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# Photoinduced reversible spin-state switching of an Fe<sup>III</sup> complex assisted by a halogen-bonded supramolecular network†

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**The organization of a molecular Fe<sup>III</sup> complex embedded in a halogen-bonded 2D network is chemically tuned to trigger temperature and light-induced spin-state switching. We attribute the associated magnetic properties and the unprecedented photoswitching effect to the optimized structural confinement provided by the presence of the supramolecular host framework.**

Photoswitching between electronic states of molecules in solid state is a continuous challenge for both technological applications and fundamental science research.<sup>1</sup> Spin crossover complexes, that undergo an equilibrium between a high-spin and a low-spin electronic configurations, can offer such an optically switchable bistability, known as light-induced excited spin state trapping or the so-called LIESST effect.<sup>2</sup> Since its discovery over 30 years ago, the reported examples exhibiting this effect have been predominantly Fe<sup>II</sup> complexes.<sup>3</sup> In contrast, the manifestation of the LIESST effect in Fe<sup>III</sup> complexes remains scarce.<sup>4–8</sup> This has been attributed to a smaller elongation of the Fe–ligand distance, relative to that of Fe<sup>II</sup>, which lead to a short lifetime of the photoinduced metastable state.<sup>4,5,9</sup>

To prevent this fast relaxation from a metastable state to the ground state, a severe distortion of the metal coordination environment, in conjunction with the introduction of extensive intermolecular  $\pi$ – $\pi$  interactions, has been proposed.<sup>5</sup> Albeit several successful examples, this strategy is not a guarantee for most Fe<sup>III</sup> complexes.<sup>10</sup> Recently, a Fe<sup>III</sup> complex inserted into a metal-oxalate network has been reported to show the LIESST effect, suggesting that a structural confinement of the Fe<sup>III</sup> complex may be an

strength, to rationally construct a supramolecular framework where Fe<sup>III</sup> molecules could be embedded. Such a network will serve not only to introduce the intermolecular elastic interactions but also to tune the spatial organization and the confinement of the guest molecules, both of which are potential important factors to activate or deactivate the LIESST effect.

We recently reported the implementation of anionic halogen-bonded supramolecular networks to drastically change the spin-state of molecular guest complexes.<sup>12</sup> The modulation of electronic and magnetic properties of the guest molecules stemmed from a significant modification of their crystal packing depending on the presence and the nature of the halogen bonding. Herein, we report the role of reaction solvent medium to control the inclusion of the molecular Fe<sup>III</sup> complex into the halogen-bonded supramolecular network and to activate the “breathing” of the host network. As a consequence, the guest molecules undergo spin crossover, light-induced spin state trapping, and photoswitching, providing an extremely rare example of a Fe<sup>III</sup> complexes whose spin state can be reversibly controlled both by light and temperature. It is important noting that the role of solvent is here a chemical tool to drastically modify the final supramolecular assembly,<sup>13</sup> which differentiates this study from previous investigations of the solvation effect on spin crossover molecules.<sup>14</sup>

