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WIGNER MEASURES AND EFFECTIVE MASS THEOREMS

VICTOR CHABU, CLOTILDE FERMANIAN-KAMMERER, AND FABRICIO MACIÀ

ABSTRACT. We study a Schrödinger equation which describes the dynamics of an electron in a crystal in the presence of impurities. We consider the regime of small wave-lengths comparable to the characteristic scale of the crystal. It is well-known that under suitable assumptions on the initial data and for highly oscillating potentials, the wave function can be approximated by the solution of a simpler equation, the effective mass equation. Using Floquet-Bloch decomposition, as it is classical in this subject, we establish effective mass equations in a rather general setting. In particular, Bloch bands are allowed to have degenerate critical points, as may occur in dimension strictly larger than one. Our analysis leads to a new type of effective mass equations which are operator-valued and of Heisenberg form and relies on Wigner measure theory and, more precisely, to its applications to the analysis of dispersion effects.

Keywords: Bloch modes, semi-classical analysis on manifolds, Wigner measures, two-microlocal measures, Effective mass theory.

1. INTRODUCTION

1.1. The dynamics of an electron in a crystal and the effective mass equation. The dynamics of an electron in a crystal in the presence of impurities is described by a wave function $\Psi(t, x)$ that solves the Schrödinger equation:

$$(1.1) \quad \begin{cases} i\hbar\partial_t\Psi(t, x) + \frac{\hbar^2}{2m}\Delta_x\Psi(t, x) - Q_{\text{per}}(x)\Psi(t, x) - Q_{\text{ext}}(t, x)\Psi(t, x) = 0, \\ \Psi|_{t=0} = \Psi_0. \end{cases}$$

The potential Q_{per} is periodic with respect to some lattice in \mathbb{R}^d and describes the interactions between the electron and the crystal. The external potential Q_{ext} takes into account the effects of impurities on the otherwise perfect crystal. Here \hbar denotes the Planck constant and m is the mass of the electrons. In many cases of physical interest, the ratio ε between the mean spacing of the lattice and the characteristic length scale of variation of Q_{ext} is very small. After performing a suitable change of units, and rescaling the external potential and the wave function (see for instance [47]) the Schrödinger equation becomes:

$$(1.2) \quad \begin{cases} i\partial_t\psi^\varepsilon(t, x) + \frac{1}{2}\Delta_x\psi^\varepsilon(t, x) - \frac{1}{\varepsilon^2}V_{\text{per}}\left(\frac{x}{\varepsilon}\right)\psi^\varepsilon(t, x) - V_{\text{ext}}(t, x)\psi^\varepsilon(t, x) = 0, \\ \psi^\varepsilon|_{t=0} = \psi_0^\varepsilon. \end{cases}$$

The potential V_{per} is periodic with respect to a fixed lattice in \mathbb{R}^d , which, for the sake of definiteness will be assumed to be \mathbb{Z}^d .

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Effective Mass Theory consists in showing that, under suitable assumptions on the initial data ψ_0^ε , the solutions of (1.2) can be approximated for ε small by those of a simpler Schrödinger equation, the *effective mass equation*, which is of the form:

$$(1.3) \quad i\partial_t\phi(t, x) + \frac{1}{2}\langle B D_x, D_x \rangle\phi(t, x) - V_{\text{ext}}(t, x)\phi(t, x) = 0,$$

where, as usual, $D_x = \frac{1}{i}\partial_x$. The approximation has to be understood in the sense that any weak limit of the density $|\psi^\varepsilon(t, x)|^2 dx dt$ is the density $|\phi(t, x)|^2 dx dt$ as ε goes to 0. In equation (1.3), B is a $d \times d$ matrix called the *effective mass tensor*; it generates the *effective Hamiltonian*

$$H_{\text{eff}}(x, \xi) = \frac{1}{2}B\xi \cdot \xi + V_{\text{ext}}(t, x).$$

The effective mass tensor is an experimentally accessible quantity that can be used to study the effect of the impurities on the dynamics of the electrons. Both the question of finding those initial conditions for which the corresponding solutions of (1.2) converge (in a suitable sense) to solutions to the effective mass equation and that of clarifying the dependence of B on the sequence of initial data have been extensively studied in the literature [11, 47, 3, 30, 9]. The effective mass tensor is related to the critical points of the *Bloch modes*. These are the eigenvalues of the operator $P(\xi)$ on $L^2(\mathbb{T}^d)$ which is canonically associated with the equation (1.2),

$$(1.4) \quad P(\xi) = \frac{1}{2}|\xi - i\nabla_y|^2 + V_{\text{per}}(y), \quad y \in \mathbb{T}^d, \quad \xi \in \mathbb{R}^d.$$

We focus here on initial data which are structurally related with one of the Blochs modes in a sense that we will make precise later; we assume that this Bloch mode is of constant multiplicity and we introduce a new method for deriving rigorously equation (1.3). The advantage of this method is that it allows to treat the case where the critical points of the considered Bloch modes are degenerate, leading to the introduction of a new family of Effective mass equations which are of Heisenberg type. Our strategy is based on the analysis of the dispersion of PDEs by a Wigner measure approach which has led us to develop global *two microlocal Wigner measures* in this specific context, while they are only defined locally in general ([19, 20]).

Note that different scaling limits for equation (1.1) have been studied in the literature: the interested reader can consult, among many others, references [25, 28, 47, 31, 10, 2, 14, 46, 16].

1.2. Floquet-Bloch decomposition. The analysis of Schrödinger operators with periodic potentials has a long history that has its origins in the seminal works by Floquet [22] on ordinary differential equations with periodic coefficients, and by Bloch [12], who developed a spectral theory of periodic Schrödinger operators in the context of solid state physics. Floquet-Bloch theory can be used to study the spectrum of the perturbed periodic Schrödinger operator:

$$-\frac{\varepsilon^2}{2}\Delta_x + V_{\text{per}}\left(\frac{x}{\varepsilon}\right) + \varepsilon^2 V_{\text{ext}}(t, x),$$

see for instance [48, 32, 33, 34] and the references therein, and [45, 24, 30] for results in the semiclassical context. The Floquet-Bloch decomposition gives as a result that the corresponding Schrödinger evolution can be decoupled in an infinite family of dispersive-type equations for the so-called *Bloch modes*. We briefly recall the basic facts that we shall need by following the approach in [24, 25].

The Floquet-Bloch decomposition is based on assuming that the solutions to (1.2) depend on both the “slow” x and the “fast” x/ε variables. The fast variables should moreover respect the symmetries of the lattice. This leads to the following Ansatz on the form of the solutions ψ^ε of (1.2):

$$(1.5) \quad \psi^\varepsilon(t, x) = U^\varepsilon\left(t, x, \frac{x}{\varepsilon}\right),$$

where $U^\varepsilon(t, x, y)$ is assumed to be \mathbb{Z}^d -periodic with respect to the variable y (and, therefore, that it can be identified to a function defined on $\mathbb{R} \times \mathbb{R}^d \times \mathbb{T}^d$, where \mathbb{T}^d denotes the torus $\mathbb{R}^d/\mathbb{Z}^d$). The function U^ε then satisfies the equation:

$$(1.6) \quad \begin{cases} i\varepsilon^2 \partial_t U^\varepsilon(t, x, y) = P(\varepsilon D_x) U^\varepsilon(t, x, y) + \varepsilon^2 V_{\text{ext}}(t, x) U^\varepsilon(t, x, y), \\ U^\varepsilon|_{t=0} = U_0^\varepsilon(x, y), \text{ such that } \psi_0^\varepsilon = L^\varepsilon U_0^\varepsilon, \end{cases}$$

where the operator L^ε maps functions F defined on $\mathbb{R}^d \times \mathbb{T}^d$ on functions on \mathbb{R}^d according to:

$$(1.7) \quad L^\varepsilon F(x) := F\left(x, \frac{x}{\varepsilon}\right),$$

and $P(\varepsilon D_x)$ denotes the operator-valued Fourier multiplier associated with the symbol $\xi \mapsto P(\varepsilon \xi)$ defined in (1.4). The initial condition in (1.6) can be interpreted in terms of the natural embedding $L^2(\mathbb{R}_x^d) \hookrightarrow L^2(\mathbb{R}_x^d \times \mathbb{T}_y^d)$ by taking $U_0^\varepsilon(x, y) = \psi_0^\varepsilon(x) \otimes \mathbb{1}(y)$. One can also have more elaborated identifications depending on the structure of the initial data, as we shall see later. Identity (1.5) makes sense, since one can check that, under suitable assumptions on the initial datum, $U^\varepsilon(t, x, \cdot)$ has enough regularity with respect to the variable y (the fact that ψ^ε must be given by (1.5) following from the uniqueness of solutions to the initial value problem (1.2)).

Assuming that the function $y \mapsto V_{\text{per}}(y)$ is smooth is enough for proving that the operator $P(\xi)$ is self-adjoint on $L^2(\mathbb{T}^d)$ (with domain $H^2(\mathbb{R}^d)$) and has a compact resolvent. For the sake of simplicity, we shall make here this assumption, even though it can be relaxed into assuming $V_{\text{per}} \in L^p(\mathbb{T}^d)$ for some convenient set of indices p which authorizes Coulombian singularity in dimension 3 (see [35]). As a consequence of the fact that $P(\xi)$ has compact resolvent, there exist a non-decreasing sequence of eigenvalues (the so-called *Bloch energies*):

$$\varrho_1(\xi) \leq \varrho_2(\xi) \leq \dots \leq \varrho_n(\xi) \leq \dots \longrightarrow +\infty,$$

and an orthonormal basis of $L^2(\mathbb{T}^d)$ consisting of eigenfunctions $(\varphi_n(\xi, \cdot))_{n \in \mathbb{N}}$ (called *Bloch waves*):

$$P(\xi) \varphi_n(y, \xi) = \varrho_n(\xi) \varphi_n(y, \xi), \quad \text{for } y \in \mathbb{T}^d.$$

Moreover, the Bloch energies $\varrho_n(\xi)$ are $2\pi\mathbb{Z}^d$ -periodic whereas the Bloch waves satisfy

$$\varphi_n(y, \xi + 2\pi k) = e^{-i2\pi k \cdot y} \varphi_n(y, \xi), \quad \text{for every } k \in \mathbb{Z}^d.$$

This follows from the fact that for every $k \in \mathbb{Z}^d$, the operator $P(\xi + 2\pi k)$ is unitarily equivalent to $P(\xi)$ since $P(\xi + 2\pi k) = e^{-i2\pi k \cdot y} P(\xi) e^{i2\pi k \cdot y}$. It is proved in [49] that the Bloch energies ϱ_n are continuous and piecewise analytic functions of $\xi \in \mathbb{R}^d$. Actually, the set $\{(\xi, \varrho_n(\xi)), n \in \mathbb{N}, \xi \in \mathbb{R}^d\}$ is an analytic set of \mathbb{R}^{2d} . Moreover, if the multiplicity of the eigenvalue $\varrho_n(\xi)$ is equal to the same constant for all $\xi \in \mathbb{R}^d$, then ϱ_n and the eigenprojector Π_n on this mode are globally analytic functions of ξ . The reader can refer to [34] for a survey on the subject.

Observing that, via the decomposition in Fourier series, any function $U \in L^2(\mathbb{R}^d_x \times \mathbb{T}^d_y)$ can be written as:

$$U(x, y) = \sum_{k \in \mathbb{Z}^d} U_k(x) e^{i2\pi k \cdot y} \quad \text{with} \quad \|U\|_{L^2(\mathbb{R}^d \times \mathbb{T}^d)}^2 = \sum_{k \in \mathbb{Z}^d} \|U_k\|_{L^2(\mathbb{R}^d)}^2,$$

we denote by $H_\varepsilon^s(\mathbb{R}^d \times \mathbb{T}^d)$, for $s \geq 0$, the Sobolev space consisting of those functions $U \in L^2(\mathbb{R}^d \times \mathbb{T}^d)$ such that there exists $C > 0$ and:

$$(1.8) \quad \forall \varepsilon > 0, \quad \|U\|_{H_\varepsilon^s(\mathbb{R}^d \times \mathbb{T}^d)}^2 := \sum_{k \in \mathbb{Z}^d} \int_{\mathbb{R}^d} (1 + |\varepsilon \xi|^2 + |k|^2)^s |\widehat{U}_k(\xi)|^2 d\xi \leq C,$$

$$\text{where } \widehat{U}_k(\xi) = \int_{\mathbb{R}^d} e^{-ix \cdot \xi} U_k(x) dx.$$

1.3. Main result. We consider the following set of assumptions.

Assumption 1.1. (1) Assume V_{per} is smooth and real-valued and that V_{ext} is a continuous function in time taking values in the set of smooth, real-valued, bounded functions on \mathbb{R}^d with bounded derivatives.

(2) Assume that ϱ_n is a Bloch mode of constant multiplicity and that the set of critical points of ϱ_n

$$\Lambda_n := \{\xi \in \mathbb{R}^d, \quad \nabla \varrho_n(\xi) = 0\}$$

is a submanifold of \mathbb{R}^d .

(3) Assume that the Hessian $d^2 \varrho_n(\xi)$ is of maximal rank above each point $\xi \in \Lambda_n$ (or equivalently that $\text{Ker } d^2 \varrho_n(\xi) = T_\xi \Lambda_n$ for all $\xi \in \Lambda_n$).

(4) Assume that the initial data $\psi_0^\varepsilon(x)$ satisfies

$$\psi_0^\varepsilon(x) = U_0^\varepsilon \left(x, \frac{x}{\varepsilon} \right) \quad \text{with} \quad \widehat{U}_0^\varepsilon(\xi, \cdot) \in \text{Ran } \Pi_n(\varepsilon \xi),$$

with U_0^ε uniformly bounded in $H_\varepsilon^s(\mathbb{R}^d \times \mathbb{T}^d)$ for some $s > d/2$.

It will be convenient to identify ϱ_n to a function defined on $(\mathbb{R}^d)^*$ rather than \mathbb{R}^d (via the standard identification given by duality). Then we define the cotangent bundle of Λ_n as the union of all cotangent spaces to Λ_n

$$(1.9) \quad T^* \Lambda_n := \{(x, \xi) \in \mathbb{R}^d \times \Lambda_n : x \in T_\xi^* \Lambda_n\},$$

each fibre $T_\xi^* \Lambda_n$ is the dual space of the tangent space $T_\xi \Lambda_n$. Note that this is well-defined, since $T_\xi^* \Lambda_n \subset (\mathbb{R}^d)^{**} = \mathbb{R}^d$. We shall denote by $\mathcal{M}_+(T^* \Lambda_n)$ the set of positive Radon measures on $T^* \Lambda_n$. We also define the normal bundle of Λ_n which is the union of those linear subspaces of \mathbb{R}^d that are normal to Λ_n :

$$(1.10) \quad N \Lambda_n := \{(z, \xi) \in \mathbb{R}^d \times \Lambda_n : z \in N_\xi \Lambda_n\},$$

where $N_\xi \Lambda_n$ consists of those $x \in (\mathbb{R}^d)^{**} = \mathbb{R}^d$ that annihilate $T_\xi \Lambda_n$. Every point $x \in \mathbb{R}^d$ can be uniquely written as $x = v + z$, where $v \in T_\xi^* \Lambda_n$ and $z \in N_\xi \Lambda_n$. Given a function $\phi \in L^\infty(\mathbb{R}^d)$ we write $m_\phi(v, \xi)$, where $v \in T_\xi^* \Lambda_n$, to denote the operator acting on $L^2(N_\xi \Lambda_n)$ by multiplication by $\phi(v + \cdot)$. Note that assumption (3) implies that the Hessian of ϱ_n defines an operator $d^2 \varrho_n(\xi) D_z \cdot D_z$ acting on $L^2(N_\xi \Lambda_n)$ for any $\xi \in \Lambda_n$.

In the statement below, the weak limit of the energy density are described by means of a time-dependent family M_n of trace-class operators acting on a certain L^2 -space. More precisely, the operators M_n depend on $t \in \mathbb{R}$ and on $\xi \in \Lambda_n$, $v \in T_\xi^* \Lambda_n$; for every choice of these parameters, $M_n(t, v, \xi)$ is a trace-class operator acting on L^2 functions of the vector space $N_\xi \Lambda$. Note that $M_n(t, \cdot)$ can also be viewed as a section of a vector bundle over $T^* \Lambda_n$, namely: $\bigsqcup_{(v, \xi) \in T^* \Lambda} \mathcal{L}_+^1(L^2(N_\xi \Lambda_n))$.

Theorem 1.2. *Assume the hypotheses of Assumption 1.1. Then, there exist a subsequence $(\varepsilon_k)_{k \in \mathbb{N}}$, a positive measure $\nu_n \in \mathcal{M}_+(T^* \Lambda_n)$, and a measurable family of self-adjoint, positive, trace-class operator*

$$M_{0,n} : T_\xi^* \Lambda_n \ni (v, \xi) \longmapsto M_{0,n}(v, \xi) \in \mathcal{L}_+^1(L^2(N_\xi \Lambda_n)), \quad \text{Tr}_{L^2(N_\xi \Lambda_n)} M_{0,n}(v, \xi) = 1,$$

such that for every $a < b$ and every $\phi \in C_c(\mathbb{R}^d)$ one has:

$$\lim_{k \rightarrow \infty} \int_a^b \int_{\mathbb{R}^d} \phi(x) |\psi^{\varepsilon_k}(t, x)|^2 dx dt = \int_a^b \int_{T^* \Lambda_n} \text{Tr}_{L^2(N_\xi \Lambda_n)} [m_\phi(v, \xi) M_n(t, v, \xi)] \nu_n(dv, d\xi) dt,$$

where $M_n(\cdot, v, \xi) \in \mathcal{C}(\mathbb{R}; \mathcal{L}_+^1(L^2(N_\xi \Lambda_n)))$ solves the Heisenberg equation:

$$(1.11) \quad \begin{cases} i\partial_t M_n(t, v, \xi) + \left[\frac{1}{2} d^2 \varrho_n(\xi) D_z \cdot D_z + m_{V_{\text{ext}}(t, \cdot)}(v, \xi), M_n(t, v, \xi) \right] = 0, \\ M_n|_{t=0} = M_{0,n}. \end{cases}$$

Remark 1.3. We point out that the measure ν_n and the family of operators $M_{0,n}$ only depend on the subsequence $\psi_0^{\varepsilon_k}$ of initial data. The way of computing them will be made clear in Section 5.

When the critical points of $\varrho_n(\xi)$ are all non degenerate, then Λ_n is discrete and $2\pi\mathbb{Z}^d$ -periodic, $T^* \Lambda_n = \Lambda_n \times \{0\}$ and $N\Lambda_n = \mathbb{R}^d$. We then have the following corollary.

