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Micro-mechanics and continuum damage mechanics

A. Burr, F. Hild, F. A. Leckie

Summary Continuum Damage Mechanics has been applied successfully to technical problems since the idea was introduced by Kachanov almost 40 years ago. In keeping with the traditions of mechanics, the formulation was based on the results of mechanical tests on specimens whose size is measured in centimeters. To model the observations which describe the deterioration of material properties it was found necessary to introduce internal variables referred to as 'damage'. The approach is phenomenological, with only a minimal attempt to provide a physical interpretation of damage. For this reason the approach has had little appeal to those whose interest is in the physical mechanisms which cause material deterioration. In this presentation a description is given of attempts to develop continuum damage mechanics so that the relationship with the physical mechanism approach is less abrupt. The procedure is illustrated with reference to ceramic matrix composites.

Key words continuum damage, micro-mechanics, ceramic matrix, composite modelling

1

Introduction

By contrast with fracture mechanics, which describes the behavior of a component weakened by the presence of a single well-defined crack, Continuum Damage Mechanics (CDM) attempts to deal with the circumstances when multiple cracking or voidage is evidently the dominant degradation and failure mechanism. The uniform distribution of voids and cracks, which was observed in high-temperature creep rupture, attracted the attention of Kachanov in 1958 [1]. Kachanov was apparently the first to attempt a description of a creep rupture model by introducing an internal damage variable into the constitutive equations describing the creep behavior. While microscopic examinations clearly identified the presence of voids and cracks, the influence of damage and its growth were deduced from the results of the traditional tests performed by mechanical engineers, and given in the form

$$\frac{\dot{\varepsilon}}{\dot{\varepsilon}_0} = \left(\frac{\sigma}{\sigma_0(1-D)} \right)^n, \quad (1)$$

$$\dot{D} = \left(\frac{\sigma}{\sigma_0(1-D)} \right)^p, \quad (2)$$

where n , p , $\dot{\varepsilon}_0$, σ_0 were selected to fit experiments, σ and ε denote stress, and strain, respectively. The damage D was assumed to be a scalar, in spite of contrary microscopic evidence which indicated the damage to have directional properties. Equations of this form have been used

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extensively [2; 3] and the term

$$\frac{\sigma}{1 - D} \quad (3)$$

is referred to as the 'effective stress' [4]. The approach has been fully described and utilized by Lemaitre and co-workers [5]. The procedure has the advantage that it is simple and can be readily related to the results of macroscopic tests, and by simple generalization can provide a means of extending the effects of damage on more complex phenomena, such as anisothermal viscoplasticity [6], by measuring stiffness loss as an indication of damage accumulation [7].

The behavior of concrete is another example which lends itself well to a CDM description. In this case, when concrete is loaded, microcracks appear at the interface of the cementitious matrix and the aggregate. Using the concept of 'effective stress', Mazars [8] has proposed simple damage laws which have proved to be useful in predicting the behavior of complex engineering components such as dams. Other applications can be found in the work of Krajcinovic [2].

In the engineering development of CDM there is little attempt to give a precise physical description of damage, since the focus is on the prediction of component behavior for which the length scale is many times that associated with the physical damage in the material. By temperament and training, design engineers tend to place greater confidence in the results of appropriate mechanical tests than they do in the results of studies of physical mechanisms. However, the inability to respond to the inquiry, 'What is this damage D ?', has limited the wider acceptance of the CDM approach and denied it the benefits to be gained from understanding the physical processes of damage growth. Ashby and Dyson [9] were able to establish the connection between the damage parameter D and the physical deterioration occurring in metals operating at elevated temperatures. A major contribution of this study was to identify the operating condition within which any particular mechanism is operative. Subsequently, Cocks and Leckie [10] recast the results of the Ashby and Dyson study in the CDM format. Using this approach, Hall and Hayhurst [11] demonstrated how CDM procedures can predict with precision the type and growth of damage occurring in complex engineering components.

Recently, the authors were involved in a study of the properties of Ceramic Matrix Composites (CMCs) consisting of a brittle matrix reinforced by continuous fibers and which, when loaded, exhibit multiple matrix cracking. The matrix cracking is known to be accompanied by slipping at the fiber/matrix interface which results in irreversible macroscopic strains. Following recent materials science studies [12], substantial knowledge exists about the damage mechanisms occurring in these materials, and the effect they have on material properties. These circumstances provided an opportunity to develop a CDM description for CMCs based on the results of materials science, which is the subject of this presentation. It is worthy of note that the CDM description was being developed at a time when little material was available to support an extensive test program. These circumstances emphasized the need of an approach, which is sufficiently general to allow for systematic improvement as material science knowledge unfolded, and the results of further tests became available.

2

The formulation of a physically-based CDM model

The general structure for the constitutive equations is obtained by following a thermodynamic approach [13–15] in identifying a finite number of internal state variables. The state of the material is then described by calculating the Helmholtz free energy density ψ , from which the associated thermodynamic forces may be derived. Once the state variables and the associated forces are defined, the formulation is then completed by defining growth laws of the state variables in terms of their associated forces. It is common at this stage [3] to define the damage growth laws by the introduction of a potential of the internal forces. This approach is not used in the present development, and the growth laws are obtained directly from experimental results.

In addition to the total strain tensor, the variables defining the state of the CMC include the size, shape and distribution of the matrix cracks. The materials science studies also indicate [12] that matrix cracking is accompanied by slipping at the fiber/matrix interface, so that the degree of interface slipping must be included in the state variables. The state of the material can also be defined in a very compact form by a scalar which is the Helmholtz free energy. The free energy is calculated by following any reversible path to the current state.

In the present study, the reversible path consists of two steps. The first step consists in introducing into the unloaded cracked material the interface slips by a series of cutting, displacing and resealing operations. This approach was introduced by Volterra [16], and used to analyze the elastic behavior of homogeneous and isotropic media by considering the elastic properties of a cut cylinder [16; 17], as well as inclusions in an infinite medium [18], or to study creeping materials [10]. In each

case, the contribution to the free energy density can be written as

$$\psi^s = \psi^s(\beta_k, \alpha), \quad (4)$$

where each β_k represents an interface slip, and α the resulting macroscopic inelastic strain tensor.

The second step consists in loading the cracked material elastically with the elastic strain ε^e so that the free energy is given by

$$\psi^e = \frac{1}{2} \varepsilon^e : \mathbf{C}(\delta_k) : \varepsilon^e, \quad (5)$$

where \mathbf{C} is the elastic stiffness tensor which is a function of the crack sizes and distributions δ_k , and ε^e is the elastic strain tensor resulting from the application of load.

The total free energy density associated with the present state is given by summing the two energy components so that

$$\psi = \frac{1}{2} \varepsilon^e : \mathbf{C}(\delta_k) : \varepsilon^e + \psi^s(\beta_k, \alpha), \quad (6)$$

where the usual summation convention is used, ':' represents the contraction with respect to two indices. The strain partition $\varepsilon^e = \varepsilon - \alpha$, where ε is the total strain tensor, permits the free energy density to be written as follows:

$$\psi = \frac{1}{2} (\varepsilon - \alpha) : \mathbf{C}(\delta_k) : (\varepsilon - \alpha) + \psi^s(\beta_k, \alpha). \quad (7)$$

The thermodynamic forces associated with the state variables ε , α , δ_k and β_k (i.e. σ , \mathbf{X} , G_k and s_k , respectively) can be found by differentiation of Eq. (7),

$$\dot{\psi} = \sigma : \dot{\varepsilon} + \mathbf{X} : \dot{\alpha} + G_k \cdot \dot{\delta}_k + s_k \cdot \dot{\beta}_k, \quad (8)$$

where

$$\sigma = \frac{\partial \psi}{\partial \varepsilon}, \quad \mathbf{X} = \frac{\partial \psi}{\partial \alpha}, \quad G_k = \frac{\partial \psi}{\partial \delta_k} \quad \text{and} \quad s_k = \frac{\partial \psi}{\partial \beta_k}, \quad (9)$$

and where s_k , G_k have the same tensorial nature as β_k , δ_k and \cdot represents contraction with respect to the number of indices equal to the rank of the tensors β_k and δ_k .

The results obtained in this section are central to the developments in the remainder of this paper. In Sect. 3 the elastic properties resulting from matrix cracking shall be analyzed. The free energy associated with interface slip is discussed in Sect. 4. The growth laws relating damage to the internal forces are discussed in Sects. 5 and 6. Instead of following the conventional procedure of identifying a scalar potential to define the growth laws as used in [13–15], a different approach is described which combines the results of micromechanics with the results of conventional mechanical tests.

