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Approximate and exact controllability of the continuity equation with a localized vector field*

Michel Duprez[†] Morgan Morancey[‡] Francesco Rossi[§]

October 18, 2017

Abstract

We study controllability of a Partial Differential Equation of transport type, that arises in crowd models. We are interested in controlling it with a control being a vector field, representing a perturbation of the velocity, localized on a fixed control set.

We prove that, for each initial and final configuration, one can steer approximately one to another with Lipschitz controls when the uncontrolled dynamics allows to cross the control set. We also show that the exact controllability only holds for controls with less regularity, for which one may lose uniqueness of the associated solution.

1 Introduction

In recent years, the study of systems describing a crowd of interacting autonomous agents has drawn a great interest from the control community (see *e.g.* the Cucker-Smale model [16]). A better understanding of such interaction phenomena can have a strong impact in several key applications, such as road traffic and egress problems for pedestrians. For a few reviews about this topic, see *e.g.* [5, 6, 9, 15, 21, 22, 26, 30].

Beside the description of interactions, it is now relevant to study problems of **control of crowds**, *i.e.* of controlling such systems by acting on few agents, or on the crowd localized in a small subset of the configuration space. The nature of the control problem relies on the model used to describe the crowd. Two main classes are widely used.

In **microscopic models**, the position of each agent is clearly identified; the crowd dynamics is described by a large dimensional ordinary differential equation, in which couplings of terms represent interactions. For control of such models, a large literature is available from the control community, under the generic name of networked control (see *e.g.* [8, 23, 24]). There are several control applications to pedestrian crowds [17, 25] and road traffic [10, 20].

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In **macroscopic models**, instead, the idea is to represent the crowd by the spatial density of agents; in this setting, the evolution of the density solves a partial differential equation of transport type. Nonlocal terms (such as convolution) model the interactions between the agents. In this article, we focus on this second approach, *i.e.* macroscopic models. To our knowledge, there exist few studies of control of this family of equations. In [28], the authors provide approximate alignment of a crowd described by the macroscopic Cucker-Smale model [16]. The control is the acceleration, and it is localized in a control region ω which moves in time. In a similar situation, a stabilization strategy has been established in [11, 12], by generalizing the Jurdjevic-Quinn method to partial differential equations.

A different approach is given by mean-field type control, *i.e.* control of mean-field equations and of mean-field games modeling crowds. See *e.g.* [1, 2, 13, 18]. In this case, problems are often of optimization nature, *i.e.* the goal is to find a control minimizing a given cost. In this article, we are mainly interested in controllability problems, for which mean-field type control approaches seem not adapted.

In this article, we study a macroscopic model, thus the crowd is represented by its density, that is a time-evolving measure $\mu(t)$ defined for positive times t on the space \mathbb{R}^d ($d \geq 1$). The natural (uncontrolled) velocity field for the measure is denoted by $v : \mathbb{R}^d \rightarrow \mathbb{R}^d$, being a vector field assumed Lipschitz and uniformly bounded. We act on the velocity field in a fixed portion ω of the space, which will be a **nonempty open connected subset** of \mathbb{R}^d . The admissible controls are thus functions of the form $\mathbb{1}_\omega u : \mathbb{R}^d \times \mathbb{R}^+ \rightarrow \mathbb{R}^d$. We then consider the following linear transport equation

$$\begin{cases} \partial_t \mu + \nabla \cdot ((v + \mathbb{1}_\omega u)\mu) = 0 & \text{in } \mathbb{R}^d \times \mathbb{R}^+, \\ \mu(0) = \mu^0 & \text{in } \mathbb{R}^d, \end{cases} \quad (1)$$

where μ^0 is the initial data (initial configuration of the crowd) and the function u is an admissible control. The function $v + \mathbb{1}_\omega u$ represents the velocity field acting on μ . System (1) is a first simple approximation for crowd modelling, since the uncontrolled vector field v is given, and it does not describe interactions between agents. Nevertheless, it is necessary to understand controllability properties for such simple equation as a first step, before dealing with velocity fields depending on the crowd itself. Thus, in a future work, we will study controllability of crowd models with a nonlocal term $v[\mu]$, based on the linear results presented here.

Even though System (1) is linear, the control acts on the velocity, thus the control problem is nonlinear, which is one of the main difficulties in this study.

The goal of this work is to study the control properties of System (1). We now recall the notion of approximate controllability and exact controllability for System (1). We say that System (1) is *approximately controllable* from μ^0 to μ^1 on the time interval $[0, T]$ if we can steer the solution to System (1) at time T as close to μ^1 as we want with an appropriate control $\mathbb{1}_\omega u$. Similarly, we say that System (1) is *exactly controllable* from μ^0 to μ^1 on the time interval $[0, T]$ if we can steer the solution to System (1) at time T exactly to μ^1 with an appropriate control $\mathbb{1}_\omega u$. In Definition 5 below, we give a formal definition of the notion of approximate controllability in terms of Wasserstein distance.

The main results of this article show that approximate and exact controllability depend on two main aspects: first, from a geometrical point of view, the uncontrolled vector field v needs to send the support of μ^0 to ω forward in time and the support of μ^1 to ω backward in time. This idea is formulated in the following Condition:

Condition 1.1 (Geometrical condition). Let μ^0, μ^1 be two probability measures on \mathbb{R}^d satisfying:

- (i) For each $x^0 \in \text{supp}(\mu^0)$, there exists $t^0 > 0$ such that $\Phi_{t^0}^v(x^0) \in \omega$, where Φ_t^v is the flow associated to v , i.e. the solution to the Cauchy problem

$$\begin{cases} \dot{x}(t) = v(x(t)) \text{ for a.e. } t > 0, \\ x(0) = x^0. \end{cases}$$

- (ii) For each $x^1 \in \text{supp}(\mu^1)$, there exists $t^1 > 0$ such that $\Phi_{-t^1}^v(x^1) \in \omega$.

This geometrical aspect is illustrated in Figure 1.

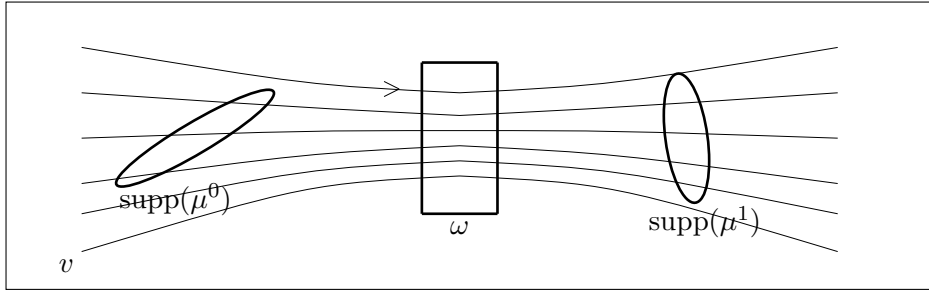


Figure 1: Geometrical condition.

Remark 1. Condition 1.1 is the minimal one that we can expect to steer any initial condition to any target. Indeed, if there exists a point x^0 of the interior of $\text{supp}(\mu^0)$ for which the first item of Condition 1.1 is not satisfied, then there exists a whole subpopulation of the measure μ^0 that never intersects the control region, thus we cannot act on it.

The second aspect that we want to highlight is the following: The measures μ^0 and μ^1 need to be sufficiently regular with respect to the flow generated by $v + \mathbb{1}_\omega u$. Three cases are particularly relevant:

a) Controllability with Lipschitz controls

If we impose the classical Carathéodory condition of $\mathbb{1}_\omega u$ being Lipschitz in space, measurable in time and uniformly bounded, then the flow $\Phi_t^{v+\mathbb{1}_\omega u}$ is an homeomorphism (see [7, Th. 2.1.1]). As a result, one can expect approximate controllability only, since for general measures there exists no homeomorphism sending one to another. For more details, see Section 4.1. We then have the following result:

Theorem 1.1 (Main result - Controllability with Lipschitz control). *Let μ^0, μ^1 be two probability measures on \mathbb{R}^d compactly supported, absolutely continuous with respect to the Lebesgue measure and satisfying Condition 1.1. Then there exists T such that System (1) is **approximately controllable** on the time interval $[0, T]$ from μ^0 to μ^1 with a control $\mathbb{1}_\omega u : \mathbb{R}^d \times \mathbb{R}^+ \rightarrow \mathbb{R}^d$ uniformly bounded, Lipschitz in space and measurable in time.*

We give a proof of Theorem 1.1 in Section 3.

b) Controllability with BV controls

To hope to obtain exact controllability of System (1), it is then necessary to search among controls $\mathbb{1}_\omega u$ with less regularity. A weaker condition on the regularity of the velocity

field for the well-posedness of System (1) has been given in [3] for vector fields v with bounded variations (we will use the abbreviation BV vector fields) satisfying in particular

$$\int_0^T \|[\operatorname{div} v]^- \|_{L^\infty(\mathbb{R}^d)} < +\infty. \quad (2)$$

As it will be explained in Section 4.1, if we choose the admissible controls satisfying the setting of [3], then it is not necessary that the support of μ^0 and μ^1 are homeomorph for exact controllability, but it is not sufficient either (see Section 4.1). For example, Condition (2) allows to steer a measure which support is connected to a measure which support is composed of two connected components, but the inverse is forbidden. Thus, even this setting does not allow to yield exact controllability.

c) Controllability with Borel controls

We then consider an even larger class of controls, that are Borel vector fields. In this setting, we have exact controllability under the geometrical Condition 1.1. The main drawback is that, in this less regular setting, System (1) is necessarily not well-posed. In particular, one has not necessarily uniqueness of the solution. For this reason, one needs to describe solutions to System (1) as pairs $(\mathbb{1}_\omega u, \mu)$, where μ is one among the admissible solutions with control $\mathbb{1}_\omega u$.

Theorem 1.2 (Main result - Controllability with Borel control). *Let μ^0, μ^1 be two probability measures on \mathbb{R}^d compactly supported and satisfying Condition 1.1. Then, there exists $T > 0$ such that System (1) is **exactly controllable** on the time interval $[0, T]$ from μ^0 to μ^1 in the following sense: there exists a couple $(\mathbb{1}_\omega u, \mu)$ composed of a Borel vector field $\mathbb{1}_\omega u : \mathbb{R}^d \times \mathbb{R}^+ \rightarrow \mathbb{R}^d$ and a time-evolving measure μ being weak solution to System (1) (see Definition 3) and satisfying*

$$\mu(T) = \mu_1.$$

A proof of Theorem 1.2 is given in Section 4.

This paper is organised as follows. In Section 2, we recall basic properties of the Wasserstein distance and the continuity equation. Section 3 is devoted to the proof of Theorem 1.1, *i.e.* the approximate controllability of System (1) with a Lipschitz localized vector field. Finally, in Section 4, we first show that exact controllability does not hold for Lipschitz or BV controls; we also prove Theorem 1.2, *i.e.* exact controllability of System (1) with a Borel localized vector field.

