



HAL
open science

Numerical reconstruction of the first band(s) in an inverse Hill's problem

Athmane Bakhta, Virginie Ehrlacher, David Gontier

► **To cite this version:**

Athmane Bakhta, Virginie Ehrlacher, David Gontier. Numerical reconstruction of the first band(s) in an inverse Hill's problem. ESAIM: Control, Optimisation and Calculus of Variations, 2020, 10.1051/cocv/2019031 . hal-01591133

HAL Id: hal-01591133

<https://hal.science/hal-01591133>

Submitted on 20 Sep 2017

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

Numerical reconstruction of the first band(s) in an inverse Hill's problem.

Athmane Bakhta* Virginie Ehrlacher† David Gontier‡

Abstract

This paper concerns an inverse band structure problem for one dimensional periodic Schrödinger operators (Hill's operators). Our goal is to find a potential for the Hill's operator in order to reproduce as best as possible some given target bands, which may not be realisable. We recast the problem as an optimisation problem, and prove that this problem is well-posed when considering singular potentials (Borel measures). We then propose different algorithms to tackle the problem numerically.

1 Introduction

The aim of this article is to present new considerations on an inverse band structure problem for periodic one-dimensional Schrödinger operators, also called Hill's operators. A Hill operator is a self-adjoint, bounded from below operator of the form $A^V := -\frac{d^2}{dx^2} + V$, acting on $L^2(\mathbb{R})$, and where V is a periodic real-valued potential. Its spectrum is composed of a reunion of intervals, which can be characterised using Bloch-Floquet theory as the reunion of the spectra of a family of self-adjoint compact resolvent operators A_q^V , indexed by an element $q \in \mathbb{R}$ called the *quasi-momentum* or *k-point* (see [22, Chapter XIII] and Section 2.1). The m^{th} band function associated to a periodic potential is the function which maps $q \in \mathbb{R}$ to the m^{th} lowest eigenvalue of A_q^V . The properties of these band functions are well-known, especially in the one-dimensional case (see e.g. [22, Chapter XIII]).

The inverse band structure problem is an interesting mathematical question of practical interest, which can be roughly formulated as follows: *is it possible to find a potential V so that its first bands are close to some target functions?*

A wide mathematical literature answers the question when the target functions are indeed the bands of some Hill's operator, corresponding to some V_{ref} . In this case, we need to *recover* a potential V that reproduces the bands of V_{ref} . We refer to [5, 6, 7, 21, 9, 24] for the case when V_{ref} is a regular potential, and to [12, 14, 15, 13, 16] when V_{ref} is singular (see also the review [18]). The main ideas of the previous references are as follows. First, the band structure of a Hill's operator can be seen as the transformation of an analytic function. In particular, the knowledge of any band on an open set is enough to recover *theoretically* the whole band structure. A potential is then reconstructed from the high energy asymptotics of the bands.

The previous methods use the knowledge of the behaviour of the high energy bands, and therefore are unsuitable for practical purpose (material design) since we usually have no accurate and numerically stable information about these high energy bands. Moreover, in practice, only the low energy bands are usually of interest. The fact that there exists no explicit characterisation of the set of the first band functions associated to a given admissible set of periodic potentials is an additional numerical difficulty. For applications, it is therefore interesting to know how to construct a potential such that only its first bands are close to some given target functions, which may not be realisable (for instance not analytic). In this present work, we therefore adopt a different point of view, which, up to our best knowledge, has not been studied: we recast the inverse problem as an optimisation problem.

*Ecole des Ponts ParisTech, athmane.bakhta@cermics.enpc.fr

†Ecole des Ponts ParisTech & INRIA, virginie.ehrlacher@enpc.fr

‡Université Paris-Dauphine, gontier@ceremade.dauphine.fr

The outline of the paper is as follows. In Section 2, we recall basic properties about Hill's operators with singular potentials. and we state our main result (Theorem 2.3). Its proof is given in Section 3. Finally, we present in Section 4 some numerical tests and propose an adaptive optimisation algorithm, which is observed to converge faster than the standard one. This adaptive algorithm relies on the use of an a posteriori error estimator for discretised eigenvalue problems, whose computation is detailed in the Appendix.

2 Spectral decomposition of periodic Schrödinger operators, and main results

In this section, we recall some properties of Hill's operators with singular potentials. Elementary notions on the Bloch-Floquet transform [22] are gathered in Section 2.1. The spectral decomposition of one-dimensional periodic Schrödinger operators with singular potentials is detailed in Section 2.2, building on the results of [17, 11, 10, 20, 4]. We state our main results in Section 2.3.

2.1 Bloch-Floquet transform

We need some notation. Let \mathcal{D}' denotes the Schwartz space of complex-valued distributions, and let $\mathcal{D}'_{\text{per}} \subset \mathcal{D}'$ be the space of distributions that are 2π -periodic. In the sequel, the unit cell is $\Gamma := [-\pi, \pi)$, and the reciprocal unit cell (or Brillouin zone) is $\Gamma^* := [-1/2, 1/2]$. For $u \in \mathcal{D}'_{\text{per}}$ and $k \in \mathbb{Z}$, the k^{th} normalised Fourier coefficient of u is denoted by $\widehat{u}(k)$. For $s \in \mathbb{R}$, we denote by

$$H_{\text{per}}^s := \left\{ u \in \mathcal{D}'_{\text{per}}, \|u\|_{H_{\text{per}}^s}^2 := \sum_{k \in \mathbb{Z}} (1 + |k|^2)^s |\widehat{u}(k)|^2 < +\infty \right\}$$

the complex-valued periodic Sobolev space, which is a Hilbert space when endowed with its natural inner product. We write $H_{\text{per},r}^s$ for the *real-valued* periodic Sobolev space, i.e.

$$H_{\text{per},r}^s := \left\{ u \in H_{\text{per}}^s, \quad \forall k \in \mathbb{Z}, \widehat{u}(-k) = \overline{\widehat{u}(k)} \right\}.$$

We also let $L_{\text{per}}^2 := H_{\text{per}}^{s=0}$. From our normalisation, it holds that

$$\forall v, w \in L_{\text{per}}^2, \langle v, w \rangle_{L_{\text{per}}^2} = \int_{\Gamma} \bar{v} w \quad \text{and} \quad \forall v, w \in H_{\text{per}}^1, \langle v, w \rangle_{H_{\text{per}}^1} = \int_{\Gamma} \frac{d\bar{v}}{dx} \frac{dw}{dx} + \int_{\Gamma} \bar{v} w.$$

Lastly, we denote by C_{per}^0 the space of 2π -periodic continuous functions, and by C_c^∞ the space of C^∞ functions over \mathbb{R} , with compact support.

To introduce the Bloch-Floquet transform, we let $\mathcal{H} := L^2(\Gamma^*, L_{\text{per}}^2)$. For any element $f \in \mathcal{H}$, we denote by $f_q(x)$ its value at the point $(q, x) \in \Gamma^* \times \Gamma$. The space \mathcal{H} is an Hilbert space when endowed with its inner product

$$\forall f, g \in \mathcal{H}, \quad \langle f, g \rangle_{\mathcal{H}} := \int_{\Gamma^*} \int_{\Gamma} \overline{f_q(x)} g_q(x) dx dq.$$

The Bloch-Floquet transform is the map $\mathcal{B} : L^2(\mathbb{R}) \rightarrow \mathcal{H}$ defined, for smooth functions $\varphi \in C_c^\infty(\mathbb{R})$, by

$$\phi_q(x) := (\mathcal{B}\varphi)_q(x) := \sum_{R \in \mathbb{Z}} \varphi(x + R) e^{-iq(R+x)}.$$

It is an isometry from $L^2(\mathbb{R})$ to \mathcal{H} , whose inverse is given by

$$(\mathcal{B}^{-1}\phi)(x) := \int_{\Gamma^*} \phi_q(x) e^{iqx} dq = \varphi(x).$$

The Bloch theorem states that if A is a self-adjoint operator on $L^2(\mathbb{R})$ with domain $D(A)$ that commutes with \mathbb{Z} -translations, then $\mathcal{B}A\mathcal{B}^{-1}$ is diagonal in the q -variable. More precisely, there exists a unique family of self-adjoint operators $(A_q)_{q \in \Gamma^*}$ on L_{per}^2 such that for all $\varphi \in L^2(\mathbb{R}) \cap D(A)$,

$$(A\varphi)(x) = \int_{\Gamma^*} (A_q \phi_q)(x) dq.$$

In this case, we write

$$A = \int_{\Gamma^*}^{\oplus} A_q dq.$$

2.2 Hill's operators with singular potentials

Giving a rigorous mathematical sense to a Hill's operator of the form $-\frac{d^2}{dx^2} + V$ on $L^2(\mathbb{R})$, when the potential V is singular is not an obvious task. In the present paper, we consider $V \in H_{\text{per},r}^{-1}$, which is a case that was first tackled in [17] (see also [11, 4, 10, 20] for recent results).

The results which are gathered in this section are direct corollaries of results which were proved in these earlier works, particularly in [11].

Proposition 2.1. *[Theorem 2.1 and Lemma 3.2 of [11]] For all $V \in H_{\text{per},r}^{-1}$, there exists $\sigma_V \in L_{\text{per}}^2$ and $\kappa_V \in \mathbb{R}$ such that*

$$V = \sigma'_V + \kappa_V \text{ in } \mathcal{D}'_{\text{per}}. \quad (2.1)$$

Moreover, if $a^V : H^1(\mathbb{R}) \times H^1(\mathbb{R}) \rightarrow \mathbb{C}$ is the sesquilinear form defined by

$$\forall v, w \in H^1(\mathbb{R}), \quad a^V(v, w) = \int_{\mathbb{R}} \frac{d\bar{v}}{dx} \frac{dw}{dx} + \int_{\mathbb{R}} \kappa_V \bar{v} w - \int_{\mathbb{R}} \sigma_V \left(\frac{d\bar{v}}{dx} w + \bar{v} \frac{dw}{dx} \right), \quad (2.2)$$

then a^V is a symmetric, continuous sesquilinear form on $H^1(\mathbb{R}) \times H^1(\mathbb{R})$, which is closed and bounded from below. Besides, a^V is independent of the choice of $\sigma_V \in L_{\text{per}}^2$ and $\kappa_V \in \mathbb{R}$ satisfying (2.1).

