



**HAL**  
open science

# Quantum ergodicity for large equilateral quantum graphs

Maxime Ingremeau, Mostafa Sabri, Brian Winn

► **To cite this version:**

Maxime Ingremeau, Mostafa Sabri, Brian Winn. Quantum ergodicity for large equilateral quantum graphs. *J. London Math. Soc.* (2), 2020, 101, pp.82-109. 10.1112/jlms.12259 . hal-01735912

**HAL Id: hal-01735912**

**<https://hal.science/hal-01735912>**

Submitted on 16 Mar 2018

**HAL** is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

# QUANTUM ERGODICITY FOR LARGE EQUILATERAL QUANTUM GRAPHS

MAXIME INGREMEAU, MOSTAFA SABRI, BRIAN WINN

ABSTRACT. Consider a sequence of finite regular graphs  $(G_N)$  converging, in the sense of Benjamini-Schramm, to the infinite regular tree. We study the induced quantum graphs with equilateral edge lengths, Kirchhoff conditions (possibly with a non-zero coupling constant  $\alpha$ ) and a symmetric potential  $U$  on the edges. We show that in the spectral regions where the infinite quantum tree has absolutely continuous spectrum, the eigenfunctions of the converging quantum graphs satisfy a quantum ergodicity theorem. In case  $\alpha = 0$  and  $U = 0$ , the limit measure is the uniform measure on the edges. In general, it has an explicit analytic density. We finally prove a stronger quantum ergodicity theorem involving integral operators, the purpose of which is to study eigenfunction correlations.

## 1. INTRODUCTION

Quantum ergodicity, one of the fundamental theorems of quantum chaos, is a result about spatial delocalization of eigenfunctions. In its original context [31, 17, 32], it says that the eigenfunctions of the Laplace-Beltrami operator on a compact Riemannian manifold become equidistributed on the manifold in the high energy limit, provided the geodesic flow is ergodic.

In their influential work [24], Kottos and Smilansky suggested that the ideas and results of quantum chaos should apply to quantum graphs. It is therefore natural to inquire whether quantum ergodicity holds on quantum graphs. By a quantum graph, we mean a metric graph, equipped with a differential operator and suitable boundary conditions at each vertex. We refer the reader to Section 2.1 for a more precise definition.

Let us briefly survey some earlier results of quantum ergodicity on quantum graphs.

The first investigation appeared in [9], where it was shown that the eigenfunctions of star graphs are not quantum ergodic in the high energy limit. More precisely, by studying the behaviour of the semi-classical measure when the number of bonds becomes large and the lengths approach a common value, one sees that it is not the uniform measure.

A positive result of quantum ergodicity later appeared in [10] for graphs associated to intervals maps. In this paper, instead of studying directly the high energy behaviour of the eigenfunctions  $\psi_n$  of the quantum graph, the authors study the eigenfunctions  $\phi_j(\lambda_n)$  of associated unitary operators  $U(\lambda_n)$  which encode the classical evolution. It is shown that both notions of quantum ergodicity (for  $\psi_n$  or  $\phi_j(\lambda_n)$ ) are intimately related if the size of the graph goes to infinity, see also [12] for a precise statement.

The previous result can serve as a motivation to consider *asymptotic quantum ergodicity* for general quantum graphs. In this case, one considers a sequence of quantum graphs, first studying the high energy behaviour, then taking the size of the graph to infinity. A general strategy for studying this regime was introduced in [20], and it is shown that the validity of asymptotic quantum ergodicity (AQE) depends on the spectral properties of the Markov transition matrix  $M_{b,b'}(\lambda) = |U_{b,b'}(\lambda)|^2$ . Namely, AQE should hold if the spectral gap of  $M(\lambda)$  does not decay too fast as the graph grows larger. This can be regarded as

---

2010 *Mathematics Subject Classification.* Primary 58J51. Secondary 34B45, 81Q10.

*Key words and phrases.* Quantum ergodicity, quantum graphs, large graphs, delocalization.

a “chaotic” assumption, similar to the ergodicity of the geodesic flow on manifolds. Note that the spectral gap of quantum star graphs closes quite rapidly.

It is natural however to study the high energy behaviour of a *fixed* compact quantum graph. This question was settled later in [18] in the case of Kirchhoff boundary conditions, where it is shown that quantum ergodicity does not hold for a generic metric on the graph, except if it is homeomorphic to an interval or a circle. The semi-classical measures were moreover characterized, and the “scars” were described.

To conclude this brief survey, we mention the paper [21] in which bounds on the entropy of the semi-classical measures were obtained for several families of quantum graphs. Finally, the paper [14] proved that asymptotic quantum ergodicity holds for sequences of quantum graphs without back-scattering, if the underlying discrete graphs are  $(q + 1)$ -regular expanders with few short loops. Under these “equi-transmitting” boundary conditions (which do not include the Kirchhoff conditions), the Markov operator  $M(\lambda)$  is just  $\frac{1}{q}\mathcal{B}$ , where  $(\mathcal{B}f)(b) = \sum_{b^+ \in \mathcal{N}_b^+} f(b^+)$  is the non-backtracking operator on the edges. It is known that if the graphs are expanders, then  $\frac{1}{q}\mathcal{B}$  has a uniform spectral gap [3, 2]. So this result of [14] confirms the general philosophy of [20] in this situation.

All the previous results were concerned with the high-energy behaviour of eigenfunctions, be it for a sequence of quantum graphs or a fixed graph. In this paper, we are interested in a different regime. Ultimately, we believe the following principle should be true : suppose that an infinite, possibly random, quantum tree has purely absolutely continuous spectrum in an interval  $I$ . Consider a sequence of quantum graphs converging to this tree in the sense of Benjamini-Schramm, and suppose the underlying discrete graphs are expanders. Then as the graphs  $\mathbf{G}_N$  grow large, the corresponding eigenfunctions whose energies *lie in the bounded interval  $I$*  will satisfy a quantum ergodicity theorem.

This kind of result can be interpreted as a delocalization result for the infinite quantum tree. It says that if the tree is *spectrally delocalized* (has pure AC spectrum in  $I$ ), then it also has a form of *spatial delocalization* (converging quantum graphs satisfy quantum ergodicity in  $I$ ).

This point of view originally stemmed from the study of quantum ergodicity for discrete graphs. The first incarnations of it appeared in the proof of quantum ergodicity for the adjacency matrix of a regular graph [3, 2, 15]. The question was later addressed in full generality in [4], in the framework of Schrödinger operators on graphs of bounded degrees. As an important application [5], it was shown that the eigenfunctions of the Anderson model on a large graph become equidistributed if the disorder is weak enough, in the region of AC spectrum previously provided by Klein [23]. This was among very few theorems of delocalization for the Anderson model, the opposite regime of localization being quite well understood today. More applications appeared in [6], where the first results of quantum ergodicity for non-regular graphs were given. We also mention the recent paper [26] which studies quantum ergodicity for sequences of compact hyperbolic surfaces, also in the bounded interval regime.

The present paper is a first step towards the proof of such a general criterion for quantum graphs. We study a simple family of quantum graphs : regular equilateral graphs, and endow them with the natural Kirchhoff conditions. We also allow for identical coupling constants  $\alpha$  on the vertices and identical symmetric potentials  $U$  on the edges. It is known [16] that the corresponding infinite tree has bands of AC spectra. We confirm the general principle by showing that eigenfunctions with energies in such bands are quantum ergodic, as the graph grow large. More precisely, when the graph  $\mathbf{G}_N$  is large enough, the probability measure  $|\psi_j^{(N)}(x)|^2 dx$  on  $\mathbf{G}_N$  approaches an explicit measure  $\frac{1}{N}\Psi_{\lambda_j}(x) dx$ , for most eigenfunctions  $\psi_j^{(N)}$ . In the special case  $U = \alpha = 0$ , the density  $\Psi_{\lambda_j}(x)$  is just a constant and we obtain the uniform measure  $\frac{2}{(q+1)N} dx$  on  $\mathbf{G}_N$  (see § 2.4 for more

precision). In general,  $\Psi_{\lambda_j}(x)$  is essentially the density of the spectral measure of the infinite quantum tree. The equidistribution in case  $U = \alpha = 0$  holds precisely because this spectral measure is constant in this case, as  $x$  varies along the tree.

Our paper is organized as follows. In Section 2 we introduce our quantum graphs and discuss the main results. In Section 3.1, we construct discrete functions on the directed bonds of  $G_N$ , which turn out to be eigenfunctions of a non-backtracking operator  $\mathcal{B}$ . We then prove our main result in Section 4. As in the case of combinatorial graphs [2, 3, 4], one can also ask about the behaviour of the eigenfunction correlators  $\overline{\psi_j^{(N)}(x)\psi_j^{(N)}(y)} dx dy$ . For this purpose, we provide a more general quantum ergodicity result, Theorem 2.3, involving integral operators. This result is proved in Section 5.

We plan to address the general criterion for quantum ergodicity in such a bounded interval regime in a future work. It should be expected that the proof will become much more difficult and technical in this case. A very interesting application would be to prove quantum ergodicity for weak random perturbations of the equilateral model we study here, whether in the edge lengths or the coupling constants. In fact, it was shown in [1] that the AC spectrum of the Laplace operator on the equilateral quantum tree  $\mathbf{T}_q$ , is stable under weak random perturbations of the edge lengths.

## 2. PRESENTATION OF THE RESULTS

**2.1. Background on quantum graphs.** Let  $G = (V, E)$  be a (combinatorial) graph, and fix  $L > 0$ . The associated equilateral metric graph, denoted by  $\mathbf{G}$ , is obtained by identifying each edge with  $[0, L]$ . In doing so, we fix an orientation on each  $e \in E$ , identifying 0 with an origin  $o_e \in V$  and  $L$  with a terminus  $t_e \in V$ . We shall write  $d_{\mathbf{G}}$  for the distance on  $\mathbf{G}$  induced by the usual distance on  $[0, L]$ .

A *quantum graph* is a self-adjoint differential operator acting on the Hilbert space  $\mathcal{H} := \bigoplus_{e \in E} L^2[0, L]$ . We consider the following Schrödinger operator.

Let  $U \in L^\infty(0, L)$  be a real-valued potential satisfying

$$U(L - x) = U(x).$$

We define an operator  $H_{\mathbf{G}}$  acting on  $\psi = (\psi_e)_{e \in E} \in \bigoplus_{e \in E} W^{2,2}(0, L) \subset \mathcal{H}$  by

$$(2.1) \quad (H_{\mathbf{G}}\psi_e)(x_e) = -\psi_e''(x_e) + U(x_e)\psi_e(x_e).$$

Thus defined,  $H_{\mathbf{G}}$  is not essentially self-adjoint: we need to impose suitable boundary conditions at each vertex. In this paper, we will consider the generalized *Kirchhoff boundary conditions* with parameter  $\alpha \in \mathbb{R}$  given by

- **Continuity:** For all  $v \in V$  and all  $e, e' \in E$ , we have  $\psi_e(t_e) = \psi_{e'}(o_{e'}) =: \psi(v)$  if  $t_e = o_{e'} = v$ .
- **Current conservation:** For all  $v \in V$ ,

$$\sum_{e: o_e=v} \psi_e'(0) - \sum_{e: t_e=v} \psi_e'(L) = \alpha\psi(v).$$

If  $G$  is finite,  $H_{\mathbf{G}}$  is self-adjoint on the domain of functions in  $\bigoplus_{e \in E} W^{2,2}(0, L)$  satisfying the Kirchhoff conditions. When  $G = \mathbb{T}_q$  is the infinite tree,  $H_{\mathbf{T}_q}$  is essentially self-adjoint on the set of compactly supported functions in  $\bigoplus_{e \in E} W^{2,2}(0, L)$  satisfying the Kirchhoff conditions [16], and can hence be uniquely extended to a self-adjoint operator. We refer the reader to [11] for more details on the construction and properties of quantum graphs.

In the sequel, we will consider sequences of finite graphs  $G_N = (V_N, E_N)$ , which we will suppose to be  $(q + 1)$ -regular for some  $q \geq 2$ , meaning that all vertices are connected to exactly  $q + 1$  neighbors. However, to lighten a bit notations and proofs, the parameters  $L$  and  $\alpha$ , as well as the potential  $U$ , will not depend on  $N$ .

*Eigenfunctions on the edges.* Fix an edge  $e \in E_N$  and consider the eigenproblem

$$(2.2) \quad -\psi_e'' + U\psi_e = \lambda\psi_e$$

for  $\lambda \in \mathbb{C}$  and  $\psi_e \in W^{2,2}(0, L)$ . On this edge  $e \equiv [0, L]$ , choose a basis of solutions  $C_\lambda(x)$ ,  $S_\lambda(x)$  of (2.2) satisfying

$$\begin{pmatrix} C_\lambda(0) & S_\lambda(0) \\ C'_\lambda(0) & S'_\lambda(0) \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}.$$

Note that if  $U \equiv 0$ , we may take  $C_\lambda(x) = \cos \sqrt{\lambda}x$  and  $S_\lambda(x) = \frac{\sin \sqrt{\lambda}x}{\sqrt{\lambda}}$ .

Denote  $c(\lambda) = C_\lambda(L)$ ,  $s(\lambda) = S_\lambda(L)$ ,  $c'(\lambda) = C'_\lambda(L)$  and  $s'(\lambda) = S'_\lambda(L)$ . Any solution  $\psi_e$  of (2.2) satisfying  $\psi_e(0, \lambda) = c_1$  and  $\psi'_e(0, \lambda) = c_2$  will have at  $L$  the values

$$(2.3) \quad \begin{pmatrix} \psi_e(L, \lambda) \\ \psi'_e(L, \lambda) \end{pmatrix} = M_\lambda \begin{pmatrix} c_1 \\ c_2 \end{pmatrix} \quad \text{where } M_\lambda = \begin{pmatrix} c(\lambda) & s(\lambda) \\ c'(\lambda) & s'(\lambda) \end{pmatrix}.$$

By the Wronskian identity, we have

$$(2.4) \quad c(\lambda)s'(\lambda) - s(\lambda)c'(\lambda) = 1.$$

Since  $U$  is symmetric, we have the following identity (cf. [16, Lemma 3.1]):

$$(2.5) \quad c(\lambda) = s'(\lambda)$$

Indeed,  $x \mapsto C_\lambda(L-x)$  and  $x \mapsto s'(\lambda)C_\lambda(x) - c'(\lambda)S_\lambda(x)$  are both solutions of (2.2), with the same value and derivative at  $x = L$ , so that they must be equal. Evaluating these functions at  $x = 0$  gives (2.5). Similarly, we have

$$(2.6) \quad s(\lambda)C_\lambda(x) - c(\lambda)S_\lambda(x) = S_\lambda(L-x).$$

**2.2. Spectral theory on the infinite tree.** Let  $H_{\mathbf{T}_q}$  be the operator (2.1) on the equilateral  $(q+1)$ -regular quantum tree  $\mathbf{G} = \mathbf{T}_q$ , with the Kirchhoff boundary conditions.

It is known [16] that

$$(2.7) \quad \sigma(H_{\mathbf{T}_q}) = \sigma_1 \cup \sigma_2,$$

where

$$(2.8) \quad \begin{aligned} \sigma_1 &= \{\lambda \in \mathbb{R} : |(q+1)c(\lambda) + \alpha s(\lambda)| \leq 2\sqrt{q}\}, \\ \sigma_2 &= \{\lambda \in \mathbb{R} : s(\lambda) = 0\}. \end{aligned}$$

Moreover, assuming  $q \geq 2$ ,

$$\sigma_{\text{ac}}(H_{\mathbf{T}_q}) = \sigma_1, \quad \sigma_{\text{pp}}(H_{\mathbf{T}_q}) = \sigma_2 \quad \text{and} \quad \sigma_1 \cap \sigma_2 = \emptyset.$$

In general,  $\sigma_2$  is a discrete set,  $\sigma_2 \subset \sigma(D)$ , where  $D$  is the operator  $D\psi = -\psi'' + U\psi$  with Dirichlet boundary conditions on  $[0, L]$  and points of  $\sigma_2$  are eigenvalues of  $H_{\mathbf{T}_q}$  with infinite multiplicity. Moreover,  $\sigma(H_{\mathbf{T}_q})$  has a band structure with infinitely many gaps [29]. In the special case  $\alpha = 0$  and  $U \equiv 0$ , we denote  $-\Delta_{\mathbf{T}_q} = H_{\mathbf{T}_q}$ , and have

$$\sigma_{\text{ac}}(-\Delta_{\mathbf{T}_q}) = \bigcup_{n=1}^{\infty} \left[ \left( \frac{(n-1)\pi + \theta}{L} \right)^2, \left( \frac{n\pi - \theta}{L} \right)^2 \right] \quad \text{and} \quad \sigma_{\text{pp}}(-\Delta_{\mathbf{T}_q}) = \left\{ \left( \frac{n\pi}{L} \right)^2 \right\}_{n=1}^{\infty},$$

where  $\theta = \arccos \frac{2\sqrt{q}}{q+1} \in (0, \frac{\pi}{2})$ .

