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WEAK ERROR FOR THE EULER SCHEME APPROXIMATION OF DIFFUSIONS WITH NON-SMOOTH COEFFICIENTS*

V. KONAKOV¹, A. KOZHINA² AND S. MENOZZI³

Abstract. We study the weak error associated with the Euler scheme of non degenerate diffusion processes with non smooth bounded coefficients. Namely, we consider the cases of Hölder continuous coefficients as well as piecewise smooth drifts with smooth diffusion matrices.

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1. INTRODUCTION

1.1. Setting.

Let $T > 0$ be a fixed given deterministic final horizon and $x \in \mathbb{R}^d$ be an initial starting point. We consider the following multidimensional SDE:

$$X_t = x + \int_0^t b(s, X_s) ds + \int_0^t \sigma(s, X_s) dW_s, \quad t \in [0, T], \quad (1.1)$$

where the coefficients $b : [0, T] \times \mathbb{R}^d \rightarrow \mathbb{R}^d$, $\sigma : [0, T] \times \mathbb{R}^d \rightarrow \mathbb{R}^d \otimes \mathbb{R}^d$ are bounded measurable in time and space and W is a Brownian motion on some filtered probability space $(\Omega, \mathcal{F}, (\mathcal{F}_t)_{t \geq 0}, \mathbb{P})$. We assume that the diffusion matrix $a(t, x) := \sigma \sigma^*(t, x)$ is uniformly elliptic and at least Hölder continuous in time and space. We will consider two kinds of assumptions for the drift coefficient b : either Hölder continuous in time and space (as for the diffusion matrix), or piecewise smooth and having at most a countable set of discontinuities. These assumptions guarantee that (1.1) admits a unique weak solution, see e.g. Bass and Perkins [BP09], [Men11] from which the uniqueness to the martingale problem for the associated generator can be derived under the current assumptions.

Define now for a given $N \in \mathbb{N}^*$ the time step $h := T/N$ and set for all $i \in \llbracket 1, N \rrbracket$, $t_i := ih$ where from now on the notation $\llbracket \cdot, \cdot \rrbracket$ is used to denote an interval of integers. Consider the continuous Euler scheme associated

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with (1.1) whose dynamics writes $X_0^h = x$ and for all $t \in [0, T]$:

$$X_t^h = x + \int_0^t b(\phi(u), X_{\phi(u)}^h) du + \int_0^t \sigma(\phi(u), X_{\phi(u)}^h) dW_u, \quad (1.2)$$

where we set $\phi(u) = \inf\{(t_i)_{i \in \llbracket 0, N-1 \rrbracket} : t_i \leq u < t_{i+1}\}$.

A useful quantity to study, arising in many applicative fields from physics to finance, is the so-called weak error which for a suitable real valued test function f writes:

$$d(f, x, T, h) := \mathbb{E}[f(X_T^{h,0,x})] - \mathbb{E}[f(X_T^{0,x})], \quad (1.3)$$

using the usual Markovian notations, i.e. $X_T^{h,0,x}$, $X_T^{t_i,x}$ respectively stand for the Euler scheme and the diffusion at time T which start at point x at 0.

There is a huge literature concerning the weak error for smooth and/or non-degenerate coefficients, from the seminal paper of Talay and Tubaro [TT90], to the extensions to the hypoelliptic framework [BT96a]. Under those conditions, the quantity $d(f, x, T, h)$ is of order h corresponding to the magnitude of the time step. In the non degenerate framework (under some uniform ellipticity or hypoellipticity conditions) it is even possible to take f to be a Dirac mass in the above expression (1.3). The associated convergence rate remains of order h for the Euler scheme, see [KM02] [BT96b] and $h^{1/2}$ in the more general case of Markov Chain approximations, see e.g [KM00] corresponding to i.i.d. sequences $(\xi_i)_{i \geq 1}$ in (1.2) that are not necessarily Gaussian. In the framework of Lipschitz coefficients we can also mention, in the scalar case, the recent work of Alfonsi *et al.* [AJKH14], who obtained bounds on the Wasserstein distances between the laws of the paths of the diffusion and its Euler scheme. Anyhow, the case of non smooth coefficients, Hölder continuous or less, has rarely been considered. Such cases might anyhow appear very naturally in many applications, when the drifts have for instance discontinuities at some given interfaces or when the diffusion coefficients are very irregular (random media).

In the framework of bounded non degenerate and Hölder continuous coefficients, let us mention the work of Mikulevičius and Platen [MP91] who obtained bounds for the weak error in (1.3) at rate $h^{\gamma/2}$ where $\gamma \in (0, 1)$ is the Hölder exponent of the coefficients b, σ in (1.1) provided $f \in C^{2+\gamma}$. This regularity is essential in that work to apply Itô's formula. Our approach permits to establish that this bound holds true, up to an additional slowly varying factor in the exponent, for the difference of the densities itself, which again corresponds to the weak error (1.3) for a δ -function. We also mention the recent work of Mikulevičius *et al.* [Mik12], [MZ15], concerning some extensions of [MP91] to jump-driven SDEs with Hölder coefficients.

Finally, concerning numerical schemes for diffusions with non-regular coefficients, we refer to the recent work of Kohatsu-Higa *et al.* [KHLY15] who investigate the weak error for possibly discontinuous drifts and diffusion coefficients that are just continuous. We are able to extend some of their controls to the densities, again up to an additional slowly varying factor in the exponent which is due to our smoothing approach.

Indeed our strategy is the following. Under the previous assumptions, both processes $(X_t)_{t \in (0, T]}$ in (1.1) and $(X_{t_i}^h)_{i \in \llbracket 1, N \rrbracket}$ in (1.2) have densities, see e.g. Sheu [She91] for the continuous process and Lemaire and Menozzi [LM10] for the scheme. Let us denote them respectively for $x \in \mathbb{R}^d$, $0 \leq i < j \leq N$, $p(t_i, t_j, x, \cdot)$ and $p^h(t_i, t_j, x, \cdot)$ for the processes starting at time t_i from point x and considered at time t_j . To study the error $(p - p^h)(t_i, t_j, x, y)$ we introduce perturbed dynamics associated with (1.1) and (1.2) respectively. Namely, for a small parameter ε , we mollify suitably the coefficients, the mollification procedure is described in its whole generality in Section 2 and depends on the two considered sets of assumptions indicated above, and consider two additional processes with dynamics:

$$\begin{aligned} X_t^{(\varepsilon)} &= x + \int_0^t b_\varepsilon(s, X_s^{(\varepsilon)}) ds + \int_0^t \sigma_\varepsilon(s, X_s^{(\varepsilon)}) dW_s, \\ X_0^{h,(\varepsilon)} &= x, \quad X_{t_{i+1}}^{h,(\varepsilon)} = X_{t_i}^{h,(\varepsilon)} + b_\varepsilon(t_i, X_{t_i}^{h,(\varepsilon)})h + \sigma_\varepsilon(t_i, X_{t_i}^{h,(\varepsilon)})(W_{t_{i+1}} - W_{t_i}), \end{aligned} \quad (1.4)$$

where $b_\varepsilon, \sigma_\varepsilon$ are mollified versions of b, σ . It is clear that both $(X_t^{(\varepsilon)})_{t \in (0, T]}$ and $(X_{t_i}^{h, (\varepsilon)})_{i \in \llbracket 1, N \rrbracket}$ have densities. The mollified coefficients indeed satisfy uniformly in the mollification parameter the previous assumptions. Let us denote those densities for $x \in \mathbb{R}^d, 0 \leq i < j \leq T$ by $p_\varepsilon(t_i, t_j, x, \cdot), p_\varepsilon^h(t_i, t_j, x, \cdot)$ respectively.

The idea is now to decompose the global error as:

$$(p - p^h)(t_i, t_j, x, y) = (p - p_\varepsilon)(t_i, t_j, x, y) + (p_\varepsilon - p_\varepsilon^h)(t_i, t_j, x, y) + (p_\varepsilon^h - p^h)(t_i, t_j, x, y). \quad (1.5)$$

The key point is that the stability of the densities with respect to a perturbation has been thoroughly investigated for diffusions and Markov Chains in Konakov *et al.* [KKM15]. The results of this work allow to control the differences $p - p_\varepsilon, p_\varepsilon^h - p^h$. On the other hand, since the coefficients $b_\varepsilon, \sigma_\varepsilon$ of $(X_t^\varepsilon)_{t \in [0, T]}, (X_{t_i}^\varepsilon)_{i \in \llbracket 0, N \rrbracket}$ are smooth the central term $p_\varepsilon - p_\varepsilon^h$ in (1.5) can be investigated thanks to the work of Konakov and Mammen [KM02] giving the error expansion at order h on the densities for the weak error. The key point is that the coefficients in the expansion depend on the derivatives of $b_\varepsilon, \sigma_\varepsilon$ which explode when ε goes to zero. This last condition is natural in order to control $p - p_\varepsilon, p_\varepsilon^h - p^h$. Thus two contributions will need to be equilibrated to derive the global error bounds. This will be done through a careful analysis of the densities (heat kernel) of the processes with dynamics described in (1.1), (1.2), (1.4). The estimates required for the error analysis will lead us to refine some bounds previously established by Il'in *et al.* [IKO62]. Let us indicate that this perturbative approach had also been considered by Kohatsu-Higa *et al.* [KHLY15] but for the weak error (1.3) involving at least a continuous function. Our approach, based on parametrix techniques, allows to handle directly the difference of the densities, and gives, up to an additional factor going to zero with the time step, the *expected* convergence rates.

1.2. Assumptions and Main Results.

Let us introduce the following assumptions.

(A1) (Boundedness of the coefficients). The components of the vector-valued function $b(t, x)$ and the matrix-valued function $\sigma(t, x)$ are bounded measurable. Specifically, there exist constants $K_1, K_2 > 0$ s.t.

$$\sup_{(t, x) \in [0, T] \times \mathbb{R}^d} |b(t, x)| \leq K_1, \quad \sup_{(t, x) \in [0, T] \times \mathbb{R}^d} |\sigma(t, x)| \leq K_2.$$

(A2) (Uniform Ellipticity). The diffusion matrix $a := \sigma\sigma^*$ is uniformly elliptic, i.e. there exists $\Lambda \geq 1, \forall (t, x, \xi) \in [0, T] \times (\mathbb{R}^d)^2$,

$$\Lambda^{-1}|\xi|^2 \leq \langle a(t, x)\xi, \xi \rangle \leq \Lambda|\xi|^2.$$

We consider two types of *smoothness* assumptions for the coefficients b, σ in (1.1).

(H) (Hölder drift and diffusion coefficient). The drift b and the diffusion coefficient σ are time-space Hölder continuous in the following sense: for some $\gamma \in (0, 1], \kappa < +\infty$, for all $(s, t) \in \mathbb{R}_+^2, (x, y) \in (\mathbb{R}^d)^2$,

$$|\sigma(s, x) - \sigma(t, y)| + |b(s, x) - b(t, y)| \leq \kappa\{|s - t|^{\gamma/2} + |x - y|^\gamma\}.$$

Observe that the last condition also readily gives, thanks to the boundedness of σ that the diffusion matrix $a = \sigma\sigma^*$ is also uniformly γ -Hölder continuous.

(PS) (Piecewise smooth drift and Smooth diffusion coefficient). The drift b is piecewise smooth, say $C^{2,4}([0, T] \times (\mathbb{R}^d \setminus \mathcal{I}), \mathbb{R}^d)$ where the set of possible discontinuities \mathcal{I} writes as a finite union of smooth, i.e. at least C^4 , closed bounded hypersurfaces with non empty interior that do not intersect. On the other hand we assume that the diffusion coefficient σ is **globally** $C^{2,4}([0, T] \times \mathbb{R}^d, \mathbb{R}^d \otimes \mathbb{R}^d)$.

From now on, we always assume conditions **(A1)**-**(A2)** to be in force. We say that assumption **(A_H)** (resp. **(A_{PS})**) holds if additionally the coefficients satisfy **(H)** (resp. **(PS)**). We will write that **(A)** holds whenever **(A_H)** or **(A_{PS})** is satisfied.

We will denote, from now on, by C a constant depending on the parameters appearing in **(A)** and T . We reserve the notation c for constants that only depend on **(A)** but not on T . The values of C, c may change from line to line. Other possible dependencies will be explicitly specified.

Theorem 1 (Error for the Euler scheme of a diffusion with Hölder coefficients). *Let $T > 0$ be fixed and consider a given time step $h := T/N$, for $N \in \mathbb{N}^*$. Set for $i \in \mathbb{N}$, $t_i := ih$. Under **(A_H)**, there exist $C \geq 1, c \in (0, 1]$ s.t. for all $0 \leq t_i < t_j \leq T$ s.t. $(t_j - t_i) \geq h^{2/(2+\gamma)}$ and $(x, y) \in (\mathbb{R}^d)^2$:*

$$p_c(t_j - t_i, y - x)^{-1} |(p - p^h)(t_i, t_j, x, y)| \leq \frac{C}{(t_j - t_i)^{(1-\gamma/2)\gamma/2}} h^{\frac{\gamma}{2} - C\psi(h)}, \quad (1.6)$$

where p, p^h respectively stand for the densities of the diffusion X and its Euler approximation X^h with time step h , for all $(t, z) \in \mathbb{R}_*^+ \times \mathbb{R}^d$, $p_c(t, z) := \frac{c^{d/2}}{(2\pi t)^{d/2}} \exp(-c \frac{|z|^2}{2t})$ and $\psi(h) = \frac{\log_3(h^{-1})}{\log_2(h^{-1})}$ where \log_k denotes for $k \in \mathbb{N}^*$ the k^{th} iterated logarithm. Let us observe that $\psi(h) \xrightarrow{h \rightarrow 0} 0$. If we are now interested in the weak error in the sense of (1.3), if a function f satisfies the growth condition:

$$\exists c_0 < c/(2T), C_0 > 0, \forall x \in \mathbb{R}^d, |f(x)| \leq C_0(1 + \exp(c_0|x|^2)), \quad (1.7)$$

and additionally $f \in C^\beta(\mathbb{R}^d, \mathbb{R})$, $\beta \in (0, 1)$, then:

$$|\mathbb{E}[f(X_{t_j}^{h, t_i, x})] - \mathbb{E}[f(X_{t_j}^{t_i, x})]| \leq C_f h^{\gamma/2}, \quad (1.8)$$

using again the usual Markovian notations, i.e. $X_{t_j}^{h, t_i, x}, X_{t_j}^{t_i, x}$ respectively stand for the Euler scheme and the diffusion at time t_j which start at point x at t_i .

Eventually, If we consider a connected Borel set A with (piecewise) smooth boundary and non empty interior when A is bounded, we also get that for all $x \in \mathbb{R}^d$ s.t. $d(x, \partial A) \geq (t_j - t_i)^{1/2} h^{\gamma/2}$ and $\eta > 0$ s.t. $\eta d(x, \partial A)^\eta \geq h^{C\psi(h)}$:

$$|\mathbb{E}[\mathbb{I}_{X_{t_j}^{h, t_i, x} \in A}] - \mathbb{E}[\mathbb{I}_{X_{t_j}^{t_i, x} \in A}]| \leq C \left\{ \frac{1}{\eta d(x, \partial A)^\eta} + 1 \right\} h^{\gamma/2}, \quad (1.9)$$

where $d(\cdot, \partial A)$ stands for the distance to the boundary of A .

Remark 1. We point out that this result is to be compared with the one obtained by Mikulevičius and Platen [MP91] for the weak error. The framework they considered is similar to ours, and their main results consists in controlling at rate $h^{\gamma/2}$ the weak error $d(f, x, T, h) = \mathbb{E}_x[f(X_T^{h, 0, x})] - \mathbb{E}[f(X_T^{0, x})]$ for a smooth function $f \in C^{2+\gamma}$. The above theorem establishes that $|d(f, x, T, h)| \leq Ch^{(\gamma-\psi(h))/2}$ as soon as f is measurable and satisfies the growth condition (1.7). This control can be useful for specific and relevant applications, like for instance quantile estimation (that would involve functions of the form $f(x) = \mathbb{I}_{|x| \leq K}$ or $f(x) = \mathbb{I}_{|x| \leq K} \exp(c|x|)$) that appear in many applications: default probabilities in mathematical finance, fatigue of structures in random mechanics. We are able to find the expected convergence rate up to a vanishing contribution. The rate $h^{\gamma/2}$ again holds, without the additional term, as soon as $f \in C^\beta(\mathbb{R}^d, \mathbb{R})$, $\beta \in (0, 1]$. The contribution in $\psi(h)$ appearing in (1.6), which slightly deteriorates the convergence, seems to be, with our approach, the price to pay to get rid of any smoothness on f .

Remark 2 (About the Convergence Rate). We also emphasize that the convergence rate in $h^{\gamma/2}$ is closer to a rate associated with a strong error. It indeed corresponds to the typical magnitude of the quantity $\mathbb{E}[|W_h|^\gamma] \leq c_\gamma h^{\gamma/2}$, which reflects the variations, on one time-step of length h , of the Euler scheme with Hölder coefficients.

Indeed, under (\mathbf{A}_H) , for all $i \in \llbracket 0, N-1 \rrbracket$:

$$\begin{aligned} \mathbb{E}\left[\sup_{u \in [t_i, t_{i+1}]} |b(u, X_u^h) - b(t_i, X_{t_i}^h)|\right] + \mathbb{E}\left[\sup_{u \in [t_i, t_{i+1}]} |\sigma(u, X_u^h) - \sigma(t_i, X_{t_i}^h)|\right] &\leq \kappa \left\{ h^{\gamma/2} + \mathbb{E}\left[\sup_{u \in [t_i, t_{i+1}]} |X_u^h - X_{t_i}^h|^\gamma\right] \right\} \\ &\leq \kappa \left\{ h^{\gamma/2} + \mathbb{E}\left[\left\{\sup_{u \in [t_i, t_{i+1}]} |\sigma(t_i, X_{t_i}^h)(W_u - W_{t_i})| + K_1 h\right\}^2\right]^{\gamma/2} \right\} \leq ch^{\gamma/2}. \end{aligned} \quad (1.10)$$

These terms typically appear in the error analysis when there is low regularity of the coefficients or of the value function $v(t, x) := \mathbb{E}[f(X_T^{t,x})]$. Under the previous assumptions, if the function f is bounded and is $C^{2+\gamma}(\mathbb{R}^d, \mathbb{R})$, $\gamma \in (0, 1)$ it is then well known, see e.g. Friedman [Fri64] or Ladyzhenskaya et al. [LSU68] that $v \in C^{1+\gamma/2, 2+\gamma}([0, T] \times \mathbb{R}^d, \mathbb{R})$. Also v satisfies the parabolic PDE $(\partial_t v + L_t v)(t, x) = 0$, $(t, x) \in [0, T] \times \mathbb{R}^d$, where L_t stands for the generator of (1.1) at time t , i.e. for all $\varphi \in C_0^2(\mathbb{R}^d, \mathbb{R})$, $x \in \mathbb{R}^d$,

$$L_t \varphi(x) = b(t, x) \cdot \nabla_x \varphi(x) + \frac{1}{2} \text{Tr}(a(t, x) D_x^2 \varphi(x)).$$

Recalling that $t_0 = 0, t_N = T$, we decompose the error as:

$$\begin{aligned} d(f, x, T, h) &:= \mathbb{E}_x[f(X_T^h)] - \mathbb{E}_x[f(X_T)] = \sum_{i=0}^{N-1} \mathbb{E}_x[v(t_{i+1}, X_{t_{i+1}}^h) - v(t_i, X_{t_i}^h)] \\ &= \sum_{i=0}^{N-1} \mathbb{E}_x \left[\int_{t_i}^{t_{i+1}} \left\{ \partial_s v(s, X_s^h) + \nabla_x v(s, X_s^h) \cdot b(t_i, X_{t_i}^h) + \frac{1}{2} \text{Tr}(D_x^2 v(s, X_s^h) a(t_i, X_{t_i}^h)) \right\} ds \right] \\ &= \sum_{i=0}^{N-1} \mathbb{E}_x \left[\int_{t_i}^{t_{i+1}} \left\{ \partial_s v + L_s v \right\} (s, X_s^h) ds \right] \\ &+ \mathbb{E}_x \left[\int_{t_i}^{t_{i+1}} \left\{ \nabla_x v(s, X_s^h) \cdot (b(t_i, X_{t_i}^h) - b(s, X_s^h)) + \frac{1}{2} \text{Tr}(D_x^2 v(s, X_s^h) (a(t_i, X_{t_i}^h) - a(s, X_s^h))) \right\} ds \right] \\ &= \sum_{i=0}^{N-1} \mathbb{E}_x \left[\int_{t_i}^{t_{i+1}} \left\{ \nabla_x v(s, X_s^h) \cdot (b(t_i, X_{t_i}^h) - b(s, X_s^h)) + \frac{1}{2} \text{Tr}(D_x^2 v(s, X_s^h) (a(t_i, X_{t_i}^h) - a(s, X_s^h))) \right\} ds \right], \end{aligned} \quad (1.11)$$

exploiting the PDE satisfied by v for the last equality. For a bounded f in $C^{2+\gamma}(\mathbb{R}^d, \mathbb{R})$, the derivatives up to order two are globally bounded. We are thus led to control in (1.11) quantities similar to those appearing in (1.10). The associated bound then precisely gives the convergence rate. The analysis extends if f is simply $C^\beta(\mathbb{R}^d, \mathbb{R})$, $\beta \in (0, 1)$. In that case the second derivatives yield an integrable singularity in time for the second order partial derivatives (see Proposition 4 which holds under the sole assumption (\mathbf{A}_H) for multi-indices α , $|\alpha| \leq 2$ and the proof of Theorem 1 in Section 3.3).