Corollary 1.4. *Assume we have Assumption 1.1 and that the critical points of $\varrho_n(\xi)$ are all non degenerate. Then the measure ν_n and the operator M_n of Theorem 1.2 above satisfy:*

- (1) *The operator $M_n(t, \xi)$ is the orthogonal projection on ψ_ξ which solves the effective mass equation:*

$$(1.12) \quad i\partial_t \psi_\xi(t, x) = \frac{1}{2} d^2 \varrho_n(\xi) D_x \cdot D_x \psi_\xi(t, x) + V_{\text{ext}}(t, x) \psi_\xi(t, x),$$

with initial data:

$$\psi_\xi|_{t=0} \text{ is the weak limit in } L^2(\mathbb{R}^d) \text{ of the sequence } \left(e^{-\frac{i}{\varepsilon_k} \xi \cdot x} \psi_0^{\varepsilon_k} \right).$$

- (2) *The measure ν_n is given by*

$$\nu_n = \sum_{\xi \in \Lambda_n} \alpha_\xi \delta_\xi, \quad \alpha_\xi = \|\psi_\xi|_{t=0}\|_{L^2(\mathbb{R}^d)}.$$

This corollary is well known and we refer to the work by Allaire and Piatnitski [3] or to [2] for similar results in a related problem; in that work homogenization and two-scale convergence techniques are used to obtain a precise description of the solution profile for data similar to ours and for Blochs mode having non-degenerated critical points. In [9], Barletti and Ben Abdallah obtained a result similar to Corollary 1.4 by following the approach initiated by Kohn and Luttinger in [37] consisting in introducing a (non-canonical) basis of modified Bloch functions.

The starting point in our approach is conceptually closer to that in [47], in the sense that we analyse the structure of Wigner measures associated to sequences of solutions. The main novelty here is the use of two-microlocal Wigner measures, that give a more explicit geometric description of the mechanism that underlies the Effective Mass Approximation, showing that it is a result of the dispersive effects associated to high-frequency solutions to the semiclassical Bloch band equations. Moreover, we are able to deal with the presence of non-isolated critical points on the Bloch energies and to prove Theorem 1.2. We believe our approach is sufficiently robust to be implemented on a Bloch band, isolated from the remainder of the spectrum, and consisting of several Bloch modes which may present crossings. We will devote further works to this specific problem. It is also interesting to notice that our result generalizes to initial data which are a finite sum of data satisfying (4) of Assumption 1.1. The weak limit of the energy density associated with the solution corresponding to this new data is the sum of weak limits of the energy densities of the solution associated with each term of the data, without any interference (see section 6.5 for a precise statement).

1.4. Strategy of the proof. The proof of Theorem 1.2 relies on the analysis of the solution U^ε to equation (1.6) with initial data U_0^ε as introduced in (4) of Assumption 1.1, and more precisely on its component U_n^ε on the n -th Bloch mode and its restriction ψ_n^ε by L^ε :

$$U_n^\varepsilon = \Pi_n(\varepsilon D_x)U^\varepsilon, \quad \psi_n^\varepsilon = L^\varepsilon U_n^\varepsilon.$$

It is shown in Section 6.3 that the family (ψ_n^ε) solves the equation

$$(1.13) \quad \begin{cases} i\varepsilon^2 \partial_t \psi_n^\varepsilon(t, x) - \varrho_n(\varepsilon D_x) \psi_n^\varepsilon(t, x) - \varepsilon^2 V_{\text{ext}}(t, x) \psi_n^\varepsilon(t, x) = \varepsilon^2 f_n^\varepsilon(t, x), \\ \psi_n^\varepsilon|_{t=0}(x) = \psi_0^\varepsilon(x) \end{cases}$$

with

$$f_n^\varepsilon = L^\varepsilon [\Pi(\varepsilon D_x), V_{\text{ext}}] U^\varepsilon,$$

There, we prove that

$$(1.14) \quad \forall T \in \mathbb{R}, \quad \exists C_T > 0, \quad \sup_{t \in [0, T]} \|\psi^\varepsilon(t, \cdot) - \psi_n^\varepsilon(t, \cdot)\|_{L^2(\mathbb{R}^d)} \leq C_T \varepsilon.$$

and

$$(1.15) \quad \exists C > 0, \quad \forall t \in \mathbb{R}, \quad \|f_n^\varepsilon(t, \cdot)\|_{L^2(\mathbb{R}^d)} \leq C \varepsilon.$$

Equation (1.14) shows that no other Bloch modes is concerned in the decomposition of U^ε and ψ^ε : the mass of ψ^ε remains above the specific mode ϱ_n because it is separated from the other ones. Therefore, a crucial step in this strategy consists in performing a detailed analysis of the dispersive equation (1.13).

1.5. Structure of the article. Sections 2 to 5 are devoted to the analysis of a dispersive equation of the form (1.13) in a more general setting. For this, we use pseudodifferential operators and semi-classical measures (Section 3) and we introduce two-microlocal tools (Section 4) that allow us to prove the main results of Section 2 in Section 5. Finally, in Section 6 we come back to the effective mass equations and prove Theorem 1.2, which requires additional results on the restriction operator L^ε , the projector $\Pi_n(\xi)$ and energy estimates for solutions to (1.6). Some Appendices are devoted to basic results about pseudodifferential calculus and trace-class operator-valued measures, and to the proof of technical lemmata.

2. QUANTIFYING THE LACK OF DISPERSION

As emphasized in the introduction, understanding the limiting behavior as $\varepsilon \rightarrow 0$ of the position densities of solutions to the Schrödinger equation (1.2) relies on a careful analysis of the solutions of equations of the form:

$$(2.1) \quad \begin{cases} i\varepsilon^2 \partial_t u^\varepsilon(t, x) = \lambda(\varepsilon D_x) u^\varepsilon(t, x) + \varepsilon^2 V_{\text{ext}}(t, x) u^\varepsilon(t, x) + \varepsilon^3 g^\varepsilon(t, x), & (t, x) \in \mathbb{R} \times \mathbb{R}^d, \\ u^\varepsilon|_{t=0} = u_0^\varepsilon, \end{cases}$$

where $(g^\varepsilon(t, \cdot))$ is locally uniformly bounded with respect to t in $L^2(\mathbb{R}^d)$.

This equation ceases to be dispersive as soon as $\lambda(\xi)$ has critical points $\xi \neq 0$, and this is always the case if λ is a Bloch energy. Heuristically, one can think that one of the consequences of a dispersive time evolution is a regularization of the high-frequency effects (that is associated to frequencies $\varepsilon\xi = c \neq 0$) caused by the sequence of initial data. These heuristics have been made precise in many cases; a presentation of our results from this point of view can be found in [15]. The reader can also find there a detailed account on the literature on the subject.

Here we show that, in the presence of critical points of λ , some of the high-frequency effects exhibited by the sequence of initial data persist after applying the time evolution (2.1). We provide a quantitative picture of this persistence by giving a complete description of the asymptotic behavior of the densities $|u^\varepsilon(t, x)|^2 dx dt$ associated to a bounded sequence (u^ε) of solutions to (2.1). We give an explicit procedure to compute all weak- \star accumulation points of the sequence of positive measures $(|u^\varepsilon(t, x)|^2 dx dt)$ in terms of quantities that can be obtained from the sequence of initial data (u_0^ε) . These results are of independent interest; we have thus chosen to present them in a more general framework than what it is necessary in our applications to Effective Mass Theory.

In order to obtain a non trivial result we must make sure that the characteristic length-scale of the oscillations carried by the sequence of initial data is of the order of ε . The following assumption is sufficient for our purposes:

H0 The sequence (u_0^ε) is uniformly bounded in $L^2(\mathbb{R}^d)$ and ε -oscillating, in the sense that its energy is concentrated on frequencies smaller or equal than $1/\varepsilon$:

$$(2.2) \quad \limsup_{\varepsilon \rightarrow 0^+} \int_{|\xi| > R/\varepsilon} |\widehat{u_0^\varepsilon}(\xi)|^2 d\xi \xrightarrow{R \rightarrow +\infty} 0.$$

We shall assume that λ is smooth and grows at most polynomially, and that its set of critical points is a submanifold of \mathbb{R}^d . More precisely, we impose the following hypotheses on λ and V :

H1 $V_{\text{ext}} \in C^\infty(\mathbb{R} \times \mathbb{R}^d)$ is bounded together with its derivatives and $\lambda \in C^\infty(\mathbb{R}^d)$, together with its derivatives, grows at most polynomially; *i.e.* there exists $N > 0$ such that, for every $\alpha \in \mathbb{N}_+^d$, one has:

$$\sup_{\xi \in \mathbb{R}^d} |\partial_\xi^\alpha \lambda(\xi)| (1 + |\xi|^N)^{-1} < \infty.$$

H2 The set

$$\Lambda := \left\{ \xi \in \mathbb{R}^d : \nabla \lambda(\xi) = 0 \right\}$$

is a connected, closed embedded submanifold of \mathbb{R}^d of codimension $0 < p \leq d$ and the Hessian $d^2\lambda$ is of maximal rank over Λ .

Hypothesis **H2** implies the existence of tubular coordinates in a neighborhood of Λ . A stronger version of **H2** is to suppose that all critical points of λ are non-degenerate (that is, the Hessian of λ , $d^2\lambda(\xi)$ is a non-degenerate quadratic form for every $\xi \in \Lambda$). This implies that $p = d$ and Λ is a discrete set in \mathbb{R}^d ; if moreover one has that λ is \mathbb{Z}^d -periodic, which is the situation when λ is a Bloch energy, this set is finite modulo \mathbb{Z}^d . We first state the main result of this section under this stronger hypothesis.

Theorem 2.1. *Suppose that the sequence of initial data (u_0^ε) verifies **H0**, denote by (u^ε) the corresponding sequence of solutions to (2.1). Suppose in addition that **H1** is satisfied and all critical points of λ are non-degenerate. Then there exists a subsequence $(u_0^{\varepsilon_k})$ such that for every $a < b$ and every $\phi \in C_c(\mathbb{R}^d)$ the following holds:*

$$(2.3) \quad \lim_{k \rightarrow \infty} \int_a^b \int_{\mathbb{R}^d} \phi(x) |u^{\varepsilon_k}(t, x)|^2 dx dt = \sum_{\xi \in \Lambda} \int_a^b \int_{\mathbb{R}^d} \phi(x) |u_\xi(t, x)|^2 dx dt,$$

where u_ξ solves the following Schrödinger equation:

$$(2.4) \quad i\partial_t u_\xi(t, x) = d^2\lambda(\xi) D_x \cdot D_x u_\xi(t, x) + V_{\text{ext}}(t, x) u_\xi(t, x),$$

with initial data:

$$u_\xi|_{t=0} \text{ is the weak limit in } L^2(\mathbb{R}^d) \text{ of the sequence } \left(e^{-\frac{i}{\varepsilon_k} \xi \cdot x} u_0^{\varepsilon_k} \right).$$

If $\Lambda = \emptyset$ then the right-hand side of (2.3) is equal to zero.

Note that u_ξ may be identically equal to zero even if the sequence (u_0^ε) oscillates in the direction ξ . For instance, if the sequence of initial data is a coherent state:

$$u_0^\varepsilon(x) = \frac{1}{\varepsilon^{d/4}} \rho \left(\frac{x - x_0}{\sqrt{\varepsilon}} \right) e^{\frac{i}{\varepsilon} \xi_0 \cdot x},$$

centered at a point (x_0, ξ_0) in phase space with $\rho \in C_0(\mathbb{R}^d)$, then $u_\xi|_{t=0} = 0$ for every $\xi \in \mathbb{R}^d$. Theorem (2.1) allows us to conclude that the corresponding solutions (u^ε) converge to zero in $L^2_{\text{loc}}(\mathbb{R} \times \mathbb{R}^d)$.

Theorem 2.1 can be interpreted as a description of the obstructions to the validity of smoothing-type estimates for the solutions to equation (2.1) in the presence of critical points of the symbol of the Fourier multiplier. We refer the reader to [15] for additional details concerning this issue and a simple proof of Theorem 2.1. Here, we obtain Theorem 2.1 as a particular case of a more general result which requires some geometric preliminaries.

As for the mode Bloch ϱ_n in the Introduction, we identify λ to a function defined on $(\mathbb{R}^d)^*$ rather than \mathbb{R}^d , and we associate with Λ its cotangent bundle $T^*\Lambda$ and its normal bundle $N\Lambda$. In the analogue of Theorem 2.1 in this context, the sum over critical points is replaced by an integral with respect to a measure over $T^*\Lambda$, and the Schrödinger equation (2.4) becomes a Heisenberg equation for a time-dependent family M of trace-class operators of $\bigsqcup_{(v, \xi) \in T^*\Lambda} \mathcal{L}^1_+(L^2(N_\xi\Lambda))$.

Theorem 2.2. *Let (u_0^ε) be a sequence of initial data satisfying **H0**, and denote by (u^ε) the corresponding sequence of solutions to (2.1). If **H1** and **H2** hold, then there exist a subsequence $(u_0^{\varepsilon_k})$, a positive measure $\nu \in \mathcal{M}_+(T^*\Lambda)$ and a measurable family of self-adjoint, positive, trace-class operators*

$$M_0 : T^*\Lambda \ni (v, \xi) \mapsto M_0(z, \xi) \in \mathcal{L}^1_+(L^2(N_\xi\Lambda)), \quad \text{Tr}_{L^2(N_\xi\Lambda)} M_0(v, \xi) = 1,$$

such that for every $a < b$ and every $\phi \in C_c(\mathbb{R}^d)$ one has:

$$(2.5) \quad \lim_{k \rightarrow \infty} \int_a^b \int_{\mathbb{R}^d} \phi(x) |u^{\varepsilon_k}(t, x)|^2 dx dt = \int_a^b \int_{T^*\Lambda} \text{Tr}_{L^2(N_\xi \Lambda)} [m_\phi(v, \xi) M_t(v, \xi)] \nu(dv, d\xi) dt,$$

where $t \mapsto M_t(v, \xi) \in \mathcal{C}(\mathbb{R}; \mathcal{L}_+^1(L^2(N_\xi \Lambda)))$ solves the following Heisenberg equation:

$$(2.6) \quad \begin{cases} i\partial_t M_t(v, \xi) = \left[\frac{1}{2} d^2 \lambda(\xi) D_z \cdot D_z + m_{V_{\text{ext}}(t, \cdot)}(v, \xi), M_t(v, \xi) \right], \\ M|_{t=0} = M_0. \end{cases}$$

Remark 2.3. When hypothesis **H2** about the rank of the Hessian $d^2 \lambda$ is dropped, then an additional term appears in (2.5) (see [15]).

When Λ consists of a set of isolated critical points, Theorems 2.1 and 2.2 are completely equivalent. Note that in this case, $T^*\Lambda = \{0\} \times \Lambda$ and the measure ν (which in this case is a measure depending on $\xi \in \mathbb{R}^d$ only) is simply

$$\nu = \sum_{\xi \in \Lambda} \alpha_\xi \delta_\xi,$$

where $\alpha_\xi = \|u_\xi|_{t=0}\|_{L^2(\mathbb{R}^d)}^2$. In addition, $N_\xi \Lambda = \mathbb{R}^d$ and the operator $M_t(\xi)$ (which again does not depend on z) is the orthogonal projection onto $u_\xi(t, \cdot)$ in $L^2(\mathbb{R}^d)$ (recall that u_ξ solves the Schrödinger equation (1.12)). These orthogonal projections satisfy the Heisenberg equation (2.6).

The proof of Theorem 2.2 follows a strategy developed in the references [39, 6, 4] in a different (though related) context. As in those references, the measure ν and the family of operators M_0 only depend on the subsequence of initial data $(u_0^{\varepsilon_k})$; we will see in Section 3 that they are defined as two microlocal Wigner measures of $(u_0^{\varepsilon_k})$ in the sense of [18, 19, 20, 39]. At this point, it might be useful to stress out that in this regime the limiting objects M, ν cannot be computed in terms of the Wigner/semiclassical measure of the sequence of initial data, as it is the case when dealing with the semiclassical limit. In [15], we have explicitly constructed sequences of initial data having the same semiclassical measure but such that their time dependent measures differ. This type of behavior was first remarked in this context in the case of the Schrödinger equation on the torus, see [38, 39].

We also emphasize that the original definition of two-microlocal Wigner measures performed in [19] and their extension to more general geometric setting [20] were only defined locally. We prove here that they extend to global objects in the geometric context of closed simply connected embedded submanifolds of \mathbb{R}^d ; related constructions were performed in the torus [39, 6, 4, 43] and the disk [5].

See also, that as soon as Λ has strictly positive dimension (*i.e.* it is not a union of isolated critical points), the measure ν may be singular with respect to the z variable, while when Λ consists in isolated points, the weak limit of the densities $|\psi^\varepsilon(t, x)|^2 dx$ are proved to be absolutely continuous with respect to the measure dx . See [15] for specific examples exhibiting this type of behavior; see also that reference for examples proving the necessity of hypothesis **H2**; it is shown there that different types of behavior can happen whenever the Hessian of λ is not of full rank on Λ .

The main idea of the proof comes from the following remark. Setting $v^\varepsilon(t, x) = u^\varepsilon(\varepsilon t, x)$, then (v^ε) solves the semi-classical equation

$$(2.7) \quad \begin{cases} i\varepsilon \partial_t v^\varepsilon(t, x) = \lambda(\varepsilon D_x) v^\varepsilon(t, x) + \varepsilon^2 V_{\text{ext}}(t, x) v^\varepsilon(t, x) + \varepsilon^3 g^\varepsilon(t, x), & (t, x) \in \mathbb{R} \times \mathbb{R}^d, \\ v^\varepsilon|_{t=0} = u_0^\varepsilon, \end{cases}$$

which means that, in the preceding analysis, we have performed the semiclassical limit $\varepsilon \rightarrow 0$ in (2.7) simultaneously with the limit $t/\varepsilon \rightarrow +\infty$. Such analysis, combining high-frequencies ($\varepsilon \rightarrow 0$) and long times ($t \sim t_\varepsilon \rightarrow +\infty$) is relevant if one wants to understand the behavior of solutions of (2.7) beyond the Ehrenfest time. This approach was followed in the case of confined geometries in the references [38, 4, 42]. Note also that in the particular case when $\lambda(\xi)$ is homogeneous of degree two, this change of time scale transforms the semiclassical equation (2.7) into the non-semiclassical one (that is, the one corresponding to $\varepsilon = 1$). Therefore, it is possible to derive results on the dynamics of the Schrödinger equation via this scaling limit [39, 8, 7, 5]. The reader can consult the survey articles [40, 6] and the introductory lecture notes [41] for additional details and references on this approach.

3. PSEUDODIFFERENTIAL OPERATORS AND SEMICLASSICAL MEASURES – PRELIMINARIES

In this section we recall some basic facts on Wigner distributions and semiclassical measures, which are the tools we are going to use to prove Theorem 2.2 and derive preliminary results about Wigner measures associated with families of solutions of equations of the form (2.1).

