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INSULATED MOLECULAR WIRES: SHEATHING SEMI-CONDUCTING POLYMERS WITH ORGANIC NANOTUBES THROUGH HETEROGENEOUS NUCLEATION

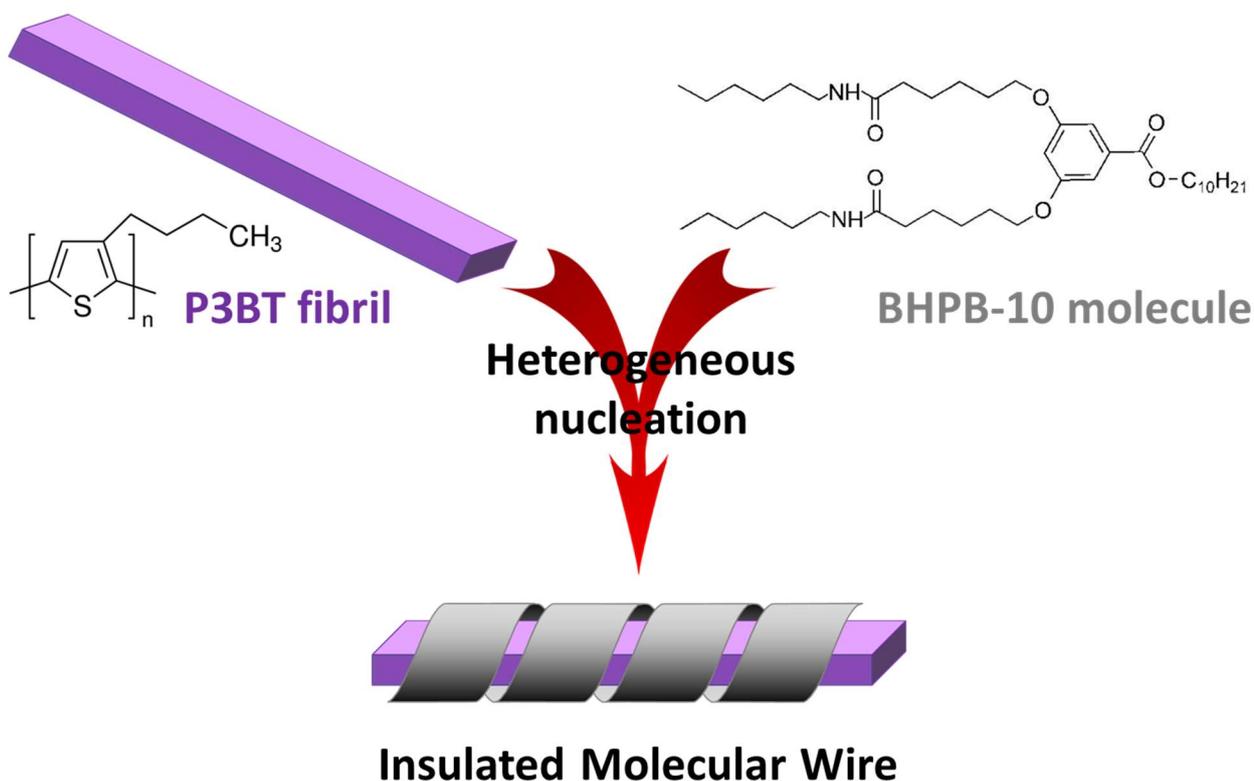
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Abstract

Insulated molecular wires formed from organic molecules may have promising applications in organic and flexible electronic devices. Here, we provide compelling evidence of the formation of insulated Molecular Wires by sheathing conducting regioregular Poly(3-butylthiophene-2,5-diyl) (P3BT) fibrils with insulating

nanotubes from 3,5-bis-(5-hexylcarbamoylpentyloxy)-benzoic acid decyl ester molecules through a nucleation and growth process. For dilute systems, conducting Atomic Force Microscopy (C-AFM) together with Force-distance curves and Current-Voltage spectroscopy are concomitantly performed to sense current from the topmost surface to the core of the composite fibrils at controlled tip indentation depths. Results show that current is sensed only when the core of the nanocomposite fibrils is reached which indicates the presence of an insulating layer around the semi-conducting P3BT fibrils. The existence of this molecular nanocomposite is further supported by neutron scattering experiments carried out on more concentrated systems at different temperatures.

Introduction

One-dimensional (1D) nanomaterials, such as nanowires, nanorods and nanotubes, show unique and fascinating properties. In order to enhance some of their properties, to combine two complementary functional materials or to increase their stability or durability, coaxial hybrid nanomaterials have been intensively studied^{1,2}. This class of materials have received growing interest owing to their essential implication in both fundamental investigation of physical phenomena (such as charge carrier transport and spin propagation) and practical application development (electronic and optoelectronic devices)²⁻⁶. Methods using top-down (microfabrication) or bottom-up (e.g. Chemical Vapour Deposition or Atomic Layer Deposition) processes have yielded inorganic core/shell nanowires used in photovoltaic devices^{2,3,5}, light emitting diodes⁷, biological and chemical species detection⁸. However, such approaches require extremely stringent methodology and does not lead to nanowires with diameter less than a few tens of nm^{3,9,10} (top-down) or require precursor materials or high temperature process⁵ (bottom-up). To overcome these limitations, intense research was focused onto a new class of materials designated as Insulated Molecular Wires (IMW)¹¹. IMW or organic coaxial nanowires have emerged as alternative materials that can be prepared by chemical synthesis¹²⁻¹⁵, surface functionalization¹⁶ or non-covalent self-assembly¹⁶⁻²¹. The latter is clearly highly promising as it is the easiest way to build up complex architectures, provided that supramolecular assembly is mastered²². In this article, we describe the making and characterization of quasi-1D insulated nanowires made from a pi-conjugated polymer and a small nanotube-forming organic

molecule. In contrast to the IMW obtained by complex chemical processes, the non-covalent assembly strategy followed here is based only on two physical processes, heterogeneous nucleation and thermo-reversible gelation. The proof of concept for making materials through this process has been recently obtained with a non-conducting polymer, isotactic polystyrene, sheathed by the same self-assembling molecule used here²³. In the present work, nanocomposite fibrils are made up of a Poly(3-butylthiophene-2,5-diyl) (P3BT) π -conjugated polymer-sheathed by a nanotube of a self-assembling molecule, namely 3,5-bis-(5-hexylcarbamoylpentyloxy)-benzoic acid decyl ester (BHPB-10)²⁴. P3BT was chosen because it belongs to the poly(3-alkylthiophene) family which is well-known in the organic electronics²⁵ and for which fibrils has already shown interesting properties^{26,27}. The length of our composites fibers is up to several microns, one order of magnitude longer than the previously reported IMWs obtained from solution process^{11,13,14,28}. Our fibers have been characterized by Small Angle Neutron Scattering (SANS) as well as Atomic Force Microscopy (AFM). Current spectroscopy has been carried out by Conducting AFM (C-AFM) to probe the conducting core/insulating shell nature of our material. In addition, cross-talk effects between the nanocomposite fibers are absent on micrometer scale, a pre-requisite to their use as IMWs in organic electronics.

The systems

The different systems observed by TEM are shown in Fig. 1. P3BT forms fibrils of square cross-section with dimensions within 5 to 15 nm; BHPB-10 forms ringlets that are locally nanotubes with outer diameter 12 nm and inner diameter 9 nm. If a homogeneous mixture of P3BT and BHPB-10 is prepared by heating above 100°C and then quenched, one can observe that BHPB-10 nucleates the growth of P3BT fibrils. Conversely, when the protocol described in Fig. 1 is implemented then the hybrid system appears quite different. Mixing the two components above 100°C so as to obtain a homogeneous solution, and then quenching to 0°C entails the nucleation of the P3BT fibrils by the BHPB-10 nanotubes. This is so because the formation kinetics of P3BT crystals is much slower than the formation of nanotubes. Conversely, when the BHPB-10 in solution in *trans*-decahydronaphthalene (BHPB-10 SOL) is added at 60°C to a suspension of fibrils and then cooled to 0°C, the hybrid system is produced.