The halogen-bonded network encapsulating an Fe<sup>III</sup> complex was synthesized following previously reported procedures but with modification of the reaction solvent.<sup>11</sup> The reaction of [Fe(qsal)<sub>2</sub>]Cl (Hqsal = *N*-(8-quinoly)salicylaldimine) with a mixture of 17 eq. of *sym*-triiodotrifluorobenzene (C<sub>6</sub>F<sub>3</sub>I<sub>3</sub>) and 1.0 eq. of KI in MeOH yielded dark plate-like crystals of [Fe(qsal)<sub>2</sub>][(C<sub>6</sub>F<sub>3</sub>I<sub>3</sub>)I] (**1**). The similar reaction in MeCN afforded dark needle-shaped crystals of a solvated compound [Fe(qsal)<sub>2</sub>][(C<sub>6</sub>F<sub>3</sub>I<sub>3</sub>)I]·MeCN (**2**), as previously reported.<sup>11</sup> Single-crystal X-ray analysis reveals that the structure of compound **1** consists of 2D honeycomb-like halogen-bonded framework [(C<sub>6</sub>F<sub>3</sub>I<sub>3</sub>)I]<sup>−</sup>, as observed in **2**. This network results from the tritopic character of the halogen bond donor, C<sub>6</sub>F<sub>3</sub>I<sub>3</sub>, and the three-fold coordination of the iodide anion as a halogen bond acceptor (Figures 1). The I⋯I<sup>−</sup> bond distances in **1** at 298 K are 3.532(8), 3.619(6), and 3.731(7) Å, that are 15, 13, and 11 % shorter than the sum of the van der Waals and ionic radii, respectively. These values indicate the employed halogen bonding interaction is more symmetric than those found in **2** that features 16, 15, and 3 % reduction. The 2D network of **1** is distorted and undulated, similar to that of **2**, with a central I<sup>−</sup> ion at *ca.* 0.9 Å off from the plane formed by three iodine atoms (Figure S1).

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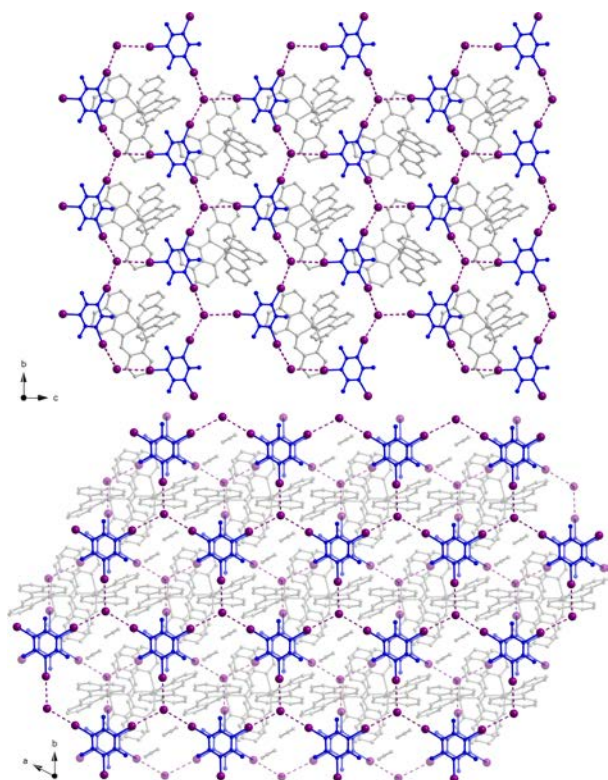
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†Electronic Supplementary Information (ESI) available: Experimental details, crystallographic data, crystal structure of **1**, additional optical and magnetic data, and crystallographic information files (CIF). CCDC [1550127-1550128](https://doi.org/10.1039/x0xx00000x) and [1568816](https://doi.org/10.1039/x0xx00000x). See DOI: 10.1039/x0xx00000x

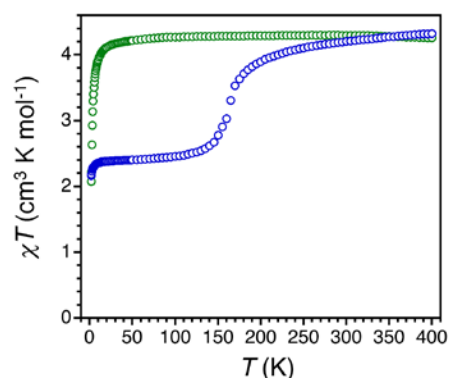
alternative strategy.<sup>7</sup> Despite these advances, the design of Fe<sup>III</sup> spin crossover systems exhibiting the photoswitching effect still represents a formidable challenge to understand key parameters governing the lifetime of the photo-induced metastable phase of Fe<sup>III</sup> compounds. Toward this end, we envision to harness halogen bonding interactions,<sup>11</sup> that feature high directionality and



**Figure 1.** X-ray crystal structure of **1** and **2**, as viewed along the crystallographic  $a$  and  $c$  axes, respectively. The  $[\text{Fe}(\text{qsal})_2]^+$  cationic complexes and the  $[(\text{C}_6\text{F}_3)_2\text{X}]^-$  anionic halogen-bonded networks are shown in grey and blue, respectively, for clarity. Halogen bonding interactions are highlighted as the purple dashed lines with I atoms in purple spheres.

