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# A simple compensation method for the accurate measurement of magnetic losses with a single strip tester

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We present a new method for the accurate characterization of soft magnetic sheets using a permeameter based on the precise compensation of the magnetomotive force (MMF) drop in the flux-closing yoke. It has been developed in order to overcome the systematic uncertainty affecting the value of the magnetic fieldstrength in single sheet testers when obtained, according to the standards, through the measurement of the magnetizing current. This phenomenon is more critical for high permeability materials, because of the reduced MMF drop across the sample. While additional sensors and auxiliary windings have been proposed in the literature, a novel approach is demonstrated here, based on the use of the permeameter upper half yoke as the MMF drop sensor and of an auxiliary winding on the lower half yoke, implementing compensation. This solution, dispensing one from dealing with the usually small signal levels of the conventional MMF drop sensors (e.g. Chattock coils), provides best results with the introduction of wedge-shaped magnetic poles, in order to accurately define the magnetic path length. The method is validated by measurements of power loss, apparent power, and hysteresis cycles on non-oriented and grain-oriented Fe-Si steel sheets, which are compared with local measurements performed on the same samples using H-coil and B-coil across a uniformly magnetized region.

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Index Terms—Magnetic loss, hysteresis cycle, compensated permeameter, apparent power.

### I. INTRODUCTION

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THE PROPER design of electromagnetic devices requires an accurate prediction of 8 electromagnetic devices requires an accurate prediction of 8 iron loss. Indeed, in modern embedded applications such a39 4 hybrid or electric vehicles [1][2], the best compromise must b40 5 found between the machine efficiency [3][4] and torqu41 6 7 density [5]. But any predictive model starts from a prope#2 8 material characterization [6][7]. The standard testing3 technique of soft magnetic steel sheets is based on the use of 4 9 10 the Epstein test frame [8]. It shows good reproducibility, a45 demonstrated by inter-laboratory comparisons of power los46 11 and apparent power measurements [9]. However, this method47 12 13 besides requiring tedious preparation of the samples, i48 14 inevitably affected by appreciable systematic deviations (up t49 15 about 10 % at high inductions) from the true values of th50 16 measured quantities, as obtained, for example, by accurat51 17 measurements using H-coils [9][10][11]. But the H-coil 2 method requires the integration of low-level signals and is3 18 19 hardly acceptable in the industrial practice. Increasing interest4 is therefore attached, at present time, to the Single Sheet<sup>5</sup> 20 Testing (SST) method, applied according to the pertaining IE $e^{6}$ 21 Standard [12][13], because of the convenient use of wide<sup>7</sup> 22 lamination samples. The STT method does not require stress  $\frac{58}{2}$ 23 relief of samples upon cutting, can be directly applied to the 24 domain-refined high-grade grain-oriented materials, and 25 shows good reproducibility of measurements [9]  $[14]_{rot}^{p_1}$ 26 Consequently, there is demand by industry for including  $SS_{3}^{\mu 2}$ 27 reference values in the material specification standards. Wit $\breve{k}_4$ 28 the SST arrangement, where the sheet sample is inserted in  $\frac{1}{85}$ 29 double-C laminated yoke, the magnetic fieldstrength is6 30 31 calculated using the measured magnetizing current and 87 32 defined magnetic path length ( $l_{\rm m} = 0.45$  m) is adopted. A maine 33 problem here is that the magnetomotive force (MMF) drop ing 34 the flux-closing yokes may not be negligible with respect to

the one across the sample, especially with high-permeability materials. This can lead to overestimated magnetic field values. At the same time, the yoke itself can provide a certain contribution to the measured loss, depending on its manufacture and the possible existence of interlaminar eddy currents at the pole faces [15].

One way to overcome this difficulty is, as discussed in [16], one of compensating the drop of the MMF in the yokes of the permeameter by an auxiliary magnetizing winding. This is driven by a Chattock coil, placed over the measuring sample region of length  $l_m$  and a very high-gain amplifier, implementing a feedback control on the auxiliary current, such as to compensate the MMF drop outside the length  $l_m$  [16]. Like all the *H*-coil measurements, this method has a weak point in the necessity of handling small signal levels.

