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TRAVELING WAVES FOR A LATTICE DYNAMICAL SYSTEM ARISING IN A DIFFUSIVE ENDEMIC MODEL

YAN-YU CHEN, JONG-SHENQ GUO, AND FRANÇOIS HAMEL

ABSTRACT. This paper is concerned with a lattice dynamical system modeling the evolution of susceptible and infective individuals at discrete niches. We prove the existence of traveling waves connecting the disease-free state to non-trivial leftover concentrations. We also characterize the minimal speed of traveling waves and we prove the non-existence of waves with smaller speeds.

1. INTRODUCTION

In this article, we consider the following lattice dynamical system (LDS)

$$(1.1) \quad \begin{cases} \frac{ds_n}{dt} = (s_{n+1} + s_{n-1} - 2s_n) + \mu - \mu s_n - \beta s_n i_n, & n \in \mathbb{Z}, \\ \frac{di_n}{dt} = d(i_{n+1} + i_{n-1} - 2i_n) - \mu i_n + \beta s_n i_n - \gamma i_n, & n \in \mathbb{Z}, \end{cases}$$

where $s_n = s_n(t)$, $i_n = i_n(t)$, $t \in \mathbb{R}$, and μ, β, γ are positive constants. Here $s_n(t)$ and $i_n(t)$ represent the population density of the susceptible individuals and the infective individuals at niches n at time t , 1 and d are the random migration coefficients for susceptible and infective population, respectively, and μ is regarded as the rate of the inflow of newborns into the susceptible population by assuming the total population of susceptible, infective and recovered individuals is normalized to be 1. The death rate of the susceptible population and the infective population are both assumed to be μ , β is the infective (transmission) coefficient and γ is the recovered/removed coefficient. Actually, as in [27], the equation for the recovered individuals $r_n(t)$ is given by

$$\frac{dr_n}{dt} = \gamma i_n - \mu r_n,$$

if there is no migration of the recovered individuals.

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For system (1.1), it is easy to see there are two constant states $(1, 0)$ and

$$(1.2) \quad (s^*, e^*) := \left(\frac{1}{\sigma}, \frac{\mu}{\beta} (\sigma - 1) \right),$$

where $\sigma := \beta/(\mu + \gamma)$. In this paper, we always assume that $\sigma > 1$, that is,

$$\beta > \mu + \gamma.$$

This means that, when the density of susceptible individuals is close to 1, the infective individuals have a positive per capita growth rate. Without migration, the steady state $(1, 0)$ is dynamically unstable with respect to perturbations whose second component are positive, while the steady state (s^*, e^*) is dynamically stable. System (1.1) is therefore called monostable.

In this paper, we are interested in the existence of traveling wave solutions of (1.1). Here a traveling wave solution is a bounded solution which can be expressed by

$$s_n(t) = \phi(n + ct) \quad \text{and} \quad i_n(t) = \psi(n + ct)$$

for $n \in \mathbb{Z}$ and $t \in \mathbb{R}$, for some nonnegative bounded functions ϕ, ψ on \mathbb{R} (the wave profiles) and some constant c (the wave speed). By setting $\xi = n + ct$ and substituting $(s_n(t), i_n(t)) = (\phi(\xi), \psi(\xi))$ into (1.1), we obtain

$$(1.3) \quad \begin{cases} -c\phi'(\xi) + D[\phi](\xi) + \mu(1 - \phi(\xi)) - \beta\phi(\xi)\psi(\xi) = 0, \\ -c\psi'(\xi) + dD[\psi](\xi) - (\mu + \gamma)\psi(\xi) + \beta\phi(\xi)\psi(\xi) = 0 \end{cases}$$

for all $\xi \in \mathbb{R}$, where

$$D[f](\xi) := f(\xi + 1) + f(\xi - 1) - 2f(\xi).$$

Furthermore, from the epidemic point of view, we are interested in traveling wave solutions connecting the trivial disease-free state $(1, 0)$ as $\xi \rightarrow -\infty$ (ahead of the front) and non-trivial states as $\xi \rightarrow +\infty$.

Define the constant c^* by

$$(1.4) \quad c^* := \inf_{\lambda > 0} \frac{d(e^\lambda + e^{-\lambda} - 2) + \beta - \mu - \gamma}{\lambda}.$$

By the assumption $\beta > \mu + \gamma$ (i.e. $\sigma > 1$), we know that $c^* = c^*(d, \beta, \mu, \gamma)$ is a well-defined real number (and the infimum in (1.4) is a minimum) and $c^* > 0$.

Our main result is the following theorem on the existence of traveling waves for (1.3) and the characterization of their minimal speed.

Theorem 1.1. *For any $c \geq c^*$, there exists a bounded classical solution (ϕ, ψ) of the system (1.3) such that*

$$(1.5) \quad 0 < \phi < 1 \text{ in } \mathbb{R}, \quad \psi > 0 \text{ in } \mathbb{R}$$

and

$$(1.6) \quad \lim_{\xi \rightarrow -\infty} (\phi(\xi), \psi(\xi)) = (1, 0),$$

together with

$$(1.7) \quad 0 < \liminf_{\xi \rightarrow +\infty} \phi(\xi) \leq s^* \leq \limsup_{\xi \rightarrow +\infty} \phi(\xi) < 1 \quad \text{and} \quad 0 < \liminf_{\xi \rightarrow +\infty} \psi(\xi) \leq e^* \leq \limsup_{\xi \rightarrow +\infty} \psi(\xi) < +\infty,$$

where (s^*, e^*) is given in (1.2). Furthermore, for any $c < c^*$, there is no classical solution (ϕ, ψ) of the system (1.3) satisfying (1.5) and (1.6).

Behind the front, as $\xi \rightarrow +\infty$, the leftover concentrations of susceptible and infective individuals are non-trivial. It is still an open question to know whether the traveling wave solutions converge to the endemic state (s^*, e^*) as $\xi \rightarrow +\infty$, but Theorem 1.1 asserts that both susceptible and infective individuals coexist behind the front and that the endemic state (s^*, e^*) is the only possible constant leftover state. To show the convergence to the endemic state as $\xi \rightarrow +\infty$, the difficulties come from the fact that (1.3) is a system and is non-local (such issues also arise for equations with non-local nonlinear interaction, see e.g. [1, 2, 4, 17, 18, 22, 20, 36, 37]). We also point out that the random migration coefficient d for the infective individuals is any arbitrary positive real number and is therefore in general different from that for the susceptible individuals. Furthermore, we mention that, due to the transmission term s with opposite signs, the systems (1.1) and (1.3) are not monotone (neither cooperative nor competitive) and therefore do not satisfy the maximum principle.

Let us finally mention some references on related problems. Actually, there is a vast literature on the study of traveling wave solutions for lattice dynamical systems or discrete versions of continuous parabolic partial differential equations. For monostable equations or monostable monotone systems, we refer to e.g. [6, 7, 8, 15, 16, 21, 23, 25, 29, 30, 34, 35, 38, 41]. Waves for bistable lattice dynamical systems have been studied in e.g. [5, 9, 10, 11, 12, 13, 14, 24, 26, 28, 31, 32, 33, 39, 40].

Remark 1.2. Notice that the necessity condition $c \geq c^*$ holds for any traveling wave (ϕ, ψ) satisfying (1.3), (1.5) and (1.6). The limiting conditions (1.7) or the boundedness of ψ are not used here.

Outline of the paper. Sections 2 and 3 are devoted to the proof of the existence of a traveling wave in case $c > c^*$, with some preliminaries on lower and upper solutions in Section 2. Approximated solutions in bounded domains are constructed and the traveling wave solving (1.3) is obtained by passing to the limit in the whole real line. Some intricate issues are to show that the limiting ψ component is bounded and that the leftover concentrations are non-trivial. Section 4 is devoted to the proof of the existence of a traveling wave for the minimal speed c^* , by passing to the limit $c_k \rightarrow (c^*)^+$, after a suitable shift of the origin and after showing that the solutions with speed c_k are uniformly bounded. Lastly, Section 5 is concerned with the proof of the necessity condition $c \geq c^*$ for any traveling wave satisfying (1.5) and (1.6).

2. PRELIMINARIES

In this section, we always assume that $c > c^*$. Then the equation

$$(2.1) \quad d(e^\lambda + e^{-\lambda} - 2) - c\lambda + \beta - \mu - \gamma = 0$$

has two positive roots λ_1 and λ_2 with $0 < \lambda_1 < \lambda_2$. Notice that

$$d(e^\lambda + e^{-\lambda} - 2) - c\lambda + \beta - \mu - \gamma < 0$$

for all $\lambda \in (\lambda_1, \lambda_2)$.

