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Spectral Signature of a Ru(II, III, IV) Complex: A Combined Experimental and Theoretical Investigation

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Abstract

The electronic and geometrical features of a recently synthesized ruthenium β -diketonato complex with protected ethynyl ligand ([Ru(dbm)₂(acac-TIPSA)], dbm = dibenzoylmethanate ion, acac-TIPSA = 3-(Trisopropylsilyl)acetylene-2,4-pentanedionate ion), building block for molecular spintronic devices, were analyzed as a function of its oxidation degree using Density Functional Theory (DFT) and Time Dependent DFT (TD-DFT), for ground and excited state properties, respectively, and compared to the available experimental data.

The UV-Vis spectral signature of native Ru(III) neutral form computed at TD-DFT level was compared to those of the oxidized Ru(IV) and reduced Ru(II) ones, allowing to clearly identify the nature of all bands occurring in this spectral region and thus providing important photophysical insights for the use of this complex as building block in polynuclear organometallic mixed valence molecular systems able to perform a swapping operation under light control.

1. Introduction

These last years, molecular spintronics ¹ has attracted the interest of many scientists emerging as an hot topic in fundamental science in particular for its potentialities related to quantum computing. ^{2, 3} Several molecular or atomic spin systems have been proposed to construct quantum gates. ⁴⁻⁹

In this context, a challenge is represented by the conception of single molecules acting as a calculating unit by themselves, that is as devices at the molecular level. ¹⁰

Very recently, one of us conceived a generic 4-center organometallic mixed valence molecular system able to perform a swapping operation controlled by light. ¹¹ This 4-center system is made of a tetranuclear ruthenium system Ru(III)₃Ru(IV). Several mononuclear precursors corresponding to the building blocks necessary to construct this target molecule have been already synthesized and characterized in detail. ¹²⁻¹⁵ Among those, a new ruthenium β -diketonato complex with protected ethynyl ligand [Ru(dbm)₂(acac-TIPSA)] (dbm = dibenzoylmethanate ion, acac-TIPSA = 3-(Trisopropylsilyl)acetylene-2,4-pentanedionate ion), schematically depicted in Scheme 1, has been recently synthesized and fully characterized at the experimental level. ¹⁵



Scheme 1. Schematic representation and labeling scheme of the systems analyzed.

The stable form of this complex is neutral, with the ruthenium atom in the oxidation state +3 since dbm and acac-TIPSA are mono-anionic ligands. This complex can be reversibly reduced or oxidized by electrolysis at constant potential (-1V vs SCE for the reduction and +1.25V vs SCE for the

oxidation). For each of these forms the UV-Vis data were collected and they are reported in Figure 1. More details about the experimental UV-Vis data can be found in our previous work.¹⁵



Figure 1. Experimental spectra recorded in CH_2Cl_2 for the different [Ru(dbm)₂(acac-TIPSA)] forms: native/Ru(III) (green line), reduced/Ru(II) (blue line) and oxidized/Ru(IV) (red line). Reduced and oxidized species have been obtained by electrolysis at constant potentials (-1 V/SCE and +1.25 V/SCE, respectively) in CH_2Cl_2 , TBAH 0.1M.

The goal of the present contribution is to clearly describe the electronic structure of the different forms of the title compound and to assign all main bands which appear in the UV-Vis spectra with the aim of identifying the spectral signature of the degree of oxidation of the ruthenium atom within the complex. This information will be of relevance when aiming at understanding or predicting the photophysical behavior of higher nuclearity compounds that shall be used as molecular spintronic devices.

Furthermore, an accurate knowledge of the electronic structure of the mononuclear ruthenium complex at the ground state is particularly important to better interpret the STM image of this single molecule on a surface, technique which may be used to better characterize the complexes whose one aims to use them as single molecule devices.

In this respect, methods rooted on Density Functional Theory (DFT) and on Time dependent DFT (TD-DFT) appear as particularly suited in order to get a clear-cut picture of the electronic structure of

both the ground and the vertical excited states, and, thus, to get the necessary understanding of the electronic features of these rather complex molecular systems.

