

Non-Linear Analytical Model for a Multi-V-Shape IPM with Concentrated Winding

P. Akiki, M. Hage Hassan, J-C Vannier, M. Bensetti, B. Dagusé, D. Prieto, M. McClelland

Abstract -- This paper presents a non-linear analytical model of a multi-V-shape Interior Permanent Magnet (IPM) motor with non-overlapping concentrated winding. The model relies on Maxwell equations to compute the flux density in the different parts of the motor. This article proposes a saturated analytical model of the stator and the rotor. The analytical model is used to calculate the average torque, the power factor and the voltage of the motor. It is compared to 2D Finite Element Analysis (FEA) and shows very good results. The developed model is 5 times faster than the FEA thus, it can be used in optimization procedures.

Index Terms – Analytical model, Concentrated winding, Electrical machines, IPM motor, Multi-V-shape magnets.

I. NOMENCLATURE

m	: Number of coils per phase per electrical period
N_s	: Number of turns per coil
N_{cs}, N_{cp}	: Number of circuits in series and in parallel
L	: Stack length
e	: Airgap length
α	: Current angle
I_C, I_B, I_C	: Currents in phases A, B and C (rms)
i_d, i_q	: d-q axis current components

II. INTRODUCTION

Recently, high performance motors have been the subject of many studies and research projects. Permanent magnet motors with rare earth materials are known for their suitability in high performance applications and have been widely used by industrials during the last decades [1], [2]. However, these magnets have reached high price, which increased the cost of the motor. In order to ensure competitiveness, varieties of motors, which do not use or with reduced rare-earth materials, have been extensively developed. Flux-switching motors offered the possibility to reduce the volume of the magnets by putting them on the stator side. The rotor is magnet-free and designed as a passive magnetic circuit with poles similar to the one used in switched reluctance motors [3]. Non-rare-earth permanent magnets such as ferrite or alnico have attracted significant interest in recent designs. The low magnetic field created by these types of magnet leads to their use in a flux concentration configuration [4] or as an additional torque source in

Paul Akiki, Maya Hage Hassan, Jean-Claude Vannier and Mohamed Bensetti are with GeePs | Group of electrical engineering-Paris, UMR CNRS 8507, CentraleSupélec, Univ.Paris-Sud, Université Paris-Saclay, Sorbonne Universités, UPMC Univ Paris 06 3 & 11 rue Joliot-Curie, Plateau de Moulon 91192 Gif-sur-Yvette, France.

Paul Akiki is also with Leroy Somer Motor, 16000 Angoulême, France. (emails :paul.akiki@centralesupelec.fr, maya.hagehassan@centralesupelec.fr, jean-claude.vannier@centralesupelec.fr, mohamed.bensetti@centralesupelec.fr)

Benjamin Dagusé, Dany Prieto and Mike McClelland are with Leroy Somer Motor, 16000 Angoulême, France. (e-mails: benjamin.daguse@emerson.com, dany.prieto@emerson.com, mike.mcclelland@emerson.com)

synchronous reluctance motors. This configuration is known as the Permanent Magnet Assisted Synchronous Reluctance Motor (PMASynRM) with distributed winding [5]–[7].

In this paper, ferrite magnets are used in the rotor with short end-winding in the stator in order to reduce the Joule losses. Therefore, the studied structure is a multi-V-shape Interior Permanent Magnet (IPM) motor with non-overlapping concentrated winding. The motor's design process usually starts by modelling the magnetic performance. The advances in computational tools made it possible to use Finite Element Analysis (FEA) to create a numerical simulation of the studied motor. FEA allows a detailed investigation of the motor behavior, with the disadvantage of generally long execution times [8]. Other modelling approaches exist, such as the reluctance network that represents the flux paths through the motor parts [9], [10] and the analytical model based on solving Maxwell equations. In [11] authors proposed an analytical model for an IPM motor in linear conditions and the method was improved in [12] to take into account the non-linear behavior of the iron in a PMASynRM with distributed windings.

In this paper, an analytical model of a multi-V-shape IPM motor with non-overlapping concentrated winding is proposed. It is an 18 slot/16 pole structure (Fig.1). It has open slots in order to facilitate the automatic insertion of the windings during the manufacturing process. We can identify the existence of an interior iron bridge and an exterior one. A flux-barrier is mainly defined by 2 parameters: δ is the barrier's tilt angle and $k = \theta/\beta$ is the ratio of the barrier's opening angle with respect to the half pole angle β . The model relies on two fundamental laws of electromagnetism: Ampere's theorem and flux conservation law [12]. This article proposes a saturated analytical model of the stator teeth, the stator yoke, the leakage flux in the slots and the saturation in the rotor magnetic circuit. The developed model is parameterized as a function of the slot number, the pole pair number and the magnet layer number.

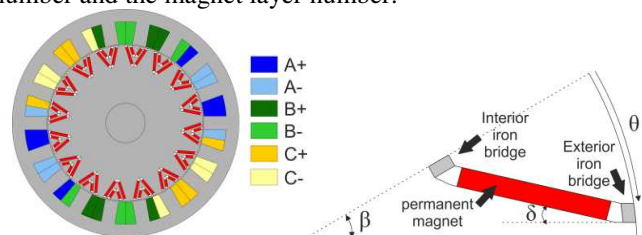


Fig.1 Studied motor structure with the rotor's main geometrical parameters

III. ANALYTICAL MODEL

The analytical model computes the flux density of the motor over one electrical period (360° elec). Due to magnetic symmetry, half of the machine is modeled.

A. Magnetomotive Force

The machine is a three phase motor with a non-overlapping concentrated tooth winding. The magnetomotive force (mmf)

is determined with respect to the winding pattern shown in Fig.2 that covers one electrical period. The *mmf* level of each tooth is given by the total current of the tooth coil as in (1). A linear variation of the *mmf* in the slots is considered.

