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# Oscillatory dynamics in Smoluchowski's coagulation equation with diagonal kernel

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October 27, 2017

## Abstract

We characterize the long-time behaviour of solutions to Smoluchowski's coagulation equation with a diagonal kernel of homogeneity  $\gamma < 1$ . Due to the property of the diagonal kernel, the value of a solution at a given cluster size depends only on a discrete set of points. As a consequence, the long-time behaviour of solutions is in general periodic, oscillating between different rescaled versions of a self-similar solution. Immediate consequences of our result are a characterization of the set of data for which the solution converges to self-similar form and a uniqueness result for self-similar profiles.

## 1 Introduction

Smoluchowski's classical coagulation equation provides a mean-field description of binary coalescence of clusters. If  $\xi$  denotes the size of a cluster and  $F(\tau, \xi)$  the corresponding number density at time  $\tau$  then the equation is given by

$$\begin{aligned} \partial_\tau F(\tau, \xi) = & \frac{1}{2} \int_0^\xi K(\xi - \eta, \eta) F(\tau, \xi - \eta) F(\tau, \eta) d\eta \\ & - \int_0^\infty K(\xi, \eta) F(\tau, \xi) F(\tau, \eta) d\eta, \quad (\tau, \xi) \in (0, \infty)^2, \end{aligned} \tag{1}$$

where the coagulation kernel  $K = K(\xi, \eta) \geq 0$  describes the rate at which clusters of size  $\xi$  and  $\eta$  coagulate. This model has been used in various applications, most prominently in aerosol physics or polymerization, but also in astrophysics or population dynamics [1, 5, 22].

We are in the following interested in kernels that grow at most sublinearly at infinity. In this case it is well-known for a large class of kernels that, if the initial condition  $F(0)$  has finite mass, i.e. finite first moment, then the solution  $F$  to (1) conserves the mass for all positive times, that is

$$\int_0^\infty \xi F(\tau, \xi) d\xi = \int_0^\infty \xi F(0, \xi) d\xi \quad \text{for all } \tau > 0, \tag{2}$$

see the survey papers [12, 13, 16] and the references therein. Since two clusters merge into one during a coalescence event, the total number of clusters, i.e. the  $L^1$ -norm of  $F(\tau)$ , is expected to decay to zero as  $\tau \rightarrow \infty$ . This property, together with (2), entails that the mean size  $s(\tau)$  at time  $\tau$  of the distribution of clusters increases without bounds with time but does not reveal much about the dynamics of the coagulation equation (1) and its possible dependence upon the coagulation kernel  $K$ . To gain further insight into these matters, the dynamical scaling hypothesis [9, 16, 23] predicts that, for large times, the solution  $F(\tau)$  to (1) behaves in a self-similar way, that is,

$$F(\tau, \xi) \sim \frac{1}{s(\tau)^2} \varphi \left( \frac{x}{s(\tau)} \right) \quad \text{as } \tau \rightarrow \infty, \quad (3)$$

for some scaling function  $\varphi$  to be determined, the prefactor of  $\varphi$  being chosen here to comply with the mass conservation (2). Though this issue has been discussed in the physics literature for homogeneous coagulation kernels with homogeneity  $\gamma \leq 1$ , that is,

$$K(\lambda\xi, \lambda\eta) = \lambda^\gamma K(\xi, \eta), \quad (\lambda, \xi, \eta) \in (0, \infty)^3, \quad (4)$$

see [9, 16, 23] and the references therein, the validity of (3) is still pending, except for the so-called solvable kernels  $K(\xi, \eta) = 2$  and  $K(\xi, \eta) = \xi + \eta$ . In these two cases, equation (1) can be solved explicitly by transform methods and a complete characterization of the long-time behaviour of solutions to (1) with solvable kernels can be found in [18]. Let us emphasize that the outcome of the analysis performed in [18] is not only that the dynamical scaling hypothesis (3) is valid for a large class of initial data but also that other self-similar behaviours are possible and involve different time and size scales, as well as different scaling profiles. For all other kernels there are so far only numerical studies available [7, 9, 11, 14] that suggest for a range of kernels convergence to self-similar form in the long-time limit. In subsequent years several results established also rigorously the existence of self-similar solutions, both with finite mass and with fat tails [6, 8, 17, 19, 20, 21], but an analysis of the long-time behaviour of solutions is still elusive. It is worth pointing out that the very first existence results of self-similar solutions for a non-solvable coagulation kernel are obtained for the so-called *diagonal kernel*

$$K(\xi, \eta) = \xi^{1+\gamma} \delta(\xi - \eta), \quad (\xi, \eta) \in (0, \infty)^2, \quad \gamma < 1, \quad (5)$$

where  $\delta$  is the Dirac mass at zero [17, 20]. Observe that the homogeneity of  $K$  is equal to  $\gamma$  as the Dirac mass is homogeneous of order  $-1$  (it is the derivative in the sense of distributions of the Heaviside function which is clearly homogeneous of order zero). In this model, only clusters of the same size can coagulate and it is fairly clear that such a model has little physical realization. Nevertheless, it can be viewed as a limit case of a coagulation kernel for which equal size interactions dominate. It has also the interesting feature of reducing the nonlocal nonlinear integral equation (1) to a nonlinear differential equation with delay which is more amenable to analysis, see (6) below, and it is commonly believed that a better understanding of its behaviour paves the way to deal with more general kernels. In the same vein, we point out that the diagonal kernel is also used

in [3, 15] to elucidate the onset of gelation, i.e. infringement of mass conservation, and the results obtained therein were for almost twenty years among the few rigorous ones available on this matter.

We thus focus in this paper on Smoluchowski's coagulation equation (1) with diagonal kernel (5) and we present the first rigorous analysis of the long-time behaviour of solutions for a non-solvable kernel. As already mentioned, it is in some sense orthogonal to the constant one since only clusters with equal sizes interact. When  $K$  is given by (5), equation (1) reduces to

$$\partial_\tau F(\tau, \xi) = \frac{\xi^{1+\gamma}}{2^{3+\gamma}} F^2(\tau, \xi/2) - \xi^{1+\gamma} F^2(\tau, \xi), \quad (\tau, \xi) \in (0, \infty)^2. \quad (6)$$

Obviously a key difference to continuous kernels is the fact that, for the diagonal one, the evolution of the number density  $F$  of clusters with mass  $\xi$  depends only on smaller values and not on larger ones. As a consequence, the equation obeys a maximum principle, a fact that we will also use in our analysis. Another key feature is that the evolution in  $\xi$  depends only on a discrete set of values of the form  $\xi/2^k$  and thus the evolution decouples in a certain sense which leads to a long-time behaviour of solutions that is different from what one expects for continuous kernels. We will explain this in more detail in Section 2.2 below, but point out already here, that our results are, to our knowledge, the first ones showing oscillatory behaviour for the solutions with finite mass. Such oscillatory behaviours can also be seen for kernels which are not diagonal, but are concentrated near the diagonal [10]. These results indicate that the self-similar behaviour is not the only possibility for the coagulation equation, although this has been often assumed and conjectured.

## 2 Main results

### 2.1 Reformulation of the problem

It turns out to be more convenient to go over to the function

$$G(\tau, \eta) = 2^{\eta(1+\gamma)} F(\tau, \xi) \quad \text{with } \xi = 2^\eta, \eta \in \mathbb{R}, \text{ and } \tau > 0.$$

Then  $G$  satisfies

$$\partial_\tau G(\tau, \eta) = \frac{1}{2^{1-\gamma}} G^2(\tau, \eta-1) - G^2(\tau, \eta), \quad (\tau, \eta) \in (0, \infty) \times \mathbb{R}, \quad (7)$$

and

$$\int_{\mathbb{R}} 2^{(1-\gamma)\eta} G(\tau, \eta) d\eta = \int_{\mathbb{R}} 2^{(1-\gamma)\eta} G(0, \eta) d\eta \quad \text{for all } \tau > 0. \quad (8)$$

Well-posedness of (7) for initial data in  $L^\infty(\mathbb{R})$  will be shown later for a reformulated problem in Lemma 3.2.

