



HAL
open science

Multiple front standing waves in the FitzHugh-Nagumo equations

Chao-Nien Chen, Eric Séré

► **To cite this version:**

Chao-Nien Chen, Eric Séré. Multiple front standing waves in the FitzHugh-Nagumo equations. *Journal of Differential Equations*, 2021, 302, pp. 895-925. 10.1016/j.jde.2021.08.005 . hal-01758786v3

HAL Id: hal-01758786

<https://hal.science/hal-01758786v3>

Submitted on 17 Sep 2021

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

Multiple front standing waves in the FitzHugh-Nagumo equations

Chao-Nien Chen ^{*} Éric Séré[†]

Abstract: There have been several existence results for the standing waves of FitzHugh-Nagumo equations. Such waves are the connecting orbits of an autonomous second-order Lagrangian system and the corresponding kinetic energy is an indefinite quadratic form in the velocity terms. When the system has two stable hyperbolic equilibria, there exist two stable standing fronts, which will be used in this paper as building blocks, to construct stable standing waves with multiple fronts in case the equilibria are of saddle-focus type. The idea to prove existence is somewhat close in spirit to [6]; however several differences are required in the argument: facing a strongly indefinite functional, we need to perform a nonlocal Lyapunov-Schmidt reduction; in order to justify the stability of multiple front standing waves, we rely on a more precise variational characterization of such critical points. Based on this approach, both stable and unstable standing waves are found.

Key words: reaction-diffusion system, FitzHugh-Nagumo equations, standing wave, stability, Hamiltonian system, heteroclinic connection, multibump solution, variational approach, Lyapunov-Schmidt reduction.

AMS subject classification: 34C37, 35J50, 35K57.

^{*}Department of Mathematics, National Tsing Hua University, Hsinchu 300, Taiwan (chen@math.nthu.edu.tw)

[†]CEREMADE, Université Paris-Dauphine, PSL Research University, CNRS, UMR 7534, Place de Lattre de Tassigny, F-75016 Paris, France (sere@ceremade.dauphine.fr)

1 Introduction

Following a fascinating idea of Turing [46], reaction-diffusion systems [1, 7, 9, 14, 27, 34, 37] serve as models for studying pattern formation and wave propagation. Significant progress [8, 10, 12, 13, 17, 19, 20, 18, 24, 26, 38, 41, 42, 47, 48] on the self-organized patterns has been made for the system of FitzHugh-Nagumo equations

$$u_t - du_{xx} = f(u) - v, \quad (1.1)$$

$$\tau v_t - v_{xx} = u - \gamma v. \quad (1.2)$$

Here $f(\xi) = \xi(\xi - \beta)(1 - \xi)$, $\beta \in (0, 1/2)$ and $d, \tau, \gamma \in (0, \infty)$. Historically the original model [27, 37] was derived as a simplification of the Hodgkin-Huxley equations [28] for nerve impulse propagation. In recent years (1.1)-(1.2) has been extensively studied as a paradigmatic activator-inhibitor system. Such systems are of great interest to the scientific community as breeding grounds for studying the generation of localized structures.

The standing wave solutions of (1.1)-(1.2) are the connecting orbits of a second order Lagrangian system

$$-du'' = f(u) - v, \quad (1.3)$$

$$-v'' = u - \gamma v. \quad (1.4)$$

Associated with (1.3)-(1.4), the Lagrangian is

$$L(u, v, u', v') = \frac{d}{2}(u')^2 - \frac{1}{2}(v')^2 + uv - \frac{\gamma}{2}v^2 - \int_0^u f(\xi)d\xi. \quad (1.5)$$

Since (1.3)-(1.4) is an autonomous Lagrangian system, the associated energy

$$E(u, v, u', v') := \frac{d}{2}(u')^2 - \frac{1}{2}(v')^2 - uv + \frac{\gamma}{2}v^2 + \int_0^u f(\xi)d\xi \quad (1.6)$$

is constant along any solution. Expressed in terms of the momenta $p = du'$ and $q = -v'$, this energy becomes the Hamiltonian

$$H(u, v, p, q) := \frac{1}{2d}p^2 - \frac{1}{2}q^2 - uv + \frac{\gamma}{2}v^2 + \int_0^u f(\xi)d\xi. \quad (1.7)$$

The Hamiltonian system associated with H in the phase space $\mathbb{R}^2 \times \mathbb{R}^2$ will be referred to as (HS). It is a reformulation of (1.3)-(1.4) as a system of first-order equations.

When γ is not too small, system (1.3)-(1.4) has two hyperbolic equilibria $(u_-, v_-) = (0, 0)$ and (u_+, v_+) with $u_+ > 0$, that are stable for (1.1)-(1.2). Since we are interested in standing front solutions joining such two equilibria, they

must have the same energy E , and this imposes the condition $\gamma = 9(2\beta^2 - 5\beta + 2)^{-1}$, as to be a basic assumption of the paper. Under this assumption, it is easy to see that $(u_+, v_+) = (2(\beta + 1)/3, 2(\beta + 1)/3\gamma)$ and the symmetry S with respect to the center $(u_+/2, v_+/2)$ preserves the Lagrangian. As a consequence, the linearized equations at (u_-, v_-) and (u_+, v_+) are the same.

Recall from [19] that \hat{v} is a C^∞ -function satisfying

$$\hat{v}(x) = \begin{cases} v_+ & \text{for } x \geq 1 \\ 0 & \text{for } x \leq -1. \end{cases} \quad (1.8)$$

To show the existence of standing front solutions of (1.3)-(1.4), we work with affine functional spaces of the form $H_w = w + H^1(\mathbb{R})$, with $w = \hat{v}, \hat{u}$, where $\hat{u} = -\hat{v}'' + \gamma\hat{v}$.

A crucial point is that (1.4) is a linear equation allowing to solve v from u . First, given any $\phi \in H^1(\mathbb{R})$, we denote $\mathcal{L}\phi$ the unique solution, in $H^1(\mathbb{R})$, of the equation

$$-g'' + \gamma g = \phi. \quad (1.9)$$

Direct calculation shows that \mathcal{L} is a bounded self-adjoint operator in $L^2(\mathbb{R})$. Then, for $u \in H_{\hat{u}}$, we denote $\mathcal{L}u := \hat{v} + \mathcal{L}\phi$ where $\phi = u - \hat{u} \in H^1(\mathbb{R})$. By construction, $\mathcal{L}u$ is the unique solution, in $H_{\hat{v}}$, of the equation

$$-v'' + \gamma v = u.$$

For $u \in H_{\hat{u}}$, define

$$J(u) = \int_{-\infty}^{\infty} L(u, \mathcal{L}u, u', \mathcal{L}u') dx. \quad (1.10)$$

It is not difficult to verify that

$$J(u) = \max_{v \in H_{\hat{v}}} \int_{-\infty}^{\infty} L(u, v, u', v') dx. \quad (1.11)$$

In [19] the action functional J has been employed, through a minimization argument, to obtain a basic type standing front solution of (1.1)-(1.2) as follows.

Theorem 1.1 *If $\beta \in (0, 1/2)$, $\gamma = 9(2\beta^2 - 5\beta + 2)^{-1}$ and $d > \gamma^{-2}$, there exists a standing front solution (u^*, v^*) of (1.1)-(1.2) with asymptotic behavior $(u^*, v^*) \rightarrow (u_-, v_-)$ as $x \rightarrow -\infty$ and $(u^*, v^*) \rightarrow (u_+, v_+)$ as $x \rightarrow \infty$. Moreover, u^* is a minimizer of J over $H_{\hat{u}}$.*

If we define $(u_*(x), v_*(x)) := (u^*(-x), v^*(-x))$, then (u_*, v_*) is also a standing front solution of (1.1)-(1.2), by reversibility of (1.3)-(1.4).

The goal of this paper is to construct multiple front solutions using (u^*, v^*) together with the reverse orbit (u_*, v_*) . We only deal with the case when the equilibria are of saddle-focus type; that is, the linearization of the Hamiltonian system (HS) at $(u_-, v_-, 0, 0)$, as well as $(u_+, v_+, 0, 0)$, has four complex eigenvalues $\pm\lambda \pm i\omega$ with $\lambda > 0$ and $\omega > 0$. As to be seen in the Appendix, this additional condition holds when $|\gamma^2 d - 2 - \beta\gamma| < 2\sqrt{1 + \beta\gamma}$. Summarizing, our assumptions on β, γ, d are the following:

$$\begin{aligned} 0 < \beta < 1/2, \quad \gamma = 9/(2\beta^2 - 5\beta + 2), \\ \frac{1}{\gamma^2} \max(1, 2 + \beta\gamma - 2\sqrt{1 + \beta\gamma}) < d < \frac{2 + \beta\gamma + 2\sqrt{1 + \beta\gamma}}{\gamma^2}. \end{aligned} \quad (1.12)$$

We now state the main existence result of the paper.

Theorem 1.2 *Assume that (1.12) is satisfied. Then there are two real numbers κ_+, κ_- , and, for each small $\sigma > 0$, a large constant $\Delta_\sigma > 0$, such that for any positive integer N and any sequence of positive integers $\mathbf{n} = (n_i)_{1 \leq i \leq N}$ with $n_i \geq \Delta_\sigma$ for every i , there exist positive numbers X_1, \dots, X_N and a solution $(\hat{u}_{\mathbf{n}}, \hat{v}_{\mathbf{n}})$ of (1.3)-(1.4) satisfying the following properties:*

(a) For i odd in $[1, N]$,

$$\|(\hat{u}_{\mathbf{n}}, \hat{v}_{\mathbf{n}})(\cdot + C_i) - (u_*, v_*)\|_{H^1(-A_i, A_{i+1})} < \sigma, \quad |X_i - 2\pi n_i/\omega - \kappa_+| < \sigma.$$

(b) For i even in $[0, N]$,

$$\|(\hat{u}_{\mathbf{n}}, \hat{v}_{\mathbf{n}})(\cdot + C_i) - (u^*, v^*)\|_{H^1(-A_i, A_{i+1})} < \sigma, \quad |X_i - 2\pi n_i/\omega - \kappa_-| < \sigma.$$

Here, $A_0 = +\infty$, $C_0 = 0$, $C_i = C_{i-1} + X_i$, $A_i = X_i/2$ for $1 \leq i \leq N$, and $A_{N+1} = +\infty$.

Let us remark that if N is odd, $(\hat{u}_{\mathbf{n}}, \hat{v}_{\mathbf{n}})$ is homoclinic to (u_-, v_-) while for N even, it is a heteroclinic connection between (u_-, v_-) and (u_+, v_+) . Such orbits are the standing waves of (1.1)-(1.2) with multiple fronts; for the Hamiltonian system they are often called multibump solutions.

As already mentioned, the range of parameters under consideration is such that the basic heteroclinics (u_*, v_*) and (u^*, v^*) connect two equilibria of saddle-focus type. In this situation, multi-bump solutions are known to exist provided the stable and unstable manifolds intersect transversally, as was proved by Devaney [25] by constructing a Smale horseshoe. Transversality condition in general is difficult to check for a given Hamiltonian although it is generically true. Instead of verifying transversality, we follow a strategy introduced in [6]. We first prove that any critical point of J is isolated up to translation invariance in the spatial variable, by solving an auxiliary boundary value problem. Then

we invoke this property to show the existence of multi-bump solutions by a variational argument.

The variational construction for multibump and chaotic solutions has a long history and the comments below are not exhaustive. To our knowledge, the earliest results were established by Bolotin [2, 3, 4] in the context of nonautonomous second order Lagrangian systems, the connecting orbits being minimizers of the action. In the case of twist maps on the annulus (also corresponding to non-autonomous Lagrangian systems), Mather [36] constructed chaotic connecting orbits by a minimization method in the region between two invariant circles. For non-autonomous first order Hamiltonian systems, multibump solutions were found by min-max methods [43, 44] under the assumption that critical points are isolated. This construction was extended to second order systems and elliptic PDEs in [22, 23, 21]. We refer to [39] and references therein for more recent development and related results in this direction. For autonomous problems of saddle-focus type a class of multi-bump solutions were obtained, in the special case of a fourth order equation related to water wave theory, by Buffoni [5] using a shooting argument. Subsequently a larger set of multi-bump solutions was constructed [6] by variational and degree arguments. This method was then adapted for studying the extended Fisher-Kolmogorov equations (of fourth order) [30]. In subsequent works [31, 32], a refined but more specific argument was introduced in order to obtain optimal results on the F-K model. As already mentioned, the present work is close in spirit to [6]. Note, however, that our system of autonomous second order Lagrangian equations is associated with a strongly indefinite variational problem and to our knowledge, it cannot be reduced to a fourth order equation which would allow a simpler variational interpretation. Instead, we use a nonlocal Lyapunov-Schmidt reduction. Moreover our approach is purely variational, contrary to [6] where degree theory was employed. Another novelty is our proof that *all* critical points are isolated up to translations in x , while in [6] the first step just consisted in showing that the basic one-bump solution is isolated. The purely variational construction and the stronger isolatedness property are needed for the sake of stability analysis, as always an important issue in considering pattern formation as well as wave propagation.

For the stationary solutions of (1.1)-(1.2), stability questions have been studied in [16, 17, 19, 18, 38, 49] by various methods. In conjunction with the strongly indefinite variational structure, the Maslov index [15, 17] and relative Morse index [16] provide useful information to determine the stability of such solutions, obtained as the critical points of the action functional. Let $\mathbb{C}^- = \{\zeta \mid \zeta \in \mathbb{C} \text{ and } \operatorname{Re}\zeta < 0\}$, where $\operatorname{Re}\zeta$ denotes the real part of ζ . Denote by Λ the linearization of (1.1)-(1.2) at a standing wave solution (u, v) . A standing wave (u, v) is said to be non-degenerate if zero is a simple eigenvalue of Λ .

Definition 1.1 *A non-degenerate standing wave (u, v) of (1.1)-(1.2) is spectrally stable if all the non-zero eigenvalues of Λ are in \mathbb{C}^- .*

The following result follows immediately from an index method developed in [11]:

Theorem 1.3 *Let (u, v) be a non-degenerate standing wave of (1.1)-(1.2). Suppose u is a local minimizer of J then (u, v) is spectrally stable, provided that $\tau < \gamma^2$.*

Note, however, that the non-degeneracy of a standing wave is equivalent to the transversality of the stable and unstable manifolds, and we are unable to prove such a property. Fortunately, we can go beyond the spectral stability analysis, thanks to a Lyapunov functional introduced in [18] in a slightly different context. In Section 5 we shall give an extension of this Lyapunov functional, which can be applied to the standing waves of (1.1)-(1.2). Let us remark that the standing waves are in affine subspaces of $H_{loc}^1(\mathbb{R}) \times H_{loc}^1(\mathbb{R})$ having $H^1(\mathbb{R}) \times H^1(\mathbb{R})$ as underlying vector space. The norm of $H^1(\mathbb{R}) \times H^1(\mathbb{R})$ induces the natural metric on such affine spaces, and we shall study the dynamical stability of the standing waves for this metric.

Theorem 1.4 *Assume that (1.12) is satisfied and let $\tau < \gamma^2$. Under the flow generated by (1.1)-(1.2) on the affine space $(\hat{u}_{\mathbf{n}} + H^1(\mathbb{R})) \times (\hat{v}_{\mathbf{n}} + H^1(\mathbb{R}))$, the standing wave $(\hat{u}_{\mathbf{n}}, \hat{v}_{\mathbf{n}})$ is asymptotically stable for the $H^1(\mathbb{R}) \times H^1(\mathbb{R})$ metric, up to a phase shift in spatial variable. More precisely, there is $\rho_{\mathbf{n}} > 0$ such that if $(u(x, t), v(x, t))$ is a solution of (1.1)-(1.2) and*

$$\|u(\cdot, 0) - \hat{u}_{\mathbf{n}}\|_{H^1(\mathbb{R})} + \|v(\cdot, 0) - \hat{v}_{\mathbf{n}}\|_{H^1(\mathbb{R})} < \rho_{\mathbf{n}},$$

then

$$\inf_{y \in \mathbb{R}} \{ \|u(\cdot, t) - \hat{u}_{\mathbf{n}}(\cdot - y)\|_{H^1(\mathbb{R})} + \|v(\cdot, t) - \hat{v}_{\mathbf{n}}(\cdot - y)\|_{H^1(\mathbb{R})} \} \xrightarrow{t \rightarrow +\infty} 0.$$

As a final remark, there are plenty of unstable standing waves; however we do not attempt to describe them all. We just state a result in the two-bump case.

Theorem 1.5 *As in Theorem 1.2, assume (1.12) and take a small $\sigma > 0$ and a large constant Δ_{σ} . For any positive integer $n \geq \Delta_{\sigma}$ there exists a solution $(\check{u}_n, \check{v}_n)$ of (1.3)-(1.4) such that, for some $X \in \mathbb{R}$ with $|X - \pi(2n+1)/\omega - \kappa_+| < \sigma$ and κ_+ as in Theorem 1.2, the following properties hold:*

- (i) $\|(\check{u}_n, \check{v}_n) - (u^*, v^*)\|_{H^1(-\infty, X/2)} < \sigma$.
- (ii) $\|(\check{u}_n, \check{v}_n) - (u_*, v_*)(\cdot - X)\|_{H^1(X/2, +\infty)} < \sigma$.

