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### OPTIMAL EXTINCTION RATES FOR THE FAST DIFFUSION EQUATION WITH STRONG ABSORPTION

### RAZVAN GABRIEL IAGAR AND PHILIPPE LAURENÇOT

ABSTRACT. Optimal extinction rates near the extinction time are derived for non-negative solutions to a fast diffusion equation with strong absorption, the power of the absorption exceeding that of the diffusion.

### 1. INTRODUCTION

Given  $m \in (0, \infty)$ ,  $q \in (0, 1)$ , and a non-negative initial condition  $u_0$  in  $BC(\mathbb{R}^N)$ ,  $u_0 \neq 0$ , it is well-known that the initial value problem

$$\partial_t u - \Delta u^m + u^q = 0$$
,  $(t, x) \in (0, \infty) \times \mathbb{R}^N$ , (1.1a)

$$u(0) = u_0 , \qquad x \in \mathbb{R}^N , \qquad (1.1b)$$

has a unique non-negative (weak) solution u which vanishes identically after a finite time, a phenomenon usually referred to as *finite time extinction* [19–21]. More precisely, introducing the extinction time

$$T_e := \sup\{t > 0 : u(t) \neq 0\} > 0 , \qquad (1.2)$$

then  $T_e$  is finite and satisfies  $T_e \leq ||u_0||_{\infty}^{(1-q)}/(1-q)$ , the latter upper bound being a straightforward consequence of (1.1) and the comparison principle. Moreover, there holds

$$u(t) \neq 0 \text{ for } t \in [0, T_e) \text{ and } u(t) \equiv 0 \text{ for } t \geq T_e .$$
 (1.3)

When q < m and  $u_0(x) \to 0$  as  $|x| \to \infty$ , finite time extinction is accompanied by an even more striking phenomenon, the *instantaneous shrinking of the support*, that is, the positivity set  $\mathcal{P}(t) :=$  $\{x \in \mathbb{R}^N : u(t,x) > 0\}$  of u at time t is a relatively compact subset of  $\mathbb{R}^N$  for all  $t \in (0, T_e)$ , even if  $\mathcal{P}(0) = \mathbb{R}^N$  initially [1, 5, 7, 19]. Observe that the inequality q < m is always satisfied when the diffusion is linear (m = 1) or slow (m > 1). Additional information on the behaviour of  $\mathcal{P}(t)$  as  $t \to T_e$  is also available when  $m \ge 1$  and N = 1 [6, 10–14].

Once finite time extinction is known to take place, gaining further insight into the underlying mechanism requires to identify the behaviour of u(t) as  $t \to T_e$ , a preliminary step being to determine

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the relevant space and time scales. Simple scaling arguments predict that, for  $r \in [1, \infty]$  and  $u_0 \in L^r(\mathbb{R}^N)$ , there is a constant  $\gamma_r > 0$  (depending on  $N, m, q, u_0$ , and r) such that

$$\|u(t)\|_r \sim \gamma_r (T_e - t)^{\alpha - (N\beta/r)} , \qquad (1.4)$$

where

$$\alpha := \frac{1}{1-q} > 0 , \qquad \beta := \frac{q-m}{2(1-q)} \in \mathbb{R} .$$
(1.5)

As already observed by several authors [9, 12, 15], a rather simple comparison argument provides a lower bound for the  $L^{\infty}$ -norm of the form (1.4). Indeed, consider  $t \in (0, T_e)$  and let  $x(t) \in \mathbb{R}^N$  be a point where u(t) reaches its maximum value, that is,  $u(t, x(t)) = ||u(t)||_{\infty}$ . Then  $u(t)^m$  also attains its maximum value at this point, so that  $\Delta u^m(t, x(t)) \leq 0$  and we infer from (1.1a) that (at least formally)

$$\frac{d}{dt} \|u(t)\|_{\infty} = \partial_t u(t, x(t)) \le -u(t, x(t))^q = -\|u(t)\|_{\infty}^q .$$

Integrating the above differential inequality over  $(t, T_e)$  gives the expected lower bound

$$||u(t)||_{\infty} \ge [(1-q)(T_e-t)]^{1/(1-q)}, \qquad t \in [0, T_e).$$
 (1.6)

The derivation of an upper bound of the form (1.4) turns out to be more involved and the results obtained so far are rather sparse: in one space dimension, the upper bound

$$||u(t)||_{\infty} \le C_{\infty}(T_e - t)^{1/(1-q)}, \quad t \in [0, T_e),$$
(1.7)

is shown in [17, Proposition 2.2] for m = 1 and in [9, Lemma 5.2, Lemma 7.2 & Lemma 9.2] for  $m \in (0, 1)$ , the latter being valid only for compactly supported initial data. The proofs are however of a completely different nature: in [17], properties of the linear heat equation are used while the approach in [9] relies on the intersection-comparison technique, which requires in particular the compactness of the support of the initial condition. Still for m = 1 but in any space dimension, the upper bound (1.7) is derived in [12, Lemma 2.1] for radially symmetric initial data  $u_0$  having a non-increasing profile and satisfying  $\Delta u_0 + \mu u_0^q \ge 0$  in  $\mathbb{R}^N$  for some  $\mu > 0$ . The last case for which (1.7) is proved corresponds to the choice m = 2 - q > 1 and the proof relies on the derivation of an Aronson-Bénilan estimate, which seems to be only available for this specific choice of the parameters m and q [15].

