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UNIFORMLY ACCURATE NUMERICAL SCHEMES FOR THE NONLINEAR DIRAC EQUATION IN THE NONRELATIVISTIC LIMIT REGIME

MOHAMMED LEMOU*, FLORIAN MÉHATS †, AND XIAOFEI ZHAO ‡

Abstract. We apply the two-scale formulation approach to propose uniformly accurate (UA) schemes for solving the nonlinear Dirac equation in the nonrelativistic limit regime. The nonlinear Dirac equation involves two small scales ε and ε^2 with $\varepsilon \rightarrow 0$ in the nonrelativistic limit regime. The small parameter causes high oscillations in time which bring severe numerical burden for classical numerical methods. We present a suitable two-scale formulation as a general strategy to tackle a class of highly oscillatory problems involving the two small scales ε and ε^2 . A numerical scheme with uniform (with respect to $\varepsilon \in (0, 1)$) second order accuracy in time and a spectral accuracy in space are proposed. Numerical experiments are done to confirm the UA property.

Key words. nonlinear Dirac equation, nonrelativistic limit, highly oscillatory equations, uniform accuracy, two-scale formulation

AMS subject classifications. 35Q55, 65M12, 74Q10

1. Introduction The nonlinear Dirac equation has been widely considered in many physical and mathematical areas, such as the electron self-interacting [15], the gravity theory [19], and recent studies in material graphene and Bose-Einstein condensates [1, 13, 17]. In this paper, we consider the nonlinear reduced Dirac equation

$$i\partial_t \Phi^\varepsilon = -\frac{i}{\varepsilon} \alpha \partial_x \Phi^\varepsilon + \frac{1}{\varepsilon^2} \beta \Phi^\varepsilon + [V_e(t) + V_m(t)\alpha] \Phi^\varepsilon + \lambda (\beta \Phi^\varepsilon, \Phi^\varepsilon) \beta \Phi^\varepsilon, \quad (1.1)$$

for $t > 0, x \in \mathbb{R}$, where $\Phi^\varepsilon = (\phi_1^\varepsilon, \phi_2^\varepsilon)^T = \Phi^\varepsilon(t, x)$ is a bi-dimensional complex-valued column vector and is subject to the initial condition

$$\Phi^\varepsilon(0, x) = \Phi_0(x), \quad x \in \mathbb{R}.$$

$V_e(t) = V_e(t, x)$ and $V_m(t) = V_m(t, x)$ are two given real-valued scalar functions representing the electrical potential and magnetic potential [2, 16], respectively, and the matrices α and β are the known as the Pauli matrices, i.e.

$$\alpha = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \beta = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}.$$

In particular, we have $(\beta \Phi^\varepsilon, \Phi^\varepsilon) = |\phi_1^\varepsilon|^2 - |\phi_2^\varepsilon|^2$. The nonlinear Dirac equation (1.1) has been widely considered in the literature [5, 4, 27, 21] as a reduced mathematical model. It was originally derived in [22] as a four-component equations in 3D for describing the spinor field with nonlinear coupling. A more complicated cubic nonlinearity could be considered as in [25, 26]. The unknown Φ^ε is interpreted as the wave function, $\lambda \in \mathbb{R}$ denotes the coupling constant and $\varepsilon \in (0, 1)$ is a dimensionless parameter which is inversely proportional to the speed of light. We refer to [2, 5] for the nondimensionalization and the model reduction to the two-component problem (1.1).

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The model (1.1) conserves the mass $M(t) := \int_{\mathbb{R}} |\Phi^\varepsilon|^2 dx$ and the energy $E(t)$, provided that the two external potentials V_e, V_m are independent of time,

$$E(t) := \int_{\mathbb{R}} \left(\frac{1}{\varepsilon^2} (\Phi^\varepsilon, \beta \Phi^\varepsilon) - \frac{i}{\varepsilon} (\Phi^\varepsilon, \alpha \partial_x \Phi^\varepsilon) + V_e |\Phi^\varepsilon|^2 + V_m (\Phi^\varepsilon, \alpha \Phi^\varepsilon) + \lambda (\beta \Phi^\varepsilon, \Phi^\varepsilon) |\Phi^\varepsilon|^2 \right) dx.$$

With a fixed $\varepsilon > 0$, the well-posedness of the Cauchy problem (1.1) has been studied and established. We refer to [4, 18, 20, 21, 3] and the references therein for detailed analytical results. Various numerical methods including finite difference time domain (FDTD) methods and operator splitting methods have also been considered in [23, 26, 27, 4] for solving the nonlinear Dirac equation in the classical regime, i.e. $\varepsilon = O(1)$.

When $\varepsilon \rightarrow 0$, which corresponds to the speed of light going to infinity and is known as the nonrelativistic limit, the nonlinear Dirac equation (1.1) has been shown to converge in the energy space to the nonlinear Schrödinger equations [2, 20]. Both the analytical results and the numerical studies [2, 20, 5, 6] show that the solution Φ^ε of (1.1) propagates waves with wavelength $O(\varepsilon^2)$ in time t and wavelength $O(1)$ in space x , as $\varepsilon \ll 1$. The small temporal wavelength makes the solution highly oscillatory in time, and as a consequence it causes severe numerical burden for classical numerical discretizations. Results in [5] show that, in order to capture the correct solution in the limit regime, one needs to use time step $\Delta t = O(\varepsilon^2)$ and mesh size $\Delta x = O(\sqrt{\varepsilon})$ for FDTD, and $\Delta t = O(\varepsilon^2)$, $\Delta x = O(1)$ for exponential integrators or time splitting spectral methods. Thus designing numerical methods that allow the use of step size independent of ε becomes a very needed and challenging issue. Recently, in study of the Klein-Gordon equation in nonrelativistic limit regime, schemes with uniformly accurate (UA) properties have been proposed in [7, 8, 10]. The UA schemes can use $\Delta t = O(1)$, $\Delta x = O(1)$ for corrected approximations when $\varepsilon \ll 1$ and are significantly better than classical methods in the limit regime and asymptotic preserving (AP) methods in the intermediate regime, as shown in [9, 12]. Among the UA schemes, there are two approaches so far. One is based on the multiscale expansion of the solution [7, 8], and the other uses the two-scale formulation approach [10, 11]. The first approach strongly relies on the pre-knowledge of the expansion and the polynomial type nonlinearity. In fact, a multiscale approach has already been considered for the linear Dirac in [6], leading to a first order UA scheme. However, this proposed scheme and its corresponding numerical analysis become rather complicated in the nonlinear case. In contrast, the two-scale formulation approach provides a general strategy to design UA schemes for a class of oscillatory problems [11]. Nevertheless, the nonlinear Dirac equation (1.1) which contains two scales ε and ε^2 , does not belong to the type of problems in [11], which means that the strategy in [10, 11] cannot be directly applied. In fact, problems involving two scales are usual in collisional kinetic equations in the diffusion limit, but the presence of these two scales (ε and ε^2) also appears in the context of oscillatory kinetic equations, as mentioned in [11, 14]. This extra scaling in the equations causes more difficulties for achieving UA property and a suitable two-scale formulation has to be found.

In this work, we are going to propose a two-scale formulation and design UA schemes for solving the nonlinear Dirac equation in the nonrelativistic limit regime. A suitable formulation is presented on which we construct a second order in time UA scheme. Extensive numerical experiments are done to show the UA property. Extension of this strategy to the case of oscillatory kinetic equations with strong magnetic field is the subject of a forthcoming paper.

The rest of the paper is organized as follows. In Section 2, we introduce the two-

scale formulation and construct the suitable initial data for this augmented problem. In Section 3, we present the numerical schemes based on this two-scale formulation. Numerical results are reported in Section 4 and some conclusions are finally drawn in Section 5.

2. Two scale formulation To filter out the main oscillation in (1.1), let us introduce the filtered data

$$u^\varepsilon = \begin{pmatrix} e^{it/\varepsilon^2} & 0 \\ 0 & e^{-it/\varepsilon^2} \end{pmatrix} \Phi^\varepsilon. \quad (2.1)$$

It solves the equation

$$\partial_t u^\varepsilon = -\frac{1}{\varepsilon} A(t/\varepsilon^2) \partial_x u^\varepsilon + F(t, t/\varepsilon^2, u^\varepsilon), \quad t > 0, x \in \mathbb{R}, \quad (2.2)$$

with

$$A(\tau) = \begin{pmatrix} 0 & e^{2i\tau} \\ e^{-2i\tau} & 0 \end{pmatrix},$$

and

$$F(t, \tau, u^\varepsilon) = -i[V_e(t) + V_m(t)A(\tau)]u^\varepsilon - i\lambda(\beta u^\varepsilon, u^\varepsilon)\beta u^\varepsilon. \quad (2.3)$$

In this paper, we assume that V_e and V_m are C^1 functions of time. For future convenience, we denote the nonlinearity without the magnetic potential by

$$F_e(t, u^\varepsilon) := -i[V_e(t)u^\varepsilon + \lambda(\beta u^\varepsilon, u^\varepsilon)\beta u^\varepsilon].$$

Now, following [10, 11], we separate the fast time variable $\tau = t/\varepsilon^2$ from the slow time variable t in the solution $u^\varepsilon(t, x)$ of (2.2) and consider the two-scale formulation for the augmented unknown $U^\varepsilon = U^\varepsilon(t, \tau, x) \in \mathbb{C}^2$:

$$\partial_t U^\varepsilon + \frac{1}{\varepsilon^2} \partial_\tau U^\varepsilon = -\frac{1}{\varepsilon} A(\tau) \partial_x U^\varepsilon + F(t, \tau, U^\varepsilon), \quad t > 0, \tau \in \mathbb{T}, x \in \mathbb{R}. \quad (2.4)$$

Here $U^\varepsilon(t, \tau, x)$ is 2π -periodic in τ and $\mathbb{T} := \mathbb{R}/(2\pi\mathbb{Z})$ denotes the torus. With initial data $U^\varepsilon(0, \tau, x)$ satisfying

$$U^\varepsilon(0, 0, x) = u^\varepsilon(0, x) = \Phi_0(x), \quad x \in \mathbb{R},$$

the solution of the two-scale problem (2.4) enables to recover the solution of the filtered equation (2.2) by setting

$$U^\varepsilon \left(t, \frac{t}{\varepsilon^2}, x \right) = u^\varepsilon(t, x), \quad t \geq 0, x \in \mathbb{R}.$$

In the two-scale formulation (2.4), τ is considered as an additional independent variable, which is periodic. The advantage one can get from this formulation is that now the initial data $U^\varepsilon(0, \tau, x)$ is only prescribed at a single point $\tau = 0$, so there is some freedom to choose the initial data. By choosing a suitable initial data, we will succeed to bound the time derivatives of U^ε uniformly with respect to ε , which allows to design uniformly accurate numerical schemes.

2.1. A toy model In order to find out a suitable initial data for problems like (2.4), we propose to analyze the following simple toy model:

$$\partial_t u + \frac{1}{\varepsilon^2} \partial_\tau u + i \frac{a(\tau)}{\varepsilon} u = 0, \quad \text{with } u(0, \tau) = u_{in}(\tau). \quad (2.5)$$

Here, $u(t, \tau)$ is a scalar unknown, $t \geq 0$ and $\tau \in \mathbb{T}$. Moreover, we assume (as for the augmented problem (2.4)) that only $u_{in}(0) = u_0$ is prescribed, that a is a real-valued smooth periodic function and that we have the property

$$\int_0^{2\pi} a(\tau) d\tau = 0.$$

Our aim is to derive expressions for $u_{in}(\tau)$ in order to ensure conditions (i) $u_{in}(0) = u_0$, and (ii) that time-derivative of u , up to some order $p \geq 1$, are uniformly bounded with respect to ε , a property which enables to construct easily uniformly accurate numerical schemes for (2.5).

