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## First Detection of Interstellar S<sub>2</sub>H

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### Abstract

We present the first detection of gas phase S<sub>2</sub>H in the Horsehead, a moderately UV-irradiated nebula. This confirms the presence of doubly sulfuretted species in the interstellar medium and opens a new challenge for sulfur chemistry. The observed S<sub>2</sub>H abundance is  $\sim 5 \times 10^{-11}$ , only a factor 4-6 lower than that of the widespread H<sub>2</sub>S molecule. H<sub>2</sub>S and S<sub>2</sub>H are efficiently formed on the UV-irradiated icy grain mantles. We performed ice irradiation experiments to determine the H<sub>2</sub>S and S<sub>2</sub>H photodesorption yields. The obtained values are  $\sim 1.2 \times 10^{-3}$  and  $< 1 \times 10^{-5}$  molecules per incident photon for H<sub>2</sub>S and S<sub>2</sub>H, respectively. Our upper limit to the S<sub>2</sub>H photodesorption yield suggests that photo-desorption is not a competitive mechanism to release the S<sub>2</sub>H molecules to the gas phase. Other desorption mechanisms such as chemical desorption, cosmic-ray desorption and grain shattering can increase the gaseous S<sub>2</sub>H abundance to some extent.

Alternatively, S<sub>2</sub>H can be formed via gas phase reactions involving gaseous H<sub>2</sub>S and the abundant ions S<sup>+</sup> and SH<sup>+</sup>. The detection of S<sub>2</sub>H in this nebula could be therefore the result of the coexistence of an active grain surface chemistry and gaseous photo-chemistry.

## Keywords

Astrochemistry; methods: laboratory: solid state; ISM: abundances; ISM: molecules; photon-dominated region (PDR); Horsehead

## 1 Introduction

Sulfur is one of the most abundant elements in the Universe ( $S/H \sim 1.3 \times 10^{-5}$ ) and plays a crucial role in biological systems on Earth, so it is important to follow its chemical history in space. Surprisingly, sulfuretted molecules are not as abundant as expected in the interstellar medium. A few sulfur compounds have been detected in diffuse clouds demonstrating that the sulfur abundance in these low density regions is close to the cosmic value (Neufeld et al. 2015). A moderate sulfur depletion (a factor of 4) is observed in the external layers of the photodissociation region (PDR) in the Horsehead nebula, as well (Goicoechea et al. 2006). In cold molecular clouds, a large depletion of sulphur is usually considered to reproduce the observations (see for instance Tieftrunk et al. 1994). However, recent models by Vidal et al. (2017) explained the observational data without or with little sulfur depletion after updating the gas and grain chemistry. In that case, HS and H<sub>2</sub>S on the grains or atomic sulphur in the gas would contain most of the sulfur. In hot cores and corinos, Wakelam et al. (2004) found that observations of S-bearing molecules would be better reproduced if sulfur was sublimated from grains in the atomic form or it is quickly converted into it. Thus far, the main solid or gas sulfur carrier is still debated.

With an adsorption energy of  $\sim 1100$  K (Hasegawa & Herbst 1993), sulfur atoms are expected to stick on the surfaces of grains with temperatures below  $\sim 22$  K. Here, because of the high hydrogen abundances and the mobility of hydrogen in the ice matrix, sulfur atoms are expected to form H<sub>2</sub>S. Indeed, gaseous H<sub>2</sub>S is the most abundant S-bearing molecule in comets, with an abundance of up to 1.5% relative to water (Bockelée-Morvan et al. 2000). A firm detection of H<sub>2</sub>S in interstellar ices has not been reported yet. A realistic upper limit of the H<sub>2</sub>S abundance in interstellar ices is 1% relative to water that is 10 times lower than the cosmic abundance (Jiménez-Escobar & Muñoz Caro 2011). One possibility to explain the low fraction of H<sub>2</sub>S in interstellar ices is that H<sub>2</sub>S is processed by UV-photons or cosmic rays in the ice leading to the formation of other S-bearing species. OCS and tentatively SO<sub>2</sub> have been detected in icy mantles but their abundances are far too low to explain the missing S budget (Geballe et al. 1985; Palumbo et al. 1995; Boogert et al. 1997).

