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Very Regular Solutions for the Landau-Lifschitz Equation with Electric Current

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Abstract

We consider a model of ferromagnetic material subject to an electric current. We prove the local in time existence of very regular solutions for this model in the scale of H^k spaces. In particular, we describe in detail the compatibility conditions at the boundary for the initial data.

Keywords: ferromagnetic materials, compatibility conditions, Existence result

1 Introduction

Ferromagnetic materials are used in several applications: radar furtivity, electromagnets, storage of digital data. They are characterized by a spontaneous magnetization represented by the magnetic moment m defined on $[0, T] \times \Omega$, where Ω is the ferromagnetic domain in which the material is confined. At low temperature (under the Curie temperature), the material is said to be saturated, that is the norm of the magnetic moment is constant, so after renormalization, we have:

$$\forall t \in [0, T], \quad \forall x \in \Omega, \quad |m(t, x)| = 1. \quad (1)$$

The magnetic moment links the magnetic field H and the magnetic induction B by the constitutive relation:

$$B = H + \overline{m},$$

where \overline{m} is the extension of m by zero outside Ω .

The most promising application of the ferromagnetic materials concerns the nano-electronics, and in particular the storage of digital informations with quick access. The ferromagnetic devices used for these applications are either nano wires, thin plates or really 3d components. For nano-electronic applications, one goal is to store informations by magnetic domains representing the bits, and transmit it very rapidly along the magnetic structures via domain walls motion. A standard way to induce this motion is by applying a magnetic field on the sample. The effects of the applied magnetic field are measured by the Zeeman energy (see [5, 11, 15]). The problem of this solution is that the information is modified since the domains can be removed by the application of an applied field. As explained in [13, 14], the good way to transport the information (to a reading head for example) without change is obtained by using of an applied electric current. This increases the access velocity to the memory without mechanical wear of the components, and preserving the data.

In this paper, we study the following Landau-Lifschitz model describing a ferromagnetic material submitted to an applied current:

$$\left\{ \begin{array}{ll} \frac{\partial m}{\partial t} = -m \times H_e(m) - m \times (m \times H_e(m)) + (v \cdot \nabla)m + m \times (v \cdot \nabla)m & \text{in } [0, T] \times \Omega, \\ H_e(m) = \Delta m + H(m), & \\ \frac{\partial m}{\partial \nu} = 0 & \text{on } [0, T] \times \partial\Omega, \\ m(0, x) = m_0(x) & \text{in } \Omega, \end{array} \right. \quad (2)$$

where:

- Ω is an open bounded set with smooth boundary and ν is the outside unit normal on $\partial\Omega$,
- H_e is the effective magnetic field including the demagnetizing field and the exchange field,
- v is the current speed. For physical reasons, the relevant boundary condition is that v satisfies $v \cdot \nu = 0$ on the boundary (the electric current remains in the domain), but we do not use this condition in our analysis.
- Δm is the exchange field,
- $H(m)$ is the demagnetizing field. It is deduced from m by the static Maxwell equations and the law of Faraday:

$$\left\{ \begin{array}{l} \text{curl } H(m) = 0 \text{ in } \mathbb{R}^3, \\ \text{div } (H(m) + \bar{m}) = 0 \text{ in } \mathbb{R}^3, \end{array} \right.$$

where \bar{m} is the extension of m by zero outside Ω .

Remark 1. *The modelisation of the applied current by the transport term $(v \cdot \nabla)m + m \times (v \cdot \nabla)m$ is explained in [13] (see also the references therein).*

We denote by $H^k(\Omega; S^2)$ the set:

$$H^k(\Omega; S^2) = \left\{ u : \Omega \rightarrow \mathbb{R}^3, \quad u \in H^k(\Omega) \text{ and } |u(x)| = 1 \text{ a.e.} \right\}.$$

In [1], [6] and [16], the existence of global weak solutions for the Landau-Lifschitz equation without electric current is obtained. With the same method as [1] or [6], G. Bonithon proves in [3] the global existence of weak solutions for (2). In the case of the Landau-Lifschitz equations without electric current when the effective field is reduced to Δm , the weak solutions are non unique (see [1]). This non uniqueness is expected in more general cases but is totally open. In addition, the Landau-Lifschitz equation does not have regularizing effects. So the existence of strong solutions is not obvious.

Existence of strong solutions for the Landau-Lifschitz equation without applied current is obtained in [7] and [8]. The solutions constructed in these papers are in $L^\infty(0, T; H^2(\Omega)) \cap L^2(0, T; H^3(\Omega))$ and are local in time.

In the present paper, following the method of [7], we construct local in time strong solutions for (2) in $L^\infty(0, T^*; H^2(\Omega)) \cap L^2(0, T^*; H^3(\Omega))$. In addition, for more regular initial data, if compatibility conditions are satisfied, we prove the existence of strong solutions with high regularity on the same maximal existence interval $[0, T^*[$. In particular, we prove that an eventual blow up for the regular solution occurs for the H^2 norm. To our knowledge, the construction of very regular solutions for the Landau-Lifschitz equation does not exist in the literature, even for models without applied current.

Our first result is the following:

Theorem 1.1. *Let $m_0 \in H^2(\Omega; S^2)$ satisfying the compatibility condition:*

$$\frac{\partial m_0}{\partial \nu} = 0 \text{ on } \partial\Omega. \quad (3)$$

Let $v \in C^0(\mathbb{R}^+; L^\infty(\Omega) \cap W^{1,3}(\Omega))$. There exist a time $T^ > 0$ depending only on the size of the initial data and a unique m such that for all $T < T^*$,*

- $m \in C^0(0, T; H^2(\Omega)) \cap L^2(0, T; H^3(\Omega))$.
- $|m(t, x)| = 1 \quad \text{in } [0, T] \times \Omega$.
- m satisfies (2).

Without additional compatibility condition, we can construct solutions in $H^3(\Omega)$:

Theorem 1.2. *Let $m_0 \in H^3(\Omega, S^2)$ such that $\frac{\partial m_0}{\partial \nu} = 0$ on $\partial\Omega$. Let $v \in C^0(\mathbb{R}^+; L^\infty(\Omega) \cap W^{1,3}(\Omega)) \cap C^1(\mathbb{R}^+; L^2(\Omega))$. Let $T^* > 0$ and $m \in L^\infty(0, T; H^2(\Omega)) \cap L^2(0, T; H^3(\Omega))$ for $T < T^*$ given by Theorem 1.1. Then, for all $T < T^*$:*

$$m \in C^0(0, T; H^3(\Omega)) \cap L^2(0, T; H^4(\Omega)).$$

The key point for the proof of Theorems 1.1 and 1.2 is that the Landau-Lifschitz equation (2) is equivalent, for solutions satisfying the saturation constraint (1), to the following problem:

$$\begin{cases} \frac{\partial m}{\partial t} = \Delta m + |\nabla m|^2 m - m \times \Delta m - m \times H(m) - m \times (m \times H(m)) \\ \quad + (v \cdot \nabla)m + m \times (v \cdot \nabla)m & \text{in } \mathbb{R}_t^+ \times \Omega, \\ \frac{\partial m}{\partial \nu} = 0 & \text{on } \mathbb{R}_t^+ \times \partial\Omega, \\ m(0, x) = m_0(x) & \text{in } \Omega. \end{cases} \quad (4)$$

In order to obtain more regularity for the solutions of (4), we must assume compatibility conditions on the initial data. Since (2) is equivalent to (4), and since our strategy consists in working with this last equation, we write the compatibility conditions derived from (4). First, we define formally the initial value of $\partial_t^k m$ at $t = 0$.

We set $V_0 = m_0$. We denote $v_i(x) := \partial_t^i v(0, x)$.

Using equation (4), if m is regular enough, we obtain the value of $\partial_t m|_{t=0}$ by taking (4) at $t = 0$ and by writing that $m|_{t=0} = m_0$. So, we define V_1 by:

$$\begin{aligned} V_1 &= \Delta m_0 + |\nabla m_0|^2 m_0 - m_0 \times (\Delta m_0 + H(m_0)) - m_0 \times (m_0 \times (H(m_0))) \\ &+ (v_0 \cdot \nabla)m_0 + m_0 \times (v_0 \cdot \nabla)m_0. \end{aligned}$$

We remark that, if $m_0 \in H^k(\Omega)$ and if v is sufficiently regular, then $V_1 \in H^{k-2}(\Omega)$ for $k \geq 2$.

In order to write the compatibility conditions, we recall the following formulas for the derivation of a product:

$$\partial_t^K(ab) = \sum_{i=0}^K C_K^i \partial_t^i a \partial_t^{K-i} b, \quad \text{with } C_K^i = \frac{k!}{(k-i)! i!}, \quad (5)$$

and

$$\partial_t^K(abc) = \sum_{\alpha \in A_K} C^\alpha \partial_t^{\alpha_1} a \partial_t^{\alpha_2} b \partial_t^{\alpha_3} c, \quad (6)$$

with:

- $A_K = \left\{ \alpha = (\alpha_1, \alpha_2, \alpha_3) \in \{0, \dots, K\}^3, \quad \alpha_1 + \alpha_2 + \alpha_3 = K \right\}$,

- $C^\alpha = C_{\alpha_1+\alpha_2+\alpha_3}^{\alpha_1+\alpha_2} \cdot C_{\alpha_1+\alpha_2}^{\alpha_1}$.

By differentiating (4) k times with respect to t , and by taking the obtained result at $t = 0$, we define recursively $V_{k+1} : \Omega \rightarrow \mathbb{R}^3$ by:

$$\begin{aligned} V_{k+1} = & \Delta V_k + \sum_{\alpha \in A_k} C^\alpha (\nabla V_{\alpha_1} \cdot \nabla V_{\alpha_2}) V_{\alpha_3} - \sum_{j=0}^k C_k^j V_j \times (\Delta V_{k-j} + H(V_{k-j})) \\ & - \sum_{\alpha \in A_k} C^\alpha \left(V_{\alpha_1} \times (V_{\alpha_2} \times H(V_{\alpha_3})) - V_{\alpha_1} \times (v_{\alpha_2} \cdot \nabla) V_{\alpha_3} \right) + \sum_{j=0}^k C_k^j (v_j \cdot \nabla) V_{k-j}. \end{aligned} \quad (7)$$

We remark that if $m_0 \in H^{2k+2+p}(\Omega)$, then $V_{k+1} \in H^p(\Omega)$.

Definition 1.3. Let $k \in \mathbb{N}$, let $m_0 \in H^{2k+2}(\Omega; S^2)$. We say that m_0 satisfies the compatibility condition at order k if :

$$\forall i \in \{0, 1, \dots, k\}, \quad \frac{\partial V_i}{\partial \nu} = 0 \text{ on } \partial\Omega.$$

Theorem 1.4. Let $k \geq 4$. Let $m_0 \in H^k(\Omega; S^2)$ and let m and T^* given in Theorem 1.1. We assume that m_0 satisfies the compatibility condition at order $\lfloor \frac{k}{2} \rfloor - 1$. In addition, we assume that

$$\forall j \leq k - \lfloor \frac{k}{2} \rfloor - 1, \quad \partial_t^j v \in C^0(0, T; H^{k-2j}(\Omega)) \cap L^2(0, T; H^{k-2j+1}(\Omega)).$$

Then, for all $T < T^*$,

$$m \in L^\infty(0, T; H^k(\Omega)) \cap L^2(0, T; H^{k+1}(\Omega)),$$

and

$$\forall j \in 1, \dots, \lfloor \frac{k}{2} \rfloor - 1, \quad \frac{\partial^j m}{\partial t^j} \in L^\infty(0, T; H^{k-2j}(\Omega)) \cap L^2(0, T; H^{k-2j+1}(\Omega)).$$

The present paper is organized as follows.

In Section 2, we recall useful lemmas. In particular, we study the regularity properties for the operator H .

Section 3 is devoted to the proof of Theorem 1.1. Basically, we follow the method of [7]. As we said before, the key point is that Equation (2) is equivalent to (4) for solutions satisfying the saturation constraint (1). In this new formulation, the dissipation due to the exchange field Δm appears completely though it only appears $m \times \Delta m$ in (2). We construct by the Galerkin method a solution of (4) in $L^\infty(0, T; H^2(\Omega)) \cap L^2(0, T; H^3(\Omega))$, and we prove that this solution satisfies in addition the saturation constraint, so that it is a solution for (2).

