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# Directional dimensions of ergodic currents on $\mathbb{C P}(2)$ 

Christophe Dupont and Axel Rogue

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#### Abstract

Let $f$ be a holomorphic endomorphism of $\mathbb{P}^{2}$ of degree $d \geq 2$. We estimate the local directional dimensions of closed positive currents $S$ with respect to ergodic dilating measures $\nu$. We infer several applications. The first one shows that the currents $S$ containing a measure of entropy $h_{\nu}>\log d$ have a directional dimension $>2$, which answers a question by de Thélin-Vigny. The second application asserts that the Dujardin's semi-extremal endomorphisms are close to suspensions of onedimensional Lattès maps. Finally, we obtain an upper bound for the dimension of the equilibrium measure, towards the formula conjectured by Binder-DeMarco.


Key words: dimension theory, positive closed currents, invariant measures, Lyapunov exponents, normal forms.

MSC 2010: 32H50, 32U40, 37C45, 37F10.

## 1 Introduction

This article concerns the ergodic properties of holomorphic endomorphisms of $\mathbb{P}^{2}$, see [16]. Let $f$ be an endomorphism of $\mathbb{P}^{2}$ of degree $d \geq 2$. The Green current $T$ is defined as $T:=\lim _{n} \frac{1}{d^{n}} f^{n *} \omega$, where $\omega$ is the Fubini-Study form of $\mathbb{P}^{2}$. The equilibrium measure is defined as $\mu:=T \wedge T$, this is an ergodic measure of entropy $\log d^{2}$, its Lyapunov exponents are $\geq \frac{1}{2} \log d$, see $[6,16]$.

We say that an ergodic measure $\nu$ is dilating if its Lyapunov exponents are positive. The ergodic measures of entropy $h_{\nu}>\log d$ are dilating: their exponents are larger than or equal to $\frac{1}{2}\left(h_{\nu}-\log d\right)$, see $[11,20]$. It is known that the support of every ergodic measure $\nu$ of entropy $h_{\nu}>\log d$ is contained in the support of $\mu$, see [10, 14]. The article [20] constructs such measures by using coding techniques.

### 1.1 Directional dimensions

Let $f$ be an endomorphism of $\mathbb{P}^{2}$ of degree $d \geq 2$. The algebraic subsets of $\mathbb{P}^{2}$ do not contain any ergodic measure $\nu$ of entropy $h_{\nu}>\log d$ : this comes from the Gromov's iterated graph argument and from the relative variational principle, see [7] and [16, Section 1.7]. In this Section, we quantify that property thanks to the Lyapunov exponents
of the measures $\nu$. We shall work in the more general setting of $(1,1)$ closed positive currents $S$. Those currents are as big as the algebraic subsets, since we have:

$$
\forall x \in \mathbb{P}^{2}, \underline{d_{S}}(x):=\liminf _{r \rightarrow 0} \frac{\log S \wedge \omega\left(B_{x}(r)\right)}{\log r} \geq 2
$$

which comes from $S \wedge \omega\left(B_{x}(r)\right) \leq c(x) r^{2}$, see [13, Chapitre 3]. We shall use the notation $\overline{d_{S}}(x)$ for the lim sup. A drawback of the trace measure $S \wedge \omega$ is that it does not distinguish any specific direction. If $Z$ is a holomorphic coordinate in the neighbourhood of $x \in \mathbb{P}^{2}$, one defines the lower local directional dimension of $S$ with respect to $Z$ by

$$
\underline{d_{S, Z}}(x):=\liminf _{r \rightarrow 0} \frac{\log \left[S \wedge\left(\frac{i}{2} d Z \wedge d \bar{Z}\right)\left(B_{x}(r)\right)\right]}{\log r}
$$

we shall denote the limsup by $\overline{d_{S, Z}}(x)$. Geometrically, the positive measure $S \wedge\left(\frac{i}{2} d Z \wedge d \bar{Z}\right)$ is the average with respect to Lebesgue measure of the slices of the current $S$ transversaly to the direction $Z$, see Proposition 9.4. If $(Z, W)$ are holomorphic coordinates near $x$, the directional dimensions of $S$ are related to the dimension of $S$ by

$$
\begin{equation*}
\underline{d_{S}}(x)=\min \left\{\underline{d_{S, Z}}(x), \underline{d_{S, W}}(x)\right\}, \overline{d_{S}}(x)=\min \left\{\overline{d_{S, Z}}(x), \overline{d_{S, W}}(x)\right\} \tag{1.1}
\end{equation*}
$$

see Proposition 9.2. We start with showing lower and upper bounds for the directional dimensions of the Green current $T$ (Theorems 1.1 and 1.2). Next we display our result concerning the general closed positive currents $S$.

## Directional dimensions of the Green current

The invariance property of the current $T$ allows to obtain estimates on the directional dimensions $\nu$-almost everywhere. In what follows, the functions $O(\epsilon)$ only depend on $\epsilon$, the degree $d$ of the endomorphism, the exponents and the entropy of $\nu$. They tend to 0 when $\epsilon$ tends to 0 and are positive for $\epsilon$ small enough. We shall say that the exponents of $\nu$ do not resonate if $\lambda_{1} \neq k \lambda_{2}$ for every $k \geq 2$.

Theorem 1.1. Let $f$ be an endomorphism of $\mathbb{P}^{2}$ of degree $d \geq 2$. Let $\nu$ be an ergodic diltating measure whose exponents $\lambda_{1}>\lambda_{2}$ do not resonate. Then there exist functions $O_{1}(\epsilon), O_{2}(\epsilon)$ satisfying the following properties. For every $\epsilon>0$ and for $\nu$-almost every $x$, there exist holomorphic coordinates $(Z, W)$ in a neighbourhood of $x$ such that

$$
\begin{gathered}
\overline{d_{T, Z}}(x) \geq 2+\frac{h_{\nu}-\log d}{\lambda_{2}}-O_{1}(\epsilon) \\
\overline{d_{T, W}}(x) \geq 2 \frac{\lambda_{2}}{\lambda_{1}}+\frac{h_{\nu}-\log d}{\lambda_{2}}-O_{2}(\epsilon)
\end{gathered}
$$

This result applies to the measure $\mu$ and allows to improve a classical lower bound concerning the dimension of the Green current $T$. This current has local $\gamma$-Hölder psh potentials for every $\gamma<\gamma_{0}:=\min \left\{1, \log d / \log d_{\infty}\right\}$, where $d_{\infty}:=\lim _{n}\left\|D f^{n}\right\|_{\infty}^{1 / n}$, see
[16, Proposition 1.18]. This implies that $T \wedge \omega\left(B_{x}(r)\right) \leq c_{\gamma}(x) r^{2+\gamma}$ for every $x \in \mathbb{P}^{2}$ and every $\gamma<\gamma_{0}$, see [24, Théorème 1.7.3]. We deduce that for every $x \in \mathbb{P}^{2}$ and for every local holomorphic coordinates $(Z, W)$ in a neighbourhood of $x$ :

$$
\begin{equation*}
\min \left\{\overline{d_{T, Z}}(x), \overline{d_{T, W}}(x)\right\}=\overline{d_{T}}(x) \geq 2+\gamma_{0} \tag{1.2}
\end{equation*}
$$

Now, since $h_{\mu}=\log d^{2}$ and $\lambda_{1} \leq \log d_{\infty}$, we get:

$$
\frac{h_{\mu}-\log d}{\lambda_{2}}=\frac{\log d}{\lambda_{2}}>\frac{\log d}{\lambda_{1}} \geq \frac{\log d}{\log d_{\infty}} \geq \gamma_{0}
$$

The lower bound in the direction $Z$ provided by Theorem 1.1 is therefore better than (1.2) when $\epsilon$ is small enough.

Now we show directional upper bounds for the current $T$ with respect to every dilating measure $\nu$ whose support is contained in the support of $\mu$.

Theorem 1.2. Let $f$ be an endomorphism of $\mathbb{P}^{2}$ of degree $d \geq 2$. Let $\nu$ be an ergodic dilating measure whose exponents $\lambda_{1}>\lambda_{2}$ do not resonate. We assume that the support of $\nu$ is contained in the support of $\mu$. Then there exist functions $O_{3}(\epsilon), O_{4}(\epsilon)$ satisfying the following properties. For every $\epsilon>0$ and for $\nu$-almost every $x$, there exist holomorphic coordinates $(Z, W)$ in a neighbourhood of $x$ such that

$$
\begin{gathered}
\underline{d_{T, Z}}(x) \leq \frac{\log d}{\lambda_{2}}+2 \frac{\lambda_{1}}{\lambda_{2}}+O_{3}(\epsilon) \\
\underline{d_{T, W}}(x) \leq \frac{\log d}{\lambda_{2}}+2+O_{4}(\epsilon)
\end{gathered}
$$

By combining these two theorems, we manage to separate coordinates $(Z, W)$ in terms of local dimensions.

Corollary 1.3. Let $f$ be an endomorphism of $\mathbb{P}^{2}$ of degree $d \geq 2$. Let $\nu$ be an ergodic dilating measure whose exponents $\lambda_{1}>\lambda_{2}$ do not resonate. For every $\epsilon>0$ and for $\nu$-almost every $x$, there exist holomorphic coordinates $(Z, W)$ in a neighbourhood of $x$ such that

$$
\underline{d_{T, W}}(x) \leq \frac{\log d}{\lambda_{2}}+2+O_{4}(\epsilon) \quad, \quad \frac{\log d}{\lambda_{2}}+2-O_{1}(\epsilon) \leq \overline{d_{T, Z}}(x)
$$

In the three preceding results, the coordinates $(Z, W)$ come from a normal form Theorem for the inverse branches of $f^{n}$, see Section 2. They depend on $\hat{x}$ in the natural extension $(\hat{f}, \hat{\nu})$, we shall denote them by $\left(Z_{\hat{x}}^{\epsilon}, W_{\hat{x}}^{\epsilon}\right)$. The coordinates $(Z, W)$ at a point $x$ are coordinates of type $\left(Z_{\hat{x}}^{\epsilon}, W_{\hat{x}}^{\epsilon}\right)$ where $\pi_{0}(\hat{x})=x$. The coordinate $W_{\hat{x}}^{\epsilon}$ is always invariant by the shift $\hat{f}, Z_{\hat{x}}^{\epsilon}$ is invariant when the exponents do not resonate. Since the current $T$ is $f$-invariant, the functions $d_{T, Z}(\hat{x})$ and $d_{T, W}(\hat{x})$ are $\hat{f}$-invariant, hence $\hat{\nu}$-almost everywhere constant, see Proposition 2.4 . We shall denote them by $\underline{d_{T, Z}}(\nu)$ and $\underline{d_{T, W}}(\nu)$. We have the same properties for the upper dimensions.

## Dimension of ergodic currents $S$

Theorem 1.4. Let $f$ be an endomorphism of $\mathbb{P}^{2}$ of degree $d \geq 2$. Let $S$ be a $(1,1)$-closed positive current on $\mathbb{P}^{2}$. We assume that the support of $S$ contains an ergodic measure $\nu$ of entropy $h_{\nu}>\log d$ whose exponents satisfy $\lambda_{1}>\lambda_{2}$ and do not resonate. Then there exists a function $O_{5}(\epsilon)$ satisfying the following properties. For every $\epsilon>0$, there exist $x \in \operatorname{Supp} \nu$ and a holomorphic coordinate $Z$ in the neighbourhood of $x$ such that:

$$
\overline{d_{S, Z}}(x) \geq 2+\frac{h_{\nu}-\log d}{\lambda_{2}}-O_{5}(\epsilon) .
$$

In particular, $S$ has a local directional dimension $>2$ at some $x \in \operatorname{Supp} \nu$.
This result is localized at a point $x$ because $S$ is not assumed to be $f$-invariant. For every closed positive current $S$ containing an ergodic dilating measure $\nu$ with exponents $\lambda_{1} \geq \lambda_{2}$, de Thélin-Vigny proved in [12] that there exists $x \in \operatorname{supp}(\nu)$ such that

$$
\begin{equation*}
\overline{d_{S}}(x) \geq 2 \frac{\lambda_{2}}{\lambda_{1}}+\frac{h_{\nu}-\log d}{\lambda_{2}} . \tag{1.3}
\end{equation*}
$$

Theorem 1.4 improves this estimate by replacing $\lambda_{2} / \lambda_{1}$ by 1 for a coordinate $Z$. When $\lambda_{1}=\lambda_{2}$, (1.1) shows that this substitution is valid for every coordinate $Z$. Our lower bound answers a question of [12] in a directional way. In the framework of invertible and meromorphic mappings, the preprint [9] gives a lower bound $>2$ for the dimensions of the Green currents $T^{ \pm}$by using coding techniques and laminar properties of $T^{ \pm}$.

### 1.2 Dimension of dilating measures

Let $f$ be an endomorphism of $\mathbb{P}^{2}$ of degree $d \geq 2$ and let $\nu$ be an ergodic measure. We refer to [23] and [26] for the beginning of this Section. The dimension of $\nu$ is defined by

$$
\operatorname{dim}_{H}(\nu):=\inf \left\{\operatorname{dim}_{H}(A), A \text { Borel set of } \mathbb{P}^{2}, \nu(A)=1\right\} .
$$

The lower and upper local dimensions of $\nu$ are

$$
\underline{d_{\nu}}:=\liminf _{r \rightarrow 0} \frac{\log \left(\nu\left(B_{x}(r)\right)\right)}{\log r}, \overline{d_{\nu}}:=\limsup _{r \rightarrow 0} \frac{\log \left(\nu\left(B_{x}(r)\right)\right)}{\log r},
$$

these limits are $\nu$-almost everywhere constant, by ergodicity of $\nu$. If $a \leq \underline{d_{\nu}} \leq \overline{d_{\nu}} \leq b$, then $a \leq \operatorname{dim}_{H}(\nu) \leq b$. If $\nu$ is dilating, we have the classical inequalities $\frac{h_{\nu}}{\lambda_{1}} \leq \underline{d_{\nu}} \leq$ $\overline{d_{\nu}} \leq \frac{h_{\nu}}{\lambda_{2}}$ where $\lambda_{1} \geq \lambda_{2}$ are the exponents of $\nu$. For the equilibrium measure $\mu$, BinderDeMarco [4] conjectured the formula:

$$
\begin{equation*}
\operatorname{dim}_{H}(\mu)=\frac{\log d}{\lambda_{1}}+\frac{\log d}{\lambda_{2}} \tag{1.4}
\end{equation*}
$$

which generalizes the one-dimensional Mañé's formula [22]. The article [4] proves the following upper bound for polynomial mappings

$$
\begin{equation*}
\operatorname{dim}_{H}(\mu) \leq 4-\frac{2\left(\lambda_{1}+\lambda_{2}\right)-\log d^{2}}{\lambda_{1}} \tag{1.5}
\end{equation*}
$$

the article [15] extends this upper bound in a meromorphic context. For every dilating measure $\nu$, we have at our disposal the following lower bound proved in [19]:

$$
\begin{equation*}
\operatorname{dim}_{H}(\nu) \geq \underline{d_{\nu}} \geq \frac{\log d}{\lambda_{1}}+\frac{h_{\nu}-\log d}{\lambda_{2}} \tag{1.6}
\end{equation*}
$$

We obtain a new upper bound which comes closer to Binder-DeMarco's conjecture (1.4).
Theorem 1.5. Let $f$ be an endomorphism of $\mathbb{P}^{2}$ of degree $d \geq 2$. Let $\nu$ be an ergodic dilating measure, of exponents $\lambda_{1}>\lambda_{2}$ and whose support is contained in the support of $\mu$. Then

$$
\underline{d_{\nu}} \leq \frac{\log d}{\lambda_{1}}+\frac{\log d}{\lambda_{2}}+2\left(1-\frac{\lambda_{2}}{\lambda_{1}}\right)
$$

Moreover, if the exponents do not resonate, then:

$$
\underline{d_{\nu}} \leq \frac{\log d}{\lambda_{1}}+\frac{\log d}{\lambda_{2}}+2 \min \left(1-\frac{\lambda_{2}}{\lambda_{1}} ; \frac{\lambda_{1}}{\lambda_{2}}-1\right)
$$

The proof extends the arguments of Theorem 1.2.

