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MEAN FIELD GAMES: A TOY MODEL ON AN ERDŐS-RENYI GRAPH.

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ABSTRACT. The purpose of this short article is to address a simple example of a game with a large number of players in mean field interaction when the graph connection between them is not complete but is of the Erdős-Renyi type. We study the quenched convergence of the equilibria towards the solution of a mean field game. To do so, we follow recent works on the convergence problem for mean field games and we heavily use the fact that the master equation of the asymptotic game has a strong solution.

1. INTRODUCTION

Mean field game theory was initiated a decade ago in independent contributions by Lasry and Lions [13, 14, 15] and by Huang, Caines and Malhamé [10, 11].

The general purpose is to address stochastic differential games with a large number of players subject to mean field interactions. Numerous works on the theory have been dedicated to the analysis of the asymptotic formulation of the game, which is precisely referred to as a mean field game. In his lectures at Collège de France, see also the lecture notes by Cardaliaguet [2], Lions [16] exposed most of the background of the analytical approach. Since then, alternative strategies, including probabilistic ones, have been suggested, see for instance [4] or the forthcoming monograph [5, 6] together with the textbook [1].

The rationale for regarding the limit under the number of players is pretty clear: Generally speaking, games with a large number of players are known to be of a high complexity; because of the underlying mean field structure, equilibria are expected to be of a somewhat simpler structure in the asymptotic regime. This is indeed a key feature of mean field particle systems that, asymptotically, particles are not only statistically identical but become also independent, this latter fact being usually known as *propagation of chaos*, see the seminal lecture notes by Sznitman [18]. Equivalently, the limiting behavior of a mean field particle system may be summed up through the dynamics of a sole representative particle interacting with its own distribution. When recast within games with a large number of players,

propagation of chaos yields the following picture. First, an equilibrium (or a solution) of the asymptotic mean field game should consist of a flow of marginal distributions $(\mu_t)_{0 \leq t \leq T}$ accounting for the statistical states of the population (when in equilibrium) at any time $t \in [0, T]$, where T is the time duration of the game. Second, it should suffice to focus on a typical player (or particle) interacting with the flow $(\mu_t)_{0 \leq t \leq T}$ instead of regarding the whole collection of players. Given $(\mu_t)_{0 \leq t \leq T}$, the player (or particle) aims at optimizing some strategy (say for instance its velocity if speaking of a particle) in order to minimize its own energy, say for instance its kinetic energy (which directly depends on the velocity) together with some potential energy (which may include the interaction with the environment). Last, the Nash condition underpinning the notion of equilibrium says that the environment $(\mu_t)_{0 \leq t \leq T}$ forms an equilibrium if the best response under the flow $(\mu_t)_{0 \leq t \leq T}$ fits $(\mu_t)_{0 \leq t \leq T}$ exactly, namely if the collection of marginal laws of the optimal state is $(\mu_t)_{0 \leq t \leq T}$ itself. In short, an equilibrium is a fixed point for an application mapping a flow of statistical distributions onto another flow of statistical distributions.

From the analytical point of view, fixed points may be characterized as solutions of a forward-backward system made of two equations, a forward Fokker-Planck equation and a backward Hamilton-Jacobi-Bellman equation. The second one accounts for the optimization part in the search of an equilibrium, whilst the first one is used to solve for the Nash condition. The probabilistic approach obeys more or less the same principle, as it also relies on a forward-backward system, but of a somewhat different structure. Precisely, this forward-backward system reads as a forward-backward stochastic differential equation of the McKean-Vlasov type. The forward component provides the form of the optimal trajectories of the stochastic optimization problem whilst the McKean-Vlasov condition enforces the fixed point constraint following from the Nash condition. Whatever the approach, the key problem is to sort out the forward-backward structure arising in the characterization. It is indeed known that the Cauchy-Lipschitz theory for forward-backward systems only applies when T is small enough. Rephrased in our setting, this says that, when T is fixed, equilibria of mean field games cannot be systematically constructed by a straightforward contraction argument. Most of the time, it requires another method; for instance, it may be based on a fixed point theorem without uniqueness. Subsequently, uniqueness for mean field games is known in very few cases; for example, it holds true when the coefficients satisfy specific monotonicity conditions, which we shall illustrate below.

Although there is no specific reason for expecting it to hold true in full generality, uniqueness is however an important question. When it fails, it may be a very difficult question to select one of the equilibria. Also, uniqueness may be very useful for justifying the convergence of equilibria of games with finitely many players towards solutions of mean field games. Proving the convergence of finite player games is indeed a difficult problem. When

the finite player equilibria are taken over open loop strategies, compactness arguments, without any need for asymptotic uniqueness, may be used, see for instance [9, 12]; however, this strategy fails when equilibria are computed over strategies in closed loop form. In the latter case, the only strategy that has been known so far for tackling the convergence problem requires uniqueness, see [3]. The idea for doing so goes back to another key object due to Lions, which is known as the *master equation*. Whenever uniqueness holds, the forward backward system used to characterize the solution of the mean field game (whatever the approach is analytical or probabilistic) may be regarded as the characteristics of a nonlinear PDE set on the space of probability measures. This latter PDE is precisely the master equation. Whenever the master equation has a classical solution, say \mathcal{U} , see for instance [3, 7] for solvability results in that direction, convergence may be proved by letting \mathcal{U} act onto the empirical distribution of the equilibria of the finite player game.

The purpose of this paper is to revisit all these facts, but in a somewhat different framework of what is usually addressed in the literature. Whilst the interactions in the finite player system that have been considered so far are always defined on the complete graph, we here assume that some of the particles do not directly interact with one another and, more precisely, that the connections between all of them are given by the realization of an Erdős-Renyi graph with a non-trivial parameter. In order to simplify, we focus on a toy model only, for which the structure of the asymptotic game is not too complex. In this framework, we prove that the limiting problem is also a mean field game. Taking advantage of the simple form of the model, we manage to simplify the master equation and to show directly, with an affordable price, that it has a classical solution. Last, we address the convergence of the equilibria of the finite player system when the realization of the graph is frozen. To simplify, we assume that equilibria are taken over open loop strategies and we make use of the master equation, as explained above, to pass to the limit.

We present the model in Section 2. The asymptotic game is investigated in Section 3. The convergence problem is addressed in Section 4.

2. A TOY MODEL

First, we introduce the finite player game we alluded to in the introduction. For the sake of simplicity, all the particles we consider below live in \mathbb{R} .

2.1. Particle system. In this paragraph, we are given a collection of N real-valued Brownian motions $((W_t^i)_{0 \leq t \leq T})_{i=1, \dots, N}$, where T is some finite time horizon and N is a fixed integer. These processes are assumed to be constructed on a probability space $(\Omega, \mathcal{F}, \mathbb{P})$. The filtration generated by the noises and possibly augmented with an initial σ -field \mathcal{F}_0 , independent of $((W_t^i)_{0 \leq t \leq T})_{i=1, \dots, N}$, is denoted by $(\mathcal{F}_t)_{0 \leq t \leq T}$.

With each index $i \in \{1, \dots, N\}$, we associate a particle (or a player). The dynamics of player i are of the following linear form:

$$dX_t^i = \alpha_t^i dt + dW_t^i, \quad t \in [0, T].$$

The process $(\alpha_t^i)_{0 \leq t \leq T}$ is called a *control*. It is required to be $(\mathcal{F}_t)_{0 \leq t \leq T}$ progressively-measurable and to satisfy the square-integrability condition

$$\mathbb{E} \int_0^T (\alpha_t^i)^2 dt < \infty.$$

Last, the initial conditions $(X_0^i)_{i=1, \dots, N}$ are assumed to be \mathcal{F}_0 -measurable, square-integrable, independent and identically distributed.

2.2. Interactions between the particles. With each particle, say $i \in \{1, \dots, N\}$, we associate a cost functional J^i depending upon all the control processes $((\alpha_t^j)_{0 \leq t \leq T})_{j=1, \dots, N}$. It takes the following form:

$$\begin{aligned} J^i(((\alpha_t^j)_{0 \leq t \leq T})_{j=1, \dots, N}) &= \mathbb{E} \left[\frac{1}{2} \left(c_g X_T^i - \int_{\mathbb{R}} g(x) d\bar{\mu}_t^{N,i} \right)^2 \right. \\ &\quad \left. + \int_0^T \left[\frac{1}{2} \left(c_f X_t^i - \int_{\mathbb{R}} f(x) d\bar{\mu}_t^{N,i} \right)^2 + \frac{1}{2} (\alpha_t^i)^2 \right] dt \right], \end{aligned}$$

where c_f and c_g are two constants, which we set to 1 for simplicity, that is $c_f = c_g = 1$. The functions f and g map \mathbb{R} into itself. Generally speaking, we assume them to be smooth enough, but we prefer not to detail this assumption at this stage of the paper and to make it clear when necessary only.

In comparison with models that have been investigated so far, the novel point is the probability measure $\bar{\mu}_t^{N,i}$. We call it the *empirical distribution of the particles connected to i* . It is defined as

$$\bar{\mu}_t^{N,i} = \frac{1}{N_i} \sum_{j=1}^N \varepsilon^{i,j} \delta_{X_t^j},$$

where $(\varepsilon^{i,j})_{1 \leq i, j \leq N}$ is a collection of bits, that it $\varepsilon^{i,j} \in \{0, 1\}$, satisfying the symmetry condition

$$\varepsilon^{i,j} = \varepsilon^{j,i}, \quad \varepsilon^{i,i} = 0.$$

In words, the family $(\varepsilon^{i,j})_{1 \leq i, j \leq N}$ forms an undirected graph between the particles: i and j are connected if $\varepsilon^{i,j} = 1$; in particular i and j are connected if j and i are connected. Above, N_i is the number of particles connected to i , namely

$$N_i = \sum_{j=1}^N \varepsilon^{i,j}.$$

Whenever $N_i = 0$, we let $\bar{\mu}_t^{N,i}$ be the null measure. In such a case this is no longer a probability measure and the action of $\bar{\mu}_t^{N,i}$ on f (or g) is zero.

