# DARBOUX INTEGRABLE SYSTEM WITH A TRIPLE POINT AND PSEUDO-ABELIAN INTEGRALS 

Aymen Braghtha

## - To cite this version:

Aymen Braghtha. DARBOUX INTEGRABLE SYSTEM WITH A TRIPLE POINT AND PSEUDOABELIAN INTEGRALS. 2016. hal-01388054v1

## HAL Id: hal-01388054 <br> https://hal.science/hal-01388054v1

Preprint submitted on 26 Oct 2016 (v1), last revised 2 Feb 2018 (v3)

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire HAL, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

# DARBOUX INTEGRABLE SYSTEM WITH A TRIPLE POINT AND PSEUDO-ABELIAN INTEGRALS 

AYMEN BRAGHTHA


#### Abstract

In this paper we consider the degeneracies of the third type. More exact, the perturbations of the Darboux integrable foliation with a triple point, i.e. the case where three of the curves $\left\{P_{i}=0\right\}$ meet at one point, are considered. Assuming that this is the only non-genericity, we prove that the number of zeros of the corresponding pseudo-abelian integrals is bounded uniformly for close Darboux integrable foliations. Let $\mathcal{F}$ denote the foliation with triple point (assume it to be at the origin), and let $\mathcal{F}_{\lambda}=\left\{M_{\lambda} \frac{d H_{\lambda}}{H_{\lambda}}=0\right\}, M_{\lambda}$ is a integrating factor, be the close foliation. The main problem is that $\mathcal{F}_{\lambda}$ can have a small nest of cycles which shrinks to the origin as $\lambda \rightarrow 0$. A particular case of this situation, namely $H_{\lambda}=(x-\lambda)^{\epsilon}(y-x)^{\epsilon}+(y+x)^{\epsilon}-\Delta$ with $\Delta$ non-vanishing at the origin (and generic in appropriate sense).


Mathematics subject classification: 34C07, 34C08
Keywords: Abelian integrals, Limit cycles, Integrable systems

## 1. Introduction and Main Result

This paper represent an amelioration of [Braghtha] in the sense to eleminate the technic condition on the perturbative one-form $\eta$. Precisely, we consider an unfolding $\omega_{\lambda}$ of the one-form $\omega_{0}$, where $\lambda$ is a small parameter and $\omega_{\lambda}$ is a family of meromorphic one-forms

$$
\begin{equation*}
\omega_{\lambda}=M_{\lambda} \frac{d H_{\lambda}}{H_{\lambda}}, \quad H_{\lambda}=P_{\lambda}^{\epsilon} \prod_{i=1}^{k} P_{i}^{\epsilon_{i}}, \quad M_{\lambda}=P_{\lambda} \prod_{i=1}^{k} P_{i} \tag{1}
\end{equation*}
$$

with $\epsilon, \epsilon_{i}>0, P_{0}, P_{\lambda}, P_{i} \in \mathbb{R}[x, y]$.
We assume that $P_{0}(0,0)=P_{1}(0,0)=P_{2}(0,0)=0$ and $P_{i}(0,0) \neq 0, i=3, \ldots, k$. Generically, the triple point unfolds into three saddles $p_{0}^{\lambda}, p_{1}^{\lambda}$ and $p_{2}^{\lambda}$ correspond to the transversal intersections of level curves $P_{1}^{-1}(0)$ and $P_{\lambda}^{-1}(0), P_{1}^{-1}(0)$ and $P_{2}^{-1}(0)$, and $P_{2}^{-1}(0)$ and $P_{\lambda}^{-1}(0)$. Here also appears a center $p_{c}^{\lambda}$ in the triangular region bounded by these levels curves.

Consider a polynomial perturbation $\omega_{\lambda, \kappa}=\omega_{\lambda}+\kappa \eta, \quad \kappa>0$ of the system $\omega_{\lambda}=M_{\lambda} \frac{\mathrm{d} H_{\lambda}}{H_{\lambda}}$, where

$$
\eta=R \mathrm{~d} x+S \mathrm{~d} y
$$

and $R, S \in \mathbb{R}[x, y]$ are a polynomials of degree $n$. The foliation $\omega_{\lambda}=0$ has a maximal nest of cycles $\gamma(\lambda, h) \subseteq\left\{H_{\lambda}(x, y)=h\right\}, h \in(0, n(\lambda))$ filling a connected component of $\mathbb{R}^{2} \backslash\left\{P_{\lambda} \prod_{i=1}^{k} P_{i}=\right.$ $0\}$, which we denote $D_{(\lambda, h)}$, whose boundary is a polycycle $\gamma(\lambda, 0)$.

