# The equivariant index of twisted dirac operators and semi-classical limits 

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# THE EQUIVARIANT INDEX OF TWISTED DIRAC OPERATORS AND SEMI-CLASSICAL LIMITS. 

PAUL-EMILE PARADAN AND MICHÈLE VERGNE


#### Abstract

Let $G$ be a compact connected Lie group with Lie algebra $\mathfrak{g}$. Let $M$ be a compact spin manifold with a $G$-action, and $\mathcal{L}$ be a $G$-equivariant line bundle on $M$. Consider an integer $k$, and let $\mathcal{Q}_{G}^{\text {spin }}\left(M, \mathcal{L}^{k}\right)$ be the equivariant index of the Dirac operator on $M$ twisted by $\mathcal{L}^{k}$. Let $m_{G}(\lambda, k)$ be the multiplicity in $\mathcal{Q}_{G}^{\text {spin }}\left(M, \mathcal{L}^{k}\right)$ of the irreducible representation of $G$ attached to the admissible coadjoint orbit $G \lambda$. We prove that the distribution $\left\langle\Theta_{k}, \varphi\right\rangle=k^{\operatorname{dim}(G / T) / 2} \sum_{\lambda} m_{G}(\lambda, k)\left\langle\beta_{\lambda / k}, \varphi\right\rangle$ has an asymptotic expansion when $k$ tends to infinity of the form $\left\langle\Theta_{k}, \varphi\right\rangle \equiv$ $k^{\operatorname{dim} M / 2} \sum_{n=0}^{\infty} k^{-n}\left\langle\theta_{n}, \varphi\right\rangle$. Here $\varphi$ is a test function on $\mathfrak{g}^{*}$ and $\left\langle\beta_{\xi}, \varphi\right\rangle$ is the integral of $\varphi$ on the coadjoint orbit $G \xi$ with respect to the canonical Liouville measure. We compute explicitly the distribution $\theta_{n}$ in terms of the graded $\hat{A}$ class of $M$ and the equivariant curvature of $\mathcal{L}$.

If $M$ is non compact, we use these asymptotic techniques to give another proof of the fact that the formal geometric quantification of a manifold with a spinc structure is functorial with respect to restriction to subgroups.


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## To the memory of Bertram Kostant.

This article is pursuing the fundamental idea of Kostant that a line bundle $\mathcal{L}$ on a $G$-manifold $M$, equipped with a $G$-invariant connection, give rise to a moment $\operatorname{map} \phi_{G}: M \rightarrow \mathfrak{g}^{*}$, and thus hopefully to a relation between the quantization of $M$ and the representation theory of $G$, since $\widehat{G}$ may be considered as a subset of $\mathfrak{g}^{*}$. We investigate further the corresponding quantization, via Dirac operators, in the case where the group $G$ is compact, and the moment map proper. Our new insight is that the asymptotic behavior when $\mathcal{L}$ is replaced by $\mathcal{L}^{k}$ is easy to compute classically, and determines completely the quantization at $k=1$. Certainly, Bertram Kostant would have found this idea obvious, but maybe also beautiful. Anyway, here it is.

## 1. Introduction

Let $M$ be a compact oriented spin manifold of even dimension $2 d$. Let $\mathcal{L}$ be a line bundle over $M$ equipped with a connection of curvature $-i \Omega$ and let $\hat{A}(M)$ be the $\hat{A}$ class of $M$ (normalized as in [2]). We do not assume the curvature $-i \Omega$ of $\mathcal{L}$ to be non-degenerate. Define the integral

$$
\mathcal{Q}^{\mathrm{geo}}(M, \mathcal{L})=\frac{1}{(-2 i \pi)^{d}} \int_{M} e^{-i \Omega} \hat{A}(M)
$$

The Dirac operator $D_{\mathcal{L}}$ twisted by $\mathcal{L}$ is an elliptic operator on $M$, and let $\mathcal{Q}^{\text {spin }}(M, \mathcal{L})=\operatorname{dim} \operatorname{Ker}\left(D_{\mathcal{L}}\right)-\operatorname{dim} \operatorname{Coker}\left(D_{\mathcal{L}}\right)$ the index of $D_{\mathcal{L}}$. By the Atiyah-Singer index theorem, $\mathcal{Q}^{\text {spin }}(M, \mathcal{L})=\mathcal{Q}^{\text {geo }}(M, \mathcal{L})$.

Assume now that a compact connected Lie group $G$, with Lie algebra $\mathfrak{g}$, acts on $(M, \mathcal{L})$. The choice of an invariant connection on the line bundle $\mathcal{L}$ determines a moment map $\phi_{G}: M \rightarrow \mathfrak{g}^{*}$. This is the Kostant moment map [12]. Let $T$ be a maximal torus of $G, \mathfrak{t}$ its Lie algebra. We identify the space $\widehat{G}$ of irreducible finite dimensional representations of $G$ to the discrete set of elements $\lambda \in \mathfrak{t}^{*}$, which are dominant admissible and regular, and we denote by $\chi_{\lambda}(g)$ the trace of the action of $g \in G$ on the irreducible representation of $G$ parametrized by $\lambda$.

If $k$ is an integer, we denote $\mathcal{L}^{k}$ the $k$-th power of the line bundle $\mathcal{L}$. Define the function $\mathcal{Q}_{G}^{\text {spin }}\left(M, \mathcal{L}^{k}\right)$ on $G$ to be the trace of the action of $g \in G$ in the virtual space $\operatorname{Ker}\left(D_{\mathcal{L}^{k}}\right)-\operatorname{Coker}\left(D_{\mathcal{L}^{k}}\right)$. Define $m_{G}(\lambda, k) \in \mathbb{Z}$ such that

$$
\mathcal{Q}_{G}^{\mathrm{spin}}\left(M, \mathcal{L}^{k}\right)=\sum_{\lambda \in \hat{G}} m_{G}(\lambda, k) \chi_{\lambda} .
$$

Consider the geometric (rescaled) analogue

$$
\Theta_{k}^{(M, \mathcal{L})}=k^{\frac{\operatorname{dim}(G / T)}{2}} \sum_{\lambda \in \hat{G}} m_{G}(\lambda, k) \beta_{\lambda / k}
$$

of weighted sums of canonical Liouville measures $\beta_{\lambda / k}$ on the coadjoint orbits $G \lambda / k$.

The aim of this article is to study the asymptotic behaviour of $\Theta_{k}^{(M, \mathcal{L})}$ when $k$ is large, and $M$ possibly non compact, and to explore a functorial consequence of this formula for reduced spaces.

We work in the spin context. The same argument would adapt to manifolds with spinc structures, provided we work with odd powers of $k$. This context is more general than the Hamiltonian context (which includes the Kähler context), and it is more natural since Weyl character formula for $\chi_{\lambda}$ is the fixed point formula for a twisted Dirac operator on the coadjoint orbit $G \lambda$. We will return to the comparison between both contexts in a forthcoming article. The article [24] determines the asymptotic expansion of $\Theta_{k}^{(M, \mathcal{L})}$ in the Hamiltonian context, when $G$ is a torus.

Let us return to the case of a spin manifold $M$ with a $G$-invariant line bundle with connection. We do not assume any more $M$ compact, but we assume the moment map $\phi_{G}$ is proper. In this introduction, to simplify, assume that all stabilizers of the action of $G$ on $M$ are abelian and connected. We also assume that $\rho$ is in the weight lattice of $T$. For every $\lambda \in \widehat{G} \subset \widehat{T} \subset \mathfrak{t}^{*}$, consider the reduced space

$$
M_{\lambda, G}:=\phi_{G}^{-1}(\lambda) / T
$$

which may be non connected. If $\lambda$ is a regular value of $\phi_{G}, M_{\lambda, G}$ is a smooth compact spin manifold. Denote by $\mathbb{C}_{[-\lambda]}$ the vector space $\mathbb{C}$ with the action of $T$ given by the character $t^{-\lambda}$ of $T$. The restriction of the line bundle $\mathbb{C}_{[-\lambda]} \otimes \mathcal{L}^{k}$ to $\phi_{G}^{-1}(\lambda)$ is the pull back of a line bundle $\mathcal{L}_{\lambda, k}$ on $M_{\lambda, G}$. So define

$$
m_{G}^{\text {geo }}(\lambda, k)=\mathcal{Q}^{\text {geo }}\left(M_{\lambda, G}, \mathcal{L}_{\lambda, k}\right)
$$

If $\lambda$ is not a regular value of $\phi_{G}$, this number can be defined by deformation. When $M$ is compact, the $[Q, R]=0$ theorem in the spin context [21] asserts that $m_{G}(\lambda, k)=m_{G}^{\text {geo }}(\lambda, k)$. This suggests to extend the
definition of $\mathcal{Q}_{G}^{\text {spin }}(M, \mathcal{L})$, when $\mathcal{L}$ is equipped with a connection (that we leave implicit in the notation) such that the moment map $\phi_{G}$ is proper, in the following way. Define

$$
\mathcal{Q}_{G}^{\text {geo }}\left(M, \mathcal{L}^{k}\right):=\sum_{\lambda} m_{G}^{\text {geo }}(\lambda, k) V_{\lambda}
$$

We call this object (in the spirit of [25]) the formal geometric quantization of $(M, \mathcal{L})$.

Here are two simple examples.
Example 1.1. - $M=T^{*} S^{1}$, with coordinates $(t, \theta)$, and the natural action of the circle group $G=S^{1}$. Consider the trivial line bundle $\mathcal{L}$ with connection $d-i t d \theta$. The moment map is $\phi_{G}(t, \theta)=t$. Thus $\mathcal{Q}_{G}^{\text {geo }}\left(M, \mathcal{L}^{k}\right)$ is independent of $k$ :

$$
\mathcal{Q}_{G}^{\mathrm{geo}}\left(M, \mathcal{L}^{k}\right)=\sum_{n \in \mathbb{Z}} e^{i n \theta}
$$

- $M=\mathbb{C}$, with coordinate $z$. We consider an integer a and the line bundle $\mathcal{L}=M \times \mathbb{C}$ with action $e^{i \theta}(z, v)=\left(e^{2 i \theta} z, e^{i a \theta} v\right)$. Take the connection $\nabla=d-\frac{i}{2} \operatorname{Im}(z d \bar{z})$. The moment map is $\phi_{G}(z)=a+|z|^{2}$ and is proper. Then

$$
\mathcal{Q}_{G}^{\text {geo }}\left(M, \mathcal{L}^{k}\right)=e^{i k a \theta} \sum_{j \geq 0} e^{i(2 j+1) \theta}
$$

As $M$ is non compact, it is not possible to define the index of $D_{\mathcal{L}}$ without introducing additional data. We use transversally elliptic operators. The use of the Kirwan vector field $\kappa_{G}$ associated to $\phi_{G}$ allows us to deform the symbol of the Dirac operator, and to obtain a $G$ invariant transversally elliptic operator $D_{\mathcal{L}, \phi_{G}}$ on $M$ if the set of zeroes of $\kappa_{G}$ is compact, which we assume in this introduction. The index of this operator $\mathcal{Q}_{G}^{-\infty}\left(M, \mathcal{L}^{k}\right)$ provides a well defined generalized function on $G$. So define $m_{G}(\lambda, k) \in \mathbb{Z}$ such that

$$
\mathcal{Q}_{G}^{-\infty}\left(M, \mathcal{L}^{k}\right)(g)=\sum_{\lambda \in \hat{G}} m_{G}(\lambda, k) \chi_{\lambda}(g)
$$

Then, we have again [10]

$$
m_{G}(\lambda, k)=m_{G}^{\text {geo }}(\lambda, k)
$$

An important consequence of this geometric relation is the fact that the function $m_{G}(\lambda, k)$ is a piecewise quasi-polynomial function. In particular, if $m_{G}(k \lambda, k)=0$ for $k$ large, then $m_{G}(k \lambda, k)=0$ for all $k \geq 1$.

We define as in the compact case the weighted sum of measures of the orbits $G \lambda / k$ :

$$
\Theta_{k}^{(M, \mathcal{L})}=k^{\operatorname{dim}(G / T) / 2} \sum_{\lambda} m_{G}(\lambda, k) \beta_{\lambda / k}
$$

As it is well known, at least in the Hamiltonian case, the measure $\frac{1}{k^{\text {dim } M / 2}} \Theta_{k}^{(M, \mathcal{L})}$ tends to the Duistermaat-Heckman measure when $k$ tends to $\infty$. Here we will prove that $\Theta_{k}^{(M, \mathcal{L})}$ has an asymptotic expansion, when $k \rightarrow \infty$, as a Laurent series (in $1 / k$ ) of distributions, and we will determine all coefficients as twisted Duistermaat-Heckman distributions related to the $\hat{A}$ class of $M$.

Recall in this spin context the definition of the Duistermaat-Heckman measure. Let $\Omega(X)=\Omega-\left\langle\phi_{G}, X\right\rangle, X \in \mathfrak{g}$. This is a closed equivariant form of degree 2 on $M$ and $e^{-i \Omega(X)}$ is the equivariant Chern character of $\mathcal{L}$. The form $\Omega^{\operatorname{dim} M / 2}$ is a density on $M$ (which might be not positive). The Duistermaat-Heckman measure is the signed measure on $\mathfrak{g}^{*}$ obtained by push-forward of $\Omega^{\operatorname{dim} M / 2}$ by the proper map $\phi_{G}$. If $M$ is compact, its Fourier transform is the function $X \mapsto \frac{1}{(-2 i \pi)^{d}} \int_{M} e^{-i \Omega(X)}$ on $\mathfrak{g}$. Similarly, if $\nu(X)$ is a closed $G$-equivariant forms on $M$, with polynomial coefficients, we can define the distribution $D H^{G}(M, \Omega, \nu)$ on $\mathfrak{g}^{*}$ by the formula:

$$
\begin{equation*}
\left\langle D H^{G}(M, \Omega, \nu), \varphi\right\rangle=\frac{1}{(-2 i \pi)^{d}} \int_{M} \int_{\mathfrak{g}} e^{-i \Omega(X)} \nu(X) \widehat{\varphi}(X) d X \tag{1.1}
\end{equation*}
$$

Here $\varphi$ is a test function on $\mathfrak{g}^{*}$, with Fourier transform $\widehat{\varphi}$. It is easy to see that this distribution is well defined if $\phi_{G}$ is proper.

Consider the equivariant $\hat{A}$ class of $M$ which belongs to the completion of the equivariant cohomology ring of $M$, and its expansion $\hat{A}(M)(X)=\sum_{n=0}^{\infty} \hat{A}_{n}(M)(X)$ in equivariant classes homogeneous of degree $2 n$. Finally, let $j_{\mathfrak{g}}^{1 / 2}(X)=\operatorname{det}_{\mathfrak{g}}^{1 / 2}\left(\frac{e^{X / 2}-e^{-X / 2}}{X}\right)$, a $G$-invariant function of $X \in \mathfrak{g}$. It determines a formal series $j_{\mathfrak{g}}^{1 / 2}(i \partial / k)$ of $G$ invariant constant coefficient differential operators on $\mathfrak{g}^{*}$.

Our main contribution is the following result (see Theorem 4.3).
Theorem 1.2. When $k$ tends to $\infty$,

$$
\begin{equation*}
\Theta_{k}^{(M, \mathcal{L})} \equiv j_{\mathfrak{g}}^{1 / 2}(i \partial / k)\left(k^{d} \sum_{n=0}^{\infty} \frac{1}{k^{n}} D H^{G}\left(M, \Omega, \hat{A}_{n}(M)\right)\right) \tag{1.2}
\end{equation*}
$$

The leading term in $k^{d}$ is the Duistermaat-Heckman measure.
Let us check this formula in our two basic examples.

Example 1.3. - For $M=T^{*} S^{1}$, the $\hat{A}$ class is identically equal to 1 , so there exists only one term in the expansion (1.2). Here Theorem 1.2 is the well known fact that

$$
\sum_{\lambda \in \mathbb{Z}} \varphi(\lambda / k) \equiv k \int_{\mathbb{R}} \varphi(\xi) d \xi
$$

- For $M=\mathbb{C}$, the equivariant $\hat{A}$ class for our action is the invariant function $\frac{\theta}{\sin (\theta)}$. Theorem 1.2 is the following variation of the EulerMacLaurin formula (see [7])

$$
\sum_{j \geq 0} \varphi(a+(2 j+1) / k) \equiv \frac{k}{2} \int_{a}^{\infty} \varphi(\xi) d \xi-\sum_{n=1}^{\infty}\left(\frac{2}{k}\right)^{n-1} \frac{B_{n}(1 / 2)}{n!} \varphi^{(n-1)}(a)
$$

Here $B_{n}(t)$ is the $n$-th Bernoulli polynomial and only even integers $n$ occur, since $B_{n}\left(\frac{1}{2}\right)=0$ for all odd $n \geq 0$.

Let us see why this formula is natural, when $M$ is compact. By Kirillov formula, the Fourier transform of $\Theta_{k}^{(M, \mathcal{L})}$ is the function

$$
X \longmapsto \mathcal{Q}_{G}^{\text {spin }}\left(M, \mathcal{L}^{k}\right)(\exp X / k) j_{\mathfrak{g}}^{1 / 2}(X / k)
$$

on $\mathfrak{g}$. The delocalized formula [2] for the index says that, for $X \in \mathfrak{g}$ small enough,

$$
\mathcal{Q}_{G}^{\mathrm{spin}}\left(M, \mathcal{L}^{k}\right)(\exp X / k)=\frac{1}{(-2 i \pi)^{d}} \int_{M} e^{-i k \Omega(X / k)} \hat{A}(M)(X / k)
$$

But, by simple inspection of the relation between the equivariant degree and the polynomial degree, this is also equal for $X$ small to

$$
\frac{1}{(-2 i \pi)^{d}} k^{d} \int_{M} e^{-i \Omega(X)}\left(\sum_{n=0}^{\infty} \frac{1}{k^{n}} \hat{A}_{n}(M)(X)\right)
$$

The formula for the asymptotic expansion of $\Theta_{k}^{(M, \mathcal{L})}$ follows easily. Remark the dichotomy between the equivariant Chern character $e^{-i \Omega(X)}$ that is unchanged in this asymptotic equivalence, and involves the exponential function $e^{i\left\langle\phi_{G}, X\right\rangle}$, and the $\hat{A}$ class that we expand as a series of homogeneous equivariant classes with polynomial coefficients. A more delicate analysis, using the delocalized formula ([20]) for the generalized function $\mathcal{Q}_{G}^{-\infty}\left(M, \mathcal{L}^{k}\right)(\exp X)$ leads us also naturally to Theorem 1.2. The Chern character $e^{-i \Omega(X)}$ has to be replaced by the Chern character $e^{-i \Omega(X)} P(X)$ where $P(X)$ is an equivariant form with generalized coefficients, supported in a compact neighborhood of the zeroes of $\kappa_{G}$, equivalent to 1 in the equivariant cohomology without compact support conditions.

