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A BOUNDARY THIN OBSTACLE PROBLEM FOR A WAVE EQUATION

BY JONG UHN KIM

Department of Mathematics Virginia Polytechnic Institute & State University Blacksburg, VA 24061

0. Introduction.

In this paper we shall discuss a hyperbolic variational inequality associated with the following initial-boundary value problem:

$$\frac{\partial^2 u}{\partial t^2} - \Delta u = f \text{ in } \Omega \times (0, T), \qquad (0-1)$$

$$u(x,0) = u_0(x), \quad \partial_t u(x,0) = u_1(x) \text{ in } \Omega,$$
 (0-2)

$$u = 0 \text{ on } \Gamma_1 \times (0, T), \tag{0-3}$$

$$u(x,t) \ge \phi(x) \text{ on } \Gamma_2 \times (0,T),$$

$$(0-4)$$

$$\frac{\partial u}{\partial \nu} \ge 0, \quad (u-\phi)\frac{\partial u}{\partial \nu} = 0 \text{ on } \Gamma_2 \times (0,T).$$
 (0-5)

Here Ω is a bounded open subset of \mathbb{R}^n , $n \geq 2$, with smooth boundary $\partial \Omega \in \mathbb{C}^2$. We assume that Γ_1 and Γ_2 are disjoint, nonempty open subsets of $\partial \Omega$ such that $\partial \Omega = \Gamma_1 \cup \overline{\Gamma}_2 = \overline{\Gamma}_1 \cup \Gamma_2$ and $\partial \Gamma_1 = \partial \Gamma_2 \in \mathbb{C}^2$. The above functions $f(x,t), u_0(x), u_1(x)$ and $\phi(x)$ are given, and $\frac{\partial}{\partial \nu}$ denotes the outward normal derivative on $\partial \Omega$. We assume that $\phi \leq 0$ and $u_0 \geq \phi$ on Γ_2 . We can put the above problem in the form of a variational inequality: find u(x,t) which satisfies (0-2), (0-3), (0-4) and

$$<\frac{\partial^2 u}{\partial t^2}, w-u>+a(u,w-u)\geq < f, w-u>$$
(0-6)

for every $w \in G$, for all $t \in (0,T)$. Here we use the notation:

$$< u, w > = \int_{\Omega} uw dx,$$

$$a(u, w) = \int_{\Omega} \nabla u \cdot \nabla w dx,$$

$$G = \{ w \in H^{1}(\Omega) : w = 0 \text{ on } \Gamma_{1} \text{ and } w \ge \phi \text{ on } \Gamma_{2} \}.$$

If a solution u of this new problem is sufficiently smooth, one can show that u also satisfies (0-1)and (0-5) so that u is a solution of the original problem (0-1) through (0-5). The corresponding stationary variational inequality: find $u \in G$ which satisfies

$$a(u, w - u) \ge \langle f, w - u \rangle, \text{ for all } w \in G, \tag{0-7}$$

is called the boundary thin obstacle problem (or the Signorini problem), and its mathematical theory has been well-established: see [6] and references therein. However, to the author's knowledge, nothing has been known about the solution of (0-1) through (0-5). The purpose of this paper is to establish the existence of a (weak) solution of (0-1) through (0-5). This problem is an example of a hyperbolic variational inequality where a constraint is imposed on the unknown function rather than on the time derivative of the unknown function. Hyperbolic variational inequalities have been extensively investigated in the case where a constraint is imposed on the time derivative of the unknown function; see [3] and [11] among others. On the other hand, when a constraint is imposed on the unknown function, existence theorems seem to be known mainly for special problems of space dimension one; see [1], [2], [5], [8], [12], [13] and [14]. In the case of several space dimensions, [10] and this paper seem to be the only known works. [10] resolved the above problem by microlocal techniques when Ω is a half-space; but the same method does not extend to a general domain, as is explained in [10].

In this paper, the basic procedure to prove the existence of a solution consists of a penalty method [11], a multiplier technique [9] and compensated compactness [4]. Finally, we note that the questions of uniqueness and energy conservation for solutions of our problem are still open.

1. Notation and Preliminaries.

$$\partial_t = \frac{\partial}{\partial_t}, \partial_{tt} = \frac{\partial^2}{\partial t^2}, \nabla g = \left(\frac{\partial g}{\partial x_1}, \dots, \frac{\partial g}{\partial x_n}\right).$$

When E is an open subset of $\mathbb{R}^N, N \ge 1, \mathcal{D}'(E)$ denotes the space of distributions over E and $C_0^m(E) = \{w \in C^m(E) : supp w \text{ is a compact subset of } E\}$. We assume that $\partial \Omega \in C^2$ and $\partial \Gamma_1 = \partial \Gamma_2 \in C^2$. For m=nonnegative integer,

$$\begin{split} H^m(\Omega) &= \{ w \in L^2(\Omega) : \left(\frac{\partial}{\partial x_1}\right)^{\alpha_1} \dots \left(\frac{\partial}{\partial x_n}\right)^{\alpha_n} w \in L^2(\Omega), 0 \le \alpha_1 + \dots + \alpha_n \le m \}, \\ H^m_0(\Omega) &= \text{the completion of } C^\infty_0(\Omega) \text{ in } H^m(\Omega), \\ H^{-m}(\Omega) &= \text{the dual of } H^m_0(\Omega), \\ H^1_{\Gamma_1}(\Omega) &= \{ w \in H^1(\Omega) : w = 0 \text{ on } \Gamma_1 \}. \end{split}$$

For s=nonnegative real, $H^{s}(\Omega)$ is defined by interpolation and $H^{-s}(\Omega)$ =the dual of $H^{s}_{0}(\Omega)$. We denote by <, > the duality pairing between $H^{-\alpha}(\Omega)$ and $H^{\alpha}(\Omega), 0 \leq \alpha \leq \frac{1}{2}$. Then, for $g \in L^{2}(\Omega), h \in H^{\alpha}(\Omega), 0 \leq \alpha \leq \frac{1}{2}$, we have < $g, h >= \int_{\Omega} ghdx$.