The purpose of this note is to contribute to the validity of (1.4) and derive optimal upper and lower bounds near the extinction time when the parameters m and q range in

$$\frac{(N-2)_+}{N} < m < q < 1 .$$
(1.8)

Recalling that a lower bound in  $L^{\infty}$  is already available, see (1.6), we begin with upper bounds.

**Theorem 1.1** (Upper bounds). Assume that m and q satisfy (1.8) and consider a non-negative initial condition  $u_0 \in BC(\mathbb{R}^N)$ ,  $u_0 \neq 0$ , for which there is  $\kappa_0 > 0$  such that

$$u_0(x) \le \kappa_0 |x|^{-2/(q-m)}$$
,  $x \in \mathbb{R}^N$ . (1.9)

Given  $r \in [1, \infty]$ , there is  $C_r > 0$  depending only on N, m, q,  $u_0$ , and r such that the solution u to (1.1a)-(1.1b) satisfies

$$|u(t)||_r \le C_r (T_e - t)^{\alpha - (N\beta/r)} , \qquad t \in (0, T_e) , \qquad (1.10)$$

the extinction time  $T_e$  being defined in (1.2).

Theorem 1.1 thus extends the validity of the upper bound (1.7) established in [9] for N = 1 and  $r = \infty$  to any space dimension  $N \ge 1$  and  $r \in [1, \infty]$ , while relaxing the assumption of compact support required in [9]. It is worth mentioning that the validity of (1.10) for  $r \in [1, \infty)$  does not seem to be a simple consequence of (1.10) for  $r = \infty$  since u(t) is positive everywhere in  $\mathbb{R}^N$  for all  $t \in (0, T_e)$  even if  $u_0$  is compactly supported, see [9, Lemma 2.5] and Proposition 1.4 below.

To be able to cope with higher space dimensions and non-compactly supported initial data, the proof of Theorem 1.1 takes a different route from that in [9] and is carried out in two steps: we first show that the algebraic decay at infinity (1.9) enjoyed by  $u_0$  remains true throughout time evolution and combine it with (1.1a) to prove (1.10) for r = 1. We next use self-similar variables and Moser's interation technique to derive (1.10) for all  $r \in (1, \infty]$ .

As a consequence of (1.6) and Theorem 1.1 for  $r = \infty$ , the correct time scale for the extinction phenomenon is identified. We now supplement the lower bound (1.6) in  $L^{\infty}$  with another one in  $L^{m+1}$ . On the one hand, it allows us to identify the right space scale. On the other hand, its derivation does not rely on the comparison principle but on energy estimates, a technique which is more likely to extend to other problems for which the former might not be available.

**Theorem 1.2** (Lower bound in  $L^{m+1}$ ). Assume that m and q satisfy (1.8) and consider a nonnegative initial condition  $u_0 \in BC(\mathbb{R}^N)$ ,  $u_0 \not\equiv 0$ , such that  $u_0 \in L^{m+1}(\mathbb{R}^N)$ . There is  $c_{m+1} > 0$ depending only on N, m, q, and  $u_0$  such that the solution u to (1.1a)-(1.1b) satisfies

$$||u(t)||_{m+1} \ge c_{m+1} (T_e - t)^{\alpha - (N\beta/(m+1))} , \qquad t \in (0, T_e) , \qquad (1.11)$$

the extinction time  $T_e$  being defined in (1.2).

Observing that Theorems 1.1 and 1.2 are shown without using the  $L^{\infty}$ -lower bound (1.6), the latter may be recovered from these two results by Hölder's inequality, with a less explicit constant though.

**Corollary 1.3.** Assume that m and q satisfy (1.8) and consider a non-negative initial condition  $u_0 \in BC(\mathbb{R}^N), u_0 \neq 0$ , enjoying the decay property (1.9). For  $r \in (m + 1, \infty]$ , there is  $c_r > 0$  depending only N, m, q,  $u_0$ , and r such that the solution u to (1.1a)-(1.1b) satisfies

$$||u(t)||_r \ge c_r (T_e - t)^{\alpha - (N\beta/r)} , \qquad t \in (0, T_e) .$$
(1.12)

Summarizing the outcome of Theorem 1.1, Theorem 1.2, and Corollary 1.3, we have shown that, for all non-negative initial data  $u_0 \in BC(\mathbb{R}^N)$ ,  $u_0 \not\equiv 0$ , enjoying the decay property (1.9), the corresponding solution u to (1.1a)-(1.1b) is bounded in  $L^r(\mathbb{R}^N)$ ,  $r \in [m+1, \infty]$ , from above and from below at time  $t \in (0, T_e)$  by the same power of  $T_e - t$ . Such estimates pave the way towards a more precise description of the behaviour of u(t) as  $t \to T_e$ , which is expected to be self-similar. That this is indeed the case is shown in [8,9] in one space dimension, another building block of the proof being the uniqueness of self-similar solutions [8].

We end up this note with the already mentioned everywhere positivity of solutions to (1.1a)-(1.1b) for positive times prior to the extinction time. As we shall see below, this property holds true for a wider range of the parameters m and q, namely  $0 < m \le q < 1$ . It is already observed in [9, Lemma 2.5] in one space dimension and we extend it herein to any space dimension. It is worth emphasizing that it includes the case q = m and contrasts markedly with the instantaneous shrinking of the support occurring when q < m.