In conclusion, Theorem 1.2 says that the formula

$$
\mathcal{Q}_{G}^{\text {spin }}(M, \mathcal{L})(\exp X)=\frac{1}{(-2 i \pi)^{d}} \int_{M} e^{i \Omega(X)} \hat{A}(M)(X)
$$

has a meaning in the asymptotic sense for a non compact manifold $M$ when $\mathcal{L}$ is replaced by $\mathcal{L}^{k}$ and $X$ by $X / k$, and is a good ersatz for the equivariant index formula of $D_{\mathcal{L}}$.

Let us now explain a consequence of this asymptotic formula for reduced spaces.

Return for a moment to the case where $M$ is compact. Plugging $g=1$ in the identity $\mathcal{Q}_{G}^{\text {spin }}\left(M, \mathcal{L}^{k}\right)(g)=\sum_{\lambda \in \hat{G}} m_{G}(\lambda, k) \chi_{\lambda}(g)$, leads to the remarkable identity

$$
\begin{equation*}
\mathcal{Q}^{\mathrm{geo}}\left(M, \mathcal{L}^{k}\right)=\sum_{\lambda \in \widehat{G}} \operatorname{vol}(G \lambda) \mathcal{Q}^{\mathrm{geo}}\left(M_{\lambda, G}, \mathcal{L}_{\lambda, k}\right) \tag{1.3}
\end{equation*}
$$

relating an integral on $M$ to a sum of integrals on the finite number of reduced spaces $M_{\lambda, G}$. We will see that this equality generalizes to reduction in stages, even when $M$ is non compact.

Let $H$ be a connected compact subgroup of $G$ with torus $T_{H}$. The $H$ action on $(M, \mathcal{L})$ leads to the moment map $\phi_{H}: M \rightarrow \mathfrak{h}^{*}$. Assume $\phi_{H}$ is still proper. In this case we can define $\mathcal{Q}_{G}^{-\infty}(M, \mathcal{L})$, and $\mathcal{Q}_{H}^{-\infty}(M, \mathcal{L})$. When $M$ is compact, from the description of $\mathcal{Q}_{G}^{\text {spin }}(M, \mathcal{L})$ as the index of the elliptic operator $D_{\mathcal{L}}$ on $M$, we see that $\mathcal{Q}_{H}^{\text {spin }}(M, \mathcal{L})$ is the restriction of the representation $\mathcal{Q}_{G}^{\text {spin }}(M, \mathcal{L})$ to $H$. When $M$ is non compact, this relation is not obvious since our $G$-transversally elliptic symbol $D_{\mathcal{L}, \phi_{G}}$ is not (usually) $H$-transversally elliptic. The following theorem

$$
\begin{equation*}
\left.\mathcal{Q}_{G}^{-\infty}(M, \mathcal{L})\right|_{H}=\mathcal{Q}_{H}^{-\infty}(M, \mathcal{L}) \tag{1.4}
\end{equation*}
$$

was proved by the first author using cutting [19]. Here we will show that this also follows from comparing the large behavior of $\mathcal{Q}_{G}^{-\infty}\left(M, \mathcal{L}^{k}\right)$ and $\mathcal{Q}_{H}^{-\infty}\left(M, \mathcal{L}^{k}\right)$ and the fact that multiplicities $m_{G}(k \lambda, k), m_{H}(k \lambda, k)$ are entirely determined by their large values.

Let $\lambda \in \widehat{G}$, and $\mu \in \widehat{H}$. Denote by $c(\lambda, \mu)$ the multiplicity of $V_{\mu}$ in $\left.V_{\lambda}\right|_{H}$. Theorem (1.4) is equivalent of the following equality of indices of Dirac operators on reduced spaces for $H$, and $G$.

For any $\mu \in \widehat{H}$, we have

$$
\mathcal{Q}^{\mathrm{geo}}\left(M_{\mu, H}, \mathcal{L}_{\mu, k}\right)=\sum_{\lambda \in \widehat{G}} c(\lambda, \mu) \mathcal{Q}^{\mathrm{geo}}\left(M_{\lambda, G}, \mathcal{L}_{\lambda, k}\right) .
$$

When $M$ is compact, we can take $H=\{1\}$, and this is the mysterious equality (1.3). In general, this equality of indices of Dirac
operators on reduced spaces for different groups is not clear, since $M_{\mu, H}=\phi_{H}^{-1}(\mu) / T_{H}$ do not carry any visible group action.

Let us briefly sketch a proof of this equality for the case where $G, H$ are torus. This simple case could be treated by considering the action of $G / H$ on the compact spin manifold $M_{\mu, H}$, but we treat it by a different method which will generalize to any pair of groups $H \subset G$.

In this case, $\widehat{G}, \widehat{H}$ are lattices in $\mathfrak{g}^{*}, \mathfrak{h}^{*}$, and we have to prove

$$
m_{H}(\mu, k)=\sum_{\lambda \in \overparen{G},\left.\lambda\right|_{\mathfrak{h}}=\mu} m_{G}(\lambda, k) .
$$

In a companion article [23], we proved the following easy result. Consider a distribution $D_{k}=\sum_{\mu \in \widehat{H}} q(\mu, k) \delta_{\mu / k}$ associated to a piecewise quasi-polynomial function $q(\mu, k)$. Assume that, for any $\zeta \in H$, an element of finite order, the distribution $\sum_{\mu \in \hat{H}} q(\mu, k) \zeta^{\mu} \delta_{\mu / k}=O\left(k^{-\infty}\right)$ when $k$ tends to $\infty$. Then $q(\mu, k)=0$.

We have computed the asymptotic expansion of both distributions $\Theta_{k}^{G}=\sum_{\lambda \in \hat{G}} m_{G}(\lambda, k) \delta_{\lambda / k}$ and $\Theta_{k}^{H}=\sum_{\mu \in \hat{H}} m_{H}(\mu, k) \delta_{\mu / k}$. One has:

$$
\begin{aligned}
\Theta_{k}^{G} & \equiv k^{d} \sum_{n=0}^{\infty} \frac{1}{k^{n}} D H^{G}\left(M, \Omega, \hat{A}_{n}(M)\right), \\
\Theta_{k}^{H} & \equiv k^{d} \sum_{n=0}^{\infty} \frac{1}{k^{n}} D H^{H}\left(M, \Omega, \hat{A}_{n}(M)\right) .
\end{aligned}
$$

Define $m_{H}^{\prime}(\mu, k)=\sum_{\lambda \in \widehat{G},\left.\lambda\right|_{\mathfrak{h}}=\mu} m_{G}(\lambda, k)$. The distribution

$$
S_{k}^{H}=\sum_{\mu \in \hat{H}} m_{H}^{\prime}(\mu, k) \delta_{\mu / k}
$$

is the push-forward $r_{*} \Theta_{k}^{G}$ of the distribution $\Theta_{k}^{G}$ under the map $r$ : $\mathfrak{g}^{*} \rightarrow \mathfrak{h}^{*}$.

The Duistermaat-Heckman distributions $D H^{G}(M, \Omega, \nu)$ behave very well under the push-forward map:

$$
r_{*} D H^{G}(M, \Omega, \nu)=D H^{H}(M, \Omega, \nu) .
$$

Indeed, at least in the compact case, the Fourier transform of $D H^{H}(M, \Omega, \nu)$ is the restriction to $\mathfrak{h}$ of the function $\frac{-1}{(2 \pi)^{d}} \int_{M} e^{-i \Omega(X)} \nu(X)$ on $\mathfrak{g}$. So we see $S_{k}^{H}$ and $\Theta_{k}^{H}$ have the same asymptotic expansion. This is not entirely sufficient to prove that $\Theta_{k}^{H}=S_{k}^{H}$. But we use the fact that both functions $m_{H}^{\prime}(\mu, k)$ and $m_{H}(\mu, k)$ are piecewise quasi-polynomials and
a similar asymptotic descent formula holds for the distribution

$$
\sum_{\mu \in \widehat{H}} m_{H}(\mu, k) \zeta^{\mu} \delta_{\mu / k}
$$

when $\zeta \in T_{H}$ is of finite order.
The proof, sketched here for $H$ and $G$ abelian, works equally well for two subgroups $H \subset G$. We just have to use the formula for the push-forward of the Liouville measure $r_{*}\left(\beta_{\lambda}\right)$ of admissible coadjoint orbits. As expected, the full series for the $\hat{A}(M)$ class, as well as the Duflo operator $j_{\mathfrak{g}}^{1 / 2}(\partial)$ plays a role in the functoriality.

Asymptotic behavior of quantization when $k$ tends to $\infty$ has been considered by many authors. Let us give a very limited and personal selection of influential works.

Asymptotic behavior of kernels of Laplacians twisted by $\mathcal{L}^{k}$ were used by Boutet de Monvel-Guillemin [6] to produce a formal star-product of functions on symplectic manifolds. In general only the few first terms of the star product formal deformation are computable (see for example [5], [13]).

Our asymptotic trace formula for the transversally operator $D_{\mathcal{L}, \phi_{G}}$ has the same flavor than Fedosov trace formula for deformation quantization [8].

The article of E. Meinrenken [15] where multiplicities $m_{G}(\lambda, k)$ are identified at the large limit to index of reduced spaces is in close relation with our setting. Asymptotic Riemann sums of values of smooth functions at equally spaced sample points of a Delzant polytope $\Delta \subset \mathfrak{g}^{*}$ were studied by Guillemin-Sternberg [9], where the full asymptotic formula is given in terms of the Todd class of the corresponding toric manifold. The natural geometric rescaling $\mu \rightarrow \mu / k$ in the computation of $\frac{1}{k^{d}} \sum_{\mu \in \mathbb{Z}^{d} \cap k \Delta} \varphi(\mu / k)$ leads to consider rescaling $X \rightarrow X / k$ in the variable $X \in \mathfrak{g}$. This was one of the inspiring examples.

Our application to restrictions to subgroups is a striking example where an obvious property at the semi-classical level (functoriality of push-forward of distributions) can lead to a proof of a subtle relation at the quantum level.

## 2. Asymptotics and representations

2.1. Fourier transforms. When $V$ is a finite dimensional real vector space, we denote $\mathcal{D}^{\prime}\left(V^{*}\right)$ the vector space of complex valued distributions on the dual vector space $V^{*}$. If $A \in \mathcal{D}^{\prime}\left(V^{*}\right)$ and $\varphi \in \mathcal{C}_{c p t}^{\infty}\left(V^{*}\right)$, we denote $\langle A, \varphi\rangle \in \mathbb{C}$ (or $\langle A(\xi), \varphi(\xi)\rangle)$ the value of $A$ on $\varphi$.

When $A \in \mathcal{D}^{\prime}\left(V^{*}\right)$ has a compact support, its Fourier transform $\mathcal{F}_{V^{*}}(A) \in \mathcal{C}^{\infty}(V)$ is defined by the relation

$$
\mathcal{F}_{V^{*}}(A)(v):=\left\langle A(\xi), e^{i\langle\xi, v\rangle}\right\rangle, \quad v \in V
$$

When $f$ is smooth and has at most polynomial growth, its Fourier transform $\mathcal{F}_{V}(f) \in \mathcal{D}^{\prime}\left(V^{*}\right)$ is defined by the relation

$$
\left\langle\mathcal{F}_{V}(f), \varphi\right\rangle=\int_{V} f(v) \widehat{\varphi}(v) d v
$$

where $\widehat{\varphi}(v)=\frac{1}{(2 \pi)^{\operatorname{dim} V}} \int_{V^{*}} \varphi(\xi) e^{-i\langle\xi, v\rangle} d \xi$. Here $d v$ and $d \xi$ are dual volume forms on $V$ and $V^{*}$.

For any $k \geq 1$, we define the operation

$$
\begin{equation*}
\operatorname{rescaling}_{V^{*}}(k): \mathcal{D}^{\prime}\left(V^{*}\right) \rightarrow \mathcal{D}^{\prime}\left(V^{*}\right) \tag{2.1}
\end{equation*}
$$

by $\left\langle\operatorname{rescaling}_{V^{*}}(k)(A), \varphi\right\rangle=\langle A(\xi), \varphi(\xi / k)\rangle$. We have also the operation rescaling ${ }_{V}(k): \mathcal{C}^{\infty}(V) \rightarrow \mathcal{C}^{\infty}(V)$ defined by rescaling ${ }_{V}(k)(f)(v)=$ $f(v / k)$.
2.2. The isomorphism $\mathcal{R}_{\mathfrak{g}}$. Let $G$ be a connected compact Lie group with Lie algebra $\mathfrak{g}$. Let $T \subset G$ be a maximal torus with Lie algebra $\mathfrak{t}$. Let $W_{G}$ be the Weyl group. We consider $\widehat{T}$ as a lattice $\Lambda$ in $\mathfrak{t}^{*}$. If $\lambda \in \Lambda$, we denote by $t^{\lambda}$ the corresponding character of $T$. If $t=\exp (X)$ with $X \in \mathfrak{t}$, then $t^{\lambda}=e^{i\langle\lambda, X\rangle}$. We choose a system $\Delta^{+}$of positive roots. In our convention, $\Delta^{+}$is contained in $\mathfrak{t}^{*}$. We choose an invariant scalar product in $\mathfrak{k}$.

Let $\mathcal{D}^{\prime}\left(\mathfrak{g}^{*}\right)^{G}$ be the space of $G$-invariant distributions on $\mathfrak{g}^{*}$ and $\mathcal{D}^{\prime}\left(\mathfrak{t}^{*}\right)^{W_{G}-\text { alt }}$ be the subspace of $W_{G}$ anti-invariant distributions on $\mathfrak{t}^{*}$. Let $\Pi_{\mathfrak{g} / \mathfrak{t}}(X)=\prod_{\alpha \in \Delta^{+}}\langle\alpha, X\rangle$ a $W_{G}$ anti-invariant function on $\mathfrak{t}$.

We will use the following classical fact.
Lemma 2.1. There exists a unique linear isomorphism

$$
\mathcal{R}_{\mathfrak{g}}: \mathcal{D}^{\prime}\left(\mathfrak{g}^{*}\right)^{G} \longrightarrow \mathcal{D}^{\prime}\left(\mathfrak{t}^{*}\right)^{W_{G}-\text { alt }}
$$

with the following two properties

- $\mathcal{R}_{\mathfrak{g}}(f A)=\left.f\right|_{\mathfrak{t}^{*}} \mathcal{R}_{\mathfrak{g}}(A)$ for all $A \in \mathcal{D}^{\prime}\left(\mathfrak{g}^{*}\right)^{G}$ and $f \in \mathcal{C}^{\infty}\left(\mathfrak{g}^{*}\right)^{G}$,
- For compactly supported distributions $A$, the Fourier transform of $\mathcal{R}_{\mathfrak{g}}(A)$ is given by

$$
\left\langle\mathcal{R}_{\mathfrak{g}}(A), e^{i\langle-, X\rangle}\right\rangle=(i)^{r} \Pi_{\mathfrak{g} / \mathfrak{t}}(X)\left\langle A, e^{i\langle-, X\rangle}\right\rangle, \quad X \in \mathfrak{t}
$$

with $r=\operatorname{dim}(G / T) / 2$.
2.3. Asymptotics. Let $\mathcal{D}^{\prime}\left(\mathfrak{g}^{*}\right)^{G}$ be the space of $G$-invariant distributions on $\mathfrak{g}^{*}$. We will study asymptotic expansions of sequences of distributions.

Definition 2.2. Let $\left(\Theta_{k}\right)_{k \geq 1}$ and $\left(\theta_{n}\right)_{n \in \mathbb{N}}$ be two sequences in $\mathcal{D}^{\prime}\left(\mathfrak{g}^{*}\right)^{G}$. We write

$$
\begin{equation*}
\Theta_{k} \equiv k^{n_{o}} \sum_{n=0}^{\infty} \frac{1}{k^{n}} \theta_{n} \tag{2.2}
\end{equation*}
$$

for some $n_{o} \in \mathbb{Z}$ if for any test function $\varphi$ and any $N \in \mathbb{N}$, we have

$$
\left\langle\Theta_{k}, \varphi\right\rangle=k^{n_{o}} \sum_{n=0}^{N} \frac{1}{k^{n}}\left\langle\theta_{n}, \varphi\right\rangle+o\left(k^{n_{o}-N}\right)
$$

In this text, we will allow each distribution $\theta_{n}$ to depends periodically of the integers $k$. This means that for each $n$ there exists $D \geq 1$ and a family of distribution $\left(\theta_{\zeta}\right)$ parametrized by the roots of unity $\left\{\zeta^{D}=1\right\}$ such that $\theta_{n}(k)=\sum_{\zeta^{D}=1} \zeta^{k} \theta_{\zeta}$ for all $k \geq 1$.
Example 2.3. Consider the distributions $\mathrm{T}_{k}=\sum_{\lambda \in \mathbb{Z}, \lambda \geq 0} \delta_{\lambda / k}(k \geq 1)$ on $\mathbb{R}$. The Euler Maclaurin formula gives the expansion

$$
\mathrm{T}_{k} \equiv k 1_{[0, \infty[ }+\frac{1}{2} \delta_{0}-\sum_{n=1}^{\infty} \frac{1}{k^{2 n-1}} \frac{b_{2 n}}{(2 n)!} \delta_{0}^{(2 n-1)}
$$

where $b_{2 n}$ are the Bernoulli numbers and $\delta_{0}^{(2 n-1)}$ is the $(2 n-1)$-th derivation of the Dirac distributions $\delta_{0}$.

Let $d$ be a $C^{\infty}$ function defined near the origin $0 \in \mathfrak{g}$. Consider the Taylor series $\sum_{n=0}^{\infty} d_{n}(X)$ of $d$ at the origin. Thus $d_{n}$ is an homogeneous polynomial function of degree $n$ on $\mathfrak{g}$. We associate to the function $d$ the formal series

$$
d(i \partial / k)=\sum_{n=0}^{\infty} k^{-n} d_{n}(i \partial)
$$

of constant coefficient differential operators $d_{n}(i \partial)$ on $\mathfrak{g}^{*}$. Thus if $Q(k)=k^{n_{o}} \sum_{n=0}^{\infty} k^{-n} \theta_{n}$ is a formal series of distributions on $\mathfrak{g}^{*}$, the series $d(i \partial / k) Q(k)$ is the formal series of distributions $k^{n_{o}} \sum_{n=0}^{\infty} k^{-n} s_{n}$ on $\mathfrak{g}^{*}$, where $s_{n}=\sum_{l+m=n} d_{l}(i \partial) \theta_{m}$.

We have the following basic fact.
Lemma 2.4. The expansion $\Theta_{k} \equiv k^{n_{o}} \sum_{n=0}^{\infty} \frac{1}{k^{n}} \theta_{n}$ holds in $\mathcal{D}^{\prime}\left(\mathfrak{g}^{*}\right)^{G}$ if and only if we have the expansion $\mathcal{R}_{\mathfrak{g}}\left(\Theta_{k}\right) \equiv k^{n_{o}} \sum_{n=0}^{\infty} \frac{1}{k^{n}} \mathcal{R}_{\mathfrak{g}}\left(\theta_{n}\right)$ in $\mathcal{D}^{\prime}\left(\mathfrak{t}^{*}\right)^{W_{G}-a l t}$.

