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Passive heaving of elliptical cylinders with active pitching – From cylinders towards flapping foils

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A B S T R A C T

This paper presents a study of the flow past elastically mounted cylinders with prescribed rotational oscillation about the cylinder centre, which are free to heave, or oscillate transverse to the flow. The configuration serves as an idealized model of a flapping-foil energy harvester. A range of geometries are tested, from the circular cylinder with an aspect ratio of 1.0 to elliptical cylinders up to aspect ratio of 6.0 approaching a flat plate. The driving frequency of the rotational oscillation is varied, while the amplitude of rotation is fixed at $\pi/2$, meaning both axes of the geometries present fully to the oncoming flow each cycle. The Reynolds number is 200. The natural frequency of the elastic-mounting is set to the Strouhal frequency for a circular cylinder. The ratio of the mass of the cylinder to the mass of the equivalent volume of displaced fluid is set to 5.0. Configurations with zero-damping reveal a rich parameter space, with increasing cross-stream oscillation with increasing geometry aspect ratio. Driving frequencies for peak oscillation amplitude are grouped around a driving frequency of 0.9 times the natural frequency of the elastic structure. The variation of the power input to actuate the rotational oscillation of the cylinder is also analysed. The fluid structure interaction is analysed for energy harvesting potential; power output is modelled by linear damping on the heave. Increasing the damping on the structure leads to optimal values of driving frequency and damping for each aspect ratio tested. For each aspect ratio, comparisons are drawn and similarities found between these optimal cases for power output and the undamped cases for maximum oscillation amplitude and velocity. The study of the parameter space serves as a useful starting point for further study of the many parameters affecting the performance of flapping-foil energy harvesting.

1. Introduction

The need to limit the growth of greenhouse gas emissions has driven the search for alternative zero- or low-emission energy sources. The feasibility of extracting energy from a flow using a pitching and heaving aerofoil was explored in [Kinsey and Dumas \(2008\)](#). Rotational and cross-stream oscillation was prescribed on a NACA 0015 aerofoil in a uniform flow, varying the pitching frequency and amplitude, while keeping the cross-stream oscillation at one chord length and the pitching axis at one third of the

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chord length from the leading edge. A mechanism to extract energy was not considered, but rather the work done by the fluid on the body. When this work is positive, power extraction is possible. For their configuration, they identified optimal rotational frequencies (normalized by chord length and flow speed) in the range $0.12 \leq f \leq 0.18$ and amplitudes in the range $70^\circ \leq A_\theta \leq 80^\circ$, where A_θ is the rotation amplitude. A key finding of the study was that the motion parameters, such as heaving amplitude and rotation frequency and amplitude, were the critical parameters for performance, and that geometry (aspect ratios of $\Gamma = 5.0, 6.67$ and 50 were tested, where Γ is the ratio of chord length to thickness) and viscous parameters (Reynolds number was varied, as well as application of turbulence modelling) were of secondary importance. These are important findings as they support the generalization of lower-Reynolds number dynamics and the convergence of behaviour with aspect ratio, albeit only for aspect ratio $\Gamma > 5.0$.

Useful reviews of the subject can be found in [Xiao and Zhu \(2014\)](#) and [Young et al. \(2014\)](#). [Xiao and Zhu \(2014\)](#) suggest that more work be done on semi-activated and self-sustaining systems, as well as on three-dimensional configuration, specifically the effect of spanwise aspect ratio. [Young et al. \(2014\)](#) suggest work on self-sustaining systems and on foil flexibility; indeed many studies have focussed on the possibility of adding flexibility or a time-dependent deformation to the aerofoil, with the intent of increasing power extraction efficiency. [Hoke et al. \(2015\)](#) examined extraction and propulsion with an aerofoil with a camber that deforms throughout the oscillation. Excluding the energy required to flex the aerofoil, this lead to an increase in efficiency. They also suggested the possibility of having the aerofoil flex passively. Using an immersed boundary-lattice Boltzmann method, [Wu et al. \(2015 a,b\)](#) investigated energy harvesting from a flapping aerofoil with either a rigid or flexible plate attached to the trailing edge, with prescribed pitching and heaving and with prescribed pitching and flow-induced heaving, finding that whether a rigid or flexible tail was used could significantly affect the Strouhal number and efficiency of the optimal configuration. [Tian et al. \(2014\)](#) investigated the effect of active and passive leading and trailing edge deformation of a flapping-plate flow energy harvester, finding that plate flexibility does not significantly affect the power-extraction capability of the plate, but active control on the leading edge of the plate can.

[Zhu \(2011\)](#) examined optimal frequencies for flapping foil energy harvesting ($f \approx 0.15$, normalized by chord length and flow speed), with prescribed pitching and heaving motions. They argue that a stable wake is better for energy harvesting and that a fully passive system is feasible, due to the near zero average power needed to actuate the foil rotation under optimal conditions. [Deng et al. \(2015\)](#) studied inertial effects on a semi-passive energy harvester with prescribed pitching and free heaving. In contrast to many other studies on the problem, power extraction was modelled using linear damping on the heave, which was included in the equation of motion. They found that efficiency decreases monotonically with mass ratio, but that inertial effects are negligible for mass ratios of less than 10. They qualify their results by noting that changes in behaviour with variation of pitching axis location, three-dimensionality and rotational moment of inertia are likely or unknown. The work of [Deng et al. \(2015\)](#) is notable in that the power extraction, via linear damping, is calculated and its effect included in the body motion dynamics; they also present their results in terms of efficiency and in terms of power output.

The current study examines primarily the effect of aspect ratio on the power output of a passively heaving cylinder with active pitching. In the works described above, generally, the efficiency of energy harvesting is described as:

$$\eta = \frac{\frac{1}{T} \int_{t=t_0}^{t=t_0+T} (P_{\text{out}} - P_{\text{in}}) dt}{\frac{1}{2} \rho U^3 b d_y}, \quad (1)$$

where P_{out} is the harvested power, P_{in} is the input power to actuate the rotation, and the term $\frac{1}{2} \rho U^3 b d_y$ represents the available energy in the flow, with ρ being the fluid density, U the freestream velocity, b the spanwise length of the body and d_y the cross-stream extent of the body's path. This is an efficiency equation that relates the performance of the energy harvester to that of more traditional wind turbines; the Betz efficiency limit of 59% is often cited ([Betz, 1966](#)), giving the limit of how much energy can be extracted from a given flow stream. Problematic to this formulation of efficiency is that some flapping foil energy harvesters will rely on large heave to extract energy. A large amplitude oscillation is an indicator of greater energy transfer from the fluid to the structure and of a greater value of path extent, b_y . This positive energy transfer, or P_{out} , also counteracts to reduce efficiency by increasing the term b_y . This measure of efficiency is important in contexts where arrays of such devices are being considered, or where a limit exists on, for example, the cross-stream space available. However, for the current study, we focus on the dynamics of a single, idealized oscillating system. Having an efficiency formulation for a system that relies on heave to generate power, but which penalizes for that same heave, can be misleading in contexts where the aim is to maximize the power output. However, in future work on arrays of oscillating cylinders or with limited cross-stream domain, the efficiency formulation shown in Eq. (1) would need to be considered. For the present work, to non-dimensionalize power, we effectively replace the cross-stream extent parameter, d_y , with, D , the major axis of the elliptical cross-section, synonymous with the chord length.

Also problematic is the treatment of the power required to drive the rotation of the body, P_{in} . The required instantaneous input power can be calculated as $P_{\text{in}} = M\dot{\theta}$, where M is the torque exerted on the body by the fluid and $\dot{\theta}$ is the rotational velocity of the body. For flapping foil energy harvesting, the time-averaged mean of P_{in} , \bar{P}_{in} , generally returns a positive value, indicating a required net input of energy to drive the rotational oscillation. However, the time-dependent term P_{in} , usually switches signs throughout the cycle, indicating a fluid torque at times pushing in the same direction as the body rotation and at others pushing against it. The term \bar{P}_{in} then represents the sum of the energies required to accelerate the body rotation (when $P_{\text{in}} > 0$) and the energy absorbed by the structure (when $P_{\text{in}} < 0$). [Section 3.3](#) deals in more detail with the calculation of the mean input power, \bar{P}_{in} .

Across the studies of flapping foil energy harvesting, a common thread is the large range of parameters, assumptions and idealizations used to analyse the problem. The current study is no different, but does have the aim of examining a more fundamental

configuration and the possibility of linking the problem with other studies, such as those of bluff-bodies and vortex-induced vibration. With that aim, the current study investigates a flapping foil energy harvester using simpler geometry, namely the circular and elliptical cylinders. The heave is flow-induced, described by an equation of motion that includes a restoring force and a damping/energy extraction force. Active control is imposed by rotation around the centre of mass of each body (at the half chord point). The rotation is of amplitude $A_\theta = \pi/2$. The aspect ratio of the cylinder is varied from the circular cylinder of aspect ratio $\Gamma = 1.0$ to the elliptical cylinder of $\Gamma = 6.0$. For the circular cylinder, the pitching contributes only viscous shear stress forcing, with the local velocity resulting from the pitching on the cylinder surface always aligned tangentially to the surface. As aspect ratio is increased beyond 1, the component of normal stress due to the pressure resulting from the pitching increases. Therefore, at higher aspect ratios the flow can potentially resist the pitching motion more forcefully.

This study seeks to examine the body and fluid dynamics of optimal configurations for power output, and examine how they vary across different cross-sectional body geometries, with the aim of presenting useful base cases and conclusions for investigation of further effects and parameter variations, such as geometry and rotation kinematics.

2. Problem definition and method

The study investigates flows for rotationally oscillating geometries ranging from the circular cylinder to the flat plate, where energy is extracted via a damper on the passive heaving motion. The geometries tested are shown in Fig. 1; namely, five ellipses of aspect ratio, $\Gamma = \frac{\text{major axis}}{\text{minor axis}} = 1.0, 1.5, 2.0, 4.0$ and 6.0 . The length scale for the system is the major axis of each elliptical cross-stream, D , while the velocity scale is the velocity of the fluid in the freestream, U .

The cylinder is elastically mounted in the cross-stream, or heave, direction. The harmonic oscillator equation governing the heaving motion is:

$$M(\ddot{Y} + (c/M)\dot{Y} + (2\pi f_N)^2 Y) = F_y, \quad (2)$$

where \ddot{Y} , \dot{Y} and Y are respectively the acceleration, velocity and displacement of the cylinder perpendicular to the freestream, F_y is the lift force on the cylinder, M is the mass of the cylinder, m^* is the ratio of the mass of the cylinder to the mass of the equivalent volume of fluid, c is the damping coefficient and $f_N = \sqrt{k/M}/(2\pi)$ is the natural structural frequency of the undamped system, where k is the spring stiffness. The natural structural frequency is chosen to be close to the vortex shedding frequency for a stationary circular cylinder, that is, $f_N^* = 0.20$; note that for the remainder of the paper, frequencies described with the superscript $*$ denote a frequency non-dimensionalized using D and U . The vortex shedding frequency of a stationary cylinder will vary with aspect ratio and with the alignment of the cylinder axes with the flow; this variation for the stationary cylinder is explored in the results.

Active control is implemented via a pitching motion, with $\theta = A_\theta \sin(2\pi f_D t)$, where θ is the angular displacement, f_D is the driving frequency and A_θ is the pitching amplitude. An amplitude of $A_\theta = \pi/2$ is chosen. The driving frequency of the pitching motion is varied in the range $0.00 \leq f_D^* \leq 0.40$. Optimal power output is explored by also varying the damping of the system, c . A constant ratio of body density to fluid density is chosen as, $m^* = \rho_{\text{body}}/\rho = 5.0$; this mass ratio is within the range suggested by Deng et al. (2015) in which inertial effects are negligible in terms of power extraction. Although Deng et al. (2015) dealt with a different geometry and parameter space to the current study, we use their finding as a guide, leaving open the possibility that our results may be sensitive to variation of mass ratio. The Reynolds number is defined as $Re = UD/\nu$, where ν is the kinematic viscosity; bearing in mind the criticality of the kinematic behaviour of the cylinder explored in Kinsey and Dumas (2008), the Reynolds number was fixed at $Re=200$ for all simulations.

