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# The jet–ISM interactions in IC 5063

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## ABSTRACT

The interstellar medium of the radio galaxy IC 5063 is highly perturbed by an AGN jet expanding in the gaseous disc of the galaxy. We model this interaction with relativistic hydrodynamic simulations and multiphase initial conditions for the interstellar medium and compare the results with recent observations. As the jets flood through the intercloud channels of the disc, they ablate, accelerate, and disperse clouds to velocities exceeding  $400 \text{ km s}^{-1}$ . Clouds are also destroyed or displaced in bulk from the central regions of the galaxy. Our models with jet powers of  $10^{44}$  and  $10^{45} \text{ erg s}^{-1}$  are capable of reproducing many of the observed features in the position velocity diagram of IC 5063, and confirm the notion that the jet is responsible for the strongly perturbed gas dynamics seen in the ionized, neutral, and molecular gas phases. In our simulations, we also see strong venting of the jet plasma perpendicular to the disc, which entrains clumps and diffuse filaments into the halo of the galaxy. Our simulations are the first 3D hydrodynamic simulations of the jet and interstellar matter of IC 5063.

**Key words:** hydrodynamics – methods: numerical – galaxies: ISM – galaxies: jets.

## 1 INTRODUCTION

Studying the interaction of AGN jets with the interstellar matter (ISM) of gas-rich galaxies can shed light on how galaxies and black holes co-evolve (Wagner & Bicknell 2011). Many low- and high-redshift radio galaxies show highly disturbed gas kinematics and outflows (Morganti, Tadhunter & Oosterloo 2005; García-Burillo et al. 2014; Collet et al. 2015; Nesvadba et al. 2017). The outflows are sometimes found to be dominated in mass by the molecular phase (Dasyra & Combes 2012). When resolved, these outflows often appear aligned with the radio jet (Nesvadba et al. 2008; Mahony et al. 2013), which suggests that the jet is directly responsible for the turbulent dispersion of the ISM and acceleration of outflows.

IC 5063 is Type 2 Seyfert in a nearby elliptical galaxy ( $z = 0.0113$ ) with powerful ( $P_{1.4\text{GHz}} = 3 \times 10^{23} \text{ W Hz}^{-1}$ ) extended radio structures, which exhibit strong interactions with the ISM (Morganti, Oosterloo & Tsvetanov 1998; Morganti et al. 2007; Tadhunter et al. 2014). The galaxy contains a gaseous disc with large-scale dust lanes (see fig. 1 of Oosterloo et al. 2017), possibly resulting from a merger with a gas-rich galaxy (Morganti et al. 1998) in the past. The galaxy is seen almost edge on in the sky and the jet appears

to be aligned parallel to the disc. This is not an uncommon feature of Seyfert galaxies, in which kiloparsec-scale radio structures are found to be randomly oriented with respect to the galaxy's major axis (Gallimore et al. 2006) or misaligned with the central gas disc (Cecil, Bland-Hawthorn & Veilleux 2002; Dopita et al. 2015).

IC 5063 contains a reservoir of molecular gas exceeding  $10^9 M_{\odot}$  (Morganti et al. 2015). Within the region containing the strong radio emission, IC 5063 exhibits kinematically perturbed gas in the warm ionized phase (Morganti et al. 2007; Dasyra et al. 2015), neutral (H I) phase (Morganti et al. 1998; Oosterloo et al. 2000), warm molecular phase (Tadhunter et al. 2014), and cold molecular phase recently observed with ALMA by Morganti et al. (2015, hereafter M15) and Dasyra et al. (2016). On the whole, the strong velocity dispersions of  $\gtrsim 500 \text{ km s}^{-1}$  seen in the various phases appear to be caused by the jet and the energy bubble it drives, with a particularly strong interaction with the ISM occurring in the western hotspot, about 500 pc from the core.<sup>1</sup>

The outflowing molecular gas was found to be optically thin, and the mass estimate of outflowing molecular gas is at least  $1.3 \times 10^6 M_{\odot}$  (Oosterloo et al. 2017). In comparison, the mass of the neutral H I outflow associated with the western hotspot is

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<sup>1</sup> We use 1 arcsec = 232 pc, as in Oosterloo et al. (2017).

approximately  $3.6 \times 10^6 M_{\odot}$  (Morganti et al. 2007). The mass outflow rate at the western hotspot in molecular gas is estimated to be at least  $4 M_{\odot} \text{ yr}^{-1}$  (Oosterloo et al. 2017).

It is not clear whether the outflowing molecular gas is mostly driven from the disc (Scannapieco 2017), or whether a substantial fraction of the molecular emission in outflow comes from molecules formed in outflowing diffuse gas, which is cooling rapidly (Richings & Faucher-Giguere 2017). Observations of the near-infrared  $\text{H}_2$  lines imply an outflowing  $\text{H}_2$  mass at the western hotspot of only approximately  $8 \times 10^2 M_{\odot}$ , suggesting that the warm molecular gas component is only a transitory phase in the process of gas cooling behind strong radiative shocks driven into the clouds accelerated by the jets (Tadhunter et al. 2014). The detailed mechanism of molecule formation in outflows is, however, not well understood.

In this paper, we present idealized relativistic hydrodynamic simulations of the inner kiloparsec region of IC 5063, concentrating on the kinematic signatures resulting from the interaction of the relativistic jets with the inhomogeneous ISM. The parameters of the simulations are chosen based on observational constraints outlined in M15.

In Section 2, we summarize the main features and interpretations of the kinematic signatures presented through the position-velocity (PV) diagram in fig. 3 in M15. We describe the equations that are solved, the numerical methods, and set-up in Section 3 and the parameters of the simulations in Section 4. The results are presented in Section 5, in which we first describe the evolution of the jet and the ISM and then focus on a comparison of the observed PV diagram (M15) with synthetic PV diagrams computed from our simulations. Additional results from our simulations are noted in a third subsection. We provide a summary and short discussion in Section 6.

## 2 KINEMATIC SIGNATURES IN CO (2–1)

The principal aim of the work presented in this paper is to investigate whether jet–ISM interactions are responsible for the enhanced velocity dispersion observed in the ISM of IC 5063. To this end, we pay particular attention to the PV diagram presented in fig. 3 of M15. The PV diagram shows several characteristics indicative of strong jet–ISM interactions.

(i) The dispersions in velocity space in the regions impacted by the jets ( $\pm 2.5$  arcsec) are broadened to a few  $100 \text{ km s}^{-1}$ . The unperturbed gas at the outer edges has a much lower dispersion ( $\sim 40 \text{ km s}^{-1}$ ).

(ii) The broadening across the region impacted by the jet is not smooth. A number of spiky features are seen. A particularly strong feature is the dispersion in the western part that coincides with the radio hotspot.

(iii) The broadening across the region impacted by the jet is not symmetric about the centre of the galaxy. The western side appears to be more strongly dispersed than the eastern side.

(iv) There is substantial emission in the regions away from the mean rotation curve (forbidden quadrants) implying strong gas motions in directions against the rotation of the disc.

(v) The gradient of the rotation profile in the region impacted by the jet appears to be shallower than the gradient expected from the gravitational potential (white line in fig. 3 of M15).

M15 outlined a model, sketched in their fig. 7, to explain some of the features described above. In their model, the jet–ISM interactions displaced clouds that were located near the centre of the galaxy to larger disc radii, leading to clouds rotating more slowly

with respect to the unperturbed gas. This model explained the first and last points in the above list, but could not directly explain the other points.

The asymmetry of the gas kinematics and the asymmetric radio morphology hint at a clumpy gas distribution in the disc (Gaibler, Khochfar & Krause 2011). Theoretical papers investigating the interactions of outflows with a turbulent gas distribution (Sutherland & Bicknell 2007; Gaibler, Khochfar & Krause 2011; Wagner & Bicknell 2011; Wagner, Bicknell & Umemura 2012; Wagner, Umemura & Bicknell 2013; Mukherjee et al. 2016) have shown that modelling the ISM as an inhomogeneous medium is crucial for correctly modelling observational characteristics and calculating feedback efficiencies.

## 3 NUMERICAL SET-UP

We use the open-source software `PLUTO` (Mignone et al. 2007) version 4.1 to numerically solve the equations of relativistic hydrodynamics, using an Eulerian finite-volume Godunov-type scheme as described in Mignone & Bodo (2005). We also solve two additional advection equations, which trace jet plasma and dense clouds, respectively, in our simulations. The code and integration schemes used are identical to those used by Mukherjee et al. (2016). We employ the third-order Runge–Kutta time-stepping scheme together with the piecewise parabolic method of spatial interpolation.

The temperature  $T$  of the gas is given by the ideal gas law,  $p = \rho k_B T / \bar{\mu}$ , where  $k_B$  is Boltzmann’s constant, and  $\bar{\mu}$  is the mean mass per particle. Atomic cooling is implemented using a tabulated non-equilibrium radiative cooling function  $\Lambda(T)$  for an optically thin plasma generated by the `MAPPINGS V` code (Sutherland & Dopita 2017) for solar metallicity gas. The cooling is dominated by emission from collisionally ionized atoms.

Between the cooling floor of 100 K and the peak of the cooling curve at just above  $10^5$  K, the mean mass per particle decreases monotonically from  $\bar{\mu} \approx 1.24$  to  $0.604$  u. The value of  $\bar{\mu}$  for a given temperature is also obtained from the pre-computed `Mappings V` cooling function, in which  $\mu$  is tabulated as a function of  $p/\rho$ . At relativistic pressures, a relativistic correction to bremsstrahlung as implemented by Krause & Alexander (2007) is employed.

Molecular cooling is not included in this work, and the cooling below  $10^4$  K down to the cooling floor of 100 K is by metal line cooling alone. In our simulations, the gas in the dense cores is quickly driven to the cooling floor even with purely atomic cooling, so any additional effects due to cooling by molecules do not influence the dynamics of the gas in this work.

## 4 SIMULATION SET-UP AND CHOICE OF PARAMETERS

The simulations of the jet–ISM interactions in IC 5063 are idealized isolated galaxy simulations. The focus of the simulations is on the interaction of the jet plasma with the hot, warm, and cold gas phases in the ISM.

A cubical region of 2 kpc width around the central black hole is initialized with a jet nozzle at the centre and a two-phase ISM representing the hot phase and warm to cold phase in IC 5063. We use a Cartesian coordinate system ( $X$ – $Y$ – $Z$ ) in which the galaxy disc is co-planar with the  $X$ – $Y$  plane; the  $Z$ -axis is the symmetry axis of the disc. The jet is aligned in the plane of the disc along the  $X$ -axis. The computational domain is divided into a grid of  $512^3$  uniform cells, providing a spatial resolution of approximately 4 pc per cell.

**Table 1.** Simulation parameters.

Parameter	Description	Value	
Jet parameters			
$P_{\text{jet}}$	Jet power	$10^{44} \text{ erg s}^{-1}$ (P44)	$10^{45} \text{ erg s}^{-1}$ (P45)
$\Gamma$	Jet bulk Lorentz factor at nozzle	4	6
$\chi$	Jet proper density parameter <sup>a</sup>	4.8	173
$r_{\text{jet}}$	Jet radius	25 pc	
$\theta_{\text{jet}}$	Jet opening angle	20°	
Potential & hot halo			
$M$	Total mass in Hernquist profile	$1.67 \times 10^{11} M_{\odot}$	
$a$	Scale length of Hernquist profile	2.8 kpc	
$\rho_{\text{h}}(0)$	Central density of hot atmosphere	$0.1 \text{ cm}^{-3}$	
$T_{\text{h}}$	Temperature of hot atmosphere	$6 \times 10^6 \text{ K}$	
Clouds and disc			
$T_{\text{c}}$	Critical temperature for warm phase	$3 \times 10^4 \text{ K}$	
$\bar{\rho}(0, 0)$	Mean central density of clouds	$200 \text{ cm}^{-3}$	
$\sigma$	Combined turbulent and thermal dispersion	$40 \text{ km s}^{-1}$	
$e_{\text{K}}$	Ratio of azimuthal speed to Keplerian speed	0.9	

<sup>a</sup>Ratio of jet rest mass energy to enthalpy (equation 8).