3 Elastic properties of isotropic and anisotropic damaged materials

3.1 Isotropic description

The effect of matrix voids and cracks on the elastic properties of materials have received considerable attention. The continuum damage description of the coupling between elasticity and damage has been presented by Lemaitre and Marquis [19], and the micromechanics for specific distributions of cracks have been completed for example by Budiansky [20; 21]. It is the expression for the Helmholtz free energy density ψ which is used to determine the associated internal forces. It proves to be simpler in some problems to use the Gibbs specific enthalpy φ which is the Legendre transformation of the Helmholtz free energy density according to the relation

$$\sigma : \varepsilon = \psi + \varphi \quad (10)$$

The Gibbs specific enthalpy φ shall be introduced later in this section when appropriate.

For an isotropic material, the initial elastic properties are defined by the shear modulus G and the bulk modulus K . If, after degradation, the material remains isotropic the elastic properties are defined by the new values of the shear modulus \tilde{G} and bulk modulus \tilde{K} . The elastic energy density, which is here identical to the Helmholtz free energy density ψ , is written as follows:

$$2\psi = 2\tilde{G}\mathbf{e}:\mathbf{e} + \tilde{K}\varepsilon_v^2, \quad (11)$$

where \mathbf{e} is the deviatoric strain tensor, and ε_v denotes the volumetric strain. The associated force to the infinitesimal strain tensor $\boldsymbol{\varepsilon}$ is the Cauchy stress tensor $\boldsymbol{\sigma}$

$$\boldsymbol{\sigma} = \frac{\partial\psi}{\partial\boldsymbol{\varepsilon}} \quad (12)$$

and the associated force(s) to the damage variable(s) D are obtained by partial differentiation of the state potential with respect to the corresponding internal variables. The exact expression depends on the details of the model, as shown in the following.

A first attempt within the context of CDM to write a state potential for isotropic damaged materials was the phenomenological description of Lemaitre and Chaboche [22]. The formulation is based on the assumption of 'equivalence of the deformations hypothesis' [23] that states that 'any strain constitutive law for a damaged material may be derived in the same way as for a virgin material except that the usual stress is replaced by the effective stress'. It is postulated that in the case of an isotropic material, only one damage variable D is needed. The consequence is that the two elastic parameters are written as a function of one scalar damage variable D

$$\frac{\tilde{G}}{G} = 1 - D, \quad (13)$$

$$\frac{\tilde{K}}{K} = 1 - D. \quad (14)$$

It is noted that with this definition, Poisson's ratio remains unaltered compared to the undamaged material. The associated force to the damage variable D is referred to as the energy release rate density Y [24], and is defined as

$$Y = -\frac{\partial\psi}{\partial D} = \frac{1}{2}(2G\mathbf{e}:\mathbf{e} + K\varepsilon_v^2). \quad (15)$$

These associated forces (here Y) are usually used as driving forces of the internal state variables (here D).

Ladevèze [25] originally proposed a general expression of the Gibbs specific enthalpy as a function of two *independent* damage variables d and δ in the so-called 'kinematics of zeroth order'. However, the same formalism can be used when the Helmholtz free energy density is analyzed, provided a partition of the strain tensor is written in terms of its deviatoric and spherical parts. The two elastic parameters are written as

$$\frac{\tilde{G}}{G} = 1 - d, \quad (16)$$

$$\frac{\tilde{K}}{K} = 1 - \delta. \quad (17)$$

The two forces associated with the damage variables d and δ are respectively given by

$$Y_d = -\frac{\partial\psi}{\partial d} = G\mathbf{e}:\mathbf{e}, \quad (18)$$

$$Y_\delta = -\frac{\partial\psi}{\partial\delta} = \frac{1}{2}K\varepsilon_v^2. \quad (19)$$

These first two models are completely phenomenological. On the other hand, the micromechanics analysis performed by Budiansky [20] predicts the effect of a dilute concentration of voids on

the macroscopic elastic properties. The new values of the shear and bulk moduli resulting from an average void density $\bar{\omega}_v$, are given by

$$\frac{\tilde{G}}{G} = 1 - 15 \frac{1-\nu}{7-5\nu} \bar{\omega}_v, \quad (20)$$

$$\frac{\tilde{K}}{K} = 1 - \frac{3(1-\nu)}{2(1-2\nu)} \bar{\omega}_v, \quad (21)$$

where ν is the Poisson's ratio of the undamaged material. This micromechanics approach gives results which are dependent on a single state variable similar to that defined by Lemaitre and Chaboche in Eq. (13-14). However, there are differences in the predictions of the two approaches as can be readily demonstrated by making the substitution

$$D = 15 \frac{1-\nu}{7-5\nu} \bar{\omega}_v. \quad (22)$$

Then Eqs. (20-21) become

$$\frac{\tilde{G}}{G} = 1 - D, \quad (23)$$

$$\frac{\tilde{K}}{K} = 1 - \frac{9K+8G}{20G} D. \quad (24)$$

The corresponding thermodynamic force or energy release rate density is given by

$$Y = - \frac{\partial \psi}{\partial D} = \frac{1}{2} \left(2G \mathbf{e} : \mathbf{e} + \frac{9K+8G}{20G} K \varepsilon_v^2 \right), \quad (25)$$

which is different from the expression given in Eq. (15).

Budiansky and O'Connell [21] have also performed similar calculations to determine the changes in the elastic properties of an isotropic distribution of penny-shaped cracks. The relationship between the shear and bulk moduli changes, and the average crack density $\bar{\omega}_c$ is given by

$$\frac{\tilde{G}}{G} = 1 - \frac{32(1-\tilde{\nu})(5-\tilde{\nu})}{45(2-\tilde{\nu})} \bar{\omega}_c, \quad (26)$$

$$\frac{\tilde{K}}{K} = 1 - \frac{16(1-\tilde{\nu}^2)}{9(1-2\tilde{\nu})} \bar{\omega}_c, \quad (27)$$

with

$$\bar{\omega}_c = \frac{45}{16(1-\tilde{\nu}^2)} \frac{(\nu-\tilde{\nu})(2-\tilde{\nu})}{[10\nu-\tilde{\nu}(1+3\nu)]}, \quad (28)$$

where $\tilde{\nu}$ denotes the Poisson's ratio of the damaged material. Setting

$$D = \frac{32(1-\tilde{\nu})(5-\tilde{\nu})}{45(2-\tilde{\nu})} \bar{\omega}_c, \quad (29)$$

shows that in this case there is no linear relationship between the damage variable D and the average crack density $\bar{\omega}_c$. It is worth remembering that Eqs. (26-29) are valid when the crack is embedded in homogeneous equivalent material (i.e. using a self-consistent scheme [26]). When $\nu=0$, then Eqs. (26-27) are simplified and are identical with the results given in Eqs. (13-14) with

$$D = \frac{16}{9} \bar{\omega}_c. \quad (30)$$

The above results have been obtained for isotropic and homogeneous damage, and are unable to model the influence of randomly distributed damage. Using a pattern approach [27] applied to a generalized self-consistent scheme [28; 29] it can be shown that the shear modulus of a material containing a distribution of voids is not only dependent upon the average concentration (i.e. the average

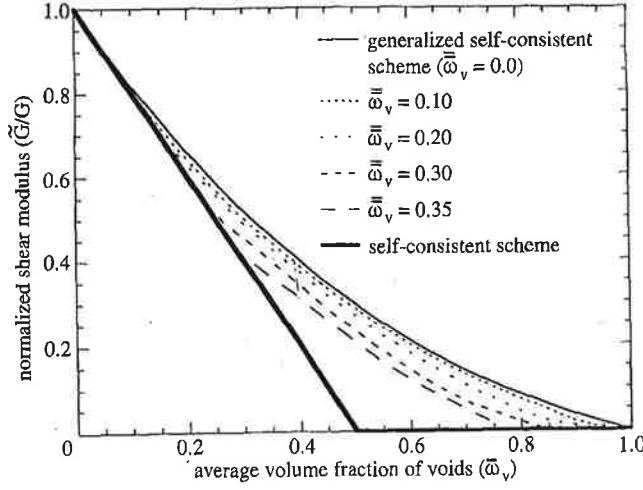


Fig. 1. Normalized shear modulus \tilde{G}/G as a function of volume fraction of voids $\bar{\omega}_v$ and the corresponding standard deviation $\bar{\omega}_v$.

void density $\bar{\omega}_v$) but also on the standard deviation ($\bar{\omega}_v$), as shown in Fig. 1. Therefore, the damage variable $D_v = \tilde{G}/G$ at least depends upon the first two moments of the void distribution. Figure 1 also shows that it is possible to describe continuously the change from the result given by a self-consistent scheme to that given by a generalized self-consistent scheme, by varying the standard deviation of the distribution of voids.