Construction of more regular solutions for (2) is totally new and entails several difficulties. In Section 4, in order to prove the H^3 regularity for the solution of (4), a direct energy estimate for the third order space derivatives is not possible because of the non local term $H(m)$ which does not satisfy the homogeneous Neumann boundary condition. Our method consists in differentiating the Galerkin approximation with respect to time and we prove by this way that $\partial_t m \in L^\infty(0, T; H^1(\Omega)) \cap L^2(0, T; H^2(\Omega))$. We recover the H^3 -regularity for m by a bootstrap argument using Equation (4). This operation is complicated by the non linear term $|\nabla m|^2 m$.

Section 5 is devoted to the proof of Theorem 1.4. As already said for the proof of the H^3 -regularity, variational estimates for high order space derivatives are not possible because of the non local term $H(m)$ which does not satisfy the homogeneous Neumann boundary condition. In addition, we are unable to perform H^2 estimates for the time derivative of m_n , solution for the Galerkin approximation of (4). Indeed, we need for that compatibility conditions for all n which are not satisfied because of the non local term $H(m_n)$.

Let us explain briefly our strategy in the simplest case: $k = 4$. We already know from Theorem 1.2 that the solution m is in $L^\infty(0, T; H^3(\Omega)) \cap L^2(0, T; H^4(\Omega))$. We derivate (4) with respect to time.

By a Galerkin process, using compatibility conditions on $\partial_t m(t=0)$, we construct a solution for the obtained problem in $L^\infty(0, T; H^2(\Omega)) \cap L^2(0, T; H^3(\Omega))$. Now, $\partial_t m$ is a solution of this problem and a uniqueness argument ensures then that $\partial_t m \in L^\infty(0, T; H^2(\Omega)) \cap L^2(0, T; H^3(\Omega))$. Finally, we improve the regularity of m by elliptic regularity theorems writing Δm in function of $\partial_t m$. In the general case, we proceed by induction with, roughly speaking, the same method.

2 Preliminary results

2.1 Equivalent norms

We use the following notations:

$$\|u\|_{2,\Omega} = \left(\|u\|_{L^2}^2 + \|\Delta u\|_{L^2}^2 \right)^{\frac{1}{2}},$$

and

$$\|\nabla u\|_{1,\Omega} = \left(\|\nabla u\|_{L^2}^2 + \|\Delta u\|_{L^2}^2 \right)^{\frac{1}{2}}.$$

We recall the following result established in [7]:

Lemma 2.1. *Let Ω be a bounded regular open set. There exist c_1 and c_2 such that for all $u \in H^2(\Omega)$ such that $\frac{\partial u}{\partial \nu} = 0$ on $\partial\Omega$,*

$$\begin{aligned} c_1 \|u\|_{2,\Omega} &\leq \|u\|_{H^2(\Omega)} \leq c_2 \|u\|_{2,\Omega}, \\ c_1 \|\nabla u\|_{1,\Omega} &\leq \|\nabla u\|_{H^1(\Omega)} \leq c_2 \|\nabla u\|_{1,\Omega}, \end{aligned}$$

and for all $u \in H^3(\Omega)$ such that $\frac{\partial u}{\partial \nu} = 0$ on $\partial\Omega$,

$$c_1 \left(\|\nabla u\|_{1,\Omega}^2 + \|\nabla \Delta u\|_{L^2}^2 \right)^{\frac{1}{2}} \leq \|\nabla u\|_{H^2(\Omega)} \leq c_2 \left(\|\nabla u\|_{1,\Omega}^2 + \|\nabla \Delta u\|_{L^2}^2 \right)^{\frac{1}{2}}.$$

From Lemma 2.1 and using the standard interpolation inequality, we rewrite Sobolev and Gagliardo-Nirenberg inequalities on the form

Lemma 2.2. *Let Ω be a regular bounded domain of \mathbb{R}^3 . There exists a constant C such that for all $u \in H^2(\Omega)$ such that $\frac{\partial u}{\partial \nu} = 0$ on $\partial\Omega$,*

$$\begin{aligned} \|u\|_{L^\infty} &\leq C \|u\|_{2,\Omega}, \\ \|\nabla u\|_{L^6} &\leq C \|u\|_{2,\Omega}, \\ \|\nabla u\|_{L^4}^2 &\leq C \|u\|_{L^\infty} \|u\|_{2,\Omega}, \end{aligned}$$

and for all $u \in H^3(\Omega)$ such that $\frac{\partial u}{\partial \nu} = 0$ on $\partial\Omega$,

$$\|D^2 u\|_{L^3} \leq C \left(\|u\|_{2,\Omega} + \|u\|_{2,\Omega}^{\frac{1}{2}} \|\nabla \Delta u\|_{L^2}^{\frac{1}{2}} \right).$$

The following lemma will be useful to estimate products of functions.

Lemma 2.3. *If $u \in H^1(\Omega)$ and $v \in H^2(\Omega)$, then $uv \in H^1(\Omega)$ and*

$$\|uv\|_{H^1(\Omega)} \leq C \|u\|_{H^1(\Omega)} \|v\|_{H^2(\Omega)}$$

Proof. Since Ω is a smooth bounded domain of \mathbb{R}^3 , by Sobolev embedding, $H^2(\Omega) \subset L^\infty(\Omega)$ so that $uv \in L^2(\Omega)$.

In addition, $\nabla(uv) = (\nabla u)v + u\nabla v$. By the same argument, $(\nabla u)v \in L^2(\Omega)$. Since $H^1(\Omega) \subset L^4(\Omega)$, then $u\nabla v \in L^2(\Omega)$. \square

2.2 Comparison lemma

We recall without proof the following standard comparison lemma:

Lemma 2.4. *Let $f : \mathbb{R}^+ \times \mathbb{R} \rightarrow \mathbb{R}$, C^0 and locally lipschitz with respect to its second variable. Let $z : [0, T^*[\rightarrow \mathbb{R}$ be the maximal solution of the Cauchy problem:*

$$\begin{cases} z' = f(t, z), \\ z(0) = z_0. \end{cases}$$

Let $y : \mathbb{R}^+ \rightarrow \mathbb{R}$, C^1 such that

$$\begin{cases} \forall t \geq 0, y'(t) \leq f(t, y(t)), \\ y(0) \leq z_0. \end{cases}$$

Then

$$\forall t \in [0, T^*[, \quad y(t) \leq z(t).$$

2.3 Galerkin basis

Let us denote by V_n the finite space spanned on the n first eigenfunctions of the operator $A = -\Delta + I$ with $D(A) = \left\{ u \in H^2(\Omega) \text{ such that } \frac{\partial u}{\partial \nu} = 0 \text{ on } \partial\Omega \right\}$, and P_n , the orthogonal projection from $L^2(\Omega)$ onto V_n .

Proposition 1. *There exists a constant C such that for all n , the orthogonal projection P_n satisfies the following properties:*

- For all $u \in H^1(\Omega)$, $\|P_n(u)\|_{H^1(\Omega)} \leq \|u\|_{H^1(\Omega)}$,
- For all $u \in H^2(\Omega)$ such that $\frac{\partial u}{\partial \nu} = 0$, $\|P_n(u)\|_{H^2(\Omega)} \leq C\|u\|_{H^2(\Omega)}$,
- For all $u \in H^3(\Omega)$ such that $\frac{\partial u}{\partial \nu} = 0$, $\|P_n(u)\|_{H^3(\Omega)} \leq C\|u\|_{H^3(\Omega)}$.

Proof. We write u on the form $u = P_n(u) + Q_n(u)$, where $Q_n(u)$ belongs to V_n^\perp . Since $\|u\|_{L^2}^2 = \|P_n(u)\|_{L^2}^2 + \|Q_n(u)\|_{L^2}^2$, we obtain

$$\|P_n(u)\|_{L^2} \leq \|u\|_{L^2}.$$

Using integration by parts and the fact that $\Delta P_n(u)$ belongs to V_n so that $\frac{\partial P_n(u)}{\partial \nu} = 0$ on $\partial\Omega$, we obtain

$$\begin{aligned} \|\nabla P_n(u)\|_{L^2}^2 &= - \int_{\Omega} \Delta P_n(u) \cdot P_n(u) \, dx \\ &= - \int_{\Omega} \Delta P_n(u) \cdot u \, dx \quad (\text{since } \Delta P_n(u) \text{ and } Q_n(u) \text{ are orthogonal}) \\ &= \int_{\Omega} \nabla P_n(u) \cdot \nabla u \, dx \quad (\text{since } P_n(u) \in V_n) \\ &\leq \|\nabla P_n(u)\|_{L^2} \|\nabla u\|_{L^2}. \end{aligned}$$

So

$$\|\nabla P_n(u)\|_{L^2} \leq \|\nabla u\|_{L^2}$$

In the same way, for the H^2 and the H^3 estimates, we remark that $\frac{\partial \Delta P_n(u)}{\partial \nu} = \frac{\partial \Delta^2 P_n(u)}{\partial \nu} = 0$ on the boundary. For the H^2 estimate, we have:

$$\begin{aligned}
\|\Delta P_n(u)\|_{L^2}^2 &= - \int_{\Omega} \nabla \Delta P_n(u) \cdot \nabla P_n(u) \, dx \\
&= \int_{\Omega} \Delta^2 P_n(u) \cdot P_n(u) \, dx \\
&= \int_{\Omega} \Delta^2 P_n(u) \cdot u \, dx && \text{(since } \Delta^2 P_n(u) \in V_n) \\
&= - \int_{\Omega} \nabla \Delta P_n(u) \cdot \nabla u \, dx \\
&= \int_{\Omega} \Delta P_n(u) \cdot \Delta u \, dx && \text{(since } \frac{\partial u}{\partial \nu} = 0 \text{ on } \partial\Omega) \\
&\leq \|\Delta P_n(u)\|_{L^2} \|\Delta u\|_{L^2}.
\end{aligned}$$

So, $\|\Delta P_n(u)\|_{L^2} \leq \|\Delta u\|_{L^2}$ and therefore:

$$\|P_n(u)\|_{H^2} \leq C \|u\|_{H^2}.$$

For the H^3 estimate,

$$\begin{aligned}
\|\nabla \Delta P_n(u)\|_{L^2}^2 &= - \int_{\Omega} \Delta^2 P_n(u) \cdot \Delta P_n(u) \, dx \\
&= \int_{\Omega} \nabla \Delta^2 P_n(u) \cdot \nabla P_n(u) \, dx \\
&= - \int_{\Omega} \Delta^3 P_n(u) \cdot P_n(u) \, dx \\
&= - \int_{\Omega} \Delta^3 P_n(u) \cdot u \, dx && \text{(since } \Delta^3 P_n(u) \in V_n) \\
&= \int_{\Omega} \nabla \Delta^2 P_n(u) \cdot \nabla u \, dx \\
&= - \int_{\Omega} \Delta^2 P_n(u) \cdot \Delta u \, dx && \text{(using that } \frac{\partial u}{\partial \nu} = 0 \text{ on } \partial\Omega), \\
&= \int_{\Omega} \nabla \Delta P_n(u) \cdot \nabla \Delta u \, dx \\
&\leq \|\nabla \Delta P_n(u)\|_{L^2} \|\nabla \Delta u\|_{L^2}.
\end{aligned}$$

Therefore, using Lemma 2.1,

$$\|P_n(u)\|_{H^3} \leq \|u\|_{H^3}.$$

□

2.4 Demagnetizing field

The operator H takes its values in the space $L^2(\mathbb{R}^3)$. We can observe that $u \longrightarrow -H(u)$ is the orthogonal projection of \bar{u} on the vector fields of gradients in $L^2(\mathbb{R}^3)$. Let us consider the restriction

of H to Ω . Classically, we have

$$\|H(u)\|_{L^p(\Omega)} \leq C\|u\|_{L^p(\Omega)} \quad \text{for } 1 < p < +\infty.$$

In the proposition below, following Ladyshenskaya [10] page 196 we can derive the following regularity result which describe the continuity of the operator H on the spaces $W^{k,p}(\Omega)$ for $p \in]1, +\infty[$ and $k \in \mathbb{N}$.

Proposition 2. *Let $p \in]1, +\infty[$, for $k \in \mathbb{N}$, then if u belongs to $W^{k,p}(\Omega)$, the restriction of $H(u)$ to Ω belongs to $W^{k,p}(\Omega)$ and there exists a constant $C_{k,p}$ such that*

$$\|H(u)\|_{W^{k,p}(\Omega)} \leq C_{k,p}\|u\|_{W^{k,p}(\Omega)}.$$

Proof. See [8]. □

3 Proof of the H^2 regularity

In [7], Carbou and Fabrie prove the existence of local strong solutions in $L^\infty(0, T; H^2) \cap L^2(0, T; H^3)$ for the Landau-Lifschitz equation without electric current. Our proof for Theorem 1.1 is basically the same. For the convenience of the reader, we give the complete proof in the present paper, emphasizing the changes due to the electric current.