Remark 3. Even though we have considered $\gamma \in (0, 1]$, our analysis should be extendable in the framework of Hölder spaces to $\gamma \in (1, 2]$. On the other hand, Theorem 1 specifies the time-singularity in small time.

Remark 4. We feel that the bounds of Theorem 1 are relevant for functions which are **truly** Hölder continuous, that is for coefficients that would involve some simple transformations of the Weierstrass functions, see e.g. [Zyg36], or of an independent Brownian sample path in order that (\mathbf{A}_H) is fulfilled. Indeed, for functions which are just locally Hölder continuous, like the mapping $x \mapsto 1 + |x|^\alpha \wedge K$, $\alpha \in (0, 1]$, we think that it would be more appropriate to study some local regularizations, close to the neighborhoods of **real** Hölder continuity (0

and $K^{1/\alpha}$ for the indicated example) and to exploit that outside of these neighborhoods, the **usual** sufficient smoothness is available. For such coefficients we think that the convergence rates might be definitely better.

Theorem 2 (Error for the Euler Scheme with Smooth Diffusion Coefficients and Piecewise Smooth Drift). *Let $T > 0$ be fixed and (\mathbf{A}_{PS}) be in force. With the notations of Theorem 1 we have that:*

- there exist $C \geq 1, c \in (0, 1]$ s.t. for all $0 \leq t_i < t_j \leq T$ s.t. $(t_j - t_i) \geq h^{1/2}$ and $(x, y) \in (\mathbb{R}^d)^2$:

$$p_c(t_j - t_i, y - x)^{-1} |(p - p^h)(t_i, t_j, x, y)| \leq C \left(\frac{h}{t_j - t_i} \right)^{1/(d+1) - C\psi(h)}. \quad (1.12)$$

- In the special case $\sigma(t, x) = \sigma$, i.e. constant diffusion coefficient¹, the previous bound improves to:

$$p_c(t_j - t_i, y - x)^{-1} |(p - p^h)(t_i, t_j, x, y)| \leq C \left(\frac{h}{(t_j - t_i)^{1/2}} \right)^{1/d - C\psi(h)}. \quad (1.13)$$

Remark 5. *This result emphasizes that as soon as the drift is irregular a true diffusion coefficient deteriorates the convergence rate. This is clear since, in that case, the difference of the densities $p - p^h$ involves higher derivatives of densities of processes with mollified coefficients which are more explosive (see Section 3.2). If the diffusion coefficient does not depend on space, we find, up to the additional term in $\psi(h)$, the usual convergence rate for the weak error if $d = 1$. However, our regularization approach clearly feels the dimension, when doing e.g. Hölder inequalities on neighborhoods of the hypersurfaces of discontinuity, and the convergence rates decrease with the dimension.*

Let us carefully mention that considering the weak error $d(f, x, t_i, t_j, h) := \mathbb{E}[f(X_{t_j}^{h, t_i, x})] - \mathbb{E}[f(X_{t_j}^{t_i, x})]$ for smooth functions f and not Dirac masses as we do, should improve the convergence rates and in particular allow to get rid of the terms in $\psi(h)$ through a careful investigation of the derivatives of the associated heat kernel, see the estimates of Proposition 4 that could be refined when considering an additional integration w.r.t. to the final variable.

1.3. On Some Related Applications.

1.3.1. Some Approximating Dynamics for Interest Rates.

A very popular model for interest rates in the financial literature is the Cox-Ingersoll-Ross process with dynamics:

$$dX_t = (a - kX_t)dt + \sigma|X_t|^{1/2}dW_t, \quad (1.14)$$

for given parameters $\sigma, k, a > 0$. From the numerical viewpoint, the behavior of the Euler scheme is not standard. For a given time-step h , the strong error was indeed proved to be, as in the usual Lipschitz case, of order $h^{1/2}$ in Berkaoui *et al.* [BBD08] provided a is not *too small*. On the other hand, numerical experiments in Alfonsi [Alf05] emphasized very slow convergence, of order $(-\ln h)^{-1}$, for small values of a . This convergence order has been established by Gyöngy and Rásonyi [GR11].

Of course the dynamics in (1.14) does not enter our framework, since it is closer to the dynamics of a Bessel-like process whose density does not satisfy Gaussian bounds. However we could introduce for positive parameters η, K which are respectively meant to be small and large enough the dynamics:

$$dX_t = (a - kX_t)dt + (\eta + \sigma|X_t|^{1/2} \wedge K)dW_t. \quad (1.15)$$

The diffusion coefficient $\tilde{\sigma}(x) = (\eta + \sigma|x|^{1/2} \wedge K)$ is then uniformly elliptic, $1/2$ Hölder continuous and bounded. On the other hand the drift is not bounded but the analysis of Theorem 1 would still hold true thanks to the

¹the case of an inhomogeneous diffusion coefficient independent of x , i.e. $\sigma(t, x) = \sigma(t)$ could also be handled provided the Gaussian part is simulated exactly in a modified Euler scheme.

work of Konakov and Markova [KM15] that allows to get rid of the linear drift through a suitable transformation. We would then derive a convergence of order $h^{1/4-C\psi(h)}$ at least for the associated Euler scheme on the densities (see also Remark 4). Even though the marginals in (1.15) enjoy Gaussian bounds, see e.g. [DM10], the expected properties for an interest rate dynamics, mean reverting and positivity, should hold with some high probability. Also, the difference between the approximate dynamics in (1.15) and the original one in (1.14) might be investigated through stochastic analysis tools (occupation times).

1.3.2. Extension to some Kinetic Models

The results of Theorems 1 and 2 should extend without additional difficulties to the case of degenerate diffusions of the form:

$$\begin{aligned} dX_t^1 &= b(t, X_t)dt + \sigma(t, X_t)dW_t, \\ dX_t^2 &= X_t^1 dt, \end{aligned} \tag{1.16}$$

denoting $X_t = (X_t^1, X_t^2)$, under the same previous assumptions (\mathbf{A}_H) or (\mathbf{A}_{PS}) on b, σ . The sensitivity analysis when we consider perturbations of the non-degenerate components, i.e. for a given $\varepsilon > 0$:

$$\begin{aligned} dX_t^{1,(\varepsilon)} &= b_\varepsilon(t, X_t^{(\varepsilon)})dt + \sigma_\varepsilon(t, X_t^{(\varepsilon)})dW_t, \\ dX_t^{2,(\varepsilon)} &= X_t^{1,(\varepsilon)} dt, \end{aligned} \tag{1.17}$$

has been performed by Kozhina [Koz16] following [KKM15]. The key point is that under (\mathbf{A}) , the required parametrix expansions of the densities associated with the solutions of equation (1.16), (1.17) were established in [KMM10]. The analysis of the derivatives of the heat kernel, that would require to extend the results of Section 3 to the considered degenerate setting will concern further research.

The paper is organized as follows. We first introduce a suitable mollification procedure of the coefficients in Section 2 and derive from the stability results of Konakov *et al.* [KKM15] how the error of the mollifying procedure is then reflected on the densities. We then give in Section 3 some pointwise controls on the derivatives of the heat-kernels with mollified coefficients. From these controls and the previous error expansion obtained for the Euler scheme with smooth coefficients by Konakov and Mammen [KM02] we establish our main estimates. Eventually, Section 4 is dedicated to the proof of the controls stated in Section 3. These proofs are based on the parametrix expansions of the underlying densities following the Mc-Kean and Singer approach [MS67].

2. MOLLIFICATION OF THE COEFFICIENTS AND STABILITY RESULTS

In order to apply the strategy described in the introduction for the error analysis we first need to *regularize* in an appropriate manner the coefficients. The mollifying procedures differ under our two sets of assumptions.

2.1. Mollification under (\mathbf{A}_H) (Hölder continuous coefficients)

In this case both coefficients b, σ need to be globally regularized in time and space. We introduce the mollified coefficients defined for all $(t, x) \in [0, T] \times \mathbb{R}^d$ and $\varepsilon > 0$ by

$$b_{\varepsilon,S}(t, x) := b(t, \cdot) * \rho_\varepsilon(x), \sigma_{\varepsilon,S}(t, x) := \sigma(t, \cdot) * \rho_\varepsilon(x), \tag{2.1}$$

where $*$ stands for the spatial convolution and for $\varepsilon > 0$, ρ_ε is a spatial mollifier, i.e. for all $x \in \mathbb{R}^d$,

$$\rho_\varepsilon(x) := \varepsilon^{-d} \rho(x/\varepsilon), \rho \in C^\infty(\mathbb{R}^d, \mathbb{R}^+), \int_{\mathbb{R}^d} \rho(y) dy = 1, |\text{supp}(\rho)| \subset K,$$

for some compact set $K \subset \mathbb{R}^d$. The subscript S in $b_{\varepsilon,S}, \sigma_{\varepsilon,S}$ appears to emphasize that the *spatial* convolution is considered. We will also need a mollification in time when the coefficients are inhomogeneous. Up to a symmetrization in time of the coefficients b, σ , i.e. we set for all $(t, x) \in [0, T] \times \mathbb{R}^d$, $b(-t, x) = b(t, x)$, $\sigma(-t, x) = \sigma(t, x)$ we can define:

$$b_\varepsilon(t, x) = b_{\varepsilon,S}(\cdot, x) \star \zeta_{\varepsilon^2}(t), \sigma_\varepsilon(t, x) = \sigma_{\varepsilon,S}(\cdot, x) \star \zeta_{\varepsilon^2}(t), \quad (2.2)$$

where \star stands for the time convolution and for $s \in \mathbb{R}$, $\zeta_{\varepsilon^2}(s) := \varepsilon^{-2} \zeta(s/\varepsilon^2)$, ζ being a scalar mollifier with compact support in $[-T, T]$. The *complete* regularization in the spatial and time variable reflects the usual parabolic scaling. This feature will be crucial to balance the singularities appearing in our analysis (see Propositions 4, 5 and their proofs below). We have the following controls.

Proposition 1 (First Controls on the Mollified Coefficients). *Assume that (\mathbf{A}_H) is in force. Then, there exists $C \geq 1$ s.t. for all $\varepsilon > 0$,*

$$\begin{aligned} \Delta_{\varepsilon,b} &:= \sup_{(t,x) \in [0,T] \times \mathbb{R}^d} |b(t, x) - b_\varepsilon(t, x)| \leq C\varepsilon^\gamma, \\ \forall \eta \in (0, \gamma), \Delta_{\varepsilon,\sigma,\eta} &:= \sup_{(t,x) \in [0,T] \times \mathbb{R}^d} |\sigma(t, x) - \sigma_\varepsilon(t, x)| + \sup_{t \in [0,T]} |(\sigma - \sigma_\varepsilon)(t, \cdot)|_\eta \\ &\leq C(\varepsilon^\gamma + \varepsilon^{\gamma-\eta}), \end{aligned} \quad (2.3)$$

where for a given function $f : \mathbb{R}^d \rightarrow \mathbb{R}$, we denote for $\eta \in (0, 1)$, $|f|_\eta := \sup_{(x,y) \in \mathbb{R}^{2d}, x \neq y} \frac{|f(x) - f(y)|}{|x - y|^\eta}$.

Proof. Write first for all $(t, x) \in [0, T] \times \mathbb{R}^d$:

$$b(t, x) - b_{\varepsilon,S}(t, x) := \int_{\mathbb{R}^d} \{b(t, x) - b(t, y)\} \rho_\varepsilon(x - y) dy = \int_{\mathbb{R}^d} \{b(t, x) - b(t, x - z\varepsilon)\} \rho(z) dz.$$

From the Hölder continuity of b assumed in **(H)** and the above equation, we deduce that $b_{\varepsilon,S}$ satisfies **(H)** as well and that:

$$\sup_{(t,x) \in [0,T] \times \mathbb{R}^d} |(b - b_{\varepsilon,S})(t, x)| \leq C_\rho \varepsilon^\gamma, \quad C_\rho := \kappa \int_K |z|^\gamma \rho(z) dz. \quad (2.4)$$

The same analysis can be performed for $\sigma_{\varepsilon,S}$, so that $\sigma_{\varepsilon,S}$ satisfies **(H)** and $|\sigma(t, \cdot) - \sigma_{\varepsilon,S}(t, \cdot)|_\infty \leq C_\rho \varepsilon^\gamma$. In particular, $b_{\varepsilon,S}, \sigma_{\varepsilon,S}$ are both $\gamma/2$ -Hölder continuous in time uniformly in $\varepsilon > 0$. Repeating the previous arguments replacing ρ_ε by ζ_{ε^2} , we therefore deduce $|b_{\varepsilon,S} - b_\varepsilon|_\infty + |\sigma_{\varepsilon,S} - \sigma_\varepsilon|_\infty \leq C_\zeta \varepsilon^\gamma$, which eventually yields:

$$|b - b_\varepsilon|_\infty + |\sigma - \sigma_\varepsilon|_\infty \leq C\varepsilon^\gamma.$$

This gives the controls concerning the sup norms in (2.3).

Let us now turn to the Hölder norm. Observe first that, for all $t \in \mathbb{R}^+$, $(x, y) \in (\mathbb{R}^d)^2$:

$$\begin{aligned} & \{\sigma(t, x) - \sigma_{\varepsilon,S}(t, x)\} - \{\sigma(t, y) - \sigma_{\varepsilon,S}(t, y)\} \\ &= \int_{\mathbb{R}^d} \left\{ [\sigma(t, x) - \sigma(t, x - z\varepsilon)] - [\sigma(t, y) - \sigma(t, y - z\varepsilon)] \right\} \rho(z) dz, \\ & \quad \{\sigma_\varepsilon(t, x) - \sigma_{\varepsilon,S}(t, x)\} - \{\sigma_\varepsilon(t, y) - \sigma_{\varepsilon,S}(t, y)\} \\ &= \int_{\mathbb{R}} \left\{ [\sigma_{\varepsilon,S}(t - \varepsilon^2 u, x) - \sigma_{\varepsilon,S}(t, x)] - [\sigma_{\varepsilon,S}(t - \varepsilon^2 u, y) - \sigma_{\varepsilon,S}(t, y)] \right\} \zeta(u) du. \end{aligned}$$

It readily follows from the γ -Hölder continuity in space (resp. $\gamma/2$ -Hölder continuity in time) of σ (resp. $\sigma_{\varepsilon,S}$) that one has the following controls:

$$\begin{aligned} |[\sigma(t,x) - \sigma_\varepsilon(t,x)] - [\sigma(t,y) - \sigma_\varepsilon(t,y)]| &\leq C(|x-y|^\gamma \wedge \varepsilon^\gamma) \leq C|x-y|^\eta \varepsilon^{\gamma-\eta}, \\ |(\sigma - \sigma_\varepsilon)(t, \cdot)|_\eta &\leq C\varepsilon^{\gamma-\eta}, \quad \eta \in (0, \gamma). \end{aligned} \quad (2.5)$$

This completes the proof. \square

We will need as well some controls on the derivatives of the mollified coefficients.

Proposition 2 (Controls on the Derivatives of the Mollified Coefficients). *Under the assumptions of Proposition 1, we have that there exists $C \geq 1$ s.t. for all $\varepsilon \in (0, 1)$ and for all multi-index α , $|\alpha| \in \llbracket 1, 4 \rrbracket$:*

$$|D_x^\alpha b_\varepsilon|_\infty + |D_x^\alpha \sigma_\varepsilon|_\infty \leq C\varepsilon^{-|\alpha|+\gamma}, \quad \sup_{t \in [0, T]} |D_x^\alpha \sigma_\varepsilon(t, \cdot)|_\gamma \leq C\varepsilon^{-|\alpha|}. \quad (2.6)$$

Also, there exists a constant C s.t.:

$$|\partial_t \sigma_\varepsilon|_\infty \leq C\varepsilon^{-2+\gamma}, \quad \sup_{t \in [0, T]} |\partial_t \sigma_\varepsilon(t, \cdot)|_\eta \leq C\varepsilon^{-2+\gamma-\eta}, \quad \forall \eta \in (0, \gamma]. \quad (2.7)$$

Proof. For all multi-index α , $|\alpha| \in \llbracket 1, 4 \rrbracket$ and $(t, x) \in [0, T] \times \mathbb{R}^d$ and all $\varepsilon > 0$:

$$D_x^\alpha \sigma_{\varepsilon,S}(t, x) = \int_{\mathbb{R}^d} \sigma(t, z) D_x^\alpha \rho_\varepsilon(x-z) dz = \int_{\mathbb{R}^d} [\sigma(t, z) - \sigma(t, x)] D_x^\alpha \rho_\varepsilon(x-z) dz.$$

Indeed, setting for all $x \in \mathbb{R}^d$, $g_\varepsilon(x) := \int_{\mathbb{R}^d} \rho_\varepsilon(x-z) dz = 1$ we have $D_x^\alpha g_\varepsilon(x) := \int_{\mathbb{R}^d} D_x^\alpha \rho_\varepsilon(x-z) dz = 0$. Thus, since $|D_x^\alpha \rho_\varepsilon(x-z)| \leq \varepsilon^{-(|\alpha|+d)} |D_w^\alpha \rho(w)|_{w=\frac{x-z}{\varepsilon}}$, we derive:

$$\begin{aligned} |D_x^\alpha \sigma_{\varepsilon,S}(t, x)| &\leq \int_{\mathbb{R}^d} |\sigma(t, z) - \sigma(t, x)| \varepsilon^{-(|\alpha|+d)} |D_w^\alpha \rho(w)|_{w=\frac{x-z}{\varepsilon}} dz \\ &\leq \kappa \varepsilon^{-|\alpha|+\gamma} \int_{\mathbb{R}^d} \left(\frac{|z-x|}{\varepsilon} \right)^\gamma \varepsilon^{-d} |D_w^\alpha \rho(w)|_{w=\frac{x-z}{\varepsilon}} dz \leq c \varepsilon^{-|\alpha|+\gamma}, \end{aligned}$$

exploiting the Hölder continuity assumption **(H)** for σ in the last but one inequality and the assumptions on ρ for the last one. Similarly, we derive for all $(t, x, y) \in [0, T] \times (\mathbb{R}^d)^2$ and all $\varepsilon > 0$:

$$\begin{aligned} |D_x^\alpha \sigma_{\varepsilon,S}(t, x) - D_x^\alpha \sigma_{\varepsilon,S}(t, y)| &\leq \int_{\mathbb{R}^d} |\sigma(t, x-z) - \sigma(t, y-z)| \varepsilon^{-(|\alpha|+d)} |D_w^\alpha \rho(w)|_{w=\frac{z}{\varepsilon}} dz \\ &\leq C \kappa \varepsilon^{-|\alpha|} |x-y|^\gamma. \end{aligned}$$

The same bounds hold for $b_{\varepsilon,S}$ as well. The previous controls readily imply (2.6) since the additional time convolution does not have any impact here.