To such perturbation one can associate the pseudo-Abelian integral

$$
\begin{equation*}
I(\lambda, h)=\int_{\gamma(\lambda, h)} \frac{\eta}{M_{\lambda}} \tag{2}
\end{equation*}
$$

which is the principal part of the Poincaré displacement function

$$
D(\kappa, \lambda, h)=\kappa h \int_{\gamma(\lambda, h)} \frac{\eta}{M_{\lambda}}+O(\kappa)
$$

of the perturbation $\omega_{\lambda, \kappa}$ along $\gamma(\lambda, h)$.
Let us impose the following generic assumptions:
$\mathbf{A}_{1}:\left.\frac{\partial P_{\lambda}}{\partial \lambda}\right|_{(0,0,0)} \neq 0$.
$\mathbf{A}_{2}: P_{1}^{-1}(0), P_{2}^{-1}(0)$ and $P_{0}^{-1}(0)$ intersect transversally two by two at the origin which is the only triple point. The level curves $P_{i}^{-1}(0), i=3, \cdots, k$ intersect transversally and two by two.

Theorem 1. Let $I(\lambda, h)$ be the family of pseudo-Abelian integrals as defined above. Under assumptions $\boldsymbol{A}_{1}$ and $\boldsymbol{A}_{2}$, there exists a bound for the number of isolated zeros of $I(\lambda, h)=$ $\int_{\gamma(\lambda, h)} \frac{\eta}{M_{\lambda}}$. The bound depends only on $n_{i}=\operatorname{deg} P_{i}, n=\max (\operatorname{deg} R, \operatorname{deg} S)$ and is uniform in the coefficients of the polynomials $P_{\lambda}, P_{i}, R, S$, the exponents $\epsilon, \epsilon_{i}, i=1, \ldots, k$ and the parameter $\lambda$.
1.1. Bobieński result. We recall also a similar result of Bobiénski [1] which he prove the existence of a local upper bound for the number of zeros of pseudo-abelian integrals. He consider a one-parameter unfolding of the singular codimension one case. The difference relies in the fact that in this work the Darboux first integrals is more general $\left(\epsilon_{1} \neq \epsilon_{2}\right)$ and the proof of our main result is purely geometric: we use the blow-up in families. This approach gives directly uniform validity of our study of the pseudo-abelian integrals.

## 2. Darboux integrable foliation

Let us a stablish a local normal form near the triple point $(0,0,0)$ for the unfolding of the degenerate polycyle $H_{0}$.

Proposition 1. Under above assumptions $\boldsymbol{A}_{1}, \boldsymbol{A}_{2}$. There exists a local analytic coordinate system $(x, y, \lambda)$ at $(0,0,0)$ such that $H_{\lambda}$ takes the form

$$
\begin{equation*}
H_{\lambda}=(x-\lambda)^{\epsilon}(y-x)^{\epsilon_{+}}(y+x)^{\epsilon-} \Delta, \quad \lambda>0 \tag{3}
\end{equation*}
$$

where $\Delta$ is an analytic unity function $\Delta(0,0,0) \neq 0$.

Let $\mathcal{F}_{1}, \mathcal{F}_{2}$ are two foliations of dimension two in space of the total complex space $\mathbb{C}^{3}$ with coordinates $(x, y, \lambda)$ such that

$$
\mathcal{F}_{1}:\left\{H(x, y, \lambda)=P_{\lambda}^{\epsilon} \prod_{i=1}^{k} P_{i}^{\epsilon_{i}}=h\right\}, \quad \mathcal{F}_{2}:\{\lambda=\text { constant }\} .
$$