Remark 2.2. *The expression (2.2) makes sense whenever $v, w \in H^1(\mathbb{R})$. This can be easily seen with the Cauchy-Schwarz inequality, and the embedding $H^1(\mathbb{R}) \hookrightarrow L^\infty(\mathbb{R})$. It is not obvious how to extend this result to higher dimension.*

A direct consequence of Proposition 2.1 is that one can consider the Friedrichs operator on $L^2(\mathbb{R})$ associated to a^V , which is denoted by A^V in the sequel. The operator A^V is thus a densely defined, self-adjoint, bounded from below operator on $L^2(\mathbb{R})$, with form domain $H^1(\mathbb{R})$ and whose domain is dense in $L^2(\mathbb{R})$. Formally, it holds that

$$A^V = -\frac{\partial^2}{\partial x^2} + V.$$

The spectral properties of the operator A^V can be studied (like in the case of regular potentials) using Bloch-Floquet theory.

The previous result, together with Bloch-Floquet theory, allows to study the operator A^V via its Bloch fibers $(A_q^V)_{q \in \Gamma^*}$. For $q \in \Gamma^*$, it holds that A_q^V is the self-adjoint extension of the operator

$$\left| -i \frac{d}{dx} + q \right|^2 + V.$$

It holds that A_q^V is a bounded from below self-adjoint operator acting on L_{per}^2 , whose form domain is H_{per}^1 , and with associated quadratic form a_q^V , defined by (recall that H_{per}^1 is an algebra)

$$\forall v, w \in H_{\text{per}}^1, \quad a_q^V(v, w) := \int_{\Gamma} \left[\overline{\left(-i \frac{d}{dx} + q \right) v} \left(-i \frac{d}{dx} + q \right) w \right] + \langle V, \bar{v} w \rangle_{H_{\text{per}}^{-1}, H_{\text{per}}^1}. \quad (2.3)$$

In other words, we have

$$A^V = \int_{\Gamma^*}^{\oplus} A_q^V dq.$$

The fact that L_{per}^2 is compactly embedded in H_{per}^1 implies that A_q^V is compact-resolvent. As a consequence, there exists a non-decreasing sequence of real eigenvalues $(\varepsilon_{q,m}^V)_{m \in \mathbb{N}^*}$ going to $+\infty$ and a corresponding orthonormal basis $(u_{q,m}^V)_{m \in \mathbb{N}^*}$ of L_{per}^2 such that

$$\forall m \in \mathbb{N}^*, \quad A_q^V u_{q,m}^V = \varepsilon_{q,m}^V u_{q,m}^V. \quad (2.4)$$

The map $\Gamma^* \ni q \mapsto \varepsilon_{q,m}^V$ is called the m^{th} band. Since the potential V is real-valued, it holds that $A_{-q}^V = \overline{A_q^V}$, so that $\varepsilon_{-q,m}^V = \varepsilon_{q,m}^V$ for all $q \in \Gamma^*$ and $m \in \mathbb{N}^*$. This implies that it is enough to study the bands on $[0, 1/2]$. Actually, we have

$$\sigma(A^V) = \bigcup_{q \in [0, 1/2]} \bigcup_{m \in \mathbb{N}^*} \{\varepsilon_{q,m}^V\}.$$

In the sequel, we mainly focus on the first band. We write $\varepsilon_q^V := \varepsilon_{q,1}^V$ for the sake of clarity. Thanks to the knowledge of the form domain of A_q^V , we know that

$$\varepsilon_q^V := \min_{\substack{v \in H_{\text{per}}^1 \\ \|v\|_{L_{\text{per}}^2} = 1}} a_q^V(v, v). \quad (2.5)$$

This characterisation will be the key to our proof. When the potential V is smooth (say $V \in L_{\text{per}}^2$), then the map $\Gamma^* \ni q \mapsto \varepsilon_{q,m}^V$ is analytic on $(-1/2, 1/2)$. Besides, it is increasing on $[0, 1/2]$ if m is odd, and decreasing if m is even (see e.g. [22, Chapter XIII]).

2.3 Main results

The goal of this article is to find a potential V so that the bands of the corresponding Hill's operator are close to some given target functions. In order to do so, we recast the problem as a minimisation one, of the form

$$V^* \in \arg \min_{V \in \mathcal{V}} \mathcal{J}(V).$$

Unfortunately, we were not able to consider the full setting where the minimisation set \mathcal{V} is the whole set $H_{\text{per},r}^{-1}$. The problem was that we were unable to control the negative part of V . To bypass this difficulty, we chose to work with potentials that are bounded from below. Such a distribution is necessary a measure (see e.g. [19]). Hence measure-valued potentials provide a natural setting for band reconstruction. We recall here some basic properties about measures.

We denote by $\mathcal{M}_{\text{per}}^+$ the space of non-negative 2π -periodic regular Borel measures on \mathbb{R} , in the sense that for all $\nu \in \mathcal{M}_{\text{per}}^+$, and all Borel set $S \in \mathcal{B}(\mathbb{R})$, it holds that $\nu(S) = \nu(S + 2\pi) \geq 0$, and $\nu(\Gamma) < \infty$. For all $\varepsilon > 0$, from the Sobolev embedding $H_{\text{per}}^{1/2+\varepsilon} \hookrightarrow C_{\text{per}}^0$, we deduce that $\mathcal{M}_{\text{per}}^+ \hookrightarrow H_{\text{per}}^{-1/2-\varepsilon} \hookrightarrow H_{\text{per}}^{-1}$, where the last embedding is compact. For $\nu \in \mathcal{M}_{\text{per}}^+$, we denote by $V_\nu \in H_{\text{per},r}^{-1}$ the unique corresponding potential, which is defined by duality through the relation:

$$\forall \phi \in H_{\text{per}}^1, \quad \int_{\Gamma} \phi d\nu = \langle V_\nu, \phi \rangle_{H_{\text{per}}^{-1}, H_{\text{per}}^1}.$$

For $B \in \mathbb{R}$, we define the set of B -bounded from below potentials

$$\mathcal{V}_B := \{V \in H_{\text{per},r}^{-1} \mid \exists \nu \in \mathcal{M}_{\text{per}}^+, \quad V = V_\nu - B\} \subset H_{\text{per},r}^{-1}.$$

This will be our minimisation space for our optimisation problem. Note that $\mathcal{V}_{B_1} \subset \mathcal{V}_{B_2}$ for $B_1 \geq B_2$.

We now introduce the functional \mathcal{J} to minimise. First, we introduce the set \mathcal{T} of allowed target functions:

$$\mathcal{T} := \{b \in C^0(\Gamma^*), \quad b \text{ is even and } b \text{ is increasing on } [0, 1/2]\}. \quad (2.6)$$

Of course, for all $V \in H_{\text{per},r}^{-1}$, it holds that $\Gamma^* \ni q \mapsto \varepsilon_q^V \in \mathcal{T}$. Finally, in order to quantify the quality of reconstruction of a band $b \in \mathcal{T}$, we introduce the error functional $\mathcal{J}_b : H_{\text{per},r}^{-1} \rightarrow \mathbb{R}$ defined by

$$\forall V \in H_{\text{per},r}^{-1}, \quad \mathcal{J}_b(V) := \frac{1}{2} \int_{\Gamma^*} |b(q) - \varepsilon_q^V|^2 dq = \int_0^{1/2} |b(q) - \varepsilon_q^V|^2 dq. \quad (2.7)$$

The main result of the present paper is the following.

Theorem 2.3. *Let $b \in \mathcal{T}$, and denote by $b^* := \int_{\Gamma^*} b(q) dq \in \mathbb{R}$. Then, for all $B > 1/4 - b^*$, there exists a solution $V_{b,B} \in \mathcal{V}_B$ to the minimisation problem*

$$V_{b,B} \in \arg \min_{V \in \mathcal{V}_B} \mathcal{J}_b(V). \quad (2.8)$$

The proof of Theorem 2.3 relies on the following proposition, which is central to our analysis. Both the proofs of Theorem 2.3 and Proposition 2.4 are provided in the next section.

Proposition 2.4. *Let $B \in \mathbb{R}$ and let $(V_n)_{n \in \mathbb{N}^*} \subset \mathcal{V}_B$. For all $n \in \mathbb{N}^*$, let $\nu_n \in \mathcal{M}_{\text{per}}^+$ such that $V_n := V_{\nu_n} - B$. Let us assume that the sequence $(\varepsilon_0^{V_n})_{n \in \mathbb{N}^*}$ is bounded and such that $\nu_n(\Gamma) \xrightarrow{n \rightarrow +\infty} +\infty$. Then, up to a subsequence (still denoted n), the functions $q \mapsto \varepsilon_q^{V_n}$ converge uniformly to a constant function $\varepsilon \in \mathbb{R}$, with $\varepsilon \geq \frac{1}{4} - B$. In other words, there is $\varepsilon \geq \frac{1}{4} - B$ such that*

$$\max_{q \in [0, 1/2]} |\varepsilon_q^{V_n} - \varepsilon| \xrightarrow{n \rightarrow \infty} 0. \quad (2.9)$$

Conversely, for all $\varepsilon \geq \frac{1}{4} - B$, there is a sequence $(V_n)_{n \in \mathbb{N}^} \subset \mathcal{V}_B$ such that (2.9) holds.*

This result implies that the first band of the sequence of operators $(A^{V_n})_{n \in \mathbb{N}^*}$, where $(V_n)_{n \in \mathbb{N}^*}$ satisfies the assumptions of Proposition 2.4, becomes flat.

Remark 2.5. *Here we have a sequence of first bands $(\varepsilon_q^{V_n})_{n \in \mathbb{N}^*}$ that converges uniformly to a constant function. However, as the first band of any Hill's operator must be increasing and analytic, the limit is not the first band of a Hill's operator.*

3 Proof of Theorem 2.3 and Proposition 2.4

3.1 Preliminary lemmas

We first prove some intermediate useful lemmas before giving the proof of Proposition 2.4 and Theorem 2.3. We start by recording a spectral convergence result.