Our asymptotic expansions will arise from Taylor series.
Let $\tau$ be a smooth function on $\mathfrak{g}$. Assume that $\tau$ and all its derivatives have at most polynomial growth. Consider the Taylor series $\sum_{n=0}^{\infty} \tau_{n}(X)$ of $\tau$ at the origin. If $b(X)$ is a smooth function on $\mathfrak{g}$ with at most polynomial growth, then $f_{k}(X)=b(X) \tau(X / k)(k \geq 1)$ defines a tempered distribution on $\mathfrak{g}$. Similarly, as $\tau_{n}(X)$ is a polynomial function, $g_{n}(X)=b(X) \tau_{n}(X)$ defines a tempered distribution on $\mathfrak{g}$. We note the following result.

Proposition 2.5. When $k$ tends to $\infty$, we have the asymptotic expansion

$$
\mathcal{F}_{\mathfrak{g}}\left(f_{k}\right) \equiv \sum_{n=0}^{\infty} \frac{1}{k^{n}} \mathcal{F}_{\mathfrak{g}}\left(g_{n}\right)=\tau(i \partial / k) \mathcal{F}_{\mathfrak{g}}(b)
$$

Later, we will need a uniform version with parameters of this proposition. So we give the proof.

Proof. Let $\varphi$ be a smooth function on $\mathfrak{g}^{*}$ with compact support. Thus $\left\langle\mathcal{F}_{\mathfrak{g}}\left(f_{k}\right), \varphi\right\rangle=\int_{\mathfrak{g}} b(X) \tau(X / k) \widehat{\varphi}(X) d X$.

Let $N$ be a positive integer, and $\tau_{\leq N}$ be the sum of the Taylor coefficients $\tau_{n}$, up to degree $N$, so $\tau=\tau_{\leq N}+\tau_{>N}$. We write $I_{k}=$ $\int_{\mathfrak{g}} b(X) \tau(X / k) \widehat{\varphi}(X) d X$ as $I_{k}^{0}+R_{k}$ with

$$
I_{k}^{0}=\int_{\mathfrak{g}} b(X) \tau_{\leq N}(X / k) \widehat{\varphi}(X) d X=\sum_{k=0}^{N} \frac{1}{k^{n}} \int_{\mathfrak{g}} b(X) \tau_{n}(X) \widehat{\varphi}(X) d X
$$

and $R_{k}=\int_{\mathfrak{g}} b(X) \tau_{>N}(X / k) \widehat{\varphi}(X) d X$.
In multi-index notation, we may write (via an integral formula depending of $\tau$ and its derivatives) $\tau_{>N}(X)=\sum_{\alpha,|\alpha|=N+1} X^{\alpha} D_{\alpha}(X)$ with $D_{\alpha}(X)$ bounded by a polynomial function of $X$. If $\tau(X)$ depends smoothly of a parameter $x$, we can construct $D_{\alpha}(X)$ depending smoothly of $x$. Thus

$$
R_{k}=\frac{1}{k^{N+1}} \sum_{\alpha,|\alpha|=N+1} \int_{\mathfrak{g}} b(X) X^{\alpha} D_{\alpha}(X / k) \widehat{\varphi}(X) d X
$$

Since $\|X / k\| \leq\|X\|$, we can bound $b(X) X^{\alpha} D_{\alpha}(X / k)$ by a polynomial function of $X$. Since $\widehat{\varphi}(X)$ is rapidly decreasing, we see that $\left|R_{k}\right| \leq$ $\frac{c_{N}}{k^{N+1}}$.

It is clear that if $b(X, k)=\sum_{m \in F} k^{m} b_{m}(X)$ is a sum over a finite set $F \subset \mathbb{Z}$ of smooth functions $b_{m}(X)$ with at most polynomial growth, and $\tau(X)$ a function with polynomial growth as well as all its derivatives, we can obtain the asymptotic expansion of $\int_{\mathfrak{g}} b(X, k) \tau(X / k) \widehat{\varphi}(X) d X$ by
summing over $m \in F$ the asymptotic expansions of $\int_{\mathfrak{g}} b_{m}(X) \tau(X / k) \widehat{\varphi}(X) d X$ multiplied by $k^{m}$. So this is given by the Laurent series (in $1 / k$ )

$$
\sum_{m \in F} \sum_{n=0}^{\infty} k^{m} \frac{1}{k^{n}} \int_{\mathfrak{g}} b_{m}(X) \tau_{n}(X) \widehat{\varphi}(X) d X
$$

We write this somewhat informally as

$$
\begin{equation*}
\int_{\mathfrak{g}} b(X, k) \tau(X / k) \widehat{\varphi}(X) d X \equiv \int_{\mathfrak{g}} b(X, k) \sum_{n=0}^{\infty} \tau_{n}(X / k) \widehat{\varphi}(X) d X \tag{2.3}
\end{equation*}
$$

In short, we replace $\tau$ by its formal Taylor series $\sum_{n=0}^{\infty} \tau_{n}(X)$ and keep $b(X, k)$ as it is.
2.4. Kirillov formula. Recall that any coadjoint orbit $\mathcal{O} \subset \mathfrak{g}^{*}$ is provided with the KKS symplectic form $\Omega_{\mathcal{O}}$ that is normalized as follows. For $\xi \in \mathcal{O}$ and $X, Y \in \mathfrak{g}$, we have $\left.\Omega_{\mathcal{O}}\right|_{\xi}(X \cdot \xi, Y \cdot \xi)=\langle\xi,[X, Y]\rangle$. Let $n_{\mathcal{O}}=\frac{1}{2} \operatorname{dim} \mathcal{O}$.

Definition 2.6. We denote $\beta_{\mathcal{O}}$ the distribution on $\mathfrak{g}^{*}$ defined by the relation

$$
\left\langle\beta_{\mathcal{O}}, \varphi\right\rangle=\int_{\xi \in \mathcal{O}} \varphi(\xi) \frac{\left(\Omega_{\mathcal{O}}\right)^{n_{\mathcal{O}}}}{(2 \pi)^{n_{\mathcal{O}}} n_{\mathcal{O}}!}
$$

for any smooth function $\varphi$ on $\mathfrak{g}^{*}$.
We have chosen a system $\Delta^{+} \subset \mathfrak{t}^{*}$ of positive roots, and let $\rho \in \mathfrak{t}^{*}$ be the corresponding element. We consider the positive Weyl chamber $\mathfrak{t}_{\geq 0}^{*}$ with interior $\mathfrak{t}_{>0}^{*}$. Identify $\widehat{G}$ to the discrete set $A_{G}=(\rho+\Lambda) \cap \mathfrak{t}_{>0}^{*}$ of positive admissible regular elements of $\mathfrak{t}^{*}$.

If $\lambda \in A_{G}$, we denote by $\chi_{\lambda}(g)$ the trace of the action of $g \in G$ on the irreducible representation $V_{\lambda}$ parametrized by $\lambda$. We have the fundamental identity

$$
\begin{equation*}
\chi_{\lambda}\left(e^{X}\right) j_{\mathfrak{g}}^{1 / 2}(X)=\left\langle\beta_{\lambda}, e^{i\langle-, X\rangle}\right\rangle \tag{2.4}
\end{equation*}
$$

where $j_{\mathfrak{g}}(X)=\operatorname{det}_{\mathfrak{g}}\left(\frac{e^{X / 2}-e^{-X / 2}}{X}\right)$, and $\beta_{\lambda}:=\beta_{G \lambda}$.
Let $\mathcal{R}_{\mathfrak{g}}: \mathcal{D}^{\prime}\left(\mathfrak{g}^{*}\right)^{G} \longrightarrow \mathcal{D}^{\prime}\left(\mathfrak{t}^{*}\right)^{W_{G}-\text { alt }}$ be the isomorphism introduced in section 2.2. If $\lambda \in A_{G}$, we have

$$
\begin{equation*}
\mathcal{R}_{\mathfrak{g}}\left(\beta_{\lambda}\right)=\sum_{w \in W_{G}} \epsilon(w) \delta_{w \lambda} \tag{2.5}
\end{equation*}
$$

where $\delta_{w \lambda}$ is the $\delta$ function at $w \lambda$.

## 3. Spin quantization

Let $M$ be an even dimensional oriented spin manifold of dimension $2 d$ and provided with an action of $G$. If $(\mathcal{L}, \nabla)$ is a $G$ equivariant line bundle equipped with a $G$ invariant hermitian connection $\nabla$, we obtain a moment map $\phi_{G}: M \rightarrow \mathfrak{g}^{*}$ and a closed two form $\Omega$ on $M$ using Kostant formula:

$$
\begin{equation*}
\nabla^{2}=-i \Omega \quad \text { and } \quad L(X)-\nabla_{X}=i\left\langle\phi_{G}, X\right\rangle \tag{3.1}
\end{equation*}
$$

Here $X \in \mathfrak{g}$ and $L(X)$ is the infinitesimal action of $X \in \mathfrak{g}$ on smooth sections of $\mathcal{L}$.

Assume $M$ compact. The spin quantization $\mathcal{Q}_{G}^{\text {spin }}(M, \mathcal{L})$ is a virtual finite dimensional representation of $G$, constructed as the index of the Dirac operator on $M$ twisted by $\mathcal{L}$. If $k$ is an integer, we denote $\mathcal{L}^{k}$ the $k$-th power of the line bundle $\mathcal{L}$ : we can consider the equivariant index

$$
\mathcal{Q}_{G}^{\mathrm{spin}}\left(M, \mathcal{L}^{k}\right)=\sum_{\lambda \in A_{G}} m_{G}(\lambda, k) V_{\lambda}
$$

Taking traces, we also write, for $g \in G$,

$$
\mathcal{Q}_{G}^{\mathrm{spin}}\left(M, \mathcal{L}^{k}\right)(g)=\sum_{\lambda \in A_{G}} m_{G}(\lambda, k) \chi_{\lambda}(g)
$$

We will need the following basic fact.
Lemma 3.1. There exist a polynomial $P(k)$ such that $\left|\mathcal{Q}_{G}^{\text {spin }}\left(M, \mathcal{L}^{k}\right)(g)\right| \leq$ $P(k)$ for all $g \in G$ and $k \geq 1$.

We associate to $(M, \mathcal{L})$ the distributions on $\mathfrak{g}^{*}$ given by

$$
\begin{equation*}
\Psi_{k}^{(M, \mathcal{L})}:=\mathcal{F}_{\mathfrak{g}} \circ \operatorname{rescaling}_{\mathfrak{g}}(k)\left(\mathcal{Q}_{G}^{\text {spin }}\left(M, \mathcal{L}^{k}\right)\left(e^{X}\right)\right) \tag{3.2}
\end{equation*}
$$

and

$$
\begin{align*}
\Theta_{k}^{(M, \mathcal{L})} & :=\mathcal{F}_{\mathfrak{g}} \circ \operatorname{rescaling}_{\mathfrak{g}}(k)\left(\mathcal{Q}_{G}^{\text {spin }}\left(M, \mathcal{L}^{k}\right)\left(e^{X}\right) j_{\mathfrak{g}}^{1 / 2}(X)\right) \\
& =\operatorname{rescaling}_{\mathfrak{g}^{*}}(k)\left(\sum_{\lambda \in A_{G}} m_{G}(\lambda, k) \beta_{\lambda}\right) \\
& =k^{r} \sum_{\lambda \in A_{G}} m_{G}(\lambda, k) \beta_{\lambda / k}, \tag{3.3}
\end{align*}
$$

where $r=\operatorname{dim}(G / T) / 2$. Thus $\Theta_{k}^{(M, \mathcal{L})}$ is the rescaled geometric analogue of $\mathcal{Q}_{G}^{\text {spin }}\left(M, \mathcal{L}^{k}\right)=\sum_{\lambda \in A_{G}} m_{G}(\lambda, k) \chi_{\lambda}$.
3.1. Quasi-polynomial behavior of the multiplicities. We recall the behaviour of the multiplicity function $m_{G}: A_{G} \times \mathbb{Z}_{>0} \rightarrow \mathbb{N}$ defined by the relation $\mathcal{Q}_{G}^{\text {spin }}\left(M, \mathcal{L}^{k}\right)=\sum_{\lambda \in A_{G}} m_{G}(\lambda, k) \chi_{\lambda}$.

We suppose in this section that the $G$-action on $M$ admits abelian infinitesimal stabilizers. Recall that we have chosen a $G$-invariant Hermitian connection on $\mathcal{L}$, and $\phi_{G}: M \rightarrow \mathfrak{g}^{*}$ is the associated moment map.

We consider the positive Weyl chamber $\mathfrak{t}_{\geq 0}^{*}$ with interior $\mathfrak{t}_{>0}^{*}$ and the following subset

$$
Y=\phi_{G}^{-1}\left(\mathfrak{t}_{>0}^{*}\right) .
$$

Then $Y$ is a $T$-invariant sub-manifold of $M$, not necessarily connected, but every connected component of $Y$ is even dimensional. We denote by $\phi_{T}: Y \rightarrow \mathfrak{t}^{*}$ the restriction of the map $\phi_{G}$ to $Y$. Thus $\phi_{T}(Y)$ is contained in $\mathfrak{t}_{>0}^{*}$.

Lemma 3.2. The Spin structure on $M$ induces a Spinc bundle $\mathcal{S}_{Y}$ on $Y$ with determinant line bundle equal to $\mathbb{C}_{-2 \rho}$.

Now we explain how we can attach a multiplicity function

$$
m_{j}: A_{G} \times \mathbb{Z}_{>0} \rightarrow \mathbb{Z}
$$

to a connected component $Y_{j}$ of $Y$. Let $\mathfrak{t}_{j} \subset \mathfrak{t}$ be the generic infinitesimal stabilizer of the $T$-action on $Y_{j}$. We denote $I_{j} \subset \mathfrak{t}^{*}$ the affine subspace with direction $\mathfrak{t}_{j}^{\perp}$ that contains $\phi_{T}\left(Y_{j}\right)$.

We consider a finite collection $\mathcal{B}_{j}$ of affine co-dimension 1 subspaces of $I_{j}$ defined as follows: $E \in \mathcal{B}_{j}$ if $E=\phi_{T}(y)+\mathfrak{t}_{y}^{\perp}$ for some $y \in Y_{j}$ such that $\operatorname{dim}\left(\mathfrak{t}_{y} / \mathfrak{t}_{j}\right)=1$.

Definition 3.3. A chamber of $I_{j}$ is a connected component of the open subset $\left(I_{j}\right)_{\text {reg }}:=\left(I_{j} \backslash \cup_{E \in \mathcal{B}_{j}} E\right) \cap \mathfrak{t}_{>0}^{*}$.

Thus the closure of a chamber is a polytope $P$ contained in $\mathfrak{t}_{\geq 0}^{*}$.
Let $\mathfrak{c}$ be a chamber of $I_{j}$ and let $\xi \in \mathfrak{c}$ be a regular value of the map $\phi_{T}: Y_{j} \rightarrow I_{j}$. We consider the orbifold

$$
\mathcal{Y}_{j, \xi}:=\left(\phi_{T}^{-1}(\xi) \cap Y_{j}\right) / T .
$$

Lemma 3.4. Suppose that $\frac{\lambda}{k} \in I_{j}$. The spinc bundle $\left.\mathcal{S}_{Y_{j}} \otimes \mathcal{L}^{k}\right|_{Y_{j}} \otimes \mathbb{C}_{-\lambda+\rho}$ on $Y_{j}$ induces a spinc bundle $\mathcal{S}_{j, \xi}^{\lambda, k}$ on $\mathcal{Y}_{j, \xi}$. We denote $\mathcal{Q}\left(\mathcal{Y}_{j, \xi}, \mathcal{S}_{j, \xi}^{\lambda, k}\right)$ the index of the corresponding Dirac operator.

The Kawasaki index theorem [11] tells us that the map

$$
\mathrm{m}_{j, \xi}:(\lambda, k) \longmapsto \mathcal{Q}\left(\mathcal{Y}_{j, \xi}, \mathcal{S}_{j, \xi}^{\lambda, k}\right)
$$

is a quasi-polynomial on the set $\left\{(\lambda, k) \in A_{G} \times \mathbb{Z}_{>0}, \frac{\lambda}{k} \in I_{j}\right\}$.

Theorem 3.5. - The quantity $\mathcal{Q}\left(\mathcal{Y}_{j, \xi}, \mathcal{S}_{j, \xi}^{\lambda, k}\right)$ does not depend on the choice of $\xi \in \mathfrak{c}$. We denote by $\mathrm{m}_{j, \mathfrak{c}}$ the corresponding quasi-polynomial function.

- Let $\mathfrak{c}_{1}, \mathfrak{c}_{2}$ be two chambers such that $\frac{\lambda}{k} \in \overline{\mathfrak{c}_{1}} \cap \overline{\mathfrak{c}_{2}}$. Then

$$
\mathrm{m}_{j, \mathfrak{c}_{1}}(\lambda, k)=\mathrm{m}_{j, \mathfrak{c}_{2}}(\lambda, k)
$$

The first point of the previous theorem tell us that $\mathrm{m}_{j, \mathrm{c}}$ is the zero map if the chamber $\mathfrak{c}$ is not contained in the image of the map $\Phi_{T}$ : $Y_{j} \rightarrow I_{j}$. It is due to the fact that we can choose a regular value $\xi \in \mathfrak{c}$ with empty reduced space $\mathcal{Y}_{j, \xi}$.

Thus we can define multiplicity functions $\left(m_{j}\right)_{j \in J}$ parametrized by the set $J$ of connected components of the manifold $Y$ as follows.

Definition 3.6. The map $m_{j}: A_{G} \times \mathbb{Z}_{>0} \rightarrow \mathbb{Z}$ is defined as follows

$$
m_{j}(\lambda, k)=\left\{\begin{array}{lr}
0 & \text { if } \frac{\lambda}{k} \notin I_{j}, \\
m_{j, \mathfrak{c}}(\lambda, k) & \text { if } \frac{\lambda}{\bar{k}} \in \overline{\mathfrak{c}} \subset I_{j} .
\end{array}\right.
$$

Recall the following theorem [21].
Theorem 3.7. The following relation

$$
m_{G}(\lambda, k)=\sum_{j \in J} m_{j}(\lambda, k)
$$

holds for any $(\lambda, k) \in A_{G} \times \mathbb{Z}_{>0}$.
In particular, this implies that the function $k \mapsto m_{G}(k \lambda, k)$ is a quasi polynomial function of $k \geq 1$.

An important consequence of this geometric description is as follows. Let $\tilde{\Lambda}$ be the lattice generated by $\Lambda$ and $\rho$. To each closure $\overline{\mathfrak{c}}$ of a chambers in $I_{j}$, we have associated a quasi-polynomial function $m_{\bar{c}}(\lambda, k)$ on the lattice $\tilde{\Lambda} \oplus \mathbb{Z}$ such that $m_{G}(\lambda, k)=m_{\bar{c}}(\lambda, k)$ if $\lambda \in k \overline{\mathfrak{c}} \cap A_{G}$.