Consider the Darboux foliation $\mathcal{F}:=\left\{\mathcal{F}_{1}, \mathcal{F}_{2}\right\}$ of dimension one in $\mathbb{C}^{3}$ with coordinates $(x, y, \lambda)$ which is given by the intersection of the leaves of $\mathcal{F}_{1}$ and $\mathcal{F}_{2}$. This foliation has a non-elementary singular point at the origin $(0,0,0)$. This foliation has a complicated singularity at the origin $(0,0,0)$. To reduice this singularity , we perform a directional blowing-up in the family $\omega_{\lambda}$. Note that blow-up in a family was introduced in [5], see also [6].
2.1. Directional Blow-up. The blow-up of $\mathbb{C}^{3}$ at the origin is defined as the incidence three dimensional manifold $W=\left\{(p, q) \in \mathbb{C P}^{2} \times \mathbb{C}^{3}: q \in p\right\}$. The blow down $\sigma: W \rightarrow \mathbb{C}^{3}$ is just the restriction to $W$ of the projection $\mathbb{C P}^{2} \times \mathbb{C}^{3}$. The inverse map $\sigma^{-1}: \mathbb{C}^{3} \rightarrow W$ is called blow-up and $\sigma^{-1}(0)=\mathbb{C P}^{2}$ is called exceptional divisor. The projective space $\mathbb{C P}^{2}$ is covered by three canonical charts: $W_{1}=\{x \neq 0\}$ with coordinates $\left(v_{1}, w_{1}\right), W_{2}=\{y \neq 0\}$ with coordinates $\left(u_{2}, w_{2}\right)$ and $W_{3}=\{\lambda \neq 0\}$ with coordinates $\left(u_{3}, v_{3}\right) . W_{1}, W_{2}$ and $W_{3}$ define canonical charts on $W$, with coordinates $\left(u_{1}, v_{1}, w_{1}\right)$,
$\left(u_{2}, v_{2}, w_{2}\right)$ and $\left(u_{3}, v_{3}, w_{3}\right)$ respectively. The blow-up $\sigma^{-1}$ is written as:

$$
\begin{array}{ccr}
\sigma_{1}^{-1}=\left.\sigma^{-1}\right|_{W_{1}}: x=u_{1} & y=u_{1} w_{1} & \lambda=w_{1} u_{1} \\
\sigma_{2}^{-1}=\left.\sigma^{-1}\right|_{W_{2}}: x=u_{2} v_{2} & y=v_{2} & \lambda=w_{2} v_{2} \\
\sigma_{3}^{-1}=\left.\sigma^{-1}\right|_{W_{3}}: x=u_{3} w_{3} & y=v_{3} w_{3} & \lambda=w_{3} \tag{6}
\end{array}
$$

Let $\sigma^{-1} \mathcal{F}$ be the lift of the foliation $\mathcal{F}$ to the complement of the exceptional divisor $\mathbb{C P}^{2}$. This foliation is regular outside of the preimage of the hypersurface $\left\{H_{\lambda}=0, \lambda=0\right\}$.

Prposition 2. Let $a=\epsilon+\epsilon_{-}+\epsilon_{+}$
(1) The foliation $\sigma^{-1} \mathcal{F}$ extends in a unique way to a holomorphic singular foliation $\sigma^{*} \mathcal{F}$ on $W$ which we call the blow-up of the original codimension two foliation $\mathcal{F}$ by the map $\sigma$.
(2) Let $\sigma_{1}^{*} \mathcal{F}$ be the restriction of the blown-up foliation $\sigma^{*} \mathcal{F}$ to the chart $W_{1}$. The singularities of $\sigma_{1}^{*} \mathcal{F}$ are located at the points $p_{+}=(0,1,0), p_{-}=(0,-1,0), q_{+}=(0,1,1)$ and $q_{-}=(0,-1,1)$. All these singular points are linearisable saddles, with eigenvalues $\mu_{+}=\left(\epsilon_{+},-a,-\epsilon_{-}\right), \mu_{-}=\left(-\epsilon_{-}, a, \epsilon_{-}\right), \nu_{+}=\left(0,-\epsilon, \epsilon_{+}\right)$and $\nu_{-}=\left(0,-\epsilon, \epsilon_{-}\right)$respectively.

Proof. We prove the second item of Proposition. Since $\sigma: W \rightarrow \mathbb{C}^{3}$ is a biholomorphism outside $\mathbb{C P}^{2}$, all singularities of $\sigma_{1}^{*} \widetilde{\mathcal{F}}$ on $W_{1} \backslash\left\{u_{1}=0\right\}$ correspond to singularities of $\mathcal{F}$.

Thus, it suffices to compute the singularities of $\sigma_{1}^{*} \mathcal{F}$ on the exceptional divisor $\left\{u_{1}=0\right\}$. On the exceptional divisor, the foliation is given by the levels of

$$
G:=\frac{\left(\pi \circ \sigma_{1}\right)^{a}}{H \circ \sigma_{1}}=w_{1}^{a}\left(1-w_{1}\right)^{-\epsilon}\left(v_{1}-1\right)^{-\epsilon_{+}}\left(v_{1}+1\right)^{-\epsilon_{-}} \tilde{\Delta}^{-1}
$$