### 1.3 Dimension of semi-extremal endomorphisms

We say that $f$ is extremal if the exponents $\lambda_{1} \geq \lambda_{2}$ of its equilibrium measure $\mu$ satisfy $\lambda_{1}=\lambda_{2}=\frac{1}{2} \log d$. The articles $[1,3,15]$ characterize these endomorphisms by the equivalent properties:

1. $\operatorname{dim}_{H}(\mu)=4$.
2. $\mu \ll L e b_{\mathbb{P}^{2}}:=\omega \wedge \omega$.
3. $T$ is a $(1,1)$ positive smooth form on an open set of $\mathbb{P}^{2}$.
4. $f$ is a Lattès map: there exist a complex torus $\mathbb{C}^{2} / \Lambda$, an affine dilation $D$ on this torus and a finite galoisian covering $\sigma: \mathbb{C}^{2} / \Lambda \rightarrow \mathbb{P}^{2}$ such that the following diagram commutes


There exist such applications for every degree $d \geq 2$, see for instance [18].
We say that $f$ is semi-extremal if the exponents $\lambda_{1} \geq \lambda_{2}$ of its equilibrium measure $\mu$ satisfy $\lambda_{1}>\lambda_{2}=\frac{1}{2} \log d$. Using (1.5) and (1.6) one sees that the Binder-DeMarco's conjecture (1.4) holds for these mappings:

$$
\begin{equation*}
\operatorname{dim}_{H}(\mu)=2+\frac{\log d}{\lambda_{1}} \tag{1.7}
\end{equation*}
$$

Classical examples of semi-extremal endomorphisms are suspensions of one-dimensional Lattès maps, they satisfy $\mu \ll T \wedge \omega$. More generally,

Theorem 1.6 (Dujardin [17]). If $\mu \ll T \wedge \omega$, then $f$ is semi-extremal.
In [17] Dujardin asked if $\mu \ll T \wedge \omega$ implies the existence of a one-dimensional Lattès factor for $f$. Theorem 1.7 below provides one step in that direction. It shows that, from a theoritical dimensional point of view, these endomorphisms look like suspensions of one-dimensional Lattès maps: the dimension of $T$ is maximal equal to 4 in a coordinate $Z$, and equal to $2+\log d / \lambda_{1}$ (the dimension of $\mu$, see (1.7)) in a coordinate $W$. The functions $O_{1}(\epsilon), O_{2}(\epsilon)$ come from Theorem 1.1.
Theorem 1.7. Let $f$ be an endomorphism of $\mathbb{P}^{2}$ of degree $d \geq 2$. We assume that $\mu \ll T \wedge \omega$ and that $\overline{d_{\mu}}=\underline{d_{\mu}}$. We also assume that the exponents $\lambda_{1}>\lambda_{2}=\frac{1}{2} \log d$ of $\mu$ do not resonate. For every $\epsilon>0$ and for $\mu$-almost every $x \in \mathbb{P}^{2}$, there exist holomorphic coordinates $(Z, W)$ in a neighbourhoord of $x$ such that

$$
4-O_{1}(\epsilon) \leq \overline{d_{T, Z}}(x) \quad \text { and } \quad 2+\frac{\log d}{\lambda_{1}}-O_{2}(\epsilon) \leq \overline{d_{T, W}}(x) \leq 2+\frac{\log d}{\lambda_{1}}
$$

### 1.4 Organization of the article

Sections 2, 3 et 4 are devoted to normal forms, the geometry of inverse branches and separated sets. Sections 5 et 6 establish Theorems 1.1, 1.4 and 1.7, the proofs rest on Theorem 5.2 which relies on the lower local dimension $\underline{d_{\nu}}$ (the lower bound (1.6) is therefore crucial to deduce these results). Theorems 1.2 and 1.5 are proved in Sections 7 and 8. Section 9 brings together technical results.

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## 2 Normal forms and Oseledec-Poincaré coordinates

### 2.1 Natural extension and normal forms

Let $\mathcal{C}_{f}$ be the critical set of $f$, this is an algebraic subset of $\mathbb{P}^{2}$. If $\nu$ is an ergodic dilating measure, then $x \mapsto \log |J a c f(x)| \in L^{1}(\nu)$, which implies $\nu\left(\mathcal{C}_{f}\right)=0$. Let $X$ be the $f$-invariant Borel set $\mathbb{P}^{2} \backslash \cup_{n \in \mathbb{Z}} f^{n}\left(\mathcal{C}_{f}\right)$ and let

$$
\hat{X}:=\left\{\hat{x}=\left(x_{n}\right)_{n \in \mathbb{Z}} \in X^{\mathbb{Z}}, \quad x_{n+1}=f\left(x_{n}\right)\right\} .
$$

Let $\hat{f}$ be the left shift on $\hat{X}$ and $\pi_{0}(\hat{x}):=x_{0}$. There exists a unique $\hat{f}$-invariant measure $\hat{\nu}$ on $\hat{X}$ such that $\left(\pi_{0}\right)_{*} \hat{\nu}=\nu$. We set $\hat{x}_{n}:=\hat{f}^{n}(\hat{x})$ for $n \in \mathbb{Z}$. A function $\left.\left.\alpha: \hat{X} \rightarrow\right] 0,+\infty\right]$ is $\epsilon$-tempered if $\alpha\left(\hat{f}^{ \pm 1}(\hat{x})\right) \geq e^{-\epsilon} \alpha(\hat{x})$. For every $\hat{x} \in X$ we denote by $f_{\hat{x}}^{-n}$ the inverse branch of $f^{n}$ defined in a neighbourhood of $x_{0}$ with values in a neighbourhood of $x_{-n}$. The articles [2] and [21] provide normal forms for these mappings.

Theorem 2.1. ([2, Proposition 4.3]) Let $f$ be an endomorphism of $\mathbb{P}^{2}$ of degree $d \geq 2$. Let $\nu$ be an ergodic dilating measure with exponents $\lambda_{1}>\lambda_{2}$. Let $\epsilon>0$.

There exists an $\hat{f}$-invariant Borel set $\hat{F} \subset \hat{X}$ such that $\hat{\nu}(\hat{F})=1$ and satisfying the following properties. There exist $\epsilon$-tempered functions $\left.\left.\eta_{\epsilon}, \rho_{\epsilon}: \hat{F} \rightarrow\right] 0,1\right]$ and $\beta_{\epsilon}, L_{\epsilon}, M_{\epsilon}$ : $\hat{F} \rightarrow[1,+\infty[$ and for every $\hat{x} \in \hat{F}$, there exists an injective holomorphic mapping

$$
\xi_{\hat{x}}^{\epsilon}: B_{x_{0}}\left(\eta_{\epsilon}(\hat{x})\right) \rightarrow \mathbb{D}^{2}\left(\rho_{\epsilon}(\hat{x})\right)
$$

such that the following diagram commutes for every $n \geq n_{\epsilon}(\hat{x})$ :

and such that

1. $\forall(p, q) \in B_{x_{0}}\left(\eta_{\epsilon}(\hat{x})\right), \frac{1}{2} d(p, q) \leq\left|\xi_{\hat{x}}^{\epsilon}(p)-\xi_{\hat{x}}^{\epsilon}(q)\right| \leq \beta_{\epsilon}(\hat{x}) d(p, q)$.
2. $\operatorname{Lip}\left(f_{\hat{x}}^{-n}\right) \leq L_{\epsilon}(\hat{x}) e^{-n \lambda_{2}+n \epsilon}$ on $B_{x_{0}}\left(\eta_{\epsilon}(\hat{x})\right)$.
3. if $\lambda_{1} \notin\left\{k \lambda_{2}, k \geq 2\right\}, R_{n, \hat{x}}(z, w)=\left(\alpha_{n, \hat{x}} z, \beta_{n, \hat{x}} w\right)$,
if $\lambda_{1}=k \lambda_{2}$ where $k \geq 2, R_{n, \hat{x}}(z, w)=\left(\alpha_{n, \hat{x}} z, \beta_{n, \hat{x}} w\right)+\left(\gamma_{n, \hat{x}} w^{k}, 0\right)$, with
(a) $e^{-n \lambda_{1}-n \epsilon} \leq\left|\alpha_{n, \hat{x}}\right| \leq e^{-n \lambda_{1}+n \epsilon}$ and $\left|\gamma_{n, \hat{x}}\right| \leq M_{\epsilon}(\hat{x}) e^{-n \lambda_{1}+n \epsilon}$,
(b) $e^{-n \lambda_{2}-n \epsilon} \leq\left|\beta_{n, \hat{x}}\right| \leq e^{-n \lambda_{2}+n \epsilon}$.

Remark 2.2. The diagram commutes for every $n \in\left\{1, \ldots, n_{\epsilon}(\hat{x})\right\}$ for the germs of the mappings, see [2]. The integer $n_{\epsilon}(\hat{x})$ is the smallest integer such that $L_{\epsilon}(\hat{x}) e^{-n \lambda_{2}+n \epsilon} \leq$ $e^{-n \epsilon}$, so that $L_{\epsilon}(\hat{x}) e^{-n \lambda_{2}+n \epsilon} \eta_{\epsilon}(\hat{x}) \leq e^{-n \epsilon} \eta_{\epsilon}(\hat{x}) \leq \eta_{\epsilon}\left(\hat{x}_{-n}\right)$. Item 2 thus ensures that $f_{\hat{x}}^{-n}\left(B_{x_{0}}\left(\eta_{\epsilon}(\hat{x})\right) \subset B_{x_{-n}}\left(\eta_{\epsilon}\left(\hat{x}_{-n}\right)\right)\right.$.

We shall need the following Lemma. Let $n_{1}(L)$ be the smallest integer $n$ satisfying $L / 4 \leq e^{n \epsilon}$. The first item uses the upper bound for $\operatorname{Lip}\left(f_{\hat{x}}^{-n}\right)$ provided by Theorem 2.1. The second item comes from [15, Proposition 3.1].

Lemma 2.3. Let $\hat{x} \in \hat{F}$ such that $\eta_{\epsilon}(\hat{x}) \geq \eta$ and $L_{\epsilon}(\hat{x}) \leq L$. If $n \geq n_{1}(L)$ and $r \leq \eta$,

1. $f_{\hat{x}_{n}}^{-n}\left(B_{x_{n}}(r / 4)\right) \subset B_{x_{0}}\left(r e^{-n \lambda_{2}+3 n \epsilon}\right)$ and $f_{\hat{x}}^{-n}\left(B_{x_{0}}(r / 4)\right) \subset B_{x_{-n}}\left(r e^{-n \lambda_{2}+3 n \epsilon}\right)$.
2. $B_{x_{0}}\left(r e^{-n \lambda_{1}-4 n \epsilon}\right) \subset f_{\hat{x}_{n}}^{-n}\left(B_{x_{n}}\left(\frac{r}{4} e^{-2 n \epsilon}\right)\right)$.

### 2.2 Oseledec-Poincaré coordinates

Let $\nu$ be an ergodic dilating measure of exponents $\lambda_{1}>\lambda_{2}$. We assume that the exponents do not resonate, which means that $\lambda_{1} \notin\left\{k \lambda_{2}, k \geq 2\right\}$. Let $\epsilon>0$, and let us apply Theorem 2.1. For every $\hat{x} \in \hat{F}$ we denote by $\left(Z_{\hat{x}}^{\epsilon}, \bar{W}_{\hat{x}}^{\epsilon}\right)$ the coordinates of $\xi_{\hat{x}}^{\epsilon}$. The commutative diagram given by Theorem 2.1 implies:

$$
\begin{equation*}
Z_{\hat{x}-n}^{\epsilon} \circ f_{\hat{x}}^{-n}=\alpha_{n, \hat{x}} \times Z_{\hat{x}}^{\epsilon}, \quad W_{\hat{x}-n}^{\epsilon} \circ f_{\hat{x}}^{-n}=\beta_{n, \hat{x}} \times W_{\hat{x}}^{\epsilon} . \tag{2.1}
\end{equation*}
$$

Hence, $f_{\hat{x}}^{-n}$ multiplies the first coordinate by $e^{-n \lambda_{1} \pm n \epsilon}$ and multiplies the second coordinate by $e^{-n \lambda_{2} \pm n \epsilon}$. Let us note that the second property holds in the resonant case $\lambda_{1} \in\left\{k \lambda_{2}, k \geq 2\right\}$. We shall name the collection of local holomorphic coordinates

$$
(Z, W)_{\epsilon}:=\left(Z_{\hat{x}}^{\epsilon}, W_{\hat{x}}^{\epsilon}\right)_{\hat{x} \in \hat{F}}
$$

Oseledec-Poincaré coordinates for $(f, \nu)$. Using (2.1) and the fact that the Green current is $f$-invariant, we obtain the following Proposition.

Proposition 2.4. Let $f$ be an endomorphism of $\mathbb{P}^{2}$ of degree $d \geq 2$ and let $T$ be its Green current. Let $\nu$ be an ergodic dilating measure of exponents $\lambda_{1}>\lambda_{2}$. Let $(Z, W)_{\epsilon}$ be Oseledec-Poincaré coordinates for $(f, \nu)$. Then there exists a $\hat{f}$-invariant Borel set $\hat{\Lambda}_{T} \subset \hat{F}$ of $\hat{\nu}$-mesure 1 such that

1. $\hat{x} \mapsto \overline{d_{T, W}}(\hat{x})$ and $\hat{x} \mapsto \underline{d_{T, W}}(\hat{x})$ are $\hat{f}$-invariant on $\hat{\Lambda}_{T}$.
2. $\hat{x} \mapsto \overline{d_{T, Z}}(\hat{x})$ and $\hat{x} \mapsto \underline{d_{T, Z}}(\hat{x})$ are $\hat{f}$-invariant on $\hat{\Lambda}_{T}$ if $\lambda_{1} \notin\left\{k \lambda_{2}, k \geq 2\right\}$.

In particular, if the exponents do not resonate, these functions are constant $\hat{\nu}$-almost everywhere. We shall denote them by

$$
\overline{d_{T, Z}}(\nu), \underline{d_{T, Z}}(\nu), \overline{d_{T, W}}(\nu), \underline{d_{T, W}}(\nu)
$$

Proof. We prove the invariance of $\overline{d_{T, W}}(\hat{x})$, the same arguments hold for the other functions. For every $z \in \mathbb{P}^{2} \backslash \mathcal{C}_{f}$ we denote

$$
a(z):=\frac{1}{2}\left\|\left(D_{z} f\right)^{-1}\right\|^{-1}, \gamma(z):=\min \left\{a(z)\|f\|_{\mathcal{C}^{2}, \mathbb{P}^{2}}^{-1}, 1\right\}
$$

Then [6, Lemme 2] asserts that $f$ is injective on $B_{z}(\gamma(z))$ and

$$
\forall r \in[0, \gamma(z)], B_{f(z)}(a(z) r) \subset f\left(B_{z}(r)\right)
$$

Let $\hat{x} \in \hat{F}$. Since $x_{n} \notin \mathcal{C}_{f}$ for every $n \in \mathbb{Z}$, we obtain for every $r \leq \gamma\left(x_{0}\right)$ :

$$
\begin{equation*}
T \wedge\left(\frac{i}{2} d W_{\hat{f}(\hat{x})}^{\epsilon} \wedge d \overline{W_{\hat{f}(\hat{x})}^{\epsilon}}\right)\left[B_{f\left(x_{0}\right)}\left(a\left(x_{0}\right) r\right)\right] \leq T \wedge\left(\frac{i}{2} d W_{\hat{f}(\hat{x})}^{\epsilon} \wedge d \overline{W_{\hat{f}(\hat{x})}^{\epsilon}}\right)\left[f\left(B_{x_{0}}(r)\right)\right] \tag{2.2}
\end{equation*}
$$

Since $f$ is injective on $B_{x_{0}}(r)$, we can change the variables to get:

$$
\begin{equation*}
T \wedge\left(\frac{i}{2} d W_{\hat{f}(\hat{x})}^{\epsilon} \wedge d \overline{W_{\hat{f}(\hat{x})}^{\epsilon}}\right)\left[f\left(B_{x_{0}}(r)\right)\right]=\int_{B_{x_{0}}(r)} f^{*} T \wedge f^{*}\left(\frac{i}{2} d W_{\hat{f}(\hat{x})}^{\epsilon} \wedge d \overline{W_{\hat{f}(\hat{x})}^{\epsilon}}\right) . \tag{2.3}
\end{equation*}
$$

Now let us recall that

$$
\begin{equation*}
f^{*} T=d T \text { and } f^{*}\left(\frac{i}{2} d W_{\hat{f}(\hat{x})}^{\epsilon} \wedge d \overline{W_{\hat{f}(\hat{x})}^{\epsilon}}\right)=|c(\hat{x})|^{2} \frac{i}{2} d W_{\hat{x}}^{\epsilon} \wedge d \overline{W_{\hat{x}}^{\epsilon}}, \tag{2.4}
\end{equation*}
$$

where the second equality comes from (2.1) by setting $c(\hat{x})^{-1}:=\beta_{1, \hat{f}(\hat{x})}$, it is valid near $x_{0}$ according to Remark 2.2. By combining (2.2), (2.3) and (2.4) we deduce:

$$
T \wedge\left(\frac{i}{2} d W_{\hat{f}(\hat{x})}^{\epsilon} \wedge d \overline{W_{\hat{f}(\hat{x})}^{\epsilon}}\right)\left[B_{f\left(x_{0}\right)}\left(a\left(x_{0}\right) r\right)\right] \leq d|c(\hat{x})|^{2} T \wedge\left(\frac{i}{2} d W_{\hat{x}}^{\epsilon} \wedge d \overline{W_{\hat{x}}^{\epsilon}}\right)\left[B_{x_{0}}(r)\right]
$$

for every $r$ small enough. Taking the logarithm and dividing by $\log \left(a\left(x_{0}\right) r\right)<0$, we get $\overline{d_{T, W}}(\hat{f}(\hat{x})) \geq \overline{d_{T, W}}(\hat{x})$ by taking limits. Since $\hat{\nu}$ is ergodic, the function $\overline{d_{T, W}}(\hat{x})$ is constant on a Borel set $\hat{\Lambda}_{T}$ of $\hat{\nu}$-measure 1 (see [25, Chapter 1.5]). One can replace it by $\bigcap_{n \in \mathbb{Z}} \hat{f}^{n}\left(\hat{\Lambda}_{T}\right)$ to obtain an invariant set.

