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# Singular perturbations for an elliptic operator in the Heisenberg group 

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#### Abstract

We study some classes of singular perturbation problems where the dynamics of the fast variables evolve in the whole space obeying to an infinitesimal operator which is subelliptic and ergodic. We prove that the corresponding ergodic problem admits a solution which is globally Lipschitz continuous and it has at most a logarithmic growth at infinity.

The main result of this paper establishes that as $\epsilon \rightarrow 0$, the value functions of the singular perturbation problems converge locally uniformly to the solution of an effective problem whose operator and data are explicitly given in terms of the invariant measure for the ergodic operator.


Keywords: Subelliptic equations, Heisenberg group, invariant measure, homogenization, singular perturbations, viscosity solutions, degenerate elliptic equations. ${ }^{1}$

## 1 Introduction

This paper is devoted to the asymptotic behaviour as $\epsilon \rightarrow 0$ of stochastic control systems of the form

$$
\begin{aligned}
d X_{t} & =\tilde{\phi}\left(X_{t}, Y_{t}, u_{t}\right) d t+\sqrt{2} \tilde{\sigma}\left(X_{t}, Y_{t}, u_{t}\right) d W_{t}, \quad X_{0}=x \in \mathbb{R}^{n} \\
d Y_{t} & =\frac{1}{\epsilon} b\left(Y_{t}\right) d t+\frac{\sqrt{2}}{\sqrt{\epsilon}} \sigma\left(Y_{t}\right) d W_{t}, \quad Y_{0}=y \in \mathbb{R}^{m}
\end{aligned}
$$

[^0]where $u_{t}$ is a control law, $W_{t}$ is a standard Brownian motion, while the coefficients $\tilde{\phi}, \tilde{\sigma}, b$ and $\sigma$ are Lipschitz continuous uniformly in $u$. Problems of this type arise from models where the variables $Y$ evolve much faster than the variables $X$. We investigate the limit of the value function
$$
V^{\epsilon}(t, x, y):=\sup _{u \in \mathcal{U}} \mathbb{E}\left[\int_{t}^{T} f\left(X_{s}, Y_{s}, u_{s}\right) d s+e^{a(t-T)} g\left(X_{T}\right)\right]
$$
where $\mathbb{E}$ denotes the expectation, $\mathcal{U}$ is the set of progressively measurable processes with values in a compact metric set $U$ and $a$ is a fixed positive parameter.

Our aim will be to characterize the limit of $V^{\epsilon}$ as the solution to an effective Cauchy problem whose operator and terminal data need to be suitably chosen.

The main issue of this paper is to tackle this problem when the coefficients are not periodic in $y$ and the diffusion matrices $\tilde{\sigma} \tilde{\sigma}^{T}$ and $\sigma \sigma^{T}$ may be degenerate and unbounded. For the sake of simplicity, we shall focus our attention to the model case where $\tilde{\sigma} \tilde{\sigma}^{T}$ is bounded and the diffusion matrix $\sigma$ is the one associated to the Heisenberg group in $\mathbb{R}^{3}$

$$
\sigma(y)=\left(\begin{array}{cc}
1 & 0 \\
0 & 1 \\
2 y_{2} & -2 y_{1}
\end{array}\right), \quad \text { for } y=\left(y_{1}, y_{2}, y_{3}\right)
$$

We note that $\sigma \sigma^{T}$ is degenerate and with unbounded coefficients.
It is well known that the ergodicity of the fast variable cannot be expected for general drift $b$. In order to overcome this issue, we consider a drift $b$ in the Ornstein-Uhlenbeck form

$$
b(y)=-\left(k_{1} y_{1}, k_{2} y_{2}, k_{3} y_{3}\right) \quad \text { for some } k_{1}>4, k_{2}>4, k_{3}>0
$$

(see A2)). This choice of the drift is reminiscent of other similar conditions about recurrence of diffusion processes in the whole space (see for example [11] and references therein).

By standard theory (see [9]), the value function $V^{\epsilon}$ is the unique (viscosity) solution to the following Cauchy problem for an Hamilton-JacobiBellman equation

$$
\begin{cases}-\partial_{t} V^{\epsilon}+H\left(x, y, D_{x} V^{\epsilon}, D_{x x}^{2} V^{\epsilon}, \frac{D_{x y}^{2} V^{\epsilon}}{\sqrt{\epsilon}}\right) &  \tag{1.1}\\ -\frac{1}{\epsilon} \mathcal{L}\left(y, D_{y} V^{\epsilon}, D_{y y} V^{\epsilon}\right)+a V^{\epsilon}=0 & \text { in }(0, T) \times \mathbb{R}^{n} \times \mathbb{R}^{3} \\ V^{\epsilon}(T, x, y)=g(x, y) & \text { on } \mathbb{R}^{n} \times \mathbb{R}^{3}\end{cases}
$$

where

$$
\begin{aligned}
H(x, y, p, X, Z) & :=\min _{u \in U}\left\{-\operatorname{tr}\left(\tilde{\sigma} \tilde{\sigma}^{T} X\right)-\tilde{\phi} \cdot p-2 \operatorname{tr}\left(\tilde{\sigma} \sigma^{T} Z\right)-f(x, y, u)\right\} \\
\mathcal{L}(y, q, Y) & :=\operatorname{tr}\left(\sigma \sigma^{T} Y\right)+b \cdot q
\end{aligned}
$$

For the sake of completeness, in order to exhibit the degeneracy and the unboundedness of the operator, we write explicitly the second order term of $\mathcal{L}$ :
(1.2) $\operatorname{tr}\left(\sigma \sigma^{T} D^{2} U\right)=U_{y_{1} y_{1}}+U_{y_{2} y_{2}}+4\left(y_{1}^{2}+y_{2}^{2}\right) U_{y_{3} y_{3}}+4 y_{2} U_{y_{1} y_{3}}-4 y_{1} U_{y_{2} y_{3}}$.

We assume without any loss of generality that $a$ is strictly positive; actually, for $a \leq 0$, the function $W^{\epsilon}(t, x, y)=e^{-A(T-t)} V^{\epsilon}(t, x, y)$, with $A>-a$, satisfies the same Cauchy problem but with a positive coefficient of the 0 -th order term.

Our aim is to establish that, as $\epsilon \rightarrow 0^{+}$, the function $V^{\epsilon}$ converges locally uniformly to a function $V=V(t, x)$ (namely, independent of $y$ ) which can be characterized as the unique (viscosity) solution to the effective Cauchy problem

$$
\begin{cases}-\partial_{t} V+\bar{H}\left(x, D_{x} V, D_{x x}^{2} V\right)+a V=0 & \text { in }(0, T) \times \mathbb{R}^{n}  \tag{1.3}\\ V(T, x)=\bar{g}(x) & \text { on } \mathbb{R}^{n}\end{cases}
$$

where, for every $(x, p, X)$, the effective Hamiltonian $\bar{H}(x, p, X)$ and the effective terminal datum are given by

$$
\begin{align*}
\bar{H}(x, p, X) & :=\int_{\mathbb{R}^{3}} H(x, y, p, X, 0) d \mu(y)  \tag{1.4}\\
\bar{g}(x) & :=\int_{\mathbb{R}^{3}} g(x, y) d \mu(y) \tag{1.5}
\end{align*}
$$

and $\mu$ is the invariant measure of the diffusion process with infinitesimal generator $-\mathcal{L}$. As a matter of facts, $\bar{H}(x, p, X)$ is the ergodic constant $\lambda$ of the cell problem

$$
\begin{equation*}
-\operatorname{tr}\left(\sigma(y) \sigma^{T}(y) D^{2} w(y)\right)-b(y) D w(y)+H(x, y, p, X, 0)=\lambda \quad y \in \mathbb{R}^{3} \tag{1.6}
\end{equation*}
$$

while $\bar{g}(x)$ is the constant obtained in the long time behaviour of the parabolic Cauchy problem

$$
\partial_{t} w^{\prime}-\mathcal{L} w^{\prime}=0 \quad \text { in }(0, \infty) \times \mathbb{R}^{3}, \quad w^{\prime}(0, y)=g(x, y) \quad \text { on } \mathbb{R}^{3}
$$

There is a large literature on singular perturbation problems: see $[1$, $12,13]$ and references therein. We shall follow a pure PDE-approach. In this framework, the singular perturbation problems are strictly related to homogenization problems (see also [18]); Alvarez and Bardi [1, 2] extended to singular perturbation problems with periodic fast variables the celebrated perturbed test function method by Evans (see also [3] for some cases in hypoelliptic periodic setting). Let us also recall that, the papers [5, 6, 11] studied singular perturbation problems of uniformly elliptic operators on the whole space.

The novelties of our results is that the variable $Y_{t}$ is unbounded and the diffusion matrix of the fast variable may be degenerate and unbounded. In other words, the main issues to overcome are the lack of periodicity and the degeneracy of the operator. The proof of our main Theorem 2.1 is not an adaptation to the subelliptic case of some arguments already known in the non degenerate case. Indeed our proof is based on the perturbed test function method suitably adapted with a Lyapunov function. Moreover our techniques shed some light on some difficult points in the literature on the whole space.

Let us recall that existence and uniqueness of the ergodic constant $\lambda$ for (1.6) (namely, that $\delta u_{\delta}$ locally converge to $\lambda$, where $u_{\delta}$ solves the approximated cell problem (3.1) below) and the stabilization to a constant have been established in our previous paper [16]. Unfortunately, by the lack of compactness for $y$, these properties seem to be not sufficient for applying the usual semilimits method for the convergence of $V^{\epsilon}$. In order to overcome this issue:

- we shall prove that the cell problem admits a corrector $w$ which is globally Lipschitz continuous and it has at most a logarithmic type growth at infinity;
- under some additional assumptions we get that the corrector $w$ is $C^{2, \alpha}\left(\mathbb{R}^{3}\right)$.
- we take advantage of the existence and uniqueness of the invariant measure and a superlinear Lyapunov function for the operator $\mathcal{L}$.

In our opinion, the proof of the global Lipschitz continuity of the corrector has its own interest because it can be extended to many other operators in unbounded domains. In this direction, let us quote the papers [11] and [15] where similar results are obtained for strictly elliptic operators.

Moreover the $C^{2}$-regularity of the corrector is not straightforward because our operator contains second order horizontal derivatives and Euclidean first derivatives as well and such a second order part of the operator does not immediately regularize the first order one.

The paper is organized as follows: in Section 2 we state the perturbation problem and our main convergence result. Section 3 is devoted to the solution of the cell problem and its properties. In Section 4, by means of these result we prove the convergence of $V^{\epsilon}$ to $V$.

## 2 The convergence result

Throughout this paper unless otherwise explicitly stated, we assume
A1) the diffusion matrix $\sigma$ has the following form:

$$
\sigma(y)=\left(\begin{array}{cc}
1 & 0 \\
0 & 1 \\
2 y_{2} & -2 y_{1}
\end{array}\right), \quad \text { for } y=\left(y_{1}, y_{2}, y_{3}\right)
$$

$A 2)$ the drift is $b(y)=-\left(k_{1} y_{1}, k_{2} y_{2}, k_{3} y_{3}\right)$ with $k_{1}>4, k_{2}>4, k_{3}>0$;
A3) the function $f=f(x, y, u)$ is Lipschitz continuous in $(x, y)$ uniformly in $u$ and, for some $C_{f}>0$, it satisfies

$$
|f(x, y, u)| \leq C_{f}(1+|x|) \quad \forall(x, y, u) \in \mathbb{R}^{n} \times \mathbb{R}^{3} \times U
$$

A4) the function $g$ is continuous in $(x, y)$ and there exits $C_{g}$ such that

$$
|g(x, y)| \leq C_{g}(1+|x|) \quad \forall(x, y) \in \mathbb{R}^{n} \times \mathbb{R}^{3} ;
$$

A5) $\tilde{\phi}(x, y, u)$ and $\tilde{\sigma}(x, y, u)$ are Lipschitz continuous and bounded in $(x, y)$ uniformly on $u:|\tilde{\phi}(x, y, u)| \leq C_{\tilde{\phi}},|\tilde{\sigma}(x, y, u)| \leq C_{\tilde{\sigma}}$;

A6) for any $(x, p, X)$ the function $F(y)=-H(x, y, p, X, 0)$ is such that $F, \frac{\partial F}{\partial y_{3}}$ and $\frac{\partial^{2} F}{\partial y_{3}^{2}}$ are bounded and globally Lipschitz.

## Remark 2.1

- Let us underline that condition (A2) is linked with the form of the second order operator given in (1.2). This assumption will play a crucial role in several main points, for instance in (3.4), Theorem 3.2, Theorem 3.4 and in (4.9).
- We note also that there holds:

$$
\begin{align*}
& \left|H(x, y, p, X, Z)-H\left(x^{\prime}, y, p^{\prime}, X^{\prime}, Z\right)\right|  \tag{2.1}\\
& \quad \leq C\left(\left|p-p^{\prime}\right|+\left|X-X^{\prime}\right|\right)+C\left|x-x^{\prime}\right|(1+|p|+|X|)
\end{align*}
$$

- We stress that assumption A6) will be only used for obtaining the regularity of the corrector.

We state now that the problem (1.1) is well posed and the solution $V^{\epsilon}$ has a sublinear growth in the slow variable.

Proposition 2.1 Under Assumptions (A1)-(A5), for any $\epsilon>0$ there exists a unique continuous viscosity solution $V^{\epsilon}$ to problem (1.1) such that

$$
\begin{equation*}
\left|V^{\epsilon}(t, x, y)\right| \leq C_{0}(1+|x|), \quad \forall(t, x, y) \in(0, T) \times \mathbb{R}^{n} \times \mathbb{R}^{n} \tag{2.2}
\end{equation*}
$$

for some positive constant $C_{0}$ independent on $\epsilon$. In particular $\left\{V^{\epsilon}\right\}_{\epsilon}$ is a family of locally equibounded functions.

