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▶ To cite this version:

Alexey Agaltsov, Roman Novikov. Error estimates for phaseless inverse scattering in the Born approximation at high energies. 2016. hal-01303885v1

HAL Id: hal-01303885 https://hal.science/hal-01303885v1

Preprint submitted on 18 Apr 2016 (v1), last revised 8 Dec 2016 (v2)

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Error estimates for phaseless inverse scattering in the Born approximation at high energies

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April 18, 2016

Abstract. We study explicit formulas for phaseless inverse scattering in the Born approximation at high energies for the Schrödinger equation with compactly supported potential in dimension $d \ge 2$. We obtain error estimates for these formulas in the configuration space.

1 Introduction

We consider the time-independent Schrödinger equation

$$-\Delta \psi + v(x)\psi = E\psi, \quad x \in \mathbb{R}^d, \ d \ge 2, \ E > 0, \tag{1}$$

where

$$v \in L^{\infty}(\mathbb{R}^d), \quad \operatorname{supp} v \subset D,$$
(2)

where D is some fixed open bounded domain in \mathbb{R}^d .

In quantum mechanics equation (1) describes an elementary particle interacting with a macroscopic object contained in D at fixed energy E.

For equation (1) we consider the classical scattering solutions $\psi^+ = \psi^+(x,k)$, $x \in \mathbb{R}^d, k \in \mathbb{R}^d, k^2 = E$. These solutions ψ^+ can be specified by the following asymptotics as $|x| \to \infty$:

$$\psi^{+}(x,k) = e^{ikx} + c(d,|k|) \frac{e^{i|k||x|}}{|x|^{(d-1)/2}} f(k,|k|\frac{x}{|x|}) + O(|x|^{-(d+1)/2}), \qquad (3)$$
$$x \in \mathbb{R}^{d}, \ k \in \mathbb{R}^{d}, \ k^{2} = E, \ c(d,|k|) = -\pi i (-2\pi i)^{(d-1)/2} |k|^{(d-3)/2},$$

for some a priori unknown f. The function f arising in (3) is defined on

$$\mathcal{M}_E = \left\{ (k,l) \in \mathbb{R}^d \times \mathbb{R}^d \colon k^2 = l^2 = E \right\},\tag{4}$$

and is known as the classical scattering amplitude for equation (1).

In quantum mechanics $|f(k, l)|^2$ describes the probability density of scattering of particle with initial impulse k into direction $l/|l| \neq k/|k|$, and is known as differential scattering cross section for equation (1); see, e.g., [11, Chapter 1, Section 6].

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The problem of finding ψ^+ and f from v is known as the direct scattering problem for equation (1). For solving this problem, one can use, in particular, the Lippmann-Schwinger integral equation for ψ^+ and an explicit integral formula for f, see, e.g., [5, 10, 26].

In turn, the problem of finding v from f is known as the inverse scattering problem (with phase information) and the problem of finding v from $|f|^2$ is known as the phaseless inverse scattering problem for equation (1).

There is a vast literature on the former inverse scattering problem with phase information; see [3, 4, 6, 7, 8, 9, 12, 13, 14, 15, 19, 20, 21, 22, 25] and references therein. In particular, it is well known that the scattering amplitude f uniquely determines v via the Born approximation formulas at high energies:

$$\widehat{v}(k-l) = f(k,l) + O(E^{-\frac{1}{2}}), \quad E \to +\infty, \ (k,l) \in \mathcal{M}_E, \tag{5}$$

$$\widehat{v}(p) = (2\pi)^{-d} \int_{\mathbb{R}^d} e^{ipx} v(x) \, dx, \quad p \in \mathbb{R}^d, \tag{6}$$

and the inverse Fourier transform; see, e.g., [9, 25].

On the other hand, the literature for the phaseless case is much more limited; see [7, 26] and references therein for the case of the aforementioned phaseless inverse problem and see [16, 17, 18, 23, 24, 26] and references therein for the case of some similar inverse problems without phase information. In addition, it is well known that the phaseless scattering data $|f|^2$ does not determine vuniquely, even if $|f|^2$ is given completely for all positive energies. In particular, it is known that

$$f_y(k,l) = e^{i(k-l)y} f(k,l),$$

$$|f_y(k,l)|^2 = |f(k,l)|^2, \quad k, l \in \mathbb{R}^d, \ k^2 = l^2 > 0,$$
(7)

where f is the scattering amplitude for v and f_y is the scattering amplitude for $v_y = v(\cdot - y)$, where $y \in \mathbb{R}^d$; see [26] and references therein.

In the present work, in view of the aforementioned non-uniqueness for the problem of finding v from $|f|^2$, we consider the modified phaseless inverse scattering problem formulated below as Problem 1. Let

$$S = \{ |f|^2, |f_1|^2, \dots, |f_m|^2 \},$$
(8)

where f is the scattering amplitude for v and f_1, \ldots, f_m are the scattering amplitudes for v_1, \ldots, v_m , where

$$v_j = v + w_j, \quad j = 1, \dots, m, \tag{9}$$

where w_1, \ldots, w_m are additional a priori known background scatterers such that $w_i \in L^{\infty}(\mathbb{R}^d)$, supp $w_i \subset \Omega_i$.

$$\Omega_{j} \text{ is an open bounded domain in } \mathbb{R}^{d}, \quad \Omega_{j} \cap D = \emptyset,$$

$$w_{j} \neq 0, \quad w_{j_{1}} \neq w_{j_{2}} \text{ if } j_{1} \neq j_{2} \text{ (in } L^{\infty}(\mathbb{R}^{d})),$$
(10)

where $j, j_1, j_2 \in \{1, \ldots, m\}$. Thus, S consists of the phaseless scattering data $|f|^2, |f_1|^2, \ldots, |f_m|^2$ measured sequentially, first, for the unknown scatterer v and then for v in the presence of known scatterer w_j disjoint from v for $j = 1, \ldots, m$.

Actually, in the present work we continue studies of [26] on the following inverse scattering problem for equation (1):

Problem 1. Reconstruct potential v from the phaseless scattering data S for some appropriate background scatterers w_1, \ldots, w_m .

Studies of Problem 1 in dimension $d \ge 2$ were started in [26]. In dimension d = 1 for m = 1 studies of Problem 1 were started earlier in [2], where phaseless scattering data was considered for all E > 0.

Actually, the key result of [26] consists in a proper extension of formula (5) for the Fourier transform \hat{v} of v to the phaseless case of Problem 1; see Section 2.

In the present work we proceed from the aforementioned result of [26] and study related approximate reconstruction of v in the configuration space. In this connection our results consist in obtaining related error estimates in the configuration space at high energies E; see Section 4.

In addition, results of the present work are necessary for extending the iterative algorithm of [25] to the phaseless case of Problem 1. The latter extension will be given in [1].

2 Extension of formula (5) to the phaseless case

Actually, the key result of [26] consists in the following formulas for solving Problem 1 in dimension $d \ge 2$ for m = 2 at high energies E:

$$\begin{pmatrix} \operatorname{Re} \widehat{v} \\ \operatorname{Im} \widehat{v} \end{pmatrix} = \frac{1}{2} \begin{pmatrix} \operatorname{Re} \widehat{w}_1 & \operatorname{Im} \widehat{w}_1 \\ \operatorname{Re} \widehat{w}_2 & \operatorname{Im} \widehat{w}_2 \end{pmatrix}^{-1} \begin{pmatrix} |\widehat{v}_1|^2 - |\widehat{v}|^2 - |\widehat{w}_1|^2 \\ |\widehat{v}_2|^2 - |\widehat{v}|^2 - |\widehat{w}_2|^2 \end{pmatrix},$$
(11)

$$\begin{aligned} |\hat{v}_{j}(p)|^{2} &= |f_{j}(k,l)|^{2} + O(E^{-\frac{1}{2}}), \quad E \to +\infty, \\ p \in \mathbb{R}^{d}, \quad (k,l) \in \mathcal{M}_{E}, \quad k-l=p, \quad j=0,1,2, \end{aligned}$$
(12)

where:

- $v_0 = v, v_j$ is defined by (9), j = 1, 2, and $f_0 = f, f_1, f_2$ are the scattering amplitudes for v_0, v_1, v_2 , respectively;
- $\hat{v} = \hat{v}(p), \ \hat{v}_j = \hat{v}_j(p), \ \hat{w}_j = \hat{w}_j(p), \ p \in \mathbb{R}^d$, are the Fourier transforms of v, $v_j, \ w_j$ (defined as in (6));
- formula (11) is considered for all $p \in \mathbb{R}^d$ such that the determinant

. .

$$\zeta_{\widehat{w}_1,\widehat{w}_2}(p) \stackrel{\text{def}}{=\!\!=} \operatorname{Re} \widehat{w}_1(p) \operatorname{Im} \widehat{w}_2(p) - \operatorname{Im} \widehat{w}_1(p) \operatorname{Re} \widehat{w}_2(p) \neq 0.$$
(13)

The point is that using formulas (12) for $d \ge 2$ with

$$k = k_E(p) = \frac{p}{2} + \left(E - \frac{p^2}{4}\right)^{1/2} \gamma(p),$$

$$l = l_E(p) = -\frac{p}{2} + \left(E - \frac{p^2}{4}\right)^{1/2} \gamma(p),$$

$$|\gamma(p)| = 1, \quad \gamma(p)p = 0,$$
(14)

where $p \in \mathbb{R}^d$, $|p| \leq 2\sqrt{E}$, one can reconstruct $|\hat{v}|^2$, $|\hat{v}_1|^2$, $|\hat{v}_2|^2$ from S at high energies for any $p \in \mathbb{R}^d$. And then using formula (11) one can reconstruct \hat{v} completely, provided that condition (13) is fulfilled for almost all $p \in \mathbb{R}^d$.

