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Paracontrolled quasilinear SPDEs

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ABSTRACT. We introduce a non-linear paracontrolled calculus and use it to renormalise a class of singular SPDEs including certain quasilinear variants of the periodic two dimensional parabolic Anderson model.

1. INTRODUCTION

We show how to renormalise a class of general quasilinear equations of which one of the simplest examples is the following parabolic SPDE:

$$\partial_t u(t, x) - a(u(t, x))\Delta u(t, x) = \xi(x), \quad u(0, x) = u_0(x), \quad x \in \mathbb{T}^2, t \geq 0, \quad (1)$$

with $a: \mathbb{R} \rightarrow [\lambda, 1]$ for $\lambda > 0$ a uniformly bounded C^3 diffusion coefficient, and $\|a^{(k)}\|_{L^\infty} \leq 1$ for $k = 0, \dots, 3$. We assume that $\xi \in \mathcal{C}^{\alpha-2}(\mathbb{T}^2)$ with $2/3 < \alpha < 1$ where $\mathcal{C}^\alpha(\mathbb{T}^2)$ is the Besov space $B_{\infty, \infty}^\alpha(\mathbb{T}^2)$. This would apply to the space white noise on \mathbb{T}^2 , for example. In this case we only expect that $u(t, \cdot) \in \mathcal{C}^\alpha(\mathbb{T}^2)$ and the term $a(u(t, \cdot))\Delta u(t, \cdot)$ is not well defined when $2\alpha - 2 < 0$. Eq. (1) is a quasilinear generalisation of the two-dimensional periodic parabolic Anderson model (PAM).

Let us remark from the start that the framework we will consider below allows to deal with a class of equations of the form

$$a_1(u(t, x))\partial_t u(t, x) - a_2(u(t, x))\Delta u(t, x) = \xi(a_3(u(t, x)), t, x), \quad x \in \mathbb{T}^2, t \geq 0, \quad (2)$$

where a_1, a_2 are sufficiently smooth non-degenerate coefficients and $\xi(z, t, x)$ is a Gaussian process with covariance

$$\mathbb{E}[\xi(z, t, x)\xi(z', t', x')] = F(z, z')Q(t - t', x - x'), \quad x, x' \in \mathbb{T}^2, t, t', z, z' \in \mathbb{R},$$

with F a smooth function and Q a distribution of parabolic regularity $\rho > -4/3$. This includes the space white noise discussed before or a time white noise with a regular dependence on the space variable or some noise mildly irregular in space and time.

Also the scalar character of the equation or of the non-linear diffusion coefficient will not play any specific role and we could consider vector-valued equations with general diffusion coefficients provided the template problem (7) below remains uniformly parabolic.

For the sake of clarity and simplicity we will discuss mainly the basic example (1) since this contains already most of the technical difficulties. The fact that one can handle models as general as (2) can be considered a direct byproduct of the techniques we will introduce below.

Recently Otto and Weber [33] and Bailleul, Debussche and Hofmanova [7] investigated quasilinear SPDEs in the context of pathwise methods and in a range of regularities compatible with the ones we will consider in this paper.

- In [33] the authors obtained a priori estimates for equations of the form

$$\partial_t u(t, x) - a(t, x)\partial_x^2 u(t, x) = f(u(t, x))\xi(t, x), \quad (t, x) \in \mathbb{T}^2$$

where both space and time variables take values in a one dimensional periodic domain and their noise can be white in time but colored in space, essentially behaving like a distribution of parabolic regularity in $(-4/3, 1)$. In order to do so they introduce a specific notion of *modelled* function and related estimates.

- Bailleul, Debussche and Hofmanová in [7] obtain local well-posedness for the generalised parabolic Anderson model equation

$$\partial_t u(t, x) - a(u(t, x))\Delta u(t, x) = g(u(t, x))\xi(x) \quad t \geq 0, x \in \mathbb{T}^2. \quad (3)$$

The authors obtain the same result as the one presented in Section 6 of our work, without the machinery of nonlinear para-products introduced here, but using only the basic tools of paracontrolled analysis and some clever transformations.

On the other hand, we remark that the apparently innocuous vectorial formulation of (3)

$$\partial_t u(t, x) - a_{ij}(u(t, x))\frac{\partial^2}{\partial x_i \partial x_j} u(t, x) = g(u(t, x))\xi \quad t \geq 0, x \in \mathbb{T}^2$$

is out of reach of the techniques used in [7], while can be treated flawlessly in our framework.

Let us state one simple result that can be obtained via the theory developed in this paper:

Theorem 1. *Fix $2/3 < \alpha < 1$. Let $\xi \in \mathcal{C}^{\alpha-2}(\mathbb{T}^2)$ be a space white noise with zero average on the torus, $u_0 \in \mathcal{C}^\alpha$ an initial condition and $a: \mathbb{R} \rightarrow [\lambda, 1]$ for some $\lambda > 0$, $a \in C^3(\mathbb{R})$ and $\|a^{(k)}\|_{L^\infty} \leq 1 \ \forall k \in 0, \dots, 3$. Let $(\xi_\varepsilon, u_{0,\varepsilon})_{\varepsilon>0}$ be a family of smooth approximations to ξ, u_0 obtained by convolution with a rescaled smoothing kernel and u_ε the classical solution to the Cauchy problem*

$$\partial_t u_\varepsilon - a(u_\varepsilon)\Delta u_\varepsilon = \xi_\varepsilon + \sigma_\varepsilon \frac{a'(u_\varepsilon)}{a(u_\varepsilon)^2}, \quad u(0) = u_{0,\varepsilon}. \quad (4)$$

Then we can choose the constants $(\sigma_\varepsilon)_{\varepsilon>0}$ and a random time $T > 0$ in such a way that the family of r.v. $(u_\varepsilon)_{\varepsilon>0} \subseteq \mathcal{L}_T^\alpha(\mathbb{T}^2)$ almost surely converge as $\varepsilon \rightarrow 0$ to a random element $u \in \mathcal{L}_T^\alpha(\mathbb{T}^2)$, where \mathcal{L}_T^α is the parabolic space $C([0, T], \mathcal{C}^\alpha(\mathbb{T}^2)) \cap C^{\alpha/2}([0, T], \mathcal{C}^0(\mathbb{T}^2))$.

This element can be characterised as the solution to a paracontrolled singular SPDE (see below for more details).

In order to devise a suitable formulation of eq. (1) and obtain a theory with $u \in \mathcal{C}^\alpha$ we decompose the non-linear diffusion term in the l.h.s. with the help of Bony's paraproduct [31] and write

$$\partial_t u - a(u) \prec \Delta u = \xi + \Phi(u) \quad (5)$$

with

$$\Phi(u) := a(u) \circ \Delta u + a(u) \succ \Delta u \quad (6)$$

where \prec, \succ are standard paraproducts and \circ denotes the resonant product (see below for precise definitions). Now the l.h.s. is always well defined irrespective of the regularity of the function u and the problem becomes that of controlling the resonant product $a(u) \circ \Delta u$ appearing in the r.h.s. . A key point of the analysis put forward below is that this term can be expected to be of regularity $2\alpha - 2 > \alpha - 2$ so better than the leading term ξ .

Our approach can be described as follows. For an equation of the form

$$\partial_t u - a_1(u)\Delta u = a_2(u)\xi,$$

we consider at first a *parametric* “template” problem with constant coefficients

$$\partial_t \vartheta(\eta, t, x) - \eta_1 \Delta \vartheta(\eta, t, x) = \eta_2 \xi(t, x), \quad (7)$$

where now $\eta = (\eta_1, \eta_2)$ are fixed numbers. A *nonlinear* paraproduct Π_{\ll} will allow us to modulate the parametric solution ϑ with the coefficient $a(u) = (a_1(u), a_2(u))$ as to capture the most irregular part of the solution u itself. As a consequence, the paracontrolled Ansatz

$$u = \Pi_{\ll}(a(u), \vartheta) + u^\sharp$$

will define a regular remainder term u^\sharp which solves a standard PDE. With this decomposition the resonant products appearing in the equation can be estimated along the lines of the standard paracontrolled arguments introduced in [15] and all the arguments introduced there can be extended in a straightforward manner to the quasilinear setting.

This approach has been inspired by the parametric controlled Ansatz of Otto and Weber [33]. At variance with their approach we use the parametric Ansatz in the context of the paradifferential calculus and consider more general noise terms.

Usefulness of paraproducts in the analysis of non-linear PDEs is by now well established: see for example the seminal paper of Meyer [31], the early review of Bony [10], the recent books of Alinhac and Gérard [1] and Bahouri, Chemin, and Danchin [3]. Let us mention also the interesting paper of Hörmander [25] where paradifferential operators allows to bypass the Nash–Moser fixpoint theorem in some applications where the loss of regularity prevents straightforward use of standard Banach fixpoint theorem. The main observation in that paper is that, with the aid of paradifferential operators, it is possible to identify a “corrected” problem for which standard Banach fixpoint applies.

Paracontrolled calculus for singular SPDEs has been introduced by Gubinelli, Imkeller and Perkowski [15] (see also the lecture notes [17]) and used to study various equations like the KPZ equation [16], the dynamic Φ_3^4 model [11] in $d = 3$ and its global well-posedness [32], the spectrum of the continuous Anderson Hamiltonian in $d = 2$ [2]. By using heat-semigroup techniques paracontrolled calculus has been extended to the manifold context by Bailleul and Bernicot [4].

Non-linear generalisation of the classic bilinear paraproducts already appeared in the notion of *paracomposition* introduced by Alinhac [10]. Non-linear versions of rough paths have been considered by one of the authors in order to study the Korteweg–de Vries equation [14]. Non-linear Young integrals were used by one of the authors in joint work with Catellier [12] to study the regularising properties of sample paths of stochastic processes. See also the related work of Hu and Le [24] on differential systems in Hölder media. Relevant to this discussion of non-linear variants of rough paths is the work of Bailleul on rough flows [8] and their application to homogenisation [6]. By looking at the composition $f(g(x))$ as the action of the distribution $\delta_{g(x)}$ on the function f , non-linear constructions can be linearised at the price of working in infinite-dimensional spaces: this is the approach chosen by Kelly and Melbourne to avoid non-linear generalisations of rough path theory in their study of homogenisation of fast–slow system with random initial conditions [26]. It is worth mentioning also Kunita’s theory of semimartingale vector fields [27] which occupy a place in stochastic analysis quite similar to that which these recent developments occupy in the landscape of rough paths/paracontrolled distributions theories.

Paracontrolled calculus is currently limited to “first order” computations. This limitation is also ubiquitous in the present work. Even if, in practice, this is not a big issue, and the calculus is still able to deal with a large class of problems, it makes the paracontrolled approach less appealing for a general theory of singular SPDEs. Let us remark that recently Bailleul and Bernicot [5] developed an higher order version of the paracontrolled calculus. However, apart from these recent development, whose impact is still to be assessed, the most general theory for singular SPDEs has been developed by Hairer [19, 20, 13] under the name of *regularity structures* theory. Regularity structures are a vast generalisation of Lyons’ rough paths [28, 30, 29] which give effective tools to describe non-linear operations acting on certain spaces of distributions, their renormalization by subtraction of local singularities and their use to solve singular SPDE. Regularity structures have been successfully applied to all the models mentioned so far [19, 18], to other models like the Sine–Gordon model [22] (which however can also be handled via paracontrolled techniques) and to study weak universality conjectures [21, 23]. In their current instantiation it does not seem possible to solve quasilinear SPDEs via regularity structures. The results of the present paper hint to the fact that a non-linear version of regularity structures is conceivable, at least in principle. Indeed one can imagine models depending on additional parameters and modelled distributions acting as evaluations of the parametric models at certain space–time dependent values of the parameters. It would be interesting to pursue further this intuition.

The structure of the paper is the following. In Section 2 we introduce our basic tools: the non-linear paraproduct decomposition and some related commutation lemmas. In Section 3 we introduce the paracontrolled Ansatz which allows to transform the singular problem (1) into a well–behaved PDE. In Section 4 we discuss the apriori estimates, the uniqueness of the solution of the transformed PDE and its continuity w.r.t. the random data and the initial condition, we introduce also the algebraic structure which allows to renormalise the model. Section 5 deals with the renormalization of the stochastic data and the construction of the enhanced noise associated to white noise. Section 6 deals with the extension of the results to more general equations, in particular with equation (3) or with noise whose law depends on the solution itself. Finally Appendix A reviews some reference material on Besov spaces and proves some technical lemmas.

Notations. We will denote $\mathcal{C}^\alpha := B_{\infty,\infty}^\alpha(\mathbb{T}^2)$ the Zygmund space of regularity $\alpha \in \mathbb{R}$ on the torus \mathbb{T}^2 . See Appendix A for the definition of the general Besov spaces $B_{p,q}^\alpha$, the Littlewood–Paley operators $(\Delta_i)_{i \geq -1}$ and the basic properties thereof needed in this paper. If V is a Banach space and $T > 0$, we denote $C_T^\alpha V$ the space of α –Hölder functions in $C_T V := C([0, T]; V)$. We introduce parabolic spaces $\mathcal{L}_T^\alpha := C_T^{\alpha/2} \mathcal{C}^0 \cap C_T \mathcal{C}^\alpha$ with norm

$$\|f\|_{\mathcal{L}_T^\alpha} = \|f\|_{C_T^{\alpha/2} \mathcal{C}^0} + \|f\|_{C_T \mathcal{C}^\alpha}. \quad (8)$$

Moreover for convenience we denote $\mathcal{C}_T^\alpha := C_T \mathcal{C}^\alpha$. We will avoid to note explicitly the time span T whenever this does not cause ambiguities. We will need also spaces for functions of (η, t, x) where η is an additional parameter in $[\lambda, 1]$ for $\lambda \in (0, 1)$ which we denote $C_\eta^k V$ with norm

$$\|F\|_{C_\eta^k V} = \sup_{\eta \in [\lambda, 1]} \sup_{n=0, \dots, k} \|\partial_\eta^n F(\eta, \cdot)\|_V, \quad (9)$$

where V is a Banach space of functions on $[0, T] \times \mathbb{T}^2$, in our case $V = \mathcal{C}_T^\alpha$ or $V = \mathcal{L}_T^\alpha$.

We will denote by $K_{i,x}(y) = 2^{2i}K(2^i(x - y))$ the kernel of the Littlewood–Paley operator Δ_i and $Q_{i,t}(s) = Q_i(t - s) = 2^{2i}Q(2^{2i}(t - s))$ a smoothing kernel at scale 2^{2i} in the time direction where Q is a smooth, positive function with compact support in \mathbb{R}_+ and mass 1. We introduce also the shortcut $P_{i,x} = K_{<i-1,x} = \sum_{j < i-1} K_{j,x}$. Another notation shortcut widely used in this article is to write $\int_{x,y}$ for integrals on \mathbb{T}^2 or \mathbb{R} with respect to the measures dx and dy without specifying the integration bounds, whenever this does not create ambiguity. Finally, we will note $\delta f_{s,y}^{t,x} = f(t, x) - f(s, y)$ and $\delta_\tau f_{s,y}^{t,x} = f(s, y) + \tau(f(t, x) - f(s, y))$ for $\tau \in [0, 1]$.

2. NONLINEAR PARAPRODUCTS

In this section we introduce the nonlinear paraproduct and related results that will be used in Section 3 to analyse equation (1).

