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# Beamforming Through Regularized Inverse Problems in Ultrasound Medical Imaging

Teodora Szasz, Adrian Basarab, and Denis Kouamé, Member, IEEE

Abstract-Beamforming (BF) in ultrasound (US) imaging has significant impact on the quality of the final image, controlling its resolution and contrast. Despite its low spatial resolution and contrast, delay-and-sum (DAS) is still extensively used nowadays in clinical applications, due to its real-time capabilities. The most common alternatives are minimum variance (MV) method and its variants, which overcome the drawbacks of DAS, at the cost of higher computational complexity that limits its utilization in real-time applications. In this paper, we propose to perform BF in US imaging through a regularized inverse problem based on a linear model relating the reflected echoes to the signal to be recovered. Our approach presents two major advantages: 1) its flexibility in the choice of statistical assumptions on the signal to be beamformed (Laplacian and Gaussian statistics are tested herein) and 2) its robustness to a reduced number of pulse emissions. The proposed framework is flexible and allows for choosing the right tradeoff between noise suppression and sharpness of the resulted image. We illustrate the performance of our approach on both simulated and experimental data, with in vivo examples of carotid and thyroid. Compared with DAS, MV, and two other recently published BF techniques, our method offers better spatial resolution, respectively contrast, when using Laplacian and Gaussian priors.

Index Terms—Adaptive beamforming (BF), basis pursuit (BP), beamspace processing, least squares (LS), linear inverse problems.

#### I. INTRODUCTION

LTRASOUND (US) imaging is one of the most fast-developing medical imaging techniques, allowing noninvasive and ultrahigh frame rate procedures at reduced costs. Cardiac, abdominal, fetal, and breast imaging are some of the applications where it is extensively used as diagnostic tool. The new advances in beam formation, signal processing, and image display enlarge the US imaging potential to other fields like brain surgery, or skin imaging (see [1], [2]).

In a classical US scanning process, short acoustic pulses are transmitted through the region of interest of the human body. The backscattered echo signals, also called raw radiofrequency (RF) data, are then processed for creating RF beamformed lines. Beamforming (BF) plays a key role in US image

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- T. Szasz and A. Basarab are with the Centre national de la recherche scientifique, Institut de Recherche en Informatique de Toulouse, University of Toulouse, 31062 Toulouse, France (e-mail: szasz@irit.fr; basarab@irit.fr).
- D. Kouamé is with the Centre national de la recherche scientifique, Institut de Recherche en Informatique de Toulouse, University of Toulouse 3, 31062 Toulouse, France (e-mail: kouame@irit.fr).

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formation, influencing the resolution and the contrast of final image. The most used BF method is the standard delay-andsum (DAS), which consists in delaying and weighting the reflected echoes before averaging them. So far, its simplicity and real-time capabilities make it extensively used in US scanners. However, its weights are independent on the echo signals, resulting in beamformed signals with a wide mainlobe width and high sidelobe level. Consequently, the resolution and the contrast of final image are relatively low [3]. Several adaptive beamformers (with weights dependent on data) from array processing literature were applied to US, the most common being the Capon or minimum variance (MV) BF [4]. It offers a very good interference rejection and better resolution than DAS, allowing higher contrast [5]. However, this method uses an estimated covariance matrix of the data and its main issue is the high computation complexity due to the calculation of the inverse covariance matrix. To overcome this, many improved versions of MV have been recently proposed (see [6], [7]), but still not adequate for real-time applications. In practice, in order to provide well-conditioned covariance matrices, diagonal loading, time, and spatial averaging approaches were investigated (see [8], [9]).

Recently, to improve the MV BF, Nilsen and Hafizovic [10] proposed a beamspace adaptive beamformer, BS-Capon, and unlike MV BF, they based their BF method on orthogonal beams formed in different directions. This technique was also applied by Jensen and Austeng [11] to develop an adaptive beamformer based on multibeam covariance matrices, called multibeam Capon beamformer. In their works, a covariance matrix is calculated for each range in the image, based on the idea that the beams were transmitted with different angles. Thus, Jensen and Austeng [11] were able to reduce the computation time of MV BF, while improving the resolution of the point-like reflectors. Following a similar idea, Jensen and Austeng [12] applied to US imaging a method initially proposed in [13], called the iterative adaptive approach (IAA). They obtained better defined cyst-like structures compared with conventional DAS and better rendering than MV.

The work presented here uses a similar idea of BF range by range. However, inspired from the source localization problems, we represent, for each range, the BF as a linear direct model relating the raw samples to the desired lateral profile of the RF image to be beamformed. This formalism allows us to invert the problem by imposing standard regularizations such as the  $\ell_1$ - or  $\ell_2$ -norm. These choices are motivated by the existing works in US image enhancement, which are typically based on Laplacian (see [14]) or Gaussian (see [15]) priors.

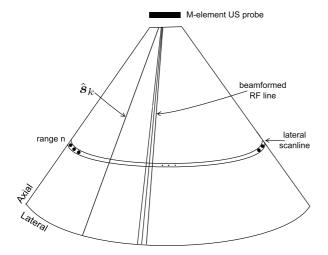


Fig. 1. Main US imaging elements used to adapt the array processing BF methods discussed in Sections II and III to medical US imaging.

Thus, the major contribution of this paper is the improvement of the existing BF techniques by combining the proposed direct model formulated in the lateral direction of the images with a regularized inversion approach. Moreover, we incorporated the proposed method with a beamspace processing technique, in order to highly reduce the number of the required US emissions.

In contrast to existing BF methods in US imaging using regularized inverse problem approaches (see [16]–[21]), our method does not use the system point spread function (PSF) in the direct model or in the inversion process. Thus, the proposed BF technique, similar to [22] applied to nondestructive evaluation, does not require any experimental measurement (see [23]) or estimation of the PSF (see [14], [24], [25]).

Laplacian and Gaussian statistics, two of the most common regularizations in such imaging problems (including US imaging applications such as deconvolution), are considered herein, allowing the reader to observe their influence on the results. Furthermore, our method opens new tracks for more complex regularization terms (see [26]–[28]) to further improve the results. The proposed approaches, generically named basis pursuit BF (BP BF) and least squares (LS) BF in this paper, were evaluated using different Field II simulated data and *in vivo* carotid and thyroid experimental data. Finally, we compared our BF techniques with four existing beamformers: the conventional DAS, MV, multibeam Capon, and IAA.

The remainder of this paper is organized as follows. First, in Section II, we summarize the theoretical background of BF. In Section III, we describe the proposed BF method from a regularized inverse problem perspective. Details about the experiments and the results are given in Section V, and Section VI concludes this paper.

#### II. BACKGROUND

The main elements used to model the BF process are depicted in Fig. 1. We consider, without loss of generality, the particular setup of an M-element US probe (M can also be the number of active elements of the probe), with the transducer's

elements denoted by  $u_m$ , with  $m = 0, \dots M - 1$ . We consider a trivial change of variable, such that the position of the mth element is

$$p_m = [m - (M-1)/2](\lambda/2), \quad m = 0, 1, \dots, M-1 \quad (1)$$

with the probe's elements positioned symmetrically around the origin. We have considered, as example, the pitch equal to  $\lambda/2$  (the spacing between elements is half of the wavelength  $\lambda = c/f_0$ , c denoting the speed of the sound through soft tissue and  $f_0$  the transducer's center frequency).