As above, the bilinear form $a(\ ,\)$ is defined by $a(g,h)=\int_\Omega
abla g\cdot
abla h dx$ and

$$G = \{ w \in H^1_{\Gamma_1}(\Omega) : w \ge \phi \text{ on } \Gamma_2 \}.$$

where $\phi(x) \in C^1(\overline{\Omega})$ and $\phi(x) \leq 0$ on Γ_2 .

Next we define the distance function d by $d(x) = dist(x, \partial \Omega)$ and set

$$V_{\delta} = \{ x \in \overline{\Omega} : d(x) < \delta \}.$$

Let δ_0 be a positive number such that $1/\delta_0$ bounds the principal curvatures of $\partial\Omega$. Since Ω is bounded and $\partial\Omega \in C^2$, it is known that $d(x) \in C^2(V_{\delta_0})$; see [7]. We also set

$$U_{\delta} = \{x \in \Omega : d(x) < \delta\},\$$

$$S_{\delta} = \{x \in \Omega : d(x) = \delta\}.$$

Then, S_{δ} is a C^2 -surface for $0 \leq \delta < \delta_0$. We shall use the following lemmas later.

LEMMA 1.1. There is a function $h(x) \in C^1(\overline{\Omega})^n$ such that h(x)=the outward unit normal vector at x on S_δ if $d(x) = \delta \leq \delta_0/2$ and h(x) = 0 if $d(x) \geq \delta_0$.

PROOF: For each $x \in S_{\delta}, 0 \leq \delta < \delta_0$, there is a unique point $y(x) \in \partial\Omega$ such that $\operatorname{dist}(x, y(x)) = \delta$ and the outward normal vector at y(x) on $\partial\Omega$ is also outward normal at x on S_{δ} . Furthermore, $y(x) \in C^1(V_{\delta_0})$. Next we denote by $\nu(y)$ the outward unit normal vector at y on $\partial\Omega$. Then, $\nu(y)$ is a C^1 function of y. Let us choose a function $\rho(t) \in C^1(R)$ such that $\rho(t) = 1$ for $t \leq \delta_0/2$ and $\rho(t) = 0$ for $t > \frac{2}{3}\delta_0$. We then set $h(x) = \nu(y(x))\rho(d(x))$, which achieves our goal.

The next lemma is known and stated without proof.

LEMMA 1.2. Let $g \in L^1(\Omega)$. Then, for $0 < \delta^* < \delta_0$,

$$\int_{U_{\delta^*}} gdx = \int_0^{\delta^*} \int_{S_{\delta}} gd\sigma_{\delta} d\delta \tag{1-1}$$

holds where $d\sigma_{\delta}$ is the surface element on S_{δ} .

Next we write a well-known consequence of the Div-Curl lemma of compensated compactness.

LEMMA 1.3. Let $\{u_k\}$ be a sequence of functions such that

$$\|u_k\|_{L^2(0,T;H^1(\Omega))} \le M,\tag{1-2}$$

$$\|\partial_t u_k\|_{L^2(0,T;L^2(\Omega))} \le M,\tag{1-3}$$

where M is a positive constant independent of k and each u_k satisfies

$$\frac{\partial^2 u_k}{\partial t^2} - \Delta u_k = g_k \tag{1-4}$$

where $\{g_k\}$ is a bounded sequence in $L^2(\Omega \times (0,T))$, and, as $k \to \infty$,

$$u_k \to u$$
 weakly in $L^2(0,T; H^1(\Omega))$. (1-5)

Then, as $k \to \infty$

$$\left(\frac{\partial u_k}{\partial t}\right)^2 - |\nabla u_k|^2 \to \left(\frac{\partial u}{\partial t}\right)^2 - |\nabla u|^2 \tag{1-6}$$

in $\mathcal{D}'(\Omega \times (0,T))$.

PROOF: Set $v_{k_i} = \frac{\partial u}{\partial x_i}, w_{k_i} = \frac{\partial u}{\partial x_i}, i = 1, \dots, n$, and $v_{k(n+1)} = -\frac{\partial u}{\partial t}, w_{k(n+1)} = \frac{\partial u}{\partial t}$. Then, apply Corollary 4.3 of [4].

Finally, we state

LEMMA 1.4. Let $\{u_k\}$ be a sequence of functions such that as $k \to \infty$,

$$u_k \to u \ weak^* \ in \ L^{\infty}(0,T; H^{\beta}(\Omega))$$
 (1-7)

$$\partial_t u_k \to \partial_t u \text{ weakly in } L^2(0,T; H^{\alpha}(\Omega))$$
 (1-8)

where $-1 \leq \alpha < \beta \leq 1$. Then, we have

$$u_k \to u \text{ in } C([0,T]; H^r(\Omega))$$
(1-9)

for any $r < \beta$.

PROOF: A special case was proved in [8]. The same argument can still be employed. Without loss of generality, we may assume $\alpha \neq -\frac{1}{2}, \beta \neq -\frac{1}{2}$, because we can take smaller α or β if necessary.

Let $r = \theta \alpha + (1 - \theta)\beta$, $0 < \theta < 1$, $r \neq -\frac{1}{2}, \frac{1}{2}$. There is $E \subset [0, T]$ such that meas(E) = 0 and for all $t_1, t_2 \in [0, T] \setminus E$, $t_1 < t_2$,

$$\begin{aligned} \|u_{k}(t_{2}) - u_{k}(t_{1})\|_{H^{r}(\Omega)} &\leq \\ &\leq M_{\theta} \|u_{k}(t_{2}) - u_{k}(t_{1})\|_{H^{\alpha}(\Omega)}^{\theta} \|u_{k}(t_{2}) - u_{k}(t_{1})\|_{H^{\beta}(\Omega)}^{1-\theta}, \\ &\leq M_{\theta} (\int_{t_{1}}^{t_{2}} \|\partial_{t}u_{k}(t)\|_{H^{\alpha}(\Omega)} dt)^{\theta} \leq M_{\theta}(t_{2} - t_{1})^{\theta/2}, \end{aligned}$$
(1-10)

holds for all k and $0 < \theta < 1$ such that $\theta \alpha + (1-\theta)\beta \neq \frac{1}{2}, -\frac{1}{2}$, where M_{θ} denotes positive constants independent of t_1, t_2 and k. Consequently, each u_k belongs to $C([0,T]; H^r(\Omega))$ for every $\alpha < r < \beta$ possibly after a modification on E. Since the embedding $H^{r_2}(\Omega) \subset H^{r_1}(\Omega)$ is compact for every $r_1 < r_2$, we can use the Ascoli theorem to arrive at the conclusion.

