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**MAXIMAL SURFACES IN ANTI-DE SITTER SPACE, WIDTH OF
CONVEX HULLS AND QUASICONFORMAL EXTENSIONS OF
QUASISYMMETRIC HOMEOMORPHISMS**

ANDREA SEPPI

ABSTRACT. We give upper bounds on the principal curvatures of a maximal surface of nonpositive curvature in three-dimensional Anti-de Sitter space, which only depend on the width of the convex hull of the surface. Moreover, given a quasisymmetric homeomorphism ϕ , we study the relation between the width of the convex hull of the graph of ϕ , as a curve in the boundary of infinity of Anti-de Sitter space, and the cross-ratio norm of ϕ .

As an application, we prove that if ϕ is a quasisymmetric homeomorphism of \mathbb{RP}^1 with cross-ratio norm $\|\phi\|$, then $\ln K \leq C\|\phi\|$, where K is the maximal dilatation of the minimal Lagrangian extension of ϕ to the hyperbolic plane.

INTRODUCTION

The study of three-dimensional Anti-de Sitter space AdS^3 was initiated by the pioneering work of Mess ([Mes07]) of 1990, and has been widely developed since then, with emphasis on its relation with Teichmüller theory, for instance in [AAW00, ABB⁺07, BBZ07, BS10, BKS11, BS12, KS07, BST17].

In particular, in [BS10] Bonsante and Schlenker studied zero mean curvature spacelike surfaces - hereafter *maximal* surfaces - with boundary contained in the boundary at infinity $\partial_\infty \text{AdS}^3$. The latter is identified in a natural way to $\mathbb{RP}^1 \times \mathbb{RP}^1$, where isometries of AdS^3 extend to projective transformations which act as elements of $\text{PSL}(2, \mathbb{R}) \times \text{PSL}(2, \mathbb{R})$. Therefore the asymptotic boundary of a maximal surface is represented by the graph of an orientation-preserving homeomorphism $\phi : \mathbb{RP}^1 \rightarrow \mathbb{RP}^1$. Bonsante and Schlenker proved that every curve in $\partial_\infty \text{AdS}^3$ corresponding to the graph of an orientation-preserving homeomorphism $\phi : \mathbb{RP}^1 \rightarrow \mathbb{RP}^1$ bounds a maximal disc S with nonpositive curvature. This result might be thought of as an asymptotic Plateau problem in Anti-de Sitter geometry. See also [BS16] for an alternative proof. By the Gauss equation in AdS^3 , nonpositivity of curvature is equivalent to the condition that the principal curvatures of S are in $[-1, 1]$.

Bonsante and Schlenker also provided a more precise description of maximal discs under the assumption that ϕ is a quasisymmetric homeomorphism of \mathbb{RP}^1 - namely, if the *cross-ratio norm*

$$\|\phi\|_{cr} = \sup_{cr(Q)=-1} |\ln |cr(\phi(Q))||$$

is finite. In this case, the maximal disc is unique and the principal curvatures are in $[-1 + \epsilon, 1 - \epsilon]$ for some ϵ . By means of a construction which associates to a maximal surface with nonpositive curvature a minimal Lagrangian diffeomorphism from \mathbb{H}^2 to \mathbb{H}^2 (a diffeomorphism of \mathbb{H}^2 is minimal Lagrangian if it is area-preserving and its graph is a minimal surface in $\mathbb{H}^2 \times \mathbb{H}^2$), the existence and uniqueness theorems on maximal surfaces led to the proof of the fact that every quasisymmetric homeomorphism of \mathbb{RP}^1 admits a unique quasiconformal minimal Lagrangian extension to \mathbb{H}^2 .

Another important ingredient introduced in [BS10] is the *width* of the convex hull. This is defined as the supremum of the length of timelike paths contained in the convex hull of the curve $gr(\phi)$. By a simple application of the maximum principle, the maximal surface

S with $\partial_\infty S = gr(\phi)$ is itself contained in the convex hull. Bonsante and Schlenker proved that for every orientation-preserving homeomorphism ϕ , the width is at most $\pi/2$, and it is strictly less than $\pi/2$ precisely when ϕ is quasimetric.

The purpose of this paper is to study the quantitative relations between the cross-ratio norm of ϕ , the width w of its convex hull, and the supremum $\|\lambda\|_\infty$ of the principal curvatures of the maximal surface S of nonpositive curvature such that $\partial_\infty S = gr(\phi)$. By the above discussion, $\|\phi\|_{cr} < +\infty$ if and only if $w < \pi/2$ if and only if $\|\lambda\|_\infty < 1$, but it is not clear whether there is a direct relation between these quantities. Using a formula proved in [KS07] which relates the differential of the minimal Lagrangian extension to the shape operator of S , our results will provide estimates on the maximal dilatation of the quasiconformal minimal Lagrangian extension, only depending on the cross-ratio norm of ϕ .

Principal curvatures of maximal surfaces. The study of the relation between the principal curvatures of a maximal surface and the width of the convex hull is split into two parts. Observe that the principal curvatures of S vanish identically when S is a totally geodesic plane, in which case the width is zero since the convex hull consists of S itself. Our first theorem describes the behavior of maximal surfaces which are close to being a totally geodesic plane:

Theorem 1.A. *There exists a constant C_1 such that, for every maximal surface S with $\|\lambda\|_\infty < 1$ and width w ,*

$$\|\lambda\|_\infty \leq C_1 \tan w.$$

This theorem provides interesting information only when w is in some neighborhood of zero, since for large w the already known bound $\|\lambda\|_\infty < 1$ is not improved. On the other hand, Bonsante and Schlenker showed that if a maximal surface of nonpositive curvature has a point where the principal curvatures are -1 and 1 , then the principal curvatures are -1 and 1 everywhere, and therefore the induced metric is flat. Moreover, the surface is a so-called *horospherical surface*, which is described explicitly and has width $\pi/2$. Our second theorem concerns surfaces which are close to this situation:

Theorem 1.B. *There exist universal constants $M > 0$ and $\delta \in (0, 1)$ such that, if S is a maximal surface in AdS^3 with $\delta \leq \|\lambda\|_\infty < 1$ and width w , then*

$$\tan w \geq \left(\frac{1}{1 - \|\lambda\|_\infty} \right)^{1/M}.$$

It is worth remarking here that an inequality going in the opposite direction can be obtained more easily, and all the necessary tools were already proved in [KS07] and [BS10]. Nevertheless, for the sake of completeness we will provide a proof of the following:

Proposition 1.C. *Let S be a maximal surface in AdS^3 with $\|\lambda\|_\infty \leq 1$ and width w . Then*

$$\tan w \leq \frac{2\|\lambda\|_\infty}{1 - \|\lambda\|_\infty^2}.$$

Since $2\|\lambda\|_\infty/(1 - \|\lambda\|_\infty^2)$ behaves like $2\|\lambda\|_\infty$ as $\|\lambda\|_\infty \rightarrow 0$, one sees that Theorem 1.A is optimal for small $\|\lambda\|_\infty$, up to determining the best possible value of the constant C_1 . On the other hand, from Proposition 1.C one obtains that $\tan w \leq 2/(1 - \|\lambda\|_\infty)$, and it remains an open question whether Theorem 1.B can be improved to an inequality of the form $\tan w \geq C_2^{-1}(1 - \|\lambda\|_\infty)^{-1}$.

Minimal Lagrangian extensions. A classical problem in Teichmüller theory concerns quasiconformal extensions to the disc of quasimetric homeomorphisms of the circle. Classical quasiconformal extensions include, for instance, the Beurling-Ahlfors extension and the Douady-Earle extension. More recently, Markovic [Mar17] proved the existence of quasiconformal harmonic extensions, where the harmonicity is referred to the complete hyperbolic metric of \mathbb{H}^2 .

Moreover, the maximal dilatation of the classical extensions has been widely studied. For instance, Beurling and Ahlfors in [BA56] proved that, if Φ_{BA} is the Beurling-Ahlfors extension of a quasymmetric homeomorphism ϕ , then the maximal dilatation $K(\Phi_{BA})$ satisfies:

$$\ln K(\Phi_{BA}) \leq 2\|\phi\|_{cr}.$$

The asymptotic behaviour was later improved in [Leh83] by

$$\ln K(\Phi_{BA}) \leq \|\phi\|_{cr} + \ln 2.$$

For the Douady-Earle extension, [DE86] proved that there exist constants δ and C such that, for every quasymmetric homeomorphism of the circle ϕ with $\|\phi\|_{cr} < \delta$, the Douady-Earle extension Φ_{DE} satisfies:

$$\ln K(\Phi_{DE}) \leq C\|\phi\|_{cr}.$$

More recently, Hu and Muzician proved in [HM12] that the following always holds:

$$\ln K(\Phi_{DE}) \leq C_1\|\phi\|_{cr} + C_2.$$

In this paper we will prove analogous results for the minimal Lagrangian extension, whose existence was proved in [BS10] as already remarked. As an application of Theorem 1.A, we will prove the following inequality:

Theorem 2.A. *There exist universal constants δ and C_1 such that, for any quasymmetric homeomorphism ϕ of \mathbb{RP}^1 with cross ratio norm $\|\phi\|_{cr} < \delta$, the minimal Lagrangian extension $\Phi_{ML} : \mathbb{H}^2 \rightarrow \mathbb{H}^2$ has maximal dilatation bounded by:*

$$\ln K(\Phi_{ML}) \leq C_1\|\phi\|_{cr}.$$

On the other hand, by an application of Theorem 1.B we will derive an asymptotic estimate of the maximal dilatation of Φ_{ML} :

Theorem 2.B. *There exist universal constants Δ and C_2 such that, for any quasymmetric homeomorphism ϕ of \mathbb{RP}^1 with cross ratio norm $\|\phi\|_{cr} > \Delta$, the minimal Lagrangian extension $\Phi_{ML} : \mathbb{H}^2 \rightarrow \mathbb{H}^2$ has maximal dilatation bounded by:*

$$\ln K(\Phi_{ML}) \leq C_2\|\phi\|_{cr}.$$

Using Proposition 1.C we also obtain an inequality in the converse direction, which holds for quasymmetric homeomorphisms with small cross-ratio norm and shows that Theorem 2.A is not improvable from a qualitative point of view.

Theorem 2.C. *There exist universal constants δ and C_0 such that, for any quasymmetric homeomorphism ϕ of \mathbb{RP}^1 with cross ratio norm $\|\phi\|_{cr} < \delta$, the minimal Lagrangian extension $\Phi : \mathbb{H}^2 \rightarrow \mathbb{H}^2$ has maximal dilatation bounded by:*

$$C_0\|\phi\|_{cr} \leq \ln K(\Phi_{ML}).$$

The constant C_0 can be taken arbitrarily close to $1/2$.

Finally, from Theorem 2.A and Theorem 2.B we will derive the following corollary.

Corollary 2.D. *There exists a universal constant C such that, for any quasymmetric homeomorphism ϕ of \mathbb{RP}^1 , the minimal Lagrangian extension $\Phi_{ML} : \mathbb{H}^2 \rightarrow \mathbb{H}^2$ has maximal dilatation $K(\Phi_{ML})$ bounded by:*

$$\ln K(\Phi_{ML}) \leq C\|\phi\|_{cr}.$$

Corollary 2.D is therefore a result for minimal Lagrangian extensions comparable to what has been proved for Beurling-Ahlfors and Douady-Earle extensions.

From the width to the cross-ratio norm. The bridge from Theorem 1.A to Theorem 2.A, and from Theorem 1.B to Theorem 2.B, is twofold. The first aspect is a direct relation between the principal curvatures of a maximal surface S and the quasiconformal dilatation of the minimal Lagrangian extension at the corresponding point. This is proved in Proposition 5.5 by using a formula of [KS07], and has as a consequence that:

$$K(\Phi_{ML}) = \left(\frac{1 + \|\lambda\|_\infty}{1 - \|\lambda\|_\infty} \right)^2.$$

On the other hand, the step from the width to the cross-ratio norm is more subtle. This is the content of the following proposition:

Proposition 3.A. *Given any quasisymmetric homeomorphism ϕ of \mathbb{RP}^1 , let w be the width of the convex hull of the graph of ϕ in $\partial_\infty \text{AdS}^3$. Then*

$$\tan w \leq \sinh \left(\frac{\|\phi\|_{cr}}{2} \right).$$

By means of these two relations and some computations, Theorem 2.A and Theorem 2.B are proved on the base of Theorem 1.A and Theorem 1.B.

To prove Proposition 3.A, assuming that the width is w , we will essentially find two support planes P_- and P_+ for the convex hull of $gr(\phi)$, on the two different sides of the convex hull, such that P_- and P_+ are connected by a timelike geodesic segment of length w . We will use the fact that the boundaries of the convex hull are pleated surfaces in order to pick four points in $\partial_\infty \text{AdS}^3$ - two in the boundary at infinity $\partial_\infty P_-$ and the other two in $\partial_\infty P_+$ - and use such four points to show that the cross-ratio norm of ϕ is large. Turning this qualitative picture into quantitative estimates, leading to the proof of Proposition 3.A, involves careful and somehow technical constructions in Anti-de Sitter space.

By using similar techniques, we will also prove an inequality in the converse direction, which is the content of the following proposition.

Proposition 3.B. *Given any quasisymmetric homeomorphism ϕ of \mathbb{RP}^1 , let w the width of the convex hull of the graph of ϕ in $\partial_\infty \text{AdS}^3$. Then*

$$\tanh \left(\frac{\|\phi\|_{cr}}{4} \right) \leq \tan w.$$

This inequality, however, is clearly not optimal, as the hyperbolic tangent tends to 1 as $\|\phi\|_{cr}$ tends to infinity. Hence the inequality is interesting only for $w < \pi/4$. Nevertheless, this inequality is used to obtain Theorem 2.C from Proposition 1.C. To prove Proposition 3.B, we will assume the cross-ratio norm is $\|\phi\|_{cr}$ and - composing with Möbius transformations in an appropriate way - construct a quadruple points in $\partial_\infty \text{AdS}^3$. Then we consider two spacelike lines connecting two pairs of points at infinity chosen in the above quadruple. By construction, those two lines are contained in the convex hull of $gr(\phi)$, hence the maximal length of a timelike geodesic segment between them provides a bound from below on the width.

Outline of the main proofs. Let us now give an outline of some technicalities involved in the proofs of Theorem 1.A and Theorem 1.B.

The starting point behind the proof of Theorem 1.A is the fact that a maximal surface S with $\partial_\infty S = gr(\phi)$ in $\partial_\infty \text{AdS}^3$ is contained in the convex hull of $gr(\phi)$. Using this fact, for every point $x \in S$ we find two timelike geodesic segments starting from x and orthogonal to two planes P_-, P_+ which do not intersect the convex hull of $gr(\phi)$. The sum of the lengths of the two segments is less than the width w . Moreover S is contained in the region bounded by P_- and P_+ .

Now the key step is to show that, heuristically, if S is contained in the region between two disjoint planes which are close to x , then the principal curvatures of S in a neighborhood of x

cannot be too large. To make this statement precise, we will apply Schauder-type estimates to the linear equation

$$(L) \quad \Delta_S u - 2u = 0,$$

where $u : S \rightarrow \mathbb{R}$ is the function which measures the sine of the (signed, timelike) distance from the plane P_- , and Δ_S is the Laplace-Beltrami operator of S (negative definite as an operator on L^2). Observe that an easy application of the maximum principle to Equation (L) proves that a maximal surface is necessarily contained in the convex hull. A more subtle study of *a priori* bounds for this equation provides the key step for Theorem 1.A.

A technical point is that the operator $\Delta_S - 2\text{id}$ depends on the maximal surface, which will be overcome by using the uniform boundedness of the coefficients, written in normal coordinates, for a class of surfaces we are interested in. The precise statement we will use is the following:

Proposition 4.A. *There exists a radius $R > 0$ and a constant $C > 0$ such that for every choice of:*

- *A maximal surface S of nonpositive curvature in AdS^3 with $\partial_\infty S$ the graph of an orientation-preserving homeomorphism;*
- *A point $x \in S$;*
- *A plane P_- disjoint from S with $d_{\text{AdS}^3}(x, P_-) \leq \pi/4$,*

the function $u(z) = \sin d_{\text{AdS}^3}(\exp_x(z), P_-)$ satisfies the Schauder-type inequality

$$\|u\|_{C^2(B(0, \frac{\pi}{2}))} \leq C \|u\|_{C^0(B(0, R))}.$$

The techniques involved here are similar to those used, in the case of minimal surfaces in three-dimensional hyperbolic geometry, in [Sep16].

To conclude, we then use an explicit expression for the shape operator of the maximal surface S in terms of the value of u , the first derivatives of u , and the second derivatives of u . Hence, using Proposition 4.A, the principal curvatures are bounded in terms of the supremum of u on a geodesic ball $B_S(x, R)$. The latter can be estimated in terms of the width w . However, in this last step it is necessary to control the size of the image of $B_S(x, R)$ under the projection to the plane P_- . To achieve this, a uniform gradient lemma is proved, to show that the maximal surface S is not too “tilted” with respect to P_- .

Similarly, the key analytical point for the proof of Theorem 1.B comes from an *a priori* estimate. Consider the function $\chi : S \rightarrow [0, +\infty)$ defined by $\chi = -\ln \lambda$, where λ is the positive principal curvature of S . It turns out that χ satisfies the equation

$$(Q) \quad \Delta_S \chi = 2(1 - e^{-2\chi}).$$

By an application of the maximum principle, one can prove that if a maximal surface of nonpositive curvature has principal curvatures -1 and 1 at some point, then the principal curvatures are identically -1 and 1 . By a careful analysis of Equation (Q), we will prove a more quantitative result, which roughly speaking shows that if the principal curvatures are “large” (i.e. close to 1) at some point, then they remain “large” on a “large” ball.

Proposition 4.B. *There exists a universal constant M such that, for every maximal surface S of nonpositive curvature in AdS^3 and every pair of points $p, q \in S$,*

$$1 - \lambda(q) \leq e^{Md_S(p, q)}(1 - \lambda(p)).$$

The proof of Proposition 4.B is based on some estimates already proved jointly by Francesco Bonsante, Jean-Marc Schlenker and Mike Wolf ([BSW]) in an unpublished work. The proof presented in this paper closely follows their arguments, which I was kindly transmitted and authorized to adapt to the purpose of this paper.

The strategy to prove Theorem 1.B is then the following. Assume there exists a point x_0 with $\lambda(x_0) = 1 - e^{-v(x_0)}$ very close to 1 . We want to show that the width w is very close to

$\pi/2$. We will first show that the line of curvature of S corresponding to the positive eigenvalue of S remains, for a certain amount of time, in the concave side of an umbilical surface U_{λ_1} tangent to S at x_0 , whose principal curvatures are both equal to $\lambda_1 = 1 - e^{-v(x_0)/4}$. Such amount of time is finite (in a unit-speed parameterization of the line of curvature), but it can be arranged to tend to infinity as $\lambda(x_0)$ tends to 1. This step is basically a maximum principle argument, but requires a technical point to show that the intrinsic acceleration of the line of curvature is also small, i.e. comparable to $e^{-v(x_0)}$.

The surface U_{λ_1} is obtained as the surface at constant timelike distance d_1 from a totally geodesic plane. As λ_1 tends to 1, d_1 tends to $\pi/4$. By following the line of curvature corresponding to the positive eigenvalue, in the two opposite directions from x_0 , for a time as indicated in the above paragraph, we obtain two points p_1 and p_2 . Moreover p_1 and p_2 converge to $\partial_\infty \text{AdS}^3$ as $\lambda(x_0) \rightarrow 1$. Of course analogous statements hold for the two points q_1, q_2 obtained by following the line of curvature from x_0 corresponding to the negative eigenvalue.

After proving quantitative versions of the above statements, we can give a lower bound for the length of the timelike geodesic segment $\overline{p_1 p_2}$ which maximizes the distance between the geodesic segments $\overline{p_1 p_2}$ and $\overline{q_1 q_2}$. Since $\overline{p_1 p_2}$ and $\overline{q_1 q_2}$ are contained in the convex hull of S , also $\overline{p_1 p_2}$ is contained in the convex hull, and therefore the lower bound on the length of $\overline{p_1 p_2}$ provides a lower bound for the width w , which only depends on $v(x_0) = -\ln(1 - \lambda(x_0))$. The reason why the obtained estimate is efficient when $\lambda(x_0)$ approaches 1 is that, in the limit configuration, the lines $\overline{p_1 p_2}$ and $\overline{q_1 q_2}$ tend to be in dual position: equivalently, in the limit every point of the first line is connected to every point of the second line by a geodesic timelike segment of length $\pi/2$. Thus as $\lambda(x_0) \rightarrow 1$, the surface S is approaching a horospherical surface in a well-quantified fashion.

Organization of the paper. In Section 1, we introduce the necessary notions on Anti-de Sitter space, maximal surfaces and the width, we collect several results proved in [BS10], and finally we give some generalities on quasisymmetric homeomorphisms. In Section 2 we discuss the relation between the cross-ratio norm and the width in Anti-de Sitter space. The main result is Proposition 3.A. Section 3 proves Theorem 1.A, while Section 4 is devoted to the proof of Theorem 1.B. Finally in Section 5 we introduce quasiconformal mappings and minimal Lagrangian extensions, and we prove Theorem 2.A, Theorem 2.B, Proposition 2.C and Corollary 2.D.

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1. ANTI-DE SITTER SPACE AND MAXIMAL SURFACES

Anti-de Sitter space AdS^3 is a pseudo-Riemannian manifold of signature $(2, 1)$ of constant curvature -1 . Consider $\mathbb{R}^{2,2}$, the vector space \mathbb{R}^4 endowed with the bilinear form of signature $(2,2)$:

$$\langle x, y \rangle = x^1 y^1 + x^2 y^2 - x^3 y^3 - x^4 y^4$$

and define

$$\widehat{\text{AdS}}^3 = \{x \in \mathbb{R}^{2,2} : \langle x, x \rangle = -1\}.$$

It turns out that $\widehat{\text{AdS}}^3$ is connected, time-orientable and has the topology of a solid torus. We define Anti-de Sitter space to be the projective domain

$$\text{AdS}^3 = P(\{\langle x, x \rangle < 0\}) \subset \mathbb{RP}^3,$$

of which $\widehat{\text{AdS}}^3$ is a double cover. The pseudo-Riemannian metric induced on $\widehat{\text{AdS}}^3$ descends to a metric on AdS^3 of constant curvature -1, again time-orientable, which will be denoted again by the product $\langle \cdot, \cdot \rangle$.

The tangent space of $\widehat{\text{AdS}}^3$ at a point x is $T_x \widehat{\text{AdS}}^3 \cong x^\perp$. Vectors in tangent spaces are classified according to their causal properties. In particular:

$$v \in T_x \mathbb{R}^{2,1} \text{ is } \begin{cases} \textit{timelike} & \text{if } \langle v, v \rangle < 0 \\ \textit{lightlike} & \text{if } \langle v, v \rangle = 0 \\ \textit{spacelike} & \text{if } \langle v, v \rangle > 0 \end{cases}.$$

Hence, given a spacelike curve $\gamma : I \rightarrow \text{AdS}^3$ (i.e. a differentiable curve whose tangent vector at every point is spacelike), we define the *length* of γ as

$$\text{length}(\gamma) = \int_I \sqrt{\langle \dot{\gamma}(t), \dot{\gamma}(t) \rangle} dt.$$

On the other hand, if γ is a timelike curve, we still define its *length*, as

$$\text{length}(\gamma) = \int_I \sqrt{-\langle \dot{\gamma}(t), \dot{\gamma}(t) \rangle} dt.$$

We will fix once and forever a time orientation, so as to talk about *future-directed* vectors and curves. Our convention is that the vector $(0, 0, 0, 1)$ (based at the point $(0, 0, 1, 0)$) is future-directed.

The group of isometries of $\widehat{\text{AdS}}^3$ which preserve the orientation and the time-orientation is $\text{SO}_+(2, 2)$, namely the connected component of the identity in the group of linear isometries of $\mathbb{R}^{2,2}$. It follows that the group of orientation-preserving, time-preserving isometries of AdS^3 is $\text{SO}_+(2, 2) / \{\pm I\}$, and will be denoted simply by $\text{Isom}(\text{AdS}^3)$.