While the topologies of the individual halogen-bonded layer in **1** and **2** are very similar, overall structures feature major differences. The honeycomb-like layers in **1** at 250 K are stacked in a perfectly eclipsed arrangement along the crystallographic  $a$  axis, in contrast to the layers in **2** revealing a staggered arrangement. This eclipsed stacking pattern in **1** gives rise to 1D channels with an average diameter of *ca.* 11.4 Å, defined by cross-pore I...F distances. In these channels, the molecular  $\text{Fe}^{\text{III}}$  complex is placed between layers and features  $\pi\cdots\pi$  and hydrogen bond interactions with the adjacent ones along the crystallographic  $a$  axis, inducing a 1D assembly (Figure S2). The individual chain is then well isolated one another inside the 1D pores formed by the halogen-bonded framework. The average Fe–N and Fe–O distances in the molecular  $[\text{Fe}(\text{qsal})_2]^+$  unit are 2.140(8) and 1.910(7) Å, respectively, unambiguously indicating a high-spin  $\text{Fe}^{\text{III}}$  configuration at 298 K.<sup>11</sup>

The observed molecular arrangement is remarkably different with that in **2**, which consisted of dimerized  $\text{Fe}^{\text{III}}$  complexes and MeCN molecules between the pairwise halogen-bonded layers (Figures 1 and S1). Furthermore, structure of **2** featured extensive intermolecular interactions between the bulky dimerized  $[\text{Fe}]_2$  units and the network, which we attributed to the absence of spin crossover by trapping a high-spin state in the 2–300 K range.<sup>11</sup> Taken together, the structural analysis indicates that the use of different reaction solvent media drastically altered the self-assembly fashion of  $\text{Fe}^{\text{III}}$  molecules during the formation of halogen-bonded supramolecular network and the guest inclusion process. In particular, the eclipsed halogen-bonded honeycomb layer in **1** (Figure 1) separates the densely packed  $[\text{Fe}]_2$  units, and therefore

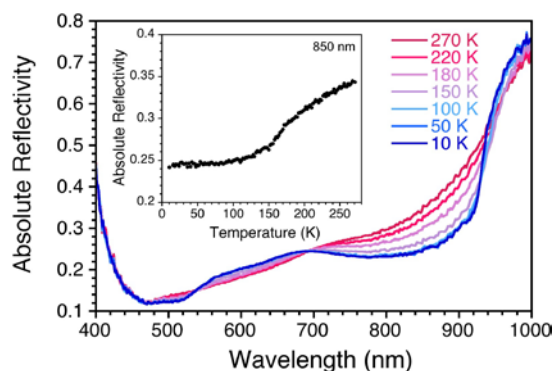


**Figure 2.** Variable-temperature dc magnetic susceptibility ( $\chi = M/H$  per mole of compound) data for **1** (blue sphere) and **2** (green sphere), shown as a plot of  $\chi T$  vs.  $T$ , collected under an applied field of 1 T from 400 to 100 K and 0.1 T below 100 K at a temperature sweep rate of 0.7 K  $\text{min}^{-1}$ .

releases the structural rigidity, relative to that in **2**, which may modulate the magnetic property.