Proof. The uniqueness follows from the comparison principle proved in Da Lio-Ley [8] (recall that they require that the diffusion matrix and the drift grow at most quadratically and respectively linearly with respect to the state). We now claim that there exist a supersolution $w^{+}$and a subsolution $w^{-}$such that $\left|w^{ \pm}(x)\right| \leq C(1+|x|)$ for $|x|$ sufficiently large. We shall prove the existence of $w^{+}$and we shall omit the analogous arguments for $w^{-}$.
Let $w_{0} \in C^{\infty}\left(\mathbb{R}^{n}\right)$ be a function in $x$ such that

$$
w_{0}=C_{1}(1+|x|) \quad \text { for }|x| \geq R \geq 1, \quad w_{0}(x) \geq g(x, y) \quad \forall(x, y),
$$

for some positive constants $C_{1}$ and $R$. For $|x| \geq R$, there holds

$$
\begin{aligned}
& -\partial_{t} w_{0}+H\left(x, y, D_{x} w_{0}, D_{x x}^{2} w_{0}, \frac{D_{x y}^{2} w_{0}}{\sqrt{\epsilon}}\right)-\frac{1}{\epsilon} \mathcal{L}\left(y, D_{y} w_{0}, D_{y y} w_{0}\right)+a w_{0}= \\
& H\left(x, y, D w_{0}, D^{2} w_{0}, 0\right)+a w_{0}= \\
& \min _{u \in U}\left\{-\operatorname{tr}\left(\tilde{\sigma} \tilde{\sigma}^{T} D^{2} w_{0}\right)-\tilde{\phi} \cdot D w_{0}-f(x, y)\right\}+a w_{0}(x) \geq \\
& \quad-C_{2} C_{1}+a C_{1}(1+|x|)-C_{f}(1+|x|)
\end{aligned}
$$

where $C_{2}$ depends on $C_{\tilde{\phi}}$ and $C_{\tilde{\sigma}}$. Choosing $C_{1}$ big enough (for instance $C_{1} \geq \frac{C_{f}+1}{a}$ ) and $R$ big enough (for instance $R \geq C_{2} C_{1}$ ) the function $w_{0}$ is a supersolution (1.1) for $|x|>R$. Eventually adding a new positive constant $C_{3}$ big enough, $w^{+}=w_{0}+C_{3}$ is a supersolution in the whole $(0, T) \times \mathbb{R}^{n} \times \mathbb{R}^{3}$ which amounts to our claim. In conclusion, applying Perron's method, we infer the existence of a solution to (1.1) verifying (2.2).

The main purpose of this paper is to prove the following
Theorem 2.1 Under assumptions A1)-A6), the solution $V^{\epsilon}$ of (1.1) converges uniformly on the compact subsets of $(0, T) \times \mathbb{R}^{n} \times \mathbb{R}^{3}$ to the unique viscosity solution $V$ of (1.3) where $\bar{H}$ and $\bar{g}$ are defined in (1.4) and respectively in (1.5).

## 3 The cell problem

In this section we prove that there exists an unique constant $\lambda$ such that the cell problem (1.6) admits solutions. We shall also prove the existence of a solution $w$ which is globally Lipschitz continuous and with log-growth at infinity. Assuming also A6) we prove that $w \in C^{2}$. This solution $w$ will play a crucial role in the proof of Theorem 2.1.

### 3.1 Approximated cell problems

In order to solve the cell problem (1.6), it is expedient to introduce the approximated problems

$$
\begin{equation*}
\delta u_{\delta}-\operatorname{tr}\left(\sigma(y) \sigma^{T}(y) D^{2} u_{\delta}\right)-b(y) D u_{\delta}=F(y) \quad \text { in } \mathbb{R}^{3}, \tag{3.1}
\end{equation*}
$$

where $\delta>0$ and $F(y)=-H(x, y, p, X, 0)$ with $(x, p, X)$ fixed. In this section the results are obtained for a general function $F(y)$ which satisfies:

$$
\begin{equation*}
F(y) \text { is continuous and bounded in } \mathbb{R}^{3} \text {. } \tag{3.2}
\end{equation*}
$$

Note that under assumptions (A1)-(A5), for $(x, p, X)$ fixed, the function $F(\cdot)=-H(x, \cdot, p, X, 0)$, satisfies assumption (3.2).

Let us recall from [16] some properties of the operator $\mathcal{L}$ and functions $u_{\delta}$ (and we refer the reader to this paper for the detailed proof).

Lemma 3.1 There exists a unique invariant measure $\mu$ associated to the operator $-\mathcal{L}$; moreover

$$
\mathcal{L}^{*} \mu=0, \quad \mu>0, \quad \mu \in C^{\infty}\left(\mathbb{R}^{3}\right)
$$

where $\mathcal{L}^{*}$ is the adjoint operator of $-\mathcal{L}$.
Remark 3.1 As a byproduct of [16], we have the following estimate on the decay of $\mu$ at infinity:

$$
\int_{\mathbb{R}^{3}}\left(y_{1}^{4}+y_{2}^{4}+y_{3}^{2}\right) d \mu(y)<+\infty
$$

Actually, by [16, Prop. 2.1 (proof)], the function $w(y):=\left(y_{1}^{4}+y_{2}^{4}\right) / 12+y_{3}^{2} / 2$ satisfies [16, eq. (2.10)] with $\phi>k\left(y_{1}^{4}+y_{2}^{4}+y_{3}^{2}\right)$ (for $k>0$ sufficiently small). Hence, for this choice of $\phi$, relations [16, equation (2.14)] and [16, equation (2.18)] hold true. Letting $\rho \rightarrow 0^{+}$we get: $\int_{\mathbb{R}^{3}} \phi d \mu<+\infty$.

Lemma 3.2 Under assumptions (A1),(A2),(3.2), there exists an unique smooth solution $u_{\delta}$ of the approximating problem (3.1) such that

$$
\left|u_{\delta}(y)\right| \leq \frac{C}{\delta} \quad \forall y \in \mathbb{R}^{3}
$$

for some positive constant $C$ independent of $\delta$. Moreover the functions $\delta u_{\delta}$ are locally uniformly Hölder continuous, i.e. there exists $\alpha \in(0,1)$ such that for every compact $K \subset \mathbb{R}^{3}$ there exists a constant $N$ such that

$$
\left|\delta u_{\delta}\left(y_{1}\right)-\delta u_{\delta}\left(y_{2}\right)\right| \leq N\left|y_{1}-y_{2}\right|^{\alpha}, \quad \forall y_{1}, y_{2} \in K, \quad \forall \delta \in(0,1)
$$

The constant $N$ only depends on $K$ and on the data of the problem (in particular is independent of $\delta$ ).

Theorem 3.1 The solution $u_{\delta}$ of problem (3.1) given in Lemma 3.2 satisfies

$$
\lim _{\delta \rightarrow 0} \delta u_{\delta}=\int_{\mathbb{R}^{3}} F(y) d \mu(y)
$$

where $\mu$ is the invariant measure of $-\mathcal{L}$ established in Lemma 3.1.

### 3.2 Global Lipschitz continuity.

In this section we derive the global Lipschitz continuity of the solution $u_{\delta}$ of (3.1) from its continuity under the weaker assumption of an at most linear growth of $F$. In our opinion, this result has its own interest. We assume:

$$
\left\{\begin{array}{l}
(A 1)-(A 2)  \tag{3.3}\\
F \text { Lipschitz continuous in } \mathbb{R}^{3} \text { with Lipschitz constant } L \\
|F(y)| \leq C_{F}(|y|+1) \quad \forall y \in \mathbb{R}^{3}
\end{array}\right.
$$

It is clear that the globally Lipschitz continuity implies that $|F(y)| \leq$ $L(|y|+1) \quad \forall y \in \mathbb{R}^{3}$, but in the following proof we want to underline separately the dependence on the Lipschitz continuity and the linear behaviour at infinity.

Lemma 3.3 Under assumptions (3.3) there exists a constant $C$ such that

$$
\begin{equation*}
\left|u_{\delta}(y)\right| \leq C\left(|y|+\frac{1}{\delta}\right), \quad y \in \mathbb{R}^{3} . \tag{3.4}
\end{equation*}
$$

Proof. The comparison principle for equation (3.1) comes from Da Lio Ley ([8]) (this is true also for elliptic operators, see [6]). The existence of a continuous viscosity solution $u_{\delta}$ in $\mathbb{R}^{3}$ comes from Perron Theorem, by finding sub- and super solution in $\mathbb{R}^{3}$. We remark that, in the case when $F$ is bounded by a constant $C_{F}$ in $\mathbb{R}^{3}$, a trivial supersolution is $\frac{C_{F}}{\delta}$ (a subsolution $-\frac{C_{F}}{\delta}$ ) and the result easily follows.
In the more general, sublinear case, let us introduce some constants:

$$
\begin{equation*}
l:=\min \left\{k_{1}-4, k_{2}-4, k_{3}\right\}>0 \tag{3.5}
\end{equation*}
$$

$$
\begin{align*}
& C_{l}:=\max \left\{1, \frac{2 C_{F}}{l}\right\}  \tag{3.6}\\
& r_{0}:=1+\frac{2 C_{l}}{C_{F}} \tag{3.7}
\end{align*}
$$

Let $U_{0}$ any regular function $U_{0} \in C^{2}\left(\mathbb{R}^{3}\right)$ such that $U_{0}(y)=|y|+1$ in $\mathbb{R}^{3} \backslash B_{r_{0}}$. (For instance $U_{0}(r)=C_{0}+C_{2} r^{2}+C_{4} r^{4}$ in $\overline{B\left(0, r_{0}\right)}$ and $U_{0}(r)=r+1$ in $\left.\mathbb{R}^{3} \backslash B_{r_{0}}\right)$.
There exists a constant $M_{0}$ independent of $\delta \in[0,1]$ such that

$$
\begin{equation*}
\delta U_{0}-\mathcal{L} U_{0}-F(y) \geq-M_{0} \quad \text { in } B_{r_{0}} \tag{3.8}
\end{equation*}
$$

Let us define the function $U \in C^{2}\left(\mathbb{R}^{3}\right)$ as

$$
U(\cdot)=C_{l}\left(U_{0}(\cdot)+\frac{2 M_{0}}{\delta}\right)
$$

We claim that $U$ is a supersolution to (3.1). Let us test the supersolution property first in $B_{r_{0}}$ then in $\mathbb{R}^{3} \backslash B_{r_{0}}$. For $y \in B_{r_{0}}$, thanks to $C_{l} \geq 1$ (see (3.6)) and the definition of $M_{0}$ in (3.8), we have:

$$
\begin{aligned}
\delta U-\mathcal{L} U-F(y) & =C_{l}\left(\delta U_{0}+2 M_{0}-\mathcal{L} U_{0}\right)-F(y) \\
& \geq\left(C_{l}-1\right) F(y)+C_{l} M_{0} \geq 0
\end{aligned}
$$

if $M_{0}$ is sufficiently large. For $y \in \mathbb{R}^{3} \backslash B_{r_{0}}$, denote $r=|y|$

$$
\begin{aligned}
& \delta U(y)-\mathcal{L}\left(y, D U(y), D^{2} U(y)\right)-F(y) \\
&= C_{l}\left(\delta\left(|y|+1+\frac{2 M_{0}}{\delta}\right)-\frac{2+4\left(y_{1}^{2}+y_{2}^{2}\right)}{r}+\frac{\left(k_{1} y_{1}^{2}+k_{2} y_{2}^{2}+k_{3} y_{3}^{2}\right)}{r}\right)+ \\
&+C_{l} \frac{1}{r^{3}}\left(y_{1}^{2}+y_{2}^{2}\right)\left(1+4 y_{3}^{2}\right)-F(y) \\
& \geq C_{l}\left(-\frac{2}{r}+\frac{\left.\left(k_{1}-4\right) y_{1}^{2}+\left(k_{2}-4\right) y_{2}^{2}+k_{3} y_{3}^{2}\right)}{r}\right)-F(y) \\
& \geq C_{l}\left(-\frac{2}{r}+l r\right)-C_{F}(r+1) \\
& \geq 0
\end{aligned}
$$

where we used (3.5), (3.6) and (3.7). From the comparison principle then $u_{\delta} \leq U(y)=C_{l}\left(U_{0}+\frac{2 M_{0}}{\delta}\right) \leq C\left(|y|+\frac{1}{\delta}\right)$. The same method applies to define a subsolution and to prove that $u_{\delta} \geq-C\left(|y|+\frac{1}{\delta}\right)$.

Theorem 3.2 Under assumptions (3.3), let $u_{\delta}$ be the unique continuous solution of (3.1) which satisfies (3.4). there holds

$$
\left|u_{\delta}\left(y^{\prime}\right)-u_{\delta}(y)\right| \leq \psi\left(\left|y^{\prime}-y\right|\right) \quad \forall y, y^{\prime} \in \mathbb{R}^{3}
$$

where $\psi \in C^{2}(\mathbb{R})$ is a concave increasing function with $\psi(0)=0$ and $\psi^{\prime}>\max \left\{\frac{L}{k_{1}-4}, \frac{L}{k_{2}-4}, \frac{L}{k_{3}}\right\}$ (recall that $L$ is the Lipschitz constant of $F$ ) and it is independent of $\delta$. In particular there holds

$$
\begin{equation*}
\left|u_{\delta}\left(y^{\prime}\right)-u_{\delta}(y)\right| \leq \bar{L}\left(\left|y^{\prime}-y\right|\right) \quad \forall y, y^{\prime} \in \mathbb{R}^{3}, \forall \delta>0 \tag{3.9}
\end{equation*}
$$

for $\bar{L}>\frac{L}{l}$, where $l$ is defined in (3.5).

Proof. For each $\eta>0$, we introduce the function

$$
\Psi(x, y)=u(x)-u(y)-\psi(|x-y|)-\eta|x|^{2}-\eta|y|^{2}
$$

where $\psi$ is a function as in the statement and for simplicity we take $u:=u_{\delta}$.
Assume for the moment that there holds

$$
\begin{equation*}
\Psi(x, y) \leq \frac{8 \eta}{\delta} \quad \forall x, y \in \mathbb{R}^{3}, \eta \in(0,1) \tag{3.10}
\end{equation*}
$$

then, for any $x, y \in \mathbb{R}^{3}$, as $\eta \rightarrow 0^{+}$, we obtain the following inequality

$$
u(x)-u(y) \leq \psi(|x-y|)
$$

which is equivalent to the statement because of the arbitrariness of $x$ and $y$.
Let us now prove inequality (3.10); to this end, we shall proceed by contradiction. Let $(\bar{x}, \bar{y})$ be the maximum point of function $\Psi$ in $\mathbb{R}^{3} \times \mathbb{R}^{3}$. This maximum does exist since from (3.4) we have that $\lim _{x \rightarrow+\infty} \frac{u(x)}{|x|^{2}}=0$.