2.1. Upper and lower solutions. First, we define the notion of upper solution $(\bar{\phi}, \bar{\psi})$ and lower solution $(\underline{\phi}, \underline{\psi})$ of (1.3) as follows.

Definition 2.1. *If $\bar{\phi}, \bar{\psi}, \underline{\phi}, \underline{\psi}$ are continuous in \mathbb{R} , of class C^1 on $\mathbb{R} \setminus \mathcal{F}$ for some finite set \mathcal{F} and if they satisfy the following inequalities*

$$(2.2) \quad D[\bar{\phi}](\xi) - c\bar{\phi}'(\xi) + \mu(1 - \bar{\phi}(\xi)) - \beta\bar{\phi}(\xi)\underline{\psi}(\xi) \leq 0,$$

$$(2.3) \quad D[\underline{\phi}](\xi) - c\underline{\phi}'(\xi) + \mu(1 - \underline{\phi}(\xi)) - \beta\underline{\phi}(\xi)\bar{\psi}(\xi) \geq 0,$$

$$(2.4) \quad dD[\bar{\psi}](\xi) - c\bar{\psi}'(\xi) - (\mu + \gamma)\bar{\psi}(\xi) + \beta\bar{\phi}(\xi)\bar{\psi}(\xi) \leq 0,$$

$$(2.5) \quad dD[\underline{\psi}](\xi) - c\underline{\psi}'(\xi) - (\mu + \gamma)\underline{\psi}(\xi) + \beta\underline{\phi}(\xi)\underline{\psi}(\xi) \geq 0$$

for all $\xi \in \mathbb{R} \setminus \mathcal{F}$, then the functions $(\bar{\phi}, \bar{\psi}), (\underline{\phi}, \underline{\psi})$ are called a pair of upper and lower solutions of (1.3).

Following [3, 19], we introduce

$$(2.6) \quad \bar{\phi}(\xi) = 1, \quad \bar{\psi}(\xi) = e^{\lambda_1 \xi}, \quad \xi \in \mathbb{R},$$

$$(2.7) \quad \underline{\phi}(\xi) = \begin{cases} 1 - \rho e^{\theta \xi}, & \xi \leq \xi_1, \\ 0, & \xi \geq \xi_1, \end{cases}$$

$$(2.8) \quad \underline{\psi}(\xi) = \begin{cases} e^{\lambda_1 \xi} - q e^{\eta \lambda_1 \xi}, & \xi \leq \xi_2, \\ 0, & \xi \geq \xi_2, \end{cases}$$

where

$$(2.9) \quad \xi_1 := -\frac{\ln \rho}{\theta} \quad \text{and} \quad \xi_2 := -\frac{\ln q}{(\eta - 1)\lambda_1}.$$

Here the constants θ, ρ, η and q are chosen *in sequence* such that the following assumptions (A1)-(A4) hold:

$$(A1) \quad \theta > 0 \text{ is small enough such that } 0 < \theta < \lambda_1 \text{ and } e^\theta + e^{-\theta} - 2 - c\theta - \mu < 0,$$

$$(A2) \quad \rho > \max \left\{ 1, \frac{\beta}{-(e^\theta + e^{-\theta} - 2 - c\theta - \mu)} \right\} \geq 1 > 0,$$

$$(A3) \quad \eta \in (1, \min\{1 + \theta/\lambda_1, \lambda_2/\lambda_1\}) \text{ such that}$$

$$d(e^{\eta \lambda_1} + e^{-\eta \lambda_1} - 2) - c\eta \lambda_1 + \beta - \mu - \gamma < 0,$$

$$(A4) \quad q > \max \left\{ e^{(1-\eta)\lambda_1 \xi_1}, \frac{\beta \rho}{-(d(e^{\eta \lambda_1} + e^{-\eta \lambda_1} - 2) - c\eta \lambda_1 + \beta - \mu - \gamma)} \right\} > 0.$$

Note that we have

$$\xi_2 = -\frac{\ln q}{(\eta - 1)\lambda_1} < \xi_1 = -\frac{\ln \rho}{\theta} < 0.$$

Also, it easy to see that

$$\max\{0, 1 - \rho e^{\theta\xi}\} \leq \underline{\phi}(\xi) \leq \overline{\phi}(\xi) = 1, \quad \max\{0, e^{\lambda_1\xi} - q e^{\eta\lambda_1\xi}\} \leq \underline{\psi}(\xi) \leq \overline{\psi}(\xi) = e^{\lambda_1\xi}$$

for all $\xi \in \mathbb{R}$.

The next lemma gives the existence of a pair of upper and lower solutions.

Lemma 2.2. *The functions $(\overline{\phi}, \overline{\psi})$ and $(\underline{\phi}, \underline{\psi})$ defined by (2.6)-(2.8) are a pair of upper and lower solutions of (1.3).*

Proof. First, the functions $\overline{\phi}$ and $\overline{\psi}$ are of class $C^1(\mathbb{R})$ and the inequalities (2.2) and (2.4) hold on \mathbb{R} , since

$$\begin{cases} D[\overline{\phi}](\xi) - c\overline{\phi}'(\xi) + \mu(1 - \overline{\phi}(\xi)) - \beta\overline{\phi}(\xi)\underline{\psi}(\xi) = -\beta\underline{\psi}(\xi) \leq 0, \\ dD[\overline{\psi}](\xi) - c\overline{\psi}'(\xi) - (\mu + \gamma)\overline{\psi}(\xi) + \beta\overline{\phi}(\xi)\overline{\psi}(\xi) \\ = e^{\lambda_1\xi} [d(e^{\lambda_1} + e^{-\lambda_1} - 2) - c\lambda_1 + \beta - \mu - \gamma] = 0 \end{cases}$$

for all $\xi \in \mathbb{R}$.

Next, the function $\underline{\phi}$ is continuous in \mathbb{R} and of class $C^1(\mathbb{R} \setminus \{\xi_1\})$ and we would like to show that (2.3) holds for $\xi \neq \xi_1$. For $\xi > \xi_1$, this is trivial since $\underline{\phi}(\xi) = 0$. When $\xi < \xi_1 (< 0)$, we have $\underline{\phi}(\xi) = 1 - \rho e^{\theta\xi}$ and so

$$\begin{aligned} & D[\underline{\phi}](\xi) - c\underline{\phi}'(\xi) + \mu(1 - \underline{\phi}(\xi)) - \beta\underline{\phi}(\xi)\overline{\psi}(\xi) \\ &= 1 - \rho e^{\theta(\xi+1)} + 1 - \rho e^{\theta(\xi-1)} - 2 + 2\rho e^{\theta\xi} + c\theta\rho e^{\theta\xi} + \mu\rho e^{\theta\xi} - \beta e^{\lambda_1\xi} + \beta\rho e^{(\theta+\lambda_1)\xi} \\ &\geq e^{\theta\xi} [-\rho(e^\theta + e^\theta - 2c\theta - \mu) - \beta e^{(\lambda_1-\theta)\xi}] \\ &\geq \beta e^{\theta\xi} [1 - e^{(\lambda_1-\theta)\xi}] \geq 0 \end{aligned}$$

by $\theta < \lambda_1$ and the choice of ρ .