Remark 1. Formulas (12) can be precised as formula (2.15) of [26]:

$$\left| |\widehat{v}_{j}(p)|^{2} - |f_{j}(k,l)|^{2} \right| \leq c(D_{j})N_{j}^{3}E^{-\frac{1}{2}},$$

$$p = k - l, \ (k,l) \in \mathcal{M}_{E}, \ E^{\frac{1}{2}} \geq \rho(D_{j},N_{j}), \ j = 0, 1, 2,$$
(15)

where $||v_j||_{L^{\infty}(D_j)} \leq N_j$, j = 0, 1, 2, and $D_0 = D$, $D_j = D \cup \Omega_j$, j = 1, 2, and formulas for the constants c, ρ are given in [26].

In addition, from the experimental point of view it seems to be, in particular, convenient to consider Problem 1 with m = 2 for the case when w_2 is just a translation of w_1 :

$$w_2(x) = w_1(x - y), \quad x \in \mathbb{R}^d, \ y \in \mathbb{R}^d.$$
(16)

In this case

$$\widehat{w}_2(p) = e^{ipy}\widehat{w}_1(p), \quad \zeta_{\widehat{w}_1,\widehat{w}_2}(p) = \sin(py)|\widehat{w}_1(p)|^2, \quad p \in \mathbb{R}^d.$$
 (17)

On the level of analysis, the principal complification of (11), (12) in comparison with (5) consists in possible zeros of the determinant $\zeta_{\widehat{w}_1,\widehat{w}_2}$ of (13). For some simplest cases, we study these zeros in the next section.

3 Zeros of the determinant $\zeta_{\widehat{w}_1, \widehat{w}_2}$

Let

$$Z_{\widehat{w}_{1},\widehat{w}_{2}} = \{ p \in \mathbb{R}^{d} : \zeta_{\widehat{w}_{1},\widehat{w}_{2}}(p) = 0 \},$$

$$Z_{\widehat{w}_{j}} = \{ p \in \mathbb{R}^{d} : \widehat{w}_{j}(p) = 0 \}, \quad j = 1, 2,$$
(18)

where ζ is defined by (13). From (13), (18) it follows that

$$Z_{\widehat{w}_1} \cup Z_{\widehat{w}_2} \subseteq Z_{\widehat{w}_1, \widehat{w}_2}.$$
(19)

In view of (19), in order to construct examples of w_1 , w_2 such that the set $Z_{\widehat{w}_1,\widehat{w}_2}$ is as simple as possible, we use the following lemma:

Lemma 1. Let

$$w(x) = |x|^{\nu} K_{\nu}(|x|) \int_{\mathbb{R}^d} q(x-y)q(y) \, dy, \quad x \in \mathbb{R}^d, \ \nu > 0, \tag{20}$$

$$K_{\nu}(s) = \frac{\Gamma(\frac{1}{2} + \nu)}{\sqrt{\pi}} \left(\frac{2}{s}\right)^{\nu} \int_{0}^{\infty} \frac{\cos(st) \, dt}{(1 + t^2)^{\frac{1}{2} + \nu}}, \quad s > 0, \tag{21}$$

$$q \in L^{\infty}(\mathbb{R}^d), \ q = \overline{q}, \ q \neq 0 \ in \ L^{\infty}(\mathbb{R}^d),$$

$$q(x) = 0 \ if \ |x| > r, \ q(x) = q(-x), \ x \in \mathbb{R}^d.$$
(22)

Then

$$w \in C(\mathbb{R}^d), \ w = \overline{w}, \ w(x) = 0 \ \text{if } |x| > 2r, \ x \in \mathbb{R}^d,$$
$$\widehat{w}(p) = \overline{\widehat{w}(p)} \ge c_1 (1+|p|)^{-\beta}, \quad p \in \mathbb{R}^d,$$
(23)

for $\beta = d + 2\nu$ and some positive constant $c_1 = c_1(q, \nu)$, where \widehat{w} is the Fourier transform of w. In addition, if $q \ge 0$, then $w \ge 0$.

We recall that K_{ν} defined by (21) is the modified Bessel function of the second kind and order ν . In addition, Γ denotes the gamma function.

Lemma 1 is proved in Section 8.

As a corollary of Lemma 1, functions

$$w_j(x) = w(x - T_j), \quad x \in \mathbb{R}^d, \ T_j \in \mathbb{R}^d,$$
(24)

where w is constructed in Lemma 1, give us examples of w_j satisfying (10) for fixed D, Ω_j and for appropriate radius r of Lemma 1 and translations T_j of (24), and such that

$$Z_{\widehat{w}_j} = \emptyset,$$

$$|\widehat{w}_j(p)| = \widehat{w}(p) \ge c_1 (1+|p|)^{-\beta}, \quad p \in \mathbb{R}^d,$$

(25)

where c_1 , β are the same as in (23). In addition,

$$\begin{aligned}
\zeta_{\widehat{w}_1,\widehat{w}_2}(p) &= \sin(py) |\widehat{w}(p)|^2, \quad y = T_2 - T_1 \neq 0, \quad p \in \mathbb{R}^d, \\
Z_{\widehat{w}_1,\widehat{w}_2} &= \left\{ p \in \mathbb{R}^d \colon \sin(py) = 0 \right\} = \left\{ p \in \mathbb{R}^d \colon py \in \pi\mathbb{Z} \right\},
\end{aligned}$$
(26)

for w_1, w_2 of (24).

As another corollary of Lemma 1, we have that

$$\begin{aligned} \zeta_{\widehat{w}_1,\widehat{w}_2}(p) &= |\widehat{w}(p)|^2 \ge c_1^2 (1+|p|)^{-2\beta}, \quad p \in \mathbb{R}^d, \\ Z_{\widehat{w}_1,\widehat{w}_2} &= \varnothing, \\ \text{if } w_1 \text{ is defined as in (24) and } w_2 &= iw_1. \end{aligned}$$
(27)

Note that complex-valued v and w_j naturally arise if we interpret equation (1) as a time-harmonic Helmholtz equation of acoustics or electrodynamics.

Finally, note that

$$Z_{\widehat{w}_1,\dots,\widehat{w}_{d+1}} = \frac{\pi}{s} \mathbb{Z}^d, \quad \text{where}$$

$$Z_{\widehat{w}_1,\dots,\widehat{w}_{d+1}} = Z_{\widehat{w}_1,\widehat{w}_2} \cap Z_{\widehat{w}_1,\widehat{w}_3} \cap \dots \cap Z_{\widehat{w}_1,\widehat{w}_{d+1}},$$

$$(28)$$

if w_1 is defined as in (24), and

$$w_2(x) = w_1(x - se_1), \dots, w_{d+1}(x) = w_1(x - se_d),$$
(29)

where (e_1, \ldots, e_d) is the standard basis of \mathbb{R}^d and s > 0.

Thus, in principle, for Problem 1 with background scatterers w_1, \ldots, w_{d+1} as in (29), for each $p \in \mathbb{R}^d \setminus \frac{\pi}{s} \mathbb{Z}^d$ formulas (11), (12) can be used with appropriate w_j in place of w_2 , where $j = 2, \ldots, d+1$.