2. Approximate Problem.

In this section we shall consider an approximate problem with a penalty parameter k > 0. The problem is: find a function $u_k(x,t)$ such that

$$\partial_{tt}u_k(x,t) \in L^{\infty}(0,T;L^2(\Omega)), \partial_t u_k \in L^{\infty}(0,T;H^1_{\Gamma_1}(\Omega))$$
(2-1)

$$u_k(x,0) = u_{0k}(x), \partial_t u_k(x,0) = u_{1k}(x)$$
(2-2)

$$<\partial_{tt}u_k, w>+a(u_k,w)=+\int_{\Gamma_2}(k(u_k-\phi)^--\frac{1}{k}\partial_tu_k)wd\sigma$$
(2-3)

holds for all $w \in H^1_{\Gamma_1}(\Omega)$, for almost all $t \in (0,T)$. Here $d\sigma$ is the surface element on $\partial\Omega$ and $g^- = \max(0, -g)$.

Our strategy consists of three steps.

- (Step 1) Establish the existence of solution of (2-1), (2-2) and (2-3).
- (Step 2) Obtain a priori estimates independent of k > 0.
- (Step 3) Pass $k \to \infty$ to arrive at a solution of the original problem.

In this section, the first two steps will be carried out.

2.1 Existence of solution

Throughout this subsection, we shall suppress the subscript k.

PROPOSITION 2.1. Let $u_0, u_1 \in C_0^{\infty}(\Omega), f \in C^1([0,T]; L^2(\Omega))$ and $\phi \in C^1(\overline{\Omega})$ with $\phi \leq 0$ on Γ_2 . Then there is a solution of (2-1), (2-2) and (2-3).

PROOF: We set $u = v + u_0 + tu_1$. Then,

$$v(x,0) \equiv 0, \partial_t v(x,0) \equiv 0, \tag{2-4}$$

and (2-3) is equivalent to

$$\langle \partial_{tt}v, w \rangle + a(v, w) = \langle f, w \rangle$$

$$+ \langle \Delta(u_0 + tu_1), w \rangle + k \int_{\Gamma_2} (v - \phi)^- w d\sigma$$

$$- \frac{1}{k} \int_{\Gamma_2} (\partial_t v) w d\sigma,$$

$$(2-5)$$

for all $w \in H^1_{\Gamma_1}(\Omega)$, for almost all $t \in (0,T)$. We shall employ the Galerkin approximation procedure. Let $\{w_j\}_{j=1}^{\infty}$ be a basis for $H^1_{\Gamma_1}(\Omega)$ where each $w_j \in C^{\infty}(\bar{\Omega}) \cap H^1_{\Gamma_1}(\Omega)$. Let us set

$$v_m = \sum_{j=1}^m a_{mj}(t) w_j(x)$$
 (2-6)

where each $a_{mj}(t)$ is determined from

$$a_{mj}(0) = \dot{a}_{mj}(0) = 0, \qquad (2-7)$$

and

$$\langle \partial_{tt}v_m, w_j \rangle + a(v_m, w_j) = \langle f, w_j \rangle$$

$$+ \langle \Delta(u_0 + tu_1), w_j \rangle + k \int_{\Gamma_2} (v_m - \phi)^- w_j d\sigma$$

$$- \frac{1}{k} \int_{\Gamma_2} (\partial_t v_m) w_j d\sigma, \quad j = 1, \dots, m.$$

$$(2-8)$$

For each *m*, there is a unique set of functions $a_{mj}(t) \in C^2([0,T]), j = 1, ..., m$, which satisfy (2-7) and (2-8). We next obtain estimates independent of *m*. Let us substitute $\partial_t v_m$ for w_j in (2-8):

$$\frac{1}{2} \frac{d}{dt} \|\partial_t v_m\|_{L^2(\Omega)}^2 + \frac{1}{2} \frac{d}{dt} a(v_m, v_m) \\ + \frac{1}{2} k \frac{d}{dt} \|(v_m - \phi)^-\|_{L^2(\Gamma_2)}^2$$

$$= \langle f, \partial_t v_m \rangle + \langle \Delta(u_0 + tu_1), \partial_t v_m \rangle - \frac{1}{k} \|\partial_t v_m\|_{L^2(\Gamma_2)}^2.$$
(2-9)