Proceeding similarly for the time convolution, exploiting the $\gamma/2$ -Hölder continuity in time of $\sigma_{\varepsilon,S}$, the bounds in (2.7) can be derived similarly. This completes the proof. \square

2.2. Mollification Under (A_{PS}) (Piecewise smooth drift and Smooth Diffusion Coefficient).

In this case we only need to regularize the drift in a neighborhood of the discontinuity points in \mathcal{I} . Let us denote by $m \in \mathbb{N}^*$, the finite number of hypersurfaces of discontinuities and write $\mathcal{I} := \cup_{i=1}^m \mathcal{S}_i$, where each \mathcal{S}_i

is a bounded hypersurface. For a given parameter $\varepsilon > 0$, define its neighborhood $\mathcal{V}_\varepsilon(\mathcal{I}) := \cup_{i=1}^m \mathcal{V}_\varepsilon(\mathcal{S}_i)$, where for $i \in \llbracket 1, m \rrbracket$, $\mathcal{V}_\varepsilon(\mathcal{S}_i) := \{z \in \mathbb{R}^d : d(z, \mathcal{S}_i) \leq \varepsilon\}$. The fact is now that we set $b_\varepsilon(t, x) = b(t, x)$ on $\mathbb{R}^d \setminus \mathcal{V}_\varepsilon(\mathcal{I})$ and perform a smooth mollification on the neighborhood $\mathcal{V}_\varepsilon(\mathcal{I})$ of the discontinuities. Of course we have that $|(b - b_\varepsilon)(t, x)| \leq C \mathbb{I}_{x \in \mathcal{V}_\varepsilon(\mathcal{I})}$ which is not necessarily small. Anyhow, for all $q > 1$, recalling the $(\mathcal{S}_i)_{i \in \mathbb{Z}}$ are bounded, we derive as well:

$$\|b - b_\varepsilon\|_{L^q([0, T] \times \mathbb{R}^d)} = \left\{ \int_0^T dt \int_{\mathbb{R}^d} |(b - b_\varepsilon)(t, x)|^q dx \right\}^{1/q} \leq C \left\{ \int_0^T dt \int_{\mathcal{V}_\varepsilon(\mathcal{I})} dx \right\}^{1/q} \leq C \varepsilon^{1/q} \sigma_g(\mathcal{I})^{1/q} T^{1/q}. \quad (2.8)$$

Observe as well that the following control holds for the derivatives of the mollified coefficient. For all multi-index α , $|\alpha| \leq 4$, there exists $C := C((\mathbf{A}))$ s.t. for all $(t, x) \in [0, T] \times \mathbb{R}^d$:

$$|\partial_x^\alpha b_\varepsilon(t, x)| \leq C \{ \varepsilon^{-|\alpha|} \mathbb{I}_{x \in \mathcal{V}_\varepsilon(\mathcal{I})} + \mathbb{I}_{x \notin \mathcal{V}_\varepsilon(\mathcal{I})} \}. \quad (2.9)$$

Under the considered assumptions it is not necessary to mollify the diffusion coefficients. Under the current assumption we thus set for all $(t, x) \in [0, T] \times \mathbb{R}^d$, $\sigma_\varepsilon(t, x) = \sigma(t, x)$, in order to keep homogeneous notations under our two running assumptions for the drift.

2.3. Stability Results

Recall now that under (\mathbf{A}_H) or (\mathbf{A}_{PS}) equation (1.1) admits a density (see e.g. [She91]), i.e. for all $0 \leq s < t \leq T$, $x \in \mathbb{R}^d$, $B \in \mathcal{B}(\mathbb{R}^d)$, $\mathbb{P}[X_t \in B | X_s = x] = \int_B p(s, t, x, y) dy$. The same holds for the Euler scheme in (1.2) (see e.g. [LM10]), for all $0 \leq i < j \leq N$, $x \in \mathbb{R}^d$, $\mathbb{P}[X_{t_j}^h \in B | X_{t_i}^h = x] = \int_B p^h(t_i, t_j, x, y) dy$. These properties remain valid for the respective perturbed diffusion and Euler scheme whose coefficients correspond to the procedures described in Section 2.1 and Section 2.2 depending on whether assumption (\mathbf{A}_H) or (\mathbf{A}_{PS}) is in force. Let us denote the densities associated with the perturbed diffusion and discretization scheme by p_ε and p_ε^h respectively.

We readily get from Theorems 1 and 2 in [KKM15] the following sensitivity results.

Lemma 1 (Sensitivity under (\mathbf{A}_H)). *Under Assumption (\mathbf{A}_H) , for $\eta \in (0, \gamma)$ there exist $C_\eta \geq 1$, $c \leq 1$ s.t. for all $0 \leq i < j \leq N$, $(x, y) \in (\mathbb{R}^d)^2$:*

$$|(p - p_\varepsilon)(t_i, t_j, x, y)| + |(p^h - p_\varepsilon^h)(t_i, t_j, x, y)| \leq C_\eta \varepsilon^{\gamma - \eta} p_c(t_j - t_i, y - x). \quad (2.10)$$

Lemma 2 (Sensitivity under (\mathbf{A}_{PS})). *Under Assumption (\mathbf{A}_{PS}) , for $q > d$ there exist $C_q \geq 1$, $c \leq 1$ s.t. for all $0 \leq i < j \leq N$, $(x, y) \in (\mathbb{R}^d)^2$:*

$$|(p - p_\varepsilon)(t_i, t_j, x, y)| + |(p^h - p_\varepsilon^h)(t_i, t_j, x, y)| \leq C_q \varepsilon^{1/q} p_c(t_j - t_i, y - x). \quad (2.11)$$

Let us mention that the constants C_η, C_q in equations (2.10) and (2.11) respectively explode when η goes to 0 and q goes to d . On the other hand it is very important to specify the explosion rates in order to equilibrate the global errors. This is the purpose of the next proposition.

Proposition 3 (Explosion of the constants in the Sensitivity analysis). *Under (\mathbf{A}_H) , there exists $C := C((\mathbf{A}_H), T) \geq 1$ s.t. for $\eta \in (0, \gamma)$ we have in equation (2.10):*

$$C_\eta \leq C \exp(C(2\eta^{-1} + 1)^{2\eta^{-1} + 1}). \quad (2.12)$$

Under (\mathbf{A}_{PS}) , there exists $C := C((\mathbf{A}_{PS}), T) \geq 1$ s.t. for $q > d$ we have in equation (2.11):

$$C_q \leq C \exp(C(\alpha(q)^{-1} + 1)^{\alpha(q)^{-1} + 1}), \quad \alpha(q) = \frac{1}{2} \left(1 - \frac{d}{q}\right). \quad (2.13)$$

Proof. Under (\mathbf{A}_H) . To obtain equation (2.10) with exponent $\varepsilon^{\gamma-\eta}$, we explicitly chose to exploit the controls from Theorem 1 and 2 in [KKM15] with a control on the difference $|\sigma - \sigma_\varepsilon|_\eta$. This possibly *small* Hölder index induces some explosions in the constants appearing in the sensitivity analysis. Those explosive contributions are quantified, for each term of the parametrix series giving the difference of the densities, in Lemma 3 of [KKM15]. The constant C_η can then, under (\mathbf{A}_H) , be bounded as follows:

$$C_\eta \leq \sum_{r \geq 1} (r+1) \frac{\bar{C}^{r+1} [\Gamma(\frac{\eta}{2})]^r}{\Gamma(1+r\frac{\eta}{2})} T^{r\frac{\eta}{2}},$$

for a constant $\bar{C} := \bar{C}((\mathbf{A}_H), T)$ which does not depend on η .

Similarly, under (\mathbf{A}_{PS}) , the controls of Lemma 3 in [KKM15] give:

$$C_q \leq \sum_{r \geq 1} (r+1) \frac{\bar{C}^{r+1} [\Gamma(\alpha(q))]^r}{\Gamma(1+r\alpha(q))} T^{r\alpha(q)}, \quad \bar{C} := \bar{C}((\mathbf{A}_{PS}), T).$$

Introducing for $\theta > 0$ the quantity:

$$I(T, \theta) := \sum_{r \geq 1} (r+1) \frac{\bar{C}^{r+1} [\Gamma(\theta)]^r}{\Gamma(1+r\theta)} T^{r\theta},$$

one actually gets that there exists constants C, \tilde{C} , independent of θ s.t.:

$$I(T, \theta) \leq C \exp(\tilde{C}(\theta^{-1} + 1)^{\theta^{-1}+1}). \quad (2.14)$$

Hence, equations (2.12) and (2.13) follow taking $\theta = \frac{\eta}{2}$ and $\theta = \alpha(q)$ respectively.

Let us now prove (2.14). One easily gets that for a given T there exists \tilde{C} independent of θ as well such that:

$$I(T, \theta) \leq \sum_{r \geq 1} (r+1) \frac{\tilde{C}^{r+1} [\Gamma(\theta)]^r}{\Gamma(1+r\frac{\theta}{2})}.$$

Set now $r_0 := \lceil \frac{1}{\theta} \rceil$ and write by monotonicity of the Γ function (see e.g. formula 8.363 (8) in Gradstein and Ryzhik [GR14]):

$$\begin{aligned} I(T, \theta) &\leq C \sum_{k \geq 0} (k+1)r_0 \sum_{kr_0 \leq r < (k+1)r_0} \frac{\{\tilde{C}\Gamma(\theta)\}^r}{\Gamma(1+k)} \\ &\leq C \sum_{k \geq 0} \frac{(k+1)r_0}{\Gamma(k+1)} \sum_{kr_0 \leq r < (k+1)r_0} \{\tilde{C}(\theta^{-1} + \exp(-1))\}^r \\ &\leq C r_0^2 \sum_{k \geq 0} \frac{(k+1)}{\Gamma(k+1)} [\tilde{C}(\theta^{-1} + \exp(-1))^{r_0}]^{k+1} \leq C \exp(\tilde{C}(\theta^{-1} + 1)^{\theta^{-1}+1}). \end{aligned}$$

This gives (2.14) and completes the proof. \square

3. ERROR ANALYSIS AND DERIVATION OF THE MAIN RESULTS

3.1. Error Decomposition and Parametrix Representation of Densities.

This Section is devoted to the proof of Theorems 1 and 2. Let $0 \leq t_i < t_j \leq T$ and $(x, y) \in (\mathbb{R}^d)^2$ be given. One writes for $\varepsilon > 0$:

$$|p(t_i, t_j, x, y) - p^h(t_i, t_j, x, y)| \leq |p - p_\varepsilon|(t_i, t_j, x, y) + |p_\varepsilon - p_\varepsilon^h|(t_i, t_j, x, y) + |p_\varepsilon^h - p^h|(t_i, t_j, x, y). \quad (3.1)$$

Now one derives from Lemma 1 that under (\mathbf{A}_H) , for all $\eta \in (0, \gamma)$:

$$|p(t_i, t_j, x, y) - p^h(t_i, t_j, x, y)| \leq C_\eta \varepsilon^{\gamma-\eta} p_c(t_j - t_i, y - x) + |(p_\varepsilon - p_\varepsilon^h)|(t_i, t_j, x, y). \quad (3.2)$$

Similarly, Lemma 2 yields that under (\mathbf{A}_{PS}) , for all $q > d$:

$$|p(t_i, t_j, x, y) - p^h(t_i, t_j, x, y)| \leq C_q \varepsilon^{1/q} p_c(t_j - t_i, y - x) + |(p_\varepsilon - p_\varepsilon^h)|(t_i, t_j, x, y). \quad (3.3)$$

To investigate and minimize the contributions in the error it thus remains from equations (3.2) and (3.3) to precisely quantify the difference $p_\varepsilon - p_\varepsilon^h$ in (3.1). Let us now recall that since the densities $p_\varepsilon, p_\varepsilon^h$ are now respectively associated with a diffusion process and its Euler scheme with *smooth* coefficients, they can be compared thanks to the results in [KM02] adapted to the current inhomogeneous setting. The only delicate, but crucial, point is that we must here specify the dependence on the derivatives of the coefficients, which explode when ε goes to 0. A key ingredient, to proceed is the *parametrix* series representation for the densities of the diffusion and its Euler scheme.

From Section 2 in [KKM15], we derive that under (\mathbf{A}) (i.e. the expansions below hold under both (\mathbf{A}_H) and (\mathbf{A}_{PS})), for all $\varepsilon \geq 0$ (the expansion below even holds for the initial coefficients taking $\varepsilon = 0$), $0 \leq s < t \leq T$, $(x, y) \in (\mathbb{R}^d)^2$:

$$p_\varepsilon(s, t, x, y) := \sum_{r \in \mathbb{N}} \tilde{p}_\varepsilon \otimes H_\varepsilon^{(r)}(s, t, x, y), \quad (3.4)$$

where for $0 \leq u < t \leq T$, $(z, y) \in (\mathbb{R}^d)^2$:

$$H_\varepsilon(u, t, z, y) := (L_u^\varepsilon - \tilde{L}_u^{\varepsilon, y}) \tilde{p}_\varepsilon(u, t, z, y), \quad (3.5)$$

and $L_u^\varepsilon, \tilde{L}_u^{\varepsilon, y}$ respectively stand for the generators at time u of the processes

$$X_t^{(\varepsilon)} = z + \int_u^t b_\varepsilon(v, X_v^{(\varepsilon)}) dv + \int_u^t \sigma_\varepsilon(v, X_v^{(\varepsilon)}) dW_v, \quad \tilde{X}_t^{(\varepsilon), y} = z + \int_u^t \sigma_\varepsilon(v, y) dW_v, \quad (3.6)$$

i.e. for all $\varphi \in C^2(\mathbb{R}^d, \mathbb{R})$, $x \in \mathbb{R}^d$,

$$L_u^\varepsilon \varphi(x) = \langle b_\varepsilon(u, x), \nabla_x \varphi(x) \rangle + \frac{1}{2} \text{Tr} \left(\sigma_\varepsilon \sigma_\varepsilon^*(u, x) D_x^2 \varphi(x) \right), \quad \tilde{L}_u^{\varepsilon, y} \varphi(x) = \frac{1}{2} \text{Tr} \left(\sigma_\varepsilon \sigma_\varepsilon^*(u, y) D_x^2 \varphi(x) \right).$$

Also $\tilde{p}_\varepsilon(u, t, z, y) := \tilde{p}_\varepsilon^y(u, t, z, w)|_{w=y}$ where $\tilde{p}_\varepsilon^y(u, t, z, \cdot)$ stands for the density at time t of the process $\tilde{X}^{(\varepsilon), y}$ starting from z at time u . We denote in (3.4), $\tilde{p}_\varepsilon \otimes H_\varepsilon^{(0)}(s, t, x, y) = \tilde{p}_\varepsilon(s, t, x, y)$ and for all $r \geq 1$, $\tilde{p}_\varepsilon \otimes H_\varepsilon^{(r)}(s, t, x, y) = \int_s^t du \int_{\mathbb{R}^d} \tilde{p}_\varepsilon(s, u, x, z) H_\varepsilon^{(r)}(u, t, z, y) dz$ where for $r \geq 2$, $H_\varepsilon^{(r)}(u, t, z, y) := H_\varepsilon \otimes H_\varepsilon^{(r-1)}(u, t, z, y) := \int_u^t dv \int_{\mathbb{R}^d} H_\varepsilon(u, v, z, w) H_\varepsilon^{(r-1)}(v, t, w, y) dw$. More generally, the symbol \otimes stands for the time-space convolution, i.e. for two real valued functions f, g defined on $[0, T]^2 \times (\mathbb{R}^d)^2$, $0 \leq s < t \leq T$, $f \otimes g(s, t, x, y) := \int_s^t du \int_{\mathbb{R}^d} f(s, u, x, z) g(u, t, z, y) dz$. We also recall that under (\mathbf{A}_{PS}) , since the diffusion coefficient is smooth we do not regularize it and denote in this case $\sigma_\varepsilon = \sigma$.

To investigate the contribution $p_\varepsilon - p_\varepsilon^h$ in (3.1) we will also use for $0 \leq i < j \leq N$, $(x, y) \in (\mathbb{R}^d)^2$ the function:

$$p_\varepsilon^d(t_i, t_j, x, y) := \sum_{r \in \mathbb{N}} \tilde{p}_\varepsilon \otimes_h H_\varepsilon^{(r)}(t_i, t_j, x, y), \quad (3.7)$$

where the quantities at hand are the same as above and the symbol \otimes_h replacing the \otimes in (3.4) denotes the *discrete* convolution, i.e. for all $r \geq 1$,

$$\tilde{p}_\varepsilon \otimes_h H_\varepsilon^{(r)}(t_i, t_j, x, y) = h \sum_{k=0}^{j-i-1} \int_{\mathbb{R}^d} \tilde{p}_\varepsilon(t_i, t_{i+k}, x, z) H_\varepsilon^{(r)}(t_{i+k}, t_j, z, y) dz.$$

Even though $p_\varepsilon^d(t_i, t_j, x, \cdot)$ is not *a priori* a density, we will call it so with a slight abuse of terminology. An important control, under **(A)**, for the terms in the parametrix series is the following:

$$|\tilde{p}_\varepsilon \otimes_h H_\varepsilon^{(r)}(s, t, x, y)| + |\tilde{p}_\varepsilon \otimes H_\varepsilon^{(r)}(s, t, x, y)| \leq \frac{((1 \vee T^{(1-\gamma)/2})c_1)^{r+1} [\Gamma(\frac{\gamma}{2})]^r}{\Gamma(1+r\frac{\gamma}{2})} p_c(t-s, y-x)(t-s)^{\frac{r\gamma}{2}}, \quad (3.8)$$

taking $\gamma = 1$ under **(A_{PS})**. We emphasize that those bounds are uniform w.r.t. $\varepsilon \geq 0$ and refer to [KM02] or Section 2 in [KKM15] for a proof.

From the same references (see also Lemma 3.6 in [KM00]), we have that the density of the Euler scheme also admits a similar *parametrix* representation. Introducing for $0 \leq i < k \leq N$, $(z, y) \in (\mathbb{R}^d)^2$, the schemes:

$$\begin{aligned} X_{t_k}^{h,(\varepsilon)} &= z + \sum_{l=i}^{k-1} (b_\varepsilon(t_l, X_{t_l}^{h,(\varepsilon)})h + \sigma_\varepsilon(t_l, X_{t_l}^{h,(\varepsilon)})(W_{t_{l+1}} - W_{t_l})), \\ \tilde{X}_{t_k}^{h,(\varepsilon),y} &= z + \sum_{l=i}^{k-1} \sigma_\varepsilon(t_l, y)(W_{t_{l+1}} - W_{t_l}), \end{aligned} \quad (3.9)$$

viewed as Markov Chains, their generators write for all $\varphi \in C^2(\mathbb{R}^d, \mathbb{R})$, $x \in \mathbb{R}^d$:

$$L_{t_i}^{h,\varepsilon} \varphi(x) := h^{-1} \mathbb{E}[\varphi(X_{t_{i+1}}^{h,(\varepsilon),t_i,x}) - \varphi(x)], \quad \tilde{L}_{t_i}^{h,\varepsilon,y} \varphi(x) = h^{-1} \mathbb{E}[\varphi(\tilde{X}_{t_{i+1}}^{h,(\varepsilon),y,t_i,x}) - \varphi(x)].$$

Define now for $0 \leq k < j \leq N$, $(z, y) \in (\mathbb{R}^d)^2$ the *Markov chain* analogue of the parametrix kernel H in (3.5) by:

$$H_\varepsilon^h(t_i, t_j, z, y) := (L_{t_i}^{h,\varepsilon} - \tilde{L}_{t_i}^{h,\varepsilon,y}) \tilde{p}_\varepsilon^h(t_i + h, t_j, x, y).$$

One gets the following parametrix representation for the density of the Euler scheme:

$$p_\varepsilon^h(t_i, t_j, x, y) := \sum_{r=0}^{j-i} \tilde{p}_\varepsilon \otimes_h H_\varepsilon^{h,(r)}(t_i, t_j, x, y). \quad (3.10)$$

Again, the subscript ε is meant to explicitly express the dependence on the mollified coefficients. Also, the terms in the above series satisfy the controls of equation (3.8) uniformly in $\varepsilon \geq 0$.