2 The Wasserstein distance and the continuity equation

In this section, we recall the definition and some properties of the Wasserstein distance and the continuity equation, which will be used all along this paper. We denote by $\mathcal{P}_c(\mathbb{R}^d)$ the space of probability measures in \mathbb{R}^d with compact support and for $\mu, \nu \in \mathcal{P}_c(\mathbb{R}^d)$, we denote by $\Pi(\mu, \nu)$ the set of *transference plans* from μ to ν , *i.e.* the probability measures on $\mathbb{R}^d \times \mathbb{R}^d$ satisfying

$$\int_{\mathbb{R}^d} d\pi(x, \cdot) = d\mu(x) \text{ and } \int_{\mathbb{R}^d} d\pi(\cdot, y) = d\nu(y).$$

Definition 1. Let $p \in [1, \infty)$ and $\mu, \nu \in \mathcal{P}_c(\mathbb{R}^d)$. Define

$$W_p(\mu, \nu) = \inf_{\pi \in \Pi(\mu, \nu)} \left\{ \left(\iint_{\mathbb{R}^d \times \mathbb{R}^d} |x - y|^p d\pi \right)^{1/p} \right\}. \quad (3)$$

The quantity is called the **Wasserstein distance**.

This is the idea of *optimal transportation*, consisting in finding the optimal way to transport mass from a given measure to another. For a thorough introduction, see *e.g.* [31].

We denote by Γ the set of Borel maps $\gamma : \mathbb{R}^d \rightarrow \mathbb{R}^d$. We now recall the definition of the *push-forward* of a measure:

Definition 2. For a $\gamma \in \Gamma$, we define the *push-forward* $\gamma\#\mu$ of a measure μ of \mathbb{R}^d as follows:

$$(\gamma\#\mu)(E) := \mu(\gamma^{-1}(E)),$$

for every subset E such that $\gamma^{-1}(E)$ is μ -measurable.

We denote by ‘‘AC measures’’ the measures which are absolutely continuous with respect to the Lebesgue measure and by $\mathcal{P}_c^{ac}(\mathbb{R}^d)$ the subset of $\mathcal{P}_c(\mathbb{R}^d)$ of AC measures. On $\mathcal{P}_c^{ac}(\mathbb{R}^d)$, the Wasserstein distance can be reformulated as follows:

Property 2.1 (see [31, Chap. 7]). *Let $p \in [1, \infty)$ and $\mu, \nu \in \mathcal{P}_c^{ac}(\mathbb{R}^d)$. It holds*

$$W_p(\mu, \nu) = \inf_{\gamma \in \Gamma} \left\{ \left(\int_{\mathbb{R}^d} |\gamma(x) - x|^p d\mu \right)^{1/p} : \gamma\#\mu = \nu \right\}. \quad (4)$$

The Wasserstein distance satisfies some useful properties:

Property 2.2 (see [31, Chap. 7]). *Let $p \in [1, \infty)$.*

- (i) *The Wasserstein distance W_p is a distance on $\mathcal{P}_c(\mathbb{R}^d)$.*
- (ii) *The topology induced by the Wasserstein distance W_p on $\mathcal{P}_c(\mathbb{R}^d)$ coincides with the weak topology.*
- (iii) *For all $\mu, \nu \in \mathcal{P}_c^{ac}(\mathbb{R}^d)$, the infimum in (4) is achieved by at least one minimizer.*

The Wasserstein distance can be extended to all pairs of measures μ, ν compactly supported with the same total mass $\mu(\mathbb{R}^d) = \nu(\mathbb{R}^d) \neq 0$, by the formula

$$W_p(\mu, \nu) = \mu(\mathbb{R}^d)^{1/p} W_p \left(\frac{\mu}{\mu(\mathbb{R}^d)}, \frac{\nu}{\nu(\mathbb{R}^d)} \right).$$

In the rest of the paper, the following properties of the Wasserstein distance will be also helpful:

Property 2.3 (see [27, 31]). *Let μ, ρ, ν, η be four positive measures compactly supported satisfying $\mu(\mathbb{R}^d) = \nu(\mathbb{R}^d)$ and $\rho(\mathbb{R}^d) = \eta(\mathbb{R}^d)$.*

- (i) *For each $p \in [1, \infty)$, it holds*

$$W_p^p(\mu + \rho, \nu + \eta) \leq W_p^p(\mu, \nu) + W_p^p(\rho, \eta). \quad (5)$$

- (ii) *For each $p_1, p_2 \in [1, \infty)$ with $p_1 \leq p_2$, it holds*

$$\begin{cases} W_{p_1}(\mu, \nu) \leq W_{p_2}(\mu, \nu), \\ W_{p_2}(\mu, \nu) \leq \text{diam}(X)^{1-p_1/p_2} W_{p_1}^{p_1/p_2}(\mu, \nu), \end{cases} \quad (6)$$

where X contains the supports of μ and ν .

We now recall the definition of the continuity equation and the associated notion of weak solutions:

Definition 3. Let $T > 0$ and μ^0 be a measure in \mathbb{R}^d . We said that a pair (μ, w) composed with a measure μ in $\mathbb{R}^d \times [0, T]$ and a vector field $w : \mathbb{R}^d \times \mathbb{R}^+ \rightarrow \mathbb{R}^d$ satisfying

$$\int_0^T \int_{\mathbb{R}^d} |w(t)| d\mu(t) dt < \infty$$

is a **weak solution** to the system, called the **continuity equation**,

$$\begin{cases} \partial_t \mu + \nabla \cdot (w\mu) = 0 & \text{in } \mathbb{R}^d \times [0, T], \\ \mu(0) = \mu^0 & \text{in } \mathbb{R}^d, \end{cases} \quad (7)$$

if for every continuous bounded function $\xi : \mathbb{R}^d \rightarrow \mathbb{R}$, the function $t \mapsto \int_{\mathbb{R}^d} \xi d\mu(t)$ is absolutely continuous with respect to t and for all $\psi \in \mathcal{C}_c^\infty(\mathbb{R}^d)$, it holds

$$\frac{d}{dt} \int_{\mathbb{R}^d} \psi d\mu(t) = \int_{\mathbb{R}^d} \langle \nabla \psi, w(t) \rangle d\mu(t)$$

for *a.e.* t and $\mu(0) = \mu^0$.

Note that $t \mapsto \mu(t)$ is continuous for the weak convergence, it then make sense to impose the initial condition $\mu(0) = \mu^0$ pointwisely in time. Before stating a result of existence and uniqueness of solutions for the continuity equation, we first recall the definition of the flow associated to a vector field.

Definition 4. Let $w : \mathbb{R}^d \times \mathbb{R}^+ \rightarrow \mathbb{R}^d$ be a vector field being uniformly bounded, Lipschitz in space and measurable in time. We define the **flow** associated to the vector field w as the application $(x^0, t) \mapsto \Phi_t^w(x^0)$ such that, for all $x^0 \in \mathbb{R}^d$, $t \mapsto \Phi_t^w(x^0)$ is the solution to the Cauchy problem

$$\begin{cases} \dot{x}(t) = w(x(t), t) \text{ for a.e. } t \geq 0, \\ x(0) = x^0. \end{cases}$$

The following property of the flow will be useful all along the present paper:

Property 2.4 (see [27]). *Let $\mu, \nu \in \mathcal{P}_c(\mathbb{R}^d)$ and $w : \mathbb{R}^d \times \mathbb{R} \rightarrow \mathbb{R}^d$ be a vector field uniformly bounded, Lipschitz in space and measurable in time with a Lipschitz constant equal to L . For each $t \in \mathbb{R}$ and $p \in [1, \infty)$, it holds*

$$W_p(\Phi_t^w \# \mu, \Phi_t^w \# \nu) \leq e^{\frac{(p+1)}{p} L |t|} W_p(\mu, \nu). \quad (8)$$

We now recall a standard result for the continuity equation:

Theorem 2.1 (see [31, Th. 5.34]). *Let $T > 0$, $\mu^0 \in \mathcal{P}_c(\mathbb{R}^d)$ and w a vector field uniformly bounded, Lipschitz in space and measurable in time. Then, System (7) admits a unique solution μ in $\mathcal{C}^0([0, T]; \mathcal{P}_c(\mathbb{R}^d))$, where $\mathcal{P}_c(\mathbb{R}^d)$ is equipped with the weak topology. Moreover:*

- (i) *If $\mu^0 \in \mathcal{P}_c^{ac}(\mathbb{R}^d)$, then the solution μ to (7) belongs to $\mathcal{C}^0([0, T]; \mathcal{P}_c^{ac}(\mathbb{R}^d))$.*
- (ii) *We have $\mu(t) = \Phi_t^w \# \mu^0$ for all $t \in [0, T]$.*

We now recall the precise notions of approximate controllability and exact controllability for System (1):

Definition 5. We say that:

- System (1) is **approximately controllable** from μ^0 to μ^1 on the time interval $[0, T]$ if for each $\varepsilon > 0$ there exists a control $\mathbb{1}_\omega u$ such that the corresponding solutions μ to System (1) satisfies

$$W_p(\mu^1, \mu(T)) \leq \varepsilon. \quad (9)$$

- System (1) is **exactly controllable** from μ^0 to μ^1 on the time interval $[0, T]$ if there exists a control $\mathbb{1}_\omega u$ such that the corresponding solution to System (1) is equal to μ^1 at time T .

It is interesting to remark that, by using properties (6) of the Wasserstein distance, estimate (9) can be replaced by:

$$W_1(\mu^1, \mu(T)) \leq \varepsilon.$$

Thus, in this work, we study approximate controllability by considering the distance W_1 only.

3 Approximate controllability with a localized Lipschitz control

In this section, we study approximate controllability of System (1) with localized Lipschitz controls. More precisely, in Sections 3.1, we consider the case where the open connected control subset ω contains the support of both μ^0 and μ^1 . We then prove Theorem 1.1 in Section 3.2.

3.1 Approximate controllability with a Lipschitz control

In this section, we prove approximate controllability of System (1) with a Lipschitz control, when the open connected control subset ω contains the support of both μ^0 and μ^1 . Without loss of generality, we can assume that the vector field v is identically zero by replacing u with $u - v$ in the control set ω . We then study approximate controllability of system

$$\begin{cases} \partial_t \mu + \operatorname{div}(u\mu) = 0 & \text{in } \mathbb{R}^d \times \mathbb{R}^+, \\ \mu(0) = \mu^0 & \text{in } \mathbb{R}^d. \end{cases} \quad (10)$$

Proposition 3.1. *Let $\mu^0, \mu^1 \in \mathcal{P}_c^{ac}(\mathbb{R}^d)$ compactly supported in ω . Then, for all $T > 0$, System (10) is approximately controllable on the time interval $[0, T]$ from μ^0 to μ^1 with a control $u : \mathbb{R}^d \times \mathbb{R}^+ \rightarrow \mathbb{R}^d$ uniformly bounded, Lipschitz in space and measurable in time. Moreover, the solution μ to System (10) satisfies*

$$\operatorname{supp}(\mu(t)) \subset \omega,$$

for all $t \in [0, T]$.