To probe the temperature dependence of the structural changes in **1**, the thermal evolution of the unit-cell parameters was investigated (Figure S3). Upon decreasing temperature, in addition to the linear thermal contraction, a noticeable abnormal jump was observed for  $a$ ,  $c$ , and  $\beta$  parameters between 150 and 180 K, implying the presence of a first-order structural transition, in agreement with the discontinuous change of  $\chi T$ . A full data collection of **1** at 100 K indeed reveals a unit-cell doubling along the crystallographic  $a$  and  $b$  directions (Table S1), displaying the characteristic superstructure Bragg peaks at  $Q = 1/2a^* + 1/2b^*$  (Figures S4 and S5). Consequently, the structure at 100 K features two distinct  $\text{Fe}^{\text{III}}$  centers with the mean Fe–N/Fe–O distances of 2.137(7)/1.907(7) Å for Fe1 and 1.961(8)/1.868(6) Å for Fe2 (Figure S6). The observed distances indicate a high-spin and a low-spin configurations for Fe1 and Fe2, respectively, which are alternating along the  $a$  and  $b$  directions (Figure S7). This is another manifestation of the formation of spin-state concentration waves (SSCW).<sup>15</sup> Accordingly, there are two independent halogen bonding environments, as shown by the I...F bond distances of 3.441(1), 3.514(1), and 3.599(2) Å for the site 1 and 3.493(1), 3.474(1) and 3.658(2) Å for the site 2 (Figure S7). As a result, the 2D honeycomb-like network is not anymore perfectly eclipsed along the  $a$  direction (Figure S8). This proves that the flexibility of the halogen-bonded supramolecular network plays a key role to accommodate the structural changes associated with the spin crossover of Fe2.

In order to investigate the magnetic behaviour of **1**, variable-temperature magnetic susceptibility data were collected on solid samples. The resulting plot of  $\chi T$  vs.  $T$  is shown in Figure 2 and compared with that of **2**, which remains in a high-spin state in the whole temperature range. At 400 K, the  $\chi T$  value for **1** is 4.32  $\text{cm}^3\text{K/mol}$ , which corresponds to the expected one for a high-spin  $S = 5/2$   $\text{Fe}^{\text{III}}$  centre with  $g = 1.99$ . As temperature is decreased,  $\chi T$  product for **1** undergoes a gradual decrease and then an abrupt drop below 200 K, to reach a value of 2.38  $\text{cm}^3\text{K/mol}$  at 20 K. This thermal dependence clearly correlates with the partial spin state conversion observed by x-ray diffraction, related to the formation of SSCW (Figure S9). The observed value at low temperature indicates *ca.* 50% of Fe centres are still in a high-spin state, which agrees well with the structural analysis at 100 K. It is worth noting



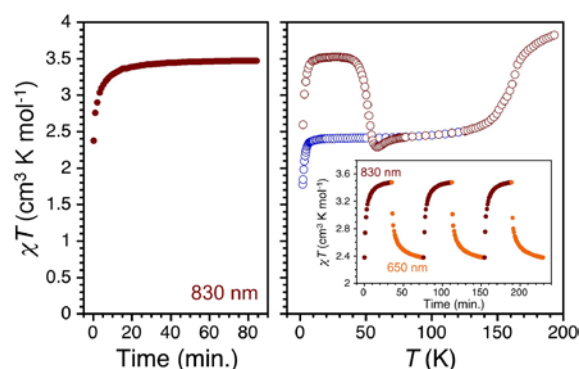
**Figure 3.** Surface optical reflectivity spectra of **1** at selected temperatures upon cooling from 270 to 10 K at 4 K/min. The experiment was performed in the dark except during the spectrum acquisition (spectroscopic white light during 1 s,  $P = 0.8 \text{ mW/cm}^2$ ). Inset: Temperature-dependence of the optical reflectivity at 850 nm of **1** between 270 and 10 K on cooling in the dark (black dots).

that the susceptibility measured between 1.8 and 400 K is completely reversible, without a thermal hysteresis, indicating that **1** and **2** are robust materials and do not degrade up to 400 K (Figure S10).