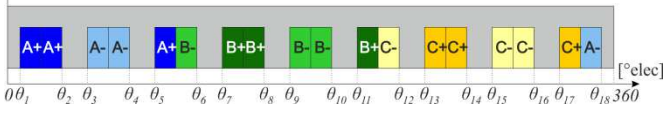


Fig.2 Winding configuration for half of the motor (360°elec)

$$mmf(\theta) = \begin{cases} +I_{A_{coil}} & \forall \theta \in [0, \theta_1] \\ -I_{A_{coil}} & \forall \theta \in [\theta_2, \theta_3] \\ +I_{A_{coil}} & \forall \theta \in [\theta_4, \theta_5] \\ +I_{B_{coil}} & \forall \theta \in [\theta_6, \theta_7] \\ -I_{B_{coil}} & \forall \theta \in [\theta_8, \theta_9] \\ +I_{B_{coil}} & \forall \theta \in [\theta_{10}, \theta_{11}] \\ +I_{C_{coil}} & \forall \theta \in [\theta_{12}, \theta_{13}] \\ -I_{C_{coil}} & \forall \theta \in [\theta_{14}, \theta_{15}] \\ +I_{C_{coil}} & \forall \theta \in [\theta_{16}, \theta_{17}] \\ +I_{A_{coil}} & \forall \theta \in [\theta_{18}, 360] \end{cases} ; \begin{cases} I_{A_{coil}} = \sqrt{2} I_A N_s \\ I_{B_{coil}} = \sqrt{2} I_B N_s \\ I_{C_{coil}} = \sqrt{2} I_C N_s \end{cases} \quad (1)$$

B. Stator Slots Additional Air-Gap

Stator slotting influences the magnetic field by affecting the distribution of the flux density in the air-gap. It increases the reluctance of the flux path and reduces the air-gap flux density in front of a slot. Therefore, it is important to take it into account. A common method for modeling the stator slotting effect uses Carter's coefficient to smooth the stator geometry but assumes that the slot width is much smaller than its height [13]. In this paper, the additional air-gap length is calculated using the mean flux path under the slot opening [14]. Thus, idealized flux paths are shown in Fig.3. Assuming quarter-circular flux lines with radii R_1 and R_2 , the arc lengths Γ_1 and Γ_2 are given by:

$$\begin{cases} \Gamma_1(v) = \frac{\pi}{2} R_{si} \sin(v) \\ \Gamma_2(v) = \frac{\pi}{2} R_{si} \sin(\theta_{slot} - v) \end{cases} \quad (2)$$

Where θ_{slot} is the slot opening angle, R_{si} is the internal radius of the stator and v is the angle under the slot opening. The additional air-gap $e_s(v)$ is the equivalent length of the two parallel paths Γ_1 and Γ_2 . The use of the first-order Taylor expansion leads to (3):

$$e_s(v) = \begin{cases} \frac{\pi}{2\theta_{slot}} R_{si} v (\theta_{slot} - v) & \text{under slot opening} \\ 0 & \text{under stator tooth} \end{cases} \quad (3)$$

The total air-gap length is given by:

$$e_{tot}(v) = e + e_s(v) \quad (4)$$

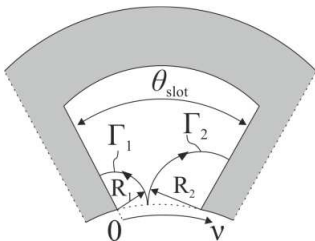


Fig.3 Idealized flux paths under stator slot

C. Flux density of the stator

The flux density of the stator is calculated in the teeth and the yoke. Additionally, the tangential leakage flux that links two adjacent teeth is taken into account. Assuming quarter-circular paths between the air-gap and the stator tooth (Fig.4), the tooth flux density is calculated using flux conservation law applied to the tooth pitch (θ_{st}) with a discretization of 1° elec.

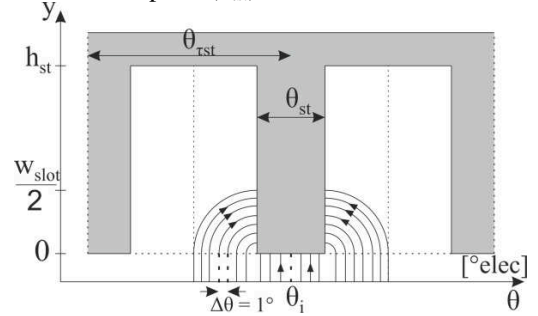


Fig.4 Flux lines between the airgap and the stator tooth

Equation (5) computes the stator tooth flux density $B_{st_{ag}}^i$ due to airgap flux density $B_{ag}(\theta)$:

$$\forall i \in \{1,9\} ; \forall n \in \left\{ \theta_i + \frac{\theta_{st}}{2}, \frac{\theta_{rst}}{2} - \frac{\theta_{st}}{2} \right\}$$

$$B_{st_{ag}n}^i = B_{st_{ag}n-1}^i + \frac{R_{si}}{w_{st}} \left(\int_{\theta_i+n}^{\theta_i+(n+1)} B_{ag}(\theta) d\theta + \int_{\theta_i-(n+1)}^{\theta_i-n} B_{ag}(\theta) d\theta \right) \quad (5)$$

Where θ_{st} is the tooth angular span and θ_i is the center angular position of the i^{th} tooth (Fig.4). The flux density for $n = 0$ is calculated using flux conservation law for the airgap flux under the stator tooth. Its expression is given by (6)

$$B_{st_{ag}0}^i = \frac{R_{si}}{w_{st}} \int_{\theta_i - \frac{\theta_{st}}{2}}^{\theta_i + \frac{\theta_{st}}{2}} B_{ag}(\theta) d\theta \quad (6)$$

Hence, we obtain $B_{st_{ag}}^i(y)$ which is the airgap flux density for each tooth along the y -axis shown in Fig.4, for $0 \leq y \leq \frac{w_{slot}}{2}$ where w_{slot} , w_{st} are respectively, the slot width and the stator tooth width.

