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Hadamard property of the *in* and *out* states for Klein-Gordon fields on asymptotically static spacetimes

Christian GÉRARD & Michał WROCHNA

ABSTRACT. We consider the massive Klein-Gordon equation on a class of asymptotically static spacetimes (in the long range sense) with Cauchy surface of bounded geometry. We prove the existence and Hadamard property of the *in* and *out* states constructed by scattering theory methods.

1. INTRODUCTION & SUMMARY

1.1. Hadamard property of *in/out* states. The construction of quantum states from scattering data is a subject that has been studied extensively in various contexts in Quantum Field Theory, including the case of the wave and Klein-Gordon equation — set either on Minkowski space, in external electromagnetic potentials [Is, Lu, Ru, Se], or on curved spacetimes with special asymptotic symmetries, to mention only the works [Wa1, DK1, DK2, DK3, Mo1]. On the physics side, the primary motivation is to give meaning to the notion of particles and anti-particles and to describe quantum scattering phenomena. From the mathematical point of view, the problems often discussed in this context in the literature involve existence of scattering and Møller operators, the question of asymptotic completeness, as well as specific properties of states such as the ground state or thermal condition with respect to an asymptotic dynamics, see e.g. [Dr, DD, DRS, GGH, I-R, Ni] for various recent developments on curved backgrounds.

In the present paper we address the question of whether the so-called *in* and *out* states on asymptotically static spacetimes satisfy the *Hadamard condition* [KW]. Nowadays regarded as an indispensable ingredient in the perturbative construction of interacting fields (see e.g. recent reviews [HW3, KM, FV2]), this property accounts for the correct short-distance behaviour of expectation values of fields. It can be conveniently formulated as a condition on the *wave front set* of the state's two-point functions [Ra] — a terminology that we explain in the paragraphs below. It is known that in the special case of the conformal wave equation, one can study the wave front set of the two-point functions quite directly in the geometrical setup of conformal scattering on asymptotically flat spacetimes [Mo2, GW3] (cf. [DMP1, DMP2, BJ] for generalizations on the allowed classes of spacetimes). Furthermore, propagation estimates in b-Sobolev spaces of variable order were used recently to show a similar result in the case of the wave equation on asymptotically Minkowski spacetimes [VW], drawing on earlier developments by Vasy et al. [BVW, HV, Va1, Va2]. The two methods being however currently limited to a special value of the mass parameter, our focus here is instead on the proof of the Hadamard property of the *in* and *out* state for the Klein-Gordon operator $P = -\square_g + m^2$ for any positive mass m , or more generally for $P = -\square_g + V$ with a real-valued potential $V \in C^\infty(M)$ satisfying an asymptotic positivity condition.

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Specifically, we first consider the special case of a $1 + d$ -dimensional globally hyperbolic spacetime (M, g) with Cauchy surface Σ and metric of the form $g = -dt^2 + h_t$, with h_t a Riemannian metric smoothly depending on t . The Klein-Gordon operator can be written in the form

$$(1.1) \quad P = \partial_t^2 + r(t)\partial_t + a(t, x, \partial_x),$$

where $r(t)$ is the multiplication operator $|h_t|^{-\frac{1}{2}}\partial_t|h_t|^{\frac{1}{2}}$ and $a(t, x, D_x) \in \text{Diff}^2(\Sigma)$ has principal symbol $k \cdot h_t^{-1}(x)k$ (where $\xi = (\tau, k)$ is the dual variable of $x = (t, x)$). Now, supposing Σ is a *manifold of bounded geometry* (see Subsect. 3.1), there exist uniform pseudodifferential operator classes $\Psi^m(\Sigma)$ due to Kordyukov and Shubin [Ko, Sh1] that generalize the well-known pseudodifferential calculus of Hörmander on \mathbb{R}^d and closed manifolds. Here in addition, in order to control decay in time, we introduce t -dependent pseudodifferential operators $\Psi_{\text{td}}^{m, \delta}(\mathbb{R}; \Sigma)$ as quantizations of t -dependent symbols $a(t, x, k)$ that satisfy

$$|\partial_t^\alpha \partial_x^\beta \partial_k^\gamma a(t, x, k)| \leq C_{\alpha\beta\gamma} \langle t \rangle^{\delta - \alpha} \langle k \rangle^{m - |\gamma|}, \quad \alpha \in \mathbb{N}, \beta, \gamma \in \mathbb{N}^d,$$

where $\langle t \rangle = (1 + t^2)^{\frac{1}{2}}$, $\langle k \rangle = (1 + |k|^2)^{\frac{1}{2}}$, and the constants $C_{\alpha\beta\gamma}$ are uniform in an appropriate sense. This allows us to state a hypothesis that accounts for asymptotic ultra-staticity of (M, g) at future and past infinity. Namely, we assume that there exists $\delta > 0$ and $a_{\text{out}}, a_{\text{in}} \in \Psi^2(\Sigma)$ elliptic and bounded from below (by a positive constant), such that on $\mathbb{R}^\pm \times \Sigma$,

$$(td) \quad \begin{aligned} a(t, x, D_x) - a_{\text{out/in}}(x, D_x) &\in \Psi_{\text{td}}^{2, -\delta}(\mathbb{R}; \Sigma), \\ r(t) &\in \Psi_{\text{td}}^{0, -1-\delta}(\mathbb{R}; \Sigma). \end{aligned}$$

In practice, in our main cases of interest $a_{\text{out/in}}(x, D_x)$ will simply be the Laplace-Beltrami operator of some asymptotic metric $h_{\text{out/in}}$ plus the mass or potential term.

Let now $\mathcal{U}(t, s)$ be the Cauchy evolution of P , i.e. the operator that maps Cauchy data of P at time s to Cauchy data at time t . In this setup, what we call *time- t covariances of the out state* are the pair of operators defined by

$$(1.2) \quad c_{\text{out}}^\pm(t) := \lim_{t_+ \rightarrow \infty} \mathcal{U}(t, t_+) c_{\text{out}}^{\pm, \text{vac}} \mathcal{U}(t_+, t)$$

whenever the limit exists (in a sense made precise later on), where $c_{\text{out}}^{\pm, \text{vac}}$ equals

$$c_{\text{out}}^{\pm, \text{vac}} = \frac{1}{2} \begin{pmatrix} \mathbf{1} & \pm a_{\text{out}}^{\frac{1}{2}} \\ \pm a_{\text{out}}^{-\frac{1}{2}} & \mathbf{1} \end{pmatrix}.$$

To elucidate the interpretation of $c_{\text{out}}^{\pm, \text{vac}}$ let us point out that $c_{\text{out}}^{\pm, \text{vac}}$ is the spectral projection on \mathbb{R}^\pm of the generator¹ of the Cauchy evolution $\mathcal{U}_{\text{out}}(t, s)$ corresponding to the asymptotic Klein-Gordon operator $P_{\text{out}} := \partial_t^2 + a_{\text{out}}$. On the other hand, to $c_{\text{out}}^{\pm, \text{vac}}$, c_{out}^\pm we can associate pairs of operators $\Lambda_{\text{out}}^{\pm, \text{vac}}, \Lambda_{\text{out}}^\pm : C_c^\infty(M) \rightarrow C^\infty(M)$ by

$$\begin{aligned} \Lambda_{\text{out}}^{\pm, \text{vac}}(t, s) &:= \mp \pi_0 \mathcal{U}_{\text{out}}(t, 0) c_{\text{out}}^{\pm, \text{vac}} \mathcal{U}_{\text{out}}(0, s) \pi_1^*, \\ \Lambda_{\text{out}}^\pm(t, s) &:= \mp \pi_0 \mathcal{U}(t, 0) c_{\text{out}}^\pm(0) \mathcal{U}(0, s) \pi_1^*, \end{aligned}$$

where we wrote $\Lambda_{\text{out}}^{\pm, \text{vac}}, \Lambda_{\text{out}}^\pm$ as operator-valued Schwartz kernels in the time variable and π_0, π_1 are the respective projections to the two pieces of Cauchy data. In QFT

¹This generator is selfadjoint for the energy scalar product.

terms (strictly speaking, using the terminology for charged fields), the operators $\Lambda_{\text{out}}^{\pm, \text{vac}}$, $\Lambda_{\text{out}}^{\pm}$ are *two-point functions*, i.e. they satisfy

$$P_{\text{out}}\Lambda_{\text{out}}^{\pm, \text{vac}} = \Lambda_{\text{out}}^{\pm, \text{vac}}P_{\text{out}} = 0, \quad \Lambda_{\text{out}}^{+, \text{vac}} - \Lambda_{\text{out}}^{-, \text{vac}} = iG_{\text{out}}, \quad \Lambda_{\text{out}}^{\pm, \text{vac}} \geq 0,$$

$$P\Lambda_{\text{out}}^{\pm} = \Lambda_{\text{out}}^{\pm}P = 0, \quad \Lambda_{\text{out}}^{+} - \Lambda_{\text{out}}^{-} = iG, \quad \Lambda_{\text{out}}^{\pm} \geq 0,$$

where G_{out} , G are the *causal propagators* for respectively P_{out} , P , i.e.

$$G_{\text{out}}(t, s) = i\pi_0 \mathcal{U}_{\text{out}}(t, s)\pi_1^*, \quad G(t, s) = i\pi_0 \mathcal{U}(t, s)\pi_1^*.$$

As a consequence, using the standard apparatus of algebraic QFT one can associate states $\omega_{\text{out}}^{\text{vac}}$, ω_{out} on the corresponding CCR C^* -algebras: $\omega_{\text{out}}^{\text{vac}}$ is then the very well studied ground state associated with P_{out} and ω_{out} is the *out* state that we study.

Our first result can be expressed as follows in terms of the two-point functions $\Lambda_{\text{out}}^{\pm}$.

Theorem 1.1. *Assume (td). Then the limit (1.2) exists and ω_{out} is a Hadamard state, i.e. the two-point functions $\Lambda_{\text{out}}^{\pm}$ satisfy the Hadamard condition:*

$$(1.3) \quad \text{WF}'(\Lambda_{\text{out}}^{\pm}) \subset \mathcal{N}^{\pm} \times \mathcal{N}^{\pm},$$

where \mathcal{N}^+ , \mathcal{N}^- are the two connected components of the characteristic set $\mathcal{N} \subset T^*M \setminus o$ of P .

Above, $\text{WF}'(\Lambda_{\text{out}}^{\pm})$ stands for the primed wave front set of $\Lambda_{\text{out}}^{\pm}$, i.e. it is the image of the wave front set of the (full) Schwartz kernel of $\Lambda_{\text{out}}^{\pm}$ by the map $(x, \xi, x', \xi') \mapsto (x, \xi, x', -\xi')$. We recall that the wave front set of a distribution characterizes the location (x, x') of its singularities, as well as the responsible directions (ξ, ξ') in Fourier space, see [Hö] for the precise definition and basic properties, cf. [BDH] for a concise introduction. The *characteristic set* of P is by definition $\mathcal{N} = p^{-1}(\{0\})$ understood as a subset of $T^*M \setminus o$ (where o is the zero section of the cotangent bundle), where $p(x, \xi) = \xi \cdot g^{-1}(x)\xi$ is the principal symbol of P .

The essential feature of the Hadamard condition (1.3) is that it constraints $\text{WF}'(\Lambda_{\text{out}}^{\pm})$ to the positive/negative frequency components $\mathcal{N}^{\pm} \times \mathcal{N}^{\pm}$ (rather than merely to $\mathcal{N} \times \mathcal{N}$, as would be the case for very general classes of bi-solutions). Thus, on a very heuristic level, the plausibility of this statement can be explained as follows. In a static situation, $c_{\text{out}}^{\pm, \text{vac}}$ are interpreted as projections that single out Cauchy data that propagate as superpositions of plane waves with positive/negative frequency, and thus with wave front set in \mathcal{N}^{\pm} . On a generic asymptotically flat spacetime it is not immediately clear what the analogous decomposition at finite times is, but instead one can try to use the decomposition given by $c_{\text{out}}^{\pm, \text{vac}}$ at *infinite times*: this is what indeed motivates the definition of $\Lambda_{\text{out}}^{\pm}$. The crucial difficulty is however to control the wave front set of the infinite time limit (1.2).

In addition to the statement of Thm. 1.1, we get in a similar vein a Hadamard state ω_{in} by taking the analogous limit with $t_- \rightarrow -\infty$ instead of $t_+ \rightarrow +\infty$; this is the so-called *in* state.

Furthermore, our results extend to a more general class of asymptotically static spacetimes $M = \mathbb{R} \times \Sigma$ with metric of the form

$$g = -c^2(x)dt^2 + (dx^i + b^i(x)dt)h_{ij}(x)(dx^j + b^j(x)dt),$$

where (Σ, h) is a manifold of bounded geometry and c, h, b as well as their inverses are bounded with all derivatives (with respect to the norm defined using a reference Riemannian metric). By *asymptotically static* we mean that there exist Riemannian metrics $h_{\text{out/in}}$ and smooth functions $c_{\text{out/in}}$ on Σ , such that on $\mathbb{R}^{\pm} \times \Sigma$,

$$\text{(ast)} \quad \begin{aligned} h(x) - h_{\text{out/in}}(x) &\in S^{-\mu}, \\ b(x) &\in S^{-\mu'}, \text{ and } c(x) - c_{\text{out/in}}(x) \in S^{-\mu} \end{aligned}$$

for some $\mu > 0$, $\mu' > 1$; in a similar vein the potential V is required to satisfy $V(x) - V_{\text{out/in}}(x) \in S^{-\mu}$ for some smooth $V_{\text{out/in}}$. Above, the notation $f \in S^{-\mu}$ means symbolic decay in time, i.e. $\partial_t^\alpha f \in O(\langle t \rangle^{-\mu-|\alpha|})$ for all $\alpha \in \mathbb{N}^{1+d}$; we refer to Subsect. 5.1 for the precise formulation.

In this more general situation, the Klein-Gordon operator is not necessarily of the form (1.1) considered so far. However, under a positivity assumption (pos) on $V_{\text{out/in}}$, it turns out that there are natural coordinates in terms of which the Klein-Gordon operator is very closely related to an operator (1.1) satisfying (td), i.e. one is obtained from the other by conjugation with some multiplication operators. This allows us to give a very similar definition of the *out/in* state $\omega_{\text{out/in}}$ and to prove a direct analogue of Thm. 1.1.

1.2. Outline of proofs. The importance of our result stems from the fact that it brings together for the first time methods from scattering theory and the analysis of Hadamard states. The main technical ingredient that we use in the proof of our theorem is an approximate diagonalization² of the Cauchy evolution by means of elliptic pseudodifferential operators, derived in detail in [GOW] and based on the strategy developed successively in the papers [Ju, JS, GW1, GW2]. Specifically, its outcome is that the Cauchy evolution of P can be written as

$$(1.4) \quad \mathcal{U}(s, t) = T(t)\mathcal{U}^{\text{ad}}(t, s)T(s)^{-1}$$

where $T(t)$ is a 2×2 matrix of pseudodifferential operators (smoothly depending on t) and $\mathcal{U}^{\text{ad}}(t, s)$ is ‘almost diagonal’ (hence the superscript ad). Namely, $\mathcal{U}^{\text{ad}}(s, t)$ is the Cauchy evolution of a time-dependent operator of the form $i\partial_t + H^{\text{ad}}(t)$, where

$$H^{\text{ad}}(t) = \begin{pmatrix} \epsilon^+(t) & 0 \\ 0 & \epsilon^-(t) \end{pmatrix}$$

modulo smooth terms (more precisely, modulo terms in $C^\infty(\mathbb{R}^2, \mathcal{W}^{-\infty}(\Sigma) \otimes \mathbb{C}^2)$, where $\mathcal{W}^{-\infty}(\Sigma)$ are the operators that map $H^{-m}(\Sigma)$ to $H^m(\Sigma)$ for each $m \in \mathbb{N}$), and $\epsilon^\pm(t)$ are elliptic pseudodifferential operators of order 1 with principal symbol $\pm(k \cdot h_t^{-1}k)^{\frac{1}{2}}$. Now, because of this particular form of the principal symbol, solutions of $(i^{-1}\partial_t + \epsilon^\pm(t))$ propagate with wave front set in \mathcal{N}^\pm . This serves one to prove that if we fix some $t_0 \in \mathbb{R}$ and set

$$c_{\text{ref}}^\pm(t_0) := T(t_0)\pi^\pm T^{-1}(t_0), \quad \text{where } \pi^+ = \begin{pmatrix} \mathbf{1} & 0 \\ 0 & 0 \end{pmatrix}, \quad \pi^- = \begin{pmatrix} 0 & 0 \\ 0 & \mathbf{1} \end{pmatrix},$$

then $\Lambda_{\text{ref}}^\pm(t, s) := \mp\pi_0 \mathcal{U}(t, t_0)c_{\text{ref}}^\pm(t_0)\mathcal{U}(t_0, s)\pi_1^*$ have wave front set only in $\mathcal{N}^\pm \times \mathcal{N}^\pm$ and therefore satisfy the Hadamard condition (1.3). As a consequence, to prove the Hadamard condition for $\Lambda_{\text{in/out}}^\pm$ it suffices to show that

$$c_{\text{in/out}}^\pm - c_{\text{ref}}^\pm \in \mathcal{W}^{-\infty}(\Sigma) \otimes B(\mathbb{C}^2).$$

To demonstrate that this is the case, we use assumption (td) to control the decay in time of various remainders in identities ‘modulo smooth’. The most crucial estimate here is

$$(1.5) \quad H^{\text{ad}}(t) - \begin{pmatrix} a(t)^{\frac{1}{2}} & 0 \\ 0 & -a(t)^{\frac{1}{2}} \end{pmatrix} \in \Psi_{\text{td}}^{0, -1-\delta}(\mathbb{R}; \Sigma) \otimes B(\mathbb{C}^2)$$

for large t , which then yields time-decay of various commutators that appear in the proofs. We obtain (1.5) by revisiting the approximate diagonalization (1.4)

²On a side note, let us mention that a different diagonalization procedure was proposed by Ruzhansky and Wirth in the context of dispersive estimates [RW, Wi]; in their method it is the (full) symbol of the generator of the Cauchy evolution that is diagonalized (rather than the Cauchy evolution itself).

using poly-homogeneous expansions of pseudodifferential operators in $\Psi_{\text{td}}^{m,-\delta}(\mathbb{R}; \Sigma)$ in both m and δ ; more details are given in Sect. 4.5. This requires us to study the classes $\Psi_{\text{td}}^{m,-\delta}(\mathbb{R}; \Sigma)$ carefully, in particular we employ a variant of Seeley's theorem on powers of pseudodifferential operators elliptic in the standard Ψ^m sense.

1.3. Plan of the paper. The paper is structured as follows.

In Sect. 2 we fix some basic terminology and recall the definition of two-point functions and covariances of states in the context of non-interacting Quantum Field Theory.

Sect. 3 contains a brief overview of the pseudodifferential calculus on manifolds of bounded geometry. We then introduce the time-dependent pseudodifferential operator classes $\Psi_{\text{td}}^{m,\delta}$ and study some of their properties.

In Sect. 4 we recall the approximate diagonalization of the Cauchy evolution used in [GOW] to construct generic Hadamard states. We then give a refinement in the setup of assumption (td) by showing time decay of various remainder terms.

Sect. 5 contains the construction of the *in* and *out* states and the proof of their Hadamard property in the case of asymptotically static spacetimes (assumptions (ast) and (pos)). The key ingredients are the reduction to the setup of assumption (td) and the estimates obtained in Sect. 4.

Various auxiliary proofs are collected in Appendix A.

2. PRELIMINARIES

2.1. Notation. The space of differential operators (of order m) over a smooth manifold M (here always without boundary) is denoted $\text{Diff}(M)$ ($\text{Diff}^m(M)$). The space of smooth functions on M with compact support is denoted $C_c^\infty(M)$.

The operator of multiplication by a function f will be denoted by f , while the operators of partial differentiation will be denoted by $\bar{\partial}_i$, so that $[\bar{\partial}_i, f] = \partial_i f$.

- If a, b are selfadjoint operators on a Hilbert space \mathcal{H} , we write $a \sim b$ if

$$a, b > 0, \quad \text{Dom } a^{\frac{1}{2}} = \text{Dom } b^{\frac{1}{2}}, \quad c^{-1}b \leq a \leq cb,$$

for some constant $c > 0$.

- Similarly, if $I \subset \mathbb{R}$ is an open interval and $\{\mathcal{H}_t\}_{t \in I}$ is a family of Hilbert spaces with $\mathcal{H}_t = \mathcal{H}$ as topological vector spaces, and $a(t), b(t)$ are two selfadjoint operators on \mathcal{H}_t , we write $a(t) \sim b(t)$ if for each $J \Subset I$ there exist constants $c_{1,J}, c_{2,J} > 0$ such that

$$(2.1) \quad a(t), b(t) \geq c_{1,J} > 0, \quad c_{2,J} b(t) \leq a(t) \leq c_{2,J}^{-1} b(t), \quad t \in J.$$

2.2. Klein-Gordon operator. Let (M, g) be a Lorentzian spacetime (we use the convention $(-, +, \dots, +)$ for the Lorentzian signature). We consider the Klein-Gordon operator with a real-valued potential $V \in C^\infty(M)$

$$P = -\square_g + V \in \text{Diff}^2(M),$$

Since V is real-valued we have $P = P^*$ in the sense of formal adjoints with respect to the $L^2(M, g)$ scalar product, canonically defined using the volume form.

