



HAL
open science

Transverse piezoelectric coefficient measurement of flexible lead zirconate titanate thin films

Thibault Dufay, Benoit Guiffard, Jean-Christophe Thomas, Raynald Seveno

► To cite this version:

Thibault Dufay, Benoit Guiffard, Jean-Christophe Thomas, Raynald Seveno. Transverse piezoelectric coefficient measurement of flexible lead zirconate titanate thin films. *Journal of Applied Physics*, 2015, 117 (20), pp.204101. 10.1063/1.4921588 . hal-01512246

HAL Id: hal-01512246

<https://hal.science/hal-01512246v1>

Submitted on 22 Jan 2020

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.



Distributed under a Creative Commons Attribution 4.0 International License

Transverse piezoelectric coefficient measurement of flexible lead zirconate titanate thin films

T. Dufay,¹ B. Guiffard,¹ J.-C. Thomas,² and R. Seveno¹

¹LUNAM Université, Université de Nantes, IETR (Institut d'Électronique et de Télécommunications de Rennes), UMR CNRS 6164, 2 rue de la Houssinière, BP 92208, 44322 Nantes Cedex 3, France

²LUNAM Université, Université de Nantes-École Centrale Nantes, GeM (Institut de Recherche en Génie Civil et Ingénierie Mécanique), UMR CNRS 6183, 2 rue de la Houssinière, BP 92208, 44322 Nantes Cedex 3, France

Highly flexible lead zirconate titanate, $\text{Pb}(\text{Zr,Ti})\text{O}_3$ (PZT), thin films have been realized by modified sol-gel process. The transverse piezoelectric coefficient d_{31} was determined from the tip displacement of bending-mode actuators made of PZT cantilever deposited onto bare or RuO_2 coated aluminium substrate ($16\ \mu\text{m}$ thick). The influence of the thickness of ruthenium dioxide RuO_2 and PZT layers was investigated for $\text{Pb}(\text{Zr}_{0.57}\text{Ti}_{0.43})\text{O}_3$. The modification of Zr/Ti ratio from 40/60 to 60/40 was done for $3\ \mu\text{m}$ thick PZT thin films onto aluminium (Al) and Al/ RuO_2 substrates. A laser vibrometer was used to measure the beam displacement under controlled electric field. The experimental results were fitted in order to find the piezoelectric coefficient. Very large tip deflections of about 1 mm under low voltage ($\sim 8\ \text{V}$) were measured for every cantilevers at the resonance frequency ($\sim 180\ \text{Hz}$). For a given Zr/Ti ratio of 58/42, it was found that the addition of a 40 nm thick RuO_2 interfacial layer between the aluminium substrate and the PZT layer induces a remarkable increase of the d_{31} coefficient by a factor of 2.7, thus corresponding to a maximal d_{31} value of 33 pC/N. These results make the recently developed PZT/Al thin films very attractive for both low frequency bending mode actuating applications and vibrating energy harvesting.

I. INTRODUCTION

Energy harvesting from ambient and low frequency mechanical vibration is actively studied because of source accessibility. Walking,¹ dancing,² or using the vibration produced by a bus³ are some examples of energy harvesting projects based on piezoelectricity.

Recent developments in low-power electronics may enable micropower energy harvesting to eliminate wiring, in order to reduce maintenance costs, or to replace electrochemical batteries.⁴ Piezoelectric materials convert efficiently the mechanical strain to an electrical charge and could generate enough power to charge batteries of some sensors (e.g., fatigue sensors).⁵ This is the reason why piezoelectric materials are widely studied for vibrating energy harvesting.⁶⁻⁹

Concerning piezoelectric thin films, several materials such as $\text{Pb}(\text{Zr,Ti})\text{O}_3$ (PZT), ZnO , and AlN are studied for actuating or sensing issues.¹⁰⁻¹² Due to their superior ferroelectric and piezoelectric properties, PZT is the most attractive candidate.¹³ However, depending on the used substrate and the composition of the PZT thin film, ferroelectric and piezoelectric properties evolve. This is why the investigation of piezoelectric coefficient, and especially the transverse piezoelectric coefficient, is an important step for the characterization of a new material. Numerous methods for measuring the transverse piezoelectric coefficient have already been reported.¹⁴⁻¹⁶ Measurement of d_{31} piezoelectric coefficient for PZT thin film as a function of the Zr/Ti ratio is interesting to determine phase transition effects on film properties. For

bulk PZT, it is well known that maximum ferroelectric and piezoelectric properties are obtained at the morphotropic phase boundary (MPB) composition.¹⁷ The MPB corresponds to the coexistence of the rhombohedral and the tetragonal phases of the PZT. For bulk PZT, even if it depends on many factors such as elaboration process, the MPB generally corresponds to Zr/Ti ratios close to 52/48.¹⁸ But, the properties of thin films show different trends than those of bulk ceramics, due to substrate, interfacial layer, and film thickness influences.^{19,20} So, the investigation of ferroelectric and piezoelectric properties as a function of Zr/Ti ratio for a new PZT thin film is a major improvement in understanding this material. In this paper, we report on the ferroelectric and piezoelectric characterization of newly developed PZT thin film ($\sim 3\ \mu\text{m}$ thick) deposited by modified sol-gel process onto commercial aluminium foil with $16\ \mu\text{m}$ in thickness.²¹ Those flexible thin films are cheap (substrate cost $< 0.1\ \$\ \text{m}^{-2}$) and could be easily used on different supports like vehicle in movement, aeration sheath, windmill, or any other structure exposed to wind, as aeroelastic energy harvester.

The study mainly focuses on the influence of both an intermediate ruthenium dioxide between PZT layer and aluminium substrate and the Zr/Ti ratio on the intrinsic properties of the thin film. More specifically, ferroelectric properties (coercive field and remnant polarization) and transverse d_{31} piezoelectric coefficients are determined for PZT/Al thin films with/without RuO_2 and as a function of the interfacial layer thickness, using bending mode cantilever configuration.