We say that  $I \subset \sigma_{\text{ac}}(H_{\mathbf{T}_q})$  is *in a fixed band* if the map  $I \ni \lambda \mapsto (q+1)c(\lambda) + \alpha s(\lambda)$  is injective.

**2.3. Presentation of our results.** Consider a sequence of  $(q+1)$ -regular combinatorial graphs  $G_N = (V_N, E_N)$ , with  $|V_N| = N$ , satisfying the following assumptions:

**(EXP)** The sequence  $(G_N)$  is a family of expanders. That is, there exists  $\beta > 0$  such that the spectrum  $\sigma(\frac{\mathcal{A}_{G_N}}{q+1}) \subset [-1 + \beta, 1 - \beta] \cup \{1\}$  for all  $N$ , where  $\mathcal{A}_{G_N}$  is the adjacency matrix of the graph  $G_N$ .

**(BST)** For all  $r > 0$ ,

$$\lim_{N \rightarrow \infty} \frac{\#\{x \in V_N : \rho_{G_N}(x) < r\}}{N} = 0$$

where  $\rho_{G_N}(x)$  is the largest  $\rho$  such that the sub-graph contained in a ball of radius  $\rho$  centered at  $x$  has no closed cycles.

These assumptions were also needed to prove quantum ergodicity for discrete graphs in [2, 4]. They are known to be “generic”, in the sense that a regular graph picked at random will typically satisfy these assumptions. There are also explicit examples of  $(G_N)$  satisfying both assumptions [27, 13]. Assumption **(BST)** means that the graphs  $(G_N)$  have few short cycles, equivalently, they converge to the combinatorial  $(q+1)$ -regular tree  $\mathbb{T}_q$  in the sense of Benjamini-Schramm [8].

We may now present our main result, which essentially says that in the regions where  $H_{\mathbf{T}_q}$  is spectrally delocalized, it also has a form of spatial delocalization. Namely, if we consider a sequence of quantum graphs converging to  $\mathbf{T}_q$ , then the eigenfunctions in any bounded interval within the AC spectrum of  $H_{\mathbf{T}_q}$  are quantum ergodic when the graph is large enough. To clarify this result, we present it in the form of three theorems of growing generality. Theorem 2.1 discusses the special case where the observables are constant on the edges. General bounded observables are considered in Theorem 2.2. Finally, Theorem 2.3 studies the case where the observables are replaced by integral kernels.

**Theorem 2.1.** *Assume  $(G_N)$  satisfies **(EXP)** and **(BST)**.*

*Let  $(\psi_n^{(N)})_{n \in \mathbb{N}}$  be an orthonormal basis of eigenfunctions of  $H_{\mathbf{G}_N}$  and let  $(\lambda_n^{(N)})$  be the corresponding sequence of eigenvalues. Let  $I$  be an open interval such that  $\bar{I} \subset \sigma_{\text{ac}}(H_{\mathbf{T}_q})$  lies in a fixed band of AC spectrum, and denote by  $N(I)$  the count of  $\lambda_n^{(N)} \in I$ .*

*Let  $a : E_N \rightarrow [-1, 1]$ . Then*

$$(2.9) \quad \lim_{N \rightarrow \infty} \frac{1}{N(I)} \sum_{\lambda_n^{(N)} \in I} \left| \langle \psi_n^{(N)}, a \psi_n^{(N)} \rangle - \langle a \rangle \right|^2 = 0,$$

where  $\langle \psi_n^{(N)}, a \psi_n^{(N)} \rangle = \sum_{e \in E_N} a(e) \int_0^L |\psi_n^{(N)}(x_e)|^2 dx_e$  and

$$\langle a \rangle = \frac{1}{|E_N|} \sum_{e \in E_N} a(e).$$

We postpone discussing the consequences of this theorem until § 2.4. In order to describe the limiting distribution of the  $|\psi_n^{(N)}(t)|^2 dt$  inside the edges, we now replace the locally constant observables  $a(x_e) = a(e)$  by general observables  $f(x_e)$  as follows.

**Theorem 2.2.** *Assume  $(G_N)$  satisfy **(EXP)** and **(BST)**.*

*Let  $(\psi_n^{(N)})_{n \in \mathbb{N}}$  be an orthonormal basis of eigenfunctions of  $H_{\mathbf{G}_N}$  and let  $I$  be an open interval such that  $\bar{I} \subset \sigma_{\text{ac}}(H_{\mathbf{T}_q})$  lies in a fixed band of AC spectrum. Then for any function  $f = (f_e) \in \bigoplus_{e \in E_N} L^2[0, L]$  satisfying  $|f(x)| \leq 1$ , we have*

$$(2.10) \quad \lim_{N \rightarrow \infty} \frac{1}{N(I)} \sum_{\lambda_n^{(N)} \in I} \left| \langle \psi_n^{(N)}, f \psi_n^{(N)} \rangle - \langle f \rangle_{\lambda_n^{(N)}} \right|^2 = 0,$$

where  $\langle \psi_n, f \psi_n \rangle = \sum_{e \in E_N} \int_0^L f_e(x_e) |\psi_n(x_e)|^2 dx_e$ ,

$$\langle f \rangle_{\lambda_n} = \frac{1}{N} \sum_{e \in E_N} \int_0^L f_e(x_e) \Psi_{\lambda_n}(x_e) dx_e,$$

$$\Psi_{\lambda}(x_e) = \frac{1}{\kappa_{\lambda}} \cdot \frac{\operatorname{Im} G_{\mathbf{T}_q}^{\lambda+i0}(\tilde{x}_e, \tilde{x}_e)}{\operatorname{Im} G_{\mathbf{T}_q}^{\lambda+i0}(o, o)},$$

and if  $w(\lambda) := (q+1)c(\lambda) + \alpha s(\lambda)$ , we have

$$(2.11) \quad \kappa_{\lambda} = \frac{q+1}{s^2(\lambda)} \int_0^L S_{\lambda}^2(t) dt + \frac{w(\lambda)}{s^2(\lambda)} \int_0^L S_{\lambda}(L-t) S_{\lambda}(t) dt.$$

Here,  $\tilde{x}_e$  is a lift of  $x_e$  to  $\mathbf{T}_q$  (the universal cover of  $\mathbf{G}_N$ ).

In the previous theorem,  $G_{\mathbf{T}_q}^{\gamma}(x, y) = (H_{\mathbf{T}_q} - \gamma)^{-1}(x, y)$  is the Green's function of  $H_{\mathbf{T}_q}$ ,  $o$  is an arbitrary vertex, and we denoted  $G^{\lambda+i0}(x, y) = \lim_{\eta \downarrow 0} G^{\lambda+i\eta}(x, y)$ . Note that this Green's function is completely explicit, see (3.9) and Appendix A.

Theorem 2.1 follows directly from Theorem 2.2 and identity (A.4), which implies that  $\int_0^L \Psi_{\lambda_n}(t) dt = \frac{2}{q+1}$ , and thus  $\frac{\int_0^L \Psi_{\lambda_n}(t) dt}{N} = \frac{1}{|E_N|}$ .

*Quantum ergodicity for integral operators.* We now present a more general quantum ergodicity result whose aim is to study the eigenfunction correlator  $\psi_n^{(N)}(x) \psi_n^{(N)}(y) dx dy$ . This also allows us to gain a better understanding of the limiting density  $\Psi_{\lambda}$ .

In general, an integral operator on  $\bigoplus_e L^2[0, L_e]$  takes the form

$$(K\psi)(x_e) = \sum_{e' \in E_N} \int_0^{L_{e'}} K_{e,e'}(x_e, y_{e'}) \psi_{e'}(y_{e'}) dy_{e'}.$$

In the following we consider integral operators whose kernels vanish if  $d(e, e')$  is large. For convenience, we consider directed edges and consider operators of the form

$$(2.12) \quad (K_k \psi)(x_b) = \sum_{(b_1, \dots, b_k)_b} \int_0^L K_{b_1, b_k}(x_{b_1}, y_{b_k}) \psi_{b_k}(y_{b_k}) dy_{b_k},$$

where  $b_1 := b$  and the sum runs over non-backtracking paths of length exactly  $k$  whose first vertices are  $(o_b, t_b)$ . Given such  $K_k$ , we have

$$(2.13) \quad 2 \langle \psi_n^{(N)}, K_k \psi_n^{(N)} \rangle = \sum_{(b_1; b_k) \in B_k} \int_0^L \int_0^L K_{b_1, b_k}(x_{b_1}, y_{b_k}) \overline{\psi_n^{(N)}(x_{b_1})} \psi_n^{(N)}(y_{b_k}) dx_{b_1} dy_{b_k},$$

where the sum now runs over all  $k$ -paths  $(b_1; b_k) \equiv (x_0; x_k)$  in  $G_N$ . The factor 2 in the left-hand side of (2.13) arises because we sum over each bond twice. We shall prove that

**Theorem 2.3.** *Under the assumptions of Theorem 2.2, if the integral kernels satisfy  $|K_{b_1, b_k}(x_{b_1}, y_{b_k})| \leq 1$ , then*

$$\lim_{N \rightarrow \infty} \frac{1}{N(I)} \sum_{\lambda_n^{(N)} \in I} \left| \langle \psi_n^{(N)}, K_k \psi_n^{(N)} \rangle - \langle K_k \rangle_{\lambda_n^{(N)}} \right|^2 = 0,$$

where

$$\langle K_k \rangle_{\lambda_n^{(N)}} = \frac{1}{N} \sum_{(b_1; b_k) \in B_k} \int_0^L \int_0^L K_{b_1, b_k}(x_{b_1}, y_{b_k}) \Psi_{\lambda_n^{(N)}, k}(x_{b_1}, y_{b_k}) dx_{b_1} dy_{b_k},$$

and for  $x_{b_1} \in b_1, y_{b_k} \in b_k$ , we have

$$(2.14) \quad \Psi_{\lambda,k}(x_{b_1}, y_{b_k}) = \frac{1}{2\kappa_\lambda} \cdot \frac{\operatorname{Im} G_{\mathbf{T}_q}^{\lambda+i0}(\tilde{x}_{b_1}, \tilde{y}_{b_k})}{\operatorname{Im} G_{\mathbf{T}_q}^{\lambda+i0}(o, o)}.$$

Here,  $\tilde{x}_{b_1}$  and  $\tilde{y}_{b_k}$  are lifts of  $x_{b_1}$  and  $y_{b_k}$  to  $\mathbf{T}_q$  such that  $d(\tilde{x}_{b_1}, \tilde{y}_{b_k}) = d(x_{b_1}, y_{b_k})$ .

**2.4. Consequences.** Using Markov's inequality, it follows from Theorem 2.1 that for any  $\epsilon > 0$ , we have

$$\lim_{N \rightarrow \infty} \frac{1}{N(I)} \# \left\{ \lambda_n^{(N)} \in I : \left| \langle \psi_n^{(N)}, a \psi_n^{(N)} \rangle - \langle a \rangle \right| > \epsilon \right\} = 0.$$

In particular, we obtain that for fixed large  $N$ , we have

$$(2.15) \quad \sum_{e \in E_N} a(e) \int_0^L |\psi_n^{(N)}(x_e)|^2 dx_e \approx \frac{1}{|E_N|} \sum_{e \in E_N} a(e)$$

for most  $\lambda_n^{(N)} \in I$ . Formula (2.15) is justified if  $a$  is supported on a macroscopic part of  $\mathbf{G}_N$ , say  $|\operatorname{supp} a| = c \cdot |\mathbf{G}_N|$  for some  $0 < c \leq 1$ . In this case, the error term, essentially an inverse power of the girth of the graph, is sufficiently small for typical graphs. However, the control we have on the remainder in the limit (2.9) is not good enough to justify (2.15) when  $a$  has a small support (for instance, when  $a$  is supported on a single edge). Similar difficulties arise for combinatorial graphs [3, 6].

For general observables  $f$ , we similarly deduce from Theorem 2.2 that

$$(2.16) \quad \int_{\mathbf{G}_N} |\psi_n^{(N)}(x)|^2 f(x) dx \approx \frac{1}{N} \int_{\mathbf{G}_N} f(x) \Psi_{\lambda_n^{(N)}}(x) dx$$

for most  $\lambda_n^{(N)} \in I$  and observables with large support. One subtlety one should not forget is that, when we say “for most  $\lambda_n^{(N)}$ ”, the set of indices for which this holds true depends on the observable  $f$ .

The limit measure turns out to be the uniform measure in the special case  $U \equiv 0$  and  $\alpha = 0$ , as we report in Appendix A.2. So in that special case we get for *any* observable  $f^{(N)} = f$  with large support that  $\int_{\mathbf{G}_N} |\psi_n^{(N)}(x)|^2 f(x) dx \approx \frac{1}{|\mathbf{G}_N|} \int_{\mathbf{G}_N} f(x) dx$  for most  $\lambda_n^{(N)} \in I$ . In other words, most eigenfunctions equidistribute in the large graph limit (in a weak sense). Such an equidistribution holds because in this special case, for any *fixed*  $\lambda$ , the density of the spectral measure of the infinite tree  $\frac{1}{\pi} \operatorname{Im} G^{\lambda+i0}(x, x)$  is constant as  $x$  moves along the tree, so that  $\Psi_{\lambda_n^{(N)}}(x)$  is a constant. In general, (2.16) says that  $|\psi_n^{(N)}(x)|^2$  varies along the tree essentially like this density does (at the fixed  $\lambda_n^{(N)}$ ). Note that by (A.4),  $x \mapsto \Psi_\lambda(x)$  is analytic (cf. [30, p. 10]) and symmetric on each edge :  $\Psi_\lambda(L - x) = \Psi_\lambda(x)$ .

Curiously, the coefficient  $\frac{1}{\kappa_{\lambda_n}}$  is the  $\ell^2$ -norm of the function  $\psi_n$  restricted to the vertices (see (4.7)).

Theorem 2.3 can similarly be interpreted as saying that when  $N$  gets large enough, the correlation  $\int \psi_n^{(N)}(x_{b_1}) \psi_n^{(N)}(y_{b_k}) dx_{b_1} dy_{b_k}$  approaches  $\frac{1}{N} \cdot \frac{1}{\kappa_\lambda} \cdot \frac{\operatorname{Im} G_{\mathbf{T}_q}^{\lambda_n}(\tilde{x}_{b_1}, \tilde{y}_{b_k})}{\operatorname{Im} G_{\mathbf{T}_q}^{\lambda_n}(o, o)} dx_{b_1} dy_{b_k}$ .

Note that Theorem 2.3 formally generalizes Theorem 2.2. In particular, the  $|\psi_n^{(N)}(t)|^2 dt$  equidistributes if and only if the diagonal  $\Psi_{\lambda_n,1}(t, t) = \text{constant}$ .

**Remark 2.4.** We can also test for quantum ergodicity in several bands of AC spectra simultaneously. In fact, if  $I$  is a finite union of open intervals  $I_r$ , each lying in a band, then the variance (2.10) of each  $I_r$  vanishes, so the same holds for the sum on  $I$ .

If we replace  $\frac{1}{N(I)}$  by  $\frac{1}{N}$  in (2.10), we can deduce moreover that the variance (2.10) also vanishes for  $I$  between the AC bands, as long as  $I \cap \sigma_2 = \emptyset$ . In this case, no real proof



is needed; this simply follows because the number of eigenvalues in such regions becomes negligible as the graph grows large. See Remark 4.4 for details.

Finally, one does not expect delocalization for  $\lambda \in \sigma_2$ . There,  $s(\lambda) = 0$ , and solutions satisfy Dirichlet boundary conditions. This means that there exist eigenfunctions supported on individual loops of the graph, in a sine-like way, and vanish identically elsewhere on the graph. See [25, Section 5] as well as [18].

**Remark 2.5.** One may wonder if we can hope for a stronger “quantum unique ergodicity” result, namely that  $\int_{\mathbf{G}_N} |\psi_n^{(N)}(x)|^2 f(x) dx \approx \frac{1}{|\mathbf{G}_N|} \int_{\mathbf{G}_N} f(x) dx$  for all  $\lambda_n^{(N)} \in I$  and all observables  $f$ .