Let us assume by contradiction that

$$
\begin{equation*}
\Psi(\bar{x}, \bar{y})=u(\bar{x})-u(\bar{y})-\psi(|\bar{x}-\bar{y}|)-\eta|\bar{x}|^{2}-\eta|\bar{y}|^{2}>\frac{8 \eta}{\delta} . \tag{3.11}
\end{equation*}
$$

Clearly, the points $\bar{x}$ and $\bar{y}$ cannot coincide, otherwise (3.11) is false. We set $\tilde{\psi}(x, y):=\psi(|x-y|)+\eta\left(|x|^{2}+|y|^{2}\right)$ and we invoke [7, Theorem 3.2]: for every $\rho>0$ there exist two symmetric $3 \times 3$ matrices $X$ and $Y$ such that

$$
\begin{gather*}
\left(p_{x}, X\right) \in \mathcal{J}^{2,+} u(\bar{x}), \quad\left(p_{y}, Y\right) \in \mathcal{J}^{2,-} u(\bar{y}),  \tag{3.12}\\
\left(\begin{array}{cc}
X & 0 \\
0 & -Y
\end{array}\right) \leq A+\rho A^{2}, \tag{3.13}
\end{gather*}
$$

where
$p_{x}:=D_{x} \tilde{\psi}(\bar{x}, \bar{y}), \quad p_{y}:=-D_{y} \tilde{\psi}(\bar{x}, \bar{y}), \quad A:=\left(\begin{array}{cc}D_{x x}^{2} \tilde{\psi}(\bar{x}, \bar{y}) & D_{x y}^{2} \tilde{\psi}(\bar{x}, \bar{y}) \\ D_{y x}^{2} \tilde{\psi}(\bar{x}, \bar{y}) & D_{y y}^{2} \tilde{\psi}(\bar{x}, \bar{y})\end{array}\right)$.
We write explicitly $p_{x}, p_{y}$ and $A$ :

$$
\begin{align*}
& p_{x}=\psi^{\prime}(|\bar{x}-\bar{y}|) \frac{\bar{x}-\bar{y}}{|\bar{x}-\bar{y}|}+2 \eta \bar{x}=\psi^{\prime}(|\bar{x}-\bar{y}|) q+2 \eta \bar{x}  \tag{3.14}\\
& p_{y}=\psi^{\prime}(|\bar{x}-\bar{y}|) \frac{\bar{x}-\bar{y}}{|\bar{x}-\bar{y}|}-2 \eta \bar{y}=\psi^{\prime}(|\bar{x}-\bar{y}|) q-2 \eta \bar{y} \tag{3.15}
\end{align*}
$$

where we defined

$$
\begin{equation*}
q:=\frac{\bar{x}-\bar{y}}{|\bar{x}-\bar{y}|} . \tag{3.16}
\end{equation*}
$$

Defining $B:=\frac{I-q \otimes q}{|\bar{x}-\bar{y}|}$, the matrix $A$ assumes the following form

$$
\begin{array}{r}
A=\psi^{\prime}(|\bar{x}-\bar{y}|)\left(\begin{array}{cc}
B & -B \\
-B & B
\end{array}\right)+\psi^{\prime \prime}(|\bar{x}-\bar{y}|)\left(\begin{array}{cc}
q \otimes q & -q \otimes q \\
-q \otimes q & q \otimes q
\end{array}\right)  \tag{3.17}\\
+2 \eta\left(\begin{array}{cc}
I & 0 \\
0 & I
\end{array}\right) .
\end{array}
$$

From the definition of sub and supersolution and $\left(p_{x}, X\right),\left(p_{y}, Y\right)$, we have

$$
\begin{aligned}
\delta u(\bar{x})-\operatorname{tr}\left(\sigma(\bar{x}) \sigma^{T}(\bar{x}) X\right)-b(\bar{x}) p_{x} & \leq F(\bar{x}) \\
\delta u(\bar{y})-\operatorname{tr}\left(\sigma(\bar{y}) \sigma^{T}(\bar{y}) Y\right)-b(\bar{y}) p_{y} & \geq F(\bar{y})
\end{aligned}
$$

Subtracting the latter inequality from the former, we infer

$$
\begin{align*}
\delta(u(\bar{x})-u(\bar{y}))-\operatorname{tr}(\sigma(\bar{x}) & \left.\sigma^{T}(\bar{x}) X-\sigma(\bar{y}) \sigma^{T}(\bar{y}) Y\right)  \tag{3.18}\\
& +\left(-b(\bar{x}) p_{x}+b(\bar{y}) p_{y}\right) \leq F(\bar{x})-F(\bar{y})
\end{align*}
$$

We want to estimate from below the three terms on the left hand side of (3.18):
i) $\mathcal{U}:=\delta(u(\bar{x})-u(\bar{y}))$,
ii) $\mathcal{T}:-\operatorname{tr}\left(\sigma(\bar{x}) \sigma^{T}(\bar{x}) X-\sigma(\bar{y}) \sigma^{T}(\bar{y}) Y\right)$,
iii) $\mathcal{G}:=-b(\bar{x}) p_{x}+b(\bar{y}) p_{y}$.
(i). The assumption by contradiction (3.11) yields

$$
\begin{equation*}
\mathcal{U}:=\delta(u(\bar{x})-u(\bar{y})) \geq \delta \psi(|\bar{x}-\bar{y}|)+\delta\left(\eta|\bar{x}|^{2}+\eta|\bar{y}|^{2}\right)+8 \eta \geq 8 \eta . \tag{3.19}
\end{equation*}
$$

(ii). Multiplying relation (3.13) by $(\zeta, \xi)$ where $\zeta$ and $\xi$ are vectors in $\mathbb{R}^{3}$ we obtain $(\zeta, \xi)\left(\begin{array}{cc}X & 0 \\ 0 & -Y\end{array}\right)(\zeta, \xi)^{T} \leq(\zeta, \xi) A(\zeta, \xi)^{T}+\rho(\zeta, \xi) A^{2}(\zeta, \xi)^{T}$. Then,
using (3.17), we have

$$
\begin{align*}
\zeta X \zeta-\xi Y \xi & \leq \psi^{\prime}(|\bar{x}-\bar{y}|)(\zeta B \zeta-\xi B \zeta-\zeta B \xi+\xi B \xi)+  \tag{3.20}\\
& +\psi^{\prime \prime}(|\bar{x}-\bar{y}|)(<\zeta-\xi, q>)^{2}+2 \eta\left(|\zeta|^{2}+|\xi|^{2}\right)+\rho a(\zeta, \xi)
\end{align*}
$$

where we denoted by $a(\zeta, \xi):=(\zeta, \xi) A^{2}(\zeta, \xi)^{T}$ and $q$ is defined in (3.16).
Recall that, for any choice of two orthonormal basis $\left\{e_{i}\right\}_{i=1,2}$ and $\left\{\tilde{e}_{i}\right\}_{i=1,2}$ in $\mathbb{R}^{2}$, (if $e_{i}$ is a orthonormal basis, $\operatorname{tr} M=\sum_{i=1}^{2} e_{i} M e_{i}^{T}$ ) we have

$$
\begin{aligned}
& \operatorname{tr}\left(\sigma \sigma^{T} X\right)=\operatorname{tr}\left(\sigma^{T} X \sigma\right)=\sum_{i=1}^{2} e_{i} \sigma^{T} X \sigma e_{i}^{T} \\
& \operatorname{tr}\left(\sigma \sigma^{T} Y\right)=\operatorname{tr}\left(\sigma^{T} Y \sigma\right)=\sum_{i=1}^{2} \tilde{e}_{i} \sigma^{T} X \sigma \tilde{e}_{i}^{T} .
\end{aligned}
$$

We choose

$$
\begin{equation*}
\zeta_{i}=e_{i} \sigma^{T}(\bar{x}), \quad \xi_{i}=\tilde{e}_{i} \sigma^{T}(\bar{y}) \tag{3.21}
\end{equation*}
$$

hence, we get
$\mathcal{T}=-\sum_{i=1}^{2} e_{i} \sigma^{T}(\bar{x}) X \sigma(\bar{x}) e_{i}^{T}+\sum_{i=1}^{2} \tilde{e}_{i} \sigma^{T}(\bar{y}) Y \sigma(\bar{y}) \tilde{e}_{i}^{T}=-\left(\sum_{i=1}^{2} \zeta_{i} X \zeta_{i}-\sum_{i=1}^{2} \xi_{i} Y \xi_{i}\right)$.
Then from inequality (3.20) we obtain

$$
\begin{aligned}
\mathcal{T} \geq & -\psi^{\prime}(|\bar{x}-\bar{y}|) \sum_{i=1}^{2}\left(\zeta_{i}-\xi_{i}\right) B\left(\zeta_{i}-\xi_{i}\right)-\psi^{\prime \prime}(|\bar{x}-\bar{y}|) \sum_{i=1}^{2}\left(\left\langle\zeta_{i}-\xi_{i}, q\right\rangle\right)^{2} \\
& -2 \eta \sum_{i=1}^{2}\left(\left|\zeta_{i}\right|^{2}+\left|\xi_{i}\right|^{2}\right)-\rho \sum_{i=1}^{2} a\left(\zeta_{i}, \xi_{i}\right) .
\end{aligned}
$$

From the definition of the matrix $B$ we have

$$
\begin{align*}
& \mathcal{T} \geq-\frac{\psi^{\prime}}{|\bar{x}-\bar{y}|} \sum_{i=1}^{2}\left(\left|\zeta_{i}-\xi_{i}\right|^{2}\right)+\left(\frac{\psi^{\prime}}{|\bar{x}-\bar{y}|}-\psi^{\prime \prime}\right) \sum_{i=1}^{2}\left(<\zeta_{i}-\xi_{i}, q>\right)^{2}  \tag{3.22}\\
&-2 \eta \sum_{i=1}^{2}\left(\left|\zeta_{i}\right|^{2}+\left|\xi_{i}\right|^{2}\right)-\rho \sum_{i=1}^{2} a\left(\zeta_{i}, \xi_{i}\right) \\
& \geq-\frac{\psi^{\prime}}{|\bar{x}-\bar{y}|} \sum_{i=1}^{2}\left(\left|\zeta_{i}-\xi_{i}\right|\right)^{2}-2 \eta \sum_{i=1}^{2}\left(\left|\zeta_{i}\right|^{2}+\left|\xi_{i}\right|^{2}\right)-\rho \sum_{i=1}^{2} a\left(\zeta_{i}, \xi_{i}\right)
\end{align*}
$$

where the last inequality was obtained taking into account that $\psi$ is increasing and concave, so $\frac{\psi^{\prime}}{|\bar{x}-\bar{y}|}-\psi^{\prime \prime} \geq 0$.
(iii). From expressions (3.14) and (3.15) of $p_{x}$ and $p_{y}$, we have

$$
\mathcal{G}=(-b(\bar{x})+b(\bar{y})) \psi^{\prime}(|\bar{x}-\bar{y}|) q+2 \eta(-b(\bar{x}) \bar{x}-b(\bar{y}) \bar{y})
$$

By our choice of $b$ (recall: $\left.b(x)=\left(-k_{1} x_{1},-k_{2} x_{2},-k_{3} x_{3}\right)\right), \mathcal{G}$ becomes

$$
\begin{array}{r}
\mathcal{G}=\left(k_{1}\left(\bar{x}_{1}-\bar{y}_{1}\right)^{2}+k_{2}\left(\bar{x}_{2}-\bar{y}_{2}\right)^{2}+k_{3}\left(\bar{x}_{3}-\bar{y}_{3}\right)^{2}\right) \frac{\psi^{\prime}(|\bar{x}-\bar{y}|)}{|\bar{x}-\bar{y}|}+  \tag{3.23}\\
2 \eta\left(k_{1}\left(\bar{x}_{1}^{2}+\bar{y}_{1}^{2}\right)+k_{2}\left(\bar{x}_{2}^{2}+\bar{y}_{2}^{2}\right)+k_{3}\left(\bar{x}_{3}^{2}+\bar{y}_{3}^{2}\right)\right)
\end{array}
$$

Now, replacing inequalities (3.19), (3.22) and (3.23) in (3.18), we obtain

$$
\begin{gathered}
L|\bar{x}-\bar{y}| \geq F(\bar{x})-F(\bar{y}) \geq \mathcal{U}+\mathcal{T}+\mathcal{G} \geq \\
8 \eta-\frac{\psi^{\prime}}{|\bar{x}-\bar{y}|} \sum_{i=1}^{2}\left(\left|\zeta_{i}-\xi_{i}\right|\right)^{2}-2 \eta \sum_{i=1}^{2}\left(\left|\zeta_{i}\right|^{2}+\left|\xi_{i}\right|^{2}\right)-\rho \sum_{i=1}^{2} a\left(\zeta_{i}, \xi_{i}\right)+ \\
\left(k_{1}\left(\bar{x}_{1}-\bar{y}_{1}\right)^{2}+k_{2}\left(\bar{x}_{2}-\bar{y}_{2}\right)^{2}+k_{3}\left(\bar{x}_{3}-\bar{y}_{3}\right)^{2}\right) \frac{\psi^{\prime}}{|\bar{x}-\bar{y}|}+ \\
2 \eta\left(k_{1}\left(\bar{x}_{1}^{2}+\bar{y}_{1}^{2}\right)+k_{2}\left(\bar{x}_{2}^{2}+\bar{y}_{2}^{2}\right)+k_{3}\left(\bar{x}_{3}^{2}+\bar{y}_{3}^{2}\right)\right)
\end{gathered}
$$

Passing to the limit as $\rho \rightarrow 0^{+}$, we obtain

$$
\begin{gathered}
\quad(3.24) \quad L \bar{x}-\bar{y} \mid \geq \\
\eta\left(8-2 \sum_{i=1}^{2}\left(\left|\zeta_{i}\right|^{2}+\left|\xi_{i}\right|^{2}\right)+2 k_{1}\left(\bar{x}_{1}^{2}+\bar{y}_{1}^{2}\right)+2 k_{2}\left(\bar{x}_{2}^{2}+\bar{y}_{2}^{2}\right)+2 k_{3}\left(\bar{x}_{3}^{2}+\bar{y}_{3}^{2}\right)\right)+ \\
\left.+\frac{\psi^{\prime}}{|\bar{x}-\bar{y}|}\left[-\sum_{i=1}^{2}\left(\left|\zeta_{i}-\xi_{i}\right|\right)^{2}+k_{1}\left(\bar{x}_{1}-\bar{y}_{1}\right)^{2}+k_{2}\left(\bar{x}_{2}-\bar{y}_{2}\right)^{2}+k_{3}\left(\bar{x}_{3}-\bar{y}_{3}\right)^{2}\right)\right]
\end{gathered}
$$