Proposition 3.1. *[Theorem 4.1 [11]] Let $(V_n)_{n \in \mathbb{N}^*} \subset H_{\text{per}, r}^{-1}$ be a sequence such that $(V_n)_{n \in \mathbb{N}^*}$ converges strongly in H_{per}^{-1} to some $V \in H_{\text{per}, r}^{-1}$. Then,*

$$\forall m \in \mathbb{N}^*, \quad \max_{q \in [0, 1/2]} |\varepsilon_{q, m}^{V_n} - \varepsilon_{q, m}^V| \xrightarrow{n \rightarrow \infty} 0.$$

In our case, since we are working with potentials that are measures, we deduce the following result.

Proposition 3.2. *Let $B \in \mathbb{R}$ and $(V_n)_{n \in \mathbb{N}^*} \subset \mathcal{V}_B$ be a bounded sequence, in the sense*

$$\sup_{n \in \mathbb{N}} \langle V_n, \mathbb{1}_\Gamma \rangle_{H_{\text{per}}^{-1}, H_{\text{per}}^1} < \infty.$$

For all $n \in \mathbb{N}^$, let $\nu_n \in \mathcal{M}_{\text{per}}^+$ such that $V_n = V_{\nu_n} - B$. Then, there exists $\nu \in \mathcal{M}_{\text{per}}^+$ such that, up to a subsequence (still denoted n), $(\nu_n)_{n \in \mathbb{N}^*}$ converges weakly- $*$ to ν in \mathcal{M}_{per} , and $(V_n)_{n \in \mathbb{N}^*}$ converges strongly in H_{per}^{-1} to $V := V_\nu - B \in \mathcal{V}_B$. Moreover, it holds that*

$$\forall m \in \mathbb{N}^*, \quad \max_{q \in [0, 1/2]} |\varepsilon_{q, m}^{V_n} - \varepsilon_{q, m}^V| \xrightarrow{n \rightarrow \infty} 0.$$

Proof. The fact that we can extract from the bounded sequence $(\nu_n)_{n \in \mathbb{N}^*}$ a weakly- $*$ convergent sequence in $\mathcal{M}_{\text{per}}^+$ is the Prokhorov's theorem applied in the torus Γ^* . The second part comes from the compact embedding $\mathcal{M}_{\text{per}} \hookrightarrow H_{\text{per}}^{-1}$. The final part is the direct application of Proposition 3.1. \square

Remark 3.3. *This proposition explains our choice to consider measure-valued potentials. Note that a similar result does not hold in the L_{per}^1 setting for instance.*

We now give a lemma which is standard in the case of regular potentials V (see [8]).

Lemma 3.4. *Let $V \in \mathcal{V}_B$ for some $B \in \mathbb{R}$. The first eigenvector $u_{q=0}^V \in H_{\text{per}}^1$ of $A_{q=0}^V$ is unique up to a global phase. It can be chosen real-valued and positive.*

Proof. We use the min-max principle (2.5), and the fact that, for $u \in H_{\text{per}}^1$, the following holds

$$\left| \frac{d}{dx} |u| \right| \leq \left| \frac{d}{dx} u \right| \quad \text{a.e.}$$

We see that if u is an eigenvector corresponding to the first eigenvalue, then so is $|u|$. We now consider a non-negative eigenvector $u \geq 0$, and prove that it is positive. The usual argument is Harnack's inequality. However, it is a priori unclear that it works in our singular setting. To prove it, we write $V = V_\nu - B$ for $\nu \in \mathcal{M}_{\text{per}}^+$, and consider the repartition function F_ν of ν : $F_\nu(x) := \nu((0, x])$. This function is not periodic, but the function $f_\nu(x) := F_\nu(x) - \nu(\Gamma) \frac{x}{|\Gamma|}$ is. Since F_ν is an non decreasing, right-continuous function, we deduce that $f_\nu \in L_{\text{per}}^\infty$. Moreover, it holds, in the H_{per}^{-1} sense, that $f'_\nu = V_\nu - |\Gamma|^{-1}\nu(\Gamma) = V + B - |\Gamma|^{-1}\nu(\Gamma)$. As a result, we see that u is solution to the minimisation problem

$$u \in \arg \min_{\substack{v \in H_{\text{per}}^1 \\ \|v\|_{L_{\text{per}}^2} = 1}} \left\{ \int_{\Gamma} \left| \frac{dv}{dx} \right|^2 + \left(\frac{\nu(\Gamma)}{|\Gamma|} - B \right) - 2 \int_{\Gamma} f_\nu \left(v \frac{dv}{dx} \right) \right\}.$$

There exists $\lambda \in \mathbb{R}$ so that the corresponding Euler-Lagrange equations can be written in the weak-form:

$$\operatorname{div} F(x, u, u') + G(x, u, u') = 0,$$

with

$$F(x, u, p) = p - f_\nu u \quad \text{and} \quad B(x, u, p) = f_\nu p + \lambda u.$$

We are now in the settings of [23, Theorem 1.1], and we deduce that $u > 0$. The rest of the proof is standard. \square

3.2 Proof of Proposition 2.4

We now prove Proposition 2.4. Let $B \in \mathbb{R}$ and let $V_n = V_{\nu_n} - B \in \mathcal{V}_B$ with $\nu_n \in \mathcal{M}_{\text{per}}^+$, be a sequence such that the sequence $(\varepsilon_{q=0}^{V_n})_{n \in \mathbb{N}^*}$ is bounded and $\nu_n(\Gamma)$ goes to $+\infty$. Since $(\varepsilon_0^{V_n})_{n \in \mathbb{N}^*}$ is bounded, then up to a subsequence (still denoted by n), there exists $\varepsilon \in \mathbb{R}$ such that $\varepsilon_0^{V_n}$ converges to ε . Our goal is to prove that the convergence also holds uniformly in $q \in \Gamma^*$.

Let $u_0^{V_n} \in H_{\text{per}}^1$ be the L_{per}^2 -normalised positive eigenvector of $A_0^{V_n}$ associated to the eigenvalue $\varepsilon_0^{V_n}$ (see Lemma 3.4). We denote by $\alpha_n := \min_{x \in \Gamma} u_0^{V_n}(x) > 0$. Let us first prove that the following convergences hold:

$$\alpha_n \int_{\Gamma} u_0^{V_n} d\nu_n \xrightarrow{n \rightarrow +\infty} 0 \quad \text{and} \quad \alpha_n^2 \nu_n(\Gamma) \xrightarrow{n \rightarrow +\infty} 0. \quad (3.1)$$

From the equality

$$\int_{\Gamma} \left| \frac{d}{dx} (u_0^{V_n}) \right|^2 + \int_{\Gamma} |u_0^{V_n}|^2 d\nu_n = \varepsilon_0^{V_n} + B,$$

we get

$$\alpha_n^2 \nu_n(\Gamma) \leq \alpha_n \int_{\Gamma} u_0^{V_n} d\nu_n \leq \int_{\Gamma} |u_0^{V_n}|^2 d\nu_n \leq \varepsilon_0^{V_n} + B. \quad (3.2)$$

As the right-hand side is bounded, and $\nu_n(\Gamma) \rightarrow +\infty$ by hypothesis, this implies $\alpha_n \rightarrow 0$. Moreover, we have

$$0 \leq \int_{\Gamma} u_0^{V_n} d\nu_n = a_0^{V_n}(u_0^{V_n}, \mathbf{1}_{\Gamma}) + B \int_{\Gamma} u_0^{V_n} = (\varepsilon_0^{V_n} + B) \int_{\Gamma} u_0^{V_n} \leq (\varepsilon_0^{V_n} + B) |\Gamma|^{1/2},$$

where we used the Cauchy-Schwarz inequality for the last part. As a result, we deduce that the sequence $(\int_{\Gamma} u_0^{V_n} d\nu_n)_{n \in \mathbb{N}^*}$ is bounded. The first convergence of (3.1) follows. The second convergence is a consequence of the first inequality in (3.2).

Let $x_n \in \Gamma = [0, 2\pi)$ be such that $\alpha_n = u_0^{V_n}(x_n)$. The fact that $\alpha_n \rightarrow 0$ implies that $l_n := \|u_0^{V_n}(x_n + \cdot) - \alpha_n\|_{L_{\text{per}}^2}^2 \rightarrow 1$ and we can thus define for n large enough

$$v_n := \frac{u_0^{V_n}(x_n + \cdot) - \alpha_n}{\|u_0^{V_n}(x_n + \cdot) - \alpha_n\|_{L_{\text{per}}^2}}.$$

It holds that $v_n \in H_{\text{per}}^1$, $\|v_n\|_{L_{\text{per}}^2} = 1$. Besides, it holds that $v_n(0) = 0$. For $q \in \Gamma^*$, we introduce the function $v_{q,n}$ defined by:

$$\forall x \in \mathbb{R}, \quad v_{q,n}(x) := v_n(x)e^{-iq[x]}, \quad \text{where we set } [x] := x \bmod 2\pi.$$

Thanks to the equality $v_n(0) = 0$, it holds that $v_{q,n} \in H_{\text{per}}^1$, and that $\|v_{q,n}\|_{L_{\text{per}}^2} = 1$. This function is therefore a valid test function for our min-max principle¹.

From the min-max principle (2.5) and the expression (2.3), we obtain

$$\begin{aligned} B + \varepsilon_q^{V_n} &\leq B + a_q^{V_n}(v_{q,n}, v_{q,n}) \\ &= \int_{\Gamma} \left| \left(-i \frac{d}{dx} + q \right) v_{q,n} \right|^2 + \int_{\Gamma} |v_{q,n}|^2 d\nu_n = \int_{\Gamma} \left| \frac{dv_n}{dx} \right|^2 + \int_{\Gamma} |v_n|^2 d\nu_n \\ &= \frac{1}{l_n} \left(\int_{\Gamma} \left| \frac{d}{dx} \left(u_0^{V_n}(x_n + \cdot) \right) \right|^2 + \int_{\Gamma} |u_0^{V_n}(x_n + \cdot) - \alpha_n|^2 d\nu_n \right) \\ &= \frac{1}{l_n} \left(\int_{\Gamma} \left| \frac{d}{dx} \left(u_0^{V_n} \right) \right|^2 + \int_{\Gamma} |u_0^{V_n}|^2 d\nu_n - 2\alpha_n \int_{\Gamma} u_0^{V_n} d\nu_n + \alpha_n^2 \nu_n(\Gamma) \right) \\ &= \frac{1}{l_n} \left(B + \varepsilon_0^{V_n} - 2\alpha_n \int_{\Gamma} u_0^{V_n} d\nu_n + \alpha_n^2 \nu_n(\Gamma) \right). \end{aligned}$$

We infer from these inequalities, and from (3.1) that

$$0 \leq \max_{q \in \Gamma^*} \left| \varepsilon_q^{V_n} - \varepsilon_0^{V_n} \right| \leq \left(B + \varepsilon_0^{V_n} \right) \left(\frac{1}{l_n} - 1 \right) + \frac{1}{l_n} \left(-2\alpha_n \int_{\Gamma} u_0^{V_n} d\nu_n + \alpha_n^2 \nu_n(\Gamma) \right) \xrightarrow{n \rightarrow +\infty} 0.$$

This already proves the convergence (2.9).

