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# OPTIMAL EXTINCTION RATES FOR THE FAST DIFFUSION EQUATION WITH STRONG ABSORPTION 

RAZVAN GABRIEL IAGAR AND PHILIPPE LAURENÇOT


#### Abstract

Optimal extinction rates near the extinction time are derived for non-negative solutions to a fast diffusion equation with strong absorption, the power of the absorption exceeding that of the diffusion.


## 1. Introduction

Given $m \in(0, \infty), q \in(0,1)$, and a non-negative initial condition $u_{0}$ in $B C\left(\mathbb{R}^{N}\right), u_{0} \not \equiv 0$, it is well-known that the initial value problem

$$
\begin{align*}
\partial_{t} u-\Delta u^{m}+u^{q} & =0, \quad(t, x) \in(0, \infty) \times \mathbb{R}^{N},  \tag{1.1a}\\
u(0) & =u_{0}, \quad x \in \mathbb{R}^{N}, \tag{1.1b}
\end{align*}
$$

has a unique non-negative (weak) solution $u$ which vanishes identically after a finite time, a phenomenon usually referred to as finite time extinction [19-21]. More precisely, introducing the extinction time

$$
\begin{equation*}
T_{e}:=\sup \{t>0: u(t) \not \equiv 0\}>0, \tag{1.2}
\end{equation*}
$$

then $T_{e}$ is finite and satisfies $T_{e} \leq\left\|u_{0}\right\|_{\infty}^{(1-q)} /(1-q)$, the latter upper bound being a straightforward consequence of (1.1) and the comparison principle. Moreover, there holds

$$
\begin{equation*}
u(t) \not \equiv 0 \text { for } t \in\left[0, T_{e}\right) \text { and } u(t) \equiv 0 \text { for } t \geq T_{e} . \tag{1.3}
\end{equation*}
$$

When $q<m$ and $u_{0}(x) \rightarrow 0$ as $|x| \rightarrow \infty$, finite time extinction is accompanied by an even more striking phenomenon, the instantaneous shrinking of the support, that is, the positivity set $\mathcal{P}(t):=$ $\left\{x \in \mathbb{R}^{N}: u(t, x)>0\right\}$ of $u$ at time $t$ is a relatively compact subset of $\mathbb{R}^{N}$ for all $t \in\left(0, T_{e}\right)$, even if $\mathcal{P}(0)=\mathbb{R}^{N}$ initially $[1,5,7,19]$. Observe that the inequality $q<m$ is always satisfied when the diffusion is linear $(m=1)$ or slow $(m>1)$. Additional information on the behaviour of $\mathcal{P}(t)$ as $t \rightarrow T_{e}$ is also available when $m \geq 1$ and $N=1[6,10-14]$.

Once finite time extinction is known to take place, gaining further insight into the underlying mechanism requires to identify the behaviour of $u(t)$ as $t \rightarrow T_{e}$, a preliminary step being to determine

[^0]the relevant space and time scales. Simple scaling arguments predict that, for $r \in[1, \infty]$ and $u_{0} \in L^{r}\left(\mathbb{R}^{N}\right)$, there is a constant $\gamma_{r}>0$ (depending on $N, m, q, u_{0}$, and $r$ ) such that
\[

$$
\begin{equation*}
\|u(t)\|_{r} \sim \gamma_{r}\left(T_{e}-t\right)^{\alpha-(N \beta / r)} \tag{1.4}
\end{equation*}
$$

\]

where

$$
\begin{equation*}
\alpha:=\frac{1}{1-q}>0, \quad \beta:=\frac{q-m}{2(1-q)} \in \mathbb{R} . \tag{1.5}
\end{equation*}
$$

As already observed by several authors [9,12,15], a rather simple comparison argument provides a lower bound for the $L^{\infty}$-norm of the form (1.4). Indeed, consider $t \in\left(0, T_{e}\right)$ and let $x(t) \in \mathbb{R}^{N}$ be a point where $u(t)$ reaches its maximum value, that is, $u(t, x(t))=\|u(t)\|_{\infty}$. Then $u(t)^{m}$ also attains its maximum value at this point, so that $\Delta u^{m}(t, x(t)) \leq 0$ and we infer from (1.1a) that (at least formally)

$$
\frac{d}{d t}\|u(t)\|_{\infty}=\partial_{t} u(t, x(t)) \leq-u(t, x(t))^{q}=-\|u(t)\|_{\infty}^{q}
$$

Integrating the above differential inequality over $\left(t, T_{e}\right)$ gives the expected lower bound

$$
\begin{equation*}
\|u(t)\|_{\infty} \geq\left[(1-q)\left(T_{e}-t\right)\right]^{1 /(1-q)}, \quad t \in\left[0, T_{e}\right) \tag{1.6}
\end{equation*}
$$

The derivation of an upper bound of the form (1.4) turns out to be more involved and the results obtained so far are rather sparse: in one space dimension, the upper bound

$$
\begin{equation*}
\|u(t)\|_{\infty} \leq C_{\infty}\left(T_{e}-t\right)^{1 /(1-q)}, \quad t \in\left[0, T_{e}\right) \tag{1.7}
\end{equation*}
$$

is shown in [17, Proposition 2.2] for $m=1$ and in [9, Lemma 5.2, Lemma $7.2 \&$ Lemma 9.2] for $m \in(0,1)$, the latter being valid only for compactly supported initial data. The proofs are however of a completely different nature: in [17], properties of the linear heat equation are used while the approach in [9] relies on the intersection-comparison technique, which requires in particular the compactness of the support of the initial condition. Still for $m=1$ but in any space dimension, the upper bound (1.7) is derived in [12, Lemma 2.1] for radially symmetric initial data $u_{0}$ having a non-increasing profile and satisfying $\Delta u_{0}+\mu u_{0}^{q} \geq 0$ in $\mathbb{R}^{N}$ for some $\mu>0$. The last case for which (1.7) is proved corresponds to the choice $m=2-q>1$ and the proof relies on the derivation of an Aronson-Bénilan estimate, which seems to be only available for this specific choice of the parameters $m$ and $q$ [15].