2.1. Numerical method

Simulations are carried out using a sharp-interface immersed boundary method to represent the moving circular and elliptical bodies on an underlying Cartesian grid, on which the flow is found by solving the incompressible Navier–Stokes equations using a second-order finite-difference scheme (Mittal et al., 2008). The pitching motion of the body is controlled, while the heaving motion is found by solving the linear oscillator Eq. (2).

Time integration is performed via a two-way time-splitting scheme. First the advection and diffusion terms of the discretized Navier–Stokes equations are integrated to an intermediate velocity field. The advection term is integrated using a second-order

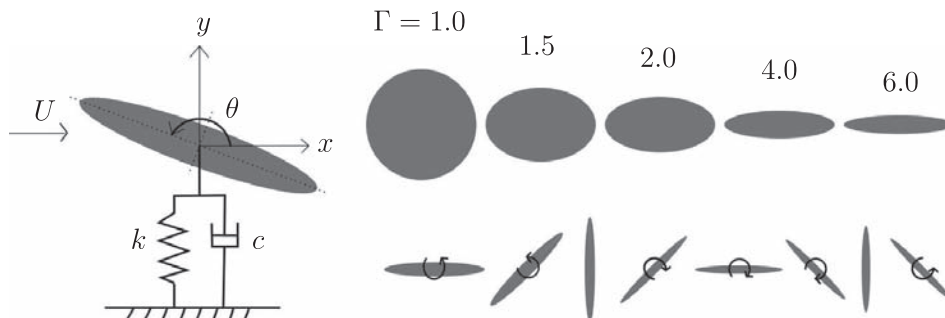


Fig. 1. A sketch of the geometries and problem setup.

Adams–Bashforth scheme, while the diffusion term is integrated using an implicit Crank–Nicolson scheme, hence the resulting equation is implicit and needs to be solved iteratively, which is done using a Gauss–Jacobi scheme.

Second, the pressure term of the Navier–Stokes equations needs to be integrated from the intermediate velocity field to obtain the velocity field at the end of the timestep. Continuity is enforced for the final velocity field, resulting in a Poisson equation for the pressure correction (the difference between the pressure at the previous timestep and the current timestep), which is solved using a custom geometric multigrid scheme. The resulting pressure field is then used to complete the integration to the velocity field at the end of the timestep.

The body surface is represented by linear finite elements – in a two-dimensional flow, this surface reduces to a curve and the finite elements reduce to linear line segments. The boundary conditions on the body surface are enforced via interpolation from the flow field and the boundary condition to the set of “ghost” points – those points on the underlying Cartesian grid which are inside the body, but which have neighbouring points outside the body. This creates a custom stencil for each ghost point, which is defined at run time and varies, as does the set of ghost points, throughout the history of the simulation as the body moves.

The forces on the body are obtained by interpolating the pressure and viscous stresses from the flow field onto the nodes defining the finite elements of the body surface, and then integrating these stresses in a finite element sense. These forces are fed to the equation of motion of the body, which is integrated forward in time along with the Navier–Stokes equations using a weakly coupled Newmark- β scheme (Newmark, 1959).

Boundary conditions for the velocity field are a Dirichlet condition on the upstream and lateral boundaries specifying the freestream condition, and a Neumann condition at the downstream boundary that all normal gradients are zero. On the immersed body boundary, the flow velocity is set to be equal to the motion of the body surface. For the pressure, a Neumann condition of the normal gradient equal to zero is prescribed at the upstream and lateral boundaries, and a Dirichlet condition of zero pressure specified at the downstream boundary. On the immersed body boundary, a Neumann condition of the pressure normal gradient equal to zero is specified.

The underlying Cartesian grid was designed with a regular grid spacing around the cylinder covering a region up to $3D$ in each direction from the origin. Grid spacings were then increased linearly to the domain boundaries. The inlet and lateral domain boundaries were located $15D$ from the origin, while the outlet boundary was located $40D$ downstream. Twice as many grid points were used in the streamwise direction than in the transverse direction; the number of grid points in each direction is restricted to the powers of 2, thereby leading to a grid size of $2^{N+1} \times 2^N$, where N is a whole number. We denote each grid as GN , such that $G10$ corresponds to a grid size of 2048×1024 .

For aspect ratio, $\Gamma = 1.0$, the circular cylinder, the results obtained using the immersed boundary method across the driving frequency range and for zero structural damping were validated against results for the same problem using a spectral-element method. The software for the spectral-element method has been tested extensively in other oscillating bluff body flows and flow-induced vibration studies (Nazarinia et al., 2009; Lo Jacono et al., 2010; Leontini et al., 2006; Griffith et al., 2011). The results for maximum non-dimensionalized oscillation amplitude, A_{\max}^* , returned by the immersed boundary method on grid $G9$ showed negligible differences with those returned by the spectral element method.

For aspect ratios, $\Gamma = 2.0$ and 4.0 , results were tested for the resolution of the underlying Cartesian grid, for the cases of maximum amplitude oscillation with zero damping, corresponding to driving frequencies $f_D^* = 0.185$ and 0.180 respectively. There is no damping, which when applied reduces the maximum amplitude and velocity of the oscillation, leading to smaller velocity gradients in the fluid. Four grids, $G8$, $G9$, $G10$ and $G11$ were tested; simulations on each grid ran with timesteps $\Delta t U/D = 0.005$, 0.003125 , 0.002 and 0.001 , respectively. Table 1 presents the grid size in the near-cylinder region, Δx_{fine} , and the maximum amplitude achieved, A_{\max}^* , for each of the two geometries.

For aspect ratio $\Gamma = 2.0$, the results for maximum amplitude, A_{\max}^* , run with grid $G9$ return an error of less than 1% compared to the value returned by grid $G11$. For aspect ratio $\Gamma = 4.0$, the results for maximum amplitude, A_{\max}^* , run with grid $G10$ return an error of less than 1% compared to the value returned by grid $G11$.

For a case with damping, Table 2 presents the same grid resolution study, but for the case of optimal power output for aspect ratio $\Gamma = 4.0$ outlined later in Section 3.4. This case features damping and smaller amplitude oscillation due to energy extraction from the heave. It shows the convergence with grid resolution of the power output, with the value of \bar{P}_{out} on grid $G9$ within 3% of the value for the higher resolution grid $G11$. For all the simulations using the geometries of aspect ratios $\Gamma = 1.0$, 1.5 and 2.0 , grid $G9$ was used. For $\Gamma = 4.0$ and 6.0 , grid $G10$ was used, except for a limited range of the most high amplitude cases for aspect ratio $\Gamma = 6.0$ and $c = 0.0$. In these cases, the amplitude was large enough to require a larger area of constant grid resolution around the origin. Another mesh was constructed, also of grid size 2048×1024 , which pushed the area of constant grid spacing out to $4D$ from the

Table 1

Parameters and maximum amplitudes for the four immersed-boundary grids tested, for $Re=200$ and $c=0$, the two cases producing maximum amplitude oscillation for aspect ratios $\Gamma = 2.0$ and 4.0 , corresponding to driving frequencies $f_D^* = 0.185$ and 0.180 , respectively.

Grid	Δx_{fine}	$A_{\max}^* (\Gamma = 2.0)$	$A_{\max}^* (\Gamma = 4.0)$
G8	$D/32$	1.4452	2.0044
G9	$D/64$	1.5561	2.1155
G10	$D/128$	1.5484	2.1628
G11	$D/256$	1.5538	2.1722

Table 2

Parameters and results of the four immersed-boundary grids tested for the case of $Re=200$, $f_D^* = 0.175$, $\Gamma = 4.0$ and $cD/MU = 0.377$.

Grid	Δx_{fine}	$\Delta t U/D$	A_{max}^*	\bar{P}_{out}
G8	$D/32$	0.0020	1.1590	0.3033
G9	$D/64$	0.0040	1.2582	0.3542
G10	$D/128$	0.0020	1.3080	0.3827
G11	$D/256$	0.0010	1.3251	0.3940

cylinder, at the expense of some far-field resolution. The comparison of results obtained for lower amplitude oscillation cases run on the two grids used for aspect ratio $\Gamma = 6.0$ returned negligible differences in maximum oscillation amplitude.

3. Results

3.1. Flow past stationary elliptical cylinders

For the vibrating cylinder cases tested here, the structural frequency is set at $f_N^* = 0.20$, which is approximately equal to the vortex-shedding frequency, f_{St}^* for the flow past a stationary circular cylinder for Reynolds number, $Re=200$. However, this result refers only to the stationary aspect ratio, $\Gamma = 1.0$, geometry; this study includes rotationally oscillating elliptical cylinders of aspect ratios up to $\Gamma = 6.0$. For background to the study, as well as context to the selection of the structural frequency adopted, the flow past the five elliptical cylinders was simulated, over a range of fixed pitch angles, θ_0 , but with no translation or rotation permitted. Fig. 2(a) plots the measured vortex shedding frequency, for each case, over the range of pitch angle, $0^\circ \leq \theta_0 \leq 90^\circ$ and aspect ratio, $1.0 \leq \Gamma \leq 6.0$. Due to the symmetry of the ellipse, this range encompasses all possible angles of the ellipse to the flow, and, effectively, the range of aspect ratios $0.167 \leq \Gamma \leq 6.0$ – an ellipse of aspect ratio $\Gamma = 0.167$ and pitch angle $\theta_0 = 90^\circ$ is identical to an ellipse of $\Gamma = 6.0$ and $\theta_0 = 0^\circ$. The results show that for the circular cylinder of aspect ratio $\Gamma = 1.0$, the pitch angle of course has no effect on the shedding frequency. For aspect ratio $\Gamma = 1.5$ and 2.0, the shedding frequency increases for all pitching angles to as high as $f_{St}^* = 0.29$ for $\Gamma = 2.0$ and $\theta_0 = 0^\circ$. This case represents an ellipse with its major axis aligned in the streamwise direction, reducing the frontal area of the cylinder and narrowing the wake. The length scale of the vortex formation is reduced, leading to higher frequency shedding than is the case for $\Gamma = 1.0$. The effect increases with aspect ratio to the point that there is no vortex-shedding for $\Gamma \geq 4.0$ and $\theta_0 \leq 15^\circ$.

At the other end, shedding frequencies as low as $f_{St}^* = 0.16$ are observed for $\Gamma = 6.0$ and $\theta_0 = 90^\circ$. This represents the high aspect ratio ellipse oriented with its minor axis aligned with the streamwise direction, retaining the frontal area of the circular $\Gamma = 1.0$ cylinder. In the near-wake, a larger formation region is present, leading to larger-scale vortex formation and lower shedding frequencies. Fig. 2(b) investigates whether the critical parameter is the width of the geometry with respect to the oncoming flow, by presenting the same data plotted in Fig. 2(a), but with the shedding frequencies normalized by the frontal area of the tilted ellipses, D_θ , rather than the ellipse major axis, D . This normalization reduces the range of shedding frequencies significantly for the four geometries of aspect ratio, $\Gamma > 1.0$.