For $K \subset M$ we denote $J_\pm(K) \subset M$ its causal future/past, see e.g. [BF2, Wa2]. Let $C_\pm^\infty(M)$ be the space of smooth functions whose support is future or past compact, that is

$$C_\pm^\infty(M) = \{f \in C^\infty(M) : \text{supp } f \subset J_\pm(K) \text{ for some compact } K \subset M\}.$$

We assume that (M, g) is globally hyperbolic, i.e. admits a foliation by Cauchy surfaces³ (in the next sections we will impose more restrictive conditions on (M, g) , but these are irrelevant for the moment). It is well known that P has then unique *advanced/retarded propagators*, i.e. operators $G_{\pm} : C_{\pm}^{\infty}(M) \rightarrow C_{\pm}^{\infty}(M)$ s.t.

$$(2.2) \quad PG_{\pm} = \mathbf{1} \quad \text{on} \quad C_{\pm}^{\infty}(M).$$

A standard duality argument that uses $P = P^*$, (2.2), and the fact that $C_{+}^{\infty}(M) \cap C_{-}^{\infty}(M) = C_c^{\infty}(M)$ on globally hyperbolic spacetimes, gives $G_{+}^* = G_{-}$ as sesquilinear forms on $C_c^{\infty}(M)$. The *causal propagator* (often also called *Pauli-Jordan commutator function*) of P is by definition $G := G_{+} - G_{-}$, interpreted here as a map from $C_c^{\infty}(M)$ to $C_{+}^{\infty}(M) + C_{-}^{\infty}(M)$, the space of space-compact smooth functions.

2.3. Symplectic space of solutions. In what follows we recall the relation between quasi-free states, two-point functions, and field quantization. The reader interested only in the analytical aspects can skip this discussion and move directly to equations (2.6)–(2.8), which can be taken as the definition of two-point functions in the present context.

By a *phase space* we will mean a pair (\mathcal{V}, q) consisting of a complex vector space \mathcal{V} and a non degenerate hermitian form q on \mathcal{V} . In our case the phase space of interest (i.e. the phase space of the classical non-interacting scalar field theory) is

$$(2.3) \quad \mathcal{V} := \frac{C_c^{\infty}(M)}{PC_c^{\infty}(M)}, \quad \bar{u}qv := i^{-1}(u|Gv),$$

where $(\cdot|\cdot)$ is the $L^2(M, g)$ pairing. The sesquilinear form q is indeed well-defined on the quotient space $C_c^{\infty}(M)/PC_c^{\infty}(M)$ because $PG = GP = 0$ on test functions. Using that $G_{+}^* = G_{-}$ one shows that q is hermitian, and it is also not difficult to show that it is non-degenerate.

Note that in contrast to most of the literature, we work with hermitian forms rather than with real symplectic ones, but the two approaches are equivalent.

2.4. States and their two-point functions. Let \mathcal{V} be a complex vector space, \mathcal{V}^* its anti-dual and $L_h(\mathcal{V}, \mathcal{V}^*)$ the space of hermitian sesquilinear forms on \mathcal{V} . If $q \in L_h(\mathcal{V}, \mathcal{V}^*)$ then we can define the polynomial CCR $*$ -algebra $\text{CCR}^{\text{pol}}(\mathcal{V}, q)$ (see e.g. [DG2, Sect. 8.3.1])⁴. It is constructed as the span of the so-called *abstract complex fields* $\mathcal{V} \ni v \mapsto \psi(v), \psi^*(v)$, which are taken to be anti-linear, resp. linear in v and are subject to the canonical commutation relations

$$[\psi(v), \psi(w)] = [\psi^*(v), \psi^*(w)] = 0, \quad [\psi(v), \psi^*(w)] = \bar{v}qw\mathbf{1}, \quad v, w \in \mathcal{V}.$$

Our main object of interests are the states⁵ on $\text{CCR}^{\text{pol}}(\mathcal{V}, q)$.

The *complex covariances* $\Lambda^{\pm} \in L_h(\mathcal{V}, \mathcal{V}^*)$ of a state ω on $\text{CCR}^{\text{pol}}(\mathcal{V}, q)$ are defined in terms of the abstract field operators by

$$(2.4) \quad \bar{v}\Lambda^{+}w = \omega(\psi(v)\psi^*(w)), \quad \bar{v}\Lambda^{-}w = \omega(\psi^*(w)\psi(v)), \quad v, w \in \mathcal{V}$$

Note that both Λ^{\pm} are positive and by the canonical commutation relations one always has $\Lambda^{+} - \Lambda^{-} = q$. We are interested in the reverse construction, namely if one has a pair of hermitian forms Λ^{\pm} such that $\Lambda^{+} - \Lambda^{-} = q$ and $\Lambda^{\pm} \geq 0$ then there is a unique *quasi-free* state ω such that (2.4) holds. We will thus further restrict our attention to quasi-free states and more specifically to their complex covariances Λ^{\pm} .

³Let us recall that a Cauchy surface is a smooth hypersurface that is intersected by every inextendible, non-spacelike (i.e. causal) curve exactly once.

⁴See also [GW1, Wr] for remarks on the transition between real and complex vector space terminology.

⁵Let us recall that a state ω is a linear functional on $\text{CCR}^{\text{pol}}(\mathcal{V}, q)$ such that $\omega(a^*a) \geq 0$ for all a in $\text{CCR}^{\text{pol}}(\mathcal{V}, q)$, and $\omega(\mathbf{1}) = 1$.

In QFT (at least for scalar fields) the phase space of interest is the one defined in (2.3). In that specific case it is convenient to consider instead of complex covariances a pair of operators $\Lambda^\pm : C_c^\infty(M) \rightarrow C^\infty(M)$ such that

$$(2.5) \quad (v|\Lambda^+w) = \omega(\psi(v)\psi^*(w)), \quad (v|\Lambda^-w) = \omega(\psi^*(w)\psi(v)), \quad v, w \in C_c^\infty(M).$$

We call Λ^\pm the *two-point functions* of the state ω and identify them with the associated complex covariances whenever possible. Note that because $(\cdot|\Lambda^\pm\cdot)$ has to induce a hermitian form on the quotient space $C_c^\infty(M)/PC_c^\infty(M)$, the two-point functions have to satisfy $P\Lambda^\pm = \Lambda^\pm P = 0$ on $C_c^\infty(M)$. By the Schwartz kernel theorem we can further identify Λ^\pm with a pair of distributions on $M \times M$, these are then bi-solutions of the Klein-Gordon equation.

In QFT on curved spacetime one is especially interested in the subclass of quasi-free *Hadamard states* [KW, Ra]. These can be defined as in the introduction, i.e. by requiring that the primed wave front set of the Schwartz kernel of Λ^\pm is contained in $\mathcal{N}^\pm \times \mathcal{N}^\pm$ (cf. e.g. [Ra, SV, Sa] for a discussion of various equivalent formulations), $\mathcal{N}^\pm \subset T^*M \setminus o$ being the two connected components of the characteristic set of P (and $o \subset T^*M$ the zero section). To sum this up, specifying a Hadamard state amounts to constructing a pair of operators $\Lambda^\pm : C_c^\infty(M) \rightarrow C^\infty(M)$ satisfying the properties:

$$(2.6) \quad P\Lambda^\pm = \Lambda^\pm P = 0, \quad \Lambda^+ - \Lambda^- = iG,$$

$$(2.7) \quad \Lambda^\pm \geq 0,$$

$$(2.8) \quad \text{WF}'(\Lambda^\pm) \subset \mathcal{N}^\pm \times \mathcal{N}^\pm.$$

Existence of two-point functions as above was proved in [FNW], and an alternative argument was given in [GW1], followed by the construction of a very large class of examples in [GOW]. The importance of Hadamard states is primarily due to their pivotal role in renormalization on curved spacetimes [BF1, HW1, HW2, Da], see [FV2, KM, HW3] for recent reviews.

Here we will be interested in showing (2.8) for specific two-point functions with prescribed asymptotic properties, motivated by the conceptual need for *distinguished* Hadamard states whenever allowed by the spacetime geometry.

2.5. Cauchy data of two-point functions. We will need a version of two-point functions acting on Cauchy data of P instead of spacetime quantities such as Λ^\pm . To this end, let $\{\Sigma_s\}_{s \in \mathbb{R}}$ be a foliation of M by Cauchy surfaces (since all Σ_s are diffeomorphic we occasionally write Σ instead). We define the map

$$\varrho_s u := (u, i^{-1}n^a \nabla_a u)|_{\Sigma_s},$$

acting on distributions u such that the restriction $|_{\Sigma_s}$ makes sense, where n^a is the unit normal vector to Σ_s . It is well-known that $\varrho_s \circ G$ maps $C_c^\infty(M)$ to $C_c^\infty(\Sigma_s)$ and that there exists an operator $G(s)$ acting on $C_c^\infty(\Sigma) \otimes \mathbb{C}^2$ (not to be confused with G) that satisfies

$$(2.9) \quad G =: (\varrho_s G)^* \circ G(s) \circ \varrho_s G,$$

where $(\varrho_s G)^*$ is the formal adjoint of $\varrho_s \circ G$ wrt. the L^2 inner product on $\Sigma_s \sqcup \Sigma_s$ respective to some density (that can depend on s , later on we will make that choice more specific). We also set

$$q(s) := i^{-1}G(s),$$

so that $q(s)^* = q(s)$.

The next result provides a Cauchy surface analogue of the two-point functions Λ^\pm , cf. [GW2] for the proof.

Proposition 2.1. *For any $s \in \mathbb{R}$ the maps:*

$$(2.10) \quad \lambda^\pm(s) \mapsto \Lambda^\pm := (\varrho_s G)^* \lambda^\pm(s) (\varrho_s G),$$

and

$$(2.11) \quad \Lambda^\pm \mapsto \lambda^\pm(s) := (\varrho_s^* G(s))^* \Lambda^\pm (\varrho_s^* G(s))$$

are bijective and inverse from one another.

It is actually convenient to make one more definition and set:

$$(2.12) \quad c^\pm(s) = \pm i^{-1} G(s) \lambda^\pm(s) : C_c^\infty(\Sigma) \otimes B(\mathbb{C}^2) \rightarrow C^\infty(\Sigma) \otimes B(\mathbb{C}^2).$$

We will simply call $c^\pm(s)$ the (*time- s*) *covariances of the state ω* . A pair of operators $c^\pm(s)$ are covariances of a state iff the operators Λ^\pm defined by (2.10) and (2.12) satisfy (2.6)-(2.7), which is equivalent to the conditions

$$(2.13) \quad c^+(s) + c^-(s) = \mathbf{1},$$

$$(2.14) \quad \lambda^\pm(s) \geq 0,$$

where we identified the operators $\lambda^\pm(s)$ with hermitian forms using the same pairing as when we took the formal adjoint in (2.9). Note that (2.13) can also be expressed as $\lambda^+(s) - \lambda^-(s) = q(s)$.

Additionally, a state (recall that we consider only quasi-free states) is *pure* iff its covariances $c^\pm(s)$ extend to projections on the completion of $C_c^\infty(\Sigma) \otimes \mathbb{C}^2$ w.r.t. the inner product given by $\lambda^+ + \lambda^-$. In practice it is sufficient to construct $c^\pm(s)$ as projections acting on a space that is big enough to contain $C_c^\infty(\Sigma) \otimes \mathbb{C}^2$, but small enough to be contained in the Hilbert space associated to $\lambda^+ + \lambda^-$.

2.6. Propagators for the Cauchy evolution. Recall that we have defined the operator $G(s)$ via the identity

$$(2.15) \quad G =: (\varrho_s G)^* \circ G(s) \circ \varrho_s G.$$

A direct consequence is that the operator $G^* \varrho_s G(s)$ assigns to Cauchy data on Σ_s the corresponding solution. Similarly, for $t, s \in \mathbb{R}$ the operator

$$(2.16) \quad \mathcal{U}(s, t) := \varrho_s G^* \varrho_t^* G(t)$$

produces Cauchy data of a solution on Σ_s given Cauchy data on Σ_t . We will call $\{\mathcal{U}(s, t)\}_{s, t \in \mathbb{R}}$ the *Cauchy evolution* of P . A straightforward computation gives the group property

$$(2.17) \quad \mathcal{U}(t, t) = \mathbf{1}, \quad \mathcal{U}(s, t') \mathcal{U}(t', t) = \mathcal{U}(s, t), \quad t' \in \mathbb{R};$$

and the conservation of the symplectic form by the evolution

$$(2.18) \quad \mathcal{U}^*(s, t) q(s) \mathcal{U}(s, t) = q(t).$$

These identities allow to conclude that the covariances $c^\pm(t)$ (and two-point functions $\lambda^\pm(t)$) at different ‘times’ of a quasi-free state are related by

$$(2.19) \quad \begin{aligned} \lambda^\pm(t) &= \mathcal{U}(s, t)^* \lambda^\pm(s) \mathcal{U}(s, t), \\ c^\pm(t) &= \mathcal{U}(t, s) c^\pm(s) \mathcal{U}(s, t). \end{aligned}$$

Notice that this induces a splitting of the evolution in two parts:

$$\mathcal{U}(s, t) = \mathcal{U}^+(s, t) + \mathcal{U}^-(s, t), \quad \text{with } \mathcal{U}^\pm(s, t) = \mathcal{U}(s, t) c^\pm(t).$$

If the state is pure then $c^\pm(t)$ are projections for all t and the operators $\mathcal{U}^\pm(s, t)$ obey the composition formula

$$\mathcal{U}^\pm(s, t') \mathcal{U}^\pm(t', t) = \mathcal{U}^\pm(s, t), \quad \mathcal{U}^\pm(s, t') \mathcal{U}^\mp(t', t) = 0, \quad t' \in \mathbb{R}.$$

Let us stress that $\mathcal{U}^\pm(t, t)$ is not the identity, but rather equals $c^\pm(t)$. Furthermore, if the state is Hadamard then $\mathcal{U}(s, t) c^\pm(t)$ propagate singularities along \mathcal{N}^\pm (see

the discussion in [GW2]). In Sect. 4 we will be interested in the reversed argument, namely we will construct covariances $c^\pm(t)$ of pure Hadamard states from a splitting of the evolution $\mathcal{U}(t, s)$ into two parts that propagate singularities along respectively \mathcal{N}^+ , \mathcal{N}^- .

3. PSEUDODIFFERENTIAL CALCULUS ON MANIFOLDS OF BOUNDED GEOMETRY

3.1. Manifolds of bounded geometry. In the present section we introduce manifolds of bounded geometry and review the pseudodifferential calculus of Kordyukov and Shubin [Ko, Sh1], making also use of some results from [GOW].

Let us denote by δ the flat metric on \mathbb{R}^d and by $B_d(y, r) \subset \mathbb{R}^d$ the open ball of center y and radius r .

If (Σ, h) is a d -dimensional Riemannian manifold and X is a (p, q) tensor on Σ , we can define the canonical norm of $X(x)$, $x \in \Sigma$, denoted by $\|X\|_x$, using appropriate tensor powers of $h(x)$ and $h^{-1}(x)$. X is *bounded* if $\sup_{x \in \Sigma} \|X\|_x < \infty$.

If $U \subset \Sigma$ is open, we denote by $\text{BT}_q^p(U, \delta)$ the Fréchet space of (p, q) tensors on U , bounded with all covariant derivatives in the above sense. Among several equivalent definitions of manifolds of bounded geometry (see [Sh1, GOW]), the one below is particularly useful in applications.

Definition 3.1. *A Riemannian manifold (Σ, h) is of bounded geometry iff for each $x \in \Sigma$, there exists an open neighborhood of x , denoted U_x , and a smooth diffeomorphism*

$$\psi_x : U_x \xrightarrow{\sim} B_d(0, 1) \subset \mathbb{R}^d$$

with $\psi_x(x) = 0$, and such that if $h_x := (\psi_x^{-1})^* h$ then:

(C1) *the family $\{h_x\}_{x \in \Sigma}$ is bounded in $\text{BT}_2^0(B_d(0, 1), \delta)$,*

(C2) *there exists $c > 0$ such that :*

$$c^{-1}\delta \leq h_x \leq c\delta, \quad x \in \Sigma.$$

A family $\{U_x\}_{x \in \Sigma}$ resp. $\{\psi_x\}_{x \in \Sigma}$ as above will be called a family of good chart neighborhoods, resp. good chart diffeomorphisms.

A known result (see [Sh1, Lemma 1.2]) says that one can find a covering $\Sigma = \bigcup_{i \in \mathbb{N}} U_i$ by good chart neighborhoods $U_i = U_{x_i}$ ($x_i \in \Sigma$) which is *uniformly finite*, i.e. there exists $N \in \mathbb{N}$ such that $\bigcap_{i \in I} U_i = \emptyset$ if $\#I > N$. Setting $\psi_i = \psi_{x_i}$, we will call the sequence $\{U_i, \psi_i\}_{i \in \mathbb{N}}$ a *good chart covering* of Σ .

Furthermore, by [Sh1, Lemma 1.3] one can associate to a good chart covering a partition of unity:

$$1 = \sum_{i \in \mathbb{N}} \chi_i^2, \quad \chi_i \in C_c^\infty(U_i)$$

such that $\{(\psi_i^{-1})^* \chi_i\}_{i \in \mathbb{N}}$ is a bounded sequence in $C_b^\infty(B_d(0, 1))$. Such a partition of unity will be called a *good partition of unity*.

3.2. Bounded tensors and bounded diffeomorphisms.

Definition 3.2. *Let (Σ, h) be of bounded geometry. We denote by $\text{BT}_q^p(\Sigma, h)$ the spaces of smooth (q, p) tensors X on Σ such that if $X_x = (\exp_x^h \circ e_x)^* X$, where $e_x : (\mathbb{R}^d, \delta) \rightarrow (T_x \Sigma, h(x))$ is an isometry, then the family $\{X_x\}_{x \in \Sigma}$ is bounded in $\text{BT}_q^p(B_d(0, \frac{r}{2}), \delta)$. We equip $\text{BT}_q^p(\Sigma, h)$ with its natural Fréchet space topology.*

We denote by $C_b^\infty(\mathbb{R}; \text{BT}_q^p(\Sigma, h))$ the space of smooth maps $\mathbb{R} \in t \mapsto X(t)$ such that $\partial_t^n X(t)$ is uniformly bounded in $\text{BT}_q^p(\Sigma, h)$ for $n \in \mathbb{N}$.

We denote by $S^\delta(\mathbb{R}; \text{BT}_q^p(\Sigma, h))$, $\delta \in \mathbb{R}$ the space of smooth maps $\mathbb{R} \in t \mapsto X(t)$ such that $\langle t \rangle^{-\delta+n} \partial_t^n X(t)$ is uniformly bounded in $\text{BT}_q^p(\Sigma, h)$ for $n \in \mathbb{N}$.

It is well known (see e.g. [GOW, Subsect. 2.3]) that we can replace in Def. 3.2 the geodesic maps $\exp_x^h \circ e_x$ by ψ_x^{-1} , where $\{\psi_x\}_{x \in \Sigma}$ is any family of good chart diffeomorphisms as in Thm. 3.1.

Definition 3.3. *Let (Σ, h) be an n -dimensional Riemannian manifold of bounded geometry and $\chi : \Sigma \rightarrow \Sigma$ a smooth diffeomorphism. One says that χ is a bounded diffeomorphism of (Σ, h) if for some some family of good chart diffeomorphisms $\{U_x, \psi_x\}_{x \in \Sigma}$, the maps*

$$\chi_x = \psi_{\chi(x)} \circ \chi \circ \psi_x^{-1}, \quad \chi_x^{-1} = \psi_{\chi^{-1}(x)} \circ \chi^{-1} \circ \psi_x : B_n(0, 1) \rightarrow B_n(0, 1)$$

are bounded in $C_b^\infty(B_n(0, 1))$ uniformly with respect to $x \in \Sigma$.

It is easy to see that if the above properties are satisfied for some family of good chart diffeomorphisms then they are satisfied for any such family, furthermore bounded diffeomorphisms are stable under composition.

3.3. Symbol classes. We recall some well-known definitions about symbol classes on manifolds of bounded geometry, following [Sh1, Ko, ALNV].

3.3.1. Symbol classes on \mathbb{R}^n . Let $U \subset \mathbb{R}^d$ be an open set, equipped with the flat metric δ on \mathbb{R}^d .