The purpose of this note is to contribute to the validity of (1.4) and derive optimal upper and lower bounds near the extinction time when the parameters $m$ and $q$ range in

$$
\begin{equation*}
\frac{(N-2)_{+}}{N}<m<q<1 \tag{1.8}
\end{equation*}
$$

Recalling that a lower bound in $L^{\infty}$ is already available, see (1.6), we begin with upper bounds.
Theorem 1.1 (Upper bounds). Assume that $m$ and $q$ satisfy (1.8) and consider a non-negative initial condition $u_{0} \in B C\left(\mathbb{R}^{N}\right), u_{0} \not \equiv 0$, for which there is $\kappa_{0}>0$ such that

$$
\begin{equation*}
u_{0}(x) \leq \kappa_{0}|x|^{-2 /(q-m)}, \quad x \in \mathbb{R}^{N} . \tag{1.9}
\end{equation*}
$$

Given $r \in[1, \infty]$, there is $C_{r}>0$ depending only on $N, m, q, u_{0}$, and $r$ such that the solution $u$ to (1.1a)-(1.1b) satisfies

$$
\begin{equation*}
\|u(t)\|_{r} \leq C_{r}\left(T_{e}-t\right)^{\alpha-(N \beta / r)}, \quad t \in\left(0, T_{e}\right) \tag{1.10}
\end{equation*}
$$

the extinction time $T_{e}$ being defined in (1.2).
Theorem 1.1 thus extends the validity of the upper bound (1.7) established in [9] for $N=1$ and $r=\infty$ to any space dimension $N \geq 1$ and $r \in[1, \infty]$, while relaxing the assumption of compact support required in [9]. It is worth mentioning that the validity of (1.10) for $r \in[1, \infty)$ does not seem to be a simple consequence of (1.10) for $r=\infty$ since $u(t)$ is positive everywhere in $\mathbb{R}^{N}$ for all $t \in\left(0, T_{e}\right)$ even if $u_{0}$ is compactly supported, see [9, Lemma 2.5] and Proposition 1.4 below.

To be able to cope with higher space dimensions and non-compactly supported initial data, the proof of Theorem 1.1 takes a different route from that in [9] and is carried out in two steps: we first show that the algebraic decay at infinity (1.9) enjoyed by $u_{0}$ remains true throughout time evolution and combine it with (1.1a) to prove (1.10) for $r=1$. We next use self-similar variables and Moser's interation technique to derive (1.10) for all $r \in(1, \infty]$.

As a consequence of (1.6) and Theorem 1.1 for $r=\infty$, the correct time scale for the extinction phenomenon is identified. We now supplement the lower bound (1.6) in $L^{\infty}$ with another one in $L^{m+1}$. On the one hand, it allows us to identify the right space scale. On the other hand, its derivation does not rely on the comparison principle but on energy estimates, a technique which is more likely to extend to other problems for which the former might not be available.

Theorem 1.2 (Lower bound in $L^{m+1}$ ). Assume that $m$ and $q$ satisfy (1.8) and consider a nonnegative initial condition $u_{0} \in B C\left(\mathbb{R}^{N}\right)$, $u_{0} \not \equiv 0$, such that $u_{0} \in L^{m+1}\left(\mathbb{R}^{N}\right)$. There is $c_{m+1}>0$ depending only on $N, m, q$, and $u_{0}$ such that the solution $u$ to (1.1a)-(1.1b) satisfies

$$
\begin{equation*}
\|u(t)\|_{m+1} \geq c_{m+1}\left(T_{e}-t\right)^{\alpha-(N \beta /(m+1))}, \quad t \in\left(0, T_{e}\right) \tag{1.11}
\end{equation*}
$$

the extinction time $T_{e}$ being defined in (1.2).
Observing that Theorems 1.1 and 1.2 are shown without using the $L^{\infty}$-lower bound (1.6), the latter may be recovered from these two results by Hölder's inequality, with a less explicit constant though.

Corollary 1.3. Assume that $m$ and $q$ satisfy (1.8) and consider a non-negative initial condition $u_{0} \in B C\left(\mathbb{R}^{N}\right), u_{0} \not \equiv 0$, enjoying the decay property (1.9). For $r \in(m+1, \infty]$, there is $c_{r}>0$ depending only $N, m, q, u_{0}$, and $r$ such that the solution $u$ to (1.1a)-(1.1b) satisfies

$$
\begin{equation*}
\|u(t)\|_{r} \geq c_{r}\left(T_{e}-t\right)^{\alpha-(N \beta / r)}, \quad t \in\left(0, T_{e}\right) \tag{1.12}
\end{equation*}
$$

Summarizing the outcome of Theorem 1.1, Theorem 1.2, and Corollary 1.3, we have shown that, for all non-negative initial data $u_{0} \in B C\left(\mathbb{R}^{N}\right), u_{0} \not \equiv 0$, enjoying the decay property (1.9), the corresponding solution $u$ to (1.1a)-(1.1b) is bounded in $L^{r}\left(\mathbb{R}^{N}\right), r \in[m+1, \infty]$, from above and from below at time $t \in\left(0, T_{e}\right)$ by the same power of $T_{e}-t$. Such estimates pave the way towards a more precise description of the behaviour of $u(t)$ as $t \rightarrow T_{e}$, which is expected to be self-similar. That this
is indeed the case is shown in $[8,9]$ in one space dimension, another building block of the proof being the uniqueness of self-similar solutions [8].

We end up this note with the already mentioned everywhere positivity of solutions to (1.1a)-(1.1b) for positive times prior to the extinction time. As we shall see below, this property holds true for a wider range of the parameters $m$ and $q$, namely $0<m \leq q<1$. It is already observed in [9, Lemma 2.5] in one space dimension and we extend it herein to any space dimension. It is worth emphasizing that it includes the case $q=m$ and contrasts markedly with the instantaneous shrinking of the support occurring when $q<m$.

