

The trunkenness of a volume-preserving vector field

Ana Rechtman, Pierre Dehornoy

▶ To cite this version:

An Rechtman, Pierre Dehornoy. The trunkenness of a volume-preserving vector field. 2016. hal-01326696v1

HAL Id: hal-01326696 https://hal.science/hal-01326696v1

Preprint submitted on 5 Jun 2016 (v1), last revised 21 Sep 2017 (v2)

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

THE TRUNKENNESS OF A VOLUME-PRESERVING VECTOR FIELD

ANA RECHTMAN, PIERRE DEHORNOY

ABSTRACT. We construct a new invariant—the trunkenness—for volume-perserving vector fields on \mathbb{S}^3 up to volume-preserving diffeomorphism. We prove that the trunkenness is independent from the helicity and that it is the limit of a knot invariant (called the trunk) computed on long pieces of orbits.

Contents

1.	Trunk of knots	3
2.	Trunkenness of measure preserving vector fields	4
3.	Independence of helicity	8
4.	Trunkenness and the existence of periodic orbits	9
5.	Trunkenness of knotted tubes	14
References		16

The problem we address here is the construction of new invariants of volume-preserving vector fields on \mathbb{S}^3 or on compact domains of \mathbb{R}^3 up to volume-preserving diffeomorphisms. This problem is motivated by at least two physical situations. First if v is the velocity field of a time-dependant ideal fluid satisfying the Euler equations (ideal hydrodynamics) then its vorticity field curl v is transported by the flow of v [Hel1858]. Second if B is the magnetic field in an incompressible plasma (ideal magnetodynamics), then B turns out to be transported by the velocity field as long as the latter does not develop singularities [Wol58]. In these contexts, invariants of curl v or B up to volume-preserving diffeomorphisms yield time-independent invariants of the system.

Not so many such invariants exist. The first one was discovered by William Thomson [Tho1867]: if the considered field has a periodic orbit or a periodic tube, then its knot type is an invariant (this remark led to the development of knot theory by Peter G. Tait [Tai1877]). However it may not be easy to find periodic orbits, and even then such an invariant only takes a small part of the field into account.

The main known invariant is called *helicity*. It is defined by the formula $\operatorname{Hel}(v) = \int v \cdot u$, where $u = \operatorname{curl}^{-1}(v)$ is an arbitrary vector-potential of v. It was discovered by Woltjer, Moreau, and Moffatt [Wol58, Mor61, Mof69]. Helicity is easy to compute or to approximate since it is enough to exhibit a vector-potential of the considered vector field, to take the scalar product and to integrate. The connection with knot theory was sketched by Moffatt [Mof69] and deepened by Arnold [Arn73] as follows. Denote by $k_X(p,t)$ a loop starting at the point p, tangent to the vector field X for a time t and closed by an arbitrary segment of bounded length. Denote by Lk the linking number of loops. Arnold showed that for almost every p_1, p_2 , the limit $\lim_{t_1, t_2 \to \infty} \frac{1}{t_1 t_2} \operatorname{Lk}(k_X(p_1, t_1), k_X(p_2, t_2))$ exists (see also [Vog02] for a corrected statement). Moreover if X is ergodic the limit coincides almost everywhere with $\operatorname{Hel}(X)$.

The idea of considering knot invariants of long pieces of orbits of the vector field was pursued by Gambaudo and Ghys [GG01] who considered ω -signatures of knots, Baader [Baa11] who considered linear saddle invariants, and Baader and Marché [BM12] who considered Vassiliev's finite type

Date: June 5, 2016.

invariants. In every case, it is shown that $\lim_{t\to\infty}\frac{1}{t^n}V(k_X(p,t))$ exists, where V is the considered invariant and n a suitable exponent called the *order* of the asymptotic invariant. However all these constructions have the drawback that they do not yield any new invariant for ergodic vector fields, as in this case the obtained limits are all functions of the helicity.

Recently, it was proved by Kudryavtseva [Kud15] for vector fields obtained by suspension of a diffeomorphism of the 2-dimensional disc and then by Enciso, Peralta-Salas and Torres de Lizaur [EPT16] for vector fields on manifolds without boundary, that the helicity is the only regular integral invariant (in the sense that its Fréchet derivative is the integral of a continuous kernel). These results partly explain why most constructions yield invariants that are functions of helicity for ergodic vector fields. However they do not rule out the existence of other invariants, but imply that such invariants cannot be too regular.

An example of such another invariant is the asymptotic crossing number considered by Freedman and He [FH91]. The advantage is that it is not proportional to helicity, but the disadvantage is that it is hard to compute, even on simple examples.

In this article we consider a less known knot invariant called the trunk (see Definition 1.1 below). It was defined by Ozawa [Ozw10], building on the concept of thin position that was introduced by Gabai [Gab87] for solving the R-conjecture. Less famous that the invariants previously studied in the context of vector fields, the trunk has the advantage that its definition relies on surfaces transverse to the considered knot, so that it is easy to transcript to the context of vector fields. The invariant depends on an invariant measure μ for the flow of the vector field that may or may not be a volume, and is invariant under diffeomorphisms that preserve the measure μ . Given a μ -preserving vector field X and a surface S, the geometric flux through S is the infinitesimal measure that crosses S in both directions (see Definition 2.1), it is denoted by Flux(X, μ, S). Our invariant is a minimax of the geometric flux, where one minimizes over all height functions and maximizes over the levels of the considered height function.

Definition 1. Assume that X is a vector field on \mathbb{S}^3 or on a compact domain of \mathbb{R}^3 that preserves a probability measure μ . Denote by ϕ_X the flow of X. The trunkenness of X with respect to μ is

$$\operatorname{Tks}(X,\mu) := \inf_{\substack{h \text{ height} \\ \text{function}}} \max_{t \in [0,1]} \operatorname{Flux}(X,\mu,h^{-1}(t)) \ = \inf_{\substack{h \text{ height} \\ \text{function}}} \max_{t \in [0,1]} \lim_{\epsilon \to 0} \frac{1}{\epsilon} \mu(\phi_X^{[0,\epsilon]}(h^{-1}(t))).$$

By a height function on \mathbb{S}^3 we refer to a function with only two singular points and whose level sets are 2-dimensional spheres. On \mathbb{R}^3 , the level sets of a height function all are planes.

From the definition, it is straightforward that the trunkenness of a vector field is invariant under diffeomorphisms that preserve the measure μ . More is true, the trunkenness is invariant under homeomorphisms that preserve μ .

Theorem A. Assume that X_1 and X_2 are vector fields on \mathbb{S}^3 or on a compact domain of \mathbb{R}^3 that preserve a probability measure μ and that there is a μ -preserving homeomorphism f that conjugates the flows of X_1 and X_2 . Then we have

$$\operatorname{Tks}(X_1, \mu) = \operatorname{Tks}(X_2, \mu).$$

What we do in this paper is to prove several properties of this new invariant. The first one is a continuity result that implies that the trunkenness of a vector field is an asymptotic invariant of order 1. For K a knot, we denote by Tk(K) its trunk (see Definition 1.1 below).

Theorem B. Suppose that $(X_n, \mu_n)_{n \in \mathbb{N}}$ is a sequence of measure-preserving vector fields such that $(X_n)_{n \in \mathbb{N}}$ converges to X and $(\mu_n)_{n \in \mathbb{N}}$ converges to μ in the weak-* sense. Then we have

$$\lim_{n\to\infty} \mathrm{Tks}(X_n,\mu_n) = \mathrm{Tks}(X,\mu).$$

In particular if X is ergodic with respect to μ then, for μ -almost every p, the limit

$$\lim_{t\to\infty}\frac{1}{t}\mathrm{Tk}(k_X(p,t))$$

exists and is equal to $Tks(X, \mu)$.

This continuity result then allows us to compute the trunkenness of some explicit vector fields on \mathbb{S}^3 called Seifert flows. These computations in turn show that the trunkenness is not dictated by helicity, thus contrasting with most previously known knot-theoretical constructions.

Theorem C. There is no function f such that for every volume-preserving vector field X on \mathbb{S}^3 one has $\operatorname{Tks}(X,\mu) = f(\operatorname{Hel}(X,\mu))$.

Finally we address the question of what happens if for a non-singular vector field on \mathbb{S}^3 there is a function that achieves the trunkenness, or in other words if the infimum in Definition 1 is a minimum.

Theorem D. Let X be a non-singular vector field on \mathbb{S}^3 preserving the measure μ and h a height function such that

$$\operatorname{Tks}(X,\mu) = \max_{t \in [0,1]} \operatorname{Flux}(X,\mu,h^{-1}(t))$$

Then X has an unknotted periodic orbit.