3.1 Equivalent system.

A dissipative term of the form $\|m \times \Delta m\|_{L^2}$ appears if we take the inner product in L^2 of (2) with Δm . This dissipation is not sufficient to obtain energy estimate in the space $H^2(\Omega)$.

We observe that, for m regular enough and $|m| = 1$ in Ω , then

$$m \times (m \times \Delta m) = (m \cdot \Delta m)m - |m|^2 \Delta m = -\Delta m - |\nabla m|^2 m.$$

So, if m is regular enough and satisfies the saturation constraint, then m is solution for the Landau-Lifschitz equation (2) if and only if m satisfies the following system:

$$\left\{ \begin{array}{l} \frac{\partial m}{\partial t} - \Delta m = \begin{array}{l} |\nabla m|^2 m - m \times \Delta m - m \times H(m) - m \times (m \times H(m)) \\ + (v \cdot \nabla)m + m \times (v \cdot \nabla)m \end{array} \quad \text{in } \mathbb{R}_t^+ \times \Omega, \\ \frac{\partial m}{\partial \nu} = 0 \quad \text{on } \mathbb{R}_t^+ \times \partial\Omega, \\ m(0, x) = m_0(x) \quad \text{in } \Omega. \end{array} \right. \quad (8)$$

This equation has the advantage to highlight the dissipative term and is more convenient to build regular solutions for (2).

3.2 Galerkin Approximation for the modified Landau-Lifschitz equation

We recall that we denote by V_n the finite space built on the n first eigenfunctions of the operator $A = -\Delta + I$ with $D(A) = \left\{ u \in H^2(\Omega) \text{ such that } \frac{\partial u}{\partial \nu} = 0 \text{ on } \partial\Omega \right\}$, and by P_n the orthogonal projection from $L^2(\Omega)$ on V_n . We aim to find a solution m_n taking its values in V_n for the following Galerkin approximation of (4):

$$\left\{ \begin{array}{l} \frac{\partial m_n}{\partial t} - \Delta m_n = \begin{array}{l} P_n (|\nabla m_n|^2 m_n - m_n \times \Delta m_n) \\ - P_n (m_n \times H(m_n) + m_n \times (m_n \times H(m_n))) \\ + P_n ((v \cdot \nabla)m_n + m_n \times (v \cdot \nabla)m_n), \end{array} \\ m_n(0) = P_n(m_0). \end{array} \right. \quad (9)$$

Cauchy-Lipschitz theorem ensures the existence of a unique solution of (9) defined on $[0, T_n[$.

3.2.1 L^2 estimate for (8)

Taking the inner product in $L^2(\Omega)$ of (9) by m_n , we obtain

$$\begin{aligned} \frac{1}{2} \frac{d}{dt} \|m_n\|_{L^2}^2 + \|\nabla m_n\|_{L^2}^2 &= \int_{\Omega} (|\nabla m_n|^2 |m_n|^2 + (v \cdot \nabla) m_n \cdot m_n) dx \\ &\leq \|\nabla m_n\|_{L^2}^2 \|m_n\|_{L^\infty}^2 + \|v\|_{L^\infty} \|\nabla m_n\|_{L^2} \|m_n\|_{L^2} \\ &\leq C \|m_n\|_{2, \Omega}^4 + C \|v\|_{L^\infty} \|m_n\|_{2, \Omega}^2. \end{aligned} \quad (10)$$

3.2.2 H^2 estimate for (8)

Now, we take the inner product in $L^2(\Omega)$ of (9) with $\Delta^2 m_n$ and we integrate by parts to get

$$\frac{1}{2} \frac{d}{dt} \|\Delta m_n\|_{L^2}^2 + \|\nabla \Delta m_n\|_{L^2}^2 \leq \sum_{i=1}^4 I_i,$$

where

$$I_1 = - \int_{\Omega} \nabla (|\nabla m_n(t)|^2 m_n(t)) \nabla \Delta m_n(t) dx,$$

$$I_2 = \int_{\Omega} \nabla (m_n(t) \times \Delta m_n(t)) \nabla \Delta m_n(t) dx,$$

$$I_3 = \int_{\Omega} \nabla \left(m_n(t) \times H(m_n(t)) + m_n(t) \times (m_n(t) \times H(m_n(t))) \right) \nabla \Delta m_n(t) dx,$$

$$I_4 = - \int_{\Omega} \nabla \left((v \cdot \nabla) m_n(t) + m_n(t) \times (v \cdot \nabla) m_n(t) \right) \nabla \Delta m_n(t) dx,$$

We bound separately each term.

The terms I_1 , I_2 and I_3 are estimated in [7]. For the convenience of the reader, we rewrite these estimates. Using Lemma 2.2, we have

- Estimate on I_1 :

$$\begin{aligned} |I_1| &\leq C \int_{\Omega} (|m_n| |\nabla m_n| |D^2 m_n| + |\nabla m_n|^3) |\nabla \Delta m_n| \\ &\leq C (\|m_n\|_{L^\infty} \|D^2 m_n\|_{L^3} \|\nabla m_n\|_{L^6} + \|\nabla m_n\|_{L^6}^3) \|\nabla \Delta m_n\|_{L^2} \\ &\leq C \|m_n\|_{2, \Omega}^{\frac{5}{2}} \|\nabla \Delta m_n\|_{L^2}^{\frac{3}{2}} + C \|m_n\|_{2, \Omega}^3 \|\nabla \Delta m_n\|_{L^2}. \end{aligned}$$

- Estimate on I_2 : we remark that $(m_n(t) \times \nabla \Delta m_n(t)) \nabla \Delta m_n(t) = 0$, so

$$\begin{aligned} |I_2| &\leq \|\Delta m_n\|_{L^3} \|\nabla m_n\|_{L^6} \|\nabla \Delta m_n\|_{L^2} \\ &\leq C \|m_n\|_{2, \Omega}^2 \|\nabla \Delta m_n\|_{L^2} + C \|m_n\|_{2, \Omega}^{\frac{3}{2}} \|\nabla \Delta m_n\|_{L^2}^{\frac{3}{2}}. \end{aligned}$$

- Estimate on I_3

$$I_3 = \int_{\Omega} (\nabla m_n(t) \times H(m_n(t)) \nabla \Delta m_n(t)) dx + \int_{\Omega} (m_n(t) \times \nabla H(m_n(t)) \nabla \Delta m_n(t)) dx.$$

$$\begin{aligned} |I_3| &\leq C \int_{\Omega} (1 + |m_n|) \left(|H(m_n)| |\nabla m_n| + |\nabla H(m_n)| |m_n| \right) |\nabla \Delta m_n| \\ &\leq C (1 + \|m_n\|_{L^\infty}) \left(\|H(m_n)\|_{L^3} \|\nabla m_n\|_{L^6} + \|\nabla H(m_n)\|_{L^6} \|m_n\|_{L^3} \right) \|\nabla \Delta m_n\|_{L^2} \\ &\leq C \|\nabla \Delta m_n\|_{L^2} (1 + \|m_n\|_{2,\Omega}^3). \end{aligned}$$

- Estimate on I_4

$$\begin{aligned} |I_4| &\leq C \int_{\Omega} \left((1 + |m_n|) (|\nabla m_n| |\nabla v| + |D^2 m_n| |v|) + |\nabla m_n|^2 |v| \right) |\nabla \Delta m_n| \\ &\leq C (1 + \|m_n\|_{L^\infty}) (\|\nabla m_n\|_{L^6} \|\nabla v\|_{L^3} + \|D^2 m_n\|_{L^2} \|v\|_{L^\infty}) \|\nabla \Delta m_n\|_{L^2} \\ &\quad + C \|\nabla m_n\|_{L^4}^2 \|v\|_{L^\infty} \|\nabla \Delta m_n\|_{L^2} \\ &\leq C (\|\nabla v\|_{L^3} + \|v\|_{L^\infty}) (1 + \|m_n\|_{2,\Omega}^2) \|\nabla \Delta m_n\|_{L^2}. \end{aligned}$$

By addition of all these estimates and using Young inequality, we obtain the following inequality

$$\frac{d}{dt} \|\Delta m_n\|_{L^2}^2 + \|\nabla \Delta m_n\|_{L^2}^2 \leq C \left(1 + \|\nabla v\|_{L^3}^2 + \|v\|_{L^\infty}^2 \right) (1 + \|m_n\|_{2,\Omega}^{10}). \quad (11)$$

3.2.3 Uniform estimate on the Galerkin approximation

Summing Inequalities (10) and (11), we obtain that there exists a constant C such that

$$\frac{d}{dt} \|m_n\|_{2,\Omega}^2 + \|\nabla m_n\|_{L^2}^2 + \|\nabla \Delta m_n\|_{L^2}^2 \leq C \left(1 + \|\nabla v\|_{L^3}^2 + \|v\|_{L^\infty}^2 \right) (1 + \|m_n\|_{2,\Omega}^{10}), \quad (12)$$

where C does not depend on n . We denote:

$$c(t) = C \left(1 + \|\nabla v(t)\|_{L^3}^2 + \|v(t)\|_{L^\infty}^2 \right).$$

From the assumptions about v , $c \in \mathcal{C}^0(\mathbb{R}^+)$.

In addition, $m_n(0) = P_n(m_0)$, and since m_0 satisfies the compatibility condition (3), by Proposition 1, we obtain that for all n :

$$\|m_n(0)\|_{2,\Omega} \leq C \|m_0\|_{2,\Omega}.$$

We set $y_n(t) = \|m_n\|_{2,\Omega}^2$. We have proven that for all n ,

$$\begin{cases} \frac{d}{dt} y_n(t) \leq c(t)(1 + y_n^5(t)), \\ y_n(0) \leq C \|m_0\|_{2,\Omega}^2. \end{cases}$$

We consider z the maximal solution of the o.d.e.:

$$\begin{cases} \frac{d}{dt} z(t) = c(t)(1 + z^5(t)), \\ z(0) = C \|m_0\|_{2,\Omega}^2, \end{cases}$$

and we denote by T^* the maximal existence time of z . By the comparison lemma, we have:

$$\forall t < T^*, \forall n, y_n(t) \leq z(t).$$

Therefore, for all $T < T^*$, there exists a constant $C(T)$ such that

$$\sup_{t \leq T} \|m_n(t)\|_{H^2(\Omega)}^2 \leq C(T), \quad (13)$$

and by integration of (12) on the interval $[0, T]$ we obtain

$$\int_0^T (\|\nabla m_n(t)\|_{L^2}^2 + \|\nabla \Delta m_n(t)\|_{L^2}^2) dt \leq C(T). \quad (14)$$

By equation (9) we conclude that

$$\sup_{t \leq T} \left\| \frac{\partial}{\partial t} m_n(t) \right\|_{L^2} \leq C(T),$$

and

$$\int_0^T \left\| \frac{\partial}{\partial t} \nabla m_n(\tau) \right\|_{L^2}^2 d\tau \leq C(T).$$

3.3 Limit when n tends to $+\infty$

These uniform estimates insure the existence of a subsequence (m_{n_k}) and a function m such that

- $m_{n_k} \rightharpoonup m$ in $L^2(0, T; H^3(\Omega))$ weak,
- $m_{n_k} \rightharpoonup m$ in $L^\infty(0, T; H^2(\Omega))$ weak *,
- $\frac{\partial m_{n_k}}{\partial t} \rightharpoonup \frac{\partial m}{\partial t}$ in $L^2(0, T; H^1(\Omega))$ weak.

Thanks to Aubin-Simon compactness lemma (see [2] and [12]), we can conclude that

$$m_{n_k} \rightarrow m \text{ in } L^\infty(0, T; H^2(\Omega)) \text{ strong,}$$

and

$$H(m_{n_k}) \rightarrow H(m) \text{ in } L^\infty(0, T; H^2(\Omega)) \text{ strong}$$

because H is a continuous map on Sobolev spaces $H^s(\Omega)$ for $s = 0, 1, 2$.