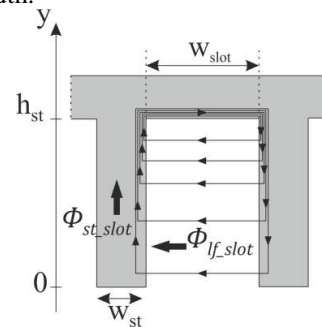


Fig.5 Tangential leakage flux in the stator slot

In order to take the slot leakage flux into consideration, the total current distribution in each slot is considered to be linear. Thus, Ampere's theorem is written on paths along the tooth height h_{st} as shown in Fig.5, and the flux conservation law is applied between the tangential leakage flux in stator slots $\phi_{lf_{slot}}(y)$ and the slot flux entering the tooth $\phi_{st_{slot}}(y)$. This yields (7) and (8).

$$B_{lf_{slot}}(y) = \frac{B_{lf_0}}{C} (e^{Ay} - e^{2Ah_{st}} e^{-Ay}) \quad (7)$$

$$B_{st_{slot}}(y) = \frac{1}{w_{st}} \int_0^y B_{l_{slot}}(u) du \quad (8)$$

$$= \frac{\mu_0 I_{slot}}{w_{slot} w_{st} A C} [e^{Ay} + e^{2Ah_{st}}(e^{-Ay} - 1) - 1]$$

$$\text{With } A = \sqrt{\frac{\mu_0}{w_{slot} w_{st}} \left(\frac{1}{\mu_{st_l}} + \frac{1}{\mu_{ds_r}} \right)}; C = 1 - e^{2Ah_{st}}$$

Where μ_0 , μ_{st_l} , μ_{st_r} , I_{slot} are respectively, the vacuum permeability, the relative permeability of the left and the right stator tooth (surrounding the considered slot) and the total current of the considered slot. Therefore the total leakage flux density and the total flux density of a tooth along its height (y-axis) are given by (9) and (10) respectively.

$$B_{st_{lf}}^i(y) = B_{st_{slot_r}}^i(y) - B_{st_{slot_l}}^i(y) \quad (9)$$

$$B_{st}^i(y) = B_{st_{ag}}^i(y) + B_{st_{lf}}^i(y) \quad (10)$$

The subscripts l and r represent the left and the right sides of a tooth. The tooth flux density is considered constant along the tooth width (θ -axis) so that the stator tooth flux density with respect to the electrical angle can be written:

$$B_{st}(\theta, y) = \begin{cases} B_{st}^i(y) & \text{under the } i^{\text{th}} \text{ tooth} \\ 0 & \text{under the slots} \end{cases} \quad (11)$$

The stator yoke flux density is calculated by applying the flux conservation equation between the stator teeth and the yoke over a variation of $\Delta\theta = 1^\circ \text{elec}$. The stator yoke flux density with respect to the electrical angle $B_{sy}(\theta)$ is expressed by (12) for $0 \leq \theta \leq 360^\circ \text{elec}$:

$$B_{sy}(\theta + \Delta\theta)h_{sy} = B_{sy}(\theta)h_{sy} + B_{st}(\theta, y) R_{yi} \Delta\theta \quad (12)$$

Where h_{sy} is the yoke height and R_{yi} is the yoke interior radius.

D. Flux density of the rotor

The flux density of the rotor has to be calculated for 8 poles, since we are dealing with the half of the machine. Fig.6 shows one rotor pole with V-shape flux-barriers that are filled with magnets. The iron piece between two flux-barriers is called "flux-guide" (green area in Fig.6).

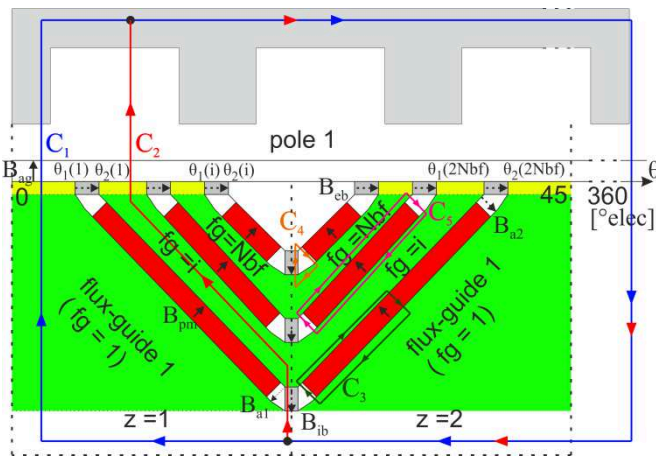


Fig.6 Motor general schematic and flux paths

Each rotor pole is divided into two zones ($z = 1, 2$) and the rotor flux densities are modeled by the flux-guide mean path. At the flux-guide boundaries, the magnetic flux is assumed to be aligned in the same direction as the flux-guide mean path (Fig.7a). Nbf is the number of flux barriers. The flux is defined as $\phi^{(z,fg,pol)}$ where fg is the flux-guide number (fg