II. EXPERIMENTAL

PZT thin films were elaborated by a modified sol-gel process commonly used in the laboratory. Lead acetate $[\text{Pb}(\text{CH}_3\text{CO}_2)_2, 3\text{H}_2\text{O}]$, zirconium n-propoxide $\text{Zr}(\text{C}_3\text{H}_7\text{O})_4$, and titanium n-propoxide $(\text{Ti}(\text{CH}_3)_2\text{CHO})_4$ are used as precursor materials. Acetic acid is used as solvent and ethylene glycol as an additive to prevent cracking.

Lead acetate is dissolved in heated acetic acid. Zirconium and titanium propoxides are mixed and then added to the lead acetate solution. The precursor solution is ready after the addition of the ethylene glycol. More technical details of the preparation of the precursor solution are reported elsewhere.²² The deposition process of the PZT thin film is given in Fig. 1. Multiple spin-coating technique at 6000 rpm was performed to obtain a film thickness between 1.8 and 4.2 μm . Each single layer successively deposited is crystallized at 650 °C during 2 min and is about 300 nm thick. The Zr/Ti ratio is modified with the proportion of precursor solutions. In both cases, i.e., bare and RuO_2 coated aluminium, the PZT composition changes from Zr/Ti = 40/60 to Zr/Ti = 60/40. Ruthenium dioxide films were prepared by mixing and stirring a ruthenium (III) nitrosyl nitrate $[\text{Ru}(\text{NO})(\text{NO}_3)_3]$ aqueous solution and 2-methoxyethanol $[\text{CH}_3\text{O}(\text{CH}_2)_2\text{OH}]$. The final solution is then spin-coated on aluminium foil and fired at 450 °C in an open air furnace.²³ This RuO_2 layer was used in a previous work to improve the dielectric properties of the thin films with stainless steel

substrate.²³ Large-area (1 cm \times 3 mm) aluminium top electrodes with 300 nm in thickness were evaporated through a shadow mask on each sample. The substrate was directly used as the bottom electrode. Top view and schematic front view of thin films are given in Fig. 2.

Scanning electron microscopy (SEM) observations were carried out to check the thickness and structure of the PZT layer. Carl Zeiss Merlin SEM was used with electron energy set to 3 keV.

Samples must be poled to exhibit macroscopic piezoelectric properties. A Sawyer-Tower circuit was employed for both poling and investigation of polarization P versus electric field E . An electric field frequency of 50 Hz was used. In order to protect samples from dielectric breakdown, the magnitude of the applied electric field was slowly increased.

Dielectric properties were investigated with an impedance meter Agilent 4294 A. Loss factor $\tan\delta$ and relative dielectric permittivity ϵ_r were measured at 1 kHz.

To determine the transverse piezoelectric coefficient d_{31} , cantilevers were clamped at one end. Dimensions of the free-part were measured with a precision of 0.5 mm. Cantilever tip displacements were evaluated with a laser vibrometer OFV-2200 (Polytec) under an applied sinusoidal voltage of magnitude V_0 at different frequencies. A mathematical model is used to fit the experimental data and extract the quasistatic displacement d_0 , the natural vibration

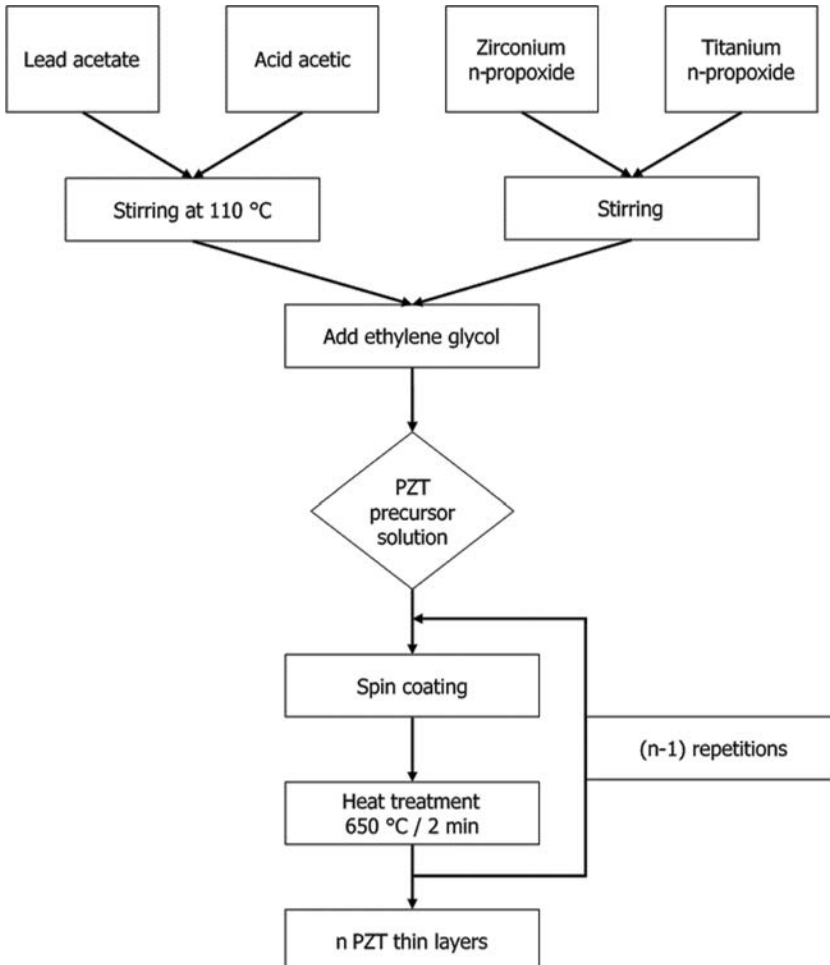


FIG. 1. Deposition process of the PZT thin film.



FIG. 2. Top view (a) and schematic front view (b) of thin films.

frequency f_0 , and the quality factor Q . Linear fit of d_0 versus applied voltage is exploited to find the d_{31} .