Proposition 1.4 (Everywhere positivity). Consider $0<m \leq q<1$. Let $u_{0} \in B C\left(\mathbb{R}^{N}\right)$ be a nonnegative initial condition, $u_{0} \not \equiv 0$, and denote the corresponding solution to (1.1a)-(1.1b) by $u$ with extinction time $T_{e}$. For $t \in\left(0, T_{e}\right)$, there holds

$$
\begin{equation*}
\mathcal{P}(t):=\left\{x \in \mathbb{R}^{N}: u(t, x)>0\right\}=\mathbb{R}^{N} \tag{1.13}
\end{equation*}
$$

Before proving the results stated above, we point out once more that the energy techniques developed herein seem to be rather flexible and are expected to have a wider range of applicability. For instance, a related approach is used in the companion paper [18], where optimal (lower and upper) bounds near the extinction time are established for a different fast diffusion equation (featuring the $p$-Laplacian operator, $p \in(1,2))$ with a gradient absorption term.

## 2. Upper bounds near the extinction time

Throughout this section, we assume that $m$ and $q$ satisfy (1.8) and consider a non-negative initial condition $u_{0} \in B C\left(\mathbb{R}^{N}\right), u_{0} \not \equiv 0$, enjoying the decay property (1.9). Let $u$ be the corresponding solution to (1.1a)-(1.1b).
2.1. $L^{1}$-estimate. We begin with the propagation throughout time evolution of the algebraic decay (1.9) and set

$$
\begin{equation*}
\kappa_{*}:=\left(\frac{2 m(m+q)}{(q-m)^{2}}\right)^{1 /(q-m)} . \tag{2.1}
\end{equation*}
$$

Lemma 2.1. For $t \in[0, \infty)$ and $x \in \mathbb{R}^{N} \backslash\{0\}$, there holds

$$
u(t, x) \leq \max \left\{\kappa_{0}, \kappa_{*}\right\}|x|^{-2 /(q-m)} .
$$

Proof. Set $\Sigma_{\kappa}(x):=\kappa|x|^{-2 /(q-m)}$ for $x \in \mathbb{R}^{N} \backslash\{0\}$, where $\kappa$ is a positive constant yet to be determined. We note that

$$
\begin{aligned}
-\Delta \Sigma^{m}(x)+\Sigma(x)^{q}= & -\kappa^{m}\left[\frac{2 m(m+q)}{(q-m)^{2}}|x|^{-2 q /(q-m)}-\frac{2 m(N-1)}{(q-m)}|x|^{-2 q /(q-m)}\right] \\
& +\kappa^{q}|x|^{-2 q /(q-m)} \\
\geq & \kappa^{m}\left(\kappa^{q-m}-\kappa_{*}^{q-m}\right)|x|^{-2 q /(q-m)}
\end{aligned}
$$

for $x \in \mathbb{R}^{N} \backslash\{0\}$, so that $\Sigma_{\kappa}$ is a supersolution to (1.1a) in $\mathbb{R}^{N} \backslash\{0\}$ for all $\kappa \geq \kappa_{*}$. We then choose $\kappa=\max \left\{\kappa_{0}, \kappa_{*}\right\}$ and use the comparison principle to complete the proof of Lemma 2.1.

We are now in a position to derive the claimed upper bound near the extinction time for $r=1$. Proof of Theorem 1.1: $r=1$. Let $t \in\left[0, T_{e}\right)$. Integrating (1.1a) over $\left(t, T_{e}\right) \times \mathbb{R}^{N}$ gives

$$
\begin{equation*}
\|u(t)\|_{1}=\int_{t}^{T_{e}} \int_{\mathbb{R}^{N}} u(s, x)^{q} d x d s \tag{2.2}
\end{equation*}
$$

Owing to (1.8), there holds $2 q /(q-m)>N$ and we infer from Lemma 2.1 and Hölder's inequality that, for $s \in\left(t, T_{e}\right)$ and $R>0$,

$$
\begin{aligned}
\int_{\mathbb{R}^{N}} u(s, x)^{q} d x & =\int_{B(0, R)} u(s, x)^{q} d x+\int_{\mathbb{R}^{N} \backslash B(0, R)} u(s, x)^{q} d x \\
& \leq C R^{N(1-q)}\|u(s)\|_{1}^{q}+\left(\max \left\{\kappa_{0}, \kappa_{*}\right\}\right)^{q}\left|\mathbb{S}^{N-1}\right| \int_{R}^{\infty} r^{N-1-(2 q /(q-m))} d r \\
& \leq C\left(R^{N(1-q)}\|u(s)\|_{1}^{q}+R^{(N(q-m)-2 q) /(q-m)}\right)
\end{aligned}
$$

We next optimize in $R$ in the previous inequality by setting $R(s):=\|u(s)\|_{1}^{-(q-m) /(N(m-q)+2)}$, which satisfies

$$
R(s)^{N(1-q)}\|u(s)\|_{1}^{q}=R(s)^{(N(q-m)-2 q) /(q-m)}=\|u(s)\|_{1}^{(N(m-q)+2 q) /(N(m-q)+2)}
$$

Consequently, taking $R=R(s)$ in the previous inequality, we obtain

$$
\int_{\mathbb{R}^{N}} u(s, x)^{q} d x \leq C\|u(s)\|_{1}^{(N(m-q)+2 q) /(N(m-q)+2)}
$$

which gives, together with (2.2), the positivity of $N(m-q)+2 q$, and the time monotonicity of $s \mapsto\|u(s)\|_{1}$,

$$
\begin{aligned}
\|u(t)\|_{1} & \leq C \int_{t}^{T_{e}}\|u(s)\|_{1}^{(N(m-q)+2 q) /(N(m-q)+2)} d s \\
& \leq C\left(T_{e}-t\right)\|u(t)\|_{1}^{(N(m-q)+2 q) /(N(m-q)+2)}
\end{aligned}
$$

from which (1.10) for $r=1$ readily follows.
2.2. Scaling variables and $L^{r}$-estimates, $r \in(1, \infty]$. The next step is to take advantage of the just derived $L^{1}$-upper bound to derive the corresponding ones in $L^{r}$ for $r \in(1, \infty]$. To this end, we introduce the scaling variables

