

Interaction of human migration and wealth distribution

Vitaly Volpert, S Petrovskii, A Zincenko

▶ To cite this version:

Vitaly Volpert, S Petrovskii, A Zincenko. Interaction of human migration and wealth distribution. Nonlinear Analysis: Hybrid Systems, 2017, 159, pp.408 - 423. 10.1016/j.na.2017.02.024 . hal-01565519

HAL Id: hal-01565519 https://hal.science/hal-01565519

Submitted on 19 Jul 2017

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers. L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

Interaction of human migration and wealth distribution

V. Volpert^{a,b,*}, S. Petrovskii^c, A. Zincenko^c

^a Institut Camille Jordan, UMR 5208 CNRS, University Lyon 1, 69622 Villeurbanne, France

^b RUDN University, ul. Miklukho-Maklaya 6, Moscow, 117198, Russia

^c Department of Mathematics, University of Leicester, Leicester, LE1 7RH, UK

ARTICLE INFO

Article history: Received 26 September 2016 Accepted 23 February 2017 Communicated by Enzo Mitidieri

Keywords: Population density Wealth distribution Cross diffusion

ABSTRACT

Dynamics of human populations depends on various economical and social factors. Their migration is partially determined by the economical conditions and it can also influence these conditions. This work is devoted to the analysis of the interaction of human migration and wealth distribution. The model consists of a system of equations for the population density and for the wealth distribution with conventional diffusion terms and with cross diffusion terms describing human migration determined by the wealth gradient and wealth flux determined by human migration. Wealth production and consumption depend on the population density while the natality and mortality rates depend on the level of wealth. In the absence of cross diffusion terms, dynamics of solutions is described by travelling wave solutions of the corresponding reaction–diffusion systems of equations. We show persistence of such solutions for sufficiently small cross diffusion coefficients. This result is based on the perturbation methods and on the spectral properties of the linearized operators.

1. Introduction

Dynamics of human population has long been a focus of research and controversy. Early insights into its uncontrolled growth resulted in a rather negative outlook and even disastrous conclusions because of the finiteness of available resources [14] (see also Section 1.1 in [18]). Although nowadays this gloomy forecast is not widely shared, good understanding of the factors controlling the growth and proliferation of human population is required in order to avoid crises, poverty and social tensions.

Due to the effect of cultural, religious, political and other phenomena specific for human societies, e.g. see [23], the population dynamics of humans is apparently much more complicated than the dynamics of animal species. Mathematical models of human dynamics allowing for the specifics and complexity of social

^{*} Corresponding author at: Institut Camille Jordan, UMR 5208 CNRS, University Lyon 1, 69622 Villeurbanne, France. Fax: +33 472431687.

E-mail address: volpert@math.univ-lyon1.fr (V. Volpert).

and cultural interactions are at their infancy. There are, however, factors and processes that are important for all living beings, and one of them is movement. As well as animal species, the dynamics of human population has a distinct spatial aspect. In particular, human migration is a ubiquitous phenomenon that has been shaping countries and societies over centuries of human history [24,5] and that takes place at various scales ranging from individuals to nations [9,20]. The direction of migration (or, more generally, population flow) is often determined by the availability of resources and/or by the quality of living.¹ That can be thought about as the effect of the resource gradient. But the effect of resource gradient is known to be a factor that affects invasion (or, more generally, dispersal) of animal and plant species too, e.g. see [12,11]. Therefore, mathematical approaches that were developed to describe spatiotemporal dynamics of ecological populations [18,13] can be expected to be a reasonable starting point to describe the human dynamic too.

Proliferation of biological populations in space occurs due to an interplay between their reproduction and the movement of individuals [22,7,21]. The density p(x,t) of the population which depends on the space xand on time t can be calculated from a mathematical model. Nowadays, a broad variety of models can be used for this purpose [10] depending on the properties of the environment, e.g. uniform or fragmented [16], and on the type of individual movement, e.g. ordinary diffusion (Brownian motion) or anomalous diffusion [15]. One commonly used type of model is given by the reaction-diffusion equations [27]. In the case where only a single species is considered explicitly, the model consists of a single partial differential equation:

$$\frac{\partial p}{\partial t} = D \frac{\partial^2 p}{\partial x^2} + f(p), \qquad (1.1)$$

where the diffusion term corresponds to random movement of the individuals (ordinary diffusion), and the reaction term represents their natality and mortality. This equation with the function f(p) = p(1-p) was introduced in [4] in order to describe propagation of the dominant gene. In a more general ecological context, $f(p) = \alpha p(K-p)$ where α is the per capita population growth rate and K is the carrying capacity of the environment [18].

A remarkable property of Eq. (1.1) is that, if considered in an unbounded uniform space, for a broad class of initial conditions its large-time asymptotical solution is given by a travelling wave, that is by the solution p(x,t) = P(x - ct), where $c = (2D\alpha)^{1/2}$ is the speed of propagation [8]. Existence, stability and the speed of propagation of travelling waves are studied in detail for the logistic function f(p) and for more general nonlinearities (see [27,29] and the references therein).

Considering Eq. (1.1) in the context of population dynamics, it therefore predicts that the propagation of the population is described by a travelling wave with the constant speed c determined by the coefficients of the equation. The properties of the propagation becomes more complicated if other species (e.g. a predator or pathogen or competitor) or components are included (e.g. resource); the model is then given by a systems of PDEs. Well known examples are given by a system of competing species and by the prey-predator system where a simple travelling population front may turn into a sequence of fronts [19] or even a more complicated pattern of spread [27,17].

In the next section we will introduce the model of economical migration studied in this work. It consists of a system of two equations with cross diffusion terms. In the case of zero cross diffusion coefficients, it reduces to a conventional reaction–diffusion system of equations. We will study the case of small cross diffusion coefficients and will show that migration of the population can be described by travelling wave solutions.

 $^{^{1}}$ Considering, for instance, the recent surge in the migration to Europe, whilst some of the migrants were apparently fleeing war or genocide, many migrants openly stated that they were in search of economic benefits and more comfortable life.

2. Model of economical migration

In order to study economical migrations, we will complement the baseline model of biological invasions (1.1) with an equation for the resource or "wealth". Let p denote the population density and u the wealth distribution. Both of them are supposed to depend on space and time. We will consider the following system of equations:

$$\frac{\partial u}{\partial t} = D_1 \frac{\partial^2 u}{\partial x^2} + D_2 \frac{\partial}{\partial x} \left(u \frac{\partial p}{\partial x} \right) + W(u, p) - S(u, p), \tag{2.2}$$

$$\frac{\partial p}{\partial t} = D_3 \frac{\partial^2 p}{\partial x^2} - D_4 \frac{\partial}{\partial x} \left(p \frac{\partial u}{\partial x} \right) + \alpha p^n (K(u) - p) - \sigma(u) p.$$
(2.3)

Here the diffusion term in the first equation describes local wealth redistribution due to the economical activity or taxes, the second diffusion term describes wealth transport due the displacement of individual, W and S are the wealth production and consumption. These functions will be specified below. The first diffusion term in the second equation describes random motion of the population, and the second diffusion term shows the migration of the population in the direction of wealth gradient. The second term in the right-hand side of this equation describes reproduction of the population where the power n is introduced in order to attain generality and to account for collective effects. The special case of n = 1 corresponds to the standard logistic population growth [18]. The last term in Eq. (2.3) describes mortality. In a general case, carrying capacity K and mortality rate σ depend on wealth. Let us note that cross diffusion terms resemble the Soret and Dufour effects in heat and mass transport, the second diffusion term in Eq. (2.3) is similar to that in chemotaxis.