**Proposition 1.4** (Everywhere positivity). Consider  $0 < m \leq q < 1$ . Let  $u_0 \in BC(\mathbb{R}^N)$  be a nonnegative initial condition,  $u_0 \not\equiv 0$ , and denote the corresponding solution to (1.1a)-(1.1b) by u with extinction time  $T_e$ . For  $t \in (0, T_e)$ , there holds

$$\mathcal{P}(t) := \{ x \in \mathbb{R}^N : u(t, x) > 0 \} = \mathbb{R}^N .$$
(1.13)

Before proving the results stated above, we point out once more that the energy techniques developed herein seem to be rather flexible and are expected to have a wider range of applicability. For instance, a related approach is used in the companion paper [18], where optimal (lower and upper) bounds near the extinction time are established for a different fast diffusion equation (featuring the *p*-Laplacian operator,  $p \in (1, 2)$ ) with a gradient absorption term.

### 2. Upper bounds near the extinction time

Throughout this section, we assume that m and q satisfy (1.8) and consider a non-negative initial condition  $u_0 \in BC(\mathbb{R}^N)$ ,  $u_0 \not\equiv 0$ , enjoying the decay property (1.9). Let u be the corresponding solution to (1.1a)-(1.1b).

2.1.  $L^1$ -estimate. We begin with the propagation throughout time evolution of the algebraic decay (1.9) and set

$$\kappa_* := \left(\frac{2m(m+q)}{(q-m)^2}\right)^{1/(q-m)} . \tag{2.1}$$

**Lemma 2.1.** For  $t \in [0, \infty)$  and  $x \in \mathbb{R}^N \setminus \{0\}$ , there holds

$$u(t,x) \le \max\{\kappa_0,\kappa_*\}|x|^{-2/(q-m)}$$

*Proof.* Set  $\Sigma_{\kappa}(x) := \kappa |x|^{-2/(q-m)}$  for  $x \in \mathbb{R}^N \setminus \{0\}$ , where  $\kappa$  is a positive constant yet to be determined. We note that

$$-\Delta \Sigma^{m}(x) + \Sigma(x)^{q} = -\kappa^{m} \left[ \frac{2m(m+q)}{(q-m)^{2}} |x|^{-2q/(q-m)} - \frac{2m(N-1)}{(q-m)} |x|^{-2q/(q-m)} \right] + \kappa^{q} |x|^{-2q/(q-m)} \geq \kappa^{m} \left( \kappa^{q-m} - \kappa_{*}^{q-m} \right) |x|^{-2q/(q-m)}$$

for  $x \in \mathbb{R}^N \setminus \{0\}$ , so that  $\Sigma_{\kappa}$  is a supersolution to (1.1a) in  $\mathbb{R}^N \setminus \{0\}$  for all  $\kappa \geq \kappa_*$ . We then choose  $\kappa = \max\{\kappa_0, \kappa_*\}$  and use the comparison principle to complete the proof of Lemma 2.1.

We are now in a position to derive the claimed upper bound near the extinction time for r = 1. Proof of Theorem 1.1: r = 1. Let  $t \in [0, T_e)$ . Integrating (1.1a) over  $(t, T_e) \times \mathbb{R}^N$  gives

$$\|u(t)\|_{1} = \int_{t}^{T_{e}} \int_{\mathbb{R}^{N}} u(s, x)^{q} \, dx ds \,.$$
(2.2)

Owing to (1.8), there holds 2q/(q-m) > N and we infer from Lemma 2.1 and Hölder's inequality that, for  $s \in (t, T_e)$  and R > 0,

$$\begin{split} \int_{\mathbb{R}^N} u(s,x)^q \, dx &= \int_{B(0,R)} u(s,x)^q \, dx + \int_{\mathbb{R}^N \setminus B(0,R)} u(s,x)^q \, dx \\ &\leq C R^{N(1-q)} \| u(s) \|_1^q + (\max\{\kappa_0,\kappa_*\})^q |\mathbb{S}^{N-1}| \int_R^\infty r^{N-1-(2q/(q-m))} \, dr \\ &\leq C \left( R^{N(1-q)} \| u(s) \|_1^q + R^{(N(q-m)-2q)/(q-m)} \right) \, . \end{split}$$

We next optimize in R in the previous inequality by setting  $R(s) := \|u(s)\|_1^{-(q-m)/(N(m-q)+2)}$ , which satisfies 2).

$$R(s)^{N(1-q)} \|u(s)\|_{1}^{q} = R(s)^{(N(q-m)-2q)/(q-m)} = \|u(s)\|_{1}^{(N(m-q)+2q)/(N(m-q)+2q)}$$

Consequently, taking R = R(s) in the previous inequality, we obtain

$$\int_{\mathbb{R}^N} u(s,x)^q \, dx \le C \|u(s)\|_1^{(N(m-q)+2q)/(N(m-q)+2)} \, ,$$

which gives, together with (2.2), the positivity of N(m-q) + 2q, and the time monotonicity of  $s \mapsto \|u(s)\|_1,$ 

$$\begin{aligned} \|u(t)\|_1 &\leq C \int_t^{T_e} \|u(s)\|_1^{(N(m-q)+2q)/(N(m-q)+2)} ds \\ &\leq C(T_e-t)\|u(t)\|_1^{(N(m-q)+2q)/(N(m-q)+2)} , \end{aligned}$$

from which (1.10) for r = 1 readily follows.