Here, the potential energy forces i to get closer to some mean state of the particles connected to i . Of course, the form we choose here may not be

of the most relevant form for practical applications, but it turns out to be tractable from the mathematical point of view.

Throughout the paper, we assume that the $(\varepsilon^{i,j})_{1 \leq i < j \leq N}$ are the realizations of independent and identically distributed Bernoulli random variables of parameter $p \in (0, 1)$. With a slight abuse of notation, we still denote by $(\varepsilon^{i,j})_{1 \leq i < j \leq N}$ the corresponding random variables. So, the graph connecting the particles with one another is a realization of an Erdős-Renyi graph. It must be stressed that the collection of random variables $(\varepsilon^{i,j})_{1 \leq i < j \leq N}$ is independent of $(X_0, (W_t^i)_{0 \leq t \leq T})$. To make it clear, we define it on another probability space $(\Xi, \mathcal{G}, \mathbf{P})$. In particular, the realization of the graph is completely frozen in the computation of the costs $(J^i)_{i=1, \dots, N}$.

2.3. Notion of equilibria. Of course, each player is willing to minimize its own energy. This prompts us to introduce the notion of Nash equilibrium. As mentioned in introduction, several forms of it are conceivable. Here, we just make use of the simplest version:

Definition 1. A tuple $((\alpha_t^i)_{0 \leq t \leq T})_{i=1, \dots, N}$ is said to be a Nash equilibrium over open loop strategies if, for any $i \in \{1, \dots, N\}$, for any other adapted and square-integrable strategy $(\beta_t)_{0 \leq t \leq T}$, it holds that

$$J^i((\alpha_t^1, \dots, \alpha_t^{i-1}, \beta_t, \alpha_t^{i+1}, \dots, \alpha_t^N)_{0 \leq t \leq T}) \geq J^i((\alpha_t^1, \dots, \alpha_t^N)_{0 \leq t \leq T}).$$

In short, a Nash equilibrium is a consensus between the players: there is no incentive for any particle to leave the consensus.

In comparison with, Nash equilibria over closed loop strategies (as we alluded to in Introduction) are taken over strategies of the form $(\alpha_t^i = \phi^i(t, X_t^1, \dots, X_t^N))_{0 \leq t \leq T}$ for a feedback function ϕ^i from $[0, T] \times \mathbb{R}^N$ to \mathbb{R} . In the latter case, the dynamics for $(X_t^1, \dots, X_t^N)_{0 \leq t \leq T}$ become a stochastic differential equation. Under appropriate assumptions on the feedback functions, it is well-solvable. If so, we have the same definition as above for a Nash equilibrium except that, therein, we have to fix the feedback function ϕ^i instead of the trajectory $(\alpha_t^i)_{0 \leq t \leq T}$, which makes a difference. In words, this means that the player may take into account the current states of the others to update his/her strategy. Although it sounds pretty much realistic, we feel better not to address this case in this note.

We now introduce the asymptotic formulation. The guess is that there should be sufficiently many edges connecting i to the rest of the population to do as if the graph was complete. So, we expect the asymptotic game to fit a standard mean field game. The solutions of this latter one are given by the following two step procedure:

First Step. For a fixed flow of probability measures $(\mu_t)_{0 \leq t \leq T}$, solve the standard optimization problem

$$\inf_{(\alpha_t)_{0 \leq t \leq T}} J((\alpha_t)_{0 \leq t \leq T}),$$

where $(\alpha_t)_{0 \leq t \leq T}$ is an adapted process with respect to the filtration generated by $(X_0, (W_t)_{0 \leq t \leq T})$, for $(X_0, (W_t)_{0 \leq t \leq T}) \sim (X_0^1, (W_t^1)_{0 \leq t \leq T})$. Above the energy functional is given by

$$J((\alpha_t)_{0 \leq t \leq T}) = \mathbb{E} \left[\frac{1}{2} \left(X_T - \int_{\mathbb{R}} g(x) d\mu_t \right)^2 + \int_0^T \left[\frac{1}{2} \left(X_t - \int_{\mathbb{R}} f(x) d\mu_t \right)^2 + \frac{1}{2} (\alpha_t)^2 \right] dt \right],$$

where $(X_t)_{0 \leq t \leq T}$ solves

$$X_t = X_0 + \int_0^t \alpha_s ds + W_t, \quad t \in [0, T].$$

Second Step. Find $(\mu_t)_{0 \leq t \leq T}$ such that the flow of marginal laws of the optimal process in the first step is $(\mu_t)_{0 \leq t \leq T}$ itself.

The reader will easily understand that, with our form of interaction, the limiting optimization problem in the first step is linear-quadratic. In this regard, the solution is very simple and is known to be, conditional on the realization of the initial condition, a Gaussian process. This is one of the main rationale for the form of the toy model addressed in this note.

Notice also that the asymptotic problem would be of the same form if we used Nash equilibria over closed loop strategies.

3. MASTER EQUATION

This section is devoted to the analysis of the mean field game.

3.1. Solvability of the mean field game. We first make use of the probabilistic approach initiated in [4]. Solutions of the mean field game may be characterized through the following forward-backward stochastic differential equation of the McKean-Vlasov type:

$$\begin{aligned} dX_t &= -Y_t dt + dW_t, \\ (1) \quad dY_t &= -(X_t - \mathbb{E}[f(X_t)]) dt + Z_t dW_t, \quad t \in [0, T], \\ Y_T &= X_T - \mathbb{E}[g(X_T)]. \end{aligned}$$

Here we assume as before that the initial condition X_0 of the forward component has been prescribed. It is assumed to be \mathcal{F}_0 measurable and $(W_t)_{0 \leq t \leq T}$ is assumed to be an $(\mathcal{F}_t)_{0 \leq t \leq T}$ -Brownian motion with values in \mathbb{R} , for a filtration $(\mathcal{F}_t)_{0 \leq t \leq T}$ that may no longer coincide with the one described in the first section. Here, the typical example for $(\mathcal{F}_t)_{0 \leq t \leq T}$ is the filtration generated by \mathcal{F}_0 and $(W_t)_{0 \leq t \leq T}$. Recall also that the solution of the backward equation is the pair $(Y_t, Z_t)_{0 \leq t \leq T}$, which is also required to be adapted. Basically, the martingale integrand $(Z_t)_{0 \leq t \leq T}$ is used to ensure the adaptedness. It is assumed to be square-integrable.

We now make use of the very simple form of the system (1). By the change of variable $(\tilde{Y}_t = Y_t - X_t, \tilde{Z}_t = Z_t - 1)_{0 \leq t \leq T}$, it may be rewritten as

$$\begin{aligned} dX_t &= -(X_t + \tilde{Y}_t)dt + dW_t, \\ d\tilde{Y}_t &= (\tilde{Y}_t + \mathbb{E}[f(X_t)])dt + \tilde{Z}_t dW_t, \quad t \in [0, T], \\ \tilde{Y}_T &= -\mathbb{E}[g(X_T)]. \end{aligned}$$

It is then well checked that $(\tilde{Y}_t)_{0 \leq t \leq T}$ is necessarily deterministic as the coefficients are purely deterministic. In particular, any solution must satisfy

$$\frac{d}{dt}(e^t \mathbb{E}(X_t)) = -e^t \tilde{Y}_t,$$

which permits to express the forward component as

$$(2) \quad e^t X_t = X_0 - \mathbb{E}(X_0) + e^t \mathbb{E}(X_t) + e^t \Gamma_t,$$

where $(\Gamma_t)_{0 \leq t \leq T}$ is the Ornstein-Uhlenbeck process

$$\Gamma_t = e^{-t} \int_0^t e^s dW_s, \quad t \in [0, T].$$

Finally, the McKean Vlasov forward backward SDE takes the form

$$\begin{aligned} dX_t &= -(X_t + \tilde{Y}_t)dt + dW_t, \\ d\tilde{Y}_t &= (\tilde{Y}_t + \mathbb{E}[f(e^{-t}(X_0 - \mathbb{E}(X_0)) + \mathbb{E}(X_t) + \Gamma_t)])dt, \quad t \in [0, T], \\ \tilde{Y}_T &= -\mathbb{E}[g(e^{-T}(X_0 - \mathbb{E}(X_0)) + \mathbb{E}(X_T) + \Gamma_T)]. \end{aligned}$$

Taking the mean and recalling (2), we deduce the the following proposition, the proof of which is straightforward.

Proposition 2. *The McKean Vlasov forward backward SDE (1) is uniquely solvable if and only if the following deterministic system is uniquely solvable:*

$$\begin{aligned} dm_t &= -(m_t + \tilde{Y}_t)dt, \\ d\tilde{Y}_t &= (\tilde{Y}_t + \mathbb{E}[f(e^{-t}(X_0 - \mathbb{E}(X_0)) + m_t + \Gamma_t)])dt, \quad t \in [0, T], \\ \tilde{Y}_T &= -\mathbb{E}[g(e^{-T}(X_0 - \mathbb{E}(X_0)) + m_T + \Gamma_T)], \end{aligned}$$

with $m_0 = \mathbb{E}(X_0)$ as initial condition, in which case $\mathbb{E}(X_t) = m_t$, for all $t \in [0, T]$.