To see that $\varepsilon \geq \frac{1}{4} - B$, we write, for $V = V_{\nu} - B$ with $\nu \in \mathcal{M}_{\text{per}}^+$ that

$$\forall q \in [-1/2, 1/2], \quad A_q^V = \left| -i \frac{d}{dx} + q \right|^2 + V_{\nu} - B \geq \left| -i \frac{d}{dx} + q \right|^2 - B \geq q^2 - B,$$

where we used the fact that the lowest eigenvalue of $\left| -i \frac{d}{dx} + q \right|^2$ is q^2 for $q \in [-1/2, 1/2]$ (this can be seen with the Fourier representation of the operator). As a consequence, for $q = \frac{1}{2}$, we obtain that for all $V \in \mathcal{V}_B$, $\varepsilon_{q=1/2}^V \geq \frac{1}{4} - B$. The result follows.

To prove the converse, we exhibit an explicit sequence of measures $(\nu_n)_{n \in \mathbb{N}^*} \subset \mathcal{M}_{\text{per}}^+$ such that $\varepsilon_q^{V_{\nu_n}} \rightarrow \frac{1}{4}$. The general result will follow by taking sequences of the form $V_n = V_{\nu_n} + \left(\varepsilon - \frac{1}{4} \right) - B$. We denote by δ_x the Dirac mass at $x \in \mathbb{R}$, and consider, for $\lambda > 0$, the measure

$$\nu_{\lambda} := \lambda \sum_{k \in \mathbb{Z}} \delta_{2\pi k} \in \mathcal{M}_{\text{per}}^+. \quad (3.3)$$

From the first part of the Proposition, it is enough to check the convergence for $q = 0$. We are looking for a solution to (we denote by $\omega_{\lambda}^2 := \varepsilon_0^{V_{\nu_{\lambda}}} \geq 0$ for simplicity)

$$-u'' + \lambda \delta_0 u(0) = \omega_{\lambda}^2 u, \quad u \geq 0, \quad u(2\pi) = u(0). \quad (3.4)$$

On $(0, 2\pi)$, u satisfies the elliptic equation $-u'' = \omega_{\lambda}^2 u$, hence is of the form

$$u(x) = Ce^{i\omega_{\lambda}x} + De^{-i\omega_{\lambda}x},$$

for some $C, D \in \mathbb{R}$. The continuity of u at 2π implies $Ce^{2i\pi\omega_{\lambda}} + De^{-2i\pi\omega_{\lambda}} = C + D$. Moreover, integrating (3.4) between 0^- and 0^+ leads to the jump of the derivative $-u'(0) + u'(2\pi) + \lambda u(0) = 0$, or

$$i\omega_{\lambda} (D - C) + i\omega_{\lambda} (Ce^{2i\pi\omega_{\lambda}} - De^{-2i\pi\omega_{\lambda}}) + \lambda(C + D) = 0.$$

¹This construction only works in one dimension. We do not know how to construct similar test functions in higher dimension.

We deduce that (C, D) is solution to the 2×2 matrix equation

$$\begin{pmatrix} 1 - e^{2i\pi\omega_\lambda} & 1 - e^{-2i\pi\omega_\lambda} \\ -i\omega_\lambda(1 - e^{2i\pi\omega_\lambda}) + \lambda & i\omega_\lambda(1 - e^{-2i\pi\omega_\lambda}) + \lambda \end{pmatrix} \begin{pmatrix} C \\ D \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \end{pmatrix}.$$

The determinant of the matrix must therefore vanish, which leads to

$$1 = \cos(2\pi\omega_\lambda) + \frac{\lambda \sin(2\pi\omega_\lambda)}{2\omega_\lambda}.$$

As $\lambda \rightarrow \infty$, one must have $\omega_\lambda \rightarrow 1/2$, or equivalently $\varepsilon_0^{V_{\nu_\lambda}} \rightarrow 1/4$. The result follows.

3.3 Proof of Theorem 2.3

We are now in position to give the proof of Theorem 2.3. Let $b \in \mathcal{T}$ and $B > 1/4 - b^*$ where $b^* := \int_{\Gamma^*} b(q) dq$. Let $V_n = V_{\nu_n} - B \subset \mathcal{V}_B$ be a minimising sequence associated to problem (2.8).

Let us first assume by contradiction that $\nu_n(\Gamma) \rightarrow \infty$. Then, according to Proposition 2.4, up to a subsequence (still denoted by n), there exists $\varepsilon \geq \frac{1}{4} - B$ such that $\varepsilon_q^{V_n}$ converges uniformly in $q \in \Gamma^*$ to the constant function ε . Also, from the second part of Proposition 2.4, the fact that $B > \frac{1}{4} - b^*$ and the fact that b^* is the unique minimiser to

$$\inf_{c \in \mathbb{R}} \mathcal{K}_b(c), \quad (3.5)$$

where $\mathcal{K}_b(c) := \int_{[0,1/2]} |b(q) - c|^2 dq$ for all $c \in \mathbb{R}$, it must hold that $\varepsilon = b^*$.

We now prove that

$$\inf_{V \in \mathcal{V}_B} \mathcal{J}_b(V) \neq \inf_{c \in \mathbb{R}} \mathcal{K}_b(c) = \mathcal{K}_b(b^*).$$

To this aim, we exhibit a potential $W \in \mathcal{V}_B$ such that $\mathcal{J}_b(W) < \mathcal{K}_b(b^*)$. Since b is continuous and increasing on $[0, 1/2]$, there exists a unique $q^* \in (0, 1/2)$ such that $b(q^*) = b^*$. We choose $\delta > 0$ small enough such that $0 < q^* - \delta < q^* + \delta < 1/2$, and set

$$\eta^{\text{ext}} := \int_0^{q^* - \delta} |b(q) - b^*|^2 dq + \int_{q^* + \delta}^{1/2} |b(q) - b^*|^2 dq \quad \text{and} \quad \eta^{\text{int}} := \int_{q^* - \delta}^{q^* + \delta} |b(q) - b^*|^2 dq,$$

so that $\mathcal{K}_b(b^*) = \eta^{\text{ext}} + \eta^{\text{int}}$. Since b is increasing and continuous, it holds that $\eta^{\text{int}} > 0$ and $\eta^{\text{ext}} > 0$, and that $b(q^* - \delta) < b^* < b(q^* + \delta)$.

We now choose a constant $\sigma > 0$ such that

$$0 < \sigma < \min \left\{ \frac{\eta^{\text{int}}}{8\delta}, B + b^* - \frac{1}{4}, b^* - b(q^* - \delta), b(q^* + \delta) - b^* \right\}.$$

Let ν_n be the measure defined in (3.3) for $\lambda = n \in \mathbb{N}$, and let

$$\widetilde{W}_n := V_{\nu_n} + b^* - \frac{1}{4}.$$

Since $\varepsilon_q^{\widetilde{W}_n}$ converges to b^* uniformly in Γ^* , there exists $n_0 \in \mathbb{N}^*$ large enough such that

$$\forall q \in \Gamma^*, \quad \left| \varepsilon_q^{\widetilde{W}_{n_0}} - b^* \right| < \sigma/2.$$

We then define

$$W := \widetilde{W}_{n_0} + b^* - \varepsilon_{q^*}^{\widetilde{W}_{n_0}} = V_{\nu_n} + \left[\left(B + b^* - \frac{1}{4} \right) - \left(\varepsilon_{q^*}^{\widetilde{W}_{n_0}} - b^* \right) \right] - B.$$

Since $\sigma < B + b^* - 1/4$, it holds that $W \in \mathcal{V}_B$. Moreover, it holds that $b^* - \sigma < \varepsilon_q^W < b^* + \sigma$ for all $q \in \Gamma^*$. Finally, for $q = q^*$, we have $\varepsilon_{q^*}^W = b^*$.

Let us evaluate $\mathcal{J}_b(W)$. We get

$$\mathcal{J}_b(W) = \int_0^{q^*-\delta} |b(q) - \varepsilon_q^W|^2 dq + \int_{q^*-\delta}^{q^*+\delta} |b(q) - \varepsilon_q^W|^2 dq + \int_{q^*+\delta}^{1/2} |b(q) - \varepsilon_q^W|^2 dq.$$

For the first part, we notice that for $0 \leq q < q^* - \delta$, we have

$$b(q) < b(q^* - \delta) < b^* - \sigma < \varepsilon_q^W < \varepsilon_{q^*}^W = b^*.$$

This yields that

$$\forall 0 \leq q < q^* - \delta, \quad |b(q) - \varepsilon_q^W| = \varepsilon_q^W - b(q) < b^* - b(q) = |b(q) - b^*|.$$

Integrating this inequality leads to

$$\int_0^{q^*-\delta} |b(q) - \varepsilon_q^W|^2 dq < \int_0^{q^*-\delta} |b(q) - b^*|^2 dq.$$

Similarly, we obtain that

$$\int_{q^*+\delta}^{1/2} |b(q) - \varepsilon_q^W|^2 dq < \int_{q^*+\delta}^{1/2} |b(q) - b^*|^2 dq.$$

Lastly, for the middle part, we have

$$\int_{q^*-\delta}^{q^*+\delta} |b(q) - \varepsilon_q^W|^2 dq < 2\delta [\varepsilon_{q^*+\delta}^W - \varepsilon_{q^*-\delta}^W] \leq 4\delta\sigma \leq \frac{\eta^{\text{int}}}{2} < \int_{q^*-\delta}^{q^*+\delta} |b(q) - b^*|^2 dq.$$

Combining all these inequalities yields that $\mathcal{J}_b(W) < \mathcal{K}_b(b^*)$. This contradicts the minimising character of the sequence $(V_n)_{n \in \mathbb{N}^*}$.