We denote by $S^m(T^*U)$, $m \in \mathbb{R}$, the space of $a \in C^\infty(U \times \mathbb{R}^d)$ such that

$$\langle k \rangle^{-m+|\beta|} \partial_x^\alpha \partial_\xi^\beta a(x, k) \text{ is bounded on } U \times \mathbb{R}^d, \quad \forall \alpha, \beta \in \mathbb{N}^n,$$

equipped with its canonical semi-norms $\|\cdot\|_{m, \alpha, \beta}$.

We set

$$S^{-\infty}(T^*U) := \bigcap_{m \in \mathbb{R}} S^m(T^*U), \quad S^\infty(T^*U) := \bigcup_{m \in \mathbb{R}} S^m(T^*U),$$

with their canonical Fréchet space topologies. If $m \in \mathbb{R}$ and $a_{m-i} \in S^{m-i}(T^*U)$ we write $a \simeq \sum_{i \in \mathbb{N}} a_{m-i}$ if for each $p \in \mathbb{N}$

$$(3.1) \quad r_p(a) := a - \sum_{i=0}^p a_{m-i} \in S^{m-p-1}(T^*U).$$

It is well-known (see e.g. [Sh2, Sect. 3.3]) that if $a_{m-i} \in S^{m-i}(T^*U)$, there exists $a \in S^m(T^*U)$, unique modulo $S^{-\infty}(T^*U)$ such that $a \simeq \sum_{i \in \mathbb{N}} a_{m-i}$.

We denote by $S_h^m(T^*U) \subset S^m(T^*U)$ the space of a such that $a(x, \lambda k) = \lambda^m a(x, k)$, for $x \in U$, $|k| \geq C$, $C > 0$ and by $S_{\text{ph}}^m(T^*U) \subset S^m(T^*U)$ the space of a such that $a \simeq \sum_{i \in \mathbb{N}} a_{m-i}$ for a sequence $a_{m-i} \in S_h^{m-i}(T^*U)$ (a is then called a *polyhomogeneous*⁶ symbol). Following [ALNV] one equips $S_{\text{ph}}^m(T^*U)$ with the topology defined by the semi-norms of a_{m-i} in $S^{m-i}(T^*U)$ and $r_p(a)$ in $S^{m-p-1}(T^*U)$, (see (3.1)). This topology is strictly stronger than the topology induced by $S^m(T^*U)$.

The space $S_{\text{ph}}^m(T^*U)/S_{\text{ph}}^{m-1}(T^*U)$ is isomorphic to $S_h^m(T^*U)$, and the image of a under the quotient map is called the *principal symbol* of a and denoted by $\sigma_{\text{pr}}(a)$.

If $U = B_n(0, 1)$ (more generally, if U is relatively compact with smooth boundary), there exists a continuous extension map $E : S^m(T^*U) \rightarrow S^m(T^*\mathbb{R}^d)$ such that $Ea|_{T^*U} = a$. Moreover E maps $S_{\text{ph}}^m(T^*U)$ into $S_{\text{ph}}^m(T^*\mathbb{R}^d)$ and is continuous for the topologies of $S_{\text{ph}}^m(T^*U)$ and $S_{\text{ph}}^m(T^*\mathbb{R}^d)$, which means that all the maps

$$a \mapsto (Ea)_{m-i}, \quad a \mapsto r_p(Ea),$$

are continuous.

⁶These are also called classical symbols in the literature.

3.3.2. *Time-dependent symbol classes on \mathbb{R}^d .* We will also need to consider various classes of *time-dependent* symbols $a(t, x, k) \in C^\infty(\mathbb{R} \times T^*U)$. First of all the space $C^\infty(\mathbb{R}; S^m(T^*U))$ is defined as the space of $a \in C^\infty(\mathbb{R} \times T^*U)$ such that

$$\langle k \rangle^{-m+|\beta|} \partial_t^\gamma \partial_x^\alpha \partial_k^\beta a(t, x, k) \text{ is bounded on } I \times U \times \mathbb{R}^d, \forall \alpha, \beta \in \mathbb{N}^n, \gamma \in \mathbb{N},$$

for any interval $I \Subset \mathbb{R}$. We denote by $C_b^\infty(\mathbb{R}; S^m(T^*U))$ the subspace of symbols which are uniformly bounded in $S^m(T^*U)$ with all time derivatives.

Furthermore, anticipating the need for some additional decay in t in Sect. 4.5, we denote by $S^\delta(\mathbb{R}; S^m(T^*U))$ the space of $a \in C^\infty(\mathbb{R} \times T^*U)$ such that

$$\langle t \rangle^{-\delta+\gamma} \langle k \rangle^{-m+|\beta|} \partial_t^\gamma \partial_x^\alpha \partial_k^\beta a(t, x, k) \text{ is bounded on } \mathbb{R} \times U \times \mathbb{R}^d, \forall \alpha, \beta \in \mathbb{N}^n, \gamma \in \mathbb{N}.$$

The notation $a \sim \sum_i a_{m-i}$ and the poly-homogeneous spaces

$$C_{(b)}^\infty(\mathbb{R}; S_{\text{ph}}^m(T^*U)), \quad S^\delta(\mathbb{R}; S_{\text{ph}}^m(T^*U)),$$

are defined analogously, by requiring estimates on the time derivatives of the a_{m-i} and r_p in (3.1).

3.3.3. *Symbol classes on Σ .* Let (Σ, h) be a Riemannian manifold of bounded geometry and $\{\psi_x\}_{x \in \Sigma}$ a family of good chart diffeomorphisms.

Definition 3.4. We denote by $S^m(T^*\Sigma)$ for $m \in \mathbb{R}$ the space of $a \in C^\infty(T^*\Sigma)$ such that for each $x \in \Sigma$, $a_x := (\psi_x^{-1})^* a \in S^m(T^*B_n(0, 1))$ and the family $\{a_x\}_{x \in \Sigma}$ is bounded in $S^m(T^*B_n(0, 1))$. We equip $S^m(T^*\Sigma)$ with the semi-norms

$$\|a\|_{m, \alpha, \beta} = \sup_{x \in \Sigma} \|a_x\|_{m, \alpha, \beta}.$$

Similarly we denote by $S_{\text{ph}}^m(T^*\Sigma)$ the space of $a \in S^m(T^*\Sigma)$ such that for each $x \in \Sigma$, $a_x \in S_{\text{ph}}^m(T^*B_n(0, 1))$ and the family $\{a_x\}_{x \in \Sigma}$ is bounded in $S_{\text{ph}}^m(T^*B_n(0, 1))$. We equip $S_{\text{ph}}^m(T^*\Sigma)$ with the semi-norms

$$\|a\|_{m, i, p, \alpha, \beta} = \sup_{x \in \Sigma} \|a_x\|_{m, i, p, \alpha, \beta}.$$

where $\|\cdot\|_{m, i, p, \alpha, \beta}$ are the semi-norms defining the topology of $S_{\text{ph}}^m(T^*B_n(0, 1))$.

We also set $S_{(\text{ph})}^\infty(T^*\Sigma) = \bigcup_{m \in \mathbb{R}} S_{(\text{ph})}^m(T^*\Sigma)$.

The definition of $S^m(T^*\Sigma)$, $S_{\text{ph}}^m(T^*\Sigma)$ and their Fréchet space topologies are independent on the choice of the family $\{\psi_x\}_{x \in \Sigma}$ of good chart diffeomorphisms.

The notation $a \simeq \sum_{i \in \mathbb{N}} a_{m-i}$ for $a_{m-i} \in S_{\text{ph}}^{m-i}(T^*\Sigma)$ is defined as before. If $a \in S_{\text{ph}}^m(T^*\Sigma)$, we denote again by a_{pr} the image of a in $S_{\text{ph}}^m(T^*\Sigma)/S_{\text{ph}}^{m-1}(T^*\Sigma)$.

The spaces $C_{(b)}^\infty(\mathbb{R}; S_{(\text{ph})}^m(T^*\Sigma))$, $S^\delta(\mathbb{R}; S_{(\text{ph})}^m(T^*\Sigma))$ are defined as in 3.3.2 and equipped with their natural Fréchet space topologies.

3.4. **Sobolev spaces and smoothing operators.** Using the metric h one defines the Sobolev spaces $H^m(\Sigma)$ as follows.

Definition 3.5. For $s \in \mathbb{R}$ the Sobolev space $H^s(\Sigma)$ is:

$$H^s(\Sigma) := \langle -\Delta_h \rangle^{-s/2} L^2(\Sigma),$$

with its natural Hilbert space topology, where $-\Delta_h$ is the Laplace-Beltrami operator on (Σ, h) , strictly speaking the closure of its restriction to $C_c^\infty(\Sigma)$.

We further set

$$H^\infty(\Sigma) := \bigcap_{m \in \mathbb{Z}} H^m(\Sigma), \quad H^{-\infty}(\Sigma) := \bigcup_{m \in \mathbb{Z}} H^m(\Sigma),$$

equipped with their Fréchet space topologies.

We denote by $\mathcal{W}^{-\infty}(\Sigma)$ the Fréchet space $B(H^{-\infty}(\Sigma), H^{\infty}(\Sigma))$ with its Fréchet space topology, given by the semi-norms

$$\|a\|_m = \|a\|_{B(H^{-m}(\Sigma), H^m(\Sigma))}, \quad m \in \mathbb{N}.$$

This allows us to define $C_{(b)}^{\infty}(\mathbb{R}; \mathcal{W}^{-\infty}(\Sigma))$, $S^{\delta}(\mathbb{R}; \mathcal{W}^{-\infty}(\Sigma))$, the latter consisting of operator-valued functions $a(t)$ such that

$$\|\partial_t^{\gamma} a(t)\|_m \in O(\langle t \rangle^{\delta-\gamma}), \quad \forall \gamma, m \in \mathbb{N}.$$

3.5. Pseudodifferential operators. Starting from the well-known Weyl quantization on open subsets of \mathbb{R}^d , one constructs a quantization map Op for symbols in $S^m(T^*\Sigma)$ using a good chart covering of Σ and good chart diffeomorphisms. More precisely let $\{U_i, \psi_i\}_{i \in \mathbb{N}}$ be a good chart covering of M and

$$\sum_{i \in \mathbb{N}} \chi_i^2 = \mathbf{1}$$

a subordinate good partition of unity, see Subsect. 3.1. If

$$(\psi_i^{-1})^* dg =: m_i dx,$$

we set

$$\begin{aligned} T_i : L^2(U_i, dg) &\rightarrow L^2(B_n(0, 1), dx), \\ u &\mapsto m_i^{\frac{1}{2}} (\psi_i^{-1})^* u, \end{aligned}$$

so that $T_i : L^2(U_i, dg) \rightarrow L^2(B_n(0, 1), dx)$ is unitary. We then fix an extension map

$$E : S_{\text{ph}}^m(T^*B_d(0, 1)) \rightarrow S_{\text{ph}}^m(T^*\mathbb{R}^d).$$

Definition 3.6. Let $a = a(t) \in C^{\infty}(\mathbb{R}; S_{\text{ph}}^m(T^*M))$. We set

$$\text{Op}(a) := \sum_{i \in \mathbb{N}} \chi_i T_i^* \circ \text{Op}^{\text{W}}(E a_i) \circ T_i \chi_i,$$

where $a_i \in S_{\text{ph}}^m(T^*B_d(0, 1))$ is the push-forward of $a|_{T^*U_i}$ by ψ_i and Op^{W} is the Weyl quantization.

If Op' is another such quantization map for different choices of U_i, ψ_i, χ_i and E then

$$\begin{aligned} S_{\text{ph}}^m(T^*\Sigma) &\rightarrow \mathcal{W}^{-\infty}(\Sigma) \\ \text{Op} - \text{Op}' : C_{(b)}^{\infty}(\mathbb{R}; S_{\text{ph}}^m(T^*\Sigma)) &\rightarrow C_{(b)}^{\infty}(\mathbb{R}; \mathcal{W}^{-\infty}(\Sigma)), \\ S^{\delta}(\mathbb{R}; S_{\text{ph}}^m(T^*\Sigma)) &\rightarrow S^{\delta}(\mathbb{R}; \mathcal{W}^{-\infty}(\Sigma)), \end{aligned}$$

are bounded. Then one defines the classes

$$\begin{aligned} \Psi^m(\Sigma) &:= \text{Op}(S_{\text{ph}}^m(T^*\Sigma)) + \mathcal{W}^{-\infty}(\Sigma), \\ C_{(b)}^{\infty}(\mathbb{R}; \Psi^m(\Sigma)) &:= \text{Op}(C_{(b)}^{\infty}(\mathbb{R}; S_{\text{ph}}^m(T^*\Sigma))) + C_{(b)}^{\infty}(\mathbb{R}; \mathcal{W}^{-\infty}(\Sigma)), \\ S^{\delta}(\mathbb{R}; \Psi^m(\Sigma)) &:= \text{Op}(S^{\delta}(\mathbb{R}; S_{\text{ph}}^m(T^*\Sigma))) + S^{\delta}(\mathbb{R}; \mathcal{W}^{-\infty}(\Sigma)). \end{aligned}$$

Thanks to including the ideal $\mathcal{W}^{-\infty}(\Sigma)$ of smoothing operators, the so-obtained pseudodifferential classes are stable under composition, for example

$$\Psi^{m_1}(\Sigma) \circ \Psi^{m_2}(\Sigma) \subset \Psi^{m_1+m_2}(\Sigma).$$

Note that $S^{\delta}(\mathbb{R}; \Psi^m(\Sigma)) = \langle t \rangle^{\delta} S^0(\mathbb{R}; \Psi^m(\Sigma))$ and similarly with $\Psi^m(\Sigma)$ replaced by $\mathcal{W}^{-\infty}(\Sigma)$ so in what follows one can assume without loss of generality that $\delta = 0$.

The spaces $\mathcal{W}^{-\infty}(\Sigma)$, $C_{(b)}^{\infty}(\mathbb{R}; \mathcal{W}^{-\infty}(\Sigma))$ and $S^{\delta}(\mathbb{R}; \mathcal{W}^{-\infty}(\Sigma))$ have natural Fréchet space topologies. If necessary we equip the spaces $\Psi^m(\Sigma)$, $C_{(b)}^{\infty}(\mathbb{R}; \Psi^m(\Sigma))$ and $S^{\delta}(\mathbb{R}; \Psi^m(\Sigma))$ with the quotient topology obtained from the map:

$$(c, R) \mapsto \text{Op}(c) + R$$

between the appropriate spaces.

If $a \in \Psi^m(\Sigma)$, the *principal symbol* $\sigma_{\text{pr}}(a) \in S_{\text{h}}^m(T^*\Sigma)$ is defined in analogy to the case $\Sigma = \mathbb{R}^d$. The operator a is *elliptic* if there exists $C > 0$ such that

$$(3.2) \quad |\sigma_{\text{pr}}(a)| \geq C|k|^m, \quad |k| \geq 1,$$

uniformly in the chart open sets. If $a \in C^{\infty}(\mathbb{R}; \Psi^m(\Sigma))$ we say that a is *elliptic* if $a(t)$ is elliptic for all $t \in \mathbb{R}$ and the constant C in (3.2) is locally uniform in t . For $a \in C_{(b)}^{\infty}(\mathbb{R}; \Psi^m(\Sigma))$ or $S^0(\mathbb{R}; \Psi^m(\Sigma))$ there is also a corresponding notion of ellipticity, where we require C to be uniform in t .

As shown in [GOW], the pseudodifferential classes $\Psi^m(\Sigma)$ fit into the general framework of Ammann, Lauter, Nistor and Vasy [ALNV], and consequently they have many convenient properties that generalize well-known facts for say, pseudodifferential operators on closed manifolds, such as the existence of complex powers for elliptic, bounded from below operators.

We state below a particular case of *Seeley's theorem* for real powers, partly proved in [GOW, Sect. 5] and based on a general result from [ALNV].

Theorem 3.7 (Seeley's theorem). *Let $a \in C^{\infty}(\mathbb{R}; \Psi^m(\Sigma))$ be elliptic, selfadjoint with $a(t) \geq c(t)\mathbf{1}$, $c(t) > 0$. Then $a^{\alpha} \in C^{\infty}(\mathbb{R}; \Psi^{m\alpha}(\Sigma))$ for any $\alpha \in \mathbb{R}$ and $\sigma_{\text{pr}}(a^{\alpha})(t) = \sigma_{\text{pr}}(a(t))^{\alpha}$.*

The same result holds replacing $C^{\infty}(\mathbb{R}; \Psi^m(\Sigma))$ by $C_{(b)}^{\infty}(\mathbb{R}; \Psi^m(\Sigma))$ or also by $S^0(\mathbb{R}; \Psi^m(\Sigma))$ if one assumes $a(t) \geq c_0\mathbf{1}$ for $c_0 > 0$.

Proof. The $C_{(b)}^{\infty}$ cases are proved in [GOW, Thm. 5.12], by checking that the general framework of [ALNV] applies to these two situations. The S^0 case can be proved similarly. The only point deserving special care is the *spectral invariance* of the ideal $S^0(\mathbb{R}; \mathcal{W}^{-\infty}(\Sigma))$, which we explain in some detail. Let $r_{-\infty} \in S^0(\mathbb{R}; \mathcal{W}^{-\infty}(\Sigma))$, considered as a bounded operator on $L^2(\mathbb{R}_t \times \Sigma_x)$. The spectral invariance property is the fact that if $\mathbf{1} - r_{-\infty}$ is invertible in $B(L^2(\mathbb{R}_t \times \Sigma_x))$ then $(\mathbf{1} - r_{-\infty})^{-1} = \mathbf{1} - r_{1,-\infty}$ for $r_{1,-\infty} \in S^0(\mathbb{R}; \mathcal{W}^{-\infty}(\Sigma))$. This can be however proved exactly as in [GOW, Lemma 5.5]. \square

3.6. Egorov's theorem. If $b(t) \in C^{\infty}(\mathbb{R}; \Psi^m(\Sigma))$ (or more generally, if $b(t)$ is a square matrix consisting of elements of $C^{\infty}(\mathbb{R}; \Psi^m(\Sigma))$ and $H^{-\infty}(\Sigma)$ is tensorized by powers of \mathbb{C} accordingly) we denote by

$$\mathcal{U}_b(t, s) : H^{-\infty}(\Sigma) \rightarrow H^{-\infty}(\Sigma)$$

the evolution generated by $b(t)$, i.e. the Cauchy evolution of $\bar{\partial}_t - ib(t)$, or put in other words, the unique solution (if it exists) of the system

$$(3.3) \quad \begin{cases} \frac{\partial}{\partial t} \mathcal{U}_b(t, s) = ib(t) \mathcal{U}_b(t, s), \\ \frac{\partial}{\partial s} \mathcal{U}_b(t, s) = -i \mathcal{U}_b(t, s) b(s), \\ \mathcal{U}_b(t, s) = \mathbf{1}. \end{cases}$$

The existence of $\mathcal{U}_b(t, s)$ can typically be established if $b(t)$ defines a differentiable family of self-adjoint operators on a Hilbert space, or a small perturbation of such

family. Specifically, consider $b(t) \in C^\infty(\mathbb{R}; \Psi^1(\Sigma))$ such that $b(t) = b_1(t) + b_0(t)$ with $b_i(t) \in C^\infty(\mathbb{R}; \Psi^i(\Sigma))$ and:

(E) $b_1(t)$ is elliptic and bounded from below on $H^\infty(\Sigma)$, locally uniformly in t .

Using [ALNV, Prop. 2.2] it follows that $b(t)$ is closed with domain $\text{Dom } b(t) = H^1(\Sigma)$. Moreover the map $\mathbb{R} \ni t \mapsto b(t) \in B(H^1(\Sigma), L^2(\Sigma))$ is norm continuous. It follows that we can define $\mathcal{U}_b(t, s)$, using for instance [RS, Thm. X.70]. In the present setup one can prove a result known generally as Egorov's theorem, we refer to [GOW] for the details and proofs.

Lemma 3.8. *Assume (E). Then:*

- (1) $\mathcal{U}_b(t, s) \in B(H^m(\Sigma))$ for $m \in \mathbb{R}$ or $m = \pm\infty$.
- (2) if $r \in \mathcal{W}^{-\infty}(\Sigma)$ then $\mathcal{U}_b(t, s)r, r\mathcal{U}_b(s, t) \in C^\infty(\mathbb{R}_{t,s}^2, \mathcal{W}^{-\infty}(\Sigma))$.
- (3) if moreover $b(t) \in S^0(\mathbb{R}; \Psi^1(\Sigma))$ and $b(t) - b^*(t) \in S^{-1-\delta}(\mathbb{R}; \Psi^0(\Sigma))$ for $\delta > 0$ then $\mathcal{U}_b(t, s)$ is uniformly bounded in $B(L^2(\Sigma))$.

Theorem 3.9 (Egorov's theorem). *Let $c \in \Psi^m(\Sigma)$ and $b(t)$ satisfying (E). Then*

$$c(t, s) := \mathcal{U}_b(t, s)c\mathcal{U}_b(s, t) \in C^\infty(\mathbb{R}_{t,s}^2, \Psi^m(\Sigma)).$$

Moreover

$$\sigma_{\text{pr}}(c)(t, s) = \sigma_{\text{pr}}(c) \circ \Phi(s, t),$$

where $\Phi(t, s) : T^*\Sigma \rightarrow T^*\Sigma$ is the flow of the time-dependent Hamiltonian $\sigma_{\text{pr}}(b)(t)$.