Taking the limit in (9), we obtain that m satisfies

$$\begin{cases} \frac{\partial m}{\partial t} - \Delta m = |\nabla m|^2 m - m \times \Delta m - m \times H(m) - m \times (m \times H(m)) \\ \quad + (v \cdot \nabla) m + m \times (v \cdot \nabla) m & \text{in } [0, T^*[\times \Omega, \\ \frac{\partial m}{\partial \nu} = 0 & \text{on } [0, T^*[\times \partial \Omega, \\ m(0, x) = m_0(x) & \text{in } \Omega. \end{cases}$$

In addition, we remark that from Theorem II.5.14 in [4], then

$$m \in C^0([0, T^*]; H^2(\Omega)).$$

3.4 Conservation of the ponctual norm

We prove now that the solution m constructed in the previous part satisfies the physical constraint $|m| = 1$, so that m satisfies Landau-Lifschitz equation. Contrary to [7], a transport term due to electric current appears in equation (4).

Using the scalar product in \mathbb{R}^3 of (4) by m , we get

$$\frac{1}{2} \frac{d}{dt} |m|^2 - m \cdot \Delta m - |\nabla m|^2 |m|^2 - m \cdot (v \cdot \nabla) m = 0 \quad \text{in } [0, T[\times \Omega.$$

For $d \leq 3$ and for all $u \in L^\infty(0, T; H^2(\Omega))$ we have

$$\Delta |u|^2 = 2u \cdot \Delta u + 2|\nabla u|^2,$$

and $m \cdot (v \cdot \nabla) m = \frac{1}{2} (v \cdot \nabla) |m|^2$, then $|m|^2$ satisfies the following equation

$$\frac{d}{dt} |m|^2 - \Delta |m|^2 - 2|\nabla m|^2 (|m|^2 - 1) - \frac{1}{2} v \cdot \nabla |m|^2 = 0.$$

Taking $a = |m|^2 - 1$, then a satisfies the system below

$$\begin{cases} \frac{\partial a}{\partial t} - \Delta a - 2|\nabla m|^2 a - \frac{1}{2} (v \cdot \nabla) a = 0 & \text{in } \Omega, \\ \frac{\partial a}{\partial \nu} = 0 & \text{on } \partial\Omega, \\ a(0) = |m_0|^2 - 1 = 0 & \text{in } \Omega. \end{cases} \quad (15)$$

Taking the inner product of this equation by a , we obtain that

$$\begin{aligned} \frac{1}{2} \frac{d}{dt} \|a\|_{L^2}^2 + \|\nabla a\|_{L^2}^2 &= 2 \int_{\Omega} |\nabla m|^2 a^2 + \frac{1}{2} \int_{\Omega} (v \cdot \nabla) a \cdot a \\ &\leq C \|\nabla m\|_{L^\infty}^2 \|a\|_{L^2}^2 + \frac{1}{2} \|v\|_{L^\infty} \|a\|_{L^2} \|\nabla a\|_{L^2}, \end{aligned}$$

and by absorbing $\|\nabla a\|_{L^2}$, we obtain that

$$\frac{d}{dt} \|a\|_{L^2}^2 + \|\nabla a\|_{L^2}^2 \leq C \left(\|\nabla m\|_{L^\infty}^2 + \|v\|_{L^\infty}^2 \right) \|a\|_{L^2}^2.$$

Since $m \in L^2(0, T; H^3(\Omega))$, the map $t \mapsto (\|\nabla m\|_{L^\infty}^2 + \|v\|_{L^\infty}^2)$ is in $L^1([0, T])$. So by using the Gronwall lemma, since $a(0, x) = 0$, we can conclude that $\|a\|_{L^2}^2 = 0$ and therefore the ponctual norm of m is conserved.

Under the assumption $|m| = 1$, the equations (2) and (4) are equivalent. Hence, we have proven the existence of a strong solution for (2) in the space $L^\infty(0, T; H^2(\Omega)) \cap L^2(0, T; H^3(\Omega))$ for $T < T^*$. It remains to prove that this solution is unique.

3.5 Uniqueness for the strong solution of (2)

Let m^1 and m^2 in $L^\infty(0, T; H^2(\Omega; S^2)) \cap L^2(0, T; H^3(\Omega))$ satisfying (2). Since (2) preserves the saturation constraint (1) for the strong solutions, they satisfy (4). We denote $w = m^1 - m^2$. Then $w \in L^\infty(0, T; H^2(\Omega; S^2)) \cap L^2(0, T; H^3(\Omega))$ and is solution to

$$\begin{cases} \frac{\partial w}{\partial t} - \Delta w = T_1 + T_2 + T_3 + T_4, \\ \frac{\partial w}{\partial \nu} = 0 \text{ on } \partial\Omega, \\ w(0, x) = 0, \end{cases} \quad (16)$$

where

- $T_1 = -m^2 \times \Delta w$,
- $T_2 = m^2(\nabla w \cdot \nabla(m^1 + m^2)) + (v \cdot \nabla)w - m^2 \times (v \cdot \nabla)w$,
- $T_3 = |\nabla m^1|^2 w - w \times \Delta m^1 + w \times (v \cdot \nabla)m^1$,
- $T_4 = -m^1 \times H(w) - w \times H(m^1) - w \times (m^1 \times H(m^1)) - m^2 \times (w \times H(m^1)) - m^2 \times (m^2 \times H(w))$.

We take the inner product of (16) with w . Since

$$\int_{\Omega} T_1 w = - \int_{\Omega} \nabla m^2 \times v \cdot \nabla w,$$

and using that $\|m^1\|_{L^\infty} = \|m^2\|_{L^\infty} = 1$, we obtain that

$$\begin{aligned} \frac{1}{2} \frac{d}{dt} \|w\|_{L^2}^2 + \|\nabla w\|_{L^2}^2 &\leq \|w\|_{L^2}^2 (\|\nabla m^1\|_{L^\infty} + 1 + \|H(m^1)\|_{L^\infty}) \\ &\quad + c \|w\|_{L^2} \|\nabla w\|_{L^2} [1 + \|\nabla m^1\|_{L^\infty} + \|\nabla m^2\|_{L^\infty} \\ &\quad + \|v\|_{L^\infty}]. \end{aligned}$$

So by using the Young inequality, we obtain that

$$\frac{d}{dt} \|w\|_{L^2}^2 + \|\nabla w\|_{L^2}^2 \leq C \left(1 + \|m^1\|_{H^3(\Omega)}^2 + \|m^2\|_{H^3(\Omega)}^2 + \|v\|_{L^\infty}^2 \right) \|w\|_{L^2}^2, \quad (17)$$

and since m^1 and m^2 are in $L^2(0, T; H^3(\Omega))$, we conclude by the Gronwall lemma that for all $t \geq 0$,

$$\|w(t)\|_{L^2}^2 \leq \|w(0)\|_{L^2}^2 \exp \left(C \int_0^t \left(1 + \|m^1(\tau)\|_{H^3(\Omega)}^2 + \|m^2(\tau)\|_{H^3(\Omega)}^2 + \|v(\tau)\|_{L^\infty}^2 \right) d\tau \right),$$

and since $w(0) = 0$, we obtain that $w = 0$ for all t , which concludes the proof of Theorem 1.1.

4 Study of the case $s = 3$

In order to obtain more regular solutions, it would be standard to multiply the Galerkin approximation (9) by $\Delta^3 m_n$ to obtain a H^3 -estimate. Unfortunately, the non local term $H(m_n)$ does not satisfy the homogeneous Neumann boundary condition, so that the necessary double integration by parts is not possible. Therefore, the H^3 regularity is obtained by derivation of (9) with respect to t in order to obtain a H^1 estimate on $\partial_t m$. We conclude the proof by a bootstrap argument to obtain that Δm and $\partial_t m$ have the same regularity.

4.1 H^1 regularity for $\partial_t m$

We differentiate the Galerkin approximation (9) with respect to t . Denoting by $w_{1,n} = \partial_t m_n$ the time derivative of m_n , we obtain that

$$\frac{\partial w_{1,n}}{\partial t} - \Delta w_{1,n} = \sum_{i=1}^6 T_i, \quad (18)$$

with

- $T_1 = -P_n(m_n \times \Delta w_{1,n})$
- $T_2 = P_n \left(|\nabla m_n|^2 w_{1,n} + 2(\nabla m_n \cdot \nabla w_{1,n})m_n - w_{1,n} \times \Delta m_n \right)$

- $T_3 = -P_n \left(w_{1,n} \times H(m_n) + m_n \times H(w_{1,n}) \right),$
- $T_4 = -P_n \left(w_{1,n} \times (m_n \times H(m_n)) + m_n \times (w_{1,n} \times H(m_n)) + m_n \times (m_n \times H(w_{1,n})) \right)$
- $T_5 = P_n \left((v \cdot \nabla) w_{1,n} + w_{1,n} \times (v \cdot \nabla) m_n + m_n \times (v \cdot \nabla) w_{1,n} \right)$
- $T_6 = P_n \left((\partial_t v \cdot \nabla) m_n + m_n \times (\partial_t v \cdot \nabla) m_n \right)$

We multiply (18) by $-\Delta w_{1,n}$ and we integrate on Ω . Since $m_n \times \Delta w_{1,n} \cdot \Delta w_{1,n} = 0$, we obtain that:

$$\frac{1}{2} \frac{d}{dt} \|\nabla w_{1,n}\|_{L^2}^2 + \|\Delta w_{1,n}\|_{L^2}^2 \leq \|\Delta w_{1,n}\|_{L^2} \sum_{i=2}^6 \|T_i\|_{L^2}.$$

Using the continuity of the operator H on $H^k(\Omega)$ for $k = 0, 1, 2$, we estimate each term:

$$\begin{aligned} \|T_2\|_{L^2} &\leq \|\nabla m_n\|_{L^6}^2 \|w_{1,n}\|_{L^6} + \|\nabla m_n\|_{L^\infty(\Omega)} \|\nabla w_{1,n}\|_{L^2(\Omega)} \|m_n\|_{L^\infty(\Omega)} \\ &\quad + \|w_{1,n}\|_{L^6(\Omega)} \|\Delta m_n\|_{L^3(\Omega)} \\ &\leq C(\|\nabla w_{1,n}\|_{L^2} + \|w_{1,n}\|_{L^2}) (1 + \|m_n\|_{H^2}^2 + \|m_n\|_{H^3(\Omega)}) \\ \|T_3 + T_4\|_{L^2} &\leq C\|w_{1,n}\|_{L^2(\Omega)} \left(1 + \|m_n\|_{H^2(\Omega)}^2\right) \\ \|T_5\|_{L^2} &\leq \|\nabla w_{1,n}\|_{L^2} (\|m_n\|_{L^\infty(\Omega)} + 1) \|v\|_{L^\infty(\Omega)} + \|w_{1,n}\|_{L^6(\Omega)} \|v\|_{L^\infty(\Omega)} \|\nabla m_n\|_{L^3(\Omega)} \\ &\leq C(\|w_{1,n}\|_{L^2} + \|\nabla w_{1,n}\|_{L^2}) (\|m_n\|_{H^2(\Omega)} + 1) \|v\|_{L^\infty(\Omega)} \\ \|T_6\|_{L^2} &\leq \|\partial_t v\|_{L^2(\Omega)} \|\nabla m_n\|_{L^\infty(\Omega)} (1 + \|m_n\|_{L^\infty(\Omega)}^2) \\ &\leq C\|\partial_t v\|_{L^2(\Omega)} (1 + \|m_n\|_{H^2(\Omega)}^2) \|m_n\|_{H^3(\Omega)}. \end{aligned}$$

Since v is sufficiently regular, by absorbing $\|\Delta w_{1,n}\|_{L^2}$, we obtain that:

$$\frac{d}{dt} \|\nabla w_{1,n}\|_{L^2}^2 + \|\Delta w_{1,n}\|_{L^2}^2 \leq g_1^n(t) \|\nabla w_{1,n}\|_{L^2}^2 + g_2^n(t), \quad (19)$$

where

$$g_1^n(t) = C \left(1 + \|m_n\|_{H^2(\Omega)}^2\right) \|m_n\|_{H^3(\Omega)}^2 + C \left(1 + \|m_n\|_{H^2(\Omega)}^4\right),$$

and

$$g_2^n(t) = C \|w_{1,n}\|_{L^2}^2 (1 + \|m_n\|_{H^3(\Omega)}^2 + \|m_n\|_{H^2(\Omega)}^4) + C \|\partial_t v\|_{L^2}^2 (1 + \|m_n\|_{H^2(\Omega)}^4) \|m_n\|_{H^3(\Omega)}.$$

Let us now estimate the initial value of $\partial_t m_n$. By Equation (9) taken at $t = 0$, we have:

$$\begin{aligned} w_{1,n}(0) &= \partial_t m_n(0) \\ &= P_n \left(\Delta m_n(0) + |\nabla m_n(0)|^2 m_n(0) - m_n(0) \times \Delta m_n(0) \right) \\ &\quad - P_n \left(m_n(0) \times H(m_n(0)) + m_n(0) \times (m_n(0) \times H(m_n(0))) \right) \\ &\quad + P_n \left((v_0 \cdot \nabla) m_n(0) + m_n(0) \times (v_0 \cdot \nabla) m_n(0) \right) \end{aligned} \quad (20)$$

(we recall that $v_0(x) = v(0, x)$).