To further assess the spin crossover process, variable-temperature optical reflectivity was measured on a solid sample of **1** (Figure 3). At 270 K, the spectrum reveals a broad absorption band in the range 400-900 nm. Upon cooling, the spectrum is smoothly modified down to 100 K, and no further significant change was observed between 100 and 10 K. The final spectrum at 10 K exhibits two distinct bands at 500 and 850 nm. To better visualize the thermal evolution of the optical properties, the absolute reflectivity at 850 nm is plotted as a function of temperature (Figure 3 inset), showing a S-shaped gradual decrease. This temperature dependence of the absolute reflectivity corresponds to the thermal evolution of  $\chi T$  product, indicating that the observed surface thermochromism mainly originates from the bulk spin crossover properties.

In order to investigate the potential photosensitivity of **1**, the solid sample was cooled down in the dark at 10 K and then irradiated with selected wavelengths using LED sources for 10 minutes ( $P = 2 \text{ mW/cm}^2$ ). The optical spectra of before and after irradiation were compared in Figure S4. The difference in the absolute reflectivity at 850 nm reveals an efficient photoconversion with LED sources of 850 and 940 nm. This spectral domain corresponds to the absorption band developed at 10 K, which indicates that this band is likely a ligand-to-metal charge transfer absorption of the low-spin  $\text{Fe}^{\text{III}}$  complex. Given the narrow selectivity in wavelength for the photoexcitation, the photoreversibility was also monitored by a successive irradiation at selected wavelength ( $P = 2 \text{ mW/cm}^2$ ), once the sample is fully excited with a 940 nm-irradiation. The photoreverse effect is efficient in a broad range of spectral domain, with an optimum wavelength around 600 nm (Figure S11). Lastly, the photoswitching was investigated using the selected 940 and 625 nm irradiations. As shown in Figure S12, the absolute surface reflectivity at 850 nm shows perfectly reversible cycles, demonstrating a good photoreversibility in **1**.

To analyze the effect of the photosensitivity on bulk magnetic properties, photomagnetic measurements were performed (Figure 4). The solid sample of **1** was cooled to 10 K in the absence of light



**Figure 4.** Dc magnetic susceptibility data for **1** collected under an applied field of 1 T. Left: A plot of  $\chi T$  vs time under an 830 nm-irradiation ( $P = 3.7 \text{ mW/cm}^2$ ) at 10 K. Right: Plots of  $\chi T$  vs temperature at 0.4 K/min before (blue) and after (brown) an 830 nm-irradiation ( $P = 3.7 \text{ mW/cm}^2$ ) at 10 K. Inset: A plot of  $\chi T$  vs time under a succession of 830 nm-irradiation (for 40 min.,  $P = 3.7 \text{ mW/cm}^2$ ) and 650 nm-irradiation (for 40 min.,  $P = 12.6 \text{ mW/cm}^2$ ) at 10 K.

before being exposed to an 830 nm-light ( $P = 3.7 \text{ mW/cm}^2$ ). As shown in the left part of Figure 4, the  $\chi T$  product quickly increases and saturates to  $3.49 \text{ cm}^3\text{K/mol}$  in less than an hour, which indicates the efficient generation of the metastable high-spin state at the bulk level. The partial conversion estimated at ca. 59% is likely due to the intense colour and the shiny nature of the microcrystalline sample of **1**.<sup>‡</sup> Upon increasing temperature, the metastable state is thermally trapped up to 40 K. Around 50 K, the photoinduced high-spin state relaxes to the thermodynamic low-spin state, confirming the presence of the LIESST effect. Note that the observed photoexcitation efficiency (at least 59%) is far superior to that reported for any  $[\text{Fe}^{\text{III}}(\text{qsal})_2]^+$  units showing a LIESST effect, which feature a weak photoconversion of less than 10%.<sup>5,6</sup> Right after the relaxation around 53 K, the  $\chi T$  product shows a drop lower than the value for the non-irradiated sample at the same temperature. The  $\chi T$  value then gradually increases to recover at 80 K the data before irradiation upon further heating. To further analyse this feature, we measured the magnetic susceptibility at a fixed temperature of 53 K after photoexcitation (Figure S13). Nevertheless, the  $\chi T$  product remains at a quasi-constant value after an hour, indicating that a combined effect of the low thermal energy and the limited structural flexibility of the compound likely prevents to access to the fully low-spin state.