where $\tilde{\Delta}$ is unit of the form $\tilde{\Delta}=c+u_{1} f, f$ is a holomorphic function. Let us compute the eigenvalues at $p_{+}, p_{-}, q_{+}$and $q_{-}$. Near the exceptional divisor $\left\{u_{1}=0\right\}$, the foliation $\sigma_{1}^{*} \mathcal{F}$ is given by

$$
\begin{aligned}
& H \circ \sigma_{1}\left(u_{1}, v_{1}, w_{1}\right)=u_{1}^{a}\left(1-w_{1}\right)^{\epsilon}\left(v_{1}-1\right)^{\epsilon+}\left(v_{1}+1\right)^{\epsilon-} \tilde{\Delta}=h, \\
& \pi \circ \sigma_{1}\left(u_{1}, v_{1}, w_{1}\right)=u_{1} w_{1}=\lambda
\end{aligned}
$$

Near $p_{+}$and after the respective changes of variable $v=\left(v_{1}-1\right)\left(v_{1}+1\right)^{\frac{\epsilon_{-}}{\epsilon_{+}}} \Delta$, the blown-up foliation $\sigma_{1}^{*} \mathcal{F}$ is given by the two first integrals $u_{1}^{a} v_{+}^{\epsilon_{+}}=h$ and $u_{1} w_{1}=s$. Then, near this point the vector field generating the foliation $\sigma_{1}^{*} \mathcal{F}$ is given by

$$
X\left(u_{1}, v, w_{1}\right)=\mu_{1}^{+} u_{1} \frac{\partial}{\partial u_{1}}+\mu_{2}^{+} v \frac{\partial}{\partial v}+\mu_{3}^{+} w_{1} \frac{\partial}{\partial w_{1}}
$$

where the vector $\mu_{ \pm}=\left(\mu_{1}^{ \pm}, \mu_{2}^{ \pm}, \mu_{3}^{ \pm}\right)$satisfies the following equations

$$
<\left(\mu_{1}^{ \pm}, \mu_{2}^{ \pm}, \mu_{3}^{ \pm}\right),\left(a, \epsilon_{ \pm}, 0\right)>=0, \quad<\left(\mu_{1}^{ \pm}, \mu_{2}^{ \pm}, \mu_{3}^{ \pm}\right),(1,0,1)>=0
$$

Here $<,>$ is the usual scalar product on $\mathbb{C}^{3}$. By simple calculations, we obtain

$$
X_{ \pm}\left(u_{1}, v_{ \pm}, w_{1}\right)= \pm \epsilon_{ \pm} u_{1} \frac{\partial}{\partial u_{1}} \mp a v_{ \pm} \frac{\partial}{\partial v_{ \pm}} \mp \epsilon_{ \pm} w_{1} \frac{\partial}{\partial w_{1}}
$$

Similar computation shows that there are local coordinates near $q_{ \pm}$in which the vector field generating the foliation is given by

$$
X_{ \pm}\left(u_{1}, v_{ \pm}, w_{ \pm}\right)=-\epsilon w_{ \pm} \frac{\partial}{v_{ \pm}}+\epsilon_{ \pm} w_{ \pm} \frac{\partial}{\partial w_{ \pm}}
$$

## 3. Proof of the theorem

Let $t=\frac{\lambda^{a}}{h}$. The blown-up foliation $\sigma_{1}^{*} \mathcal{F}$ is given by the two first integrals

$$
\begin{aligned}
& G=w_{1}^{a}\left(1-w_{1}\right)^{-\epsilon}\left(v_{1}-1\right)^{-\epsilon_{+}}\left(v_{1}+1\right)^{-\epsilon_{-}} \tilde{\Delta}^{-1}=t \\
& F=u_{1} w_{1}=\lambda
\end{aligned}
$$

Consider the two-dimensional square $Q \subset \mathbb{C P}^{2}$ with vertices $p_{+}, p_{-}, q_{+}$and $q_{-}$. All levels curves $\{G=t\}$ inside $Q$ correspond to values of $t \in[0,+\infty]$. We consider the family of hyperbolic polycycles $\delta^{t}$, i.e. at each intersection of two consecutive curve we have a saddle point,.

$$
\delta^{t}=\left(\sigma_{1}^{-1}(\gamma(0,0) \backslash(0,0,0)) \cup(Q \cap\{G=t\})\right)^{\mathbb{R}}, t \in[0,+\infty]
$$

where $(\ldots)^{\mathbb{R}}$ denotes the real part of a complex analytic set. All polycycle $\delta^{t}$ satisfies the generecity assumptions from [2].