We expect to have as special solutions to (6) self-similar solutions of the form

$$(\tau, \xi) \mapsto \frac{1}{\tau^{2/(1-\gamma)}} \phi\left(\frac{\xi}{\tau^{1/(1-\gamma)}}\right), \quad (\tau, \xi) \in (0, \infty)^2,$$

which means for (7) rescaled traveling wave solutions of the form

$$G(\tau, \eta) = \frac{1}{\tau} H \left( \eta - \frac{\ln \tau}{(1-\gamma) \ln 2} \right), \quad (\tau, \eta) \in (0, \infty) \times \mathbb{R}.$$

In order to study the existence of such special solutions and the large time behaviour of solutions to (7) we introduce the new variables

$$\alpha := (1-\gamma) \ln 2, \quad t := \frac{\ln \tau}{\alpha}, \quad x := \eta - \frac{\ln \tau}{\alpha}, \quad \text{and} \quad e^{-\alpha t} h(t, x) := \alpha G(\tau, \eta), \quad (9)$$

and we deduce from (7) and (8) that  $h$  solves

$$\partial_t h(t, x) = \partial_x h(t, x) + \alpha h(t, x) + e^{-\alpha} h^2(t, x-1) - h^2(t, x), \quad (t, x) \in \mathbb{R}^2, \quad (10)$$

and

$$\int_{\mathbb{R}} e^{\alpha x} h(t, x) dx = \int_{\mathbb{R}} e^{\alpha x} h(0, x) dx, \quad t \in \mathbb{R}. \quad (11)$$

A self-similar solution to (6) corresponds now to a stationary solution  $\bar{h}$  of (10), that is, to a solution of

$$0 = \partial_x \bar{h}(x) + \alpha \bar{h}(x) + e^{-\alpha} \bar{h}^2(x-1) - \bar{h}^2(x), \quad x \in \mathbb{R}. \quad (12)$$

Notice that the equation (12) is invariant under a shift, that is, if  $\bar{h}$  is a solution to (12), then so is  $x \mapsto \bar{h}(x + \lambda)$  for all  $\lambda \in \mathbb{R}$ . Recall that we are interested in solutions  $\bar{h}$  to (12) that are nonnegative, locally integrable, and such that  $\bar{h} \in L^1(\mathbb{R}; e^{\alpha x} dx)$ . By shifting  $\bar{h}$  we can normalize the mass so that

$$\int_{\mathbb{R}} e^{\alpha x} \bar{h}(x) dx = 1.$$

Notice also that it follows directly from (12) and the local integrability of  $\bar{h}$  that  $\bar{h} \in C^1(\mathbb{R})$ . It is proved in [17] by a shooting method that for any  $a > 0$  there exists a nonnegative solution  $\bar{h}_a \in L^1(\mathbb{R}; e^{\alpha x} dx)$  to (12) that is bounded, decreasing and satisfies

$$\lim_{x \rightarrow -\infty} 2^{-\sigma x} \left( \bar{h}_a(x) - \frac{\alpha}{1-e^{-\alpha}} \right) = -a, \quad (13)$$

where  $\sigma$  is the only positive root of the equation  $(1 + \ln 2\sigma/\alpha)(1 - e^{-\alpha}) = 2(1 - e^{-\alpha} 2^{-\sigma})$ . The constant  $a$  is just a normalization and corresponds to rescaling the mass  $\int_{\mathbb{R}} e^{\alpha x} \bar{h}_a(x) dx$ . Furthermore, the function  $\bar{h}_a$  satisfies

$$0 < \bar{h}_a(x) \leq C_a \exp \left( -L_a 2^x \right) \quad \text{for some } C_a > 0 \text{ and } L_a > 0. \quad (14)$$

Our goal in this paper is to study the behaviour of solutions to (10) as  $t \rightarrow \infty$  and to figure out what is the role played by the family  $(\bar{h}_a)_{a>0}$ , if any. As a consequence of our results we will actually show that stationary solutions are unique up to rescaling.

## 2.2 Main results

A key property of the diagonal kernel is that the evolution of  $h(t, x)$  depends only on the discrete set of values  $x - k$ ,  $k \in \mathbb{N}$ . Correspondingly, one of the key points of the argument used in this paper is a decomposition of the plane  $\{(t, x) \in \mathbb{R}^2\}$  in a family of lines whose evolution under (10) is mutually decoupled.

More specifically, given  $\theta \in [0, 1)$ , we define the following family of lines:

$$\mathcal{S}_\theta := \{(t, x) \in \mathbb{R}^2 : x + t - \theta \in \mathbb{Z}\} .$$

Obviously for given  $\theta_1, \theta_2 \in [0, 1)$  with  $\theta_1 \neq \theta_2$  we have  $\mathcal{S}_{\theta_1} \cap \mathcal{S}_{\theta_2} = \emptyset$ , while  $\bigcup_{\theta \in [0, 1)} \mathcal{S}_\theta = \mathbb{R}^2$ . Hence, the following function  $\Theta$  is well-defined.

**Definition 2.1.** *We define the function  $\Theta: \mathbb{R} \times \mathbb{R} \rightarrow [0, 1)$  as the unique  $\theta = \Theta(t, x)$  such that  $(t, x) \in \mathcal{S}_{\Theta(t, x)}$ , that is,  $\Theta(t, x) = t + x - \lfloor t + x \rfloor$ .*

Notice that the function  $\Theta$  is 1-periodic in both variables, that is

$$\Theta(t + 1, x) = \Theta(t, x) , \quad \Theta(t, x) = \Theta(t, x - 1) , \quad (t, x) \in \mathbb{R} \times \mathbb{R} . \quad (15)$$

We also define the function

$$\psi(t, \theta) := t - \theta - \lfloor t - \theta \rfloor \in [0, 1) , \quad (t, \theta) \in [0, \infty) \times [0, 1) . \quad (16)$$

Notice that in particular  $\psi(0, \theta) = 1 - \theta$  for  $\theta \in (0, 1)$  and that  $t \mapsto \psi(t, \theta)$  is right-continuous and jumps from 1 to 0 at times  $n + \theta$ ,  $n \in \mathbb{N}$ ,

$$\psi((n + \theta)^-, \theta) := \lim_{t \nearrow n + \theta} \psi(t, \theta) = 1 \quad \text{and} \quad \psi(n + \theta, \theta) = 0 . \quad (17)$$

We now fix a nonnegative solution  $\bar{h} \in C^1(\mathbb{R}) \cap L^1(\mathbb{R}; e^{\alpha x} dx)$  to (12) enjoying the properties (13) and (14) and with mass normalized to one, that is

$$1 = \int_{\mathbb{R}} e^{\alpha x} \bar{h}(x) dx = \int_0^1 \left\{ \sum_{k \in \mathbb{Z}} e^{\alpha(k + \theta)} \bar{h}(k + \theta) \right\} d\theta . \quad (18)$$

Introducing

$$\nu(\theta) := \sum_{k \in \mathbb{Z}} e^{\alpha(k + \theta)} \bar{h}(k + \theta) , \quad \theta \in [0, 1) ,$$

we infer from (12) and (14) that

$$0 = \sum_{k \in \mathbb{Z}} [\partial_x \bar{h}(k + \theta) + \alpha \bar{h}(k + \theta)] e^{\alpha(k + \theta)} = \sum_{k \in \mathbb{Z}} \frac{d}{d\theta} \left[ e^{\alpha(k + \theta)} \bar{h}(k + \theta) \right] = \frac{d\nu}{d\theta}(\theta) .$$

Consequently,  $\nu(\theta) = \text{const.}$  for  $\theta \in [0, 1)$  and the normalization (18) entails that

$$\sum_{k \in \mathbb{Z}} e^{\alpha(k + \theta)} \bar{h}(k + \theta) = 1 , \quad \theta \in [0, 1) . \quad (19)$$

Then, in order to obtain from  $\bar{h}$  a stationary solution to (10) with mass  $M > 0$ , we need to shift  $\bar{h}$  and see that the shift  $\lambda$  is determined by the following relation:

$$M = \sum_{k \in \mathbb{Z}} e^{\alpha k} \bar{h}(k - \lambda) \quad \text{if and only if} \quad \lambda = \frac{\ln M}{\alpha}. \quad (20)$$

As pointed out above, the evolution of (10) decouples into an evolution for each ‘fibre’  $\mathcal{S}_\theta$ ,  $\theta \in [0, 1)$ . We shall see that, given a nonnegative initial condition  $h_0 \in L^\infty(\mathbb{R})$ , the long-time behaviour of the corresponding solution to (10) is determined by the mass of the initial condition in each fibre, given by

$$m_0(\theta) := \sum_{k \in \mathbb{Z}} e^{\alpha(k+\theta)} h_0(k + \theta), \quad \theta \in [0, 1). \quad (21)$$

More precisely, we will see that a solution of (10) is for large times approximated in each fibre  $\theta \in [0, 1)$  by the stationary solution with the same mass, that is, according to (20), the shift  $\bar{h}(\cdot - \lambda(\theta))$  of  $\bar{h}$  such that

$$m_0(\theta) = \sum_{k \in \mathbb{Z}} e^{\alpha k} \bar{h}(k - \lambda(\theta)), \quad \lambda(\theta) := \frac{\ln m_0(\theta)}{\alpha}. \quad (22)$$

We can now formulate our main results.