Production of wealth is often considered in the form of Cobb–Douglas production function:

$$W(L,Q,M) = bL^{\alpha}Q^{\beta}M^{\gamma}, \qquad (2.4)$$

where L_i represents the labour input, Q_i the capital, machinery, transportation, infrastructure, M_i natural resources and materials [3]; a positive coefficient *b* represents a measure of technology, α, β and γ are positive constants. The condition that doubling the inputs leads to doubling the output signifies that $\alpha + \beta + \gamma = 1$. This model satisfies the law of diminishing marginal productivity [6]. In some models, this sum is considered different from 1. Assuming that *M* is constant, capital *Q* is proportional to wealth and labour *L* is proportional to the population density, we get the wealth production rate in the form W(u, p) = f(u)g(p). We will consider the functions f(u) and g(p) in a more general form. Typically these are growing functions with saturation.

Next, we specify the rate of wealth consumption. At the level of individual consumption c, we will consider Keynes consumption function, c = r + sy, where y is income, r and s are some positive coefficients. Assuming that average income is proportional to the wealth, we will consider the rate of wealth consumption in the form S(u, p) = au + (r + su)p, where the first term corresponds to its consumption independent of the population level and the second term depending on it.

Let us now estimate the dependence of carrying capacity K(u) on wealth u. Suppose that the number of children N in the family is proportional to the family income R and inversely proportional to the child related expenses $E, N \sim R/E$. Then there are two limiting cases. In the first one, E is fixed, and R is proportional to the wealth u. In this case, K(u) is a growing function of u. In the second case, R is still proportional to u, but E also depends on u and it grows faster than R. In particular, E dependence on ucan be related to education and various commodities required for children in modern society. In this second case, K(u) is a decreasing function of u. Both patterns are known. When modern technologies come to undeveloped countries, they lead to a rapid population growth. On the other hand, it is also known that natality rate decreases with growth of the education level. Finally, the mortality rate decreases as a function of wealth, $\sigma(u) = \sigma_0 - \sigma_1 u/(1+u)$.

In this work we study the existence of travelling waves for system (2.2), (2.3) in the case of small cross diffusion coefficients. If $D_2 = D_4 = 0$, then the wave existence follows from the properties of the monotone systems [26,29]. In order to study the case where these parameters are positive and sufficiently small, we apply the perturbation technique based on the implicit function theorem. We begin with the spectral properties of the corresponding linear operators and establish the conditions under which they satisfy the Fredholm property (Section 3). It allows us to formulate the solvability conditions and prove the invertibility of linear operators on properly chosen subspaces. The invertibility is used in the implicit function theorem. In Section 4 we employ it to prove the wave existence with a variable speed in the bistable case (Theorem 4.1) and with a fixed speed in the monostable case (Theorem 4.2). These results are applied to study the migration model.

3. Spectral properties of linear operators

We consider the system (2.2), (2.3)

$$\frac{\partial u}{\partial t} = D_1 \frac{\partial^2 u}{\partial x^2} + D_2 \frac{\partial}{\partial x} \left(u \frac{\partial p}{\partial x} \right) + F(u, p), \tag{3.1}$$

$$\frac{\partial p}{\partial t} = D_3 \frac{\partial^2 p}{\partial x^2} - D_4 \frac{\partial}{\partial x} \left(p \frac{\partial u}{\partial x} \right) + G(u, p) \tag{3.2}$$

with more general functions F(u, p) and G(u, p) in view of other possible applications. We will consider travelling wave solutions of this systems, that is solutions of the form u(x, t) = U(x - ct), p(x, t) = P(x - ct), which satisfy the system of equations

$$D_1 U'' + D_2 (UP')' + cU' + F(U, P) = 0, (3.3)$$

$$D_3 P'' - D_4 (PU')' + cP' + G(U, P) = 0$$
(3.4)

considered on the whole axis. Here prime denotes the derivative with respect to the variable z = x - ct. Consider the operator A corresponding to the left-hand side of system (3.3), (3.4):

$$A\begin{pmatrix} U\\ P \end{pmatrix} = \begin{cases} D_1 U'' + D_2 (UP')' + cU' + F(U, P)\\ D_3 P'' - D_4 (PU')' + cP' + G(U, P) \end{cases}$$

and the operator L linearized about some functions U_0, P_0 :

$$L\begin{pmatrix}U\\P\end{pmatrix} = \begin{cases} a_{11}U'' + a_{12}P'' + b_{11}U' + b_{12}P' + c_{11}U + c_{12}P\\ a_{21}U'' + a_{22}P'' + b_{21}U' + b_{22}P' + c_{21}U + c_{22}P \end{cases},$$

where

$$a_{11} = D_1, \qquad a_{12} = D_2 U_0(z), \qquad a_{21} = -D_4 P_0(z), \qquad a_{22} = D_3,$$
(3.5)

$$b_{11} = c + D_2 P_0'(z), \qquad b_{12} = D_2 U_0'(z), \qquad b_{21} = -D_4 P_0'(z), \qquad b_{22} = c - D_4 U_0'(z), \tag{3.6}$$

$$c_{11} = F'_U(U_0, P_0) + D_2 P''_0, \qquad c_{12} = F'_P(U_0, P_0), c_{21} = G'_U(U_0, P_0), \qquad c_{22} = G'_P(U_0, P_0) - D_4 U''_0.$$
(3.7)

We suppose that the diffusion coefficients D_i , i = 1, 2, 3, 4 are positive, the functions F(U, P) and G(U, P)have continuous first partial derivatives with respect to U and P, and these derivatives satisfy the Lipschitz condition in every bounded set in \mathbb{R}^2 . Moreover, we assume that $U_0, P_0 \in C^{2+\alpha}(\mathbb{R})$ and $U_0(z), P_0(z) > 0$ for all $z \in \mathbb{R}$. Then the coefficients of the operator L belong to the space $C^{\alpha}(\mathbb{R})$. We consider the operator L as acting from the space $E_1 = C^{2+\alpha}(\mathbb{R})$ into the space $E_2 = C^{\alpha}(\mathbb{R})$. Let us recall that the operator L satisfies the Fredholm property if its kernel has a finite dimension, its image is close and the codimension of the image has also a finite dimension. The latter means that the nonhomogeneous equation Lw = f is solvable if and only if a finite number of solvability conditions $\phi_i(f) = 0, i = 1, ..., N$ are satisfied. Here $w = (U, P), \phi_i$ are some functionals from the dual space E_2^* . In order to study the Fredholm property of this operator, we suppose that the functions U_0 and P_0 have limits at infinity and introduce the limiting operators

$$L^{\pm} \begin{pmatrix} U \\ P \end{pmatrix} = \begin{cases} a_{11}^{\pm}U'' + a_{12}^{\pm}P'' + b_{11}^{\pm}U' + b_{12}^{\pm}P' + c_{11}^{\pm}U + c_{12}^{\pm}P \\ a_{21}^{\pm}U'' + a_{22}^{\pm}P'' + b_{21}^{\pm}U' + b_{22}^{\pm}P' + c_{21}^{\pm}U + c_{22}^{\pm}P, \end{cases}$$

where the superscripts in the coefficients signify their limiting values at $\pm \infty$.

