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# ITERATED BOOLEAN RANDOM VARIETIES AND APPLICATION TO FRACTURE STATISTICS MODELS 

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#### Abstract

Models of random sets and of point processes are introduced to simulate some specific clustering of points, namely on random lines in $\mathbb{R}^{2}$ and $\mathbb{R}^{3}$ and on random planes in $\mathbb{R}^{3}$. The corresponding point processes are special cases of Cox processes. The generating distribution function of the probability distribution of the number of points in a convex set $K$ and the Choquet capacity $T(K)$ are given. A possible application is to model point defects in materials with some degree of alignment. Theoretical results on the probability of fracture of convex specimens in the framework of the weakest link assumption are derived, and are used to compare geometrical effects on the sensitivity of materials to fracture.


## 1. Introduction

Point processes showing clustering effects are interesting models to simulate nonhomogenous location of points in space, as seen for instance for some defects in materials: for polycrystals modellled by random tessellations, defects can be located on the grain boundaries; in composite materials, they can appear on fibers of a network. The aim of this paper is to introduce some random sets models based on point processes reproducing these kinds of situation, and to study some of their theroretical probabilistic properties.

After a reminder on random sets obtained from Boolean random varieties in $\mathbb{R}^{n}$, two-steps varieties in $\mathbb{R}^{n}$ are introduced and characterized. The cases of point processes in $\mathbb{R}^{2}$ and in $\mathbb{R}^{3}$ are detailed to generate random points on lines and on planes. A three-steps Poisson points in $\mathbb{R}^{3}$ enables us to take into account alignements in Poisson planes. All these point processes are particular cases of Cox point processes, for which the generating function of the probability distribution of the number of points in a convex set $K$ and the Choquet capacity $T(K)$ are calculated.

In a last section, we make use of iterated Boolean varieties to propose new probabilistic models of fracture based on the weakest link assumption, that can be applied to model the intergranular fracture of polycrystals or the fiber fracture of composites.

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## 2. Reminder on Boolean random varieties

In this section are given the construction of Boolean random sets based on random varieties in $\mathbb{R}^{n}$, and their main probabilistic properties, namely their Choquet capacity.
2.1. Construction and properties of the linear Poisson varieties model in $\mathbb{R}^{n}$. A geometrical introduction of the Poisson linear varieties is as follows (Matheron 1975): a Poisson point process $\left\{x_{i}(\omega)\right\}$, with intensity $\theta_{k}(d \omega)$ is considered on the varieties of dimension $(n-k)$ containing the origin $O$, and with orientation $\omega$. On every point $x_{i}(\omega)$ is located a variety with dimension $k, V_{k}(\omega)_{x_{i}}$, orthogonal to the direction $\omega$. By construction, we have $V_{k}=\cup_{x_{i}(\omega)} V_{k}(\omega)_{x_{i}}$. For instance in $R^{3}$ can be built a network of Poisson hyperplanes $\Pi_{\alpha}$ (orthogonal to the lines $D_{\omega}$ containing the origin) or a network of Poisson lines in every plane $\Pi_{\omega}$ containing the origin.

Definition 1. In $\mathbb{R}^{n}$, $n$ Poisson linear varieties of dimension $k(k=0,1, \ldots, n-1)$ $V_{k}$, can be built: the Poisson point process for $k=0$, and the Poisson hyperplanes for $k=n-1$. For $k \geq 1$, a network of Poisson linear varieties of dimension $k$ can be considered as a Poisson point process in the space $S_{k} \times \mathbb{R}^{n-k}$, with intensity $\theta_{k}(d \omega) \mu_{n-k}(d x) ; \theta_{k}$ is a positive Radon measure for the set of subspaces of dimension $k, S_{k}$, and $\mu_{n-k}$ is the Lebesgue measure of $\mathbb{R}^{n-k}$.

If $\theta_{k}(d \omega)$ is any Radon measure, the obtained varieties are anisotropic. When $\theta_{k}(d \omega)=\theta_{k} d \omega$, the varieties are isotropic. If the Lebesgue measure $\mu_{n-k}(d x)$ is replaced by a measure $\theta_{n-k}(d x)$, non stationary random varieties are obtained.

The probabilistic properties of the Poisson varieties are easily derived from their definition as a Poisson point process.

Theorem 1. The number of varieties of dimension $k$ hit by a compact set $K$ is a Poisson variable, with parameter $\theta(K)$ :

$$
\begin{equation*}
\theta(K)=\int \theta_{k}(d \omega) \int_{K(\omega)} \theta_{n-k}(d x)=\int \theta_{k}(d \omega) \theta_{n-k}(K(\omega)) \tag{2.1}
\end{equation*}
$$

where $K(\omega)$ is the orthogonal projection of $K$ on the orthogonal space to $V_{k}(\omega)$, $V_{k^{\perp}}(\omega)$. For the stationary case,

$$
\begin{equation*}
\theta(K)=\int \theta_{k}(d \omega) \mu_{n-k}(K(\omega)) \tag{2.2}
\end{equation*}
$$

The Choquet capacity $T(K)=P\left\{K \cap V_{k} \neq \varnothing\right\}$ of the varieties of dimension $k$ is given by

$$
\begin{equation*}
T(K)=1-\exp \left(-\int \theta_{k}(d \omega) \int_{K(\omega)} \theta_{n-k}(d x)\right) \tag{2.3}
\end{equation*}
$$

In the stationary case, the Choquet capacity is

$$
\begin{equation*}
T(K)=1-\exp \left(-\int \theta_{k}(d \omega) \mu_{n-k}(K(\omega))\right) \tag{2.4}
\end{equation*}
$$

Proof. By construction, the random varieties $V_{k}(\omega)$ induce by intersection on every orthogonal variety of dimension $n-k, V_{k^{\perp}}(\omega)$, a Poisson point process with dimension $n-k$ and with intensity $\theta_{k}(d \omega) \theta_{n-k}(d x)$. Therefore, the contribution of the direction $\omega$ to $N(K)$, is the Poisson variable $N(K, \omega)$ with intensity $\theta_{n-k}(K(\omega))$.

Since the contributions of the various directions are independent, Eq. (2.1) results immediately.

Proposition 2. Consider now the isotropic ( $\theta_{k}$ being constant) and stationary case, and a convex set $K$. Due to the symmetry of the isotropic version, we can consider $\theta_{k}(d \omega)=\theta_{k} d \omega$ as defined on the half unit sphere (in $\mathbb{R}^{k+1}$ ) of the directions of the varieties $V_{k}(\omega)$. The number of varieties of dimension $k$ hit by a compact set $K$ is a Poisson variable, with parameter $\theta(K)$ given by:

$$
\begin{equation*}
\theta(K)=\theta_{k} \int \mu_{n-k}(K(\omega)) d \omega=\theta_{k} \frac{b_{n-k} b_{k+1}}{b_{n}} \frac{k+1}{2} W_{k}(K) \tag{2.5}
\end{equation*}
$$

where $b_{k}$ is the volume of the unit ball in $\mathbb{R}^{k}\left(b_{k}=\frac{\pi^{k / 2}}{\Gamma\left(1+\frac{k}{2}\right)}\right)\left(b_{1}=2, b_{2}=\pi, b_{3}=\right.$
$\left.\frac{4}{3} \pi\right)$, and $W_{k}(K)$ is the Minkowski's functional of $K$, homogeneous and of degree $n-k$ (Matheron 1975).

The following examples are useful for applications:

- When $k=n-1$, the varieties are Poisson planes in $\mathbb{R}^{n}$; in that case, $\theta(K)=\theta_{n-1} n W_{n-1}(K)=\theta_{n-1} \mathcal{A}(K)$, where $\mathcal{A}(K)$ is the norm of $K$ (average projected length over orientations).
- In the plane $\mathbb{R}^{2}$ are obtained the Poisson lines, with $\theta(K)=\theta L(K), L$ being the perimeter.
- In the three-dimensional space $\mathbb{R}^{3}$ are obtained Poisson lines for $k=1$ and Poisson planes for $k=2$. For Poisson lines, $\theta(K)=\frac{\pi}{4} \theta S(K)$ and for Poisson planes, $\theta(K)=\theta M(K)$, where $S$ and $M$ are the surface area and the integral of the mean curvature.
2.2. Boolean random varieties. Boolean random sets can be built, starting
from Poisson varieties and a random primary grain (Jeulin 1991; Jeulin 1991a, Jeulin 2015a).