**Theorem 2.2.** *Consider the solution  $h$  of (10) with a nonnegative initial condition  $h_0 \in L^\infty(\mathbb{R})$  satisfying  $m_0 \in L^1(0, 1)$ , where  $m_0$  is defined in (21). Introducing*

$$\mu(t, x) := \frac{\ln(m_0(\Theta(t, x)))}{\alpha}, \quad (t, x) \in (0, \infty) \times \mathbb{R},$$

there holds

$$\lim_{t \rightarrow \infty} \int_{\mathbb{R}} e^{\alpha x} |h(t, x) - \bar{h}(x - \mu(t, x))| dx = 0.$$

Theorem 2.2 provides a detailed description of the asymptotic behaviour of solutions to (10) as  $t \rightarrow \infty$  and in particular implies that the long-time behaviour is in general periodic. We note that periodic long-time behaviour is also observed in a growth-fragmentation model where the fragmentation kernel is diagonal [2]. Theorem 2.2 also characterizes the class of initial data with finite mass that yield a traveling wave solution (resp. self-similar solution in the original variables) as  $t \rightarrow \infty$ .

**Corollary 2.3.** *We have convergence to a stationary solution if and only if the function  $m_0$  as defined in (21) is constant almost everywhere in  $[0, 1)$ .*

Corollary 2.3 applies in particular to functions  $h_0$  given by  $h_0(x) := H_k e^{-\alpha(x-k)}$  for  $x \in [k, k+1)$  and  $k \in \mathbb{Z}$ , where  $(H_k)_{k \in \mathbb{Z}}$  is a sequence of positive real numbers which is bounded and such that  $\sum H_k e^{\alpha k}$  converges. For instance,  $H_k = e^{-\beta k_+}$ ,  $\beta > \alpha$ , or  $H_k = (k_+)^{-q} e^{-\alpha k_+}$ ,  $q > 1$ , recalling that  $k_+ := \max\{k, 0\}$ . Let us nevertheless emphasize here

that, according to Theorem 2.2, solutions emanating from arbitrarily small perturbations of such initial data may have a drastically different behaviour in the long term, as soon as the corresponding function  $m_0$  in (21) is not constant.

Another consequence of Theorem 2.2 is the fact that stationary solutions of (10) are unique up to rescaling.

**Proposition 2.4.** *Let  $\bar{h}_1$  and  $\bar{h}_2$  be two nonnegative solutions to (12) in  $L^1(\mathbb{R}; e^{\alpha x} dx)$ . Then there exists  $\Lambda \in \mathbb{R}$  such that  $\bar{h}_1(x) = \bar{h}_2(x - \Lambda)$  for  $x \in \mathbb{R}$ .*

The remainder of the paper is organized as follows. Section 3, which is the core of our paper, contains the proof of Theorem 2.2 and consists of several lemmas. The key idea is that we can construct a Lyapunov functional  $L^\theta$  for the evolution in each fibre that resembles a contractivity property of solutions to scalar conservation laws and nonlinear diffusion equations [4, 24], see Lemma 3.4. We also identify its limit for large times in Lemma 3.6. Another key auxiliary result is a tightness property that we prove in Lemma 3.3. This in turn allows us to obtain a lower bound for the dissipation functional in Lemma 3.5. All these results are proved pointwise in the sense that we prove them for each fibre. Corollary 2.3 is then a direct consequence of Theorem 2.2. Finally, in order to apply Theorem 2.2 to deduce the uniqueness result stated in Proposition 2.4 we need an *a priori*  $L^\infty$  bound for arbitrary nonnegative solutions to (12). The corresponding result provided in Lemma 4.1 and the conclusion are the contents of Section 4.

### 3 Proof of Theorem 2.2

Consider  $h_0 \in L^\infty(\mathbb{R})$ ,  $h_0 \geq 0$ , such that  $m_0$  defined in (21) belongs to  $L^1(0, 1)$ . In particular,  $m_0(\theta)$  is finite for almost every  $\theta \in [0, 1)$ . Let  $h$  be the solution to (10) with initial condition  $h(0) = h_0$  at time  $t = 0$ .

We introduce further suitable variables. For  $\theta \in [0, 1)$  such that  $m_0(\theta) \in (0, \infty)$  we define the right-continuous functions

$$\varphi_k^\theta(t) := h(t, k + 1 - \psi(t, \theta)) \quad \text{and} \quad \bar{\varphi}_k^\theta(t) := \bar{h}(k + 1 - \lambda(\theta) - \psi(t, \theta)) \quad (23)$$

for  $(t, k) \in [0, \infty) \times \mathbb{Z}$ , where  $\psi$  is defined in (16) and  $\lambda(\theta) = \ln m_0(\theta)/\alpha$ . Observe that  $\bar{\varphi}_k^\theta$  is 1-periodic.

The motivation for introducing these variables is as follows. In order to get rid of the transport term  $\partial_x h$  in (10) we as usual go over to characteristics. However, the solution is then continuously shifted to the right by the coagulation term. Thus, to keep the solution at scales of order one, we shift the solution back to the left at times  $n + \theta$ ,  $n \in \mathbb{N}$ . In a more quantitative way, we report the following identity:

**Lemma 3.1.** *For  $t > 0$  there holds*

$$\int_{\mathbb{R}} e^{\alpha x} \left| h(t, x) - \bar{h}(x - \mu(t, x)) \right| dx = \int_0^1 e^{\alpha(1-\psi(t,\theta))} \sum_{k \in \mathbb{Z}} e^{\alpha k} \left| \varphi_k^\theta(t) - \bar{\varphi}_k^\theta(t) \right| d\theta.$$

*Proof.* Clearly

$$\begin{aligned} & \int_{\mathbb{R}} e^{\alpha x} \left| h(t, x) - \bar{h}(x - \mu(t, x)) \right| dx \\ &= \int_0^1 e^{\alpha y} \sum_{k \in \mathbb{Z}} e^{\alpha k} \left| h(t, y + k) - \bar{h}(y + k - \mu(t, y + k)) \right| dy. \end{aligned}$$

We next make the change of variables  $y = \theta - t + [t]$  for  $y \in (t - [t], 1)$  and  $y = 1 + \theta - t + [t]$  for  $y \in (0, t - [t])$  and use the definition of  $\mu$  and the properties of  $\Theta$  and  $\psi$  to complete the proof.  $\square$

Define next

$$\mathcal{J}^\theta := (0, \theta) \cup \bigcup_{n \in \mathbb{N}} (n + \theta, n + 1 + \theta). \quad (24)$$

It follows from (10) that

$$\frac{d}{dt} \varphi_k^\theta(t) = \alpha \varphi_k^\theta(t) + e^{-\alpha} (\varphi_{k-1}^\theta)^2(t) - (\varphi_k^\theta)^2(t) \quad \text{for } t \in \mathcal{J}^\theta \quad (25)$$

and, by (17) and (23),

$$\varphi_k^\theta((n + \theta)^-) = \lim_{t \nearrow n + \theta} \varphi_k^\theta(t) = h(t, n + \theta) = \lim_{t \searrow n + \theta} \varphi_{k-1}^\theta(t) = \varphi_{k-1}^\theta((n + \theta)). \quad (26)$$

Also,

$$\sum_{k \in \mathbb{Z}} e^{\alpha(k + \theta)} \varphi_k^\theta(0) = m_0(\theta), \quad (27)$$

and, since

$$\frac{d}{dt} \sum_{k \in \mathbb{Z}} e^{\alpha k} \varphi_k^\theta(t) = \alpha \sum_{k \in \mathbb{Z}} e^{\alpha k} \varphi_k^\theta(t), \quad t \in \mathcal{J}^\theta,$$

by (25), we deduce that

$$\sum_{k \in \mathbb{Z}} e^{\alpha k} \varphi_k^\theta(t) = e^{\alpha(t - \theta)} m_0(\theta), \quad t \in [0, \theta),$$

and, for  $n \in \mathbb{N}$ ,

$$\sum_{k \in \mathbb{Z}} e^{\alpha k} \varphi_k^\theta(t) = e^{\alpha(t - n - \theta)} \sum_{k \in \mathbb{Z}} e^{\alpha k} \varphi_k^\theta(n + \theta), \quad t \in [n + \theta, n + 1 + \theta).$$

As

$$\sum_{k \in \mathbb{Z}} e^{\alpha k} \varphi_k^\theta((n + \theta)^-) = e^\alpha \sum_{k \in \mathbb{Z}} e^{\alpha k} \varphi_k^\theta(n + \theta)$$

for  $n \in \mathbb{N}$  by (26), an induction argument leads us to the identity

$$\sum_{k \in \mathbb{Z}} e^{\alpha k} \varphi_k^\theta(t) = e^{\alpha(\psi(t, \theta) - 1)} m_0(\theta), \quad t \geq 0. \quad (28)$$

Next, thanks to (12), (22), and (23), we easily check that  $(\bar{\varphi}_k^\theta)_k$  satisfy also (25), (26), and (28). Due to (14) we also have that  $\bar{\varphi}_k^\theta(t) \leq C \exp(-L2^k)$  as  $k \rightarrow \infty$  uniformly for  $t \in [0, \infty)$ .