Let $\mathcal{A}$ be the collection of closed polytopes formed by the faces of the closures $\overline{\mathfrak{c}}$ of the chambers in $I_{j}$. To each $P \in \mathcal{A}$, we associate

- the quasi polynomial $m_{P}$ on $\tilde{\Lambda} \oplus \mathbb{Z}$ which is equal to $m_{\overline{\mathfrak{c}}}$ for a polytope $\overline{\mathfrak{c}}$ containing $P$ as face,
- the following cone in $\mathfrak{t}^{*} \times \mathbb{R}$

$$
C_{P}:=\{(t \xi, t), t \geq 0, \xi \in P\}
$$

and we denote $\left[C_{P}\right]$ the characteristic function of $C_{P}$.
By inclusion-exclusion, we see that the multiplicity function $m_{G}$ admits the decomposition

$$
\begin{equation*}
m_{G}=\sum_{P \in \mathcal{A}} \alpha_{P} m_{P}\left[C_{P}\right] \tag{3.4}
\end{equation*}
$$

for some appropriate choice of constants $\alpha_{P} \in \mathbb{Z}$.
Example 3.8. Let us give a simple example for the group $G=S U(2)$. Consider $\mathfrak{t}^{*}$ with basis $\rho$. In this basis $A_{G}=\{\lambda ; \lambda \in \mathbb{Z}, \lambda>0\}$. We consider the line bundle $\mathcal{L}$ associated to $(\rho, \rho)$ on the spin manifold $M=G / T \times G / T$. We see that $\mathcal{Q}_{G}^{\text {spin }}(M, \mathcal{L})$ is the tensor product $V(k \rho) \otimes V(k \rho)$, and $V(k \rho)$ is the irreducible representation of dimension $k$ of $G$. So $\mathcal{Q}_{G}^{\text {spin }}(M, \mathcal{L})$ is the sum of the representations $V(j \rho)$, with $j$ odd and less than $2 k$. If we write $\mathcal{Q}_{G}^{\text {spin }}\left(M, \mathcal{L}^{k}\right)=\sum_{\lambda>0} m_{G}(\lambda, k) V(\lambda \rho)$,

$$
m_{G}(\lambda, k)=\frac{1}{2}\left(1-(-1)^{\lambda}\right), \text { for } 1 \leq \lambda \leq 2 k
$$

3.2. Quasi-polynomials and asymptotics. We consider the families of distributions on $\mathfrak{g}^{*}$ associated to piecewise quasi-polynomial functions.
Definition 3.9. A function $m: \tilde{\Lambda} \times \mathbb{Z} \rightarrow \mathbb{C}$ is called a piecewise quasi-polynomial functions if there is a collection $\mathcal{A}$ of closed rational polytopes in $\mathfrak{t}^{*}$ and a collection of quasi-polynomial functions $\left(m_{P}\right)_{P \in \mathcal{A}}$ such that

$$
\begin{equation*}
m=\sum_{P \in \mathcal{A}} m_{P}\left[C_{P}\right] \tag{3.5}
\end{equation*}
$$

Here $\mathcal{A}$ is not necessarily finite but, in order that (3.5) makes sense, we assume that for any compact $\mathcal{K} \subset \mathfrak{t}^{*}$ the set $\{P \in \mathcal{A}, P \cap \mathcal{K} \neq \emptyset\}$ is finite.

To a piecewise quasi-polynomial function $m: \tilde{\Lambda} \times \mathbb{Z} \rightarrow \mathbb{C}$ we associate the family of invariant distributions on $\mathfrak{g}^{*}$

$$
\Theta_{k}(m)=k^{r} \sum_{\lambda \in A_{G}} m(\lambda, k) \beta_{\lambda / k}
$$

In a companion article [23], we proved the following easy result.
Proposition 3.1. The family $\Theta_{k}(m)$ admits an asymptotic expansion

$$
\Theta_{k}(m) \equiv \mathrm{AS}_{k}(m):=\sum_{n \in \mathbb{Z}} \frac{1}{k^{n}} \theta_{n}(k)
$$

Here $\left(\theta_{n}(k)\right)_{n \in \mathbb{Z}}$ is a collection of invariant distributions on $\mathfrak{g}^{*}$, depending periodically of $k$, and such that $\theta_{n}(k)=0$ if $n<n_{0}^{m}$.

For our computation we need also a variation of Proposition 3.1. Let $\mathrm{j}: \mathfrak{g} \rightarrow \mathbb{C}$ be an invariant analytic function such that its Fourier transform $B_{\mathrm{j}}:=\mathcal{F}_{\mathfrak{g}}(\mathrm{j})$ is a compactly supported measure. We denote $B_{\mathrm{j}}^{k}=\operatorname{rescaling}_{\mathfrak{g}^{*}}(k) B_{\mathrm{j}}$.

We consider now the family of distributions $B_{\mathrm{j}}^{k} \star \Theta_{k}(m)$ where $\star$ denotes the convolution.

Proposition 3.2. The family $B_{\mathrm{j}}^{k} \star \Theta_{k}(m)$ admits the asymptotic expansion

$$
B_{\mathrm{j}}^{k} \star \Theta_{k}(m) \equiv \mathrm{j}(i \partial / k) \mathrm{AS}_{k}(m)
$$

Proof. Let $\varphi$ be a test function on $\mathfrak{g}^{*}$. Let $R>0$ such that the ball $\left\{\xi \in \mathfrak{g}^{*},\|\xi\| \leq R / 2\right\}$ contains the compact supports of $\varphi$ and $B_{\mathrm{j}}$.

We start with the relation

$$
\left\langle B_{\mathrm{j}}^{k} \star \Theta_{k}(m), \varphi\right\rangle=\left\langle\Theta_{k}(m), \varphi_{k}\right\rangle
$$

where

$$
\varphi_{k}(\xi)=\int_{\mathfrak{g}^{*}} B_{\mathrm{j}}\left(\xi^{\prime}\right) \varphi\left(\xi+\xi^{\prime} / k\right), \quad k \geq 1
$$

are smooth functions supported on the ball $\left\{\xi \in \mathfrak{g}^{*},\|\xi\| \leq R\right\}$. Let us write the Taylor series of $\varphi$ at $\xi: \varphi(\xi+\eta)=\varphi(\xi)+\sum_{n=1}^{N} d_{n} \varphi(\xi)(\eta)+$ $r_{N}(\xi, \eta)$. Here $\eta \mapsto d_{n} \varphi(\xi)(\eta)$ is an homogeneous polynomial of degree $n$ depending smoothly of the variable $\xi$, and $r_{N}$ is a smooth function of $(\xi, \eta)$ such that

$$
\begin{equation*}
\left|r_{N}(\xi, \eta)\right| \leq c_{N}\|\eta\|^{N+1}, \quad \forall(\xi, \eta) \in \mathfrak{g}^{*} \times \mathfrak{g}^{*} \tag{3.6}
\end{equation*}
$$

for some $c_{N}>0$.
If we write the Taylor series of j at $0, \mathrm{j}(X)=\sum_{n=0}^{\infty} \mathrm{j}_{n}(X)$, we see that

$$
\int_{\mathfrak{g}^{*}} B_{\mathrm{j}}(\eta) d_{n} \varphi(\xi)(\eta)=\mathrm{j}_{n}(-i \partial) \varphi(\xi)
$$

So the functions $\varphi_{k}$ admit the following description

$$
\varphi_{k}=\sum_{n=0}^{N} \frac{1}{k^{n}} \mathrm{j}_{n}(-i \partial) \varphi+R_{k}^{N}
$$

where $R_{k}^{N}(\xi)=\int_{\mathfrak{g}^{*}} B_{\mathrm{j}}\left(\xi^{\prime}\right) r_{N}\left(\xi, \xi^{\prime} / k\right)$. Thanks to (3.6), we see that there exists $C_{N}$ such that $\left|R_{k}^{N}(\xi)\right| \leq C_{N} / k^{N+1}, \forall \xi \in \mathfrak{g}^{*}$. We check also that for any $k \geq 1$ the support of $R_{N}^{k}$ is contained in the ball $\left\{\xi \in \mathfrak{g}^{*},\|\xi\| \leq R\right\}$.

Finally we obtain

$$
\begin{equation*}
\left\langle B_{\mathrm{j}}^{k} \star \Theta_{k}(m), \varphi\right\rangle=\sum_{n=0}^{N} \frac{1}{k^{n}}\left\langle\Theta_{k}(m), \mathrm{j}_{n}(-i \partial) \varphi\right\rangle+\operatorname{Rest}_{k}^{N} \tag{3.7}
\end{equation*}
$$

where

$$
\operatorname{Rest}_{k}^{N}=\left\langle\Theta_{k}(m), R_{k}^{N}\right\rangle=k^{r} \sum_{\|\lambda\| / k \leq R} m(\lambda, k)\left\langle\beta_{\lambda / k}, R_{k}^{N}\right\rangle
$$

The term $\mid$ Rest $_{k}^{N} \mid$ is bounded by $\frac{C_{N}}{k^{N+1}} \sum_{\|\lambda\| / k \leq R}|m(\lambda, k)| \operatorname{vol}(G \lambda)$. Since $m(\lambda, k)$ is a piecewise quasi-polynomial function we see that for any $p \geq 1$ there exist $N \geq 1$ such that

$$
\begin{equation*}
\operatorname{Rest}_{k}^{N}=O\left(k^{-p}\right) \tag{3.8}
\end{equation*}
$$

Identities (3.7) and (3.8) show that the family $B_{\mathrm{j}}^{k} \star \Theta_{k}(m)$ admits the asymptotic expansion $\mathrm{j}(i \partial / k) \mathrm{AS}_{k}(m)$.
3.3. Equivariant cohomology. Let $N$ be a $G$-manifold and let $\mathcal{A}(N)$ be the space of differential forms on $N$, graded by its exterior degree. Following [3] and [26], an equivariant form is a $G$-invariant smooth function $\alpha: \mathfrak{g} \rightarrow \mathcal{A}(N)$, thus $\alpha(X)$ is a differential form on $N$ depending smoothly of $X \in \mathfrak{g}$. Consider the operator

$$
\begin{equation*}
d_{\mathfrak{g}} \alpha(X)=d \alpha(X)-\iota\left(v_{X}\right) \alpha(X) \tag{3.9}
\end{equation*}
$$

where $\iota\left(v_{X}\right)$ is the contraction by the vector field $v_{X}$ generated by the action of $-X$ on $N$. Then $d_{\mathfrak{g}}$ is an odd operator with square 0 , and the equivariant cohomology is defined to be the cohomology space of $d_{\mathfrak{g}}$. It is important to note that the dependence of $\alpha$ on $X$ may be $C^{\infty}$. If the dependence of $\alpha$ in $X$ is polynomial, we denote by $H_{G}^{*}(N)$ the corresponding $\mathbb{Z}$-graded algebra. By definition, the grading of $P(X) \otimes \mu, P$ a homogeneous polynomial and $\mu$ a differential form on $N$, is the exterior degree of $\mu$ plus twice the polynomial degree in $X$.

The line bundle $\mathcal{L} \rightarrow M$ determines the closed equivariant form $\Omega(X)=\Omega-\left\langle\phi_{G}, X\right\rangle$ on $M$. Here we will not assume any non degeneracy condition on $\Omega$.

Choose a $G$-invariant Riemannian metric on $M$. Let

$$
J(A)=\operatorname{det}_{\mathbb{R}^{2 d}}\left(\frac{e^{A / 2}-e^{-A / 2}}{A}\right)
$$

an invariant function of $A$. Then $J(0)=1$. Consider $\frac{1}{J^{1 / 2}(A)}$ and its Taylor expansion at 0 :

$$
\frac{1}{J^{1 / 2}(A)}=\operatorname{det}_{\mathbb{R}^{2 d}}^{1 / 2}\left(\frac{A}{e^{A / 2}-e^{-A / 2}}\right)=\sum_{n=0}^{\infty} B_{n}(A)
$$

Each function $B_{n}(A)$ is an invariant polynomial of degree $n$ (in fact $B_{n}$ is non zero only for $n$ even) and by the Chern Weil construction, and choice of connections on $T M$, the function $B_{n}$ determines a closed equivariant form $\hat{A}_{n}(M)(X)$ on $M$ of homogeneous equivariant degree
$2 n$. Remark that $\hat{A}_{0}(M)=1$. We define the formal series of equivariant forms:

$$
\hat{A}(M)(X)=\sum_{n=0}^{\infty} \hat{A}_{n}(M)(X) .
$$

As $M$ is compact, we can find a positive constant $r_{M}$ such that, for $\|X\|<r_{M}$, these series of equivariant forms is convergent. In particular $\hat{A}(M)(0)$ is a closed differential form on $M$ which represents the usual $\hat{A}$ class of $M$.
3.4. The equivariant index. Recall the "delocalized index formula" (see [2]). For $X \in \mathfrak{g}$ such that $\|X\|<r_{M}$, so that $\hat{A}(M)(X)$ is well defined, we have

$$
\begin{equation*}
\mathcal{Q}_{G}^{\text {spin }}(M, \mathcal{L})(\exp X)=\frac{1}{(-2 i \pi)^{d}} \int_{M} e^{-i \Omega(X)} \hat{A}(M)(X) \tag{3.10}
\end{equation*}
$$

Here $d=\operatorname{dim} M / 2$. Note that Kostant relations (3.1) implies that the equivariant form $\Omega(X)=\Omega-\left\langle\phi_{G}, X\right\rangle$ is closed. In other words $\iota(X) \Omega+d\left\langle\phi_{G}, X\right\rangle=0$ for any $X \in \mathfrak{g}$.

For each integer $n$, consider the analytic function on $\mathfrak{g}$ given by

$$
\mathcal{I}_{n}(X)=\frac{1}{(-2 i \pi)^{d}} \int_{M} e^{-i \Omega(X)} \hat{A}_{n}(M)(X) .
$$

There is a remarkable relation between the character associated to $\mathcal{L}^{k}$ and the dilation $X \rightarrow X / k$ on $\mathfrak{g}$.

Lemma 3.10. When $X \in \mathfrak{g}$ is such that $\|X\|<r_{M}$, then for any $k \geq 1$ one has

$$
\mathcal{Q}_{G}^{\text {spin }}\left(M, \mathcal{L}^{k}\right)(\exp (X / k))=k^{d} \sum_{n=0}^{\infty} \frac{1}{k^{n}} \mathcal{I}_{n}(X) .
$$

Proof. Write $\hat{A}(M)(X)=\sum_{n=0}^{\infty} \tau_{n}(X)$ as a sum of forms with coefficients homogeneous polynomials in $X$ of degree $n$. Thus

$$
\hat{A}_{n}(M)(X)=\sum_{q \leq n}\left(\tau_{q}(X)\right)_{[n-2 q]}
$$

where $\omega_{[s]}$ is the component of exterior degree $s$ of a differential form $\omega$.

For $\|X\|<r_{M}, \sum_{n=0}^{\infty} \tau_{n}(X)$ is a convergent series with sum the equivariant $\hat{A}$ class. We obtain

$$
\mathcal{Q}_{G}^{\mathrm{spin}}\left(M, \mathcal{L}^{k}\right)(\exp (X / k))=\frac{1}{(-2 i \pi)^{d}} \int_{M} e^{-i k \Omega(X / k)}\left(\sum_{n=0}^{\infty} \tau_{n}(X / k)\right)
$$

Let $b(X, k)=e^{-i k \Omega(X / k)}=e^{-i k \Omega} e^{i\left\langle\phi_{G}, X\right\rangle}$. Remark that $b(X, k)$ depends polynomially of $k$. We rewrite

$$
\begin{equation*}
\mathcal{Q}_{G}^{\mathrm{spin}}\left(M, \mathcal{L}^{k}\right)(\exp (X / k))=\frac{1}{(-2 i \pi)^{d}} \int_{M} b(X, k)\left(\sum_{n=0}^{\infty} \frac{1}{k^{n}} \tau_{n}(X)\right) \tag{3.11}
\end{equation*}
$$

Only the top exterior degree term contribute to the integral. We compute it in terms of the $\hat{A}_{n}(M)$.

## Lemma 3.11.

$$
\left[b(X, k)\left(\sum_{n=0}^{\infty} \frac{1}{k^{n}} \tau_{n}(X)\right)\right]_{[2 d]}=k^{d}\left[e^{-i \Omega(X)}\left(\sum_{n=0}^{\infty} \frac{1}{k^{n}} \hat{A}_{n}(M)(X)\right)\right]_{[2 d]}
$$

Proof. The coefficient of $k^{d} \frac{1}{k^{n}}$ in the left hand side is

$$
\sum_{a=0}^{d} \frac{(-i \Omega)^{a}}{a!}\left(\tau_{n+a-d}(X)\right)_{[2 d-2 a]} e^{i\left\langle\phi_{G}, X\right\rangle}
$$

This is the term of exterior degree $2 d$ of $e^{-i \Omega(X)} \hat{A}_{n}(M)(X)$.
This identity implies the lemma.
3.5. Twisted Duistermaat-Heckman distributions. Let $\nu$ be a closed equivariant form on $M$ with polynomial coefficients. Let $\Omega(X)=$ $\Omega-\left\langle\phi_{G}, X\right\rangle$ our equivariant curvature. As $\nu(X)$ depends polynomially on $X, \int_{M} e^{-i \Omega(X)} \nu(X)$ is a function on $\mathfrak{g}$ of at most polynomial growth.
Definition 3.12. We denote $D H^{G}(M, \Omega, \nu) \in \mathcal{D}^{\prime}\left(\mathfrak{g}^{*}\right)^{G}$ the Fourier transform of $\frac{1}{(-2 i \pi)^{d}} \int_{M} e^{-i \Omega(X)} \nu(X)$.

We thus have for $\varphi$ smooth with compact support,

$$
\begin{aligned}
\left\langle D H^{G}(M, \Omega, \nu), \varphi\right\rangle & =\frac{1}{(-2 i \pi)^{d}} \int_{M} \int_{\mathfrak{g}} e^{-i \Omega(X)} \nu(X) \widehat{\varphi}(X) d X \\
& =\frac{1}{(-2 i \pi)^{d}} \int_{M} e^{-i \Omega}[\nu(-i \partial) \varphi]\left(\phi_{G}(m)\right) .
\end{aligned}
$$

If $\nu(X)=\sum_{a} p_{a}(X) \nu_{a}$ with polynomial functions $p_{a}$ on $\mathfrak{g}$ and differential forms $\nu_{a}$ on $M$, the expression

$$
[\nu(-i \partial) \varphi]\left(\phi_{G}(m)\right)=\sum_{a}\left[p_{a}(-i \partial)(\varphi)\right]\left(\phi_{G}(m)\right) \nu_{a}
$$

is a differential form on $M$ supported on the pull-back by $\phi_{G}$ of the support of $\varphi$. This shows, and this will be used in our application to a line bundle with proper moment map, that $D H^{G}(M, \Omega, \nu)$ is a well defined distribution supported on $\phi_{G}(M)$, provided $\phi_{G}$ is proper.