Proposition 3.1. Suppose that the matrix $\mathcal{A}(z) = (a_{ij}(z))$ is such that det $\mathcal{A}(z) \ge a_0$ for all $z \in \mathbb{R}$ and for some positive number a_0 . Then the operator L is normally solvable with a finite dimensional kernel if and only if the limiting equations

$$L^{\pm}w = 0 \tag{3.8}$$

have only zero solutions in E_1 .

The proof of the proposition follows from more general results on elliptic operators in unbounded domains [25, Chapter 4, Theorems 3.1, 3.2]. Let us note that by virtue of the condition on the determinant of the matrix \mathcal{A} , the operator L is elliptic in the Douglis–Nirenberg sense. Taking into account the explicit form (3.5) of the coefficients, we get

$$\det \mathcal{A} = D_1 D_3 + D_2 D_4 U_0 P_0 \ge D_1 D_3 > 0.$$

Therefore the condition of the proposition is satisfied.

Since the limiting operators have constant coefficients, then we can apply the Fourier transform to equations (3.8). These equations have only zero solutions if and only if

$$\det\left(-\xi^2 A^{\pm} + i\xi B^{\pm} + C^{\pm}\right) \neq 0 \quad \forall \xi \in \mathbb{R},$$
(3.9)

where

$$A^{\pm} = \begin{pmatrix} a_{11}^{\pm} & a_{12}^{\pm} \\ a_{21}^{\pm} & a_{22}^{\pm} \end{pmatrix}, \qquad B^{\pm} = \begin{pmatrix} b_{11}^{\pm} & b_{12}^{\pm} \\ b_{21}^{\pm} & b_{22}^{\pm} \end{pmatrix}, \qquad C^{\pm} = \begin{pmatrix} c_{11}^{\pm} & c_{12}^{\pm} \\ c_{21}^{\pm} & c_{22}^{\pm} \end{pmatrix}.$$

Set $M^{\pm}(\xi) = -\xi^2 A^{\pm} + i\xi B^{\pm} + C^{\pm}$. Taking into account the explicit form of the coefficients (3.5)–(3.7), we get

$$M^{\pm}(\xi) = \begin{pmatrix} -\xi^2 D_1 + s_{11}^{\pm} & , & -\xi^2 D_2 U_0^{\pm} + s_{12}^{\pm} \\ \xi^2 D_4 P_0^{\pm} + s_{21}^{\pm} & , & -\xi^2 D_3 + s_{22}^{\pm} \end{pmatrix} + ci\xi I,$$

where I is the identity matrix and

$$s_{11}^{\pm} = F'_U(U_0^{\pm}, P_0^{\pm}), \qquad s_{12}^{\pm} = F'_P(U_0^{\pm}, P_0^{\pm}), \qquad s_{21}^{\pm} = G'_U(U_0^{\pm}, P_0^{\pm}), \qquad s_{22}^{\pm} = G'_P(U_0^{\pm}, P_0^{\pm})$$

In what follows, the superscripts \pm will be omitted for brevity. Consider first the case c = 0. Then

$$\det M(\xi) = \alpha \xi^4 + \beta \xi^2 + \gamma,$$

where

$$\alpha = D_1 D_3 + D_2 D_4 U_0 P_0, \qquad \beta = -D_1 s_{22} - D_3 s_{11} + D_2 U_0 s_{21} - D_4 P_0 s_{12}, \qquad \gamma = s_{11} s_{22} - s_{12} s_{21}.$$

Theorem 3.2. The operator L with the coefficients (3.5)-(3.7) satisfies the Fredholm property if and only if for c = 0 one of the conditions

$$\beta^2 < 4\alpha\gamma, \tag{3.10}$$

or

$$\beta^2 \ge 4\alpha\gamma \quad \text{and} \quad -\beta + \sqrt{\beta^2 - 4\alpha\gamma} < 0$$
(3.11)

is satisfied and for $c \neq 0$ conditions

$$\gamma \neq 0 \quad \text{and} \quad \det M(\xi_0) \neq 0,$$
(3.12)

where $\xi_0^2 = (s_{11} + s_{22})/(D_1 + D_3)$, are satisfied.

The proof of this theorem is given in Appendix A.

Remark 3.3. Let us recall that the essential spectrum of the operator L is the set of all complex numbers λ for which the operator $L - \lambda$ does not satisfy the Fredholm property. According to the results presented above, it is determined by the equality

$$\det(M(\xi) - \lambda I) = 0, \quad \xi \in \mathbb{R}.$$
(3.13)

If the essential spectrum of the operator L lies completely in the left-half plane, then the index of the operator equals 0. Indeed, the operator $L - \lambda$ is invertible for large positive λ [25], and the index is constant in each connected set at the complex plane where the operator satisfies the Fredholm property.

If det $M(\xi) = 0$ for some real ξ , then the essential spectrum of the operator L crosses the origin, and the operator does not satisfy the Fredholm property. This is also related to the instability of solutions by means of spatial structures emerging at infinity. Suppose that the matrix $S = (s_{ij})$ has both eigenvalues with the negative real part. Then $\gamma > 0$ and $s_{11} + s_{22} < 0$. This condition is not sufficient to provide that det $M(\xi) \neq 0$. Indeed, if the coefficient β is negative and sufficiently large in the absolute value, then det $M(\xi)$ can vanish. In the case of diagonal diffusion matrix $(D_2 = D_4 = 0)$, if one of the coefficients s_{11} and s_{22} is positive and the corresponding diffusion coefficient is sufficiently large, then det $M(\xi) = 0$ for some ξ . This is related to the Turing instability. If both coefficients s_{11} and s_{22} are negative, then the equality det $M(\xi) = 0$ can be achieved due to the cross diffusion coefficients if $s_{12} > 0$ or $s_{21} < 0$.

We also observe that conditions (3.10), (3.11) are not satisfied for $\gamma \leq 0$. Indeed, this is obvious for condition (3.10). As to condition (3.10), since $\alpha > 0$, then $-\beta + \sqrt{\beta^2 - 4\alpha\gamma} \geq 0$.

4. Small cross diffusion coefficients

4.1. Perturbations of waves

In the case of small cross diffusion coefficients we set $D_2 = \epsilon d_2$, $D_4 = \epsilon d_4$, where ϵ , d_2 and d_4 are positive numbers. System (3.3), (3.4) writes

$$D_1 U'' + \epsilon d_2 (UP')' + cU' + F(U, P) = 0, \qquad (4.1)$$

$$D_3 P'' - \epsilon d_4 (PU')' + cP' + G(U, P) = 0.$$
(4.2)

We will prove the existence of solutions of this system for sufficiently small ϵ assuming that the solution exists for $\epsilon = 0$. We will consider below applications to the migration model where the existence for $\epsilon = 0$ can be verified.