Trying to provide new insights into the ice composition, experimental simulations of the irradiation of interstellar ices containing H<sub>2</sub>S under astrophysically relevant conditions have been performed in laboratory using UV photons (Jiménez-Escobar & Muñoz Caro 2011; Jiménez-Escobar et al. 2014), X-rays (Jiménez-Escobar et al. 2012), or ions (Moore et al. 2007; Ferrante et al. 2008; Garozzo et al. 2010). Energetic processing of H<sub>2</sub>S-bearing ices readily generates sulfur-sulfur bonds, and the main S-bearing products in these experiments

are  $\text{H}_2\text{S}_2$  and  $\text{S}_2\text{H}$  that were detected by Jiménez-Escobar & Muñoz Caro (2011) through their infrared absorption bands. The molecule  $\text{H}_2\text{S}_2$  could subsequently photodissociate forming  $\text{S}_2$  and  $\text{S}_3$  depending on the irradiation time. These molecules with two S atoms and even more could thus contain a significant fraction of the missing sulfur in dense clouds. In line with this work, Druard & Wakelam (2012) suggested that polysulphanes could be a sulfur reservoir in the ice and are rapidly converted into atomic sulfur once in the gas phase. Martín-Doménech et al. (2016a) unsuccessfully searched for  $\text{S}_2\text{H}$  and  $\text{H}_2\text{S}_2$  in the gas phase toward the well-known hot corino, IRAS 16293–2422. The lack of gaseous  $\text{S}_2\text{H}$  and  $\text{H}_2\text{S}_2$  was interpreted as the consequence of the rapid destruction of these species once sublimated in such a warm and dense environment (Martín-Doménech et al. 2016a; Fortenberry & Francisco 2017).

In this Letter, we report the first interstellar detection of  $\text{S}_2\text{H}$  in the prototypical photo-dissociation region, the Horsehead. In Sect. 4, we discuss the possible grain-surface and gas-phase  $\text{S}_2\text{H}$  formation routes. New measurements of the photodesorption yields of  $\text{S}_2\text{H}$  and  $\text{H}_2\text{S}$  are presented in Sect. 5.

## 2 Observations and Data Reduction

The data used in this work are from the Horsehead WHISPER (Wide-band High-resolution Iram-30m Surveys at two Positions with Emir Receivers, PI: J. Pety) project and the Director's Discretionary Time project D11-16. The Horsehead WHISPER project is a complete unbiased line survey of the 3, 2, and 1 mm bands using the IRAM 30m telescope. Two positions are observed: i) the HCO peak (RA= $5^{\text{h}}40^{\text{m}}53^{\text{s}}.936$ , Dec= $2^{\circ}28'00''$ , J2000), which is characteristic of the photo-dissociation region at the UV-illuminated surface of the Horsehead nebula (Gerin et al. 2009) (also referred to as PDR position), and ii) the DCO<sup>+</sup> peak (RA= $5^{\text{h}}40^{\text{m}}55^{\text{s}}.61$ , Dec= $2^{\circ}27'38''$ , J2000), which corresponds to a cold and UV-shielded condensation located less than  $40''$  away from the PDR edge (Pety et al. 2007). During the observations we used the Position-Switching procedure with the reference position located at an offset ( $-100'', 0$ ) relative to RA:  $05^{\text{h}}40^{\text{m}}54^{\text{s}}.27$  Dec:  $-02^{\circ}28'00''.0$ . Several lines of  $\text{S}_2\text{H}$  were tentatively detected towards the two positions observed in the WHISPER survey. In order to confirm the  $\text{S}_2\text{H}$  detection, we requested Director's Discretionary Time (D11-16) to observe a single setup covering the frequencies listed in Table 1. The merged  $\text{S}_2\text{H}$  spectra are shown in Fig. 1. Line intensities are given in main brightness temperature ( $T_{\text{MB}}$ ) and the lines were observed with a frequency resolution of 49 kHz.

In order to have a deeper insight into the  $\text{S}_2\text{H}$  chemistry, we compare the new  $\text{S}_2\text{H}$  observations with the  $\text{H}_2\text{S}$   $1_{1,0} \rightarrow 1_{0,1}$  map observed during April 2006 with the IRAM 30m telescope. These observations were done using the frequency switching mode and a spectral resolution of 40 kHz. Averaged noise level per resolution element at 168 GHz is  $\text{rms}(T_{\text{MB}})=170$  mK. The integrated intensity emission of the  $\text{H}_2\text{S}$  line varies between  $1.0^{-1.5}$  K km s<sup>-1</sup> across the molecular cloud with an abrupt border in the west (see Fig. 1). The  $\text{H}_2\text{S}$  emission presents a local minimum towards the DCO<sup>+</sup> peak, similar to the morphology observed in other species such as  $\text{CH}_3\text{OH}$  (Guzmán et al. 2011, 2013), suggesting gaseous  $\text{H}_2\text{S}$  depletion towards this cold dense core.

### 3 Column Densities and Abundances

The rotational spectrum of S<sub>2</sub>H was calculated by Tanimoto et al. (2000). The spectroscopic data can be found in the CDMS catalogue (Müller et al. 2005). We have detected eight S<sub>2</sub>H lines located at 94526.1508, 94526.3208, 94731.0115, 94731.2080, 110294.0282, 110294.1530, 110497.9666, and 110498.1104 MHz. The S<sub>2</sub>H hyperfine transitions are forming doublets very close in frequency (~0.12 MHz) that remain unresolved in our data (see Fig. 1). In Table 1, we show the Gaussian fits to the observed line features, each one clearly detected with S/N>5. We have adopted as central frequency the one of the most intense component of the doublet. For this reason the central velocity shown in Table 1 is systematically shifted by ~0.2–0.5 km s<sup>-1</sup> from the Horsehead systemic velocity, 10.5 km s<sup>-1</sup>. We have checked possible contamination by other compounds using the CDMS and JPL catalogues. There is no other good candidate to be a carrier of these lines. The large linewidth of the 94.731 GHz line towards the DCO<sup>+</sup> peak position, ~1.7 km s<sup>-1</sup>, is more likely due to the poor baseline around this feature.