$\in \{1, Nbf\}$), and pol is the pole number ($pol \in \{1, 8\}$). For the flux-guide between poles ($fg = 1$), (13) and (14) are written:

$$\phi_{fg_f}^{(1,1,pol)} = \phi_{ag}^{(1,1,pol)} + \phi_{ag}^{(2,1,pol-1)} + \phi_{eb_0}^{(1,1,pol-1)} - \phi_{eb_f}^{(2,1,pol-1)} \quad (13)$$

$$\phi_{fg_0}^{(1,1,pol)} = \phi_{fb}^{(1,1,pol)} + \phi_{fb}^{(2,1,pol-1)} + \phi_{fg_f}^{(1,1,pol)} \quad (14)$$

With $pol = 1$ if $pol-l = 8$. For flux-guides between flux-barriers ($fg \in \{2, Nbf\}$), (15) and (16) are written with $z \in \{1, 2\}$:

$$\phi_{fg_f}^{(z,fg,pol)} = \phi_{ag}^{(z,fg,pol)} - (2-z)\phi_{eb_f}^{(z,fg-1,pol)} + (2-z)\phi_{eb_0}^{(z,fg,pol)} + (z-1)\phi_{eb_0}^{(z,fg-1,pol)} - (z-1)\phi_{eb_f}^{(z,fg,pol)} \quad (15)$$

$$\phi_{fg_0}^{(z,fg,pol)} = \phi_{ib}^{(fg-1,pol)} - \phi_{ib}^{(fg,pol)} + \phi_{fb}^{(z,fg,pol)} - \phi_{fb}^{(z,fg-1,pol)} + \phi_{fg_f}^{(z,fg,pol)} \quad (16)$$

Where ϕ_{fg_f} , ϕ_{fg_0} are the boundary fluxes of the rotor flux-guides. ϕ_{eb_f} , ϕ_{eb_0} are the boundary fluxes of the exterior iron bridges. ϕ_{ib} and ϕ_{ag} are the fluxes of the interior iron bridge and the airgap respectively. ϕ_{fb} is the flux crossing the flux-barrier.

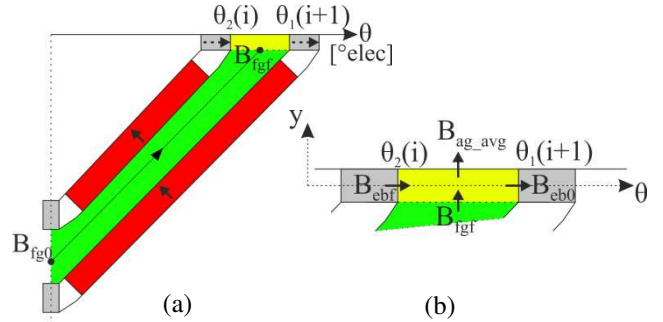


Fig.7 Flux density mean path in a rotor flux-guide (a). Local saturation near the exterior iron bridge (b)

The boundary flux densities are then given by (17) - (20)

$$B_{fg_f}^{(1,1,pol)} = \frac{\phi_{fg_f}^{(1,1,pol)}}{2w_{fg_f}(1)} = B_{fg_f}^{(2,1,pol-1)} \quad (17)$$

$$B_{fg_0}^{(1,1,pol)} = \frac{\phi_{fg_0}^{(1,1,pol)}}{2w_{fg_0}(1)} = B_{fg_0}^{(2,1,pol-1)} \quad (18)$$

$$B_{fg_f}^{(z,fg,pol)} = \frac{\phi_{fg_f}^{(z,fg,pol)}}{w_{fg_f}(fg)} \quad (19)$$

$$B_{fg_0}^{(z,fg,pol)} = \frac{\phi_{fg_0}^{(z,fg,pol)}}{w_{fg_0}(fg)} \quad (20)$$

w_{fg_f} and w_{fg_0} are the boundary widths of the flux-guides. A linear interpolation is used between the boundary flux densities in order to compute the interior flux densities of each flux-guide.

The iron bridges are generally saturated. Hence, the mmf drop of the rotor flux-guide near the iron bridge is higher than other areas. To take into account this local saturation phenomenon, the area near the exterior iron bridges is

modeled separately (Fig.7b). We assume a linear variation of the flux density between two bridges in the θ -axis ($B_x(\theta)$). The boundary flux density of the flux-guide B_{fgf} and the average value of the airgap flux density B_{agavg} are used to calculate an average value of the flux-guide flux density in the y-axis:

$$B_y = \frac{B_{agavg} + B_{fgf}}{2} \quad (21)$$

Thus, for the electrical angle between $\theta_2(i)$ and $\theta_1(i+1)$ the flux-guide local saturation flux density is given by:

$$B_{local_{sat}}(\theta) = \sqrt{B_x(\theta)^2 + B_y^2} \quad (22)$$

E. Equations of the global system

The motor general scheme is shown in Fig.6. The magnets are inserted in the V-shape flux-barriers and are surrounded by air from both sides. The model considers the non-linearity of the ferromagnetic material. The electrical angular position θ is discretized from 0 to 360° along the airgap. Ampere's theorem is written on five contours (C1 to C5) and expressed by (23) to (27), respectively.

$$H_{ag}^{(\theta)} e_{tot}^{(\theta)} - H_{ag}^{(360)} e_{tot}^{(360)} + At_s^{(\theta)} - At_s^{(360)} + At_r^{(\theta)} - At_r^{(360)} + mmf^{(\theta)} - mmf^{(360)} = 0 \quad (23)$$

$$H_{ag}^{(\theta)} e_{tot}^{(\theta)} - H_{ag}^{(360)} e_{tot}^{(360)} - \sum_{m=1}^{fg-1} (H_{ib}^{(m,pol)} h_{ib}^{(m)}) + At_s^{(\theta)} - At_s^{(360)} + At_r^{(\theta)} - At_r^{(360)} + mmf^{(\theta)} - mmf^{(360)} = 0 \quad (24)$$

$$-H_{a1}^{(z,nb,pol)} w_{fb}^{(nb)} + (-1)^{pol} H_{pm}^{(z,nb,pol)} w_{pm}^{(nb)} + At_r^{(z,nb,pol)} = 0 \quad (25)$$

$$-H_{ib}^{(nb,pol)} h_{ib}^{(nb)} + H_{a1}^{(z,nb,pol)} w_{fb}^{(nb)} = 0 \quad (26)$$

$$H_{a2}^{(z,nb,pol)} w_{fb}^{(nb)} - H_{a1}^{(z,nb,pol)} w_{fb}^{(nb)} + At_r^{(z,nb,pol)} = 0 \quad (27)$$