Note that $v_m - \phi \in C^1([0,T] \times \overline{\Omega})$ and hence, $(v_m - \phi)^-$ is Lipschitz continuous and

$$\frac{d}{dt} \| (v_m - \phi)^- \|_{L^2(\Gamma_2)}^2 \in L^{\infty}(0, T).$$

It follows from (2-7) and (2-9) that

$$\|\partial_t v_m\|_{L^2(\Omega)} \le M, \text{ for all } t \in [0,T]$$
(2-10)

$$\|v_m\|_{H^1_{\Gamma_1}(\Omega)} \le M$$
, for all $t \in [0,T]$ (2-11)

$$\|(v_m - \phi)^-\|_{L^2(\Gamma_2)} \le M, \text{ for all } t \in [0, T]$$
(2-12)

$$\int_0^T \|\partial_t v_m\|_{L^2(\Gamma_2)}^2 dt \le M,$$
(2-13)

where M denotes positive constants independent of m. Next we differentiate (2-8) with respect to t and substitute $\partial_{tt}v_m$ for w_j :

$$< \partial_{ttt} v_m, \delta_{tt} v_m > + a(\partial_t v_m, \partial_{tt} v_m)$$

$$= < \partial_t f, \partial_{tt} v_m > + < \Delta u_1, \partial_{tt} v_m >$$

$$+ k \int_{\Gamma_2} \partial_t \{ (v_m - \phi)^- \} \partial_{tt} v_m d\sigma$$

$$- \frac{1}{k} \int_{\Gamma_2} (\partial_{tt} v_m)^2 d\sigma,$$

$$(2-14)$$

which yields

$$\frac{1}{2} \frac{d}{dt} \|\partial_{tt} v_m\|_{L^2(\Omega)}^2 + \frac{1}{2} \frac{d}{dt} a(\partial_t v_m, \partial_t v_m) \\
\leq \|\partial_t f\|_{L^2(\Omega)} \|\partial_{tt} v_m\|_{L^2(\Omega)} + \|\Delta u_1\|_{L^2(\Omega)} \|\partial_{tt} v_m\|_{L^2(\Omega)} \\
+ \frac{1}{2} k^3 \|\partial_t \{(v_m - \phi)^-\}\|_{L^2(\Gamma_2)}^2 - \frac{1}{2k} \|\partial_{tt} v_m\|_{L^2(\Gamma_2)}^2.$$
(2-15)

We consider (2-8) at t = 0 and substitute $\partial_{tt} v_m(0)$ for w_j :

$$\|\partial_{tt}v_m(0)\|_{L^2(\Omega)}^2 = \langle f(0), \partial_{tt}v_m(0) \rangle + \langle \Delta u_0, \partial_{tt}v_m(0) \rangle.$$
(2-16)

We have used the fact that $v_m(0) = \partial_t v_m(0) = 0$. Now (2-16) implies

$$\left\|\partial_{tt}v_m(0)\right\|_{L^2(\Omega)} \le M,\tag{2-17}$$

M being a positive constant independent of m. Combining (2-13), (2-15) and (2-17), we derive

$$\|\partial_{tt}v_m\|_{L^2(\Omega)} \le M, \text{ for all } t \in [0,T],$$

$$(2-18)$$

$$\|\partial_t v_m\|_{H^1_{\Gamma_1}(\Omega)} \le M$$
, for all $t \in [0, T]$, (2-19)

where M denotes positive constants independent of m. By means of (2-11), (2-18) and (2-19), we can extract a subsequence still denoted by $\{v_m\}$ such that

$$v_m \to v \ weak * \ in \ L^{\infty}(0,T; H^1_{\Gamma_1}(\Omega))$$
 (2-20)

$$\partial_t v_m \to \partial_t v \text{ weak } * \text{ in } L^{\infty}(0,T; H^1_{\Gamma_1}(\Omega))$$
 (2-21)

$$\partial_{tt} v_m \to \partial_{tt} v \ weak * \ in \ L^{\infty}(0,T;L^2(\Omega))$$
 (2-22)

for some function v. By virtue of Lemma 1.4, we have

$$(v_m - \phi)^- \to (v - \phi)^- \text{ in } C([0, T]; L^2(\partial\Omega)).$$

$$(2-23)$$

We then make use of (2-20) through (2-23) to derive (2-4) and (2-5) through a standard procedure.

Remark 2.2. The purpose of the term $\frac{1}{k}\partial_t u_k$ in (2-3) is to obtain approximate solutions which are regular enough to justify manipulations in using a multiplier technique in the next subsection. The utility of that term is seen in (2-15).

2.2 A priori estimates independent of k.

We still suppress the subscript k and let u(x, t) be a solution constructed in Proposition 2.1 for a fixed k > 0. We can substitute $\partial_t u$ for w in (2-3):

$$\begin{split} &\frac{1}{2}\frac{d}{dt}\|\partial_{t}u\|_{L^{2}(\Omega)}^{2}+\frac{1}{2}\frac{d}{dt}a(u,u)=\\ &-\frac{1}{2}k\frac{d}{dt}\int_{\Gamma_{2}}\{(u-\phi)^{-}\}^{2}d\sigma-\frac{1}{k}\int_{\Gamma_{2}}(\partial_{t}u)^{2}d\sigma, \end{split} \tag{2-24}$$

for almost all $t \in (0,T)$. We now derive from (2-24)

$$\|\partial_t u\|_{L^2(\Omega)} \le M, \text{ for all } t \in [0, T]$$

$$(2-25)$$

$$\|u\|_{H^1_{\Gamma_1}(\Omega)} \le M, \text{ for all } t \in [0,T]$$

$$(2-26)$$

$$\|(u-\phi)^{-}\|_{L^{2}(\Gamma_{2})} \leq M/\sqrt{k}, \text{ for all } t \in [0,T],$$
(2-27)

$$\int_{0}^{T} \|\partial_{t}u\|_{L^{2}(\Gamma_{2})}^{2} dt \leq Mk, \qquad (2-28)$$

where M denotes positive constants which are independent of k, and depend only on the number L such that $\|u_0\|_{H^1_0(\Omega)} + \|u_1\|_{L^2(\Omega)} + \|f\|_{L^2(\Omega \times [0,T))} \leq L$.