The WEEDS software has been used to simulate the S<sub>2</sub>H spectrum in the whole frequency coverage of the WHISPER survey assuming LTE conditions. We have fitted the detections and upper-limits of 97 lines with upper level energies lower than 75 K found in the frequency range of the full survey using the Bayesian method described by Majumdar et al. (2017). The whole spectrum can be fitted assuming that the emission uniformly fills the beam and the rotation temperatures and S<sub>2</sub>H column densities listed in Table 1. The fitted line-widths are 0.68±0.12 km s<sup>-1</sup> for the core and 0.63±0.1 km s<sup>-1</sup> for the PDR. The observed S<sub>2</sub>H linewidths are consistent with the emission coming from the UV irradiated gas. Species that are more abundant in the cold and UV shielded gas of the core as DCO<sup>+</sup> and H<sup>13</sup>CO<sup>+</sup>, present narrower line-widths towards the DCO<sup>+</sup> peak than towards the PDR (Goicoechea et al. 2009). However, others PDR-like species such as HCO present similar linewidths towards both positions (Gerin et al. 2009). This suggests that even towards the core position, the S<sub>2</sub>H emission is mainly coming from the UV-illuminated layers of the cloud along the line of sight.

The S<sub>2</sub>H rotation temperatures reveal subthermal excitation and are similar to those derived for other high dipole moment compounds like o-H<sub>2</sub>CO (Guzmán et al. 2011). The estimated abundance (wrt hydrogen nuclei) is ~5×10<sup>-11</sup> towards the DCO<sup>+</sup> peak and about a factor of 2 larger towards the HCO peak.

From the chemical point of view, it is interesting to compare the S<sub>2</sub>H abundance with those of the related species H<sub>2</sub>S. Unfortunately there is only one transition of H<sub>2</sub>S that is easily observable with the 30m telescope given the physical conditions in the Horsehead, the o-H<sub>2</sub>S 1<sub>1,0</sub>→1<sub>0,1</sub> line. Thus, we need to assume a rotation temperature to derive the H<sub>2</sub>S column density. Since the H<sub>2</sub>S dipole moment ( $\mu_b=0.978$  D; Viswanathan et al. 1984) is similar to those of S<sub>2</sub>H ( $\mu_a=1.161$  D,  $\mu_b=0.827$  D; Peterson et al. 2008), we assume the same rotation temperature for both molecules. With these assumptions and adopting an ortho-to-para ratio of 3, we derive a H<sub>2</sub>S abundance of ~3×10<sup>-10</sup> towards the two positions. This would imply that [S<sub>2</sub>H]/[H<sub>2</sub>S]=0.15±0.09 in the DCO<sup>+</sup> peak and [S<sub>2</sub>H]/[H<sub>2</sub>S]=0.27±0.14, toward the HCO peak. These numbers are consistent with the [S<sub>2</sub>H]/[H<sub>2</sub>S] ice ratio obtained

by Jiménez-Escobar et al. (2012) in their simulations of UV irradiation of H<sub>2</sub>S ices. In the following, we qualitatively explore the possible surface and gas-phase formation routes of S<sub>2</sub>H.

## 4 S<sub>2</sub>H Formation

The formation of S<sub>2</sub>H is a intricate problem due to the low H-SS energy bonding. Evidences for the formation of S<sub>2</sub>H during irradiation of pure H<sub>2</sub>S and H<sub>2</sub>S:H<sub>2</sub>O ice mixtures were provided by Jiménez-Escobar & Muñoz Caro (2011) using the same experimental setup as the one described here. One way of forming S<sub>2</sub>H could be the grain surface reactions: s-H atom (hereafter, 's-' is used to refer to the solid phase) addition on s-S<sub>2</sub>, and s-S + s-HS reaction, followed by chemical desorption. However, the low exothermicity of the first reaction should prevent efficient chemical desorption (Minissale et al. 2016; Wakelam et al. 2017). The second reaction should not be efficient in cold cores because, below 15 K, S atom and HS radical are not mobile on ice considering the adsorption energies given by Wakelam et al. (2017). Moreover, S<sub>2</sub>H is a very reactive species in the gas phase, reacting with H, N, C and O atoms without barrier, so likely also on surface. An alternative surface induced S<sub>2</sub>H production may be s-H<sub>2</sub>S<sub>2</sub> photo-dissociation-desorption: s-H<sub>2</sub>S<sub>2</sub> + hν → S<sub>2</sub>H + H. s-H<sub>2</sub>S<sub>2</sub> may be efficiently produced by the s-HS + s-HS reaction but it needs mobile HS on ice and so a high grain temperature.