(iii) If $\tau < \gamma^2$ then $(\check{u}_n, \check{v}_n)$ is unstable in the following sense: there is $\epsilon_n > 0$ such that, for any $\rho > 0$, there exist $\tau_n(\rho) > 0$ and a solution $(u(x, t), v(x, t))$ of (1.1)-(1.2) satisfying

$$\|u(\cdot, 0) - \check{u}_n\|_{H^1(\mathbb{R})} + \|v(\cdot, 0) - \check{v}_n\|_{H^1(\mathbb{R})} < \rho,$$

while for all $t \geq \tau_n(\rho)$, one has

$$\inf_{y \in \mathbb{R}} \left\{ \|u(\cdot, t) - \check{u}_n(\cdot - y)\|_{H^1(\mathbb{R})} + \|v(\cdot, t) - \check{v}_n(\cdot - y)\|_{H^1(\mathbb{R})} \right\} \geq \epsilon_n.$$

The solution $(\check{u}_n, \check{v}_n)$ will be found in Section 6 by a mountain-pass type mini-max method. In the proof of (iii) (as well as in the proof of Theorem 1.4), a crucial ingredient is Proposition 3.1, stating that *any* critical point of J is isolated up to translation.

2 Preliminaries

In this section we recall the variational setting [19] used to study (u^*, v^*) and discuss related properties, including a reduced functional J which is bounded from below. In the sequel, we work with affine functional spaces of the form $H_w = w + H^1(\mathbb{R})$, with $w = 0, u_+, v_+, u_*, u^*, v_*$ or v^* . For $a_u = 0, u_+, u_*, u^*$ respectively, and $u \in H_{a_u}$, we also denote $\mathcal{L}u := a_v + \mathcal{L}(u - a_u)$, with $a_v = 0, v_+, v_*, v^*$ respectively. Let us remark that $\mathcal{L}u$ is the unique solution, in H_{a_v} , of the equation

$$-v'' + \gamma v = u.$$

The operator \mathcal{L} has a simple expression:

Lemma 2.1 *Let $G(x) = \frac{1}{2\sqrt{\gamma}} e^{-\sqrt{\gamma}|x|}$. If $u \in H_{a_u}$ with $a_u \in \{0, u_+, u^*, u_*\}$, then $\mathcal{L}u = G * u$.*

Proof. The function G is in $W^{1,1}(\mathbb{R})$ and is a fundamental solution of the operator $-\frac{d^2}{dx^2} + \gamma$. If $u \in H_{a_u}$ then it is continuous and bounded on \mathbb{R} , so $G * u(x) = \int_{\mathbb{R}} G(x-y)u(y)dy$ is a well-defined function of x . This function is of class C^2 , and solves $-v'' + \gamma v = u$. In particular, $G * a_u$ solves $-v'' + \gamma v = a_u$. In addition, $\lim_{x \rightarrow \pm\infty} G * a_u(x) = \lim_{x \rightarrow \pm\infty} a_v(x)$, since $\int_{\mathbb{R}} G = \frac{1}{\gamma}$. So, by the maximum principle, $G * a_u = a_v$. Finally, since $u - a_u \in H^1(\mathbb{R})$, $G * (u - a_u) \in H^1(\mathbb{R})$. We conclude that $G * u = \mathcal{L}u$. \blacksquare

Recall from (1.5) that the Lagrangian associated with (1.3)-(1.4) is $L(u, v, u', v')$. Note that the main difference with [6] is that the present system does not seem to be reducible to an almost linear, fourth-order system having a simple

variational interpretation. So one has to deal with an indefinite Lagrangian (1.5). Fortunately, this Lagrangian is concave in (v, v') . We exploit this property as follows: observe that for $u \in H_{a_u}$ and $v = \mathcal{L}u \in H_{a_v}$,

$$\begin{aligned} \left\| \frac{d}{dx}(v - a_v) \right\|_{L^2(\mathbb{R})}^2 + \gamma \|v - a_v\|_{L^2(\mathbb{R})}^2 &= \int_{-\infty}^{\infty} (u - a_u)(v - a_v) dx \\ &\leq \|u - a_u\|_{L^2(\mathbb{R})} \|v - a_v\|_{L^2(\mathbb{R})}. \end{aligned}$$

Hence there is a $C_0 > 0$ such that

$$\|\mathcal{L}u - a_v\|_{H^1(\mathbb{R})} \leq C_0 \|u - a_u\|_{L^2(\mathbb{R})}. \quad (2.1)$$

Given $\phi \in H^1(\mathbb{R})$, define, for all $\psi \in H^1(\mathbb{R})$,

$$I(\psi) = \int_{-\infty}^{\infty} \left(\frac{1}{2} |\psi'|^2 + \frac{\gamma}{2} \psi^2 - \phi \psi \right) dx$$

Lemma 2.2 *Let $\phi \in H^1(\mathbb{R})$. Then*

$$I(\psi) - I(\mathcal{L}\phi) = \frac{1}{2} \int_{-\infty}^{\infty} \{ (\psi' - (\mathcal{L}\phi)')^2 + \gamma (\psi - \mathcal{L}\phi)^2 \} dx$$

for all $\psi \in H^1(\mathbb{R})$.

Proof. It follows from straightforward calculation, by making use of

$$\int_{-\infty}^{\infty} ((\mathcal{L}\phi')^2 + \gamma (\mathcal{L}\phi)^2) dx = \int_{-\infty}^{\infty} \phi \mathcal{L}\phi dx$$

and

$$\int_{-\infty}^{\infty} (\psi' \mathcal{L}\phi' + \gamma \psi \mathcal{L}\phi) dx = \int_{-\infty}^{\infty} \phi \psi dx.$$

■

For $a_u \in \{0, u_+, u^*, u_*\}$ and all $u \in H_{a_u}$, define

$$J(u) = \int_{-\infty}^{\infty} L(u, \mathcal{L}u, u', \mathcal{L}u') dx. \quad (2.2)$$

The next lemma is an immediate consequence of Lemma 2.2.

Lemma 2.3 *Taking (a_u, a_v) in $\{(0, 0); (u_+, v_+); (u^*, v^*); (u_*, v_*)\}$, if $u \in H_{a_u}$ then:*

$$J(u) = \max_{v \in H_{a_v}} \int_{-\infty}^{\infty} L(u, v, u', v') dx. \quad (2.3)$$

The following proposition, which was proved in [19], implies that all the critical values of J on H_{u^*} must be positive, since $d - \frac{1}{\gamma^2} > 0$.

Proposition 2.4 *If $u \in H_w$ (with $w = u_-, u_+, u_*$ or u^*) and $v = \mathcal{L}u$, then*

$$J(u) = \int_{-\infty}^{\infty} \left\{ \frac{1}{2} \left[\left(d - \frac{1}{\gamma^2} \right) (u')^2 + \left(v' - \frac{u'}{\gamma} \right)^2 + \gamma \left(v - \frac{u}{\gamma} \right)^2 \right] + \frac{1}{4} u^2 (u - u_+)^2 \right\} dx.$$

If $J(u) \leq C$ then Proposition 2.4 gives an upper bound on $\|u'\|_{L^2}$, $\|(\mathcal{L}u)'\|_{L^2}$, $\|\gamma \mathcal{L}u - u\|_{L^2}$ and $\|\text{dist}(u, \{0, u_+\})\|_{L^2}$. If, in addition, u is a critical point of J , then, invoking (1.3)-(1.4), we obtain an L^∞ bound for $(u, \mathcal{L}u, u', \mathcal{L}u')$:

Corollary 2.5 *Assume that $u \in H_w$ (with $w = u_-, u_+, u_*$ or u^*) is a critical point of J and that $J(u) \leq C$ for some $C > 0$. Then there is a $C' > 0$, depending only on C , such that*

$$\|(u - w, \mathcal{L}(u - w), u' - w', \mathcal{L}(u' - w'))\|_{L^\infty} \leq C'.$$

Moreover, in the case $w \in \{u_-, u_+\}$ then $\lim_{C \rightarrow 0} C' = 0$, while if $w \in \{u_*, u^*\}$, the upper bound C cannot approach zero.

The choice of parameters made in this paper implies that the Hamiltonian system (HS) given by (1.7) has the same linearization at $z_- = (u_-, v_-, 0, 0)$ and $z_+ = (u_+, v_+, 0, 0)$. Moreover these two points are equilibria of saddle-focus type. In other words, the linearized Hamiltonian system at z_\pm possesses four complex eigenvalues $\mu e^{i\varpi}, \mu e^{-i\varpi}, -\mu e^{i\varpi}, -\mu e^{-i\varpi}$ with $\mu > 0$ and $\varpi \in (0, \frac{\pi}{2})$. Now, we may consider the corresponding linearization of the Lagrangian system (1.3)-(1.4) at (u_-, v_-) or (u_+, v_+) . This linearization takes the form $h'' = Bh$, where $h(x)$ is a column vector with two components and B is a square matrix of order 2 having two simple eigenvalues: $\mu^2 e^{2i\varpi}$ and $\mu^2 e^{-2i\varpi}$. As a consequence, there exists a change-of-basis matrix P such that, taking $Y_\pm = P(u - u_\pm, v - v_\pm)^T$, (1.3)-(1.4) becomes $Y''_\pm = \mu^2 R_{2\varpi} Y_\pm + O(|Y_\pm|^2)$. Here, $R_{2\varpi}$ is the matrix of a rotation with angle 2ϖ in \mathbb{R}^2 and the notation T means transposition. In the sequel we denote λ and ω the real and imaginary parts of $\mu e^{i\varpi}$.

With $z_\pm := (u_\pm, v_\pm, 0, 0)$ being hyperbolic equilibria, a non-constant solution of (1.3)-(1.4) cannot stay near them for all values of x . This together with Corollary 2.5 gives a positive lower bound ℓ for the value of J at non-constant critical points. Let $|\cdot|$ be the euclidean norm in \mathbb{R}^2 . We list below a number of properties for the Hamiltonian system (HS).

Proposition 2.6 *There exist positive numbers ρ_0 and ℓ such that the following properties hold.*

(i) *Suppose $z(x) = (u(x), v(x), p(x), q(x))$ is a non-constant maximal solution of (HS). Then:*

-If $|P(u(x) - u_{\pm}, v(x) - v_{\pm})^T| < \rho_0$ for all $x \leq 0$, then $z(0) \in W^u(z_{\pm})$, and there is an $a_1 > 0$ such that $|P(u(a_1) - u_{\pm}, v(a_1) - v_{\pm})^T| = \rho_0$ and

$$\frac{d}{dx}|P(u(x) - u_{\pm}, v(x) - v_{\pm})^T| \geq \frac{\lambda}{2}|P(u(x) - u_{\pm}, v(x) - v_{\pm})^T| \quad (2.4)$$

for all $x \leq a_1$.

-If $|P(u(x) - u_{\pm}, v(x) - v_{\pm})^T| < \rho_0$ for all $x \geq 0$, then $z(0) \in W^s(z_{\pm})$, and there is an $a_2 < 0$ such that $|P(u(a_2) - u_{\pm}, v(a_2) - v_{\pm})^T| = \rho_0$ and

$$\frac{d}{dx}|P(u(x) - u_{\pm}, v(x) - v_{\pm})^T| \leq -\frac{\lambda}{2}|P(u(x) - u_{\pm}, v(x) - v_{\pm})^T| \quad (2.5)$$

for all $x \geq a_2$.

(ii) *Let u_c be a non-constant critical point of J and $u_{s_{\pm}} = \lim_{x \rightarrow \pm\infty} u_c(x)$, with $s_{\pm} \in \{-, +\}$. If $0 < \rho \leq \rho_0$, there exist $x_1 < x_2$ such that for any $x \in (-\infty, x_1)$, $P(u_c, \mathcal{L}u_c)^T(x)$ lies in the open disk D_- of center $P(u_{s_-}, v_{s_-})^T$ with radius ρ and $P(u_c, \mathcal{L}u_c)^T(x_1)$ sits on the boundary of D_- , while for any $x \in (x_2, \infty)$, $P(u_c, \mathcal{L}u_c)^T(x)$ lies in the open disk D_+ of center $P(u_{s_+}, v_{s_+})^T$ with radius ρ and $P(u_c, \mathcal{L}u_c)^T(x_2)$ sits on the boundary of D_+ .*

(iii) *If u_c is a non-constant critical point of J in H_w with $w = u_-, u_+, u_*$ or u^* , then $J(u_c) \geq \ell$. Moreover $\|u_c - u_{\pm}\|_{H^1} \geq \ell$ in case $w = u_{\pm}$.*

Proof. We prove (i) when $|P(u(x) - u_{\pm}, v(x) - v_{\pm})^T| < \rho_0$ for all $x \leq 0$; the other case immediately follows from time reversal. With $z_{\pm} = (u_{\pm}, v_{\pm}, 0, 0)$ being hyperbolic equilibria, if ρ_0 is small enough, then, by the Hartman-Grossman Theorem, $z(0)$ must lie in the local unstable manifold $W_{\text{loc}}^u(z_{\pm})$, which is an embedded submanifold tangent to the unstable space $E^u(z_{\pm})$, by the Stable Manifold Theorem (see e.g. [33]).

Let $\dot{z}(x) = (\dot{u}(x), \dot{v}(x), \dot{p}(x), \dot{q}(x))$ be a solution of the linearization of (HS) at z_{\pm} . If \dot{z} lies in $E^u(z_{\pm})$, then $\dot{Y}(x) := P(\dot{u}(x), \dot{v}(x))^T$ satisfies the first-order equation $\dot{Y}' = \mu R_{\infty} \dot{Y}$ and this implies that $\frac{d}{dx}|\dot{Y}| = \lambda|\dot{Y}|$. Then, for the nonlinear flow of (HS), (2.4) follows for small ρ_0 . Note that in (2.4), $\frac{\lambda}{2}$ could be replaced by any number $\lambda' < \lambda$, at the expense of choosing a smaller ρ_0 when λ' is closer to λ .

Next we prove (ii). For fixed $0 < \rho \leq \rho_0$, let $X \in \mathbb{R}$ be such that $|P(u_c(x) - u_{s_-}, \mathcal{L}u_c(x) - v_{s_-})^\top| < \rho$ for all $x \leq -X$ and $z(x) := (u_c(x - X), \mathcal{L}u_c(x - X), du'_c(x - X), -\mathcal{L}u'_c(x - X))$.

Since z satisfies the assumptions of (i) corresponding to the case $z(0) \in W^u(z_{s_-})$, we may set

$$x_1 = \max\{x \leq a_1 - X : \forall y < x, |P(u_c(y) - u_{s_-}, \mathcal{L}u_c(y) - v_{s_-})^\top| < \rho\},$$

where a_1 was defined by (i). Clearly x_1 satisfies the required properties and x_2 can be found in a similar way.

Finally, since $P(u_c(x_1), \mathcal{L}u_c(x_1))^\top$ has reached the boundary of D_- , Corollary 2.5 implies an estimate of the form $J(u_c) \geq \ell$ for some $\ell > 0$. For the same reason, if $u_c \in H_{u_\pm}$, $\|u_c - u_\pm\|_{H^1}$ cannot be too small: indeed, the H^1 norm controls the L^∞ norm. This proves (iii). \blacksquare

Next we analyze the behavior of Palais-Smale sequences of J . Recall the classical notion of Palais-Smale sequence.

Definition 2.1 *A sequence $\{u_m\}$ in H_w , with $w \in \{u_\pm, u_*, u^*\}$, is called a Palais-Smale sequence for J if $\{J(u_m)\}$ is convergent and $J'(u_m) \rightarrow 0$ in $H^{-1}(\mathbb{R})$ as $m \rightarrow \infty$.*

Due to translation invariance in x , a Palais-Smale sequence does not necessarily have a convergent subsequence for the $H^1(\mathbb{R})$ metric. However, adapting the arguments of [19], we obtain the following result in the spirit of the concentration-compactness theory by Pierre-Louis Lions [35].

