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# **BEM Simulation Of 3D Updated Resin Front For LCM Processes**

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Abstract. The use of composite materials in large structures has increased in recent years. Aircraft industry has recently begun to investigate the field of Liquid Composite Molding (LCM) through research programs because of its ability to produce large parts at low cost. The present paper focuses on modeling a 3D radial impregnation through an anisotropic fibrous perform. As a preliminary work, it is assumed an isothermal flow and no hydro-mechanical coupling. Governing equations are Darcy's law and mass conservation. Simulation is performed combining Boundary Element Method (BEM) with a lagrangian moving mesh method. An analytical solution is developed to assess the numerical model.

#### Introduction

Composite materials have been used in the aircraft industry over the past three decades for their excellent specific mechanical properties. Large structural parts made of resin/fiber composites require specific forming processes. The traditional prepreg autoclave process produces high quality parts, but involves heavy manufacturing cost due to autoclave investment and use of high value added raw materials. As an alternative, recent technologies in Liquid Composite Molding (LCM) have become established recently under the designation of infusion processes. Liquid Resin Infusion (LRI) is one of them. It consists in forcing the liquid resin to circulate transversally into the dry preform using a distribution medium, consecutively to the vacuum application [1].

Prediction of resin flow during the infusion filling stage requires a full three-dimensional model in order to take into account the lag due to the distribution medium. Previous works revealed the adequacy of Boundary Element Method (BEM) solving the Darcy's flow [2-4]. It was performed assuming a two-dimensional flow, neglecting the impregnation through the thickness.

The present paper is focused on a three-dimensional numerical modeling of the Newtonian isothermal flow using constant Boundary Elements. Combining Darcy's law with incompressibility equation leads to Laplace equation which is solved using BEM at each step time on an updated lagrangian mesh. Preliminary results are presented, i.e. BEM simulations performed in the case of central injection. Computed results are compared with an analytical solution.

### **Governing equations**

**Modeling Assumptions.** Impregnation is modeled as a Darcy's flow. Fibrous perform is regarded as a porous medium. Hydro-mechanical couplings are neglected, so that the medium is assumed to be not deformable and of a constant permeability. The flow is considered as isothermal

and resin viscosity is assumed to be constant, neglecting chemical reactions effects. Consequently, governing equations reduce to Darcy's law and incompressibity equation.

**Darcy's law.** Darcy's law [5] states that the average velocity  $\vec{v}$  is proportional to the pressure gradient according to the following equation:

$$\vec{v} = -\frac{[K]}{\mu} \vec{\nabla} p \tag{1}$$

where [K] is the permeability tensor, p the resin pressure and  $\mu$  the liquid resin viscosity. Combining Eq. 1 and incompressibility equation gives:

$$\vec{\nabla} \cdot \left( -\frac{[K]}{\mu} \vec{\nabla} p \right) = 0 \tag{2}$$

For sake of simplicity, we assume in this paper that the principal orthogonal axis of the reinforcement coincides with the reference coordinates system, such as [K] reduces to diagonal matrix. Eq. 2 can then be reduced to Laplace equation in the isotropic equivalent domain by stretching the coordinates. In the Cartesian coordinate system, the equivalent coordinates  $\bar{x}_e = (x_{e_1}, x_{e_2})$  may be deduced from the physical coordinates  $\bar{x} = (x_1, x_2)$  using the following transformation:

$$x_{e_i} = \sqrt{\frac{K_e}{K_i}} x_i \tag{3}$$

where  $K_i$  is a principal value of [K] and  $K_e$  the equivalent permeability given as  $K_e = \sqrt[3]{K_1 K_2 K_3}$ . In the isotropic equivalent domain, Eq. 1 and Eq. 2 reduce simply to:

$$\vec{v}_e = -\frac{K_e}{\mu} \vec{\nabla} p \tag{4}$$

$$\Delta p = 0 \tag{5}$$

**Analytical solution.** Following the approach introduced by Adams et al. [6], an analytical solution is used to calculate the evolution of the flow front versus time. Computations are performed in the isotropic equivalent system, using the equivalent permeability  $K_e$  and the spherical injection gate of radius  $r_{0_e}$ . A constant pressure injection  $p_{0_e}$  is applied, so that the macroscopically resin front shape is spherical of radius  $r_{f_e}$ . In such configuration the resin pressure varies only in the radial direction according to:

$$\frac{p - p_f}{p_0 - p_f} = 1 - \frac{\frac{1}{r_{0e}} - \frac{1}{r}}{\frac{1}{r_{0e}} - \frac{1}{r_{fe}}}$$
(6)

Substituting the expression of the gradient pressure into Darcy's law and introducing the system porosity  $\mathcal E$  gives the radial component of the front velocity  $v_{f_e}$ :

$$v_{f_e} = \frac{K_e}{\mu \varepsilon} \frac{(p_0 - p_f)}{r_{f_e}^2 \left(\frac{1}{r_{0_e}} - \frac{1}{r_{f_e}}\right)}$$
(7)