Next we shall adapt a multiplier technique to derive some uniform estimates in a neighborhood of $\partial\Omega$. Let us set

$$\Omega_{\delta} = \{x \in \Omega : \operatorname{dist}(x, \partial \Omega) > \delta\}$$

for $0 < \delta \leq \delta_0$. Obviously, $\partial \Omega_{\delta} = S_{\delta}$, which was defined in Section 1. We need to observe that solution u of (2-1), (2-2) and (2-3) has sufficient regularity to justify manipulations in using a multiplier technique. First of all, (2-3) implies that

$$\partial_{tt}u - \Delta u = f \tag{2-29}$$

holds in $\mathcal{D}'(\Omega \times (0,T))$. By virtue of (2-1) and the assumption that $f \in C^1([0,T]; L^2(\Omega))$, we infer that

$$\Delta u \in L^{\infty}(0,T;L^{2}(\Omega)). \tag{2-30}$$

In fact, (2-29) holds in $L^{\infty}(0,T; L^{2}(\Omega))$. Now (2-30) implies

$$u \in L^{\infty}(0,T; H^2(\Omega_{\delta})), \tag{2-31}$$

for each $\delta > 0$. Because of the mixed boundary condition on $\partial\Omega$, one cannot claim that $u \in L^{\infty}(0,T; H^2(\Omega))$. We also note that (2-1) implies

$$u \in C([0,T]; H^1_{\Gamma_1}(\Omega)), \partial_t u \in C([0,T]; L^2(\Omega)),$$
 (2-32)

possibly after a modification on a set of measure zero in t.

We now proceed to obtain the estimates. Using the function h(x) constructed in Lemma 1.1, we multiply both sides of (2-29) by $(h \cdot \nabla)u$ and integrate over $\Omega_{\delta} \times (0,T), 0 < \delta \leq \delta_0/2$:

$$\int_0^T \int_{\Omega_{\delta}} (\partial_{tt} u)((h \cdot \nabla) u) dx dt$$

=
$$\int_{\Omega_{\delta}} (\partial_{t} u(T))((h \cdot \nabla) u(T)) dx - \int_{\Omega_{\delta}} u_1((h \cdot \nabla) u_0) dx$$

$$- \int_0^T \int_{\Omega_{\delta}} (\partial_{t} u)((h \cdot \nabla) \partial_{t} u) dx dt =$$

(applying the divergence theorem to the last integral)

$$= \int_{\Omega_{\delta}} (\partial_{t} u(T))((h \cdot \nabla)u(T))dx - \int_{\Omega_{\delta}} u_{1}((h \cdot \nabla)u_{0})dx$$
$$- \frac{1}{2} \int_{0}^{T} \int_{S_{\delta}} (\partial_{t} u)^{2} d\sigma_{\delta} dt + \frac{1}{2} \int_{0}^{T} \int_{\Omega_{\delta}} (\nabla \cdot h)(\partial_{t} u)^{2} dx dt.$$
(2-33)

$$\begin{split} &\int_{0}^{T} \int_{\Omega_{\delta}} (-\Delta u)((h \cdot \nabla)u) dx dt \\ &= -\int_{0}^{T} \int_{S_{\delta}} \left(\frac{\partial u}{\partial \nu}\right)^{2} d\sigma_{\delta} dt + \int_{0}^{T} \int_{\Omega_{\delta}} \nabla u \cdot \nabla \{(h \cdot \nabla)u\} dx dt \\ &= -\int_{0}^{T} \int_{S_{\delta}} \left(\frac{\partial u}{\partial \nu}\right)^{2} d\sigma_{\delta} dt + \sum_{i,j=1}^{n} \int_{0}^{T} \int_{\Omega_{\delta}} \frac{\partial u}{\partial z_{i}} \frac{\partial h_{j}}{\partial z_{i}} \frac{\partial u}{\partial z_{j}} dx dt \\ &- \int_{0}^{T} \int_{\Omega_{\delta}} \frac{1}{2} (\nabla \cdot h) |\nabla u|^{2} dx dt + \frac{1}{2} \int_{0}^{T} \int_{S_{\delta}} |\nabla u|^{2} d\sigma_{\delta} dt, \end{split}$$
(2-34)

where $\frac{\partial}{\partial \nu}$ is the unit outward normal derivative on S_{δ} , i.e., $\frac{\partial}{\partial \nu} = h \cdot \nabla$ on S_{δ} . Combining (2-33) and (2-34), we get

$$\frac{1}{2} \int_{0}^{T} \int_{S_{\delta}} (\partial_{t} u)^{2} d\sigma_{\delta} dt - \frac{1}{2} \int_{0}^{T} \int_{S_{\delta}} |\nabla u|^{2} d\sigma_{\delta} dt + \int_{0}^{T} \int_{S_{\delta}} \left(\frac{\partial u}{\partial \nu}\right)^{2} d\sigma_{\delta} dt = C(\delta),$$
(2-35)

where

$$C(\delta) = -\int_0^T \int_{\Omega_{\delta}} f((h \cdot \nabla)u) dx dt$$

+ obvious terms in (2-33) and (2-34). (2-36)

Recalling (2-25) and (2-26), it is evident that

$$|C(\delta)| \le M, \text{ for all } 0 < \delta \le \delta_0/2, \tag{2-37}$$

where M is a positive constant which is independent of k and depends only on the number L such that

$$||u_0||_{H^1_0(\Omega)} + ||u_1||_{L^2(\Omega)} + ||f||_{L^2(\Omega \times (0,T))} \le L.$$

We next choose any $\rho(s) \in C(R)$ such that for all $s, 0 \le |\rho(s)| \le 1$, and set $\psi(x) = \rho(\operatorname{dist}(x, \partial\Omega))$. Then, $\psi(x)$ is a continuous function and is constant on each $S_{\delta}, 0 < \delta \le \delta_0$. Hence, it follows from (2-35) that

$$\begin{split} &\frac{1}{2}\int_0^T \int_0^{\delta^*} \int_{S_{\delta}} \psi(x)(\partial_t u)^2 d\sigma_{\delta} d\delta dt \\ &-\frac{1}{2}\int_0^T \int_0^{\delta^*} \int_{S_{\delta}} \psi(x) |\nabla u|^2 d\sigma_{\delta} d\delta dt \\ &+\int_0^T \int_0^{\delta^*} \int_{S_{\delta}} \psi(x)((h\cdot\nabla)u)^2 d\sigma_{\delta} d\delta dt = \end{split}$$

(using Lemma 1.2)

$$= \frac{1}{2} \int_0^T \int_{U_{\delta^*}} \psi(x) (\partial_t u)^2 dx dt - \frac{1}{2} \int_0^T \int_{U_{\delta^*}} \psi(x) |\nabla u|^2 dx dt$$
$$+ \int_0^T \int_{U_{\delta^*}} \psi(x) ((h \cdot \nabla) u)^2 dx dt = 0(\delta^*), \tag{2-38}$$

for each $0 < \delta^* \leq \delta_0/2$, where $0(\delta^*)$ is a number such that

$$|0(\delta^*)| \le M\delta^*, \tag{2-39}$$

where M is the same positive number as in (2-37). (2-38) will be used in the next section.