4 Error estimates in the configuration space

We recall that for inverse scattering with phase information the scattering amplitude f on \mathcal{M}_E processed by (5) and the inverse Fourier transform yield the approximate reconstruction

$$u(\cdot, E) = v + O(E^{-\alpha})$$
 in $L^{\infty}(\mathbb{R}^d)$ as $E \to +\infty$, $\alpha = \frac{n-d}{2n}$, (30)

if $v \in W^{n,1}(\mathbb{R}^d)$, n > d (in addition to the initial assumption (2)), where $W^{n,1}(\mathbb{R}^d)$ denotes the standard Sobolev space of *n*-times differentiable functions in $L^1(\mathbb{R}^d)$:

$$W^{n,1}(\mathbb{R}^d) = \left\{ u \in L^1(\mathbb{R}^d) \colon \|u\|_{n,1} < \infty \right\},$$

$$\|u\|_{n,1} = \max_{|J| \le n} \left\| \frac{\partial^{|J|} u}{\partial x^J} \right\|_{L^1(\mathbb{R}^d)}, \quad n \in \mathbb{N} \cup \{0\}.$$
 (31)

More precisely, the approximation $u(\cdot, E)$ in (30) is defined by

$$u(x,E) = \int_{\mathcal{B}_{r(E)}} e^{-ipx} f(k_E(p), l_E(p)) \, dp, \quad x \in \mathbb{R}^d,$$

$$r(E) = 2\tau E^{\frac{\alpha}{n-d}} \quad \text{for some fixed } \tau \in (0,1],$$

(32)

where

$$\mathcal{B}_r = \left\{ p \in \mathbb{R}^d \colon |p| \le r \right\},\tag{33}$$

 α is defined in (30), and $k_E(p)$, $l_E(p)$ are defined as in (14) with some piecewise continuous vector-function γ on \mathbb{R}^d ; see, e.g., [25].

Analogs of $u(\cdot, E)$ for the phaseless case are given below in this section. In particular, related formulas depend on the zeros of determinant $\zeta_{\widehat{w}_1,\widehat{w}_2}$ of (13).

We consider

$$U_{\widehat{w}_{1},\widehat{w}_{2}} = \operatorname{Re} U_{\widehat{w}_{1},\widehat{w}_{2}} + i \operatorname{Im} U_{\widehat{w}_{1},\widehat{w}_{2}}, \\ \begin{pmatrix} \operatorname{Re} U_{\widehat{w}_{1},\widehat{w}_{2}} \\ \operatorname{Im} U_{\widehat{w}_{1},\widehat{w}_{2}} \end{pmatrix} = \frac{1}{2} M_{\widehat{w}_{1},\widehat{w}_{2}}^{-1}(p) b_{\widehat{w}_{1},\widehat{w}_{2}}(p, E),$$
(34)

$$M_{\widehat{w}_1,\widehat{w}_2}(p) = \begin{pmatrix} \operatorname{Re} \widehat{w}_1(p) & \operatorname{Im} \widehat{w}_1(p) \\ \operatorname{Re} \widehat{w}_2(p) & \operatorname{Im} \widehat{w}_2(p) \end{pmatrix},$$
(35)

$$M_{\widehat{w}_1,\widehat{w}_2}^{-1}(p) = \frac{1}{\zeta_{\widehat{w}_1,\widehat{w}_2}(p)} \begin{pmatrix} \operatorname{Im} \widehat{w}_2(p) & -\operatorname{Im} \widehat{w}_1(p) \\ -\operatorname{Re} \widehat{w}_2(p) & \operatorname{Re} \widehat{w}_1(p) \end{pmatrix},$$
(36)

$$b_{\widehat{w}_1,\widehat{w}_2}(p,E) = \begin{pmatrix} |f_1(p,E)|^2 - |f(p,E)|^2 - |\widehat{w}_1(p)|^2 \\ |f_2(p,E)|^2 - |f(p,E)|^2 - |\widehat{w}_2(p)|^2 \end{pmatrix},$$
(37)

$$f(p, E) = f(k_E(p), l_E(p)), \ f_j(p, E) = f_j(k_E(p), l_E(p)), \ j = 1, 2,$$
(38)

where \widehat{w}_1 , \widehat{w}_2 , f, f_1 , f_2 are the same as in (11), (12), $\zeta_{\widehat{w}_1,\widehat{w}_2}$ is defined by (13), $k_E(p)$, $l_E(p)$ are the same as in (14), (32), and $p \in \mathcal{B}_{2\sqrt{E}}$, $d \geq 2$.

For Problem 1 for $d \ge 2$, m = 2, and for the case when $\zeta_{\widehat{w}_1,\widehat{w}_2}$ has no zeros (the case of (27) in Section 3) we have the following result:

Theorem 1. Let v satisfy (2) and $v \in W^{n,1}(\mathbb{R}^d)$ for some n > d. Let w_1, w_2 be the same as in (27). Let

$$u(x,E) = \int_{\mathcal{B}_{r_1(E)}} e^{-ipx} U_{\widehat{w}_1,\widehat{w}_2}(p,E) \, dp, \quad x \in \mathbb{R}^d,$$

$$r_1(E) = 2\tau E^{\frac{\alpha_1}{n-d}}, \quad \alpha_1 = \frac{n-d}{2(n+\beta)}, \quad \text{for some fixed } \tau \in (0,1],$$
(39)

where $U_{\widehat{w}_1,\widehat{w}_2}$ is defined by (34), \mathcal{B}_r is defined by (33), β is the number of (23), (27). Then

$$u(\cdot, E) = v + O(E^{-\alpha_1}) \quad in \ L^{\infty}(\mathbb{R}^d), \ E \to +\infty.$$
(40)

Theorem 1 is proved in Section 5. Next, we set

$$Z_{\widehat{w}_1,\widehat{w}_2}^{\varepsilon} = \left\{ p \in \mathbb{R}^d \colon py \in (-\varepsilon,\varepsilon) + \pi\mathbb{Z} \right\}, \quad y \in \mathbb{R}^d \setminus 0, \quad 0 < \varepsilon < 1,$$
(41)

where \widehat{w}_1 , \widehat{w}_2 and y are the same as in (24)–(26). One can see that $Z_{\widehat{w}_1,\widehat{w}_2}^{\varepsilon}$ is the open $\frac{\varepsilon}{|y|}$ -neighborhood of $Z_{\widehat{w}_1,\widehat{w}_2}$ defined in (26).

Note that

for any
$$p \in Z_{\widehat{w}_1, \widehat{w}_2}^{\varepsilon}$$
 there exists
the unique $z(p) \in \mathbb{Z}$ such that $|py - \pi z(p)| < \varepsilon$. (42)

In addition to $U_{\widehat{w}_1,\widehat{w}_2}$ of (34), we define

$$U_{\widehat{w}_{1},\widehat{w}_{2}}^{\varepsilon}(p,E) = \frac{1}{2} \left(U_{\widehat{w}_{1},\widehat{w}_{2}}(p_{-}^{\varepsilon},E) + U_{\widehat{w}_{1},\widehat{w}_{2}}(p_{+}^{\varepsilon},E) \right),$$

$$p_{\pm}^{\varepsilon} = p_{\perp} + \pi z(p) \frac{y}{|y|^{2}} \pm \varepsilon \frac{y}{|y|^{2}}, \quad p_{\perp} = p - (py) \frac{y}{|y|^{2}}, \quad p \in \mathcal{B}_{2\sqrt{E}} \cap Z_{\widehat{w}_{1},\widehat{w}_{2}}^{\varepsilon},$$

(43)

where z(p) is the integer number of (42).

For Problem 1 for $d \ge 2$, m = 2, and for the case when $\zeta_{\widehat{w}_1,\widehat{w}_2}$ has zeros on hyperplanes (the case of (26) in Section 3) we have the following result:

Theorem 2. Let v satisfy (2) and $v \in W^{n,1}(\mathbb{R}^d)$ for some n > d. Let w_1, w_2 be the same as in (24)–(26). Let

$$u(x, E) = u_{1}(x, E) + u_{2}(x, E), \quad x \in \mathbb{R}^{d},$$

$$u_{1}(x, E) = \int_{\mathcal{B}_{r_{2}(E)} \setminus Z_{\widehat{w}_{1}, \widehat{w}_{2}}^{\varepsilon_{2}(E)}} e^{-ipx} U_{\widehat{w}_{1}, \widehat{w}_{2}}(p, E) \, dp,$$

$$u_{2}(x, E) = \int_{\mathcal{B}_{r_{2}(E)} \cap Z_{\widehat{w}_{1}, \widehat{w}_{2}}^{\varepsilon_{2}(E)}} e^{-ipx} U_{\widehat{w}_{1}, \widehat{w}_{2}}^{\varepsilon}(p, E) \, dp,$$

$$r_{2}(E) = 2\tau E^{\frac{\alpha_{2}}{n-d}}, \quad \varepsilon_{2}(E) = E^{-\frac{\alpha_{2}}{2}},$$

$$\alpha_{2} = \frac{n-d}{2\left(n+\beta+\frac{n-d}{2}\right)}, \quad for \ some \ fixed \ \tau \in (0,1],$$

$$(44)$$

where $U_{\widehat{w}_1,\widehat{w}_2}$ and $U_{\widehat{w}_1,\widehat{w}_2}^{\varepsilon}$ are defined by (34), (43), \mathcal{B}_r and $Z_{\widehat{w}_1,\widehat{w}_2}^{\varepsilon}$ are defined by (33), (41), and β is the number of (23). Then

$$u(\cdot, E) = v + O(E^{-\alpha_2}) \quad in \ L^{\infty}(\mathbb{R}^d), \ E \to +\infty.$$
(45)

Theorem 2 is proved in Section 6. Next, we set

$$Z^{\varepsilon}_{\widehat{w}_1,\dots,\widehat{w}_{d+1}} = \mathcal{B}'_{\varepsilon/s} + \frac{\pi}{s} \mathbb{Z}^d, \quad 0 < \varepsilon < 1, \\ \mathcal{B}'_r = \mathcal{B}_r \setminus \partial \mathcal{B}_r, \quad r > 0,$$

$$(46)$$

where $\widehat{w}_1, \ldots, \widehat{w}_{d+1}$ are the same as in (28), (29), and \mathcal{B}_r is defined by (33). One can see that $Z_{\widehat{w}_1,\ldots,\widehat{w}_{d+1}}^{\varepsilon}$ is the open $\frac{\varepsilon}{s}$ -neighborhood of $Z_{\widehat{w}_1,\ldots,\widehat{w}_{d+1}}$ defined in (28).