A series of K focused beams is transmitted with different incident angles  $\theta_k$ ,  $k=1,\ldots,K$ . The returning echoes are recorded using the same US probe, being time delayed, such that the time of flight is compensated for, so the backscatter from the point of interest is summed up coherently. If we consider that each of the recorded raw signal after the time-delay compensation has N time samples, the size of the recorded data from all the K directions will be  $M \times N \times K$ . Finally, the final RF US image is a collection of RF beamformed lines, each of which being the result of BF the raw RF signals coming from an emission in the direction  $\theta_k$ ,  $k \in \{1, \ldots, K\}$ , using M elements of the transducer.

The classical DAS BF can be expressed as

$$\hat{s}_{k}(n) = \sum_{m=1}^{M} w_{m} y_{m}^{(k)}(n - \Delta_{m}(n)), \quad n = 1, \dots, N$$

$$k = 1, \dots, K \quad (2)$$

where  $\Delta_m(n)$  is the time delay for focusing at the point of interest sample, being dependent on the distance between the mth element and the point of interest,  $w_m$  are the BF weights, and  $y_m^{(k)} \in \mathbb{C}^{N \times 1}$  is the raw data received by the mth element of the US probe, corresponding to the emission steered at angle  $\theta_k$ . A simplified form of (2) can be formulated as

$$\hat{\mathbf{s}}_k = \mathbf{w}^H \mathbf{y}_k \tag{3}$$

where  $y_k \in \mathbb{C}^{M \times N}$  is the time-compensated version of  $y_m^{(k)}$  in (2) for the kth emission (for the sake of generality, we consider  $y_k$  to be complex-valued data), w is the vector of the beamformer weights of size  $M \times 1$ , and  $(\cdot)^H$  represents the conjugate transpose. DAS BF selects the weights independent on data, solving

$$\min_{\mathbf{w}} \mathbf{w}^H \mathbf{w}, \quad \text{such that } \mathbf{w}^H \mathbf{1} = 1 \tag{4}$$

where 1 is a length M column vector of ones since the raw data were focused using time delays. The solution of (4) is

$$w_{\rm DAS} = \frac{1}{M}.$$
 (5)

If we replace (5) in (3), we get

$$\hat{\mathbf{s}}_k = \frac{1}{M} \mathbf{1}^T \mathbf{y}_k \tag{6}$$

where  $\{\cdot\}^T$  denotes the transpose. A common technique used in US is to apply weighting functions such as Hanning or Hamming apodizations to (6) to further reduce the sidelobes of  $\hat{s}_k$ , resulting in improved contrast of the beamformed image, at the cost of a slight lateral spatial degradation.

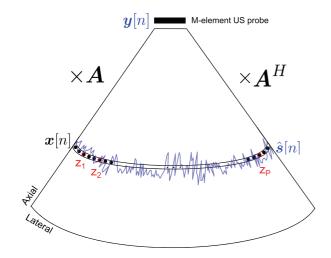


Fig. 2. Elements used to form the proposed model.

Further details about adaptive BF in US imaging (i.e., the methods used for comparison purposes) and about beamspace processing are provided in Appendixes A–D.

## III. PROPOSED METHOD: BEAMFORMING THROUGH REGULARIZED INVERSE PROBLEMS

#### A. Model Formulation

The main elements used to model the proposed method are depicted in Fig. 2. For the sake of simplicity, let us focus our problem at a time sample (range) n. The proposed BF method is sequentially applied in the same manner to each range. If  $y_k[n] \in \mathbb{C}^{M \times 1}$  is the raw data after the compensation of the time of flight for the kth steering direction  $\theta_k$ , we can form the steering vectors as in (26). Let A be the  $M \times K$  steering matrix containing the steering vectors in (26) for all  $\theta_k$  directions,  $k = 1, \ldots, K$ 

$$A = [a_{\theta_1}, a_{\theta_2}, \dots, a_{\theta_K}]. \tag{7}$$

Note that A is known and depends on the positions of the probe elements and on the locations to beamform. Thus, it is independent on the actual positions of the reflectors.

For each range n, we want to estimate the signal corresponding to a reflector as a function of its location that will contain dominant peaks at reflector positions. Thus, the main difference from the multibeam Capon BF method is that instead of calculating the values of the weights  $\mathbf{w}_{\theta n}$  as in (28), that are further used to calculate the reflector's signal, we are directly estimating the corresponding signal by considering the raw data  $y_{k}[n]$  as observations. In other words, we want to obtain an estimate of the reflected echo  $x[n] \in \mathbb{C}^{K \times 1}$ through the observations  $y_k[n]$ . Unfortunately, one difficulty arises: since  $y_k[n]$  is corresponding to only one emission, modeling our problem using raw data as observations to estimate the reflectors requires high computational cost, since we are dealing with multiple directions. We recall that the size of raw data in our problem at a range n is  $M \times K$ . To overcome this issue, motivated by the results in [29] and [30], we propose to use the DAS beamformed data instead of the original

raw data. In addition to data dimensionality reduction, it was shown in [29] and [30] that more accurate results may be achieved by proceeding in this way. Thus, we can formulate our model as follows:

$$\hat{\mathbf{s}}[n] = (\mathbf{A}^H \mathbf{A}) \mathbf{x}[n] + \mathbf{g}[n] \tag{8}$$

where  $\hat{s}[n] \in \mathbb{C}^{K \times 1}$  is a lateral scanline of the DAS beamformed image formed as discussed in Section II, A is the steering matrix formed with (7), and g[n] is an additive white Gaussian noise. If we denote by  $\hat{S}$  the DAS beamformed image of size  $K \times N$ , formed by juxtaposing the DAS RF lines  $\hat{s}_k$  expressed in (6) for all K directions, we consider  $\hat{s}[n]$  the lateral scanline from  $\hat{S}$  taken at the time sample (range) n. Note that, since the transducers are emitting the same pulse, we assumed that x[n], which is the unknown signal, is the same for all K emissions and for all transmitters (e.g., [31]).

The role of the multiplication of the steering matrix A with its conjugate transpose  $A^H$  in (8) is to relate the position of the elements with the position of all K reflectors on a scanline. This relation is a result of considering on the one hand that the elements are impinging to the reflectors situated on the lateral scanline (the multiplication of A with x[n]), while on the other hand, the reflectors are impinging to the elements through their reflected pulses (the multiplication with  $A^H$ ). Hence, the result of the DAS beamformed scanline  $\hat{s}[n]$  is related to the unknown signal x[n] through a direct linear model. Fig. 2 offers a schematic of our model in (8). Thus, after the compensation of the time of flight, the received raw data y[n] at a range n are formed by multiplying the steering matrix A with the desired signal x[n], y[n] = Ax[n]. This multiplication could be sufficient for describing the proposed model if we are considering the raw data y[n], as observations. However, since we are using the DAS beamformed data instead of the original raw data, we further take into account the geometrical relationship between the potential sources and the elements of the probe (through the multiplication with  $A^H$ ).

#### B. Beamspace Processing

In order to solve (8), we first apply beamspace processing, a common tool used in source localization approaches that reduces the computational complexity, while improving the resolution and reducing the sensitivity to the position of the sensor (see [30], [32]). Its main purpose in US is to reduce the number of the US emissions, thus reducing the acquisition time and the computational load required by the BF process. We should note that our method of transforming the data into beamspace domain is totally different from the technique resumed in Appendix B. The main reason is that, using the beamspace processing presented in [10], we need all the acquired raw data for applying beamspace processing as described in (22). Hence, even if on the one hand, the computational complexity required by the estimated covariance matrix inversion is reduced, on the other hand, it is increased by the operations required to transform the entire set of the raw data into its beamspaced correspondents.