When cracks and cavities are present, the simplest result to be derived is by assuming no interaction between the two defect populations as well as a Taylor scheme. The overall behavior is still isotropic, but now it depends on the statistical distributions of the two populations. From the previous results, the average free energy density ψ is then related to the average free energy density ψ_c of a cracked system, and the average free energy density ψ_v of a system containing voids by

$$\psi = (1 - f)\psi_c(\bar{\omega}_c) + f\psi_v(\bar{\omega}_v), \quad (31)$$

so that

$$\tilde{G}(\bar{\omega}_c, \bar{\omega}_v, \nu) = (1 - f)\tilde{G}_c(\bar{\omega}_c, \nu) + f\tilde{G}_v(\bar{\omega}_v, \nu) = G(1 - d),$$

$$\tilde{K}(\bar{\omega}_c, \bar{\omega}_v, \nu) = (1 - f)\tilde{K}_c(\bar{\omega}_c, \nu) + f\tilde{K}_v(\bar{\omega}_v, \nu) = K(1 - d),$$

where f is the volume proportion of material embedding voids. The average void density and the average crack density are the two independent damage variables required to model the behavior. The forces associated with the average crack density $\bar{\omega}_c$ and the average void density $\bar{\omega}_v$ are defined by

$$Y_c = -\frac{\partial\psi}{\partial\bar{\omega}_c} = \frac{1}{2}(1 - f) \left(2\frac{\partial\tilde{G}_c(\bar{\omega}_c, \nu)}{\partial\bar{\omega}_c} \mathbf{e}:\mathbf{e} + \frac{\partial\tilde{K}_c(\bar{\omega}_c, \nu)}{\partial\bar{\omega}_c} \varepsilon_v^2 \right), \quad (32)$$

$$Y_v = -\frac{\partial\psi}{\partial\bar{\omega}_v} = \frac{1}{2}f \left(2\frac{\partial\tilde{G}_v(\bar{\omega}_v, \nu)}{\partial\bar{\omega}_v} \mathbf{e}:\mathbf{e} + \frac{\partial\tilde{K}_v(\bar{\omega}_v, \nu)}{\partial\bar{\omega}_v} \varepsilon_v^2 \right). \quad (33)$$

If the internal variables selected are the damage variables d and δ , then the associated Y_d and Y_δ defined in Eqs. (18–19) are functions of the average crack density as well as the void density. From this it can be appreciated that the formulation of the free energy from micromechanics has the advantage that it allows more physical insight.

Even though the initial behavior of a damaged material may be isotropic, loading conditions shall result in an evolution of cracks which introduce anisotropic behavior. For example, this is the case in materials experiencing multiple cracking such as concrete, rocks, and brittle matrix composites. In this case, the previous descriptions are not valid, and it is necessary to take the crack orientation into account.

3.2

Anisotropic description

When studying the effect of a distribution of crack orientation, it is easier initially to limit the study to plane stress conditions. This is also a circumstance when it is an advantage to introduce the Gibbs specific enthalpy ϕ following the Legendre transformation of Eq. (10). Figure 2 shows a sheet subjected

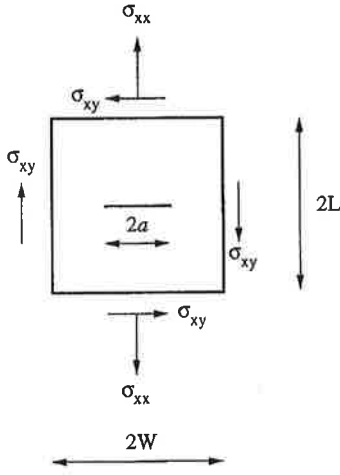


Fig. 2. Elementary cell of size $2L \times 2W$ containing a crack of size $2a$ submitted to normal and shear stresses

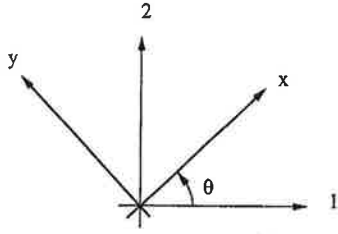


Fig. 3. The initial principal directions of orthotropy, or material directions, 1 and 2 often do not coincide with the loading directions x and y . The angle θ measures their respective orientation

to a tensile stress σ_{xx} and shear stress σ_{xy} with a crack of length $2a$ (σ_{xx} , σ_{yy} and σ_{xy} are the components of the stress tensor σ of the material expressed in the x - y frame, Fig. 3). The change of potential energy density due to the presence of a crack is given by [21]

$$\Delta\varphi = \frac{\sigma_{xx}^2 + \sigma_{xy}^2}{E} \omega_c(a), \quad (34)$$

where E is the Young's modulus of the virgin material, and $\omega_c(a)$ is the crack density assumed to be independent of the crack orientation n . The corresponding compliance variation ΔS at constant stress level is given by

$$\Delta S = \frac{\partial^2 \Delta\varphi}{\partial \sigma \partial \sigma}. \quad (35)$$

Since the energy density variation written in Eq. (34) is proportional to $\sigma_{xx}^2 + \sigma_{xy}^2$, Eq. (35) shows that the compliance variation affects the Young's modulus along the x -direction as well as the shear modulus, so that the expression of the Gibbs specific enthalpy is

$$2\varphi = \frac{\sigma_{xx}^2}{E(1-D_x)} + \frac{2\nu\sigma_{xx}\sigma_{yy}}{E} + \frac{\sigma_{yy}^2}{E} + \frac{\sigma_{xy}^2}{\tilde{G}(D_x)}, \quad (36)$$

with

$$\tilde{G}(D_x) = \frac{G}{1 + \left(\frac{D_x}{1-D_x}\right) \frac{1}{2(1+\nu)}} \quad (37)$$

and

$$D_x = \frac{2\omega_c(a)}{1 + 2\omega_c(a)} \quad (38)$$

where E , ν are the initial elastic properties of the matrix, \tilde{G} is the shear modulus which is a function of damage variable D_x .

The associated force to the damage variable D_x corresponds to the energy release rate density Y_x

$$Y_x = \frac{\partial \varphi}{\partial D_x} = \frac{1}{2} \frac{\sigma_{xx}^2 + \sigma_{xy}^2}{E(1 - D_x)^2} \quad (39)$$

The corresponding Helmholtz free energy density is obtained as a Legendre–Fenchel transformation of Eq. (36)

$$2\psi = \frac{E((1 - D_x)\varepsilon_{xx}^2 + 2\nu(1 - D_x)\varepsilon_{xx}\varepsilon_{yy} + \varepsilon_{yy}^2)}{1 - \nu^2(1 - D_x)} + 4\tilde{G}(D_x)\varepsilon_{xy}^2 \quad (40)$$

and the energy release rate density Y_x is now defined as

$$Y_x = -\frac{\partial \psi}{\partial D_x} \quad (41)$$

and has the same significance and expression as that given in Eq. (39).

When interactions between cracks are accounted for, it is shown that within the framework of the Nonlocal Damage Theory (NDT) [30], the damage variable measured as loss of stiffness not only depends on the crack density ω_c but also on the interaction between cracks. For instance, in tension it is shown that the damage variable D of cracked systems along one direction, measured as a loss of stiffness in the perpendicular direction, is given by

$$D = \frac{H\omega_c}{1 + H\omega_c} \quad (42)$$

with

$$\omega_c = \frac{\pi a^2}{4LW} \quad (43)$$

where $2L$ and $2W$ denote the cell size as well as the distance between two horizontal and two vertical cracks. The interaction coefficient H represents the interaction between the cracks in the considered array [31]. It is worth noting that Eq. (38) is similar to Eq. (43). However, depending on the micromechanical model, the relationship between the damage variable and the crack density is different.

