

### Atomic layer deposition of LaxZr1-xO2-delta (x=0.25) high-k dielectrics for advanced gate stacks

D. Tsoutsou, L. Lamagna, S. N. Volkos, A. Molle, S. Baldovino, Sylvie Schamm-Chardon, Pierre-Eugène Coulon, M. Fanciulli

#### ► To cite this version:

D. Tsoutsou, L. Lamagna, S. N. Volkos, A. Molle, S. Baldovino, et al.. Atomic layer deposition of LaxZr1-xO2-delta (x=0.25) high-k dielectrics for advanced gate stacks. Applied Physics Letters, 2009, 94 (5), pp.53504 - 53504. 10.1063/1.3075609 . hal-01745031

#### HAL Id: hal-01745031 https://hal.science/hal-01745031

Submitted on 9 Apr 2018  $\,$ 

**HAL** is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers. L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

## Atomic layer deposition of $La_x Zr_{1-x}O_{2-\delta}$ (x = 0.25) high-k dielectrics for advanced gate stacks

D. Tsoutsou, L. Lamagna, S. N. Volkos, A. Molle, S. Baldovino, S. Schamm, P. E. Coulon, and M. Fanciulli

Citation: Appl. Phys. Lett. **94**, 053504 (2009); doi: 10.1063/1.3075609 View online: https://doi.org/10.1063/1.3075609 View Table of Contents: http://aip.scitation.org/toc/apl/94/5 Published by the American Institute of Physics

#### Articles you may be interested in

Structural and electrical properties of atomic layer deposited Al-doped ZrO<sub>2</sub> films and of the interface with TaN electrode Journal of Applied Physics **112**, 014107 (2012); 10.1063/1.4731746

Stabilization of very high-k tetragonal phase in Ge-doped  $ZrO_2$  films grown by atomic oxygen beam deposition Journal of Applied Physics **106**, 024107 (2009); 10.1063/1.3182636

High-κ gate dielectrics: Current status and materials properties considerations Journal of Applied Physics **89**, 5243 (2001); 10.1063/1.1361065

## **Sci**light

Sharp, quick summaries illuminating the latest physics research

Sign up for **FREE**!

# Atomic layer deposition of $La_x Zr_{1-x}O_{2-\delta}$ (x=0.25) high-*k* dielectrics for advanced gate stacks

D. Tsoutsou,<sup>1</sup> L. Lamagna,<sup>1,a)</sup> S. N. Volkos,<sup>1</sup> A. Molle,<sup>1</sup> S. Baldovino,<sup>1</sup> S. Schamm,<sup>2</sup> P. E. Coulon,<sup>2</sup> and M. Fanciulli<sup>1,3</sup>

<sup>1</sup>Laboratorio Nazionale MDM, CNR-INFM, Via C. Olivetti 2, Agrate Brianza (MI) 20041, Italy <sup>2</sup>CEMES-CNRS, Université de Toulouse, 29 rue J. Marvig, BP 94347, 31055 Toulouse Cedex 4, France

<sup>3</sup>Dipartimento di Scienza dei Materiali, Università degli studi di Milano-Bicocca, Milano, Italy

(Received 16 September 2008; accepted 22 December 2008; published online 3 February 2009)

Thin  $La_x Zr_{1-x}O_{2-\delta}$  (*x*=0.25) high permittivity (*k*) films are grown on Si(100) by atomic layer deposition at 300 °C using (<sup>i</sup>PrCp)<sub>3</sub>La, (MeCp)<sub>2</sub>ZrMe(OMe) and O<sub>3</sub> species. Their properties are studied by grazing incidence x-ray diffraction, high resolution transmission electron microscopy, electron energy loss spectroscopy, x-ray photoelectron spectroscopy, and electrical measurements on the as-grown films and after vacuum annealing at 600 °C. Annealed films feature resistance to hygroscopicity, a large *k* value of around 30 and an acceptable leakage current density. A low-*k* silica-rich interlayer is also evidenced at both pristine and annealed high-*k*/Si interfaces. © 2009 American Institute of Physics. [DOI: 10.1063/1.3075609]

