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LETTER TO THE EDITOR

The Milky Way rotation curve revisited

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ABSTRACT

The *Herschel* survey of the Galactic Plane (Hi-GAL) is a continuum Galactic plane survey in five wavebands at 70, 160, 250, 350 and 500 μm . From such images, about 150 000 sources have been extracted for which the distance determination is a challenge. In this context the velocity of these sources has been determined thanks to a large number of molecular data cubes. But to convert the velocity to kinematic distance, one needs to adopt a rotation curve for our Galaxy. For three different samples of tracers, we test different analytical forms. We find that the power-law expression, $\theta(R)/\theta_0 = 1.022 (R/R_0)^{0.0803}$ with $R_0, \theta_0 = 8.34 \text{ kpc}, 240 \text{ km s}^{-1}$ is a good and easily manipulated expression for the distance determination process.

Key words. Galaxy: disk – Galaxy: kinematics and dynamics

1. Introduction

In this paper we propose to update the rotation curve of our Galaxy. This is done to determine the distance of the sources extracted from the Hi-GAL survey (e.g. Molinari et al. 2016). This open time key-project of the *Herschel* Space Observatory (Pilbratt et al. 2010) is a 5-band photometric imaging survey (with the SPIRE and PACS photometric cameras) at 70, 160, 250, 350, and 500 μm of a $|b| \leq 1^\circ$ wide strip of the Milky Way Galactic plane.

In the context of the Vialactea Project¹ a large number of molecular data cubes, covering the full galactic plane, have been compiled (Molinari et al. 2016) from which velocity for the HiGAL sources are extracted (Mège et al., in prep.). To convert this velocity into kinematic distance we need to adopt a rotation curve for our Galaxy.

Up to now, two main rotation curves have generally been used: the Clemens (1985) and the Brand & Blitz (1993) curves. To determine the global rotation curve it is important to have access to data from all the galactic quadrants and to combine the three usual methods to link the velocity to the distance. Usually CO and/or HI observations of the interstellar medium are used to establish the inner (inner the Solar circle) rotation curve (assuming circular rotation) thanks to the tangent-point method (which used the terminal line-of-sight velocities). Such an approach is complemented adding the H II regions (or star-forming complexes) to trace the external part (outside the Solar circle) of the rotation curve because for H II regions one can independently measure the velocity of the gas and the distance of the exciting stars. Recently, a third method based on the maser parallaxes produced a new set of data. Indeed, maser(s) observed in star-forming regions are used to determine their parallactic distance and are very useful to determine the distance of embedded

star-formation regions for which the classical exciting OB stars are not observable. Recently, Reid et al. (2009, 2014) used maser parallax distances of star-forming regions to trace the rotation curve. However, these new results are determined from a “small” number (100) of high-mass star-forming regions and from the northern part of the galactic plane ($l \sim 0^\circ$ to 240°). Because of the long-known asymmetry (Kerr 1964; Georgelin & Georgelin 1976; Blitz & Spergel 1991; Levine et al. 2008) between the southern and the northern rotation curves, to determine the distance of any source in our Galaxy it is better to use a rotation curve established from the full range of longitudes (Gómez 2006). It is in this context that we revisit the Milky Way rotation curve.

2. The sample

The best method to determine the rotation curve of our galaxy is to independently measure the velocity and the distance (exciting star or maser parallax distances) of the objects. However, this method limits probing the local Galaxy only (the stellar distance can be evaluated only up to about 6 kpc due to extinction, and maser parallax in star-forming regions requires long time baseline). To probe the inner rotation of our Galaxy on a larger scale the H I/CO tangent method is usually used. If this method can have uncertainties due to local motions, Chemin et al. (2015) show that it is adequate for galactocentric distances larger than 4.5 kpc.

We propose here to update the rotation curve of our Galaxy by combining different samples:

1. Sample 1: the H II regions/complexes-stellar distance cataloged by Brand & Blitz (1993). Their H II regions catalog has been downloaded from Vizier² and provides, for every

¹ <http://vialactea-sg.oact.inaf.it:8080/web/guest/home>

² <http://vizier.u-strasbg.fr/viz-bin/VizieR?-source=J/A+A/275/67>

object, the l, b coordinates, the V_{LSR} , and the stellar distance. This catalog provides a sample of 152 objects (select to have $|b| < 3^\circ$) distributed in the four galactic quadrants.