Let $\delta(\lambda, t)=\sigma^{-1}(\gamma(\lambda, h)) \subset W$ be the pull-back of the cycle $\gamma(\lambda, h)$ by the blowing-up map. Let $\delta^{t}$ be the polycycle corresponding to the cycle $\delta(\lambda, t)$.

Let

$$
J(\lambda, t)=\int_{\delta(\lambda, t)} \sigma_{1}^{*} \frac{\eta}{M_{\lambda}}
$$

The integral $J(\lambda, t)$ can be analytically continued to the universal cover of $\mathbb{C}^{2} \backslash\{\lambda t=0\}$.
3.1. Variation operator. Given any multivalued function $F$ defined in a neighborhood of the origin in $\mathbb{C}$ i.e. a holomorphic function defined on the universal covering $\widetilde{\mathbb{C}^{*}}$ of $\mathbb{C}^{*}$. We define the rescaled monodromy as

$$
\mathcal{M o n}_{(t, \alpha)} F(t)=F\left(t e^{i \pi \alpha}\right)
$$

The variation is given as the difference between the counterclockwise and clockwise continuation

$$
\begin{aligned}
\mathcal{V a r}_{(t, \alpha} F(t) & =\mathcal{M o n}_{(t, \alpha)} J(t)-\mathcal{M o n}_{(t,-\alpha)} J(t) \\
& =F\left(t e^{i \pi \alpha}\right)-F\left(t e^{-i \pi \alpha}\right) .
\end{aligned}
$$

Now, let $G$ be a multivalued function in two variables $\lambda$ and $t$ defined in universal covering $\mathbb{C}^{2} \backslash\{s t=0\}$ of $\mathbb{C}^{2} \backslash\{s t=0\}$. We define the mixed variation as

$$
\begin{aligned}
& \mathcal{V a r}_{(\beta, \lambda)} \circ \mathcal{V a r}_{\alpha, t} G(\lambda, t)=\operatorname{Var}_{(\beta, \lambda)}\left(G\left(\lambda, t e^{i \pi \alpha}\right)-G\left(\lambda, t e^{-i \pi \alpha}\right)\right)= \\
& G\left(\lambda e^{i \pi \beta}, t e^{i \pi \alpha}\right)-G\left(\lambda e^{-i \pi \beta}, t e^{i \pi \alpha}\right)-G\left(\lambda e^{i \pi \beta}, t e^{-i \pi \alpha}\right)+G\left(\lambda e^{-i \pi \beta}, t e^{-i \pi \alpha}\right) .
\end{aligned}
$$

Lemma 1 The variations $\mathcal{V}_{(\lambda, \beta)}$ and $\mathcal{V}_{(t, \alpha)}$ commute

$$
\mathcal{V}_{(\lambda, \beta)} \circ \mathcal{V}^{2} r_{(t, \alpha)}=\mathcal{V}^{2} r_{(t, \alpha)} \circ \mathcal{V}^{2} r_{(\lambda, \beta)}
$$

Proof. The proof is a direct consequence of monodromy theorem.
3.2. The analytic proporties of the integral $J$ near the polycycle $\delta^{t}$. Using the partition of unity of the blown-up space we can decompose our cycle of integration $\delta(\lambda, t)$ in a relative cycles and we can check that any relative cycle can be chosen as a lift of a base path (for more details see [3]). Then, we have

Proposition 3 . The integral $J(\lambda, t)$ satisfies the rescaled iterated variation equations

$$
\begin{align*}
& \mathcal{V a r}_{\left(t, \beta_{1}\right)} \circ \ldots \circ \operatorname{Var}_{\left(t, \beta_{k}\right)} J(\lambda, t)=0,  \tag{7}\\
& \operatorname{Var}_{(\lambda, 1)} \circ \operatorname{Var}_{(\lambda, 1)} J(\lambda, t)=0, \tag{8}
\end{align*}
$$

where $\beta_{i}$ are analytic functions in $\epsilon, \epsilon_{+}, \epsilon_{-}, \ldots, \epsilon_{k}$.
3.3. Increment of argument of $J$. The estimation of a local bound of the number of isolated zeros of the pseudo-abelian integral $I(\lambda, h)=\int_{\gamma(\lambda, h)} \frac{\eta}{M_{\lambda}}$ is analoguous to estimate a local $\lambda$ uniform bound of isolated zeros of the integral $J(\lambda, t)=\int_{\delta(\lambda, t)} \sigma_{1}^{*} \frac{\eta}{M_{\lambda}}$. As consequence of the equations (7) and (8) the integral $J(s, t)$ has the following expansion

$$
\begin{equation*}
J(\lambda, t)=J_{1}(\lambda, t)+J_{2}(\lambda, t) \log \lambda \tag{9}
\end{equation*}
$$