$$
\begin{equation*}
s:=\ln \left(\frac{T_{e}}{T_{e}-t}\right), \quad y:=x\left(T_{e}-t\right)^{\beta}, \quad(t, x) \in\left[0, T_{e}\right) \times \mathbb{R}^{N} \tag{2.3}
\end{equation*}
$$

and the new unknown function $v$ defined by

$$
\begin{equation*}
u(t, x)=\left(T_{e}-t\right)^{\alpha} v\left(\ln \left(T_{e}\right)-\ln \left(T_{e}-t\right), x\left(T_{e}-t\right)^{\beta}\right), \quad(t, x) \in\left[0, T_{e}\right) \times \mathbb{R}^{N} \tag{2.4}
\end{equation*}
$$

or, equivalently,

$$
\begin{equation*}
v(s, y)=T_{e}^{-\alpha} e^{\alpha s} u\left(T_{e}\left(1-e^{-s}\right), y T_{e}^{-\beta} e^{\beta s}\right), \quad(s, y) \in[0, \infty) \times \mathbb{R}^{N} \tag{2.5}
\end{equation*}
$$

It readily follows from (1.1a)-(1.1b) that $v$ solves

$$
\begin{align*}
\partial_{s} v(s, y) & =\alpha v(s, y)+\beta y \cdot \nabla v(s, y)+\Delta v^{m}(s, y)-v(s, y)^{q}, \quad(s, y) \in(0, \infty) \times \mathbb{R}^{N},  \tag{2.6}\\
v(0, y) & =v_{0}(y):=T_{e}^{-\alpha} u_{0}\left(y T_{e}^{-\beta}\right), \quad y \in \mathbb{R}^{N} . \tag{2.7}
\end{align*}
$$

Since

$$
\begin{equation*}
\|u(t)\|_{r}=\left(T_{e}-t\right)^{\alpha-(N \beta / r)}\|v(s)\|_{r}, \quad t \in\left(0, T_{e}\right) \tag{2.8}
\end{equation*}
$$

for all $r \in[1, \infty]$, we realize that an upper bound such as (1.10) on $\|u(t)\|_{r}$ for $t \in\left(0, T_{e}\right)$ obviously follows from a uniform upper bound on $\|v(s)\|_{r}$ for $s \geq 0$, the converse being true as well. In particular, it follows from (2.8) and Theorem 1.1 for $r=1$ that

$$
\begin{equation*}
\|v(s)\|_{1} \leq C_{1}, \quad s \geq 0 \tag{2.9}
\end{equation*}
$$

and we may assume without loss of generality that $C_{1} \geq 1$.
We now aim at using a bootstrap argument to deduce from (2.6) and (2.9) that $v$ belongs to $L^{\infty}\left(0, \infty ; L^{r}\left(\mathbb{R}^{N}\right)\right)$ for all $r \in(1, \infty]$. To this end, Moser's iteration technique is a suitable tool and the way we apply it is inspired from [2, Theorem 3.1]. But since [2, Theorem 3.1] is devoted to the slow diffusion case $m>1$, some technical aspects of its proof do not seem to apply directly here and we borrow additional arguments from the proof of [3, Proposition 2].

Lemma 2.2. Let $r \in(0, \infty]$. There is $C_{r+1}>0$ depending only on $N, m, q, u_{0}$, and $r$ such that

$$
\|v(s)\|_{r+1} \leq C_{r+1}, \quad s \geq 0
$$

Proof. Let $r \in[2-m, \infty)$. Multiplying (2.6) by $v^{r}$, integrating over $\mathbb{R}^{N}$, and using integration by parts, we obtain

$$
\begin{aligned}
& \frac{1}{r+1} \frac{d}{d s}\|v\|_{r+1}^{r+1}+r m \int_{\mathbb{R}^{N}} v^{r+m-2}|\nabla v|^{2} d y+\int_{\mathbb{R}^{N}} v^{r+q} d y=\left(\alpha-\frac{N \beta}{r+1}\right)\|v\|_{r+1}^{r+1}, \\
& \frac{d}{d s}\|v\|_{r+1}^{r+1}+\frac{4 m r(r+1)}{(m+r)^{2}}\left\|\nabla v^{(m+r) / 2}\right\|_{2}^{2} \leq \alpha(r+1)\|v\|_{r+1}^{r+1} .
\end{aligned}
$$

Since $4 m r(r+1) \geq 2 m(m+r)^{2}$, we end up with

$$
\begin{equation*}
\frac{d}{d s}\|v\|_{r+1}^{r+1}+2 m\left\|\nabla v^{(m+r) / 2}\right\|_{2}^{2} \leq \alpha(r+1)\|v\|_{r+1}^{r+1} \tag{2.10}
\end{equation*}
$$

We next fix $\zeta \in\left(2 / m, 2^{*}\right)$ where $2^{*}:=2 N /(N-2)_{+}\left(\right.$with $2^{*}=\infty$ for $\left.N=1,2\right)$. On the one hand, it follows from the Gagliardo-Nirenberg inequality that

$$
\begin{equation*}
\left\|v^{(m+r) / 2}\right\|_{\zeta} \leq C\left\|\nabla v^{(m+r) / 2}\right\|_{2}^{\theta}\left\|v^{(m+r) / 2}\right\|_{1}^{1-\theta} \tag{2.11}
\end{equation*}
$$

with

$$
\theta:=\frac{2 N(\zeta-1)}{(N+2) \zeta}
$$

On the other hand, since $(m+r) / 2 \in[1, \zeta(m+r) / 2]$ for $r \geq 2-m$, we infer from Hölder's inequality that

$$
\begin{equation*}
\|v\|_{(m+r) / 2}^{(m+r) / 2} \leq\|v\|_{\zeta(m+r) / 2}^{\zeta(m+r)(m+r-2) / 2[\zeta(m+r)-2]}\|v\|_{1}^{(\zeta-1)(m+r) /[\zeta(m+r)-2]} . \tag{2.12}
\end{equation*}
$$