The duration of the simulations is set by time required for the jet head to advance to at least 0.5 kpc from the centre of the galaxy. This time-scale is of the order of 1 Myr (see Appendix A). We can therefore ignore the self-gravity of the warm phase gas and the effects of star formation and supernova feedback in these simulations. We also ignore the effects of magnetic fields, which are expected to make only a small difference in terms of the energy transfer efficiency from jets to the ISM (Asahina, Nomura & Ohsuga 2017).

The aim of the simulations is to provide a fully hydrodynamic model that qualitatively, and to some extent quantitatively, explains the observations of IC 5063. We have therefore not performed a costly parameter space study, but have chosen parameters informed from observations. We present two models with differing jet powers. The parameters for the simulations are listed in Table 1. In the following section, we explain in more detail the set-up of the simulations and the parameters chosen for the gravitational potential, the different phases of the ISM, and the jet.

#### 4.1 Gravitational potential and initial gas disc

In modelling the gravitational potential of the galaxy, we do not need to take the gravitational field of the black hole into account since, for a mass  $M_{\text{bh}} \approx 2.8 \times 10^8 M_{\odot}$  (Nicastro, Martocchia & Matt 2003) and stellar velocity dispersion  $\sigma \approx 160 \text{ km s}^{-1}$  (Woo & Urry 2002), the radius of influence of the black hole is  $\sim GM_{\text{bh}}/\sigma^2 \approx 46 \text{ pc}$ , comparable to the radius of injection of the jet ( $\sim 40 \text{ pc}$ ) in the simulation. For a comparison with the jet powers employed in the simulations, the Eddington luminosity of the black hole is important and this is approximately  $3.5 \times 10^{46} \text{ erg s}^{-1}$ .

We use the analytic Hernquist potential (Hernquist 1990) to model the gravitational potential of IC 5063. The parameters of this potential are the total mass,  $M$ , and the scale length,  $a$ , where the mass,  $M(r)$ , and the gravitational potential,  $\Phi(r)$ , as functions of spherical radius,  $r$ , are given by

$$M(r) = M \frac{r^2}{(r+a)^2} \quad (1)$$

$$\Phi(r) = -\frac{GM}{r+a}. \quad (2)$$

We use the inferred effective radius  $R_{\text{e}} = 21.5 \text{ arcsec} = 5139 \text{ pc}$  (Kulkarni et al. 1998; M15) and the central velocity dispersion  $160 \text{ km s}^{-1}$  to evaluate the parameters  $M$  and  $a$ , as follows. Since the gravitational potential is most likely dominated by baryonic matter within the effective radius, we use the Hernquist (1990) relation between half-light radius and scale length, viz,  $R_{\text{e}} = 1.8153 a$ , to estimate  $a \approx 11.8 \text{ arcsec} \approx 2.8 \text{ kpc}$ .

In the Hernquist model, the projected isotropic velocity dispersion rises rapidly to a maximum  $\approx 0.319 (GM/a)^{1/2}$  at a projected radius  $x \approx 0.224 a \approx 2 \text{ arcsec} \approx 477 \text{ pc}$  for IC 5063. We equate the maximum velocity dispersion with the observed central value of  $160 \text{ km s}^{-1}$  to obtain a mass  $M \approx 1.67 \times 10^{11} M_{\odot}$ .

These estimates can be checked for consistency using the observed major axis H I rotational velocity beyond the region of influence of the jet. Let  $\theta \approx 74^\circ$  be the angle between the disc normal and the line of sight (M15). Then, in the Hernquist model, the observed rotational velocity is given by

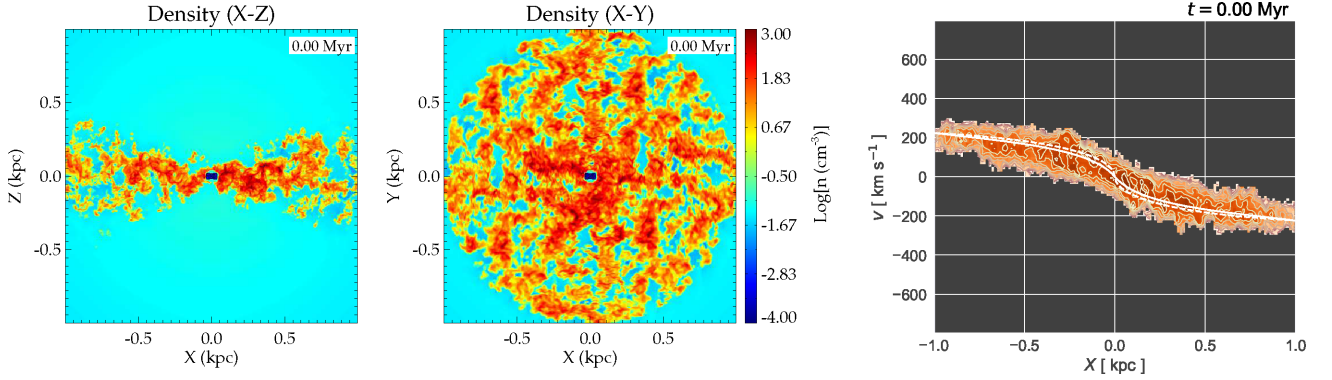
$$v_{\text{rot}} = \pm \left( \frac{GM}{a} \right)^{1/2} \frac{(r/a)^{1/2}}{1+r/a} \sin \theta. \quad (3)$$

The region just beyond the influence of the jet is at an angular distance of 4 arcsec, implying an observed rotational velocity  $\approx \pm 209 \text{ km s}^{-1}$ . This compares favourably with the observed speed of  $\pm 190 \text{ km s}^{-1}$  (M15). The initial distribution of gas in the galaxy is defined by a hot isothermal atmosphere in hydrostatic equilibrium. Assuming the central density at  $r = 0$  to be  $\rho_{\text{h}}(0)$ , for gas with temperature  $T_{\text{h}}$ , mean molecular weight  $\mu = 0.6$ , and atomic mass unit  $m_{\text{a}}$ , the halo density profile is given by

$$\rho_{\text{h}} = \rho_{\text{h}}(0) \left[ \exp \left( \frac{-\phi(r)}{k_{\text{B}} T_{\text{h}} / (\mu m_{\text{a}})} \right) \right]. \quad (4)$$

The temperature of the hot gas is initialized to  $T_{\text{h}} = 6 \times 10^6 \text{ K}$  (as in Table 1). Superimposed on the halo is a turbulent disc prescribed using the approach described in Sutherland & Bicknell (2007), which is summarized below. A number of cells in the computational volume are replaced by denser, cooler gas using the following approach.

Let the cool gas density in cylindrical coordinates be  $\rho(R, z)$ , with  $\rho(0, 0)$  being the value at the centre of the galaxy,  $\sigma_{\text{t}}$  the turbulent velocity dispersion, and  $T$  the mean temperature. The azimuthal velocity of the disc is defined to be a constant fraction,  $e_{\text{K}}$  of the



**Figure 1.** Density ( $\log(n[\text{cm}^{-3}])$ ) in the  $X$ - $Z$  (left) and  $X$ - $Y$  (centre) planes at  $t = 0$ . Right: synthetic PV diagram based on CO (2–1) emission at  $t = 0$  (based on equation 9). The dashed and solid lines trace the circular velocity according to equation (3) and the projected circular velocity weighted by the mean density profile of the turbulent disc given by equation (5).

Keplerian velocity. Then the *mean* density of the disc is given by

$$\frac{\bar{\rho}(R, z)}{\bar{\rho}(0, 0)} = \exp - \left[ \frac{\Phi(R, z) - e_K^2 \Phi(R, 0) - (1 - e_K^2) \Phi(0, 0)}{\sigma^2} \right], \quad (5)$$

where

$$\sigma^2 = \sigma_t^2 + \frac{k\tilde{T}}{\bar{\mu}}. \quad (6)$$

The distribution of dense gas is constructed by multiplying a fractal cube by the mean density described by equation (5). The fractal cube consists of a lognormal density distribution with unit mean and a power-law spectrum in Fourier space. The temperature of the gas is defined by pressure equilibrium with the hot atmosphere. This gas distribution then replaces a number of cells in the hot atmosphere, except when the temperature in a cell exceeds a critical value  $T_c$ , corresponding to a low density  $\rho_c = \bar{\mu} p_{\text{hot}}/kT_c$ . While the gas temperature, in principle, contributes to the overall dispersion parameter  $\sigma$ , in these simulations the contribution is very small since  $\tilde{T} \sim 100$  K because of the cooling floor implemented. Hence, we define the cool gas density by neglecting the temperature contribution and assigning the parameter  $\sigma_t$  alone.

A Gaussian turbulent velocity field, with velocity dispersion  $\sigma$ , is imposed on the dense gas. This is in addition to the Sutherland & Bicknell (2007) Ansatz and was first implemented in Mukherjee et al. (2016). The parameters of the distribution are summarized in Table 1 together with the other parameters of the simulation. The initial conditions are shown by two orthogonal mid-plane density slices through the simulation box at time  $t = 0$  and the PV diagram of the rotating disc (line of sight along  $Y$ ) in Fig. 1.

## 4.2 The relativistic jet

We present two models which have identical set-ups and parameters, except for the jet power. The jet power in model P44 is  $10^{44}$  erg s<sup>-1</sup>, and that in model P45 is  $10^{45}$  erg s<sup>-1</sup>. The Eddington ratios of these jets are approximately 0.015 and 0.15, respectively.

The choice of jet powers is partly motivated by our previous studies of jet–ISM interactions (Wagner & Bicknell 2011), in which only jets with powers exceeding  $10^{44}$  erg s<sup>-1</sup> were found to be capable of powering velocity dispersions greater than  $\sim 400$  km s<sup>-1</sup>. Using the scaling relation by Cavagnolo et al. (2010) with  $P_{1.4\text{GHz}} = 3 \times 10^{23} \text{ W Hz}^{-1}$  gives a jet power of  $9 \times 10^{43}$  erg s<sup>-1</sup>;

the scaling relation by Willott et al. (1999) yields  $5 \times 10^{43}$  erg s<sup>-1</sup> (M15). We discuss these estimates further in Section 6.