One of the main motivations for constructing topological invariants of a vector field X is to find lower bound on the energies $E_p(X) := \int |X|^p d\mu$. Indeed since a topological invariant yields a time-independent invariant of the physical system, an energy bound in term of a topological invariant will also be time-independent, although the energy may vary when the vector field is transported under (volume-preserving) diffeomorphisms. Such energy bounds exist for the helicity and for the asymptotic crossing number. We do not know whether the trunkenness bounds the energy.

The plan of the article is as follows. First we recall in Section 1 the definition of the trunk of a knot in order to make the definition for vector fields natural. Then we define the trunkenness of a measure-preserving vector field and prove Theorems A and B in Section 2. Using this result we compute of trunkenness of Seifert vector fields and prove Theorem C in Section 3. We prove Theorem D in Section 4. Finally in Section 5 we compute the trunkenness of some vector fields supported in the tubular neighborhood of a link.

Acknowledgements. We thank Michel Boileau who suggested to study the trunk instead of the genus for vector fields during a visit of P.D. to Toulouse in 2013.

1. Trunk of knots

For K a knot, we denote by K the set of all embeddings of K into \mathbb{R}^3 . The standard height function on \mathbb{R}^3 is the function $h_z: \mathbb{R}^3 \to \mathbb{R}, (x,y,z) \mapsto z$. Every level $h_z^{-1}(t)$ is a 2-dimensional plane. An embedding $k \in K$ is said to be in Morse position with respect to h_z if the restriction of h_z to k is a Morse function. In this case there are only finitely many points at which k is tangent to a level of h_z .

Definition 1.1. Assume that k is an embedded knot in \mathbb{R}^3 that is in Morse position with respect to h_z . The trunk of the curve k relatively to h_z is

$$\operatorname{tk}_{h_z}(k) := \max_{t \in \mathbb{R}} \sharp \{k \cap h_z^{-1}(t)\}.$$

The $trunk \ of \ a \ knot \ K$ is then defined by

$$\operatorname{Tk}(K) := \min_{k \in \mathcal{K}} \operatorname{tk}_{h_z}(k) = \min_{k \in \mathcal{K}} \max_{t \in \mathbb{R}} \sharp \{k \cap h_z^{-1}(t)\}.$$

The trunk of a knot was defined by M. Ozawa [Ozw10] and motivated by D. Gabai's definition of the waist of a knot [Gab87].

Example 1.2. A knot is trivial if and only if its trunk equals 2. Indeed the embedding as the boundary of a vertical disc shows that the trunk is less than or equal to 2, and every embedding in Morse position of the trivial knot has to intersect some horizontal plane in at least two points. Conversely, if the trunk of a knot is equal to 2, then it admits an embedding that intersects every horizontal plane in at most two points. The union of the segments that connect these pairs of points is a disc bounded by the knot, implying that the knot is trivial.

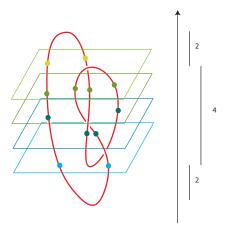


FIGURE 1. The trunk of the trefoil knot: the maximal number of intersection points between a horizontal level and the proposed embedding is 4. This number cannot be reduced under isotopy, hence the trunk of the trefoil is 4.

Example 1.3. For p, q in \mathbb{N} , the torus knot T(p, q) can be realized as the closure of a braid with q strands, yielding $\operatorname{Tk}(T(p,q)) \leq 2q$. By symmetry one also gets $\operatorname{Tk}(T(p,q)) \leq 2p$. Actually, one can prove $\operatorname{Tk}(T(p,q)) = 2\min(p,q)$, see Remark 1.2 in $[\operatorname{Ozw} 10]$.

Instead of fixing the function and changing the knot up to isotopy, one can fix the knot and change the function up to orientation preserving diffeomorphism. With this in mind, one defines a height function on \mathbb{R}^3 as a function obtained by precomposing h_z by a diffeomorphism, that is, a function of the form

$$h: \mathbb{R}^3 \to \mathbb{R}$$

 $(x, y, z) \mapsto h_z(\phi(x, y, z))$

for ϕ an orientation preserving diffeomorphism of \mathbb{R}^3 . In particular, a height function is a function whose levels are smooth planes. For K a knot and k a fixed embedding of K in \mathbb{R}^3 , one can then define

$$\operatorname{tk}_h(k) := \max_{t \in \mathbb{R}} \sharp \{k \cap h^{-1}(t)\},\,$$

so that we have the alternative definition

(1)
$$\operatorname{Tk}(K) = \min_{\substack{h \text{ height} \\ \text{function}}} \operatorname{tk}_h(k) = \min_{\substack{h \text{ height} \\ \text{function}}} \max_{t \in \mathbb{R}} \sharp \{k \cap h^{-1}(t)\}.$$

2. Trunkenness of measure preserving vector fields

We use the definition of Equation (1) to define the trunkenness of a vector field with respect to an invariant measure. The main question then concerns the analog of the number of intersection points of a surface with a curve when the curve is replaced by a vector field. A natural answer is the geometric flux. If X is a vector field that preserves a measure μ given by a volume element Ω , one can then consider the 2-form $\iota_X\Omega$. For S a piece of oriented surface that is positively transverse to X, the integral $\int_S \iota_X\Omega$ computes the instantaneous volume that crosses S. In other words, by Fubini Theorem we have $\mu(\phi^{[0,t]}(S)) = (\int_S \iota_X\Omega) \cdot t$. On the other hand if S is negatively transverse to X we have $\mu(\phi^{[0,t]}(S)) = -(\int_S \iota_X\Omega) \cdot t$. Therefore in this case, for any surface S, the instantaneous volume crossing S is given by $\int_S |\iota_X\Omega|$. Now if the measure μ is not given by integrating a volume form one cannot consider the above integral, but the quantity $\mu(\phi^{[0,t]}(S))$ still makes sense for any piece of surface S (see Figure 2).

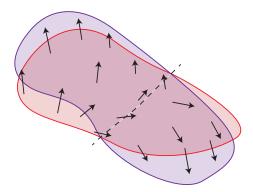


FIGURE 2. For S a piece of a surface (in red), the domain $\phi^{[0,\epsilon]}(S)$ is roughly the part of the space located between S and $\phi^{\epsilon}(S)$. It is pinched around those points where the vector field is tangent to S. Its volume is $\epsilon \cdot \operatorname{Flux}(X, \mu, S)$ at the first order.

Definition 2.1. For X a vector field that preserves a measure μ and for S a surface, the *geometric flux* of (X, μ) through S is

$$\operatorname{Flux}(X,\mu,S) := \lim_{\epsilon \to 0} \frac{1}{\epsilon} \, \mu(\phi^{[0,\epsilon]}(S)).$$

This definition generalizes the number of intersection points of a knot with a surface. Indeed one can see an embedding k of a knot K as a vector field with a particular invariant measure in the following way: consider a non-singular vector field X_k that is tangent to k at every point and denote by ϕ_k^t the induced flow. Since k is closed, ϕ_k is T_k -periodic for some $T_k > 0$. The Dirac linear measure associated to X_k is defined by

$$\mu_k(A) = \text{Leb}(\{t \in [0, T_k], \phi_k^t(x) \in A\})$$

where A is a measurable set and x an arbitrary point on k. The measure μ_k is X_k -invariant and has total mass T_k . In this setting, for S a surface that intersects k in finitely many points, a point p in the set $k \cap S$ has μ_k -measure zero and thus cannot be detected by the measure. But by definition of μ_k the set $\phi_k^{[0,\epsilon]}(p)$ is an arc of k of μ_k -measure ϵ and since $k \cap S$ is made of finitely many points, for ϵ small enough, the set $\phi_k^{[0,\epsilon]}(k \cap S)$ has μ_k -measure exactly $\epsilon \cdot \sharp \{k \cap S\}$. In other words, one has

$$\sharp\{k\cap S\}=\lim_{\epsilon\to 0}\frac{1}{\epsilon}\,\mu_k(\phi_k^{[0,\epsilon]}(k\cap S)).$$

As μ_k is concentrated on k, we thus have

(2)
$$\sharp\{k \cap S\} = \lim_{\epsilon \to 0} \frac{1}{\epsilon} \mu_k(\phi_k^{[0,\epsilon]}(S)) = \operatorname{Flux}(X_k, \mu_k, S),$$

so the geometric flux indeed generalizes the number of intersection points.