In this paper, we consider a single strip double-C yoke permeameter, applied to annealed Epstein samples, where a new compensation method, simpler and more effective than previous literature solutions [16], is implemented. Epstein strip samples are used in the present experiments for practical convenience, but the method could be easily adapted to standard SST permeameters. The idea is one of using the upper half of the yoke as a zero MMF indicator and the lower half for accommodating the compensation circuit. It is a simple measuring arrangement, where the signal to be controlled is relatively large and much easier to handle than the weak and noise-prone signal generated by a Chattock coil. As a further advantage, there is no geometrical discontinuity in the magnetizing winding, as required instead, with ensuing inhomogeneity of the applied field, by the insertion of the Chattock coil [16]. The method is validated by comparison with accurate measurements performed upon a relatively restricted median region of the strip sample by the H-coil method. Non-oriented (NO) and high-permeability grainoriented (GO) Fe-Si samples have been tested, with very good 103

70 agreement between the results obtained with the magnetizing 0 71 current and those obtained by localized H and B coils. 101

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72 II. THE COMPENSATED PERMEAMETER

In this section we describe the geometry of the system and
 the circuit for the compensation of the MMF drop.

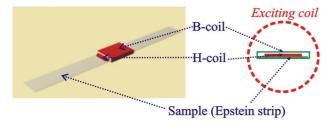
## 75 A. System geometry and magnetic path length

The sample is a conventional annealed Epstein strip. The induction derivative is measured by means of a 500-turn  $\frac{100}{100}$ mm long pickup coil placed at the center of the strip. Inside mm long pickup coil placed at the center of the strip. Inside the *B*-coil, a many-turn calibrated flat *H*-coil (1 mm thick turn-area 2.25  $\cdot 10^{-2}$  m<sup>2</sup>) of same length is placed upon the sample surface (see

Fig. 1). The *H*-coil provides, after integration of the measured voltage, the tangential field upon the measuring area. Since in the measuring region the applied field and the induction are verified to be highly uniform, tangential field and internal field are bound to coincide. The measured local to be applied for the measured local set the induction of the measurement of the measurement

87 loss is thus identified with the true loss figure of the material,

88 that is, the reference quantity for the results obtained with the 89 compensated permeameter.



92 Fig. 1 - Sample (Epstein strip), with enwrapping *B*-coil and tangential *H*- coil

93 (3D view and front view).

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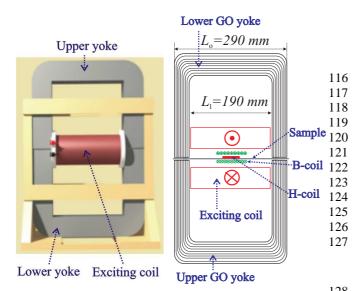


Fig. 2 - a) 3D view of the permeameter. b) 2D view of magnetizer, exciting circuit, and measuring coils.

The developed permeameter, shown in Fig. 2, consists of 30
double-C laminated yoke of 50 mm × 50 mm cross-section 31
area, made of 0.30 mm thick high-permeability GO strips. 32
uniformly wound magnetizing solenoid covers the distance 33
= 190 mm between the pole faces of the yoke. It is endow 34
with series connected additional narrow windings at its end 35
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by which maximum uniformity of the applied field across the whole distance  $L_i$  is obtained ([11], p. 109).

According to Ampère's law, the magnetomotive force NI generated by a current I flowing into the N-turn solenoid is related to the magnetic field H along a closed path L by the equation

$$N \cdot I = \oint_{I} H \cdot \mathrm{d}l. \tag{1}$$

If a method is found by which the drop of the MMF in the flux-closing yoke is made either negligible or fully compensated, (1) simply becomes  $N \cdot I = H \cdot L_s$ , where  $L_s$  is the length of the mean magnetic flux path in the sample.  $L_s$  has value intermediate between  $L_i$ , the distance between the pole faces, and  $L_o$ , the width of the permeameter (Fig. 2b), as the obvious result of flux channeling from the sample into the yokes (see Fig. 3a). It is actually not well defined and it depends on the magnetization level in the sheet sample. It will appreciably change on approaching the material saturation.

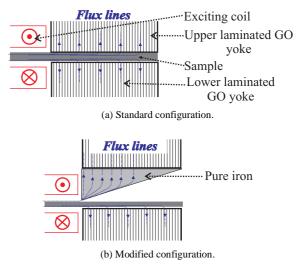


Fig. 3 - Defining the magnetic flux path length in the sample.