3.7. Some auxiliary results. For the sake of having a slightly more short-hand notation, for (Σ, h) of bounded geometry we set:

$$\Psi_{\text{td}}^{m, \delta}(\mathbb{R}; \Sigma) := S^\delta(\mathbb{R}; \Psi^m(\Sigma)).$$

for pseudodifferential operator classes with time decay (td) of the symbols.

3.7.1. Difference of fractional powers. We now state an auxiliary result about fractional powers of elliptic operators that will be needed later on.

Proposition 3.10. *Let $a_i \in \Psi_{\text{td}}^{2,0}(\mathbb{R}; \Sigma)$, $i = 1, 2$ elliptic with $a_i = a_i^*$ and $a_i(t) \geq c_0 \mathbf{1}$ for some $c_0 > 0$. Assume that $a_1 - a_2 \in \Psi_{\text{td}}^{2, -\delta}(\mathbb{R}; \Sigma)$ with $\delta > 0$. Then for each $\alpha \in \mathbb{R}$ one has:*

$$a_1^\alpha - a_2^\alpha \in \Psi_{\text{td}}^{2\alpha, -\delta}(\mathbb{R}; \Sigma).$$

Prop. 3.10 is proved in Subsect. A.1.

3.7.2. Resummation of symbols. We now examine the resummation of symbols. One can think of this as a statement about the uniform symbol classes on \mathbb{R}^d , after applying a chart diffeomorphism.

We denote $\Psi_{\text{td}}^{-\infty, -\delta}(\mathbb{R}; \Sigma) := \bigcap_{m \in \mathbb{R}} \Psi_{\text{td}}^{m, -\delta}(\mathbb{R}; \Sigma)$.

Lemma 3.11. *Let $\delta \in \mathbb{R}$ and let (m_j) be a real sequence decreasing to $-\infty$. Then if $a_j \in \Psi_{\text{td}}^{m_j, -\delta}(\mathbb{R}; \Sigma)$ there exists $a \in \Psi_{\text{td}}^{m_0, -\delta}(\mathbb{R}; \Sigma)$, unique modulo $\Psi_{\text{td}}^{-\infty, -\delta}(\mathbb{R}; \Sigma)$, such that*

$$a \sim \sum_{j=0}^{\infty} a_j, \text{ i.e. } a - \sum_{j=0}^N a_j \in \Psi_{\text{td}}^{m_{N+1}, -\delta}(\mathbb{R}; \Sigma), \forall N \in \mathbb{N}.$$

Proof. By introducing the new variable $s = \int_0^t \langle \sigma \rangle^{-1} d\sigma$ (so that $\langle t \rangle \partial_t = \partial_s$) and putting the extra variable s together with the x variables we can reduce ourselves to the situation covered by the standard proof (see e.g. [Sh2, Prop. 3.5]). \square

4. PARAMETRIX FOR THE CAUCHY EVOLUTION AND HADAMARD STATES

4.1. Model Klein-Gordon equation. We fix a d -dimensional manifold Σ equipped with a reference Riemannian metric k such that (Σ, k) is of bounded geometry. We equip $M = \mathbb{R} \times \Sigma$, the elements of which are denoted by $x = (t, \mathbf{x})$, with a Lorentzian metric g and a real function V such that:

$$(4.1) \quad \begin{aligned} g &= -dt^2 + h_{ij}(t, \mathbf{x}) dx^i dx^j, \\ h &\in C^\infty(\mathbb{R}; \text{BT}_2^0(\Sigma, k)), \quad h^{-1} \in C^\infty(\mathbb{R}; \text{BT}_0^2(\Sigma, k)), \\ V &\in C^\infty(\mathbb{R}; \text{BT}_0^0(\Sigma, k)). \end{aligned}$$

Although the first assumption may look restrictive, we will give in Subsect. 5.2 a reduction procedure that will allow us to treat more general cases.

In this setup, the Klein-Gordon operator $P = -\square_g + V$ equals

$$(4.2) \quad \begin{aligned} P &= |h|^{-\frac{1}{2}} \bar{\partial}_t |h|^{\frac{1}{2}} \bar{\partial}_t - |h|^{-\frac{1}{2}} \bar{\partial}_i h^{ij} |h|^{\frac{1}{2}} \bar{\partial}_j + V \\ &= \bar{\partial}_t^2 + r(t, \mathbf{x}) \bar{\partial}_t + a(t, \mathbf{x}, \bar{\partial}_\mathbf{x}), \end{aligned}$$

where

$$a(t, \mathbf{x}, \bar{\partial}_\mathbf{x}) = -|h|^{-\frac{1}{2}} \bar{\partial}_i h^{ij} |h|^{\frac{1}{2}} \bar{\partial}_j + V(t, \mathbf{x})$$

is formally self-adjoint with respect to the t -dependent $L^2(\Sigma, |h|^{\frac{1}{2}} dx)$ -inner product and

$$r(t, \mathbf{x}) = |h|^{-\frac{1}{2}} \partial_t (|h|^{\frac{1}{2}})(t, \mathbf{x}).$$

Note that the above function is closely related to the extrinsic curvature of Σ in M .

In the sequel we will often abbreviate $a(t, \mathbf{x}, \bar{\partial}_\mathbf{x})$ by $a(t)$ or a .

4.2. Construction of parametrix. Following [GOW] we now explain how one obtains a parametrix for the Cauchy evolution for the model Klein-Gordon operator (and a splitting of it) by means of an approximate time-dependent diagonalization. We will then adapt it to the setup of scattering theory.

The first step consists of observing that the Klein-Gordon equation $(\partial_t^2 + r(t)\partial_t + a(t))\phi(t) = 0$ is equivalent to

$$(4.3) \quad \mathbf{i}^{-1} \partial_t \psi(t) = H(t) \psi(t), \quad \text{where } H(t) = \begin{pmatrix} 0 & \mathbf{1} \\ a(t) & \mathbf{i}r(t) \end{pmatrix},$$

by setting

$$(4.4) \quad \psi(t) = \begin{pmatrix} \phi(t) \\ \mathbf{i}^{-1} \partial_t \phi(t) \end{pmatrix}.$$

Let us denote by $\mathcal{U}(s, t)$ the evolution generated by $H(t)$, cf. (3.3). Recall that on Cauchy data on $\Sigma_s = \{s\} \times \Sigma$, we have a symplectic form induced from an operator $G(s)$, defined by:

$$G = (\varrho_s G)^* \circ G(s) \circ (\varrho_s G).$$

Here the formal adjoint will always be taken with respect to the $L^2(\Sigma, |h|^{\frac{1}{2}} dx)$ -inner product. We have also introduced the hermitian operator $q(s) = \mathbf{i}G(s)$. It is well known that with these choices, $q(s)$ equals specifically

$$(4.5) \quad q(s) = \begin{pmatrix} 0 & \mathbf{1} \\ \mathbf{1} & 0 \end{pmatrix},$$

in particular it does not depend on s (we will thus simply write q instead). Furthermore,

$$(4.6) \quad \mathcal{U}^*(t, s) q \mathcal{U}(t, s) = q,$$

(the Cauchy evolution is symplectic).

4.2.1. *Riccati equation.* The approximate diagonalization of $\mathcal{U}(s, t)$ will be based on solving the Riccati equation

$$(4.7) \quad i\partial_t b - b^2 + a + irb = 0,$$

modulo smoothing terms, where the unknown is $b(t) \in C^\infty(\mathbb{R}; \Psi^1(\Sigma))$. By repeating the arguments in [GW1, GW2] this can be solved modulo terms in $C^\infty(\mathbb{R}; \mathcal{W}^{-\infty}(\Sigma))$. Concretely, supposing for the moment that $a(t) \geq c(t)\mathbf{1}$ for $c(t) > 0$, upon setting $\epsilon = a^{\frac{1}{2}}$, $b = \epsilon + b_0$ one obtains the equations:

$$b_0 = \frac{i}{2}(\epsilon^{-1}\partial_t\epsilon + \epsilon^{-1}r\epsilon) + F(b_0),$$

$$F(b_0) = \frac{1}{2}\epsilon^{-1}(i\partial_t b_0 + [\epsilon, b_0] + irb_0 - b_0^2).$$

These can be solved by substituting a poly-homogeneous expansion of the symbol of b_0 , yielding an approximate solution of (4.7) in the sense that

$$(4.8) \quad i\partial_t b - b^2 + a + irb = r_{-\infty} \in C^\infty(\mathbb{R}; \mathcal{W}^{-\infty}(\Sigma)).$$

Set

$$(4.9) \quad b^+ = b, \quad b^- = -b^*.$$

Taking the adjoint of both sides of (4.8) with respect to the t -dependent inner product $L^2(\Sigma, |h|^{\frac{1}{2}}dx)$ and using that

$$(\partial_t b)^* = \partial_t(b^*) + rb^* - b^*r,$$

we obtain

$$(4.10) \quad i\partial_t b^\pm - b^{\pm 2} + a + irb^\pm = r_{-\infty}^\pm,$$

with $r_{-\infty}^+ = r_{-\infty}$, $r_{-\infty}^- = r_{-\infty}^* \in C^\infty(\mathbb{R}; \mathcal{W}^{-\infty}(\Sigma))$.

In general we can, using the locally uniform ellipticity of $a(t)$, find a cutoff function $\varphi \in C_c^\infty(\mathbb{R})$ such that $a(t) + \varphi(a(t)) \geq c(t)\mathbf{1}$ for $c(t)$ as above. Since $\varphi(a(t))$ is a smoothing operator, replacing $a(t)$ by $a(t) + \varphi(a(t))$ is a harmless modification.

A redefinition of $b(t)$ involving a cutoff in low frequencies as in [GW2, GOW] gives then control of the norm sufficient to obtain in addition

$$(4.11) \quad (b^+(t) - b^-(t))^{-1} \geq C(t)\epsilon(t)^{-1}$$

for some $C(t) > 0$, while keeping the property that $b^\pm(t) = \pm\epsilon(t) + C^\infty(\mathbb{R}; \Psi^0(\Sigma))$, and with (4.10) still valid for some $r_{-\infty}^\pm \in C^\infty(\mathbb{R}; \mathcal{W}^{-\infty}(\Sigma))$.

Observe now that the Riccati equation (4.10) implies the following approximate factorization of the Klein-Gordon operator:

$$(4.12) \quad (\bar{\partial}_t + ib^\pm(t) + r(t)) \circ (\bar{\partial}_t - ib^\pm(t)) = \bar{\partial}_t^2 + r\bar{\partial}_t + a - r_{-\infty}^\pm.$$

The benefits of having such a factorization were already recognized by Junker and Schrohe [Ju, JS] in the context of Hadamard states (although it was obtained only in the special case of FRLW spacetimes with compact Cauchy hypersurface). Here we use (4.12) to diagonalize (4.3) by setting

$$\tilde{\psi}(t) := \begin{pmatrix} \bar{\partial}_t - ib^-(t) \\ \bar{\partial}_t - ib^+(t) \end{pmatrix} \phi(t).$$

A direct computation yields then $\tilde{\psi}(t) = S^{-1}(t)\psi(t)$ with

$$(4.13) \quad S^{-1}(t) = i \begin{pmatrix} -b^-(t) & \mathbf{1} \\ -b^+(t) & \mathbf{1} \end{pmatrix}, \quad S(t) = i^{-1} \begin{pmatrix} \mathbf{1} & -\mathbf{1} \\ b^+(t) & -b^-(t) \end{pmatrix} (b^+(t) - b^-(t))^{-1},$$

where well-definiteness and invertibility of $S(t)$ rely on the fact that $b^+(t) - b^-(t)$ is invertible by (4.11). We obtain from (4.12) that

$$\begin{aligned} \begin{pmatrix} \bar{\partial}_t + ib^- + r & 0 \\ 0 & \bar{\partial}_t + ib^+ + r \end{pmatrix} \tilde{\psi}(t) &= \begin{pmatrix} \bar{\partial}_t^2 + a + r\bar{\partial}_t - r_{-\infty}^- \\ \bar{\partial}_t^2 + a + r\bar{\partial}_t - r_{-\infty}^+ \end{pmatrix} \phi(t) \\ &= \begin{pmatrix} r_{-\infty}^- & 0 \\ r_{-\infty}^+ & 0 \end{pmatrix} S(t) \tilde{\psi}(t) = i^{-1} \begin{pmatrix} r_{-\infty}^- & -r_{-\infty}^- \\ r_{-\infty}^+ & -r_{-\infty}^+ \end{pmatrix} (b^+ - b^-)^{-1} \tilde{\psi}(t). \end{aligned}$$

Therefore, $\tilde{\psi}(t)$ solves a diagonal matrix equation modulo smooth terms. More precisely, we have $\tilde{\psi}(t) = \mathcal{U}_B(t, s) \tilde{\psi}(s)$ for

$$(4.14) \quad B(t) = \tilde{B}(t) + R_{-\infty}(t),$$

$$(4.15)$$

$$\tilde{B}(t) = \begin{pmatrix} -b^- + ir & 0 \\ 0 & -b^+ + ir \end{pmatrix}, \quad R_{-\infty}(t) = - \begin{pmatrix} r_{-\infty}^- & -r_{-\infty}^- \\ r_{-\infty}^+ & -r_{-\infty}^+ \end{pmatrix} (b^+ - b^-)^{-1},$$

Ultimately, we can thus conclude that

$$(4.16) \quad \mathcal{U}(t, s) = S(t) \mathcal{U}_B(t, s) S(s)^{-1}.$$

4.3. Improved approximate diagonalization. It is convenient to modify $S(t)$ to obtain a simple formula for the symplectic form $S^*(t)q(t)S(t)$ preserved by the almost diagonalized evolution. Namely, setting

$$(4.17) \quad \begin{aligned} T(t) &:= S(t)(b^+ - b^-)^{\frac{1}{2}}(t) = i^{-1} \begin{pmatrix} \mathbf{1} & -\mathbf{1} \\ b^+ & -b^- \end{pmatrix} (b^+ - b^-)^{-\frac{1}{2}}, \\ T^{-1}(t) &= i(b^+ - b^-)^{-\frac{1}{2}} \begin{pmatrix} -b^- & \mathbf{1} \\ -b^+ & \mathbf{1} \end{pmatrix}, \end{aligned}$$

we find that for q defined in (4.5) one has:

$$T^*(t)qT(t) = \begin{pmatrix} \mathbf{1} & 0 \\ 0 & -\mathbf{1} \end{pmatrix} =: q^{\text{ad}}.$$

We now define

$$(4.18) \quad \mathcal{U}(t, s) =: T(t) \mathcal{U}^{\text{ad}}(t, s) T(s)^{-1},$$

and we obtain that $\mathcal{U}^{\text{ad}}(t, s)^* q^{\text{ad}} \mathcal{U}^{\text{ad}}(t, s) = q^{\text{ad}}$. Furthermore, the generator of the evolution group $\{\mathcal{U}^{\text{ad}}(t, s)\}_{t, s \in \mathbb{R}}$ is:

$$(4.19)$$

$$\begin{aligned} H^{\text{ad}}(t) &= (b^+ - b^-)^{-\frac{1}{2}} B(t) (b^+ - b^-)^{\frac{1}{2}} - i\partial_t (b^+ - b^-)^{-\frac{1}{2}} (b^+ - b^-)^{\frac{1}{2}} \\ &= \begin{pmatrix} -b^- + r_b^- & 0 \\ 0 & -b^+ + r_b^+ \end{pmatrix} - (b^+ - b^-)^{-\frac{1}{2}} \begin{pmatrix} r_{-\infty}^- & -r_{-\infty}^- \\ r_{-\infty}^+ & -r_{-\infty}^+ \end{pmatrix} (b^+ - b^-)^{-\frac{1}{2}}, \end{aligned}$$

where $r_{-\infty}^{\pm} \in C^\infty(\mathbb{R}; \mathcal{W}^{-\infty}(\Sigma))$ are the remainder terms from (4.10), and

$$(4.20) \quad r_b^\pm = ir + [(b^+ - b^-)^{-\frac{1}{2}}, b^\pm] - i\partial_t (b^+ - b^-)^{-\frac{1}{2}} (b^+ - b^-)^{\frac{1}{2}} \in \Psi^0(\Sigma).$$

This way, denoting by H^{d} the diagonal part of $H^{\text{ad}}(t)$, using that $H^{\text{ad}}(t)^* q^{\text{ad}} = q^{\text{ad}} H^{\text{ad}}(t)$ we have:

$$H^{\text{d}}(t) = H^{\text{d}*}(t), \quad H^{\text{d}}(t) = \begin{pmatrix} \epsilon^+(t) & 0 \\ 0 & \epsilon^-(t) \end{pmatrix},$$

where

$$\epsilon^\pm = -b^\mp + r_b^\mp + C^\infty(\mathbb{R}; \mathcal{W}^{-\infty}(\Sigma)),$$

and $H^{\text{ad}}(t) = H^{\text{d}}(t) + V_{-\infty}^{\text{ad}}(t)$, where $V_{-\infty}^{\text{ad}}(t) \in C^\infty(\mathbb{R}; \mathcal{W}^{-\infty}(\Sigma) \otimes B(\mathbb{C}^2))$. The evolution $\mathcal{U}^{\text{d}}(t, s)$ generated by $H^{\text{d}}(t)$ is diagonal, in fact:

$$(4.21) \quad \mathcal{U}^{\text{d}}(t, s) = \begin{pmatrix} \mathcal{U}_{\epsilon^+}(t, s) & 0 \\ 0 & \mathcal{U}_{\epsilon^-}(t, s) \end{pmatrix}.$$

Moreover:

$$(4.22) \quad \begin{aligned} \mathcal{U}(t, s) &= T(t)\mathcal{U}^{\text{ad}}(t, s)T(s)^{-1} \\ &= T(t)\mathcal{U}^{\text{d}}(t, s)T(s)^{-1} + C^\infty(\mathbb{R}^2; \mathcal{W}^{-\infty}(\Sigma)). \end{aligned}$$

This is shown by an ‘interaction picture’ argument explained in detail in [GOW]; we omit the proof here.

Remark 4.1. *One easily sees that $S(t)$ is an isomorphism from $L^2(\Sigma) \oplus L^2(\Sigma)$ to $H^1(\Sigma) \oplus L^2(\Sigma)$ (the so-called energy space of Cauchy data of (4.3)), while $T(t)$ is an isomorphism from $L^2(\Sigma) \oplus L^2(\Sigma)$ to $H^{\frac{1}{2}}(\Sigma) \oplus H^{-\frac{1}{2}}(\Sigma)$ (this is the charge space that appears naturally in the quantization of the Klein-Gordon equation).*

4.4. Splitting of the parametrix and of the Cauchy evolution. Let us set

$$(4.23) \quad \pi^+ = \begin{pmatrix} \mathbf{1} & 0 \\ 0 & 0 \end{pmatrix}, \quad \pi^- = \begin{pmatrix} 0 & 0 \\ 0 & \mathbf{1} \end{pmatrix}.$$

Since $\mathcal{U}^{\text{d}}(t, s)$ is diagonal we have:

$$\mathcal{U}^{\text{d}}(t, s) = \mathcal{U}^{\text{d}}(t, s)\pi^+ + \mathcal{U}^{\text{d}}(t, s)\pi^-,$$

with $\mathcal{U}^{\text{d}}(t, s)\pi^\pm$ propagating with wave front set contained in \mathcal{N}^\pm (this follows from b^\pm being $\pm\epsilon$ modulo terms of lower order). This suggests that at least modulo smoothing terms, the splitting of $\mathcal{U}(t, s)$ at time s should be given by a pair of operators $c_{\text{ref}}^\pm(s)$ defined as follows. We first fix a reference time $t_0 \in \mathbb{R}$.

Definition 4.2. *We set:*

$$c_{\text{ref}}^\pm(t_0) := T(t_0)\pi^\pm T^{-1}(t_0) = \begin{pmatrix} \mp(b^+ - b^-)^{-1}b^\mp & \pm(b^+ - b^-)^{-1} \\ \mp b^\pm(b^+ - b^-)^{-1}b^- & \pm b^\pm(b^+ - b^-)^{-1} \end{pmatrix}(t_0).$$

Then $c_{\text{ref}}^\pm(t_0)$ is a 2×2 matrix of pseudodifferential operators and

$$c_{\text{ref}}^\pm(t_0)^2 = c_{\text{ref}}^\pm(t_0), \quad c_{\text{ref}}^+(t_0) + c_{\text{ref}}^-(t_0) = \mathbf{1}.$$

We set:

$$(4.24) \quad \mathcal{U}^\pm(t, s) := \mathcal{U}(t, t_0)c_{\text{ref}}^\pm(t_0)\mathcal{U}(t_0, s),$$

so that

$$(4.25) \quad \mathcal{U}(t, s) = \mathcal{U}^+(t, s) + \mathcal{U}^-(t, s).$$

This splitting has the following properties (see [GOW]):

Proposition 4.3.

$$(4.26) \quad \begin{aligned} i) & \mathcal{U}^\pm(t, s)\mathcal{U}^\pm(s, t') = \mathcal{U}^\pm(t, t'), \\ ii) & (\bar{\partial}_t - iH(t))\mathcal{U}^\pm(t, s) = \mathcal{U}^\pm(t, s)(\bar{\partial}_s + iH(s)) = 0, \\ iii) & \text{WF}(\mathcal{U}^\pm(t, s))' = \{(X, X') \in T^*\Sigma \times T^*\Sigma : X = \Phi^\pm(t, s)(X')\}, \end{aligned}$$

where $\Phi^\pm(t, s) : T^*\Sigma \rightarrow T^*\Sigma$ is the symplectic flow generated by the time-dependent Hamiltonian $\pm(h^{ij}(t, x)k_i k_j)^{\frac{1}{2}}$.