2.2. Scaling variables and L<sup>r</sup>-estimates,  $r \in (1, \infty]$ . The next step is to take advantage of the just derived  $L^1$ -upper bound to derive the corresponding ones in  $L^r$  for  $r \in (1, \infty]$ . To this end, we introduce the scaling variables

$$s := \ln\left(\frac{T_e}{T_e - t}\right) , \qquad y := x(T_e - t)^{\beta} , \qquad (t, x) \in [0, T_e) \times \mathbb{R}^N , \qquad (2.3)$$

and the new unknown function v defined by

$$u(t,x) = (T_e - t)^{\alpha} v \left( \ln(T_e) - \ln(T_e - t), x(T_e - t)^{\beta} \right) , \qquad (t,x) \in [0,T_e) \times \mathbb{R}^N , \qquad (2.4)$$

or, equivalently,

$$v(s,y) = T_e^{-\alpha} e^{\alpha s} u \left( T_e(1-e^{-s}), y T_e^{-\beta} e^{\beta s} \right) , \qquad (s,y) \in [0,\infty) \times \mathbb{R}^N .$$
 (2.5)

It readily follows from (1.1a)-(1.1b) that v solves

$$\partial_s v(s,y) = \alpha v(s,y) + \beta y \cdot \nabla v(s,y) + \Delta v^m(s,y) - v(s,y)^q , \qquad (s,y) \in (0,\infty) \times \mathbb{R}^N , \qquad (2.6)$$

$$v(0,y) = v_0(y) := T_e^{-\alpha} u_0 \left( y T_e^{-\beta} \right) , \qquad y \in \mathbb{R}^N .$$
 (2.7)

Since

$$||u(t)||_r = (T_e - t)^{\alpha - (N\beta/r)} ||v(s)||_r , \qquad t \in (0, T_e) , \qquad (2.8)$$

for all  $r \in [1, \infty]$ , we realize that an upper bound such as (1.10) on  $||u(t)||_r$  for  $t \in (0, T_e)$  obviously follows from a uniform upper bound on  $||v(s)||_r$  for  $s \ge 0$ , the converse being true as well. In particular, it follows from (2.8) and Theorem 1.1 for r = 1 that

$$\|v(s)\|_1 \le C_1 , \qquad s \ge 0 , \tag{2.9}$$

and we may assume without loss of generality that  $C_1 \geq 1$ .

We now aim at using a bootstrap argument to deduce from (2.6) and (2.9) that v belongs to  $L^{\infty}(0,\infty; L^r(\mathbb{R}^N))$  for all  $r \in (1,\infty]$ . To this end, Moser's iteration technique is a suitable tool and the way we apply it is inspired from [2, Theorem 3.1]. But since [2, Theorem 3.1] is devoted to the slow diffusion case m > 1, some technical aspects of its proof do not seem to apply directly here and we borrow additional arguments from the proof of [3, Proposition 2].

# **Lemma 2.2.** Let $r \in (0, \infty]$ . There is $C_{r+1} > 0$ depending only on N, m, q, $u_0$ , and r such that $\|v(s)\|_{r+1} \leq C_{r+1}$ , $s \geq 0$ .

*Proof.* Let  $r \in [2 - m, \infty)$ . Multiplying (2.6) by  $v^r$ , integrating over  $\mathbb{R}^N$ , and using integration by parts, we obtain

$$\frac{1}{r+1}\frac{d}{ds}\|v\|_{r+1}^{r+1} + rm\int_{\mathbb{R}^N} v^{r+m-2}|\nabla v|^2 \, dy + \int_{\mathbb{R}^N} v^{r+q} \, dy = \left(\alpha - \frac{N\beta}{r+1}\right)\|v\|_{r+1}^{r+1} \, ,$$
$$\frac{d}{ds}\|v\|_{r+1}^{r+1} + \frac{4mr(r+1)}{(m+r)^2}\left\|\nabla v^{(m+r)/2}\right\|_2^2 \le \alpha(r+1)\|v\|_{r+1}^{r+1} \, .$$

Since  $4mr(r+1) \ge 2m(m+r)^2$ , we end up with

$$\frac{d}{ds} \|v\|_{r+1}^{r+1} + 2m \left\|\nabla v^{(m+r)/2}\right\|_2^2 \le \alpha(r+1) \|v\|_{r+1}^{r+1}.$$
(2.10)

We next fix  $\zeta \in (2/m, 2^*)$  where  $2^* := 2N/(N-2)_+$  (with  $2^* = \infty$  for N = 1, 2). On the one hand, it follows from the Gagliardo-Nirenberg inequality that

$$\left\| v^{(m+r)/2} \right\|_{\zeta} \le C \left\| \nabla v^{(m+r)/2} \right\|_{2}^{\theta} \left\| v^{(m+r)/2} \right\|_{1}^{1-\theta} , \qquad (2.11)$$

with

$$\theta := \frac{2N(\zeta - 1)}{(N+2)\zeta} \; .$$

On the other hand, since  $(m+r)/2 \in [1, \zeta(m+r)/2]$  for  $r \geq 2-m$ , we infer from Hölder's inequality that

$$\|v\|_{(m+r)/2}^{(m+r)/2} \le \|v\|_{\zeta(m+r)/2}^{\zeta(m+r)(m+r-2)/2[\zeta(m+r)-2]} \|v\|_{1}^{(\zeta-1)(m+r)/[\zeta(m+r)-2]} .$$
(2.12)