3.1. Wigner transform and Wigner measures. Wigner distributions provide a useful way for computing weak- \star accumulation points of a sequence of densities $|f^\varepsilon(x)|^2 dx$ constructed from a L^2 -bounded sequence (f^ε) of solutions of a semiclassical (pseudo) differential equation. They provide a joint physical/Fourier space description of the energy distribution of functions in \mathbb{R}^d . The Wigner distribution of a function $f \in L^2(\mathbb{R}^d)$ is defined as:

$$W_f^\varepsilon(x, \xi) := \int_{\mathbb{R}^d} f\left(x - \frac{\varepsilon v}{2}\right) \overline{f\left(x + \frac{\varepsilon v}{2}\right)} e^{i\xi \cdot v} \frac{dv}{(2\pi)^d},$$

and has several interesting properties (see, for instance, [23]).

- $W_f^\varepsilon \in L^2(\mathbb{R}^d \times \mathbb{R}^d)$.
- Projecting W_f^ε on x or ξ gives the position or momentum densities of f respectively :

$$\int_{\mathbb{R}^d} W_f^\varepsilon(x, \xi) d\xi = |f(x)|^2, \quad \int_{\mathbb{R}^d} W_f^\varepsilon(x, \xi) dx = \frac{1}{(2\pi\varepsilon)^d} \left| \widehat{f}\left(\frac{\xi}{\varepsilon}\right) \right|^2.$$

Note that despite this, W_f^ε is not positive in general.

- For every $a \in \mathcal{C}_c^\infty(\mathbb{R}^d \times \mathbb{R}^d)$ one has:

$$(3.1) \quad \int_{\mathbb{R}^d \times \mathbb{R}^d} a(x, \xi) W_f^\varepsilon(x, \xi) dx d\xi = (\text{op}_\varepsilon(a)f, f)_{L^2(\mathbb{R}^d)},$$

where $\text{op}_\varepsilon(a)$ is the semiclassical pseudodifferential operator of symbol a obtained through the Weyl quantization rule:

$$\text{op}_\varepsilon(a)f(x) = \int_{\mathbb{R}^d \times \mathbb{R}^d} a\left(\frac{x+y}{2}, \varepsilon\xi\right) e^{i\xi \cdot (x-y)} f(y) dy \frac{d\xi}{(2\pi)^d}.$$

If (f^ε) is a bounded sequence in $L^2(\mathbb{R}^d)$ then $(W_{f^\varepsilon}^\varepsilon)$ is a bounded sequence of tempered distributions in $\mathcal{S}'(\mathbb{R}^d \times \mathbb{R}^d)$. This is proved using identity (3.1) combined with the fact that the operators $\text{op}_\varepsilon(a)$ are uniformly bounded by a suitable semi-norm in $\mathcal{S}(\mathbb{R}^d \times \mathbb{R}^d)$, see (A.1). Appendix A contains additional facts on the theory of pseudodifferential operators, as well as references to the literature.

In addition, every accumulation point of $(W_{f^\varepsilon}^\varepsilon)$ in $\mathcal{S}'(\mathbb{R}^d \times \mathbb{R}^d)$ is a positive distribution and therefore, by Schwartz's theorem, a positive measure on $\mathbb{R}^d \times \mathbb{R}^d$. These measures are called *semiclassical* or *Wigner measures*. See references [25, 36, 27, 28] for different proofs of the results we have presented so far.

Now, if $\mu \in \mathcal{M}_+(\mathbb{R}^d \times \mathbb{R}^d)$ is an accumulation point of $(W_{f^\varepsilon}^\varepsilon)$ along some subsequence (ε_k) and $(|f^{\varepsilon_k}|^2)$ converges weakly- \star towards a measure $\nu \in \mathcal{M}_+(\mathbb{R}^d)$ then one has:

$$(3.2) \quad \int_{\mathbb{R}^d} \mu(\cdot, d\xi) \leq \nu.$$

Equality holds if and only if (f^ε) is ε -oscillating:

$$(3.3) \quad \limsup_{\varepsilon \rightarrow 0^+} \int_{|\xi| > R/\varepsilon} |\widehat{f^\varepsilon}(\xi)|^2 d\xi \xrightarrow{R \rightarrow +\infty} 0,$$

see [25, 27, 28]. The hypothesis **H0** that we made on the initial data for equation (2.1), is this ε -oscillating property. Note also that (3.2) implies that μ is always a finite measure of total mass bounded by $\sup_\varepsilon \|f^\varepsilon\|_{L^2(\mathbb{R}^d)}^2$.

Remark 3.1. If $\|\langle \varepsilon D_x \rangle^s f^\varepsilon\|_{L^2(\mathbb{R}^d)}$ is uniformly bounded for some constant $s > 0$, then the family f^ε is ε -oscillating.

3.2. Wigner measure and family of solutions of dispersive equations. We will now consider Wigner distributions associated to solutions of the evolution equation (2.1) where V_{ext} and λ satisfy hypothesis **H1** and $(g^\varepsilon(t, \cdot))$ is locally uniformly bounded with respect to t in $L^2(\mathbb{R}^d)$.

When the sequence (u_0^ε) of initial data is uniformly bounded in $L^2(\mathbb{R}^d)$, so is the corresponding sequence $(u^\varepsilon(t, \cdot))$ of solutions to (2.1) for every $t \in \mathbb{R}$. Therefore the sequence of Wigner distributions $(W_{u^\varepsilon(t, \cdot)}^\varepsilon)$ is bounded in $\mathcal{C}(\mathbb{R}; \mathcal{S}'(\mathbb{R}^d \times \mathbb{R}^d))$. Nevertheless, its time derivatives are unbounded and, in general, one cannot hope to find a subsequence that converges pointwise (or even almost everywhere) in t (see Proposition 3.4 below). This difficulty can be overcome if one considers the time-average of the Wigner distributions.

Proposition 3.2. *Let (u^ε) be a sequence of solutions to (2.1) issued from an $L^2(\mathbb{R}^d)$ -bounded family of initial data (u_0^ε) . Then there exist a subsequence (ε_k) tending to zero as $k \rightarrow \infty$ and a t -measurable family $\mu_t \in \mathcal{M}_+(\mathbb{R}^d \times \mathbb{R}^d)$ of finite measures, with total mass essentially uniformly bounded in $t \in \mathbb{R}$, such that, for every $\theta \in L^1(\mathbb{R})$ and $a \in \mathcal{C}_c^\infty(\mathbb{R}^d \times \mathbb{R}^d)$:*

$$\lim_{k \rightarrow \infty} \int_{\mathbb{R} \times \mathbb{R}^d \times \mathbb{R}^d} \theta(t) a(x, \xi) W_{u^{\varepsilon_k}(t, \cdot)}^{\varepsilon_k}(x, \xi) dx d\xi dt = \int_{\mathbb{R} \times \mathbb{R}^d \times \mathbb{R}^d} \theta(t) a(x, \xi) \mu_t(dx, d\xi) dt.$$

If moreover, the families (u_0^ε) and $g^\varepsilon(t, \cdot)$ are ε -oscillating, then for every $\theta \in L^1(\mathbb{R})$ and $\phi \in \mathcal{C}_c^\infty(\mathbb{R}^d)$:

$$\lim_{k \rightarrow \infty} \int_{\mathbb{R}} \int_{\mathbb{R}^d} \theta(t) \phi(x) |u^{\varepsilon_k}(t, x)|^2 dx dt = \int_{\mathbb{R}} \int_{\mathbb{R}^d \times \mathbb{R}^d} \theta(t) \phi(x) \mu_t(dx, d\xi) dt.$$

This result is proved in [38], Theorem 1; see also Appendix B in [42]. Note that its proof uses the following observation.

Remark 3.3. Let $(u^\varepsilon(t, \cdot))$ be a sequence of solutions to (2.1) with ε -oscillating sequence of initial data (u_0^ε) and assume $g^\varepsilon(t, \cdot)$ is ε -oscillating for all time $t \in \mathbb{R}$. Then, $u^\varepsilon(t, \cdot)$ also is ε -oscillating for all $t \in \mathbb{R}$.

3.3. Localisation of Wigner measures on the critical set. The fact that $(u^\varepsilon(t, \cdot))$ is a sequence of solutions to (2.1) imposes restrictions on the measures μ_t that can be attained as a limit of their Wigner functions. In the region in the phase space $\mathbb{R}_x^d \times \mathbb{R}_\xi^d$ where equation (2.1) is dispersive (*i.e.* away from the critical points of λ) the energy of the sequence $(u^\varepsilon(t, \cdot))$ is dispersed at infinite speed to infinity. These heuristics are made precise in the following result.

Proposition 3.4. *Let $(u^\varepsilon(t, \cdot))$ be a sequence of solutions to (2.1) issued from an $L^2(\mathbb{R}^d)$ -bounded and ε -oscillating sequence of initial data (u_0^ε) , and suppose that the measures μ_t are given by Proposition 3.2. Then, for almost every $t \in \mathbb{R}$ the measure μ_t is supported above the set of critical points of λ :*

$$\text{supp } \mu_t \subset \Lambda = \{(x, \xi) \in \mathbb{R}^d \times \mathbb{R}^d : \nabla \lambda(\xi) = 0\}.$$

The result of Proposition 3.4 follows from a geometric argument : the fact that u^ε are solutions to (2.1) translates in an invariance property of the measures μ_t .

Lemma 3.5. *For almost every $t \in \mathbb{R}$, the measure μ_t is invariant by the flow*

$$\phi_s^1 : \mathbb{R}^d \times \mathbb{R}^d \ni (x, \xi) \mapsto (x + s\nabla \lambda(\xi), \xi) \in \mathbb{R}^d \times \mathbb{R}^d, \quad s \in \mathbb{R}.$$

This means that for every function a on $\mathbb{R}^d \times \mathbb{R}^d$ that is Borel measurable one has:

$$\int_{\mathbb{R}^d \times \mathbb{R}^d} a \circ \phi_s^1(x, \xi) \mu_t(dx, d\xi) = \int_{\mathbb{R}^d \times \mathbb{R}^d} a(x, \xi) \mu_t(dx, d\xi), \quad s \in \mathbb{R}.$$

This result is part of Theorem 2 in [38]. We reproduce the argument here for the reader's convenience, since we are going to use similar techniques in the sequel.

Proof of Lemma 3.5. It is enough to show that, for all $a \in C_c^\infty(\mathbb{R}^d \times \mathbb{R}^d)$ and $\theta \in C_c^\infty(\mathbb{R})$, the quantity

$$R^\varepsilon(\theta, a) := \int_{\mathbb{R} \times \mathbb{R}^d \times \mathbb{R}^d} \theta(t) \frac{d}{ds} (a \circ \phi_s^1(x, \xi)) \Big|_{s=0} W_{u^{\varepsilon_k}(t, \cdot)}^{\varepsilon_k}(x, \xi) dx d\xi dt$$

tends to 0 for the subsequence ε_k of Proposition 3.2. Note that

$$\frac{d}{ds} (a \circ \phi_s^1) \Big|_{s=0} = \nabla_\xi \lambda \cdot \nabla_x a = \{\lambda, a\};$$

therefore, by the symbolic calculus of semiclassical pseudodifferential operators, Proposition A.1:

$$\text{op}_\varepsilon \left(\frac{d}{ds} (a \circ \phi_s^1) \Big|_{s=0} \right) = \frac{i}{\varepsilon} [\lambda(\varepsilon D), \text{op}_\varepsilon(a)] + O_{\mathcal{L}(L^2(\mathbb{R}^d))}(\varepsilon)$$

and, using the fact that u^ε solves (2.1):

$$\frac{i}{\varepsilon} \int_{\mathbb{R}} \theta(t) ([\lambda(\varepsilon D), \text{op}_\varepsilon(a)] u^\varepsilon(t, \cdot), u^\varepsilon(t, \cdot)) dt + O(\varepsilon)$$

$$= -\varepsilon \int_{\mathbb{R}} \theta(t) \frac{d}{dt} (\text{op}_\varepsilon(a) u^\varepsilon(t, \cdot), u^\varepsilon(t, \cdot)) dt = \varepsilon \int_{\mathbb{R}} \theta'(t) (\text{op}_\varepsilon(a) u^\varepsilon(t, \cdot), u^\varepsilon(t, \cdot)) dt = O(\varepsilon).$$

This estimate together with identity (3.1) show that $R^\varepsilon(\theta, a) = O(\varepsilon)$, which gives the result that we wanted to prove. \square

Proposition 3.4 follows easily from Lemma 3.5 and the following elementary fact.

Lemma 3.6. *Let $\Omega \subset \mathbb{R}^d$ and $\Phi_s : \mathbb{R}^d \times \Omega \rightarrow \mathbb{R}^d \times \Omega$ a flow satisfying: for every compact $K \subset \mathbb{R}^d \times \Omega$ such that K contains no stationary points of Φ there exist constants $\alpha, \beta > 0$ such that:*

$$\alpha|s| - \beta \leq |\Phi_s(x, \xi)| \leq \alpha|s| + \beta, \quad \forall (x, \xi) \in K.$$

Let μ be a finite, positive Radon measure on $\mathbb{R}^d \times \Omega$ that is invariant by the flow Φ_s . Then μ is supported on the set of stationary points of Φ_s .

Proof. It suffices to show that $\mu(K) = 0$ for every compact set $K \subset \mathbb{R}^d \times \Omega$ as in the statement of the lemma. By the assumption made on Φ_s , it is possible to find $s_k \rightarrow +\infty$ such that $\Phi_{s_k}(K)$, $k \in \mathbb{N}$, are mutually disjoint. The invariance property of μ implies that $\mu(\Phi_{s_k}(K)) = \mu(K)$ and therefore, for every $N > 0$:

$$\mu \left(\bigcup_{k=1}^N \Phi_{s_k}(K) \right) = N\mu(K).$$

Since μ is finite, we must have $\mu(K) = 0$. \square

4. TWO-MICROLOCAL WIGNER DISTRIBUTIONS

The localization result for semiclassical measures that we obtained in the preceding section is still very far from the conclusions of Theorems 2.1 and 2.2. In particular, Proposition 3.4 does not explain how the measures μ_t depend on the sequence of initial data of the sequence of solutions $(u^\varepsilon(t, \cdot))$. For obtaining more information, we use two-microlocal tools that we introduce in a rather general framework in this section.

From now on, we assume that X is a connected, closed embedded submanifold of $(\mathbb{R}^d)^*$ with codimension $p > 0$. Given any $\sigma \in X$, $T_\sigma X$ and $N_\sigma X$ will stand for the cotangent and normal spaces of X at σ respectively (as defined in (1.9) and (1.10)). The tubular neighborhood theorem (see for instance [29]) ensures that there exists an open neighborhood U of $\{(\sigma, 0) : \sigma \in X\} \subseteq NX$ such that the map:

$$U \ni (\sigma, v) \mapsto \sigma + v \in (\mathbb{R}^d)^*,$$

is a diffeomorphism onto its image V . Its inverse is given by:

$$V \ni \xi \mapsto (\sigma(\xi), \xi - \sigma(\xi)) \in U,$$

for some smooth map $\sigma : V \rightarrow X$. When $X = \{\xi_0\}$ consists of a single point, the function σ is constant, identically equal to ξ_0 .

We extend the phase space $T^*\mathbb{R}^d := \mathbb{R}_x^d \times (\mathbb{R}^d)_\xi^*$ with a new variable $\eta \in \overline{\mathbb{R}^d}$, where $\overline{\mathbb{R}^d}$ is the compactification of \mathbb{R}^d obtained by adding a sphere \mathbf{S}^{d-1} at infinity. The test functions associated with this extended phase space are functions $a \in C^\infty(T^*\mathbb{R}_{x,\xi}^d \times \mathbb{R}_\eta^d)$ which satisfy the two following properties:

- (1) There exists a compact $K \subset T^*\mathbb{R}^d$ such that, for all $\eta \in \mathbb{R}^p$, the map $(x, \xi) \mapsto a(x, \xi, \eta)$ is a smooth function compactly supported in K .
- (2) There exists a smooth function a_∞ defined on $T^*\mathbb{R}^d \times \mathbf{S}^{d-1}$ and $R_0 > 0$ such that, if $|\eta| > R_0$, then $a(x, \xi, \eta) = a_\infty(x, \xi, \eta/|\eta|)$.

We denote by \mathcal{A} the set of such functions and for $a \in \mathcal{A}$ we write:

$$(4.1) \quad a_\varepsilon(x, \xi) := a\left(x, \xi, \frac{\xi - \sigma(\xi)}{\varepsilon}\right).$$

Given $f \in L^2(\mathbb{R}^d)$, we define the two-microlocal Wigner distribution $W_f^{X, \varepsilon} \in \mathcal{D}'(\mathbb{R}^d \times V \times \overline{\mathbb{R}^d})$ by:

$$(4.2) \quad \langle W_f^{X, \varepsilon}, a \rangle := (\text{op}_\varepsilon(a_\varepsilon)f|f)_{L^2(\mathbb{R}^d)}, \quad \forall a \in \mathcal{A}.$$

Since $a_\varepsilon(x, \varepsilon\xi) = a\left(x, \varepsilon\xi, \frac{\varepsilon\xi - \sigma(\varepsilon\xi)}{\varepsilon}\right)$ has derivatives that are uniformly bounded in ε , the Calderón-Vaillancourt theorem (see Appendix A) gives the uniform boundedness of the family of operators $(\text{op}_\varepsilon(a_\varepsilon))_{\varepsilon>0}$ in $L^2(\mathbb{R}^d)$. In addition, any function $a \in \mathcal{C}_c^\infty(\mathbb{R}^d \times V)$ can be naturally identified to a function in \mathcal{A} which does not depend on the last variable. For such a , one clearly has

$$\langle W_f^{X, \varepsilon}, a \rangle = \int_{\mathbb{R}^d \times \mathbb{R}^d} a(x, \xi) W_f^\varepsilon(x, \xi) dx d\xi.$$

Putting the above remarks together, one obtains the following.

Proposition 4.1. *Let $(f^\varepsilon)_{\varepsilon>0}$ be bounded in $L^2(\mathbb{R}^d)$; suppose in addition that this sequence has a semiclassical measure μ . Then, $(W_{f^\varepsilon}^{X, \varepsilon})_{\varepsilon>0}$ is a bounded sequence in $\mathcal{D}'(\mathbb{R}^d \times V \times \overline{\mathbb{R}^d})$ whose accumulation points μ^X satisfy:*

$$\langle \mu^X, a \rangle = \int_{\mathbb{R}^d \times \mathbb{R}^d} a(x, \xi) \mu(dx, d\xi), \quad \forall a \in \mathcal{C}_c^\infty(\mathbb{R}^d \times V).$$

The distributions μ^X turn out to have additional structure (they are not positive measures on $\mathbb{R}^d \times V \times \overline{\mathbb{R}^d}$, though) and can be used to give a more precise description of the restriction $\mu|_{\mathbb{R}^d \times X}$ of semiclassical measures. The measure μ^X decomposes into two parts: a *compact* part, which is essentially the restriction of μ^X to the interior $\mathbb{R}^d \times V \times \mathbb{R}^d$ of $\mathbb{R}^d \times V \times \overline{\mathbb{R}^d}$, and a *part at infinity*, which corresponds to the restriction to the sphere at infinity $\mathbb{R}^d \times V \times \mathbf{S}^{d-1}$.