Note that

for any
$$p \in Z^{\varepsilon}_{\widehat{w}_1,...,\widehat{w}_{d+1}}$$
 there exists
the unique $z(p) \in \mathbb{Z}^d$ such that $|sp - \pi z(p)| < \varepsilon$. (47)

In addition, we consider i' such that

$$i' = i'(p, s), \quad p = (p_1, \dots, p_d) \in \mathbb{R}^d \setminus \frac{\pi}{s} \mathbb{Z}^d, \ s > 0,$$

$$i' \text{ take values in } \{2, \dots, d+1\},$$

$$|\sin(sp_{i'-1})| \ge |\sin(sp_{i-1})| \quad \text{for all } i \in \{2, \dots, d+1\}.$$
(48)

$$U_{\widehat{w}_1,\dots,\widehat{w}_{d+1}}(p,E) = U_{\widehat{w}_1,\widehat{w}_{i'}}(p,E), \quad p \in \mathbb{R}^d \setminus \frac{\pi}{s} \mathbb{Z}^d, \qquad (49)$$
$$U_{\widehat{w}_1,\dots,\widehat{w}_{d+1}}^{\varepsilon}(p,E) =$$

$$\frac{1}{|\mathbb{S}^{d-1}|} \int_{\mathbb{S}^{d-1}} U_{\widehat{w}_1,\dots,\widehat{w}_{d+1}}(\frac{\varepsilon}{s}\vartheta + \frac{\pi}{s}z(p), E) \, d\vartheta, \quad p \in Z^{\varepsilon}_{\widehat{w}_1,\dots,\widehat{w}_{d+1}}, \qquad (50)$$
$$\mathbb{S}^{d-1} = \{ p \in \mathbb{R}^d \colon |p| = 1 \}, \qquad (51)$$

where $|\mathbb{S}^{d-1}|$ denotes the standard Euclidean volume of \mathbb{S}^{d-1} , $U_{\widehat{w}_1,\widehat{w}_{i'}}$ is defined in a similar way with $U_{\widehat{w}_1,\widehat{w}_2}$ of (34), $\widehat{w}_1,\ldots,\widehat{w}_{d+1}$ are the same as in (28), (29), and i' is the same as in (48).

For Problem 1 for $d \ge 2$, m = d + 1, we also have the following result:

Theorem 3. Let v satisfy (2) and $v \in W^{n,1}(\mathbb{R}^d)$ for some n > d. Let w_1, \ldots, w_{d+1} be the same as in (29). Let

$$u(x, E) = u_{1}(x, E) + u_{2}(x, E), \quad x \in \mathbb{R}^{d},$$

$$u_{1}(x, E) = \int_{\mathcal{B}_{r_{3}(E)} \setminus Z_{\hat{w}_{1}, \dots, \hat{w}_{d+1}}^{\varepsilon_{3}(E)}} e^{-ipx} U_{\hat{w}_{1}, \dots, \hat{w}_{d+1}}(p, E) \, dp,$$

$$u_{2}(x, E) = \int_{\mathcal{B}_{r_{3}(E)} \cap Z_{\hat{w}_{1}, \dots, \hat{w}_{d+1}}^{\varepsilon_{3}(E)}} e^{-ipx} U_{\hat{w}_{1}, \dots, \hat{w}_{d+1}}^{\varepsilon}(p, E) \, dp,$$

$$r_{3}(E) = 2\tau E^{\frac{\alpha_{3}}{n-d}}, \quad \varepsilon_{2}(E) = E^{-\frac{\alpha_{3}}{d+1}},$$

$$\alpha_{3} = \frac{n-d}{2\left(n+\beta+\frac{n-d}{d+1}\right)}, \quad for \ some \ fixed \ \tau \in (0,1],$$
(52)

where $U_{\widehat{w}_1,\ldots,\widehat{w}_{d+1}}$ and $U_{\widehat{w}_1,\ldots,\widehat{w}_{d+1}}^{\varepsilon}$ are defined by (49), (50), \mathcal{B}_r and $Z_{\widehat{w}_1,\ldots,\widehat{w}_{d+1}}^{\varepsilon}$ are defined by (33), (46), and β is the number of (23). Then

$$u(\cdot, E) = v + O(E^{-\alpha_3}) \quad in \ L^{\infty}(\mathbb{R}^d), \ E \to +\infty.$$
(53)

Theorem 3 is proved in Section 7.

5 Proof of Theorem 1

Proposition 1. Let v satisfy (2) and w_1 , w_2 be the same as in (27), $d \ge 2$. Then:

$$\begin{aligned} \left| \widehat{v}(p) - U_{\widehat{w}_1, \widehat{w}_2}(p, E) \right| &\leq c_2 |\widehat{w}(p)|^{-1} E^{-\frac{1}{2}} \quad for \ p \in \mathcal{B}_{2\sqrt{E}}, \ E \geq \rho_1, \\ c_2 &= 2c(D_0) N_0^3 + 2c(D_1) N_1^3, \\ \rho_1 &= \max_{j=0,1} \rho(D_j, N_j), \end{aligned}$$
(54)

where w is the function of (23), (24), and c, ρ , N_j , D_j , j = 0, 1, 2, are the same as in estimates (15).

Let

Proposition 1 follows from formulas (11), (14), estimates (15), definitions (34)–(38), and the properties that

$$\Omega_2 = \Omega_1, \ D_2 = D_1, \ N_2 = N_1. \tag{55}$$

In turn, properties (55) follow from (9) for j = 1, 2, (10) for j = 1, and from the equality $w_2 = iw_1$ assumed in (27).

Next, we represent v as follows:

$$v(x) = v^{+}(x, r) + v^{-}(x, r), \quad x \in \mathbb{R}^{d}, \ r > 0,$$
$$v^{+}(x, r) = \int_{\mathcal{B}_{r}} e^{-ipx} \widehat{v}(p) \ dp,$$
$$v^{-}(x, r) = \int_{\mathbb{R}^{d} \setminus \mathcal{B}_{r}} e^{-ipx} \widehat{v}(p) \ dp.$$
(56)

Since $v \in W^{n,1}(\mathbb{R}^d)$, n > d, we have

$$|v^{-}(x,r)| \leq c_{3} ||v||_{n,1} r^{d-n}, \quad x \in \mathbb{R}^{d}, \ r > 0,$$

$$c_{3} = |\mathbb{S}^{d-1}| \frac{(2\pi)^{-d} d^{n}}{n-d},$$
(57)

where $\|\cdot\|_{n,1}$ is defined in (31), and $|\mathbb{S}^{d-1}|$ is the standard Euclidean volume of \mathbb{S}^{d-1} . Indeed,

$$|p_1^{k_1} \cdots p_d^{k_d} \hat{v}(p)| \le (2\pi)^{-d} ||v||_{n,1}, \quad p = (p_1, \dots, p_d) \in \mathbb{R}^d,$$

for any $k_1, \dots, k_d \in \mathbb{N} \cup \{0\}, \, k_1 + \dots + k_d \le n,$ (58)

assuming also that $p_j^0 = 1$. Taking an appropriate sum in (58) over all such k_1 , ..., k_d with $k_1 + \cdots + k_d = m \le n$, we get

$$|p|^{m}|\hat{v}(p)| \le (|p_{1}| + \dots + |p_{d}|)^{m}|\hat{v}(p)| \le (2\pi)^{-d}d^{m}||v||_{n,1}, \quad p \in \mathbb{R}^{d}.$$
 (59)

The definition of v^- of (56) and inequalities (59) for m = n imply (57).