We now mimic for vector fields the definition of the trunk of a knot. In order to have a well-defined maximum, in what follows we assume the vector fields are on a compact domain $D^3 \subset \mathbb{R}^3$ or on the 3-sphere $\mathbb{S}^3 = \mathbb{R}^3 \cup \infty$. In the later case, we define the *standard height function*

$$h_0: \mathbb{S}^3 \to [0,1]$$

 $(x,y,z) \mapsto 1 - \frac{1}{1+x^2+y^2+z^2}.$

The levels $h_0^{-1}(0)$ and $h_0^{-1}(1)$ consist of the points (0,0,0) and ∞ respectively and every other level $h_0^{-1}(t)$ is a 2-dimensional sphere of radius $\sqrt{t/(1-t)}$ centered at the origin. A *height function* on \mathbb{S}^3 is then a function obtained by precomposing by an orientation preserving diffeomorphism ϕ of \mathbb{S}^3 , that is, a function of the form $h: \mathbb{S}^3 \to [0,1], (x,y,z) \mapsto h_0(\phi(x,y,z))$.

Definition 2.2. Let X be a vector field whose flow preserves a measure μ on a compact domain of \mathbb{R}^3 or on \mathbb{S}^3 , and h a height function. We set

$$\operatorname{tks}_h(X,\mu) := \max_{t \in [0,1]} \operatorname{Flux}(X,\mu,h^{-1}(t)) = \max_{t \in [0,1]} \lim_{\epsilon \to 0} \frac{1}{\epsilon} \mu(\phi^{[0,\epsilon]}(h^{-1}(t))).$$

The trunkenness of (X, μ) is defined as

$$\operatorname{Tks}(X,\mu) := \inf_{\substack{h \text{ height} \\ \text{function}}} \operatorname{tks}_h(X,\mu) = \inf_{\substack{h \text{ height} \\ \text{function}}} \max_{t \in [0,1]} \operatorname{Flux}(X,\mu,h^{-1}(t))$$
$$= \inf_{\substack{h \text{ height} \\ \text{function}}} \max_{t \in [0,1]} \lim_{\epsilon \to 0} \frac{1}{\epsilon} \mu(\phi^{[0,\epsilon]}(h^{-1}(t))).$$

Note that we can only consider an infimum instead of a minimum as in the case of knots. In Section 4 we prove that for non-singular vector fields, if the trunkenness is a minimum then the vector field possesses an unknotted periodic orbit.

If the invariant measure μ is given by the integration of a volume form Ω , we get the alternative definitions

$$\operatorname{tks}_h(X,\Omega) = \max_{t \in [0,1]} \int_{h^{-1}(t)} |\iota_X \Omega|, \quad \text{and} \quad \operatorname{Tks}(X,\Omega) = \inf_{\substack{h \text{ height} \\ \text{function}}} \max_{t \in [0,1]} \int_{h^{-1}(t)} |\iota_X \Omega|.$$

We now prove the invariance under homeomorphism.

Proof of Theorem A. Assume without loss of generality that there is a homeomorphism f such that $f \cdot \phi_{X_1}^t = \phi_{X_2}^t \cdot f$, with $\phi_{X_i}^t$ the flow of X_i , and assume by contradiction that $\mathrm{Tks}(X_1, \mu) < \mathrm{Tks}(X_2, \mu)$. Let $0 < \delta = \mathrm{Tks}(X_2, \mu) - \mathrm{Tks}(X_1, \mu)$. Take h_n to be a sequence of (differentiable) height functions such that

$$\lim_{n\to\infty} \operatorname{tks}_{h_n}(X_1,\mu) = \operatorname{Tks}(X_1,\mu).$$

Consider the continuous functions $h_n = h_n \cdot f^{-1}$ and for each n consider a differentiable function H_n that is arbitrarily C^0 -close to h_n . Hence, for each n, the level sets of H_n and h_n are arbitrarily close. Observe that we can assume that H_n is a height function (by taking H_n such that its level sets are all spheres or all planes, depending on the domain of the vector field). Thus we can choose H_n so that, for every $t \in [0,1]$ and for ϵ small enough, we have

$$|\mu(\phi_{X_2}^{[0,\epsilon]}(H_n^{-1}(t))) - \mu(\phi_{X_2}^{[0,\epsilon]}(\widetilde{h}_n^{-1}(t)))| < \frac{\delta}{4}.$$

Fix n large enough so that $\operatorname{tks}_{h_n}(X_1,\mu) - \operatorname{Tks}(X_1,\mu) < \frac{\delta}{4}$, and for such n set T in [0,1] so that

$$\operatorname{tks}_{H_n}(X_2, \mu) = \operatorname{Flux}(X_2, \mu, H_n^{-1}(T)).$$

We have two possibilities:

(1) If $\operatorname{Flux}(X_2, \mu, H_n^{-1}(T)) \geq \operatorname{Flux}(X_2, \mu, \widetilde{h}_n^{-1}(T))$, then for ϵ small enough

$$\mu(\phi_{X_2}^{[0,\epsilon]}(H_n^{-1}(T))) \le \frac{\delta}{4} + \mu(\phi_{X_2}^{[0,\epsilon]}(\widetilde{h}_n^{-1}(T))).$$

Hence

$$\begin{aligned} \operatorname{Tks}(X_{2}, \mu) & \leq & \operatorname{tks}_{H_{n}}(X_{2}, \mu) = \lim_{\epsilon \to 0} \frac{1}{\epsilon} \mu(\phi_{X_{2}}^{[0, \epsilon]}(H_{n}^{-1}(T))) \\ & \leq & \frac{\delta}{4} + \lim_{\epsilon \to 0} \frac{1}{\epsilon} \mu(\phi_{X_{2}}^{[0, \epsilon]}(\widetilde{h}_{n}^{-1}(T))) \\ & = & \frac{\delta}{4} + \lim_{\epsilon \to 0} \frac{1}{\epsilon} \mu(\phi_{X_{1}}^{[0, \epsilon]}(h_{n}^{-1}(T))) \\ & \leq & \frac{\delta}{4} + \operatorname{tks}_{h_{n}}(X_{1}, \mu) < \frac{\delta}{2} + \operatorname{Tks}(X_{1}, \mu). \end{aligned}$$

Thus, $\delta = \text{Tks}(X_2, \mu) - \text{Tks}(X_1, \mu) < \frac{\delta}{2}$, a contradiction proving Theorem A in this case.

(2) If
$$\operatorname{Flux}(X_2, \mu, H_n^{-1}(T)) < \operatorname{Flux}(X_2, \mu, \widetilde{h}_n^{-1}(T))$$
, then

$$\operatorname{tks}_{H_n}(X_2, \mu) = \operatorname{Flux}(X_2, \mu, H_n^{-1}(T))$$

$$< \operatorname{Flux}(X_2, \mu, \widetilde{h}_n^{-1}(T)) = \operatorname{Flux}(X_1, \mu, h_n^{-1}(T))$$

$$\leq \operatorname{tks}_{h_n}(X_1, \mu) < \operatorname{Tks}(X_1, \mu) + \frac{\delta}{4} < \operatorname{Tks}(X_2, \mu).$$

Thus, $\operatorname{tks}_{H_n}(X_2,\mu) < \operatorname{Tks}(X_2,\mu)$, a contradiction to the definition of the trunkenness. This finishes the proof of Theorem A.

Computing the trunkenness of a vector field is not easy in general. Considering a given height function gives an upper bound on the trunkenness, but lower bounds are harder to find. Theorem B is a continuity result that provides a useful tool.

Proof of Theorem B. We begin with the first part of the theorem, namely we prove that if $(X_n, \mu_n)_{n \in \mathbb{N}}$ is a sequence of measure-preserving vector fields such that $(X_n)_{n \in \mathbb{N}}$ converges to X and $(\mu_n)_{n \in \mathbb{N}}$ converges to μ in the weak-* sense, then we have $\lim_{n \to \infty} \operatorname{Tks}(X_n, \mu_n) = \operatorname{Tks}(X, \mu)$.

Fix $\epsilon > 0$. The convergence hypothesis implies that for every surface S, for every $\delta > 0$ small and for n big enough, we have

(3)
$$|\mu_n(\phi_{X_n}^{[0,\delta]}(S)) - \mu(\phi_X^{[0,\delta]}(S))| \le \epsilon,$$

where ϕ_X^t and $\phi_{X_n}^t$ denote the flows of the vector fields X and X_n , respectively.