Definition 2. A Boolean model with primary grain $A^{\prime}$ is built on Poisson linear varieties in two steps: i) we start from a network $V_{k}$; ii) every variety $V_{k \alpha}$ is dilated by an independent realization of the primary grain $A^{\prime}$. The Boolean RACS A is given by

$$
A=\cup_{\alpha} V_{k \alpha} \oplus A^{\prime}
$$

By construction, this model induces on every variety $V_{k^{\perp}}(\omega)$ orthogonal to $V_{k}(\omega)$ a standard Boolean model with dimension $n-k$, with random primary grain $A^{\prime}(\omega)$ and with intensity $\theta(\omega) d \omega$. The Choquet capacity of this model immediately follows, after averaging over the directions $\omega$; it can also be deduced from Eq. (2.4), after replacing $K$ by $A^{\prime} \oplus \check{K}$ and averaging.

Theorem 3. The Choquet capacity of the Boolean model built on Poisson linear varieties of dimension $k$ is given by

$$
\begin{equation*}
T(K)=1-\exp \left(-\int \theta_{k}(d \omega) \bar{\mu}_{n-k}\left(A^{\prime}(\omega) \oplus \check{K}(\omega)\right)\right) \tag{2.6}
\end{equation*}
$$

For isotropic varieties, the Choquet capacity of Boolean varieties is given by

$$
\begin{equation*}
T(K)=1-\exp \left(-\theta_{k} \frac{b_{n-k} b_{k+1}}{b_{n}} \frac{k+1}{2} \bar{W}_{k}\left(A^{\prime} \oplus \check{K}\right)\right) \tag{2.7}
\end{equation*}
$$

Particular cases of Eq. (2.6) are obtained when $K=\{x\}$ (giving the probability $q=P\left\{x \in A^{c}\right\}=\exp \left(-\int \theta_{k}(d \omega) \mu_{n-k}\left(A^{\prime}(\omega)\right)\right)$ and when $K=\{x, x+h\}$, giving the covariance of $A^{c}, Q(h)$ :

$$
\begin{equation*}
Q(h)=q^{2} \exp \left(\int \theta_{k}(d \omega) K_{n-k}(\omega, \vec{h} \cdot \vec{u}(\omega))\right) \tag{2.8}
\end{equation*}
$$

where $K_{n-k}(\omega, h)=\bar{\mu}_{n-k}\left(A^{\prime}(\omega) \cap A_{-h}^{\prime}(\omega)\right)$ and $\vec{u}(\omega)$ is the unit vector with the direction $\omega$. For a compact primary grain $A^{\prime}$, there exists for any $h$ an angular sector where $K_{n-k}(\omega, h) \neq 0$, so that the covariance generally does not reach its sill, at least in the isotropic case, and the integral range, obtained by integration of the correlation function, is infinite. Consider now some examples in $\mathbb{R}^{2}$ and in $\mathbb{R}^{3}$.
2.2.1. Fibers in 2D. In the plane can be built a Boolean model on Poisson lines.

For an isotropic lines network (figure 1), and if $A^{\prime} \oplus \check{K}$ is a convex set, we have, from equation (2.7):

$$
\begin{equation*}
T(K)=1-\exp \left(-\theta \bar{L}\left(A^{\prime} \oplus \check{K}\right)\right) \tag{2.9}
\end{equation*}
$$

If $A^{\prime} \oplus \check{K}$ is not a convex set, the integral of projected lengths over a line with the orientation varying between 0 and $\pi$ must be taken. If $A^{\prime}$ and $K$ are convex sets, we have $\bar{L}\left(A^{\prime} \oplus \check{K}\right)=\bar{L}\left(A^{\prime}\right)+L(K)$. Consider now the isotropic case. Using for $A^{\prime}$ a random disc with a random radius $R$ (with expectation $\bar{R}$ ) and for $K$ a disc with radius $r$, equation 2.9 becomes:

$$
\begin{aligned}
& T(r)=1-\exp (-2 \pi \theta(\bar{R}+r)) \\
& T(0)=P\{x \in A\}=1-\exp (-2 \pi \theta \bar{R})
\end{aligned}
$$

which can be used to estimate $\theta$ and $\bar{R}$, and to validate the model.
In $\mathbb{R}^{3}$, can be built a Boolean model on Poisson planes or on Poisson lines.
2.2.2. Boolean model on Poisson planes in $\mathbb{R}^{3}$. A Boolean model built on Poisson planes generates a structure with strata. On isotropic Poisson planes, we have for a convex set $A^{\prime} \oplus \check{K}$ by application of equation (2.7):

$$
\begin{equation*}
T(K)=1-\exp \left(-\theta \bar{M}\left(A^{\prime} \oplus \check{K}\right)\right) \tag{2.10}
\end{equation*}
$$

When $A^{\prime}$ and $K$ are convex sets, we have $\bar{M}\left(A^{\prime} \oplus \check{K}\right)=\bar{M}\left(A^{\prime}\right)+M(K)$. If $A^{\prime} \oplus \check{K}$ is not convex, $T(K)$ is expressed as a function of the length $l$ of the projection over the lines $D_{\omega}$ by $T(K)=1-\exp \left(-\theta \int_{2 \pi s t e r} \bar{l}\left(A^{\prime}(\omega) \oplus \check{K}(\omega)\right) d \omega\right)$. For instance if $A^{\prime}$ is a random sphere with a random radius $R$ (with expectation $\bar{R}$ ) and $K$ is a sphere with radius $r$, equation 2.10 becomes:

$$
\begin{aligned}
& T(r)=1-\exp (-4 \pi \theta(\bar{R}+r)) \\
& T(0)=P\{x \in A\}=1-\exp (-4 \pi \theta \bar{R})
\end{aligned}
$$

which can be used to estimate $\theta$ and $\bar{R}$, and to validate the model.
2.2.3. Boolean model on Poisson lines in $\mathbb{R}^{3}$. A Boolean model built on Poisson lines generates a fiber network, with possible overlaps of fibers. On isotropic Poisson lines, we have for a convex set $A^{\prime} \oplus \check{K}$

$$
\begin{equation*}
T(K)=1-\exp \left(-\theta \frac{\pi}{4} \bar{S}\left(A^{\prime} \oplus \check{K}\right)\right) \tag{2.11}
\end{equation*}
$$

If $A^{\prime} \oplus \check{K}$ is not a convex set, $T(K)$ is expressed as a function of the area $A$ of the projection over the planes $\Pi_{\omega}$ by

$$
\begin{equation*}
T(K)=1-\exp \left(-\frac{1}{2} \theta \int_{2 \pi s t e r} \bar{A}\left(A^{\prime}(\omega) \oplus \check{K}(\omega)\right) d \omega\right) \tag{2.12}
\end{equation*}
$$

If $A^{\prime}$ is a random sphere with a random radius $R$ (with expectation $\bar{R}$ and second moment $E\left(R^{2}\right)$ ) and $K$ is a sphere with radius $r$, equation 2.11 becomes:

$$
\begin{aligned}
& T(r)=1-\exp \left(-\pi^{2} \theta\left(E\left(R^{2}\right)+2 r \bar{R}+r^{2}\right)\right) \\
& T(0)=P\{x \in A\}=1-\exp \left(-\pi^{2} \theta E\left(R^{2}\right)\right)
\end{aligned}
$$

which can be used to estimate $\theta, E\left(R^{2}\right)$ and $\bar{R}$, and to validate the model. A model of Poisson fibers parallel to a plane, and with a uniform distribution of orientations in the plane was used to model cellulosic fiber materials (Delisée et al. 2001). In (Schladitz et al. 2006), non isotropic dilated Poisson lines were used to model and to optimize the acoustic absorption of nonwoven materials.

## 3. Two steps Boolean varieties

It is possible to generate further Boolean models by iteration of Poisson varieties. For instance in $\mathbb{R}^{2}$, we first consider Poisson lines, and in a second step Poisson points on every line. These points are germs to locate primary grains $A^{\prime}$ to generate a Boolean model. Compared to the standard Boolean model, this one shows alignments of grains. Similarly in $\mathbb{R}^{3}$ we can start from Poisson planes $V_{2 \alpha}$ and use Poisson lines $V_{1 \beta}$ in every plane to generate a Boolean model with fibers. In contrast with Poisson fibers in $\mathbb{R}^{3}$, this model generates a random set with some coplanar fibers. Such long range random sets could mimic specific microstructures with an order in a lower dimension subspace of $\mathbb{R}^{n}$, such as preferred germination of objects on specific planes or lines.