III. RESULTS AND DISCUSSION

SEM images of PZT 40/60 are shown in Fig. 3. PZT layer thickness is measured on the front view. A granular structure is seen and grains size is comprised between 100 and 300 nm as it could be observed on surface view. SEM study was done for three Zr/Ti ratios: 40/60, 50/50, and 60/40 and for each ratio, two samples were observed. It appeared that Zr/Ti ratio has no influence on the morphological structure.

A. Influence of RuO₂ layer thickness on the electrical properties of PZT thin films

Measurement of electrical properties was done for PZT thin films with three different ruthenium dioxide (RuO₂) layer thicknesses: 40, 80, and 120 nm. Dielectric and ferroelectric properties were investigated to evaluate the influence of RuO₂ thickness. For this study, the Zr/Ti ratio was fixed to 57/43, which is a commonly used ratio at the laboratory and in the vicinity of the MPB. The study of the influence of Zr/Ti on the electrical/piezoelectric properties will be presented in Sec. III C.

The P - E hysteresis loops, remnant polarization (P_r), coercive field (E_c), dielectric permittivity (ϵ_r), and loss factor ($\tan\delta$) of 3 μm -thick PZT 57/43 thin films with different thicknesses of RuO₂ are given in Fig. 4. At an applied electric field of 167 kV/cm, samples with a deliberately added intermediate oxide layer show a reasonably square well-saturated hysteresis loop, whereas without RuO₂, the hysteresis loop is not saturated. Here, the choice was made to apply

the same electric field to the samples, even if the hysteresis loop without RuO₂ is not saturated.

It has been shown that preparing PZT on base metals resulted in the formation of interfacial layers, yielding the degradation of ferroelectric properties.²⁴ In particular, aluminium oxide formation could occur between the substrate and the first PZT layer during the heat treatment, leading to oxygen vacancies in the first deposited PZT layers. That could explain the asymmetric loop—observed during poling of samples without RuO₂—with the trapping of electrons when positive voltage is applied.^{25,26} Besides, this is also at the origin of the relative high value of coercive field. When RuO₂ layer is used, it protects PZT layer from substrate oxidation during crystallization and limits the defects (e.g., oxygen vacancies) at the interface between the substrate and the PZT. Without the RuO₂ layer, P_r is weak ($\sim 8 \mu\text{C}/\text{cm}^2$) and E_c is high ($\sim 87 \text{ kV}/\text{cm}$), for the same applied electric field (167 kV/cm) at which the hysteresis loop is not saturated. As seen in Fig. 4(b), the remnant polarization increases up to $\sim 25 \mu\text{C}/\text{cm}^2$ when 40 nm thick RuO₂ layer is added and P_r remains almost constant for higher RuO₂ thicknesses. The coercive field is also strongly affected since it is divided by almost 2 with the presence of a 120 nm thick ruthenium oxide layer. Thus, for these two parameters, RuO₂ has a very interesting beneficial effect.

RuO₂ influence is also observed on the relative dielectric permittivity, as shown in Fig. 4(c). The ϵ_r value is ~ 310 for PZT on aluminium and increases up to about 550 for the three PZT thin films on RuO₂-coated aluminium. These quite low values come from the annealing time and temperature, respectively, shorter and weaker than those commonly used in the literature.^{7,27,28} Higher annealing time or temperature could improve crystallization and so ferroelectric properties, but it would degrade the aluminium foil substrate owing to

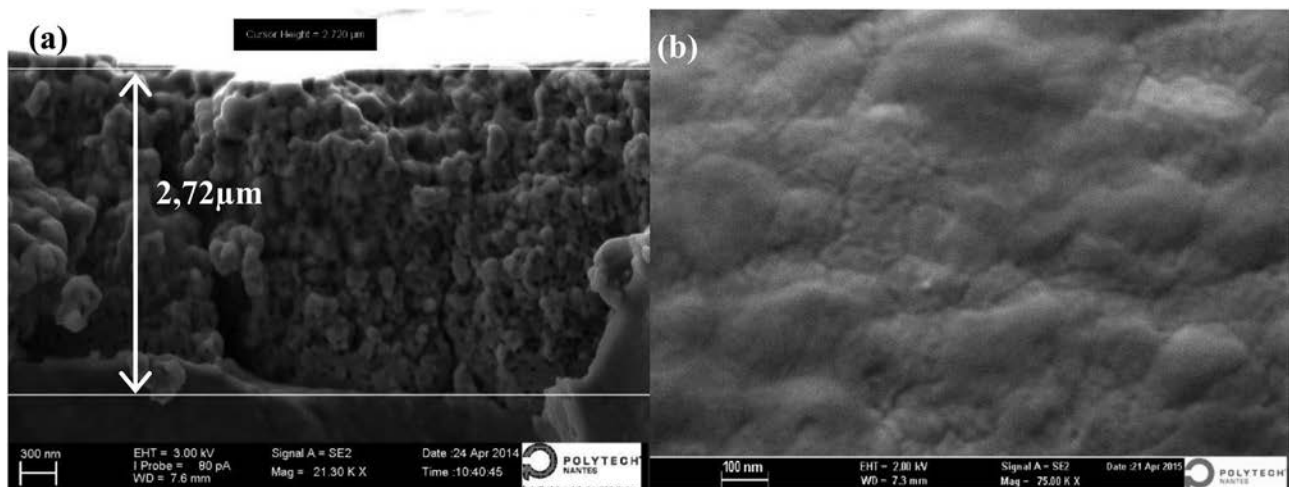


FIG. 3. SEM images of PZT 40/60 deposited onto aluminium substrate. (a) Cross-section and (b) top surface.

