

Forced-ferromagnetic state in a Tm2Fe17H5 single crystal

E. A. Tereshina, M. D. Kuz'Min, Y. Skourski, M. Doerr, W. Iwasieczko, J.

Wosnitza, I. S. Tereshina

▶ To cite this version:

E. A. Tereshina, M. D. Kuz'Min, Y. Skourski, M. Doerr, W. Iwasieczko, et al.. Forced-ferromagnetic state in a Tm2Fe17H5 single crystal. Journal of Physics: Condensed Matter, 2017, 29 (24), 10.1088/1361-648X/aa70a6. hal-01694223

HAL Id: hal-01694223 https://hal.science/hal-01694223

Submitted on 10 May 2018

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers. L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

Letter Forced-ferromagnetic state in a Tm2Fe17H5 single crystal

E A Tereshina¹, M D Kuz'min², Y Skourski³, M Doerr⁴, W Iwasieczko⁵, J Wosnitza^{3,4} and I S Tereshina⁶

¹ Institute of Physics CAS, 18221 Prague, Czechia

² Aix-Marseille Université, IM2NP, UMR CNRS 7334, 13397 Marseille, France

³ Dresden High Magnetic Field Laboratory (HLD-EMFL), Helmholtz-Zentrum Dresden-

Rossendorf, D-01314 Dresden, Germany

⁴ Technische Universität Dresden, Institut für Festkörperphysik, D-01062 Dresden, Germany

⁵ Institute of Low Temperature and Structure Research PAS, 50-950 Wroclaw, Poland

⁶ Faculty of Physics, Lomonosov Moscow State University, 119991 Moscow, Russia

E-mail: tereshina@physics.msu.ru

Received 23 March 2017, revised 30 April 2017 Accepted for publication 3 May 2017 Published 17 May 2017

Abstract

We report the attainment of the ferromagnetic state in an interstitially modified heavy rare-earthiron intermetallic compound in an external magnetic field. The starting composition is RE2Fe17, which is the RE–Fe binary richest in iron. We concentrate on the Tm–Fe compound, which is the most sensitive to magnetic field. The maximum possible amount of hydrogen (5 at.H/f.u.) is

inserted into a Tm2Fe17 single crystal. We demonstrate that in a magnetic field of 57 T Tm2Fe

17H5 reaches the ferromagnetic state with an enviably high polarization of 2.25 T.

Keywords: single crystal, hydride, forced-ferromagnetic state, rare-earth-iron intermetallic compound

(Some figures may appear in colour only in the online journal)

1. Introduction

Magnetism of the rare-earth (RE)—iron intermetallic compounds is of special importance in view of applications. Suffice it to mention the celebrated Nd2Fe14B magnet, in which light boron stabilizes the crystal structure. The Sm2Fe17N_x system, interstitially modified with nitrogen (here x ranges between 2 and 3), has some advantages over Nd2Fe14B due to its high fundamental characteristics (such as magnetic polarization, Curie temperature, anisotropy constant [1, 2]), but its poten-tial has not been realized yet [3]. Improved characteristics can rationally be reached only with a better understanding of the magnetic properties of such interstitially modified RE2Fe17 materials and for that the use of bulk, nondecrepitated single crystals is highly desirable [4]. The RE and Fe sublattices are ordered ferro- or ferrimagnetically depending on the rare-earth atom [5–11]. The total magnetic moment of light rare earths (LRE) is usually small (~3 μ B/atom for Nd, ~0.7 μ B/atom for Sm) and is parallel to the total magnetic moment of the iron, which results in a mod-erate enhancement of the net polarization. It is for this reason that LRE-Fe intermetallics are used as magnetic materials [3, 6, 12– 14]. The situation in heavy rare earth (HRE)-Fe com-pounds is different. The total magnetic moment of HREs is as large as 10 μ B/atom (for Dy or Ho). Unfortunately, the large total moment is of no immediate benefit, since it is antiparal-lel to the moment of the iron sublattice. From a fundamental point of view, the magnetic behavior of HRE-Fe compounds (ferrimagnets) is more diverse than that of ferromagnets. While a ferromagnet simply saturates in external magnetic field, a ferrimagnet demonstrates a significant growth of mag-netization. It goes through a sequence of spinreorientation phase transitions until it reaches the fieldinduced ferromagn- etic state.