These models are based on doubly stochastic Poisson random variables for which the Choquet capacity can be obtained.

Definition 3. Two steps random varieties are defined as follows: starting from Poisson linear varieties $V_{k}$ of dimension $k$ and with intensity $\theta_{k}(d \omega)$ in $\mathbb{R}^{n}$, Poisson linear varieties $V_{k^{\prime} \beta}$ with dimension $0 \leq k^{\prime}<k$ and with intensity $\theta_{k^{\prime}}(d \omega)$ are implanted on each $V_{k \alpha}$. Then each $V_{k^{\prime} \beta}$ is dilated by independent realizations of a random compact primary grain $A^{\prime} \subset \mathbb{R}^{n}$ to generate the Boolean RACS A:

$$
A=\cup_{\beta} V_{k^{\prime} \beta} \oplus A^{\prime}
$$

Remark 1. By construction, when $k^{\prime}=0$ the varieties $V_{k^{\prime} \beta}$ are a particular case of a Cox process driven by the random set $V_{k}$, and the derived random set $A$ is a Cox Boolean model (Jeulin 2012).

In what follows the purpose is restricted to the stationary and isotropic case, with the two intensities $\theta_{k}$ and $\theta_{k^{\prime}}$.

Theorem 4. The number $N(K)$ of varieties of dimension $k^{\prime}<k$ hit by the compact set $K$ is a random variable with generating function

$$
\begin{equation*}
G_{k^{\prime}}(s, K)=E\left\{s^{N(K)}\right\}=\exp \left[\theta_{k} a_{k} W_{k}(K)\left[\varphi_{k^{\prime}}\left(\theta_{k^{\prime}} a_{k^{\prime}}(1-s), K\right)-1\right]\right] \tag{3.1}
\end{equation*}
$$

where $a_{k^{\prime}}=\frac{b_{n-k^{\prime}} b_{k^{\prime}+1}}{b_{n}} \frac{k^{\prime}+1}{2}$ and $\varphi_{k^{\prime}}(\lambda, K)$ is the Laplace transform of the random variable $W_{k^{\prime}}\left(K \cap V_{k \alpha}\right)$, $W_{k^{\prime}}$ being the Minkowski functional homogeneous with degree $k-k^{\prime}$ in $\mathbb{R}^{k}$ :

$$
\begin{equation*}
\varphi_{k^{\prime}}(\lambda, K)=E\left\{\exp \left[-\lambda W_{k^{\prime}}\left(K \cap V_{k}\right)\right]\right\} \tag{3.2}
\end{equation*}
$$

the mathematical expectation being taken over the realizations $V_{k \alpha}$. As a consequence, the Choquet capacity of the Boolean RACS A built on the Poisson linear varieties $V_{k^{\prime}}$ using a deterministic primary grain $A^{\prime}$ is derived from $G_{k^{\prime}}\left(0, A^{\prime} \oplus \check{K}\right)$, $E\left\}\right.$ being the expectation with respect to the random variety $V_{k \alpha}$ :

$$
\begin{align*}
& 1-T(K)  \tag{3.3}\\
& =\exp \left[-\theta_{k} a_{k} W_{k}\left(A^{\prime} \oplus \check{K}\right)\left[1-E\left\{\exp \left[-\theta_{k^{\prime}} a_{k^{\prime}} W_{k^{\prime}}\left(A^{\prime} \oplus \check{K} \cap V_{k}\right)\right]\right\}\right]\right]
\end{align*}
$$

Proof. The random number $N_{k}$ of varieties $V_{k \alpha}$ hit by $K$ is a Poisson variable with expectation $\theta_{k} a_{k} W_{k}(K)$. On each $V_{k \alpha}$ are generated $N_{k^{\prime}}$ varieties $V_{k^{\prime} \beta}, N_{k^{\prime}}$ being a Poisson variable with expectation $\theta_{k^{\prime}} a_{k^{\prime}} W_{k^{\prime}}\left(K \cap V_{k}\right)$. For a random section $K \cap V_{k \alpha}$, the generating function of $N_{k^{\prime}}$ is

$$
\begin{equation*}
\Gamma(s)=\exp \left[-\theta_{k^{\prime}} a_{k^{\prime}} W_{k^{\prime}}\left(K \cap V_{k \alpha}\right)(1-s)\right] \tag{3.4}
\end{equation*}
$$

Taking the expectation of (3.4) with respect to $W_{k^{\prime}}\left(K \cap V_{k \alpha}\right)$ and then of $\Gamma(s)^{N_{k}}$ gives (3.1).

The Choquet capacity requires the use of the Laplace transform $\varphi_{k^{\prime}}\left(\lambda, A^{\prime} \oplus \check{K}\right)$. It is not easy to expressed them in a closed form for specific compact sets $K$ and $A^{\prime}$. However the required distribution functions and their Laplace transforms can be estimated by simulation of the random variables obtained from random variables $W_{k^{\prime}}\left(A^{\prime} \oplus \check{K} \cap V_{k \alpha}\right)$ obtained from random sections. Examples of closed form expressions are given now for two-steps Poisson points in $\mathbb{R}^{2}$ and in $\mathbb{R}^{3}$.
3.1. Poisson points on Poisson lines in the plane. Starting from Poisson lines in the plane, a 1D Poisson point process is independently generated on each lines.

Proposition 5. The generating function $G_{K}(s)$ of the random number of points $N_{P}(K)$ contained in a convex set $K$ in $\mathbb{R}^{2}$ with perimeter $\mathcal{L}(K)$, random intercept length $L(K)$ (with Laplace transform $\varphi_{L}(\lambda, K)$ ), is given by

$$
\begin{equation*}
G_{K}(s)=\exp \left(-\theta_{1} \mathcal{L}(K)\left(1-\varphi_{L}(\theta(1-s), K)\right)\right) \tag{3.5}
\end{equation*}
$$

We have

$$
\begin{equation*}
1-T(K)=Q(K)=\exp \left(-\theta_{1} \mathcal{L}(K)\left(1-\varphi_{L}(\theta, K)\right)\right) \tag{3.6}
\end{equation*}
$$

The Choquet capacity of the corresponding Boolean model for convex sets $K$ and $A^{\prime}$ is obtained by replacing $K$ by $\left(A^{\prime} \oplus \check{K}\right)$ in equation 3.6.

Proof. $K$ hits a Poisson random number of lines $N_{D}$ with parameter $\theta_{1} \mathcal{L}(K)$. Each chord with random length $L(K)$ contains a Poisson number of points with parameter $\theta L(K)$ and generating function

$$
\Gamma(s, L(K))=\exp (-\theta L(K)(1-s))
$$

After deconditioning with respect to $L(K)$, with Laplace transform $\varphi_{L}(\lambda, K)$, we obtain the generating function $\Gamma(s, K)=E_{L}\{\Gamma(s, L(K))\}=\varphi_{L}(\theta(1-s), K)$. Then we consider the sum of $N_{D}$ realizations of the random variable $L(K)$ to obtain equation 3.5 by expectation with respect to $N_{D}$ :

$$
\begin{aligned}
G_{K}(s) & =E_{N_{D}}\left\{\varphi_{L}(\theta(1-s), K)^{N_{D}}\right\} \\
& =\exp \left(\theta_{1} \mathcal{L}(K)\left(\varphi_{L}(\theta(1-s), K)-1\right)\right)
\end{aligned}
$$

When $K$ is the disc $C(r)$ with radius $r$, the generating function of the random number of points $N_{P}(r)$ in $C(r)$ is obtained by

$$
G(s, r)=\exp \left[-2 \pi r \theta_{1}\left(1-\varphi_{L}(\theta(1-s), r)\right)\right]
$$

with $\varphi_{L}(\lambda, r)$ given by equation 7.1 or 7.2 . We have

$$
Q(r)=\exp \left[-2 \pi r \theta_{1}\left(1-\varphi_{L}(\theta, r)\right)\right]
$$

3.2. Poisson points on Poisson planes in $\mathbb{R}^{3}$. This point process is obtained in two steps:
(1) We start with Poisson planes in $\mathbb{R}^{3}$ (consider here the isotropic case), with intensity $\theta_{2}$
(2) On each Poisson plane, is generated a 2D Poisson point process, with intensity $\theta$.

