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A fictitious domain method for unilateral contact problems in non-destructive testing.

E. Bécache^{†¶} P. Joly^{‡¶} G. Scarella^{§¶}

Abstract

In this work, we present a numerical method for solving the diffraction of transient elastic waves by cracks of arbitrary shapes in complex media, with Signorini's boundary conditions on the crack. We use a fictitious domain method based on a mixed displacement-stress formulation for elastodynamics. We propose an off-centered time discretisation scheme for enforcing the stability.

Keywords: elastodynamics, unilateral contact, fictitious domain method, non-destructive testing, cracks

1. Introduction

In this paper, we are interested in solving the diffraction of transient elastic waves by cracks of arbitrary shapes in complex media, with Signorini's boundary conditions on the crack. This is the continuation of a previous work [1] done on the linear problem, that is when the boundary condition on the crack is a free surface boundary condition. To get an efficient method, we want to use regular meshes and at the same time respect the geometry of the crack. This is possible thanks to the fictitious domain method, which takes into account the boundary condition via a Lagrange multiplier defined on the crack, which can

[†](eliane.becache@inria.fr)

[‡](patrick.joly@inria.fr)

[§](gilles.scarella@inria.fr)

[¶]INRIA, Domaine de Voluceau-Rocquencourt, BP 105, F-78153 Le Chesnay Cédex

be interpreted as the jump of the displacement through the crack. This allows to work with a uniform mesh in the whole domain and an independent mesh on the crack. In order to consider the unilateral contact boundary condition as a constraint, we are led to use the mixed displacement-stress formulation for elastodynamics. We will present a fictitious domain formulation of this problem in which the boundary conditions are taken into account by a variational inequality for the Lagrange multiplier. For the space discretisation of this problem, we propose to use the mixed finite element using spaces of symmetric tensors for the stress[1]. This choice was shown to allow the obtention of an explicit time discretization scheme (mass-lumping) in the linear case. In the non-linear case, we cannot use a centered difference scheme for the time discretisation which would lead to an unconditionally unstable scheme. That is why we propose an off-centered scheme that we show to be stable. This scheme is explicit in the volume unknowns (displacement and stress) but implicit in the Lagrange multiplier: one has to solve an optimisation problem with bound constraints at each time step.

2. Presentation of the dynamic unilateral contact problem

We want to solve the displacement-stress formulation of elastodynamics in a domain $\Omega = C \setminus \Gamma \in \mathbb{R}^2$, where C is a rectangle (see Fig. 1):

$$\begin{cases} \rho \frac{\partial^2 u}{\partial t^2} - \operatorname{div} \sigma &= f & \text{in } \Omega \times]0, T[, \\ A\sigma &= \varepsilon(u) & \text{in } \Omega \times]0, T[, \end{cases} \quad (1)$$

where u is the displacement field, σ the stress tensor, and $\varepsilon(u)$ the strain tensor defined as $\varepsilon_{ij}(u) = (\partial_i u_j + \partial_j u_i)/2$. We add to (1) Dirichlet boundary conditions on the exterior boundary: $u = 0$ on $\partial C \times]0, T[$, Signorini's boundary conditions without friction on

the crack Γ , [3]:

$$\left\{ \begin{array}{ll} [u_N]_\Gamma & \geq 0 \text{ on } \Gamma, \\ \sigma_N & \leq 0 \text{ on } \Gamma, \\ \sigma_N [u_N]_\Gamma & = 0 \text{ on } \Gamma, \\ [\sigma n]_\Gamma & = 0 \text{ on } \Gamma, \\ \sigma_T & = 0 \text{ on } \Gamma. \end{array} \right. \quad (2)$$

where, if n denotes the unit normal to Γ , we set: $u_N = u \cdot n$, $\sigma_N = \sigma n \cdot n$, $\sigma_T = \sigma n - \sigma_N n$ and prescribed initial datas that we will systematically omitt.