2. Sample 2: the maser parallax distance catalogs of star-forming regions from Reid et al. (2014) and Honma et al. (2012). We retrieved the catalogs from Reid et al. (2014) and Honma et al. (2012). These catalogs give, for each object, the l, b coordinates, the V_{LSR} (CO velocity) of the associated molecular cloud, and the parallax π (which is then converted into distance as $d_\pi = 1/\pi$). Both catalogs give a final sample of 101 objects (selected to have distance larger than 1 kpc) located mainly in quadrants 1 and 2 of our Galaxy.
3. Sample 3: the HI tangent+CO-HII regions catalog from Sofue et al. (2009). They compiled HI tangent point data from Burton & Gordon (1978), Clemens (1985) and Fich et al. (1989), HI-disk thickness method from Honma & Sofue (1997a,b), CO and HII regions from Fich et al. (1989) and Blitz et al. (1982). This data provides, for every object, the rotation velocity and the galactocentric distance (calculated with $R_0, \theta_0 = 8$ kpc, 200 km s^{-1}). Selecting data with velocity uncertainty less than 50 km s^{-1} and rotation velocity between 150 and 350 km s^{-1} gives a sample of 408 measurements.
4. Sample 4: the HI tangent measurements from McClure-Griffiths & Dickey (2007) and McClure-Griffiths & Dickey (2016). They are based on the Southern Galactic Plane survey (McClure-Griffiths et al. 2005) and the VLA galactic plane survey (Stil et al. 2006). The northern and southern survey cover $18^\circ \leq l \leq 67^\circ$ (with latitude varying from $|b| < 1.3^\circ$ to $|b| < 2.3^\circ$), and $253^\circ \leq l \leq 358^\circ$ ($|b| \leq 1.5^\circ$), respectively. Selecting only the data with galactocentric distance larger than 4 kpc, gives a sample of 1243 measurements. In their data tables, no velocity uncertainty is given, so we adopt a 10 km s^{-1} uncertainty.

3. The adopted local standard of rest and solar motion parameters

Up to now, to establish the V_{LSR} , and to determine the distance, assumptions have been made concerning the Solar parameters: first the Solar motions ($U_\odot, V_\odot, W_\odot$) to the local standard of rest (LSR) and the LSR parameters, which are the distance of the Sun to the galactic center (R_0) and the rotation velocity (θ_0). The IAU standard values for these quantities are $R_0 = 8.5$ kpc, $\theta_0 = 220 \text{ km s}^{-1}$, $U_\odot = 10.27 \text{ km s}^{-1}$, $V_\odot = 15.32 \text{ km s}^{-1}$, and $W_\odot = 7.74 \text{ km s}^{-1}$. One can recall that $U_\odot, V_\odot, W_\odot$ are exclusively used to calculate the V_{LSR} from the measured heliocentric radial velocity while R_0 and θ_0 are used in the kinematic distance determination. From maser parallaxes, new R_0, θ_0 values were also determined as being 8.05 ± 0.45 kpc, $238 \pm 14 \text{ km s}^{-1}$ and 8.34 ± 0.16 kpc, $240 \pm 8 \text{ km s}^{-1}$ by Honma et al. (2012) and Reid et al. (2014), respectively. Such a low R_0 value (between 7.7 and 8.27 kpc) is also found from independent measurements (Meyer et al. 2012; Gillessen et al. 2013; Chatzopoulos et al. 2015) as recently underlined by Boehle et al. (2016) who find $R_0 = 7.86$ kpc. In parallel, ω_0 , determined from different approaches (Feast & Whitelock 1997; Reid & Brunthaler 2004; Reid et al. 2014; Bobylev 2017), is between 27.19 and $29.45 \text{ km s}^{-1} \text{ kpc}^{-1}$ which implies a θ_0 value larger than the IAU one.

Reid et al. (2009, 2014) suggested also from the 3D motion measured for masers that U_\odot, V_\odot , and W_\odot must be updated and that the particular motion of the sources must be taken into account. Unfortunately, for Hi-GAL sources, we will have no

information about their own U, V, W . We therefore assume them to be null. Several other authors suggested alternative values to the standard ones for $U_\odot, V_\odot, W_\odot$, and R_0, θ_0 (we refer to Hou & Han 2014, for a revue on this). If two main sets of “ R_0, θ_0 ” emerge (the IAU standard one and the $R_0, \theta_0 = 8.34$ kpc, 240 km s^{-1}) for U_\odot, V_\odot , and W_\odot , no general agreement, especially for V_\odot , is brought out. However, because in the frame of the Hi-GAL survey we use source velocity extracted from different l, b, V_{LSR} data cubes (Molinaro et al. 2016) we can expect that, by default, their V_{LSR} is calculated with the $U_\odot, V_\odot, W_\odot$ IAU standard.

4. The updated version of the Galactic rotation curve

To produce an updated version of the rotation curve, we used the data listed above and fitted different analytical expressions. From the data, we define three distinct sub-samples: one combining samples 1, 2 and 3 (“Sub 123”), one combining samples 1 and 2 (“Sub 12”) and one combining samples 1, 2 and 4 (“Sub 124”). To avoid redundancy, in “Sub 123” we only add the 74 HII regions of sample 1 not in common with the ones already used in sample 3. In “Sub 12” the HII regions (all the regions from sample 1) and masers are put together to probe the rotation curve as traced by a similar method (velocity independent of the distance calculation). To avoid redundancy, we do not combine samples 3 and 4. However, as mentioned by McClure-Griffiths & Dickey (2016) comparing CO and HI data, we checked the good agreement between them.