If we set for $t \in \mathbb{R}$:

$$(4.27) \quad \mathcal{U}^\pm(t, t) =: c_{\text{ref}}^\pm(t) = \mathcal{U}(t, t_0) c_{\text{ref}}^\pm(t_0) \mathcal{U}(t_0, t),$$

then

$$c_{\text{ref}}^\pm(t)^2 = c_{\text{ref}}^\pm(t), \quad c_{\text{ref}}^+(t) + c_{\text{ref}}^-(t) = \mathbf{1}, \quad c_{\text{ref}}^\pm(t) = \mathcal{U}(t, s) c_{\text{ref}}^\pm(s) \mathcal{U}(s, t).$$

As a consequence, one gets that $c_{\text{ref}}^\pm(t)$ are the time- t covariances of a Hadamard state [GOW]. In general, we say that a state is a *regular Hadamard state* if its time- t covariances differ from $c_{\text{ref}}^\pm(t)$ by terms in $\mathcal{W}^{-\infty}(\Sigma) \otimes B(\mathbb{C}^2)$, and one can show that it suffices to check that property for one value of t [GOW]. In summary:

Theorem 4.4 ([GOW]). *The pair of operators $c_{\text{ref}}^\pm(t)$ defined in (4.27) are the covariances of a pure, regular Hadamard state.*

We stress that in general $c_{\text{ref}}^\pm(t)$ are not ‘canonical’ nor ‘distinguished’, because they depend on the choice of the reference time t_0 and on the precise choice of the operators $b^\pm(t)$ (to which one can always add suitable regularizing terms). On the other hand, in Sect. 4.5 we will construct covariances $c_{\text{in}}^\pm(t)$ and $c_{\text{out}}^\pm(t)$ of the distinguished *in* and *out* states, and the operators $c_{\text{ref}}^\pm(t)$ will play an important role in the proof of their Hadamard property: a suitable sufficient condition for that is in fact that

$$(4.28) \quad c_{\text{out/in}}^\pm(t) - c_{\text{ref}}^\pm(t) \in \mathcal{W}^{-\infty}(\Sigma) \otimes B(\mathbb{C}^2)$$

for some (and hence all) $t \in \mathbb{R}$.

4.5. Further estimates in scattering settings. In what follows we give a refinement of the constructions in Sect. 4 for the model Klein-Gordon equation in a scattering situation, corresponding to a situation when the metric g , resp. the potential V converge to ultra-static metrics $g_{\text{out/in}} = -dt^2 + h_{\text{out/in},ij}(\mathbf{x}) dx^i dx^j$, resp. to time-independent potentials $V_{\text{out/in}}$ as $t \rightarrow \pm\infty$. We start by fixing two classes of assumptions on the model Klein-Gordon equation (4.2).

We will often abbreviate the classes $\Psi_{\text{td}}^{m,\delta}(\mathbb{R}; \Sigma)$ (introduced in Subsect. 3.5-3.7) by $\Psi_{\text{td}}^{m,\delta}$. We make the following assumption:

$$(td) \quad \begin{aligned} a(t, \mathbf{x}, D_{\mathbf{x}}) &= a_{\text{out/in}}(\mathbf{x}, D_{\mathbf{x}}) + \Psi_{\text{td}}^{2,-\delta}(\mathbb{R}; \Sigma) \text{ on } \mathbb{R}^\pm \times \Sigma, \quad \delta > 0, \\ r(t) &\in \Psi_{\text{td}}^{0,-1-\delta}(\mathbb{R}; \Sigma), \\ a_{\text{out/in}}(\mathbf{x}, D_{\mathbf{x}}) &\in \Psi^2(\Sigma) \text{ elliptic, } a_{\text{out/in}}(\mathbf{x}, D_{\mathbf{x}}) = a_{\text{out/in}}(\mathbf{x}, D_{\mathbf{x}})^* \geq C_\infty > 0. \end{aligned}$$

The assumption $\delta > 0$ corresponds to a *long-range* setup (as opposed to the more narrow *short-range* case when $\delta > 1$).

Below, we give estimates on the solution of the Riccati equation, taking now into account the decay in time that follows from (td). To simplify notation we simply write $a_1(t) = a_2(t) + \Psi_{\text{td}}^{m,\delta}(\mathbb{R}^\pm; \Sigma)$ when $a_1(t) = a_2(t) + \Psi_{\text{td}}^{m,\delta}(\mathbb{R}; \Sigma)$ in $\mathbb{R}^\pm \times \Sigma$. We also abbreviate $\Psi_{\text{td}}^{m,\delta}(\mathbb{R}^\pm; \Sigma)$ by $\Psi_{\text{td}}^{m,\delta}$ when it is clear from the context whether the future or past case is meant.

From hypothesis (td) we deduce that there exists $c(t) \in C_c^\infty(\mathbb{R})$ such that $a(t) + c(t)\mathbf{1} \sim a_{\text{out/in}}$, uniformly in $t \in \mathbb{R}^\pm$. By functional calculus we can find $\varphi \in C_c^\infty(\mathbb{R})$ such that $a(t) + \varphi(\frac{a(t)}{c(t)}) \sim a_{\text{out/in}}$, uniformly in $t \in \mathbb{R}^\pm$. The error term $\varphi(\frac{a(t)}{c(t)})$ belongs to $C_c^\infty(\mathbb{R}; \mathcal{W}^{-\infty}(\Sigma))$.

We can hence replace $a(t)$ by $a(t) + \varphi(\frac{a(t)}{c(t)})$ in the Riccati equation (4.7) and assume that

$$a(t) \sim a_{\text{out/in}} \text{ uniformly in } t \in \mathbb{R}^\pm.$$

If $\epsilon_{\text{out/in}} := a_{\text{out/in}}^{\frac{1}{2}}$, then from Prop. 3.10 we deduce that if (td) holds then

$$(4.1) \quad \epsilon(t) := a(t)^{\frac{1}{2}} = \epsilon_{\text{out/in}} + \Psi_{\text{td}}^{1,-\delta}.$$

Proposition 4.5. *There exists $b(t) = \epsilon(t) + \Psi_{\text{td}}^{0,-1-\delta}(\mathbb{R}; \Sigma) = \epsilon_{\text{out/in}} + \Psi_{\text{td}}^{1,-\delta}(\mathbb{R}^{\pm}; \Sigma)$ that solves*

$$i\partial_t b - b^2 + a + irb \in \Psi_{\text{td}}^{-\infty,-1-\delta}(\mathbb{R}; \Sigma).$$

The proof is given in Appendix A.2.

Proposition 4.6. *Assume (td) and let r_b^{\pm} be defined in (4.20) and $r_{-\infty}^{\pm}$ in (4.10). Then*

$$r_b^{\pm} \in \Psi_{\text{td}}^{0,-1-\delta}(\mathbb{R}; \Sigma), \quad r_{-\infty}^{\pm} \in \Psi_{\text{td}}^{-\infty,-1-\delta}(\mathbb{R}; \Sigma).$$

The proof is given in Appendix A.3.

5. THE out/in STATES ON ASYMPTOTICALLY STATIC SPACETIMES

5.1. Assumptions. In what follows we introduce a class of asymptotically static spacetimes on which we will construct the *in* and *out* states and prove their Hadamard property. One of the key ingredients is the reduction to a model Klein-Gordon operator that satisfies the assumptions (td) considered in Subsect. 4.5.

We will use the framework of manifolds and diffeomorphisms of bounded geometry introduced in Defs. 3.2, 3.3.

We fix a d -dimensional manifold Σ equipped with a reference Riemannian metric k such that (Σ, k) is of bounded geometry, and consider $M = \mathbb{R}_t \times \Sigma_y$, setting $y = (t, y)$, $n = 1 + d$. We equip M with a Lorentzian metric g of the form

$$(5.1) \quad g = -c^2(y)dt^2 + (dy^i + b^i(y)dt)h_{ij}(y)(dy^j + b^j(y)dt),$$

where we assume:

$$\begin{aligned} h_{ij} &\in C_b^\infty(\mathbb{R}; \text{BT}_2^0(\Sigma, k)), \quad h_{ij}^{-1} \in C_b^\infty(\mathbb{R}; \text{BT}_0^2(\Sigma, k)), \\ (\text{bg}) \quad b &\in C_b^\infty(\mathbb{R}; \text{BT}_0^1(\Sigma, k)), \\ c, \quad c^{-1} &\in C_b^\infty(\mathbb{R}; \text{BT}_0^0(\Sigma, k)). \end{aligned}$$

We recall that $\tilde{t} \in C^\infty(M)$ is called a *time function* if $\nabla \tilde{t}$ is a timelike vector field. It is called a *Cauchy time function* if its level sets are Cauchy hypersurfaces. By [CC, Thm. 2.1] we know that (M, g) is globally hyperbolic and t is a Cauchy time function.

We will consider the Klein-Gordon operator on (M, g) :

$$(5.2) \quad P = -\square_g + V,$$

with $V \in C_b^\infty(\mathbb{R}; \text{BT}_0^0(\Sigma, k))$ a smooth real-valued function. We consider two static metrics

$$g_{\text{out/in}} = -c_{\text{out/in}}^2(y)dt^2 + h_{\text{out/in}}(y)dy^2$$

and time-independent potentials $V_{\text{out/in}}$ and assume the following conditions:

$$\begin{aligned} (\text{ast}) \quad h(y) - h_{\text{out/in}}(y) &\in S^{-\mu}(\mathbb{R}^{\pm}; \text{BT}_2^0(\Sigma, k)), \\ b(y) &\in S^{-\mu'}(\mathbb{R}; \text{BT}_0^1(\Sigma, k)), \\ c(y) - c_{\text{out/in}}(y) &\in S^{-\mu}(\mathbb{R}^{\pm}; \text{BT}_0^0(\Sigma, k)), \\ V(y) - V_{\text{out/in}}(y) &\in S^{-\mu}(\mathbb{R}^{\pm}; \text{BT}_0^0(\Sigma, k)), \\ (\text{pos}) \quad \frac{n-2}{4(n-1)} (R_{c_{\text{out/in}}^{-2} h_{\text{out/in}}} - c_{\text{out/in}}^2 R_{g_{\text{out/in}}}) + c_{\text{out/in}}^2 V_{\text{out/in}} &\geq m^2, \end{aligned}$$

for some $\mu > 0$, $\mu' > 1$ and $m > 0$. Above, R_g , resp. R_h denotes the scalar curvature of g , resp. h .

Condition (ast) means that g , resp. V are asymptotic to the static metrics $g_{\text{out/in}}$, resp. to the time-independent potentials $V_{\text{out/in}}$ as $t \rightarrow \pm\infty$. Condition (pos) means that the asymptotic Klein-Gordon operators $\bar{\partial}_t^2 + a_{\text{out/in}}(x, \bar{\partial}_x)$ introduced in Lemma 5.2 below are *massive*.

It follows from (bg) that $h_{\text{out/in}} \in \text{BT}_2^0(\Sigma, k)$, $h_{\text{out/in}}^{-1} \in \text{BT}_0^2(\Sigma, k)$, and that $V_{\text{out/in}}, V_{\text{out/in}}^{-1} \in \text{BT}_0^0(\Sigma, k)$.

5.2. Reduction to the model case. In this subsection we perform the reduction of the Klein-Gordon operator P to the model case considered in Sect. 4.5. We start with the well-known orthogonal decomposition of g associated with the time function t . Namely, we set

$$v := \frac{g^{-1}dt}{dt \cdot g^{-1}dt} = \bar{\partial}_t + b^i \bar{\partial}_{y^i},$$

which using (bg) is a complete vector field. Furthermore, we denote by ϕ_t its flow, so that

$$\phi_t(x) = (t, y(t, 0, x)), \quad t \in \mathbb{R}, \quad x \in \Sigma,$$

where $y(t, s, \cdot)$ is the flow of the time-dependent vector field b on Σ . We also set

$$(5.3) \quad \chi : \mathbb{R} \times \Sigma \ni (t, x) \mapsto (t, y(t, 0, x)) \in \mathbb{R} \times \Sigma.$$

Lemma 5.1. *Assume (bg), (ast). Then*

$$\hat{g} := \chi^* g = -\hat{c}^2(t, x) dt^2 + \hat{h}(t, x) dy^2, \quad \chi^* V = \hat{V},$$

where:

$$\hat{c}, \hat{c}^{-1}, \hat{V} \in C_b^\infty(\mathbb{R}; \text{BT}_0^0(\Sigma, k)),$$

$$\hat{h} \in C_b^\infty(\mathbb{R}; \text{BT}_2^0(\Sigma, k)), \quad \hat{h}^{-1} \in C_b^\infty(\mathbb{R}; \text{BT}_0^2(\Sigma, k)).$$

Moreover there exist bounded diffeomorphisms $y_{\text{out/in}}$ of (Σ, k) such that if:

$$\hat{h}_{\text{out/in}} := y_{\text{out/in}}^* h_{\text{out/in}},$$

$$\hat{c}_{\text{out/in}} := y_{\text{out/in}}^* c_{\text{out/in}}, \quad \hat{V}_{\text{out/in}} := y_{\text{out/in}}^* V_{\text{out/in}},$$

then we have:

$$\hat{h}_{\text{out/in}} \in \text{BT}_2^0(\Sigma, k), \quad \hat{h}_{\text{out/in}}^{-1} \in \text{BT}_0^2(\Sigma, k),$$

$$\hat{c}_{\text{out/in}}, \hat{c}_{\text{out/in}}^{-1}, \hat{V}_{\text{out/in}} \in \text{BT}_0^0(\Sigma, k),$$

and furthermore,

$$\hat{h} - \hat{h}_{\text{out/in}} \in S^{-\min(1-\mu', \mu)}(\mathbb{R}^\pm, \text{BT}_2^0(\Sigma, k)),$$

$$\hat{c} - \hat{c}_{\text{out/in}} \in S^{-\min(1-\mu', \mu)}(\mathbb{R}^\pm, \text{BT}_0^0(\Sigma, k)),$$

$$\hat{V} - \hat{V}_{\text{out/in}} \in S^{-\mu}(\mathbb{R}^\pm, \text{BT}_0^0(\Sigma, k)).$$

Lemma 5.1 is proved in Appendix A.4.

Writing P as $-\square_g + \frac{n-2}{4(n-1)} R_g + W$ for $W = V - \frac{n-2}{4(n-1)} R_g$, and using the conformal invariance of $-\square_g + \frac{n-2}{4(n-1)} R_g$ and the estimates in Lemma 5.1, we obtain the following result, which completes the reduction to the model case.

If $P \in \text{Diff}(\mathbb{R} \times \Sigma)$ we denote by $\chi^* P$ the pullback of P by χ defined by $(\chi^* P)u \circ \chi = (Pu) \circ \chi$.

Lemma 5.2. *Assume (bg), (ast), (pos) and consider the Klein-Gordon operator P in (5.2). Let $\hat{h}, \hat{c}, \hat{V}$ be as in Lemma 5.1 and set:*

$$\hat{P} = \chi^* P, \quad \tilde{P} = \hat{c}^{1-n/2} \hat{P} \hat{c}^{1+n/2}, \quad \tilde{g} = \hat{c}^{-2} \hat{g}, \quad \tilde{h} = \hat{c}^{-2} \hat{h}.$$

Then

$$\tilde{P} = \bar{\partial}_t^2 + r(t, \mathbf{x}) \bar{\partial}_t + a(t, \mathbf{x}, \bar{\partial}_\mathbf{x}),$$

for

$$a(t, \mathbf{x}, \bar{\partial}_\mathbf{x}) = -\Delta_{\tilde{h}} + \tilde{V}, \quad r = |\tilde{h}_t|^{-\frac{1}{2}} \partial_t |\tilde{h}_t|^{\frac{1}{2}},$$

$$\tilde{V} = \frac{n-2}{4(n-1)} (R_{\tilde{g}} - \hat{c}^2 R_{\hat{g}}) + \hat{c}^2 \hat{V}.$$

Moreover a, r satisfy (td) with $\delta = \min(\mu, \mu' - 1)$ and

$$a_{\text{out/in}}(\mathbf{x}, \bar{\partial}_\mathbf{x}) = -\Delta_{\tilde{h}_{\text{out/in}}} + \tilde{V}_{\text{out/in}}(\mathbf{x}),$$

where

$$\tilde{V}_{\text{out/in}} = \left(\frac{n-2}{4(n-1)} (R_{c_{\text{out/in}}^{-2} h_{\text{out/in}}} - c_{\text{out/in}}^2 R_{g_{\text{out/in}}}) + c_{\text{out/in}}^2 V_{\text{out/in}} \right) \circ Y_{\text{out/in}}.$$

Note that condition (pos) simply means that $\tilde{V}_{\text{out/in}} \geq m^2 > 0$.

5.3. Cauchy evolutions. In this subsection we relate the Cauchy evolutions of P and of the model Klein-Gordon operator \tilde{P} .

The trace operator for P associated to the time function t is given by:

$$(5.4) \quad \varrho_t \phi = \begin{pmatrix} u(t, \cdot) \\ i^{-1} n \cdot \nabla \phi(t, \cdot) \end{pmatrix},$$

where n is the future directed unit normal to Σ_t . The corresponding trace operator for $\hat{P} = \chi^* P$ is:

$$\hat{\varrho}_t \phi = \begin{pmatrix} \phi(t, \cdot) \\ i^{-1} \hat{c}^{-1} \partial_t \phi(t, \cdot) \end{pmatrix},$$

so that denoting $\chi^* \phi = \phi \circ \chi$, we have:

$$\hat{\varrho}_t \chi^* \phi = \chi_t^* \varrho_t \phi \quad \text{for } \chi_t^* \begin{pmatrix} u_0 \\ u_1 \end{pmatrix} = \begin{pmatrix} u_0 \circ \chi_t \\ u_1 \circ \chi_t \end{pmatrix},$$

and $\chi_t(\mathbf{x}) = y(t, 0, \mathbf{x})$, see (5.3). Finally the trace operator for \tilde{P} as in Lemma 5.2 is

$$\tilde{\varrho}_t \phi = \begin{pmatrix} \phi(t, \cdot) \\ i^{-1} \phi(t, \cdot) \end{pmatrix}$$

so that if $\tilde{\phi} = \hat{c}^{n/2-1} \phi$ is the conformal transformation in Lemma 5.2, we have:

$$\tilde{\varrho}_t \tilde{\phi} = R(t) \hat{\varrho}_t \phi, \quad \text{for } R(t) = \hat{c}^{n/2-1} \begin{pmatrix} 1 & 0 \\ -i(n/2-1) \partial_t \ln(\hat{c}) & 1 \end{pmatrix}.$$

Let us denote by $\mathcal{U}(t, s)$ the Cauchy evolution for P associated to ϱ_t and by $\mathcal{U}^{\text{ad}}(t, s)$ the almost diagonal Cauchy evolution introduced in Subsect. 4.3 for the model Klein-Gordon operator \tilde{P} . The following lemma follows from the above computations and (4.18).