Assume that $Tks(X_n, \mu_n)$ does not converge to $Tks(X, \mu)$, then for any $N \in \mathbb{N}$ there exists $n \geq N$ such that

$$|\operatorname{Tks}(X_n, \mu_n) - \operatorname{Tks}(X, \mu)| > 3\epsilon.$$

Fix n big enough satisfying the last inequality and (3). Then either

$$\operatorname{Tks}(X_n, \mu_n) < \operatorname{Tks}(X, \mu)$$
 or $\operatorname{Tks}(X, \mu) < \operatorname{Tks}(X_n, \mu_n)$.

Next we analyse these two cases and deduce a contradiction for each.

If $\operatorname{Tks}(X,\mu) < \operatorname{Tks}(X_n,\mu_n)$, consider a sequence of height functions h_k such that

$$\lim_{k \to \infty} \operatorname{tks}_{h_k}(X, \mu) = \operatorname{Tks}(X, \mu).$$

Modulo extracting a subsequence we can assume that for all k we have that $0 \leq \operatorname{tks}_{h_k}(X,\mu) - \operatorname{Tks}(X,\mu) \leq \epsilon$. Observe that this difference is always non-negative by the definition of the trunkenness. Then

$$3\epsilon < \operatorname{Tks}(X_n, \mu_n) - \operatorname{Tks}(X, \mu)$$

=
$$\operatorname{Tks}(X_n, \mu_n) - \operatorname{tks}_{h_k}(X, \mu) + \operatorname{tks}_{h_k}(X, \mu) - \operatorname{Tks}(X, \mu),$$

hence,

$$2\epsilon < \operatorname{Tks}(X_n, \mu_n) - \operatorname{tks}_{h_k}(X, \mu) + \operatorname{tks}_{h_k}(X_n, \mu_n) - \operatorname{tks}_{h_k}(X_n, \mu_n) < \operatorname{Tks}(X_n, \mu_n) - \operatorname{tks}_{h_k}(X_n, \mu_n) + |\operatorname{tks}_{h_k}(X_n, \mu_n) - \operatorname{tks}_{h_k}(X, \mu)|$$

For k fixed, since μ_n tends to μ in the weak-* sense, we can assume that the term $|\operatorname{tks}_{h_k}(X_n,\mu_n) - \operatorname{tks}_{h_k}(X,\mu)|$ is smaller than ϵ by possibly taking n larger. Then we get $2\epsilon < \operatorname{Tks}(X_n,\mu_n) - \operatorname{tks}_{h_k}(X_n,\mu_n) + \epsilon$. We deduce $\epsilon + \operatorname{tks}_{h_k}(X_n,\mu_n) < \operatorname{Tks}(X_n,\mu_n)$, a contradiction to the definition since we have

$$\epsilon + \operatorname{tks}_{h_h}(X_n, \mu_n) < \operatorname{Tks}(X_n, \mu_n) \leq \operatorname{tks}_{h_h}(X_n, \mu_n).$$

The other case is similar. Assume now $\operatorname{Tks}(X_n, \mu_n) < \operatorname{Tks}(X, \mu)$ and for each n consider a sequence of height functions $h_{n,k}$ such that

$$\lim_{k \to \infty} \operatorname{tks}_{h_{n,k}}(X_n, \mu_n) = \operatorname{Tks}(X_n, \mu_n).$$

As in the previous case, we assume that for all k we have $0 \le \operatorname{tks}_{h_{n,k}}(X_n, \mu_n) - \operatorname{Tks}(X_n, \mu_n) \le \epsilon$. Then, for n large enough,

$$3\epsilon < \operatorname{Tks}(X,\mu) - \operatorname{Tks}(X_n,\mu_n)$$

$$= \operatorname{Tks}(X,\mu) - \operatorname{tks}_{h_{n,k}}(X_n,\mu_n) + \operatorname{tks}_{h_{n,k}}(X_n,\mu_n) - \operatorname{tks}_{h_{n,k}}(X,\mu)$$

$$+ \operatorname{tks}_{h_{n,k}}(X,\mu) - \operatorname{Tks}(X_n,\mu_n)$$

$$\leq \operatorname{Tks}(X,\mu) - \operatorname{tks}_{h_{n,k}}(X,\mu) + \epsilon + |\operatorname{tks}_{h_{n,k}}(X,\mu) - \operatorname{tks}_{h_{n,k}}(X_n,\mu_n)|$$

$$\leq \operatorname{Tks}(X,\mu) - \operatorname{tks}_{h_{n,k}}(X,\mu) + 2\epsilon.$$

We conclude that $\epsilon + \operatorname{tks}_{h_{n,k}}(X,\mu) < \operatorname{Tks}(X,\mu)$, a contradiction to the definition of $\operatorname{Tks}(X,\mu)$. Thus, for n large enough, $|\operatorname{Tks}(X_n,\mu_n) - \operatorname{Tks}(X,\mu)| < 3\epsilon$. This proves the first part of Theorem B.

The second part of Theorem B gives an asymptotic interpretation to the trunkenness of a vector field. Consider a sequence of knots $(K_n)_{n\in\mathbb{N}}$ that support vector fields X_n and denote by $(t_n)_{n\in\mathbb{N}}$ the respective periods. From Equation (2) we get that for each n

$$\sharp \{K_n \cap h^{-1}(t)\} = \lim_{\epsilon \to 0} \frac{1}{\epsilon} \mu_n(\phi_n^{[0,\epsilon]}(h^{-1}(t))),$$

where μ_n are the mesures supported by the knots and ϕ_n^t is the flow of X_n .

If we suppose that X_n tends to X and that the normalized linear Dirac measures $(\frac{1}{t_n}\mu_n)_{n\in\mathbb{N}}$ converge to μ in the weak-* sense, then we can reformulate the previous result into

$$\lim_{n \to \infty} \frac{1}{t_n} \operatorname{Tk}(K_n) = \operatorname{Tks}(X, \mu).$$

Now if X is ergodic with respect to μ then, for almost every p and for every sequence $t_n \to \infty$, the linear Dirac masses concentrated on the knots $K_n := k_X(p, t_n)$ tend to μ in the weak-* sense. Recall that $k_X(p, t_n)$ is the knot obtained by following the orbit of p for a time t_n concatenated with a geodesic from $\phi_X^{t_n}(p)$ to p, as explained in the introduction. Then the convergence implies that $\lim_{t\to\infty} \frac{1}{t} \operatorname{Tk}(k_X(p,t))$ exists and is equal to $\operatorname{Tks}(X,\mu)$. This proves the second part of Theorem B. \square

3. Independence of helicity

As mentioned in the introduction, the helicity is a well-known invariant of vector fields up to volume-preserving diffeomorphism. In this section, all vector fields are on the sphere \mathbb{S}^3 and preserve a volume form, that we denote by Ω . For X such a vector field, Cartan's formula implies that $\iota_X\Omega$ is a closed 2-form, and since the ambient manifold is simply connected it is exact. We may then write $\iota_X\Omega=d\alpha$, for α some differential 1-form. The helicity of X is defined as

$$\operatorname{Hel}(X) = \int_{\mathbb{S}^3} \alpha \wedge d\alpha,$$

and does not depends on the choice of the primitive α .

As we recalled in the introduction most known asymptotic invariants are in fact proportional to a power of helicity [Arn73, GG01, Baa11, BM12]. The goal of this section it to prove that the trunkenness of a vector field is not a function of its helicity. In order to do so we compute the trunkenness and the helicity of a vector field that preserves the invariant tori of a Hopf fibration of \mathbb{S}^3 .

Considering \mathbb{S}^3 as the unit sphere $\{(z_1, z_2) \in \mathbb{C}^2, |z_1|^2 + |z_2|^2 = 1\}$, the Seifert flow of slope (α, β) is the flow $\phi_{\alpha,\beta}$ given by

$$\phi_{\alpha,\beta}^t(z_1, z_2) := (z_1 e^{i2\pi\alpha t}, z_2 e^{i2\pi\beta t}).$$

This flow preserves the standard volume form, that is, the volume form Ω_{Haar} associated to the Haar measure of \mathbb{S}^3 . The flow has two distinctive periodic orbits corresponding to $z_1 = 0$ and $z_2 = 0$ that are trivial knots in \mathbb{S}^3 . The tori $|z_1/z_2| = r$ for $0 < r < \infty$ are invariant and the flow on each of them is the linear flow of slope α/β . If α/β is rational, put $\alpha/\beta = p/q$ with

 $p, q \in \mathbb{N}$ coprime. Then every orbit of $\phi_{\alpha,\beta}$, different from the two trivial ones, is a torus knot of type T(p,q).