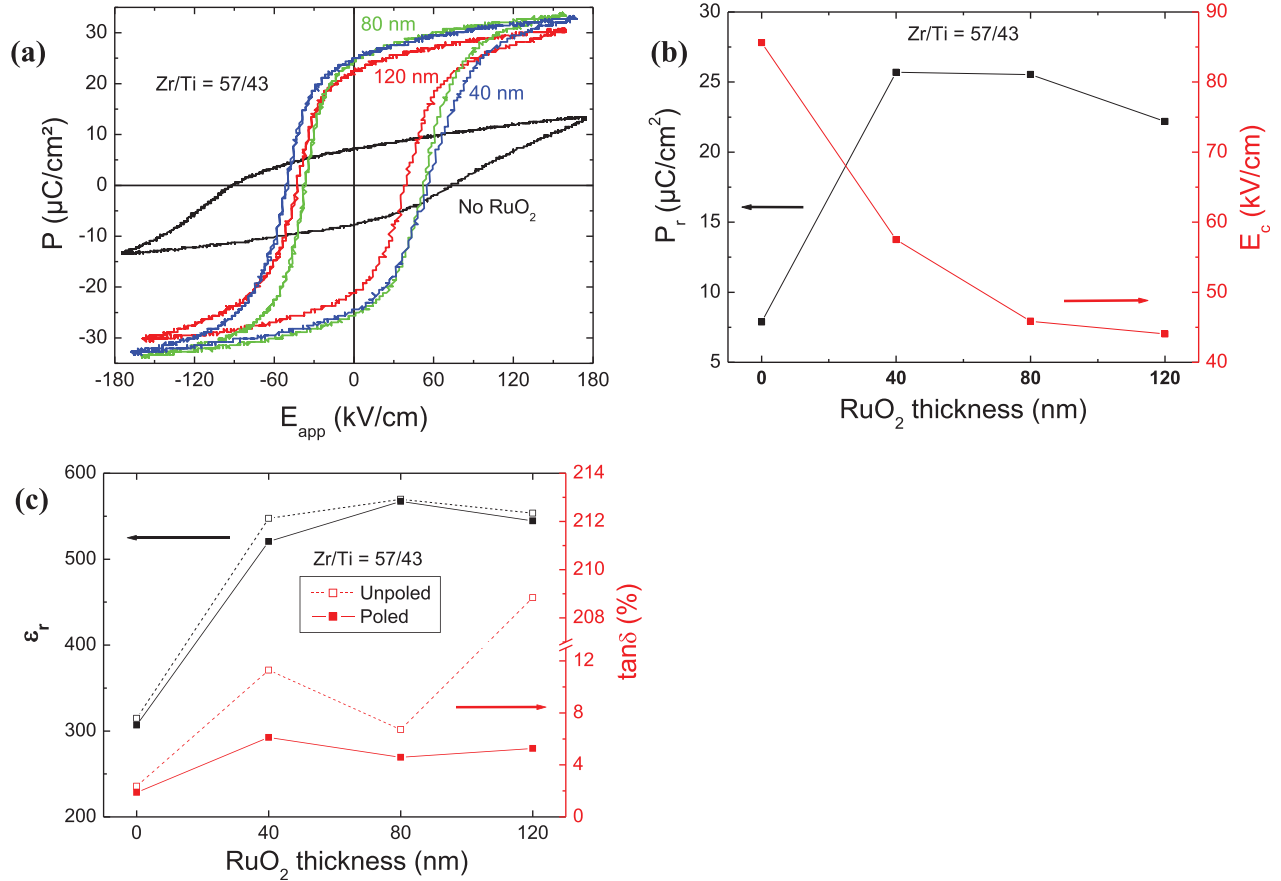


FIG. 4. Ferroelectric and dielectric properties of a 3 μm -thick PZT 57/43 thin film as a function of RuO_2 thickness: (a) hysteresis loops, (b) remnant polarization P_r and coercive field E_c , and (c) relative dielectric permittivity ϵ_r and loss factor $\tan\delta$: comparison between poled and unpoled material.

its low thickness (16 μm) and melting point ($\sim 660^\circ\text{C}$). In addition, a more crystallized thin film will lose flexibility and will favour the appearance of cracks. The loss factor value remains roughly the same after the poling step—around 5%—with and without the oxide layer. During the deposition of PZT, RuO_2 may diffuse and create some micrometer-sized diffusion channels through the PZT layer acting as electrical conduction paths, which could explain the $\tan\delta$ increase observed for PZT on 120 nm thick RuO_2 layer. Consequently, the drastic $\tan\delta$ decrease after the poling step could be explained by a self-healing effect at the top electrode.²⁹

This effect seems to be confirmed by the optical microscopy photographs of the top electrode before and after poling (Fig. 5). In fact, those micro-channels are sealed by the local

melting of the top surface of the thin film due to electrical breakdown when the electric field is applied during the poling process, as shown in Fig. 5(b). Consequently, the dielectric loss factor is decreased thanks to the self-healing of the electrode.

Finally, this study shows that it is not necessary to coat a 120 nm thick RuO_2 layer, a 40 nm thick layer is enough to enhance the dielectric and ferroelectric properties. In the following experiments, a RuO_2 layer of 40 nm will be used.

In order to estimate the piezoelectric coefficient, a laser vibrometer was used for measuring the displacement of the free-end of the beam under an applied voltage. The experimental plots of the displacement as a function of frequency and their fits are given in Fig. 6(a) for three different applied voltages. The experimental data are fitted with Eq. (1)³⁰

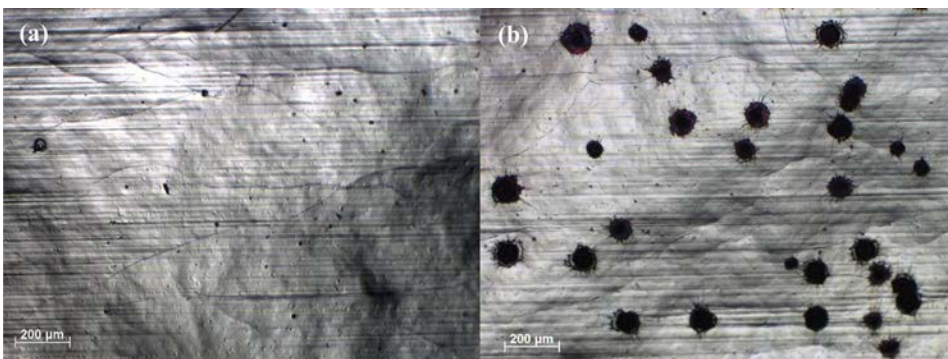


FIG. 5. Optical microscopy photographs of Al top electrode (a) before poling and (b) after poling. Black holes on the electrode after poling correspond to the local melting of the electrode near the conduction channel.