In the gas-phase, S<sub>2</sub>H may be produced by the electronic dissociative recombination of H<sub>2</sub>S<sub>2</sub><sup>+</sup>. Even if there is no data on this reaction, the loss of one H atom is always an important exit channel on dissociative recombination (Plessis et al. 2010). There are two known H<sub>2</sub>S<sub>2</sub><sup>+</sup> production pathways: the S<sup>+</sup> + H<sub>2</sub>S → H<sub>2</sub>S<sub>2</sub><sup>+</sup> + hν reaction (Anicich, V.G. 2003), despite the fact that the reference is an unpublished work and previous experimental studies did not identify this channel (Smith et al. 2004), and the SH<sup>+</sup> + H<sub>2</sub>S → H<sub>2</sub>S<sub>2</sub><sup>+</sup> + H reaction which is well characterized (Anicich, V.G. 2003). We note that S<sup>+</sup> and SH<sup>+</sup> are only abundant in the UV-irradiated gas (Gerin et al. 2016). Therefore, in spite of the large uncertainties in the reaction rates, we can conclude that these formation routes are only efficient in the UV-illuminated cloud surfaces.

## 5 Experimental Study of the Photodesorption of S<sub>2</sub>H and H<sub>2</sub>S

Jiménez-Escobar & Muñoz Caro (2011) showed that sulfur-sulfur bonds, in particular H<sub>2</sub>S<sub>2</sub> and S<sub>2</sub>H, are formed in irradiated ices. Here we focus on the determination of the S<sub>2</sub>H and H<sub>2</sub>S photodesorption yields which are key to determine the origin (surface vs gas phase chemistry) of the observed S<sub>2</sub>H. For this aim, we performed experimental simulations under astrophysically relevant conditions using the ISAC setup (Muñoz Caro et al. 2010), an ultra-high vacuum chamber with a work pressure on the order of 4×10<sup>-11</sup> mbar, corresponding to the pressure found in the interior of the pre-stellar cores. Sulfur is expected to be locked on the icy mantles in these regions, H<sub>2</sub>S being the most abundant S-bearing molecule in cometary ices. Pure amorphous H<sub>2</sub>S ice samples with thicknesses of about 40×10<sup>15</sup> molecules cm<sup>-2</sup> were deposited from the gas phase (H<sub>2</sub>S gas, Praxair, 99.8%) onto a KBr

substrate at 8 K, and subsequently irradiated using an F-type microwave-discharged hydrogen flow lamp with a vacuum-ultraviolet flux of  $2 \times 10^{14}$  photons  $\text{cm}^2 \text{s}^{-1}$  at the sample position (Muñoz Caro et al. 2010). The emission spectrum of the lamp (reported in Chen et al. 2014, and Cruz-Díaz et al. 2014) resembles that of the secondary UV field in dense cloud interiors, calculated by Gredel et al. (1989). A Pfeiffer Prisma quadrupole mass spectrometer (QMS) was used during irradiation of the ice samples to monitor the mass fragments  $m/z = 34$  (corresponding to photodesorbing  $\text{H}_2\text{S}$  molecules), and  $m/z = 64$  (corresponding to any desorbing photoproduct with a sulfur-sulfur bond, observed to form in Jiménez-Escobar et al. (2012)). In our experiment, we did not monitor  $\text{S}_2\text{H}$  directly. However, if  $\text{H}_2\text{S}_2$  or  $\text{S}_2\text{H}$  were desorbed, we would expect to detect all the fragments derived from these species, in particular  $\text{S}_2^+$ . While photodesorption of  $\text{H}_2\text{S}$  was detected, no gaseous  $\text{S}_2^+$  was observed (see Fig. 2). The measured ion current was converted into a photodesorption yield following calibration of the QMS (see Martín-Doménech et al. 2015). Photodesorption of  $\text{H}_2\text{S}$  took place with a decreasing yield, reaching a steady-state value of  $1.2 \times 10^{-3}$  molecules per incident photon after  $\sim 30$  minutes of irradiation, which corresponds to the fluence experienced by ice mantles during the typical cloud lifetime (Shen et al. 2004). A factor of 2 is assumed as the error in the photodesorption yield values due to the uncertainties in the calibration process, see Martín-Doménech et al. 2016b. Following the non-detection of any sulfur-sulfur photo-product, an upper limit of  $1 \times 10^{-5}$  molecules per incident photon (the sensitivity limit of our QMS) was assumed for the photodesorption of  $\text{S}_2\text{H}$ . Direct  $\text{S}_2\text{H}$  photodesorption or  $\text{H}_2\text{S}_2$  photo-dissociation-desorption are therefore not expected to be the origin of the gaseous  $\text{S}_2\text{H}$ .