Let $g: [0, T] \times \mathbb{T}^2 \rightarrow \mathbb{R}$, and $h: \mathbb{R} \times [0, T] \times \mathbb{T}^2 \rightarrow \mathbb{R}$ be smooth functions. We can decompose the composition $h(g(\cdot), \cdot)$ via *nonlinear* paraproducts as follows. Define

$$\Pi_{<}(g, h)(t, x) := \sum_q \int_{y,z} P_{q,x}(y) K_{q,x}(z) h(g(t, y), t, z) \quad (10)$$

$$\Pi_{\circ}(g, h)(t, x) := \sum_{k \sim q} \int_{y,z} K_{k,x}(y) K_{q,x}(z) h(g(t, y), t, z) \quad (11)$$

$$\Pi_{>}(g, h)(t, x) := \sum_k \int_{y,z} K_{k,x}(z) P_{k,x}(y) h(g(t, z), t, y). \quad (12)$$

This gives a map

$$(g, h) \mapsto \Pi_{\diamond}(g, h) := \Pi_{<}(g, h) + \Pi_{\circ}(g, h) + \Pi_{>}(g, h) = h(g(\cdot), \cdot) \quad (13)$$

that can be uniquely extended to

$$\Pi_{\diamond}: \mathcal{C}_T^\rho \times C_\eta^2 \mathcal{C}_T^\gamma \rightarrow \mathcal{C}_T^{\gamma \wedge \rho} \quad \rho \in (0, 1), \gamma \in \mathbb{R}, \rho + \gamma > 0,$$

thanks to the following bounds:

Lemma 2. (Nonlinear paraproduct estimates) *Let $g \in \mathcal{C}_T^\rho$ for some $\rho \in (0, 1)$, $g \in [\lambda, 1]$, and $h \in C_\eta^2 \mathcal{C}_T^\gamma$ for any $\gamma \in \mathbb{R}$. Then*

$$\|\Pi_{<}(g, h)\|_{\mathcal{C}_T^\gamma} \lesssim \|h\|_{C_\eta \mathcal{C}_T^\gamma}, \quad \|\Pi_{>}(g, h)\|_{\mathcal{C}_T^{\rho \wedge (\rho + \gamma)}} \lesssim \|g\|_{\mathcal{C}_T^\rho} \|h\|_{C_\eta^1 \mathcal{C}_T^\gamma},$$

and

$$\|\Pi_{<}(g_1, h) - \Pi_{<}(g_2, h)\|_{\mathcal{C}_T^\gamma} \lesssim \|g_1 - g_2\|_{C_T L^\infty} \|h\|_{C_\eta^1 \mathcal{C}_T^\gamma},$$

$$\begin{aligned} \|\Pi_{>}(g_1, h) - \Pi_{>}(g_2, h)\|_{\mathcal{C}_T^{\rho \wedge (\rho + \gamma)}} &\lesssim \|g_1 - g_2\|_{C_T L^\infty} (\|g_1\|_{\mathcal{C}_T^\rho} + \|g_2\|_{\mathcal{C}_T^\rho}) \|h\|_{C_\eta^2 \mathcal{C}_T^\gamma} \\ &\quad + \|g_1 - g_2\|_{\mathcal{C}_T^\rho} \|h\|_{C_\eta^1 \mathcal{C}_T^\gamma}. \end{aligned}$$

Moreover if $\rho + \gamma > 0$ we have also

$$\|\Pi_{\circ}(g, h)\|_{\mathcal{C}_T^{\rho + \gamma}} \lesssim \|h\|_{C_\eta^1 \mathcal{C}_T^\gamma} \|g\|_{\mathcal{C}_T^\rho},$$

$$\begin{aligned} \|\Pi_{\circ}(g_1, h) - \Pi_{\circ}(g_2, h)\|_{\mathcal{C}_T^{\rho + \gamma}} &\lesssim \|g_1 - g_2\|_{C_T L^\infty} (\|g_1\|_{\mathcal{C}_T^\rho} + \|g_2\|_{\mathcal{C}_T^\rho}) \|h\|_{C_\eta^2 \mathcal{C}_T^\gamma} \\ &\quad + \|g_1 - g_2\|_{\mathcal{C}_T^\rho} \|h\|_{C_\eta^1 \mathcal{C}_T^\gamma}. \end{aligned}$$

In particular if $\rho + \gamma > 0$ the composition $\Pi_{\diamond}(g, h) = h(g(\cdot), \cdot)$ is linear in h and locally Lipschitz in g :

$$\begin{aligned} \|\Pi_{\diamond}(g, h)\|_{\mathcal{E}_T^{\gamma}} &\lesssim \|h\|_{C_{\eta}^1 \mathcal{E}_T^{\gamma}} \|g\|_{\mathcal{E}_T^{\rho}}, \\ \|\Pi_{\diamond}(g_1, h) - \Pi_{\diamond}(g_2, h)\|_{\mathcal{E}_T^{\gamma}} &\lesssim \|g_1 - g_2\|_{\mathcal{E}_T^{\rho}} (1 + \|g_1\|_{\mathcal{E}_T^{\rho}} + \|g_2\|_{\mathcal{E}_T^{\rho}}) \|h\|_{C_{\eta}^2 \mathcal{E}^{\gamma}}. \end{aligned}$$

Proof. Using the fact that

$$\begin{aligned} \|\Delta_k h(g(t, y), t, \cdot)\|_{L^{\infty}} &\lesssim 2^{-\gamma k} \|h\|_{C_{\eta} \mathcal{E}_T^{\gamma}}, \\ \|\Delta_k h(g(t, y), t, \cdot) - \Delta_k h(g(t, y'), t, \cdot)\|_{L^{\infty}} &\lesssim 2^{-\gamma k} \|h\|_{C_{\eta}^1 \mathcal{E}_T^{\gamma}} \|g\|_{C_T \mathcal{E}^{\rho}} |y - y'|^{\rho}, \end{aligned}$$

and

$$\|\Delta_k h(g_1(t, y), t, \cdot) - \Delta_k h(g_2(t, y), t, \cdot)\|_{L^{\infty}} \lesssim 2^{-\gamma k} \|h\|_{C_{\eta}^1 \mathcal{E}_T^{\gamma}} \|g_1 - g_2\|_{C_T L^{\infty}},$$

we obtain the bounds on $\Pi_{\prec}(g, h)$, $\Pi_{\succ}(g, h)$, $\Pi_{\circ}(g, h)$ and $\Pi_{\prec}(g_1, h) - \Pi_{\prec}(g_2, h)$. We proceed to estimate the term $\Pi_{\succ}(g_1, h) - \Pi_{\succ}(g_2, h)$. We will use the following notation for brevity:

$$\delta g_{2z}^{1y} := g_1(t, y) - g_2(t, z) \quad \text{and} \quad \delta_{\tau} g_{2z}^{1y} := g_2(t, z) + \tau(g_1(t, y) - g_2(t, z)).$$

Then

$$\begin{aligned} &\left| \int_{y,z} K_{k,x}(z) P_{k,x}(y) [h(g_1(t, z), t, y) - h(g_2(t, z), t, y)] \right| \\ &= \left| \int_{y,z,\tau \in [0,1]} K_{k,x}(z) P_{k,x}(y) [\partial_{\eta} h(\delta_{\tau} g_{2z}^{1z}, t, y) \delta g_{2z}^{1z} - \partial_{\eta} h(\delta_{\tau} g_{2x}^{1x}, t, y) \delta g_{2x}^{1x}] \right| \\ &\lesssim \left| \int_{y,z,t \in [0,1], \sigma \in [0,1]} K_{k,x}(z) P_{k,x}(y) \partial_{\eta}^2 h(\delta_{\sigma}(\delta_{\tau} g_{2x}^{1z}), t, y) (\delta g_{1x}^{1z} - \delta g_{2x}^{2z}) \delta g_{2z}^{1z} \right| \\ &\quad + \left| \int_{y,z,\tau \in [0,1]} K_{k,x}(z) P_{k,x}(y) \partial_{\eta} h(\delta_{\tau} g_{2x}^{1x}, t, y) (\delta g_{2z}^{1z} - \delta g_{2x}^{1x}) \right| \\ &\lesssim \|g_1 - g_2\|_{C_T L^{\infty}} (\|g_1\|_{C_T \mathcal{E}^{\rho}} + \|g_2\|_{C_T \mathcal{E}^{\rho}}) \|h\|_{C_{\eta}^2 C_T \mathcal{E}^{\gamma}} 2^{-\rho k} \sum_{q < k-1} 2^{-\gamma q} \\ &\quad + \|g_1 - g_2\|_{C_T \mathcal{E}^{\rho}} \|h\|_{C_{\eta}^1 C_T \mathcal{E}^{\gamma}} 2^{-\rho k} \sum_{q < k-1} 2^{-\gamma q}. \end{aligned}$$

With the same reasoning we can bound the norm of $\Pi_{\circ}(g_1, h) - \Pi_{\circ}(g_2, h)$. \square

We will need the following time-smoothed nonlinear paraproduct

$$\Pi_{\ll}(g, h)(t, x) := \sum_i \int_{y,s} Q_{i,t}(s) P_{i,x}(y) (\Delta_i h(g(s, y), t, \cdot))(x), \quad (14)$$

with $Q \in C_c^{\infty}(\mathbb{R})$ with total mass 1, and $Q_{i,t}(s) := 2^{2i} Q(2^i(t - s))$ as specified in the introduction. In (14) we use the convention that a continuous function $t \mapsto g(t)$ on \mathbb{R}_+ is extended to \mathbb{R} by defining $g(t) = g(0)$ for $t \leq 0$. This preserves the Hölder norms of index in $[0, 1]$. The modified nonlinear paraproduct enjoys similar bounds to the regular one.

Lemma 3. *Let $g \in C_T L^{\infty}$, $g \in [\lambda, 1]$ and $h \in C_{\eta}^1 \mathcal{E}_T^{\gamma}$ for $\gamma \in (0, 2)$. Then*

$$\|\Pi_{\ll}(g, h)\|_{\mathcal{E}_T^{\gamma}} \lesssim \|h\|_{C_{\eta} \mathcal{E}_T^{\gamma}} \quad \text{and} \quad \|\Pi_{\ll}(g, h)\|_{\mathcal{L}_T^{\gamma}} \lesssim \|h\|_{C_{\eta} \mathcal{E}_T^{\gamma}}.$$

Moreover, $\Pi_{\ll}(g, h)$ is linear in h and:

$$\begin{aligned} \|\Pi_{\ll}(g_1, h) - \Pi_{\ll}(g_2, h)\|_{\mathcal{E}_T^{\gamma}} &\lesssim \|g_1 - g_2\|_{C_T L^{\infty}} \|h\|_{C_{\eta}^1 \mathcal{E}_T^{\gamma}}, \\ \|\Pi_{\ll}(g_1, h) - \Pi_{\ll}(g_2, h)\|_{\mathcal{L}_T^{\gamma}} &\lesssim \|g_1 - g_2\|_{C_T L^{\infty}} \|h\|_{C_{\eta}^1 \mathcal{E}_T^{\gamma}}. \end{aligned}$$

Proof. The norm $\|\Pi_{\leftarrow}(g, h)\|_{C_T \mathcal{C}^\gamma}$ can be treated in the same way as in Lemma 2. We estimate $\|\Pi_{\leftarrow}(g, h)\|_{C_T^{\gamma/2} \mathcal{C}^0}$ as follows:

$$\begin{aligned} & \|\Delta_j \Pi_{\leftarrow}(g, h)(t_1) - \Delta_j \Pi_{\leftarrow}(g, h)(t_2)\|_{L^\infty} \\ & \lesssim \sup_x \left| \int_z K_{j,x}(z) \sum_{i \sim j} \int_{y,s} Q_{i,t_1}(s) P_{i,z}(y) [\Delta_i h(g(s, y), t_1, z) - \Delta_i h(g(s, y), t_2, z)] \right| \\ & \quad + \sup_x \left| \int_z K_{j,x}(z) \sum_{i \sim j} \int_{y,s} [Q_{i,t_1}(s) - Q_{i,t_2}(s)] P_{i,z}(y) \Delta_i h(g(s, y), t_2, z) \right| \\ & \lesssim \|h(\cdot, t_1, \cdot) - h(\cdot, t_2, \cdot)\|_{C_\eta \mathcal{C}^0} + |t_1 - t_2|^{\gamma/2} \|h\|_{C_\eta \mathcal{C}_T^\gamma}. \end{aligned}$$

The second inequality can be obtained easily with the same techniques used so far. \square

Remark 4. Using the Fourier support properties of the kernel $P_{q,x}$ it is easy to see that it has mass 1. Therefore, the paraproducts (10) and (14) when $g(t, x) = \bar{g}$ is a constant smooth function become $\forall (t, x) \in \mathbb{R} \times \mathbb{T}^2$:

$$\begin{aligned} \Pi_{\leftarrow}(\bar{g}, h)(t, x) &= h(\bar{g}, t, x) = \Pi_{\diamond}(\bar{g}, h), \\ \Pi_{\leftarrow}(\bar{g}, h)(t, x) &= h(\bar{g}, t, x). \end{aligned}$$

2.1. Nonlinear commutator.

The next technical ingredient is a commutator lemma between the non-linear paraproduct of (14) and the standard resonant product. It will be needed below to analyse a term of the form $\Pi_{\leftarrow}(g, h) \circ \Delta \Pi_{\leftarrow}(g, h)$, then we will specialise our discussion to this specific structure. Notice that in the following the various space-time operators act pointwise in the parameter η , in the sense that, for example:

$$(h \circ \Delta h)(\eta, t, x) = (h(\eta, t, \cdot) \circ \Delta h(\eta, t, \cdot))(x).$$

Lemma 5. We define the map $\Lambda: C^\infty([0, T], \mathbb{T}^2) \times C_\eta^2 C^\infty([0, T], \mathbb{T}^2) \rightarrow C^\infty([0, T], \mathbb{T}^2)$ by

$$\Lambda(g, h) := [\Pi_{\leftarrow}(g, h) \circ \Delta \Pi_{\leftarrow}(g, h)] - \Pi_{\diamond}(g, h \circ \Delta h).$$

Then for all $\rho \in (0, 1)$, $\gamma < \mathbb{R}$, $\varepsilon > 0$ such that $2\gamma - 2 + \rho - \varepsilon > 0$ and $g \in [\lambda, 1]$, we have

$$\|\Lambda(g, h)\|_{\mathcal{C}_T^{2\gamma-2+\rho-\varepsilon}} \lesssim \|g\|_{\mathcal{L}_T^\rho} \|h\|_{C_\eta^1 \mathcal{C}_T^\gamma}^2$$

and

$$\begin{aligned} \|\Lambda(g_1, h) - \Lambda(g_2, h)\|_{\mathcal{C}_T^{2\gamma-2+\rho-\varepsilon}} &\lesssim \|g_1 - g_2\|_{C_T L^\infty} (\|g_1\|_{\mathcal{L}_T^\rho} + \|g_2\|_{\mathcal{L}_T^\rho}) \|h\|_{C_\eta^2 \mathcal{C}_T^\gamma}^2 \\ &\quad + \|g_1 - g_2\|_{\mathcal{L}_T^\rho} \|h\|_{C_\eta^1 \mathcal{C}_T^\gamma}^2. \end{aligned}$$

As a consequence Λ can be uniquely extended to a locally Lipschitz function

$$\Lambda: \mathcal{L}_T^\rho \times C_\eta^2 \mathcal{C}_T^\gamma \rightarrow \mathcal{C}_T^{2\gamma-2+\rho-\varepsilon}.$$

Proof. Let $A(t, z) := \Pi_{\leftarrow}(g, h) \circ \Delta \Pi_{\leftarrow}(g, h)(t, z)$. We can approximate $A(t, z)$ with its value for a fixed $g = g(t, z)$, to obtain

$$\begin{aligned} \Delta_q A(t, x) &= \int_z K_{q,x}(z) (\Pi_{\leftarrow}(g, h) \circ \Delta \Pi_{\leftarrow}(g, h))(t, z) \\ &= \int_z K_{q,x}(z) (\Pi_{\leftarrow}(g(t, z), h) \circ \Delta \Pi_{\leftarrow}(g(t, z), h))(t, z) \\ &\quad + \int_z K_{q,x}(z) ((\Pi_{\leftarrow}(g, h) - \Pi_{\leftarrow}(g(t, z), h)) \circ \Delta \Pi_{\leftarrow}(g(t, z), h))(t, z) \quad (15) \\ &\quad + \int_z K_{q,x}(z) (\Pi_{\leftarrow}(g, h) \circ \Delta (\Pi_{\leftarrow}(g, h) - \Pi_{\leftarrow}(g(t, z), h)))(t, z). \quad (16) \end{aligned}$$

with the remainders (15) and (16). To estimate the first term notice that

$$\Delta_j[\Delta\Pi_{\leftarrow}(g(t, z), h)](t, z) = \Delta_j\Pi_{\leftarrow}(g(t, z), \Delta h)(t, z)$$

and that, by Remark 4

$$\begin{aligned} \Delta_i\Pi_{\leftarrow}(g(t, z), h)(t, z) &= \int_{z'} K_{i,z}(z')\Pi_{\leftarrow}(g(t, z), h)(t, z') \\ &= \int_{z'} K_{i,z}(z')h(g(t, z), t, z') \\ &= \Delta_i h(g(t, z), t, \cdot)(z). \end{aligned}$$

This yields

$$\begin{aligned} (\Pi_{\leftarrow}(g(t, z), h) \circ \Delta\Pi_{\leftarrow}(g(t, z), h))(t, z) &= \sum_{i \sim j} \Delta_i\Pi_{\leftarrow}(g(t, z), h) \\ &= \sum_{i \sim j} \Delta_i h(g(t, z), t, \cdot)(z) \Delta_j[\Delta h(g(t, z), t, \cdot)](z) \\ &= \Pi_{\diamond}(g, h \circ \Delta h)(t, z). \end{aligned}$$

We proceed then to estimate (15) and (16). We obtain

$$\begin{aligned} &\int_z K_{q,x}(z)[(\Pi_{\leftarrow}(g, h) - \Pi_{\leftarrow}(g(t, z), h)) \circ \Delta\Pi_{\leftarrow}(g(t, z), h)](t, z) \\ &= \int_z K_{q,x}(z) \sum_{i \sim j \gtrsim q} (\Delta_i\Pi_{\leftarrow}(g, h)(t, z) - \Delta_i\Pi_{\leftarrow}(g(t, z), h)(t, z)) \Delta_j \Delta\Pi_{\leftarrow}(g(t, z), h)(t, z). \end{aligned}$$

Using Lemma 3 we have

$$|\Delta \Delta_j \Pi_{\leftarrow}(g(t, z), h)(t, z)| \lesssim 2^{(2-\gamma)j} \|h\|_{C_\eta \mathcal{E}_T^\gamma}.$$

Lemma 6 gives

$$|\Delta_i(\Pi_{\leftarrow}(g, h) - \Pi_{\leftarrow}(g(t, z), h))(t, z)| \lesssim 2^{-(\gamma+\rho-\varepsilon)i} \|g\|_{\mathcal{L}_T^\rho} \|h\|_{C_\eta^1 \mathcal{E}_T^\gamma},$$

and thus (15) is bounded by $2^{-(2\gamma+\rho-2-\varepsilon)q} \|g\|_{\mathcal{L}_T^\rho} \|h\|_{\mathcal{E}_T^\gamma}^2$.