4.1. The compact part. On the neighborhood of any point $\sigma \in X$, one may find a system of p equations on \mathbb{R}^d for which X is the zero set. Let $\varphi(\xi) = 0$ be such a system in an open set Ω that we can assume included in the set V where the map σ is defined. Then, a parametrization of $N_\sigma X$ associated to this system of equations is

$$N_\sigma X = \{ {}^t d\varphi(\sigma)z, \quad z \in \mathbb{R}^p \}.$$

For $\sigma \in X$, we define functions of $L^2(N_\sigma X)$ as square integrable functions

$$\mathbb{R}^p \ni z \mapsto f(z),$$

where z is the parameter of a parametrization of $N_\sigma X$ that we fixed *a priori*.

Besides, one associates with the system $\varphi(\xi) = 0$ a smooth map $\xi \mapsto B(\xi)$ from the neighborhood Ω of σ into the set of $d \times p$ matrices such that

$$(4.3) \quad \xi - \sigma(\xi) = B(\xi)\varphi(\xi), \quad \xi \in \Omega.$$

Given a function $a \in \mathcal{C}_c^\infty(\mathbb{R}^d \times \Omega \times \mathbb{R}^d)$ and a point $(\sigma, v) \in TX$, we can use the system of coordinates $\varphi(\xi) = 0$ to define an operator acting on $f \in L^2(N_\sigma X)$ given by:

$$Q_a^\varphi(\sigma, v)f(z) = \int_{\mathbb{R}^p \times \mathbb{R}^p} a\left(v + {}^t d\varphi(\sigma) \frac{z+y}{2}, \sigma, B(\sigma)\eta\right) f(y) e^{i\eta \cdot (z-y)} \frac{d\eta dy}{(2\pi)^p}.$$

In other words, $Q_a^\varphi(\sigma, v)$ is obtained from a by applying the non-semiclassical Weyl quantization to the symbol

$$(z, \eta) \mapsto a\left(v + {}^t d\varphi(\sigma)z, \sigma, B(\sigma)\eta\right) \in \mathcal{C}_c^\infty(\mathbb{R}^p \times \mathbb{R}^p).$$

We write

$$Q_a^\varphi(\sigma, v) = a^W\left(v + {}^t d\varphi(\sigma)z, \sigma, B(\sigma)D_z\right).$$

If one changes the system of coordinates into $\tilde{\varphi}(\xi) = 0$ on some open neighborhood $\tilde{\Omega}$ of σ , then, there exists a smooth map $R(\xi)$ defined on the open set $\Omega \cap \tilde{\Omega}$ (where both system of coordinates can be used), and valued in the set of invertible $p \times p$ matrices, such that $\tilde{\varphi}(\xi) = R(\xi)\varphi(\xi)$. One then observe that the matrix $\tilde{B}(\xi)$ associated with the choice of $\tilde{\varphi}$ is given by $\tilde{B}(\xi) = B(\xi)R(\xi)^{-1}$. Besides, for $a \in \mathcal{C}_c^\infty(\mathbb{R}^d \times (\Omega \cap \tilde{\Omega}) \times \mathbb{R}^d)$,

$$\begin{aligned} Q_a^{\tilde{\varphi}}(\sigma, v) &= \int_{\mathbb{R}^p \times \mathbb{R}^p} a\left(v + {}^t d\tilde{\varphi}(\sigma) \frac{z+y}{2}, \sigma, \tilde{B}(\sigma)\eta\right) f(y) e^{i\eta \cdot (z-y)} \frac{d\eta dy}{(2\pi)^p} \\ &= \int_{\mathbb{R}^p \times \mathbb{R}^p} a\left(v + {}^t d\varphi(\sigma) {}^t R(\sigma) \frac{z+y}{2}, \sigma, B(\sigma)R(\sigma)^{-1}\eta\right) f(y) e^{i\eta \cdot (z-y)} \frac{d\eta dy}{(2\pi)^p}. \end{aligned}$$

We obtain

$$Q_a^{\tilde{\varphi}}(\sigma, v) = U(\sigma)Q_a^\varphi(\sigma, v)U^*(\sigma),$$

where $U(\sigma)$ is the unitary operator of $L^2(N_\sigma X) \sim L^2(\mathbb{R}^p)$ associated with the linear map from \mathbb{R}^p into itself: $z \mapsto {}^t R(\sigma)z$. More precisely,

$$\forall f \in L^2(\mathbb{R}^p), \quad U(\sigma)f(z) = |\det R(\sigma)|^{\frac{p}{2}} f({}^t R(\sigma)z).$$

This map is the one associated with the change of parametrization on $N_\sigma X$ induced by turning φ into $\tilde{\varphi}$, and the map $(z, \zeta) \mapsto ({}^t R(\sigma)z, R(\sigma)^{-1}\zeta)$ is a symplectic transform of the cotangent of \mathbb{R}^p . This is the standard rule of transformation of pseudodifferential operators through linear change of variables (see [1] for an example or any textbook about pseudodifferential calculus).

Because of this invariance property with respect to the change of system of coordinates, we shall say that a defines an operator $Q_a(\sigma, v)$ on $L^2(N_\sigma X)$. Clearly, $Q_a(\sigma, v)$ is smooth and compactly supported in (σ, v) ; moreover, $Q_a(\sigma, v) \in \mathcal{K}(L^2(N_\sigma X))$, for every $(\sigma, v) \in TX$, where $\mathcal{K}(L^2(N_\sigma X))$ stands for the space of compact operators on $L^2(N_\sigma X)$.

Proposition 4.2. *Let μ^X be given by Proposition 4.1. Then there exist a positive measure ν on T^*X and a measurable family:*

$$M : T^*X \ni (\sigma, v) \mapsto M(\sigma, v) \in \mathcal{L}_+^1(L^2(N_\sigma X)),$$

satisfying

$$\mathrm{Tr}_{L^2(N_\sigma X)} M(\sigma, v) = 1, \quad \text{for } \nu\text{-a.e. } (\sigma, v) \in T^*X,$$

and such that, for every $a \in C_c^\infty(\mathbb{R}^d \times V \times \mathbb{R}^d)$ one has:

$$\langle \mu^X, a \rangle = \int_{T^*X} \text{Tr}_{L^2(N_{\sigma X})}(Q_a(\sigma, v)M(\sigma, v))\nu(d\sigma, dv).$$

Proof. We suppose that we are given a local system of p equations of X by $\varphi(\xi) = 0$. Put $\xi = (\xi', \xi'') \in \mathbb{R}^p \times \mathbb{R}^{d-p}$. Without loss of generality, we may assume that $d_{\xi'}\varphi(\xi)$ is invertible. We consider the smooth valued function B satisfying $\xi - \sigma(\xi) = B(\xi)\varphi(\xi)$ and we introduce the local diffeomorphism

$$\Phi : (\varphi(\xi), \xi'') \mapsto \xi.$$

Note that if $\xi = \Phi(\zeta)$, $\zeta = (\zeta', \zeta'')$, we have $\zeta' = \varphi(\xi) = \varphi(\Phi(\zeta))$ and $\zeta'' = \xi''$. We use this diffeomorphism according to the next lemma.

Lemma 4.3. *For all $f \in L^2(\mathbb{R}^d)$ and $a \in \mathcal{A}$,*

$$(\text{op}_\varepsilon(a_\varepsilon)f, f) = \left(\text{op}_\varepsilon \left(a \left({}^t d\Phi(\xi)^{-1}x, \Phi(\xi), B(\Phi(\xi)) \frac{\xi'}{\varepsilon} \right) \right) \mathcal{U}_\varepsilon f, \mathcal{U}_\varepsilon f \right) + O(\varepsilon)\|f\|^2$$

where $f \mapsto \mathcal{U}_\varepsilon f$ is an isometry of $L^2(\mathbb{R}^d)$.

The proof of this lemma is in the Appendix C. This lemma reduces the problem to the analysis of the concentration of the bounded family $\tilde{f}^\varepsilon = (\mathcal{U}_\varepsilon f)$ on the submanifold $\Lambda_0 = \{\xi' = 0\}$ which has the additional property to be a vector space. This special case has been studied in [15] where it is proved (see pages 96-97, Proposition 2) that up to a subsequence, there exist a positive measure ν_0 on $T^*\mathbb{R}^{d-p}$ and a measurable family of trace 1 operators:

$$M_0 : T^*\mathbb{R}^{d-p} \ni (\sigma, v) \mapsto M_0(\sigma, v) \in \mathcal{L}_+^1(L^2(\mathbb{R}^p)),$$

satisfying for any $b \in C_c^\infty(\mathbb{R}^{2d+p})$,

$$\begin{aligned} & \lim_{\varepsilon \rightarrow 0} \left(\text{op}_\varepsilon(b_\varepsilon)\tilde{f}^\varepsilon, \tilde{f}^\varepsilon \right) \\ &= \int_{\mathbb{R}^{d-p} \times \mathbb{R}^{d-p}} \text{Tr}_{L^2(\mathbb{R}^p)} \left(b^W((z, u''), (0, \theta''), D_z) M_0(u'', \theta'') \right) d\nu_0(du'', d\theta''). \end{aligned}$$

The reader will find in Appendix B comments on the operator-valued families. Therefore, for compactly supported $a \in \mathcal{A}$, and choosing $b(x, \xi, \eta') = a({}^t d\Phi(\xi)^{-1}x, \Phi(\xi), B(\Phi(\xi))\eta')$, one obtains

$$\begin{aligned} & \lim_{\varepsilon \rightarrow 0} (\text{op}_\varepsilon(a_\varepsilon)f^\varepsilon, f^\varepsilon) \\ &= \int_{\mathbb{R}^{d-p} \times \mathbb{R}^{d-p}} \text{Tr}_{L^2(\mathbb{R}^p)} \left(a^W({}^t d\Phi(0, \theta'')^{-1}(z, u''), \Phi(0, \theta''), B(\Phi(0, \theta''))D_z) \right. \\ & \quad \left. \times M_0(u'', \theta'') \right) d\nu_0(du'', d\theta''). \end{aligned}$$

Note that the map $\theta'' \mapsto \sigma = \Phi(0, \theta'')$ is a parametrization of X with associated parametrization of T^*X ,

$$(\theta'', u'') \mapsto (\sigma, v) = (\Phi(0, \theta''), {}^t d\Phi(0, \theta'')^{-1}(0, u'')).$$

Since the Jacobian of this mapping is 1, after a change of variable, we obtain an operator valued measurable family M on T^*X and a measure ν on T^*X such that

$$\begin{aligned} & \lim_{\varepsilon \rightarrow 0} (\text{op}_\varepsilon(a_\varepsilon)f, f) \\ &= \int_{T^*X} \text{Tr}_{L^2(\mathbb{R}^p)} (a^W ({}^t d\Phi(0, \theta''(\sigma))^{-1}(z, 0) + v, \sigma, B(\sigma)D_z) M(\sigma, v)) \, d\nu(d\sigma, dv). \end{aligned}$$

We now take advantage of the fact that $\varphi(\Phi(\zeta)) = \zeta'$ for all $\zeta \in \mathbb{R}^d$ in order to write

$$d\varphi(\Phi(\zeta))d\Phi(\zeta) = (\text{Id}, 0).$$

We deduce

$$\forall z \in \mathbb{R}^p, \quad {}^t d\Phi(\zeta) {}^t d\varphi(\Phi(\zeta))z = (z, 0),$$

which implies

$$\forall z \in \mathbb{R}^p, \quad {}^t d\varphi(\Phi(\zeta))z = {}^t d\Phi(\zeta)^{-1}(z, 0).$$

Therefore,

$$\begin{aligned} \lim_{\varepsilon \rightarrow 0} (\text{op}_\varepsilon(a_\varepsilon)f, f) &= \int_{T^*X} \text{Tr}_{L^2(\mathbb{R}^p)} (a^W ({}^t d\varphi(\sigma)z + v, \sigma, B(\sigma)D_z) M(\sigma, v)) \, d\nu(d\sigma, dv) \\ &= \int_{T^*X} \text{Tr}_{L^2(N_\sigma X)} (Q_a(\sigma)M(\sigma, v)) \, d\nu(d\sigma, dv). \end{aligned}$$

□

4.2. Measure structure of the part at infinity. To analyze the part at infinity, we use a cut-off function $\chi \in C_c^\infty(\mathbb{R}^d)$ such that $0 \leq \chi \leq 1$, $\chi(\eta) = 1$ for $|\eta| \leq 1$ and $\chi(\eta) = 0$ for $|\eta| \geq 2$, and we write

$$\left\langle W_f^{X,\varepsilon}, a \right\rangle = \left\langle W_f^{X,\varepsilon}, a_R \right\rangle + \left\langle W_f^{X,\varepsilon}, a^R \right\rangle,$$

with

$$(4.4) \quad a_R(x, \xi, \eta) := a(x, \xi, \eta) \chi\left(\frac{\eta}{R}\right) \quad \text{and} \quad a^R(x, \xi, \eta) := a(x, \xi, \eta) \left(1 - \chi\left(\frac{\eta}{R}\right)\right).$$

Observe that a_R is compactly supported in all variables. We thus focus on the second part, and more precisely on the quantity

$$\limsup_{R \rightarrow \infty} \limsup_{\varepsilon \rightarrow 0^+} \left\langle W_f^{X,\varepsilon}, a^R \right\rangle.$$

We denote by $S\Lambda$ the compactified normal bundle to Λ , viewed as a submanifold of $\mathbb{R}^d \times \mathbb{R}^d$, the fiber of which is $T_\sigma^* \mathbb{R}^d \times S_\sigma \Lambda$ above σ with $S_\sigma \Lambda$ being obtained by taking the quotient of $N_\sigma \Lambda$ by the action of \mathbb{R}_+^* by homotheties.

Proposition 4.4. *Let (f^ε) be a bounded family of $L^2(\mathbb{R}^d)$. There exists a subsequence ε_k and a measure γ on $S\Lambda$ such that for all $a \in \mathcal{A}$,*

$$\begin{aligned} \lim_{R \rightarrow \infty} \lim_{k \rightarrow +\infty} \left\langle W_{f^{\varepsilon_k}}^{X,\varepsilon_k}, a^R \right\rangle &= \int_{\mathbb{R}^d \times X \times \mathbf{S}^{d-1}} a_\infty(x, \sigma, \omega) \gamma(dx, d\sigma, d\omega) \\ &\quad + \int_{\mathbb{R}^d \times X^c \times \mathbf{S}^{d-1}} a_\infty\left(x, \xi, \frac{\xi - \sigma(\xi)}{|\xi - \sigma(\xi)|}\right) \mu(dx, d\xi), \end{aligned}$$

where X^c denotes the complement of the set X in \mathbb{R}^d .

Proof. We begin by recalling the arguments that prove the existence of the measure γ , which are the same that the one developed in the vector case in [15]. Since $a = a_\infty$ for $|\eta|$ large enough, we have $a^R = a_\infty^R$ as soon as R is large enough and the quantity

$$\limsup_{R \rightarrow \infty} \limsup_{\varepsilon \rightarrow 0^+} \left\langle W_{f^\varepsilon}^{X,\varepsilon}, a^R \right\rangle$$

will only depend on a_∞ . Therefore, by considering a dense subset of $\mathcal{C}_c(T^*\mathbb{R}^d \times \mathbf{S}^{d-1})$, we can find a subsequence (ε_k) by a diagonal extraction process such that the following linear form on $\mathcal{C}_c(T^*\mathbb{R}^d \times \mathbf{S}^{d-1})$ is well-defined

$$\ell : a_\infty \mapsto \lim_{R \rightarrow \infty} \lim_{k \rightarrow +\infty} \left\langle W_{f^{\varepsilon_k}}^{X,\varepsilon_k}, a^R \right\rangle.$$

We then observe that

$$\forall \alpha, \beta \in \mathbb{N}^d, \exists C_{\alpha,\beta} > 0, \sup_{\mathbb{R}^{2d}} \left| \partial_x^\alpha \partial_\xi^\beta (a^R)_\varepsilon \right| \leq C_{\alpha,\beta} (\varepsilon^{|\beta|} + R^{-|\beta|}).$$

This implies that the symbolic calculus on symbols $(a^R)_\varepsilon$ is semiclassical with respect to the small parameter $\sqrt{\varepsilon^2 + R^{-2}}$. To be precise, one has the following weak Gårding inequality: if $a \geq 0$, then, for all $\kappa > 0$, there exists a constant C_κ such that

$$\left\langle W_{f^\varepsilon}^{X,\varepsilon}, a^R \right\rangle \geq - \left(\kappa + C_\kappa \left(\varepsilon + \frac{1}{R} \right) \right) \|f^\varepsilon\|_{L^2(\mathbb{R}^d)}^2.$$

We then conclude that the linear form ℓ defined above is positive and defines a positive Radon measure $\tilde{\rho}$. It remains to compute $\tilde{\rho}$ outside X . In this purpose, we set

$$a^R = a_\delta^R + a^{R,\delta} \quad \text{with} \quad a_\delta^R(x, \xi, \eta) = a^R(x, \xi, \eta) (1 - \chi) \left(\frac{\xi - \sigma(\xi)}{\delta} \right)$$

and we observe that, by the definition of μ :

$$\lim_{\delta \rightarrow 0} \lim_{R \rightarrow \infty} \lim_{\varepsilon \rightarrow 0} \left\langle W_{f^{\varepsilon_k}}^{X,\varepsilon_k}, a_\delta^R \right\rangle = \int_{\mathbb{R}^d \times X^c \times \mathbf{S}^{d-1}} a_\infty \left(x, \xi, \frac{\xi - \sigma(\xi)}{|\xi - \sigma(\xi)|} \right) \mu(dx, d\xi),$$

which concludes the proof of the existence of the measure γ .