Where nb is the number of the flux-barriers varying from 1 to Nbf . H , w and h are respectively the magnetic field, the width and the height. The subscripts pm , fb , ib , $a1$ and $a2$ represent the permanent magnet, the flux-barrier, the interior iron bridge and the air next to the magnet sides, respectively. At_s and At_r are the mmf drop of the stator path (tooth + yoke) and the rotor path respectively. The above equations are completed with (28) that express the flux conservation law in the airgap.

$$\langle B_{ag} \rangle = 0 \quad (28)$$

Equations (23) and (24) are not valid for the exterior iron bridge angular positions. The equation system above the exterior iron bridges is given in the next section (III-F).

The flux of the interior iron bridges is calculated by applying the flux conservation law in the central iron piece and between the flux-barriers for $pol \in \{1,8\}$:

Central iron piece:

$$\sum_{z=1}^2 \left\{ (-1)^{pol} \phi_{pm}^{(z,Nbf,pol)} + \phi_{fa1}^{(z,Nbf,pol)} + \phi_{fa2}^{(z,Nbf,pol)} \right\} + \phi_{ag}^{(1,Nbf+1,pol)} + \phi_{eb0}^{(2,Nbf,pol)} - \phi_{ebf}^{(1,Nbf,pol)} + \phi_{ib}^{(Nbf)} = 0 \quad (29)$$

Between two flux-barriers: $fg \in \{2, Nbf\}$

$$\sum_{z=1}^2 \left\{ \begin{aligned} &(-1)^{pol} \phi_{pm}^{(z,fg-1,pol)} + \phi_{a1}^{(z,fg-1,pol)} \\ &+ \phi_{a2}^{(z,fg-1,pol)} + \phi_{ag}^{(z,fg,pol)} \\ &- (-1)^{pol} \phi_{pm}^{(z,fg,pol)} - \phi_{a1}^{(z,fg,pol)} - \phi_{a2}^{(z,fg,pol)} \end{aligned} \right\} + \phi_{eb0}^{(1,fg,pol)} + \phi_{eb0}^{(2,fg-1,pol)} - \phi_{ebf}^{(1,fg-1,pol)} - \phi_{ebf}^{(2,fg,pol)} - \phi_{ib}^{(fg)} + \phi_{ib}^{(fg-1)} = 0 \quad (30)$$

It is worth mentioning that (29) and (30) take into account the alternating magnetic orientation of the permanent magnet among poles.

F. Flux density of the exterior iron bridge

To complete the global system, the airgap flux density above the exterior iron bridge has to be computed. The flux conservation law (31) and Ampere's theorem (32) are used on the iron bridge as shown in Fig.8. The mmf drop in the stator iron is included when applying Ampere's theorem.

$$B_{ag}(\theta_i) = \frac{h_{eb}}{R_{rot} \Delta\theta} (B_{eb}(\theta_i) - B_{eb}(\theta_{i+1})) \quad (31)$$

$$1 \leq i \leq n+1$$

$$R_{rot} \int_{\theta_1}^{\theta_i} H_{eb}^{(\theta)} d\theta + \frac{B_{ag}^{(\theta_i)}}{\mu_0} e_{tot}^{(\theta_i)} - \frac{B_{ag}^{(\theta_1)}}{\mu_0} e_{tot}^{(\theta_1)} + At_s^{(\theta_i)} - At_s^{(\theta_1)} + mmf^{(\theta)} - mmf^{(\theta_1)} = 0 \quad (32)$$

$$2 \leq i \leq n+2$$

Where $(n+2)$ is the electrical angular span of the iron bridge. R_{rot} is the rotor radius and h_{eb} is the height of the exterior iron bridge. Equation (32) is completed with the boundary value $B_{ag}(\theta_1)$. The airgap flux density above the iron bridge is then deduced by (31).

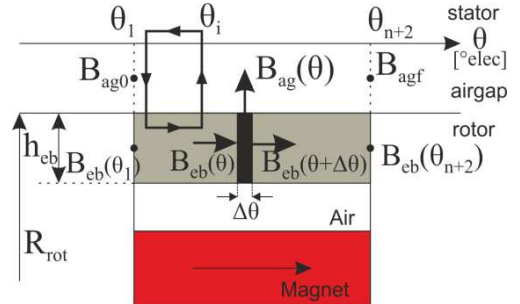


Fig.8 Exterior iron bridge: Ampere's theorem and flux conservation law

G. Performance calculation

The airgap flux density is used to calculate the stator phase flux linkages (ϕ_a , ϕ_b , ϕ_c). The slot leakage flux is added to obtain the total flux linked by the stator coils. Equation (33) gives the expression of phase-A flux linkage. For the other phases it is obtained by replacing the subscript A by B and C.

$$\phi_a = \frac{N_{cs}}{N_{cp}} N_s \sum_{k=1}^m R_{si} L \int_{\theta_{A-k}}^{\theta_{A+k}} B_{ag}(\theta) d\theta + \phi_{st_{lf}} \quad (33)$$