To overcome this, we based our idea on the beamspace processing techniques proposed in array processing

localization). (notably in source More specifically, Malioutov et al. [29] used a method that maps the data from full dimension space of the directions (DS) into a lower dimension beamspace (BS) through a linear transform prior to source localization processing. In our case, for each range n, we project the data resulted by applying DAS BF,  $\hat{s}[n]$ , in BS before BF it through regularized inverse problems. To emphasize,  $\hat{s}[n] \in \mathbb{C}^{K \times 1}$  is projected on a lateral sampled grid of  $P \ll K$  locations. In other words, the proposed BF method, contrarily to all the other discussed BF methods, uses only P focused emitted beams among all the K transmissions to beamform a particular lateral scanline of K samples. Thus, the number of emissions is reduced by a factor of (K/P). This will result in a reduced dimensionality of the data compared with the other BF methods and an improved computational complexity compared with MV, multibeam Capon, and IAA.

Let  $z[n] \in \mathbb{C}^{P \times 1}$  be the beamspace transformed vector formed by sampling the DAS beamformed lateral scanline  $\hat{s}[n]$  on a grid of P locations (see Fig. 2)

$$z[n] = \boldsymbol{D}^H \hat{\boldsymbol{s}}[n] \tag{9}$$

where D of size  $K \times P$  is the beamspace decimation matrix, which will reduce the dimensionality of a vector from  $K \times 1$  to  $P \times 1$ . Hence, since the decimation factor is (K/P), D has all elements zero, except the elements  $d_{i,j}$  with j = (K/P)i, which will get the value 1. Similarly, the beamspaced steering matrix  $A_{\rm BS}^H$  of size  $P \times M$  is formed, composed of the P beamspaced steering vectors

$$A_{\rm BS}^H = \boldsymbol{D}^H A^H. \tag{10}$$

Concretely, we form  $A_{\mathrm{BS}} \in \mathbb{C}^{M \times P}$  by taking from  $A^H$  each (K/P)th steering vector. Thus, the model formed by (8) after applying beamspace processing becomes

$$z[n] = \boldsymbol{D}^{H}\hat{\boldsymbol{s}}[n] = (\boldsymbol{A}_{\mathrm{BS}}^{H}\boldsymbol{A})\boldsymbol{x}[n] + \boldsymbol{D}^{H}\boldsymbol{g}[n]$$
 (11)

where x[n] of size  $K \times 1$  is the lateral profile at range n to be estimated. Thus, we can see (11) as an inverse problem, where z[n] is the DAS beamformed data corresponding to P < K emissions, and considered as the observation data.

#### C. Beamforming Through Regularized Inverse Problems

Given the ill-posedness of the direct model in (11), we propose hereafter to invert it using standard regularization techniques. For achieving this, a cost function, denoted by J(x[n]), consisting into the linear combination of two terms is considered. The first term, denoted by  $J_1(x[n])$ , represents the data attachment, while the second, denoted by  $J_2(x[n])$ , is the regularization prior

$$J(x[n]) = J_1(x[n]) + \lambda J_2(x[n])$$
 (12)

where  $\lambda$  is a scalar, called regularization parameter, that adjusts the tradeoff between the fidelity to the data and the regularization term. A large  $\lambda$  will strongly favor the *a priori* about x[n], while a small  $\lambda$  gives a high confidence to the observations. Keeping in mind that the additional noise in (11)

is Gaussian, the data attachment term is expressed by an  $\ell_2$ -norm, giving the following cost function:

$$J(x[n]) = \|z[n] - (A_{BS}^{H}A)x[n]\|_{2}^{2} + \lambda J_{2}(x[n]).$$
 (13)

In this paper, the choice of the regularization term  $J_2(x[n])$  is guided by the existing literature on statistical modeling of US images, previously applied to various applications such as image deconvolution or segmentation (see [14], [33]). It has thus been shown that Laplacian and Gaussian statistics are well adapted to model US images. For this reason, we give in the two following paragraphs the mathematical derivations and BF results using  $\ell_1$  norm (the sum of absolute difference) and  $\ell_2$  norm (or the Euclidean norm, that is, the sum of squared difference) regularization terms. We should note that while the first will promote sparse solutions, the latter will promote smoother results.

The choice of a quadratic data fidelity term is related to the additive zero-mean Gaussian assumption on the noise. We emphasize that the noise considered in our paper is different from the multiplicative speckle noise generally assumed to affect envelope images in US imaging. Instead, the additive noise considered in our paper affects the RF data and is caused by the acquisition process. The same model has been previously used by several authors (see [14], [31]).

1) Laplacian Statistics Through Basis Pursuit: Considering that the signal x[n] to be beamformed follows Laplacian statistics, the minimization of the cost function in (13) turns into the optimization procedure in the following, usually called BP in the literature [34]:

$$x_{BP}[n] = \underset{x[n]}{\operatorname{argmin}} \left( \| z[n] - \left( A_{BS}^{H} A \right) x[n] \|_{2}^{2} + \lambda \| x[n] \|_{1} \right)$$
(14)

where  $\|\cdot\|_1$  denotes the  $l_1$  norm. The minimization problem (14) is convex, hence continuous, and has one global minimum for any  $\lambda > 0$ .

Herein, we used the well-known YALL1 to solve (14) [35], a software package that contains implementation of alternating direction method that also solves BP. A comparison of the six mostly used BP implementations is done in [36], and three of them were also compared in [37] with application in underwater acoustics, where it is shown that YALL1 gives the best performances for real-time applications.

2) Gaussian Statistics Through Least Squares: To achieve smooth solutions of the proposed BF method, we modeled our problem with an  $\ell_2$ -norm-based minimization function, and we used the Tikhonov regularized least-squared method (or rigid regression) for solving it [38]. The cost function is of the form

$$\mathbf{x}_{LS}[n] = \underset{\mathbf{x}[n]}{\operatorname{argmin}} \left( \|\mathbf{z}[n] - \left(\mathbf{A}_{BS}^{H} \mathbf{A}\right) \mathbf{x}[n] \|_{2}^{2} + \lambda \|\mathbf{x}[n]\|_{2}^{2} \right)$$
 (15)

where  $\|\cdot\|_2$  denotes the  $\ell_2$ -norm. For solving (15), we used its analytical solution

$$\mathbf{x}_{LS}[n] = \left( \left( \mathbf{A}_{BS}^{H} \mathbf{A} \right)^{H} \left( \mathbf{A}_{BS}^{H} \mathbf{A} \right) + \lambda \mathbf{I}_{K \times K} \right)^{-1} \left( \mathbf{A}_{BS}^{H} \mathbf{A} \right)^{H} \mathbf{z}[n]$$
(16)

where  $I_{K\times K}$  denotes the identity matrix of size  $K\times K$ .

In order to obtain the BP and LS beamformed images, for each time sample n, with  $n \in \{1, ..., N\}$ , we estimate its corresponding lateral scanline  $x_{BP}[n]$  (using BP BF method) or  $x_{LS}[n]$  (using LS BF method), and we are juxtaposing all the obtained scanlines, in the axial direction of the image.