Let us now analyze the geometric properties of this measure. We choose a system of local coordinates of Λ and introduce the matrix B as in (4.3). By Lemma 4.3 and the result of [15] for vector spaces: up to a subsequence, there exists a measure $\tilde{\gamma}_0$ on $\mathbb{R}^d \times \mathbb{R}^{d-p} \times \mathbf{S}^{p-1}$ such that

$$\begin{aligned} & \lim_{\delta \rightarrow 0^+} \lim_{R \rightarrow \infty} \lim_{\varepsilon \rightarrow 0^+} \left\langle W_f^{X,\varepsilon}, a^{R,\delta} \right\rangle \\ &= \int_{\mathbb{R}^d \times \mathbb{R}^{d-p} \times \mathbf{S}^{p-1}} a_\infty \left({}^t d\Phi(0, \xi'')^{-1} x, \Phi(0, \xi''), \frac{B(\Phi(0, \xi'')) \omega}{|B(\Phi(0, \xi'')) \omega|} \right) \tilde{\gamma}_0(dx, d\xi, d\omega). \end{aligned}$$

The mapping $\xi'' \mapsto \Phi(0, \xi'')$ is a parametrization of X and the mapping

$$(x, \xi) \mapsto ({}^t d\Phi(0, \xi'')^{-1} x, \Phi(0, \xi''))$$

is the associated mapping of $T_X^* \mathbb{R}^d$. Therefore, this relation defines a measure $\tilde{\gamma}$ on $T^*X \times \mathbf{S}^{p-1}$ such that

$$(4.5) \quad \lim_{\delta \rightarrow 0^+} \lim_{R \rightarrow \infty} \lim_{\varepsilon \rightarrow 0^+} \left\langle W_f^{X,\varepsilon}, a^{R,\delta} \right\rangle = \int_{T^*X \times \mathbf{S}^{p-1}} a_\infty \left(x, \sigma, \frac{B(\sigma) \omega}{|B(\sigma) \omega|} \right) \tilde{\gamma}(dx, d\xi, d\omega).$$

Besides, using that

$$(4.6) \quad \text{Id} = d\sigma(\sigma_0) + B(\sigma_0)d\varphi(\sigma_0)$$

for any $\sigma_0 \in X$, we deduce that for any $\zeta \in T_{\sigma_0}\mathbb{R}^d$, we have the decomposition

$$\zeta = d\sigma(\sigma_0)\zeta + B(\sigma_0)d\varphi(\sigma_0)\zeta, \quad \text{with } d\sigma(\sigma_0)\zeta \in T_{\sigma_0}X \quad \text{and } B(\sigma_0)d\varphi(\sigma_0)\zeta \in N_{\sigma_0}X.$$

Now, since $d\varphi$ is of rank p , one can write any $\omega \in \mathbf{S}^{p-1}$ as $\omega = d\varphi(\sigma_0)\zeta$ and the points $B(\sigma_0)\omega$ are in $N_{\sigma_0}X$. By identification of γ in (4.5), we deduce that $\gamma(x, \sigma, \cdot)$ is a measure on the set

$$\left\{ \frac{B(\sigma)\omega}{|B(\sigma)\omega|}, \quad \omega \in \mathbf{S}^{p-1} \right\} = N_{\sigma}X / \mathbb{R}_+^* = S_{\sigma}X,$$

which completes the proof of the proposition. \square

5. TWO MICROLOCAL WIGNER MEASURES AND FAMILIES OF SOLUTIONS TO DISPERSIVE EQUATIONS

We now consider families of solutions to equation (2.1). As proved in Proposition 3.4, the Wigner measure of the family $(u^\varepsilon(t, \cdot))$ concentrates on the set $\Lambda = \{\nabla\lambda(\xi) = 0\}$. In order to analyze μ^t above Λ , we perform a second microlocalization above the set $X = \Lambda$, with average in time. We consider for $\theta \in L^1(\mathbb{R})$ the quantities

$$\int_{\mathbb{R}} \theta(t) \left\langle W_{u^\varepsilon(t, \cdot)}^{\Lambda, \varepsilon}, a \right\rangle dt$$

for symbols $a \in \mathcal{A}$. Up to extracting a subsequence ε_k , we construct L^∞ maps

$$t \mapsto \gamma_t(dx, d\sigma, d\omega), \quad t \mapsto \nu_t(d\sigma, dv), \quad t \mapsto M_t(\sigma, v)$$

valued respectively on the set of positive Radon measures on $\mathbb{R}^d \times \Lambda \times \mathbf{S}^{d-1}$, on the set of positive Radon measures on $T^*\Lambda$ and finally on the set of measurable families from $T^*\Lambda$ onto the set of positive trace class operators on $L^2(N\Lambda)$, such that for all $\theta \in L^1(\mathbb{R})$ and for all $a \in \mathcal{A}$:

$$\begin{aligned} \int_{\mathbb{R}} \theta(t) \left\langle W_{u^{\varepsilon_k}(t, \cdot)}^{\Lambda, \varepsilon_k}, a \right\rangle dt &\xrightarrow{k \rightarrow +\infty} \int_{\mathbb{R}} \int_{\mathbb{R}^d \times \Lambda \times \mathbf{S}^{d-1}} \theta(t) a_\infty(x, \sigma, \omega) \gamma_t(dx, d\sigma, d\omega) dt \\ &\quad + \int_{\mathbb{R}} \int_{T^*\Lambda} \theta(t) \text{Tr}_{L^2(N_\sigma\Lambda)}(Q_a(\sigma, v) M_t(\sigma, v) \nu_t(d\sigma, dv)) dt. \end{aligned}$$

The measures γ^t and ν^t , and the map M^t satisfy additional properties coming from the fact that the family $(u^\varepsilon(t, \cdot))$ solves a time-dependent equation. These properties are discussed in the next two sections. We shall see that the measures γ_t are invariant under a linear flow and that we can choose the sequence ε_k such that the map $t \mapsto M_t$ is continuous (and even \mathcal{C}^1).

5.1. Transport properties of the compact part. Since Λ is the set of critical points of λ , the matrix $d^2\lambda$ is intrinsically defined above points of Λ . Thus, using the formalism of the preceding sections,

$$Q_{d^2\lambda(\sigma)\eta \cdot \eta} = d^2\lambda(\sigma) D_z \cdot D_z.$$

Proposition 5.1. *The map $t \mapsto \nu_t$ is constant and the map*

$$t \mapsto M_t(\sigma, v) \in \mathcal{C}(\mathbb{R}; \mathcal{L}_+^1(L^2(N_\sigma\Lambda)))$$

solves the Heisenberg equation (2.6).

Proof. We analyze for $a \in \mathcal{C}_c^\infty(\mathbb{R}^{3d})$ the time evolution of the quantity $\left\langle W_{u^\varepsilon(t, \cdot)}^{\Lambda, \varepsilon}, a \right\rangle$. We have

$$\frac{d}{dt} \left\langle W_{u^\varepsilon(t, \cdot)}^{\Lambda, \varepsilon}, a \right\rangle = \frac{1}{i\varepsilon^2} ([\text{op}_\varepsilon(a_\varepsilon), \lambda(\varepsilon D)] u^\varepsilon(t, \cdot), u^\varepsilon(t, \cdot)) + \frac{1}{i} ([\text{op}_\varepsilon(a_\varepsilon), V_{\text{ext}}] u^\varepsilon(t, \cdot), u^\varepsilon(t, \cdot)) + O(\varepsilon).$$

By standard symbolic calculus for Weyl quantization, we have in $\mathcal{L}(L^2(\mathbb{R}^d))$

$$\frac{1}{i\varepsilon^2} [\text{op}_\varepsilon(a_\varepsilon), \lambda(\varepsilon D)] = \frac{1}{\varepsilon} \text{op}_\varepsilon(\nabla \lambda(\xi) \cdot \nabla_x a_\varepsilon) + O(\varepsilon).$$

Besides, by Taylor formula and by use of $\nabla \lambda(\sigma(\xi)) = 0$, we have

$$(5.1) \quad \nabla \lambda(\xi) = d^2 \lambda(\sigma(\xi)) (\xi - \sigma(\xi)) + \Gamma(\xi) (\xi - \sigma(\xi)) \cdot (\xi - \sigma(\xi)),$$

where Γ is a smooth matrix. This yields

$$\frac{1}{\varepsilon} \nabla \lambda(\xi) \cdot \nabla_x a_\varepsilon(x, \xi) = b_\varepsilon(x, \xi)$$

with

$$b(x, \xi, \eta) = d^2 \lambda(\sigma(\xi)) \eta \cdot \nabla_x a(x, \xi, \eta) + \Gamma(\xi) (\xi - \sigma(\xi)) \cdot \eta \nabla_x a(x, \xi, \eta).$$

At this stage of the proof, we see that $\frac{d}{dt} \left\langle W_{u^\varepsilon(t, \cdot)}^{\Lambda, \varepsilon}, a \right\rangle$ is uniformly bounded in ε , thus using a suitable version of Ascoli's theorem and a standard diagonal extraction argument, we can find a sequence (ε_k) such that the limit exists for all $a \in \mathcal{C}_c^\infty(\mathbb{R}^{3d})$ and all time $t \in [0, T]$ (for some $T > 0$ fixed) with a limit that is a continuous map in time. The transport equation that we are now going to prove shall guarantee the independence of the limit from $T > 0$.

We observe that for any local system of equations of Λ , $\varphi(\xi) = 0$, the operator Q_b^φ satisfies for $(\sigma, v) \in T\Lambda$,

$$\begin{aligned} Q_b^\varphi(\sigma, v) &= b^W(v + {}^t d\varphi(\sigma)z, \sigma, B(\sigma)D_z) \\ &= \text{op}_1(d^2 \lambda(\sigma)B(\sigma)\eta \cdot \nabla_x a(v + {}^t d\varphi(\sigma)z, \sigma, B(\sigma)\eta)). \end{aligned}$$

On the other hand, we observe that, setting

$$\theta(\xi, \eta) = \frac{1}{2} d^2 \lambda(\xi) \eta \cdot \eta,$$

we have

$$(5.2) \quad \begin{aligned} i [Q_\theta^\varphi(\sigma), Q_a^\varphi(\sigma, v)] &= i [{}^t B(\sigma) d^2 \lambda(\sigma) B(\sigma) D_z \cdot D_z, Q_a^\varphi(\sigma, v)] \\ &= \text{op}_1({}^t d\varphi(\sigma) {}^t B(\sigma) d^2 \lambda(\sigma) B(\sigma) \eta \cdot \nabla_x a(v + {}^t d\varphi(\sigma)z, \sigma, B(\sigma)\eta)), \end{aligned}$$

and we now focus on the matrix ${}^t d\varphi(\sigma) {}^t B(\sigma) d^2 \lambda(\sigma) B(\sigma)$, and thus on the properties of the hessian $d^2 \lambda(\sigma)$.

For $\xi \in \Lambda$, the bilinear form $d^2 \lambda(\xi)$ is defined intrinsically on $T_\xi \mathbb{R}^d$ and $d^2 \lambda(\xi) = 0$ on $T_\xi \Lambda$. We deduce from (4.6) that any $\zeta \in T_\xi \mathbb{R}^d$ satisfies

$$\zeta = d\sigma(\xi)\zeta + B(\xi)d\varphi(\xi)\zeta \quad \text{with } d\sigma(\xi)\zeta \in T_\sigma \Lambda.$$

Therefore,

$$\forall \xi \in \Lambda, \quad d^2 \lambda(\xi) = d^2 \lambda(\xi) B(\xi) d\varphi(\xi).$$

Taking into account this information, equation (5.2) becomes

$$i [Q_\theta^\varphi(\sigma), Q_a^\varphi(\sigma, v)] = \text{op}_1(d^2 \lambda(\sigma)B(\sigma)\eta \cdot \nabla_x a(v + {}^t d\varphi(\sigma)z, \sigma, B(\sigma)\eta)).$$

We conclude

$$Q_b^\varphi(\sigma, v) = i [Q_\theta^\varphi(\sigma), Q_a^\varphi(\sigma, v)].$$

This implies that

$$i\partial_t(M_t(\sigma, v)\nu_t(d\sigma, dv)) = \left[\frac{1}{2}d^2\lambda(\sigma)D_z \cdot D_z + m_{V_{\text{ext}}(t, \cdot)}(v, \sigma), M_t(\sigma, v) \right] \nu_t(d\sigma, dv).$$

Taking the trace, we get $\partial_t\nu_t = 0$, thus ν_t is equal to some constant measure ν and M_t satisfies equation (2.6), which proves the proposition. \square

5.2. Invariance and localization of the measure at infinity. We are concerned with the property of the L^∞ -map $t \mapsto \gamma^t(dx, d\sigma, d\omega)$ valued in the set of positive Radon measures on $S\Lambda$. We now define a flow on $S\Lambda$ by setting for $s \in \mathbb{R}$

$$\phi_2^s : (x, \sigma, \omega) \mapsto (x + s d^2\lambda(\sigma)\omega, \sigma, \omega).$$

Proposition 5.2. *The measure γ^t is invariant by the flow ϕ_2^s .*

Proof. The proof essentially follows the lines of the proof of Theorem 2.5 in [4]. We use the cut-off function χ introduced before and set

$$a^{R, \delta}(x, \xi, \eta) = a(x, \xi, \eta) \chi\left(\frac{\xi - \sigma(\xi)}{\delta}\right) \left(1 - \chi\left(\frac{\eta}{R}\right)\right);$$

we introduce the smooth symbol

$$b_s^{R, \delta}(x, \xi, \eta) = a^{R, \delta}\left(x + s d^2\lambda(\xi) \frac{\eta}{|\eta|}, \xi, \eta\right),$$

which satisfies $(b_s^R)_\infty = a_\infty \circ \phi_2^s$. Using equation (5.1), we obtain

$$\left(b_s^{R, \delta}\right)_\varepsilon(x, \xi) = a^{R, \delta}\left(x + \frac{s}{|\xi - \sigma(\xi)|} \nabla\lambda(\xi), \xi, \frac{\xi - \sigma(\xi)}{\varepsilon}\right) + \delta r_\varepsilon^{R, \delta}(x, \xi)$$

where for all multi-index $\alpha, \beta \in \mathbb{N}^d$, there exists a constant $C_{\alpha, \beta} > 0$ such that $r_\varepsilon^{R, \delta}$ satisfies:

$$\sup_{x, \xi \in \mathbb{R}^d} \left| \partial_x^\alpha \partial_\xi^\beta r_\varepsilon^{R, \delta} \right| \leq C_{\alpha, \beta}.$$

As a consequence, $\langle W_{u^\varepsilon(t, \cdot)}^{\Lambda, \varepsilon}, r_\varepsilon^{R, \delta} \rangle$ is uniformly bounded in R, δ, ε and:

$$\langle W_{u^\varepsilon(t, \cdot)}^{\Lambda, \varepsilon}, b_s^{R, \delta} \rangle = \langle W_{u^\varepsilon(t, \cdot)}^{\Lambda, \varepsilon}, \tilde{b}_s^{R, \delta} \rangle + O(\delta),$$

uniformly with respect to R and ε , with

$$\tilde{b}_s^{R, \delta}(x, \xi, \eta) = a^{R, \delta}\left(x + \frac{s}{|\xi - \sigma(\xi)|} \nabla\lambda(\xi), \xi, \eta\right).$$

Note that this symbol is smooth because $|\xi - \sigma(\xi)| > R\varepsilon$ on the support of $a^{R, \delta}$. We are going to prove that for all $\theta \in \mathcal{C}_c^\infty(\mathbb{R})$,

$$\lim_{\delta \rightarrow 0^+} \lim_{R \rightarrow \infty} \lim_{\varepsilon \rightarrow 0^+} \int_{\mathbb{R}} \theta(t) \frac{d}{ds} \langle W_{u^\varepsilon(t, \cdot)}^{\Lambda, \varepsilon}, \tilde{b}_s^{R, \delta} \rangle dt = 0.$$

Indeed, by the calculus of the preceding section, we have

$$\frac{d}{ds} \langle W_{u^\varepsilon(t, \cdot)}^{\Lambda, \varepsilon}, \tilde{b}_s^{R, \delta} \rangle = \langle W_{u^\varepsilon(t, \cdot)}^{\Lambda, \varepsilon}, \nabla\lambda \cdot \nabla_x c_s^{R, \delta} \rangle$$

with

$$c_s^{R,\delta}(x, \xi, \eta) = \frac{1}{|\xi - \sigma(\xi)|} a^{R,\delta} \left(x + \frac{s}{|\xi - \sigma(\xi)|} \nabla \lambda(\xi), \xi, \eta \right).$$

The symbol $c_s^{R,\delta}$ is such that for all multi-index $\alpha \in \mathbb{N}^d$, there exists a constant $C_\alpha > 0$ for which:

$$\sup_{x, \xi \in \mathbb{R}^d} \left| \partial_x^\alpha (c_s^{R,\delta})_\varepsilon \right| \leq C_\alpha (R\varepsilon)^{-1}.$$

This implies in particular:

$$\left\| \text{op}_\varepsilon((c_s^{R,\delta})_\varepsilon) \right\|_{\mathcal{L}(L^2(\mathbb{R}^d))} \leq \frac{C}{R\varepsilon}.$$

By symbolic calculus, we have

$$\frac{1}{i\varepsilon} \left[\text{op}_\varepsilon((c_s^{R,\delta})_\varepsilon), \lambda(\varepsilon D) \right] = \text{op}_\varepsilon \left(\nabla \lambda(\xi) \cdot \nabla_x (c_s^{R,\delta})_\varepsilon \right) + O\left(\frac{\varepsilon}{R}\right).$$

We deduce that for all $\theta \in C_c^\infty(\mathbb{R})$,

$$\begin{aligned} & \int_{\mathbb{R}} \theta(t) \frac{d}{ds} \langle W_{u^\varepsilon(t,\cdot)}^{\Lambda,\varepsilon}, \tilde{b}_s^{R,\delta} \rangle dt \\ &= \int_{\mathbb{R}} \theta(t) \left(\frac{1}{i\varepsilon} \left[\text{op}_\varepsilon((c_s^{R,\delta})_\varepsilon), \lambda(\varepsilon D) \right] u^\varepsilon(t,\cdot), u^\varepsilon(t,\cdot) \right) dt + O\left(\frac{\varepsilon}{R}\right) \\ &= \int_{\mathbb{R}} \theta(t) \left(\frac{1}{i\varepsilon} \left[\text{op}_\varepsilon((c_s^{R,\delta})_\varepsilon), \lambda(\varepsilon D) + \varepsilon^2 V_{\text{ext}}(t, x) \right] u^\varepsilon(t,\cdot), u^\varepsilon(t,\cdot) \right) dt + O\left(\frac{1}{R}\right) \\ &= -\varepsilon \int_{\mathbb{R}} \theta(t) \frac{d}{dt} \left(\text{op}_\varepsilon((c_s^{R,\delta})_\varepsilon) u^\varepsilon(t,\cdot), u^\varepsilon(t,\cdot) \right) dt + O\left(\frac{1}{R}\right) \\ &= O(\varepsilon) + O\left(\frac{1}{R}\right). \end{aligned}$$

As a conclusion,

$$\begin{aligned} \langle W_{u^\varepsilon(t,\cdot)}^{\Lambda,\varepsilon}, b_s^{R,\delta} \rangle &= \langle W_{u^\varepsilon(t,\cdot)}^{\Lambda,\varepsilon}, \tilde{b}_s^{R,\delta} \rangle + O(\delta) \\ &= \langle W_{u^\varepsilon(t,\cdot)}^{\Lambda,\varepsilon}, \tilde{b}_0^{R,\delta} \rangle + O(|s|\varepsilon) + O(|s|R^{-1}) + O(\delta) \\ &= \langle W_{u^\varepsilon(t,\cdot)}^{\Lambda,\varepsilon}, b_0^{R,\delta} \rangle + O(|s|\varepsilon) + O(|s|R^{-1}) + O(\delta), \end{aligned}$$

which implies the Proposition. \square

5.3. Proofs of Theorems 2.1 and 2.2. Remind that Theorem 2.2 implies Theorem 2.1, thus we focus on Theorem 2.2. We first observe that the measure γ_t is zero. Indeed, by **H2**; for $\sigma \in \Lambda$, $d^2\lambda(\sigma)$ is one to one on $N_\sigma\Lambda$. Therefore, since γ_t is a measure on $S\Lambda$, the invariance property of Proposition 5.2 and an argument similar to the one of Lemma 3.6 yields that $\gamma_t = 0$. As a consequence, the semi-classical measure μ_t is only given by the compact part and one has for any $a \in C_c^\infty(\mathbb{R}^{2d})$ and $\theta \in L^1(\mathbb{R})$,

$$\int_{\mathbb{R}} \theta(t) \int_{\mathbb{R}^{2d}} a(x, \xi) \mu^t(dx, d\xi) = \int_{\mathbb{R}} \theta(t) \int_{T^*\Lambda} \text{Tr}_{L^2(N_\sigma\Lambda)} (Q_a(\sigma, v) M_t(d\sigma, dv)) dt.$$

Then, taking $\theta = \mathbb{1}_{[a,b]}$ for $a, b \in \mathbb{R}$, $a < b$, and in view of Proposition 3.2 and of Lemma 3.3, we deduce that for every $\phi \in \mathcal{C}_c(\mathbb{R}^d)$ one has for the subsequence defining M_t and ν_t :

$$\lim_{\varepsilon \rightarrow 0} \int_a^b \int_{\mathbb{R}^d} \phi(x) |u^\varepsilon(t, x)|^2 dx dt = \int_a^b \int_{T^*\Lambda} \text{Tr}_{L^2(N_{\xi\Lambda})} [Q_\phi(v, \xi) M_t(v, \xi)] \nu(dv, d\xi) dt,$$

where M_t satisfies (2.6). This concludes the proof of Theorem 2.2. We emphasize that the measure ν and the operator valued family M_0 are utterly determined by the initial data.