3.2. Error Expansion for The Euler Scheme: Controls on the Densities.

From Theorem 1.1, Theorem 2.1 and their proofs in [KM02] we have with the notations of the previous paragraph:

$$(p_\varepsilon - p_\varepsilon^h)(t_i, t_j, x, y) = (p_\varepsilon - p_\varepsilon^d)(t_i, t_j, x, y) + h \int_0^1 (1-\tau) \left\{ p_\varepsilon^d \otimes_h (\tilde{L}_{\cdot,*}^\varepsilon - \tilde{L}_{\cdot,*}^{*,\varepsilon})^2 p_\varepsilon^{\tau,h}(t_i, t_j, x, y) \right\} d\tau, \quad (3.11)$$

where we denote for $0 \leq t_i < t_j \leq T, \tau \in [0, 1]$:

$$\begin{aligned} p_\varepsilon^{\tau, h}(t_i, t_j, x, y) &:= \sum_{r=0}^{j-i} \tilde{p}_\varepsilon^\tau \otimes_h H_\varepsilon^{h, (r)}(t_i, t_j, x, y), \\ \forall (k, z) \in (i, j] \times \mathbb{R}^d, \tilde{p}_\varepsilon^\tau(t_i, t_k, x, z) &:= \int_{\mathbb{R}^d} \tilde{p}_\varepsilon^x(t_i, t_i + \tau h, x, w) \tilde{p}_\varepsilon^z(t_i + \tau h, t_k, w, z) dw. \end{aligned}$$

Also, for $k \in \{1, 2\}, t \in (t_i, t_j), (\tilde{L}_{t, *})^k \phi(x, y) := (L_{t, \xi}^\varepsilon)^k \phi(x, y)|_{\xi=x}, (\tilde{L}_t^{*, \varepsilon})^k \phi(x, y) := (\bar{L}_{t, \xi}^\varepsilon)^k \phi(x, y)|_{\xi=y}$ for

$$\begin{aligned} L_{t, \xi}^\varepsilon \phi(x, y) &= \langle b_\varepsilon(t, \xi), D_x \phi(x, y) \rangle + \frac{1}{2} \text{Tr}(a_\varepsilon(t, \xi) D_x^2 \phi(x, y)), \\ \bar{L}_{t, \xi}^\varepsilon \phi(x, y) &= \frac{1}{2} \text{Tr}(a_\varepsilon(t, \xi) D_x^2 \phi(x, y)). \end{aligned}$$

Observe that $L_t^\varepsilon \phi(x, y) = L_{t, *}^\varepsilon \phi(x, y)$, but more generally the operators do not coincide anymore when iterated. Also, we indicate that the operators involved slightly differ from [KM02] since we chose to use a Gaussian process without drift as proxy, see (3.6) and (3.9). Another difference is the fact that we deal with inhomogeneous coefficients, and the notations $\tilde{L}_{t, *}, \tilde{L}_t^{*, \varepsilon}$ in (3.11) are used to emphasize the time dependence of the operators in the discrete convolution \otimes_h . Anyhow, reproducing the proof of [KM02] taking into account the indicated differences leads to the expression in (3.11).

We mention carefully that in order to analyze the contribution of the last two terms in the r.h.s. of (3.11) no smoothness in time of the coefficients is needed. On the other hand, such smoothness is clearly required to derive some convergence rates, since to control $p_\varepsilon - p_\varepsilon^d$ we need to investigate the difference between time integrals and Riemann sums (see Proposition 5 and its proof below).

The term $\int_0^1 (1 - \tau) \{p_\varepsilon^d \otimes_h (\tilde{L}_{t, *}^\varepsilon - \tilde{L}_t^{*, \varepsilon})^2 p_\varepsilon^{\tau, h}(t_i, t_j, x, y)\} d\tau$ involves derivatives of the coefficients and *heat kernels* up to order 4. The point is again that the derivatives of the coefficients and kernels explode with ε going to 0 (see equation (2.6)). It is precisely this aspect that deteriorates the convergence rate w.r.t. the usual *smooth* case. We carefully mention that if $\sigma(t, x) = \sigma$, the previous contributions involve lower derivatives of the heat kernel (up to order 2).

The key elements are now the following Propositions. The first one gives bounds for the derivatives of the densities involved in the parametrix series (3.4), (3.7). The second one controls the difference between the discrete and continuous convolutions in (3.11).

Proposition 4 (Controls for the derivatives of the densities). *Let $\alpha, |\alpha| \leq 4$ be a multi-derivation index. Under (\mathbf{A}_H) , there exist constants C, c s.t. for all $0 \leq s < t \leq T, (x, y) \in (\mathbb{R}^d)^2$:*

$$\begin{aligned} |D_x^\alpha \bar{p}_\varepsilon(s, t, x, y)| &\leq \frac{C}{(t-s)^{|\alpha|/2}} p_c(t-s, y-x), |\alpha| \leq 2, \\ |D_x^\alpha \bar{p}_\varepsilon(s, t, x, y)| &\leq \frac{C}{(t-s)^{|\alpha|/2}} p_c(t-s, y-x) (1 + \varepsilon^{-|\alpha|+2} (t-s)^{\gamma/2}), |\alpha| \in \llbracket 3, 4 \rrbracket, \\ |D_y^\alpha \bar{p}_\varepsilon(s, t, x, y)| &\leq \frac{C \varepsilon^{-|\alpha|+\gamma}}{(t-s)^{|\alpha|/2}} p_c(t-s, y-x). \end{aligned} \tag{3.12}$$

Under (\mathbf{A}_{PS}) , for all $q > d$, $\eta \in (0, \alpha(q))$, $\alpha(q) = \frac{1}{2}(1 - \frac{d}{q})$, there exist constants C, c s.t. for all $0 \leq s < t \leq T$, $(x, y) \in (\mathbb{R}^d)^2$:

$$\begin{aligned} |D_x \bar{p}_\varepsilon(s, t, x, y)| &\leq \frac{C}{(t-s)^{1/2}} p_c(t-s, y-x), \\ |D_x^\alpha \bar{p}_\varepsilon(s, t, x, y)| &\leq \frac{C}{(t-s)^{|\alpha|/2}} p_c(t-s, y-x) (1 + C_{\eta, q} \varepsilon^{-|\alpha|+2-\eta+1/q} \mathbb{1}_{|\alpha| \geq 3} (t-s)^{\eta/2}), \quad |\alpha| \in \llbracket 2, 4 \rrbracket, \\ |D_y \bar{p}_\varepsilon(s, t, x, y)| &\leq \frac{C}{(t-s)^{1/2}} (1 + \varepsilon^{-\eta} C_\eta (t-s)^{\eta/2}) p_c(t-s, y-x), \\ |D_y^\alpha \bar{p}_\varepsilon(s, t, x, y)| &\leq \frac{C(1 + C_{\eta, q} \varepsilon^{-|\alpha|+1-\eta} (t-s)^{\eta/2})}{(t-s)^{|\alpha|/2}} p_c(t-s, y-x), \quad |\alpha| \in \llbracket 2, 4 \rrbracket, \end{aligned} \quad (3.13)$$

where $C_{\eta, q} = C_\eta \times C_q$ and C_η, C_q are as in Proposition 3.

In the above expressions \bar{p}_ε can be any of the densities $p_\varepsilon, p_\varepsilon^d, p_\varepsilon^{\tau, h}$ uniformly in $\tau \in [0, 1]$. For $p_\varepsilon^d, p_\varepsilon^{\tau, h}$, the time variables s, t are taken on the time grid.

Remark 6. We point out that, under (\mathbf{A}_H) , the previous controls for $\bar{p}_\varepsilon = p_\varepsilon$ improve in some sense those of [IKO62], since we do not exploit any smoothness in time of the coefficients and we get the same pointwise controls for the derivatives of the non degenerate heat-kernel with Hölder coefficients in space up to order 2, uniformly in $\varepsilon \in [0, 1]$.

Proposition 5 (Bounds for the difference between continuous and discrete time convolutions). *Under (\mathbf{A}_H) , there exist C, c s.t. for all $0 \leq t_i < t_j \leq T$, $(x, y) \in (\mathbb{R}^d)^2$, $\eta \in (0, \gamma)$:*

$$|(p_\varepsilon - p_\varepsilon^d)(t_i, t_j, x, y)| \leq C \left[\frac{\varepsilon^{-2+\gamma} h}{(t_j - t_i)^{1-\gamma/2}} + \frac{h^{(\gamma-\eta)/2}}{\eta} \right] p_c(t_j - t_i, y - x). \quad (3.14)$$

Under (\mathbf{A}_{PS}) , there exist C, c s.t. for all $0 \leq t_i < t_j \leq T$, $(x, y) \in (\mathbb{R}^d)^2$, $q > d$, $\eta \in (0, \alpha(q))$:

$$|(p_\varepsilon - p_\varepsilon^d)(t_i, t_j, x, y)| \leq C \left\{ C_{\eta, q} \frac{h \varepsilon^{-(1+\eta)}}{(t_j - t_i)^{1-\eta/2}} + \frac{h^{1-\eta/2}}{\eta} \right\} p_c(t_j - t_i, y - x), \quad (3.15)$$

with $C_{\eta, q}, \alpha(q)$ as in Propositions 3.12, 3 respectively. If now $\sigma(t, x) = \sigma$, i.e. constant diffusion term, the previous bound improves to:

$$|(p_\varepsilon - p_\varepsilon^d)(t_i, t_j, x, y)| \leq C \left\{ \frac{1}{\alpha(q)} \frac{h \varepsilon^{-(1+\eta)+1/q}}{(t_j - t_i)^{1/2+\alpha(q)}} + \frac{h^{1-\eta/2}}{\eta} \right\} p_c(t_j - t_i, y - x). \quad (3.16)$$

We postpone the proof of Propositions 4 and 5 to Section 4 for clarity. It now remains to exploit Propositions 4, 5 and (3.11) to specifically control how the weak error for the densities depends on the explosive norms of the mollified coefficients.

3.3. Proof of The Main Results for Hölder Coefficients (Theorem 1 under (\mathbf{A}_H))

Observe from Proposition 4 that, for all $k \in \llbracket 1, j-1 \rrbracket$, $(z, y) \in (\mathbb{R}^d)^2$, $\tau \in [0, 1]$:

$$|(\tilde{L}_{t_k, *}^\varepsilon - \tilde{L}_{t_k}^{*, \varepsilon}) p_\varepsilon^{\tau, h}(t_k, t_j, z, y)| \leq \frac{C}{(t_j - t_k)^{1-\gamma/2}} p_c(t_j - t_k, y - z).$$

We analyze the contribution $[p_\varepsilon^d \otimes_h (\tilde{L}_{\cdot,*}^\varepsilon - \tilde{L}_{\cdot,*}^{*,\varepsilon})^2 p_\varepsilon^{\tau,h}](t_i, t_j, x, y)$ in (3.11) thanks to Proposition 4 as follows:

$$\begin{aligned} & |[p_\varepsilon^d \otimes_h (\tilde{L}_{\cdot,*}^\varepsilon - \tilde{L}_{\cdot,*}^{*,\varepsilon})^2 p_\varepsilon^{\tau,h}](t_i, t_j, x, y)| \leq Ch \\ & \times \left\{ \sum_{k \in \llbracket i+1, \lceil (i+j)/2 \rrbracket \rrbracket} \int_{\mathbb{R}^d} \frac{\varepsilon^{-1+\gamma}}{(t_k - t_i)^{1/2}} p_c(t_k - t_i, z - x) \left[\frac{1 + \varepsilon^{-1}(t_j - t_k)^{\gamma/2}}{(t_j - t_k)^{3/2}} \right] p_c(t_j - t_k, y - z) dz \right. \\ & \left. + \sum_{k \in \llbracket \lceil (i+j)/2 \rceil + 1, j - 1 \rrbracket} \int_{\mathbb{R}^d} \frac{\varepsilon^{-2+\gamma}}{(t_k - t_i)} p_c(t_k - t_i, z - x) \frac{1}{(t_j - t_k)^{1-\gamma/2}} p_c(t_j - t_k, y - z) dz \right\}, \end{aligned} \quad (3.17)$$

where we perform one integration by part w.r.t. z for the first integral and two for the second one (taking once the adjoints). We thus get:

$$|[p_\varepsilon^d \otimes_h (\tilde{L}_{\cdot,*}^\varepsilon - \tilde{L}_{\cdot,*}^{*,\varepsilon})^2 p_\varepsilon^{\tau,h}](t_i, t_j, x, y)| \leq C \frac{\varepsilon^{-1+\gamma}}{(t_j - t_i)} \left\{ 1 + \varepsilon^{-1}(t_j - t_i)^{\gamma/2} \right\} p_c(t_j - t_i, y - x). \quad (3.18)$$

We thus finally derive from (3.11), (3.14), (3.18):

$$\begin{aligned} & |(p_\varepsilon - p_\varepsilon^h)(t_i, t_j, x, y)| \\ & \leq C \left\{ \left(\frac{h\varepsilon^{-2+\gamma}}{(t_j - t_i)^{1-\gamma/2}} + \frac{h^{(\gamma-\eta)/2}}{\eta} \right) + \frac{h\varepsilon^{-1+\gamma}}{t_j - t_i} (1 + \varepsilon^{-1}(t_j - t_i)^{\gamma/2}) \right\} p_c(t_j - t_i, y - x). \end{aligned} \quad (3.19)$$

We can assume for the rest of the proof that $0 \leq t_j - t_i \leq T \leq 1$, which is the only case leading to include the possibly explosive contributions $\left((t_j - t_i)^{-(1-\beta\gamma/2)} \right)_{\beta \in \{0,1\}}$ in (3.19) for the optimization over ε . For $T \geq 1$ those contributions can indeed be bounded by 1.

Equation (3.19) together with (3.1), (3.2) now yields for $\varepsilon \leq (t_j - t_i)^{\gamma/2}$, recalling as well from Proposition 3 that $\eta^{-1} \leq C_\eta$:

$$|(p - p^h)(t_i, t_j, x, y)| \leq C \left\{ C_\eta (\varepsilon^{\gamma-\eta} + h^{(\gamma-\eta)/2}) + \frac{h\varepsilon^{-2+\gamma}}{(t_j - t_i)^{1-\gamma/2}} \right\} p_c(t_j - t_i, y - x), \quad \eta \in (0, \gamma).$$

Taking now $C_\eta \varepsilon^{\gamma-\eta} = \frac{h\varepsilon^{-2+\gamma}}{(t_j - t_i)^{1-\gamma/2}}$ leads to:

$$\varepsilon^{2-\eta} = \frac{h}{(t_j - t_i)^{1-\gamma/2} C_\eta} \iff C_\eta \varepsilon^{\gamma-\eta} = C_\eta^{(2-\gamma)/(2-\eta)} \left\{ \frac{h}{(t_j - t_i)^{1-\gamma/2}} \right\}^{\frac{\gamma-\eta}{2-\eta}}. \quad (3.20)$$

Observe now that, since $(t_j - t_i) \leq T \leq 1$:

$$\varepsilon \leq (t_j - t_i)^{\gamma/2} \iff h \leq (t_j - t_i)^{(\gamma/2+1)}.$$

Hence, the condition $(t_j - t_i) \geq h^{2/(2+\gamma)}$ indeed guarantees that $\varepsilon \leq (t_j - t_i)^{\gamma/2}$ as we assumed above. Thus, for this choice of ε , we derive:

$$\begin{aligned} & |(p - p^h)(t_i, t_j, x, y)| \leq C \left\{ C_\eta h^{(\gamma-\eta)/2} + C_\eta^{\frac{2-\gamma}{2-\eta}} \left(\frac{h}{(t_j - t_i)^{1-\gamma/2}} \right)^{\frac{\gamma-\eta}{2-\eta}} \right\} p_c(t_j - t_i, y - x) \\ & \leq C \exp(C(2\eta^{-1} + 1)^{2\eta^{-1}+1}) \left\{ h^{(\gamma-\eta)/2} + \left(\frac{h}{(t_j - t_i)^{1-\gamma/2}} \right)^{\frac{\gamma-\eta}{2-\eta} \frac{1-\gamma/2}{2-\eta}} \right\} p_c(t_j - t_i, y - x), \end{aligned} \quad (3.21)$$

using the bounds of Proposition 3 for the last inequality. The point is now to carefully choose $\eta := \eta(h)$. Let us consider the specific sequence $\eta = \eta(h) := 2^{\frac{\log_3(h^{-1})}{\log_2(h^{-1})}}$, where we recall that for $k \in \mathbb{N}$, $\log_k(x)$ stands for the k^{th} iterated logarithm of x . Setting $\beta_h := h^{-\eta}$ and $\alpha_h := \exp\left(C(2\eta^{-1} + 1)^{2\eta^{-1}+1}\right)$, we get that:

$$\begin{aligned} \log_2(\beta_h) &= \log(\eta \log(h^{-1})) = \log(2) + \log_4(h^{-1}) - \log_3(h^{-1}) + \log_2(h^{-1}), \\ \log_2(\alpha_h) &= \log(C(2\eta^{-1} + 1)^{2\eta^{-1}+1}) = \log(C) + (2\eta^{-1} + 1) \log(2\eta^{-1} + 1) \\ &= \log(C) + \left(\frac{\log_2(h^{-1})}{\log_3(h^{-1})} + 1\right) \log(2\eta^{-1}(1 + \frac{\eta}{2})) \\ &= \log(C) + \left(\frac{\log_2(h^{-1})}{\log_3(h^{-1})} + 1\right) \{\log(2\eta^{-1}) + \log(1 + \frac{\eta}{2})\} \\ &= \log(C) + \left(\frac{\log_2(h^{-1})}{\log_3(h^{-1})} + 1\right) \{\log_3(h^{-1}) - \log_4(h^{-1}) + \log(1 + \frac{\eta}{2})\} \\ &= \log_2(h^{-1}) - \frac{\log_2(h^{-1}) \log_4(h^{-1})}{\log_3(h^{-1})} + \log_3(h^{-1}) - \log_4(h^{-1}) + R_h, \\ R_h &:= \log(C) + \log(1 + \frac{\log_3(h^{-1})}{\log_2(h^{-1})}) \left\{ \frac{\log_2(h^{-1})}{\log_3(h^{-1})} + 1 \right\}. \end{aligned}$$

It is easily seen that there exists a finite constant $\bar{C} > 0$ s.t. for all h small enough, $R_h \leq \bar{C}$ and that $\log_2(\beta_h) \geq \log_2(\alpha_h) - \bar{C}$. By monotonicity of the exponential, recalling as well that $\eta \in (0, \gamma)$, we thus derive:

$$\left(\beta_h^{\frac{1}{2}} + \beta_h^{\frac{1-\gamma/2}{2-\eta}}\right) \alpha_h = \left(h^{-\frac{\eta}{2}} + h^{-\eta \frac{1-\gamma/2}{2-\eta}}\right) \exp(C(2\eta^{-1} + 1)^{2\eta^{-1}+1}) \leq 2h^{-\eta(1/2 + \exp(\bar{C}))}. \quad (3.22)$$

Plugging (3.22) into (3.21) we complete the proof of equation (1.6) in Theorem 1.

To prove the other statements concerning the weak error let us first point out that we do not need to regularize the coefficients as above. Let $t_j = jh \in [0, T]$ be fixed. Observe that setting for all $(t, x) \in [0, t_j] \times \mathbb{R}^d$, $v(t, x) := \mathbb{E}[f(X_{t_j}^{t,x})]$, as soon as f satisfies the integrability condition (1.7) and from Proposition 4, which is valid for $\varepsilon = 0$ (no regularization), we have that for all $(t, x) \in [0, t_j] \times \mathbb{R}^d$:

$$|\nabla_x v(t, x)| \leq \frac{C}{(t_j - t)^{1/2}}. \quad (3.23)$$

On the other hand, we get that for $f \in C^\beta(\mathbb{R}^d, \mathbb{R})$, $\beta \in (0, 1]$, we have for a multi-index α , $|\alpha| = 2$, $(t, x) \in [0, t_j] \times \mathbb{R}^d$:

$$\partial_x^\alpha v(t, x) = \int_{\mathbb{R}^d} \partial_x^\alpha p(t, t_j, x, y) f(y) dy = \int_{\mathbb{R}^d} \partial_x^\alpha p(t, t_j, x, y) (f(y) - f(x)) dy,$$

recalling that $\partial_x^\alpha \int_{\mathbb{R}^d} p(t, t_j, x, y) dy = 0$ for the last identity. Then, under (\mathbf{A}_H) we derive from the above equation, (3.12) (for $\varepsilon = 0$) and the Hölder continuity of f , that there exists $C \geq 1$ s.t. for all $(t, x) \in [0, t_j] \times \mathbb{R}^d$:

$$|\partial_x^\alpha v(t, x)| \leq \frac{C}{(t_j - t)^{1-\beta/2}}, \quad (3.24)$$

yielding an integrable singularity in time.