To illustrate the subtlety of this question, let us discuss the very easy case where  $q = 1$  and  $U = \alpha = 0$ . In other words,  $G_N$  is just an  $N$ -cycle. Under Kirchhoff conditions, the Laplacian on  $\mathbf{G}_N$  is the same as the Laplacian on the interval  $\mathbf{I}_N = [0, NL]$  with periodic boundary conditions. Note that this case does not enter the framework of our theorem, since these graphs are not expanders.

An orthonormal basis of eigenfunctions of  $-\Delta_{\mathbf{I}_N}$  is given by  $e_k(x) = \frac{\exp(\frac{2\pi i k x}{NL})}{\sqrt{NL}}$ ,  $k \in \mathbb{Z}$ . For this basis, quantum ergodicity, and even quantum unique ergodicity, are trivial; in fact  $\int |e_k(x)|^2 f(x) dx = \frac{1}{NL} \int f(x) dx$ .

A more interesting basis is given by  $(v_j)_{j \in \mathbb{N}}$ , where

$$v_0(x) = \frac{1}{\sqrt{NL}}, \quad v_{2j}(x) = \sqrt{\frac{2}{NL}} \sin\left(\frac{2j\pi x}{NL}\right), \quad \text{and} \quad v_{2j+1}(x) = \sqrt{\frac{2}{NL}} \cos\left(\frac{2j\pi x}{NL}\right).$$

In this case we have,  $\int |v_{2j}(x)|^2 f(x) dx = \frac{1}{NL} \int f(x) dx - \frac{1}{NL} \int f(x) \cos(\frac{4j\pi x}{NL}) dx$  as well as  $\int |v_{2j+1}(x)|^2 f(x) dx = \frac{1}{NL} \int f(x) dx + \frac{1}{NL} \int f(x) \cos(\frac{4j\pi x}{NL}) dx$ . Hence,

$$\frac{1}{N(I)} \sum_{\lambda_j \in I} |\langle v_j, f v_j \rangle - \langle f \rangle|^2 = \frac{1}{N(I)} \sum_{\lambda_j \in I} \left| \frac{1}{NL} \int f(x) \cos\left(\frac{4j\pi x}{NL}\right) dx \right|^2.$$

But  $\sum_{j \in \mathbb{N}} \left| \sqrt{\frac{2}{NL}} \int \cos(\frac{4j\pi x}{NL}) f(x) dx \right|^2 = \sum_{j \in \mathbb{N}} |\langle v_{4j+1}, f \rangle|^2 \leq \|f\|_{L^2(\mathbf{I}_N)}^2 \leq NL$ , assuming  $|f(x)| \leq 1$ . Hence,

$$\frac{1}{N(I)} \sum_{\lambda_j \in I} |\langle v_j, f v_j \rangle - \langle f \rangle|^2 \leq \frac{1}{N(I)} \cdot \frac{1}{NL} \cdot NL = \frac{1}{N(I)} \rightarrow 0$$

as  $N \rightarrow \infty$ , since  $N(I) \rightarrow \infty$ . In fact,  $\lambda_j \in [a, b]$  iff  $(\frac{2\pi j}{NL})^2 \in [a, b]$ , i.e.  $j \in [\frac{NL}{2\pi} \sqrt{a}, \frac{NL}{2\pi} \sqrt{b}]$ , so  $N(I) = c_I \cdot N$ . This proves (2.10) for this basis.

However, note that for  $\lambda_j \in I$ , so that  $j = \frac{NL}{2\pi} c_j$  for some  $c_j \in [\sqrt{a}, \sqrt{b}]$ , we have  $\int |v_{2j}(x)|^2 f(x) dx = \frac{2}{NL} \int f(x) \sin^2(c_j x) dx$ . We may choose observables such that this is not close to the uniform average. In fact, taking  $f(x) = \cos(2c_j x)$  yields the error  $\frac{1}{NL} \int f(x) \cos(2c_j x) dx = \frac{1}{2}$ . Therefore, in this case, we do not have  $\int_{\mathbf{G}_N} |\psi_n^{(N)}(x)|^2 f(x) dx \approx \frac{1}{|\mathbf{G}_N|} \int_{\mathbf{G}_N} f(x) dx$  for all  $\lambda_n^{(N)} \in I$  and all observables  $f$ .

The moral of this discussion is that, if one wants to show an equidistribution result holding for all eigenfunctions, one should probably restrict a lot the space of observables. In particular, the observables should not oscillate at a speed much smaller than the size of the graph  $\mathbf{G}_N$ , like they do in our theorems.

### 3. RELATIONSHIP BETWEEN THE CONTINUOUS AND DISCRETE PROBLEMS

**3.1. Non-backtracking eigenfunctions.** Given  $\gamma \in \mathbb{C}$ , we set

$$(3.1) \quad w(\gamma) := (q+1)c(\gamma) + \alpha s(\gamma).$$

As in [16], we consider

$$(3.2) \quad \mu^\pm(\gamma) := \frac{w(\gamma) \pm \sqrt{w(\gamma)^2 - 4q}}{2q}.$$

This function arises in the study of transfer matrices on the quantum tree. Note that for  $\lambda \in \sigma_{\text{ac}}(H_{\mathbf{T}_q})$ , the  $\mu^\pm(\lambda)$  are complex conjugate (see (2.8)), with  $|\mu^\pm(\lambda)|^2 = \frac{1}{q}$ . In general,

$$(3.3) \quad \mu^+(\gamma)\mu^-(\gamma) = \frac{1}{q}, \quad q(\mu^+(\gamma) + \mu^-(\gamma)) = w(\gamma).$$

If  $(\lambda_j^{(N)})_{j \in \mathbb{N}} \subset \mathbb{R}$  is the discrete spectrum of  $H_{\mathbf{G}_N}$ , we shall denote  $\mu_j^\pm := \mu^\pm(\lambda_j^{(N)})$ . We drop the superscripts  $(N)$  from notations where it will not lead to confusion.

Fix  $G = G_N$ . Given an orthonormal basis  $(\psi_j)$  of eigenfunctions of  $H_{\mathbf{G}}$  with eigenvalues  $\lambda_j$ , we define the discrete function  $\mathring{\psi}_j$  on  $V_N$  by

$$(3.4) \quad \mathring{\psi}_j(v) = \psi_j(v).$$

This is well-defined as  $\psi_j$  satisfies Kirchhoff conditions.

Let  $B$  be the set of directed edges of  $G_N$ , so  $|B| = 2|E_N| = N(q+1)$ . Each edge  $b \in B$  has an origin  $o_b \in V_N$  and a terminus  $t_b \in V_N$ . We denote by  $\hat{b}$  the reversal of  $b$ .

Consider the non-backtracking operator  $\mathcal{B}$  defined on  $\ell^2(B)$  by

$$(\mathcal{B}f)(b) = \sum_{b^+ \in \mathcal{N}_b^+} f(b^+),$$

where  $\mathcal{N}_b^+$  is the set of edges  $b'$  with  $o_{b'} = t_b$  and  $b' \neq \hat{b}$ , i.e. the set of outgoing edges from  $b$ . Define  $\tau_\pm : \mathbb{C}^{V_N} \rightarrow \mathbb{C}^B$  by

$$(\tau_+\psi)(b) = \psi(t_b) \quad \text{and} \quad (\tau_-\psi)(b) = \psi(o_b).$$

Inspired by [2], we define

$$(3.5) \quad f_j = \tau_+\mathring{\psi}_j - \mu_j^-\tau_-\mathring{\psi}_j \quad \text{and} \quad f_j^* = \tau_-\mathring{\psi}_j - \mu_j^-\tau_+\mathring{\psi}_j.$$

**Lemma 3.1.** *The following properties hold:*

$$(3.6) \quad \mathcal{A}_G \mathring{\psi}_j = w(\lambda_j) \mathring{\psi}_j,$$

$$(3.7) \quad \mu_j^-\mathcal{B}f_j = f_j \quad \text{and} \quad \mu_j^-\mathcal{B}^*f_j^* = f_j^*.$$

*Proof.* Let  $v \in V_N$  and  $b \in B$  be such that  $v = t_b$ . We write

$$(\mathcal{A}_G \mathring{\psi}_j)(t_b) = \psi_j(o_b) + \sum_{b^+ \in \mathcal{N}_b^+} \psi_j(t_{b^+}).$$

Recall that  $\psi_j(t_{b^+}) = \psi_j(o_{b^+})c(\lambda_j) + \psi_j'(o_{b^+})s(\lambda_j)$ . By the Kirchhoff conditions,  $\psi_j(o_{b^+}) = \psi_j(t_b)$  and  $\sum_{b^+ \in \mathcal{N}_b^+} \psi_j'(o_{b^+}) = \psi_j'(t_b) + \alpha\psi_j(t_b)$ . We thus get

$$(\mathcal{A}_G \mathring{\psi}_j)(t_b) = \psi_j(o_b) + [qc(\lambda_j) + \alpha s(\lambda_j)]\psi_j(t_b) + s(\lambda_j)\psi_j'(t_b).$$

By (2.3),  $\begin{pmatrix} \psi_j(o_b) \\ \psi_j'(o_b) \end{pmatrix} = M_{\lambda_j}^{-1} \begin{pmatrix} \psi_j(t_b) \\ \psi_j'(t_b) \end{pmatrix}$ . It follows that  $\psi_j(o_b) = c(\lambda_j)\psi_j(t_b) - s(\lambda_j)\psi_j'(t_b)$ .

Recalling (3.1), this proves (3.6). Next,

$$\begin{aligned} (\mathcal{B}f_j)(b) &= \sum_{b^+ \in \mathcal{N}_b^+} [\psi_j(t_{b^+}) - \mu_j^-\psi_j(o_{b^+})] \\ &= (\mathcal{A}\mathring{\psi}_j)(t_b) - \psi_j(o_b) - q\mu_j^-\psi_j(t_b) \\ &= [w(\lambda_j) - q\mu_j^-]\psi_j(t_b) - \psi_j(o_b). \end{aligned}$$

Using (3.3), we get  $(\mathcal{B}f_j)(b) = q\mu_j^+ \psi_j(t_b) - \psi_j(o_b) = \frac{1}{\mu_j^-} f_j(b)$ .

Finally, to prove that  $\mathcal{B}^* f_j^* = \frac{1}{\mu_j^-} f_j^*$ , note that if  $\iota$  is the edge reversal, i.e.  $(\iota f)(x_0, x_1) = f(x_1, x_0)$ , then  $\mathcal{B}^* = \iota \mathcal{B} \iota$  and  $f_j^* = \iota f_j$ , so we deduce that  $\mathcal{B}^* f_j^* = \frac{1}{\mu_j^-} f_j^*$ .  $\square$

If  $k \in \mathbb{N}$ , we denote by  $\mathcal{H}_k = \mathbb{C}^{B_k}$  the set of complex-valued functions on the set  $B_k$  of (discrete) non-backtracking paths of length  $k$  in  $G_N$ . Note that  $B_0 = V_N$  and  $B_1 = B$ .

In Section 4 below we are going to express the quantum variance in (2.10) using the discrete graph function  $\psi_j$  which is an eigenvector of  $\mathcal{A}_{G_N}$ . In this reduction, energy-dependent observables  $K^\lambda$  arise. Experience has shown us that the *non-backtracking quantum variance* is more adapted to handle such observables. This is defined using the eigenvectors  $f_j, f_j^*$  of  $\mathcal{B}$  and  $\mathcal{B}^*$  as follows :

$$(3.8) \quad \text{Var}_{\text{nb}}^I(K^\lambda) = \frac{1}{N(I)} \sum_{\lambda_j \in I} \left| \langle f_j^*, K_B^{\lambda_j} f_j \rangle \right|^2$$

for  $K^{\lambda_j} \in \mathcal{H}_k$ , where we denoted  $\langle f, K_B g \rangle = \sum_{(x_0; x_k) \in B_k} \overline{f(x_0, x_1)} K(x_0; x_k) g(x_{k-1}, x_k)$ .

**3.2. Relationship between the continuous and discrete resolvents.** Let  $G_{\mathbf{T}_q}^\gamma(x, y)$  be the integral kernel of the resolvent of  $H_{\mathbf{T}_q}$  at energy  $\gamma \in \mathbb{C}$ , and let  $G_{\mathcal{A}_{\mathbf{T}_q}}^\gamma(v, w)$  be the Green's function  $(\mathcal{A}_{\mathbf{T}_q} - \gamma)^{-1}(v, w)$ . This is well-defined, at least when  $\gamma \in \mathbb{C}^+ = \{\text{Im } z > 0\}$ , so that  $\gamma$  is away from the spectrum. On the real line, we denote

$$G^\lambda(x, y) := G^{\lambda+i0}(x, y) := \lim_{\eta \downarrow 0} G^{\lambda+i\eta}(x, y), \quad \lambda \in \mathbb{R}$$

when this limit exists. We use this notation for Green's functions of different objects, indicated with subscripts.

Explicit expressions for the Green's function on the quantum tree were worked-out by Carlson in [16]. One fixes  $b_0 \in \mathbf{T}_q$ ,  $b_0 \equiv [0, L]$ , and define for  $x \in b_0$ ,

$$U_\gamma(x) = -s(\gamma)C_\gamma(x) + [c(\gamma) - q\mu^+(\gamma)] S_\gamma(x),$$

$$V_\gamma(x) = s(\gamma)C_\gamma(L-x) - [c(\gamma) - q\mu^+(\gamma)] S_\gamma(L-x).$$

Given  $x_0 \in b_0$ , using the notations of [1], we let  $\mathbf{T}_{x_0}^+$  and  $\mathbf{T}_{x_0}^-$  be the subtrees produced by cutting  $\mathbf{T}$  at  $x_0$ . Any  $b^+ \in \mathbf{T}_{x_0}^+$  has the form  $\mathcal{B}^n b_0$ , which is an abusive notation to say that  $b^+$  can be reached by a non-backtracking path of length  $n$ . Similarly, any  $b^- \in \mathbf{T}_{x_0}^-$  takes the form  $b^- = \mathcal{B}^{*m} b_0$  for some  $m$ . We then define the functions

$$\begin{cases} U_{\gamma; x_0}^-(y) = (\mu^-(\gamma))^m U_\gamma(\tilde{y}) & \text{if } y \in \mathcal{B}^{*m} b_0, \\ V_{\gamma; x_0}^+(y) = (\mu^-(\gamma))^n V_\gamma(\tilde{y}) & \text{if } y \in \mathcal{B}^n b_0, \end{cases}$$

where  $\tilde{y}$  is the point in  $b_0$  at the corresponding position to  $y$ . Then  $U_{\gamma; x_0}^-$  and  $V_{\gamma; x_0}^+$  are solutions to the Kirchhoff problem on  $\mathbf{T}_{x_0}^-$  and  $\mathbf{T}_{x_0}^+$ , respectively. Given  $x, y \in \mathbf{T}$ , choose  $o, v$  such that  $x, y \in \mathbf{T}_o^+ \cap \mathbf{T}_v^-$ . Define

$$(3.9) \quad G_{\mathbf{T}_q}^\gamma(x, y) = \begin{cases} \frac{U_{\gamma; v}^-(x) V_{\gamma; o}^+(y)}{W_{v, o}^\gamma(x)} & \text{if } y \in \mathbf{T}_x^+, \\ \frac{U_{\gamma; v}^+(y) V_{\gamma; o}^-(x)}{W_{v, o}^\gamma(x)} & \text{if } y \in \mathbf{T}_x^-, \end{cases}$$

where  $W_{v, o}^\gamma(x)$  is the Wronskian

$$W_{v, o}^\gamma(x) = V_{\gamma; o}^+(x) (U_{\gamma; v}^-)'(x) - (V_{\gamma; o}^+)'(x) U_{\gamma; v}^-(x).$$

Then it is shown in [16] that  $G_{\mathbf{T}_q}^\gamma(x, y)$  is precisely the Green kernel of  $H_{\mathbf{T}_q}$ .