Proposition 6. The generating function $G_{K}(s)$ of the random number of points $N_{P}(K)$ contained in a convex set $K$ is given by

$$
\begin{equation*}
G_{K}(s)=\exp \left[-\theta_{2} M(K)\left(1-\psi_{A}(\theta(1-s), K \cap \pi)\right)\right] \tag{3.7}
\end{equation*}
$$

and we get

$$
\begin{equation*}
1-T(K)=Q(K)=\exp \left[-\theta_{2} M(K)\left(1-\psi_{A}(\theta, K \cap \pi)\right)\right] \tag{3.8}
\end{equation*}
$$

with: $M(K)$ : integral of mean curvature of $K ; A(K \cap \pi)$ : area of sections of $K$ by a random plane $\pi$, with Laplace transform $\psi_{A}(\lambda, K \cap \pi)$

The Choquet capacity of the corresponding Boolean model for convex sets $K$ and $A^{\prime}$ is obtained by replacing $K$ by $\left(A^{\prime} \oplus \check{K}\right)$ in equation 3.8.

Proof. The random number of planes $N_{\pi}(K)$ hit by $K$ is a Poisson variable with parameter $\theta_{2} M(K)$. Each plane $\pi$ cuts $K$ according to a convex random set with area $A(K \cap \pi)$, containing a Poisson number of points, with parameter $\theta A(K)$ and generating function

$$
\Gamma(s, A(K))=\exp (-\theta A(K)(1-s))
$$

After deconditioning with respect to $A(K)$, with Laplace transform $\psi_{A}(\lambda, K \cap \pi)$, we obtain the generating function $\Gamma(s, K)=E_{A}\{\Gamma(s, A(K))\}=\psi_{A}(\theta(1-s), K \cap \pi)$. Then we consider the sum of $N_{\pi}$ realizations of the random variable $A(K)$ to obtain equation 3.7 by expectation with respect to $N_{\pi}$ :

$$
\begin{aligned}
G_{K}(s) & =E_{N_{\pi}}\left\{\psi_{A}(\theta(1-s), K \cap \pi)^{N_{\pi}}\right\} \\
& =\exp \left[\theta_{2} M(K)\left(\psi_{A}(\theta(1-s), K \cap \pi)-1\right)\right]
\end{aligned}
$$

The generating function $G(s, r)$ of the random number of points $N_{P}(r)$ in the sphere with radius $r$ is given by

$$
G(s, r)=\exp \left[-4 \pi r \theta_{2}(1-\psi(\theta \pi(1-s), r))\right]
$$

and

$$
1-T(r)=Q(r)=\exp \left[-4 \pi r \theta_{2}(1-\psi(\theta \pi, r))\right]
$$

with (see equation 7.4)

$$
\psi(\lambda, r)=\frac{\exp \left(-\lambda r^{2}\right) \int_{0}^{r \sqrt{\lambda}} \exp \left(y^{2}\right) d y}{r \sqrt{\lambda}}
$$

3.3. Poisson points on Poisson lines in $\mathbb{R}^{3}$. This point process is obtained in two steps:
(1) We start from isotropic Poisson lines in $\mathbb{R}^{3}$, with intensity $\theta_{1}$
(2) On each Poisson line, a 1D Poisson point process with intensity $\theta$ is generated.

Proposition 7. The generating function $G_{K}(s)$ of the random number of points $N_{D}(K)$ contained in a convex set $K$ is given by

$$
\begin{equation*}
G_{K}(s)=\exp \left[-\frac{\pi}{4} \theta_{1} S(K)\left(1-\varphi_{L}(\theta(1-s), K)\right)\right] \tag{3.9}
\end{equation*}
$$

so that

$$
\begin{equation*}
1-T(K)=Q(K)=\exp \left[-\frac{\pi}{4} \theta_{1} S(K)\left(1-\varphi_{L}(\theta, K)\right)\right] \tag{3.10}
\end{equation*}
$$

where $S(K)$ is the surface area of $K$, and $\varphi_{L}(\lambda, K)$ the Laplace transform of $a$ random chord $L(K)$ in $K$.

The Choquet capacity of the corresponding Boolean model for convex sets $K$ and $A^{\prime}$ is obtained by replacing $K$ by $\left(A^{\prime} \oplus \check{K}\right)$ in equation 3.10.
Proof. The random number of lines $N_{D}(K)$ hit by $K$ is a Poisson variable with parameter $\frac{\pi}{4} \theta_{1} S(K)$. Each line cuts $K$ according to a random chord with length $L(K)$, containing a Poisson number of points with parameter $\theta L(K)$ and with generating function

$$
\Gamma(s, L(K))=\exp (\theta L(K)(s-1))
$$

After deconditioning with respect to $L(K)$, with Laplace transform $\varphi_{L}(\lambda, K)$, we obtain the generating function $\Gamma(s, K)=E_{L}\{\Gamma(s, L(K))\}=\varphi_{L}(\theta(1-s), K)$. Then we consider the sum of $N_{D}$ realizations of the random variable $L(K)$ to obtain equation 3.9 by expectation with respect to $N_{D}$ :

$$
\begin{aligned}
G_{K}(s) & =E_{N_{D}}\left\{\varphi_{L}(\theta(1-s), K)^{N_{D}}\right\} \\
& =\exp \left[-\frac{\pi}{4} \theta_{1} S(K)\left(1-\varphi_{L}(\theta(1-s), K)\right)\right]
\end{aligned}
$$

The generating function $G(s, r)$ of the random number of points $N_{P}(r)$ in the sphere with radius $r$ is expressed from

$$
\begin{aligned}
\log (G(s, r))= & \\
& -\pi^{2} \theta_{1} r^{2}\left(1-\frac{2}{(2 r \theta(1-s))^{2}}[1-(1+2 r \theta(1-s)) \exp (-2 r \theta(1-s))]\right)
\end{aligned}
$$

and

$$
1-T(r)=Q(r)=\exp \left[-\pi^{2} \theta_{1} r^{2}\left(1-\frac{2}{(2 r \theta)^{2}}[1-(1+2 r \theta) \exp (-2 r \theta)]\right)\right]
$$

## 4. Three-steps Poisson points in $\mathbb{R}^{3}$

A new point process is obtained by a three steps iteration: Poisson points on 2D Poisson lines contained in Poisson planes
(1) We start from Poisson planes in $\mathbb{R}^{3}$ (isotropic case), with intensity $\theta_{2}$
(2) On each Poisson plane, a 2 D Poisson lines process with intensity $\theta_{1}$ is generated
(3) On each line, a 1D Poisson point process with intensity $\theta$ is given

Proposition 8. Consider a convex compact set $K$, with random planar sections $K \cap \pi$. The generating function $G_{K}(s)$ of the random number of points $N_{D}(K)$ contained in the convex set $K$ is given from

$$
\begin{align*}
& \log \left(G_{K}(s)\right)  \tag{4.1}\\
= & -\theta_{2} M(K)\left(E_{\pi}\left\{\exp \left[\theta_{1} \mathcal{L}(K \cap \pi)\left(\varphi_{L}(\theta(1-s), K \cap \pi)-1\right)\right]\right\}-1\right)
\end{align*}
$$

where $E_{\pi}$ is the mathematical expectation over random sections. We get

$$
\begin{aligned}
(4 \sqrt[1 D g]{g}(1-T(K)) & = \\
\log (Q(K)) & =
\end{aligned}
$$

$$
-\theta_{2} M(K)\left(E_{\pi}\left\{\exp \left[\theta_{1} \mathcal{L}(K \cap \pi)\left(\varphi_{L}(\theta, K \cap \pi)-1\right)\right]\right\}-1\right)
$$

with perimeter $\mathcal{L}(K \cap \pi)$ of sections of $K$, with Laplace transform $\psi_{\mathcal{L}}(\lambda, K \cap \pi)$, random chord of each planar section $L(K \cap \pi)$, with Laplace transform $\varphi_{L}(\lambda, K \cap \pi)$.

The Choquet capacity of the corresponding Boolean model for convex sets $K$ and $A^{\prime}$ is obtained by replacing $K$ by $\left(A^{\prime} \oplus \check{K}\right)$ in equation 4.2.