Figure A.1 shows the different sub-samples and the typical error bars. By default, the tangent method always gives very small error bars and smaller scattering with respect to the HII regions and masers. In parallel, the galactocentric range around 8.5 kpc is naturally well populated by HII regions because at larger distance from the Sun the extinction no longer allows stellar distance determination. In the $R \sim 6$ kpc to 8 kpc the tangent method data points and HII regions agree, while below 6 kpc, the masers show a clear offset from them.

In the literature, several expressions for the rotation curve are used:

- A polynomial form: $\theta(R)/\theta_0 = a_1 + a_2 r + a_3 r^2$ with $r = (R/R_0) - 1$ used by Reid et al. (2014).
- Power law forms: Honma et al. (2012) and Brand & Blitz (1993) used the following power law forms $\theta(R)/\theta_0 = a_1 (R/R_0)^{a_2}$ and $\theta(R)/\theta_0 = a_1 (R/R_0)^{a_2} + a_3$, respectively.
- A universal form: Persic et al. (1996) suggest a more universal form $\theta(R)/\theta_0 = a_1 [1 + a_2 ((R/a_3) - 1)]$ based on a sample of extragalactic rotation curves.
- The Polyex model: Giovanelli & Haynes (2002) used, to fit rotation curves for 2246 galaxies, another universal analytical expression (known as the “polyex” model) with the form $\theta(R)/\theta_0 = (1 - e^{-R/a_1}) \times (1 + (a_2 R/a_1))$.

Before performing the fit we scale the data to the chosen R_0, θ_0 set, following, for example, Xin & Zheng (2013). To compare with the old and new results, we performed the fit on the three sub-samples with both $R_0 = 8.5$ kpc, $\theta_0 = 220 \text{ km s}^{-1}$ and $R_0 = 8.34$ kpc, $\theta_0 = 240 \text{ km s}^{-1}$ sets. In practice, following Fich et al. (1989), the rotation curves are fitted in ω versus R , because they are observationally independent quantities. We requested also that the cataloged objects have $|b| < 3^\circ$ (objects with larger latitude are probably close objects for which the systemic velocity can be distorted by local motions) and

$R > 4$ kpc because closer to the Galactic center the contribution of the bulge and the bar to the kinematics can become important (e.g. Chemin et al. 2015; Reid et al. 2014). The fits are done minimizing the normalised weighted χ^2 expression (where the weight is the inverse of the squared uncertainty) using the “Minit” subroutine (Nelder & Mead 1965). The different fitted solutions are shown in Figs. A.2 and A.3 for the both adopted R_0 , θ_0 sets, while the results are listed in Table A.1. Whatever the model, the value of the standard normalised χ^2 we found is small (between 0.09 and 0.45) while it is expected around unity. A value less than one does not necessary indicate a better fit but underlines uncertainty in the determination of the variance (Bevington & Robinson 2003). Despite the fact that it gives indication about the data dispersion around the fitted curve, Fich et al. (1989) show that the χ^2 numerical value cannot be easily used to sort the goodness of fit analysis. We then also used the evaluated parameter uncertainties³ to compare the fits because they are related to the width of the minimized expression minimum.

Looking at the figures we note the all the fitted curves give similar results except the polynomial form which, for several configurations, departs from the others after $R = 8$ kpc. From “Sub12”, we note also that all the models fit well the data in the range 6 to 10 kpc, while below 6 kpc they are not able to fit the regions offset from the tangent data. This underlines the fact that adding the tangent data will not impact too much the fitted curve in such inner parts and that fixing R_0 , θ_0 strongly forces the fit in this distance range.

5. Discussion

To compare our results with previous studies we have to keep in mind that we fit with fixed R_0 and θ_0 , as did McClure-Griffiths & Dickey (2016) and Levine et al. (2008), for example, while Reid et al. (2014) and Honma et al. (2012), for example, have them as free parameters. In addition, Reid & Dame (2016) found that a slightly curved rotation curve with $\theta_0 = 240 \text{ km s}^{-1}$ can mimic a flat rotation curve with $\theta_0 = 220 \text{ km s}^{-1}$, convincing us to fit with fixed R_0 and θ_0 . However fixing R_0 and θ_0 implies some expected relation between the fitted parameters. For example, $a_1 \sim 1$ is expected for Polynomial and Power-law models while a_3 and a_1 close to R_0 is expected for the Universal and Polyex models, respectively. Any large departure from these expected values underlines a least good fit, however, for most of the fitted parameters and the expected values are in agreement.