For the following results for flow-induced vibration from pitching ellipses, a constant structural frequency of $f_N^* = 0.20$ is selected, which is near the shedding frequency for the $Re=200$ flow past a stationary cylinder and is also near or within the range of

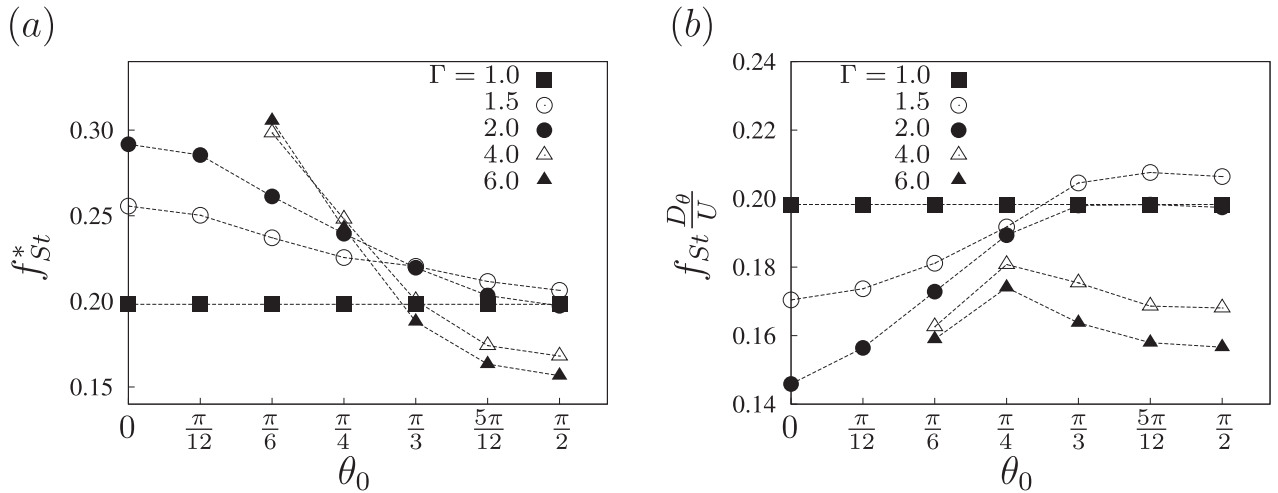


Fig. 2. For stationary cylinders, (a) the variation of vortex shedding frequency, f_{St}^* , with pitch angle, θ_0 , for the five geometries tested in the aspect ratio range, $1.0 \leq \Gamma \leq 6.0$. For aspect ratios $\Gamma = 4.0$ and 6.0, there is no vortex-shedding for pitch angles $\theta_0 = 0^\circ$ and 15° . Note, that due to the symmetry of the ellipse, this figure also encompasses results for aspect ratios in the range, $0.167 \leq \Gamma < 1.0$. In (b), the variation of vortex shedding frequency but with the frequencies normalized by the frontal area of the ellipse, rather than the major axis length D .

f_{St}^* exhibited by the other ellipses across the θ_0 range. The natural structural frequency, $f_N^* = 0.20$, serves as a starting point for exploring the parameter space, leaving open the possibility that other behaviours and responses may result from different selected values of f_N^* .

3.2. Flow-induced heave with active pitching

The results in this section are for cylinders free to heave, according to Eq. (2). An active pitching of amplitude $A_\theta = \pi/2$ and variable frequency is also applied. To begin with, results are shown for undamped configurations, $c=0$. For aspect ratio $\Gamma = 4.0$, Fig. 3 plots, at top, the variation of maximum displacement with driving frequency. Note that some of the simulations produce a non-symmetric heave response, therefore the maximum displacement is calculated as the largest displacement from the mean.

For driving frequencies $0.15 < f_D^* < 0.35$ the vortex shedding is synchronized with the driving frequency, while it is de-synchronized outside of this range. The most distinct variation with driving frequency is the increase in maximum amplitude near $f_D^* = 0.15$. This change in behaviour is also observed in the second panel of the figure plotting the maximum value of the coefficient of lift force, $C_L = 2F_y/\rho DU_\infty^2$. $C_{L,max}$ reaches its maximum value near the same driving frequency where the step change in maximum amplitude is observed. Similarly, in the fourth panel plotting the phase lag, ϕ_θ , between the heave response of the cylinder to the pitch oscillation, a step change from a phase lag of approximately $\pm \pi$ to approximately $-\pi/2$ at the same driving frequencies can be observed, clearly indicating a significant change in dynamic behaviour.

The third panel plots the absolute value of the mean of the lift force, $\overline{C_L}$; this quantity gives an indication of the mean symmetry of the flow and the cylinder dynamics. For $\overline{C_L} \neq 0$, the mean flow is non-symmetric around the y -axis. In relation to the driving frequencies where the large step changes in A_{max}^* and ϕ_θ are seen in the first and fourth panels, the value of $\overline{C_L}$ is greater than zero for isolated cases (for $f_D^* = 0.120$ and 0.150). This indicates a non-symmetric flow. For the $f_D^* = 0.150$ case, a small increase in driving frequency shifts the shedding to a different symmetric mode.

Fig. 4 plots the same data shown in Fig. 3, but for the five elliptical cross-sections of aspect ratios, $\Gamma = 1.0, 1.5, 2.0, 4.0$ and 6.0 . For each of the aspect ratios, the variation of the maximum displacement peaks in the region $0.14 \leq f_D^* \leq 0.20$. For aspect ratio $\Gamma = 1.0$, the peak displacement occurs for $f_D^* = 0.16$. For the elliptical cylinders of $\Gamma = 1.5, 2.0$ and 4.0 , the peak shifts to around $f_D^* = 0.18$, while for $\Gamma = 6.0$, the peak is at $f_D^* = 0.15$. Increasing aspect ratio, even by a small amount, increases the amplitude of oscillation; the stronger the heave the more potential for extracting energy.

Similar to the $\Gamma = 4.0$ case, the $\Gamma = 6.0$ case exhibits a distinct step change in A_{max}^* near a driving frequency $f_D^* = 0.150$, as does the $\Gamma = 2.0$ case, although to a lesser extent, near $f_D^* = 0.180$. For $\Gamma = 2.0, 4.0$ and 6.0 , $C_{L,max}$ reaches its maximum value near the same driving frequency where the step change in maximum amplitude is observed. The third panel also shows relatively isolated

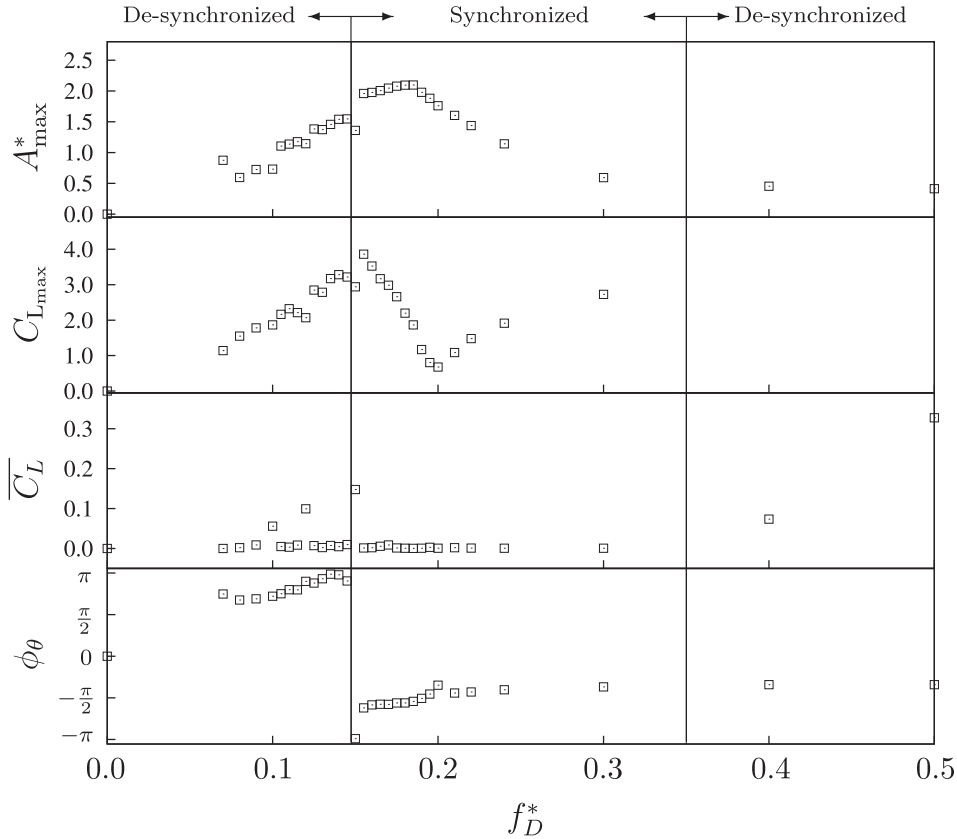


Fig. 3. For $Re=200$, $m^* = 5.0$ and $\Gamma = 4.0$, variation with driving frequency, f_D^* , of, from top, the maximum amplitude of vibration, A_{max}^* , maximum coefficient of lift, $C_{L,max}$, the mean coefficient of lift, $\overline{C_L}$, and the phase lag, ϕ_θ , between the heave response, Y , and the pitch, θ . Damping in all cases is set to $c=0$.

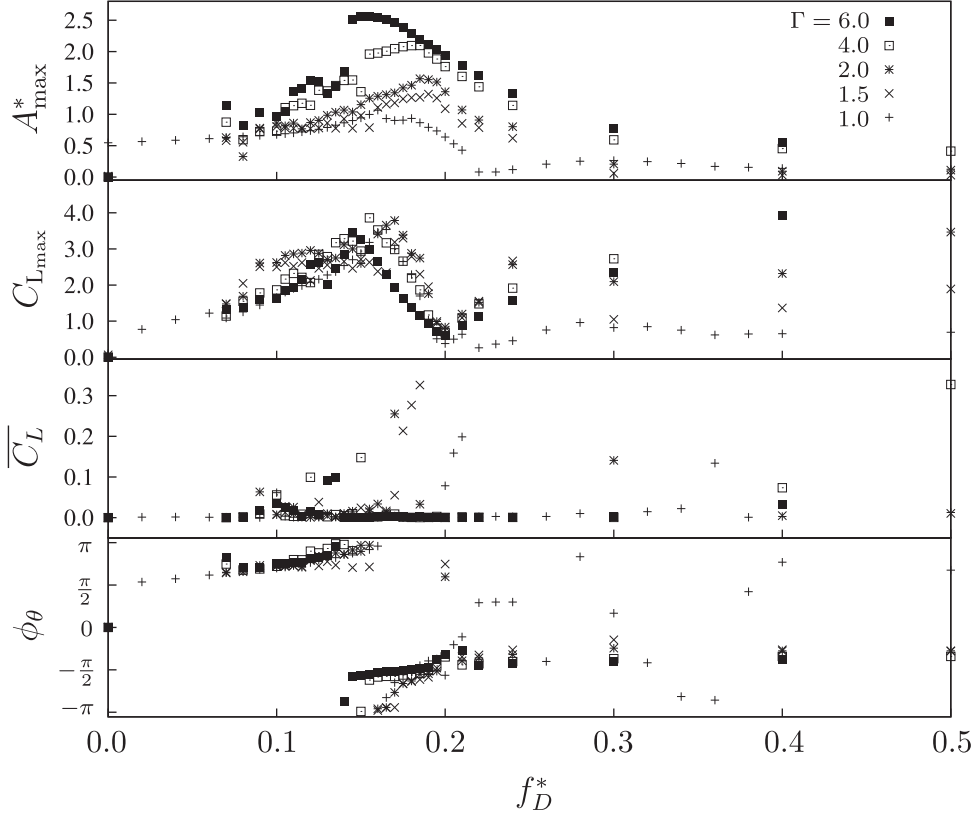


Fig. 4. For $Re=200$, $m^* = 5.0$ and $1.0 \leq \Gamma \leq 6.0$, variation with driving frequency, f_D^* , of, from top, the maximum amplitude of vibration, A_{Lmax}^* , maximum coefficient of lift, C_{Lmax} , the mean coefficient of lift, $\overline{C_L}$, and the phase lag, ϕ_θ , between the heave response, Y , and the pitch, θ . Damping in all cases is set to $c=0$.

cases of asymmetric flow near the transitions observable in the other panels. In the fourth panel plotting the phase lag, ϕ_θ , the variations with driving frequency show a general similarity across the aspect ratio range, with the de-synchronized cases of $f_D^* \lesssim 0.15$ trending upwards towards a phase lag of π and then a shift in the synchronized region to a phase lag of $-\pi/2$. [Nazarinia et al. \(2009\)](#) investigated the vortex shedding downstream of a circular cylinder with prescribed rotation and transverse oscillation, at $Re=1322$ (experimental) and 225 (numerical), varying the phase difference between the two oscillations. In contrast to the results of [Fig. 4](#), a desynchronized region, characterized by quasi-periodic and chaotic vortex-shedding, was observed for phase differences in the approximate range $-3\pi/4 \leq \phi \leq -\pi/4$. Results outside of this range, including the whole positive side of the phase difference range, exhibited locked vortex-shedding. The flow field of a locked case from the work of [Nazarinia et al. \(2009\)](#) (their Fig. 4b) bears strong similarity to a locked case from the current work for aspect ratio $\Gamma = 1.0$ (see the second case of [Fig. 6](#)). However, the current results show desynchronized vortex-shedding over a positive phase lag range, and locked vortex-shedding over a negative. In terms of phase, the two sets of results do not match up well, but there are key differences in configurations to be considered. Firstly, the amplitude of rotational oscillation in the current work is 1.5 times greater than in [Nazarinia et al. \(2009\)](#), which will significantly affect the vorticity generation on the cylinder. Furthermore, in the current work the prescribed transverse oscillation is replaced with a passive heave. Also, the results of [Nazarinia et al. \(2009\)](#) are for oscillation frequencies equal to the Strouhal frequency; they also found that small shifts in the driving frequency away from the Strouhal frequency could shift the boundaries between locked and unlocked shedding cases. In the current work, we see unlocked shedding for positive phase lags for driving frequencies of a half or three quarters of the Strouhal and structural frequencies, so different behaviour should be expected. The key point is that the wakes of the locked cases for $\Gamma = 1.0$ are qualitatively similar, even though there may be differences in phase between the pitching and heaving and in the frequencies of the systems.