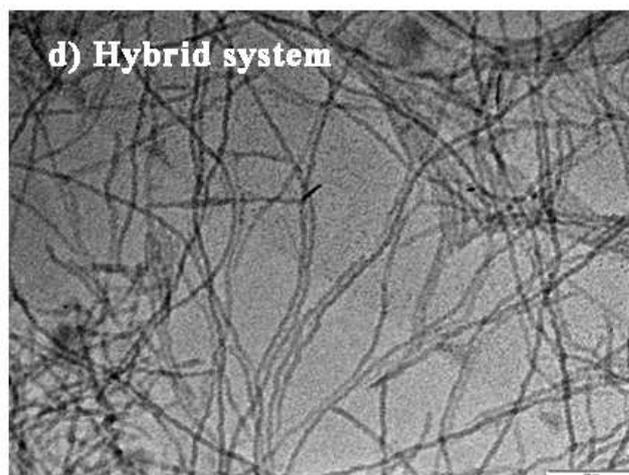
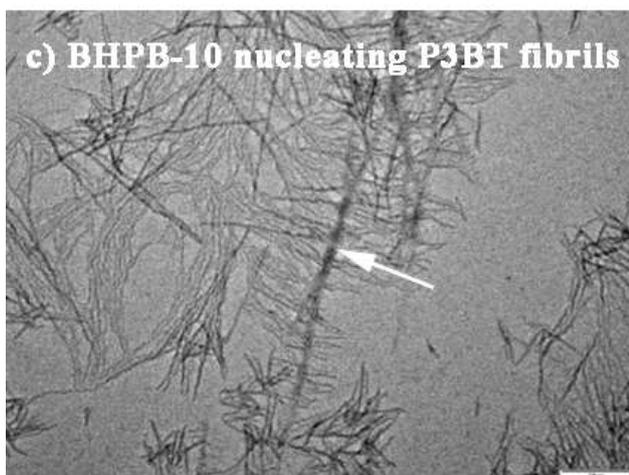
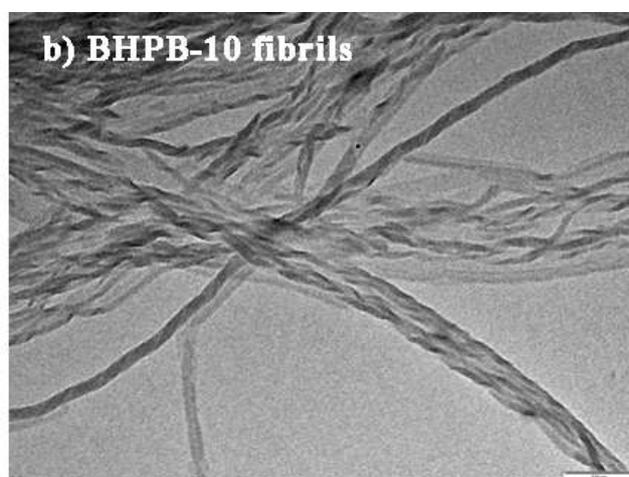
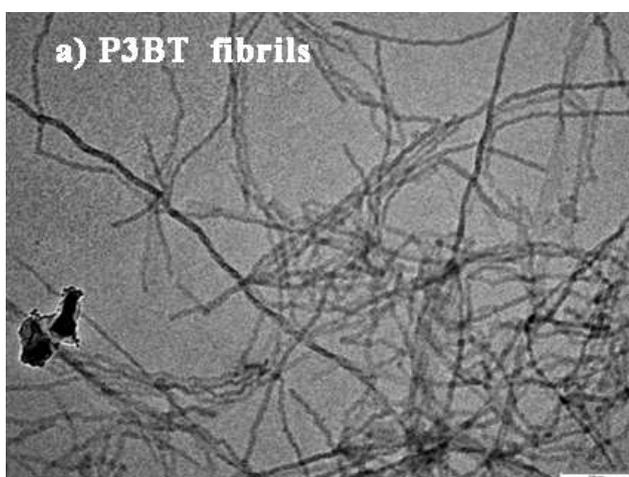
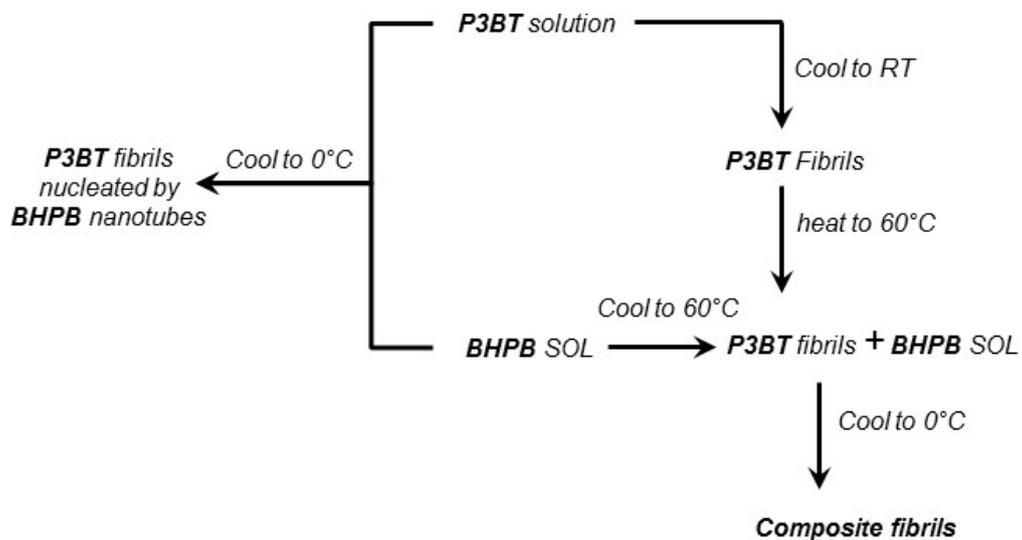


Figure 1: The preparation protocol, top; bottom TEM pictures obtained on: a) purified P3BT fibrils prepared in trans-decahydronaphthalene at room temperature; b) BHPB-10 ringlets produced by quenching at 0°C solutions in trans-decahydronaphthalene; c) P3BT fibrils nucleated onto BHPB-10 nanotubes (arrow) by cooling a homogeneous mixture ($0.05 \times 10^2 \text{g/cm}^3 \text{P3BT} + 0.01 \times 10^2 \text{g/cm}^3 \text{BHPB-10}$); d) hybrid system obtained by mixing P3BT fibrils grown beforehand with BHPB-10 SOL at 60°C and then quench to 0°C. White scale bar= 200 nm. Note that the systems are diluted to 10 to 100 times to allow for TEM observations.

Current spectroscopy using C-AFM technique at controlled indentation depths

Conducting Atomic Force Microscopy (C-AFM) has been used to probe the electrical properties of the molecular nanostructures under investigation, namely, P3BT fibrils, BHPB-10 fibrils and the hybrid nanomaterial. Current-Voltage (I-V) spectroscopy on C-AFM setups is a tool to have access to the electronic nature of the investigated material deposited on a conducting substrate²⁹: no collected current if the material is electrically insulating, non-linear I-V response if the material is semi-conducting and a linear I-V response if the material is conducting. Force-distance (FZ) experiments are generally used to measure intrinsic properties such as stiffness, hydrophobicity or surface charge density of a given material through the adhesion forces felt by the AFM tip upon retraction^{30,31}. The precise vertical Z-displacement of the piezoelectric scanner allows controlling the tip penetration with respect to the material surface. In the present work, FZ curves and I-V spectroscopy are concomitantly used to sense current from the topmost surface of the fiber as well as from the fiber core at controlled indentation depths (figure 2A).

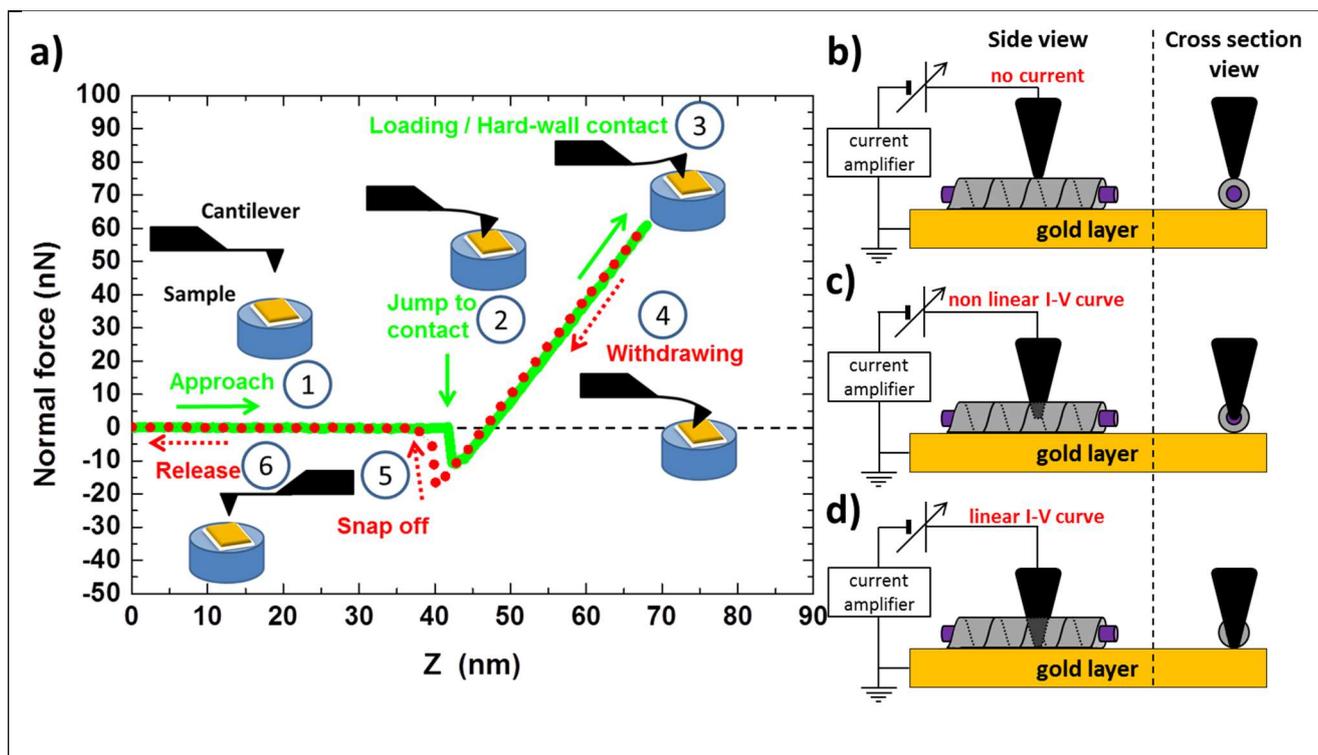


Figure 2| A) Force-distance AFM experiment principle. a) Typical force-distance FZ curve showing different regions of tip-sample interaction, b) a conducting AFM tip is used to probe current on the surface of the hybrid fiber (no current, insulating shell), c) indenting to reach the fiber core (non linear I-V curve, semiconducting core), and finally d) touching the gold substrate (linear I-V curve, conducting substrate).

The organic material is spin-coated on an 80nm thick gold layer which is grounded using a conductive silver paint. The current is probed by means of a conducting platinum silicide (PtSi) AFM tip (Nanosensors, model: PtSi-FM) as a second terminal. I-V curves recorded on a bare grounded gold electrode give a linear relation with a steep slope as expected in the case of a conducting material (Supplementary Fig. 2).