Proof of Proposition 3.1. We assume that $d := 2$, but the reader will see that the proof can be clearly adapted to dimension one or to any other space dimension. In view to simplify the computations, we suppose that $T := 1$ and $\operatorname{supp}(\mu^i) \subset (0, 1)^2 \subset \subset \omega$ for $i = 1, 2$. We first partition $(0, 1)^2$. Let $n \in \mathbb{N}^*$, consider $a_0 := 0$, $b_0 := 0$ and define

the points a_i, b_i for all $i \in \{1, \dots, n\}$ by induction as follows: suppose that for a given $i \in \{0, \dots, n-1\}$ the points a_i and b_i are defined, then the points a_{i+1} and b_{i+1} are the smallest values such that

$$\int_{(a_i, a_{i+1}) \times \mathbb{R}} d\mu^0 = \frac{1}{n} \quad \text{and} \quad \int_{(b_i, b_{i+1}) \times \mathbb{R}} d\mu^1 = \frac{1}{n}.$$

Again, for each $i \in \{0, \dots, n-1\}$, we consider $a_{i,0} := 0, b_{i,0} := 0$ and supposing that for a given $j \in \{0, \dots, n-1\}$ the points $a_{i,j}$ and $b_{i,j}$ are already defined, $a_{i,j+1}$ and $b_{i,j+1}$ are the smallest values such that

$$\int_{A_{ij}} d\mu^0 = \frac{1}{n^2} \quad \text{and} \quad \int_{B_{ij}} d\mu^1 = \frac{1}{n^2},$$

where $A_{ij} := (a_i, a_{i+1}) \times (a_{ij}, a_{i(j+1)})$ and $B_{ij} := (b_i, b_{i+1}) \times (b_{ij}, b_{i(j+1)})$. Since μ^0 and μ^1 have a mass equal to 1 and are supported in $(0, 1)^2$, then $a_n, b_n \leq 1$ and $a_{i,n}, b_{i,n} \leq 1$ for all $i \in \{0, \dots, n-1\}$. We give in Figure 2 an example of such partition.

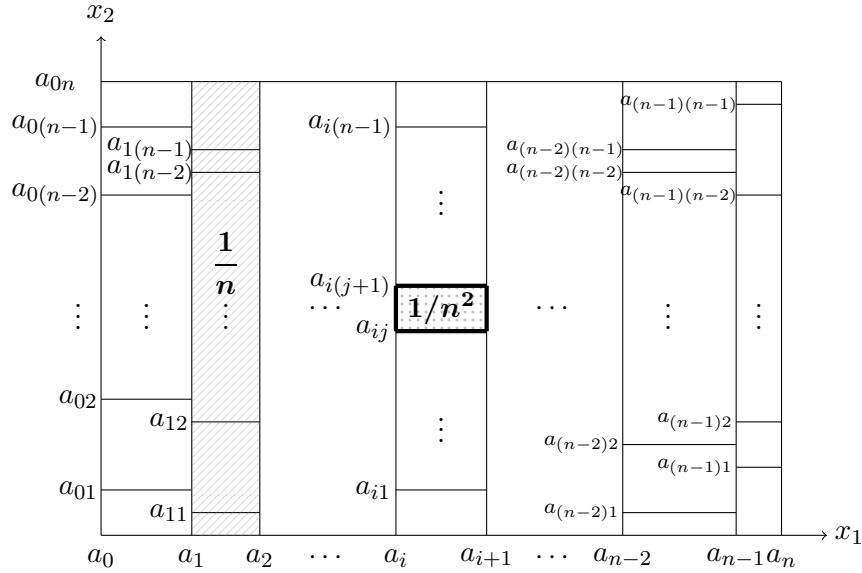


Figure 2: Example of a partition for μ^0 .

If one aims to define a vector field sending each A_{ij} to B_{ij} , then some shear stress is naturally introduced, as described in Remark 2. To overcome this problem, we first define sets $\tilde{A}_{ij} \subset\subset A_{ij}$ and $\tilde{B}_{ij} \subset\subset B_{ij}$ for all $i, j \in \{0, \dots, n-1\}$. We then send the mass of μ^0 from each \tilde{A}_{ij} to \tilde{B}_{ij} , while we do not control the mass contained in $A_{ij} \setminus \tilde{A}_{ij}$. More precisely, for all $i, j \in \{0, \dots, n-1\}$, we define, as in Figure 3, $a_i^-, a_i^+, a_{ij}^-, a_{ij}^+$ the smallest values such that

$$\int_{(a_i^-, a_i^+) \times (a_{ij}^-, a_{i(j+1)}^-)} d\mu^0 = \int_{(a_i^+, a_{i+1}^+) \times (a_{ij}^+, a_{i(j+1)}^+)} d\mu^0 = \frac{1}{n^3}$$

and

$$\int_{(a_i^-, a_i^+) \times (a_{ij}^-, a_{i(j+1)}^-)} d\mu^0 = \int_{(a_i^-, a_i^+) \times (a_{ij}^+, a_{i(j+1)}^+)} d\mu^0 = \frac{1}{n} \times \left(\frac{1}{n^2} - \frac{2}{n^3} \right).$$

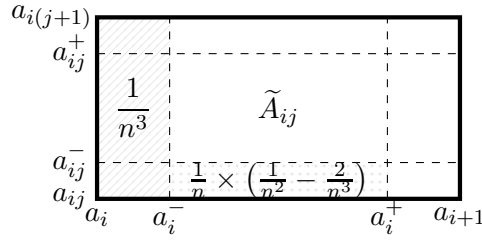


Figure 3: Example of cell.

We similarly define b_i^+ , b_i^- , b_{ij}^+ , b_{ij}^- and finally define

$$\tilde{A}_{ij} := (a_i^-, a_i^+) \times (a_{ij}^-, a_{ij}^+) \text{ and } \tilde{B}_{ij} := (b_i^-, b_i^+) \times (b_{ij}^-, b_{ij}^+).$$

The goal is to build a solution to System (10) such that the corresponding flow Φ_t^u satisfies

$$\Phi_T^u(\tilde{A}_{ij}) = \tilde{B}_{ij}, \quad (11)$$

for all $i, j \in \{0, \dots, n-1\}$. We observe that we do not take into account the displacement of the mass contained in $A_{ij} \setminus \tilde{A}_{ij}$. We will show that the mass of the corresponding term tends to zero when n goes to infinity. The rest of the proof is divided into two steps. In a first step, we build a flow satisfying (11), then the corresponding vector field. In a second step, we compute the Wasserstein distance between μ^1 and $\mu(T)$, showing that it converges to zero when n goes to infinity.

Step 1: We first build a flow satisfying (11). We recall that $T := 1$. For each $i \in \{0, \dots, n-1\}$, we denote by c_i^- and c_i^+ the linear functions equal to a_i^- and a_i^+ at time $t = 0$ and equal to b_i^- and b_i^+ at time $t = T = 1$, respectively, *i.e.* the functions defined for all $t \in [0, T]$ by:

$$c_i^-(t) = (b_i^- - a_i^-)t + a_i^- \quad \text{and} \quad c_i^+(t) = (b_i^+ - a_i^+)t + a_i^+.$$

Similarly, for all $i, j \in \{0, \dots, n-1\}$, we denote by c_{ij}^- and c_{ij}^+ the linear functions equal to a_{ij}^- and a_{ij}^+ at time $t = 0$ and equal to b_{ij}^- and b_{ij}^+ at time $t = T = 1$, respectively, *i.e.* the functions defined for all $t \in [0, T]$ by:

$$c_{ij}^-(t) = (b_{ij}^- - a_{ij}^-)t + a_{ij}^- \quad \text{and} \quad c_{ij}^+(t) = (b_{ij}^+ - a_{ij}^+)t + a_{ij}^+.$$

Consider the application being the following linear combination of c_i^- , c_i^+ and c_{ij}^- , c_{ij}^+ on \tilde{A}_{ij} , *i.e.*

$$x(x^0, t) := \begin{pmatrix} x_1(x^0, t) \\ x_2(x^0, t) \end{pmatrix} = \begin{pmatrix} \frac{a_i^+ - x_1^0}{a_i^+ - a_i^-} c_i^-(t) + \frac{x_1^0 - a_i^-}{a_i^+ - a_i^-} c_i^+(t) \\ \frac{a_{ij}^+ - x_2^0}{a_{ij}^+ - a_{ij}^-} c_{ij}^-(t) + \frac{x_2^0 - a_{ij}^-}{a_{ij}^+ - a_{ij}^-} c_{ij}^+(t) \end{pmatrix}, \quad (12)$$

where $x^0 = (x_1^0, x_2^0) \in \tilde{A}_{ij}$. Let us prove that an extension of the application $(x^0, t) \mapsto x(x^0, t)$ is a flow associated to a vector field u . After some computations, we obtain

$$\begin{cases} \frac{dx_1}{dt}(x^0, t) = \alpha_i(t)x_1(x^0, t) + \beta_i(t) & \forall t \in [0, T], \\ \frac{dx_2}{dt}(x^0, t) = \alpha_{ij}(t)x_2(x^0, t) + \beta_{ij}(t) & \forall t \in [0, T], \end{cases}$$

where for all $t \in [0, T]$,

$$\begin{cases} \alpha_i(t) = \frac{b_i^+ - b_i^- + a_i^- - a_i^+}{c_i^+(t) - c_i^-(t)}, & \beta_i(t) = \frac{a_i^+ b_i^- - a_i^- b_i^+}{c_i^+(t) - c_i^-(t)}, \\ \alpha_{ij}(t) = \frac{b_{ij}^+ - b_{ij}^- + a_{ij}^- - a_{ij}^+}{c_{ij}^+(t) - c_{ij}^-(t)}, & \beta_{ij}(t) = \frac{a_{ij}^+ b_{ij}^- - a_{ij}^- b_{ij}^+}{c_{ij}^+(t) - c_{ij}^-(t)}. \end{cases}$$

The last quantities are well defined since for all $i, j \in \{0, \dots, n-1\}$ and $t \in [0, T]$

$$\begin{cases} |c_i^+(t) - c_i^-(t)| \geq \max\{|a_i^+ - a_i^-|, |b_i^+ - b_i^-|\}, \\ |c_{ij}^+(t) - c_{ij}^-(t)| \geq \max\{|a_{ij}^+ - a_{ij}^-|, |b_{ij}^+ - b_{ij}^-|\}. \end{cases}$$

For all $t \in [0, T]$, consider the set

$$\tilde{C}_{ij}(t) := (c_i^-(t), c_i^+(t)) \times (c_{ij}^-(t), c_{ij}^+(t)).$$

We remark that $\tilde{C}_{ij}(0) = \tilde{A}_{ij}$ and $\tilde{C}_{ij}(T) = \tilde{B}_{ij}$. On

$$\tilde{C}_{ij} := \{(x, t) : t \in [0, T], x \in \tilde{C}_{ij}(t)\},$$

we then define the vector field u by

$$\begin{cases} u_1(x, t) = \alpha_i(t)x_1 + \beta_i(t), \\ u_2(x, t) = \alpha_{ij}(t)x_2 + \beta_{ij}(t), \end{cases}$$

for all $(x, t) \in \tilde{C}_{ij}$ ($x = (x_1, x_2)$). We extend u by a uniform bounded C^∞ function outside $\cup_{ij} \tilde{C}_{ij}$, then u is a C^∞ function and it satisfies $\text{supp}(u) \subset \omega$. Then, System (1) admits an unique solution and the flow on \tilde{C}_{ij} is given by (12).