The exact solution of (2.5) is given by

$$u(t, \tau) = \exp(-i\varepsilon b(\tau)) \exp(i\varepsilon b(\tau - t/\varepsilon^2)) u_{in}(\tau - t/\varepsilon^2),$$

where $b(\tau) = \int_0^\tau a(s) ds$. Note that b is periodic, since the average of a vanishes. It is clear that

$$\frac{\partial^k u}{dt^k}(t, \tau) = \exp(-i\varepsilon b(\tau)) \frac{(-1)^k}{\varepsilon^{2k}} \frac{d^k}{ds^k} \left(\exp(i\varepsilon b(s)) u_{in}(s) \right) \Big|_{s=\tau-t/\varepsilon^2}.$$

Therefore, in order to get uniformly bounded derivatives for $k = 0, \dots, p$, we must have

$$\frac{d^p}{ds^p} \left(\exp(i\varepsilon b(s)) u_{in}(s) \right) = \mathcal{O}(\varepsilon^{2p}),$$

which is satisfied as soon as

$$u_{in}(\tau) = Q_{p-1}(\tau) \exp(-i\varepsilon b(\tau)) + \mathcal{O}(\varepsilon^{2p}).$$

Here Q_{p-1} is any polynomial of degree $\leq p-1$. Due to the smoothness and periodicity of u_{in} , we have necessarily that Q is a constant polynomial, thus

$$u_{in}(\tau) = C \exp(-i\varepsilon b(\tau)) + \mathcal{O}(\varepsilon^{2p}).$$

Therefore, the following initial data is suitable to ensure (i) and (ii):

$$u_{in}(\tau) = u_0 + u_0 \sum_{k=1}^{2p-1} \frac{(-ib(\tau))^k}{k!} \varepsilon^k.$$

In particular, one observes that one needs to choose the initial data as an expansion in powers ε up to the order ε^{2p-1} in order to ensure the boundedness of the time derivatives of u up to the order p . This crucial property will guide our analysis in the next subsection. Of course, for (2.4), we do not have an exact solution, but we will obtain iteratively the expansion in powers of ε of the suitable initial data by using Chapman-Enskog techniques.

2.2. Suitable initial data for the augmented problem (2.4)

In order to define the initial data $U^\varepsilon(0, \tau)$ where here and after we omit the space variable x for simplicity, let us perform *formally* the Chapman-Enskog expansion of U^ε [10, 11, 12]. To this aim, we introduce the operators L and Π for a periodic function $h(\tau) : \mathbb{T} \rightarrow \mathbb{C}^2$ as

$$Lh := \partial_\tau h, \quad \Pi h := \frac{1}{2\pi} \int_0^{2\pi} h(\tau) d\tau,$$

and when $\Pi h = 0$, the operator L is invertible with

$$(L^{-1}h)(\tau) = (I - \Pi) \int_0^\tau h(\theta) d\theta.$$

We perform the Chapman-Enskog expansion by setting

$$U^\varepsilon(t, \tau) = \underline{U}(t) + h(t, \tau), \quad \text{with } \underline{U} = \Pi U^\varepsilon.$$

We observe that $\Pi A = 0$ and write the micro-macro formulation of (2.4) as

$$\partial_t \underline{U} = -\frac{1}{\varepsilon} \Pi(A(\tau) \partial_x h) + \Pi(F(t, \tau, \underline{U} + h)), \quad (2.6)$$

$$\partial_t h = -\frac{1}{\varepsilon^2} Lh - \frac{1}{\varepsilon} A(\tau) \partial_x \underline{U} - \frac{1}{\varepsilon} (I - \Pi)(A(\tau) \partial_x h) + (I - \Pi)F(t, \tau, \underline{U} + h). \quad (2.7)$$

Based on this formulation, it appears that, when $\varepsilon \rightarrow 0$, $U^\varepsilon(t, \tau)$ converges to $U(t)$ satisfying

$$\partial_t U = C \partial_x^2 U + \Pi(F(t, \tau, U)) = F_e(t, U), \quad (2.8)$$

where the matrix C is given below by (2.12). In particular, the solution Φ^ε of the nonlinear Dirac equation (1.1) will formally be close to Φ defined by

$$\Phi = \begin{pmatrix} e^{-it/\varepsilon^2} & 0 \\ 0 & e^{it/\varepsilon^2} \end{pmatrix} U.$$

The above limit model is a system of coupled nonlinear Schrödinger equations. When there is no magnetic potential, i.e. $V_m = 0$, the convergence has been proved rigorously in some energy space in [20].

2.2.1. Uniform boundedness of time derivatives up to order 2 In this subsection, we derive the expression of $U^\varepsilon(0, \tau)$ such that the time derivatives of U^ε up to order 2 are uniformly bounded. Following the previous subsection on the toy model, we need an expansion up to the order ε^3 . In the following formal calculations, we thus assume that derivatives of h until order 2 are bounded.

By applying the inverse of L to (2.7), we get

$$h = -\varepsilon L^{-1} A \partial_x \underline{U} - \varepsilon L^{-1} (I - \Pi) A \partial_x h + \varepsilon^2 L^{-1} (I - \Pi) F(t, \tau, \underline{U} + h) - \varepsilon^2 L^{-1} \partial_t h. \quad (2.9)$$

Then we further have

$$\begin{aligned} \partial_x h &= -\varepsilon L^{-1} A \partial_x^2 \underline{U} - \varepsilon L^{-1} (I - \Pi) A \partial_x^2 h + \varepsilon^2 L^{-1} (I - \Pi) \partial_x F(t, \tau, \underline{U} + h) \\ &\quad - \varepsilon^2 L^{-1} \partial_{tx} h, \end{aligned} \quad (2.10)$$

$$\begin{aligned} \partial_t h &= -\varepsilon L^{-1} A \partial_{tx} \underline{U} - \varepsilon L^{-1} (I - \Pi) A \partial_{tx} h + \varepsilon^2 L^{-1} (I - \Pi) [\partial_t F(t, \tau, \underline{U} + h) \\ &\quad + \partial_u F(t, \tau, \underline{U} + h) (\partial_t \underline{U} + \partial_t h)] - \varepsilon^2 L^{-1} \partial_t^2 h. \end{aligned} \quad (2.11)$$

Hence, we can expand the unknowns in powers of ε and find

$$h = -\varepsilon L^{-1} A \partial_x \underline{U} + \varepsilon^2 L^{-1} (I - \Pi) A L^{-1} A \partial_x^2 \underline{U} + \varepsilon^2 L^{-1} (I - \Pi) F(t, \tau, \underline{U}) + \mathcal{O}(\varepsilon^3).$$

Therefore,

$$\begin{aligned} U^\varepsilon(0, \tau) &= \underline{U} - \varepsilon L^{-1} A \partial_x \underline{U} + \varepsilon^2 L^{-1} (I - \Pi) A L^{-1} A \partial_x^2 \underline{U} \\ &\quad + \varepsilon^2 L^{-1} (I - \Pi) F(t, \tau, \underline{U}) + \mathcal{O}(\varepsilon^3). \end{aligned}$$

Direct computations yield

$$B(\tau) := L^{-1} A(\tau) = -\frac{i}{2} \begin{pmatrix} 0 & e^{2i\tau} \\ -e^{-2i\tau} & 0 \end{pmatrix}$$

and

$$C \equiv C(\tau) := A(\tau) B(\tau) = \frac{i}{2} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \quad (2.12)$$

so $(I - \Pi)C(\tau) = 0$ and

$$U^\varepsilon(0, \tau) = \underline{U}(0) - \varepsilon B(\tau) \partial_x \underline{U}(0) + \varepsilon^2 L^{-1} (I - \Pi) F(0, \tau, \underline{U}) + \mathcal{O}(\varepsilon^3).$$

Now we use that $U^\varepsilon(0, 0) = \Phi_0$ to get

$$\begin{aligned} \underline{U}(0) &= \Phi_0 + \varepsilon B(0) \partial_x \underline{U}(0) - \varepsilon^2 f_0(0) + \mathcal{O}(\varepsilon^3) \\ &= \Phi_0 + \varepsilon B(0) \partial_x \Phi_0 + \varepsilon^2 B(0)^2 \partial_x^2 \Phi_0 - \varepsilon^2 f_0(0) + \mathcal{O}(\varepsilon^3), \end{aligned} \quad (2.13)$$

where

$$f_0(\tau) := L^{-1} (I - \Pi) F(0, \tau, \Phi_0) = -i V_m(0) B(\tau) \Phi_0,$$

and then

$$\begin{aligned} U^\varepsilon(0, \tau) &= \Phi_0 - \varepsilon (B(\tau) - B(0)) \partial_x \Phi_0 - \varepsilon^2 (B(\tau) - B(0)) B(0) \partial_x^2 \Phi_0 \\ &\quad + \varepsilon^2 (f_0(\tau) - f_0(0)) + \mathcal{O}(\varepsilon^3). \end{aligned}$$

We thus get an expression U_2^ε for the suitable initial data up to $\mathcal{O}(\varepsilon^3)$, i.e. $U^\varepsilon(0, \tau) = U_2^\varepsilon(\tau) + \mathcal{O}(\varepsilon^3)$, as follows:

$$\begin{aligned} U_2^\varepsilon(\tau) &:= \Phi_0 + \frac{i\varepsilon}{2} \begin{pmatrix} 0 & e^{2i\tau} - 1 \\ 1 - e^{-2i\tau} & 0 \end{pmatrix} \partial_x \Phi_0 + \frac{\varepsilon^2}{4} \begin{pmatrix} 1 - e^{2i\tau} & 0 \\ 0 & 1 - e^{-2i\tau} \end{pmatrix} \partial_x^2 \Phi_0 \\ &\quad - \frac{\varepsilon^2}{2} V_m(0) \begin{pmatrix} 0 & e^{2i\tau} - 1 \\ 1 - e^{-2i\tau} & 0 \end{pmatrix} \Phi_0. \end{aligned}$$

For future use, we denote $U_2(\tau) = U_2^\varepsilon(\tau) - \varepsilon^2 (f_0(\tau) - f_0(0))$, and also define

$$U_1^\varepsilon(\tau) := \Phi_0 + \frac{i\varepsilon}{2} \begin{pmatrix} 0 & e^{2i\tau} - 1 \\ 1 - e^{-2i\tau} & 0 \end{pmatrix},$$

which gives $U^\varepsilon(0, \tau) = U_1^\varepsilon(\tau) + \mathcal{O}(\varepsilon^2)$.