Clearly, there are several modes of vibration present in the limited parameter space plotted in [Fig. 4](#). For the section of the parameter space corresponding to aspect ratios $1.0 \leq \Gamma \leq 6.0$ and driving frequencies $0.11 \leq f_D^* \leq 0.20$, [Fig. 5](#) plots contours of spanwise vorticity for each case, at intervals of $0.01f_D^*$. The focus is on this part of the parameter space as it encompasses the high-amplitude oscillation synchronized regime – of most interest for power extraction – and also the distinct variations with driving frequency observed in [Figs. 3 and 4](#). Each vorticity field in [Fig. 5](#) is plotted at the point of maximum heave. The aspect ratio increases from left, $\Gamma = 1.0$, to right $\Gamma = 6.0$, and the driving frequency increases from bottom, $f_D^* = 0.11$, to top, $f_D^* = 0.20$. The figure encompasses several changes in behaviour; however broadly speaking, at the lower end of the driving frequency range presented, the vortex-shedding is disordered and unsynchronized with the driving frequency. Cross-referencing with [Fig. 4](#), these cases exhibit smaller heave responses than for the synchronized cases. For the cases of $\Gamma > 1.0$, the variation with driving frequency of the phase lag between the heave response and the pitching actuation can be observed in the orientation of the cylinder at its peak heave. For the desynchronized cases, the ellipses, at maximum heave, are oriented with their minor axes aligned with the streamwise direction, while for the larger amplitude synchronized oscillations, the ellipses are oriented with the major axis aligned with the flow.

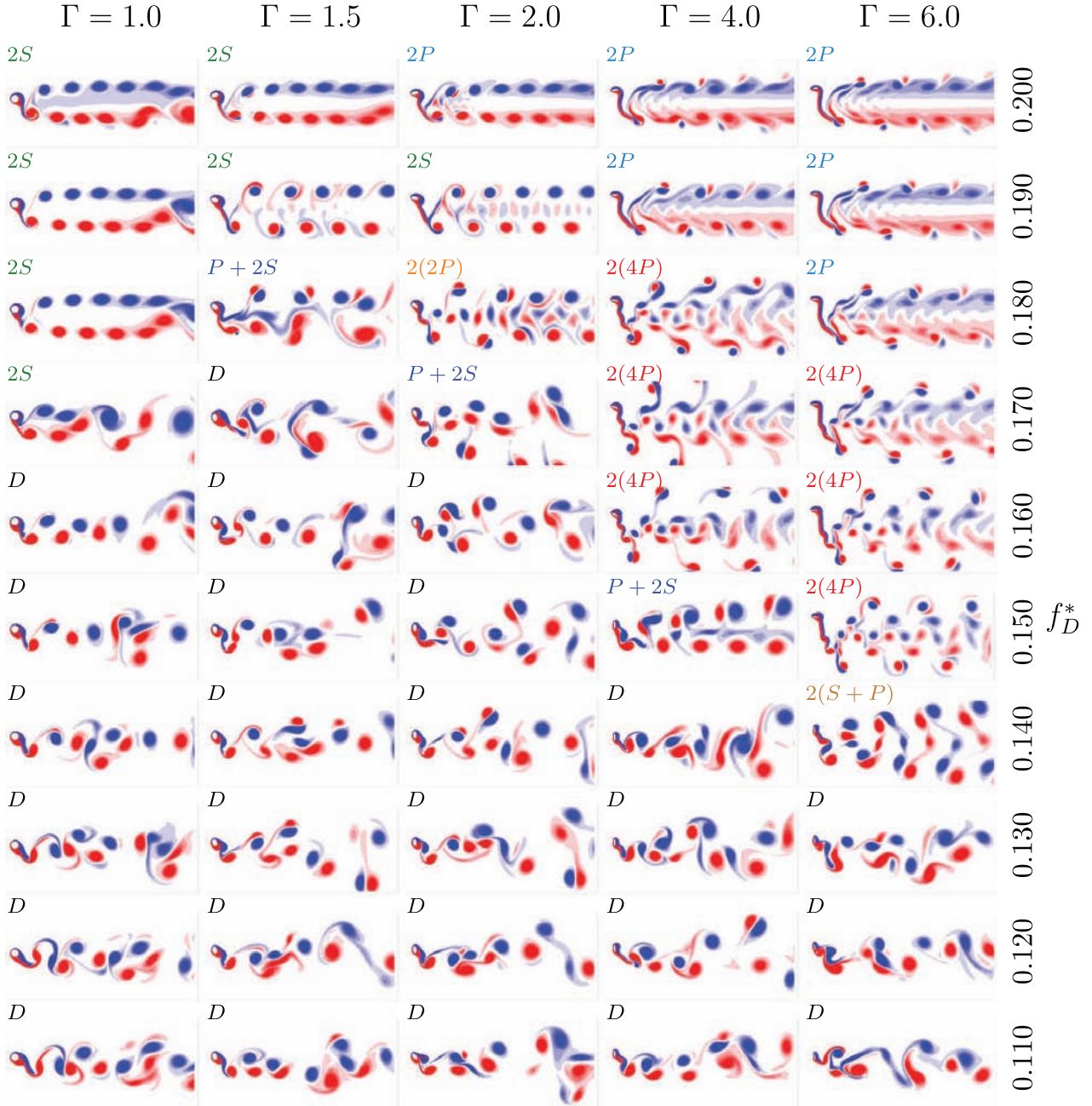


Fig. 5. Vorticity fields in the parameter space range, $0.110 \leq f_D^* \leq 0.200$ and $1.0 \leq \Gamma \leq 6.0$. Contours range $-1 \leq \omega_z \frac{D}{U} \leq 1$.

There are several modes of oscillation and vortex shedding, and an attempt has been made to categorize them on Fig. 5 using seven separate labels, each label describing the quantity of vortex pairs (P 's) and single vortices (S 's) shed each cycle of the pitching oscillation, following the original naming scheme devised by Williamson and Roshko (1988). Fig. 6 presents examples of each of the modes. The label D mode refers to cases where the vortex shedding is not synchronized, or desynchronized, with the driving frequency, f_D^* . In the example of the D mode in Fig. 6, there is no clear regular arrangement of the vortices evident in the vorticity contours. This disorder is reflected in the time series of the lift coefficient and the cylinder displacement. The frequency spectrum shows many peaks besides the primary frequency of $f/f_D = 1$. For these desynchronized cases, these appear for driving frequencies less than $f_D^* \simeq 0.15$, the frequency associated with the elastic structure, f_N^* , plays a role. In the case shown in Fig. 6, for aspect ratio $\Gamma = 1.0$ and driving frequency $f_D^* = 0.13$, the structural frequency is equal to $1.538f_D^*$. For this case – and for other D mode shedding cases – the frequency spectra reveal peaks at the structural frequency, $f = f_N$, and the driving frequency, $f = f_D$, but also at $f = f_N + (f_N - f_D)$. Such combinations imply some frequency modulation of the body oscillation (and presumably the vortex shedding) at f_N^* by the rotational oscillation at f_D^* (Chowning, 1973). This is similar to the frequency modulation observed in the wakes of streamwise oscillating cylinders in a free stream, where vortex shedding at one frequency is modulated by the driving frequency (Leontini et al., 2011, 2013).

For all the aspect ratios, for $f_D^* < 0.14$, desynchronized vortex-shedding is observed. As f_D^* increases, for all aspect ratios, ordered vortex-shedding modes appear. For the aspect ratio, $\Gamma = 1.0$, case, a $2S$ mode emerges, characterized by two single vortices of

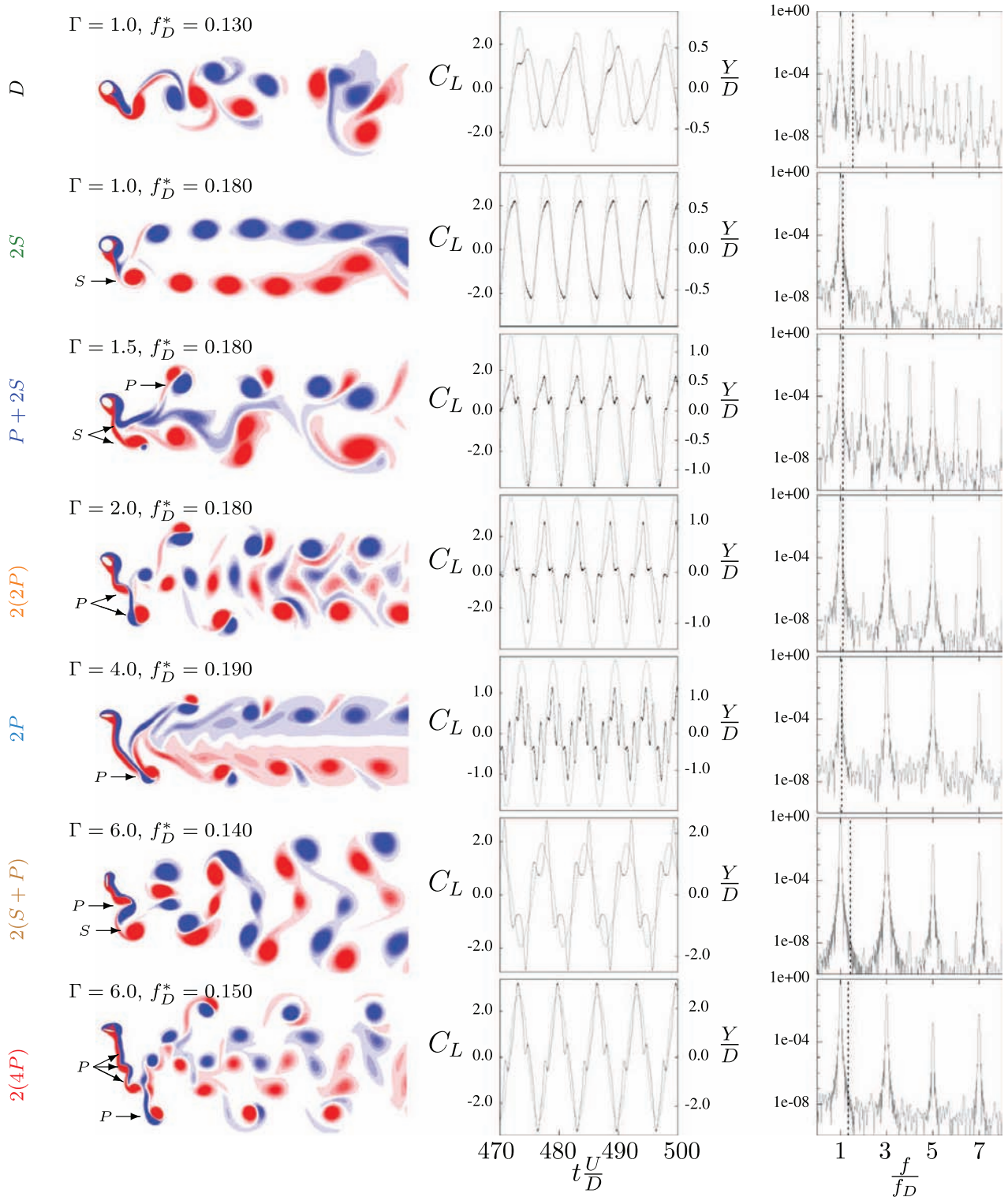


Fig. 6. Further details of each of the oscillation modes identified in Fig. 5, presenting at left contours of spanwise vorticity ($-1 \leq \omega_z \frac{D}{U} \leq 1$), in the middle, plots of the time signal of cylinder displacement and lift coefficient and at right, frequency spectra of the lift signal, normalized by the peak power. The dotted vertical line on the frequency spectra corresponds to $f = f_N$.

opposite sign shed each cycle of the pitching oscillation. In the example of Fig. 6, the ordered vortex-shedding can be observed in the times series and the frequency spectrum. The spectrum reveals peaks at the driving frequency, f_D^* , as well as monotonically decreasing peaks at the 3rd, 5th and 7th harmonics of the driving frequency. This arrangement of peaks is indicative of the spatio-temporal symmetry of this flow, at least in the near wake – only odd harmonics maintain this symmetry. Advanced one half-cycle, reflected around the x -axis, with the sign of the vorticity reversed, the vorticity topology would be essentially identical. The $2S$ mode only appears for the aspect ratio, $\Gamma = 1.0, 1.5$ and 2.0 geometries. For the cases of the mode for $\Gamma = 1.5$ and 2.0 , the flow is close to

transitioning to the $2P$ mode.