A possible approach followed by several authors to study the rotation curve is to simply fit a linear expression of the form $\theta(R)/\theta_0 = a_1 + a_2 (R/R_0)$. Some authors fit only the inner part of the rotation curve (3 to 8 kpc), because this is traced from the tangent velocity method. This is the case for Fich et al. (1989) for quadrant I, Levine et al. (2008) for quadrants I and IV and McClure-Griffiths & Dickey (2016) for quadrant I, who find $(a_1, a_2) = (0.887, 0.186)$, $(0.855, 0.024)$, $(0.829, 0.026)$, $(0.82, 0.026)$, respectively. Others fit a linear rotation curve to data within 4 and 16 kpc (Reid et al. 2014) or even within 8 and 11 kpc (Huang et al. 2016) finding $(a_1, a_2) = (1.007, -8.3 \times 10^{-3})$ and $(1.23, -0.023)$, respectively. We note that focussing on the inner rotation curve gives a positive slope while fitting up to a larger radius changes the slope to negative (but close to zero value).

We first test the polynomial model. In this model $a_1 \times \theta_0$ gives the overall amplitude of the rotation curve, and a_2 and a_3

describe the position of the curve extremum (with respect to R_0) and the curve inflection, respectively. For a decreasing curve, a_3 must be negative while the smaller $|a_3|$ is, the flatter the curve. Reid et al. (2014), fitting such a form to masers, found $a_1, a_2, a_3 = 1, 0.002, -0.06$ (following our parameter definitions). Focussing on the two curves of “Sub12”, we find a similar a_1 but a systematically larger $|a_3|$ suggesting our curves are more decreasing. However, we can note that the polynomial form consistently gives the worst χ^2 with respect to the other forms as can also be seen in Figs. A.2 and A.3.

In the “polyex” expression, Giovanelli & Haynes (2002) describe a_1 as the scale length for the inner steep rise (the radial distance at which θ is 0.63 of the asymptotic velocity for a flat rotation curve) while a_2 sets the slope of the rotation curve’s outer part. From approximately 2200 low-redshift galaxies, Catinella et al. (2006) show that $0.002 < a_2 < 0.087$. However, for galaxies with maximum velocity around 220–240 km s^{-1} , a_2 is expected to be between approximately 0.003 and 0.006. Similar values are found for “Sub 123” and “Sub 124” while for “Sub 12” it is found negative but still close to zero. In Catinella et al. (2006), a_1 is in units of exponential disk scale length (R_d) and for galaxies with maximum velocity around 220–240 km s^{-1} , it is estimated around 0.45. With a disk scale length for our Galaxy between 2.15 kpc (Bovy & Rix 2013) and 3.19 kpc (Sofue 2012) we expect a_1 to be between 1.07 and 1.43 kpc. But whatever our sample, a_1 is found around 1 kpc and even smaller for “Sub 12” (with the IAU R_0, θ_0).

In the Universal law expression, $a_1 \times \theta_0$ is the maximum velocity, a_3 is the radius where this maximum velocity (also noted R_{max}) is reached and a_2 is the velocity variation between a_3 and R_{opt} . As expected, a_1 is close to 1 whatever the sample. Persic et al. (1996) find $-0.1 \leq a_2 \leq 0.6$. We find such values for all our fitting configurations. From Persic et al. (1996), usually $a_3 \sim 2.2 \times R_d$, which suggest a_3 between 4.7 kpc and 7 kpc for our Galaxy. For $R_0 = 8.5$ kpc, a_3 is found very close to R_0 while for $R_0 = 8.34$ kpc a departure is noted reaching 40% for “Sub 123”. Reid et al. (2014) also fit such a Universal law to masers and find $a_1, a_2, a_3 = 1, 0.003, 12.13$ kpc (following our parameter definitions and with their $R_0, \theta_0 = 8.31$ kpc and 241 km s^{-1}) which suggests a flatter rotation curve.

In the power law forms, a_2 , the exponent, describes how quickly the curve decreases/increases while a_3 is the deviation term, which represents a simple way for observations to deviate from the power-law function. We note that $a_1 + a_3 \sim 1$, as expected for the Brand & Blitz (1993) form, is well recovered for IAU R_0, θ_0 while it is slightly smaller for $R_0, \theta_0 = 8.34$ kpc and 240 km s^{-1} . In addition to their HII regions catalog, Brand & Blitz (1993) added HI tangent velocities to compute a rotation curve (with $R_0, \theta_0 = 8.5$ kpc, 220 km s^{-1}) with the power law form and found $a_1, a_2, a_3 = 1.00767, 0.0394, 0.00712$ similar to our results obtained for “Sub 123” and “Sub 124”.