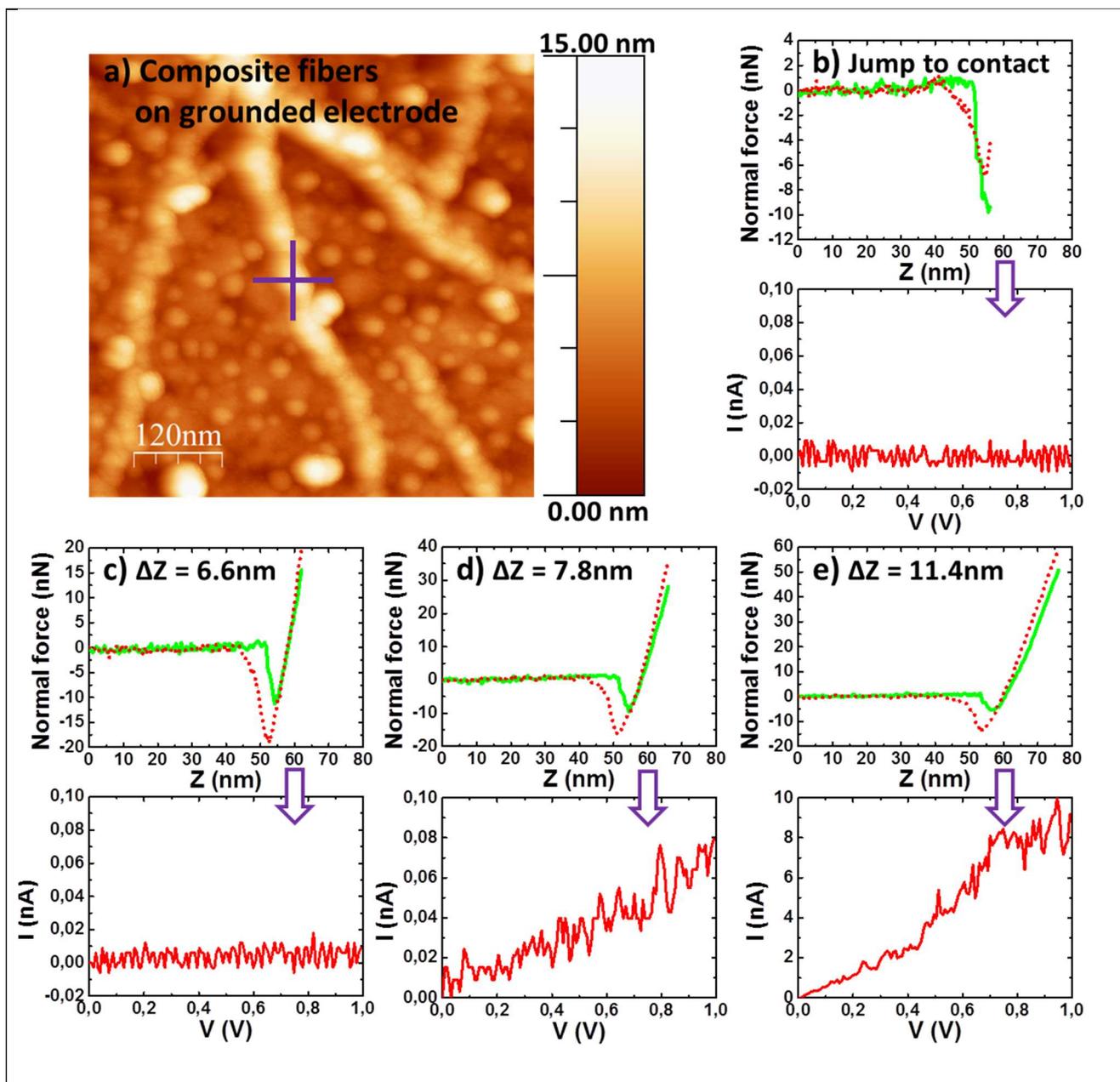


Figure 2 | B) Current-Voltage spectroscopy on hybrid fibers at controlled indentation depth. a) Topography of P3BT-BHPB-10 composite fibers on a grounded Au electrode. b-e) FZ curves at different tip depth position (ΔZ) into the fiber (at the point located by the purple cross) and the related I-V curves measured for each ΔZ , showing the change in the electronic properties (from insulating shell to semiconducting core).

The principle of the FZ approach is illustrated on Fig. 2A: the force felt by the cantilever is recorded as the function of the piezoelectric vertical displacement Z of the sample holder. When the tip is far from the surface, the force is zero. Upon further approach, a jump to contact occurs when a negative deflection is seen in the FZ curve arising from the attraction of the tip to the surface on account of long-range attractive forces. On further loading the force becomes proportional to the displacement while the penetration depth of the tip into the material depends on the sample stiffness. The Z distance after the 'jump to contact' is the sum of cantilever deflection and the tip penetration/indentation depth³⁰. A reference sample made of freshly cleaved mica is used to calculate accurately the penetration depth ΔZ ($=Z_{\text{material}}-Z_{\text{mica}}$). Upon tip withdrawing (Z decreasing), the sharp decrease to zero force corresponds to the contact snap off. The electronic properties of P3BT, BHPB-10, and the hybrid fibers are characterized by I-V curves registered at several tip depths into the different organic materials. Experimental details can be found in the supplementary material.

On pure P3BT fibrils, current is measured instantly at the jump to contact of the tip with the fiber surface. Non-linear I-V curves indicate the semi-conducting nature of pure P3BT as expected (Supplementary Fig.3). The current saturation level (at 100 nA) is reached when the tip reaches the gold layer, which corresponds to $\Delta Z= 4.5$ nm. This value is in excellent agreement with the height of the fiber derived from standard AFM cross-section analysis (4.2 nm).

In the case of pure BHPB-10 nanotubes, no current is sensed upon jump into contact or further indentation, confirming its insulating nature (Supplementary Fig. 4), until the tip touches the grounded electrode, leading to current saturation.

For the composite fibrils, no current is sensed upon jump to contact (Fig. 2B) and down to a distance of 6.6 nm from the top surface of the fiber (see Fig. 2B(c)). Thus the conducting AFM tip is first sensing a layer of insulating material on the top of the fiber surface, as illustrated on Fig. 2 A(b). A small value of current is being sensed when tip penetrates 7.8 nm into the composite fiber (Fig. 2B (d)). Here the AFM tip is just approaching the conducting part of the composite fiber. A further increment in the penetration depth results in a clear jump of current (≈ 10 nA for 1V) as the tip establishes subsequent contact with the semi-

conducting domain (in our case P3BT) as ascertained by the non-linear variation of the I-V curve on Fig. 2 B€ and as illustrated on Fig. 2 A (c). Allowing the tip to penetrate deeper into the fibril entails a current saturation which corresponds to the reaching of the gold grounded electrode (Fig. 2 A(d)). This experiment reveals that the I-V response of a P3BT-BHPB-10 composite fiber is very different when compared to that of a pure P3BT fiber, in the later where current is sensed at the very instant when the tip is in contact with the fiber. The fact that current is sensed only from the core of the fiber and not from the topmost layers is indeed direct evidence to the core-shell structure composed of a conducting P3BT core sheathed by an insulating BHPB-10.

Here, the question may arise as to whether there is a real sheathing through a heterogeneous nucleation process or simply P3BT fibrils being accidentally covered by BHPB-10 material. To test this point, films were prepared by first spin coating P3BT fibrils on a gold electrode, followed by the spin coating of BHPB-10 material, and drying under vacuum. Cross-section analysis of the sequentially deposited film show two types of fibers with 5 nm and 30 nm heights that correspond to the heights of P3BT fibrils and BHPB-10 nanotubes, respectively. This clearly indicates that P3BT and BHPB-10 fibrils are well-separated entities under these conditions. FZ and I-V measurements on BHPB-10 bundles show that no current is sensed at the jump to contact region. When the tip traveled a distance equivalent to the BHPB-10 fiber bundle height, current is sensed from underlying P3BT/Au electrode (See Supplementary Fig. 4).

Further investigations into the electrical properties of the fibers were performed by mapping current flowing across ground and isolated Au electrodes separated by an insulating (SiO₂) 5µm-wide channel. The idea is to highlight the proof of concept of these composite fibers as IMWs: conduction of current by the inner part along the fiber axis and electrical insulation by the outer shell. On this purpose, we used contact mode C-AFM, where current maps are generated concomitantly with topography images. A Pt/Ir coated tip (Nanosensors, model: PPP-CONTPt) is scanned over the sample surface with an applied bias voltage (fixed at 1V) between the ground electrode and the tip end (acting as the second terminal): for each pixel of the C-AFM image the value of the collected current at this position is recorded. The 1V voltage was chosen to unambiguously distinguish between the insulating or conducting nature of the materials under consideration. The principle of the experiment is drawn in Fig. 2 C(a) which represents the five possible

AFM tip positions in the case of the hybrid fibers: i) directly in contact with the grounded gold electrode (current flows), ii) on top of a hybrid fiber on the grounded electrode (no current, configuration similar to Fig. 2-b), iii) on top of a hybrid fiber in the insulating SiO₂ channel (no current), iv) on top of a hybrid fiber on the isolated electrode (no current) and v) directly in contact with the isolated gold electrode (current flows thanks to the hybrid fibers acting as IMWs). Current map of the hybrid fibers is shown on Fig. 2 C, while that of the pure P3BT and pure BHPB-10 fibers can be found in the Supplementary Materials (Supplementary Fig. 5). In the case of hybrid fibers, current appears on both grounded (case i) and isolated (case v) electrodes revealing that these electrodes are electrically connected. In addition, i) the channel area is devoid of any current despite topography image showing fibers across the channel and ii) these hybrid fibers appear dark (i.e. no current) on the grounded and isolated electrodes: these two points clearly point out that the hybrid fibers have an insulating outer part, which impedes the current flow (cases ii, iii and iv). Though current appears at the isolated electrode, the absence of current in the channel is at variance with the high current value observed in the channel for pure P3BT fibrils. Indeed, when pure P3BT fibrils are deposited, current appears with the same intensity on both electrodes as well as across the insulating channel between them, indicating the conducting nature of P3BT fibers (Supplementary Fig. 5a). In the case of pure BHPB-10 fibers, current appears only on the ground electrode while no current is sensed on the channel or at the isolated electrode. It is interesting to note here that areas corresponding to BHPB-10 fibers on the ground electrode appear dark in the current map in agreement with the insulating character of BHPB-10 (See Supplementary Fig. 5b). These C-AFM results confirm the sheathing of BHPB-10 over P3BT that renders the composite fiber insulating properties perpendicular to the fiber axis (absence of channel current, hybrid fibers appearing dark even on the grounded electrode), and conducting properties in the co-axial direction (current appears on the isolated electrode). From these experiments, it is also clear that the electrical cross-talk is prevented even on large fiber bundles and over large distances (5µm-wide channel). This confirms the use of our composite material as efficient IMWs.