Geodesics of $\widehat{\text{AdS}}^3$ are the intersection of $\widehat{\text{AdS}}^3$ with linear planes of $\mathbb{R}^{2,2}$. Therefore, geodesics of AdS^3 are projective lines which intersect the projective domain $\text{AdS}^3 \subset \mathbb{RP}^3$. It is easy to see that a unit speed parameterization of a spacelike geodesic with initial point p and initial tangent (spacelike) vector v is:

$$(1) \quad \gamma(t) = [\cosh(t)p + \sinh(t)v],$$

where the square brackets denote the class in AdS^3 of a point of $\widehat{\text{AdS}}^3$. On the other hand, when v is a timelike vector, the parameterization of the timelike geodesic is

$$(2) \quad \gamma(t) = [\cos(t)p + \sin(t)v].$$

Hence timelike geodesics are closed and have length π . Equations (1) and (2) enable to derive immediately the formulae for the length of a spacelike geodesic segment:

$$(3) \quad \cosh(\text{length}(\overline{pq})) = |\langle p, q \rangle|,$$

while for a timelike geodesic segment one gets:

$$(4) \quad \cos(\text{length}(\overline{pq})) = |\langle p, q \rangle|.$$

Analogously, totally geodesic planes of AdS^3 are projective planes and are the projection of the intersection of $\widehat{\text{AdS}}^3$ with three-dimensional linear subspaces of $\mathbb{R}^{2,2}$. There is a duality

between totally geodesic planes and points of AdS^3 , which is given by associating to a point $x \in \text{AdS}^3$ the *dual plane* $P = x^\perp$. One defines

$$P = x^\perp \text{ is } \begin{cases} \textit{timelike} & \text{if } \langle x, x \rangle > 0 & (\Leftrightarrow \text{the induced metric is Lorentzian}) \\ \textit{lightlike} & \text{if } \langle x, x \rangle = 0 & (\Leftrightarrow \text{the induced metric is degenerate}) \\ \textit{spacelike} & \text{if } \langle x, x \rangle < 0 & (\Leftrightarrow \text{the induced metric is Riemannian}) \end{cases} .$$

Spacelike totally geodesic planes are isometric to the hyperbolic plane \mathbb{H}^2 . Using Equation (4), it is easy to check that the dual plane of a point x coincides with

$$x^\perp = \{\gamma(\pi/2) | \gamma : [0, \pi] \rightarrow \text{AdS}^3 \text{ is a timelike geodesic, } \gamma(0) = \gamma(\pi) = x\}.$$

In the affine chart $\{x_4 \neq 0\}$, AdS^3 fills the domain $\{x^2 + y^2 < 1 + z^2\}$, interior of a one-sheeted hyperboloid; however AdS^3 is not contained in a single affine chart, hence in this description we are missing a totally geodesic plane at infinity P_∞ , which is the dual plane of the origin. Since geodesics in AdS^3 are intersections of AdS^3 with linear planes in $\mathbb{R}^{2,2}$, in the affine chart geodesics are represented by straight lines. See Figure 1.1 for a picture in the affine chart $\{x_4 \neq 0\}$.

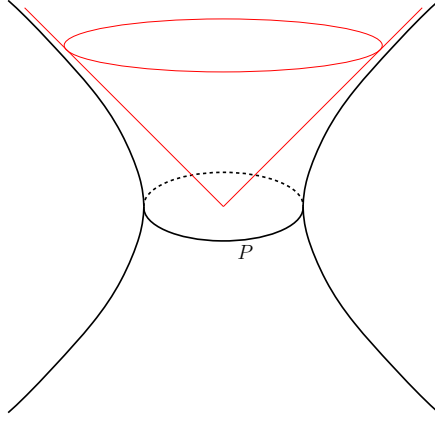


Figure 1.1. The light-cone of future null geodesic rays from a point and a totally geodesic spacelike plane P .

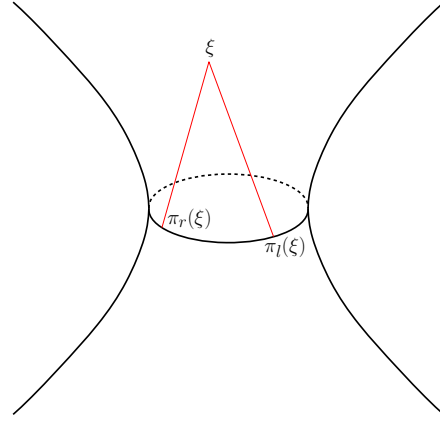


Figure 1.2. Left and right projection from a point $\xi \in \partial_\infty \text{AdS}^3$ to the plane $P = \{x_3 = 0\}$

The boundary at infinity of AdS^3 is defined as the topological frontier of AdS^3 in \mathbb{RP}^3 , namely the doubly ruled quadric

$$\partial_\infty \text{AdS}^3 = P(\{\langle x, x \rangle = 0\}).$$

It is naturally endowed with a conformal Lorentzian structure, for which the null lines are precisely the left and right ruling. Given a spacelike plane P , which we recall is obtained as intersection of AdS^3 with a linear hyperplane of \mathbb{RP}^3 and is a copy of \mathbb{H}^2 , P has a natural boundary at infinity $\partial_\infty(P)$ which coincides with the usual boundary at infinity of \mathbb{H}^2 . Moreover, P intersects each line in the left or right ruling in exactly one point. If a spacelike plane P is chosen, $\partial_\infty \text{AdS}^3$ can be identified with $\partial_\infty \mathbb{H}^2 \times \partial_\infty \mathbb{H}^2$ by means of the following description: $\xi \in \partial_\infty \text{AdS}^3$ corresponds to $(\pi_l(\xi), \pi_r(\xi))$, where π_l and π_r are the projection to $\partial_\infty(P)$ following the left and right ruling respectively (compare Figure 1.2). Under this identification, the isometry group of AdS^3 acts on $\partial_\infty \text{AdS}^3$ by projective transformations, and it turns out that

$$\text{Isom}(\text{AdS}^3) \cong \text{PSL}(2, \mathbb{R}) \times \text{PSL}(2, \mathbb{R}).$$

Given an orientation-preserving homeomorphism $\phi : \partial_\infty \mathbb{H}^2 \rightarrow \partial_\infty \mathbb{H}^2$, by means of this identification, the graph of ϕ can be thought of as a curve in $\partial_\infty \text{AdS}^3$, denoted simply by $gr(\phi)$.

1.1. Maximal surfaces. This paper is concerned with spacelike embedded surfaces in Anti-de Sitter space. A smooth embedded surface $\sigma : S \rightarrow \text{AdS}^3$ is called spacelike if the first fundamental form $I(v, w) = \langle d\sigma(v), d\sigma(w) \rangle$ is a Riemannian metric on S . Equivalently, the tangent plane is a spacelike plane at every point. Let N be a unit normal vector field to the embedded surface S . We denote by ∇ and ∇^S the ambient connection and the Levi-Civita connection of the surface S , respectively. The *second fundamental form* of S is defined as

$$\nabla_{\tilde{v}} \tilde{w} = \nabla_{\tilde{v}}^S \tilde{w} + II(v, w)N$$

if \tilde{v} and \tilde{w} are vector fields extending v and w . The *shape operator* is the $(1, 1)$ -tensor defined as $B(v) = \nabla_v N$. It satisfies the property

$$II(v, w) = \langle B(v), w \rangle.$$

Definition 1.1. A smooth embedded spacelike surface S in AdS^3 is *maximal* if $\text{tr}B = 0$.

The shape operator is symmetric with respect to the first fundamental form of the surface S ; hence the condition of maximality amounts to the fact that the principal curvatures (namely, the eigenvalues of B) are opposite at every point.

Definition 1.2. We say that an embedded spacelike surface S in AdS^3 is *entire* if it is a compression disc, i.e. it is a topological disc and its frontier is contained in $\partial_\infty \text{AdS}^3$.

The condition that S is entire is equivalent to the fact that S can be expressed as a graph over \mathbb{H}^2 in a suitable coordinate system, see [BS10].

An existence result for maximal surfaces in AdS^3 was given by Bonsante and Schlenker.

Theorem 1.3 ([BS10]). *Given any orientation-preserving homeomorphism $\phi : \mathbb{RP}^1 \rightarrow \mathbb{RP}^1$, there exists an entire maximal surface with nonpositive curvature S in AdS^3 such that $\partial_\infty S = gr(\phi)$.*

Observe that, by the Gauss equation in AdS^3 , the curvature of the induced metric on the maximal surface S is $-1 + \lambda^2$, where λ and $-\lambda$ are the principal curvatures. Hence the condition of nonpositive curvature corresponds to the fact that the principal curvatures of S are in $[-1, 1]$.

1.2. Width of convex hulls. Since AdS^3 is a projective geometry, we have a well-defined notion of convexity. In particular, we can give the definition of *convex hull* of a curve in $\partial_\infty \text{AdS}^3$.

Definition 1.4. Given a curve $\Gamma = gr(\phi)$ in $\partial_\infty \text{AdS}^3$, the convex hull of Γ , which we denote by $\mathcal{CH}(\Gamma)$, is the intersection of half-spaces bounded by planes P such that $\partial_\infty P$ does not intersect Γ , and the half-space is taken on the side of P containing Γ .

It can be proved that the convex hull of Γ , which is well-defined in \mathbb{RP}^3 , is contained in $\text{AdS}^3 \cup \partial_\infty \text{AdS}^3$, and is actually contained in an affine chart.

Let us fix a totally geodesic spacelike plane Q (for instance, the plane $\{x_4 = 0\}$, which is the plane at infinity in Figure 1.1). We will denote by $d_{\text{AdS}^3}(\cdot, \cdot)$ the *timelike distance* in $\text{AdS}^3 \setminus Q$, which is defined as follows.

Definition 1.5. Given points p and q which are connected by a timelike curve, the timelike distance between p and q is

$$d_{\text{AdS}^3}(p, q) = \sup_{\gamma} \{\text{length}(\gamma) \mid \gamma : [0, 1] \rightarrow \text{AdS}^3 \setminus Q \text{ is a timelike curve, } \gamma(0) = p, \gamma(1) = q\}.$$

The distance between two such points p, q is achieved along the timelike geodesic connecting p and q . The timelike distance satisfies the reverse triangle inequality, meaning that,

$$d_{\mathbb{A}dS^3}(p, r) \geq d_{\mathbb{A}dS^3}(p, q) + d_{\mathbb{A}dS^3}(q, r),$$

provided both pairs (p, q) and (q, r) are connected by a timelike curve. In a completely analogous way, we define the distance of a point x from a totally geodesic spacelike plane P as the supremum of the length of a timelike curve connecting x to P .

We are now ready to introduce the notion of *width* of the convex hull, as defined in [BS10].

Definition 1.6. Given a curve $\Gamma = gr(\phi)$ in $\partial_\infty \mathbb{A}dS^3$, the width w of the convex hull $\mathcal{CH}(\Gamma)$ is the supremum of the length of a timelike geodesic contained in $\mathcal{CH}(\Gamma)$.

Remark 1.7. Note that the distance $d_{\mathbb{A}dS^3}(p, q)$ is achieved along the geodesic timelike segment connecting p and q . Hence Definition 1.6 is equivalent to

$$(5) \quad w = \sup_{p \in \partial_- \mathcal{C}, q \in \partial_+ \mathcal{C}} d_{\mathbb{A}dS^3}(p, q),$$

where $\mathcal{C} = \mathcal{CH}(\Gamma)$ and $\partial_\pm \mathcal{C}$ denote the two components (one future and one past component) of the boundary of the convex hull of Γ in $\mathbb{A}dS^3$.

In particular, we note that

$$(6) \quad w = \sup_{x \in \mathcal{C}} (d_{\mathbb{A}dS^3}(x, \partial_- \mathcal{C}) + d_{\mathbb{A}dS^3}(x, \partial_+ \mathcal{C})).$$

To stress once more the meaning of this equality, note that the supremum in (6) cannot be achieved on a point x such that the two segments realizing the distance from x to $\partial_- \mathcal{C}$ and $\partial_+ \mathcal{C}$ are not part of a unique geodesic line. Indeed, if at x the two segments form an angle, the piecewise geodesic can be made longer by avoiding the point x , as in Figure 1.3. We also remark that if the distance between a point x and $\partial_\pm \mathcal{C}$ is achieved along a geodesic segment l , then the maximality condition imposes that l must be orthogonal to a support plane to $\partial_\pm \mathcal{C}$ at $\partial_\pm \mathcal{C} \cap l$.

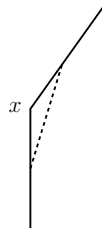


Figure 1.3. A path through x which is not geodesic does not achieve the maximum distance.

1.3. An application of the maximum principle. A key property used in this paper is that maximal surfaces with boundary at infinity a curve $\Gamma = gr(\phi)$ are contained in the convex hull of Γ . Although this fact is known, we prove it here by applying maximum principle to a simple linear PDE describing maximal surfaces.

Hereafter $\text{Hess } u$ denotes the Hessian of a smooth function u on the surface S , i.e. the (1,1) tensor

$$\text{Hess } u(v) = \nabla_v^S \text{grad } u.$$

Finally, Δ_S denotes the Laplace-Beltrami operator of S , which can be defined as

$$\Delta_S u = \text{tr}(\text{Hess } u).$$

Proposition 1.8 was proved in [BS10]. We give a proof here for the sake of completeness. We first observe that, given a point x and a totally geodesic spacelike plane P , it is easy to

check (as for Equations (3) and (4)) that the timelike distance $d_{\text{AdS}^3}(x, P)$ of x from $P = p^\perp$ satisfies

$$\sin d_{\text{AdS}^3}(x, P) = |\langle x, p \rangle|.$$

Proposition 1.8. *Given a maximal surface $S \subset \text{AdS}^3$ and a plane P , let $u : S \rightarrow \mathbb{R}$ be the function $u(x) = \sin d_{\text{AdS}^3}(x, P)$, where $d_{\text{AdS}^3}(x, P)$ is considered as a signed distance. Let N be the future unit normal to S and $B = \nabla N$ the shape operator. Then*

$$(7) \quad \text{Hess } u - uE = \sqrt{1 - u^2 + \|\text{grad } u\|^2} B,$$

where E denotes the identity operator. As a consequence, u satisfies the linear equation

$$(L) \quad \Delta_S u - 2u = 0.$$

Proof. Let us assume that P is the plane dual to the point $p \in \text{AdS}^3$. We will perform the computation in the double cover $\widehat{\text{AdS}^3}$. Then u is the restriction to S of the function U defined by:

$$(8) \quad U(x) = \sin d_{\text{AdS}^3}(x, P) = \langle x, p \rangle.$$

Let N be the unit normal vector field to S ; we compute $\text{grad } u$ by projecting the gradient ∇U of U to the tangent plane to S :

$$(9) \quad \nabla U = p + \langle p, x \rangle x$$

$$(10) \quad \text{grad } u(x) = p + \langle p, x \rangle x + \langle p, N \rangle N$$

Now $\text{Hess } u(v) = \nabla_v^S \text{grad } u$, where ∇^S is the Levi-Civita connection of S , namely the projection of the flat connection of $\mathbb{R}^{2,2}$, and so for any $v \in T_x \widehat{\text{AdS}^3}$ one gets:

$$\text{Hess } u(v) = \langle p, x \rangle v + \langle p, N \rangle \nabla_v^S N = u(x)v + \langle \nabla U, N \rangle B(v).$$

Moreover, $\nabla U = \text{grad } u + \langle \nabla U, N \rangle N$ and thus

$$(11) \quad \langle \nabla U, N \rangle^2 = \|\text{grad } u\|^2 - \langle \nabla U, \nabla U \rangle = 1 - u^2 + \|\text{grad } u\|^2,$$

which proves (7). By taking the trace, (L) follows. \square

Corollary 1.9. *Let S be an entire maximal surface in AdS^3 . Then S is contained in the convex hull of $\partial_\infty S$.*

Proof. If $\Gamma = \partial_\infty S$ is a circle, then S is a totally geodesic plane which coincides with the convex hull of Γ . Hence we can suppose Γ is not a circle. Consider a plane P_- which does not intersect Γ and the function u defined as in Equation (8) in Proposition 1.8, with respect to P_- . Suppose their mutual position is such that $u \geq 0$ in the region of S close to the boundary at infinity (i.e. in the complement of a large compact set). If there exists some point where $u < 0$, then by Equation (L) at a minimum point $\Delta_S u = 2u < 0$, which gives a contradiction. The proof is analogous for a plane P_+ on the other side of Γ , by switching the signs. This shows that every convex set containing Γ contains also S . \square

1.4. Two opposite examples. We report here the two examples to have in mind for our study of maximal surface. The first is a very simple example, namely a totally geodesic plane, for which the principal curvatures vanish at every point. To some extent, the second example can be considered as the opposite of a totally geodesic plane, since in the second case the principal curvatures are 1 and -1 at every point.

Example 1.10. (Totally geodesic planes) By definition, a totally geodesic plane P has shape operator $B \equiv 0$. Hence a totally geodesic plane is a maximal surface, and $\partial_\infty P = gr(A)$, where $A \in \text{PSL}(2, \mathbb{R})$ is the trace on $\mathbb{RP}^1 = \partial_\infty \mathbb{H}^2$ of an isometry of \mathbb{H}^2 . Hence the width of the convex hull vanishes. It is also easy to see that, given a curve in $\partial_\infty \text{AdS}^3$, if $w = 0$ then the curve is necessarily the boundary of a totally geodesic plane. Therefore the unique maximal surface with zero width (up to isometries of AdS^3) is a totally geodesic plane.

Example 1.11. (Horospherical surfaces) Consider a spacelike line l in AdS^3 . The dual line l^\perp is obtained as the intersection of all totally geodesic planes dual to points of l . Recall that timelike geodesics in AdS^3 are closed and have length π . Hence one equivalently has:

$$l^\perp = \{\gamma(\pi/2) | \gamma : [0, \pi] \rightarrow \text{AdS}^3 \text{ is a timelike geodesic}, \gamma(0) = \gamma(\pi) \in l\}.$$

Let us define the smooth surface

$$H = \{\gamma(\pi/4) | \gamma : [0, \pi/2] \rightarrow \text{AdS}^3 \text{ is a future-directed timelike geodesic}, \gamma(0) \in l, \gamma(\pi/2) \in l^\perp\}.$$

See Figure 1.4 for a schematic picture.

The group of isometries which preserve l (and thus preserves also l^\perp) is isomorphic to $\mathbb{R} \times \mathbb{R}$, where $\mathbb{R} \times \{0\}$ fixes l pointwise and acts on l^\perp by translation (it actually acts as a rotation around the spacelike line l), while $\{0\} \times \mathbb{R}$ does the opposite. The induced metric on H is flat, and thus H is isometric to the Euclidean plane. Moreover, for every point $x \in H$, the surface H has an orientation-reversing, time-reversing isometry obtained by reflection in a plane P tangent to a point $x \in H$, followed by rotation of angle $\pi/2$ around the timelike geodesic orthogonal to P at x . This basically shows that the principal curvatures λ_1, λ_2 of H are necessarily opposite to one another, hence H is a maximal surface. Moreover, since by the Gauss equation

$$0 = K_H = -1 - \lambda_1 \lambda_2 = -1 + \lambda^2,$$

it follows that the principal curvatures are necessarily ± 1 at every point. Finally, by construction, the width of the convex hull of H is precisely $\pi/2$.

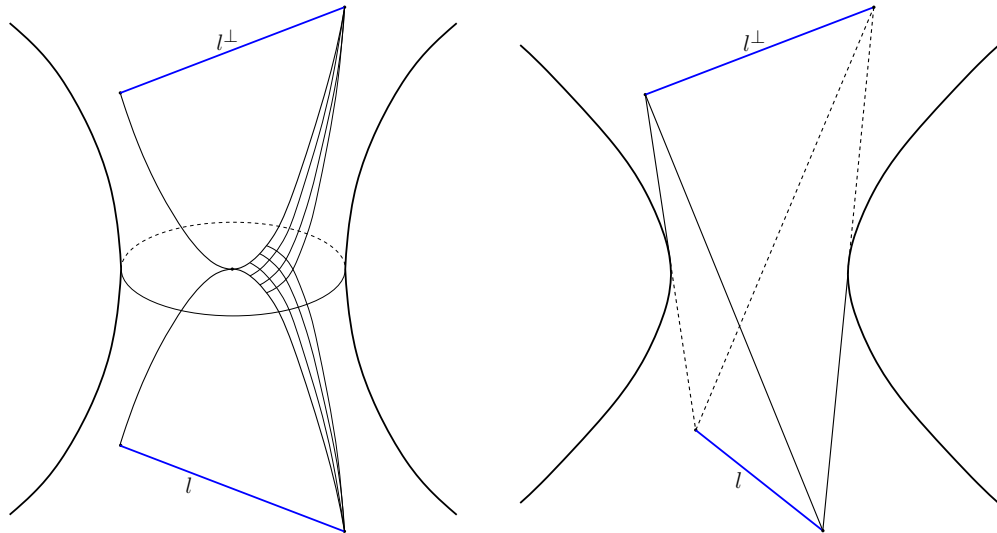


Figure 1.4. A subset of a horospherical surface. The lines of curvature provide Euclidean coordinates for the induced metric, and are planar curves parallel to l or l^\perp . The boundary at infinity is composed of four lightlike segments.

Bonsante and Schlenker proved an important property of rigidity of maximal surfaces with large principal curvatures.

Lemma 1.12 ([BS10, Lemma 5.5]). *Given a maximal surface S in AdS^3 with nonpositive curvature, if the curvature is 0 at some point, then S is a subset of a horospherical surface.*

The proof of Lemma 1.12 follows from applying the maximum principle to Equation (Q), which is stated in the following lemma. Recall that an umbilical point on a surface is a point where the two principal curvatures are equal. Hence for a maximal surface, umbilical points are the points where the principal curvatures vanish.

Lemma 1.13 ([KS07, Lemma 3.11]). *Given a maximal surface S in AdS^3 , with principal curvatures $\pm\lambda$, let $\chi : S \rightarrow \mathbb{R}$ be the function defined (in the complement of umbilical points) by $\chi = -\ln|\lambda|$. Then χ satisfies the quasi-linear equation*

$$(Q) \quad \Delta_S \chi = 2(1 - e^{-2\chi}).$$

1.5. Uniformly negative curvature. As a warm-up for what will come next, we give here a proof of Proposition 1.C. Our proof was basically already implicit in [BS10, Claim 3.21], and will be a consequence of the following easy lemma, which we prove here by completeness. See also [KS07].

Lemma 1.14. *Given a smooth spacelike surface S in AdS^3 , let S_ρ be the surface at timelike distance ρ from S , obtained by following the normal flow. Then the pull-back to S of the induced metric on the surface S_ρ is given by*

$$(12) \quad I_\rho = I((\cos(\rho)E + \sin(\rho)B)\cdot, (\cos(\rho)E + \sin(\rho)B)\cdot).$$

The second fundamental form and the shape operator of S_ρ are given by

$$(13) \quad II_\rho = I((-\sin(\rho)E + \cos(\rho)B)\cdot, (\cos(\rho)E - \sin(\rho)B)\cdot),$$

$$(14) \quad B_\rho = (\cos(\rho)E + \sin(\rho)B)^{-1}(-\sin(\rho)E + \cos(\rho)B).$$

Proof. Let σ be a smooth embedding of the maximal surface S , with oriented unit normal N . The geodesics orthogonal to S at a point $x = \sigma(y)$ can be written as

$$\gamma_x(\rho) = \cos(\rho)\sigma(y) + \sin(\rho)N(x).$$

Then we compute

$$\begin{aligned} I_\rho(v, w) &= \langle d\gamma_x(\rho)(v), d\gamma_x(\rho)(w) \rangle \\ &= \langle \cos(\rho)d\sigma_x(v) + \sin(\rho)d(N \circ \sigma)_x(v), \cos(\rho)d\sigma_x(w) + \sin(\rho)d(N \circ \sigma)_x(w) \rangle \\ &= I(\cos(\rho)v + \sin(\rho)B(v), \cos(\rho)w + \sin(\rho)B(w)). \end{aligned}$$

We have used the equation $B = \nabla N$. The formula for the second fundamental form follows from the fact that $II_\rho = \frac{1}{2} \frac{dI_\rho}{d\rho}$, and the formula for B_ρ from equating $B_\rho = I_\rho^{-1}II_\rho$. \square

It follows that, if the principal curvatures of a maximal surface S are $\lambda_1 = \lambda \in [0, 1)$ and $\lambda_2 = -\lambda$, then the principal curvatures of S_ρ are

$$\lambda_\rho = \frac{\lambda - \tan(\rho)}{1 + \lambda \tan(\rho)} = \tan(\rho_0 - \rho),$$

where $\tan \rho_0 = \lambda$, and

$$\lambda'_\rho = \frac{-\lambda - \tan(-\rho)}{1 - \lambda \tan(\rho)} = \tan(-\rho_0 - \rho).$$

In particular λ_ρ and λ'_ρ are non-singular for every ρ between $-\pi/4$ and $\pi/4$. It then turns out that S_ρ is convex at every point for $\rho < -\|\rho_0\|_\infty$, and concave for $\rho > \|\rho_0\|_\infty$. This proves that the width w is less than $2\|\rho_0\|_\infty = 2 \arctan \|\lambda\|_\infty$. Therefore

$$\tan\left(\frac{w}{2}\right) \leq \|\lambda\|_\infty,$$

and the statement of Proposition 1.C follows.

Proposition 1.C. *Let S be an entire maximal surface in AdS^3 with $\|\lambda\|_\infty < 1$ and let w be the width of the convex hull. Then*

$$\tan w \leq \frac{2\|\lambda\|_\infty}{1 - \|\lambda\|_\infty^2}.$$

A direct consequence is that, if S is an entire maximal surface with uniformly negative curvature (equivalently, with $\|\lambda\|_\infty < 1$), then $w < \pi/2$, see also [BS10, Corollary 3.22]. Also the converse holds:

Proposition 1.15 ([BS10]). *Given any entire maximal surface S of nonpositive curvature in AdS^3 , the width w of the convex hull of S is at most $\pi/2$. Moreover, $w < \pi/2$ if and only if S has uniformly negative curvature.*

The proof of the converse implication in [BS10] used quasisymmetric homeomorphisms (see below) and minimal Lagrangian extensions (which will be discussed in Section 5). Observe that our Theorem 1.B gives a quantitative version of the converse implication of Proposition 1.15.