We can easily bound (16) in the same way, and this proves the first inequality. For the second inequality, Lemma 3 yields

$$|\Delta_j \Delta \Pi_{\leftarrow}(g_1, h)(t, z) - \Delta_j \Delta \Pi_{\leftarrow}(g_2, h)(t, z)| \lesssim 2^{(2-\gamma)j} \|g_1 - g_2\|_{C_T L^\infty} \|h\|_{C_\eta^1 \mathcal{E}_T^\gamma},$$

and using the second inequality of Lemma 6 we obtain the desired bound.

The extension of Λ to $\mathcal{L}_T^\rho \times C_\eta^2 \mathcal{L}_T^\gamma$ is standard (see e.g. the proof of the commutator lemma in [15], Lemma 2.4). \square

Lemma 6. *With the same assumptions of Lemma 5 we have*

$$|\Delta_i \Pi_{\leftarrow}(g, h)(t, z) - \Delta_i \Pi_{\leftarrow}(g(t, z), h)(t, z)| \lesssim 2^{(\varepsilon-\rho-\gamma)i} \|g\|_{\mathcal{L}_T^\rho} \|h\|_{C_\eta^1 \mathcal{E}_T^\gamma}$$

and

$$\begin{aligned} &|\Delta_i \Pi_{\leftarrow}(g_1, h)(t, z) - \Delta_i \Pi_{\leftarrow}(g_1(t, z), h)(t, z) - \Delta_i \Pi_{\leftarrow}(g_2, h)(t, z) + \Delta_i \Pi_{\leftarrow}(g_2(t, z), h)(t, z)| \\ &\lesssim 2^{(\varepsilon-\rho-\gamma)i} [\|g_1 - g_2\|_{\mathcal{L}_T^\rho} \|h\|_{C_\eta^1 \mathcal{E}_T^\gamma} + \|g_1 - g_2\|_{C_T L^\infty} (\|g_1\|_{\mathcal{L}_T^\rho} + \|g_2\|_{\mathcal{L}_T^\rho}) \|h\|_{C_\eta^2 \mathcal{E}_T^\gamma}] \end{aligned}$$

Proof.

$$\begin{aligned}
& [\Delta_i \Pi_{\leftarrow}(g, h) - \Delta_i \Pi_{\leftarrow}(g(t, z), h)](t, z) \\
&= \sum_{k \sim i} \int_{x, y, s, \tau} K_{i, z}(x) Q_{k, t}(s) P_{k-1, x}(y) \partial_\eta \Delta_k h(\delta_\tau g_{tz}^{sy}, t, x) (\delta g_{ty}^{sy} + \delta g_{tz}^{ty}) \\
&\lesssim \sum_{k \sim i} \int_{x, y, s, \tau} |K_{i, z}(x) Q_{k, t}(s) P_{k-1, x}(y)| \|\partial_\eta \Delta_k h\|_{C_T L^\infty} |t - s|^{(\rho - \varepsilon)/2} \|g\|_{C_T^{\rho/2 - \varepsilon/2} L^\infty} \\
&\quad + \sum_{k \sim i} \int_{x, y, s, \tau} |K_{i, z}(x) Q_{k, t}(s) P_{k-1, x}(y)| \|\partial_\eta \Delta_k h\|_{C_T L^\infty} |y - z| \|g\|_{\mathcal{E}_T^\rho} \\
&\lesssim 2^{-(\rho - \varepsilon)i} 2^{-\gamma i} \|\partial_\eta h\|_{\mathcal{E}_T^\gamma} \left(\|g\|_{C_T^\rho \mathcal{E}^0} + \|g\|_{\mathcal{E}_T^\rho} \right)
\end{aligned}$$

where we used the notation $\delta g_{tz}^{sy} = g(s, y) - g(t, z)$, $\delta_\tau g_{tz}^{sy} = g(t, z) + \tau(g(s, y) - g(t, z))$ and Lemma 22. This proves the first bound.

The second inequality can be obtained in the same way with the techniques already used here and in Lemma 2. \square

2.2. Approximate paradifferential problem.

In this section we construct an approximate solution to the equation

$$(\partial_t - g \prec \Delta)u = f, \quad u(0, \cdot) = 0, \quad (17)$$

with data $f \in \mathcal{E}^{\gamma-2}$ and $g \in \mathcal{L}_T^\rho$, for some fixed $\rho, \gamma \in (0, 1)$. The idea is to obtain it via a certain class of paradifferential operators. We introduce the operator \mathcal{L} acting on functions of (η, t, x) by

$$(\mathcal{L}U)(\eta, t, x) := \partial_t U(\eta, t, x) - \eta \Delta U(\eta, t, x). \quad (18)$$

We will also use the notation $\mathcal{L}_{\eta_1} := \partial_t - \eta_1 \Delta$ with $\eta_1 \in \mathbb{R}$.

Observe that if u does not depend on η we can define

$$\Pi_{\prec}(g, \mathcal{L})u := \Pi_{\prec}(g, \mathcal{L}u) \quad (19)$$

and from definition (10) with $h = \mathcal{L}u$ we obtain $\Pi_{\prec}(g, \mathcal{L})u = \partial_t u - g \prec \Delta u$.

We can describe the commutation between the differential operator \mathcal{L} and the paraproduct $\Pi_{\leftarrow}(g, \cdot)$ via the following estimate:

Lemma 7. *Let $\rho \in (0, 1)$, $\gamma \in \mathbb{R}$. Let $U \in C_\eta^2 \mathcal{E}_T^\gamma$ and $g \in \mathcal{L}_T^\rho$ such that $g \in [\lambda, 1]$. Define*

$$\Psi(g, U) := R_1 + R_2$$

with R_1 and R_2 as in (22), (23). Then for every $\varepsilon > 0$

$$\|\Psi(g, U)\|_{\mathcal{E}_T^{\rho+\gamma-2-\varepsilon}} \lesssim (1 + \|g\|_{C_T L^\infty}) \|g\|_{\mathcal{L}_T^\rho} \|U\|_{C_\eta^1 \mathcal{E}_T^\gamma}. \quad (20)$$

Moreover, $\Psi(g, U)$ is linear in U and

$$\|\Psi(g_1, U) - \Psi(g_2, U)\|_{\mathcal{E}_T^{\rho+\gamma-2-\varepsilon}} \lesssim \|g_1 - g_2\|_{\mathcal{L}_T^\rho} (1 + \|g_1\|_{\mathcal{L}_T^\rho} + \|g_2\|_{\mathcal{L}_T^\rho}) \|U\|_{C_\eta^2 \mathcal{E}_T^\gamma}.$$

In particular, we have

$$\Psi(g, U) = \Pi_{\leftarrow}(g, \mathcal{L}U) - \Pi_{\prec}(g, \mathcal{L})\Pi_{\leftarrow}(g, U) \in C_T \mathcal{E}^{\rho+\gamma-2-\varepsilon} \quad (21)$$

whenever this expression makes sense.

Proof. We start considering $g \in C^\infty([0, T], \mathbb{T}^2)$ and $U \in C_\eta^2 C^\infty([0, T], \mathbb{T}^2)$, and prove (21) in this setting. Notice that $\Pi_{\leftarrow}(g(t, y), \mathcal{L}_{g(t, y)}U) = \mathcal{L}U(g(t, y))$. As a consequence, we can estimate

$$\begin{aligned}
& \Pi_{\leftarrow}(g, \mathcal{L}U)(t, x) - \Pi_{\leftarrow}(g, \mathcal{L}\Pi_{\leftarrow}(g, U))(t, x) \\
&= \Pi_{\leftarrow}(g, \mathcal{L}U)(t, x) - \sum_k \int_y P_{k, x}(y) (\mathcal{L}_{g(t, y)} \Delta_k \Pi_{\leftarrow}(g, U))(t, x) \\
&= \Pi_{\leftarrow}(g, \mathcal{L}U)(t, x) - \sum_k \int_y P_{k, x}(y) (\partial_t \Delta_k \Pi_{\leftarrow}(g, U))(t, x) \\
&\quad + \sum_k \int_y P_{k, x}(y) g(t, y) (\Delta \Delta_k \Pi_{\leftarrow}(g, U))(t, x) \\
&= \Pi_{\leftarrow}(g, \mathcal{L}U - \partial_t U)(t, x) + \sum_k \int_y P_{k, x}(y) g(t, y) (\Delta_k \Pi_{\leftarrow}(g, \Delta U))(t, x) \\
&\quad + \sum_k \int_y P_{k, x}(y) g(t, y) (\Delta_k [\Delta, \Pi_{\leftarrow}(g, \cdot)] U)(t, x) - \sum_k \int_y P_{k, x}(y) (\Delta_k [\partial_t, \Pi_{\leftarrow}(g, \cdot)] U)(t, x)
\end{aligned}$$

with the commutators

$$\begin{aligned}
[\Delta, \Pi_{\leftarrow}(g, \cdot)]U &:= \Delta \Pi_{\leftarrow}(g, U) - \Pi_{\leftarrow}(g, \Delta U), \\
[\partial_t, \Pi_{\leftarrow}(g, \cdot)]U &:= \partial_t \Delta_k \Pi_{\leftarrow}(g, U) - \Delta_k \Pi_{\leftarrow}(g, \partial_t U).
\end{aligned}$$

We have

$$\Pi_{\leftarrow}(g, \mathcal{L}U - \partial_t U)(t, x) + \sum_k \int_y P_{k, x}(y) g(t, y) (\Delta_k \Pi_{\leftarrow}(g, \Delta U))(t, x) = R_1(t, x)$$

with the definition

$$\begin{aligned}
R_1(t, x) &:= \\
& \sum_{k, i} \int_{y', z} P_{k, x}(y) K_{k, x}(z) P_{i, z}(y') Q_{i, t}(s) [g(t, y) - g(s, y')] \Delta \Delta_i U(g(s, y'), t, z) \quad (22)
\end{aligned}$$

and

$$\sum_k \int_y P_{k, x}(y) [g(t, y) (\Delta_k [\Delta, \Pi_{\leftarrow}(g, \cdot)] U)(t, x) - (\Delta_k [\partial_t, \Pi_{\leftarrow}(g, \cdot)] U)(t, x)] = R_2(t, x)$$

with the definition

$$\begin{aligned}
R_2(t, x) &:= \sum_{k, i} \int_{y, y', s} P_{k, x}(y) K_{k, x}(z) Q_{i, t}(s) g(t, y) \Delta P_{i, z}(y') \Delta_i U(g(s, y'), t, z) \\
&\quad + 2 \sum_{k, i} \int_{y, y', s} P_{k, x}(y) K_{k, x}(z) Q_{i, t}(s) g(t, y) \nabla P_{i, z}(y') \nabla \Delta_i U(g(s, y'), t, z) \\
&\quad - \sum_{k, i} \int_{y, y', s} P_{k, x}(y) K_{k, x}(z) \partial_t Q_{i, t}(s) P_{i, z}(y') \Delta_i U(g(s, y'), t, z). \quad (23)
\end{aligned}$$

Indeed:

$$\begin{aligned}
([\partial_t, \Pi_{\leftarrow}(g, \cdot)]U)(t, x) &= \sum_i \int_{y, s} (\partial_t Q_{i, t})(s) P_{i, x}(y) (\Delta_i U(g(s, y), t, x)), \\
([\Delta, \Pi_{\leftarrow}(g, \cdot)]U)(t, x) &= \sum_i \int_{y, s} Q_{i, t}(s) \Delta P_{i, x}(y) (\Delta_i U(g(s, y), t, x)) \quad (24) \\
&\quad + 2 \sum_i \int_{y, s} Q_{i, t}(s) \nabla P_{i, x}(y) (\nabla \Delta_i U(g(s, y), t, x)).
\end{aligned}$$

This shows that (21) holds for smooth functions.

With the techniques used in Lemma 6 we can estimate

$$|\Delta_q R_1(t, x)| \lesssim \sum_{k \sim q} \left(2^{-(\rho-\varepsilon)k} \|g\|_{C_T^{\rho/2} \mathcal{C}^0} + 2^{-\rho k} \|g\|_{\mathcal{C}_T^\rho} \right) 2^{(2-\gamma)k} \|U\|_{C_\eta \mathcal{C}_T^\gamma}.$$

By the spectral support properties of the commutators we have that

$$\begin{aligned} \|[\Delta, \Pi_{\leftarrow}(g, \cdot)]U\|_{\mathcal{C}_T^{\gamma+\rho-2}} &\lesssim \|g\|_{\mathcal{C}_T^\rho} \|U\|_{C_\eta^1 \mathcal{C}_T^\gamma}, \\ \|\Delta_q [\partial_t, \Pi_{\leftarrow}(g, \cdot)]U\|_{C_T L^\infty} &\lesssim \left(2^{(2+\varepsilon-\rho-\gamma)q} \|g\|_{C_T^{\rho/2} \mathcal{C}^0} + 2^{(2-\rho-\gamma)q} \|g\|_{\mathcal{C}_T^\rho} \right) \|U\|_{C_\eta^1 \mathcal{C}_T^\gamma}. \end{aligned}$$

This yields

$$\|R_2\|_{\mathcal{C}_T^{\gamma+\rho-2-\varepsilon}} \lesssim (1 + \|g\|_{C_T L^\infty}) \|g\|_{\mathcal{C}_T^\rho} \|U\|_{C_\eta^1 \mathcal{C}_T^\gamma}.$$

We have so far proved (20) and then (21) follows by continuity. The local Lipschitz dependence on g can be obtained via similar computations. \square

Remark 8. If f does not depend on η we consider the parametric problem

$$(\partial_t - \eta \Delta) U_f(\eta, t) = f, \quad U_f(\eta, 0) = 0, \quad \eta \in [\lambda, 1], \quad (25)$$

which is solved by

$$U_f(\eta, t) = \int_0^t e^{\eta \Delta(t-s)} f ds.$$

Remark that

$$\partial_\eta U_f(\eta, t) = \int_0^t e^{\eta \Delta(t-s)} (t-s) \Delta f ds \quad \text{and} \quad \partial_\eta^2 U_f(\eta, t) = \int_0^t e^{\eta \Delta(t-s)} (t-s)^2 \Delta^2 f ds.$$

We have, thanks to the well-known Schauder estimates of Lemma 21 (since $\eta \geq \lambda$):

$$\|U_f\|_{C_\eta^2 \mathcal{C}_T^\gamma} := \sup_{n=0,1,2} \sup_{\eta \in [\lambda, 1]} \|\partial_\eta^n U_f(\eta)\|_{\mathcal{C}_T^\gamma} \lesssim \|f\|_{\mathcal{C}_T^{\gamma-2}} \quad (26)$$

We define then

$$u(t, x) := \Pi_{\leftarrow}(g, U_f)(t, x) \quad (27)$$

and observe that $u(t, x)$ is an approximate solution of equation (17), indeed

$$(\partial_t - g \prec \Delta)u = \Pi_{\prec}(g, \mathcal{L} \Pi_{\leftarrow}(g, U_f)) = \Pi_{\leftarrow}(g, \mathcal{L} U_f) - \Psi(g, U_f) = f - \Psi(g, U_f)$$

and the estimation in Lemma 7 together with the bound (26) yield immediately the following inequality:

$$\|\Psi(g, U_f)\|_{\mathcal{C}_T^{\rho+\gamma-2-\varepsilon}} \lesssim \|g\|_{\mathcal{C}_T^\rho} (1 + \|g\|_{C_T L^\infty}) \|f\|_{\mathcal{C}_T^{\gamma-2}}. \quad (28)$$