Hence the sequence $(\nu_n(\Gamma))_{n \in \mathbb{N}^*}$ is bounded. The proof of Theorem 2.3 then follows from Proposition 3.2.

4 Numerical tests

In this section, we present some numerical results obtained on different toy inverse band structure problems. We propose an adaptive optimisation algorithm in which the different discretisation parameters are progressively increased. Such an approach, although heuristic, shows a significant gain in computational time on the presented test cases in comparison to a naive optimisation approach.

In Section 4.1, we present the discretised version of the inverse band problem for multiple target bands. We present the different optimisation procedures used for this problem (direct and adaptive) in Section 4.2. Numerical results on different test cases are given in Section 4.3. The reader should keep in mind that although the proof given in the previous section only works for the reconstruction of the first band, it is possible to numerically look for methods that reproduce several bands.

4.1 Discretised inverse band structure problem

For $k \in \mathbb{Z}$, we let $e_k(x) := \frac{1}{\sqrt{2\pi}} e^{ikx}$ be the k -th Fourier mode. For $s \in \mathbb{N}^*$, we define by

$$X_s := \text{Span} \{e_k, k \in \mathbb{Z}, |k| \leq s\} \quad (4.1)$$

the finite dimensional space of L_{per}^2 consisting of the $N_s := 2s + 1$ lowest Fourier modes. We denote by $\Pi_{X_s} : L_{\text{per}}^2 \rightarrow X_s$ the L_{per}^2 orthogonal projector onto X_s . In practice, the solutions of the eigenvalue problem (2.4) are approximated using a Galerkin method in X_s . We denote by $\varepsilon_{q,1}^{V,s} \leq \dots \leq \varepsilon_{q,N_s}^{V,s}$ the eigenvalues (ranked in increasing order, counting multiplicity) of the operator $A_q^{V,s} := \Pi_{X_s} A_q^V \Pi_{X_s}^*$. We also denote by $(u_{q,1}^{V,s}, \dots, u_{q,N_s}^{V,s})$ an orthonormal basis of X_s composed of eigenvectors associated to these eigenvalues so that

$$\forall 1 \leq j \leq N_s, \quad A_q^{V,s} u_{q,j}^{V,s} = \varepsilon_{q,j}^{V,s} u_{q,j}^{V,s}. \quad (4.2)$$

An equivalent variational formulation of (4.2) is the following:

$$\forall 1 \leq j \leq N_s, \quad \forall v \in X_s, \quad a_q^V \left(u_{q,j}^{V,s}, v \right) = \varepsilon_{q,j}^{V,s} \left\langle u_{q,j}^{V,s}, v \right\rangle_{L_{\text{per}}^2}.$$

As s goes to $+\infty$, it holds that $\varepsilon_{q,m}^{V,s} \xrightarrow{s \rightarrow +\infty} \varepsilon_{q,m}^V$.

In order to perform the integration in (2.7), we discretise the Brillouin zone. We use a regular grid of size $Q \in \mathbb{N}^*$, and set

$$\Gamma_Q^* := \left\{ -\frac{1}{2} + \frac{j}{Q}, \quad j \in \{0, \dots, Q-1\} \right\}.$$

We emphasise that since the maps $q \mapsto \varepsilon_{q,m}$ are analytic and periodic, the discretisation error coming from the integration will be exponentially small with respect to Q . In practice, we fix $Q \in \mathbb{N}^*$.

Let $M \in \mathbb{N}^*$ be a desired number of targeted bands and $b_1, \dots, b_M \in C_{\text{per}}^0$ be real-valued even functions, and such that b_m is increasing when m is odd and decreasing when m is even. Our cost functional is therefore $\mathcal{J} : H_{\text{per},r}^{-1} \rightarrow \mathbb{R}$, defined by

$$\forall V \in H_{\text{per},r}^{-1}, \quad \mathcal{J}(V) := \frac{1}{Q} \sum_{q \in \Gamma_Q^*} \sum_{m=1}^M |b_m(q) - \varepsilon_{q,m}^V|^2.$$

Its discretised version, when the eigenvalues problems are solved with a Galerkin approximation, is

$$\forall s \in \mathbb{N}^*, \quad \forall V \in H_{\text{per},r}^{-1}, \quad \mathcal{J}^s(V) := \frac{1}{Q} \sum_{q \in \Gamma_Q^s} \sum_{m=1}^M |b_m(q) - \varepsilon_{q,m}^{V,s}|^2.$$

Recall that our goal is to find a potential $V \in H_{\text{per},r}^{-1}$ which minimise the functional \mathcal{J}^s . In practice, an element $V \in H_{\text{per},r}^{-1}$ is approximated with a finite set of Fourier modes. For $p \in \mathbb{N}^*$, we denote by

$$Y_p := \text{Span} \left\{ \sum_{k \in \mathbb{Z}, |k| \leq p} \widehat{V}_k e_k, \quad \forall k \in \mathbb{Z}, |k| \leq p, \quad \overline{\widehat{V}_{-k}} = \widehat{V}_k \right\}. \quad (4.3)$$

Altogether, we want to solve

$$V^{s,p} := \arg \min_{V \in Y_p} \mathcal{J}^s(V).$$

4.2 Algorithms for optimisation procedures

4.2.1 Naive algorithm

We first present a naive optimisation procedure, using a gradient descent method, where the parameters s and p are fixed beforehand. We tested three different versions of the gradient descent algorithm: steepest descent (**SD**), conjugate gradient with Polak Ribiere formula (**PR**) and quasi Newton with the Broyden-Fletcher-Goldfarb-Shanno formula (**BFGS**). We do not detail here these classical descents and corresponding line search routines for the sake of conciseness and refer the reader to [1, 3].

For all $V \in H_{\text{per},r}^{-1}$, there exists real-valued coefficients $(c_k^V)_{k \in \mathbb{N}}$ and $(d_k^V)_{k \in \mathbb{N}^*}$ such that

$$V(x) = c_0^V + \sum_{k \in \mathbb{N}^*} c_k^V \cos(kx) + d_k^V \sin(kx), \quad \text{and} \quad \sum_{k \in \mathbb{N}^*} (1 + |k|^2)^{-1} (|c_k^V|^2 + |d_k^V|^2) < +\infty.$$

For all $k \in \mathbb{N}$ (respectively $k \in \mathbb{N}^*$), we can express the derivative $\partial_{c_k^V} \mathcal{J}^s(V)$ (respectively $\partial_{d_k^V} \mathcal{J}^s(V)$) exactly in terms of the Bloch eigenvectors $u_{q,m}^{V,s}$. Indeed, it holds that

$$\partial_{c_k^V} \mathcal{J}^s(V) = \frac{1}{Q} \sum_{q \in \Gamma_Q^s} \sum_{m=1}^M 2 (\varepsilon_{q,m}^{V,s} - b_m(q)) \partial_{c_k^V} (\varepsilon_{q,m}^{V,s}).$$

On the other hand, from the Hellman-Feynman theorem, it holds that

$$\partial_{c_k^V} (\varepsilon_{q,m}^{V,s}) = \left\langle u_{q,m}^{V,s}, \partial_{c_k^V} A_q^V, u_{q,m}^{V,s} \right\rangle = \langle u_{q,m}^{V,s}, \cos(k \cdot) u_{q,m}^{V,s} \rangle_{L_{\text{per}}^2}.$$

Similarly, for all $k \in \mathbb{N}^*$,

$$\partial_{d_k^V} (\varepsilon_{q,m}^{V,s}) = \left\langle u_{q,m}^{V,s}, \partial_{d_k^V} A_q^V, u_{q,m}^{V,s} \right\rangle = \langle u_{q,m}^{V,s}, \sin(k \cdot) u_{q,m}^{V,s} \rangle_{L_{\text{per}}^2}.$$

In the rest of the article, for all $p \in \mathbb{N}^*$, we will denote by $\nabla \mathcal{J}^s(V)|_{Y^p}$ the $2p + 1$ -dimensional real-valued vector so that

$$\nabla \mathcal{J}^s(V)|_{Y^p} = \left(\partial_{d_p^V} \mathcal{J}^s(V), \partial_{d_{p-1}^V} \mathcal{J}^s(V), \dots, \partial_{d_1^V} \mathcal{J}^s(V), \partial_{c_p^V} \mathcal{J}^s(V), \partial_{c_1^V} \mathcal{J}^s(V), \dots, \partial_{c_p^V} \mathcal{J}^s(V) \right).$$

In order for the reader to better compare our adaptive algorithm with this naive one, we provide its pseudo-code below (Algorithm 1).

<p>Algorithm 1: Naive optimisation algorithm</p> <p>Input: $p, s \in \mathbb{N}^*$; $W_0 \in Y_p$: initial guess; $\varepsilon > 0$: prescribed global precision; $\nu > 0$: tolerance for the norm of the gradient;</p> <p>Output: $W_* \in Y_p$ such that $\ \nabla \mathcal{J}^s(W_*) _{Y_p}\ \leq \nu$;</p> <p>Instructions: $n = 0, W = W_0$; while $\ \nabla \mathcal{J}^s(W) _{Y_p}\ > \nu$ do compute a descent direction $D \in Y_p$ at $\mathcal{J}^s(W)$ (using SD / PR / BFGS); choose $t \in \mathbb{R}$ so that $t \in \underset{\tilde{t} \in \mathbb{R}}{\text{argmin}} \mathcal{J}^s(W + \tilde{t}D)$; set $W \leftarrow W + tD$; return $W_* = W$.</p>

Although this method gives satisfactory numerical optimisers as shown in Section 4.3, its computational time grows very quickly with the discretisation parameters p and s . Besides, it is not clear how these parameters should be chosen a priori, given some target bands. This motivates the design of an adaptive algorithm.

4.2.2 Adaptive algorithm

In order to improve on the efficiency of the numerical optimisation procedure, we propose an adaptive algorithm, where the discretisation parameters s or p are increased during the optimisation process. To describe this procedure, we introduce two criteria to determine whether s or p need to be increased during the algorithm.