#### IV. EXPERIMENTS

For evaluating the proposed BP BF and LS BF approaches, we considered different types of simulated and experimental data. We compared our BF results with DAS (Section II), MV (Appendix A), multibeam Capon (Appendix C), and IAA (Appendix D) BF methods. The simulations were made using the Field II program (see [39], [40]). The first simulation includes a sparse medium, the second one contains a circular hypoechoic cyst in a medium with speckle, and the third one contains a simulation of the short axis (SAx) view of a cardiac image, as suggested in [41]. The first experimental data consist in a carotid that was recorded with an Ultrasonix MDP research platform. Finally, the second experimental data contain a thyroid medium with a malignant tumor. The thyroid data were recorded with the Sonoline Elegra clinical scanner, modified for research purpose. The parameters of the simulated and experimental data are presented in Table I. Note that for MV and multibeam Capon beamformers, spatial and temporal averaging, as well as diagonal loading technique are used for the estimation of the covariance matrix, as discussed in Section VI.

An important aspect is that, when applying the proposed BF methods, five times [(K/P) = 5] less emissions were used in the BF process, by applying the beamspace processing presented in Section III-B. This hangs on for all the examples we are presenting in this paper. For these examples, reducing five times the US transmissions is optimal in terms of gain in resolution, while reducing computational time.

The values of the regularization parameter  $\lambda$  for all the presented examples are grouped in Table I. The optimal value of  $\lambda$  was chosen manually. We emphasize that this was the case for all hyperparameters of all comparative methods. Several studies exist in the literature for automatic estimation of the regularization hyperparameter (see [42]–[44]) that can improve the robustness of the proposed methods, at increased computational cost. Nonetheless, their implementation is beyond the scope of this paper.

#### A. Parameters for the Comparative Methods

The results of MV BF were obtained using the implementation described in [9]. The length of the spatial averaging window, L, was defined as half of the number of the probe's elements, i.e., L=(M/8). A temporal window of ten samples was used in our examples, and the diagonal loading parameter was fixed to  $\Delta=(1/10L)$ . The adaptive coherence method was applied to the MV BF method.

The results of multibeam Capon were obtained using the multibeam approach suggested in [11]. The K emissions were uniformly distributed between  $\pm 30^{\circ}$ . The beamspace transform down to 33 dimensions was applied, able to retain

the variance for incoming narrowband far-field signals. The diagonal loading factor was set to 0.01.

Finally, for IAA implementation, we used the source code provided in [12]. The number of iterations was set to 15 for our examples.

Note that for all the comparative methods, several parameters need to be carefully tuned in order to obtain acceptable results. However, using the proposed approach, only the regularization hyperparameter  $\lambda$  needs to be set.

#### B. Simulated Point Reflectors

The medium contains five point reflectors, four of them aligned in pairs of two and separated by 4 mm, and the other laterally centered at 0 mm. They are located at axial depths ranging from 63 to 68 mm, with a transmit focus at 65 mm and a dynamic receive focalization.

#### C. Simulated Phantom Data

To evaluate the accuracy, contrast and resolution of the aforementioned beamformers, a hypoechoic cyst of radius 5 mm, located at the depth 80 mm, in a speckle pattern. The speckle pattern contains 50 000 randomly placed scatterers, with Gaussian-distributed amplitudes. This example was inspired from the simulation of a synthetic kidney example included in Field II software. The attenuation was not taken into account.

#### D. Simulated Cardiac Image

The SAx view is the cross-sectional view of the heart and is a well-exploited perspective in echocardiography, containing information about the left ventricle (LV) and right ventricle. In our simulation, we visualize the LV. The transmit focus point is set at 65 mm. The final image is ultrarealistic, the amplitudes being related to an *in vivo* cardiac image [41]. The number of scatterers was sufficiently large to produce fully developed speckle.

#### E. In Vivo Data: Carotid

The carotid US is a common procedure used to detect strokes or the risk of strokes due to the narrowing of the carotid arteries. The data were acquired from a healthy subject, with the Ultrasonix MDP research platform, attached with the parallel channel acquisition system, SonixDAQ. The linear ultrasonic probe L14-W/60 Prosonic (Korea) of 128 elements was used.

#### F. In Vivo Data: Thyroid

The thyroid US is done to visualize the thyroid gland to detect possible tumors or deformations. Two sets of data were acquired: first one, from a subject with a tumor and the other one from a healthy subject. For both acquisitions, we used the clinical Sonoline Elegra US system modified for research purposes, and a 7.5L40 P/N 5260281-L0850 linear array transducer of Siemens Medical Systems, having the characteristics described in Table I.

TABLE I
PARAMETERS OF SIMULATED AND EXPERIMENTAL IMAGES

Parameters for simulation of:	Point scatterers Fig. 3	Cyst simulation <b>Fig. 5</b>	Simulation of cardiac Image Fig. 9	Experimental carotid <b>Fig. 10</b>	In vivo thyroid Fig. 12
		Trans	ducer		
Transducer type			Linear array		
Transducer element pitch $[\mu m]$	256	256	192.5	110	120
Transducer element kerf $[\mu m]$	20	20	38.5	25	
Transducer element height [mm]	5	5	14	4	
Central frequency, $f_0$ [MHz]	3	3	4	7	7.2
Sampling frequency, $f_s$ [MHz]	100	100	40	40	40
Speed of sound, c [m/s]			1540		
Wavelength [ $\mu$ m]	513.3	513.3	385	220	
Excitation pulse			Two-cycle sinusoidal at $f_0$		
		Synthetic Ape	rture Emission		
Receive Apodization			Hanning		
Number of transmitting elements	64	64	64	128	128
Number of receiving elements	64	64	64	128	128
Number of emissions $(K)$	260	260	204	192	312
	The values	of $\lambda$ for the simul	ated and experimental images		
BP	0.5	0.5	0.2	0.5	5
LS	0.7	1	0.5	1	1

#### G. Image Quality Measures

Three image quality metrics were evaluated: the contrast-to-noise ratio (CNR), the signal-to-noise ratio (SNR), and the resolution gain (RG). They were computed based on the envelope-detected signals independent of image display range.

Based on two regions  $R_1$  and  $R_2$  belonging to two different structures, CNR is defined as [45]

$$CNR = \frac{|\mu_{R_1} - \mu_{R_2}|}{\sqrt{\sigma_{R_1}^2 + \sigma_{R_2}^2}}$$
 (17)

where  $\mu_{R_1}$  and  $\mu_{R_2}$  are the mean values in the regions  $R_1$  and  $R_2$ , respectively, and  $\sigma_{R_1}$  and  $\sigma_{R_2}$  are the standard deviations of intensities in  $R_1$  and  $R_2$ , respectively.

The SNR is defined as the ratio between the mean value  $\mu$  and the standard deviation  $\sigma$  in homogeneous regions [12]

$$SNR = \frac{\mu}{\sigma}.$$
 (18)

The RG is defined in [46] as the ratio between the normalized autocorrelation function with values higher than 3 dB (computed for the DAS beamformed image in our case), over the normalized autocorrelation function (higher than 3 dB) of the images formed using the other aforementioned BF methods (MV, multibeam Capon, IAA, BP, and LS). Note that a value of RG > 1 needs to be achieved for achieving a better resolution than DAS beamformer.