The helicity of $\phi_{\alpha,\beta}$ is equal to $\alpha\beta$. To compute it in the rational case $(\alpha,\beta) = (p,q)$ with p,q coprime, observe that all the orbits except two are periodic of period 1. The linking number of an arbitrary pair of such orbits is pq. Therefore the asymptotic linking number (also called asymptotic Hopf invariant) equals pq and, by Arnold's Theorem [Arn73], so does the helicity.

Proposition 3.1. The trunkenness of the Seifert flow $\phi_{\alpha,\beta}$ with respect to the standard volume form Ω_{Haar} is equal to $2\min(\alpha,\beta)$.

Proof. Let us first prove $\mathrm{Tks}(X_{\alpha,\beta},\Omega_{\mathrm{Haar}}) \leq 2\beta$. For this it is enough to exhibit a height function h that yields $\mathrm{tks}_h(X_{\alpha,\beta},\Omega_{\mathrm{Haar}}) = 2\beta$. First define $\infty = (0,1)$ and 0 = (0,-1) in $\mathbb{S}^3 \subset \mathbb{C}^2$ and take the stereographic projection to identify

$$\{(z_1, z_2) \in \mathbb{C}^2, |z_1|^2 + |z_2|^2 = 1\} \simeq \mathbb{R}^3 \cup \{\infty\}.$$

Take now as h the standard height function h_0 of $\mathbb{R}^3 \cup \{\infty\}$. The spheres are centered at $0 \in \mathbb{R}^3$ that corresponds to the point $(0,-1) \in \mathbb{S}^3 \subset \mathbb{C}^2$, hence the orbit $z_1 = 0$ intersects twice each level sphere $h_0^{-1}(t)$. The middle sphere, $S = h_0^{-1}(1/2)$, contains the other special orbit $z_2 = 0$ and is the only sphere that intersects all the orbits of $\phi_{\alpha,\beta}$. Then the function $t \mapsto \int_{h_0^{-1}(t)} |\iota_{X_{\alpha,\beta}}\Omega_{\text{Haar}}|$ has a maximum for t = 1/2.

For computing $\int_S |\iota_{X_{\alpha,\beta}}\Omega_{\mathrm{Haar}}|$, we remark that the 2-sphere S has the orbit $(e^{i2\pi\alpha t},0)$ as an equator, that the flow is positively transverse to the northern hemisphere and negatively transverse to the southern hemisphere. Then the integral $\int_S |\iota_{X_{\alpha,\beta}}\Omega_{\mathrm{Haar}}|$ is equal to twice the flux of $X_{\alpha,\beta}$ through any disc bounded by the curve $(e^{i2\pi\alpha t},0)$. Consider the flat disc D in \mathbb{S}^3 bounded by $(e^{i2\pi\alpha t},0)$. The first return time to D is constant and equal to $1/\beta$, so the flux multiplied by $1/\beta$ gives the total volume of \mathbb{S}^3 , that is 1. Therefore $\mathrm{Flux}(X_{\alpha,\beta},\Omega_{\mathrm{Haar}},D)$ is equal to β , and we obtain $\mathrm{tks}_{h_0}(X_{\alpha,\beta},\Omega_{\mathrm{Haar}}) = \int_S |\iota_{X_{\alpha,\beta}}\Omega_{\mathrm{Haar}}| = 2\beta$. By symmetry, we then have $\mathrm{Tks}(X_{\alpha,\beta},\Omega_{\mathrm{Haar}}) \leqslant 2\min(\alpha,\beta)$.

For proving the converse inequality $\operatorname{Tks}(X_{\alpha,\beta},\Omega_{\operatorname{Haar}}) \geqslant 2\min(\alpha,\beta)$, we approximate $X_{\alpha,\beta}$ by a sequence $(X_{p_n/r_n,q_n/r_n})_{n\in\mathbb{N}}$, where p_n,q_n,r_n are integer numbers. Theorem B yields

$$\operatorname{Tks}(X_{\alpha,\beta},\Omega_{\operatorname{Haar}}) = \lim_{n \to \infty} \operatorname{Tks}(X_{p_n/r_n,q_n/r_n},\Omega_{\operatorname{Haar}}).$$

As the trunkenness is an order-1 invariant (it is multiplied by λ if the vector field is multiplied by λ), we only have to prove $\text{Tks}(X_{p,q},\Omega_{\text{Haar}}) = 2\min(p,q)$ for p,q two coprime natural numbers.

Since every orbit of $X_{p,q}$ is periodic, we can consider a sequence of $(K_n)_{n\in\mathbb{N}}$ of collections of periodic orbits whose induced linear Dirac measures converge to Ω_{Haar} . We take K_n to be an n-component link all of whose components are torus knots T(p,q). Actually K_n is a cabling with n strands on T(p,q), so by Zupan's theorem [Zup12], the trunk of K_n is $2n \min(p,q)$. Since the period of each component of K_n is 1, the total length of K_n is n, and we get $\text{Tks}(X_{p,q}, \Omega_{\text{Haar}}) \geqslant 2 \min(p,q)$.

Proof of Theorem C. The previous computations show that for a Seifert flow $X_{\alpha,\beta}$ on \mathbb{S}^3 we have $\operatorname{Hel}(X_{\alpha,\beta},\Omega_{\operatorname{Haar}}) = \alpha\beta$ and $\operatorname{Tks}(X_{\alpha,\beta},\Omega_{\operatorname{Haar}}) = 2\min(\alpha,\beta)$. There is no real function g such that $\min(\alpha,\beta) = g(\alpha\beta)$, so there is no function g such that $\operatorname{Tks}(X_{\alpha,\beta},\Omega_{\operatorname{Haar}}) = g(\operatorname{Hel}(X_{\alpha,\beta},\Omega_{\operatorname{Haar}}))$. \square

The formula $\operatorname{Tks}(X_{\alpha,\beta},\Omega_{\operatorname{Haar}})=2\min(\alpha,\beta)$ is also interesting to compare with Kudryavtseva's and Encisco-Peralata-Salas-Torres de Lizaur's theorems: the function $(\alpha,\beta)\mapsto 2\min(\alpha,\beta)$ is continuous but not differentiable, so that trunkenness is a continuous vector field invariant, but it is not *integral regular* in the sense of [Kud15, EPT16].

4. Trunkenness and the existence of periodic orbits

In this section we address the question of what happens when the infimum in the definition of the trunkenness is a minimum, for non-singular vector fields on \mathbb{S}^3 with an invariant mesure μ . We deduce that the vector field must posses an unknotted periodic orbit by finding a periodic orbit

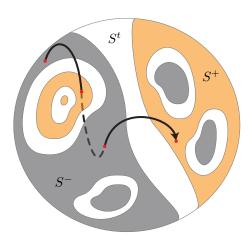


FIGURE 3. The decomposition of the level S of the height function into S^+ where X is positively transverse to S (orange), S^- where X is negatively transverse (black), and S^t where X is tangent to T (white). A piece of orbit of X that intersects S four times is also shown.

tangent to a level of the function. The proof of Theorem D in particular implies that there is height function h such that

$$\operatorname{tks}_h(X, \mu) = \operatorname{Flux}(X, \mu, h^{-1}(t_{\max}))$$

for some (not necessarily unique) $t_{\text{max}} \in [0,1]$ and such that all the connected components of $S = h^{-1}(t_{\text{max}})$ along which X is tangent to S are bounded by periodic orbits of S. The existence of vector fields on S without periodic orbits [Kup94] implies that there are vector fields for which we cannot consider a minimum to define the trunkenness.

Proof of Theorem D. For f a height function we define

$$F_f: [0,1] \rightarrow \mathbb{R}$$

 $t \mapsto \operatorname{Flux}(X, \mu, f^{-1}(t))$

Set h the height function such that $\operatorname{Tks}(X,\mu) = \operatorname{tks}_h(X,\mu)$ and let $t_{\max} \in [0,1]$ be a maximum of F_h , hence

$$\operatorname{Tks}(X,\mu) = \operatorname{tks}_h(X,\mu) = \lim_{\epsilon \to 0} \frac{1}{\epsilon} \mu(\phi_X^{[0,\epsilon]}(h^{-1}(t_{\max}))).$$

Denote by S the level set $h^{-1}(t_{\text{max}})$. Observe that the value t_{max} is not necessarily unique.