Step 2: We now prove that the refinement of the grid provides convergence to the target μ^1 , *i.e.*

$$W_1(\mu^1, \mu(T)) \xrightarrow{n \rightarrow \infty} 0.$$

We remark that

$$\int_{\tilde{B}_{ij}} d\mu(T) = \int_{\tilde{B}_{ij}} d\mu^1 = \frac{1}{n^2} - \frac{2}{n^3} - \frac{2}{n} \left(\frac{1}{n^2} - \frac{2}{n^3} \right) = \frac{(n-2)^2}{n^4}.$$

Hence, by defining

$$R := (0, 1)^2 \setminus \bigcup_{ij} \tilde{B}_{ij},$$

we also have

$$\int_R d\mu(T) = \int_R d\mu^1 = 1 - \frac{(n-2)^2}{n^2}.$$

Using (5), it holds

$$W_1(\mu^1, \mu(T)) \leq \sum_{i,j=1}^n W_1(\mu^1_{|\tilde{B}_{ij}}, \mu(T)_{|\tilde{B}_{ij}}) + W_1(\mu^1_{|R}, \mu(T)_{|R}). \quad (13)$$

We now estimate each term in the right-hand side of (13). Since we deal with AC measures, using Properties 2.2, there exist measurable maps $\gamma_{ij} : \mathbb{R}^2 \rightarrow \mathbb{R}^2$, for all $i, j \in \{0, \dots, n-1\}$, and $\bar{\gamma} : \mathbb{R}^2 \rightarrow \mathbb{R}^2$ such that

$$\begin{cases} \gamma_{ij} \# (\mu^1_{|\tilde{B}_{ij}}) = \mu(T)_{|\tilde{B}_{ij}}, \\ W_1(\mu^1_{|\tilde{B}_{ij}}, \mu(T)_{|\tilde{B}_{ij}}) \\ \quad = \int_{\tilde{B}_{ij}} |x - \gamma_{ij}(x)| d\mu^1(x) \end{cases} \quad \text{and} \quad \begin{cases} \bar{\gamma} \# (\mu^1_{|R}) = \mu(T)_{|R}, \\ W_1(\mu^1_{|R}, \mu(T)_{|R}) \\ \quad = \int_R |x - \bar{\gamma}(x)| d\mu^1(x). \end{cases}$$

In the first term in the right hand side of (13), observe that γ_{ij} moves masses inside \tilde{B}_{ij} only. Thus, for all $i, j \in \{0, \dots, n-1\}$, using the triangle inequality,

$$\begin{aligned} W_1(\mu^1_{|\tilde{B}_{ij}}, \mu(T)_{|\tilde{B}_{ij}}) &= \int_{\tilde{B}_{ij}} |x - \gamma_{ij}(x)| d\mu^1(x) \\ &\leq [(b_i^+ - b_i^-) + (b_{ij}^+ - b_{ij}^-)] \int_{\tilde{B}_{ij}} d\mu^1(x) \leq (b_i^+ - b_i^- + b_{ij}^+ - b_{ij}^-) \frac{(n-2)^2}{n^4}. \end{aligned} \quad (14)$$

For the second term in the right-hand side of (13), observe that $\bar{\gamma}$ moves a small mass in the bounded set $(0, 1)$. Thus it holds

$$W_1(\mu^1_{|R}, \mu(T)_{|R}) = \int_R |x - \bar{\gamma}(x)| d\mu^1(x) \leq 2 \left(1 - \frac{(n-2)^2}{n^2}\right) = 8 \frac{n-1}{n^2}. \quad (15)$$

Combining (13), (14) and (15), we obtain

$$\begin{aligned} W_1(\mu^1, \mu(T)) &\leq \left(\sum_{i,j=1}^n (b_i^+ - b_i^- + b_{ij}^+ - b_{ij}^-) \frac{(n-2)^2}{n^4} \right) + 8 \frac{n-1}{n^2} \\ &\leq 2n \frac{(n-2)^2}{n^4} + 8 \frac{n-1}{n^2} \xrightarrow{n \rightarrow \infty} 0. \end{aligned}$$

□

Remark 2. It is not possible in general to build a Lipschitz vector field sending directly each A_{ij} to B_{ij} using the strategy developed in the proof of Proposition 3.1. Indeed, we would obtain discontinuous velocities on the lines c_i . Figure 4 illustrates this phenomenon in the case $n = 2$.

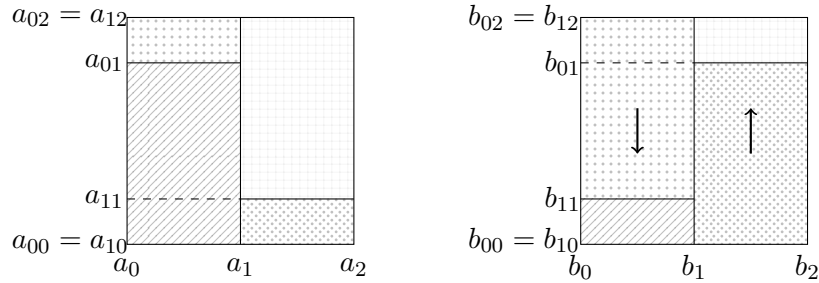


Figure 4: Shear stress (left: μ^0 , right: μ^1)

3.2 Approximate controllability with a localized regular control

This section is devoted to prove Theorem 1.1: we aim to prove approximate controllability of System (1) with a Lipschitz localized control. This means that we remove the constraints $\text{supp}(\mu^0) \subset \omega$, $\text{supp}(\mu^1) \subset \omega$ and $v := 0$, that we used in Section 3.1. On the other side, we impose Condition 1.1.

Before the main proof, we need three useful results. First of all, we give a consequence of Condition 1.1:

Condition 3.1. There exist two real numbers T_0^* , $T_1^* > 0$ and a nonempty open set $\omega_0 \subset\subset \omega$ such that

- (i) For each $x^0 \in \text{supp}(\mu^0)$, there exists $t^0 \in [0, T_0^*]$ such that $\Phi_{t^0}^v(x^0) \in \omega_0$, where Φ_t^v is the flow associated to v .
- (ii) For each $x^1 \in \text{supp}(\mu^1)$, there exists $t^1 \in [0, T_1^*]$ such that $\Phi_{-t^1}^v(x^1) \in \omega_0$.

Lemma 3.1. *If Condition 1.1 is satisfied for $\mu^0, \mu^1 \in \mathcal{P}_c(\mathbb{R}^d)$, then Condition 3.1 is satisfied too.*

Proof. We use a compactness argument. Let $\mu^0 \in \mathcal{P}_c(\mathbb{R}^d)$ and assume that Condition 1.1 holds. Let $x^0 \in \text{supp}(\mu^0)$. Using Condition 1.1, there exists $t^0(x^0) > 0$ such that $\Phi_{t^0(x^0)}^v(x^0) \in \omega$. Choose $r(x^0) > 0$ such that $B_{r(x^0)}(\Phi_{t^0(x^0)}^v(x^0)) \subset\subset \omega$, where $B_r(x^0)$ denotes the open ball of radius $r > 0$ centered at point x^0 in \mathbb{R}^d . Such $r(x^0)$ exists, since ω is open. By continuity of the application $x^1 \mapsto \Phi_{t^0(x^0)}^v(x^1)$ (see [7, Th. 2.1.1]), there exists $\hat{r}(x^0)$ such that

$$x^1 \in B_{\hat{r}(x^0)}(x^0) \Rightarrow \Phi_{t^0(x^0)}^v(x^1) \in B_{r(x^0)}(\Phi_{t^0(x^0)}^v(x^0)).$$

Since μ^0 is compactly supported, we can find a set $\{x_1^0, \dots, x_{N_0}^0\} \subset \text{supp}(\mu^0)$ such that

$$\text{supp}(\mu^0) \subset \bigcup_{i=1}^{N_0} B_{\hat{r}(x_i^0)}(x_i^0).$$

We similarly build a set $\{x_1^1, \dots, x_{N_1}^1\} \subset \text{supp}(\mu^1)$. Thus Condition 3.1 is satisfied for

$$T_k^* := \max\{t^k(x_i^k) : i \in \{1, \dots, N_k\}\},$$

with $k = 0, 1$ and

$$\omega_0 := \left(\bigcup_{i=1}^{N_0} B_{r(x_i^0)}(\Phi_{t^0(x_i^0)}^v(x_i^0)) \right) \cup \left(\bigcup_{i=1}^{N_1} B_{r(x_i^1)}(\Phi_{t^1(x_i^1)}^v(x_i^1)) \right) \subset\subset \omega.$$

□

The second useful result is the following proposition, showing that we can store a large part of the mass of μ^0 in ω , under Condition 3.1.

Proposition 3.2. *Let $\mu^0 \in \mathcal{P}_c^{ac}(\mathbb{R}^d)$ satisfying the first item of Condition 3.1. Then, for all $\varepsilon > 0$, there exists a space-dependent vector field $\mathbb{1}_{\omega} u$ Lipschitz and uniformly bounded and a Borel set $A \subset \mathbb{R}^d$ such that*

$$\mu^0(A) = \varepsilon \text{ and } \text{supp}(\Phi_{T_0^*}^{v+\mathbb{1}_{\omega} u} \# \mu_{|A^c}^0) \subset \omega. \quad (16)$$

Proof. For each $k \in \mathbb{N}^*$, we denote by ω_k the closed set defined by

$$\omega_k := \{x^0 \in \mathbb{R}^d : d(x^0, \omega_0^c) \geq 1/k\}$$

and a cutoff function $\theta_k \in \mathcal{C}^\infty(\mathbb{R}^d)$ satisfying

$$\begin{cases} 0 \leq \theta_k \leq 1, \\ \theta_k = 1 \text{ in } \omega_0^c, \\ \theta_k = 0 \text{ in } \omega_k. \end{cases}$$

For all $x^0 \in \text{supp}(\mu^0)$, we define

$$t_0(x^0) := \inf\{t \in \mathbb{R}^+ : \Phi_t^v(x^0) \in \omega_0\} \text{ and } t_k(x^0) := \inf\{t \in \overline{\mathbb{R}}^+ : \Phi_t^v(x^0) \in \omega_k\}.$$

For all $k \in \mathbb{N}^*$, we consider

$$u_k := (\theta_k - 1)v \quad (17)$$

and

$$S_k := \{x^0 \in \text{supp}(\mu^0) \setminus \omega_0 : \exists s \in (t_0(x^0), t_k(x^0)), \text{ s.t. } \Phi_s^v(x^0) \in \overline{\omega_0^c}\}.$$

The rest of the proof is divided into three steps:

- In Step 1, we prove that the range of the flow associated to x^0 with the control u_k is included in the range of the flow associated to x^0 without control.
- In Step 2, we show that S_k is a Borel set for all $k \in \mathbb{N}^*$.
- In Step 3, we prove that for a K large enough we have

$$\mu^0(\omega \setminus \omega_K) + \mu^0(S_K) \leq \varepsilon. \quad (18)$$

Step 1: Consider the flow $y(t) := \Phi_t^v(x^0)$ associated to x^0 without control, *i.e.* the solution to

$$\begin{cases} \dot{y}(t) = v(y(t)), & t \geq 0, \\ y(0) = x^0 \end{cases}$$

and the flow $z_k(t) := \Phi_t^{v+u_k}(x^0)$ associated to x^0 with the control u_k given in (17), *i.e.* the solution to

$$\begin{cases} \dot{z}_k(t) = (v + u_k)(z_k(t)) = \theta_k(z_k(t)) \times v(z_k(t)), & t \geq 0, \\ z_k(0) = x^0. \end{cases} \quad (19)$$