Now we push the asymptotic expansion to the next order. Inserting (2.10) and (2.11) into (2.9) and letting $t=0$, we get

$$\begin{aligned}
h(0, \tau) &= -\varepsilon B \partial_x \underline{U} - \varepsilon L^{-1} (I - \Pi) A [-\varepsilon B \partial_x^2 \underline{U} + \varepsilon^2 L^{-1} (I - \Pi) C \partial_x^3 \underline{U} \\
&\quad + \varepsilon^2 L^{-1} (I - \Pi) \partial_x F(0, \tau, \underline{U})] + \varepsilon^2 L^{-1} (I - \Pi) F(0, \tau, \underline{U} + h) \\
&\quad + \varepsilon^3 L^{-1} B \partial_{tx} \underline{U} + O(\varepsilon^4) \\
&= -\varepsilon B \partial_x \underline{U} + \varepsilon^2 L^{-1} (I - \Pi) F(0, \tau, U_1^\varepsilon) + \varepsilon^3 L^{-1} B \partial_{tx} \underline{U} \\
&\quad - \varepsilon^3 L^{-1} (I - \Pi) A L^{-1} (I - \Pi) \partial_x F(0, \tau, \underline{U}) + O(\varepsilon^4).
\end{aligned} \tag{2.14}$$

From (2.6) and (2.10), we have

$$\partial_t \underline{U} = C \partial_x^2 \underline{U} + \Pi F(t, \tau, \underline{U}) + O(\varepsilon). \tag{2.15}$$

Noting $F_e(t, \underline{U}) = \Pi F(t, \tau, \underline{U})$, we find

$$\begin{aligned}
h(0, \tau) &= -\varepsilon B \partial_x \underline{U} + \varepsilon^2 L^{-1} (I - \Pi) F(0, \tau, U_1^\varepsilon) + \varepsilon^3 L^{-1} B [C \partial_x^3 \underline{U} + \partial_x F_e(0, \underline{U})] \\
&\quad - \varepsilon^3 L^{-1} (I - \Pi) A L^{-1} (I - \Pi) \partial_x F(0, \tau, \underline{U}) + O(\varepsilon^4).
\end{aligned}$$

Using (2.13), we then get

$$\begin{aligned}
h(0, \tau) &= -\varepsilon B(\tau) [\partial_x \Phi_0 + \varepsilon B(0) \partial_x^2 \Phi_0 + \varepsilon^2 B^2(0) \partial_x^3 \Phi_0 - \varepsilon^2 \partial_x f_0(0)] + \varepsilon^2 f_1(\tau) \\
&\quad + \varepsilon^3 L^{-1} B [C \partial_x^3 \Phi_0 + \partial_x F_e(0, \Phi_0)] - \varepsilon^3 L^{-1} (I - \Pi) A \partial_x f_0 + O(\varepsilon^4),
\end{aligned} \tag{2.16}$$

where

$$f_1(\tau) := L^{-1} (I - \Pi) F(0, \tau, U_1^\varepsilon).$$

Based on

$$U^\varepsilon(0, \tau) = \underline{U}(0) + h(0, \tau) = \Phi_0 + h(0, \tau) - h(0, 0),$$

and remarking that $L^{-1} B = -A/4$ and $L^{-1} (I - \Pi) A \partial_x f_0 \equiv 0$, we can update the expansion of $\underline{U}(0)$ as

$$\begin{aligned}
\underline{U}(0) &= \Phi_0 + \varepsilon B(0) \partial_x \Phi_0 + \varepsilon^2 B^2(0) \partial_x^2 \Phi_0 - \varepsilon^2 f_1(0) + \varepsilon^3 B^3(0) \partial_x^3 \Phi_0 \\
&\quad - \varepsilon^3 B(0) \partial_x f_0(0) + \frac{\varepsilon^3}{4} A(0) [C \partial_x^3 \Phi_0 + \partial_x F_e(0, \Phi_0)] + O(\varepsilon^4),
\end{aligned} \tag{2.17}$$

and then get an expression U_3^ε of the prepared initial data up to $O(\varepsilon^4)$, i.e. $U^\varepsilon(0, \tau) = U_3^\varepsilon(\tau) + O(\varepsilon^4)$, as follows:

$$U_3^\varepsilon(\tau) = U_2(\tau) + \varepsilon^2 [f_1(\tau) - f_1(0)] - \varepsilon^3 [B(\tau) - B(0)] B^2(0) \partial_x^3 \Phi_0 \tag{2.18}$$

$$\begin{aligned}
&\quad - \frac{\varepsilon^3}{4} (A(\tau) - A(0)) [C \partial_x^3 \Phi_0 + \partial_x F_e(0, \Phi_0)] + \varepsilon^3 (B(\tau) - B(0)) \partial_x f_0(0) \\
&= U_2(\tau) + \varepsilon^2 [f_1(\tau) - f_1(0)] + \frac{i\varepsilon^3}{4} \begin{pmatrix} 0 & e^{2i\tau} - 1 \\ 1 - e^{-2i\tau} & 0 \end{pmatrix} \partial_x^3 \Phi_0 \\
&\quad + \frac{i\varepsilon^3}{4} \begin{pmatrix} 1 - e^{2i\tau} & 0 \\ 0 & 1 - e^{-2i\tau} \end{pmatrix} \partial_x (V_m(0) \Phi_0) \\
&\quad + \frac{\varepsilon^3}{4} \begin{pmatrix} 0 & 1 - e^{2i\tau} \\ 1 - e^{-2i\tau} & 0 \end{pmatrix} \partial_x F_e(0, \Phi_0).
\end{aligned} \tag{2.19}$$

2.2.2. Uniform boundedness of time derivatives up to order 3 Repeating the procedure above, we are going to get the next order expansion. Here we assume that derivatives of h until order 3 are bounded. From (2.11), we have

$$\partial_t h = -\varepsilon B \partial_{tx} \underline{U} + \varepsilon^2 L^{-1} (I - \Pi) [\partial_t F(t, \tau, \underline{U}) + \partial_u F(t, \tau, \underline{U}) \partial_t \underline{U}] + O(\varepsilon^3).$$

Plugging the above expansion and (2.14) into (2.9) and letting $t=0$, we get

$$\begin{aligned} h(0, \tau) &= -\varepsilon B \partial_x \underline{U} + \varepsilon L^{-1} (I - \Pi) A [\varepsilon B \partial_x^2 \underline{U} - \varepsilon^2 f_1(\tau) - \varepsilon^3 L^{-1} B \partial_{txx} \underline{U}] \\ &\quad + \varepsilon^2 L^{-1} (I - \Pi) F(0, \tau, U_2^\varepsilon) + \varepsilon^3 L^{-1} B \partial_{tx} \underline{U} - \varepsilon^4 f_t(\tau) + O(\varepsilon^5) \\ &= -\varepsilon B \partial_x \underline{U} + \varepsilon^2 L^{-1} (I - \Pi) F(0, \tau, U_2^\varepsilon) - \varepsilon^3 L^{-1} (I - \Pi) A f_1 \\ &\quad - \frac{\varepsilon^3}{4} A \partial_{tx} \underline{U} - \varepsilon^4 f_t(\tau) + O(\varepsilon^5), \end{aligned} \quad (2.20)$$

where

$$\begin{aligned} f_t(\tau) &:= L^{-2} (I - \Pi) [\partial_t F(0, \tau, \Phi_0) + \partial_u F(0, \tau, \Phi_0) (C \partial_x^2 \Phi_0 + F_e(0, \Phi_0))] \\ &= \frac{i}{4} A(\tau) [\partial_t V_m(0) \Phi_0 + V_m(0) (C \partial_x^2 \Phi_0 + F_e(0, \Phi_0))]. \end{aligned}$$

From (2.6), we find

$$\partial_t \underline{U}(0) = C \partial_x^2 \underline{U} + \Pi F(0, \tau, U_1^\varepsilon) - \varepsilon \Pi A \partial_x f_1 + O(\varepsilon^2).$$

Then (2.20) becomes

$$\begin{aligned} h(0, \tau) &= -\varepsilon B(\tau) \partial_x \underline{U} + \varepsilon^2 L^{-1} (I - \Pi) F(0, \tau, U_2^\varepsilon) - \varepsilon^3 L^{-1} (I - \Pi) A f_1 \\ &\quad - \frac{\varepsilon^3}{4} A(\tau) [C \partial_x^3 \underline{U} + \Pi \partial_x F(0, \tau, U_1^\varepsilon) - \varepsilon \Pi A \partial_x^2 f_1] - \varepsilon^4 f_t(\tau) + O(\varepsilon^5). \end{aligned} \quad (2.21)$$

Combining (2.21) with (2.17) and noting $A(0)C = 4B^3(0)$, we get

$$\begin{aligned} h(0, \tau) &= -\varepsilon B(\tau) [\partial_x \Phi_0 + \varepsilon B(0) \partial_x^2 \Phi_0 + \varepsilon^2 B^2(0) \partial_x^3 \Phi_0 - \varepsilon^2 \partial_x f_1(0) + 2\varepsilon^3 B^3(0) \partial_x^4 \Phi_0 \\ &\quad - \varepsilon^3 B(0) \partial_x^2 f_0(0) + \frac{\varepsilon^3}{4} A(0) \partial_x^2 F_e(0, \Phi_0)] + \varepsilon^2 L^{-1} (I - \Pi) F(0, \tau, U_2^\varepsilon) \\ &\quad - \varepsilon^3 L^{-1} (I - \Pi) A f_1 - \frac{\varepsilon^3}{4} A(\tau) C [\partial_x^3 \Phi_0 + \varepsilon B(0) \partial_x^4 \Phi_0] \\ &\quad - \frac{\varepsilon^3}{4} A(\tau) [\Pi \partial_x F(0, \tau, U_1^\varepsilon) - \varepsilon \Pi A \partial_x^2 f_1] - \varepsilon^4 f_t(\tau) + O(\varepsilon^5). \end{aligned} \quad (2.22)$$

Then we have

$$\begin{aligned} \underline{U}(0) &= \Phi_0 + \varepsilon B(0) \partial_x \Phi_0 + \varepsilon^2 B^2(0) \partial_x^2 \Phi_0 - \varepsilon^2 f_2^\varepsilon(0) + 2\varepsilon^3 B^3(0) \partial_x^3 \Phi_0 - \varepsilon^3 B(0) \partial_x f_1(0) \\ &\quad + \frac{\varepsilon^3}{4} A(0) \Pi \partial_x F(0, \tau, U_1^\varepsilon) + \varepsilon^3 g^\varepsilon(0) + 3\varepsilon^4 B^4(0) \partial_x^4 \Phi_0 - \varepsilon^4 B^2(0) \partial_x^2 f_0(0) \\ &\quad + \frac{\varepsilon^4}{4} B(0) A(0) \partial_x^2 F_e(0, \Phi_0) - \frac{\varepsilon^4}{4} A(0) \Pi A \partial_x^2 f_1 + \varepsilon^4 f_t(0) + O(\varepsilon^5), \end{aligned} \quad (2.23)$$

where

$$f_2(\tau) := L^{-1} (I - \Pi) F(0, \tau, U_2^\varepsilon), \quad g_1(\tau) := L^{-1} (I - \Pi) A f_1.$$

Thus, we get the prepared initial data $U^\varepsilon(0, \tau) = U_4^\varepsilon(\tau) + O(\varepsilon^5)$, with

$$\begin{aligned}
U_4^\varepsilon(\tau) &= U_2(\tau) + \varepsilon^2 [f_2(\tau) - f_2(0)] + \frac{i\varepsilon^3}{4} \begin{pmatrix} 0 & e^{2i\tau} - 1 \\ 1 - e^{-2i\tau} & 0 \end{pmatrix} \partial_x^3 \Phi_0 \\
&+ \frac{\varepsilon^3}{4} \begin{pmatrix} 0 & 1 - e^{2i\tau} \\ 1 - e^{-2i\tau} & 0 \end{pmatrix} \Pi \partial_x F(0, \tau, U_1^\varepsilon) - \varepsilon^3 [g_1(\tau) - g_1(0)] \\
&- \frac{i\varepsilon^3}{2} \begin{pmatrix} 0 & e^{2i\tau} - 1 \\ 1 - e^{-2i\tau} & 0 \end{pmatrix} \partial_x f_1(0) - \frac{3\varepsilon^4}{16} \begin{pmatrix} e^{2i\tau} - 1 & 0 \\ 0 & e^{-2i\tau} - 1 \end{pmatrix} \partial_x^4 \Phi_0 \\
&- \frac{\varepsilon^4}{8} \begin{pmatrix} 0 & e^{2i\tau} - 1 \\ 1 - e^{-2i\tau} & 0 \end{pmatrix} \partial_x^2 (V_m(0) \Phi_0) + \frac{\varepsilon^4}{4} \begin{pmatrix} 0 & e^{2i\tau} - 1 \\ e^{-2i\tau} - 1 & 0 \end{pmatrix} \Pi A \partial_x^2 f_1 \\
&- \frac{i\varepsilon^4}{8} \begin{pmatrix} 1 - e^{2i\tau} & 0 \\ 0 & e^{-2i\tau} - 1 \end{pmatrix} \partial_x^2 F_e(0, \Phi_0) - \varepsilon^4 (f_t(\tau) - f_t(0)).
\end{aligned} \tag{2.24}$$