Proof. The random number of planes $N_{\pi}(K)$ hit by $K$ is a Poisson variable with parameter $\theta_{2} M(K)$. Each plane $\pi$ cuts $K$ according to a convex random set with random perimeter $\mathcal{L}(K \cap \pi)$ hitting a Poisson number of lines, with parameter $\theta_{1} \mathcal{L}(K \cap \pi)$. Each line cuts $K \cap \pi$ according to a random chord $L(K \cap \pi)$ containing a Poisson number of points with parameter $\theta L(K)$ and generating function

$$
\Gamma(s, L(K \cap \pi))=\exp (-\theta L(K \cap \pi)(1-s))
$$

For a given section $K \cap \pi$, the generating function of the number of points on a line is obtained by deconditioning over $L(K \cap \pi)$, so that

$$
\Gamma(s, K \cap \pi)=\varphi_{L}(\theta(1-s), K \cap \pi)
$$

The generating function of the random number of points on $K \cap \pi$ is given by the expectation of $\Gamma(s, K \cap \pi)^{N}$, $N$ being the Poisson variable with parameter $\theta_{1} \mathcal{L}(K \cap \pi)$ and therefore

$$
E\left\{\Gamma(s, K \cap \pi)^{N}\right\}=\exp \left[\theta_{1} \mathcal{L}(K \cap \pi)\left(\varphi_{L}(\theta(1-s), K \cap \pi)-1\right)\right]
$$

Deconditioning with respect to the random section $K \cap \pi$ gives

$$
\Gamma(s, K)=E_{\pi}\left\{\exp \left[\theta_{1} \mathcal{L}(K \cap \pi)\left(\varphi_{L}(\theta(1-s), K \cap \pi)-1\right)\right]\right\}
$$

Deconditioning now with respect to the Poisson number of planes $N_{\pi}(K)$, we take the expectation of $\Gamma(s, K)^{N_{\pi}(K)}$ to get equation 4.1.

The generating function of the number of points of the process inside a sphere with radius $r$ is given by

$$
\begin{aligned}
& \log (G(s, r)) \\
= & -4 \pi \theta_{2} r\left(1-E_{R}\left\{\exp \left[2 \pi \theta_{1} R\left(\varphi_{L}(\theta(1-s), R)-1\right)\right]\right\}\right)
\end{aligned}
$$

with

$$
\begin{aligned}
& E_{R}\left\{\exp \left[2 \pi \theta_{1} R\left(\varphi_{L}(\theta(1-s), R)-1\right)\right]\right\} \\
= & \int_{0}^{r} \exp \left[2 \pi \theta_{1} u\left(\varphi_{L}(\theta(1-s), u)-1\right)\right] f(u, r) d u
\end{aligned}
$$

- $f(u, r)$ : distribution function of the radius $R$ of random sections of a sphere
- $\varphi_{L}(\lambda, R)$ : Laplace transform of random chords of the disc with radius $R$ The Choquet capacity for a sphere with radius $r$ is given by

$$
\begin{aligned}
\log (1-T(r))= & \log (Q(r)) \\
& =-4 \pi \theta_{2} r\left(1-\int_{0}^{r} \exp \left[2 \pi \theta_{1} u\left(\varphi_{L}(\theta, u)-1\right)\right] f(u, r) d u\right)
\end{aligned}
$$

## 5. Use of Iterated Boolean varieties for probabilistic models of FRACTURE BASED ON THE WEAKEST LINK ASSUMPTION

The standard weakest link model is based on the assumption that fracture in a brittle material is initiated on the most critical defect, that controls the full fracture process. For this model, it means that when there is at least one point $x$ in a specimen where the applied principal stress component $\sigma(x)$ is larger than the local critical stress $\sigma_{c}(x)$, the specimen is broken. Usually it is assumed that the occurrence or absence of critical defects (generating fracture) of any volume elements generate a set of independent events. After a decomposition of the volume $V$ into links $v_{i}$ and assuming that there is a fracture of the volume $V$ when a single link $v_{i}$ is broken, a classical computation for independent events gives:

$$
P\{\text { Non fracture }\}=\prod_{i} P\left\{\text { Non fracture of } v_{i}\right\}
$$

For $v_{i} \rightarrow 0, P\{$ fracture $\} \simeq \Phi((\sigma(x)) d x$, with $\Phi$ increasing with the loading $\sigma$ and $P\{$ Non fracture of $d x\} \simeq 1-\Phi((\sigma(x)) d x$. Therefore with these assumptions,

$$
P\{\text { Non fracture of } V\}=\exp \left(-\int_{V} \Phi(\sigma(x)) d x\right)=\exp \left(-V \Phi\left(\sigma_{e q}\right)\right)
$$

where the equivalent stress is defined from

$$
\Phi\left(\sigma_{e q}\right)=\frac{1}{V} \int_{V} \Phi((\sigma(x)) d x
$$

This assumption is equivalent to a distribution of point defects in a matrix with $\sigma_{c}=\infty$, according to a Poisson point process in space, with intensity $\Phi(\sigma)$, where $\Phi(\sigma)$ is the average number per unit volume of defects with a critical stress $\sigma_{c}$ lower than $\sigma$.

For a homogeneous applied stress field $\sigma(x)=\sigma$,

$$
P\{\text { Non fracture of } V\}=\exp (-V \Phi(\sigma))
$$

For the Weibull model, the function $\Phi(\sigma)$ is power law in $\sigma: \Phi(\sigma)=\theta\left(\sigma-\sigma_{0}\right)^{m}$ and $P\{$ Non fracture of $V\}$ follows a 3 parameters Weibull distribution.

The weakest link model corresponds to a "Infimum" Boolean random function with point support primary random functions (PRF), and is immediately extended to any PRF with a support having almost surely compact sections (Jeulin 1991, 2001, 2012, 2015a).

In what follows the weakest link model is extended to the case of the various point processes introduced in the first parts of this paper. It allows for clustering of defects on Poisson varieties. A comparison is made with the standard Poissonbased model and between the various models, when using the same function $\Phi(\sigma)$ for point defects.
5.1. Fracture statistics for Poisson point defects on Poisson lines in $\mathbb{R}^{2}$. As seen before, a two steps point process can be used to locate random defects:
(1) Poisson lines in $\mathbb{R}^{2}$ (isotropic case), with intensity $\theta_{1}$
(2) On each Poisson plane, a 1D Poisson point process of point defects acting in fracture, with intensity $\theta$ replaced by $\Phi(\sigma)$ in equation 3.6.
We get

$$
\begin{equation*}
P\left\{\sigma_{R} \geq \sigma\right\}_{L}=\exp \left(-\theta_{1} \mathcal{L}(K)\left(1-\varphi_{L}(\Phi(\sigma), K)\right)\right) \tag{5.1}
\end{equation*}
$$

When $K$ is the disc with radius $r$,

$$
\begin{equation*}
P\left\{\sigma_{R} \geq \sigma\right\}_{L}=Q(r, \sigma)=\exp \left[2 \pi r \theta_{1}\left(\varphi_{L}(\Phi(\sigma), r)-1\right)\right] \tag{5.2}
\end{equation*}
$$

5.2. Comparison of Fracture statistics for Poisson points and for points on lines in $\mathbb{R}^{2}$. In the plane, the average number of Poisson points contained in the disc of radius $r$ is

$$
E\left\{N_{P}(r)\right\}=\pi r^{2} \theta_{2}
$$

The average number of Poisson lines hit by the disc is $2 \pi r \theta_{1}$. Therefore, the average number of points of the two-step process on lines is $2 \pi r \theta_{1} \theta E\{L\}, E\{L\}$ being the average chord of the disc. We have $-\pi K^{\prime}(0)=2 \pi r$ and then $-K^{\prime}(0)=2 r$, so that $E\{L\}=\frac{\pi r 2}{2 r}=\pi \frac{r}{2}$. The average number of points on lines is given by :

$$
E\left\{N_{P}(r)\right\}=2 \pi r \theta_{1} \theta \pi \frac{r}{2}=\pi^{2} r^{2} \theta_{1} \theta
$$