Suppose that $F(U^{\pm}, P^{\pm}) = G(U^{\pm}, P^{\pm}) = 0$ for some (U^{\pm}, P^{\pm}) and set

$$S^{\pm} = \begin{pmatrix} F'_U(U^{\pm}, P^{\pm}) & F'_P(U^{\pm}, P^{\pm}) \\ G'_U(U^{\pm}, P^{\pm}) & G'_P(U^{\pm}, P^{\pm}) \end{pmatrix}.$$

Theorem 4.1. Suppose that system (4.1), (4.2) has a solution $(U_0(z), P_0(z))$ with the limits

$$U_0(\pm\infty) = U^{\pm}, \qquad P_0(\pm\infty) = P^{\pm}$$

for $\epsilon = 0$ and some $c = c_0$, and one of the conditions (3.10), (3.11) is satisfied if $c_0 = 0$. If the matrices S^+ and S^- have all eigenvalues in the left-half plane and the zero eigenvalue of the operator

$$L_0 \begin{pmatrix} U \\ P \end{pmatrix} = \begin{cases} D_1 U'' + c_0 U' + F'_U (U^{\pm}, P^{\pm}) U + F'_P (U^{\pm}, P^{\pm}) P \\ D_3 P'' + c_0 P' + G'_U (U^{\pm}, P^{\pm}) U + G'_P (U^{\pm}, P^{\pm}) P \end{cases}$$

is simple, then for all ϵ sufficiently small system (4.1), (4.2) has a solution for some $c = c_{\epsilon}$.

Proof. We consider the operator

$$A_{\epsilon} \begin{pmatrix} U \\ P \end{pmatrix} = \begin{cases} D_1 U'' + \epsilon d_2 (UP')' + cU' + F(U, P) = 0, \\ D_3 P'' - \epsilon d_4 (PU')' + cP' + G(U, P) = 0 \end{cases}$$

corresponding to system (4.1), (4.2). It acts from the space $E_1 \times R$ into the space E_2 , where $E_1 = C^{2+\alpha}(\mathbb{R}), E_2 = C^{\alpha}(\mathbb{R})$ $(U, P \in E_1, c \in \mathbb{R})$. According to the conditions of the theorem, equation $A_0(w) = 0$ has a solution $w = w_0$. Here $w = (U, P), w_0 = (U_0, P_0)$. We will use the implicit function theorem in order to prove the existence of solutions of the equation $A_{\epsilon}(w) = 0$ for all sufficiently small $\epsilon > 0$. It can be easily verified that the operator $A'_{\epsilon}(w_0)$ is continuous with respect to (ϵ, w_0) (see, e.g., [1]). It remains to verify the invertibility of the operator L_0 .

Let us begin with the spectral properties of the operator L_0 . The essential spectrum of this operator is given by the equality (3.13) where

$$\det(M^{\pm}(\xi) - \lambda I) = \alpha^{\pm}\xi^{4} + (\beta^{\pm} - c_{0}^{2})\xi^{2} + \gamma^{\pm} + \lambda^{2} + (c_{0}i\xi - \lambda)\left(s_{11}^{\pm} + s_{22}^{\pm} - (D_{1} + D_{3})\xi^{2}\right).$$

Since the matrices S^{\pm} have eigenvalues with negative real parts, then $\gamma^{\pm} = \det S^{\pm} > 0$, $s_{11}^{\pm} + s_{22}^{\pm} < 0$. It can be directly verified that the essential spectrum $\lambda(\xi)$ of the operator L_0 determined as a solution of Eq. (3.13) lies in the left-half plane of the complex plane for all real ξ . Therefore the operator L_0 satisfies the Fredholm property and it has the zero index (see Remark 3.3).

Differentiating Eqs. (4.1), (4.2), we obtain that the operator L_0 has a zero eigenvalue with the eigenfunction (U'_0, P'_0) . Therefore it is not invertible and we cannot apply the implicit function theorem directly. We will show that it is applicable on a properly chosen subspace.

Equation $L_0 w = f$ is solvable if and only if

$$\int_{-\infty}^{\infty} f(z) w_*(z) dz = 0,$$

where w_* is the eigenfunction corresponding to the zero eigenvalue of the formally adjoint operator L^* .

Let us consider the subspace E_1^0 of the space $E_1 = C^{2+\alpha}(\mathbb{R})$ which consists of functions from E_1 for which

$$\int_{-\infty}^{\infty} f(z)w_0'(z)dz = 0,$$

where $w'_0 = (U'_0, P'_0)$. Since the eigenfunction w'_0 is exponentially decaying at infinity, then this integral is well defined. We consider the operator A_{ϵ} as acting from $E_1^0 \times R$ $(U, P \in E_1^0, c \in R)$ into $E_2 = C^{\alpha}(\mathbb{R})$. The constant c is unknown and it should be taken into account in the linearized operator:

$$A'_{\epsilon}(w_0)w = L_0w + cw'_0,$$

where $w_0 = (U_0, P_0), w = (U, P).$

From the condition that the zero eigenvalue of the operator L_0 is simple, it follows that

$$\int_{-\infty}^{\infty} w_0'(z) w_*(z) dz \neq 0.$$
(4.3)

Indeed, if this integral equals zero, then the equation $L_0w = w'_0$ has a solution w_1 . Therefore $L_0w_1 \neq 0, L^2w_1 = 0$, and the zero eigenvalue is not simple.

Hence equation $A'_{\epsilon}(w_0)w = f$ is solvable for any $f \in E_2$. Indeed, if we write it in the form

$$L_0 w = f - c w'_0$$

then we determine c from the equality

$$\int_{-\infty}^{\infty} f(z)w_*(z)dz = c \int_{-\infty}^{\infty} w_0'(z)w_*(z)dz$$

(cf. (4.3)), and the solvability condition is satisfied.

Thus, the operator $A'_{\epsilon}(w_0) : E_1^0 \times \mathbb{R} \to E_2$ is invertible. The assertion of the theorem follows from the implicit function theorem. \Box

Theorem 4.2. Suppose that system (4.1), (4.2) has a solution $(U_0(z), P_0(z))$ with the limits

$$U_0(\pm\infty) = U^{\pm}, \qquad P_0(\pm\infty) = P^{\pm}$$

for $\epsilon = 0$ and some $c = c_0 \neq 0$. Suppose, next, the matrix S^+ has one positive and one negative eigenvalue, and the matrix S^- has both eigenvalues in the left-half plane. If the dimension of the kernel of the operator L_0 equals 1 and the codimension of its image equals 0, then for all ϵ sufficiently small system (4.1), (4.2) has a solution for the same value $c = c_0$.

Proof. Similar to the proof of the previous theorem we consider the operator $A_{\epsilon} : E_1 \times \mathbb{R} \to E_2$. According to the assumption of the theorem, equation $A_0(w) = 0$ has a solution w_0 ($w = (U, P), w_0 = (U_0, P_0)$). We will use the implicit function theorem in order to prove that equation $A_{\epsilon}(w) = 0$ has a solution for all $\epsilon > 0$ sufficiently small.