By the new change of variable $(\bar{m}_t = e^t m_t - \mathbb{E}(X_0), \bar{Y}_t = e^{-t} \tilde{Y}_t)_{0 \leq t \leq T}$, the above system is equivalent to

$$(3) \quad \begin{aligned} d\bar{m}_t &= -e^{2t} \bar{Y}_t dt, \quad \bar{m}_0 = 0, \\ d\bar{Y}_t &= e^{-t} \mathbb{E}[f(e^{-t}(X_0 + \bar{m}_t) + \Gamma_t)]dt, \\ \bar{Y}_T &= -e^{-T} \mathbb{E}[g(e^{-T}(X_0 + \bar{m}_T) + \Gamma_T)], \end{aligned}$$

in which case $\mathbb{E}(X_t) = e^{-t}(\bar{m}_t + \mathbb{E}(X_0))$. The following proposition provides conditions under which existence and uniqueness hold true.

Proposition 3. *Assume that f and g are non-increasing and Lipschitz-continuous functions. Then, for any square-integrable initial condition X_0 , there exists a unique solution to the system (3).*

Proof. The result follows from a standard fact in the theory of forward-backward systems. It suffices to observe that, when X_0 is given, the functions

$$\begin{aligned}\mathbb{R} \ni x &\mapsto -\mathbb{E}[f(e^{-t}(X_0 + x) + \Gamma_t)], \quad t \in [0, T], \\ \mathbb{R} \ni x &\mapsto -\mathbb{E}[g(e^{-T}(X_0 + x) + \Gamma_T)],\end{aligned}$$

are non-decreasing and Lipschitz continuous. This permits to regard the system (3) as a deterministic forward-backward system with non-decreasing and Lipschitz continuous coefficients. It is then known to be uniquely solvable, see for instance [17]. \square

Actually, we can even prove a more general stability result.

Proposition 4. *Assume that f and g satisfy the same assumptions as above. Then, there exists a constant C such that, for any two initial conditions X_0 and X'_0 of the system (1),*

$$|\bar{Y}_0 - \bar{Y}'_0| \leq C\mathbb{E}[|X_0 - X'_0|],$$

where $(\bar{m}_t, \bar{Y}_t)_{0 \leq t \leq T}$ and $(\bar{m}'_t, \bar{Y}'_t)_{0 \leq t \leq T}$ are the corresponding solutions to (3).

Proof. Again, the proof is pretty standard. It consists in regarding the product $((\bar{m}'_t - \bar{m}_t)(\bar{Y}'_t - \bar{Y}_t))_{0 \leq t \leq T}$. We compute

$$\begin{aligned}\frac{d}{dt}[(\bar{m}'_t - \bar{m}_t)(\bar{Y}'_t - \bar{Y}_t)] \\ = -e^{2t}(\bar{Y}'_t - \bar{Y}_t)^2 \\ + e^{-t}(\bar{m}'_t - \bar{m}_t)\left(\mathbb{E}[f(e^{-t}(X'_0 + \bar{m}'_t) + \Gamma_t) - \mathbb{E}[f(e^{-t}(X_0 + \bar{m}_t) + \Gamma_t)]\right).\end{aligned}$$

Since f is Lipschitz-continuous and non-increasing, we can find a constant C , only depending on the coefficients, such that

$$\begin{aligned}\frac{d}{dt}[(\bar{m}'_t - \bar{m}_t)(\bar{Y}'_t - \bar{Y}_t)] &\leq -e^{2t}(\bar{Y}'_t - \bar{Y}_t)^2 \\ &\quad + C|\bar{m}'_t - \bar{m}_t|\mathbb{E}[|X'_0 - X_0|],\end{aligned}$$

which yields

$$\begin{aligned}-(\bar{m}'_t - \bar{m}_t)(\bar{Y}'_t - \bar{Y}_t) + \int_t^T e^{2s}(\bar{Y}'_s - \bar{Y}_s)^2 ds \\ \leq -(\bar{m}'_T - \bar{m}_T)(\bar{Y}'_T - \bar{Y}_T) + C\mathbb{E}[|X'_0 - X_0|] \int_t^T |\bar{m}'_s - \bar{m}_s| ds.\end{aligned}$$

Now, using the fact that g is Lipschitz-continuous and non-increasing,

$$\begin{aligned}
& -(\bar{m}'_T - \bar{m}_T)(\bar{Y}'_T - \bar{Y}_T) \\
& = (\bar{m}'_T - \bar{m}_T)e^{-T} \left(\mathbb{E}[g(e^{-T}(X'_0 + \bar{m}'_T) + \Gamma_T)] \right. \\
& \quad \left. - \mathbb{E}[g(e^{-T}(X_0 + \bar{m}_T) + \Gamma_T)] \right) \\
& \leq C|\bar{m}'_T - \bar{m}_T| \mathbb{E}[|X'_0 - X_0|].
\end{aligned}$$

Recalling that $\bar{m}_0 = \bar{m}'_0 = 0$, we end up with

$$\begin{aligned}
& \int_0^T e^{2s} (\bar{Y}'_s - \bar{Y}_s)^2 ds \\
& \leq C|\bar{m}'_T - \bar{m}_T| \mathbb{E}[|X'_0 - X_0|] + C \mathbb{E}[|X'_0 - X_0|] \int_t^T |\bar{m}'_s - \bar{m}_s| ds.
\end{aligned}$$

Returning to the forward equation, we deduce that

$$\begin{aligned}
& \sup_{0 \leq t \leq T} |\bar{m}_t - \bar{m}'_t|^2 \\
& \leq C|\bar{m}'_T - \bar{m}_T| \mathbb{E}[|X'_0 - X_0|] + C \mathbb{E}[|X'_0 - X_0|] \int_t^T |\bar{m}'_s - \bar{m}_s| ds.
\end{aligned}$$

By a standard convexity inequality, we obtain

$$\sup_{0 \leq t \leq T} |\bar{m}_t - \bar{m}'_t| \leq C \mathbb{E}[|X'_0 - X_0|].$$

Plugging into the backward equation, we easily complete the proof. \square

3.2. Master field. The combination of Propositions 3 and 4 permits to define a Lipschitz continuous mapping

$$L^2(\Omega, \mathcal{F}, \mathbb{P}; \mathbb{R}) \ni X_0 \mapsto \bar{Y}_0 \in \mathbb{R},$$

which maps the initial condition of the system (1) onto the initial value of the backward component of the system (3). Remarkably, the Lipschitz property holds true with respect to the L^1 norm. Importantly, the right-hand side only depends on X_0 through its distribution. So, the above mapping may be regarded as a mere mapping from $\mathcal{P}_2(\mathbb{R})$ into \mathbb{R} , namely

$$\mathcal{V}(0, \cdot) : \mathcal{P}_2(\mathbb{R}) \ni \mu_0 \mapsto \bar{Y}_0 \in \mathbb{R},$$

where the right-hand side is computed for some $X_0 \sim \mu_0$. This mapping is Lipschitz continuous for the 1-Wasserstein distance \mathcal{W}_1 (and so for the 2-Wasserstein distance as well).

Of course, we can easily extend the argument to any initial time $t_0 \in [0, T]$ by initializing the system (1) at time t_0 . In words, this amounts to initialize the system (3) at time t_0 or, equivalently, in shifting it to $[0, T - t_0]$. The mapping $\mathcal{V}(t_0, \cdot)$ is then defined as

$$\mathcal{V}(t_0, \cdot) : \mathcal{P}_2(\mathbb{R}) \ni \mu_0 \mapsto \bar{Y}_{t_0} \in \mathbb{R},$$

where the right-hand side is the solution of the backward equation in (3), when this latter one is set on $[0, T - t_0]$ in lieu of $[0, T]$, e^T and Γ_T in the terminal condition being replaced by e^{T-t_0} and Γ_{T-t_0} , and X_0 being some \mathcal{F}_0 -measurable random variable with the same distribution as the initial condition at time t_0 of the system (1).

A crucial fact for the analysis of the finite player game equilibria is to prove that the mapping \mathcal{V} is smooth with respect to the measure argument, or, say to simplify, that the mapping $\mathcal{V}(0, \cdot)$ is smooth. This requires to give first a meaning to the underlying notion of derivative. To do so, we shall use Lions' approach to the so-called *Wasserstein derivative*, as introduced in [16], see also [2, 5]. This requires to define first the lifting of $\mathcal{V}(0, \cdot)$ to $L^2(\Omega, \mathcal{F}_0, \mathbb{P}; \mathbb{R})$, namely

$$\hat{\mathcal{V}}(0, X_0) = \mathcal{V}(0, \mathcal{L}(X_0)),$$

where $\mathcal{L}(X_0)$ is the law of X_0 . We then say that $\mathcal{V}(0, \cdot)$ is continuously differentiable if the lift $\hat{\mathcal{V}}(0, \cdot)$ is Fréchet continuously differentiable. Of course, this makes sense as long as the space $(\Omega, \mathcal{F}_0, \mathbb{P})$ is rich enough so that, for any $\mu_0 \in \mathcal{P}_2(\mathbb{R})$, there exists a random variable $X_0 \in L^2(\Omega, \mathcal{F}_0, \mathbb{P}; \mathbb{R})$ such that $X_0 \sim \mu_0$. Here, $\mathcal{P}_2(\mathbb{R})$ is the space of probability measures on \mathbb{R} with a finite second-order moment. So, in the sequel, $(\Omega, \mathcal{F}_0, \mathbb{P})$ is assumed to be atomless, which makes it rich enough. A crucial point with Lions' approach to Wasserstein differential calculus is that the Fréchet derivative of $\hat{\mathcal{V}}(0, \cdot)$, which can be identified with a square-integrable random variable, may be represented at point X_0 as $\partial_\mu \mathcal{V}(0, \mu_0)(X_0)$ for a mapping $\partial_\mu \mathcal{V}(0, \mu_0)(\cdot) : \mathbb{R} \ni v \mapsto \partial_\mu \mathcal{V}(0, \mu_0)(v) \in \mathbb{R}$. This latter function plays the role of Wasserstein derivative of $\mathcal{V}(0, \cdot)$ in the measure argument.