3. A fictitious domain formulation

The fictitious domain method consists in extending the two unknowns (u, σ) in the whole domain C and in introducing Lagrange multipliers to take into account the Signorini's boundary conditions (2). Setting $\mathcal{G} = H_{00}^{1/2}(\Gamma)$, we introduce

$$\begin{aligned} M &= \{u \in [L^2(C)]^2\}, \\ X &= \{\tau \in [L^2(C)]^4, \operatorname{div} \tau \in [L^2(C)]^2 / \tau_{ij} = \tau_{ji}\}, \\ L_N &= H_{00+}^{1/2}(\Gamma) = \{\mu_N \in \mathcal{G} / \mu_N \geq 0 \text{ a.e. on } \Gamma\}, \\ L_T &= [H_{00T}^{1/2}(\Gamma)]^2 = \{\mu_T \in \mathcal{G}^2 / \mu_T \cdot n = 0\} \end{aligned}$$

The fictitious domain formulation consists in finding $(\sigma, u, \lambda_N, \lambda_T) :]0, T[\rightarrow X \times M \times L_N \times L_T$

$$\left\{ \begin{array}{ll} a(\sigma, \tau) + d(\tau, u) + b_T(\tau, \lambda_T) + b_N(\tau, \lambda_N) & = 0 \quad \forall \tau \in X, \\ (\rho \frac{\partial^2 u}{\partial t^2}, v) - d(\sigma, v) & = (f, v) \quad \forall v \in M, \\ b_T(\sigma, \mu_T) & = 0 \quad \forall \mu_T \in L_T, \\ b_N(\sigma, \mu_N - \lambda_N) & \leq 0 \quad \forall \mu_N \in L_N. \end{array} \right. \quad (3)$$

with

$$\begin{aligned} a(\sigma, \tau) &= \int_C A\sigma : \tau \, dx, \quad d(\tau, w) = \int_C w \cdot \operatorname{div} \tau \, dx, \\ b_T(\tau, \mu_T) &= \langle \tau_T, \mu_T \rangle_{L'_T, L_T}, \\ b_N(\tau, \mu_N) &= \langle \tau_N, \mu_N \rangle_{\mathcal{G}', \mathcal{G}}. \end{aligned}$$

The Lagrange multipliers can be interpreted as $\lambda_N = [u_N]$ and $\lambda_T = [u_T]$, with $u_T = u - u_N n$.

4. Discretisation

4.1. Semi-discretisation in space

We now introduce some finite dimensional spaces $X_h \subset X$, $M_h \subset M$, $\mathcal{G}_H \subset \mathcal{G}$ and $(L_T)_H \subset L_T$ of dimensions respectively $N_X, N_M, N_{\mathcal{G}}, N_L^T$. The semi-discretisation in space can then be written in a matricial form as: finding $(U, \Sigma, \Lambda_N, \Lambda_T) \in \mathbb{R}^{N_M} \times \mathbb{R}^{N_X} \times \mathbb{R}_+^{N_{\mathcal{G}}} \times \mathbb{R}^{N_L^T}$ such that

$$\begin{cases} M_\sigma \Sigma + D^* U + B_T^* \Lambda_T + B_N^* \Lambda_N &= 0, & (i) \\ M_v \frac{d^2 U}{dt^2} - D \Sigma &= F, & (ii) \\ B_T \Sigma &= 0, & (iii) \\ (B_N \Sigma, \mu_N - \Lambda_N) &\leq 0, \quad \forall \mu_N \in \mathbb{R}_+^{N_{\mathcal{G}}}. & (iv) \end{cases} \quad (4)$$

In practice, and this is the interesting point in the fictitious domain method, we introduce two meshes: the volumic unknowns U and Σ are defined on a regular grid, \mathcal{T}_h made of squares K_j of size h while the surfacic unknowns Λ_N and Λ_T are computed on a nonuniform mesh on Γ , \mathcal{T}_H made of segments S_j of size H_j , $H = \sup_j H_j$, see Fig.2.

Remark 1 *The inequality (4)-(iv) can be reinterpreted as*

$$\Lambda_N = \Pi(\Lambda_N + B_N \Sigma) \quad (5)$$

where Π is the orthogonal projection on $\mathbb{R}_+^{N_{\mathcal{G}}}$.

Choice of the finite elements. We intend to use the same discretisation than for the linear problem (see [1]). For the lowest order element, this choice corresponds to :

$$\begin{aligned} X_h &= \{\sigma_h \in X / \forall K \in \mathcal{T}_h, \sigma_h|_K \in (Q_1(K))^4\} \\ M_h &= \{v_h \in M / \forall K \in \mathcal{T}_h, v_h|_K \in (Q_0(K))^2\} \\ \mathcal{G}_H &= \{\mu_H \in \mathcal{G} / \forall S \in \mathcal{T}_H, \mu_H|_S \in P_1(S)\} \end{aligned}$$

Its main interest is that it leads to block diagonal mass matrices (even diagonal for M_v) so that M_v and M_σ are very easy to invert.