where

$$
\begin{align*}
& J_{2}(\lambda, t)=\mathcal{V a r}_{(\lambda, 1)} J(\lambda, t)=\int_{\text {eight loop }} \sigma_{1}^{*} \frac{\eta}{M_{\lambda}},  \tag{10}\\
& \mathcal{V a r}_{\left(t, \beta_{1}\right)} \circ \ldots \circ \mathcal{V a r}_{\left(t, \beta_{k}\right)} J_{i}(\lambda, t)=0, \quad i=1,2,  \tag{11}\\
& \operatorname{Var}_{(\lambda, 1)} J_{i}(\lambda, t)=0, i=1,2 . \tag{12}
\end{align*}
$$

Theorem 2. Let $\epsilon>0$ be sufficiently small. Then, for all $|\lambda|<\epsilon$ the number of zeros

$$
\#\{t \in[0,+\infty]: J(\lambda, t)=0\}
$$

is uniformly bounded with respect to $\lambda$.
The integral $J(\lambda, t)$ has analytic prolongation to the complex argument $t$. This is a multivalued function with unique ramification point $t=0$.

Let $\partial \Omega$ be the boundary of a domain complex $\Omega$ which consists of a big circular arc $C_{R_{1}}=$ $\left\{|t|=R_{1},|\arg t| \leq \alpha \pi\right\}$, a two segments $C^{ \pm}=\left\{r_{1} \leq|t| \leq R_{1},|\arg t|= \pm \alpha \pi\right\}$ and the small circular $\operatorname{arc} C_{r_{1}}=\left\{|t|=r_{1},|\arg t| \leq \alpha \pi\right\}$.

To count the number of zeros of the function $J(s, t)$, we estimate the increment of argument of the function $J(\lambda, t)$ we apply the argument principle which says that

$$
\# Z\left(J_{\Omega}\right) \leq \frac{1}{2 \pi} \Delta \arg \left(J_{\partial \Omega}\right)=\frac{1}{2 \pi}\left(\Delta \arg \left(J_{C_{R_{1}}}\right)+\Delta \arg \left(J_{C_{r_{1}}}\right)+\Delta \arg \left(J_{C^{ \pm}}\right)\right)
$$

The increment of arguments $\Delta \arg \left(J_{C_{R_{1}}}\right)$ and $\Delta \arg \left(J_{C^{ \pm}}\right)$are locally uniformly bounded, for more details see [3]. The problem consist to estimate the increment of argument of the integral $J(s, t)$ along the small circular arc $C_{r_{1}}$.

Let

$$
\begin{aligned}
\Lambda\left(\lambda ; n_{0}, n_{1}, \cdots, n_{k} ; n\right)= & \left\{M_{\lambda} \frac{d H_{\lambda}}{H_{\lambda}}+\kappa(R d x+R d y): H_{\lambda}=P_{\lambda} \prod_{i=1}^{k} P_{i}^{\epsilon_{i}}\right. \\
& \left.\operatorname{deg} P_{\lambda} \leq n_{0}, \operatorname{deg} P_{i} \leq n_{i}, \operatorname{deg}(R, S) \leq n\right\}
\end{aligned}
$$

be the parameters space. Consider the following functional space $\mathcal{P}$

$$
\mathcal{P}\left(v, V ; \alpha_{1}, \ldots, \alpha_{k} ; \lambda\right):=\left\{\sum \sum c_{j l n}(\lambda) t^{\alpha_{j} n} \log ^{n} t, c_{j l n} \in \mathbb{C}, v \leq \alpha_{j} n \leq V, 0 \leq l \leq k\right\}
$$

3.3.1. Finite order approximation to $J(\lambda, t)$. Our goal is to obtain an asymptotic expansion for $J(\lambda, t)$. This will allow us to prove the existence of an upper bound for the increment of the $\operatorname{argument} \Delta \arg \left(J_{C_{r_{1}}}\right)$. The problem is proposed in the fact that the existence of the term $\log (\lambda) \rightarrow \infty$ as $\lambda \rightarrow 0$ in the expression of the function $J(\lambda, t)$. To resolve the problem we make a weighted quasi-homogeneous blowing-up. Let

$$
c=(\log \lambda)^{-1}, \quad a=J_{1}(\lambda, t), \quad b=J_{2}(\lambda, t)
$$

Then, we have

$$
J(\lambda, t)=J_{1}(\lambda, t)+\log (\lambda) J_{2}(\lambda, t)=c^{-1}(c a+b)=c^{-1} \psi(a, b, c)
$$