The contradiction is easily obtained choosing as the two orthonormal basis the canonical basis in $\mathbb{R}^{2}, e_{1}=\tilde{e}_{1}=(1,0)$ and $e_{2}=\tilde{e}_{2}=(0,1)$. Then the vectors $\zeta_{i}$ and $\xi_{i}$ (see (3.21)), with $\bar{x}=\left(\bar{x}_{1}, \bar{x}_{2}, \bar{x}_{3}\right)$ and $\bar{y}=\left(\bar{y}_{1}, \bar{y}_{2}, \bar{y}_{3}\right)$ become $\zeta_{1}=\left(1,0,2 \bar{x}_{2}\right), \zeta_{2}=\left(1,0,-2 \bar{x}_{1}\right), \xi_{1}=\left(1,0,2 \bar{y}_{2}\right), \xi_{2}=\left(1,0,-2 \bar{y}_{1}\right)$, and

$$
\begin{gathered}
\left|\zeta_{1}\right|^{2}=1+4 \bar{x}_{2}^{2},\left|\zeta_{2}\right|^{2}=1+4 \bar{x}_{1}^{2},\left|\xi_{1}\right|^{2}=1+4 \bar{y}_{2}^{2},\left|\xi_{2}\right|^{2}=1+4 \bar{y}_{1}^{2} \\
\left|\zeta_{1}-\xi_{1}\right|^{2}=4\left(\bar{x}_{2}-\bar{y}_{2}\right)^{2},\left|\zeta_{2}-\xi_{2}\right|^{2}=4\left(\bar{x}_{1}-\bar{y}_{1}\right)^{2}
\end{gathered}
$$

Hence, relation (3.24) becomes

$$
\begin{aligned}
& L|\bar{x}-\bar{y}| \geq \eta\left[8-8-8\left(\bar{x}_{1}^{2}+\bar{y}_{1}^{2}+\bar{x}_{2}^{2}+\bar{y}_{2}^{2}\right)+\right. \\
& \left.2 k_{1}\left(\bar{x}_{1}^{2}+\bar{y}_{1}^{2}\right)+2 k_{2}\left(\bar{x}_{2}^{2}+\bar{y}_{2}^{2}\right)+2 k_{3}\left(\bar{x}_{3}^{2}+\bar{y}_{3}^{2}\right)\right]+ \\
& \left.\frac{\psi^{\prime}}{|\bar{x}-\bar{y}|}\left[\left(k_{1}-4\right)\left(\bar{x}_{1}-\bar{y}_{1}\right)^{2}+\left(k_{2}-4\right)\left(\bar{x}_{2}-\bar{y}_{2}\right)^{2}\right)+k_{3}\left(\bar{x}_{3}-\bar{y}_{3}\right)^{2}\right] \\
& \geq 2 \eta\left[\left(k_{1}-4\right)\left(\bar{x}_{1}^{2}+\bar{y}_{1}^{2}\right)+\left(k_{2}-4\right)\left(\bar{x}_{2}^{2}+\bar{y}_{2}^{2}\right)+k_{3}\left(\bar{x}_{3}^{2}+\bar{y}_{3}^{2}\right)\right]+ \\
& \frac{\psi^{\prime}}{|\bar{x}-\bar{y}|}\left[\left(k_{1}-4\right)\left(\bar{x}_{1}-\bar{y}_{1}\right)^{2}+\left(k_{2}-4\right)\left(\bar{x}_{2}-\bar{y}_{2}\right)^{2}+k_{3}\left(\bar{x}_{3}-\bar{y}_{3}\right)^{2}\right] .
\end{aligned}
$$

By our choice of $k_{1}, k_{2}$ and $k_{3}$ in (A2) (namely, $k_{1}, k_{2}>4, k_{3}>0$ ) we get $L|\bar{x}-\bar{y}|^{2} \geq \psi^{\prime}(|\bar{x}-\bar{y}|)\left[\left(k_{1}-4\right)\left(\bar{x}_{1}-\bar{y}_{1}\right)^{2}+\left(k_{2}-4\right)\left(\bar{x}_{2}-\bar{y}_{2}\right)^{2}+k_{3}\left(\bar{x}_{3}-\bar{y}_{3}\right)^{2}\right]$, thus we obtain a contradiction provided that we choose a function $\psi$ such that

$$
\psi^{\prime}>\max \left\{\frac{L}{k_{1}-4}, \frac{L}{k_{2}-4}, \frac{L}{k_{3}}\right\} .
$$

Hence, the proof of our claim (3.10) is accomplished. The second statement of the theorem easily follows by taking $\psi(z)=\bar{L} z$, with $\bar{L}>\max \left\{\frac{L}{k_{1}-4}, \frac{L}{k_{2}-4}, \frac{L}{k_{3}}\right\}$.

Remark 3.2 Similar arguments can be applied to other matrices still related to degenerate elliptic operators as, for example, in dimension 2:

$$
\sigma(y):=\left(\sigma_{i j}(y)\right)_{i, j} \text { with } \sigma_{i j}(y)=a_{i j} y_{1}+b_{i j} y_{2}+c_{i j}
$$

which in particular encompasses the Ornstein-Uhlenbeck operator and Grushin operator, respectively

$$
\sigma_{O U}=\left(\begin{array}{cc}
1 & 0 \\
0 & 0
\end{array}\right), \quad \sigma_{G}=\left(\begin{array}{cc}
1 & 0 \\
0 & y
\end{array}\right)
$$

For the Grushin operator in a forthcoming paper [17] we will obtain a local Hölder continuity uniform in $\delta$ using a technique introduced in [10].

### 3.3 A key estimate on the growth of the approximate corrector

The aim of this section is to establish that the solution to the approximating cell problem has a logarithmic growth at infinity. Our arguments are borrowed form [11, Proposition 3.2].

Lemma 3.4 Assume (A1) and (A2). Let $u_{\delta}(y)$ be the solution of equation (3.1) with (3.2). There exists $C>0$ such that

$$
\left|u_{\delta}(y)-u_{\delta}(0)\right| \leq C\left[1+\log \left(\left(y_{1}^{2}+y_{2}^{2}\right)^{2}+y_{3}^{2}+1\right)\right] \quad \forall y \in \mathbb{R}^{3}, \delta \in(0,1)
$$

Proof. We can argue as in [11, Proposition 3.2], replacing its Lemma 3.3 with our Theorem 3.2; to this end, our first step is to claim that, for $C_{1}$ and $R$ sufficiently large, the function $g(y):=C_{1} \log \left(\left(y_{1}^{2}+y_{2}^{2}\right)^{2}+y_{3}^{2}\right)$ is a supersolution to (3.1) in $\mathbb{R}^{3} \backslash B_{R}$. Indeed, by equality (1.2) there holds:

$$
\begin{aligned}
\operatorname{tr}\left(\sigma \sigma^{T} D^{2} g(y)\right) & =\frac{8 C_{1}\left(y_{1}^{2}+y_{2}^{2}\right)}{\left(y_{1}^{2}+y_{2}^{2}\right)^{2}+y_{3}^{2}} \\
b(y) \cdot D g(y) & =-C_{1} \frac{4\left(y_{1}^{2}+y_{2}^{2}\right)\left(k_{1} y_{1}^{2}+k_{2} y_{2}^{2}\right)+2 k_{3} y_{3}^{2}}{\left(y_{1}^{2}+y_{2}^{2}\right)^{2}+y_{3}^{2}}
\end{aligned}
$$

By these identities, we get

$$
\delta g(y)-\operatorname{tr}\left(\sigma(y) \sigma^{T}(y) D^{2} g\right)-b(y) D g \geq F(y), y \in \mathbb{R}^{3} \backslash B_{R}
$$

provided that $C$ and $R$ are sufficiently large. Now if $\max _{\bar{B}_{R}} u_{\delta} \leq 0$ then we have $\max _{\bar{B}_{R}} u_{\delta} \leq g(y)$ for any $y \in \partial \bar{B}_{R}$. By the comparison principle established in [8], we obtain $u_{\delta} \leq g$ in $\mathbb{R}^{3}$. If $\max _{\bar{B}_{R}} u_{\delta}>0$, we note that $g_{1}(\cdot):=g(\cdot)+\max _{\bar{B}_{R}} u_{\delta}$ is still a supersolution of (3.1) in $\mathbb{R}^{3} \backslash B_{R}$. Hence, again by the comparison principle we have $u_{\delta} \leq g_{1}$ in $\mathbb{R}^{3}$. By Theorem 3.2 we infer: $u_{\delta}(y)-u_{\delta}(0) \leq g_{1}(y)+\bar{L} R$ which gives one of the two inequalities of the statement. The proof of the other one is similar and we shall omit it.

### 3.4 The cell problem

Theorem 3.3 Under assumptions A1)-A5) of Section 2, for every ( $x, p, X$ ) the constant $\lambda=-\int_{\mathbb{R}^{3}} H(\bar{x}, y, \bar{p}, \bar{X}, 0) d \mu(y)$ ( $\mu$ is the invariant measure defined in Lemma 3.1) is the unique constant such that the cell problem (1.6)
admits a solution $w(y)$ which is globally Lipschitz continuous and satisfies the following estimate:

$$
\begin{equation*}
|w(y)-w(0)| \leq C\left[1+\log \left(\left(y_{1}^{2}+y_{2}^{2}\right)^{2}+y_{3}^{2}+1\right)\right] \quad \forall y \in \mathbb{R}^{3} \tag{3.25}
\end{equation*}
$$

Moreover, the solution $w$ is unique up to an additive constant.

Proof. To prove the existence of such a $\lambda$ we argue as in [11, Proposition 3.2], replacing its Lemma 3.3 with our Theorem 3.2. We consider the solution $u_{\delta}$ of the approximated cell problem (3.1), recalling that, from A3), $F(y)$ is bounded in $\mathbb{R}^{3}$; then $w_{\delta}(y):=u_{\delta}(y)-u_{\delta}(0)$ satisfies

$$
\delta w_{\delta}(y)-\operatorname{tr}\left(\sigma(y) \sigma^{T}(y) D^{2} w_{\delta}\right)-b(y) D w_{\delta}=F(y)-\delta u_{\delta}(0)
$$

From the Lipschitz continuity of $u_{\delta}(y)$ in (3.9) we have that

$$
\left|w_{\delta}(y)\right|=\left|u_{\delta}(y)-u_{\delta}(0)\right| \leq \bar{L}|y|
$$

and

$$
\left|w_{\delta}(y)-w_{\delta}(z)\right|=\left|u_{\delta}(y)-u_{\delta}(z)\right| \leq \bar{L}|y-z|
$$

hence $w_{\delta}(y)$ are locally equibounded and equicontinuous. Then by AscoliArzela theorem and standard diagonal argument we can conclude that there exists a function $w$ with the desired properties. Moreover from Theorem 3.1 we know that

$$
\delta u_{\delta} \rightarrow \int H(x, y, p, X, 0) d \mu(y)=-\lambda
$$

To prove the uniqueness of $\lambda$ and the uniqueness up to a constant of $w$, we use the arguments of [10, Thm.4.5]. For the sake of completeness, let us recall them briefly.
Assume by contradiction that there exist two constants $\lambda_{1} \neq \lambda_{2}$ and two Lipschitz functions $w_{1}, w_{2}$ such that $\left(\lambda_{1}, w_{1}\right)$ and $\left(\lambda_{2}, w_{2}\right)$ are both solutions to problem (1.6). Without any loss of generality we assume $\lambda_{1}<\lambda_{2}$. By standard arguments, $w_{1}$ and $w_{2}$ are regular solutions. We set $u(\cdot):=w_{1}(\cdot)-$ $w_{2}(\cdot)$ and $U_{1}(y):=y_{1}^{4}+y_{2}^{4}+y_{3}^{2}$. Without any loss of generality (eventually adding a constant), we assume $\sup _{\mathbb{R}^{3}} u>0$. We observe that, for $\gamma>0$ sufficiently small and $\beta>0$ sufficiently large, there hold

$$
\begin{align*}
-\operatorname{tr}\left(\sigma \sigma^{T} D^{2} u\right)-b D u & =\lambda_{1}-\lambda_{2}  \tag{3.26}\\
-\operatorname{tr}\left(\sigma \sigma^{T} D^{2} U_{1}\right)-b D U_{1} & \geq \gamma U_{1}-\beta \tag{3.27}
\end{align*} \quad \text { in } \mathbb{R}^{3} .
$$

(For instance an explicit and tedious calculation gives: $\gamma<2 k_{3}, \gamma<4 k_{1}$, $\gamma<4 k_{2}, \beta \geq \frac{100}{4 k_{1}-\gamma}+\frac{100}{4 k_{2}-\gamma}$.)
Consider $\rho>0$ so small to have $\rho \beta<\lambda_{2}-\lambda_{1}$ and set $U(\cdot):=\rho U_{1}(\cdot)$. By the global Lipschitz continuity of $w_{1}$ and $w_{2}$, we have

$$
\begin{equation*}
\lim _{|y| \rightarrow \infty}(u(y)-U(y))=-\infty \tag{3.28}
\end{equation*}
$$

Hence, there exists an open bounded set $\Omega \subset \mathbb{R}^{3}$ such that $u \leq U$ in $\mathbb{R}^{3} \backslash \Omega$. By linearity of the operator, relations (3.26) and (3.27) entail that the function $\eta(\cdot):=u(\cdot)-U(\cdot)$ satisfies

$$
-\operatorname{tr}\left(\sigma \sigma^{T} D^{2} \eta\right)-b D \eta \leq-\rho \gamma U_{1}-\left(\lambda_{2}-\lambda_{1}-\rho \beta\right)<0 \quad \text { in } \mathbb{R}^{3}
$$

where the last inequality is due to our choice of $\rho$ and to $U_{1} \geq 0$. Applying the maximum principle to $\eta$ on the domain $\Omega$ we obtain: $\eta \leq 0$ in $\Omega$. Hence, we have: $\eta \leq 0$ in $\mathbb{R}^{3}$, namely

$$
u(y) \leq \rho U_{1}(y) \quad \forall y \in \mathbb{R}^{3}
$$

Letting $\rho \rightarrow 0^{+}$, we get $u \leq 0$ in $\mathbb{R}^{3}$ which gives the desired contradiction.
Let us now pass to prove that if $\left(\lambda, w_{1}\right)$ and $\left(\lambda, w_{2}\right)$ are both solutions to (1.6) then $w_{1}=w_{2}+C$, for some constant $C$. By (3.27), there exists $R>0$ such that

$$
-\operatorname{tr}\left(\sigma \sigma^{T} D^{2} U_{1}\right)-b D U_{1}>0 \quad \text { for }|y|>R
$$

For $u=w_{1}-w_{2}$ as before, we claim that there holds

$$
\begin{equation*}
\sup _{\mathbb{R}^{3}} u=\frac{\max }{B(0, R)} u . \tag{3.29}
\end{equation*}
$$

Actually, for any $\rho>0$, for $\eta(\cdot)=u(\cdot)-\rho U_{1}(\cdot)$,

$$
-\operatorname{tr}\left(\sigma \sigma^{T} D^{2} \eta\right)-b D \eta<0 \quad \text { for }|y|>R
$$

As before (see (3.28)) $\lim _{|y| \rightarrow \infty} \eta(y)=-\infty$ and this implies that $\eta$ attains his maximum on $\mathbb{R}^{3}$.
By the maximum principle, $\eta$ cannot attain its maximum over $\overline{B(0, R)}$ C at any point in its interior. Then

$$
u(y)-\rho U_{1}(y) \leq \max _{\left|y^{\prime}\right|=R}\left(u\left(y^{\prime}\right)-\rho U_{1}\left(y^{\prime}\right)\right) \quad \forall|y| \geq R
$$

letting $\rho \rightarrow 0^{+}$, we obtain our claim (3.29).
By (3.29), for any $r>R$, the strong maximum principle on $u$ over $B(0, r)$ ensures that $u$ is a constant function on $B(0, r)$. By the arbitrariness of $r$, we obtain the desired result.