We can now also suppose, without loss of generality that $\gamma \geq \beta$. Indeed, if $f \in C^\beta(\mathbb{R}^d, \mathbb{R})$ for $\beta > \gamma$, it is locally $C^\gamma(\mathbb{R}^d, \mathbb{R})$. Observe in that case that we have $h^{\gamma/\beta} \leq h \leq 1$. Plugging (3.24) and (3.23) in an expansion

similar to (1.11), recalling from Proposition 4 that $v \in C^{\beta/2, \beta}([0, t_j] \times \mathbb{R}^d, \mathbb{R}) \cap C^{1,2}([0, t_j] \times \mathbb{R}^d, \mathbb{R})$, gives:

$$\begin{aligned}
& |\mathbb{E}[f(X_{t_j}^{h, t_i, x}) - f(X_{t_j}^{t_i, x})]| \leq |\mathbb{E}[f(X_{t_j}^{h, t_i, x}) - v(t_{j-1}, X_{t_{j-1}}^{h, t_i, x})]| \\
& + C \sum_{k=i}^{j-2} \int_{t_k}^{t_{k+1}} ds \mathbb{E}[\{|\nabla_x v(s, X_s^{h, t_i, x})| + |D_x^2 v(s, X_s^{h, t_i, x})|\} \{|s - t_k|^{\gamma/2} + |X_s^{h, t_i, x} - X_{t_k}^{h, t_i, x}|^\gamma\}] \\
& \leq \mathbb{E}[|f(X_{t_j}^{h, t_i, x}) - f(X_{t_{j-h\gamma/\beta}}^{h, t_i, x})|] + \mathbb{E}[|v(t_j, X_{t_{j-h\gamma/\beta}}^{h, t_i, x}) - v(t_j - h\gamma/\beta, X_{t_{j-h\gamma/\beta}}^{h, t_i, x})|] \\
& + |\mathbb{E}[v(t_j - h\gamma/\beta, X_{t_{j-h\gamma/\beta}}^{h, t_i, x}) - v(t_{j-1}, X_{t_{j-1}}^{h, t_i, x})]| + Ch^{\gamma/2} \sum_{k=i}^{j-2} \int_{t_k}^{t_{k+1}} \left(1 + \frac{1}{(t_j - s)^{1-\beta/2}}\right) ds \\
& \leq \mathbb{E}[|f(X_{t_j}^{h, t_i, x}) - f(X_{t_{j-h\gamma/\beta}}^{h, t_i, x})|] + \mathbb{E}[|v(t_j, X_{t_{j-h\gamma/\beta}}^{h, t_i, x}) - v(t_j - h\gamma/\beta, X_{t_{j-h\gamma/\beta}}^{h, t_i, x})|] \\
& + \int_{t_{j-1}}^{t_j - h\gamma/\beta} \mathbb{E}[|\nabla_x v(s, X_s^{h, t_i, x})| |b(s, X_s^{h, t_i, x}) - b(t_{j-1}, X_{t_{j-1}}^{h, t_i, x})| + |D_x^2 v(s, X_s^{h, t_i, x})| |a(s, X_s^{h, t_i, x}) - a(t_{j-1}, X_{t_{j-1}}^{h, t_i, x})|] ds \\
& \quad + Ch^{\gamma/2} \sum_{k=i}^{j-2} \int_{t_k}^{t_{k+1}} \left(1 + \frac{1}{(t_j - s)^{1-\beta/2}}\right) ds \\
& \leq C\mathbb{E}[|X_{t_j}^{h, t_i, x} - X_{t_{j-h\gamma/\beta}}^{h, t_i, x}|^\beta] + Ch^{\gamma/2} \left\{1 + \int_{t_i}^{t_j - h\gamma/\beta} \left(1 + \frac{1}{(t_j - s)^{1-\beta/2}}\right) ds\right\} \leq Ch^{\gamma/2},
\end{aligned} \tag{3.25}$$

expanding as in (1.11) the term $|\mathbb{E}[v(t_j - h\gamma/\beta, X_{t_{j-h\gamma/\beta}}^{h, t_i, x}) - v(t_{j-1}, X_{t_{j-1}}^{h, t_i, x})]|$ with Itô's formula and using (1.10) for the last two inequalities. This gives the required control in (1.8).

Let us now prove (1.9). We write for $0 \leq i < j \leq N$, $x \in \mathbb{R}^d$:

$$\begin{aligned}
& \mathbb{E}[\mathbb{I}_{X_{t_j}^{h, t_i, x} \in A}] - \mathbb{E}[\mathbb{I}_{X_{t_j}^{t_i, x} \in A}] \\
& = \{\mathbb{E}[\mathbb{I}_{X_{t_j}^{h, t_i, x} \in A}] - \mathbb{E}[f_\delta(X_{t_j}^{h, t_i, x})]\} + \{\mathbb{E}[f_\delta(X_{t_j}^{h, t_i, x})] - \mathbb{E}[f_\delta(X_{t_j}^{t_i, x})]\} + \{\mathbb{E}[f_\delta(X_{t_j}^{t_i, x})] - \mathbb{E}[\mathbb{I}_{X_{t_j}^{t_i, x} \in A}]\} =: \sum_{i=1}^3 T_i^\delta,
\end{aligned} \tag{3.26}$$

where $\delta > 0$ is such that A satisfies the interior ball condition with radius 2δ , f_δ stands for a smooth approximation of the mapping $x \mapsto \mathbb{I}_{x \in A}$ which is equal to 1 on $A_\delta := \{y \in A : d(y, \partial A) \geq \delta\}$ where $d(\cdot, \partial A)$ stands for the distance to the boundary and to 0 outside of $A \cup V_\delta(A)$, $V_\delta(A) := \{y \in \mathbb{R}^d : d(y, \partial A) \leq \delta\}$ (neighborhood of size δ of the boundary). In particular $|\nabla f_\delta|_\infty = \sup_{x \in V_\delta(A)} |\nabla f_\delta(x)| \leq C\delta^{-1}$.

The terms T_1^δ and T_3^δ can be handled similarly thanks to the Gaussian upper bound that is satisfied, under (\mathbf{A}_H) , by the density of both the diffusion and its Euler scheme, see Proposition 4 or again [She91], [LM10]. Precisely, with the notations of (3.26) and provided that $\delta \leq (t_j - t_i)^{1/2}$:

$$|T_1^\delta + T_3^\delta| \leq \mathbb{E}[\mathbb{I}_{X_{t_j}^{t_i, x} \in V_\delta(A)}] + \mathbb{E}[\mathbb{I}_{X_{t_j}^{h, t_i, x} \in V_\delta(A)}] \leq \frac{C\delta}{(t_j - t_i)^{1/2}} \exp\left(-c \frac{d(x, \partial A)^2}{t_j - t_i}\right). \tag{3.27}$$

Indeed, we have that locally, up to a change of coordinate, only one variable is orthogonal to the hypersurface ∂A . We can thus integrate the Gaussian bounds w.r.t. the other ones. This yields the above control.

Observe that to find the indicated convergence rate this imposes $\delta \leq (t_j - t_i)^{1/2} h^{\gamma/2}$ which specifies the admissible magnitude for the parameter δ . On the other hand, to analyze T_2^δ we have that setting for all $(t, x) \in [0, t_j] \times \mathbb{R}^d$, $v_\delta(t, x) := \mathbb{E}[f_\delta(X_{t_j}^{t, x})]$ the terminal function f_δ is η -Hölder continuous, for all $\eta \in (0, \gamma]$,

with Hölder modulus of continuity bounded by $\delta^{-\eta}$ on $V_{2\delta}(A)$. We will now establish, similarly to (3.24), that for all multi-index α , $|\alpha| \leq 2$, $(t, x) \in [0, t_j] \times \mathbb{R}^d$:

$$|D_x^\alpha v_\delta(t, x)| \leq \frac{C}{(\delta \vee d(x, \partial A))^\eta} \frac{1}{(t_j - t)^{1-\eta/2}}. \quad (3.28)$$

Recall indeed that

$$|D_x^\alpha v_\delta(t, x)| = \left| \int_{\mathbb{R}^d} D_x^\alpha p(t, t_j, x, y) (f_\delta(y) - f_\delta(x)) dy \right| \leq \frac{C}{(t_j - t)} \int_{\mathbb{R}^d} p_c(t_j - t, y - x) |f_\delta(y) - f_\delta(x)| dy, \quad (3.29)$$

exploiting Proposition 4 for the last inequality. Thus

- if both $x, y \notin V_\delta(A)$, then $f_\delta(x) = \mathbb{I}_{x \in A}$, $f_\delta(y) = \mathbb{I}_{y \in A}$. If $x \in (A \cup V_\delta(A))^C$, $y \in A \setminus V_\delta(A)$, or by symmetry $y \in (A \cup V_\delta(A))^C$, $x \in A \setminus V_\delta(A)$, then $|x - y| \geq \delta \vee d(x, \partial A)$. If now $x, y \in (A \cup V_\delta(A))^C$ or $x, y \in A \setminus V_\delta(A)$ then $f_\delta(x) = f_\delta(y)$ yielding a trivial contribution in (3.29).
- if $x, y \in V_\delta(A)$, then the control of the Hölder modulus gives: $|f_\delta(x) - f_\delta(y)| \leq C\delta^{-\eta}|x - y|^\eta = C(\delta \vee d(x, \partial A))^{-\eta}|x - y|^\eta$.
- if $x \in V_\delta(A)$, $y \notin V_\delta(A)$ (resp. $y \in V_\delta(A)$, $x \notin V_\delta(A)$) we can exploit the Hölder continuity for $y \in V_{2\delta}(A)$ (resp. $x \in V_{2\delta}(A)$) and the fact that $|x - y| \geq \delta \vee d(x, \partial A)$ for $y \notin V_{2\delta}(A)$ (resp. $x \notin V_{2\delta}(A)$).

In all cases, we have established that $|f_\delta(x) - f_\delta(y)| \leq C(\delta \vee d(x, \partial A))^{-\eta}|x - y|^\eta$, which plugged into (3.29) yields the control (3.28). Recall now that, again from Proposition 4, we have $v_\delta \in C^{\eta/2, \eta}([0, t_j] \times \mathbb{R}^d) \cap C^{1,2}([0, t_j] \times \mathbb{R}^d)$. In particular, v_δ has the same Hölder continuity modulus as f_δ . We can as well assume w.l.o.g. that $\gamma/\eta \geq 1$ so that $h^{\gamma/\eta} \leq h \leq 1$.

Exploiting now (3.28) in an expansion similar to (1.11) and (3.25), we get:

$$\begin{aligned} |T_2^\delta| &\leq |\mathbb{E}[f_\delta(X_{t_j}^{h, t_i, x}) - f_\delta(X_{t_j - h^{\gamma/\eta}}^{h, t_i, x})]| + |\mathbb{E}[v_\delta(t_j, X_{t_j - h^{\gamma/\eta}}^{h, t_i, x}) - v_\delta(t_j - h^{\gamma/\eta}, X_{t_j - h^{\gamma/\eta}}^{h, t_i, x})]| \\ &\quad + |\mathbb{E}[v_\delta(t_j - h^{\gamma/\eta}, X_{t_j - h^{\gamma/\eta}}^{h, t_i, x}) - v_\delta(t_{j-1}, X_{t_{j-1}}^{h, t_i, x})]| \\ &\quad + C \sum_{k=i}^{j-2} \int_{t_k}^{t_{k+1}} ds \mathbb{E}[\{|\nabla_x v_\delta(s, X_s^{h, t_i, x})| + |D_x^2 v_\delta(s, X_s^{h, t_i, x})|\} \{ |s - t_k|^{\gamma/2} + |X_s^{h, t_i, x} - X_{t_k}^{h, t_i, x}|^\gamma \}] \\ &\leq Ch^{\gamma/2} \left\{ \mathbb{E}[(\delta \vee d(X_{t_j}^{h, t_i, x}, \partial A))^{-2\eta}]^{1/2} + 1 + \int_{t_i}^{t_j - h^{\gamma/\eta}} \left(1 + \frac{1}{(t_j - s)^{1-\eta/2}} \mathbb{E}[(\delta \vee d(X_s^{h, t_i, x}, \partial A))^{-2\eta}]^{1/2} \right) ds \right\}, \end{aligned} \quad (3.30)$$

where the term $|\mathbb{E}[v_\delta(t_j - h^{\gamma/\eta}, X_{t_j - h^{\gamma/\eta}}^{h, t_i, x}) - v_\delta(t_{j-1}, X_{t_{j-1}}^{h, t_i, x})]|$ is again expanded with Itô's formula which yields bounds similar to those appearing for the contributions associated with the indexes $k \in \llbracket i, j-2 \rrbracket$.

Recalling as well that the Euler scheme satisfies the Aronson Gaussian bounds (see Proposition 4 and [LM10] for details) we obtain:

$$\mathbb{E}[(\delta \vee d(X_s^{h, t_i, x}, \partial A))^{-2\eta}] \leq C \left\{ (\delta \vee d(x, \partial A))^{-2\eta} + \int_{\frac{1}{2}d(x, \partial A) \geq d(y, \partial A)} \frac{\exp(-c \frac{|x-y|^2}{s-t_i})}{(\delta \vee d(y, \partial A))^{2\eta}} \frac{dy}{(s-t_i)^{d/2}} \right\}.$$

Since on $\{\frac{1}{2}d(x, \partial A) \geq d(y, \partial A)\}$ we have $|x - y| \geq |x - \pi_{\partial A}(y)| - |\pi_{\partial A}(y) - y| \geq \frac{d(x, \partial A)}{2} \geq d(y, \partial A)$, where $\pi_{\partial A}(y)$ denotes the projection of y on the boundary ∂A , we get:

$$\begin{aligned} \mathbb{E}[\{\delta \vee d(X_s^{h, t_i, x}, \partial A)\}^{-2\eta}] &\leq C \left\{ (\delta \vee d(x, \partial A))^{-2\eta} + \int_{\frac{1}{2}d(x, \partial A) \geq d(y, \partial A)} \frac{\exp(-c \frac{d(y, \partial A)^2}{s-t_i})}{(\delta \vee d(y, \partial A))^{2\eta}} \frac{dy}{(s-t_i)^{d/2}} \right\} \\ &\leq C \left\{ (\delta \vee d(x, \partial A))^{-2\eta} + 1 \right\}. \end{aligned}$$

Hence, since $d(x, \partial A) \geq (t_j - t_i)^{1/2} h^{\gamma/2} \geq \delta$, we get from (3.30)

$$|T_2^\delta| \leq Ch^{\gamma/2} \left\{ \frac{1}{\eta d(x, \partial A)^\eta} + 1 \right\},$$

which together with (3.27) gives (1.9). We conclude emphasizing that the condition $\eta d(x, \partial A)^\eta \geq h^{C\psi(h)}$, indeed guarantees that $\frac{h^{\gamma/2}}{\eta d(x, \partial A)^\eta} \leq h^{\gamma/2 - C\psi(h)}$. This means that the bound in (1.9) is then better than the more general one appearing in (1.6).

3.4. Proof of The Main Results for piecewise smooth coefficients (Theorem 2 under (\mathbf{A}_{PS}))

The idea is to proceed as in the previous section from equation (3.11) and (3.17). To emphasize the specificity of Assumptions (\mathbf{A}_{PS}) , due to the approximation of the piecewise smooth drift, we begin with the special case $\sigma(t, x) = \sigma$. In that framework, a global integration by part, associated with the controls of (2.9) and Proposition 4, yields:

$$\begin{aligned} &| [p_\varepsilon^d \otimes_h (\tilde{L}_{\cdot, *}^\varepsilon - \tilde{L}^{*, \varepsilon})^2 p_\varepsilon^{\tau, h}](t_i, t_j, x, y) | \\ &\leq Ch \sum_{k \in \llbracket i+1, j-1 \rrbracket} \left| \int_{\mathbb{R}^d} \operatorname{div}_z \left(p_\varepsilon^d(t_i, t_k, x, z) b_\varepsilon(t_k, z) \right) \langle b_\varepsilon(t_k, z), \nabla_z p_\varepsilon^{\tau, h}(t_k, t_j, z, y) \rangle dz \right| \\ &\leq Ch \sum_{k \in \llbracket i+1, j-1 \rrbracket} \int_{\mathbb{R}^d} (\varepsilon^{-\eta} + (1 + \varepsilon^{-1} \mathbb{1}_{z \in \mathcal{V}_\varepsilon(\mathcal{I})})) p_c(t_k - t_i, z - x) \times \frac{1}{(t_j - t_k)^{1/2}} p_c(t_j - t_k, y - z) dz. \end{aligned}$$

The point is now to use the Hölder inequality to exploit that the set on which $\nabla_z b_\varepsilon$ gives an explosive bound is *small*. We get:

$$\begin{aligned} &| [p_\varepsilon^d \otimes_h (\tilde{L}_{\cdot, *}^\varepsilon - \tilde{L}^{*, \varepsilon})^2 p_\varepsilon^{\tau, h}](t_i, t_j, x, y) | \\ &\leq Ch \sum_{k \in \llbracket i+1, j-1 \rrbracket} \frac{1}{(t_j - t_k)^{1/2}} \left(\varepsilon^{-\eta} p_c(t_j - t_i, y - x) + \varepsilon^{-1+1/q} \left(\int_{\mathbb{R}^d} p_c(t_k - t_i, z - x)^{\bar{q}} p_c(t_j - t_k, y - z)^{\bar{q}} dz \right)^{1/\bar{q}} \right), \end{aligned}$$

denoting by $\bar{q} > 1$ the conjugate of q , $q^{-1} + \bar{q}^{-1} = 1$. Recall now that:

$$\left(\int_{\mathbb{R}^d} p_c(t_k - t_i, z - x)^{\bar{q}} p_c(t_j - t_k, y - z)^{\bar{q}} dz \right)^{1/\bar{q}} = \left(\frac{c(t_j - t_i)}{(2\pi)(t_k - t_i)(t_j - t_k)} \right)^{d/2(1-1/\bar{q})} \bar{q}^{-d/(2\bar{q})} p_c(t_j - t_i, y - x).$$

This yields:

$$| [p_\varepsilon^d \otimes_h (\tilde{L}_{\cdot, *}^\varepsilon - \tilde{L}^{*, \varepsilon})^2 p_\varepsilon^{\tau, h}](t_i, t_j, x, y) | \leq \frac{C}{\alpha(q)} \frac{\varepsilon^{-1+1/q}}{(t_j - t_i)^{1/2+\alpha(q)}} p_c(t_j - t_i, y - x). \quad (3.31)$$

Performing now in the general case, involving derivatives of the heat kernel up to order 4, an integration by part similar to the one described for (3.17) and using the Hölder inequality as above for the terms involving

derivatives of b_ε , we derive from Proposition 4 that for all $q > d$, $\eta \in (0, \alpha(q))$:

$$|[p_\varepsilon^d \otimes_h (\tilde{L}_{\cdot, \cdot}^\varepsilon - \tilde{L}^{*, \varepsilon})^2 p_\varepsilon^{\tau, h}](t_i, t_j, x, y)| \leq \frac{C}{(t_j - t_i)} \left\{ 1 + C_{\eta, q} \varepsilon^{-(1+\eta)} (t_j - t_i)^{\eta/2} \right\} p_c(t_j - t_i, y - x). \quad (3.32)$$

We thus get, from (3.3), (3.11) and Proposition 5:

$$\begin{aligned} |p - p^h(t_i, t_j, x, y)| &\leq C \left\{ C_q \varepsilon^{1/q} + \left\{ 1 + (t_j - t_i)^{\eta/2} \right\} C_{\eta, q} \frac{h \varepsilon^{-(1+\eta)}}{(t_j - t_i)} \right. \\ &\quad \left. + \left\{ C_{\eta, q} \frac{h \varepsilon^{-(1+\eta)}}{(t_j - t_i)^{1-\eta/2}} + \frac{h^{1-\eta/2}}{\eta} \right\} \right\} p_c(t_j - t_i, y - x), \end{aligned}$$

using equations (3.15), (3.32) in the general case and

$$|p - p^h(t_i, t_j, x, y)| \leq C \left\{ C_q \varepsilon^{1/q} + \frac{1}{\alpha(q)} \frac{h \varepsilon^{-1+1/q}}{(t_j - t_i)^{1/2+\alpha(q)}} + \left\{ \frac{1}{\alpha(q)} \frac{h \varepsilon^{-(1+\eta)+1/q}}{(t_j - t_i)^{1/2+\alpha(q)}} + \frac{h^{1-\eta/2}}{\eta} \right\} \right\} p_c(t_j - t_i, y - x),$$

from (3.16), (3.31), if $\sigma(t, x) = \sigma$. We then set $C_q \varepsilon^{1/q} = C_{\eta, q} \frac{h \varepsilon^{-(1+\eta)}}{(t_j - t_i)^{1-\eta/2}}$ in the general case and $C_q \varepsilon^{1/q} = \frac{1}{\alpha(q)} \frac{h \varepsilon^{-(1+\eta)+1/q}}{(t_j - t_i)^{1/2+\alpha(q)}}$ if $\sigma(t, x) = \sigma$. The results can be derived as in the previous section choosing $\eta := \eta(h) = \psi(h)$, $q := q(h)$ s.t. $\alpha(q) = \psi(h)$.