As $c^{-1}=\log \lambda \in \mathbb{R}$, we have

$$
\arg (J(\lambda, t))=\arg \left(c^{-1} \psi(a, b, c)\right)=\arg (\psi(a, b, c))
$$

Proposition 4. There exist a some meromorphic function $\phi$ such that

$$
\Delta \arg (\psi(a, b, c))=\Delta \arg (\phi)
$$

Proof. To estimate the increment of $\operatorname{argument} \arg (q(a, b, c))$ of polynomial $q$ we make a quasihomogeneous blowing-up of weight $\left(\frac{1}{2}, 1, \frac{1}{2}\right)$ with the center $C_{1}=\{a=b=c=0\}$. The explicit formulae of the blowing-up $\pi_{1}^{-1}$ in the affine charts $\tau_{1}=\{a \neq 0\}, \tau_{2}=\{b \neq 0\}$ and $\tau_{3}=\{c \neq 0\}$ is written respectively as

$$
\begin{aligned}
& \left.\pi_{1}^{-1}\right|_{\tau_{1}}=\pi_{11}^{-1}: a=\sqrt{a_{1}}, \quad b=b_{1} a_{1}, \quad c=c_{1} \sqrt{a_{1}}, \\
& \left.\pi_{1}^{-1}\right|_{\tau_{2}}=\pi_{12}^{-1}: a=a_{2} \sqrt{b_{2}}, \quad b=b_{2}, \quad c=c_{2} \sqrt{b_{2}}, \\
& \left.\pi_{1}^{-1}\right|_{\tau_{3}}=\pi_{13}^{-1}: a=a_{3} \sqrt{c_{3}}, \quad b=b_{3} c_{3}, \quad c=\sqrt{c_{3}} .
\end{aligned}
$$

The total transform $\pi_{1}^{*} q(a, b, c)$ of $q(a, b, c)$ is given, in different charts, by

$$
\begin{aligned}
& \pi_{11}^{*}(c a+b)=a_{1}\left(c_{1}+b_{1}\right)=a_{1} P_{1}\left(a_{1}, b_{1}, c_{1}\right) \\
& \pi_{12}^{*}(c a+b)=b_{2}\left(a_{2} c_{2}+1\right)=b_{2} P_{2}\left(a_{2}, b_{2}, c_{2}\right) \\
& \pi_{13}^{*}(c a+b)=c_{3}\left(a_{3}+b_{3}\right)=c_{3} P_{3}\left(a_{3}, b_{3}, c_{3}\right)
\end{aligned}
$$

where $\left\{a_{1}=0\right\},\left\{b_{2}=0\right\}$ and $\left\{c_{3}=0\right\}$ are local equations of the exceptional divisor and $\left\{P_{1}=0\right\},\left\{P_{2}=0\right\}$ and $\left\{P_{3}=0\right\}$ local equations of strict transform of $\{c a+b=0\}$.

We observe that the exceptional divisor $\left\{a_{1}=0\right\}$ ( resp $\left\{c_{3}=0\right\}$ ) has not a normal crossing with the strict transform $\left\{P_{1}=0\right\}$ (resp $\left\{P_{3}=0\right\}$ ). Conversely in the chart $\tau_{2}$, the exceptional divisor $\left\{b_{2}=0\right\}$ has a normal crossing with the strict transform $\left\{P_{2}=0\right\}$. To resolve this problem we make a second blowing-up $\pi_{2}^{-1}$ with center $C_{2}$ such that
(1) in the chart $\tau_{1}$ with coordinates $\left(a_{1}, b_{1}, c_{1}\right)$, the center $C_{2}$ is given by $C_{2}=\left\{c_{1}=b_{1}=0\right\}$,
(2) in the chart $\tau_{2}$ the blowing-up $\pi_{2}$ is a biholomorphism,
(3) in the chart $\tau_{3}$ with coordinates $\left(a_{3}, b_{3}, c_{3}\right)$, the center $C_{2}$ is given by $C_{2}=\left\{a_{3}=b_{3}=0\right\}$, and
(1) in the chart $\tau_{1}$, we have $\pi_{2}^{-1}\left(C_{2}\right)$ is covred by two coordinates charts $V_{b_{1}}$ and $V_{c_{1}}$ with coordinate $\left(\tilde{a}_{1}, \tilde{b}_{1}, \tilde{c}_{1}\right)$ such that the blowing-up $\pi_{2}^{-1}$ is given in the charts $V_{b_{1}}$ and $V_{c_{1}}$ respectively, by $a_{1}=\tilde{a}_{1}, b_{1}=\tilde{b}_{1}, c_{1}=\tilde{b}_{1} \tilde{c}_{1}$ and $a_{1}=\tilde{a}_{1}, b_{1}=\tilde{b}_{1} \tilde{c}_{1}, c_{1}=\tilde{c}_{1}$,
(2) in the chart $\tau_{3}$, we have $\pi_{2}^{-1}\left(C_{2}\right)$ is covred by two coordinates charts $V_{a_{3}}$ and $V_{b_{3}}$ with coordinates $\left(\tilde{a}_{3}, \tilde{b}_{3}, \tilde{c}_{3}\right)$ such that the blowing-up $\pi_{2}^{-1}$ is given in $V_{a_{3}}$ and $V_{b_{3}}$ respectively, by $a_{3}=\tilde{a}_{3}, b_{3}=\tilde{a}_{3} \tilde{b}_{3}, c_{3}=\tilde{c}_{3}$ and $a_{3}=\tilde{a}_{3} \tilde{b}_{3}, b_{3}=\tilde{a}_{3}, c_{3}=\tilde{c}_{3}$.
Then, we have
(1) in the chart $V_{b_{1}}$,