3. PARACONTROLLED ANSATZ

In order to give a meaning to the PDE in (5) with initial condition $u_0 \in \mathcal{C}^\alpha$, our initial goal will be to get informations on solutions $\theta = \theta(g)$ of the equation

$$\partial_t \theta - g \prec \Delta \theta = \xi,$$

for a fixed $g \in \mathcal{C}_T^\alpha$, $2/3 < \alpha < 1$, $g \in [\lambda, 1]$. Using the results of Section 2.2, we consider to this effect the parametric problem

$$(\partial_t - \eta \Delta) \vartheta(\eta, t) = \xi,$$

for $\eta \in [\lambda, 1]$. We will consider the stationary solution of this problem which has the form

$$\vartheta(\eta, x) = \int_0^\infty e^{\eta \Delta s} \xi \, ds = (-\eta \Delta)^{-1} \xi \quad (29)$$

and in order for (29) to be well defined we impose that the noise ξ has zero mean on \mathbb{T}^2 (this is a simplifying assumption which can be easily removed, e.g. at the price of adding a linear term to the equation). We can control (29) by bounding its Littlewood-Paley blocks with a Bernstein lemma for distributions with compactly supported Fourier transform ([3], Lemma 2.1) to obtain:

$$\|\vartheta\|_{C_\eta^2 \mathcal{L}_T^\alpha} = \|\vartheta\|_{C_\eta^2 \mathcal{L}_T^\alpha} \lesssim \|\xi\|_{\mathcal{C}^{\alpha-2}}. \quad (30)$$

We define now for every $t \in [0, T]$

$$\theta(t, x) := \Pi_{\prec}(a(u), \vartheta).$$

Thanks to Lemma 3 we have the bound $\|\theta\|_{\mathcal{L}_T^\alpha} \lesssim \|\vartheta\|_{C_\eta \mathcal{L}_T^\alpha} \lesssim \|\xi\|_{\mathcal{C}^{\alpha-2}}$. We observe that this definition together with Lemma 7 gives

$$\partial_t \theta - a(u) \prec \Delta \theta = \xi - \Psi(a(u), \vartheta)$$

with $\|\Psi(a(u), \vartheta)\|_{\mathcal{L}_T^{2\alpha-2-\varepsilon}} \lesssim \|a(u)\|_{\mathcal{L}_T^\alpha}^2 \|\xi\|_{\mathcal{C}_T^{\alpha-2}}$. We expect then $\Psi(a(u), \vartheta)$ to be bounded in $\mathcal{L}_T^{2\alpha-2-\varepsilon}$ for any $\varepsilon > 0$. At this point let us introduce the Ansatz

$$u = \theta + u^\sharp. \quad (31)$$

Remark 9. Notice that we are not making any assumption on the existence of such u , which is the subject of Section 4. Our aim here is to find the equation that a couple $(u, u^\sharp) \in \mathcal{L}_T^\alpha \times \mathcal{L}_T^{2\alpha}$ verifying (31) must solve, in order for u to solve (5).

Observe that

$$\partial_t u - a(u) \prec \Delta u = (\partial_t - a(u) \prec \Delta) \theta + (\partial_t - a(u) \prec \Delta) u^\sharp = \xi + (\partial_t - a(u) \prec \Delta) u^\sharp - \Psi(a(u), \vartheta).$$

It follows that u^\sharp must solve

$$\begin{cases} (\partial_t - a(u) \prec \Delta) u^\sharp = \Phi(u) + \Psi(a(u), \vartheta) \\ u^\sharp(t=0) = u_0^\sharp := u_0 - \Pi_{\prec}(a(u_0), \vartheta)(t=0) \in \mathcal{L}^\alpha \end{cases} \quad (32)$$

with $\Phi(u) = a(u) \circ \Delta u + a(u) \succ \Delta u$, and if we can make sense of the resonant term $a(u) \circ \Delta u$, it is reasonable to expect $u^\sharp(t, \cdot) \in \mathcal{L}^{2\alpha} \forall t \in (0, T]$. Indeed, take $U^\sharp := U_Q$ to be the solution of

$$\mathcal{L} U^\sharp(\eta) := (\partial_t - \eta \Delta) U^\sharp(\eta) = Q \quad U^\sharp(\eta, t=0) = 0 \quad (33)$$

for some $Q = Q(u^\sharp)$ to be determined and $\eta \in [\lambda, 1]$. Using again Lemma 7 as shown in Remark 8 we have

$$(\partial_t - a(u) \prec \Delta) \Pi_{\prec}(a(u), U^\sharp) = Q(u^\sharp) - \Psi(a(u), U^\sharp).$$

For $\eta \in [\lambda, 1]$ we define $\mathcal{P}_t u_0^\sharp(\eta) := e^{\eta \Delta t} u_0^\sharp$ so that $\mathcal{L}(\mathcal{P}_t u_0^\sharp) = 0$, with \mathcal{L} as in (18).

We set

$$u^\sharp := \Pi_{\prec}(a(u), U^\sharp) + \Pi_{\prec}(a(u), \mathcal{P} u_0^\sharp). \quad (34)$$

Taking

$$Q(u^\sharp) := \Phi(u) + \Psi(a(u), \vartheta) + \Psi(a(u), U^\sharp) + \Psi(a(u), \mathcal{P} u_0^\sharp),$$

we obtain that U^\sharp solves equation (33) if and only if u^\sharp solves equation (32). As we will see, $Q(u^\sharp)(t)$ belongs to $\mathcal{C}^{2\alpha-2} \forall t \in (0, T]$ but not uniformly as $t \rightarrow 0$. However it belongs to $\mathcal{C}^{\alpha-2}$ uniformly as $t \rightarrow 0$.

It remains to control the resonant term $a(u) \circ \Delta u$ appearing in $\Phi(u)$. We have

$$a(u) \circ \Delta u = a(u) \circ \Delta \theta + a(u) \circ \Delta u^\sharp.$$

By parilinearization (see Theorem 25) $a(u) = a'(u) \prec u + R_a(u)$ with $\|R_a(u)\|_{\mathcal{C}_T^{2\alpha}} \lesssim 1 + \|u\|_{\mathcal{C}_T^\alpha}^2$, and then

$$a(u) \circ \Delta \theta = (a'(u) \prec u) \circ \Delta \theta + R_a(u) \circ \Delta \theta.$$

In order to use the commutator lemma (Lemma 26) we can estimate $a'(u)$, recalling that $\alpha \in (0, 1)$, as

$$\|a'(u)\|_{\mathcal{C}_T^\alpha} \lesssim \|a''\|_{L^\infty} \|u\|_{\mathcal{C}_T^\alpha}$$

and write

$$a(u) \circ \Delta \theta = a'(u)(u \circ \Delta \theta) + C(a'(u), u, \Delta \theta) + R_a(u) \circ \Delta \theta.$$

Then, Ansatz (31) gives

$$a(u) \circ \Delta \theta = a'(u)(\theta \circ \Delta \theta) + a'(u)(u^\sharp \circ \Delta \theta) + C(a'(u), u, \Delta \theta) + R_a(u) \circ \Delta \theta.$$

Summarizing, we have:

$$\begin{aligned} \Phi(u) &= a'(u)(\theta \circ \Delta \theta) + a(u) \succ \Delta u + a'(u)(u^\sharp \circ \Delta \theta) \\ &\quad + C(a'(u), u, \Delta \theta) + R_a(u) \circ \Delta \theta + a(u) \circ \Delta u^\sharp \end{aligned}$$

Thanks to the nonlinear commutator (Lemma 5), we can decompose the resonant term $\theta \circ \Delta \theta$ to obtain

$$\begin{aligned} \Phi(u) &= a(u) \succ \Delta u + a'(u)(u^\sharp \circ \Delta \theta) + C(a'(u), u, \Delta \theta) + R_a(u) \circ \Delta \theta + a(u) \circ \Delta u^\sharp \\ &\quad + a'(u)\Lambda(a(u), \vartheta) + a'(u)\Pi_\diamond(a(u), \Theta_2) \end{aligned}$$

and $\Lambda(a(u), \vartheta) \in \mathcal{C}_T^{3\alpha-2-\varepsilon}$ if $u \in \mathcal{L}_T^\alpha$. Here we defined

$$\Theta_2(\eta, x) := (\vartheta \circ \Delta \vartheta)(\eta, x) = \sum_{i \sim j} \Delta_i \vartheta(\eta, \cdot)(x) \Delta_j [\Delta \vartheta(\eta, \cdot)](x) \quad (35)$$

Finally, recalling the decomposition of u^\sharp in two terms (34) we obtain

$$\Phi(u) = a'(u)\Pi_\diamond(a(u), \Theta_2) + \Phi_1(u) + \Phi_2(u)$$

where

$$\begin{aligned} \Phi_1(u) &:= a(u) \succ \Delta u + C(a'(u), u, \Delta \theta) + R_a(u) \circ \Delta \theta + a'(u)\Lambda(a(u), \vartheta) \\ &\quad + a'(u)(\Pi_{\prec}(a(u), U^\sharp) \circ \Delta \theta) + a(u) \circ \Delta \Pi_{\prec}(a(u), U^\sharp), \\ \Phi_2(u) &:= a'(u)(\Pi_{\prec}(a(u), \mathcal{P}u_0^\sharp) \circ \Delta \theta) + a(u) \circ \Delta \Pi_{\prec}(a(u), \mathcal{P}u_0^\sharp). \end{aligned}$$

Thanks to Lemma 2 the terms $a'(u)\Pi_\diamond(a(u), \Theta_2)$ and $\Phi_1(u)$ can be estimated in $\mathcal{C}_T^{2\alpha-2}$, provided $\Theta_2 \in C_\eta^2 \mathcal{C}_T^{2\alpha-2}$ (see Section 5). On the other hand the term $\Phi_2(u)(t)$ can be estimated in $\mathcal{C}^{2\alpha-2}$ only for strictly positive times $t > 0$ due to the lack of regularity of the initial condition u_0^\sharp which a priori lives only in \mathcal{C}^α .

Note moreover that the specific form of Φ allows to deduce that if we replace Θ_2 by $\tilde{\Theta}_2 = \Theta_2 - H$ with $H \in C_\eta^2 \mathcal{C}_T^{2\alpha-2}$ then this is equivalent to consider an equation for u of the form

$$\partial_t u(t, x) - a(u(t, x))\Delta u(t, x) = \xi(x) - a'(u(t, x))H(a(u(t, x)), t, x).$$

Let us resume this long discussion in the following theorem:

Theorem 10. *Assume that $\xi \in \mathcal{C}^0$, $u_0 \in \mathcal{C}^2$, $H \in C_\eta^2 \mathcal{C}_T^0$. $u \in C_T^1 \mathcal{C}^2$ is the classical solution to the equation*

$$\partial_t u(t, x) - a(u(t, x)) \Delta u(t, x) = \xi(x) - a'(u(t, x)) H(a(u(t, x))), t, x, \quad u(0) = u_0, \quad (36)$$

up to time $T > 0$ if

$$u = \Pi_{\leftarrow}(a(u), \vartheta + U^\sharp + \mathcal{P} u_0^\sharp),$$

where ϑ is the solution to eq. (29) and U^\sharp is the solution to the PDE

$$(\partial_t - \eta \Delta) U^\sharp(\eta) = F(u, U^\sharp, u_0^\sharp) \quad U^\sharp(\eta, 0) = 0 \quad \eta \in [\lambda, 1] \quad (37)$$

with

$$\begin{aligned} F(u, U^\sharp, u_0) &= a'(u) \Pi_{\diamond}(a(u), \Theta_2) + \Phi_1(u) + \Phi_2(u) + \Psi(a(u), \vartheta) + \Psi(a(u), U^\sharp) \\ &\quad + \Psi(a(u), \mathcal{P} u_0^\sharp) \end{aligned}$$

and $\Theta_2 = \vartheta \circ \Delta \vartheta - H$.

Definition 11. *For any $\alpha \in \mathbb{R}$ we define $\mathcal{X}^\alpha \subseteq C_\eta^2 \mathcal{C}^\alpha \times C_\eta^2 \mathcal{C}^{2\alpha-2}$ the closure of the image of the map*

$$(\rho, H) \in C_\eta^2 \mathcal{C}^2 \times C_\eta^2 \mathcal{C}^0 \mapsto J(\rho, H) = (\rho, \rho \circ \Delta \rho - H) \in C_\eta^2 \mathcal{C}^2 \times C_\eta^2 \mathcal{C}^0$$

(in the topology of $C_\eta^2 \mathcal{C}^\alpha \times C_\eta^2 \mathcal{C}^{2\alpha-2}$).

We call the elements in \mathcal{X}^α *enhanced noises*. In the next section we will exploit the space \mathcal{X}^α for $2/3 < \alpha < 1$ to solve equations (37) and (31).

4. LOCAL WELLPOSEDNESS

The main result of this section is the local well-posedness for equations (31) and (37) when $(\vartheta, \Theta_2) \in \mathcal{X}^\alpha$ and $u_0 \in \mathcal{C}^\alpha$ for $2/3 < \alpha < 1$. This yields a unique solution to (36), thanks to Theorem 10.

Theorem 12. *Let $\alpha > 2/3$. Then for any $(\vartheta, \Theta_2) \in \mathcal{X}^\alpha$ and $u_0 \in \mathcal{C}^\alpha$ there exists a time $T > 0$ depending only on $\|(\vartheta, \Theta_2)\|_{\mathcal{X}^\alpha}$ and $\|u_0\|_\alpha$ up to which the system of equations (31) and (37) has a unique solution $(u, U^\sharp) \in \mathcal{L}_T^\alpha \times C_\eta^2 \mathcal{L}_T^{2\delta}$ for all $\delta < \alpha$ such that $2\delta + \alpha > 2$. For any fixed $\tau > 0$ there exist a ball $B_\tau \subseteq \mathcal{C}^\alpha \times \mathcal{X}^\alpha$ such that the solution map*

$$\Sigma_\tau: (u_0, \vartheta, \Theta_2) \in B_\tau \mapsto (u, U^\sharp) \in \mathcal{L}_\tau^\alpha \times C_\eta^2 \mathcal{L}_\tau^{2\delta}$$

is well defined and Lipschitz continuous in the data.

Remark 13. The proof is based on a Picard fixed point argument. In order to obtain a small-time contraction we carry on our estimates of U^\sharp in the space $C_\eta^2 \mathcal{L}_T^{2\delta} \supset C_\eta^2 \mathcal{L}_T^{2\alpha}$.

Proof. (Theorem 12) Let $\mathcal{G}_T = \mathcal{L}_T^\alpha \times C_\eta^2 \mathcal{L}_T^{2\delta}$. We introduce the map

$$\Gamma: (u, U^\sharp) \in \mathcal{G}_T \mapsto (\Gamma_u(u, U^\sharp), \Gamma_{U^\sharp}(u, U^\sharp)) \in \mathcal{G}_T$$

by

$$\Gamma_u(u, U^\sharp) := \Pi_{\leftarrow}(a(u), \vartheta) + \Pi_{\leftarrow}(a(u), \Gamma_{U^\sharp}(u, U^\sharp)) + \Pi_{\leftarrow}(a(u), \mathcal{P} u_0^\sharp)$$

and

$$(\partial_t - \eta \Delta) \Gamma_{U^\sharp}(u, U^\sharp)(\eta) = F(u, U^\sharp, u_0^\sharp), \quad \Gamma_{U^\sharp}(u, U^\sharp)(\eta)(0) = 0, \quad \eta \in [\lambda, 1],$$

We will establish that this map is a contraction in the space \mathcal{G}_T .

First, we have to show that there exists a ball $B \subset \mathcal{G}_T$ such that $\Gamma(B) \subseteq B$. We have the bound $\|\mathcal{P}u_0^\sharp\|_{C_\eta^1 \mathcal{L}_T^\alpha} \lesssim \|u_0^\sharp\|_{\mathcal{L}_T^\alpha}$. It is easy to obtain, using the estimates of Section 2 and Lemma 21:

$$\begin{aligned} & \left\| \int_0^T e^{-\eta\Delta(t-s)} [\Phi_1(u) + \Psi(a(u), \vartheta) + \Psi(a(u), U^\sharp) + \Psi(a(u), \mathcal{P}u_0^\sharp)]_s ds \right\|_{C_\eta^2 \mathcal{L}_T^{2\delta}} \\ & \lesssim T^\kappa (1 + \|u\|_{\mathcal{L}_T^\alpha})^4 (1 + \|\xi\|_{\mathcal{L}^{\alpha-2}})^2 \|u_0^\sharp\|_{\mathcal{L}^\alpha} (1 + \|U^\sharp\|_{C_\eta^2 \mathcal{L}_T^{2\delta}}) \end{aligned}$$

for some $\kappa > 0$.