As the parameter s is increased, the approximated eigenvalues $\varepsilon_{q,m}^{V,s}$ becomes more accurate, and the discretised cost functional \mathcal{J}^s gets closer to the true one \mathcal{J} . Our criterion for s relies on the use of an a posteriori error estimator for the eigenvalue problem (4.2). More precisely, assume we can calculate at low numerical cost an estimator $\Delta_{q,m}^{V,s} \in \mathbb{R}_+$ such that

$$|\varepsilon_{m,q}^V - \varepsilon_{m,q}^{V,s}| \leq \Delta_{q,m}^{V,s},$$

(see Appendix A), then we would have that

$$\begin{aligned}
|\mathcal{J}(V) - \mathcal{J}^s(V)| &= \left| \frac{1}{Q} \sum_{q \in \Gamma_Q^*} \sum_{m=1}^M (|b_m(q) - \varepsilon_{q,m}^V|^2 - |b_m(q) - \varepsilon_{q,m}^{V,s}|^2) \right| \\
&= \left| \frac{1}{Q} \sum_{q \in \Gamma_Q^*} \sum_{m=1}^M (2b_m(q) - \varepsilon_{q,m}^V - \varepsilon_{q,m}^{V,s}) (\varepsilon_{q,m}^{V,s} - \varepsilon_{q,m}^V) \right| \\
&\leq \frac{1}{Q} \sum_{q \in \Gamma_Q^*} \sum_{m=1}^M (2|b_m(q) - \varepsilon_{q,m}^{V,s}| + \Delta_{q,m}^{V,s}) \Delta_{q,m}^{V,s} =: \mathcal{S}_V^s.
\end{aligned}$$

The quantity \mathcal{S}_V^s estimates the error between $\mathcal{J}(V)$ and $\mathcal{J}^s(V)$ and therefore gives information on the necessity to adapt the value of the discretisation parameter s .

We now derive a criterion for the parameter p . When this parameter is increased, the minimisation space Y_p gets larger. A natural way to decide whether or not to increase p is therefore to consider the gradient of \mathcal{J}^s , at the current minimisation point $W \in Y_p$, but calculated on a larger subspace $Y_{p'} \supset Y_p$ with $p' > p$.

In practice, the natural choice $p' = p + 1$ is inefficient. This is not a surprise, as there is no reason a priori to expect a sudden change at exactly the next Fourier mode. We therefore took the heuristic choice $p' = 2p$. More specifically, we define

$$\mathcal{P}_V^p := \left\| \nabla_V \mathcal{J}^s(V) \Big|_{Y_{2p}} \right\|.$$

Note that this estimator needs to be computed only when V is a local minimum of \mathcal{J}^s on Y_p . When this estimator is larger than some threshold, we increase p so that the new space Y_p contains the Fourier mode which provides the highest contribution in $(\nabla_V \mathcal{J}^s(V)) \Big|_{Y_{2p}}$.

The adaptive procedure we propose is described in details in Algorithm 2:

Algorithm 2: Adaptive optimisation algorithm

Input:

$p_0, s_0 \in \mathbb{N}^*$: initial discretisation parameters;
 $W_0 \in Y_{p_0}$: initial guess;
 $\eta > 0$: global discretisation precision;
 $\nu > 0$: gradient norm precision;

Output:

$p \geq p_0, s \geq s_0$: final discretisation parameters;
 $W_* \in Y_p$ such that $\|\nabla \mathcal{J}^s(W_*)\|_{Y_p} \leq \nu$, $\mathcal{S}_{W_*}^s \leq \eta$ and $\mathcal{P}_{W_*}^p \leq \eta$;

Instructions:

$n = 0, W = W_0$;

while $\|\nabla \mathcal{J}^s(W)\|_{Y_p} > \nu$ **or** $\mathcal{S}_W^s > \eta$ **or** $\mathcal{P}_W^p > \eta$ **do**

while $\|\nabla \mathcal{J}_p^s(W)\|_{Y_p} > \nu$ **do**

 compute a descent direction $D \in Y_p$ at $\mathcal{J}^s(W)$ (using **SD** / **PR** / **BFGS**);
 choose $t \in \mathbb{R}$ so that $t \in \arg \min_{\bar{t} \in \mathbb{R}} \mathcal{J}^s(W + \bar{t}D)$;

 set $W \leftarrow W + tD$;

if $\mathcal{S}_W^s > \eta$ **then**

 set $s \leftarrow s + 1$;

else if $\mathcal{P}_W^p > \eta$ **then**

 set $p \leftarrow \arg \max_{p < \bar{p} \leq 2p} \max \left(\left| \partial_{d_{\bar{p}}}^V \mathcal{J}^s(W) \right|, \left| \partial_{c_{\bar{p}}}^V \mathcal{J}^s(W) \right| \right)$;

return $W_* = W$.

4.3 Numerical results

In this section, we illustrate the different algorithms presented above.

We consider the case where the target functions come from a target potential $V_t \in Y_{p_t}$, whose Fourier coefficients are randomly chosen for some $p_t \in \mathbb{N}^*$. We therefore take $b_m(q) := \varepsilon_{q,m}^{V_t, s_t}$, and try to recover the first M functions b_m . The numerical parameters are $M = 3$, $Q = 25$, $\nu = 10^{-5}$, $\eta = 10^{-6}$ and $s_t = 20$. The initial guess is $W_0 = 0$. The naive algorithms are run with $s = s_t$ and $p = p_t$, while the adaptive algorithms start with $s_0 = p_0 = 1$. In addition, the a posteriori estimator is obtained with $s_{\text{ref}} = 250$ and $\theta = 0.01$ (see Appendix A). All tests are done with the naive and adaptive algorithms, with steepest descent (**SD**), conjugate gradient with Polak Ribiere formula (**PR**) and quasi Newton with the Broyden-Fletcher-Goldfarb-Shanno formula (**BFGS**).

In our first test, we try to recover a simple shifted cosine function (i.e. $p_t = 1$). Results are shown in Figure 1. We observe that the bands and the potential are well reconstructed. We also notice that the adaptive algorithm takes more iterations to converge. However, as we will see later, most iterations are performed for low values of the parameters s and p , and therefore are usually faster in terms of CPU time (see Table 1 below).

In the second test case, we try to recover a more complex potential with $p_t = 8$ (see Figure 2). In this case, all the algorithms reproduce well the first bands, but fail to recover the potential. Actually, we see how different methods can lead to different local minima for the functional \mathcal{J} . This reflects the complex landscape of this function.

We end this section by reporting results obtained with the different algorithms, and for different target potential $V_t \in Y_{p_t}$ with $p_t = 1, 4, 8, 12$ (see Table 1). In this table, N denotes the number of iterations, s_N and p_N are the values of the parameters s and p at the last iteration (in particular, for the naive algorithms, we have $s_N = s_t = 20$ and $p_N = p_t$). Lastly, for each algorithm **algo**, we define a relative CPU time

$$\tau_{\text{algo}} = \frac{t_{\text{algo}}}{t_{\text{SD}}},$$

where t_{algo} is the CPU time consumed by the algorithm **algo** and t_{SD} is the CPU time consumed by the classical steepest descent. In particular, $\tau_{\text{SD}} = 1$.

p_t	-	BFGS		PR		SD	
	-	naive	adaptive	naive	adaptive	naive	adaptive
1	τ	0.259	1.176	0.929	1.320	1	1.255
	N	8	31	21	154	24	90
	s_N	20	3	20	4	20	3
	p_N	1	3	1	2	1	3
4	τ	0.070	0.009	0.464	0.281	1	0.259
	N	54	1424	1927	7091	8453	19095
	s_N	20	8	20	7	20	5
	p_N	4	5	4	3	4	3
8	τ	0.470	0.151	1.090	0.144	1	0.519
	N	553	1041	1023	1515	7326	26783
	s_N	20	6	20	7	20	6
	p_N	8	4	8	4	8	4
12	τ	0.007	0.001	0.054	0.004	1	0.044
	N	765	2474	2413	2727	50312	34865
	s_N	20	9	20	9	20	9
	p_n	12	8	12	8	12	8

Table 1: Results for recovery test with different algorithms. Red values are reference values.

We notice that although the adaptive approach requires more iterations to converge, it is usually faster than the naive one. As we already mentioned, this is due to the fact that most of the iterations are performed with small values of p and s , and are therefore faster. Moreover, we notice that the adaptive algorithms tend to find an optimised potential which $p_N \leq p_t$, i.e. a less oscillatory potential than the target one.

Acknowledgements

The authors heartily thank Éric Cancès, Julien Vidal, Damiano Lombardi and Antoine Levitt for their great help in this work and for inspiring discussions. The IRDEP institute is acknowledged for funding.

A A posteriori error estimator for the eigenvalue problem

We present in this appendix the a posteriori error estimator for eigenvalue problems that we use in Section 4.3. More details about this estimator are given in [2].

Let \mathcal{H} be a finite dimensional space of size N_{ref} and let A be a self-adjoint operator on \mathcal{H} . In our case, \mathcal{H} is some $X_{s_{\text{ref}}}$ for some large $s_{\text{ref}} \gg 1$, and $A = A_q^{V, s_{\text{ref}}}$. The eigenvalues of A , counting multiplicities are denoted by $\varepsilon_1 \leq \varepsilon_2 \leq \dots \leq \varepsilon_{N_{\text{ref}}}$.

For $N \ll N_{\text{ref}}$, we consider X_N a finite dimensional subspace of \mathcal{H} . We denote by Π_{X_N} the orthogonal projection on X_N , and by $A^N := \Pi_{X_N} A \Pi_{X_N}^*$. The eigenvalues of A^N are denoted by $\varepsilon_1^N \leq \varepsilon_2^N \leq \dots \leq \varepsilon_N^N$. Let us also denote by $(u_m^N)_{1 \leq m \leq N}$ a corresponding orthogonal basis of X^N , so that

$$\forall 1 \leq m \leq N, \quad A^N u_m^N = \varepsilon_m^N u_m^N.$$

We recall that, from the min-max principle, it holds that $\varepsilon_m \leq \varepsilon_m^N$. A certified a posteriori error estimator for the m -th eigenvalue is a non-negative real number $\Delta_m^N \in \mathbb{R}_+$ such that

$$\varepsilon_m^N - \varepsilon_m \leq \Delta_m^N.$$

We also require that the expression of Δ_m^N only involves the approximated eigenpair ε_m^N and u_m^N (and not ε_m).