Several works have already demonstrated the capability and the accuracy of DFT for the description of the electronic and geometrical features of complexes containing transition metal ions and, more particularly of Ru complexes. As a matter of fact several studies reported in literature, both fully theoretical and mixed theoretical and experimental, focused on Ru(II) complexes of technological (such as dyes for Dye Sensitized Solar Cells)¹⁶⁻¹⁸, biological (such as DNA intercalators)^{19, 20} and fundamental interest (such as compounds for photoinduced charge separation) $^{21-24}$, testified that, using sufficiently accurate exchange correlation potentials, a quantitative description of the experimental photophysical properties can be obtained. In particular, hybrid exchange and correlation functionals appeared as particularly well suited to obtained reliable structural and electronic features of Ru complexes at the ground state.²⁵ Analogously several recent works demonstrated the remarkable accuracy of TD-DFT in conjunction with the same exchange and correlation functionals for the description of the spectral features of Ru(II) containing metal complexes and, especially of the transitions of Metal to Ligand Character (MLCT) with limited though space character occurring in the visible spectral range. ²⁶ Moreover, the use of hybrid exchange and correlation functionals allows for a balanced description of MLCT and Ligand Centered (LC) states thus allowing for a trustable assignment of intermingled experimental spectra or the use of TD-DFT for the in-silico design of new systems.

Since TD-DFT methods and computationally more expensive ab-initio approaches, such as post-Hartree Fock methods, provide a similar description of the photophysical properties of Ru containing complexes, not surpringly the former are currently extensively used for the investigation of the electronic properties at both the ground and the excited state of large d-transition metal complexes.²⁵⁻²⁶ Remarkably, theoretical works concerning Ru(III) and Ru(IV) complexes are less numerous than those devoted to Ru(II) complexes and mainly related to the investigation of catalytically active compounds (both of synthetic or biomimetic relevance).²⁷⁻³¹

The paper is structured as follows: after a brief description of the computational approaches used (section 2), the electronic and geometrical features are discussed in section 3.1 together with the absorption properties of each compound (section 3.2). Finally, some general conclusions, concerning the accuracy of the methods and the spectroscopic signature of the Ru centers, are drawn in last section.

2. Computational details

All calculations were performed with the Gaussian09 program package. ³² The hybrid PBE0 functional³³ casting 25% of Hartree-Fock exchange in the GGA PBE functional,^{34, 35} was used throughout both for structural optimization and for excited state calculations performed at TD-DFT level. For both structural optimization and excited state calculations, all atoms were described using a

double zeta quality basis set (LANL2DZ).³⁶ For the Ru atom the same basis set and the corresponding pseudopotential were applied.³⁷⁻³⁹

In the case of the Ru(IV) complex, structural optimization were performed both for the singlet and triplet spin state. This last one (triplet) resulted as the most stable spin state (of roughly 10 kcal/mol). Optical transitions were computed in order to cover the full UV-Vis range. Simulated spectra reported in Figures were obtained from the computed vertical transition energies and oscillator strengths by a using a convolution with Gaussian functions (half height width: 40 nm) with Gabedit software. ⁴⁰ The MO drawings were obtained with GaussView 5.0.⁴¹

Solvent effects (dichloromethane, DCM) were included using the Polarizable Continuum Model (PCM) both for ground and excited state calculations, as available in the Gaussian code (CPCM) ^{42, 43} only in the case of the Ru(II) complex. Structural and absorption data computed at this level of theory are reported in Supported Information for comparison purposes.

3. Results and discussion

3.1 Geometrical and electronic structural features

The computed geometrical parameters for Ru complex in its native Ru(III), oxidized Ru(IV) and reduced Ru(II) form together with the experimental Xray data available for the Ru(III) species are reported in Table 1, while the labeling scheme adopted is reported in Scheme 1.