The $2P$ mode is similar to the $2S$ mode, but involves a tendril of the opposite sign vorticity joining the single vortex of the $2S$ mode and then rolling up to form the other half of a vortex pair. In the example of Fig. 6, the frequency spectra is similar to the spectra of the $2S$ mode, retaining the spatio-temporal symmetry, but with greater relative energy in the 3rd and 5th harmonics of the driving frequency.

For the geometries of aspect ratio $\Gamma = 2.0, 4.0$ and 6.0 , as the driving frequency is reduced from $f_D^* = 0.20$, where the $2P$ mode is present, the amplitude of oscillation increases. The two modes which appear, the $2(2P)$ and $2(4P)$ modes, are large amplitude types of the $2P$, retaining the same spatio-temporal symmetry. The extra vortices present in the wake are formed via an instability of the shear layers that are shed from the body.

The large vertical motion of the ellipse, with its major axis essentially aligned with this vertical motion, creates two shear layers with an increasing cross-stream orientation as the amplitude increases – essentially a vertical wake. As the amplitude of oscillation increases, the velocity of the cylinder increases, which in turn results in a narrower vertical wake. This vertical wake is susceptible to the same classic Kármán instability as any bluff-body wake, and it could roll up into a series of vortices behind the main vortex pair. There will be a critical vertical velocity for the onset of the Kármán instability, and further increases of the velocity beyond this will also influence the frequency, and wavelength, of the instability.

For the $2(2P)$ mode a single extra vortex pair (with respect to the $2P$ mode) is created per half cycle, while for the $2(4P)$ mode, three extra pairs are created. An examination of the $2(4P)$ mode is presented in Fig. 7. Each half cycle, four vortex pairs are shed. In the first panel of Fig. 7, the trailing edge vortex on the $y < 0$ side is labelled TEP_- . This pair forms when the ellipse reaches its minimum heave and the rotation of the ellipse is such that the major axis is aligned in the streamwise direction. On the $y > 0$, the opposite sign TEP_+ vortex pair can be seen forming. In the wake of each of these pairs, the two shear layers form a vertical Kármán street, which exhibits three distinct roll-ups. The first pair (labelled P_1 in Fig. 7) is also reinforced by a vortex formation

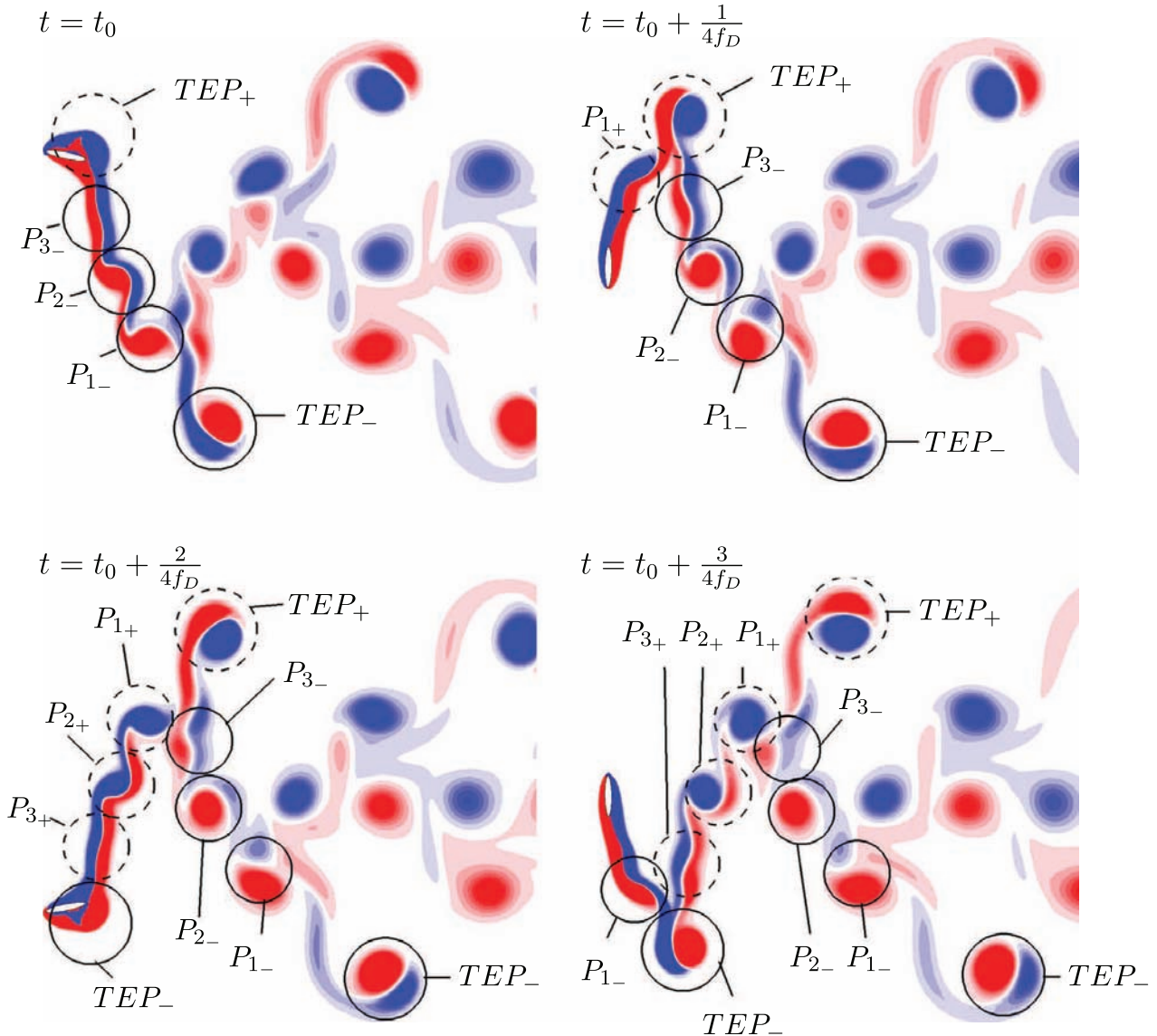


Fig. 7. A depiction of the $2(4P)$ shedding mode with vortex pairs labelled for the case of aspect ratio $\Gamma = 6.0$ and $f_D^* = 0.150$.

from the leading edge as the body sweeps from an essentially horizontal configuration at maximum displacement to a vertical configuration near zero displacement. Due to the complexity of the vorticity topology, some of the pairs are not as distinct as others; the vorticity of one sign of a pair dissipates, is destroyed, or interacts with another pair. The key point here is that the vorticity created on the cylinder surface rolls into distinct vortices and vortex pairs both in the cylinder wake and on the cylinder surface. In this way, the $2(4P)$ mode is similar to the $2P$ and $2(2P)$ modes in terms of vortex generation from the cylinder, differing principally in the vortex dynamics away from the cylinder. The frequency spectrum of the lift coefficient for the example of mode $2(4P)$ shows a strong peak on the 7th harmonic of the driving frequency, which reflects the multiple vortex pair formation each half cycle.

This appears to be related to the frequency of the instability of the vertical wake, and the symmetry of the flow. A whole number of wavelengths or vortex pairs shed per half cycle, combined with the main vortex, or pair of vortices, shed at the end of the half cycle will maintain the spatio-temporal symmetry of the flow. Adding one pair of vortices per half cycle will see one-and-a-half shedding “events” per half cycle, and three events per cycle – therefore a significant frequency component at the third harmonic. Adding two pairs sees five vortices per half cycle, and a contribution at the fifth harmonic. Adding three pairs sees seven vortices per half cycle and a contribution at the seventh harmonic, as observed for the $2(4P)$ wake and confirmed from observation of the spectrum of the lift force for this case in Fig. 6.

This vertical wake shedding also plays a role in determining the amplitude. Since the preferred modes see a whole number of vortex shedding cycles in the vertical wake, the amplitude needs to adjust to allow this to occur. This adjustment changes the number of shedding cycles in two ways. First, the larger amplitude provides more space to fit vortex pairs. Second, the increase in amplitude implies an increase in the cylinder velocity, which sees a decrease in the vertical wake width. This decrease in wake width sees a decrease in the vortex shedding wavelength (coupled with an increase in frequency), hence increasing the number of vortex pairs in a given length. The flow adjusts to set the amplitude close to an integer multiple of this vortex shedding wavelength.

As well as this vertical vortex shedding, the wake configuration is set by whether or not a pair or a single vortex is shed when the amplitude is maximum. This change is related to a change in phase between the rotation and heave ϕ_θ , which causes a change in the mechanism of vortex formation. The cases at $\Gamma = 6.0$, $f_D^* = 0.140$ ($2(S + P)$) and $\Gamma = 6.0$, $f_D^* = 0.150$ ($2(4P)$) illustrate these two mechanisms.

These two cases sit either side of the large step change in maximum cylinder displacement observed for $\Gamma = 6.0$ in Fig. 4. The mode at $f_D^* = 0.140$ is differentiated from the similarly named $2(P + S)$ mode because of its distinct vortex formation mechanism at the top of its travel related to its rotational phase.

The case at $f_D^* = 0.140$ has a phase $\phi_\theta \simeq \pi/2$, as shown in Fig. 4. This means that the ellipse is essentially pointing straight up and down at the top of the motion, as shown in Fig. 6. As the ellipse approaches the top of its travel it decelerates, and the negative vorticity in the shear layer on the back of the body is carried off by inertia and rolls up into a single vortex. This is the first mechanism of vortex formation. As the ellipse moves down towards the centre it also rotates, and the sweeping motion of the leading edge forms a positive vortex, which peels off a negative vortex from the shear layer now forming on the front of the body, hence forming a pair. This same process is repeated in the next half cycle.

The case at $f_D^* = 0.150$ has a phase $\phi_\theta \simeq -\pi/2$, as shown in Fig. 4. This means the ellipse is essentially pointing straight into the flow at the top of the motion, as shown in Fig. 6. As the ellipse approaches the top of its travel, it rotates from pointing straight up and down, and effectively “flicks” a negative-signed vortex from the trailing edge of the ellipse. This rotation also forms positive vorticity on the top side of the ellipse, which the newly formed negative vortex peels off, forming a vortex pair. This is the second mechanism of vortex formation. As the ellipse moves down towards the centre, it forms the essentially vertical wake that can break into subsequent vortex pairs due to the Kármán instability, in this case seeing three wavelengths, and three vortex pairs formed in the travel from top to bottom. This process is then repeated in the next half cycle. This complicated case is further illustrated in Fig. 7.