#### V. RESULTS AND DISCUSSION

#### A. Individual Point Reflectors

With this simulation, we evaluate the potential of the proposed methods in sparse media. The resulted beamformed images are illustrated in Fig. 3. The result of DAS BF is

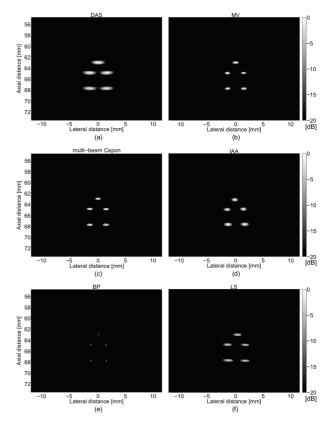


Fig. 3. (a) DAS, (b) MV, (c) multibeam Capon, (d) IAA, (e) BP, and (f) LS BF results of the simulation of individual point scatterers.

shown in Fig. 3(a). Using the MV BF, the lateral resolution is improved compared with DAS, IAA, and LS BF [see Fig. 3(b)], and it is comparable with the result of multibeam

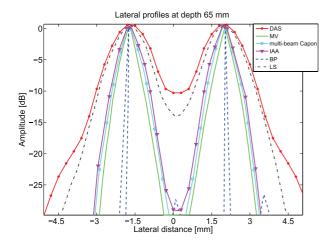


Fig. 4. Lateral profiles at 65-mm depth of the point reflectors represented in Fig. 3.

 $\label{eq:table II}$  CNR, SNR, and RG Values for the Simulated Phantom in Fig. 5

BF Method	CNR	SNR	RG
DAS	4.8	0.4	1
MV	5.3	0.61	3.64
multi-beam Capon	5.4	0.58	4.87
IAA	3.5	0.63	3.57
BP	6.5	0.62	8.72
LS	7.4	0.68	2.65

Capon BF [Fig. 3(c)]. Concerning the IAA beamformed result, as stated in [12], it gives better point-target resolvability than DAS [Fig. 3(d)]. The proposed BP BF has the best resolution of the point-like reflectors, being able to perfectly detect the five reflectors, by obtaining the most narrower mainlobes, due to the fact that BP results in a sparse representation of the beamformed signals [Fig. 3(e)]. As expected, LS beamformer results in solutions that tend to be smooth and regular, as in Fig. 3(f).

Fig. 4 presents the lateral profiles of the compared BF methods at 65 mm. We can observe that multibeam Capon and MV are comparable in terms of lateral profiles, but MV offers better delimitation of the two points. As observed, BP BF outperforms the other BF methods, being able to perfectly resolve the two points, also suppressing the sidelobes. Finally, LS BF gives the smoothest result.

#### B. Simulated Hypoechoic Cyst

The BF results of a hypoechoic cyst in a speckle pattern are shown in Fig. 5. We have highlighted with white circle the true borders of the cyst, in order to show the accuracy of the proposed methods regarding the dimensionality of the scanned structures.

The image quality metrics are detailed in Table II. To calculate the CNR, we have considered the region  $R_2$  inside

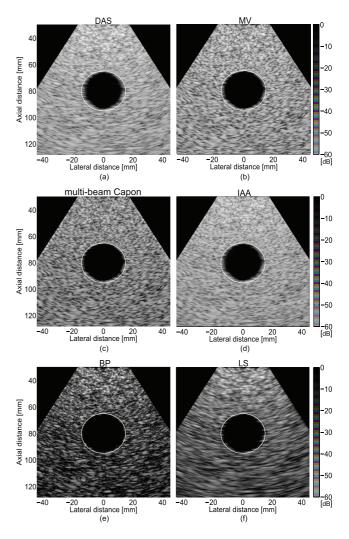


Fig. 5. (a) DAS, (b) MV, (c) multibeam Capon, (d) IAA, (e) BP, and (f) LS BF results of the hypoechoic cyst simulation.

the hypoechoic cyst (the black region), and the region  $R_1$ inside the homogeneous speckle, at the same depth and with the same dimension as the region  $R_2$ , as suggested in [5]. The SNR was computed for  $R_1$ . For calculating RG, the whole image was considered. As expected, with DAS, the cyst appears more narrow due to the low resolution and its low capability of resolving cyst-like structures inside the speckle pattern [Fig. 5(a)]. Using MV, we slightly increase the contrast and the resolution in the final image compared with DAS, the dimension of the cyst being closer to its real dimension, as shown in Fig. 5(b). Better resolution is obtained when the multibeam Capon is used, the RG being increased by a factor of almost 1.4. The improvement in resolution can also be observed in the delimitation of the cyst region compared with the white circle that represents the real dimension of the cyst [see Fig. 5(c)]. Compared with DAS, IAA increases the resolution of the beamformed image, but not as much as MV or multibeam Capon [Fig. 5(d)]. However, a contrast degradation can be observed from Table II. Finally, the proposed methods are reflecting more correctly the real dimension of the cyst, especially when using the BP BF [Fig. 5(e)], this being

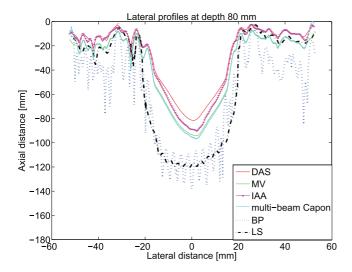


Fig. 6. Lateral profiles at 80-mm depth of the cyst phantom represented in Fig. 3.

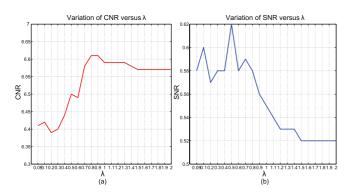


Fig. 7. Variation of (a) CNR and (b) SNR versus  $\lambda$  when the BP method was applied to the hypoechoic cyst simulation [Fig. 5(e)].

in concordance with the high increase in resolution (with a factor of two compared with MV) and contrast. As expected, LS tends to favor continuity and smoothness, especially when dealing with the speckle pattern [see Fig. 5(f)], the gain in resolution being less important. However, even so, it is more precise in reflecting the dimensionality of the cyst. Note that in terms of SNR, in comparison with DAS, all the other beamformers give better SNR, the best improvement being obtained with LS BF, which is also outperforming the other beamformers in terms of contrast.

Fig. 6 presents the lateral profiles of the results presented in Fig. 5, where the previous observations are confirmed. The curves in Fig. 6 are computed by averaging 15 lateral profiles around a depth of 80 mm. The proposed methods, BP and LS, have larger mainlobes than the other BF methods, which correspond to the true dimension of the hypoechoic cyst. We can also confirm the increase in contrast presented in Table II in the case of BP and LS BF approaches.

Fig. 7 presents the variation of CNR and SNR parameters as a function of  $\lambda$  hyperparameter in the case of BP BF. We can observe that a favorable compromise between CNR and SNR is reached when  $\lambda=0.5$ . The value of CNR can

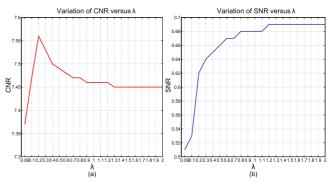


Fig. 8. Variation of (a) CNR and (b) SNR versus  $\lambda$  when the LS method was applied to the hypoechoic cyst simulation [Fig. 5(f)].

TABLE III
CNR, SNR, AND RG VALUES FOR THE SIMULATED US CARDIAC
BEAMFORMED IMAGES IN FIG. 9

BF Method	CNR	SNR	RG
DAS	1.12	0.47	1
MV	0.90	0.56	3.56
multi-beam Capon	0.61	0.54	4.23
IAA	1.30	0.62	5.3
BP	1.45	1.75	9.89
LS	1.55	1.88	2.06

be improved by increasing the value of  $\lambda$ . For example, when  $\lambda=0.9$ , CNR = 6.61, but the value of SNR is reduced, SNR = 0.55. Similarly, decreasing the value of  $\lambda$  will increase the value of SNR, while losing in CNR. Similarly, with Fig. 7, Fig. 8 presents the variation of CNR and SNR function of  $\lambda$  hyperparameter in the case of LS BF.