The efficient photoexcitation and the thermal trapping of the metastable phase enabled us to test the photoreverse effect on a bulk sample of **1** by photomagnetic measurements. At 10 K, the crystalline sample was first exposed during 40 minutes to a 830 nm-light ( $P = 3.7 \text{ mW/cm}^2$ ), and then a 650 nm-light ( $P = 12.6 \text{ mW/cm}^2$ ) was used for another 40 minutes. During this sequence of irradiations, the magnetic susceptibility was measured and the resulting plot of  $\chi T$  vs. time is shown in inset of Figure 4. Remarkably the data shows excellent reversibility and reproducibility during three successive cycles of irradiations, highlighting the efficient solid-state photoswitching of spin-state in **1**. To our knowledge, this is the first example of a reversible photoconversion in compounds containing  $[\text{Fe}^{\text{III}}(\text{qsal})_2]^+$  units.

To have an insight into the unprecedented reversible photoswitching effect in **1**, we have compared the structural differences with the few literature examples containing the  $[\text{Fe}^{\text{III}}(\text{qsal})_2]^+$  unit and exhibiting a photomagnetic effect.<sup>5</sup> For the



mean Fe–N and Fe–O bond lengths between high-spin and low-spin states, the reported compounds show a contraction by 0.150–0.167 and 0.003–0.018 Å, respectively. In contrast, the corresponding shortening in **1** amounts to 0.187 and 0.043 Å, and accordingly is much larger than that observed in the reported examples. As such, the more significant structural change in the photoexcited high-spin state is anticipated for **1**, which may help to avoid the fast relaxation of the metastable phase. Such a structural distortion effect has been previously proposed in other Fe<sup>III</sup> systems as well.<sup>4,7</sup> Furthermore, the rare observation of the LIESST effect in known [Fe<sup>III</sup>(qsal)<sub>2</sub>]<sup>+</sup> complexes, in conjunction with its low efficiency, has been attributed to an overlap of the LIESST and the reverse LIESST bands.<sup>5</sup> In contrast, compound **1** shows a distinct difference between the spectral bands for the photoexcitation and the reverse photoconversion, which likely stems from a combined effect of the substantial structural distortion and the significant intermolecular interactions with the host supramolecular network. Taken together, the optimized structural confinement and complexity provided by halogen-bonded network play a key role to favourably access and trap the metastable phase, which enables the reversible photoswitching.

The foregoing results describe a drastic solvent-induced modification in the organization of a guest Fe<sup>III</sup> complex inside a halogen-bonded supramolecular network. In particular, the induced structural relief triggers the breathing of the host network to accommodate the spin crossover and the reversible photoswitching of the spin-state of the guest molecule. Work is underway to further control the inclusion process of guest functional molecules into the host framework, for example by increasing the pore size of the halogen-bonded networks using larger neutral tris-iodinated spacers.<sup>16</sup>

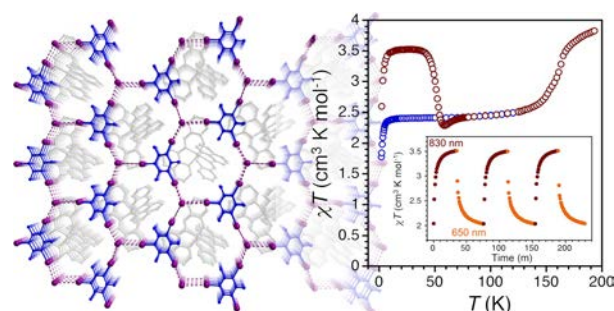
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## Notes and references

‡ The conversion percentage is calculated using a following equation:  $(\chi T_{400K} - \chi T_{20K}) / (\chi T_{20K, hv} - \chi T_{20K}) \times 100$ .

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