1.6. Quasisymmetric homeomorphisms of the circle. Given an orientation-preserving homeomorphism $\phi : \mathbb{RP}^1 \rightarrow \mathbb{RP}^1$, we define the *cross-ratio norm* of ϕ as

$$\|\phi\|_{cr} = \sup_{cr(Q)=-1} |\ln |cr(\phi(Q))||,$$

where $Q = (z_1, z_2, z_3, z_4)$ is any quadruple of points on \mathbb{RP}^1 and we use the following definition of cross-ratio:

$$cr(z_1, z_2, z_3, z_4) = \frac{(z_4 - z_1)(z_3 - z_2)}{(z_2 - z_1)(z_3 - z_4)}.$$

According to this definition, a quadruple $Q = (z_1, z_2, z_3, z_4)$ is symmetric (i.e. the hyperbolic geodesics connecting z_1 to z_3 and z_2 to z_4 intersect orthogonally) if and only if $cr(Q) = -1$.

Observe that $\|\phi\|_{cr} \in [0, \infty]$ and $\|\phi\|_{cr} = 0$ if and only if ϕ is a projective transformation, i.e. $\phi \in \text{PSL}(2, \mathbb{R})$. Indeed, by post-composing with a projective transformation, one can assume that ϕ fixes three points of \mathbb{RP}^1 , and then the conclusion is straightforward. A homeomorphism is *quasisymmetric* when it has finite cross-ratio norm:

Definition 1.16. An orientation-preserving homeomorphism $\phi : \mathbb{RP}^1 \rightarrow \mathbb{RP}^1$ is quasisymmetric if and only if $\|\phi\|_{cr} < +\infty$.

Quasisymmetric homeomorphisms arise naturally in the context of quasiconformal mappings and universal Teichmüller space, which is actually one of the main themes of this paper. However, we defer the introduction of this point of view to Section 5, where some applications are given.

The first intimate correlation between the cross-ratio norm of a quasisymmetric homeomorphism ϕ and the width of the convex hull of $gr(\phi)$ was proved in [BS10]:

Theorem 1.17 ([BS10, Theorem 1.12]). *Given any orientation-preserving homeomorphism $\phi : \mathbb{RP}^1 \rightarrow \mathbb{RP}^1$, let w be the width of the convex hull of $gr(\phi)$. Then $w < \pi/2$ if and only if ϕ is quasisymmetric.*

Again, Proposition 3.A will provide a more precise version of Theorem 1.17, giving a quantitative inequality between the width and the cross-ratio norm.

Remark 1.18. Actually, Theorem 1.17 holds under a more general hypothesis, namely for more general curves in $\partial_\infty \text{AdS}^3$, which, roughly speaking, may also contain lightlike segments. We will not be interested in these objects in this paper.

A refinement of Theorem 1.3 was given in [BS10] under the assumption that the curve at infinity is the graph of a quasisymmetric homeomorphism.

Theorem 1.19 ([BS10]). *Given any **quasisymmetric** homeomorphism $\phi : \mathbb{RP}^1 \rightarrow \mathbb{RP}^1$, there exists a **unique** entire maximal surface S in AdS^3 with **uniformly** negative curvature such that $\partial_\infty S = gr(\phi)$.*

1.7. Compactness properties. One of the main tools used in [BS10] is a result of compactness for maximal surfaces in AdS^3 . To conclude the preliminaries of this paper, we briefly discuss some properties of compactness for maximal surfaces and quasisymmetric homeomorphisms.

Given a spacelike plane P_0 in AdS^3 and a point $x_0 \in P_0$, let l be the timelike geodesic through x_0 orthogonal to P_0 . We define the solid cylinder $Cl(x_0, P_0, R_0)$ of radius R_0 above P_0 centered at x_0 as the set of points $x \in \text{AdS}^3$ which lie on a spacelike plane P_x orthogonal to l such that $d_{P_x}(x, l \cap P_x) \leq R_0$. See also Figure 1.5.

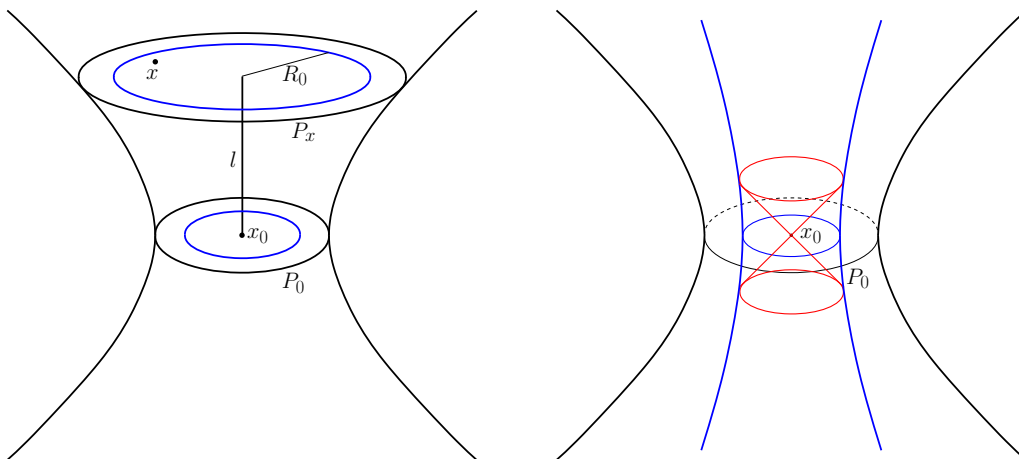


Figure 1.5. The definition of the solid cylinder $Cl(x_0, P_0, R_0)$. On the right, its intersection with the future and past lightcone over x_0 .

The solid cylinder $Cl(x_0, P_0, R_0)$ can also be conveniently described in the following way. Assuming (in the double cover $\widetilde{\text{AdS}}^3$, for one moment) that $x_0 = (0, 0, 1, 0)$ and the normal vector to P_0 is $N_0 = (0, 0, 0, 1)$, let us consider the following coordinate system:

$$(15) \quad (r, \theta, \zeta) \mapsto [\cos \theta \sinh r, \sin \theta \sinh r, \cos \zeta \cosh r, \sin \zeta \cosh r],$$

defined for $r \in \mathbb{R}, \theta \in S^1, \zeta \in S^1$. This means that the level sets with $\zeta = c$ are totally geodesic planes orthogonal to the timelike line

$$\{[0, 0, \cos \zeta, \sin \zeta] : \zeta \in S^1\},$$

which passes through the point $x_0 = (0, 0, 1, 0)$ with future-directed normal $N_0 = (0, 0, 0, 1)$. It is easy to see that, by this description, the solid cylinder $Cl(x_0, P_0, R_0)$ is determined by the relation $r \leq R_0$.

From the tools in the paper [BS10], the following lemma is proved:

Lemma 1.20. *Given a spacelike plane P_0 and a point $x_0 \in P_0$, every sequence S_n of entire maximal surfaces with nonpositive curvature, tangent to P_0 at x_0 , admits a subsequence converging C^∞ to an entire maximal surface on the solid cylinders $Cl(x_0, P_0, R)$, for every $R > 0$.*

A somehow similar property of compactness for quasisymmetric homeomorphisms will be used in several occasions. See also [BZ06] for a discussion.

Theorem 1.21. *Let $k > 0$ and $\phi_n : \mathbb{RP}^1 \rightarrow \mathbb{RP}^1$ be a family of orientation-preserving quasisymmetric homeomorphisms of the circle, with $\|\phi_n\|_{cr} \leq k$. Then there exists a subsequence ϕ_{n_k} for which one of the following holds:*

- *The homeomorphisms ϕ_{n_k} converge uniformly to a quasisymmetric homeomorphism $\phi : \mathbb{RP}^1 \rightarrow \mathbb{RP}^1$, with $\|\phi\|_{cr} \leq k$;*

- The homeomorphisms ϕ_{n_k} converge uniformly on the complement of any open neighborhood of a point of \mathbb{RP}^1 to a constant map $c : \mathbb{RP}^1 \rightarrow \mathbb{RP}^1$.

Remark 1.22. It is also not difficult to prove that, given a sequence of entire maximal surfaces S_n which converges uniformly on compact cylinders $Cl(x_0, P_0, R)$ to an entire maximal surface S_∞ , then the asymptotic boundaries $\partial_\infty S_n$ converge (in the Hausdorff convergence, for instance) to the asymptotic boundary $\partial_\infty S_\infty$ of the limit surface.

2. CROSS-RATIO NORM AND THE WIDTH OF THE CONVEX HULL

The purpose of this section is to investigate the relation between the width of the convex hull of $gr(\phi)$ and the cross-ratio norm of ϕ when ϕ is a quasymmetric homeomorphism. The main result is thus Proposition 3.A. To some extent, Proposition 3.A is a quantitative version of Theorem 1.17.

Proposition 3.A. *Given any quasymmetric homeomorphism ϕ of \mathbb{RP}^1 , let w be the width of the convex hull of the graph of ϕ in $\partial_\infty \text{AdS}^3$. Then*

$$(16) \quad \tan(w) \leq \sinh\left(\frac{\|\phi\|_{cr}}{2}\right).$$

Proof. To prove the upper bound on the width, suppose the width of the convex hull \mathcal{C} of $gr(\phi)$ is $w \in (0, \pi/2)$. Let $k = \|\phi\|_{cr}$. We can find a sequence of pairs (p_n, q_n) such that $d_{\text{AdS}^3}(p_n, q_n) \nearrow w$, with $p_n \in \partial_- \mathcal{C}$, $q_n \in \partial_+ \mathcal{C}$. We can assume the geodesic connecting p_n and q_n is orthogonal to $\partial_- \mathcal{C}$ at p_n ; indeed one can replace p_n with a point in $\partial_- \mathcal{C}$ which maximizes the distance from q_n , if necessary (see Remark 1.7). Let us now apply isometries T_n so that $T_n(p_n) = p = [\hat{p}] \in \text{AdS}^3$, for $\hat{p} = (0, 0, 1, 0) \in \widehat{\text{AdS}^3}$, and $T_n(q_n)$ lies on the timelike geodesic through p orthogonal to $P_- = (0, 0, 0, 1)^\perp$.

The curve at infinity $gr(\phi)$ is mapped by T_n to a curve $gr(\phi_n)$, where ϕ_n is obtained by pre-composing and post-composing ϕ with Möbius transformations (this is easily seen from the description of $\text{Isom}(\text{AdS}^3)$ as $\text{PSL}(2, \mathbb{R}) \times \text{PSL}(2, \mathbb{R})$). Hence ϕ_n is still quasymmetric with norm $\|\phi_n\|_{cr} = \|\phi\|_{cr} = k$.

It is easy to see that ϕ_n cannot converge to a map sending the complement of a point in \mathbb{RP}^1 to a single point of \mathbb{RP}^1 . Indeed, the curves $gr(\phi_n)$ are all contained between P_- and a spacelike plane P_n disjoint from P_- , which contains the point $T_n(q_n)$. Moreover the

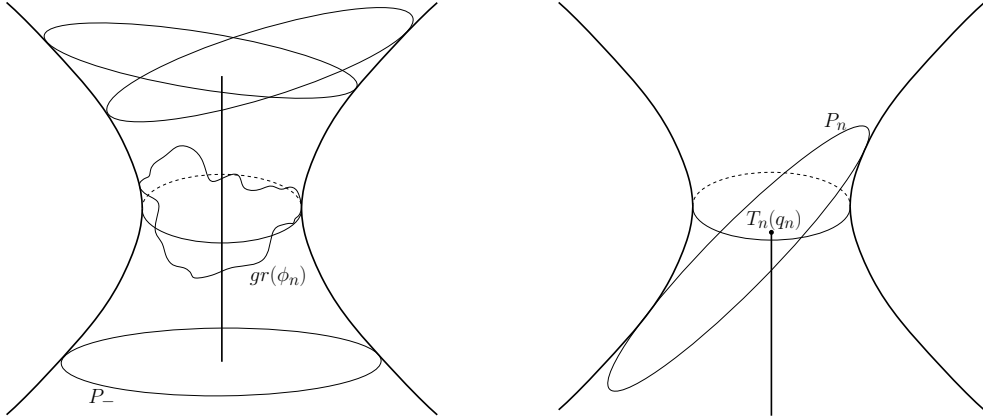


Figure 2.1. The curves $gr(\phi_n)$ are contained in a bounded region in an affine chart, hence they cannot diverge to a constant map. This is easily seen, for instance, by putting the plane P_- at infinity and observing that the plane P_n has to be spacelike, disjoint from P_- , and to intersect the point $T_n(q_n)$ which is in the lower half-plane.

distance of p from $T_n(q_n) \in P_n$ is at most w . This shows that the curves $gr(\phi_n)$ all lie in a bounded region in an affine chart of AdS^3 ; this would not be the case if ϕ_n were converging on the complement of one point to a constant map. See Figure 2.1.

Hence, by the convergence property of k -quasisymmetric homeomorphisms (Theorem 1.21), ϕ_n converges to a k -quasisymmetric homeomorphism ϕ_∞ , so that w equals the width of the convex hull of $gr(\phi_\infty)$. Let us denote by \mathcal{C}_∞ the convex hull of $gr(\phi_\infty)$.

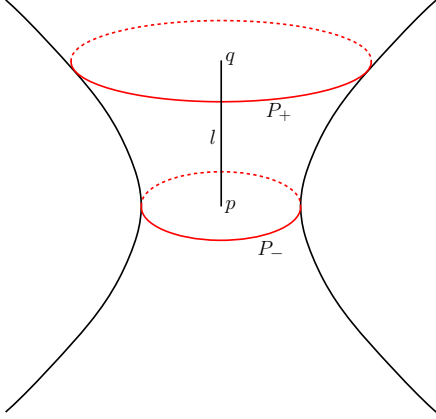


Figure 2.2. The setting of the proof of Proposition 3.A.

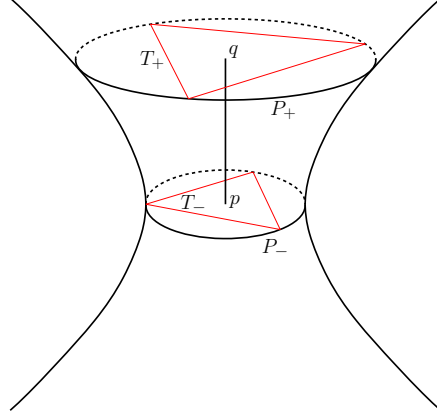


Figure 2.3. The point p is contained in the convex envelope of three (or two) points in $\partial_\infty(P_-)$; analogously q in P_+ .

We will mostly refer to the coordinates in the affine chart $\{x^3 \neq 0\}$, namely $(x, y, z) = (x^1/x^3, x^2/x^3, x^4/x^3)$. Our assumption is that the point p has coordinates $(0, 0, 0)$ and $P_- = \{(x, y, 0) : x^2 + y^2 < 1\}$ is the totally geodesic plane through p which is a support plane for $\partial_- \mathcal{C}_\infty$. The geodesic line l through p orthogonal to P_- is $\{(0, 0, z)\}$. By construction, the width of \mathcal{C}_∞ equals $d_{\text{AdS}^3}(p, q)$, where $q = (0, 0, h) = l \cap \partial_+ \mathcal{C}_\infty$. It is then an easy computation to show that $h = \tan w$. Hence the plane

$$P_+ = \{(x, y, h) : (x, y) \in \mathbb{R}^2, x^2 + y^2 < 1 + h^2\},$$

which is the plane orthogonal to l through q , is a support plane for $\partial_+ \mathcal{C}_\infty$. See Figure 2.2.

Since $\partial_- \mathcal{C}_\infty$ and $\partial_+ \mathcal{C}_\infty$ are pleated surfaces, $\partial_- \mathcal{C}_\infty$ contains an ideal triangle T_- , such that $p \in T_-$ (possibly p is on the boundary of T_-). The ideal triangle might also be degenerate if p is contained in an entire geodesic, but this will not affect the argument. Hence we can find three geodesic half-lines in P_- connecting p to $\partial_\infty \text{AdS}^3$ (or an entire geodesic connecting p to two opposite points in the boundary, if T_- is degenerate). Analogously we have an ideal triangle T_+ in P_+ , compare Figure 2.3. The following sublemma will provide constraints on the position the half-geodesics in P_+ can assume. See Figure 2.4 and 2.5 for a picture of the ‘‘sector’’ described in Lemma 2.1.

Sublemma 2.1. *Suppose $\partial_- \mathcal{C}_\infty \cap P_-$ contains a half-geodesic*

$$g = \{t(\cos \theta, \sin \theta, 0) : t \in [0, 1)\}$$

from p , asymptotic to the point at infinity $\eta = (\cos \theta, \sin \theta, 0)$. Then $\partial_+ \mathcal{C}_\infty \cap P_+$ must be contained in $P_+ \setminus S(\eta)$, where $S(\eta)$ is the sector $\{x \cos \theta + y \sin \theta > 1\}$.

Proof. The computation will be carried out using the coordinates of the double cover $\widehat{\text{AdS}^3}$ of AdS^3 . It suffices to check the assertion when $\theta = \pi$, since in the statement there is a rotational symmetry along the vertical axis. The half-geodesic g is parametrized by

$$g(t) = [\sinh(t), 0, \cosh(t), 0],$$

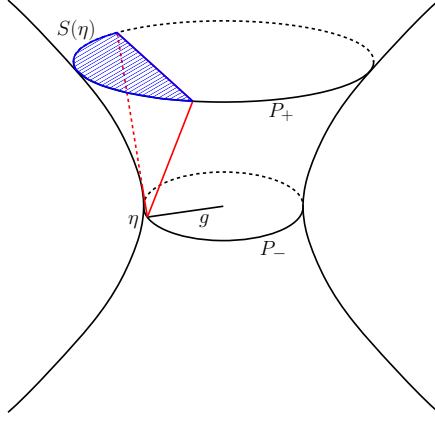


Figure 2.4. The sector $S(\eta)$ as in Sublemma 2.1.

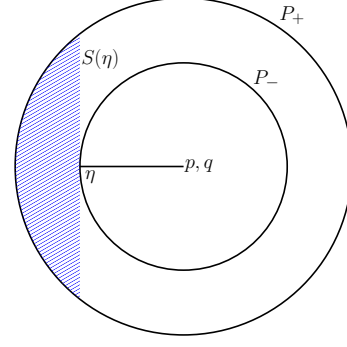


Figure 2.5. The (x, y) -plane seen from above. The sector $S(\eta)$ is bounded by the chord in P_+ tangent to the concentric circle, which projects vertically to P_- .

for $t \in (-\infty, 0]$. Since the width is less than $\pi/2$, every point in $\partial_+\mathcal{C}_\infty \cap P_+$ must lie in the region bounded by P_- and the dual plane $g(t)^\perp$. Indeed for every t , $g(t)^\perp$ is the locus of points at timelike distance $\pi/2$ from $g(t)$. We have

$$P_+ = \{[\cos(\alpha) \sinh(r), \sin(\alpha) \sinh(r), \cos(w) \cosh(r), \sin(w) \cosh(r)] : r \geq 0, \alpha \in [0, 2\pi)\}.$$

Hence the intersection $P_+ \cap g(t)^\perp$ is given by imposing that a point of P_+ has zero product with points $g(t)$, which gives the condition

$$\sinh(t) \cos(\alpha) \sinh(r) = \cosh(t) \cos(w) \cosh(r).$$

Thus points in the intersection are of the form (in the affine coordinates of $\{x^3 \neq 0\}$):

$$\left(\frac{1}{\tanh(t)}, \frac{\tan(\alpha)}{\tanh(t)}, \tan(w) \right).$$

Therefore, points in $\partial_+\mathcal{C}_\infty \cap P_+$ need to have $x \geq 1/\tanh(t)$, and since this holds for every $t \leq 0$, we have $x \geq -1$. \square

By the Sublemma 2.1, if p is contained in the convex envelope of three points η_1, η_2, η_3 in $\partial_\infty(P_-)$, then any point at infinity of $\partial_+\mathcal{C}_\infty \cap P_+$ is necessarily contained in $P_+ \setminus (S(\eta_1) \cup S(\eta_2) \cup S(\eta_3))$. We will use this fact to choose two pairs of points, η, η' in $\partial_\infty(P_-)$ and ξ, ξ' in $\partial_\infty(P_+)$, in a convenient way. This is the content of next sublemma. See Figure 2.6.

Sublemma 2.2. *Suppose p is contained in the convex envelope of three points η_1, η_2, η_3 in $\partial_\infty(P_-)$. Then $gr(\phi_\infty)$ must contain (at least) two points ξ, ξ' of $\partial_\infty(P_+)$ which lie in different connected components of $\partial_\infty(P_+) \setminus (S(\eta_1) \cup S(\eta_2) \cup S(\eta_3))$.*

Proof. The proof is simple 2-dimensional Euclidean geometry. Recall that the point q , which is the “center” of the plane P_+ , is in the convex hull of $gr(\phi_\infty)$. If the claim were false, then one connected component of $\partial_\infty(P_+) \setminus (S(\eta_1) \cup S(\eta_2) \cup S(\eta_3))$ would contain a sector S_0 of angle $\geq \pi$. But then the points η_1, η_2, η_3 would all be contained in the complement of S_0 . This contradicts the fact that p is in the convex hull of η_1, η_2, η_3 . \square

Remark 2.3. If p is in the convex envelope of only two points at infinity, which means that P_- contains an entire geodesic, the previous statement is simplified, see Figure 2.7.

Let us now choose two points $\eta, \eta' \in \partial_\infty(P_-)$ among η_1, η_2, η_3 , and $\xi, \xi' \in \partial_\infty(P_+)$ in such a way that ξ and ξ' lie in two different connected components of $\partial_\infty(P_+) \setminus (S(\eta_1) \cup S(\eta_2))$. The strategy will be to use this quadruple to show that the cross-ratio distortion of ϕ_∞ is

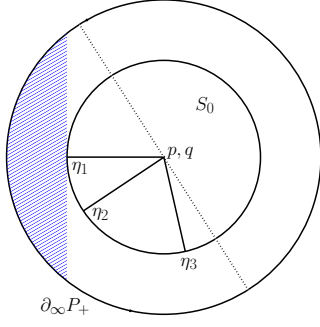


Figure 2.6. The proof of Sublemma 2.2. Below, the choice of points η, η', ξ, ξ' .

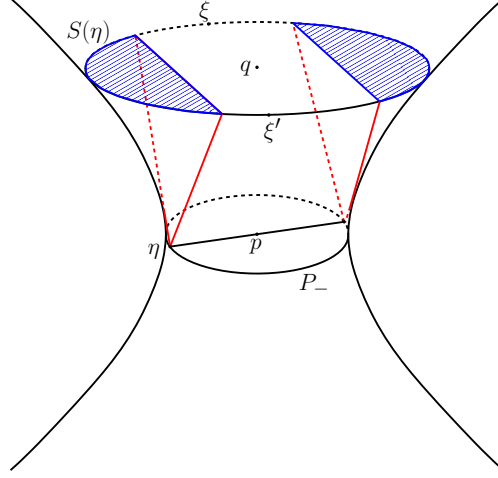
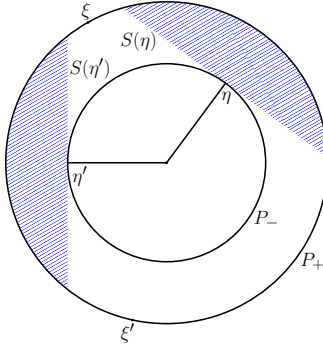


Figure 2.7. The same statement of Sublemma 2.2 is simpler if p is contained in an entire geodesic line contained in P_- .

not too small, depending on the width w . However, such quadruple is not symmetric in general. Hence ξ' will be replaced later by another point ξ'' . First we need some tool to compute the left and right projections to $\partial_\infty \mathbb{H}^2$ of the chosen points.

We use the plane P_- to identify $\partial_\infty \text{AdS}^3$ with $\partial_\infty \mathbb{H}^2 \times \partial_\infty \mathbb{H}^2$. Let π_l and π_r denote left and right projection to $\partial_\infty(P_-)$, following the left and right ruling of $\partial_\infty \text{AdS}^3$. In what follows, angles like θ_l, θ_r and similar symbols will always be considered in $(-\pi, \pi]$.