In Appendix A.1, we prove the following lemma :

**Lemma 3.2.** *Let  $\gamma \in \mathbb{C}^+$  and  $v, w$  be vertices in  $\mathbb{T}_q$ . Then we have*

$$G_{\mathbb{T}_q}^\gamma(v, w) = -s(\gamma)G_{\mathcal{A}_{\mathbb{T}_q}}^{w(\gamma)}(v, w).$$

*If  $\lambda \in \sigma_{ac}(H_{\mathbb{T}_q})$ , we also have  $G_{\mathbb{T}_q}^{\lambda+i0}(v, w) = -s(\lambda)G_{\mathcal{A}_{\mathbb{T}_q}}^{w(\lambda)+i0}(v, w)$ .*

**3.3. Spherical functions.** The spherical function of parameter  $\lambda \in \mathbb{R}$  is given, for  $d \in \mathbb{N}_0$ , by

$$(3.10) \quad \Phi_\lambda(d) = q^{-d/2} \left( \frac{2}{q+1} P_d \left( \frac{\lambda}{2\sqrt{q}} \right) + \frac{q-1}{q+1} Q_d \left( \frac{\lambda}{2\sqrt{q}} \right) \right),$$

where  $P_r(\cos \theta) = \cos r\theta$  and  $Q_r(\cos \theta) = \frac{\sin(r+1)\theta}{\sin \theta}$  are the Chebyshev polynomials of the first and second kinds, respectively.

As is well-known, for  $\lambda \in (-2\sqrt{q}, 2\sqrt{q})$ , we have

$$\Phi_\lambda(d(v, w)) = \frac{\operatorname{Im} G_{\mathcal{A}_{\mathbb{T}_q}}^\lambda(v, w)}{\operatorname{Im} G_{\mathcal{A}_{\mathbb{T}_q}}^\lambda(o, o)},$$

where  $o \in \mathbb{T}_q$  is any vertex. This follows from the fact that  $\Phi_{\lambda, v}(w) := \Phi_\lambda(d(v, w))$  is the unique eigenfunction of  $\mathcal{A}_{\mathbb{T}_q}$  which is radial from  $v$  and normalized as  $\Phi_{\lambda, v}(v) = 1$ . So one just checks that the function on the right-hand side satisfies these properties.

Using Lemma 3.2, we deduce that for  $\lambda \in \sigma_{ac}(H_{\mathbb{T}_q})$ ,

$$(3.11) \quad \Phi_{w(\lambda)}(k) = \frac{\operatorname{Im} G_{\mathbb{T}_q}^\lambda(v, w)}{\operatorname{Im} G_{\mathbb{T}_q}^\lambda(o, o)}$$

for vertices  $v, w$  separated by a distance  $k$  on the tree.

The spherical function is also the unique solution of the recursive relation

$$(3.12) \quad \Phi_\lambda(d) = \frac{1}{q} (\lambda \Phi_\lambda(d-1) - \Phi_\lambda(d-2))$$

with initial conditions

$$(3.13) \quad \Phi_\lambda(1) = \frac{\lambda}{q+1} \quad \text{and} \quad \Phi_\lambda(0) = 1.$$

#### 4. PROOF OF THE MAIN THEOREM

**4.1. Reduction to a discrete quantum variance.** Recall that the functions  $\psi_n^\circ$  were defined in (3.4). The following lemma tells us that  $\|\psi_n\|_{L^2(\mathbf{G}_N)}$  can be expressed in terms of  $\|\psi_n^\circ\|_{\ell^2(V_N)}$ .

**Lemma 4.1.** *Let  $(\psi_n)$  be an orthonormal basis of eigenfunctions with corresponding eigenvalues  $(\lambda_n)$ . For any  $n$  such that  $\lambda_n \notin \sigma_2$ , we have*

$$(4.1) \quad 2\|\psi_n\|^2 = \frac{2(q+1)}{s^2(\lambda_n)} \left( \int_0^L S_{\lambda_n}^2(x) dx \right) \|\psi_n^\circ\|^2 \\ + \frac{2}{s^2(\lambda_n)} \left( \int_0^L S_{\lambda_n}(L-x) S_{\lambda_n}(x) dx \right) w(\lambda_n) \|\psi_n^\circ\|^2.$$

*Proof.* We may write  $\psi_n(x_b) = \psi_n(o_b)C_{\lambda_n}(x_b) + \psi_n'(o_b)S_{\lambda_n}(x_b)$ , expressing  $\psi_n$  in the basis  $C_{\lambda_n}(x)$  and  $S_{\lambda_n}(x)$ . Since  $\psi_n(t_b) = \psi_n(o_b)c(\lambda_n) + \psi_n'(o_b)s(\lambda_n)$ , and since  $\lambda_n \notin \sigma_2$  implies  $s(\lambda_n) \neq 0$ , we get using (2.6),

$$(4.2) \quad \psi_n(x_b) = \frac{S_{\lambda_n}(L-x_b)}{s(\lambda_n)} \psi_n(o_b) + \frac{S_{\lambda_n}(x_b)}{s(\lambda_n)} \psi_n(t_b)$$

Thus,

$$\begin{aligned}
\int_0^L f_b(x) |\psi_n(x)|^2 dx &= \frac{|\psi_n(o_b)|^2}{s^2(\lambda_n)} \int_0^L f_b(x) S_{\lambda_n}^2(L-x) dx \\
&\quad + \frac{\overline{\psi_n(o_b)} \psi_n(t_b) + \psi_n(o_b) \overline{\psi_n(t_b)}}{s^2(\lambda_n)} \int_0^L f_b(x) S_{\lambda_n}(L-x) S_{\lambda_n}(x) dx \\
(4.3) \quad &\quad + \frac{|\psi_n(t_b)|^2}{s^2(\lambda_n)} \int_0^L f_b(x) S_{\lambda_n}^2(x) dx.
\end{aligned}$$

Consider the special case where  $f_b \equiv 1$  for all  $b$ . Since  $\|\psi_n\|^2 = \frac{1}{2} \sum_{b \in B} \int_0^L |\psi_n(x)|^2 dx$  (the half is due to directed edges), and since  $\int_0^L S_{\lambda_n}^2(L-x) dx = \int_0^L S_{\lambda_n}^2(x) dx$ , we get

$$\begin{aligned}
2\|\psi_n\|^2 &= \frac{1}{s^2(\lambda_n)} \left( \int_0^L S_{\lambda_n}^2(x) dx \right) \left( \sum_{b \in B} |\psi_n(o_b)|^2 + \sum_{b \in B} |\psi_n(t_b)|^2 \right) \\
&\quad + \frac{1}{s^2(\lambda_n)} \left( \int_0^L S_{\lambda_n}(L-x) S_{\lambda_n}(x) dx \right) \sum_{b \in B} (\overline{\psi_n(o_b)} \psi_n(t_b) + \psi_n(o_b) \overline{\psi_n(t_b)}).
\end{aligned}$$

Now  $\sum_b |\psi_n(o_b)|^2 = \sum_b |\psi_n(t_b)|^2 = (q+1)\|\dot{\psi}_n\|^2$ ,  $\sum_b \overline{\psi_n(o_b)} \psi_n(t_b) = \langle \dot{\psi}_n, \mathcal{A}_{G_N} \dot{\psi}_n \rangle$  and  $\sum_b \psi_n(t_b) \overline{\psi_n(o_b)} = \langle \mathcal{A}_{G_N} \dot{\psi}_n, \dot{\psi}_n \rangle$ .

The result then follows from (3.6), which tells us that  $\mathcal{A}_{G_N} \dot{\psi}_n = w(\lambda_n) \dot{\psi}_n$ .  $\square$

Recall that  $B_k$  is the set of (discrete) non-backtracking paths of length  $k$  in  $G = G_N$ .

We now define the following discrete quantum variance. Given  $K^\lambda \in \mathcal{H}_k$  a family of operators depending on  $\lambda$ , we set

$$(4.4) \quad \text{Var}^I(K^\lambda) = \frac{1}{N(I)} \sum_{\lambda_n \in I} \left| \langle \dot{\psi}_n, K_G^{\lambda_n} \dot{\psi}_n \rangle \right|^2,$$

where

$$(4.5) \quad \langle \varphi, K_G \psi \rangle = \sum_{(x_0; x_k) \in B_k} \overline{\varphi(x_0)} K(x_0; x_k) \psi(x_k).$$

We also introduce

$$\langle K \rangle = \frac{1}{N} \sum_{(x_0; x_k) \in B_k} K(x_0; x_k) \quad \text{and} \quad \mathcal{H}_k^o = \{K \in \mathcal{H}_k : \langle K \rangle = 0\}$$

and notice that if  $S_k \in \mathcal{H}_k$  are the constant functions

$$(4.6) \quad S_0(x_0) = 1 \quad \text{and} \quad S_k(x_0; x_k) = \frac{1}{(q+1)q^{k-1}}$$

for  $k \geq 1$ , then for any  $K \in \mathcal{H}_k$ , we have  $K - \langle K \rangle S_k \in \mathcal{H}_k^o$ .

The following lemma allows us to reduce the statement of Theorem 2.2 to a bound on the variances of three operators. Recall that  $\kappa_{\lambda_n}$  and  $\langle f \rangle_{\lambda_n}$  were defined in Theorem 2.2.

**Lemma 4.2.** *Let  $(\psi_n)$  be an orthonormal basis of eigenfunctions of  $H_{\mathbf{G}_N}$  with eigenvalues  $(\lambda_n)$ . Assume  $\lambda_n \in \bar{I} \subset \sigma_{\text{ac}}(H_{\mathbf{T}_q})$ . Then*

(1) *We have*

$$(4.7) \quad \|\dot{\psi}_n\|^2 = \frac{1}{\kappa_{\lambda_n}} \leq C_{L,I}$$

(2) *Given a function  $f = (f_e)$  on the quantum graph  $\mathbf{G}_N$ , define  $K_{f,n}, J_{f,n} \in \mathcal{H}_0$  and  $M_{f,n} \in \mathcal{H}_1$  by*

$$K_{f,n}(x_0) = \sum_{x_1 \sim x_0} \frac{1}{s^2(\lambda_n)} \int_0^L f_{(x_0, x_1)}(t) S_{\lambda_n}^2(L-t) dt,$$

$$J_{f,n}(x_1) = \sum_{x_0 \sim x_1} \frac{1}{s^2(\lambda_n)} \int_0^L f_{(x_0, x_1)}(t) S_{\lambda_n}^2(t) dt,$$

$$M_{f,n}(x_0, x_1) = \frac{2}{s^2(\lambda_n)} \int_0^L f_{(x_0, x_1)}(t) S_{\lambda_n}(L-t) S_{\lambda_n}(t) dt.$$

Then

$$(4.8) \quad 2 \langle \psi_n, f \psi_n \rangle = \langle \dot{\psi}_n, (K_{f,n} + J_{f,n} + M_{f,n})_G \dot{\psi}_n \rangle.$$

Moreover, we have

$$(4.9) \quad 2(\langle \psi_n, f \psi_n \rangle - \langle f \rangle_{\lambda_n}) = \langle \dot{\psi}_n, (K_{f,n} - \langle K_{f,n} \rangle S_0)_G \dot{\psi}_n \rangle + \langle \dot{\psi}_n, (J_{f,n} - \langle J_{f,n} \rangle S_0) \dot{\psi}_n \rangle \\ + \langle \dot{\psi}_n, (M_{f,n} - \langle M_{f,n} \rangle S_1)_G \dot{\psi}_n \rangle.$$

- (3) If  $f$  is locally constant, i.e.  $f_e(x) = c_e$  and  $|c_e| \leq 1$ , define  $K_f \in \mathcal{H}_0$  and  $M_f \in \mathcal{H}_1$  by  $K_f(x_0) = \sum_{x_1 \sim x_0} c_{(x_0, x_1)}$  and  $M_f(x_0, x_1) = c_{(x_0, x_1)}$ . Then we get for  $I \subset \sigma_{ac}(H_{\mathbf{T}_q})$  and  $\langle f \rangle$  as in Theorem 2.1,

$$\frac{1}{N(I)} \sum_{\lambda_n \in I} |\langle \psi_n, f \psi_n \rangle - \langle f \rangle|^2 \leq C'_{L,I} (\text{Var}^I(K_f - \langle K_f \rangle S_0) + \text{Var}^I(M_f - \langle M_f \rangle S_1)).$$

*Proof.* (1) The equality in (4.7) follows from (4.1), by the definition of  $\kappa_{\lambda_n}$ . To see the uniform bound, first note by Cauchy-Schwarz that  $|\int_0^L S_{\lambda_n}(L-x) S_{\lambda_n}(x) dx| \leq \int_0^L S_{\lambda_n}^2(x) dx$ . We therefore have

$$\kappa_{\lambda_n} \geq \frac{q+1-|w(\lambda_n)|}{s^2(\lambda_n)} \int_0^L S_{\lambda_n}^2(x) dx.$$

As  $\lambda_n \in \sigma_{ac}(H_{\mathbf{T}_q})$ , we have  $|w(\lambda_n)| \leq 2\sqrt{q}$ . Hence,  $\kappa_{\lambda_n} \geq \frac{q+1-2\sqrt{q}}{s^2(\lambda_n)} \int_0^L S_{\lambda_n}^2(x) dx$ . Finally, note that  $I \ni \lambda \mapsto \frac{1}{s^2(\lambda)} \int_0^L S_{\lambda}^2(x) dx$  is continuous, so it reaches its minimum on  $\bar{I} \subset \sigma_{ac}(H_{\mathbf{T}_q})$ , say at  $\lambda_*$ . This minimum cannot be zero, otherwise  $S_{\lambda_*}^2(x)$  would be identically zero. Hence,  $\frac{1}{s^2(\lambda_n)} \int_0^L S_{\lambda_n}^2(x) dx \geq c_{L,I} > 0$ , proving the claim.

(2) Equation (4.8) readily follows from (4.3). To deduce (4.9), note that

$$2(\langle \psi_n, f \psi_n \rangle - \langle f \rangle_{\lambda_n}) = \langle \dot{\psi}_n, (K_{f,n} + J_{f,n} + M_{f,n})_G \dot{\psi}_n \rangle - 2 \langle f \rangle_{\lambda_n} \kappa_{\lambda_n} \cdot \|\dot{\psi}_n\|^2.$$

On the other hand,

$$\langle \dot{\psi}_n, (\langle K_{f,n} \rangle S_0 + \langle J_{f,n} \rangle S_0 + \langle M_{f,n} \rangle S_1)_G \dot{\psi}_n \rangle = \left( \langle K_{f,n} \rangle + \langle J_{f,n} \rangle + \frac{\langle M_{f,n} \rangle}{q+1} w(\lambda_n) \right) \cdot \|\dot{\psi}_n\|^2,$$

where we used (3.6), that  $\mathcal{A}_{G_N} \dot{\psi}_n = w(\lambda_n) \dot{\psi}_n$ . But

$$\langle K_{f,n} \rangle + \langle J_{f,n} \rangle + \frac{\langle M_{f,n} \rangle}{q+1} w(\lambda_n) = \frac{1}{N} \sum_{(x_0, x_1) \in B} \frac{1}{s^2(\lambda_n)} \int_0^L f_{(x_0, x_1)}(t) \\ \cdot \left[ S_{\lambda_n}^2(L-t) + S_{\lambda_n}^2(t) + \frac{2S_{\lambda_n}(L-t)S_{\lambda_n}(t)}{q+1} w(\lambda_n) \right] dt \\ = \frac{\kappa_{\lambda_n}}{N} \sum_{(x_0, x_1) \in B} \int_0^L f_{(x_0, x_1)}(t) \Psi_{\lambda_n}(t) dt = 2 \kappa_{\lambda_n} \langle f \rangle_{\lambda_n},$$

where we used (A.4). Hence, (4.9) is proved by comparing the expressions.

(3) If  $f$  is locally constant,  $f_{(x_0, x_1)}(t) = c_{(x_0, x_1)}$ , we get  $\langle f \rangle_{\lambda_n} = \langle f \rangle$ . Also,

$$\langle \dot{\psi}_n, (K_{f,n})_G \dot{\psi}_n \rangle = \langle \dot{\psi}_n, (J_{f,n})_G \dot{\psi}_n \rangle = a_{n,L} \langle \dot{\psi}_n, (K_f)_G \dot{\psi}_n \rangle, \\ \langle \dot{\psi}_n, (M_{f,n})_G \dot{\psi}_n \rangle = b_{n,L} \langle \dot{\psi}_n, (M_f)_G \dot{\psi}_n \rangle \\ \langle K_{f,n} \rangle = \langle J_{f,n} \rangle = a_{n,L} \langle K_f \rangle \quad \text{and} \quad \langle M_{f,n} \rangle = b_{n,L} \langle M_f \rangle.$$

where  $a_{n,L} = \frac{1}{s^2(\lambda_n)} \int_0^L S_{\lambda_n}^2(t) dt$  and  $b_{n,L} = \frac{2}{s^2(\lambda_n)} \int_0^L S_{\lambda_n}(L-t)S_{\lambda_n}(t) dt$ . Since  $|a_{n,L}| \leq C_I L$  and  $|b_{n,L}| \leq C'_I L$  for  $\lambda_n \in \bar{I} \subset \sigma_{ac}(H_{\mathbf{T}_q})$ , the proof is complete.  $\square$

**4.2. From the variance to the non-backtracking variance.** The aim of this Section is to express (4.9) in terms of non-backtracking variances, following [2].