To get the next order expansion, we recall that we assume $\partial_t h, \partial_t^2 h, \partial_t^3 h = \mathcal{O}(1)$ as $\varepsilon \rightarrow 0$, which indicates that

$$\begin{aligned}
\partial_t h &= -\varepsilon B \partial_{tx} \underline{U} + \varepsilon^2 L^{-1} (I - \Pi) \frac{d}{dt} F(t, \tau, \underline{U}(t) + h(t)) + \varepsilon^3 L^{-1} B \partial_{ttx} \underline{U} \\
&- \varepsilon^3 L^{-1} (I - \Pi) A L^{-1} (I - \Pi) \frac{d}{dt} \partial_x F(t, \tau, \underline{U}(t)) + O(\varepsilon^4).
\end{aligned}$$

Using

$$\begin{aligned}
\frac{d}{dt} F(t, \tau, \underline{U}(t) + h(t)) &= \partial_t F(t, \tau, \underline{U} + h) + \partial_u F(t, \tau, \underline{U} + h) (\partial_t \underline{U} + \partial_t h) \\
&= \partial_t F(t, \tau, \underline{U} - \varepsilon B \partial_x \underline{U}) + \partial_u F(t, \tau, \underline{U} - \varepsilon B \partial_x \underline{U}) (\partial_t \underline{U} - \varepsilon B \partial_{tx} \underline{U}) + O(\varepsilon^2)
\end{aligned}$$

and denoting

$$\begin{aligned}
H_0(\tau) &:= L^{-1} (I - \Pi) [\partial_t F(0, \tau, \Phi_0) + \partial_u F(0, \tau, \Phi_0) C \partial_x^2 \Phi_0 + \partial_u F(0, \tau, \Phi_0) F_e(0, \Phi_0)], \\
H_1(\tau) &:= \partial_t F(0, \tau, U_1^\varepsilon) + \partial_u F(0, \tau, U_1^\varepsilon) C (\partial_x^2 \Phi_0 + \varepsilon B(0) \partial_x^3 \Phi_0) \\
&+ \partial_u F(0, \tau, U_1^\varepsilon) [\Pi F(0, \tau, U_1^\varepsilon) - \varepsilon \Pi A \partial_x f_1] \\
&- \varepsilon \partial_u F(0, \tau, U_1^\varepsilon) B(\tau) [C \partial_x^3 \Phi_0 + \partial_x F_e(0, \Phi_0)],
\end{aligned}$$

we get

$$\begin{aligned}
\partial_t h(0, \tau) &= -\varepsilon B(\tau) \partial_{tx} \underline{U} + \varepsilon^2 L^{-1} (I - \Pi) H_1 + \varepsilon^3 L^{-1} B \partial_{ttx} \underline{U} \\
&- \varepsilon^3 L^{-1} (I - \Pi) A \partial_x H_0 + O(\varepsilon^4).
\end{aligned} \tag{2.25}$$

A detailed computation gives

$$H_0(\tau) = -iB(\tau) [\partial_t V_m(0) \Phi_0 + V_m(0) (C \partial_x^2 \Phi_0 + F_e(0, \Phi_0))].$$

Plugging (2.16) into (2.6) at $t=0$, we have

$$\begin{aligned}
\partial_t \underline{U}(0) &= C \partial_x^2 \Phi_0 + \Pi F(0, \tau, U_2^\varepsilon) + \varepsilon C B(0) \partial_x^3 \Phi_0 + \frac{\varepsilon^2}{2} C \partial_x^4 \Phi_0 - \varepsilon \Pi A \partial_x f_1 \\
&- \varepsilon^2 C \partial_x^2 f_0(0) + \frac{\varepsilon^2}{4} \partial_x^2 F_e(0, \Phi_0) + O(\varepsilon^3),
\end{aligned}$$

and

$$\partial_t^2 \underline{U}(0) = C^2 \partial_x^4 \Phi_0 + C \partial_x^2 F_e(0, \Phi_0) + Z_e + O(\varepsilon),$$

where

$$Z_e = \partial_u F_e(0, \Phi_0) (C \partial_x^2 \Phi_0 + F_e(0, \Phi_0)) + \partial_t F_e(0, \Phi_0).$$

Combining the above two identities with (2.25), and inserting them together with (2.23) and (2.22) into (2.9), we finally obtain the prepared initial data $U^\varepsilon(0, \tau) = U_5^\varepsilon(\tau) + O(\varepsilon^6)$, with

$$\begin{aligned} U_5^\varepsilon(\tau) &= U_2(\tau) + \varepsilon^2 (f_3(\tau) - f_3(0)) + \frac{i\varepsilon^3}{4} \begin{pmatrix} 0 & e^{2i\tau} - 1 \\ 1 - e^{-2i\tau} & 0 \end{pmatrix} \partial_x^3 \Phi_0 \\ &+ \frac{\varepsilon^3}{4} \begin{pmatrix} 0 & 1 - e^{2i\tau} \\ 1 - e^{-2i\tau} & 0 \end{pmatrix} \Pi \partial_x F(0, \tau, U_2^\varepsilon) - \varepsilon^3 (g_2(\tau) - g_2(0)) \\ &- \frac{i\varepsilon^3}{2} \begin{pmatrix} 0 & e^{2i\tau} - 1 \\ 1 - e^{-2i\tau} & 0 \end{pmatrix} \partial_x f_2(0) + \frac{\varepsilon^4}{4} \begin{pmatrix} e^{2i\tau} - 1 & 0 \\ 0 & e^{-2i\tau} - 1 \end{pmatrix} \partial_x^2 f_1(0) \\ &- \frac{i\varepsilon^4}{8} \begin{pmatrix} 1 - e^{2i\tau} & 0 \\ 0 & e^{-2i\tau} - 1 \end{pmatrix} \Pi \partial_x^2 F(0, \tau, U_1^\varepsilon) - \frac{3\varepsilon^4}{16} \begin{pmatrix} e^{2i\tau} - 1 & 0 \\ 0 & e^{-2i\tau} - 1 \end{pmatrix} \partial_x^4 \Phi_0 \\ &+ \frac{i\varepsilon^4}{2} \begin{pmatrix} 0 & e^{2i\tau} - 1 \\ 1 - e^{-2i\tau} & 0 \end{pmatrix} \partial_x g_1(0) + \varepsilon^4 (v(\tau) - v(0)) - \varepsilon^4 (w_1(\tau) - w_1(0)) \\ &+ \frac{\varepsilon^4}{4} \begin{pmatrix} 0 & e^{2i\tau} - 1 \\ e^{-2i\tau} - 1 & 0 \end{pmatrix} \Pi A \partial_x^2 f_1 + \frac{3i\varepsilon^5}{16} \begin{pmatrix} 0 & e^{2i\tau} - 1 \\ 1 - e^{-2i\tau} & 0 \end{pmatrix} \partial_x^5 \Phi_0 \\ &+ \varepsilon^5 (w_0(\tau) - w_0(0)) - \frac{3\varepsilon^5}{16} \begin{pmatrix} 0 & e^{2i\tau} - 1 \\ e^{-2i\tau} - 1 & 0 \end{pmatrix} \partial_x^3 F_e(0, \Phi_0) \\ &+ \frac{i\varepsilon^5}{8} \begin{pmatrix} 1 - e^{2i\tau} & 0 \\ 0 & e^{-2i\tau} - 1 \end{pmatrix} \Pi A \partial_x^3 f_1 - \frac{\varepsilon^5}{8} \begin{pmatrix} e^{2i\tau} - 1 & 0 \\ 0 & 1 - e^{-2i\tau} \end{pmatrix} \partial_x Z_m \\ &- \frac{i\varepsilon^5}{8} \begin{pmatrix} e^{2i\tau} - 1 & 0 \\ 0 & e^{-2i\tau} - 1 \end{pmatrix} \partial_x^3 (V_m(0) \Phi_0) - \frac{i\varepsilon^5}{8} \begin{pmatrix} 0 & e^{2i\tau} - 1 \\ 1 - e^{-2i\tau} & 0 \end{pmatrix} \partial_x Z_e \end{aligned} \quad (2.26)$$

where

$$\begin{aligned} f_3(\tau) &:= L^{-1}(I - \Pi)F(0, U_3^\varepsilon), & g_2(\tau) &:= L^{-1}(I - \Pi)A \partial_x f_2, \\ v(\tau) &:= L^{-1}(I - \Pi)A L^{-1}(I - \Pi)A \partial_x f_1, \\ Z_m &:= \partial_t V_m(0) \Phi_0 + V_m(0)(C \partial_x^2 \Phi_0 + F_e(0, \Phi_0)), \\ w_0(\tau) &:= L^{-2}(I - \Pi)A \partial_x H_0, & w_1(\tau) &:= L^{-2}(I - \Pi)H_1. \end{aligned}$$

The augmented problem (2.4), endowed with one of the initial data U_{2k-1}^ε ($k \in \{1, 2, 3\}$), provides a solution which has uniformly bounded time derivatives up to the order k , respectively. The proof of this property, which will not be done here, can be derived in an analogous way as in [11]. The detailed proof would be lengthy and tedious, so we omit it for brevity.

3. Numerical method From the analytical results in [2, 20], the solution of the Dirac equation (1.1) decays very fast at infinity for localized initial data. Thus, to give the numerical discretization, we truncate the Dirac equation (1.1) in the whole

space to a bounded interval (a, b) which is large enough, and impose periodic boundary conditions as done in [5, 6, 26, 27], i.e.

$$\begin{aligned} i\partial_t\Phi^\varepsilon &= -\frac{i}{\varepsilon}\alpha\partial_x\Phi^\varepsilon + \frac{1}{\varepsilon^2}\beta\Phi^\varepsilon + [V_e(t) + V_m(t)\alpha]\Phi^\varepsilon + \lambda(\beta\Phi^\varepsilon, \Phi^\varepsilon)\beta\Phi^\varepsilon, \quad x \in (a, b), \\ \Phi^\varepsilon(t, a) &= \Phi^\varepsilon(t, b), \quad \partial_x\Phi^\varepsilon(t, a) = \partial_x\Phi^\varepsilon(t, b), \quad t \geq 0; \quad \Phi^\varepsilon(0, x) = \Phi_0(x). \end{aligned} \quad (3.1)$$

3.1. Uniformly accurate scheme Based on the two-scale formulation, we give the numerical discretization to solve the nonlinear Dirac equation. Due to the truncation (3.1), the corresponding two-scale problem (2.4) reads

$$\partial_t U^\varepsilon + \frac{1}{\varepsilon^2}\partial_\tau U^\varepsilon = -\frac{1}{\varepsilon}A(\tau)\partial_x U^\varepsilon + F(t, \tau, U^\varepsilon), \quad t > 0, \quad \tau \in \mathbb{T}, \quad x \in (a, b), \quad (3.2a)$$

$$U^\varepsilon(t, \tau, x) = U^\varepsilon(t, \tau + 2\pi, x), \quad t \geq 0, \quad \tau \in \mathbb{T}, \quad x \in \overline{(a, b)}, \quad (3.2b)$$

$$U^\varepsilon(t, \tau, a) = U^\varepsilon(t, \tau, b), \quad t \geq 0, \quad \tau \in \mathbb{T}. \quad (3.2c)$$