The previous results can be generalized to materials with a distribution of crack size a and orientations \mathbf{n} . This distribution is characterized by a probability density function $f(a, \mathbf{n})$. The average energy variation $\Delta\bar{\varphi}$ is then expressed as

$$\Delta\bar{\varphi} = \frac{1}{2\pi} \int_{\mathcal{C}} \int_0^{+\infty} \Delta\varphi(a, \mathbf{n}) f(a, \mathbf{n}) da d\mathbf{l} \quad (44)$$

where \mathcal{C} is the circle of unit directions \mathbf{n} . If the interactions are neglected, the energy variation for a crack of size a and orientation \mathbf{n} is given by

$$\Delta\varphi(a, \mathbf{n}) = \omega_c(a) \frac{\mathbf{n} \cdot \boldsymbol{\sigma} \cdot \boldsymbol{\sigma} \cdot \mathbf{n}}{E} \quad (45)$$

The average compliance variation ΔS at constant stress level is then given by

$$\Delta S = \frac{\partial^2 \Delta\bar{\varphi}}{\partial \boldsymbol{\sigma} \partial \boldsymbol{\sigma}} = \frac{1}{2\pi} \frac{1}{E} \int_{\mathcal{C}} \langle \omega_c(\mathbf{n}) \rangle \mathbf{N}(\mathbf{n}) d\mathbf{l} \quad (46)$$

with

$$\langle \omega_c(\mathbf{n}) \rangle = \int_0^{+\infty} \omega_c(a) f(a, \mathbf{n}) da,$$

$$N_{\alpha\beta\gamma\delta}(\mathbf{n}) = \frac{1}{2} (n_\alpha n_\gamma \delta_{\beta\delta} + n_\alpha n_\delta \delta_{\gamma\beta} + n_\beta n_\gamma \delta_{\alpha\delta} + n_\beta n_\delta \delta_{\alpha\gamma}) \quad \alpha, \beta, \gamma, \delta = 1, 2$$

where δ is the Kronecker symbol. The average crack density $\langle \omega_c(\mathbf{n}) \rangle$ along a normal \mathbf{n} can be expanded through a modified Fourier representation [32]

$$\langle \omega_c(\mathbf{n}) \rangle = \omega_0 1 + \omega_{\alpha\beta} f_{\alpha\beta} + \omega_{\alpha\beta\gamma\delta} f_{\alpha\beta\gamma\delta} + \dots, \quad (47)$$

where $1, f_{\alpha\beta}, f_{\alpha\beta\gamma\delta}, \dots$ are the basis functions, and $\omega_0, \omega_{\alpha\beta}, \omega_{\alpha\beta\gamma\delta}, \dots$ are the Fourier coefficients. The basis functions $1, f_{\alpha\beta}$, and $f_{\alpha\beta\gamma\delta}, \dots$ are orthogonal to each other so that

$$\int_{\mathcal{C}} 1 f_{\alpha\beta} d\mathbf{l} = \int_{\mathcal{C}} 1 f_{\alpha\beta\gamma\delta} d\mathbf{l} = \int_{\mathcal{C}} f_{\alpha\beta} f_{\gamma\delta\chi\varphi} d\mathbf{l} = \dots = 0. \quad (48)$$

The tensors $f_{\alpha\beta}$ and $f_{\alpha\beta\gamma\delta}$ are completely symmetric with respect to their indices, traceless, and can be defined as follows:

$$f_{\alpha\beta}(\mathbf{n}) = n_\alpha n_\beta - \frac{1}{2} \delta_{\alpha\beta}, \quad (49)$$

$$f_{\alpha\beta\gamma\delta}(\mathbf{n}) = n_\alpha n_\beta n_\gamma n_\delta - \frac{1}{6} (n_\alpha n_\beta \delta_{\gamma\delta} + n_\gamma n_\delta \delta_{\alpha\beta} + n_\alpha n_\gamma \delta_{\beta\delta} + n_\alpha n_\delta \delta_{\gamma\beta} + n_\beta n_\gamma \delta_{\alpha\delta} + n_\beta n_\delta \delta_{\alpha\gamma}) \\ + \frac{1}{6 \times 4} (\delta_{\alpha\beta} \delta_{\gamma\delta} + \delta_{\alpha\gamma} \delta_{\beta\delta} + \delta_{\alpha\delta} \delta_{\beta\gamma}).$$

...

The expansion coefficients $\omega_{\alpha\beta}, \omega_{\alpha\beta\gamma\delta}, \dots$ are also completely symmetric and traceless with respect to their indices. They are expressed as the inner product of $\langle \omega_c(\mathbf{n}) \rangle$ and the basis functions

$$\omega_0 = \frac{1}{2\pi} \int_{\mathcal{C}} \langle \omega_c(\mathbf{n}) \rangle d\mathbf{l}, \quad (50)$$

$$\omega_{\alpha\beta} = \frac{1}{2\pi} \frac{2 \times 4}{2} \int_{\mathcal{C}} \langle \omega_c(\mathbf{n}) \rangle f_{\alpha\beta}(\mathbf{n}) d\mathbf{l},$$

$$\omega_{\alpha\beta\gamma\delta} = \frac{1}{2\pi} \frac{2 \times 4 \times 6 \times 8}{2 \times 3 \times 4} \int_{\mathcal{C}} \langle \omega_c(\mathbf{n}) \rangle f_{\alpha\beta\gamma\delta}(\mathbf{n}) d\mathbf{l},$$

...

These tensors are written for plane problems, the same formalism can be used in 3D problems as shown in [32; 33].

According to Onat and Leckie [32–34], a fourth-order tensor, e.g. the stiffness tensor \mathbf{C} , with the following symmetries

$$C_{\alpha\beta\gamma\delta} = C_{\beta\alpha\gamma\delta} = C_{\alpha\beta\delta\gamma} = C_{\gamma\delta\alpha\beta}, \quad (51)$$

can be written as a unique linear combination of the previous basis functions. The fourth-order tensor $\mathbf{N}(\mathbf{n})$ can be expressed as a combination of the first two basis functions

$$N_{\alpha\beta\gamma\delta}(\mathbf{n}) = \frac{1}{2} (\delta_{\alpha\gamma} \delta_{\beta\delta} + \delta_{\alpha\delta} \delta_{\beta\gamma}) + \frac{1}{2} (f_{\alpha\gamma}(\mathbf{n}) \delta_{\beta\delta} + f_{\alpha\delta}(\mathbf{n}) \delta_{\gamma\beta} + f_{\beta\gamma}(\mathbf{n}) \delta_{\alpha\delta} + f_{\beta\delta}(\mathbf{n}) \delta_{\alpha\gamma}), \quad (52)$$

and the compliance variation depends upon the first two coefficients of the series expansion (47)

$$\Delta \bar{S}_{\alpha\beta\gamma\delta} = \frac{1}{E} \left[\frac{1}{2} \omega_0 (\delta_{\alpha\gamma} \delta_{\beta\delta} + \delta_{\alpha\delta} \delta_{\beta\gamma}) + \frac{1}{8} (\omega_{\alpha\gamma} \delta_{\beta\delta} + \omega_{\alpha\delta} \delta_{\gamma\beta} + \omega_{\beta\gamma} \delta_{\alpha\delta} + \omega_{\beta\delta} \delta_{\alpha\gamma}) \right]. \quad (53)$$

Equation (53) shows that in the case of a random distribution of cracks, the first two coefficients of the series expansion (47) define the damage state variables. To get the expression of the Helmholtz free energy density, a Legendre–Fenchel transformation of the Gibbs specific enthalpy, has to be used.

These results are written for plane problems, and the same formalism can be used in 3D problems related to random distributions of penny-shaped cracks as shown in [32–34]. In that case, the first three coefficients of the series expansion of the average crack density $\langle \omega_c(\mathbf{n}) \rangle$ along a normal \mathbf{n} define the damage state variables.

Inelastic properties of damaged materials

The inelastic strains which occur in CMCs are the result of slip at interfaces. The inelastic strains resulting from the slip and the elastic energy introduced into the broken part are the subjects of this section. The unit cell in Fig. 4 shows a crack in a unidirectionally reinforced CMC accompanied by interface slip occurring over a debond length l_f . The volume fraction of the broken part is f , and the elastic moduli of the broken part and the broken part are E_u , G_u and E_b , G_b , respectively. There are two components of slip whose effects are discussed separately.