We now prove a well-posedness result for (25)-(26) that in particular gives a uniform bound on  $(\varphi_k^\theta)_k$  that is independent of  $\theta$ . Let  $\mathcal{Y}$  be the space

$$\mathcal{Y} := \left\{ \phi = (\phi_k)_{k \in \mathbb{Z}} : \phi \in l_\infty(\mathbb{Z}) \text{ and } \sum_{k \in \mathbb{Z}} e^{\alpha k} |\phi_k| < \infty \right\}$$

and denote its positive cone by  $\mathcal{Y}_+$ .

**Lemma 3.2.** *Fix  $\theta \in [0, 1)$  such that  $m_0(\theta) < \infty$ . Let  $\varphi^\theta(0) = (\varphi_k^\theta(0))_k \in \mathcal{Y}_+$  such that*

$$\left\| \varphi^\theta(0) \right\|_{l_\infty} \leq C_0 \text{ and } \sum_{k \in \mathbb{Z}} e^{\alpha k} \varphi_k^\theta(0) = e^{-\alpha \theta} m_0(\theta) \quad (29)$$

for some  $C_0 > 0$ . Then there exists a unique function  $\varphi^\theta \in C^0([0, \infty); \mathcal{Y}_+)$  with  $\varphi_k^\theta \in C^1(\mathcal{J}^\theta)$  for all  $k \in \mathbb{Z}$  which solves (25) and (26) with initial condition  $\varphi^\theta(0)$  and satisfies (28) as well as the following property:

$$\sup_{t \geq 0} \left\| \varphi^\theta(t) \right\|_{l_\infty} \leq c_0 := \max \left\{ C_0, \frac{\alpha}{1 - e^{-\alpha}} \right\}. \quad (30)$$

*Proof.* Throughout the proof we omit the dependence on  $\theta$  in the notation. It turns out to be convenient to go over to the unknown function

$$\rho_k(t) := e^{-\alpha(\psi(t)-1)} \varphi_k(t), \quad (t, k) \in [0, \infty) \times \mathbb{Z}.$$

Then we are looking for a solution of

$$\frac{d}{dt} \rho_k(t) = e^{\alpha(\psi(t)-1)} \left( e^{-\alpha} \rho_{k-1}^2(t) - \rho_k^2(t) \right), \quad t \in \mathcal{J}, \quad (31)$$

satisfying

$$\rho_k((n + \theta)^-) = e^{-\alpha} \rho_{k-1}(n + \theta), \quad n \in \mathbb{N}, \quad k \in \mathbb{Z}, \quad (32)$$

and

$$\sum_{k \in \mathbb{Z}} e^{\alpha k} \rho_k(t) = m_0(\theta) \quad \text{for all } t. \quad (33)$$

Introducing  $R(\rho) := (R_k(\rho))_{k \in \mathbb{Z}}$  with  $R_k(\rho) := e^{-\alpha} \rho_{k-1}^2 - \rho_k^2$ , one has

$$\|R(\rho) - R(\sigma)\|_{l_1(e^{\alpha k})} \leq 2 (\|\rho\|_{l_\infty} + \|\sigma\|_{l_\infty}) \|\rho - \sigma\|_{l_1(e^{\alpha k})}$$

for any  $\rho = (\rho_k)_k$  and  $\sigma = (\sigma_k)_k$  in  $\mathcal{Y}$ . Consequently, a standard fixed point argument implies that there is  $T \in (0, \theta)$  small enough such that there is a unique solution to (31) with initial condition  $(e^{\alpha \theta} \varphi_k(0))_k$  in

$$\mathcal{M}_T := \{ \rho = (\rho_k)_k \in C^0([0, T]; l_1(e^{\alpha k})) : \|\rho\|_{l_\infty} \leq 2c_0 \}.$$

Nonnegativity of each component  $\rho_k$ ,  $k \in \mathbb{Z}$ , of the solution follows from the maximum principle. In order to show that  $\rho$  satisfies (33) in  $[0, T]$  we use (31) to compute, for  $M \geq 1$ ,

$$\frac{d}{dt} \sum_{k=-M}^M e^{\alpha k} \rho_k(t) = e^{\alpha(\psi(t)-1)} \left( e^{-\alpha(M+1)} \rho_{-(M+1)}^2(t) - e^{\alpha M} \rho_M^2(t) \right),$$

so that

$$\frac{d}{dt} \sum_{k=-M}^M e^{\alpha k} \rho_k(t) \geq -e^{\alpha M} \rho_M^2(t), \quad (34)$$

and, since  $\rho \in \mathcal{M}_T$ ,

$$\frac{d}{dt} \sum_{k=-M}^M e^{\alpha k} \rho_k(t) \leq e^{\alpha(\psi(t)-1)} e^{-\alpha(M+1)} \rho_{-(M+1)}^2(t) \leq 4c_0^2 e^{-\alpha(M+1)}.$$

Integrating the previous inequality and letting  $M \rightarrow \infty$  we obtain that

$$\sum_{k \in \mathbb{Z}} e^{\alpha k} \rho_k(t) \leq m_0 \quad \text{for all } t \in [0, T].$$

In particular,  $\rho_M(t) \leq m_0 e^{-\alpha M}$  for  $t \in [0, T]$  and we improve the lower bound (34) to

$$\frac{d}{dt} \sum_{k=-M}^M e^{\alpha k} \rho_k(t) \geq -e^{-\alpha M} m_0^2.$$

Integrating the above inequality and letting  $M \rightarrow \infty$  give

$$\sum_{k \in \mathbb{Z}} e^{\alpha k} \rho_k(t) \geq m_0 \quad \text{for all } t \in [0, T],$$

and, as a consequence, the solution  $\rho$  satisfies (33) in  $[0, T]$ .

It remains to show that (30) is satisfied in  $[0, T]$  from which it follows that we can extend the solution to the interval  $[0, \theta]$ . We can then repeat the argument in all subsequent intervals  $[n + \theta, n + \theta + 1]$ ,  $n \in \mathbb{N}$ .

To prove (30) we use a maximum principle argument and note that, given  $\gamma \in C^1([0, T])$  and  $M \in \mathbb{N}$ , it follows from (31) that

$$\begin{aligned} \frac{d}{dt} \sum_{k=-\infty}^M e^{\alpha k} (\rho_k - \gamma)_+ &= \sum_{k=-\infty}^M e^{\alpha(k-1)} (\rho_{k-1}^2 - \gamma^2) \text{sign}_+(\rho_k - \gamma) e^{\alpha(\psi-1)} \\ &\quad + \sum_{k=-\infty}^M e^{\alpha k} (e^{-\alpha} - 1) \gamma^2 \text{sign}_+(\rho_k - \gamma) e^{\alpha(\psi-1)} \\ &\quad - \sum_{k=-\infty}^M e^{\alpha k} (\rho_k^2 - \gamma^2)_+ e^{\alpha(\psi-1)} \\ &\quad - \sum_{k=-\infty}^M e^{\alpha k} \text{sign}_+(\rho_k - \gamma) \frac{d\gamma}{dt} \end{aligned}$$

Choosing  $\gamma$  such that

$$\frac{d\gamma}{dt} = -e^{\alpha(\psi-1)}(1 - e^{-\alpha})\gamma^2, \quad \gamma(0) = e^{\alpha\theta}C_0,$$

we realize that

$$\begin{aligned} \frac{d}{dt} \sum_{k=-\infty}^M e^{\alpha k} (\rho_k - \gamma)_+ &\leq \sum_{k=-\infty}^{M-1} e^{\alpha k} (\rho_k^2 - \gamma^2)_+ e^{\alpha(\psi-1)} \\ &\quad - \sum_{k=-\infty}^M e^{\alpha k} (\rho_k^2 - \gamma^2)_+ e^{\alpha(\psi-1)} \leq 0, \end{aligned}$$

hence

$$\sum_{k=-\infty}^M e^{\alpha k} (\rho_k - \gamma)_+(t) \leq \sum_{k=-\infty}^M e^{\alpha k} (\rho_k - \gamma)_+(0) = 0, \quad t \in [0, T].$$

As  $M \in \mathbb{N}$  is arbitrary we conclude that

$$\rho_k(t) \leq \gamma(t), \quad (t, k) \in [0, T] \times \mathbb{Z}.$$

Since  $\psi(t, \theta) - 1 = t - \theta$  for  $t \in (0, \theta)$ , the function  $\gamma$  is given by

$$\gamma(t) = \frac{\alpha C_0 e^{\alpha\theta}}{\alpha + C_0(e^{\alpha t} - 1)(1 - e^{-\alpha})}, \quad t \in [0, T],$$

and we end up with

$$\varphi_k(t) = e^{\alpha(t-\theta)} \rho_k(t) \leq e^{\alpha(t-\theta)} \gamma(t) \leq \sup_{\tau \geq 0} \left\{ \frac{\alpha C_0 e^{\alpha\tau}}{\alpha + C_0(e^{\alpha\tau} - 1)(1 - e^{-\alpha})} \right\} = c_0$$

for  $(t, k) \in [0, T] \times \mathbb{Z}$ , which proves the claim.  $\square$

To proceed further, we set

$$L^\theta(t) := \sum_k e^{\alpha k} \left( \varphi_k^\theta(t) - \bar{\varphi}_k^\theta(t) \right)_+, \quad t \geq 0, \quad (35)$$

which is a right-continuous function of time, and notice that (23) and (26) give

$$L^\theta((n + \theta)^-) = e^\alpha L^\theta((n + \theta)). \quad (36)$$

Our goal is to show that  $L^\theta(t) \rightarrow 0$  as  $t \rightarrow \infty$ . This will follow from the fact that  $L^\theta$  is almost a Lyapunov functional and that we can provide a lower bound on the dissipation functional. Towards that aim we first prove a tightness result.