Proposition 2.7 *Let ρ_0 be the number introduced in Proposition 2.6, $w \in \{u_\pm, u_*, u^*\}$ and $u_{s_-} := \lim_{x \rightarrow -\infty} w(x)$, $s_- \in \{-, +\}$. Suppose that $\{u_m\}$ is a Palais-Smale sequence for J in H_w such that $\lim_{m \rightarrow \infty} J(u_m) > 0$. Then there is a $\rho_1 \in (0, \rho_0)$ independent of $\{u_m\}$ and such that, for any fixed $\rho \in (0, \rho_1)$, the following properties hold.*

(i) *For m large enough, there exists $x_m \in \mathbb{R}$ such that, for all $x \in (-\infty, x_m)$, $P(u_m, \mathcal{L}u_m)^\top(x)$ lies in the open disk D_- of center $P(u_{s_-}, \mathcal{L}u_{s_-})^\top$ with radius ρ and $P(u_m, \mathcal{L}u_m)^\top(x_m)$ sits on the boundary of D_- . Moreover, after extraction, $\tilde{u}_m := u_m(\cdot + x_m)$ converges in H_{loc}^1 to a non-constant critical point $u^{(1)}$ of J , with $\lim_{x \rightarrow -\infty} u^{(1)}(x) = u_{s_-}$, while $u_{s_+} := \lim_{x \rightarrow +\infty} u^{(1)}(x)$ may differ from $\lim_{x \rightarrow +\infty} w(x)$.*

(ii) *If $\liminf \|\tilde{u}_m - u^{(1)}\|_{H^1} > 0$, i.e. the convergence of $\{\tilde{u}_m\}$ does not hold for the $H^1(\mathbb{R})$ metric, there exist two sequences $\underline{t}_m \rightarrow \underline{t} \in \mathbb{R}$ and $\bar{t}_m \rightarrow +\infty$ such that for all $x \in (\underline{t}_m, \bar{t}_m)$, $P(\tilde{u}_m, \mathcal{L}\tilde{u}_m)^\top(x)$ lies in the open*

disk \tilde{D} of center $P(u_{s_+}, v_{s_+})^T$ with radius ρ , while $P(\tilde{u}_m, \mathcal{L}\tilde{u}_m)^T(\underline{t}_m)$ and $P(\tilde{u}_m, \mathcal{L}\tilde{u}_m)^T(\bar{t}_m)$ both sit on the boundary of \tilde{D} . Furthermore along a subsequence, $\{\tilde{u}_m(\cdot + \bar{t}_m)\}$ converges in H_{loc}^1 to a non-constant critical point $u^{(2)}$ of J . Then $u_{s_+} = \lim_{x \rightarrow +\infty} u^{(1)}(x) = \lim_{x \rightarrow -\infty} u^{(2)}(x)$. Moreover

$$\ell \leq J(u^{(i)}) \leq \lim J(u_m) - \ell \quad \text{for } i \in \{1, 2\}, \quad \text{and } \limsup \|\tilde{u}_m - u^{(1)}\|_{H^1} \geq \ell.$$

Proof. Since $f(\xi) + \beta\xi = o(\xi)$ as $\xi \rightarrow 0$, there exists $\rho_1 \in (0, \rho_0)$ such that $|P(u, v)^T| \leq \rho_1$ implies $-f(u)u \geq \frac{\beta u^2}{2}$. Then, by symmetry of the Lagrangian with respect to $(u_+/2, v_+/2)$, $|P(u - u_+, v - v_+)^T| \leq \rho_1$ implies $(v_+ - f(u))(u - u_+) \geq \frac{\beta(u - u_+)^2}{2}$.

To show the existence of x_m for m large enough, we assume by contradiction that for any integer m_0 , one can find $m \geq m_0$ such that for all $x \in \mathbb{R}$, $P(u_m, \mathcal{L}u_m)^T(x)$ lies in D_- . This is possible only if $w = u_{s_-}$. Taking ρ_1 small enough, one has $\text{dist}(P^{-1}D_-, (u_{s_-}, v_{s_-})^T) > 0$, so the boundedness of $\{J(u_m)\}$ together with Proposition 2.4 implies that

$$\|u_m - u_{s_-}\|_{H^1} + \|\mathcal{L}u_m - v_{s_-}\|_{H^1} = O(1).$$

Now, by definition of \mathcal{L} , $\hat{e}_m := -(\mathcal{L}u_m)'' + \gamma\mathcal{L}u_m - u_m = 0$. Since $J'(u_m) \rightarrow 0$, one also has $e_m := -du_m'' - f(u_m) + \mathcal{L}u_m \rightarrow 0$ in H^{-1} . Combining these estimates, one gets

$$\epsilon_m := \langle e_m, u_m - u_{s_-} \rangle + \langle \hat{e}_m, \mathcal{L}u_m - v_{s_-} \rangle = o(1). \quad (2.6)$$

On the other hand, by direct calculation,

$$\epsilon_m = \int_{-\infty}^{\infty} \{d(u_m')^2 + (\mathcal{L}u_m')^2 + \gamma(\mathcal{L}u_m - v_{s_-})^2 + (v_{s_-} - f(u_m))(u_m - u_{s_-})\} \quad (2.7)$$

Since we assumed that $P(u_m, \mathcal{L}u_m)^T(x)$ lies in D_- for all $x \in \mathbb{R}$, invoking (2.6)-(2.7) yields

$$\int_{-\infty}^{\infty} \{d(u_m')^2 + (\mathcal{L}u_m')^2 + \gamma(\mathcal{L}u_m - v_{s_-})^2 + \frac{\beta}{2}(u_m - u_{s_-})^2\} \leq o(1).$$

This implies that as $m_0 \rightarrow \infty$, $\|u_m - u_{s_-}\|_{H^1} \rightarrow 0$ and consequently $J(u_m) \rightarrow 0$, which violates the assumption $\liminf_{m \rightarrow \infty} J(u_m) > 0$. We have thus proved the existence of x_m for m large.

Next, set $\tilde{u}_m := u_m(\cdot + x_m)$. Since $J(\tilde{u}_m)$ is bounded and $J'(\tilde{u}_m) \rightarrow 0$, using an argument from [19] yields a critical point $u^{(1)}$ of J and a subsequence, still denoted by $\{\tilde{u}_m\}$, such that $\{(\tilde{u}_m, \mathcal{L}\tilde{u}_m)\}$ converges to $(u^{(1)}, \mathcal{L}u^{(1)})$ in $H_{\text{loc}}^1(\mathbb{R})$. Note that $u^{(1)}$ satisfies $|P(u^{(1)}(x) - u_{s_-}, \mathcal{L}u^{(1)}(x) - v_{s_-})^T| \leq \rho$ for $x \in (-\infty, 0)$ and $|P(u^{(1)}(0) - u_{s_-}, \mathcal{L}u^{(1)}(0) - v_{s_-})^T| = \rho$. By Proposition 2.6(i), it is clear that $(u^{(1)}, \mathcal{L}u^{(1)})(x)$ converges to (u_{s_-}, v_{s_-}) as $x \rightarrow -\infty$.

Assuming that the convergence of \tilde{u}_m to $u^{(1)}$ does not hold for the $H^1(\mathbb{R})$ metric, we now prove (ii). Let $u_{s_+} := \lim_{x \rightarrow +\infty} u^{(1)}(x)$ ($s_+ \in \{-, +\}$). Let $b_1 > 0$ be such that

$$|P(u^{(1)}(x) - u_{s_+}, \mathcal{L}u^{(1)}(x) - v_{s_+})^\top| < \rho/2 \text{ for all } x \geq b_1. \quad (2.8)$$

By Proposition 2.6(ii), there is $b_2 < b_1$ such that $|P(u^{(1)}(b_2) - u_{s_+}, \mathcal{L}u^{(1)}(b_2) - v_{s_+})^\top| = \rho_0$. Since $\rho_0 > \rho_1 > \rho$ and $\{(\tilde{u}_m, \mathcal{L}\tilde{u}_m)\}$ uniformly converges to $(u^{(1)}, \mathcal{L}u^{(1)})$ on $[b_2, b_1]$, it follows that $|P(\tilde{u}_m(b_2) - u_{s_+}, \mathcal{L}\tilde{u}_m(b_2) - v_{s_+})^\top| > \rho$ and $|P(\tilde{u}_m(b_1) - u_{s_+}, \mathcal{L}\tilde{u}_m(b_1) - v_{s_+})^\top| < \rho$ for m large enough. Set

$$\underline{t}_m := \max\{x \leq b_1 : |P(\tilde{u}_m(x) - u_{s_+}, \mathcal{L}\tilde{u}_m(x) - v_{s_+})^\top| \geq \rho\}.$$

Then $b_2 < \underline{t}_m < b_1$ and we may assert that $\underline{t}_m \rightarrow \underline{t}$ after extracting a subsequence. By uniform convergence of $\{(\tilde{u}_m, \mathcal{L}\tilde{u}_m)\}$ on compact sets, we conclude that $|P(u^{(1)}(\underline{t}) - u_{s_+}, \mathcal{L}u^{(1)}(\underline{t}) - v_{s_+})^\top| = \rho$ and $|P(u^{(1)}(x) - u_{s_+}, \mathcal{L}u^{(1)}(x) - v_{s_+})^\top| \leq \rho$ for all $x \geq \underline{t}$.

Set $u_m^{(2)}(x) := \tilde{u}_m(x) - u^{(1)}(x) + u_{s_+}$. Then $\lim_{x \rightarrow -\infty} u_m^{(2)}(x) = u_{s_+}$, $\{u_m^{(2)}\}$ is a Palais-Smale sequence for J , and $\liminf J(u_m^{(2)}) > 0$ follows from the fact that $\{\tilde{u}_m\}$ fails to converge for the $H^1(\mathbb{R})$ metric. Let \hat{D} be the disk of center $P(u_{s_+}, v_{s_+})^\top$ with radius $(\rho_1 + \rho)/2$. Replacing u_m by $u_m^{(2)}$ and D_- by \hat{D} in the contradiction argument given above, we find that for m sufficiently large $P(u_m^{(2)}, \mathcal{L}u_m^{(2)})^\top(x)$ has to exit from \hat{D} . But $(u_m^{(2)}, \mathcal{L}u_m^{(2)})$ converges uniformly to $(u_{s_+}, \mathcal{L}u_{s_+})$ on any bounded interval, so if we denote $x_m^{(2)}$ the first exit date, we have $\lim_{m \rightarrow \infty} x_m^{(2)} = +\infty$. As a consequence, $P(\tilde{u}_m, \mathcal{L}\tilde{u}_m)^\top(x)$ cannot stay in \tilde{D} for all the values of x in the interval $[b_1, \infty)$, and the first exit date

$$\bar{t}_m := \min\{x \geq b_1 : |P(\tilde{u}_m(x) - u_{s_+}, \mathcal{L}\tilde{u}_m(x) - v_{s_+})^\top| \geq \rho\}$$

satisfies $\bar{t}_m \rightarrow \infty$.

Then, along a subsequence, $\{u_m^{(2)}(\cdot + \bar{t}_m)\}$ converges in $H_{\text{loc}}^1(\mathbb{R})$ to a non-constant critical point $u^{(2)}$. By Proposition 2.6(i), $\lim_{x \rightarrow -\infty} (u^{(2)}, \mathcal{L}u^{(2)})(x) = (u_{s_+}, v_{s_+})$, since $P(u^{(2)}, \mathcal{L}u^{(2)})^\top(x)$ stays in the closed disk of center $P(u_{s_+}, v_{s_+})^\top$ with radius ρ for all $x \leq 0$. This completes the proof of (ii).

By the construction of $u_m^{(2)}$, $J(u^{(1)}) + \lim J(u_m^{(2)}) = \lim J(u_m)$ and $J(u^{(2)}) \leq \lim J(u_m^{(2)})$, so $J(u^{(1)}) + J(u^{(2)}) \leq \lim J(u_m)$. We also have $\|u^{(2)}\|_{H^1} \leq \limsup \|\tilde{u}_m - u^{(1)}\|_{H^1}$. Combining these inequalities with Proposition 2.6(iii), we conclude that $\ell \leq J(u^{(i)}) \leq \lim J(u_m) - \ell$ for $i = 1, 2$ and $\limsup \|\tilde{u}_m - u^{(1)}\|_{H^1} \geq \ell$. This ends the proof of Proposition 2.7. \blacksquare

A consequence of Proposition 2.7 is the following local Palais-Smale compactness property:

Corollary 2.8 *Let ℓ be the number as defined in Proposition 2.6. Suppose a Palais-Smale sequence satisfies $\|u_m - u_n\|_{H^1} \leq \ell/2$ for all m, n , then it has a convergent subsequence for the $H^1(\mathbb{R})$ metric.*

Proof. The corollary will be proved by contradiction. The idea behind the proof is the following: a Palais-Smale sequence $\{u_m\}$ is always precompact for the H^1_{loc} topology. If it is not precompact for the global $H^1(\mathbb{R})$ metric, then some “mass” has to escape at infinity. But Proposition 2.6 and Proposition 2.7 imply that this “mass” has an H^1 norm at least equal to ℓ , so the inequality $\|u_m - u_n\|_{H^1} \leq \ell/2$ cannot hold for all m, n .

More precisely, for the Palais-Smale sequence $\{u_m\}$, we pick out the point x_m from each u_m as in Proposition 2.7(i) and we consider a non-constant critical point $u^{(1)}$ of J that is the limit (after extraction if necessary) of $\{\tilde{u}_m\} := \{u_m(\cdot + x_m)\}$ for the $H^1_{\text{loc}}(\mathbb{R})$ topology. Let us assume that $\{u_m\}$ has no convergent subsequence in the H^1 metric. Then three cases may occur:

If $\{x_m\}$ has a convergent subsequence of limit \bar{x} , then the convergence of $\{\tilde{u}_m\}$ to $u^{(1)}$ fails in the $H^1(\mathbb{R})$ metric. Hence $\liminf \|\tilde{u}_m - u^{(1)}\|_{H^1} > 0$. Applying Proposition 2.7(ii), we conclude that $\limsup \|u_m - u^{(1)}(\cdot - \bar{x})\|_{H^1} \geq \ell$.

Next we turn to the case when, after extraction, $\lim x_m = +\infty$. Taking the same definition of the constant u_{s_-} as in Proposition 2.7, we see that $\{u_m\}$ converges to u_{s_-} in $H^1_{\text{loc}}(\mathbb{R})$, and $\limsup_{m \rightarrow \infty} \|u_m - u_{s_-}\|_{H^1} \geq \|u^{(1)} - u_{s_-}\|_{H^1} \geq \ell$.

The third case is when, after extraction, $\lim x_m = -\infty$. After a new extraction if necessary, $\{u_m\}$ converges to a critical point u of J , which can be constant or non-constant. Let $u_s := \lim_{x \rightarrow -\infty} u(x)$, with $s \in \{-, +\}$. Then $\limsup_{m \rightarrow \infty} \|u_m - u\|_{H^1} \geq \|u^{(1)} - u_s\|_{H^1} \geq \ell$.

In each of the three cases, we have found, after extraction, that $\{u_m\}$ converges to a limit u in $H^1_{\text{loc}}(\mathbb{R})$, but $\limsup_{m \rightarrow \infty} \|u_m - u\|_{H^1(\mathbb{R})} \geq \ell$. Now, let m_0 be such that $\|u_{m_0} - u\|_{H^1(\mathbb{R})} \geq 3\ell/4$ and let R_0 be such that $\|u_{m_0} - u\|_{H^1(-R_0, R_0)} \geq 2\ell/3$. By local convergence, for n large enough we have $\|u_n - u\|_{H^1(-R_0, R_0)} < \ell/6$, hence $\|u_{m_0} - u_n\|_{H^1(-R_0, R_0)} > \ell/2$. This contradicts the assumption of the corollary. \blacksquare

In the construction of multi-front solutions, the trajectories between two fronts will need to be in good control. Such trajectories are very close to one of the two stable equilibria with asymptotical behavior being dominated by the linearized equations. Recall from (1.6) that if (u, v) is a solution of (1.3)-(1.4) then $E(u', v', u, v)$ is constant along the trajectory. In particular, for any critical point u of J , the energy $E(u', \mathcal{L}u', u, \mathcal{L}u)$ is identically zero. We now state a lemma in the same spirit as Lemmas 3.1, A.1, A.2 and A.3 of [6].