Lemma 5.3. *Let $Z(t) := (\chi_t^*)^{-1} R(t) T(t)$, where $T(t)$ is defined in (4.17). Then*

$$(5.5) \quad \mathcal{U}(t, s) = Z(t) \mathcal{U}^{\text{ad}}(t, s) Z^{-1}(s).$$

We have a similar reduction for the asymptotic Klein-Gordon operators:

$$P_{\text{out/in}} = -\square_{g_{\text{out/in}}} + \frac{n-2}{4(n-1)} R_{g_{\text{out/in}}} + V_{\text{out/in}},$$

for $g_{\text{out/in}} = -c_{\text{out/in}}^2(y)dt^2 + h_{\text{out/in}}(y)dy^2$, where $h_{\text{out/in}}$, $c_{\text{out/in}}$, $V_{\text{out/in}}$ were introduced in (ast). The associated trace operator is

$$\varrho_{t,\text{out/in}}\phi = \left(\begin{array}{c} \phi(t, \cdot) \\ i^{-1} c_{\text{out/in}}^{-1} \partial_t \phi(t, \cdot) \end{array} \right).$$

We also set

$$\chi_{\text{out/in}}^* \begin{pmatrix} u_0 \\ u_1 \end{pmatrix} = \begin{pmatrix} u_0 \circ Y_{\text{out/in}} \\ u_1 \circ Y_{\text{out/in}} \end{pmatrix}, \quad R_{\text{out/in}} = \hat{c}_{\text{out/in}}^{(d-1)/2} \mathbf{1},$$

and for $\epsilon_{\text{out/in}} = a_{\text{out/in}}^{\frac{1}{2}}$:

$$T_{\text{out/in}} = (i\sqrt{2})^{-1} \begin{pmatrix} \epsilon_{\text{out/in}}^{-\frac{1}{2}} & -\epsilon_{\text{out/in}}^{-\frac{1}{2}} \\ \epsilon_{\text{out/in}}^{\frac{1}{2}} & \epsilon_{\text{out/in}}^{\frac{1}{2}} \end{pmatrix}, \quad Z_{\text{out/in}} = (\chi_{\text{out/in}}^*)^{-1} R_{\text{out/in}} T_{\text{out/in}},$$

so that the Cauchy evolution of $P_{\text{out/in}}$ is given by

$$(5.6) \quad \mathcal{U}_{\text{out/in}}(t, s) = Z_{\text{out/in}} \circ \mathcal{U}_{\text{out/in}}^{\text{ad}}(t, s) \circ Z_{\text{out/in}}^{-1},$$

where $\mathcal{U}_{\text{out/in}}^{\text{ad}}$ stands for the evolution generated by

$$(5.7) \quad H_{\text{out/in}}^{\text{ad}} = \begin{pmatrix} \epsilon_{\text{out/in}} & 0 \\ 0 & \epsilon_{\text{out/in}} \end{pmatrix}.$$

The following fact will be needed in the sequel.

Lemma 5.4. *We have:*

$$Z^{-1}(t)Z_{\text{out/in}} - \mathbf{1}, \quad Z_{\text{out/in}}^{-1}Z(t) - \mathbf{1} \rightarrow 0 \text{ in } B(L^2(\Sigma) \otimes \mathbb{C}^2) \text{ as } t \rightarrow \pm\infty.$$

Proof. From Prop. 4.5 we obtain that $T_{\text{out/in}}^{-1}T(t) - \mathbf{1}$ tends to 0 in norm as $t \rightarrow \pm\infty$. By Lemma 5.1, $R(t)$ tends to $R_{\text{out/in}}$ in norm. Finally, from the proof of Lemma 5.1, see in particular (A.12), we obtain that $(\chi_{\text{out/in}}^*)^{-1}\chi_t^*$ tends to $\mathbf{1}$ in norm. This implies the lemma. \square

5.4. Construction of Hadamard states by scattering theory. In this subsection we construct the *out/in* states $\omega_{\text{out/in}}$ for the Klein-Gordon operator P and show that they are Hadamard states. We assume hypotheses (bg), (ast), (pos).

By the positivity condition (pos), the asymptotic Klein-Gordon operators $P_{\text{out/in}}$ admit *vacuum states* (that is, *ground states* for the dynamics $\mathcal{U}_{\text{out/in}}$) $\omega_{\text{out/in}}^{\text{vac}}$. In terms of $t = 0$ Cauchy data their covariances are the projections:

$$c_{\text{out/in}}^{\pm, \text{vac}} = Z_{\text{out/in}} \pi^{\pm} Z_{\text{out/in}}^{-1},$$

where π^{\pm} are defined in (4.23). Clearly we have

$$\mathcal{U}_{\text{out/in}}(t, s) c_{\text{out/in}}^{\pm, \text{vac}} \mathcal{U}_{\text{out/in}}(s, t) = c_{\text{out/in}}^{\pm, \text{vac}},$$

i.e. $\omega_{\text{out/in}}^{\text{vac}}$ are invariant under the asymptotic dynamics. For $t \in \mathbb{R}$ we now consider the projections:

$$(5.8) \quad \begin{aligned} c_{\text{out/in}}^{\pm, t}(0) &:= \mathcal{U}(0, t) c_{\text{out/in}}^{\pm, \text{vac}} \mathcal{U}(t, 0) \\ &= \mathcal{U}(0, t) \mathcal{U}_{\text{out/in}}(t, 0) c_{\text{out/in}}^{\pm, \text{vac}} \mathcal{U}_{\text{out/in}}(0, t) \mathcal{U}(t, 0). \end{aligned}$$

By taking the $t \rightarrow \pm\infty$ limit of $c_{\text{out/in}}^{\pm,t}(0)$ we obtain the time-0 covariances $c_{\text{out/in}}^{\pm}(0)$ of a state $\omega_{\text{out/in}}$ (for the Klein-Gordon operator P) that ‘equals $\omega_{\text{out/in}}^{\text{vac}}$ asymptotically’ at $t = \pm\infty$. The main new result that we prove is that $\omega_{\text{out/in}}$ are Hadamard states.

Before stating the theorem let us recall that the Sobolev spaces $H^m(\Sigma)$ are naturally defined using the reference Riemannian metric k on Σ . The *charge space* $H^{\frac{1}{2}}(\Sigma) \oplus H^{-\frac{1}{2}}(\Sigma)$ is the natural space of Cauchy data in connection with quantized Klein-Gordon fields.

Theorem 5.5. *Assume hypotheses (bg), (ast), (pos). Then*

$$(5.9) \quad \lim_{t \rightarrow \pm\infty} c_{\text{out/in}}^{\pm,t}(0) =: c_{\text{out/in}}^{\pm}(0) = c_{\text{ref}}^{\pm}(0) + \mathcal{W}^{-\infty}(\Sigma), \quad \text{in } B(H^{\frac{1}{2}}(\Sigma) \oplus H^{-\frac{1}{2}}(\Sigma)),$$

where $c_{\text{ref}}^{\pm}(0) = Z(0)\pi^{\pm}Z^{-1}(0)$. The operators $c_{\text{out/in}}^{\pm}(0)$ are pairs of projections defining a pure state $\omega_{\text{out/in}}$ for the Klein-Gordon operator P . Moreover $\omega_{\text{out/in}}$ is a Hadamard state.

Proof. From (5.5), (5.6) we obtain:

$$(5.10) \quad \begin{aligned} \mathcal{U}_{\text{out/in}}(0,t)\mathcal{U}(t,0) &= Z_{\text{out/in}}(0)\mathcal{U}_{\text{out/in}}^{\text{ad}}(0,t)Z_{\text{out/in}}^{-1}Z(t)\mathcal{U}^{\text{ad}}(t,0)Z^{-1}(0), \\ \mathcal{U}(0,t)\mathcal{U}_{\text{out/in}}(t,0) &= Z(0)\mathcal{U}^{\text{ad}}(0,t)Z^{-1}(t)Z_{\text{out/in}}\mathcal{U}_{\text{out/in}}^{\text{ad}}(t,0)Z_{\text{out/in}}^{-1}. \end{aligned}$$

It follows that:

$$(5.11) \quad \begin{aligned} c_{\text{out/in}}^{\pm,t}(0) &= Z(0)\mathcal{U}^{\text{ad}}(0,t)Z^{-1}(t)Z_{\text{out/in}}\mathcal{U}_{\text{out/in}}^{\text{ad}}(t,0) \\ &\quad \times \pi^{\pm}\mathcal{U}_{\text{out/in}}^{\text{ad}}(0,t)Z_{\text{out/in}}^{-1}Z(t)\mathcal{U}^{\text{ad}}(t,0)Z^{-1}(0). \end{aligned}$$

Since $Z(0) : L^2(\Sigma) \otimes \mathbb{C}^2 \rightarrow H^{\frac{1}{2}}(\Sigma) \oplus H^{-\frac{1}{2}}(\Sigma)$ is boundedly invertible it suffices to show the existence of the limit

$$\begin{aligned} d_{\text{out/in}}^{\pm} &= \lim_{t \rightarrow \pm\infty} \mathcal{U}^{\text{ad}}(0,t)Z^{-1}(t)Z_{\text{out/in}}\mathcal{U}_{\text{out/in}}^{\text{ad}}(t,0)\pi^{\pm} \\ &\quad \times \mathcal{U}_{\text{out/in}}^{\text{ad}}(0,t)Z_{\text{out/in}}^{-1}Z(t)\mathcal{U}^{\text{ad}}(t,0) \end{aligned}$$

in $B(L^2(\Sigma) \otimes \mathbb{C}^2)$.

By Prop. 5.6 (1) below we know that $\mathcal{U}^{\text{ad}}(t,s)$, $\mathcal{U}_{\text{out/in}}^{\text{ad}}(t,s)$ are uniformly bounded in $B(L^2(\Sigma) \otimes \mathbb{C}^2)$. Hence using Lemma 5.4 we can replace $Z^{-1}(t)Z_{\text{out/in}}$ and $Z_{\text{out/in}}^{-1}Z(t)$ by $\mathbf{1}$ in the rhs of (5.11), modulo an error of size $o(t^0)$ in $B(L^2(\Sigma) \otimes \mathbb{C}^2)$, i.e. we are reduced to prove the existence of the limit

$$\begin{aligned} d_{\text{out/in}}^{\pm} &:= \lim_{t \rightarrow \pm\infty} \mathcal{U}^{\text{ad}}(0,t)\mathcal{U}_{\text{out/in}}^{\text{ad}}(t,0)\pi^{\pm}\mathcal{U}_{\text{out/in}}^{\text{ad}}(0,t)\mathcal{U}^{\text{ad}}(t,0) \\ &= \lim_{t \rightarrow \pm\infty} W_{\text{out/in}}(t)\pi^{\pm}W_{\text{out/in}}^{-1}(t), \end{aligned}$$

where $W_{\text{out/in}}(t) = \mathcal{U}^{\text{ad}}(0,t)\mathcal{U}_{\text{out/in}}^{\text{ad}}(t,0)$. By Prop. 5.6 the limit exists in $B(L^2(\Sigma) \otimes \mathbb{C}^2)$ and equals $\pi^{\pm} + \mathcal{W}^{-\infty}(\Sigma)$. The limit operators $d_{\text{out/in}}^{\pm}$ are projections as norm limits of projections. It follows that

$$(5.12) \quad c_{\text{out/in}}^{\pm}(0) = Z(0)d_{\text{out/in}}^{\pm}Z(0)^{-1} + \mathcal{W}^{-\infty}(\Sigma) = c_{\text{ref}}^{\pm}(0) + \mathcal{W}^{-\infty}(\Sigma)$$

is a projection. The conditions (2.13), (2.14) are satisfied by $c_{\text{out/in}}^{\pm}(0)$ since they are satisfied by $c_{\text{out/in}}^{\pm,t}(0)$ for each finite t . Therefore $c_{\text{out/in}}^{\pm}$ are the covariances of two pure states $\omega_{\text{out/in}}$ for P . Finally as in [GOW] we obtain from (5.12) that $\omega_{\text{out/in}}$ are Hadamard states. \square

In the proof of Thm. 5.5, the crucial ingredient is the following proposition.

Proposition 5.6. *Let $H^{\text{ad}}(t), H_{\text{out/in}}^{\text{ad}}$ be as in (4.19), (5.7). Then:*

- (1) $\mathcal{W}_{\text{out/in}}^{\text{ad}}(t, s)$ and $\mathcal{W}^{\text{ad}}(t, s)$ are uniformly bounded in $B(H^m(\Sigma) \otimes \mathbb{C}^2)$, for all $m \in \mathbb{R}$.
- (2) Let $W_{\text{out/in}}(t) = \mathcal{W}^{\text{ad}}(0, t)\mathcal{W}_{\text{out/in}}^{\text{ad}}(t, 0)$. Then

$$\lim_{t \rightarrow +\infty} W_{\text{out/in}}(t)\pi^+W_{\text{out/in}}(t)^{-1} = \pi^+ + \mathcal{W}^{-\infty}(\Sigma) \otimes L(\mathbb{C}^2), \text{ in } B(L^2(\Sigma) \otimes \mathbb{C}^2).$$

Proof. *Proof of (1):* we can assume without loss of generality that $s = 0$. The statement for $\mathcal{W}_{\text{out/in}}^{\text{ad}}(t, 0)$ is obvious since $H_{\text{out/in}}^{\text{ad}} = \begin{pmatrix} \epsilon_{\text{out/in}} & 0 \\ 0 & -\epsilon_{\text{out/in}} \end{pmatrix}$. Let us prove it for $\mathcal{W}^{\text{ad}}(t, 0)$. We have:

$$(5.13) \quad \begin{aligned} H^{\text{ad}}(t) &= \begin{pmatrix} -b^-(t) + ir_b^-(t) & 0 \\ 0 & -b^+(t) + ir_b^+(t) \end{pmatrix} + \Psi_{\text{td}}^{-\infty, -1-\delta}(\mathbb{R}; \Sigma) \otimes B(\mathbb{C}^2) \\ &= \begin{pmatrix} \epsilon(t) & 0 \\ 0 & -\epsilon(t) \end{pmatrix} + \Psi_{\text{td}}^{0, -1-\delta}(\mathbb{R}; \Sigma) \otimes B(\mathbb{C}^2), \end{aligned}$$

by Props. 4.5, 4.6. Since $\epsilon(t)$ is selfadjoint, this implies that $\mathcal{W}^{\text{ad}}(t, 0)$ is uniformly bounded in $B(L^2(\Sigma))$, which proves (1) for $m = 0$.

We now note that $\|u\|_{H^m(\Sigma)} \sim \|\epsilon^m(t)u\|_{L^2(\Sigma)}$, uniformly for $t \in \mathbb{R}$, since $\epsilon(t)$ is elliptic uniformly for $t \in \mathbb{R}$. Therefore to prove (1) it suffices, using the uniform boundedness of $\mathcal{W}^{\text{ad}}(t, 0)$ in $B(L^2(\Sigma))$, to show that

(5.14) $\mathcal{W}^{\text{ad}}(0, t) (\epsilon(t)^m \otimes \mathbf{1}_{\mathbb{C}^2}) \mathcal{W}^{\text{ad}}(t, 0) (\epsilon(0)^{-m} \otimes \mathbf{1}_{\mathbb{C}^2})$ is uniformly bounded in $B(L^2(\Sigma))$. We have by (5.13):

$$\begin{aligned} &\partial_t \mathcal{W}^{\text{ad}}(0, t) (\epsilon(t)^m \otimes \mathbf{1}_{\mathbb{C}^2}) \mathcal{W}^{\text{ad}}(t, 0) (\epsilon(0)^{-m} \otimes \mathbf{1}_{\mathbb{C}^2}) \\ &= \mathcal{W}^{\text{ad}}(0, t) (\partial_t \epsilon^m(t) \otimes \mathbf{1}_{\mathbb{C}^2} + i[H^{\text{ad}}(t), \epsilon^m(t) \otimes \mathbf{1}_{\mathbb{C}^2}]) \mathcal{W}^{\text{ad}}(t, 0) (\epsilon(0)^{-m} \otimes \mathbf{1}_{\mathbb{C}^2}) \\ &= \mathcal{W}^{\text{ad}}(0, t) (\partial_t \epsilon^m(t) \otimes \mathbf{1}_{\mathbb{C}^2} + i[H^{\text{ad}}(t), \epsilon^m(t) \otimes \mathbf{1}_{\mathbb{C}^2}]) (\epsilon(t)^{-m} \otimes \mathbf{1}) \mathcal{W}^{\text{ad}}(t, 0) \\ &\quad \times \mathcal{W}^{\text{ad}}(0, t) (\epsilon(t)^m \otimes \mathbf{1}_{\mathbb{C}^2}) \mathcal{W}^{\text{ad}}(t, 0) (\epsilon(0)^{-m} \otimes \mathbf{1}_{\mathbb{C}^2}) \\ &=: M(t) \mathcal{W}^{\text{ad}}(0, t) (\epsilon(t)^m \otimes \mathbf{1}_{\mathbb{C}^2}) \mathcal{W}^{\text{ad}}(t, 0) (\epsilon(0)^{-m} \otimes \mathbf{1}_{\mathbb{C}^2}). \end{aligned}$$

By (td) and Prop. 3.10 we see that $\partial_t \epsilon^m(t) \in \Psi_{\text{td}}^{m, -1-\delta}$, and by (5.13) that $[H^{\text{ad}}(t), \epsilon^m(t) \otimes \mathbf{1}_{\mathbb{C}^2}] \in \Psi_{\text{td}}^{m, -1-\delta}$. Therefore $\|M(t)\|_{B(L^2(\Sigma) \otimes \mathbb{C}^2)} \in O(\langle t \rangle^{-1-\delta})$. Hence, setting

$$f(t) := \|\mathcal{W}^{\text{ad}}(0, t) (\epsilon(t)^m \otimes \mathbf{1}_{\mathbb{C}^2}) \mathcal{W}^{\text{ad}}(t, 0) (\epsilon(0)^{-m} \otimes \mathbf{1}_{\mathbb{C}^2})\|_{B(L^2(\Sigma))},$$

we have $f(0) = 1$, $|\partial_t f(t)| \in O(\langle t \rangle^{-1-\delta})f(t)$. If $f(t) \neq +\infty$ for each t , an application of Gronwall's inequality would immediately imply (5.14). If $m \leq 0$ the use of Gronwall's inequality is justified by applying the above time dependent operator to a vector $u \in H^m(\Sigma)$. If $m > 0$ we replace the unbounded operator $A = \epsilon(t) \otimes \mathbf{1}_{\mathbb{C}^2}$ by the bounded operator $A_\delta = A(1 + i\delta A)$, for $\delta > 0$. For the corresponding function $f_\delta(t)$ we obtain that $f_\delta(0) \leq 1$, $|\partial_t f_\delta(t)| \in O(\langle t \rangle^{-1-\delta})f_\delta(t)$ uniformly for $0 < \delta \leq 1$. Then (5.14) follows using that $\|A^m u\| = \sup_{0 < \delta \leq 1} \|A_\delta^m u\|$.

Proof of (2): note first that $[\pi^+, A] = 0$ for any diagonal operator A . Therefore:

$$W_{\text{out/in}}(t)\pi^+W_{\text{out/in}}(t)^{-1} = \mathcal{W}(0, t)\pi^+\mathcal{W}(t, 0),$$

and by (5.13)

$$(5.15) \quad \begin{aligned} &\partial_t (W_{\text{out/in}}(t)\pi^+W_{\text{out/in}}(t)^{-1}) = -i\mathcal{W}(0, t)[H^{\text{ad}}(t), \pi^+]\mathcal{W}(t, 0) \\ &= \mathcal{W}(0, t)[R_{-\infty}(t), \pi^+]\mathcal{W}(t, 0), \quad R_{-\infty} \in \Psi_{\text{td}}^{-\infty, -1-\delta}(\mathbb{R}; \Sigma) \otimes B(\mathbb{C}^2). \end{aligned}$$

By (1), this implies that $\partial_t(W_{\text{out/in}}(t)\pi^+W_{\text{out/in}}(t)^{-1}) \in \Psi_{\text{td}}^{-\infty, -1-\delta}(\mathbb{R}; \Sigma) \otimes B(\mathbb{C}^2)$, hence:

$$\begin{aligned} & \lim_{t \rightarrow +\infty} W_{\text{out/in}}(t)\pi^+W_{\text{out/in}}(t) \\ &= \pi^+ + \int_0^{+\infty} \partial_t(W_{\text{out/in}}(t)\pi^+W_{\text{out/in}}(t)^{-1})dt \text{ in } B(L^2(\Sigma) \otimes \mathbb{C}^2). \end{aligned}$$

The integral term belongs to $\mathcal{W}^{-\infty}(\Sigma)$. \square

APPENDIX A

A.1. Proof of Prop. 3.10. To prove Prop. 3.10 we first need an auxiliary lemma about parameter-dependent pseudodifferential calculus.

We start by introducing parameter dependent versions of the spaces $\Psi^m(\Sigma)$, $S^0(\mathbb{R}; \Psi^m(\Sigma))$.

We define the symbol classes $\tilde{S}^m(T^*\Sigma)$ for $m \in \mathbb{R}$ as the space of functions $c(x, k, \lambda) \in C^\infty(T^*\Sigma \times \mathbb{R})$ such that:

$$\partial_\lambda^\gamma \partial_x^\alpha \partial_k^\beta c(x, k, \lambda) \in O(\langle k \rangle + \langle \lambda \rangle)^{m-|\beta|-\gamma}, \quad \alpha, \beta \in \mathbb{N}^d, \quad \gamma \in \mathbb{N},$$

as usual understood after fixing a good chart cover and good chart diffeomorphisms, with uniformity of the constants with respect to the element of the cover. The standard example of such a symbol is $c(x, k, \lambda) = (a(x, k) + \langle \lambda \rangle^m)$, for $a \in S^m(T^*\Sigma)$ elliptic and positive.