Second-order derivatives may be defined accordingly. The function $\mathcal{V}(0, \cdot)$ is said to be twice continuously differentiable if we can find, for each $\mu \in \mathcal{P}_2(\mathbb{R})$, a version of $\mathbb{R} \ni v \mapsto \partial_\mu \mathcal{V}(0, \mu)(v)$ (which is *a priori* defined as a function in $L^2(\mathbb{R}, \mu; \mathbb{R})$) such that $\mathcal{P}_2(\mathbb{R}) \times \mathbb{R} \ni (\mu, v) \mapsto \partial_\mu \mathcal{V}(0, \mu)(v)$ is continuous, differentiable in v when μ is fixed, and differentiable in μ when v is fixed, with jointly continuous derivatives, namely $\mathcal{P}_2(\mathbb{R}) \times \mathbb{R} \ni (\mu, v) \mapsto \partial_v \partial_\mu \mathcal{V}(0, \mu)(v)$ is continuous and $\mathcal{P}_2(\mathbb{R}) \times \mathbb{R} \times \mathbb{R} \ni (\mu, v, v') \mapsto \partial_\mu^2 \mathcal{V}(0, \mu)(v, v')$ is continuous (here the variable v' accounts for the second-order derivative in the direction μ).

3.3. Proving the smoothness of \mathcal{V} . Generally speaking, the smoothness of \mathcal{V} at order two (which the one we need for writing down the master equation, see Subsection 3.4) is addressed in the two papers [3, 7] and in the monograph [6]. The general method consists in proving first that $\mathcal{V}(t, \cdot)$ is a smooth function of the measure argument for t close to T , or equivalently, in establishing the smoothness of $\mathcal{V}(t, \cdot)$ for $t \in [0, T]$ when T is less than some small $\delta > 0$. In order to so, f and g must be smooth enough, say bounded and three times differentiable, with bounded derivatives. The idea is then to show that the solution to (3) varies smoothly with $X_0 \in L^2(\Omega, \mathcal{F}, \mathbb{P}; \mathbb{R})$

when $T \leq \delta$; this approach should be regarded as a flow method, but with an initial condition in a space of infinite dimension. In this regard, the condition $T \leq \delta$ plays a key role, as it permits to define the successive derivatives of the flow as the solutions of differentiated forward-backward systems. We feel better not to detail the analysis and to refer the reader to the aforementioned citations. The crux is then to pass from a small time horizon T to an arbitrary time horizon. This is the point when the Lipschitz control of \mathcal{V} comes in. In [6] and [7], it is shown that, for a given $T > 0$, the analysis of the smoothness of \mathcal{V} could be extended from the neighborhood $[T - \delta, T]$ of the terminal time T to the whole $[0, T]$ by iterating backwards the small time analysis along a sequence of small intervals of the form $[T - (n + 1)\delta, T - n\delta]$, with n such that $T - n\delta > 0$. Importantly, δ can be assumed to be the same throughout the induction. To make it clear, it must be stressed that the time length δ , as defined right above, depends in fact on the Lipschitz constants of f and g with respect to the law of X_0 for the distance \mathcal{W}_1 , with the slight abuse that f and g are seen as functions of the law of X_0 . When iterating the small time analysis as we just described, the terminal boundary condition has to be updated on any new interval, namely g has to be replaced by $\mathcal{V}(T - n\delta, \cdot)$ when working on the interval $[T - (n + 1)\delta, T - n\delta]$.

For the sake of illustration, we provide an explicit example.

Proposition 5. *Assume that f is null and that g is non-increasing, bounded and smooth, with bounded derivatives. Then, the function \mathcal{V} is jointly continuous. It is differentiable in time and $[\partial\mathcal{V}/\partial t]$ is jointly continuous. Also, for any $t \in [0, T]$, the function $\mathcal{V}(t, \cdot)$ is differentiable with respect to the measure argument and the function*

$$[0, T] \times \mathcal{P}_2(\mathbb{R}) \times \mathbb{R} \ni (t, \mu, v) \mapsto \partial_\mu \mathcal{V}(t, \mu)(v)$$

is bounded and continuous. It is differentiable in the variable v and in the variable μ and the functions

$$[0, T] \times \mathcal{P}_2(\mathbb{R}) \times \mathbb{R} \ni (t, \mu, v) \mapsto \partial_v \partial_\mu \mathcal{V}(t, \mu)(v),$$

$$[0, T] \times \mathcal{P}_2(\mathbb{R}) \times \mathbb{R} \times \mathbb{R} \ni (t, \mu, v, v') \mapsto \partial_\mu^2 \mathcal{V}(t, \mu)(v, v'),$$

are bounded and continuous.

Proof. If $f = 0$, the solution of the forward-backward system (3) is given by the fixed point condition

$$\begin{aligned} (4) \quad \bar{m}_T &= \int_0^T e^{2t-T} \mathbb{E}[g(e^{-T}(X_0 + \bar{m}_T) + \Gamma_T)] dt \\ &= \sinh(T) \mathbb{E}[g(e^{-T}(X_0 + \bar{m}_T) + \Gamma_T)]. \end{aligned}$$

In order to compute the derivative of $\mathcal{V}(0, \cdot)$, we proceed as follows. We consider two random variables X_0 and Y in $L^2(\Omega, \mathcal{F}, \mathbb{P}; \mathbb{R})$, with $\mathbb{E}[|Y|^2] \leq 1$ and with (X_0, Y) independent of $(W_t)_{0 \leq t \leq T}$. For $\varepsilon \in \mathbb{R}$, we call $(\bar{m}_t^\varepsilon, \bar{Y}_t^\varepsilon)_{0 \leq t \leq T}$

the solution to (3) when (1) is initialized with $X_0 + \varepsilon Y$. Since $(\bar{Y}_t^\varepsilon)_{0 \leq t \leq T}$ is constant in time,

$$(5) \quad \bar{m}_T^\varepsilon = -\frac{1}{2}(e^{2T} - 1)\mathcal{V}(0, \mathcal{L}(X_0 + \varepsilon Y)).$$

By Proposition 4,

$$|\bar{m}_T^\varepsilon - \bar{m}_T^0| \leq C\varepsilon \mathbb{E}[|Y|].$$

Letting $e_T = e^{-T} \sinh(T)$, we have

$$\begin{aligned} \bar{m}_T^\varepsilon &= \sinh(T) \mathbb{E}[g(e^{-T}(X_0 + \varepsilon Y + \bar{m}_T^\varepsilon) + \Gamma_T)] \\ &= \sinh(T) \mathbb{E}[g(e^{-T}(X_0 + \bar{m}_T^0 + \varepsilon Y + (\bar{m}_T^\varepsilon - \bar{m}_T^0)) + \Gamma_T)] \\ &= \sinh(T) \mathbb{E}[g(e^{-T}(X_0 + \bar{m}_T^0) + \Gamma_T)] \\ &\quad + \varepsilon e_T \mathbb{E}[g'(e^{-T}(X_0 + \bar{m}_T^0) + \Gamma_T)Y] \\ &\quad + e_T \mathbb{E}[g'(e^{-T}(X_0 + \bar{m}_T^0) + \Gamma_T)](\bar{m}_T^\varepsilon - \bar{m}_T^0) \\ &\quad + O(\varepsilon^2), \end{aligned}$$

where the Landau symbol $O(\varepsilon^2)$ is independent of X_0 and Y . Recalling that $\sinh(T) \mathbb{E}[g(e^{-T}(X_0 + \bar{m}_T^0) + \Gamma_T)] = \bar{m}_T^0$, we get

$$\begin{aligned} &\left(1 - e_T \mathbb{E}[g'(e^{-T}(X_0 + \bar{m}_T^0) + \Gamma_T)]\right)(\bar{m}_T^\varepsilon - \bar{m}_T^0) \\ &= \varepsilon e_T \mathbb{E}[g'(e^{-T}(X_0 + \bar{m}_T^0) + \Gamma_T)Y] + O(\varepsilon^2), \end{aligned}$$

that is, since $g' \leq 0$,

$$\bar{m}_T^\varepsilon - \bar{m}_T^0 = \varepsilon e_T \frac{\mathbb{E}[g'(e^{-T}(X_0 + \bar{m}_T^0) + \Gamma_T)Y]}{1 - e_T \mathbb{E}[g'(e^{-T}(X_0 + \bar{m}_T^0) + \Gamma_T)]} + O(\varepsilon^2).$$

Recall from (5) that

$$\mathcal{V}(0, \mathcal{L}(X_0)) = -e^{-2T} e_T^{-1} \bar{m}_T^0,$$

which shows that $\mathcal{V}(0, \cdot)$ is differentiable with respect to the measure argument; also, for any $\mu \in \mathcal{P}_2(\mathbb{R})$ and $v \in \mathbb{R}$,

$$\partial_\mu \mathcal{V}(0, \mu)(v) = -\frac{e^{-2T} \mathbb{E}[g'(e^{-T}v - e^T e_T \mathcal{V}(0, \mu) + \Gamma_T)]}{1 - e_T \mathbb{E}[g'(e^{-T}X_0 - e^T e_T \mathcal{V}(0, \mu) + \Gamma_T)]}.$$