By the assumption that $(\vartheta, \Theta_2) \in \mathcal{X}^\alpha$ we deduce that there exists $M > 0$ such that $\|\Theta_2\|_{C_\eta^2 \mathcal{L}_T^{2\alpha-2}} \leq M$. We have

$$\left\| \int_0^T e^{-\eta\Delta(t-s)} [a'(u)\Pi_\diamond(a(u), \Theta_2)]_s ds \right\|_{C_\eta^2 \mathcal{L}_T^{2\delta}} \lesssim T^{\alpha-\delta} (1 + \|u\|_{\mathcal{L}_T^\alpha})^2 \|\Theta_2\|_{C_\eta^2 \mathcal{L}_T^{2\alpha-2}}.$$

To bound the term $\Phi_2(u)$ we observe that $\|\mathcal{P}tu_0^\sharp\|_{C_\eta^2 \mathcal{L}_T^{2\alpha}} \lesssim t^{-\frac{\alpha}{2}} \|u_0^\sharp\|_{\mathcal{L}^\alpha}$ thanks to Lemma 21. This gives

$$\left\| \int_0^T e^{-\eta\Delta(t-s)} \Phi_2(u)_s ds \right\|_{C_\eta^2 \mathcal{L}_T^{2\delta}} \lesssim T^{\alpha-\delta} (1 + \|u\|_{\mathcal{L}_T^\alpha}) (1 + \|\xi\|_{\mathcal{L}^{\alpha-2}}) \|u_0^\sharp\|_{\mathcal{L}^\alpha}$$

and then $\Gamma_{U^\sharp}(u, U^\sharp)$ is bounded in $C_\eta^2 \mathcal{L}_T^{2\delta}$ for T small enough. We have also

$$\begin{aligned} \|\Gamma_u(u, U^\sharp)\|_{\mathcal{L}_T^\alpha} & \lesssim \|\xi\|_{\mathcal{L}^{\alpha-2}} + \|u_0^\sharp\|_{\mathcal{L}^\alpha} + \|\Gamma_{U^\sharp}(u, U^\sharp)\|_{C_\eta \mathcal{L}_T^\alpha} \\ & \lesssim \|\xi\|_{\mathcal{L}^{\alpha-2}} + \|u_0^\sharp\|_{\mathcal{L}^\alpha} + T^{\frac{2\delta-\alpha}{2}} \|\Gamma_{U^\sharp}(u, U^\sharp)\|_{C_\eta^2 \mathcal{L}_T^{2\delta}} \end{aligned}$$

and these bounds show that $\Gamma(B) \subseteq B$. The contractivity of $\Gamma_{U^\sharp}(u, U^\sharp)$ can be obtained in the same way. Now consider $\Gamma_u(u, U^\sharp)$: we have

$$\begin{aligned} & \|\Pi_{\leftarrow}(a(u_1), U_1^\sharp) - \Pi_{\leftarrow}(a(u_2), U_2^\sharp)\|_{\mathcal{L}_T^\alpha} \\ & \lesssim T^{\frac{2\delta-\alpha}{2}} (\|U_1^\sharp - U_2^\sharp\|_{C_\eta \mathcal{L}_T^{2\delta}} + \|u_1 - u_2\|_{C_T L^\infty} \|U_2^\sharp\|_{C_\eta^1 \mathcal{L}_T^{2\delta}}) \end{aligned}$$

while for the other terms in $\Gamma_u(u_1, U_1^\sharp) - \Gamma_u(u_2, U_2^\sharp)$ we remark that

$$\sup_{s \in [0, t]} \|u_{1,s} - u_{2,s} + u_0\|_{L^\infty} \lesssim t^{\varepsilon/2} \|u_1 - u_2\|_{C_{[0, t]}^{\varepsilon/2} L^\infty}.$$

Then $\forall 0 < \varepsilon < \alpha$, using Lemma 3 and Lemma 22:

$$\begin{aligned} \|\Pi_{\leftarrow}(a(u_1), \vartheta) - \Pi_{\leftarrow}(a(u_2), \vartheta)\|_{\mathcal{L}_T^\alpha} & \lesssim \|a(u_1) - a(u_2)\|_{C_T L^\infty} \|\vartheta\|_{C_\eta^1 \mathcal{L}_T^\alpha} \\ & \lesssim \|u_1 - u_2\|_{C_T L^\infty} \|\xi\|_{\mathcal{L}^{\alpha-2}} \\ & \lesssim T^{\varepsilon/2} \|u_1 - u_2\|_{C_T^{\varepsilon/2} L^\infty} \|\xi\|_{\mathcal{L}^{\alpha-2}} \\ & \lesssim T^{\varepsilon/2} \|u_1 - u_2\|_{\mathcal{L}_T^\alpha} \|\xi\|_{\mathcal{L}^{\alpha-2}}. \end{aligned}$$

With the same reasoning we estimate

$$\begin{aligned} \|\Pi_{\leftarrow}(a(u_1), \mathcal{P}u_0^\sharp) - \Pi_{\leftarrow}(a(u_2), \mathcal{P}u_0^\sharp)\|_{\mathcal{L}_T^\alpha} & \lesssim T^{\varepsilon/2} \|u_1 - u_2\|_{C_T^{\varepsilon/2} L^\infty} \|\mathcal{P}u_0^\sharp\|_{C_\eta^1 \mathcal{L}_T^\alpha} \\ & \lesssim T^{\varepsilon/2} \|u_1 - u_2\|_{\mathcal{L}_T^\alpha} \|u_0^\sharp\|_{\mathcal{L}^\alpha} \end{aligned}$$

and then Γ is a contraction for small times.

The uniqueness of the solution $(u, U^\sharp) \in \mathcal{L}_T^\alpha \times C_\eta^2 \mathcal{L}_T^{2\delta}$ and the Lipschitz continuity of the localized solution map Σ_τ can be proved along the same lines via standard arguments. \square

5. RENORMALIZATION

At this point we want to construct an enhanced noise Ξ associated to the white noise ξ . Already in the standard setting of the generalised PAM model with constant diffusion matrix, the construction of the enhancement requires a renormalization since the resonant product $\vartheta \circ \Delta \vartheta$ is not well defined.

Let $\psi \in \mathcal{S}(\mathbb{T}^2)$ be a cutoff function and let $\psi_\varepsilon(x) = \varepsilon^{-2}\psi(x/\varepsilon)$. Then define a regularised noise by $\xi_\varepsilon = \psi_\varepsilon * \xi$ and let $\vartheta_\varepsilon = (-\eta\Delta)^{-1}\xi_\varepsilon$. Notice that

$$H_\varepsilon(\eta) := \mathbb{E}[\vartheta_\varepsilon(\eta, x) \circ \Delta \vartheta_\varepsilon(\eta, x)] = \mathbb{E}[\vartheta_\varepsilon(\eta, x) \Delta \vartheta_\varepsilon(\eta, x)] = - \sum_{k \in \mathbb{Z}^2 \setminus \{0\}} \frac{\hat{\psi}_\varepsilon(k)^2}{\eta^2 |k|^2} = - \frac{\sigma_\varepsilon}{\eta^2}$$

where

$$\sigma_\varepsilon := \sum_{k \in \mathbb{Z}^2 \setminus \{0\}} \frac{\hat{\psi}_\varepsilon(k)^2}{|k|^2} \simeq |\log \varepsilon|$$

as $\varepsilon \rightarrow 0$. Subtracting the diverging quantity H_ε to $\vartheta_\varepsilon \circ \Delta \vartheta_\varepsilon$ and then taking the limit as $\varepsilon \rightarrow 0$ delivers a finite result.

Theorem 14. *Take $\alpha < 1$ and let $\Xi_\varepsilon = (\xi_\varepsilon, \Xi_{2,\varepsilon}) := (\xi_\varepsilon, \vartheta_\varepsilon \circ \Delta \vartheta_\varepsilon - H_\varepsilon)$. Then the family $(\Xi_\varepsilon)_\varepsilon \subseteq \mathcal{X}^\alpha$ converges a.s. and in L^p to a random element $\Xi = (\xi, \Xi_2) \in \mathcal{X}^\alpha$.*

Proof. The proof is a mild modification of the proof for PAM [15]. In order to establish the required $C_\eta^2 \mathcal{C}_T^{2\alpha-2}$ regularity for Ξ_2 we follow the computations for the case where the diffusion coefficient is constant. We have only to discuss the additional regularity in the η parameter. In order to do so observe that

$$\Xi_{2,\varepsilon}(\eta) = \sum_{i \sim j} \llbracket \Delta_i \vartheta_\varepsilon(\eta) \Delta_j \Delta \vartheta_\varepsilon(\eta) \rrbracket$$

where $\llbracket \cdot \rrbracket$ denotes the Wick product with respect to the Gaussian structure of ξ . Then we have

$$\partial_\eta \Xi_{2,\varepsilon}(\eta) = \sum_{i \sim j} \llbracket \Delta_i \partial_\eta \vartheta_\varepsilon(\eta) \Delta_j \Delta \vartheta_\varepsilon(\eta) \rrbracket + \sum_{i \sim j} \llbracket \Delta_i \vartheta_\varepsilon(\eta) \Delta_j \Delta \partial_\eta \vartheta_\varepsilon(\eta) \rrbracket,$$

and

$$\begin{aligned} \partial_\eta^2 \Xi_{2,\varepsilon}(\eta) &= \sum_{i \sim j} \llbracket \Delta_i \partial_\eta^2 \vartheta_\varepsilon(\eta) \Delta_j \Delta \vartheta_\varepsilon(\eta) \rrbracket + \sum_{i \sim j} \llbracket \Delta_i \vartheta_\varepsilon(\eta) \Delta_j \Delta \partial_\eta^2 \vartheta_\varepsilon(\eta) \rrbracket \\ &\quad + \sum_{i \sim j} 2 \llbracket \Delta_i \partial_\eta \vartheta_\varepsilon(\eta) \Delta_j \Delta \partial_\eta \vartheta_\varepsilon(\eta) \rrbracket. \end{aligned}$$

Now the computations relative to the regularities of these additional stochastic objects are equivalent to those for the term $\Xi_{2,\varepsilon}$ where one or two instances of $\vartheta_\varepsilon(\eta)$ are replaced by Gaussian fields of similar regularities of the form $\partial_\eta \vartheta_\varepsilon(\eta)$ and $\partial_\eta^2 \vartheta_\varepsilon(\eta)$, a direct inspection of the proof allows us to deduce that we have almost sure $\mathcal{C}^{2\alpha-2}$ regularity for these terms and also for random fields $\partial_\eta^n \Xi_{2,\varepsilon}$ for any finite n . This allows also to deduce that the random field is a.s. smooth in the parameter η . Similar computations allow to prove continuity in ε for $\varepsilon > 0$. The rest of the proof is standard. \square

In conclusion we see that in order to be able to use this convergence result we need to modify our approximate PDE and consider instead

$$\partial_t u_\varepsilon - a(u_\varepsilon) \Delta u_\varepsilon = \xi_\varepsilon - a'(u_\varepsilon) H_\varepsilon(a(u_\varepsilon))$$

which gives the renormalised equation (4).

Our well-posedness results for the paracontrolled formulation of this equation together with the convergence result in Theorem 14 allow to deduce that $u_\varepsilon \rightarrow u$ in \mathcal{C}_T^δ for any $2/3 < \delta < \alpha < 1$ and that the limiting process u satisfies a modified version of eq. (1), namely

$$\partial_t u - a(u) \diamond \Delta u = \xi, \quad u(0) = u_0,$$

where $a(u) \diamond \Delta u$ denotes a *renormalized* diffusion term given by

$$a(u) \diamond \Delta u := a(u) \prec \Delta u + a'(u) \Pi_\diamond(a(u), \Xi_2) + \Phi_1(u) + \Phi_2(u). \quad (38)$$

6. NONLINEAR SOURCE TERMS

Let us start by discussing the presence of a u dependent r.h.s. in eq. (1). We want to solve

$$\partial_t u - a_1(u) \Delta u = a_2(u) \xi$$

where a_1 is a non-linear diffusion coefficient as before and $a_2: \mathbb{R} \rightarrow \mathbb{R}$ is another bounded function with sufficiently many bounded derivatives. We rewrite this equation as

$$\Pi_\prec(a(u), \mathcal{L})u = a_2(u) \prec \xi + a_1(u) \circ \Delta u + a_2(u) \circ \xi + a_1(u) \succ \Delta u + a_2(u) \succ \xi$$

where now $a(u) = (a_1(u), a_2(u))$ is a vector valued non-linearity. Since we don't need u to depend on any parameter $\eta = (\eta_1, \eta_2)$, we have defined \mathcal{L} as

$$\mathcal{L}(\eta) := \partial_t - \eta_1 \Delta$$

and used the identity $\Pi_\prec(a(u), \mathcal{L})u = (\partial_t - a_1(u) \prec \Delta)u$, similarly to what we have done in (19).

Notice that the non-linear paraproduct can be extended trivially to the vector valued case in such a way that, for example,

$$\Pi_\prec((g_1, g_2), h)(t, x) = \sum_i \int_{y,s} Q_{i,t}(s) P_{i,x}(y) (\Delta_i h((g_1(s, y), g_2(s, y)), t, \cdot))(x).$$

As before we make the Ansatz

$$u = \Pi_\prec(a(u), \vartheta) + u^\sharp$$

where now ϑ solves

$$\mathcal{L}(\eta) \vartheta(\eta) = (\partial_t - \eta_1 \Delta) \vartheta(\eta) = \eta_2 \xi,$$

for $\eta = (\eta_1, \eta_2) \in [\lambda, 1] \times [-L, L]$ where L is a large but fixed constant. The bounded domain is important to be able to have uniform estimates and reuse the estimates proved above in the simple situation of $\eta_2 = 1$. The solution of this equation is

$$\vartheta(\eta, \cdot) = \eta_2 \int_0^\infty e^{\eta_1 \Delta s} \xi ds = -\frac{\eta_2}{\eta_1} \Delta^{-1} \xi.$$

Observe that

$$\Pi_\prec(a(u), \mathcal{L})u = \Pi_\prec(a(u), \mathcal{L}) \Pi_\prec(a(u), \vartheta) + \Pi_\prec(a(u), \mathcal{L})u^\sharp$$

and recall that (Lemma 7) $\Pi_\prec(a(u), \mathcal{L}) \Pi_\prec(a(u), \vartheta) = \Pi_\prec(a(u), \mathcal{L} \vartheta) + \Psi(a(u), \vartheta)$. Now

$$(\mathcal{L} \vartheta)(\eta) = (\partial_t - \eta_1 \Delta) \vartheta(\eta, t, x) = \Xi(\eta), \quad \eta = (\eta_1, \eta_2) \in [\lambda, 1] \times [-L, L]$$

with $\Xi(\eta)(t, x) = \eta_2 \xi(x)$ and then

$$\Pi_\prec(a(u), \mathcal{L} \vartheta) = \Pi_\prec(a(u), \Xi) = a_2(u) \prec \xi.$$

In conclusion

$$\Pi_\prec(a(u), \mathcal{L}) \Pi_\prec(a(u), \vartheta) = a_2(u) \prec \xi + \Psi(a(u), \vartheta)$$

and the equation for u^\sharp reads

$$\begin{aligned} \Pi_{\prec}(a(u), \mathcal{L}u^\sharp) &= a_1(u) \circ \Delta u + a_2(u) \circ \xi + [a_2(u) \prec \xi - a_2(u) \ll \xi] \\ &\quad + a_1(u) \succ \Delta u + a_2(u) \succ \xi - \Psi(a(u), \vartheta) \end{aligned}$$

where now all the terms on the r.h.s. can be treated as remainder terms. Let us just remark that the commutation term $a_2(u) \prec \xi - a_2(u) \ll \xi$ has a standard treatment via Lemma 27. Of course, the first two terms require to be treated as resonant terms. Note that, modulo terms of order $\mathcal{E}_T^{3\alpha-2}$ (or $\mathcal{E}^{\alpha/2}\mathcal{E}^{2\alpha-2}$ as defined in Lemma 21) the terms $a_1(u) \circ \Delta u + a_2(u) \circ \xi$ are equivalent to

$$a_1'(u)\Pi_{\diamond}(a(u), \vartheta \circ \Delta\vartheta) + a_2'(u)(\Pi_{\ll}(a(u), \vartheta) \circ \xi)$$

and that by computations similar to those of the previous sections one can prove that

$$(\Pi_{\ll}(a(u), \vartheta) \circ \xi) = \Pi_{\diamond}(a(u), \vartheta \circ \xi) + \mathcal{E}_T^{3\alpha-2}$$

so the resonant terms are comparable to the sum of the two terms

$$a_1'(u)\Pi_{\diamond}(a(u), \vartheta \circ \Delta\vartheta) + a_2'(u)\Pi_{\diamond}(a(u), \vartheta \circ \xi)$$

which require renormalization of the form

$$\frac{a_1'(u)a_2(u)^2}{a_1(u)^2}\sigma_\varepsilon - \frac{a_2'(u)a_2(u)}{a_1(u)}\sigma_\varepsilon \quad (39)$$

and the convergence follows with the same arguments of Section 5.