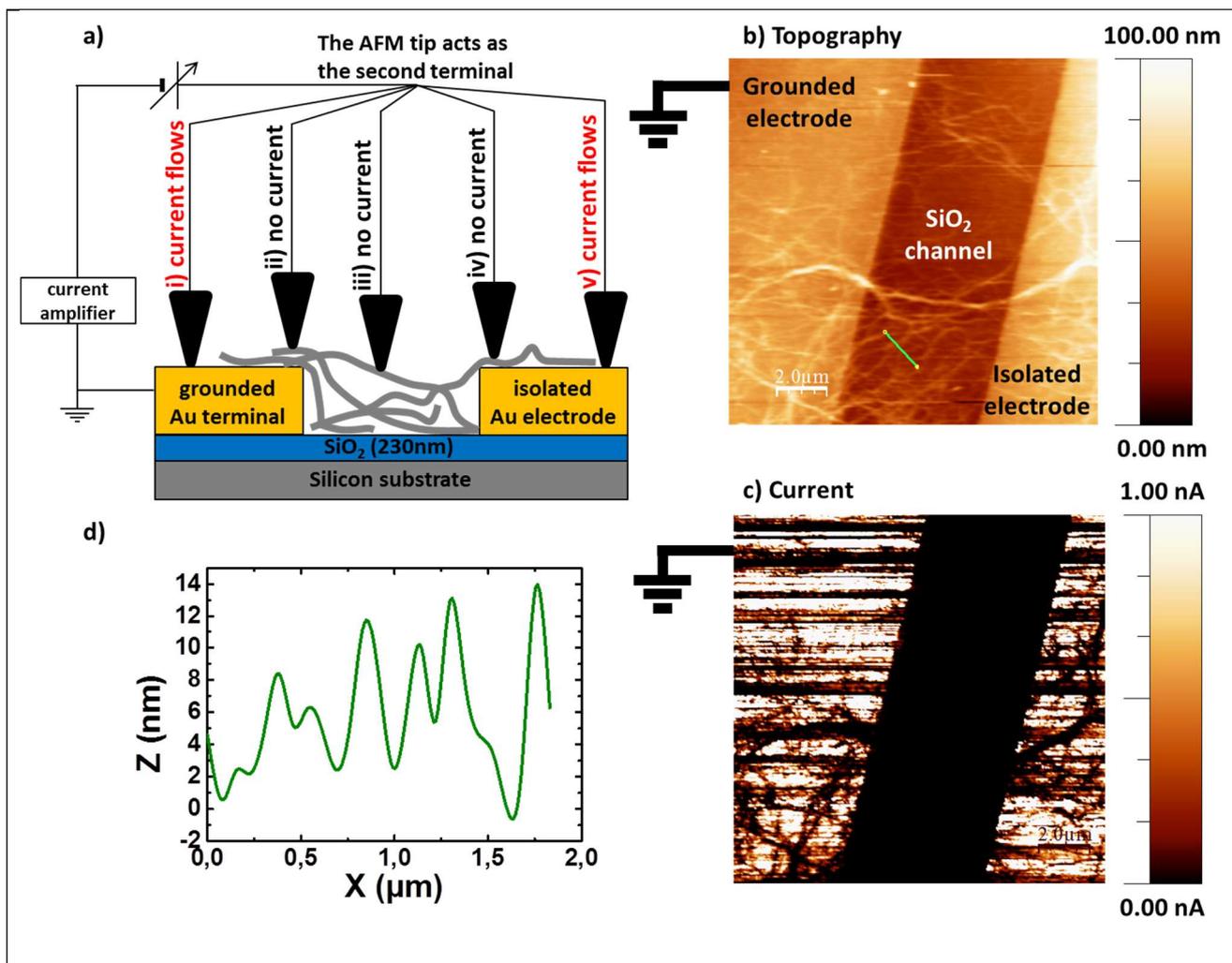


Figure 2 | C) C-AFM measurements performed on hybrid fibers deposited in-between gold electrodes.

a) Schematic representation. One electrode is grounded while the other one is isolated and separated by an insulating SiO₂ channel. Hybrid fibers are spin-coated on the substrate. The C-AFM current mapping principle is the following: the current flows across the grounded electrode and the AFM tip which acts as the second terminal. The five possible AFM tip positions are represented: i) directly in contact with the grounded gold electrode (current flows), ii) on top of a hybrid fiber on the grounded electrode (no current, configuration similar to Fig. 1b), iii) on top of a hybrid fiber in the insulating SiO₂ channel (no current), iv) on top of a hybrid fiber on the isolated electrode (no current) and v) directly in contact with the isolated gold electrode (current flows thanks to the hybrid fibers acting as IMWs).

b) and c) Topography images and current mapping are recorded simultaneously. The level of color in the current mapping indicates if the current flows or not at this tip position in contact with the material: a bright pixel indicates that a conduction path exist between the spot position and the grounded electrode (e.g. tip directly on top of the grounded electrode), a dark pixel indicates that no conduction path exists between the tip and the grounded electrode.

These P3BT-BHPB-10 composite fibers show current flowing between both electrodes (case v) but the channel appears dark despite the presence of fibers present on the channel (case iii), thus confirming the efficient insulating sheathing.

d) AFM surface profile of the composite fibers.

Small-Angle Neutron Scattering experiments

The small-angle neutron scattering were carried out on three samples: BHPB-10/*trans*-decahydronaphthaleneD, $C_{\text{BHPB-10}} = 0.01 \text{ g/cm}^3$, P3BT/*trans*-decahydronaphthaleneD, $C_{\text{P3BT}} = 0.01 \text{ g/cm}^3$, and the ternary system BHPB-10/P3BT/*trans*-decahydronaphthaleneD, $C_{\text{BHPB-10}} = 0.01 \text{ g/cm}^3 + C_{\text{P3BT}} = 0.01 \text{ g/cm}^3$. These systems were studied at 25°C, 47°C, and 60°C, respectively. The results are presented in Fig.3 by means of a Kratky-representation ($q^2I(q)$ vs q). The scattering curve of BHPB-10 can be well-fitted with the scattering function of a hollow cylinder:

$$I(qr) = \frac{\pi\mu_L}{q} \left[\frac{2}{(1-\gamma^2)qr} \times (J_1(qr) - \gamma J_1(\gamma qr)) \right]^2 \quad [1]$$

where r is the outer radius of the cylinder, γ the ratio between the inner and the outer radii, and μ_L the mass per unit length.

The fit yields $r_{\text{out}} = 12 \text{ nm}$ and $\gamma = 0.75$, results in good agreement with previous findings^{23,24} on this molecule. As P3BT forms fibrils with rectangular cross-sections, a good fit can be obtained for the P3BT scattering curve by considering straight ribbons whose theoretical form factor reads:

$$\varphi(qr) = \frac{2}{\pi} \int_0^{\pi/2} \left[\frac{\sin qa/2 \cos \theta}{qa/2 \cos \theta} \times \frac{\sin qb/2 \sin \theta}{qb/2 \sin \theta} \right]^2 \sin \theta d\theta \quad [2]$$

The fit presented in Fig. 5a corresponds to $a = 4 \text{ nm}$ and $b = 14 \text{ nm}$. The diagonal of the cross-section amounts therefore to about 15 nm. These values correspond to values that we have observed by TEM and AFM]. A better fit could certainly be obtained by taking into account a slight cross-sectional polydispersity but this is not the direct concern of this paper.

As highlighted in Fig. 3, the scattering curve for the ternary system, $I_{\text{Bternary}}^{\text{exp}}$, is not the sum of the scattering curves determined for the binary systems as it should have been if no particular interaction were taking place. It also shows oscillations, yet the maximum at $q = 0.31 \text{ nm}^{-1}$ observed for this sum is shifted to $q = 0.25 \text{ nm}^{-1}$, and the magnitude of the intensities do not match. Either intramolecular terms have to be considered or the occurrence of a hybrid system involving the P3BT fibrils and the BHPB-10 nanotubes ought to be considered. As will be highlighted below, the results obtained at 60°C definitely allows one to disregard intermolecular terms.