As the complementary metal oxide semiconductor (CMOS) device dimensions continue to decrease in order to improve chip performance, high permittivity (high-k) oxides are being introduced into the CMOS technology as alternative dielectrics to  $SiO_xN_y$  with the principal aim to suppress detrimental gate leakage current effects.<sup>1</sup> Among the various rare earth oxides<sup>2</sup> currently considered for the sub-45 nm high-k CMOS era is  $La_2O_3$  mainly due to its high k value  $(\sim 27 \text{ when crystallized into the hexagonal } P6_3/mmc \text{ phase})^3$ and wide band gap energy (~6 eV).<sup>4</sup> However, La(OH)<sub>3</sub> formation due to OH<sup>-</sup> groups absorption when La<sub>2</sub>O<sub>3</sub> is exposed to air significantly degrades its k value.<sup>5</sup>  $ZrO_2$  could be another option for the sub-45 nm high-k CMOS nodes because it satisfies technology requirements in terms of k value and band gap energy but the material is marginally chemically/thermally stable on Si,<sup>6</sup> therefore, restricting an aggressive scaling down of the equivalent oxide thickness (EOT). Alloying  $ZrO_2$  with  $La_2O_3$  to produce a ternary  $La_x Zr_{1-x}O_{2-\delta}$  (LZO) compound might be a promising route to engineer a material suitable for high-k dielectric applications. Among the various LZO growth techniques,<sup>7–12</sup> atomic layer deposition (ALD) appears to be the most attractive one for integration into the CMOS flow as it provides the capability to grow smooth and conformal films at low temperatures, allowing an accurate film thickness control.<sup>13</sup>

In this article, we report on the growth of LZO/Si interfaces by ALD using O<sub>3</sub>, a more powerful oxidizer than H<sub>2</sub>O tested by Gaskell *et al.*,<sup>10,12</sup> in combination with the same La [i.e., (<sup>i</sup>PrCp)<sub>3</sub>La] and Zr [i.e., (MeCp)<sub>2</sub>ZrMe(OMe); also known as ZrD-04] precursors as in Refs. 10 and 12. Contrary to Refs. 10–12, which are almost exclusively concerned with the properties of as-grown LZO films, our study merely focuses on annealed films since CMOS device processing involves postdeposition thermal treatments. LZO films were deposited in the Savannah 200 (Cambridge Nanotech Inc.) ALD reactor on H-terminated *n*-Si(100) at a growth temperature of 300 °C. (<sup>i</sup>PrCp)<sub>3</sub>La and ZrD-04 complexes were kept at 150 and 110 °C, respectively. O<sub>3</sub> was injected into the reactor at a concentration ~200 g/Nm<sup>3</sup>. ALD cycle number varied between 12 and 50, resulting in a stack thickness within the 5–30 nm range. La atomic fraction x = La/(La+Zr)=0.25, which was tuned by altering the ALD pulse number in the growth cycle, was estimated from the La 3d and Zr 3d core-level x-ray photoelectron spectroscopy (XPS) lines relying on the relevant atomic sensitivity factors. Postgrowth rapid thermal annealing (RTA) was performed for 60 s in vacuum (~1 mbar) at 600 °C. As opposed to RTA in forming gas or N<sub>2</sub>, annealing in vacuum permits monitoring of thermally induced structural, compositional, and electrical changes minimizing any contribution to them from the ambient. Moreover, vacuum should enhance the removal of the OH<sup>-</sup> groups from the La-based film.

LZO films were characterized by a combination of (i) grazing incidence x-ray diffraction (GIXRD), (ii) transmission electron microscopy in the high resolution (HRTEM) and electron energy loss spectroscopy (EELS) modes, (iii) XPS, (iv) capacitance-voltage (C-V), parallel conductance-voltage  $(G_p V)$  and current density-voltage (J-V)measurements. GIXRD spectra were collected with a position sensitive detector using the Cu  $K\alpha$  radiation. The local structural and chemical properties around the film/substrate interface and in the film were analyzed by HRTEM and EELS with a field emission transmission electron microscope (FEI Tecnai<sup>TM</sup> F20) operated at 200 kV, equipped with a spherical aberration corrector and, for local EELS studies, with a scanning stage (STEM) and an imaging filter (Gatan GIF TRIDIEM). Cross-sectional TEM samples (bulk is cut perpendicular to the film/substrate interface) were prepared by classical procedures. XPS spectra were recorded at a takeoff angle of 90° in a semispherical spectrometer employing a standard Al  $K\alpha$  source (1486.6 eV) with an energy resolution of 0.8 eV and a pass energy of 20 eV. A multiple function fit was applied to the XPS data. Briefly, a Shirley baseline was removed from each XPS spectrum and all the Si oxidation state contributions were then fitted by single Gaussian-Lorentzian functions. Room temperature electrical measurements were performed on metal insulator semiconductor

<sup>&</sup>lt;sup>a)</sup>Author to whom correspondence should be addressed. Electronic mail: luca.lamagna@mdm.infm.it.