We deduce from (2.11) and (2.12) that

$$\begin{aligned} \|v\|_{\zeta(m+r)/2}^{(m+r)/2} &= \left\|v^{(m+r)/2}\right\|_{\zeta} \le C \|\nabla v^{(m+r)/2}\|_{2}^{\theta} \left(\|v\|_{(m+r)/2}^{(m+r)/2}\right)^{1-\theta} \\ &\le C \left\|\nabla v^{(m+r)/2}\right\|_{2}^{\theta} \left[\|v\|_{\zeta(m+r)/2}^{\zeta(m+r)(m+r-2)/2[\zeta(m+r)-2]}\|v\|_{1}^{(\zeta-1)(m+r)/[\zeta(m+r)-2]}\right]^{1-\theta} ,\end{aligned}$$

hence

$$\|v\|_{\zeta(m+r)/2}^{\zeta(m+r)[N(m+r)+2-N]/N[\zeta(m+r)-2]} \le C \|\nabla v^{(m+r)/2}\|_2^2 \|v\|_1^{[2N-(N-2)\zeta](m+r)/N[\zeta(m+r)-2]} .$$
(2.13)

Moreover, since  $\zeta > 2/m$  and m < 1, we have  $2r \le m[\zeta(m+r) - 2]$ , hence

$$\frac{[2N - (N-2)\zeta](m+r)}{N[\zeta(m+r) - 2]} \le \frac{[2N - (N-2)\zeta]}{N} \frac{m(m+r)}{2r} \le \frac{m[2N - (N-2)\zeta]}{N} ,$$

so that

$$\begin{aligned} \|v\|_{1}^{[2N-(N-2)\zeta](m+r)/N[\zeta(m+r)-2]} &\leq C_{1}^{[2N-(N-2)\zeta](m+r)/N[\zeta(m+r)-2]} \\ &\leq C_{1}^{m[2N-(N-2)\zeta]/N} . \end{aligned}$$
(2.14)

Also,

$$1 - \frac{N[\zeta(m+r) - 2]}{\zeta[N(m+r) + 2 - N]} = \frac{2N - \zeta(N - 2)}{\zeta[N(m+r) + 2 - N]} > 0 ,$$

and we infer from (2.13), (2.14), and Young's inequality that

$$\begin{aligned} \|v\|_{\zeta(m+r)/2}^{m+r} &\leq \frac{N[\zeta(m+r)-2]}{\zeta[N(m+r)+2-N]} \|v\|_{\zeta(m+r)/2}^{\zeta(m+r)[N(m+r)+2-N]/N[\zeta(m+r)-2]} \\ &+ \frac{2N-\zeta(N-2)}{\zeta[N(m+r)+2-N]} \\ &\leq C \left\|\nabla v^{(m+r)/2}\right\|_{2}^{2} + 1 \ . \end{aligned}$$

Therefore, there is  $\nu \in (0, 1)$  depending only on N, m, q, and  $u_0$  such that

$$\nu\left(\|v\|_{\zeta(m+r)/2}^{m+r}-1\right) \le \left\|\nabla v^{(m+r)/2}\right\|_{2}^{2} .$$
(2.15)

Moreover, since  $r + 1 \in [1, \zeta(m + r)/2]$ , it follows from (2.9) and Hölder's and Young's inequalities that

$$\begin{aligned} \|v\|_{r+1}^{r+1} &\leq \|v\|_{\zeta(m+r)/2}^{\zeta(r(r+m)/[\zeta(r+m)-2]} \|v\|_{1}^{[(\zeta-2)r+\zeta m-2]/(\zeta r+\zeta m-2)} \\ &\leq C_{1}^{[(\zeta-2)r+\zeta m-2]/(\zeta r+\zeta m-2)} \|v\|_{\zeta(m+r)/2}^{\zeta r(r+m)/[\zeta(r+m)-2]} \\ &\leq C_{1} \|v\|_{\zeta(m+r)/2}^{\zeta r(r+m)/[\zeta(r+m)-2]} \\ &\leq \frac{\zeta m-2}{\zeta(m+r)-2} C_{1}^{[\zeta(m+r)-2]/(\zeta m-2)} + \frac{\zeta r}{\zeta(m+r)-2} \|v\|_{\zeta(m+r)/2}^{m+r} \\ &\leq C_{1}^{[\zeta(m+r)-2]/(\zeta m-2)} + \|v\|_{\zeta(m+r)/2}^{m+r} \end{aligned}$$
(2.16)

Next, let  $\sigma > 1$  to be chosen appropriately later on and set

$$\mathcal{I}_{r}(s) := \int_{\mathbb{R}^{N}} v(s, y)^{\zeta[r+m+\sigma(1-m)]/[\sigma(\zeta-2)+2]} \, dy \, , \qquad s \ge 0$$

Since  $\sigma(\zeta - 2) + 2 \in [\zeta, \sigma\zeta]$  and

$$r+1 = \frac{\sigma-1}{\sigma}(m+r) + \frac{m+r+\sigma(1-m)}{\sigma}$$

we deduce from (2.9) and Hölder's and Young's inequalities that, for  $\delta > 0$ ,

$$\|v\|_{r+1}^{r+1} \le \|v\|_{\zeta(m+r)/2}^{(\sigma-1)(m+r)/\sigma} \mathcal{I}_{r}^{[\sigma(\zeta-2)+2]/\sigma\zeta} \le \frac{(\sigma-1)\delta}{\sigma} \|v\|_{\zeta(m+r)/2}^{m+r} + \frac{1}{\sigma\delta^{\sigma-1}} \mathcal{I}_{r}^{[\sigma(\zeta-2)+2]/\zeta} \le \delta \|v\|_{\zeta(m+r)/2}^{m+r} + \frac{1}{\delta^{\sigma-1}} \mathcal{I}_{r}^{[\sigma(\zeta-2)+2]/\zeta} .$$
(2.17)