To compare the two fracture statistics, we consider the same average number of defects in the disc, so that we have to use

$$
\theta_{2}=\pi \theta_{1} \theta
$$

We have

$$
\begin{aligned}
& \log \left(P\left\{\sigma_{R} \geq \sigma\right\}_{P}\right)-\log \left(P\left\{\sigma_{R} \geq \sigma\right\}_{L}\right) \\
= & 2 \pi r \theta_{1}\left(1-\varphi_{L}(\theta, r)-\frac{\pi \theta_{1} \theta}{2 \pi r \theta_{1}} \pi r^{2}\right) \\
= & 2 \pi r \theta_{1}\left(1-\varphi_{L}(\theta, r)-\theta \frac{\pi}{2} r\right)
\end{aligned}
$$

Using the parameter $\alpha=2 \theta r$ we have to compute

$$
2 \pi r \theta_{1}\left(1-\frac{\pi}{2}(S t r u v e L(-1, \alpha)-\operatorname{BesselI}(1, \alpha))-\frac{\pi}{4} \alpha\right)
$$

From numerical calculation, it turns out that this expression remains negative for any $\alpha$ and then

$$
\left(P\left\{\sigma_{R} \geq \sigma\right\}_{P}\right)<P\left\{\sigma_{R} \geq \sigma\right\}_{L}
$$

This result is satisfied for any intensity $\Phi(\sigma)$. In 2 D , it is easier to break a specimen with Poisson point defects than with point defects on Poisson lines.
5.3. Fracture statistics for Poisson point defects on Poisson planes in $\mathbb{R}^{3}$. As earlier, we locate point defects according to a two steps point process:
(1) Poisson planes in $\mathbb{R}^{3}$ (isotropic case), with intensity $\theta_{2}$
(2) On each Poisson plane, 2D Poisson point process of point defects, with intensity $\theta$ replaced by $\Phi(\sigma)$ in equation 3.8.
Considering the Poisson tessellation generated by Poisson planes, this model figures out point defects located on grain boundaries, generating intergranular fracture. We get

$$
\begin{equation*}
P\left\{\sigma_{R} \geq \sigma\right\}_{\pi}=\exp \left[-\theta_{2} M(K)\left(1-\psi_{A}(\Phi(\sigma), K \cap \pi)\right)\right] \tag{5.3}
\end{equation*}
$$

In the case of a spherical specimen with radius $r$,

$$
\begin{equation*}
P\left\{\sigma_{R} \geq \sigma\right\}_{\pi}=\exp \left[-4 \pi r \theta_{2}(1-\psi(\pi \Phi(\sigma), r))\right] \tag{5.4}
\end{equation*}
$$

with $\psi(\lambda, r)$ given by equation 7.4.
5.4. Fracture statistics for Poisson point defects on Poisson lines in $\mathbb{R}^{3}$. A model of long fiber network with point defects is obtained from Poisson lines, where we replace $\theta$ by $\Phi(\sigma)$ in equation 3.10:

$$
\begin{equation*}
P\left\{\sigma_{R} \geq \sigma\right\}_{D}=\exp \left[-\frac{\pi}{4} \theta_{1} S(K)\left(1-\varphi_{L}(\Phi(\sigma), K)\right)\right] \tag{5.5}
\end{equation*}
$$

In the case of a spherical specimen with radius $r$,

$$
\begin{aligned}
\left(5 . 6 P \left\{\sigma_{R} \geq\right.\right. & \sigma\}_{D}= \\
& \exp \left[-\pi^{2} \theta_{1} r^{2}\left(1-\frac{2}{(2 r \Phi(\sigma))^{2}}[1-(1+2 r \Phi(\sigma)) \exp (-2 r \Phi(\sigma))]\right]\right.
\end{aligned}
$$

5.5. Comparison of Fracture statistics for Poisson points and for points on planes. With consider the fracture of a sphere of sphere of radius $r$ containing a random number of points $N_{P}(r)$ with a given average.

For the standard Poisson point process,

$$
E\left\{N_{P}(r)\right\}=\frac{4}{3} \pi r^{3} \theta_{3}
$$

For Poisson points on Poisson planes,

$$
E\left\{N_{P}(r)\right\}=\frac{8}{3} \pi^{2} r^{3} \theta_{2} \theta
$$

For a fixed average number of defects in the sphere of radius $r$, we get

$$
\theta_{3}=2 \pi \theta_{2} \theta
$$

Using the same intensity $\Phi(\sigma)=\theta$ for the two processes

$$
\begin{aligned}
\log \left(P \left\{\sigma_{R}\right.\right. & \left.\geq \sigma\}_{P}\right)-\log \left(P\left\{\sigma_{R} \geq \sigma\right\}_{\pi}\right) \\
& =4 \pi r \theta_{2}\left(1-\psi(\theta \pi, r)-\frac{\theta r^{2}}{2}\right)
\end{aligned}
$$

and

$$
P\left\{\sigma_{R} \geq \sigma\right\}_{\pi}<P\left\{\sigma_{R} \geq \sigma\right\}_{P} \text { for } r^{2} \Phi(\sigma)<1.8
$$

Given the same statistics of defects $\Phi(\sigma)$, for low applied stresses, or at a small scale, the "intergranular" fracture probability is higher than the standard probability of fracture. For high applied stresses the reverse is true, and the material is less sensitive to "intergranular" fracture. The two probability curves cross for $r^{2} \Phi(\sigma) \simeq 1.8$.
5.6. Comparison of Fracture statistics for Poisson points and for points on lines. We consider again the fracture of a sphere of sphere of radius $r$ containing a random number of points $N_{P}(r)$ with a given average.

For the standard Poisson point process,

$$
E\left\{N_{P}(r)\right\}=\frac{4}{3} \pi r^{3} \theta_{3}
$$

For Poisson points on Poisson lines,

$$
E\left\{N_{P}(r)\right\}=\frac{4}{3} \pi^{2} r^{3} \theta_{1} \theta
$$

Given the average number of defects in the sphere of radius $r$, we have

$$
\theta_{3}=\pi \theta_{1} \theta
$$

Using the same intensity $\Phi(\sigma)=\theta$ for the two processes, and the auxiliary variable $\alpha=2 r \theta$, we have

$$
\begin{aligned}
\log \left(P \left\{\sigma_{R}\right.\right. & \left.\geq \sigma\}_{D}\right)-\log \left(P\left\{\sigma_{R} \geq \sigma\right\}_{P}\right) \\
& =\pi^{2} \theta_{1} r^{2}\left(\frac{2}{\alpha^{2}}(1-(1+\alpha) \exp (-\alpha))-1+\frac{2}{3} \alpha\right)
\end{aligned}
$$

and

$$
P\left\{\sigma_{R} \geq \sigma\right\}_{D}<P\left\{\sigma_{R} \geq \sigma\right\}_{P}
$$

Therefore, given the same statistics of defects $\Phi(\sigma)$, the "fiber" fracture probability is higher than the standard probability of fracture. The material is more sensitive to point defects on fibers.

### 5.7. Comparison of Fracture statistics for Poisson points on planes and

 for points on lines. For a given average number of defects in the sphere of radius $r$, we have$$
\theta_{2} \theta_{\pi}=\frac{1}{2} \theta_{1} \theta_{D}
$$

Taking $2 r \theta_{D}=\pi r^{2} \theta_{\pi}=\alpha$, we get $\pi^{2} \theta_{1} r^{2}=4 \pi r \theta_{2}$. Using the same intensity $\Phi(\sigma)=\theta=\theta_{\pi}=\theta_{D}$ for the two processes, we get

$$
\begin{aligned}
\log \left(P \left\{\sigma_{R}\right.\right. & \left.\geq \sigma\}_{\pi}\right)-\log \left(P\left\{\sigma_{R} \geq \sigma\right\}_{D}\right) \\
& =\pi^{2} \theta_{1} r^{2}\left(\psi(\theta \pi, r)-\frac{2}{\alpha^{2}}(1-(1+\alpha) \exp (-\alpha))\right)>0
\end{aligned}
$$

and therefore $P\left\{\sigma_{R} \geq \sigma\right\}_{\pi}>P\left\{\sigma_{R} \geq \sigma\right\}_{D}$ for any distribution $\Phi(\sigma)$ of defects and it is easier to break a specimen with defects on fibers than with defects on planes.
5.8. Fracture statistics for defects obtained in the three steps iteration. We consider now a model of a long fibers in random planes, with point defects located on the fibers, where $\theta$ is replaced by $\Phi(\sigma)$ in equation 4.2. We have

$$
\begin{aligned}
\log \left(P \left\{\sigma_{R} \geq\right.\right. & \left.\sigma\}_{3 \text { iterations }}\right)= \\
& \theta_{2} M(K)\left(E_{\pi}\left\{\exp \left[\theta_{1} \mathcal{L}(K \cap \pi)\left(\varphi_{L}(\Phi(\sigma), K \cap \pi)-1\right)\right]\right\}-1\right)
\end{aligned}
$$

For fracture statistics of the sphere with radius $r$,

$$
\begin{aligned}
\log \left(P \left\{\sigma_{R} \geq\right.\right. & \left.\sigma\}_{3 \text { iterations }}\right)= \\
& 4 \pi \theta_{2} r\left(\int_{0}^{r} \exp \left[2 \pi \theta_{1} u\left(\varphi_{L}(\Phi(\sigma), u)-1\right)\right] f(u, r) d u-1\right)
\end{aligned}
$$

5.9. Comparison of Fracture statistics for Poisson points and for the three steps iteration. We study the fracture statistics of a sphere of radius $r$ containing a random number of points $N_{P}(r)$ with a given average.