### 3.5 Regularity of the corrector $w$

In this subsection we prove the $C^{2}$-regularity of the corrector $w$. This result seems not straightforward. Actually, since our operator $\mathcal{L}$ contains second order horizontal derivatives and Euclidean first derivatives as well, the second order part of the operator does not immediately regularize the first order one. On the other hand is worth to observe that, for $H(x, \cdot, p, X, 0) \in$ $C^{\infty}$, the solution $w$ of (1.6) is $C^{\infty}$ by hypoellipticity.

We start with a lemma which states the equivalence between solution in the sense of distributions and continuous viscosity solutions, under a growth condition at infinity.

Lemma 3.5 Consider the equation

$$
\begin{equation*}
-\operatorname{tr}\left(\sigma \sigma^{T} D^{2} \chi\right)-b(y) D \chi+K \chi=R(y), y \in \mathbb{R}^{3} \tag{3.30}
\end{equation*}
$$

where $R$ is a bounded globally Lipschitz continuous function and $K$ is a strictly positive constant. Then

1) there exists a unique bounded and globally Lipschitz continuous viscosity solution $\chi$;
2) $\chi$ is a solution in the sense of distributions;
3) any bounded solution of (3.30) in the sense of distributions coincides with $\chi$.

Proof. 1) Follows from Lemma 3.2 and Theorem 3.2.
2) Follows from [14, Theorem 1].
3) If $\chi_{1}$ is a solution of (3.30) in the sense of distributions, we define $\bar{\chi}:=$ $\chi_{1}-\chi,(\bar{\chi}$ is bounded). By linearity, $\bar{\chi}$ solves in the sense of distributions

$$
-\operatorname{tr}\left(\sigma \sigma^{T} D^{2} \bar{\chi}\right)-b D \bar{\chi}+K \bar{\chi}=0 .
$$

By the hypoellipticity of the operator, $\bar{\chi}$ is smooth. Hence $\chi_{1}=\bar{\chi}+\chi$ is continuous and, by [14, Theorem 2] is also a viscosity solution. By the uniqueness of sublinear viscosity solutions of (3.30), we get $\chi_{1}=\chi$.

Theorem 3.4 Under assumptions (A1)-(A6), let $w$ be the solution of the cell problem (1.6) found in Theorem 3.3. Then $w \in C_{\text {loc }}^{2, \alpha}\left(\mathbb{R}^{3}\right)$, for some $\alpha \in(0,1)$.

Proof. Let us denote by $X_{1}=\left(\partial_{y_{1}}+2 y_{2} \partial_{y_{3}}\right)$ and $X_{2}=\left(\partial_{y_{2}}-2 y_{1} \partial_{y_{3}}\right)$ the two vector fields associated to the two columns of the matrix $\sigma$. Recall that
these two vectors are the generators of the Heisenberg group and span all $\mathbb{R}^{3}$ because their commutator is $\left[X_{1}, X_{2}\right]=-4 \frac{\partial}{\partial y_{3}}$.
Along the proof $\alpha$ is a strictly positive constant which may change line to line.
The corrector $w$ solves

$$
\begin{equation*}
-\operatorname{tr}\left(\sigma(y) \sigma^{T}(y) D^{2} w(y)\right)-b(y) D w(y)=G(y) \tag{3.31}
\end{equation*}
$$

with

$$
G(y):=\lambda-H\left(\bar{x}, y, D_{x} \psi(\bar{t}, \bar{x}), D_{x x}^{2} \psi(\bar{t}, \bar{x}), 0\right)
$$

First let us get the Lipschitz continuity of $\frac{\partial w}{\partial y_{3}}$. Derivating equation (3.31) with respect to $y_{3}$ we obtain that the function $u:=\frac{\partial w}{\partial y_{3}}$ solves in the sense of distributions

$$
\begin{equation*}
-\operatorname{tr}\left(\sigma \sigma^{T} D^{2} u\right)-b D u+k_{3} u=\frac{\partial G}{\partial y_{3}} \tag{3.32}
\end{equation*}
$$

with $k_{3}>0$ by assumption A2). Note that $u$ is bounded by Theorem 3.2, then by Lemma 3.5 we get that $u$ is Lipschitz continuous and it is also a viscosity solution; hence

$$
\begin{equation*}
u=\frac{\partial w}{\partial y_{3}} \in B \operatorname{Lip}\left(\mathbb{R}^{3}\right) \tag{3.33}
\end{equation*}
$$

Deriving equation (3.32) with respect to $y_{3}$, we get that the function $z:=$ $\partial^{2} w / \partial y_{3}^{2}$ solves

$$
\begin{equation*}
-\operatorname{tr}\left(\sigma \sigma^{T} D^{2} z\right)-b D z+2 k_{3} z=\frac{\partial^{2} G}{\partial y_{3}^{2}} \tag{3.34}
\end{equation*}
$$

By assumption A6), we can apply Lemma 3.5 also to (3.34) and we get that the function $z$ is globally Lipschitz continuous, i.e.

$$
\begin{equation*}
\frac{\partial^{2} w}{\partial y_{3}^{2}} \in B \operatorname{Lip}\left(\mathbb{R}^{3}\right) \tag{3.35}
\end{equation*}
$$

Now we study the regularity of $w$ with respect to $y_{1}$ and $y_{2}$; to this end, let us come back to (3.31). From the Lipschitz continuity of $w$ (see Theorem 3.2), we get

$$
-\operatorname{tr}\left(\sigma \sigma^{T} D^{2} w\right) \in L_{l o c}^{\infty}\left(\mathbb{R}^{3}\right)
$$

By (3.33), we have

$$
\frac{\partial^{2} w}{\partial y_{1} \partial y_{3}}, \frac{\partial^{2} w}{\partial y_{2} \partial y_{3}}, \frac{\partial^{2} w}{\partial y_{3}^{2}} \in L^{\infty}\left(\mathbb{R}^{3}\right)
$$

Taking into account the explicit expression of $-\operatorname{tr}\left(\sigma \sigma^{T} D^{2} w\right)$ we have that

$$
\frac{\partial^{2} w}{\partial y_{1}^{2}}+\frac{\partial^{2} w}{\partial y_{2}^{2}} \in L_{l o c}^{\infty}\left(\mathbb{R}^{3}\right)
$$

This relation and (3.35) imply $\Delta w \in L_{l o c}^{\infty}\left(\mathbb{R}^{3}\right),(\Delta$ is the Euclidean Laplacian). Hence from classical results on uniformly elliptic equations we obtain

$$
D w \in C_{l o c}^{0, \alpha}\left(\mathbb{R}^{3}\right)
$$

Now we can replace $w$ with $u$, (3.33) with (3.35) and (3.31) with (3.32), using the same arguments we get:

$$
\begin{equation*}
D u \in C_{l o c}^{0, \alpha}\left(\mathbb{R}^{3}\right) \tag{3.36}
\end{equation*}
$$

In particular

$$
\frac{\partial^{2} w}{\partial y_{1} \partial y_{3}}, \frac{\partial^{2} w}{\partial y_{2} \partial y_{3}}, \frac{\partial^{2} w}{\partial y_{3}^{2}} \in C_{l o c}^{0, \alpha}\left(\mathbb{R}^{3}\right)
$$

As before we have that $\Delta w \in C_{l o c}^{0, \alpha}\left(\mathbb{R}^{3}\right)$. Hence, from classical results on uniformly elliptic equations we obtain the statement.

Remark 3.3 We remark that in this proof the structure of the operator $\mathcal{L}$ and Theorem 3.2 play a crucial role, this allows us to overcome the application of some deep results on the hypoelliptic theory.

## 4 Proof of Theorem 2.1.

In this section we prove the convergence of the solution $V^{\epsilon}$ of (1.1) to the solution of the effective equation (1.3).

Let us recall from Proposition 2.1 that, for every compact $K \subset \mathbb{R}^{n}$, the solutions $V^{\epsilon}$ are equibounded in $(0, T) \times K \times \mathbb{R}^{3}$, hence the following semilimits

$$
\begin{aligned}
\bar{V}(t, x, y) & =\limsup _{\epsilon \rightarrow 0^{+}, t^{\prime} \rightarrow t, x^{\prime} \rightarrow x, y^{\prime} \rightarrow y} V^{\epsilon}\left(t^{\prime}, x^{\prime}, y^{\prime}\right) \quad \text { for } t<T \\
\bar{V}(T, x, y) & =\limsup _{t^{\prime} \rightarrow T^{-}, x^{\prime} \rightarrow x, y^{\prime} \rightarrow y} \bar{V}\left(t^{\prime}, x^{\prime}, y^{\prime}\right) \quad \text { for } t=T
\end{aligned}
$$

(and similarly for $\underline{\underline{V}}$ replacing limsup by liminf) are well defined. This two step definition of $\bar{V}$ is needed to overcome an expected initial layer.

For the sake of clarity we shall divide the proof in several steps, as follows:

Step 1. $\bar{V}$ and $\underline{V}$ are independent of $y$;
Step 2. $\bar{V}$ and $\underline{V}$ are respectively a subsolution and a supersolution of the parabolic equation (1.3);

Step 3. $\bar{V}(T, x) \leq \bar{g}(x) \leq \underline{V}(T, x)$, where $\bar{g}(x)$ is defined in (1.4);
Step 4. $\bar{V}=\underline{V}=: V$ and $V^{\epsilon} \rightarrow V$ locally uniformly.

### 4.1 Step 1

Lemma 4.1 Under assumptions A1)-A5), $\bar{V}$ and $\underline{V}$ are independent of $y$.
Proof. Let us observe that $\bar{V}(t, x, y)$ and $\underline{V}(t, x, y)$ are respectively BUSC and BLSC. We prove that $\bar{V}(t, x, y)$ is independent of $y$; being similar, the proof for $\underline{V}(t, x, y)$ is omitted.

We claim that for $\left(t_{0}, x_{0}\right) \in(0, T) \times \mathbb{R}^{n}$ fixed, $\bar{V}\left(t_{0}, x_{0}, y\right)$ is a subsolution for $y \in \mathbb{R}^{3}$ to equation

$$
\begin{equation*}
-\operatorname{tr}\left(\sigma \sigma^{T}(y) D_{y y} V\right)-b(y) \cdot D_{y} V=0 \tag{4.1}
\end{equation*}
$$

Assuming for the moment that is true, since $\bar{V}\left(t_{0}, x_{0}, y\right)$ is BUSC in $y$, we can apply the Liouville theorem proved in [16, Proposition 3.1] to deduce that the function $\bar{V}\left(t_{0}, x_{0}, y\right)$ does not depend on $y$.

In order to prove that $\bar{V}\left(t_{0}, x_{0}, y\right)$ is a subsolution to equation (4.1) we follow the same arguments as in Step 2 of [5, Theorem 3.2], which we write for the sake of completeness.