The other mode categorized in Fig. 5, the $P+2S$ mode, is an example of synchronized asymmetric shedding. This mode exists at the transition between the desynchronized modes of the cases of lower driving frequencies and the spatio-temporally symmetric $2P$ and related modes (note that the flow for $\Gamma = 6.0$ and $f_D^* = 0.135$ is also of the $P+2S$ mode, but is not shown in Fig. 5. The field is qualitatively similar to the flow for $\Gamma = 4.0$ and $f_D^* = 0.150$, but biased in the opposite direction in y). The mode is synchronized with

$$\omega_z = \frac{\partial v}{\partial x} - \frac{\partial u}{\partial y}$$

$$t = t_0$$

$$\omega_z = -\left(\frac{\partial v}{\partial x} - \frac{\partial u}{\partial y}\right)$$

$$t = t_0 + \frac{1}{2f_D}$$

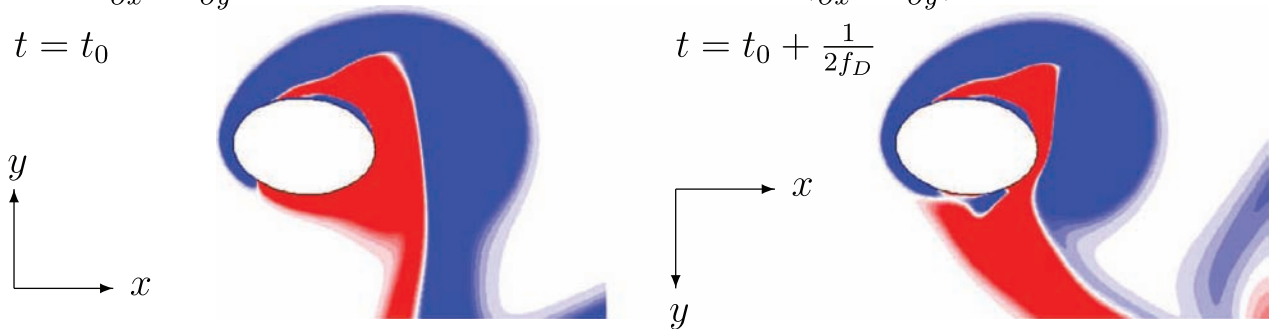


Fig. 8. For the case of aspect ratio $\Gamma = 1.5$ from Fig. 4, at left, a snapshot of vorticity near maximum positive displacement, with the ellipse aligned with the flow. At right, the same view of the flow, shifted in time half a period, reflected around the y -axis and the sign of vorticity reversed.

the driving frequency, yet has non-zero mean lift. This bias to one side is unusual, given the symmetry of the geometry and motion. A similar phenomenon has been observed for the vortex-shedding flow past square cylinders (Sheard et al., 2009) and also in the onset of $P+S$ wakes in the flow past oscillating circular cylinders (Leontini et al., 2006, 2007). This asymmetric phenomenon is illustrated in Fig. 8, for $\Gamma = 1.5$ and $f_D^* = 0.18$, which shows a closeup of the vorticity in the near wake region at the instant where the elliptical body is aligned with the oncoming flow, shortly after the snapshot at maximum displacement shown in Fig. 6 for the same case. The second image is of the same snapshot half a period later, with the flow reflected about the y -axis and the sign of the vorticity reversed. In a flow with complete spatio-temporal symmetry around the centreline, the images would be identical. In the case shown here, on one side of the oscillation, the formation of the vortex is cut by the trailing edge of the ellipse, while on the other, a larger vortex forms. It was noted in Sheard et al. (2009), as well as in Sohankar et al. (1999), that the asymmetry occurs in two-dimensional flows above the critical Reynolds number for absolute stability to spanwise perturbations and that the phenomenon does not persist in fully three-dimensional simulations. The appearance of the asymmetry suggests the two-dimensional flow may alter significantly when simulated in three-dimensions. The persistence of the findings in this study, and similar ones, to three-dimensional flows is a possible area of further interest, particularly as much of the literature on flapping aerofoil power generation to date has not considered three-dimensional effects.

The mode that generally exhibits the greatest cylinder oscillation amplitude and velocity is the $2(P+2S)$ mode. A greater heave indicates the potential for power extraction. However, the actuation of the pitching motion requires energy input. Although higher aspect ratio ellipses produce greater heave, as shown in Fig. 4, the increasing alignment of the surface of the ellipse normal to the tangential direction of the pitch, leads to an increasing pressure component to the torque applied by the fluid on the body. This increased resistance needs to be overcome, proving a negative energy input to any final net power output, meaning this rotational actuation power input, P_{in} , needs to be considered.

3.3. Calculation of P_{in}

To rotationally oscillate the cylinder in the fluid, an input power is required, which will reduce the efficiency of any energy extraction from the heave. Therefore, some consideration of its variation and magnitude is needed. The equation of motion for the cylinder rotation is given by

$$\tau = I\alpha \quad (3)$$

where τ is the torque, I is the rotational moment of inertia and α is the rotational acceleration. The total torque is the sum of the torque needed to drive the oscillation, τ_d , and the torque that the fluid exerts on the cylinder, τ_f . Multiplying through by the rotational velocity, ω , Eq. (3) can be cast as a balance of power, such that

$$\tau_d\omega + \tau_f\omega = I\alpha\omega \quad (4)$$

In this way, the work done by any external control such as a motor over one oscillation cycle to rotate the cylinder is given by

$$\int_{t=t_0}^{t=t_0+T} \tau_d\omega \, dt, \quad (5)$$

and the (rotational) work done by the fluid on the cylinder over one oscillation cycle is given by

$$\int_{t=t_0}^{t=t_0+T} \tau_f\omega \, dt. \quad (6)$$

I , α and ω are known or prescribed, so an expression for the instantaneous power needed to rotate the cylinder is

$$P_{in} = \tau_d\omega = I\alpha\omega - \tau_f\omega. \quad (7)$$

The mean power input over an oscillation cycle is given by

$$\bar{P}_{in} = \frac{1}{T} \int_{t=t_0}^{t=t_0+T} P_{in} \, dt. \quad (8)$$

Dimensional input and output power terms (which are per unit of spanwise length), when expressed in this study, are assumed non-dimensionalized by $\rho U^3 D$.

The instantaneous torque on the cylinder exerted by the fluid can be calculated by integrating the product of the surface stress component normal to the vector from the axis of rotation and the length of this vector over the surface of the cylinder, i.e.,

$$\tau_f = \int_S \sigma \cdot ((\hat{\mathbf{r}} \times \hat{\mathbf{k}}) / |(\hat{\mathbf{r}} \times \hat{\mathbf{k}})|) |\mathbf{r}| \, ds, \quad (9)$$

where σ is the fluid stress vector from pressure (acting normally) and the traction vector from shear stress, S is the surface of the cylinder and \mathbf{r} is the vector from the axis of rotation to the cylinder surface, and $\hat{\mathbf{k}}$ is the unit vector aligned with the axis of rotation.

For aspect ratio $\Gamma = 1.0$, the only torque exerted by the fluid on the circular cylinder comes from the fluid viscosity, which generally acts in the opposite direction to the rotation. The entire surface of the circular cylinder is arrayed tangentially to the rotation, so there is no pressure component to fluid torque, τ_f . The amplitude of the oscillation is constant, so as the driving frequency, f_D^* , increases, so too does the magnitude of α and hence of P_{in} . The positive value of \bar{P}_{in} is an energy cost and reduces the

efficiency of the power extraction. However, the behaviour does not hold for higher aspect ratios. As aspect ratio increases from $\Gamma = 1.0$, the elliptical surface is increasingly perpendicular to the direction of the rotation, leading to the growth of the pressure component contribution to τ_f . What does this mean for any final sum of extracted energy?

3.3.1. Consideration of energy extraction scenarios affecting \bar{P}_{in}

A representative case is chosen from Fig. 4, for $\Gamma = 4.0$ and $f_D^* = 0.180$, which returns mean power input $\bar{P}_{in} = 0.156$. Fig. 9(a) plots the variation of P_{in} over two oscillation cycles. From the plot, the positive value of \bar{P}_{in} is apparent; however, the sign of P_{in} changes several times during each cycle; sometimes extra power is needed to drive the cylinder, sometimes power needs to be extracted or absorbed to resist the acceleration coming from the fluid. But to have the value of $\bar{P}_{in} = 0.156$ included in the final calculation of the efficiency of energy extraction from the heave, assumes that the energy absorbed during those parts of the cycle where the fluid is acting to accelerate the rotation of the cylinder beyond its prescribed path, is fully collected. This is a best case scenario. In an application of the energy-harvesting method presented in this paper, the driving mechanism will have its own efficiency. Investigating the efficiency of a driving mechanism is beyond the scope of the current paper. However, the following results are included to give an understanding of the effect on net power of different input power scenarios. The three input power scenarios presented effectively cover a range of input power efficiencies.

Fig. 9(b) plots the same variation as in (a), only with the negative parts of the series set to zero, P_{in1} . This imagines a scenario where energy is supplied to drive the cylinder, but the work done by the fluid on the cylinder is dissipated by the structure, not collected and not included in the final calculation of total power. For this variation, \bar{P}_{in} is greater, with a value of 0.235. Fig. 9(c) plots the absolute value of the variation of Fig. 9(a), P_{in2} . This is a worst case scenario, where energy is spent accelerating the rotation, but also spent in braking. Clearly, this is more expensive again, giving a \bar{P}_{in} of 0.313. For aspect ratio $\Gamma = 4.0$, the variation of \bar{P}_{in} with f_D^* for the three scenarios plotted in Figs. 9(a), (b) and (c), is shown in Fig. 9(d), illustrating the range of possible values of \bar{P}_{in} .

In the following section, the variation of the mean power extracted from the heave, \bar{P}_{out} , is explored independently of \bar{P}_{in} ; following that, the importance of \bar{P}_{in} to the optimal cases for \bar{P}_{out} will be considered, accounting for the three possible scenarios outlined above.

3.4. Damping to include energy extraction from the heave in the model

To simulate power extraction from the cylinder heave, a linear damping term, c , is applied to Eq. (2). The instantaneous power extracted from the heave is given as

$$P_{out} = c\dot{Y}^2, \quad (10)$$

where \dot{Y} is the velocity of the cylinder transverse to the oncoming flow. The average power output is given by

$$\bar{P}_{out} = \frac{1}{T} \int_{t=t_0}^{t=t_0+T} P_{out} dt. \quad (11)$$

Again, the power term per unit of spanwise length is non-dimensionalized by $\rho U^3 D$. The power output scales with the damping, c , however the higher the damping the stronger the effect of the second term of Eq. (2). The damping increases the resistance to motion, reducing the magnitude of \dot{Y} and of P_{in} . Thus power extraction affects the motion of the cylinder and thereby the flow. This interaction means that an optimal damping term for power extraction will exist for each aspect ratio.

Fig. 10 plots the same scalars as for the undamped cases presented in Fig. 4, but for the $\Gamma = 4.0$ case for various values of the damping, as a function of the driving frequency f_D^* . The first feature that is most apparent is that the overall behaviour of the system is quite similar for all the values of damping, with only an apparent scaling of the amplitude. As the damping is increased from

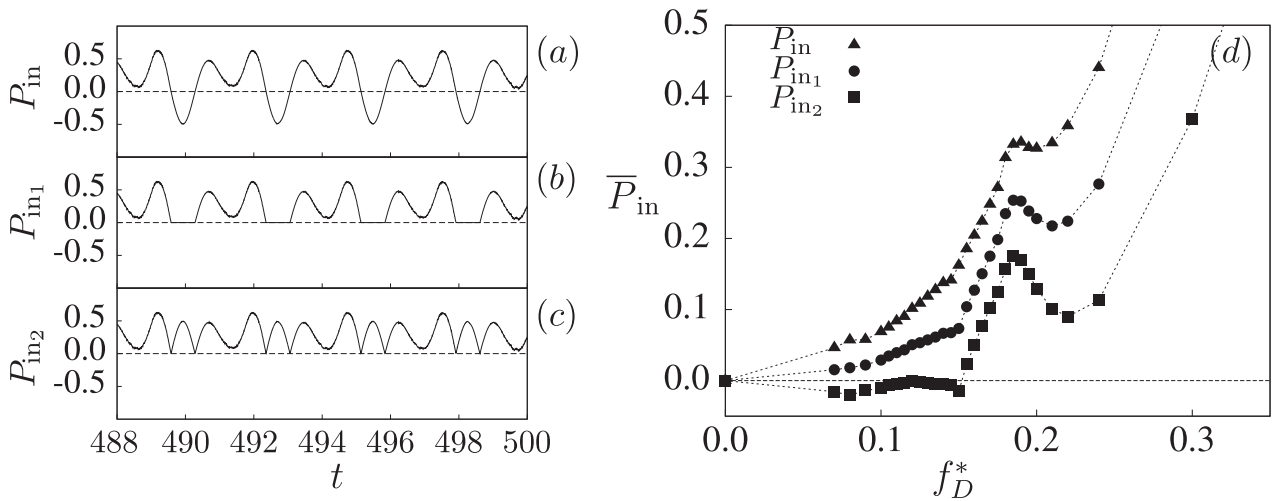


Fig. 9. For the case of aspect ratio $\Gamma = 4.0$, $c=0.0$ and $f_D^* = 0.180$, the variation with time of input power with (a) complete rotational energy extraction, P_{in} , (b) no rotational energy extraction, P_{in1} , and (c) energy input for braking, P_{in2} . In (d), the variation of \bar{P}_{in} with f_D^* for P_{in} , P_{in1} and P_{in2} .