From the conditions of the theorem it follows that the linearized operator $L_0 = A'_0(w_0)$ satisfies the Fredholm property, the dimension of its kernel equals 1, and its image coincides with the whole space E_2 . Let E_1^0 be the subspace of the space $E_1 = C^{2+\alpha}(\mathbb{R})$ which consists of functions from E_1 orthogonal in L^2 to the eigenfunction (U'_0, P'_0) . We consider the operator A_{ϵ} as acting from E_1^0 into $E_2 = C^{\alpha}(\mathbb{R})$. Since the operator $L_0: E_1^0 \to E_2$ is invertible, then we can apply the implicit function theorem. Therefore equation $A_{\epsilon}(w) = 0$ has a solution $w \in E_1^0$ for all $\epsilon > 0$ sufficiently small. \Box

Remark 4.3. Existence of waves is proved in the bistable case (Theorem 4.1) and in the monostable case (Theorem 4.2). In the former, the wave speed depends on ϵ , while in the latter, the wave speed is fixed. This difference is determined by the index of the linearized operators. It is 0 in the bistable case and 1 in the monostable case. We will consider the applications of these theorems in the next section.

4.2. Applications to the migration model

In order to apply the results of the previous section we begin with the case of zero cross diffusion coefficients:

$$D_1 U'' + cU' + F(U, P) = 0, (4.4)$$

$$D_3 P'' + cP' + G(U, P) = 0, (4.5)$$

where

$$F(U,P) = f(U)g(P) - au, \qquad G(U,P) = \alpha P^n(K(u) - P) - \sigma P.$$

We assume here for simplicity that a and σ are constant. More general case presented in the introduction can also be considered. We have also introduced the power n of the population density in the reproduction term. If n = 0 (and eventually $\sigma = 0$), then we neglect the reproduction and mortality and suppose that the evolution of the population is determined by the carrying capacity K(u). In the case n = 1 we obtain conventional logistic equation where the reproduction rate depends only on female density, if n = 2 then it depends on the product of male and female densities assumed to be equal to each other. In what follows we suppose that f'(U) > 0, g'(P) > 0, K'(U) > 0, that is wealth production increases with wealth and population density, carrying capacity increases as a function of wealth. Therefore

$$\frac{\partial F}{\partial P} > 0, \qquad \frac{\partial G}{\partial U} > 0,$$
(4.6)

system (4.4), (4.5) is monotone, and we can apply the results on wave existence [29].

Bistable case without reproduction. For simplicity, we begin with the case n = 0 and set $\alpha = 1, \sigma = 0$. Note that this case is apparently not entirely biologically realistic, because it assumes that the local population growth rate can be positive for P = 0 as $G(U, 0) \neq 0$ (which can as well mean that the system is not closed and there are other sources of population migration that we do not discuss here). However, this case is a convenient example to demonstrate the mathematical technique and to reveal the principle properties of the system.

The zeros of the nonlinearities can be found from the equations

$$f(U)g(P) - aU = 0,$$
 $K(U) - P = 0.$

We suppose that f(U) is an S-shape function qualitatively shown in Fig. 1. The functions g(P) and K(U) have a similar behaviour with possibly different inflection points. Set h(U) = f(U)g(K(U)). Equation h(U) = aU can have different numbers of solutions. Two possible situations are shown in Fig. 1. In the first case, there are three solutions, the points (U^{\pm}, P^{\pm}) are stable in the sense that the corresponding matrices (Theorem 4.1) have eigenvalues in the left-half plane. The intermediate point is unstable, the corresponding matrix has a positive eigenvalue. There exists a unique solution $(U_0(z), P_0(z))$ of system (4.4), (4.5). All conditions of Theorem 4.1 are satisfied, so there exists a solution for small cross diffusion coefficients.

In the second case, there are five solutions. Besides the points (U^{\pm}, P^{\pm}) , there is one more stable point $(U^*, P^* = K(U^*))$. In this case, system (4.4), (4.5) has a solution $(U_1(z), P_1(z))$ with the limits

$$U_1(-\infty) = U^-, \qquad P_1(-\infty) = P^-, \qquad U_1(+\infty) = U^*, \qquad P_1(+\infty) = P^*$$
(4.7)

and another solution $(U_2(z), P_2(z))$ with the limits

$$U_1(-\infty) = U^*, \qquad P_1(-\infty) = P^*, \qquad U_1(+\infty) = U^+, \qquad P_1(+\infty) = P^+.$$
 (4.8)

The corresponding values of the wave speed are c_1 and c_2 , respectively. If $c_1 > c_2$, then there exists a solution $(U_0(z), P_0(z))$ with the limits (U^{\pm}, P^{\pm}) . This solution describes the dynamics of the original system (2.2),



Fig. 1. Qualitative form of the functions f(U) and h(U). Solutions of the equation h(U) = U determine the stationary points and their stability.



Fig. 2. A sketch of the two wave system arising in the bistable case if $c_1 \leq c_2$ (solid line corresponds to the *P* component of the solution, dashed line to the *U* component).

(2.3) (without the cross-diffusion terms) in the large time limit. Similar solutions exist in the case of small nonzero cross diffusion coefficients. In case $c_1 \leq c_2$, such solution does not exist [28,26]. For intermediate time, the dynamics of the system can more generally be given by a system of two travelling front; see Fig. 2. Bistable case with reproduction. In the case n = 2 we will suppose for simplicity that g(P) and K(u) are linear functions, g(P) = P, K(U) = U, and $\alpha = 1$. This assumption allows us to get the analytical expressions for the zero line of the functions F and G:

$$P = \frac{aU}{f(U)}, \qquad U = P + \frac{\sigma}{P} \quad (P \neq 0) \text{ and } P = 0$$

(Fig. 3). The number of stationary points of the kinetic system of equations

$$\frac{dU}{dt} = F(U, P), \qquad \frac{dP}{dt} = G(U, P)$$



Fig. 3. Two specific examples of zero lines of the functions F and G. There are two stable points in the left figure (marked by "s") and three of them in the right figure.

depends on the parameters. Two examples are shown in Fig. 3. There are two stable stationary points in the left figure, $(U^+, P^+) = (0, 0)$ and $(U^-, P^-) > 0$. Other points are unstable. Therefore there exists a unique (up to a shift) monotonically decreasing solution of system (4.4), (4.5) with the limits (U^{\pm}, P^{\pm}) at infinity.

There are three stable stationary points in the example in the right figure. Denote the intermediate point by (U^*, P^*) . Then there are waves with the limits (4.7), (4.8). Denote their speeds by c_1 and c_2 , respectively. As above in the case n = 0, the wave with the limits (U^{\pm}, P^{\pm}) exists if and only if $c_1 > c_2$. In all these cases, the waves also exist for sufficiently small cross diffusion coefficients.

Monostable case with reproduction. In the case n = 1, the zero lines of the functions F and G are determined by the system of equations:

$$f(U)g(P) - aU = 0,$$
 $P = K(U) - \frac{\sigma}{\alpha}$ and $P = 0.$

The number and stability of the corresponding stationary points depend on the parameters and on the form of the functions f(U), g(P) and K(U). We will not analyse here all possible cases. Let us note that the point with P = 0 is unstable if $\alpha K'(U_0) > \sigma$, where U_0 is a solution of the equation F(U)g(0) = aU.