First of all, one can note a fairly good agreement between the computed and the experimental data available for the Ru(III) complex. The coordination sphere of the Ru(III) is rather well reproduced by the calculations since the difference between computed and experimental Ru-O bonds is always below 2% (that is 0.03 Å). Furthermore, the shorter Ru-O bonds experimentally observed in the case of the oxygen atoms of the acac-TIPSA ligand (Ru-O₅ and Ru-O₆ in Table 1) are overall correctly reproduced by the calculations. As expected the Ru ion keeps a pseudo octahedral coordination, only slightly distorted due the presence of non-equivalent ligands. The computed distances are also in rather good agreement with the experimental ones previously reported in literature for similar compounds.^{14, 15, 44-51} From a structural point of view, not surprisingly oxidation seems mainly to affect Ru coordination sphere, following the expected contraction of the Ru-O bonds when one goes from Ru(II)_{S=0} to Ru(IV)_{S=1}, the average Ru-O bond length varying from 2.042 Å (computed for Ru(II)) to 1.986 Å (computed for Ru(IV)). On the other hand, since the oxidation is essentially centered on the metal, the ligands are not subject to very relevant structural changes, especially for what concerns the TIPSA moiety.

	Ru(II)	Ru(III)	Ru(III)	Ru(IV)	Ru(IV)
Ru ₁ -O ₄	2.039	2.007	1.994	1.988	1.986
Ru_1-O_1	2.039	2.008	2.003	2.015	1.986
Ru_1-O_3	2.033	2.027	1.998	1.983	1.976
Ru_1-O_2	2.033	2.03	2.024	1.951	1.976
Ru_1-O_6	2.053	2.019	1.998	2.02	1.997
Ru_1-O_5	2.053	2.019	1.996	1.966	1.996
<ru-o> average</ru-o>	2.042	2.018	2.002	1.987	1.986
C≡C	1.237	1.233	1.208	1.235	1.234
O_6 - Ru_1 - O_5	90.2	91.5	92.7	87.0	86.9
O_4 - Ru_1 - O_3	92.2	89.9	91.9	93.3	88.9
O_1 - Ru_1 - O_2	92.1	89.6	91.9	87.9	88.9
O_1 - Ru_1 - O_6	89.6	89.3	92.7	89.9	90.7
O_5 - Ru_1 - O_4	88.9	88.3	88.5	87.5	90.9
O_3 - Ru_1 - O_2	89.2	90.1	92.9	88.7	89.5

 Table 1 Computed geometrical parameters (bond lengths in Å, angles in degrees) together with the available experimental data.

The MO isodensity surfaces and energies computed for the d^6 (i.e. reduced Ru(II)) species are reported in Figure 2.

The computed orbital splitting corresponds to what expected for a Ru(II) ion in octahedral environment : the three highest occupied MOs are very closely lying in energy and display a large contribution on the metal center. These orbitals basically correspond to the t_{2g} set of the d orbitals of a d^6 ion in octahedral environment while the loss of perfect degeneracy is due to the presence of different ligands in the Ru coordination sphere. Below these orbitals, a doubly occupied MO with large contribution from atomic orbitals centered on the TIPSA moiety is found, thus underlying the donor character of this ligand. On the other hand the lowest lying unoccupied orbitals (from the LUMO to the LUMO+3) show contributions mainly arising from the dbm moieties.



Figure 2. Orbital energies (in a.u.) and isodensity plots (contour value 0.025 a.u.) computed for the reduced [Ru(II)(dbm)₂(acac-TIPSA)]⁻ species.