The first component of slip occurs when tensile stress is applied in the fiber direction. The faces of the crack are open by an amount Δ and cause slip over the debond length l_f , Fig. 4. When the slip is introduced and locked in position, a self-equilibrating state of stress ρ^μ is introduced into the composite and the total compatible strains consist of two components

$$\varepsilon^\mu = \varepsilon^{\mu e} + \varepsilon^{\mu in}, \quad (54)$$

where $\varepsilon^{\mu e}$ is the elastic strain tensor which satisfies the elastic stress/strain condition

$$\varepsilon^{\mu e} = \mathbf{C}^\mu : \rho^\mu, \quad (55)$$

and $\varepsilon^{\mu in}$ are the inelastic strains. The macroscopic inelastic strains may be calculated using the relationship derived in [10]

$$V \hat{\sigma} : \alpha + \int_A \hat{\mathbf{T}}^+ \cdot \mathbf{u} dS = \int_\Omega \hat{\sigma}^\mu : (\varepsilon^{\mu e} + \varepsilon^{\mu in}) dV, \quad (56)$$

where $\hat{\sigma}^\mu$ is the elastic stress in the composite corresponding to a macro stress $\hat{\sigma}$, and Ω an integration domain of volume V . The body Ω contains a plane of surface A with outward normals \mathbf{n}^+ and \mathbf{n}^- . The positive side of the plane moves an amount \mathbf{u} with respect to the negative side. The body is treated as if it contained an internal surface S that is subjected to tractions $\hat{\mathbf{T}}^+ = -\hat{\sigma} \cdot \mathbf{n}^+$ along S^+ and $\hat{\mathbf{T}}^- = -\hat{\sigma} \cdot \mathbf{n}^-$ along S^- . When $\hat{\sigma}_{kl} = \hat{\sigma}_{11} = 1$, the internal elastic stresses are

$$\hat{\sigma}_{11}^{\mu u} = \frac{E_u}{E} \hat{\sigma}_{11}, \quad (57)$$

$$\hat{\sigma}_{11}^{\mu b} = \frac{E_b}{E} \hat{\sigma}_{11}, \quad (58)$$

where $E = (1 - f)E_u + fE_b$, then

$$\alpha_{11} = \frac{\Delta f E_b}{L E}. \quad (59)$$

Similarly,

$$\alpha_{22} = \alpha_{12} = 0. \quad (60)$$

It might be noted that the same result applies irrespective of the crack angle.

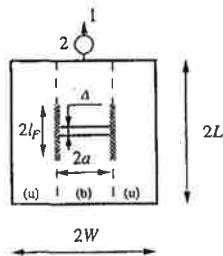


Fig. 4. Elementary cell in tension of size $2L \times 2W$ containing a crack of size $2a$. A friction zone is characterized by the friction length l_f .

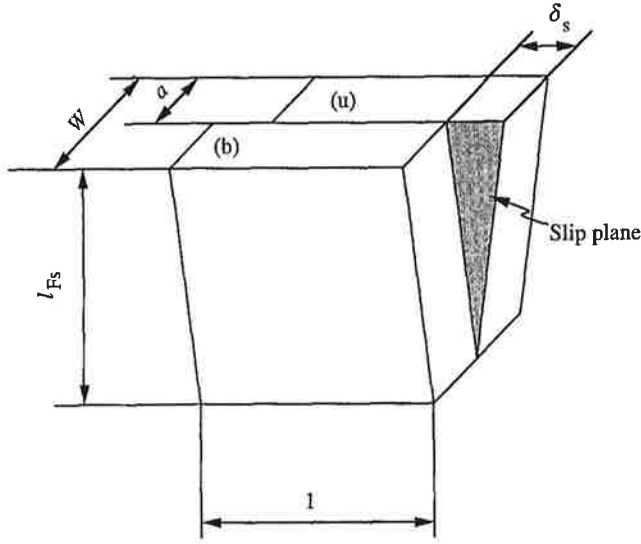


Fig. 5. Elementary cell in shear containing a crack of size $2a$. A friction zone is characterized by the friction length l_{Fs} .

The second type of slip which is the result of shear is shown in Fig. 5. Applying again the previous results, and using the shear field shown in Fig. 5, when $\hat{\sigma}_{k1} = \hat{\sigma}_{12} = \hat{\sigma}_{21} = 1$, for which the elastic stress distribution is

$$\hat{\sigma}_{12}^{uu} = \frac{G_u}{G} \hat{\sigma}_{12}, \quad \hat{\sigma}_{12}^{ub} = \frac{G_b}{G} \hat{\sigma}_{12}, \quad (61)$$

gives the result

$$\alpha_{12} = \frac{\delta_s f G_b}{L G}, \quad (62)$$

$$\alpha_{11} = \alpha_{22} = 0, \quad (63)$$

where $G = (1-f)G_u + fG_b$. With the assumption on the shear stress at the slipping interface, it is possible to show that the elastic energy resulting from the slip Δ is

$$\psi^s(\Delta) = \frac{2 f E_b (1-f) E_u}{3 E} \left(\frac{\Delta}{l_F} \right)^2 \frac{l_F}{L}, \quad (64)$$

and the corresponding value for the shear slip δ_s is

$$\psi^s(\delta) = \frac{1 f G_b (1-f) G_u}{2 G} \left(\frac{\delta}{l_{Fs}} \right)^2 \frac{l_{Fs}}{L}. \quad (65)$$

The results derived in Sect. 3 and 4 will be used in the following sections to derive a 2D constitutive law of CMCs under monotonic loading conditions.

5 CMCs with multidirectional fibers systems

This section is concerned with the behavior of ceramics reinforced by CMCs. It has been demonstrated by Evans and Marshall [35] that following matrix cracking, sliding occurs at the fiber/matrix interface which results in inelastic deformations. The presence of matrix cracks and inelastic deformations may impart to the material the ability to redistribute stresses. The results of experiments on notched panels on SiC/CAS composites [36] suggest the capacity of the material to redistribute stress is sufficiently high for this material to be notch-insensitive. The ability to redistribute stress is an important property since design studies indicate that working stresses are sufficiently high for matrix cracking to be unavoidable in regions of stress concentration.

The micromechanics which describes interface bedonding and sliding has been established by Hutchinson and Jensen [37] and Evans et al. [38]. In contrast to the early phenomenological studies [25; 39], the intention of the present study is to develop a continuum description of the damage processes which is mechanism-based, and which may be used to describe the behavior of CMCs under the conditions of multiaxial stresses occurring in practice. Since crack spacing at saturation is small in CMCs [12], CDM is an appropriate means of describing damage. Changes in elastic moduli, measured on a macroscopic level, provide a simpler and more robust means of measuring damage than does microscopic measurement of crack density. The latter requires the average of many readings before reliable values are established [40].

By combining CDM with the micromechanical studies referred to previously, constitutive equations are developed which lend themselves to the finite element procedures commonly used in practice [41; 42]. Following established procedures, the properties of each layer are first derived, and those of the composite are calculated by ensuring compatibility conditions.

The components of each layer consist of the matrix, the fiber and the interface, with f being the fiber volume fraction. The fiber direction defines the 1–2 axes. The axes x – y correspond to the principal axes of the strains in the ceramic matrix. The definition of the axes used at the constituent, layer and composite levels are shown in Fig. 3. Following Sect. 3, the loss of stiffness due to matrix cracking and fiber breakage is first established, and this is followed by the influence of the slip at the interface studies in Sect. 4.

5.1

Elastic energies of the composite associated with matrix cracking and fiber breakage

5.1.1

Constituent level: matrix and fiber

The initial behavior of the matrix is assumed to be isotropic. The presence of cracks leads to an anisotropic behavior. The assumption is made that cracking occurs in the x -direction (e.g. the maximal principal strain direction) in the matrix, and that this direction is constant during a monotonic load history. It is worth noting that this assumption excludes nonproportional loading conditions as well as equibiaxial straining. Under this hypothesis, only one damage variable for matrix cracking is needed, and is denoted by D_{mx} . The crack distribution is different from zero only along one direction. The results of Sect. 3 show that the Young's modulus along each direction as well as the shear modulus are altered [21; 43]. The expression of the elastic energy density of the matrix is (see Eq. (40))

$$\psi_m = \frac{1}{2} \frac{E_m((1 - D_{mx})\varepsilon_{mxx}^2 + 2\nu_m(1 - D_{mx})\varepsilon_{mxx}\varepsilon_{myy} + \varepsilon_{myy}^2)}{1 - \nu_m^2(1 - D_{mx})} + 2\tilde{G}_m(D_{mx})\varepsilon_{mxy}^2, \quad (66)$$

with

$$\tilde{G}_m(D_{mx}) = \frac{G_m}{1 + \frac{D_{mx}}{1 - D_{mx}} \frac{1}{2(1 + \nu_m)}}, \quad (67)$$

where E_m , ν_m are the initial elastic properties of the matrix, \tilde{G}_m is the shear modulus which is a function of damage variables D_{mx} , and ε_{mxx} , ε_{myy} , ε_{mxy} are the components of the strain tensor ε_m of the matrix expressed in the x - y frame.