First of all we prove that $\bar{V}(t, x, y)$ is a subsolution to equation (4.1) for $(t, x, y) \in(0, T) \times \mathbb{R}^{3} \times \mathbb{R}^{3}$. To do this we fix a point $(\bar{t}, \bar{x}, \bar{y})$ and a smooth function $\psi$ such that $\bar{V}-\psi$ has a local strict maximum at $(\bar{t}, \bar{x}, \bar{y})=\bar{P}$ in $\overline{B(\bar{P}, r)}=\{(t, x, y):|(t, x, y)-(\bar{t}, \bar{x}, \bar{y})| \leq r\}$, for some $r>0$. Using the definition of the half relaxed semilimit it is possible to prove (see [4]) that there exists $\epsilon_{n} \rightarrow 0$ and $\left(t_{n}, x_{n}, y_{n}\right) \in \overline{B(\bar{P}, r)}$ such that $\left(t_{n}, x_{n}, y_{n}\right) \rightarrow$ $(\bar{t}, \bar{x}, \bar{y}),\left(t_{n}, x_{n}, y_{n}\right)$ are maxima for $V^{\epsilon_{n}}-\psi$ in $\overline{B(\bar{P}, r)}$ and $V^{\epsilon_{n}}\left(t_{n}, x_{n}, y_{n}\right) \rightarrow$
$\bar{V}(\bar{t}, \bar{x}, \bar{y})$. Since $V^{\epsilon_{n}}$ solves (1.1) then

$$
\begin{aligned}
& -\partial_{t} \psi+H\left(x_{n}, y_{n}, D_{x} \psi, D_{x x}^{2} \psi, \frac{D_{x y}^{2} \psi}{\sqrt{\epsilon_{n}}}\right) \\
& -\frac{1}{\epsilon_{n}} \operatorname{tr}\left(\sigma\left(y_{n}\right) \sigma\left(y_{n}\right)^{T} D_{y y} \psi\right)-\frac{1}{\epsilon}_{n} b\left(y_{n}\right) \cdot D_{y} \psi+a V^{\epsilon_{n}}\left(t_{n}, x_{n}, y_{n}\right) \leq 0
\end{aligned}
$$

Then

$$
\begin{aligned}
& -\operatorname{tr}\left(\sigma\left(y_{n}\right) \sigma\left(y_{n}\right)^{T} D_{y y} \psi\right)-b\left(y_{n}\right) \cdot D_{y} \psi \leq \\
& \epsilon_{n}\left[\partial_{t} \psi-H\left(x_{n}, y_{n}, D_{x} \psi, D_{x x}^{2} \psi, \frac{D_{x y}^{2} \psi}{\sqrt{\epsilon}}\right)-a V^{\epsilon_{n}}\right]
\end{aligned}
$$

From the regularity of $\psi$, the continuity of $H$ (obtained from A1), A3), A5)) and the local uniform boundedness of $V^{\epsilon_{n}}$, the part in the brackets on the right hand side is uniformly bounded with respect to $n$ in $B(\bar{P}, r)$, then passing to the limit as $\epsilon_{n} \rightarrow 0$ we get

$$
-\operatorname{tr}\left(\sigma(\bar{y}) \sigma(\bar{y})^{T} D_{y y} \psi\right)-b(\bar{y}) \cdot D_{y} \psi \leq 0
$$

i.e. $\bar{V}(t, x, y)$ is a subsolution to equation (4.1) for $(t, x, y) \in(0, T) \times \mathbb{R}^{n} \times \mathbb{R}^{3}$.

We now show that, arguing as in [4, Lemma II.5.17], for every fixed $\left(t_{0}, x_{0}\right) \in(0, T) \times \mathbb{R}^{n}$ the function $\bar{V}\left(t_{0}, x_{0}, y\right)$ is a subsolution to equation (4.1). We fix $\bar{y}$ and $\phi(y)$, a smooth function such that $\bar{V}\left(t_{0}, x_{0}, y\right)-\phi(y)$ has a strict local maximum at $\bar{y}$ in $B(\bar{y}, \delta)$ and such that $\phi(y) \geq 1$ in $B(\bar{y}, \delta)$. Let us chose $\delta>0$ small enough s.t. $t_{0}-\delta>0$. We define, for $\eta>0$

$$
\phi_{\eta}(t, x, y)=\phi(y)\left(1+\frac{\left|x-x_{0}\right|^{2}+\left|t-t_{0}\right|^{2}}{\eta}\right)
$$

and we consider $\left(t_{\eta}, x_{\eta}, y_{\eta}\right)$ a maximum point of $\bar{V}-\phi_{\eta}$ in $\overline{B\left(\left(t_{0}, x_{0}, \bar{y}\right), \delta\right)}$. We remark that

$$
\begin{gather*}
\bar{V}\left(t_{\eta}, x_{\eta}, y_{\eta}\right)-\phi_{\eta}\left(t_{\eta}, x_{\eta}, y_{\eta}\right) \geq \bar{V}\left(t_{0}, x_{0}, \bar{y}\right)-\phi(\bar{y})  \tag{4.2}\\
\bar{V}\left(t_{\eta}, x_{\eta}, y_{\eta}\right)-\phi\left(y_{\eta}\right) \geq \bar{V}\left(t_{0}, x_{0}, \bar{y}\right)-\phi(\bar{y})
\end{gather*}
$$

and we can prove that, eventually passing to subsequences (as $\eta \rightarrow 0$ ) first that $\left(t_{\eta}, x_{\eta}\right) \rightarrow\left(t_{0}, x_{0}\right)$, then that $y_{\eta} \rightarrow \bar{y}$ using the strict maximum
property of $\bar{y}$.
Using (4.2) and the upper semicontinuity of $\bar{V}$

$$
K_{\eta}:=\left(1+\frac{\left|x_{\eta}-x_{0}\right|^{2}+\left|t_{\eta}-t_{0}\right|^{2}}{\eta}\right) \rightarrow K>0
$$

Now, using the fact that $\bar{V}(t, x, y)$ is a subsolution of (4.1) in $\left(t_{\eta}, x_{\eta}, y_{\eta}\right)$ we get

$$
-\operatorname{tr}\left(\sigma\left(y_{\eta}\right) \sigma\left(y_{\eta}\right)^{T} D_{y y}^{2} \phi_{\eta}\right)-b\left(y_{\eta}\right) \cdot D_{y} \phi_{\eta} \leq 0
$$

which gives, passing to the limit as $\eta \rightarrow 0$

$$
-\operatorname{tr}\left(\sigma(\bar{y}) \sigma(\bar{y})^{T} D_{y y}^{2} \phi\right)-b(\bar{y}) \cdot D_{y} \phi \leq 0
$$

Remark 4.1 Using (4.2) and the upper semicontinuity of $\bar{V}$ it is possible to prove that $K_{\eta} \rightarrow 1$. This property in not used in our proof but can be useful in more general and nonlinear cases.

### 4.2 Step 2

Proposition 4.1 Under the assumptions A1)-A5), $\bar{V}$ and $\underline{V}$ are respectively a subsolution and a supersolution of the parabolic equation in (1.3).

Proof. We prove that $\bar{V}$ is a viscosity subsolution of (1.3) in $] 0, T\left[\times \mathbb{R}^{n}\right.$. The proof that $\underline{V}$ is a viscosity supersolution is analogous, so we shall omit it.

We take a smooth test function $\psi(t, x)$ such that $(\bar{t}, \bar{x}) \in] 0, T\left[\times \mathbb{R}^{n}\right.$ is a strict local maximum point for $\bar{V}-\psi$. We have to prove that

$$
-\partial_{t} \psi(\bar{t}, \bar{x})+\bar{H}\left(\bar{x}, D_{x} \psi(\bar{t}, \bar{x}), D_{x x}^{2} \psi(\bar{t}, \bar{x})\right)+a \psi(\bar{t}, \bar{x}) \leq 0
$$

Without any loss of generality we can assume that:

1. $\bar{V}(\bar{t}, \bar{x})=\psi(\bar{t}, \bar{x})$;
2. $\psi$ is coercive in $x$ uniformly in t , i.e.

$$
\begin{equation*}
\lim _{|x| \rightarrow \infty} \inf _{t \in[0, T]} \psi(t, x)=+\infty \tag{4.3}
\end{equation*}
$$

3. there holds

$$
\begin{equation*}
\inf _{x \in \mathbb{R}^{n}} \psi\left(\frac{\bar{t}}{2}, x\right)>M+1, \quad \inf _{x \in \mathbb{R}^{n}} \psi\left(\frac{\bar{t}+T}{2}, x\right)>M+1 \tag{4.4}
\end{equation*}
$$

where $M$ is a constant such that $\mid V^{\epsilon}(t, x, y) \leq M$;
4. $\sup _{(t, x) \in K}\left|\partial_{t} \psi(t, x)\right| \leq C_{K}$ for any $K$ compact in $[0, T] \times \mathbb{R}^{n}$.

For any fixed $\eta \in] 0,1]$, let us consider now the "perturbed test function":

$$
\psi_{\epsilon \eta}(t, x, y):=\psi(t, x)+\epsilon(w(y)+\eta \chi(y))
$$

where $w(y)$ is the viscosity solution of the cell problem (1.6) founded in Theorem 3.3 associated to $\left(\bar{x}, D_{x} \psi(\bar{t}, \bar{x}), D_{x x}^{2} \psi(\bar{t}, \bar{x})\right)$ and $\chi(y)$ is the Lyapunov function

$$
\begin{equation*}
\chi(y)=y_{1}^{2}+y_{2}^{2}+y_{3}^{2} \tag{4.5}
\end{equation*}
$$

Note that, from (3.25) and the definition of $\chi$ in (4.5), we have

$$
\begin{equation*}
w(y)+\eta \chi(y) \rightarrow+\infty, \text { if }|y| \rightarrow+\infty \tag{4.6}
\end{equation*}
$$

and there exists a constant $k_{0}$ independent of $\eta$ such that

$$
\begin{equation*}
w(y)+\eta \chi(y) \geq-k_{0}(1+\log (\eta)) \tag{4.7}
\end{equation*}
$$

Let consider the function

$$
\Psi(t, x, y):=V^{\epsilon}(t, x, y)-\psi_{\epsilon \eta}(t, x, y)
$$

Thanks to the equi-boundedness of $V^{\epsilon},(4.6)$ and (4.3) we have:

$$
\Psi(t, x, y) \rightarrow-\infty, \text { if }(x, y) \rightarrow+\infty
$$

and there exists a point $\left(t_{\epsilon, \eta}, x_{\epsilon, \eta}, y_{\epsilon, \eta}\right) \in\left[\frac{\bar{t}}{2}, \frac{\bar{t}+T}{2}\right] \times \mathbb{R}^{n} \times \mathbb{R}^{3}$ which is a global maximum point of $\Psi$ in $\left[\frac{\bar{t}}{2}, \frac{\bar{t}+T}{2}\right] \times \mathbb{R}^{n} \times \mathbb{R}^{3}$.

Claim $4.1\left(t_{\epsilon, \eta}, x_{\epsilon, \eta}\right)$ is bounded uniformly in $\epsilon$.

The points $t_{\epsilon, \eta}$ are obviously bounded. Now using the maximum property of ( $\left.t_{\epsilon, \eta}, x_{\epsilon, \eta}, y_{\epsilon, \eta}\right)$, we have:

$$
V^{\epsilon}\left(t_{\epsilon, \eta}, x_{\epsilon, \eta}, y_{\epsilon, \eta}\right)-\psi_{\epsilon \eta}\left(t_{\epsilon, \eta}, x_{\epsilon, \eta}, y_{\epsilon, \eta}\right) \geq V^{\epsilon}(\bar{t}, 0,0)-\psi(\bar{t}, 0)-\epsilon(w(0)+\eta \chi(0) ;
$$

then from (4.7)

$$
\begin{aligned}
& K \geq V^{\epsilon}\left(t_{\epsilon, \eta}, x_{\epsilon, \eta}, y_{\epsilon, \eta}\right)-V^{\epsilon}(\bar{t}, 0,0)+\psi(\bar{t}, 0)+\epsilon(w(0)+\eta \chi(0)) \geq \\
& \quad \psi\left(t_{\epsilon, \eta}, x_{\epsilon, \eta}\right)+\epsilon\left(w\left(y_{\epsilon, \eta}\right)+\eta \chi\left(y_{\epsilon, \eta}\right)\right) \geq \psi\left(t_{\epsilon, \eta}, x_{\epsilon, \eta}\right)-\epsilon k_{0}(1+\log (\eta))
\end{aligned}
$$

and we end the proof of Claim 4.1 using and (4.3).
Claim 4.2 If $t_{\epsilon, \eta}=\frac{\bar{t}}{2}$ or $t_{\epsilon, \eta}=\frac{\bar{t}+T}{2}$, then for any $\left(t^{\prime}, x^{\prime}, y^{\prime}\right) \in\left[\frac{\bar{t}}{2}, \frac{\bar{t}+T}{2}\right] \times$ $\mathbb{R}^{n} \times \mathbb{R}^{3}$

$$
\Psi\left(t^{\prime}, x^{\prime}, y^{\prime}\right) \leq-1+\epsilon k_{0}(1+\log (\eta))
$$

Thanks to (4.7)

$$
\Psi\left(t^{\prime}, x^{\prime}, y^{\prime}\right) \leq V^{\epsilon}\left(t_{\epsilon, \eta}, x_{\epsilon, \eta}, y_{\epsilon, \eta}\right)-\psi\left(t_{\epsilon, \eta}, x_{\epsilon, \eta}\right)+\epsilon k_{0}(1+\log (\eta)) .
$$

Using now (4.4) and the definition of $M$

$$
\begin{aligned}
\Psi\left(t^{\prime}, x^{\prime}, y^{\prime}\right) & \leq V^{\epsilon}\left(t_{\epsilon, \eta}, x_{\epsilon, \eta}, y_{\epsilon, \eta}\right)-(M+1)+\epsilon k_{0}(1+\log (\eta)) \\
& \leq-1+\epsilon k_{0}(1+\log (\eta)) .
\end{aligned}
$$

Claim 4.3 If $\left.t_{\epsilon, \eta} \in\right] \frac{\bar{t}}{2}, \frac{\bar{t}+T}{2}[$, then

$$
\begin{align*}
& -\partial_{t} \psi\left(t_{\epsilon, \eta}, x_{\epsilon, \eta}\right)+\bar{H}\left(\bar{x}, D_{x} \psi(\bar{t}, \bar{x}), D_{x x}^{2} \psi(\bar{t}, \bar{x})\right)  \tag{4.8}\\
& -\eta \mathcal{L}\left(y_{\epsilon, \eta}, D_{y} \chi\left(y_{\epsilon, \eta}\right), D_{y y}^{2} \chi\left(y_{\epsilon, \eta}\right)\right)+\mathcal{F}_{\epsilon, \eta}+a V^{\epsilon}\left(t_{\epsilon, \eta}, x_{\epsilon \eta}, y_{\epsilon, \eta}\right) \leq 0,
\end{align*}
$$

where

$$
\begin{aligned}
\mathcal{F}_{\epsilon, \eta}=H\left(x_{\epsilon, \eta}, y_{\epsilon, \eta}, D_{x} \psi\left(t_{\epsilon, \eta}, x_{\epsilon, \eta}\right)\right. & \left., D_{x x}^{2} \psi\left(t_{\epsilon, \eta}, x_{\epsilon, \eta}\right), 0\right) \\
& -H\left(\bar{x}, y_{\epsilon, \eta}, D_{x} \psi(\bar{t}, \bar{x}), D_{x x}^{2} \psi(\bar{t}, \bar{x}), 0\right) .
\end{aligned}
$$