Using Proposition 1 we can estimate the H^1 norm of $w_{1,n}(0)$ without compatibility condition on the following way:

$$\begin{aligned}
\|w_{1,n}(0)\|_{H^1(\Omega)} &\leq \left\| \Delta m_n(0) + |\nabla m_n(0)|^2 m_n(0) - m_n(0) \times \Delta m_n(0) \right\|_{H^1(\Omega)} \\
&\quad + \left\| m_n(0) \times H(m_n(0)) + m_n(0) \times (m_n(0) \times H(m_n(0))) \right\|_{H^1(\Omega)} \\
&\quad + \left\| (v_0 \cdot \nabla) m_n(0) + m_n(0) \times (v_0 \cdot \nabla) m_n(0) \right\|_{H^1(\Omega)} \\
&\leq C(1 + \|m_n(0)\|_{H^3(\Omega)}^3) \\
&\leq C(1 + \|m_0\|_{H^3(\Omega)}^3),
\end{aligned}$$

using the compatibility condition (3) and Proposition 1.

Let us consider $z_n(t) = \|\nabla w_{1,n}\|_{H^1(\Omega)}^2$ then we have

$$\begin{cases} z_n'(t) \leq g_1^n(t) z_n(t) + g_2^n(t), \\ z_n(0) \leq C(1 + \|m_0\|_{H^3(\Omega)}^3). \end{cases}$$

Let ξ_n be the solution of

$$\begin{cases} \xi_n'(t) = g_1^n(t) \xi_n(t) + g_2^n(t), \\ \xi_n(0) = C(1 + \|m_0\|_{H^3(\Omega)}^3). \end{cases}$$

We have:

$$\xi_n(t) = \xi_n(0) \exp\left(\int_0^t g_1^n(\tau) d\tau\right) + \int_0^t g_2^n(\tau) \exp\left(\int_\tau^t g_1^n(s) ds\right) d\tau.$$

From the estimates (13) and (14), we obtain that for all $T < T^*$, there exists a constant $C(T)$ independent of n such that:

$$\int_0^T g_1^n(\tau) d\tau \leq C(T) \quad \text{and} \quad \int_0^T g_2^n(\tau) d\tau \leq C(T),$$

so that we deduce that

$$\forall T < T^*, \exists C(T), \forall n, \xi_n(T) \leq C(T).$$

Therefore, by the Comparison Lemma 2.4, we obtain that for every $T < T^*$, there exists a constant $C(T)$ independent of n such that

$$\sup_{t \leq T} \|w_{1,n}(t)\|_{H^1(\Omega)}^2 \leq C,$$

and by integration of (19) on the interval $[0, T]$ we obtain

$$\int_0^T \|\Delta w_{1,n}(t)\|_{L^2}^2 dt \leq C.$$

From these inequalities, we can conclude that there exist a subsequence (w_{1,n_k}) such that

- $w_{1,n_k} \rightharpoonup \partial_t m$ in $L^2(0, T; H^2(\Omega))$ weak,
- $w_{1,n_k} \rightharpoonup \partial_t m$ in $L^\infty(0, T; H^1(\Omega))$ weak*.

Therefore we obtain that $\partial_t m$ is bounded in the space $L^\infty(0, T; H^1(\Omega)) \cap L^2(0, T; H^2(\Omega))$ for all $T < T^*$. Let us show that this estimate yields that $m \in L^\infty(0, T; H^3(\Omega)) \cap L^2(0, T; H^4(\Omega))$.

4.2 Regularity of Δm

In this paragraph, we will prove that $\Delta m \in L^\infty(0, T; H^1(\Omega)) \cap L^2(0, T; H^2(\Omega))$, so that $m \in L^\infty(0, T; H^3(\Omega)) \cap L^2(0, T; H^4(\Omega))$. By the equation, we have:

$$\Delta m - m \times \Delta m = \partial_t m - |\nabla m|^2 m + m \times (H(m) + m \times H(m)) - (v \cdot \nabla)m - m \times (v \cdot \nabla)m.$$

Let us consider g_m the map defined by $\xi \longrightarrow \xi - m \times \xi$. This map is linear bijective which inverse $\psi_m = g_m^{-1}$ is given by

$$\psi_m(y) = \xi = \frac{1}{2}(y + m \times y + (m \cdot y)m).$$

Since $|m| = 1$ then $\psi_m(\partial_t m) = \frac{1}{2}(\partial_t m + m \times \partial_t m)$, $\psi_m(m) = m$ and we obtain that:

$$\Delta m = \frac{1}{2}(\partial_t m + m \times \partial_t m) - |\nabla m|^2 m + W, \quad (21)$$

where $W = m \times (m \times H(m)) - m \times (v \cdot \nabla)m$.

Lemma 4.1. *For all $T < T^*$, $\psi_m(\partial_t m) \in L^\infty(0, T; H^1(\Omega)) \cap L^2(0, T; H^2(\Omega))$.*

Proof. We recall that

$$\psi_m(\partial_t m) = \frac{1}{2}(\partial_t m + m \times \partial_t m).$$

In the previous section, we have proven that $\partial_t m \in L^\infty(0, T; H^1(\Omega))$ and $m \in L^\infty(0, T; H^2(\Omega))$. So by Lemma 2.3,

$$\psi_m(\partial_t m) \in L^\infty(0, T; H^1(\Omega)).$$

Now, $H^2(\Omega)$ is an algebra, so since $m \in L^\infty(0, T; H^2(\Omega))$ and $\partial_t m \in L^2(0, T; H^2(\Omega))$, we obtain that

$$\psi_m(\partial_t m) \in L^2(0, T; H^2(\Omega)).$$

□

Lemma 4.2. *For all $T < T^*$, $W \in L^\infty(0, T; H^1(\Omega)) \cap L^2(0, T; H^2(\Omega))$.*

Proof. We know that $m \in L^\infty(0, T; H^2(\Omega))$. From Proposition 1, the same holds for $H(m)$, so since $L^\infty(0, T; H^2(\Omega))$ is an algebra,

$$m \times (m \times H(m)) \in L^\infty(0, T; H^2(\Omega)).$$

So *a fortiori* this term belongs to $L^\infty(0, T; H^1(\Omega)) \cap L^2(0, T; H^2(\Omega))$.

Concerning the other term, m and v are bounded in $L^\infty(0, T; H^2(\Omega))$ and ∇m is bounded in $L^\infty(0, T; H^1(\Omega))$. So by Lemma 2.3,

$$m \times (v \cdot \nabla)m \in L^\infty(0, T; H^1(\Omega)).$$

On the other hand, m and v are bounded in $L^\infty(0, T; H^2(\Omega))$ and ∇m is bounded in $L^2(0, T; H^2(\Omega))$, thus, since $H^2(\Omega)$ is an algebra,

$$m \times (v \cdot \nabla)m \in L^2(0, T; H^2(\Omega)).$$

□

We aim to deduce from (21) more regularity for m . We have:

$$|\nabla(|\nabla m|^2 m)| \leq |\nabla m|^3 + 2|m||\nabla m||D^2 m|$$

so

$$\begin{aligned} \|\nabla(|\nabla m|^2 m)\|_{L^2} &\leq C (\|\nabla m\|_{L^6}^3 + \|m\|_{L^\infty} \|\nabla m\|_{L^6} \|D^2 m\|_{L^3}) \\ &\leq C \|m\|_{H^2}^3 + C \|m\|_{H^2}^{\frac{5}{2}} \|m\|_{H^3}^{\frac{1}{2}} \end{aligned}$$

by Proposition 2.1. Since $m \in L^\infty(0, T; H^2(\Omega)) \cap L^2(0, T; H^3(\Omega))$, then

$$|\nabla m|^2 m \in L^4(0, T; H^1(\Omega)).$$

Using Lemmas 4.1 and 4.2, we obtain then from (21) that:

$$\Delta m \in L^4(0, T; H^1(\Omega)), \text{ so } m \in L^4(0, T; H^3(\Omega)).$$

Now, $\nabla m \in L^4(0, T; H^2(\Omega))$ and $m \in L^\infty(0, T; H^2(\Omega))$. So

$$|\nabla m|^2 m \in L^2(0, T; H^2(\Omega)).$$

By using lemmas 4.1 and 4.2, we obtain then that

$$\Delta m \in L^2(0, T; H^2(\Omega)),$$

and so

$$m \in L^2(0, T; H^4(\Omega)).$$

By Theorem II.5.14 in [4], since we know that $\partial_t m \in L^2(0, T; H^2(\Omega))$, we obtain that

$$m \in C^0(0, T; H^3(\Omega)).$$

This concludes the proof of Theorem 1.2.

5 Very regular solutions

It is not possible to obtain H^2 estimates on $w_{1,n} = \partial_t m_n$ using Equation (18). Indeed, we would need a uniform estimate on the initial value $w_{1,n}(0)$ in the H^2 norm, and using Proposition 1, it would be necessary to check a compatibility condition of the form:

$$\begin{aligned} &\frac{\partial}{\partial \nu} \left(\Delta m_n(0) + |\nabla m_n(0)|^2 m_n(0) - m_n(0) \times \Delta m_n(0) - m_n(0) \times H(m_n(0)) \right. \\ &\left. + m_n(0) \times (m_n(0) \times H(m_n(0))) + (v_0 \cdot \nabla) m_n(0) + m_n(0) \times (v_0 \cdot \nabla) m_n(0) \right) = 0. \end{aligned}$$

Because of the non local term $H(m_n(0))$, this condition cannot be satisfied for all n .

We assume that $m_0 \in H^4(\Omega; S^2)$ and satisfies the compatibility condition at order one (see Definition 1.3). In order to obtain more regular solutions, our strategy is the following.

1. We compute the equation (23) (see below) satisfied by the time derivative $\partial_t m$.
2. We construct a solution w_1 for this equation in $L^\infty(0, T; H^2) \cap L^2(0, T; H^3)$. At this step we need a compatibility condition at the boundary for $\partial_t m(t=0)$.
3. We already know that $\partial_t m \in L^\infty(0, T; H^1) \cap L^2(0, T; H^2)$ and satisfies (23). In addition, we show a uniqueness result for the solutions of (23), so that $w_1 = \partial_t m$.
4. Writing Δm in function of $\partial_t m$, we prove that $\Delta m \in L^\infty(0, T; H^2) \cap L^2(0, T; H^3)$, so that $m \in L^\infty(0, T; H^4) \cap L^2(0, T; H^5)$.

If $m_0 \in H^5(\Omega; S^2)$ satisfies the compatibility condition at order one, we obtain the desired regularity by differentiating the Galerkin approximation of (*) with respect to time (in the same spirit we obtained the H^3 regularity for m in the previous part).

This strategy will be used at any order to obtain very regular solutions for the Landau-Lifschitz equation.

We will detail this process at any order by proving by induction the following property $\mathcal{P}(k)$:

$$\mathcal{P}(k) \left\{ \begin{array}{l} \text{If } m_0 \in H^{2k}(\Omega; S^2) \text{ satisfies the compatibility condition at order } k-1, \text{ if for all } j \leq k-1, \\ \partial_t^j v \in \mathcal{C}^0(\mathbb{R}^+; H^{2k-2j}(\Omega)), \text{ then the solution } m \text{ of the Landau-Lifschitz equation (2) with} \\ \text{initial data } m_0 \text{ satisfies:} \\ \\ \forall T < T^*, \forall i \in \{0, \dots, k\}, \partial_t^i m \in L^\infty(0, T; H^{2k-2i}(\Omega)) \cap L^2(0, T; H^{2k-2i+1}(\Omega)). \\ \\ \text{If in addition } m_0 \in H^{2k+1}(\Omega; S^2), \text{ if for all } j \in \{0, \dots, k\}, \partial_t^j v \in \mathcal{C}^0(\mathbb{R}^+; H^{2k-2j+1}(\Omega)), \\ \text{then } m \text{ satisfies:} \\ \\ \forall T < T^*, \forall i \in \{0, \dots, k\}, \partial_t^i m \in L^\infty(0, T; H^{2k-2i+1}(\Omega)) \cap L^2(0, T; H^{2k-2i+2}(\Omega)). \end{array} \right.$$

This property is proven for $k = 1$ in Sections 3 and 4.