This orbital picture doesn't change a lot in the case of the native Ru(III) species (whose MO diagram is reported in Supporting Information, Figure SI.1) although the metal character of the highest occupied MO decreases as a consequence of the mainly metal centered oxidation that stabilize the metal centered MO. The fact that redox processes take place essentially at the Ru atom is qualitatively confirmed by the inspection of the Mulliken charges (and Mulliken Atomic Spin population in the case of open shell systems, that is the Ru(III) and Ru(IV)_{S=1}). Indeed, computed charges varies from 0.44 in the case of the Ru(II) species to 0.62 and 0.82 for the Ru(III) and Ru(IV), species, respectively. The localization of the hole(s) on the Ru atom for the Ru(III) complex after its oxidation from the Ru(II) d⁶ closed shell electronic configuration, is confirmed by the computed Mulliken Atomic Spin Density (0.84 e-) actually corresponding to the presence of an unpaired electron on the Ru atom. Analogously the oxidation from Ru(III) to Ru(IV) seems mainly to take place on the Ru center, the spin density computed on this latter being, for the Ru(IV)_{S=1} complex, of 1.25 e-.

Oxidation of the Ru(III) species yields to the Ru(IV) species for which the intermediate spin triplet state (S=1) is indeed computed to be more stable than the corresponding singlet state of roughly 10 kcal/mol.

Several examples of six-coordinated Ru(IV) complexes displaying either a diamagnetic or a paramagnetic ground state exist in literature.²⁷⁻³¹ The recent theoretical paper of Gonzalez and collaborators shows how DFT is indeed able to provide a realistic picture of the spin state ordering in the case of Ru(IV) compounds.²⁷ The results obtained in this work are coherent with the conclusions of Gonzalez work concerning the stabilization of the intermediate spin state (S=1) in presence of O or S containing ligands and to the MO diagram computed for the Ru(II) species (Figure 2) displaying to pseudo-degenerate orbitals of $d\pi$ character as highest occupied MOs. A careful analysis of MOs (Figure SI.2) computed for the Ru(IV) compound actually shows that the highest occupied MOs display a more marked ligand character, the d center orbitals being shifted more significantly to lower energies as a consequence of the metal centered oxidation.

More interestingly the presence of a diamagnetic (S=0) or paramagnetic (S=1) ground state is associated with important structural reorganization. In particular, the angle between O-Ru-O plane and O-C-C-C-O acac-TIPSA one is computed to be 29° for diamagnetic (S=0) compound while is reduced to 0° in the case of the (S=1) intermediate spin state (Figure SI.3). Analysis of computed Mulliken atomic charges qualitatively shows that in the case of the diamagnetic $Ru(IV)_{S=0}$ species the second oxidation involves to a slightly larger extent than in the case of the $Ru(IV)_{S=1}$ also the ligands (and in particular the acac-TIPSA moiety) which may explain the difference in structure computed for the $Ru(IV)_{S=1}$ and $Ru(IV)_{S=0}$ complexes.

3.2 Absorption properties

As previously reported, ¹⁵ the experimental spectra of the neutral, Ru(III), species is dominated by three intense bands, two in the UV domain (at 258 nm and 330 nm, respectively) and one in the visible range (at 413 nm), while a weaker band appears at lower energies (at 570 nm).

If reduction (and oxidation) of this species has small impact in the UV domain, the most intense bands being only slightly shifted in energy although being significantly changed in relative intensity, in the visible and near IR domains the redox signature of both oxidation and reduction is clear.

Upon reduction to Ru(II), a rise in absorbance in the visible range is observed due to the emergence of two new bands at 554 and 656 nm , respectively. On the other hand upon oxidation a broad band appears in the near-IR region, centered at 980 nm.

The experimentally measured spectra, together with the corresponding computed transition energies and simulated spectra for the reduced, native and oxidized species are reported in Figure 3-5, respectively.

An overall quantitative agreement between the observed and computed absorption energies can be remarked. The difference in energy between the maxima in experimental spectra and the computed transition energies are all below the 0.15 eV threshold, which is indeed considered the accuracy that

one can reasonably expect at this level of theory. Beside the relative intensity of the picks is well reproduced both within the same species as well as when comparing the complex at different degree of oxidation. Thus, we can conclude that TD-DFT calculations are able to capture all main spectral features related to the oxidation or reduction processes.