By definition of viscosity subsolution of (1.1) and using the regularity of $w$ (proved in Theorem 3.4), $\psi$ and $\chi$ :

$$
\begin{aligned}
-\partial_{t} \psi\left(t_{\epsilon, \eta}, x_{\epsilon, \eta}\right)+H\left(x_{\epsilon, \eta}, y_{\epsilon, \eta}, D_{x} \psi\left(t_{\epsilon, \eta}, x_{\epsilon, \eta}\right)\right. & \left., D_{x x}^{2} \psi\left(t_{\epsilon, \eta}, x_{\epsilon, \eta}\right), 0\right) \\
-\mathcal{L}\left(y_{\epsilon, \eta}, D_{y} w\left(y_{\epsilon, \eta}\right), D_{y y}^{2} w\left(y_{\epsilon, \eta}\right)\right)-\eta \mathcal{L}\left(y_{\epsilon, \eta},\right. & \left.D_{y} \chi\left(y_{\epsilon, \eta}\right), D_{y y}^{2} \chi\left(y_{\epsilon, \eta}\right)\right) \\
& +a V^{\epsilon}\left(t_{\epsilon, \eta}, x_{\epsilon \eta}, y_{\epsilon, \eta}\right) \leq 0
\end{aligned}
$$

Now we use (1.6) for $y=y_{\epsilon, \eta}$

$$
\begin{aligned}
& -\partial_{t} \psi\left(t_{\epsilon, \eta}, x_{\epsilon, \eta}\right)+H\left(x_{\epsilon, \eta}, y_{\epsilon, \eta}, D_{x} \psi\left(t_{\epsilon, \eta}, x_{\epsilon, \eta}\right), D_{x x}^{2} \psi\left(t_{\epsilon, \eta}, x_{\epsilon, \eta}\right), 0\right) \\
& \quad+\bar{H}\left(\bar{x}, D_{x} \psi(\bar{t}, \bar{x}), D_{x x}^{2} \psi(\bar{t}, \bar{x})\right)-H\left(\bar{x}, y_{\epsilon, \eta}, D_{x} \psi(\bar{t}, \bar{x}), D_{x x}^{2} \psi(\bar{t}, \bar{x}), 0\right) \\
& \quad-\eta \mathcal{L}\left(y_{\epsilon, \eta}, D_{y} \chi\left(y_{\epsilon, \eta}\right), D_{y y}^{2} \chi\left(y_{\epsilon, \eta}\right)\right)+a V^{\epsilon}\left(t_{\epsilon, \eta}, x_{\epsilon \eta}, y_{\epsilon, \eta}\right) \leq 0
\end{aligned}
$$

which is equivalent to (4.8).
Claim 4.4 If $\left.t_{\epsilon, \eta} \in\right] \frac{\bar{t}}{2}, \frac{\bar{t}+T}{2}\left[\right.$, then $\left\{y_{\epsilon, \eta}\right\}_{\epsilon}$ is uniformly bounded.
By contradiction: let us assume that there exists a sequence $\left\{\epsilon_{n}\right\}_{n}$ with $\epsilon_{n} \rightarrow 0$ such that

$$
\left.t_{\epsilon_{n}, \eta} \in\right] \frac{\bar{t}}{2}, \frac{\bar{t}+T}{2}\left[, \quad\left|y_{\epsilon_{n}, \eta}\right| \rightarrow \infty \text { for } n \rightarrow \infty\right.
$$

Then, there holds, calling $y_{\epsilon_{n}, \eta}=\left(y_{\epsilon_{n}, \eta, 1}, y_{\epsilon_{n}, \eta, 2}, y_{\epsilon_{n}, \eta, 3}\right)$ :

$$
\begin{align*}
& -\mathcal{L}\left(y_{\epsilon_{n}, \eta}, D_{y} \chi\left(y_{\epsilon_{n}, \eta}\right), D_{y y} \chi\left(y_{\epsilon_{n}, \eta}\right)\right)=  \tag{4.9}\\
& \quad 2\left(-2+\left(k_{1}-4\right) y_{\epsilon_{n}, \eta, 1}^{2}+\left(k_{2}-4\right) y_{\epsilon_{n}, \eta, 2}^{2}+k_{3} y_{\epsilon_{n}, \eta, 3}^{2}\right) \rightarrow \infty .
\end{align*}
$$

Moreover, by Claim 4.1 and 2.1 we get

$$
\left|\mathcal{F}_{\epsilon_{n}, \eta}\right| \leq K(\eta)=K .
$$

Coming back to (4.8), using again Claim 4.1 and the uniform boundedness of $V^{\epsilon}$ we have a contradiction.

Claim 4.5 There holds

$$
\left(t_{\epsilon, \eta}, x_{\epsilon, \eta}\right) \rightarrow(\bar{t}, \bar{x}) \quad \text { for } \epsilon \rightarrow 0 .
$$

There exists $(\tilde{t}, \tilde{x}) \in\left[\frac{\bar{t}}{2}, \frac{\bar{t}+T}{2}\right] \times \mathbb{R}^{n}$ such that (possibly passing to a subsequence)

$$
\left(t_{\epsilon, \eta}, x_{\epsilon, \eta}\right) \rightarrow(\tilde{t}, \tilde{x}) \quad \text { as } \epsilon \rightarrow 0 .
$$

By definition of $\left(t_{\epsilon, \eta}, x_{\epsilon, \eta}, y_{\epsilon, \eta}\right)$ we have: $\forall\left(t^{\prime}, x^{\prime}, y^{\prime}\right) \in\left[\frac{\bar{t}}{2}, \frac{\bar{t}+T}{2}\right] \times \mathbb{R}^{n} \times \mathbb{R}^{3}$

$$
\begin{align*}
& V^{\epsilon}\left(t_{\epsilon, \eta}, x_{\epsilon, \eta}, y_{\epsilon, \eta}\right)-\psi\left(t_{\epsilon, \eta}, x_{\epsilon, \eta}\right)-\epsilon\left(w\left(y_{\epsilon, \eta}\right)+\eta \chi\left(y_{\epsilon, \eta}\right)\right) \geq  \tag{4.10}\\
& V^{\epsilon}\left(t^{\prime}, x^{\prime}, y^{\prime}\right)-\psi\left(t^{\prime}, x^{\prime}\right)-\epsilon\left(w\left(y^{\prime}\right)+\eta \chi\left(y^{\prime}\right)\right)
\end{align*}
$$

Let $(t, x, y) \in\left[\frac{\bar{t}}{2}, \frac{\bar{t}+T}{2}\right] \times \mathbb{R}^{n} \times \mathbb{R}^{3}$. Passing to $\limsup _{\left(\epsilon, t^{\prime}, x^{\prime}, y^{\prime}\right) \rightarrow(0, t, x, y)}$ in the previous inequality (4.10):

$$
\begin{align*}
\limsup _{\left(\epsilon, t^{\prime}, x^{\prime}, y^{\prime}\right) \rightarrow(0, t, x, y)}\left[V^{\epsilon}\left(t_{\epsilon, \eta}, x_{\epsilon, \eta}, y_{\epsilon, \eta}\right)-\psi\left(t_{\epsilon, \eta}, x_{\epsilon, \eta}\right)-\epsilon\left(w\left(y_{\epsilon, \eta}\right)+\eta \chi\left(y_{\epsilon, \eta}\right)\right)\right] & \geq  \tag{4.11}\\
\bar{V}(t, x)-\psi(t, x) &
\end{align*}
$$

Moreover there is a sequence $\epsilon_{n} \rightarrow 0$ such that

$$
\begin{align*}
& \lim _{\left(\epsilon, t^{\prime}, x^{\prime}, y^{\prime}\right) \rightarrow(0, t, x, y)}\left[V^{\epsilon}\left(t_{\epsilon, \eta}, x_{\epsilon, \eta}, y_{\epsilon, \eta}\right)-\psi\left(t_{\epsilon, \eta}, x_{\epsilon, \eta}\right)-\epsilon\left(w\left(y_{\epsilon, \eta}\right)+\eta \chi\left(y_{\epsilon, \eta}\right)\right)\right]=  \tag{4.12}\\
& \lim _{n \rightarrow \infty}\left[V^{\epsilon_{n}}\left(t_{\epsilon_{n}, \eta}, x_{\epsilon_{n}, \eta}, y_{\epsilon_{n}, \eta}\right)-\psi\left(t_{\epsilon_{n}, \eta}, x_{\epsilon, \eta}\right)-\epsilon_{n}\left(w\left(y_{\epsilon, \eta}\right)+\eta \chi\left(y_{\epsilon_{n}, \eta}\right)\right)\right]
\end{align*}
$$

We split the argument according to the case that $t_{\epsilon_{n}, \eta}$ belongs to the interior or to the boundary of $\left[\frac{\bar{t}}{2}, \frac{\bar{t}+T}{2}\right]$.

- If $t_{\epsilon_{n}, \eta}=\frac{\bar{t}}{2}$ or $t_{\epsilon_{n}, \eta}=\frac{\bar{t}+T}{2}$ we apply claim 4.2

$$
\begin{array}{r}
V^{\epsilon_{n}}\left(t_{\epsilon_{n}, \eta}, x_{\epsilon_{n}, \eta}, y_{\epsilon_{n}, \eta}\right)-\psi\left(t_{\epsilon_{n}, \eta}, x_{\epsilon, \eta}\right)-\epsilon_{n}\left(w\left(y_{\epsilon, \eta}\right)+\eta \chi\left(y_{\epsilon_{n}, \eta}\right)=\right. \\
\Psi\left(t_{\epsilon_{n}, \eta}, x_{\epsilon_{n}, \eta}, y_{\epsilon_{n}, \eta}\right) \leq-1+\epsilon_{n} k_{0}(1+\log (\eta))
\end{array}
$$

If $t_{\epsilon_{n}, \eta}=\frac{\bar{t}}{2}$ or $t_{\epsilon_{n}, \eta}=\frac{\bar{t}+T}{2}$ for an infinite sequence of indices $\epsilon_{n}$ we have:

$$
-1 \geq \bar{V}(t, x)-\psi(t, x)
$$

and this is a contradiction since, for $(t, x)=(\bar{t}, \bar{x}), \bar{V}(\bar{t}, \bar{x})-\psi(\bar{t}, \bar{x})=$ 0.

- If $\left.t_{\epsilon, \eta} \in\right] \frac{\bar{t}}{2}, \frac{\bar{t}+T}{2}\left[\right.$, then by Claim $4.4\left\{y_{\epsilon, \eta}\right\}_{\epsilon}$ is uniformly bounded and we can assume (possibly passing to a subsequence) that

$$
y_{\epsilon_{n}, \eta} \rightarrow \tilde{y}_{\eta}
$$

Using (4.11), (4.12) and the upper-semicontinuity of $\bar{V}$

$$
\begin{aligned}
\bar{V}(\tilde{t}, \tilde{x})-\psi(\tilde{t}, \tilde{x}) \geq & \limsup _{\left(\epsilon^{\prime}, t^{\prime}, x^{\prime}, y^{\prime}\right) \rightarrow\left(0, \tilde{t}, \tilde{,}, \tilde{y}_{\eta}\right)}\left[V^{\epsilon^{\prime}}\left(t^{\prime}, x^{\prime}, y^{\prime}\right)\right. \\
& \left.\quad-\psi\left(t^{\prime}, x^{\prime}\right)-\epsilon^{\prime}\left(w\left(y^{\prime}\right)+\eta \chi\left(y^{\prime}\right)\right)\right] \\
\geq & \lim _{n \rightarrow \infty}\left[V^{\epsilon_{n}}\left(t_{\epsilon_{n}, \eta}, x_{\epsilon_{n}, \eta}, y_{\epsilon_{n}, \eta}\right)\right. \\
& \left.\quad-\psi\left(\epsilon_{\epsilon_{n}, \eta}, x_{\epsilon, \eta}\right)-\epsilon_{n}\left(w\left(y_{\epsilon, \eta}\right)+\eta \chi\left(y_{\epsilon_{n}, \eta}\right)\right)\right] \\
\geq & \bar{V}(t, x)-\psi(t, x) .
\end{aligned}
$$

Using the strict maximum property of $(\bar{t}, \bar{x})$, we get $(\bar{t}, \bar{x})=(\tilde{t}, \tilde{x})$.
Let us remark that the previous inequalities imply also:

$$
\begin{equation*}
\lim _{n \rightarrow \infty} V^{\epsilon_{n}}\left(t_{\epsilon_{n}, \eta}, x_{\epsilon_{n}, \eta}, y_{\epsilon_{n}, \eta}\right)=\bar{V}(\bar{t}, \bar{x}) \tag{4.13}
\end{equation*}
$$

Claim 4.6 There holds

$$
\begin{align*}
-\partial_{t} \psi(\bar{t}, \bar{x})+\bar{H}(\bar{x}, & \left.D_{x} \psi(\bar{t}, \bar{x}), D_{x x}^{2} \psi(\bar{t}, \bar{x})\right)  \tag{4.14}\\
& -\eta \mathcal{L}\left(\tilde{y}_{\eta}, D_{y} \chi\left(\tilde{y}_{\eta}\right), D^{2} \chi\left(\tilde{y}_{\eta}\right)\right)+a \bar{V}(\bar{t}, \bar{x}) \leq 0 .
\end{align*}
$$

Using Claim 4.3, we get

$$
\begin{aligned}
& \quad-\partial_{t} \psi\left(t_{\epsilon_{n}, \eta}, x_{\epsilon_{n}, \eta}\right)+\bar{H}\left(\bar{x}, D_{x} \psi(\bar{t}, \bar{x}), D_{x x}^{2} \psi(\bar{t}, \bar{x})\right) \\
& -\eta \mathcal{L}\left(\tilde{y}_{\eta}, D_{y} \chi\left(y_{\epsilon_{n}, \eta}\right), D_{y y}^{2} \chi\left(y_{\epsilon_{n}, \eta}\right)\right)+\mathcal{F}_{\epsilon_{n}, \eta}+a V^{\epsilon_{n}}\left(t_{\epsilon_{n}, \eta}, x_{\epsilon_{n}, \eta}, y_{\epsilon_{n}, \eta}\right) \leq 0 .
\end{aligned}
$$

Thanks to the regularity properties of $H$ (see (2.1)) and $\psi$, it is easy to get

$$
\mathcal{F}_{\epsilon_{n}, \eta} \rightarrow 0
$$

From (4.13), the statement follows easily.
Claim 4.7 There holds

$$
-\partial_{t} \psi(\bar{t}, \bar{x})+\bar{H}\left(\bar{x}, D_{x} \psi(\bar{t}, \bar{x}), D_{x x}^{2} \psi(\bar{t}, \bar{x})\right)+a \bar{V}(\bar{t}, \bar{x}) \leq 0 .
$$

We split the argument according to the fact that $\left\{\tilde{y}_{\eta}\right\}_{\eta}$ is uniformly bounded or not. If $\left\{\tilde{y}_{\eta}\right\}_{\eta}$ is uniformly bounded, then passing to the limit as $\eta \rightarrow 0$ in inequality (4.14) we get the statement. If $\left\{\tilde{y}_{\eta}\right\}_{\eta}$ is unbounded, eventually passing to a subsequence, we can assume that $\left|\tilde{y}_{\eta}\right| \rightarrow+\infty$ as $\eta \rightarrow 0$. Arguing as in (4.9), we get

$$
-\mathcal{L}\left(\tilde{y}_{\eta}, D_{y} \chi\left(\tilde{y}_{\eta}\right), D_{y y}^{2} \chi\left(\tilde{y}_{\eta}\right)\right) \rightarrow+\infty \quad \text { as } \eta \rightarrow 0 .
$$

In particular, for $\eta$ sufficiently small, there holds

$$
-\mathcal{L}\left(\tilde{y}_{\eta}, D_{y} \chi\left(\tilde{y}_{\eta}\right), D_{y y}^{2} \chi\left(\tilde{y}_{\eta}\right)\right) \geq 0 .
$$

Replacing this inequality in relation (4.14), we get the statement.
Conclusion of the proof. By the arbitrariness of the test function $\psi$ and of the point $(\bar{t}, \bar{x})$, we obtain that $\bar{V}$ is a (viscosity) subsolution of the parabolic equation in (1.3).