Sufficiently high external magnetic fields have to be employed in order to observe a complete magnetization pro-cess along different crystallographic directions of the sample. Mutual orientation of the RE spin and orbital moments is gov-erned by the spin-orbit coupling, $H_{so} = \lambda S_{RE} \cdot L_{RE}$, where the coupling constant λ is largely an atomic property [15]. All HRE have λ negative and large, $|\lambda| \sim 10^3$ K. As against that, the RE-Fe exchange interaction, $H_{ex} = J_{RE}$ -Fe $S_{RE} \cdot S_{Fe}$, is two orders of magnitude weaker, J_{RE} -Fe ~ 10 K [16]. In a favorable case, a magnetic field of several dozen Tesla should suffice to align S_{RE} and S_{Fe} parallel. More importantly, J_{RE} -Fe is not an atomic constant. It varies from compound to com-pound, depending in particular on interatomic distances [6].

Here, we report the attainment of the ferromagnetic state in HRE-Fe single-crystalline hydrides, observed in high magn- etic fields. It is advantageous to perform such experiments on Tm–Fe compounds. Since thulium is the HRE with a Landé factor closest to unity, it takes the smallest applied magnetic field to reach the ferromagnetic state (for a given $J_{\text{RE-Fe}}$). For the same reason, the field in our experiments was always applied along the easy magnetization direction of single-crystalline samples.

Earlier, a Tm₂Fe₁₇ single crystal and its polycrystalline deuteride Tm₂Fe₁₇D_{3.2} were studied in fields up to 74 and 60 T, respectively [17]. Unfortunately, Isnard *et al* did not reach the ferromagnetic state in both compounds. In the present work, samples with various hydrogen content, including the maximum possible amount of 5 at.H/f.u., were used. We showed that Tm₂Fe₁₇H₅ reaches the ferromagnetic state in a magnetic field of 57 T and attains an enviably high polarization of 2.25 T. The study is performed on single-crystalline Tm₂Fe₁₇H_x samples that allowed us to obtain reliable data on the type and behavior of the magnetocrystalline anisotropy in these materials.

2. Materials synthesis and experimental details

Initial samples were prepared by arc-melting of 99.9% pure elements under helium atmosphere. To stimulate growth of single-crystalline grains, the ingots were re-melted and cooled down slowly in a resistance furnace. Several 1 mm-large grains were extracted and checked by x-ray Laue tech-nique. Usually single crystals decrepitate when exposed to hydrogen. To prevent this from happening, the hydrogenation was carried out under particularly mild conditions—elevated temperature and low partial pressure of H2—both of which slow down the reaction. The higher hydrogen contents were achieved by subsequently raising the pressure. The amount of absorbed hydrogen in Tm2Fe17H_x was determined from the pressure drop

in the reaction chamber. Phase composi-tion of the samples was controlled by powder x-ray diffrac-tion. Magnetization measurements were performed by using a pulsed-field induction magnetometer [18] at the Dresden High

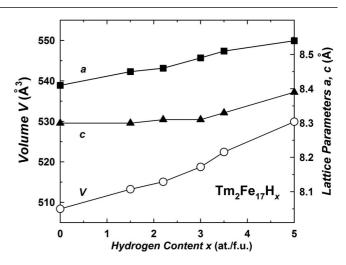


Figure 1. The lattice parameters *a* and *c* and the unit cell volume *V* of $Tm_2Fe_{17}H_x$ (lines are guides to the eye).

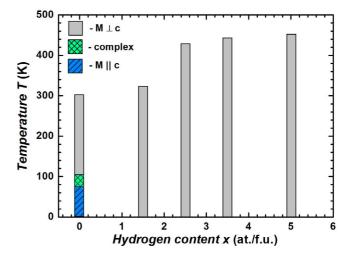


Figure 2. Magnetic phase diagram demonstrating the orientation of the magnetic moment with respect to the crystal axis *c* for Tm₂Fe₁₇ and some hydrides Tm₂Fe₁₇H_x. The width of the bars corresponds to the experimental error in *x*.

Magnetic Field Laboratory. The maximum field was 60 T with a total pulse duration of 25 ms [19]. A commercial magnetom-eter (Quantum Design PPMS-14) was used for magnetization measurements in steady magnetic fields up to 14 T at temper- atures between 2 and 500 K.