6. Conclusion

Testing different analytical forms to different samples we find that all the forms, except the polynomial one, give satisfactory fitting results. The two power-law forms are often superimposed. The models used in extra-galactic studies (universal and polyex forms) are also tested. However they implement a radius scaling parameter which is difficult to relate to R_0 as is the case for the other form. Using a stellar/maser distance sample gives more departures, even if the fits are good, between the different forms; in particular in the outer part of the rotation curve. In addition,

³ See the Minit reference manual at <https://root.cern.ch/sites/d35c7d8c.web.cern.ch/files/minuit.pdf>

in the inner part, the fitted curve is not able to pass through the data points but passes at the expected location of the tangent point as plotted by “Sub 123” and “Sub 124”. The power-law appears then as the simplest and easiest form describing the rotation curve, and because the two power forms are often superimposed, we favour the simplest form (with no a_3). In the frame of the kinematic distance determination of the Hi-GAL sources, we then adopt the power law form $\theta(R)/\theta_0 = 1.022 (R/R_0)^{0.0803}$ with $R_0, \theta_0 = 8.34$ kpc, 240 km s^{-1} .

To improve the Galactic rotation curve it appears important to better sample its outer part ($R > 10$ kpc). In this framework, we expect that the incoming ESA-*Gaia* database and maser parallactic distances will provide such information. Indeed, the ESA-*Gaia* database should provide better distance and rotation velocity determination (and velocity field) for the OB stars exciting the H II regions and identify and quantify the circular velocity departures of such regions (in the detection limits). It will also allow one to trace the galactic rotation curve (and the velocity field) given by the stellar background potential and to compare it to the observed one, as it is expected (Gómez 2006) that the observed rotation curve is systematically above the true one. Maser parallax distances appear also as a very accurate and promising tool (and complementary to *Gaia*) for directly determining the distance of the star-forming regions in which Hi-GAL sources are located.

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Appendix A: Figures and table

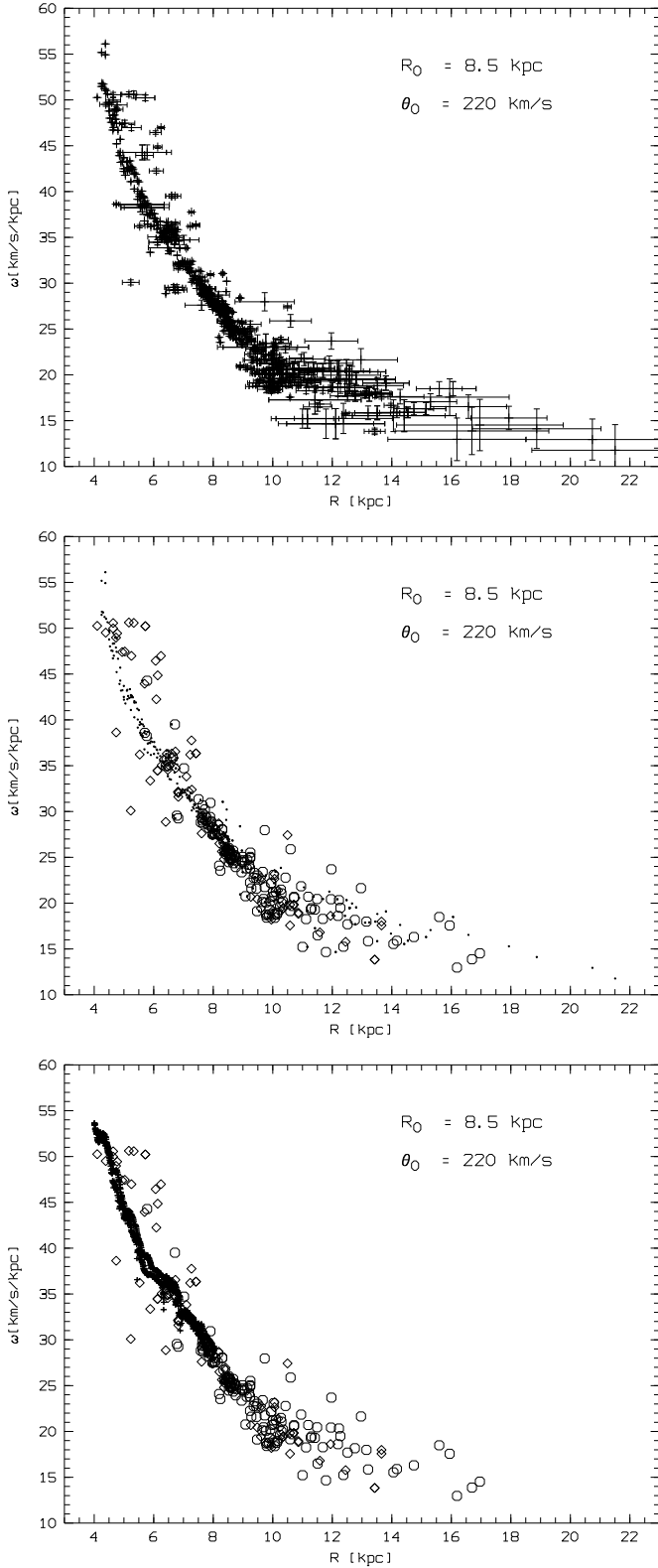


Fig. A.1. Data sample used for the rotation curve fitting. The *upper panel* shows the error bars. The *middle panel* shows sources from sample 1 (circles), sample 2 (diamonds), and sample 3 (dots), respectively. The *lower panel* shows sample 4 (dots) instead of sample 3 while the other symbols are similar as in middle panel.