Consider the solution γ_k to the following system

$$\begin{cases} \dot{\gamma}_k(t) = \theta_k(y(\gamma_k(t))), & t \geq 0, \\ \gamma_k(0) = 0. \end{cases} \quad (20)$$

Since θ_k and y are Lipschitz, then System (20) admits a solution defined for all times. We remark that $\xi_k := y \circ \gamma_k$ is solution to System (19). Indeed, for all $t \geq 0$ it holds

$$\begin{cases} \dot{\xi}_k(t) = \dot{\gamma}_k(t) \times \dot{y}(\gamma_k(t)) = \theta_k(\xi_k(t)) \times v(\xi_k(t)), & t \geq 0, \\ \xi_k(0) = y(\gamma_k(0)) = y(0). \end{cases}$$

By uniqueness of the solution to System (19), we obtain

$$y(\gamma_k(t)) = z_k(t) \text{ for all } t \geq 0.$$

Using the fact that $0 \leq \theta \leq 1$ and the definition of γ_k , we have

$$\begin{cases} \gamma_k \text{ increasing,} \\ \gamma_k(t) \leq t & \forall t \in [0, t_k(x^0)], \\ \gamma_k(t) \leq t_k(x^0) & \forall t \geq t_k(x^0). \end{cases}$$

We deduce that, for all $x^0 \in \text{supp}(\mu^0)$, it holds

$$\{z_k(t) : t \geq 0\} \subset \{y(s) : s \in [0, t_k(x^0)]\}.$$

Step 2: We now prove that S_k is a Borel set by showing that the set

$$R_k := \{x^0 \in \mathbb{R}^d : t_0(x^0) < \infty \text{ and } \exists s \in (t_0(x^0), t_k(x^0)) \text{ s.t. } \Phi_s^v(x^0) \in \overline{\omega_0^c}\}$$

is open. Let $k \in \mathbb{N}^*$, x^0 be an element of R_k and search $r(x^0) > 0$ such that $B_{r(x^0)}(x^0) \subset R_k$. There exists $s \in (t_0(x^0), t_k(x^0))$ such that $\Phi_s^v(x^0) \in \overline{\omega_0^c}$. Since $\overline{\omega_0^c}$ is open, for a $\beta > 0$, we have $B_\beta(\Phi_s^v(x^0)) \subset \overline{\omega_0^c}$. By continuity of the application $x^1 \mapsto \Phi_s^v(x^1)$, there exists $r(x^0) > 0$ such that

$$x^1 \in B_{r(x^0)}(x^0) \Rightarrow \Phi_s^v(x^1) \in B_\beta(\Phi_s^v(x^0)).$$

Thus, for all $k \in \mathbb{N}^*$, R_k is open and S_k is a Borel set.

Step 3: We now prove that (18) holds for a K large enough. Since we deal with AC measure, there exists $K_0 \in \mathbb{N}^*$ such that for all $k \geq K_0$

$$\mu^0(\omega_0 \setminus \omega_k) \leq \varepsilon/2.$$

Argue now by contradiction to prove that there exists $K_1 \geq K_0$ such that

$$\mu^0(S_{K_1}) \leq \varepsilon/2.$$

Assume that $\mu^0(S_k) > \varepsilon/2$ for all $k \geq K_0$. Using the inclusion $S_{k+1} \subset S_k$, we deduce that

$$\mu^0\left(\bigcap_{k \in \mathbb{N}^*} S_k\right) \geq \varepsilon/2.$$

Since μ^0 is absolute continuous with respect to λ (the Lebesgue measure), there exists $\alpha > 0$ such that

$$\lambda\left(\bigcap_{k \in \mathbb{N}^*} S_k\right) \geq \alpha.$$

We deduce that the intersection of the set S_k is nonempty. Let $\overline{x}^0 \in \text{supp}(\mu^0) \setminus \overline{\omega_0}$ be an element of this intersection. By definition of S_k , for all $k \geq K_0$, there exists s_k satisfying

$$\begin{cases} s_k \in (t_0(\overline{x}^0), t_k(\overline{x}^0)), \\ \Phi_{s_k}^v(\overline{x}^0) \in \overline{\omega_0^c}. \end{cases} \quad (21)$$

Moreover, the convergence of $t_k(\overline{x}^0)$ to $t_0(\overline{x}^0)$, implies that

$$s_k \rightarrow t_0(\overline{x}^0). \quad (22)$$

Using the continuity of $x^1 \mapsto \Phi_t^v(x^1)$ and the definition of $t_0(x^0)$, there exists $\beta > 0$ such that

$$\Phi_t^v(\overline{x}^0) \in \omega_0 \text{ for all } t \in (t_0, t_0 + \beta). \quad (23)$$

We deduce that (23) contradicts (21) and (22). Thus there exists $K \in \mathbb{N}^*$ such that

$$\mu^0(S_K) + \mu^0(\omega \setminus \omega_K) \leq \varepsilon.$$

Since we deal with AC measures, we add a Borel set to have the equality in (16), *i.e.* there exists a Borel set S such that

$$\mu^0(S_K \cup \omega \setminus \omega_K \cup S) = \varepsilon.$$

We conclude that, for u defined by

$$u(t) := u^1 := u_K \text{ for all } t \in [0, T_0^*],$$

and $A := S_K \cup \omega \setminus \omega_K \cup S$, Properties (16) are satisfied. \square

The third useful result for the proof of Theorem 1.1 allows to approximately steer a measure contained in ω to a measure contained in an open hypercube $S \subset\subset \omega$.

Proposition 3.3. *Let $\mu^0 \in \mathcal{P}_c^{ac}(\mathbb{R}^d)$ satisfying $\text{supp}(\mu^0) \subset \omega$. Define an open hypercube S strictly included in $\omega \setminus \text{supp}(\mu^0)$ and choose $\delta > 0$. Then, for all $\varepsilon > 0$, there exists a vector field $\mathbb{1}_\omega u$, Lipschitz and uniformly bounded and a Borel set A such that*

$$\mu^0(A) = \varepsilon \text{ and } \text{supp}(\Phi_\delta^{v+\mathbb{1}_\omega u} \# \mu_{|A^c}^0) \subset S.$$

Proof. Consider S_0 a nonempty open set of \mathbb{R}^d of class \mathcal{C}^∞ strictly included in S and ω_1 an open set of \mathbb{R}^d of class \mathcal{C}^∞ satisfying

$$\text{supp}(\mu^0) \cup S \subset\subset \omega_1 \subset\subset \omega.$$

An example is given in Figure 5. From [19, Lemma 1.1, Chap. 1] (see also [14, Lemma

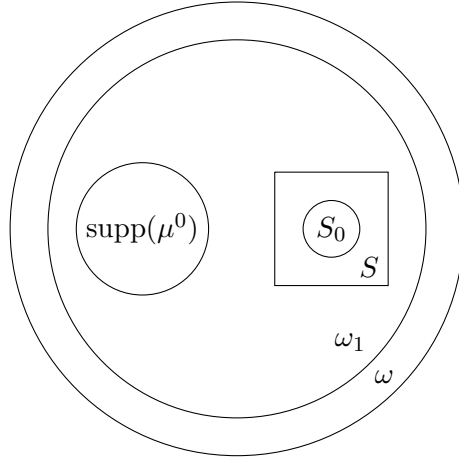


Figure 5: Construction of ω_1

2.68, Chap. 2]), there exists a function $\eta \in \mathcal{C}^2(\overline{\omega_1})$ satisfying

$$\kappa_0 \leq |\nabla \eta| \leq \kappa_1 \text{ in } \omega_1 \setminus S_0, \quad \eta > 0 \text{ in } \omega_1 \quad \text{and} \quad \eta = 0 \text{ on } \partial \omega_1, \quad (24)$$

with $\kappa_0, \kappa_1 > 0$. Let $k \in \mathbb{N}^*$. Consider $u_k : \mathbb{R}^d \rightarrow \mathbb{R}^d$ Lipschitz and uniformly bounded satisfying

$$u_k := \begin{cases} k \nabla \eta - v & \text{in } \omega_1, \\ 0 & \text{in } \omega^c. \end{cases}$$

Let $x^0 \in \text{supp}(\mu^0)$. Consider the flow $z_k(t) = \Phi_t^{v+u_k}(x^0)$ associated to x^0 with the control u_k , *i.e.* the solution to system

$$\begin{cases} \dot{z}_k(t) = v(z_k(t)) + u_k(z_k(t)), & t \geq 0, \\ z_k(0) = x^0. \end{cases} \quad (25)$$

The different conditions in (24) imply that

$$n \cdot \nabla \eta < C < 0 \text{ on } \partial\omega_1, \quad (26)$$

where n represents the outward unit normal to $\partial\omega_1$. Since $\text{supp}(\mu^0) \subset \omega_1$, it holds $z_k(t) \in \omega_1$ for all $t \geq 0$, otherwise, by taking the scalar product of (25) and n on $\partial\omega_1$, we obtain a contradiction with (26).

We now prove that there exists $K(x^0) \in \mathbb{N}^*$ such that for all $k \geq K(x^0)$ there exists $t_k(x^0) \in (0, \delta)$ such that $z_k(t_k(x^0))$ belongs to S_0 . By contradiction, assume that there exists a sequences $\{k_n\}_{n \in \mathbb{N}^*} \subset \mathbb{N}^*$ such that for all $t \in (0, \delta)$

$$z_{k_n}(t) \in S_0^c. \quad (27)$$

Consider the function f_n defined for all $t \in [0, \delta]$ by

$$f_n(t) := k_n \eta(z_{k_n}(t)). \quad (28)$$

Its time derivative is given for all $t \in [0, \delta]$ by

$$\dot{f}_n(t) = k_n \dot{z}_{k_n}(t) \cdot \nabla \eta(z_{k_n}(t)) = k_n^2 |\nabla \eta(z_{k_n}(t))|^2$$

Then, using (27), properties (24) of η and definition (28) of f_n , it holds

$$f_n(\delta) \geq k_n^2 \kappa_0^2 \delta \quad \text{and} \quad f_n(\delta) \leq k_n \|\eta\|_\infty.$$

We observe that the two last inequalities are in contradiction for n large enough. Then there exists $K(x^0) \in \mathbb{N}^*$ such that for all $k \geq K(x^0)$ there exists $t_k(x^0) \in (0, \delta)$ such that $z_k(t_k(x^0))$ belongs to S_0 .

By continuity, there exists $r(x^0) > 0$ such that $\Phi_{t_{K(x^0)}(x^0)}^{v+u_{K(x^0)}}(x^1)$ belongs to S_0 for all $x^1 \in B_{r(x^0)}(x^0)$. Since $v + u_k$ is linear with respect to k in ω_1 , then, using the same argument as in Step 1 of the proof of Proposition 3.2, the range of the flow Φ^{v+u_k} is independent of k . Thus, for all $k \geq K(x^0)$ there exists $t_k^0(x^0) \in (0, \delta)$ such that $\Phi_{t_k^0(x^0)}^{v+u_k}(x^1) \in S_0$ for all $x^1 \in B_{r(x^0)}(x^0)$.