,

Combining (2.10), (2.15), and (2.17) leads us to

$$\begin{aligned} \frac{d}{ds} \|v\|_{r+1}^{r+1} &+ 2m\nu \left( \|v\|_{\zeta(m+r)/2}^{m+r} - 1 \right) \leq \frac{d}{ds} \|v\|_{r+1}^{r+1} + 2m \left\| \nabla v^{(m+r)/2} \right\|_{2}^{2} \\ &\leq \alpha(r+1) \|v\|_{r+1}^{r+1} \\ &\leq \alpha\delta(r+1) \|v\|_{\zeta(m+r)/2}^{m+r} + \frac{\alpha(r+1)}{\delta^{\sigma-1}} \mathcal{I}_{r}^{[\sigma(\zeta-2)+2]/\zeta} . \end{aligned}$$

We then choose  $\delta=m\nu/\alpha(r+1)$  in the above inequality to obtain

$$\frac{d}{ds} \|v\|_{r+1}^{r+1} + m\nu \|v\|_{\zeta(m+r)/2}^{m+r} \le 2m\nu + \frac{\alpha^{\sigma}(r+1)^{\sigma}}{(m\nu)^{\sigma-1}} \mathcal{I}_r^{[\sigma(\zeta-2)+2]/\zeta} .$$

We finally use (2.16) to estimate from below the second term of the left-hand side of the previous inequality and end up with

$$\frac{d}{ds} \|v\|_{r+1}^{r+1} + m\nu\|v\|_{r+1}^{r+1} \le 2m\nu + m\nu C_1^{[\zeta(m+r)-2]/(\zeta m-2)} + \frac{\alpha^{\sigma}(r+1)^{\sigma}}{(m\nu)^{\sigma-1}} \mathcal{I}_r^{[\sigma(\zeta-2)+2]/\zeta} .$$
(2.18)

We first choose

$$\sigma = \frac{\zeta(m+r) - 2}{\zeta m - 2} > 1$$

in (2.18) and observe that this choice guarantees that

$$\zeta[m+r+\sigma(1-m)] = \sigma(\zeta-2)+2 \; .$$

Consequently,  $\mathcal{I}_r = ||v||_1$  and we deduce from (2.9) and (2.18) that there is C(r) > 0 depending on  $N, m, q, u_0$ , and r such that

$$\frac{d}{ds} \|v\|_{r+1}^{r+1} + m\nu \|v\|_{r+1}^{r+1} \le C(r)$$

Integrating the previous differential inequality entails that

$$\sup_{s \ge 0} \|v(s)\|_{r+1} < \infty .$$
(2.19)

The validity of (2.19) extends to all  $r \in (0, 2 - m)$  by (2.9) and Hölder's inequality.

To complete the proof of Lemma 2.2, we are left to check the boundedness of v in  $L^{\infty}(\mathbb{R}^N)$ . To this end, we take  $\sigma = \sigma_0 := 2(\zeta - 1)/(\zeta - 2) > 1$  in (2.18) and obtain, after integration with respect to time,

$$\begin{aligned} \|v(s)\|_{r+1}^{r+1} &\leq \|v_0\|_{r+1}^{r+1} e^{-m\nu s} + 2 + C_1^{[\zeta(m+r)-2]/(\zeta m-2)} \\ &+ \left(\frac{\alpha(r+1)}{m\nu}\right)^{\sigma_0} \left[\sup_{s_* \in [0,s]} \mathcal{I}_r(s_*)\right]^2 \\ &\leq \|v_0\|_1 \|v_0\|_{\infty}^r + 2 + C_1^{[\zeta(m+r)-2]/(\zeta m-2)} \\ &+ \left(\frac{\alpha(r+1)}{m\nu}\right)^{\sigma_0} \left[\sup_{s_* \in [0,s]} \mathcal{I}_r(s_*)\right]^2 ,\end{aligned}$$

and

$$\mathcal{I}_r = \|v\|_{[(r+m)+\sigma_0(1-m)]/2}^{[(r+m)+\sigma_0(1-m)]/2}$$

Therefore, there are  $K_0 > 0$  and  $K_1 > 0$  depending only on N, m, q, and  $u_0$  such that

$$\sup_{s \ge 0} \left\{ \|v(s)\|_{r+1}^{r+1} \right\} \le K_0 \left( K_1^{r+1} + (1+r)^{\sigma_0} \sup_{s \ge 0} \left\{ \|v(s)\|_{[(r+m)+\sigma_0(1-m)]/2}^{[(r+m)+\sigma_0(1-m)]/2} \right\} \right) .$$
(2.20)

We now define the sequence  $(r_j)_{j\geq 0}$  by

$$1 + r_{j+1} = 2(1 + r_j) - (1 - m)(\sigma_0 - 1) , \qquad j \ge 0 , \qquad r_0 := 2 - m ,$$

and set

$$V_j := \sup_{s \ge 0} \left\{ \|v(s)\|_{r_j+1}^{r_j+1} \right\} , \qquad j \ge 0 .$$

For  $j \ge 0$ , we take  $r = r_{j+1}$  in (2.20) and realize that

$$V_{j+1} \le K_0 \left( K_1^{1+r_{j+1}} + (1+r_{j+1})^{\sigma_0} V_j^2 \right)$$
  
$$\le K_0 (1+r_{j+1})^{\sigma_0} \max \left\{ K_1^{1+r_{j+1}}, V_j^2 \right\} , \qquad j \ge 0 .$$

Since  $\sigma_0 - 1 < 1/(1-m)$  thanks to the constraint  $\zeta > 2/m$ , one has  $1 + r_0 - (1-m)(\sigma_0 - 1) > 0$  and we are in a position to apply [22, Lemma A.1], which we recall in Lemma 2.3 below for completeness, to conclude that there is  $K_2 > 0$  depending only on m,  $\zeta$ ,  $K_0$ , and  $K_1$  such that

$$V_j^{1/(1+r_j)} \le K_2 , \qquad j \ge 0 .$$

Equivalently,

$$\sup_{s \ge 0} \{ \|v(s)\|_{1+r_j} \} \le K_2 , \qquad j \ge 0 ,$$

and letting  $j \to \infty$  entails that  $||v(s)||_{\infty} \leq K_2$  for all  $s \geq 0$ , thereby completing the proof of Lemma 2.2.