Finally, the function $\underline{\psi}$ is continuous in \mathbb{R} and of class $C^1(\mathbb{R} \setminus \{\xi_2\})$ and we claim that (2.5) holds for $\xi \neq \xi_2$. Clearly, (2.5) holds for $\xi > \xi_2$. For the case $\xi < \xi_2$, due to $\xi_2 < \xi_1 < 0$, we know that $\underline{\phi}(\xi) = 1 - \rho e^{\theta\xi}$ and $\underline{\psi}(\xi) = e^{\lambda_1\xi} - q e^{\eta\lambda_1\xi}$. Then we obtain

$$\begin{aligned} & dD[\underline{\psi}](\xi) - c\underline{\psi}'(\xi) - (\mu + \gamma)\underline{\psi}(\xi) + \beta\underline{\phi}(\xi)\underline{\psi}(\xi) \\ &\geq d[-q e^{\eta\lambda_1(\xi+1)} - q e^{\eta\lambda_1(\xi-1)} + 2q e^{\eta\lambda_1\xi}] + c\eta\lambda_1 q e^{\eta\lambda_1\xi} - (\beta - \mu - \gamma)q e^{\eta\lambda_1\xi} - \beta\rho e^{(\theta+\lambda_1)\xi} \\ &= e^{\eta\lambda_1\xi} \left\{ -q[d(e^{\eta\lambda_1} + e^{-\eta\lambda_1} - 2) - c\eta\lambda_1 + \beta - \mu - \gamma] - \beta\rho e^{[\theta+(1-\eta)\lambda_1]\xi} \right\} \\ &\geq \beta\rho e^{\eta\lambda_1\xi} (1 - e^{(\theta+(1-\eta)\lambda_1)\xi}) \geq 0 \end{aligned}$$

by the choices of η and q . Therefore, the proof of this lemma has been completed. \square

2.2. An auxiliary truncated problem. Now, given $l > -\xi_2 (> 0)$, we consider the following truncated problem

$$(2.10) \quad \begin{cases} D[\phi] - c\phi' + \mu(1 - \phi) - \beta\phi\psi = 0 & \text{in } [-l, l], \\ dD[\psi] - c\psi' - (\mu + \gamma)\psi + \beta\phi\psi = 0 & \text{in } [-l, l], \\ (\phi, \psi) = (\overline{\phi}, \overline{\psi}) & \text{on } (-\infty, -l), \\ (\phi, \psi) = (\phi(l), \psi(l)) & \text{on } (l, +\infty), \end{cases}$$

where

$$\begin{cases} \phi'(-l) := \lim_{h \searrow 0} \frac{\phi(-l+h) - \phi(-l)}{h}, & \psi'(-l) := \lim_{h \searrow 0} \frac{\psi(-l+h) - \psi(-l)}{h}, \\ \phi'(l) := \lim_{h \searrow 0} \frac{\phi(l) - \phi(l-h)}{h}, & \psi'(l) := \lim_{h \searrow 0} \frac{\psi(l) - \psi(l-h)}{h}. \end{cases}$$

Next, we give some notations. Set $\mathcal{C}^l := C([-l, l]) \times C([-l, l])$ and

$$\mathcal{S}^l := \{(\phi, \psi) \in \mathcal{C}^l \mid \underline{\phi} \leq \phi \leq \bar{\phi}, \underline{\psi} \leq \psi \leq \bar{\psi} \text{ in } [-l, l] \text{ and } (\phi, \psi)(-l) = (\bar{\phi}, \bar{\psi})(-l)\}.$$

From the definition of $\bar{\phi}, \underline{\phi}, \bar{\psi}, \underline{\psi}$, we know that $0 \leq \phi \leq 1$ and $0 \leq \psi \leq e^{\lambda_1 l}$ in $[-l, l]$ for any $(\phi, \psi) \in \mathcal{S}^l$. Hence \mathcal{S}^l is a nonempty bounded closed convex set in $(\mathcal{C}^l, \|\cdot\|)$, where $\|\cdot\|$ is the usual sup norm. For any $(\phi, \psi) \in \mathcal{S}^l$, we extend (ϕ, ψ) by continuity outside the interval $[-l, l]$ as in (2.10) and we introduce the continuous functions $H_1^l(\phi, \psi)$ and $H_2^l(\phi, \psi)$ defined in \mathbb{R} by

$$\begin{cases} H_1^l(\phi, \psi)(\xi) &= \alpha \phi(\xi) + D[\phi](\xi) + \mu(1 - \phi(\xi)) - \beta \phi(\xi) \psi(\xi), \\ H_2^l(\phi, \psi)(\xi) &= \alpha \psi(\xi) + d D[\psi](\xi) - (\mu + \gamma) \psi(\xi) + \beta \phi(\xi) \psi(\xi), \end{cases}$$

where $\alpha = \alpha^l$ is a positive constant such that

$$\alpha > \max \{2 + \mu + \beta e^{\lambda_1 l}, 2d + \mu + \gamma\}.$$

For $(\phi_i, \psi_i) \in \mathcal{S}^l$, $i = 1, 2$, with $\phi_1 \leq \phi_2$ and $\psi_1 \leq \psi_2$ in $[-l, l]$, we have

$$(2.11) \quad H_1^l(\phi_1, \psi_2)(\xi) \leq H_1^l(\phi_1, \psi_1)(\xi) \leq H_1^l(\phi_2, \psi_1)(\xi) \quad \text{and} \quad H_2^l(\phi_1, \psi_1)(\xi) \leq H_2^l(\phi_2, \psi_2)(\xi)$$

for all $\xi \in [-l, l]$. Finally, we define the operator $F^l = (F_1^l, F_2^l)$ from \mathcal{S}^l into \mathcal{C}^l as follows

$$\begin{cases} F_1^l(\phi, \psi)(\xi) &= e^{\alpha(-l-\xi)/c} \bar{\phi}(-l) + \int_{-l}^{\xi} \frac{e^{\alpha(z-\xi)/c}}{c} H_1^l(\phi, \psi)(z) dz, \quad \xi \in [-l, l], \\ F_2^l(\phi, \psi)(\xi) &= e^{\alpha(-l-\xi)/c} \bar{\psi}(-l) + \int_{-l}^{\xi} \frac{e^{\alpha(z-\xi)/c}}{c} H_2^l(\phi, \psi)(z) dz, \quad \xi \in [-l, l]. \end{cases}$$

Note that a fixed point (ϕ, ψ) of the operator F^l , extended outside the interval $[-l, l]$ as in (2.10), gives a solution of (2.10) which is continuous in \mathbb{R} and of class $C^1(\mathbb{R} \setminus \{-l, l\})$.

To show the existence of such a fixed point, we apply Schauder's fixed point theorem in the next lemma.

Lemma 2.3. *Given $l > -\xi_2$, there exists a $C(\mathbb{R}) \times C(\mathbb{R})$ and $C^1(\mathbb{R} \setminus \{-l, l\}) \times C^1(\mathbb{R} \setminus \{-l, l\})$ solution (ϕ, ψ) of (2.10) such that*

$$(2.12) \quad 0 \leq \underline{\phi} \leq \phi \leq 1 \quad \text{and} \quad 0 \leq \underline{\psi} \leq \psi \leq \bar{\psi} \quad \text{in } (-\infty, l].$$

Proof. First, we claim that $F^l(\mathcal{S}^l) \subset \mathcal{S}^l$. By (2.11), for any $(\phi, \psi) \in \mathcal{S}^l$, we have

$$F_1^l(\underline{\phi}, \bar{\psi}) \leq F_1^l(\phi, \psi) \leq F_1^l(\bar{\phi}, \underline{\psi}) \quad \text{and} \quad F_2^l(\underline{\phi}, \underline{\psi}) \leq F_2^l(\phi, \psi) \leq F_2^l(\bar{\phi}, \bar{\psi}) \quad \text{in } [-l, l].$$

By Lemma 2.2 and the definition of the upper and lower solutions, we also derive that

$$\underline{\phi} \leq F_1^l(\underline{\phi}, \bar{\psi}), \quad F_1^l(\bar{\phi}, \underline{\psi}) \leq \bar{\phi}, \quad \underline{\psi} \leq F_2^l(\underline{\phi}, \underline{\psi}) \quad \text{and} \quad F_2^l(\bar{\phi}, \bar{\psi}) \leq \bar{\psi} \quad \text{in } [-l, l].$$

Hence $F^l(\mathcal{S}^l) \subset \mathcal{S}^l$.

By using Arzela-Ascoli Theorem, the operator $F^l : \mathcal{S}^l \rightarrow \mathcal{S}^l$ is completely continuous with respect to the sup norm. With the help of Schauder's fixed point theorem, we conclude that there exists a pair $(\phi, \psi) \in \mathcal{S}^l$ such that $(\phi, \psi) = F^l(\phi, \psi)$. Therefore, (ϕ, ψ) , extended outside the interval $[-l, l]$ as in (2.10), solves (2.10) and satisfies the properties stated in Lemma 2.3. \square

3. EXISTENCE OF A TRAVELING WAVE FOR $c > c^*$

3.1. Proof of Theorem 1.1 for $c > c^*$. In this section, we show Theorem 1.1 for any fixed real number $c \in (c^*, +\infty)$. Namely, we show the existence of a bounded solution (ϕ, ψ) of (1.3) satisfying $0 < \phi < 1$ in \mathbb{R} , $\psi > 0$ in \mathbb{R} , and such that (1.6) and (1.7) hold.