We wish to impose a flux path such that the length  $L_s$  coincides with the distance between the pole faces  $L_i$ . To this end, the following modification of the upper yoke is proposed. Two identical wedge-shaped pure iron poles are placed beneath the two limbs of the upper yoke, as shown in Fig. 3b and Fig. 5, with the contact lines between the sheet surface and the poles placed exactly at the distance  $L_i$ . As discussed in the following Section, this is expected to force, under the action of a feedback system employing auxiliary windings on the yokes, the flux path length in the sample to be equal to  $L_i$ . Under these conditions, the magnetic field H in the sample will be obtained according to the relation

$$N \cdot I = H \cdot L_i \,. \tag{2}$$

independent of the way the flux lines enter the lower yoke.

## B. The principle of MMF compensation.

The idea here developed follows to some extent from the method using a Chattock coil sensor to cancel the MMF on a known length, as described in [16]. *Here we take the yoke itself as a zero MMF indicator*. A few-turn secondary coil wound around the upper yoke provides the derivative of the magnetic flux  $\varphi_1$  flowing in it. According to the simplified reluctance description of the system shown in Fig. 4, where

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137 the reluctance of the sample is  $R_{\rm S}$ , the reluctance of the upped 4 and lower yokes is  $R_{\rm Y}$ , and that of the wedge-shaped pole 185 138  $R_{\rm P}$ , the MMF drop E pertaining to the magnetic circuit outside 139 the sample is proportional to the flux  $\varphi_1$ . Since there is  $\eta_{0,7}$ 140 MMF source in the upper yoke, E will be reduced to zero  $b\chi_8$ 141 canceling the flux  $\varphi_1$ , that is, bringing to zero the voltage 142 correspondingly induced in the secondary coil. This can  $b_{70}$ 143 accomplished by adding a compensation winding on the low  $\frac{1}{11}$ 144 yoke and controlling it in such a way that the corresponding  $\frac{1}{2}$ 145 generated MMF  $N_c I_c$  leads to the condition E = 0. The control  $\frac{1}{73}$ 146 147 loop is schematically shown in Fig. 5. An analog control card (PID controller) keeps the voltage  $v_1 \propto d\varphi_1/dt$  equal to zero,  $b\sqrt[3]{4}$ 148 149 properly supplying, via a high-gain linear amplifier, the5 150 compensation winding on the lower yoke. Consequently, 4176 the flux  $\varphi_{\rm S}$  in the sample is made to entirely flow in the lower 151 152 yoke.

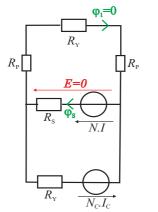


Fig. 4. Reluctance network of the compensated permeameter.

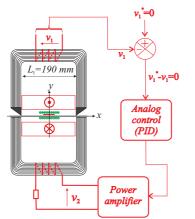


Fig. 5. Compensated permeameter and control circuit.

184 185 153 This compensation method is simple and sensitive, sincel 86 does not require any specific MMF sensor and a high signal 87 154 to-noise ratio is ensured by the high permeability of the  $GQ_8$ 155 laminations employed in the yokes. The voltage  $v_1$  is, forg 156 example, always much larger than the one achievable  $by_{190}$ 157 Chattock coil, and its control around the zero value 181 158 correspondingly easier and more precise. To be remarked that  $\frac{1}{192}$ 159 imperfect contact between the wedge shaped poles and the 160sample, which can be lumped in the pole reluctance  $R_{\rm P}$ , has  $r_{\rm P}$ 161 little detrimental effect on the permeameter performance, 194 162 163 because the flux in the upper yoke is made to vanish. 196

#### **III. EXPERIMENTAL RESULTS**

The novel compensated permeameter has been tested on nonoriented Fe-(3 wt%)Si sheets (thickness 0.35 mm) and high-permeability (HiB) grain-oriented sheets (thickness 0.28 mm). Energy loss, apparent power, and hysteresis loops were measured under sinusoidal induction waveform at f = 100 Hz and peak polarization values  $0.2 \text{ T} \le J_p \le 1.5 \text{ T}$ . The value of the magnetic field *H* was obtained both by measuring the magnetizing current and by integrating the signal induced in the *H*-coil placed at the center of the Epstein strip.

#### A. Non-oriented Fe-Si sheets

Fig. 6 shows the measured energy loss vs. peak polarization obtained through the *H*-coil and current methods.

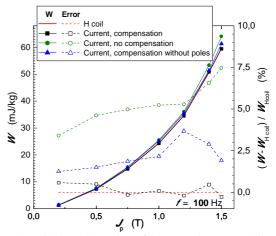


Fig. 6. NO Fe-Si sheets. Comparison of the energy loss measured by the *H*-coil method and the current measurement method. The permeameter can be compensated as described above, non-compensated (the compensation circuit is not switched on), or compensated without the iron poles. The dashed lines provide the deviations of the loss figures measured with the current method with respect to the *H*-coil method.