The proof of Theorem 1.1 for  $r \in (1, \infty]$  is now a straightforward consequence of (2.8) and Lemma 2.2.

**Lemma 2.3.** Let a > 1,  $b \ge 0$ ,  $c \in \mathbb{R}$ ,  $C_0 \ge 1$ ,  $C_1 \ge 1$ , and  $p_0 > 0$  be given such that  $p_0(a-1)+c > 0$ . We define the sequence  $(p_k)_{k\ge 0}$  of positive real numbers by  $p_{k+1} = ap_k + c$  for  $k \ge 0$  and assume that  $(Q_k)_{k\ge 0}$  is a sequence of positive real numbers satisfying

$$Q_0 \le C_1^{p_0}$$
,  $Q_{k+1} \le C_0 p_{k+1}^b \max\left\{C_1^{p_{k+1}}, Q_k^a\right\}, \ k \ge 0$ .

Then the sequence  $\left(Q_k^{1/p_k}\right)_{k\geq 0}$  is bounded.

### 3. Lower bound near the extinction time

We now turn to the lower bound near the extinction time in  $L^{m+1}(\mathbb{R}^N)$ .

Proof of Theorem 1.2. For  $t \in [0, T_e]$ , we define

$$X(t) := ||u(t)||_{m+1}^{m+1}$$
 and  $Y(t) := \int_{\mathbb{R}^N} u(t, x)^{m+q} dx$ 

Let  $t \in (0, T_e)$ . It follows from (1.1a) that

$$\frac{1}{m+1}\frac{dX}{dt}(t) + \|\nabla u^m(t)\|_2^2 + Y(t) = 0.$$
(3.1)

Since

$$1 < \frac{m+q}{m} < \frac{m+1}{m} < 2^* := \frac{2N}{(N-2)_+}$$

by (1.8) we infer from the Gagliardo-Nirenberg inequality that

$$X(t)^{m/(m+1)} = \|u(t)^m\|_{(m+1)/m} \le C \|\nabla u^m(t)\|_2^{\theta} \|u(t)^m\|_{(m+q)/m}^{1-\theta} \le CY(t)^{m(1-\theta)/(m+q)} \|\nabla u^m(t)\|_2^{\theta} ,$$

where

$$\theta := \frac{2Nm(1-q)}{(m+1)[m(N+2) - q(N-2)]} \ .$$

Consequently, since  $u(t) \neq 0$  as  $t \in (0, T_e)$ ,

$$\|\nabla u^m(t)\|_2^2 \ge CX(t)^{2m/\theta(m+1)}Y(t)^{-2m(1-\theta)/\theta(m+q)} ,$$

which gives, together with (3.1),

$$\frac{dX}{dt}(t) + CX(t)^{2m/\theta(m+1)}Y(t)^{-2m(1-\theta)/\theta(m+q)} + (m+1)Y(t) \le 0.$$
(3.2)

Setting

$$\xi := 1 + \frac{2m(1-\theta)}{\theta(m+q)} > 1$$
 and  $\gamma := \frac{2m}{\theta\xi(m+1)}$ ,

it follows from Young's inequality that

$$X(t)^{\gamma} = X(t)^{\gamma} Y(t)^{-(\xi-1)/\xi} Y(t)^{(\xi-1)/\xi} \le \frac{1}{\xi} X(t)^{\xi\gamma} Y(t)^{1-\xi} + \frac{\xi-1}{\xi} Y(t) \le X(t)^{2m/\theta(m+1)} Y(t)^{-2m(1-\theta)/\theta(m+q)} + Y(t)$$

Combining this inequality with (3.2) leads us to the differential inequality

$$\frac{dX}{dt}(t) + CX(t)^{\gamma} \le 0 , \qquad t \in (0, T_e) .$$
(3.3)

Now,

$$\begin{split} \gamma &= \frac{2m}{\theta\xi(m+1)} = \frac{2m(m+q)}{(2m+\theta(q-m))(m+1)} \\ &= \frac{2m(m+q)[m(N+2)-q(N-2)]}{2m\{(m+1)[m(N+2)-q(N-2)]+N(q-m)(1-q)\}} \\ &= \frac{m(N+2)-q(N-2)}{m(N+2)-qN+2} \in (0,1) \;, \end{split}$$

and we integrate (3.3) over  $(t, T_e)$  to obtain

$$-X(t)^{1-\gamma} + (1-\gamma)C(T_e - t) \le 0 , \qquad t \in (0, T_e) .$$

Noticing that

$$(m+1)\alpha - N\beta = \frac{m+1}{1-q} - \frac{N(q-m)}{2(1-q)} = \frac{m(N+2) - qN + 2}{2(1-q)} = \frac{1}{1-\gamma} ,$$

the lower bound (1.11) readily follows from the previous inequality.