Proposition 2.5. Let $f$ be an endomorphism of $\mathbb{P}^{2}$ and let $T$ be its Green current. Let $\nu$ be an ergodic measure. Then the functions $x \mapsto \underline{d_{T}}(x)$ and $x \mapsto \overline{d_{T}}(x)$ are invariant, hence $\nu$-almost everywhere constant. We denote them by $\underline{d_{T}}(\nu)$ and $\overline{d_{T}}(\nu)$.

Proof. The arguments follow the proof of Proposition 2.4. In this case we study the measure $T \wedge \omega$, and we replace the second equality in (2.4) by $f^{*} \omega \leq \rho\left(x_{0}\right) \omega$ on $B_{x_{0}}\left(\gamma\left(x_{0}\right)\right)$, where $\rho\left(x_{0}\right)>0$ is a large enough positive constant.

## 3 Geometry of the inverse branches and uniformizations

Let $\nu$ be an ergodic dilating measure of exponents $\lambda_{1}>\lambda_{2}$. Let $\epsilon>0$ and let $(Z, W)_{\epsilon}$ be Oseledec-Poincaré coordinates for $(f, \nu)$. Our aim is to construct, for every $\delta>0$, a Borel set $\hat{\Lambda}_{\epsilon} \subset \hat{X}$ satisfying $\hat{\nu}\left(\hat{\Lambda}_{\epsilon}\right) \geq 1-\delta / 2$ which provides convenient uniformizations.

### 3.1 Dynamical balls

The dynamical distance is defined by $d_{n}(x, y):=\max _{0 \leq k \leq n} d\left(f^{k}(x), f^{k}(y)\right)$. We denote by $B_{n}(x, r)$ the ball centered at $x$ and of radius $r$ for the distance $d_{n}$.

Lemma 3.1. There exist $r_{0}>0, n_{2} \geq 1$ and $C \subset \mathbb{P}^{2}$ such that $\nu(C) \geq 1-\delta / 8$ and satisfying the following properties: for every $x \in C$ and every $n \geq n_{2}$ :

$$
\begin{gathered}
\nu\left(B_{n}\left(x, r_{0} / 8\right)\right) \geq e^{-n h_{\nu}-\epsilon n} . \\
\forall r \leq r_{0}, \nu\left(B_{n}(x, 5 r)\right) \leq \nu\left(B_{n}\left(x, 5 r_{0}\right)\right) \leq e^{-n h_{\nu}+\epsilon n} .
\end{gathered}
$$

Proof. Brin-Katok Theorem [8] ensures that there exists $C_{1} \subset \mathbb{P}^{2}$ of full $\nu$-measure such that for every $x \in C_{1}$ :

$$
\lim _{r \rightarrow 0}\left(\liminf _{n \rightarrow+\infty} \frac{-1}{n} \log \nu\left(B_{n}(x, r)\right)\right)=\lim _{r \rightarrow 0}\left(\limsup _{n \rightarrow+\infty} \frac{-1}{n} \log \nu\left(B_{n}(x, r)\right)\right)=h_{\nu}
$$

Hence, for every $x \in C_{1}$ there exists $r_{0}(x)$ such that $r \leq r_{0}(x)$ implies

$$
\liminf _{n \rightarrow+\infty} \frac{-1}{n} \log \left(\nu\left(B_{n}(x, 5 r)\right)\right) \geq h_{\nu}-\epsilon / 2 \text { and } \limsup _{n \rightarrow+\infty} \frac{-1}{n} \log \left(\nu\left(B_{n}(x, r / 8)\right)\right) \leq h_{\nu}+\epsilon / 2
$$

Let $r_{0}$ such that $C_{2}:=\left\{x \in C_{1}, r_{0}(x) \geq r_{0}\right\}$ satisfies $\nu\left(C_{2}\right) \geq 1-\delta / 16$. For every $x \in C_{2}$, there exists $n_{2}(x)$ such that $n \geq n_{2}(x)$ implies

$$
\begin{gathered}
\nu\left(B_{n}\left(x, r_{0} / 8\right)\right) \geq e^{-n h_{\nu}-\epsilon n} \\
\forall r \leq r_{0}, \nu\left(B_{n}(x, 5 r)\right) \leq \nu\left(B_{n}\left(x, 5 r_{0}\right)\right) \leq e^{-n h_{\nu}+\epsilon n}
\end{gathered}
$$

Let $n_{2} \geq 1$ such that $C:=C_{2} \cap\left\{x \in C_{1}, n_{2}(x) \leq n_{2}\right\}$ satisfies $\nu(C) \geq 1-\delta / 8$.
For every $L>0$, let $m_{L}$ be the smallest integer $m$ such that $L e^{-m\left(\lambda_{2}+\epsilon\right)} \leq 1$ and let $n_{3}(L)$ be the smallest integer larger than $m_{L}$ such that $e^{-n \epsilon} \leq M^{-m_{L}}$, where $M:=$ $\max \left\{\|D f\|_{\infty, \mathbb{P}^{2}}, 1\right\}$.

Lemma 3.2. Let $\hat{x} \in \hat{F}$ such that $\eta_{\epsilon}(\hat{x}) \geq \eta$ and $L_{\epsilon}(\hat{x}) \leq L$. For every $n \geq n_{3}(L)$ and every $r \leq \eta$,

$$
f_{\hat{x}_{n}}^{-n}\left(B_{x_{n}}\left(r e^{-2 n \epsilon}\right)\right) \subset B_{n}\left(x_{0}, r\right)
$$

Proof. Let us observe that for every $0 \leq k \leq n, f^{k}$ is injective on $f_{\hat{x}_{n}}^{-n}\left(B_{x_{n}}\left(r e^{-2 n \epsilon}\right)\right)$ and that $f^{k} f_{\hat{x}_{n}}^{-n}=f_{\hat{x}_{n}}^{-n+k}$. By setting $p=n-k$, it suffices to show that

$$
\begin{equation*}
\forall p \in \llbracket 0, n \rrbracket \quad f_{\hat{x}_{n}}^{-p}\left(B_{x_{n}}\left(r e^{-2 n \epsilon}\right)\right) \subset B_{x_{n-p}}(r) . \tag{3.1}
\end{equation*}
$$

To simplify let us set $m:=m_{L}$ et $n_{3}:=n_{3}(L)$. We immediately have

$$
\begin{equation*}
\forall n \geq n_{3}, \forall p \in \llbracket 0, n \rrbracket, \quad f_{\hat{x}_{n}}^{-p}\left(B_{x_{n}}\left(r e^{-2 n \epsilon}\right)\right) \subset f_{\hat{x}_{n}}^{-p}\left(B_{x_{n}}\left(\frac{r}{M^{m}} e^{-n \epsilon}\right)\right) \tag{3.2}
\end{equation*}
$$

To verify (3.1), we shall consider separately the cases $p \leq m$ and $p>m$. We know that for every $p, \operatorname{Lip} f_{\hat{x}_{n}}^{-p} \leq L\left(\hat{x}_{n}\right) e^{-p \lambda_{2}+p \epsilon} \leq L e^{n \epsilon} e^{-p \lambda_{2}+p \epsilon}$ on $B_{x_{n}}\left(\eta_{\epsilon}\left(\hat{x}_{n}\right)\right)$, which contains $B_{x_{n}}\left(\eta e^{-n \epsilon}\right)$. Hence for every $n \geq n_{3} \geq m, p \in \llbracket m, n \rrbracket$ and $r \leq \eta$ :

$$
f_{\hat{x}_{n}}^{-p}\left(B_{x_{n}}\left(r e^{-n \epsilon}\right)\right) \subset B_{x_{n-p}}\left(r e^{-n \epsilon} L e^{n \epsilon} e^{-p \lambda_{2}+p \epsilon}\right)=B_{x_{n-p}}\left(r L e^{-p \lambda_{2}+p \epsilon}\right) \subset B_{x_{n-p}}(r)
$$

Since $M^{m} \geq 1$ this implies for every $n \geq n_{3} \geq m, p \in \llbracket m, n \rrbracket$ and $r \leq \eta$ :

$$
\begin{equation*}
f_{\hat{x}_{n}}^{-p}\left(B_{x_{n}}\left(\frac{r}{M^{m}} e^{-n \epsilon}\right)\right) \subset B_{x_{n-p}}\left(\frac{r}{M^{m}}\right) \tag{3.3}
\end{equation*}
$$

Thus, by using (3.2) and $M^{m} \geq 1$ :

$$
\begin{equation*}
\forall p \in \llbracket m, n \rrbracket, \quad f_{\hat{x}_{n}}^{-p}\left(B_{x_{n}}\left(r e^{-2 n \epsilon}\right)\right) \subset B_{x_{n-p}}(r) \tag{3.4}
\end{equation*}
$$

We have proved (3.1) for $p \in \llbracket m, n \rrbracket$. Let us show this inclusion for $p \in \llbracket 0, m \rrbracket$. For every $p \in \llbracket 0, m \rrbracket$, let us set $p=m-p^{\prime}$ where $p^{\prime} \in \llbracket 0, m \rrbracket$. Then

$$
f_{\hat{x}_{n}}^{-p}\left(B_{x_{n}}\left(\frac{r}{M^{m}} e^{-n \epsilon}\right)\right)=f^{p^{\prime}}\left(f_{\hat{x}_{n}}^{-m}\left(B_{x_{n}}\left(\frac{r}{M^{m}} e^{-n \epsilon}\right)\right) \subset f^{p^{\prime}}\left(B_{x_{n-m}}\left(\frac{r}{M^{m}}\right)\right)\right.
$$

where the inclusion comes from (3.3) with $p=m$. We deduce:

$$
\begin{equation*}
\forall p \in \llbracket 0, m \rrbracket, \quad f_{\hat{x}_{n}}^{-p}\left(B_{x_{n}}\left(r e^{-2 n \epsilon}\right)\right) \subset B_{x_{n-m+p^{\prime}}}\left(\frac{r}{M^{m}} M^{p}\right) \subset B_{x_{n-p}}(r) \tag{3.5}
\end{equation*}
$$

We finally obtain (3.1) by combining (3.4) and (3.5).

### 3.2 Pullback of the Fubini-Study form $\omega$

Let $\nu$ be an ergodic dilating measure of exponents $\lambda_{1}>\lambda_{2}$. Let $(Z, W)_{\epsilon}$ be OseledecPoincaré coordinates for $(f, \nu)$. Let $n_{4}(\beta)$ be the smallest integer $n$ such that $e^{-n \epsilon} \leq \beta^{-1}$.
Proposition 3.3. Let $\hat{x} \in \hat{F}$ such that $\eta_{\epsilon}(\hat{x}) \geq \eta$ and $\beta_{\epsilon}(\hat{x}) \leq \beta$. If $n \geq \max \left\{n_{4}(\beta), n_{\epsilon}\left(\hat{x}_{n}\right)\right\}$ and if $r \leq \eta$, then we have on $f_{\hat{x}_{n}}^{-n}\left(B_{x_{n}}\left(r e^{-n \epsilon}\right)\right)$ :

1. $\left(f^{n}\right)^{*} \omega \geq e^{-4 n \epsilon+2 n \lambda_{1}}\left(\frac{i}{2} d Z_{\hat{x}}^{\epsilon} \wedge d \overline{Z_{\hat{x}}^{\epsilon}}\right)$ if the exponents do not resonate.
2. $\left(f^{n}\right)^{*} \omega \geq e^{-4 n \epsilon+2 n \lambda_{2}}\left(\frac{i}{2} d W_{\hat{x}}^{\epsilon} \wedge d \overline{W_{\hat{x}}^{\epsilon}}\right)$.

The remainder of this Section is devoted to the proof. Theorem 2.1 gives

$$
\left(f^{n}\right)^{*} \omega=\left(\xi_{\hat{x}}^{\epsilon}\right)^{*}\left(\left(R_{n, \hat{x}_{n}}\right)^{-1}\right)^{*}\left(\left(\xi_{\hat{x}_{n}}\right)^{-1}\right)^{*} \omega
$$

on $f_{\hat{x}_{n}}^{-n}\left(B_{x_{n}}\left(r e^{-n \epsilon}\right)\right)$. Let $\omega_{0}:=\frac{i}{2} d z \wedge d \bar{z}+\frac{i}{2} d w \wedge d \bar{w}$ be the standard form on $\mathbb{D}^{2}$.
Lemma 3.4. Let $\hat{x} \in \hat{F}$ such that $\eta_{\epsilon}(\hat{x}) \geq \eta$ and $\beta_{\epsilon}(\hat{x}) \leq \beta$. For every $n \geq n_{4}(\beta)$ and $r \leq \eta$, we have on $\xi_{\hat{x}_{n}}^{\epsilon}\left(B_{x_{n}}\left(r e^{-n \epsilon}\right)\right)$ :

$$
e^{-2 n \epsilon} \omega_{0} \leq\left(\left(\xi_{\hat{x}_{n}}^{\epsilon}\right)^{-1}\right)^{*} \omega \leq 2 \omega_{0}
$$

Proof. For every $p=(z, w)$ and $p^{\prime}=\left(z^{\prime}, w^{\prime}\right)$ in $\xi_{\hat{x}_{n}}^{\epsilon}\left(B_{x_{n}}\left(r e^{-n \epsilon}\right)\right)$, we have

$$
e^{-n \epsilon} \beta^{-1} d\left(p, p^{\prime}\right) \leq\left|\left(\xi_{\hat{x}_{n}}^{\epsilon}\right)^{-1}(p)-\left(\xi_{\hat{x}_{n}}^{\epsilon}\right)^{-1}\left(p^{\prime}\right)\right| \leq 2 d\left(p, p^{\prime}\right)
$$

This implies for every $n \geq n_{4}(\beta)$ and $(z, w) \in \xi_{\hat{x}_{n}}^{\epsilon}\left(B_{x_{n}}\left(r e^{-n \epsilon}\right)\right)$ :

$$
\forall u \in \mathbb{C}^{2}, \quad e^{-2 n \epsilon}|u| \leq\left|D_{(z, w)}\left(\xi_{\hat{x}_{n}}^{\epsilon}\right)^{-1}(u)\right| \leq 2|u|
$$

This provides the desired estimates.

Lemma 3.5. Let $\hat{x} \in \hat{F}$. If $n \geq n_{\epsilon}\left(\hat{x}_{n}\right)$, then

1. $\left(\left(R_{n, \hat{x}_{n}}\right)^{-1}\right)^{*} \omega_{0} \geq e^{2\left(n \lambda_{1}-n \epsilon\right)} \frac{i}{2} d z \wedge d \bar{z}$ if the exponents do not resonate.
2. $\left(\left(R_{n, \hat{x}_{n}}\right)^{-1}\right)^{*} \omega_{0} \geq e^{2\left(n \lambda_{2}-n \epsilon\right)} \frac{i}{2} d w \wedge d \bar{w}$.

Proof. We use the fact that the linear part of $R_{n, \hat{x}_{n}}$ is diagonal with coefficients $e^{-n \epsilon-n \lambda_{1}} \leq$ $\left|\alpha_{n, \hat{x}_{n}}\right| \leq e^{n \epsilon-n \lambda_{1}}$ and $e^{-n \epsilon-n \lambda_{2}} \leq\left|\beta_{n, \hat{x}_{n}}\right| \leq e^{n \epsilon-n \lambda_{2}}$ (see Theorem 2.1) and the fact that the (1,1)-forms $\frac{i}{2} d z \wedge d \bar{z}$ and $\frac{i}{2} d w \wedge d \bar{w}$ are positive.

To end the proof of Proposition 3.3, we observe that for every $\hat{x} \in \hat{F}$ :

$$
\left(\xi_{\hat{x}}^{\epsilon}\right)^{*}\left(\frac{i}{2} d z \wedge d \bar{z}\right)=\left(\frac{i}{2} d Z_{\hat{x}}^{\epsilon} \wedge d \overline{Z_{\hat{x}}^{\epsilon}}\right) \quad\left(\xi_{\hat{x}}^{\epsilon}\right)^{*}\left(\frac{i}{2} d w \wedge d \bar{w}\right)=\left(\frac{i}{2} d W_{\hat{x}}^{\epsilon} \wedge d \overline{W_{\hat{x}}^{\epsilon}}\right)
$$

which follows from the definitions of $Z_{\hat{x}}^{\epsilon}$ and $W_{\hat{x}}^{\epsilon}$.