Let $k \geq 1$. Let us assume that $\mathcal{P}(k)$ is true, and let us establish $\mathcal{P}(k+1)$.

Let $m_0 \in H^{2(k+1)}(\Omega; S^2)$ satisfying the compatibility condition at order k . By Property $\mathcal{P}(k)$ we already know that

$$\forall T < T^*, \forall i \in \{0, \dots, k\}, \partial_t^i m \in L^\infty(0, T; H^{2k-2i+1}(\Omega)) \cap L^2(0, T; H^{2k-2i+2}(\Omega)).$$

We denote by $w_j = \partial_t^j m$ for $j \in \{0, \dots, k\}$. By differentiating (2) j times with respect to t , we obtain that w_j satisfies:

$$\left\{ \begin{array}{l} \frac{\partial w_j}{\partial t} - \Delta w_j = -m \times \Delta w_j + K_m(\nabla w_j) + L_m(w_j) + F_j, \\ \frac{\partial w_j}{\partial \nu} = 0 \quad \text{on } \mathbb{R}^+ \times \partial\Omega, \\ w_j(t=0) = V_j, \end{array} \right. \quad (22)$$

where

$$\begin{aligned} K_m(\nabla w) &= 2(\nabla m \cdot \nabla w)m + (v \cdot \nabla)w + m \times (v \cdot \nabla)w, \\ L_m(w) &= |\nabla m|^2 w - w \times \Delta m - w \times H(m) - m \times H(w) - w \times (m \times H(m)) \\ &\quad - m \times (w \times H(m)) - m \times (m \times H(w)) + w \times (v \cdot \nabla)m, \\ F_j &= \sum_{\alpha \in \bar{A}_j} C^\alpha ((\nabla w_{\alpha_1} \cdot \nabla w_{\alpha_2})w_{\alpha_3} - w_{\alpha_1} \times (w_{\alpha_2} \times H(w_{\alpha_3}))) \\ &\quad - \sum_{i=1}^{j-1} C_j^i (w_i \times \Delta w_{j-i} + w_i \times H(w_{j-i})) \\ &\quad + \sum_{i=0}^{j-1} C_j^i (\partial_t^{j-i} v \cdot \nabla)w_i + \sum_{\alpha \in \bar{A}_j} C^\alpha w_{\alpha_1} \times (\partial_t^{\alpha_2} v \cdot \nabla)w_{\alpha_3}, \end{aligned}$$

with

$$\bar{A}_j = \{(\alpha_1, \alpha_2, \alpha_3) \in \{0, \dots, j-1\}^3, \alpha_1 + \alpha_2 + \alpha_3 = j\}$$

and

$$\tilde{A}_j = \{(\alpha_1, \alpha_2, \alpha_3) \in \{0, \dots, j-1\} \times \{0, \dots, j\} \times \{0, \dots, j-1\}, \alpha_1 + \alpha_2 + \alpha_3 = j\}.$$

In particular, $w_k = \partial_t^k w$ is in $L^\infty(0, T; H^1(\Omega)) \cap L^2(0, T; H^2(\Omega))$ and satisfies:

$$\begin{cases} \frac{\partial w_k}{\partial t} - \Delta w_k = -m \times \Delta w_k + K_m(\nabla w_k) + L_m(w_k) + F_k, \\ \frac{\partial w_k}{\partial \nu} = 0 & \text{on } \mathbb{R}^+ \times \partial\Omega, \\ w_k(t=0) = V_k. \end{cases} \quad (23)$$

We aim to construct a more regular solution for (23). We establish before the following estimates.

Lemma 5.1. *For all $j \in \{1, \dots, k\}$, for all $T < T^*$,*

$$F_j \in L^\infty(0, T; H^{2k-2j+1}(\Omega)) \cap L^2(0, T; H^{2k-2j+2}(\Omega)).$$

Proof. We fix $j \leq k$. From Property $\mathcal{P}(k)$, for all $i \in \{0, \dots, j-1\}$, $w_i \in L^\infty(0, T; H^{2k-2j+3}(\Omega)) \cap L^2(0, T; H^{2k-2j+4}(\Omega))$.

For all $\alpha \in \bar{A}_j$, ∇w_{α_1} , ∇w_{α_2} , w_{α_1} , w_{α_2} , w_{α_3} and $H(w_{\alpha_3})$ are in $L^\infty(0, T; H^{2k-2j+2}(\Omega))$ which is an algebra since $2k - 2j + 2 \geq 2$. So

$$\sum_{\alpha \in \bar{A}_j} C^\alpha ((\nabla w_{\alpha_1} \cdot \nabla w_{\alpha_2}) w_{\alpha_3} - w_{\alpha_1} \times (w_{\alpha_2} \times H(w_{\alpha_3}))) \in L^\infty(0, T; H^{2k-2j+2}(\Omega)),$$

and *a fortiori* this quantity is in $L^\infty(0, T; H^{2k-2j+1}(\Omega)) \cap L^2(0, T; H^{2k-2j+2}(\Omega))$.

In the same way,

$$\sum_{i=1}^{j-1} C_j^i w_i \times H(w_{j-i}) \in L^\infty(0, T; H^{2k-2j+1}(\Omega)) \cap L^2(0, T; H^{2k-2j+2}(\Omega)).$$

On the other hand, for $i \in \{1, \dots, j-1\}$, $w_i \in L^\infty(0, T; H^{2k-2j+2}(\Omega))$ and $\Delta w_{j-i} \in L^\infty(0, T; H^{2k-2j+1}(\Omega)) \cap L^2(0, T; H^{2k-2j+2}(\Omega))$. So, since $H^s(\Omega)$ is an algebra for $s \geq 2$ or by applying Lemma 2.3, we obtain that:

$$\sum_{i=1}^{j-1} C_j^i w_i \times \Delta w_{j-i} \in L^\infty(0, T; H^{2k-2j+1}(\Omega)) \cap L^2(0, T; H^{2k-2j+2}(\Omega)).$$

Since for all $i \leq j$, $\partial_t^i v \in C^0(0, T; H^{2k-2j+1}(\Omega)) \cap L^2(0, T; H^{2k-2j+2}(\Omega))$, with the same arguments as below, we obtain that the remainder terms are in $L^\infty(0, T; H^{2k-2j+1}(\Omega)) \cap L^2(0, T; H^{2k-2j+2}(\Omega))$. \square

5.1 Construction of more regular solution for (24)

We consider the following Cauchy problem:

$$\begin{cases} \frac{\partial w}{\partial t} - \Delta w = -m \times \Delta w + K_m(\nabla w) + L_m(w) + F_k, \\ \frac{\partial w}{\partial \nu} = 0 & \text{on } \mathbb{R}^+ \times \partial\Omega, \\ w(t=0) = V_k. \end{cases} \quad (24)$$

We aim to construct a solution for (24) in $L^\infty(0, T; H^2(\Omega)) \cap L^2(0, T; H^3(\Omega))$

We recall that, from the compatibility condition at order k , $V_k \in H^2(\Omega)$ and $\frac{\partial V_k}{\partial \nu} = 0$ on $\partial\Omega$.

We consider the following Galerkin approximation for (24):

$$\begin{cases} w^n \in \mathcal{C}^1([0, T_n]; V_n) \\ \frac{\partial w^n}{\partial t} - \Delta w^n = P_n \left(-m \times \Delta w^n + K_m(\nabla w^n) + L_m(w^n) + F_k \right), \\ w^n(t=0) = P_n(V_k). \end{cases} \quad (25)$$

Since the coefficients of this ordinary differential equation are continuous on $[0, T^*]$, since the equation is linear, the maximal existence time T_n equals T^* . Let us obtain uniform estimates on w^n .

Bound for the initial data. Using the compatibility condition at order k , we know that $V_k \in H^2(\Omega)$ and that $\frac{\partial V_k}{\partial \nu} = 0$ on $\partial\Omega$, so we can apply Proposition 1:

$$\exists K, \forall n, \|w^n(0)\|_{H^2} \leq K. \quad (26)$$

L^2 estimate. We multiply (25) by w^n and we obtain that:

$$\frac{1}{2} \frac{d}{dt} \|w^n\|_{L^2}^2 + \|\nabla w^n\|_{L^2}^2 = T_1 + T_2 + T_3 + T_4$$

where

$$\begin{aligned} T_1 &= - \int_{\Omega} m \times \Delta w^n \cdot w^n = \int_{\Omega} m \times w^n \cdot \Delta w^n \\ &= - \int_{\Omega} \nabla(m \times w^n) \cdot \nabla w^n \quad (\text{using the homogeneous boundary conditions}) \\ &= \int_{\Omega} (\nabla m) \times w^n \cdot \nabla w^n \\ &\leq \|\nabla m\|_{L^\infty} \|w^n\|_{L^2} \|\nabla w^n\|_{L^2}, \\ T_2 &= \int_{\Omega} K_m(\nabla w^n) w^n \\ &\leq \|K_m(\nabla w^n)\|_{L^2} \|w^n\|_{L^2} \\ &\leq (2\|\nabla m\|_{L^\infty} + 2\|v\|_{L^\infty}) \|w^n\|_{L^2} \|\nabla w^n\|_{L^2} \quad (\text{since } \|m\|_{L^\infty} = 1), \\ T_3 &= \int_{\Omega} L_m(w^n) w^n \\ &\leq \|L_m(w^n)\|_{L^2} \|w^n\|_{L^2} \\ &\leq (\|\nabla m\|_{L^6}^2 + \|\Delta m\|_{L^3} + \|\nabla m\|_{L^6} \|v\|_{L^\infty}) \|w^n\|_{L^6} \|w^n\|_{L^2} \\ &\quad + C(\|H(m)\|_{L^\infty} + 1) \|w^n\|_{L^2}^2, \\ T_4 &= \int_{\Omega} F_k \cdot w^n \\ &\leq \|F_k\|_{L^2} \|w^n\|_{L^2}. \end{aligned}$$

H^2 estimate. We multiply (25) by $-\Delta^2 w^n$ and we obtain that:

$$\frac{1}{2} \frac{d}{dt} \|\Delta w^n\|_{L^2}^2 + \|\nabla \Delta w^n\|_{L^2}^2 \leq I_1 + I_2 + I_3 + I_4,$$

where

$$\begin{aligned} I_1 &= \|(\nabla m) \times \Delta w^n\|_{L^2} \|\nabla \Delta w^n\|_{L^2} \\ &\leq \|\nabla m\|_{L^\infty} \|w^n\|_{L^2} \|\nabla \Delta w^n\|_{L^2}, \\ I_2 &= \left| \int_{\Omega} \nabla(K_m(\nabla w^n)) \nabla \Delta w^n \right| \\ &\leq C \int_{\Omega} (|D^2 m| |\nabla w^n| + |\nabla m| |D^2 w^n| + |\nabla m|^2 |\nabla w^n| + |\nabla v| |\nabla w^n|) |\nabla \Delta w^n| \\ &\quad + \int_{\Omega} (|v| |D^2 w^n| + |\nabla m| |v| |\nabla w^n|) |\nabla \Delta w^n| \\ &\leq C \left(\|D^2 m\|_{L^3} \|\nabla w^n\|_{L^6} + \|\nabla m\|_{L^\infty} \|D^2 w^n\|_{L^2} + \|\nabla m\|_{L^6}^2 \|\nabla w^n\|_{L^6} \right. \\ &\quad \left. + \|\nabla v\|_{L^\infty} \|\nabla w^n\|_{L^2} + \|v\|_{L^\infty} \|D^2 w^n\|_{L^2} + \|\nabla m\|_{L^\infty} \|v\|_{L^\infty} \|\nabla w^n\|_{L^2} \right) \|\nabla \Delta w^n\|_{L^2}, \\ I_3 &= \left| \int_{\Omega} \nabla(L_m(w^n)) \cdot \nabla \Delta w^n \right| \\ &\leq C \int_{\Omega} (|D^2 m| |\nabla w^n| + |\nabla m|^2 |\nabla w^n| + |\Delta m| |\nabla w^n| + |\nabla \Delta m| |w^n| + |w^n| |\nabla H(m)|) |\nabla \Delta w^n| \\ &\quad + C \int_{\Omega} (|H(m)| |\nabla w^n| + |\nabla m| |w^n| + |\nabla w^n| |v| |\nabla m| + |w^n| |\nabla v| |\nabla m| + |w^n| |v| |D^2 m|) |\nabla \Delta w^n| \\ &\leq C (\|D^2 m\|_{L^3} + \|\nabla m\|_{L^6}^2 + \|H(m)\|_{L^3} + \|\nabla m\|_{L^3} + \|v\|_{L^\infty}) \|\nabla w^n\|_{L^6} \|\nabla \Delta w^n\|_{L^2} \\ &\quad + C (\|\nabla \Delta m\|_{L^2} + \|\nabla H(m)\|_{L^2} + \|\nabla m\|_{L^2} + \|\nabla v\|_{L^4} \|\nabla m\|_{L^4}) \|w^n\|_{L^\infty} \|\nabla \Delta w^n\|_{L^2} \\ &\quad + \|v\|_{L^\infty} \|D^2 m\|_{L^4} \|w^n\|_{L^\infty} \|\nabla \Delta w^n\|_{L^2}, \\ I_4 &\leq \|\nabla F_k\|_{L^2} \|\nabla \Delta w^n\|_{L^2}. \end{aligned}$$