Lemma 2.9 *Let $\nu \in (0, \frac{\pi}{2\omega})$ and δ be small positive numbers. For η and ζ in \mathbb{R}^2 let us denote (r_η, θ_η) and (r_ζ, θ_ζ) the polar coordinates of $P(\eta_1 - u_+, \eta_2 - v_+)^T$ and $P(\zeta_1 - u_+, \zeta_2 - v_+)^T$, respectively. There exists a small radius $\bar{r} > 0$ such*

that, if r_η and r_ζ are smaller than \bar{r} , then the boundary value problem

$$-du'' = f(u) - v, \quad (2.9)$$

$$-v'' = u - \gamma v, \quad (2.10)$$

$$(u(0), v(0)) = (\eta_1, \eta_2), \quad (u(T), v(T)) = (\zeta_1, \zeta_2), \quad (2.11)$$

has a unique solution in a small neighborhood of (u_+, v_+) . Denoting this solution by $(\bar{U}, \bar{V})_{T, \eta_1, \eta_2, \zeta_1, \zeta_2}(\cdot)$, the following estimate holds on the interval $(0, T)$:

$$|P((\bar{U}, \bar{V})_{T, \eta_1, \eta_2, \zeta_1, \zeta_2}(x) - (u_+, v_+))^T| < \min\{r_\eta, r_\zeta\}.$$

Moreover, if $E_{\eta_1, \eta_2, \zeta_1, \zeta_2}(T)$ denotes the energy of the solution $(\bar{U}, \bar{V})_{T, \eta_1, \eta_2, \zeta_1, \zeta_2}$, then the following properties hold (recalling that the four eigenvalues of the linearization of (HS) at the saddle-focus equilibrium $(u_+, v_+, 0, 0)$ have $\pm\omega$ as imaginary parts):

- (i) For T large enough, $E_{\eta_1, \eta_2, \zeta_1, \zeta_2}(T) > 0$ if $\cos(\theta_\zeta - \theta_\eta - T\omega + \varphi_+) > \delta$, while $E_{\eta_1, \eta_2, \zeta_1, \zeta_2}(T) < 0$ if $\cos(\theta_\zeta - \theta_\eta - T\omega + \varphi_+) < -\delta$. Here, φ_+ is a phase independent of the parameters.
- (ii) There is a real number κ_+ and, for each $r \leq \bar{r}/2$, a smaller radius $\epsilon(r)$ proportional to r , such that, if $|P(u^*(z) - u_+, v^*(z) - v_+)^T| = r$, $|(\eta_1 - u^*(z), \eta_2 - v^*(z))| < \epsilon$, $|(\zeta_1 - u^*(z), \zeta_2 - v^*(z))| < \epsilon$ and $n \geq 1/\epsilon$ with n an integer, then

$$E_{\eta_1, \eta_2, \zeta_1, \zeta_2}(\kappa_+ - 2z + 2\pi n/\omega - \nu) > 0,$$

$$E_{\eta_1, \eta_2, \zeta_1, \zeta_2}(\kappa_+ - 2z + 2\pi n/\omega + \nu) < 0.$$

Similar assertions hold with the same \bar{r} and $\epsilon(r)$ when replacing (u_+, v_+) by (u_-, v_-) , (u^*, v^*) by (u_*, v_*) and φ_+, κ_+ by possibly different phases φ_-, κ_- .

We refer to [6] where a very similar statement is proved; there the existence and local uniqueness of (\bar{U}, \bar{V}) follow from Lemma A.3. The sign property (i) of the energy is a consequence of Lemma A.2, and see Lemma 3.1 for the detail. Finally, (ii) is a consequence of (i). Note that the formula for the sign of E in (i) or (ii) is correct for a suitable choice of P , for other choices the sign in front of $T\omega$ has to be changed.

3 Isolated critical points

The aim of this section is to show that all critical points of J are isolated, up to translations in x . As already mentioned in the introduction, for the construction of multi-front solutions we would just need to know that the basic heteroclinic u^* is isolated up to translations, as in [6]. Showing that all critical points are isolated is considerably more difficult, however this stronger property will be needed for working out the stability analysis in the last part of this paper.

Proposition 3.1 *Any critical point u_c of J is an isolated critical point for the $H^1(\mathbb{R})$ metric, up to translations in x .*

Proposition 3.1 will follow from two facts. The first one is an alternative. Its proof is delicate and relies crucially on the results of Section 2 and the real-analyticity of the Hamiltonian:

Lemma 3.2 *Either the stable and unstable manifolds of the hyperbolic points $z_- = (u_-, v_-, 0, 0)$ and $z_+ = (u_+, v_+, 0, 0)$ are bounded, or every critical point of J is isolated up to translation in x .*

Note that the four manifolds mentioned in Lemma 3.2 are all bounded (resp. unbounded) if at least one of them is bounded (resp. unbounded), since the Lagrangian $L(u, v, u', v')$ is invariant under time reversal and under the symmetry of center $(u_+/2, v_+/2)$.

The second fact is stated in the next proposition, in which we find trajectories which are either on the unstable manifold of $(u_-, v_-, 0, 0)$ or on the unstable manifold of $(u_+, v_+, 0, 0)$, and that reach a point of arbitrarily large size. This shows that the stable and unstable manifolds mentioned in Lemma 3.2 are unbounded, and thus Proposition 3.1 is established:

Proposition 3.3 *For any $b \in \mathbb{R}$, there is a solution (u, v) of (1.3)-(1.4) which satisfies $u(0) = \frac{2}{\gamma}v(0) = b$ and one of the following conditions:*

$$(i) \lim_{x \rightarrow -\infty} (u(x), v(x)) = (u_-, v_-);$$

$$(ii) \lim_{x \rightarrow -\infty} (u(x), v(x)) = (u_+, v_+).$$

Proof. We use an argument similar to the proof of Theorem 1.1 of [19]. For $s \in \{+, -\}$, let \hat{v}_s be a C^∞ -function on $(-\infty, 0]$ such that

$$\hat{v}_s(0) = \frac{2b}{\gamma}, -\hat{v}_s''(0) + \gamma\hat{v}_s(0) = b \text{ and } \hat{v}_s(x) = v_s \text{ if } x \leq -1. \quad (3.1)$$

For convenience in notation, we define $\mathbf{H}_w = w + H_0^1(-\infty, 0)$. Set $\hat{u}_s = \gamma\hat{v}_s - \hat{v}_s''$. For $w \in \mathbf{H}_{\hat{u}_s}$, let

$$\hat{J}(w) = \int_{-\infty}^0 \left[\frac{d}{2}(w')^2 + \frac{1}{2}w\hat{\mathcal{L}}w - \int_0^w f(\xi)d\xi \right] dx, \quad (3.2)$$

where $\hat{\mathcal{L}}w$ denotes the unique solution of

$$-v'' + \gamma v = w, \quad v \in \mathbf{H}_{\hat{v}_s}. \quad (3.3)$$

The formula of Proposition 2.4 has an easy extension to the present situation, which can be proved in exactly the same way: if $w \in \mathbf{H}_{\hat{u}_\varepsilon}$ then

$$\begin{aligned} \hat{J}(w) &= \int_{-\infty}^0 \frac{1}{2} \left[\left(d - \frac{1}{\gamma^2} \right) (w')^2 + \left((\hat{\mathcal{L}}w)' - \frac{w'}{\gamma} \right)^2 + \gamma \left(\hat{\mathcal{L}}w - \frac{w}{\gamma} \right)^2 \right] dx \\ &\quad + \int_{-\infty}^0 \frac{1}{4} (w - u_-)^2 (w - u_+)^2 dx. \end{aligned} \quad (3.4)$$

Set $\hat{c}_s = \inf_{w \in \mathbf{H}_{\hat{u}_s}} \hat{J}(w)$. We distinguish two cases:

Case 1: $\hat{c}_- \leq \hat{c}_+$. Pick a sequence $\{u_m\} \subset \mathbf{H}_{\hat{u}_-}$ such that $\lim_{m \rightarrow \infty} \hat{J}(u_m) = \hat{c}_-$. With a priori bounds on u_m derived from (3.4) and by passing to a subsequence if necessary, we may assert, as in the proof of Theorem 1.1 of [19], that $\{u_m\}$ converges to a function u in $H_{loc}^1(-\infty, 0]$. Moreover $\hat{\mathcal{L}}u_m$ has a limit v in H_{loc}^1 , and $u(0) = \frac{\gamma}{2}v(0) = b$. Invoking (3.4) again, we then prove that $u \in \mathbf{H}_{\hat{u}_s}$ for some $s \in \{+, -\}$, and $v = \hat{\mathcal{L}}u$, $\hat{J}(u) \leq c_-$. Now, since $u \in \mathbf{H}_{\hat{u}_s}$ and $\hat{c}_- \leq \hat{c}_+$, we must have $\hat{J}(u) \geq c_s \geq c_-$. Hence $\hat{J}(u) = c_-$, $\{u_m\}$ strongly converges to u in $H^1(-\infty, 0)$ and u is a minimizer of \hat{J} in $\mathbf{H}_{\hat{u}_-}$. Thus (u, v) satisfies (1.3)-(1.4) and (i) holds.

Case 2: $\hat{c}_- > \hat{c}_+$. The argument is similar, taking a sequence $\{u_m\} \subset \mathbf{H}_{\hat{u}_+}$ such that $\lim_{m \rightarrow \infty} \hat{J}(u_m) = \hat{c}_+$. As in case 1, after extraction, $\{u_m\}$ strongly converges to a minimizer u of \hat{J} in $\mathbf{H}_{\hat{u}_+}$ and (ii) holds. Now the proof of Proposition 3.3 is complete. \blacksquare

Proof of Lemma 3.2 and Proposition 3.1. We actually just need to prove Lemma 3.2, since Proposition 3.1 then immediately follows.

Let u_c be a non-constant critical point of J and $\lim_{z \rightarrow -\infty} u_c(x) = u_{s_-}$, $\lim_{z \rightarrow +\infty} u_c(x) = u_{s_+}$, $s_-, s_+ \in \{-, +\}$. Note that $z_{s_-} = (u_{s_-}, v_{s_-}, 0, 0)$ is a hyperbolic equilibrium of the first order Hamiltonian system (HS) with Hamiltonian function $H(u, v, p, q)$, which is associated to the second order Lagrangian system (1.3)-(1.4). Here, p denotes the momentum conjugate to u and q the momentum conjugate to v . With the Hamiltonian H being real analytic, the local unstable manifold $W_{loc}^u(z_{s_-})$ of z_{s_-} is a real analytic submanifold of $\mathbb{R}^2 \times \mathbb{R}^2$. Moreover the unstable space of the linearization of (HS) at z_{s_-} is the graph of a linear map from \mathbb{R}^2 to itself, hence $W_{loc}^u(z_{s_-})$ is the graph of a real analytic map φ from a small neighborhood \mathcal{U}_{s_-} of (u_{s_-}, v_{s_-}) into \mathbb{R}^2 , and $\varphi(u_{s_-}, v_{s_-}) = (0, 0)$. Similarly, $z_{s_+} = (u_{s_+}, v_{s_+}, 0, 0)$ is hyperbolic and its local stable manifold $W_{loc}^s(z_{s_+})$ is the graph of a real analytic map $\psi = (\psi_p, \psi_q)$ from a small neighborhood \mathcal{U}_{s_+} of (u_{s_+}, v_{s_+}) into \mathbb{R}^2 , and $\psi(u_{s_+}, v_{s_+}) = (0, 0)$.

Recall that we use the same ρ_1 as introduced in Proposition 2.7. Also, in Lemma 2.9, we choose $\delta = 1/2$ with the associated $\bar{r} > 0$. To employ Proposition 2.7, we pick $0 < \rho \leq \min\{\rho_1/2, \bar{r}\}$ such that $|P(u - u_{s_\pm}, v - v_{s_\pm})| \leq \rho$

implies $(u, v) \in \mathcal{U}_{s_{\pm}}$. Given $\theta \in \mathbb{R}$, we denote $(\mathbf{u}, \mathbf{v})_{\theta}^{\text{T}} := P^{-1}(\rho \cos \theta, \rho \sin \theta)^{\text{T}} + (u_{s_{-}}, v_{s_{-}})^{\text{T}}$ and $\mathbf{z}_{\theta} := ((\mathbf{u}, \mathbf{v})_{\theta}, \varphi((\mathbf{u}, \mathbf{v})_{\theta}))$.

Let $v_c = \mathcal{L}u_c$. Remembering that $\rho_1 < \rho_0$, we may take x_1, x_2 associated to u_c and ρ as in Proposition 2.6(ii). Let $\Phi = (U, V, \Pi_U, \Pi_V)$ be the flow of (HS) at time $T = x_2 - x_1 + 1$ (both T and Φ depend on the critical point u_c under consideration). Then we may write $P(u_c(x_1) - u_{s_{-}}, v_c(x_1) - v_{s_{-}})^{\text{T}} = (\rho \cos \theta_c, \rho \sin \theta_c)^{\text{T}}$ for some angle θ_c , and $P(U(\mathbf{z}_{\theta_c}), V(\mathbf{z}_{\theta_c}))^{\text{T}}$ lies in the open disk D_+ of center $P(u_{s_{+}}, v_{s_{+}})^{\text{T}}$ with radius ρ .

From now on in this proof, *let us assume that u_c is not isolated up to translations*, for the H^1 metric in the set of critical points of J . Under this assumption, we are going to show that at least one of the two unstable manifolds $W^u(z_{-}), W^u(z_{+})$ is bounded. By time reversibility and symmetry with respect to $(u_{+}/2, v_{+}/2)$, this will automatically imply that both unstable manifolds are bounded, as well as the two stable manifolds $W^s(z_{\pm})$.

Let u be a critical point of J sufficiently close to u_c in the $H^1(\mathbb{R})$ topology but not equal to a translate of u_c . Then the trajectory parametrized by $(u, \mathcal{L}u, du', -\mathcal{L}u')$ in the phase space must contain a point \mathbf{z}_{θ} with θ arbitrarily close, but not equal, to the angle θ_c . Moreover, since $P(U(\mathbf{z}_{\theta_c}), V(\mathbf{z}_{\theta_c}))^{\text{T}}$ lies in the open disk D_+ of center $P(u_{s_{+}}, v_{s_{+}})^{\text{T}}$ with radius ρ , the same is true for $P(U(\mathbf{z}_{\theta}), V(\mathbf{z}_{\theta}))^{\text{T}}$ if u is close enough to u_c in $H^1(\mathbb{R})$ topology. So θ is a zero of each of the functions $\chi_1(\theta) = \Pi_U(\mathbf{z}_{\theta}) - \psi_p(U(\mathbf{z}_{\theta}), V(\mathbf{z}_{\theta}))$ and $\chi_2(\theta) = \Pi_V(\mathbf{z}_{\theta}) - \psi_q(U(\mathbf{z}_{\theta}), V(\mathbf{z}_{\theta}))$.

The above argument shows that θ_c is not an isolated zero of the real-analytic functions χ_1 and χ_2 , which are defined in a small open interval \mathcal{I} containing θ_c . So these functions are identically zero on \mathcal{I} , which means that the flow Φ sends all the points \mathbf{z}_{θ} near \mathbf{z}_{θ_c} to points of $W_{loc}^s(z_{s_{+}})$.

To each θ in \mathcal{I} , we associate the solution Z_{θ} of the Hamiltonian system with initial value $Z_{\theta}(0) = \mathbf{z}_{\theta}$. Let $u_{\theta}(x)$ be the first component of the vector $Z_{\theta}(x)$ and let $j(\theta) := J(u_{\theta})$. Then it is not hard to see that the function $\theta \in \mathcal{I} \rightarrow u_{\theta}$ is continuous, and even of class C^1 , for the $H^1(\mathbb{R})$ metric on the target space. Thus, by the chain rule, $\frac{d}{d\theta}j(\theta) = 0$, and $j(\theta) = J(\theta_c)$. Moreover, all the angles in \mathcal{I} correspond to critical points u_{θ} that are not isolated, up to translation, in $H^1(\mathbb{R})$ topology.

Denote by Θ the set of angles θ associated to all non-isolated critical points J which converge to $u_{s_{-}}$ as $x \rightarrow -\infty$. We have shown that Θ is an open subset of \mathbb{R} , and by assumption it contains θ_c . Clearly Θ is 2π -periodic, and the above argument shows that j is a 2π -periodic and locally constant function on Θ . Let \mathcal{J} be the maximal open interval in Θ containing θ_c and let $\hat{\theta} := \sup \mathcal{J}$. If $\hat{\theta} = +\infty$, then $\mathcal{J} = \Theta = \mathbb{R}$ and j is constant on \mathbb{R} , so Corollary 2.5 gives a uniform L^{∞} estimate on *all* the solutions Z_{θ} , so $W^u(z_{s_{-}}) = \bigcup_{(\theta, x) \in [0, 2\pi) \times \mathbb{R}} Z_{\theta}(x)$ is bounded. Thus, to complete the proof of Lemma 3.2, we just need to study

the remaining case when $\hat{\theta}$ is a finite number.

Then, as $\theta \rightarrow \hat{\theta}$, $\theta \in \mathcal{J}$, Z_θ converges in the C_{loc}^1 topology to $Z_{\hat{\theta}}$, by continuous dependence of the solutions of the Hamiltonian system with respect to initial data. Since $\hat{\theta} \notin \mathcal{J}$, the convergence of u_θ to $u_{\hat{\theta}}$ does not hold for the $H^1(\mathbb{R})$ metric. As a consequence, from Proposition 2.7(iii), $u_{\hat{\theta}}$ is a critical point of j satisfying the estimate $\ell \leq j(u_{\hat{\theta}}) \leq j(\theta_c) - \ell$. From now on, set $(u_{s_1}, v_{s_1}) = \lim_{x \rightarrow +\infty} (u_{\hat{\theta}}, \mathcal{L}u_{\hat{\theta}})(x)$, $s_1 \in \{-, +\}$. Let $\underline{t} \in \mathbb{R}$ be such that for all $x > \underline{t}$, $P(u_{\hat{\theta}}, \mathcal{L}u_{\hat{\theta}})^{\text{T}}(x)$ lies in the open disk D_1 of center $P(u_{s_1}, v_{s_1})^{\text{T}}$ with radius ρ , while $P(u_{\hat{\theta}}, \mathcal{L}u_{\hat{\theta}})^{\text{T}}(\underline{t})$ sits on the boundary of D_1 . Then, for θ close enough to $\hat{\theta}$, there are two numbers $\underline{t}_\theta < \bar{t}_\theta$ such that for all $x \in (\underline{t}_\theta, \bar{t}_\theta)$, $P(u_\theta, \mathcal{L}u_\theta)^{\text{T}}(x)$ lies in the open disk D_1 , while $P(u_\theta, \mathcal{L}u_\theta)^{\text{T}}(\underline{t}_\theta)$ and $P(u_\theta, \mathcal{L}u_\theta)^{\text{T}}(\bar{t}_\theta)$ both sit on the boundary of D_1 . Moreover $\lim_{\theta \rightarrow \hat{\theta}} \underline{t}_\theta = \underline{t}$ and $\lim_{\theta \rightarrow \hat{\theta}} \bar{t}_\theta = +\infty$. The existence of \underline{t}_θ is due to the convergence of Z_θ in C_{loc}^1 , and the existence of \bar{t}_θ is due to the lack of convergence in the $H^1(\mathbb{R})$ metric.