First, we consider a positive increasing sequence $\{l_k\}_{k \in \mathbb{N}}$ such that $l_k \rightarrow \infty$ as $k \rightarrow \infty$, and $l_k > -\xi_2$ for all $k \in \mathbb{N}$, where $\xi_2 < 0$ is as in (2.9). By Lemma 2.3, for each $k \in \mathbb{N}$, there exists a $C(\mathbb{R}) \cap C^1(\mathbb{R} \setminus \{-l_k, l_k\})$ solution (ϕ_k, ψ_k) of (2.10) and (2.12) for $l = l_k$. For each $K \in \mathbb{N}$ such that $l_K \geq 2$, since $\bar{\psi}$ is bounded above in $[-l_K, l_K]$, it follows from (2.12) that the sequences

$$\{\phi_k\}_{k \geq K}, \quad \{\psi_k\}_{k \geq K}, \quad \{\phi_k \psi_k\}_{k \geq K}$$

are uniformly bounded on $[-l_K, l_K]$. Also, the sequences $\{\phi'_k\}_{k \geq K}$ and $\{\psi'_k\}_{k \geq K}$ are uniformly bounded in $[-l_K + 1, l_K - 1]$, due to (2.10) and (2.12). Since $\phi''_k(\xi)$ and $\psi''_k(\xi)$ can be expressed in terms of $\phi_k(\xi)$, $\psi_k(\xi)$, $\phi_k(\xi \pm 1)$, $\psi_k(\xi \pm 1)$, $\phi_k(\xi \pm 2)$, $\psi_k(\xi \pm 2)$, $\phi'_k(\xi)$ and $\psi'_k(\xi)$ in $[-l_K + 2, l_K - 2]$, one infers that the sequences $\{\phi''_k\}_{k \geq K}$ and $\{\psi''_k\}_{k \geq K}$ are uniformly bounded in $[-l_K + 2, l_K - 2]$. By using Arzela-Ascoli theorem on $[-l_K + 2, l_K - 2]$ for every $K \in \mathbb{N}$ large enough, we obtain a subsequence $\{(\phi_{k_j}, \psi_{k_j})\}$ of $\{(\phi_k, \psi_k)\}$ through the diagonal process such that

$$\phi_{k_j} \rightarrow \phi, \quad \psi_{k_j} \rightarrow \psi, \quad \phi'_{k_j} \rightarrow \phi', \quad \psi'_{k_j} \rightarrow \psi' \quad \text{as } j \rightarrow +\infty$$

uniformly in any compact subinterval of \mathbb{R} , for some functions $\phi \in C^1(\mathbb{R})$ and $\psi \in C^1(\mathbb{R})$. Then (ϕ, ψ) is a solution of the system (1.3) with

$$(3.1) \quad 0 \leq \underline{\phi} \leq \phi \leq 1 \quad \text{and} \quad 0 \leq \underline{\psi} \leq \psi \leq \bar{\psi} \quad \text{in } \mathbb{R}.$$

By the definitions of $\underline{\phi}$, $\bar{\psi}$ and $\underline{\psi}$, it easy to check that

$$(\phi, \psi)(-\infty) = (1, 0).$$

Notice also that, by differentiating the equations (1.3), one infers by induction that the functions ϕ and ψ are of class C^∞ in \mathbb{R} .

Lemma 3.1. *The functions ϕ and ψ are non-trivial, in the sense that*

$$0 < \phi < 1 \quad \text{and} \quad \psi > 0 \quad \text{in } \mathbb{R}.$$

Proof. Firstly, owing to the definition of $\underline{\psi}$, we have $\psi > 0$ in $(-\infty, \xi_2)$. For contradiction, we assume that there exists a real number $\bar{\xi}_0 \in [\xi_2, +\infty)$ such that $\psi(\bar{\xi}_0) = 0$ and $\psi(\xi) > 0$ for all $\xi < \bar{\xi}_0$. Since $\psi \geq 0$ in \mathbb{R} , we also have $\psi'(\bar{\xi}_0) = 0$. From the second equation of (1.3), we get that $\psi(\bar{\xi}_0 - 1) = \psi(\bar{\xi}_0 + 1) = 0$, a contradiction to the definition of $\bar{\xi}_0$.

Let us now show that $\phi > 0$ over \mathbb{R} . Indeed, if $\phi(\xi^*) = 0$ for some real number ξ^* , then

$$0 = -c \phi'(\xi^*) + D[\phi](\xi^*) + \mu (1 - \phi(\xi^*)) - \beta \phi(\xi^*) \psi(\xi^*) = -c \phi'(\xi^*) + D[\phi](\xi^*) + \mu > 0,$$

since $\phi'(\xi^*) = 0$, $D[\phi](\xi^*) \geq 0$ and $\mu > 0$. This contradiction leads to the inequality $\phi > 0$ in \mathbb{R} .

Similarly, we claim that $\phi < 1$ in \mathbb{R} by a contradiction argument. If there exists a real number $\tilde{\xi}$ such that $\phi(\tilde{\xi}) = 1$, then

$$0 = -c\phi'(\tilde{\xi}) + D[\phi](\tilde{\xi}) + \mu(1 - \phi(\tilde{\xi})) - \beta\phi(\tilde{\xi})\psi(\tilde{\xi}) = -c\phi'(\tilde{\xi}) + D[\phi](\tilde{\xi}) - \beta\psi(\tilde{\xi}) < 0,$$

since $\phi'(\tilde{\xi}) = 0$, $D[\phi](\tilde{\xi}) \leq 0$ and $\psi(\tilde{\xi}) > 0$. This contradiction leads to the inequality $\phi < 1$ in \mathbb{R} . \square

The next main step consists in showing that the function ψ is actually bounded. A first key-point is the following Harnack type property for equations of the type (1.3) satisfied by the second component ψ . We state this property in a more general framework.

Lemma 3.2. *Let M be a positive real number. Then there exists a constant $C = C(M) > 0$ such that, for any continuous functions a and b with $M^{-1} \leq a(\xi) \leq M$ and $b(\xi) \geq -M$ for all $\xi \in \mathbb{R}$ and for any positive $C^1(\mathbb{R})$ function u satisfying*

$$u'(\xi) \geq a(\xi)u(\xi + 1) + b(\xi)u(\xi) \quad \text{for all } \xi \in \mathbb{R},$$

there holds

$$C^{-1} \leq \frac{u(\xi + 1)}{u(\xi)} \leq C \quad \text{for all } \xi \in \mathbb{R}.$$

In order not to lengthen too much the main line of the proof of Theorem 1.1 with $c > c^*$, the proof of Lemma 3.2 is postponed in Section 3.2.

Coming back to our solutions (ϕ, ψ) of (1.3), since $c > 0$ and ϕ is nonnegative, it follows from Lemma 3.2 applied to the positive function $u = \psi$ solving $\psi'(\xi) \geq (d/c)\psi(\xi + 1) - (2d/c + \mu/c + \gamma/c)\psi(\xi)$ that the functions $\xi \mapsto \psi(\xi \pm 1)/\psi(\xi)$ are bounded in \mathbb{R} . Hence, from the equation (1.3) itself and since ϕ is bounded, the function

$$\xi \mapsto \frac{\psi'(\xi)}{\psi(\xi)}$$

is therefore bounded too.

The following two lemmas deal with the behavior of ϕ and ψ at $+\infty$ if $\limsup_{\xi \rightarrow +\infty} \psi(\xi) = +\infty$. The first one says that ϕ is small when ψ is large. This property actually holds locally uniformly with respect to the speed c . It is stated in this more general framework since it will be used again in Section 4 to get the existence of a bounded solution (ϕ, ψ) of (1.3) with speed c^* .

Lemma 3.3. *Let $0 < \underline{c} \leq \bar{c}$ be two given positive real numbers. Let $\{c_k\}$ be a sequence of real numbers in $[\underline{c}, \bar{c}]$ and let $\{(\phi_k, \psi_k)\}$ be a sequence of solutions of (1.3) with speed c_k and satisfying (1.5). If $\{\xi_k\}$ is a sequence of real numbers such that $\psi_k(\xi_k) \rightarrow +\infty$ as $k \rightarrow +\infty$, then $\phi_k(\xi_k) \rightarrow 0$ as $k \rightarrow +\infty$.*

Since this lemma is concerned with general sequences of solutions with different speeds, and in order not to lengthen too much the main line of the proof of Theorem 1.1 with given speed $c > c^*$, the proof of Lemma 3.3 is postponed in Section 3.2.