The relations between the jet power and proper density to the jet pressure,  $p_{\text{jet}}$ , internal energy,  $\varepsilon_{\text{jet}}$ , density,  $\rho_{\text{jet}}$ , bulk Lorentz factor,  $\Gamma$ , and velocity,  $\beta_{\text{jet}}$ , are

$$P_{\text{jet}} = \frac{\gamma}{\gamma - 1} p_{\text{jet}} \Gamma^2 c \beta_{\text{jet}} A_{\text{jet}} \left( 1 + \frac{\Gamma - 1}{\Gamma} \chi \right) \quad (7)$$

$$\chi = \frac{\rho_{\text{jet}} c^2}{\varepsilon_{\text{jet}} + p_{\text{jet}}} = \frac{(\gamma - 1) \rho_{\text{jet}} c^2}{\gamma p_{\text{jet}}}. \quad (8)$$

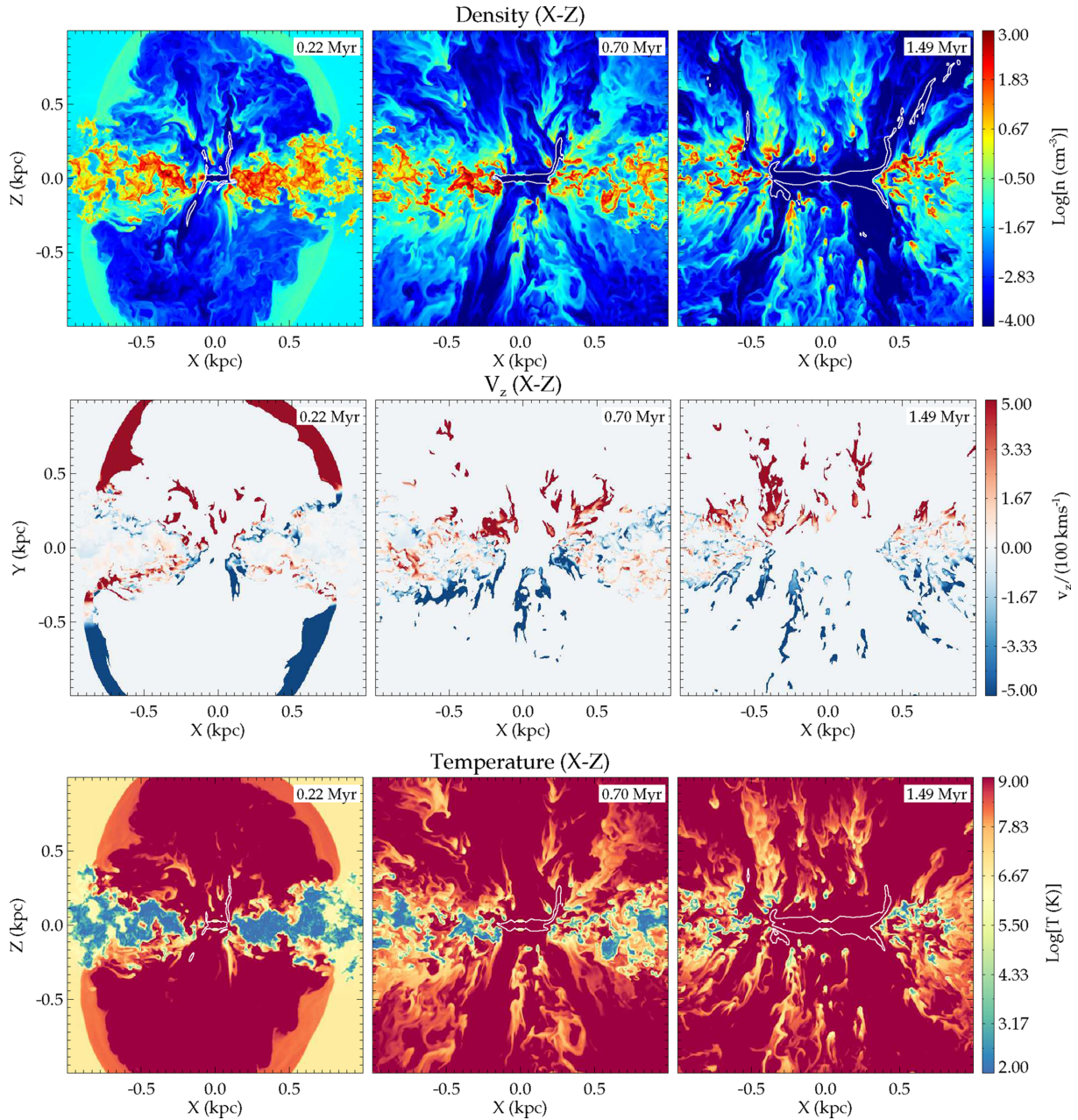
The adiabatic index is assumed to be  $\gamma = 5/3$  for the jet plasma, corresponding to that of an ideal equation of state. Such an assumption is appropriate for  $\chi > 1$  (Mignone & McKinney 2007), as used in our simulation. An ideal equation of state is also a better description for the dense thermal gas included in the non-homogeneous ISM. More on this has been discussed in section 2.3 of Mukherjee et al. (2016). The parameters for the jet are included in Table 1.

A small spherical region with a radius of 40 pc around the central black hole is clear of any warm phase gas and is set as a fixed internal boundary from which the jet emerges conically and bidirectionally with an opening angle of 20°. The jet radius at the inlet is 25 pc. These initial conditions are prescribed to represent a freely expanding conical jet.

The orientation of the jet with respect to the line of sight is not well constrained observationally, but the jet axis is thought to lie close to the plane of the disc and perpendicular to the line of sight. For simplicity, we orient the jet exactly along the disc axis. The evolution of the jet is dominated by its interaction with the warm phase of the ISM and the random and fractal clumpiness of the warm phase ensures that this choice of jet orientation does not lead to a vastly different solution than cases where the jet is slightly inclined with respect to the disc.

## 5 RESULTS

We first describe the general features of the evolution of the jet and the ISM in the simulations. We then compare the kinematics of the jet-perturbed gas in the disc as seen in our simulations with the CO (2–1) observations of M15.



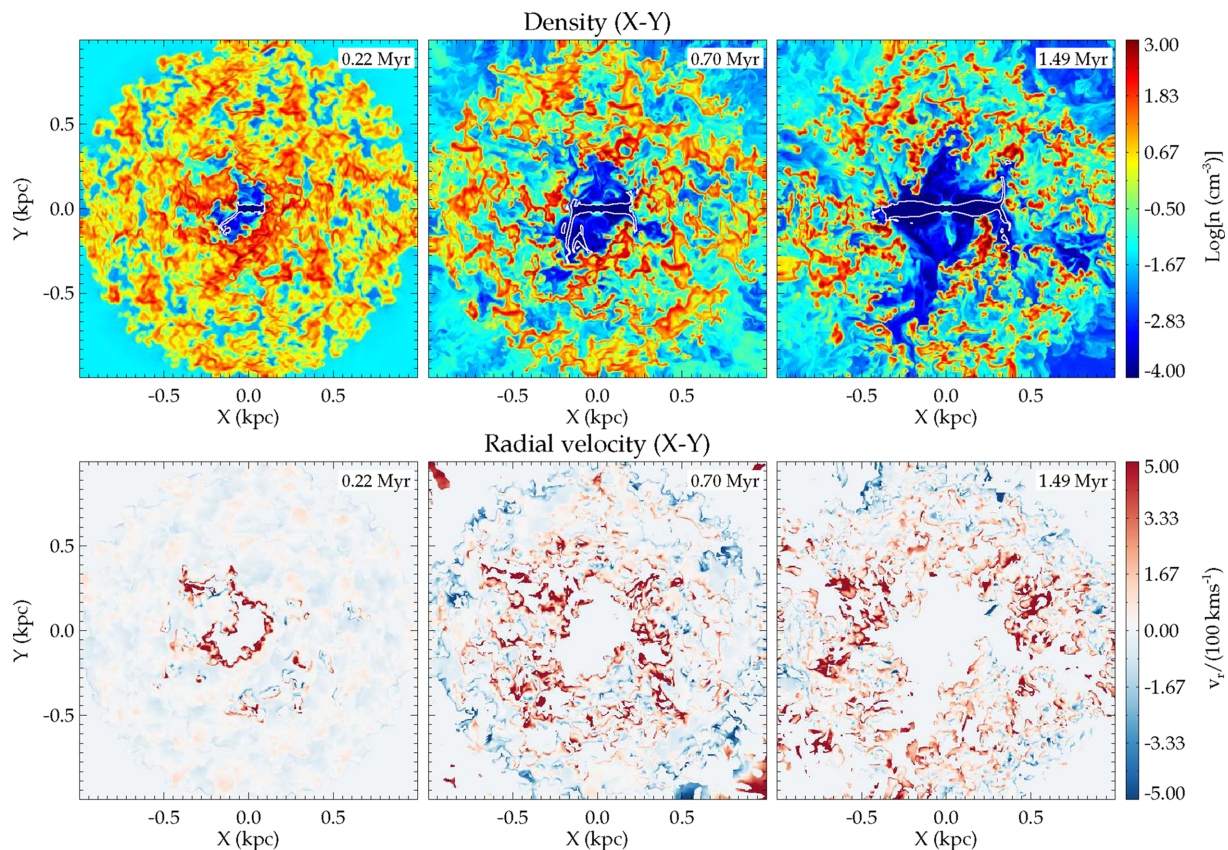
**Figure 2.** Mid-plane slices in the  $X$ - $Z$  plane ( $Y = 0$ ) of the evolution of simulation P44, for which the jet power is  $P_{\text{jet}} = 10^{44} \text{ erg s}^{-1}$ . The last column of panels with  $t = 1.49 \text{ Myr}$  corresponds roughly to the current extent of the jets in IC 5063. The top row of panels shows the density ( $\log(n[\text{cm}^{-3}])$ ), the middle row the vertical component of the velocity ( $v_z$ ) normalized to  $100 \text{ km s}^{-1}$ , for gas with density  $n > 0.1 \text{ cm}^{-3}$ , and the bottom row shows the temperature. The white line corresponds to the contours of jet tracer with value 0.1 (maximum being 1), demarcating the region with non-thermal plasma.

### 5.1 Evolution of jet and ISM

The evolution of the simulation with a jet power of  $10^{44} \text{ erg s}^{-1}$  (P44) is shown in three snapshots in Figs 2 (side-on) and 3 (face-on), and the evolution of the simulation with a jet power of  $10^{45} \text{ erg s}^{-1}$  (P45) is shown in Figs 4 (side-on) and 5 (face-on). Each figure contains three rows of panels showing mid-plane density slices, mid-plane temperature slices, and mid-plane velocity slices, respectively. For P44, the snapshot at time  $t = 1.49 \text{ Myr}$  corresponds roughly to the observed conditions in IC 5063, based mainly on the extents of the jet streams, while for P45 it is the snapshot at time  $t = 0.24 \text{ Myr}$ . A 3D rendering of simulation P45 at  $t = 0.24 \text{ Myr}$  from different perspectives is shown in Fig. 6.

The evolution of the jet and ISM in simulations of AGN feedback in a multiphase ISM on kiloparsec galaxy scales is well documented in the literature (Saxton et al. 2005; Sutherland & Bicknell 2007; Wagner et al. 2012; Mukherjee et al. 2016; Bieri et al. 2017). As in the simulations described by, e.g. Mukherjee et al. (2016), the interaction between the jet and ISM is very strong because the interaction depth of the jet (the column density of clouds that the jets impact; Wagner et al. 2012) is very high if the jet lies in the plane of the gas-rich galactic disc.

The thermal and ram pressures in the jet plasma drive an energy bubble into the ambient ISM. The jet plasma in the bubble, being orders of magnitudes lighter than the dense phases of the



**Figure 3.** Mid-plane slices in the  $X$ – $Y$  plane ( $Z = 0$ ) of the evolution of simulation P44, for which the jet power is  $P_{\text{jet}} = 10^{44} \text{ erg s}^{-1}$ . The last column of panels for which  $t = 1.49 \text{ Myr}$  corresponds approximately to the current extent of the jets in IC 5063. The top row shows the density ( $\log[n \text{ cm}^{-3}]$ ) and the bottom row shows the cylindrically radial component of the velocity ( $v_{\text{cyl}} = \sqrt{V_x^2 + V_y^2}$ ) normalized to  $100 \text{ km s}^{-1}$ , for gas with density  $n > 0.1 \text{ cm}^{-3}$ .

ISM, percolates through paths of least resistance between clouds, a phase termed the ‘flood-and-channel’ phase (Sutherland & Bicknell 2007). In the process, the clouds are compressed but they also experience strong hydrodynamic ablation driven through the Kelvin–Helmholtz instability. The ablated material forms cometary tails and partly mixes with the intercloud flow, which consists of swept up hot-phase ISM, and jet plasma.

The clouds and filaments of ablated cloud material are long-lived because of efficient radiative cooling. Radiative cooling does not strongly affect the global evolution of the AGN bubble as the outflow remains largely energy-driven, at least on kpc scales. The dynamics of the warm and cold phases are, however, strongly affected as clouds fragment and their cores and filaments are efficiently accelerated by the surrounding ram pressure and thermal pressure gradients.