### 3.3 Uniformizations

Let $\nu$ be an ergodic dilating measure of exponents $\lambda_{1}>\lambda_{2}$ and let $\delta>0$. Let $\epsilon>0$ and let $(Z, W)_{\epsilon}$ be Oseledec-Poincaré coordinates for $(f, \nu)$.

## Measure of dynamical balls

We apply Lemma 3.1. There exist $r_{0}>0, n_{2} \geq 1$ and $C \subset \mathbb{P}^{2}$ such that $\nu(C) \geq$ $1-\delta / 8$ and for every $x \in C$ and $n \geq n_{2}$ :

$$
\begin{gathered}
\nu\left(B_{n}\left(x, r_{0} / 8\right)\right) \geq e^{-n h_{\nu}-\epsilon n}, \\
\forall r \leq r_{0}, \nu\left(B_{n}(x, 5 r)\right) \leq \nu\left(B_{n}\left(x, 5 r_{0}\right)\right) \leq e^{-n h_{\nu}+\epsilon n} .
\end{gathered}
$$

We denote $\Lambda^{(1)}:=\pi_{0}^{-1}(C) \cap \hat{F}$.

## Control of the functions $n_{\epsilon}, \rho_{\epsilon}, L_{\epsilon}, \eta_{\epsilon}, \beta_{\epsilon}$ of Theorem 2.1

Let $n_{0}, \rho_{0}>0, L_{0}>0, \eta_{0}>0$ and $\beta_{0}>0$ such that

$$
\begin{equation*}
\Lambda^{(2)}:=\left\{\hat{x} \in \hat{F}, n_{\epsilon}(\hat{x}) \leq n_{0}, \rho_{\epsilon}(\hat{x}) \geq \rho_{0}, L_{\epsilon}(\hat{x}) \leq L_{0}, \eta_{\epsilon}(\hat{x}) \geq \eta_{0}, \beta_{\epsilon}(\hat{x}) \leq \beta_{0}\right\} \tag{3.6}
\end{equation*}
$$

satisfies $\hat{\nu}\left(\Lambda^{(2)}\right) \geq 1-\delta / 8$.

## Uniformization of the dimension of the current.

Let $S$ be a positive closed current on $\mathbb{P}^{2}$ whose support contains the support of $\nu$. Let $r_{1}>0$ such that

$$
\Lambda^{(3)}:=\left\{\begin{array}{cc}
\hat{x} \in \hat{F}, \quad \forall r \leq r_{1}, & r^{\overline{d_{S, Z}}(\hat{x})+\epsilon} \leq\left(S \wedge\left(\frac{i}{2} d Z_{\hat{\hat{x}}}^{\epsilon} \wedge d \overline{\bar{Z}_{\hat{\hat{\epsilon}}}}\right)\right)\left(B_{x_{0}}(r)\right) \leq r^{d_{S, Z}}(\hat{x})-\epsilon \\
\left(S \wedge\left(\frac{i}{2} d W_{\hat{x}}^{\epsilon} \wedge d \overline{W_{\hat{x}}^{\epsilon}}\right)\right)\left(B_{x_{0}}(r)\right) \leq r^{d_{S, W}}(\hat{x})-\epsilon
\end{array}\right\}
$$

satisfies $\hat{\nu}\left(\Lambda^{(3)}\right) \geq 1-\delta / 8$. In the case of the Green current $T$, the functions $\overline{d_{T, Z}}, d_{T, Z}$ and $\underline{d_{T, W}}$ are $\hat{\nu}$-almost everywhere constant and denoted $\overline{d_{T, Z}}(\nu), \underline{d_{T, Z}}(\nu)$ and $\underline{d_{T, W}} \overline{(\nu)}$.

## Uniformization of the dimension of the measure

The lower dimension $\underline{d_{\nu}}$ is defined in Section 1.4. Let $r_{2}>0$ such that

$$
D:=\left\{x \in \mathbb{P}^{2}, \forall r \leq r_{2}, \nu\left(B_{x}(r)\right) \leq r \underline{d_{\nu}}-\epsilon\right\}
$$

satisfies $\nu(D) \geq 1-\delta / 8$. We set $\Lambda^{(4)}:=\pi_{0}^{-1}(D) \cap \hat{F}$.
Definition of $\hat{\Lambda}_{\epsilon}, \eta_{1}$ and $N_{\epsilon}$.
The integers $n_{1}(L), n_{3}(L)$ and $n_{4}(\beta)$ were defined before Lemma 2.3, 3.2 and Proposition 3.3. Let $n_{5}$ be the smallest integer $n$ such that $e^{-n \epsilon} \leq 1 / 2$ and $2 e^{-n\left(\lambda_{1}+\epsilon\right)}<1$.

$$
\begin{gathered}
\hat{\Lambda}_{\epsilon}:=\Lambda^{(1)} \cap \Lambda^{(2)} \cap \Lambda^{(3)} \cap \Lambda^{(4)}, \\
\eta_{1}:=\min \left\{\eta_{0}, r_{0}, r_{1}, r_{2}\right\}, \\
N_{\epsilon}:=\max \left\{n_{0}, n_{1}\left(L_{0}\right), n_{2}, n_{3}\left(L_{0}\right), n_{4}\left(\beta_{0}\right), n_{5}\right\}
\end{gathered}
$$

We have $\hat{\nu}\left(\hat{\Lambda}_{\epsilon}\right) \geq 1-\delta / 2$.

## Definition of $\hat{\Delta}_{\epsilon}^{n}$.

We set

$$
\forall n \geq N_{\epsilon}, \quad \hat{\Delta}_{\epsilon}^{n}:=\hat{F} \cap \hat{f}^{-n}\left\{n_{\epsilon}(\hat{x}) \leq n\right\}=\left\{\hat{x} \in \hat{F}, n_{\epsilon}\left(\hat{x}_{n}\right) \leq n\right\}
$$

Since $\hat{\nu}$ is $\hat{f}$-invariant and $\Lambda^{(3)} \subset\left\{n_{\epsilon}(\hat{x}) \leq n\right\}$, we have $\hat{\nu}\left(\hat{\Delta}_{\epsilon}^{n}\right) \geq \hat{\nu}\left(\Lambda^{(3)}\right) \geq 1-\delta / 8$. Hence

$$
\forall n \geq N_{\epsilon}, \quad \hat{\nu}\left(\hat{\Lambda}_{\epsilon} \cap \hat{\Delta}_{\epsilon}^{n}\right) \geq 1-\delta
$$

## 4 Separated sets

A subset $\left\{x_{1}, \ldots, x_{N}\right\} \subset \mathbb{P}^{2}$ is $r$-separated if $d\left(x_{i}, x_{j}\right) \geq r$ for every $i \neq j$. For $A \subset \mathbb{P}^{2}$, a subset $\left\{x_{1}, \ldots, x_{N}\right\} \subset A$ is maximal $r$-separated with respect to $A$ if it is $r$-separated and if for every $y \in A$, there exists $i \in\{1, \ldots N\}$ such that $d\left(y, x_{i}\right)<r$. We use similar definitions for the distance $d_{n}$, in which case we say that the subsets are $(n, r)$ separated.

### 4.1 Elementary separation

Lemma 4.1. Let $f$ be an endomorphism of $\mathbb{P}^{2}$ of degree $d \geq 2$ and let $\nu$ be an ergodic measure. Let $A \subset \pi_{0}\left(\hat{\Lambda}_{\epsilon}\right)$ such that $\nu(A)>0$ and let $\left.\left.c \in\right] 0,1\right]$. Let $n \geq N_{\epsilon}$ and let $\left\{x_{1}, \ldots x_{N_{n}}\right\} \subset A$ be maximal $\left(n, c \eta_{1}\right)$-separated with respect to $A$. Then

1. for every $i \neq j, B_{n}\left(x_{i}, c \eta_{1} / 2\right) \cap B_{n}\left(x_{j}, c \eta_{1} / 2\right)=\emptyset$.
2. $A \subset \cup_{i=1}^{N_{n}} B_{n}\left(x_{i}, c \eta_{1}\right)$.
3. $\nu\left(B_{n}\left(x_{i}, c \eta_{1}\right)\right) \leq e^{-n h_{\nu}+n \epsilon}$.
4. $e^{-n h_{\nu}-n \epsilon} \leq \nu\left(B_{n}\left(x_{i}, c \eta_{1}\right)\right)$ si $c \geq 1 / 8$.
5. $N_{n} \geq \nu(A) e^{n h_{\nu}-n \epsilon}$.

Proof. Item 1 comes from separation, Item 2 from the maximal property, Items 3 and 4 from Section 3.3, because $n \geq N_{\epsilon}, c \eta_{1} \leq \eta_{1}$ and $x_{i} \in C$. Items 2 and 3 then imply $\nu(A) \leq \sum_{i=1}^{N_{n}} \nu\left(B_{n}\left(x_{i}, c \eta_{1}\right)\right) \leq N_{n} e^{-n h_{\nu}+n \epsilon}$, which gives Item 5.

### 4.2 Concentrated separation

Lemma 4.1 applied with $c=1 / 4$ gives $\nu\left(B_{n}\left(x_{i}, \eta_{1} / 4\right)\right) \geq e^{-n h_{\nu}-n \epsilon}$ for every $x_{i}$ in a maximal ( $n, \eta_{1} / 4$ )-separated subset of $A$. We shall see that it is possible to select a large number of $x_{i}$ such that

$$
\nu\left(B_{n}\left(x_{i}, \eta_{1} / 4\right) \cap A\right) \geq e^{-n h_{\nu}-2 n \epsilon} .
$$

We take the arguments of de Thélin-Vigny in $\left[12\right.$, Section 6]. Let $n_{\delta}$ be the smallest integer $n$ such that $e^{-n \epsilon} \leq \delta / 2$.

Lemma 4.2. Let $A \subset \pi_{0}\left(\hat{\Lambda}_{\epsilon}\right)$ such that $\nu(A) \geq \delta$. For every $n \geq \max \left\{N_{\epsilon}, n_{\delta}\right\}$, there exists a ( $n, \eta_{1} / 4$ )-separated subset $\left\{x_{1}, \ldots, x_{N_{n, 2}}\right\}$ of $A$ such that

1. for every $i \neq j, B_{n}\left(x_{i}, \eta_{1} / 8\right) \cap B_{n}\left(x_{j}, \eta_{1} / 8\right)=\emptyset$.
2. for every $1 \leq i \leq N_{n, 2}, \nu\left(B_{n}\left(x_{i}, \eta_{1} / 4\right) \cap A\right) \geq e^{-n h_{\nu}-2 \epsilon n}$.
3. $N_{n, 2} \geq \nu(A) e^{n h_{\nu}-2 n \epsilon}$.
4. $e^{-n h_{\nu}-n \epsilon} \leq \nu\left(B_{n}\left(x_{i}, c \eta_{1}\right)\right)$ si $c \geq 1 / 8$.

Proof. Let us apply Lemma 4.1 with $c=1 / 4$ and $n \geq \max \left\{N_{\epsilon}, n_{\delta}\right\}$. There exists a maximal ( $n, \eta_{1} / 4$ )-separated subset $\left\{x_{1}, \ldots, x_{N_{n, 1}}\right\}$ of $A$ satisfying:

- for every $i \neq j, B_{n}\left(x_{i}, \eta_{1} / 8\right) \cap B_{n}\left(x_{j}, \eta_{1} / 8\right)=\emptyset$,
$-A \subset \cup_{i=1}^{N_{n, 1}} B_{n}\left(x_{i}, \eta_{1} / 4\right)$,
- $e^{-n h_{\nu}-n \epsilon} \leq \nu\left(B_{n}\left(x_{i}, \eta_{1} / 8\right)\right)$,
- $N_{n, 1} \geq \nu(A) e^{n h_{\nu}-n \epsilon}$.

Let us set $I:=\left\{1 \leq i \leq N_{n, 1}, \nu\left(B_{n}\left(x_{i}, \eta_{1} / 4\right) \cap A\right) \geq e^{-n h_{\nu}-2 \epsilon n}\right\}$. Let $N_{n, 2}$ be the cardinality of $I$, and assume that $I=\llbracket 1, N_{n, 2} \rrbracket$ (we may adapt the sums below if $N_{n, 2}=$ $0)$. We want to bound $N_{n, 2}$ from below. We know that $A \subset \cup_{i=1}^{N_{n, 1}} B_{n}\left(x_{i}, \eta_{1} / 4\right)$, hence

$$
\nu(A) \leq \sum_{i=1}^{N_{n, 2}} \nu\left(B_{n}\left(x_{i}, \eta_{1} / 4\right) \cap A\right)+\sum_{i=N_{n, 2}+1}^{N_{n, 1}} \nu\left(B_{n}\left(x_{i}, \eta_{1} / 4\right) \cap A\right) .
$$

If $i \notin \llbracket 1, N_{n, 2} \rrbracket$, we have $\nu\left(B_{n}\left(x_{i}, \eta_{1} / 4\right) \cap A\right)<e^{-n h_{\nu}-2 \epsilon n}$ by definition of $I$. Otherwise, $\nu\left(B_{n}\left(x_{i}, \eta_{1} / 4\right)\right) \leq e^{-n h_{\nu}+\epsilon n}$ since $x_{i} \in C$. This implies

$$
\begin{equation*}
\nu(A) \leq N_{n, 2} e^{-n h_{\nu}+\epsilon n}+\left(N_{n, 1}-N_{n, 2}\right) e^{-n h_{\nu}-2 \epsilon n} . \tag{4.1}
\end{equation*}
$$

Let us give an upper bound for $N_{n, 1}$. Since the balls $B_{n}\left(x_{i}, \eta_{1} / 8\right)$ are pairwise disjoint and since $\nu\left(B_{n}\left(x_{i}, \eta_{1} / 8\right)\right) \geq e^{-n h_{\nu}-\epsilon n}$, we get

$$
e^{n h_{\nu}+\epsilon n} \geq N_{n, 1} \geq N_{n, 1}-N_{n, 2}
$$

Combining this and (4.1), we obtain

$$
\nu(A) \leq N_{n, 2} e^{-n h_{\nu}+\epsilon n}+e^{-\epsilon n} .
$$

Since $n \geq n_{\delta}$, we have $e^{-n \epsilon} \leq \delta / 2 \leq \nu(A) / 2$, and hence $N_{n, 2} \geq \nu(A) e^{n h_{\nu}-\epsilon n} / 2$. Finally $N_{n, 2} \geq \nu(A) e^{n h_{\nu}-2 \epsilon n}$ since $n \geq N_{\epsilon} \geq n_{5}$.

Now we put in $B_{n}\left(x, \eta_{1} / 2\right)$ a lot of balls whose centers are in $B_{n}\left(x, \eta_{1} / 4\right) \cap A$.
Lemma 4.3. Let $A \subset \pi_{0}\left(\hat{\Lambda}_{\epsilon}\right)$ such that $\nu(A)>0$. Let $x \in A$ and let $n \geq N_{\epsilon}$ such that

$$
\nu\left(B_{n}\left(x, \eta_{1} / 4\right) \cap A\right) \geq e^{-n h_{\nu}-2 n \epsilon} .
$$

Let $\left\{y_{1}, \ldots, y_{M_{n}}\right\}$ be a maximal $2 \eta_{1} e^{-n \lambda_{1}-4 n \epsilon}$-separated subset in $B_{n}\left(x, \eta_{1} / 4\right) \cap A$.

1. for every $i \neq j, B\left(y_{i}, \eta_{1} e^{-n \lambda_{1}-4 n \epsilon}\right) \cap B\left(y_{j}, \eta_{1} e^{-n \lambda_{1}-4 n \epsilon}\right)=\emptyset$.
2. $B_{n}\left(x, \eta_{1} / 4\right) \cap A \subset \cup_{i=1}^{M_{n}} B\left(y_{i}, 2 \eta_{1} e^{-n \lambda_{1}-4 n \epsilon}\right)$.
3. $B\left(y_{i}, \eta_{1} e^{-n \lambda_{1}-4 n \epsilon}\right) \subset B_{n}\left(x, \eta_{1} / 2\right)$.
4. $M_{n} \geq e^{-n h_{\nu}-2 n \epsilon}\left(\frac{1}{2 \eta_{1}} e^{n \lambda_{1}+4 n \epsilon}\right)^{\frac{d_{\nu}}{}-\epsilon}$.