Coming back to our solution (ϕ, ψ) of (1.3) satisfying (1.5) and (1.6), the following result shows the convergence of ψ to $+\infty$ at $+\infty$ if it were not bounded.

Lemma 3.4. *If $\limsup_{\xi \rightarrow +\infty} \psi(\xi) = +\infty$, then $\lim_{\xi \rightarrow +\infty} \psi(\xi) = +\infty$.*

Proof. Assume by way of contradiction that $\limsup_{\xi \rightarrow +\infty} \psi(\xi) = +\infty$ and $\liminf_{\xi \rightarrow +\infty} \psi(\xi) < +\infty$. Since ψ'/ψ is globally bounded, there are then $M \in \mathbb{R}$ and two sequences $\{\theta_k\}$ and $\{\xi_k\}$ converging to $+\infty$ and such that

$$\psi(\theta_k) \leq M, \quad \theta_k < \xi_k - 1 < \xi_k < \xi_k + 1 < \theta_{k+1}, \quad \psi(\xi_k) = \max_{[\theta_k, \theta_{k+1}]} \psi \quad \left(= \max_{[\xi_k - 1, \xi_k + 1]} \psi \right)$$

for all $k \in \mathbb{N}$ and $\lim_{k \rightarrow +\infty} \psi(\xi_k) = +\infty$. Therefore, $\psi'(\xi_k) = 0$ and $dD[\psi](\xi_k) \leq 0$. Hence, by (1.3), one infers that $(\mu + \gamma - \beta \phi(\xi_k)) \psi(\xi_k) \leq 0$ for all $k \in \mathbb{N}$. This is clearly impossible for large k since $\psi(\xi_k) > 0$, and $\phi(\xi_k) \rightarrow 0$ as $k \rightarrow +\infty$ by Lemma 3.3. The proof is thereby complete. \square

To proceed further, we recall the following useful fundamental theory from [8] (or [6]) in dealing with the asymptotic tail behavior of wave profiles for a lattice dynamical system.

Proposition 3.5. [8] *Let $\varsigma > 0$ be a positive constant, let $B : \mathbb{R} \rightarrow \mathbb{R}$ be a continuous function having finite $B(\pm\infty) := \lim_{x \rightarrow \pm\infty} B(x)$ and let z be a continuous function such that*

$$(3.2) \quad \varsigma z(x) = e^{\int_x^{x+1} z(s) ds} + e^{\int_x^{x-1} z(s) ds} + B(x), \quad \forall x \in \mathbb{R}.$$

Then z is uniformly continuous and bounded in \mathbb{R} . In addition, the limits $\omega^\pm = \lim_{x \rightarrow \pm\infty} z(x)$ exist and are real roots of the characteristic equations

$$\varsigma \omega = e^\omega + e^{-\omega} + B(\pm\infty).$$

With this result and the previous lemmas in hand, we can show that ψ is bounded in \mathbb{R} .

Lemma 3.6. *The function ψ is bounded.*

Proof. Assume not. Then $\limsup_{\xi \rightarrow +\infty} \psi(\xi) = +\infty$, since ψ is continuous, positive, and $\psi(-\infty) = 0$. Therefore, Lemmas 3.3 and 3.4 imply that $\psi(\xi) \rightarrow +\infty$ and $\phi(\xi) \rightarrow 0$ as $\xi \rightarrow +\infty$. From (1.3), the continuous function $z := \psi'/\psi$ satisfies

$$\frac{c}{d} z(x) = e^{\int_x^{x+1} z(s) ds} + e^{\int_x^{x-1} z(s) ds} - 2 - \frac{\mu + \gamma}{d} + \frac{\beta \phi(x)}{d}$$

for all $x \in \mathbb{R}$. Since ϕ has finite limits at $\pm\infty$ and $\phi(+\infty) = 0$, it then follows from Proposition 3.5 that, in particular, z has a finite limit ω at $+\infty$, with

$$(3.3) \quad d(e^\omega + e^{-\omega} - 2) = c\omega + \mu + \gamma.$$

Since μ and γ are positive, this equation has a negative and a positive root. The function $z = \psi'/\psi$ cannot converge to the negative root at $+\infty$, since $\psi(+\infty) = +\infty$. Therefore, ψ'/ψ converges at $+\infty$ to the positive root ω of (3.3). Remember now that $\lambda_1 < \lambda_2$ are the two positive roots of equation (2.1). Since $\beta > 0$, one infers immediately that $\lambda_1 < \lambda_2 < \omega$. But $\lim_{\xi \rightarrow +\infty} \psi'(\xi)/\psi(\xi) = \omega > 0$ yields $\ln \psi(\xi) \sim \omega \xi$ as $\xi \rightarrow +\infty$, while (3.1) implies that $\psi(\xi) \leq \bar{\psi}(\xi) = e^{\lambda_1 \xi}$ for all $\xi \in \mathbb{R}$. One gets a contradiction, since $\lambda_1 < \omega$. As a conclusion, the function ψ is bounded and the proof of Lemma 3.6 is complete. \square

To complete the proof of Theorem 1.1 in case $c > c^*$, we show in the following lemmas that none of the components ϕ and ψ can be trivial at $+\infty$.

Lemma 3.7. *There holds $\inf_{\mathbb{R}} \phi > 0$.*

Proof. Remember that the C^∞ function ϕ satisfies $0 < \phi < 1$ in \mathbb{R} and $\phi(-\infty) = 1$. Assume by contradiction that $\inf_{\mathbb{R}} \phi = 0$. Then there exists a sequence $\{\xi_k\}$ converging to $+\infty$ such that $\phi(\xi_k) \rightarrow 0$ as $k \rightarrow +\infty$. On the other hand, since both functions ϕ and ψ are bounded, the equations (1.3) guarantee that the functions ϕ and ψ have bounded derivatives at any order. Therefore, by Arzela-Ascoli theorem, the functions $\xi \mapsto \phi(\xi + \xi_k)$ and $\xi \mapsto \psi(\xi + \xi_k)$ converge in $C_{loc}^\infty(\mathbb{R})$ as $k \rightarrow +\infty$, up to extraction of a subsequence, to some nonnegative C^∞ functions ϕ_∞ and ψ_∞ . Furthermore,

$$(3.4) \quad -c\phi'_\infty + D[\phi_\infty] + \mu(1 - \phi_\infty) - \beta\phi_\infty\psi_\infty = 0$$

in \mathbb{R} and $\phi_\infty(0) = 0$. Since 0 is a global minimum of ϕ_∞ , one has $\phi'_\infty(0) = 0$ and the above equality at 0 leads to a contradiction, since $\phi_\infty \geq 0$ and $\mu > 0$. Therefore, $\inf_{\mathbb{R}} \phi > 0$. \square

To show that ψ cannot approach 0 at $+\infty$, even for a sequence, the key-step is the following lemma saying that ψ is increasing when it is small. The property actually holds locally uniformly with respect to the speed c and we state the lemma in this slightly more general framework, since it will be used as such in Section 4.

Lemma 3.8. *Let $0 < \underline{c} \leq \bar{c}$ be two given positive real numbers. There is $\varepsilon > 0$ such that, for any $\Gamma \in [\underline{c}, \bar{c}]$ and for any solution (Φ, Ψ) of (1.3) (with speed Γ in place of c) satisfying (1.5), there holds*

$$\forall \xi \in \mathbb{R}, \quad (\Psi(\xi) \leq \varepsilon) \implies (\Psi'(\xi) > 0).$$

In order to conclude now the proof of Theorem 1.1 with $c > c^*$, the proof of Lemma 3.8 is postponed in Section 3.2. Coming back to our solution (ϕ, ψ) , we immediately get from Lemma 3.8 and the positivity of ψ in \mathbb{R} that

$$(3.5) \quad \liminf_{\xi \rightarrow +\infty} \psi(\xi) > 0.$$

We also claim that

$$(3.6) \quad \limsup_{\xi \rightarrow +\infty} \phi(\xi) < 1.$$

Indeed, otherwise, there exists a sequence of real numbers $\{\xi_k\}$ converging to $+\infty$ such that $\phi(\xi_k) \rightarrow 1$ as $k \rightarrow +\infty$. As in the proof of Lemma 3.7, up to extraction of a subsequence, the functions $\xi \mapsto \phi(\xi + \xi_k)$ and $\xi \mapsto \psi(\xi + \xi_k)$ converge as $k \rightarrow +\infty$ in $C_{loc}^\infty(\mathbb{R})$ to some nonnegative C^∞ functions ϕ_∞ and ψ_∞ solving (1.3). Furthermore, $0 < \phi_\infty \leq 1$ and $\psi_\infty > 0$ in \mathbb{R} from Lemma 3.7 and (3.5). Since $\phi_\infty(0) = 1$, one has $\phi'_\infty(0) = 0$. The equation (3.4) satisfied by ϕ_∞ at 0 leads to a contradiction, since $D[\phi_\infty](0) \leq 0$ and $-\beta\phi_\infty(0)\psi_\infty(0) = -\beta\psi_\infty(0) < 0$. Therefore, the claim (3.6) holds.