#### C. Simulated Cardiac Image

The results of BF the cardiac medium are shown in Fig. 9. With this example, we are interested in visualizing the LV region (hypoechoic), which is surrounded by the hyperechoic regions containing the anterior and posterior walls of the heart as well as the septum. The small echoic regions inside the LV region are the papillary muscles, which due to the low contrast and resolution of the DAS beamformed image are hard to be distinguished [Fig. 9(a)]. A better visualization of the walls is obtained with MV [Fig. 9(b)] and multibeam Capon [Fig. 9(c)], also resulting in an improved resolution, confirmed with a higher RG value (see Table III). For the calculation of the CNR, we considered  $R_1$  the region inside the white square, situated at approximately 18 (laterally) and around 55 mm (axially), while  $R_2$  is delimited by the black square, around -20 (laterally) and 55 mm (axially). To compute SNR, the  $R_1$  region was considered.

An interesting observation is that the value of the CNR in the case of MV and multibeam Capon is not improved compared with DAS. This is explained by the results in [5], where it has been shown that the improvement of the contrast

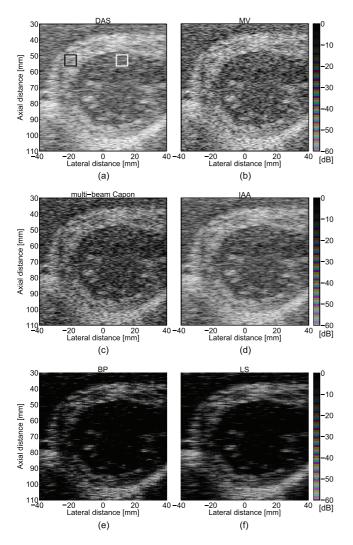


Fig. 9. (a) DAS, (b) MV, (c) multibeam Capon, (d) IAA, (e) BP, and (f) LS BF results of the simulation of a cardiac image.

directly depends on the high definition of the regions (the LV, the septum, and the walls in our example). Since the amplitude of the reflectors from the walls and septum are not so high compared with the region of LV that contains speckle, the contrast of the final image is affected. However, IAA improves both the contrast and the resolution of the image, presenting more defined regions, as shown in Fig. 9(d). Yet, the best improvement of the resolution is obtained when we promote Laplacian BF solutions, with BP BF [see Fig. 9(e)], resulting in an improvement by a factor of two in RG compared with MV and multibeam Capon and by a factor of almost ten compared with DAS. Of course, as expected, LS BF is highly improving the contrast and the SNR of the resulted image, while the RG is lower than when using the other BF approaches [Fig. 9(f)].

#### D. In Vivo Data: Carotid

Fig. 10 presents the BF results of the studied beamformers, and in Table IV, we calculated their corresponding CNR,

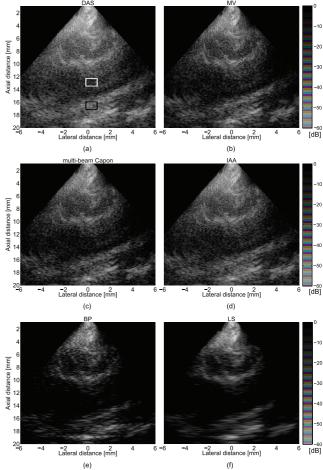


Fig. 10. (a) DAS, (b) MV, (c) multibeam Capon, (d) IAA, (e) BP, and (f) LS BF results of experimental carotid data.

TABLE IV
CNR, SNR, RG, AND COMPUTATIONAL TIME VALUES FOR THE
EXPERIMENTAL CAROTID BEAMFORMED IMAGES FROM FIG. 10

BF Method	CNR	SNR	RG	Computation time [s]
DAS	1.84	1.46	1	4.5552
MV	2.24	1.43	1.25	122
multi-beam Capon	1.32	1.49	1.25	368
IAA	1.48	1.47	1.34	8.9266
BP	1.85	1.49	1.48	60.4320
LS	1.94	1.55	1.17	8.8692

SNR, RG, and computational time values. In this example, the carotid is placed between 8 and 15 mm in the axial direction. In this region, the interior of the carotid artery is the hypoechoic structure surrounded by the arterial walls (which are hyperechoic). To calculate the CNR, we have considered region  $R_2$  inside the carotid (the white rectangle positioned at 0 mm laterally) and the region  $R_1$  inside the region of speckle (the black rectangle positioned at 0 mm laterally). The SNR for  $R_1$  was computed.

As observed, using DAS BF is hard to distinguish between the interior of the carotid and its walls [Fig. 10(a)]. This can also be explained by the fact that DAS BF result represents the lower RG. A better visualization of the structures of interest is obtained with MV and multibeam Capon, which have similar RG values. However, the contrast of the MV beamformed image is better, increasing the value of CNR by a factor of  $\approx 1$  compared with multibeam Capon. We can observe that multibeam Capon is clearly defining the region inside the carotid, by reducing the level of speckle inside it [see Fig. 10(c)]. The IAA beamformed image is comparable with the one of multibeam Capon, but it conserves better the speckle inside the carotid, offering a better resolution and a better contrast of the image. With the proposed approaches, however, we are able to better distinguish the interior of the carotid artery, as well as its walls, with a high gain in contrast and resolution resulted by applying BP BF. A loss in resolution can be observed when using LS, compared with BP, due to the use of the  $\ell_2$ -norm regularization. Note that, due to the formulation of the proposed direct model (8) that includes an additive noise, the proposed method is intrinsically denoising the signal (e.g., the noise inside the carotid is reduced) through the inversion process (see Table IV). The denoising effect obtained by our BF approach does not suffer from any spatial resolution loss, as it could be the case if the raw data or the beamformed images were low-pass filtered.

Regarding the computational time, note that it is highly dependent on the length of the acquired raw data. For this case, the number of ranges was around N=2500, and the proposed approaches were applied without a previous decimation of the raw data. Of course, the standard parallel computing methods could additionally improve the computational complexity, since the BF process is done for each lateral scanline. All the discussed methods were implemented with MATLAB R2013b, on an Intel i7 2600 CPU working at 3.40 GHz. Note that even so, LS BF is approaching the time capabilities of DAS, being just twice slower than DAS. Moreover, BP is also faster than MV. Thus, using the discussed techniques for improving the computational expense makes the two proposed methods good candidates for real-time applications.

#### E. In Vivo Data: Healthy Thyroid

Fig. 11 presents the BF results of healthy thyroid data. The thyroid (echoic region) is situated between the trachea and the carotid artery (laterally, between -20 and 30 mm approximately). Fig. 11(a) illustrates the result obtained with DAS BF. As expected, the contrast of the image is low and it is hard to distinguish the thyroid structure from the trachea, especially in the top-left part of the thyroid. However, when BP [Fig. 11(b)] and LS [Fig. 11(c)] are used, the thyroid region is easy to be identified, and the contrast of the image is increased.

The values of CNR, SNR, and RG are depicted in Table V. To compute CNR, we considered the region  $R_2$  inside the thyroid (the black circle positioned at approximately -10 mm laterally) and the region  $R_1$  inside trachea (the white circle positioned at approximately -40 mm laterally). The SNR for  $R_1$  was computed. We can observe that the best values of

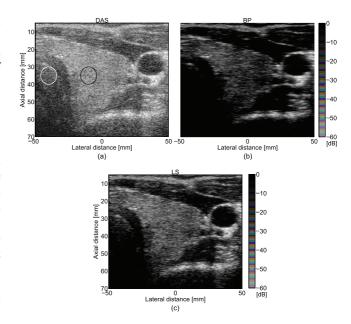


Fig. 11. (a) DAS, (b) BP, and (c) LS BF results of in vivo healthy thyroid data.