By compactness, there exists $\{x_1^0, \dots, x_{N_0}^0\}$ such that

$$\text{supp}(\mu^0) \subset \bigcup_{i=1}^{N_0} B_{r(x_i^0)}(x_i^0).$$

We deduce that for $K := \max_i \{K(x_i^0)\}$, for all $x^0 \in \text{supp}(\mu^0)$ there exists $t^0(x^0)$ for which $\Phi_{t^0(x^0)}^{v+u_K}(x^0)$ belongs to S_0 . We remark that the first item of Condition 3.1 holds replacing ω , ω_0 and T_0^* by S , S_0 and δ , respectively. We conclude applying Proposition 3.2 replacing ω , ω_0 , T_0^* and v by S , S_0 , δ and $v + u_K$, respectively. \square

Remark 3. An alternative method to prove Proposition 3.3 involves building an explicit flow composed with straight lines as in the proof of Proposition 3.1. However, for such method we need to assume that ω is convex, contrarily to the more general approach developed in the proof of Proposition 3.3.

We now have all the tools to prove Theorem 1.1.

Proof of Theorem 1.1. Consider μ^0, μ^1 satisfying Condition 1.1. By Lemma 3.1, there exist T_0^*, T_1^*, ω_0 for which μ^0, μ^1 satisfy Condition 3.1. Let $\delta, \varepsilon > 0$ and $T := T_0^* + T_1^* + \delta$. We now prove that we can construct a Lipschitz uniformly bounded and control $\mathbb{1}_\omega u$ such that the corresponding solution μ to System (1) satisfies

$$W_1(\mu(T), \mu^1) \leq \varepsilon.$$

Denote by $T_0 := 0, T_1 := T_0^*, T_2 := T_0^* + \delta/3, T_3 := T_0^* + 2\delta/3, T_4 := T_0^* + \delta$ and $T_5 := T_0^* + T_1^* + \delta$. Also fix an open hypercube $S \subset\subset \omega \setminus \omega_0$. There exist $R > 0$ and $\bar{x} \in \mathbb{R}^d$ such that the supports of μ^0 and μ^1 are strictly included in a hypercube with edges of length R . Define

$$\bar{R} := R + T \times \sup_{\mathbb{R}^d} |v|.$$

Applying Proposition 3.2 on $[T_0, T_1] \cup [T_4, T_5]$ and Proposition 3.3 on $[T_1, T_2] \cup [T_3, T_4]$, we can construct some space-dependent controls u^1, u^2, u^4, u^5 Lipschitz and uniformly bounded, with $\text{supp}(u^i) \subset \omega$, and two Borel sets A_0 and A_1 such that

$$\mu^0(A_0) = \mu^1(A_1) = \frac{\varepsilon}{2d\bar{R}}.$$

Moreover, the solution forward in time to

$$\begin{cases} \partial_t \rho_0 + \nabla \cdot ((v + \mathbb{1}_\omega u^1) \rho_0) = 0 & \text{in } \mathbb{R}^d \times [T_0, T_1], \\ \partial_t \rho_0 + \nabla \cdot ((v + \mathbb{1}_\omega u^2) \rho_0) = 0 & \text{in } \mathbb{R}^d \times [T_1, T_2], \\ \rho_0(T_0) = \mu_{|A_0}^0 & \text{in } \mathbb{R}^d \end{cases}$$

and the solution backward in time to

$$\begin{cases} \partial_t \rho_1 + \nabla \cdot ((v + \mathbb{1}_\omega u^5) \rho_1) = 0 & \text{in } \mathbb{R}^d \times [T_4, T_5], \\ \partial_t \rho_1 + \nabla \cdot ((v + \mathbb{1}_\omega u^4) \rho_1) = 0 & \text{in } \mathbb{R}^d \times [T_3, T_4], \\ \rho_1(T_5) = \mu_{|A_1}^1 & \text{in } \mathbb{R}^d \end{cases}$$

satisfy $\text{supp}(\rho_0(T_2)) \subset S$ and $\text{supp}(\rho_1(T_3)) \subset S$.

We remark that $|\rho_0(T_2)| = |\rho_1(T_3)| = 1 - \varepsilon/2d\bar{R}$. We now apply Proposition 3.1 to approximately steer $\rho_0(T_2)$ to $\rho_1(T_3)$ inside S as follows: we find a control u^3 on the time interval $[T_2, T_3]$ satisfying $\text{supp}(u^3) \subset S$ such that the solution ρ to

$$\begin{cases} \partial_t \rho + \nabla \cdot ((v + \mathbb{1}_\omega u^3) \rho) = 0 & \text{in } \mathbb{R}^d \times [T_2, T_3], \\ \rho(T_2) = \rho_0(T_2) & \text{in } \mathbb{R}^d \end{cases}$$

satisfies

$$W_1(\rho(T_3), \rho_1(T_3)) \leq \frac{\varepsilon}{2e^{2L(T_5-T_3)}},$$

where L is the uniform Lipschitz constant for u^4 and u^5 . Thus, denoting by u the concatenation of u^1, u^2, u^3, u^4, u^5 on the time interval $[0, T]$, we approximately steer $\mu_{|A_0}^0$ to $\mu_{|A_1}^1$, since by (8) the solution μ to

$$\begin{cases} \partial_t \mu + \nabla \cdot ((v + \mathbb{1}_\omega u^i) \mu) = 0 & \text{in } \mathbb{R}^d \times [T_{i-1}, T_i], i \in \{1, \dots, 5\}, \\ \mu(0) = \mu_{|A_0}^0 & \text{in } \mathbb{R}^d \end{cases}$$

satisfies

$$W_1(\Phi_T^{v+u} \# \mu_{|A_0}^0, \mu_{|A_1}^1) = W_1(\mu(T_5), \mu_{|A_1}^1) \leq e^{2L(T_5-T_3)} \frac{\varepsilon}{2e^{2L(T_5-T_3)}} = \frac{\varepsilon}{2}. \quad (29)$$

Since we deal with AC measures, using Properties 2.2, there exists a measurable map $\gamma : \mathbb{R}^d \rightarrow \mathbb{R}^d$ such that

$$\begin{cases} \gamma \# \mu_{|A_1}^1 = \Phi_T^{v+u} \# \mu_{|A_0}^0, \\ W_1(\Phi_T^{v+u} \# \mu_{|A_0}^0, \mu_{|A_1}^1) = \int_{\mathbb{R}^d} |x - \gamma(x)| d\mu_{|A_1}^1(x). \end{cases}$$

We deduce that

$$W_1(\Phi_T^{v+u} \# \mu_{|A_0}^0, \mu_{|A_1}^1) = \int_{\mathbb{R}^d} |x - \gamma(x)| d\mu_{|A_1}^1(x) \leq d\bar{R} \times \frac{\varepsilon}{2d\bar{R}} = \frac{\varepsilon}{2}. \quad (30)$$

Inequalities (5), (29) and (30) leads to the conclusion:

$$W_1(\Phi_T^{v+u} \# \mu^0, \mu^1) \leq W_1(\Phi_T^{v+u} \# \mu_{|A_0^c}^0, \mu_{|A_1^c}^1) + W_1(\Phi_T^{v+u} \# \mu_{|A_0}^0, \mu_{|A_1}^1) \leq \varepsilon.$$

□

4 Exact controllability

In this section, we study exact controllability for System (1). In Section 4.1, we show that exact controllability of System (1) does not hold for Lipschitz or BV controls. In Section 4.2, we prove Theorem 1.2, *i.e.* exact controllability of System (1) with a Borel localized control under some geometrical conditions.

4.1 Negative results for exact controllability

In this section, we show that exact controllability does not hold in general for Lipschitz or BV controls. We will see that topological aspects play a crucial role at this level.

a) Non exact controllability with Lipschitz controls

As explained in the introduction, if we impose the classical Carathéodory condition of $\mathbb{1}_\omega u : \mathbb{R}^d \times \mathbb{R}^+ \rightarrow \mathbb{R}^d$ being uniformly bounded, Lipschitz in space and measurable in time, then the flow $\Phi_t^{v+1_\omega u}$ is a homeomorphism (see [7, Th. 2.1.1]). More precisely, the flow and its inverse are locally Lipschitz. This implies that the support of μ^0 and $\mu(T)$ are homeomorph. Thus, if the support of μ^0 and μ^1 are not homeomorph, then exact controllability does not hold with Lipschitz controls. In particular, we cannot steer a measure which support is connected to a measure which support is composed of two connected components with Lipschitz controls and conversely.

b) Non exact controllability with BV controls

To hope to obtain exact controllability of System (1), it is then necessary to search for a control with less regularity. A weaker condition on the regularity of the vector field for the well-posedness of System (1) has been given in [3]. Consider an initial data μ^0 in $L^\infty \cap \mathcal{P}_c^{ac}(\mathbb{R}^d)$, a target μ^1 in $L^\infty \cap \mathcal{P}_c^{ac}(\mathbb{R}^d)$ and a vector field u satisfying:

- (i) $u(\cdot, t) \in BV_{loc}(\mathbb{R}^d)$ for a.e. $t \in (0, T)$.

(ii) For all $R > 0$,

$$\begin{aligned} \sup_{\mathbb{R}^d} |u| + \int_0^T \|[\operatorname{div} u]^- \|_{L^\infty(B_R(0))} + \int_0^T \|u\|_{BV(B_R(0))} dt \\ + \int_0^T \int_{B_R(0)} |\operatorname{div} u(t)| dx dt < +\infty. \end{aligned}$$

Under these assumptions, there exists a unique solution to System (1) in the sense of [3] (see [3, Theorems 4.1 and 6.2]).