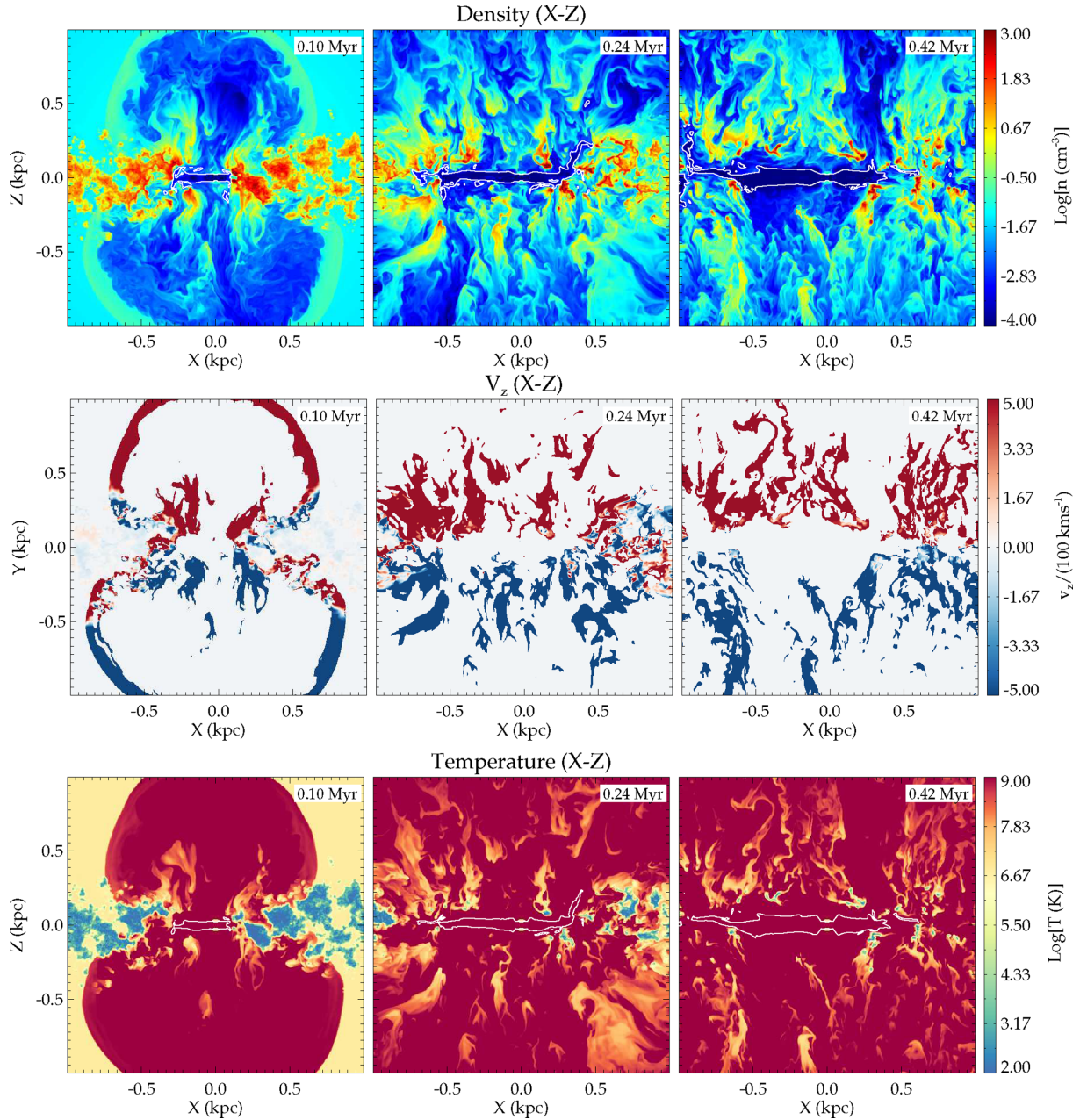
We see in Figs 2–5 that cloud cores are accelerated to  $100$ – $200 \text{ km s}^{-1}$ , while diffuse filaments exceed velocities of  $500 \text{ km s}^{-1}$ . Some cloud cores are displaced and most of the ablated gas is carried along intercloud channels to larger disc radii. In these simulations, the jet never fully breaks out of the disc within the duration of the simulations. In the face-on velocity maps of Figs 3 and 5, we see that the highest gas velocities occur near the locations of the jet head (i.e. the radio hot spots) at  $(X, Y) \approx (\pm 0.5, 0) \text{ kpc}$ . The corresponding density maps also show that this is where the strongest cloud compression occurs. We also note that, as seen in the side-on density slices of Figs 2 and 4, the jet plasma quickly sweeps away the hot halo gas, replacing it with jet plasma and entrained ISM

from the disc. The jet also excavates a cavity in the central parts of the galaxy devoid of dense gas.

The third panel in Figs 2 and 4 shows the temperature in the  $X$ – $Z$  plane. The jet–ISM interaction forms a multiphase ISM as discussed previously in Mukherjee et al. (2016, 2017). Shocks driven by the energy bubble heat the initially cold disc to  $T \gtrsim 10^5 \text{ K}$ , with shocked clouds being driven vertically as discussed earlier. The tenuous cavity filled with relativistic plasma is at a higher temperature ( $T \gtrsim 10^8 \text{ K}$ ).

The main difference between the two simulations presented here (P44 and P45) is that the higher powered jet penetrates through the disc much faster (almost 10 times faster) than the lower powered jet. As a result, the higher powered jet ends up excavating a roughly cylindrical cavity about the jet axis, rather than a spherical cavity about the central core. Of course, the resulting velocity dispersion caused by the higher powered jet is much stronger, as expected from the scaling relation of feedback efficiency and jet power (Wagner & Bicknell 2011). Accordingly, the more powerful jet in simulation P45 is able to clear its path faster and reach an extent of  $0.5 \text{ kpc}$  within  $0.24 \text{ Myr}$ , while the lower powered jet in simulation P44 requires  $1.49 \text{ Myr}$  to reach the same size. The difference in velocity dispersion between the two simulations is discussed more quantitatively in Section 5.2.

In order to compare the simulations with observations, we generate a number of synthetic radio surface brightness images. Since the simulation is purely hydrodynamic and we do not track the evolution of relativistic electrons, we estimate the synchrotron emissivity



**Figure 4.** Same as Fig. 2 but for simulation P45, for which the jet power is  $P_{\text{jet}} = 10^{45} \text{ erg s}^{-1}$ . The central column of panels for which  $t = 0.24 \text{ Myr}$  corresponds approximately to the current extent of the jets in IC 5063.

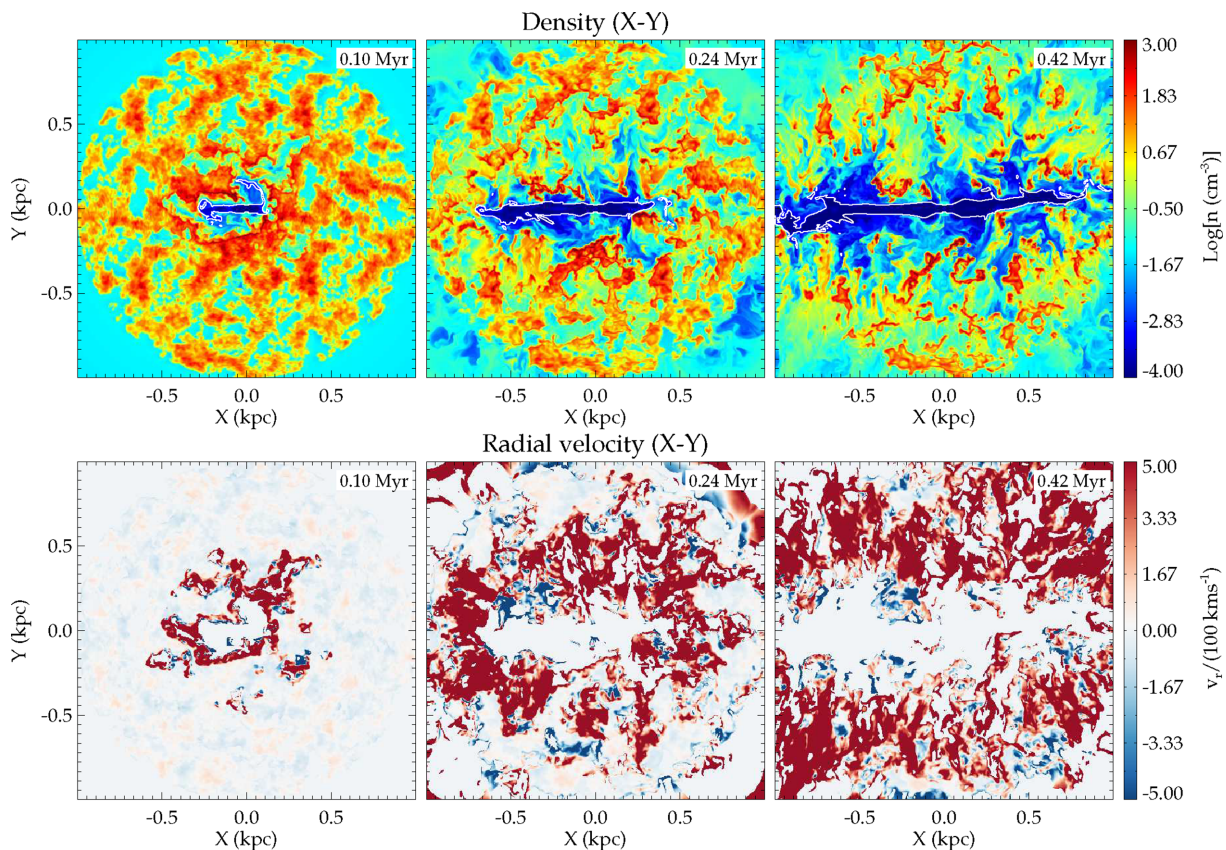
by assuming that the energy densities of the magnetic field and relativistic electrons are fixed fractions, respectively,  $f_e$  and  $f_B$ , of the internal energy density. We adopt nominal values of  $f_e = f_B = 0.1$  but the resulting images can be readily scaled to other values. The surface brightness is calculated by integrating the radiative transfer equation through the volume, taking into account free-free absorption by the dense clouds. The method of generating the synthetic images is presented in more detail in Bicknell et al. (2018).

In Fig. 7, we show a synthetic synchrotron surface brightness image at an observed frequency of 8 GHz generated from our simulation at  $t = 704 \text{ kyr}$ , corresponding approximately to the observed extent of the radio jet. The left-hand panel in the figure shows the surface brightness at the resolution of our simulation, and the right-hand panel shows the surface brightness convolved with a beam of

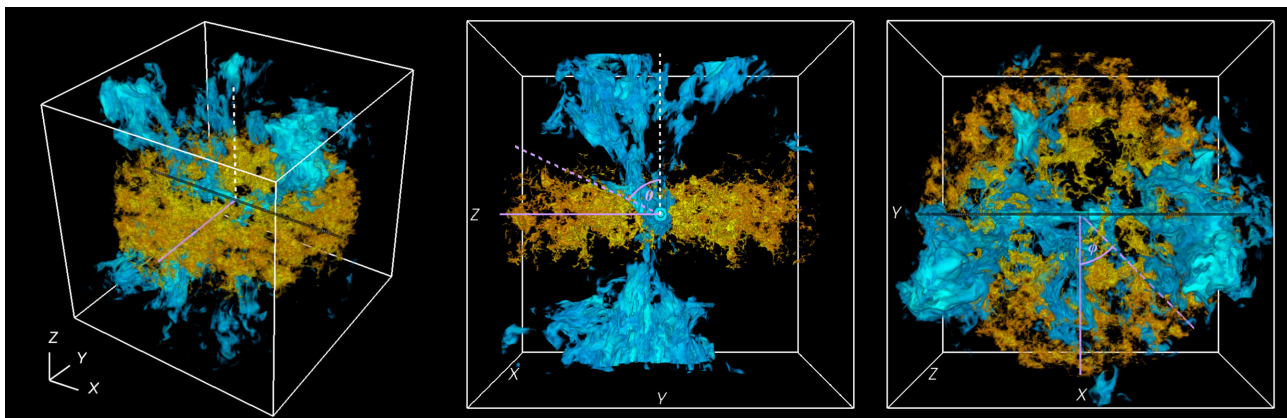
size  $234 \text{ pc} \times 234 \text{ pc}$ , which corresponds to the resolution of the image at 8 GHz in Morganti et al. (1998). Our synthetic radio image resembles the observed radio image, in particular at the western ‘hotspot’, where the radio surface brightness is particularly high due to the strong interaction of the jet with a large dense cloud. Although the distribution of the dense gas in the disc is statistically symmetric, the asymmetric features arise from interactions with the local fluctuations in the density field resulting in the clumpy inhomogeneous ISM. Note that there is no core component because we do not model the subparsec scale nuclear region of the galaxy.