TABLE V
CNR, SNR, and RG Values for the *In Vivo* Healthy
Thyroidal Beamformed Images From Fig. 11

BF Method	CNR	SNR	RG
DAS	0.55	0.22	1
BP	1.13	0.32	3
LS	1.56	0.64	2.5

the CNR and SNR are obtained when the LS method was applied, the thyroid region being obvious to be discerned. The boundaries of the carotid artery are also well defined [Fig. 11(c)].

#### F. In Vivo Data: Thyroid With Tumor

The beamformed results of the thyroid data with tumor are presented in Fig. 12. The malignant tumor with an irregular structure can be seen between the left lobe of the thyroid (the hyperechoic structure situated near the trachea) and the carotid artery [the hypoechoic circular structure with the center at approximately 33 mm (axially) and 40 mm (laterally)]. We can observe that, contrarily to DAS beamformed image, where the tumor is hard to be distinguished [see Fig. 12(a)], both our methods improve the visualization of the main structures, enhancing the edges of the tumor. The values of CNR, SNR, and RG are depicted in Table VI, where a gain in resolution with a factor of almost three can be observed when using BP, compared with DAS, while with LS, we obtain a higher improvement in contrast and SNR than with BP BF. CNR was computed by considering region  $R_2$  inside the tumor (the black circle positioned at 0 mm laterally) and the region  $R_1$ inside the left lobe of the thyroid (the white circle positioned

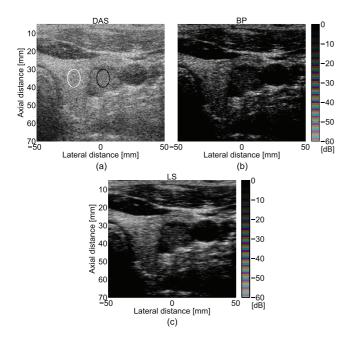


Fig. 12. (a) DAS, (b) BP, and (c) LS BF results of in vivo thyroid data with tumor.

TABLE VI
CNR, SNR, AND RG VALUES FOR THE *In Vivo* THYROIDAL
BEAMFORMED IMAGES FROM FIG. 12

BF Method	CNR	SNR	RG
DAS	0.71	0.62	1
BP	1.16	0.79	2.9
LS	1.32	0.86	1.5

at approximately -20 mm laterally). The SNR for  $R_1$  was computed.

#### VI. CONCLUSION

We have presented a new BF approach in US medical imaging that solves a regularized inverse problem based on a linear model relating, for each depth, the US reflected data to the signal of interest. Contrarily to existing techniques that use adaptive or nonadaptive weights to average the raw data in order to form RF lines, we directly recover, for each depth, the desired signals using Laplacian or Gaussian statistical assumptions. The proposed regularization-based BF allows us to take advantage of the beamspace processing that enables to highly reduce the number of US transmissions (by a factor of five in our examples), while improving the quality of the beamformed images compared with four existing beamformers. Multiple simulated and experimental examples were presented, which compare our approach with DAS, MV, multibeam Capon, and IAA beamformers. We showed that our BF approaches, based on Laplacian and Gaussian prior information, although based on the same model, are complementary in terms of result quality. Thus, Laplacian statistics are favoring sparse results, while the Gaussian law is offering

more regular and smooth images. We also proved through RG, CNR, and SNR image quality metrics that we obtained an important gain in spatial resolution and/or in contrast, while maintaining a reasonable computational time compared with other existing techniques. As a future work, we will consider other statistical assumptions, such as the generalized Gaussian distribution, resulting in  $\ell_p$ -norm minimization with the parameter p between 0 and 2. Following the choice of p, this should guarantee a better compromise between the gain in contrast and the improve of the spatial resolution (see [27], [47], [48]). Another way to obtain this compromise could be to combine the two regularization terms (through Laplacian and Gaussian statistics) used in our approach, resulting in an elastic net regularization (see [26], [49]).

Another interesting perspective offered by our BF direct linear model is the possibility to combine it with existing postprocessing techniques, aiming to enhance the quality of US images, such as deconvolution or super-resolution.

Including the knowledge of the PSF in our beamformer will also facilitate the comparison with existing regularized BF methods in medical US imaging.

## APPENDIX A MINIMUM VARIANCE BEAMFORMING

MV (or Capon filter) BF [50] consists in minimizing the array output power by maintaining a unit gain at the focal point. It adaptively calculates the weights, by solving

$$\min_{\boldsymbol{w}} \boldsymbol{w}^{H} \boldsymbol{R}_{k} \boldsymbol{w}, \quad \text{such that } \boldsymbol{w}^{H} \boldsymbol{1} = 1$$
 (19)

with the analytical solution

$$w_{MV} = \frac{R_k^{-1} \mathbf{1}}{\mathbf{1}^T R_k^{-1} \mathbf{1}} \tag{20}$$

where  $\mathbf{R}_k = E[\mathbf{y}_k \mathbf{y}_k^H]$  is the covariance matrix of  $\mathbf{y}_k$  and  $\mathbf{1}$  is a length M column vector of ones. These weights are then used to calculate the desired RF beamformed lines in a similar way as with DAS. In practice,  $\mathbf{R}_k$  is unavailable and the estimated covariance matrix  $\hat{\mathbf{R}}_k$  is used as alternative, derived from L received samples

$$\hat{R}_k = \sum_{l=1}^{L} y_k(l) y_k^H(l).$$
 (21)

Since the received focused raw data are coherent, several methods were proposed to decorrelate the data as much as possible: subaperture (or subarray) averaging (also called spatial smoothing), time averaging, and diagonal loading significantly improve the standard MV BF (see [51]).

## APPENDIX B BEAMSPACE BEAMFORMING

Starting from the MV BF method presented in Section VI, named element-space-based Capon (ES-Capon), Nilsen and Hafizovic [10] proposed a beamspace beamformer (BS-Capon) that allowed reducing the computational complexity of the MV BF by a ratio of three. Basically, they reduce the size of the covariance matrix in (21) by replacing it with a smaller covariance matrix of orthogonal beams. The expression of

the orthogonal beams is detailed in [10], and the beamspace transformation is expressed as follows:

$$y_{k_{\rm BS}} = B y_k \tag{22}$$

where  $B = [b_1, ..., b_M]^T$  is the  $M \times M$  Butler matrix whose elements are defined as

$$b_{mn} = \frac{1}{\sqrt{M}} e^{j\frac{2\pi}{M}\left(m - \frac{1}{2}\right)n}.$$
 (23)

 $\boldsymbol{B}$  is a unitary matrix  $(\boldsymbol{B}\boldsymbol{B}^H = \boldsymbol{B}^H\boldsymbol{B} = \boldsymbol{I}_{M\times M})$ , equivalent to an M-point discrete Fourier transform matrix.  $\boldsymbol{I}_{M\times M}$  is the identity matrix of size  $M\times M$ .