FIG. 1. GIXRD spectra for the 24 nm ZrO<sub>2</sub> and LZO films. Incidence angle was  $1^\circ$  and angular resolution was  $0.03^\circ.$ 

(MIS) capacitors of  $7.8 \times 10^{-4}$  cm<sup>2</sup> area. Capacitance equivalent oxide thickness (CET) was determined from the device accumulation capacitance at 800 kHz.

Figure 1 displays GIXRD spectra for the 24 nm thick ZrO<sub>2</sub> and LZO films. Although the monoclinic phase is the most stable phase<sup>14</sup> for pure ZrO<sub>2</sub> below 1100 °C, the asgrown  $ZrO_2$  film [Fig. 1(a)], which was deposited at 300 °C using the ZrD-04 and O<sub>3</sub> recipe,<sup>15</sup> is found to adopt a cubic (c-) or tetragonal (t-)  $ZrO_2$  structure<sup>16</sup> owing to ALD growth under nonequilibrium conditions. An XRD-based delineation between the c- and t-ZrO<sub>2</sub> phases is difficult since the respective reflections<sup>16</sup> appear at vicinal angles. La incorporation into the host ZrO<sub>2</sub> matrix through the ternary ALD process mostly generates thin amorphous as-grown LZO films as evidenced from the low signal in Fig. 1(b) and subsequently confirmed by the TEM image in Fig. 2(a). As illustrated in Fig. 1(c), LZO films crystallize after RTA into the c- or t-ZrO<sub>2</sub> phases. Thinner (20 and 12 nm) LZO films were found (data not shown) to exhibit the same structural behavior as the 24 nm thick one. Contrary to La<sub>2</sub>O<sub>3</sub> films, LZO films are not hygroscopic as demonstrated by both Fig. 1(d)



FIG. 2. (Color online) Cross-section HRTEM images for the 12 nm thick LZO film. (a) As-grown state. (b) After RTA for 60 s in vacuum at 600  $^{\circ}$ C. Dotted lines are guide to the eye. (c) STEM-EELS profiles of elemental La/O, Zr/O, and Si/O ratios for the annealed state.



FIG. 3. (Color online) XPS Si 2p and 2s shape profiles at a take-off angle of 90° for a 5 nm thick LZO film. [(a)–(b)] Si 2p and 2s lines for the as-grown state. [(c)–(d)] Si 2p and 2s lines after RTA for 60 s in vacuum at 600 °C. Chemical shifts are 1.3, 2.1, 3.1, and 4.0 eV for 1+, 2+, 3+, and 4+ oxidation states, respectively.

and XPS O 1s line (data not shown) taken after 90 days long exposure to air.

HRTEM data were recorded on the 12 nm thick asgrown and annealed LZO stacks. As depicted in Figs. 2(a) and 2(b), an amorphous interface layer (IL) of clear contrast, which suggests a silica-rich nature, appears just above the Si substrate for both cases with the IL thickness increasing after RTA. Atop the IL, a predominantly amorphous high-k layer for the as-grown stack and a crystalline one for the annealed stack are observed. Crystallographic parameters measured on the projected planes of the HRTEM images and on selected area electron diffraction patterns (data not shown) correspond to those for the *c*- or *t*-ZrO<sub>2</sub> phases.<sup>16</sup> From the EELS spectra collected across the film/substrate interface towards the film and encompassing characteristic La, O, and Si ionization edges, the relative elemental La/O, Zr/O, and Si/O ratios were extracted and plotted as a function of the distance from the Si surface. For the annealed stack, Fig. 2(c) displays that past a Si- and O-rich near-interface region, the La/O and Zr/O ratios remain fairly constant not only along the film thickness but also in the direction parallel to the film/ substrate interface. Differently from Ref. 11, no cluster nucleation, as revealed by TEM-EELS analysis, occurs at the interface or in the film.

To unveil the IL composition, an XPS analysis was conducted on a 5 nm thick LZO stack. Since the La 4d line overlaps with the high binding energy (BE) tail of the Si 2pline, we considered the Si 2s line for the pertinent analysis. Figure 3 compares the Si 2s lines for the as-grown and annealed stacks. The Si 2s peak is decomposed into six components: elemental Si (Si<sup>0</sup>), four oxidation state components (Si<sup>1+</sup>, Si<sup>2+</sup>, Si<sup>3+</sup>, and Si<sup>4+</sup> corresponding to Si<sub>2</sub>O, SiO, Si<sub>2</sub>O<sub>3</sub>, and  $SiO_2$  species, respectively),<sup>17</sup> and an additional component (Si1-) related to a minor Si-metal bonding. In agreement with HRTEM analysis, XPS data manifest a RTA-induced increase in the IL thickness because the broad peak due to the overall oxide contribution is larger after RTA. They also corroborate the STEM-EELS data on the silica-rich nature of the IL since the Si<sup>4+</sup> component, which implies a SiO<sub>2</sub> bonding arrangement, dominates after RTA.