## 6 Discussion and Conclusions

At a distance of 400 pc, the Horsehead is a PDR viewed nearly edge-on and illuminated by the O9.5V star  $\sigma$ Ori at a projected distance of  $\sim 3.5$  pc. The intensity of the incident FUV radiation field is  $\chi=60$  relative to the interstellar radiation field in Draine units. This PDR presents a differentiated chemistry from others associated with nearby HII regions such as the Orion Bar. One main difference is that the dust temperature is around  $\sim 20$ -30 K in the PDR (Goicoechea et al. 2009), i.e. below or close to the sublimation temperature of many species, allowing a rich surface chemistry on the irradiated surfaces. Our unbiased line survey has provided valuable hints on the chemistry of this region. The detection of the molecular ions  $\text{CF}^+$  and  $\text{HOC}^+$  towards the HCO peak are well understood in terms of gas-phase photochemistry (Guzmán et al. 2012). We learned that there is an efficient top-down chemistry in the PDR, in which large polyatomic molecules or small grains are photo-destroyed into smaller hydrocarbon molecules/precursors, such as  $\text{C}_2\text{H}$ ,  $\text{C}_3\text{H}_2$ ,  $\text{C}_3\text{H}$  and  $\text{C}_3\text{H}^+$  (Pety et al. 2012; Guzmán et al. 2015). The detection of several complex organic molecules (COMs) towards the warm ( $T_{\text{kin}} \sim 60$  K) PDR and its associated cold ( $T_{\text{kin}} \sim 20$  K) core was unexpected. In fact, the chemical complexity reached in the Horsehead is extraordinarily high with COMs of up to 7 atoms:  $\text{HCOOH}$ ,  $\text{H}_2\text{CCO}$ ,  $\text{CH}_3\text{CHO}$  and  $\text{CH}_3\text{CCH}$  (Guzmán et al. 2014). Current pure gas-phase models cannot reproduce the inferred  $\text{H}_2\text{CO}$ ,  $\text{CH}_3\text{OH}$  and COMs abundances in the Horse-head PDR (Guzmán et al. 2011, 2013), which supports the grain surface origin of these molecules. Le Gal et al. (2017) was able to reproduce the observed COMs abundances using a chemical model with grain

surface chemistry and found that chemical desorption, instead of photodesorption, is the main process to release COMs to the gas phase.  $\text{CH}_3\text{CN}$  and  $\text{CH}_3\text{NC}$ , key species for the formation of prebiotic molecules, seem to have a very specific formation pathway in the PDR (Gratier et al. 2013). The Horsehead is therefore an excellent site to study the influence of UV radiation on the grain surface chemistry and its subsequent impact on the gas phase.

We present the first detection of  $\text{S}_2\text{H}$  in the Horse-head. The observed  $\text{S}_2\text{H}$  abundance is  $\sim 5 \times 10^{-11}$ , only a factor 4-6 lower than that of  $\text{H}_2\text{S}$ . Our laboratory experiments show that the  $\text{H}_2\text{S}$  and  $\text{S}_2\text{H}$  photodesorption yields are  $1.2 \times 10^{-3}$  and  $< 1 \times 10^{-5}$  molecules per incident photon, respectively. Although  $\text{S}_2\text{H}$  can be formed on warm ( $T_d > 15$  K) grains, our upper limit to the  $\text{S}_2\text{H}$  photodesorption yield suggest that this mechanism is not efficient to release the  $\text{S}_2\text{H}$  molecules from the grain mantles. Other desorption mechanisms such as chemical desorption, cosmic-ray desorption and grain shattering could increase the  $\text{S}_2\text{H}$  abundance in gas phase.  $\text{S}_2\text{H}$  can also be formed in gas-phase by reactions involving  $\text{H}_2\text{S}$  and the ions  $\text{S}^+$  and  $\text{SH}^+$ . These ions are expected to be abundant in the external layers of the PDR (Goicoechea et al. 2006). The photodesorption of  $\text{H}_2\text{S}$  could hence boost the  $\text{S}_2\text{H}$  production in gas phase. We conclude that the abundance of  $\text{S}_2\text{H}$  in the Horsehead is more likely the consequence of the favorable physical conditions prevailing in this nebula where grain mantles irradiated by UV photons coexist with the ions  $\text{S}^+$  and  $\text{SH}^+$  that are only abundant in PDRs.