The evolution of the flow front is governed by the following differential equation:

$$\left(\frac{dr}{dt}\right)_{f_e} = \frac{K_e}{\mu\varepsilon} \frac{\left(p_0 - p_f\right)}{r_{f_e}^2 \left(\frac{1}{r_{0_e}} - \frac{1}{r_{f_e}}\right)} \tag{8}$$

Finally, the time dependant flow front position can be deduced by integrating the previous equation during the filling stage. This leads to the following solution:

$$\left[\frac{r_{f_e}}{r_{0_e}}\right]^2 \cdot \left[2\frac{r_{f_e}}{r_{0_e}} - 3\right] + 1 = 6\frac{K_e(p_0 - p_f)}{\mu \varepsilon r_{0_e}^2} t \tag{9}$$

Eq. 9 gives immediately the flow front which is spherical in shape for an isotropic reinforcement. For the orthotropic case, the reciprocal transformation is applied to give the three semi-axis of the flow front which is ellipsoidal in shape.

#### **Numerical Simulation**

**Boundary Element Method.** The calculation domain is delimited by its boundary  $\Gamma = \Gamma_p \bigcup \Gamma_q$  A know pressure p (resp. normal pressure gradient q) is prescribed on the boundary surfaces  $\Gamma_p$  (resp.  $\Gamma_q$ ). Eq. 4 is multiplied by the Green function  $p^*$ . Integrated twice over the calculation domain leads to the well-known Somigiana's equation [7,8]:

$$c_i p_i + \int_{\Gamma} p q^* d\Gamma = \int_{\Gamma} p^* q \, d\Gamma \quad \text{with } c_i = \frac{\theta}{2\pi}$$
 (10)

Where  $p_i$  is the value of the pressure at a point  $M_i$  located on the boundary (which is supposed to be smooth),  $q^*$  the pressure gradient associated with  $p^*$  and  $\theta$  is the internal angle of the corner in radians. For a three-dimensional domain,  $p^*$  and  $q^*$  are given as:

$$p^* = \frac{1}{4\pi} \left(\frac{1}{r}\right) \text{ and } q^* = \frac{1}{4\pi} \frac{\vec{r} \cdot \vec{n}}{r^3}$$
 (11)

where  $\vec{r} = \overrightarrow{M_iP}$ ,  $M_i$  the point of application of the Dirac delta function and P the point under consideration. Meshing the boundary into N constant boundary elements and applying Eq. 10 gives:

$$c_i p_i + \sum_{j=1}^N p_j \int_{\Gamma_i} q^* d\Gamma = \sum_{j=1}^N q_j \int_{\Gamma_i} p^* d\Gamma$$
(12)

Introducing the  $N^2$  matrix H and G , Eq. 12 can be transformed into:

$$\sum_{i=1}^{N} p_{j} H_{ij} = \sum_{i=1}^{N} q_{j} G_{ij}$$
 (13)

where  $H_{ij} = \frac{1}{2} \delta_{ij} + \int_{\Gamma_j} q^* d\Gamma$  and  $G_{ij} = \int_{\Gamma_j} p^* d\Gamma$ . Finally Eq. 13 is reordered to take the form of a

linear system as:

$$AX = F ag{14}$$

Where X is the vector of N unknowns, A is a  $N^2$  matrix obtained by reordering H and G. F is the known vector computed from the boundary conditions, H and G matrix.

**Fluid domain updating.** The fluid domain is continuously expanding as it flows, constraining to update the resin flow front. The filling stage is regarded as a succession of quasi-steady states, so that the geometry is redefined at each time step. Unknown pressures p and normal pressure gradients q are calculated using Eq. 14. Nodes on the moving boundary are relocated using a backward Euler integration scheme:

$$\vec{x}(t + \Delta t) = \vec{x}(t) + \Delta t(\vec{v}.\vec{n})\vec{n} \tag{15}$$

Combining Eq. 15 and Eq. 4 yields to:

$$\vec{x}(t + \Delta t) = \vec{x}(t) + \Delta t \left(K_e \frac{q}{\mu} \vec{n}\right) \tag{16}$$

The procedure updates the resin flow front, but induces confined mesh degeneration. To overcome this difficulty, an additional adjustment is performed relocating the nodes in contact with the mold.

**Algorithm.** The implemented algorithm is summarized below:

- Initial computation domain is meshed.
- Stationary boundary conditions p or q are applied on each BEM element.
- Nodes coordinates are transformed according to Eq. 3.
- Normal pressure gradients q are calculated according to Eq. 14.
- Normal pressure gradients are prescribed on each node averaging values from surrounding elements.
- Nodes on the moving boundary are relocated according to Eq. 16. Nodes in contact with the upper mold wall are relocated according to a homogenization procedure.
- New normal pressures p and normal pressure gradients q are calculated next steps. The three last steps are repeated until the allotted mold-filling time is reached.
- Nodes coordinates are transformed according to the reciprocal transformation.