We remark that the structure of the second renormalisation term, which is due to the r.h.s. in the equation, is the same of that found by Bailleul, Debussche and Hofmanova in [7].

Remark 15. Our approach works straightforwardly for the equation

$$\partial_t u(t, x) - a_{ij}(u(t, x))\partial_{ij}^2 u(t, x) = g(u(t, x))\xi$$

with $a: \mathbb{R} \rightarrow M_2(\mathbb{R})$ such that $\sum_{i,j} a(u)_{ij} x_i x_j \geq C|x|^2 \forall x \in \mathbb{R}^2$ for $C > 0$ and $\partial_{ij}^2 := \frac{\partial^2}{\partial x_i \partial x_j}$.

To see that, let $a(u) := (a_{ij}(u), g(u)) \in \mathbb{R}^5$ and $\eta = (\eta_{i,j}, \eta_g) \in \mathbb{R}^5$. Let $\mathcal{L}(\eta) := \partial_t - \eta_{ij}\partial_{ij}^2$ and $\Xi(\eta) := \eta_g \xi$ with the uniform ellipticity condition $\sum_{i,j} \eta_{ij} x_i x_j \geq C|x|^2 \forall x \in \mathbb{R}^2$. It is easy to verify that Lemma 7 and Lemma 5 hold within this setting, just considering nonlinear para-products for functions depending on 5 parameters. We have then:

$$u = \Pi_{\ll}(a(u), \vartheta + U^\sharp + \mathcal{P}u_0^\sharp)$$

with $\vartheta(\eta)$ stationary solution of $\mathcal{L}\vartheta(\eta) = \Xi(\eta)$, $\mathcal{P}_t u_0^\sharp := e^{\eta_{ij}\partial_{ij}^2 t} u_0^\sharp$ and $U^\sharp(\eta)$ which solves

$$\mathcal{L}U^\sharp(\eta) = \Pi_{\diamond}((a(u), a'(u)), \Theta_1) + \Pi_{\diamond}((a(u), a'(u)), \Theta_2) + Q(u, U^\sharp)$$

with $Q(u, U^\sharp) \in \mathcal{E}^{2\alpha-2-\varepsilon}$, $\Theta_1 \in C_\eta^2 C_\eta^2 \mathcal{E}^{2\alpha-2} = \vartheta(\eta) \circ \eta'_{ij} \partial_{ij}^2 \vartheta(\eta)$, $\Theta_2(\eta, \eta') = \vartheta(\eta) \circ \eta'_g \xi$ and $U^\sharp(t=0) = 0$. Note that we can write ϑ as

$$\vartheta(\eta) = \eta_g \int_0^\infty e^{t\eta_{ij}\partial_{ij}^2} \xi dt \quad \hat{\vartheta}(k) = \eta_g \frac{\hat{\xi}(k)}{\eta_{ij} k_i k_j}, \quad k \in \mathbb{Z}^2 \setminus \{0\}.$$

From the uniform ellipticity condition we have that $\|\vartheta\|_{C_\eta^k \mathcal{E}^\alpha} \lesssim \|\xi\|_{\mathcal{E}^\alpha}$, and Schauder estimates analogous to those of Lemma 21 hold as well.

Now consider the renormalization. We have

$$\begin{aligned} H_1^\varepsilon(\eta, \eta') &:= \mathbb{E}(\Theta_1(\eta, \eta')) = -\eta_g^2 \sum_{k \in \mathbb{Z}^2 \setminus \{0\}} \hat{\psi}_\varepsilon(k)^2 \frac{\sum_{i,j} \eta'_{ij} k_i k_j}{(\sum_{i,j} \eta_{ij} k_i k_j)^2}, \\ H_2^\varepsilon(\eta, \eta') &:= \mathbb{E}(\Theta_2(\eta, \eta')) = \eta_g \eta'_g \sum_{k \in \mathbb{Z}^2 \setminus \{0\}} \frac{\hat{\psi}_\varepsilon(k)^2}{\sum_{i,j} \eta_{ij} k_i k_j}. \end{aligned}$$

The convergence of $\Theta_1^\varepsilon - H_1^\varepsilon$, $\Theta_2^\varepsilon - H_2^\varepsilon$ in $C_{(\eta, \eta')}^k \mathcal{C}^{2\alpha-2}(\mathbb{T}^2)$ can be obtained with the techniques used in [15], Section 5.2.

7. FULL GENERALITY

Within the framework of the present work we are actually able to treat equations of the form

$$\partial_t u(t, x) - a_1(u(t, x)) \Delta u(t, x) = \xi(a_2(u(t, x)), x) \quad (40)$$

where $\xi(\eta_2, x)$ is a Gaussian process with covariance

$$\mathbb{E}[\xi(\eta_2, x) \xi(\tilde{\eta}_2, \tilde{x})] = F(\eta_2, \tilde{\eta}_2) \delta(x - \tilde{x})$$

where F is a smooth covariance function. Let as before $2/3 < \alpha < 1$. In this case we can take as a parametric equation

$$\mathcal{L}(\eta) \vartheta := \partial_t \vartheta(\eta, t, x) - \eta_1 \Delta \vartheta(\eta, t, x) = \xi(\eta_2, x)$$

whose solution ϑ is a Gaussian process, smooth with respect to the variable $\eta = (\eta_1, \eta_2)$ which we assume taking value in a compact subset of \mathbb{R}^2 for which $\eta_1 \geq \lambda > 0$ with fixed λ . Letting $a(u) = (a_1(u), a_2(u))$ we can rewrite the l.h.s. of eq. (40) in the form

$$\partial_t u - a_1(u) \Delta u = \Pi_{\diamond}(a(u), \mathcal{L}u)$$

and the r.h.s. as

$$\xi(a_2(u(t, x)), x) = \Pi_{\diamond}(a(u), \Xi)$$

where $\Xi(\eta, x) = \xi(\eta_2, x)$. Now we perform the paraproduct decomposition to get

$$\begin{aligned} \Pi_{\prec}(a(u), \mathcal{L}u) - \Pi_{\prec}(a(u), \Xi) &= \Pi_{\circ}(a(u), \Xi) + \Pi_{\circ}(a(u), \mathfrak{D}u) \\ &\quad + \Pi_{\succ}(a(u), \Xi) + \Pi_{\succ}(a(u), \mathfrak{D}u). \end{aligned}$$

with $\mathfrak{D}(\eta) := \eta_1 \Delta$. Let $\mathcal{P}_t(\eta) := e^{t\eta_1 \Delta}$ as before, and invoke the paracontrolled Ansatz in the usual form

$$u = \Pi_{\prec}(a(u), \vartheta + U^\sharp + \mathcal{P}u_0^\sharp).$$

Using that

$$\begin{aligned} \Pi_{\prec}(a(u), \mathcal{L} \Pi_{\prec}(a(u), \vartheta + U^\sharp + \mathcal{P}u_0^\sharp)) &= \Pi_{\prec}(a(u), \mathcal{L}(\vartheta + U^\sharp + \mathcal{P}u_0^\sharp)) \\ &\quad + \Psi(a(u), \vartheta + U^\sharp + \mathcal{P}u_0^\sharp) \end{aligned}$$

and observing that we can take $\mathcal{L}\vartheta = \Xi$ and that $\mathcal{L}\mathcal{P}u_0^\sharp = 0$ to get

$$\Pi_{\prec}(a(u), \mathcal{L}U^\sharp) = F(u, U^\sharp)$$

where

$$\begin{aligned} F(u, U^\sharp) &= \Pi_{\circ}(a(u), \Xi) + \Pi_{\circ}(a(u), \mathfrak{D}u) + \Pi_{\succ}(a(u), \Xi) + \Pi_{\succ}(a(u), \mathfrak{D}u) \\ &\quad + [\Pi_{\prec}(a(u), \Xi) - \Pi_{\prec}(a(u), \Xi)] - \Psi(a(u), \vartheta + U^\sharp + \mathcal{P}u_0^\sharp) \end{aligned}$$

which is solved by U^\sharp satisfying

$$\mathcal{L}U^\sharp = F(u, U^\sharp).$$

Indeed $\Pi_{\prec}(a(u), F(u, U^\sharp)) = F(u, U^\sharp)$, since $F(u, U^\sharp)$ does not depend on the additional parameter. Remark that the term $\Pi_{\prec}(a(u), \Xi) - \Pi_{\leftarrow}(a(u), \Xi)$, which does not appear in the simpler case, can be treated with Lemma 28.

It remains now to discuss the handling of the resonant products under the paracontrolled assumption, namely $\Pi_{\circ}(a(u), \Xi)$ and $\Pi_{\circ}(a(u), \mathfrak{D}u)$. Next lemma is a parilinearization result adapted to our non-linear context.

Lemma 16. *Assume that $u \in \mathcal{C}_T^\rho$ and $Z \in C_\eta^2 \mathcal{C}_T^\gamma$ then if $\gamma + 2\rho > 0$ we have*

$$C(u, Z) := \Pi_{\circ}(a(u), Z) - u \circ \Pi_{\prec}((a(u), a'(u)), \mathcal{D}Z) \in \mathcal{C}_T^{\gamma+2\rho}$$

where $\mathcal{D}Z((\eta, \eta'), t, x) := \sum_i \eta_i' \partial_{\eta_i} Z(\eta, t, x)$.

Proof.

$$\begin{aligned} \Pi_{\circ}(a(u), Z)(t, x) &= \sum_{i \sim j} \int_{y, z} K_{i, x}(y) K_{j, x}(z) Z(a(u(t, y)), t, z) \\ &= \sum_{i \sim j} \sum_k \int_{y, z, z', z''} K_{i, x}(y) K_{j, x}(z) P_{k, z}(z'') K_{k, z}(z') Z(a(u(t, y)), t, z') \\ &= \sum_{i \sim j} \sum_k \int_{y, z, z', z''} K_{i, x}(y) K_{j, x}(z) P_{k, z}(z'') K_{k, z}(z') [Z(a(u(t, y)), t, z') - Z(a(u(t, z'')), t, z')] \\ &= \sum_{i \sim j} \sum_k \int_{y, z, z', z''} K_{i, x}(y) K_{j, x}(z) P_{k, z}(z'') K_{k, z}(z') \times \\ &\quad \times \left[\sum_{\ell} a_{\ell}'(u(t, z'')) \delta u_{t, z''}^{\ell y} \partial_{a_{\ell}} Z(a(u(t, z'')), t, z') \right] \\ &\quad + \sum_{i \sim j} \sum_{k \sim j} \int_{\substack{z', z'' \\ y, z}} K_{i, x}(y) K_{j, x}(z) P_{k, z}(z'') K_{k, z}(z') O((\delta u_{t, z''}^{\ell y})^2) \partial_{\eta}^2 Z(a(u(t, y)), t, z') \end{aligned}$$

and observe that the first term is equal to $u \circ \Pi_{\prec}((a(u), a'(u)), \mathcal{D}Z)$ while the second term can be easily estimated in $\mathcal{C}_T^{\gamma+2\rho}$. \square

Using this result and Lemma 28 we can expand

$$\begin{aligned} \Pi_{\circ}(a(u), \Xi) &= u \circ \Pi_{\prec}((a(u), a'(u)), \mathcal{D}\Xi) + \mathcal{C}_T^{3\alpha-2} \\ &= \Pi_{\leftarrow}(a(u), \vartheta) \circ \Pi_{\prec}((a(u), a'(u)), \mathcal{D}\Xi) + \mathcal{C}_T^{3\alpha-2} \\ &= \Pi_{\diamond}((a(u), a'(u)), \vartheta \circ \mathcal{D}\Xi) + \mathcal{C}_T^{3\alpha-2} \end{aligned}$$

and similarly, noting that

$$\begin{aligned} \Pi_{\prec}((a(u), a'(u)), (\mathcal{D}\mathfrak{D})u) &= \Pi_{\prec}((a(u), a'(u)), (\mathcal{D}\mathfrak{D})\Pi_{\leftarrow}(a(u), \vartheta)) + \mathcal{C}_T^{3\alpha-2} \\ &= \Pi_{\prec}((a(u), a'(u)), (\mathcal{D}\mathfrak{D})\vartheta) + \mathcal{C}_T^{3\alpha-2} \end{aligned}$$

where $(\mathcal{D}\mathfrak{D})(\eta, \eta') = \eta_1' \Delta$, we have

$$\begin{aligned} \Pi_{\circ}(a(u), \mathcal{L}u) &= u \circ \Pi_{\prec}((a(u), a'(u)), (\mathcal{D}\mathfrak{D})u) + \mathcal{C}_T^{3\alpha-2} \\ &= \Pi_{\leftarrow}(a(u), \vartheta) \circ \Pi_{\prec}((a(u), a'(u)), (\mathcal{D}\mathfrak{D})\vartheta) + \mathcal{C}_T^{3\alpha-2} \\ &= \Pi_{\diamond}((a(u), a'(u)), \vartheta \circ (\mathcal{D}\mathfrak{D})\vartheta) + \mathcal{C}_T^{3\alpha-2} \end{aligned}$$

Finally the equation for U^\sharp reads

$$\mathcal{L}U^\sharp = \Pi_\diamond((a(u), a'(u)), \vartheta \circ \mathcal{D}\Xi + \vartheta \circ (\mathcal{D}\mathcal{D})\vartheta) + \mathcal{E}_T^{3\alpha-2}.$$

This equation can be solved essentially as we did in the simpler context. We see that the general enhancement has the form

$$(\xi, \vartheta \circ \mathcal{D}\Xi + \vartheta \circ (\mathcal{D}\mathcal{D})\vartheta)$$

which of course will require renormalization like we did before. In particular

$$\begin{aligned} (\vartheta \circ \mathcal{D}\Xi + \vartheta \circ (\mathcal{D}\mathcal{D})\vartheta)(\eta, \eta') &= \vartheta(\eta) \circ \eta'_2 \partial_{\eta_2} \xi(\eta_2, \cdot) + \vartheta(\eta) \circ \eta'_1 \Delta \vartheta(\eta) \\ &= -\frac{\eta'_2}{\eta_1} (\Delta^{-1} \xi(\eta_2, \cdot)) \circ \partial_{\eta_2} \xi(\eta_2, \cdot) + \frac{\eta'_1}{\eta_1} (\Delta^{-1} \xi(\eta_2, \cdot)) \circ \xi(\eta_2, \cdot) \end{aligned}$$

where we used that $\eta_1 \Delta \vartheta(\eta) = -\xi(\eta_2, \cdot)$. Now observe that

$$\mathbb{E}[(\Delta^{-1} \xi_\varepsilon(\eta_2, \cdot)) \circ \xi_\varepsilon(\eta_2, \cdot)] = -F(\eta_2, \eta_2) \sigma_\varepsilon$$

and that

$$\mathbb{E}[(\Delta^{-1} \xi_\varepsilon(\eta_2, \cdot)) \circ \partial_{\eta_2} \xi_\varepsilon(\eta_2, \cdot)] = -(\partial_1 F)(\eta_2, \eta_2) \sigma_\varepsilon$$

with $\partial_1 F$ denoting the derivative with respect to the first entry.

In the end the renormalized enhanced noise is obtained as the limit in \mathcal{X}^α of $(\xi_\varepsilon, \Xi_{2,\varepsilon})$ where

$$\Xi_{2,\varepsilon}(\eta, \eta') = -\frac{\eta'_2}{\eta_1} (\Delta^{-1} \xi_\varepsilon(\eta_2, \cdot)) \circ \partial_{\eta_2} \xi_\varepsilon(\eta_2, \cdot) + \frac{\eta'_1}{\eta_1} (\Delta^{-1} \xi_\varepsilon(\eta_2, \cdot)) \circ \xi_\varepsilon(\eta_2, \cdot) - H_\varepsilon(\eta, \eta')$$

with

$$H_\varepsilon(\eta, \eta') = \frac{\eta'_2}{\eta_1} (\partial_1 F)(\eta_2, \eta_2) \sigma_\varepsilon - \frac{\eta'_1}{\eta_1^2} F(\eta_2, \eta_2) \sigma_\varepsilon.$$

We remark that if we take $F(\eta_2, \tilde{\eta}_2) = \eta_2 \tilde{\eta}_2$ we reobtain the situation treated in Section 6, indeed in this case

$$\Pi_\diamond((a(u), a'(u)), H_\varepsilon) = \frac{a'_2(u) a_2(u)}{a_1(u)} \sigma_\varepsilon - \frac{a'_1(u) a_2(u)^2}{a_1(u)^2} \sigma_\varepsilon.$$

which coincides with (39).