In order to complete the proof of (1.7), let us finally show that

$$(3.7) \quad \liminf_{\xi \rightarrow +\infty} \phi(\xi) \leq s^* \leq \limsup_{\xi \rightarrow +\infty} \phi(\xi) \quad \text{and} \quad \liminf_{\xi \rightarrow +\infty} \psi(\xi) \leq e^* \leq \limsup_{\xi \rightarrow +\infty} \psi(\xi).$$

Call $\phi_- = \liminf_{\xi \rightarrow +\infty} \phi(\xi)$, $\phi_+ = \limsup_{\xi \rightarrow +\infty} \phi(\xi)$, $\psi_- = \liminf_{\xi \rightarrow +\infty} \psi(\xi)$ and $\psi_+ = \limsup_{\xi \rightarrow +\infty} \psi(\xi)$. One already knows from (3.5), (3.6) and Lemmas 3.6 and 3.7 that

$$0 < \phi_- \leq \phi_+ < 1 \quad \text{and} \quad 0 < \psi_- \leq \psi_+ < +\infty.$$

Consider now a sequence $\{\xi_k\}$ converging to $+\infty$ such that $\psi(\xi_k) \rightarrow \psi_-$ as $k \rightarrow +\infty$. Up to extraction of a subsequence (as for instance in the proof of Lemma 3.7), the functions $\xi \mapsto \phi(\xi + \xi_k)$ and $\xi \mapsto \psi(\xi + \xi_k)$ converge in $C_{loc}^\infty(\mathbb{R})$ to some bounded functions $0 < \phi_\infty < 1$ and $\psi_\infty > 0$ satisfying (1.3). Furthermore, $0 < \psi_- = \psi_\infty(0) = \min_{\mathbb{R}} \psi_\infty$. Therefore, $\psi'_\infty(0) = 0$ and $D[\psi_\infty](0) \geq 0$. Hence

$$-(\mu + \gamma) \psi_- + \beta \phi_\infty(0) \psi_- \leq 0,$$

that is, $\beta \phi_\infty(0) \leq \mu + \gamma$. This yields $\phi_- = \liminf_{\xi \rightarrow +\infty} \phi(\xi) \leq (\mu + \gamma)/\beta = 1/\sigma = s^*$. Similarly, it follows that $\phi_+ = \limsup_{\xi \rightarrow +\infty} \phi(\xi) \geq s^*$. Consider also a sequence $\{\zeta_k\}$ converging to $+\infty$ such that $\phi(\xi_k) \rightarrow \phi_-$ as $k \rightarrow +\infty$. As above, up to extraction of a subsequence, the functions $\xi \mapsto \phi(\xi + \zeta_k)$ and $\xi \mapsto \psi(\xi + \zeta_k)$ converge in $C_{loc}^\infty(\mathbb{R})$ to some bounded functions $0 < \Phi_\infty < 1$ and $\Psi_\infty > 0$ satisfying (1.3). Furthermore, $0 < \phi_- = \Phi_\infty(0) = \min_{\mathbb{R}} \Phi_\infty$. Therefore, $\Phi'_\infty(0) = 0$ and $D[\Phi_\infty](0) \geq 0$. Hence

$$\mu(1 - \phi_-) - \beta \phi_- \Psi_\infty(0) \leq 0.$$

Since $0 < \phi_- \leq s^* = 1/\sigma$, one gets immediately that $\Psi_\infty(0) \geq (\mu/\beta)(\sigma - 1) = e^*$, whence $\psi_+ = \limsup_{\xi \rightarrow +\infty} \psi(\xi) \geq e^*$. Similarly, it follows that $\psi_- = \liminf_{\xi \rightarrow +\infty} \psi(\xi) \leq e^*$.

As a conclusion, (1.7) is proved and the proof of Theorem 1.1 in case $c > c^*$ is thereby complete.

As explained after the statement of Theorem 1.1 in Section 1, the question of the existence of a limit of (ϕ, ψ) at $+\infty$ is unclear. However, we can say that the a priori existence of a limit of one of these two functions guarantees the convergence of both, and that the endemic state (s^*, e^*) defined in (1.2) is the only possible limit.

Lemma 3.9. *Let (ϕ, ψ) be a bounded classical solution of (1.3) satisfying (1.5), (1.6) and (1.7), with speed $c \geq c^*$. If $\phi(+\infty)$ or $\psi(+\infty)$ exists, then they both exist and*

$$(\phi(+\infty), \psi(+\infty)) = (s^*, e^*).$$

Proof. Assume first that $l = \lim_{\xi \rightarrow +\infty} \phi(\xi)$ exists. Property (1.7) yields $0 < l = s^* < 1$. Consider now any sequence $\{\xi_k\}$ converging to $+\infty$. Up to extraction of a subsequence, the functions $\xi \mapsto \phi(\xi + \xi_k)$ and $\xi \mapsto \psi(\xi + \xi_k)$ converge in $C_{loc}^\infty(\mathbb{R})$ to some functions $\phi_\infty = l = s^*$ and ψ_∞ such that

$$\mu(1 - s^*) - \beta s^* \psi_\infty(\xi) = 0 \quad \text{for all } \xi \in \mathbb{R}.$$

Therefore, the function ψ_∞ is identically equal to the constant $\mu(1 - s^*)/(\beta s^*) = e^*$. Since the limit does not depend on the sequence $\{\xi_k\}$, one infers that $\lim_{\xi \rightarrow +\infty} \psi(\xi) = e^*$.

Conversely, if $L = \lim_{\xi \rightarrow +\infty} \psi(\xi)$ exists, property (1.7) yields $0 < L = e^*$. For any sequence $\{\xi_k\}$ converging to $+\infty$, the functions $\xi \mapsto \phi(\xi + \xi_k)$ and $\xi \mapsto \psi(\xi + \xi_k)$ converge in $C_{loc}^\infty(\mathbb{R})$,

up to a subsequence, to some functions ϕ_∞ and $\psi_\infty = L = e^*$ such that

$$-(\mu + \gamma) e^* + \beta e^* \phi_\infty(\xi) = 0 \quad \text{for all } \xi \in \mathbb{R}.$$

Therefore, the function ϕ_∞ is identically equal to the constant $(\mu + \gamma)/\beta = s^*$. Since the limit does not depend on the sequence $\{\xi_k\}$, one infers that $\lim_{\xi \rightarrow +\infty} \phi(\xi) = s^*$.

Therefore, if the limit $l = \phi(+\infty)$ or the limit $L = \psi(+\infty)$ exists, then they both exist such that $(\phi(+\infty), \psi(+\infty)) = (s^*, e^*)$. \square

Remark 3.10. The condition $c \geq c^*$ in Lemma 3.9 is not a restriction, since we shall prove in Section 5 that, for any solution (ϕ, ψ) of (1.3) satisfying (1.5) and (1.6) with speed c , there holds $c \geq c^*$.