Figure 4(a) illustrates *C-V* curves for as-grown and annealed 13 nm thick LZO stacks. For both cases, accumulation capacitance exhibits minor frequency dispersion within



FIG. 4. Electrical characteristics for Al-gate MIS capacitors from a 13 nm thick LZO stack. (a) Multifrequency C-V curves for as-grown (empty symbols) and RTA-treated (filled symbols) devices. (b) Stack CET vs LZO film physical thickness plot for RTA-treated devices. (c) J-V curves for as-grown and RTA-treated devices. RTA was carried out for 60 s in vacuum at 600 °C. MIS capacitors received no postmetallization annealing.

the 5–800 kHz range. Table I summarizes electrical figures of merit derived from the data of Fig. 4(a). Evidently, RTA notably lowers the density of oxide trapped charge and interface defects  $(D_{it})$  while markedly enhancing the positive fixed charge density near the dielectric/Si interface. From the slope of a linear fit to a stack CET versus LZO physical thickness plot in Fig. 4(b), the k value for annealed LZO films is estimated to be  $\sim$ 30. Based on the theoretically calculated k value for different  $ZrO_2$  crystalline phases,<sup>14,20</sup> our result points more to a c- rather than t-ZrO<sub>2</sub> phases. In good agreement with the HRTEM and STEM-EELS analyses, the fit to the above plot intercepts with the y-axis at a point corresponding to an IL with a k value close to that of SiO<sub>2</sub>. Our as-grown LZO films feature a k value that varies with the stack thickness such that it is as low as 17.5 for a 12.5 nm thick stack and as high as 24 for a 22.5 nm thick stack. In the light of dark field TEM images (data not shown), this trend is ascribed to the fact that as-grown LZO films are not entirely amorphous but contain a small fraction of isolated nanocrystals, which contribute to a different extent to the film k value because their size is found to increase with increasing stack thickness. Consequently, it is inferred that the thermally induced crystallization of the as-grown amorphous LZO films is accompanied by a moderate k value enhancement also with respect to the parental ZrO<sub>2</sub> films.<sup>15</sup> Given the connection between phonon vibrational spectra and k value,<sup>18</sup> the above finding can be tentatively attributed to softening, through partial substitution of Zr with La atoms,<sup>19</sup> of the c- or t-ZrO<sub>2</sub> infrared-active transverse optical phonon mode with minimum frequency.<sup>20</sup> Finally, it is ascertained from Fig. 4(c) that although RTA significantly degrades, compared to the as-grown state, the stack hardness to leakage paths, an acceptable leakage current density still flows through the MIS capacitor when pulsed in strong accumulation and inversion.

In conclusion, thin and uniform LZO films were grown on Si(100) by ALD at 300  $^{\circ}$ C using (<sup>i</sup>PrCp)<sub>3</sub>La, ZrD-04, and

TABLE I. Figures of merit for 13 nm as-grown and annealed LZO stacks (W=hysteresis width,  $V_{FB}$ =flatband voltage).

Stack state	W <sup>a</sup> (mV)	$V_{\mathrm{FB}}^{}\mathrm{b}}$ (V)	$D_{\rm it}^{\rm c}$ (×10 <sup>12</sup> eV <sup>-1</sup> cm <sup>-2</sup> )
as-grown	$170(\pm 10) \\ 40(\pm 10)$	$-0.46(\pm 0.07)$	3.8(±0.4)
RTA in vacuum		$-1.47(\pm 0.15)$	2.1(±0.2)

<sup>a</sup>Referenced to  $C_{max}/2$  for a bidirectional  $-2 \leftrightarrow +2$  V sweep at 100 kHz. <sup>b</sup>Referenced to the experimental 800 kHz *C-V* curve; theoretical value derived from the Al-Si work function difference is -0.21 V.

<sup>c</sup>Hill–Coleman method (Ref. 21) applied to the  $G_pV$  peak at 100 kHz.