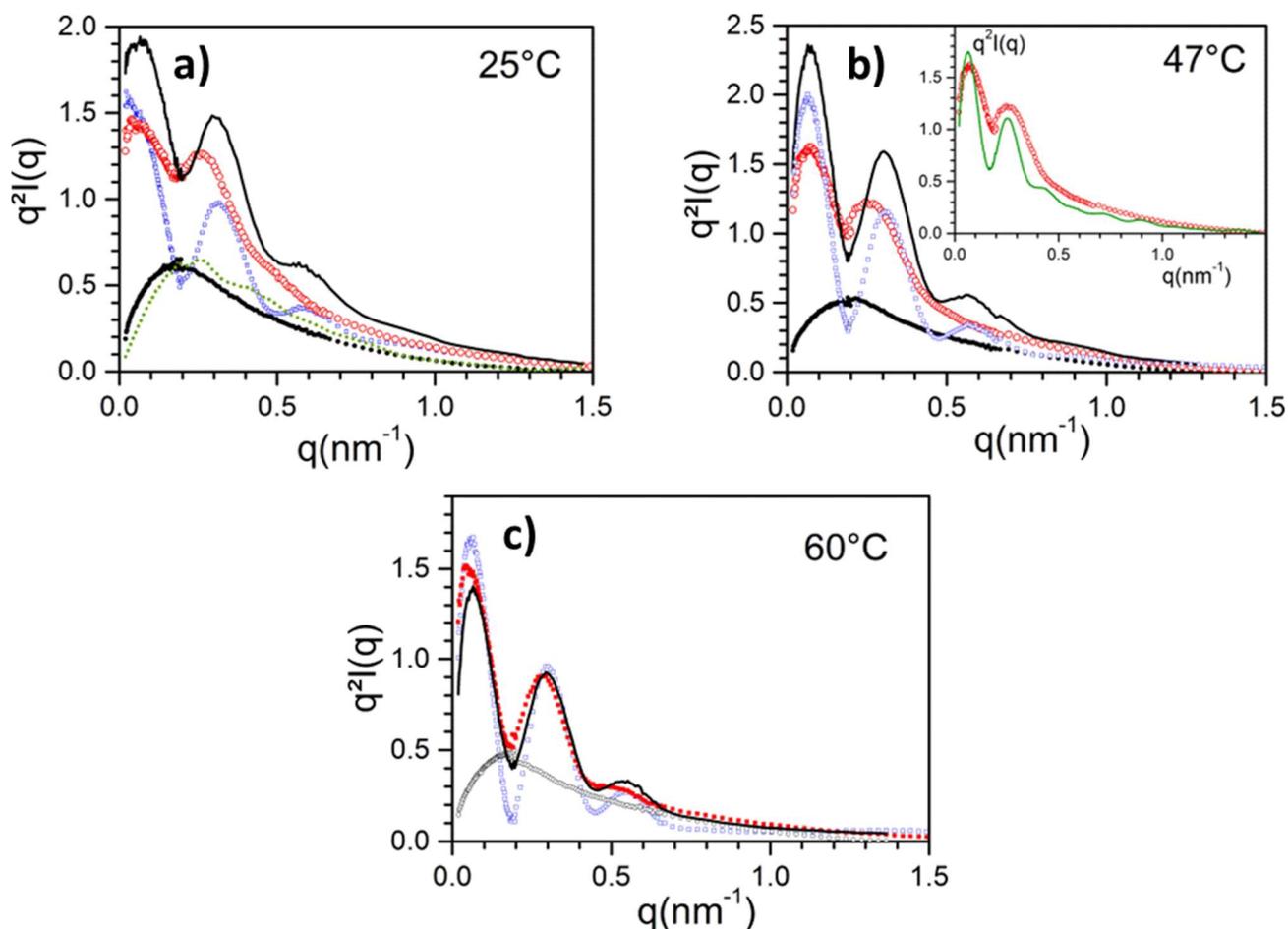


Figure 3 | Small-angle scattering curves plotted by means of a Kratky-plot ($q^2I(q)$ vs q) for the binary system P3BT/trans-decahydronaphthalene gels, $C_{\text{P3BT}} = 0.01 \text{ g/cm}^3$ (black dots); the binary system BHPB-10/trans-decahydronaphthalene gels, $C_{\text{BHPB-10}} = 0.01 \text{ g/cm}^3$ (blue open squares); the ternary system P3BT/BHPB-10/trans-decahydronaphthalene, $C_{\text{BHPB-10}} = 0.01 \text{ g/cm}^3 + C_{\text{P3BT}} = 0.01 \text{ g/cm}^3$ (red open circles); the solid curve stands for the sum of the intensities of each binary system, except at 60°C where it represents 70% of this sum. a) $T = 25^\circ\text{C}$ b) $T = 47^\circ\text{C}$, inset shows the best fit in solid green curve (see text for details) and c) $T = 60^\circ\text{C}$.

The same fits basically work at 47°C although the scattered intensity of the P3BT in the binary mixture has slightly decreased; possibly due to the melting of some of the fibrils and thermal expansion of the solvent (Fig. 3b).

The results obtained at 60°C rather support the existence of a hybrid system against the existence of intermolecular terms (Fig. 3c). As can be seen, the scattered intensity is this time the sum of the scattered intensity of each component in the binary systems scaled by a factor 0.7. This disregards the effect of intermolecular terms for the lower temperatures in the explored q -range. Clearly, a fraction of the material has vanished that is neither pure P3BT fibrils nor pure BHPB-10 nanotubes. Note that melted material,

particularly BHPB-10, scatters very little compared to the polymer fibrils and the nanotubes so that its contribution to the intensity can be ignored to a good approximation.

If one contemplates the possible occurrence of a hybrid system where BHPB-10 nanotubes sheathe P3BT fibrils, then this can be modelled with two concentric cylinders assuming in a first approximation a cylinder-like structure for the P3BT fibrils. In view of the lower scattering of P3BT fibrils with respect to BHPB-10's this is a very good approximation (see supp. Information). In other word the scattering is dominated by the BHPB-10 moiety.

This approach is definitely based on the C-AFM outcome that definitely allows one to narrow down the possible library of models. One has to keep in mind that each of these concentric cylinders possesses a different neutron contrast factors in *trans*-decahydronaphthalene (see supplementary information) so that the theoretical expression for the scattered intensity reads³²:

$$I_{\text{hybrid}}^{\text{theo}}(q) \sim \frac{\pi \mu_{\text{LHyb}}}{q} \left[\frac{2\gamma A_{\text{in}}}{A_m q r_{\text{out}}} J_1(q\gamma r_{\text{out}}) + \frac{2A_{\text{out}}}{A_m q r_{\text{out}}} \times (J_1(qr_{\text{out}}) - \gamma J_1(\gamma q r_{\text{out}})) \right]^2 \quad [3]$$

where A_{in} and A_{out} are the scattering amplitude of the inner and outer cylinder, and with

$$A_m = \gamma^2 A_{\text{in}} + (1 - \gamma^2) A_{\text{out}} \quad [4]$$

From the results at 60°C we infer that about 0.3 of this hybrid material was formed. Therefore, the final expression for the expected intensity at 25°C should read:

$$I(q) = 0.35 \times I_{\text{P3BT}}^{\text{exp}}(q) + 0.35 \times I_{\text{BHPB-10}}^{\text{exp}}(q) + 0.3 \times I_{\text{hybrid}}^{\text{theo}}(q) \quad [5]$$

where $I_{\text{P3BT}}^{\text{exp}}$ and $I_{\text{BHPB-10}}^{\text{exp}}$ are the experimental intensities measured for the binary systems.

Equal fraction of P3BT and BHPB-10 are considered since they correspond to the starting material.

In order to match the positions of the first two maximum the outer cylinder should possess an external radius of $r_{\text{out}} = 20$ nm and an internal radius, being also the radius of the inner solid cylinder, of $r_{\text{in}} = 9$ nm (figure 4b inset). This fit should be regarded indicative, although there are not so many possibilities in view of the constraints imposed by the position of the second maximum together with the relative ratios of the first two maxima.

The inner radius, $r_{\text{in}} = 9$ nm, is consistent with the half-value of the diagonal of the P3BT ribbons (7.5 nm).

Conversely, the outer radius is larger by a factor about 1.6 with respect to the outer radius of the pure

nanotubes. This gives a layer of about 11 nm, which suggests that approximately 3 layers of BHPB-10 would sheathe the P3BT fibrils since the nanotube thickness is about 3 nm. This value, 11nm, is consistent with the findings from the C-AFM observations, where the tip has to penetrate by about 8 to 11nm to establish contact with the semi-conducting P3BT. Note that the scattering function cannot differentiate between closed nanotubes and helical ribbons at this level of resolution. The sheathing shell may well be made up with one closed nanotube onto which helical ribbons further grow.

Under these conditions a molecular mixture reminiscent of an eutectic compound may be formed which can account the hybrid system melting at a temperature below that the pure BHPB-10 nanotubes²³ ($T_{\text{BHPB-10}} = 67^\circ\text{C}$ for $C = 0.01 \text{ g/cm}^3$) and of the pure P3BT fibrils ($T_{\text{P3BT}} = 100^\circ\text{C}$). The formation of an eutectic system also implies that the degree of sheathing may vary with the concentrations of each species and the composition. Determination of the ternary phase diagram would be a prerequisite for gaining a deeper knowledge on these systems.

Additional information

See supporting information for details on C-AFM experiments.

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Insulated molecular wires: sheathing semi-conducting polymers with organic nanotubes through heterogeneous nucleation

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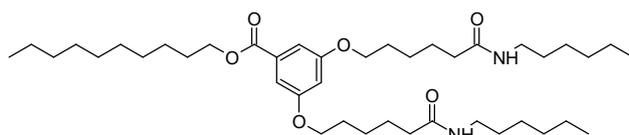
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METHODS

Samples preparation

The molecule used in this study for its propensity of producing nanotubes in some organic solvents is 3,5-Bis-(5-hexylcarbamoylpentyloxy)-benzoic acid decyl ester (abbreviated as BHPB-10 in what follows) whose chemical structure is shown below.



The synthesis of BHPB-10 is described in details in reference²⁴.