Proposition A.1. *Assume that ε_m (resp. ε_m^N) is a non-degenerate eigenvalue of A (resp. A^N), and that*

$$0 < \varepsilon_m^N - \varepsilon_m < \text{dist}(\varepsilon_m^N, \sigma(A) \setminus \{\varepsilon_m\}). \quad (\text{A.1})$$

Let $\lambda_m < \varepsilon_m$. Then there exists $\delta_m > 0$ such that, for all $0 \leq \delta < \delta_m$, we have

$$\varepsilon_m^N - \varepsilon_m \leq \left\langle r_m^N, (A - c_\delta)^{-1} (A - d_\delta) (A - c_\delta)^{-1} r_m^N \right\rangle, \quad (\text{A.2})$$

where we set $c_\delta := \varepsilon_m^N + \delta$, $d_\delta := \lambda_m + \delta$, and where $r_m^N := (A - \varepsilon_m^N) u_m^N$ is the residual.

Proof. Assumption (A.1) implies that $\varepsilon_m^N \notin \sigma(A)$, so that $(A - \varepsilon_m^N)$ is invertible. From the fact that $\langle u_m^N, A u_m^N \rangle = \varepsilon_m^N$, and the definition of the residual, it holds that

$$\varepsilon_m^N - \varepsilon_m = \left\langle r_m^N, (A - \varepsilon_m^N)^{-1} (A - \varepsilon_m) (A - \varepsilon_m^N)^{-1} r_m^N \right\rangle. \quad (\text{A.3})$$

Thus, a sufficient condition for (A.2) to hold is that

$$(A - c_\delta)^{-1} (A - d_\delta) (A - c_\delta)^{-1} \geq (A - \varepsilon_m^N)^{-1} (A - \varepsilon_m) (A - \varepsilon_m^N)^{-1}.$$

Thanks to the spectral decomposition of A , this is the case if and only if,

$$\forall 1 \leq \tilde{m} \leq N_{\text{ref}}, \quad \frac{\varepsilon_{\tilde{m}} - d_\delta}{(\varepsilon_{\tilde{m}} - c_\delta)^2} \geq \frac{\varepsilon_{\tilde{m}} - \varepsilon_m}{(\varepsilon_{\tilde{m}} - \varepsilon_m^N)^2}.$$

Denoting by $\eta := \text{dist}(\varepsilon_m^N, \sigma(A) \setminus \{\varepsilon_m\}) - (\varepsilon_m^N - \varepsilon_m)$, this holds true as soon as $\delta \leq \delta_m := \min(\varepsilon_m - \lambda_m, \eta)$. The result follows. \square

In order to use the left-side of (A.2) as an a posteriori estimator, we need to choose $\lambda_m < \varepsilon_m$ and $\delta_m > 0$. For the choice of λ_m , we follow [25], and notice that

$$\varepsilon_m \geq \lambda_m := \mu - \left(\frac{N_{\text{ref}} - m - 1}{m + 1} \right)^{1/2} \sigma,$$

where we set

$$\mu := \frac{1}{N_{\text{ref}}} \text{Tr } A \quad \text{and} \quad \sigma^2 := \frac{1}{N_{\text{ref}}} \text{Tr } A^2 - \mu^2.$$

For the choice of δ_m , we chose the simple rule

$$\delta_m = \theta (\varepsilon_m^N - \kappa) \quad \text{with} \quad 0 < \theta \ll 1 \quad \text{and} \quad \kappa \in \mathbb{R} \quad \text{independent of } m.$$

The real number κ is chosen to be an a priori lower bound of the lowest eigenvalue ε_1 of A . This choice is heuristic in the sense that we cannot guarantee that the assumptions of Proposition A.1 are satisfied. However, the encouraging numerical results we obtain below motivated our choice to use such an estimator (see Section A).

Numerical test To illustrate the efficiency of our heuristic, we tested it to compute the first bands of the Hill's operator A^V with

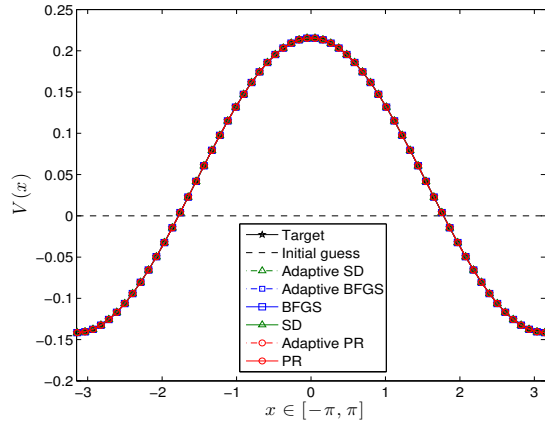
$$V(x) = \sum_{k=-3}^3 \hat{V}_k e_k, \quad \text{where} \quad \hat{V}_0 = 2 \quad \text{and} \quad \overline{\hat{V}_{-1}} = \overline{\hat{V}_{-2}} = \hat{V}_1 = \hat{V}_2 = 1 + 0.5i.$$

The reference operator is $A := A_q^{V, s_{\text{ref}}}$ with $s_{\text{ref}} = 250$, and the first three bands are computed on the space X^s defined in 4.1 with $s = 6$. We plot in Figure 3 the true error $\varepsilon_{q,m}^{V,s} - \varepsilon_{q,m}^{V,s_{\text{ref}}}$ for $m = 1, 2, 3$, and the corresponding a posteriori error with $\kappa = 0$ and different values of θ (namely $\theta = 0.1, 0.5, 1$). We observe that our estimator is sharp for a large range of θ .

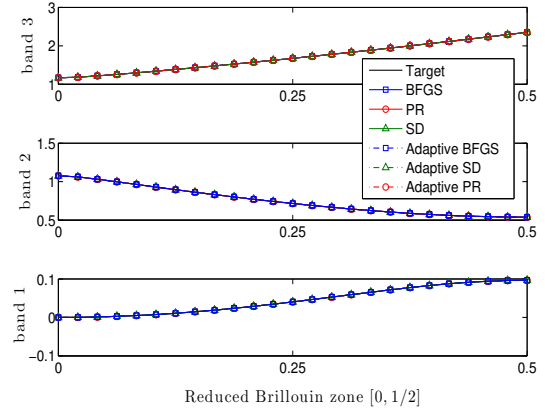
References

- [1] A. Bakhta. *Mathematical models and numerical simulation of photovoltaic materials*. PhD thesis, Université Paris Est - Ecole des Ponts ParisTech, 2017.
- [2] A. Bakhta and D. Lombardi. An a posteriori error estimator based on shifts for positive hermitian eigenvalue problems. <https://hal.inria.fr/hal-01584180/>, 2017.
- [3] J.F. Bonnans, J.Ch. Gilbert, C. Lemaréchal, and C. Sagastizábal. *Numerical Optimization*. Springer Verlag, 2003.
- [4] N.C. Dias, C. Jorge, and J.N. Prata. One-dimensional Schrödinger operators with singular potentials: A Schwartz distributional formulation. *J. of Differential Equations*, 260(8):6548–6580, 2016.
- [5] G. Eskin. Inverse spectral problem for the Schrödinger equation with periodic vector potential. *Commun. Math. Phys.*, 125(2):263–300, 1989.
- [6] J. Eskin, J. Ralston, and E. Trubowiz. On isospectral periodic potential in \mathbb{R}^n . I. *Commun. Pure Appl. Maths.*, 37(5):647–676, 1984.
- [7] J. Eskin, J. Ralston, and E. Trubowiz. On isospectral periodic potential in \mathbb{R}^n . II. *Commun. Pure Appl. Maths.*, 37(6):715–753, 1984.
- [8] L.C. Evans. *Partial Differential Equations*. Graduate studies in mathematics. American Mathematical Society, 1998.
- [9] G. Freiling and V. Yurko. *Inverse Sturm-Liouville problems and their applications*. Nova Science Publishers, 2001.
- [10] F. Gesztesy and M. Zinchenko. On spectral theory for Schrödinger operators with strongly singular potentials. *Math. Nachr.*, 279(9-10):1041–1082, 2006.
- [11] R.O. Hryniv and Y.V. Mykytyuk. 1-D Schrödinger operators with periodic singular potentials. *Methods Funct. Anal. Topology*, 7(4):31–42, 2001.
- [12] R.O. Hryniv and Y.V. Mykytyuk. Inverse spectral problems for Sturm-Liouville operators with singular potentials. *Inverse Problems*, 19(3):665, 2003.

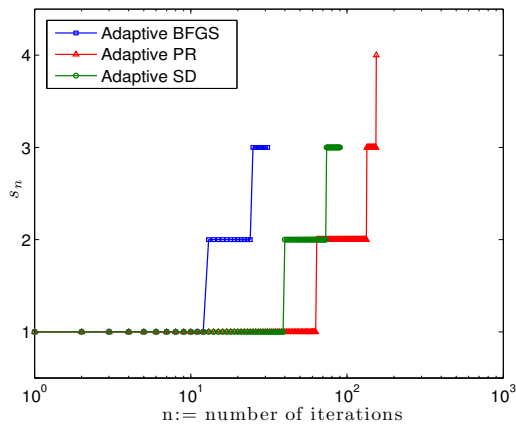
- [13] R.O. Hryniv and Y.V. Mykytyuk. Inverse spectral problems for Sturm-Liouville operators with singular potentials. III. Reconstruction by three spectra. *J. Math. Anal. Appl.*, 284(2):626–646, 2003.
- [14] R.O. Hryniv and Y.V. Mykytyuk. Half-inverse spectral problems for Sturm-Liouville operators with singular potentials. *Inverse Problems*, 20(5):1423, 2004.
- [15] R.O. Hryniv and Y.V. Mykytyuk. Inverse spectral problems for Sturm-Liouville operators with singular potentials. II. Reconstruction by two spectra. *North-Holland Mathematics Studies*, 197:97–114, 2004.
- [16] R.O. Hryniv and Y.V. Mykytyuk. Inverse spectral problems for Sturm-Liouville operators with singular potentials. IV. Potentials in the Sobolev space scale. *Proceedings of the Edinburgh Mathematical Society*, 49(2):309–329, 2006.
- [17] T. Kato. Schrödinger operators with singular potentials. *Israel Journal of Mathematics*, 13(1):135–148, 1972.
- [18] P. Kuchment. An overview of periodic elliptic operators. *Bull. Amer. Math. Soc.*, 53(3):343–414, 2016.
- [19] E.H. Lieb and M. Loss. *Analysis*, volume 14 of *Graduate studies in mathematics*. 2001.
- [20] V. Mikhaelets and V. Molyboga. One-dimensional Schrödinger operators with singular periodic potentials. *Methods Funct. Anal. Topology*, 14(2):184–200, 2008.
- [21] J. Pöschel and E. Trubowiz. *Inverse spectral theory. Pure and applied mathematics*. Academic Press, 1987.
- [22] M. Reed and B. Simon. *Methods of modern mathematical physics. IV: Analysis of operators*. Elsevier, 1978.
- [23] N.S. Trudinger. On Harnack type inequalities and their application to quasilinear elliptic equations. *Commun. Appl. Math.*, 20(4):721–747, 1967.
- [24] O. Veliev. *Multidimensional periodic Schrödinger operator. Perturbation theory and applications*. Academic Press, 2015.
- [25] H. Wolkowicz and G.P.H. Styan. Bounds on eigenvalues using traces. *Linear Algebra Appl.*, 29:471–506, 1980.