For the standard Poisson point process,

$$
E\left\{N_{P}(r)\right\}=\frac{4}{3} \pi r^{3} \theta_{3}
$$

For Poisson points on Poisson lines on Poisson planes,

$$
E\left\{N_{P}(r)\right\}=\frac{4}{3} \pi r^{3}\left(\theta_{2} \theta_{1} \theta 2 \pi^{2}\right)
$$

Given an average number of defects in the sphere of radius $r$, we have

$$
\theta_{3}=2 \pi^{2} \theta_{2} \theta_{1} \theta
$$

To compare fracture statistics of Poisson points and of the three iterations case, we use the ratio

$$
\frac{4}{3} \frac{\theta_{3} r^{3}}{4 \theta_{2} r}=\frac{4}{3} \frac{2 \pi^{2} \theta_{2} \theta_{1} \theta}{4 \theta_{2}} r^{2}=\frac{2}{3} \pi^{2} \theta_{1} \theta r^{2}
$$

With auxilliary variables $2 \theta r=\alpha$ and $\theta_{1} r=\beta$, we have to compare $\frac{1}{3} \pi^{2} \theta_{1} \alpha r=$ $\frac{1}{3} \pi^{2} \alpha \beta$ to

$$
1-\int_{0}^{r} \exp \left[-2 \pi \theta_{1} u\left(1-\varphi_{L}\left(\varphi_{L}(\theta, u)\right)\right] f(u, r) d r\right.
$$

Using $\frac{u}{r}=y$ and $d u=r d y$, we get

$$
\begin{aligned}
& 1-\int_{0}^{1} \exp \left[-2 \pi \theta_{1} r y\left(1-\varphi_{L}(\theta, r y)\right)\right] \frac{y}{\sqrt{1-y^{2}}} d y \\
= & 1-\int_{0}^{1} \exp \left[-2 \pi \beta y\left(1-\varphi_{L}(\theta, r y)\right)\right] \frac{y}{\sqrt{1-y^{2}}} d y
\end{aligned}
$$

We make a comparison by numerical calculation of the integral over $\alpha$, for a given $\beta$. For $\beta=0.01,0.1,1$ and 10 ,

$$
P\left\{\sigma_{R} \geq \sigma\right\}_{3 \text { iterations }}>P\left\{\sigma_{R} \geq \sigma\right\}_{P}
$$

5.10. Comparison of Fracture statistics for Poisson points on Poisson planes and for the three steps iteration. We fix the average number of points in the sphere with radius $r$. For Poisson points on Poisson planes

$$
E\left\{N_{\pi}(r)\right\}=\frac{8}{3} \pi^{2} r^{3} \theta_{2 \pi} \theta_{\pi}
$$

and for 3 iterations

$$
E\left\{N_{P}(r)\right\}_{3 \text { iterations }}=\frac{4}{3} \pi r^{3}\left(\theta_{2} \theta_{1} \theta 2 \pi^{2}\right)
$$

To keep the same average values, we fix

$$
2 \theta_{2 \pi} \theta_{\pi}=2 \pi \theta_{2} \theta_{1} \theta
$$

Taking $\theta_{\pi}=\theta$ to get the same statistics over points, and $\theta_{2}$ identical for the two models, in order to keep the same scale for the Poisson polyedra, we get $\pi \theta_{1}=1$ and $\theta_{1}=1 / \pi$.

With auxilliary variables $2 \theta r=\alpha$ and $\theta_{1} r=\beta$, we have to compare

$$
1-\psi(\theta \pi, r)
$$

and

$$
1-\int_{0}^{1} \exp \left[-2 \pi \beta y\left(1-\varphi_{L}(\theta, r y)\right)\right] \frac{y}{\sqrt{1-y^{2}}} d y
$$

For $\beta=0.01,0.1,0.5,0.75$, numerical calculations give

$$
P\left\{\sigma_{R} \geq \sigma\right\}_{3 \text { iterations }}>P\left\{\sigma_{R} \geq \sigma\right\}_{\pi}
$$

For $\beta=1$,

$$
\begin{aligned}
& P\left\{\sigma_{R} \geq \sigma\right\}_{3 \text { iterations }}<P\left\{\sigma_{R} \geq \sigma\right\}_{\pi} \text { when } \alpha<1.99 \\
& P\left\{\sigma_{R} \geq \sigma\right\}_{3 \text { iterations }}>P\left\{\sigma_{R} \geq \sigma\right\}_{\pi} \text { when } \alpha>1.99
\end{aligned}
$$

For $\beta=2,10$,

$$
P\left\{\sigma_{R} \geq \sigma\right\}_{3 \text { iterations }}<P\left\{\sigma_{R} \geq \sigma\right\}_{\pi}
$$

## 6. Conclusion

The models of random sets and point processes studied in this paper were designed to simulate some specific clustering of points, namely on random lines in $\mathbb{R}^{2}$ and $\mathbb{R}^{3}$ and on random planes in $\mathbb{R}^{3}$. A possible application is to model point defects in materials with some degree of alignment. We derived general theoretical results, useful to compare geometrical effects on the sensitivity of materials to fracture. Based on the presented theoretical results, applications can be looked for from statistical experimental data.
7. Appendix: Laplace transforms for sections of discs and of spheres
7.1. Random chords in a disc of radius $r$. Let $L$ the random length of a chord obtained by random sections of the disc with radius $r$ (this means chords given by intersections of the disc by lines with a uniform location along its diameter).
Proposition 9. The cumulative distribution of the random variable $L$ is given by

$$
P\{L<l\}=1-\sqrt{1-\left(\frac{l}{2 r}\right)^{2}}
$$

for $l<2 r$.

Proof. For any convex set $A$ with geometrical covariogram $K(h)$, the distribution $P\{L<l\}$ is given by the expression $1-\frac{K^{\prime}(h)}{K^{\prime}(0)}$, where $K^{\prime}(h)$ is also the projection of $A \cap A_{h}$ in the direction of vector $h$. For a disc with radius $r$ in $\mathbb{R}^{2}$ we get:

$$
\left(\frac{h}{2}\right)^{2}+\left(\frac{K^{\prime}(h)}{2}\right)^{2}=(r)^{2}
$$

and then

$$
(1-P\{L<l\})^{2}=\left(\frac{K^{\prime}(l)}{K^{\prime}(0)}\right)^{2}=1-\left(\frac{l}{2 r}\right)^{2}
$$

The Laplace transform of the distribution of random chords $L$ in a disc with radius $r$ is given by

$$
\varphi_{L}(\lambda, r)=\frac{1}{4 r^{2}} \int_{0}^{2 r} \frac{l \exp (-\lambda l)}{\sqrt{1-\left(\frac{l}{2 r}\right)^{2}}} d l
$$

Using $y=\frac{l}{2 r}$, so that $d l=2 r d y, \varphi_{L}(\lambda, r)$ can be derived using the formal computation software Mathematica

$$
\begin{align*}
\varphi_{L}(\lambda, r) & =\frac{1}{4 r^{2}} \int_{0}^{1} \exp (-\lambda 2 r y) \frac{2 r y 2 r}{\sqrt{1-y^{2}}} d y \\
& =\int_{0}^{1} \exp (-\lambda 2 r y) \frac{y d y}{\sqrt{1-y^{2}}} \\
& =\pi / 2[-\operatorname{Bessel}(1,2 \lambda r)+\operatorname{Struve} L(-1,2 \lambda r)] \tag{7.1}
\end{align*}
$$