3. Convergence of approximate solutions.

We assume that $u_0 \in H^1_0(\Omega), u_1 \in L^2(\Omega), f \in L^2(\Omega \times (0,T)), \phi \in C^1(\overline{\Omega})$ with $\phi \leq 0$ on Γ_2 , and define a solution of (0-1) through (0-5) by

<u>Definition 3.1</u> A function u(x,t) is a solution of (0-1) through (0-5) if

$$u \in L^{\infty}(0,T; H^{1}_{\Gamma_{1}}(\Omega)), \partial_{t}u \in L^{\infty}(0,T; L^{2}(\Omega)) \cap C([0,T]; H^{-1/2}(\Omega)),$$
(3-1)

$$u(x,0) = u_0(x), \partial_t u(x,0) = u_1(x), \tag{3-2}$$

 $u(t) \in G$, for almost all $t \in (0, T)$ (3-3)

$$<\partial_t u(T), w(T) - u(T) > - < u_1, w(0) - u_0 >$$

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$$-\int_{0}^{T} \langle \partial_{t}u, \partial_{t}w - \partial_{t}u \rangle dt + \int_{0}^{T} a(u, w - u)dt$$

$$\geq \int_{0}^{T} \langle f, w - u \rangle dt, \qquad (3-4)$$

for every w such that $w \in L^{\infty}(0,T; H^{1}_{\Gamma_{1}}(\Omega)), \partial_{t}w \in L^{\infty}(0,T; L^{2}(\Omega)), w(t) \in G$, for almost all $t \in (0,T)$.

We shall justify this definition. Suppose that u is a solution according to the above definition. By taking $w = u \pm \xi, \xi \in C_0^{\infty}(\Omega \times (0,T))$, in (3-4), we find that (0-1) is satisfied in $\mathcal{D}'(\Omega \times (0,T))$. (0-3) and (0-4) are satisfied in the sense $u(t) \in G$, for almost all $t \in (0,T)$. Now we assume further regularity on u:

$$\partial_{tt} u \in L^2(0,T; L^2(\Omega \setminus E_\delta)) \tag{3-5}$$

$$u \in L^2(0,T; H^2(\Omega \setminus E_\delta))$$
(3-6)

for each $\delta > 0$, where we set

$$E_{\delta} = \{ x \in \overline{\Omega} : \operatorname{dist}(x, \partial \Gamma_1) \leq \delta \}.$$

Recall that $\partial \Gamma_1 = \partial \Gamma_2$ is an (n-2)-dimensional, C^2 -manifold by assumption. Then, (0-1) holds in a stronger sense:

$$\int_{0}^{T} <\partial_{tt}u, w > dt - \int_{0}^{T} <\Delta u, w > dt = \int_{0}^{T} < f, w > dt$$
(3-7)

holds for every $w \in L^2(\Omega \times (0,T))$ such that $supp \ w(x,t) \cap E_{\delta}$ is empty, for almost all t, for some $\delta > 0$.

Next choose any nonnegative function $\xi(x) \in C_0^1(\Gamma_2)$. Then, we can extend $\xi(x)$ so that $\xi(x) \in C^1(\bar{\Omega})$ and $supp \ \xi \cap \{E_\delta \cup \Gamma_1\}$ is empty for some $\delta > 0$. Choose any nonnegative function $\eta(t) \in C_0^1((0,T))$ and take $w = u + \xi(x)\eta(t)$ in (3-4). By means of integration by parts and the divergence theorem (this procedure is justified by (3-5) and (3-6)), we find

$$\int_{0}^{T} <\partial_{tt}u, \xi(x)\eta(t) > dt - \int_{0}^{T} <\Delta u, \xi(x)\eta(t) > dt$$
$$+ \int_{0}^{T} \int_{\Gamma_{2}} \left(\frac{\partial u}{\partial \nu}\right) \xi(x)\eta(t) d\sigma dt \ge \int_{0}^{T} < f, \xi(x)\eta(t) > dt.$$
(3-8)

It follows from (3-7) and (3-8) that

$$\int_{0}^{T} \int_{\Gamma_{2}} \left(\frac{\partial u}{\partial \nu} \right) \xi(x) \eta(t) d\sigma dt \ge 0,$$
(3-9)

which is true for all nonnegative functions $\xi(x) \in C_0^1(\Gamma_2)$ and $\eta(t) \in C_0^1((0,T))$. Hence we can conclude that

$$\frac{\partial u}{\partial \nu} \ge 0$$
, almost everywhere in $\Gamma_2 \times (0,T)$. (3-10)

It remains to show the second part of (0.5). Let us choose a function $\xi_{\delta}(x) \in C^1(\overline{\Omega})$ such that $0 \leq \xi_{\delta}(x) \leq 1$, for all $x \in \overline{\Omega}, \xi_{\delta} \equiv 1$ on $\Gamma_2 \setminus E_{\delta}$ and $supp \ \xi_{\delta} \cap \{E_{\delta/2} \cup \Gamma_1\}$ is empty. We also choose any $\eta(t) \in C_0^1((0,T))$ such that $-1 \leq \eta(t) \leq 1$ for all t. Then, we set

$$w = u + \xi_{\delta}(x)\eta(t)(u - \phi)$$

in (3.4). As above, we use integration by parts, the divergence theorem and (3-7) to arrive at

$$\int_{0}^{T} \int_{\Gamma_{2}} \xi_{\delta}(x) \eta(t) \left(\frac{\partial u}{\partial \nu}\right) (u - \phi) d\sigma dt = 0.$$
(3-11)

Here, we have the equality, because $-\eta(t)$ can be also chosen in place of $\eta(t)$. Since $u - \phi \ge 0$ almost everywhere in $\Gamma_2 \times (0,T)$ and $\delta > 0$ is arbitrary, (3-10) and (3-11) imply

$$\left(\frac{\partial u}{\partial \nu}\right)(u-\phi) = 0, \text{ almost everywhere in } \Gamma_2 \times (0,T).$$
(3-12)

This completes the justification of the above definition of a solution.