4. PROOF OF THE TECHNICAL RESULTS FROM SECTION 3.

4.1. Proof of Proposition 4.

4.1.1. Proof under (\mathbf{A}_H) .

Let us establish the result for p_ε . We start from the parametrix representation of p_ε obtained in (3.4). In all cases, we can readily derive from (3.6) (recall that $\tilde{X}^{\varepsilon, y}$ is a non degenerate Gaussian process) and (2.6) that for the *main* term in the expansion:

$$|D_x^\alpha \tilde{p}_\varepsilon(s, t, x, y)| \leq \frac{C}{(t-s)^{|\alpha|/2}} p_c(t-s, y-x), \quad |D_y^\alpha \tilde{p}_\varepsilon(s, t, x, y)| \leq \frac{C \varepsilon^{-|\alpha|+\gamma}}{(t-s)^{|\alpha|/2}} p_c(t-s, y-x). \quad (4.1)$$

Let us now concentrate on the remainder term:

$$R_\varepsilon(s, t, x, y) := \sum_{i \geq 1} \tilde{p}_\varepsilon \otimes H_\varepsilon^{(i)}(s, t, x, y) = \tilde{p}_\varepsilon \otimes \Phi_\varepsilon(s, t, x, y), \quad \Phi_\varepsilon(s, t, x, y) := \sum_{i \geq 1} H_\varepsilon^{(i)}(s, t, x, y).$$

We focus on the first two inequalities in (3.12), the last one can be proved similarly. The ideas are close to those in [IKO62], but we need to adapt them since they considered the “forward” version of the parametrix expansions. The key point is that, for Hölder coefficients we have bounded controls for the derivatives of the remainder in the backward variable up to order two. It is first easily seen for the first derivatives, since the first order derivation gives an integrable singularity in time in the previous expansions. Indeed, from (4.1) and (3.8), one readily gets the statement if $|\alpha| = 1$. The case $|\alpha| \geq 2$ is much more subtle and needs to be discussed

thoroughly. Write indeed:

$$\begin{aligned}
D_x^\alpha R_\varepsilon(s, t, x, y) &= \lim_{\tau \rightarrow 0} \int_{s+\tau}^{(t+s)/2} du \int_{\mathbb{R}^d} D_x^\alpha \tilde{p}_\varepsilon(s, u, x, z) \Phi_\varepsilon(u, t, z, y) dz + \\
&\quad \int_{(t+s)/2}^t du \int_{\mathbb{R}^d} D_x^\alpha \tilde{p}_\varepsilon(s, u, x, z) \Phi_\varepsilon(u, t, z, y) dz \\
&=: \lim_{\tau \rightarrow 0} D_x^\alpha R_\varepsilon^\tau(s, t, x, y) + D_x^\alpha R_\varepsilon^f(s, t, x, y).
\end{aligned} \tag{4.2}$$

The contribution $D_x^\alpha R_\varepsilon^f(s, t, x, y)$ does not exhibit time singularities in the integral, since on the considered integration set $u - s \geq \frac{1}{2}(t - s)$. Let us now recall the usual control on the parametrix kernel under (\mathbf{A}_H) , see e.g. Section 2 in [KKM15]. There exist c, c_1 s.t. for all $0 \leq u < t \leq T, (z, y) \in (\mathbb{R}^d)^2$:

$$|H_\varepsilon(u, t, z, y)| \leq \frac{c_1(1 \vee T^{(1-\gamma)/2})}{(t-u)^{1-\gamma/2}} p_c(t-u, z-y). \tag{4.3}$$

Equation (4.3) for H_ε then yields for all $0 \leq s < t \leq T, (x, y) \in (\mathbb{R}^d)^2$:

$$|H_\varepsilon^{(r)}(s, t, x, y)| \leq ((1 \vee T^{(1-\gamma)/2})c_1)^r \prod_{i=1}^{r-1} B\left(\frac{\gamma}{2}, 1 + (i-1)\frac{\gamma}{2}\right) p_c(t-s, y-x)(t-s)^{-1+\frac{r\gamma}{2}}, \tag{4.4}$$

with the convention $\prod_{i=1}^0 = 1$. We thus derive that for all $0 \leq s < t \leq T, (x, y) \in (\mathbb{R}^d)^2$:

$$|\Phi_\varepsilon(s, t, x, y)| \leq \frac{C}{(t-s)^{1-\gamma/2}} p_c(t-s, y-x) \tag{4.5}$$

$$|\Phi_\varepsilon(u, t, z, y)| \leq \frac{C}{(t-u)^{1-\gamma/2}} p_c(t-u, y-z) \text{ and } |\Phi_\varepsilon(u, t, x, y)| \leq \frac{C}{(t-u)^{1-\gamma/2}} p_c(t-u, y-x).$$

Thus, from equations (4.1) and (4.5):

$$|D_x^\alpha R_\varepsilon^f(s, t, x, y)| \leq \frac{C}{(t-s)^{|\alpha|/2}} p_c(t-s, y-x). \tag{4.6}$$

The delicate contribution is indeed $D_x^\alpha R_\varepsilon^\tau(s, t, x, y)$ for which we need to be more careful. If $|\alpha| = 2$ we exploit some cancellation properties of the derivatives of the Gaussian kernels. Recall now that for an arbitrary $w \in \mathbb{R}^d$, setting for $0 \leq s < u \leq T, \Sigma_\varepsilon(s, u, w) := \int_s^u \sigma_\varepsilon \sigma_\varepsilon^*(v, w) dv$,

$$\begin{aligned}
\tilde{p}_\varepsilon^w(s, u, x, z) &= \frac{1}{(2\pi)^{d/2} \det(\Sigma_\varepsilon(s, u, w))^{1/2}} \exp\left(-\frac{1}{2} \langle \Sigma_\varepsilon(s, u, w)^{-1}(z-x), z-x \rangle\right), \\
D_{x_i x_j}^2 \tilde{p}_\varepsilon^w(s, u, x, z) &= \left\{ (\Sigma_\varepsilon^{-1}(s, u, w)(z-x))_i (\Sigma_\varepsilon^{-1}(s, u, w)(z-x))_j \right. \\
&\quad \left. - \delta_{ij} (\Sigma_\varepsilon^{-1}(s, u, w))_{ii} \right\} \tilde{p}_\varepsilon^w(s, u, x, z), \quad \forall (i, j) \in \llbracket 1, d \rrbracket^2,
\end{aligned} \tag{4.7}$$

where for $q \in \mathbb{R}^d$, we denote for $i \in \llbracket 1, d \rrbracket$ by q_i its i^{th} entry. Hence, for all multi-index $\alpha, |\alpha| = 2$:

$$\int_{\mathbb{R}^d} D_x^\alpha \tilde{p}_\varepsilon^w(s, u, x, z) dz = 0. \tag{4.8}$$

Introducing the centering function $c_\varepsilon^\alpha(s, u, x, z) := (D_x^\alpha \tilde{p}_\varepsilon^w(s, u, x, z))|_{w=x}$, we rewrite:

$$\begin{aligned} D_x^\alpha R_\varepsilon^\tau(s, t, x, y) &= \int_{s+\tau}^{(s+t)/2} du \int_{\mathbb{R}^d} (D_x^\alpha \tilde{p}_\varepsilon - c_\varepsilon^\alpha)(s, u, x, z) \Phi_\varepsilon(u, t, z, y) dz \\ &\quad + \int_{s+\tau}^{(s+t)/2} du \int_{\mathbb{R}^d} c_\varepsilon^\alpha(s, u, x, z) (\Phi_\varepsilon(u, t, z, y) - \Phi_\varepsilon(u, t, x, y)) dz \\ &:= (R_\varepsilon^{\tau,1} + R_\varepsilon^{\tau,2})(s, t, x, y), \end{aligned} \quad (4.9)$$

exploiting the centering condition (4.8) to introduce the last term of the first equality. On the one hand, the terms $D_x^\alpha \tilde{p}_\varepsilon(s, u, x, z)$, $c_\varepsilon^\alpha(s, u, x, z)$ only differ in their frozen coefficients (respectively at point z and x). Exploiting the Hölder property in space of the mollified coefficients, it is then easily seen that:

$$|(D_x^\alpha \tilde{p}_\varepsilon - c_\varepsilon^\alpha)(s, u, x, z)| \leq \frac{C|x-z|^\gamma}{(u-s)} p_c(u-s, z-x) \leq \frac{C}{(u-s)^{1-\gamma/2}} p_c(u-s, z-x),$$

yielding an integrable singularity in time so that, from (4.5):

$$|R_\varepsilon^{\tau,1}(s, t, x, y)| \leq \frac{C}{(t-s)^{1-\gamma}} p_c(t-s, y-x). \quad (4.10)$$

Let us now control the other contribution. The key idea is now to exploit the *smoothing property* of the kernel Φ_ε . Assume indeed that for $A := \{z \in \mathbb{R}^d : |x-z| \leq c(t-s)^{1/2}\}$ (recall as well that $u \in [s, \frac{s+t}{2}]$) one has:

$$|\Phi_\varepsilon(u, t, x, y) - \Phi_\varepsilon(u, t, z, y)| \leq C \frac{|x-z|^{\gamma/2}}{(t-u)^{1-\gamma/4}} p_c(t-u, y-z). \quad (4.11)$$

Then, we can derive from (4.1), (4.9) and (4.11):

$$\begin{aligned} |R_\varepsilon^{\tau,2}(s, t, x, y)| &\leq C^2 \int_{s+\tau}^{(s+t)/2} du \int_A \frac{|x-z|^{\gamma/2}}{(u-s)} p_c(u-s, z-x) \frac{1}{(t-u)^{1-\gamma/4}} p_c(t-u, y-z) dz \\ &\quad + \frac{C}{(t-s)^{\gamma/4}} \int_{s+\tau}^{(s+t)/2} du \int_{A^c} \frac{|x-z|^{\gamma/2}}{(u-s)} p_c(u-s, z-x) \{|\Phi_\varepsilon(u, t, z, y)| + |\Phi_\varepsilon(u, t, x, y)|\} dz. \end{aligned} \quad (4.12)$$

From (4.5), we finally get on the considered time set:

$$\begin{aligned} |R_\varepsilon^{\tau,2}(s, t, x, y)| &\leq C p_c(t-s, y-x) \int_{s+\tau}^{(s+t)/2} du \frac{1}{(u-s)^{1-\gamma/4}} \frac{1}{(t-u)^{1-\gamma/4}} \\ &\leq \frac{C}{(t-s)^{1-\gamma/2}} p_c(t-s, y-x), \end{aligned}$$

which together with (4.10), (4.9), (4.6) and (4.2) gives the statement. It remains to establish (4.11). From the definition of Φ_ε and the smoothing effect of the kernel H_ε in (4.4), it suffices to prove that on the set $\bar{A} := \{z \in \mathbb{R}^d : |z-x| \leq c(u'-u)^{1/2}\}$:

$$|H_\varepsilon(u, u', x, w) - H_\varepsilon(u, u', z, w)| \leq C \frac{|x-z|^{\gamma/2}}{(u'-u)^{1-\gamma/4}} p_c(u'-u, w-z), \quad (4.13)$$

for $u' \in (u, t]$, $u \in [s, (s+t)/2]$. Observe that $\bar{A} \subset A$. Indeed, recalling that we want to establish (4.11) on A if $z \notin \bar{A}$, we get from (4.4):

$$\begin{aligned} & \int_u^t du' \int_{\bar{A}^c} |H_\varepsilon(u, u', x, w) - H_\varepsilon(u, u', z, w)| \left(\sum_{i \geq 2} H_\varepsilon^{(i)}(u', t, w, y) \right) dw \\ & \leq \int_u^t du' \int_{\bar{A}^c} \frac{C}{(u' - u)^{1-\gamma/2}} (p_c(u' - u, w - x) + p_c(u' - u, w - z)) \\ & \times \frac{|x - z|^{\gamma/2}}{(u' - u)^{\gamma/4}} \frac{C}{(t - u')^{1-\gamma}} p_c(t - u', y - w) dw \leq C \frac{|x - z|^{\gamma/2}}{(t - u)^{1-5\gamma/4}} p_c(t - u, y - z) \leq C \frac{|x - z|^{\gamma/2}}{(t - u)^{1-\gamma/4}} p_c(t - u, y - z), \end{aligned}$$

exploiting that $z \in A$, $t - u \geq \frac{1}{2}(t - s)$, and the usual convexity inequality $\frac{|y-x|^2}{t-u} \geq \frac{|y-z|^2}{2(t-u)} - \frac{|z-x|^2}{t-u} \geq \frac{|y-z|^2}{2(t-u)} - c$ for the last but one inequality. On the other hand, on \bar{A} we get (4.11) from (4.13) and (4.4).

Let us turn to the proof of (4.13). We concentrate on the second derivatives in H_ε which yield the most singular contributions:

$$\begin{aligned} & \text{Tr}((a_\varepsilon(u, x) - a_\varepsilon(u, w))D_x^2 \tilde{p}_\varepsilon(u, u', x, w)) - \text{Tr}((a_\varepsilon(u, z) - a_\varepsilon(u, w))D_x^2 \tilde{p}_\varepsilon(u, u', z, w)) \\ & = \text{Tr}((a_\varepsilon(u, x) - a_\varepsilon(u, z))D_x^2 \tilde{p}_\varepsilon(u, u', x, w)) - \text{Tr}((a_\varepsilon(u, z) - a_\varepsilon(u, w))(D_x^2 \tilde{p}_\varepsilon(u, u', z, w) - D_x^2 \tilde{p}_\varepsilon(u, u', x, w))) \\ & =: I + II. \end{aligned} \quad (4.14)$$

Then, from (4.1),

$$|I| \leq C \frac{|x - z|^\gamma}{(u - u')^\gamma} p_c(u' - u, w - x) \leq \frac{C|x - z|^{\gamma/2}}{(u - u')^{1-\gamma/4}} p_c(u' - u, w - x) \leq \frac{C|x - z|^{\gamma/2}}{(u - u')^{1-\gamma/4}} p_c(u' - u, w - z), \quad (4.15)$$

using that $z \in \bar{A}$ for the second inequality, again combined with the convexity inequality $\frac{|x-w|^2}{u'-u} \geq \frac{|z-w|^2}{2(u'-u)} - \frac{|x-z|^2}{u'-u} \geq \frac{|z-w|^2}{2(u'-u)} - c$ for the last one. Now, from the explicit expression of the second order derivatives in (4.7), **(A2)** and usual computations we also derive:

$$|II| \leq \frac{C|z - w|^\gamma}{(u' - u)} \frac{|z - x|^{\gamma/2}}{(u' - u)^{\gamma/4}} p_c(u' - u, w - z) \leq \frac{C|z - x|^{\gamma/2}}{(u' - u)^{1-\gamma/4}} p_c(u' - u, w - z). \quad (4.16)$$

This gives (4.13) and completes the proof for $|\alpha| \leq 2$.

Let us now turn to $|\alpha| \geq 3$. In those cases the singularities induced by the derivatives are not integrable in short time, even if we exploit cancellations. We are thus led to perform integration by parts, deteriorating the bounds since these operations make the derivatives of the mollified coefficients appear.

Recalling $\alpha \in \mathbb{N}^d$, denote by l a multi-index s.t. $|l| = 2$ and $\alpha - l \geq 0$ (where the inequality is to be understood componentwise). From equations (4.2), (4.6), we only have to consider the contribution $D_x^\alpha R_\varepsilon^T(s, t, x, y)$. Write:

$$\begin{aligned} D_x^\alpha R_\varepsilon^T(s, t, x, y) & = D_x^{\alpha-l} \int_{s+\tau}^{(t+s)/2} du \int_{\mathbb{R}^d} D_x^l \tilde{p}_\varepsilon(s, u, x, z) \Phi_\varepsilon(u, t, z, y) dz \\ & = D_x^{\alpha-l} \int_{s+\tau}^{(t+s)/2} du \int_{\mathbb{R}^d} g^{l, \varepsilon}(s, u, x, z) \Phi_\varepsilon(u, t, z, y) dz, \end{aligned} \quad (4.17)$$

where $g^{l,\varepsilon}(s, u, x, z) := D_x^l \tilde{p}_\varepsilon(s, u, x, z)$. Let us write introducing the cancellation term c_ε^l :

$$\begin{aligned}
D_x^\alpha R_\varepsilon^\tau(s, t, x, y) &= D_x^{\alpha-l} \int_{s+\tau}^{(s+t)/2} du \int_{\mathbb{R}^d} (g^{l,\varepsilon} - c_\varepsilon^l)(s, u, x, z) \Phi_\varepsilon(u, t, z, y) dz \\
&\quad + D_x^{\alpha-l} \int_{s+\tau}^{(s+t)/2} du \int_{\mathbb{R}^d} c_\varepsilon^l(s, u, x, z) (\Phi_\varepsilon(u, t, z, y) - \Phi_\varepsilon(u, t, x, y)) dz \\
&= D_x^{\alpha-l} \int_{s+\tau}^{(s+t)/2} du \int_{\mathbb{R}^d} (g^{l,\varepsilon} - c_\varepsilon^l)(s, u, x, x+z) \Phi_\varepsilon(u, t, x+z, y) dz \\
&\quad + D_x^{\alpha-l} \int_{s+\tau}^{(s+t)/2} du \int_{\mathbb{R}^d} c_\varepsilon^l(s, u, x, x+z) (\Phi_\varepsilon(u, t, x+z, y) - \Phi_\varepsilon(u, t, x, y)) dz.
\end{aligned} \tag{4.18}$$

The purpose of that change of variable, already performed in [KM02], is that we get integrable time singularities in the contributions $D_x^{\alpha-l}(g^{l,\varepsilon} - c_\varepsilon^l)(s, u, x, x+z)$. Anyhow, the mollified coefficients $b_\varepsilon, \sigma_\varepsilon$ have explosive derivatives. From the definition of $g^{l,\varepsilon}$ and (2.6) one easily gets that there exists c, C s.t. for all $\alpha, |\alpha| \leq 4$:

$$\begin{aligned}
|D_x^{\alpha-l}(g^{l,\varepsilon} - c_\varepsilon^l)(s, u, x, x+z)| &\leq \frac{C\varepsilon^{-|\alpha-l|}}{(u-s)^{1-\gamma/2}} p_c(u-s, z), \\
|D_x^{\alpha-l}c_\varepsilon^l(s, u, x, x+z)| &\leq \frac{C\varepsilon^{-|\alpha-l|+\gamma}}{(u-s)} p_c(u-s, z).
\end{aligned} \tag{4.19}$$

From (4.18) and (4.19) it thus remains to control the terms $D_x^{\alpha-l}\Phi_\varepsilon(u, t, z+x, y), D_x^{\alpha-l}(\Phi_\varepsilon(u, t, x+z, y) - \Phi_\varepsilon(u, t, x, y))$ which are the most singular ones in $D_x^\alpha R_\varepsilon^\tau(s, t, x, y)$. To this end, we will establish by induction that the following control holds:

$$\begin{aligned}
\exists c, C, \forall 0 \leq s < t \leq T, (x, y) \in (\mathbb{R}^d)^2, \forall \beta, |\beta| \leq 3, |D_x^\beta H_\varepsilon^{(i)}(s, t, x, y)| &\leq \\
\frac{C^i \varepsilon^{-|\beta|}}{(t-s)^{|\beta|/2}} (t-s)^{-1+i\gamma/2} \prod_{j=1}^{i-1} B(\gamma/2, j\gamma/2) p_c(t-s, y-x), &\tag{4.20}
\end{aligned}$$

with the convention that $\prod_{j=1}^0 = 1$. Observe first that for $|\beta| = 0$ (no derivation), estimate (4.20) readily follows from (4.3). Let us now suppose $|\beta| > 0$. Observe from the definition of H_ε that (4.20) is satisfied for $i = 1$. Let us assume it holds for a given i and let us prove it for $i + 1$. Write again:

$$\begin{aligned}
D_x^\beta H_\varepsilon^{(i+1)}(s, t, x, y) &= \int_{(s+t)/2}^t du \int_{\mathbb{R}^d} D_x^\beta H_\varepsilon(s, u, x, z) H_\varepsilon^{(i)}(u, t, z, y) dz \\
+ D_x^\beta \int_s^{(s+t)/2} du \int_{\mathbb{R}^d} H_\varepsilon(s, u, x, x+z) H_\varepsilon^{(i)}(u, t, x+z, y) dz &=: (R_1^{i,\beta} + R_2^{i,\beta})(s, t, x, y).
\end{aligned}$$

The term $R_1^{i,\beta}$ is easily controlled by (4.20) for $\beta = 0$ and the induction hypothesis. Observe also that, from Proposition 2 one derives similarly to (4.19) that:

$$|D_x^\beta H_\varepsilon(s, u, x, x+z)| \leq \frac{C\varepsilon^{-|\beta|}}{(u-s)^{1-\gamma/2}} p_c(u-s, z).$$

Together with the induction hypothesis and the Leibniz rule for differentiation, this allows to control $R_2^{i,\beta}$. The controls on $\{R_j^{i,\beta}\}_{j \in \{1,2\}}$ give (4.20) for $i+1$. We eventually derive (reminding that $|l|=2$):

$$|D_x^{\alpha-l} \Phi_\varepsilon(u, t, x+z, y)| \leq \frac{C}{(t-u)^{(|\alpha|-2)/2}} \frac{\varepsilon^{-|\alpha|+2}}{(t-u)^{1-\gamma/2}} p_c(t-u, y-(x+z)). \quad (4.21)$$

The spatial Hölder continuity of the derivatives of the kernel Φ_ε could be checked following the previous steps performed respectively to get the spatial Hölder continuity of the kernel and the controls on its derivatives. One gets, on $|z| \leq c(t-u)^{1/2}$:

$$|D_x^{\alpha-l} \Phi_\varepsilon(u, t, x+z, y) - D_x^{\alpha-l} \Phi_\varepsilon(u, t, x, y)| \leq \frac{C|z|^{\gamma/2}}{(t-u)^{(|\alpha|-2)/2}} \frac{\varepsilon^{-|\alpha|+2}}{(t-u)^{1-\gamma/4}} p_c(t-u, y-(x+z)),$$

which together with (4.21), (4.19), (4.18) gives (proceeding as above for $|z| \geq c(t-u)^{1/2}$):

$$|D_x^\alpha R_\varepsilon^\tau(s, t, x, y)| \leq \frac{C\varepsilon^{-|\alpha|+2}(t-s)^{\gamma/2}}{(t-s)^{|\alpha|/2}} p_c(t-s, y-x).$$

The second equation of (3.12) follows for $\bar{p}_\varepsilon = p_\varepsilon$ from the above control and (4.6), (4.2). Observe that the control for the derivative w.r.t. y has additional singularity in ε . This is clear since we directly differentiate the frozen mollified coefficients. Now the statements readily hold for p_ε^d , since the *integration in time* played no role in the previous computations. For p_ε^h , the only point that should be totally justified is the smoothing property and Hölder continuity of the discrete Kernel $\Phi_\varepsilon^h(t_i, t_j, x, y) := \sum_{r=1}^{j-i} H_\varepsilon^{h,(r)}(t_i, t_j, x, y)$. The smoothing property, equivalent of (4.20), has been investigated in [LM10]. The spatial Hölder continuity can be derived as above.