We deduce from (2.11) and (2.12) that

$$
\begin{aligned}
\|v\|_{\zeta(m+r) / 2}^{(m+r) / 2} & =\left\|v^{(m+r) / 2}\right\|_{\zeta} \leq C\left\|\nabla v^{(m+r) / 2}\right\|_{2}^{\theta}\left(\|v\|_{(m+r) / 2}^{(m+r) / 2}\right)^{1-\theta} \\
& \leq C\left\|\nabla v^{(m+r) / 2}\right\|_{2}^{\theta}\left[\|v\|_{\zeta(m+r) / 2}^{\zeta(m+r)(m+r-2) / 2[\zeta(m+r)-2]}\|v\|_{1}^{(\zeta-1)(m+r) /[\zeta(m+r)-2]}\right]^{1-\theta},
\end{aligned}
$$

hence

$$
\begin{equation*}
\|v\|_{\zeta(m+r) / 2}^{\zeta(m+r)[N(m+r)+2-N] / N[\zeta(m+r)-2]} \leq C\left\|\nabla v^{(m+r) / 2}\right\|_{2}^{2}\|v\|_{1}^{[2 N-(N-2) \zeta](m+r) / N[\zeta(m+r)-2]} \tag{2.13}
\end{equation*}
$$

Moreover, since $\zeta>2 / m$ and $m<1$, we have $2 r \leq m[\zeta(m+r)-2]$, hence

$$
\begin{aligned}
\frac{[2 N-(N-2) \zeta](m+r)}{N[\zeta(m+r)-2]} & \leq \frac{[2 N-(N-2) \zeta]}{N} \frac{m(m+r)}{2 r} \\
& \leq \frac{m[2 N-(N-2) \zeta]}{N}
\end{aligned}
$$

so that

$$
\begin{align*}
\|v\|_{1}^{[2 N-(N-2) \zeta](m+r) / N[\zeta(m+r)-2]} & \leq C_{1}^{[2 N-(N-2) \zeta](m+r) / N[\zeta(m+r)-2]} \\
& \leq C_{1}^{m[2 N-(N-2) \zeta] / N} . \tag{2.14}
\end{align*}
$$

Also,

$$
1-\frac{N[\zeta(m+r)-2]}{\zeta[N(m+r)+2-N]}=\frac{2 N-\zeta(N-2)}{\zeta[N(m+r)+2-N]}>0,
$$

and we infer from (2.13), (2.14), and Young's inequality that

$$
\begin{aligned}
\|v\|_{\zeta(m+r) / 2}^{m+r} \leq & \frac{N[\zeta(m+r)-2]}{\zeta[N(m+r)+2-N]}\|v\|_{\zeta(m+r) / 2}^{\zeta(m+r)[N(m+r)+2-N] / N[\zeta(m+r)-2]} \\
& +\frac{2 N-\zeta(N-2)}{\zeta[N(m+r)+2-N]} \\
\leq & C\left\|\nabla v^{(m+r) / 2}\right\|_{2}^{2}+1 .
\end{aligned}
$$

Therefore, there is $\nu \in(0,1)$ depending only on $N, m, q$, and $u_{0}$ such that

$$
\begin{equation*}
\nu\left(\|v\|_{\zeta(m+r) / 2}^{m+r}-1\right) \leq\left\|\nabla v^{(m+r) / 2}\right\|_{2}^{2} \tag{2.15}
\end{equation*}
$$

Moreover, since $r+1 \in[1, \zeta(m+r) / 2]$, it follows from (2.9) and Hölder's and Young's inequalities that

$$
\begin{align*}
\|v\|_{r+1}^{r+1} & \leq\|v\|_{\zeta(m+r) / 2}^{\zeta r(r+m)[\zeta(r+m)-2]}\|v\|_{1}^{[(\zeta-2) r+\zeta m-2] /(\zeta r+\zeta m-2)} \\
& \leq C_{1}^{[(\zeta-2) r+\zeta m-2] /(\zeta r+\zeta m-2)}\|v\|_{\zeta(m+r) / 2}^{\zeta r(r+m) /[\zeta(r+m)-2]} \\
& \leq C_{1}\|v\|_{\zeta(m+r) / 2}^{\zeta r(r+m) /[(r+m)-2]} \\
& \leq \frac{\zeta m-2}{\zeta(m+r)-2} C_{1}^{[\zeta(m+r)-2] /(\zeta m-2)}+\frac{\zeta r}{\zeta(m+r)-2}\|v\|_{\zeta(m+r) / 2}^{m+r} \\
& \leq C_{1}^{[\zeta(m+r)-2] /(\zeta m-2)}+\|v\|_{\zeta(m+r) / 2}^{m+r} \tag{2.16}
\end{align*}
$$

Next, let $\sigma>1$ to be chosen appropriately later on and set

$$
\mathcal{I}_{r}(s):=\int_{\mathbb{R}^{N}} v(s, y)^{\zeta[r+m+\sigma(1-m)] /[\sigma(\zeta-2)+2]} d y, \quad s \geq 0
$$

Since $\sigma(\zeta-2)+2 \in[\zeta, \sigma \zeta]$ and

$$
r+1=\frac{\sigma-1}{\sigma}(m+r)+\frac{m+r+\sigma(1-m)}{\sigma}
$$

we deduce from (2.9) and Hölder's and Young's inequalities that, for $\delta>0$,

$$
\begin{align*}
\|v\|_{r+1}^{r+1} & \leq\|v\|_{\zeta(m+r) / 2}^{(\sigma-1)(m+r) / \sigma} \mathcal{I}_{r}^{[\sigma(\zeta-2)+2] / \sigma \zeta} \\
& \leq \frac{(\sigma-1) \delta}{\sigma}\|v\|_{\zeta(m+r) / 2}^{m+r}+\frac{1}{\sigma \delta^{\sigma-1}} \mathcal{I}_{r}^{[\sigma(\zeta-2)+2] / \zeta} \\
& \leq \delta\|v\|_{\zeta(m+r) / 2}^{m+r}+\frac{1}{\delta^{\sigma-1}} \mathcal{I}_{r}^{[\sigma(\zeta-2)+2] / \zeta} \tag{2.17}
\end{align*}
$$