Proof. Item 1 comes from separation, Item 2 from the maximal property. Lemmas 2.3 then 3.2 give

$$
B\left(y_{i}, \eta_{1} e^{-n \lambda_{1}-4 n \epsilon}\right) \subset f_{\hat{y}_{i, n}}^{-n}\left(B_{y_{i, n}}\left(\eta_{1} e^{-2 n \epsilon} / 4\right) \subset B_{n}\left(y_{i}, \eta_{1} / 4\right) .\right.
$$

Since $y_{i} \in B_{n}\left(x, \eta_{1} / 4\right)$, we have $B_{n}\left(y_{i}, \eta_{1} / 4\right) \subset B_{n}\left(x, \eta_{1} / 2\right)$, which yields Item 3. Item 2 implies

$$
\nu\left(B_{n}\left(x, \eta_{1} / 4\right) \cap A\right) \leq \sum_{i=1}^{M_{n}} \nu\left(B\left(y_{i}, 2 \eta_{1} e^{-n \lambda_{1}-4 n \epsilon}\right)\right) .
$$

By assumption, the left hand side is larger than $e^{-n h_{\nu}-2 n \epsilon}$. For the right hand side, since $n \geq N_{\epsilon} \geq n_{5}$, we have $2 \eta_{1} e^{-n \lambda_{1}-n \epsilon}<\eta_{1} \leq r_{2}$ and thus

$$
\left.\nu\left(B\left(y_{i}, 2 \eta_{1} e^{-n \lambda_{1}-4 n \epsilon}\right)\right) \leq\left(2 \eta_{1} e^{-n \lambda_{1}-4 n \epsilon}\right)\right)^{d_{\nu}-\epsilon}
$$

by using $y_{i} \in A \subset \pi_{0}\left(\hat{\Lambda}_{\epsilon}\right) \subset D$. This shows $e^{-n h_{\nu}-2 n \epsilon} \leq M_{n}\left(2 \eta_{1} e^{-n \lambda_{1}-4 n \epsilon}\right) \underline{d_{\nu}-\epsilon}$.

## 5 Lower bounds for the directional dimensions of $T$

Let $\nu$ be an ergodic dilating measure whose exponents $\lambda_{1}>\lambda_{2}$ do not resonate. Let $\epsilon>0$ and let $(Z, W)_{\epsilon}$ be Oseledec-Poincaré coordinates for $(f, \nu)$. We have $\overline{d_{T, Z}}(\nu):=\overline{d_{T, Z}}(\hat{x})$ and $\overline{d_{T, W}}(\nu)=\overline{d_{T, W}}(\hat{x})$ for $\hat{\nu}$-almost every $\hat{x}$ (see Proposition 2.4). In this Section we prove Theorems 5.1 and 5.2. These results specify in a directional way Theorems A and B of de Thélin-Vigny [12] concerning the dimension of $T$. We use below arguments of [12] by replacing the lower bound (obtained in [12] by slicing arguments)

$$
\left(f^{n}\right)^{*} \omega \geq e^{2 n \lambda_{2}} e^{-n \epsilon} \omega
$$

by the two lower bounds (obtained by normal forms Theorem 2.1)

$$
\left(f^{n}\right)^{*} \omega \geq e^{2 n \lambda_{1}} e^{-n \epsilon} d Z \wedge d \bar{Z} \text { and }\left(f^{n}\right)^{*} \omega \geq e^{2 n \lambda_{2}} e^{-n \epsilon} d W \wedge d \bar{W}
$$

Theorem 5.1 uses elementary separation (Lemma 4.1), Theorem 5.2 uses concentrated separation (Lemmas 4.2 and 4.3). Theorem 5.2 implies Theorem 1.1 (via the lower bound (1.6) for $\underline{d_{\nu}}$ ) and Theorem 1.7.

### 5.1 First lower bounds for the upper dimensions $d_{T, Z}$ et $d_{T, W}$

Theorem 5.1. Let $f$ be an endomorphism of $\mathbb{P}^{2}$ of degree $d \geq 2$. Let $\nu$ be an ergodic dilating measure whose exponents $\lambda_{1}>\lambda_{2}$ do not resonate. There exist functions $O_{5}(\epsilon), O_{6}(\epsilon)$ satisfying the following properties. Let $\epsilon>0$ and let $(Z, W)_{\epsilon}$ be OseledecPoincaré coordinates for $(f, \nu)$. Then

$$
\begin{gathered}
\overline{d_{T, Z}}(\nu) \geq 2+\frac{h_{\nu}-\log d}{\lambda_{1}}-O_{5}(\epsilon) \\
\overline{d_{T, W}}(\nu) \geq 2 \frac{\lambda_{2}}{\lambda_{1}}+\frac{h_{\nu}-\log d}{\lambda_{1}}-O_{6}(\epsilon)
\end{gathered}
$$

Proof. Let us denote $\overline{d_{T, Z}}:=\overline{d_{T, Z}}(\nu)$. For first estimate, we are going to show

$$
\begin{equation*}
\left(\lambda_{1}+4 \epsilon\right)\left(\overline{d_{T, Z}}-2+\epsilon\right)+13 \epsilon \geq h_{\nu}-\log d \tag{5.1}
\end{equation*}
$$

which provides

$$
\overline{d_{T, Z}} \geq 2+\frac{h_{\nu}-\log d-13 \epsilon}{\lambda_{1}+4 \epsilon}-\epsilon=: 2+\frac{h_{\nu}-\log d}{\lambda_{1}}-O_{5}(\epsilon)
$$

Let $\delta>0$. Let $\hat{\Lambda}_{\epsilon}$ and $N_{\epsilon}$ be given by Section 3.3. For every $n \geq N_{\epsilon}$ we set $A_{n}:=$ $\pi_{0}\left(\hat{\Lambda}_{\epsilon} \cap \hat{\Delta}_{\epsilon}^{n}\right)$, it satisfies $\nu\left(A_{n}\right) \geq \hat{\nu}\left(\hat{\Lambda}_{\epsilon} \cap \hat{\Delta}_{\epsilon}^{n}\right) \geq 1-\delta>0$. Lemma 4.1 applied with $c=1$ yields a maximal $\left(n, \eta_{1}\right)$-separated subset $\left\{x_{1}, \ldots, x_{N_{n}}\right\}$ of $A_{n}$ with

$$
\begin{equation*}
N_{n} \geq \nu\left(A_{n}\right) e^{n h_{\nu}-n \epsilon} \geq(1-\delta) e^{n h_{\nu}-n \epsilon} \tag{5.2}
\end{equation*}
$$

For every $i$, let us choose $\hat{x}_{i} \in \hat{\Lambda}_{\epsilon} \cap \hat{\Delta}_{\epsilon}^{n}$ such that $\pi_{0}\left(\hat{x}_{i}\right)=x_{i}$. From Proposition 9.3, we get $d^{n}=\int_{\mathbb{P}^{2}}\left(\left(f^{n}\right)_{*} T\right) \wedge \omega$. Therefore

$$
d^{n} \geq \sum_{i=1}^{N_{n}} \int_{\mathbb{P}^{2}}\left(f^{n}\right)_{*}\left(1_{B_{n}\left(x_{i}, \eta_{1} / 2\right)} T\right) \wedge \omega \geq \sum_{i=1}^{N_{n}} \int_{\mathbb{P}^{2}}\left(1_{B_{n}\left(x_{i}, \eta_{1} / 2\right)} T\right) \wedge\left(f^{n}\right)^{*} \omega .
$$

By using Lemmas 2.3 and 3.2 with $\hat{x}_{i} \in \hat{\Lambda}_{\epsilon}$, we get

$$
\begin{equation*}
B_{x_{i}}\left(\frac{\eta_{1}}{2} e^{-n \lambda_{1}-4 n \epsilon}\right) \subset f_{\hat{x}_{i, n}}^{-n}\left(B_{x_{i, n}}\left(\frac{\eta_{1}}{8} e^{-2 n \epsilon}\right)\right) \subset B_{n}\left(x_{i}, \frac{\eta_{1}}{8}\right) . \tag{5.3}
\end{equation*}
$$

Since $T \wedge\left(f^{n}\right)^{*} \omega$ is a positive measure, we deduce

$$
d^{n} \geq \sum_{i=1}^{N_{n}} \int_{\mathbb{P}^{2}}\left(1_{B_{x_{i}}\left(\frac{\eta_{1}}{2} e^{-n \lambda_{1}-4 n \epsilon}\right)} T\right) \wedge\left(f^{n}\right)^{*} \omega .
$$

Thanks to the first inclusion of (5.3) and $\hat{x}_{i} \in \hat{\Delta}_{\epsilon}^{n}$ (which implies $n \geq n_{\epsilon}\left(\hat{x}_{i, n}\right)$ ), we can use Proposition 3.3 to bound $\left(f^{n}\right)^{*} \omega$ from below. We obtain

$$
d^{n} \geq \sum_{i=1}^{N_{n}} e^{2 n \lambda_{1}-4 n \epsilon}\left(T \wedge \frac{i}{2} d Z_{\hat{x}_{i}}^{\epsilon} \wedge d \overline{Z_{\hat{x}_{i}}^{\epsilon}}\right)\left(B_{x_{i}}\left(\frac{\eta_{1}}{2} e^{-n \lambda_{1}-4 n \epsilon}\right)\right) .
$$

Since $\hat{x_{i}} \in \hat{\Lambda}_{\epsilon} \subset \Lambda^{(3)}$ and $\eta_{1} \leq r_{1}$, we deduce

$$
d^{n} \geq \sum_{i=1}^{N_{n}} e^{2 n \lambda_{1}-4 n \epsilon}\left(\frac{\eta_{1}}{2} e^{-n \lambda_{1}-4 n \epsilon}\right)^{\overline{d_{T, Z}}+\epsilon} .
$$

Finally, we use the estimate (5.2):

$$
d^{n} \geq \nu\left(A_{n}\right)\left(\frac{\eta_{1}}{2}\right)^{\overline{d_{T, Z}^{\epsilon}}+\epsilon} e^{n\left(h_{\nu}-\left(\lambda_{1}+4 \epsilon\right)\left(\overline{d_{T, Z}^{\epsilon}}-2+\epsilon\right)-13 \epsilon\right)}
$$

where $\nu\left(A_{n}\right) \geq(1-\delta)$. By taking logarithm and dividing by $n$, we get (5.1) when $n \rightarrow \infty$. The second estimate concerning the coordinate $W$ is proved in a similar way, by using Proposition 3.3 to bound $\left(f^{n}\right)^{*} \omega$ from below. We precisely get

$$
\begin{equation*}
\left(\lambda_{1}+4 \epsilon\right)\left(\overline{d_{T, W}}+\epsilon\right)+5 \epsilon \geq h_{\nu}-\log d+2 \lambda_{2} \tag{5.4}
\end{equation*}
$$

which yields

$$
\overline{d_{T, W}} \geq \frac{2 \lambda_{2}}{\lambda_{1}+4 \epsilon}+\frac{h_{\nu}-\log d-5 \epsilon}{\lambda_{1}+4 \epsilon}-\epsilon=: 2 \frac{\lambda_{2}}{\lambda_{1}}+\frac{h_{\nu}-\log d}{\lambda_{1}}-O_{6}(\epsilon) .
$$

This completes the proof of Theorem 5.1.

### 5.2 Proof of Theorem 1.1

Theorem 1.1 is a consequence of Theorem 5.2 below and of the lower bound (1.6).
Theorem 5.2. Let $f$ be an endomorphism of $\mathbb{P}^{2}$ of degree $d \geq 2$. Let $\nu$ be an ergodic dilating measure whose exponents $\lambda_{1}>\lambda_{2}$ do not resonate. There exist functions $O_{1}(\epsilon), O_{2}(\epsilon)$ satisfying the following properties. Let $\epsilon>0$ and let $(Z, W)_{\epsilon}$ be OseledecPoincaré coordinates for $(f, \nu)$. Then

$$
\begin{gathered}
\overline{d_{T, Z}}(\nu) \geq 2+\underline{d_{\nu}}-\frac{\log d}{\lambda_{1}}-O_{1}(\epsilon) \\
\overline{d_{T, W}}(\nu) \geq 2 \frac{\lambda_{2}}{\lambda_{1}}+\underline{d_{\nu}}-\frac{\log d}{\lambda_{1}}-O_{2}(\epsilon)
\end{gathered}
$$

Proof. Let us set $\overline{d_{T, Z}}:=\overline{d_{T, Z}}(\nu)$. We are going to show

$$
\begin{equation*}
\left(\lambda_{1}+4 \epsilon\right)\left(\overline{d_{T, Z}}-\underline{d_{\nu}}+2 \epsilon\right)+8 \epsilon \geq 2 \lambda_{1}-\log d \tag{5.5}
\end{equation*}
$$

This yields as desired

$$
\begin{equation*}
\overline{d_{T, Z}}-\underline{d_{\nu}} \geq \frac{2 \lambda_{1}}{\lambda_{1}+4 \epsilon}-\frac{\log d-8 \epsilon}{\lambda_{1}+4 \epsilon}-2 \epsilon=: 2-\frac{\log d}{\lambda_{1}}-O_{1}(\epsilon) \tag{5.6}
\end{equation*}
$$

Let $\delta>0$. Let $\hat{\Lambda}_{\epsilon}$ and $N_{\epsilon}$ be given by Section 3.3. For every $n \geq N_{\epsilon}$ we set $A_{n}:=$ $\pi_{0}\left(\hat{\Lambda}_{\epsilon} \cap \hat{\Delta}_{\epsilon}^{n}\right)$, it satisfies $\nu\left(A_{n}\right) \geq 1-\delta$. Let $\left\{x_{1}, \ldots, x_{N_{n, 2}}\right\}$ be a $\left(n, \eta_{1} / 4\right)$-separated subset of $A_{n}$ provided by Lemma 4.2. Then for every $x_{i}$, we set a $2 \eta_{1} e^{-n \lambda_{1}-4 n \epsilon}$-separated subset $\left\{y_{1}^{i}, \ldots, y_{M_{n}}^{i}\right\}$ given by Lemma 4.3. For every $i$ we choose $\hat{x}_{i} \in \hat{\Lambda}_{\epsilon} \cap \hat{\Delta}_{\epsilon}^{n}$ such that $\pi_{0}\left(\hat{x}_{i}\right)=x_{i}$, and for every $j$ we choose $\hat{y}_{j}^{i} \in \hat{\Lambda}_{\epsilon} \cap \hat{\Delta}_{\epsilon}^{n}$ such that $\pi_{0}\left(\hat{y}_{j}^{i}\right)=y_{j}^{i}$. According to Proposition 9.3, we have $d^{n}=\int_{\mathbb{P}^{2}}\left(f^{n}\right)_{*} T \wedge \omega$, thus

$$
d^{n} \geq \sum_{i=1}^{N_{n, 2}} \sum_{j=1}^{M_{n}} \int_{\mathbb{P}^{2}}\left(f^{n}\right)_{*}\left(1_{B\left(y_{j}^{i}, \eta_{1} e^{-n \lambda_{1}-4 n \epsilon}\right)} T\right) \wedge \omega=\sum_{i=1}^{N_{n, 2}} \sum_{j=1}^{M_{n}} \int_{B\left(y_{j}^{i}, \eta_{1} e^{-n \lambda_{1}-4 n \epsilon}\right)} T \wedge\left(f^{n}\right)^{*} \omega
$$

Lemma 2.3 with $\hat{y}_{j}^{i} \in \hat{\Lambda}_{\epsilon}$ implies

$$
\begin{equation*}
B\left(y_{j}^{i}, \eta_{1} e^{-n \lambda_{1}-4 n \epsilon}\right) \subset f_{\hat{y}_{j, n}^{i}}^{-n}\left(B\left(y_{j, n}^{i}, \frac{\eta_{1}}{4} e^{-2 n \epsilon}\right)\right) \tag{5.7}
\end{equation*}
$$

Since $\hat{y}_{j}^{i} \in \hat{\Delta}_{\epsilon}^{n}$, we can apply Proposition 3.3 to bound $\left(f^{n}\right)^{*} \omega$ from below:

$$
\begin{equation*}
d^{n} \geq \sum_{i=1}^{N_{n, 2}} \sum_{j=1}^{M_{n}} e^{2 n \lambda_{1}-4 n \epsilon}\left(T \wedge \frac{i}{2} d Z_{\hat{y}_{j}^{i}}^{\epsilon} \wedge d \overline{Z_{\hat{y}_{j}^{i}}^{\epsilon}}\right)\left(B_{y_{j}^{i}}\left(\eta_{1} e^{-n \lambda_{1}-4 n \epsilon}\right)\right) \tag{5.8}
\end{equation*}
$$

Now $\hat{y}_{j}^{i} \in \hat{\Lambda}_{\epsilon} \subset \Lambda^{(3)}$ and $n \geq N_{\epsilon}$, hence

$$
d^{n} \geq \sum_{i=1}^{N_{n, 2}} \sum_{j=1}^{M_{n}} e^{2 n \lambda_{1}-4 n \epsilon}\left(\eta_{1} e^{-n \lambda_{1}-4 n \epsilon}\right)^{\overline{d_{T, Z}}+\epsilon}
$$

Finally, we use the lower bounds for $M_{n}$ (Lemma 4.3) and for $N_{n, 2}$ (Lemma 4.2). We obtain for every $n \geq \max \left\{N_{\epsilon}, n_{1-\delta}\right\}$ :

$$
d^{n} \geq(1-\delta) e^{n h_{\nu}-2 n \epsilon} \cdot e^{-n h_{\nu}-2 n \epsilon}\left(\frac{1}{2 \eta_{1}} e^{n \lambda_{1}+4 n \epsilon}\right)^{\underline{d_{\nu}-\epsilon}} \cdot e^{2 n \lambda_{1}-4 n \epsilon}\left(\eta_{1} e^{-n \lambda_{1}-4 n \epsilon}\right)^{\overline{d_{T, Z}}+\epsilon}
$$

Let us note that the entropy $h_{\nu}$ disappear for the benefit of $\underline{d_{\nu}}$, and we get

$$
\begin{equation*}
d^{n} \geq c e^{-8 n \epsilon}\left(e^{n \lambda_{1}+4 n \epsilon}\right)^{\frac{d_{\nu}}{}-\overline{d_{T, Z}^{\epsilon}}-2 \epsilon} e^{2 n \lambda_{1}} \tag{5.9}
\end{equation*}
$$

where $c:=(1-\delta) \eta_{1}^{\overline{d_{T, Z}}+\epsilon} /\left(2 \eta_{1}\right) \underline{d_{\nu}}-\epsilon$. Taking logarithm and then dividing by $n$, we obtain (5.5) when $n \rightarrow+\infty$. Similarly, we can prove

$$
\left(\lambda_{1}+4 \epsilon\right)\left(\overline{d_{T, W}}-\underline{d_{\nu}}+2 \epsilon\right)+8 \epsilon \geq 2 \lambda_{2}-\log d
$$

by using again Proposition 3.3 to bound $\left(f^{n}\right)^{*} \omega$ from below.