4.2. The fully discretised scheme.

It would be tempting to discretize (4) using centered finite difference operators, for instance:

$$\frac{\Lambda_N^{n+1} + \Lambda_N^{n-1}}{2} = \Pi(\Lambda_N^n + B_N \Sigma^n)$$

which would give an explicit scheme. However, one can show that this choice leads to an unconditionally unstable scheme! That is why we propose the following off-centered scheme:

$$\left\{ \begin{array}{ll} M_\sigma \Sigma^n + D^* U^n + B_T^* \Lambda_T^n + B_N^* \Lambda_N^n &= 0, \quad (i) \\ M_v \frac{U^{n+1} - 2U^n + U^{n-1}}{\Delta t^2} - D \Sigma^n &= F^n, \quad (ii) \\ B_T \Sigma^n &= 0, \quad (iii) \\ \Lambda_N^{n+1} = \Pi(\Lambda_N^{n+1} + B_N \frac{\Sigma^n + \Sigma^{n+1}}{2}) & \quad (iv) \end{array} \right. \quad (6)$$

Note that if the mass matrices are block diagonal, this scheme is only implicit in Λ_N and is explicit in the other unknowns. We can show a stability result:

Theorem 1 *The scheme (6) is stable under the usual CFL stability condition*

$$\frac{\Delta t^2}{4} \|D^* D\| \leq 1, \text{ with } \|D^* D\| = \sup_{\Sigma} \frac{(D\Sigma, D\Sigma)}{(M_\sigma \Sigma, \Sigma)} \quad (7)$$

For proving this result we show the decay of an energy. The precise result is the following.

We set

$$\begin{aligned} V^{n+1/2} &= \frac{U^{n+1} - U^n}{\Delta t}, \\ E_1^{n+1} &= \frac{1}{2} \left((M_\sigma \Sigma^{n+1}, \Sigma^{n+1}) + (M_v V^{n+3/2}, V^{n+1/2}) \right). \end{aligned} \tag{8}$$

This quantity is an energy under the CFL condition (7), and one has the identity

$$E_1^{n+1} - E_1^n = \Delta t \left(\frac{\Lambda_N^n - \Lambda_N^{n+1}}{\Delta t}, B_N \frac{\Sigma^{n+1} + \Sigma^n}{2} \right),$$

which shows that $E_1^{n+1} \leq E_1^n$ thanks to 6-(iv).

5. Numerical aspects

The implementation of the method amounts to combine an explicit scheme for the unknowns U and Σ with an optimisation problem (quadratic functional with bounds constraints) to be solved at each time step for the unknown Λ . The algorithm has been tested in 1D (comparison with analytic solutions). We are currently developping a 2D code, the optimisation procedure being handled by an algorithm combining the active set method with gradient projection method, [2]. Numerical results will be presented at the conference.

References

- [1] E. Bécache, P. Joly, and C. Tsogka. Fictitious domains, mixed finite elements and perfectly matched layers for 2d elastic wave propagation. *to appear in J. Comp. Acous.* (Tech. Report INRIA 3889, 2000).
- [2] J. Nocedal and S. J. Wright. *Numerical Optimization*. Springer, 1999.
- [3] J.R. Willis and V.P. Smyshlyaev. Effective relations for nonlinear dynamics of cracked solids. *J.Mech.Phys.Solids*, 44(1):49–75, 1996.

Figures

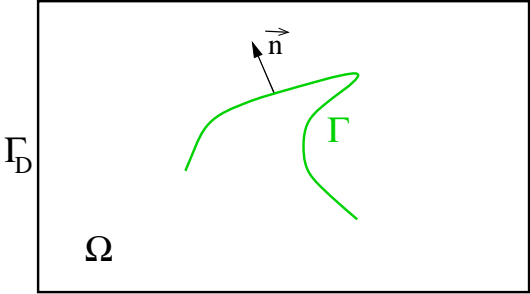


Figure 1: The geometry of the problem

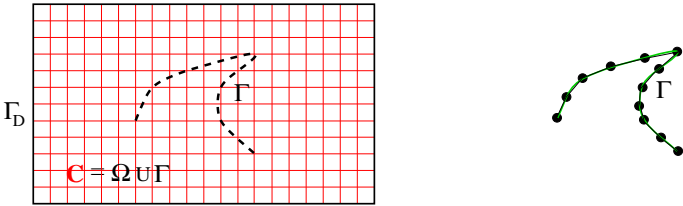


Figure 2: The two meshes