Where θ_{A+} and θ_{A-} are the center angular position of the coils A+ and A- (Fig.2). The d-q axis flux linkages (ϕ_d , ϕ_q) are then computed using Park's transformation. The average torque (T_{avg}), the power factor (PF) and the voltage (V) are

computed at 4 static positions for an accurate estimation [15]. They are expressed by (34), (35) and (36) respectively.

$$T_{avg} = \frac{3p}{2}(\phi_d i_q - \phi_q i_d) \quad (34)$$

$$PF = \cos\left(\tan^{-1}\left(-\frac{\phi_d}{\phi_q}\right) - \tan^{-1}\left(\frac{i_q}{i_d}\right)\right) \quad (35)$$

$$V = \omega \sqrt{\phi_d^2 + \phi_q^2} \quad (36)$$

IV. FINITE ELEMENT COMPARISON

A 2D finite element model of the studied motor is created and its results are compared with the results of the analytical model. The first motor geometry (M_1) is reported in the appendix. At first, the magnetomotive force (mmf) is verified and the airgap flux density is shown for two pairs of current values and current angles. Then, the flux density of the stator teeth is shown and the flux density of the rotor flux-guides and exterior iron bridges are reported. Finally, the comparison of the average torque, the power factor and the voltage is conducted for M_1 and another geometry M_2 in order to verify the robustness of the analytical model.

The studied motors have 2 V-shape flux-barriers filled with ferrite magnets and rotate at 500 rpm.

A. Magnetomotive force

In order to verify the validity of the magnetomotive force model (1) and the stator permeance function (3), the system formed by (23) and (28) was solved under linear conditions with a passive rotor disc. Thus, the terms $At_s^{(\theta)}$ and $At_r^{(\theta)}$ do not exist. Fig.9 shows the airgap flux density for a current value of 43A (rms) and a current angle of 0° .

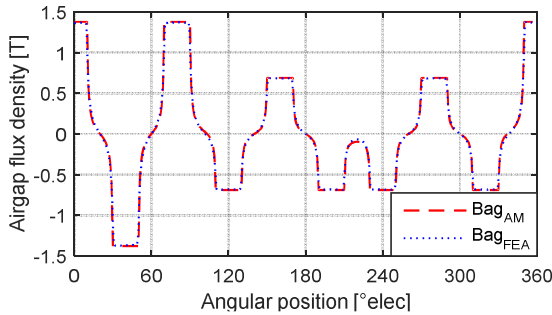


Fig.9 Airgap flux density under linear conditions and passive rotor iron disc

In this particular configuration, the airgap flux density is the image of the mmf since no Ampere-turns are consumed in the motor's iron parts. The analytical model (AM) reproduces almost the exact shape of the Finite Element Analysis (FEA) result. This validates the mmf and the stator saliency models for the non-overlapping concentrated tooth winding.

B. Airgap flux density

The airgap flux density of the motor M_1 is obtained by resolving the equation system (23) to (28). A comparison between the analytical model and the finite element model is made. The airgap flux density is given for two current values, 28A (Fig.10) and 43A (Fig.11) with a current angle of 10° elec. The analytical model gives satisfactory results and is accurate for both current values at two saturation levels of the ferromagnetic material (lower saturation at 28A and higher saturation at 43A).

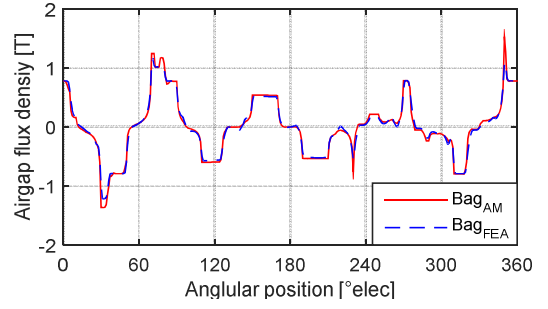


Fig.10 Comparison of the airgap flux density between the AM and the FEA (Irms=28A, $\alpha=10^\circ$ elec)

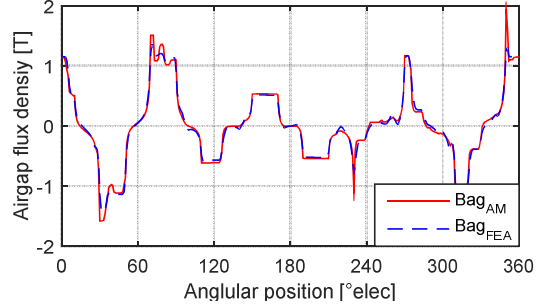


Fig.11 Comparison of the airgap flux density between the AM and the FEA (Irms=43A, $\alpha=10^\circ$ elec)

C. Flux density of the stator teeth

The developed model of the stator teeth that takes into consideration the airgap flux density and the slot tangential leakage flux (section III-C) is computed and compared to the results of the FEA. The comparison is carried out for a current of 43A and a current angle of 10° elec. Since only half of the motor is modeled, the number of teeth is 9. Fig.12 shows the flux density of the 9 teeth along the tooth height.

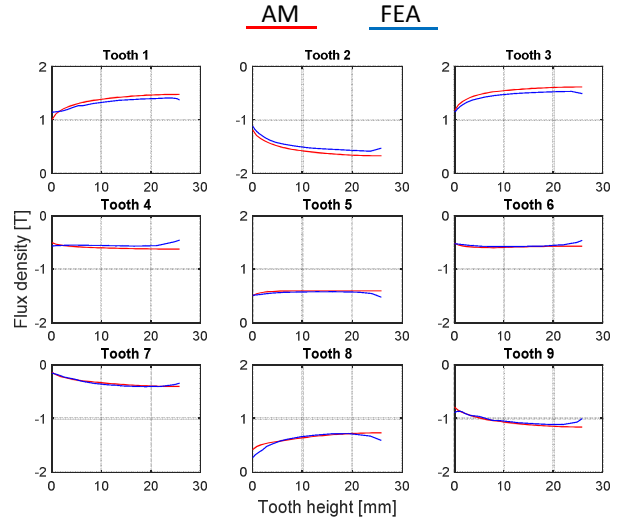


Fig.12 Comparison of the stator teeth flux density between the AM and the FEA (Irms=43A, $\alpha=10^\circ$ elec)

Fig.12 shows that the stator tooth flux density is not constant along the tooth. The analytical model shows a good concordance with the finite element results.