One interesting issue is to compare the sulfur and oxygen chemistry. We have not detected  $\text{H}_2\text{S}_2$ ,  $\text{HSO}$ ,  $\text{H}_2\text{O}_2$  and  $\text{HO}_2$  in the Horsehead with the upper limits shown in Table 3. We find interesting that the column densities of  $\text{HSO}$  and  $\text{HO}_2$  are lower than that of  $\text{S}_2\text{H}$ , although the oxygen elemental abundance is 30 times greater than that of sulfur. In gas phase,  $\text{S}_2\text{H}$  is mainly formed through  $\text{S}^+ + \text{H}_2\text{S} \rightarrow \text{H}_2\text{S}_2^+ + h\nu$  and  $\text{SH}^+ + \text{H}_2\text{S} \rightarrow \text{H}_2\text{S}_2^+ + \text{H}$  followed by dissociative recombination of  $\text{H}_2\text{S}_2^+$ . Oxygen and sulfur have indeed similar reactivity but, due to their different ionization potentials,  $\text{O}^+$  is expected less abundant than  $\text{S}^+$  and then the  $\text{O}^+$  and  $\text{OH}^+$  reactions play a smaller role. We have also compared the  $\text{SH}^+ + \text{H}_2\text{O}$  and  $\text{SH}^+ + \text{H}_2\text{S}$  gas-phase reactions which may be intermediate paths at work for producing  $\text{SOH}$  and  $\text{S}_2\text{H}$ , respectively. The channel towards  $\text{HSO}^+ + \text{H}_2$  reaction is endothermic in opposite to the channels towards  $\text{S}_2\text{H}^+ + \text{H}_2$  and  $\text{S}_2\text{H}_2^+ + \text{H}$ . Therefore, in gas phase the formation of  $\text{S}_2\text{H}$  is favored relative to  $\text{HSO}$ .  $\text{HSO}$  and related species have not been observed in space thus far (Cazzoli et al. 2016; Fortenberry & Francisco 2017). Laboratory experiments demonstrate that grain surface chemistry involving  $\text{H}_2\text{O}$  and  $\text{H}_2\text{S}$  also present different pathways. Photo-desorption experiments reported by Cruz-Diaz et al. (2017) show that  $\text{H}_2\text{O}_2$  is not formed in UV irradiated water ice. In contrast, Jiménez-Escobar & Muñoz Caro (2011) showed that  $\text{H}_2\text{S}_2$  is formed when a  $\text{H}_2\text{S}$  and  $\text{H}_2\text{S}-\text{H}_2\text{O}$  ices are irradiated, providing a path to form species with two sulfur atoms. Summarizing, sulfur and oxygen are not analogues in the gas-phase and surface chemistry, and the comparison of their related species requires the full chemical modelling of the region.