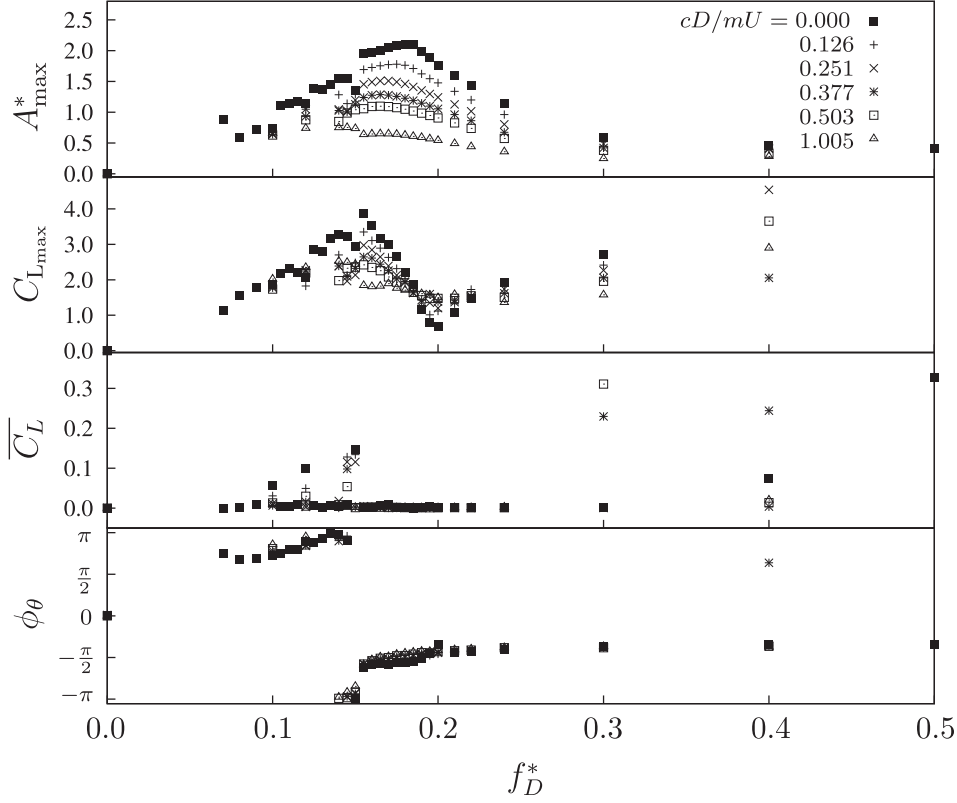


Fig. 10. For $Re=200$, $m^* = 5.0$ and $\Gamma = 4.0$, variation with driving frequency, f_D^* , of, from the top, the maximum amplitude of vibration, A_{\max}^* , the maximum lift coefficient, $C_{L_{\max}}$, the absolute value of the mean lift, \overline{C}_L , and the phase lag between the rotation and the heave, ϕ_θ , across a range of damping coefficients, cD/mU .

$c = 0.0$, the amplitude of oscillation decreases monotonically across the range of driving frequency. This reflects the increasing effect of the $(c/M)\dot{Y}$ term in the harmonic oscillation described by Eq. (2).

The two branches of the amplitude response tied to the phase between the controlled pitching and the heaving motion (which is

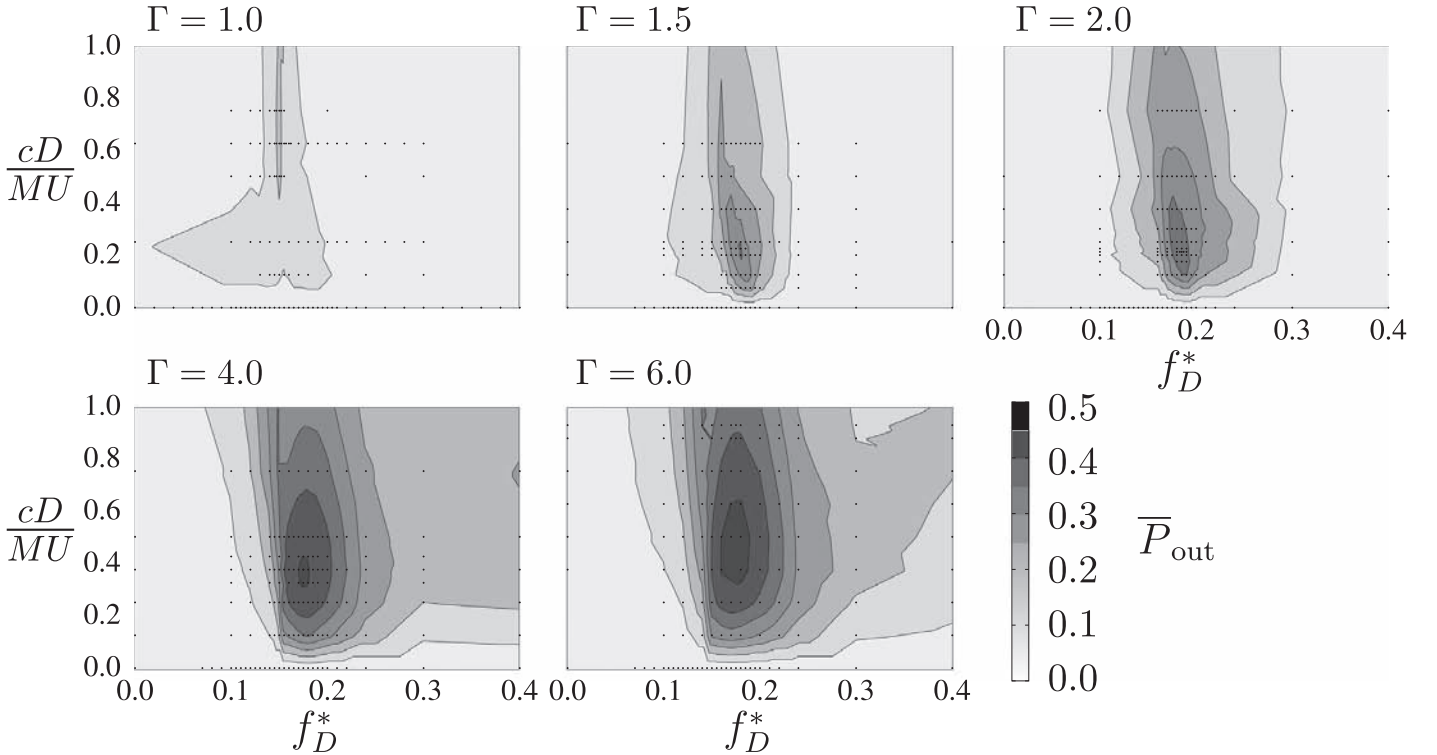


Fig. 11. Contour maps of mean output power, $\overline{P}_{\text{out}}$, for aspect ratios, $\Gamma = 1.0, 1.5, 2.0, 4.0$ and 6.0 , over the driving frequency, f_D^* , normalized damping coefficient, cD/mU , parameter space. Axis ranges and contour scales are constant across the five maps.

also related to a change in main vortex production mechanism as explained in Section 3.2) still exist. The asymmetric modes, apparent due to a non-zero mean lift, are also recovered at low driving frequencies.

A focus of this study is the impact this damping has on the power extracted from the flow, and therefore the energy harvesting potential. For the zero damping set of results, the mean power output, \bar{P}_{out} , calculated using Eq. (10), is zero. As damping increases, \bar{P}_{out} initially increases, but the velocity of the cross-stream oscillation is also reduced. This will eventually reduce the power extracted. Thus an optimal damping occurs for power output. For the cases shown in Fig. 10, of aspect ratio $\Gamma = 4.0$, the optimal configuration is for a damping $cD/mU = 0.377$ and driving frequency $f_D^* = 0.175$, producing a mean power output of $\bar{P}_{out} = 0.354$.

For all of the aspect ratios tested, Fig. 11 presents contour maps of mean power output, as function of both driving frequency, f_D^* , and the damping coefficient, cD/mU , for the aspect ratios, $\Gamma = 1.0, 1.5, 2.0, 4.0$ and 6.0 . The axis ranges and contour scale are constant across the five plots. The limiting case $f_D^* = 0.00$ represents the case with no rotational forcing. The plots show that across the geometries an increase in power output from the unforced case is possible. For the five geometries, the cases of maximum power output are grouped mostly around a driving frequency of $f_D^* = 0.180$, with the exception of the $\Gamma = 1.0$ case, where the optimal case is at $f_D^* = 0.150$. All these driving frequencies for optimal power output are less than the natural structural frequency of the system, $f_N^* = 0.20$.

With the darkening colour contours with aspect ratio, an increase in the magnitude of power output with aspect ratio is apparent. The sharpest change in behaviour is from the circular cylinder of aspect ratio, $\Gamma = 1.0$, to the higher aspect ratio cases. Fig. 12(a) plots the variation of the different formulations of the input power, discussed in Section 3.3.1, \bar{P}_{in} , \bar{P}_{in1} and \bar{P}_{in2} across the optimal cases for the five aspect ratios tested.

Interestingly, for the non-idealized \bar{P}_{in} scenarios, the lower aspect ratio optimal cases perform poorly, the required input power increasing markedly for the P_{in} formulations of Figs. 9(b) and (c). For the case of aspect ratio $\Gamma = 6.0$, the worst case scenario for input power requires a mean input power of more than half of the mean output power. This power input constitutes a significant effect on the net power output, illustrating the importance of optimizing the actuation mechanism for the cylinder rotation. Fig. 12(b) plots the net power ($\bar{P}_{out} - \bar{P}_{in}$) for the three different mean power input formulations. The highest values for \bar{P}_{net} are for the higher aspect ratio cases with marginal differences between the geometries of $\Gamma = 4.0$ and 6.0 .

The effect on net power output of the different scenarios of mean input power calculation is presented in Fig. 13. For the $\Gamma = 4.0$ geometry, the contour maps show the results for \bar{P}_{out} alone, $\bar{P}_{out} - \bar{P}_{in}$, $\bar{P}_{out} - \bar{P}_{in1}$ and $\bar{P}_{out} - \bar{P}_{in2}$. The net power output reduces as expected with the inclusion of the input power, but the location in the parameter space of the optimal cases does not significantly shift. This is to be expected given, firstly, the third panel of Fig. 10, which plots the variation with driving frequency of the \bar{P}_{in} for a range of damping values. The plot shows that as the damping increases, for a given driving frequency, the variation of the input power is strongly asymptotic. Secondly, Fig. 9(b), plotting the variation of \bar{P}_{in} , \bar{P}_{in1} and \bar{P}_{in2} with driving frequency for $\Gamma = 4.0$, shows that the differences between the three input power scenarios described are near invariant with driving frequency. Thus, the variations in the contour maps of Fig. 13 consist largely in the contour levels, with little variation with driving frequency or damping.

To relate the flows for optimal power output to the undamped flows of maximum oscillation velocity, Fig. 14 plots the vorticity fields for each case of maximum power output from Fig. 11, with the cylinder in each case at maximum displacement. For all aspect ratios, the case of optimal power output exhibits the 2S mode of vortex shedding. However, the fields for aspect ratios $\Gamma = 2.0, 4.0$ and 6.0 are showing signs of transitioning to a 2P mode, with a tendril of opposite sign vorticity trailing each single vortex, but not yet forming a distinct vortex to complete the pair of the 2P mode. Also, the more complicated vortex wakes such as the 2(P+S) and 2(P+2S) modes are not recovered; the peaks of the frequency spectra decrease monotonically on the 3rd, 5th and 7th harmonics. The optimal flows exhibit the same, or nearly the same, driving frequencies as for the undamped cases of maximum oscillation amplitude.