D. Flux density of the exterior iron bridge

Each rotor barrier has two exterior iron bridges. Their flux densities are obtained by the system of equations (32). For a current value of 43A (rms) and a current angle of 10° elec, the comparison of the exterior iron bridges flux density of motor M_1 between the AM and the FEA is given in Fig.13.

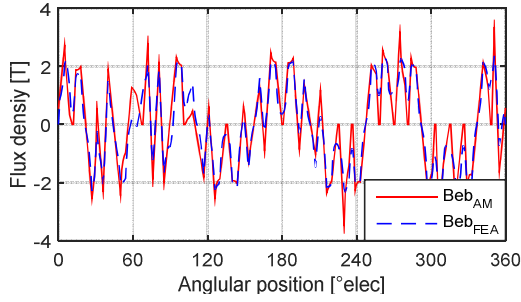


Fig.13 Comparison of the exterior iron bridges flux densities between the AM and the FEA ($I_{rms}=43A$, $\alpha=10^\circ elec$)

Fig.13 shows a good agreement between the AM and the FEA results. The flux density of the iron bridges can reach high saturation levels and causes the peaks observed in the AM. Despite these local peaks, the airgap flux density given in Fig.11, reflects the saturation of the motor and shows very good results when compared to the FEA.

E. Flux density of the rotor

The motor M_1 has 2 flux barriers therefore it has 2 flux-guides per pole. Fig.14 shows the flux density of the flux-guides for the first pole. It was computed along the mean flux path of the flux-guide, with a current of 43A (rms) and a current angle of $10^\circ elec$. The x-axis represents the length of the flux-guide.

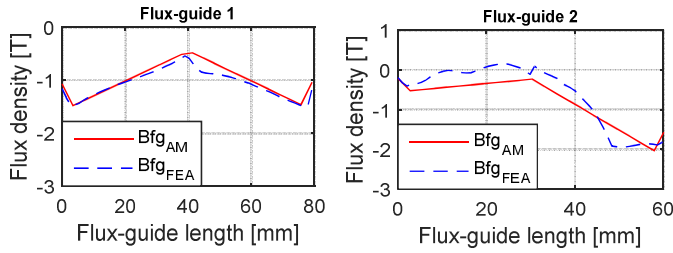


Fig.14 Comparison of the flux density of the rotor flux-guides for the 1st pole between the AM and the FEA ($I_{rms}=43A$, $\alpha=10^\circ elec$)

The analytical model shows good results when compared to the finite element model. The local variations of the flux density along the flux-guide shown in the FEA are close to the linear approximation of the analytical model.

F. Average torque

The average torque is computed for the motor M_1 at two current levels: 28A and 43A. Another motor geometry is added to the comparison. The motor M_2 has the same exterior dimensions as M_1 but differs in its rotor parameters (Fig.1). The barriers parameters of both motors are given in per unit and reported in the appendix. T_{avg} of M_2 is also computed for a current of 28 A and 43A (rms). Fig.15 shows the comparison of the AM and the FEA for the two motor structures. It can be seen that the curves are very close at all current levels for the two structures. The discrepancy between the two models at the maximum average torque is given in TABLE I and is around 2% (for the two current values) which confirms the accuracy of the analytical model.

TABLE I
COMPARISON BETWEEN AM AND FEA AT MAXIMUM AVERAGE TORQUE

	M1		M2	
Current rms (A)	28	43	28	43
AM: T_{avg} (Nm)	104.4	163.4	128	192.8
FEA: T_{avg} (Nm)	102.1	160.2	128.6	194.5
Discrepancy %	2.2	2	0.4	0.9

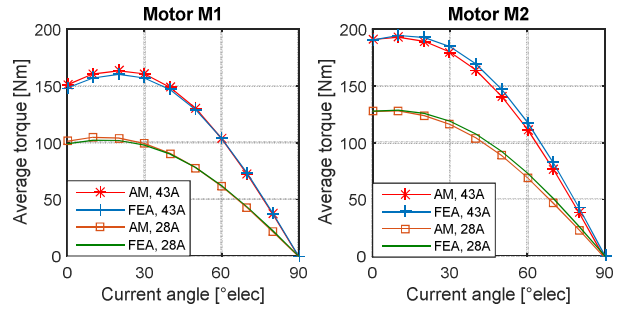


Fig.15 Comparison of the average torque between the AM and the FEA for two motor structures

G. Power factor

The power factor (PF) is also computed for the two motors under the same conditions of the torque. The analytical model curves and the finite element curves show good agreement. Fig.16 shows the PF versus the current angle.