6. BLOCH PROJECTORS AND SEMICLASSICAL MEASURES

In this section we prove Theorem 1.2, as a result of the analysis in Section 4. We shall use properties of the operator of restriction L^ε defined in (1.7) and of the projector $\Pi_n(\varepsilon D_x)$. Then, we prove *a priori* estimates for solutions of equation (1.6) and use them to reduce the dynamics of our original problem to those of equation (1.13) (Corollary 6.8).

Note that, modulo adding a positive constant to equation (1.2), we may assume that $P(\varepsilon D_x)$ is a non-negative operator. With this in mind, the following estimates, that will be repeatedly used in what follows, hold.

Remark 6.1. There exists a constant $c > 0$ such that:

$$c^{-1} \|U\|_{H_\varepsilon^s(\mathbb{R}^d \times \mathbb{T}^d)} \leq \| \langle \varepsilon D_x \rangle^s U \|_{L^2(\mathbb{R}^d \times \mathbb{T}^d)} + \|P(\varepsilon D_x)^{s/2} U\|_{L^2(\mathbb{R}^d \times \mathbb{T}^d)} \leq c \|U\|_{H_\varepsilon^s(\mathbb{R}^d \times \mathbb{T}^d)},$$

for every $U \in L^2(\mathbb{R}^d \times \mathbb{T}^d)$ and $\varepsilon > 0$, where, as usual, $\langle \xi \rangle = (1 + |\xi|^2)^{1/2}$ and where the sets H_ε^s have been defined in (1.8).

6.1. High frequency behavior of the operator of restriction to the diagonal and of the Bloch projectors. We first focus on the properties of the operator of restriction to the diagonal L^ε and prove its boundedness in appropriate functional spaces.

Lemma 6.2. *Suppose $s > d/2$, then the operator*

$$L^\varepsilon : L^2(\mathbb{R}_x^d; H^s(\mathbb{T}_y^d)) \longrightarrow L^2(\mathbb{R}^d)$$

is uniformly bounded in ε .

Moreover, if $U^\varepsilon \in L^2(\mathbb{R}_x^d; H^s(\mathbb{T}_y^d))$ satisfies the estimate:

$$(6.1) \quad \limsup_{\varepsilon \rightarrow 0^+} \|\mathbb{1}_R(\varepsilon D_x) U^\varepsilon\|_{L^2(\mathbb{R}^d; H^s(\mathbb{T}^d))} \xrightarrow{R \rightarrow \infty} 0,$$

where $\mathbb{1}_R$ is the characteristic function of $\{|\xi| > R\}$, then the sequence $(L^\varepsilon U^\varepsilon)$ is bounded in $L^2(\mathbb{R}^d)$ and ε -oscillating.

Remark 6.3. Suppose that (U^ε) is bounded in $H_\varepsilon^r(\mathbb{R}^d \times \mathbb{T}^d)$ for some $r > d/2$. Then condition (6.1) is satisfied for every $d/2 < s < r$. This follows from the bound:

$$\|\mathbb{1}_R(\varepsilon D_x) U^\varepsilon\|_{L^2(\mathbb{R}^d; H^s(\mathbb{T}^d))} \leq R^{s-r} \|U^\varepsilon\|_{H_\varepsilon^r(\mathbb{R}^d \times \mathbb{T}^d)}.$$

In particular, if ψ_0^ε satisfies item (3) of Assumptions 1.1, then (ψ_0^ε) is ε -oscillating.

Proof. Let $U^\varepsilon \in L^2(\mathbb{R}_x^d; H^s(\mathbb{T}_y^d))$ and write

$$U^\varepsilon(x, y) = \sum_{k \in \mathbb{Z}^d} U_k^\varepsilon(x) e^{i2\pi k \cdot y},$$

and

$$\|U^\varepsilon\|_{L^2(\mathbb{R}_x^d; H^s(\mathbb{T}_y^d))}^2 = \sum_{k \in \mathbb{Z}^d} \langle k \rangle^{2s} \|U_k^\varepsilon\|_{L^2(\mathbb{R}^d)}^2.$$

Then there exist constants $C, C_{d,s} > 0$ such that

$$\sum_{k \in \mathbb{Z}^d} \|U_k^\varepsilon\|_{L^2(\mathbb{R}^d)} \leq C \left(\sum_{k \in \mathbb{Z}^d} |k|^{2s} \|U_k^\varepsilon\|_{L^2(\mathbb{R}^d)}^2 \right)^{1/2} \leq C_{d,s} \|U^\varepsilon\|_{L^2(\mathbb{R}_x^d; H^s(\mathbb{T}_y^d))},$$

and therefore:

$$(6.2) \quad \|L^\varepsilon U^\varepsilon\|_{L^2(\mathbb{R}^d)} \leq \sum_{k \in \mathbb{Z}^d} \|U_k^\varepsilon\|_{L^2(\mathbb{R}^d)} \leq C_{d,s} \|U^\varepsilon\|_{L^2(\mathbb{R}_x^d; H^s(\mathbb{T}_y^d))}.$$

Let us now show that, under the hypothesis of the proposition, $v^\varepsilon := L^\varepsilon U^\varepsilon$ defines an ε -oscillating sequence. Given $\delta > 0$, since $s > d/2$, there exists $N_\delta > 0$ such that

$$\sum_{|k| > N_\delta} |k|^{-2s} < \delta^2.$$

Define:

$$v_\delta^\varepsilon(x) = \sum_{|k| \leq N_\delta} U_k^\varepsilon(x) e^{i2\pi k \cdot \frac{x}{\varepsilon}}.$$

Clearly,

$$\|v^\varepsilon - v_\delta^\varepsilon\|_{L^2(\mathbb{R}^d)} \leq \delta \|U^\varepsilon\|_{L^2(\mathbb{R}_x^d; H^s(\mathbb{T}_y^d))}.$$

Therefore, it suffices to show that for any $\delta > 0$ the sequence (v_δ^ε) is ε -oscillating. The Fourier transform of v_δ^ε is:

$$\widehat{v}_\delta^\varepsilon(\xi) = \sum_{|k| \leq N_\delta} \widehat{U}_k^\varepsilon \left(\xi - \frac{2\pi k}{\varepsilon} \right).$$

Therefore,

$$\|\mathbb{1}_R(\varepsilon D_x) v_\delta^\varepsilon\|_{L^2(\mathbb{R}^d)} \leq \sum_{|k| \leq N_\delta} \|\mathbb{1}_R(\varepsilon D_x + 2\pi k) U_k^\varepsilon\|_{L^2(\mathbb{R}^d)}.$$

If $R > R_0$ for $R_0 > 0$ large enough, one has $\mathbb{1}_R(\cdot + 2\pi k) \leq \mathbb{1}_{R/2}$ for every $|k| \leq N_\delta$. This allows us to conclude that for $R > R_0$:

$$\|\mathbb{1}_R(\varepsilon D_x) v_\delta^\varepsilon\|_{L^2(\mathbb{R}^d)} \leq \sum_{|k| \leq N_\delta} \|\mathbb{1}_{R/2}(\varepsilon D_x) U_k^\varepsilon\|_{L^2(\mathbb{R}^d)} \leq C_{d,s} \|\mathbb{1}_R(\varepsilon D_x) U^\varepsilon\|_{L^2(\mathbb{R}_x^d; H^s(\mathbb{T}_y^d))}$$

and the conclusion follows. \square

We shall also need information on the derivatives with respect to ξ of the operator $\Pi_n(\xi)$. We recall the formula

$$\Pi_n(\xi) = -\frac{1}{2i\pi} \sum_{j=1}^N \chi_j(\xi) \oint_{\mathcal{C}_j} (P(\xi) - z)^{-1} dz$$

where the functions $\chi_j \in \mathcal{C}^\infty(\mathbb{R}^d/2\pi\mathbb{Z}^d)$ form a partition of unity and, for $j = 1, \dots, N$, \mathcal{C}_j is a contour in the complex plane separating $\varrho_n(\xi)$, for $\xi \in \text{supp } \chi_j$, from the remainder of the spectrum. The existence of such contours is guaranteed by the fact that $\varrho_n(\xi)$ is of constant

multiplicity for all $\xi \in \mathbb{R}^d$ and, thus, is separated from the remainder of the spectrum. As a consequence of this formula, of Lemma 6.1 and of the relation

$$[\Pi_n(\varepsilon D_x), P(\varepsilon D_x)^{s/2}] = [\Pi_n(\varepsilon D_x), \langle \varepsilon D_x \rangle^s] = 0,$$

we deduce the following result.

Lemma 6.4. *The map $\xi \mapsto \Pi_n(\xi)$ is a smooth bounded map from \mathbb{R}^d into $\mathcal{L}(L^2(\mathbb{T}^d))$. In addition, the operator $\Pi_n(\varepsilon D_x)$ maps the space $H_\varepsilon^s(\mathbb{R}^d \times \mathbb{T}^d)$ into itself.*

6.2. A priori estimates on $U^\varepsilon(t, \cdot)$. In order to derive the desired properties of $\psi_n^\varepsilon(t, x)$, the solution to (1.13), we need to prove some *a priori* estimates for the solutions of equation (1.6). We will use them for reducing the analysis of $\psi^\varepsilon(t, \cdot)$ (the solution to our original problem (1.2)) to that of $\psi_n^\varepsilon(t, \cdot)$.

Lemma 6.5. *Given $s \geq 0$, there exists a constant $C_s > 0$ such that any solution U^ε to (1.6) with initial datum $U_0^\varepsilon \in H^s(\mathbb{R}^d \times \mathbb{T}^d)$ satisfies:*

$$(6.3) \quad \|U^\varepsilon(t, \cdot)\|_{H_\varepsilon^s(\mathbb{R}^d \times \mathbb{T}^d)} \leq \|U_0^\varepsilon\|_{H_\varepsilon^s(\mathbb{R}^d \times \mathbb{T}^d)} + C_s \varepsilon |t|,$$

uniformly in $\varepsilon > 0$.

Corollary 6.6. *Lemma 6.5 and Remark 6.3 imply that for all $t \in \mathbb{R}$, the family $(\psi^\varepsilon(t, \cdot))$ is ε -oscillating.*

Proof. In view of Remark 6.1, we are first going to study the families

$$(\langle \varepsilon D_x \rangle U^\varepsilon) \quad \text{and} \quad (P(\varepsilon D_x)^{1/2} U^\varepsilon).$$

Start noticing that $\langle \varepsilon D_x \rangle U^\varepsilon$ satisfies the equation

$$(6.4) \quad i\varepsilon^2 \partial_t (\langle \varepsilon D_x \rangle U^\varepsilon) = P(\varepsilon D_x) (\langle \varepsilon D_x \rangle U^\varepsilon) + \varepsilon^2 V_{\text{ext}} \langle \varepsilon D_x \rangle U^\varepsilon - \varepsilon^2 [V_{\text{ext}}, \langle \varepsilon D_x \rangle] U^\varepsilon.$$

As a consequence, using the boundedness of $\nabla_x V_{\text{ext}}$ on $\mathbb{R} \times \mathbb{R}^d$, we obtain by the symbolic calculus of semiclassical pseudodifferential operators, that the source term can be estimated by:

$$\|[V_{\text{ext}}(t, \cdot), \langle \varepsilon D_x \rangle] U^\varepsilon(t, \cdot)\|_{L^2(\mathbb{R}^d \times \mathbb{T}^d)} \leq C \varepsilon \|U^\varepsilon(t, \cdot)\|_{L^2(\mathbb{R}^d \times \mathbb{T}^d)},$$

for some constant $C > 0$ independent of $\varepsilon > 0$ and $t \in \mathbb{R}$. Using standard energy estimates, we deduce the existence of a constant $C_1 > 0$ such that for all $t \in \mathbb{R}$,

$$\|\langle \varepsilon D_x \rangle U^\varepsilon(t, \cdot)\|_{L^2(\mathbb{R}^d \times \mathbb{T}^d)} \leq \|\langle \varepsilon D_x \rangle U_0^\varepsilon\|_{L^2(\mathbb{R}^d \times \mathbb{T}^d)} + C_1 \varepsilon |t|.$$

A completely analogous argument yields the estimate:

$$\|P(\varepsilon D_x)^{1/2} U^\varepsilon(t, \cdot)\|_{L^2(\mathbb{R}^d \times \mathbb{T}^d)} \leq \|P(\varepsilon D_x)^{1/2} U_0^\varepsilon\|_{L^2(\mathbb{R}^d \times \mathbb{T}^d)} + C_1 \varepsilon |t|.$$

A standard recursive argument gives, for all $s \in \mathbb{N}$, the existence of a constant $C_s > 0$ such that for all $t \in \mathbb{R}$,

$$\begin{aligned} & \|\langle \varepsilon D_x \rangle^s U^\varepsilon(t, \cdot)\|_{L^2(\mathbb{R}^d \times \mathbb{T}^d)} + \|P(\varepsilon D_x)^{s/2} U^\varepsilon(t, \cdot)\|_{L^2(\mathbb{R}^d \times \mathbb{T}^d)} \leq \\ & \|\langle \varepsilon D_x \rangle^s U_0^\varepsilon\|_{L^2(\mathbb{R}^d \times \mathbb{T}^d)} + \|P(\varepsilon D_x)^{s/2} U_0^\varepsilon\|_{L^2(\mathbb{R}^d \times \mathbb{T}^d)} + C_s \varepsilon |t|, \end{aligned}$$

and the result follows for any $s \in \mathbb{R}^+$ by interpolation. \square

We now focus on the case where the initial data U_0^ε belongs to a particular Bloch eigenspace: $U_0^\varepsilon = \Pi_n(\varepsilon D_x)U_0^\varepsilon$. We set

$$\tilde{U}^\varepsilon(t, \cdot) = \Pi_n(\varepsilon D_x)U^\varepsilon(t, \cdot).$$

Note that by Lemma 6.4, for any $t \in \mathbb{R}$, the family $\tilde{U}^\varepsilon(t, \cdot)$ is uniformly bounded in $H_\varepsilon^s(\mathbb{R}^d \times \mathbb{T}^d)$.

Lemma 6.7. *Assume $U_0^\varepsilon = \Pi_n(\varepsilon D_x)U_0^\varepsilon$ and consider $\tilde{U}^\varepsilon(t, \cdot)$ as defined above. Then, for all $T > 0$, there exists $C_T > 0$ such that*

$$\sup_{t \in [0, T]} \left\| U^\varepsilon(t, \cdot) - \tilde{U}^\varepsilon(t, \cdot) \right\|_{H_\varepsilon^s(\mathbb{T}^d \times \mathbb{R}^d)} \leq C_T \varepsilon.$$

Let us prove now Lemma 6.7.

Proof. Note first that, in view of Remark 6.1, it is enough to prove the uniform boundedness in $L^2(\mathbb{T}^d \times \mathbb{R}^d)$ of

$$U^\varepsilon(t, \cdot) - \tilde{U}^\varepsilon(t, \cdot), \quad P(\varepsilon D_x)^{s/2}(U^\varepsilon(t, \cdot) - \tilde{U}^\varepsilon(t, \cdot)) \quad \text{and} \quad \langle \varepsilon D_x \rangle^s (U^\varepsilon(t, \cdot) - \tilde{U}^\varepsilon(t, \cdot)).$$

We have $U^\varepsilon(0, \cdot) = \tilde{U}^\varepsilon(0, \cdot)$ and \tilde{U}^ε solves

$$(6.5) \quad i\varepsilon^2 \partial_t \tilde{U}^\varepsilon(t, x) = P(\varepsilon D_x) \tilde{U}^\varepsilon(t, x) + \varepsilon^2 V_{\text{ext}}(t, x) \tilde{U}^\varepsilon(t, x) + \varepsilon^2 B^\varepsilon(t) U^\varepsilon(t, x),$$

with

$$B^\varepsilon(t) = [\Pi_n(\varepsilon D_x), V_{\text{ext}}(t, \cdot)].$$

The symbolic calculus of semiclassical pseudodifferential operators implies that:

$$\|B^\varepsilon(t)U^\varepsilon(t, \cdot)\|_{L^2(\mathbb{R}^d \times \mathbb{T}^d)} = O(\varepsilon), \quad \text{locally uniformly in } t.$$

As for $\langle \varepsilon D_x \rangle \tilde{U}^\varepsilon$ one has:

$$i\varepsilon^2 \partial_t (\langle \varepsilon D_x \rangle \tilde{U}^\varepsilon) = P(\varepsilon D_x) \langle \varepsilon D_x \rangle \tilde{U}^\varepsilon + \varepsilon^2 V_{\text{ext}} \langle \varepsilon D_x \rangle \tilde{U}^\varepsilon + \varepsilon^2 C^\varepsilon \langle \varepsilon D_x \rangle U^\varepsilon - \varepsilon^2 [V_{\text{ext}}, \langle \varepsilon D_x \rangle] \tilde{U}^\varepsilon,$$

with,

$$C^\varepsilon = [\Pi_n(\varepsilon D_x), \langle \varepsilon D_x \rangle V_{\text{ext}} \langle \varepsilon D_x \rangle^{-1}].$$

Again, the symbolic calculus gives that $\|C^\varepsilon(t) \langle \varepsilon D_x \rangle U^\varepsilon(t, \cdot)\|_{L^2(\mathbb{R}^d \times \mathbb{T}^d)} = O(\varepsilon)$ locally uniformly in t . Taking into account that $\langle \varepsilon D_x \rangle U^\varepsilon$ satisfies equation (6.4) and is bounded in $L^2(\mathbb{R}^d \times \mathbb{T}^d)$, one concludes that:

$$\|\langle \varepsilon D_x \rangle (U^\varepsilon(t, \cdot) - \tilde{U}^\varepsilon(t, \cdot))\|_{L^2(\mathbb{R}^d \times \mathbb{T}^d)} \leq C\varepsilon |t|.$$

An analogous reasoning holds for $P(\varepsilon D_x)^{1/2}(U^\varepsilon(t, \cdot) - \tilde{U}^\varepsilon(t, \cdot))$. One concludes using an inductive argument following the lines of the end of the proof of Lemma 6.5. \square

6.3. Analysis of the Bloch component ψ_n^ε . By the definition of $\psi_n^\varepsilon(t, x)$, we have

$$\psi_n^\varepsilon(t, \cdot) = L^\varepsilon \tilde{U}^\varepsilon(t, \cdot);$$

and the family is bounded in $L^2(\mathbb{R}^d)$ for all $t \in \mathbb{R}$. Moreover, as a corollary of Lemma 6.7, the following holds.