Figure 3. Experimental (black line) normalized absorption spectra of the reduced (Ru(II)) $[Ru(II)(dbm)_2(acac-TIPSA)]^{-1}$ complex together with the corresponding computed vertical transition energies and simulated spectra (red line)



Figure 4. Experimental (black line) normalized absorption spectra of the native (Ru(III)) $[Ru(III)(dbm)_2(acac-TIPSA)]^0$ complex together with the corresponding computed vertical transition energies and simulated spectra (red line)



Figure 5. Experimental (black line) normalized absorption spectra of the oxidized (Ru(IV)) $[Ru(IV)(dbm)_2(acac-TIPSA)]^{+1}$ complex together with the corresponding computed vertical transition energies and simulated spectra for Ru(IV)_{S=1} (red line)

The band experimentally measured at 330 nm in the case of the neutral compound (Figure 4, Ru(III)) is practically unchanged upon reduction (Figure 3) while slightly (roughly 10 nm) red shifted upon oxidation (Figure 5).

This band is indeed computed at 309 nm in the case of the Ru(III) species and red shifted at 316 and 329 nm for the Ru(II) and Ru(IV) complexes, respectively. Therefore both the sign and magnitude of the shifts observed upon both reduction and oxidation are almost quantitatively reproduced by the calculations.

Analyzing the nature molecular orbitals involved in all these transitions, these latter can be unambiguously assigned as ligand centred (LC) excitations, with contribution stemming from orbitals centered both on the dbm ligands and a non-negligible Inter Ligand Charge Transfer (ILCT) character associated to mono-electronic excitations from orbitals centered on the acac-TIPSA ligand to orbitals of the dbm moieties, which should, indeed, be not very sensitive to the redox processes occurring at the metal center.

On the other hand, since most of the transitions involving metal orbitals occur in the visible and near-IR spectral range it is not surprising that this latter contain the spectral signature of the oxidation degree of the Ru atom within the complex.

In the visible range, the neutral Ru(III) complex experimentally shows a rather broad band at 570 nm and a more intense band at 413 nm, while no absorption is present in the near-IR region. From our calculations (Figure 4) is rather clear that the lowest energy transition (computed at 580 nm in very good agreement with the experimental maximum observed at 570 nm) has a significant Ligand to Metal Charge Transfer Character (LMCT). From an orbital point of view, this transition mainly corresponds to a monoelectronic excitation from a doubly occupied orbital centered mainly on the acac and ethynyl moieties of the acac-TIPSA ligand to an essentially metal centered orbital.

Two 'classical' Metal to Ligand Charge Transfer Transitions (MLCT) are indeed computed at 399 and 378 nm in rather good agreement with the experimental spectra showing a single MCLT band at 413 nm, together with a number of other, less intense transitions always of MLCT character in the same spectral region. The oscillator strength associated to the most intense MLCT transitions is higher than that computed for the LMCT transition (0.05 vs 0.03 for the MCLTs and the LMCT, respectively), in agreement with the observed relative intensity of the bands associated to such transitions. The MLCT transitions occur from doubly occupied molecular orbitals centered on the metal to ligands (dbm) centered orbitals. Beside an overestimation of the absolute transition energy ascribable to the lack of solvent effects in our simulation, which is indeed still compatible with the level of theory here applied, most probably the energetic splitting between the two computed MLCT (of 0.174 eV) is slightly overestimated.

Upon reduction to a negatively charged Ru(II) complex (Figure 3), the LMCT band disappears and two transitions are computed in the visible range, at 500 and 582 nm, with an intensity ratio of 2:3.

These two transitions, of MCLT character involving mainly ligand orbitals centered on the dbm moieties fit very well the experimental spectra consisting, in this spectral region, of two transitions at 554 and 656 nm, respectively. Indeed, in the experimental spectra, a broad band in the near IR region (i.e. around 750 nm) is observed. This broad band can be ascribed to a HOMO-LUMO transition of MCLT character, which is computed at 748 nm with quite low intensity as depicted in Figure 3 or a to spin forbidden transition (S=0 to S=1) or to a combination of both. In order to qualitative estimate the energy of the spin forbidden transition, this latter was computed using a Δ SCF approach at 979 nm (in the gas phase) and 898 nm when simulating solvent effects (here DCM) by a continuum polarizable (PCM approach), ⁵² that is at significantly lower energies that what experimentally observed.