It is assumed that the fibers are aligned along the 1-direction and that fiber breakage is perpendicular to the fiber direction. Therefore the elastic energy is given by

$$\psi_f = \frac{1}{2}[E_f(1 - D_{f1})\varepsilon_{f11}^2 + E_f\varepsilon_{f22}^2] + 2G_f(D_{f1})\varepsilon_{f12}^2, \quad (68)$$

where E_f is the Young's modulus of the uncracked fiber, $G_f(D_{f1})$ is the shear modulus of the cracked fiber embedded in the matrix, ε_{f11} , ε_{f22} and ε_{f12} the components of the strain tensor ε_f of the fiber expressed in the 1–2 frame.

The expression of the stresses in the matrix σ_m and in the fibers σ_f are obtained by partial differentiation of the elastic energy density with respect to the strain tensors ε_m and ε_f , respectively,

$$\sigma_m = \frac{\partial \psi_m}{\partial \varepsilon_m}, \quad (69)$$

$$\sigma_f = \frac{\partial \psi_f}{\partial \varepsilon_f}, \quad (70)$$

and the associated forces are defined as

$$Y_{mx} = - \frac{\partial \psi_m}{\partial D_{mx}}, \quad (71)$$

$$Y_{f1} = - \frac{\partial \psi_f}{\partial D_{f1}}. \quad (72)$$

These generalized forces are the energy release rate densities associated with matrix cracking and fiber breakage, respectively.

5.1.2

Layered composite

When the composite consists of layers of unidirectional fibers with different orientations, the laminate properties are determined by applying laminate theory to the properties of individual layers.

5.1.2.1

Layer level

A layer consists of fibers aligned along one orientation embedded in a matrix. To determine the behavior of this layer, micro-interface compatibility conditions are written in terms of the strains $\boldsymbol{\varepsilon}^L$ and stresses $\boldsymbol{\sigma}^L$ on the layer level. Those conditions are the compatibility and the equilibrium between a fiber and the surrounding matrix, which takes place in that system. Therefore, it is more convenient to write the conditions in the 1-2 material frame as follows:

$$\varepsilon_{m11} = \varepsilon_{f11} = \varepsilon_{11}^L, \quad (73)$$

$$f_m \sigma_{m11} + f_f \sigma_{f11} = \sigma_{11}^L, \quad (74)$$

$$f_m \varepsilon_{m22} + f_f \varepsilon_{f22} = \varepsilon_{22}^L, \quad (75)$$

$$\sigma_{m22} = \sigma_{f22} = \sigma_{22}^L, \quad (76)$$

$$f_m \varepsilon_{m12} + f_f \varepsilon_{f12} = \varepsilon_{12}^L, \quad (77)$$

$$\sigma_{m12} = \sigma_{f12} = \sigma_{12}^L, \quad (78)$$

where f_f denotes the fiber volume fraction, and f_m denotes the matrix volume fraction. When the principal strain directions do not coincide with the material frame, Eqs. (73-78) have to be rewritten in the 1-2 frame. The application of these equations then defines the elastic properties of the layer,

$$\boldsymbol{\sigma}^L = \mathbf{E}^L(D_{mx}, D_{f1}) : \boldsymbol{\varepsilon}^L \quad (79)$$

where $\mathbf{E}^L(D_{mx}, D_{f1})$ is the fourth order elastic tensor of the layer level, and is a function of all damage variables at the constituent level. From Eq. (79), the elastic energy density associated with matrix cracking and fiber breakage on the layer level can be written as

$$\psi^L = \frac{1}{2} \boldsymbol{\varepsilon}^L : \mathbf{E}^L(D_{mx}, D_{f1}) : \boldsymbol{\varepsilon}^L \quad (80)$$

5.1.2.2

Composite level

For simplicity, the case of two layers at 0 and 90 degrees is considered. The micromechanical quantities associated with the 0 degree layer are superscripted by ⁰⁰, and those at 90 degrees by ⁹⁰. The elastic behavior of the composite system is determined by applying the classical laminate theory. The compatibility condition and the overall stresses $\boldsymbol{\sigma}$ are

$$\boldsymbol{\varepsilon} = \boldsymbol{\varepsilon}^{00} = \boldsymbol{\varepsilon}^{90}, \quad (81)$$

$$\boldsymbol{\sigma} = f^{00} \boldsymbol{\sigma}^{00} + f^{90} \boldsymbol{\sigma}^{90}, \quad (82)$$

where f^{00} and f^{90} denote, respectively, the volume fraction of the first and the second layer, σ^{00} and σ^{90} denote, respectively, the stress tensor of the first and the second layer. By solving Eqs. (73–78) and using Eqs. (81–82), the overall behavior of the composite is defined as

$$\sigma = \mathbf{E}(D_{mx}^{00}, D_{f1}^{00}, D_{mx}^{90}, D_{f1}^{90}) : \varepsilon, \quad (83)$$

with

$$\mathbf{E}(D_{mx}^{00}, D_{f1}^{00}, D_{mx}^{90}, D_{f1}^{90}) = f^{00} \mathbf{E}^{00}(D_{mx}^{00}, D_{f1}^{00}) + f^{90} \mathbf{E}^{90}(D_{mx}^{90}, D_{f1}^{90}), \quad (84)$$

where $\mathbf{E}(D_{mx}^{00}, D_{f1}^{00}, D_{mx}^{90}, D_{f1}^{90})$ is the fourth order elastic tensor on the composite level, which is a function of all damage variables $D_{mx}^{00}, D_{f1}^{00}, D_{mx}^{90}, D_{f1}^{90}$ on the constituent level for the two layers. From Eqs. (81–82), the elastic energy associated with matrix cracking and fiber breakage on the composite level can be written as

$$\psi^D = f^{00} \psi^{00} + f^{90} \psi^{90}. \quad (85)$$

5.2

State potential associated with debonding and fiber Pull out

Inelasticity is essentially due to sliding at the interface between the fiber and the matrix. Sliding is involved in debonding as well as fiber pullout. From a micromechanical point of view, this sliding can take place as soon as a crack is bridged by fibers. In a CDM formulation, only the equivalent homogenous sliding and the associated forces are considered. Therefore, the cell model used to describe cracking and sliding is that shown in Figs. 4, 5. The analysis that has been done in Sect. 4 to give the expression of the elastic energy due to sliding of a layered composite along the 1–2 directions can be extended as follows:

$$\psi^S = \frac{1}{2} E \left(\frac{\alpha_{11}^2}{d_{11}} + \frac{\alpha_{22}^2}{d_{22}} \right) + \frac{1}{2} G \left(\frac{\alpha_{12}^2}{d_{12}} \right), \quad (86)$$

with

$$E = \frac{4}{3} \frac{f^{00} E^{00} f^{90} E^{90}}{f^{00} E^{00} + f^{90} E^{90}}, \quad G = \frac{f^{00} G^{00} f^{90} G^{90}}{f^{00} G^{00} + f^{90} G^{90}},$$

where E^{00} is the Young's modulus of the 0-degree layer in the fiber direction 00, E^{90} is the Young's modulus of the 90 degree layer in the fiber direction 90, G^{00} is the shear modulus of the 0-degree layer, G^{90} is the shear modulus of the 90-degree layer, d_{11} , d_{22} and d_{12} are damage quantities related to sliding, and α_{11} , α_{22} and α_{12} are the inelastic strains.