We assume that \mathbb{S}^3 and S are oriented and thus we can distinguish three subsets of S (see Figure 3):

- S^t the closed set along wich X is tangent to S;
- S^+ the open set along wich X is positively transverse to S:
- S^- the open set along wich X is negatively transverse to S.

We claim that none of these sets is empty. First assume that S^+ is empty, then $S = S^t \cup S^-$. Observe that $\mathbb{S}^3 \setminus S$ has two connected components that are diffeomorphic to open 3-dimensional balls. Hence there is one of these connected components, denote its closure by D, that is invariant under the diffeomorphism of \mathbb{S}^3 defined by ϕ_X^t for any t > 0, and since X is non-singular there exists $t_0 > 0$ such that $\Phi := \phi_X^{t_0}$ has no fixed points. Thus Φ maps D to D, and D is the closed 3-dimensional disc, hence by Brouwer fixed point theorem Φ has a fixed point, a contradiction. Thus S^+ is non-empty, and the same argument proves that S^- is non-empty. Now S^+ and S^- are open subsets of S and have empty intersection, hence $S^t = S \setminus (S^+ \cup S^-)$ is not empty.

Decompose $B := \partial S^t$ as $B = B^+ \cup B^-$, where B^+ and B^- are in the boundary of S^+ and S^- respectively. Observe that there is at least one connected component of S^t whose boundary has both positive and negative parts.

Assume that X has no periodic orbits, then B is made of circles and each one of these circles is tranverse to X in at least one point.

Lemma 4.1. If a point p in S^t is such that either its positive or negative orbit is contained in S, then X has an unknotted periodic orbit.

Proof. Assume without loss of generality that the positive orbit of p is contained in S, then it limits to some invariant set of X contained in S. By Poincaré-Bendixson Theorem such a set has to be a periodic orbit of X. Since the periodic orbit is contained in the sphere S it has to be unknotted.

Then under the assumption that X has no unknotted periodic orbits, we have that for every $p \in S^t$ the positive and negative orbit of p have to leave S at some point. Consider now two points $p, q \in B$ that are in the same orbit and such that the orbit segment connecting them is contained in S^t . We can thus assume that there exists $\tau \geq 0$ such that $\phi_X^{\tau}(p) = q$ and $\phi_X^s(p) \in S^t$ for all $0 \leq s \leq \tau$. Observe that by allowing τ to be zero, we consider the case p = q.

Take $\epsilon > 0$ and consider the orbit segment

$$\mathcal{O} = \{ \phi_X^s(p) \mid -\epsilon < s < \tau + \epsilon \}.$$

The flowbox Theorem implies that there is a neighborhood U of $\mathcal O$ and a diffeomorphism

$$\varphi : \mathbb{B}(1) \times (-\epsilon, \tau + \epsilon) \to U$$

so that the flowlines are the image under φ of the vertical segments $\{\cdot\} \times (-\epsilon, \tau + \epsilon)$. Here $\mathbb{B}(1)$ denotes the 2-dimensional open disc of radius 1.

Proposition 4.2. There exists $p, q \in B$ as above such that $p \in B^-$ and $q \in B^+$ or $p \in B^+$ and $q \in B^-$.

The idea of the proof is that if for every pair of points p, q as above that are both in B^+ or B^- , we can change the function h for another height function h_1 such that

$$\operatorname{tks}_h(X,\mu) = \operatorname{tks}_{h_1}(X,\mu)$$

and such that the level of h_1 realizing the trunkenness has no tangent part that separates S^+ from S^- , which is impossible.

Proof. Consider a pair of points $p, q \in B^+$ as above (see Figure 4). Let V_0 be the neighborhood of \mathcal{O} defined as $\varphi(\mathbb{D}(1/3) \times (-\epsilon, \tau + \epsilon))$, where $\mathbb{D}(1/3)$ is the 2-dimensional closed disc of radius 1/3 and $V = \varphi(\mathbb{B}(2/3) \times (-\epsilon, \tau + \epsilon))$. Hence $V_0 \subset V \subset U$. We will deform the levels of h intersecting V without changing the trunkenness.

The image of S by φ^{-1} defines a surface Σ that is positively transverse to the vertical direction in $\mathbb{B}(1) \times (-\epsilon, 0)$ and $\mathbb{B}(1) \times (\tau, \tau + \epsilon)$, as in Figure 5 left. If we project Σ to $\mathbb{B}(1) \times \{0\}$, the flux through Σ and through the image of Σ under the projection are the same. Thus if we change Σ (or any surface) for another surface whose projection, counted with multiplicities and signs, is the same as the one of Σ , the flux remains constant.

Inside $\mathbb{D}(1/3) \times (-\epsilon, \tau + \epsilon)$ we change the images under φ^{-1} of the levels of h to obtain a family of surfaces that are always positively transverse to X and are C^{∞} -close to the original ones, in such a way that the projection to $\mathbb{B}(1) \times \{0\}$ is preserved (as in Figure 5 right). Using the neighborhood V, we can paste the deformed surfaces of V_0 with the original surfaces in $U \setminus V$.

Using φ this deformation can be pushed forward to the manifold so that the level surfaces of h that intersected V are modified and X is positively transverse to the surfaces inside V_0 . Let \widetilde{S} be the surface obtained from S after the deformation and \widetilde{h} be a height function whose level sets correspond to the deformed surfaces. Then

$$\operatorname{tks}_h(X,\mu) = \operatorname{tks}_{\widetilde{h}}(X,\mu),$$

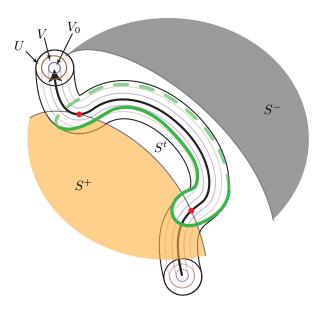


FIGURE 4. A neighborhood U of a piece of orbit \mathcal{O} of X that connects two points p, q in B^+ . The intersection of ∂U with S is bold and green.

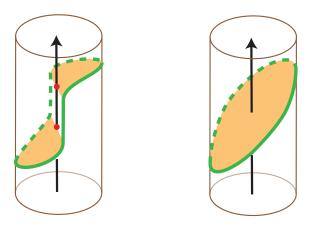


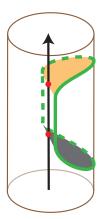
FIGURE 5. Modifying the levels of h in U around the arc pq in the case $p, q \in B^+$.

since $\operatorname{Flux}(X,\mu,\widetilde{S})=\operatorname{Flux}(X,\mu,S)$, and for every surface intersection V the corresponding equation holds. Clearly, the same proof works if $p,q\in B^-$.

Observe that the level \widetilde{S} of \widetilde{h} realizing the trunkenness coincides with S outside V. Moreover, the tangent part \widetilde{S}^t of \widetilde{S} is strictly smaller than S^t : the connected component A of S^t containing p and q got transformed into $A \setminus (V_0 \cap A)$. In other words there is now a strip accross A.

Assume now that for any couple of points $p, q \in B$ that is joined by an orbit segment contained in S^t we have either $p, q \in B^+$ or $p, q \in B^-$. The sets \widetilde{S}^t and B are compact, we can thus find a finite number of pairs of points $p, q \in \partial B$ such that after deforming the levels h as above along each of the corresponding orbit segments, we obtain a height function h_1 such that

$$\operatorname{tks}_h(X,\mu) = \operatorname{tks}_{h_1}(X,\mu) = \operatorname{Flux}(X,\mu,S_1),$$



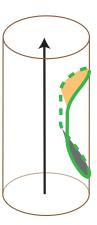


FIGURE 6. Modifying the levels of h in U around the arc pq in the case $p \in B^+$, $q \in B^-$.

where S_1 is the level of h_1 realizing the trunkenness and such that S_1^t is formed by discs. Hence none of the connected components of S_1^t separates the positively transverse set S_1^+ from the negatively transverse set S_1^- , a contradiction.

Consider now a pair of points p, q as in Proposition 4.2 and assume that $p \in B^-$ and $q \in B^+$. We now construct a height function h_2 such that

$$\operatorname{tks}_h(X,\mu) > \operatorname{tks}_{h_2}(X,\mu),$$

a contradiction that implies that X has unknotted periodic orbits.

Consider the function F_h and let M_h be the set of $t \in [0,1]$ realizing the trunkenness, that is $M_h = \{t \in [0,1] \mid F_h(t) = \operatorname{tks}_h(X,\mu)\}$. Observe that M_h is closed. We distinguish two cases:

- (1) all the elements of M_h are isolated;
- (2) there exist a closed interval I_{max} contained in M_h .