Sublemma 2.4. *Suppose $\xi \in \partial_\infty(P_+)$, where the length of the timelike geodesic segment orthogonal to P_- and P_+ is w . If $\pi_l(\xi) = (\cos(\theta_l), \sin(\theta_l), 0)$, then $\pi_r(\xi) = (\cos(\theta_l - 2w), \sin(\theta_l - 2w), 0)$.*

Proof. By the description of the left ruling (see Section 1), recalling $h = \tan(w)$, it is easy to check that

$$\begin{aligned} \xi &= (\cos(\theta_l), \sin(\theta_l), 0) + h(\sin(\theta_l), -\cos(\theta_l), 1) = (\cos(\theta_l) + h \sin(\theta_l), \sin(\theta_l) - h \cos(\theta_l), h) \\ &= (\sqrt{1+h^2} \cos(\theta_l - w), \sqrt{1+h^2} \sin(\theta_l - w), h). \end{aligned}$$

By applying the same argument to the right projection, the claim follows. \square

We can assume $\eta' = (-1, 0, 0)$. We shall adopt in this part the complex notation, i.e. \mathbb{H}^2 is thought of in the disc model as a subset of \mathbb{C} , where \mathbb{C} is identified to the plane $\{z = 0\}$ in the affine chart. In this way, η' corresponds to $(-1, -1) \in \partial_\infty \mathbb{H}^2 \times \partial_\infty \mathbb{H}^2$. Let $\eta = (e^{i\theta_0}, e^{i\theta_0})$; by symmetry, we can assume $\theta_0 \in [0, \pi)$; in this case we need to consider the point $\xi = (e^{i\theta_l}, e^{i\theta_r})$ constructed above, with $\theta_r \in [\theta_0, \pi)$. More precisely, Sublemma 2.4 shows $\theta_r = \theta_l - 2w$; by Sublemma 2.1 we must have $\theta_l - w \notin (\theta_0 - w, \theta_0 + w) \cup (\pi - w, \pi) \cup (-\pi, -\pi + w)$ and thus, by choosing ξ in the correct connected component (i.e. switching ξ and ξ' if necessary), necessarily $\theta_l \in [\theta_0 + 2w, \pi]$ (see Figure 2.8).

We remark again that the quadruple $Q = \pi_l(\xi', \eta, \xi, \eta')$ will not be symmetric in general, so we need to consider a point ξ'' instead of ξ' so as to obtain a symmetric quadruple. However, if $\theta_0 \in (-\pi, 0)$, then one would consider the point ξ' in the connected component having $\theta_r \in (-\pi, \theta_0)$ - and then a point ξ'' in the other connected component so as to have a symmetric quadruple - and obtain the same final estimate.

So let $\xi'' = (e^{i\theta_l''}, e^{i\theta_r''})$ be a point on $gr(\phi)$ so that the quadruple $Q = \pi_l(\xi'', \eta, \xi, \eta')$ is symmetric; we are going to compute the cross-ratio of $\phi(Q) = \pi_r(\xi'', \eta, \xi, \eta')$. However, in order to avoid dealing with complex numbers, we first map $\partial_\infty \mathbb{H}^2 = \partial_\infty(P_-)$ to $\mathbb{R}P^1 \cong \mathbb{R} \cup \{\infty\}$ using the Möbius transformation

$$z \mapsto \frac{z-1}{i(z+1)}$$

which maps $e^{i\theta}$ to $\tan(\theta/2) \in \mathbb{R}$ if $\theta \neq \pi$, and -1 to ∞ . We need to compute

$$(17) \quad |\ln |cr(\phi(Q))|| = \left| \ln \left| \frac{\tan(\theta_r/2) - \tan(\theta_0/2)}{\tan(\theta_0/2) - \tan(\theta_r''/2)} \right| \right|$$

and in particular we want to show this is uniformly away from 1. By construction $\theta_r < \theta_l$ (see also Figure 2.9), and since P_- does not disconnect $gr(\phi)$, also $\theta_r'' < \theta_l''$. Hence we have

$$(18) \quad \tan(\theta_0/2) - \tan(\theta_r''/2) \geq \tan(\theta_0/2) - \tan(\theta_l''/2).$$

The condition that $(\theta_l'', \theta_0, \theta_l, \infty)$ forms a symmetric quadruple translates on \mathbb{R} to the condition that

$$(19) \quad \tan(\theta_0/2) - \tan(\theta_l''/2) = \tan(\theta_l/2) - \tan(\theta_0/2).$$

Using (18) and (19) in the argument of the logarithm in (17), we obtain:

$$\frac{\tan(\theta_r/2) - \tan(\theta_0/2)}{\tan(\theta_0/2) - \tan(\theta_r''/2)} \leq \frac{\tan((\theta_l/2) - w) - \tan(\theta_0/2)}{\tan(\theta_l/2) - \tan(\theta_0/2)} =: S(\theta_l).$$

Note that $S(\theta_l) < 1$ on $[\theta_0 + 2w, \pi]$ and $S(\theta_l) \rightarrow 0$ when $\theta_l \rightarrow \theta_0 + 2w$ or $\theta_l \rightarrow \pi$: this corresponds to the fact that $gr(\phi_\infty)$ tends to contain a lightlike segment. On the other hand

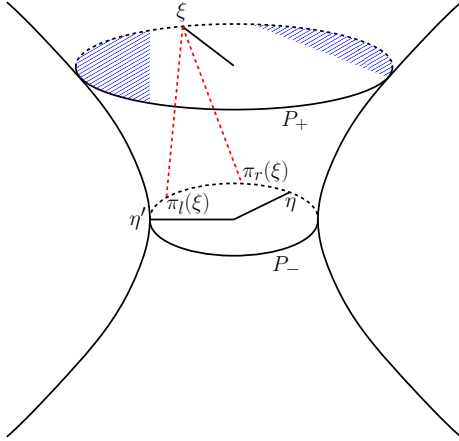


Figure 2.8. The choice of points η, ξ, η' in $\partial_\infty \text{AdS}^3$, endpoints at infinity of geodesic half-lines in the boundary of the convex hull.

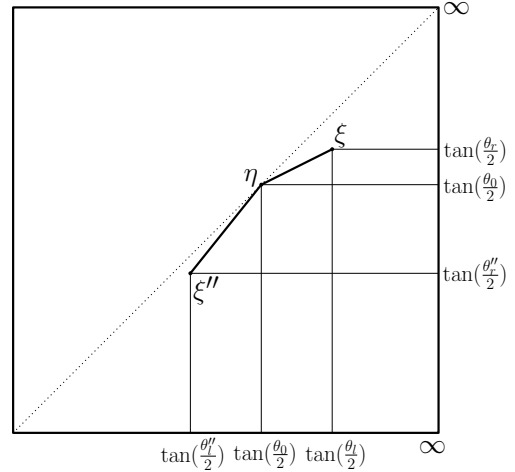


Figure 2.9. We give an upper bound on the ratio between the slopes of the two thick lines. The dotted line represents the plane P_- .

$S(\theta_l)$ is positive on $[\theta_0 + 2w, \pi]$ and the maximum S_{max} is achieved at some interior point of the interval. A computation gives

$$|cr(\phi(Q))| \leq S_{max} = \left(\frac{\cos(\theta_0/2 + w)}{\cos(\theta_0/2) + \sin(w)} \right)^2.$$

The RHS quantity depends on θ_0 , but is maximized on $[0, \pi - 2w]$ for $\theta_0 = 0$, where it assumes the value $(1 - \sin(w))/(1 + \sin(w))$. This gives

$$e^{\|\phi_\infty\|_{cr}} \geq \left| \frac{1}{cr(\phi(Q))} \right| \geq \frac{1 + \sin(w)}{1 - \sin(w)}.$$

From this we deduce

$$\sin(w) \leq \frac{e^{\|\phi_\infty\|_{cr}} - 1}{e^{\|\phi_\infty\|_{cr}} + 1} = \tanh \frac{\|\phi_\infty\|_{cr}}{2}$$

or equivalently

$$\tan(w) \leq \sinh \frac{\|\phi_\infty\|_{cr}}{2}.$$

Since $\|\phi_\infty\|_{cr} \leq \|\phi\|_{cr}$, the proof is concluded. \square

By using a very similar analysis, though simpler, we can prove an inequality in the converse direction.

Proposition 3.B. *Given any quasisymmetric homeomorphism ϕ of \mathbb{RP}^1 , let w be the width of the convex hull of the graph of ϕ in $\partial_\infty \text{AdS}^3$. Then*

$$\tanh \left(\frac{\|\phi\|_{cr}}{4} \right) \leq \tan w.$$

Proof. Suppose $\|\phi\|_{cr} > k$. Then we can find a quadruple of symmetric points Q such that $|cr(\phi(Q))| = e^k$. Consider the points ξ', η, ξ, η' on $\partial_\infty \text{AdS}^3$ such that their left and right projection are Q and $\phi(Q)$, respectively.

Recall that the isometries of AdS^3 act on $\partial_\infty \text{AdS}^3 \cong \mathbb{RP}^1 \times \mathbb{RP}^1$ as a pair of Möbius transformations, therefore they preserve the cross-ratio of both Q and $\phi(Q)$. Thus we can suppose $Q = (-1, 0, 1, \infty)$ and $\phi(Q) = (-e^{k/2}, 0, e^{-k/2}, \infty)$ when the quadruples are regarded as composed of points on $\mathbb{R} \cup \{\infty\}$.

Passing to the coordinates in S^1 (by the map $\theta \in S^1 \mapsto \tan(\theta/2) \in \mathbb{R}$) for this quadruple of points at infinity, it is easy to see that - in the affine chart $\{x^3 \neq 0\}$ - the position of the four points has an order 2 symmetry obtained by rotation around the z -axis. See Figure 2.10. This is ensured by the special renormalization chosen for Q and $\phi(Q)$.

Hence the geodesic line g_1 with endpoints at infinity η and η' is contained in the plane P_- as in the first part of the proof. More precisely, in the usual affine chart $\{x^3 \neq 0\}$,

$$g_1 = \{(\tanh(t), 0, 0) : t \in \mathbb{R}\}.$$

The geodesic line g_2 connecting ξ and ξ' has the form

$$g_2(s) = \left\{ \left(\frac{\cos(\alpha) \tanh(s)}{\cos(w')}, \frac{\sin(\alpha) \tanh(s)}{\cos(w')}, \tan(w') \right) : s \in \mathbb{R} \right\}.$$

The lines g_1 and g_2 are in the convex hull of $gr(\phi)$ and have the common orthogonal segment l which lies in the z -axis in the usual affine chart (Figure 2.10), the feet of l being achieved for $t = 0$ and $s = 0$.

The distance between g_1 and g_2 is achieved along this common orthogonal geodesic and its value is w' . Recalling Sublemma 2.4 and the computation in its proof, we find $\alpha = \theta_l - w' = \pi/2 - w'$ and $\theta_r = \theta_l - 2w'$. Since $\tan(\theta_r/2) = e^{-k/2}$ and $\theta_l = \pi/2$, one can compute

$$w' = \pi/4 - \arctan(e^{-k/2}).$$

It follows that

$$\tan w \geq \tan w' = \frac{1 - e^{-k/2}}{1 + e^{-k/2}} = \tanh\left(\frac{k}{4}\right).$$

Since this is true for an arbitrary $k \leq \|\phi\|_{cr}$, the inequality

$$\tan w \geq \tanh\left(\frac{\|\phi\|_{cr}}{4}\right)$$

holds. □

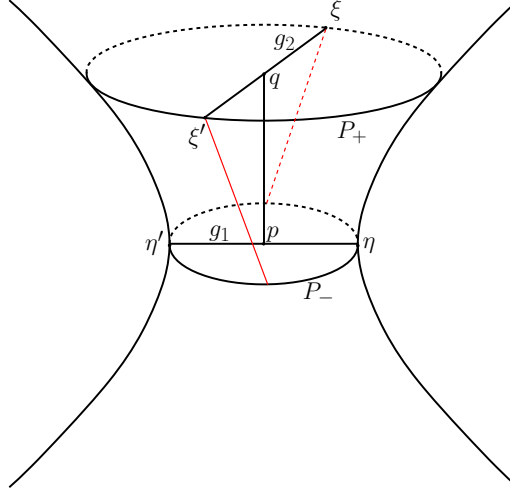


Figure 2.10. The distance between the two lines g and g' is achieved along the common orthogonal geodesic.

3. MAXIMAL SURFACES WITH SMALL PRINCIPAL CURVATURES

Let S be a maximal surface in AdS^3 . Let P_- be a spacelike plane which does not intersect the convex hull. We want to use the fact that the function $u(x) = \sin d_{\text{AdS}^3}(x, P_-)$, satisfies the equation

$$(L) \quad \Delta_S u - 2u = 0.$$

given in Proposition 1.8. This will enable us to use Equation (7) to give estimates on the principal curvatures of S .

3.1. Uniform gradient estimates. We start by obtaining some technical estimates on the gradient of u , which will have as a consequence that a maximal surface cannot be very “tilted” with respect to a plane outside the convex hull.

Lemma 3.1. *The universal constant $L = 2(1 + \sqrt{2})$ is such that, for every point x on a maximal surface of nonpositive curvature in AdS^3 , $\|\text{grad } u\|^2 < L$.*

Proof. Let γ be a path on S obtained by integrating the gradient vector field; more precisely, we impose $\gamma(0) = x$ and

$$\gamma'(t) = -\frac{\text{grad } u}{\|\text{grad } u\|}.$$

Observe that

$$u(\gamma(t)) - u(x) = \int_0^t du(\gamma'(s))ds = \int_0^t -\langle \text{grad } u(s), \frac{\text{grad } u(s)}{\|\text{grad } u(s)\|} \rangle ds = -\int_0^t \|\text{grad } u(s)\| ds.$$

We denote $y(t) = \|\text{grad } u(\gamma(t))\|$. We will show that $y(0)$ is bounded by a universal constant, using the fact that $u(\gamma(t))$ cannot become negative on S (recall Corollary 1.9). We have

$$(20) \quad \frac{d}{dt} \Big|_{t=0} y(t)^2 = 2\langle \nabla_{\gamma'(t)} \text{grad } u(\gamma(t)), \text{grad } u(\gamma(t)) \rangle = 2\langle \text{Hess}u(\gamma'(t)), \text{grad } u(\gamma(t)) \rangle.$$

Since, by equation (7), $\text{Hess}u - uE = \sqrt{1 - u^2 + \|\text{grad } u\|^2}B$ and $\|B(v)\| \leq \|v\|$,

$$\begin{aligned} -\frac{d}{dt} \Big|_{t=0} y(t)^2 &\leq \left| \frac{d}{dt} \Big|_{t=0} y(t)^2 \right| \leq 2\|\text{Hess}u(\gamma'(t))\| \|\text{grad } u(\gamma(t))\| \\ &\leq 2 \left(u(\gamma(t)) + \sqrt{1 - u(\gamma(t))^2 + y(t)^2} \right) y(t). \end{aligned}$$

Using that

$$\left(u(\gamma(t)) + \sqrt{1 - u(\gamma(t))^2 + y(t)^2} \right)^2 \leq 2(u(\gamma(t))^2 + 1 - u(\gamma(t))^2 + y(t)^2) = 2(1 + y(t)^2)$$

we obtain

$$(21) \quad -\frac{d}{dt} \Big|_{t=0} y(t) \leq \sqrt{2} \sqrt{1 + y(t)^2}.$$

It follows that

$$\int_0^t \frac{y'(s)}{\sqrt{1 + y(s)^2}} ds = \text{arcsinh } y(t) - \text{arcsinh } y(0) \geq \sqrt{2}t,$$

and therefore by a direct computation,

$$(22) \quad y(t) \geq y(0) \cosh(\sqrt{2}t) - \sqrt{1 + y(0)^2} \sinh(\sqrt{2}t).$$

Now

$$u(\gamma(t)) - u(x) = -\int_0^t y(s) ds \leq \frac{1}{\sqrt{2}} \left(-y(0) \sinh(\sqrt{2}t) + \sqrt{1 + y(0)^2} (\cosh(\sqrt{2}t) - 1) \right) =: F(t).$$

We must have $u(\gamma(t)) \geq 0$ for every t ; so we impose that $F(t) \geq -u(x)$ for every t . The minimum of F is achieved for

$$\tanh(\sqrt{2}t_{min}) = \frac{y(0)}{\sqrt{1 + y(0)^2}}.$$

Therefore

$$F(t_{min}) = -\frac{1}{\sqrt{2}} \left(1 + \sqrt{1 + y(0)^2} \right) \geq -u(x)$$

which is equivalent to $y(0)^2 \leq 2(u(x)^2 + \sqrt{2}u(x))$. Recalling $u \in [-1, 1]$, $\|\text{grad } u(x)\|^2 \leq 2(1 + \sqrt{2})$ independently on the maximal surface S and on the support plane P_- . \square

We now apply the above uniform gradient estimate to prove a fact which will be of use shortly. Given two unit timelike vectors $v, v' \in T_x \text{AdS}^3$ (both future-directed, or both past-directed), we define the hyperbolic angle between v and v' as the number $\alpha \geq 0$ such that $\cosh \alpha = |\langle v, v' \rangle|$. Compare with Figure 3.2 below.

Lemma 3.2. *There exists a constant $\bar{\alpha}$ such that the following holds for every maximal surface S in AdS^3 and every totally geodesic plane P_- in the past of S which does not intersect S . Let l be a geodesic line orthogonal to P_- and let $x = l \cap S$. Suppose x is at timelike distance less than $\pi/4$ from P_- . Then the hyperbolic angle α at x between l and the normal vector to S is bounded by $\alpha \leq \bar{\alpha}$.*

Proof. We use the same notation as Proposition 1.8. It is clear that the tangent direction to l is given by the vector ∇U , where $U(x) = \sin d_{\mathbb{A}\text{dS}^3}(x, P_-) = \langle x, p \rangle$ is defined on the entire $\mathbb{A}\text{dS}^3$ and p is the point dual to P_- . Recall u is the restriction of U to S . From Equations (9) and (11) of Proposition 1.8, we have

$$\langle \nabla U, \nabla U \rangle = -1 + u^2 = \|\text{grad } u\|^2 - \langle \nabla U, N \rangle^2.$$

It follows that the angle α at x between the normal to the maximal surface S and the geodesic l can be computed as

$$(\cosh \alpha)^2 = \left\langle \frac{\nabla U(x)}{\|\nabla U(x)\|}, N \right\rangle^2 = \frac{1 - u(x)^2 + \|\text{grad } u(x)\|^2}{1 - u(x)^2}$$

and so α is bounded by Lemma 3.1 and the assumption that $u(x)^2 \leq 1/2$. \square

3.2. Schauder estimates. We move to proving Schauder-type estimates on the derivatives of the function $u = \sin d_{\mathbb{A}\text{dS}^3}(\cdot, P_-)$, expressed in suitable coordinates, of the form

$$\|u\|_{C^2(B_0(0, \frac{R}{2}))} \leq C \|u\|_{C^0(B_0(0, R))}$$

where the constant does not depend on S and P_- . Here $B_0(0, R)$ denotes the Euclidean ball centered at 0 of radius R .

Denote by w the width of the convex hull of $\partial_\infty S = gr(\phi)$; recall that $w \leq \pi/2$. Let x be a point of S . By Remark 1.7, we have that $d_{\mathbb{A}\text{dS}^3}(x, \partial_- \mathcal{C}) + d_{\mathbb{A}\text{dS}^3}(x, \partial_+ \mathcal{C}) \leq w$, therefore one among $d_{\mathbb{A}\text{dS}^3}(x, \partial_- \mathcal{C})$ and $d_{\mathbb{A}\text{dS}^3}(x, \partial_+ \mathcal{C})$ must be smaller than $\pi/4$. Composing with an isometry of $\mathbb{A}\text{dS}^3$ (which possibly reverses time-orientation), we can assume $d_{\mathbb{A}\text{dS}^3}(x, \partial_- \mathcal{C}) \leq d_{\mathbb{A}\text{dS}^3}(x, \partial_+ \mathcal{C})$, which implies that x has distance less than $\pi/4$ from P_- . This assumption will be important in the following.

Recall from Subsection 1.7 that, if l is a timelike line orthogonal to the plane P_0 at x_0 , the solid cylinder $Cl(x_0, P_0, R_0)$ is the set of points $x \in \mathbb{A}\text{dS}^3$ which lie on a spacelike plane P_x orthogonal to l such that $d_{P_x}(x, l \cap P_x) \leq R_0$. See also Figure 1.5.

Proposition 4.A. *There exists a radius $R > 0$ and a constant $C > 0$ such that for every choice of:*

- *A maximal surface $S \subset \mathbb{A}\text{dS}^3$ with $\partial_\infty S$ the graph of an orientation-preserving homeomorphism;*
- *A point $x \in S$;*
- *A plane P_- disjoint from S with $d_{\mathbb{A}\text{dS}^3}(x, P_-) \leq \pi/4$,*

the function $u(\cdot) = \sin d_{\mathbb{A}\text{dS}^3}(\cdot, P_-)$ expressed in terms of normal coordinates centered at x , namely

$$u(z) = \sin d_{\mathbb{A}\text{dS}^3}(\exp_x(z), P_-)$$

where $\exp_x : \mathbb{R}^2 \cong T_x S \rightarrow S$ denotes the exponential map, satisfies the Schauder-type inequality

$$(23) \quad \|u\|_{C^2(B_0(0, \frac{R}{2}))} \leq C \|u\|_{C^0(B_0(0, R))}.$$

Proof. Fix a radius $R_0 > 0$. First, we show that there exists a radius $R > 0$ such that the image of the Euclidean ball $B_0(0, R)$ under the exponential map at every point $x \in S$, for every surface S , is contained in the solid cylinder $Cl(x, T_x S, R_0)$. Indeed, suppose this does not hold, namely

$$(24) \quad \inf_{x \in S} \sup \{R : \exp_x(B_0(0, R)) \subset Cl(x, T_x S, R_0)\} = 0.$$

Then one can find a sequence S_n of maximal surfaces and points x_n such that, if R_n is the supremum of those radii R for which $\exp_{x_n}(B_0(0, R))$ is contained in the respective cylinder of radius R_0 , then R_n goes to zero. We can compose with isometries of $\mathbb{A}\text{dS}^3$ so that all points x_n are sent to the same point x_0 and all surfaces are tangent at x_0 to the same plane P_0 . By Lemma 1.20, there exists a subsequence converging inside $Cl(x_0, P_0, R_0)$ to a maximal

surface S_∞ . Therefore the infimum in the LHS of Equation (24) cannot be zero, since for the limiting surface S_∞ there is a radius R_∞ such that $\exp_x(B_0(0, R_\infty)) \subset Cl(x, T_x S, R_0)$.

We use a similar argument to prove the main statement. We can consider P_- a fixed plane, and a point $x \in S$ lying on a fixed geodesic l orthogonal to P_- . Suppose the claim does not hold, namely there exists a sequence of surfaces S_n in the future of P_- such that for the function $u_n(z) = \sin d_{\text{AdS}^3}(\exp_{x_n}(z), P_n)$,

$$\|u_n\|_{C^2(B_0(0, \frac{R}{2}))} \geq n \|u\|_{C^0(B_0(0, R))}.$$

Let us compose each S_n with an isometry $T_n \in \text{Isom}(\text{AdS}^3)$ so that $S'_n = T_n(S_n)$ is tangent at a fixed point x_0 to a fixed plane P_0 , whose normal unit vector is N_0 .

We claim that the sequence of isometries T_n is bounded in $\text{Isom}(\text{AdS}^3)$, since T_n^{-1} maps the element (x_0, N_0) of the tangent bundle $T\text{AdS}^3$ to a bounded region of $T\text{AdS}^3$. Indeed, by our assumptions, $T_n^{-1}(x_0) = x_n$ lies on a geodesic l orthogonal to P_- and has distance less than $\pi/4$ (in the future) from P_- ; moreover by Lemma 3.2 the vector $(dT_n)^{-1}(N_0)$ forms a bounded angle with l .

By Lemma 1.20, up to extracting a subsequence, we can assume $S'_n \rightarrow S'_\infty$ on $Cl(x_0, P_0, R_0)$ with all derivatives. Since we can also extract a converging subsequence from T_n , we assume $T_n \rightarrow T_\infty$, where T_∞ is an isometry of AdS^3 . Therefore $T_n(P_-)$ converges to a totally geodesic plane P_∞ .

Using the first part of this proof and Lemma 1.20, on the image of the ball $B_0(0, R)$ under the exponential map of S'_n , the coefficients of the Laplace-Beltrami operators $\Delta_{S'_n}$ (in normal coordinates on $B_0(0, R)$) converge to the coefficients of $\Delta_{S'_\infty}$. Hence the operators $\Delta_{S'_n} - 2$ are uniformly strictly elliptic with uniformly bounded coefficients. By classical Schauder estimates (see [GT83]), using the fact that u_n solves the equation $\Delta_{S'_n}(u_n) - 2u_n = 0$, there exists a constant c such that

$$\|u_n\|_{C^2(B_0(0, \frac{R}{2}))} \leq c \|u_n\|_{C^0(B_0(0, R))},$$

for every n . This gives a contradiction. \square

Remark 3.3. The statements of Lemma 3.2 and Proposition 4.A could be improved so as to be stated in terms of the choice of any radius $R > 0$, any number $w_0 < \pi/2$ (replacing $\pi/4$), where the constant C would depend on such choices. Similarly for Proposition 3.4 below. However, these details would not improve the final statement of Theorem 1.A and thus are not pursued here.

We are therefore in a good point to obtain an estimate of the second derivatives of u (and thus for the principal curvatures of S) in terms of the width w . However, let us remark that in Anti-de Sitter space the projection from a spacelike curve or surface to a totally geodesic spacelike plane is not distance-contracting. Hence we need to give an additional computation in order to ensure (by substituting the radius R in Proposition 4.A by a smaller one if necessary) that the projection from the geodesic balls $B_S(x, R)$ to P_- has image contained in a uniformly bounded set. This is proved in the next Proposition, see also Figure 3.1.