Notice that the distribution $D H^{G}(M, \Omega, \nu)$ depends only of the class $[\nu] \in H_{G}^{*}(M)$.

Recall that we have associated to $(M, \mathcal{L})$ the distributions on $\mathfrak{g}^{*}$ given by

$$
\begin{aligned}
\Psi_{k}^{(M, \mathcal{L})} & :=\mathcal{F}_{\mathfrak{g}} \circ \operatorname{rescaling}_{\mathfrak{g}}(k)\left(\mathcal{Q}_{G}^{\text {spin }}\left(M, \mathcal{L}^{k}\right)\left(e^{X}\right)\right) \\
\Theta_{k}^{(M, \mathcal{L})} & :=\mathcal{F}_{\mathfrak{g}} \circ \operatorname{rescaling}_{\mathfrak{g}}(k)\left(\mathcal{Q}_{G}^{\text {spin }}\left(M, \mathcal{L}^{k}\right)\left(e^{X}\right) j_{\mathfrak{g}}^{1 / 2}(X)\right)
\end{aligned}
$$

The distribution $\Theta_{k}^{(M, \mathcal{L})}$ is the weighted sum $k^{r} \sum_{\lambda \in A_{G}} m_{G}(\lambda, k) \beta_{\lambda / k}$ of the canonical measure on the orbits $G \lambda / k$.

Theorem 3.13. Let $M$ be a compact spin even dimensional oriented manifold with a $G$ equivariant line bundle $\mathcal{L}$. When $k$ tends to $\infty$, we have the asymptotic expansions

$$
\begin{equation*}
\Psi_{k}^{(M, \mathcal{L})} \equiv k^{d} \sum_{n=0}^{\infty} \frac{1}{k^{n}} D H^{G}\left(M, \Omega, \hat{A}_{n}(M)\right) \tag{3.12}
\end{equation*}
$$

and

$$
\begin{equation*}
\Theta_{k}^{(M, \mathcal{L})} \equiv j_{\mathfrak{g}}^{1 / 2}(i \partial / k)\left(k^{d} \sum_{n=0}^{\infty} \frac{1}{k^{n}} D H^{G}\left(M, \Omega, \hat{A}_{n}(M)\right)\right) \tag{3.13}
\end{equation*}
$$

Proof. We present a detailed proof since we will have to adapt this proof to the case of a non compact manifold $M$.

We fix a $G$-invariant function $h: \mathfrak{g} \rightarrow \mathbb{R}$ equal to 1 on a small ball $\|X\| \leq r$ and with compact support contained in $\|X\|<r_{M}$, so that $h(X) \hat{A}(M)(X)$ is a smooth compactly supported function on $\mathfrak{g}$ with value differential forms on $M$.

Let $\varphi$ be a smooth compactly supported function on $\mathfrak{g}^{*}$. We have $\left\langle\Psi_{k}^{(M, \mathcal{L})}, \varphi\right\rangle=\mathrm{I}_{k}+\mathrm{J}_{k}$ where

$$
\mathrm{I}_{k}=\int_{\mathfrak{g}} \mathcal{Q}_{G}^{\mathrm{spin}}\left(M, \mathcal{L}^{k}\right)\left(e^{X / k}\right) h(X / k) \widehat{\varphi}(X) d X
$$

and

$$
\mathrm{J}_{k}=\int_{\mathfrak{g}} \mathcal{Q}_{G}^{\mathrm{spin}}\left(M, \mathcal{L}^{k}\right)\left(e^{X / k}\right)(1-h(X / k)) \widehat{\varphi}(X) d X
$$

We will see that $\mathrm{J}_{k}=O\left(k^{-\infty}\right)$. For estimating $\mathrm{I}_{k}$, we will use the delocalized index formula (3.10).

Using the fact that $\left|\mathcal{Q}_{G}^{\text {spin }}\left(M, \mathcal{L}^{k}\right)(g)\right|$ is uniformly bounded by a polynomial in $k$, that the support of $(1-h(X / k))$ is contained in the set $\|X\| \geq r k, r>0$, and that the function $\widehat{\varphi}$ is rapidly decreasing, we see that $\mathrm{J}_{k}=O\left(k^{-\infty}\right)$.

We now analyse $\mathrm{I}_{k}$. We can use Formula (3.10) since $\|X / k\| \leq r_{M}$. Thus

$$
\mathrm{I}_{k}=\frac{1}{(-2 i \pi)^{d}} \int_{M} \int_{\mathfrak{g}} e^{-i k \Omega(X / k)} \hat{A}(M)(X / k) h(X / k) \widehat{\varphi}(X) d X .
$$

Proposition 2.5 extends to the case of vector valued functions on $\mathfrak{g}$. Let $E$ be a finite dimensional space, and let $\tau: \mathfrak{g} \rightarrow \Lambda E$ be a smooth function with at most polynomial growth, as well as all its derivatives. The coefficients $\tau_{n}(X)$ of its Taylor series $\sum_{n=0}^{\infty} \tau_{n}(X)$ are $\Lambda E$ valued polynomials functions on $\mathfrak{g}$ homogeneous of degree $n$. Let $b(X, k)=\sum_{m=0}^{d} k^{m} b_{m}(X)$ where $b_{m}(X)$ are smooth functions of $X$ with value in $\Lambda E$ with at most polynomial growth. Thus

$$
b(X, k) \sum_{n=0}^{\infty} \tau_{n}(X / k)=\left(\sum_{m=0}^{d} k^{m} b_{m}(X)\right)\left(\sum_{n=0}^{\infty} \frac{1}{k^{n}} \tau_{n}(X)\right)
$$

is a Laurent series in $\frac{1}{k}$ of functions on $\mathfrak{g}$ with values in $\Lambda E$.
We have

$$
\begin{equation*}
\int_{\mathfrak{g}} b(X, k) \tau(X / k) \widehat{\varphi}(X) d X \equiv \int_{\mathfrak{g}} b(X, k)\left(\sum_{n=0}^{\infty} \tau_{n}(X / k)\right) \widehat{\varphi}(X) d X \tag{3.14}
\end{equation*}
$$

In short, we replace $\tau$ by its Taylor series.
Consider the differential form $h(X) \hat{A}(M)(X)$. For each $x \in M$, $\tau_{x}^{h}(X)=\left.h(X) \hat{A}(M)\right|_{x}(X)$ is a smooth compactly supported function on $\mathfrak{g}$ with values in $\Lambda T_{x}^{*} M$. Let

$$
b_{x}(X, k)=e^{i\left\langle\phi_{G}(x), X\right\rangle} e^{-i k \Omega_{x}} .
$$

It depends polynomially of $k$ and each coefficient of this polynomial is a bounded function of $X$ (as it is proportional to $e^{i\left\langle\phi_{G}(x), X\right\rangle}$ ). So we can certainly apply the asymptotic formula (3.14).

Let $\mathrm{I}_{k}^{x}=\frac{1}{(-2 i \pi)^{d}} \int_{\mathfrak{g}} b_{x}(X, k) \tau_{x}^{h}(X / k) \widehat{\varphi}(X) d X$, with values in $\Lambda T_{x}^{*} M$. So $x \mapsto \mathrm{I}_{k}^{x}$ is a differential form and $\mathrm{I}_{k}=\int_{x \in M} \mathrm{I}_{k}^{x}$.

The manifold $M$ is compact. Proposition 2.5 gives us an asymptotic expansion for $\mathrm{I}_{k}^{x}$, and the proof shows that the rest is bounded uniformly. So the asymptotic expansion of $\mathrm{I}_{k}$ is obtained by replacing (at each $x \in M) \tau_{x}^{h}(X / k)$ by its Taylor series. As $h$ is identically 1 near zero, we obtain in the notations of the proof of Lemma 3.10,

$$
\mathrm{I}_{k} \equiv \frac{1}{(-2 i \pi)^{d}} \int_{M} \int_{\mathfrak{g}} b(X, k)\left(\sum_{n=0}^{\infty} \frac{1}{k^{n}} \tau_{n}(X)\right) \widehat{\varphi}(X) d X
$$

By Lemma 3.11, this is

$$
k^{d} \frac{1}{(-2 i \pi)^{d}} \int_{M} \int_{\mathfrak{g}} e^{-i \Omega(X)}\left(\sum_{n=0}^{\infty} \frac{1}{k^{n}} \hat{A}_{n}(M)(X)\right) \widehat{\varphi}(X) d X
$$

By definition of the twisted Duistermaat-Heckman measure, we obtain the expansion of Formula (3.12).

The proof of Expansion formula (3.13) is identical.
3.6. Asymptotics localized at $g \in G$. For applications to formal geometric quantization, we state an analogous asymptotic descent formula. Now let $g \in G$. Let $G(g)$ be the centralizer of $g$ in $G$, and $\mathfrak{g}(g)$ its Lie algebra.

Consider the manifold $M^{g}$ of fixed points of the action of $g$ on $M$ : it may have several connected components $M_{a}^{g}$, all of even dimension. Since $M$ is spin, $M^{g}$ admits an orientation (that is all connected components are oriented). Let $\mathcal{N}^{g}$ be the normal bundle of $M^{g}$ in $M$ and let

$$
R_{\mathcal{N}^{g}}(X)=R_{\mathcal{N}^{g}}+\mu_{\mathcal{N}^{g}}(X)
$$

be its equivariant curvature [2].
We denote by $\Omega_{g}(X)$ the restriction of $\Omega(X)$ on $M^{g}$. We consider the following equivariant form

$$
D_{g}(M)(X)=\operatorname{det}_{\mathcal{N}^{g}}^{1 / 2}\left(I-\left(g^{-1}\right)^{\mathcal{N}^{g}} e^{-R_{\mathcal{N}^{g}}(X)}\right)
$$

Here the square root is chosen such that $\operatorname{det}_{\mathcal{N}^{g}}^{1 / 2}\left(I-\left(g^{-1}\right)^{\mathcal{N}^{g}}\right)>0$.
For $X \in \mathfrak{g}(g)$ small enough, we have

$$
\begin{equation*}
\mathcal{Q}_{G}^{\text {spin }}(M, \mathcal{L})(g \exp (X))=\int_{M^{g}} c_{g} \frac{\widehat{A}\left(M^{g}\right)(X)}{D_{g}(M)(X)} g^{\mathcal{L}} e^{-i \Omega_{g}(X)} \tag{3.15}
\end{equation*}
$$

where

- $g^{\mathcal{L}}$ is the locally constant function on $M^{g}$ with value a complex number of modulus 1 given by the action $g$ on the fiber of $\left.\mathcal{L}\right|_{M^{g}}$. We write $u_{a}$ for the restriction of $g^{\mathcal{L}}$ to a connected component $M_{a}^{g}$. This is a complex number of modulus 1 . If $g$ is of finite order, $u_{a}$ is a root of unity.
- $c_{g}=\epsilon_{g}(-2 i \pi)^{-\operatorname{dim} M / 2}(2 \pi)^{\mathrm{rk} \mathcal{N}^{g} / 2}$ is a locally constant function independent of $\mathcal{L}$. Here the sign $\epsilon_{g}= \pm$ depends only of the action of $g$ on the spinor bundle $\left.\mathcal{S}\right|_{M^{g}}$.

Consider the equivariant form $V_{g}(\mathcal{L}, k)(X)=c_{g}\left(g^{\mathcal{L}}\right)^{k} \frac{\widehat{A}\left(M^{g}\right)(X)}{D_{g}(M)(X)}$ on $M^{g}$. Then we have $V_{g}(\mathcal{L}, k)(X)=\sum_{n=0}^{\infty} V_{g}(n, k)(X)$ where $V_{g}(n, k)(X)$ are $G(g)$-equivariant forms on $M^{g}$ homogeneous of even degree $2 n$.

Thus, for $X \in \mathfrak{g}(g)$ small, we have

$$
\mathcal{Q}_{G}^{\mathrm{spin}}\left(M, \mathcal{L}^{k}\right)(g \exp (X))=\sum_{n=0}^{\infty} \int_{M^{g}} e^{-i k \Omega_{g}(X)} V_{g}(n, k)(X)
$$

Here Lemma 3.10 becomes
Lemma 3.14. When $X \in \mathfrak{g}(g)$ is sufficiently small, then for any $k \geq 1$, one has ${ }^{1}$

$$
\mathcal{Q}_{G}^{\mathrm{spin}}\left(M, \mathcal{L}^{k}\right)(g \exp (X / k))=\sum_{n=0}^{\infty} \frac{1}{k^{n}} k^{\frac{\operatorname{dim} M^{g}}{2}} \mathcal{I}_{g}(n, k)(X)
$$

where $\mathcal{I}_{g}(n, k)(X):=\int_{M^{g}} e^{-i \Omega_{g}(X)} V_{g}(n, k)(X)$.
Now consider the character $\chi_{\lambda}$ of the irreducible representation attached to the admissible orbit $\mathcal{O}_{\lambda}$. If $g \in G$, the fixed point set $\mathcal{O}_{\lambda}^{g}$ is a union of coadjoint orbits of the connected component $G(g)^{\circ}$ of $G(g)$. They are all of the same even dimension $r_{g}=\operatorname{dim} G(g)-\operatorname{dim} T$.

Then we have the formula, for $X \in \mathfrak{g}(g)$,

$$
\begin{equation*}
\chi_{\lambda}\left(g e^{X}\right) j_{\mathfrak{g}(g)}^{1 / 2}(X) \operatorname{det}_{\mathfrak{g} / \mathfrak{g}(g)}^{1 / 2}\left(1-g^{-1} e^{-X}\right)=\left\langle\beta(g, \lambda)(\xi), e^{i\langle\xi, X\rangle}\right\rangle \tag{3.16}
\end{equation*}
$$

Here the square root is determined such that $\operatorname{det}_{\mathfrak{g} / \mathfrak{g}(g)}^{1 / 2}\left(1-g^{-1}\right)>0$ and $\beta(g, \lambda)$ is a measure on $\mathcal{O}_{\lambda}^{g}$ entirely determined by this equation.

Let $\tilde{T} \rightarrow T$ the two-fold cover such that any $\lambda \in A_{G}$ defines a character $\tilde{t} \in \tilde{T} \rightarrow \tilde{t}^{\lambda}$. Suppose that $g \in T$ and denote $W_{G(g)} \subset W_{G}$ the Weyl group of $G(g)^{\circ}$. Then

$$
\mathcal{O}_{\lambda}^{g}=\bigcup_{\bar{w} \in W_{G(g)} \backslash W_{G}} G(g)^{o} w \lambda .
$$

Let $\mathcal{R}_{\mathfrak{g}(g)}: \mathcal{D}^{\prime}(\mathfrak{g}(g))^{G(g)^{o}} \rightarrow \mathcal{D}^{\prime}(\mathfrak{t})^{W_{G(g)}-\text { alt }}$ the canonical isomorphism.
Lemma 3.15. Choose $\tilde{g} \in \tilde{T}$ above $g \in T$.

- The measure $\beta(g, \lambda)$ is defined by the relation ${ }^{2}$

$$
\beta(g, \lambda)=\gamma_{\tilde{g}} \sum_{\bar{w} \in W_{G(g)} \backslash W_{G}} \tilde{g}^{w \lambda} \beta_{G(g)^{o} w \lambda}
$$

where $\gamma_{\tilde{g}}$ is a complex number of modulus 1 .

- We have $\mathcal{R}_{\mathfrak{g}(g)}(\beta(g, \lambda))=\gamma_{\tilde{g}} \sum_{w \in W_{G}} \epsilon(w) \tilde{g}^{w \lambda} \delta_{w \lambda}$.

[^0]Proof. The first point follows from the character formula and the second is a direct consequence of the first.

Definition 3.16. We denote $D H^{G(g)}\left(M^{g}, \Omega_{g}, V_{g}(n, k)\right) \in \mathcal{D}^{\prime}\left(\mathfrak{g}(g)^{*}\right)^{G(g)}$ the Fourier transform of the function $\mathcal{I}_{g}(n, k)$ (see Lemma 3.14).

Then for a test function $\varphi,\left\langle D H^{G(g)}\left(M^{g}, \Omega_{g}, V_{g}(n, k)\right), \varphi\right\rangle$ is of the form $\sum_{a} u_{a}^{k}\left\langle s_{a, n}, \varphi\right\rangle$ with $\left\langle s_{a, n}, \varphi\right\rangle=\int_{M_{a}^{g}} \int_{\mathfrak{g}(g)} e^{-i \Omega_{g}(Y)} \alpha_{a, n}(Y) \widehat{\varphi}(Y) d Y$ : here $\alpha_{a, n}(Y)$ is an equivariant form on $M_{a}^{g}$ of degree $2 n$, and $u_{a}$ is the restriction of $g^{\mathcal{L}}$ to $M_{a}^{g}$.

Asymptotic expansion of distributions of the form $\sum_{a} \sum_{n=0}^{\infty} u_{a}^{k} k^{-n} h_{n, a}$ are thus well defined and the coefficients $H(n, k)=\sum_{a} u_{a}^{k} h_{n, a}$ are uniquely determined.

With the same proof than Theorem 3.13 we obtain
Theorem 3.17. Let $M$ be a compact spin even dimensional oriented manifold with a $G$ equivariant line bundle $\mathcal{L}$. Let $g \in G$ of finite order. When $k$ tends to $\infty$, we have the asymptotic equivalence ${ }^{3}$

$$
\begin{gathered}
\operatorname{rescaling}_{\mathfrak{g}(g)^{*}}(k)\left(\sum_{\lambda \in \Lambda} m(\lambda, k) \beta(g, \lambda)\right) \equiv \\
j_{\mathfrak{g}(g)}^{\frac{1}{2}}(i \partial / k) \operatorname{det}_{\mathfrak{g} / \mathfrak{g}(g)}^{\frac{1}{2}}\left(1-g^{-1} e^{-i \partial / k}\right) \sum_{n=0}^{\infty} \frac{1}{k^{n}} k^{\frac{\mathrm{dim} M^{g}}{2}} D H^{G(g)}\left(M^{g}, \Omega_{g}, V_{g}(n, k)\right)
\end{gathered}
$$

## 4. Spin Quantization in the non-compact setting

In this section we work with a even dimensional oriented spin manifold $M$ of dimension $2 d$ provided with an action of $G$. We do not assume that $M$ is compact. Let $(\mathcal{L}, \nabla)$ be a $G$ equivariant line bundle equipped with a $G$-invariant hermitian connection $\nabla$.

We assume that the moment map $\phi_{G}: M \rightarrow \mathfrak{g}^{*}$ defined by the Kostant relation (3.1) is a proper map. In the next section we explain how is defined the formal geometric quantization of the data $\left(M, \phi_{G}, \mathcal{L}\right)$.