Remark 17. Consider the more general equation (2), where the noise depends explicitly on time, e.g. with a covariance

$$\mathbb{E}[\xi(\eta, t, x) \xi(\eta', t', x')] = F(\eta, \eta') Q(t - t', x - x')$$

with F a smooth function and Q a distribution of parabolic regularity $\rho > -4/3$. First note that the coefficient $a_1(u) \in [\lambda, 1]$ in front of the time derivative can be eliminated trivially by dividing.

In order to handle the time dependence of the noise, the framework of this paper will still apply, provided we consider space–time para-products instead of para-products which act only on the space variable. This can be done exactly following the lines of the paper [15] where time para-products were considered in the paracontrolled approach to solutions to SDE driven by gaussian signals.

The constraint of regularity $\rho > -4/3$ does allow to treat a noise which is white in time and smooth in space, but not a space–time white noise. It is well known that the first order paracontrolled approach on which the present paper is based does not allow to treat this kind of irregular signals in full generality.

Appendix A. Besov spaces

In this Appendix we collect some classical results from harmonic analysis needed in the paper. For a gentle introduction to Littlewood-Paley theory and Besov spaces see the recent monograph [3], where most of our results are taken from. There the case of tempered distributions on \mathbb{R}^d is considered. The Schauder estimates for the heat semigroup are classical and can be found in [15, 17].

Fix $d \in \mathbb{N}$ and denote by $\mathbb{T}^d = (\mathbb{R}/(2\pi\mathbb{Z}))^d$ the d -dimensional torus. We focus here on distributions and SPDEs on the torus, but everything in this Appendix applies mutatis mutandis on the full space \mathbb{R}^d , see [15]. The space of distributions $\mathcal{D}' = \mathcal{D}'(\mathbb{T}^d)$ is defined as the set of linear maps f from $C^\infty = C^\infty(\mathbb{T}^d, \mathbb{C})$ to \mathbb{C} , such that there exist $k \in \mathbb{N}$ and $C > 0$ with

$$|\langle f, \varphi \rangle| := |f(\varphi)| \leq C \sup_{|\mu| \leq k} \|\partial^\mu \varphi\|_{L^\infty(\mathbb{T}^d)}$$

for all $\varphi \in C^\infty$. In particular, the Fourier transform $\mathcal{F}f: \mathbb{Z}^d \rightarrow \mathbb{C}$, $\mathcal{F}f(k) = \langle f, e^{-ik\cdot} \rangle$, is defined for all $f \in \mathcal{D}'$, and it satisfies $|\mathcal{F}f(k)| \leq |P(k)|$ for a suitable polynomial P . We will also write $\hat{f}(k) = \mathcal{F}f(k)$. Conversely, if $(g(k))_{k \in \mathbb{Z}^d}$ is at most of polynomial growth, then its inverse Fourier transform

$$\mathcal{F}^{-1}g = (2\pi)^{-d} \sum_{k \in \mathbb{Z}^d} e^{i(k,\cdot)} g(k)$$

defines a distribution, and we have $\mathcal{F}^{-1}\mathcal{F}f = f$ as well as $\mathcal{F}\mathcal{F}^{-1}g = g$. To see this, it suffices to note that the Fourier transform of $\varphi \in C^\infty$ decays faster than any rational function (we say that it is of *rapid decay*). Indeed, for $\mu \in \mathbb{N}_0^d$ we have $|k^\mu \hat{g}(k)| = |\mathcal{F}(\partial^\mu g)(k)| \leq \|\partial^\mu g\|_{L^1(\mathbb{T}^d)}$ for all $k \in \mathbb{Z}^d$. As a consequence we get the Parseval formula $\langle f, \bar{\varphi} \rangle = (2\pi)^{-d} \sum_k \hat{f}(k) \overline{\hat{\varphi}(k)}$ for $f \in \mathcal{D}'$ and $\varphi \in C^\infty$.

Linear maps on \mathcal{D}' can be defined by duality: if $A: C^\infty \rightarrow C^\infty$ is such that for all $k \in \mathbb{N}$ there exists $n \in \mathbb{N}$ and $C > 0$ with $\sup_{|\mu| \leq k} \|\partial^\mu(A\varphi)\|_{L^\infty} \leq C \sup_{|\mu| \leq n} \|\partial^\mu \varphi\|$, then we set $\langle {}^t A f, \varphi \rangle = \langle f, A\varphi \rangle$. Differential operators are defined by $\langle \partial^\mu f, \varphi \rangle = (-1)^{|\mu|} \langle f, \partial^\mu \varphi \rangle$. If $\varphi: \mathbb{Z}^d \rightarrow \mathbb{C}$ grows at most polynomially, then it defines a *Fourier multiplier*

$$\varphi(D)f = \mathcal{F}^{-1}(\varphi \mathcal{F}f),$$

which gives us a distribution $\varphi(D)f \in \mathcal{D}'$ for every $f \in \mathcal{D}'$.

Littlewood-Paley blocks give a decomposition of any distribution on \mathcal{D}' into an infinite series of smooth functions.

Definition 18. A *dyadic partition of unity* consists of two nonnegative radial functions $\chi, \rho \in C^\infty(\mathbb{R}^d, \mathbb{R})$, where ρ is supported in a ball $\mathcal{B} = \{|x| \leq c\}$ and ρ is supported in an annulus $\mathcal{A} = \{a \leq |x| \leq b\}$ for suitable $a, b, c > 0$, such that

1. $\chi + \sum_{j \geq 0} \rho(2^{-j}\cdot) \equiv 1$ and
2. $\chi \rho(2^{-j}\cdot) \equiv 0$ for $j \geq 1$ and $\rho(2^{-i}\cdot)\rho(2^{-j}\cdot) \equiv 0$ for all $i, j \geq 0$ with $|i - j| > 1$.

We will often write $\rho_{-1} = \chi$ and $\rho_j = \rho(2^{-j}\cdot)$ for $j \geq 0$.

Dyadic partitions of unity exist, see [3]. The reason for considering smooth partitions rather than indicator functions is that indicator functions do not have good Fourier properties. We fix a dyadic partition of unity (χ, ρ) and define the dyadic blocks

$$\Delta_j f = \rho_j(D)f = \mathcal{F}^{-1}(\rho_j \hat{f}), \quad j \geq -1.$$

We also use the notation

$$S_j f = \sum_{i < j-1} \Delta_i f$$

and notice that

$$\Delta_j f(x) = \int K_{j,x}(y) f(y) dy, \quad S_j f(x) = \int P_{j,x}(y) f(y) dy$$

with $K_{j,x}(y) = 2^{dj} K(2^j(x-y))$, $P_{j,x}(y) = \sum_{i < j-1} K_{i,x}(y)$, K radial with zero mean.

Every dyadic block has a compactly supported Fourier transform and it belongs therefore to C^∞ . It is easy to see that $f = \sum_{j \geq -1} \Delta_j f = \lim_{j \rightarrow \infty} S_j f$ for all $f \in \mathcal{D}'$.

For $\alpha \in \mathbb{R}$, the Hölder-Besov space \mathcal{E}^α is given by $\mathcal{E}^\alpha = B_{\infty, \infty}^\alpha(\mathbb{T}^d, \mathbb{R})$, where for $p, q \in [1, \infty]$ we define

$$B_{p,q}^\alpha = B_{p,q}^\alpha(\mathbb{T}^d, \mathbb{R}) = \left\{ f \in \mathcal{D}' : \|f\|_{B_{p,q}^\alpha} = \left(\sum_{j \geq -1} (2^{j\alpha} \|\Delta_j f\|_{L^p})^q \right)^{1/q} < \infty \right\},$$

with the usual interpretation as ℓ^∞ norm in case $q = \infty$. Then $B_{p,q}^\alpha$ is a Banach space and while the norm $\|\cdot\|_{B_{p,q}^\alpha}$ depends on (χ, ρ) , the space $B_{p,q}^\alpha$ does not, and any other dyadic partition of unity corresponds to an equivalent norm.

If $\alpha \in (0, \infty) \setminus \mathbb{N}$, then \mathcal{E}^α is the space of $[\alpha]$ times differentiable functions whose partial derivatives of order $[\alpha]$ are $(\alpha - [\alpha])$ -Hölder continuous (see page 99 of [3]). Note however, that for $k \in \mathbb{N}$ the space \mathcal{E}^k is strictly larger than C^k , the space of k times continuously differentiable functions.

The following lemma gives useful characterisation of Besov regularity for functions that can be decomposed into pieces which are localized in Fourier space.

Lemma 19.

1. Let \mathcal{A} be an annulus, let $\alpha \in \mathbb{R}$, and let (u_j) be a sequence of smooth functions such that $\mathcal{F}u_j$ has its support in $2^j \mathcal{A}$, and such that $\|u_j\|_{L^\infty} \lesssim 2^{-j\alpha}$ for all j . Then

$$u = \sum_{j \geq -1} u_j \in \mathcal{E}^\alpha \quad \text{and} \quad \|u\|_\alpha \lesssim \sup_{j \geq -1} \{2^{j\alpha} \|u_j\|_{L^\infty}\}.$$

2. Let \mathcal{B} be a ball, let $\alpha > 0$, and let (u_j) be a sequence of smooth functions such that $\mathcal{F}u_j$ has its support in $2^j \mathcal{B}$, and such that $\|u_j\|_{L^\infty} \lesssim 2^{-j\alpha}$ for all j . Then

$$u = \sum_{j \geq -1} u_j \in \mathcal{E}^\alpha \quad \text{and} \quad \|u\|_\alpha \lesssim \sup_{j \geq -1} \{2^{j\alpha} \|u_j\|_{L^\infty}\}.$$

The Bernstein inequalities of the next lemma are extremely useful when dealing with functions with compactly supported Fourier transform.

Lemma 20. Let \mathcal{A} be an annulus and let \mathcal{B} be a ball. For any $k \in \mathbb{N}_0$, $\lambda > 0$, and $1 \leq p \leq q \leq \infty$ we have that

1. if $u \in L^p$ is such that $\text{supp}(\mathcal{F}u) \subseteq \lambda \mathcal{B}$, then

$$\max_{\mu \in \mathbb{N}^d: |\mu|=k} \|\partial^\mu u\|_{L^q} \lesssim_k \lambda^{k+d(\frac{1}{p}-\frac{1}{q})} \|u\|_{L^p};$$

2. if $u \in L^p$ is such that $\text{supp}(\mathcal{F}u) \subseteq \lambda \mathcal{A}$, then

$$\lambda^k \|u\|_{L^p} \lesssim_k \max_{\mu \in \mathbb{N}^d: |\mu|=k} \|\partial^\mu u\|_{L^p}.$$

We recall the following standard heat kernel estimations (see [3], Chapter 2).

Lemma 21. (Schauder estimates) *Let $V_t = \int_0^t e^{\eta(t-s)\Delta} v_s ds$ and $\mathcal{P}_t u_0 = e^{\eta\Delta t} u_0$, with $\eta \geq \lambda$. We define \mathcal{L}_T^α and $C_\eta^k \mathcal{L}_T^\alpha$, $C_\eta^k \mathcal{C}_T^\alpha$ for $k \in \mathbb{N}$ as in (8), (9) and introduce the norm*

$$\|v\|_{\mathcal{E}_T^\delta \mathcal{C}^\alpha} = \sup_{t \in [0, T]} t^\delta \|v(t, \cdot)\|_{\mathcal{C}^\alpha}.$$

Then for any $\gamma \in [0, 1)$ and $\alpha \in \mathbb{R}$:

$$\begin{aligned} \sup_{t \in [0, T]} t^\gamma \|V_t\|_{C_\eta^k \mathcal{C}^{\alpha-2\beta}} &\lesssim \sup_{t \in [0, T]} t^{\gamma+\beta} \|v_t\|_{\mathcal{C}^{\alpha-2}} \quad \forall \beta \in [0, 1) \\ \|V\|_{C_\eta^k \mathcal{L}_T^{\alpha-2\beta}} &\lesssim T^\beta \|v\|_{\mathcal{C}_T^{\alpha-2}} \quad \forall \beta \in [0, 1) \\ \|V\|_{C_\eta^k \mathcal{L}_T^{\alpha-2\beta}} &\lesssim T^{\beta-\delta} \|v\|_{\mathcal{E}_T^\delta \mathcal{C}^{\alpha-2}} \quad \forall \beta \in [0, 1), \forall \delta \in [0, \beta] \\ \|V\|_{C_\eta^k \mathcal{L}_T^\gamma} &\lesssim T^{\frac{\rho-\gamma}{2}+1-\delta} \|v\|_{\mathcal{E}_T^\delta \mathcal{C}^\rho} \quad \forall \rho \in [\gamma-2, \gamma), \forall \delta \in \left[0, \frac{\rho-\gamma}{2}+1\right] \\ \|\mathcal{P}_t u_0\|_{C_\eta^k \mathcal{C}^\alpha} &\lesssim t^{-\frac{\alpha-\beta}{2}} \|u_0\|_{\mathcal{C}^\beta} \quad \forall \beta < \alpha \\ \|\mathcal{P}_t u_0 - \mathcal{P}_s u_0\|_{C_\eta^k \mathcal{C}^\alpha} &\lesssim s^{-\delta} |t-s|^{\frac{\beta-\alpha}{2}+\delta} \|u_0\|_{\mathcal{C}^\beta} \\ &\quad \forall t \neq s \in \mathbb{R}^+, \beta \leq \alpha+2, \delta \in \left[0, 1 - \frac{\beta-\alpha}{2}\right] \end{aligned}$$

We need the following interpolation lemma:

Lemma 22. *Let $\gamma \in (0, 2)$, $0 < \varepsilon < \gamma$ and $u \in \mathcal{L}_T^\gamma$. Then*

$$\|u\|_{C_T^{\gamma/2-\varepsilon/2} L^\infty} \lesssim \|u\|_{\mathcal{L}_T^\gamma}$$

Proof.

$$\sup_{s \neq t} \frac{\|u_t - u_s\|_{L^\infty}}{|t-s|^{\gamma/2-\varepsilon/2}} \leq \sup_{s \neq t} \left[\sum_{i \leq n} \frac{\|\Delta_i u_t - \Delta_i u_s\|_{L^\infty}}{|t-s|^{\gamma/2-\varepsilon/2}} + \sum_{i > n} \frac{\|\Delta_i u_t - \Delta_i u_s\|_{L^\infty}}{|t-s|^{\gamma/2-\varepsilon/2}} \right]$$

and choosing $2^{-n-1} \leq |t-s|^{1/2} \leq 2^{-n}$ we obtain

$$\begin{aligned} \sum_{i < n} \frac{\|\Delta_i u_t - \Delta_i u_s\|_{L^\infty}}{|t-s|^{\gamma/2-\varepsilon/2}} &\lesssim \|u\|_{C_T^{\gamma/2} \mathcal{C}^0} \sum_{i \leq n} |t-s|^{\varepsilon/2} \\ \sum_{i \geq n} \frac{\|\Delta_i u_t - \Delta_i u_s\|_{L^\infty}}{|t-s|^{\gamma/2-\varepsilon/2}} &\lesssim \|u\|_{\mathcal{L}_T^\gamma} \sum_{i > n} 2^{-\gamma i} 2^{-(\gamma+\varepsilon)n} \end{aligned}$$

and this gives the result. \square

Terms of the type $\|a(u(t, x))\|_{\mathcal{C}^0}$ with $a: \mathbb{R} \rightarrow \mathbb{R}$ cannot be estimated directly with their Hölder norm. In the following lemma we note some bounds used in Section 4.