We emphasize that the radio morphology at the observed resolution belies the complex distribution of the jet plasma in the galaxy, which bears little resemblance to classical radio sources. For





**Figure 5.** Same as Fig. 3 but for simulation P45, for which the jet power is  $P_{\text{jet}} = 10^{45} \text{ erg s}^{-1}$ . The central column of panels for which  $t = 0.24 \text{ Myr}$  corresponds roughly to the current extent of the jets in IC 5063.

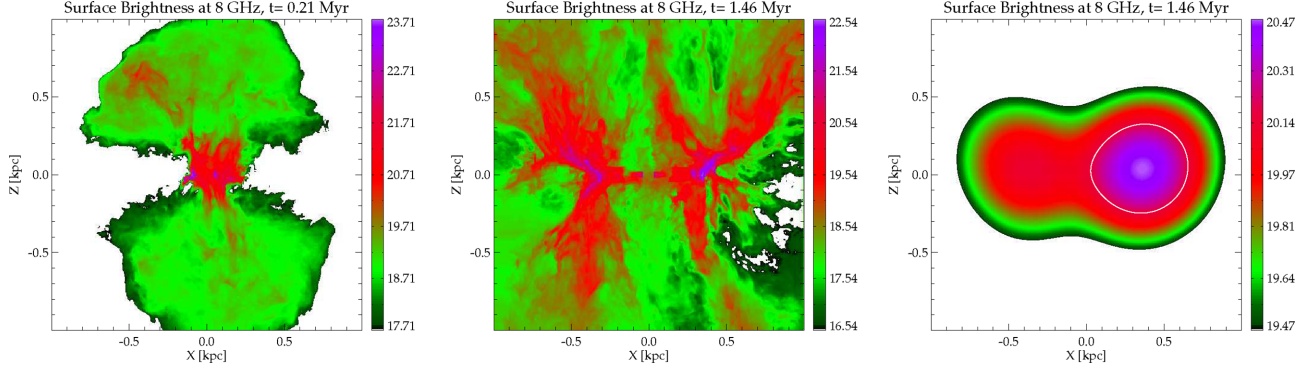


**Figure 6.** Volume render of the simulation snapshot at time  $t = 0.24 \text{ Myr}$  of simulation P45. The yellow–orange material represents dense disc gas and the blue material the jet plasma. Diffuse warm and hot gas and gas mixed with jet plasma are not shown as these phases have a large volume filling factor. The grey dotted line traces the jet axis, the white dashed line is the rotational symmetry axis of the disc, and the purple dashed line is the default view normal used in most of the plots in the paper. Left: a perspective from an angle. Center: a side-on view along the jet axis, with the front half of the galaxy clipped away. The purple dashed line shows a line of sight that is tilted by an angle  $\theta$  with respect to the disc normal, but is perpendicular to the jet axis. Right: a face-on view on to the disc along the  $Z$ -axis. The purple dashed line here represents a line of sight along the plane of the disc, but rotated by an angle  $\phi$  from the default view normal.

example, the jets, as they propagate through the disc, do not develop well-defined lobes that arise when overpressured jet plasma expands against a smooth ambient medium. The regions of bright radio emission are not classical hotspots in the sense of knots or terminal shocks, but are best described as a ‘splatter region’, where the jet head interacts strongly with a dense cloud in its path (see

fig. 5 in Oosterloo et al. 2000), splitting and splattering the main jet stream in all directions.

Because the jet is so light compared to the clouds, it is easily deflected (or even split), as we see in the white contours of the western jet in the panels in Fig. 2 for  $t = 1.49$  and  $2.10 \text{ Myr}$ . Throughout the simulation, the jet plasma also escapes perpendicular to the disc,



**Figure 7.** Synthetic 8 GHz surface brightness from simulation P44 taking into account free–free absorption. The ray tracing is performed with rays parallel to the  $Y$ -axis, in the direction of  $-y$ . Left: unconvolved image at the resolution of the simulation at  $t = 0.21$  Myr. Center: unconvolved image at the resolution of the simulation at  $t = 1.49$  Myr. Right: the image convolved with a beam of size  $234 \text{ pc} \times 234 \text{ pc}$ . The white line denotes the contour for 10 per cent of peak brightness. The power at 8 GHz of the left and middle panels are  $6.6 \times 10^{21}$  and  $2.2 \times 10^{21} \text{ W Hz}^{-1}$ .

entraining with it cloudlets that are accelerated to several  $100 \text{ km s}^{-1}$ . Although most of the mass of the cold and warm gas remains contained within the disc, the acceleration of ablated gas is also efficient perpendicular to the disc. This effect is a somewhat serendipitous discovery, to which we return in Section 5.3.

## 5.2 PV diagrams

From the velocity field in the simulations, we construct synthetic PV diagrams and compare them to the PV diagrams of CO (2–1) published by M15. In our simulations, we do not directly track the molecular gas, nor do we follow the chemistry of molecule formation, so that we utilize temperature, cloud tracer, and density thresholds to infer the possible presence of molecular gas and to calculate the emissivity.

Let  $\phi_w$  be the cloud tracer variable ( $0 \leq \phi_w \leq 1$ ),  $n_{\text{crit}}$  the critical density for a molecular transition, and  $n$  and  $f_{\text{mol}}$  be the particle number density and the molecule fraction by mass in a cell, respectively. We take the emissivity,  $\epsilon$ , to be

$$\epsilon = \begin{cases} A\phi_w^2 n^2 & \text{if } n < n_{\text{crit}} \text{ and } f_{\text{mol}} > 0 \\ B\phi_w n & \text{if } n \geq n_{\text{crit}} \text{ and } f_{\text{mol}} > 0 \\ 0 & \text{otherwise,} \end{cases} \quad (9)$$

where

$$f_{\text{mol}} = \begin{cases} \phi_w & \text{if } T < 5000 \text{ K and } n \geq 10 \text{ cm}^{-3} \text{ and } \phi_w \geq 0.98 \\ 0 & \text{otherwise.} \end{cases} \quad (10)$$

The values for the temperature, density, and jet tracer thresholds in equation (10) were chosen to select a sufficient number of dense, cool cells in the simulations that are likely to contain molecular gas. The normalizations  $A$  and  $B$  are arbitrary, but satisfy  $B = A\phi_w n_{\text{crit}}$ . We took the critical density  $n_{\text{crit}} = 2 \times 10^4 \text{ cm}^{-3}$  for CO (2–1), although the PV diagrams are not very sensitive to the choice of  $n_{\text{crit}}$  since only a small fraction of cells in our simulations have densities above the critical density.

The emissivity in each cell is then integrated along the line of sight through a slit of width  $444 \text{ pc}$  across a length of  $1 \text{ kpc}$  along the mid-plane of the galaxy. The alignment and width of the slit are the same as those used by M15. We bin the integrated emission into 128 bins along the slit and 128 bins in velocity space between  $-800$  and  $800 \text{ km s}^{-1}$ . This results in somewhat finer spatial and velocity resolutions than those of the M15 data.

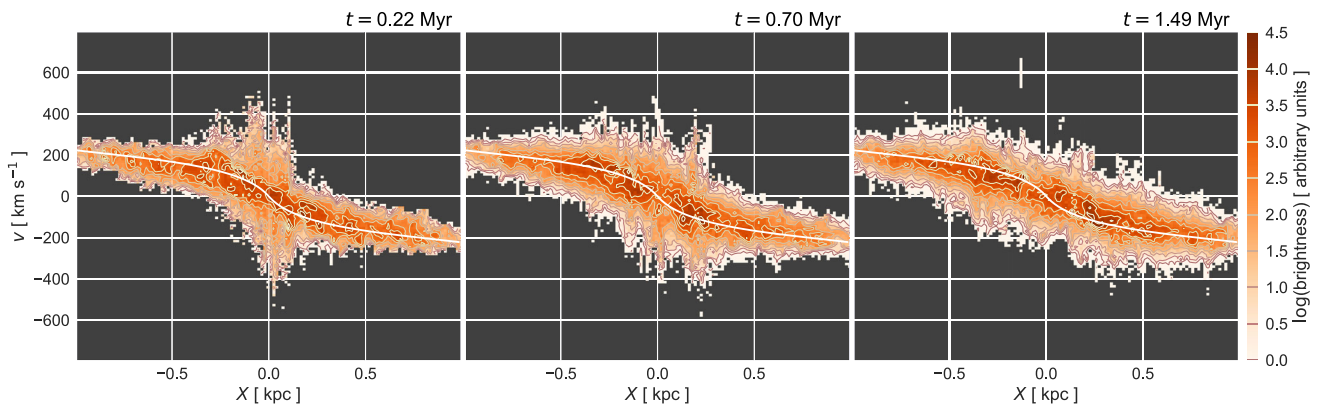
When producing the synthetic PV diagrams, we assume that the gas is optically thin. All the disturbed and outflowing gas in IC 5063 is known to be optically thin (Dasyra et al. 2016; Oosterloo et al. 2017). Furthermore, our simulations show that strong velocity dispersions exist throughout the disc, suggesting that the ISM in the disc may be a mixture of gas that is optically thin<sup>2</sup> and optically thick. In any case, our synthetic PV diagrams likely overestimate the surface brightness along the rotation curve relative to the surface brightness of the dispersed gas. Modelling the detailed distribution and radiative transfer of CO transitions for the disc is beyond the scope of this work and we therefore do not attempt to directly compare the range of the modelled and observed CO surface brightness in the PV diagram, focusing instead only on the gas kinematics. However, the fact that the kinematics of the neutral gas (Oosterloo et al. 2000) and the PV diagrams of ionized gas phases (Morganti et al. 2007) are similar to the PV diagram of CO (2–1) (M15) indicates that the underlying gas motions are physically related and that our model, therefore, properly captures the jet–ISM interactions that result in the observed velocity dispersions.

Given the above simplifications and keeping in mind that the dynamical range in density is limited by the spatial resolution of the simulations, the synthetic PV diagrams presented here are only meant to show qualitative trends and qualitative agreement with the ALMA data obtained by M15.