Poly(3-butylthiophene-2,5-diyl) was purchased from Merck. Its weigh-averaged molecular weight is $M_w=32$ kDa with a polydispersity $M_w/M_n= 2.0$ and a regioregularity of 77%. The amorphous fraction was removed from the suspensions of crystallized material through successive centrifugations (see supplementary information). Only the P3BT fraction of highest crystallinity was used thereafter.

The hybrid materials prepared through a four-step process (shown in Fig. 1): 1) preparing separately a suspension of P3BT fibrils at room temperature in *trans*-decahydronaphthalene , 2) preparing separately a homogeneous solution of BHPB-10 in *trans*-decahydronaphthalene at $T= 70^\circ\text{C}$, 3) mixing at 60°C the P3BT suspension and the BHPB-10 solution, 4) quenching the system at 0°C .

Mixing at 60°C allows one to keep a homogeneous solution of BHPB-10 in *trans*-decahydronaphthalene since nanotube formation occurs only at $T= 45^\circ\text{C}$, meanwhile P3BT fibrils are prevented from melting. Quenching at 0°C triggers the rapid growth of the BHPB-10 nanotubes.

Experiments

TEM

After preparation of the samples for the C-AFM experiments, the solutions are further diluted to allow for observation in TEM (typically 10 to 100 times). A drop of the resulting solution is deposited onto a carbon-

coated copper grid, and the excess of solvent is carefully removed by means of blot-paper. This avoids additional aggregation due to solvent evaporation. TEM was performed in bright field s using a Philips CM12 microscope equipped with a MVIII CCD camera (Soft Imaging System) operating at 120 kV. **C-AFM experiments** were performed using the Nanotec Cervantes AFM with the Dulcinea controller (Nanotec Electronica, Spain). Images of the fibers deposited on the grounded electrode were acquired in Tapping mode AFM under ambient conditions using PtSi tips (Nanosensors, model: PtSi-FM). Once stable and drift free image of the fibers are recorded, FZ and I-V cures were measured at several position on selected fibers with controlled Z increments. Time to record a single spectral curve is ~ 2 s with number of points at least 200. I-V curves on the grounded electrode give a straight line with a steep slope that is characteristic of a highly conducting material (Supplementary Fig. 1). New tip was used for each set of samples. Tips are then used to perform FZ curves on freshly cleaved mica in order to determine the actual tip penetration distance on the sample.

Small-angle neutron scattering. The experiments were performed on D11 camera located at the Institut Laue-Langevin (ILL, Grenoble, France). A wavelength of $\lambda_m = 0.6$ nm was used with a wavelength distribution characterized by a full width at half maximum of about 10%. A built-in two-dimensional position sensitive detector composed of 64x64 cells was used (details available on <http://www.ill.eu>). The sample-detector distance was varied (3, 10, 34 m) in order to span a large q-range ($0.05 < q \text{ (nm}^{-1}) < 2.25$) with $q = (4\pi/\lambda) \sin(\vartheta/2)$, where ϑ is the scattering angle. The position sensitive detector was calibrated with light water for correcting for cell efficiency and obtaining absolute intensity, $I_{abs}(q)$. The absolute intensity was finally obtained after the usual signal processing and using $d\Sigma/d\Omega = 0.985 \text{ cm}^{-1}$ for the water cross-section as determined experimentally for D11 at $\lambda_m = 0.6$ nm, and the contrast factor for the labelled species.

Per deuterated *trans*-decahydronaphthalene, $C_{10}D_{18}$, purchased from Cambridge Isotopes was used throughout. The scattering amplitudes for P3BT and BHPB-10 in this deuterated solvent used for equation 3 are calculated by means of:

$$A_{out} = \left[\frac{A_{BHPB-10}}{v_{BHPB-10}} - \frac{A_{TdecaD}}{v_{TdecaD}} \right] \text{ and } A_{in} = \left[\frac{A_{P3BT}}{v_{P3BT}} - \frac{A_{TdecaD}}{v_{TdecaD}} \right]$$

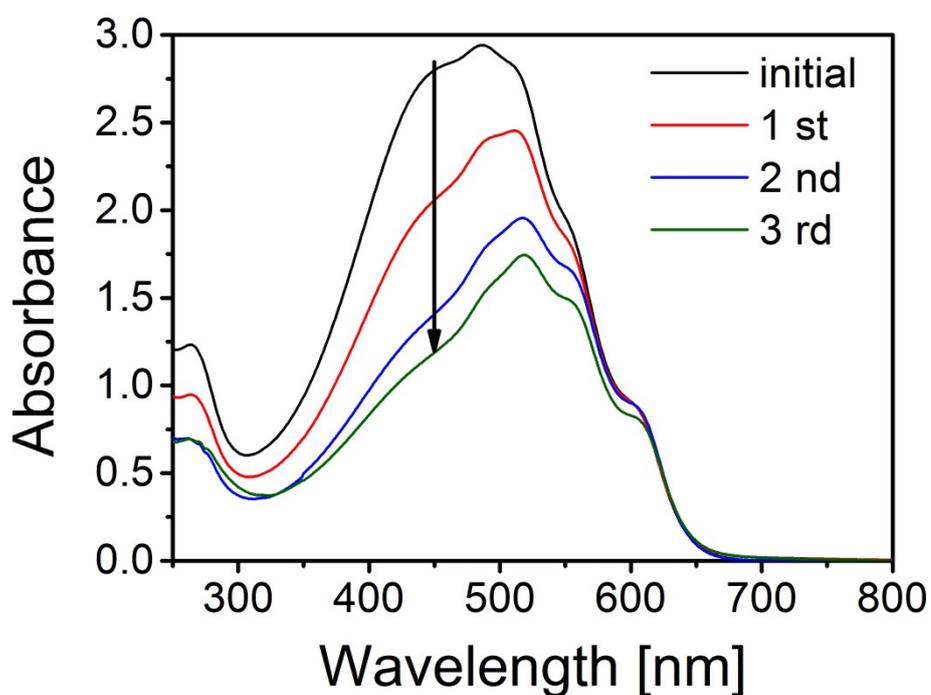
where A and v , with the appropriate subscript, are the scattering amplitudes and the molar volumes, respectively, and with $A_{\text{BHPB-10}}/v_{\text{BHPB-10}} = 0.01096$, $A_{\text{P3BT}}/v_{\text{P3BT}} = 0.0362$, and $A_{\text{TdecaD}}/v_{\text{TdecaD}} = 0.1176$. As a reminder, for neutrons A is the sum of the scattering lengths of the atoms constituting the molecule.

SUPPLEMENTARY FIGURES AND FIGURE CAPTIONS

Purification of the P3BT samples

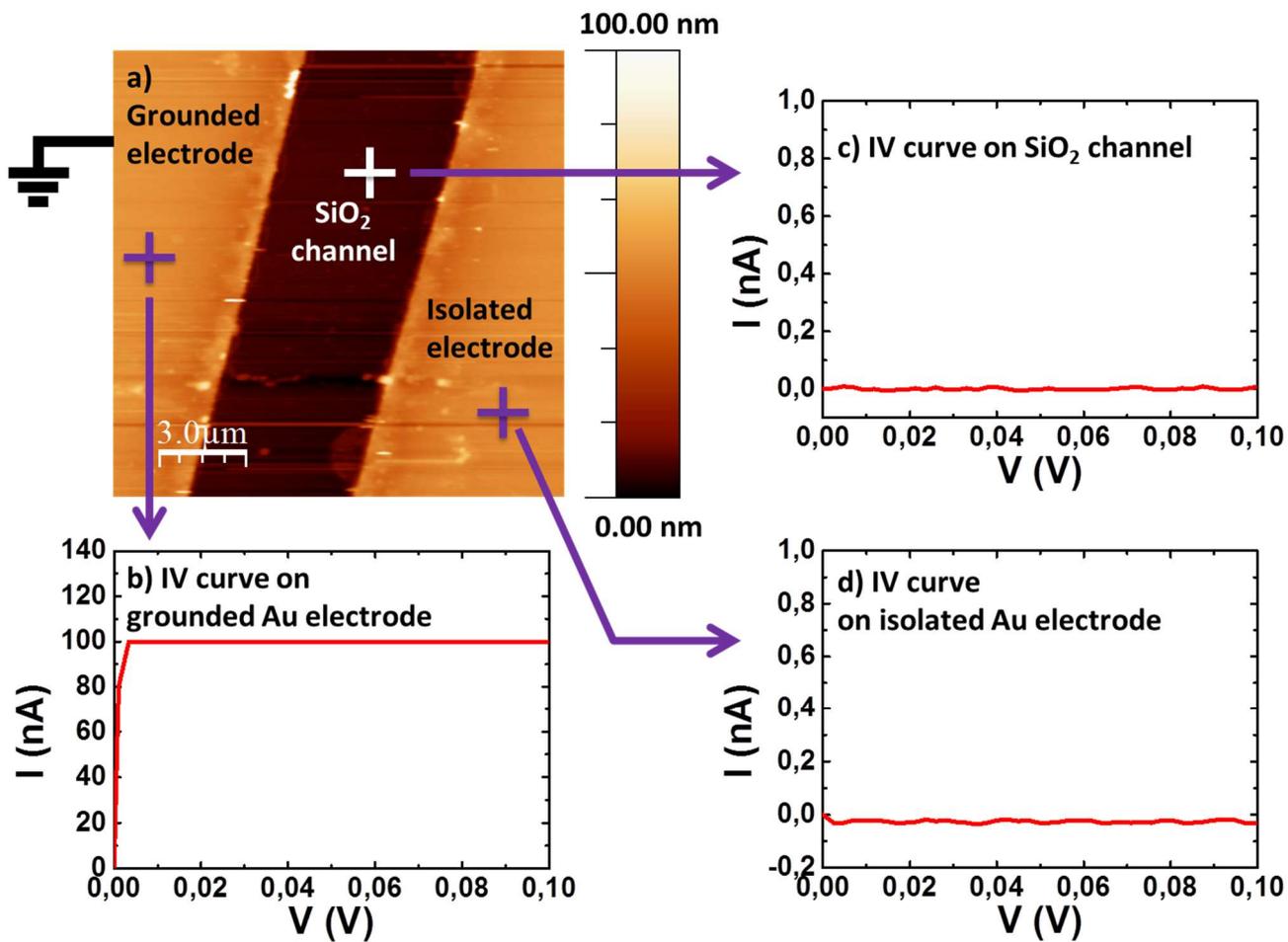
Trans-decahydronaphthalene was purified through long columns of silica gel prior to use.