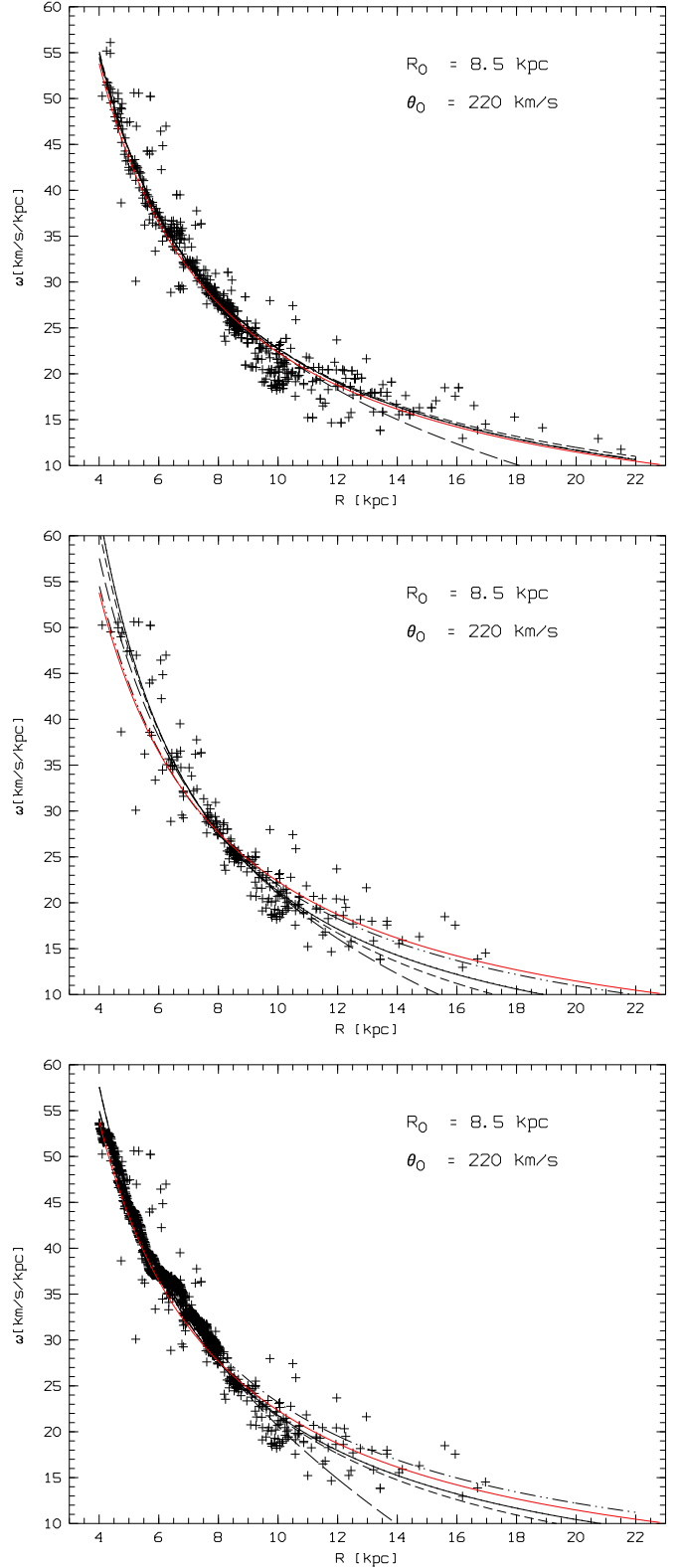


Fig. A.2. Fitted rotation curves with R_0 , θ_0 set to 8.5 kpc and 220 km s^{-1} , respectively. The fitted Brand & Blitz (1993), power law, polynomial, Universal and polyex forms are displayed as solid, dotted, long dash, short dash, and dash-dot lines, respectively. The Brand & Blitz (1993) rotation curve is superimposed (red line).