5.2.1

State laws

The following development deals with ceramic matrix layered composites. The total elastic energy density of the composite is the sum of the elastic energy density of the damaged composite ψ^D and the elastic energy density due to sliding ψ^S

$$\psi = \frac{1}{2} (\varepsilon - \alpha) : \mathbf{E}(D_{mx}^{00}, D_{mx}^{90}, D_{f1}^{00}, D_{f1}^{90}) : (\varepsilon - \alpha) + \psi^S, \quad (87)$$

where D_m^{00} , D_m^{90} , D_{f1}^{00} , and D_{f1}^{90} are the damage variables modeling matrix cracking and fiber breakage in the 0 and 90 degree plies. The forces associated with total strains are

$$\sigma = \frac{\partial \psi}{\partial \varepsilon} = \mathbf{E}(D_{mx}^{00}, D_{mx}^{90}, D_{f1}^{00}, D_{f1}^{90}) : (\varepsilon - \alpha), \quad (88)$$

and correspond to the macroscopic stresses. The associated forces to the damage variables modeling matrix cracking and fiber breakage are

$$Y = - \frac{\partial \psi}{\partial D}, \quad (89)$$

where $D = \{D_{mx}^{00}; D_{mx}^{90}; D_{f1}^{00}; D_{f1}^{90}\}$, and $Y = \{Y_{mx}^{00}; Y_{mx}^{90}; Y_{f1}^{00}; Y_{f1}^{90}\}$ represent the energy release rate densities due to matrix cracking and fiber breakage. The associated forces to the damage variables modeling sliding are

$$y = - \frac{\partial \psi}{\partial \mathbf{d}}, \quad (90)$$

where $\mathbf{d} = \{d_{11}; d_{22}; d_{12}\}$ and y correspond to the energy release rate densities due to sliding. The associated forces to the anelastic strains are

$$X = \frac{\partial \psi}{\partial \alpha}, \quad (91)$$

and represent the back stresses in the sliding zone.

5.2.2

Evolution laws

5.2.2.1

1D evolution law of a CMC

The aim of this section is to derive the evolution laws from a 1D analysis. Postmortem analysis of broken specimens indicates the presence of arrays of microcracks in the matrix accompanied by debonding or friction at the fiber/matrix interface. Hutchinson and Jensen [37] and Evans et al. [38] have analyzed the behavior of a unidirectional CMC in tension by considering the unit cell shown in Fig. 4, when matrix cracking of spacing $2L$ is accompanied by sliding at the interface over a debond length $2l_f$. The micromechanics can predict the 1D macroscopic stress-strain response shown in Fig. 6. The microcracks are usually aligned with the principal stress or strain directions. By studying a cracked panel with crack of length $2a$ in a cell of area $4LW$ (Fig. 2), the reduction in stiffness may be estimated. If the initial behavior of the elementary cell is isotropic, and the Young's modulus is E , as shown earlier, the stiffness loss depends on the crack density defined as $\pi a^2/4LW$. By assuming plane stress conditions, and that the crack density is so small that crack interaction can be neglected, a first approximation for the reduced elastic modulus \tilde{E} can be written as (see Sect. 2)

$$\frac{\tilde{E}}{E} = \frac{1}{1 + 2 \frac{\pi a^2}{4LW}} = 1 - D. \quad (92)$$

In the case of constituents with different elastic properties, D depends upon the elastic properties of the two constituents, as well as the ratios a/W and a/L . The stress/strain relationship becomes

$$\sigma = \frac{\bar{\sigma}}{1 - D} = E \bar{\epsilon}, \quad (93)$$

where σ is the stress on a microscopic level, $\bar{\sigma}$ and $\bar{\epsilon}$ are the macroscopic stress and strain, respectively. The model illustrated in Fig. 4 is now analyzed and uses an analytic approach which follows the thermodynamic developments summarized in Sect. 2. This is done by calculating the internal elastic density ψ^s in the unit cell [44] caused by debonding and sliding at the interface. The first

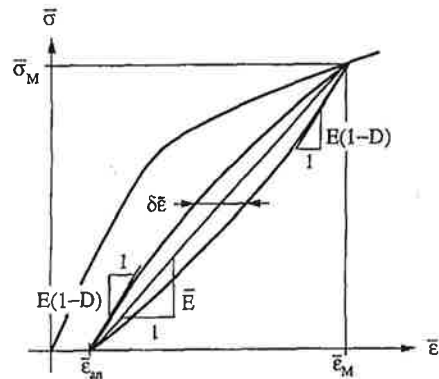


Fig. 6. Stress, $\bar{\sigma}$, versus strain, $\bar{\epsilon}$, during a loading-unloading-reloading sequence

step consists in moving the unbroken part (u) with respect to the broken part (b) with no external load by an amount Δ over a length l_F (Fig. 7). The displacement Δ gives rise to a self-balanced linear stress field along a length l_F in parts (b) and (u). By integration over l_F , and then by averaging over the total length L , the elastic energy density associated with this process is given by

$$\psi^s = \frac{2}{3} \frac{fE_b(1-f)E_u}{E} \left(\frac{\Delta}{l_F} \right)^2 \frac{l_F}{L}. \quad (94)$$

As shown in Eq. (59), the crack opening displacement Δ induces an irreversible strain α expressed as

$$\alpha = \frac{fE_b \Delta}{E L}. \quad (95)$$

The second step consists of an elastic loading of the damaged system so that the elastic energy density is given by

$$\psi^e = \frac{1}{2} E (1 - D) (\bar{\epsilon} - \alpha)^2. \quad (96)$$

The total elastic energy density is the sum of the two elements of the energy densities so that

$$\psi = \frac{1}{2} E (1 - D) (\bar{\epsilon} - \alpha)^2 + \frac{2}{3} \frac{fE_b(1-f)E_u}{E} \left(\frac{\Delta}{l_F} \right)^2 \frac{l_F}{L}. \quad (97)$$

For convenience, the energy density can be expressed in a more compact form by using state variables which are the total strain $\bar{\epsilon}$, the damage variable D modeling the loss of stiffness due to the cracking mechanism, the damage variable $d = fE_b l_F / (1-f)E_u L$ which defines the size of the slip zone related to the crack spacing, and the crack opening strain α . The elastic energy density in terms of the new internal variables is

$$\psi = \frac{1}{2} E (1 - D) (\bar{\epsilon} - \alpha)^2 + \frac{2}{3} E \left(\frac{\alpha^2}{d} \right). \quad (98)$$

The forces associated with the state variables ($\bar{\epsilon}, D, d, \alpha$) are respectively given by

$$\bar{\sigma} = \frac{\partial \psi}{\partial \bar{\epsilon}} = E (1 - D) (\bar{\epsilon} - \alpha), \quad (99)$$

$$Y = - \frac{\partial \psi}{\partial D} = \frac{E}{2} (\bar{\epsilon} - \alpha)^2, \quad (100)$$

$$y = - \frac{\partial \psi}{\partial d} = \frac{2E}{3} \left(\frac{\alpha}{d} \right)^2, \quad (101)$$

$$X = \frac{\partial \psi}{\partial \alpha} = - \bar{\sigma} + \frac{4E}{3} \frac{\alpha}{d}, \quad (102)$$

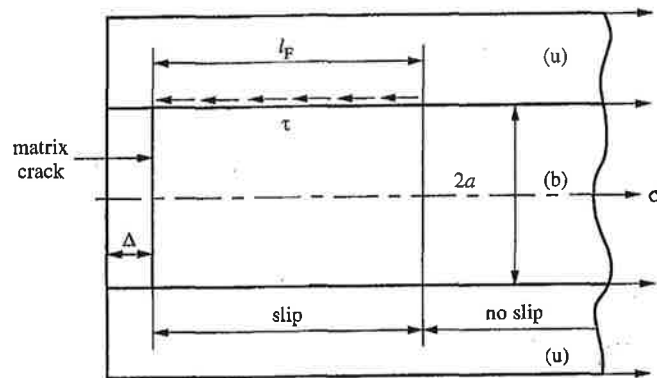


Fig. 7. Motion of the unbroken part (u) with respect to the broken part (b) with no external load by an amount Δ over a length l_F

where $\bar{\sigma}$ is the applied stress, Y is the energy released rate available to form matrix cracks, y the energy release rate available for interface debonding and X is the back-stress induced by the slipping mechanism.

In the present approach the growth laws of the state variables (D, d, α) are established from macroscopic quantities measured in the course of unloading and reloading experiments. To this end, use is made of the solution [44] of the response of the unit cell (Fig. 4) when subjected to an unloading/reloading cycle during which the magnitude of the shear stress remains constant.