Let $\Delta t > 0$ be the time step, and denote $t_n = n\Delta t$ for $n = 0, 1, \dots$. Let us choose mesh sizes $\Delta x = (b - a)/N$ and $\Delta\tau = 2\pi/N_\tau$ with N and N_τ two positive even integers. We also denote the grid points in x and τ by

$$x_j := a + j\Delta x, \quad j = 0, 1, \dots, N; \quad \tau_j = j\Delta\tau, \quad j = 0, 1, \dots, N_\tau.$$

For first order time discretization, inspired by the scheme in [10], we propose the following semi-implicit Euler method:

$$\frac{U^{n+1}(\tau, x) - U^n(\tau, x)}{\Delta t} + \frac{1}{\varepsilon^2}\partial_\tau U^{n+1}(\tau, x) = -\frac{1}{\varepsilon}A(\tau)\partial_x U^{n+1}(\tau, x) + F^n(\tau, x), \quad (3.3)$$

where we denote $U^n(\tau, x) \approx U^\varepsilon(t_n, \tau, x)$ and $F^n(\tau, x) \approx F(t_n, \tau, U^\varepsilon(t_n, \tau, x))$, for $n = 0, 1, \dots$

3.1.1. Uniform accuracy of the semi-implicit scheme (3.3) We recall that the initial data $U_3^\varepsilon(\tau)$ given by (2.19) has been obtained by Chapman-Enskog expansions, in order to ensure the uniform boundedness of U^ε , $\partial_t U^\varepsilon$ and $\partial_t^2 U^\varepsilon$ with respect to ε . Although we omit the rigorous proof of this property, we assume that this holds and use it to prove the uniform accuracy of the semi-discrete scheme (3.3) initialized with $U^0 = U_3^\varepsilon$. In fact, we have the following result.

THEOREM 3.1. *Assume that the solution $U^\varepsilon(t, \tau, x)$ of (2.4) has uniformly bounded (with respect to ε) derivatives for $t \in [0, T]$, up to order two. Let $U^n(\tau, x)$ be the solution of the numerical scheme (3.3), initialized with $U^0 = U^\varepsilon(0)$, and $\Delta t > 0$. Then, for $n\Delta t \leq T$, we have*

$$\|U^\varepsilon(n\Delta t) - U^n\|_{L_\tau^\infty L_x^2} \leq C\Delta t,$$

where C is a constant independent of n , Δt and ε .

Proof. Taking the Fourier transform of (3.3) with respect to the variable x , we get

$$\frac{\widehat{U}^{n+1} - \widehat{U}^n}{\Delta t} + \frac{1}{\varepsilon^2}\partial_\tau \widehat{U}^{n+1} = -\frac{1}{\varepsilon}A(\tau)i\mu \widehat{U}^{n+1} + \widehat{F}^n,$$

where μ is the Fourier variable associated to x . In terms of the following quantity

$$Q(V) = V + \frac{\Delta t}{\varepsilon}A(\tau)i\mu V + \frac{\Delta t}{\varepsilon^2}\partial_\tau V,$$

the scheme reads

$$Q\left(\widehat{U^{n+1}}\right) = \widehat{U^n} + \Delta t \widehat{F^n}. \quad (3.4)$$

Step 1: invertibility property of Q . We claim that Q is invertible on $L_\tau^\infty L_\mu^2$ and that

$$\|Q^{-1}(W)\|_{L_\tau^\infty L_\mu^2} \leq \|W\|_{L_\tau^\infty L_\mu^2}. \quad (3.5)$$

Indeed, consider the equation $Q(V) = W$ and set

$$Z(\tau) := \begin{pmatrix} e^{-i\tau} & 0 \\ 0 & e^{i\tau} \end{pmatrix} V.$$

We get the equation

$$Z + i\Omega Z + \frac{\Delta t}{\varepsilon^2} \partial_\tau Z = \begin{pmatrix} e^{-i\tau} & 0 \\ 0 & e^{i\tau} \end{pmatrix} W,$$

with

$$\Omega = \begin{pmatrix} \Delta t/\varepsilon^2 & \mu\Delta t/\varepsilon \\ \mu\Delta t/\varepsilon & -\Delta t/\varepsilon^2 \end{pmatrix}.$$

Then,

$$\frac{d}{d\tau} \left(\exp\left(\frac{\varepsilon^2}{\Delta t}(I + i\Omega)\tau\right) Z \right) = \frac{\Delta t}{\varepsilon^2} \exp\left(\frac{\varepsilon^2}{\Delta t}(I + i\Omega)\tau\right) \begin{pmatrix} e^{-i\tau} & 0 \\ 0 & e^{i\tau} \end{pmatrix} W.$$

We integrate this expression between τ and $\tau + 2\pi$ and use the periodicity of $Z(\tau)$ to obtain

$$\begin{aligned} & \left(e^{\frac{\tau+2\pi}{\Delta t}\varepsilon^2} e^{i\frac{\tau+2\pi}{\Delta t}\varepsilon^2\Omega} - e^{\frac{\tau}{\Delta t}\varepsilon^2} e^{i\frac{\tau}{\Delta t}\varepsilon^2\Omega} \right) Z(\tau) \\ &= \frac{\Delta t}{\varepsilon^2} \int_\tau^{\tau+2\pi} e^{\frac{s}{\Delta t}\varepsilon^2} e^{i\frac{s}{\Delta t}\varepsilon^2\Omega} \begin{pmatrix} e^{-is} & 0 \\ 0 & e^{is} \end{pmatrix} W(s) ds \end{aligned}$$

Using that Ω is real-valued and symmetric, we get, with the notation $\|Z\| = \sqrt{|Z_1|^2 + |Z_2|^2}$ for all $Z \in \mathbb{C}^2$, that

$$\begin{aligned} & \left\| e^{\frac{\tau+2\pi}{\Delta t}\varepsilon^2} e^{i\frac{\tau+2\pi}{\Delta t}\varepsilon^2\Omega} Z(\tau) - e^{\frac{\tau}{\Delta t}\varepsilon^2} e^{i\frac{\tau}{\Delta t}\varepsilon^2\Omega} Z(\tau) \right\| \\ & \geq e^{\frac{\tau+2\pi}{\Delta t}\varepsilon^2} \left\| e^{i\frac{\tau+2\pi}{\Delta t}\varepsilon^2\Omega} Z(\tau) \right\| - e^{\frac{\tau}{\Delta t}\varepsilon^2} \left\| e^{i\frac{\tau}{\Delta t}\varepsilon^2\Omega} Z(\tau) \right\| \\ & = \left(e^{\frac{\tau+2\pi}{\Delta t}\varepsilon^2} - e^{\frac{\tau}{\Delta t}\varepsilon^2} \right) \|Z(\tau)\|, \end{aligned}$$

and that

$$\begin{aligned} & \left\| \frac{\Delta t}{\varepsilon^2} \int_\tau^{\tau+2\pi} e^{\frac{s}{\Delta t}\varepsilon^2} e^{i\frac{s}{\Delta t}\varepsilon^2\Omega} \begin{pmatrix} e^{-is} & 0 \\ 0 & e^{is} \end{pmatrix} W(s) ds \right\|_{L^2_\mu} \\ & \leq \frac{\Delta t}{\varepsilon^2} \int_\tau^{\tau+2\pi} e^{\frac{s}{\Delta t}\varepsilon^2} \left\| \begin{pmatrix} e^{-is} & 0 \\ 0 & e^{is} \end{pmatrix} W(s) \right\|_{L^2_\mu} \\ & \leq \left(e^{\frac{\tau+2\pi}{\Delta t}\varepsilon^2} - e^{\frac{\tau}{\Delta t}\varepsilon^2} \right) \|W\|_{L_\tau^\infty L_\mu^2} \end{aligned}$$

These two inequalities give $\|Z(\tau)\|_{L_\mu^2} \leq \|W\|_{L_\tau^\infty L_\mu^2}$, which proves the claim.

Step 2: Taylor expansion of the exact solution. The exact solution of (2.4) satisfies

$$U(t^n) - U(t^{n+1}) + \Delta t \partial_t U(t^{n+1}) = R^n \quad \text{with} \quad R^n = \int_{t^n}^{t^{n+1}} (s - t^n) \partial_t^2 U(s) ds$$

which yields

$$Q(\widehat{U}(t^{n+1})) = \widehat{U}(t^n) + \Delta t \widehat{F}(t^{n+1}, \tau, \widehat{U}(t^{n+1})) - \widehat{R}^n.$$

Here we have denoted $\widehat{G}(\widehat{U}) = \widehat{G}(U)$.

We now introduce the error $E^n = \widehat{U}(t^n) - \widehat{U}^n$ and subtracting (3.4) from this equation, we get

$$Q(E^{n+1}) = E^n + \Delta t \left(\widehat{F}(t^{n+1}, \tau, \widehat{U}(t^{n+1})) - \widehat{F}(t^n, \tau, \widehat{U}^n) \right) - \widehat{R}^n. \quad (3.6)$$

Step 3: error estimate. We observe from the expression (2.3) of F and the assumptions on V_e and V_m that F is locally Lipschitz continuous with respect to t and U . Let us fix $M > 0$ such that $\|U\|_{L_t^\infty L_\tau^\infty L_x^2} \leq M$ for $t \in [0, T]$. Let N be the number of time discretization points and $\Delta t = T/N$. Let n_0 be the largest integer $n \leq N$ such that $\|U^n\|_{L_\tau^\infty L_x^2} \leq 2M$. Our aim is to prove that $n_0 = N$ and to estimate the error.

Let us proceed by contradiction and assume that $n_0 \leq N - 1$. Then we have

$$\left\| \widehat{F}(t^{n+1}, \tau, \widehat{U}(t^{n+1})) - \widehat{F}(t^n, \tau, \widehat{U}^n) \right\|_{L_\tau^\infty L_\mu^2} \leq C \left(\Delta t + \|E^n\|_{L_\tau^\infty L_\mu^2} \right),$$

where C , here and after, is a generic constant which only depends on M and T . Applying (3.4) to (3.6), we deduce that

$$\|E^{n+1}\|_{L_\tau^\infty L_\mu^2} \leq (1 + C\Delta t) \|E^n\|_{L_\tau^\infty L_\mu^2} + C\Delta t^2 + \|R^n\|_{L_\tau^\infty L_\mu^2}.$$

Now, in order to estimate R^n , we apply the assumption made in Theorem 3.1 saying that $\partial_t^2 U$ is uniformly bounded with respect to ε . Thus $\|R^n\|_{L_\tau^\infty L_\mu^2} \leq C\Delta t^2$. Therefore

$$\|E^{n+1}\|_{L_\tau^\infty L_\mu^2} \leq (1 + C\Delta t) \|E^n\|_{L_\tau^\infty L_\mu^2} + C\Delta t^2,$$

for all $n \leq n_0$, which implies

$$\|E^{n+1}\|_{L_\tau^\infty L_\mu^2} \leq C\Delta t.$$

In particular, if $n_0 \Delta t < T$ and if we choose $C\Delta t < M$ then we deduce from this estimate that

$$\|\widehat{U}^{n_0+1}\|_{L_\tau^\infty L_\mu^2} \leq \|\widehat{U}(t^{n_0+1})\|_{L_\tau^\infty L_\mu^2} + \|E^{n_0+1}\|_{L_\tau^\infty L_\mu^2} \leq 2M.$$