It is absolutely obvious to check that $\partial_\mu \mathcal{V}(0, \mu)(v)$ can be differentiated once more with respect to μ and v , namely

$$\partial_v \partial_\mu \mathcal{V}(0, \mu)(v) = -\frac{e^{-3T} \mathbb{E}[g''(e^{-T}v - e^T e_T \mathcal{V}(0, \mu) + \Gamma_T)]}{1 - e_T \mathbb{E}[g'(e^{-T}X_0 - e^T e_T \mathcal{V}(0, \mu) + \Gamma_T)]},$$

whilst the chain rule yields

$$\begin{aligned} & \partial_\mu^2 \mathcal{V}(0, \mu)(v, v') \\ &= e^{-T} e_T \partial_\mu \mathcal{V}(0, \mu)(v') \left(\frac{\mathbb{E}[g''(e^{-T}v - e^T e_T \mathcal{V}(0, \mu) + \Gamma_T)]}{1 - e_T \mathbb{E}[g'(e^{-T}X_0 - e^T e_T \mathcal{V}(0, \mu) + \Gamma_T)]} \right. \\ & \quad \left. + \frac{e_T \mathbb{E}[g'(e^{-T}v - e^T e_T \mathcal{V}(0, \mu) + \Gamma_T)] \mathbb{E}[g''(e^{-T}v - e^T e_T \mathcal{V}(0, \mu) + \Gamma_T)]}{(1 - e_T \mathbb{E}[g'(e^{-T}X_0 - e^T e_T \mathcal{V}(0, \mu) + \Gamma_T)]^2} \right). \end{aligned}$$

The above formulas can be extended to any time $t \in [0, T]$ (instead of the sole $t = 0$). It suffices to replace T by $T - t$, e_T by e_{T-t} , Γ_T by Γ_{T-t} and then $\mathcal{V}(0, \mu)$ by $\mathcal{V}(t, \mu)$. By (4) and (5), it holds that

$$\mathcal{V}(t, \mu) = -e^{-(T-t)} \int_{\mathbb{R}} \mathbb{E}[g(e^{-(T-t)}v - e^{T-t}e_{T-t}\mathcal{V}(t, \mu) + \Gamma_{T-t})] d\mu(v),$$

This prompts us to let

$$\Phi(t, y) = -e^{-(T-t)} \int_{\mathbb{R}} \mathbb{E}[g(e^{-(T-t)}v - e^{T-t}e_{T-t}y + \Gamma_{T-t})] d\mu(v),$$

for $t \in [0, T]$ and $y \in \mathbb{R}$. Obviously, Φ is continuously differentiable with respect to y and

$$\frac{\partial \Phi}{\partial y}(t, y) = e_{T-t} \int_{\mathbb{R}} \mathbb{E}[g'(e^{-(T-t)}v - e^{T-t}e_{T-t}y + \Gamma_{T-t})] d\mu(v) \leq 0.$$

Notice also from Itô's formula that

$$\begin{aligned} & \Phi(t, y) \\ &= -e^{-(T-t)} \left(\int_{\mathbb{R}} \mathbb{E}[g(e^{-(T-t)}v - e^{T-t}e_{T-t}y)] d\mu(v) \right. \\ & \quad \left. + \frac{1}{2} \int_0^{T-t} \int_{\mathbb{R}} \mathbb{E}[g''(e^{-(T-t)}v - e^{T-t}e_{T-t}y + e^{s-(T-t)}\Gamma_s)] e^{2(s-(T-t))} d\mu(v) \right), \end{aligned}$$

so that Φ is continuously differentiable with respect to t . Observing that

$$\mathcal{V}(t, \mu) = \Phi(t, \mathcal{V}(t, \mu)),$$

we deduce from the implicit function theorem that, for any $\mu \in \mathcal{P}_2(\mathbb{R})$, the function $\mathcal{V}(\cdot, \mu)$ is continuously differentiable. Also, for any $t \in [0, T]$,

$$\frac{\partial}{\partial t} \mathcal{V}(t, \mu) = \left(1 - \frac{\partial \Phi}{\partial y}(t, \mathcal{V}(t, \mu)) \right)^{-1} \frac{\partial \Phi}{\partial t}(t, \mathcal{V}(t, \mu)).$$

Since \mathcal{V} is Lipschitz-continuous in the measure argument, we deduce that it is in fact jointly continuous. Also, all the derivatives we computed above are jointly continuous in all their arguments (including time). \square

3.4. Form of the master equation. Returning to (1) and recalling that $(\tilde{Y}_t = e^t \bar{Y}_t)_{0 \leq t \leq T}$, and that $(Y_t = X_t + \tilde{Y}_t)_{0 \leq t \leq T}$, we finally end up with the fact that

$$Y_0 = X_0 + \mathcal{V}(0, \mathcal{L}(X_0)),$$

and more generally, by regarding X_t as the initial condition when the system is restricted to the interval $[t, T]$, we also have

$$(6) \quad Y_t = X_t + \tilde{Y}_t = X_t + \mathcal{V}(t, \mathcal{L}(X_t)),$$

which prompts us to let

$$\mathcal{U}(t, x, \mu) = x + \mathcal{V}(t, \mu), \quad (t, x, \mu) \in [0, T] \times \mathbb{R} \times \mathcal{P}_2(\mathbb{R}).$$

We then claim:

Proposition 6. *Assume that \mathcal{V} satisfies the conclusion of Proposition 5, then the function \mathcal{U} satisfies the PDE, set on $[0, T] \times \mathbb{R} \times \mathcal{P}(\mathbb{R})$,*

$$(7) \quad \begin{aligned} & \partial_t \mathcal{U}(t, x, \mu) - \mathcal{U}(t, x, \mu) \partial_x \mathcal{U}(t, x, \mu) + \frac{1}{2} \partial_x^2 \mathcal{U}(t, x, \mu) \\ & - \int_{\mathbb{R}} \mathcal{U}(t, v, \mu) \partial_\mu \mathcal{U}(t, x, \mu)(v) d\mu(v) + \frac{1}{2} \int_{\mathbb{R}} \partial_v \partial_\mu \mathcal{U}(t, x, \mu)(v) d\mu(v) \\ & + x - \int_{\mathbb{R}} f(v) d\mu(v) = 0, \end{aligned}$$

for $(t, x, \mu) \in [0, T] \times \mathbb{R} \times \mathcal{P}_2(\mathbb{R})$, with the terminal boundary condition

$$\mathcal{U}(T, x, \mu) = x - \int_{\mathbb{R}} g(v) d\mu(v).$$

Equation (7) is called the master equation of the system (1). Notice that, because of the special form of \mathcal{U} , the partial derivatives of \mathcal{U} in x here simplify.

Proof. The proof of Proposition 6 consists in expanding (6) by means of a suitable version of the chain rule for flows of marginal distributions of a diffusion process, see Lemma 7, and by identifying the expansion with the backward equation in (1). \square

The following lemma is taken from [5, 7].

Lemma 7. *Let $(\chi_t)_{0 \leq t \leq T}$ be a real-valued Itô process of the form*

$$d\chi_t = \beta_t dt + dW_t,$$

where $(\beta_t)_{0 \leq t \leq T}$ is a real-valued adapted process satisfying

$$\mathbb{E} \int_0^T \beta_t^2 dt < \infty.$$

Let $\mathcal{U} : [0, T] \times \mathbb{R} \times \mathcal{P}(\mathbb{R}) \rightarrow \mathbb{R}$ satisfy the assumption of Proposition 6. Then, the process $(\mathcal{U}(t, X_t, \mathcal{L}(X_t)))_{0 \leq t \leq T}$ expands as an Itô process, namely

$$\begin{aligned} d[\mathcal{U}(t, X_t, \mathcal{L}(X_t))] &= \left(\partial_t \mathcal{U}(t, X_t, \mathcal{L}(X_t)) + \partial_x \mathcal{U}(t, X_t, \mathcal{L}(X_t)) \beta_t + \frac{1}{2} \partial_x^2 \mathcal{U}(t, X_t, \mathcal{L}(X_t)) \right. \\ &\quad \left. + \mathbb{E}[\partial_\mu \mathcal{U}(t, x, \mathcal{L}(X_t))(X_t) \beta_t]_{x=X_t} + \frac{1}{2} \mathbb{E}[\partial_v \partial_\mu \mathcal{U}(t, x, \mathcal{L}(X_t))(X_t)]_{x=X_t} \right) dt \\ &\quad + \partial_x \mathcal{U}(t, X_t, \mathcal{L}(X_t)) dW_t, \end{aligned}$$

for $t \in [0, T]$.

4. CONVERGENCE OF THE FINITE GAMES

We now address the validity of the limiting formulation. To do so, we focus on the case when the equilibria in the finite player game are taken over open loop strategies, see Definition 1.

Throughout this section, we assume that the conclusion of Proposition 5 holds true.

4.1. First-order condition. Actually, we are going to simplify once more the problem. Instead of addressing the convergence of the equilibria, we shall address the convergence of the solutions to the first-order condition that any equilibria must satisfy. Equivalently, this amounts to say that we shall not address the existence of the equilibria.

We use the same setting as in Section 2. Namely, we are given a collection of N real-valued Brownian motions $((W_t^i)_{0 \leq t \leq T})_{i=1, \dots, N}$, where T is some finite time horizon and N is a fixed integer. These processes are constructed on a probability space $(\Omega, \mathcal{F}, \mathbb{P})$ and the filtration generated by the noises and possibly augmented with an initial σ -field \mathcal{F}_0 , independent of $((W_t^i)_{0 \leq t \leq T})_{i=1, \dots, N}$, is denoted by $(\mathcal{F}_t)_{0 \leq t \leq T}$. Also, the initial conditions $(X_0^i)_{i=1, \dots, N}$ are assumed to be \mathcal{F}_0 -measurable, square-integrable, independent and identically distributed.

By the game version of the Pontryagin principle, see for instance [5, Chap. 2], any equilibria must satisfy the following forward-backward system:

$$\begin{aligned} dX_t^i &= -Y_t^i dt + dW_t^i, \\ (8) \quad dY_t^i &= -\left(X_t^i - \int_{\mathbb{R}} f(x) d\bar{\mu}_t^{N,i}(x) \right) dt + \sum_{j=1}^N Z_t^{i,j} dW_t^j, \end{aligned}$$

for $t \in [0, T]$ and $i = 1, \dots, N$, and with the terminal boundary condition

$$Y_T^i = X_T^i - \int_{\mathbb{R}} g(x) d\bar{\mu}_T^{N,i}(x).$$

As already explained, the solutions of the backward equations are required to be adapted and the martingale integrands to be square-integrable. By using the same kind of change of variable as in Section 2, we can “remove”

the linear term in x in the backward component; thanks to [8], we can deduce that the above system is uniquely solvable when f and g are bounded and Lipschitz continuous.