3.2. Proof of Lemmas 3.2, 3.3 and 3.8. In this section, we prove some technical lemmas stated in Section 3.1.

Proof of Lemma 3.2. Although the idea of the proof is similar to the one given in [8], we provide the details here for completeness. Up to multiplication of u by a positive constant and up to a shift in space, one can assume without loss of generality that $u(0) = 1$ and it is sufficient to show that $u(\pm 1) \leq C = C(M)$. Firstly, since $u'(\xi) \geq -M u(\xi)$ for all $\xi \in \mathbb{R}$, the function $\xi \mapsto v(\xi) := u(\xi) e^{M\xi}$ is nondecreasing, hence

$$u(-1) \leq e^M u(0) = e^M.$$

Secondly, for all $\xi \in [0, 1]$, one has

$$v'(\xi) = (u'(\xi) + M u(\xi)) e^{M\xi} \geq a(\xi) u(\xi + 1) e^{M\xi} \geq \frac{v(\xi + 1) e^{-M}}{M} \geq \frac{v(1) e^{-M}}{M} = \frac{u(1)}{M}.$$

Hence, $v(\xi) \geq v(0) + u(1) \xi/M = 1 + u(1) \xi/M$ for all $\xi \in [0, 1]$. In other words,

$$u(\xi) \geq \left(1 + \frac{u(1) \xi}{M}\right) e^{-M\xi} \quad \text{for all } \xi \in [0, 1].$$

Finally, for all $\xi \in [-1/2, 0]$,

$$v'(\xi) \geq a(\xi) u(\xi + 1) e^{M\xi} \geq \frac{e^{M\xi}}{M} \times \left(1 + \frac{u(1) (\xi + 1)}{M}\right) e^{-M(\xi+1)} \geq \frac{e^{-M}}{M} \times \left(1 + \frac{u(1)}{2M}\right).$$

Therefore,

$$1 = v(0) \geq \underbrace{v(-1/2)}_{\geq 0} + \frac{e^{-M}}{2M} \times \left(1 + \frac{u(1)}{2M}\right) \geq \frac{e^{-M}}{2M} \times \left(1 + \frac{u(1)}{2M}\right).$$

Hence

$$u(1) \leq 2M(2Me^M - 1)$$

and the proof of Lemma 3.2 is thereby complete with $C(M) = \max\{e^M, 2M(2Me^M - 1)\}$. \square

Proof of Lemma 3.3. Let $0 < \underline{c} \leq \bar{c}$, $\{c_k\}$, $\{(\phi_k, \psi_k)\}$ and $\{\xi_k\}$ be as in the statement and assume by way of contradiction that there are $\varepsilon > 0$ and a subsequence, still denoted with the same index k , such that $\psi_k(\xi_k) \rightarrow +\infty$ as $k \rightarrow +\infty$ and $\phi_k(\xi_k) \geq \varepsilon$ for all $k \in \mathbb{N}$.

Since $0 < \phi_k < 1$ and $\psi_k > 0$ in \mathbb{R} , the equation (1.3) for ϕ_k (with $c_k \in [\underline{c}, \bar{c}]$) implies that $\phi'_k \leq (2 + \mu)/\underline{c}$ in \mathbb{R} . Hence

$$(3.8) \quad \phi_k(\xi) \geq \frac{\varepsilon}{2} \quad \text{for all } \xi \in [\xi_k - \delta, \xi_k] \text{ and for all } k \in \mathbb{N},$$

where $\delta = \varepsilon \underline{c}/(4 + 2\mu) > 0$. On the other hand, since

$$\psi'_k(\xi) \geq \frac{d}{\bar{c}} \psi_k(\xi + 1) - \frac{2d + \mu + \gamma}{\underline{c}} \psi_k(\xi) \quad \text{for all } \xi \in \mathbb{R} \text{ and for all } k \in \mathbb{N},$$

Lemma 3.2 applied to the positive functions ψ_k implies that the functions $\xi \mapsto \psi_k(\xi \pm 1)/\psi_k(\xi)$ are globally bounded independently of $k \in \mathbb{N}$. Hence, the functions ψ'_k/ψ_k are globally bounded in \mathbb{R} independently of $k \in \mathbb{N}$. Therefore, the limit $\lim_{k \rightarrow +\infty} \psi_k(\xi_k) = +\infty$ implies that $0 < M_k := \min_{[\xi_k - \delta, \xi_k]} \psi_k \rightarrow +\infty$ as $k \rightarrow +\infty$. Now, equation (1.3) and the inequalities $0 < \phi_k < 1$ and (3.8) yield

$$\max_{[\xi_k - \delta, \xi_k]} \phi'_k \leq \frac{2 + \mu}{\underline{c}} - \frac{\beta \varepsilon M_k}{2\bar{c}} \rightarrow -\infty \quad \text{as } k \rightarrow +\infty.$$

This contradicts the global boundedness of the functions ϕ_k . The proof of Lemma 3.3 is thereby complete. \square

Proof of Lemma 3.8. Assume by way of contradiction that there is no such ε . Then there exist a sequence of real numbers $\{c_k\}$ in $[\underline{c}, \bar{c}]$, a sequence of solutions $\{(\phi_k, \psi_k)\}$ of (1.3) with speed $c = c_k$ and $0 < \phi_k < 1$, $\psi_k > 0$ in \mathbb{R} , and a sequence of real numbers $\{\xi_k\}$ such that

$$(3.9) \quad \psi_k(\xi_k) \rightarrow 0 \text{ as } k \rightarrow +\infty \quad \text{and} \quad \psi'_k(\xi_k) \leq 0 \text{ for all } k \in \mathbb{N}.$$

Up to a shift of the origin, one can assume without loss of generality that

$$(3.10) \quad \xi_k = 0$$

for all $k \in \mathbb{N}$. Up to extraction of a subsequence, one can also assume that $c_k \rightarrow c_\infty \in [\underline{c}, \bar{c}]$ as $k \rightarrow +\infty$.

Notice first that Lemma 3.2 and the equations (1.3) satisfied by (ϕ_k, ψ_k) with $c_k \in [\underline{c}, \bar{c}] \subset (0, +\infty)$ imply that the sequence $\{\psi'_k/\psi_k\}$ is bounded in $L^\infty(\mathbb{R})$, that is, there is $C > 0$ such that $|\psi'_k(\xi)| \leq C \psi_k(\xi)$ for all $k \in \mathbb{N}$ and $\xi \in \mathbb{R}$. Since $\psi_k(0) \rightarrow 0^+$ as $k \rightarrow +\infty$, it follows that

$$\psi_k \rightarrow 0 \text{ locally uniformly in } \mathbb{R} \text{ as } k \rightarrow +\infty.$$

As a consequence, there also holds that $\psi'_k \rightarrow 0$ locally uniformly in \mathbb{R} as $k \rightarrow +\infty$.

Furthermore, by differentiating the equation (1.3) satisfied by ϕ_k , one gets that the functions ϕ'_k and ϕ''_k are locally bounded (and the functions ϕ_k are globally bounded). Therefore, the functions ϕ_k converge in $C^1_{loc}(\mathbb{R})$, up to extraction of a subsequence, to a function $0 \leq \phi_\infty \leq 1$ solving (1.3) with speed c_∞ and with $\psi = 0$, that is,

$$(3.11) \quad c_\infty \phi'_\infty = D[\phi_\infty] + \mu(1 - \phi_\infty) \quad \text{in } \mathbb{R}.$$

Call $\alpha = \inf_{\mathbb{R}} \phi_\infty$ and let $\{\zeta_m\}$ be sequence of real numbers such that $\phi_\infty(\zeta_m) \rightarrow \alpha$ as $m \rightarrow +\infty$. Up to extraction of a subsequence, the functions $\xi \mapsto \phi_\infty(\xi + \zeta_m)$ converge as $m \rightarrow +\infty$ in $C^\infty_{loc}(\mathbb{R})$ to a function Φ_∞ solving $c_\infty \Phi'_\infty = D[\Phi_\infty] + \mu(1 - \Phi_\infty)$ in \mathbb{R} ,

$\alpha \leq \Phi_\infty \leq 1$ in \mathbb{R} and $\Phi_\infty(0) = \alpha$. Consequently, $\Phi'_\infty(0) = 0$ and $D[\Phi_\infty](0) \geq 0$, whence $\mu(1 - \alpha) = \mu(1 - \Phi_\infty(0)) \leq 0$. Thus, $\alpha \geq 1$. Since $\alpha = \inf_{\mathbb{R}} \phi_\infty$ and $\phi_\infty \leq 1$ in \mathbb{R} , one concludes that

$$\phi_\infty = 1 \quad \text{in } \mathbb{R}.$$

Now set

$$\Psi_k(\xi) = \frac{\psi_k(\xi)}{\psi_k(0)}$$

for $k \in \mathbb{N}$ and $\xi \in \mathbb{R}$. Since the sequence $\{\psi'_k/\psi_k\}$ is bounded in $L^\infty(\mathbb{R})$, the positive functions Ψ_k are locally bounded, in the sense that $\sup_{k \in \mathbb{N}, |\xi| \leq R} \Psi_k(\xi) < +\infty$ for all $R > 0$. Therefore, the functions

$$\Psi'_k(\xi) = \frac{\psi'_k(\xi)}{\psi_k(0)} = \frac{\psi'_k(\xi)}{\psi_k(\xi)} \times \Psi_k(\xi)$$