This contradicts the fact that n_0 is the largest integer $n \leq N$ such that $\|U^n\|_{L_\tau^\infty L_x^2} \leq 2M$. We conclude that $n_0 = N$. It is also clear from the above estimates that we have

$$\|E^n\|_{L_\tau^\infty L_\mu^2} \leq C\Delta t,$$

for all $n \leq N$, which concludes the proof. \square

3.1.2. The fully discretized scheme In the x -direction, we apply the Fourier transform and obtain for $l = -N/2, \dots, N/2 - 1$,

$$\frac{\widehat{U}_l^{n+1}(\tau) - \widehat{U}_l^n(\tau)}{\Delta t} + \frac{1}{\varepsilon^2} \partial_\tau \widehat{U}_l^{n+1}(\tau) = -\frac{1}{\varepsilon} A(\tau) i\mu_l \widehat{U}_l^{n+1}(\tau) + \widehat{F}_l^n(\tau), \quad (3.7)$$

where $\mu_l = \frac{2\pi l}{b-a}$ and

$$U^n(\tau, x) = \sum_{l=-N/2}^{N/2-1} \widehat{U}_l^n(\tau) e^{i\mu_l(x-a)}, \quad F^n(\tau, x) = \sum_{l=-N/2}^{N/2-1} \widehat{F}_l^n(\tau) e^{i\mu_l(x-a)},$$

with

$$\widehat{U}_l^n(\tau) = \frac{1}{N} \sum_{j=0}^{N-1} U^n(\tau, x_j) e^{i\mu_l(x_j-a)}, \quad \widehat{F}_l^n(\tau) = \frac{1}{N} \sum_{j=0}^{N-1} F^n(\tau, x_j) e^{i\mu_l(x_j-a)}.$$

In terms of the two components $U^\varepsilon = (u_1, u_2)^T$ and $F = (f_1, f_2)^T$, the numerical scheme (3.7) reads

$$\frac{(\widehat{u}_1^{n+1})_l(\tau) - (\widehat{u}_1^n)_l(\tau)}{\Delta t} + \frac{1}{\varepsilon^2} \partial_\tau (\widehat{u}_1^{n+1})_l(\tau) = -\frac{i\mu_l}{\varepsilon} e^{2i\tau} (\widehat{u}_2^{n+1})_l(\tau) + (\widehat{f}_1^n)_l(\tau), \quad (3.8a)$$

$$\frac{(\widehat{u}_2^{n+1})_l(\tau) - (\widehat{u}_2^n)_l(\tau)}{\Delta t} + \frac{1}{\varepsilon^2} \partial_\tau (\widehat{u}_2^{n+1})_l(\tau) = -\frac{i\mu_l}{\varepsilon} e^{-2i\tau} (\widehat{u}_1^{n+1})_l(\tau) + (\widehat{f}_2^n)_l(\tau). \quad (3.8b)$$

Then we discretize the τ -direction as follows

$$\begin{aligned} (\widehat{u}_j^n)_l &= \left((\widehat{u}_j^n)_l(\tau_0), \dots, (\widehat{u}_j^n)_l(\tau_{N_\tau-1}) \right)^T, \quad l = -N/2, \dots, N/2 - 1, \\ (\widehat{f}_j^n)_l &= \left((\widehat{f}_j^n)_l(\tau_0), \dots, (\widehat{f}_j^n)_l(\tau_{N_\tau-1}) \right)^T, \quad j = 1, 2. \end{aligned}$$

Introducing the Fourier pseudo-differential matrix D_τ [24]:

$$D_\tau = (d_{n,m}) \in \mathbb{C}^{N_\tau \times N_\tau}, \quad d_{n,m} = \frac{i}{N_\tau} \sum_{l=-N_\tau/2}^{N_\tau/2-1} l e^{il(\tau_n - \tau_m)}, \quad n, m = 1, \dots, N_\tau,$$

we derive the full discretization of (3.8) as

$$\frac{(\widehat{u}_1^{n+1})_l - (\widehat{u}_1^n)_l}{\Delta t} + \frac{1}{\varepsilon^2} D_\tau (\widehat{u}_1^{n+1})_l = -\frac{i\mu_l}{\varepsilon} e^{2i\tau} (\widehat{u}_2^{n+1})_l + (\widehat{f}_1^n)_l, \quad (3.9a)$$

$$\frac{(\widehat{u}_2^{n+1})_l - (\widehat{u}_2^n)_l}{\Delta t} + \frac{1}{\varepsilon^2} D_\tau (\widehat{u}_2^{n+1})_l = -\frac{i\mu_l}{\varepsilon} e^{-2i\tau} (\widehat{u}_1^{n+1})_l + (\widehat{f}_2^n)_l, \quad n \geq 0, \quad (3.9b)$$

for $l = -\frac{N}{2}, \dots, \frac{N}{2} - 1$, where here and after $e^{\pm 2i\tau}$ are interpreted as diagonal matrices. Solving (3.9), we get the detailed scheme of the first order method (UA1) for $n = 0, 1, \dots$, and $l = -N/2, \dots, N/2 - 1$,

$$A_\tau (\widehat{u}_1^{n+1})_0 = \frac{1}{\Delta t} (\widehat{u}_1^n)_0 + (\widehat{f}_1^n)_0, \quad A_\tau (\widehat{u}_2^{n+1})_0 = \frac{1}{\Delta t} (\widehat{u}_2^n)_0 + (\widehat{f}_2^n)_0, \quad (3.10a)$$

$$B_\tau^l (\widehat{u}_1^{n+1})_l = A_\tau^l \left[\frac{1}{\Delta t} (\widehat{u}_1^n)_l + (\widehat{f}_1^n)_l \right] - \frac{1}{\Delta t} (\widehat{u}_2^n)_l - (\widehat{f}_2^n)_l, \quad l \neq 0, \quad (3.10b)$$

$$\widehat{u}_2^{n+1})_l = \frac{\varepsilon e^{-2i\tau}}{i\mu_l} \left[-A_\tau (\widehat{u}_1^{n+1})_l + \frac{1}{\Delta t} (\widehat{u}_1^n)_l + (\widehat{f}_1^n)_l \right], \quad l \neq 0, \quad (3.10c)$$

where

$$A_\tau = \frac{Id}{\Delta t} + \frac{D_\tau}{\varepsilon^2}, \quad A_\tau^l = \frac{\varepsilon A_\tau e^{-2i\tau}}{i\mu_l}, \quad B_\tau^l = A_\tau^l A_\tau - \frac{i\mu_l e^{-2i\tau}}{\varepsilon}, \quad l \neq 0,$$

and where Id denotes the identity matrix. The inverse of the above matrices A_τ, B_τ^l can be computed numerically once for all, so the UA1 method (3.10) is explicit in practise. Also noticing that $B_\tau^l = -B_\tau^{-l}$, we only need to store and compute the inverse of B_τ^l for $l = -N/2, \dots, -1$.

3.1.3. The uniformly accurate second order scheme Again inspired by [10], our second order method starts with a prediction step from the first order method

$$\frac{U^{n+1/2}(\tau, x) - U^n(\tau, x)}{\Delta t} + \frac{1}{\varepsilon^2} \partial_\tau U^{n+1/2}(\tau, x) = -\frac{1}{\varepsilon} A(\tau) \partial_x U^{n+1/2}(\tau, x) + F^n(\tau, x),$$

followed by the correction step

$$\begin{aligned} & \frac{U^{n+1}(\tau, x) - U^n(\tau, x)}{2\Delta t} + \frac{1}{2\varepsilon^2} [\partial_\tau U^{n+1}(\tau, x) + \partial_\tau U^n(\tau, x)] \\ & = -\frac{1}{2\varepsilon} A(\tau) [\partial_x U^{n+1}(\tau, x) + \partial_x U^n(\tau, x)] + F^{n+1/2}(\tau, x), \end{aligned}$$

where $F^{n+1/2}(\tau, x) := F(t_{n+1/2}, \tau, U^{n+1/2}(\tau, x))$. Then by applying the Fourier pseudo-spectral discretization similarly as before, we end up with the following detailed second order scheme (UA2). For $n=0, 1, \dots$, and $l = -N/2, \dots, N/2 - 1$,

$$A_\tau^+(\widehat{u_1^{n+1}})_0 = A_\tau^-(\widehat{u_1^n})_0 + (\widehat{f_1^{n+1/2}})_0, \quad A_\tau^+(\widehat{u_2^{n+1}})_0 = A_\tau^-(\widehat{u_2^n})_0 + (\widehat{f_2^{n+1/2}})_0, \quad (3.11a)$$

$$\begin{aligned} B_\tau^{+,l}(\widehat{u_1^{n+1}})_l &= A_\tau^{+,l} \left[A_\tau^-(\widehat{u_1^n})_l + (\widehat{f_1^{n+1/2}})_l \right] - A_\tau^-(\widehat{u_2^n})_l - (\widehat{f_2^{n+1/2}})_l \\ &+ \frac{i\mu_l e^{-2i\tau}}{2\varepsilon} (\widehat{u_1^n})_l - A_\tau^+(\widehat{u_2^n})_l, \quad l \neq 0, \end{aligned} \quad (3.11b)$$

$$(\widehat{u_2^{n+1}})_l = -(\widehat{u_2^n})_l + \frac{2\varepsilon e^{-2i\tau}}{i\mu_l} \left[A_\tau^-(\widehat{u_1^n})_l - A_\tau^+(\widehat{u_1^{n+1}})_l + (\widehat{f_1^{n+1/2}})_l \right], \quad l \neq 0, \quad (3.11c)$$

with

$$A_\tau^{1/2}(\widehat{u_1^{n+1/2}})_0 = \frac{1}{2\Delta t} (\widehat{u_1^n})_0 + (\widehat{f_1^n})_0, \quad A_\tau^{1/2}(\widehat{u_2^{n+1/2}})_0 = \frac{1}{2\Delta t} (\widehat{u_2^n})_0 + (\widehat{f_2^n})_0,$$

$$B_\tau^{1/2,l}(\widehat{u_1^{n+1/2}})_l = A_\tau^{1/2,l} \left[\frac{1}{2\Delta t} (\widehat{u_1^n})_l + (\widehat{f_1^n})_l \right] - \frac{1}{2\Delta t} (\widehat{u_2^n})_l - (\widehat{f_2^n})_l, \quad l \neq 0,$$

$$(\widehat{u_2^{n+1/2}})_l = \frac{\varepsilon e^{-2i\tau}}{i\mu_l} \left[-A_\tau^{1/2}(\widehat{u_1^{n+1/2}})_l + \frac{1}{2\Delta t} (\widehat{u_1^n})_l + (\widehat{f_1^n})_l \right], \quad l \neq 0,$$

where

$$A_\tau^\pm = \frac{Id}{\Delta t} \pm \frac{D_\tau}{2\varepsilon^2}, \quad A_\tau^{+,l} = \frac{2\varepsilon A_\tau^+ e^{-2i\tau}}{i\mu_l}, \quad B_\tau^{+,l} = A_\tau^{+,l} A_\tau^+ - \frac{i\mu_l e^{-2i\tau}}{2\varepsilon}, \quad l \neq 0,$$

$$A_\tau^{1/2} = \frac{Id}{2\Delta t} + \frac{D_\tau}{\varepsilon^2}, \quad A_\tau^{1/2,l} = \frac{\varepsilon A_\tau^{1/2} e^{-2i\tau}}{i\mu_l}, \quad B_\tau^{1/2,l} = A_\tau^{1/2,l} A_\tau^{1/2} - \frac{i\mu_l e^{-2i\tau}}{\varepsilon}.$$

Similarly, we remark that the inverse of matrices $A_\tau^+, A_\tau^{1/2}, B_\tau^{+,l}, B_\tau^{1/2,l}$ can be computed once for all, and by noticing that $B_\tau^{1/2,l} = -B_\tau^{1/2,-l}$ and $B_\tau^{+,l} = -B_\tau^{+,-l}$, we can only store the inverse of $B_\tau^{+,l}$ and $B_\tau^{1/2,l}$ for $l = -N/2, \dots, -1$.