We now proceed to present the main result.

THEOREM 3.2. For $u_0 \in H^1_0(\Omega), u_1 \in L^2(\Omega)$ and $f \in L^2(\Omega \times (0,T))$, there is a solution of (0-1) through (0-5).

PROOF: Choose sequences $\{u_{0k}\}, \{u_{1k}\}$ in $C_0^{\infty}(\Omega)$ and f_k in $C^1([0,T]; L^2(\Omega))$ such that as $k \to \infty$,

$$u_{0k} \to u_0 \text{ in } H^1_0(\Omega), \tag{3-13}$$

$$u_{1k} \to u_1 \text{ in } L^2(\Omega), \qquad (3-14)$$

$$f_k \to f \text{ in } L^2(\Omega \times (0,T)).$$
 (3-15)

Using these, for each k, let u_k be a solution constructed in Proposition 2.1. Then, each u_k satisfies

$$\int_{0}^{T} < \partial_{tt} u_{k}, w - u_{k} > dt + \int_{0}^{T} a(u_{k}, w - u_{k}) dt$$

= $\int_{0}^{T} < f, w - u_{k} > dt + k \int_{0}^{T} \int_{\Gamma_{2}} (u_{k} - \phi)^{-} (w - u_{k}) d\sigma dt$
 $- \frac{1}{k} \int_{0}^{T} \int_{\Gamma_{2}} (\partial_{t} u_{k}) (w - u_{k}) d\sigma dt,$ (3-16)

for every w(x,t) such that $w \in L^{\infty}(0,T; H^{1}_{\Gamma_{1}}(\Omega)), \partial_{t}w \in L^{\infty}(0,T; L^{2}(\Omega))$ and $w(t) \in G$, for almost all $t \in (0,T)$. Using integration by parts and

$$\int_{0}^{T} \int_{\Gamma_{2}} (u_{k} - \phi)^{-} (w - u_{k}) d\sigma dt$$

=
$$\int_{0}^{T} \int_{\Gamma_{2}} (u_{k} - \phi)^{-} (w - \phi) d\sigma dt - \int_{0}^{T} \int_{\Gamma_{2}} (u_{k} - \phi)^{-} (u_{k} - \phi) d\sigma dt$$

$$\geq 0, \qquad 12 \qquad (3-17)$$

we obtain

$$<\partial_{t}u_{k}(T), w(T) - u_{k}(T) > - < u_{1k}, w(0) - u_{0k} >$$

$$-\int_{0}^{T} <\partial_{t}u_{k}, \partial_{t}w > dt + \int_{0}^{T} \int_{\Omega} (\partial_{t}u_{k})^{2}dxdt$$

$$+\int_{0}^{T} a(u_{k}, w)dt - \int_{0}^{T} \int_{\Omega} |\nabla u_{k}|^{2}dxdt$$

$$\geq \int_{0}^{T} < f, w - u_{k} > dt - \frac{1}{k} \int_{0}^{T} \int_{\Gamma_{2}} (\partial_{t}u_{k})wd\sigma dt$$

$$+ \frac{1}{2k} \int_{\Gamma_{2}} (u_{k}(T))^{2}d\sigma.$$

$$(3-18)$$

By virtue of (2-25) and (2-26), we can extract a subsequence still denoted by $\{u_k\}$ such that as $k \to \infty$,

$$u_k \to u \text{ weak } * \text{ in } L^{\infty}(0,T; H^1_{\Gamma_1}(\Omega)), \qquad (3-19)$$

$$\partial_t u_k \to \partial_t u \ weak * \ in \ L^{\infty}(0,T;L^2(\Omega)).$$
 (3-20)

Since each u_k satisfies (2-29), we infer from (3-15) and (3-19) that

$$\partial_{tt} u_k \rightarrow \partial_{tt} u \text{ weakly in } L^2(0,T; H^{-1}(\Omega)).$$
 (3-21)

Hence, by Lemma 1.4, we find that

$$\partial_t u_k \to \partial_t u \text{ in } C([0,T]; H^{-\frac{1}{2}}(\Omega)),$$
 (3-22)

$$u_k \rightarrow u \text{ in } C([0,T]; H^{\frac{1}{2}}(\Omega)),$$

$$(3-23)$$

so that

$$<\partial_t u_k(T), w(T) - u_k(T) > \rightarrow <\partial_t u(T), w(T) - u(T) >$$
(3-24)

as $k \to \infty$, for each w. In the meantime, we notice that

$$\begin{split} &\frac{1}{k} |\int_0^T \int_{\Gamma_2} (\partial_t u_k) w d\sigma dt| \\ &\leq \frac{1}{k} (\int_0^T \int_{\Gamma_2} (\partial_t u_k)^2 d\sigma dt)^{\frac{1}{2}} (\int_0^T \int_{\Gamma_2} w^2 d\sigma dt)^{\frac{1}{2}}, \\ &(\text{using (2-28))}, \\ &\leq M (\int_0^T \int_{\Gamma_2} w^2 d\sigma dt)^{\frac{1}{2}} / \sqrt{k} \to 0, \end{split}$$
(3-25)

as $k \to \infty$, for each w. Next we shall show that

$$\int_0^T \int_\Omega ((\partial_t u)^2 - |\nabla u|^2) dx dt \ge \limsup_{k \to \infty} \sup \int_0^T \int_\Omega ((\partial_t u_k)^2 - |\nabla u_k|^2) dx dt$$
(3-26)