The as-received P3BT was first washed in a soxhlet by acetone and hexane. Then, the P3BT chains that poorly or even do not crystallize have been removed by successive centrifugations. P3BT is first allowed to crystallize from dilute solutions in *trans*-decahydronaphthalene ($0.075 \times 10^{-2} \text{ g/cm}^3$). Centrifugation separates the P3BT fibrils from chains that have not crystallized, the latter being located in the supernatant. After an ageing time of 24H to allow for the growth of P3BT fibrils, the samples have undergone three centrifugation processes at 1000 rpm for 2 hrs. As can be seen in supplementary figure 1, where the UV-vis absorption spectra are displayed, a very high degree of crystallinity is thus obtained.



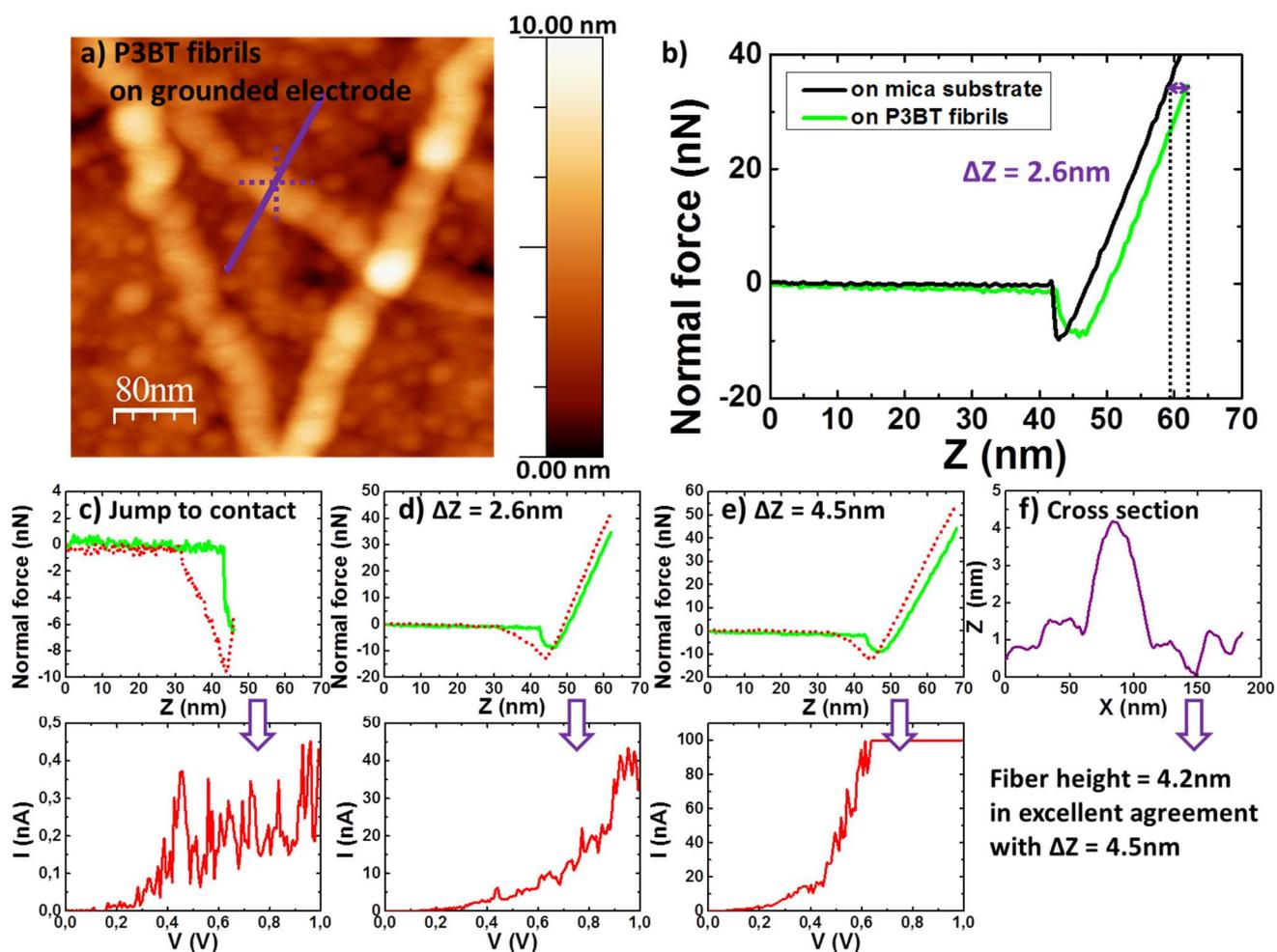
Supplementary Figure 1. UV-Vis absorption spectra after successive centrifugations of a P3BT/*trans*-decahydronaphthalene suspension. The decrease of the maximum at 450 nm indicates gradual disappearance of the non-crystallizable chains.

Supplementary Figure 2



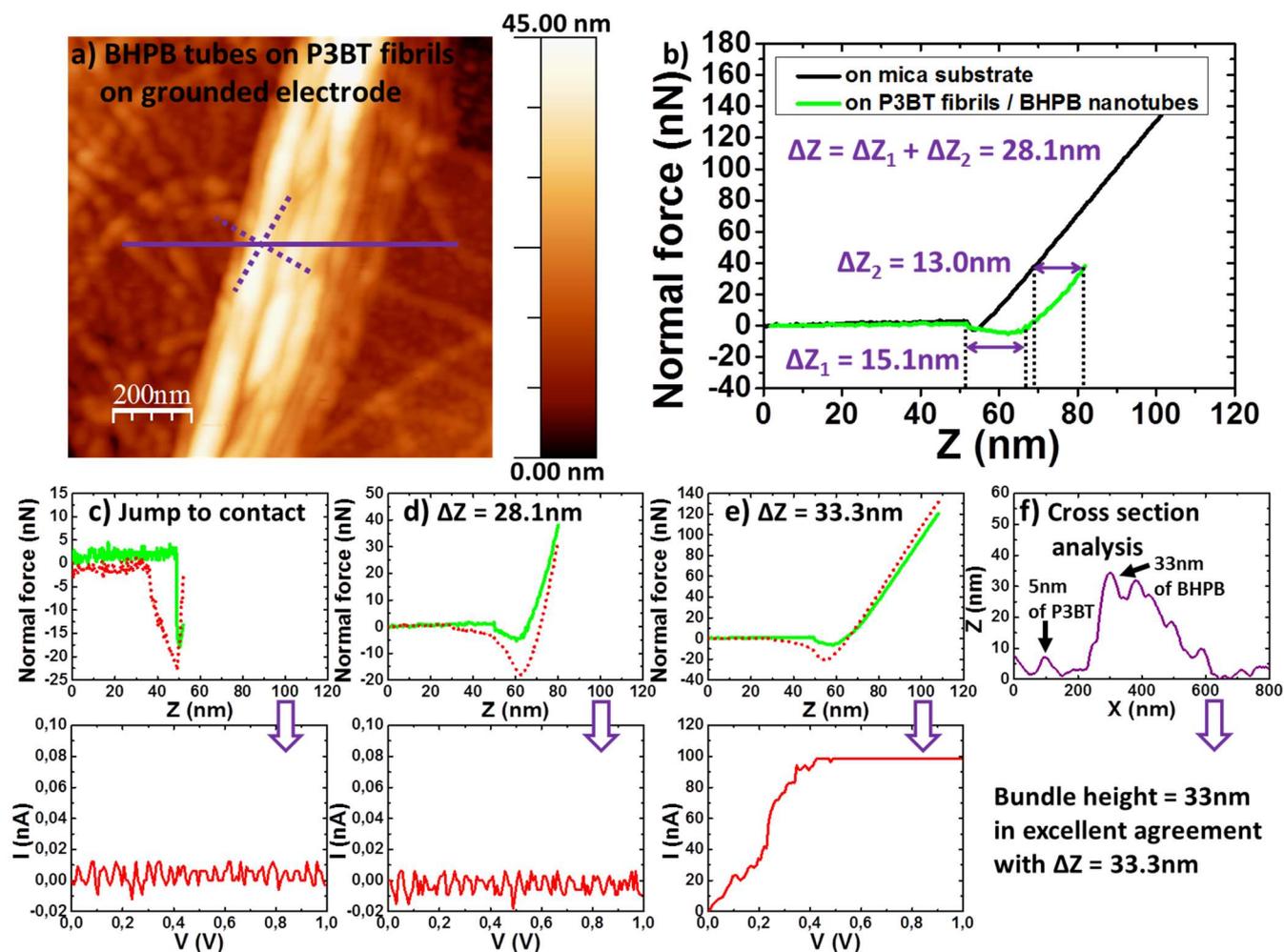
Supplementary Figure 2. IV curves on the different parts of the bare substrate. a) AFM topographic image of interdigitated Au electrodes acquired using PtSi tip in tapping mode; b) IV curves measured using the PtSi tip on the ground electrode show a linear relation with steep slope - characteristics of a conducting material; c) and d) IV curves show no current on the SiO₂ channel and the isolated electrode respectively.