3. Results and discussion

Pure Tm₂Fe₁₇ has the hexagonal Th₂Ni₁₇-type crystal structure (space group *P*63/*mmc*) with lattice parameters a = 8.409(3) Å and c = 8.300(3) Å. Hydrogenation increases the unit-cell volume of Tm₂Fe₁₇ without changing its crystal structure. The lattice parameter *a* of Tm₂Fe₁₇H_{*x*} grows linearly with hydrogen content *x*, see figure 1. However, *c* remains practically constant as hydrogen fills the octahedral interstitial sites (at x < 3), and increases in the interval 3 = x - 5, as the tetrahe-dral sites are filled [20]. For the composition Tm₂Fe₁₇H₅, the lattice parameters are a = 8.540(3) Å and c = 8.381(3) Å; the relative volume change upon hydrogenation V/V is ~4.2%.

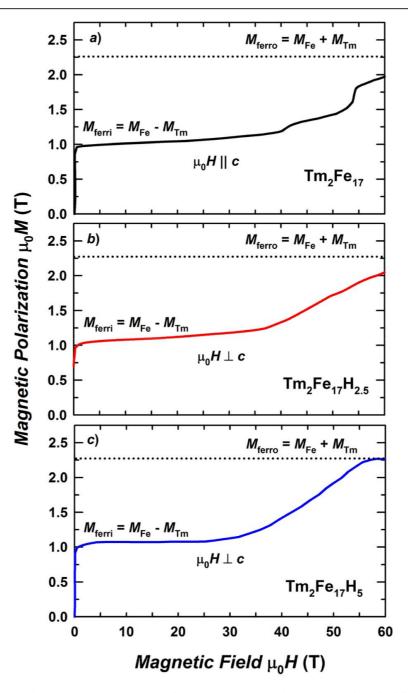


Figure 3. Magnetization curves of Tm₂Fe₁₇, Tm₂Fe₁₇H_{2.5} and Tm₂Fe₁₇H₅ measured at T = 4.2 K in pulsed fields applied along the easy axes.

Tm₂Fe₁₇ is a ferrimagnet with a Curie temperature of $T_{\rm C}$ = 295 K and a spontaneous moment of 22 $\mu_{\rm B}$ /f.u. ($\mu_0 M_{\rm ferri}$ = 1.0 T) at T = 4.2 K. Figure 2 shows the magnetic phase diagram of Tm₂Fe₁₇, as well as of its hydrides Tm₂Fe₁₇H_x, deduced from magnetization data measured on single crystals along different crystallographic directions in a wide temperature range up to 500 K (not shown here). At low temperatures, the anisotropy of pure Tm₂Fe₁₇ is of the 'easy-axis' type; this changes to 'easy plane' as temperature increases. The change proceeds via an intermediate 'easy-cone' phase. Our results agree well with data presented in [17].

Hydrogenation is known [4] to affect strongly (i) the type of magnetic anisotropy (all investigated Tm₂Fe₁7H_x hydrides are easy-plane magnets), and (ii) the Curie temperature (seen to rise from 295 to 450 K as x grows from 0 to 5). It also modi-fies significantly the magnetization curves of Tm₂Fe₁7H_x, see figure 3. Here, the field was applied along the easy magnetization direction, namely, along the *c* axis for Tm₂Fe₁7H_{2.5} (figure 2(a)) and perpendicular to the *c* axis for Tm₂Fe₁7H_{2.5} (figure 2(b)) and Tm₂Fe₁7H₅ (figure 2(c)) (the basal-plane magnetocrystalline anisotropy of the studied compounds is close to zero). For Tm₂Fe₁7, one can observe two rather sharp field-induced transitions, where the magnetization increases

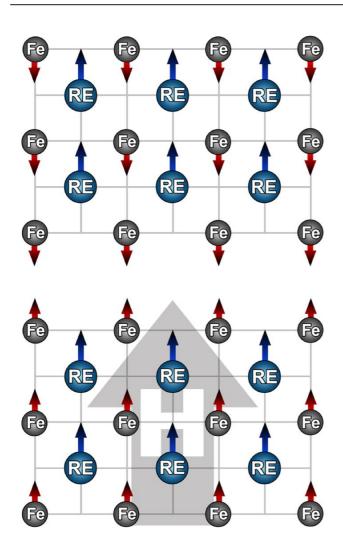


Figure 4. Ferrimagnetic (top) and field-induced ferromagnetic (bottom) states of an HRE-Fe compound.

abruptly. As against that, the magnetization of Tm₂Fe₁₇H_{2.5}

and Tm₂Fe₁₇H₅ shows a continuous growth above a certain threshold field. For the final composition with x = 5 the magn- etic field of 57 T was sufficient for attaining the ferromagnetic state (figure 3(c)).