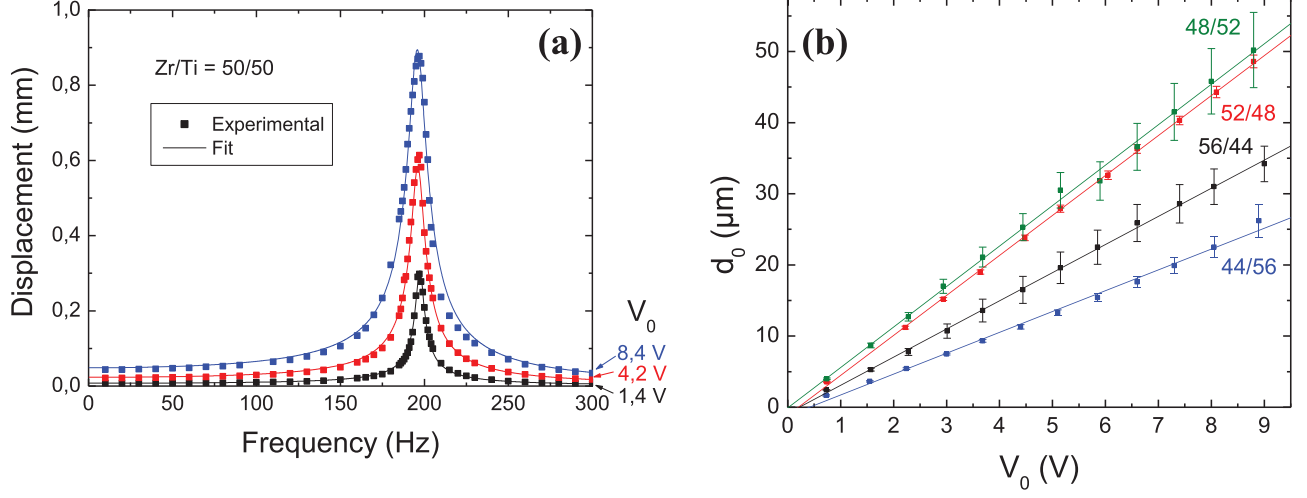


FIG. 6. (a) Tip displacement for PZT 50/50 at different applied voltages. (b) Linear fit of quasistatic displacement d_0 as a function of the applied voltage for various Zr/Ti ratios.

$$d = \frac{d_0}{\sqrt{\left(1 - \left(\frac{f}{f_0}\right)^2\right)^2 + \frac{1}{Q^2} \left(\frac{f}{f_0}\right)^2}}, \quad (1)$$

where d_0 is the quasi-static displacement, f_0 , the natural frequency, and Q , the quality factor.

A good agreement is observed between experimental data and the fits for each applied voltage. A remarkable point is the order of magnitude of the tip displacement (deflection) at the resonance frequency—around 1 mm—obtained under 8.4 V. This transverse displacement value is very uncommon for bending mode actuators made of PZT thin films and it shows that thin films have a good flexibility. With very similar PZT thicknesses and dimensions of cantilever beam (Table I), Morimoto *et al.* and Harigai *et al.* obtained less than 10 μm tip displacement.^{9,31}

In order to calculate the transverse piezoelectric coefficient d_{31} , a linear fit of the quasistatic displacement d_0 as a function of applied voltage was done, as shown in Fig. 6(b). The fit is realized with Eq. (2)^{32,33}

$$d_0 = d_{31} \times \frac{6s_a s_p t_a (t_a + t_p) \times \left(l_e l_b - \frac{l_e^2}{2}\right)}{s_a^2 t_p^4 + 4s_a s_p t_a t_p^3 + 6s_a s_p t_a^2 t_p^2 + 4s_a s_p t_a^3 t_p + s_p^2 t_a^4} \times \frac{w_e}{w_b} \times V_0 = d_{31} \times \beta \times V_0 \quad (2)$$

TABLE I. Comparison of PZT cantilever beam tip displacement under an applied voltage around 8.5 V.

	This work:		
	with/without RuO ₂	Morimoto <i>et al.</i> ⁹	Harigai <i>et al.</i> ³¹
Zr/Ti	40/60 to 60/40	52/48	53/47
PZT thickness (μm)	3	2.8	3.6
Substrate (thickness)	Aluminium foil (16 μm)	Stainless steel (50 μm)	Silicium (120 μm)
Length (cm)	1	1.7	0.77
Width (mm)	3	4.5	1.8
Tip displacement around 8.5 V (mm)	0.5–0.9/0.8–1.2	0.008	0.004

where s is the inverse of the Young's modulus, t the thickness, l_b , l_e the lengths and w_b , w_e the widths of the beam and the electrode, respectively, and V_0 is the amplitude of the applied voltage. The “ p ” and “ a ” subscripts correlate with PZT layer and aluminium substrate, respectively. The slope of the linear fits is divided by β to calculate the d_{31} value. This calculation is used in the following to find the d_{31} . Here again, the fits are in a good agreement with experimental data.

B. Influence of the PZT thickness on the electrical properties of PZT thin films

Thickness of PZT 57/43 thin film deposited onto RuO₂/Al was modified from 1.8 to 4.2 μm . Poling was done as explained before, with an electric field of 167 kV/cm. Both P_r and E_c are plotted as a function of PZT thickness in Fig. 7(a) and the piezoelectric coefficient of the PZT is given for all the tested thicknesses in Fig. 7(b). It is found that both P_r and E_c are roughly independent upon the PZT thickness. It seems that the thickness of the PZT layer has no impact on ferroelectric properties. Whatever, the electromechanical conversion depends on piezoelectric properties so the thickness dependence of d_{31} is investigated. For lower thicknesses, i.e., 1.8 and 2.4 μm , the measured d_{31} values are around 12 pC/N, while for thicknesses above 3 μm , d_{31} is multiplied at least by 2 and reaches 27 pC/N at 4.2 μm . An abrupt increase in d_{31} is observed when the PZT thickness changes from 2.4 to 3.0 μm . But for thicknesses under 3.0 μm , some difficulties appeared in the poling step. Few samples of less than 3.0 μm -thick were in short-circuit, which could mean that the insulating properties of the first PZT layers are medium. So the weak piezoelectric properties of the low thickness PZT thin layer beam may be explained by the medium quality of the material.