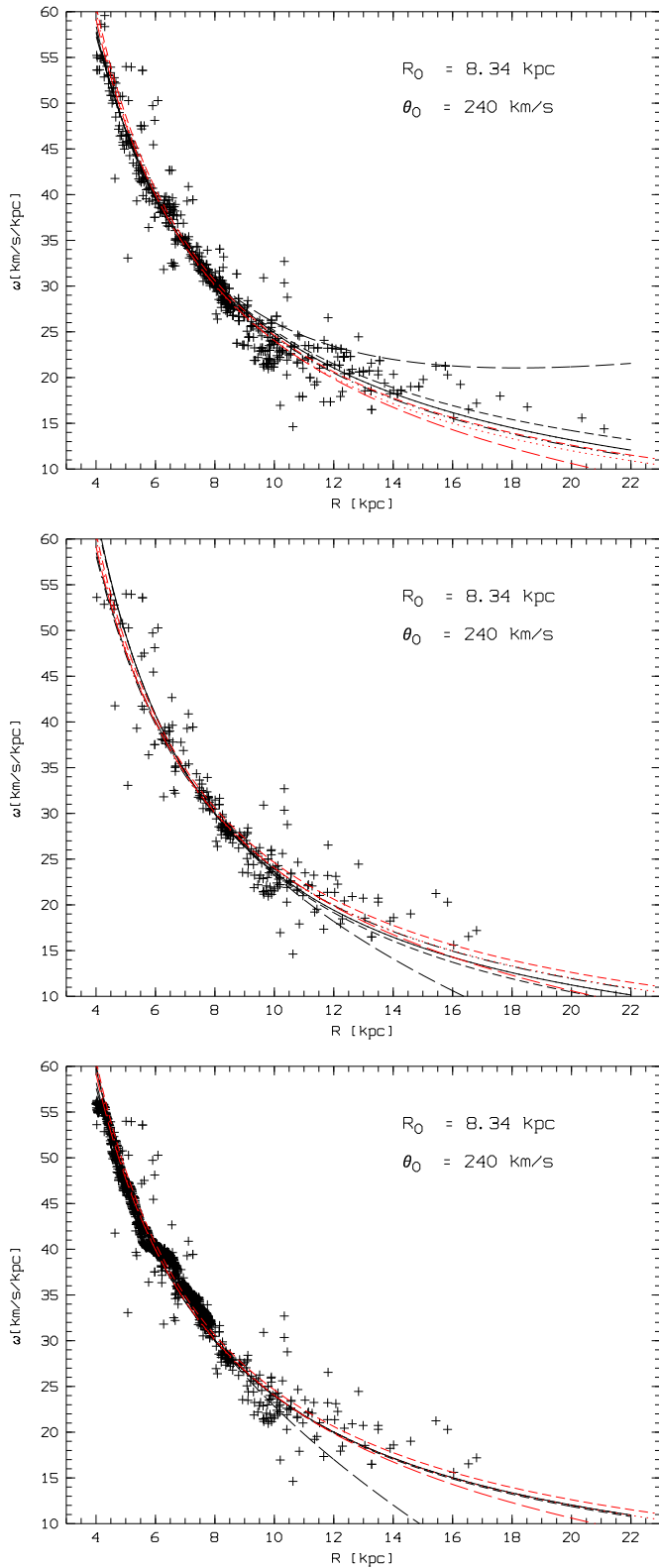


Fig. A.3. As in Fig. A.2 but for R_0 , θ_0 set to 8.34 kpc and 240 km s⁻¹ respectively. The Reid et al. (2014) polynomial (long dashes), power-law (dotted line) and universal (short dashes) rotation curves are superimposed (in red).

Table A.1. Fitting results.