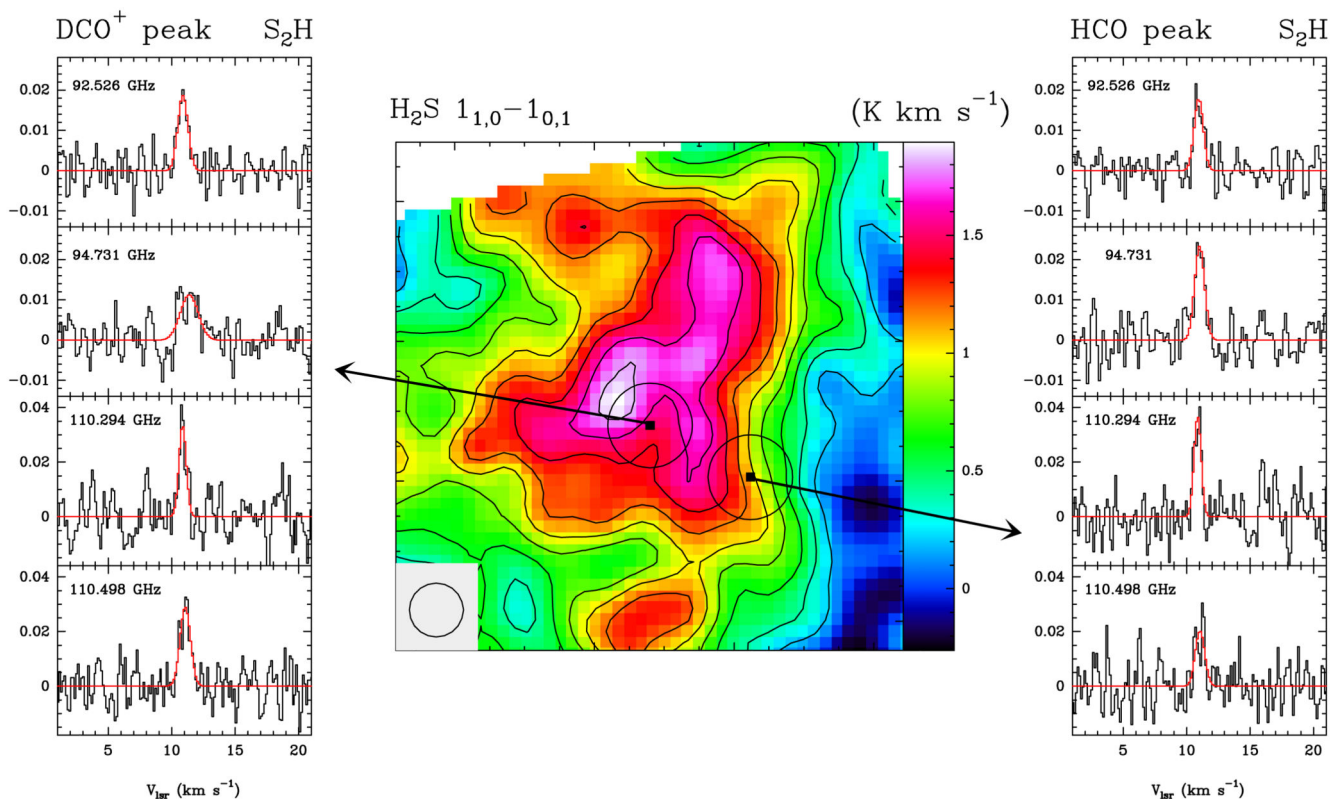
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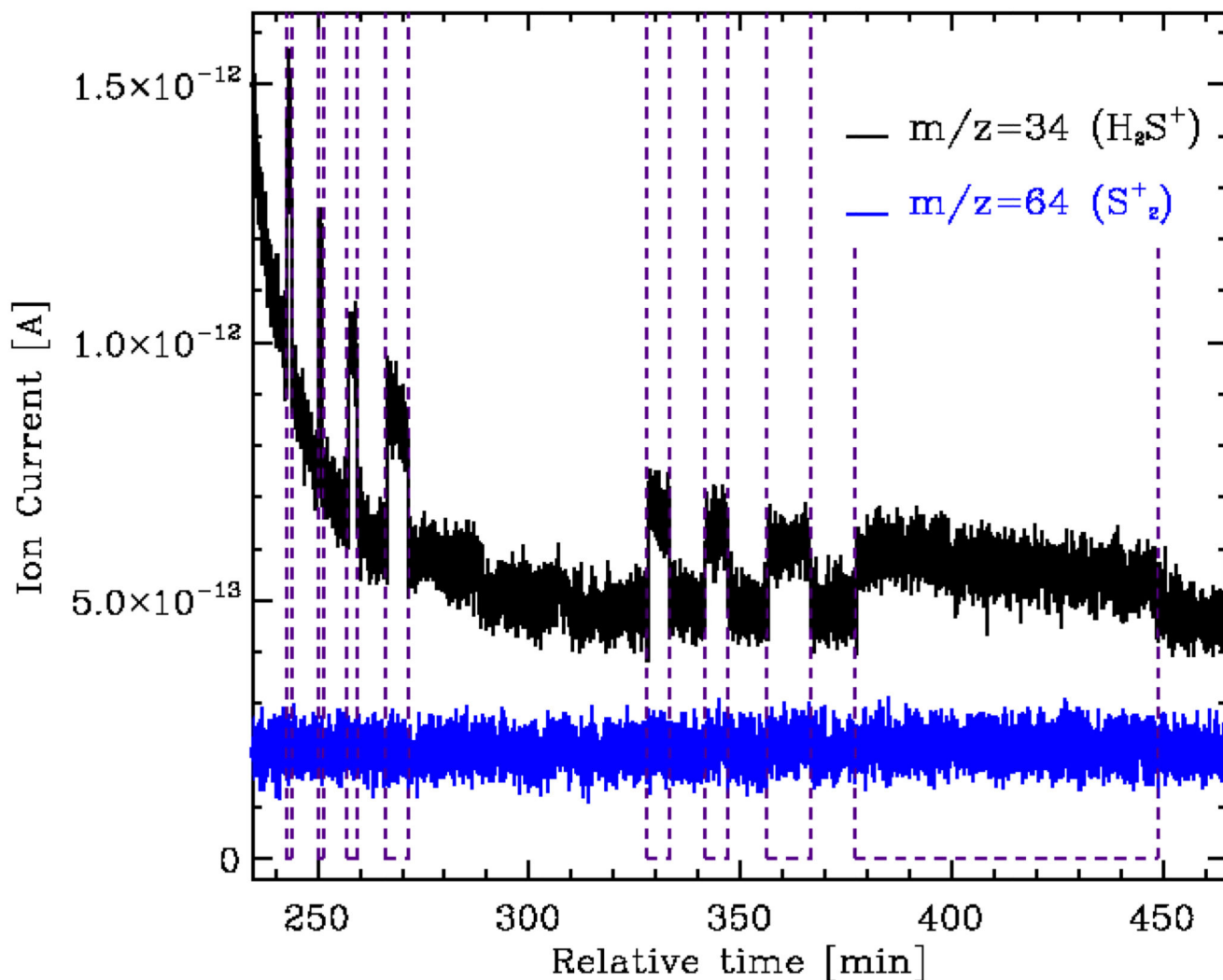
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**Figure 1.**

In the central panel, we show the integrated intensity map of the  $\text{H}_2\text{S } 1_{1,0} \rightarrow 1_{0,1}$  line (168.763 GHz). UV-illumination from  $\sigma\text{Ori}$  comes from the west (right). The beam is drawn in the bottom-left corner. Black circles around the surveyed positions indicate the beam of the  $\text{S}_2\text{H}$  detections. Spectra of the four  $\text{S}_2\text{H}$  lines detected towards the two positions targeted in the Whisper spectral are plotted in the left ( $\text{DCO}^+$  peak) and right ( $\text{HCO}$  peak) panels. The frequency in GHz is indicated in the top-left corner. In red, the Gaussian fits shown in Table 1.