Now, it follows from the Hartman-Grobman theorem that the distance between $Z_\theta(\underline{t}_\theta)$ and $W_{\text{loc}}^s(z_{s_1})$ tends to zero as $\theta \rightarrow \hat{\theta}$, since $(\bar{t}_\theta - \underline{t}_\theta) \rightarrow \infty$. As a consequence, remembering that (2.5) holds on $W_{\text{loc}}^s(z_{s_1})$, we obtain a similar estimate when θ is sufficiently close to $\hat{\theta}$:

$$\frac{d}{dx} |P(u_\theta(\underline{t}_\theta) - u_{s_1}, \mathcal{L}u_\theta(\underline{t}_\theta) - v_{s_1})^{\text{T}}| \leq -\frac{\lambda}{4} |P(u_\theta(\underline{t}_\theta) - u_{s_1}, \mathcal{L}u_\theta(\underline{t}_\theta) - v_{s_1})^{\text{T}}|.$$

Such an inequality, together with the implicit function theorem gives the continuity, and even the real-analyticity, of \underline{t}_θ as a function of θ on a small interval $(\hat{\theta} - \varepsilon, \hat{\theta}) \subset \mathcal{J}$. Similarly, the distance between $Z_\theta(\bar{t}_\theta)$ and $W_{\text{loc}}^u(z_{s_1})$ tends to zero as $\theta \rightarrow \hat{\theta}$, and the continuity of \bar{t}_θ can be proved using an inequality analogous to (2.4).

So we can define two continuous functions $\underline{\alpha}, \bar{\alpha} : (\hat{\theta} - \varepsilon, \hat{\theta}) \rightarrow \mathbb{R}$ which satisfy

$$\begin{aligned} P(u_\theta - u_{s_1}, \mathcal{L}u_\theta - v_{s_1})^{\text{T}}(\underline{t}_\theta) &= \rho (\cos \underline{\alpha}(\theta), \sin \underline{\alpha}(\theta))^{\text{T}}, \\ P(u_\theta - u_{s_1}, \mathcal{L}u_\theta - v_{s_1})^{\text{T}}(\bar{t}_\theta) &= \rho (\cos \bar{\alpha}(\theta), \sin \bar{\alpha}(\theta))^{\text{T}}. \end{aligned}$$

Since $\lim_{\theta \rightarrow \hat{\theta}} (u_\theta, \mathcal{L}u_\theta)(\underline{t}_\theta) = (u_{\hat{\theta}}, \mathcal{L}u_{\hat{\theta}})(\underline{t})$, $\underline{\alpha}(\theta)$ has a finite limit $\underline{\alpha}(\hat{\theta})$ as $\theta \rightarrow \hat{\theta}$. In order to study the limit of $\bar{\alpha}$, we are going to use Lemma 2.9, remembering that we have fixed $\delta = 1/2$ and chosen $\rho \leq \bar{r}$. Since u_θ is a critical point of J , its energy is zero, so $E_{u_\theta(\underline{t}_\theta), \mathcal{L}u_\theta(\underline{t}_\theta), u_\theta(\bar{t}_\theta), \mathcal{L}u_\theta(\bar{t}_\theta)}(\bar{t}_\theta - \underline{t}_\theta) = 0$. Then, from (i) in Lemma 2.9, we find a phase φ_1 independent of the parameters, such that, for θ close enough to $\hat{\theta}$,

$$\cos(\bar{\alpha}(\theta) - \underline{\alpha}(\theta) - (\bar{t}_\theta - \underline{t}_\theta)\omega + \varphi_1) \in [-1/2, 1/2].$$

But $(\bar{t}_\theta - \underline{t}_\theta) \rightarrow \infty$ and $\underline{\alpha}(\theta) \rightarrow \underline{\alpha}(\hat{\theta})$ as $\theta \rightarrow \hat{\theta}$. Moreover $\bar{\alpha}$ depends continuously on θ . So we must have $\lim_{\theta \rightarrow \hat{\theta}} \bar{\alpha}(\theta) = \infty$. As a consequence, given any $\alpha \in [0, 2\pi)$, there is a sequence $\{\theta_\alpha^{(n)}\}$ such that $\bar{\alpha}(\theta_\alpha^{(n)}) = \alpha + 2n\pi$ for all

n large enough, and $\theta_\alpha^{(n)} \rightarrow \hat{\theta}$ as $n \rightarrow \infty$. Passing to a subsequence if necessary, it follows from Proposition 2.7 that $\{u_{\theta_\alpha^{(n)}}(\cdot + \bar{t}_{\theta_\alpha^{(n)}})\}$ converges for the local H^1 topology to a non-constant critical point of J , denoted by $u_\alpha^{(2)}$, with $P(u_\alpha^{(2)} - u_{s_1}, \mathcal{L}u_\alpha^{(2)} - v_{s_1})^\top(0) = \rho(\cos \alpha, \sin \alpha)^\top$ and $\ell \leq J(u_\alpha^{(2)}) \leq J(u_c) - \ell$. Then Corollary 2.5 gives an L^∞ estimate on $Z_\alpha^{(2)} = (u_\alpha^{(2)}, \mathcal{L}u_\alpha^{(2)}, d(u_\alpha^{(2)})', -\mathcal{L}(u_\alpha^{(2)})')$, which is independent of α . Moreover, since the distance between $Z_{\theta_\alpha^{(n)}}(\bar{t}_{\theta_\alpha^{(n)}})$ and $W_{\text{loc}}^u(z_{s_1})$ tends to zero as $n \rightarrow \infty$, $Z_\alpha^{(2)}(0)$ belongs to $W_{\text{loc}}^u(z_{s_1})$. But $W_{\text{loc}}^u(z_{s_1})$ is a graph in $\mathbb{R}^2 \times \mathbb{R}^2$, so we conclude that $W^u(z_{s_1}) = \bigcup_{(\alpha, x) \in [0, 2\pi) \times \mathbb{R}} Z_\alpha^{(2)}(x)$. As a consequence, this unstable manifold is bounded. This completes the proof of Lemma 3.2. So Proposition 3.1 is true. \blacksquare

Recall that the Lagrangian $L(u, v, u', v')$ is autonomous and u_* is a non-constant solution. By taking a small translation in x if necessary, from now on we always assume that $u'_*(0) \neq 0$. This condition also holds for u^* , since it is the reverse orbit of u_* .

As a consequence of Proposition 3.1, we obtain the following result.

Corollary 3.4 *There exist $h_0, \sigma_0 > 0$ and, for any $0 < h < h_0$, a radius $\bar{\sigma}(h) > 0$ with $\lim_{h \rightarrow 0} \bar{\sigma}(h) = 0$, such that the local sublevel set*

$$\mathcal{V}_h = \{u \in H_{u^*} : u(0) = u^*(0), \|u - u^*\|_{H^1(\mathbb{R})} \leq \sigma_0 \text{ and } J(u) \leq J(u^*) + h\}$$

satisfies the following property:

$$u \in \mathcal{V}_h \Rightarrow \|u - u^*\|_{H^1(\mathbb{R})} < \bar{\sigma}(h).$$

Proof. From Corollary 2.8, for σ_1 small enough, the functional J satisfies the Palais-Smale condition on the closed ball of center u_* with radius σ_1 of $H^1(\mathbb{R})$ -norm. Since $u'_*(0) \neq 0$, by Proposition 3.1 there exists $\sigma_0 \leq \sigma_1$ such that u_* is the unique critical point of J on the closed ball of center u_* with radius σ_0 , so does u^* . Consequently in this closed ball, Theorem 1.1 tells that u^* is the unique minimizer of J , which completes the proof. \blacksquare

Consider a sufficiently large number z and define

$$\mathcal{V}_{h,z} := \{u \in H^1(-z, z) : u \equiv \hat{u} \text{ on } [-z, z] \text{ for some } \hat{u} \in \mathcal{V}_h\}. \quad (3.5)$$

Now, for $u \in \mathcal{V}_{h,z}$ with h small and z large, the functional J is C^2 and strictly convex on

$$\mathcal{C}_u := \{\tilde{u} \in H_{u^*} : \tilde{u} \equiv u \text{ on } [-z, z] \text{ and } \|\tilde{u} - u^*\|_{H^1(\mathbb{R})} \leq \bar{\sigma}(h)\},$$

which is a closed, bounded and convex subset of H_{u^*} . Indeed, if $\tilde{u} \in \mathcal{C}_u$, any other element of \mathcal{C}_u near \tilde{u} is of the form $\tilde{u} + w$ with $w \equiv 0$ on $[-z, z]$, while

$\|(\tilde{u} - u_-)(\tilde{u} - u_+)\|_{L^\infty(\mathbb{R} \setminus [-z, z])}$ is small. Thus a direct calculation gives, for some $\hat{k} > 0$:

$$D^2 J(\tilde{u}) \cdot w \cdot w = \int_{\mathbb{R}} \{d(w')^2 - f'(\tilde{u})w^2 + (\mathcal{L}w')^2 + \gamma(\mathcal{L}w)^2\} \geq \hat{k}\|w\|_{H^1(\mathbb{R})}^2$$

Moreover, if $\tilde{u} \in \mathcal{C}_u$ satisfies $\|\tilde{u} - u^*\|_{H^1(\mathbb{R})} = \bar{\sigma}(h)$, then $J(\tilde{u}) > J(u^*) + h \geq \min_{\mathcal{C}_u} J$. So the minimizer $b(u)$ of J on \mathcal{C}_u does not saturate the constraint $\|\tilde{u} - u^*\|_{H^1(\mathbb{R})} \leq \bar{\sigma}(h)$. As a consequence, $(b(u), \mathcal{L}b(u))$ solves the system (1.3)-(1.4) outside the interval $[-z, z]$, and by the implicit function theorem the map $b : \mathcal{V}_{h,z} \rightarrow H_{u^*}$ constructed in this way is smooth. This provides a Lyapunov-Schmidt reduction $J_z = J \circ b$ of J defined on $\mathcal{V}_{h,z}$, and the following corollary holds.

Corollary 3.5 *For h_0 small enough, there is $z_0 > 0$ such that if $h \in (0, h_0)$ and $z > z_0$, then*

$$\rho(h) := \inf\{\|J'_z(u)\|_{(H^1(-z,z))^*} : u \in \mathcal{V}_{h,z} \text{ and } J_z(u) = J(u^*) + h\} > 0.$$

Proof. Suppose that the assertion of the corollary is false. Then $\rho(h) = 0$ for small h and Corollary 2.8 implies that a Palais-Smale sequence converges to a critical point of J in a small ball of center u^* at the critical level $J(u^*) + h$. Hence there would exist critical points of J in any small neighborhood of u^* . This is contrary to Proposition 3.1. \blacksquare

4 Construction of multi-front waves

With $u'_*(0) \neq 0$, we now get into details about how to construct the multi-front solutions. Let $h > 0$ be small and $\Delta > 0$ large (to be determined later as depending on h). Pick an arbitrary integer $N \geq 1$ and an arbitrary finite sequence of positive integers $\mathbf{n} = (n_i)_{1 \leq i \leq N}$ such that $n_i \geq \Delta$ for all i . Take $z > 0$ large enough so that $r_- := |P(u^*(-z) - u_-, v^*(-z) - v_-)^T| \leq \bar{r}/2$ and $r_+ := |P(u^*(z) - u_+, v^*(z) - v_+)^T| \leq \bar{r}/2$, where \bar{r} is the small radius considered in Lemma 2.9.

Recall $\mathcal{V}_{h,z}$ from (3.5) and define as follows a smooth map $b_{\mathbf{n}}$ from $(\mathcal{V}_{h,z})^{N+1} \times [-\nu, \nu]^N$ into H_w , with $w = 0$ for N odd, $w = u^*$ for N even. To any $(\mathbf{u}, \mathbf{x}) = ((u_i)_{0 \leq i \leq N}, (x_i)_{1 \leq i \leq N})$ in $(\mathcal{V}_{h,z})^{N+1} \times [-\nu, \nu]^N$, we associate a unique function

$u = b_{\mathbf{n}}(\mathbf{u}, \mathbf{x}) \in H_w$ such that:

- (S₁) $\forall i \in [0, N] \cap 2\mathbb{Z}$, $u \equiv u_i(\cdot - C_i)$ on $(C_i - z, C_i + z)$,
- (S₂) $\forall i \in [0, N] \cap (2\mathbb{Z} + 1)$, $u \equiv u_i(C_i - \cdot)$ on $(C_i - z, C_i + z)$,
- (S₃) $\|u - u_-\|_{H^1(-\infty, -z)} \leq Kr_-$,
- (S₄) $\forall i \in [0, N - 1] \cap 2\mathbb{Z}$, $\|u - u_+\|_{H^1(C_i+z, C_{i+1}-z)} \leq Kr_+$,
- (S₅) $\forall i \in [0, N - 1] \cap (2\mathbb{Z} + 1)$, $\|u - u_-\|_{H^1(C_i+z, C_{i+1}-z)} \leq Kr_-$,
- (S₆) $\|u - u_{\pm}\|_{H^1(C_N+z, \infty)} \leq Kr_{\pm}$, where $\pm = +$ for N even, $\pm = -$ for N odd,
- (S₇) $C_0 = 0$, $C_i = C_{i-1} + X_i$ ($1 \leq i \leq N$),
- (S₈) $X_{2j} = x_{2j} + \kappa_- + \frac{2\pi n_{2j}}{\omega}$,
- (S₉) $X_{2j+1} = x_{2j+1} + \kappa_+ + \frac{2\pi n_{2j+1}}{\omega}$,
- (S₁₀) $(u, \mathcal{L}u)$ satisfies (1.3)-(1.4) on each of the intervals $(-\infty, -z]$, $[C_i + z, C_{i+1} - z]$, $[C_N + z, +\infty)$.

Taking a large K independent of r_{\pm} , we claim that for h small and z large enough, conditions (S₁)-(S₁₀) determine u in a unique way, and explain why the corresponding function $b_{\mathbf{n}}(\mathbf{u}, \mathbf{x})$ is smooth. Observe that one can define the set $\mathcal{U}_{(\mathbf{u}, \mathbf{x})}$ consisting of all functions u satisfying conditions (S₁)-(S₆). This set is convex, bounded, closed in the H^1 topology. Moreover, the controls (S₃)-(S₆) on u imply the strict convexity of J restricted to $\mathcal{U}_{(\mathbf{u}, \mathbf{x})}$. Indeed, if $u \in \mathcal{U}_{(\mathbf{u}, \mathbf{x})}$, any other element of $\mathcal{U}_{(\mathbf{u}, \mathbf{x})}$ is of the form $u + w$ with $w \equiv 0$ on $A := \bigcup_{1 \leq i \leq N} [C_i - z, C_i + z]$, while $\|(u - u_-)(u - u_+)\|_{L^\infty(\mathbb{R} \setminus A)}$ is small for Kr_{\pm} small. Then a direct calculation gives

$$D^2 J(u) \cdot w \cdot w = \int_{\mathbb{R}} \{d(w')^2 - f'(u)w^2 + (\mathcal{L}w')^2 + \gamma(\mathcal{L}w)^2\} \geq \bar{k} \|w\|_{H^1(\mathbb{R})}^2$$

for some $\bar{k} > 0$, exactly as in the proof of Corollary 3.4.

So J has a unique minimizer in $\mathcal{U}_{(\mathbf{u}, \mathbf{x})}$. Moreover for K large enough, if a function u belongs to $\mathcal{U}_{(\mathbf{u}, \mathbf{x})}$, and saturates at least one of the constraints (S₃)-(S₆) then $J(u) > \min_{\mathcal{U}_{(\mathbf{u}, \mathbf{x})}} J$. In conclusion, the minimizer does not saturate any of the constraints, so it is the only solution of (S₁₀) in $\mathcal{U}_{(\mathbf{u}, \mathbf{x})}$ and the implicit function theorem gives a smooth function $b_{\mathbf{n}}$ of (\mathbf{u}, \mathbf{x}) in the H^1 topology.