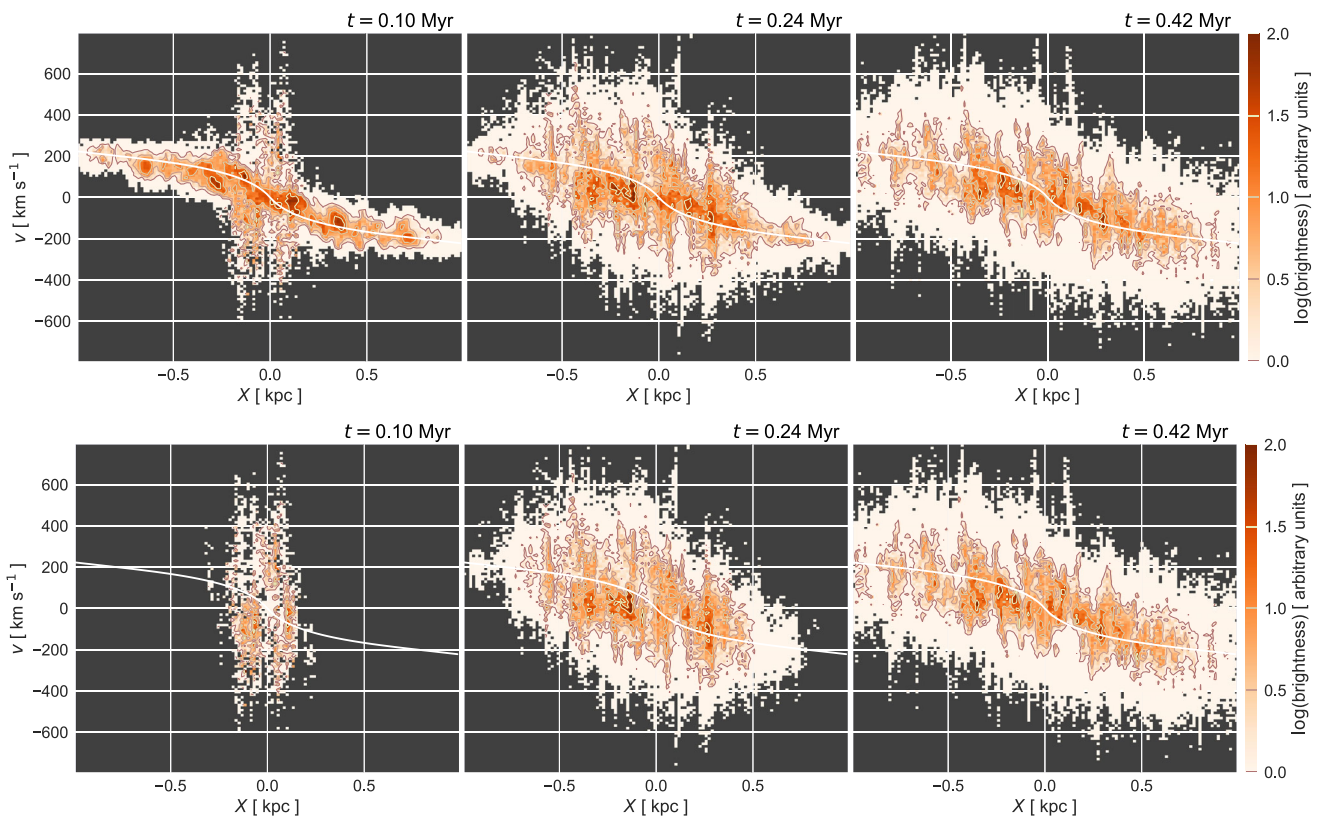
In the simulations, we can easily distinguish between gas impacted by the jet and gas in the unperturbed disc using the tracer variable for the jet plasma and selecting cells for which the jet tracer is greater than a small but non-zero value, typically  $10^{-8}$ . The small value ensures that most cells in the interior of clouds are selected, even if the mixing of the jet plasma and cloud material is dominated by numerical diffusion.

In Figs 8 and 9, we show PV diagrams for models P44 and P45, respectively. The solid white curve traces the projected circular velocity (equation 3) weighted by the mean density profile of the turbulent disc given by equation (5). It approximates the mean rotation of the disc gas represented in the PV plots at  $t = 0$  (see also Fig. 1). The brightness is in arbitrary units, but limited to a dynamical range of 4.5 dex, although for the PV diagrams of simulation P45, we limit the coloured contours to 2 dex, because

<sup>2</sup> The optical depth of line emission in a medium with velocity dispersion depends on the velocity gradient as  $\tau \propto \frac{1}{|dv/dz|}$  (see e.g. Castor 1970; Neufeld & Kaufman 1993).



**Figure 8.** PV diagrams generated from the P44 simulation for three different snapshots. The right-hand panel corresponds approximately to the current state of IC 5063 with regard to the extent of the jets.



**Figure 9.** PV diagrams for the model in which the jet power is  $10^{45}$  erg  $s^{-1}$ , a factor of 10 higher than that for the model shown in Fig. 8. The top row shows the PV diagrams for all gas according to equation (9), and the bottom row shows the PV diagram only for gas that is perturbed by the jet. The dynamic range covered by the colourmap and contours is set to 2 dex, although the white area shows the extent of the data down to 4.5 dex in dynamic range to aid comparison with Fig. 8.

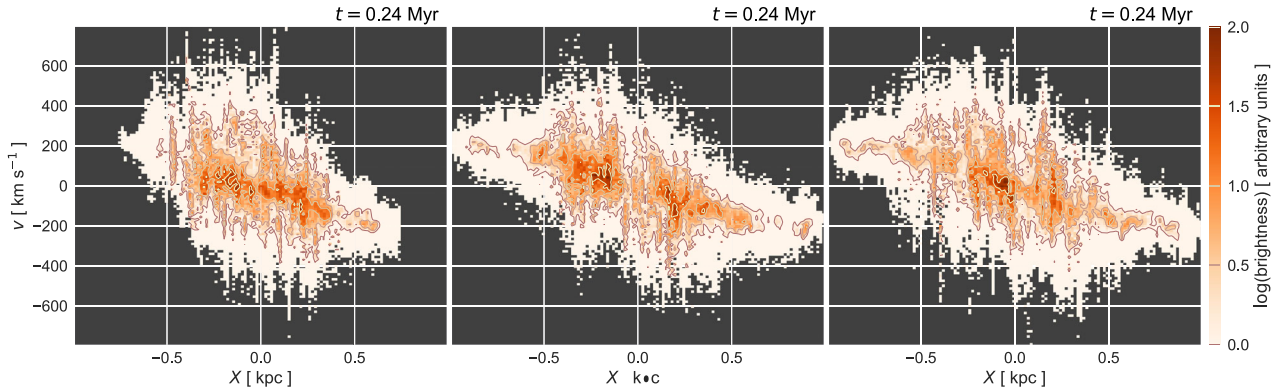
the velocity dispersions attained by the more powerful jet are higher. In fact, the velocity dispersions in simulation P45 are almost a factor of 2 higher compared to those in simulation P44, as expected from the scaling  $P_{\text{jet}} \propto \sigma^{1/4}$  found by Wagner & Bicknell (2011).

The top row in Fig. 9 shows the PV diagrams for all gas for which  $f_{\text{mol}} > 0$ , and the bottom row shows the PV diagrams only for gas perturbed by the jet. A visual comparison of the two rows of PV diagrams underlines how strong the effect of the jet on the dynamics of the clouds throughout the entire disc is.

We recall that there are five notable features in the observed PV diagrams of IC 5063 presented in M15: (1) broad velocity

dispersions of the order of a few 100  $\text{km s}^{-1}$  across the spatial region overlapping with the jet; (2) spikiness; (3) asymmetry in the velocity dispersions; (4) strong dispersion in the forbidden quadrants (i.e. motions in the counter-rotational direction); and (5) straightening of the rotation profile compared to that predicted by the gravitational potential.

To a greater or lesser extent, we recover all these features in the synthetic PV diagrams generated from the simulations. The broadening of the velocity dispersion and the spikiness of the broadening are particularly clear, especially for simulation P45, but also for simulation P44 while  $t < 0.7$  Myr. In the splattering regions



**Figure 10.** PV plots for simulation P45 for the case in which the plane of the disc is inclined by  $26^\circ$  to the line of sight, that is,  $\theta = 74^\circ$  (left) and for the cases in which the line of sight is parallel to the disc but oriented by an azimuthal angle of  $\phi = \pm 45^\circ$  towards the jet axis (centre and right).

of strong jet–ISM interactions, the broadening is by as much as  $600 \text{ km s}^{-1}$ ; the clouds here are dispersed nearly isotropically, giving rise to emission from the forbidden quadrants. At  $t = 1.49 \text{ Myr}$ , however, the jet in simulation P44 does not appear capable of producing the observed velocity dispersion. Simulation P45, on the other hand, exhibits strong velocity dispersion even at  $t = 0.24 \text{ Myr}$  when the jet has evolved roughly to the extents observed in the 8 GHz radio image. The asymmetry in the broadening is not as pronounced as in the data in M15 for either simulation, which may indicate an even more irregularly clumped ISM than we have assumed.

On the other hand, the straightening of the rotation profile along the central  $0.3 \text{ pc}$  region seen in the PV diagram, which reduces the steepening of the gradient compared to the solid white line, is clearly seen in both simulations. At  $t = 1.49$  and  $0.24 \text{ Myr}$ , simulations P44 and P45 show shallower rotation profiles, compared to the rotation profile at time  $t = 0$ . Simulation P45 displays a considerably stronger straightening of the rotation profile, compared to simulation P44. In simulation P45, the dense gas in the disc also appears to be more strongly broken up, as evidenced by the gaps between regions delineated by the highest brightness contour levels.

The degree to which the rotation profile is straightened due to jet–ISM interactions depends on how much of the inner dense clouds is displaced to larger radii. For simulation P45 at  $t = 0.24 \text{ Myr}$ , the gradient along the brightest parts is almost constant, meaning that most of the gas is redistributed to large disc radii. However, the process may be more complicated. Bulk cloud displacement does occur, but in the inner few kpc, most cloud material is ablated and clouds are dispersed in their entirety. We must also always bear in mind that we may not be capturing cloud survival properly – magnetic fields and turbulence may play a role, but we are also limited by numerical diffusion at the grid resolution and unresolved turbulent dissipation. We cannot, therefore, make a strong statement as to whether simulation P44 or simulation P45 represents a closer match to the observed data as far as the straightening of the rotation profile is concerned.

M15 outlined a model that explained the dispersions seen in the PV diagrams through radial mass transport by the pressure of the jet-driven, expanding radio bubbles. This model, although confined to a simple geometry, was able to explain the dispersion in the gas and the straightening of the rotation curve. Our simulations indicate that the physics in their model is correct, and in addition capture the stochastic nature of jet–cloud interactions that lead to the asymmetric and jagged features in the PV diagrams.

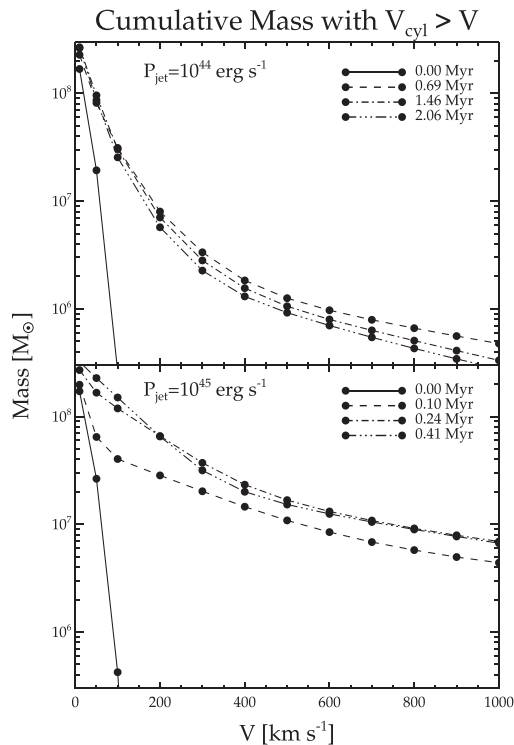
There is some freedom in the choice for the direction of line of sight with respect to the plane of the disc and the jet axis. We have tested various viewing angles and the PV diagrams are fairly similar for line-of-sight inclinations of less than  $\pm 45^\circ$  with respect to either the plane of the disc or the jet axis. Note that the pairs of viewing angles  $(\theta, \phi)$  and  $(-\theta, \phi)$  for any  $\theta$  and  $\phi$  are symmetric as far as the average density and velocity fields are concerned. However, pairs of viewing angles  $(\theta, \phi)$  and  $(\theta, -\phi)$  differ intrinsically because the orientation of the jet flux relative to the velocity field of the rotating disc gas along lines of sight offset from the centre is different.

In Fig. 10, we show the PV plot for simulation P45 for the case in which the line of sight is inclined by  $\theta = 74^\circ$  to the disc normal and for the cases in which the line of sight is oriented by an azimuthal angle of  $45^\circ$  with respect to the normal to the jet axis along the plane of the disc ( $\phi = \pm 45^\circ$ ). In all three cases, the main features of the PV diagram for simulation P45 along  $(\theta, \phi) = (0, 0)$  (Fig. 9, middle panel) are retained, but individual regions of strong dispersion vary. For example, when the viewing angle is tilted along  $\phi$  so that the component of the jet flux parallel to the line of sight acts in the opposite direction to the component of the rotating disc gas parallel to the line of sight, the emission from the forbidden quadrants becomes enhanced (right-hand panel in Fig. 10). Broadly speaking, however, the features of the synthetic PV diagrams are insensitive to the choice of line of sight within a few tens of degrees, and we attribute this to the isotropic nature of the gas dispersions caused by the jets.