**Figure 2.** Photodesorption of  $\text{H}_2\text{S}$  (black) detected by the QMS during irradiation of a pure  $\text{H}_2\text{S}$  ice sample. No increase of the measured ion current for the mass fragment  $m/z = 64$  (blue, corresponding to any sulfur-sulfur photo-product) was detected. Irradiation intervals are indicated with vertical dashed lines. Signals are shifted for clarity.

**Table 1**

## Gaussian fits

<b>Freq(MHz)</b>	<b>Area(K kms<sup>-1</sup>)</b>	<b>v<sub>lsr</sub>(km s<sup>-1</sup>)</b>	<b>v(km s<sup>-1</sup>)</b>	<b>T<sub>MB</sub>(K)</b>	<b>rms (K)</b>
DCO <sup>+</sup> peak (core)					
94526.32	0.019 (0.002)	10.90 (0.05)	0.9 (0.1)	0.018	0.004
94731.21	0.020 (0.003)	11.39 (0.13)	1.7 (0.2)	0.011	0.004
110294.15	0.024 (0.004)	10.86 (0.05)	0.7 (0.1)	0.033	0.008
110498.11	0.029 (0.004)	11.09 (0.06)	0.9 (0.1)	0.029	0.007
HCO peak (PDR)					
94526.32	0.016 (0.002)	10.97 (0.06)	0.9 (0.1)	0.018	0.004
94731.21	0.023 (0.003)	11.03 (0.05)	0.9 (0.1)	0.023	0.004
110294.15	0.026 (0.003)	10.87 (0.04)	0.7 (0.1)	0.037	0.008
110498.11	0.019 (0.004)	11.07 (0.09)	0.9 (0.2)	0.020	0.007

**Table 2**

Summary of column densities and fractional abundances

Molecule	DCO <sup>+</sup> peak				HCO peak		
	HPBW (")	T <sub>rot</sub> (K)	N(X) (cm <sup>-2</sup> )	N(X)/N <sub>H</sub>	T <sub>rot</sub> (K)	N(X) (cm <sup>-2</sup> )	N(X)/N <sub>H</sub>
H <sub>2</sub>	12		2.9×10 <sup>22</sup>	0.5		1.9×10 <sup>22</sup>	0.5
S <sub>2</sub> H	22 - 26	8.73 <sup>+1.36</sup> <sub>-1.10</sub>	3.0 <sup>+0.9</sup> <sub>-0.6</sub> × 10 <sup>12</sup>	5.2 <sup>+1.5</sup> <sub>-1.0</sub> × 10 <sup>-11</sup>	12.69 <sup>+1.78</sup> <sub>-1.54</sub>	3.3 <sup>+1.2</sup> <sub>-0.7</sub> × 10 <sup>12</sup>	8.7 <sup>+3.1</sup> <sub>-1.9</sub> × 10 <sup>-11</sup>
H <sub>2</sub> S <sup>I</sup>	14	9 <sup>+1</sup> <sub>-1</sub>	1.9 <sup>+0.4</sup> <sub>-0.3</sub> × 10 <sup>13</sup>	3.3 <sup>+0.7</sup> <sub>-0.6</sub> × 10 <sup>-10</sup>	12 <sup>+2</sup> <sub>-2</sub>	1.2 <sup>+0.1</sup> <sub>-0.2</sub> × 10 <sup>13</sup>	3.1 <sup>+0.3</sup> <sub>-0.5</sub> × 10 <sup>-10</sup>

<sup>(1)</sup>We assume the rotation temperatures derived from S<sub>2</sub>H.

**Table 3**

## Column density upper limits

Molecule	Freq (GHz)	rms <sup>1</sup> (mK)	N <sub>X</sub> <sup>2</sup> (cm <sup>-2</sup> )
S <sub>2</sub> H <sub>2</sub>	139.885	9	<8.5×10 <sup>11</sup>
HSO	158.391	30	<1.5×10 <sup>12</sup>
HO <sub>2</sub>	130.260	11	<4.7×10 <sup>11</sup>
H <sub>2</sub> O <sub>2</sub>	90.365	4	<1.0×10 <sup>12</sup>

(1) The rms has been calculated for a channel width of  $\approx 0.3 \text{ km s}^{-1}$ . The obtained rms is similar in the two surveyed positions.

(2)  $3\sigma$  upper limits assuming LTE,  $T_{\text{rot}}=10 \text{ K}$  and a linewidth of  $0.6 \text{ km s}^{-1}$