As in [2, §3], we introduce the stochastic operators  $\mathcal{S} : \mathcal{H}_k \rightarrow \mathcal{H}_k$ ,

$$\mathcal{S} := \frac{\mathcal{A}_G}{q+1} \quad \text{if } k=0 \quad \text{and} \quad \mathcal{S} := \frac{1}{q}\mathcal{B} \quad \text{for } k \geq 1.$$

Here,  $(\mathcal{B}f)(x_0; x_k) = \sum_{x_{k+1} \in \mathcal{N}_{x_k} \setminus \{x_{k-1}\}} f(x_1, x_{k+1})$ . We then define

$$\mathcal{S}_T := \frac{1}{T} \sum_{r=0}^{T-1} (T-r)\mathcal{S}^r \quad \text{and} \quad \tilde{\mathcal{S}}_T := \frac{1}{T} \sum_{r=1}^T \mathcal{S}^r.$$

We then observe that

$$K^{\lambda_j} = (I - \mathcal{S})(\mathcal{S}_T K^{\lambda_j}) + \tilde{\mathcal{S}}_T K^{\lambda_j}.$$

On the other hand, define  $i_k : \mathcal{H}_k \rightarrow \mathcal{H}_{k+1}$  and  $\nabla : \mathcal{H}_k \rightarrow \mathcal{H}_{k+1}$  by  $(i_k J)(x_0; x_{k+1}) = J(x_1; x_{k+1})$  and  $(\nabla K)(x_0; x_{k+1}) = K(x_1; x_{k+1}) - K(x_0; x_k)$ . We then have

$$\nabla^* i_k K = q(I - \mathcal{S})K.$$

Hence, using the fact that  $\text{Var}^I(K_1 + K_2) \leq 2\text{Var}^I(K_1) + 2\text{Var}^I(K_2)$ , we obtain

$$(4.10) \quad \text{Var}^I(K^\lambda - \langle K^\lambda \rangle S_k) \leq 2q^{-2} \text{Var}^I(\nabla^* i_k \mathcal{S}_T K^\lambda) + 2\text{Var}^J(\tilde{\mathcal{S}}_T K^\lambda - \langle K^\lambda \rangle S_k),$$

where  $S_k$  is as in (4.6).

Next, using [2, Lemma 6.4] and the idea in the proof of [2, Lemma 7.8], we express

$$2i \text{Im } \mu_j^- \cdot \langle \psi_j, (\nabla^* K)_G \psi_j \rangle = \langle f_j^*, K_B f_j \rangle - \langle g_j^*, K_B g_j \rangle,$$

where  $g_j, g_j^*$  are defined like  $f_j, f_j^*$ , except that  $\mu_j^-$  are replaced by  $\overline{\mu_j^-}$ . It follows that<sup>1</sup>

$$\text{Var}^I(\nabla^* i_k \mathcal{S}_T K^\lambda) \leq \max_j \frac{1}{|\text{Im } \mu_j^-|^2} \text{Var}_{\text{nb}}^I(i_k \mathcal{S}_T K^\lambda).$$

Assuming  $\bar{I} \subset \sigma_{ac}(H_{\mathbf{T}_q})$  is in a fixed band, so that  $\overline{w(\bar{I})} \subset (-2\sqrt{q}, 2\sqrt{q})$ , this maximum is bounded by some  $C_I$ . Back to (4.10), note that for any  $K$ , we have

$$\tilde{\mathcal{S}}_T \langle K \rangle S_k = \langle K \rangle S_k,$$

since  $\mathcal{S}\mathbf{1} = \mathbf{1}$ . We thus finally get

$$(4.11) \quad \text{Var}^I(K^\lambda - \langle K^\lambda \rangle S_k) \leq 2q^{-2} C_I \text{Var}_{\text{nb}}^I(i_k \mathcal{S}_T K^\lambda) + 2\text{Var}^I(\tilde{\mathcal{S}}_T [K^\lambda - \langle K^\lambda \rangle S_k]).$$

*Taking advantage of the invariance.* Let us introduce the operators  $\mathcal{R}_{n,r}^{\lambda_j} : \mathcal{H}_k \rightarrow \mathcal{H}_{n+k}$  defined by

$$(\mathcal{R}_{n,r}^{\lambda_j} K^{\lambda_j})(x_0; x_{n+k}) = \overline{(\mu_j^-)^{n-r} (\mu_j^-)^r} K^{\lambda_j}(x_{n-r}; x_{n-r+k}).$$

<sup>1</sup>Actually, instead of  $\text{Var}_{\text{nb}}^I(J)$ , we should have  $\frac{1}{2} \text{Var}_{\text{nb}}^I(J) + \frac{1}{2} \widetilde{\text{Var}_{\text{nb}}^I(J)}$ , where  $\widetilde{\text{Var}_{\text{nb}}^I}$  is defined using  $g_j, g_j^*$  instead. Since this variance is controlled exactly like  $\text{Var}_{\text{nb}}^I(J)$ , except for replacing  $\mu_j^-$  by  $\overline{\mu_j^-}$  in the arguments, we omitted it for transparency.

Using (3.7), we have  $\text{Var}_{\text{nb}}^I(K^\lambda) = \text{Var}_{\text{nb}}^I(\frac{1}{n} \sum_{r=1}^n \mathcal{R}_{n,r}^\lambda K^\lambda)$ . Summarizing and using (4.9), we see that

$$\begin{aligned}
 & \frac{1}{N(I)} \sum_{\lambda_j \in I} \left| \langle \psi_j, f\psi_j \rangle - \langle f \rangle_{\lambda_j} \right|^2 \\
 (4.12) \quad & \leq C_I \text{Var}_{\text{nb}}^I \left( \frac{1}{n} \sum_{r=1}^n \mathcal{R}_{n,r}^\lambda i_0 \mathcal{S}_T K_f^\lambda \right) + C_I \text{Var}_{\text{nb}}^I \left( \frac{1}{n} \sum_{r=1}^n \mathcal{R}_{n,r}^\lambda i_0 \mathcal{S}_T J_f^\lambda \right) \\
 & \quad + C_I \text{Var}_{\text{nb}}^I \left( \frac{1}{n} \sum_{r=1}^n \mathcal{R}_{n,r}^\lambda i_1 \mathcal{S}_T M_f^\lambda \right) + 2 \text{Var}^I \left( \tilde{\mathcal{S}}_T [K_f^\lambda - \langle K_f^\lambda \rangle S_0] \right) \\
 & \quad + 2 \text{Var}^I \left( \tilde{\mathcal{S}}_T [J_f^\lambda - \langle J_f^\lambda \rangle S_0] \right) + 2 \text{Var}^I \left( \tilde{\mathcal{S}}_T [M_f^\lambda - \langle M_f^\lambda \rangle S_1] \right),
 \end{aligned}$$

where  $K_f^{\lambda_j} := K_{f,j}$ ,  $J_f^{\lambda_j} := J_{f,j}$  and  $M_f^{\lambda_j} := M_{f,j}$ .

**4.3. Controlling the variances.** We next try to bound the variances in terms of Hilbert-Schmidt norms. For clarity, we start with energy-independent observables  $K$ .

**Lemma 4.3.** *Suppose that  $I$  is in a fixed band of AC spectrum. Then for any  $K \in \mathcal{H}_k$ ,*

$$\text{Var}^I(K) \leq C_{L,I} \frac{N}{N(I)} \|K\|_{\mathcal{H}}^2 + c_{k,q,L,I} \frac{N}{N(I)} \frac{\#\{x : \rho_{G_N}(x) < k\}}{N} \|K\|_\infty^2.$$

where  $\|K\|_{\mathcal{H}}^2 = \frac{1}{N} \sum_{(x_0; x_k) \in B_k} |K(x_0; x_k)|^2$  and  $\|K\|_\infty = \max_{(x_0; x_k)} |K(x_0; x_k)|$ .

*Proof.* As we showed in Lemma 3.1, if  $\lambda$  is an eigenvalue of the quantum graph with eigenvector  $\psi$ , then  $w(\lambda)$  is an eigenvalue of  $\mathcal{A}_{G_N}$  with eigenvector  $\dot{\psi}$ . By assumption, if  $\lambda_n, \lambda_j \in I$  and  $\lambda_n \neq \lambda_j$ , then  $w(\lambda_n) \neq w(\lambda_j)$ , so the eigenvectors  $\dot{\psi}_n$  and  $\dot{\psi}_j$  are orthogonal. If  $\lambda_n = \lambda_j = \lambda$  and  $\psi_n \perp \psi_j$ , we find as in (4.1) that  $\langle \dot{\psi}_n, \dot{\psi}_j \rangle_{\kappa_\lambda} = 0$ , so  $\dot{\psi}_n$  and  $\dot{\psi}_j$  are again orthogonal. If  $\phi_n := \dot{\psi}_n / \|\dot{\psi}_n\|$ , we may thus complete  $(\phi_n)_{\lambda_n \in I}$  to an orthonormal basis  $(\phi_n)_{n=1}^N$  of  $\ell^2(G_N)$ . Hence, recalling (4.4) and using (4.7), we get

$$\begin{aligned}
 \text{Var}^I(K) & \leq \frac{C_{L,I}}{N(I)} \sum_{\lambda_n \in I} \|K_G \dot{\psi}_n\|^2 \leq \frac{C_{L,I}^2}{N(I)} \sum_{\lambda_n \in I} \|K_G \phi_n\|^2 \\
 & \leq \frac{C_{L,I}^2}{N(I)} \sum_{n=1}^N \|K_G \phi_n\|^2 = C_{L,I}^2 \frac{N}{N(I)} \|K_G\|_{HSN}^2,
 \end{aligned}$$

where  $\|K_G\|_{HSN}^2 = \frac{1}{N} \sum_{x,y \in V_N} |K_G(x,y)|^2$ . We may then apply [2, Proposition 2], which tells us that

$$\|K_G\|_{HSN}^2 \leq \|K\|_{\mathcal{H}}^2 + c_{k,q,L,I} \frac{\#\{x : \rho_{G_N}(x) < k\}}{N} \|K\|_\infty^2$$

for some  $c_{k,q,L,I} > 0$ . This concludes the proof.  $\square$

**Remark 4.4.** The fraction  $\frac{N}{N(I)}$  arising in Lemma 4.3 is asymptotically bounded. In fact, as  $w$  is injective on  $I$ , we have  $N(I) = \#\{w(\lambda_j) \in w(I)\}$ . Moreover,  $w(\lambda_j)$  are all eigenvalues of  $\mathcal{A}_{G_N}$ . Conversely, if  $\mu_j \in w(I)$ ,  $\mu_j = w(t_j)$  is an eigenvalue of  $\mathcal{A}_{G_N}$ , we may construct an eigenfunction of  $H_{\mathbf{G}_N}$  with eigenvalue  $t_j$  via formula (4.2). Hence,  $N(I)$  is the number of eigenvalues of  $\mathcal{A}_{G_N}$  in  $w(I)$ , so  $\frac{N(I)}{N} \rightarrow \mu_o(w(I)) := \langle \delta_o, \chi_{w(I)}(\mathcal{A}_{\mathbb{T}_q}) \delta_o \rangle$  by the law of Kesten-McKay [22, 28]. Since  $w(I) \subset (-2\sqrt{q}, 2\sqrt{q})$ , we have  $\mu_o(w(I)) \geq C_I > 0$ . Hence<sup>2</sup>,  $\limsup_{N \rightarrow \infty} \frac{N}{N(I)} \leq C_I^{-1}$ .

<sup>2</sup>Note that this argument continues to hold on any  $I$  on which  $w$  is injective and  $I \cap \sigma_2 = \emptyset$ . If moreover  $I \cap \sigma_1 = \emptyset$ , we see that  $\frac{N(I)}{N} \rightarrow \mu_o(w(I)) = 0$  because  $w(I) \cap (-2\sqrt{q}, 2\sqrt{q}) = \emptyset$ .



From here, one may adapt the “ultra-short” proof of [2, §4], see also [6, §2], to obtain a simple proof of Theorem 2.1. Indeed, in this case, the observables are independent of the energy (cf. Lemma 4.2.(3)). Since our main concern is Theorem 2.2, we consider the general case directly. As the observables now depend on the energy (Lemma 4.2.(2)), we use the non-backtracking quantum variance and give the following analogue of Lemma 4.3.

**Lemma 4.5.** *Let  $K^\lambda \in \mathcal{H}_k$ ,  $k \geq 1$ , depend on  $\lambda$  as in (4.12), and let  $I$  be in a fixed band of the AC spectrum. Then*

$$(4.13) \quad \limsup_{N \rightarrow \infty} \text{Var}_{\text{nb}}^I(K^\lambda) \leq C(L, I) \limsup_{N \rightarrow \infty} \int_I \|K^\lambda\|_{\mathcal{H}}^2 d\lambda,$$

where  $\|K^\lambda\|_{\mathcal{H}}^2 = \frac{1}{N} \sum_{(x_0; x_k)} |K^\lambda(x_0; x_k)|^2$ . A similar bound holds for  $\text{Var}^I(K^\lambda)$  if  $k = 0, 1$ .

*Proof.* As in Lemma 4.3, if  $\phi_j := \psi_j / \|\psi_j\|$ , we may complete  $(\phi_j)_{w(\lambda_j) \in w(I)}$  to an orthonormal basis  $(\phi_j)_{j=1}^N$  of  $\ell^2(G_N)$ . Denote  $\tilde{f}_j = f_j / \|\psi_j\|$ . We then have for  $m_j = w(\lambda_j)$ ,

$$\text{Var}_{\text{nb}}^I(K^\lambda) \leq \frac{C_{L,I}^2}{N(I)} \sum_{m_j \in w(I)} \|K_B^{w^{-1}(m_j)} \tilde{f}_j\|^2.$$

As previously mentioned,  $m_j = w(\lambda_j)$  are eigenvalues of  $\mathcal{A}_{G_N}$  and we denote the rest of the eigenvalues corresponding to the completed base  $(\phi_j)_{j=1}^N$  also by  $\{m_j\}_{j=1}^N$ .

Since  $w$  is injective and continuous on  $I$ , it is either strictly increasing or strictly decreasing. Assuming the former without loss of generality, we now define  $J^{m_j} = K^{w^{-1}(m_j)}$  if  $m_j \in w(I)$ ,  $J^{m_j} = K^b$  if  $m_j \geq w(b)$  and  $J^{m_j} = K^a$  if  $m_j \leq w(a)$ , where  $[a, b] := \bar{I}$ .

Let  $\chi$  be a continuous function, with  $0 \leq \chi \leq 1$ ,  $\chi = 1$  on  $w(I)$ ,  $\chi = 0$  outside a  $\delta$ -enlargement  $w(I)_\delta$  of  $w(I)$ . We have

$$\begin{aligned} \text{Var}_{\text{nb}}^I(K^\lambda) &\leq \frac{C_{L,I}^2}{N(I)} \sum_{m_j \in w(I)} \|J_B^{m_j} \tilde{f}_j\|^2 \leq \frac{C_{L,I}^2}{N(I)} \sum_{j=1}^N \chi(m_j) \|J_B^{m_j} \tilde{f}_j\|^2 \\ &= \frac{C^2}{N(I)} \sum_{j=1}^N \chi(m_j) \sum_{(x_0, x_1) \in B} \sum_{(x_2; x_k), (y_2; y_k)} J^{m_j}(x_0; x_k) \overline{J^{m_j}(x_0; y_k)} \tilde{f}_j(x_{k-1}, x_k) \overline{\tilde{f}_j(y_{k-1}, y_k)}, \end{aligned}$$

where the last sum runs over paths  $(x_2; x_k), (y_2; y_k)$  with  $x_2, y_2 \in \mathcal{N}_{x_1} \setminus \{x_0\}$ , and  $(x_0; y_k) := (x_0, x_1, y_2, \dots, y_k)$ . To keep the notations clear, we now assume that

$$(4.14) \quad J^m(x_0; x_k) = F(m)K(x_0; x_k)$$

for some continuous function  $F$  on  $\mathbb{R}$ ,  $|F(m)| \leq c_I$ . We comment on the case of the operators in (4.12) at the end of the proof.