Combining (2.10), (2.15), and (2.17) leads us to

$$
\begin{aligned}
\frac{d}{d s}\|v\|_{r+1}^{r+1} & +2 m \nu\left(\|v\|_{\zeta(m+r) / 2}^{m+r}-1\right) \leq \frac{d}{d s}\|v\|_{r+1}^{r+1}+2 m\left\|\nabla v^{(m+r) / 2}\right\|_{2}^{2} \\
& \leq \alpha(r+1)\|v\|_{r+1}^{r+1} \\
& \leq \alpha \delta(r+1)\|v\|_{\zeta(m+r) / 2}^{m+r}+\frac{\alpha(r+1)}{\delta^{\sigma-1}} \mathcal{I}_{r}^{[\sigma(\zeta-2)+2] / \zeta}
\end{aligned}
$$

We then choose $\delta=m \nu / \alpha(r+1)$ in the above inequality to obtain

$$
\frac{d}{d s}\|v\|_{r+1}^{r+1}+m \nu\|v\|_{\zeta(m+r) / 2}^{m+r} \leq 2 m \nu+\frac{\alpha^{\sigma}(r+1)^{\sigma}}{(m \nu)^{\sigma-1}} \mathcal{I}_{r}^{[\sigma(\zeta-2)+2] / \zeta}
$$

We finally use (2.16) to estimate from below the second term of the left-hand side of the previous inequality and end up with

$$
\begin{equation*}
\frac{d}{d s}\|v\|_{r+1}^{r+1}+m \nu\|v\|_{r+1}^{r+1} \leq 2 m \nu+m \nu C_{1}^{[\zeta(m+r)-2] /(\zeta m-2)}+\frac{\alpha^{\sigma}(r+1)^{\sigma}}{(m \nu)^{\sigma-1}} \mathcal{I}_{r}^{[\sigma(\zeta-2)+2] / \zeta} \tag{2.18}
\end{equation*}
$$

We first choose

$$
\sigma=\frac{\zeta(m+r)-2}{\zeta m-2}>1
$$

in (2.18) and observe that this choice guarantees that

$$
\zeta[m+r+\sigma(1-m)]=\sigma(\zeta-2)+2
$$

Consequently, $\mathcal{I}_{r}=\|v\|_{1}$ and we deduce from (2.9) and (2.18) that there is $C(r)>0$ depending on $N, m, q, u_{0}$, and $r$ such that

$$
\frac{d}{d s}\|v\|_{r+1}^{r+1}+m \nu\|v\|_{r+1}^{r+1} \leq C(r)
$$

Integrating the previous differential inequality entails that

$$
\begin{equation*}
\sup _{s \geq 0}\|v(s)\|_{r+1}<\infty \tag{2.19}
\end{equation*}
$$

The validity of (2.19) extends to all $r \in(0,2-m)$ by (2.9) and Hölder's inequality.
To complete the proof of Lemma 2.2, we are left to check the boundedness of $v$ in $L^{\infty}\left(\mathbb{R}^{N}\right)$. To this end, we take $\sigma=\sigma_{0}:=2(\zeta-1) /(\zeta-2)>1$ in (2.18) and obtain, after integration with respect to time,

$$
\begin{aligned}
\|v(s)\|_{r+1}^{r+1} \leq & \left\|v_{0}\right\|_{r+1}^{r+1} e^{-m \nu s}+2+C_{1}^{[\zeta(m+r)-2] /(\zeta m-2)} \\
& +\left(\frac{\alpha(r+1)}{m \nu}\right)^{\sigma_{0}}\left[\sup _{s_{*} \in[0, s]} \mathcal{I}_{r}\left(s_{*}\right)\right]^{2} \\
\leq & \left\|v_{0}\right\|_{1}\left\|v_{0}\right\|_{\infty}^{r}+2+C_{1}^{[\zeta(m+r)-2] /(\zeta m-2)} \\
& +\left(\frac{\alpha(r+1)}{m \nu}\right)^{\sigma_{0}}\left[\sup _{s_{*} \in[0, s]} \mathcal{I}_{r}\left(s_{*}\right)\right]^{2}
\end{aligned}
$$

and

$$
\mathcal{I}_{r}=\|v\|_{\left[(r+m)+\sigma_{0}(1-m)\right] / 2}^{\left[(r+m)+\sigma_{0}(1-m) / 2\right.}
$$

Therefore, there are $K_{0}>0$ and $K_{1}>0$ depending only on $N, m, q$, and $u_{0}$ such that

$$
\begin{equation*}
\sup _{s \geq 0}\left\{\|v(s)\|_{r+1}^{r+1}\right\} \leq K_{0}\left(K_{1}^{r+1}+(1+r)^{\sigma_{0}} \sup _{s \geq 0}\left\{\|v(s)\|_{\left[(r+m)+\sigma_{0}(1-m)\right] / 2}^{\left[(r+m)+\sigma_{0}(1-m)\right]}\right\}\right) . \tag{2.20}
\end{equation*}
$$

We now define the sequence $\left(r_{j}\right)_{j \geq 0}$ by

$$
1+r_{j+1}=2\left(1+r_{j}\right)-(1-m)\left(\sigma_{0}-1\right), \quad j \geq 0, \quad r_{0}:=2-m
$$

and set

$$
V_{j}:=\sup _{s \geq 0}\left\{\|v(s)\|_{r_{j}+1}^{r_{j}+1}\right\}, \quad j \geq 0
$$

For $j \geq 0$, we take $r=r_{j+1}$ in (2.20) and realize that

$$
\begin{aligned}
V_{j+1} & \leq K_{0}\left(K_{1}^{1+r_{j+1}}+\left(1+r_{j+1}\right)^{\sigma_{0}} V_{j}^{2}\right) \\
& \leq K_{0}\left(1+r_{j+1}\right)^{\sigma_{0}} \max \left\{K_{1}^{1+r_{j+1}}, V_{j}^{2}\right\}, \quad j \geq 0
\end{aligned}
$$

Since $\sigma_{0}-1<1 /(1-m)$ thanks to the constraint $\zeta>2 / m$, one has $1+r_{0}-(1-m)\left(\sigma_{0}-1\right)>0$ and we are in a position to apply [22, Lemma A.1], which we recall in Lemma 2.3 below for completeness, to conclude that there is $K_{2}>0$ depending only on $m, \zeta, K_{0}$, and $K_{1}$ such that

$$
V_{j}^{1 /\left(1+r_{j}\right)} \leq K_{2}, \quad j \geq 0
$$

Equivalently,

$$
\sup _{s \geq 0}\left\{\|v(s)\|_{1+r_{j}}\right\} \leq K_{2}, \quad j \geq 0
$$

and letting $j \rightarrow \infty$ entails that $\|v(s)\|_{\infty} \leq K_{2}$ for all $s \geq 0$, thereby completing the proof of Lemma 2.2.