In addition, using Proposition 1 and the estimate on \hat{w} of (23), we obtain:

As a corollary of (56), (57), (60), we have

$$\left| v(x) - \int_{\mathcal{B}_r} e^{-ipx} U_{\widehat{w}_1, \widehat{w}_2}(p, E) \, dp \right| \le c_3 \|v\|_{n, 1} r^{d-n} + c_4 E^{-\frac{1}{2}} r^{d+\beta}, \tag{61}$$

where $x \in \mathbb{R}^d$, $1 \leq r \leq 2\sqrt{E}$, $E \geq \rho_1$. In addition, if $r = r_1(E)$, where $r_1(E)$ is defined in (39), then

$$r^{d-n} = (2\tau)^{d-n} E^{-\alpha_1},$$

$$E^{-\frac{1}{2}} r^{d+\beta} = (2\tau)^{d+\beta} E^{-\alpha_1}.$$
(62)

Formula (40) follows from formulas (61) and (62).

Theorem 1 is proved.

6 Proof of Theorem 2

Proposition 2. Let v satisfy (2) and w_1 , w_2 be the same as in (24)–(26), $d \ge 2$. Then:

$$\begin{aligned} \left| \widehat{v}(p) - U_{\widehat{w}_{1},\widehat{w}_{2}}(p,E) \right| &\leq c_{5}\varepsilon^{-1}(1+|p|)^{\beta}E^{-\frac{1}{2}}, \\ p \in \mathcal{B}_{2\sqrt{E}} \setminus Z_{\widehat{w}_{1},\widehat{w}_{2}}^{\varepsilon}, E \geq \rho_{2}, \ 0 < \varepsilon < 1, \\ c_{5} &= \frac{\pi}{2}(2c(D_{0})N_{0}^{3} + c(D_{1})N_{1}^{3} + c(D_{2})N_{2}^{3})c_{1}, \\ \rho_{2} &= \max_{j=0,1,2}\rho(D_{j},N_{j}), \end{aligned}$$
(63)

in addition, if $v \in W^{n,1}(\mathbb{R}^d)$, $n \ge 0$, then:

$$\begin{aligned} |\widehat{v}(p) - U_{\widehat{w}_{1},\widehat{w}_{2}}^{\varepsilon}(p,E)| \\ &\leq 2^{\beta}c_{5}\varepsilon^{-1}(1+|p|)^{\beta}E^{-\frac{1}{2}} + c_{6}\varepsilon(1+\frac{\pi}{|y|}|z(p)|+|p_{\perp}|)^{-n}, \qquad (64) \\ &p \in \mathcal{B}_{2\sqrt{E}} \cap Z_{\widehat{w}_{1},\widehat{w}_{2}}^{\varepsilon}, \ E \geq \rho_{2}, \ 0 < \varepsilon < \min\{1,\frac{1}{2}|y|\}; \\ &c_{6} = 2^{n}\frac{(d+1)^{n+1}}{(2\pi)^{d}|y|} \max_{j=1,...,d} \|x_{j}v\|_{n,1}, \end{aligned}$$

where $c, \rho, N_j, D_j, j = 0, 1, 2$, are the same as in estimates (15), c_1, β are the same as in Lemma 1; $z(p), p_{\perp}$ are defined in (42), (43), $x_jv = x_jv(x)$, and $\|\cdot\|_{n,1}$ is defined in (31).

Proof of Proposition 2. It follows from formulas (24), (26) and (35), (36) that

$$M_{\widehat{w}_1,\widehat{w}_2}(p) = \widehat{w}(p) \begin{pmatrix} \cos(T_1p) & \sin(T_1p) \\ \cos(T_2p) & \sin(T_2p) \end{pmatrix}, \quad p \in \mathbb{R}^d,$$

$$M_{\widehat{w}_1,\widehat{w}_2}^{-1}(p) = \frac{1}{\sin(py)\widehat{w}(p)} \begin{pmatrix} \sin(T_2p) & -\sin(T_1p) \\ -\cos(T_2p) & \cos(T_1p) \end{pmatrix}, \quad p \in \mathbb{R}^d \setminus Z_{\widehat{w}_1,\widehat{w}_2}^{\varepsilon}.$$
(65)

Also note that

$$|\sin(py)| \ge \frac{2\varepsilon}{\pi}, \quad p \in \mathbb{R}^d \setminus Z^{\varepsilon}_{\widehat{w}_1, \widehat{w}_2}, \ 0 < \varepsilon < 1.$$
(66)

The estimate (63) follows from (11), (15), (23), (34), (37), (38) and from (65), (66).

It remains to prove (64). Using definition (43), one can write

$$U_{\widehat{w}_{1},\widehat{w}_{2}}^{\varepsilon}(p,E) - \widehat{v}(p) = \varphi_{1}^{\varepsilon}(p,E) + \varphi_{2}^{\varepsilon}(p), \quad p \in \mathcal{B}_{2\sqrt{E}} \cap Z_{\widehat{w}_{1},\widehat{w}_{2}}^{\varepsilon},$$

$$\varphi_{1}^{\varepsilon}(p,E) = \frac{1}{2} \left(U_{\widehat{w}_{1},\widehat{w}_{2}}^{\varepsilon}(p_{-}^{\varepsilon},E) - \widehat{v}(p_{-}^{\varepsilon}) \right) + \frac{1}{2} \left(U_{\widehat{w}_{1},\widehat{w}_{2}}^{\varepsilon}(p_{+}^{\varepsilon},E) - \widehat{v}(p_{+}^{\varepsilon}) \right), \quad p \in \mathcal{B}_{2\sqrt{E}} \cap Z_{\widehat{w}_{1},\widehat{w}_{2}}^{\varepsilon},$$

$$\varphi_{2}^{\varepsilon}(p) = \frac{1}{2} \left(\widehat{v}(p_{-}^{\varepsilon}) + \widehat{v}(p_{+}^{\varepsilon}) \right) - \widehat{v}(p), \quad p \in Z_{\widehat{w}_{1},\widehat{w}_{2}}^{\varepsilon}.$$
(67)

Using estimate (63), formula (67) and the definitions of p_{\pm}^{ε} in (43), we get

$$\begin{aligned} |\varphi_1^{\varepsilon}(p,E)| &\leq \frac{1}{2}c_5\varepsilon^{-1}E^{-\frac{1}{2}}\left((1+|p_-^{\varepsilon}|)^{\beta}+(1+|p_+^{\varepsilon}|)^{\beta}\right) \\ &\leq c_5\varepsilon^{-1}(1+|p|+2\frac{\varepsilon}{|y|})^{\beta}E^{-\frac{1}{2}} \leq 2^{\beta}c_5\varepsilon^{-1}(1+|p|)^{\beta}E^{-\frac{1}{2}}, \\ &\text{for } \varepsilon \text{ as in } (64), \ p \in \mathcal{B}_{2\sqrt{E}} \cap Z_{\widehat{w}_1,\widehat{w}_2}^{\varepsilon}. \end{aligned}$$
(68)

Next, using the definition of φ_2^{ε} in (67) and the mean value theorem, we obtain

$$|\varphi_2^{\varepsilon}(p)| \le \frac{\varepsilon}{|y|} \max\{|\frac{y}{|y|} \nabla \widehat{v}(\xi)| \colon \xi \in [p_-^{\varepsilon}, p_+^{\varepsilon}]\}, \quad p \in Z_{\widehat{w}_1, \widehat{w}_2}^{\varepsilon}, \tag{69}$$

where $[p_{-}^{\varepsilon}, p_{+}^{\varepsilon}]$ denotes the segment joining p_{-}^{ε} to p_{+}^{ε} . Here, the mean value theorem was used for $\hat{v}(\xi)$ on $[p_{-}^{\varepsilon}, p]$ and on $[p, p_{+}^{\varepsilon}]$.

Note also that

$$|\nabla \widehat{v}(\xi)| \le d \max_{j=1,\dots,d} \left| \frac{\partial \widehat{v}}{\partial \xi_j}(\xi) \right|, \quad \xi = (\xi_1,\dots,\xi_d) \in [p_-^\varepsilon, p_+^\varepsilon]. \tag{70}$$

In addition, the following estimates hold:

$$\left|\frac{\partial \widehat{v}}{\partial \xi_j}(\xi)\right| \le \frac{(1+d)^n}{(2\pi)^d (1+|\xi|)^n} \|x_j v\|_{n,1}, \quad \xi \in [p_-^\varepsilon, p_+^\varepsilon], \ j = 1, \dots, d.$$
(71)

Indeed, taking the sum in (59) over all m = 0, ..., n with the binomial coefficients, we get

$$(1+|p|)^{n}|\widehat{v}(p)| \leq (1+|p_{1}|+\dots+|p_{d}|)^{n}|\widehat{v}(p)| \leq (2\pi)^{-d}(1+d)^{n}||v||_{n,1}, \quad p \in \mathbb{R}^{d}.$$

$$(72)$$

Estimates (71) follow from (72), where we replace v by $x_j v$ and use that v belongs to $W^{n,1}(\mathbb{R}^d)$ and is compactly supported.