**Lemma 3.3.** *Given  $\varepsilon > 0$  and  $\theta \in [0, 1)$  such that  $m_0(\theta) < \infty$ , let  $M = M(\theta, \varepsilon) > 0$  be such that*

$$\sum_{k \geq M} e^{\alpha k} h_0(k + \theta) < \varepsilon. \quad (37)$$

*Then there exists  $N = N(\varepsilon, M, m_0(\theta), \|h_0\|_{L^\infty})$  such that*

$$\sum_{|k| \geq N} e^{\alpha k} \varphi_k^\theta(t) \leq 2e^\alpha \varepsilon \quad \text{for all } t \in [0, \infty), \quad (38)$$

*and*

$$\sum_{|k| \geq N} e^{\alpha k} \bar{\varphi}_k^\theta(t) \leq 2e^\alpha \varepsilon \quad \text{for all } t \in [0, \infty). \quad (39)$$

*Proof.* Within this proof we again omit the dependence on  $\theta$  in the notation.

We first notice that the uniform bound (30) and the boundedness of  $\bar{h}$  imply that there exists  $N_1 \in \mathbb{N}$  such that

$$\sum_{k \leq -N_1} e^{\alpha k} \varphi_k(t) + \sum_{k \leq -N_1} e^{\alpha k} \bar{\varphi}_k(t) \leq \varepsilon, \quad t \in [0, \infty).$$

In order to control the mass at large positive  $k$  we construct a supersolution for the quantiles. More precisely, for  $t \geq 0$ , we define

$$Q_k(t) := \sum_{\ell \geq k} e^{\alpha \ell} \varphi_\ell(t) \quad \text{and} \quad \bar{Q}_k(t) := \sum_{\ell \geq k} e^{\alpha \ell} \bar{\varphi}_{\ell - \ell_0}(t) + \varepsilon e^{\alpha \psi(t)}$$

with some sufficiently large  $\ell_0 \in \mathbb{N}$  that will be determined later. Owing to (25) we note that  $Q_k$  solves

$$\frac{d}{dt} Q_k = \alpha Q_k + e^{-\alpha(k-1)} (Q_{k-1} - Q_k)^2, \quad t \in \mathcal{J},$$

and  $\bar{Q}_k$  solves the same equation for  $k \in \mathbb{Z}$ . Furthermore, by (23) and (26), we have

$$Q_k((n + \theta)^-) = e^\alpha Q_{k-1}((n + \theta)) \quad \text{and} \quad \bar{Q}_k((n + \theta)^-) = e^\alpha \bar{Q}_{k-1}((n + \theta))$$

for  $n \in \mathbb{N}$  and  $k \in \mathbb{Z}$ . We also observe that, by (28),

$$\lim_{k \rightarrow -\infty} Q_k(t) = e^{\alpha(\psi(t)-1)} m_0 \quad \text{and} \quad \lim_{k \rightarrow -\infty} \bar{Q}_k(t) = e^{\alpha(\ell_0 + \psi(t))} m_0 + \varepsilon e^{\alpha \psi(t)}, \quad t \geq 0.$$

For the difference  $W_k := Q_k - \bar{Q}_k$  we obtain

$$\frac{d}{dt} W_k = \alpha W_k + e^{-\alpha(k-1)} (W_{k-1} - W_k) (\bar{Q}_{k-1} - \bar{Q}_k + Q_{k-1} - Q_k), \quad (t, k) \in \mathcal{J} \times \mathbb{Z}. \quad (40)$$

Since

$$\bar{Q}_{k-1} - \bar{Q}_k + Q_{k-1} - Q_k = e^{\alpha(k-1)} \bar{\varphi}_{k-1-\ell_0} + e^{\alpha(k-1)} \varphi_{k-1},$$

we infer from (23) and (30) that

$$0 \leq \bar{Q}_{k-1} - \bar{Q}_k + Q_{k-1} - Q_k \leq (c_0 + \|\bar{h}\|_{L^\infty}) e^{\alpha(k-1)}, \quad k \in \mathbb{Z}. \quad (41)$$

In addition,

$$\lim_{k \rightarrow -\infty} W_k(t) = m_0 e^{\alpha(\psi(t)-1)} (1 - e^{\alpha \ell_0}) - \varepsilon e^{\alpha \psi(t)} \leq -\varepsilon, \quad t \in [0, \infty),$$

so that

$$\sum_{k=-\infty}^K (W_k)_+ = \sum_{k=-K}^K (W_k)_+$$

for  $K \in \mathbb{N}$  large enough. We then deduce from (40) and (41) that

$$\begin{aligned} \frac{d}{dt} \sum_{k=-\infty}^K (W_k)_+ &\leq \alpha \sum_{k=-\infty}^K (W_k)_+ + \sum_{k=-\infty}^K e^{-\alpha(k-1)} [\bar{Q}_{k-1} - \bar{Q}_k + Q_{k-1} - Q_k] (W_{k-1})_+ \\ &\leq (\alpha + c_0 + \|\bar{h}\|_{L^\infty}) \sum_{k=-\infty}^K (W_k)_+, \end{aligned}$$

and thus, after integration,

$$\sum_{k=-\infty}^K (W_k)_+(t) \leq \sum_{k=-\infty}^K (W_k)_+((n-1+\theta)_+), \quad t \in [(n-1+\theta)_+, n+\theta], \quad n \in \mathbb{N}. \quad (42)$$

We now choose  $\ell_0 = \ell_0(M, m_0)$  so large such that

$$\bar{Q}_k(0) \geq 2m_0 \geq Q_k(0) \quad \text{for all } k \leq M.$$

Also, by (37), we have

$$\bar{Q}_k(0) \geq \varepsilon \geq Q_k(0) \quad \text{for all } k > M.$$

With this choice of  $\ell_0$ , we conclude that  $W_k(0) \leq 0$  for all  $k \in \mathbb{Z}$ , hence  $W_k(t) \leq 0$  for  $t \in [0, \theta)$  and  $k \in \mathbb{Z} \cap (-\infty, K]$  by (42). Since  $K \in \mathbb{N}$  is arbitrary we have thus proved that  $Q_k(t) \leq \bar{Q}_k(t)$  for  $t \in [0, \theta)$  and  $k \in \mathbb{Z}$ . In particular,

$$e^\alpha Q_{k-1}(\theta) = Q_k(\theta_-) \leq \bar{Q}_k(\theta_-) = e^\alpha \bar{Q}_{k-1}(\theta),$$

which allows us to iterate the above argument and end up with

$$Q_k(t) \leq \bar{Q}_k(t), \quad (t, k) \in [0, \infty) \times \mathbb{Z}.$$

Now, according to (23),  $\bar{Q}_k$  is 1-periodic and, for  $t \in [0, 1)$ , we infer from (14) that

$$\bar{Q}_k(t) \leq C \sum_{\ell \geq k} e^{\alpha \ell} \exp\left(-L 2^{\ell - \ell_0 + 1 - \lambda - \psi(t)}\right) \leq (2e^\alpha - 1)\varepsilon$$

and

$$\sum_{\ell \geq k} e^{\alpha \ell} \bar{\varphi}_k(t) \leq C \sum_{\ell \geq k} e^{\alpha \ell} \exp\left(-L2^{\ell+1-\lambda-\psi(t)}\right) \leq (2e^\alpha - 1)\varepsilon$$

for  $k \geq N_2(\theta)$  sufficiently large. Choosing  $N = \max\{N_1, N_2\}$  finishes the proof.  $\square$

**Lemma 3.4.** Consider  $\theta \in [0, 1)$  such that  $m_0(\theta) < \infty$ . For  $t \in \mathcal{J}^\theta$  we have

$$\frac{d}{dt} L^\theta(t) = \alpha L^\theta(t) - D^\theta(t),$$

where

$$D^\theta := \sum_{k \in \mathbb{Z}} e^{\alpha k} (\varphi_k^\theta + \bar{\varphi}_k^\theta) w_k^\theta [\text{sign}_+(w_k^\theta) - \text{sign}_+(w_{k+1}^\theta)] \geq 0,$$

and  $w_k^\theta := \varphi_k^\theta - \bar{\varphi}_k^\theta$ ,  $k \in \mathbb{Z}$ .