By adding up the previous L^2 and H^2 estimates, after absorbing $\|\nabla \Delta w^n\|_{L^2}$, we obtain that

$$\frac{d}{dt} \|w^n\|_{2,\Omega}^2 + \|\nabla w^n\|_{2,\Omega}^2 \leq g_1(t) \|w^n\|_{2,\Omega}^2 + g_2^n(t), \quad (27)$$

where

$$g_1(t) = C(1 + \|m\|_{H^2}^2)(1 + \|v\|_{H^2}) + C\|m\|_{H^3}^2,$$

$$g_2^n(t) = C\|F_k\|_{H^1}^2.$$

We already know that for all $T < T^*$, $m \in L^\infty(0, T; H^2(\Omega)) \cap L^2(0, T; H^3(\Omega))$, so $g_1 \in L^1(0, T)$ for all $T < T^*$. In addition, by lemma 5.1, g_2^n is uniformly bounded (with respect to n) in $L^1(0, T)$ for all $T < T^*$. Since (26) yields a uniform bound for $\|w^n(0)\|_{2,\Omega}^2$, by the comparison lemma, we deduce that for all $T < T^*$, there exists a constant $C(T)$ such that for all n ,

$$\|w^n\|_{L^\infty(0, T; H^2(\Omega))} + \|w^n\|_{L^2(0, T; H^3(\Omega))} \leq C(T). \quad (28)$$

By the equation (25), we obtain in addition the following estimate on the time derivatives:

$$\forall T < T^*, \exists C(T), \forall n, \|\partial_t w^n\|_{L^\infty(0,T;L^2(\Omega))} + \|\partial_t w^n\|_{L^2(0,T;H^1(\Omega))} \leq C(T). \quad (29)$$

By standard arguments (see section 3.3) we obtain by extracting subsequences that there exists w satisfying (24) such that

$$\forall T < T^*, w \in L^\infty(0, T; H^2(\Omega)) \cap L^2(0, T; H^3(\Omega)).$$

5.2 Uniqueness for (24)

Let w_1 and w_2 two solutions of (24) in $L^\infty(0, T; H^1(\Omega)) \cap L^2(0, T; H^2(\Omega))$. We denote $W = w_1 - w_2$, and we remark that w satisfies the following problem:

$$\begin{cases} \frac{\partial W}{\partial t} - \Delta W = -m \times \Delta W + K_m(\nabla W) + L_m(W), \\ \frac{\partial W}{\partial \nu} = 0 & \text{on } \mathbb{R}^+ \times \partial\Omega, \\ W(t=0) = 0. \end{cases} \quad (30)$$

Since $W \in L^\infty(0, T; H^1(\Omega)) \cap L^2(0, T; H^2(\Omega))$ and since $\partial_t W \in L^2(0, T; L^2(\Omega))$, we can perform the following L^2 estimate, taking the inner product of the equation with W :

$$\begin{aligned} \frac{d}{dt} \|W\|_{L^2}^2 + 2\|\nabla W\|_{L^2}^2 &\leq C(\|\nabla m\|_{L^\infty} + \|v\|_{L^\infty}) \|W\|_{L^2} \|\nabla W\|_{L^2} \\ &\quad + (\|\nabla m\|_{L^6}^2 + \|\Delta m\|_{L^3} + \|\nabla m\|_{L^6} \|v\|_{L^\infty}) \|W\|_{L^6} \|W\|_{L^2} \\ &\quad + C(\|H(m)\|_{L^\infty} + 1) \|W\|_{L^2}^2. \end{aligned}$$

We multiply Equation (30) by $-\Delta W$ and we integrate on $[0, T] \times \Omega$ for all $T < T^*$. We can do it because ΔW and $\frac{\partial W}{\partial t}$ are in $L^2([0, T] \times \Omega)$. Using that $W \in C^0(0, T; H^1(\Omega))$ and that $W(0) = 0$, we obtain that:

$$\|\nabla W(T)\|_{L^2}^2 + 2 \int_0^T \int_\Omega |\Delta W|^2 dx dt = -2 \int_0^T \int_\Omega (K_m(\nabla W) + L_m(W)) \Delta W.$$

We have

$$\|K_m(\nabla W)\|_{L^2} \leq C\|\nabla m\|_{L^\infty} \|\nabla W\|_{L^2} + C\|v\|_{L^\infty} \|W\|_{L^2},$$

and

$$\|L_m(W)\|_{L^2} \leq C\|\nabla m\|_{L^6}^2 \|W\|_{L^6} + C\|W\|_{L^6} \|\Delta m\|_{L^3} + C\|W\|_{L^2} (\|H(m)\|_{L^\infty} + \|\nabla m\|_{L^\infty} \|v\|_{L^\infty}).$$

We integrate the L^2 estimate from 0 to T . Adding up with the H^1 estimate and using Young formula, we obtain that:

$$\|W(T)\|_{L^2}^2 + \|\nabla W(T)\|_{L^2}^2 + \int_0^T \int_\Omega (|\nabla W|^2 + |\Delta W|^2) dx dt \leq C \int_0^T \left(\|\nabla W(t)\|_{L^2}^2 + \|W(t)\|_{L^2}^2 \right) g(t) dt,$$

where

$$g(t) = 1 + \|v(t)\|_{L^\infty}^2 + \|m(t)\|_{2,\Omega}^4 + \|\nabla m(t)\|_{L^\infty}^2 (1 + \|v(t)\|_{L^\infty}^2).$$

By properties of m and v , $g \in L^1(0, T)$ for all $T < T^*$ so that we can use the Gronwall lemma to conclude that

$$W = 0 \text{ on } [0, T^*] \times \Omega.$$

Now, w and $\partial_t^k m$ are solutions for (24) in the space $C^0([0, T^*]; H^1(\Omega)) \cap L^2(0, T; H^2(\Omega))$, so by the previous uniqueness result, $w = \partial_t^k m$, so that

$$\forall T < T^*, \partial_t^k m \in L^\infty(0, T; H^2(\Omega)) \cap L^2(0, T; H^3(\Omega)). \quad (31)$$

5.3 Regularity for m

From Property $\mathcal{P}(k)$ we know that

$$\forall j \in \{0, \dots, k\}, \quad w_j := \partial_t^j m \in L^\infty(0, T; H^{2k-2j+1}(\Omega)) \cap L^2(0, T; H^{2k-2j+2}(\Omega)).$$

In addition, we have proven in the previous subsection that $\partial_t^k m \in L^\infty(0, T; H^2(\Omega)) \cap L^2(0, T; H^3(\Omega))$.

Let us prove by induction the following claim:

Claim. For all $T < T^*$, for $j \in \{0, \dots, k-1\}$, $w_{k-j} \in L^\infty(0, T; H^{2+2j}(\Omega)) \cap L^2(0, T; H^{3+2j}(\Omega))$.

Proof. This property is true for $j = 0$ by (31).

Let $j \in \{1, \dots, k-1\}$. Let us assume that the property is true at the rank $j-1$. Then, from Equation (22), replacing j by $k-j$, we have:

$$\Delta w_{k-j} - m \times \Delta w_{k-j} = \partial_t w_{k-j} - K_m(\nabla w_{k-j}) - L_m(w_{k-j}) - F_{k-j},$$

that is, since $\partial_t w_{k-j} = w_{k-j+1}$,

$$\Delta w_{k-j} = \psi_m \left(w_{k-j+1} - K_m(\nabla w_{k-j}) - L_m(w_{k-j}) - F_{k-j} \right)$$

(where ψ_m is defined on page 16). We have:

- $m \in L^\infty(0, T; H^{2k+1}(\Omega))$
- $\nabla m \in L^\infty(0, T; H^{2k}(\Omega))$
- $\Delta m \in L^\infty(0, T; H^{2k-1}(\Omega))$
- $w_{k-j} \in L^\infty(0, T; H^{2j+1}(\Omega))$
- $v \in L^\infty(0, T; H^{2k}(\Omega))$.

On the one hand, we recall that

$$K_m(\nabla w_{k-j}) = 2(\nabla m \cdot \nabla w_{k-j})m + (v \cdot \nabla)w_{k-j} + m \times (v \cdot \nabla)w_{k-j}.$$

Since $2j \leq 2k$, ∇m , ∇w_{k-j} , v and m are in $L^\infty(0, T; H^{2j}(\Omega))$, which is an algebra (since $j \geq 1$), then

$$K_m(\nabla w_{k-j}) \in L^\infty(0, T; H^{2j}(\Omega)).$$

Since $2j+1 \leq 2k$, ∇m , v and m are in $L^\infty(0, T; H^{2j+1}(\Omega))$ and since $\nabla w_{k-j} \in L^2(0, T; H^{2j+1}(\Omega))$, then

$$K_m(\nabla w_{k-j}) \in L^2(0, T; H^{2j+1}(\Omega)).$$

On the other hand, we recall that

$$\begin{aligned} L_m(w_{k-j}) = & |\nabla m|^2 w_{k-j} - w_{k-j} \times \Delta m - w_{k-j} \times H(m) - m \times H(w_{k-j}) - w_{k-j} \times (m \times H(m)) \\ & - m \times (w_{k-j} \times H(m)) - m \times (m \times H(w_{k-j})) + w_{k-j} \times (v \cdot \nabla)m. \end{aligned}$$

From Proposition 2, $H(w_{k-j})$ and w_{k-j} have the same regularity, *i.e.* are in $L^\infty(0, T; H^{2j+1}(\Omega))$. We remark that $j \leq k-1$ so that $2j+1 \leq 2k-1$. So ∇m , $H(m)$, m and Δm are in $L^\infty(0, T; H^{2j+1}(\Omega))$. The same holds for v . Since this space is an algebra, we obtain that

$$L_m(w_{k-j}) \in L^\infty(0, T; H^{2j+1}(\Omega)).$$

From Lemma 5.1,

$$F_{k-j} \in L^\infty(0, T; H^{2j+1}(\Omega)).$$

From the property at rank $j - 1$,

$$w_{k-j+1} \in L^\infty(0, T; H^{2j}(\Omega)) \cap L^2(0, T; H^{2j+1}(\Omega)).$$

In addition, since $m \in L^\infty(0, T; H^{2j+1}(\Omega))$, we keep the same regularity when we compose with ψ_m , so:

$$\Delta w_{k-j} \in L^\infty(0, T; H^{2j}(\Omega)) \cap L^2(0, T; H^{2j+1}(\Omega)),$$

and by elliptic regularity results,

$$w_{k-j} \in L^\infty(0, T; H^{2j+2}(\Omega)) \cap L^2(0, T; H^{2j+3}(\Omega)).$$

This concludes the proof of the claim. \square

In particular, we have proven that

$$w_1 = \partial_t m \in L^\infty(0, T; H^{2k}(\Omega)) \cap L^2(0, T; H^{2k+1}(\Omega)).$$

Now, we have:

$$\Delta m = \frac{1}{2} (\partial_t m + m \times \partial_t m) - |\nabla m|^2 m + W,$$

where $W = m \times (m \times H(m)) - m \times (v \cdot \nabla) m$.

Since ∇m , m and $H(m)$ are in $L^\infty(0, T; H^{2k}(\Omega))$, the same holds for $-|\nabla m|^2 m + W$.

Since $m \in L^\infty(0, T; H^{2k+1}(\Omega))$, $\partial_t m + m \times \partial_t m \in L^\infty(0, T; H^{2k}(\Omega)) \cap L^2(0, T; H^{2k+1}(\Omega))$.