Estimates (69)-(71) imply

$$|\varphi_2^{\varepsilon}(p)| \le 2^{-n} c_6 \varepsilon \max\left\{ \left(1 + |\xi|\right)^{-n} \colon \xi \in [p_-^{\varepsilon}, p_+^{\varepsilon}] \right\}, \quad p \in Z_{\widehat{w}_1, \widehat{w}_2}^{\varepsilon}.$$
(73)

Using also that

$$\xi = \tau \frac{y}{|y|} + p_{\perp}, \text{ where } |\tau - \frac{\pi}{|y|} z(p)| \le \frac{\varepsilon}{|y|}, \text{ if } \xi \in [p_{-}^{\varepsilon}, p_{+}^{\varepsilon}], \tag{74}$$

and that $\varepsilon < |y|$, we obtain

$$\begin{aligned} |\varphi_2^{\varepsilon}(p)| &\leq 2^{-n} c_6 \varepsilon \left(1 + \frac{1}{2} \left(\frac{\pi}{|y|} |z(p)| - \frac{\varepsilon}{|y|} + |p_{\perp}|\right)\right)^{-n} \\ &\leq c_6 \left(1 + \frac{\pi}{|y|} |z(p)| + |p_{\perp}|\right)^{-n}, \quad p \in Z_{\widehat{w}_1, \widehat{w}_2}^{\varepsilon}. \end{aligned}$$
(75)

Estimate (64) follows from (68) and (75).

Proposition 2 is proved.

The final part of the proof of Theorem 2 is as follows. In a similar way with (56), we represent v as follows:

$$v(x) = v_{1}^{+}(x,r) + v_{2}^{+}(x,r) + v^{-}(x,r), \quad x \in \mathbb{R}^{d}, \ r > 0,$$

$$v_{1}^{+}(x,r) = \int_{\mathcal{B}_{r} \setminus Z_{\widehat{w}_{1},\widehat{w}_{2}}^{\varepsilon}} e^{-ipx} \widehat{v}(p) \ dp,$$

$$v_{2}^{+}(x,r) = \int_{\mathcal{B}_{r} \cap Z_{\widehat{w}_{1},\widehat{w}_{2}}^{\varepsilon}} e^{-ipx} \widehat{v}(p) \ dp,$$

$$v^{-}(x,r) = \int_{\mathbb{R}^{d} \setminus \mathcal{B}_{r}} e^{-ipx} \widehat{v}(p) \ dp.$$
(76)

Since $v \in W^{n,1}(\mathbb{R}^d)$, estimate (57) holds.

Using estimates (63), (64), we get:

$$\left| v_1^+(x,r) - \int_{\mathcal{B}_r \setminus Z_{\hat{w}_1,\hat{w}_2}^{\varepsilon}} e^{-ipx} U_{\hat{w}_1,\hat{w}_2}(p,E) \, dp \right|$$

$$(77)$$

$$+v_{+}^{2}(x,r) - \int_{\mathcal{B}_{r} \cap Z_{\hat{w}_{1},\hat{w}_{2}}^{\varepsilon}} e^{-ipx} U_{\hat{w}_{1},\hat{w}_{2}}^{\varepsilon}(p,E) \, dp \, \leq I_{1} + I_{2},$$

$$I_1 = 2^{\beta} c_5 \varepsilon^{-1} E^{-\frac{1}{2}} \int_{\mathcal{B}_r} (1+|p|)^{\beta} dp,$$
(78)

$$I_{2} = c_{6}\varepsilon \int_{\mathcal{B}_{r} \cap Z_{\hat{w}_{1},\hat{w}_{2}}^{\varepsilon}} (1 + \frac{\pi}{|y|} |z(p)| + |p_{\perp}|)^{-n} dp,$$
(79)

$$x \in \mathbb{R}^d, \ 1 \le r \le 2\sqrt{E}, \ E \ge \rho_2,$$

where ρ_2 is the same as in Proposition 2. In addition:

$$I_1 \le c_7 \varepsilon^{-1} E^{\frac{1}{2}} r^{d+\beta}, \quad c_7 = |\mathbb{S}^{d-1}| \frac{2^{d+2\beta}}{d+\beta} c_5; \tag{80}$$

$$I_{2} = c_{6}\varepsilon \sum_{z\in\mathbb{Z}} \int_{\substack{\tau^{2}+p_{\perp}^{2}\leq\tau^{2}\\|\tau-\frac{\pi}{|y|}z|\leq\frac{\varepsilon}{|y|}}} \left(1+\frac{\pi}{|y|}|z|+|p_{\perp}|\right)^{-n}d\tau dp_{\perp},$$

where $\tau \in \mathbb{R}, \, p_{\perp} \in \mathbb{R}^d, \, p_{\perp} y = 0,$

$$I_{2} \leq 2c_{6} \frac{\varepsilon^{2}}{|y|} \sum_{z \in \mathbb{Z}} \int_{\xi \in \mathbb{R}^{d-1}, |\xi| \leq r} \left(1 + \frac{\pi}{|y|} |z| + |\xi|\right)^{-n} d\xi \leq c_{6} c_{8} \varepsilon^{2},$$

$$c_{8} = \frac{2}{|y|} \frac{|\mathbb{S}^{d-2}|}{n-d+1} \sum_{z \in \mathbb{Z}} (1 + \frac{\pi}{|y|} |z|)^{d-n-1}.$$
(81)

In addition, if $r = r_2(E)$, $\varepsilon = \varepsilon_2(E)$, where $r_2(E)$, $\varepsilon_2(E)$ are defined in (44), then $r_2d-n = (2-)d-n r_2-\alpha_2$

$$r^{d-n} = (2\tau)^{d-n} E^{-\alpha_2},$$

$$\varepsilon^{-1} E^{-\frac{1}{2}} r^{d+\beta} = (2\tau)^{d+\beta} E^{-\alpha_2},$$

$$\varepsilon^2 = E^{-\alpha_2}.$$
(82)

Formula (45) follows from representation (76), estimates (57), (77), (80), (81) and formulas (82).

Theorem 2 is proved.

7 Proof of Theorem 3

Proposition 3. Let v satisfy (2) and w_1, \ldots, w_{d+1} be the same as in (24), (29), $d \ge 2$. Then:

$$\begin{aligned} |\widehat{v}(p) - U_{\widehat{w}_{1},...,\widehat{w}_{d+1}}(p,E)| &\leq c_{9}\varepsilon^{-1}(1+|p|)^{\beta}E^{-\frac{1}{2}}, \\ p \in \mathcal{B}_{2\sqrt{E}} \setminus Z_{\widehat{w}_{1},...,\widehat{w}_{d+1}}^{\varepsilon}, E \geq \rho_{3}, \ 0 < \varepsilon < 1, \\ c_{9} &= \frac{\pi\sqrt{d}}{2} \big(2c(D_{0})N_{0}^{3} + c(D_{1})N_{1}^{3} + \max_{j=2,...,d+1} c(D_{j})N_{j}^{3} \big) c_{1}, \\ \rho_{3} &= \max_{j=0,...,d+1} \rho(D_{j},N_{j}), \end{aligned}$$
(83)

in addition, if $v \in W^{n,1}(\mathbb{R}^d)$, $n \ge 0$, then:

$$\begin{aligned} |\widehat{v}(p) - U_{\widehat{w}_{1},...,\widehat{w}_{d+1}}^{\varepsilon}(p,E)| \\ &\leq 2^{\beta}c_{9}\varepsilon^{-1}(1+|p|)^{\beta}E^{-\frac{1}{2}} + 2c_{6}\varepsilon\left(1+2\frac{\pi}{s}\|z(p)\|_{2}\right)^{-n}, \\ p \in \mathcal{B}_{2\sqrt{E}} \cap Z_{\widehat{w}_{1},...,\widehat{w}_{d+1}}^{\varepsilon}, E \geq \rho_{3}, \ 0 < \varepsilon < \min\{1,\frac{1}{2}s\}, \end{aligned}$$
(84)

where $c, \rho, D_j, N_j, j = 0, ..., d+1$, are defined as in (15); c_1, β are the same as in Lemma 1, c_6 is the same as in Proposition 2, z(p) is defined in (47) and $||z(p)||_2$ is the standard Euclidean norm of z(p).