*Proof.* This is a simple explicit computation.  $\square$

**Lemma 3.5.** Consider  $\theta \in [0, 1)$  such that  $m_0(\theta) < \infty$  and  $\varepsilon > 0$ . Let  $N \in \mathbb{N}$  be such that (38) holds and let  $\mathcal{T}_{\varepsilon, N}^\theta$  be the set of solutions  $q = (q_k)_{k \in \mathbb{Z}}$  to (25) on  $(\theta, \theta + 1)$  satisfying the mass equation (28), the uniform bound (30), the tightness estimate (38), and the lower bound

$$\sum_{k \in \mathbb{Z}} e^{\alpha k} (q_k(t) - \bar{\varphi}_k(t))_+ \geq 16e^\alpha \varepsilon \quad (43)$$

on the interval  $[\theta, \theta + 1)$ .

Then there exists  $\nu = \nu(\theta, N, \varepsilon) > 0$  such that

$$\int_\theta^{\theta+1} \sum_{k \in \mathbb{Z}} e^{\alpha k} (q_k(t) + \bar{\varphi}_k(t)) \omega_k(t) [\text{sign}_+(\omega_k(t)) - \text{sign}_+(\omega_{k+1}(t))] dt \geq \nu,$$

where  $\omega_k := q_k - \bar{\varphi}_k^\theta$ ,  $k \in \mathbb{Z}$ .

*Proof.* Throughout this proof we again omit the dependence on  $\theta$ .

Assume for contradiction that there exists a sequence  $(p^m)_{m \geq 1}$  in  $\mathcal{T}_{\varepsilon, N}$  such that

$$\lim_{m \rightarrow \infty} \int_\theta^{\theta+1} D_m(t) dt = 0, \quad (44)$$

where  $p^m = (p_k^m)_{k \in \mathbb{Z}}$ ,  $W_k^m := p_k^m - \bar{\varphi}_k$ , and

$$D_m := \sum_{k \in \mathbb{Z}} e^{\alpha k} (p_k^m + \bar{\varphi}_k) W_k^m [\text{sign}_+(W_k^m) - \text{sign}_+(W_{k+1}^m)].$$

Owing to (25) and (30),  $(p_k^m)_{m \geq 1}$  is bounded in  $C^1([\theta, \theta + 1])$  for all  $k \in \mathbb{Z}$  and we can extract a subsequence, again denoted by  $(p^m)_{m \geq 1}$  such that  $(p_k^m)_{m \geq 1}$  converges uniformly in  $[\theta, \theta + 1)$  to a function  $p_k$  for all  $k \in \mathbb{Z}$ . Setting  $p := (p_k)_{k \in \mathbb{Z}}$ , we easily see that  $p$

is a solution to (25) and satisfies the uniform bound (30) and the tightness bound (38) in  $[\theta, \theta + 1)$ . In addition, due to the tightness property (38), we deduce from the mass equation (28) and the lower bound (43) which are valid for  $p^m$  that

$$\left| \sum_{k \in \mathbb{Z}} e^{\alpha k} p_k(t) - e^{\alpha(\psi(t)-1)} m_0 \right| \leq 4e^\alpha \varepsilon \quad \text{and} \quad \left| \sum_{|k| \leq N} e^{\alpha k} W_k(t) \right| \leq 4e^\alpha \varepsilon, \quad (45)$$

with  $W_k := p_k - \bar{\varphi}_k$ , as well as

$$\sum_{|k| \leq N} e^{\alpha k} (W_k)_+(t) \geq 8e^\alpha \varepsilon, \quad (46)$$

for  $t \in [\theta, \theta + 1)$ .

Now, since each term in the sum  $D_m$  is non-negative, there holds

$$\lim_{m \rightarrow \infty} \int_{\theta}^{\theta+1} b_k^m(t) dt = 0, \quad k \in \mathbb{Z}, \quad (47)$$

with

$$b_k^m := e^{\alpha k} (p_k^m + \bar{\varphi}_k) W_k^m (\text{sign}_+(W_k^m) - \text{sign}_+(W_{k+1}^m)), \quad k \in \mathbb{Z}. \quad (48)$$

It remains to take the limit in (47). Fix  $k \in \mathbb{Z} \cap [-N, N]$ . Introducing

$$b_k := e^{\alpha k} (p_k + \bar{\varphi}_k) W_k [\text{sign}_+(W_k) - \text{sign}_+(W_{k+1})], \quad k \in \mathbb{Z},$$

and the sets

$$\begin{aligned} \mathcal{B}_k &:= \{t \in (\theta, \theta + 1) : W_k(t) > 0 \text{ and } W_{k+1}(t) = 0\}, \\ \mathcal{C}_k &:= \{t \in (\theta, \theta + 1) : W_k(t) < 0 \text{ and } W_{k+1}(t) = 0\}, \\ \mathcal{G}_k &:= (\theta, \theta + 1) \setminus (\mathcal{B}_k \cup \mathcal{C}_k), \end{aligned}$$

we have  $(\theta, \theta + 1) = \mathcal{B}_k \cup \mathcal{C}_k \cup \mathcal{G}_k$  and we notice that

$$\lim_{m \rightarrow \infty} b_k^m(t) = b_k(t) \quad \text{for } t \in \mathcal{G}_k,$$

while

$$b_k(t) \leq \liminf_{m \rightarrow \infty} b_k^m(t) \quad \text{for } t \in \mathcal{C}_k.$$

Then, (47) implies that

$$0 = \int_{\mathcal{C}_k \cup \mathcal{G}_k} B_k(t) dt + \lim_{m \rightarrow \infty} \int_{\mathcal{B}_k} B_k^m(t) dt. \quad (49)$$

We now claim that  $|\mathcal{B}_k| = 0$ . To prove this we use the equation satisfied by  $W_{k+1}$  which reads, due to (25),

$$\frac{d}{dt} W_{k+1}(t) = \alpha W_{k+1}(t) + e^{-\alpha} [p_k(t) + \bar{\varphi}_k(t)] W_k(t) - [p_{k+1}(t) + \bar{\varphi}_{k+1}(t)] W_{k+1}(t) \quad (50)$$

for  $t \in (\theta, \theta + 1)$ . Let  $t_0 \in \mathcal{B}_k$ . By (50)

$$\frac{d}{dt} W_{k+1}(t_0) = e^{-\alpha} [p_k(t_0) + \bar{\varphi}_k(t_0)] W_k(t_0) > 0.$$

Consequently, there is  $\delta > 0$  such that  $\mathcal{B}_k \cap [t_0 - \delta, t_0 + \delta] = \{t_0\}$  and  $\mathcal{B}_k$  contains only isolated points. This implies that  $|\mathcal{B}_k| = 0$  and therefore (49) reduces to

$$0 = \int_{\theta}^{\theta+1} b_k(t) dt. \quad (51)$$

This in turn implies that

$$W_k [\text{sign}_+(W_k) - \text{sign}_+(W_{k+1})] = 0 \quad \text{a.e. in } (\theta, \theta + 1), \quad k \in \mathbb{Z} \cap [-N, N]. \quad (52)$$

We will now show that the functions  $\varphi_k$  and the corresponding  $W_k$  constructed above cannot exist. Indeed, assume first that there is  $t_* \in [\theta, \theta + 1)$  such that  $W_{-N}(t_*) > 0$ . Due to (45) and (46),  $W_k(t_*)$  cannot be positive for all  $k \in \mathbb{Z} \cap [-N, N]$  and we define

$$k_* + 1 := \min\{k \in \mathbb{Z} \cap [-N, N] : W_k(t_*) \leq 0\} \leq N.$$

Then  $W_{k_*}(t_*) > 0$  and  $W_{k_*+1}(t_*) \leq 0$ . Since

$$\frac{d}{dt} W_{k_*+1} = \alpha W_{k_*+1} + e^{-\alpha} [p_{k_*} + \bar{\varphi}_{k_*}] W_{k_*} - [p_{k_*+1} + \bar{\varphi}_{k_*+1}] W_{k_*+1}$$

by (50), we realize that

$$\frac{d}{dt} [W_{k_*+1}(t) e^{-\alpha t}] \Big|_{t=t_*} > 0,$$

so that

$$W_{k_*+1}(t) e^{-\alpha t} < W_{k_*+1}(t_*) e^{-\alpha t_*} \leq 0, \quad t \in [t_* - \delta, t_*],$$

for some  $\delta > 0$ . Reducing  $\delta$  if necessary, we may also assume that  $W_{k_*}(t) \geq W_{k_*}(t_*)/2$  for  $t \in [t_* - \delta, t_*]$ . This implies that

$$W_{k_*}(t) [\text{sign}_+(W_{k_*}(t)) - \text{sign}_+(W_{k_*+1}(t))] = W_{k_*}(t) \geq \frac{W_{k_*}(t_*)}{2} > 0$$

for  $t \in [t_* - \delta, t_*)$  and contradicts (52).