We now give an example of non exact controllability of System (1) with a velocity field satisfying these assumptions and $\omega = \mathbb{R}^d$. Consider

$$\mu^0 := \frac{1}{2} \mathbb{1}_{(-1,0)}(x) dx + \frac{1}{2} \mathbb{1}_{(1,2)}(x) dx \quad \text{and} \quad \mu^1 := \frac{1}{2} \mathbb{1}_{(-1,1)} dx.$$

Suppose that there exists u satisfying (i), (ii) and $\Phi_T^u \# \mu^0 = \mu^1$. The solution to System (1) is then unique. Consider y and z solutions to

$$\begin{cases} \dot{y}(t) = u(y(t), t) \text{ for a.e. } t \in (0, T), \\ y(0) = 0 \end{cases} \quad \text{and} \quad \begin{cases} \dot{z}(t) = u(z(t), t) \text{ for a.e. } t \in (0, T), \\ z(0) = 1. \end{cases}$$

Since $\Phi_T^u((-1, 0)) = (-1, 0)$ and $\Phi_T^u((1, 2)) = (0, 1)$, it holds $y(T) = z(T)$. Thus there exists an infimum time $t_1 \in [0, T]$ such that $y(t_1) = z(t_1)$. Since $\Phi_T^u \# \mu^0 = \mu^1$, we obtain a contradiction with (ii) by observing that there exists $t_0 \in (0, t_1)$ such that

$$\int_0^T \|[\partial_x u(t)]^-\|_\infty dt \geq \int_{t_0}^{t_1} \left| \frac{u(y(t), t) - u(z(t), t)}{y(t) - z(t)} \right| dt = \int_{t_0}^{t_1} \left| \frac{\dot{y}(t) - \dot{z}(t)}{y(t) - z(t)} \right| dt = +\infty.$$

Thus, we cannot steer a measure which support is composed of two connected components to a measure which support is connected with BV controls satisfying (i) and (ii), hence general results on exact controllability cannot hold. However, the inverse is possible. For example, if we denote by

$$\mu^0 := \mathbb{1}_{(-1,1)} dx \quad \text{and} \quad u(x) := \begin{cases} \sqrt{x} & \text{if } x \geq 0, \\ 0 & \text{otherwise,} \end{cases}$$

then u satisfies (i) and (ii) and the unique solution μ to System (1) is given by

$$\mu(t) = \mathbb{1}_{(-1,0)}(x) dx + \left(1 - \frac{t}{2\sqrt{x}}\right) \times \mathbb{1}_{(\frac{t^2}{4}, (\frac{t}{2}+1)^2)}(x) dx.$$

4.2 Exact controllability with Borel controls

In this section, we prove Theorem 1.2, *i.e.* exact controllability of System (1) in the following sense: there exists a couple $(\mathbb{1}_\omega u, \mu)$ solution to System (1) satisfying $\mu(T) = \mu^1$. Before proving Theorem 1.2, we need three useful results.

The first one is the following proposition, showing that we can store the whole mass of μ^0 in ω , under Condition 3.1. It is the analogue of Proposition 3.2. In this case, we control the whole mass, but we do not have necessarily uniqueness of the solution to System (1).

Proposition 4.1. *Let $\mu^0 \in \mathcal{P}_c(\mathbb{R}^d)$ satisfying the first item of Condition 3.1. Then there exists a couple $(\mathbb{1}_\omega u, \mu)$ composed of a Borel vector field $\mathbb{1}_\omega u : \mathbb{R}^d \times \mathbb{R}^+ \rightarrow \mathbb{R}^d$ and a time-evolving measure μ being weak solution to System (1) and satisfying*

$$\operatorname{supp}(\mu(T_0^*)) \subset \omega.$$

Proof. For each $x^0 \in \mathbb{R}^d$, we denote by

$$\tilde{t}^0(x^0) := \inf\{t \geq 0 : \Phi_t^v(x^0) \in \bar{\omega}_0\}$$

and consider the application $\Psi_t(x^0)$ defined for all $t \geq 0$ by

$$\Psi_t(x^0) = \begin{cases} \Phi_t^v(x^0) & \text{if } t \leq \tilde{t}^0(x^0), \\ \Phi_{\tilde{t}^0(x^0)}^v(x^0) & \text{otherwise.} \end{cases}$$

For all $t \geq 0$, the application Ψ_t is a Borel map. Consider μ defined for all $t \geq 0$ by

$$\mu(t) := \Psi_t \# \mu^0.$$

We remark that, for all $t, s \in [0, T_0^*]$ such that $t \geq s$,

$$\mu(t) = \Psi_{t-s} \# \mu(s). \quad (31)$$

Since $\Phi^v(x^0)$ is Lipschitz, for all $x^0 \in \mathbb{R}^d$ and $t \in [0, T_0^*]$, it holds

$$|\Psi_t(x^0) - x^0| \leq C \min\{t, \tilde{t}^0(x^0)\} \leq Ct. \quad (32)$$

Combining (31) and (32), we deduce for all $t, s \in [0, T_0^*]$ with $s \leq t$

$$W_2^2(\mu(s), \mu(t)) \leq \int_{\mathbb{R}^d} |\Psi_{t-s}(x) - x|^2 d\mu(s) \leq \sup_{x \in \mathbb{R}^d} |\Psi_{t-s}(x) - x|^2 \leq C|t - s|^2.$$

We deduce that the metric derivative $|\mu'|$ of μ defined for all $t \in [0, T_0^*]$ by

$$|\mu'| (t) := \lim_{s \rightarrow t} \frac{W_2(\mu(t), \mu(s))}{|t - s|} \quad (33)$$

is uniformly bounded on $[0, T_0^*]$. Then μ is an absolute continuous curve on $\mathcal{P}_c(\mathbb{R}^d)$ (see [4, Def. 1.1.1]). Using [4, Th. 8.3.1], there exists a Borel vector $w : \mathbb{R}^d \times (0, T_0^*) \rightarrow \mathbb{R}^d$ satisfying

$$\|w(t)\|_{L^2(\mu(t); \mathbb{R}^d)} \leq |\mu'| (t) \text{ a.e. } t \in [0, T_0^*]$$

and the couple (w, μ) is a weak solution to

$$\begin{cases} \partial_t \mu + \nabla \cdot (w\mu) = 0 & \text{in } \mathbb{R}^d \times [0, T_0^*], \\ \mu(0) = \mu^0 & \text{in } \mathbb{R}^d. \end{cases} \quad (34)$$

Moreover, for all $t \in [0, T_0^*]$, it holds

$$w(t) \in \text{Tan}_{\mu(t)}(\mathcal{P}_c(\mathbb{R}^d)) := \overline{\{\nabla \varphi : \varphi \in \mathcal{C}_c^\infty(\mathbb{R}^d)\}}^{L^2(\mu(t); \mathbb{R}^d)}.$$

Consider an open set ω_1 of class \mathcal{C}^∞ satisfying $\omega_0 \subset\subset \omega_1 \subset\subset \omega$. We now prove that $w(t)$ coincides with $v(t)$ in $\text{supp}(\mu(t)) \setminus \bar{\omega}_1$ a.e. $t \in [0, T_0^*]$, i.e. we can choose $u = 0$ outside ω . Fix $t \in [0, T_0^*]$ and consider $x \in \text{supp}(\mu(t)) \cap \omega_1^c$. There necessarily exists $x^0 \in \text{supp}(\mu^0)$ such that $\Phi_t^v(x^0) = x$, otherwise $x \in \partial\omega_0$. Moreover for a $B := B_r(x^0)$ with $r > 0$ $\Phi_s^v(B) \subset\subset \omega_0^c$ for all $s \in [0, t]$, otherwise there exists $s \in [0, t]$ for which $\Phi_s^v(x^0) \in \partial\omega_0$. Thus

$$\Phi_t^v = \Psi_t \text{ in } B. \quad (35)$$

We denote by $A := \Phi_t^v(B)$. We have now prove that

$$\Psi_t^{-1}(A) = (\Phi_t^v)^{-1}(A). \quad (36)$$

Consider $x \in (\Phi_t^v)^{-1}(A)$. Equality (35) implies $\Phi_t^v(x) = \Psi_t(x)$. Then $x \in \Psi_t^{-1}(A)$. Consider now $x \in \Psi_t^{-1}(A)$, which means $\Psi_t(x) \in A$. Using the fact that $A \cap \overline{\omega}_0 \neq \emptyset$, $t < \tilde{x}^0(x)$. Then $\Psi_t(x) = \Phi_t^v(x)$ and $x \in (\Phi_t^v)^{-1}(A)$. Thus (36) holds. By definition of the push forward,

$$\mu|_A(t) = \Psi_t \# (\mu|_{\Psi_t^{-1}(A)}^0) \text{ and } (\Phi_t^v \# \mu^0)|_A = \Phi_t^v \# (\mu|_{\Phi_t^{-1}(A)}^0).$$

Since $\Psi_t = \Phi_t^v$ on the set $B = (\Phi_t^v)^{-1}(A) = \Psi_t^{-1}(A)$, this implies

$$\mu|_A(t) = \Phi_t^v \# \mu|_A^0.$$

By compactness of $\text{supp}(\mu(t) \cap \omega_1^c)$, it holds

$$\mu(t)|_{\omega_1^c} = (\Phi_t^v \# \mu^0)|_{\omega_1^c}.$$

We deduce that, for all $\varphi \in \mathcal{C}_c^\infty(\mathbb{R}^d)$ such that $\text{supp}(\varphi) \subset \subset \omega_1^c$,

$$\frac{d}{dt} \int_{\mathbb{R}^d} \varphi d\mu(t) = \int_{\mathbb{R}^d} \langle \nabla \varphi, w \rangle d\mu(t) \text{ and } \frac{d}{dt} \int_{\mathbb{R}^d} \varphi d\mu(t) = \int_{\mathbb{R}^d} \langle \nabla \varphi, v \rangle d\mu(t).$$

If it holds $v \in \text{Tan}_{\mu(t)}(\mathcal{P}_c(\mathbb{R}^d))$, then $w(t) = v$, $\mu(t)$ a.e. in $\overline{\omega}_1^c$, and we conclude by taking $u := w - v$ which is supported in ω .

If now $v \notin \text{Tan}_{\mu(t)}(\mathcal{P}_c(\mathbb{R}^d))$, we can write $v = v_1 + v_2$ with $v_1 \in \text{Tan}_{\mu(t)}(\mathcal{P}_c(\mathbb{R}^d))$ and $v_2 \in \text{Tan}_{\mu(t)}(\mathcal{P}_c(\mathbb{R}^d))^\perp$, where

$$\text{Tan}_{\mu(t)}(\mathcal{P}_c(\mathbb{R}^d))^\perp = \{\nu \in L^2(\mu(t) : \mathbb{R}^d) : \nabla \cdot (\nu \mu(t)) = 0\}$$

(see for instance [4, Prop. 8.4.3]). In other terms, v_2 plays no role in the weak formulation of the continuity equation. Thus, with the same argument, we can prove that $w(t) = v_1$, $\mu(t)$ a.e. in $\overline{\omega}_1^c$ and we conclude by tacking $u := w - v_1$. \square

The second useful result for the proof of Theorem 1.2 allows to exactly steer a measure contained in ω to a nonempty open convex set $S \subset \subset \omega$. It is the analogue of Proposition 3.3. In this case, as in Proposition 4.1, we control the whole mass, but we do not have necessarily uniqueness of the solution to System (1).