| | Sub 123 | Sub 12 | Sub 124 |
|----------------------------------|--------------------------------------------------------------------------------------------------------------------------------------------------------------|---------------------------------------------------------------------------------------------------------------------------|--------------------------------------------------------------------------------------------------------------------------------------|
| 8.5/220. Poly. | $a_1 = 1.025 \pm 4 \times 10^{-6}$ $a_2 = -0.013 \pm 6 \times 10^{-5}$ $a_3 = -0.147 \pm 1 \times 10^{-4}$ $\chi^2 = 0.213$ | $a_1 = 0.996 \pm 6 \times 10^{-4}$ $a_2 = -0.204 \pm 0.011$ $a_3 = -0.198 \pm 0.043$ $\chi^2 = 0.422$ | $a_1 = 0.998 \pm 5 \times 10^{-4}$ $a_2 = -0.259 \pm 8 \times 10^{-3}$ $a_3 = -0.477 \pm 2 \times 10^{-2}$ $\chi^2 = 0.186$ |
| 8.5/220. Power law | $a_1 = 1.027 \pm 3 \times 10^{-6}$ $a_2 = 0.0352 \pm 2 \times 10^{-5}$ $\chi^2 = 0.175$ | $a_1 = 0.995 \pm 6 \times 10^{-4}$ $a_2 = -0.179 \pm 0.011$ $\chi^2 = 0.289$ | $a_1 = 1.000 \pm 5 \times 10^{-4}$ $a_2 = -0.062 \pm 4 \times 10^{-3}$ $\chi^2 = 0.099$ |
| 8.5/220. BB Power law | $a_1 = 1.156 \pm 1 \times 10^{-3}$ $a_2 = 0.031 \pm 3 \times 10^{-5}$ $a_3 = -0.128 \pm 1 \times 10^{-3}$ $\chi^2 = 0.185$ | $a_1 = 1.342 \pm 0.078$ $a_2 = -0.134 \pm 0.010$ $a_3 = -0.346 \pm 0.052$ $\chi^2 = 0.288$ | $a_1 = 1.091 \pm 3 \times 10^{-2}$ $a_2 = -0.057 \pm 5 \times 10^{-3}$ $a_3 = -0.09 \pm 8 \times 10^{-2}$ $\chi^2 = 0.099$ |
| 8.5/220. Univ. law | $a_1 = 1.028 \pm 7 \times 10^{-5}$ $a_2 = 0.044 \pm 7 \times 10^{-5}$ $a_3 = 8.5 \pm 0.013$ $\chi^2 = 0.175$ | $a_1 = 0.995 \pm 0.033$ $a_2 = -0.206 \pm 0.040$ $a_3 = 8.49 \pm 1.36$ $\chi^2 = 0.316$ | $a_1 = 0.999 \pm 2 \times 10^{-2}$ $a_2 = -0.095 \pm 2 \times 10^{-2}$ $a_3 = 8.49 \pm 1.68$ $\chi^2 = 0.099$ |
| 8.5/220. Polyex Univ. law | $a_1 = 1.0091 \pm 2 \times 10^{-4}$ $a_2 = 0.0032 \pm 9 \times 10^{-7}$ $\chi^2 = 0.174$ | $a_1 = 0.785 \pm 0.203$ $a_2 = -0.00015 \pm 6 \times 10^{-5}$ $\chi^2 = 0.292$ | $a_1 = 1.053 \pm 0.400$ $a_2 = 0.0058 \pm 1.5 \times 10^{-3}$ $\chi^2 = 0.075$ |
| 8.34/240. Poly | $a_1 = 1.032 \pm 2 \times 10^{-5}$ $a_2 = 0.221 \pm 2 \times 10^{-4}$ $a_3 = 0.216 \pm 4 \times 10^{-4}$ $\chi^2 = 0.215$ | $a_1 = 0.996 \pm 6 \times 10^{-4}$ $a_2 = -0.090 \pm 0.011$ $a_3 = -0.244 \pm 0.042$ $\chi^2 = 0.452$ | $a_1 = 0.998 \pm 5 \times 10^{-4}$ $a_2 = -0.149 \pm 8 \times 10^{-3}$ $a_3 = -0.422 \pm 2 \times 10^{-2}$ $\chi^2 = 0.160$ |
| 8.34/240. Power law | $a_1 = \mathbf{1.022} \pm \mathbf{2} \times \mathbf{10}^{-5}$ $a_2 = \mathbf{0.0803} \pm \mathbf{4} \times \mathbf{10}^{-5}$ $\chi^2 = \mathbf{0.169}$ | $a_1 = 0.996 \pm 5 \times 10^{-4}$ $a_2 = -0.069 \pm 0.010$ $\chi^2 = 0.303$ | $a_1 = 0.999 \pm 5 \times 10^{-4}$ $a_2 = 0.032 \pm 3 \times 10^{-3}$ $\chi^2 = 0.095$ |
| 8.34/240. BB Power law | $a_1 = 0.838 \pm 3 \times 10^{-4}$ $a_2 = 0.0989 \pm 5 \times 10^{-5}$ $a_3 = 0.184 \pm 3 \times 10^{-4}$ $\chi^2 = 0.169$ | $a_1 = 1.000 \pm 8 \times 10^{-3}$ $a_2 = -0.069 \pm 0.010$ $a_3 = -0.004 \pm 8 \times 10^{-3}$ $\chi^2 = 0.301$ | $a_1 = 0.937 \pm 7 \times 10^{-3}$ $a_2 = 0.003 \pm 4 \times 10^{-3}$ $a_3 = 0.063 \pm 7 \times 10^{-3}$ $\chi^2 = 0.095$ |
| 8.34/240. Univ. law | $a_1 = 1.071 \pm 6 \times 10^{-4}$ $a_2 = 0.148 \pm 5 \times 10^{-4}$ $a_3 = 11.69 \pm 0.04$ $\chi^2 = 0.167$ | $a_1 = 0.992 \pm 0.026$ $a_2 = -0.092 \pm 0.030$ $a_3 = 8.70 \pm 2.5$ $\chi^2 = 0.313$ | $a_1 = 0.999 \pm 4 \times 10^{-3}$ $a_2 = -0.629 \pm 5 \times 10^{-3}$ $a_3 = 7.73 \pm 4.98$ $\chi^2 = 0.097$ |
| 8.34/240. Polyex Univ. law | $a_1 = 1.35 \pm 2 \times 10^{-4}$ $a_2 = 0.0028 \pm 3 \times 10^{-6}$ $\chi^2 = 0.173$ | $a_1 = 1.12 \pm 0.11$ $a_2 = -0.00028 \pm 8 \times 10^{-5}$ $\chi^2 = 0.291$ | $a_1 = 1.125 \pm 3 \times 10^{-2}$ $a_2 = 1.23 \times 10^{-3} \pm 7 \times 10^{-5}$ $\chi^2 = 0.090$ |

Notes. Values in bold is our final choice.