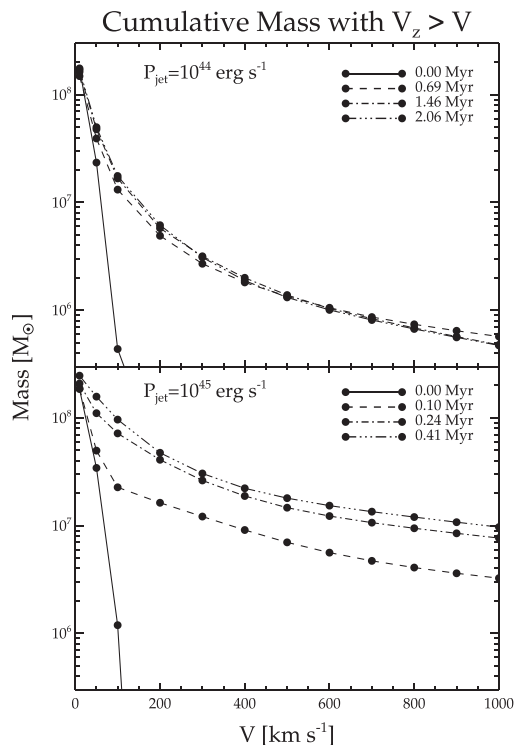
### 5.3 Other results

We measure the mass in the outflowing gas in our simulations, distinguishing between gas flowing along the plane of the disc, and gas flowing perpendicularly away from the galactic plane. In Fig. 11, we show the evolution of the mass of gas with cylindrical radial velocity  $v_{\text{cyl}} > v$  as a function of  $v$ , at different times for both simulations.

In Fig. 12, we show a similar plot of the mass in gas with vertical velocity  $v_z > v$  as a function of  $v$ . For simulation P44, gas of mass of  $8 \times 10^6 M_\odot$  is outflowing along the disc plane with a velocity greater than  $200 \text{ km s}^{-1}$  while in the case of simulation P45, it is  $7 \times 10^7 M_\odot$ . The outflowing gas masses perpendicular to the disc are similar. This is likely due to the fact that the jet plasma propagates quasi-isotropically within the confines of the disc, and, as a consequence, the mechanical advantage is equally effective in the  $z$ -direction as it is directed radially in the plane of the disc. In addition, dense gas that is accelerated in the  $z$ -direction experiences



**Figure 11.** Mass of gas with cylindrical velocity ( $v_{\text{cyl}}$ ) greater than a given value, at different times. After the jet has evolved significantly,  $8 \times 10^6 M_{\odot}$  of gas has an outflow velocity of  $\gtrsim 200 \text{ km s}^{-1}$  for simulation P44. For simulation P45, the total mass of gas with  $v_{\text{cyl}} \gtrsim 200 \text{ km s}^{-1}$  is  $\sim 7 \times 10^7 M_{\odot}$ .



**Figure 12.** Mass of gas with vertical velocity ( $v_z$ ) greater than a given value, at different times. After the jet has evolved significantly,  $6 \times 10^6 M_{\odot}$  of gas has an outflow velocity of  $\gtrsim 200 \text{ km s}^{-1}$  for simulation P44. For simulation P45, the total mass of gas with  $v_z \gtrsim 200 \text{ km s}^{-1}$  is  $\sim 5 \times 10^7 M_{\odot}$ . The outflowing masses in the plane of the disc are similar to the outflowing masses in the vertical direction (see Fig. 12).

less hindrance by other clouds located ahead. The outflowing gas masses translate to total mass outflow rates of  $\sim 8$  and  $70 M_{\odot} \text{ yr}^{-1}$ , for simulations P44 and P45, respectively, assuming a characteristic dynamical time-scale of  $t_{\text{dyn}} \approx 1 \text{ Myr}$ . Oosterloo et al. (2017) find total outflowing gas masses of at least  $5 \times 10^6 M_{\odot}$ , and a mass outflow rate of around  $12 M_{\odot} \text{ yr}^{-1}$ . Thus, in simulation P44, we find a mass outflow rate similar to that found observationally by Oosterloo et al. (2017), but the mass outflow rate for simulation P45 is about an order of magnitude larger.

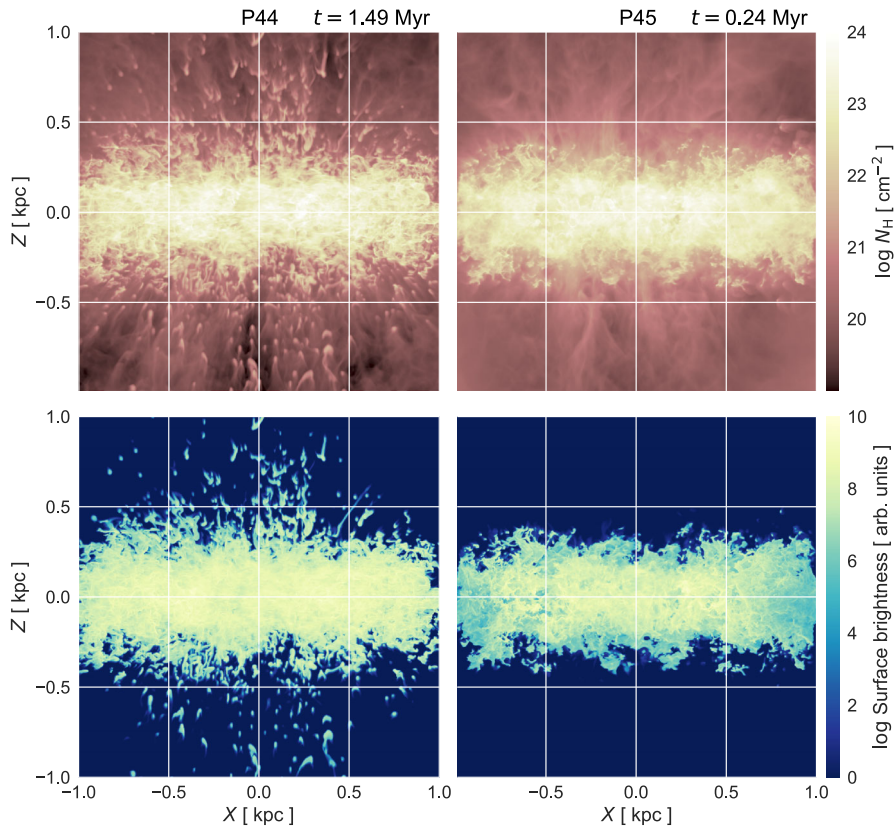
In Section 5.1, we mentioned that the jet plasma is also escaping perpendicular to the disc, entraining with it small clumps and filaments of disc material. While the 8 GHz radio images do not reveal the jet plasma escaping perpendicular to the disc, the 1.4 GHz continuum image presented in fig. 2 of Morganti et al. (1998) shows radio emission aligned with the minor axis of the galaxy. This radio structure may be the result of a starburst wind, but we propose the alternative explanation that it is radio plasma escaping perpendicular to the disc, as seen in our simulations. In fact, the galaxy-wide star formation rate of IC 5063 based on a far-infrared luminosity of  $L_{\text{FIR}} \approx 1.6 \times 10^{10} L_{\odot}$  is approximately  $5 M_{\odot} \text{ yr}^{-1}$  (Satyapal et al. 2005) with an associated 1.4 GHz luminosity of  $P_{1.4 \text{ GHz}} \approx 8 \times 10^{21} \text{ W Hz}^{-1}$  (Murphy et al. 2011). It appears unlikely that the radio structures of almost 10 kpc in extent are bubbles associated with a star formation activity of less than  $5 M_{\odot} \text{ yr}^{-1}$ . A direct comparison between the observed 1.4 GHz morphology and a synthetic radio map of the extended jet plasma would, however, require a simulation box size at least five times larger than that we have employed.

In Fig. 13, we show in the top two panels the column densities along  $Y$  for the simulations P44 and P45. In the second row of the same figure, we show the  $\text{H}\alpha$  surface brightness viewed along  $Y$  in the optically thin limit. The  $\text{H}\alpha$  emissivity is obtained directly from the cooling function. The snapshots are those for which the extents of the jets are in agreement with the currently observed extents of the jets. The striking difference between the two simulations is the presence of clumps lifted off the disc in simulation P44, whereas simulation P45 shows only diffuse filaments entrained in the jet plasma outflow perpendicular to the disc.

The difference between the two cases is a direct consequence of the power of the jets and the resulting confinement time. In simulation P45, the higher power jet heats and strips the outer layers of clouds on a shorter time-scale, dispersing the inner region of the ISM. The regions at larger disc radii remain more intact. In simulation P44, on the other hand, the jet plasma takes more time to percolate through the ISM, shocking and fragmenting the clouds. The jet plasma escaping perpendicular from the disc has more time to entrain and accelerate cloudlets in bulk. The acceleration of cloudlets is also aided by the vertical thermal pressure gradient. A similar result of low-power jets affecting a wider extent of the ISM has been shown in our previous papers (Mukherjee et al. 2016, 2017).

This interpretation of the difference in evolution is supported by maps of the line-of-sight velocity dispersion and mean temperature, shown in Fig. 14. For these plots, we weighted the velocity and temperature by density, and applied a cloud tracer threshold of  $\phi_w > 0.98$ . We clearly see that simulation P45 results in a higher velocity dispersion and higher mean temperature within the disc than simulation P44, and produces a higher velocity, hotter, more diffuse vertical outflow than simulation P44.

Also, as noted earlier in Fig. 12, the more powerful jet drives out more gas than the weaker jet. However, that gas is mostly in the hot diffuse form not visible in, e.g.  $\text{H}\alpha$ . Thus, the morphology



**Figure 13.** Comparison of side-on column density and  $H\alpha$  surface brightness between simulations P44 and P45. Top panels are the column density, and bottom panels are  $H\alpha$  surface brightness. The left-hand panels are simulation P44, and the right are P45. Simulation P44 shows cometary clumps lifted off the disc of the galaxy due to energy deposited in the disc by the jet during its long confinement time. The much shorter confinement time of the jet in simulation P45 does not lead to such an effect and the vertical outflow is mainly hot gas beyond  $10^5$  K.

of shocked gas observed from emission lines, and the phase of the outflowing gas (cold, warm, or hot), can help distinguish between different models of feedback.

Finally, we consider the possibility that molecules form in the dispersed or outflowing gas. In particular, if diffuse outflowing gas can condense again to cool dense filaments that form molecules, the velocity dispersions seen in the molecular phase may be achievable with a lower jet power. Since we do not follow the formation of molecules in our simulations, we test the plausibility of this scenario by comparing the cooling time to the dynamical time of the dispersed or outflowing gas. We show mid-plane slices for simulations P44 and P45 of the ratio of cooling time to dynamical time in Fig. 15. For the dynamical time, we take  $t_{\text{dyn}} = 1 \text{ kpc}/v_{\text{rad}}$  and the cooling time is  $t_{\text{cool}} = \varepsilon/\Lambda(\rho, T)$ , with a warm phase tracer cut of  $\phi_w > 0.98$ .

In both simulations, the cooling time in dense, slow-moving cloud surfaces and clump cores is shorter than the dynamical time, but the cooling times of gas filaments ablated from clouds exceed their dynamical time. The comparatively long cooling time in fast, diffuse filaments may indicate that molecule formation in these is unlikely, although a simulation including the chemistry of molecule formation, such as that performed by Richings & Faucher-Giguere (2017), is required to make a more conclusive statement.