Proof of Proposition 3. In a similar way with formulas (65), one can write

$$M_{\widehat{w}_{1},\widehat{w}_{i'}}(p) = \widehat{w}_{1}(p) \begin{pmatrix} 1 & 0\\ \cos(sp_{i'-1}) & \sin(sp_{i'-1}) \end{pmatrix}, \quad p \in \mathbb{R}^{d} \setminus \frac{\pi}{s} \mathbb{Z}^{d},$$

$$M_{\widehat{w}_{1},\widehat{w}_{i'}}^{-1}(p) = \frac{1}{\sin(sp_{i'-1})\widehat{w}_{1}(p)} \begin{pmatrix} \sin(sp_{i'-1}) & 0\\ -\cos(sp_{i'-1}) & 1 \end{pmatrix}, \quad p \in \mathbb{R}^{d} \setminus Z_{\widehat{w}_{1},...,\widehat{w}_{d+1}}^{\varepsilon},$$
(85)

where i' = i'(p, s) is defined in (48). Also note that

,

$$|\sin(sp_{i'-1})| \ge \frac{2\varepsilon}{\pi\sqrt{d}}, \quad p \in \mathbb{R}^d \setminus Z^{\varepsilon}_{\widehat{w}_1,\dots,\widehat{w}_{d+1}}, \ 0 < \varepsilon < 1.$$
(86)

Estimate (83) follows from (11), (15), (23), (34), (37), (38), (48), (49) and from (85), (86).

It remains to prove (84). Using definition (50), we represent

$$U_{\widehat{w}_{1},...,\widehat{w}_{d+1}}^{\varepsilon}(p,E) - \widehat{v}(p) = \varphi_{1}^{\varepsilon}(p,E) + \varphi_{2}^{\varepsilon}(p), \quad p \in \mathcal{B}_{2\sqrt{E}} \cap Z_{\widehat{w}_{1},...,\widehat{w}_{d+1}}^{\varepsilon}$$
$$\varphi_{1}^{\varepsilon}(p,E) = \frac{1}{|\mathbb{S}^{d-1}|} \int_{\mathbb{S}^{d-1}} \left(U_{\widehat{w}_{1},...,\widehat{w}_{d+1}}(\eta,E) - \widehat{v}(\eta) \right) \Big|_{\eta = \frac{\varepsilon}{s}\vartheta + \frac{\pi}{s}z(p)} d\vartheta, \qquad (87)$$
$$\varphi_{2}^{\varepsilon}(p) = \frac{1}{|\mathbb{S}^{d-1}|} \int_{\mathbb{S}^{d-1}} \left(\widehat{v} \left(\frac{\varepsilon}{s}\vartheta + \frac{\pi}{s}z(p) \right) - \widehat{v}(p) \right) d\vartheta,$$

where z(p) is defined in (47).

Using formulas (83), (87), we obtain

$$\begin{aligned} |\varphi_1^{\varepsilon}(p,E)| &\leq c_9 \varepsilon^{-1} E^{-\frac{1}{2}} \frac{1}{|\mathbb{S}^{d-1}|} \int_{\mathbb{S}^{d-1}} \left(1 + \left|\frac{\varepsilon}{s}\vartheta + \frac{\pi}{s}z(p)\right|\right)^{\beta} d\vartheta \\ &\leq c_9 \varepsilon^{-1} E^{-\frac{1}{2}} \frac{1}{|\mathbb{S}^{d-1}|} \int_{\mathbb{S}^{d-1}} \left(1 + |p| + 2\frac{\varepsilon}{s}\right)^{\beta} d\vartheta \leq 2^{\beta} c_9 \varepsilon^{-1} (1 + |p|)^{\beta} E^{-\frac{1}{2}}, \quad (88) \\ & \text{for } \varepsilon \text{ as in } (84). \end{aligned}$$

Next, using the definition of φ_2^{ε} in formula (87) and the mean value theorem, we get the following estimate:

$$|\varphi_2^{\varepsilon}(p)| \le 2\frac{\varepsilon}{s} \max\{|\nabla \widehat{v}(\xi)| \colon \xi \in \mathbb{R}^d, \ |\xi - \frac{\pi}{s}z(p)| \le \frac{\varepsilon}{s}\}, \ p \in Z_{\widehat{w}_1, \dots, \widehat{w}_{d+1}}^{\varepsilon}.$$
 (89)

Here, the mean value theorem was used for $\hat{v}(\xi)$ on $[p, \frac{\varepsilon}{s}\vartheta + \frac{\pi}{s}z(p)], \vartheta \in \mathbb{S}^{d-1}$. One can see that

estimates (70) and (71) hold for all
$$\xi \in \mathbb{R}^d$$

such that $|\xi - \frac{\pi}{s} z(p)| \le \frac{\varepsilon}{s}$, where $p \in Z^{\varepsilon}_{\widehat{w}_1, \dots, \widehat{w}_{d+1}}$. (90)

It follows from (89), (90) and from the upper estimate on ε of (84), that

$$\begin{aligned} |\varphi_{2}^{\varepsilon}(p)| &\leq 2^{1-n} c_{6} \varepsilon \max\left\{ (1+|\xi|)^{-n} \colon |\xi - \frac{\pi}{s} z(p)| \leq \frac{\varepsilon}{s} \right\} \\ &\leq 2^{1-n} c_{6} \varepsilon \left(1 + \frac{\pi}{s} ||z(p)||_{2} - \frac{\varepsilon}{s} \right)^{-n} \\ &\leq 2 c_{6} \varepsilon \left(1 + 2\frac{\pi}{s} ||z(p)||_{2} \right)^{-n}, \ p \in Z_{\widehat{w}_{1}, \dots, \widehat{w}_{d+1}}^{\varepsilon}. \end{aligned}$$
(91)

Estimate (84) follows from estimates (88) and (91).

Proposition 3 is proved.

The final part of the proof of Theorem 3 is as follows. In a similar way with (76), we represent v as follows:

$$v(x) = v_{1}^{+}(x,r) + v_{2}^{+}(x,r) + v^{-}(x,r), \quad x \in \mathbb{R}^{d}, \ r > 0,$$

$$v_{1}^{+}(x,r) = \int_{\mathcal{B}_{r} \setminus Z_{\hat{w}_{1},...,\hat{w}_{d+1}}^{\varepsilon}} e^{-ipx} \hat{v}(p) \, dp,$$

$$v_{2}^{+}(x,r) = \int_{\mathcal{B}_{r} \cap Z_{\hat{w}_{1},...,\hat{w}_{d+1}}^{\varepsilon}} e^{-ipx} \hat{v}(p) \, dp,$$

$$v^{-}(x,r) = \int_{\mathbb{R}^{d} \setminus \mathcal{B}_{r}} e^{-ipx} \hat{v}(p) \, dp.$$
(92)

Since v belongs to $W^{n,1}(\mathbb{R}^d)$, estimate (57) is valid.

Using estimates (83), (84) we obtain

$$\begin{vmatrix}
v_{1}^{+}(x,r) - \int_{\mathcal{B}_{r} \setminus Z_{\hat{w}_{1},...,\hat{w}_{d+1}}} e^{-ipx} U_{\hat{w}_{1},...,\hat{w}_{d+1}}(p,E) \, dp \\
+ v_{2}^{+}(x,r) - \int_{\mathcal{B}_{r} \cap Z_{\hat{w}_{1},...,\hat{w}_{d+1}}} e^{-ipx} U_{\hat{w}_{1},...,\hat{w}_{d+1}}^{\varepsilon}(p,E) \, dp \\
\int_{1} = 2^{\beta} c_{9} \varepsilon^{-1} E^{-\frac{1}{2}} \int_{\mathcal{B}_{r}} (1+|p|)^{\beta} dp, \\
J_{2} = 2c_{6} \varepsilon \int_{\mathcal{B}_{r} \cap Z_{\hat{w}_{1},...,\hat{w}_{d+1}}} (1+2\frac{\pi}{s} \|z(p)\|_{2})^{-n} dp, \\
x \in \mathbb{R}^{d}, \ 1 \leq r \leq 2\sqrt{E}, \ E \geq \rho_{3},
\end{cases}$$
(93)

where ρ_3 is the same as in Proposition 3. In addition,

$$J_{1} \leq c_{10}\varepsilon^{-1}E^{-\frac{1}{2}}r^{d+\beta}, \quad c_{10} = |\mathbb{S}^{d-1}|\frac{2^{d+\beta}}{d+\beta}c_{9},$$

$$J_{2} \leq c_{11}\varepsilon^{d+1}, \quad c_{11} = \left(\frac{1}{s}\right)^{d}|\mathcal{B}_{1}|\sum_{z\in\mathbb{Z}^{d}}\left(1+2\frac{\pi}{s}||z||_{2}\right)^{-n},$$
(94)

where $|\mathcal{B}_1|$ is the standard Euclidean volume of \mathcal{B}_1 . Finally, if $r = r_3(E)$, $\varepsilon = \varepsilon_3(E)$, where $r_3(E)$, $\varepsilon_3(E)$ are defined in (52), then

$$r^{d-n} = (2\tau)^{d-n} E^{-\alpha_3},$$

$$\varepsilon^{-1} E^{-\frac{1}{2}} r^{d+\beta} = (2\tau)^{d+\beta} E^{-\alpha_3},$$

$$\varepsilon^{d+1} = E^{-\alpha_3}.$$
(95)

Formula (53) follows from representation (92), estimates (57), (93), (94) and formulas (95).