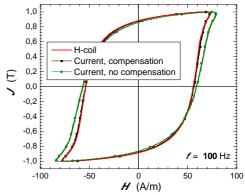


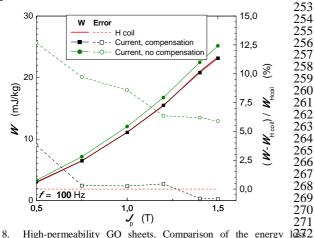
Fig. 7. NO Fe-Si sheets.  $J_p = 1$  T. Comparison of the hysteresis loops measured by the *H*-coil method and by the current measurement method, with and without compensation.

It appears that the results by the *H*-coil method and the current method with compensation and wedge-shaped iron poles on the upper yoke show remarkable agreement. On the other hand, because of the additional loss contribution by the yokes, overestimated figures are obtained by use of the uncompensated permeameter, the higher  $J_p$  the higher the loss deviation (up to about 7 %). If the compensation procedure is applied to the standard permeameter configuration without wedge-shaped poles, the loss value is still overestimated (from 2% to 4%), because the magnetic path cannot be fully

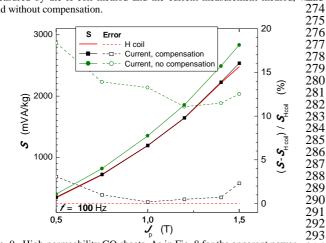
198constrained to the sample length  $L_i$  and is slightly longer.231199The good agreement between the *H*-coil and the compensated2200current methods is confirmed by the corresponding33201measured hysteresis loops shown in Fig. 7.234

## 202 B. HiB grain-oriented sheets 235

Measurements of high permeability GO materials are 203 demanding, because the MMF drop in the yokes may be  $f_{447}$ 204from negligible with respect to the one in the sample. At  $th_{38}$ 205 same time, the localized measurements using the *H*-coil  $\tilde{ars_{99}}$ 206 difficult, especially at low frequencies, because the signal  $c \tilde{a} \tilde{\mu}_0$ 207 be very small. This adds to the interest for the here proposed  $\frac{1}{2}$ 208 solution, which includes also the measurement of the apparent 2209 power. Fig. 8 compares the energy loss values obtained with  $\frac{1}{43}$ 210 the *H*-coil method and the magnetizing current method, with  $\frac{1}{44}$ 211 and without compensation, in the GO sheet. It is apparent that  $\frac{1}{45}$ 212 the compensated permeameter and the *H*-coil measurement  $\frac{1}{246}$ 213 provide close results. Under uncompensated conditions,  $one_{247}$ 214 215 finds instead that the current method overestimates the magnetic loss by a substantial extent, especially at  $\log_{10}$ 216inductions (about 10% for  $J_p = 0.5$  T, 6% for  $J_p = 1.5$  T). This 249 217 is expected, because the loss in the yoke depends only on  $\frac{4^{+7}}{250}$ 218 not on the type of material under test, and its ratio to  $tb\tilde{s}_1$ 219 220 sample loss increases with better materials 252



 $\overline{222}$  Fig. 8. High-permeability GO sheets. Comparison of the energy  $12\overline{332}$ measured by the *H*-coil method and the current measurement method, where 323and without compensation. 274



225 226 Fig. 9. High-permeability GO sheets. As in Fig. 8 for the apparent power. 293 294

Fig. 9 shows the behavior of the measured apparent pow295 Again, we find good agreement between the results provided7 by the two methods: *H*-coil and compensated permeameter.208 is noted that part of the small discrepancies occurring between 300

the *H*-coil and compensated current measurements could be attributed to the difficulties intrinsic to the *H*-coil method (small signal, integration problems...). This further stresses the merits of the here proposed approach.

#### IV. CONCLUSIONS

A permeameter has been developed, which applies a simple and effective magnetomotive force compensation method for the accurate characterization of soft magnetic steel sheets. It does not require specific sensors, except wedge shaped pole faces for the precise definition of the magnetic path length. It works on the principle of using the flux-closing yoke itself for both sensing and compensation.

The performances of this permeameter have been validated by measurements on non-oriented and grain-oriented Fe-Si sheets, whose results excellently compare with the results provided by measurements performed using the tangential *H*coil method.

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