## 6 Currents $S$ and semi-extremal endomorphisms

### 6.1 Proof of Theorem 1.4

Let $S$ be a $(1,1)$ closed positive current of $\mathbb{P}^{2}$. If $S$ does not satisfy $f^{*} S=d S$, the directional dimensions may be not $\hat{\nu}$-almost everywhere constant (see Proposition 2.4). In this case, in the manner of de Thélin-Vigny [12], we take on an adapted definition and obtain the following result. The functions $O_{5}(\epsilon), O_{6}(\epsilon)$ were defined in Theorem 5.1.

Theorem 6.1. Let $f$ be an endomorphism of $\mathbb{P}^{2}$ of degree $d \geq 2$. Let $S$ be a $(1,1)$ closed positive current of $\mathbb{P}^{2}$, of mass 1 . Let $\nu$ be an ergodic dilating measure whose exponents $\lambda_{1}>\lambda_{2}$ do not resonate. We assume that $\operatorname{Supp}(\nu) \subset \operatorname{Supp} S$. Let $\epsilon>0$ and let $(Z, W)_{\epsilon}$ be Oseledec-Poincaré coordinates for $(f, \nu)$. For every $\hat{\Lambda} \subset \hat{F}$ such that $\hat{\nu}(\hat{\Lambda})>0$, we set

$$
\overline{d_{S, Z}}(\hat{\Lambda}):=\sup _{\hat{x} \in \hat{\Lambda}} \overline{d_{S, Z}}(\hat{x}), \quad \overline{d_{S, W}}(\hat{\Lambda}):=\sup _{\hat{x} \in \hat{\Lambda}} \overline{d_{S, W}}(\hat{x}) .
$$

Then

$$
\begin{gathered}
\overline{d_{S, Z}}(\hat{\Lambda}) \geq 2+\frac{h_{\nu}-\log d}{\lambda_{1}}-O_{5}(\epsilon) \\
\overline{d_{S, W}}(\hat{\Lambda}) \geq 2 \frac{\lambda_{2}}{\lambda_{1}}+\frac{h_{\nu}-\log d}{\lambda_{1}}-O_{6}(\epsilon)
\end{gathered}
$$

Proof. Let $2 \delta:=\hat{\nu}(\hat{\Lambda})$. We construct $\hat{\Lambda}_{\epsilon}$ et $N_{\epsilon}$ for the current $S$ as in Section 3.3. We have $\hat{\nu}\left(\hat{\Lambda}_{\epsilon} \cap \hat{\Delta}_{\epsilon}^{n}\right) \geq 1-\delta$ for every $n \geq N_{\epsilon}$, thus $\hat{\nu}\left(\hat{\Lambda} \cap \hat{\Lambda}_{\epsilon} \cap \hat{\Delta}_{\epsilon}^{n}\right) \geq \delta>0$. We follow the arguments of Theorem 5.1. Lemma 4.1 applied to $A_{n}:=\pi_{0}\left(\hat{\Lambda} \cap \hat{\Lambda}_{\epsilon} \cap \hat{\Delta}_{\epsilon}^{n}\right)$ and $c=1$ provides a maximal $\left(n, \eta_{1}\right)$-separated subset $\left\{x_{1}, \ldots, x_{N_{n}}\right\}$ of $A_{n}$. For every $i$, let us
choose $\hat{x}_{i} \in \hat{\Lambda} \cap \hat{\Lambda}_{\epsilon} \cap \hat{\Delta}_{\epsilon}^{n}$ such that $\pi_{0}\left(\hat{x}_{i}\right)=x_{i}$. According to Proposition 9.3 we have $d^{n}=\int_{\mathbb{P}^{2}}\left(f^{n}\right)_{*} S \wedge \omega$, hence

$$
d^{n}=\int_{\mathbb{P}^{2}}\left(\left(f^{n}\right)_{*} S\right) \wedge \omega \geq \sum_{i=1}^{N_{n}} \int_{\mathbb{P}^{2}}\left(1_{B_{n}\left(x_{i}, \eta_{1} / 2\right)} S\right) \wedge\left(f^{n}\right)^{*} \omega
$$

Then we use the inclusions and the lower bound for $\left(f^{n}\right)^{*} \omega$ to obtain:

$$
d^{n} \geq \sum_{i=1}^{N_{n}} e^{2 n \lambda_{1}-4 n \epsilon}\left(S \wedge \frac{i}{2} d Z_{\hat{x}_{i}}^{\epsilon} \wedge d \overline{Z_{\hat{x}_{i}}^{\epsilon}}\right)\left(B_{x_{i}}\left(\frac{\eta_{1}}{2} e^{-n \lambda_{1}-4 n \epsilon}\right)\right)
$$

Since $\hat{x_{i}} \in \hat{\Lambda}_{\epsilon}$ and $n \geq N_{\epsilon}$, we get

$$
d^{n} \geq \sum_{i=1}^{N_{n}} e^{2 n \lambda_{1}-4 n \epsilon}\left(\frac{\eta_{1}}{2} e^{-n \lambda_{1}-4 n \epsilon}\right)^{\overline{d_{S, Z}}\left(\hat{x}_{i}\right)+\epsilon}
$$

Now we use the adapted definition of $\overline{d_{S, Z}}(\hat{\Lambda})$ and the lower estimate (5.2) to obtain

$$
\begin{equation*}
d^{n} \geq \nu\left(A_{n}\right)\left(\frac{\eta_{1}}{2}\right)^{\overline{d_{S, Z}}(\hat{\Lambda})+\epsilon} e^{n h_{\nu}-13 n \epsilon-n\left(\lambda_{1}+4 \epsilon\right)\left(\overline{d_{S, Z}}(\hat{\Lambda})-2+\epsilon\right)} \tag{6.1}
\end{equation*}
$$

where $\nu\left(A_{n}\right) \geq \delta$. The lower bound concerning $W$ is proved in a similar way.
Remark 6.2. Theorem 6.1 is the counterpart of Theorem 5.1 for currents S. Similarly, the counterpart of Theorem 5.2 can be proved with $n \geq \max \left\{N_{\epsilon}, n_{\delta}\right\}$ in the proof.

### 6.2 Proof of Theorem 1.7

Since $T$ is $f$-invariant and $\nu$ is ergodic, we have $\overline{d_{T}}(x)=\overline{d_{T}}(\mu)$ for $\mu$-almost every $x$, see Proposition 2.5. According to Proposition 9.1, $\mu \ll \sigma_{T}$ implies

$$
\begin{equation*}
\overline{d_{T}}(\mu) \leq \overline{d_{\mu}} \tag{6.2}
\end{equation*}
$$

Let us analyse these quantities. On the one hand, Proposition 9.2 yields $\overline{d_{T}}(\mu)=$ $\min \left\{\overline{d_{T, Z}}(x), \overline{d_{T, W}}(x)\right\}$ for $\mu$-almost every $x \in \mathbb{P}^{2}$ and for every holomorphic coordinates $(Z, W)$ near $x$. On the other hand, since $\underline{d_{\mu}}=\overline{d_{\mu}}$, then $\underline{d_{\mu}}=\overline{d_{\mu}}=\operatorname{dim}_{H}(\mu)$, which is equal to $2+\frac{\log d}{\lambda_{1}}$ by (1.7). One deduces from $(6.2)$ that if $\bar{\mu} \ll \sigma_{T}$, then

$$
\begin{equation*}
\min \left\{\overline{d_{T, Z}}(x), \overline{d_{T, W}}(x)\right\} \leq 2+\frac{\log d}{\lambda_{1}} \tag{6.3}
\end{equation*}
$$

Now we use Theorem 5.2. Let $\epsilon>0$ such that $4-O_{1}(\epsilon)>2+\frac{\log d}{\lambda_{1}}$, where the function $O_{1}(\epsilon)$ is defined by (5.6). Let $(Z, W)_{\epsilon}$ Oseledec-Poincaré coordinates for $(f, \mu)$.

First we bound $\overline{\bar{d}_{T, Z}}(\mu)$ from below, then we establish the formula for $\overline{d_{T, W}}(\mu)$ modulo the function $O_{2}(\epsilon)$. If $\mu \ll \sigma_{T}$, then $\lambda_{2}=\frac{1}{2} \log d$ by Theorem 1.6. We deduce from (1.6) that $2+\frac{\log d}{\lambda_{1}} \leq \underline{d_{\mu}}$. Theorem 5.2 then provides

$$
\overline{d_{T, Z}}(\mu) \geq 4-O_{1}(\epsilon) \quad, \quad \overline{d_{T, W}}(\mu) \geq 2+\frac{\log d}{\lambda_{1}}-O_{2}(\epsilon) .
$$

Finally, using $4-O_{1}(\epsilon)>2+\frac{\log d}{\lambda_{1}}$ and (6.3) with Oseledec-Poincaré coordinates $(Z, W)_{\epsilon}$, we get $\overline{d_{T, W}}(\mu) \leq 2+\frac{\log d}{\lambda_{1}}$ as desired.

## $7 \quad$ Upper bounds for the directional dimensions of $T$

In this Section we show Theorem 1.2. Let $\nu$ be an ergodic dilating measure such that $\operatorname{Supp}(\nu) \subset \operatorname{Supp}(\mu)$ and whose exponents $\lambda_{1}>\lambda_{2}$ do not resonate. Let $\epsilon>0$ and let $(Z, W)_{\epsilon}$ be Oseledec-Poincaré coordinates for $(f, \nu)$. We want to prove

$$
\underline{d_{T, Z}}(\nu) \leq \frac{\log d}{\lambda_{2}}+2 \frac{\lambda_{1}}{\lambda_{2}}+O_{3}(\epsilon) \quad \text { and } \quad \underline{d_{T, W}}(\nu) \leq \frac{\log d}{\lambda_{2}}+2+O_{4}(\epsilon) .
$$

We shall directly obtain these upper bounds for $\hat{\nu}$-almost every $\hat{x}$, by using the jacobians of $T \wedge d Z_{\hat{x}}^{\epsilon} \wedge d \overline{Z_{\hat{x}}^{\epsilon}}$ and $T \wedge d W_{\hat{x}}^{\epsilon} \wedge d \overline{W_{\hat{x}}^{\epsilon}}$ with respect to $f$. In particular, we shall not use separated subsets. The Monge-Ampère equation $\mu=T \wedge T$ will be crucial.

### 7.1 Dimensions of the Green current on the equilibrium measure

The following Proposition is proved in Section 7.3.
Proposition 7.1. Let $f$ be an endomorphism of $\mathbb{P}^{2}$ of degree $d \geq 2$ and let $T$ be its Green current. Let $x \in \operatorname{Supp} \mu$ and let $Z$ be a local holomorphic coordinate (submersion) in a neighbourhood $V$ of $x$. Then $T \wedge\left(\frac{i}{2} d Z \wedge d \bar{Z}\right)$ is not the zero measure on $V$.

This implies:
Proposition 7.2. Let $f$ be an endomorphism of $\mathbb{P}^{2}$ of degree $d \geq 2$. Let $\nu$ be an ergodic dilating measure of exponents $\lambda_{1}>\lambda_{2}$ and whose support is contained in the support of $\mu$. Let $\epsilon>0$ and let $(Z, W)_{\epsilon}$ be Oseledec-Poincaré coordinates for $(f, \nu)$. We recall that for every $\hat{x} \in \hat{F},\left(Z_{\hat{x}}^{\epsilon}, W_{\hat{x}}^{\epsilon}\right)$ is defined on $B_{x_{0}}\left(\eta_{\epsilon}(\hat{x})\right)$. Then, for every $0<r<\eta_{\epsilon}(\hat{x})$,

$$
\left[T \wedge \frac{i}{2} d Z_{\hat{x}}^{\epsilon} \wedge d \overline{Z_{\hat{x}}^{\epsilon}}\right]\left(B_{x}(r)\right)>0 \quad \text { and } \quad\left[T \wedge \frac{i}{2} d W_{\hat{x}}^{\epsilon} \wedge d \overline{W_{\hat{x}}^{\epsilon}}\right]\left(B_{x}(r)\right)>0 .
$$

In particular, for every $\delta>0$, there exist $m_{0} \geq 1, L_{0} \geq 1$ and $q_{0} \geq 1$ such that

$$
\hat{\Omega}_{\epsilon}:=\left\{\begin{array}{ll}
\eta_{\epsilon}(\hat{x}) \geq \frac{1}{4 m_{0}}, L(\hat{x}) \leq L_{0}, & {\left[T \wedge \frac{i}{2} d Z_{\hat{x}}^{\epsilon} \wedge d \overline{Z_{\hat{x}}^{\epsilon}}\right]\left(B_{x}\left(\frac{1}{4 m_{0}}\right)\right) \geq \frac{1}{q_{0}}} \\
{\left[T \wedge \frac{i}{2} d W_{\hat{x}}^{\epsilon} \wedge d \overline{W_{\hat{x}}^{\epsilon}}\right]\left(B_{x}\left(\frac{1}{4 m_{0}}\right)\right) \geq \frac{1}{q_{0}}}
\end{array}\right\}
$$

satisfies $\hat{\nu}\left(\hat{\Omega}_{\epsilon}\right) \geq 1-\delta$.

Proof. The first part immediately follows from Proposition 7.1. To prove the second part, let $m_{0} \geq 1$ and $L_{0} \geq 1$ be such that $\hat{\nu}\left\{\eta_{\epsilon} \geq \frac{1}{4 m_{0}}\right\} \cap\left\{L \leq L_{0}\right\} \geq 1-\delta / 2$. Then we choose $q_{0}$ large enough so that $\hat{\nu}\left(\hat{\Omega}_{\epsilon}\right) \geq 1-\delta$.