4.1.2. Proof under (\mathbf{A}_{PS})

Let us now turn to the proof of the heat kernel bounds for p_ε under (\mathbf{A}_{PS}) , which almost follows the same lines. Observe first that the result for $|\alpha|=1$ still follows from (4.1) and (3.8). The key point is again that the derivative of the Gaussian kernel yields an integrable singularity. For $|\alpha|=2$, we still separate the contribution $R_\varepsilon(s, t, x, y)$ as in (4.2) and again focus on $\lim_{\tau \rightarrow 0} D_x^\alpha R_\varepsilon^\tau(s, t, x, y)$ which is the only term yielding a potential singularity. With then notations of (4.9), it is sufficient to investigate $R_\varepsilon^{\tau,2}(s, t, x, y)$. Indeed, under (\mathbf{A}_{PS}) , equation (4.10) actually holds with $\gamma=1$. We recall that to control $R_\varepsilon^{\tau,2}(s, t, x, y)$, the key estimate was (4.11). We aim at proving the different control, for all $u \in [s, t]$, for all $\eta \in (0, 1]$:

$$|\Phi_\varepsilon(u, t, x, y) - \Phi_\varepsilon(u, t, z, y)| \leq C\varepsilon^{-\eta} \frac{|x-z|^\eta}{(t-u)^{3/4}} p_c(t-u, y-z), \quad (4.22)$$

on $A := \{z \in \mathbb{R}^d : |x-z| \leq (t-s)^{1/2} \wedge \varepsilon\}$. Then, we can derive from (4.1), (4.9) and (4.22):

$$\begin{aligned} |R_\varepsilon^{\tau,2}(s, t, x, y)| &\leq C\varepsilon^{-\eta} \int_{s+\tau}^{(s+t)/2} du \int_A \frac{|x-z|^\eta}{(u-s)} p_c(u-s, z-x) \frac{1}{(t-u)^{3/4}} p_c(t-u, y-z) dz \\ &+ C((t-s)^{1/2} \wedge \varepsilon)^{-\eta} \int_{s+\tau}^{(s+t)/2} du \int_{A^c} \frac{|x-z|^\eta}{(u-s)} p_c(u-s, z-x) \{|\Phi_\varepsilon(u, t, z, y)| + |\Phi_\varepsilon(u, t, x, y)|\} dz. \end{aligned}$$

Since the drift b_ε is uniformly bounded, uniformly in $\varepsilon \in [0, 1]$, we have under (\mathbf{A}_{PS}) the following usual control on the parametrix kernel (see e.g. Section 2 in [KKM15]):

$$|H_\varepsilon(u, t, z, y)| \leq \frac{c_1}{(t-u)^{1/2}} p_c(t-u, z-y). \quad (4.23)$$

Equation (4.23) for H_ε then yields

$$|H_\varepsilon^{(r)}(s, t, x, y)| \leq c_1^r \prod_{i=1}^{r-1} B\left(\frac{1}{2}, 1 + (i-1)\frac{1}{2}\right) p_c(t-s, y-x)(t-s)^{-1+\frac{r}{2}}, \quad (4.24)$$

again with the convention $\prod_{i=1}^0 = 1$. We thus derive $|\Phi_\varepsilon(u, t, z, y)| \leq \frac{C}{(t-u)^{1/2}} p_c(t-u, y-z)$ and $|\Phi_\varepsilon(u, t, x, y)| \leq \frac{C}{(t-u)^{1/2}} p_c(t-u, y-x)$. We finally get on the considered time set:

$$\begin{aligned} |R_\varepsilon^{\tau, 2}(s, t, x, y)| &\leq C((t-s)^{1/2} \wedge \varepsilon)^{-\eta} p_c(t-s, y-x) \int_{s+\tau}^{(s+t)/2} du \frac{1}{(u-s)^{1-\eta/2}} \frac{1}{(t-u)^{3/4}} \\ &\leq \frac{C((t-s)^{1/2} \wedge \varepsilon)^{-\eta}}{\eta(t-s)^{3/4-\eta/2}} p_c(t-s, y-x). \end{aligned}$$

It remains to establish (4.22). From the definition of Φ_ε and the smoothing effect of the kernel H_ε in (4.24), it suffices to prove that on $\bar{A} := \{z \in \mathbb{R}^d : |x-z| \leq c[(u'-u)^{1/2} \wedge \varepsilon]\}$:

$$|H_\varepsilon(u, u', x, w) - H_\varepsilon(u, u', z, w)| \leq C \left\{ \frac{|x-z|^\eta}{(u'-u)^{3/4}} ((u'-u)^{1/2} \wedge \varepsilon)^{-\eta} \right\} p_c(u'-u, w-z), \quad (4.25)$$

for $u' \in (u, t]$, $u \in [s, (s+t)/2]$. The contributions associated with $z \in \bar{A}^C$ can be handled as above. To establish the above control we focus on the first order terms involving the regularized coefficient with initial discontinuities. Indeed the second order contribution can be analyzed as in (4.14), (4.15), (4.16), taking $\gamma = 1$ in those expressions. In particular, the time singularity in $(u-u')^{3/4}$ in (4.25) comes precisely from those terms. Recalling that under (\mathbf{A}_{PS}) the driftless proxy does not depend on ε (since the diffusion is smooth, see (3.9) in which one has $\sigma_\varepsilon = \sigma$ under (\mathbf{A}_{PS})), we denote its density by \tilde{p} and write:

$$\begin{aligned} &\langle b_\varepsilon(u, x), D_x \tilde{p}(u, u', x, w) \rangle - \langle b_\varepsilon(u, z), D_x \tilde{p}(u, u', z, w) \rangle \\ &= \langle b_\varepsilon(u, x) - b_\varepsilon(u, z), D_x \tilde{p}(u, u', x, w) \rangle + \langle b_\varepsilon(u, z), D_x \tilde{p}(u, u', x, w) - D_x \tilde{p}(u, u', z, w) \rangle := I + II. \end{aligned}$$

On the one hand, from the mean value theorem and recalling that $|D_x b_\varepsilon|_\infty \leq C\varepsilon^{-1} \leq C((u'-u)^{1/2} \wedge \varepsilon)^{-1}$ we get:

$$\begin{aligned} |I| &\leq \frac{C}{(u'-u)^{1/2}} \left\{ 2|b|_\infty \left(\frac{|x-z|}{(u'-u)^{1/2} \wedge \varepsilon} \right)^\eta \mathbb{I}_{|x-z| > (u'-u)^{1/2} \wedge \varepsilon} + \varepsilon^{-1} |x-z| \mathbb{I}_{|x-z| \leq (u'-u)^{1/2} \wedge \varepsilon} \right\} p_c(u'-u, w-x) \\ &\leq C((u'-u)^{1/2} \wedge \varepsilon)^{-\eta} \frac{|x-z|^\eta}{(u'-u)^{1/2}} p_c(u'-u, w-x) \leq C((u'-u)^{1/2} \wedge \varepsilon)^{-\eta} \frac{|x-z|^\eta}{(u'-u)^{1/2}} p_c(u'-u, w-z), \end{aligned}$$

using again a convexity inequality for the last control, recalling that $z \in \bar{A}$. On the other hand still from the mean value Theorem and usual controls on the derivatives of the Gaussian density:

$$|II| \leq \frac{C|x-z|}{(u'-u)} \int_0^1 p_c(u'-u, w - \{z + \lambda(x-z)\}) d\lambda \leq \frac{C|x-z|^\eta}{(u'-u)^{\eta/2}} p_c(u'-u, w-z) \leq \frac{C|x-z|^\eta}{(u'-u)^{1/2}} p_c(u'-u, w-z).$$

The above estimates give (4.25) and concludes the proof for $|\alpha| = 2$.

Let us turn to $|\alpha| \geq 3$. The idea is again to proceed as under (\mathbf{A}_H) , up to a suitable modification of the key estimate (4.20) which can now be localized and becomes for all $q > d$:

$$\begin{aligned} & \exists c, C, \forall 0 \leq s < t \leq T, (x, y) \in (\mathbb{R}^d)^2, \forall \beta, |\beta| \leq 3, |D_x^\beta H_\varepsilon^{(i)}(s, t, x, y)| \leq \\ & \frac{C^i (\varepsilon^{-|\beta|} \mathbb{1}_{x \in V_\varepsilon(\mathcal{I})} + \varepsilon^{-|\beta|+1/q})}{(t-s)^{|\beta|/2}} (t-s)^{-1+i\alpha(q)} \prod_{j=1}^{i-1} B(\alpha(q), \alpha(q)j) p_c(t-s, y-x), \alpha(q) = \frac{1}{2} \left(1 - \frac{d}{q}\right), \end{aligned} \quad (4.26)$$

with $\prod_{j=1}^0 = 1$. We again proceed by induction. Observe first that for $|\beta| = 0$ (no derivation), estimate (4.26) readily follows from (4.23). Let us now suppose $|\beta| > 0$. Observe as well from the definition of H_ε that (4.26) is satisfied for $i = 1$. Let us assume it holds for a given i and let us prove it for $i + 1$. Write again:

$$\begin{aligned} D_x^\beta H_\varepsilon^{(i+1)}(s, t, x, y) &= \int_{(s+t)/2}^t du \int_{\mathbb{R}^d} D_x^\beta H_\varepsilon(s, u, x, z) H_\varepsilon^{(i)}(u, t, z, y) dz \\ &+ D_x^\beta \int_s^{(s+t)/2} du \int_{\mathbb{R}^d} H_\varepsilon(s, u, x, x+z) H_\varepsilon^{(i)}(u, t, x+z, y) dz =: (R_1^{i,\beta} + R_2^{i,\beta})(s, t, x, y). \end{aligned}$$

The term $R_1^{i,\beta}$ is easily controlled by (4.26) that holds from the induction hypothesis for $j = 0$ (direct differentiation of H_ε) and $\beta = 0$ for the considered i (no differentiation of $H_\varepsilon^{(i)}$). Observe also that, similarly to (4.19), one has:

$$|D_x^\beta H_\varepsilon(s, u, x, x+z)| \leq \frac{C(\varepsilon^{-|\beta|} \mathbb{1}_{x \in V_\varepsilon(\mathcal{I})} + 1)}{(u-s)^{1/2}} p_c(u-s, z). \quad (4.27)$$

Now, from the Leibniz rule for differentiation, (4.27) and the induction hypothesis, we have:

$$\begin{aligned} |R_2^{i,\beta}(s, t, x, y)| &\leq C^{i+1} \prod_{j=1}^{i-1} B(\alpha(q), \alpha(q)j) \left\{ \sum_{\tilde{\beta}, |\tilde{\beta}| \leq |\beta|} C_{|\tilde{\beta}|}^{|\tilde{\beta}|} \int_s^{(s+t)/2} \frac{du (t-u)^{-1+i\alpha(q)}}{(u-s)^{1/2} (t-u)^{(|\beta|-|\tilde{\beta}|)/2}} \right. \\ &\times \int_{\mathbb{R}^d} p_c(u-s, z) (\varepsilon^{-|\tilde{\beta}|} \mathbb{1}_{x \in V_\varepsilon(\mathcal{I})} + 1) (\varepsilon^{-(|\beta|-|\tilde{\beta}|)} \mathbb{1}_{x+z \in V_\varepsilon(\mathcal{I})} + \varepsilon^{-(|\beta|-|\tilde{\beta}|)+1/q}) p_c(t-u, y-x-z) dz \left. \right\} \\ &\leq C^{i+1} \prod_{j=1}^{i-1} B(\alpha(q), \alpha(q)j) \left\{ \sum_{\tilde{\beta}, |\tilde{\beta}| \leq |\beta|} C_{|\tilde{\beta}|}^{|\tilde{\beta}|} \int_s^{(s+t)/2} \frac{du (t-u)^{-1+i\alpha(q)}}{(u-s)^{1/2} (t-u)^{(|\beta|-|\tilde{\beta}|)/2}} \right. \\ &\times [p_c(t-s, y-x) (\varepsilon^{-|\beta|} \mathbb{1}_{x \in V_\varepsilon(\mathcal{I})} + \varepsilon^{-|\beta|+1/q}) + \varepsilon^{-|\beta|+1/q} \left(\int_{\mathbb{R}^d} p_c(u-s, z)^{\bar{q}} p_c(t-u, y-x-z)^{\bar{q}} dz \right)^{1/\bar{q}} \left. \right\}, \end{aligned} \quad (4.28)$$

denoting by $\bar{q} > 1$ the conjugate of q , $q^{-1} + \bar{q}^{-1} = 1$. Recall now that:

$$\begin{aligned} \left(\int_{\mathbb{R}^d} p_c(u-s, z)^{\bar{q}} p_c(t-u, y-x-z)^{\bar{q}} dz \right)^{1/\bar{q}} &= \left(\frac{c(t-s)}{(2\pi)(u-s)(t-u)} \right)^{d/2(1-1/\bar{q})} \bar{q}^{-d/(2\bar{q})} p_c(t-s, y-x) \\ &\leq C(u-s)^{-d/(2q)} p_c(t-s, y-x), \end{aligned}$$

for $u \in [s, (s+t)/2]$. Hence,

$$\begin{aligned} |R_2^{i,\beta}(s, t, x, y)| &\leq \frac{C^{i+1}}{(t-s)^{|\beta|/2}} \prod_{j=1}^{i-1} B(\alpha(q), \alpha(q)j) \left\{ \int_s^{(s+t)/2} \frac{du(t-u)^{-1+i\alpha(q)}}{(u-s)^{1/2(1+d/q)}} \right\} \\ &\quad \times p_c(t-s, y-x) (\varepsilon^{-|\beta|} \mathbb{I}_{x \in V_\varepsilon(\mathcal{I})} + \varepsilon^{-|\beta|+1/q}), \\ &\leq \frac{C^{i+1}}{(t-s)^{|\beta|/2}} \prod_{j=1}^{i-1} B(\alpha(q), \alpha(q)j) (t-s)^{-1+(i+1)\alpha(q)} \int_0^{1/2} (1-u)^{-1+i\alpha(q)} u^{-1+\alpha(q)} du \\ &\quad \times p_c(t-s, y-x) (\varepsilon^{-|\beta|} \mathbb{I}_{x \in V_\varepsilon(\mathcal{I})} + \varepsilon^{-|\beta|+1/q}). \end{aligned}$$

The controls on $\{R_j^{i,\beta}\}_{j \in \{1,2\}}$ give (4.26) for $i+1$.

Estimate (4.26) yields for every multi-index l , $|l| = 2$:

$$|D_x^{\alpha-l} \Phi_\varepsilon(u, t, x+z, y)| \leq \frac{C_q}{(t-u)^{(|\alpha|-2)/2}} \frac{\mathbb{I}_{x+z \in V_\varepsilon(\mathcal{I})} \varepsilon^{-|\alpha|+2} + \varepsilon^{-|\alpha|+2+1/q}}{(t-u)^{1-\alpha(q)}} p_c(t-u, y-(x+z)). \quad (4.29)$$

The spatial Hölder continuity of the derivatives of the kernel Φ_ε could be checked following the previous steps performed respectively to get the spatial Hölder continuity of the kernel and the controls on its derivatives. One gets, on $|z| \leq c\{(t-u)^{1/2} \wedge \varepsilon\}$ for all $\eta \in (0, 1]$:

$$\begin{aligned} &|D_x^{\alpha-l} \Phi_\varepsilon(u, t, x+z, y) - D_x^{\alpha-l} \Phi_\varepsilon(u, t, x, y)| \\ &\leq \frac{C_q \varepsilon^{-\eta} |z|^\eta}{(t-u)^{(|\alpha|-2)/2}} \frac{\varepsilon^{-|\alpha|+2} \mathbb{I}_{x+z \in V_\varepsilon(\mathcal{I})} + \varepsilon^{-|\alpha|+2+1/q}}{(t-u)^{1-\alpha(q)+\eta/2}} p_c(t-u, y-(x+z)). \end{aligned} \quad (4.30)$$

Now, equation (4.18) still holds under (\mathbf{A}_{PS}) , with $g^{l,\varepsilon} = g^l$, $c_\varepsilon^l = c^l$, i.e. the driftless proxy does not depend on ε . Also, the smoothness assumption on σ allows to improve (4.19). Precisely, there exist c, C s.t. for all α , $|\alpha| \leq 4$:

$$|D_x^{\alpha-l} (g^l - c^l)(s, u, x, x+z)| \leq \frac{C}{(u-s)^{1/2}} p_c(u-s, z), \quad |D_x^{\alpha-l} c^l(s, u, x, x+z)| \leq \frac{C}{(u-s)} p_c(u-s, z),$$

which together with (4.29), (4.30), (4.18) and choosing $\alpha(q) > \eta$ gives (proceeding as above for $|z| \geq c\{(t-u)^{1/2} \wedge \varepsilon\}$):

$$|D_x^\alpha R_\varepsilon^\tau(s, t, x, y)| \leq \frac{C_{\eta,q} \varepsilon^{-(\eta+|\alpha|)+2+1/q} (t-s)^{\eta/2}}{(t-s)^{|\alpha|/2}} p_c(t-s, y-x).$$

The controls on the derivatives w.r.t. to the forward variables are derived similarly. We here simply illustrate on the first term $\tilde{p}_\varepsilon \otimes H^\varepsilon(s, t, x, y)$ of the parametrix series how the derivatives must be handled. The stated controls would follow from inductions similar to the previous ones. Write for a given multi-index β :

$$\begin{aligned} &D_y^\beta (\tilde{p}_\varepsilon \otimes H_\varepsilon(s, t, x, y)) \\ &= \int_s^{(s+t)/2} du \int_{\mathbb{R}^d} \tilde{p}(s, u, x, z) D_y^\beta \{ (b_\varepsilon(u, z), D_z \tilde{p}(u, t, z, y)) \} + \frac{1}{2} \text{Tr} \{ (a(u, z) - a(u, y)) D_z^2 \tilde{p}(u, t, z, y) \} dz + \\ &\quad \lim_{\tau \downarrow 0} \int_{(s+t)/2}^{t-\tau} du \int_{\mathbb{R}^d} D_y^\beta (\tilde{p}(s, u, x, z)) [(b_\varepsilon(u, z), D_z \tilde{p}(u, t, z, y)) \\ &\quad + \frac{1}{2} \text{Tr} \{ (a(u, z) - a(u, y)) D_z^2 \tilde{p}(u, t, z, y) \}] dz := (D_1^\beta + D_2^\beta)(s, t, x, y). \end{aligned}$$