C. Influence of the Zr/Ti ratio on the electrical properties of PZT thin films

Different Zr/Ti ratios were tested for samples with PZT layer of 3 μm in thickness deposited onto bare and 40 nm

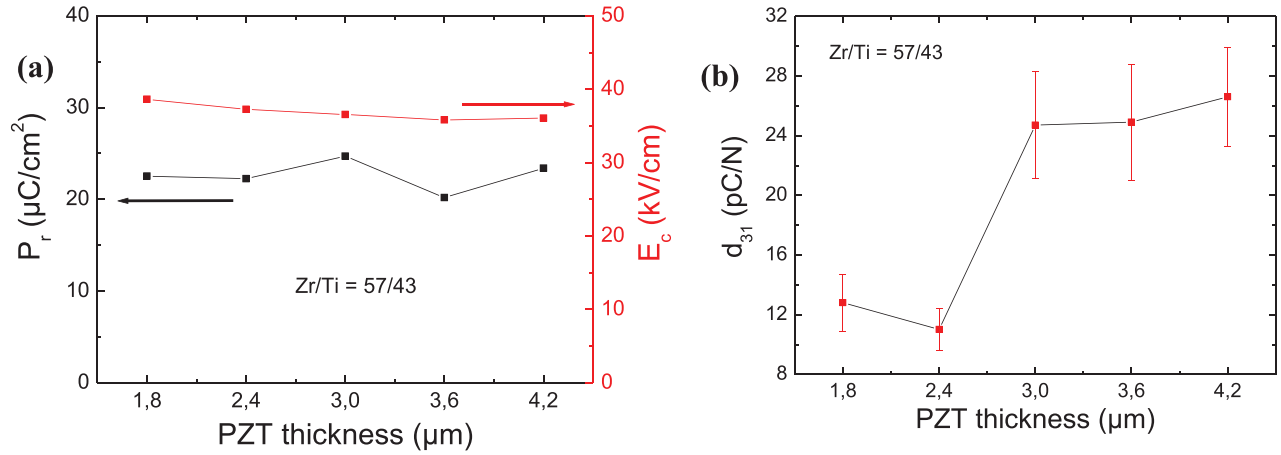


FIG. 7. (a) P_r and E_c and (b) d_{31} evolution of PZT 57/43 thin film deposited onto 40 nm RuO_2 coated aluminium substrate as a function of PZT layer thickness.

RuO_2 coated aluminium foil. This study was carried out to show the influence of Zr/Ti ratio on both ferroelectric properties and the transverse piezoelectric coefficient d_{31} . To maximize the evaluated properties of PZT thin films with and without RuO_2 , it was chosen to apply electric field values of 167 kV/cm 267 kV/cm , respectively. Under these appropriate electric field strengths, saturated hysteresis loops were obtained.

As shown in Fig. 8(a), for both PZT thin films with/without RuO_2 , there is no obvious tendency of the P_r and E_c variation versus the Zr content. However, ferroelectric properties are significantly improved with the RuO_2 layer, whatever the Zr/Ti ratio. As a matter of fact, ferroelectric saturation was reached for all samples with RuO_2 . Besides, saturation of P-E loops is obtained with an applied electric field of 167 kV/cm , whereas 267 kV/cm is necessary for the samples without the oxide layer. Moreover, with the intermediate oxide layer, the hysteresis loops are symmetric, and the coercive field is at least divided by 2. There is a quite improvement of the remnant polarization. From an average value of 14 $\mu\text{C}/\text{cm}^2$ without RuO_2 for the studied Zr contents, it rises to 24 $\mu\text{C}/\text{cm}^2$ with a 40 nm oxide layer. Ferroelectric properties are improved with the intermediate

RuO_2 layer and no clear dependence of this improvement upon the Zr/Ti ratio could be evidenced, but measurement of the d_{31} remains the principal objective of this study.

The obtained results of d_{31} calculation are plotted in Fig. 8(b). Without RuO_2 , d_{31} increases for $0.4 < x < 0.48$ and reaches a maximum of 14 pC/N for PZT compositions with %Zr = 52–54. It is corresponding typically to PZT ceramics compositions of the MPB (Zr/Ti = 52/48) where both tetragonal and rhombohedral phases coexist at room temperature in the phase diagram.¹⁷ In addition, a shift of the MPB to the high Zr content is observed for PZT thin films.³⁴ When the molar percentage of zirconium gets higher, the transverse coefficient value seems to decrease. This behaviour was already observed in the literature for both thick³⁵ and thin PZT films.³⁶ However, experimental d_{31} values of PZT on aluminium substrate are lower than those reported in the literature for PZT on silicon (Si, 118.9 pC/N (Ref. 37)) or stainless steel (SS, 55 pC/N (Ref. 38)). But the difference could be explained by the originality of the flexible aluminium substrate. Indeed, the PZT thin layer (3 μm , 92.5 GPa evaluated from Kanno *et al.*³⁹) is not constrained by the uncommon low thickness (16 μm) and the weak Young's modulus (65 GPa) of the substrate (commercial aluminium