The expressions obtained from the analysis for the residual stress ρ_b and the internal variables D, d, α in terms of the macroscopic quantities shown in Fig. 6 are given respectively by

$$\frac{-\rho_b}{E} = \left(\sqrt{\frac{\bar{\epsilon}_{in} + 2\delta\bar{\epsilon}}{4\delta\bar{\epsilon}}} - 1 \right) (\bar{\epsilon}_M - \bar{\epsilon}_{in} - 2\delta\bar{\epsilon}), \quad (103)$$

$$D = \frac{\bar{\epsilon}_M \bar{D} - \bar{\epsilon}_{in} \bar{D} - 2\delta\bar{\epsilon}}{\bar{\epsilon}_M - \bar{\epsilon}_{in} - 2\delta\bar{\epsilon}}, \quad (104)$$

$$\frac{d}{4} = \frac{\sqrt{(\bar{\epsilon}_{in} + 2\delta\bar{\epsilon}) \delta\bar{\epsilon}}}{\bar{\epsilon}_M - \bar{\epsilon}_{in} - 2\delta\bar{\epsilon}}, \quad (105)$$

$$\alpha = \frac{d \bar{\sigma}_M - \rho_b (1 - D)}{2 E (1 - D)}, \quad (106)$$

where $-\rho_b E/E_b$ is the thermal mismatch residual stress in the broken layer, $\bar{\epsilon}_{in}$ is the inelastic strain upon complete unloading, $\delta\bar{\epsilon}$ is the maximum hysteresis loop width, E is the Young's modulus of the composite, $\bar{\epsilon}_M$ ($\bar{\sigma}_M$) is the maximum applied strain (stress), D and \bar{D} are, respectively, the microscopic and macroscopic damage associated with matrix cracking, d is the damage related to interface debonding and α is the inelastic strain at the current state point. Eq. (106) is only valid for monotonic loading conditions.

By performing a series of unloading/reloading cycles, the internal variables can be determined from experiment using Eqs. (103–106). The residual stress ρ_b is calculated from Eq. (103), and it is a test of the effectiveness of the model that the same value of the residual stress is obtained for each loading sequence. The values of D and d are given by applying Eqs. (104) and (105), respectively. The information is now available to complete the calculation for Eq. (106). The corresponding associated forces are given by Eqs. (99–102). The relationship between the internal variables and the associated quantities can then be investigated. It is this method which is proposed to model the behavior of the CMC laminate [44].

5.2.2.2

Identification of a 2D evolution law of a CMC

The identification procedure is performed on a [0/90]_s laminate architecture of CMCs. The first step is to define all the internal state variables needed to model the material behavior. The three strain variables, $\{\epsilon_{11}, \epsilon_{22}, \epsilon_{12}\}$, are given either from experiment or as input from a FE calculation. The four damage variables, i.e. $\{D_{mx}^{00}, D_{mx}^{90}, D_{f1}^{00}, D_{f1}^{90}\}$ are used to define the change in the elastic properties, where D_m is matrix-cracking damage, and fiber breakage damage is D_f . Assuming the damage evolution laws, $D_m(Y_m)$ and $D_f(Y_f)$, are functions of the associated forces Y_m and Y_f , respectively, then only two evolution laws, one for each mechanism, are sufficient to compute the four components of damage. The three damage variables, $\{d_{11}; d_{22}; d_{12}\}$, define the sliding distances, with d_{11} or d_{22} being associated with sliding in the fiber directions, and d_{12} associated with shear sliding. Consequently, only two evolution laws are needed, $d_{11}(y)$ or $d_{22}(y)$, and $d_{12}(y)$.

For the same reasons, two evolution laws for the inelastic strains, $\alpha_{11}(X)$ or $\alpha_{22}(X)$, and $\alpha_{12}(X)$ give the evolution of the three inelastic strains, $\{\alpha_{11}, \alpha_{22}, \alpha_{12}\}$. In conclusion, the model has 13 state variables, three of which are strain inputs and the remaining ten micromechanical state variables are derived from six evolution laws.

The second step is to define the relevant tests required to identify the growth of the six state variables. This is achieved from unloading tests performed at regular intervals and measuring the macroscopic inelastic strain upon complete unloading, $\bar{\epsilon}_{in}$, the macroscopic damage, of the composite, \bar{D} , and the maximum hysteresis loop width, $\delta\bar{\epsilon}$ (Fig. 6). Using Eqs. (88–91), the internal states can be calculated.

When tension is applied at 45 degrees on a [0/90]_s layered composite, only matrix cracking occurs. The macroscopic damage, \bar{D}^{45} , is related to the microscopic damage variables taking place in the matrix alone, $\{D_{mx}^{00}; D_{mx}^{90}\}$. Moreover, the loading condition is such that the two damage variables, D_{mx}^{00} and D_{mx}^{90} , have the same value. Therefore, the evolution law, $D_m(Y_m)$, is directly given by the evolution of the macroscopic damage of the composite \bar{D}^{45} .

The evolution law of the damage variable associated with fiber breakage, $D_f(Y_f)$, is found from measurement of the macroscopic damage, \bar{D}^{00} , in a tension test at 0 degrees on the [0/90]_s composite. This has been compensated by the contribution of matrix cracking which may be calculated from the results of the first test. Similarly, the evolution laws of the state variables related to sliding, $\alpha_{11}(X)$ or $\alpha_{22}(X)$, and $d_{11}(y)$ or $d_{22}(y)$, are known from the evolution of the macroscopic inelastic strain upon complete unloading, $\bar{\epsilon}_{in}^{00}$, the macroscopic damage of the composite, \bar{D}^{00} , and the maximum hysteresis loop width, $\delta\bar{\epsilon}^{00}$, using micromechanics relations (103–106)

$$\frac{-\rho_b^{00}}{E^{00}} = \left(\sqrt{\frac{\bar{\epsilon}_{in}^{00} + 2\delta\bar{\epsilon}^{00}}{4\delta\bar{\epsilon}^{00}}} - 1 \right) (\bar{\epsilon}_m^{00} - \bar{\epsilon}_{in}^{00} - 2\delta\bar{\epsilon}^{00}), \quad (107)$$

$$D^{00} = \frac{\bar{\epsilon}_M^{00}\bar{D}^{00} - \bar{\epsilon}_{in}^{00}\bar{D}^{00} - 2\delta\bar{\epsilon}^{00}}{\bar{\epsilon}_M^{00} - \bar{\epsilon}_{in}^{00} - 2\delta\bar{\epsilon}^{00}}, \quad (108)$$

$$\frac{d_{11}}{4} = \frac{\sqrt{(\bar{\epsilon}_{in}^{00} + 2\delta\bar{\epsilon}^{00})\delta\bar{\epsilon}^{00}}}{\bar{\epsilon}_M^{00} - \bar{\epsilon}_{in}^{00} - 2\delta\bar{\epsilon}^{00}}, \quad (109)$$

$$\alpha_{11} = \frac{d_{11}\bar{\sigma}_M^{00} - \rho_b^{00}(1 - D^{00})}{2 \frac{E^{00}(1 - D^{00})}{2}}, \quad (110)$$

where ρ_b^{00} is proportional to the residual stress in the broken part of the 0 degree layer. Finally, returning to the results of the tension test at 45 degrees on a [0/90]_s layered composite, the evolution laws of the state variables related to sliding, $\alpha_{12}(X)$ and $d_{12}(y)$ are given by the following relationships similar to those calculated by the micromechanical analysis:

$$\frac{d_{12}}{4} = \frac{\bar{\epsilon}_{in}^{45}}{\bar{\epsilon}_M^{45} - 2\bar{\epsilon}_{in}^{45}}, \quad (111)$$

$$\alpha_{12} = \frac{d_{12} \frac{\bar{\epsilon}_M^{45}}{2}}{1 + \frac{d_{12}}{2}}. \quad (112)$$

In this analysis, it can be noticed that the residual stress $-\rho_b^{45}$ in the tensile direction is equal to zero, and, therefore, $\bar{\epsilon}_{in}^{45} = 2\delta\bar{\epsilon}^{45}$, which leads to the above results.

In summary, only two tests, on the same architecture, enable us to extract all six evolution laws that define the behavior of the material as shown in [45] on experiments performed on SiC/SiC specimens [46].

6

Conclusions

CDM has grown out of the need to predict the behavior of engineering components composed of materials whose properties degrade because of damage. The traditional approach, which is now well established as a design procedure, has been formulated with only minimal attention to the internal mechanisms occurring in the material. In the above presentation, the formulation of classical CDM has been extended so that the description of the physical mechanisms which cause damage can be included. This approach has the advantage that the studies of material scientists are integrated into a formulation retaining the proven advantages of CDM. The illustrative examples of the application of the procedure is concerned with the properties of CMC which can suffer matrix cracking and fiber/matrix sliding resulting in a loss of stiffness, the appearance of inelastic deformations, as well as hysteresis loops.

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