$$
\pi_{2}^{*} \circ \pi_{1}^{*} \psi(a, b, c) \stackrel{0}{\approx} \tilde{a}_{1} \tilde{b}_{1}=J_{2}(\lambda, t),
$$

(2) in the chart $V_{c_{1}}$,

$$
\pi_{2}^{*} \circ \pi_{1}^{*} \psi(a, b, c) \stackrel{0}{\approx} \tilde{a}_{1} \tilde{c}_{1}=\frac{J_{1}(\lambda, t)}{\log \lambda}
$$

(3) in the chart $V_{a_{3}}$,

$$
\pi_{2}^{*} \circ \pi_{1}^{*} \psi(a, b, c) \stackrel{0}{\approx} \tilde{a}_{3} \tilde{c}_{3}=\frac{J_{1}(\lambda, t)}{\log \lambda}
$$

(4) in the chart $V_{b_{3}}$,

$$
\pi_{2}^{*} \circ \pi_{1}^{*} \psi(a, b, c) \stackrel{0}{\approx} \tilde{b}_{3} \tilde{c}_{3}=\frac{J_{1}(\lambda, t)}{\log \lambda}
$$

(5) in the chart $\tau_{2}$, the exceptinal divisor $\left\{b_{2}=0\right\}$ has a normal crossings with the strict transform $P_{2}=0$ and the blowing-up $\pi_{2}^{-1}$ is a biholomophism, so the function

$$
\frac{J_{1}(\lambda, t)}{\log \lambda}+J_{2}(\lambda, t)=\frac{J(\lambda, t)}{\log \lambda}
$$

is meromorphic function.
Then, we conclude that $\phi \in\left\{\frac{J(\lambda, t)}{\log \lambda}, \frac{J_{1}(\lambda, t)}{\log \lambda}, J_{2}(\lambda, t)\right\}$.
3.3.2. Proof of the theorem 2. Using lemma 4.8 of [2] for $\phi \in\left\{\frac{J_{1}(\lambda, t)}{\log \lambda}, J_{2}(\lambda, t)\right\}$, there exist a $\tilde{\phi} \in \mathcal{P}(\ldots)$ such that $|\phi(\lambda, t)-\tilde{\phi}(\lambda, t)| \leq N$. For $\phi=\frac{J(\lambda, t)}{\log \lambda}$, we can conclude using Gabrielov's theorem [4].

## References

[1] Bobieński, Marcin Pseudo-Abelian integrals along Darboux cycles a codimension one case. J. Differential Equations 246 (2009), no. 3, 1264-1273.
[2] Bobieński, Marcin; Mardešić, Pavao Pseudo-Abelian integrals along Darboux cycles. Proc. Lond. Math. Soc. (3) 97 (2008), no. 3, 669-688.
[3] Braghtha Aymen (2013) Les zéros des intégrales pseudo-abeliennes: cas non générique. PhD thesis. Université de Bourgogne: France.
[4] Gabrièlov, A. M Projections of semianalytic sets. (Russian) Funktsional. Anal. i Priložen. 2 1968 no. 4, 18-30

## Burgundy University, Burgundy Institue of Mathematics, <br> U.M.R. 5584 du C.N.R.S., B.P. 47870, 21078 Dijon

Cedex - France.
E-mail adress: aymenbraghtha@yahoo.fr