Suppose that the point (U^+, P^+) is unstable and the point (U^-, P^-) is stable in the sense that the corresponding matrix has a positive eigenvalue in the first case, and both of them are negative in the second case. Moreover, we assume that $U^- > U^+, P^- > P^+$, and there are no other stationary points in the rectangle $U^+ \leq U \leq U^-, P^+ \leq P \leq P^-$. Then monotonically decreasing solutions of system (4.4), (4.5) with the limits (U^{\pm}, P^{\pm}) exist for all values of speed $c \geq c_0$, where c_0 is some positive number.

The index of the linearized operator L_0 equals 1 [2]. Therefore, if the dimension of the kernel of the operator is 1, then the codimension of the image is 0, and we can apply Theorem 4.2. In order to determine the dimension of the kernel of the operator L_0 , let us recall that the derivative $(U'_0(z), P'_0(z))$ of the solution is the eigenfunction corresponding to the zero eigenvalue. Assuming that all solutions λ^i_{\pm} of the equations

$$\det T_{\pm}(\lambda) = 0$$

where $T_{\pm}(\lambda) = D\lambda^2 + c\lambda + S^{\pm}$, D is the diagonal matrix with the diagonal elements D_1 and D_3 , are different from each other, we get

$$\begin{pmatrix} U_0'(z) \\ P_0'(z) \end{pmatrix} \sim \sum_{i=1}^4 k_{\pm}^i q_{\pm}^i e^{\lambda_{\pm}^i z}, \quad z \to \pm \infty.$$
(4.9)

Here q_{\pm}^i are some vectors, k_{\pm}^i are some constants.

It can be shown that there are three negative values λ_{+}^{i} which determine behaviour of solutions at $+\infty$ (Appendix B). Assuming for simplicity that they are different from each other, $\lambda_{+}^{3} < \lambda_{+}^{2} < \lambda_{+}^{1} < 0$, we can verify that the waves with the speed $c > c_{0}$ decay with the maximal exponent λ_{+}^{1} . In this case, the kernel of the linearized operator has dimension 1, Theorem 4.2 is applicable, and we prove the wave existence for small cross diffusion coefficients.

If $c = c_0$, then solution $(U_0(z), P_0(z))$ of system (4.4), (4.5) decays with the exponent $\lambda_+^2 < \lambda_+^1$. In this case we can introduce an exponential weight at ∞ with some weight $\mu, \lambda_+^2 < \mu < \lambda_+^1$. The essential spectrum of the linearized operator L_0 considered in the weighted space is completely in the left-half plane, and its index equals zero. The derivative $(U'_0(z), P'_0(z))$ of the solution belongs to the weighted space. Hence the dimension of the kernel of the operator and the codimension of its image equal 1. Similar to the bistable case, the solution persists under small perturbations of the system but the speed c_0 can change.

5. Discussion

In this work we study the interaction of human population with wealth distribution. This interaction concerns the fluxes and the production terms for both variables. Conventional diffusion term in the equation for the population density describes random motion of the individuals while the second diffusion terms corresponds to the migration of human population along the gradient of wealth distribution. Similar terms in the equation for wealth distribution describe diffusive flux and wealth transport due to human migration. Wealth transfer from rich to poor regions occurs by redistribution mechanisms and by various types of economic activity (trade, investment, delocalization, and so on). Production terms depend on both variables. Natality and mortality rates in the equation for the population density depend on wealth, production and consumption of wealth depend on the population density.

In the case of zero cross diffusion coefficients we obtain a reaction–diffusion system of equations. It satisfies the conditions of monotonicity (positive off-diagonal elements in the Jacobian of production rates). Such systems satisfy the maximum and comparison principle allowing the investigation of the existence and stability of travelling waves.

In the bistable case, if there are only two stable stationary points, then there exists a unique travelling wave connecting them. It is globally asymptotically stable. If there are three stable stationary points, then there are two consecutive waves connecting them. The relation between their speeds determines the existence of the merged wave. We prove the wave existence in the case of small cross diffusion coefficients using the implicit function theorem. Its application is based on spectral properties of the linearized operators. The perturbation of the problem influences the wave speed.

In the monostable case with zero cross diffusion coefficients the waves exist for all values of the speed greater than or equal to the minimal speed. The index of the corresponding linear operator equals 1. We show that its kernel has the same dimension. Therefore the image of the operators coincides with the whole space. This property of the operator allows us to prove wave existence for a fixed wave speed. Its value does not change for small nonzero cross diffusion coefficients; see Theorem 4.2 and Corollary 7 (Appendix B).

Note that, in this paper, it was not our intention to provide an exhausting study of human migrations. Our goal is more modest, i.e. to introduce a new spatially explicit model of coupled population-wealth dynamics and to reveal some of its basic mathematical properties, thus laying the foundation for future research in this direction. Correspondingly, we have only studied in this work dynamics of solutions under some particular conditions on diffusion and production terms. We can expect new effects and more complicated dynamics for other values of parameters and production functions. In particular, cross diffusion terms can lead to the emergence of patterns (cf. Remark 3.3 in Section 3). These questions will be studied in the subsequent works.

Acknowledgements

The first author was partially supported by the Ministry of Education and Science of the Russian Federation (the agreement number 02.a03.21.0008) and by the program PICS CNRS 6583 Matbio.

Appendix A

We will present here the proof of Theorem 3.2. We begin with auxiliary lemmas.

Lemma 1. If c = 0, then the determinant of the matrix $M(\xi)$ is different from 0 for all real ξ if and only if one of the following two conditions is satisfied:

$$\beta^2 < 4\alpha\gamma,\tag{A.1}$$

or

$$\beta^2 \ge 4\alpha\gamma \quad \text{and} \quad -\beta + \sqrt{\beta^2 - 4\alpha\gamma} < 0.$$
 (A.2)

Indeed, we consider det $M(\xi)$ as a second-order polynomial with respect to ξ^2 . If condition (A.1) is satisfied, then it does not have real roots. For the second condition, these roots are negative, and the corresponding ξ are not real.

If $c \neq 0$, then

$$\det M(\xi) = \alpha \xi^4 + (\beta - c^2)\xi^2 + \gamma + ci\xi \left(s_{11} + s_{22} - (D_1 + D_3)\xi^2\right)$$

Lemma 2. If $c \neq 0$, then the determinant of the matrix $M(\xi)$ is different from 0 for all real ξ if and only if both of the following conditions are satisfied:

$$\gamma \neq 0 \quad \text{and} \quad \det M(\xi_0) \neq 0,$$
 (A.3)

where $\xi_0^2 = (s_{11} + s_{22})/(D_1 + D_3)$.

The proof of this lemma is straightforward. We note that the real part of det $M(\xi)$ can vanish with its imaginary part different from 0. Then the conditions of the lemma are satisfied and, as it follows from the next theorem, the operator can satisfy the Fredholm property but a part of its essential spectrum can be in the right-half plane resulting in the instability of solutions. We can now prove the Fredholm property of the operator L.