Proposition 3.4. *There exist constant radii R'_0 and R' such that for every maximal surface of nonpositive curvature S in AdS^3 , every point $x_0 \in S$ and every totally geodesic plane P_- which does not intersect S , such that the distance of x_0 from P_- is at most $\pi/4$, the orthogonal projection $\pi|_S : S \rightarrow P_-$ maps $S \cap Cl(x_0, T_{x_0}S, R'_0)$ to $B_{P_-}(\pi(x_0), R')$.*

Proof. We can suppose $T_{x_0}S$ is the intersection of the plane $\{x_4 = 0\}$ with AdS^3 and $x_0 = [0, 0, 1, 0]$. As in the coordinate system (15), the points x in $Cl(x_0, T_{x_0}S, R'_0)$ have coordinates

$$x = [\cos \theta \sinh r, \sin \theta \sinh r, \cos \zeta \cosh r, \sin \zeta \cosh r],$$

for $r \leq R'_0$. Let us denote by $I^+(p)$ (resp. $I^-(x_0)$) the cone of points connected to p by a future-directed (resp. past-directed) timelike path in $\text{AdS}^3 \setminus Q$, where Q is the plane at infinity in the affine chart.

Since S is spacelike, $S \cap Cl(x_0, T_{x_0}S, R'_0)$ is contained in $Cl(x_0, T_{x_0}S, R'_0) \setminus (I^+(x_0) \cup I^-(x_0))$. See also Figure 1.5. Hence $|\langle x, x_0 \rangle| > 1$ (recall Equation (3) in Section 1), which is equivalent to

$$(25) \quad |\cos \zeta| > \frac{1}{\cosh r}.$$

Let l be the geodesic through x_0 orthogonal to P_- . We will conduct the computation in the double cover $\widehat{\text{AdS}}^3 \subset \mathbb{R}^{2,2}$. We can assume l has tangent vector at x_0 given by $l'(0) = (\sinh \alpha, 0, 0, \cosh \alpha)$, where of course α is the angle between l and the normal to S at x_0 . Therefore

$$l(t) = (\cos t)x_0 + (\sin t)l'(0) = (\sin t \sinh \alpha, 0, \cos t, \sin t \cosh \alpha).$$

Let $w_1 = d_{\widehat{\text{AdS}}^3}(x_0, P_-)$, so $P_- = p^\perp$, where p is the point

$$p = l'(-w_1) = (\cos w_1 \sinh \alpha, 0, \sin w_1, \cos w_1 \cosh \alpha).$$

The projection of x to P_- is given by

$$\pi(x) = \frac{x + \langle x, p \rangle p}{\sqrt{1 - \langle x, p \rangle^2}}$$

provided $\langle x, p \rangle^2 < 1$, which is the condition for x to be in the domain of dependence of P_- . (We say that x is in the domain of dependence of P_- if the dual plane of x is disjoint from P_- .) The distance d between $\pi(x)$ and $\pi(x_0) = l(-w_1)$ is given by the expression

$$(26) \quad \cosh d = |\langle \pi(x), l(-w_1) \rangle| = \left| \frac{\langle x, l(-w_1) \rangle}{\sqrt{1 - \langle x, p \rangle^2}} \right|.$$

Now, we have

$$\begin{aligned} |\langle x, p \rangle| &= |\cos \theta \sinh r \cos w_1 \sinh \alpha - \cos \zeta \cosh r \sin w_1 - \sin \zeta \cosh r \cos w_1 \cosh \alpha| \\ &\leq \sinh r \sinh \alpha + \frac{\sqrt{2}}{2} \cosh r + \sinh r \cosh \alpha = \frac{\sqrt{2}}{2} \cosh r + (\sinh r)e^\alpha. \end{aligned}$$

In the last line, we have used that $|\sin \zeta| = \sqrt{1 - (\cos \zeta)^2} \leq \tanh r$, by Equation (25), and that $\sin w_1 < \sqrt{2}/2$. Since the hyperbolic angle α is uniformly bounded by Lemma 3.2 (Figure 3.2), it follows that if $r \leq R'_0$ for R'_0 sufficiently small, $\sqrt{1 - \langle x, p \rangle^2}$ is uniformly bounded from below. Moreover,

$$\begin{aligned} |\langle x, l(-w_1) \rangle| &= |-\cos \theta \sinh r \sin w_1 \sinh \alpha - \cos \zeta \cosh r \cos w_1 + \sin \zeta \cosh r \sin w_1 \cosh \alpha| \\ &\leq \sinh r \sinh \alpha + \cosh r + \cosh r \cosh \alpha \end{aligned}$$

is uniformly bounded. This shows, from Equation (26), that $\cosh d \leq \cosh R'$ for some constant radius R' (depending on R'_0). This concludes the proof. \square

Therefore, replacing R_0 in Lemma 1.20 with $\min\{R_0, R'_0\}$, we have that the geodesic balls of radius R (R as in Proposition 4.A) on S centered at x project to P_- with image contained in $B_{P_-}(\pi(x), R')$. The radii R and R' are fixed, not depending on S .

3.3. Principal curvatures. In this subsection we finally prove the estimate on the supremum of the principal curvatures of S in terms of the width. Recall the statement of the main theorem of this section:

Theorem 1.A. *There exists a constant C_1 such that, for every maximal surface S with $\|\lambda\|_\infty < 1$ and width w ,*

$$\|\lambda\|_\infty \leq C_1 \tan w.$$

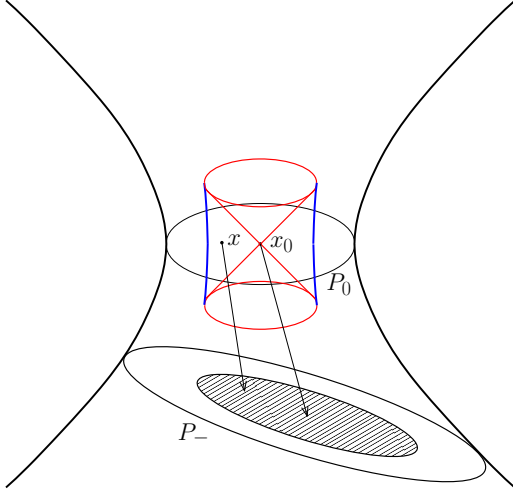


Figure 3.1. Projection from points in $Cl(x_0, T_{x_0}S, R'_0)$ which are connected to x_0 by a spacelike geodesic have bounded image.

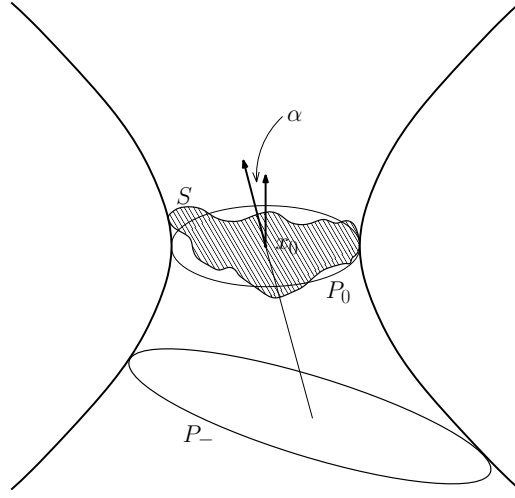


Figure 3.2. The key point is that the hyperbolic angle α is uniformly bounded, by Lemma 3.2.

We take an arbitrary point $x \in S$. By Remark 1.7, we know that there are two disjoint planes P_- and P_+ with $d_{\text{AdS}^3}(x, P_-) + d_{\text{AdS}^3}(x, P_+) = w_1 + w_2 \leq w$ where w is the width. As in the previous subsection, we will assume P_- is a fixed plane in AdS^3 , upon composing with an isometry. Figure 3.3 gives a picture of the situation of the following lemma.

Lemma 3.5. *Let $p \in P_-$, $q \in P_+$ be the endpoints of geodesic segments l_1 and l_2 from $x \in S$ orthogonal to P_- and P_+ , of length w_1 and w_2 , with $w_1 \leq w_2$. Let $p' \in P_-$ a point at distance R' from p and let $d = d_{\text{AdS}^3}((\pi|_{P_+})^{-1}(p'), P_-)$. Then*

$$(27) \quad \tan d \leq (1 + \sqrt{2}) \cosh R' \tan(w_1 + w_2).$$

Proof. As in the previous proof, we do the computation in $\widehat{\text{AdS}^3}$. We assume $x = (0, 0, 1, 0)$ and l_1 is the geodesic segment parametrized by $l_1(t) = (\cos t)x - (\sin t)(0, 0, 0, 1)$, so that the plane P_- is dual to $p_- = (0, 0, \sin w_1, \cos w_1)$. Points on the plane P_- at distance R' from $\pi(x) = l_1(w_1) = (0, 0, \cos w_1, -\sin w_1)$ have coordinates

$$p' = (\cos \theta \sinh R', \sin \theta \sinh R', \cosh R' \cos w_1, -\cosh R' \sin w_1).$$

We also assume l_2 has initial tangent vector $l'_2(0) = (\sinh \alpha, 0, 0, \cosh \alpha)$, where α is the hyperbolic angle between $(0, 0, 0, 1)$ and $l'_2(0)$, so that $l_2(t) = (\cos t)x + (\sin t)(\sinh \alpha, 0, 0, \cosh \alpha)$. Note that $l'_2(w_2) = (\cos w_2 \sinh \alpha, 0, -\sin w_2, \cos w_2 \cosh \alpha) =: p_+$ is the unit vector orthogonal to P_+ , by construction.

We derive a condition which must necessarily be satisfied by α , because P_- and P_+ are disjoint. Indeed, we must have

$$|\langle p_-, p_+ \rangle| = -\sin w_1 \sin w_2 + \cos w_1 \cos w_2 \cosh \alpha \leq 1,$$

which is equivalent to

$$(28) \quad \cosh \alpha < \frac{1 + \sin w_1 \sin w_2}{\cos w_1 \cos w_2}.$$

Let us now write

$$(\tanh \alpha)^2 = \left(1 + \frac{1}{\cosh \alpha}\right) \left(1 - \frac{1}{\cosh \alpha}\right) \leq 2(\cosh \alpha - 1),$$

and therefore, using (28),

$$(29) \quad (\tanh \alpha)^2 < 2 \left(\frac{1 - \cos(w_1 + w_2)}{\cos w_1 \cos w_2} \right) \leq 2 \left(\frac{1 - (\cos(w_1 + w_2))^2}{\cos w_1 \cos w_2} \right) \leq 2 \frac{(\sin(w_1 + w_2))^2}{\cos w_1 \cos w_2}.$$

To compute d , we now write explicitly the geodesic γ starting from p' and orthogonal to P_- . We find d such that $\gamma(d) \in P_+$ and this will give the expected inequality. We have

$$\gamma(d) = (\cos d)p' + (\sin d)(0, 0, \sin w_1, \cos w_1)$$

and $\gamma(d) \in P_+$ if and only if $\langle \gamma(d), p_+ \rangle = 0$, which gives the condition

$$\begin{aligned} \cos d(\cosh R'(\cos w_1 \sin w_2 + \cos w_2 \sin w_1 \cosh \alpha) + \sinh R'(\cos \theta \cos w_2 \sinh \alpha)) \\ + \sin d(\sin w_1 \sin w_2 - \cos w_1 \cos w_2 \cosh \alpha) = 0. \end{aligned}$$

We express

$$\begin{aligned} \tan d = \cosh R' \frac{\cos w_1 \sin w_2 + \cos w_2 \sin w_1 \cosh \alpha}{\cos w_1 \cos w_2 \cosh \alpha - \sin w_1 \sin w_2} \\ + \sinh R' \frac{\cos \theta \cos w_2 \sinh \alpha}{\cos w_1 \cos w_2 \cosh \alpha - \sin w_1 \sin w_2}. \end{aligned}$$

The first term in the RHS is easily seen to be less than $\cosh R' \tan(w_1 + w_2)$. We turn to the second term. Using (29), it is bounded by

$$\sinh R' \tanh \alpha \frac{\cos w_2}{\cos(w_1 + w_2)} \leq \sqrt{2} \sinh R' \tan(w_1 + w_2) \left(\frac{\cos w_2}{\cos w_1} \right)^{\frac{1}{2}}.$$

In conclusion, having assumed $w_1 \leq w_2$, we can put $\cos(w_2)/\cos(w_1) \leq 1$, sum the two terms and get

$$\tan d \leq (1 + \sqrt{2}) \cosh R' \tan(w_1 + w_2).$$

□

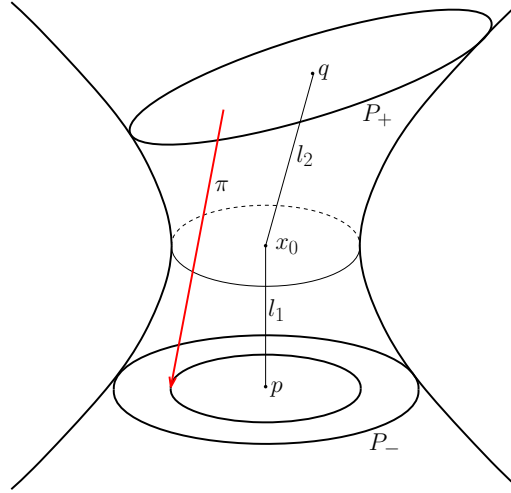


Figure 3.3. The setting of Lemma 3.5. We assume $w_1 = d_{\text{AdS}^3}(x_0, p) < d_{\text{AdS}^3}(x_0, q) = w_2$.

Proof of Theorem 1.A. Let $x \in S$ and consider the point x_- of $\partial_- \mathcal{C}$ which minimizes the distance from x , where \mathcal{C} is the convex hull of S . Let P_- be the plane through x_- orthogonal to the geodesic line containing x and x_- (recall Remark 1.7). The plane P_- is then a support plane of $\partial_- \mathcal{C}$. We construct analogously the support plane P_+ for $\partial_+ \mathcal{C}$. As discussed in Remark 1.7,

$$d_{\text{AdS}^3}(x, P_-) + d_{\text{AdS}^3}(x, P_+) \leq w.$$

Moreover, we can assume (upon composing with a time-orientation-reversing isometry, if necessary) that $d_{\text{AdS}^3}(x, P_-) \leq d_{\text{AdS}^3}(x, P_+)$. As a consequence, $d_{\text{AdS}^3}(x, P_-) \leq \pi/4$.

Let us now consider the function u , defined in normal coordinates:

$$u = \sinh d_{\text{AdS}^3}(\exp_x(\cdot), P_-).$$

By Equation (7), we have the following expression for the shape operator of S :

$$B = \frac{1}{\sqrt{1 - u^2 + \|\text{grad } u\|^2}} (\text{Hess } u - u E).$$

In normal coordinates at x the Hessian of u is given just by the second derivatives of u ; in Proposition 4.A we showed the second derivatives of u are bounded, up to a factor, by $\|u\|_{C^0(B_S(x, R))}$. By Proposition 3.4, $\|u\|_{C^0(B_S(x, R))}$ is smaller than the supremum of the hyperbolic sine of the distance d from P_- of points of S which project to $B_{P_-}(\pi(x), R')$. Therefore we have the following estimate for the principal curvatures at x :

$$|\lambda| \leq C' \frac{\|u\|}{\sqrt{1 - \|u\|^2}} \leq C' \tan(\sup\{d_{\text{AdS}^3}(p, P_-) : p \in (\pi_S)^{-1}(B_{P_-}(\pi(x), R'))\}).$$

The constant C' involves the constant which appears in Equation (23) in Proposition 4.A. The quantity in brackets in the RHS is certainly less than

$$(\sup\{d_{\text{AdS}^3}(p, P_-) : p \in (\pi_{P_+})^{-1}(B_{P_-}(\pi(x), R'))\}).$$

Thus, applying Lemma 3.5 we obtain:

$$\|\lambda\|_\infty \leq C_1 \tan w.$$

The constant C then involves C' and $\cosh R'$. Such inequality holds independently on the point x and thus concludes the proof. \square

4. MAXIMAL SURFACES WITH LARGE PRINCIPAL CURVATURES

The main goal of this section is to prove Theorem 1.B. Recall that we denote by $\|\lambda\|_\infty$ the supremum of the principal curvatures of a maximal surface S with $\partial_\infty S = gr(\phi) \subset \partial_\infty \text{AdS}^3$, and by w the width of the convex hull of $\partial_\infty S$.

Theorem 1.B. *There exist universal constants $M > 0$ and $\delta \in (0, 1)$ such that, if S is an entire maximal surface in AdS^3 with $\delta \leq \|\lambda\|_\infty < 1$ and width w , then*

$$\tan w \geq \left(\frac{1}{1 - \|\lambda\|_\infty} \right)^{1/M}.$$

Observe that it clearly suffices to prove that for every $\lambda_0 \geq \delta$, if there exists a point $x_0 \in S$ such that $\lambda(x_0) = \lambda_0$, then

$$(30) \quad \tan w \geq \left(\frac{1}{1 - \lambda_0} \right)^{1/M}.$$

Indeed, if $\|\lambda\|_\infty$ is not achieved on S , one has that Equation (30) holds for every $\lambda_0 < \|\lambda\|_\infty$, and by continuity (30) holds for $\lambda_0 = \|\lambda\|_\infty$ as well.

4.1. Uniform gradient estimates. We define the function $v : S \rightarrow \mathbb{R}$ by

$$v(x) = \ln \left(\frac{1}{1 - \lambda(x)} \right),$$

where $\lambda(x)$ is the positive eigenvalue of the shape operator. Observe that $v(x) \rightarrow \infty$ as $\lambda(x)$ approaches 1, while $v(x) = 0$ if x is an umbilical point, and the statement of Theorem 1.B amounts to proving that $\tan w \geq e^{v(x_0)/M}$. The first important step is to show that, if v is large at some point, then it remains large on a geodesic ball of S with large radius.

Proposition 4.1. *There exists a universal constant M such that $\|\text{grad } v\| \leq M$, for every maximal surface S in AdS^3 with uniformly negative curvature.*

We will prove some *a priori* estimates for the function $\chi = -\ln|\lambda|$, which is defined in the complement of umbilical points and it takes values in $[0, +\infty)$. Recall that by Lemma 1.13 the function χ satisfies the quasi-linear equation:

$$(Q) \quad \Delta_S \chi = 2(1 - e^{-2\chi}),$$

while by the Gauss equation the curvature of S is given by:

$$(31) \quad K = -1 + e^{-2\chi}.$$

We start by computing the gradient and the Laplacian of the function $\|\text{grad } \chi\|^2$.

Lemma 4.2. *The following identities hold:*

$$(32) \quad \text{grad}(\|\text{grad } \chi\|^2) = 2\text{Hess}\chi(\text{grad } \chi),$$

$$(33) \quad \Delta_S(\|\text{grad } \chi\|^2) = (10e^{-2\chi} - 2)\|\text{grad } \chi\|^2 + 2\|\text{Hess}\chi\|^2.$$

Proof. In order to compute the gradient, notice that if ξ is any vector tangent to S , then

$$\xi(\|\text{grad } \chi\|^2) = 2\langle \nabla_\xi(\text{grad } \chi), \text{grad } \chi \rangle = 2\langle \text{Hess}\chi(\xi), \text{grad } \chi \rangle = 2\langle \xi, \text{Hess}\chi(\text{grad } \chi) \rangle.$$

This proves the first identity. By taking the covariant derivative of $\text{grad } \|\text{grad } \chi\|^2$ we get

$$\begin{aligned} \text{Hess}(\|\text{grad } \chi\|^2) &= 2\nabla(\text{Hess}\chi)(\text{grad } \chi) + 2(\text{Hess}\chi)^2 \\ &= 2\nabla_\bullet(\nabla \text{grad } \chi)(\text{grad } \chi) + 2(\text{Hess}\chi)^2 \\ &= 2\nabla_{\text{grad } \chi}(\text{Hess}\chi) + 2R(\bullet, \text{grad } \chi) \text{grad } \chi + 2(\text{Hess}\chi)^2. \end{aligned}$$

Taking the trace we get

$$\Delta_S(\|\text{grad } \chi\|^2) = 2\text{grad } \chi(\Delta_S \chi) + 2K\|\text{grad } \chi\|^2 + 2\|\text{Hess}\chi\|^2,$$

and using (Q) and (31) we deduce that

$$\begin{aligned} \Delta_S(\|\text{grad } \chi\|^2) &= 4\langle \text{grad } \chi, \text{grad}(1 - e^{-2\chi}) \rangle + 2(-1 + e^{-2\chi})\|\text{grad } \chi\|^2 + 2\|\text{Hess}\chi\|^2 \\ &= 4 \cdot 2e^{-2\chi}\|\text{grad } \chi\|^2 + 2(-1 + e^{-2\chi})\|\text{grad } \chi\|^2 + 2\|\text{Hess}\chi\|^2 \\ &= (10e^{-2\chi} - 2)\|\text{grad } \chi\|^2 + 2\|\text{Hess}\chi\|^2. \end{aligned}$$

This concludes the proof. □

Lemma 4.3. *Let $g : (0, +\infty) \rightarrow \mathbb{R}$ any smooth function and consider the function $\psi : S \rightarrow \mathbb{R}$ defined by*

$$\psi = e^{g(\chi)}\|\text{grad } \chi\|^2.$$

Then we have

$$(34) \quad \Delta_S \psi \geq a(\chi)\psi^2 + b(\chi)\psi + c(\chi),$$

where

$$(35) \quad \begin{aligned} a(\chi) &= g''e^{-g}, \\ b(\chi) &= 6g'(\chi)(1 - e^{-2\chi}) + (10e^{-2\chi} - 2), \\ c(\chi) &= e^g(\Delta_S \chi)^2. \end{aligned}$$

Proof. Let us put $\alpha = e^g$, then (using (32)) we have

$$\text{grad } \psi = \alpha' \|\text{grad } \chi\|^2 \text{grad } \chi + 2\alpha \text{Hess}\chi(\text{grad } \chi).$$

Differentiating again we get

$$\begin{aligned} \text{Hess } \psi &= \alpha'' \|\text{grad } \chi\|^2 \text{grad } \chi \otimes \text{grad } \chi + \alpha' \|\text{grad } \chi\|^2 \text{Hess}\chi + 2\alpha' \text{grad } \chi \otimes (\text{Hess}\chi(\text{grad } \chi)) \\ &\quad + 2\alpha' (\text{Hess}\chi(\text{grad } \chi)) \otimes \text{grad } \chi + \alpha \text{Hess}(\|\text{grad } \chi\|^2), \end{aligned}$$

so by using (Q) and (33) we get

$$\begin{aligned} \Delta_S \psi &= \alpha'' \|\text{grad } \chi\|^4 + 2\alpha'(1 - e^{-2x}) \|\text{grad } \chi\|^2 + 4\alpha' \langle \text{grad } \chi, \text{Hess}\chi(\text{grad } \chi) \rangle \\ &\quad + \alpha(10e^{-2x} - 2) \|\text{grad } \chi\|^2 + 2\alpha \|\text{Hess}\chi\|^2 \\ &= \alpha'' \|\text{grad } \chi\|^4 + (2\alpha'(1 - e^{-2x}) + \alpha(10e^{-2x} - 2)) \|\text{grad } \chi\|^2 + \\ &\quad + 4\alpha' \langle \text{grad } \chi, \text{Hess}\chi(\text{grad } \chi) \rangle + 2\alpha \|\text{Hess}\chi\|^2. \end{aligned}$$

The main term to be estimated is the scalar product $\langle \text{grad } \chi, \text{Hess}\chi(\text{grad } \chi) \rangle$. Let μ_1, μ_2 be the eigenvalues of $\text{Hess}\chi$ and x_1, x_2 be the coordinates of $\text{grad } \chi$ with respect to an orthonormal basis of eigenvectors of $\text{Hess}\chi$. Then we have

$$\langle \text{grad } \chi, \text{Hess}\chi(\text{grad } \chi) \rangle = \mu_1 x_1^2 + \mu_2 x_2^2.$$

Notice that $\Delta_S \chi = \mu_1 + \mu_2$, so if we put $\xi = \mu_1 - \mu_2$ we deduce that

$$\langle \text{grad } \chi, \text{Hess}\chi(\text{grad } \chi) \rangle = \frac{\Delta_S \chi}{2} \|\text{grad } \chi\|^2 + \frac{\xi}{2} (x_1^2 - x_2^2).$$

So we get

$$(36) \quad 4\alpha' \langle \text{grad } \chi, \text{Hess}\chi(\text{grad } \chi) \rangle = 4\alpha'(1 - e^{-2x}) \|\text{grad } \chi\|^2 + 2\alpha' \xi (x_1^2 - x_2^2).$$

Putting (36) into the equality we obtained previously yields

$$\begin{aligned} \Delta_S \psi &= \alpha'' \|\text{grad } \chi\|^4 + (2\alpha'(1 - e^{-2x}) + \alpha(10e^{-2x} - 2)) \|\text{grad } \chi\|^2 \\ &\quad + 4\alpha'(1 - e^{-2x}) \|\text{grad } \chi\|^2 + 2\alpha' \xi (x_1^2 - x_2^2) + 2\alpha \|\text{Hess}\chi\|^2 \\ &= \alpha'' \|\text{grad } \chi\|^4 + (6\alpha'(1 - e^{-2x}) + \alpha(10e^{-2x} - 2)) \|\text{grad } \chi\|^2 \\ &\quad + 2\alpha' \xi (x_1^2 - x_2^2) + 2\alpha \|\text{Hess}\chi\|^2. \end{aligned}$$

Now observe that

$$|2\alpha' \xi (x_1^2 - x_2^2)| \leq 2|\alpha'| |\xi| \|\text{grad } \chi\|^2 \leq \frac{(\alpha')^2}{\alpha} \|\text{grad } \chi\|^4 + \alpha \xi^2,$$

hence we obtain

$$\begin{aligned} \Delta_S \psi &\geq \alpha'' \|\text{grad } \chi\|^4 + (6\alpha'(1 - e^{-2x}) + \alpha(10e^{-2x} - 2)) \|\text{grad } \chi\|^2 \\ &\quad - \frac{(\alpha')^2}{\alpha} \|\text{grad } \chi\|^4 - \alpha \xi^2 + 2\alpha \|\text{Hess}\chi\|^2 = \\ &= \left(\alpha'' - \frac{(\alpha')^2}{\alpha} \right) \|\text{grad } \chi\|^4 + (6\alpha'(1 - e^{-2x}) + \alpha(10e^{-2x} - 2)) \|\text{grad } \chi\|^2 \\ &\quad + \alpha(2\|\text{Hess}\chi\|^2 - \xi^2). \end{aligned}$$

But $\|\text{Hess}\chi\|^2 = \mu_1^2 + \mu_2^2$, so that

$$2\|\text{Hess}\chi\|^2 - \xi^2 = 2(\mu_1^2 + \mu_2^2) - (\mu_1 - \mu_2)^2 = (\mu_1 + \mu_2)^2 = (\Delta_S \chi)^2.$$

We finally obtain

$$\Delta_S \psi \geq \left(\alpha'' - \frac{(\alpha')^2}{\alpha} \right) \|\text{grad } \chi\|^4 + (6\alpha'(1 - e^{-2x}) + \alpha(10e^{-2x} - 2)) \|\text{grad } \chi\|^2 + \alpha(\Delta_S \chi)^2.$$

Observing that $\alpha' = g'e^g$ and $\alpha'' = (g'' + (g')^2)e^g$, we get $\alpha'' - (\alpha')^2/\alpha = g''e^g$ and thus

$$\begin{aligned}\Delta_S \psi &\geq g''e^g \|\text{grad } \chi\|^4 + (6g'(1 - e^{-2x}) + (10e^{-2x} - 2))e^g \|\text{grad } \chi\|^2 + \alpha(\Delta_S \chi)^2 \\ &= g''e^{-g}\psi^2 + (6g'(1 - e^{-2x}) + (10e^{-2x} - 2))\psi + e^g(\Delta_S \chi)^2,\end{aligned}$$

as in the statement. \square

Remark 4.4. In the hypothesis of Lemma 4.3, suppose that g is a convex function so that $a(\chi) > 0$. If $\sup_{\chi>0}(-b(\chi)/a(\chi)) = M_1 < +\infty$ we have that $\Delta_S \psi > 0$ whenever $\psi > M_1$.