Corollary 6.8. *Suppose that ψ^ε and ψ_n^ε are the respective solutions of equations (1.2) and (1.13) with the same initial datum $L^\varepsilon U_0^\varepsilon$, where $U_0^\varepsilon = \Pi_n(\varepsilon D_x)U_0^\varepsilon$. Then for every $T > 0$ there exist $C_T > 0$ such that, uniformly in ε ,*

$$\sup_{t \in [0, T]} \|\psi^\varepsilon(t, \cdot) - \psi_n^\varepsilon(t, \cdot)\|_{L^2(\mathbb{R}^d)} \leq C_T \varepsilon.$$

The proof is a direct consequence of Lemma 6.7, since Lemma 6.2 ensures that

$$\|\psi^\varepsilon(t, \cdot) - \psi_n^\varepsilon(t, \cdot)\|_{L^2(\mathbb{R}^d)} \leq C \|U^\varepsilon(t, \cdot) - \tilde{U}^\varepsilon(t, \cdot)\|_{L^2(\mathbb{R}^d, H^s(\mathbb{T}^d))}.$$

We now conclude our analysis of the Bloch component $\psi_n^\varepsilon(t, \cdot)$. The following result gathers the remaining information that we will need in order to conclude, together with Corollary 6.8, the proof of Theorem 1.2.

Proposition 6.9. *The family ψ_n^ε solves equation (1.13)*

$$\begin{cases} i\varepsilon^2 \partial_t \psi_n^\varepsilon(t, x) - \varrho_n(\varepsilon D_x) \psi_n^\varepsilon(t, x) - \varepsilon^2 V_{\text{ext}}(t, x) \psi_n^\varepsilon(t, x) = \varepsilon^2 f_n^\varepsilon(t, x), \\ \psi_n^\varepsilon|_{t=0}(x) = \psi_0^\varepsilon(x) \end{cases}$$

with (1.15): $\|f_n^\varepsilon(t, \cdot)\|_{L^2(\mathbb{R}^d)} \leq C\varepsilon$ for all $t \in \mathbb{R}$, $\varepsilon > 0$.

Proof. Let us first prove that ψ_n^ε solves (1.13). We denote by J the set of the indexes of the Bloch eigenfunctions $\varphi_j(\cdot, \xi)$ which form an orthonormal basis of $\text{Ran } \Pi_n(\xi)$. Define for $j \in J$,

$$u_j^\varepsilon(t, x) := \int_{\mathbb{T}^d} \bar{\varphi}_j(y, \varepsilon D_x) \tilde{U}^\varepsilon(t, x, y) dy,$$

and notice that:

$$\psi_n^\varepsilon(t, x) = (L^\varepsilon \tilde{U}^\varepsilon)(t, x) = \sum_{j \in J} \varphi_j\left(\frac{x}{\varepsilon}, \varepsilon D_x\right) u_j^\varepsilon(t, x).$$

Since \tilde{U}^ε solves (6.5) and $P(\xi)\varphi_j(\cdot, \xi) = \varrho_n(\xi)\varphi_j(\cdot, \xi)$ for all $\xi \in \mathbb{R}^d$, the family u_j^ε solves:

$$i\varepsilon^2 \partial_t u_j^\varepsilon(t, x) = \varrho_n(\varepsilon D_x) u_j^\varepsilon(t, x) + \varepsilon^2 V_{\text{ext}}(t, x) u_j^\varepsilon(t, x) + \varepsilon^2 g_j^\varepsilon(t, x),$$

where:

$$g_j^\varepsilon(t, x) := \int_{\mathbb{T}^d} [\bar{\varphi}_j(y, \varepsilon D_x), V_{\text{ext}}(t, x)] U^\varepsilon(t, x, y) dy.$$

Since $\varrho_n(\xi)$ is $2\pi\mathbb{Z}^d$ -periodic, it is easy to check that:

$$[L^\varepsilon \varphi_j(\cdot, \varepsilon D_x), \varrho_n(\varepsilon D_x)] = 0.$$

Summing the relations over $j \in J$, this implies (1.13) with $f_n^\varepsilon = L^\varepsilon[\Pi_n(\varepsilon D_x), V_{\text{ext}}]U^\varepsilon$. Now, Lemma 6.2 and the symbolic calculus of pseudodifferential operators gives, for any $t \in \mathbb{R}$:

$$\|f_n^\varepsilon(t, \cdot)\|_{L^2(\mathbb{R}^d)} \leq C \|[\Pi_n(\varepsilon D_x), V_{\text{ext}}(t, \cdot)]U^\varepsilon(t, \cdot)\|_{L^2(\mathbb{R}^d; H^s(\mathbb{T}^d))} \leq C'\varepsilon \|U^\varepsilon(t, \cdot)\|_{L^2(\mathbb{R}^d; H^s(\mathbb{T}^d))},$$

which concludes the proof. \square

6.4. Proofs of Theorems 1.2. The proof of Theorem 1.2 (which implies Corollary 1.4) easily follows from our results so far.

Proof. By Corollary 6.6, the family $(\psi^\varepsilon(t, \cdot))$ is ε -oscillating. Therefore, the weak limits of $|\psi^\varepsilon(t, x)|^2 dx$ are the projection on \mathbb{R}_x^d of the Wigner measures associated with $(\psi^\varepsilon(t, \cdot))$. By Corollary 6.8, the Wigner measures of $(\psi^\varepsilon(t, \cdot))$ coincide with those of $(\psi_n^\varepsilon(t, \cdot))$. Finally, Proposition 6.9 allows us to use the results of Theorem 2.1 for determining the Wigner measure of $(\psi_n^\varepsilon(t, \cdot))$. \square

6.5. Some comments on initial data that are a finite superposition of Bloch modes. Our results also apply to initial data that are a finite linear combination of the form:

$$(6.6) \quad \psi_0^\varepsilon = \sum_{n \in \mathcal{N}} L^\varepsilon U_{0,n}^\varepsilon$$

with \mathcal{N} a finite subset of \mathbb{N} such that for all $n \in \mathcal{N}$, $P(\varepsilon D_x)U_{0,n}^\varepsilon = \varrho_n(\varepsilon D_x)U_{0,n}^\varepsilon$, for distinct ϱ_n of constant multiplicity and $U_{0,n}^\varepsilon$ uniformly bounded in $H_\varepsilon^s(\mathbb{R}^d \times \mathbb{T}^d)$ for all $n \in \mathcal{N}$.

Proposition 6.10. *Assume we turn (3) into (6.6) in the hypothesis of Assumption 1.1 and that item (2) of Assumption 1.1 holds for every ϱ_n with $n \in \mathcal{N}$. Then, there exist a subsequence $(\varepsilon_k)_{k \in \mathbb{N}}$, positive measures $\nu_n \in \mathcal{M}_+(T^*\Lambda_n)$, and measurable families of self-adjoint, positive, trace-class operators*

$$M_{0,n} : T_\xi^*\Lambda_n \ni (v, \xi) \longmapsto M_{0,n}(v, \xi) \in \mathcal{L}_+^1(L^2(N_\xi\Lambda_n)), \quad \text{Tr}_{L^2(N_\xi\Lambda_n)} M_{0,n}(v, \xi) = 1,$$

such that for every $a < b$ and $\phi \in \mathcal{C}_c(\mathbb{R}^d)$ one has:

$$\begin{aligned} \lim_{k \rightarrow \infty} \int_a^b \int_{\mathbb{R}^d} \phi(x) |\psi^{\varepsilon_k}(t, x)|^2 dx dt \\ = \sum_{n \in \mathcal{N}} \int_a^b \int_{T^*\Lambda_n} \text{Tr}_{L^2(N_\xi\Lambda_n)} [m_\phi(v, \xi) M_n(t, v, \xi)] \nu_j(dv, d\xi) dt, \end{aligned}$$

where $M_n(\cdot, v, \xi) \in \mathcal{C}(\mathbb{R}; \mathcal{L}_+^1(L^2(N_\xi\Lambda_n)))$ solves the Heisenberg equation (1.11) with initial data $M_{0,n}$ associated with the concentration of ψ_0^ε on Λ_n .

Proof. We associate to any $n \in \mathbb{N}$ their respective Bloch components $\psi_n^\varepsilon(t, \cdot)$ of $\psi^\varepsilon(t, \cdot)$ as we previously did. We just have to prove that for all $n, n' \in \mathcal{N}$, $n \neq n'$,

$$\int_{\mathbb{R}} \theta(t) (\text{op}_\varepsilon(a) \psi_n^\varepsilon(t, \cdot), \psi_{n'}^\varepsilon(t, \cdot)) \xrightarrow{\varepsilon \rightarrow 0} 0,$$

which implies that the Wigner measure of $\sum_{n \in \mathcal{N}} \psi_n^\varepsilon$ is the sum of the Wigner measures of the ψ_n^ε . We take $a \in \mathcal{C}_c^\infty(\mathbb{R}^{2d})$ and $\tilde{a} = (\varrho_n - \varrho_{n'})^{-1}a \in \mathcal{C}_c^\infty(\mathbb{R}^{2d})$; then, for all $t \in \mathbb{R}$,

$$\begin{aligned} (\text{op}_\varepsilon(a) \psi_n^\varepsilon(t, \cdot), \psi_{n'}^\varepsilon(t, \cdot)) &= (\text{op}_\varepsilon(\tilde{a}) \varrho_n(\varepsilon D_x) \psi_n^\varepsilon(t, \cdot), \psi_{n'}^\varepsilon(t, \cdot)) \\ &\quad - (\text{op}_\varepsilon(\tilde{a}) \psi_n^\varepsilon(t, \cdot), \varrho_{n'}(\varepsilon D_x) \psi_{n'}^\varepsilon(t, \cdot)) + O(\varepsilon) \end{aligned}$$

from which we deduce:

$$(\text{op}_\varepsilon(a) \psi_n^\varepsilon(t, \cdot), \psi_{n'}^\varepsilon(t, \cdot)) = i\varepsilon^2 \frac{d}{dt} (\text{op}_\varepsilon(\tilde{a}) \psi_n^\varepsilon(t, \cdot), \psi_{n'}^\varepsilon(t, \cdot)) + O(\varepsilon).$$

Therefore, if $\theta \in \mathcal{C}_c^\infty(\mathbb{R})$, then

$$\begin{aligned} \int_{\mathbb{R}} \theta(t) (\text{op}_\varepsilon(a) \psi_n^\varepsilon(t, \cdot), \psi_{n'}^\varepsilon(t, \cdot)) dt &= O(\varepsilon) + i\varepsilon^2 \int_{\mathbb{R}} \theta(t) \frac{d}{dt} (\text{op}_\varepsilon(\tilde{a}) \psi_n^\varepsilon(t, \cdot), \psi_{n'}^\varepsilon(t, \cdot)) dt \\ &= O(\varepsilon) - i\varepsilon^2 \int_{\mathbb{R}} \theta'(t) (\text{op}_\varepsilon(\tilde{a}) \psi_n^\varepsilon(t, \cdot), \psi_{n'}^\varepsilon(t, \cdot)) dt = O(\varepsilon). \end{aligned}$$

□

APPENDIX A. SEMICLASSICAL PSEUDODIFFERENTIAL OPERATORS

In this appendix we recall a few basic notions on the theory of pseudodifferential operators that we use through this article. The reader can consult the references [1, 17, 21, 44, 50] for additional background and for proofs of the results that follow.

Recall that given a function $a \in C^\infty(\mathbb{R}^d \times \mathbb{R}^d)$ that is bounded together with its derivatives (we denote the space of all such functions by S), one defines the semiclassical pseudodifferential operator of symbol a obtained through the Weyl quantization rule to be the operator $\text{op}_\varepsilon(a)$ that acts on functions $f \in \mathcal{S}(\mathbb{R}^d)$ by:

$$\text{op}_\varepsilon(a)f(x) = \int_{\mathbb{R}^d \times \mathbb{R}^d} a\left(\frac{x+y}{2}, \varepsilon\xi\right) e^{i\xi \cdot (x-y)} f(y) dy \frac{d\xi}{(2\pi)^d}.$$

These operators are bounded in $L^2(\mathbb{R}^d)$. The Calderón-Vaillancourt theorem [13] ensures the existence of a constant $C_d > 0$ such that for every $a \in S$ one has

$$(A.1) \quad \|\text{op}_\varepsilon(a)\|_{\mathcal{L}(L^2(\mathbb{R}^d))} \leq C_d N(a),$$

where

$$N_d(a) := \sum_{\alpha \in \mathbb{N}^{2d}, |\alpha| \leq J_0} \sup_{\mathbb{R}^d \times \mathbb{R}^d} |\partial_{x,\xi}^\alpha a|$$

for some $J_0 \in \mathbb{N}$ depending only on d . We make use repeatedly of the following result, known as the symbolic calculus for pseudodifferential operators.

Proposition A.1. *Let $a, b \in S$, then*

$$\text{op}_\varepsilon(a)\text{op}_\varepsilon(b) = \text{op}_\varepsilon(ab) + \frac{\varepsilon}{2i}\text{op}_\varepsilon(\{a, b\}) + \varepsilon^2 R_\varepsilon^{(2)},$$

with $\{a, b\} = \nabla_\xi a \cdot \nabla_x b - \nabla_x a \cdot \nabla_\xi b$ and

$$[\text{op}_\varepsilon(a), \text{op}_\varepsilon(b)] = \frac{\varepsilon}{i}\text{op}_\varepsilon(\{a, b\}) + \varepsilon^3 R_\varepsilon^{(3)},$$

$$\|R_\varepsilon^{(j)}\|_{\mathcal{L}(L^2(\mathbb{R}^d))} \leq C \sup_{|\alpha|+|\beta|=j} N_d(\partial_\xi^\alpha \partial_x^\beta a) N_d(\partial_\xi^\beta \partial_x^\alpha b), \quad j = 1, 2,$$

for some constant $C > 0$ independent of a, b and ε .

APPENDIX B. TRACE OPERATOR-VALUED MEASURES

In this appendix we recall general considerations on operator-valued measures. Let X be a complete metric space and (Y, σ) a measure space; write $\mathcal{H} := L^2(Y, \sigma)$ and denote by $\mathcal{L}^1(\mathcal{H})$, $\mathcal{K}(\mathcal{H})$ and $\mathcal{L}(\mathcal{H})$ the spaces of trace-class, compact and bounded operators on \mathcal{H} respectively. A trace-operator valued Radon measure on X is a linear functional:

$$M : \mathcal{C}_0(X) \longrightarrow \mathcal{L}^1(\mathcal{H})$$

satisfying the following boundedness condition. For every compact $K \subset X$ there exist a constant $C_K > 0$ such that:

$$\text{Tr} |M(\phi)| \leq C_K \sup_K |\phi|, \quad \forall \phi \in \mathcal{C}_0(K).$$

Such an operator-valued measure is positive if for every $\phi \geq 0$, $M(\phi)$ is an Hermitian positive operator. Let M be a positive trace operator-valued measure on X , denote by $\nu \in \mathcal{M}_+(X)$ the positive real measure defined by:

$$\int_X \phi(x) \nu(dx) = \text{Tr } M(\phi), \quad \forall \phi \in \mathcal{C}_0(X).$$

The Radon-Nikodym theorem for operator valued measures (see, for instance, the appendix in [26]) ensures the existence of a ν -locally integrable function:

$$Q : X \mapsto \mathcal{L}^1(\mathcal{H}), \quad \text{Tr } Q(x) = 1, \quad Q(x) \text{ positive Hermitian for } \nu\text{-a.e. } x \in X,$$

such that:

$$M(\phi) = \int_X \phi(x) Q(x) \nu(dx), \quad \forall \phi \in \mathcal{C}_0(X).$$

Note that this formula implies that M can be identified to a positive element of the dual of $\mathcal{C}_0(X; \mathcal{K}(\mathcal{H}))$ via:

$$\langle M, T \rangle \equiv \int_X \text{Tr}[T(x)M(dx)] := \int_X \text{Tr}(T(x)Q(x))\nu(dx), \quad T \in \mathcal{C}_0(X; \mathcal{K}(\mathcal{H})).$$

It can be also shown that every such positive functional arises in this way. Let $(e_j(x))_{j \in \mathbb{N}}$ denote an orthonormal basis of \mathcal{H} consisting of eigenfunctions of $Q(x)$:

$$Q(x)e_j(x) = \varrho_j(x)e_j(x), \quad \sum_{j=1}^{\infty} \varrho_j(x) = 1, \quad \nu\text{-a.e.}$$

Clearly, both ϱ_j and e_j , $j \in \mathbb{N}$, are locally ν -integrable and

$$Q(x) = \sum_{j=1}^{\infty} \varrho_j(x) |e_j(x)\rangle \langle e_j(x)|, \quad \nu\text{-a.e.},$$

where, as usual, $|e_j(x)\rangle \langle e_j(x)|$ denotes the orthogonal projection in \mathcal{H} onto $e_j(x)$. Moreover, as a consequence of the monotone convergence theorem, the following result easily follows.