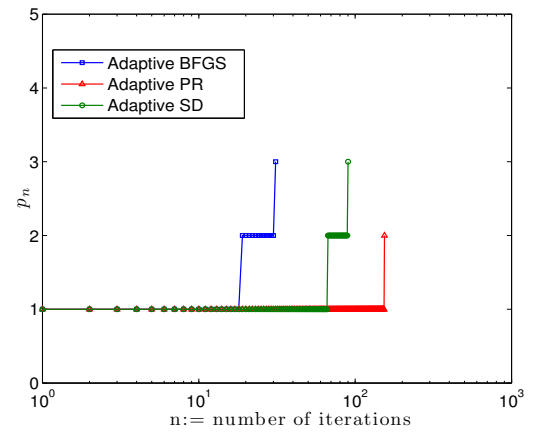
(a) Potentials



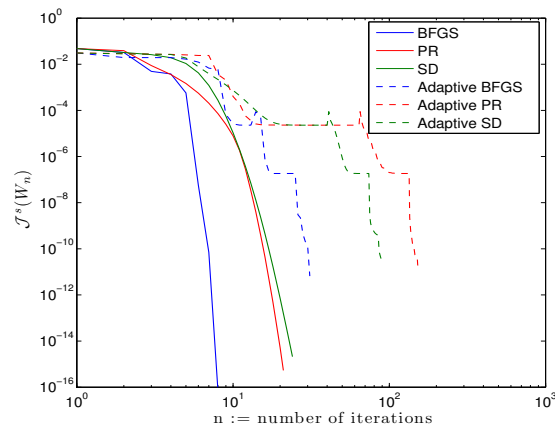
(b) Bands



(c) Evolution of s

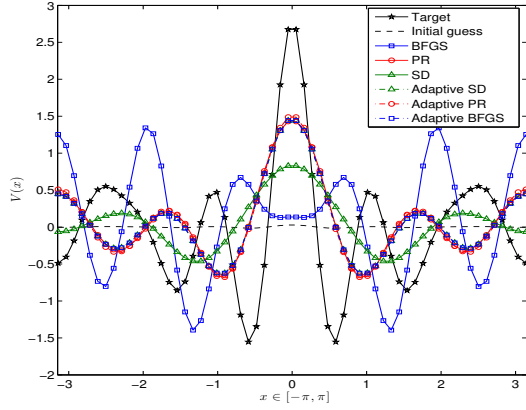


(d) Evolution of p

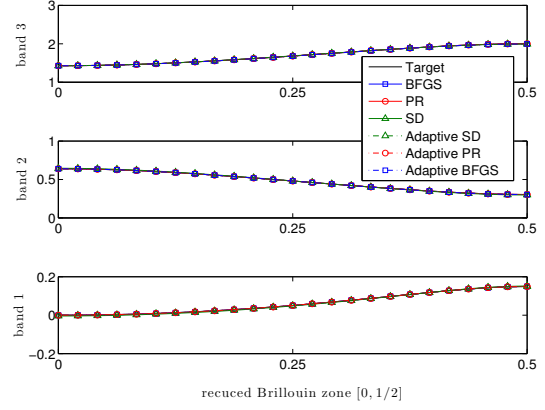


(e) Convergence of the algorithms

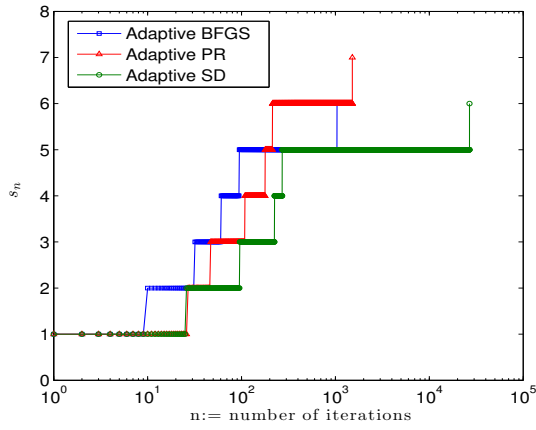
Figure 1: Recovery of the cosine potential



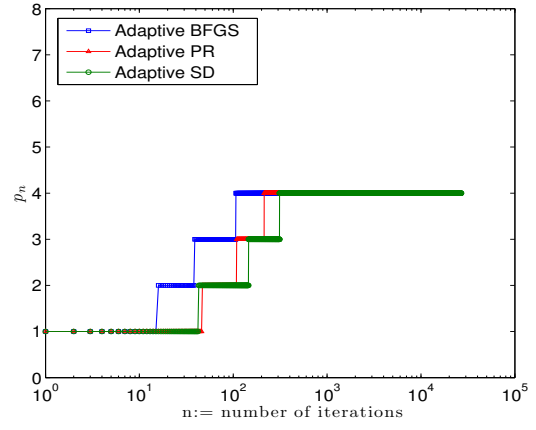
(a) Potentials



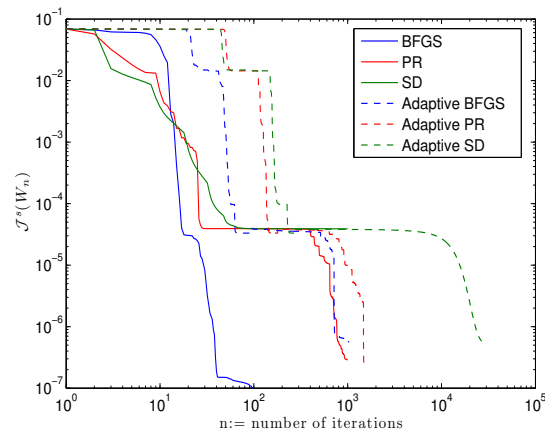
(b) Bands



(c) Evolution of s



(d) Evolution of p



(e) Convergence of the algorithms

Figure 2: Recovery of an oscillating potential.

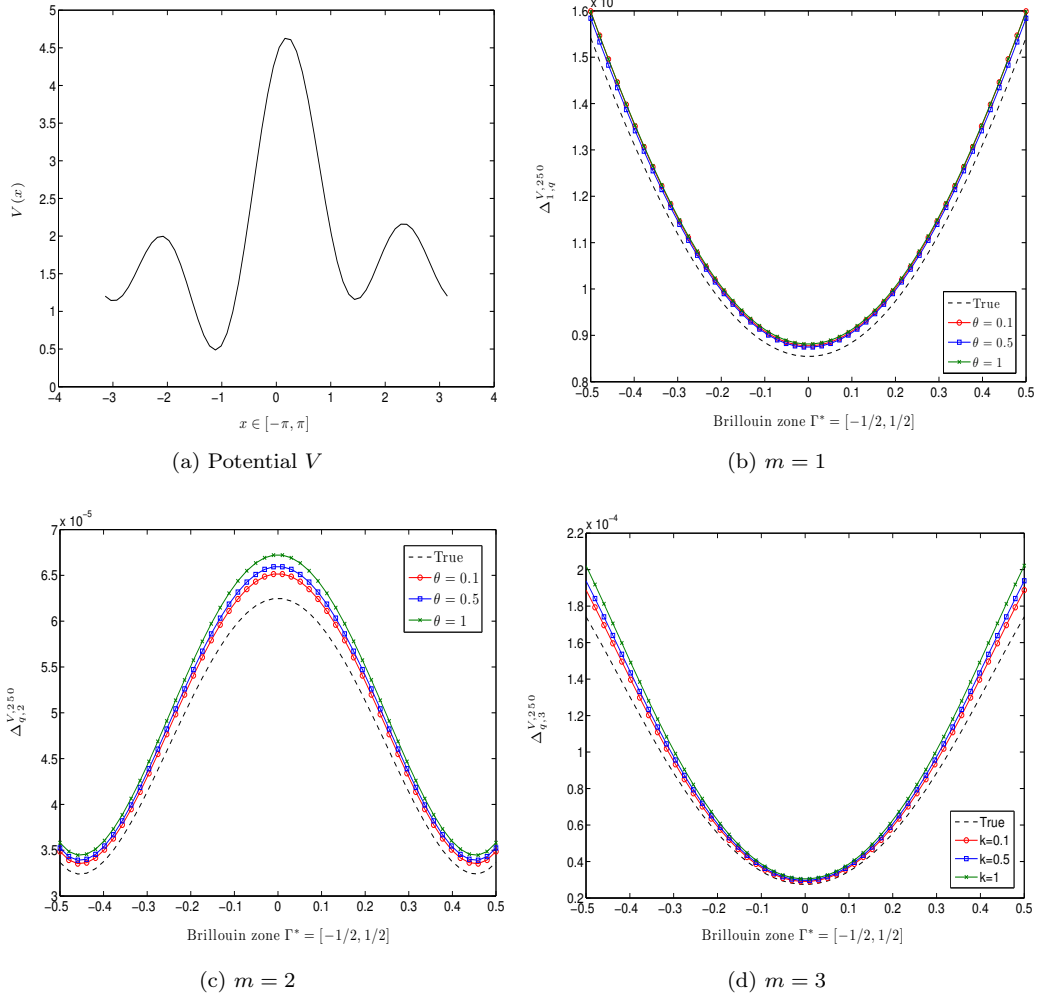


Figure 3: Numerical validation of the a posteriori error estimator proposed in Appendix A.