Lemma 4.5. *There is a constant M_2 such that $\|\text{grad } \lambda\| < M_2$ for any maximal surface of nonpositive curvature in AdS^3 .*

Proof. Take a point $p \in S$ and consider normal coordinates x, y centered at p . We have that

$$B(\partial_x) = a(x, y)\partial_x + b(x, y)\partial_y,$$

where a, b are smooth functions in a neighborhood of p . Since B is traceless self-adjoint, it turns out that

$$\lambda = \sqrt{a^2 + b^2},$$

so that

$$\text{grad } \lambda = \frac{1}{\sqrt{a^2 + b^2}}(a \text{grad } a + b \text{grad } b).$$

In particular

$$\|\text{grad } \lambda\|^2 \leq \|\text{grad } a\|^2 + \|\text{grad } b\|^2.$$

On the other hand, at the point p we have

$$(\nabla B)(\partial_x) = \nabla(B(\partial_x)) = (\text{grad } a) \otimes (\partial_x) + (\text{grad } b) \otimes (\partial_y),$$

so we deduce that at the point p

$$\|\text{grad } \lambda\|^2 \leq \|\nabla B\|^2.$$

By applying Lemma 1.20, it is not difficult to show that there exists a universal constant M_2 such that $\|\nabla B\| \leq M_2$ for any point of any maximal surface with nonpositive curvature. \square

Proof of Proposition 4.1. We will derive an a-priori bound on $\|\text{grad } v\|$ where

$$v = -\ln(1 - |\lambda|) = -\ln(1 - e^{-x}).$$

Indeed we have that

$$(37) \quad \|\text{grad } v\|^2 = \frac{1}{(1 - e^{-x})^2} \|\text{grad } \lambda\|^2 = \frac{e^{-2x}}{(1 - e^{-x})^2} \|\text{grad } \chi\|^2,$$

hence in particular $\|\text{grad } v\| \leq M_2$ at umbilical points, where M_2 is the constant given by Lemma 4.5. On the other hand, we have $\|\text{grad } v\|^2 = e^g \|\text{grad } \chi\|^2$ where

$$g(\chi) = -2(\chi + \ln(1 - e^{-\chi})).$$

By a direct computation

$$g' = -2 - 2\frac{e^{-x}}{1 - e^{-x}} = -\frac{2}{1 - e^{-x}},$$

$$g'' = \frac{2e^{-x}}{(1 - e^{-x})^2}.$$

Applying Lemma 4.3 we have

$$\Delta_S(\|\text{grad } v\|^2) \geq a(\chi)\|\text{grad } v\|^4 + b(\chi)\|\text{grad } v\|^2 + c(\chi),$$

where $c(\chi) \geq 0$ and

$$a(\chi) = \left(\frac{2e^{-\chi}}{(1-e^{-\chi})^2} \right) (1-e^{-\chi})^2 e^{2\chi} = 2e^\chi,$$

$$b(\chi) = \left(-12 \left(\frac{1}{1-e^{-\chi}} \right) (1-e^{-2\chi}) + (10e^{-2\chi} - 2) \right) = 10e^{-2\chi} - 12e^{-\chi} - 14.$$

Observe that $-b(\chi)/a(\chi) \rightarrow 0$ as $\chi \rightarrow +\infty$ and $-b(\chi)/a(\chi) \rightarrow 8$ as $\chi \rightarrow 0$. By Remark 4.4, taking $M = \sup_{\chi>0} (-b(\chi)/a(\chi))$ we have that $\Delta_S(\|\text{grad } v\|^2) > 0$ whenever $\|\text{grad } v\|^2 > M$.

Summarizing we have that

- $\|\text{grad } v\|^2$ is bounded by a constant M_2 at umbilical points;
- if at some non-umbilical point $\|\text{grad } v\|^2 > M_1$ then $\Delta_S(\|\text{grad } v\|^2) > 0$.

We now claim that

$$(38) \quad \sup \|\text{grad } v\|^2 \leq M^2 := \max(M_1, M_2).$$

To prove the claim, suppose to have a sequence of points x_n such that $\|\text{grad } v\|^2(x_n)$ converges to $\sup_S \|\text{grad } v\|^2$. The latter is finite by Equation (37), since $\|\text{grad } \lambda\|$ is bounded and $e^{-\chi}$ stays uniformly away from 1 by the hypothesis of uniformly negative curvature of S . Take a sequence of isometries T_n of AdS^3 so that $T_n(x_n)$ is a fixed point x_0 and the maximal surface $S_n = T_n(S)$ is tangent to a fixed spacelike plane through x_0 .

By Lemma 1.20, up to a subsequence the surfaces S_n converge C^∞ on compact sets to a maximal surface S_∞ .

If x_0 is an umbilical point for S_∞ , then by Equation (37) $\|\text{grad } v_n(x_0)\|^2$ is bounded by $(1+\epsilon)M_2$ for n large. Hence $\|\text{grad } v(x_n)\|$ is bounded by $(1+\epsilon)M_2$, and thus $\sup_S \|\text{grad } v\|^2$ is bounded by $(1+\epsilon)M_2$. (It is actually bounded by M_2 itself, since in the argument ϵ is arbitrary.) On the other hand, if x_0 is not an umbilical point for S_∞ , by the C^∞ convergence

$$\|\text{grad } v_\infty\|^2(x_0) = \sup_S \|\text{grad } v\|^2 = \sup_{S_\infty} \|\text{grad } v_\infty\|^2,$$

therefore x_0 is an interior maximum point for $\|\text{grad } v_\infty\|^2$. Hence $\Delta_S \|\text{grad } v_\infty\|^2(x_0) \leq 0$ and $\|\text{grad } v_\infty\|^2 \leq M_1$ by the first part of the proof. \square

Proposition 4.B. *There exists a universal constant M such that, for every maximal surface S of nonpositive curvature in AdS^3 and every pair of points $p, q \in S$,*

$$1 - \lambda(q) \leq e^{Md_S(p,q)}(1 - \lambda(p)).$$

Proof. Using Proposition 4.1, let $\gamma : [0, d_S(p, q)]$ be a unit speed parameterization of the geodesic segment connecting p and q , and get

$$|v(p) - v(q)| = \left| \int_{[0, d_S(p,q)]} dv(\dot{\gamma}(t)) dt \right| \leq \int_{[0, d_S(p,q)]} \|\text{grad } v\| dt \leq Md_S(p, q).$$

Therefore

$$(39) \quad v(q) \geq v(p) - Md_S(p, q),$$

from which the statements follows, by recalling that v is the function on S defined so that $e^{v(x)} = 1/(1 - \lambda(x))$. \square

In particular, we will apply Proposition 4.B in the following form.

Corollary 4.6. *There exists a universal constant M such that, for every maximal surface S in AdS^3 and every pair $x_0 \in S$, $v(x) \geq v(x_0)/2$ for every point x in the geodesic ball $B_S(x_0, v(x_0)/2M)$.*

Proof. Follows directly from Equation (39) in the proof of Proposition 4.B, with the choice $d(x_0, x) \leq v(x_0)/2M$. \square

4.2. Barriers for the lines of curvature. We need also to deduce that, on a large ball $B_S(x_0, v(x_0)/2M)$ as estimated in Proposition 4.B, the lines of curvature of a maximal surface S are closer and closer to being geodesics, in the sense that their intrinsic acceleration is small, as $\lambda(x_0)$ tends to 1.

Corollary 4.7. *For every $\delta \in (0, 1)$ there exists a constant $M(\delta)$ such that, for every maximal surface S in AdS^3 and every point $x_0 \in S$ with $\lambda(x_0) \geq \delta$, the lines of curvature of S have intrinsic acceleration, inside the ball $B_S(x_0, v(x_0)/2M)$ for the intrinsic metric of S , bounded by:*

$$\|\nabla_{\dot{\gamma}_c}^S \dot{\gamma}_c(t)\| \leq M(\delta)e^{-v(\gamma_c(t))},$$

where γ_c is a unit-speed parametrization of any portion of line of curvature of S contained inside the ball $B_S(x_0, v(x_0)/2M)$.

Proof. It turns out (see for instance [KS07]) that the second fundamental form \mathbb{I} of S is the real part of a holomorphic quadratic differential for the complex structure underlying the induced metric on S . Let us denote by q this holomorphic quadratic differential, so that $\mathbb{I} = \text{Re}(q)$. By Corollary 4.6, assuming $\lambda(x_0) \geq \delta > 0$, the eigenvalues of S are nonzero on the geodesic ball $B_S(x_0, v(x_0)/2M)$, and thus also $q \neq 0$ on the same ball. Hence one can find a conformal chart for $B_S(x_0, v(x_0)/2M)$ for which $q = dz^2$.

In this coordinate the first fundamental form I of S has the form $e^\varphi |dz|^2$, for some real function φ . We claim that φ coincides with the function $\chi = -\ln|\lambda|$. Indeed, observe that the shape operator of S has the form

$$B = e^{-\varphi} \mathbb{I} = e^{-\varphi} \text{Re}(q) = e^{-\varphi} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} = \begin{pmatrix} e^{-\varphi} & 0 \\ 0 & -e^{-\varphi} \end{pmatrix}.$$

Since the eigenvalues of B are $\pm\lambda$, assuming $\lambda > 0$, we get $\lambda = e^{-\varphi}$ and therefore $\varphi = \chi$ as claimed.

In such coordinates, the lines of curvature of S are the lines with constant coordinates $x = \Re(z)$ or $y = \Im(z)$. Let us denote by e_x, e_y the orthonormal frame given by such lines of curvature. By a direct computation, one checks that

$$\nabla_{e_x}^S e_x = \left(-\frac{1}{2} e^{-\frac{\chi}{2}} \partial_y \chi \right) e_y, \quad \nabla_{e_y}^S e_y = \left(-\frac{1}{2} e^{-\frac{\chi}{2}} \partial_x \chi \right) e_x.$$

Hence one has

$$\|\nabla_{e_x}^S e_x\|^2 = \frac{1}{4} e^{-\chi} (\partial_y \chi)^2, \quad \|\nabla_{e_y}^S e_y\|^2 = \frac{1}{4} e^{-\chi} (\partial_x \chi)^2.$$

Observe that the gradient of χ , for the induced metric on S , has squared norm

$$\|\text{grad } \chi\|^2 = e^{-\chi} ((\partial_x \chi)^2 + (\partial_y \chi)^2),$$

and thus one directly obtains

$$\|\nabla_{e_x}^S e_x\|^2 \leq \frac{1}{4} \|\text{grad } \chi\|^2.$$

On the other hand, by Equation (37), we have

$$\|\text{grad } \chi\|^2 = e^{2\chi} e^{-2v} \|\text{grad } v\|^2.$$

Since by hypothesis $\lambda(x_0)$ is bounded away from zero by δ , by Corollary 4.6 e^χ is uniformly bounded by some constant $C = C(\delta)$ on $B_S(x_0, v(x_0)/2M)$. Hence

$$\|\nabla_{e_x}^S e_x\| \leq C(\delta) M e^{-v},$$

where M is the constant of Proposition 4.1, and the same holds for $\|\nabla_{e_y}^S e_y\|$. Upon relabeling the constant M , this concludes the proof. \square

In the following, we will always fix $\delta \in (0, 1)$ and denote by M a larger constant satisfying the statement of both Proposition 4.1 and Corollary 4.7.

Observe that, given a totally geodesic plane P , the surface at timelike distance d from P (in the past, say) is a complete convex constant mean curvature umbilical surface with shape operator $(\tan d)E$ at every point. This follows for instance by applying Equation (14) of Lemma 1.14. Therefore, given a surface S with future unit normal vector N_0 at the point $x_0 \in S$, consider the totally geodesic plane P which contains the point $y_0 = \exp_{x_0}((\arctan \bar{\lambda})N_0)$ and is orthogonal to the timelike line through x_0 and y_0 . Thus the surface at distance $\bar{d} = \arctan \bar{\lambda}$ in the past from P , which we denote by $U_{\bar{\lambda}}(x_0, N_0)$, is an umbilical constant mean curvature surface tangent to S at x_0 . The shape operator of $U_{\bar{\lambda}}(x_0, N_0)$ is $\bar{\lambda}E$. See Figure 4.1.

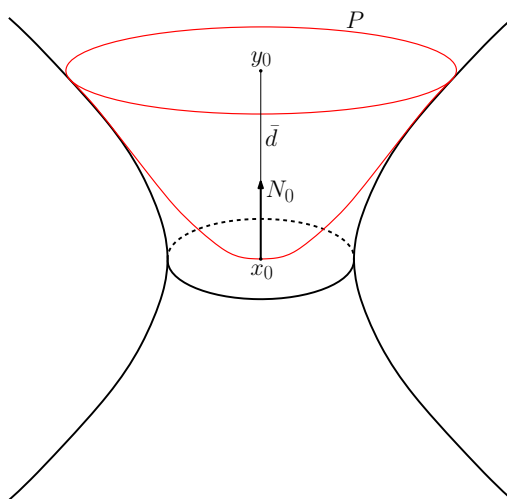


Figure 4.1. The umbilical surface $U_{\bar{\lambda}}(x_0, N_0)$, for $\bar{\lambda} = \tan \bar{d}$, constructed as a parallel surface of the totally geodesic plane P .

We shall denote by $l_c^+(x_0, a)$ (resp. $l_c^-(x_0, a)$) the segment of the line of curvature of S for the positive (resp. negative) eigenvalue, which contains x_0 and whose extrema are at distance a from x_0 for the induced metric.

The following lemma is a subtle application of a maximum principle argument. See also Figure 4.2 (for the statement) and Figure 4.3 (for the proof).

Lemma 4.8. *There exists a constant $\delta \in (0, 1)$ as follows. Suppose S is a maximal surface with future unit normal vector N_0 at x_0 and with $\lambda(x_0) = 1 - e^{-v(x_0)} \geq \delta$. Then $l_c^+(x_0, a)$ is entirely contained in the convex side of the surface $U_{\lambda_1}(x_0, N_0)$, for $a = v(x_0)/2M$ and $\lambda_1 = 1 - e^{-v(x_0)/4}$.*

Proof. Choose $\delta > 0$ as in Corollary 4.7, and suppose *ab absurdum* that $p \in l_c^+(x_0, a)$ is strictly in the past of $U_{\lambda_1}(x_0, N_0)$.

Recall that, for a spacelike curve γ in a Lorentzian manifold, the curvature of γ is defined as $\kappa = \sqrt{|\langle \nabla_{\dot{\gamma}}^S \dot{\gamma}, \nabla_{\dot{\gamma}}^S \dot{\gamma} \rangle|}$. If $\gamma_c^+ : [-a, a] \rightarrow l_c^+(x_0, a)$ is a unit-speed parameterization of $l_c^+(x_0, a)$, we have

$$\nabla_{\dot{\gamma}_c^+} \dot{\gamma}_c^+ = \nabla_{\dot{\gamma}_c^+}^S \dot{\gamma}_c^+ + \lambda N,$$

where N is the unit future-directed normal vector field of the maximal surface S . Hence

$$\kappa^2 = \lambda^2 - \|\nabla_{\dot{\gamma}_c^+}^S \dot{\gamma}_c^+\|^2.$$

On the other hand, let T be the timelike plane spanned by $\dot{\gamma}_c^+$ and N_0 , and let γ_1 be a unit-speed parameterization of the spacelike curve $U_{\lambda_1}(x_0, N_0) \cap T$. Such curve is a geodesic of $U_{\lambda_1}(x_0, N_0)$, by a classical argument of symmetry.

We first claim that $\gamma_c^+(-\epsilon, \epsilon)$ is contained in the future of $U_{\lambda_1}(x_0, N_0)$ for some $\epsilon > 0$. Indeed, if this were not the case, the curvature of γ_1 at x_0 should be larger than the curvature of γ_c^+ . Since γ_1 is geodesic for $U_{\lambda_1}(x_0, N_0)$, the curvature of γ_1 is

$$\kappa_1 = \lambda_1.$$

By Corollary 4.6 we have $v(x) \geq v(x_0)/2$ on $B_S(x_0, v(x_0)/2M)$, and by Corollary 4.7, the intrinsic acceleration of $l_c^+(x_0, a)$ is bounded by $\|\nabla_{\dot{\gamma}_c^+}^S \dot{\gamma}_c^+\| \leq Me^{-v} \leq Me^{-v(x_0)/2}$. Hence

$$\begin{aligned} \kappa^2 &= \lambda^2 - \|\nabla_{\dot{\gamma}_c^+}^S \dot{\gamma}_c^+\|^2 \geq (1 - e^{-v(x_0)/2})^2 - Me^{-v(x_0)/2} \\ &= 1 - (2 + M)e^{-v(x_0)/2} + e^{-v(x_0)} \\ &\geq 1 - e^{-v(x_0)/4} \geq (1 - e^{-v(x_0)/4})^2 = \lambda_1^2 = \kappa_1^2. \end{aligned}$$

We have replaced δ by a larger number if necessary, and used the assumption $\lambda(x_0) \geq \delta$. This gives a contradiction and concludes the claim.

Now consider the function $d: [-a, a] \rightarrow \mathbb{R}$, where $d(t)$ is the signed distance of $\gamma_c^+(t)$ from the surface $U_{\lambda_1}(x_0, N_0)$. The function d is positive in the interval $(-\epsilon, \epsilon)$ by the previous claim, and negative at the point t_0 such that $\gamma_c^+(t_0) = p$. Hence d must achieve a maximum $d_{\max} = d(t_{\max}) > 0$. At the point $\gamma_c^+(t_{\max})$, the curve γ_c^+ is therefore tangent to the surface V at distance d_{\max} from $U_{\lambda_1}(x_0, N_0)$. Again by Lemma 1.14, if d_1 is such that $\lambda_1 = \tan d_1$, then V is an umbilical constant mean curvature convex surface, whose shape operator at every point is $\tan(d_1 - d_{\max})E$. Denote $\lambda_2 = \tan(d_1 - d_{\max})$ and observe that $\lambda_2 < \lambda_1$.

By a similar argument as the previous claim, we compare the curve γ_c^+ to the curve γ_2 , which parameterizes the intersection of V with the timelike plane spanned by $\dot{\gamma}_c^+(t_{\max})$ and $N(\gamma_c^+(t_{\max}))$. We remark that in this case, γ_2 is not a geodesic for V . Since t_{\max} is a maximum point, the curvature κ of γ_c^+ at t_{\max} needs to be smaller than the curvature κ_2 of γ_2 . But by the same computation,

$$\kappa^2 \geq \lambda_1^2 \geq \lambda_2^2 \geq \lambda_2^2 - \|\nabla_{\dot{\gamma}_2}^V \dot{\gamma}_2\|^2 = \kappa_2^2.$$

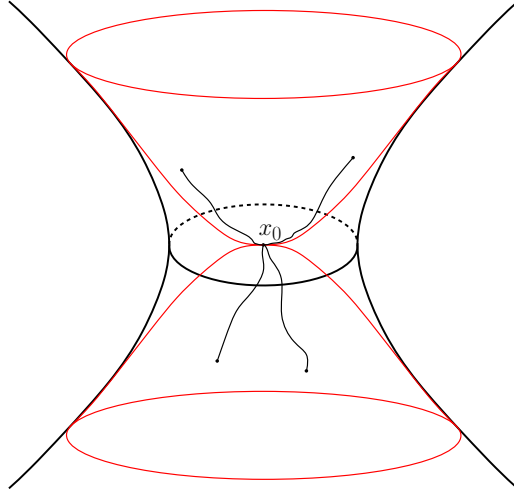


Figure 4.2. Lines of curvature through x_0 and the surfaces $U_{\lambda_1}^+(x_0, N_0)$ and $U_{\lambda_1}^-(x_0, N_0)$

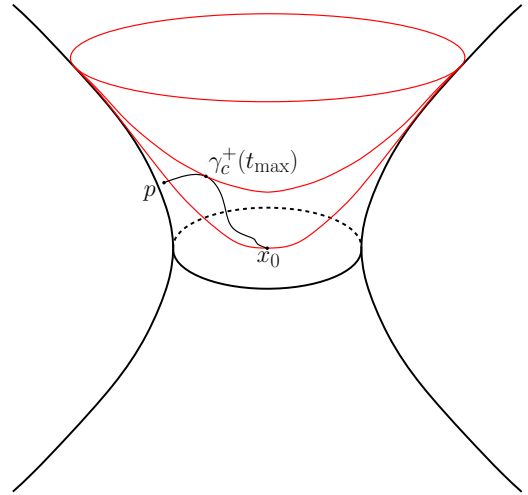


Figure 4.3. The proof of Lemma 4.8 is an argument by contradiction.

This gives a contradiction and thus concludes the proof. \square

Remark 4.9. Clearly a statement analogous to Lemma 4.8 holds for the line of curvature $l_c^-(x_0, a)$ which, under the same assumptions, is entirely contained below the concave umbilical surface $U_{\lambda_1}^-(x_0, N_0)$, tangent to S at x_0 , obtained as the surface at distance $\tan \lambda_1$ in the future of a totally geodesic plane. See Figure 4.2.

4.3. Estimating the width from below. Observe that, in the extreme situation $\lambda_1 = 1$, the umbilical surfaces $U_1^+(x_0, N_0)$ and $U_1^-(x_0, N_0)$ have the following good property. Take two timelike planes T_1 and T_2 which intersect orthogonally in the timelike geodesic through x_0 , directed by the vector N_0 . Then $U_1^+(x_0, N_0) \cap T_1$ is a geodesic of $U_1^+(x_0, N_0)$, and $U_1^-(x_0, N_0) \cap T_2$ is a geodesic of $U_1^-(x_0, N_0)$. The endpoints at infinity of $U_1^+(x_0, N_0) \cap T_1$ determine a spacelike entire line of AdS^3 , which is dual to the line determined by the endpoints at infinity of $U_1^-(x_0, N_0) \cap T_2$. Hence the width of the convex hull of the four points at infinity is $\pi/2$. Indeed, the curves $U_1^+(x_0, N_0) \cap T_1$ and $U_1^-(x_0, N_0) \cap T_2$ are lines of curvature for a horospherical surface. See Figures 4.4 and 4.5. Roughly speaking, in this subsection we want to quantify “how close” we get to this situation when $\lambda(x_0)$ (and thus also λ_1) approaches 1.