We end up this section with the derivation of the lower bound for  $r \in (m+1, \infty]$  from Theorem 1.1 for r = 1 and Theorem 1.2.

Proof of Corollary 1.3. We first note that, owing to (1.8), there holds 2/(q-m) > N and (1.9) entails that  $u_0 \in L^1(\mathbb{R}^N)$ . Since  $u_0$  also belongs to  $L^{\infty}(\mathbb{R}^N)$ , we conclude that  $u_0 \in L^{m+1}(\mathbb{R}^N)$ .

Let  $r \in (m + 1, \infty]$  and  $t \in (0, T_e)$ . We infer from Theorem 1.1, Theorem 1.2, and Hölder's inequality that

$$c_{m+1}^{m+1}(T_e - t)^{(m+1)\alpha - N\beta} \le \|u(t)\|_{m+1}^{m+1} \le \|u(t)\|_r^{m/(r-1)} \|u(t)\|_1^{(r-1-m)/(r-1)} \le C_1(T_e - t)^{(\alpha - N\beta)(r-1-m)/(r-1)} \|u(t)\|_r^{m/(r-1)},$$

from which (1.12) readily follows.

### 

### 4. Everywhere positivity

In this section, we assume that  $0 < m \leq q < 1$  and consider a non-negative initial condition  $u_0 \in BC(\mathbb{R}^N)$ ,  $u_0 \not\equiv 0$ . We denote the corresponding solution to (1.1a)-(1.1b) by u and define its extinction time by (1.2). As in [9], the proof relies on an upper bound for  $\partial_t u$  which we establish now.

**Lemma 4.1.** For t > 0 there holds

$$\partial_t u(t) \le \frac{u(t)}{(1-m)t}$$
 in  $\mathbb{R}^N$ .

When q = m, Lemma 4.1 is a consequence of [4, Theorem 2], the proof relying on an homogeneity argument. Though the operator  $-\Delta u^m + u^q$  is not homogeneous, we may still adapt the proof of [4, Theorem 2] when  $q \ge m$ .

*Proof.* Given a non-negative initial condition  $u_0 \in BC(\mathbb{R}^N)$ , we denote the corresponding solution to (1.1a)-(1.1b) at time  $t \ge 0$  by  $S(t)u_0$ . Recall that, if  $u_0$  and  $v_0$  are two non-negative functions in  $BC(\mathbb{R}^N)$  satisfying  $u_0 \ge v_0$ , then the comparison principle entails  $S(t)u_0 \ge S(t)v_0$  for all  $t \ge 0$ .

**Step 1.** We first claim that, for  $\lambda \geq 1$ ,

$$S(\lambda t)u_0 \le \lambda^{1/(1-m)} S(t) \left(\lambda^{1/(m-1)} u_0\right) , \qquad t \ge 0 .$$
(4.1)

Indeed, setting  $u(t) := S(t)u_0$  for  $t \ge 0$ , the function v defined by  $v(t) := \lambda^{1/(m-1)}S(\lambda t)u_0$  satisfies

$$\partial_t v(t,x) - \Delta v^m(t,x) + v(t,x)^q = \lambda^{m/(m-1)} \partial_t u(\lambda t,x) - \lambda^{m/(m-1)} \Delta u^m(\lambda t,x) + \lambda^{q/(m-1)} u(\lambda t,x)^q = \left(\lambda^{q/(m-1)} - \lambda^{m/(m-1)}\right) u(\lambda t,x)^q \le 0 .$$

Since  $v(0) = \lambda^{1/(m-1)} u_0 \leq u_0$ , we infer from the comparison principle that (4.1) holds true.

Step 2. Now, fix t > 0 and consider h > 0. Since  $\lambda = (1 + h/t) > 1$  and  $m \in (0, 1)$ , we infer from (4.1) and the comparison principle that

$$S(t+h)u_0 - S(t)u_0 = S(\lambda t)u_0 - S(t)u_0$$
  

$$\leq \lambda^{1/(1-m)}S(t) \left(\lambda^{1/(m-1)}u_0\right) - S(t)u_0$$
  

$$\leq \left[\left(1 + \frac{h}{t}\right)^{1/(1-m)} - 1\right]S(t)u_0.$$

Dividing the above inequality by h and passing to the limit as  $h \to 0$  complete the proof.

We now argue as in the proof of [9, Lemma 2.5] to complete the proof of Proposition 1.4.

Proof of Proposition 1.4. Fix  $t \in (0, T_e)$  and assume for contradiction that  $u(t, x_0) = 0$  for some  $x_0 \in \mathbb{R}^N$ . By (1.1a) and Lemma 4.1, there holds

$$-\Delta u^m(t) + u(t)^q + \frac{u(t)}{(1-m)t} \ge 0 \quad \text{in} \quad \mathbb{R}^N ,$$

so that  $u(t)^m$  is a supersolution to

$$-\Delta w + dw = 0$$
 in  $\mathbb{R}^N$ ,

with  $d(x) := u(t, x)^{q-m} + u(t, x)^{1-m}/((1-m)t)$  for  $x \in \mathbb{R}^N$ . Since t > 0 and  $m \le q < 1$ , the function d is non-negative and bounded and we infer from the strong maximum principle [16, Theorem 8.19] that  $u(t)^m \equiv 0$  in  $\mathbb{R}^N$ , contradicting  $t < T_e$ . Consequently,  $u(t)^m$  is positive everywhere in  $\mathbb{R}^N$  and the proof of Proposition 1.4 is complete.

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