Notice that the twisted Duistermaat-Heckman distributions $D H^{G}(M, \Omega, \nu) \in \mathcal{D}^{\prime}\left(\mathfrak{g}^{*}\right)^{G}$ are still defined by the relation

$$
\begin{equation*}
\left\langle D H^{G}(M, \Omega, \nu), \varphi\right\rangle=\frac{1}{(-2 i \pi)^{d}} \int_{M} e^{-i \Omega}[\nu(-i \partial) \varphi]\left(\phi_{G}(m)\right) . \tag{4.1}
\end{equation*}
$$

Here $\nu$ is a closed equivariant form on $M$ with polynomial coefficients and $\varphi$ is a test function on $\mathfrak{g}^{*}$. The integral in (4.1) is well-defined since the differential form $e^{-i \Omega}[\nu(-i \partial) \varphi]\left(\phi_{G}(m)\right)$ has a compact support.

[^1]4.1. Formal geometric quantization: definition. The invariant scalar product in $\mathfrak{k}^{*}$ provides an identification $\mathfrak{k} \simeq \mathfrak{k}^{*}$.

Definition 4.1. - The Kirwan vector field associated to $\phi_{G}$ is defined by

$$
\begin{equation*}
\kappa_{G}(m)=-\phi_{G}(m) \cdot m, \quad m \in M . \tag{4.2}
\end{equation*}
$$

- We denote by $Z_{G}$ the set of zeroes of $\kappa_{G}$. Thus $Z_{G}$ is a $G$-invariant closed subset of $M$.

The set $Z_{G}$, which is not necessarily smooth, admits the following description. Consider the closed Weyl chamber $\mathfrak{t}_{\geq 0}^{*}$. We see that

$$
\begin{equation*}
Z_{G}=\coprod_{\gamma \in \mathcal{B}_{G}} Z_{\gamma} \tag{4.3}
\end{equation*}
$$

where $Z_{\gamma}$ corresponds to the compact set $G\left(M^{\gamma} \cap \phi_{G}^{-1}(\gamma)\right)$, and $\mathcal{B}_{G}=$ $\phi_{G}\left(Z_{G}\right) \cap \mathfrak{t}_{\geq 0}^{*}$. The properness of $\phi_{G}$ insures that, for any compact subset $C \subset \mathfrak{t}^{*}$, the intersection $\mathcal{B}_{G} \cap C$ is finite. Here $M^{\gamma}$ is the zeroes of the vector field on $M$ defined by the infinitesimal action of $\gamma$.

Let $\mathcal{S} \rightarrow M$ be the spinor bundle on $M$. The principal symbol of the Dirac operator $D_{\mathcal{S}}$ is the bundle map $\sigma(M) \in \Gamma\left(\mathrm{T}^{*} M, \operatorname{hom}\left(\mathcal{S}^{+}, \mathcal{S}^{-}\right)\right)$ defined by the Clifford action

$$
\sigma(M)(m, \nu)=\mathbf{c}_{m}(\tilde{\nu}):\left.\left.\mathcal{S}\right|_{m} ^{+} \rightarrow \mathcal{S}\right|_{m} ^{-}
$$

where $\nu \in \mathrm{T}^{*} M \simeq \tilde{\nu} \in \mathrm{~T} M$ is an identification associated to an invariant Riemannian metric on $M$.

Definition 4.2. The symbol $\sigma\left(M, \phi_{G}\right)$ shifted by the vector field $\kappa_{G}$ is the symbol on $M$ defined by

$$
\sigma\left(M, \phi_{G}\right)(m, \nu)=\sigma(M)\left(m, \tilde{\nu}-\kappa_{G}(m)\right)
$$

for any $(m, \nu) \in \mathrm{T}^{*} M$.
For any $G$-invariant open subset $\mathcal{U} \subset M$ such that $\mathcal{U} \cap Z_{G}$ is compact in $M$, we see that the restriction $\left.\sigma\left(M, \phi_{G}\right)\right|_{\mathcal{U}}$ is a transversally elliptic symbol on $\mathcal{U}$, and so its equivariant index is a well defined element in $\hat{R}(G)$ (see $[1,22]$ ).

Thus we can define the following localized equivariant indices.
Definition 4.3. - A closed invariant subset $Z \subset Z_{G}$ is called a component of $Z_{G}$ if it is a union of connected components of $Z_{G}$.

- If $Z$ is a compact component of $Z_{G}$, and $E \rightarrow M$ is an equivariant vector bundle, we denote by

$$
\mathcal{Q}_{G}^{\mathrm{spin}}(M, E, Z) \in \hat{R}(K)
$$

the equivariant index of $\left.\sigma\left(M, \phi_{G}\right) \otimes E\right|_{\mathcal{U}}$ where $\mathcal{U}$ is an invariant neighbourhood of $Z$ so that $\mathcal{U} \cap Z_{G}=Z$.

By definition, $Z=\emptyset$ is a component of $Z_{\mathcal{S}}$ and $\mathcal{Q}_{G}^{\text {spin }}(M, \mathcal{S}, \emptyset)=0$. For any $\gamma \in \mathcal{B}_{\mathcal{S}}, Z_{\gamma}$ is a compact component of $Z_{\mathcal{S}}$.

When the manifold $M$ is compact, the set $\mathcal{B}_{G}$ is finite and we have the decomposition $\mathcal{Q}_{G}^{\text {spin }}\left(M, \mathcal{L}^{k}\right)=\sum_{\gamma \in \mathcal{B}_{G}} \mathcal{Q}_{G}^{\text {spin }}\left(M, \mathcal{L}^{k}, Z_{\gamma}\right)$.

Definition 4.1. When the moment map $\phi_{G}$ is proper, we define the formal geometric quantization of the data $\left(M, \phi_{G}, \mathcal{L}^{k}\right)$ as

$$
\begin{equation*}
\mathcal{Q}_{G}^{-\infty}\left(M, \mathcal{L}^{k}\right):=\sum_{\gamma \in \mathcal{B}_{G}} \mathcal{Q}_{G}^{\text {spin }}\left(M, \mathcal{L}^{k}, Z_{\gamma}\right) \tag{4.4}
\end{equation*}
$$

The sum of the right hand side is not necessarily finite but it converges in $\hat{R}(G)$ (see $[10,14,18,19]$ ).

In the following example, for any $\lambda \in \mathbb{Z}$, we denote $\mathbb{C}_{[\lambda]}$ the vector space $\mathbb{C}$ with the action of $S^{1}: t \cdot z=t^{\lambda} z$, for $(t, z) \in S^{1} \times \mathbb{C}_{[\lambda]}$.

Example 4.2. Consider the $S^{1}$-manifold $M=\mathbb{C}_{[2]}$. The $S^{1}$-equivariant spinor bundle on $M$ is $\mathcal{S}=M \times\left(\mathbb{C}_{[-1]} \oplus \mathbb{C}_{[1]}\right)$. The Clifford action of $v \in \mathrm{~T}_{z} M$ is $\mathbf{c}(v) w=v w$ if $w \in \mathbb{C}_{[-1]}$ and $\mathbf{c}(v) w=-\bar{v} w$ if $w \in \mathbb{C}_{[1]}$.

Fix $a \in \mathbb{N}$. Consider the equivariant line bundle $\mathcal{L}(a)=M \times \mathbb{C}_{[a]}$ with connection $\nabla=d-\frac{i}{2} \operatorname{Im}(z d \bar{z})$. The two-form is $\Omega^{a}=\frac{1}{2} \operatorname{Im}(d z d \bar{z})$, the moment map $\phi_{S^{1}}^{a}(z)=a+|z|^{2}$ is proper, and the corresponding critical set $Z_{S^{1}}^{a}$ is reduced to $\{0\} \subset M$.
$A$ small computation shows that $\mathcal{Q}_{S^{1}}^{-\infty}\left(M, \mathcal{L}(a)^{k}\right)=\mathcal{Q}_{S^{1}}^{\text {spin }}\left(M, \mathcal{L}(a)^{k},\{0\}\right)$ is equal to $\mathbb{C}_{[k a]} \otimes \sum_{j \geq 0} \mathbb{C}_{[2 j+1]}$ (see [22]).

Let $m_{G}(\lambda, k)$ be the multiplicity of $V_{\lambda}$ in $\mathcal{Q}_{G}^{-\infty}\left(M, \mathcal{L}^{k}\right)$. In other words,

$$
\mathcal{Q}_{G}^{-\infty}\left(M, \mathcal{L}^{k}\right)=\sum_{\lambda \in A_{G}} m_{G}(\lambda, k) V_{\lambda}
$$

In this context, the multiplicities $m_{G}(\lambda, k)$ have still an interpretation in terms of reduced spaces (see [10]). Hence, when the generic infinitesimal stabilizer is abelian, Theorem 3.7 still holds, so $m_{G}(\lambda, k) \neq 0$ only if $\lambda / k \in \phi_{G}(M)$.

As in the previous section, we are interested to the asymptotic behaviour of the following family of distributions

$$
\Theta_{k}^{(M, \mathcal{L})}:=k^{r} \sum_{\lambda \in A_{G}} m_{G}(\lambda, k) \beta_{\lambda / k}
$$

where $r=\operatorname{dim}(G / T) / 2$. Note that for any test function $\varphi$ with compact support the identity

$$
\left\langle\Theta_{k}^{(M, \mathcal{L})}, \varphi\right\rangle:=k^{r} \sum_{\lambda \in A_{G}} m_{G}(\lambda, k)\left\langle\beta_{\lambda / k}, \varphi\right\rangle
$$

is well defined since there exists only a finite number of terms such that $\left\langle\beta_{\lambda / k}, \varphi\right\rangle \neq 0$.

Let us recall that we can associate the twisted Duistermaat-Heckman distributions $D H^{G}\left(M, \Omega, \hat{A}_{n}(M)\right) \in \mathcal{D}^{\prime}\left(\mathfrak{g}^{*}\right)^{G}$ to any equivariant form $\hat{A}_{n}(M)($ see (4.1)).

The aim of this section is to prove the following extension of Theorem 3.13 .

Theorem 4.3. Let $M$ be a spin even dimensional oriented manifold with a $G$-equivariant line bundle $(\mathcal{L}, \nabla)$. Suppose that the moment map $\phi_{G}$ is proper. When $k$ tends to $\infty$, we have the asymptotic expansion

$$
\begin{equation*}
\Theta_{k}^{(M, \mathcal{L})} \equiv j_{\mathfrak{g}}^{1 / 2}(i \partial / k)\left(k^{d} \sum_{n=0}^{\infty} \frac{1}{k^{n}} D H^{G}\left(M, \Omega, \hat{A}_{n}(M)\right)\right) . \tag{4.5}
\end{equation*}
$$

Example 4.4. We continue Example 4.2 with $M=\mathbb{C}_{[2]}$ and $\mathcal{L}(a)=$ $M \times \mathbb{C}_{[a]}$. Then $\Theta_{k}^{(M, \mathcal{L}(a))}=\sum_{j \geq 0} \delta_{a+(2 j+1) / k}$. The equivariant $\hat{A}$-class is $\widehat{A}(M)(X)=\frac{X}{\sin (X)}$ for $X \in \operatorname{Lie}\left(S^{1}\right) \simeq \mathbb{R}$. Identity (4.5) says that

$$
\Theta_{k}^{(M, \mathcal{L}(a))} \equiv \frac{k}{2} \frac{i \partial / k}{\sin (i \partial / k)} 1_{[a, \infty[.} .
$$

This is the formula given in Example 1.3.
With the same proof than Theorem 4.3 we obtain
Theorem 4.5. Let $M$ be a spin even dimensional oriented manifold with a $G$-equivariant line bundle $\mathcal{L}$. Suppose that the moment map $\phi_{G}$ is proper. Let $g \in G$ of finite order. When $k$ tends to $\infty$, we have the asymptotic equivalence

$$
\begin{gathered}
\operatorname{rescaling}_{\mathfrak{g}(g)^{*}}(k)\left(\sum_{\lambda \in \Lambda} m(\lambda, k) \beta(g, \lambda)\right) \equiv \\
j_{\mathfrak{g}(g)}^{\frac{1}{2}}(i \partial / k) \operatorname{det}_{\mathfrak{g} / \mathfrak{g}(g)}^{\frac{1}{2}}\left(1-g^{-1} e^{-i \partial / k}\right) \sum_{n=0}^{\infty} \frac{1}{k^{n}} k^{\frac{\operatorname{dim} M^{g}}{2}} D H^{G(g)}\left(M^{g}, \Omega_{g}, V_{g}(n, k)\right)
\end{gathered}
$$

We prove Theorem 4.3 in the next two subsections.
4.2. Formal geometric quantization: delocalized formulas. We consider the generalized character

$$
\mathcal{Q}_{G}^{\mathrm{spin}}\left(M, \mathcal{L}^{k}, Z_{\gamma}\right)=\sum_{\lambda \in A_{G}} m_{\gamma}(\lambda, k) V_{\lambda}
$$

Here $\mathcal{Q}_{G}^{\text {spin }}\left(M, \mathcal{L}^{k}, Z_{\gamma}\right)(g)=\sum_{\lambda \in A_{G}} m_{\gamma}(\lambda, k) \chi_{\lambda}(g)$ defines a distribution on $G$ because the multiplicity function $\lambda \mapsto m_{\gamma}(\lambda, k)$ has at most a polynomial growth.

The function $m_{\gamma}(\lambda, k)$ is in fact with at most polynomial growth in both variables $(\lambda, k)$. This follows for example from the general multiplicity formula for transversally elliptic operators given in [4]. In this formula, the multiplicity formula is given for a single general transversally elliptic symbol $\sigma$ and multiplicities are obtained as values on $A_{G}$ of of a certain piecewise quasi polynomial functions on $A_{G}$. If we consider a family of the form $\sigma_{k}=\sigma \otimes \mathcal{L}^{k}$, the formula is locally piecewise polynomial on a certain finite number of affine cones in $\mathfrak{t}_{\geq 0}^{*} \oplus \mathbb{R}$ intersected with $A_{G} \oplus \mathbb{Z}$.

We consider the corresponding distribution on $\mathfrak{g}$

$$
\mathcal{Q}_{G}^{\mathrm{spin}}\left(M, \mathcal{L}^{k}, Z_{\gamma}\right)\left(e^{X}\right)=\sum_{\lambda \in A_{G}} m_{\gamma}(\lambda, k) \chi_{\lambda}\left(e^{X}\right)
$$

The distribution on $\mathfrak{g}^{*}$ defined by

$$
\Theta_{k}^{\gamma}:=k^{r} \sum_{\lambda \in A_{G}} m_{\gamma}(\lambda, k) \beta_{\lambda / k}
$$

is the Fourier transform of $\mathcal{Q}_{G}^{\text {spin }}\left(M, \mathcal{L}^{k}, Z_{\gamma}\right)\left(e^{X / k}\right) j_{\mathfrak{g}}^{1 / 2}(X / k)$. It is tempered also because the multiplicity function $\lambda \mapsto m_{\gamma}(\lambda, k)$ has at most a polynomial growth.

In order to give a formula for the asymptotic of $\Theta_{k}^{\gamma}$ we need to introduced an equivariant form $P_{\gamma}(X)$ with generalized coefficient on $M$ that is supported in a small neighbourhood of $Z_{\gamma}$.

Take $\chi: M \rightarrow \mathbb{R}$ a $G$-invariant function equal to 1 in a neighbourhood of $Z_{\gamma}$, compactly supported, and such that $\operatorname{Support}(\chi) \cap Z_{G}=Z_{\gamma}$. Consider the invariant 1-form $\theta:=\left(\kappa_{G},-\right)$ where $\kappa_{G}$ is the Kirwan vector field associated to $\phi_{G}$, and $(-,-)$ is an invariant Riemannian metric.

We consider the equivariant form $D \theta(X)=d \theta-\left\langle\Phi_{\theta}, X\right\rangle$, and the equivariant form with generalized coefficients

$$
P_{\gamma}(X):=\chi+i d \chi \theta \int_{0}^{\infty} e^{-i t D \theta(X)} d t
$$

Thus, if $f(X)$ is in the Schwartz space of $\mathfrak{g}$, the integral

$$
\int_{\mathfrak{g}} P_{\gamma}(X) f(X) d X:=\chi \hat{f}(0)+i d \chi \theta \int_{0}^{\infty} e^{-i t d \theta} \hat{f}\left(t \Phi_{\theta}\right) d t
$$

is convergent. Indeed $\left\langle\Phi_{\theta}, \phi_{G}\right\rangle(m)=\left\|\kappa_{G}(m)\right\|^{2}$ and so $\Phi_{\theta} \neq 0$ on the support of $d \chi$. It defines a differential form with compact support since it is equal to 0 outside the support of $\chi$. So if $f(X)$ is a function on $\mathfrak{g}$ to differential forms on $M$ such that $X \rightarrow f(X)$ is rapidly decreasing, then we can define $\int_{M} \int_{\mathfrak{g}} P_{\gamma}(X) f(X) d X$. Let $\varphi$ a test function on $\mathfrak{g}^{*}$, and $\nu(X)$ an equivariant form. Then $e^{-i \Omega(X)} \nu(X) \widehat{\varphi}(X)$ is rapidly decreasing if $\nu(X)$ is an equivariant form with polynomial coefficients. So we define $D H^{G}\left(M, \Omega, \nu P_{\gamma}\right) \in \mathcal{D}^{\prime}\left(\mathfrak{g}^{*}\right)^{G}$ by

$$
\left\langle D H^{G}\left(M, \Omega, \nu P_{\gamma}\right), \varphi\right\rangle=\frac{1}{(-2 i \pi)^{d}} \int_{M} \int_{\mathfrak{g}} e^{-i \Omega(X)} \nu(X) P_{\gamma}(X) \widehat{\varphi}(X) .
$$

An improved version of Witten non abelian localization theorem is the following formula

$$
\begin{equation*}
D H^{G}(M, \Omega, \nu)=\sum_{\gamma \in \mathcal{B}_{G}} D H^{G}\left(M, \Omega, \nu P_{\gamma}\right) \tag{4.6}
\end{equation*}
$$

which is somehow a consequence of the fact that $\sum_{\gamma \in \mathcal{B}_{G}} P_{\gamma}$ is equal to 1 in cohomology (see [16]).

We will prove the following theorem.
Theorem 4.4. We have the following relation

$$
\Theta_{k}^{\gamma} \equiv k^{d} j_{\mathfrak{g}}^{1 / 2}(i \partial / k) \sum_{n=0}^{\infty} \frac{1}{k^{n}} D H^{G}\left(M, \Omega, \hat{A}_{n}(M) P_{\gamma}\right) .
$$

We see that $P_{\gamma}(X)=\chi+\sum_{j=1}^{d} \omega_{j} \int_{0}^{\infty} t^{j-1} e^{i t\left\langle\Phi_{\theta}, X\right\rangle} d t$ where $\omega_{j}$ is a differential form, compactly supported, of degree $2 j$. Then

$$
E(X, k):=P_{\gamma}(X / k)=\chi+\sum_{j=1}^{d} k^{j} \omega_{j} \int_{0}^{\infty} t^{j-1} e^{i t\left\langle\Phi_{\theta}, X\right\rangle} d t, \quad k \geq 1
$$

is a polynomial in $k$ with value equivariant forms with generalized coefficients.