## 6 SUMMARY AND DISCUSSION

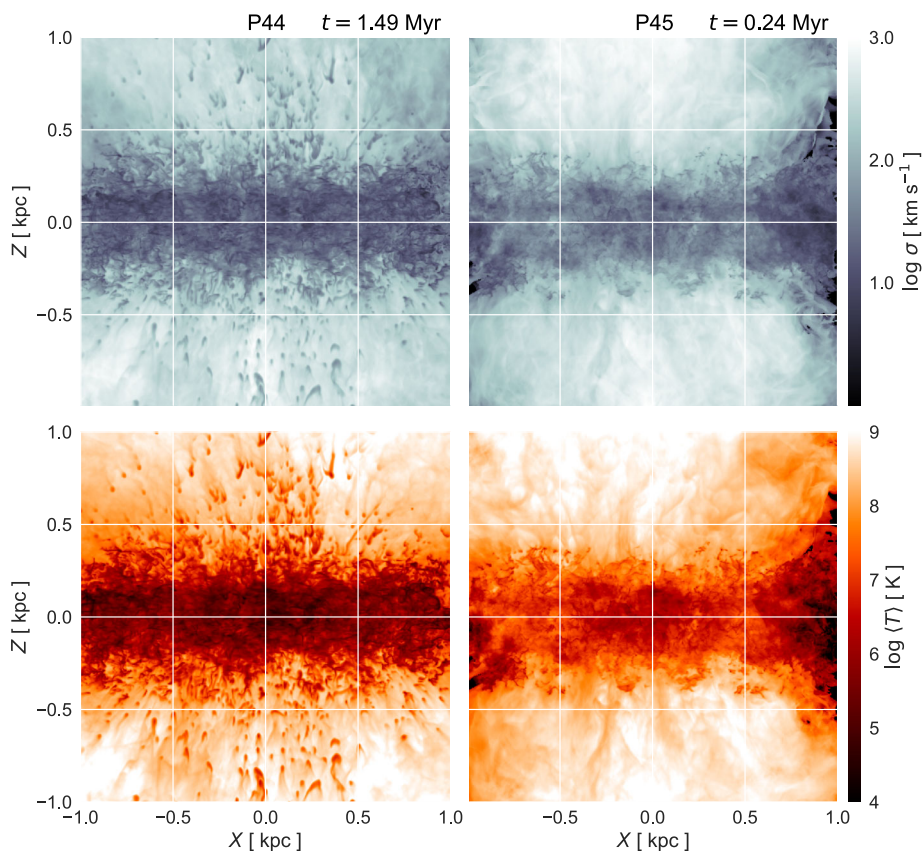
We have performed 3D relativistic hydrodynamic simulations of the jet–ISM interactions in IC 5063, a radio galaxy in which the jet is

aligned with the plane of the kpc-scale disc, and in which strong velocity dispersions are observed in all gas phases. We have set up a fractal, turbulent, rotating disc with parameters informed by observations, and have presented two simulations with different jet powers:  $10^{44}$  (P44) and  $10^{45} \text{ erg s}^{-1}$  (P45).

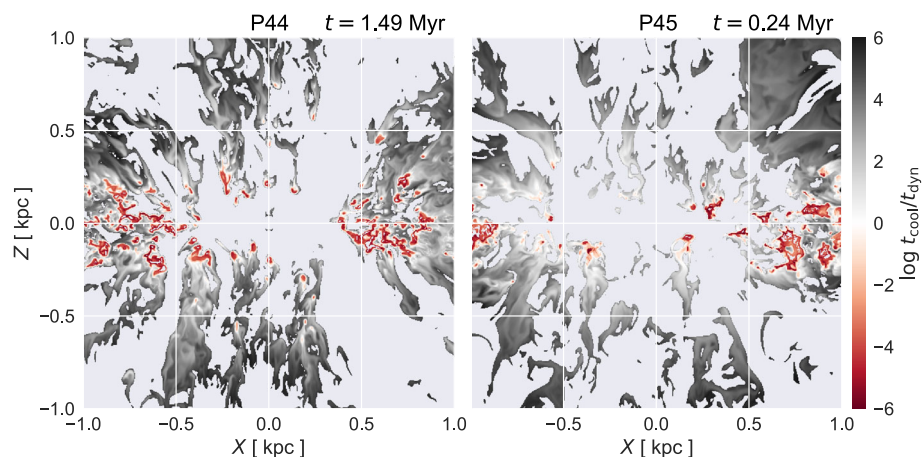
The jets in the simulation evolve very differently to the classical ‘dentist’s drill’ (Scheuer 1982). In the simulations, we recover the evolution of the flood-and-channel phase typical of jet–ISM interactions, in which clouds are shocked and gas is ablated from the clouds in the ISM. In this geometry, with the jet axis lying in the plane of the gaseous disc, the jet is strongly confined and jet plasma vents perpendicular to the disc.

The jet plasma that vents perpendicular to the disc entrains gas ablated from the clouds. In the case of simulation P44, in which the jet is confined six times longer than in the case of simulation P45, clumps of dense gas are lifted off the disc; these may possibly be visible in emission from a warm ionized medium (e.g.  $H\alpha$  as we show in Fig. 13).

In order to directly compare the gas kinematics resulting from jet–ISM interactions with gas kinematics inferred from recent ALMA observations by M15, we have constructed synthetic PV diagrams from the density, temperature, velocity, and cloud tracer fields in our simulations. We compare our synthetic PV diagrams to the PV diagram presented in M15 for CO (2–1) and find that our simulations broadly reproduce the key features in the observational data, including an increase in velocity dispersion to a few  $100 \text{ km s}^{-1}$  in the regions within approximately 0.6 kpc, in which the jet has



**Figure 14.** Comparison of side-on line-of-sight velocity dispersion and mean temperature along the line of sight between simulations P44 and P45. Top panels are the velocity dispersion, and bottom panels are mean temperature. The left-hand panels are from simulation P44, and the right are for simulation P45. The velocity dispersion and mean temperature are density-weighted and a dense phase tracer threshold of  $\phi_w > 0.98$  is applied. Simulation P44 shows cometary clumps lifted off the disc of the galaxy due to energy deposited in the disc by the jets during their long confinement time. The much shorter confinement time of the jets in simulation P45 does not lead to such an effect and the vertical outflow is mainly hot gas beyond  $10^6$  K.



**Figure 15.** Mid-plane slices perpendicular to the  $Y$ -axis of the ratio of cooling time to dynamical time for simulations P44 (left) and P45 (right). The dynamical time is defined as  $1 \text{ kpc}/v_{\text{rad}}$ , where  $v_{\text{rad}}$  is the radial velocity. Dense, slow-moving cloud surfaces and clump cores have short cooling times relative to their dynamical time, while the cooling times of gas filaments ablated from clouds exceed their dynamical time.

interacted strongly with the dense clouds in the ISM. In this region, the CO surface brightness is enhanced and coincident with the 8 GHz radio image.

The asymmetry in the PV diagrams between the eastern and western sides of the galaxy, as well as the jagged features in the PV plots, is also reproduced. These features are largely explained by

the clumpy, inhomogeneous nature of the ISM, through which the jet requires a substantial amount of time to penetrate. The forbidden quadrants in the PV diagram show a strong signal because the jet, especially near the jet head, disperses the gas nearly isotropically, resulting in gas being accelerated against the direction of the disc rotation.

The peaks in the radio surface brightness correspond to these locations, where the main jet stream impacts a massive dense cloud in its path and splitting into secondary jet streams while compressing, fragmenting, and dispersing the cloud. Such a ‘splatter region’ is particularly prominent, both in the observations and in our simulations, in the western part of the galaxy, where a significant velocity dispersion is observed in the molecular gas as well as in the ionized and neutral phases.

We have presented two models with jet powers  $10^{44}$  and  $10^{45} \text{ erg s}^{-1}$ , respectively. As expected, a higher jet power leads to stronger velocity dispersion in a much shorter time (1.49 Myr for simulation P44, and 0.24 Myr for simulation P45). Our results indicate that the ISM in our simulation is a reasonable representation of the ISM in IC 5063.  $P_{\text{jet}} = 10^{44} \text{ erg s}^{-1}$  is a lower limit to the jet power because jets of even lower power would likely not result in the observed velocity dispersion, unless molecules are efficiently produced in the diffuse outflowing gas over the short dynamical time of the outflow. A simple estimate of the dynamical time versus cooling time for the outflowing gas does not favour this scenario. The simulations by Richings & Faucher-Giguere (2017) also suggest that time-scales larger than a few  $10^5 \text{ yr}$  are required to obtain a significant amount of molecular gas in an AGN-driven outflow, although their simulations employed a smooth ISM.

The values of the jet power of IC 5063 quoted in the literature, including by M15, are of the order of  $10^{43}$ – $10^{44} \text{ erg s}^{-1}$ . They are obtained through scaling relations between radio power and cavity power derived for classical evolved radio jets doing work against smooth haloes of galaxies or clusters (Cavagnolo et al. 2010). Such scaling relations are known to be unreliable (Godfrey & Shabala 2016), and would not apply to jets propagating into the clumpy ISM of a galactic disc, as is the case for IC 5063.

From our simulations, we find the radio power at 1.4 GHz for the P44 simulation at time  $\sim 1.49 \text{ Myr}$  to be  $\sim 5.9 \times 10^{21} \text{ W Hz}^{-1}$  and that for the P45 simulation at  $\sim 0.24 \text{ Myr}$  to be  $\sim 4 \times 10^{23} \text{ W Hz}^{-1}$ . The above estimates assume a nominal value of  $f_e = f_B = 0.1$ , as discussed in Section 5. For an electron–positron plasma, the values of these fractions may be 0.5 each, which would increase the simulated radio powers by a factor of 92 (discussed in more detail in Bicknell et al. 2018). The observed radio power at 1.4 GHz was estimated to be  $3 \times 10^{23} \text{ W Hz}^{-1}$  (M15). Thus, the actual mechanical power of the jet is likely to lie in between the two values considered in this work  $10^{44}$  and  $10^{45} \text{ erg s}^{-1}$ . This is an order of magnitude higher than that estimated from the simple scaling relations obtained as an empirical fit to the estimates of radio power and energy in the radio cavities inferred from radio and X-ray data, respectively (Bîrzan et al. 2008; Cavagnolo et al. 2010). However, although the mechanical power of the jet in IC 5063 implied by our simulations differs from the mean relations derived in Cavagnolo et al. (2010), the values are consistent with the scatter in the measurements of the jet power (e.g. fig. 1 of Cavagnolo et al. 2010) at values of radio powers similar to that of IC 5063.

The kinematics of the gas affected by the jets as represented by the PV diagrams in M15 and Oosterloo et al. (2017) favour a jet power closer to  $10^{45} \text{ erg s}^{-1}$ , but the mass outflow rates favour a jet power closer to  $10^{44} \text{ erg s}^{-1}$ . Currently, we are not able to put stronger constraints on the jet power, but future studies that include the kinematics of the emission line gas, and an estimate of the jet age may help in this regard. Our work has shown, however, that the power of jets and the nature of the intervening ISM with which it

interacts can be constrained by careful modelling of the perturbed velocity fields that the jet–ISM interactions generate.

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## APPENDIX A: PROPAGATION TIME OF THE JET

We approximate the dynamical time of the jet propagating through the disc with the advance speed of the jet head calculated using equations (20) and (21) in Safouris et al. (2008, see also equations 23 and 24 in Sutherland & Bicknell 2007). Let  $\zeta$  be the ratio between jet density and the density of the ambient medium into which the jet is propagating. The time for the jet head to traverse a distance  $L$  is then

$$t_{\text{jet}} = \frac{L}{c} \left( 1 + \frac{1}{\Gamma} \sqrt{\frac{\chi}{(1+\chi)\zeta}} \right) \left( 1 - \frac{1}{\Gamma^2} \right)^{-1/2}. \quad (\text{A1})$$

Strictly, the equations mentioned above were derived for the advance speed of the terminal hotspot for jets propagating into a homogeneous medium. However, by using the average density of the multiphase interstellar medium in our simulations, we obtain a reasonable estimate for the jet head advance speed. Using  $L = 0.5$  kpc,  $\zeta = 10^{-5}$ ,  $\chi = 0.48$ , and  $\Gamma = 4$ , we obtain  $t_{\text{jet}} \approx 0.4$  Myr. This estimate agrees within a factor of a few with the duration of our simulation up to the point where the jets reach a distance of 0.5 kpc from the centre.

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