Following the strategy introduced in [3], we let

$$\bar{Y}_t^i = \mathcal{U}(t, X_t^i, \bar{\mu}_t^N), \quad \bar{Z}_t^i = \partial_x \mathcal{U}(t, X_t^i, \bar{\mu}_t^N) = 1, \quad t \in [0, T],$$

for $i \in \{1, \dots, N\}$, with

$$(9) \quad \bar{\mu}_t^N = \frac{1}{N} \sum_{j=1}^N \delta_{X_t^j}, \quad t \in [0, T].$$

The main strategy for addressing the convergence of the solution to (8) (as N tends to ∞) is to consider the distance between $(\bar{Y}_t^i)_{0 \leq t \leq T}$ and $(Y_t^i)_{0 \leq t \leq T}$, which is the purpose of the following claim.

Proposition 8. *There exists a constant C , independent of N and of the realization of the Erdős-Renyi graph, such that*

$$\begin{aligned} & \sup_{0 \leq t \leq T} \left(\frac{1}{N} \sum_{i=1}^N \mathbb{E}[|\bar{Y}_t^i - Y_t^i|^2] \right) \\ & \leq \frac{C}{N} + C \frac{1}{N} \sum_{i=1}^N \mathbb{E} \left[\left(\frac{1}{N_i} \mathbf{1}_{N_i \geq 1} \sum_{j=1}^N (\varepsilon^{i,j} g(X_T^j)) - \frac{1}{N} \sum_{j=1}^N g(X_T^j) \right)^2 \right] \\ & \quad + C \int_0^T \frac{1}{N} \sum_{i=1}^N \mathbb{E} \left[\left(\frac{1}{N_i} \mathbf{1}_{N_i \geq 1} \sum_{j=1}^N (\varepsilon^{i,j} f(X_t^j)) - \frac{1}{N} \sum_{j=1}^N f(X_t^j) \right)^2 \right] dt. \end{aligned}$$

As made clear below, the proof is based on the standard Itô formula, applied to the function:

$$u^i : [0, T] \times \mathbb{R}^d \ni (t, x^1, \dots, x^N) \mapsto \mathcal{U}(t, x^i, \bar{\mu}_x^N), \quad \bar{\mu}_x^N = \frac{1}{N} \sum_{j=1}^N \delta_{x^j}.$$

To do so, we must explain how to compute the (Euclidean) derivatives of u^i in terms of the Wasserstein derivative of \mathcal{U} . The following is taken from [3, 7] or [5].

$$\begin{aligned} \partial_{x^i} u^i(t, x^1, \dots, x^N) &= \partial_x \mathcal{U}(t, x^i, \bar{\mu}_x^N) + \frac{1}{N} \partial_{\mu} \mathcal{U}(t, x^i, \bar{\mu}_x^N)(x^i) \\ \partial_{x^j} u^i(t, x^1, \dots, x^N) &= \frac{1}{N} \partial_{\mu} \mathcal{U}(t, x^i, \bar{\mu}_x^N)(x^j), \quad j \neq i, \\ \partial_{x^i}^2 u^i(t, x^1, \dots, x^N) &= \partial_x^2 \mathcal{U}(t, x^i, \bar{\mu}_x^N) + \frac{1}{N} \partial_v \partial_{\mu} \mathcal{U}(t, x^i, \bar{\mu}_x^N)(x^i) \\ &\quad + \frac{1}{N^2} \partial_{\mu}^2 \mathcal{U}(t, x^i, \bar{\mu}_x^N)(x^i, x^i) \\ &= \frac{1}{N} \partial_v \partial_{\mu} \mathcal{U}(t, x^i, \bar{\mu}_x^N)(x^i) + \frac{1}{N^2} \partial_{\mu}^2 \mathcal{U}(t, x^i, \bar{\mu}_x^N)(x^i, x^i), \end{aligned}$$

and

$$\begin{aligned}\partial_{x^j}^2 u^i(t, x^1, \dots, x^N) &= \frac{1}{N} \partial_v \partial_\mu \mathcal{U}(t, x^i, \bar{\mu}_x^N)(x^j) \\ &\quad + \frac{1}{N^2} \partial_\mu^2 \mathcal{U}(t, x^i, \bar{\mu}_x^N)(x^j, x^j), \quad j \neq i,\end{aligned}$$

where, to compute the second-order derivatives, we used the fact that, in our case, $\partial_x \mathcal{U}(t, x, \mu)$ is constant (equal to 1) and $\partial_\mu \mathcal{U}(t, x, \mu)(v)$ is independent of x .

Proof. By applying Itô's formula and by implementing the PDE satisfied by \mathcal{U} at point $(t, x^i, \bar{\mu}_x^N)$, we get

$$\begin{aligned}d\bar{Y}_t^i &= \partial_x \mathcal{U}(t, X_t^i, \bar{\mu}_t^N) \left(\mathcal{U}(t, X_t^i, \bar{\mu}_t^N) - Y_t^i \right) dt \\ &\quad + \frac{1}{N} \sum_{j=1}^N \partial_\mu \mathcal{U}(t, X_t^i, \bar{\mu}_t^N)(X_t^j) \left(\mathcal{U}(t, X_t^j, \bar{\mu}_t^N) - Y_t^j \right) dt \\ &\quad + \frac{1}{2N^2} \sum_{j=1}^N \partial_\mu^2 \mathcal{U}(t, X_t^i, \bar{\mu}_t^N)(X_t^j, X_t^j) dt \\ &\quad - \left(X_t^i - \int_{\mathbb{R}} f(x) d\bar{\mu}_t^N \right) dt \\ &\quad + \bar{Z}_t^i dW_t^i + \frac{1}{N} \sum_{j=1}^N \partial_\mu \mathcal{U}(t, X_t^i, \bar{\mu}_t^N)(X_t^j) dW_t^j.\end{aligned}$$

Inserting the backward equation satisfied by $(Y_t^i)_{0 \leq t \leq T}$ and recalling that $(\bar{Y}_t^i = \mathcal{U}(t, X_t^i, \bar{\mu}_t^N))_{0 \leq t \leq T}$, we obtain

$$\begin{aligned}d(\bar{Y}_t^i - Y_t^i) &= \partial_x \mathcal{U}(t, X_t^i, \bar{\mu}_t^N) (\bar{Y}_t^i - Y_t^i) dt \\ &\quad + \frac{1}{N} \sum_{j=1}^N \partial_\mu \mathcal{U}(t, X_t^i, \bar{\mu}_t^N)(X_t^j) (\bar{Y}_t^j - Y_t^j) dt \\ &\quad + \frac{1}{2N^2} \sum_{j=1}^N \partial_\mu^2 \mathcal{U}(t, X_t^i, \bar{\mu}_t^N)(X_t^j, X_t^j) dt \\ &\quad - \left(\frac{1}{N_i} \sum_{j=1}^N \varepsilon^{i,j} f(X_t^j) - \frac{1}{N} \sum_{j=1}^N f(X_t^j) \right) dt \\ &\quad + \sum_{j=1}^N \left(\bar{Z}_t^i \mathbf{1}_{i=j} + \frac{1}{N} \partial_\mu \mathcal{U}(t, X_t^i, \bar{\mu}_t^N)(X_t^j) - Z_t^{i,j} \right) dW_t^j.\end{aligned}$$

Taking the square and then the expectation, we deduce that there exists a constant C , independent of N and i , such that, for all $t \in [0, T]$,

$$\begin{aligned} d\mathbb{E}[|\bar{Y}_t^i - Y_t^i|^2] &\leq C\mathbb{E}[|\bar{Y}_t^i - Y_t^i|^2]dt \\ &\quad + C\frac{1}{N}\sum_{j=1}^N\mathbb{E}[|\bar{Y}_t^j - Y_t^j|^2]dt \\ &\quad + C\left\{\frac{1}{N} + \mathbb{E}\left[\left(\frac{1}{N_i}\sum_{j=1}^N\varepsilon^{i,j}f(X_t^j) - \frac{1}{N}\sum_{j=1}^Nf(X_t^j)\right)^2\right]\right\}dt. \end{aligned}$$

where we used the fact that the derivatives of \mathcal{U} are bounded.