Let $U_{\gamma} \subset M$ be a relatively compact invariant subset containing the support of $\chi$. So we can choose $r_{\gamma}>0$ so that $\hat{A}(M)(X)$ is well defined on $U_{\gamma}$ when $\|X\|<r_{\gamma}$.

We will deduce Theorem 4.4 from the following result proved in [20].

Proposition 4.5. If $\|X\|<r_{\gamma}$, we have the relation

$$
\mathcal{Q}_{G}^{\mathrm{spin}}\left(M, \mathcal{L}^{k}, Z_{\gamma}\right)\left(e^{X}\right)=\frac{1}{(-2 i \pi)^{d}} \int_{M} \hat{A}(M)(X) P_{\gamma}(X) e^{-i k \Omega(X)}
$$

So for $k \geq 1$, the generalized function $\mathcal{Q}_{G}^{\text {spin }}\left(M, \mathcal{L}^{k}, Z_{\gamma}\right)\left(e^{X / k}\right) j_{\mathfrak{g}}^{1 / 2}(X / k)$ coincides with

$$
\frac{1}{(-2 i \pi)^{d}} \int_{M} j_{\mathfrak{g}}^{1 / 2}(X / k) \hat{A}(M)(X / k) P_{\gamma}(X / k) e^{-i k \Omega(X / k)}
$$

on the ball $\left\{\|X\| / k<r_{\gamma}\right\}$.
Let us compute the Laurent series

$$
\left[\hat{A}(M)(X / k) j_{\mathfrak{g}}^{1 / 2}(X / k) P_{\gamma}(X / k) e^{-i k \Omega(X / k)}\right]_{[2 d]}
$$

where $[-]_{[2 d]}$ means the component of maximal degree $2 d$ in $\mathcal{A}^{\bullet}(M)$. We write $j_{\mathfrak{g}}^{1 / 2}(X) \hat{A}(M)(X)=\sum_{n=0}^{\infty} d_{n}(X)$ as a sum of closed equivariant forms of equivariant degree $n$. Recall that $E(X, k)=P_{\gamma}(X / k)$ and $b(X, k)=e^{-i k \Omega(X / k)}$ depend polynomially of $k$.

We write also $\hat{A}(M)(X) j_{\mathfrak{g}}^{1 / 2}(X)=\sum_{n=0}^{\infty} \tau_{n}(X)$ where $\tau_{n}(X)$ is an equivariant form with coefficients homogeneous polynomial functions of $X$ of degree $n$. We proceed as in the proof of Lemma 3.10 comparing the terms of top exterior degree $2 d$. With same proof as Lemma 3.11, we have the following formula.

## Lemma 4.6.

$$
\left[b(X, k)\left(\sum_{n=0}^{\infty} \frac{1}{k^{n}} \tau_{n}(X)\right) E(X, k)\right]_{[2 d]}=\left[e^{-i \Omega(X)}\left(\sum_{n=0}^{\infty} \frac{1}{k^{n}} d_{n}(X)\right) P_{\gamma}(X)\right]_{[2 d]}
$$

Thus we can write formally

$$
\mathcal{Q}_{G}^{\text {spin }}\left(M, \mathcal{L}^{k}, Z_{\gamma}\right)\left(e^{X / k}\right) j_{\mathfrak{g}}^{1 / 2}(X)=k^{d} \sum_{n=0}^{\infty} \frac{1}{k^{n}} g_{n}(X)
$$

where $g_{n}(X)$ is the distribution

$$
g_{n}(X):=\frac{1}{(-2 i \pi)^{d}} \int_{M} d_{n}(X) P_{\gamma}(X) e^{-i \Omega(X)}
$$

4.3. Proof of Theorem 4.4. We fix a $G$-invariant function $h: \mathfrak{g} \rightarrow \mathbb{R}$ equal to 1 in a neighbourhood of 0 and with support contained in $\|X\|<r_{\gamma}$. If $\varphi$ is a smooth function with compact support on $\mathfrak{g}^{*}$, we have $\left\langle\Theta_{k}^{\gamma}, \varphi\right\rangle=\mathrm{I}_{k}+\mathrm{J}_{k}$ where

$$
\mathrm{I}_{k}=\int_{\mathfrak{g}} \mathcal{Q}_{G}^{\text {spin }}\left(M, \mathcal{L}^{k}, Z_{\gamma}\right)\left(e^{X / k}\right) j_{\mathfrak{g}}^{1 / 2}(X / k) h(X / k) \widehat{\varphi}(X) d X
$$

and

$$
\mathrm{J}_{k}=\int_{\mathfrak{g}} \mathcal{Q}_{G}^{\mathrm{spin}}\left(M, \mathcal{L}^{k}, Z_{\gamma}\right)\left(e^{X / k}\right) j_{\mathfrak{g}}^{1 / 2}(X / k)(1-h(X / k)) \widehat{\varphi}(X) d X
$$

For estimating $\mathrm{J}_{k}$, we will use estimation of the Fourier coefficients $m_{\gamma}(\lambda, k)$ and we will prove that $\mathrm{J}_{k}=O\left(k^{-\infty}\right)$. For estimating $\mathrm{I}_{k}$, we will use Proposition (4.5).

To analyse these expressions, we need the following technical lemma.
Lemma 4.7. Let $u(X)$ be a $C^{\infty}$ function of $X$. Assume that $u(X)$ and all its derivatives are of at most polynomial growth and that $u(X)$ vanishes at order $N$ at $X=0$.

Let $R$ be an integer. Then there exists a constant $c_{R}$ such that

$$
\left|\int_{\mathfrak{g}} u(X / k) e^{i\langle\xi, X\rangle} \widehat{\varphi}(X) d X\right| \leq \frac{1}{k^{N+1}} \frac{c_{R}}{\left(1+\|\xi\|^{2}\right)^{R}}
$$

for all $(\xi, k) \in \mathfrak{g}^{*} \times \mathbb{N} \backslash\{0\}$.
Proof. We may write in multi-index notation

$$
u(X)=\sum_{\alpha,|\alpha|=N+1} X^{\alpha} v_{\alpha}(X)
$$

with $v_{\alpha}(X)$ (and derivatives) bounded by polynomial functions of $X$. We compute $I_{R}=k^{N+1}\left(1+\|\xi\|^{2}\right)^{R} \int_{\mathfrak{g}} u(X / k) e^{i\langle\xi, X\rangle} \widehat{\varphi}(X) d X$. Thus

$$
I_{R}=\left(1+\|\xi\|^{2}\right)^{R} \int_{\mathfrak{g}} \sum_{|\alpha|=N+1} X^{\alpha} v_{\alpha}(X / k) e^{i\langle\xi, X\rangle} \widehat{\varphi}(X) d X
$$

Let $\Delta=-\sum_{j} \partial_{X_{j}}^{2}$ be the Laplacian on $\mathfrak{g}$. Then $I_{R}$ is equal to $\sum_{|\alpha|=N+1} \int_{\mathfrak{g}} X^{\alpha} v_{\alpha}(X / k) \widehat{\varphi}(X)\left(\left(1+\Delta^{2}\right)^{R} \cdot\left(e^{i\langle\xi, X\rangle}\right) d X=\int_{\mathfrak{g}} L_{R}(X, k) e^{i\langle\xi, X\rangle} d X\right.$ where $L_{R}(X, k)=\left(1+\Delta^{2}\right)^{R} \cdot\left(\sum_{|\alpha|=N+1} X^{\alpha} v_{\alpha}(X / k) \widehat{\varphi}(X)\right)$.

Using the fact that $\widehat{\varphi}(X)$ is rapidly decreasing, that the derivatives of the functions $v_{\alpha}$ are bounded by polynomials, and that $1 / k \leq 1$ for $k \geq 1$, we see that $L_{R}(X, k)$ can be bounded by a rapidly decreasing function of $X$ independent of $k$. So $\left|\int_{\mathfrak{g}} L_{R}(X, k) e^{i\langle\xi, X\rangle} d X\right| \leq c_{R}$.

We return to our proof. We start by checking that $\mathrm{J}_{k}=O\left(k^{-\infty}\right)$. For this computation, we can assume that $\varphi$ is $G$-invariant. Let

$$
c(\lambda, k)=\int_{\mathfrak{g}} \chi_{\lambda}\left(e^{X / k}\right) j_{\mathfrak{g}}^{1 / 2}(X / k)(1-h(X / k)) \widehat{\varphi}(X) d X
$$

This is equal to $\operatorname{vol}(G \lambda) \int_{\mathfrak{g}} e^{i\langle\lambda / k, X\rangle}(1-h(X / k)) \widehat{\varphi}(X) d X$.

By definition $\mathcal{Q}_{G}^{\text {spin }}\left(M, \mathcal{L}^{k}, Z_{\gamma}\right)\left(e^{X}\right)=\sum_{\lambda \in A_{G}} m_{\gamma}(\lambda, k) \chi_{\lambda}\left(e^{X}\right)$, so we get

$$
\mathrm{J}_{k}=\sum_{\lambda \in A_{G}} m_{\gamma}(\lambda, k) \operatorname{vol}(G \lambda) \int_{\mathfrak{g}} e^{i\langle\lambda / k, X\rangle}(1-h(X / k)) \widehat{\varphi}(X) d X
$$

Now consider $u(X)=1-h(X)$. It vanishes identically in a neighbourhood of 0 . As $h(X)$ is compactly supported, $u(X)$ and all derivatives are bounded. So we can apply Lemma 4.7 and obtain $|c(\lambda, k)| \leq$ $\operatorname{vol}(G \lambda) \frac{1}{k^{N+1}} \frac{c_{R, N}}{\left(1+\|\lambda / k\|^{2}\right)^{R}}$ for any integers $(R, N)$. Since the multiplicities $m_{\gamma}(\lambda, k)$ have at most a polynomial growth in the variable $(\lambda, k)$, we can conclude that $\mathrm{J}_{k}=O\left(k^{-\infty}\right)$.

We now estimate

$$
\mathrm{I}_{k}=\int_{\mathfrak{g}} \mathcal{Q}_{G}^{\text {spin }}\left(M, \mathcal{L}^{k}, Z_{\gamma}\right)\left(e^{X / k}\right) j_{\mathfrak{g}}^{1 / 2}(X / k) h(X / k) \widehat{\varphi}(X) d X
$$

As $h(X / k)=0$ when $\|X\| / k \geq r_{\gamma}$, we may use Proposition 4.5. Let $\tau^{h}(X)=h(X) \hat{A}(M)(X) j_{\mathfrak{g}}^{1 / 2}(X)$, a smooth compactly supported function from $\mathfrak{g}$ to differential form on $U_{\gamma}$. We fix $x \in U_{\gamma}$ and consider $\tau_{x}^{h}(X)$ a function from $X \in \mathfrak{g}$ to $\Lambda T_{x}^{*} M$. Define

$$
\mathrm{I}_{k}^{x}=\int_{\mathfrak{g}} \tau_{x}^{h}(X / k) E_{x}(X, k) b_{x}(X, k) \widehat{\varphi}(X) d X
$$

with

$$
E_{x}(X, k)=\chi(x)+\left.\sum_{j=1}^{d} k^{j}\left(\omega_{j}\right)\right|_{x} \int_{0}^{\infty} t^{j-1} e^{i t\langle\zeta, X\rangle} d t
$$

if $\zeta=\Phi_{\theta}(x)$ and $b_{x}(X, k)=e^{-i k \Omega_{x}} e^{i\langle\xi, X\rangle}$ if $\xi=\phi_{G}(x)$. So

$$
\mathrm{I}_{k}^{x}=\int_{M} \int_{\mathfrak{g}} \tau_{x}^{h}(X / k) E_{x}(X, k) b_{x}(X, k) \widehat{\varphi}(X) d X
$$

In view of Lemma 4.6, we only need to prove that the asymptotic expansion of $\mathrm{I}_{k}^{x}$ is obtained by replacing $\tau_{x}^{h}(X)$ by its Taylor series (we need to care of uniform estimates in $x$ in the compact support of $\chi$ ).

Thus we fix $x$ and write $E_{x}(X, k)=\chi(x)+R_{x}(X, k)$. So $\mathrm{I}_{k}^{x}=\mathrm{S}_{k}^{x}+\mathrm{T}_{k}^{x}$ with

$$
\mathrm{S}_{k}^{x}=\chi(x) \int_{\mathfrak{g}} \tau_{x}^{h}(X / k) b_{x}(X, k) \widehat{\varphi}(X) d X
$$

and

$$
\mathrm{T}_{k}^{x}=\int_{\mathfrak{g}} R_{x}(X, k) \tau_{x}^{h}(X / k) b_{x}(X, k) \widehat{\varphi}(X) d X
$$

The expression for $S_{k}^{x}$ has been already analysed when proving Proposition 3.13, and indeed it admits an asymptotic expansion obtained by replacing $\tau_{x}(X)$ by its Taylor series.

Consider $a(X, t)=e^{i\langle\xi, X\rangle} e^{i t\langle\zeta, X\rangle} P(t)$ where $P(t)$ is a polynomial function on $t$ with values in $\Lambda T_{x}^{*} M$. We only need to prove

$$
\mathrm{W}_{k}^{x}=\int_{t=0}^{\infty} \int_{\mathfrak{g}} a(X, t) \tau_{x}^{h}(X / k) \widehat{\varphi}(X) d X
$$

admits an asymptotic expansion also obtained by replacing $\tau_{x}^{h}(X)$ by its Taylor series. We fix $x \in M$, write $\tau_{x}^{h}(X)=\tau_{\leq N}(X)+\tau_{>N}(X)$. Then $\tau_{\leq N}(X)$ vanishes at order $N$ at $X=0$. As $\tau(\bar{X})$ was compactly supported, $\tau_{\leq N}(X)$ and $\tau_{>N}(X)$ are of at most polynomial growth, as well as derivatives. So we use Lemma 4.7 and obtain, for any positive integer $R$

$$
\int_{\mathfrak{g}} a(X, t) \tau_{>N}(X / k) \widehat{\varphi}(X) d X \leq P(t) \frac{1}{k^{N+1}} \frac{c_{N, R}}{\left(1+|\xi+t \zeta|^{2}\right)^{R}}
$$

If $R$ is sufficiently large, $\int_{t=0}^{\infty} P(t) \frac{c_{N, R}}{\left(1+|\xi+t \zeta|^{2}\right)^{R}}<\infty$. So we obtain our estimate for the rest. All our estimates can be done uniformly in $x$ when $x$ runs in the compact support of $\chi$. This ends the proof of Theorem 4.4.
4.4. Proof of Theorem 4.3. The decomposition $\Theta_{k}^{(M, \mathcal{L})}=\sum_{\gamma} \Theta_{k}^{\gamma}$ is well-defined since the distribution $\Theta_{k}^{\gamma}$ is supported in $\{\|\xi\| \geq\|\gamma\|\}$ (see [18]). If $\varphi$ is a test function, then

$$
\left\langle\Theta_{k}^{(M, \mathcal{L})}, \varphi\right\rangle=\sum_{\gamma}\left\langle\Theta_{k}^{\gamma}, \varphi\right\rangle
$$

where the sum in the right hand side has only a finite number of nonzero terms.

Thanks to Theorem 4.4, we have the asymptotic expansion

$$
\Theta_{k}^{\gamma} \equiv k^{d} j_{\mathfrak{g}}^{1 / 2}(i \partial / k) \sum_{n=0}^{\infty} \frac{1}{k^{n}} D H^{G}\left(M, \Omega, \hat{A}_{n}(M) P_{\gamma}\right)
$$

for any $\gamma$. Hence $\Theta_{k}^{(M, \mathcal{L})} \equiv k^{d} j_{\mathfrak{g}}^{1 / 2}(i \partial / k) \sum_{n=0}^{\infty} k^{-n} \theta_{n}$ with $\theta_{n}$ is equal to $\sum_{\gamma} D H^{G}\left(M, \Omega, \hat{A}_{n}(M) P_{\gamma}\right)$. The proof of Theorem 4.3 is complete since $\sum_{\gamma} D H^{G}\left(M, \Omega, \hat{A}_{n}(M) P_{\gamma}\right)=D H^{G}\left(M, \Omega, \hat{A}_{n}(M)\right)$ by Equation (4.6).

## 5. Functoriality

5.1. $H$ and $G$. Let $H$ be a connected compact subgroup of $G$. Let $r: \mathfrak{g}^{*} \rightarrow \mathfrak{h}^{*}$ be the projection. For $\theta$ a distribution with compact support on $\mathfrak{g}^{*}$, the push-forward $\left\langle r_{*} \theta, \varphi\right\rangle=\langle\theta, \varphi \circ r\rangle$ is well defined. The Fourier transform $\mathcal{F}_{\mathfrak{h}^{*}}\left(r_{*} \theta\right)$ is the restriction to $\mathfrak{h}$ of the Fourier transform $\mathcal{F}_{\mathfrak{g}^{*}}(\theta)$. We can define $r_{*} \theta$ more generally whenever $\theta$ is compactly supported along the fibers of $r$.

The twisted Duistermaat-Heckman distributions behave very well under the push-forward map $r_{*}$.

From (4.1) we get immediately the following
Proposition 5.1. Suppose that the moment map $\phi_{H}:=r \circ \phi_{G}$ is proper. Then for any $\nu \in H_{G}^{*}(M)$ we have

$$
r_{*} D H^{G}(M, \Omega, \nu)=D H^{H}(M, \Omega, \nu) .
$$

Rescaling behave also very well under the push-forward map $r_{*}$ :

$$
\operatorname{rescaling}_{\mathfrak{h}^{*}}(k) \circ r_{*}=r_{*} \circ \operatorname{rescaling}_{\mathfrak{g}^{*}}(k) .
$$

In the rest of the article, objects associated to $\lambda$ are associated to $G$, while objects associated to $\mu$ are associated to $H$.

For $\mu \in A_{H}$ and $\lambda \in A_{G}$, let $c(\mu, \lambda)$ be the multiplicity of the representation $V_{\mu}$ of $H$ in the restriction of $V_{\lambda}$ to $H$. In other words, for $h \in H, \chi_{\lambda}(h)=\sum_{\mu \in A_{H}} c(\mu, \lambda) \chi_{\mu}(h)$.