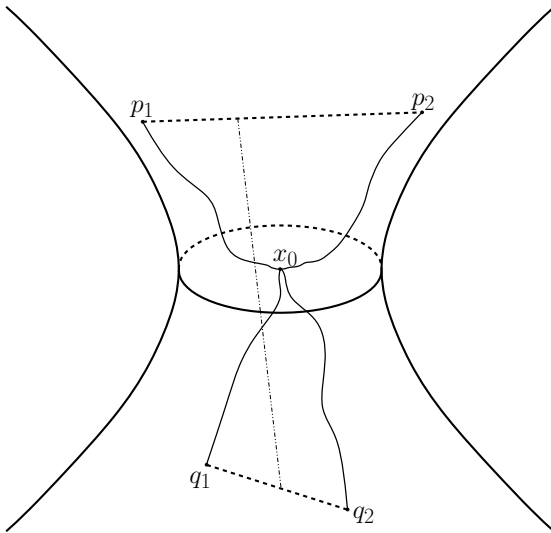


Figure 4.4. We want to estimate the timelike distance between the lines $\overline{p_1 p_2}$ and $\overline{q_1 q_2}$, where p_1, p_2, q_1, q_2 are endpoints of segments on the lines of curvature.

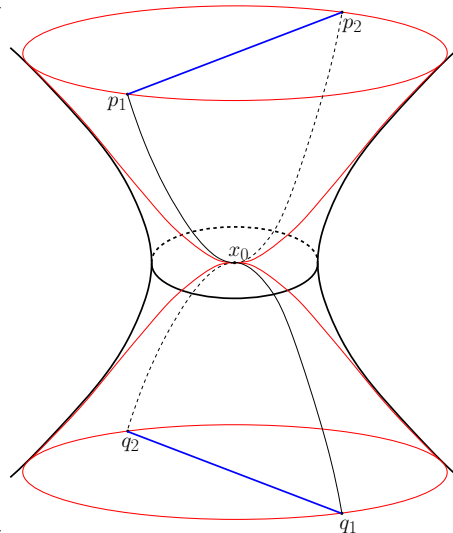


Figure 4.5. The configuration is optimal on a horospherical surface, when the lines of curvature are planar curves (and geodesics), and $\overline{p_1 p_2}$ and $\overline{q_1 q_2}$ are dual lines.

Given a timelike totally geodesic plane $T = w^\perp$ (where $w \in \mathbb{R}^{2,2}$ is a spacelike vector of unit norm) and a point $p \in \text{AdS}^3$, it is easy to see that $|\langle p, w \rangle|$ represents the hyperbolic sine of the length of the spacelike segment $\overline{pp_0}$, such that $p_0 \in T$ and the line containing $\overline{pp_0}$ is orthogonal to T at p_0 . If T_r denotes the (timelike) surface composed of points for which this (signed) length is r , the following lemma gives an estimate on how the lines of curvature of a maximal surface escape from the surfaces T_r . Here $T = T_0$ is chosen as the timelike plane orthogonal to the line of curvature at the base point x_0 .

Lemma 4.10. *There exist constants $\delta \in (0, 1)$ and $t_0 \geq 0$ as follows. Let S be any maximal surface with $\lambda(x_0) = 1 - e^{-v(x_0)} \geq \delta$, and let $\gamma_c : [0, a] \rightarrow \text{AdS}^3$ be a unit-speed*

parameterization of a line of curvature of S with $\gamma_c(0) = x_0$, where $a = v(x_0)/2M$. If $\varphi(t) = \langle \gamma_c(t), \dot{\gamma}_c(0) \rangle$, then

$$\varphi(t) \geq e^t$$

for all $t \in [t_0, a]$.

Proof. Let us compute

$$\dot{\varphi}(t) = \langle \dot{\gamma}(t), \dot{\gamma}(0) \rangle.$$

We now think $\gamma(t)$ as a point of $\widehat{\text{AdS}}^3 \subset \mathbb{R}^{2,2}$. Denote $\psi(t) = \langle N(t), \dot{\gamma}(0) \rangle$, where we write $N(t) = N(\gamma(t))$ by a slight abuse of notation. Hence

$$\ddot{\gamma}(t) = \gamma(t) + \lambda(\gamma(t))N(\gamma(t)) + \nabla_{\dot{\gamma}(t)}^S \dot{\gamma}(t),$$

and therefore

$$\ddot{\varphi}(t) = \varphi(t) + \lambda(t)\psi(t) + \langle \nabla_{\dot{\gamma}(t)}^S \dot{\gamma}(t), \dot{\gamma}(0) \rangle.$$

If we denote $\rho(t) = \dot{\varphi}(t) = \langle \dot{\gamma}(t), \dot{\gamma}(0) \rangle$ and $\alpha(t) = \langle \nabla_{\dot{\gamma}(t)}^S \dot{\gamma}(t), \dot{\gamma}(0) \rangle$, the triple (φ, ψ, ρ) solves the (non-linear, non-autonomous) system of ODEs

$$(40) \quad \begin{cases} \dot{\varphi}(t) = \rho(t) \\ \dot{\psi}(t) = \lambda(t)\rho(t) \\ \dot{\rho}(t) = \varphi(t) + \lambda(t)\psi(t) + \alpha(t) \end{cases},$$

with the initial conditions

$$(41) \quad \begin{cases} \varphi(0) = 0 \\ \psi(0) = 0 \\ \rho(0) = 1 \end{cases}.$$

Sublemma 4.11. *There exists $\delta \in (0, 1)$ such that, if $\lambda(x_0) \geq \delta$, then*

$$|\alpha(t)| < e^{-v(x_0)/4} \sqrt{1 + \varphi(t)^2 + \psi(t)^2 - \rho(t)^2}$$

for every $t \in [0, v(x_0)/2M]$. In particular, setting $\epsilon = e^{-v(x_0)/4}$,

$$|\alpha(t)| < \epsilon(1 + |\varphi(t)| + |\psi(t)|).$$

Proof. Since γ is a unit-speed parameterization, $(\gamma(t), N(t), \dot{\gamma}(t), \nabla_{\dot{\gamma}(t)}^S \dot{\gamma}(t) / \|\nabla_{\dot{\gamma}(t)}^S \dot{\gamma}(t)\|)$ is an orthonormal frame for every t , provided $\nabla_{\dot{\gamma}(t)}^S \dot{\gamma}(t) \neq 0$. Hence

$$\begin{aligned} \dot{\gamma}(0) &= -\langle \dot{\gamma}(0), \gamma(t) \rangle \gamma(t) - \langle \dot{\gamma}(0), N(t) \rangle N(t) + \langle \dot{\gamma}(0), \dot{\gamma}(t) \rangle \dot{\gamma}(t) + \langle \dot{\gamma}(0), \nabla_{\dot{\gamma}(t)}^S \dot{\gamma}(t) \rangle \frac{\nabla_{\dot{\gamma}(t)}^S \dot{\gamma}(t)}{\|\nabla_{\dot{\gamma}(t)}^S \dot{\gamma}(t)\|^2} \\ &= -\varphi(t)\gamma(t) - \psi(t)N(t) + \rho(t)\dot{\gamma}(t) + \alpha(t) \frac{\nabla_{\dot{\gamma}(t)}^S \dot{\gamma}(t)}{\|\nabla_{\dot{\gamma}(t)}^S \dot{\gamma}(t)\|^2}. \end{aligned}$$

Therefore one gets

$$1 = -\varphi(t)^2 - \psi(t)^2 + \rho(t)^2 + \frac{\alpha(t)^2}{\|\nabla_{\dot{\gamma}(t)}^S \dot{\gamma}(t)\|^2}.$$

Recalling that, from Corollary 4.7, $\|\nabla_{\dot{\gamma}}^S \dot{\gamma}\| \leq Me^{-v(x_0)/2} < e^{-v(x_0)/4}$ for δ sufficiently large, one concludes the claim. \square

By Corollary 4.6, $\lambda(t) \geq \eta := 1 - e^{-v(x_0)/2}$ for $t \in [0, v(x_0)/2M]$. We will compare the solution of the system (40) with the solution $(\varphi_1(t), \psi_1(t), \rho_1(t))$ of the following system:

$$(42) \quad \begin{cases} \dot{\varphi}_1(t) = \rho_1(t) \\ \dot{\psi}_1(t) = \eta\rho_1(t) \\ \dot{\rho}_1(t) = (1 - \epsilon)\varphi_1(t) + (\eta - \epsilon)\psi_1(t) - \epsilon \end{cases},$$

with the same initial conditions

$$(43) \quad \begin{cases} \varphi_1(0) = 0 \\ \psi_1(0) = 0 \\ \rho_1(0) = 1 \end{cases} .$$

Sublemma 4.12. *Let $(\varphi(t), \psi(t), \rho(t))$ be the solution of (40),(41) and $(\varphi_1(t), \psi_1(t), \rho_1(t))$ be the solution of (42),(43). For every interval $t \in [0, t_0)$ where $\varphi(t), \psi(t) > 0$, one has $\varphi(t) \geq \varphi_1(t)$.*

Proof. The system (40) can be written in the form of integro-differential equation:

$$\dot{\rho}(t) = \int_0^t \rho(s) ds + \lambda(t) \int_0^t \lambda(s) \rho(s) ds + \alpha(t),$$

while system (42) takes the form

$$\dot{\rho}_1(t) = (1 - \epsilon) \int_0^t \rho_1(s) ds + (\eta - \epsilon) \int_0^t \eta \rho_1(s) ds - \epsilon.$$

By Sublemma 4.11, as soon as $\varphi(t), \psi(t) > 0$ and $t \in [0, v(x_0)/2M]$,

$$\dot{\rho}(t) > (1 - \epsilon) \int_0^t \rho(s) ds + (\eta - \epsilon) \int_0^t \eta \rho(s) ds - \epsilon.$$

Hence one gets

$$\dot{\rho}(t) - \dot{\rho}_1(t) > (1 - \epsilon) \int_0^t (\rho(s) - \rho_1(s)) ds + (\eta - \epsilon) \eta \int_0^t (\rho(s) - \rho_1(s)) ds.$$

This is enough to conclude that $\rho(t) > \rho_1(t)$ for all $t \in [0, t_0)$. Indeed, if $t = t_{\max}$ were a maximal point for which $\rho(t) > \rho_1(t)$, then $\dot{\rho}(t_{\max}) - \dot{\rho}_1(t_{\max})$ would still be strictly positive, thus giving a contradiction. As a direct consequence, $\varphi(t) > \varphi_1(t)$ for all $t \in [0, t_0)$. \square

To conclude the proof of Lemma 4.10, it suffices to check by a direct computation that the solution of (42) with initial conditions (43) is:

$$\begin{pmatrix} \varphi_1(t) \\ \psi_1(t) \\ \rho_1(t) \end{pmatrix} = \begin{pmatrix} \frac{\epsilon}{1 - \epsilon + \eta(\eta - \epsilon)} \\ \frac{\eta \epsilon}{1 - \epsilon + \eta(\eta - \epsilon)} \\ 0 \end{pmatrix} + c_1 \begin{pmatrix} \frac{1}{\sqrt{1 - \epsilon + \eta(\eta - \epsilon)}} \\ \frac{\eta}{\sqrt{1 - \epsilon + \eta(\eta - \epsilon)}} \\ -1 \end{pmatrix} e^{-t\sqrt{1 - \epsilon + \eta(\eta - \epsilon)}} + c_2 \begin{pmatrix} \frac{1}{\sqrt{1 - \epsilon + \eta(\eta - \epsilon)}} \\ \frac{\eta}{\sqrt{1 - \epsilon + \eta(\eta - \epsilon)}} \\ 1 \end{pmatrix} e^{t\sqrt{1 - \epsilon + \eta(\eta - \epsilon)}},$$

where the constants are

$$c_1 = \frac{1}{2} \left(-1 - \frac{\epsilon}{\sqrt{1 - \epsilon + \eta(\eta - \epsilon)}} \right)$$

and

$$c_2 = \frac{1}{2} \left(1 - \frac{\epsilon}{\sqrt{1 - \epsilon + \eta(\eta - \epsilon)}} \right).$$

This shows in particular that

$$\rho_1(t) = \cosh(t\sqrt{1 - \epsilon + \eta(\eta - \epsilon)}) - \frac{\epsilon}{\sqrt{1 - \epsilon + \eta(\eta - \epsilon)}} \sinh(t\sqrt{1 - \epsilon + \eta(\eta - \epsilon)}),$$

and is therefore positive for all $t \geq 0$ (since ϵ is small). Hence also $\varphi_1(t) \geq 0$ and therefore $\varphi(t)$ remains positive as well. Therefore the assumptions of Sublemma 4.12 are satisfied for all $t \in [0, v(x_0)/2M]$. Hence there exists $t_0 > 0$ such that

$$\varphi(t) \geq \varphi_1(t) \geq e^t$$

for $t \in [t_0, a]$ as claimed. \square

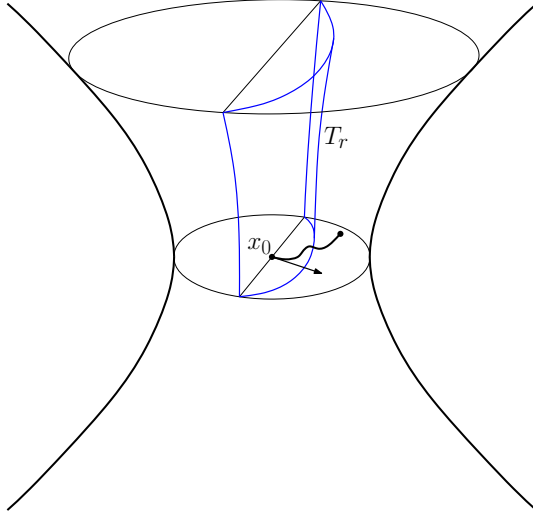


Figure 4.6. Lemma 4.10 asserts that the lines of curvature escape from the region bounded by the surface T_r .

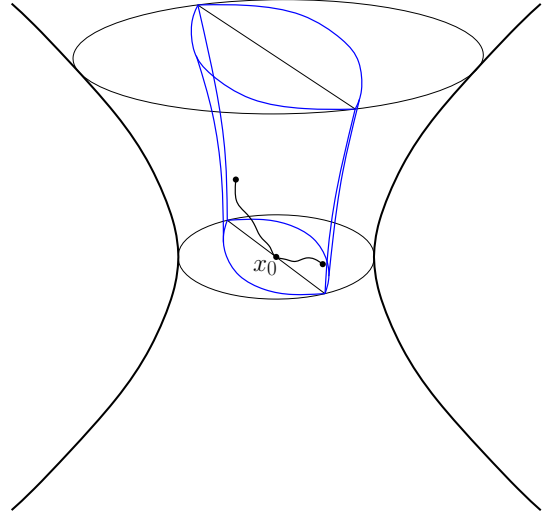


Figure 4.7. Lemma 4.14 instead quantifies how the lines of curvature remain inside a thin slice bounded by two timelike surfaces.

Remark 4.13. Using the same techniques as in Lemma 4.10, one can consider the function $\varphi(t) = \langle \gamma_c(t), w \rangle$, where w is the unit spacelike vector orthogonal to both $\dot{\gamma}_c(0)$ and to $N(x_0)$. One then similarly defines $\rho(t) = \dot{\varphi}(t) = \langle \dot{\gamma}(t), w \rangle$ and $\psi(t) = \langle \gamma(t), N(t) \rangle$. Hence the triple (φ, ψ, ρ) solves the system (40), now with the initial conditions

$$(44) \quad \begin{cases} \varphi(0) = 0 \\ \psi(0) = 0 \\ \rho(0) = 0 \end{cases} .$$

Since $\lambda \leq 1$, we consider the supersolution $(\varphi_2(t), \psi_2(t), \rho_2(t))$ which solves the system:

$$(45) \quad \begin{cases} \dot{\varphi}_2(t) = \rho_2(t) \\ \dot{\psi}_2(t) = \rho_2(t) \\ \dot{\rho}_2(t) = (1 + \epsilon)\varphi_2(t) + (1 + \epsilon)\psi_2(t) + \epsilon \end{cases} ,$$

with the same initial conditions

$$(46) \quad \begin{cases} \varphi_2(0) = 0 \\ \psi_2(0) = 0 \\ \rho_2(0) = 0 \end{cases} .$$

The latter system is easily solved, since it is equivalent to the equation $\ddot{\rho}(t) = 2(1 + \epsilon)\rho(t)$, and therefore one gets

$$\rho_2(t) = \frac{\epsilon}{\sqrt{2(1 + \epsilon)}} \sinh(t\sqrt{2(1 + \epsilon)})$$

and

$$\varphi_2(t) = \psi_2(t) = \frac{\epsilon}{2(1 + \epsilon)} (\cosh(t\sqrt{2(1 + \epsilon)}) - 1)$$

as a solution to (45), (46). Recalling that, under the usual assumptions, we have $\epsilon = e^{-v(x_0)/4}$. Thus for $t = a = v(x_0)/2M$, $\varphi_2(t)$ is dominated by an exponential of exponent $(v(x_0)/4)(-1 + (2\sqrt{2}(1 + \epsilon))/M)$, which is estimated from above by $e^{-v(x_0)/8}$ provided M is sufficiently large (we can always replace M by a larger constant, as we did several times

before). In conclusion, we get $\varphi(a) \leq \varphi_2(a) \leq e^{-t/8}$. Of course the situation is symmetric, and one can prove the same upper bound for $-\varphi(t)$.

We report here the statement of a lemma, which was discussed in Remark 4.13.

Lemma 4.14. *There exists constants $\delta \in (0, 1)$ and $t_0 \geq 0$ as follows. Let S be any maximal surface with future unit normal vector N_0 at x_0 and with $\lambda(x_0) = 1 - e^{-v(x_0)} \geq \delta$, and let $\gamma_c : [0, a] \rightarrow \text{AdS}^3$ be a unit-speed parameterization of a line of curvature of S with $\gamma_c(0) = x_0$, with $a = v(x_0)/2M$. If $\varphi(t) = \langle \gamma_c(t), w \rangle$, where w is the unit spacelike vector orthogonal to both $\dot{\gamma}_c(0)$ and to $N(x_0)$, then*

$$|\varphi(a)| \leq e^{-v(x_0)/8}.$$

To give an estimate from below for the width, we will consider a maximal surface with large principal curvatures at the point x_0 , which we will assume to be the point $x_0 = [0, 0, 1, 0]$. We are going to use again the coordinate system (15), which we write here again:

$$(r, \theta, \zeta) \mapsto [\cos \theta \sinh r, \sin \theta \sinh r, \cos \zeta \cosh r, \sin \zeta \cosh r].$$

We are assuming the maximal surface S is tangent at x_0 to the plane $\zeta = 0$. Hence the level sets $\zeta = c$ are totally geodesic planes orthogonal to the timelike like which starts x_0 with initial tangent vector $N(x_0)$.

Lemma 4.15. *Let $a = v(x_0)/2M$ and $\tan d_1 = \lambda_1 = 1 - e^{-v(x_0)/4}$. Let $p_1 = \gamma_c^+(-a)$ and $p_2 = \gamma_c^+(a)$ be the endpoints of the segment $l_c^+(x_0, a)$ of a line of curvature through x_0 . Then*

$$|\sin(d_1 - \zeta(p_i))| \leq \frac{\sin d_1}{\cosh(r(p_i))}.$$

In other words, $\zeta(p_i) \geq \bar{\zeta}$ where $\bar{\zeta} \leq d_1$ satisfies

$$\sin(d_1 - \bar{\zeta}) = \frac{\sin d_1}{\cosh(r(p_i))}.$$

Proof. We know from Lemma 4.8 that $l_c^+(x_0, a)$ is entirely contained in the future-directed side of the surface $U_{\lambda_1}(x_0, N_0)$, for $a = v(x_0)/2M$ and $\lambda_1 = 1 - e^{-v(x_0)/4}$. Recall that the surface $U_{\lambda_1}(x_0, N_0)$ is obtained as the surface at distance d_1 (past-directed) from the plane $\zeta = d_1$, where $\tan d_1 = \lambda_1$. This plane is also defined by

$$\{\zeta = d_1\} = [0, 0, \sin d_1, -\cos d_1]^T.$$

Observe that the product of p_i and $[0, 0, \sin d_1, -\cos d_1]$, in absolute value, is the sine of the timelike distance of p_i from the plane $[0, 0, \sin d_1, -\cos d_1]^T$. Hence if $(r(p_i), \theta(p_i), \zeta(p_i))$ are the coordinates of p_i , then

$$\begin{aligned} \sin d_1 &\geq |\langle p_i, [0, 0, \sin d_1, -\cos d_1] \rangle| \\ &= |\sin d_1 \cos(\zeta(p_i)) \cosh(r(p_i)) - \cos d_1 \sin(\zeta(p_i)) \cosh(r(p_i))| \\ &= |\sin(d_1 - \zeta(p_i))| \cosh(r(p_i)). \end{aligned}$$

from which the statement follows straightforwardly. \square

Proof of Theorem 1.B. As observed earlier, we can assume that there exists a point x_0 (and we set $x_0 = [0, 0, 1, 0]$) with $\lambda(x_0) \geq \delta$. Composing with an isometry, we also assume that (in the double cover $\widehat{\text{AdS}^3}$) $N_0 = (0, 0, 0, 1)$ and that the tangent vectors to the lines of curvature at x_0 are $\dot{\gamma}_c^+(0) = (1, 0, 0, 0)$ and $\dot{\gamma}_c^-(0) = (0, 1, 0, 0)$. Let as usual $p_1 = \gamma_c^+(-a)$ and $p_2 = \gamma_c^+(a)$ be the endpoints of the segment $l_c^+(x_0, a)$, and $q_1 = \gamma_c^-(-a)$ and $q_2 = \gamma_c^-(a)$ be the endpoints of the segment $l_c^-(x_0, a)$, for $a = v(x_0)/2M$. Recall also Figure 4.4.

The width w of the convex hull of S is at least the supremum of the length of geodesic timelike segments which connect the spacelike segments $\overline{p_1 p_2}$ and $\overline{q_1 q_2}$, which we denote by

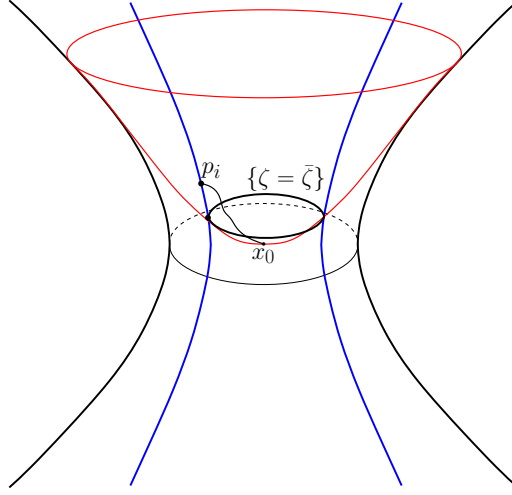


Figure 4.8. In Lemma 4.15, the intersection of the surface $U_{\lambda_1}(x_0, N_0)$ with the cylinder which contains p_i lies in the plane $\{\zeta = \bar{\zeta}\}$.

$d_{\mathbb{A}dS^3}(\overline{p_1 p_2}, \overline{q_1 q_2})$. Let $(r(p_i), \theta(p_i), \zeta(p_i))$ the coordinates of p_i and $(r(q_i), \theta(q_i), \zeta(q_i))$ the coordinates of q_i . By Lemma 4.15, $\zeta(p_i) \geq \bar{\zeta}$ and $\zeta(q_i) \leq -\bar{\zeta}$, for $i = 1, 2$.

Hence $d_{\mathbb{A}dS^3}(\overline{p_1 p_2}, \overline{q_1 q_2})$ is certainly larger than $d_{\mathbb{A}dS^3}(\overline{p'_1 p'_2}, \overline{q'_1 q'_2})$, where p'_i has coordinates $(r(p_i), \theta(p_i), \bar{\zeta})$ and q'_i has coordinates $(r(q_i), \theta(q_i), -\bar{\zeta})$. Compare also Figure 4.9. Indeed, every timelike segment connecting $\overline{p_1 p_2}$ and $\overline{q_1 q_2}$ can be continued to a timelike segment connecting $\overline{p'_1 p'_2}$ and $\overline{q'_1 q'_2}$ of larger timelike length.