Let us assume for the ground state of Tm₂Fe₁₇ an antiparallel configuration of the Tm and Fe moments

 $(M_{\text{ferri}} = M_{\text{Fe}} - M_{\text{Tm}} = 22 \ \mu_{\text{B}}/\text{f.u.})$ and adopt the free-ion value for the magnetic moment of Tm, $M_{\text{Tm}} = 14 \ \mu_{\text{B}}/\text{f.u.}$

This yields for the iron-sublattice moment $M_{\text{Fe}} = 36 \ \mu\text{B/f.u.}$, which is in good agreement with what is known for RE2Fe17 [13]. In the forced ferromagnetic state, Tm2Fe17 should have

 $M_{\text{ferro}} = M_{\text{Fe}} + M_{\text{Tm}} = 50 \ \mu_{\text{B}}/\text{f.u.}$ ($\mu_0 M_{\text{ferro}} = 2.29 \ \text{T}$). For Tm₂Fe₁₇H₅, $M_{\text{ferri}} = 23 \ \mu_{\text{B}}/\text{f.u.}$ and $M_{\text{ferro}} = 51 \ \mu_{\text{B}}/\text{f.u.}$, that is, $\mu_0 M_{\text{ferri}} = 1.02 \ \text{T}$ and $\mu_0 M_{\text{ferro}} = 2.25 \ \text{T}$. In both compounds, the magnetization more than doubles as a result of the ferri-

to-ferro transition. Since hydrogenation weakens the Tm–Fe exchange we are able to observe the ferromagnetic state for Tm₂Fe₁₇H₅ already at 57 T (figures 3(c) and 4). Neglecting the weak Tm–Tm interaction, one can consider the Tm–Fe exchange integral J_{TmFe} to be proportional to the molecular

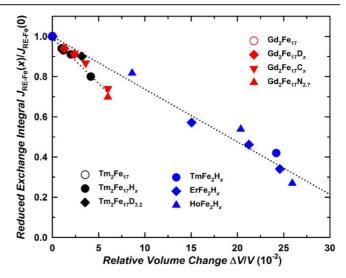


Figure 5. Reduced exchange integral $J_{RE-Fe}(x)/J_{RE-Fe}(0)$ for

RE2Fe17(H, N, C)_{*x*} (RE = Gd [22, 23], Tm) and REFe2H_{*x*} (RE = Ho, Er, Tm [25–27]) as a function of the relative volume change V/V.

field on Tm. The latter can be determined by several meth-ods [16, 6, 17] In [17], it was shown that the molecular field is just given by the abscissa of the midpoint of the sloping part of the experimental magnetization curve. Since the pro-portionality coefficient is not accurately known, it is conveni-

ent to describe the effect of hydrogen on $J_{Tm}Fe$ by means of the reduced quantity,- $J_{Tm}Fe(x)/J_{Tm}Fe(0)$, i.e. the exchange integral or the molecular field normalized to unity for the parent compound. This is plotted for Tm₂Fe₁₇H_x versus the relative volume change, V/V, as determined from the lattice parameters (figure 5). In agreement with earlier observations [21], $J_{RE-Fe}(x)/J_{RE-Fe}(0)$ decreases as a result of hydrogenation. Moreover, in figure 5 we present similar data for

 $Gd_2Fe_{17}(D,N,C)_x$ compounds, whose molecular fields were determined by use of inelastic neutron scattering [22, 23]. All data can be approximated by a straight line.

Interestingly, the J_{RE-Fe} data plotted against relative volume change in HREFe₂-H systems (figure 5) also demonstrate a linear trend for the HRE Ho, Er and Tm [24–26]. Hydrogenation of the cubic Laves phases REFe₂ can well produce values of V/V between 25 and 30%, while in RE₂Fe₁₇ by way of hydrogenation, nitrogenation, and/or carbonation a volume increase of just under 9–10% appears [1]. For REFe₂, the exchange fields weaken by up to about 70%, whereas for RE₂Fe₁₇, it weakens by 20–30%. Despite the considerable decrease of J_{TmFe} , high magnetic fields are required in order to attain the forced-ferromagnetic state in TmFe₂H_x due to the increased magnetocrystalline anisotropy upon hydrogenation caused by the structural phase transition from the parent cubic to the rhombohedral structure.