Therefore  $W_{-N}(t) \leq 0$  for all  $t \in [\theta, \theta + 1)$ . We fix any value  $t_* \in [\theta, \theta + 1)$  and define

$$k_1 := \min\{k \in \mathbb{Z} \cap [-N, N] : W_k(t_*) \geq 0\}.$$

Then  $k_1 > -N$  and (46) guarantees that  $k_1 \leq N$ . Since  $W_{k_1-1}(t_*) < 0$  and  $W_{k_1}(t_*) \geq 0$ , we use once more (50) to obtain

$$\frac{d}{dt} [W_{k_1}(t) e^{-\alpha t}] \Big|_{t=t_*} < 0.$$

We may then find  $\delta > 0$  sufficiently small such that

$$W_{k_1-1}(t) < \frac{W_{k_1-1}(t_*)}{2} < 0 < W_{k_1}(t), \quad t \in [t_* - \delta, t_*].$$

This implies that

$$W_{k_1-1}(t) [\text{sign}_+(W_{k_1-1}(t)) - \text{sign}_+(W_{k_1}(t))] = -W_{k_1-1}(t) \geq -\frac{W_{k_1-1}(t_*)}{2} > 0$$

for  $t \in [t_* - \delta, t_*]$  and contradicts again (52). This concludes the proof.  $\square$

**Lemma 3.6.** *Let  $\theta \in [0, 1)$  be such that  $m_0(\theta) < \infty$ . The function  $L^\theta$  defined in (35) satisfies*

$$\lim_{t \rightarrow \infty} L^\theta(t) = \lim_{t \rightarrow \infty} \sum_{k \in \mathbb{Z}} e^{\alpha k} \left| \varphi_k^\theta(t) - \bar{\varphi}_k^\theta(t) \right| = 0.$$

*Proof.* We first note that, for  $t \in \mathcal{J}^\theta$ , Lemma 3.4 ensures that

$$\frac{d}{dt} \left( L^\theta(t) e^{-\alpha t} \right) = -e^{\alpha t} D^\theta(t). \quad (53)$$

After integration we find, for  $n \in \mathbb{N}$ ,

$$L^\theta((n+1+\theta)^-) + \int_{n+\theta}^{n+1+\theta} D^\theta(s) ds \leq e^\alpha L^\theta(n+\theta),$$

which, together with (36), gives

$$e^\alpha L^\theta((n+1+\theta)) + \int_{n+\theta}^{n+1+\theta} D^\theta(s) ds \leq e^\alpha L^\theta(n+\theta). \quad (54)$$

In particular,  $(L^\theta(n+\theta))_{n \in \mathbb{N}}$  is a nonincreasing sequence and

$$\sum_{n \in \mathbb{N}} \int_{n+\theta}^{n+1+\theta} D^\theta(t) dt \leq e^\alpha L^\theta(\theta). \quad (55)$$

Assume for contradiction that

$$\eta := \inf_{n \in \mathbb{N}} L^\theta(n+\theta) > 0.$$

Then, owing to (36),  $L^\theta((n+\theta)^-) \geq e^\alpha \eta$  while we infer from (53) that

$$\begin{aligned} L^\theta(\theta^-) e^{\alpha(t-\theta)} &\leq L^\theta(t), & t \in [0, \theta), \\ L^\theta((n+1+\theta)^-) e^{\alpha(t-n-1-\theta)} &\leq L^\theta(t), & t \in [n+\theta, n+1+\theta), \end{aligned}$$

for  $n \in \mathbb{N}$ . Combining these estimates gives

$$L^\theta(t) \geq \eta, \quad t \in [0, \infty). \quad (56)$$

Consider now  $\varepsilon \in (0, \eta/(16e^\alpha))$ . Since  $m_0(\theta) < \infty$  there is  $M = M(\theta) \in \mathbb{N}$  such that (37) holds true and it follows from Lemma 3.3 that there is  $N = N(\theta, \varepsilon)$  such that

$$\sum_{|k| \geq N} e^{\alpha k} \varphi_k^\theta(t) \leq 2e^\alpha \varepsilon, \quad t \in [0, \infty). \quad (57)$$

We next define

$$p_k^m(s) := \varphi_k^\theta(m+s), \quad (s, m, k) \in [\theta, \theta+1) \times \mathbb{N} \times \mathbb{Z}.$$

Then, for each  $m \in \mathbb{N}$ ,  $p^m := (p_k^m)_{k \in \mathbb{Z}}$  is a solution to (25) in  $(\theta, \theta+1)$  which satisfies the mass equation (28) in  $[\theta, \theta+1)$ . It also enjoys the uniform estimate (30) by Lemma 3.2 as well as the tightness property (38) by (57), still in  $[\theta, \theta+1)$ . Finally, for  $t \in [\theta, \theta+1)$ ,

$$\sum_{k \in \mathbb{Z}} e^{\alpha k} \left( p_k^m(t) - \bar{\varphi}_k^\theta(t) \right)_+ = L^\theta(m+s) \geq \eta \geq 16e^\alpha \varepsilon$$

by (56). In other words,  $p^m$  belongs to the set  $\mathcal{T}_{\varepsilon, N}^\theta$  defined in Lemma 3.5 and we infer from Lemma 3.5 that there is  $\nu = \nu(\varepsilon, \theta) > 0$  such that

$$\int_\theta^{\theta+1} \sum_{k \in \mathbb{Z}} e^{\alpha k} \left( p_k^m(t) + \bar{\varphi}_k^\theta(t) \right) W_k^m(t) \left[ \text{sign}_+(W_k^m(t)) - \text{sign}_+(W_{k+1}^m(t)) \right] dt \geq \nu,$$

where  $W_k^m := p_k^m - \bar{\varphi}_k^\theta$ ,  $(m, k) \in \mathbb{N} \times \mathbb{Z}$ . In terms of  $\varphi^\theta$ ,

$$\int_{m+\theta}^{m+1+\theta} D^\theta(t) dt \geq \nu, \quad m \in \mathbb{N}.$$

Summing up with respect to  $m \in \mathbb{N}$  clearly contradicts (55). Therefore  $\eta = 0$  and, since

$$L^\theta(t) \leq L^\theta(n+\theta)e^\alpha, \quad t \in [n+\theta, n+1+\theta), \quad n \in \mathbb{N},$$

by (53), we conclude that  $L^\theta(t) \rightarrow 0$  as  $t \rightarrow \infty$ .

Finally, let  $t \geq 0$ . We recall that

$$\sum_{k \in \mathbb{Z}} e^{\alpha k} \varphi_k^\theta(t) = m_0(\theta) e^{\alpha(\psi(t, \theta) - 1)} = \sum_{k \in \mathbb{Z}} e^{\alpha k} \bar{\varphi}_k^\theta(t)$$

by (28), so that

$$\begin{aligned} \sum_{k \in \mathbb{Z}} e^{\alpha k} \left| \varphi_k^\theta(t) - \bar{\varphi}_k^\theta(t) \right| &= \sum_{k \in \mathbb{Z}} e^{\alpha k} \left[ 2 \left( \varphi_k^\theta(t) - \bar{\varphi}_k^\theta(t) \right)_+ - \left( \varphi_k^\theta(t) - \bar{\varphi}_k^\theta(t) \right) \right] \\ &= 2L^\theta(t), \end{aligned}$$

and the proof is complete.  $\square$

**Proof of Theorem 2.2.** By Lemma 3.1

$$\int_{\mathbb{R}} e^{\alpha x} \left| h(t, x) - \bar{h}(x - \mu(t, x)) \right| dx = \int_0^1 e^{\alpha(1-\psi(t,\theta))} \sum_{k \in \mathbb{Z}} e^{\alpha k} \left| \varphi_k^\theta(t) - \bar{\varphi}_k^\theta(t) \right| d\theta.$$

On the one hand, since  $m_0 \in L^1(0, 1)$ , then  $m_0(\theta)$  is finite for almost every  $\theta \in (0, 1)$  and, according to Lemma 3.6,

$$\lim_{t \rightarrow \infty} \sum_{k \in \mathbb{Z}} e^{\alpha k} \left| \varphi_k^\theta(t) - \bar{\varphi}_k^\theta(t) \right| = 0$$

for almost every  $\theta \in (0, 1)$ . On the other hand, thanks to (28),

$$e^{\alpha(1-\psi(t,\theta))} \sum_{k \in \mathbb{Z}} e^{\alpha k} \left| \varphi_k^\theta(t) - \bar{\varphi}_k^\theta(t) \right| \leq 2m_0(\theta).$$

The Lebesgue dominated convergence then entails that

$$\lim_{t \rightarrow \infty} \int_{\mathbb{R}} e^{\alpha x} \left| h(t, x) - \bar{h}(x - \mu(t, x)) \right| dx = 0,$$

which proves the claim. □

## 4 A priori estimates for stationary solutions

**Lemma 4.1.** *Let  $\bar{h} \in C^1(\mathbb{R}) \cap L^1(\mathbb{R}; e^{\alpha x} dx)$  be a nonnegative solution to (12). Then  $\bar{h} \in L^\infty(\mathbb{R})$ .*

*Proof. Step 1:* We first claim that the function  $\bar{h}$  satisfies

$$e^{\alpha x} \bar{h}(x) = \int_{x-1}^x \bar{h}^2(t) e^{\alpha t} dt \quad \text{for all } x \in \mathbb{R}. \quad (58)$$

To prove (58), we first multiply (12) by  $e^{\alpha x}$  and integrate to find

$$e^{\alpha x} \bar{h}(x) = \int_{x-1}^x \bar{h}^2(t) e^{\alpha t} dt + J \quad (59)$$

for some  $J \in \mathbb{R}$ . In particular,  $e^{\alpha x} \bar{h}(x) \geq J$  and the integrability of  $x \mapsto e^{\alpha x} \bar{h}(x)$  implies that  $J \leq 0$ .