Lemma B.1. *Let M be a positive trace operator-valued measure on X . Then there exist a non-negative function $\rho \in L^1_{\text{loc}}(X, \nu; L^1(Y, \sigma))$ such that, for every $a \in \mathcal{C}_0(X; L^\infty(Y, \sigma))$ one has:*

$$\int_X \text{Tr}[m_a(x)M(dx)] = \int_X \int_Y a(x, y) \rho(x, y) \sigma(dy) \nu(dx),$$

where $m_a(x)$ denotes the operator acting on \mathcal{H} by multiplication by $a(x, \cdot)$. The density ρ is given by:

$$\rho(x, y) = \sum_{j=1}^{\infty} \varrho_j(x) |e_j(x, y)|^2.$$

APPENDIX C. PROOF OF LEMMA 4.3

We denote by \mathcal{F}_ε the semi-classical Fourier transform defined for $f \in L^2(\mathbb{R}^d)$ by

$$\mathcal{F}_\varepsilon f(\xi) = (2\pi\varepsilon)^{-d/2} \widehat{f}\left(\frac{\xi}{\varepsilon}\right)$$

and we observe that for $a \in \mathcal{C}_c^\infty(\mathbb{R}^{3d})$,

$$(\text{op}_\varepsilon(a_\varepsilon)f, f) = (2\pi\varepsilon)^{-d} \int_{\mathbb{R}^{3d}} a_\varepsilon\left(-x, \frac{\xi + \xi'}{2}\right) e^{i_\varepsilon x \cdot (\xi - \xi')} \mathcal{F}_\varepsilon f(\xi') \overline{\mathcal{F}_\varepsilon f(\xi)} d\xi d\xi' dx,$$

where a_ε is associated with a according to (4.1). We consider a smooth cut-off function χ which is equal to 1 on the support of a so that we have $a(x, \xi)\chi(\xi) = a(x, \xi)$ and we write

$$\begin{aligned} (\text{op}_\varepsilon(a_\varepsilon)f, f) &= \\ & (2\pi\varepsilon)^{-d} \int_{\mathbb{R}^{3d}} a_\varepsilon\left(-x, \frac{\xi + \xi'}{2}\right) e^{i_\varepsilon x \cdot (\xi - \xi')} \mathcal{F}_\varepsilon f(\xi') \overline{\mathcal{F}_\varepsilon f(\xi)} \chi(\xi) \chi(\xi') d\xi d\xi' dx + O(\varepsilon). \end{aligned}$$

The rest term $O(\varepsilon)$ comes from Taylor formula close to $\frac{\xi + \xi'}{2}$, the observation that

$$(\xi_j - \xi'_j) e^{i_\varepsilon x \cdot (\xi - \xi')} = \frac{\varepsilon}{i} \partial_{x_j} \left(e^{i_\varepsilon x \cdot (\xi - \xi')} \right), \quad 1 \leq j \leq d,$$

and the use of integration by parts in x . Similarly, we just need to consider vectors (ξ, ξ') which are close to the diagonal and if we introduce a smooth function Θ compactly supported on $|\xi| \leq 1$ and equal to 1 close to 0, then for some $\delta > 0$ (that will be chosen small enough later), we have

$$\begin{aligned} (\text{op}_\varepsilon(a_\varepsilon)f, f) &= (2\pi\varepsilon)^{-d} \int_{\mathbb{R}^{3d}} a_\varepsilon\left(-x, \frac{\xi + \xi'}{2}\right) \\ & \times e^{i_\varepsilon x \cdot (\xi - \xi')} \mathcal{F}_\varepsilon f(\xi') \overline{\mathcal{F}_\varepsilon f(\xi)} \Theta\left(\frac{\xi - \xi'}{\delta}\right) \chi(\xi) \chi(\xi') d\xi d\xi' dx + O(\varepsilon). \end{aligned}$$

We are left with the integral

$$\begin{aligned} I_\varepsilon &= (2\pi\varepsilon)^{-d} \int_{\mathbb{R}^{3d}} a_\varepsilon\left(-x, \frac{\xi + \xi'}{2}\right) e^{i_\varepsilon x \cdot (\xi - \xi')} \mathcal{F}_\varepsilon f(\xi') \overline{\mathcal{F}_\varepsilon f(\xi)} \chi(\xi) \chi(\xi') \\ & \times \Theta\left(\frac{\xi - \xi'}{\delta}\right) d\xi d\xi' dx \\ &= (2\pi\varepsilon)^{-d} \int_{\mathbb{R}^{3d}} a_\varepsilon\left(-x, \frac{\Phi(\zeta) + \Phi(\zeta')}{2}\right) e^{i_\varepsilon x \cdot (\Phi(\zeta) - \Phi(\zeta'))} \mathcal{F}_\varepsilon f(\Phi(\zeta')) \\ & \times \overline{\mathcal{F}_\varepsilon f(\Phi(\zeta))} J_\Phi(\zeta) J_\Phi(\zeta') \chi \circ \Phi(\zeta) \chi \circ \Phi(\zeta') \Theta\left(\frac{\Phi(\zeta) - \Phi(\zeta')}{\delta}\right) d\zeta d\zeta' dx \end{aligned}$$

where $\zeta \mapsto J_\Phi(\zeta)$ is the Jacobian of the diffeomorphism Φ . Setting

$$\zeta = \theta + \varepsilon \frac{v}{2} \quad \text{and} \quad \zeta' = \theta - \varepsilon \frac{v}{2},$$

we have for $t \in \mathbb{R}$,

$$\Phi(\theta + \varepsilon tv) = \Phi(\theta) + \varepsilon t d\Phi(\theta)v + \varepsilon^2 \int_0^1 d^2\Phi(\theta + \varepsilon tsv)[v, v](1-s) ds,$$

whence

$$\begin{aligned}\frac{1}{2}(\Phi(\zeta) + \Phi(\zeta')) &= \Phi(\theta) + \frac{\varepsilon^2}{2}B_\varepsilon^+(\theta, v)[v, v], \\ \Phi(\zeta) - \Phi(\zeta') &= \varepsilon d\Phi(\theta)v + \varepsilon^2B_\varepsilon^-(\theta, v)[v, v],\end{aligned}$$

with

$$B_\varepsilon^\pm(\theta, v) = \int_0^1 d^2 \left(\Phi \left(\theta + \varepsilon s \frac{v}{2} \right) \pm \Phi \left(\theta - \varepsilon s \frac{v}{2} \right) \right) (1-s) ds.$$

Note that the functions B_ε^\pm are smooth, bounded and with bounded derivatives, uniformly in ε , as soon as the variables θ and εv are in a compact. We obtain

$$\begin{aligned}I_\varepsilon &= (2\pi)^{-d} \int_{\mathbb{R}^{3d}} a_\varepsilon \left(-x, \Phi(\theta) + \frac{\varepsilon^2}{2}B_\varepsilon^+(\theta, v)[v, v] \right) e^{ix \cdot (d\Phi(\theta)v + \varepsilon B_\varepsilon^-(\theta, v)[v, v])} \\ &\quad \times \mathcal{F}_\varepsilon f \left(\Phi \left(\theta - \varepsilon \frac{v}{2} \right) \right) \overline{\mathcal{F}_\varepsilon f} \left(\Phi \left(\theta + \varepsilon \frac{v}{2} \right) \right) J_\Phi \left(\theta + \varepsilon \frac{v}{2} \right) J_\Phi \left(\theta - \varepsilon \frac{v}{2} \right) d\theta d\theta' dx,\end{aligned}$$

where we have omitted the localization functions in $\theta + \varepsilon \frac{v}{2}$ and $\theta - \varepsilon \frac{v}{2}$, which makes that the integral is compactly supported in θ and εv . Moreover, we have $\varepsilon|v| \leq \delta$ on the domain of integration. We shall crucially use this information later.

The change of variable $x = {}^t d\Phi(\theta)^{-1}u$ gives

$$\begin{aligned}I_\varepsilon &= (2\pi)^{-d} \int_{\mathbb{R}^{3d}} a_\varepsilon \left(-{}^t d\Phi(\theta)^{-1}u, \Phi(\theta) + \frac{\varepsilon^2}{2}B_\varepsilon^+(\theta, v)[v, v] \right) \\ &\quad \times e^{iu \cdot v + i\varepsilon u \cdot {}^t d\Phi(\theta)^{-1}B_\varepsilon^-(\theta, v)[v, v]} \mathcal{F}_\varepsilon f \left(\Phi \left(\theta - \varepsilon \frac{v}{2} \right) \right) \overline{\mathcal{F}_\varepsilon f} \left(\Phi \left(\theta + \varepsilon \frac{v}{2} \right) \right) \\ &\quad \times J_\Phi \left(\theta + \varepsilon \frac{v}{2} \right) J_\Phi \left(\theta - \varepsilon \frac{v}{2} \right) J_\Phi^{-1}(\theta) d\theta d\theta' du,\end{aligned}$$

with the same property on the domain of integration (θ in a compact and $\varepsilon|v| < \delta$). Note that

$$\begin{aligned}&a_\varepsilon \left(-{}^t d\Phi(\theta)^{-1}u, \Phi(\theta) + \frac{\varepsilon^2}{2}B_\varepsilon^+(\theta, v)[v, v] \right) \\ &= a \left(-{}^t d\Phi(\theta)^{-1}u, \Phi(\theta) + \frac{\varepsilon^2}{2}B_\varepsilon^+(\theta, v)[v, v], \right. \\ &\quad \left. \frac{1}{\varepsilon}B \left(\Phi(\theta) + \frac{\varepsilon^2}{2}B_\varepsilon^+(\theta, v)[v, v] \right) \varphi \left(\Phi(\theta) + \frac{\varepsilon^2}{2}B_\varepsilon^+(\theta, v)[v, v] \right) \right) \\ &= a \left(-{}^t d\Phi(\theta)^{-1}u, \Phi(\theta), B(\Phi(\theta)) \frac{\theta'}{\varepsilon} \right) + \varepsilon r_\varepsilon^{(2)}(\theta, u, v)[v, v].\end{aligned}$$

The matrix $r_\varepsilon^{(2)}$ is supported in a compact independent of ε in the variables (u, θ) . Besides, the matrix $r_\varepsilon^{(2)}$ is smooth, bounded, and with bounded derivatives, uniformly in ε , as soon as the variable εv is in a compact, which is the case on the domain of integration of the integral I_ε . Using Taylor formula on the Jacobian terms, we write

$$\begin{aligned}&a_\varepsilon \left(-{}^t d\Phi(\theta)^{-1}u, \Phi(\theta) + \frac{\varepsilon^2}{2}B_\varepsilon^+(\theta, v)[v, v] \right) J_\Phi \left(\theta + \varepsilon \frac{v}{2} \right) J_\Phi \left(\theta - \varepsilon \frac{v}{2} \right) \\ &= a \left(-{}^t d\Phi(\theta)^{-1}u, \Phi(\theta), B(\Phi(\theta)) \frac{\theta'}{\varepsilon} \right) J_\Phi(\theta)^2 + \varepsilon r_\varepsilon^{(2)}(\theta, u, v)[v, v] + \varepsilon r_\varepsilon^{(1)}(\theta, u, v) \cdot v,\end{aligned}$$

where the vector $r_\varepsilon^{(1)}$ is supported in a compact independent of ε in the variables (u, θ) and, as $r_\varepsilon^{(2)}$, is smooth, bounded, and with bounded derivatives, uniformly in ε on the domain of integration of the integral I_ε (where θ is in a compact and $\varepsilon|v| \leq \delta$, δ to be chosen later).

Denote by \mathcal{U}_ε the isometry of $L^2(\mathbb{R}^d)$:

$$f^\varepsilon \mapsto J_\Phi(\cdot)^{\frac{d}{2}} \mathcal{F}_\varepsilon f(\Phi(\cdot)),$$

then

$$(\text{op}_\varepsilon(a_\varepsilon)f, f) = (\text{op}_\varepsilon(\tilde{a}_\varepsilon)\mathcal{U}_\varepsilon f, \mathcal{U}_\varepsilon f) + \varepsilon (R_\varepsilon \mathcal{U}_\varepsilon f, \mathcal{U}_\varepsilon f),$$

with

$$\tilde{a}_\varepsilon(u, \theta) = a\left(-{}^t d\Phi(\theta)^{-1}u, \Phi(\theta), B(\Phi(\theta))\frac{\theta'}{\varepsilon}\right),$$

and where R_ε is the operator of kernel

$$(\theta, \theta') \mapsto (2\pi\varepsilon)^{-d} K_\varepsilon\left(\frac{\theta + \theta'}{2}, \frac{\theta - \theta'}{\varepsilon}\right),$$

with $K_\varepsilon = K_\varepsilon^{(1)} + K_\varepsilon^{(2)}$,

$$\begin{aligned} K_\varepsilon^{(1)}(\theta, v) &= \int_{\mathbb{R}^d} \left(r_\varepsilon^{(1)}(\theta, u, v) \cdot v + r_\varepsilon^{(2)}(\theta, u, v)[v, v] \right) e^{iu \cdot v + i\varepsilon u \cdot {}^t d\Phi(\theta)^{-1} B_\varepsilon^-(\theta, v)[v, v]} du, \\ K_\varepsilon^{(2)}(\theta, v) &= \int_{\mathbb{R}^d} a_\varepsilon\left({}^t d\Phi(\theta)^{-1}u, \Phi(\theta) + \frac{\varepsilon^2}{2} B_\varepsilon^+(\theta, v)[v, v]\right) e^{iu \cdot v} \\ &\quad \times \frac{1}{\varepsilon} \left[e^{i\varepsilon u \cdot {}^t d\Phi(\theta)^{-1} B_\varepsilon^-(\theta, v)[v, v]} - 1 \right] du. \end{aligned}$$

The proof concludes by Schur lemma and the next result.

Lemma C.1. *Let us fix δ small enough. Then, for any $j \in \{1, 2\}$, there exists a constant $C_j > 0$ such that for all $\varepsilon > 0$,*

$$\int_{\mathbb{R}^d} \sup_{\theta \in \mathbb{R}^d} |K_\varepsilon^{(j)}(\theta, v)| dv \leq C_j.$$

Indeed, by this Lemma, we obtain that for all $\varepsilon > 0$

$$\begin{aligned} (2\pi\varepsilon)^{-d} \int_{\mathbb{R}^d} \sup_{\theta \in \mathbb{R}^d} \left| K_\varepsilon\left(\frac{\theta + \theta'}{2}, \frac{\theta - \theta'}{\varepsilon}\right) \right| d\theta' &= (2\pi)^{-d} \int_{\mathbb{R}^d} \sup_{\theta \in \mathbb{R}^d} |K_\varepsilon(\theta - \varepsilon v, v)| dv \\ &\leq (2\pi)^{-d} \int_{\mathbb{R}^d} \sup_{\theta \in \mathbb{R}^d} |K_\varepsilon(\theta, v)| dv \leq C_1 + C_2, \end{aligned}$$

and similarly

$$\begin{aligned} (2\pi\varepsilon)^{-d} \int_{\mathbb{R}^d} \sup_{\theta' \in \mathbb{R}^d} \left| K_\varepsilon\left(\frac{\theta + \theta'}{2}, \frac{\theta - \theta'}{\varepsilon}\right) \right| d\theta &= (2\pi)^{-d} \int_{\mathbb{R}^d} \sup_{\theta' \in \mathbb{R}^d} |K_\varepsilon(\theta' + \varepsilon v, v)| dv \\ &\leq (2\pi)^{-d} \int_{\mathbb{R}^d} \sup_{\theta' \in \mathbb{R}^d} |K_\varepsilon(\theta', v)| dv \leq C_1 + C_2, \end{aligned}$$

By Schur Lemma, these two inequalities yield the boundedness of R_ε uniformly in ε as an operator on $L^2(\mathbb{R}^d)$.

Let us now prove Lemma C.1.

Proof. Note first that the functions $K_\varepsilon^{(j)}$ are compactly supported in the variable θ , uniformly in ε . We are going to prove that for any $N > 0$, there exists a constant $C_{N,j}$ such that, for $|v| > 1$,

$$(1 + |v|^2)^N \left| K_\varepsilon^{(j)}(\theta, v) \right| \leq C_{N,j}.$$

These inequalities are enough to conclude as in the lemma. For proving these inequalities, we crucially use that the domain of integration in u is compact and we shall gain the decrease in v by using the oscillations inside the integral.

Let us first focus on $K_\varepsilon^{(1)}$. Since θ is in a compact and B_ε^- is bounded, we have

$$|v + \varepsilon {}^t d\Phi(\theta)^{-1} B_\varepsilon^-(\theta, v)[v, v]| \geq |v| - M\delta|v|$$

for some constant M . Therefore, if $\delta M < 1/2$, we have

$$|v + \varepsilon {}^t d\Phi(\theta)^{-1} B_\varepsilon^-(\theta, v)[v, v]| > \frac{1}{2}|v|,$$

and, for $|v| > 1$, integration by parts give

$$K_\varepsilon^{(1)}(\theta, v) = \int_{\mathbb{R}^d} |v + \varepsilon {}^t d\Phi(\theta)^{-1} B_\varepsilon^-(\theta, v)[v, v]|^{-2N} \left(\Delta_u^N r_\varepsilon^{(1)}(\theta, u, v) \cdot v + \Delta_u^N r_\varepsilon^{(2)}[v, v] \right) \times e^{iu \cdot v + i\varepsilon u \cdot {}^t d\Phi(\theta)^{-1} B_\varepsilon^-(\theta, v)[v, v]} du.$$

Since $r_\varepsilon^{(1)}$ and $r_\varepsilon^{(2)}$ have smooth compactly supported derivatives in u , uniformly bounded in ε , we obtain the existence of a constant $C_{N,1}$ such that

$$|K_\varepsilon^{(1)}(\theta, v)| \leq |v|^{-2N} C_{N,1}.$$

Let us now study $K_\varepsilon^{(2)}$ that we turn into

$$K_\varepsilon^{(2)}(\theta, v) = i \int_0^1 \int_{\mathbb{R}^d} u {}^t d\Phi(\theta)^{-1} B_\varepsilon^-(\theta, v) a_\varepsilon \left({}^t d\Phi(\theta)^{-1} u, \Phi(\theta) + \frac{\varepsilon^2}{2} B_\varepsilon^+(\theta, v)[v, v] \right) \times e^{iu \cdot v + i\varepsilon u \cdot {}^t d\Phi(\theta)^{-1} B_\varepsilon^-(\theta, v)[v, v]} du dt.$$

Once written on this form, one can see that the arguments developed for $K_\varepsilon^{(1)}$ apply again since the function

$$u \mapsto u {}^t d\Phi(\theta)^{-1} B_\varepsilon^-(\theta, v) a_\varepsilon \left({}^t d\Phi(\theta)^{-1} u, \Phi(\theta) + \frac{\varepsilon^2}{2} B_\varepsilon^+(\theta, v)[v, v] \right)$$

is compactly supported in the variable u , smooth and bounded with derivatives that are bounded uniformly in ε . \square

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