We readily get from the controls of (4.1) that:

$$|D_1^\beta(s, t, x, y)| \leq \frac{C}{(t-s)^{(|\beta|-1)/2}} p_c(t-s, y-x), \quad (4.31)$$

which is the expected control. Since a is smooth the terms involving the second derivatives w.r.t. z in D_2^β can be handled performing the change of variables $z' = z + y$ as above (see also [KM02] under the current smoothness assumption on the diffusion coefficient). Let us thus focus on the contribution:

$$D_{21}^\beta(s, t, x, y) := \lim_{\tau \downarrow 0} \int_{(s+t)/2}^{t-\tau} du \int_{\mathbb{R}^d} D_y^\beta \left(\tilde{p}(s, u, x, z) \langle b_\varepsilon(u, z), D_z \tilde{p}(u, t, z, y) \rangle \right) dz.$$

Consider first the case $|\beta| = 1$. Write:

$$\begin{aligned} D_{21}^{\beta, \tau}(s, t, x, y) &:= \int_{(s+t)/2}^{t-\tau} du \int_{\mathbb{R}^d} D_y^\beta \left(\tilde{p}(s, u, x, z) \langle b_\varepsilon(u, z), D_z \tilde{p}(u, t, z, y) \rangle \right) dz \\ &= \int_{(s+t)/2}^{t-\tau} du \int_{\mathbb{R}^d} \tilde{p}(s, u, x, z) \langle b_\varepsilon(u, z), D_y^\beta D_z \tilde{p}(u, t, z, y) \rangle dz \\ &= \int_{(s+t)/2}^{t-\tau} du \int_{\mathbb{R}^d} [\tilde{p}(s, u, x, z) - \tilde{p}(s, u, x, y)] \langle b_\varepsilon(u, z), D_y^\beta D_z \tilde{p}(u, t, z, y) \rangle dz \\ &\quad + \int_{(s+t)/2}^{t-\tau} du \int_{\mathbb{R}^d} \tilde{p}(s, u, x, y) \langle b_\varepsilon(u, z) - b_\varepsilon(u, y), D_y^\beta D_z \tilde{p}(u, t, z, y) \rangle dz \\ &=: [D_{211}^{\beta, \tau} + D_{212}^{\beta, \tau}](s, t, x, y), \end{aligned} \quad (4.32)$$

recalling that for all $y \in \mathbb{R}^d$, $\int_{\mathbb{R}^d} D_z \tilde{p}(u, t, z, y) dz = 0$, so that $D_y^\beta \int_{\mathbb{R}^d} D_z \tilde{p}(u, t, z, y) dz = 0$, for the last but one equality. Still from the controls of (4.1), we readily get:

$$\begin{aligned} |D_{211}^{\beta, \tau}| &\leq \frac{C}{(t-s)^{1/2}} \int_{(s+t)/2}^{t-\tau} du \int_{\mathbb{R}^d} |z-y| \left\{ \int_0^1 p_c(u-s, z-x + \lambda(y-z)) d\lambda \mathbb{1}_{|z-y| \leq (t-s)^{1/2}} \right. \\ &\quad \left. + (p_c(u-s, z-x) + p_c(u-s, y-x)) \mathbb{1}_{|z-y| > (t-s)^{1/2}} \right\} \frac{1}{(t-u)} p_c(t-u, y-z) dz \leq C p_c(t-s, y-x). \end{aligned}$$

On the other hand:

$$\begin{aligned} |D_{212}^{\beta, \tau}|(s, t, x, y) &\leq C p_c(t-s, y-x) \int_{(s+t)/2}^{t-\tau} du \int_{\mathbb{R}^d} \left\{ \mathbb{1}_{|z-y| \leq \varepsilon} \frac{|z-y|}{\varepsilon} + \mathbb{1}_{|z-y| > \varepsilon} \frac{|z-y|^\eta}{\varepsilon^\eta} \right\} \frac{1}{(t-u)} p_c(t-u, y-z) dz \\ &\leq \frac{C}{\eta} \varepsilon^{-\eta} (t-s)^{\eta/2} p_c(t-s, y-x), \quad \eta \in (0, 1]. \end{aligned}$$

We therefore eventually derive from the above controls, (4.32) and (4.31) that for $|\beta| = 1$

$$D_y^\beta \left(\tilde{p} \otimes H^\varepsilon(s, t, x, y) \right) \leq C p_c(t-s, y-x) \left\{ 1 + \frac{C}{\eta} \varepsilon^{-\eta} (t-s)^{\eta/2} \right\}.$$

Take now $|\beta| \geq 2$, and let l be a multi-index s.t. $|l| = 1$ and $\beta - l \geq 0$. Set for all $0 \leq u < t \leq T$, $(z, y) \in \mathbb{R}^d$, $g_l(u, t, z, y) := D_y^l D_z \tilde{p}(u, t, z, y)$. Observe that there exists $C \geq 1$, $|g_l(u, t, z, y)| \leq C(t-u)^{-1} p_c(t-u, y-z)$ and also, similarly to (4.27), for all multi-index $\tilde{\beta}$, $|\tilde{\beta}| \leq 4$, $|D_y^{\tilde{\beta}} g_l(u, t, z, y)| \leq C(t-u)^{-1} p_c(t-u, z)$. Rewrite

now from (4.32):

$$\begin{aligned} D_{21}^{\beta, \tau}(s, t, x, y) &= D_y^{\beta-l} \int_{(s+t)/2}^{t-\tau} du \int_{\mathbb{R}^d} (\tilde{p}(s, u, x, z) - \tilde{p}(s, u, x, y)) \langle b_\varepsilon(u, z), g_l(u, t, z, y) \rangle dz \\ &\quad + D_y^{\beta-l} \int_{(s+t)/2}^{t-\tau} du \int_{\mathbb{R}^d} \tilde{p}(s, u, x, y) \langle b_\varepsilon(u, z) - b_\varepsilon(u, y), g_l(u, t, z, y) \rangle dz, \end{aligned}$$

recalling that $\int_{\mathbb{R}^d} g_l(u, t, z, y) dz = 0$ for the last equality. Now,

$$\begin{aligned} |D_{21}^{\beta, \tau}(s, t, x, y)| &= \left| D_y^{\beta-l} \int_{(s+t)/2}^{t-\tau} du \int_{\mathbb{R}^d} (\tilde{p}(s, u, x, z+y) - \tilde{p}(s, u, x, y)) \langle b_\varepsilon(u, z+y), g_l(u, t, z+y, y) \rangle dz \right. \\ &\quad \left. + D_y^{\beta-l} \int_{(s+t)/2}^{t-\tau} du \int_{\mathbb{R}^d} \tilde{p}(s, u, x, y) \langle b_\varepsilon(u, z+y) - b_\varepsilon(u, y), g_l(u, t, z+y, y) \rangle dz \right| \\ &\leq \sum_{\substack{\beta_1, \beta_2, \beta_3, \\ \sum_{i=1}^3 |\beta_i| = |\beta| - 1}} (|\beta_1|, |\beta_2|, |\beta_3|)! \int_{(s+t)/2}^{t-\tau} du \int_{\mathbb{R}^d} |D_y^{\beta_1} \tilde{p}(s, u, x, z+y) - D_y^{\beta_1} \tilde{p}(s, u, x, y)| |D_y^{\beta_2} b_\varepsilon(u, z+y)| |D_y^{\beta_3} g_l(u, t, z+y, y)| dz \\ &\quad + \int_{(s+t)/2}^{t-\tau} du \int_{\mathbb{R}^d} |D_y^{\beta_1} \tilde{p}(s, u, x, y)| |D_y^{\beta_2} b_\varepsilon(u, z+y) - D_y^{\beta_2} b_\varepsilon(u, y)| |D_y^{\beta_3} g_l(u, t, z+y, y)| dz, \end{aligned}$$

where $(|\beta_1|, |\beta_2|, |\beta_3|)! = \frac{(\sum_{i=1}^3 |\beta_i|)!}{\prod_{i=1}^3 (|\beta_i|)!}$ stands for the multinomial coefficients with entries $(|\beta_i|)_{i \in [1,3]}$. Recall as well from (2.9) that we have the following control:

$$\begin{aligned} |D_y^{\beta_2} b_\varepsilon(u, z+y) - D_y^{\beta_2} b_\varepsilon(u, y)| &\leq C \left((1 + \varepsilon^{-|\beta|+1} (\mathbb{I}_{y+z \in V_\varepsilon(\mathcal{I})} + \mathbb{I}_{y \in V_\varepsilon(\mathcal{I})})) \frac{|z|}{(t-s)^{1/2}} \mathbb{I}_{|z| > (t-s)^{1/2}} \right. \\ &\quad \left. + ((1 + \varepsilon^{-|\beta|} \mathbb{I}_{y \in V_{2\varepsilon}(\mathcal{I})}) |z| \mathbb{I}_{|z| \leq \varepsilon} + (1 + \varepsilon^{-|\beta|+1} (\mathbb{I}_{y+z \in V_\varepsilon(\mathcal{I})} + \mathbb{I}_{y \in V_\varepsilon(\mathcal{I})})) \left(\frac{|z|}{\varepsilon}\right)^\eta \mathbb{I}_{|z| > \varepsilon}) \mathbb{I}_{|z| \leq (t-s)^{1/2}} \right) \\ &\leq \left((1 + \varepsilon^{-|\beta|+1} (\mathbb{I}_{y+z \in V_\varepsilon(\mathcal{I})} + \mathbb{I}_{y \in V_\varepsilon(\mathcal{I})})) \frac{|z|}{(t-s)^{1/2}} \mathbb{I}_{|z| > (t-s)^{1/2}} \right. \\ &\quad \left. + ((1 + \varepsilon^{-|\beta|+1} \mathbb{I}_{y \in V_{2\varepsilon}(\mathcal{I})}) \mathbb{I}_{|z| \leq \varepsilon} + (1 + \varepsilon^{-|\beta|+1} (\mathbb{I}_{y+z \in V_\varepsilon(\mathcal{I})} + \mathbb{I}_{y \in V_\varepsilon(\mathcal{I})})) \mathbb{I}_{|z| > \varepsilon}) \left(\frac{|z|}{\varepsilon}\right)^\eta \mathbb{I}_{|z| \leq (t-s)^{1/2}} \right). \end{aligned}$$

Thus,

$$\begin{aligned} |D_{21}^{\beta, \tau}(s, t, x, y)| &\leq C \sum_{\substack{\beta_1, \beta_2, \beta_3, \\ \sum_{i=1}^3 |\beta_i| = |\beta| - 1}} (|\beta_1|, |\beta_2|, |\beta_3|)! \int_{(s+t)/2}^{t-\tau} du \int_{\mathbb{R}^d} \left\{ \frac{|z|}{(u-s)^{(|\beta_1|+1)/2}} p_c(u-s, x-y+\lambda z) \mathbb{I}_{|z| \leq (t-s)^{1/2}} \right. \\ &\quad \left. + \frac{|z|}{(t-s)^{1/2}} (p_c(u-s, y+z-x) + p_c(u-s, y-x)) \mathbb{I}_{|z| > (t-s)^{1/2}} \right\} (1 + \varepsilon^{-|\beta|+1} \mathbb{I}_{y+z \in V_\varepsilon(\mathcal{I})}) \frac{1}{t-u} p_c(t-u, z) dz \\ &\quad + \int_{(s+t)/2}^{t-\tau} du \int_{\mathbb{R}^d} \frac{1}{(t-s)^{|\beta_1|/2}} p_c(t-s, y-x) \left\{ (1 + \varepsilon^{-|\beta|+1} (\mathbb{I}_{y+z \in V_\varepsilon(\mathcal{I})} + \mathbb{I}_{y \in V_\varepsilon(\mathcal{I})})) \frac{|z|}{(t-s)^{1/2}} \mathbb{I}_{|z| > (t-s)^{1/2}} \right. \\ &\quad \left. + \{(1 + \varepsilon^{-|\beta|+1} \mathbb{I}_{y \in V_{2\varepsilon}(\mathcal{I})}) \mathbb{I}_{|z| \leq \varepsilon} + (1 + \varepsilon^{-|\beta|+1} (\mathbb{I}_{y \in V_\varepsilon(\mathcal{I})} + \mathbb{I}_{y+z \in V_\varepsilon(\mathcal{I})})) \mathbb{I}_{|z| > \varepsilon}\} \mathbb{I}_{|z| \leq (t-s)^{1/2}} \left(\frac{|z|}{\varepsilon}\right)^\eta \right\} \\ &\quad \frac{1}{t-u} p_c(t-u, z) dz \leq \frac{C_{|\beta|}}{(t-s)^{|\beta|}} p_c(t-s, y-x) (1 + \varepsilon^{-|\beta|+1} (\frac{\varepsilon^{-\eta}}{\eta} + \frac{\varepsilon^{1/q}}{\alpha(q)})), \end{aligned}$$

recalling that the contribution in $\frac{\varepsilon^{1/q}}{\alpha(q)}$ comes from the terms involving $\frac{\mathbb{1}_{y+z \in V_\varepsilon(\mathcal{I})}}{(t-u)^{1/2}}$ that can be handled using Hölder inequalities similarly to (4.28). This gives the stated control.

4.2. Proof of Proposition 5

Write similarly to the proof of Theorem 2.1 in [KM02]:

$$\begin{aligned} (p_\varepsilon - p_\varepsilon^d)(t_i, t_j, x, y) &= (p_\varepsilon \otimes H_\varepsilon - p_\varepsilon \otimes_h H_\varepsilon)(t_i, t_j, x, y) + (p_\varepsilon - p_\varepsilon^d) \otimes_h H_\varepsilon(t_i, t_j, x, y) \\ &= \sum_{r \geq 0} (p_\varepsilon \otimes H_\varepsilon - p_\varepsilon \otimes_h H_\varepsilon) \otimes_h H_\varepsilon^{(r)}(t_i, t_j, x, y), \end{aligned} \quad (4.33)$$

where we apply iteratively the first equality to get the second one. From (4.4), that holds for $\gamma = 1$ under (\mathbf{A}_{PS}) , the key point is thus to control $p_\varepsilon \otimes H_\varepsilon - p_\varepsilon \otimes_h H_\varepsilon$. Write:

$$\begin{aligned} & (p_\varepsilon \otimes H_\varepsilon - p_\varepsilon \otimes_h H_\varepsilon)(t_i, t_j, x, y) \\ &= \sum_{k=0}^{j-i-1} \int_{t_i+kh}^{t_i+(k+1)h} du \int_{\mathbb{R}^d} dz \{ p_\varepsilon(t_i, u, x, z) H_\varepsilon(u, t_j, z, y) - p_\varepsilon(t_i, t_{i+k}, x, z) H_\varepsilon(t_{i+k}, t_j, z, y) \} \\ &= \sum_{k=0}^{j-i-1} \left\{ \int_{t_i+kh}^{t_i+(k+1)h} du \int_{\mathbb{R}^d} dz \{ [p_\varepsilon(t_i, u, x, z) - p_\varepsilon(t_i, t_{i+k}, x, z)] H_\varepsilon(u, t_j, z, y) \} \right. \\ &+ \left. \int_{t_i+kh}^{t_i+(k+1)h} du \int_{\mathbb{R}^d} dz \{ p_\varepsilon(t_i, t_{i+k}, x, z) [H_\varepsilon(u, t_j, z, y) - H_\varepsilon(t_{i+k}, t_j, z, y)] \} \right\} =: (D_\varepsilon^{d,1} + D_\varepsilon^{d,2})(t_i, t_j, x, y). \end{aligned} \quad (4.34)$$

For the term $D_\varepsilon^{d,1}$ we first write:

$$\begin{aligned} p_\varepsilon(t_i, u, x, z) - p_\varepsilon(t_i, t_{i+k}, x, z) &= (u - t_{i+k}) \int_0^1 (\partial_v p_\varepsilon(t_i, v, x, z))_{v=t_{i+k}+\lambda(u-t_{i+k})} d\lambda \\ &= (u - t_{i+k}) \int_0^1 (L_v^{\varepsilon,*} p_\varepsilon(t_i, v, x, z))_{v=t_{i+k}+\lambda(u-t_{i+k})} d\lambda. \end{aligned}$$

Under (\mathbf{A}_H) , reproducing the integration by parts strategy that led to (3.18), we then derive:

$$|D_\varepsilon^{d,1}|(t_i, t_j, x, y) \leq C \frac{h\varepsilon^{-2+\gamma}}{(t_j - t_i)^{1-\gamma/2}} p_c(t_j - t_i, y - x). \quad (4.35)$$

Similarly, under (\mathbf{A}_{PS}) , from the computations described in Section 3.4, we derive in the general case that for all $q > d, \eta \in (0, \alpha(q))$:

$$|D_\varepsilon^{d,1}|(t_i, t_j, x, y) \leq C_{\eta,q} \frac{h\varepsilon^{-1-\eta}}{(t_j - t_i)^{1-\eta/2}} p_c(t_j - t_i, y - x). \quad (4.36)$$

On the other hand introduce:

$$\begin{aligned} & (\bar{D}_\varepsilon^{d,21} + \bar{D}_\varepsilon^{d,22})(t_i, t_{i+k}, u, t_j, x, y) := \\ & C \int_{\mathbb{R}^d} p_c(t_{i+k} - t_i, z - x) |a_\varepsilon(u, z) - a_\varepsilon(u, y) - (a_\varepsilon(t_{i+k}, z) - a_\varepsilon(t_{i+k}, y))| \frac{1}{t_j - t_{i+k}} p_c(t_j - t_{i+k}, y - z) dz \\ & + \left| \int_{\mathbb{R}^d} p_\varepsilon(t_i, t_{i+k}, x, z) (a_\varepsilon(u, z) - a_\varepsilon(u, y)) [D_z^2 \tilde{p}_\varepsilon(u, t_j, z, y) - D_z^2 \tilde{p}_\varepsilon(t_{i+k}, t_j, z, y)] \right|, \end{aligned}$$

that correspond to the most singular contributions in $D_\varepsilon^{d,2}$ as far as the time singularity is concerned. For $\bar{D}_\varepsilon^{d,22}$ we can again perform Taylor expansion in time, use the Kolmogorov equations and integrate by parts as above to derive under (\mathbf{A}_H) :

$$\sum_{k=0}^{j-i-1} \int_{t_i+kh}^{t_i+(k+1)h} du |\bar{D}_\varepsilon^{d,22}|(t_i, t_{i+k}, u, t_j, x, y) \leq C \frac{h\varepsilon^{-2+\gamma}}{(t_j - t_i)^{1-\gamma/2}} p_c(t_j - t_i, y - x). \quad (4.37)$$

On the other hand, using the $\gamma/2$ -Hölder continuity in time of a we get:

$$\begin{aligned} & |\bar{D}_\varepsilon^{d,21}(t_i, t_{i+k}, u, t_j, x, y)| \\ & \leq C \int_{\mathbb{R}^d} p_c(t_{i+k} - t_i, z - x) |u - t_{i+k}|^{\gamma/2} \frac{1}{t_j - t_{i+k}} p_c(t_j - t_{i+k}, y - z) dz \\ & \leq Ch^{(\gamma-\eta)/2} p_c(t_j - t_i, y - x) (t_j - t_{i+k})^{-1+\eta/2}, \end{aligned}$$

for $\eta \in (0, \gamma)$. Plugging now the above control, (4.37), (4.35) in (4.34) we derive the result under (\mathbf{A}_H) from (4.4) and (4.33).

Now under (\mathbf{A}_{PS}) , the previous strategy yields:

$$\sum_{k=0}^{j-i-1} \int_{t_i+kh}^{t_i+(k+1)h} du |\bar{D}_\varepsilon^{d,22}|(t_i, t_{i+k}, u, t_j, x, y) \leq C_{\eta,q} \frac{h\varepsilon^{-1-\eta}}{(t_j - t_i)^{1-\eta/2}} p_c(t_j - t_i, y - x). \quad (4.38)$$

Also, the smoothness in time (Lipschitz continuity) of the diffusion coefficients gives for $\eta \in (0, 1]$,

$$|\bar{D}_\varepsilon^{d,21}(t_i, t_{i+k}, u, t_j, x, y)| \leq Ch^{1-\eta/2} p_c(t_j - t_i, y - x) (t_j - t_{i+k})^{-1+\eta/2},$$

for $\eta \in (0, 1]$. Again, the results under (\mathbf{A}_{PS}) follows from (4.4) and (4.33) plugging the above control, (4.38), (4.36) in (4.34). In the particular case $\sigma(t, x) = \sigma$, the bound of the proposition is derived similarly, using again Proposition 4, equation (3.13), since in that case at most one integration by part is needed.

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