Proposition 4.2. *Let $\mu^0 \in \mathcal{P}_c(\mathbb{R}^d)$ satisfying $\text{supp}(\mu^0) \subset \omega$. Define a nonempty open convex set S strictly included in $\omega \setminus \text{supp}(\mu^0)$ and choose $\delta > 0$. Then there exists a couple $(\mathbf{1}_\omega u, \mu)$ composed of a Borel vector field $\mathbf{1}_\omega u : \mathbb{R}^d \times \mathbb{R}^+ \rightarrow \mathbb{R}^d$ and a time-evolving measure μ being weak solution to System (1) satisfying*

$$\text{supp}(\mu(\delta)) \subset S.$$

Proof. Consider S_0 a nonempty open set of \mathbb{R}^d of class \mathcal{C}^∞ strictly included in S and ω_1 an open set of \mathbb{R}^d of class \mathcal{C}^∞ satisfying

$$\text{supp}(\mu^0) \cup S \subset \subset \omega_1 \subset \subset \omega.$$

An example is given in Figure 5. Consider $\eta \in \mathcal{C}^2(\overline{\omega}_1)$ defined in the proof of Proposition 3.3 satisfying (24). For all $k \in \mathbb{N}^*$, we consider a Lipschitz vector field v_k satisfying

$$v_k := \begin{cases} k \nabla \eta & \text{in } \omega_1, \\ v & \text{in } \omega^c. \end{cases}$$

We denote by

$$\tilde{t}_k^0(x^0) := \inf\{t \geq 0 : \Phi_t^{v_k}(x^0) \in \overline{S_0}\}.$$

For all $x^0 \in \mathbb{R}^d$ and all $k \in \mathbb{N}^*$, consider the application $\Psi_{k,\cdot}(x^0)$ defined for all $t \geq 0$ by

$$\Psi_{k,t}(x^0) = \begin{cases} \Phi_t^{v_k}(x^0) & \text{if } t \leq \tilde{t}_k^0(x^0), \\ \Phi_{\tilde{t}_k^0(x^0)}^{v_k}(x^0) & \text{otherwise.} \end{cases}$$

Using the same argument as in the proof of Proposition 3.3, for K large enough, $\Psi_{K,\delta}(x^0)$ belongs to S for all $x^0 \in \text{supp}(\mu^0)$. Consider μ defined for all $t \in (0, \delta)$ by $\mu(t) := \Psi_{K,t} \# \mu^0$. As in the proof of Proposition 4.1, there exists a vector field u_K such that (u_K, μ) is a weak solution to System (34). Moreover $u_K(t) = v_K$, $\mu(t)$ a.e. in $\overline{S^c}$ and a.e. $t \in [0, \delta]$. Thus, we conclude that $(\mathbb{1}_\omega(u_K - v_K), \mu)$ is solution to System (1) and $\text{supp}(\mu(\delta)) \subset S$. \square

The third useful result for the proof of Theorem 1.2 allows to exactly steer a measure contained in a nonempty open convex set $S \subset\subset \omega$ to a given measure contained in S . It is the analogue of Proposition 3.1. In this situation, we obtain exact controllability of System (1), but, again, we do not have necessarily uniqueness of the solution to System (1).

Proposition 4.3. *Let $\mu^0, \mu^1 \in \mathcal{P}_c(\mathbb{R}^d)$ satisfying $\text{supp}(\mu^0) \subset S$ and $\text{supp}(\mu^1) \subset S$ for a nonempty open convex set S strictly included in ω . Choose $\delta > 0$. Then there exists a couple $(\mathbb{1}_\omega u, \mu)$ composed of a Borel vector field $\mathbb{1}_\omega u : \mathbb{R}^d \times \mathbb{R}^+ \rightarrow \mathbb{R}^d$ and a time-evolving measure μ being weak solution to System (1) and satisfying*

$$\text{supp}(\mu) \subset S \text{ and } \mu(\delta) = \mu^1.$$

Proof. Let π be the optimal plan given in (3) associated to the Wasserstein distance between μ^0 and μ^1 . For $i \in \{1, 2\}$, we denote by $p_i : \mathbb{R}^d \times \mathbb{R}^d \rightarrow \mathbb{R}^d$ the projection operator defined by

$$p_i : (x_1, x_2) \mapsto x_i.$$

Consider the time-evolving measure μ defined for all $t \in [0, \delta]$ by

$$\mu(t) := \frac{1}{\delta} [(\delta - t)p_1 + tp_2] \# \pi. \quad (37)$$

Using [4, Th. 7.2.2], μ is a constant speed geodesic connecting μ^0 and μ^1 in $\mathcal{P}_c(\mathbb{R}^d)$, i.e. for all $s, t \in [0, \delta]$

$$W_2(\mu(t), \mu(s)) = \frac{(t - s)}{\delta} W_2(\mu^0, \mu^1).$$

We deduce that the metric derivative $|\mu'|$ of μ (see (33)) is uniformly bounded on $[0, \delta]$. Then μ is an absolute continuous curve on $\mathcal{P}_c(\mathbb{R}^d)$ (see [4, Def. 1.1.1]). Thus, using [4, Th. 8.3.1], there exists a Borel vector field $w : \mathbb{R}^d \times (0, \delta) \rightarrow \mathbb{R}^d$ such that

$$\|w(t)\|_{L^2(\mu(t); \mathbb{R}^d)} \leq |\mu'| (t) \text{ a.e. } t \in [0, \delta]$$

and the couple (w, μ) is a weak solution to

$$\begin{cases} \partial_t \mu + \nabla \cdot (w\mu) = 0 & \text{in } \mathbb{R}^d \times [0, \delta], \\ \mu(0) = \mu^0 & \text{in } \mathbb{R}^d. \end{cases}$$

Consider $\theta \in \mathcal{C}_c^\infty(\mathbb{R}^d)$ such that

$$0 \leq \theta \leq 1, \quad \theta = 1 \text{ in } S \quad \text{and} \quad \theta = 0 \text{ in } \omega^c.$$

We remark that μ is supported in S , then the couple $(\mathbb{1}_\omega u, \mu)$ with

$$u := \theta \times (w - v)$$

is solution to

$$\begin{cases} \partial_t \mu + \nabla \cdot ((v + \mathbb{1}_\omega u)\mu) = 0 & \text{in } \mathbb{R}^d \times [0, \delta], \\ \mu(0) = \mu^0 & \text{in } \mathbb{R}^d. \end{cases}$$

□

Remark 4. It is well known (see for instance [29]) that for two measures compactly supported with the same total mass μ^0, μ^1 , the Wasserstein distance can be expressed as follows

$$W_2(\mu^0, \mu^1) = \min_{(v, \mu) \in \mathcal{B}} \left\{ \left(\int_0^1 \int_{\mathbb{R}^d} |v(t)|^2 d\mu(t) dt \right)^{1/2} : \begin{aligned} & \partial_t \mu + \nabla \cdot (v\mu) = 0, \quad \mu(0) = \mu^0, \quad \mu(1) = \mu^1 \end{aligned} \right\}, \quad (38)$$

where \mathcal{B} is the set of couples (v, μ) composed of a time evolving measure $\mu(t)$ and a Borel vector field $v : \mathbb{R}^d \times \mathbb{R} \rightarrow \mathbb{R}^d$ satisfying

$$\int_0^1 \int_{\mathbb{R}^d} |v(t)|^2 d\mu(t) dt < \infty.$$

Equality (38) is called the *Benamou-Brenier Formula*. In the proof of the Proposition 4.3, it is possible to replace the definition in (37) of μ by the minimizer of (38) which already satisfies the continuity equation.

We now have all the tools to prove Theorem 1.2.

Proof of Theorem 1.2. Consider μ^0 and μ^1 satisfying Condition 1.1. Applying Lemma 3.1, Condition 3.1 holds for some ω_0, T_0^* and T_1^* . Let $T := T_0^* + T_1^* + \delta$ with $\delta > 0$ and $T_0, T_1, T_2, T_3, T_4, T_5$ be the times given in the proof of Theorem 1.1. Using Proposition 4.1 on $[T_0, T_1] \cup [T_4, T_5]$, there exist $\rho_1 \in \mathcal{C}^0([T_0, T_1], \mathcal{P}_c(\mathbb{R}^d))$, $\rho_5 \in \mathcal{C}^0([T_4, T_5], \mathcal{P}_c(\mathbb{R}^d))$ and some space-dependent Borel controls u^1, u^5 with

$$\text{supp}(u^1) \cup \text{supp}(u^5) \subset \omega$$

such that $(\mathbb{1}_\omega u^1, \rho_1)$ is a weak solution forward in time to

$$\begin{cases} \partial_t \rho_1 + \nabla \cdot ((v + \mathbb{1}_\omega u^1)\rho_1) = 0 & \text{in } \mathbb{R}^d \times [T_0, T_1], \\ \rho_1(T_0) = \mu^0 & \text{in } \mathbb{R}^d \end{cases}$$

and $(\mathbb{1}_\omega u^5, \rho_5)$ is a weak solution backward in time to

$$\begin{cases} \partial_t \rho_5 + \nabla \cdot ((v + \mathbb{1}_\omega u^5)\rho_5) = 0 & \text{in } \mathbb{R}^d \times [T_4, T_5], \\ \rho_5(T_5) = \mu^1 & \text{in } \mathbb{R}^d. \end{cases}$$

Moreover $\text{supp}(\rho_1(T_1)) \subset \omega$ and $\text{supp}(\rho_5(T_4)) \subset \omega$. Consider a nonempty open convex set S strictly included in $\omega \setminus \omega_0$. Using Proposition 4.2 on $[T_1, T_2] \cup [T_3, T_4]$, there exist $\rho_2 \in \mathcal{C}^0([T_1, T_2], \mathcal{P}_c(\mathbb{R}^d))$, $\rho_4 \in \mathcal{C}^0([T_3, T_4], \mathcal{P}_c(\mathbb{R}^d))$ and some space-dependent Borel controls u^2, u^4 with

$$\text{supp}(u^2) \cup \text{supp}(u^4) \subset \omega$$

such that $(\mathbb{1}_\omega u^2, \rho_2)$ is a weak solution forward in time to

$$\begin{cases} \partial_t \rho_2 + \nabla \cdot ((v + \mathbb{1}_\omega u^2) \rho_2) = 0 & \text{in } \mathbb{R}^d \times [T_1, T_2], \\ \rho_2(T_1) = \rho_1(T_1) & \text{in } \mathbb{R}^d \end{cases}$$

and $(\mathbb{1}_\omega u^4, \rho_4)$ is a weak solution backward in time to

$$\begin{cases} \partial_t \rho_4 + \nabla \cdot ((v + \mathbb{1}_\omega u^4) \rho_4) = 0 & \text{in } \mathbb{R}^d \times [T_3, T_4], \\ \rho_4(T_4) = \rho_5(T_4) & \text{in } \mathbb{R}^d. \end{cases}$$

Moreover $\text{supp}(\rho_2(T_2)) \subset S$ and $\text{supp}(\rho_4(T_3)) \subset S$. Using Proposition 4.3 on $[T_2, T_3]$, there exist $\rho_3 \in \mathcal{C}^0([T_2, T_3], \mathcal{P}_c(\mathbb{R}^d))$ satisfying $\text{supp}(\rho_3) \subset S$ and a Borel control u^3 with

$$\text{supp}(u^3) \subset \omega$$

such that $(\mathbb{1}_\omega u^3, \rho_3)$ is a weak solution forward in time to

$$\begin{cases} \partial_t \rho_3 + \nabla \cdot ((v + \mathbb{1}_\omega u^3) \rho_3) = 0 & \text{in } \mathbb{R}^d \times [T_2, T_3], \\ \rho_3(T_2) = \rho_2(T_2) & \text{in } \mathbb{R}^d \end{cases}$$

and satisfies $\rho_3(T_3) = \rho_4(T_3)$. Thus the couple $(\mathbb{1}_\omega u, \mu)$ defined by

$$(\mathbb{1}_\omega u, \mu) = (\mathbb{1}_\omega u^i, \rho_i) \text{ in } \mathbb{R}^d \times [T_{i-1}, T_i], \quad i \in \{1, \dots, 5\}$$

is a weak solution to System (1) and satisfies $\mu(T) = \mu^1$. \square

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