Therefore, $\Delta m \in L^\infty(0, T; H^{2k}(\Omega)) \cap L^2(0, T; H^{2k+1}(\Omega))$, so

$$m \in L^\infty(0, T; H^{2k+2}(\Omega)) \cap L^2(0, T; H^{2k+3}(\Omega)).$$

5.4 Regularity for $\partial_t^k m$ in the case $m_0 \in H^{2k+3}(\Omega)$

We assume now in addition that $m_0 \in H^{2k+3}(\Omega; S^2)$ and satisfies the compatibility condition at order k . From the previous case, we already know that

$$\forall j \in \{0, \dots, k\}, \forall T < T^*, \partial_t^j m \in L^\infty(0, T; H^{2k+2-2j}(\Omega)) \cap L^2(0, T; H^{2k+3-2j}(\Omega)).$$

In addition, we know that $\partial_t^k m$ is the unique solution for Equation (24). We will prove that we can improve the regularity of $\partial_t^k m$ using the Galerkin formulation (25).

We derivate (25) with respect to time. Denoting by $\alpha_n := \partial_t w^n$, we obtain that:

$$\begin{cases} \frac{\partial \alpha^n}{\partial t} - \Delta \alpha^n = -P_n \left(m \times \Delta \alpha^n + K_m(\nabla \alpha^n) + L_m(\alpha^n) + \tilde{F}_k \right) \text{ on } [0, T_n[, \\ \alpha^n(0) = P_n \left(\Delta(P_n(V_k)) - m_0 \times \Delta(P_n(V_k)) + K_{m_0}(\nabla(P_n(V_k))) + L_{m_0}(P_n(V_k)) + F_k(t=0) \right), \end{cases} \quad (32)$$

where

$$\tilde{F}_k = -\partial_t m \times \Delta w^n + \partial_t(K_m)(\nabla w^n) + \partial_t(L_m)(w^n) + \partial_t F_k.$$

We multiply (32) by $\Delta \alpha^n$. After integration on Ω and using Young formula, we obtain that:

$$\frac{d}{dt} \|\nabla \alpha^n\|_{L^2}^2 + \|\Delta \alpha^n\|_{L^2}^2 \leq C \left(\|K_m(\nabla \alpha^n)\|_{L^2}^2 + \|L_m(\alpha^n)\|_{L^2}^2 + \|\tilde{F}_k\|_{L^2}^2 \right).$$

We have:

$$\begin{aligned}
\|K_m(\nabla\alpha^n)\|_{L^2} &= \|2(\nabla m \cdot \nabla\alpha^n)m + (v \cdot \nabla)\alpha^n + m \times (v \cdot \nabla)\alpha^n\|_{L^2} \\
&\leq C(\|\nabla m\|_{L^\infty} + \|v\|_{L^\infty})\|\nabla\alpha^n\|_{L^2}, \\
\|L_m(\alpha^n)\|_{L^2} &\leq \| |\nabla m|^2\alpha^n - \alpha^n \times \Delta m - \alpha^n \times H(m) - m \times H(\alpha^n) - \alpha^n \times (m \times H(m)) \|_{L^2} \\
&\quad + \|m \times (\alpha^n \times H(m)) + m \times (m \times H(\alpha^n))\|_{L^2} + \|\alpha^n \times (v \cdot \nabla)m\|_{L^2} \\
&\leq C\left(\|\nabla m\|_{L^\infty}^2 + \|\Delta m\|_{L^\infty} + \|H(m)\|_{L^\infty} + \|v\|_{L^\infty}\|\nabla m\|_{L^\infty} + 1\right)\|\alpha^n\|_{L^2} \\
&\leq C(T) \text{ for all } T < T^*,
\end{aligned}$$

since $m \in L^\infty(0, T; H^4(\Omega))$ and by (29).

Concerning \tilde{F}_k , we have:

$$\begin{aligned}
\|\partial_t m \times \Delta w^n\|_{L^2} &\leq C\|\partial_t m\|_{L^\infty}\|w^n\|_{H^2} \\
&\leq C(T) \text{ for all } T < T^*,
\end{aligned}$$

since $\partial_t m \in L^\infty(0, T; H^{2k}(\Omega))$ and by (28),

$$\begin{aligned}
\partial_t(K_m)(\nabla w^k) &= 2(\nabla\partial_t m \cdot \nabla w^n)m + 2(\nabla m \cdot \nabla w^n)\partial_t m + (\partial_t v \cdot \nabla)w^n + \partial_t m \times (v \cdot \nabla)w^n \\
&\quad + m \times (\partial_t v \cdot \nabla)w^n
\end{aligned}$$

$$\|\partial_t(K_m)(\nabla w^k)\|_{L^2} \leq C\|\nabla w^n\|_{L^6}(\|\nabla\partial_t m\|_{L^3} + \|\nabla m\|_{L^\infty}\|\partial_t m\|_{L^3} + \|\partial_t v\|_{L^3} + \|v\|_{L^\infty}\|\partial_t m\|_{L^3}).$$

Since $k \geq 1$, since $\partial_t m \in L^\infty(0, T; H^2(\Omega))$, then $\nabla\partial_t m \in L^\infty(0, T; H^1(\Omega))$ so by embedding, for all $t \leq T < T^*$, $\|\nabla\partial_t m(t)\|_{L^3} \leq K(T)$.

In addition,

$$\|\nabla w^n\|_{L^6} \leq C\|w^n\|_{H^2} \leq K(T) \text{ by (28)}.$$

So,

$$\|\partial_t(K_m)(\nabla w^k)(t)\|_{L^2} \leq K(T) \text{ for all } t \leq T < T^*.$$

$$\begin{aligned}
\partial_t(L_m)(w_k) &= 2(\nabla m \cdot \nabla\partial_t m)w^n - w^n \times \Delta\partial_t m - w^n \times H(\partial_t m) - \partial_t m \times H(w^n) \\
&\quad - w^n \times (\partial_t m \times H(m)) - w^n \times (m \times H(\partial_t m)) - \partial_t m \times (w^n \times H(m)) \\
&\quad - m \times (w^n \times H(\partial_t m)) - \partial_t m \times (m \times H(w^n)) - m \times (\partial_t m \times H(w^n)) \\
&\quad + w^n \times (\partial_t v \cdot \nabla)m + w^n \times (v \cdot \nabla)\partial_t m
\end{aligned}$$

$$\begin{aligned}
\|\partial_t(L_m)(w_k)\|_{L^2} &\leq C\|w^n\|_{L^\infty}(\|\nabla m\|_{L^\infty}\|\nabla\partial_t m\|_{L^2} + \|\Delta\partial_t m\|_{L^2} + \|H(\partial_t m)\|_{L^2}) \\
&\quad + C\|w^n\|_{L^\infty}(\|\partial_t m\|_{L^2}\|H(m)\|_{L^\infty} + \|\partial_t m\|_{L^2} + \|v\|_{L^\infty}\|\nabla\partial_t m\|_{L^2}) \\
&\quad + C\|w^n\|_{L^\infty}\|\partial_t v\|_{L^\infty}\|\nabla m\|_{L^2} + C\|H(w^n)\|_{L^2}\|\partial_t m\|_{L^2}.
\end{aligned}$$

Using Proposition 2, using (28) and the known bounds on $\partial_t m$, we obtain that:

$$\forall T < T^*, \exists C(T), \forall n, \frac{d}{dt}\|\alpha^n\|_{L^2}^2 + \|\Delta\alpha^n\|_{L^2}^2 \leq C(T)\|\nabla\alpha^n\|_{L^2}^2 + C(T). \quad (33)$$

Now from Proposition 1, we have:

$$\begin{aligned} \|\nabla\alpha^n(0)\|_{L^2} &\leq \|\Delta(P_n(V_k)) - m_0 \times \Delta(P_n(V_k)) + K_{m_0}(\nabla(P_n(V_k))) + L_{m_0}(P_n(V_k)) + F_k(t=0)\|_{H^1} \\ &\leq C(1 + \|P_n(V^k)\|_{H^3}). \end{aligned}$$

From the compatibility condition at order k , $\frac{\partial V^k}{\partial \nu} = 0$ on $\partial\Omega$, so from Proposition 1,

$$\|P_n(V^k)\|_{H^3} \leq C\|V^k\|_{H^3} \text{ for all } n.$$

Therefore, there exists a constant C such that for all n ,

$$\|\nabla\alpha^n(0)\|_{L^2} \leq C. \quad (34)$$

Estimates (33) and (34) coupled with the comparison lemma 2.4 yield that:

$$\forall T < T^*, \exists C(T), \|\alpha^n\|_{L^\infty(0,T;H^1(\Omega))} + \|\alpha^n\|_{L^2(0,T;H^2(\Omega))} \leq C(T).$$

Since $\alpha^n \rightharpoonup \partial_t^{k+1}m$ in the distribution sense, we obtain that

$$\forall T < T^*, \partial_t^{k+1}m \in L^\infty(0,T;H^1(\Omega)) \cap L^2(0,T;H^2(\Omega)). \quad (35)$$

5.5 Regularity for m in the case $m_0 \in H^{2k+3}(\Omega)$

We know from Section 5.3 that

$$\forall j \in \{0, \dots, k\}, \partial_t^j m \in L^\infty(0,T;H^{2k-2j+2}(\Omega)) \cap L^2(0,T;H^{2k-2j+3}(\Omega)).$$

So coupling this estimate with (35), we gain one rank for the regularity of each w_j , and exactly with the same method as in Section 5.3, we prove that

$$\forall j \in \{0, \dots, k+1\}, \partial_t^j m \in L^\infty(0,T;H^{2k-2j+3}(\Omega)) \cap L^2(0,T;H^{2k-2j+4}(\Omega)),$$

so that the proof of Property $\mathcal{P}(k+1)$ is complete.

References

- [1] François Alouges and Alain Soyeur, *On global weak solutions for Landau-Lifshitz equations: existence and nonuniqueness*, Nonlinear Anal. **18** (1992), no. 11, 1071–1084.
- [2] T. Aubin, *Un théorème de compacité*, C. R. Acad. Sci. Paris **256** (1963), 5242–5044.
- [3] Gaël Bonithon, *Landau-Lifschitz-Gilbert equation with applied electric current*. Discrete Contin. Dyn. Syst. 2007, Dynamical Systems and Differential Equations. Proceedings of the 6th AIMS International Conference, suppl., 138-144.
- [4] F. Boyer, P. Fabrie, *Éléments d'analyse pour l'étude de quelques modèles d'écoulements de fluides visqueux incompressibles*, Mathématiques & Applications (Berlin) 52. Springer-Verlag, Berlin, 2006.
- [5] William F. Brown, *Micromagnetics*, Wiley, New York (1963).
- [6] Gilles Carbou and Pierre Fabrie, *Time average in micromagnetism*, J. Differential Equations **147** (1998), no. 2, 383–409.
- [7] G. Carbou, P. Fabrie, *Regular Solutions for Landau-Lifschitz equations in a bounded domain*, Differential Integral Equations **14** (2001), 213–229.

- [8] G. Carbou, P. Fabrie, O. Guès, *On the ferromagnetism equations in the non static case*, Commun. Pure Appl. Anal. **3** (2004), no. 3, 367–393.
- [9] G. Foias and R. Temam, *Remarques sur les équations de Navier-Stokes stationnaires et les phénomènes successifs de bifurcation*, An. Sc. Norm. Super. Pisa IV, **5** (1978), 29–63.
- [10] O.A. Ladyzhenskaya, *The boundary value problem of mathematical physics*, Springer Verlag Applied Math. Sciences, **49**, 1985.
- [11] L. Landau et E. Lifschitz, *Electrodynamique des milieux continus, cours de physique théorique, tome VIII* (ed. Mir) Moscou, 1969.
- [12] J. Simon, *Compact sets in the space $L^p(0, T; B)$* , Ann. Mat. Pura Appli. **146** (1987), no. 4, 65–96.
- [13] A.Thiaville, Y.Nakatani, J.Miltat and Y.Susuki, *Micromagnetic understanding of current driven domain wall motion in patterned nanowires*, Europhys. Lett. **69** (2005), no. 6, 990-996.
- [14] A.Thiaville, Y.Nakatani, J.Miltat and N. Vernier, *Domain wall motion by spin-polarized current: a micromagnetic study*, J. Appl. Phys., Part 2, **95** (2004), no. 11, 7049-7051.
- [15] A.Thiaville, J.M.Garcia and J.Miltat, *Domain wall dynamics in nanowires*. Journal of Magnetism and Magnetic Materials **242-245** (2002) 1061-1063.
- [16] Augusto Visintin, *On Landau Lifschitz equation for ferromagnetism*, Japan Journal of Applied Mathematics **1** (1985), no. 2, 69–84.