We define for every $n \geq 1$ :

$$
\hat{\Omega}_{\epsilon}^{n}:=\hat{\Omega}_{\epsilon} \cap \hat{f}^{-n}\left(\hat{\Omega}_{\epsilon}\right) .
$$

Since $\hat{\nu}$ is invariant, we have:

$$
\begin{equation*}
\hat{\nu}\left(\hat{\Omega}_{\epsilon}^{n}\right) \geq 1-2 \delta . \tag{7.1}
\end{equation*}
$$

The following Proposition will be useful to prove Theorems 1.2 and 1.5. $L_{0}$ is defined in Proposition 7.2 and $n_{1}\left(L_{0}\right) \geq 1$ is defined before Lemma 2.3.
Proposition 7.3. Let $f$ be an endomorphism of $\mathbb{P}^{2}$ of degree $d \geq 2$. Let $\nu$ be a ergodic dilating measure of exponents $\lambda_{1}>\lambda_{2}$ and such that $\operatorname{Supp}(\nu) \subset \operatorname{Supp}(\mu)$. For every $n \geq n_{1}\left(L_{0}\right)$ and $\hat{x} \in \hat{\Omega}_{\epsilon}^{n}$, we have:

$$
\begin{aligned}
& {\left[T \wedge \frac{i}{2} d Z_{\hat{x}}^{\epsilon} \wedge d \overline{Z_{\hat{x}}^{\epsilon}}\right]\left(B_{x}\left(\frac{1}{m_{0}} e^{-n \lambda_{2}+3 n \epsilon}\right)\right) \geq \frac{1}{d^{n}} e^{-2 n \lambda_{1}-2 n \epsilon} \frac{1}{q_{0}} \quad \text { if } \lambda_{1} \notin\left\{k \lambda_{2}, k \geq 2\right\},} \\
& {\left[T \wedge \frac{i}{2} d W_{\hat{x}}^{\epsilon} \wedge d \overline{W_{\hat{x}}^{\epsilon}}\right]\left(B_{x}\left(\frac{1}{m_{0}} e^{-n \lambda_{2}+3 n \epsilon}\right)\right) \geq \frac{1}{d^{n}} e^{-2 n \lambda_{2}-2 n \epsilon} \frac{1}{q_{0}} \quad \text { for every } \lambda_{1}>\lambda_{2} .}
\end{aligned}
$$

Proof. Let $\hat{x} \in \hat{\Omega}_{\epsilon}^{n}$ and let

$$
E_{n}:=f_{\hat{x}_{n}}^{-n}\left(B_{x_{n}}\left(\frac{1}{4 m_{0}}\right)\right) .
$$

The inverse branch $f_{\hat{x}_{n}}^{-n}$ is well defined on $B_{x_{n}}\left(\frac{1}{4 m_{0}}\right)$ since $\hat{x}_{n} \in \hat{\Omega}_{\epsilon}$. Let $g_{n}$ be the restriction of $f^{n}$ on $E_{n}$. By using $f_{\hat{x}_{n}}^{-n} \circ g_{n}=I d_{E_{n}}$ and $T=\frac{1}{d^{n}} g_{n}^{*} T$ on $E_{n}$, we obtain

$$
\begin{aligned}
T \wedge \frac{i}{2} d Z_{\hat{x}}^{\epsilon} \wedge d \overline{Z_{\hat{x}}^{\epsilon}} & =\frac{1}{d^{n}} g_{n}^{*} T \wedge g_{n}^{*}\left(f_{\hat{x}_{n}}^{-n}\right)^{*}\left(\frac{i}{2} d Z_{\hat{x}}^{\epsilon} \wedge d \overline{Z_{\hat{x}}^{\epsilon}}\right) \\
& =\frac{1}{d^{n}} g_{n}^{*}\left[T \wedge \frac{i}{2}\left(d Z_{\hat{x}}^{\epsilon} \circ\left(f_{\hat{x}_{n}}^{-n}\right)\right) \wedge d\left(\overline{\left.Z_{\hat{x}}^{\epsilon} \circ\left(f_{\hat{x}_{n}}^{-n}\right)\right)}\right]\right.
\end{aligned}
$$

on the open subset $E_{n}$. Now we use (2.1) to write $Z_{\hat{x}}^{\epsilon} \circ\left(f_{\hat{x}_{n}}^{-n}\right)=\alpha_{n, \hat{x}_{n}} Z_{\hat{x}_{n}}$. Since $\left|\alpha_{n, \hat{x}_{n}}\right|^{2} \geq e^{-2 n \lambda_{1}-2 n \epsilon}$, we get on $E_{n}$ :

$$
\begin{equation*}
T \wedge \frac{i}{2} d Z_{\hat{x}}^{\epsilon} \wedge d \overline{Z_{\hat{x}}^{\epsilon}} \geq \frac{1}{d^{n}} e^{-2 n \lambda_{1}-2 n \epsilon} g_{n}^{*}\left[T \wedge \frac{i}{2} d Z_{\hat{x}_{n}}^{\epsilon} \wedge d \overline{Z_{\hat{x}_{n}}^{\epsilon}}\right] . \tag{7.2}
\end{equation*}
$$

We are going to bound from above the left hand side and to bound from below the right hand side (applied to $E_{n}$ ). Using Lemma 2.3 with $r=1 / m_{0}$ and $n \geq L_{0}$ we obtain $E_{n} \subset B_{x}\left(\frac{1}{m_{0}} e^{-n \lambda_{2}+3 n \epsilon}\right)$, hence

$$
\begin{equation*}
T \wedge \frac{i}{2} d Z_{\hat{x}}^{\epsilon} \wedge d \overline{Z_{\hat{x}}^{\epsilon}}\left(B_{x}\left(\frac{1}{m_{0}} e^{-n \lambda_{2}+3 n \epsilon}\right)\right) \geq T \wedge \frac{i}{2} d Z_{\hat{x}}^{\epsilon} \wedge d \overline{Z_{\hat{x}}^{\epsilon}}\left(E_{n}\right) \tag{7.3}
\end{equation*}
$$

For the right hand side, since $g_{n}$ is injective on $E_{n}$ and $g_{n}\left(E_{n}\right)=B_{x_{n}}\left(\frac{1}{4 m_{0}}\right)$, we get

$$
\begin{equation*}
g_{n}^{*}\left[T \wedge \frac{i}{2}\left(d Z_{\hat{x}_{n}}^{\epsilon} \wedge \overline{d Z_{\hat{x}_{n}}^{\epsilon}}\right)\right]\left(E_{n}\right)=\left[T \wedge \frac{i}{2}\left(d Z_{\hat{x}_{n}}^{\epsilon} \wedge\left(\overline{d Z_{\hat{x}_{n}}^{\epsilon}}\right)\right]\left(B_{x_{n}}\left(\frac{1}{4 m_{0}}\right)\right) \geq \frac{1}{q_{0}},\right. \tag{7.4}
\end{equation*}
$$

where the inequality comes from $\hat{x}_{n} \in \hat{\Omega}_{\epsilon}$. By combining (7.2), (7.3) and (7.4) we obtain

$$
\left[T \wedge \frac{i}{2} d Z_{\hat{x}}^{\epsilon} \wedge d \overline{Z_{\hat{x}}^{\epsilon}}\right]\left(B_{x}\left(\frac{1}{m_{0}} e^{-n \lambda_{2}+3 n \epsilon}\right)\right) \geq \frac{1}{d^{n}} e^{-2 n \lambda_{1}-2 n \epsilon} \frac{1}{q_{0}} .
$$

We use $W_{\hat{x}}^{\epsilon} \circ\left(f_{\hat{x}_{n}}^{-n}\right)=\beta_{n, \hat{x}} W_{\hat{x}_{n}}$ et $\left|\beta_{n, \hat{x}}\right|^{2} \geq e^{-2 n \lambda_{2}-2 n \epsilon}$ to prove the other lower bound.

### 7.2 Proof of Theorem 1.2

We take the notations of Section 7.1. Let

$$
\hat{\Omega}_{\epsilon}:=\limsup _{n \in \mathbb{N}} \hat{\Omega}_{\epsilon} \cap \hat{f}^{-n}\left(\hat{\Omega}_{\epsilon}\right)=\limsup _{n \in \mathbb{N}} \hat{\Omega}_{\epsilon}^{n} .
$$

We have $\hat{\nu}\left(\hat{\Omega}_{\epsilon}\right) \geq 1-2 \delta$ according to (7.1). Let $\hat{x} \in \hat{\Omega}_{\epsilon}$. Then there exists an increasing sequence of intergers $\left(l_{p}\right)_{p}$ such that

$$
\hat{x} \in \hat{\Omega}_{\epsilon} \cap \hat{f}^{-l_{p}}\left(\hat{\Omega}_{\epsilon}\right)=\hat{\Omega}_{\epsilon}^{l_{p}}
$$

for every $p \geq 0$. Proposition 7.3 then asserts for $p$ large enough:

$$
\left[T \wedge \frac{i}{2} d Z_{\hat{x}}^{\epsilon} \wedge d \overline{Z_{\hat{x}}^{\epsilon}}\right]\left(B_{x}\left(\frac{1}{m_{0}} e^{-l_{p} \lambda_{2}+3 l_{p} \epsilon}\right)\right) \geq \frac{1}{d^{l_{p}}} e^{-2 l_{p} \lambda_{1}-2 l_{p} \epsilon} \frac{1}{q_{0}} .
$$

If $p$ is also large enough so that $e^{l_{p} \epsilon} \geq \frac{1}{m_{0}}$ and $\frac{1}{q_{0}} \geq e^{-l_{p} \epsilon}$, we obtain with $r_{p}:=e^{-l_{p}\left(\lambda_{2}-4 \epsilon\right)}$ :

$$
\left[T \wedge \frac{i}{2} d Z_{\hat{x}}^{\epsilon} \wedge d \overline{Z_{\hat{x}}^{\epsilon}}\right]\left(B_{x}\left(r_{p}\right)\right) \geq e^{-l_{p}\left(\log d+2 \lambda_{1}+3 \epsilon\right)}=r_{p}^{\left(\log d+2 \lambda_{1}+3 \epsilon\right) /\left(\lambda_{2}-4 \epsilon\right)} .
$$

Since $\left(r_{p}\right)_{p}$ tends to 0 and $\hat{\nu}\left(\hat{\Omega}_{\epsilon}\right)>0$, we get

$$
\underline{d_{T, Z}}(\nu) \leq \frac{\log d+2 \lambda_{1}+3 \epsilon}{\lambda_{2}-4 \epsilon}=: \frac{\log d}{\lambda_{2}}+2 \frac{\lambda_{1}}{\lambda_{2}}+O_{3}(\epsilon) .
$$

One can prove

$$
\underline{d_{T, W}}(\nu) \leq \frac{\log d+2 \lambda_{2}+3 \epsilon}{\lambda_{2}-4 \epsilon}=: \frac{\log d}{\lambda_{2}}+2+O_{4}(\epsilon)
$$

in a similar way.

### 7.3 Monge-Ampère mass

We prove Proposition 7.1. Let $x \in \operatorname{Supp} \mu$, let $V$ be a neighbourhood of $x$ and let $Z: V \rightarrow \mathbb{C}$ be a holomorphic coordinate (submersion) on $V$. We want to prove that the positive measure $T \wedge \frac{i}{2} d Z \wedge d \bar{Z}$ is not the zero measure on $V$. With no loss of generality, we can assume that $x=(0,0), V=\mathbb{D}(2) \times \mathbb{D}(2)$ and $Z(z, w)=z$. Let also $T=2 i \partial \bar{\partial} G$ on $V$, where $G$ is a continuous psh function. We denote $\sigma_{z}(u):=(z, u)$.

Lemma 7.4. If $\left(T \wedge \frac{i}{2} d Z \wedge d \bar{Z}\right)(\mathbb{D}(2) \times \mathbb{D}(2))=0$, then $G \circ \sigma_{z}$ is harmonic on $\mathbb{D}$ for every $z \in \mathbb{D}$.

Proof. Let $z_{0} \in \mathbb{D}$ and let $\varphi \in C_{0}^{\infty}(\mathbb{D})$ be a test function. Let $\psi \in C_{0}^{\infty}\left(\mathbb{D}^{2}\right)$ such that $\psi \circ \sigma_{z_{0}}=\varphi$ on $\mathbb{D}$. According to Proposition 9.4, we have

$$
\left(T \wedge \frac{i}{2} d Z \wedge d \bar{Z}\right)(\psi)=\int_{z \in \mathbb{D}}\left(\int_{w \in \mathbb{D}}\left(G \circ \sigma_{z}\right)(w) \times \Delta\left(\psi \circ \sigma_{z}\right)(w) \mathrm{d} \operatorname{Leb}(w)\right) \mathrm{d} \operatorname{Leb}(z)
$$

which is equal to zero by our assumption. Since the measurable function

$$
z \mapsto \int_{w \in \mathbb{D}}\left(G \circ \sigma_{z}\right)(w) \times \Delta\left(\psi \circ \sigma_{z}\right)(w) \mathrm{d} \operatorname{Leb}(w)
$$

is non negative, there exists $A \subset \mathbb{D}$ such that $\operatorname{Leb}(A)=\operatorname{Leb}(\mathbb{D})$ and

$$
\begin{equation*}
\forall z \in A, \quad \int_{w \in \mathbb{D}}\left(G \circ \sigma_{z}\right)(w) \times \Delta\left(\psi \circ \sigma_{z}\right)(w) \mathrm{d} \operatorname{Leb}(w)=0 \tag{7.5}
\end{equation*}
$$

Let us extend this property to every $z \in \mathbb{D}$. Since $A$ is a dense subset of $\mathbb{D}$, there exists a sequence $\left(z_{n}\right)_{n}$ of points in $A$ which converges to $z$. Using (7.5), we get

$$
\begin{equation*}
\forall n \geq 1, \quad \int_{w \in \mathbb{D}}\left(G \circ \sigma_{z_{n}}\right)(w) \times \Delta\left(\psi \circ \sigma_{z_{n}}\right)(w) \mathrm{d} \operatorname{Leb}(w)=0 \tag{7.6}
\end{equation*}
$$

Since $G$ is continuous on $\overline{\mathbb{D}^{2}}$ and $\psi$ is smooth on $\overline{\mathbb{D}^{2}}, G$ and $2 i \partial \bar{\partial} \psi$ are uniformly continuous on $\overline{\mathbb{D}^{2}}$. This implies that $G \circ \sigma_{z_{n}}$ uniformly converges to $G \circ \sigma_{z}$ on $\mathbb{D}$ and that $\Delta\left(\psi \circ \sigma_{z_{n}}\right)$ uniformly converges to $\Delta\left(\psi \circ \sigma_{z}\right)$ on $\mathbb{D}$. Taking the limits in (7.6), we get

$$
\forall z \in \mathbb{D}, \quad \int_{w \in \mathbb{D}}\left(G \circ \sigma_{z}\right)(w) \times \Delta\left(\psi \circ \sigma_{z}\right)(w) \mathrm{d} \operatorname{Leb}(w)=0
$$

In particular, we obtain using $\psi \circ \sigma_{z_{0}}=\varphi$ :

$$
\int_{w \in \mathbb{D}}\left(G \circ \sigma_{z_{0}}\right)(w) \times \Delta \varphi(w) \mathrm{d} \operatorname{Leb}(w)=0
$$

This holds for every $\varphi \in C_{0}^{\infty}(\mathbb{D})$, hence the function $G \circ \sigma_{z_{0}}$ is harmonic on $\mathbb{D}$.
Now we use the following result, see [5, Lemme IV.1.1] and [24, Section A.10].
Theorem 7.5 (Briend). Let $G$ be a continuous psh function on $\mathbb{D}(2) \times \mathbb{D}(2)$. Let $E$ be the set of points $p \in \mathbb{D}\left(\frac{1}{4}\right) \times \mathbb{D}\left(\frac{1}{4}\right)$ such that there exists a holomorphic disc $\sigma_{p}: \mathbb{D} \rightarrow$ $\mathbb{D}(2) \times \mathbb{D}(2)$ satisfying

1. the boundary of $\sigma_{p}$ is outside $\mathbb{D}\left(\frac{1}{2}\right) \times \mathbb{D}\left(\frac{1}{2}\right)$,
2. $G \circ \sigma_{p}$ is harmonic $\mathbb{D}$.

Then $(2 i \partial \bar{\partial} G \wedge 2 i \partial \bar{\partial} G)(E)=0$.
In our situation, $\mathbb{D}\left(\frac{1}{4}\right) \times \mathbb{D}\left(\frac{1}{4}\right)=E$ since one can take for $\sigma_{p}$ the discs $\sigma_{z}: \mathbb{D} \rightarrow$ $\mathbb{D} \times \mathbb{D}, u \mapsto(z, u)$. Indeed, the boundary of $\sigma_{z}$ is contained in $\{z\} \times \partial \mathbb{D}$ and $G \circ \sigma_{z}$ is harmonic on $\mathbb{D}$ according to Lemma 7.4. Theorem 7.5 then gives:

$$
\begin{equation*}
(2 i \partial \bar{\partial} G \wedge 2 i \partial \bar{\partial} G)\left(\mathbb{D}\left(\frac{1}{4}\right) \times \mathbb{D}\left(\frac{1}{4}\right)\right)=0 \tag{7.7}
\end{equation*}
$$

which contradicts $x=0 \in \operatorname{Supp} \mu=\operatorname{Supp}(2 i \partial \bar{\partial} G \wedge 2 i \partial \bar{\partial} G)$.