We start with the first case. Take $t_{\text{max}} \in M_h$ and set $S = h^{-1}(t_{\text{max}})$. To obtain h_2 we will deform the levels near S, if M_h has more than one element the following deformation has to be done near each corresponding level set. Thus we assume without loss of generality that t_{max} is the only element of M_h .

Let \mathcal{O} , φ , U, V and V_0 be defined as above. Let $[a,b] \subset [0,1]$ be such that $t_{\max} \in [a,b]$ and if $h^{-1}(t) \cap V_0 \neq \emptyset$ then $t \in [a,b]$. Shrinking V if necessary, we can assume that $\varphi(\mathbb{B}(2/3) \times (-\epsilon, \tau + \epsilon))$ is contained in $h^{-1}([a,b])$ and that for any $t \in [0,1] \setminus [a,b]$ we have

$$\operatorname{Flux}(X, \mu, h^{-1}(t)) < \operatorname{tks}_h(X, \mu) - \delta,$$

for some $\delta > 0$. We assume also that $\varphi(\mathbb{B}(2/3) \times [0,\tau])$ is contained in $h^{-1}([a,b])$.

Consider now the surface Σ obtained as the image under φ^{-1} of S in U. Then either Σ is positively transverse to the vertical direction in $\mathbb{B}(1) \times (-\epsilon, 0)$ and negatively transverse to the vertical direction in $\mathbb{B}(1) \times (\tau, \tau + \epsilon)$, or it is negatively transverse to the vertical direction in $\mathbb{B}(1) \times (-\epsilon, 0)$ and positively transverse to the vertical direction in $\mathbb{B}(1) \times (\tau, \tau + \epsilon)$. We assume without loss of generality that Σ is negatively transverse to the vertical direction in $\mathbb{B}(1) \times (-\epsilon, 0)$ and positively transverse to the vertical direction in $\mathbb{B}(1) \times (\tau, \tau + \epsilon)$, as represented in Figure 6.

We want to deform the levels of h intersecting V_0 in such a way that their trace under the projection to $\mathbb{B}(1) \times \{0\}$ is reduced.

Shrinking V if necessary, we can assume that the level sets are:

- tangent or negatively transverse to X in $\varphi(\mathbb{B}(2/3) \times (-\epsilon, 0))$,
- tangent or positively transverse to X in $\varphi(\mathbb{B}(2/3) \times (\tau, \tau + \epsilon))$.

Consider the circle $C = \Sigma \cap \partial(\mathbb{D}(1/3) \times (-\epsilon, \tau + \epsilon))$. Since $\partial(\mathbb{D}(1/3) \times (-\epsilon, \tau + \epsilon))$ is topologically a sphere, C divides it into two discs. Let W be the disc that is entirely tangent to the vertical direction. There is a continuous deformation from Σ to the continuous surface $\widetilde{\Sigma}$ that coincides

with Σ outside $\mathbb{D}(1/3) \times (-\epsilon, \tau + \epsilon)$ and coincides with W in this set. Observe that the projection of $\widetilde{\Sigma}$ to $\mathbb{B}(1) \times \{0\}$ is a proper subset of the projection of Σ . We can now approximate $\widetilde{\Sigma}$ in $\mathbb{B}(2/3) \times (-\epsilon, \tau + \epsilon)$ with a smooth surface Σ_1 , that can be obtained from Σ by a continuous deformation that is the identity near the boundary of $\mathbb{B}(2/3) \times (-\epsilon, \tau + \epsilon)$.

Apply this deformation to nearby levels by pushing (to the right as in Figure 6) all the level sets intersecting $\mathbb{B}(1/3) \times (0,\tau)$ with a continuous map that is the identity near the boundary of $\mathbb{B}(2/3) \times (-\epsilon, \tau + \epsilon)$. By construction the projection to $\mathbb{B}(1) \times \{0\}$ of the image under φ^{-1} of any level set intersecting V is a subset of the original one and for levels sufficiently near $h^{-1}(t_{\text{max}})$ it is a proper subset. Thus the flux of X gets reduced on any of the deformed surfaces.

As before, we use φ to push forward the deformation to \mathbb{S}^3 and obtain the levels of a new height function h_2 such that $\operatorname{tks}_{h_2}(X,\mu) < \operatorname{tks}_h(X,\mu)$.

The proof of the case where M_h has only isolated elements ends here and we are left with the case where the set M_h contains an interval $I_{\text{max}} = [a, b]$. Take $t_{\text{max}} = a$ and use the previous procedure to obtain a new height function h_2 such that the interval I_{max} is reduced to $I_2 = [a_2, b]$ for some $a_2 > a$. This process can be applied as long as the boundary of the tangent part of $h^{-1}(t)$ for any $t \in [a, b]$ is not composed by periodic orbits of X. Recursively either we find an unknotted periodic orbit or we obtain a height function f such that $\text{tks}_f(X, \mu) < \text{tks}_h(X, \mu)$ as needed. \square

The proof of Theorem D gives the following result.

Corollary 4.3. Let X be a non-singular vector field on \mathbb{S}^3 , μ an invariant measure and h a height function such that

$$\operatorname{Tks}(X,\mu) = \operatorname{tks}_h(X,\mu).$$

Then there is a height function h_1 such that:

- $\operatorname{tks}_h(X,\mu) = \operatorname{tks}_{h_1}(X,\mu);$
- for every $s \in M_{h_1}$ and for every point p in the boundary of the tangent part of $S = h_1^{-1}(s)$, either p belongs to a periodic orbit or its orbit limits to a periodic orbit. In both cases the periodic orbit is contained in S and is thus unknotted.

5. Trunkenness of knotted tubes

In this section we compute the trunkenness with respect to a volume form of some vector fields supported in a tubular neighborhood of a link or knot. Our statement is reminiscent from Zupan's theorem [Zup12] concerning the trunk of the cable of a knot.

Proposition 5.1. (see Figure 7 right) Suppose that X is a Ω -preserving vector field supported on tubes T_1, \ldots, T_j that are tubular neighborhoods of knots k_1, \ldots, k_j , each knot k_i being non-trivial, and such that X is transverse to all canonical meridian discs and the flux of X through each of them is constant equal to some $F \neq 0$, then we have

$$\operatorname{Tks}(X,\Omega) = F \cdot \operatorname{Tk}(k_1 \cup \cdots \cup k_i).$$

Note that the trunkenness in this case is independent from the dynamics of X inside each tube, and in particular of the first-return map on a meridian disc, exactly like the trunk of a cable link is independent of the twist number of the cabling. However the restriction that all the knots k_i are non-trivial is necessary. Indeed the trunkenness of a vector field supported on a tube that forms a trivial knot may be smaller than 2F. For example take an unknotted solid torus in \mathbb{R}^3 that supports a vector field obtained by the suspension of the identity (with no extra twist) and assume that the vector field is zero outside the solid torus (see Figure 7 top left). The trunkenness of such vector field is zero: it is enough to consider a height function whose levels are always tangent to the vector field. Observe that the levels of such a height function intersect the solid torus along circles or annuli. As this example shows, the trunkenness of a vector field supported in the neighborhood of a link with at least one unknotted component seems hard to determine (see Figure 7).

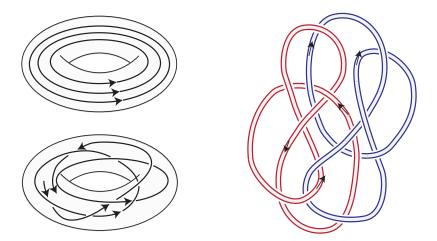


FIGURE 7. On the left, two vector fields supported on the unlink and both transverse to the canonical meridian discs. The top one has all its orbits horizontal, hence its trunkenness is 0 while the bottom one has linked orbits. Its trunkenness is at most 2 · Flux (X, Ω, D) but it might be smaller. On the right a vector field supported on a link L both of which components are figure-eight knots. If the flux of the vector field through the canonical meridian discs of both components is equal to F, then the trunkenness of this vector field is $F \cdot \text{Tk}(L) = 8F$ by Proposition 5.1.