The subspaces of symbols poly-homogeneous in (k, λ) are denoted by $\tilde{S}_{\text{ph}}^m(T^*\Sigma)$. We define $\tilde{\mathcal{W}}^{-\infty}(\Sigma)$ as the set of smooth maps $\mathbb{R} \ni \lambda \mapsto a(\lambda) \in \mathcal{W}^{-\infty}(\Sigma)$ such that:

$$\|\partial_\lambda^\gamma a(\lambda)\|_{B(H^{-m}(\Sigma), H^m(\Sigma))} \in O(\langle \lambda \rangle^{-n}), \quad \forall m, n, \gamma \in \mathbb{N},$$

and we set

$$\tilde{\Psi}^m(\Sigma) := \text{Op}(\tilde{S}_{\text{ph}}^m(T^*\Sigma)) + \tilde{\mathcal{W}}^{-\infty}(\Sigma).$$

We also define the time-dependent versions:

$$S^0(\mathbb{R}; \tilde{S}_{\text{ph}}^m(T^*\Sigma)), \quad S^0(\mathbb{R}; \tilde{\mathcal{W}}^{-\infty}(\Sigma)), \quad S^0(\mathbb{R}; \tilde{\Psi}^m(\Sigma)),$$

in analogy with Subsect. 3.3. For example $c(t, x, k, \lambda) \in S^0(\mathbb{R}; \tilde{S}^m(T^*\Sigma))$ if

$$\partial_t^n \partial_\lambda^\gamma \partial_x^\alpha \partial_k^\beta c(t, x, k, \lambda) \in O(\langle t \rangle^{-n} (\langle k \rangle + \langle \lambda \rangle)^{m-|\beta|-\gamma}), \quad \alpha, \beta \in \mathbb{N}^d, \quad \gamma, n \in \mathbb{N}.$$

Lemma A.1. *Let $a(t) \in S^0(\mathbb{R}; \Psi^2(\Sigma))$ such that $a(t)$ is elliptic, selfadjoint on $L^2(\Sigma)$ with $a(t) \geq c_0 \mathbf{1}$, $c_0 > 0$. Then $(a(t) + \lambda^2)^{-1} \in S^0(\mathbb{R}; \tilde{\Psi}^{-2}(\Sigma))$.*

Proof. The proof is based on a reduction to the situation without the parameter λ . We first present the argument in the time-independent case.

Step 1. Let us denote by $l \in \mathbb{R}$ the dual variable to λ . We consider the manifold of bounded geometry $\Sigma_x \times \mathbb{R}_l$ equipped with the metric $h_{ij}(x)dx^i dx^j + dl^2$. As good chart covering we can take $\tilde{U}_i = U_i \times \mathbb{R}$, $\tilde{\psi}_i(x, l) = (\psi_i(x), l)$ where $\{U_i, \psi_i\}_{i \in \mathbb{N}}$ is a good chart covering for (Σ, h) . A subordinate good partition of unity is $\tilde{\chi}_i(x, l) = \chi_i(x)$.

The classes $S_{\text{ph}}^m(T^*(\Sigma \times \mathbb{R}))$ are then defined as in Subsect. 3.3. Denoting $\text{ad}_A B := [A, B]$, one sets and one sets:

$$(A.1) \quad \mathcal{W}^{-\infty}(\Sigma \times \mathbb{R}) = \{A : \text{ad}_l^\gamma A \in B(H^{-m}(\Sigma \times \mathbb{R}), H^m(\Sigma \times \mathbb{R})), m, \gamma \in \mathbb{N}\},$$

This choice of the ideal $\mathcal{W}^{-\infty}(\Sigma \times \mathbb{R})$ is dictated by the definition of $\tilde{\mathcal{W}}^{-\infty}(\Sigma)$. We set then $\Psi^m(\Sigma \times \mathbb{R}) = \tilde{\text{Op}}(S_{\text{ph}}^m(T^*(\Sigma \times \mathbb{R}))) + \mathcal{W}^{-\infty}(\Sigma \times \mathbb{R})$, where $\tilde{\text{Op}}$ is defined as in Subsect. 3.5 with Σ replaced by $\Sigma \times \mathbb{R}$. Note that because of our choice of the chart covering $\tilde{\text{Op}}$ is the usual Weyl quantization w.r.t. the (l, λ) variables.

Step 2. In step 2 we describe the link between $\Psi^m(\Sigma \times \mathbb{R})$ and $\tilde{\Psi}^m(\Sigma)$. We note that

$$\tilde{S}_{\text{ph}}^m(T^*\Sigma) = \{c \in S_{\text{ph}}^m(T^*(\Sigma \times \mathbb{R})) : \partial_l c = 0\},$$

and denoting by T_l the group of translations in l we have

$$[T_l, \widetilde{\text{Op}}(c)] = 0, \forall l \in \mathbb{R} \Leftrightarrow c \in \tilde{S}_{\text{ph}}^m(T^*\Sigma).$$

Equivalently, if \mathcal{F} is the Fourier transform in l we have

$$\begin{aligned} (A.2) \quad & c \in S_{\text{ph}}^m(T^*(\Sigma \times \mathbb{R})), [T_l, \widetilde{\text{Op}}(c)] = 0 \\ & \Leftrightarrow \mathcal{F}\widetilde{\text{Op}}(c)\mathcal{F}^{-1} = \int_{\mathbb{R}}^{\oplus} \text{Op}(c(\lambda))d\lambda, \text{ for } c(\lambda) \in \tilde{S}_{\text{ph}}^m(T^*\Sigma). \end{aligned}$$

Let now $w \in \mathcal{W}^{-\infty}(\Sigma \times \mathbb{R})$ with $[w, T_l] = 0$. We have:

$$(A.3) \quad \mathcal{F}w\mathcal{F}^{-1} = \int_{\mathbb{R}}^{\oplus} w(\lambda)d\lambda.$$

Since $\text{ad}_l^\gamma w \in \bigcap_{m \in \mathbb{N}} B(H^{-m}(\Sigma \times \mathbb{R}), H^m(\Sigma \times \mathbb{R}))$ we obtain that:

$$\int_{\mathbb{R}} \langle \lambda \rangle^n \|\partial_\lambda^\gamma w(\lambda)u(\lambda)\|_{H^p(\Sigma)}^2 d\lambda \leq C_{n,p} \int_{\mathbb{R}} \langle \lambda \rangle^{-n} \|u(\lambda)\|_{H^{-p}(\Sigma)}^2 d\lambda, \forall \gamma, n, p \in \mathbb{N},$$

or equivalently

$$\int_{\mathbb{R}}^{\oplus} \langle \lambda \rangle^n (-\Delta_h + 1)^{p/2} \partial_\lambda^\gamma w(\lambda) (-\Delta_h + 1)^{p/2} d\lambda \in B(L^2(\Sigma \times \mathbb{R})).$$

By Sobolev's embedding theorem this implies that

$$\|\partial_\lambda^\gamma w(\lambda)\|_{B(H^{-p/2}(\Sigma), H^{p/2}(\Sigma))} \in O(\langle \lambda \rangle^{-n}) \forall \gamma, n, p \in \mathbb{N},$$

hence $w(\lambda) \in \widetilde{\mathcal{W}}^{-\infty}(\Sigma)$. Conversely, if $w(\lambda) \in \widetilde{\mathcal{W}}^{-\infty}(\Sigma)$ it is immediate that w defined by (A.3) belongs to $\mathcal{W}^{-\infty}(\Sigma \times \mathbb{R})$. Hence we have shown

$$w \in \mathcal{W}^{-\infty}(\Sigma \times \mathbb{R}), [w, T_l] = 0$$

$$(A.4) \quad \Leftrightarrow \mathcal{F}w\mathcal{F}^{-1} = \int_{\mathbb{R}}^{\oplus} w(\lambda)d\lambda, \text{ for } w(\lambda) \in \widetilde{\mathcal{W}}^{-\infty}(\Sigma).$$

Let us now consider the time-dependent situation. If we define the time-dependent classes $C_b^\infty(\mathbb{R}; \tilde{S}^m(T^*\Sigma))$, $C_b^\infty(\mathbb{R}; \widetilde{\mathcal{W}}^{-\infty}(\Sigma))$ and $C_b^\infty(\mathbb{R}; \tilde{\Psi}^m(\Sigma))$ in the obvious way, then

$$\begin{aligned} (A.5) \quad & c \in C_b^\infty(\mathbb{R}; S_{\text{ph}}^m(T^*(\Sigma \times \mathbb{R}))), [T_l, \widetilde{\text{Op}}(c)(t)] = 0 \\ & \Leftrightarrow \mathcal{F}\widetilde{\text{Op}}(c)(t)\mathcal{F}^{-1} = \int_{\mathbb{R}}^{\oplus} \text{Op}(c(t, \lambda))d\lambda, c(t, \lambda) \in C_b^\infty(\mathbb{R}; \tilde{S}^m(T^*\Sigma)), \\ & w \in C_b^\infty(\mathbb{R}; \mathcal{W}^{-\infty}(\Sigma \times \mathbb{R})), [w(t), T_l] = 0 \\ & \Leftrightarrow \mathcal{F}w(t)\mathcal{F}^{-1} = \int_{\mathbb{R}}^{\oplus} w(t, \lambda)d\lambda, w(t, \lambda) \in C_b^\infty(\mathbb{R}; \widetilde{\mathcal{W}}^{-\infty}(\Sigma)). \end{aligned}$$

The same results hold also if we replace $C_b^\infty(\mathbb{R}; A)$ by $S^\delta(\mathbb{R}; A)$ for $A = S_{\text{ph}}^m(T^*(\Sigma \times \mathbb{R}))$, $\tilde{S}^m(T^*\Sigma)$ etc. In fact it suffices to note that $c(t) \in S^\delta(\mathbb{R}; A)$ iff $\langle t \rangle^{-\delta+n} \partial_t^n c(t) \in C_b^\infty(\mathbb{R}; A)$ for all $n \in \mathbb{N}$.

Step 3. Let now $a(t) \in S^0(\mathbb{R}; \Psi^2(\Sigma))$ be as in the lemma and let $A(t) = a(t) + D_t^2$ acting on $L^2(\Sigma \times \mathbb{R})$. The operator $A(t)$ is elliptic in $S^0(\mathbb{R}; \Psi^2(\Sigma \times \mathbb{R}))$, selfadjoint on $H^2(\Sigma \times \mathbb{R})$ with $A(t) \geq c_0 \mathbf{1}$ for c_0 as in the lemma.

We would like to apply Thm. 3.7 for $\alpha = -1$ to the class $S^0(\mathbb{R}; \Psi^m(\Sigma \times \mathbb{R}))$ to obtain that $A(t)^{-1}$ belongs to $S^0(\mathbb{R}; \Psi^{-2}(\Sigma \times \mathbb{R}))$. Note that the ideal $S^0(\mathbb{R}; \mathcal{W}^{-\infty}(\Sigma \times \mathbb{R}))$ is smaller than the one used in Subsect. 3.5 for the manifold $\Sigma \times \mathbb{R}$, because multi-commutators $\text{ad}_l^\gamma A$ appear in (A.1).

However we can still apply the abstract framework in [ALNV] to this situation, see [GOW, Subsect. 5.3.4] for a concise summary. We choose as Hilbert space $\mathcal{H} = L^2(\mathbb{R}_t \times \Sigma_x \times \mathbb{R}_l)$. As injective operator in $S^0(\mathbb{R}; \mathcal{W}^{-\infty}(\Sigma \times \mathbb{R}))$ we choose $e^{\Delta_h + D_l^2 + 1}$. The only point which differs a little from the situation in Subsect. 3.5 is the *spectral invariance* of $S^0(\mathbb{R}; \mathcal{W}^{-\infty}(\Sigma \times \mathbb{R}))$, see the proof of Thm. 3.7: if $R \in S^0(\mathbb{R}; \mathcal{W}^{-\infty}(\Sigma \times \mathbb{R}))$ and $(\mathbf{1} + R)$ is boundedly invertible in $B(\mathcal{H})$, then we have

$$(\mathbf{1} + R)^{-1} = \mathbf{1} + R_1, \quad R_1 = -R + R(\mathbf{1} + R)^{-1}R.$$

We have $R_1 = \int_{\mathbb{R}}^{\oplus} R_{-1}(t) dt$, for

$$(A.6) \quad R_1(t) = -R(t) + R(t)(\mathbf{1} + R(t))^{-1}R(t).$$

We have to check that

$$\|\partial_t^n \text{ad}_l^\gamma R_{-1}(t)\|_{B(H^{-p}(\Sigma \times \mathbb{R}), H^p(\Sigma \times \mathbb{R}))} \in O(\langle t \rangle^{-n}), \quad \forall \gamma, n, p \in \mathbb{N}.$$

This follows from (A.6) using the Leibniz rule for ∂_t and ad_l and the identities

$$\begin{aligned} \partial_t(\mathbf{1} + R(t))^{-1} &= -(\mathbf{1} + R(t))^{-1} \partial_t R(t) (\mathbf{1} + R(t))^{-1} \\ \text{ad}_l(\mathbf{1} + R(t))^{-1} &= -(\mathbf{1} + R(t))^{-1} \text{ad}_l R(t) (\mathbf{1} + R(t))^{-1}. \end{aligned}$$

In conclusion we can apply Seeley's theorem and obtain that $A(t)^{-1}$ belongs to $S^0(\mathbb{R}; \Psi^{-m}(\Sigma \times \mathbb{R}))$. We have then:

$$\mathcal{F} A(t)^{-1} \mathcal{F}^{-1} = \int_{\mathbb{R}}^{\oplus} (a(t) + \lambda^2)^{-1} d\lambda,$$

which by (A.5) implies that $(a(t) + \lambda^2)^{-1} \in S^0(\mathbb{R}; \tilde{\Psi}^{-2}(\Sigma))$. This concludes the proof of the lemma. \square

Proof of Prop. 3.10. In view of the identity

$$a_1^{1+\alpha} - a_2^{1+\alpha} = (a_1 - a_2)a_1^\alpha + a_2(a_1^\alpha - a_2^\alpha),$$

we see that it suffices to prove the proposition for $0 < \alpha < 1$. We will use the following formula, valid for example if a is a selfadjoint operator on a Hilbert space \mathcal{H} with $a \geq c\mathbf{1}$, $c > 0$:

$$(A.7) \quad a^\alpha = C_\alpha \int_0^{+\infty} (a + s)^{-1} s^\alpha ds = C_\alpha \int_{\mathbb{R}} (a + \lambda^2)^{-1} \lambda^{2\alpha+1} d\lambda, \quad \alpha \in \mathbb{R},$$

where the integrals are norm convergent in say, $B(\text{Dom } a^m, \mathcal{H})$ for m large enough.

We have for $r(t) = a_1(t) - a_2(t)$:

$$\begin{aligned} (a_1(t) + \lambda^2)^{-1} &= (a_2(t) + \lambda^2)^{-1} (\mathbf{1} + r(t)(a_1(t) + \lambda^2)^{-1}) \\ &= (a_2(t) + \lambda^2)^{-1} + (a_2(t) + \lambda^2)^{-2} (a_2(t) + \lambda^2) r(t) (a_1(t) + \lambda^2)^{-1} \\ &= (a_2(t) + \lambda^2)^{-1} + (a_2(t) + \lambda^2)^{-2} a_2(t) c_1(t, \lambda) \\ &= (a_2(t) + \lambda^2)^{-1} + a_2(t) c_2(t, \lambda), \end{aligned}$$

where using Lemma A.1, $c_1(t, \lambda) \in S^{-\delta}(\mathbb{R}; \tilde{\Psi}^0(\Sigma))$ and $c_2(t, \lambda) \in S^{-\delta}(\mathbb{R}; \tilde{\Psi}^{-4}(\Sigma))$. From (A.7) we obtain that:

$$(A.8) \quad a_1^\alpha(t) - a_2^\alpha(t) = C_\alpha a_2(t) \int_{\mathbb{R}} c_2(t, \lambda) \lambda^{2\alpha+1} d\lambda.$$

We now write $c_2(t, \lambda)$ as $\text{Op}(d_2(t, \lambda)) + w_2(t, \lambda)$, for $d_2 \in S^{-\delta}(\mathbb{R}; \widetilde{S}_{\text{ph}}^{-4}(T^*\Sigma))$ and $w_2(t, \lambda) \in S^{-\delta}(\mathbb{R}; \widetilde{\mathcal{W}}^{-\infty}(\Sigma))$. Using that

$$\int_{\mathbb{R}} ((\xi) + \langle \lambda \rangle)^{-4-k} \lambda^{2\alpha+1} d\lambda \sim \langle \xi \rangle^{2\alpha-2-k},$$

we first obtain that

$$\int_{\mathbb{R}} d_2(t, \lambda) \lambda^{2\alpha+1} d\lambda \in S^{-\delta}(\mathbb{R}; S_{\text{ph}}^{2\alpha-2}(\Sigma)).$$

Similarly we obtain that $\int_{\mathbb{R}} w_2(t, \lambda) \lambda^{2\alpha+1} d\lambda \in S^{-\delta}(\mathbb{R}; \mathcal{W}^{-\infty}(\Sigma))$. Using (A.8) this implies that $a_1^\alpha(t) - a_2^\alpha(t) \in S^{-\delta}(\mathbb{R}; \Psi^{2\alpha}(\Sigma))$, as claimed.

A.2. Proof of Prop. 4.5. We follow the proof in [GW1]. The out and in cases are treated similarly. We set $a_0 = \frac{i}{2}(\epsilon^{-1}\partial_t\epsilon + \epsilon^{-1}r\epsilon)$,

$$F(c) := \frac{1}{2}\epsilon^{-1}(\partial_t c + [\epsilon, c] + irc - c^2) = G(c) - \frac{1}{2}\epsilon^{-1}c^2.$$

and look for $b(t)$ as $\epsilon(t) + b_0$, where $b_0 = a_0 + F(b_0)$. Let us start by studying some properties of the map F . First if $c \in \Psi_{\text{td}}^{0, -\mu}$ then:

$$\begin{aligned} G(c) &\in \Psi_{\text{td}}^{-1,0} \times \Psi_{\text{td}}^{0,-1-\mu} + \Psi_{\text{td}}^{-1,0} \times \Psi_{\text{td}}^{0,-\mu} + \Psi_{\text{td}}^{-1,0} \Psi_{\text{td}}^{0,-1-\delta} \times \Psi_{\text{td}}^{0,-\mu}, \\ \epsilon^{-1}c^2 &\in \Psi_{\text{td}}^{-1,-2\mu}, \end{aligned}$$

hence

$$(A.9) \quad c \in \Psi_{\text{td}}^{0,-\mu} \Rightarrow F(c) \in \Psi_{\text{td}}^{-1,-\mu}.$$

Secondly, if $c_1, c_2 \in \Psi_{\text{td}}^{0,-\mu}$ and $c_1 - c_2 \in \Psi_{\text{td}}^{-j,-\mu}$ then:

$$\begin{aligned} G(c_1) - G(c_2) &= G(c_1 - c_2) \\ &\in \Psi_{\text{td}}^{-1,0} \times \Psi_{\text{td}}^{-j,-1-\mu} + \Psi_{\text{td}}^{-1,0} \times \Psi_{\text{td}}^{-j,-\mu} + \Psi_{\text{td}}^{-1,0} \times \Psi_{\text{td}}^{0,-1-\delta} \times \Psi_{\text{td}}^{-j,-\mu}, \\ \epsilon^{-1}(c_1^2 - c_2^2) &= \epsilon^{-1}c_1(c_1 - c_2) + \epsilon^{-1}(c_1 - c_2)c_2 \\ &\in \Psi_{\text{td}}^{-1,0} \times \Psi_{\text{td}}^{0,-\mu} \times \Psi_{\text{td}}^{-j,-\mu} + \Psi_{\text{td}}^{-1,0} \times \Psi_{\text{td}}^{-j,-\mu} \times \Psi_{\text{td}}^{0,-\mu}, \end{aligned}$$

hence

$$(A.10) \quad c_1, c_2 \in \Psi_{\text{td}}^{0,-\mu}, c_1 - c_2 \in \Psi_{\text{td}}^{-j,-\mu} \Rightarrow F(c_1) - F(c_2) \in \Psi_{\text{td}}^{-j-1,-\mu}.$$

We also have

$$\begin{aligned} a_0 &= \frac{i}{2}(\epsilon^{-1}\partial_t\epsilon + \epsilon^{-1}r\epsilon) \\ &\in \Psi_{\text{td}}^{-1,0} \times \Psi_{\text{td}}^{1,-1-\delta} + \Psi_{\text{td}}^{-1,0} \times \Psi_{\text{td}}^{0,-1-\delta} \times \Psi_{\text{td}}^{1,0} \subset \Psi_{\text{td}}^{0,-1-\delta}. \end{aligned}$$