### 4.3 Step 3

Proposition 4.2 There holds

$$
\bar{V}(T, x) \leq \bar{g}(x) \leq \underline{V}(T, x) \quad \forall x \in \mathbb{R}^{n} .
$$

Proof. We shall prove only the former inequality since the latter is analogous. We first recall from Lemma 4.1 that, for $t>0$, the function $\bar{V}$ is independent of $y$; hence, also $\bar{V}(T, x, y)$ is independent of $y$. Fix $\bar{x} \in \mathbb{R}^{n}$; for every $r>0$ sufficiently small, we define

$$
g^{r}(y):=\sup _{|x-\bar{x}| \leq r} g(x, y)
$$

and we observe that assumption (A4) ensures that $g^{r}$ is a continuous bounded function with

$$
\left|g^{r}(y)-g(x, y)\right| \leq \omega(r) \quad \forall x \in B_{r}(\bar{x}) .
$$

This implies that

$$
\begin{equation*}
-\omega(r)+g^{r}(y) \leq g(\bar{x}, y) \leq \omega(r)+g^{r}(y) \quad \forall y \in \mathbb{R}^{3} . \tag{4.15}
\end{equation*}
$$

We also introduce the parabolic Cauchy problem

$$
\begin{cases}\partial_{t} w^{r}-\operatorname{tr}\left(\sigma \sigma^{T} D_{y y}^{2} w^{r}\right)-b \cdot D_{y} w^{r}=0 & \text { in }(0, \infty) \times \mathbb{R}^{3}  \tag{4.16}\\ w^{r}(0, y)=g^{r}(y) & \text { on } \mathbb{R}^{3} ;\end{cases}
$$

by standard arguments, it admits exactly one bounded solution. On the other hand, let us also consider the problem

$$
\begin{cases}\partial_{t} w^{\prime}-\operatorname{tr}\left(\sigma \sigma^{T} D_{y y}^{2} w^{\prime}\right)-b \cdot D_{y} w^{\prime}=0 & \text { in }(0, \infty) \times \mathbb{R}^{3}  \tag{4.17}\\ w^{\prime}(0, y)=g(\bar{x}, y) & \text { on } \mathbb{R}^{3} ;\end{cases}
$$

we recall from [16, Theorem4.2] that there holds $\lim _{t \rightarrow \infty} w^{\prime}(t, y)=\bar{g}(\bar{x})$ locally uniformly in $y$; in particular, for every $\eta>0$ and $R>1$, there exists $\tau>0$ such that

$$
\begin{equation*}
\left|w^{\prime}(t, y)-\bar{g}(\bar{x})\right| \leq \eta \quad \forall(t, y) \in(\tau, \infty) \times B_{R}(0) \tag{4.18}
\end{equation*}
$$

Moreover, by relation (4.15), one can easily show that $w^{r}(t, y) \pm \omega(r)$ are respectively a supersolution and a subsolution to problem (4.17); hence, the comparison principle yields

$$
\left|w^{\prime}(t, y)-w^{r}(t, y)\right| \leq \omega(r) \quad \forall(t, y) \in(0, \infty) \times \mathbb{R}^{3}
$$

By the last inequality and (4.18) we deduce that: for every $\eta>0$ and $R>1$, there exists $\tau>0$ such that

$$
\begin{equation*}
\left|w^{r}(t, y)-\bar{g}(\bar{x})\right| \leq \eta+\omega(r) \quad \forall(t, y) \in(\tau, \infty) \times B_{R}(0) \tag{4.19}
\end{equation*}
$$

For later use, we introduce some notations; we set $Q^{r}:=(T-r, T) \times$ $B_{r}(\bar{x})$ and let $M \in \mathbb{R}$ be sufficiently large that for every $(t, x, y) \in(0, T) \times$ $B_{r}(\bar{x}) \times \mathbb{R}^{3}$ :

$$
\begin{align*}
\left\|w^{\prime}\right\|_{\infty} & \leq M  \tag{4.20}\\
\left\|w^{r}\right\|_{\infty} & \leq M  \tag{4.21}\\
\left|V^{\epsilon}(t, x, y)\right| & \leq M \tag{4.22}
\end{align*}
$$

Consider also a smooth function $\psi_{0}=\psi_{0}(x)$ (namely, it is independent of $t$ and of $y$ ) such that

$$
\left\{\begin{array}{l}
\psi_{0}(\bar{x})=0,  \tag{4.23}\\
\psi_{0}(x) \geq 0 \quad \forall x \in B_{r}(\bar{x}), \\
\psi_{0}(x) \geq M-\inf _{(z, y) \in B_{r}(\bar{x}) \times \mathbb{R}^{3}} g(z, y) \quad \forall x \in \partial B_{r}(\bar{x}) .
\end{array}\right.
$$

Let $C>0$ be a constant such that

$$
\begin{equation*}
\left|H\left(x, y, D_{x} \psi_{0}(x), D_{x x}^{2} \psi_{0}(x), 0\right)\right| \leq C \quad \forall(x, y) \in B_{r}(\bar{x}) \times \mathbb{R}^{3} \tag{4.24}
\end{equation*}
$$

For $(t, x, y) \in Q^{r} \times \mathbb{R}^{3}$, we define

$$
\psi^{\epsilon}(t, x, y):=w^{r}\left(\frac{T-t}{\epsilon}, y\right)+\psi_{0}(x)+C_{1}(T-t)
$$

with $C_{1}:=C+a M$. We claim that the function $\psi^{\epsilon}$ is a supersolution to the following initial-boundary value problem

$$
\left\{\begin{array}{c}
(i) \quad-\partial_{t} \Psi+H\left(x, y, D_{x} \Psi, D_{x x}^{2} \Psi, D_{x y}^{2} \Psi / \sqrt{\epsilon}\right)-\frac{1}{\epsilon} \mathcal{L}\left(y, D_{y} \Psi, D_{y y}^{2} \Psi\right)  \tag{4.25}\\
\quad+a \Psi=0 \quad \text { in } Q^{r} \times \mathbb{R}^{3} \\
(i i) \quad \Psi(T, x, y)=g(x, y) \quad \text { on } B_{r}(\bar{x}) \times \mathbb{R}^{3} \\
(i i i) \quad \Psi(t, x, y)=M \quad \text { on }(T-r, T) \times \partial B_{r}(\bar{x}) \times \mathbb{R}^{3} .
\end{array}\right.
$$

Assume for the moment that this claim is true. On the other hand, the function $V^{\epsilon}$ is a subsolution to problem (4.25); therefore, by comparison principle (see, for instance, $[2$, Proposition 1 (proof)]), we get

$$
V^{\epsilon}(t, x, y) \leq \psi^{\epsilon}(t, x, y) \quad \forall(t, x, y) \in(T-r, T) \times B_{r}(\bar{x}) \times \mathbb{R}^{3}
$$

For $y \in B_{R / 2}(0)$ ( $R$ is defined in (4.19)) and $t \in(T-r, T)$, we get

$$
\begin{aligned}
\bar{V}(t, x, y) & =\limsup _{\epsilon \rightarrow 0^{+}, t^{\prime} \rightarrow t, x^{\prime} \rightarrow x, y^{\prime} \rightarrow y} V^{\epsilon}\left(t^{\prime}, x^{\prime}, y^{\prime}\right) \\
& \leq \limsup _{\epsilon \rightarrow 0^{+}, t^{\prime} \rightarrow t, x^{\prime} \rightarrow x, y^{\prime} \rightarrow y} \psi^{\epsilon}\left(t^{\prime}, x^{\prime}, y^{\prime}\right) \\
& \leq \limsup _{\epsilon \rightarrow 0^{+}, t^{\prime} \rightarrow t, y^{\prime} \rightarrow y} w^{r}\left(\frac{T-t^{\prime}}{\epsilon}, y^{\prime}\right)+\limsup _{t^{\prime} \rightarrow t, x^{\prime} \rightarrow x}\left[\psi_{0}\left(x^{\prime}\right)+C_{1}\left(T-t^{\prime}\right)\right] \\
& \leq \bar{g}(\bar{x})+\eta+\omega(r)+\psi_{0}(x)+C_{1}(T-t)
\end{aligned}
$$

where the last inequality is due to relation (4.19) (observe that definitely $y^{\prime} \in B_{R}(0)$ and $\left.\frac{T-t^{\prime}}{\epsilon}>\tau\right)$ and to the continuity of $\psi_{0}$. Since $\bar{V}$ is independent of $y$, we deduce

$$
\bar{V}(t, x) \leq \bar{g}(\bar{x})+\eta+\omega(r)+\psi_{0}(x)+C_{1}(T-t) .
$$

Passing to the $\limsup _{t^{\prime} \rightarrow T^{-}, x^{\prime} \rightarrow \bar{x}, y^{\prime} \rightarrow y}\left(\right.$ recall $\left.\psi_{0}(\bar{x})=0\right)$, we infer

$$
\bar{V}(T, \bar{x}) \leq \bar{g}(\bar{x})+\eta+\omega(r)
$$

By the arbitrariness of $\eta$ and of $r$, we get

$$
\bar{V}(T, \bar{x}) \leq \bar{g}(\bar{x})
$$

which is equivalent to our statement.
Let us now pass to prove the claim: $\psi^{\epsilon}$ is a supersolution to problem (4.25). First we check the initial-boundary conditions (ii) and (iii).

In order to prove (ii), it suffices to note that the definition of $g^{r}$ and the second property in (4.23) entail

$$
\psi^{\epsilon}(T, x, y)=w^{r}(0, y)+\psi_{0}(x) \geq w^{r}(0, y) \geq g(x, y) \quad \forall(x, y) \in B_{r}(\bar{x}) \times \mathbb{R}^{3}
$$

In order to prove (iii), we observe that $g^{r}(y) \geq \inf _{(x, z) \in B_{r}(\bar{x}) \times \mathbb{R}^{3}} g(x, z)$; hence, the comparison principle for (4.16) yields

$$
w^{r}(t, y) \geq \inf _{(x, y) \in B_{r}(\bar{x}) \times \mathbb{R}^{3}} g(x, y) \quad \forall(t, y) \in(0, \infty) \times \mathbb{R}^{3}
$$

Taking also into account the third property in (4.23), we conclude

$$
\psi^{\epsilon}(t, x, y)=w^{r}\left(\frac{T-t}{\epsilon}, y\right)+\psi_{0}(x)+C_{1}(T-t) \geq M
$$

for every $(t, x, y) \in(T-r, T) \times \partial B_{r}(\bar{x}) \times \mathbb{R}^{3}$.
Now we prove (i). Let us assume that $w^{r}$ is a classical solution to (4.16). In this case, in $(T-r, T) \times B_{r}(\bar{x}) \times \mathbb{R}^{3}$ there holds

$$
\begin{aligned}
& -\partial_{t} \psi^{\epsilon}+H\left(x, y, D_{x} \psi^{\epsilon}, D_{x x}^{2} \psi^{\epsilon}, D_{x y}^{2} \psi^{\epsilon} / \sqrt{\epsilon}\right)-\frac{1}{\epsilon} \mathcal{L}\left(y, D_{y} \psi^{\epsilon}, D_{y y}^{2} \psi^{\epsilon}\right)+a \psi^{\epsilon} \\
& =\frac{1}{\epsilon}\left[\partial_{t} w^{r}-\mathcal{L}\left(y, D_{y} w^{r}, D_{y y}^{2} w^{r}\right)\right]+C_{1}+H\left(x, y, D_{x} \psi_{0}, D_{x x}^{2} \psi_{0}, 0\right)+a \psi^{\epsilon} \\
& =C_{1}+H\left(x, y, D_{x} \psi_{0}, D_{x x}^{2} \psi_{0}, 0\right)+a w^{r}\left(\frac{T-t}{\epsilon}, y\right)+a \psi_{0}(x)+a C_{1}(T-t) \\
& =C_{1}+a w^{r}\left(\frac{T-t}{\epsilon}, y\right)+H\left(x, y, D_{x} \psi_{0}, D_{x x}^{2} \psi_{0}, 0\right) \\
& \geq C_{1}-C-a M \\
& \geq 0
\end{aligned}
$$

where we used the definition of $C_{1}:=C+a M$, the definition of $C$ in (4.24) and the definition of $M$ (4.20).
In the case when $w^{r}$ is only a viscosity solution to (4.16), we can accomplish the proof following the same arguments of [2, Theorem 3].

### 4.4 Step 4

We proved that $\bar{V}$ and $\underline{V}$ are respectively a subsolution and a supersolution of (1.3). WE can apply the comparison principle to equation (1.3) that holds since $\bar{H}$ defines a degenerate elliptic equation and $H$ is Lipschitz continuous
on $x$ (see (2.1)). Then $\bar{V}(t, x) \leq \underline{V}(t, x)$; by definition the reverse inequality is obvious. Then $V=\bar{V}(t, x)=\underline{V}(t, x)$ is the unique continuous solution of the parabolic equation (1.3) and the local uniform convergence follows from standards arguments.

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