Recalling that $\bar{Y}_T^i = \mathcal{U}(T, X_T^i, \bar{\mu}_T^N) = X_T^i - \int_{\mathbb{R}} g(x) d\bar{\mu}_T^N(x)$, taking the mean over $i \in \{1, \dots, N\}$ and appealing to Gronwall's lemma, we obtain

$$\begin{aligned} (10) \quad &\sup_{0 \leq t \leq T} \left(\frac{1}{N} \sum_{i=1}^N \mathbb{E}[|\bar{Y}_t^i - Y_t^i|^2] \right) \\ &\leq \frac{C}{N} + C\frac{1}{N}\sum_{i=1}^N\mathbb{E}\left[\left(\frac{1}{N_i}\sum_{j=1}^N(\varepsilon^{i,j}g(X_T^j)) - \frac{1}{N}\sum_{j=1}^Ng(X_T^j)\right)^2\right] \\ &\quad + C\int_0^T\frac{1}{N}\sum_{i=1}^N\mathbb{E}\left[\left(\frac{1}{N_i}\sum_{j=1}^N(\varepsilon^{i,j}f(X_t^j)) - \frac{1}{N}\sum_{j=1}^Nf(X_t^j)\right)^2\right]dt, \end{aligned}$$

which completes the proof. \square

4.2. Almost sure analysis on the graph. We now compute (with a similar computation for the term driven by f):

$$\begin{aligned} (11) \quad &\frac{1}{N}\sum_{i=1}^N\mathbb{E}\left[\left(\frac{1}{N_i}\sum_{j=1}^N(\varepsilon^{i,j}g(X_T^j)) - \frac{1}{N}\sum_{j=1}^Ng(X_T^j)\right)^2\right] \\ &= \mathbb{E}\left[\frac{1}{N}\sum_{i=1}^N\left(\frac{\mathbf{1}_{N_i \geq 1}}{N_i^2}\sum_{j,\ell=1}^N(\varepsilon^{i,j}\varepsilon^{i,\ell}g(X_T^j)g(X_T^\ell)) \right. \right. \\ &\quad \left. \left. + \frac{1}{N^2}\sum_{j,\ell=1}^Ng(X_T^j)g(X_T^\ell)\right)\right] \\ &\quad - \frac{2}{N^2}\mathbb{E}\left[\sum_{i=1}^N\frac{\mathbf{1}_{N_i \geq 1}}{N_i}\sum_{j,\ell=1}^N(\varepsilon^{i,j}g(X_T^j)g(X_T^\ell))\right]. \end{aligned}$$

Observe that

$$\begin{aligned}
 (12) \quad & \frac{1}{N} \sum_{i=1}^N \frac{\mathbf{1}_{N_i \geq 1}}{N_i^2} \sum_{j,\ell=1}^N (\varepsilon^{i,j} \varepsilon^{i,\ell} g(X_T^j) g(X_T^\ell)) \\
 &= \frac{1}{N} \sum_{j,\ell=1}^N \left[\left(\sum_{i=1}^N \frac{\mathbf{1}_{N_i \geq 1}}{N_i^2} \varepsilon^{i,j} \varepsilon^{i,\ell} \right) g(X_T^j) g(X_T^\ell) \right] \\
 &= \frac{1}{N^2} \sum_{j,\ell=1}^N \left[\left(\frac{1}{N} \sum_{i=1}^N \frac{N^2}{N_i^2} \mathbf{1}_{N_i \geq 1} \varepsilon^{i,j} \varepsilon^{i,\ell} \right) g(X_T^j) g(X_T^\ell) \right],
 \end{aligned}$$

and,

$$\begin{aligned}
 (13) \quad & \frac{2}{N^2} \sum_{i=1}^N \frac{\mathbf{1}_{N_i \geq 1}}{N_i} \sum_{j,\ell=1}^N (\varepsilon^{i,j} g(X_T^j) g(X_T^\ell)) \\
 &= \frac{2}{N^2} \sum_{j,\ell=1}^N \left[\left(\sum_{i=1}^N \frac{\mathbf{1}_{N_i \geq 1}}{N_i} \varepsilon^{i,j} \right) g(X_T^j) g(X_T^\ell) \right] \\
 &= \frac{2}{N^2} \sum_{j,\ell=1}^N \left[\left(\frac{1}{N} \sum_{i=1}^N \frac{N}{N_i} \mathbf{1}_{N_i \geq 1} \varepsilon^{i,j} \right) g(X_T^j) g(X_T^\ell) \right].
 \end{aligned}$$

We claim:

Proposition 9. *On $(\Xi, \mathcal{G}, \mathbf{P})$, consider a sequence $(\varepsilon^{i,j})_{j \geq i \geq 1}$ of independent and identically distributed Bernoulli random variables of parameter $p \in (0, 1)$. Let $\varepsilon^{i,i} = 0$ and $\varepsilon^{i,j} = \varepsilon^{j,i}$ if $i > j$.*

Then, \mathbf{P} almost surely,

$$\begin{aligned}
 \lim_{N \rightarrow \infty} \frac{1}{N^2} \sum_{j,\ell=1}^N \left| 1 - \frac{1}{N} \sum_{i=1}^N \frac{N^2}{N_i^2} \mathbf{1}_{N_i \geq 1} \varepsilon^{i,j} \varepsilon^{i,\ell} \right| &= 0, \\
 \lim_{N \rightarrow \infty} \frac{1}{N} \sum_{j=1}^N \left| 1 - \frac{1}{N} \sum_{i=1}^N \frac{N}{N_i} \mathbf{1}_{N_i \geq 1} \varepsilon^{i,j} \right| &= 0,
 \end{aligned}$$

where

$$N_i = \frac{1}{N} \sum_{j=1}^N \varepsilon^{i,j}.$$

Proof. Fix $N \geq 2$ and observe, for a given $i \in \{1, \dots, N\}$, that the random variables $(\varepsilon^{i,j})_{j \neq i}$ are independent and identically distributed. By Hoeffding's inequality,

$$(14) \quad \forall \varepsilon > 0, \quad \mathbf{P} \left[\left| \frac{N_i}{N-1} - p \right| \geq \varepsilon \right] \leq \exp(-2(N-1)\varepsilon^2).$$

First Step. By (14), we have

$$\forall t > 0, \quad \mathbf{P} \left[\sqrt{N-1} \left| \frac{N_i}{N-1} - p \right| \geq t \right] \leq \exp(-2t^2),$$

from which we deduce that, for any integer $q \geq 1$,

$$\mathbf{E}\left[\left|\frac{N_i}{N-1} - p\right|^q\right] \leq \frac{C_q}{N^{q/2}},$$

for a constant C_q only depending on q . Hence,

$$(15) \quad \begin{aligned} \mathbf{E}\left[\left|\frac{N_i}{N} - p\right|^q\right] &\leq 2^{q-1} \mathbf{E}\left[\left|\frac{N_i}{N-1} - p\right|^q\right] + 2^{q-1} \mathbf{E}\left[\left|\frac{N_i}{N} - \frac{N_i}{N-1}\right|^q\right] \\ &\leq \frac{C_q}{N^{q/2}}, \end{aligned}$$

where we allowed the constant C_q to vary from line to line.

Second Step. By (14), we also have, for any large enough real $k \geq 1$, precisely for $1/k \leq p/2$,

$$\begin{aligned} \mathbf{P}\left[\frac{N-1}{N_i} \mathbf{1}_{N_i \geq 1} \geq k\right] &\leq \mathbf{P}\left[\frac{N_i}{N-1} \leq \frac{1}{k}\right] \\ &\leq \mathbf{P}\left[\left|\frac{N_i}{N-1} - p\right| \geq \frac{p}{2}\right] \\ &\leq \exp\left(-\frac{(N-1)p^2}{2}\right). \end{aligned}$$

In particular, for any integer $q \geq 1$,

$$\begin{aligned} \mathbf{E}\left[\left(\frac{N-1}{N_i}\right)^q \mathbf{1}_{N_i \geq 1}\right] &\leq q \int_0^N k^{q-1} \mathbf{P}\left[\frac{N-1}{N_i} \mathbf{1}_{N_i \geq 1} \geq k\right] dk \\ &\leq q \int_0^{2/p} k^{q-1} dk + q \int_0^N k^{q-1} \exp\left(-\frac{(N-1)p^2}{2}\right) dk \\ &\leq \frac{2^q}{p^q} + N^q \exp\left(-\frac{(N-1)p^2}{2}\right). \end{aligned}$$

Allowing the constant C_q to depend on p , we obtain

$$(16) \quad \mathbf{E}\left[\left(\frac{N}{N_i}\right)^q \mathbf{1}_{N_i \geq 1}\right] \leq C_q.$$

By (15) and (16) and by Cauchy-Schwarz inequality,

$$\begin{aligned} \mathbf{E}\left[\left|\frac{N}{N_i} \mathbf{1}_{N_i \geq 1} - \frac{1}{p}\right|^q\right] &= \mathbf{E}\left[\left(\frac{N \mathbf{1}_{N_i \geq 1}}{p N_i}\right)^q \left|p - \frac{N_i}{N}\right|^q\right] + \frac{1}{p^q} \mathbf{P}[N_i = 0] \\ &\leq \mathbf{E}\left[\left(\frac{N \mathbf{1}_{N_i \geq 1}}{p N_i}\right)^q \left|p - \frac{N_i}{N}\right|^q\right] + \frac{1}{p^q} \mathbf{P}\left[\left|\frac{N_i}{N-1} - p\right| \geq p\right] \\ &\leq \frac{C_q}{N^{q/2}}, \end{aligned}$$

where we made use of (14) again to derive the last term in the second line.

Third Step. We now prove the second claim in the statement of Proposition 9.

$$\begin{aligned}
& \left| \frac{1}{N} \sum_{j=1}^N \left| 1 - \frac{1}{N} \sum_{i=1}^N \frac{N}{N_i} \mathbf{1}_{N_i \geq 1} \varepsilon^{i,j} \right| \right| \\
& \leq \frac{1}{N} \sum_{j=1}^N \left| \frac{1}{N} \sum_{i=1}^N \left(\frac{N}{N_i} \mathbf{1}_{N_i \geq 1} - \frac{1}{p} \right) \varepsilon^{i,j} \right| + \frac{1}{pN} \sum_{j=1}^N \left| \frac{N_j}{N} - p \right| \\
& \leq \frac{1}{N} \sum_{i=1}^N \left| \frac{N}{N_i} \mathbf{1}_{N_i \geq 1} - \frac{1}{p} \right| + \frac{1}{pN} \sum_{j=1}^N \left| \frac{N_j}{N} - p \right|.
\end{aligned}$$

Taking the power q and then the expectation under \mathbf{P} and appealing to the conclusions of the two first steps, we get

$$\mathbf{E} \left[\left(\frac{1}{N} \sum_{j=1}^N \left| 1 - \frac{1}{N} \sum_{i=1}^N \frac{N}{N_i} \mathbf{1}_{N_i \geq 1} \varepsilon^{i,j} \right| \right)^q \right] \leq \frac{C_q}{N^{q/2}},$$

which suffices to conclude.