Thanks to the semi-implicit approximations in time, the time step Δt of UA1 and UA2 methods are free from any CFL-type conditions on $\Delta\tau$, Δx or stability condition on ε . After obtaining $U^n(\tau, x)$, then we take $\tau = t_n/\varepsilon^2$ to get the approximation of $\Phi^\varepsilon(t_n, x)$, i.e. $\Phi^\varepsilon(t_n, x) \approx U^n(t_n/\varepsilon^2, x)$.

4. Numerical results For comparison purpose, we define the initial data

$$U_0^\varepsilon := \Phi_0,$$

along with $U_1^\varepsilon, U_2^\varepsilon, U_3^\varepsilon, U_4^\varepsilon$ and U_5^ε . We shall test the error of the proposed UA1 (3.10) and UA2 (3.11) under those choices of initial data by performing the following three numerical experiments. The computational domain is chosen as $(a, b) = (-8, 8)$. The ‘exact’ solution is obtained numerically by UA2 with very small step sizes, e.g. $\Delta t = 10^{-6}, \Delta x = 1/64, N_\tau = 64$. We shall focus on the time discretization error.

Example I: (nonlinear without magnetic potential) We take the potential and initial data in (1.1) as

$$V_e = \frac{1-x}{2+2x^2}, \quad V_m = 0, \quad \phi_1^\varepsilon(t=0) = \frac{e^{-x^2}}{\sqrt{2}}, \quad \phi_2^\varepsilon(t=0) = e^{-\sqrt{2}x^2},$$

and choose $\lambda = 0.5$. We solve the Dirac equation (3.2) till $t = 0.5$, and show the error of the numerical solution (ϕ_1^n, ϕ_2^n) with $n = t/\Delta t$ as

$$\|\phi_1^\varepsilon(t) - \phi_1^n\|_{l^\infty} + \|\phi_2(t) - \phi_2^n\|_{l^\infty}.$$

The time discretization errors of UA1 and UA2 under different ε are shown in Fig. 4.1 and Fig. 4.2, respectively, where we take the spatial mesh size of the numerical method small enough, e.g. $\Delta x = 1/64, N_\tau = 32$.

Example II: (linear with magnetic potential) We use the same setup in Example I but without nonlinearity and with a non-zero magnetic potential, i.e.

$$V_m = \frac{(x+1)^2}{1+x^2}, \quad \lambda = 0.$$

The temporal errors of UA2 at $t = 0.5$ under different ε are shown in Fig. 4.3.

Example III: (nonlinear with magnetic potential) We keep both the nonlinearity and magnetic potential, i.e.

$$V_m = \frac{(x+1)^2}{1+x^2}, \quad \lambda = 0.5.$$

The time discretization error of UA2 at $t = 0.5$, with initial data U_5^ε , is shown in Fig. 4.4. The corresponding spatial errors of UA2 with respect to Δx and $\Delta\tau$ are shown in Fig. 4.5. The behavior of the spatial errors of UA1 and of the other two numerical examples are similar, so the results are omitted here for brevity.

The error between the solution of the nonlinear Dirac equation (1.1) and the solution of the limit model (2.8), i.e.

$$\|\phi_1^\varepsilon(t) - \phi_1(t)\|_{L^\infty} + \|\phi_2^\varepsilon(t) - \phi_2(t)\|_{L^\infty},$$

are shown in Fig. 4.6 at $t=0.5$.

Based on the numerical results, we can have the following observation:

(i) The UA1 scheme with initial data U_3^ε shows uniformly first order temporal accuracy and the UA2 scheme with U_5^ε shows uniformly second order temporal accuracy in all the three cases. In space discretization, the UA schemes have uniformly spectral accuracy in both x and the artificial τ .

(ii) The convergence rate from the nonlinear Dirac equation to the limit model is of order $O(\varepsilon)$, as $\varepsilon \rightarrow 0$.

5. Conclusion We proposed some uniformly accurately (UA) schemes for solving the nonlinear Dirac equation in the nonrelativistic limit regimes. Our approach is based on a suitable two-scale formulation which offers a general strategy for constructing UA schemes for a class of highly oscillatory problems involving two small scales. We derive correct initial data for this augmented formulation, using a Chapman-Enskog expansion. This allows us to construct a UA scheme with second order accuracy in time and spectral accuracy in space. Numerical tests were done to show the UA property. Our approach can also be applied to solve the oscillatory kinetic equations with diffusion scaling. This the subject of a work in progress.

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REFERENCES

- [1] M. Ablowitz, Y. Zhu, *Nonlinear waves in shallow honeycomb lattices*, SIAM J. Appl. Math., 72, 240–260, 2012.
- [2] P. Bechouche, N. Mauser, F. Poupaud, *Semi-nonrelativistic limits of the Dirac equation with external time-dependent electromagnetic field*, Commun. Math. Phys., 197, 405–425, 1998.
- [3] Th. Bartsch, Y. Ding, *Solutions of nonlinear Dirac equations*, J. Differential Equations, 226, 210–249 (2006).
- [4] N. Bournaveas, G.E. Zouraris, *Theory and numerical approximations for a nonlinear 1+1 Dirac system*, ESAIM: M2AN, 46, 841–874, 2012.
- [5] W. Bao, Y. Cai, X. Jia, Q. Tang, *Numerical methods and comparison for the Dirac equation in the nonrelativistic limit regime*, preprint, 2015.
- [6] W. Bao, Y. Cai, X. Jia, Q. Tang, *A uniformly accurate (UA) multiscale time integrator pseudospectral method for the Dirac equation in the nonrelativistic limit regime*, preprint, arXiv:1507.04103v1 [math.NA], 2015.
- [7] W. Bao, Y. Cai, X. Zhao, *A uniformly accurate multiscale time integrator pseudospectral method for the Klein-Gordon equation in the nonrelativistic limit regime*, SIAM J. Numer. Anal., 52, 2488–2511, 2014.
- [8] W. Bao, X. Dong, X. Zhao, *Uniformly accurate multiscale time integrators for highly oscillatory second order differential equations*, J. Math. Study, 47, 111–150, 2014.
- [9] W. Bao, X. Zhao, *A numerical study and comparison of numerical methods for the nonlinear Klein-Gordon equation in the nonrelativistic limit regime*, preprint, 2016.
- [10] P. Chartier, N. Crouseilles, M. Lemou, F. Méhats, *Uniformly accurate numerical schemes for highly oscillatory Klein-Gordon and nonlinear Schrödinger equations*, Numer. Math. 129, 211–250, 2015.
- [11] N. Crouseilles, M. Lemou, F. Méhats, *Asymptotic Preserving schemes for highly oscillatory Vlasov-Poisson equations*, J. Comp. Phys., 248, 287–308, 2013.
- [12] N. Crouseilles, M. Lemou, F. Méhats, X. Zhao, *Uniformly accurate forward semi-Lagrangian methods for highly oscillatory Vlasov-Poisson equations*, preprint, hal-01286947, 2016.
- [13] C.L. Fefferman, M.I. Weinstein, *Honeycomb lattice potentials and Dirac points*, J. Amer. Math. Soc., 25, 1169–1220, 2012.
- [14] E. Frénod, S.A. Hirstoaga, M. Lutz, E. Sonnendrücker, *Long time behaviour of an exponential integrator for a Vlasov-Poisson system with strong magnetic field*, Communications in Comput. Phys., 18, 263–296, 2015.

- [15] D.J. Gross, A. Neveu, *Dynamical symmetry breaking in asymptotically free field theories*, Phys. Rev. D, 10, 3235–3253, 1974.
- [16] Z. Huang, S. Jin, P. Markowich, Ch. Sparber, C. Zheng, *A time-splitting spectral scheme for the Maxwell-Dirac system*, J. Comput. Phys., 208, 761–789, 2005.
- [17] L.H. Haddad, L.D. Carr, *The nonlinear Dirac equation in Bose-Einstein condensates: superfluid fluctuations and emergent theories from relativistic linear stability equations*, New J. Phys., 17, 093037, 2015.
- [18] H. Huh, *Global strong solution to the Thirring model in critical space*, J. Math. Anal. Appl., 381, 513–520, 2011.
- [19] T.W.B. Kibble, *Lorentz invariance and the gravitational field*, J. Math. Phys., 2, 212–221, 1961.
- [20] B. Najman, *The nonrelativistic limit of the nonlinear Dirac equation*, Ann. Inst. Henri Poincaré, 9, 3–12, 1992.
- [21] H. Pecher, *Local well-posedness for the nonlinear Dirac equation in two space dimensions*, Commun. Pure Appl. Anal., 13, 673–685, 2014.
- [22] M. Soler, *Classical, stable, nonlinear spinor field with positive rest energy*, Phys. Rev. D, 1, 2766–2769, 1970.
- [23] S. Shao, H. Tang, *Interaction of solitary waves with a phase shift in a nonlinear Dirac model*, Commun. Comput. Phys., 3, 950–967, 2008.
- [24] J. Shen, T. Tang, L. Wang, *Spectral Methods: algorithms, analysis and applications*, Springer-Verlag, Berlin Heidelberg, 2011.
- [25] W. Thirring, *A soluble relativistic field theory*, Ann. Physics 3 (1958), pp. 91–112.
- [26] J. Xu, S. Shao, H. Tang, *Numerical methods for nonlinear Dirac equation*, J. Comput. Phys., 245, 131–149, 2013.
- [27] H. Wang, H. Tang, *An efficient adaptive mesh redistribution method for a nonlinear Dirac equation*, J. Comp. Phys., 222, 176–193, 2007.

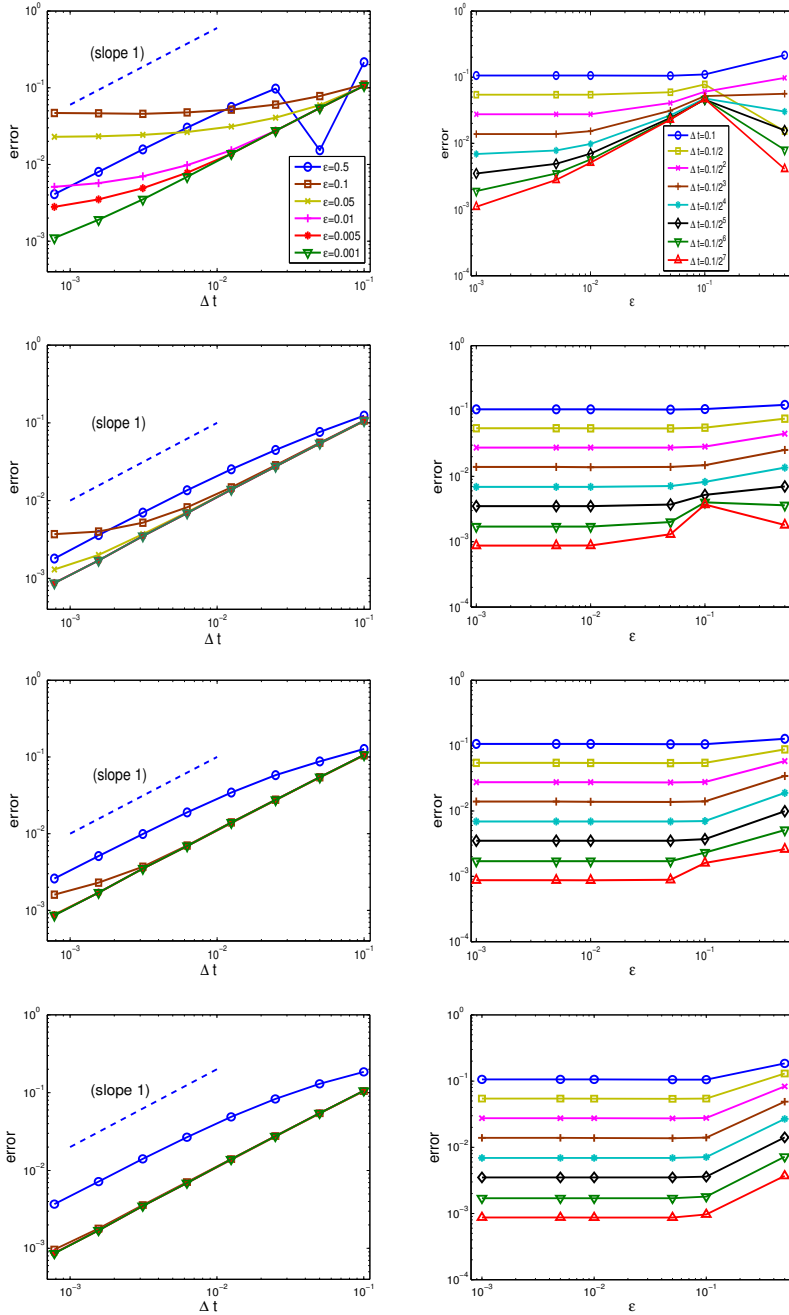


Fig. 4.1: Temporal error of UA1 in Example I with respect to Δt and ε : results of using U_0^ε (first row); results of using U_1^ε (second row) results of using U_2^ε (third row); results of using U_3^ε (last row).