O<sub>3</sub>. Their structural, chemical, and electrical properties were studied after 60 s vacuum annealing at 600 °C. Annealed films crystallize into a cubic or tetragonal  $ZrO_2$  phases with a k value of ~30. They are also immune to hygroscopicity and exhibit an acceptable leakage current density. Although such a high k value is desirable for EOT scaling down, process optimization is still required to minimize the sizable low-k silica-rich IL at the high-k/Si junction and (near) dielectric/Si interface defect density.

This work has been done as a part of the European FP6-Program "REALISE" (Grant No. IST-NMP 016172).

- <sup>1</sup>J. Robertson, Rep. Prog. Phys. **69**, 327 (2006).
- <sup>2</sup>G. Scarel, A. Svane, and M. Fanciulli, in *Rare Earth Oxide Thin Films*, edited by M. Fanciulli and G. Scarel (Springer, Berlin, 2007), Vol. 106, Chap. 1, pp. 1–14.
- <sup>3</sup>G. Scarel, A. Debernardi, D. Tsoutsou, S. Spiga, S. C. Capelli, L. Lamagna, S. N. Volkos, M. Alia, and M. Fanciulli, Appl. Phys. Lett. **91**, 102901 (2007).
- <sup>4</sup>A. V. Prokofiev, A. I. Shelyakh, and B. T. Melekh, J. Alloys Compd. **242**, 41 (1996).
- <sup>5</sup>Y. Zhao, M. Toyama, K. Kita, K. Kyuno, and A. Toriumi, Appl. Phys. Lett. **88**, 072904 (2006).
- <sup>6</sup>M. C. Zeman, C. C. Fulton, G. Lucovsky, R. J. Nemanich, and W.-C. Yang, J. Appl. Phys. **99**, 023519 (2006).
- <sup>7</sup>K. Bobzin, E. Lugscheider, and N. Bagcivan, High Temp. Mater. Processes **10**, 103 (2006).
- <sup>8</sup>J. W. Seo, J. Fompeyrine, A. Guiller, G. Norga, C. Marchiori, H. Siegwart, and J.-P. Locquet, Appl. Phys. Lett. 83, 5211 (2003).
- <sup>9</sup>J. Fompeyrine, G. Norga, A. Guiller, C. Marchiori, J. W. Seo, J.-P. Locquet, H. Siegwart, D. Halley, and C. Rossel, Proceedings of the WODIM 2002 (IMEP, Grenoble, France, 2002), p. 65.
- <sup>10</sup>J. M. Gaskell, A. C. Jones, H. C. Aspinall, S. Taylor, P. Taechakumput, P. R. Chalker, P. N. Heys, and R. Odedra, Appl. Phys. Lett. **91**, 112912 (2007).
- <sup>11</sup>K. B. Jinesh, W. F. Besling, E. Tois, J. H. Klootwijk, W. Dekkers, M. Kaiser, F. Bakker, M. Tuominen, and F. Roozeboom, Appl. Phys. Lett. **93**, 062903 (2008).
- <sup>12</sup>J. M. Gaskell, A. C. Jones, P. R. Chalker, M. Werner, H. C. Aspinall, S. Taylor, P. Taechakumput, and P. N. Heys, Chem. Vap. Deposition 13, 684 (2007).
- <sup>13</sup>R. L. Puurunen, J. Appl. Phys. **97**, 121301 (2005).
- <sup>14</sup>G.-M. Rignanese, J. Phys.: Condens. Matter 17, R357 (2005).
- <sup>15</sup>A. Lamperti, L. Lamagna, G. Congedo, S. Spiga, and M. Fanciulli (unpublished).
- <sup>16</sup>G. Katz, J. Am. Ceram. Soc. 54, 531 (1971); G. Teufer, Acta Crystallogr. Sect. A: Cryst. Phys., Diffr., Theor. Gen. Crystallogr. 15, 1187 (1962).
- <sup>17</sup>A. Yoshigoe and Y. Teraoka, Surf. Interface Anal. 34, 432 (2002).

<sup>20</sup>X. Zhao and D. Vanderbilt, Phys. Rev. B **65**, 075105 (2002).

<sup>&</sup>lt;sup>18</sup>E. Bonera, G. Scarel, M. Fanciulli, P. Delugas, and V. Fiorentini, Phys. Rev. Lett. **94**, 027602 (2005).

<sup>&</sup>lt;sup>19</sup>S. A. Shevlin, A. Curioni, and W. Andreoni, Phys. Rev. Lett. **94**, 146401 (2005).

<sup>&</sup>lt;sup>21</sup>W. A. Hill and C. C. Coleman, Solid-State Electron. 23, 987 (1980).