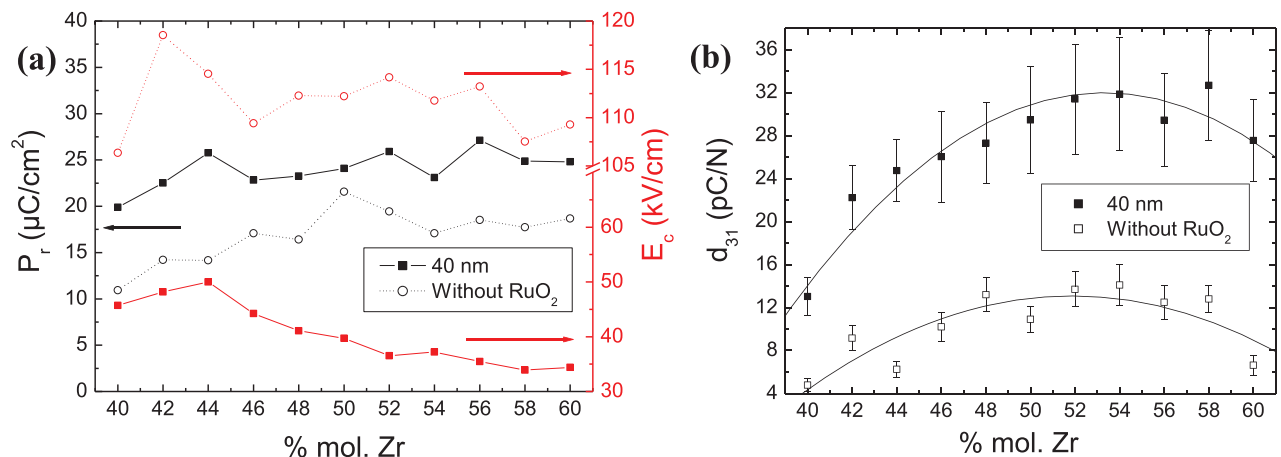


FIG. 8. (a) Variation of the remnant polarization and coercive field as a function of Zr content. (b) Variation of the transverse piezoelectric coefficient d_{31} as a function of Zr content. Solid lines are guides for the eyes. Comparison between PZT thin film with and without RuO_2 .

foil) like it would be with Si (130.2 GPa for $\langle 100 \rangle$ oriented silicon) or SS (~ 200 GPa).

Besides, a large increase of d_{31} is obtained with the intermediate oxide layer. As a matter of fact, it is multiplied by an average factor of 2.7, for the different studied Zr contents.

As seen before, when the PZT crystallizes in a tetragonal phase ($x < 0.48$), the transverse piezoelectric coefficient increases with increasing Zr content. With RuO₂, the maximum is not observed in the vicinity of the MPB but for the PZT thin film of composition 58/42 ($d_{31} = 33$ pC/N). This maximum is still a little bit lower than d_{31} values of thin films given in the literature.^{9,40} However, in view of potential applications, the low d_{31} values of PZT/Al are compensated by the remarkable flexibility, the associated low working frequencies, and the low cost of fabrication, which render these lightweight PZT thin films promising materials for some applications like gas flow sensing, wearable electronics, or micro-energy harvesting.

Nevertheless, concerning the Zr/Ti ratio, it is difficult to choose an optimal composition with the uncertainty of the d_{31} measurement for Zr content between 52% and 58%. If we pay attention to the trend of the curve, it seems the maximum d_{31} could be near the MPB, which is a similar trend previously observed without RuO₂. However, the optimal composition seems to be shifted towards the rhombohedral phase.

Finally, the 40 nm RuO₂ layer has a real positive impact on ferroelectric and piezoelectric properties of PZT/Al thin films. In fact, this layer acts as a diffusion barrier at the PZT-aluminium interface. The density of oxygen vacancies—which have trapped electrons when a positive field is applied—is reduced with RuO₂. So, the real field applied to the PZT layer, for the same amplitude, may be higher when the intermediate RuO₂ layer is used.

IV. CONCLUSIONS

PZT thin films of various compositions and thicknesses were deposited onto bare and RuO₂ coated aluminium. Dielectric, ferroelectric, and piezoelectric characteristics were measured on bending mode cantilever beams. Results show the influence of the intermediate oxide layer (RuO₂) which contributes to reduce the amount of oxygen vacancies in the first deposited PZT layers and improve almost all the studied properties. For a 3 μm thick PZT thin film with different Zr/Ti ratios, an interfacial RuO₂ layer of 40 nm in thickness yields a large increase of average remnant polarization P_r is around 14 $\mu\text{C}/\text{cm}^2$ to 24 $\mu\text{C}/\text{cm}^2$, while the average coercive field E_c is divided by at least 2 and d_{31} is increased with an average factor of 2.7.

It was found that the optimal material parameters for measurement of d_{31} coefficient are the presence of the 40 nm thick RuO₂ layer and a PZT thickness of 3 μm . Besides, the Zr/Ti ratio must lie between 52/48 and 58/42, which means that PZT compositions on the rhombohedral side of the MPB give the best results. A maximal d_{31} value of 33 pC/N was obtained for Zr/Ti = 58/42.

Finally, the optimal properties of the thin film are of the same order of magnitude than the observed values in the

literature. However, a better flexibility—with displacement around 1 mm under low voltage (~ 8 V)—is noticed for our films due to the low aluminium substrate thickness and Young's modulus, which is asset for electromechanical applications requiring large displacements induced by low voltage appliance. Such PZT thin film cantilevers may also be of interest for aeroelastic energy harvesting from low frequency vibration. In this sense, future work will be devoted to the direct piezoelectric effect in these recently developed PZT/Al thin films.

¹J. Kymissis, C. Kendall, J. Paradiso, and N. Gershenfeld, in *Second International Symposium on Wearable Computing 1998, Digest of Papers* (IEEE, 1998), pp. 132–139.

²J. J. H. Paulides, J. W. Jansen, L. Encica, E. A. Lomonova, and M. Smit, in *IEEE International Electric Machines and Drives Conference 2009, IEMDC09* (IEEE, 2009), pp. 439–444.

³R. L. Harne, *Mech. Syst. Signal Process.* **36**, 604 (2013).

⁴A. Chandrakasan, R. Amirtharajah, J. Goodman, and W. Rabiner, in *Proceedings of 1998 IEEE International Symposium on Circuits and Systems 1998, ISCAS98* (IEEE, 1998), pp. 604–607.