We now follow the proof in [GW1, Lemma A.1], setting $b_0 = a_0$, $b_n = a_0 + F(b_{n-1})$, and obtain by induction that $b_n - b_{n-1} \in \Psi_{\text{td}}^{-n,-1-\delta}$. We set

$$b_0 \sim a_0 + \sum_{n=1}^{\infty} b_n - b_{n-1} \in \Psi_{\text{td}}^{0,-1-\delta}$$

by Lemma 3.11. We obtain that

$$i\partial_t b - b^2 + a + irb \in \Psi_{\text{td}}^{-\infty,-1-\delta}.$$

By construction we have $b(t) = \epsilon(t) + \Psi_{\text{td}}^{0,-1-\delta}(\mathbb{R}; \Sigma)$. Applying Prop. 3.10 we get

$$\epsilon(t) = \epsilon_{\text{out/in}} + \Psi_{\text{td}}^{1,-\delta}(\mathbb{R}; \Sigma) \text{ on } \mathbb{R}^{\pm} \times \Sigma.$$

□

A.3. Proof of Prop. 4.6. From Prop. 4.5 we first obtain that $b^+ - b^- = (b + b^*) = 2\epsilon + \Psi_{\text{td}}^{0,-1-\delta}(\mathbb{R}; \Sigma)$. It follows first that $(b^+ - b^-)^2 = 4a + \Psi_{\text{td}}^{1,-1-\delta}(\mathbb{R}; \Sigma)$ and then by Prop. 4.5 that

$$(b^+ - b^-)^\alpha = ((b^+ - b^-)^2)^{\alpha/2} = \begin{cases} (2\epsilon)^{\frac{1}{2}} + \Psi_{\text{td}}^{0,-1-\delta}(\mathbb{R}; \Sigma), & \alpha = \frac{1}{2} \\ (2\epsilon)^{-\frac{1}{2}} + \Psi_{\text{td}}^{-3/2,-1-\delta}(\mathbb{R}; \Sigma), & \alpha = -\frac{1}{2}. \end{cases}$$

We obtain again by Prop. 4.5 that:

$$\begin{aligned} [(b^+ - b^-)^{-\frac{1}{2}}, b^\pm] &= [(2\epsilon)^{-\frac{1}{2}} + \Psi_{\text{td}}^{-3/2,-1-\delta}(\mathbb{R}; \Sigma), \pm\epsilon + \Psi_{\text{td}}^{0,-1-\delta}(\mathbb{R}; \Sigma)] \\ &\in \Psi_{\text{td}}^{-3/2,-1-\delta}(\mathbb{R}; \Sigma), \\ \partial_t(b^+ - b^-)^{-\frac{1}{2}}(b^+ - b^-)^{\frac{1}{2}} &= (\partial_t(2\epsilon)^{-\frac{1}{2}} + \Psi_{\text{td}}^{-3/2,-2-\delta}(\mathbb{R}; \Sigma)) \times \Psi_{\text{td}}^{\frac{1}{2},0}(\mathbb{R}; \Sigma) \\ &= \partial_t(2\epsilon)^{-\frac{1}{2}} \times \Psi_{\text{td}}^{\frac{1}{2},0}(\mathbb{R}; \Sigma) + \Psi_{\text{td}}^{-1,-2-\delta}(\mathbb{R}; \Sigma). \end{aligned}$$

Since by Prop. 3.10 $(2\epsilon)^{-\frac{1}{2}} = (2\epsilon_{\text{out/in}})^{-\frac{1}{2}} + \Psi_{\text{td}}^{-3/2,-\delta}(\mathbb{R}^\pm; \Sigma)$, we have

$$\partial_t(2\epsilon)^{-\frac{1}{2}} \in \Psi_{\text{td}}^{-3/2,-1-\delta}(\mathbb{R}; \Sigma) \Rightarrow \partial_t(b^+ - b^-)^{-\frac{1}{2}}(b^+ - b^-)^{\frac{1}{2}} \in \Psi_{\text{td}}^{-1,-1-\delta}(\mathbb{R}; \Sigma).$$

Since by hypothesis (td), $r \in \Psi_{\text{td}}^{0,-1-\delta}(\mathbb{R}; \Sigma)$, we obtain that $r_b^\pm \in \Psi_{\text{td}}^{0,-1-\delta}(\mathbb{R}; \Sigma)$. Finally we obtain immediately from Prop. 4.5 that $r_\infty^\pm = i\partial_t b^\pm - (b^\pm)^2 + a + irb^\pm \in \Psi_{\text{td}}^{-\infty,-1-\delta}(\mathbb{R}; \Sigma)$. \square

A.4. Proof of Lemma 5.1. Let us fix two good chart coverings $\{U_i, \psi_i\}_{i \in \mathbb{N}}$ and $\{\tilde{U}_i, \tilde{\psi}_i\}_{i \in \mathbb{N}}$ with $U_i \Subset \tilde{U}_i$. Since $b \in C_b^\infty(\mathbb{R}; \text{BT}_0^1(\Sigma, k))$, we obtain easily by transporting b to $B_n(0, 1)$ using ψ_i that there exists $t_{+,\epsilon} > 0$ such that $y(t, s, \cdot)$ is a bounded diffeomorphism of (Σ, k) , uniformly for $|t - s| \leq t_{+,\epsilon}$. By the group property of the flow we can replace $t_{+,\epsilon}$ by any $t_+ > 0$, keeping the above uniformity property.

Moreover if $b_i := (\psi_i^{-1})^* b$ we obtain from (ast) that $b_i \in S^{-\delta}(\mathbb{R}; \text{BT}_0^1(B_n(0, 1)))$, uniformly in $i \in \mathbb{N}$. If $y_i(t, s, \cdot)$ denotes the flow of b_i we obtain that:

$$y_i(t, s, x) = x + \int_s^t b_i(\sigma, y_i(\sigma, s, x)) d\sigma.$$

From this we obtain that there exists $t_+ \gg 1$ such that

$$y_i(\pm t, \pm t_+, \cdot) : B_n(0, \frac{1}{2}) \rightarrow B_n(0, 1)$$

for all $t \geq t_+$ and moreover

$$\lim_{t \rightarrow \pm\infty} y_i(t, \pm t_+, y) = \int_{\pm t_+}^{\pm\infty} b_i(\sigma, y_i(\sigma, \pm t_+, x)) ds =: y_i(\pm\infty, t_+, x).$$

We can also choose t_+ large enough so that if we set

$$(A.11) \quad y(\pm\infty, \pm t_+, x) := \psi_i^{-1} \circ y_i(\pm\infty, \pm t_+, \cdot) \circ \psi_i(y), \quad x \in U_i$$

then $y(\pm\infty, \pm t_+, \cdot)$ is well defined, and is a bounded diffeomorphism of (Σ, k) . We now set:

$$y_{\text{out/in}} := y(\pm\infty, \pm t_+, \cdot) \circ y(\pm t_+, 0, \cdot),$$

which is also a bounded diffeomorphism of (Σ, k) . We also obtain from (A.11) and the previous estimates on $y(t, s, \cdot)$ for $|t - s| \leq t_+$ that $\{y(t, 0, \cdot)\}_{t \in \mathbb{R}}$ is a bounded family of bounded diffeomorphisms of (Σ, k) . Moreover from (A.11) we obtain that

$$(A.12) \quad y_i(t, 0, x) - y_{i,\text{out/in}}(x) \in S^{1-\delta'}(\mathbb{R}; C_b^\infty(B_n(0, 1))), \quad \text{uniformly in } i \in \mathbb{N}.$$

Let us now consider the metric $\chi^* g$.

Since $v \cdot dt = 0$, $\chi^* g = {}^t D\chi(g \circ \chi)D\chi = -\hat{c}^2(t, x)dt^2 + \hat{h}(t, x)dx^2$. Using (A.12) we obtain that

$$\begin{aligned}\hat{c}(t, x) &= c(t, y(t, x)) + S^{-2\delta'}(\mathbb{R}; \text{BT}_0^0(\Sigma)) \\ &= c_{\text{out/in}}(y(t, x)) + S^{-\min(2\delta', \delta)}(\mathbb{R}^\pm; \text{BT}_0^0(\Sigma)) \\ &= c_{\text{out/in}}(y_{\text{out/in}}(x)) + S^{-\min(1-\delta', \delta)}(\mathbb{R}^\pm; \text{BT}_0^0(\Sigma)).\end{aligned}$$

Similarly,

$$\begin{aligned}\hat{h}(t, x) &= {}^t D y(t, x) h(t, y(t, x)) D y(t, x) \\ &= {}^t D y(t, x) h_{\text{out/in}}(y(t, x)) D y(t, x) + S^{-\delta}(\mathbb{R}^\pm; \text{BT}_2^0(\Sigma)) \\ &= {}^t D y_{\text{out/in}}(x) h_{\text{out/in}}(y_{\text{out/in}}(x)) D y_{\text{out/in}}(x) + S^{-\min(1-\delta', \delta)}(\mathbb{R}^\pm; \text{BT}_2^0(\Sigma)), \\ \chi^* V &= r(t, y(t, x)) = r_{\text{out/in}}(y(t, x)) + S^{-\delta}(\mathbb{R}^\pm; \text{BT}_0^0(\Sigma)) \\ &= V_{\text{out/in}}(y_{\text{out/in}}(x)) + S^{-\min(1-\delta', \delta)}(\mathbb{R}^\pm; \text{BT}_0^0(\Sigma)).\end{aligned}$$

Since by definition

$$\hat{h}_{\text{out/in}} = y_{\text{out/in}}^* h_{\text{out/in}}, \quad \hat{c}_{\text{out/in}} = y_{\text{out/in}}^* c_{\text{out/in}}, \quad \hat{V}_{\text{out/in}} = y_{\text{out/in}}^* V_{\text{out/in}},$$

we obtain the assertion. \square

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REFERENCES

- [Ar] H. Araki, M. Shiraishi, *On quasi-free states of canonical commutation relations I*, Publ. RIMS Kyoto Univ. 7 (1971/72), 105–120.
- [ALNV] B. Ammann, R. Lauter, V. Nistor, A. Vasy, *Complex powers and non-compact manifolds*, Comm. in PDE **29** (2004), 671–705.
- [BDH] C. Brouder, N.V. Dang, F. Hélein, *A smooth introduction to the wavefront set*, J. Phys. A: Math. Theor. 47, no. 44 (2014).
- [BGP] C. Bär, N. Ginoux, F. Pfäffle, *Wave equation on Lorentzian manifolds and quantization*, *ESI Lectures in Mathematics and Physics*, EMS (2007).
- [BF1] R. Brunetti, K. Fredenhagen, *Microlocal Analysis and Interacting Quantum Field Theories: Renormalization on Physical Backgrounds*, Comm. Math. Phys. **208** (2000), 623–661.
- [BF2] C. Bär (ed.), K. Fredenhagen (ed.), *Quantum Field Theory on Curved Spacetimes*, Lect. Notes Phys. 786 (2009).
- [BJ] M. Brum, S. E. Jorás, *Hadamard state in Schwarzschild-de Sitter spacetime*, Class. Quantum Grav. 32, no. 1 (2014).
- [BVW] D. Baskin, A. Vasy, J. Wunsch, *Asymptotics of radiation fields in asymptotically Minkowski space*, to appear in Am. J. Math., [arXiv:1212.5141](https://arxiv.org/abs/1212.5141) (2012).
- [CC] Y. Choquet-Bruhat, Y. Cotsakis, *Global hyperbolicity and completeness*, J. Geom. and Phys. **43**, (2002), 345–350.
- [Co] H.O. Cordes, *A global parametrix for pseudo-differential operators over \mathbb{R}^n , with applications*, preprint no. 90, SFB 72, Bonn
- [Da] N.V. Dang, *Renormalization of quantum field theory on curved spacetimes, a causal approach*, Ph.D. thesis, Paris Diderot University, (2013).
- [DD] C. Dappiaggi, N. Drago, *A new deformation argument for Hadamard states via an extended Møller operator*, preprint [arXiv:1506.09122](https://arxiv.org/abs/1506.09122) (2015).
- [DG1] J. Dereziński, C. Gérard, *Scattering Theory of Classical and Quantum N-particle Systems*, Texts and Monographs in Physics, Springer, 1997.

- [DG2] J. Dereziński, C. Gérard, *Mathematics of Quantization and Quantum Fields*, Cambridge Monographs in Mathematical Physics, Cambridge University Press, 2012.
- [DH] J.J. Duistermaat J.J., L. Hörmander, *Fourier integral operators II*, Acta Math. **128** (1972), 183–269.
- [DK1] J. Dimock, B.S. Kay, *Classical wave operators and asymptotic quantum field operators on curved space-times*, Annales de l’I.H.P. A 37 no. 2, (1982), 93–114.
- [DK2] J. Dimock, B.S. Kay, *Classical and quantum scattering theory for linear scalar fields on the Schwarzschild metric. I*, Ann. Phys. 175 (2), (1987), 366–426.
- [DK3] J. Dimock, B.S. Kay, *Classical and quantum scattering theory for linear scalar fields on the Schwarzschild metric. II*, J. Math. Phys. 27, (1986), 2520.
- [DMP1] C. Dappiaggi, V. Moretti, N. Pinamonti, *Distinguished quantum states in a class of cosmological spacetimes and their Hadamard property*, J. Math. Phys. 50 (2009) 062304.
- [DMP2] C. Dappiaggi, V. Moretti, N. Pinamonti, *Rigorous construction and Hadamard property of the Unruh state in Schwarzschild spacetime*, Adv. Theor. Math. Phys. 15 (2011) 355.
- [Dr] A. Drouot, *A quantitative version of Hawking radiation*, arXiv:1412.8379 (2014).
- [DRS] M. Dafermos, I. Rodnianski, Y. Shlapentokh-Rothman, *A scattering theory for the wave equation on Kerr black hole exteriors*, preprint arXiv:1412.8379 (2014).
- [FNW] S.A. Fulling, F.J. Narcowich, R.M. Wald, *Singularity structure of the two-point function in quantum field theory in curved space-time, II*, Annals of Physics, **136** (1981), 243–272.
- [FV1] C.J. Fewster, R. Verch, *The necessity of the Hadamard condition*, Class. Quant. Grav. 30, 235027 (2013).
- [FV2] C.J. Fewster, R. Verch, *Algebraic quantum field theory in curved spacetimes*, in: Advances in Algebraic Quantum Field Theory, Springer (2015).
- [KW] B.S. Kay, R.M. Wald, *Theorems on the uniqueness and thermal properties of stationary, nonsingular, quasifree states on spacetimes with a bifurcate Killing horizon*, Phys. Rep. 207, 49 (1991)
- [GGH] V. Georgescu, C. Gérard, D. Häfner, *Asymptotic completeness for superradiant Klein-Gordon equations and applications to the De Sitter Kerr metric*, preprint arXiv:1405.5304, (2014).
- [GHV] J. Gell-Redman, N. Haber, A. Vasy, *The Feynman propagator on perturbations of Minkowski space*, Comm. Math. Phys., 342, 1, (2016), 333–384.
- [GOW] C. Gérard, O. Oulghazi, M. Wrochna, *Hadamard states for the Klein-Gordon equation on Lorentzian manifolds of bounded geometry*, preprint arXiv:1602.00930 (2016).
- [GW1] C. Gérard, M. Wrochna, *Construction of Hadamard states by pseudo-differential calculus*, Comm. Math. Phys. **325** (2) (2014), 713–755.
- [GW2] C. Gérard, M. Wrochna, *Hadamard states for the linearized Yang-Mills equation on curved spacetime*, Comm. Math. Phys. 337 (1) (2015), 253–320.
- [GW3] C. Gérard, M. Wrochna, *Construction of Hadamard states by characteristic Cauchy problem*, Anal. PDE 9 (1), 111–149 (2016).
- [Ho] S. Hollands, *The Hadamard Condition for Dirac Fields and Adiabatic States on Robertson-Walker space-times*, Comm. Math. Phys. **216** (2001), 635–661.
- [Hö] L. Hörmander, *The Analysis of Linear Partial Differential Operators I. Distribution Theory and Fourier Analysis*, Springer (1985).
- [HV] P. Hintz, A. Vasy, *Semilinear wave equations on asymptotically de Sitter, Kerr-de Sitter and Minkowski spacetimes*, to appear in Anal. PDE, arXiv:1306.4705 (2013).
- [HW1] S. Hollands, R.M. Wald, *Existence of local covariant time ordered products of quantum fields in curved spacetime*, Comm. Math. Phys. 231, no. 2 (2002), 309–345.
- [HW2] S. Hollands, R.M. Wald, *Conservation of the stress tensor in perturbative interacting quantum field theory in curved spacetimes*, Rev. Math. Phys. 17.3 (2005), 277–311.
- [HW3] S. Hollands, R.M. Wald, *Quantum fields in curved spacetime*, in: General Relativity and Gravitation: A Centennial Perspective, Cambridge University Press (2015).
- [I-R] G. Idelon–Riton, *Scattering theory for the Dirac equation in Schwarzschild-Anti-de Sitter space-time*, preprint arXiv:1412.0869 (2014).
- [Is] H. Isozaki, *QFT for scalar particles in external fields on Riemannian manifolds*, Rev. Math. Phys. 13, No. 6 (2001), 767–798.
- [Ju] W. Junker, *Adiabatic Vacua and Hadamard States for Scalar Quantum Fields on Curved Spacetime*, PhD thesis, University of Hamburg 1995.
- [JS] W. Junker, E. Schrohe, *Adiabatic Vacuum States on General Space-time Manifolds: Definition, Construction, and Physical Properties*, Ann. Henri Poincaré, **3** (2002), 1113–1181.
- [Ko] Y. Kordyukov, *L^p -theory of elliptic differential operators on manifolds of bounded geometry*, Acta Appl. Math. 23 (1991), 223–260.
- [KM] I. Khavkine, V. Moretti, *Algebraic QFT in curved spacetime and quasifree Hadamard states: an introduction*, in: Advances in Algebraic Quantum Field Theory, Springer (2015).

- [Lu] L.-E. Lundberg, *Relativistic quantum theory for charged spinless particles in external vector fields*, Comm. Math. Phys. 31, (1973), 295–316.
- [Mo1] V. Moretti, *Uniqueness theorem for BMS-invariant states of scalar QFT on the null boundary of asymptotically flat spacetimes and bulk-boundary observable algebra correspondence*, Comm. Math. Phys. **268** (2006), 727–756.
- [Mo2] V. Moretti, *Quantum out-states holographically induced by asymptotic flatness: invariance under space-time symmetries, energy positivity and Hadamard property*, Comm. Math. Phys. **279** (2008), 31–75.
- [Ni] J-P. Nicolas, *Conformal scattering on the Schwarzschild metric*, preprint [arXiv:1312.1386](https://arxiv.org/abs/1312.1386), to appear in Annales de l’Institut Fourier, (2013).
- [Ra] M. Radzikowski, *Micro-local approach to the Hadamard condition in quantum field theory on curved space-time*, Comm. Math. Phys. **179** (1996), 529–553.
- [Ru] S.N.M. Ruijsenaars, *Charged particles in external fields I. Classical theory*, J. Math. Phys. 18 (4), (1977), 720–737.
- [RS] M. Reed, B. Simon, *Methods of modern mathematical physics II : Fourier analysis, self-adjointness*, Academic Press, 1975.
- [RW] M. Ruzhansky, J. Wirth, *Dispersive estimates for t -dependent hyperbolic systems*, J. Differ. Equations 251, (2011), 941–969.
- [Sa] K. Sanders, *Equivalence of the (generalized) Hadamard and microlocal spectrum condition for (generalized) free fields in curved space-time*, Comm. Math. Phys. **295** (2010), 485–501.
- [Se] R. Seiler, *Quantum theory of particles with spin zero and one half in external fields*, Comm. Math. Phys. 25, (1972), 127–151.
- [Sh1] M.A. Shubin, *Spectral theory of elliptic operators on non-compact manifolds*, Astérisque 207 (1992), 37–108.
- [Sh2] M.A. Shubin, *Pseudo-differential operators and spectral theory*, Springer (2001).
- [SV] H. Sahlmann, R. Verch, *Microlocal spectrum condition and Hadamard form for vector-valued quantum fields in curved spacetime*, Rev. Math. Phys., 13(10) (2001), 1203-1246.
- [Va1] A. Vasy, *Microlocal analysis of asymptotically hyperbolic and Kerr-de Sitter spaces, (With an appendix by S. Dyatlov)*, Inventiones Math., (2013), 194.2: 381–513.
- [Va2] A. Vasy, *On the positivity of propagator differences*, [arxiv:1411.7242](https://arxiv.org/abs/1411.7242) (2014).
- [VW] A. Vasy, M. Wrochna, *Quantum fields from global propagators on asymptotically Minkowski and extended de Sitter spacetimes*, preprint [arXiv:1512.08052](https://arxiv.org/abs/1512.08052), (2015).
- [Wa1] R.M. Wald, *Existence of the S-matrix in quantum field theory in curved space-time*, Ann. Phys. (N. Y.), t. 118, (1979), 490–510.
- [Wa2] R.M. Wald, *General Relativity*, University of Chicago Press (1984).
- [Wi] J. Wirth, *On t -dependent hyperbolic systems. Part 2*, preprint [arXiv:1508.02635](https://arxiv.org/abs/1508.02635) (2015).
- [Wr] M. Wrochna, *Singularities of two-point functions in Quantum Field Theory*, PhD thesis, University of Göttingen (2013).

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