It is convenient to introduce the notation:

$$\begin{split} A_k &= \frac{1}{2} (\partial_t u_k)^2 - \frac{1}{2} |\nabla u_k|^2 + ((h \cdot \nabla) u_k)^2, \\ B_k &= ((h \cdot \nabla) u_k)^2, \\ A &= \frac{1}{2} (\partial_t u)^2 - \frac{1}{2} |\nabla u|^2 + ((h \cdot \nabla) u)^2, \\ B &= ((h \cdot \nabla) u)^2. \end{split}$$

First of all, by Lemma 1.3, as $k \to \infty$,

$$A_k - B_k \to A - B \text{ in } \mathcal{D}'(\Omega \times (0,T)).$$
 (3-27)

Since the sequence $\{A_k - B_k\}$ is bounded in $L^{\infty}(0,T;L^1(\Omega))$, we observe that as $k \to \infty$,

$$\int_0^T \int_\Omega \psi(x) (A_k - B_k) dx dt \to \int_0^T \int_\Omega \psi(x) (A - B) dx dt$$
(3-28)

for each $\psi(x) \in C_0(\Omega)$. Let us define a function $\rho_{\delta}(s) \in C(R)$ such that $\rho_{\delta}(s)$ is nondecreasing in s and

$$\rho_{\delta}(s) = \begin{cases} 0 & \text{for all } s \le \delta/2 \\ 1 & \text{for all } s \ge \delta. \end{cases}$$
(3-29)

We then set $\psi_{\delta}(x) = \rho_{\delta}(\operatorname{dist}(x,\partial\Omega))$. Fix any $0 < \delta \leq \delta_0/2$. According to (3-28),

$$\int_0^T \int_\Omega \psi_\delta(x) (A_k - B_k) dx dt \to \int_0^T \int_\Omega \psi_\delta(x) (A - B) dx dt$$
(3-30)

as $k \to \infty$. Furthermore, with the aid of (2-38), we notice that

$$\int_{0}^{T} \int_{\Omega} (\psi_{\delta}(x) - 1) A_{k} dx dt = \int_{0}^{T} \int_{U_{\delta}} (\psi_{\delta}(x) - 1) A_{k} dx dt = 0_{1}(\delta),$$
(3-31)

where we can estimate $0_1(\delta)$ by using the number M in (2-37), and (3-13), (3-14) and (3-15),

$$|0_1(\delta)| \le M\delta$$
, for all k. (3-32)

Meanwhile we also have

$$\int_{0}^{T} \int_{\Omega} B_{k} dx dt \geq \int_{0}^{T} \int_{\Omega} \psi_{\delta}(x) B_{k} dx dt, \qquad (3-33)$$

for each k, since B_k is nonnegative and integrable, and $0 \le \psi_{\delta}(x) \le 1$, for all x. Since $A - B \in L^{\infty}(0,T; L^1(\Omega))$, it is apparent that

$$\int_0^T \int_\Omega (1 - \psi_\delta(x))(A - B) dx dt = \int_0^T \int_{U_\delta} (1 - \psi_\delta(x))(A - B) dx dt$$

= $0_2(\delta) \to 0$ as $\delta \to 0$. (3.34)

Choose any $\epsilon > 0$. Then, by (3-32) and (3-34), there is $0 < \tilde{\delta}(\epsilon) \le \delta_0/2$ such that

$$|\mathbf{0}_1(\delta)| + |\mathbf{0}_2(\delta)| \le \epsilon, \tag{3-35}$$

for all $0 < \delta \leq \tilde{\delta}(\epsilon)$ and all k. Using (3-30) and (3-34), we have for each $0 < \delta \leq \delta_0/2$,

$$\int_0^T \int_\Omega (A-B) dx dt = \int_0^T \int_\Omega \psi_\delta(x) (A-B) dx dt + 0_2(\delta)$$
$$= 0_2(\delta) + \lim_{k \to \infty} \int_0^T \int_\Omega \psi_\delta(x) (A_k - B_k) dx dt$$

$$= 0_{2}(\delta) + \lim_{k \to \infty} \left\{ \int_{0}^{T} \int_{\Omega} (A_{k} - B_{k}) dx dt + \int_{0}^{T} \int_{\Omega} (\psi_{\delta}(x) - 1) A_{k} dx dt - \int_{0}^{T} \int_{\Omega} \psi_{\delta}(x) B_{k} dx dt + \int_{0}^{T} \int_{\Omega} B_{k} dx dt \right\}.$$
(3-36)

In particular, for $0 < \delta \leq \tilde{\delta}(\epsilon)$, we use (3-31), (3-33) and (3-35) to arrive at

$$\int_0^T \int_\Omega (A-B) dx dt \ge -\epsilon + \limsup_{k \to \infty} \sup_{k \to \infty} \int_0^T \int_\Omega (A_k - B_k) dx dt.$$
(3.37)

Since ϵ is arbitrary, it follows that

$$\int_0^T \int_{\Omega} (A - B) dx dt \ge \lim_{k \to \infty} \sup \int_0^T \int_{\Omega} (A_k - B_k) dx dt,$$
(3-38)

which shows (3-26). Combining (3-18), (3-19), (3-20), (3-24), (3-25) and (3-26), we conclude that (3-4) is satisfied. Finally, (3.23) implies

$$(u_k - \phi)^- \to (u - \phi)^- \text{ in } C([0, T]; L^2(\Gamma_2)).$$
 (3-39)

But, each u_k satisfies (2-27) and thus

$$u \ge \phi$$
, almost everywhere in Γ_2 , for each t , (3-40)

from which (3-3) follows. Now the proof of Theorem 3.2 is complete.

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