Now, the segment $\overline{p'_1 p'_2}$ is clearly contained in the plane $\zeta = \bar{\zeta}$, and it contains a point i_+ with coordinates $(r(i_+), \theta(i_+) = \pi/2, \bar{\zeta})$. See Figure 4.10. Therefore

$$i_+ = [0, \sinh r(i_+), \cos \bar{\zeta} \cosh r(i_+), \sin \bar{\zeta} \cosh r(i_+)].$$

Analogously, the segment $\overline{q'_1 q'_2}$ contains the point

$$i_- = [\sinh r(i_-), 0, \cos \bar{\zeta} \cosh r(i_-), -\sin \bar{\zeta} \cosh r(i_-)].$$

Hence one can give the following bound for the width w :

$$\cos w \leq |\langle i_+, i_- \rangle| = 2 \cosh r(i_-) \cosh r(i_+) \cos \bar{\zeta} \sin \bar{\zeta} = \cosh r(i_-) \cosh r(i_+) \cos(2\bar{\zeta}),$$

which is equivalent to

$$(\tan w)^2 = \frac{1}{(\cos w)^2} - 1 \geq \frac{1}{(\cos(2\bar{\zeta}))^2} \frac{1}{(\cosh r(i_+))^2} \frac{1}{(\cosh r(i_-))^2} - 1.$$

We now want to estimate the factors $\cos(2\bar{\zeta})$ and $\cosh r(i_+)$, $\cosh r(i_-)$. For the former, let us write $\bar{\zeta} = d_1 - (d_1 - \bar{\zeta})$ and compute

$$(47) \quad \cos(2\bar{\zeta}) = \cos(2d_1) \cos(2(d_1 - \bar{\zeta})) + \sin(2d_1) \sin(2(d_1 - \bar{\zeta})).$$

Recall that $\tan d_1 = \lambda_1 = 1 - e^{-v(x_0)/4}$, hence $(\cos d_1)^2 = 1/(1 + \lambda_1^2)$ and

$$(48) \quad \cos(2d_1) = (\cos d_1)^2 - (\sin d_1)^2 = \frac{1}{1 + \lambda_1^2} - \frac{\lambda_1^2}{1 + \lambda_1^2} \leq 2(1 - \lambda_1) = 2e^{-v(x_0)/4}.$$

For the second term of the RHS of Equation (47), using Lemma 4.15, we have

$$(49) \quad \sin(2(d_1 - \bar{\zeta})) = 2 \sin(d_1 - \bar{\zeta}) \cos(d_1 - \bar{\zeta}) \leq \frac{2 \sin d_1}{\cosh r(p_i)}.$$

Observe that

$$\cosh r(p_i) \geq |\sinh r(p_i)| \geq |\langle p_i, (1, 0, 0, 0) \rangle| = |\langle p_i, \dot{\gamma}_c^+(0) \rangle| \geq e^{v(x_0)/2M},$$

where in the last step we have used Lemma 4.10 (up to changing the orientation of the parameterization γ_c^+ for one of the two points p_i). Therefore using the inequalities of Equations (48) and (49) in (47), and relabeling M by a yet larger constant, we get

$$\frac{1}{\cos(2\zeta)} \geq e^{v(x_0)/M},$$

provided $\lambda(x_0)$ is larger than the constant δ .

On the other hand, observe that $r(i_+)$ is the distance in the hyperbolic plane $\zeta = \bar{\zeta}$ of the point i_+ from the geodesic defined by $\theta = 0$. By the convexity of the distance function, $r(i_+)$ is less than the maximum between the distance of p'_1 and p'_2 from the line $\theta = 0$, which remains bounded by Lemma 4.14 (see Remark 4.13). Actually, $r(i_+)$ tends to zero as $v(x_0) \rightarrow \infty$, and the same holds for $r(i_-)$. Hence the factors $\cosh r(i_+)$ and $\cosh r(i_-)$ remain bounded, and this concludes the proof that

$$\tan w \geq e^{v(x_0)/M} = \left(\frac{1}{1 - \lambda(x_0)} \right)^{1/M}.$$

Hence, by a continuity argument, also the inequality

$$\tan w \geq \left(\frac{1}{1 - \|\lambda\|_\infty} \right)^{1/M}$$

holds. \square

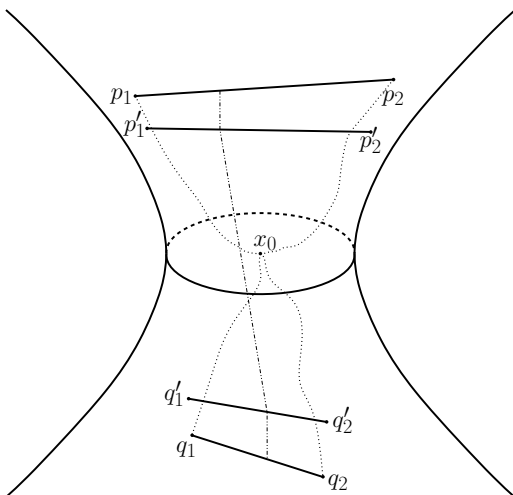


Figure 4.9. The width is not increased when replacing the points p_1, p_2, q_1, q_2 by p'_1, p'_2, q'_1, q'_2 .

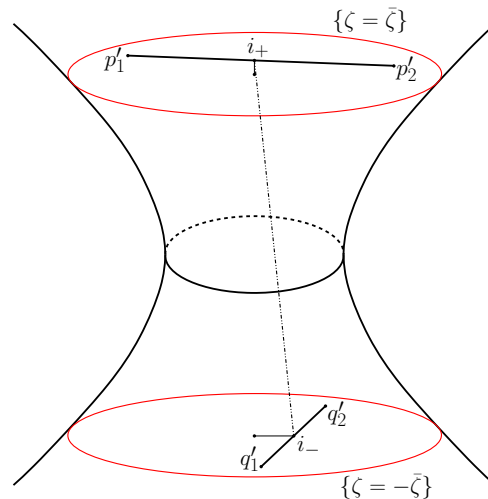


Figure 4.10. The position of the points i_+ and i_- , which are used to derive the lower bound on the width.

5. APPLICATION: MINIMAL LAGRANGIAN QUASICONFORMAL EXTENSIONS

We begin this section by briefly introducing the theory of quasiconformal mappings and universal Teichmüller space. Useful references are [Gar87, GL00, Ahl06, FM07]. Next, we discuss the applications of Theorem 1.A, Theorem 1.B and Proposition 1.C, proving Theorem 2.A, Theorem 2.B, Theorem 2.C and Corollary 2.D.

5.1. Quasiconformal mappings. We recall the definition of quasiconformal map.

Definition 5.1. Given a domain $\Omega \subset \mathbb{C}$, an orientation-preserving homeomorphism

$$f : \Omega \rightarrow f(\Omega) \subset \mathbb{C}$$

is *quasiconformal* if f is absolutely continuous on lines and there exists a constant $k < 1$ such that

$$|\partial_{\bar{z}}f| \leq k|\partial_zf|.$$

Let us denote $\mu_f = \partial_{\bar{z}}f/\partial_zf$, which is called *complex dilatation* of f . This is well-defined almost everywhere, hence it makes sense to take the L_∞ norm. Thus a homeomorphism $f : \Omega \rightarrow f(\Omega) \subset \mathbb{C}$ is quasiconformal if $\|\mu_f\|_\infty < 1$. Moreover, a quasiconformal map as in Definition 5.1 is called *K -quasiconformal*, where

$$K = \frac{1+k}{1-k}.$$

It turns out that the best such constant $K \in [1, +\infty)$ represents the *maximal dilatation* of f , i.e. the supremum over all $z \in \Omega$ of the ratio between the major axis and the minor axis of the ellipse which is the image of a unit circle under the differential d_zf .

It is known that a 1-quasiconformal map is conformal, and that the composition of a K_1 -quasiconformal map and a K_2 -quasiconformal map is K_1K_2 -quasiconformal. Hence composing a quasiconformal map with conformal maps does not change the maximal dilatation.

Actually, there is an explicit formula for the complex dilatation of the composition of two quasiconformal maps f, g on Ω :

$$(50) \quad \mu_{g \circ f^{-1}} = \frac{\partial_z f \mu_g - \mu_f}{\partial_z f (1 - \overline{\mu_f} \mu_g)}.$$

Using Equation (50), one can see that f and g differ by post-composition with a conformal map if and only if $\mu_f = \mu_g$ almost everywhere.

The connection between quasiconformal homeomorphisms of \mathbb{H}^2 and quasisymmetric homeomorphisms of the boundary of \mathbb{H}^2 is made evident by the following classical theorem (see [BA56]).

Ahlfors-Beuring Theorem. Every quasiconformal map $\Phi : \mathbb{H}^2 \rightarrow \mathbb{H}^2$ extends to a quasisymmetric homeomorphism of $\mathbb{R}P^1 = \partial\mathbb{H}^2$. Conversely, any quasisymmetric homeomorphism $\phi : \mathbb{R}P^1 \rightarrow \mathbb{R}P^1$ admits a quasiconformal extension to \mathbb{H}^2 .

5.2. Minimal Lagrangian extension. Our purpose is to give a quantitative description of *minimal Lagrangian* extensions of a quasisymmetric homeomorphism ϕ .

Definition 5.2. A diffeomorphism $\Phi : \mathbb{H}^2 \rightarrow \mathbb{H}^2$ is minimal Lagrangian if Φ is area-preserving and the graph of Φ is a minimal surface in $\mathbb{H}^2 \times \mathbb{H}^2$.

The following characterization of minimal Lagrangian diffeomorphisms is well-known. A proof can be found in [Tou15, Proposition 1.2.6].

Proposition 5.3. A diffeomorphism $\Phi : \mathbb{H}^2 \rightarrow \mathbb{H}^2$ is minimal Lagrangian if and only if $\Phi^*(g_{\mathbb{H}^2}) = g_{\mathbb{H}^2}(b, \cdot)$, where $b \in \Gamma(\text{End}(TM))$ is a bundle morphism such that:

- b is self-adjoint for $g_{\mathbb{H}^2}$;
- $\det b = 1$;
- $d^\nabla b = 0$.

Here d^∇ is the exterior derivative, hence $d^\nabla b$ is the two-form defined by:

$$d^\nabla b(v, w) = \nabla_{\tilde{v}}(b(\tilde{w})) - \nabla_{\tilde{w}}(b(\tilde{v})) - b[\tilde{v}, \tilde{w}],$$

where \tilde{v}, \tilde{w} are vector fields which extend the vectors v, w in a neighborhood of the base point. The vanishing of $d^\nabla b$ is the so-called *Codazzi condition*.

In [BS10], entire maximal surfaces of uniformly negative curvature were used to prove the following theorem:

Theorem 5.4 ([BS10, Theorem 1.4]). *For every quasimetric homeomorphism $\phi : \mathbb{RP}^1 \rightarrow \mathbb{RP}^1$, there exists a unique quasiconformal minimal Lagrangian extension $\Phi_{ML} : \mathbb{H}^2 \rightarrow \mathbb{H}^2$.*

The key observation is that the maximal surface with $\partial_\infty S = gr(\phi)$ corresponds to the minimal Lagrangian extension Φ_{ML} of ϕ . The extension is given geometrically in the following way. Fix a totally geodesic plane P in AdS^3 , which is a copy of hyperbolic plane. Given a point $x \in S$, we define two isometries $\Phi_l^x, \Phi_r^x \in \text{Isom}(\text{AdS}^3)$ which map the tangent plane $T_x S$ to P . The first isometry Φ_l^x is obtained by following the left ruling of $\partial_\infty \text{AdS}^3$. Analogously Φ_r^x for the right ruling. This gives two diffeomorphisms Φ_l and Φ_r from S to P , by

$$\Phi_l(x) = \Phi_l^x(x), \quad \Phi_r(x) = \Phi_r^x(x).$$

The diffeomorphism Φ is then defined as

$$\Phi_{ML} = (\Phi_l)^{-1} \circ \Phi_r.$$

In [KS07, Lemma 3.16] it is shown that the pull-back of the hyperbolic metric h of P on S by means of Φ_r and Φ_l is given by

$$(51) \quad \Phi_l^* h = I((E + JB) \cdot, (E + JB) \cdot),$$

and

$$(52) \quad \Phi_r^* h = I((E - JB) \cdot, (E - JB) \cdot),$$

where I is the first fundamental form of S , J is the almost-complex structure of S , B the shape operator and E the identity. We are now ready to give a relation between the principal curvatures of S and the quasiconformal distortion of Φ :

Proposition 5.5. *Given a maximal surface S in AdS^3 , the quasiconformal distortion of the minimal Lagrangian map $\Phi : \mathbb{H}^2 \rightarrow \mathbb{H}^2$ at a point $y \in \mathbb{H}^2$ is given by*

$$K(y) = \left(\frac{1 + |\lambda(x)|}{1 - |\lambda(x)|} \right)^2,$$

where $y = (\Phi_l)(x)$. Therefore, by taking $K(\Phi_{ML}) = \sup_y K(y)$, namely $K(\Phi_{ML})$ is the maximal dilatation of Φ , the following holds:

$$K(\Phi_{ML}) = \left(\frac{1 + \|\lambda\|_\infty}{1 - \|\lambda\|_\infty} \right)^2.$$

Proof. Let h be the hyperbolic metric of P ; it follows from Equations (51) and (52) that

$$\Phi^* h = h((E + JB)^{-1}(E - JB) \cdot, (E + JB)^{-1}(E - JB) \cdot).$$

The quasiconformal distortion of Φ at a fixed point x can be computed as the ratio between $\sup \|\Phi_*(v)\|$ and $\inf \|\Phi_*(v)\|$ where the supremum and the infimum are taken over all tangent vectors $v \in T_x P$ with $\|v\| = 1$. Since B is diagonalizable with eigenvalues $\pm\lambda$, $(E + JB)^{-1}(E - JB)$ can be diagonalized to the form

$$\begin{pmatrix} \frac{1-\lambda}{1+\lambda} & 0 \\ 0 & \frac{1+\lambda}{1-\lambda} \end{pmatrix}.$$

Hence, assuming $0 < \lambda(x) < 1$, the quasiconformal distortion is given by

$$K(\Phi_l(x)) = \left(\frac{1 + \lambda(x)}{1 - \lambda(x)} \right)^2,$$

as claimed. \square

We are now ready to prove Theorem 2.A.

Theorem 2.A. *There exist universal constants δ and C_1 such that, for any quasisymmetric homeomorphism ϕ of \mathbb{RP}^1 with cross ratio norm $\|\phi\|_{cr} < \delta$, the minimal Lagrangian extension $\Phi_{ML} : \mathbb{H}^2 \rightarrow \mathbb{H}^2$ has maximal dilatation bounded by:*

$$\ln K(\Phi_{ML}) \leq C_1 \|\phi\|_{cr}.$$

Proof. Putting together the inequalities in Proposition 3.A, Theorem 1.A and Proposition 5.5, we obtain the following inequality:

$$(53) \quad K(\Phi) \leq \left(\frac{1 + C_1 \tan w}{1 - C_1 \tan w} \right)^2 \leq \left(\frac{1 + C_1 \sinh\left(\frac{\|\phi\|_{cr}}{2}\right)}{1 - C_1 \sinh\left(\frac{\|\phi\|_{cr}}{2}\right)} \right)^2.$$

Clearly, the inequality (53) holds provided $\|\phi\|_{cr}$ is sufficiently small so that $1 - C \sinh\left(\frac{\|\phi\|_{cr}}{2}\right) > 0$. Since

$$\left. \frac{d}{dx} \right|_{x=0} \ln \left(\frac{1 + C_1 \sinh\left(\frac{x}{2}\right)}{1 - C_1 \sinh\left(\frac{x}{2}\right)} \right)^2 = 2C_1,$$

the claim is proved, by appropriately choosing the constant (which depends on the choice of δ , and is still called C_1). \square

We now move to the proof of Theorem 2.B.

Theorem 2.B. *There exist universal constants Δ and C_2 such that, for any quasisymmetric homeomorphism ϕ of \mathbb{RP}^1 with cross ratio norm $\|\phi\|_{cr} > \Delta$, the minimal Lagrangian extension $\Phi_{ML} : \mathbb{H}^2 \rightarrow \mathbb{H}^2$ has maximal dilatation bounded by:*

$$\ln K(\Phi_{ML}) \leq C_2 \|\phi\|_{cr}.$$

Proof. By Theorem 1.B, we have

$$\frac{1}{1 - \|\lambda\|_\infty} \leq (\tan w)^M$$

and therefore

$$\|\lambda\|_\infty \leq 1 - \frac{1}{(\tan w)^M}.$$

By Proposition 5.5, if K is the maximal dilatation of Φ_{ML} , then

$$(54) \quad K^{1/2} = \frac{1 + \|\lambda\|_\infty}{1 - \|\lambda\|_\infty} \leq \frac{2 - \frac{1}{(\tan w)^M}}{\frac{1}{(\tan w)^M}} = 2(\tan w)^M - 1.$$

Finally, using Proposition 3.A,

$$K^{1/2} \leq 2 \left(\sinh \left(\frac{\|\phi\|_{cr}}{2} \right) \right)^M - 1.$$

Therefore one gets

$$\ln K \leq 2 \left((M/2) \|\phi\|_{cr} + \ln 2 \right) \leq C_2 \|\phi\|_{cr},$$

choosing the constant C_2 sufficiently large, under the (repeatedly used) assumption that $\|\lambda\|_\infty$ is larger than a constant δ (and thus K is larger than some universal constant $K_0 > 1$). \square

Finally, by using the inequalities in Proposition 3.A, Proposition 1.C and Proposition 5.5, we obtain the following estimate:

$$\|\phi\|_{cr} \leq 2 \ln \left(\frac{(\sqrt{K} + 1 - \sqrt{2})(\sqrt{K} + 1 + \sqrt{2})}{(\sqrt{K} - 1 + \sqrt{2})(1 + \sqrt{2} - \sqrt{K})} \right),$$

which clearly holds if the quasiconformal coefficient $K = K(\Phi_{ML})$ of the minimal Lagrangian extension $\Phi_{ML} : \mathbb{D} \rightarrow \mathbb{D}$ is in $[1, (1 + \sqrt{2})^2)$. Let us observe that the function

$$K \mapsto 2 \ln \left(\frac{(\sqrt{K} + 1 - \sqrt{2})(\sqrt{K} + 1 + \sqrt{2})}{(\sqrt{K} - 1 + \sqrt{2})(1 + \sqrt{2} - \sqrt{K})} \right)$$

is differentiable with derivative at 0 equal to 2. Hence the following holds:

Theorem 2.C. *There exist universal constants δ and C_0 such that, for any quasisymmetric homeomorphism ϕ of \mathbb{RP}^1 with cross ratio norm $\|\phi\|_{cr} < \delta$, the minimal Lagrangian extension $\Phi : \mathbb{H}^2 \rightarrow \mathbb{H}^2$ has maximal dilatation bounded by:*

$$C_0 \|\phi\|_{cr} \leq \ln K(\Phi_{ML}).$$

The constant C_0 can be taken arbitrarily close to $1/2$.

In particular, any constant C satisfying the statement of Theorem 2.A cannot be smaller than $1/2$.

Corollary 2.D. *There exists a universal constant C such that, for any quasisymmetric homeomorphism ϕ of \mathbb{RP}^1 , the minimal Lagrangian extension $\Phi_{ML} : \mathbb{H}^2 \rightarrow \mathbb{H}^2$ has maximal dilatation $K(\Phi_{ML})$ bounded by:*

$$\ln K(\Phi_{ML}) \leq C \|\phi\|_{cr}.$$

Proof. In light of Theorem 2.A and Theorem 2.B, it will be sufficient to prove that there exists a constant K_0 such that $\ln K(\Phi_{ML}) \leq K_0$ for all quasisymmetric homeomorphisms ϕ with $\delta \leq \|\phi\|_{cr} \leq \Delta$. To prove this, suppose by contradiction that there exists a sequence ϕ_n with $\delta \leq \|\phi_n\|_{cr} \leq \Delta$ such that the corresponding minimal Lagrangian extensions have maximal dilatation $K_n \rightarrow \infty$. Therefore one can pick a sequence of points x_n on the maximal surface S_n (where of course $\partial_\infty S_n = gr(\phi_n)$) such that the (positive) principal curvature $\lambda_n(x_n)$ tends to 1 as $n \rightarrow \infty$.

By composing with isometries of AdS^3 , we can assume x_n is a fixed point x_0 and all surfaces S_n are tangent to the same plane through x_0 . Indeed, composing with elements of $\text{PSL}(2, \mathbb{R})$ does not change the cross-ratio norm of ϕ_n . By Lemma 1.20 there exists a subsequence n_k such that the maximal surfaces S_{n_k} converge C^∞ on compact sets to an entire maximal surface S_∞ of nonpositive curvature. Using Theorem 1.21 there exists a further subsequence n_{k_j} such that $\phi_{n_{k_j}}$ converges to a quasisymmetric homeomorphism ϕ_∞ . Moreover (see also Remark 1.22), S_∞ has asymptotic boundary $\partial_\infty S_\infty = gr(\phi_\infty)$. But by the C^∞ convergence, the principal curvatures of S_∞ at x_0 are 1 and -1 . By Lemma 1.12, S_∞ is a horospherical surfaces and this gives a contradiction. \square

REFERENCES

- [AAW00] Reiko Aiyama, Kazuo Akutagawa, and Tom Y. H. Wan. Minimal maps between the hyperbolic discs and generalized Gauss maps of maximal surfaces in the anti-de Sitter 3-space. *Tohoku Math. J. (2)*, 52(3):415–429, 2000.
- [ABB⁺07] Lars Andersson, Thierry Barbot, Riccardo Benedetti, Francesco Bonsante, William M. Goldman, François Labourie, Kevin P. Scannell, and Jean-Marc Schlenker. Notes on: “Lorentz spacetimes of constant curvature” [Geom. Dedicata **126** (2007), 3–45; mr2328921] by G. Mess. *Geom. Dedicata*, 126:47–70, 2007.
- [Ahl06] Lars V. Ahlfors. *Lectures on quasiconformal mappings*, volume 38 of *University Lecture Series*. American Mathematical Society, Providence, RI, second edition, 2006. With supplemental chapters by C. J. Earle, I. Kra, M. Shishikura and J. H. Hubbard.
- [BA56] A. Beurling and L. Ahlfors. The boundary correspondence under quasiconformal mappings. *Acta Math.*, 96:125–142, 1956.
- [BBZ07] Thierry Barbot, François Béguin, and Abdelghani Zeghib. Constant mean curvature foliations of globally hyperbolic spacetimes locally modelled on AdS₃. *Geom. Dedicata*, 126:71–129, 2007.
- [BKS11] Francesco Bonsante, Kirill Krasnov, and Jean-Marc Schlenker. Multi-black holes and earthquakes on Riemann surfaces with boundaries. *Int. Math. Res. Not. IMRN*, (3):487–552, 2011.
- [BS10] Francesco Bonsante and Jean-Marc Schlenker. Maximal surfaces and the universal Teichmüller space. *Invent. Math.*, 182(2):279–333, 2010.
- [BS12] Francesco Bonsante and Jean-Marc Schlenker. Fixed points of compositions of earthquakes. *Duke Math. J.*, 161(6):1011–1054, 2012.
- [BS16] Francesco Bonsante and Andrea Seppi. Area-preserving diffeomorphisms of hyperbolic plane and K-surfaces in Anti-de Sitter space. *Preprint, ArXiv: 1610.05701*, 2016.
- [BST17] Francesco Bonsante, Andrea Seppi, and Andrea Tamburelli. On the volume of Anti-de Sitter maximal globally hyperbolic three-manifolds. *Geometric and Functional Analysis*, 2017. <https://doi.org/10.1007/s00039-017-0423-x>.
- [BSW] Francesco Bonsante, Jean-Marc Schlenker, and Michael Wolf. Personal communication.
- [BZ06] Ara Basmajian and Mahmoud Zeinalian. Möbius transformations of the circle form a maximal convergence group. In *The geometry of Riemann surfaces and abelian varieties*, volume 397 of *Contemp. Math.*, pages 1–6. Amer. Math. Soc., Providence, RI, 2006.
- [DE86] Adrien Douady and Clifford J. Earle. Conformally natural extension of homeomorphisms of the circle. *Acta Math.*, 157(1-2):23–48, 1986.
- [FM07] A. Fletcher and V. Markovic. *Quasiconformal maps and Teichmüller theory*, volume 11 of *Oxford Graduate Texts in Mathematics*. Oxford University Press, Oxford, 2007.
- [Gar87] F. P. Gardiner. *Teichmüller theory and quadratic differentials*. Wiley, 1987.
- [GL00] F. P. Gardiner and N. Lakic. *Quasiconformal Teichmüller Theory*. American Mathematical Soc., 2000.
- [GT83] David Gilbarg and Neil S. Trudinger. *Elliptic partial differential equations of second order*, volume 224 of *Grundlehren der Mathematischen Wissenschaften [Fundamental Principles of Mathematical Sciences]*. Springer-Verlag, Berlin, second edition, 1983.
- [HM12] Jun Hu and Oleg Muzician. Cross-ratio distortion and Douady-Earle extension: I. A new upper bound on quasiconformality. *J. Lond. Math. Soc. (2)*, 86(2):387–406, 2012.
- [KS07] Kirill Krasnov and Jean-Marc Schlenker. Minimal surfaces and particles in 3-manifolds. *Geom. Dedicata*, 126:187–254, 2007.
- [Leh83] Matti Lehtinen. The dilatation of Beurling-Ahlfors extensions of quasimetric functions. *Ann. Acad. Sci. Fenn. Ser. A I Math.*, 8(1):187–191, 1983.
- [Mar17] Vladimir Markovic. Harmonic maps and the Schoen conjecture. *J. Amer. Math. Soc.*, 30(3):799–817, 2017.
- [Mes07] Geoffrey Mess. Lorentz spacetimes of constant curvature. *Geom. Dedicata*, 126:3–45, 2007.
- [Sep16] Andrea Seppi. Minimal discs in hyperbolic space bounded by a quasicircle at infinity. *Comment. Math. Helv.*, 91(4):807–839, 2016.
- [Tou15] Jérémy Toulisse. *Minimal Lagrangian diffeomorphisms between hyperbolic cone surfaces and Anti-de Sitter geometry*. PhD thesis, University of Luxembourg, Luxembourg, 2015.

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