Recalling (3.5) and the formula  $f(\mathcal{A}_G)(x, y) = \sum_{j=1}^N f(m_j) \phi_j(x) \overline{\phi_j(y)}$ , we get

$$\begin{aligned} \text{Var}_{\text{nb}}^I(K^\lambda) &\leq \frac{C^2}{N(I)} \sum_{j=1}^N |F(m_j)|^2 \chi(m_j) \sum_{(x_0, x_1) \in B} \sum_{(x_2; x_k), (y_2; y_k)} K(x_0; x_k) \overline{K(x_0; y_k)} \\ &\quad \times \left( \left( \phi_j(x_{k-1}) - \nu(m_j) \phi_j(x_k) \right) \overline{\left( \phi_j(y_{k-1}) - \nu(m_j) \phi_j(y_k) \right)} \right) \\ &= \frac{C^2}{N(I)} \sum_{(x_0, x_1) \in B} \sum_{(x_2; x_k), (y_2; y_k)} K(x_0; x_k) \overline{K(x_0; y_k)} \left[ (\chi|F|^2)(\mathcal{A}_G)(x_k, y_k) \right. \\ &\quad \left. - (\bar{\nu}\chi|F|^2)(\mathcal{A}_G)(x_k, y_{k-1}) - (\nu\chi|F|^2)(\mathcal{A}_G)(x_{k-1}, y_k) \right. \\ &\quad \left. + \frac{1}{q} (\chi|F|^2)(\mathcal{A}_G)(x_{k-1}, y_{k-1}) \right], \end{aligned}$$

where  $\nu(m) = \mu^-(w^{-1}(m))$ , and we used the fact that  $|\mu^-(\lambda)|^2 = \frac{1}{q}$  for  $\lambda \in I_\delta$ .

As in [2, Proposition 3], see also [4, Theorem 4.1], we approximate each continuous function  $\chi(m)|F(m)|^2$ ,  $\nu(m)\chi(m)|F(m)|^2$ ,  $\bar{\nu}(m)\chi(m)|F(m)|^2$  by polynomials  $Q_1, Q_2, Q_3$  up to an  $\epsilon$  error. Let  $d_\epsilon = \max(\deg(Q_i))$ .

Recall that  $\rho_{G_N}(x)$  is the largest  $\rho$  such that the sub-graph contained in a ball of radius  $\rho$  centered at  $x$  has no closed cycles. If  $\rho_G(x_0) \geq d_\epsilon + k =: d_{\epsilon,k}$ , then we have  $Q_i(\mathcal{A}_G)(x, y) = Q_i(\mathcal{A}_{\mathbb{T}_q})(\tilde{x}, \tilde{y})$  for  $x, y$  in a  $k$ -ball around  $x_0$  where  $\tilde{x}, \tilde{y} \in \mathbb{T}_q$  are lifts of  $x, y \in G$ . The ‘‘bad’’  $x_0$  with  $\rho_G(x_0) < d_{\epsilon,k}$  induce an error term. Finally, we replace the polynomials  $Q_i$  by the original functions, yielding an  $\epsilon$  error. In the end, we get

$$\begin{aligned} \text{Var}_{\text{nb}}^I(K^\lambda) &\leq \frac{C^2}{N(I)} \sum_{\rho_G(x_0) \geq d_{\epsilon,k}} \sum_{x_1 \sim x_0} \sum_{(x_2; x_k), (y_2; y_k)} K(x_0; x_k) \overline{K(x_0; y_k)} \\ &\quad \left[ (\chi|F|^2)(\mathcal{A}_{\mathbb{T}_q})(\tilde{x}_k, \tilde{y}_k) - (\bar{\nu}\chi|F|^2)(\mathcal{A}_{\mathbb{T}_q})(\tilde{x}_k, \tilde{y}_{k-1}) - (\nu\chi|F|^2)(\mathcal{A}_{\mathbb{T}_q})(\tilde{x}_{k-1}, \tilde{y}_k) \right. \\ &\quad \left. + \frac{1}{q}(\chi|F|^2)(\mathcal{A}_{\mathbb{T}_q})(\tilde{x}_{k-1}, \tilde{y}_{k-1}) \right] + \frac{C'N(q+1)q^{2k}}{N(I)} \epsilon \|K\|_\infty^2 \\ &\quad + C' \frac{\#\{\rho_G(x) < d_{\epsilon,k}\}}{N(I)} (q+1)q^{2k} \|K\|_\infty^2 \end{aligned}$$

for some  $C' > 0$ .

Since  $\mathcal{A}_{\mathbb{T}_q}$  has purely absolutely continuous spectrum, it follows that  $f(\mathcal{A}_{\mathbb{T}_q})(v, w) = \frac{1}{\pi} \int f(m) \text{Im } G^m(v, w) dm$ , where  $G^m(v, w) = \lim_{\eta \downarrow 0} (\mathcal{A}_{\mathbb{T}_q} - (m + i\eta))^{-1}(v, w)$  (see for instance [7, Lemma 3.6]). Hence,

$$\begin{aligned} \text{Var}_{\text{nb}}^I(K^\lambda) &\leq \frac{c}{N(I)} \sum_{\rho_G(x_0) \geq d_{\epsilon,k}} \sum_{x_1 \sim x_0} \sum_{(x_2; x_k), (y_2; y_k)} \left( K(x_0; x_k) \overline{K(x_0; y_k)} \right. \\ &\quad \left. \int_{w(I)_\delta} \chi(m)|F(m)|^2 \left[ \text{Im } G^m(\tilde{x}_k, \tilde{y}_k) - \overline{\mu^-(w^{-1}m)} \text{Im } G^m(\tilde{x}_k, \tilde{y}_{k-1}) \right. \right. \\ &\quad \left. \left. - \mu^-(w^{-1}m) \text{Im } G^m(\tilde{x}_{k-1}, \tilde{y}_k) + \frac{1}{q} \text{Im } G^m(\tilde{x}_{k-1}, \tilde{y}_{k-1}) \right] dm \right) \\ &\quad + \frac{C'N(q+1)q^{2k}}{N(I)} \epsilon \|K\|_\infty^2 + C' \frac{\#\{\rho_G(x) < d_{\epsilon,k}\}}{N(I)} (q+1)q^{2k} \|K\|_\infty^2. \end{aligned}$$

If  $d(v, w) = s$ , then  $\text{Im } G^m(v, w) = c_m \Phi_m(s)$ , where  $\Phi_m$  is the spherical function defined in (3.10) and  $c_m = \text{Im } G^m(o, o)$ . Let  $e_k = (x_{k-1}, x_k)$ ,  $e'_k = (y_{k-1}, y_k)$  and suppose  $e_k \neq e'_k$ . Then there is a path  $(v_0, \dots, v_s)$  with  $v_0 = \tilde{x}_k$ ,  $v_1 = \tilde{x}_{k-1}$ ,  $v_{s-1} = \tilde{y}_{k-1}$ ,  $v_s = \tilde{y}_k$ . So the term in square brackets takes the form

$$(4.15) \quad c_m \left( \Phi_m(s) - \left[ \overline{\mu^-(w^{-1}m)} + \mu^-(w^{-1}m) \right] \Phi_m(s-1) + \frac{1}{q} \Phi_m(s-2) \right) = 0,$$

where we used the fact that  $q(\overline{\mu^-(t)} + \mu^-(t)) = w(t)$ , as well as (3.12). Since  $\rho_G(x_0) \geq d_{\epsilon,k}$ ,  $e_k \neq e'_k$  iff  $(y_2; y_k) \neq (x_2; x_k)$ . We thus see that the sum vanishes if  $(y_2; y_k) \neq (x_2; x_k)$ , so only the diagonal terms may be nonzero. Recalling (3.13), we get

$$\begin{aligned} \text{Var}_{\text{nb}}^I(K^\lambda) &\leq \frac{c}{N(I)} \sum_{\rho_G(x_0) \geq d_{\epsilon,k}} \sum_{(x_1; x_k)} \\ (4.16) \quad &\int_{w(I)_\delta} \left( |F(m)|^2 |K(x_0; x_k)|^2 \cdot \chi(m) c_m \left[ \frac{q+1}{q} - \frac{m^2}{q(q+1)} \right] dm \right) \\ &\quad + \frac{C'N(q+1)q^{2k}}{N(I)} \epsilon \|K\|_\infty^2 + C' \frac{\#\{\rho_G(x) < d_{\epsilon,k}\}}{N(I)} (q+1)q^{2k} \|K\|_\infty^2. \end{aligned}$$

Recall that  $F(m)K(x_0; x_k) = J^m(x_0; x_k) = K^{w^{-1}(m)}(x_0; x_k)$ . We finally change the variables  $m = w(\lambda)$ . Note that  $|w'(\lambda)| \leq C$  (in fact,  $\lambda \mapsto w(\lambda)$  is analytic [30, p. 10]) and

$c_m \leq c_I$ . Using Remark 4.4, we see that the first term of (4.16) is bounded by

$$C(L, I) \limsup_{N \rightarrow \infty} \int_{I_{\delta'}} \|K^\lambda\|_{\mathcal{H}}^2 d\lambda$$

where  $\delta'$  goes to zero with  $\delta$ . The last term in (4.16) goes to zero as  $N \rightarrow \infty$  thanks to **(BST)**. Therefore, the proof is now complete for  $\text{Var}_{\text{nb}}^I(K^\lambda)$  by first taking  $N \rightarrow \infty$ , then  $\epsilon \downarrow 0$  and  $\delta \downarrow 0$ .

The control of  $\text{Var}^I(K^\lambda)$  for  $k = 0$  is similar. For  $k = 1$ , we find that

$$\text{Var}^I(K^\lambda) \leq \frac{C}{N(I)} \int_{w(I)} \chi(m) \sum_{x_0 \in V} \sum_{x_1, y_1 \in \mathcal{N}_{x_0}} J^m(x_0, x_1) \overline{J^m(x_0, y_1)} \text{Im} G^m(\tilde{x}_1, \tilde{y}_1) dm.$$

We write the sums  $\sum_{x \in V} \sum_{x_1, y_1 \in \mathcal{N}_{x_0}}$  in the form

$$\sum_{(x_0, x_1)} |J^m(x_0, x_1)|^2 \text{Im} G^m(\tilde{x}_1, \tilde{x}_1) + \sum_{(x_{-1}, x_0, x_1)} J^m(x_0, x_1) \overline{J^m(x_0, x_{-1})} \text{Im} G^m(\tilde{x}_1, \tilde{x}_{-1}).$$

The first term has the required form, since  $\text{Im} G^m(v, v)$  is bounded. For the second term, we have  $\text{Im} G^m(\tilde{x}_1, \tilde{x}_{-1}) = c_m \Phi_m(2)$ . Using the Cauchy-Schwarz inequality, we have  $|\sum_{(x_{-1}, x_0, x_1)} J^m(x_0, x_1) \overline{J^m(x_0, x_{-1})}| \leq q \sum_{(x_0, x_1)} |J^m(x_0, x_1)|^2$ . This completes the proof.

Finally, in the previous proof we assumed (4.14) that  $J^m(x_0; x_k) = F(m)K(x_0; x_k)$ . In (4.12), for  $k = 1, 2$ , we have operators of the form

$$J^m(x_0; x_{n+k}) = \frac{1}{n} \sum_{r=1}^n \sum_{s=0}^{T-1} \sum_{y_1, \dots, y_{s+1} \in B(x_{n-r+1}, s+1)} \int_0^L F_{T,t,r,s}(m) f_{(y_s, y_{s+1})}(t) dt,$$

where  $F_{T,t,r,s}$  is continuous and uniformly bounded  $|F_{T,t,r,s}(m)| \leq C_I$ . Looking back at the proof, we see that all arguments continue to apply; the notations simply become more cumbersome. Here, instead of  $|F(m)|^2$ , we have  $F_{T,t,r,s}(m) \overline{F_{T,t',r',s'}(m)}$ , which we approximate by a polynomial as before.  $\square$

**4.4. End of the proof.** We may now conclude the proof of Theorem 2.2. From (4.12) and Lemma 4.5, we have

$$(4.17) \quad \begin{aligned} & \limsup_{N \rightarrow \infty} \frac{1}{N(I)} \sum_{\lambda_j \in I} |\langle \psi_j, f \psi_j \rangle - \langle f \rangle_{\lambda_j}|^2 \\ & \leq C'(L, I) \limsup_{N \rightarrow \infty} \int_I \left[ \left\| \frac{1}{n} \sum_{r=1}^n \mathcal{R}_{n,r}^\lambda i_0 \mathcal{S}_T K_f^\lambda \right\|_{\mathcal{H}}^2 + \left\| \frac{1}{n} \sum_{r=1}^n \mathcal{R}_{n,r}^\lambda i_0 \mathcal{S}_T J_f^\lambda \right\|_{\mathcal{H}}^2 \right. \\ & \quad + \left\| \frac{1}{n} \sum_{r=1}^n \mathcal{R}_{n,r}^\lambda i_1 \mathcal{S}_T M_f^\lambda \right\|_{\mathcal{H}}^2 + \left\| \tilde{\mathcal{S}}_T [K_f^\lambda - \langle K_f^\lambda \rangle S_0] \right\|_{\mathcal{H}}^2 \\ & \quad \left. + \left\| \tilde{\mathcal{S}}_T [J_f^\lambda - \langle J_f^\lambda \rangle S_0] \right\|_{\mathcal{H}}^2 + \left\| \tilde{\mathcal{S}}_T [M_f^\lambda - \langle M_f^\lambda \rangle S_1] \right\|_{\mathcal{H}}^2 \right] d\lambda, \end{aligned}$$

The first three terms may be controlled in exactly the same way as in [2, p. 661–662], and give vanishing contributions as  $N \rightarrow \infty$ , followed by  $n \rightarrow \infty$ , for any  $T$ .

The remaining terms are of the form  $\|\tilde{\mathcal{S}}_T J^\lambda\|_{\mathcal{H}}$ , for  $J^\lambda = K^\lambda - \langle K^\lambda \rangle S_k$ . By [2, Remark 2.2], we know that  $\|\tilde{\mathcal{S}}_T J^\lambda\|_{\mathcal{H}_k} \leq \frac{c_{k,\beta}}{T} \|J^\lambda\|_{\mathcal{H}_k}$ . The result follows by letting  $T \rightarrow \infty$ .

## 5. THE CASE OF INTEGRAL OPERATORS

The aim of this Section is to prove Theorem 2.3.