#### Validation of numerical simulation

**Boundary conditions.** The simulation consists in propagating the resin front into an anisotropic reinforcement characterized by its permeability tensor [K] and its porosity  $\mathcal{E}$ . A constant pressure  $p_0$  is prescribed at the spherical inlet of radius  $r_0$ . On the flow front, the pressure  $p_f$  is constant and set by the vacuum pump. Processing and rheological parameters assigned are referenced in Table 1. They keep close to the LRI process conditions but the distribution medium is not taken into account as a first approach.

Fig. 1 shows the boundary conditions prescribed on the boundary of the initial fluid domain. The non penetration condition is assumed assigning  $\overline{q} = 0$  at the upper mold wall.

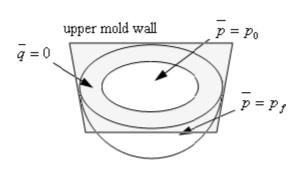


Figure 1: Boundary conditions

[K] (10 <sup>-11</sup>	$m^2$ ) $\mu$	(Pa.s)	${\cal E}$
$\begin{bmatrix} 4 & 0 & 0 \end{bmatrix}$	) ]	0.1	0.5
$\begin{bmatrix} 4 & 0 & 0 \\ 0 & 2 & 0 \\ 0 & 0 & 1 \end{bmatrix}$	)		
$\begin{bmatrix} 0 & 0 & 1 \end{bmatrix}$			
$p_0$ (Pa)	$p_f$ (Pa	$r_0$	(m)
1.10 <sup>-5</sup>	-1.10 <sup>-5</sup>	3	3.10 <sup>-2</sup>

Table 1: Processing parameters

Comparison with analytical solution. For symmetry reasons, the simulation is performed on a quarter domain meshed using 540 elements as shown in Figure 2. The step time is set at  $10^{-2}$  s. CPU time for the complete simulation is around one hour on Intel® Core<sup>TM</sup>2 Duo CPU (2.26 GHz, 1.93 Go of RAM) using Matlab® environment. The computed flow fronts are in fair agreement with the analytical one as shown in Figure 3. At 10 s elapsed time, the deviations between the numerical and analytical flow fronts positions do not exceed 1%. Prolongating the simulation until 60 s gives deviations around 3%, due to the gradual mesh degeneration as the iterations go along. Figure 4 illustrates the flow front evolution during the filling stage, comparing the front shapes at t=0.5 s and t=10s.

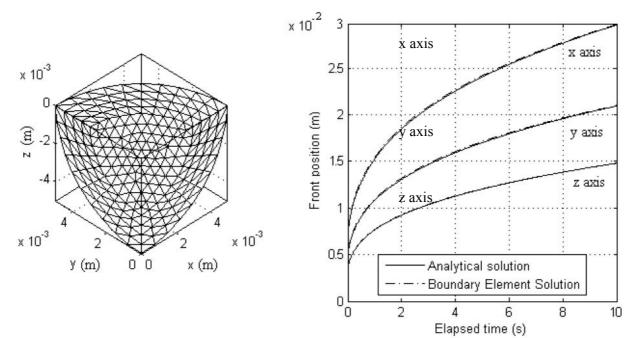


Figure 2: Initial mesh domain

Figure 3: Comparison with analytical solution

The initial mesh generation is performed on the equivalent isotropic domain. In order to assess the mesh effect on the numerical method, some simulations have been carried out meshing the physical domain. Deviations are all the more important that the permeability values are different. This is due to the transformation from Eq. 3 which stretches the mesh, degenerating elements before any calculation.

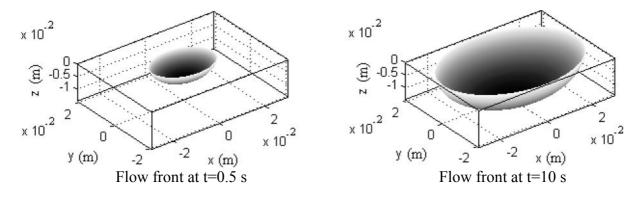


Figure 4: Flow front computed at t=0.5 s and t=10 s

#### Conclusion

The 3D radial front propagation is accurately simulated. The proposed technique is able to give a high accuracy of the front position, making it suitable for front tracking in permeability measurement [9]. It is a great advantage comparing to Volume Of Fluid techniques.

However, the proposed method is not directly applicable to complex shapes, because there is no contact algorithm implemented. In order to perform it, current works [10] focus on fixed grid approach, more suitable to handle contact (in particular merging and dividing fronts).

Future works will include hydro-mechanical couplings occurring during Liquid Composite Molding process, such as fibrous perform deformations.

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