Supplementary Figure 3



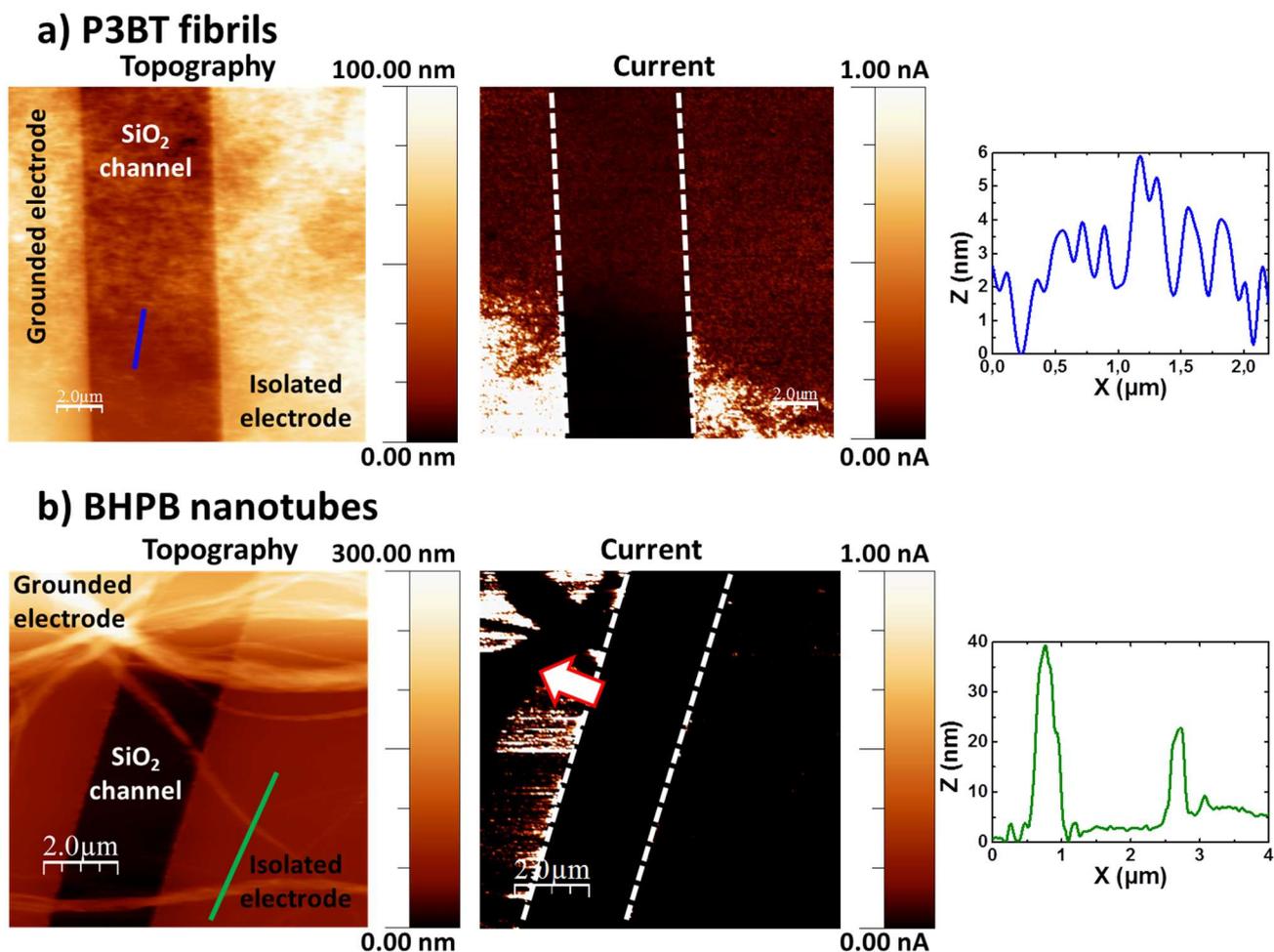
Supplementary Figure 3. Current-Voltage spectroscopy on P3BT fibrils at controlled indentation depth. a) Topography of P3BT fibrils on a grounded Au electrode; b) ΔZ calculation of P3BT fibril tip penetration by comparing FZ curves on mica; c-e) FZ curves at different tip depth positions (ΔZ) into the fiber (at the point located by the dotted purple cross) and the related I-V curves measured for each ΔZ , confirming the semiconducting character of the P3BT fibrils (non linear I-V curve); f) cross-section analysis of P3BT fibrils along the straight line in panel a.

Supplementary Figure 4



Supplementary Figure 4. Current-Voltage spectroscopy on sequentially deposited BHPB-10 fibers on P3BT fibrils at controlled indentation depth. a) Topography of BHPB-10 nanotubes on P3BT fibrils deposited on a grounded Au electrode; b) Calculation of tip penetration depth. The total tip penetration (ΔZ) here is obtained by the addition of the distance traveled by the tip through the spacing between fibrils, indicated by zero force from jump to contact until the point of linearly increasing load, (ΔZ_1), to the tip penetration distance into the fiber obtained by comparing FZ curves on mica (ΔZ_2); c-e) FZ curves at different tip depth positions (ΔZ) into the material (at the point located by the dotted purple cross) and the related I-V curves measured for each ΔZ , confirming the fact that P3BT and BHPB-10 form separated entities under these conditions; f) cross-section analysis of P3BT fibrils and BHPB-10 nanotubes along the straight line in panel a.

Supplementary Figure 5.



Supplementary Figure 5. AFM topography and C-AFM current mapping (see Fig. 3a for the experiment principle):

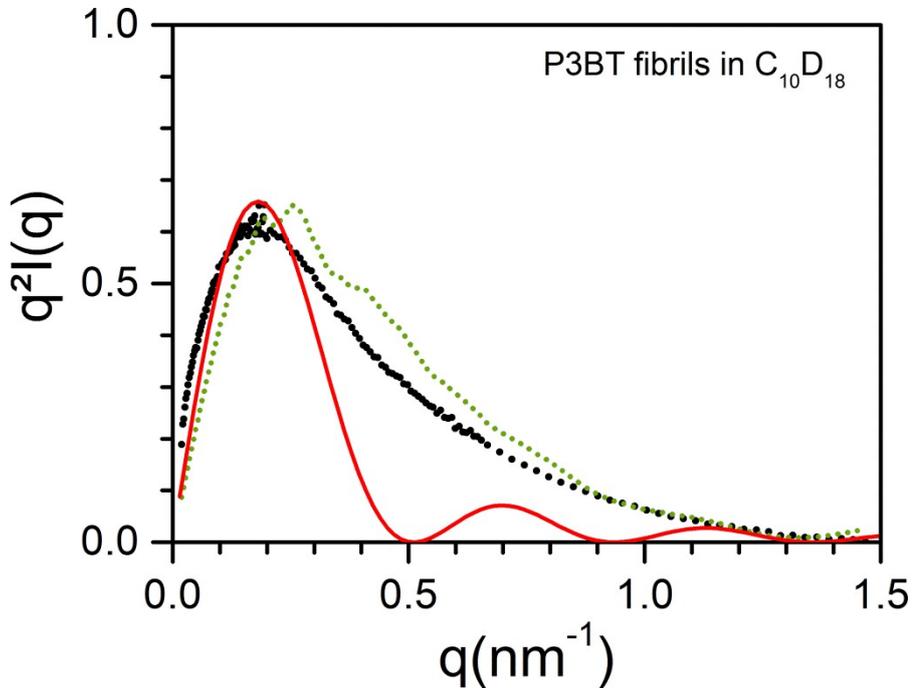
a) on P3BT fibrils: the current flows between both electrodes and channel indicating the conducting nature of P3BT;

b) on BHPB-10 nanotubes: the current flows only at the grounded electrode. No current appears in the channel and the isolated electrode. Red arrow indicates dark areas that correspond to BHPB-10 nanotubes on the ground electrode where no current is recorded due to the insulating nature of BHPB-10.

Cross-section analysis along the straight line in the topography panel are shown besides each image.

Neutron experiments

A better fit for the P3BT fibrils is obtained by considering a rectangular cross-section. Considering a cylinder cross-section fits the low q domain but fails at larger q .



Supplementary Figure 6. Scattering by P3BT fibrils in deuterated *trans*-decahydronaphthalene. Comparison between a fit with a solid cylinder of radius 7.5 nm (red curve) and a ribbon of rectangular cross-section $a=4$ nm and $b=14$ nm. Note that 7.5 nm is the value of the half-diagonal of such a rectangle.

Neutron contrast

The contrast of P3BT with respect to deuterated *trans*-decahydronaphthalene (TDecaD) is:

$$A_{P3BT}^2 = \left(B_{P3BT} - \frac{v_{P3BT}}{V_{TDecaD}} B_{TDecaD} \right)^2$$

Where v with the appropriate subscripts are the molar volume of deuterated *trans*-decahydronaphthalene and P3BT respectively,

The scattering amplitude are calculated by summing the scattering amplitudes a_i of all the atoms composing the molecules, namely

$$B_{P3BT} = 8a_C + 8a_H + a_S$$

$$B_{TDecaD} = 10a_C + 18a_D$$

Similarly, for BHPB-10

$$A_{BHPB-}^2 = \left(B_{BHPB-10} - \frac{v_{BHPB-10}}{V_{TDecaD}} B_{TDecaD} \right)^2$$

With

$$B_{BHPB-} = 39a_C + 68a_H + 6a_O + 2a_N$$

$$a_C = 0.662, a_H = -0.375, a_D = 0.670, a_O = 0.575, a_N = 0.94, a_S = 0.285$$

$$d_{P3BT} = 1.24 \text{ g/cm}^3, d_{BHPB-10} = 1.1 \text{ g/cm}^3$$

$$A_{P3BT}^2 = 113 \times 10^{-24} \text{ cm}^2$$

$$A_{BHPB-10}^2 = 4453 \times 10^{-24} \text{ cm}^2$$