We start by decomposing the integral in (2.13). Using (4.2), and the notation  $S_n(x) := S_{\lambda_n}(x)$ , we have

$$\int_0^L \overline{\psi_n(x_b)} (K_k \psi_n)(x_b) dx_b = \sum_{(b_1; b_k)_b} \int_0^L \int_0^L K_{b, b_k}(x_b, y_{b_k}) \left[ \frac{S_n(L - x_b) \overline{\psi_n(o_b)}}{s(\lambda_n)} + \frac{S_n(x_b) \overline{\psi_n(t_b)}}{s(\lambda_n)} \right] \cdot \left[ \frac{S_n(L - y_{b_k}) \psi_n(o_{b_k})}{s(\lambda_n)} + \frac{S_n(y_{b_k}) \psi_n(t_{b_k})}{s(\lambda_n)} \right] dx_b dy_{b_k},$$

so denoting  $b_1 = (x_0, x_1)$  and  $b_k = (x_{k-1}, x_k)$ , we have

$$\begin{aligned} & 2 \langle \psi_n, K_k \psi_n \rangle \\ &= \sum_{(x_0; x_k) \in B_k} \overline{\psi_n(x_0)} \psi_n(x_{k-1}) \int_0^L \int_0^L K_{b_1, b_k}(r_{b_1}, s_{b_k}) \frac{S_n(L - r_{b_1}) S_n(L - s_{b_k})}{s^2(\lambda_n)} dr_{b_1} ds_{b_k} \\ & \quad + \sum_{(x_0; x_k) \in B_k} \overline{\psi_n(x_1)} \psi_n(x_k) \int_0^L \int_0^L K_{b_1, b_k}(r_{b_1}, s_{b_k}) \frac{S_n(r_{b_1}) S_n(s_{b_k})}{s^2(\lambda_n)} dr_{b_1} ds_{b_k} \\ & \quad + \sum_{(x_0; x_k) \in B_k} \overline{\psi_n(x_0)} \psi_n(x_k) \int_0^L \int_0^L K_{b_1, b_k}(r_{b_1}, s_{b_k}) \frac{S_n(L - r_{b_1}) S_n(s_{b_k})}{s^2(\lambda_n)} dr_{b_1} ds_{b_k} \\ & \quad + \sum_{(x_0; x_k) \in B_k} \overline{\psi_n(x_1)} \psi_n(x_{k-1}) \int_0^L \int_0^L K_{b_1, b_k}(r_{b_1}, s_{b_k}) \frac{S_n(r_{b_1}) S_n(L - s_{b_k})}{s^2(\lambda_n)} dr_{b_1} ds_{b_k}. \end{aligned}$$

Assume  $k \geq 2$  (the case  $k = 1$  is similar to Theorem 2.2). If  $J_{K,n} \in \mathcal{H}_k$ ,  $M_{K,n} \in \mathcal{H}_{k-1}$  and  $P_{K,n} \in \mathcal{H}_{k-2}$  are defined by

$$J_{K,n}(x_0; x_k) = \int_0^L \int_0^L K_{(x_0; x_k)}(r_{b_1}, s_{b_k}) \frac{S_n(L - r_{b_1}) S_n(s_{b_k})}{s^2(\lambda_n)} dr_{b_1} ds_{b_k},$$

$$\begin{aligned} & M_{K,n}(x_0; x_{k-1}) \\ &= \sum_{x_k \in \mathcal{N}_{x_{k-1}} \setminus \{x_{k-2}\}} \int_0^L \int_0^L K_{(x_0; x_k)}(r_{b_1}, s_{b_k}) \frac{S_n(L - r_{b_1}) S_n(L - s_{b_k})}{s^2(\lambda_n)} dr_{b_1} ds_{b_k} \\ & \quad + \sum_{x_{-1} \in \mathcal{N}_{x_0} \setminus \{x_1\}} \int_0^L \int_0^L K_{(x_{-1}; x_{k-1})}(r_{b_0}, s_{b_{k-1}}) \frac{S_n(r_{b_0}) S_n(s_{b_{k-1}})}{s^2(\lambda_n)} dr_{b_0} ds_{b_{k-1}} \end{aligned}$$

$$\begin{aligned} & P_{K,n}(x_1; x_{k-1}) \\ &= \sum_{x_0 \in \mathcal{N}_{x_1} \setminus \{x_2\}, x_k \in \mathcal{N}_{x_{k-1}} \setminus \{x_{k-2}\}} \int_0^L \int_0^L K_{(x_0; x_k)}(r_{b_1}, s_{b_k}) \frac{S_n(r_{b_1}) S_n(L - s_{b_k})}{s^2(\lambda_n)} dr_{b_1} ds_{b_k}, \end{aligned}$$

where  $K_{(x_0; x_k)} := K_{(x_0, x_1), (x_{k-1}, x_k)}$ , we thus get

$$(5.1) \quad 2 \langle \psi_n, K_k \psi_n \rangle = \langle \dot{\psi}_n, (J_{K,n} + M_{K,n} + P_{K,n})_G \dot{\psi}_n \rangle.$$

As before, we can aim to control  $\text{Var}^J(K - \langle K \rangle S_k)$  for  $K \in \mathcal{H}_k$ , with  $S_k$  as in (4.6). This is why we consider

$$\begin{aligned} & \langle \dot{\psi}_n, (\langle J_{K,n} \rangle S_k + \langle M_{K,n} \rangle S_{k-1} + \langle P_{K,n} \rangle S_{k-2})_G \dot{\psi}_n \rangle \\ (5.2) \quad &= [\langle J_{K,n} \rangle \Phi_{w(\lambda_n)}(k) + \langle M_{K,n} \rangle \Phi_{w(\lambda_n)}(k-1) + \langle P_{K,n} \rangle \Phi_{w(\lambda_n)}(k-2)] \cdot \|\dot{\psi}_n\|^2 \\ &=: 2 \langle K_k \rangle_{n,L,\alpha} \end{aligned}$$

where in the first equality, we used (3.6) and the fact that if  $\tilde{\Phi}_k(\lambda) := \Phi_\lambda(k)$ , then<sup>3</sup>

$$(5.3) \quad (S_k)_G = \tilde{\Phi}_k(\mathcal{A}_G).$$

Using (4.7), this yields the formula

$$\langle K_k \rangle_{n,L,\alpha} = \frac{1}{N} \sum_{(b_1; b_k) \in B_k} \int_0^L \int_0^L K_{b_1, b_k}(x_{b_1}, y_{b_k}) \Psi_{\lambda_n, k}(x_{b_1}, y_{b_k}) dx_{b_1} dy_{b_k}$$

where, for  $k \geq 2$ ,

$$(5.4) \quad \Psi_{\lambda_n, k}(x_{b_1}, y_{b_k}) = \frac{1}{2\kappa_{\lambda_n}} \left[ \frac{S_n(L - x_{b_1})S_n(y_{b_k})}{s^2(\lambda_n)} \Phi_{w(\lambda_n)}(k) \right. \\ \left. + \frac{S_n(L - x_{b_1})S_n(L - y_{b_k}) + S_n(x_{b_1})S_n(y_{b_k})}{s^2(\lambda_n)} \Phi_{w(\lambda_n)}(k - 1) \right. \\ \left. + \frac{S_n(x_{b_1})S_n(L - y_{b_k})}{s^2(\lambda_n)} \Phi_{w(\lambda_n)}(k - 2) \right].$$

For  $k = 1$ , we have for  $b = (x_0, x_1)$ ,

$$(5.5) \quad \Psi_{\lambda_n, 1}(x_b, y_b) = \frac{1}{2\kappa_{\lambda_n}} \left[ \frac{S_n(L - x_b)S_n(L - y_b) + S_n(x_b)S_n(y_b)}{s^2(\lambda_n)} \right. \\ \left. + \frac{S_n(L - x_b)S_n(y_b) + S_n(x_b)S_n(L - y_b)}{s^2(\lambda_n)} \Phi_{w(\lambda_n)}(1) \right].$$

Using (3.11), we may express (5.4) as

$$\Psi_{\lambda_n, k}(x_{b_1}, y_{b_k}) = \frac{1}{2\kappa_{\lambda_n} \operatorname{Im} G_{\mathbf{T}_q}^{\lambda_n}(o, o) s^2(\lambda_n)} \\ \times \left[ S_n(L - x_{b_1})S_n(y_{b_k}) \operatorname{Im} G_{\mathbf{T}_q}^{\lambda_n}(\tilde{o}_{b_1}, \tilde{t}_{b_k}) + S_n(L - x_{b_1})S_n(L - y_{b_k}) \operatorname{Im} G_{\mathbf{T}_q}^{\lambda_n}(\tilde{o}_{b_1}, \tilde{o}_{b_k}) \right. \\ \left. + S_n(x_{b_1})S_n(y_{b_k}) \operatorname{Im} G_{\mathbf{T}_q}^{\lambda_n}(\tilde{t}_{b_1}, \tilde{t}_{b_k}) + S_n(x_{b_1})S_n(L - y_{b_k}) \operatorname{Im} G_{\mathbf{T}_q}^{\lambda_n}(\tilde{t}_{b_1}, \tilde{o}_{b_k}) \right].$$

Noting that  $\tilde{y}_{b_k} \mapsto \operatorname{Im} G_{\mathbf{T}_q}^{\lambda_n}(v, \tilde{y}_{b_k})$  is an eigenfunction of  $H_{\mathbf{T}_q}$ , we see as in (4.2) that

$$\operatorname{Im} G_{\mathbf{T}_q}^{\lambda_n}(v, \tilde{y}_{b_k}) = \frac{S_n(L - y_{b_k}) \operatorname{Im} G_{\mathbf{T}_q}^{\lambda_n}(v, \tilde{o}_{b_k}) + S_n(y_{b_k}) \operatorname{Im} G_{\mathbf{T}_q}^{\lambda_n}(v, \tilde{t}_{b_k})}{s(\lambda_n)},$$

so we deduce that for  $\tilde{x}_{b_1} \in \tilde{b}_1$ ,

$$(5.6) \quad \Psi_{\lambda_n, k}(x_{b_1}, y_{b_k}) \\ = \frac{1}{2\kappa_{\lambda_n} \operatorname{Im} G_{\mathbf{T}_q}^{\lambda_n}(o, o)} \frac{S_n(L - x_{b_1}) \operatorname{Im} G_{\mathbf{T}_q}^{\lambda_n}(\tilde{x}_0, \tilde{y}_{b_k}) + S_n(x_{b_1}) \operatorname{Im} G_{\mathbf{T}_q}^{\lambda_n}(\tilde{x}_1, \tilde{y}_{b_k})}{s(\lambda_n)} \\ = \frac{1}{2\kappa_{\lambda_n}} \cdot \frac{\operatorname{Im} G_{\mathbf{T}_q}^{\lambda_n}(\tilde{x}_{b_1}, \tilde{y}_{b_k})}{\operatorname{Im} G_{\mathbf{T}_q}^{\lambda_n}(o, o)},$$

<sup>3</sup>This is well-known : by definitions (4.5) and (4.6), we have  $[(S_k)_G \psi](x) = \frac{1}{(q+1)q^{k-1}} \sum_{(x_0; x_k), x_0=x} \psi(x_k)$  for  $k \geq 1$  and  $[(S_0)_G \psi](x) = \psi(x)$ . Since  $\tilde{\Phi}_0(\lambda) = 1$  and  $\tilde{\Phi}_1(\lambda) = \frac{\lambda}{q+1}$ , then (5.3) holds for  $k = 0, 1$ . Next,  $[(S_{k+1})_G \psi](x) = \frac{1}{(q+1)q^k} \sum_{(x_0; x_{k+1}), x_0=x} \psi(x_{k+1}) = \frac{1}{(q+1)q^k} \sum_{(x_0; x_k), x_0=x} [(\mathcal{A}_G \psi)(x_k) - \psi(x_{k-1})] = \frac{1}{q} [(\mathcal{S}_k \mathcal{A}_G \psi)(x) - (\mathcal{S}_{k-1} \psi)(x)]$ . If (5.3) holds for  $j \leq k$ , then using (3.12), this yields  $[\tilde{\Phi}_{k+1}(\mathcal{A}_G) \psi](x)$ , as asserted.

which is the expression in (2.14). Similarly,  $\Psi_{\lambda_n,1}(x_b, y_b) = \frac{1}{2\kappa_{\lambda_n}} \cdot \frac{\operatorname{Im} G_{\mathbf{T}_q}^{\lambda_n}(x_b, y_b)}{\operatorname{Im} G_{\mathbf{T}_q}^{\lambda_n}(o, o)}$ . In other words,  $\langle K_k \rangle_{n,L,\alpha} = \langle K_k \rangle_{\lambda_n}$ . Recalling (5.1) and (5.2), we thus get

$$\begin{aligned} \frac{1}{N(I)} \sum_{\lambda_n \in I} |\langle \psi_n, K_k \psi_n \rangle - \langle K_k \rangle_{\lambda_n}|^2 &\leq \operatorname{Var}^I(J_K^\lambda - \langle J_K^\lambda \rangle S_k) \\ &\quad + \operatorname{Var}^I(M_K^\lambda - \langle M_K^\lambda \rangle S_{k-1}) + \operatorname{Var}^I(P_K^\lambda - \langle P_K^\lambda \rangle S_{k-2}), \end{aligned}$$

where  $J_K^{\lambda_n} := J_{K,n}$ ,  $M_K^{\lambda_n} := M_{K,n}$  and  $P_K^{\lambda_n} := P_{K,n}$ .

From here, we argue as in the discussion following Remark 4.4. Namely, we use non-backtracking quantum variances and control them as before using the operators  $\mathcal{R}_{n,r}^\lambda$ . The only difference is that the variance of the remainders  $\operatorname{Var}^I(\tilde{\mathcal{S}}_T[J - \langle J \rangle S_k])$  should now be controlled for general  $k$  in Lemma 4.5, not just  $k = 0, 1$ . Using the same ideas of the proof of this lemma, we arrive at a bound of the form

$$\operatorname{Var}^I(K^\lambda) \lesssim \frac{c}{\pi N(I)} \int_{w(I)} \chi(m) \sum_{x_0 \in V} \sum_{x_k, y_k \in S(x_0, k)} J^m(x_0; x_k) \overline{J^m(x_0; y_k)} \operatorname{Im} G^m(\tilde{x}_k, \tilde{y}_k) dm,$$

where  $S(x_0, k) = \{d(y, x_0) = k\}$ . We may just bound this using  $|\operatorname{Im} G^m(\tilde{x}_k, \tilde{y}_k)| \leq c_k$ , so the sum becomes  $\sum_{x_0} |\sum_{x_k \in S(x_0, k)} |J^m(x_0; x_k)||^2 \leq (q+1)q^{k-1} \sum_{(x_0; x_k) \in \mathcal{B}_k} |J^m(x_0; x_k)|^2$ . Of course, such a crude bound would be problematic if we needed to take  $k \rightarrow \infty$  as in the control of the non-backtracking quantum variance (where  $k$  is replaced by  $n+k$ ,  $n \rightarrow \infty$ ). In this case the expression should be bounded with more caution (in Lemma 4.5 for example, we used the recursive formula (4.15) to rule out off-diagonal terms, thereby avoiding Cauchy-Schwarz). Here however,  $k$  is fixed and is the same as the one in  $K_k$ .

**Acknowledgements:** M.I. was funded by the LabEx IRMIA, and partially supported by the Agence Nationale de la Recherche project GeRaSic (ANR-13-BS01-0007-01).

M.S. was supported by a public grant as part of the *Investissement d'avenir* project, reference ANR-11-LABX-0056-LMH, LabEx LMH.

We thank Nalini Anantharaman for suggesting this problem to us.

## APPENDIX A. SOME GREEN'S FUNCTION IDENTITIES

### A.1. Relationship between resolvents on the quantum and combinatorial trees.

Lemma 3.2 can be in principle derived from the results of [16, Section 4]. Given that we use it in a crucial way, we present here a self-contained proof.

The first step is to calculate  $G_{\mathbf{T}_q}^\gamma(o_b, o_b)$  for some bond  $b \in \mathbf{T}_q$ . Indeed by (3.9) we have

$$G_{\mathbf{T}_q}^\gamma(o_b, o_b) = \frac{U_{\gamma;v}^-(o_b) V_{\gamma;o}^+(o_b)}{W_{v,o}^\gamma(o_b)}$$

for  $\gamma \notin \sigma(H_{\mathbf{T}_q})$ . If  $b \in \mathcal{B}^n b_o \cap \mathcal{B}^{*m} b_v$ , where  $b_o$  and  $b_v$  are the edges containing  $o, v$  respectively, we get  $U_{\gamma;v}^-(o_b) = \mu^-(\gamma)^m U_\gamma(0)$ ,  $V_{\gamma;o}^+(o_b) = \mu^-(\gamma)^n V_\gamma(0)$  and  $W_{v,o}^\gamma(o_b) = (\mu^-(\gamma))^{m+n} W^\gamma$ , where  $W^\gamma = V^\gamma(0) U_\gamma'(0) - V_\gamma'(0) U_\gamma(0) = s(\gamma)(1 - (q\mu^+(\gamma))^2)$ . Hence,

$$\begin{aligned} \text{(A.1)} \quad G_{\mathbf{T}_q}^\gamma(o_b, o_b) &= \frac{U_\gamma(0) V_\gamma(0)}{W^\gamma} = \frac{-s(\gamma) q \mu^+(\gamma) s(\gamma)}{s(\gamma)[1 - (q\mu^+(\gamma))^2]} \\ &= \frac{-s(\gamma)}{\mu^-(\gamma) - q\mu^+(\gamma)} = \frac{-s(\gamma)}{(q+1)\mu^-(\gamma) - w(\gamma)}, \end{aligned}$$

where we used (3.3). More generally, if  $(v_0, \dots, v_k)$  is a non-backtracking path of vertices in  $\mathbb{T}_q$ , then we can take  $o = v_0$ ,  $v = v_k$  and

$$G_{\mathbf{T}_q}^\gamma(v_0, v_k) = \frac{U_{\gamma;v}^-(v_0)V_{\gamma;o}^+(v_k)}{W_{v,o}^\gamma(v_0)} = (\mu^-(\gamma))^k \frac{U_\gamma(v_0)V_\gamma(v_0)}{W_{v,o}^\gamma(v_0)} = \frac{-s(\gamma)(\mu^-(\gamma))^k}{(q+1)\mu^-(\gamma) - w(\gamma)}$$

We thus showed that for any  $\gamma \in \mathbb{C}^+$ ,

$$(A.2) \quad G_{\mathbf{T}_q}^\gamma(v, w) = -s(\gamma) \frac{(\mu^-(\gamma))^{d(v,w)}}{(q+1)\mu^-(\gamma) - w(\gamma)} = -s(\gamma)G_{\mathcal{A}_{\mathbf{T}_q}}^{w(\gamma)}(v, w),$$

where the last equality is by the classic form of the Green's function of  $\mathcal{A}_{\mathbf{T}_q}$ , see e.g. [19, Chapter 1] or [7, Lemma 2.1]. Note that from (3.1),  $\mu^\pm(\gamma)$  coincides with the quantity  $\epsilon_\pm(w(\gamma))$  in [2, Section 6], both being the solutions of  $qx^2 - w(\gamma)x + 1 = 0$ .