Generally, the damping seems to only limit the oscillation amplitude of the body. This limitation in the amplitude also restricts the wake modes observed, as the more complicated patterns are formed by the flow trying to “lock” a whole number of vortex pairs

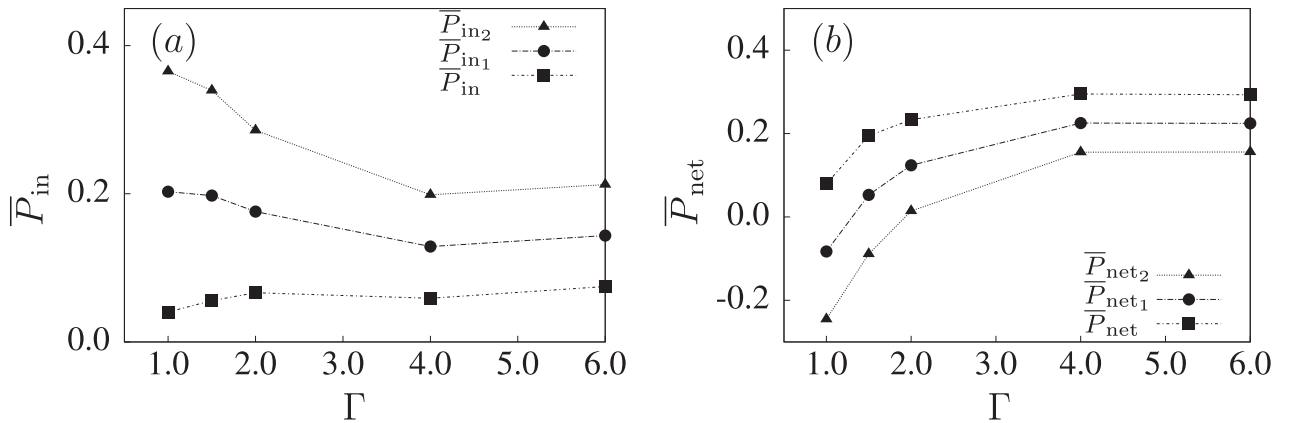


Fig. 12. For the optimal power output cases of Fig. 11. (a) The variation with aspect ratio of \bar{P}_{in} , \bar{P}_{in1} and \bar{P}_{in2} and (b) the variation of \bar{P}_{net} calculated using the three different mean power input formulations.

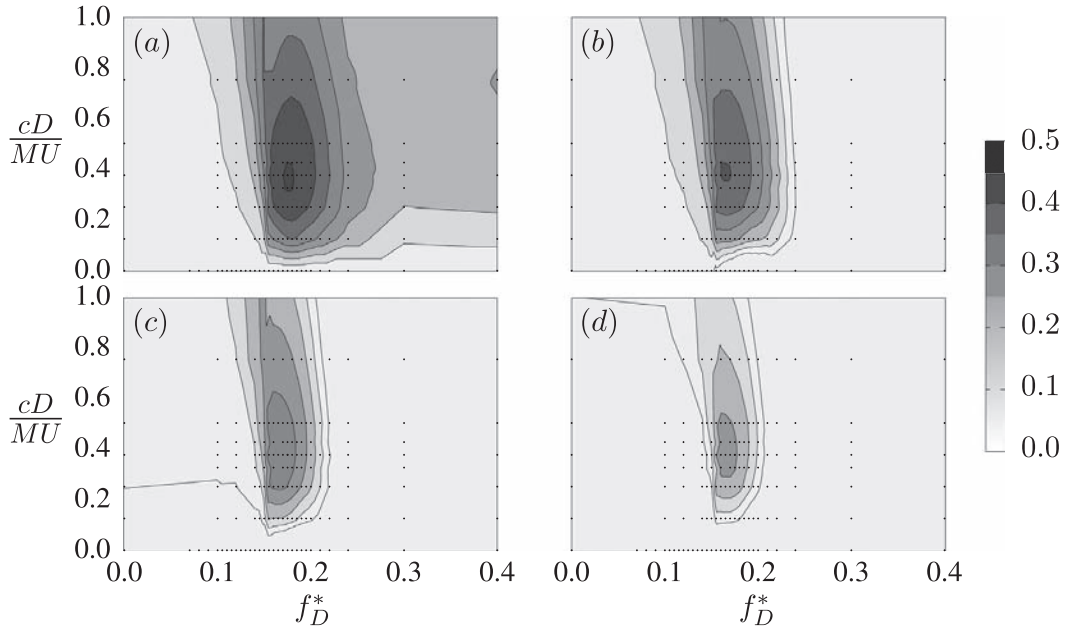


Fig. 13. For aspect ratio, $\Gamma = 4.0$, contour maps of (a) output power, \bar{P}_{out} , (b) output power minus input power, $\bar{P}_{\text{out}} - \bar{P}_{\text{in}}$, (c) output power minus input power with no rotational energy-harvesting, $\bar{P}_{\text{out}} - \bar{P}_{\text{in}_1}$, and (d) output power minus input power with braking energy, $\bar{P}_{\text{out}} - \bar{P}_{\text{in}_2}$. Axis ranges and contour scales are constant across the five maps.

into one oscillation half cycle – lower amplitude mean less pairs, and therefore simpler vortex wakes. The flows of the optimal cases seem to be lower amplitude variants of the undamped cases. A useful comparison could be made with the undamped cases exhibiting the same amplitude oscillation amplitude, but at different frequency.

4. Conclusions

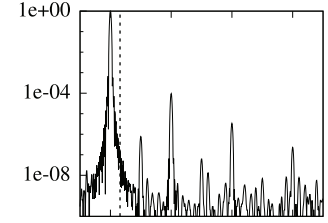
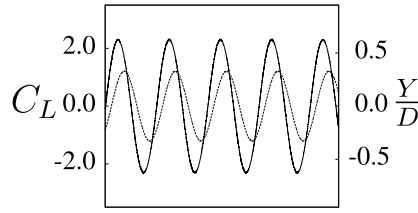
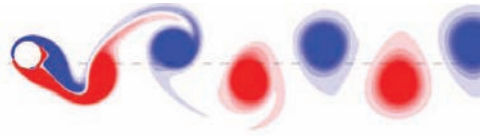
The large parameter space covering the flow past elastically mounted rotationally oscillating elliptical bluff bodies reveals a rich range of flow phenomena, such as the complex shedding and vortex roll-ups seen in the flow past undamped high-aspect ratio ellipses. As aspect ratio is increased from the circular cylinder to the ellipse, converging towards a flat plate, the amplitude and velocity of the cross-stream oscillation increases significantly. It was shown that for the class of problem investigated here, what begins as a typical vortex-shedding wake for the circular cylinder, $\Gamma = 1.0$, transitions quickly to a more complex field at higher aspect ratio, with essentially a vertical wake aligned laterally to the oncoming flow rolling up and shedding at higher frequency than the rotation driving frequency. Also analysed was the input power for the actuation of the cylinder rotation; it was shown that the cylinder rotation requires a torque working both with and against the torque exerted by the fluid on the body at different periods of the cycle. Therefore, the efficiency of any power output from energy harvesters should be considered against the accelerating and decelerating mechanisms of the rotational actuation.

Power output was modelled as a damping on the velocity of the cross-stream oscillation, producing optimal cases of damping and driving frequency. The frequencies of the optimal cases for power output corresponded closely with the frequencies of the maximum amplitude undamped cases, the flows exhibiting similarities, except for the case of aspect ratio, $\Gamma = 1.5$.

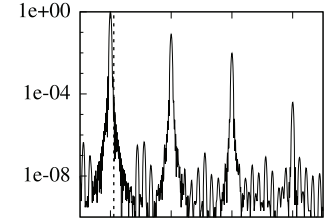
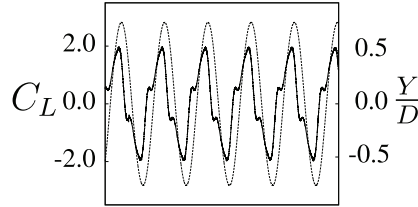
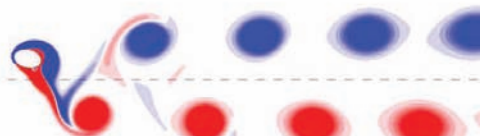
The number of parameters available to test for a theoretical flapping foil energy harvester is extremely large and only a small portion of the potential parameter space has been explored in the current paper. Many similar studies use a NACA aerofoil geometry, rather than the ellipses of the current study. Although the current study find a convergence with increasing thinness (or aspect ratio) of the cylinder, many aerofoils more naturally rotate around a pitching axis near their centre of mass, generally around a third of a chord downstream of the leading edge. The effect of pitching axis has not been examined here. Including another cylinder geometry, such as a NACA aerofoil of aspect ratio 4 or 6 would also yield interesting comparisons.

Also of interest is the extension to three dimensions. Whether the addition of oscillatory pitching and/or passive heave will make the flow less or more unstable to three-dimensional modes is not clear. While the added energy from the cylinder pitch could make the flow more unstable, there exists in the literature examples of body motion or oscillation having a stabilizing effect for three-dimensional stability of bluff-body flows, most notably [Lo Jacono et al. \(2010\)](#), who found that rotationally oscillating a circular cylinder in a freestream damped the three-dimensional instability modes usually observed. For the higher aspect ratio cases presented in the current work, it is possible that they are unstable to three-dimensional perturbation, although this would have to be tested. Future study could analyse the stability of the two-dimensional flows, as well as simulate the flows in three dimensions, both with infinite spanwise cylinder length and a finite spanwise aspect ratio.

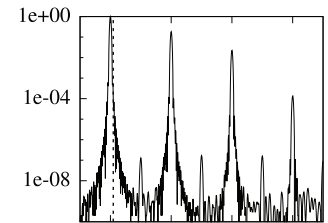
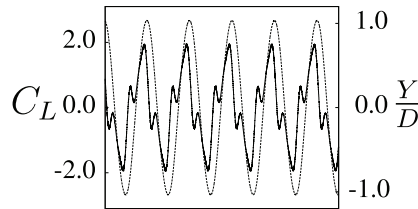
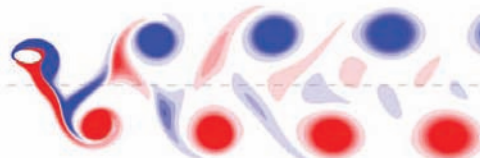
$$\Gamma = 1.0, \frac{cD}{MU} = 0.628, f_D^* = 0.152$$



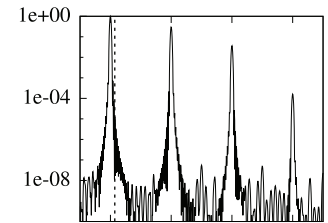
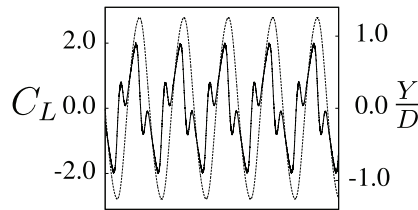
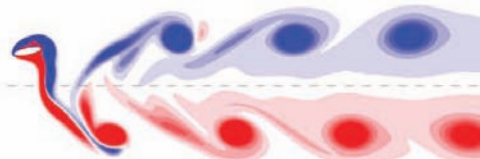
$$\Gamma = 1.5, \frac{cD}{MU} = 0.251, f_D^* = 0.180$$



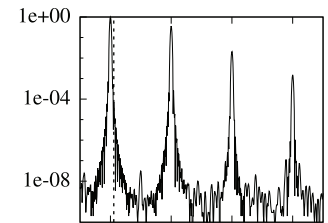
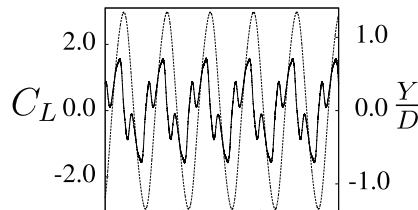
$$\Gamma = 2.0, \frac{cD}{MU} = 0.214, f_D^* = 0.183$$



$$\Gamma = 4.0, \frac{cD}{MU} = 0.377, f_D^* = 0.175$$



$$\Gamma = 6.0, \frac{cD}{MU} = 0.503, f_D^* = 0.180$$



$t \frac{U}{D}$

$\frac{f}{f_D}$

Fig. 14. For each aspect ratio, the vorticity fields at the moment of maximum body displacement for the optimal cases taken from Fig. 11. The time signals of the lift force and the cylinder displacement are plotted. At right, the frequency spectra of the lift signal are plotted, with the frequency normalized by the driving frequency, f_D . The dotted vertical line on the frequency spectra corresponds to $f = f_N$. The parameters for each optimal case are given with each vorticity field.

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