The proof of Theorem 1.1 for $r \in(1, \infty]$ is now a straightforward consequence of (2.8) and Lemma 2.2.

Lemma 2.3. Let $a>1, b \geq 0, c \in \mathbb{R}, C_{0} \geq 1, C_{1} \geq 1$, and $p_{0}>0$ be given such that $p_{0}(a-1)+c>0$. We define the sequence $\left(p_{k}\right)_{k \geq 0}$ of positive real numbers by $p_{k+1}=a p_{k}+c$ for $k \geq 0$ and assume that $\left(Q_{k}\right)_{k \geq 0}$ is a sequence of positive real numbers satisfying

$$
Q_{0} \leq C_{1}^{p_{0}}, \quad Q_{k+1} \leq C_{0} p_{k+1}^{b} \max \left\{C_{1}^{p_{k+1}}, Q_{k}^{a}\right\}, k \geq 0
$$

Then the sequence $\left(Q_{k}^{1 / p_{k}}\right)_{k \geq 0}$ is bounded.

## 3. Lower bound near the extinction time

We now turn to the lower bound near the extinction time in $L^{m+1}\left(\mathbb{R}^{N}\right)$.
Proof of Theorem 1.2. For $t \in\left[0, T_{e}\right]$, we define

$$
X(t):=\|u(t)\|_{m+1}^{m+1} \quad \text { and } \quad Y(t):=\int_{\mathbb{R}^{N}} u(t, x)^{m+q} d x .
$$

Let $t \in\left(0, T_{e}\right)$. It follows from (1.1a) that

$$
\begin{equation*}
\frac{1}{m+1} \frac{d X}{d t}(t)+\left\|\nabla u^{m}(t)\right\|_{2}^{2}+Y(t)=0 \tag{3.1}
\end{equation*}
$$

Since

$$
1<\frac{m+q}{m}<\frac{m+1}{m}<2^{*}:=\frac{2 N}{(N-2)_{+}}
$$

by (1.8) we infer from the Gagliardo-Nirenberg inequality that

$$
\begin{aligned}
X(t)^{m /(m+1)} & =\left\|u(t)^{m}\right\|_{(m+1) / m} \leq C\left\|\nabla u^{m}(t)\right\|_{2}^{\theta}\left\|u(t)^{m}\right\|_{(m+q) / m}^{1-\theta} \\
& \leq C Y(t)^{m(1-\theta) /(m+q)}\left\|\nabla u^{m}(t)\right\|_{2}^{\theta}
\end{aligned}
$$

where

$$
\theta:=\frac{2 N m(1-q)}{(m+1)[m(N+2)-q(N-2)]}
$$

Consequently, since $u(t) \not \equiv 0$ as $t \in\left(0, T_{e}\right)$,

$$
\left\|\nabla u^{m}(t)\right\|_{2}^{2} \geq C X(t)^{2 m / \theta(m+1)} Y(t)^{-2 m(1-\theta) / \theta(m+q)}
$$

which gives, together with (3.1),

$$
\begin{equation*}
\frac{d X}{d t}(t)+C X(t)^{2 m / \theta(m+1)} Y(t)^{-2 m(1-\theta) / \theta(m+q)}+(m+1) Y(t) \leq 0 \tag{3.2}
\end{equation*}
$$

Setting

$$
\xi:=1+\frac{2 m(1-\theta)}{\theta(m+q)}>1 \quad \text { and } \quad \gamma:=\frac{2 m}{\theta \xi(m+1)}
$$

it follows from Young's inequality that

$$
\begin{aligned}
X(t)^{\gamma} & =X(t)^{\gamma} Y(t)^{-(\xi-1) / \xi} Y(t)^{(\xi-1) / \xi} \leq \frac{1}{\xi} X(t)^{\xi \gamma} Y(t)^{1-\xi}+\frac{\xi-1}{\xi} Y(t) \\
& \leq X(t)^{2 m / \theta(m+1)} Y(t)^{-2 m(1-\theta) / \theta(m+q)}+Y(t)
\end{aligned}
$$

Combining this inequality with (3.2) leads us to the differential inequality

$$
\begin{equation*}
\frac{d X}{d t}(t)+C X(t)^{\gamma} \leq 0, \quad t \in\left(0, T_{e}\right) \tag{3.3}
\end{equation*}
$$

Now,

$$
\begin{aligned}
\gamma & =\frac{2 m}{\theta \xi(m+1)}=\frac{2 m(m+q)}{(2 m+\theta(q-m))(m+1)} \\
& =\frac{2 m(m+q)[m(N+2)-q(N-2)]}{2 m\{(m+1)[m(N+2)-q(N-2)]+N(q-m)(1-q)\}} \\
& =\frac{m(N+2)-q(N-2)}{m(N+2)-q N+2} \in(0,1)
\end{aligned}
$$

and we integrate (3.3) over $\left(t, T_{e}\right)$ to obtain

$$
-X(t)^{1-\gamma}+(1-\gamma) C\left(T_{e}-t\right) \leq 0, \quad t \in\left(0, T_{e}\right)
$$

Noticing that

$$
(m+1) \alpha-N \beta=\frac{m+1}{1-q}-\frac{N(q-m)}{2(1-q)}=\frac{m(N+2)-q N+2}{2(1-q)}=\frac{1}{1-\gamma}
$$

the lower bound (1.11) readily follows from the previous inequality.