Fourth Step. In order to tackle the first claim, we first observe that

$$\mathbf{E} \left[\left| \frac{N^2}{N_i^2} \mathbf{1}_{N_i \geq 1} - \frac{1}{p^2} \right|^q \right] \leq \mathbf{E} \left[\left| \frac{N}{N_i} \mathbf{1}_{N_i \geq 1} + \frac{1}{p} \right|^{2q} \right]^{1/2} \mathbf{E} \left[\left| \frac{N}{N_i} \mathbf{1}_{N_i \geq 1} - \frac{1}{p} \right|^{2q} \right]^{1/2}.$$

By (16) and the conclusion of the second step, we deduce that

$$\mathbf{E} \left[\left| \frac{N^2}{N_i^2} \mathbf{1}_{N_i \geq 1} - \frac{1}{p^2} \right|^q \right] \leq \frac{C_q}{N^{q/2}}.$$

We also note that, for any given $j \neq \ell$, the random variables $(\varepsilon^{i,\ell} \varepsilon^{i,j})_{i \neq j, i \neq \ell}$ are independent and follow the same Bernoulli distribution of parameter p^2 . Therefore, proceeding as in the first step, we have

$$\mathbf{E} \left[\left| \frac{1}{N} \sum_{i=1}^N \varepsilon^{i,\ell} \varepsilon^{i,j} - p^2 \right|^q \right] \leq \frac{C_q}{N^{q/2}}.$$

In order to complete the proof, it remains to see, in analogy with the third step, that

$$\begin{aligned}
& \left| \frac{1}{N^2} \sum_{j,\ell=1}^N \left| 1 - \frac{1}{N} \sum_{i=1}^N \frac{N^2}{N_i^2} \mathbf{1}_{N_i \geq 1} \varepsilon^{i,j} \varepsilon^{i,\ell} \right| \right| \\
& \leq \frac{1}{N} \sum_{i=1}^N \left| \frac{N^2}{N_i^2} \mathbf{1}_{N_i \geq 1} - \frac{1}{p^2} \right| + \frac{1}{N^2 p^2} \sum_{j,\ell=1}^N \left| p^2 - \frac{1}{N} \sum_{i=1}^N \varepsilon^{i,j} \varepsilon^{i,\ell} \right|.
\end{aligned}$$

Taking the power q and then the expectation under \mathbf{P} ,

$$\mathbf{E} \left[\left(\frac{1}{N^2} \sum_{j,\ell=1}^N \left| 1 - \frac{1}{N} \sum_{i=1}^N \frac{N^2}{N_i^2} \mathbf{1}_{N_i \geq 1} \varepsilon^{i,j} \varepsilon^{i,\ell} \right| \right)^q \right] \leq \frac{C_q}{N^{q/2}},$$

which yields the result. \square

4.3. Final statement. From now on, we assume that $(\Omega, \mathcal{F}, \mathbb{P})$ is equipped with an infinite sequence of independent Brownian motions $((W_t^i)_{0 \leq t \leq T})_{i \geq 1}$ and with an infinite sequence of identically distributed and independent initial conditions $(X_0^i)_{i \geq 1}$, both sequences being required to be independent.

Here then is the final statement.

Theorem 10. *Assume that f and g are bounded and Lipschitz continuous and that the conclusion of Proposition 5 holds true.*

On $(\Xi, \mathcal{G}, \mathbf{P})$, consider a sequence $(\varepsilon^{i,j})_{j > i \geq 1}$ of independent and identically distributed Bernoulli random variables of parameter $p \in (0, 1)$. Let $\varepsilon^{i,i} = 0$ and $\varepsilon^{i,j} = \varepsilon^{j,i}$ if $i > j$.

Then, \mathbf{P} almost surely,

$$\lim_{N \rightarrow \infty} \sup_{0 \leq t \leq T} \mathbb{E}[\mathcal{W}_2(\bar{\mu}_t^N, \mu_t)^2] = 0,$$

where \mathcal{W}_2 denotes the 2-Wasserstein distance, $(\bar{\mu}_t^N)_{0 \leq t \leq T}$ is as in (9) and $(\mu_t)_{0 \leq t \leq T}$ is equal to $(\mathcal{L}(X_t))_{0 \leq t \leq T}$, $(X_t)_{0 \leq t \leq T}$ denoting the forward component of the solution to (1).

Proof. By (11), (12) and (13), and by Proposition 9, we deduce that there exists a sequence of random variables $(\delta_N)_{N \geq 1}$ constructed on $(\Xi, \mathcal{G}, \mathbf{P})$ such that

$$\mathbf{P}\left[\lim_{N \rightarrow \infty} \delta_N = 0\right] = 1,$$

and

$$\frac{1}{N} \sum_{i=1}^N \mathbb{E}\left[\left(\frac{1}{N_i} \sum_{j=1}^N (\varepsilon^{i,j} g(X_T^j)) - \frac{1}{N} \sum_{j=1}^N g(X_T^j)\right)^2\right] \leq \delta_N.$$

Proceeding similarly with the terms driven by f in (10), we deduce that

$$(17) \quad \sup_{0 \leq t \leq T} \left(\frac{1}{N} \sum_{i=1}^N \mathbb{E}[|\bar{Y}_t^i - Y_t^i|^2] \right) \leq \delta_N,$$

the details of $(\delta_N)_{N \geq 1}$ being allowed to change from line to line.

For any $i \geq 1$, call $((\hat{X}_t^i)_{0 \leq t \leq T})_{i \geq 1}$ the copy of $(X_t)_{0 \leq t \leq T}$ when driven by $(X_0^i, (W_t^i)_{0 \leq t \leq T})$ instead of $(X_0, (W_t)_{0 \leq t \leq T})$, namely

$$d\hat{X}_t^i = -\mathcal{U}(t, \hat{X}_t^i, \mathcal{L}(\hat{X}_t^i))dt + dW_t^i, \quad t \in [0, T]; \quad \hat{X}_0^i = X_0^i.$$

Then, recalling that $(\bar{Y}_t^i = \mathcal{U}(t, X_t^i, \bar{\mu}_t^N))_{0 \leq t \leq T}$ and using the Lipschitz property of \mathcal{U} in x and μ , we have

$$\mathbb{E}[|\hat{X}_t^i - X_t^i|^2] \leq C \int_0^t \mathbb{E}[|\bar{Y}_s^i - Y_s^i|^2]ds + C \int_0^t \mathbb{E}[\mathcal{W}_2(\bar{\mu}_s^N, \mu_s)^2]ds, \quad t \in [0, T],$$

for a constant C independent of N and t . Above, we used the identity $(\mu_t = \mathcal{L}(\hat{X}_t^i))_{0 \leq t \leq T}$. Taking the mean over i and using (17), we obtain

$$\mathbb{E}[\mathcal{W}_2(\bar{\mu}_t^N, \hat{\mu}_t^N)^2] \leq \delta_N + C \int_0^t \mathbb{E}[\mathcal{W}_2(\bar{\mu}_s^N, \mu_s)^2]ds, \quad t \in [0, T],$$

where we let

$$\hat{\mu}_t^N = \frac{1}{N} \sum_{i=1}^N \delta_{\hat{X}_t^i}, \quad t \in [0, T].$$

This yields

$$\sup_{0 \leq t \leq T} \mathbb{E}[\mathcal{W}_2(\bar{\mu}_t^N, \hat{\mu}_t^N)^2] \leq \delta_N + \int_0^T \mathbb{E}[\mathcal{W}_2(\hat{\mu}_t^N, \mu_t)^2] dt.$$

As $\hat{\mu}_t^N$ in the right-hand side is the empirical distribution of an independent and identically distributed sample of law μ_t , for each $t \in [0, T]$, the term $\mathbb{E}[\mathcal{W}_2(\hat{\mu}_t^N, \mu_t)^2]$ tends to 0 as N tends to ∞ . As $\sup_{0 \leq t \leq T} \mathbb{E}[\mathcal{W}_2(\hat{\mu}_t^N, \mu_t)^2] \leq 2 \sup_{0 \leq t \leq T} \mathbb{E}[|X_t|^2] < \infty$, Lebesgue's dominated convergence theorem shows that the integral tends to 0 as well. Hence,

$$\lim_{N \rightarrow \infty} \sup_{0 \leq t \leq T} \mathbb{E}[\mathcal{W}_2(\bar{\mu}_t^N, \hat{\mu}_t^N)^2] = 0.$$

It remains to observe that, in fact,

$$\lim_{N \rightarrow \infty} \sup_{0 \leq t \leq T} \mathbb{E}[\mathcal{W}_2(\hat{\mu}_t^N, \mu_t)^2] = 0,$$

which follows from a uniform continuity argument. Indeed, for any $0 \leq s < t \leq T$,

$$\begin{aligned} \mathbb{E} \left[\left| \mathcal{W}_2(\hat{\mu}_t^N, \mu_t)^2 - \mathcal{W}_2(\hat{\mu}_s^N, \mu_s)^2 \right| \right] &\leq C \mathbb{E} \left[\left| \mathcal{W}_2(\hat{\mu}_t^N, \mu_t) - \mathcal{W}_2(\hat{\mu}_s^N, \mu_s) \right|^2 \right]^{1/2} \\ &\leq C \mathbb{E} \left[\mathcal{W}_2(\hat{\mu}_t^N, \hat{\mu}_s^N)^2 + \mathcal{W}_2(\mu_t, \mu_s)^2 \right]^{1/2} \\ &\leq C \mathbb{E}[|X_t - X_s|^2] \\ &\leq C(t - s), \end{aligned}$$

the constant C being allowed to vary from line to line. \square

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