4. Conclusions

The study of the field-induced transitions from a ferri- to a ferromagnetic state in the intermetallic compounds is an interesting fundamental problem. Up to now very few forced-ferromagnetic states in rare-earth intermetallic compounds

with iron were reported [24, 27]. In the present work, the measurements were carried out on single-crystalline Tm2Fe17 hydrogenated to the maximum hydrogen content of 5 H atoms per formula unit along the easy magnetization direction in fields up to 60 T. Hydrogenation allowed us to decrease the Tm–Fe exchange and thus to lower the transition field to the forced-ferromagnetic state down to 57 T as compared to about 80 T expected for the parent compound Tm2Fe17 [17]. In the ferromagnetic state, Tm2Fe 17H5 attained an enviably high polarization of 2.25 T. The larger value $\mu 0M_{\text{ferro}} = 2.65$

T can be obtained for Ho₂Fe₁₇ in the even higher applied magnetic field (84 T [18]). (Other numbers for comparison: the celebrated hard magnetic material Nd₂Fe₁₄B has 1.85 T at T = 0 K while this value in the industrial magnets does not surpass 1.5 T).

Although it may take a lower magnetic field to stabilize ferromagnetism in Co-based compounds, e.g. 39 T sufficed for Tm₂Co₁₇ [28], polarization in the ferromagnetic state is lower, too (μ 0*M*ferro = 1.96 T for Tm₂Co₁₇). The search for new multicomponent compounds with substitutional or/and interstitial atoms where a transition into a high-magnetization state takes place at even lower magnetic field is of practical importance. Further progress in this field can be anticipated.

Acknowledgments

We wish to thank Prof D Givord and Prof N V Mushnikov for useful discussion, O Chistyakov and A Hackemer for help in producing single-phase parent samples. This work was supported by RFBR, project No. 16-03-00612. We gratefully acknowledge the support of the HLD at HZDR, member of the European Magnetic Field Laboratory (EMFL). Part of the measurements has been performed in MLTL (http://mltl. eu) and was supported by Material Growth and Measurement Laboratory MGML (http://mgml.eu/).

References

- Fujii H and Sun H 1995 Interstitially modified intermetallics of rare earth and 3d elements *Handbook of Magnetic Materials* vol 9, ed K H J Buschow (Amsterdam: North-Holland) ch 3
- [2] Nakayama H, Takagi K, Ozaki K and Kobayashi K 2012 Correlation between microstructure and magnetic properties in Sm₂Fe₁₇N₃ magnet prepared by pulsed current sintering *Mater. Trans.* 53 1962–6
- [3] Gutfleisch O, Willard M A, Brück E, Chen C H, Sankar S G and Liu J P 2011 Magnetic materials and devices for the 21st century: stronger, lighter, and more energy efficient Adv. Mater. 23 821–42
- [4] Tereshina E A, Drulis H, Skourski Y and Tereshina I S 2013 Strong room-temperature easy-axis anisotropy in Tb₂Fe₁₇H₃: an exception among R₂Fe₁₇ hydrides *Phys. Rev.* B 87 214425
- [5] Duc N H 1997 Intersublattice exchange coupling in the lanthanide-transition metal intermetallics *Handbook on the Physics and Chemistry of Rare Earths* vol 24, ed K A Gschneider Jr and L Eyring (Amsterdam: Elsevier) ch 163
- [26] Mushnikov N V 2004 Magnetovolume effects and magnetic anisotropy in the itinerant and localized subsystems of f–d intermetallics *Thesis of Doctor of Sciences* Ekaterinburg (in Russian)