Up to this stage, a Lyapunov-Schmidt reduction has been performed, and the next task is to minimize the reduced functional $\mathbf{J} = J \circ b_{\mathbf{n}}$. Note that the set \mathcal{V}_h is a bounded, closed sublevel set of the weakly lower semicontinuous functional J , thus $\mathcal{V}_{h,z}$ is weakly compact in $H^1(-z, z)$. Moreover one easily checks that \mathbf{J} is weakly lower semicontinuous on $(\mathcal{V}_{h,z})^{N+1} \times [-\nu, \nu]^N$, which leads to the existence of a minimizer $(\bar{\mathbf{u}}, \bar{\mathbf{x}})$ of \mathbf{J} on that set.

Lemma 4.1 *Given z large, h small and choosing Δ large enough independently of N , if $n_i \geq \Delta$ for every $1 \leq i \leq N$, then $\hat{u}_{\mathbf{n}} := b_{\mathbf{n}}(\bar{\mathbf{u}}, \bar{\mathbf{x}})$ is a local minimizer of J .*

To prove Lemma 4.1, we introduce the set $\mathcal{O} = \bigcup_{(\mathbf{u}, \mathbf{x}) \in (\mathcal{V}_{h,z})^{N+1} \times [-\nu, \nu]^N} \mathcal{U}_{(\mathbf{u}, \mathbf{x})}$ consisting of functions u satisfying (S_1) - (S_6) for some $(\mathbf{u}, \mathbf{x}) \in (\mathcal{V}_{h,z})^{N+1} \times [-\nu, \nu]^N$. The next lemma shows that \mathcal{O} contains a small ball for the $H^1(\mathbb{R})$ metric, with center at $\hat{u}_{\mathbf{n}}$. Clearly $\hat{u}_{\mathbf{n}}$ minimizes J on \mathcal{O} , by virtue of the construction used in the variational argument, and thus Lemma 4.1 is an immediate consequence.

Lemma 4.2 *Choose z large and h small, then Δ large enough, independently of N . Assume that $n_i \geq \Delta$ for all $1 \leq i \leq N$. If $(\bar{\mathbf{u}}, \bar{\mathbf{x}}) = ((\bar{u}_i)_{0 \leq i \leq N}, (\bar{x}_i)_{1 \leq i \leq N})$ is a minimizer of \mathbf{J} in $(\mathcal{V}_{h,z})^{N+1} \times [-\nu, \nu]^N$ then*

- (i) $J_z(\bar{u}_i) < J(u^*) + h$ for all $0 \leq i \leq N$,
- (ii) $-\nu < \bar{x}_i < \nu$ for all $1 \leq i \leq N$.

The following lemma will be used to prove Lemma 4.2.

Lemma 4.3 *Let z and h be small enough, both being independent of N . For any $\alpha > 0$ there exists $\bar{\Delta}(\alpha)$, not depending on N , such that if $n_i \geq \bar{\Delta}(\alpha) \forall 1 \leq i \leq N$ then*

$$\|J'_z(u_i) - \partial_{u_i} \mathbf{J}\|_{(H^1(-z,z))^*} < \alpha, \quad \forall (\mathbf{u}, \mathbf{x}) \in (\mathcal{V}_{h,z})^{N+1} \times [-\nu, \nu]^N, \quad 0 \leq i \leq N.$$

The proof of Lemma 4.3 is standard (see e.g. [6]). We omit it.

Proof of Lemma 4.2. We argue indirectly. Suppose that $J_z(\bar{u}_l) = J(u^*) + h$ for some $l \in (0, N)$, applying Lemma 4.3 yields

$$\langle \partial_{u_l} \mathbf{J}(\bar{\mathbf{u}}, \bar{\mathbf{x}}), \nabla_{H^1(-z,z)} J_z(\bar{u}_l) \rangle \geq \frac{\rho(h)}{2},$$

with $\rho(h)$ given by Corollary 3.5. Then moving u_l slightly in the direction of $-\nabla_{H^1(-z,z)} J_z(\bar{u}_l)$ would decrease $\mathbf{J}(\mathbf{u}, \mathbf{x})$, which contradicts the minimality of $\mathbf{J}(\bar{\mathbf{u}}, \bar{\mathbf{x}})$. The proof of (i) is complete.

We next apply Lemma 2.9 to prove (ii). Remembering that we have fixed z so that

$$r_{\pm} := |P(u^*(\pm z), v^*(\pm z))^T - P(u_{\pm}, v_{\pm})^T| \leq \bar{r}/2,$$

taking h small enough and Δ large, we may impose

$$|(\bar{u}_i(\pm z), \mathcal{L}b_{\mathbf{n}}(\bar{\mathbf{u}}, \bar{\mathbf{x}})(C_i \pm (-1)^i z)) - (u^*(\pm z), v^*(\pm z))| < \epsilon(r_{\pm})$$

with $\epsilon(r_{\pm})$ as in Lemma 2.9 and the condition $n_i \geq 1/\min\{\epsilon(r_-), \epsilon(r_+)\}$ being imposed. Suppose $\bar{x}_l = -\nu$ for some $l \in (1, N-1)$, it follows from Lemma 2.9 that

$$\partial_{x_l} \mathbf{J}(\bar{\mathbf{u}}, \bar{\mathbf{x}}) = -E_{\eta_1, \eta_2, \zeta_1, \zeta_2}(\kappa_s - 2z + 2\pi n_l/\omega - \nu) < 0,$$

where $s = +$ if l is odd, $s = -$ if l is even, $(\eta_1, \eta_2) = (b_{\mathbf{n}}(\bar{\mathbf{u}}, \bar{\mathbf{x}}), \mathcal{L}b_{\mathbf{n}}(\bar{\mathbf{u}}, \bar{\mathbf{x}}))(C_{l-1} + z)$ and $(\zeta_1, \zeta_2) = (b_{\mathbf{n}}(\bar{\mathbf{u}}, \bar{\mathbf{x}}), \mathcal{L}b_{\mathbf{n}}(\bar{\mathbf{u}}, \bar{\mathbf{x}}))(C_l - z)$. Then increasing x_l slightly would decrease \mathbf{J} , which again contradicts the minimality of $\mathbf{J}(\bar{\mathbf{u}}, \bar{\mathbf{x}})$. Likewise, if $\bar{x}_l = \nu$ we could decrease \mathbf{J} by slightly decreasing x_l . Now the proof of Lemma 4.2 is complete. \blacksquare

We are now ready to prove the existence result of multi-front solutions, stated in Theorem 1.2. The stability of such solutions will be investigated in the next section.

Proof of Theorem 1.2. Impose that z be large enough so that $Kr_{\pm} < \tilde{\sigma}$, where $\tilde{\sigma}$ is to be chosen later. Then pick z large and h small enough so that the small number $\bar{\sigma}(h)$, defined in Corollary 3.4, is less than $\tilde{\sigma}$, and so that the conditions of Lemma 4.1 are satisfied.

Then the critical point $\hat{u}_{\mathbf{n}}$ satisfies the following estimates:

- $\|\hat{u}_{\mathbf{n}} - u^*\|_{H^1(-\infty, X_1 - z)} \leq 3\tilde{\sigma}$,
- $\forall i \in [1, N-1] \cap 2\mathbb{Z}$, $\|\hat{u}_{\mathbf{n}}(\cdot + C_i) - u^*\|_{H^1(-X_i + z, X_{i+1} - z)} \leq 3\tilde{\sigma}$,
- $\forall i \in [1, N-1] \cap (2\mathbb{Z} + 1)$, $\|\hat{u}_{\mathbf{n}}(\cdot + C_i) - u_*\|_{H^1(-X_i + z, X_{i+1} - z)} \leq 3\tilde{\sigma}$,
- $\|\hat{u}_{\mathbf{n}}(\cdot + C_N) - u_{\pm}\|_{H^1(-X_N + z, \infty)} \leq 3\tilde{\sigma}$, where $\pm = +$ for N even, $\pm = -$ for N odd.

Now we define $\hat{v}_{\mathbf{n}} = \mathcal{L}\hat{u}_{\mathbf{n}}$. Since $\hat{u}_{\mathbf{n}}$ is a critical point of J , $(\hat{u}_{\mathbf{n}}, \hat{v}_{\mathbf{n}})$ is a solution of (1.3)-(1.4). From Lemma 2.1 we also have $\hat{v}_{\mathbf{n}} = G * \hat{u}_{\mathbf{n}}$ with $G(x) = \frac{1}{2\sqrt{\gamma}} e^{-\sqrt{\gamma}|x|}$. From this and the above estimates on $\hat{u}_{\mathbf{n}}$, one can conclude that for $\tilde{\sigma}$ small enough and Δ large enough (both depending on σ), $(\hat{u}_{\mathbf{n}}, \hat{v}_{\mathbf{n}})$ satisfies the properties (a) and (b) of Theorem 1.2. This theorem is thus proved. \blacksquare

5 Stability

From the proof of Theorem 1.2, we know that $\hat{u}_{\mathbf{n}}$ is a local minimizer of J and $\hat{v}_{\mathbf{n}} = \mathcal{L}\hat{u}_{\mathbf{n}}$. To prove Theorem 1.4, we introduce the functional

$$\mathcal{E}(u, v) := J(u) + \frac{\gamma}{2(1 + \hat{\delta})} \|v - \mathcal{L}u\|^2 \quad (5.1)$$

defined for $u \in H_{\hat{u}_n}$, $v \in H_{\hat{v}_n}$. Here and in the sequel, $\|\cdot\|$ means the $L^2(\mathbb{R})$ -norm, and the parameter $\hat{\delta} > 0$ will be chosen later. Proposition 3.1 shows that (\hat{u}_n, \hat{v}_n) is a local minimizer of \mathcal{E} for the natural topology of the affine space $H_{\hat{u}_n} \times (\hat{v}_n + L^2(\mathbb{R}))$, and it is an isolated critical point of \mathcal{E} up to translation in the spatial variable. Also, due to Corollary 2.8, \mathcal{E} satisfies the Palais-Smale condition in a small neighborhood of (\hat{u}_n, \hat{v}_n) .

Consider the Cauchy problem:

$$u_t - du_{xx} = f(u) - v, \quad (5.2)$$

$$\tau v_t - v_{xx} = u - \gamma v, \quad (5.3)$$

with the initial data in the function space $Y = Y_u \times Y_v$. Here $Y_u = H_{\hat{u}_n} \cap C_b(\mathbb{R})$, $Y_v = H_{\hat{v}_n} \cap C_b(\mathbb{R})$ and $C_b(\mathbb{R})$ is the set of bounded uniformly continuous functions on \mathbb{R} . For $(u_1, v_1), (u_2, v_2) \in Y$, define $\|(u_1, v_1) - (u_2, v_2)\|_Y = \|u_1 - u_2\|_{H^1(\mathbb{R})} + \|u_1 - u_2\|_{L^\infty(\mathbb{R})} + \|v_1 - v_2\|_{H^1(\mathbb{R})} + \|v_1 - v_2\|_{L^\infty(\mathbb{R})}$. Then with minor modification, Theorem 2.1 of [40] (see also [45], Theorem 14.2) shows that, for given initial data in Y , (5.2)-(5.3) has a unique solution $(u(\cdot, t), v(\cdot, t))$. This solution exists globally in time and $(u(\cdot, t), v(\cdot, t)) \in C([0, \infty), Y)$. For the proofs, we refer to [40, 45] for the detail, including the method of contracting rectangles. Similar results hold if we work on different function spaces; for instance, take $Y_u = \hat{u}_n + C_0^2(\mathbb{R})$ and $Y_v = \hat{v}_n + C_0^2(\mathbb{R})$ with the natural topology inherited from $C_0^2(\mathbb{R})$. Here

$$C_0^2(\mathbb{R}) = \{w \mid \frac{d^j w}{d^j x} \in C_b(\mathbb{R}) \text{ and } \lim_{|x| \rightarrow \infty} \frac{d^j w}{d^j x}(x) = 0 \text{ for } 0 \leq j \leq 2\}.$$

The above results will be used in the proofs of Theorem 1.4 and Theorem 1.5. In the next proposition, we extend a result of [18] to show that $\mathcal{E}(u, v)$ is a Lyapunov functional for the evolution flow generated by (1.1)-(1.2), for $\hat{\delta}$ chosen small enough.

Proposition 5.1 *Assume that $0 < \tau < \gamma^2$. Let $\hat{\delta} > 0$ such that $1 + \hat{\delta}/2 < \gamma^2/\tau$. Then for any smooth solution $(u(x, t), v(x, t))$ of (1.1)-(1.2),*

$$\begin{aligned} & \frac{d}{dt} \mathcal{E}(u(\cdot, t), v(\cdot, t)) \\ & \leq -\frac{\hat{\delta}}{2(1+\hat{\delta})} \|u_t\|^2 - \frac{1}{1+\hat{\delta}} \left(\frac{\gamma^2}{\tau} - 1 - \frac{\hat{\delta}}{2} \right) \|v - \mathcal{L}u\|^2 - \frac{\gamma}{(1+\hat{\delta})\tau} \left\| \frac{\partial}{\partial x} (v - \mathcal{L}u) \right\|^2 \\ & \leq 0. \end{aligned}$$

Proof. Let $w = v - \mathcal{L}u$. If $u \in Y_u$ and $v \in Y_v$ then $v - \mathcal{L}u \in H^1(\mathbb{R})$. By Theorem 2.3 of [40] (see also [45], Theorem 14.3) and density argument,

it suffices to treat the case $Y_u = \hat{u}_n + C_0^2(\mathbb{R})$ and $Y_v = \hat{v}_n + C_0^2(\mathbb{R})$. Clearly (1.1)-(1.2) is equivalent to

$$u_t = du_{xx} + u(u - \beta)(1 - u) - \mathcal{L}u - w, \quad (5.4)$$

$$\tau(w_t + \mathcal{L}u_t) = w_{xx} - \gamma w. \quad (5.5)$$

In terms of (u, w) , we rewrite (5.1) as

$$\mathcal{E}_1(u, w) := J(u) + \frac{\gamma}{2(1 + \hat{\delta})} \|w\|^2. \quad (5.6)$$

Let $(u(x, t), w(x, t))$ be a solution of (5.4)-(5.5). Multiplying (5.5) by w and integrating by parts, we obtain

$$\tau[(w, w_t)_{L^2} + (w, \mathcal{L}u_t)_{L^2}] + \|w_x\|^2 + \gamma \|w\|^2 = 0. \quad (5.7)$$

Recall that \mathcal{L} is a self-adjoint operator in $L^2(\mathbb{R})$. Hence

$$(w, \mathcal{L}u_t)_{L^2} = (\mathcal{L}w, u_t)_{L^2}. \quad (5.8)$$

By making use of (5.7)-(5.8), a direct calculation gives

$$\begin{aligned} & \frac{d}{dt} \mathcal{E}_1(u(\cdot, t), w(\cdot, t)) \\ &= - \int_{-\infty}^{\infty} (du_{xx} + u(u - \beta)(1 - u) - \mathcal{L}u)u_t dx + \frac{\gamma}{1 + \hat{\delta}} (w, w_t)_{L^2} \\ &= - \|u_t\|^2 - (w, u_t)_{L^2} - \frac{\gamma}{1 + \hat{\delta}} \left((w, \mathcal{L}u_t)_{L^2} + \frac{1}{\tau} (\|w_x\|^2 + \gamma \|w\|^2) \right) \\ &\leq -\frac{1}{2} \|u_t\|^2 + \frac{1}{2} \|w\|^2 - \frac{\gamma}{(1 + \hat{\delta})} (\mathcal{L}w, u_t)_{L^2} \\ &\quad - \frac{\gamma^2}{(1 + \hat{\delta})\tau} \|w\|^2 - \frac{\gamma}{(1 + \hat{\delta})\tau} \|w_x\|^2 \\ &\leq -\frac{1}{2} \|u_t\|^2 + \frac{1}{2} \|w\|^2 + \frac{1}{(1 + \hat{\delta})} \|w\| \|u_t\| - \frac{\gamma^2}{(1 + \hat{\delta})\tau} \|w\|^2 - \frac{\gamma}{(1 + \hat{\delta})\tau} \|w_x\|^2 \\ &\leq -\frac{1}{2} \left(1 - \frac{1}{1 + \hat{\delta}} \right) \|u_t\|^2 + \left(\frac{1}{2} + \frac{1}{2(1 + \hat{\delta})} - \frac{\gamma^2}{(1 + \hat{\delta})\tau} \right) \|w\|^2 \\ &\quad - \frac{\gamma}{(1 + \hat{\delta})\tau} \|w_x\|^2 \\ &\leq -\frac{\hat{\delta}}{2(1 + \hat{\delta})} \|u_t\|^2 - \frac{1}{1 + \hat{\delta}} \left(\frac{\gamma^2}{\tau} - 1 - \frac{\hat{\delta}}{2} \right) \|w\|^2 - \frac{\gamma}{(1 + \hat{\delta})\tau} \|w_x\|^2 \leq 0. \end{aligned}$$

Since $\mathcal{E}_1(u(\cdot, t), w(\cdot, t)) = \mathcal{E}(u(\cdot, t), v(\cdot, t))$, the proposition is proved. \blacksquare

Proof of Theorem 1.4. By the construction, $\hat{u}_{\mathbf{n}}$ is a local minimizer of J . From Proposition 3.1, it is an isolated critical point of J , up to translations. Moreover, from Corollary 2.8 and the translation invariance of J , there is $\ell > 0$ such that any Palais-Smale sequence $\{u_j\}$ for J satisfying $\inf_{y \in \mathbb{R}} \|u_j - \hat{u}_{\mathbf{n}}(\cdot - y)\|_{H^1} \leq \ell/4$ is precompact up to translations.