Theorem 3 is proved.

8 Proof of Lemma 1

Note that

$$\widehat{w}(p) = \int_{\mathbb{R}^d} |\widehat{q}(\xi)|^2 \widehat{\omega}_{\nu}(p-\xi) \, d\xi, \quad p \in \mathbb{R}^d, \tag{96}$$

$$\omega_{\nu}(x) = |x|^{\frac{\nu}{2}} K_{\nu}(|x|), \quad x \in \mathbb{R}^d,$$
(97)

where $\hat{q}, \hat{\omega}_{\nu}$ are the Fourier transforms of q, ω_{ν} . The Fourier transform $\hat{\omega}_{\nu}$ can be computed explicitly:

$$\widehat{\omega}_{\nu}(p) = \frac{c_{12}}{(1+|p|^2)^{\frac{d}{2}+\nu}}, \quad c_{12} = \frac{\Gamma(\frac{d}{2}+\nu)2^{\nu-1}}{\pi^{\frac{d}{2}}}.$$
(98)

Indeed, formula (98) follows from the Fourier inversion theorem and the following computations:

$$\begin{split} \int_{\mathbb{R}^d} \frac{e^{-ipx} \, dp}{(1+|p|^2)^{\frac{d}{2}+\nu}} &= \int_{\mathbb{R}} \int_{\mathbb{R}^{d-1}} \frac{e^{-i|x|t} \, dt \, d\xi}{(1+t^2+|\xi|^2)^{\frac{d}{2}+\nu}} \\ &= |\mathbb{S}^{d-2}| \int_{\mathbb{R}} \int_{\mathbb{R}} \frac{e^{-i|x|t} r^{d-2} \, dt \, dr}{(1+t^2+r^2)^{\frac{d}{2}+\nu}} \\ \frac{r=\sqrt{1+t^2\tau}}{|\mathbb{S}^{d-2}|} \int_{\mathbb{R}} \frac{e^{-i|x|t} \, dt}{(1+t^2)^{\frac{1}{2}+\nu}} \int_{0}^{+\infty} \frac{\tau^{d-2} \, d\tau}{(1+\tau^2)^{\frac{d}{2}+\nu}} \\ &= c_{12}^{-1} |x|^{\nu} K_{\nu}(|x|), \ x \in \mathbb{R}^d. \end{split}$$

Here, it was used that

$$\int_{0}^{+\infty} \frac{\tau^{d-2} d\tau}{(1+\tau^2)^{\frac{d}{2}+\nu}} = \frac{1}{2} B(\frac{d-1}{2}, \nu + \frac{1}{2}) = \frac{1}{2} \frac{\Gamma(\frac{d-1}{2})\Gamma(\nu + \frac{1}{2})}{\Gamma(\frac{d}{2}+\nu)},$$
$$|\mathbb{S}^{d-2}| = \frac{2\pi^{\frac{d-1}{2}}}{\Gamma(\frac{d-1}{2})},$$

where B and Γ denote the beta and gamma functions.

Using (96), (98), we obtain the estimates

$$\widehat{w}(p) \ge \int_{|\xi| \le 1} \frac{c_{12} |\widehat{q}(\xi)|^2}{(1+|p-\xi|)^{d+2\nu}} d\xi \ge \frac{c_1(q,\nu)}{(1+|p|)^{d+2\nu}}, \quad p \in \mathbb{R}^d,$$

$$c_1(q,\nu) = \frac{c_{12}}{2^{d+2\nu}} \int_{|\xi| \le 1} |\widehat{q}(\xi)|^2 d\xi.$$
(99)

Properties (23) follow from (20), (22), (96) and (99). Lemma 1 is proved.

References

- [1] A. D. Agaltsov, T. Hohage, and R. G. Novikov, "An iterative approach to monochromatic phaseless inverse scattering," in preparation.
- [2] T. Aktosun and P. E. Sacks, "Inverse problems on the line without phase information," *Inverse Problems*, vol. 14, pp. 211–224, 1998.
- [3] N. V. Alexeenko, V. A. Burov, and O. D. Rumyantseva, "Solution of the three-dimensional acoustical inverse scattering problem. The modified Novikov algorithm," *Acoustical Physics*, vol. 54, no. 3, pp. 407–419, 2008.
- [4] J. A. Barceló, C. Castro, and J. M. Reyes, "Numerical approximation of the potential in the two-dimesional inverse scattering problem," *Inverse Problems*, vol. 32, no. 1, 2016, 015006 (19pp).

- [5] F. A. Berezin and M. A. Shubin, *The Schrödinger equation*, ser. Mathematics and its applications. Dordrecht: Kluwer Academic, 1991, vol. 66.
- [6] V. A. Burov, N. V. Alekseenko, and O. D. Rumyantseva, "Multifrequency generalization of the Novikov algorithm for the two-dimensional inverse scattering problem," *Acoustical Physics*, vol. 56, no. 6, pp. 843–856, 2009.
- [7] K. Chadan and P. C. Sabatier, *Inverse problems in quantum scattering theory*, 2nd ed. Berlin: Springer, 1989.
- [8] G. Eskin, *Lectures on linear partial differential equations*, ser. Graduate studies in mathematics. American mathematical society, 2011, vol. 123.
- [9] L. D. Faddeev, "Uniqueness of the solution of the inverse scattering problem," Vest. Leningrad Univ., vol. 7, pp. 126–130, 1956, (in Russian).
- [10] —, "Inverse problem of quantum scattering theory. II." Journal of Soviet Mathematics, vol. 5, no. 3, pp. 334–396, 1976.
- [11] L. D. Faddeev and S. P. Merkuriev, Quantum scattering theory for several particle systems, ser. Mathematical physics and applied mathematics. Dordrecht: Kluwer Academic, 1993, vol. 11.
- [12] P. G. Grinevich, "The scattering transform for the two-dimensional Schrödinger operator with a potential that decreases at infinity at fixed nonzero energy," *Russian Math. Surveys*, vol. 55, no. 6, pp. 1015–1083, 2000.
- [13] P. Hähner and T. Hohage, "New stability estimates for the inverse acoustic inhomogeneous medium problem and applications," SIAM J. Math. Anal., vol. 33, no. 3, pp. 670–685, 2001.
- [14] M. I. Isayev, "Exponential instability in the inverse scattering problem on the energy interval," *Funct. Anal. Appl.*, vol. 47, no. 3, pp. 187–194, 2013.
- [15] M. I. Isayev and R. G. Novikov, "New global stability estimates for monochromatic inverse acoustic scattering," SIAM J. Math. Analysis, vol. 43, no. 5, pp. 1495–1504, 2013.
- [16] M. V. Klibanov, "Phaseless inverse scattering problems in three dimensions," SIAM J. Appl. Math., vol. 74, pp. 392–410, 2014.
- [17] M. V. Klibanov and V. G. Romanov, "Reconstruction procedures for two inverse scattering problems without the phase information," SIAM J. Appl. Math., vol. 76, no. 1, pp. 178–196, 2016.
- [18] —, "Two reconstruction procedures for a 3D phaseless inverse scattering problem for the generalized Helmholtz equation," *Inverse Problems*, vol. 32, no. 1, 2016, 015005 (16pp).

- [19] R. G. Novikov, "Multidimensional inverse spectral problem for the equation $-\Delta \psi + (v(x) Eu(x))\psi = 0$," Funct. Anal. Appl., vol. 22, pp. 263–272, 1988.
- [20] —, "Rapidly converging approximation in inverse quantum scattering in dimension 2," *Physics Letters A*, vol. 238, pp. 73–78, 1998.
- [21] —, "The ∂-approach to approximate inverse scattering at fixed energy in three dimensions," Int. Math. Res. Pap., no. 6, pp. 287–349, 2005.
- [22] —, "Approximate Lipschitz stability for non-overdetermined inverse scattering at fixed energy," J. Inverse Ill-Posed Probl., vol. 26, no. 6, pp. 813– 823, 2013.
- [23] —, "Inverse scattering without phase information," in Séminaire Laurent Schwartz — EDP et applications. Ecole Polytechnique, 2014-2015, exp. no. 16 (13pp).
- [24] —, "Formulas for phase recovering from phaseless scattering data at fixed frequency," Bull. Sci. Math., vol. 139, no. 8, pp. 923–936, 2015.
- [25] —, "An iterative approach to non-overdetermined inverse scattering at fixed energy," *Sbornik: Mathematics*, vol. 206, no. 1, pp. 120–134, 2015.
- [26] —, "Explicit formulas and global uniqueness for phaseless inverse scattering in multidimensions," *The Journal of Geometric Analysis*, vol. 26, no. 1, pp. 346–359, 2016.