- [6] Campbell I A 1972 Indirect exchange for rare earths in metals J. Phys. F: Met. Phys. 2 L47
- [7] France J J M and Radwanski R J 1993 Magnetic properties of binary rare-earth 3d-transition-metal intermetallic compounds *Handbook of Magnetic Materials* vol 7, ed K H J Buschow (Amsterdam: North-Holland) ch 5, p 307
- [8] Buschow K H J 1988 Permanent magnet materials based on 3d-rich ternary compounds *Handbook of Magnetic Materials* vol 4, ed E P Wohlfarth and K H J Buschow (Amsterdam: North-Holland) ch 1
- [9] Brooks M S S, Gasche T, Auluck S, Nordström L, Severin L, Trygg J and Johansson B 1992 Density functional theory of molecular fields in R-M systems J. Magn. Magn. Mater. 104–7 1381–2
- [10] Belorizky E, Fremy M E, Gavigan J P, Givord D and Li H S 1987 Evidence in rare-earth (R)-transition metal (M) intermetallics for a systematic dependence of R-M exchange interactions on the nature of the R atom J. Appl. Phys. 61 3971–3
- [11] Verhoef R, de Boer F R, Sinnema S, Franse J J M, Tomiyama F, Ono M, Date M and Yamagishi A 1992
 Multiple field-induced transitions in Er₂Fe₁₇ in fields up to 50 T *Physica* B 177 223–6
- [12] Skomski R and Coey J M D 1998 Permanent Magnetism (Bristol: Institute of Physics)
- [13] Herbst J F 1991 R2Fe14B materials: intrinsic properties and technological aspects *Rev. Mod. Phys.* 63 819–98
- [14] Coey J M D 2010 Magnetism and Magnetic Materials (Cambridge: Cambridge University Press)
- [15] Skomski R 2008 *Simple Models of Magnetism* (Oxford: Oxford University Press)
- [16] Liu J P, de Boer F R, de Châtel P F, Coehoorn R and Buschow K H J 1994 On the 4f–3d exchange interaction in intermetallic compounds J. Magn. Magn. Mater. 132 159–79
- [17] Isnard O et al 2013 High magnetic field study of the Tm2Fe17 and Tm2Fe17D3.2 compounds Phys. Rev. B 88 174406
- [18] Skourski Y, Kuz'min M D, Skokov K P, Andreev A V and Wosnitza J 2011 High-field magnetization of Ho₂Fe₁₇ Phys. Rev. B 83 214420
- [19] Zherlitsyn S, Wustmann B, Herrmannsdörfer T and Wosnitza J 2012 Status of the pulsed-magnet-development program at the Dresden high magnetic field laboratory *IEEE Trans. Appl. Supercond.* 22 4300603
- [20] Isnard O, Miraglia S, Soubeyroux J L, Fruchart D and I'Héritier P 1994 A structural analysis and some magnetic properties of the R₂Fe₁₇H_x series *J. Magn. Magn. Mater.* 137 151–6
- [21] Hong N M 1996 Volume dependence of magnetic interactions in R₂Fe₁₇ and R₂Fe₁₄B and their interstitial compounds with H, C and N *Physica* B 226 391–8
- [22] Isnard O, Sippel A, Loewenhaupt M and Bewley R 2001 A high energy inelastic neutron scattering investigation of the Gd–Fe exchange coupling in Gd2Fe17D_x (x = 0, 3 and 5) J. Phys.: Condens. Matter 13 3533–43
- 5) J. Phys.: Condens. Matter 13 3533–43
 [23] Loewenhaupt M, Tils P, Buschow K H J and Eccleston R S 1995 Nitrogen- and carbon-induced reduction of Gd– Fe exchange coupling in Gd2Fe17 doped samples J. Alloys Compd. 222 39–43
 [24] Mushnikov N V, Zajkov N K, Bartashevich M I, Goto T,
- [24] Mushnikov N V, Zajkov N K, Bartashevich M I, Goto T, Aruga-Katori H, Yamaguchi M and Yamamoto I 1997 Magnetization process and magnetic structure of the TmFe₂H_{3.4} hydride in ultrahigh magnetic fields up to 100 T J. Magn. Magn. Mater. 167 229–36
- [25] Deryagin A V, Kudrevatykh N V, Moskalev V N and Mushnikov N V 1984 Phys. Met. Metall. 58 96–100
- [27] Kato H, Lim D W, Yamada M, Nakagawa Y, Aruga Katori H and Goto T 1995 Field-induced phase transitions in

ferrimagnetic R₂Fe₁₄B in ultra-high magnetic fields *Physica* B **211** 105 [28] Andreev A V, Kuz'min M D, Narumi Y, Skourski Y, Kudrevatykh N V, Kindo K, de Boer F R and Wosnitza J

2010 High-field magnetization study of a Tm₂Co₁₇ single crystal *Phys. Rev.* B **81** 134429