Let us take $0 < \sigma_{\mathbf{n}} \leq \ell/4$ such that $\hat{u}_{\mathbf{n}}$ minimizes J on the set

$$\mathcal{X}_{\mathbf{n}} := \{u \in H_{\hat{u}_{\mathbf{n}}} : \inf_{y \in \mathbb{R}} \|u - \hat{u}_{\mathbf{n}}(\cdot - y)\|_{H^1} \leq \sigma_{\mathbf{n}}\},$$

and such that any critical point of J in $\mathcal{X}_{\mathbf{n}}$ is a translate $\hat{u}_{\mathbf{n}}(\cdot - y)$ of $\hat{u}_{\mathbf{n}}$. Since $\sigma_{\mathbf{n}} \leq \ell/4$, J satisfies the Palais-Smale condition up to translation in the set $\mathcal{X}_{\mathbf{n}}$, and there exists $h_{\mathbf{n}} > 0$ such that the local sublevel set

$$\mathcal{Y}_{\mathbf{n}} := \{u \in \mathcal{X}_{\mathbf{n}} : J(u) \leq J(\hat{u}_{\mathbf{n}}) + h_{\mathbf{n}}\}$$

is a subset of

$$\mathcal{X}'_{\mathbf{n}} := \left\{ u \in H_{\hat{u}_{\mathbf{n}}} : \inf_{y \in \mathbb{R}} \|u - \hat{u}_{\mathbf{n}}(\cdot - y)\|_{H^1(\mathbb{R})} \leq \frac{\sigma_{\mathbf{n}}}{2} \right\}.$$

Next we choose $0 < \rho_{\mathbf{n}} \leq \sigma_{\mathbf{n}}$ such that $\mathcal{E}(u, v) \leq J(\hat{u}_{\mathbf{n}}) + h_{\mathbf{n}}$ for all $(u, v) \in H_{\hat{u}_{\mathbf{n}}} \times H_{\hat{v}_{\mathbf{n}}}$ satisfying $\inf_{y \in \mathbb{R}} \{\|u - \hat{u}_{\mathbf{n}}(\cdot - y)\|_{H^1(\mathbb{R})} + \|v - \hat{v}_{\mathbf{n}}(\cdot - y)\|_{H^1(\mathbb{R})}\} < \rho_{\mathbf{n}}$.

Let $(u(x, t), v(x, t))$ be a solution of (1.1)-(1.2) such that

$$\|u(\cdot, 0) - \hat{u}_{\mathbf{n}}\|_{H^1(\mathbb{R})} + \|v(\cdot, 0) - \hat{v}_{\mathbf{n}}\|_{H^1(\mathbb{R})} < \rho_{\mathbf{n}}.$$

Then $J(u(\cdot, 0)) \leq \mathcal{E}(u(\cdot, 0), v(\cdot, 0)) \leq J(\hat{u}_{\mathbf{n}}) + h_{\mathbf{n}}$, and $u(\cdot, 0) \in \mathcal{X}_{\mathbf{n}}$, so $u(\cdot, 0) \in \mathcal{Y}_{\mathbf{n}}$. Since \mathcal{E} is a Lyapunov functional, we know that $\mathcal{E}(u(\cdot, t), v(\cdot, t)) \leq J(\hat{u}_{\mathbf{n}}) + h_{\mathbf{n}}$ for all $t \geq 0$. Hence

$$J(u(\cdot, t)) \leq J(\hat{u}_{\mathbf{n}}) + h_{\mathbf{n}} \text{ for all } t \geq 0. \quad (5.9)$$

By continuity of the flow $t \rightarrow u(\cdot, t)$ for the H^1 topology, let us prove by contradiction that $u(\cdot, t)$ stays in $\mathcal{Y}_{\mathbf{n}}$ for all $t \geq 0$. Otherwise, there would exist a maximal time T such that $u(\cdot, t) \in \mathcal{Y}_{\mathbf{n}}$ ($\forall 0 \leq t \leq T$). This together with (5.9) would imply $\inf_{y \in \mathbb{R}} \|u(\cdot, T) - \hat{u}_{\mathbf{n}}(\cdot - y)\|_{H^1(\mathbb{R})} = \sigma_{\mathbf{n}}$, which is not possible since $\mathcal{Y}_{\mathbf{n}} \subset \mathcal{X}'_{\mathbf{n}}$.

Now, since \mathcal{E} is bounded from below, it follows from Proposition 5.1 that

$$\int_0^\infty \|u_t(\cdot, t)\|_{L^2(\mathbb{R})}^2 dt < \infty, \quad (5.10)$$

$$\int_0^\infty \|v(\cdot, t) - \mathcal{L}u(\cdot, t)\|_{L^2(\mathbb{R})}^2 dt < \infty \quad (5.11)$$

and

$$\int_0^\infty \left\| \frac{\partial}{\partial x} (v(\cdot, t) - \mathcal{L}u(\cdot, t)) \right\|_{L^2(\mathbb{R})}^2 dt < \infty. \quad (5.12)$$

Since $u(\cdot, t)$ stays in \mathcal{Y}_n for all $t \geq 0$, the flow $t \rightarrow (u(\cdot, t), v(\cdot, t))$ is uniformly continuous with respect to the H^1 metric, and $t \rightarrow u_t$ is uniformly continuous with respect to the L^2 metric. Consequently

$$\lim_{t \rightarrow \infty} \{ \|u_t(\cdot, t)\|_{L^2(\mathbb{R})} + \|v(\cdot, t) - \mathcal{L}u(\cdot, t)\|_{H^1(\mathbb{R})} \} = 0$$

and $J'(u(\cdot, t)) \rightarrow 0$ in $H^{-1}(\mathbb{R})$. Recall that the Palais-Smale condition holds in \mathcal{Y}_n up to translation. Hence the set $\{u(\cdot, t), t \geq 0\}$ is precompact up to translation for the $H^1(\mathbb{R})$ metric, and its limit points are necessarily critical points of J in \mathcal{Y}_n . Furthermore such a critical point must be a translate of \hat{u}_n . As a conclusion,

$$\liminf_{t \rightarrow \infty} \inf_{y \in \mathbb{R}} \{ \|u(\cdot, t) - \hat{u}_n(\cdot - y)\|_{H^1(\mathbb{R})} + \|v(\cdot, t) - \hat{u}_n(\cdot - y)\|_{H^1(\mathbb{R})} \} = 0.$$

■

6 Unstable waves

In this section we are going to prove Theorem 1.5. Some notations will be taken from the proof of Theorem 1.2. The Lyapunov-Schmidt reduction will be very similar, so the details will be omitted. But note that there will be a shift of π/ω in the condition on the distance between bumps: see (S'_5) below, compared with (S_9) . Then a localized mountain-pass principle will be employed instead of minimization, in order to construct unstable two-bump solutions.

Starting with the set $\mathcal{V}_{h,z}$ defined by (3.5), we construct, for h small and K, z, n large enough, a smooth map \hat{b}_n from $(\mathcal{V}_{h,z})^2 \times [-\nu, \nu]$ to $H^1(\mathbb{R})$; here to each (u_0, u_1, x) in $(\mathcal{V}_{h,z})^2 \times [-\nu, \nu]$, we associate the unique function $u \in H^1(\mathbb{R})$ satisfying the following conditions:

- (S'_1) $u \equiv u_0$ on $(-z, z)$,
- (S'_2) $u \equiv u_1(X - \cdot)$ on $(X - z, X + z)$,
- (S'_3) $\|u - u_-\|_{H^1(-\infty, -z)} \leq Kr_-$,
- (S'_4) $\|u - u_-\|_{H^1(X+z, \infty)} \leq Kr_+$,
- (S'_5) $X = x + \kappa_+ + \frac{\pi(2n+1)}{\omega}$,
- (S'_6) $(u, \mathcal{L}u)$ satisfies (1.3)-(1.4) on each of the intervals $(-\infty, -z]$, $[z, X - z]$, $[X + z, +\infty)$.

With this definition of \check{b}_n , we define $\check{J} := J \circ \check{b}_n$. If (u_0, u_1, x) is a critical point of \check{J} in the interior of $\mathcal{V}_{h,z}^2 \times [-\nu, \nu]$, then $\check{b}_n(u_0, u_1, x)$ is a critical point of J .

Moreover, adapting Lemma 2.9 to the present situation, we see that, for each n large enough, there is a small $\mu_n > 0$ such that if $x \in [-\nu, -\nu/2]$ then

$$\partial_x \check{J}(u_0, u_1, x) \geq \mu_n, \quad (6.1)$$

while for $x \in [\nu/2, \nu]$,

$$\partial_x \check{J}(u_0, u_1, x) \leq -\mu_n. \quad (6.2)$$

Denoting u_z^* the restriction of u^* to $[-z, z]$, we set

$$\Gamma := \{g \in C([-1, 1], (\mathcal{V}_{h,z})^2 \times [-\nu, \nu]) : g(\pm 1) = (u_z^*, u_z^*, \pm \nu)\}$$

and for each large n assigned in (S'_5) , we define

$$c_n := \inf_{g \in \Gamma} \max\{\check{J} \circ g(\tau) : \tau \in [-1, 1] \text{ and } g(\tau) \in (\mathcal{V}_{h,z})^2 \times [-\nu/2, \nu/2]\}.$$

We also introduce the following notations:

$$\check{J}^a = J^{-1}((-\infty, a]), \quad \check{J}_a = J^{-1}([a, +\infty)), \quad \check{J}^{\dot{a}} = J^{-1}((-\infty, a)),$$

$$\text{Cr}(\check{J}, a) = \{(u_0, u_1, x) \in (\mathcal{V}_{h,z})^2 \times [-\nu, \nu] : \check{J}(u_0, u_1, x) = a \text{ and } D\check{J}(u_0, u_1, x) = 0\}.$$

For any $g \in \Gamma$, there is $\tau_0 \in [-1, 1]$ such that $g(\tau_0) \in (\mathcal{V}_{h,z})^2 \times \{0\}$. As a consequence, $c_n \geq \inf\{\check{J}(u_0, u_1, 0) : (u_0, u_1) \in (\mathcal{V}_{h,z})^2\}$. Moreover, when n tends to infinity, $\check{J}(u_0, u_1, x)$ converges uniformly to $J_z(u_0) + J_z(u_1)$ on $(\mathcal{V}_{h,z})^2 \times [-\nu, \nu]$, so $c_n \geq 2J(u^*) + o(1)_{n \rightarrow \infty}$.

On the other hand, the function $g_0 : \tau \in [-1, 1] \mapsto (u_z^*, u_z^*, \nu\tau)$ belongs to Γ . As a consequence, $c_n \leq \sup_{\tau \in [-1/2, 1/2]} \check{J} \circ g_0(\tau) = 2J(u^*) + o(1)_{n \rightarrow \infty}$. Thus, taking n large enough, we may impose $\check{J}^{c_n + h/4} \subset (\mathcal{V}_{3h/4, z})^2 \times [-\nu, \nu]$.

Remembering that \check{J} is a C^1 functional on $F = (\mathcal{V}_{h,z})^2 \times [-\nu, \nu]$ satisfying the Palais-Smale condition and that its critical points are isolated, we obtain the following deformation lemma:

Lemma 6.1 *For $\bar{\varepsilon} > 0$ small enough, if V, W are open neighborhoods of $\text{Cr}(\check{J}, c_n)$ in F such that $\bar{V} \subset W$ and $\text{dist}(V, F \setminus W) > 0$, there exist $\varepsilon \in (0, \bar{\varepsilon})$ and a deformation $\eta \in C([0, 1] \times F, F)$ such that:*

- (i) $\eta(0, \cdot) = \text{id}_F$,
- (ii) $\eta(1, (\check{J}^{c_n + \varepsilon} \setminus V) \cup ((\mathcal{V}_{h,z})^2 \times ([-\nu, -\nu/2] \cup [\nu/2, \nu])))$
 $\subset \check{J}^{c_n - \varepsilon} \cup ((\mathcal{V}_{h,z})^2 \times ([-\nu, -\nu/2] \cup [\nu/2, \nu]))$,
- (iii) $\eta([0, 1] \times \bar{V}) \subset W$,
- (iv) $\eta(t, u_0, u_1, x) = (u_0, u_1, x)$,
 $\forall (t, u_0, u_1, x) \in [0, 1] \times (\check{J}^{c_n + \varepsilon} \cup \check{J}^{c_n - \varepsilon} \cup ((\mathcal{V}_{h,z})^2 \times ([-\nu, -3\nu/4] \cup [3\nu/4, \nu])))$.

Proof. This result is a variant of Lemma 1 in [29], and its proof is similar. Using (6.1)-(6.2), one builds a locally Lipschitz pseudo-gradient vector field Ξ for \check{J} on $\check{J}^{c_n+2\bar{\varepsilon}} \cap \check{J}_{c_n-2\bar{\varepsilon}} \setminus \text{Cr}(\check{J}, c_n)$ such that, for any $(u_0, u_1, x) \in \check{J}^{c_n+2\bar{\varepsilon}} \cap \check{J}_{c_n-2\bar{\varepsilon}} \cap ((\mathcal{V}_{h,z})^2 \times ([-\nu, -\nu/2] \cup [\nu/2, \nu]))$, the x -component of $\Xi(u_0, u_1, x)$ has the sign of x . Then one builds a locally Lipschitz cut-off function χ equal to 1 on $\check{J}^{c_n+\varepsilon} \cap \check{J}_{c_n-\varepsilon} \cap ((\mathcal{V}_{h,z})^2 \times [-\nu/2, \nu/2]) \setminus V$ and vanishing on $V' \cup \check{J}_{c_n+\bar{\varepsilon}} \cup \check{J}^{c_n-\bar{\varepsilon}} \cup ((\mathcal{V}_{h,z})^2 \times ([-\nu, -3\nu/4] \cup [3\nu/4, \nu]))$ with V' an open neighborhood of $\text{Cr}(\check{J}, c_n)$ such that $\bar{V}' \subset V$. The deformation η is obtained as the flow of $-\chi\Xi$, after a suitable time reparametrization. \blacksquare

Now, using Lemma 6.1, we can easily adapt Hofer's arguments in [29], and find a "mountain-pass type" critical point $(u_0^\#, u_1^\#, x^\#)$ of \check{J} at level c_n .

Setting $(\check{u}_n, \check{v}_n) := (\check{b}_n(u_0^\#, u_1^\#, x^\#), \mathcal{L}\check{b}_n(u_0^\#, u_1^\#, x^\#))$, we have thus found a two-bump solution of (1.3)-(1.4): for h small and n large enough, $(\check{u}_n, \check{v}_n)$ satisfies Theorem 1.5(i), (ii).

Moreover, "mountain-pass type" means that for any neighborhood Ω of $(u_0^\#, u_1^\#, x^\#)$, the local sublevel set $\Omega \cup \check{J}^{c_n}$ is non-empty and not path-connected. We are now going to prove the instability of $(\check{u}_n, \check{v}_n)$ for (1.1)-(1.2).