We end up this section with the derivation of the lower bound for $r \in(m+1, \infty]$ from Theorem 1.1 for $r=1$ and Theorem 1.2.

Proof of Corollary 1.3. We first note that, owing to (1.8), there holds $2 /(q-m)>N$ and (1.9) entails that $u_{0} \in L^{1}\left(\mathbb{R}^{N}\right)$. Since $u_{0}$ also belongs to $L^{\infty}\left(\mathbb{R}^{N}\right)$, we conclude that $u_{0} \in L^{m+1}\left(\mathbb{R}^{N}\right)$.

Let $r \in(m+1, \infty]$ and $t \in\left(0, T_{e}\right)$. We infer from Theorem 1.1, Theorem 1.2, and Hölder's inequality that

$$
\begin{aligned}
c_{m+1}^{m+1}\left(T_{e}-t\right)^{(m+1) \alpha-N \beta} & \leq\|u(t)\|_{m+1}^{m+1} \leq\|u(t)\|_{r}^{r m /(r-1)}\|u(t)\|_{1}^{(r-1-m) /(r-1)} \\
& \leq C_{1}\left(T_{e}-t\right)^{(\alpha-N \beta)(r-1-m) /(r-1)}\|u(t)\|_{r}^{r m /(r-1)},
\end{aligned}
$$

from which (1.12) readily follows.

## 4. Everywhere positivity

In this section, we assume that $0<m \leq q<1$ and consider a non-negative initial condition $u_{0} \in B C\left(\mathbb{R}^{N}\right), u_{0} \not \equiv 0$. We denote the corresponding solution to (1.1a)-(1.1b) by $u$ and define its extinction time by (1.2). As in [9], the proof relies on an upper bound for $\partial_{t} u$ which we establish now.

Lemma 4.1. For $t>0$ there holds

$$
\partial_{t} u(t) \leq \frac{u(t)}{(1-m) t} \quad \text { in } \quad \mathbb{R}^{N}
$$

When $q=m$, Lemma 4.1 is a consequence of [4, Theorem 2], the proof relying on an homogeneity argument. Though the operator $-\Delta u^{m}+u^{q}$ is not homogeneous, we may still adapt the proof of [4, Theorem 2] when $q \geq m$.

Proof. Given a non-negative initial condition $u_{0} \in B C\left(\mathbb{R}^{N}\right)$, we denote the corresponding solution to (1.1a)-(1.1b) at time $t \geq 0$ by $S(t) u_{0}$. Recall that, if $u_{0}$ and $v_{0}$ are two non-negative functions in $B C\left(\mathbb{R}^{N}\right)$ satisfying $u_{0} \geq v_{0}$, then the comparison principle entails $S(t) u_{0} \geq S(t) v_{0}$ for all $t \geq 0$.
Step 1. We first claim that, for $\lambda \geq 1$,

$$
\begin{equation*}
S(\lambda t) u_{0} \leq \lambda^{1 /(1-m)} S(t)\left(\lambda^{1 /(m-1)} u_{0}\right), \quad t \geq 0 \tag{4.1}
\end{equation*}
$$

Indeed, setting $u(t):=S(t) u_{0}$ for $t \geq 0$, the function $v$ defined by $v(t):=\lambda^{1 /(m-1)} S(\lambda t) u_{0}$ satisfies

$$
\begin{aligned}
\partial_{t} v(t, x)-\Delta v^{m}(t, x)+v(t, x)^{q}= & \lambda^{m /(m-1)} \partial_{t} u(\lambda t, x)-\lambda^{m /(m-1)} \Delta u^{m}(\lambda t, x) \\
& +\lambda^{q /(m-1)} u(\lambda t, x)^{q} \\
= & \left(\lambda^{q /(m-1)}-\lambda^{m /(m-1)}\right) u(\lambda t, x)^{q} \leq 0
\end{aligned}
$$

Since $v(0)=\lambda^{1 /(m-1)} u_{0} \leq u_{0}$, we infer from the comparison principle that (4.1) holds true.

Step 2. Now, fix $t>0$ and consider $h>0$. Since $\lambda=(1+h / t)>1$ and $m \in(0,1)$, we infer from (4.1) and the comparison principle that

$$
\begin{aligned}
S(t+h) u_{0}-S(t) u_{0} & =S(\lambda t) u_{0}-S(t) u_{0} \\
& \leq \lambda^{1 /(1-m)} S(t)\left(\lambda^{1 /(m-1)} u_{0}\right)-S(t) u_{0} \\
& \leq\left[\left(1+\frac{h}{t}\right)^{1 /(1-m)}-1\right] S(t) u_{0}
\end{aligned}
$$

Dividing the above inequality by $h$ and passing to the limit as $h \rightarrow 0$ complete the proof.
We now argue as in the proof of [9, Lemma 2.5] to complete the proof of Proposition 1.4.
Proof of Proposition 1.4. Fix $t \in\left(0, T_{e}\right)$ and assume for contradiction that $u\left(t, x_{0}\right)=0$ for some $x_{0} \in \mathbb{R}^{N}$. By (1.1a) and Lemma 4.1, there holds

$$
-\Delta u^{m}(t)+u(t)^{q}+\frac{u(t)}{(1-m) t} \geq 0 \quad \text { in } \quad \mathbb{R}^{N}
$$

so that $u(t)^{m}$ is a supersolution to

$$
-\Delta w+d w=0 \quad \text { in } \quad \mathbb{R}^{N},
$$

with $d(x):=u(t, x)^{q-m}+u(t, x)^{1-m} /((1-m) t)$ for $x \in \mathbb{R}^{N}$. Since $t>0$ and $m \leq q<1$, the function $d$ is non-negative and bounded and we infer from the strong maximum principle [16, Theorem 8.19] that $u(t)^{m} \equiv 0$ in $\mathbb{R}^{N}$, contradicting $t<T_{e}$. Consequently, $u(t)^{m}$ is positive everywhere in $\mathbb{R}^{N}$ and the proof of Proposition 1.4 is complete.

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