## 8 Upper bound for the dimension of dilating measures

We prove Theorem 1.5. We shall take the proof of Theorem 5.2 and use Proposition 7.3. Let $\epsilon>0$ and let $\hat{\Lambda}_{\epsilon}$ and $\hat{\Delta}_{\epsilon}^{n}$ be the sets defined in Section 3.3. The set $\hat{\Omega}_{\epsilon}^{n}$ has been defined in Section 7.1, it satisfies $\hat{\nu}\left(\hat{\Omega}_{\epsilon}^{n}\right) \geq 1-2 \delta$. Hence we have $\hat{\nu}\left(\hat{\Lambda}_{\epsilon} \cap \hat{\Delta}_{\epsilon}^{n} \cap \hat{\Omega}_{\epsilon}^{n}\right) \geq 1-3 \delta$ for every $n \geq N_{\epsilon}$. Now let $K_{n}$ be the unique integer satisfying

$$
\begin{equation*}
\eta_{1} e^{-n \lambda_{1}-4 n \epsilon} e^{-\lambda_{2}+3 \epsilon} \leq \frac{1}{m_{0}} e^{-K_{n} \lambda_{2}+3 K_{n} \epsilon} \leq \eta_{1} e^{-n \lambda_{1}-4 n \epsilon} \tag{8.1}
\end{equation*}
$$

so that $K_{n} \simeq n \lambda_{1} / \lambda_{2}$. Let $\left\{x_{1}, \cdots, x_{N_{n, 2}}\right\}$ be a $\left(n, \eta_{1} / 4\right)$-separated subset of $A_{n}:=$ $\pi_{0}\left(\hat{\Lambda}_{\epsilon} \cap \hat{\Delta}_{\epsilon}^{n} \cap \hat{\Omega}_{\epsilon}^{K_{n}}\right)$ provided by Lemma 4.2. We have for every $n \geq \max \left\{N_{\epsilon}, n_{1-3 \delta}\right\}:$

$$
\begin{equation*}
N_{n, 2} \geq \nu\left(A_{n}\right) e^{n h_{\nu}-2 n \epsilon} \geq(1-3 \delta) e^{n h_{\nu}-2 n \epsilon} \tag{8.2}
\end{equation*}
$$

Then, for every $x_{i}$, let $\left\{y_{1}^{i}, \ldots, y_{M_{n}}^{i}\right\}$ be a $2 e^{-n \lambda_{1}-4 n \epsilon}$-separated subset of $B_{n}\left(x_{i}, \eta_{1} / 4\right) \cap$ $A_{n}$ provided by Lemma 4.3. The cardinality of this set satisfies:

$$
\begin{equation*}
M_{n} \geq e^{-n h_{\nu}-2 n \epsilon}\left(\frac{1}{2 \eta_{1}} e^{n \lambda_{1}+4 n \epsilon}\right)^{\underline{d_{\nu}-\epsilon}} \tag{8.3}
\end{equation*}
$$

For every $j \in\left\{1, \ldots, M_{n}\right\}$, we set $\hat{y}_{j}^{i} \in \hat{\Lambda}_{\epsilon} \cap \hat{\Delta}_{\epsilon}^{n} \cap \hat{\Omega}_{\epsilon}^{K_{n}}$ such that $y_{j}^{i}=\pi_{0}\left(\hat{y}_{j}^{i}\right)$. Then we follow the proof of Theorem 5.2 until the inequality (5.8):

$$
\begin{equation*}
d^{n} \geq \sum_{i=1}^{N_{n, 2}} \sum_{j=1}^{M_{n}} e^{2 n \lambda_{1}-4 n \epsilon}\left[T \wedge\left(\frac{i}{2} d Z_{\hat{y}_{j}^{i}}^{\epsilon} \wedge d \overline{Z_{\hat{y}_{j}^{i}}^{\epsilon}}\right)\right]\left(B_{y_{j}^{i}}\left(\eta_{1} e^{-n \lambda_{1}-4 n \epsilon}\right)\right) \tag{8.4}
\end{equation*}
$$

We want to apply Proposition 7.3. According to (8.1),

$$
B_{y_{j}^{i}}\left(\eta_{1} e^{-n \lambda_{1}-4 n \epsilon}\right) \supset B_{y_{j}^{i}}\left(\frac{1}{m_{0}} e^{-K_{n} \lambda_{2}+3 K_{n} \epsilon}\right)
$$

We apply the positive measure $T \wedge\left(\frac{i}{2} d Z_{\hat{y}_{j}^{i}}^{\epsilon} \wedge d \overline{Z_{\hat{y}_{j}^{i}}^{\epsilon}}\right)$ to this inclusion. Since $\hat{y}_{j}^{i} \in \hat{\Omega}_{\epsilon}^{K_{n}}$, we deduce from Proposition 7.3 that for every $n$ satisfying $n \geq N_{\epsilon}$ and $K_{n} \geq N_{\epsilon}$ :

$$
\left[T \wedge\left(\frac{i}{2} d Z_{\hat{y}_{i, j}}^{\epsilon} \wedge d \overline{Z_{\hat{y}_{i, j}}^{\epsilon}}\right)\right]\left(B_{y_{j}^{i}}\left(\eta_{1} e^{-n \lambda_{1}-4 n \epsilon}\right)\right) \geq \frac{1}{d^{K_{n}}} e^{-2 K_{n} \lambda_{1}-2 K_{n} \epsilon} \frac{1}{q_{0}} .
$$

We infer from (8.4) that for every $n$ satisfying $n \geq N_{\epsilon}$ and $K_{n} \geq N_{\epsilon}$,

$$
\begin{equation*}
d^{n} \geq N_{n, 2} \cdot M_{n} \cdot e^{2 n \lambda_{1}-4 n \epsilon} \frac{1}{d^{K_{n}}} e^{-2 K_{n} \lambda_{1}-2 K_{n} \epsilon} \frac{1}{q_{0}} . \tag{8.5}
\end{equation*}
$$

Now we use the upper bounds for $N_{n, 2}$ and $M_{n}$ given by (8.2) and (8.3):

$$
d^{n+K_{n}} \geq(1-3 \delta) e^{n h_{\nu}-2 n \epsilon} \cdot e^{-n h_{\nu}-2 n \epsilon}\left(\frac{1}{2 \eta_{1}} e^{n \lambda_{1}+4 n \epsilon}\right)^{\frac{d_{\nu}-\epsilon}{}} \cdot e^{2 n \lambda_{1}-4 n \epsilon} e^{-2 K_{n} \lambda_{1}-2 K_{n} \epsilon} \frac{1}{q_{0}} .
$$

If $C_{1}(\epsilon):=(1-3 \delta) / q_{0}\left(2 \eta_{1}\right)^{\underline{d_{\nu}}-\epsilon}$, we get:

$$
\log d+\frac{K_{n}}{n} \log d \geq \frac{1}{n} \log C_{1}(\epsilon)-8 \epsilon+\left(\lambda_{1}+4 \epsilon\right)\left(\underline{d_{\nu}}-\epsilon\right)+2 \lambda_{1}-2 \frac{K_{n}}{n}\left(\lambda_{1}+\epsilon\right) .
$$

By using (8.1), we have

$$
\log d+\frac{\lambda_{1}+4 \epsilon}{\lambda_{2}-3 \epsilon} \log d \geq \frac{1}{n} \log C_{2}(\epsilon)-8 \epsilon+\left(\lambda_{1}+4 \epsilon\right)\left(\underline{d_{\nu}}-\epsilon\right)+2 \lambda_{1}-2 \frac{\lambda_{1}+4 \epsilon}{\lambda_{2}-3 \epsilon}\left(\lambda_{1}+\epsilon\right)
$$

where $C_{2}(\epsilon)$ is another constant. Letting $n$ tend to $+\infty$ and then $\epsilon$ to 0 , we get

$$
\underline{d_{\nu}} \leq \frac{\log d}{\lambda_{1}}+\frac{\log d}{\lambda_{2}}+2\left(\frac{\lambda_{1}}{\lambda_{2}}-1\right) .
$$

To obtain the other upper bound, we use the analogue of (8.4) for $W$. Applying Proposition 7.3 with respect to $W$, we obtain instead of (8.5):

$$
d^{n} \geq N_{n, 2} \cdot M_{n} \cdot e^{2 n \lambda_{2}-4 n \epsilon} \frac{1}{d^{K_{n}}} e^{-2 K_{n} \lambda_{2}-2 K_{n} \epsilon} \frac{1}{q_{0}} .
$$

Then we get

$$
\underline{d_{\nu}} \leq \frac{\log d}{\lambda_{1}}+\frac{\log d}{\lambda_{2}}+2\left(1-\frac{\lambda_{2}}{\lambda_{1}}\right),
$$

which completes the proof of Theorem 1.5.

## 9 Appendix

### 9.1 Dimension of measures

Proposition 9.1. Let $\nu_{1}$ and $\nu_{2}$ be two probability measures on $\mathbb{P}^{2}$ such that $\nu_{1} \ll \nu_{2}$. Then for $\nu_{1}$-almost every $x \in \mathbb{P}^{2}$, we have:

$$
\underline{d_{\nu_{1}}}(x) \geq \underline{d_{\nu_{2}}}(x) \quad \text { and } \quad \overline{d_{\nu_{1}}}(x) \geq \overline{d_{\nu_{2}}}(x) .
$$

Proof. Let $\varphi \in L^{1}\left(\nu_{2}\right)$ such that $\nu_{1}(A)=\int_{A} \varphi \nu_{2}$ for every Borel set $A$ of $\mathbb{P}^{2}$. Using the dominated convergence Theorem,

$$
\lim _{M \rightarrow \infty} \int_{\mathbb{P}^{2}} 1_{\{\varphi \leq M\}} \varphi \mathrm{d} \nu_{2}=\int_{\mathbb{P}^{2}} \varphi \mathrm{~d} \nu_{2}=1
$$

For every $n \geq 1$, we let $M_{n}$ satisfy $\int_{\mathbb{P}^{2}} 1_{\left\{\varphi \leq M_{n}\right\}} \varphi d \nu_{2} \geq 1-\frac{1}{n}$. By the Lebesgue density Theorem, for $\nu_{1}$-almost every $x$ in $\left\{\varphi \leq M_{n}\right\}$, we have

$$
\lim _{r \rightarrow 0} \frac{\nu_{1}\left(B_{x}(r) \cap\left\{\varphi \leq M_{n}\right\}\right)}{\nu_{1}\left(B_{x}(r)\right)}=1
$$

Then for every $r$ small enough, we have

$$
\frac{1}{2} \nu_{1}\left(B_{x}(r)\right) \leq \nu_{1}\left(B_{x}(r) \cap\left\{\varphi \leq M_{n}\right\}\right)=\int_{B_{x}(r) \cap\left\{\varphi \leq M_{n}\right\}} \varphi d \nu_{2} \leq M_{n} \int_{B_{x}(r)} d \nu_{2}
$$

And thus $\nu_{1}\left(B_{x}(r)\right) \leq 2 M_{n} \nu_{2}\left(B_{x}(r)\right)$. We deduce that

$$
\underline{d_{\nu_{1}}}(x) \geq \underline{d_{\nu_{2}}}(x) \text { and } \overline{d_{\nu_{1}}}(x) \geq \overline{d_{\nu_{2}}}(x)
$$

for $\nu_{1}$-almost every $x \in\left\{\varphi \leq M_{n}\right\}$. We end with $\nu_{1}\left(\cup_{n \in \mathbb{N}}\left\{\varphi \leq M_{n}\right\}\right)=1$.
Now we take the notations of Section 1.1.
Proposition 9.2. Let $S$ be a $(1,1)$-closed positive current on $\mathbb{P}^{2}$. Let $x \in \mathbb{P}^{2}$ and let $(Z, W)$ be holomorphic coordinates near $x$. Then

$$
\underline{d_{S}}(x)=\min \left\{\underline{d_{S, Z}}(x), \underline{d_{S, W}}(x)\right\}, \quad \overline{d_{S}}(x)=\min \left\{\overline{d_{S, Z}}(x), \overline{d_{T, W}}(x)\right\} .
$$

Proof. Let us set $\sigma_{S, Z}=S \wedge\left(\frac{i}{2} d Z \wedge d \bar{Z}\right)$ and $\sigma_{S, W}=S \wedge\left(\frac{i}{2} d W \wedge d \bar{W}\right)$. There exists $c>0$ such that $\frac{1}{c}\left(\sigma_{S, Z}+\sigma_{S, W}\right) \leq \sigma_{S} \leq c\left(\sigma_{S, Z}+\sigma_{S, W}\right)$ on a neighbourhood of $x$, see [13] Chapter III, §3. We deduce for every $r$ small enough

$$
\frac{1}{c} \max \left[\sigma_{S, Z}\left(B_{x}(r)\right), \sigma_{S, W}\left(B_{x}(r)\right)\right] \leq \sigma_{S}\left(B_{x}(r)\right) \leq 2 c \max \left[\sigma_{S, Z}\left(B_{x}(r)\right), \sigma_{S, W}\left(B_{x}(r)\right)\right]
$$

We finish by observing that the local dimension of the maximum of two measures is equal to the minimum of these two dimensions, since one divides by $\log r$ which is negative.

### 9.2 Cohomology and slices

We refer to Sections 1.2 and A. 3 of Dinh-Sibony's book [16].
Proposition 9.3. Let $S$ be a $(1,1)$-closed positive current of $\mathbb{P}^{2}$ of mass 1. Let $\omega$ be the Fubini-Study form on $\mathbb{P}^{2}$ and let $f: \mathbb{P}^{2} \rightarrow \mathbb{P}^{2}$ be an endomorphism of degree $d$. Then,

$$
\int_{\mathbb{P}^{2}}\left(f^{n}\right)_{*} S \wedge \omega=\int_{\mathbb{P}^{2}} S \wedge\left(f^{n}\right)^{*} \omega=d^{n}
$$

Proof. The first equality comes from the definition of duality. We show the second one. By using $f^{*} \omega=d \cdot \omega+2 i \partial \bar{\partial} u$, where $u$ is a smooth function on $\mathbb{P}^{2}$, we obtain by induction

$$
\left(f^{n}\right)^{*} \omega=d^{n} \omega+2 i \partial \bar{\partial} v_{n}
$$

where $v_{n}:=\left(d^{n-1} \cdot u+\cdots+d \cdot u \circ f^{n-2}+u \circ f^{n-1}\right)$. Hence

$$
\int_{\mathbb{P}^{2}} S \wedge\left(f^{n}\right)^{*} \omega=\int_{\mathbb{P}^{2}} S \wedge\left(d^{n} \omega+2 i \partial \bar{\partial} v_{n}\right) .
$$

Since $S$ is a closed current of mass 1 , we have $\int_{\mathbb{P}^{2}} S \wedge 2 i \partial \bar{\partial} v_{n}=0$ and $\int_{\mathbb{P}^{2}} S \wedge d^{n} \omega=d^{n}$.
Proposition 9.4. Let $G$ be a continuous psh function on $\mathbb{D}^{2}$ and let $S=2 i \partial \bar{\partial} G$. Let $(Z, W)$ be the coordinates on $\mathbb{D}^{2}$ and let $\phi \in C_{0}^{\infty}\left(\mathbb{D}^{2}\right)$. Then
$S \wedge \frac{i}{2} d Z \wedge d \bar{Z}(\phi)=\int_{z \in \mathbb{D}}\left(\int_{w \in \mathbb{D}} G_{z}(w) \times \Delta \phi_{z}(w) d \operatorname{Leb}(w)\right) d \operatorname{Leb}(z)=\int_{z \in \mathbb{D}}\left(\sigma_{z}^{*} S\right)\left(\phi_{z}\right) d \operatorname{Leb}(z)$, where $\sigma_{z}: u \mapsto(z, u), G_{z}:=G \circ \sigma_{z}$ and $\phi_{z}:=\phi \circ \sigma_{z}$.

Proof. By definition,

$$
S \wedge \frac{i}{2} d Z \wedge d \bar{Z}(\phi)=2 i \partial \bar{\partial} G\left(\phi \frac{i}{2} d Z \wedge d \bar{Z}\right)=\int_{\mathbb{D}^{2}} G .2 i \partial \bar{\partial}\left(\phi \times \frac{i}{2} d Z \wedge d \bar{Z}\right) .
$$

The computation

$$
2 i \partial \bar{\partial}\left(\phi \times \frac{i}{2} d Z \wedge d \bar{Z}\right)=4\left(\frac{\partial^{2} \phi}{\partial w \partial \bar{w}}\right) \frac{i}{2} d W \wedge d \bar{W} \wedge \frac{i}{2} d Z \wedge d \bar{Z}=4\left(\frac{\partial^{2} \phi}{\partial w \partial \bar{w}}\right) \mathrm{d} \operatorname{Leb}(z, w)
$$

allows to write

$$
\begin{aligned}
S \wedge \frac{i}{2} d Z \wedge d \bar{Z}(\phi) & =\int_{(z, w) \in \mathbb{D}^{2}} G(z, w) \times 4 \frac{\partial^{2} \phi}{\partial w \partial \bar{w}}(z, w) \mathrm{d} \operatorname{Leb}(z, w) \\
& =\int_{z \in \mathbb{D}}\left(\int_{w \in \mathbb{D}} G_{z}(w) \times \Delta \phi_{z}(w) \mathrm{d} \operatorname{Leb}(w)\right) \mathrm{d} \operatorname{Leb}(z) .
\end{aligned}
$$

Finally, the quantity in brackets is equal to $\left(\Delta G_{z}\right)\left(\phi_{z}\right)=\left(\sigma_{z}^{*} S\right)\left(\phi_{z}\right)$.

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