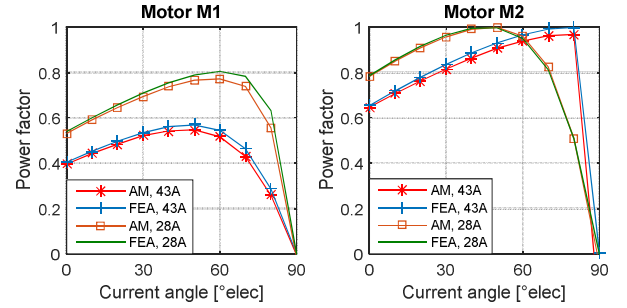


Fig.16 Comparison of the power factor between the AM and the FEA for two motor structures

At maximum average torque, for 28A the PF of M_1 is 0.59 (AM) and 0.6 (FEA). For 43A it is 0.48 (AM) and 0.49 (FEA). For M_2 at 28A the PF is 0.85 (AM) and 0.855 (FEA). For 43A it is 0.7 (AM) and 0.71 (FEA). The difference between the two models remains acceptable.

H. Voltage

The two motors rotate at 500 rpm. Fig.17 shows the results for the voltage curves computed with the analytical and finite element models. TABLE II shows the discrepancy between the two models at the current angle that corresponds to the maximum average torque. It is less than 5% which confirms the validity of the analytical model.

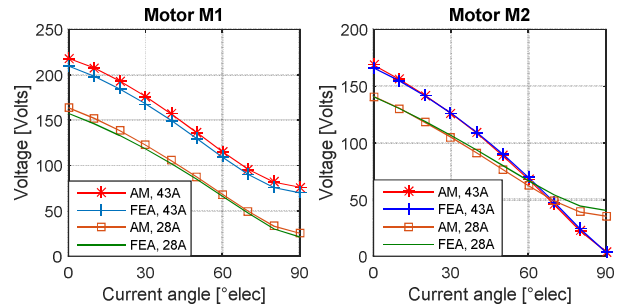


Fig.17 Comparison of the voltage between the AM and the FEA for two motor structures

TABLE II
COMPARISON OF VOLTAGE BETWEEN AM AND FEA AT MAXIMUM AVERAGE TORQUE

	M1		M2	
Current rms (A)	28	43	28	43
AM: Voltage (V)	152.1	193.2	130.5	156
FEA: Voltage (V)	146.1	184.4	130.2	154.7
Discrepancy %	4.1	4.7	0.2	0.8

I. Computational time

In order to obtain accurate values of the average torque, the PF and the voltage, 4 static computations are needed [15]. The average computational time of one calculation step is 3.1s for the AM and 17s for the FEA. This makes the analytical model about 5 times faster than the finite element model.

V. CONCLUSION

This paper presented a complete analytical model for a saturated multi-V-shape IPM with non-overlapping concentrated winding. The equations are based on Ampere's theorem and flux conservation law. The proposed model took into consideration the evolution of the flux density along the height of the stator teeth with the tangential slot leakage flux. The saturation in the rotor flux-guide is modeled and the local saturation in the area near the exterior iron bridges is included for a more accurate computation of the airgap flux density.

The analytical model was compared to a 2D finite element model. The results of the flux density in the different parts of the motor showed a very good agreement with the finite element analysis. The curves of the average torque, the power factor and the voltage for two motor structures showed also good agreement with the finite element model.

The analytical model is about 5 times faster than the finite element model and gives accurate results. Thus, in future work, it will be used for the electromagnetic calculation in a multi-physics optimization routine.

VI. APPENDIX

MOTOR SPECIFICATIONS

	M1	M2
Stator exterior radius (mm)	130	130
Axial length (mm)	200	200
Barriers width (pu)	1	1.3
Barriers tilt angle (pu)	1	0.1
Barriers opening angle (pu)	1	1.5
Specific torque (Nm/liter)	15.08	18.32

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VIII. BIOGRAPHIES

Paul Akiki received his Engineering degree and Master's degree from Ecole Supérieure d'Electricité (Supélec), France, in 2014. He is currently working towards the Ph.D degree in Electrical Engineering with the Group of electrical engineering – Paris (GeePs) at CentraleSupélec, France.

Maya Hage Hassan received her B.S. degree in mechanical engineering from the LU-Faculty of engineering, Lebanon, and her M.S. degree from the Ecole Centrale de Nantes, France, and a Ph.D. degree in electrical engineering from University Paris-Sud in 2014. She is now an Assistant Professor at CentraleSupélec and with the Group of electrical engineering – Paris, CNRS UMR 8507. Her current research interests include optimization and design of electric machines and drives.

Jean-Claude Vannier is a Professor and head of "Department of Energy and Power Systems" at CentraleSupélec, France. His primary research interest is in energy conversion systems (motors, actuators, generators) and concerns the modeling, the design, and the optimization of these equipments.

Mohamed Bensetti received his Master's degree (DEA) in 2001 and his Ph.D. in Electrical Engineering in 2004 from the University of Paris-Sud, France. He is now an Associate Professor at CentraleSupélec. He is currently with GeePs (Group of electrical engineering -Paris). His current research interests are Electromagnetic modeling, ElectroMagnetic Compatibility (EMC) and Power Electronics, including modelling, simulation and instrumentation.

Benjamin Dagusé received a degree in Electrical Engineering from École Polytechnique de l'Université de Nantes, in 2008, the M.Sc. degree in 2009 and the Ph.D degree from Supélec in 2013. He is now working with Leroy-Somer as R&D Electrical Design Manager. His main research interests include design of electrical motors with their power electronic supply.

Dany Prieto received the M.S. degree in Electrical Engineering from the University of Nantes, France, in 2011 and the Ph.D. degree from CentraleSupélec in 2015. Since 2015, he has been working with Leroy-Somer Motors as R&D Electrical Engineer. His research interests include PMSM and induction motors with their power electronic supply.

Mike McClelland obtained his engineering degree from Leeds University in 1985 and PhD in 1994. He has been working for Emerson ever since on the electromagnetic, mechanical and thermal design of brushless machines and is now employed as Technical Director for Leroy Somer Motors & Drives division. He is a Fellow of the Institute of Engineering and Technology (UK).