Using a power series for $\exp (-\lambda 2 r y)=1+\sum_{i-1}^{i=\infty}(-1)^{n} \frac{(2 \lambda r)^{n}}{n!}$ we get

$$
\int_{0}^{1} \frac{y^{n} d y}{\sqrt{1-y^{2}}}=\frac{\sqrt{\pi}}{2} \frac{\Gamma\left(\frac{1+n}{2}\right)}{\Gamma\left(1+\frac{n}{2}\right)}
$$

and therefore

$$
\begin{align*}
\varphi_{L}(\lambda, r) & =1+\frac{\sqrt{\pi}}{2} \sum_{i=1}^{i=\infty}(-1)^{n} \frac{(2 \lambda r)^{n}}{n!} \frac{\Gamma\left(\frac{1+n+1}{2}\right)}{\Gamma\left(1+\frac{n+1}{2}\right)} \\
& =1+\frac{\sqrt{\pi}}{2} \sum_{i=1}^{i=\infty}(-1)^{n} \frac{(2 \lambda r)^{n}}{n!} \frac{\Gamma\left(\frac{n+2}{2}\right)}{\Gamma\left(\frac{n+3}{2}\right)} \tag{7.2}
\end{align*}
$$

7.2. Random radius of sections of a sphere with radius $R_{3}$. The cumulative distribution function of the radius $R$ of a random section of the sphere with radius $R_{3}$ is given by

$$
P\{R<r\}=1-\sqrt{1-\left(\frac{r}{R_{3}}\right)^{2}}
$$

The probability density function is given by

$$
\begin{equation*}
f\left(r, R_{3}\right)=\frac{r}{R_{3}^{2}} \frac{1}{\sqrt{1-\left(\frac{r}{R_{3}}\right)^{2}}} \tag{7.3}
\end{equation*}
$$

By using $y=\frac{r}{R_{3}}$ we have $R_{3} d y=d r$. The Laplace transform of the distribution of $R$ is given by

$$
\begin{aligned}
\varphi_{R}\left(\lambda, R_{3}\right) & =\int_{0}^{R_{3}} \frac{r}{R_{3}^{2}} \frac{1}{\sqrt{1-\left(\frac{r}{R_{3}}\right)^{2}}} \exp (-\lambda r) d r \\
& =\int_{0}^{1} \frac{R_{3} y}{R_{3}^{2}} \frac{1}{\sqrt{1-y^{2}}} R_{3} \exp \left(-\lambda R_{3} y\right) d y \\
& =\int_{0}^{1} \frac{y}{\sqrt{1-y^{2}}} \exp \left(-\lambda R_{3} y\right) d y \\
& =\pi / 2\left[-\operatorname{BesselI}\left(1, \lambda R_{3}\right)+\operatorname{StruveL}\left(-1, \lambda R_{3}\right)\right]
\end{aligned}
$$

From $\varphi_{R}\left(\lambda, R_{3}\right)$ is obtained the Laplace transform of the perimeter of sections, $\psi_{\mathcal{L}}\left(\lambda, B\left(R_{3}\right) \cap \pi\right)$, replacing $\lambda$ by $2 \pi \lambda$.

From the distribution of the radius $R$ can be computed the Laplace transform of $R^{2}$ :

$$
\begin{aligned}
E_{R}\left\{\exp \left(-\lambda R^{2}\left(R_{3}\right)\right)\right\} & =\psi\left(\lambda, R_{3}\right) \\
& =\int_{0}^{R_{3}} \exp \left(-\lambda u^{2}\right) \frac{u}{r^{2} \sqrt{1-\left(\frac{u}{R_{3}}\right)^{2}}} d u
\end{aligned}
$$

Using $\frac{u}{R_{3}}=y$ and $d u=R_{3} d y$,

$$
\begin{align*}
& E_{R}\left\{\exp \left(-\lambda R^{2}\left(R_{3}\right)\right)\right\} \\
= & \int_{0}^{1} \exp \left(-\lambda R_{3}^{2} y^{2}\right) \frac{R_{3} y R_{3} d y}{r^{2} \sqrt{1-y^{2}}}=\int_{0}^{1} \exp \left(-\lambda R_{3}^{2} y^{2}\right) \frac{y d y}{\sqrt{1-y^{2}}} \\
= & \frac{\exp \left(-\lambda R_{3}^{2}\right) \int_{0}^{R_{3} \sqrt{\lambda}} \exp \left(y^{2}\right) d y}{R_{3} \sqrt{\lambda}} \tag{7.4}
\end{align*}
$$

7.3. Random chords in a sphere with radius $R_{3}$. We start with the geometrical covariogram of the sphere with diameter $a=2 R_{3}$. For $h<a$,

$$
\begin{gathered}
K(h)=V\left(1-\frac{3}{2} \frac{h}{a}+\frac{1}{2}\left(\frac{h}{a}\right)^{3}\right) \\
K^{\prime}(h)=V\left(-\frac{3}{2} \frac{1}{a}+\frac{3}{2} \frac{h^{2}}{a^{3}}\right)
\end{gathered}
$$

The cumulative distribution function is obtained by:

$$
1-F(l)=\frac{K^{\prime}(l)}{K^{\prime}(0)}=1-\frac{l^{2}}{a^{2}}, \text { for } l<a
$$

and

$$
F(l)=\frac{l^{2}}{a^{2}}, \text { for } l<a
$$

with density $f(l)$

$$
f(l)=\frac{2 l}{a^{2}}, \text { for } l<a
$$

For a sphere with radius $R_{3}$

$$
f(l)=\frac{l}{2 R_{3}^{2}}, \text { for } l<2 R_{3}
$$

The Laplace transform of $f l$ ) is obtained from

$$
\varphi_{S L}\left(\lambda, R_{3}\right)=\frac{1}{2 R_{3}^{2}} \int_{0}^{2 R_{3}} l \exp (-\lambda l) d l
$$

Using $y=\frac{l}{2 R_{3}}$ with $d y=\frac{d l}{2 R_{3}}$

$$
\begin{aligned}
\varphi_{S L}\left(\lambda, R_{3}\right) & =\frac{1}{2 R_{3}^{2}} \int_{0}^{2 R_{3}} \exp (-\lambda l) 2 R_{3} y\left(2 R_{3}\right) d y \\
& =2 \int_{0}^{1} \exp \left(-\lambda 2 R_{3} y\right) y d y
\end{aligned}
$$

Since

$$
\begin{aligned}
& \int_{0}^{1} y \exp (-\alpha y) d y \\
= & \frac{1-(1+\alpha) \exp (-\alpha)}{\alpha^{2}}
\end{aligned}
$$

we get

$$
\varphi_{L}\left(\lambda, R_{3}\right)=\frac{2}{\left(2 R_{3} \lambda\right)^{2}}\left[1-\left(1+2 R_{3} \lambda\right) \exp \left(-2 R_{3} \lambda\right)\right]
$$

Using a power series expansion of $\psi(\lambda, r), \exp \left(-\lambda r^{2} y^{2}\right)=1+\sum_{1}^{\infty}(-1)^{n} \frac{\left(\lambda r^{2} y^{2}\right)^{n}}{n!}$. We get:

$$
\begin{aligned}
\int_{0}^{1} \frac{y^{2 n+1} d y}{\sqrt{1-y^{2}}} & =\frac{\sqrt{\pi}}{2} \frac{\Gamma\left(\frac{2 n+2}{2}\right)}{\Gamma\left(1+\frac{2 n+1}{2}\right)} \\
& =\frac{\sqrt{\pi}}{2} \frac{n!}{\Gamma\left(n+1+\frac{1}{2}\right)}
\end{aligned}
$$

and

$$
\begin{aligned}
\psi(\lambda, r) & =\int_{0}^{1} \exp \left(-\lambda r^{2} y^{2}\right) \frac{y d y}{\sqrt{1-y^{2}}} \\
& =1+\frac{\sqrt{\pi}}{2} \sum_{1}^{\infty}(-1)^{n} \frac{\left(\lambda r^{2}\right)^{n}}{\Gamma\left(n+1+\frac{1}{2}\right)}
\end{aligned}
$$

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