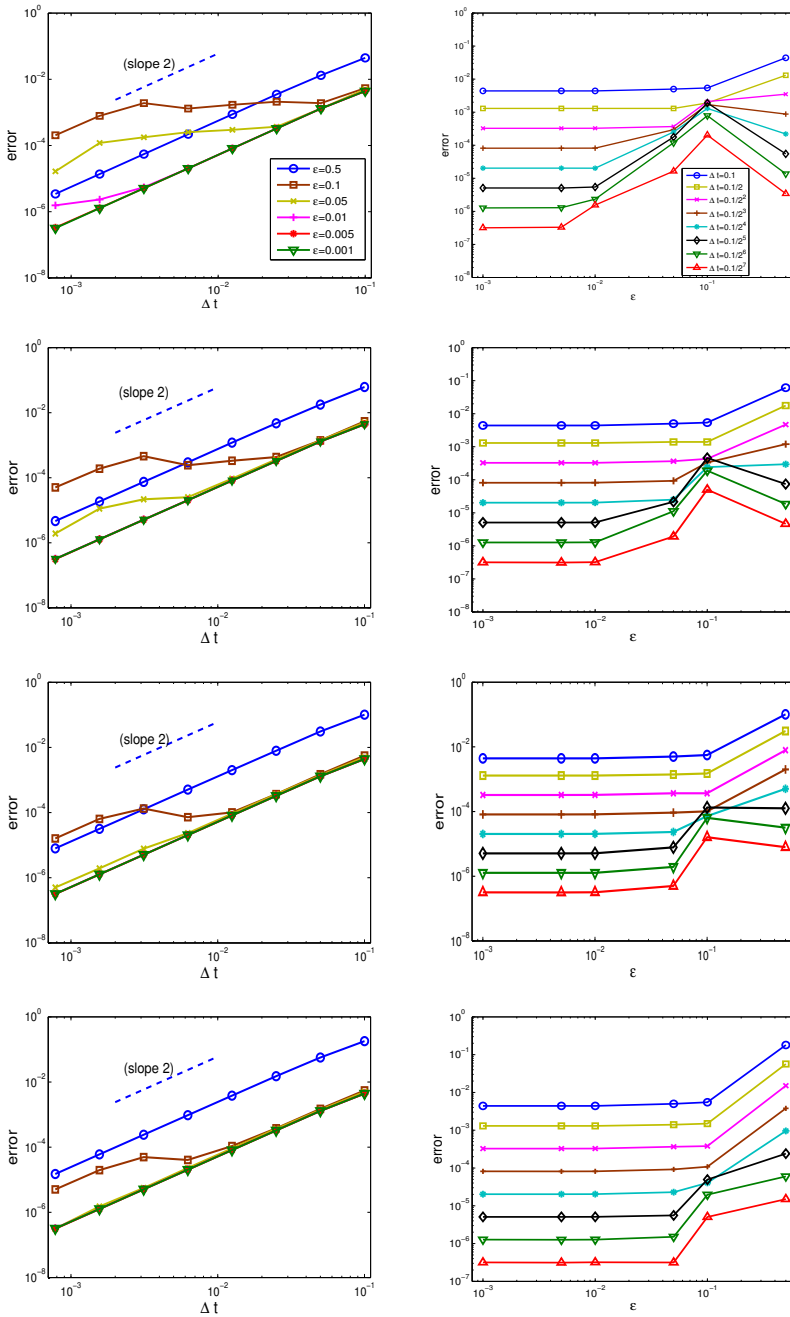


Fig. 4.2: Temporal error of UA2 in Example I with respect to Δt and ε : results of using U_2^ε (first row); results of using U_3^ε (second row) results of using U_4^ε (third row); results of using U_5^ε (last row).

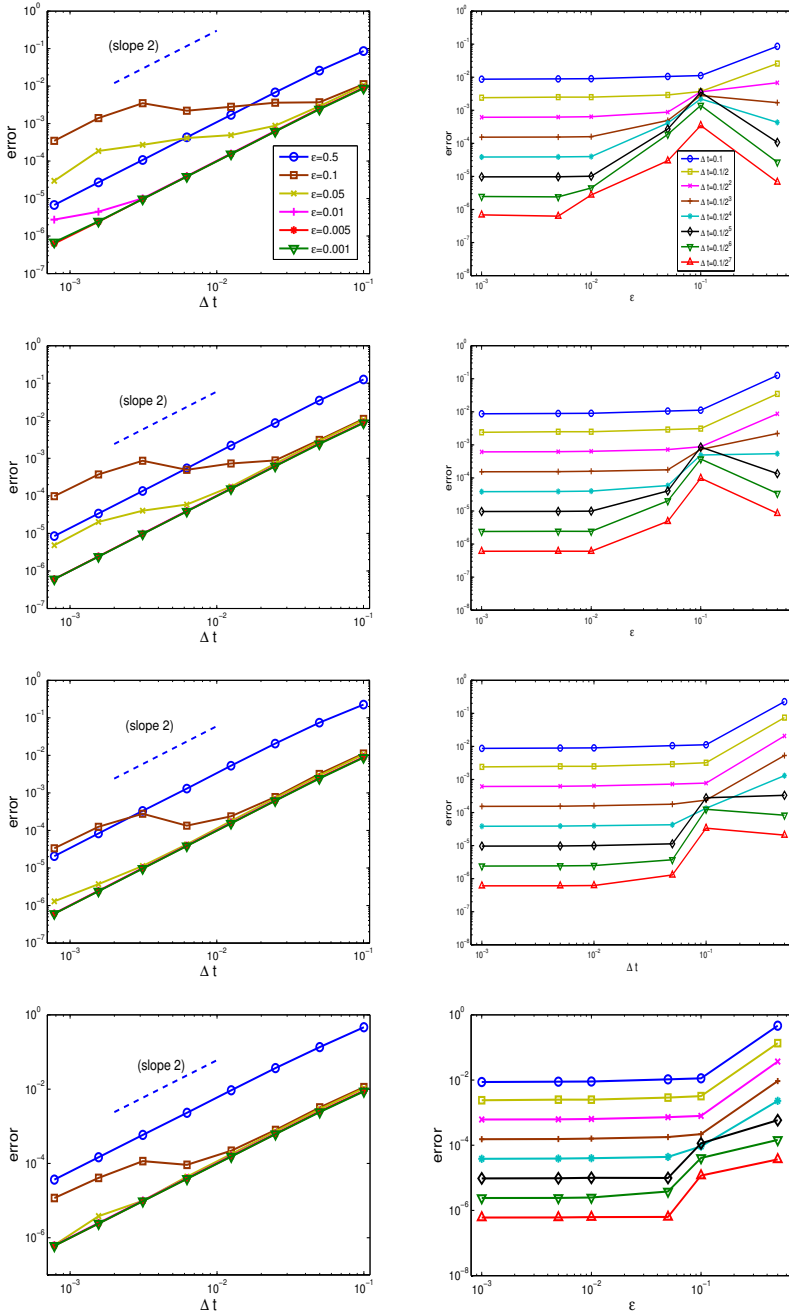


Fig. 4.3: Temporal error of UA2 in Example II with respect to Δt and ε : results of using U_2^ε (first row); results of using U_3^ε (second row) results of using U_4^ε (third row); results of using U_5^ε (last row).

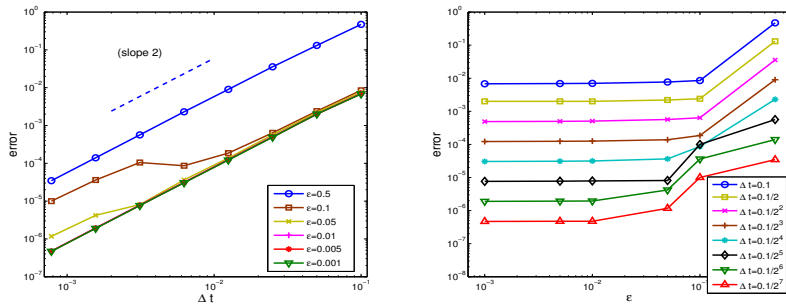


Fig. 4.4: Temporal error of UA2 with U_ε^ξ in Example III with respect to Δt and ε : results of the nonlinear Dirac equation with magnetic potential case.

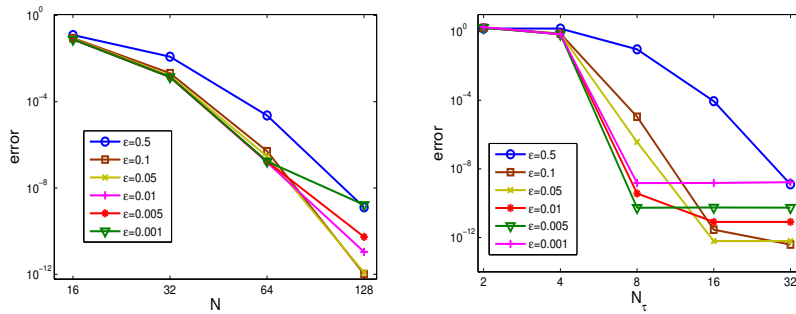


Fig. 4.5: Spatial error of UA2 in Example III with respect to $N(=(b-a)/\Delta x)$ and N_τ .

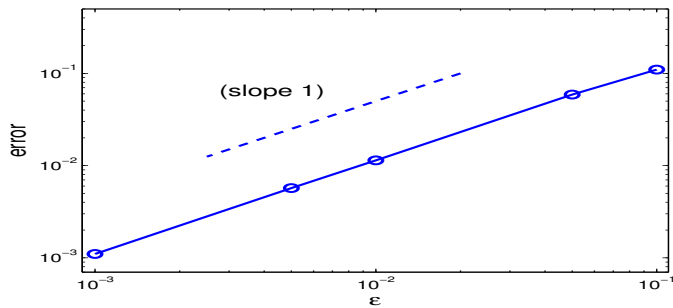


Fig. 4.6: The maximum error in solution between the nonlinear Dirac equation and the limit model (2.8) in Example III.