⁵G. Poulin, E. Sarraute, and F. Costa, *Sens. Actuators Phys.* **116**, 461 (2004).

⁶J. C. Park and J. Y. Park, *Ceram. Int.* **39**, S653 (2013).

⁷Y. B. Jeon, R. Sood, J.-H. Jeong, and S.-G. Kim, *Sens. Actuators Phys.* **122**, 16 (2005).

⁸P. Murali, M. Marzencki, B. Belgacem, F. Calame, and S. Basrour, *Proc. Chem.* **1**, 1191 (2009).

⁹K. Morimoto, I. Kanno, K. Wasa, and H. Kotera, *Sens. Actuators Phys.* **163**, 428 (2010).

¹⁰S. Trolier-McKinstry and P. Murali, *J. Electroceram.* **12**, 7 (2004).

¹¹C. Dagdeviren, B. D. Yang, Y. Su, P. L. Tran, P. Joe, E. Anderson, J. Xia, V. Doraiswamy, B. Dehdashti, X. Feng, B. Lu, R. Poston, Z. Khalpey, R. Ghaffari, Y. Huang, M. J. Slepian, and J. A. Rogers, *Proc. Natl. Acad. Sci.* **111**, 1927 (2014).

¹²M. D. Nguyen, M. Dekkers, H. N. Vu, and G. Rijnders, *Sens. Actuators Phys.* **199**, 98 (2013).

¹³M.-A. Dubois and P. Murali, *Sens. Actuators Phys.* **77**, 106 (1999).

¹⁴J. F. Shepard, Jr., P. J. Moses, and S. Trolier-McKinstry, *Sens. Actuators Phys.* **71**, 133 (1998).

¹⁵J. E. A. Southin, S. A. Wilson, D. Schmitt, and R. W. Whatmore, *J. Phys. Appl. Phys.* **34**, 1456 (2001).

¹⁶I. Kanno, H. Kotera, and K. Wasa, *Sens. Actuators Phys.* **107**, 68 (2003).

¹⁷B. Jaffe, W. R. Cook, and H. L. Jaffe, *Piezoelectric Ceramics* (Academic Press, London, 1971).

¹⁸B. Noheda, D. E. Cox, G. Shirane, J. A. Gonzalo, L. E. Cross, and S.-E. Park, *Appl. Phys. Lett.* **74**, 2059 (1999).

¹⁹A. K. Tagantsev and I. A. Stolichnov, *Appl. Phys. Lett.* **74**, 1326 (1999).

²⁰S.-H. Kim, J.-S. Yang, C. Y. Koo, J.-H. Yeom, E. Yoon, C. S. Hwang, J.-S. Park, S.-G. Kang, D.-J. Kim, and J. Ha, *Jpn. J. Appl. Phys., Part 1* **42**, 5952 (2003).

²¹R. Seveno and D. Averty, *J. Sol-Gel Sci. Technol.* **68**, 175 (2013).

²²R. Seveno, P. Limousin, D. Averty, J.-L. Chartier, R. Le Bihan, and H. W. Gundel, *J. Eur. Ceram. Soc.* **20**, 2025 (2000).

²³J. H. Yi, R. Seveno, and H. W. Gundel, *Integr. Ferroelectr.* **23**, 199 (1999).

²⁴M. D. Losego, L. H. Jimison, J. F. Ihlefeld, and J.-P. Maria, *Appl. Phys. Lett.* **86**, 172906 (2005).

²⁵B. G. Chae, C. H. Park, Y. S. Yang, and M. S. Jang, *Appl. Phys. Lett.* **75**, 2135 (1999).

²⁶Q. Zou, H. E. Ruda, B. G. Yacobi, K. Saegusa, and M. Farrell, *Appl. Phys. Lett.* **77**, 1038 (2000).

²⁷M. H. M. Zai, A. Akiba, H. Goto, M. Matsumoto, and E. M. Yeatman, *Thin Solid Films* **394**, 96 (2001).

²⁸Y.-L. Tu, M. L. Calzada, N. J. Phillips, and S. J. Milne, *J. Am. Ceram. Soc.* **79**, 441 (1996).

²⁹N. Klein, *IEEE Trans. Electron Devices* **13**, 788 (1966).

³⁰F. K. Kneubühl, *Oscillations and Waves* (Springer, Berlin, 1997).

³¹T. Harigai, H. Adachi, and E. Fujii, *J. Appl. Phys.* **107**, 096101 (2010).

³²J. G. Smits and W.-S. Choi, *IEEE Trans. Ultrason. Ferroelectr. Freq. Control* **38**, 256 (1991).

- ³³J. G. E. Gardeniers, A. G. B. J. Verholen, N. R. Tas, and M. C. Elwenspoek, *J. Korean Phys. Soc.* **32**, S1573 (1998).
- ³⁴S. Hoon Oh and H. Jang, *Phys. Rev. B* **62**, 14757 (2000).
- ³⁵E. Boucher, B. Guiffard, L. Lebrun, and D. Guyomar, *Ceram. Int.* **32**, 479 (2006).
- ³⁶N. Ledermann, P. Murali, J. Baborowski, S. Gentil, K. Mukati, M. Cantoni, A. Seifert, and N. Setter, *Sens. Actuators Phys.* **105**, 162 (2003).
- ³⁷M. Dekkers, H. Boschker, M. van Zalk, M. Nguyen, H. Nazeer, E. Houwman, and G. Rijnders, *J. Micromech. Microeng.* **23**, 025008 (2013).
- ³⁸S. Umemiya, M. Hida, T. Aoki, and M. Kondo, in *15th IEEE International Symposium on Applications of Ferroelectrics 2006, ISAF06* (IEEE, 2006), pp. 318–321.
- ³⁹I. Kanno, S. Fujii, T. Kamada, and R. Takayama, *Appl. Phys. Lett.* **70**, 1378 (1997).
- ⁴⁰D.-M. Chun, M. Sato, and I. Kanno, *J. Appl. Phys.* **113**, 044111 (2013).