Let us choose $\epsilon_n > 0$ such that J satisfies the Palais-Smale condition, up to translation, in the set

$$\mathcal{N}_n := \{u \in H^1(\mathbb{R}) : \inf_{y \in \mathbb{R}} \|u - \check{u}_n(\cdot - y)\|_{H^1(\mathbb{R})} \leq 2\epsilon_n\},$$

and such that any critical point of J in \mathcal{N}_n is a translate of \check{u}_n . Since $(u_0^\#, u_1^\#, x^\#)$ is of mountain-pass type, for any $0 < \rho < \epsilon_n$ there exists $u_\rho \in H^1(\mathbb{R})$ such that

$$\|u_\rho - \check{u}_n\|_{H^1(\mathbb{R})} + \|\mathcal{L}u_\rho - \check{v}_n\|_{H^1(\mathbb{R})} < \rho \quad \text{and} \quad J(u_\rho) < c_n = J(\check{u}_n). \quad (6.3)$$

Let $\{(u(x, t), v(x, t)) : (x, t) \in \mathbb{R} \times \mathbb{R}_+\}$ be the solution of (1.1)-(1.2) with initial datum $(u(\cdot, 0), v(\cdot, 0)) = (u_\rho, \mathcal{L}u_\rho)$. Using again the Lyapunov functional \mathcal{E} defined in Section 5, we see that $\mathcal{E}(u_\rho, \mathcal{L}u_\rho) = J(u_\rho)$ and thus

$$J(u(\cdot, t)) \leq \mathcal{E}(u(\cdot, t), v(\cdot, t)) \leq J(u_\rho) < J(\check{u}_n), \quad \forall t \geq 0.$$

To complete the proof, we argue indirectly to show the existence of $\tau_n(\rho) \geq 0$ such that

$$\inf_{y \in \mathbb{R}} \{\|u(\cdot, t) - \check{u}_n(\cdot - y)\|_{H^1(\mathbb{R})} + \|v(\cdot, t) - \check{v}_n(\cdot - y)\|_{H^1(\mathbb{R})}\} \geq \epsilon_n, \quad \forall t \geq \tau_n(\rho).$$

Assuming that $\tau_n(\rho)$ does not exist, we find a sequence $\{t_j\}$ of times, with $t_j \rightarrow \infty$ as $j \rightarrow \infty$, and a sequence of translations $\{y_j\}$ such that

$$\|u(\cdot, t_j) - \check{u}_n(\cdot - y_j)\|_{H^1(\mathbb{R})} + \|v(\cdot, t_j) - \check{v}_n(\cdot - y_j)\|_{H^1(\mathbb{R})} < \epsilon_n.$$

By the continuity properties of the flow, we have

$$\|u(\cdot, t) - \check{u}_n(\cdot - y_j)\|_{H^1(\mathbb{R})} + \|v(\cdot, t) - \check{v}_n(\cdot - y_j)\|_{H^1(\mathbb{R})} \leq 2\epsilon_n, \quad \forall t_j \leq t \leq t_j + \theta$$

with $\theta > 0$ independent of j . Since \mathcal{E} is bounded from below, applying Proposition 5.1 yields

$$\lim_{j \rightarrow \infty} \int_{t_j}^{t_j + \theta} \|u_t(\cdot, t)\|_{L^2(\mathbb{R})}^2 dt = 0, \quad (6.4)$$

$$\lim_{j \rightarrow \infty} \int_{t_j}^{t_j + \theta} \|v(\cdot, t) - \mathcal{L}u(\cdot, t)\|_{L^2(\mathbb{R})}^2 dt = 0 \quad (6.5)$$

and

$$\lim_{j \rightarrow \infty} \int_{t_j}^{t_j + \theta} \left\| \frac{\partial}{\partial x} (v(\cdot, t) - \mathcal{L}u(\cdot, t)) \right\|_{L^2(\mathbb{R})}^2 dt = 0. \quad (6.6)$$

Pick $t'_j \in [t_j, t_j + \theta]$ such that

$$\lim_{j \rightarrow \infty} \{ \|u_t(\cdot, t'_j)\|_{L^2(\mathbb{R})}^2 + \|v(\cdot, t'_j) - \mathcal{L}u(\cdot, t'_j)\|_{H^1(\mathbb{R})}^2 \} = 0.$$

Then $u^{(j)} := u(\cdot, t'_j)$ is a Palais-Smale sequence for J in \mathcal{N}_n , and $\lim J(u^{(j)}) \leq J(u_\rho) < J(\check{u}_n)$. So, using the Palais-Smale condition up to translation in \mathcal{N}_n , we get a critical point in that set, at a critical level less than $J(\check{u}_n)$. This is absurd since the only critical points of J in \mathcal{N}_n are the translates of \check{u}_n . Theorem 1.5 is thus proved. \square

7 Appendix

In this Appendix, we clarify the conditions on the parameters ensuring that both (u_-, v_-) and (u_+, v_+) are saddle-focus equilibria at the same energy level.

First of all, given $\beta \in (0, 1)$ and $d, \gamma > 0$, the equilibria of (1.3)-(1.4) are of the form $(u, \frac{u}{\gamma})$ where u is solution of the polynomial equation $u(u - \beta)(1 - u) = \frac{u}{\gamma}$. An obvious solution is 0, and this gives a first equilibrium $(u_-, v_-) = (0, 0)$ having energy zero. If $\gamma > \frac{4}{(1-\beta)^2}$ there are two other equilibria, obtained by solving $u^2 - (1 + \beta)u + (\beta + \frac{1}{\gamma}) = 0$. By an elementary computation, one easily checks that one of these additional equilibria has zero energy if and only if $\gamma = \frac{9}{2\beta^2 - 5\beta + 2}$, with $\beta < 1/2$. Then this equilibrium is $(u_+, v_+) = (\frac{2(\beta+1)}{3}, \frac{2(\beta+1)}{3\gamma})$. The other one is $(u_+/2, v_+/2)$ and it is a center of symmetry of our problem. This is better seen if one sets $k = \frac{1}{3}(\beta^2 - \beta + 1)$ and makes the change of variables

$$\begin{cases} u = (\beta + 1)/3 + \sqrt{k}U, \\ v = (\beta + 1)/3\gamma + \sqrt{k}V. \end{cases} \quad (7.1)$$

Then, with our choice $\gamma = \frac{9}{2\beta^2 - 5\beta + 2}$, (1.3)-(1.4) is equivalent to the system

$$-dU'' = k(U - U^3) - V, \quad (7.2)$$

$$-V'' = U - \gamma V. \quad (7.3)$$

The new system is symmetric with respect to the origin. Its nonzero equilibria, which correspond to (u_{\pm}, v_{\pm}) in the new coordinates, are $\pm \left(\sqrt{1 - \frac{1}{k\gamma}}, \frac{1}{\gamma} \sqrt{1 - \frac{1}{k\gamma}} \right)$.

So, if we impose $0 < \beta < 1/2$ and $\gamma = \frac{9}{2\beta^2 - 5\beta + 2}$, it is clear that (1.3)-(1.4) has the same linearization at (u_-, v_-) and (u_+, v_+) . Linearizing this system at $(u_-, v_-) = (0, 0)$, we get an equation of the form $h'' = Bh$ where $h(x)$ is a column vector with two components, and

$$B = \begin{bmatrix} d^{-1}\beta & d^{-1} \\ -1 & \gamma \end{bmatrix}.$$

The equilibria (u_{\pm}, v_{\pm}) are of saddle-focus type when B has no real eigenvalue, which means that $(\text{tr } B)^2 - 4\det B < 0$. This condition may be written as follows:

$$\gamma^2 d^2 - 2(2 + \beta\gamma)d + \beta^2 < 0.$$

It holds when

$$\frac{2 + \beta\gamma - 2\sqrt{1 + \beta\gamma}}{\gamma^2} < d < \frac{2 + \beta\gamma + 2\sqrt{1 + \beta\gamma}}{\gamma^2}.$$

But we also need the functional J to be bounded from below, to ensure the existence of the minimizer u^* . This imposes the additional condition $d > \frac{1}{\gamma^2}$, as can be seen from Proposition 2.4.

In summary, to guarantee that (u_-, v_-) and (u_+, v_+) are saddle-focus equilibria having the same energy and that J is bounded from below, we require

$$0 < \beta < 1/2, \quad \gamma = 9/(2\beta^2 - 5\beta + 2), \\ \frac{1}{\gamma^2} \max(1, 2 + \beta\gamma - 2\sqrt{1 + \beta\gamma}) < d < \frac{2 + \beta\gamma + 2\sqrt{1 + \beta\gamma}}{\gamma^2},$$

as stated in (1.12).

Acknowledgments The authors thank an anonymous referee who gave valuable comments and suggestions to improve the paper. Research is supported in part by MOST 105-2115-M-007-009-MY3, the Ministry of Science and Technology, Taiwan. Part of the work was done when Chen was visiting the Université Paris-Dauphine, and Séré was visiting the National Tsing Hua University and National Center for Theoretical Sciences, Taiwan.

References

- [1] M. BODE, A. W. LIEHR, C. P. SCHENK AND H. -G. PURWINS, *Interaction of dissipative solitons: particle-like behaviour of localized structures in a three-component reaction-diffusion system*, Physica D 161 (2002), 45-66.
- [2] S. BOLOTIN, *The effect of singularities of the potential energy on the integrability of mechanical systems*, J. Appl. Math. Mech.(1.1)-(1.2) 48 (1984), 255-260.
- [3] S. BOLOTIN, *Variational methods of constructing chaotic motions in rigid-body dynamics*, J. Appl. Math. Mech. 56 (1992), 198-205.
- [4] S. BOLOTIN, *Variational criteria for nonintegrability and chaos in Hamiltonian systems*, pp.173-179 in: Hamiltonian Mechanics, J. Seimenis, ed., Proceedings of a NATO Advanced Research Workshop, NATO Adv. Sci. Inst. Ser. B Phys. 331, 1994.
- [5] B. BUFFONI, *Infinitely many large amplitude homoclinic orbits for a class of autonomous Hamiltonian systems*, J. Differential Equations 121 (1995), 109-120.
- [6] B. BUFFONI AND E. SÉRÉ, *A global condition for quasi-random behavior in a class of conservative systems*, Comm. Pure Appl. Math. 49 (1996), 285-305.
- [7] C.-N. CHEN, C.-C. CHEN AND C.-C. HUANG, *Traveling waves for the FitzHugh-Nagumo system on an infinite channel*, J. Differential Equations 261 (2016), 3010-3041.
- [8] C.-N. CHEN AND Y. S. CHOI, *Standing pulse solutions to FitzHugh-Nagumo equations*, Arch. Rational Mech. Anal. 206 (2012), 741-777.
- [9] C.-N. CHEN AND Y. S. CHOI, *Traveling pulse solutions to FitzHugh-Nagumo equations*, Calculus of Variations and Partial Differential Equations 54 (2015), 1-45.
- [10] C.-N. CHEN, Y. S. CHOI AND N. FUSCO, *The Γ -limit of traveling waves in the FitzHugh-Nagumo system*, J. Differential Equations 267 (2019), 1805-1835.
- [11] C.-N. CHEN, Y. S. CHOI AND X. HU, *An index method for stability analysis of traveling and standing waves*, preprint.
- [12] C.-N. CHEN, Y. S. CHOI, Y. HU AND X. REN, *Higher dimensional bubble profiles in a sharp interface limit of the FitzHugh-Nagumo system*, SIAM J. Math. Anal. 50 (2018), 5072-5095.

- [13] C.-N. CHEN, Y. S. CHOI AND X. REN, *Bubbles and droplets in a singular limit of the FitzHugh-Nagumo system*, Interfaces and Free Boundaries, 20 (2018), 165-210.
- [14] C.-N. CHEN, S.-I. EI AND S.-Y. TZENG, *Heterogeneity-induced effects for pulse dynamics in FitzHugh-Nagumo type systems*, Physica D: Nonlinear Phenomena, 382 (2018), 22-32.
- [15] C.-N. CHEN AND X. HU, *Maslov index for homoclinic orbits of Hamiltonian systems*, Ann. Inst. H. Poincaré Anal. Non Linéaire 24 (2007), 589-603.
- [16] C.-N. CHEN AND X. HU, *Stability criteria for reaction-diffusion systems with skew-gradient structure*, Comm. Partial Differential Equations 33 (2008), 189-208.
- [17] C.-N. CHEN AND X. HU, *Stability analysis for standing pulse solutions to FitzHugh-Nagumo equations*, Calculus of Variations and Partial Differential Equations, 49 (2014), 827-845.
- [18] C.-N. CHEN, S. JIMBO AND Y. MORITA, *Spectral comparison and gradient-like property in the FitzHugh-Nagumo type equations*, Nonlinearity 28 (2015), 1003-1016.
- [19] C.-N. CHEN, S.-Y. KUNG AND Y. MORITA, *Planar standing wavefronts in the FitzHugh-Nagumo equations*, SIAM J. Math. Anal. 46 (2014), 657-690.
- [20] C.-N. CHEN AND K. TANAKA, *A variational approach for standing waves of FitzHugh-Nagumo type systems*, J. Differential Equations 257 (2014), 109-144.
- [21] C.-N. CHEN AND S.-Y. TZENG, *Existence and multiplicity results for heteroclinic orbits of second order Hamiltonian systems*, J. Differential Equations 158 (1999), 211-250.
- [22] V. COTI ZELATI AND P.H. RABINOWITZ, *Homoclinic orbits for second order Hamiltonian systems possessing superquadratic potentials*, J. Amer. Math. Soc. 4 (1991), 693-727.
- [23] V. COTI ZELATI AND P.H. RABINOWITZ, *Homoclinic type solutions for a semilinear elliptic PDE on \mathbb{R}^n* , Comm. Pure Appl. Math. 45 (1992), 1217-1269.
- [24] E. N. DANCER AND S. YAN, *A minimization problem associated with elliptic systems of FitzHugh-Nagumo type*, Ann. Inst. H. Poincaré Anal. Nonlinéaire 21 (2004), 237-253.

- [25] R. L. DEVANEY, *Homoclinic orbits in Hamiltonian systems*, J. Differential Equations 21 (1976), 431-438.
- [26] A. DOELMAN, P. VAN HEIJSTER AND T. KAPER, *Pulse dynamics in a three-component system: existence analysis*, J. Dynam. Differential Equations 21 (2008), 73-115.
- [27] R. FITZHUGH, *Impulses and physiological states in theoretical models of nerve membrane*, Biophys. J. 1 (1961), 445-466.
- [28] A. L. HODGKIN AND A. F. HUXLEY, *A quantitative description of membrane current and its application to conduction and excitation in nerve*, J. Physiol. 117 (1952), 500-544.
- [29] H. HOFER, *A geometric description of the neighbourhood of a critical point given by the mountain-pass theorem*. J. London Math. Soc. 31 (1985), 566-570.
- [30] W. D. KALIES AND R. C. A. M. VANDERVORST, *Multitransition homoclinic and heteroclinic solutions of the extended Fisher-Kolmogorov equation*. J. Differential Equations 131 (1996), 209-228.
- [31] W. D. KALIES, J. KWAPISZ AND R. C. A. M. VANDERVORST, *Homotopy classes for stable connections between Hamiltonian saddle-focus equilibria*. Comm. Math. Phys. 193 (1998), 337-371.
- [32] W. D. KALIES, J. KWAPISZ, J. B. VANDENBERG AND R. C. A. M. VANDERVORST, *Homotopy classes for stable periodic and chaotic patterns in fourth-order Hamiltonian systems*, Comm. Math. Phys. 214 (2000), 573-592.
- [33] A. KATOK AND B. HASSELBLATT, *Introduction to the modern theory of dynamical systems*, Encyclopedia of Mathematics and Its Applications 54, Cambridge University Press, 1995.
- [34] A. W. LIEHR, *Dissipative Solitons in Reaction-Diffusion Systems*, Springer Series in Synergetics 70, Springer-Verlag, Berlin, 2013.
- [35] P. L. LIONS, *The concentration-compactness principle in the calculus of variations. The locally compact case, part 1*, Annales de l'IHP (C) 1 (2), 109-145 (1984).
- [36] J. N. MATHER, *Variational construction of orbits of twist diffeomorphisms*, J. Amer. Math. Soc. 4 (1991), 207-263.
- [37] J. NAGUMO, S. ARIMOTO AND S. YOSHIZAWA, *An active pulse transmission line simulating nerve axon*, Proc. I. R. E. 50 (1962), 2061-2070.

- [38] Y. OSHITA, *On stable nonconstant stationary solutions and mesoscopic patterns for FitzHugh-Nagumo equations in higher dimensions*, J. Differential Equations 188 (2003), 110-134.
- [39] P. H. RABINOWITZ AND E. W. STREDULINSKY, *Extensions of Moser-Bangert theory: Local minimal solutions. Progress in Nonlinear Differential Equations and their Applications*, 81. Birkhauser/Springer, New York, 2011.
- [40] J. RAUCH AND J. SMOLLER, *Qualitative theory of the FitzHugh-Nagumo equations*, Advances in Math. 27 (1978), 12-44.
- [41] C. REINECKE AND G. SWEERS, *A positive solution on \mathbb{R}^n to a system of elliptic equations of FitzHugh-Nagumo type*, J. Differential Equations 153 (1999), 292-312.
- [42] X. REN AND J. WEI, *Nucleation in the FitzHugh-Nagumo system: Interface-spike solutions*, J. Differential Equations 209 (2005), 266-301.
- [43] E. SÉRÉ, *Existence of infinitely many homoclinic orbits in Hamiltonian systems*, Math. Z. 209 (1992), 27-42.
- [44] E. SÉRÉ, *Looking for the Bernoulli shift*, Ann. Inst. H. Poincaré C, Anal. non Linéaire 10 (1993), 561-590.
- [45] J. SMOLLER, *Shock Waves and Reaction Diffusion Equations*, Second edition, Springer-Verlag, New York, 1994.
- [46] A. M. TURING, *The chemical basis of morphogenesis*, Phil. Trans. R. Soc. Lond. B 237 (1952), 37-72.
- [47] PETER VAN HEIJSTER, C. -N. CHEN, YASUMASA NISHIURA AND TAKASHI TERAMOTO, *Localized patterns in a three-component FitzHugh-Nagumo model revisited via an action functional*, J. Dyn. Differ. Equ. 30 (2018), 521-555.
- [48] J. WEI AND M. WINTER, *Clustered spots in the FitzHugh-Nagumo system*, J. Differential Equations 213 (2005), 121-145.
- [49] E. YANAGIDA, *Standing pulse solutions in reaction-diffusion systems with skew-gradient structure*, J. Dyn. Differ. Equ. 14 (2002), 189-205.