It next follows from the integrability of  $x \mapsto e^{\alpha x} \bar{h}(x)$  that there is a sequence  $(x_n)_{n \geq 1}$  in  $(1, \infty)$  such that

$$\lim_{n \rightarrow \infty} x_n = \infty \quad \text{and} \quad \bar{h}(x_n - 1) \leq e^{\alpha(x_n - 1)} \bar{h}(x_n - 1) \leq \frac{1}{2}, \quad n \geq 1. \quad (60)$$

We are going to show that the differential equation (12) for  $\bar{h}$  guarantees that  $\bar{h}$  does not exceed one in  $(x_n - 1, x_n)$ . Indeed, by (12),  $\partial_x \bar{h} \leq \bar{h}^2$  and integrating this equation gives, thanks to (60),

$$2 - \frac{1}{\bar{h}(x)} \leq \frac{1}{\bar{h}(x_n - 1)} - \frac{1}{\bar{h}(x)} \leq x - x_n + 1 \leq 1, \quad x \in [x_n - 1, x_n],$$

hence

$$\bar{h}(x) \leq 1, \quad x \in [x_n - 1, x_n]. \quad (61)$$

Combining (59) (with  $x = x_n$ ) and (61) gives

$$0 \leq -J = \int_{x_n-1}^{x_n} \bar{h}^2(y) e^{\alpha y} dy - e^{\alpha x_n} \bar{h}(x_n) \leq \int_{x_n-1}^{x_n} \bar{h}(y) e^{\alpha y} dy.$$

Since the right-hand side of the above inequality converges to zero as  $n \rightarrow \infty$ , we conclude that  $J = 0$  which proves (58).

*Step 2:* To complete the proof we argue by contradiction and assume that there is  $M \geq 8e^{2\alpha}$  and  $x_0 \in \mathbb{R}$  such that

$$\bar{h}(x_0) \geq M. \quad (62)$$

According to (58) we have either

$$\int_{x_0-1/2}^{x_0} \bar{h}^2(y) e^{\alpha y} dy \geq \frac{\bar{h}(x_0) e^{\alpha x_0}}{2} \quad (63)$$

or

$$\int_{x_0-1}^{x_0-1/2} \bar{h}^2(y) e^{\alpha y} dy \geq \frac{\bar{h}(x_0) e^{\alpha x_0}}{2}. \quad (64)$$

In case that (63) holds, consider  $y \in [x_0, x_0 + 1/2]$ . Then  $(x_0 - 1/2, x_0) \subset (y - 1, y)$  and we deduce from (58), (62), and (63) that

$$\bar{h}(y) = e^{-\alpha y} \int_{y-1}^y \bar{h}^2(z) e^{\alpha z} dz \geq e^{-\alpha y} \int_{x_0-1/2}^{x_0} \bar{h}^2(z) e^{\alpha z} dz \geq \frac{M}{2e^{\alpha/2}}.$$

Combining this lower bound with (58) gives

$$\bar{h}\left(x_0 + \frac{1}{2}\right) \geq e^{-\alpha(x_0+1/2)} \int_{x_0}^{x_0+1/2} \bar{h}^2(y) e^{\alpha y} dy \geq \frac{M^2}{8e^{2\alpha}} \geq M.$$

Similarly, if (64) holds, then  $(x_0 - 1, x_0 - 1/2) \subset (y - 1, y)$  for  $y \in [x_0 - 1/2, x_0]$  and it follows from (58), (62), and (64) that  $\bar{h}(y) \geq M/(2e^{\alpha/2})$ . In particular, we find

$$\bar{h}\left(x_0 + \frac{1}{2}\right) \geq e^{-\alpha(x_0+1/2)} \int_{x_0-1/2}^{x_0} \bar{h}^2(z) e^{\alpha z} dz \geq \frac{M^2}{8e^{2\alpha}} \geq M.$$

We have thus proved that

$$\bar{h}\left(x_0 + \frac{1}{2}\right) \geq M. \quad (65)$$

Arguing by induction, we actually conclude that

$$\bar{h}\left(x_0 + \frac{k}{2}\right) \geq M, \quad k \in \mathbb{N}. \quad (66)$$

Furthermore, arguing as in the proof of (65), we infer from (58) and (66) that

$$\bar{h}(y) \geq \frac{M}{2e^{\alpha/2}}, \quad y \in I_k,$$

where either  $I_k := (x_0 + (k-1)/2, x_0 + k/2)$  or  $I_k := (x_0 + k/2, x_0 + (k+1)/2)$ . Define

$$\mathcal{L} := \left\{ k \in \mathbb{N} : I_k = \left(x_0 + \frac{k-1}{2}, x_0 + \frac{k}{2}\right) \right\}$$

and

$$\mathcal{R} := \left\{ k \in \mathbb{N} : I_k = \left(x_0 + \frac{k}{2}, x_0 + \frac{k+1}{2}\right) \right\}.$$

Then  $\mathcal{L} \cup \mathcal{R} = \mathbb{N}$  and

$$\begin{aligned} 1 &= \int_{\mathbb{R}} e^{\alpha x} \bar{h}(x) dx \geq \frac{1}{2} \sum_{k \in \mathbb{N}} \int_{x_0 + (k-1)/2}^{x_0 + k/2} e^{\alpha x} \bar{h}(x) dx + \frac{1}{2} \sum_{k \in \mathbb{N}} \int_{x_0 + k/2}^{x_0 + (k+1)/2} e^{\alpha x} \bar{h}(x) dx \\ &\geq \frac{1}{2} \sum_{k \in \mathcal{L}} \int_{I_k} e^{\alpha x} \bar{h}(x) dx + \frac{1}{2} \sum_{k \in \mathcal{R}} \int_{I_k} e^{\alpha x} \bar{h}(x) dx \\ &\geq \frac{M}{8e^{\alpha/2}} \left[ \sum_{k \in \mathcal{L}} e^{\alpha(x_0 + (k-1)/2)} + \sum_{k \in \mathcal{L}} e^{\alpha(x_0 + k/2)} \right]. \end{aligned}$$

Since  $\mathcal{L}$  or  $\mathcal{R}$  is infinite, the right-hand side of the above inequality cannot be bounded, which gives a contradiction and hence proves the statement of the lemma.  $\square$

*Proof of Proposition 2.4.* Let  $\tilde{h} \in C^1(\mathbb{R}) \cap L^1(\mathbb{R}; e^{\alpha x} dx)$  be a nonnegative solution to (12). By Lemma 4.1 we have  $\tilde{h} \in L^\infty(\mathbb{R})$ . Then  $\tilde{h}$  is also a solution to the evolution equation (10) and, by Theorem 2.2,

$$\lim_{t \rightarrow \infty} \int_{\mathbb{R}} e^{\alpha x} \left| \tilde{h}(x) - \bar{h}(x - \tilde{\mu}(t, x)) \right| = 0,$$

where

$$\tilde{\mu}(t, x) := \frac{\ln(\tilde{m}_0(\Theta(t, x)))}{\alpha}, \quad (t, x) \in [0, \infty) \times \mathbb{R},$$

with  $\Theta$  defined in Definition 2.1 and

$$\tilde{m}_0(\theta) := \sum_{k \in \mathbb{Z}} e^{\alpha(k+\theta)} \tilde{h}(k + \theta), \quad \theta \in [0, 1).$$

Arguing as in the proof of (19), there exists  $\tilde{\nu} \in [0, \infty)$  such that  $\tilde{m}_0(\theta) = \tilde{\nu}$  for a.e.  $\theta \in (0, 1)$ , that is,  $\tilde{\mu}$  is a constant and the claim is proved.  $\square$

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