Proof of Theorem 3.2. It follows from the conditions of the theorem, Proposition 3.1 and Lemmas 1, 2 that the operator L is normally solvable with a finite dimensional kernel. It remains to prove that the codimension of its image is finite. Consider an operator L_{τ} such that its dependence on the parameters $\tau \in [0, 1]$ is continuous in the operator norm, it is normally solvable with a finite dimensional kernel for all $\tau, L_0 = L$ and the operator L_1 has a finite codimension of the image. Since the index of the operator L_{τ} does not change during such deformation, then the operator L_0 also has a finite codimension of the image.

Thus, in order to prove the theorem, we need to construct a continuous deformation of the operator L to some model operator L_1 in such a way that conditions of Lemma 1 or 2 are satisfied, and the operator L_1 satisfies the Fredholm property. We will reduce the operator L to the operator with zero cross diffusion coefficients ($D_2 = D_4 = 0$). In this case, the Fredholm property of such operators is known [26].

Let us begin with the case where c = 0 and condition (A.1) is satisfied. Consider the operator L_{τ} obtained from the operator L if we replace D_2, D_4 by $(1 - \tau)D_2, (1 - \tau)D_4$, respectively, D_1, D_3 by $(1 + k(1 - \tau))D_1, (1 + k(1 - \tau))D_3$, and s_{ij} by $(1 + k(1 - \tau))s_{ij}$. Here k is a positive number which will be chosen below. Then $L_0 = L$ and the operator L_1 has zero cross diffusion coefficients. Moreover,

$$\begin{aligned} \alpha_{\tau} &= (1+k(1-\tau))^2 D_1 D_3 + (1-\tau)^2 D_2 D_4 U_0 P_0, \\ \beta_{\tau} &= (1+k(1-\tau)) \left(-(D_1 s_{22} + D_3 s_{11}) + (1-\tau) (D_2 U_0 s_{21} - D_4 P_0 s_{12}) \right), \\ \gamma_{\tau} &= (1+k(1-\tau))^2 (s_{11} s_{22} - s_{12} s_{21}). \end{aligned}$$

Hence for k sufficiently large we have $\beta_{\tau}^2 < 4\alpha_{\tau}\gamma_{\tau}$. Hence the operator L_{τ} is normally solvable with a finite dimensional kernel for all $\tau \in [0, 1]$. Its index is independent of τ .

A similar construction can be done for c = 0 with condition (A.2) and for $c \neq 0$. Let us note that during the deformation condition (A.2) can change to (A.1). Then we proceed as above. In order to consider behaviour at $+\infty$ and $-\infty$ independently of each other we can introduce variable coefficients D_1 and D_3 with different limits at infinities. \Box

Appendix B

Lemma 3. Equation det $T_{-}(\lambda) = 0$ has two positive solutions λ_{-}^{1} and λ_{-}^{2} . If $c > c_{0}$, then they are different, and a positive vector q_{-}^{i} in (4.9) at $-\infty$ corresponds to the minimal of them.

Proof. If $D_1 = D_3$, then the solutions of the equation det $T_-(\lambda) = 0$ can be found explicitly from the equations

$$D_1 \lambda^2 + c\lambda + \mu_{-}^{1,2} = 0,$$

where $\mu_{-}^{1,2} < 0$ are the eigenvalues of the matrix S^{-} . In this case there are two positive and two negative values of λ . When we change the value of D_3 , these solutions cannot cross the imaginary axis since det $T_{\pm}(i\xi) \neq 0$ for any real ξ provided that the essential spectrum does not cross the origin. Hence there are two values of λ with positive real parts. By virtue of the existence of monotonically decreasing solution $(U^0(z), P^0(z))$ these solutions are real. They do not coincide because such solutions persist for a small change of c due to the assumption of the lemma.

It remains to verify that positive vector q_{-}^{i} in (4.9) corresponds to the minimal of these solutions. Since $T_{-}(\lambda)q_{-} = 0$, then $q_{-} > 0$ is the eigenvector corresponding to the zero eigenvalue of the matrix $T_{-}(\lambda)$. Hence 0 is the maximal eigenvalue of this matrix, and another one is negative. The matrix $T_{-}(0)$ has two negative eigenvalues. Therefore, $T_{-}(\lambda)$ has one negative and one zero eigenvalue for the minimal solution of the equation det $T_{-}(\lambda) = 0$, and one positive and one zero eigenvalue for the maximal solution. \Box

Lemma 4. Equation det $T_+(\lambda) = 0$ has three solutions with negative real parts. If they are real and different, then positive vector q^i_+ in (4.9) at ∞ corresponds to one of the two maximal solutions.

The proof of the lemma is similar to the previous one.

Proposition 5. Suppose that all solutions of the equations det $T_{\pm}(\lambda) = 0$ are real and different. If the positive vector q_{+}^{i} in (4.9) at ∞ corresponds to the maximal negative solution of the equation det $T_{+}(\lambda) = 0$, then the kernel of the operator L_{0} has dimension 1, and $c > c_{0}$.

Proof. The vector-function $w(z) = -(U'_0(z), P'_0(z))$ is a positive eigenfunction corresponding to the zero eigenvalue of the operator L_0 . According to Lemmas 3 and 4, we have

$$\begin{split} & w(z) \sim k_{+}^{1} q_{+}^{1} e^{\lambda_{+}^{1}} + k_{+}^{2} q_{+}^{2} e^{\lambda_{+}^{2}} + k_{+}^{3} q_{+}^{3} e^{\lambda_{+}^{3}}, \quad z \to \infty, \\ & w(z) \sim k_{-}^{1} q_{-}^{1} e^{\lambda_{-}^{1}} + k_{-}^{2} q_{-}^{2} e^{\lambda_{-}^{2}}, \quad z \to -\infty, \end{split}$$

where $k^1_\pm \neq 0, q^1_\pm > 0, 0 > \lambda^1_+ > \lambda^2_+ > \lambda^3_+, 0 < \lambda^1_- < \lambda^2_-.$

Suppose that there is another bounded solution $w_*(z)$ of the equation $L_0 v = 0$. Then

$$\begin{split} w_*(z) &\sim c_+^1 q_+^1 e^{\lambda_+^1} + c_+^2 q_+^2 e^{\lambda_+^2} + c_+^3 q_+^3 e^{\lambda_+^3}, \quad z \to \infty, \\ w_*(z) &\sim c_-^1 q_-^1 e^{\lambda_-^1} + c_+^2 q_-^2 e^{\lambda_-^2}, \quad z \to -\infty \end{split}$$

with some constants c_{\pm}^i . Without loss of generality we can assume that $c_{\pm}^1 = 0$. Otherwise we can take a linear combination of the functions $w_*(z)$ and w(z) which satisfies this condition. Therefore $w_*(x)$ decays faster at ∞ than w(z). We can also assume that the vector-function $w_*(z)$ has some positive values for at least one of its components and for some z. Indeed, otherwise we can multiply it by -1.

Consider the function $\tau w(z)$ with some positive number τ . Since the matrix S^- has negative eigenvalues, then there exists such $z = z_0$ that from the inequality $\tau w(z_0) \ge w_*(z_0)$ it follows that $\tau w(z) \ge w_*(z)$ for all $z \le z_0$ [29]. The inequalities between the vectors are understood component-wise. On the other hand, $\tau w(z) > w_*(z)$ for any $\tau > 0$ and all z sufficiently large (possibly depending on τ). Hence

$$\tau w(z) > w_*(z), \quad z \in \mathbb{R} \tag{B.1}$$

for τ large enough.