Lemma 23. *Let $a \in C_b^3$ uniformly bounded and $u \in \mathcal{L}_T^\alpha = \mathcal{C}_T^\alpha \cap C_T^{\alpha/2} \mathcal{C}^0$, then*

$$\begin{aligned} \|a(u)\|_{\mathcal{L}_T^\alpha} &\lesssim \|a\|_{L^\infty} + \|a'\|_{L^\infty} \left(\|u\|_{C_T^{\alpha/2} \mathcal{C}^0} + \|u\|_{\mathcal{C}_T^\alpha} \right) + \|a''\|_{L^\infty} \|u\|_{C_T^{\alpha/2} \mathcal{C}^0} \|u\|_{\mathcal{C}_T^\alpha} \\ &\lesssim 1 + \|u\|_{\mathcal{L}_T^\alpha} + \|u\|_{\mathcal{L}_T^\alpha}^2, \end{aligned}$$

$$\begin{aligned} \|a(u_1) - a(u_2)\|_{\mathcal{L}_T^\alpha} &\lesssim \|u_1 - u_2\|_{\mathcal{L}_T^\alpha} (1 + \|u_1\|_{\mathcal{C}_T^\alpha} + \|u_2\|_{\mathcal{C}_T^\alpha})^2, \\ \|a(u_1) - a(u_2)\|_{\mathcal{C}_T^0} &\lesssim \|u_1 - u_2\|_{\mathcal{C}_T^0} (1 + \|u_1\|_{\mathcal{C}_T^\alpha} + \|u_2\|_{\mathcal{C}_T^\alpha}). \end{aligned}$$

Proof. The bound on $\|a(u)\|_{\mathcal{E}_T^\alpha}$ is trivial. We estimate $\|a(u_t) - a(u_s)\|_{\mathcal{E}^0}$ as

$$\begin{aligned} & \left| \int_z K_{i,x}(z) [a(u(t,z)) - a(u(s,z))] \right| = \left| \int_{z,\tau \in [0,1]} K_{i,x}(z) a'(\delta_\tau u_{sz}^{tz}) [u(t,z) - u(s,z)] \right| \\ & \leq \left| \int_{z,w,\tau \in [0,1]} K_{i,x}(z) \sum_{j \lesssim k \sim i} [\Delta_j u(t,z) - \Delta_j u(s,z)] K_{k,z}(w) a'(\delta_\tau u_{sw}^{tw}) \right| \\ & \quad + \left| \int_{z,w,\tau \in [0,1]} K_{i,x}(z) \sum_{k \lesssim j \sim i} [\Delta_j u(t,z) - \Delta_j u(s,z)] K_{k,z}(w) a'(\delta_\tau u_{sw}^{tw}) \right| \\ & \quad + \left| \int_{z,w,\tau \in [0,1]} K_{i,x}(z) \sum_{k \sim j \gtrsim i} [\Delta_j u(t,z) - \Delta_j u(s,z)] K_{k,z}(w) a'(\delta_\tau u_{sw}^{tw}) \right|. \end{aligned}$$

If $k > -1$ we have

$$\int_w K_{k,z}(w) a'(\delta_\tau u_{sw}^{tw}) = \int_w K_{k,z}(w) [a'(\delta_\tau u_{sw}^{tw}) - a'(\delta_\tau u_{sz}^{tz})]$$

and then the first term above becomes

$$\begin{aligned} & \left| \int_{z,w,\tau,\sigma} K_{i,x}(z) \sum_{j \lesssim k \sim i} [\Delta_j u(t,z) - \Delta_j u(s,z)] K_{k,z}(w) a''(\delta_\sigma(\delta_\tau u_s^t)_z^w) [\delta u_{tz}^{tw} - \delta u_{sz}^{sw}] \right| \\ & \lesssim \|a''\|_{L^\infty} \|u_t - u_s\|_{\mathcal{E}^0} \sum_{j \lesssim k \sim i} \int_w |K_{k,z}(w)| |w - z|^\alpha \|u\|_{\mathcal{E}_T^\alpha} \lesssim i 2^{-\alpha i} \|a''\|_{L^\infty} \|u_t - u_s\|_{\mathcal{E}^0} \|u\|_{\mathcal{E}_T^\alpha}. \end{aligned}$$

The second term is

$$\begin{aligned} & \left| \int_{z,w,\tau \in [0,1]} K_{i,x}(z) \sum_{j \sim i} [\Delta_j u(t, \cdot) - \Delta_j u(s, \cdot)](z) P_{j,z}(w) a'(\delta_\tau u_{sw}^{tw}) \right| \\ & \lesssim \|u_t - u_s\|_{\mathcal{E}^0} \|P_{j,z}(w) a'(\delta_\tau u_{sw}^{tw})\|_{L^\infty} \lesssim \|u_t - u_s\|_{\mathcal{E}^0} \|a'\|_{L^\infty} \end{aligned}$$

The third term can be estimated as the first one when $k > -1$. Otherwise we just bound it as

$$\left| \int_{z,w,\tau \in [0,1]} K_{i,x}(z) \sum_{j \leq N} [\Delta_j u(t, \cdot) - \Delta_j u(s, \cdot)](z) K_{-1,z}(w) a'(\delta_\tau u_{sw}^{tw}) \right| \lesssim \|u_t - u_s\|_{\mathcal{E}^0} \|a'\|_{L^\infty}.$$

For the three terms together we have the bound

$$\|a(u_t) - a(u_s)\|_{\mathcal{E}^0} \lesssim \|u_t - u_s\|_{\mathcal{E}^0} (\|a'\|_{L^\infty} + \|a''\|_{L^\infty}) \|u\|_{\mathcal{E}_T^\alpha}$$

With the same technique we obtain

$$\begin{aligned} \|a(u_1) - a(u_2)\|_{C_T^{\alpha/2} \mathcal{E}^0} & \lesssim \|a'\|_{L^\infty} \|u_1 - u_2\|_{C_T^{\alpha/2} \mathcal{E}^0} \\ & \quad + \|a''\|_{L^\infty} \|u_1 - u_2\|_{C_T^{\alpha/2} \mathcal{E}^0} \|u_1 - u_2\|_{\mathcal{E}_T^\alpha} \\ & \quad + \|a'''\|_{L^\infty} \|u_1 - u_2\|_{C_T^{\alpha/2} \mathcal{E}^0} \|u_1 - u_2\|_{\mathcal{E}_T^\alpha}^2 \end{aligned}$$

and this gives the second estimate. The third one can be obtained easily. \square

A.1. Bony's paraproduct.

In terms of Littlewood–Paley blocks, the product fg can be decomposed as

$$fg = \sum_{j \geq -1} \sum_{i \geq -1} \Delta_i f \Delta_j g = f \prec g + f \succ g + f \circ g.$$

Where the *paraproducts* $f \prec g$ and $f \succ g$ and the *resonant product* $f \circ g$ are defined as

$$f \prec g = g \succ f := \sum_{j \geq -1} \sum_{i=-1}^{j-2} \Delta_i f \Delta_j g \quad \text{and} \quad f \circ g := \sum_{|i-j| \leq 1} \Delta_i f \Delta_j g.$$

We will often use the shortcuts $\sum_{i \sim j}$ for $\sum_{|i-j| \leq 1}$ and $\sum_{i \lesssim j}$ for $\sum_{i < j-1}$. Of course, the decomposition depends on the dyadic partition of unity used to define the blocks Δ_j , and also on the particular choice of the pairs (i, j) in the diagonal part. The choice of taking all (i, j) with $|i-j| \leq 1$ into the diagonal part corresponds to property 2 in our definition of dyadic partitions of unity.

Bony’s crucial observation [9, 31] is that the paraproduct $f \prec g$ (and thus $f \succ g$) is always a well-defined distribution. Heuristically, $f \prec g$ behaves at large frequencies like g (and thus retains the same regularity), and f provides only a frequency modulation of g . The only difficulty in constructing fg for arbitrary distributions lies in handling the resonant product $f \circ g$. The basic result about these bilinear operations is given by the following estimates.

Theorem 24. (*Paraproduct estimates*) For any $\beta \in \mathbb{R}$ and $f, g \in \mathcal{D}'$ we have

$$\|f \prec g\|_{\mathcal{C}^\beta} \lesssim_\beta \|f\|_{L^\infty} \|g\|_{\mathcal{C}^\beta}, \quad (41)$$

and for $\alpha < 0$ furthermore

$$\|f \prec g\|_{\mathcal{C}^{\alpha+\beta}} \lesssim_{\alpha, \beta} \|f\|_{\mathcal{C}^\alpha} \|g\|_{\mathcal{C}^\beta}. \quad (42)$$

For $\alpha + \beta > 0$ we have

$$\|f \circ g\|_{\mathcal{C}^{\alpha+\beta}} \lesssim_{\alpha, \beta} \|f\|_{\mathcal{C}^\alpha} \|g\|_{\mathcal{C}^\beta}. \quad (43)$$

Bony proved also a basic *paralinearisation* result, soon after improved by Meyer. We give here a particular version suited to our purposes.

Theorem 25. Let $\alpha \in (0, 1)$, $f \in (\mathcal{C}^\alpha)^k$ and $F \in C^3(\mathbb{R}^k; \mathbb{R})$ then

$$R_F(f) := F(f) - F'(f) \prec f \in \mathcal{C}^{2\alpha}$$

with

$$\|R_F(f)\|_{\mathcal{C}^{2\alpha}} \lesssim \|F\|_{C^2} (1 + \|f\|_{\mathcal{C}^\alpha})^2.$$

Moreover the map $f \mapsto R_F(f)$ is locally Lipschitz and

$$\|R_F(f) - R_F(\tilde{f})\|_{\mathcal{C}^{2\alpha}} \lesssim \|F\|_{C^3} (1 + \|f\|_{\mathcal{C}^\alpha} + \|\tilde{f}\|_{\mathcal{C}^\alpha})^2 \|\tilde{f} - f\|_{\mathcal{C}^\alpha}.$$

The additional key ingredient at the core of the paracontrolled approach is the following commutation result proved in [15], Lemma 2.4:

Lemma 26. Assume that $\alpha, \beta, \gamma \in \mathbb{R}$ are such that $\alpha + \beta + \gamma > 0$ and $\beta + \gamma < 0$. Then for $f, g, h \in C^\infty$ the trilinear operator

$$C(f, g, h) := ((f \prec g) \circ h) - f(g \circ h)$$

allows for the bound

$$\|C(f, g, h)\|_{\mathcal{E}^{\beta+\gamma}} \lesssim \|f\|_{\mathcal{E}^\alpha} \|g\|_{\mathcal{E}^\beta} \|h\|_{\mathcal{E}^\gamma}, \quad (44)$$

and can thus be uniquely extended to a bounded trilinear operator

$$C: \mathcal{E}^\alpha \times \mathcal{E}^\beta \times \mathcal{E}^\alpha \rightarrow \mathcal{E}^{\beta+\gamma}.$$

We will need the following two lemmas to compare standard and time-smoothed para-products. The first one has essentially the same proof as [15], Lemma 5.1.

Lemma 27. *Let $\rho \in (0, 2)$, $\gamma \in \mathbb{R}$. Then for every $\varepsilon > 0$ we have the bound*

$$\|g \prec h - g \llcorner h\|_{\mathcal{E}_T^{\rho+\gamma-\varepsilon}} \lesssim \|g\|_{C_T^{\rho/2} \mathcal{E}_0} \|h\|_{\mathcal{E}_T^\gamma}.$$

The second lemma has a standard proof.

Lemma 28. *Let $g \in \mathcal{L}_T^\rho$, $h \in C_\eta^1 \mathcal{E}_T^\gamma$ with $\rho \in (0, 1)$, $\gamma \in \mathbb{R}$. We have, $\forall \varepsilon > 0$*

$$\|\Pi_\prec(g, h) - \Pi_\llcorner(g, h)\|_{\mathcal{E}_T^{\rho+\gamma-\varepsilon}} \lesssim \|g\|_{\mathcal{L}_T^\rho} \|h\|_{C_\eta^1 \mathcal{E}_T^\gamma}.$$

BIBLIOGRAPHY

- [1] S. Alinhac and P. Gérard. *Opérateurs pseudo-différentiels et théorème de Nash-Moser*. EDP Sciences, 1991.
- [2] R. Allez and K. Chouk. The continuous Anderson hamiltonian in dimension two. *ArXiv:1511.02718 [math]*, nov 2015. ArXiv: 1511.02718.
- [3] H. Bahouri, J.-Y. Chemin, and R. Danchin. *Fourier Analysis and Nonlinear Partial Differential Equations*. Springer, jan 2011.
- [4] I. Bailleul and F. Bernicot. Heat semigroup and singular PDEs. *ArXiv:1501.06822 [math]*, jan 2015. ArXiv: 1501.06822.
- [5] I. Bailleul and F. Bernicot. Higher order paracontrolled calculus. *ArXiv:1609.06966 [math]*, sep 2016. ArXiv: 1609.06966.
- [6] I. Bailleul and R. Catellier. Rough flows and homogenization in stochastic turbulence. *ArXiv:1601.01902 [math-ph]*, jan 2016. ArXiv: 1601.01902.
- [7] Ismael Bailleul, Arnaud Debussche, and Martina Hofmanova. Quasilinear generalized parabolic anderson model equation. *ArXiv:1610.06726 [math]*, Oct 2016. ArXiv: 1610.06726.
- [8] Ismaël Bailleul. Flows driven by rough paths. *Revista Matemática Iberoamericana*, 31(3):901–934, 2015.
- [9] J.-M. Bony. Calcul symbolique et propagation des singularités pour les équations aux dérivées partielles non linéaires. *Ann. Sci. Éc. Norm. Supér. (4)*, 14:209–246, 1981.
- [10] J.-M. Bony, G. Grubb, L. Hörmander, H. Komatsu, J. Sjöstrand, L. Cattabriga, and L. Rodino. *Microlocal Analysis and Applications Lectures given at the 2nd Session of the Centro Internazionale Matematico Estivo (C.I.M.E.) held at Montecatini Terme, Italy, July 3-11, 1989*. Springer Berlin Heidelberg, Berlin, Heidelberg, 1991.
- [11] R. Catellier and K. Chouk. Paracontrolled Distributions and the 3-dimensional Stochastic Quantization Equation. *ArXiv:1310.6869 [math-ph]*, oct 2013. ArXiv: 1310.6869.
- [12] R. Catellier and M. Gubinelli. Averaging along irregular curves and regularisation of ODEs. *Stochastic Processes and their Applications*, 126(8):2323–2366, aug 2016.
- [13] P. K. Friz and M. Hairer. *A Course on Rough Paths: With an Introduction to Regularity Structures*. Springer, aug 2014.
- [14] M. Gubinelli. Rough solutions for the periodic Korteweg–de Vries equation. *Communications on Pure and Applied Analysis*, 11(2):709–733, 2012.
- [15] M. Gubinelli, P. Imkeller, and N. Perkowski. Paracontrolled distributions and singular PDEs. *Forum of Mathematics. Pi*, 3:0, 2015.
- [16] M. Gubinelli and N. Perkowski. KPZ reloaded. *ArXiv:1508.03877 [math-ph]*, aug 2015. ArXiv: 1508.03877.
- [17] M. Gubinelli and N. Perkowski. Lectures on singular stochastic PDEs. *ArXiv:1502.00157 [math]*, jan 2015. ArXiv: 1502.00157.

- [18] M. Hairer. Solving the KPZ equation. *Annals of Mathematics*, 178(2):559–664, jun 2013.
- [19] M. Hairer. A theory of regularity structures. *Inventiones mathematicae*, 198(2):269–504, mar 2014.
- [20] M. Hairer. Singular stochastic PDEs. *ArXiv:1403.6353 [math]*, mar 2014. ArXiv: 1403.6353.
- [21] M. Hairer and J. Quastel. A class of growth models rescaling to KPZ. *ArXiv:1512.07845 [math-ph]*, dec 2015. ArXiv: 1512.07845.
- [22] M. Hairer and H. Shen. The dynamical sine-Gordon model. *ArXiv:1409.5724 [math-ph]*, sep 2014. ArXiv: 1409.5724.
- [23] M. Hairer and W. Xu. Large scale behaviour of 3d phase coexistence models. *ArXiv:1601.05138 [math-ph]*, jan 2016. ArXiv: 1601.05138.
- [24] Y. Hu and K. Le. Nonlinear Young integrals and differential systems in Hölder media. *ArXiv:1404.7582 [math]*, apr 2014.
- [25] L. Hörmander. The Nash-Moser theorem and paradifferential operators. In *Analysis, et cetera*, pages 429–449. Academic Press, Boston, MA, 1990.
- [26] D. Kelly and I. Melbourne. Deterministic homogenization for fast-slow systems with chaotic noise. *ArXiv:1409.5748 [math]*, sep 2014. ArXiv: 1409.5748.
- [27] H. Kunita. Stochastic differential equations and stochastic flows of diffeomorphisms. In P. L. Hennequin, editor, *École d'Été de Probabilités de Saint-Flour XII - 1982*, number 1097 in Lecture Notes in Mathematics, pages 143–303. Springer Berlin Heidelberg, 1984. DOI: 10.1007/BFb0099433.
- [28] T. Lyons. Differential equations driven by rough signals. *Revista Matemática Iberoamericana*, pages 215–310, 1998.
- [29] T. J. Lyons, M. J. Caruana, and T. Lévy. *Differential Equations Driven by Rough Paths: Ecole d'Été de Probabilités de Saint-Flour XXXIV-2004*. Springer, 1 edition, jun 2007.
- [30] T. Lyons and Z. Qian. *System Control and Rough Paths*. Oxford University Press, 2002.
- [31] Y. Meyer. Remarques sur un théorème de J.-M. Bony. In *Rendiconti del Circolo Matematico di Palermo. Serie II*, pages 1–20. 1981.
- [32] J.-C. Mourrat and H. Weber. Global well-posedness of the dynamic Φ^4_3 model on the torus. *ArXiv:1601.01234 [math]*, jan 2016. ArXiv: 1601.01234.
- [33] F. Otto and H. Weber. Quasilinear SPDEs via rough paths. *ArXiv:1605.09744 [math]*, may 2016. ArXiv: 1605.09744.