If  $\gamma = \lambda + i\eta$  with  $\lambda \in \sigma_{\text{ac}}(H_{\mathbf{T}_q})$ , then taking  $\eta \downarrow 0$  in (A.2), we get  $G_{\mathbf{T}_q}^{\lambda+i0}(v, w) = -s(\lambda) \frac{(\mu^-(\lambda))^{d(v,w)}}{(q+1)\mu^-(\lambda) - w(\lambda)}$ . This coincides with  $-s(\lambda)G_{\mathcal{A}_{\mathbf{T}_q}}^{w(\lambda)+i0}(v, w)$ , proving the lemma<sup>4</sup>.

**A.2. Limiting measures and Green's functions.** If we re-write  $U_\gamma(x)$  and  $V_\gamma(x)$  in the basis  $\{S_\gamma(x), S_\gamma(L-x)\}$ , we find as in (4.2),

$$\begin{aligned} V_\gamma(x) &= q\mu^+(\gamma)S_\gamma(L-x) + S_\gamma(x), \\ U_\gamma(x) &= -(S_\gamma(L-x) + q\mu^+(\gamma)S_\gamma(x)), \end{aligned}$$

If  $x, y$  belong to the same bond  $b = (o, v)$ , and supposing that  $y \in \mathbf{T}_x^+$ , we have by (3.9),

$$\begin{aligned} G_{\mathbf{T}_q}^\gamma(x, y) &= \frac{\mu^-(\gamma)}{\mu^-(\gamma)} \cdot \frac{U_\gamma(x)V_\gamma(y)}{W^\gamma} \\ &= -\frac{(\mu^-(\gamma)S_\gamma(L-x) + S_\gamma(x))(q\mu^+(\gamma)S_\gamma(L-y) + S_\gamma(y))}{s(\gamma)[(q+1)\mu^-(\gamma) - w(\gamma)]}, \end{aligned}$$

using (3.3). Taking  $\gamma = \lambda + i0$  with  $\lambda \in \sigma_{\text{ac}}(H_{\mathbf{T}_q})$  and recalling that  $G_{\mathbf{T}_q}^\lambda(o, o) = \frac{-s(\lambda)}{(q+1)\mu^-(\lambda) - w(\lambda)}$  by (A.1), we get

$$\begin{aligned} G_{\mathbf{T}_q}^\lambda(x, y) &= \frac{G_{\mathbf{T}_q}^\lambda(o, o)}{s^2(\lambda)} \left( S_\lambda(L-x)S_\lambda(L-y) + S_\lambda(x)S_\lambda(y) \right. \\ &\quad \left. + \mu^-(\lambda)S_\lambda(L-x)S_\lambda(y) + q\mu^+(\lambda)S_\lambda(x)S_\lambda(L-y) \right). \end{aligned}$$

Hence,

$$\begin{aligned} \text{Im } G_{\mathbf{T}_q}^\lambda(x, y) &= \frac{\text{Im } G_{\mathbf{T}_q}^\lambda(o, o)}{s^2(\lambda)} \left( S_\lambda(L-x)S_\lambda(L-y) + S_\lambda(x)S_\lambda(y) \right. \\ &\quad \left. + \frac{w(\lambda)}{2q}S_\lambda(L-x)S_\lambda(y) + \frac{w(\lambda)}{2}S_\lambda(x)S_\lambda(L-y) \right) \\ &\quad + \frac{\text{Re } G_{\mathbf{T}_q}^\lambda(o, o)}{s^2(\lambda)} \left( \text{Im } \mu^-(\lambda)S_\lambda(L-x)S_\lambda(y) + \text{Im } q\mu^+(\lambda)S_\lambda(x)S_\lambda(L-y) \right). \end{aligned}$$

---

<sup>4</sup>Note that  $-s(\lambda) \frac{(\mu^-(\lambda))^{d(v,w)}}{(q+1)\mu^-(\lambda) - w(\lambda)}$  is well-defined for  $\lambda \in \sigma_{\text{ac}}(H_{\mathbf{T}_q})$ . In fact, the denominator has a non-zero imaginary part in this case, so in particular it never vanishes. In general, it is easy to see that this fraction can have a pole only if  $w(\lambda) = \pm(q+1)$ . If  $\alpha = 0$ , this only occurs for  $\lambda \in \sigma_2$ . Hence,  $G_{\mathbf{T}_q}^{\lambda+i0}(v, w)$  exists more generally for  $\lambda \in \mathbb{R} \setminus \sigma_2$  in this case and equals  $-s(\lambda)G_{\mathcal{A}_{\mathbf{T}_q}}^{w(\lambda)+i0}(v, w)$ .

Now  $\operatorname{Im} G_{\mathbf{T}_q}^\lambda(o, o) = \frac{(q+1)\operatorname{Im}\mu^-(\lambda)s(\lambda)}{|(q+1)\mu^-(\lambda)-w(\lambda)|^2}$  and  $\operatorname{Re} G_{\mathbf{T}_q}^\lambda(o, o) = \frac{q-1}{2q} \frac{w(\lambda)s(\lambda)}{|(q+1)\mu^-(\lambda)-w(\lambda)|^2}$ . Hence,  $\operatorname{Re} G_{\mathbf{T}_q}^\lambda(o, o) = \frac{q-1}{2q} \frac{w(\lambda)}{(q+1)\operatorname{Im}\mu^-(\lambda)} \operatorname{Im} G_{\mathbf{T}_q}^\lambda(o, o)$ . We thus get

$$\begin{aligned}
 \text{(A.3)} \quad \operatorname{Im} G_{\mathbf{T}_q}^\lambda(x, y) &= \frac{\operatorname{Im} G_{\mathbf{T}_q}^\lambda(o, o)}{s(\lambda)^2} \left( S_\lambda(L-x)S_\lambda(L-y) + S_\lambda(x)S_\lambda(y) \right. \\
 &\quad \left. + \frac{w(\lambda)}{2q} S_\lambda(L-x)S_\lambda(y) + \frac{w(\lambda)}{2} S_\lambda(x)S_\lambda(L-y) \right) \\
 &\quad + \frac{(q-1)w(\lambda)}{2q(q+1)} S_\lambda(L-x)S_\lambda(y) - \frac{q-1}{2(q+1)} w(\lambda) S_\lambda(x)S_\lambda(L-y) \\
 &= \frac{\operatorname{Im} G_{\mathbf{T}_q}^\lambda(o, o)}{s(\lambda)^2} \left( S_\lambda(L-x)S_\lambda(L-y) + S_\lambda(x)S_\lambda(y) \right. \\
 &\quad \left. + \frac{w(\lambda)}{q+1} [S_\lambda(L-x)S_\lambda(y) + S_\lambda(x)S_\lambda(L-y)] \right),
 \end{aligned}$$

which recovers  $\Psi_{\lambda_n, 1}(x, y) = \frac{1}{2\kappa_{\lambda_n}} \cdot \frac{\operatorname{Im} G_{\mathbf{T}_q}^{\lambda_n}(x, y)}{\operatorname{Im} G_{\mathbf{T}_q}^{\lambda_n}(o, o)}$  from (5.5). Because the final expression in (A.3) is symmetric in  $x$  and  $y$ , it holds for  $y \in \mathbf{T}_x^-$  too.

By a similar method we may also recover  $\Psi_{\lambda_n, k}(x_{b_1}, y_{b_k})$  with the help of identity (3.11), thus obtaining an alternative proof of (5.6).

If we specialize to  $y = x$  in (A.3), we get

$$\text{(A.4)} \quad \frac{\operatorname{Im} G_{\mathbf{T}_q}^{\lambda+i0}(x, x)}{\operatorname{Im} G_{\mathbf{T}_q}^{\lambda+i0}(o, o)} = \frac{1}{s(\lambda)^2} \left( S_\lambda^2(L-x) + S_\lambda^2(x) + \frac{2w(\lambda)}{q+1} S_\lambda(L-x)S_\lambda(x) \right).$$

If we consider the special case where  $\alpha = 0$ ,  $U \equiv 0$  we find the interesting property that  $\operatorname{Im} G_{\mathbf{T}_q}^{\lambda+i0}(x, x)$  is constant as a function of  $x$ . Indeed we have then  $S_\lambda(x) = \frac{\sin(\sqrt{\lambda}x)}{\sqrt{\lambda}}$  and  $w(\lambda) = (q+1)\cos(\sqrt{\lambda}L)$ , so

$$\begin{aligned}
 \frac{\operatorname{Im} G_{\mathbf{T}_q}^{\lambda+i0}(x, x)}{\operatorname{Im} G_{\mathbf{T}_q}^{\lambda+i0}(o, o)} &= \frac{1}{\sin^2(\sqrt{\lambda}L)} \left( (\sin(\sqrt{\lambda}L)\cos(\sqrt{\lambda}x) - \sin(\sqrt{\lambda}x)\cos(\sqrt{\lambda}L))^2 + \sin^2(\sqrt{\lambda}x) \right. \\
 &\quad \left. + 2\cos(\sqrt{\lambda}L) (\sin(\sqrt{\lambda}L)\cos(\sqrt{\lambda}x) - \sin(\sqrt{\lambda}x)\cos(\sqrt{\lambda}L)) \sin(\sqrt{\lambda}x) \right) \\
 &= \frac{1}{\sin^2(\sqrt{\lambda}L)} \left( \sin^2(\sqrt{\lambda}L)\cos^2(\sqrt{\lambda}x) - \sin^2(\sqrt{\lambda}x)\cos^2(\sqrt{\lambda}L) + \sin^2(\sqrt{\lambda}x) \right) \\
 &= 1.
 \end{aligned}$$

On the other hand, by (2.11), (A.4), we have  $\kappa_{\lambda_n} = \frac{q+1}{2} \int_0^L \frac{\operatorname{Im} G_{\mathbf{T}_q}^{\lambda_n}(x, y)}{\operatorname{Im} G_{\mathbf{T}_q}^{\lambda_n}(o, o)} dx$ . So  $\kappa_{\lambda_n} = \frac{q+1}{2}L$  in this case. Hence, we deduce that the limiting measure in the special case  $\alpha = 0$ ,  $U \equiv 0$  is the uniform measure :

$$\text{(A.5)} \quad \Psi_{\lambda_n}(x) = \frac{2}{L(q+1)}.$$

## REFERENCES

- [1] M. Aizenman, R. Sims and S. Warzel, *Absolutely continuous spectra of quantum tree graphs with weak disorder*, Comm. Math. Phys. **264** (2006) 371–389.
- [2] N. Anantharaman, *Quantum ergodicity on regular graphs*, Comm. Math. Phys. **353** (2017) 633–690.
- [3] N. Anantharaman, E. Le Masson, *Quantum ergodicity on large regular graphs*, Duke Math. J. **164** (2015) 723–765.
- [4] N. Anantharaman, M. Sabri, *Quantum ergodicity on graphs : from spectral to spatial delocalization*, arXiv:1704.02766. Preprint 2017.



- [5] N. Anantharaman, M. Sabri, *Quantum ergodicity for the Anderson model on regular graphs*, J. Math. Phys. **58** (2017), 091901.
- [6] N. Anantharaman, M. Sabri, *Recent results of quantum ergodicity on graphs and further investigation*, arXiv:1711.07666. Preprint 2017.
- [7] N. Anantharaman, M. Sabri, *Poisson kernel expansions for Schrödinger operators on trees*, arXiv:1610.05907. To appear in J. Spectral Theory.
- [8] I. Benjamini, O. Schramm, *Recurrence of distributional limits of finite planar graphs*, Electron. J. Probab. **6** (2001) 13 pp.
- [9] G. Berkolaiko, J. P. Keating and B. Winn, *No quantum ergodicity for star graphs*, Comm. Math. Phys. **250** (2004) 259–285.
- [10] G. Berkolaiko, J. P. Keating and U. Smilansky, *Quantum ergodicity for graphs related to interval maps*, Comm. Math. Phys. **273** (2007) 137–159.
- [11] G. Berkolaiko and P. Kuchment, *Introduction to quantum graphs*, SURV 186, AMS 2013.
- [12] G. Berkolaiko, B. Winn, *Relationship between scattering matrices and spectrum of quantum graphs*, Trans. Amer. Math. Soc. **362** (2010) 6261–6277.
- [13] J. Bourgain, A. Gamburd *Uniform expansion bounds for Cayley graphs of  $SL_2(\mathbb{F}_p)$* , Ann. of Math. **167(2)** (2008) 625642.
- [14] M. Brammall and B. Winn, *Quantum ergodicity for quantum graphs without back-scattering*, Ann. Henri Poincaré **17** (2016) 1353–1382.
- [15] S. Brooks, E. Le Masson, E. Lindenstrauss, *Quantum ergodicity and averaging operators on the sphere*, Int. Math. Res. Not. **19** (2016) 6034–6064.
- [16] R. Carlson, *Hill's equation for a homogeneous tree*, Electron. J. Differen. Equations **23** (1997) 1–30.
- [17] Y. Colin de Verdière, *Ergodicité et fonctions propres du laplacien*, Comm. Math. Phys. **102** (1985) 497–502.
- [18] Y. Colin de Verdière, *Semi-classical measures on quantum graphs and the Gauß map of the determinant manifold*, Ann. Henri Poincaré **16** (2015) 347–364.
- [19] A. Figà-Talamanca, T. Steger, *Harmonic Analysis for Anisotropic Random Walks on Homogeneous Trees*, Memoirs of the AMS no. 531, 1994.
- [20] S. Gnuzmann, J. P. Keating, F. Pietot *Eigenfunction statistics on quantum graphs*, Ann. of Phys. **325** (2010) 2595–2640.
- [21] L. Kamení and R. Schubert *Entropy of eigenfunctions on quantum graphs*, arXiv:1405.5871. Preprint 2014.
- [22] H. Kesten, *Symmetric random walks on groups*, Trans. Amer. Math. Soc. **92** (1959) 336–354.
- [23] A. Klein, *Extended states in the Anderson model on the Bethe lattice*, Adv. Math. **133** (1998) 163–184.
- [24] T. Kottos, U. Smilansky, *Quantum chaos on graphs*, Phys. Rev. Lett. **79(24)** (1997) 4794–4797.
- [25] P. Kuchment, *Quantum graphs: II. Some spectral properties of quantum and combinatorial graphs*, J. Phys. A, **38** (2005) 4887–4900.
- [26] E. Le Masson, T. Sahlsten, *Quantum ergodicity and Benjamini-Schramm convergence of hyperbolic surfaces*, Duke Math. J. **166** (2017) 3425–3460.
- [27] A. Lubotzky, R. Phillips, P. Sarnak. *Ramanujan graphs* Combinatorica, **8(3)** (1988) 261–277.
- [28] B. D. McKay, *The expected eigenvalue distribution of a large regular graph*, Linear Algebra Appl. **40** (1981) 203–216.
- [29] K. Pankrashkin, *Spectra of Schrödinger operators on equilateral quantum graphs*, Lett. Math. Phys. **77** (2006) 139–154.
- [30] J. Pöschel and E. Trubowitz, *Inverse Spectral Theory*, Academic Press 1987.
- [31] A. I. Šnirel'man, *Ergodic properties of eigenfunctions*, Uspehi Mat. Nauk, **29** (1974) 181–182.
- [32] S. Zelditch, *Uniform distribution of eigenfunctions on compact hyperbolic surfaces*, Duke Math. J. **55** (1987) 919–941.

UNIVERSITÉ DE STRASBOURG, CNRS, IRMA UMR 7501, F-67000 STRASBOURG, FRANCE.  
*E-mail address:* `ingremeau@math.unistra.fr`

DEPARTMENT OF MATHEMATICS, FACULTY OF SCIENCE, CAIRO UNIVERSITY, CAIRO 12613, EGYPT.

LABORATOIRE DE MATHÉMATIQUES D'ORSAY, UNIV. PARIS-SUD, CNRS, UNIVERSITÉ PARIS-SACLAY, 91405 ORSAY, FRANCE.  
*E-mail address:* `mostafa.sabri@math.u-psud.fr`

DEPARTMENT OF MATHEMATICAL SCIENCES, LOUGHBOROUGH UNIVERSITY, LEICESTERSHIRE, LE11 3TU, UNITED KINGDOM.  
*E-mail address:* `b.winn@lboro.ac.uk`