Denote by τ_0 infimum of all τ for which $\tau w(z) > w_*(z)$ for $z \ge z_0$. Then $\tau_0 > 0$ since $w_*(z)$ has positive values. Inequality (B.1) holds for such values of τ . Clearly, $\tau_0 w(z) \ge w_*(z)$ for $z \ge z_0$ and $\tau_0 w(z_1) = w_*(z_1)$ for some $z_1 \ge z_0$ and for one of the components of the vector-function. Therefore, $\tau_0 w(z) \ge w_*(z)$ for all $z \in \mathbb{R}$. Thus, we obtain a contradiction with the positiveness theorem (or maximum principle) at $z = z_1$.

We proved that w(z) is a unique up to a factor solution of equation $L_0 v = 0$. Since the index of this operator equals 1, then the image of the operator coincides with the whole space E_2 . The operator L_0 is invertible as acting from the subspace E_1^0 of functions orthogonal to w_0 into E_2 . From the implicit function theorem it follows that solution of system (4.4), (4.5) persists under a small variation of c. Hence $c > c_0$. \Box

Corollary 6. A monotonically decreasing solution of system (4.4), (4.5), which decays at ∞ with the exponent $\lambda^1_+ < 0$, (the maximal negative solution of the equation det $T_+(\lambda) = 0$) is unique up to a shift and it corresponds to a speed $c > c_0$. Solution with the minimal speed c_0 decays with the exponent $\lambda^2_+ < \lambda^1_+$, which is the second solution of this equation.

Corollary 7. System (4.1), (4.2) has solutions for all $c > c_0$ and $\epsilon > 0$ sufficiently small.

References

- N. Apreutesei, A. Ducrot, V. Volpert, Travelling waves for integro-differential equations in population dynamics, Discrete Contin. Dyn. Syst. Ser. B 11 (3) (2009) 541–561.
- [2] J.F. Collet, V. Volpert, Computation of the index of linear elliptic operators in unbounded cylinders., J. Funct. Anal. 164 (1999) 34–59.
- [3] The digital economist. http://www.digitaleconomist.org/d_macro.html.
- [4] R. Fisher, The wave of advance of advantageous genes, Ann. Eugen. 7 (1937) 355–369.
- [5] K.E. Graf, Hunter–gatherer dispersals in the mammoth-steppe: technological provisioning and land-use in the Enisei River valley, south-central Siberia, J. Archaeol. Sci. 37 (2010) 210–223.
- [6] C. Groth, Lecture notes in macroeconomics. Mimeo, 2015.
- [7] R. Hengeveld, Dynamics of Biological Invasions, Chapman and Hall, London, 1989.
- [8] A.N. Kolmogorov, I.G. Petrovsky, N.S. Piskunov, A study of the diffusion equation with increase in the amount of substance, and its application to a biological problem, Bull. Moscow Univ. Math. Mech. 1 (6) (1937) 1–26. In: Selected Works of A.N. Kolmogorov, Vol. 1, V.M. Tikhomirov (Ed.), Kluwer, London, 1991.
- [9] M. Kupiszewski, Extra-union migration: the east-west perspective, in: P.P. Rees, J.S.C. Stillwell, A. Convey, M. Kupiszewski (Eds.), Population Migration in the European Union, John Wiley and Sons, London, 1996, pp. 11–35.
- [10] M.A. Lewis, S.V. Petrovskii, J. Potts, The Mathematics Behind Biological Invasions, in: Interdisciplinary Applied Mathematics, vol. 44, Springer, 2016.
- [11] B. Li, W.F. Fagan, K.I. Meyer, Success, failure, and spreading speeds for invasions on spatial gradients, J. Math. Biol. 70 (2015) 265–287.
- [12] C.J. Lortie, J.H. Cushman, Effects of a directional abiotic gradient on plant community dynamics and invasion in a coastal dune system, J. Ecol. 95 (2007) 468–481.
- [13] H. Malchow, S.V. Petrovskii, E. Venturino, Spatiotemporal Patterns in Ecology and Epidemiology: Theory, Models, Simulations, Chapman & Hall / CRC Press, 2008.

- [14] T.R. Malthus, An essay on the Principal of Population, Penguin Books, 1970, Originally published in 1798.
- [15] V. Mendez, D. Campos, F. Bartumeus, Stochastic Foundations in Movement Ecology, in: Springer Series in Synergetics, Springer, Berlin, 2014.
- [16] D.C. Mistro, L.A.D. Rodrigues, S.V. Petrovskii, Spatiotemporal complexity of biological invasion in a space- and timediscrete predator-prey system with the strong Allee effect, Ecol. Complex. 9 (2012) 16–32.
- [17] A.Y. Morozov, S.V. Petrovskii, B.-L. Li, Spatiotemporal complexity of the patchy invasion in a predator-prey system with the Allee effect, J. Theoret. Biol. 238 (2006) 18–35.
- [18] J.D. Murray, Mathematical Biology I. An Introduction, in: Interdisciplinary Applied Mathematics, vol. 17, Springer, Berlin, 2002.
- [19] S.V. Petrovskii, K. Kawasaki, F. Takasu, N. Shigesada, Diffusive waves, dynamical stabilization and spatio-temporal chaos in a community of three competitive species, Jpn. J. Ind. Appl. Math. 18 (2001) 459–481.
- [20] A. Rogers, Introduction to Multiregional Mathematical Demography, Wiley, New York, 1975.
- [21] N. Shigesada, K. Kawasaki, Biological Invasions: Theory and Practice, Oxford University Press, Oxford, 1997.
- [22] P. Turchin, Quantitative Analysis of Movement, Sinauer, Sunderland, 1998.
- [23] P. Turchin, T.E. Currie, E.A.L. Turner, S. Gavrilets, War, space, and the evolution of Old World complex societies, Proc. Natl. Acad. Sci. 110 (2013) 16384–16389.
- [24] C.S.M. Turney, H. Brown, Catastrophic early Holocene sea level rise, human migration and the Neolithic transition in Europe, Quat. Sci. Rev. 26 (2007) 2036–2041.
- [25] V. Volpert, Elliptic Partial Differential Equations. Volume 1. Fredholm Theory of Elliptic Problems in Unbounded Domains, Birkhäuser, 2011.
- [26] V. Volpert, Elliptic Partial Differential Equations. Volume 2. Reaction-diffusion Equations, Birkhäuser, 2014.
- [27] V. Volpert, S. Petrovskii, Reaction-diffusion waves in biology, Phys. Life Rev. 6 (2009) 267–310.
- [28] V.A. Volpert, A.I. Volpert, Wave trains described by monotone parabolic systems, Nonlinear World 3 (1996) 159–181.
- [29] A. Volpert, Vit. Volpert, Vl. Volpert, Traveling Wave Solutions of Parabolic Systems, in: Translation of Mathematical Monographs, vol. 140, Amer. Math. Society, Providence, 1994.