Proof. Let $L = k_1 \cup k_2 \cup \ldots \cup k_n \subset \mathbb{S}^3$ be a link all of whose components are non-trivial knots. Let h be a height function such that $Tk(L) = tk_h(L)$. Observe that we can assume that h exists since for a knot or link the trunk is defined by a minimum (see Definition 1.1). The trunk of L is then realized by one or more level sets of h, that is, there exist t_1, t_2, \dots, t_k such that for any $1 \le i \le k$ there is a critical value c_i of $h|_L$ so that $t_i < c_i < t_{i+1}$ and $Tk(L) = \sharp\{L \cap h^{-1}(t_i)\}$. Let

$$m_j(t_i) = \sharp \{k_j \cap h^{-1}(t_i)\}, \quad \text{for} \quad 1 \le j \le n \quad \text{and} \quad 1 \le i \le k.$$

For any $0 < \epsilon < 1$ there is an Ω -preserving diffeomorphism f_{ϵ} of \mathbb{S}^3 that makes the solid tori T_j longer and of radius ϵr_j for any $1 \leq j \leq n$, with r_j the radius of the tubular neighborhood T_j . Let X_{ϵ} be the vector field obtained from X via f_{ϵ} , then $\mathrm{Tks}(X,\Omega) = \mathrm{Tks}(X_{\epsilon},\Omega)$ for every ϵ .

Assume that there is a height function g and an index i such that $\operatorname{tks}_g(X,\Omega) < F \cdot \sum_{j=1}^n m_j(t_i)$. Consider the height function $g_{\epsilon} = g \circ f_{\epsilon}^{-1}$, then Theorem A implies $\operatorname{tks}_g(X,\Omega) = \operatorname{tks}_{g_{\epsilon}}(X,\Omega)$ and

$$\operatorname{tks}_{g_{\epsilon}}(X_{\epsilon}, \Omega) < F \cdot \sum_{j=1}^{n} m_{j}(t_{i}).$$

We claim that $\operatorname{tk}_{g_{\epsilon}}(L) < \sum_{i=1}^{n} m_{i}(t_{i}) = \operatorname{Tk}(L)$ holds, which is absurd.

Assume first that g_{ϵ} restricted to the core of the tori $f_{\epsilon}(T_i)$ is a Morse function. The core of the tori form the link L thus $\operatorname{tk}_{g_{\epsilon}}(L) \geq \operatorname{Tk}(L)$ in one hand. But since

$$\operatorname{tks}_{g_{\epsilon}}(X_{\epsilon}, \Omega) = \operatorname{Flux}(X_{\epsilon}, \Omega, g_{\epsilon}^{-1}(T)) < F \cdot \min_{1 \le i \le k} \sum_{j=1}^{n} m_{j}(t_{i})$$

for some $T \in [0,1]$, taking ϵ arbitrarily small tells us that the number of discs in the intersection of $g_{\epsilon}^{-1}(T)$ and the tori is smaller than $\mathrm{Tk}(L)$. In other words we conclude that $\mathrm{tk}_{g_{\epsilon}}(L) < \mathrm{Tk}(L)$, which is impossible. We are left with the case where g_{ϵ} is not a Morse function when restricted to L, seen as the core of the tori $f_{\epsilon}(T_i)$. The following lemma finishes the proof of the proposition.

Lemma 5.2. In the previous context, there exists another height function g''_{ϵ} so that:

- g''_ε|_L is a Morse function;
 tks_{g_ε}(X_ε, Ω) = tks_{g''}(X_ε, Ω).

Proof. Let $S = g_{\epsilon}^{-1}(t)$ be a level that is tangent to at least one of the components of L and let k_1 be one of the tangent components. We will modify g_{ϵ} in a neighborhood of k_1 , this modification taking place near the part of k_1 that is tangent to S.

Observe that S is tangent to a closed subinterval of k_1 , since if it were tangent to all of k_1 then k_1 would be the trivial knot, in contrary to our assumption on L. For $\delta < \epsilon r_1$, consider $N(\delta, k_1)$ the δ -tubular neighborhood of k_1 that is contained in $f_{\epsilon}(T_1)$. For δ sufficiently small, the intersection $S \cap N(\delta, k_1)$ is composed by a finite number of discs. Let D_1 be one of these discs, then ∂D_1 is included in $\partial N(\delta, k_1)$ and is either a contractible circle in $\partial N(\delta, k_1)$ or not.

If ∂D_1 is contractible in $\partial N(\delta, k_1)$, then we can locally deform the levels of g_{ϵ} near D_1 , to obtain a new height function g'_{ϵ} such that:

- the disc D'_1 obtained from D_1 is tangent to k_1 at one point;
- near D'_1 all levels are transverse to k_1 ;
- $\operatorname{tks}_{g_{\epsilon}}(X_{\epsilon}, \Omega) = \operatorname{tks}_{g'_{\epsilon}}(X_{\epsilon}, \Omega).$

The support of the deformation is contained in $N(\delta, k_1)$. Performing this deformation near any disc in $S \cap N(\delta, k_1)$ whose boundary is contractible in $\partial N(\delta, k_1)$, we can assume now that the discs of S that are tangent to k_1 have non-contractible boundary in $\partial N(\delta, k_1)$. Thus these discs are homologous to meriodional discs $\mathbb{D}^2 \times \{\theta\}$ of T_1 . Again, we can locally deform the levels, so that this discs are transverse to k_1 . The deformed height function g''_{ϵ} is Morse when restricted to L.

Thus g''_{ϵ} is Morse when restricted to k_1 . Repeating this process if necessary, we obtain a that g''_{ϵ} is Morse when restricted to L.

References

[Arn73] ARNOLD Vladimir I.: The asymptotic Hopf invariant and its applications. *Proc. Summer School in Diff. Equations at Dilizhan*, 1973 (1974), Evevan (in Russian); English transl. *Sel. Math. Sov.* 5 (1986), 327–345.

[AK98] Arnold Vladimir I. and Khesin Boris: Topological methods in hydrodynamics, Appl. Math. Sci. 125, Springer (1998).

[Baa11] Baader Sebastian: Asymptotic concordance invariants for ergodic vector fields, Comment. Math. Helv. 86 (2011), 1–12.

[BM12] BAADER Sebastian and MARCHÉ Julien: Asymptotic Vassiliev invariants for vector fields, Bull. Soc. Math. France 140 (2012), 569–582.

[EPT16] Enciso Alberto, Peralta-Salas Daniel and Torres de Lizaur Francisco: Helicity is the only integral invariant of volume-preserving transformations: arXiv:1602.04745.

[FH91] FREEDMAN Michael H. and HE Zheng-Xu: Divergence-Free Fields: Energy and Asymptotic Crossing Number, Ann. of Math. (2) 134 (1991), 189–229.

[Gab87] Gabai David: Foliations and the topology of 3-manifolds III, J. Differential Geom. 26 (1987), 479–536.

[GG01] GAMBAUDO Jean-Marc and GHYS Étienne, Signature asymptotique d'un champ de vecteurs en dimension 3, Duke Math. J. 106 (2001), 41–79.

[Hel1858] VON HELMHOLTZ Hermann: Über Integrale der hydrodynamischen Gleichungen, welche den Wirbelbewegungen entsprechen, J. Reine Angew. Math. 55 (1858), 25–55.

[Kud14] KUDRYAVTSEVA Elena A.: Conjugation invariants on the group of area-preserving diffeomorphisms of the disk, Math. Notes 95 (2014), 877–880.

[Kud15] Kudryavtseva Elena A.: Helicity is the only invariant of incompressible flows whose derivative is continuous in C^1 -topology, arXiv:1511.03746.

[Kup94] Kuperberg Krystyna, A smooth counterexample to the Seifert conjecture, Ann. of Math. (2) 140 (1994), 723–732.

[Mof69] MOFFATT Keith: The degree of knottedness of tangle vortex lines, J. Fluid. Mech. 106 (1969), 117–129.

[Mor61] MOREAU Jean-Jacques: Constantes d'un îlot tourbillonnaire en fluide parfait barotrope., C. R. Acad. Sci. Paris 252 (1961), 2810–2812.

[Ozw10] Ozawa Makoto: Waist and trunk of knots, Geom. Dedicata 149 (2010), 8594.

[Tai1877] Tait Peter Guthrie: On knots, Trans. Roy. Soc. Edin. 28 (1877), 145–190.

[Tho1867] THOMSON William: On Vortex Atoms, Proc. Roy. Soc. Edinburgh 6 (1867), 94–105; reprinted in Phil. Mag. 34 (1867), 15–24.

[Vog02] Vogel Thomas: On the asymptotic linking number, Proc. Amer. Math. Soc. 131 (2002), 2289–2297.

[Wol58] Woltjer Lodewijk: A Theorem on Force-free magnetic Fields, Proc. Natl. Acad. Sci. USA 44 (1958), 489–491.

[Zup12] Zupan Alexander: A lower bound on the width of satellite knots, Topology Proc. 40 (2012), 179–188.