The band experimentally observed at 413 nm, and practically unaffected by reduction from Ru(III) to Ru(II), is still ascribed to a series of low intensity transitions of MLCT character computed around 400 nm and mostly involving dbm centered MOs as vacant orbitals.

Overall, one can thus conclude that reduction to the Ru(II) complex implies an increase of absorption in the visible range mostly associated to the rising of MCLT type transitions, while the UV part of the spectra, involving ILCT transitions, is not strongly perturbed.

The situation is clearly different when the Ru(III) complex is oxidized to yield the positively charged Ru(IV) species (Figure 5). Indeed, both from experimental and computed spectra, a drastic decrease of absorption in the visible range and increase in the near-IR region are observed. The decrease is associated to the absence of the MLCT transitions occurring at lower energy (that is the one computed around 600 nm for the Ru(II)) mainly due to the shift at lower energy of metal centered orbitals, as previously discussed for the Ru(III) and Ru(IV) species.

Indeed, the positively charged Ru(IV) species is characterized by two transitions occurring in the near-IR domain (at 761 and 1140 nm, respectively) both of marked LMCT character. These two transitions are both contributing to the broad band observed in the near-IR region with absorption maximum at around 980 nm. If both the computed transitions possess an LMCT character (Figure 5) the ligand centered orbitals involved are indeed significantly different. In particular, the most intense LMCT transition, that is the one computed at 1140 nm shows an acac-TIPSA to Ru character while the ligand orbitals involved in the higher energy LMCT (i.e. the one at 761 nm) are mostly centered on the dbm moieties. Finally, it is worth noting that the singlet Ru(IV) species, computed to be less stable of roughly 10kcal/mol , also shows a very similar spectrum (reported in Figure SI.5 and SI.6) to the triplet Ru(IV) complex only the low lying LMCT transition being perturbed as a consequence of the different oxidation pattern already discussed for the two species. Therefore, the concomitant presence of the S=0 and S=1 Ru(IV) species cannot be ruled out on the basis of the sole analysis of their spectral features.

4. Conclusions

The electronic and geometrical properties of a recently synthesized ruthenium β -diketonato complex with protected ethynyl ligand ([Ru(dbm)₂(acac-TIPSA)], dbm = dibenzoylmethanate ion, acac-TIPSA = 3-(Trisopropylsilyl)acetylene-2,4-pentanedionate ion) were interpreted as a function of its oxidation degree using Density Functional Theory (DFT) and Time Dependent DFT (TD-DFT), for ground and excited state properties, respectively, and compared to the available experimental data.

The UV-Vis spectral signature of native Ru(III) neutral form computed at TD-DFT level was compared to those of the oxidized (Ru(IV)) and reduced (Ru(II)) species. The results obtained show how the increase in absorption in the near-IR region observed upon oxidation of the native (Ru(III)) system is related to the presence of low lying LMCT states involving both the TIPSA and the dbm ligands.

LMCT transitions mainly involving the TIPSA ligand are indeed present also in the case of the native system although shifted to higher energy (in the visible region around 550 nm), while, as expected completely absent in the case of the reduced Ru(II) compound. A sizable absorption in the visible domain dominated by classical MLCT transitions is indeed observed for this latter (Ru(II)).

Finally, in the near UV domain all forms show a very similar spectral behavior dominated by LC transitions, not sensible to the metal degree of oxidation.

The clear identification of the nature of all bands occurring in these systems obtained by comparison of the present TD-DFT results with the experimental data provides important photophysical insights of relevance in the case of the use of this complex as building block in polynuclear organometallic mixed valence molecular systems able to perform a swapping operation under light control

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