The transformation in (22) is applied to all signals and weights vectors in the ES to find their beamspaced version. Therefore, the weights of ES-Capon BF are formed by solving

$$\min_{\boldsymbol{w}_{BS}} \boldsymbol{w}_{BS}^{H} \boldsymbol{R}_{BS} \boldsymbol{w}_{BS}, \text{ such that } \boldsymbol{w}_{BS}^{H} e_{1} = 1.$$
 (24)

The solution of (24) is

$$\mathbf{w}_{\rm BS} = \frac{\mathbf{R}_{\rm BS}^{-1} e_1}{e_1^T \mathbf{R}_{\rm BS}^{-1} e_1} \tag{25}$$

where  $R_{\rm BS} = E[y_{k_{\rm BS}}y_{k_{\rm BS}}^H]$  is the covariance matrix of  $y_{k_{\rm BS}}$  and  $e_m$  is an  $M \times 1$  vector having the value 1 in the mth position and zero in all other positions. Finally, we can state that BS-Capon BF can be seen as the description of the Capon filter from (20) in the Fourier domain.

As stated before in this section, in the case of DAS, MV, and BS-Capon BF, the final RF US image is a collection of RF beamformed lines, each of which being the result of BF the raw RF signals coming from a focused wave emission in the direction  $\theta_k$ ,  $k \in \{1, ..., K\}$ , using the M elements of the transducer.

The most commonly used visualization mode in US medical imaging is the B (brightness)-mode. It is obtained by applying envelope detection and log-compression techniques to each beamformed RF line. Finally, all the RF lines are juxtaposed in the lateral direction to form the final 2-D US image, as shown in Fig. 1.

## APPENDIX C MULTIBEAM CAPON BEAMFORMING

Jensen and Austeng [11] used beamspace processing for reducing the dimensionality of the data, and proposed a new approach of Capon BF, called multibeam Capon BF. For more convenience, let us briefly recall their approach.

For a given range n, let us select its corresponding lateral scanline, as illustrated in Fig. 1. Since the signals  $y_{k_{\rm BS}}$  have been focused in axial direction (by applying time delays) before being beamformed, we just need to compensate for the phase shifts based on the distances from the samples of the lateral scanline (equivalent, in our case, with K, the number of beam directions).

The compensation of the phase shifts,  $\Delta \phi_m$ , with  $M=0,\ldots,M-1$ , is depicted in Fig. 13, assuming that the time-compensated data reach the elements at angle  $\theta_k$ . We consider the first elements as reference, so its phase shift is zero. Thus, based on the far-field assumption, we can

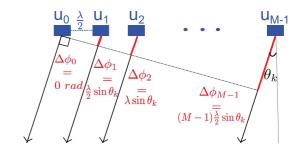


Fig. 13. Phase shift compensation of the focused raw data.

formulate the complex exponential version of the manifold vector for a given direction k, which corresponds to the incident angle  $\theta_k$ , as (see [11], [52])

$$\mathbf{a}_{\theta_k} = [1 \ e^{-j\pi \sin(\theta_k)} \cdots e^{-j(M-1)\pi \sin(\theta_k)}]^T.$$
 (26)

Thus, using phase shifts, for focusing along a lateral scanline, contrarily to the matrix  $R_{BS}$  used in BS-Capon, for a given range n, the covariance matrix R[n] will cover all the directions  $\theta_k$ , k = 1, ..., K. Therefore, the weights corresponding to a given direction  $\theta$  and a range n will be formed by solving

$$\min_{\mathbf{w}} \mathbf{w}^H \mathbf{R}[n] \mathbf{w}, \quad \text{such that } \mathbf{w}^H \mathbf{a}_{\theta,n} = 1$$
 (27)

having the solution

$$\boldsymbol{w}_{\theta,n} = \frac{\boldsymbol{R}^{-1}[n]\boldsymbol{a}_{\theta,n}}{\boldsymbol{a}_{\theta,n}^T \boldsymbol{R}^{-1}[n]\boldsymbol{a}_{\theta,n}}.$$
 (28)

These weights are then applied to calculate the signal corresponding to a lateral scanline, at a range n.

#### APPENDIX D

#### ITERATIVE ADAPTIVE APPROACH BEAMFORMING

Based on the beamspace processing technique and on the calculation of the multibeam covariance matrix discussed in Section VI, Jensen and Austeng [12] applied IAA [13] to US medical imaging. Following this recent work, a covariance matrix,  $\bar{R}[n]$  based on  $\bar{K}$  potential reflectors placed across a considered lateral scanline, was defined as:

$$\bar{\mathbf{R}}[n] = \sum_{k=1}^{K} |\mathbf{y}_{BS}[n]|^2 \mathbf{a}_{\theta} \mathbf{a}_{\theta}^T = \mathbf{A}_{BS} \mathbf{P} \mathbf{A}_{BS}^T$$
 (29)

with  $y_{BS}[n] \in \mathbb{C}^{N_b \times \bar{K}}$  the beamspaced time-delayed raw data at a given range n, before applying the phase-shift transform. A is the matrix containing the manifold column vectors defined in (26), and P is a diagonal matrix with the elements of  $|y_{BS}[n]|^2$  along its diagonal. The values of P are then iteratively updated and calculated by taking into account the weights corresponding to a lateral scanline, by following (28). Finally, P is used to estimate the amplitude of each reflector of the IAA BF result.

Contrarily of DAS, MV, and BS-Capon BF where, to form the final beamformed image, the RF lines are juxtaposed in the lateral direction, multibeam Capon and IAA BF are axially juxtaposing the beamformed lateral scanlines to form the final beamformed image.

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**Teodora Szasz** received the M.Sc. degree from the Technical University of Cluj-Napoca, Cluj-Napoca, Romania, in 2013. She is currently pursuing the Ph.D. degree with the Toulouse Institute of Computer Science Research, University Paul Sabatier Toulouse 3, Toulouse, France.

Her current research interests include the adaptive beamforming of medical ultrasound imaging and inverse problems.



Adrian Basarab received the M.S. and Ph.D. degrees in signal and image processing from the National Institute for Applied Sciences, Lyon, France, in 2005 and 2008, respectively.

Since 2009, he has been an Assistant Professor with University Paul Sabatier Toulouse 3, Toulouse, France, where he was also a member of the Institut de Recherche en Informatique de Toulouse Laboratory. He is an Associate Editor of the book Digital Signal Processing. His current research interests include medical imaging, particularly in motion

estimation, inverse problems, and ultrasound image formation.



**Denis Kouamé** (M'97) received the Ph.D. and Habilitation to supervise research works (HDR) degrees in signal processing and medical ultrasound imaging from the University of Tours, Tours, France, in 1996 and 2004, respectively.

From 1996 to 1998, he was a Senior Engineer with GIP Tours, Tours. From 1998 to 2008, he was initially an Assistant and then an Associate Professor with the University of Tours. He was the Head of the Signal and Image Processing Group from 2000 to 2006 and the Ultrasound Imaging Group

from 2006 to 2008 at the Ultrasound and Signal Laboratory, University of Tours. He is currently a Professor with University Paul Sabatier Toulouse 3, Toulouse, France, and a member of the Institut de Recherche en Informatique de Toulouse (IRIT) Laboratory (UMR 5505 of the CNRS), Toulouse. He also leads the Image Comprehension and Processing Group, IRIT. His current research interests include signal and image processing with applications to medical imaging and particularly ultrasound imaging, including high-resolution imaging, image resolution enhancement, Doppler signal processing, detection and estimation with applications to cerebral emboli detection, multidimensional parametric modeling, spectral analysis, inverse problems related to compressed sensing and restoration.

Dr. Kouamé has been involved in organizing several conferences and also led a number of invited conferences, special sessions, and tutorials at several IEEE conferences and workshops. He has been serving as an Associate Editor of the IEEE TRANSACTIONS ON ULTRASONICS, FERROELECTRICS, AND FREQUENCY CONTROL.