Consider the $H$-invariant function $j_{\mathfrak{g} / \mathfrak{h}}^{1 / 2}(Y)=\operatorname{det}_{\mathfrak{g} / \mathfrak{h}}^{1 / 2}\left(\frac{e^{Y / 2}-e^{-Y / 2}}{Y}\right)$ on $\mathfrak{h}$. Its Fourier transform is a compactly supported measure $B_{\mathfrak{g} / \mathfrak{h}}$ on $\mathfrak{h}^{*}$. We have an exact relation between the push-forward of the measure $d \beta_{\lambda}$ on $\mathfrak{h}^{*}$ and measures of $H$-admissible coadjoint orbits. We denote $B_{\mathfrak{g} / \mathfrak{h}}^{k}=\operatorname{rescaling}_{\mathfrak{h}^{*}}(k) B_{\mathfrak{g} / \mathfrak{h}}$.
Lemma 5.2. For $\lambda \in A_{G}$, we have the relations

$$
\begin{aligned}
r_{*}\left(\beta_{\lambda}\right) & =B_{\mathfrak{g} / \mathfrak{h}} \star\left(\sum_{\mu \in A_{H}} c(\mu, \lambda) \beta_{\mu}\right) \\
\operatorname{rescaling}_{\mathfrak{h}^{*}}(k)\left(r_{*}\left(\beta_{\lambda}\right)\right) & =B_{\mathfrak{g} / \mathfrak{h}}^{k} \star\left(\operatorname{rescaling}_{\mathfrak{h}^{*}}(k)\left(\sum_{\mu \in A_{H}} c(\mu, \lambda) \beta_{\mu}\right)\right) .
\end{aligned}
$$

Here the $\star$ sign denotes the convolution.
Proof. The first identity follows immediately by Fourier transform of the formula: $\chi_{\lambda}\left(e^{Y}\right) j_{\mathfrak{g}}^{1 / 2}(Y)=j_{\mathfrak{g} / \mathfrak{h}}^{1 / 2}(Y) \sum_{\mu \in A_{H}} c(\mu, \lambda) \chi_{\mu}\left(e^{Y}\right) j_{\mathfrak{h}}^{1 / 2}(Y)$, for $Y \in \mathfrak{h}$.

We get the second identity by applying the operator rescaling $\mathfrak{h}^{*}(k)$ to the first one.

We now prove a similar descent formula. We write $\mathfrak{g}=\mathfrak{h} \oplus \mathfrak{q}$ where $\mathfrak{q}$ is $H$-invariant. Let $g \in H$ and consider the tempered function

$$
d(Y):=\chi_{\lambda}\left(g e^{Y}\right) j_{\mathfrak{g}(g)}^{1 / 2}(Y) \operatorname{det}_{\mathfrak{g} / \mathfrak{g}(g)}^{1 / 2}\left(1-g^{-1} e^{-Y}\right), \quad Y \in \mathfrak{g}(g)
$$

In one hand $d(Y)=\left\langle\beta(g, \lambda), e^{i\langle-, Y\rangle}\right\rangle$ and for $Y \in \mathfrak{h}(g)$ we have

$$
d(Y)=j_{\mathfrak{g}(g) / \mathfrak{h}(g)}^{1 / 2}(Y) \operatorname{det}_{\mathfrak{q} / \mathfrak{q}(g)}^{1 / 2}\left(1-g^{-1} e^{-Y}\right) \sum_{\mu} c(\lambda, \mu)\left\langle\beta(g, \mu), e^{i\langle-, Y\rangle}\right\rangle
$$

If $A$ is a compactly supported distribution on $\mathfrak{g}(g)^{*}$ we still denote $r_{*}(A)$ the push-forward distribution on $\mathfrak{h}(g)^{*}$. We denote $B_{\mathfrak{g}(g) / \mathfrak{h}(g)}$ the compactly supported measure on $\mathfrak{h}(g)^{*}$ which is the Fourier transform of the $H(g)$-invariant function $j_{\mathfrak{g}(g) / \mathfrak{h}(g)}^{1 / 2}$. We denote $B_{\mathfrak{g}(g) / \mathfrak{h}(g)}^{k}:=$ rescaling $_{\mathfrak{h}(g)^{*}}(k) B_{\mathfrak{g}(g) / \mathfrak{h}(g)}$. Let $C_{\mathfrak{q} / \mathfrak{q}(g)}$ be the Fourier transform of the $H(g)$-invariant function $Y \mapsto \operatorname{det}_{\mathfrak{q} / \mathfrak{q}(g)}^{1 / 2}\left(1-g^{-1} e^{-Y}\right)$. We denote $C_{\mathfrak{q} / \mathfrak{q}(g)}^{k}:=$ rescaling $_{\mathfrak{h}(g)^{*}}(k) C_{\mathfrak{q} / \mathfrak{q}(g)}$.

The previous identities give the following
Lemma 5.3. Let $g \in H$. Then rescaling $_{\mathfrak{h}(g)^{*}}(k)\left(r_{*} \beta(g, \lambda)\right)$ is equal to

$$
B_{\mathfrak{g}(g) / \mathfrak{h}(g)}^{k} \star C_{\mathfrak{q} / \mathfrak{q}(g)}^{k} \star\left(\operatorname{rescaling}_{\mathfrak{h}(g)^{*}}(k)\left(\sum_{\mu \in A_{H}} c(\mu, \lambda) \beta(g, \mu)\right)\right) .
$$

5.2. Functoriality relatively to restrictions. Here we consider an oriented spin manifold $M$ of even dimension with an action of a compact connected Lie group $G$. Let $\mathcal{L} \rightarrow M$ be a $G$-equivariant line bundle equipped with an invariant hermitian connection $\nabla$. We assume here that $\phi_{G}$ is proper. In this case we can define $\mathcal{Q}_{G}^{-\infty}\left(M, \mathcal{L}^{k}\right) \in \hat{R}(G)$ for any $k \geq 1$.

The main result of this section is the following result
Theorem 5.1. Suppose that $H \subset G$ is closed connected subgroup such that $\phi_{H}$ is proper. Then
(1) $\mathcal{Q}_{G}^{-\infty}\left(M, \mathcal{L}^{k}\right)$ is $H$-admissible,
(2) $\left.\mathcal{Q}_{G}^{-\infty}\left(M, \mathcal{L}^{k}\right)\right|_{H}=\mathcal{Q}_{H}^{-\infty}\left(M, \mathcal{L}^{k}\right)$.

We will prove Theorem 5.1 under the assumption that the generic infinitesimal stabilizer of the $G$-action on $M$ is abelian. It can be shown easily that it implies the general case (see Lemma 4.2 in [19]).

Let us denote $c(\lambda, \mu)$ the multiplicity of $V_{\mu}$ in $\left.V_{\lambda}\right|_{H}$. If we consider the generalized characters $\mathcal{Q}_{G}^{-\infty}\left(M, \mathcal{L}^{k}\right)=\sum_{\lambda \in \hat{G}} m_{G}(\lambda, k) V_{\lambda}$ and
$\mathcal{Q}_{H}^{-\infty}\left(M, \mathcal{L}^{k}\right)=\sum_{\mu \in \hat{H}} m_{H}(\mu, k) V_{\mu}$, we see that Theorem 5.1 is equivalent of the following

Theorem 5.2. For any $\mu \in A_{H}$ we have

$$
m_{H}(\mu, k)=\sum_{\lambda \in A_{G}} m_{G}(\lambda, k) c(\lambda, \mu),
$$

where the right hand side is a finite sum.
We consider the sequences of distributions $\Theta_{k}^{G}=k^{r} \sum_{\lambda} m_{G}(\lambda, k) \beta_{\lambda / k}$ and $\Theta_{k}^{H}=k^{r^{\prime}} \sum_{\mu} m_{H}(\mu, k) \beta_{\mu / k}$ associated to formal geometric quantizations $\mathcal{Q}_{G}^{-\infty}\left(M, \mathcal{L}^{k}\right)$ and $\mathcal{Q}_{H}^{-\infty}\left(M, \mathcal{L}^{k}\right)$.

We start with
Proposition 5.3. - We have $B_{\mathfrak{g} / \mathfrak{h}}^{k} \star \Theta_{k}^{H} \equiv j_{\mathfrak{g} / \mathfrak{b}}^{1 / 2}(i \partial / k) \Theta_{k}^{H}$.

- The push-forward $r_{*}\left(\Theta_{k}^{G}\right)$ is a well defined sequence of distributions on $\mathfrak{h}^{*}$. We have

$$
r_{*}\left(\Theta_{k}^{G}\right) \equiv j_{\mathfrak{g} / \mathfrak{h}}^{1 / 2}(i \partial / k) \Theta_{k}^{H} .
$$

Proof. Since $m_{H}(\mu, k)$ is a piecewise quasi-polynomial function (see Section 3.1), the first point is a consequence of Proposition 3.2.

We now prove the second point. Let $\varphi \in \mathcal{C}^{\infty}\left(\mathfrak{h}^{*}\right)$ be a function with compact support. Let $R_{\varphi}>0$ such that the compact set $\mathcal{K}_{\varphi}:=$ $\phi_{G}\left(\phi_{H}^{-1}(\operatorname{Support}(\varphi))\right.$ is contained in $\left\{\xi \in \mathfrak{g}^{*},\|\xi\| \leq R_{\varphi}\right\}$. The expression

$$
\left\langle r_{*}\left(\Theta_{k}^{G}\right), \varphi\right\rangle:=k^{r} \sum_{\lambda} m_{G}(\lambda, k) \varphi(r(\lambda / k)) .
$$

is well-defined since the term $m_{G}(\lambda, k) \varphi(r(\lambda / k))$ is non-zero for a finite number of $\lambda$. More precisely, since $m_{G}(\lambda, k) \neq 0$ only if $\lambda / k \in \phi_{G}(M)$, we see $m_{G}(\lambda, k) \varphi(r(\lambda / k)) \neq 0$ only if $\lambda / k \in \mathcal{K}_{\varphi}$.

If $\epsilon \in \mathcal{C}^{\infty}\left(\mathfrak{g}^{*}\right)$ is a function with compact support such that $\epsilon(\xi)=1$ if $\|\xi\| \leq R_{\varphi}$, we defined $\tilde{\varphi}:=\varphi \circ r \times \epsilon$. We see that

$$
\left\langle r_{*}\left(\Theta_{k}^{G}\right), \varphi\right\rangle=\left\langle\Theta_{k}^{G}, \tilde{\varphi}\right\rangle=k^{d} \sum_{n=0}^{N} \frac{1}{k^{n}}\left\langle\theta_{n}^{G}, \tilde{\varphi}\right\rangle+o\left(k^{d-N}\right)
$$

where $\left(\theta_{n}^{G}\right)$ is the family of distributions on $\mathfrak{g}^{*}$ such that

$$
j_{\mathfrak{g}}^{1 / 2}(i \partial / k) \sum_{n=0}^{\infty} \frac{1}{k^{n}} D H^{G}\left(M, \Omega, \hat{A}_{n}(M)\right)=\sum_{n=0}^{\infty} \frac{1}{k^{n}} \theta_{n}^{G} .
$$

We check easily that for any polynomial $P$ on $\mathfrak{g}$, we have

$$
\left\langle D H^{G}\left(M, \Omega, \hat{A}_{n}(M)\right), P(\partial) \tilde{\varphi}\right\rangle=\left\langle D H^{H}\left(M, \Omega, \hat{A}_{n}(M)\right),\left.P\right|_{\mathfrak{h}}(\partial) \varphi\right\rangle .
$$

Hence the formal series $\left\langle j_{\mathfrak{g}}^{1 / 2}(i \partial / k) \sum_{n=0}^{\infty} \frac{k^{d}}{k^{n}} D H^{G}\left(M, \Omega, \hat{A}_{n}(M)\right), \tilde{\varphi}\right\rangle$ is equal to

$$
\left\langle j_{\mathfrak{g}}^{1 / 2}(i \partial / k) \sum_{n=0}^{\infty} \frac{k^{d}}{k^{n}} D H^{H}\left(M, \Omega, \hat{A}_{n}(M)\right), \varphi\right\rangle \equiv\left\langle j_{\mathfrak{g} / \mathfrak{h}}^{1 / 2}(i \partial / k) \Theta_{k}^{H}, \varphi\right\rangle
$$

Consider the function $m_{H}^{\prime}(\mu, k):=\sum_{\lambda \in A_{G}} m_{G}(\lambda, k) c(\lambda, \mu)$ and the difference $d(\mu, k):=m_{H}(\mu, k)-m_{H}^{\prime}(\mu, k)$. Our aim is to show that $d=0$.

We start with the following
Corollary 5.4. We have rescaling $\mathfrak{h}^{*}(k)\left(\sum_{\mu \in A_{H}} d(\mu, k) \beta_{\mu}\right) \equiv 0$.
Proof. The previous proposition tells us that

$$
r_{*}\left(\Theta_{k}^{G}\right) \equiv B_{\mathfrak{g} / \mathfrak{h}}^{k} \star \operatorname{rescaling}_{\mathfrak{h}^{*}}(k)\left(\sum_{\mu \in A_{H}} m_{H}(\mu, k) \beta_{\mu}\right) .
$$

If we use Lemma 5.2, we can compute $r_{*}\left(\Theta_{k}^{G}\right)$ in another manner:

$$
\begin{aligned}
r_{*}\left(\Theta_{k}^{G}\right) & =\operatorname{rescaling}_{\mathfrak{h}^{*}}(k)\left(\sum_{\lambda \in A_{G}} m_{G}(\lambda, k) r_{*}\left(\beta_{\lambda}\right)\right) \\
& =B_{\mathfrak{g} / \mathfrak{h}}^{k} \star \operatorname{rescaling}_{\mathfrak{h}^{*}}(k)\left(\sum_{\mu \in A_{H}} m_{H}^{\prime}(\mu, k) \beta_{\mu}\right) .
\end{aligned}
$$

At this stage we have that $B_{\mathfrak{g} / \mathfrak{h}}^{k} \star$ rescaling $_{\mathfrak{h}^{*}}(k)\left(\sum_{\mu \in A_{H}} d(\mu, k) \beta_{\mu}\right) \equiv 0$, or equivalently $B_{\mathfrak{g} / \mathfrak{h}}^{k} \star\left(\sum_{\mu \in A_{H}} d(\mu, k) \beta_{\mu / k}\right) \equiv 0$. Since the function $d(\mu, k)$ is a piecewise quasi-polynomial function, we know that $\Theta_{k}:=$ $\sum_{\mu \in A_{H}} d(\mu, k) \beta_{\mu / k}$ admits an asymptotic expansion

$$
\Theta_{k} \equiv k^{n_{o}} \sum_{n \geq 0} k^{-n} \theta_{n}(k)
$$

where the distributions $\theta_{n}(k)$ depends periodically on $k$ (see Proposition 3.1). Thanks to Proposition 3.2, we know that $B_{\mathfrak{g} / \mathfrak{h}}^{k} \star \Theta_{k}$ admits the asymptotic expansion

$$
0 \equiv B_{\mathfrak{g} / \mathfrak{h}}^{k} \star \Theta_{k} \equiv j_{\mathfrak{g} / \mathfrak{h}}^{1 / 2}(i \partial / k)\left(k^{n_{o}} \sum_{n \geq 0} k^{-n} \theta_{n}(k)\right)
$$

If we write the Taylor series $j_{\mathfrak{g} / \mathfrak{h}}^{1 / 2}(X)=\sum_{n \geq 0} d_{n}(X)$, we obtain for any $n \geq 0$ the relation $\sum_{m+l=n} d_{m}(i \partial) \theta_{l}(k)=0$. As $d_{0}=1$, we see that $\theta_{n}(k)=0$ for all $n \geq 0$.

If we use Lemma 5.3 and Theorem 4.5, we can prove similarly the following extension of Corollary 5.4.

Proposition 5.5. For any $g \in H$ of finite order, we have

$$
\text { rescaling }_{\mathfrak{h}(g)^{*}}(k)\left(\sum_{\mu \in A_{H}} d(\mu, k) \beta(g, \mu)\right) \equiv 0
$$

The function $d(\mu, k)$ is defined for $\mu \in A_{H}=\left(\rho_{H}+\Lambda_{H}\right) \cap \mathfrak{t}_{H,>0}^{*}$. Let $\tilde{\Lambda}_{H}$ be the lattice generated by $\Lambda_{H}$ and $\rho_{H}$. Let $W_{H}$ the Weyl group of $H$. We can then extend the function $d(\mu, k)$ to a $W_{H}$-anti-invariant function $\tilde{d}$ on $\tilde{\Lambda}_{H}: \tilde{d}(w \mu, k)=\epsilon(w) d(\mu, k)$ for $w \in W_{H}, \mu \in A_{H}$ and $\tilde{d}(\mu, k)=0$ if $\mu \notin \rho_{H}+\Lambda_{H}$.

Let $T_{H}$ be a Cartan subgroup of $H$. We have a covering $\tilde{T}_{H} \rightarrow T_{H}$ such that $\tilde{\Lambda}_{H}$ is the weight lattice of torus $\tilde{T}_{H}$. Take $\tilde{g} \in \tilde{T}_{H}$ and its image $g \in T_{H}$.

We now identify invariant distribution on $\mathfrak{h}(g)^{*}$ to $W_{H(g)}$-anti-invariant distributions on $\mathfrak{t}_{H}^{*}$ through the isomorphism $\mathcal{R}_{\mathfrak{h}(g)}$. Using Lemma 3.15, we have $\mathcal{R}_{\mathfrak{h}(g)}(\beta(g, \mu))=\gamma_{\tilde{g}} \sum_{w \in W_{G}} \epsilon(w) \tilde{g}^{w \mu} \delta_{w \mu}$.

We thus obtain
Proposition 5.6. For any $\tilde{g} \in \tilde{T}_{H}$ of finite order, we have

$$
\sum_{\nu \in \tilde{\Lambda}_{H}} \tilde{d}(\nu, k) \tilde{g}^{\nu} \delta_{\nu / k} \equiv 0
$$

Since $\tilde{d}(\nu, k)$ is a piecewise quasi-polynomial function $\tilde{\Lambda}_{H} \times \mathbb{Z}$, Proposition 5.6 implies that $\tilde{d}=0$ (see Proposition 3.1 in [23]). Hence $d=0$ : the proof of Theorem 5.1 is complete.

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[^0]:    ${ }^{1}$ The term $k \frac{\operatorname{dim} M^{g}}{2} \mathcal{I}_{g}(n, k)(X)$ must be understood as the sum $\sum_{a} k^{\frac{\operatorname{dim} M_{a}^{g}}{2}} \int_{M_{a}^{g}} e^{-i \Omega_{g}(X)} V_{g}(n, k)(X)$.
    ${ }^{2}$ We can also define $\beta(g, \lambda)$ by the relation $\beta(g, \lambda)=\gamma_{g} \sum_{\bar{w}} g^{w \lambda-\rho} \beta_{G(g)^{\circ} w \lambda}$.

[^1]:    ${ }^{3}$ The term $k^{\frac{\operatorname{dim} M^{g}}{2}} D H^{G(g)}\left(M^{g}, \Omega_{g}, V_{g}(n, k)\right)$ must be understood as the sum $\sum_{a} k^{\frac{\operatorname{dim} M_{a}^{g}}{2}} D H^{G(g)}\left(M_{a}^{g}, \Omega_{g}, V_{g}(n, k)\right)$.