are locally bounded too. Since each Ψ_k satisfies

$$-c_k \Psi''_k(\xi) + d D[\Psi'_k](\xi) - (\mu + \gamma) \Psi'_k(\xi) + \beta \phi'_k(\xi) \Psi_k(\xi) + \beta \phi_k(\xi) \Psi'_k(\xi) = 0$$

in \mathbb{R} and the sequence $\{\phi_k\}$ is bounded in $C^1_{loc}(\mathbb{R})$, one infers that the functions Ψ''_k are locally bounded too. By Arzela-Ascoli theorem, it follows that, up to extraction of a subsequence, the positive functions Ψ_k converge in $C^1_{loc}(\mathbb{R})$ to a nonnegative solution Ψ_∞ of

$$(3.12) \quad c_\infty \Psi'_\infty = d D[\Psi_\infty] + (\beta - \mu - \gamma) \Psi_\infty \quad \text{in } \mathbb{R},$$

where one used the fact that $\phi_k(\xi) \rightarrow \phi_\infty(\xi) = 1$ as $k \rightarrow +\infty$ for all $\xi \in \mathbb{R}$. Furthermore, we claim that $\Psi_\infty > 0$ in \mathbb{R} . Otherwise, there is $\xi_0 \in \mathbb{R}$ such that $\Psi_\infty(\xi_0) = 0$, and $\Psi'_\infty(\xi_0) = 0$. It follows from (3.12) applied at ξ_0 that $\Psi_\infty(\xi_0 + 1) = \Psi_\infty(\xi_0 - 1) = 0$, and then $\Psi_\infty(\xi_0 + m) = 0$ for all $m \in \mathbb{Z}$ by immediate induction. Since $c_\infty \Psi'_\infty \geq (\beta - \mu - \gamma - 2) \Psi_\infty$ in \mathbb{R} , the nonnegative function $\xi \mapsto \Psi_\infty(\xi) e^{-(\beta - \mu - \gamma - 2)\xi/c_\infty}$ is nondecreasing. Since it vanishes at $\xi_0 + m$ for all $m \in \mathbb{Z}$, one concludes that it is identically equal to 0, whence $\Psi_\infty = 0$ in \mathbb{R} . This contradicts the fact that $\Psi_\infty(0) = 1$. Therefore,

$$\Psi_\infty(\xi) > 0$$

for all $\xi \in \mathbb{R}$.

The continuous function $z := \Psi'_\infty/\Psi_\infty$ obeys

$$(3.13) \quad \frac{c_\infty}{d} z(\xi) = e^{\int_\xi^{\xi+1} z(s) ds} + e^{\int_\xi^{\xi-1} z(s) ds} - 2 + \frac{\beta - \mu - \gamma}{d} \quad \text{in } \mathbb{R}.$$

Therefore, by Proposition 3.5, $z(\xi) = \Psi'_\infty(\xi)/\Psi_\infty(\xi)$ has finite limits ω_\pm as $\xi \rightarrow \pm\infty$, which are roots of the characteristic equation

$$c_\infty \omega_\pm = d(e^{\omega_\pm} + e^{-\omega_\pm} - 2) + \beta - \mu - \gamma.$$

Since $c_\infty \geq \underline{c} > 0$ and $\beta > \mu + \gamma$, the roots of the previous equation are necessarily positive. In particular, Ψ'_∞ is positive at $\pm\infty$. Furthermore, by differentiating (3.13), one gets that

$$(3.14) \quad c_\infty z'(\xi) = d(z(\xi + 1) - z(\xi)) \frac{\Psi_\infty(\xi + 1)}{\Psi_\infty(\xi)} + d(z(\xi - 1) - z(\xi)) \frac{\Psi_\infty(\xi - 1)}{\Psi_\infty(\xi)} \quad \text{in } \mathbb{R}.$$

Therefore, if z has a minimum $\underline{\xi}$ in \mathbb{R} , then $z'(\underline{\xi}) = 0$ and $z(\underline{\xi} + 1) = z(\underline{\xi} - 1) = z(\underline{\xi})$, whence $z(\underline{\xi} + m) = z(\underline{\xi})$ for all $m \in \mathbb{Z}$ by immediate induction. As a consequence,

$$\inf_{\mathbb{R}} z \geq \min\{z(-\infty), z(+\infty)\} > 0.$$

Finally, $\Psi'_\infty > 0$ in \mathbb{R} , hence $0 < \Psi'_\infty(0) = \lim_{k \rightarrow +\infty} \Psi'_k(0) = \lim_{k \rightarrow +\infty} \psi'_k(0)/\psi_k(0)$ and $\psi'_k(0) > 0$ for all k large enough. This contradicts the fact that $\psi'_k(0) \leq 0$ for all $k \in \mathbb{N}$ (remember (3.9) and (3.10)).

As a conclusion, there is $\varepsilon > 0$ such that $\psi'(\xi) > 0$ for any $\xi \in \mathbb{R}$ with $\psi(\xi) \leq \varepsilon$ for any solution (ϕ, ψ) of (1.3) with $c \in [\underline{c}, \bar{c}]$, $0 < \phi < 1$ and $\psi > 0$ in \mathbb{R} . The proof of Lemma 3.8 is thereby complete. \square

4. THE CASE $c = c^*$

This section is devoted to the proof of the existence of a traveling wave (ϕ, ψ) of (1.3) satisfying (1.5), (1.6) and (1.7) with speed $c = c^*$. To do so, we consider a sequence $\{c_k\}$ of real numbers such that $c_k \in (c^*, c^* + 1]$ for each $k \in \mathbb{N}$, and

$$c_k \rightarrow c^* \quad \text{as } k \rightarrow +\infty.$$

For each $k \in \mathbb{N}$, Section 3 provides the existence of a traveling wave (ϕ_k, ψ_k) of (1.3) (with speed c_k) satisfying (1.5), (1.6) and (1.7). The natural strategy is to pass to the limit as $k \rightarrow +\infty$, in order to get the existence of a traveling wave with the limiting speed c^* . To achieve this goal, we need some a priori bounds for the functions ψ_k in order to get a non-trivial solution at the limit. We also point out that the inequalities (3.1) satisfied by the approximated waves (ϕ_k, ψ_k) do not carry over at the limit $c_k \rightarrow c^*$ (since the coefficients in the definitions of the lower solutions depend on c_k and degenerate at the limit $c_k \rightarrow c^*$). Therefore, we will have to suitably shift and renormalize the approximated waves (ϕ_k, ψ_k) before passing to the limit as $k \rightarrow +\infty$.

The first a priori bound asserts that the functions ψ_k do not converge to 0 uniformly as $k \rightarrow +\infty$.

Lemma 4.1. *There holds $\liminf_{k \rightarrow +\infty} \|\psi_k\|_{L^\infty(\mathbb{R})} > 0$.*

Proof. Assume that the conclusion does not hold. Then, up to extraction of a subsequence, one can assume without loss of generality that $\|\psi_k\|_{L^\infty(\mathbb{R})} \rightarrow 0$ as $k \rightarrow +\infty$. Since $c_k \in [c^*, c^* + 1] \subset (0, +\infty)$ for each $k \in \mathbb{N}$, Lemma 3.8 implies that $\psi'_k > 0$ in \mathbb{R} for all k large enough. Since each ψ_k is bounded, it follows that the limit $\psi_k(+\infty)$ exists in \mathbb{R} , for all k large enough. Since each (ϕ_k, ψ_k) satisfies the assumptions of Lemma 3.9, one then infers in particular that, for all k large enough,

$$\psi_k(+\infty) = e^* = \frac{\mu}{\beta} (\sigma - 1) > 0.$$

This contradicts the fact that $\lim_{k \rightarrow +\infty} \|\psi_k\|_{L^\infty(\mathbb{R})} = 0$. Thus, the conclusion of Lemma 4.1 holds. \square

The second key-point is the boundedness of the sequence $\{\psi_k\}$ in $L^\infty(\mathbb{R})$.

Lemma 4.2. *There holds $\limsup_{k \rightarrow +\infty} \|\psi_k\|_{L^\infty(\mathbb{R})} < +\infty$.*

Proof. Assume that the conclusion does not hold. Then, up to extraction of a subsequence, one has $\|\psi_k\|_{L^\infty(\mathbb{R})} \rightarrow +\infty$ as $k \rightarrow +\infty$. For each $k \in \mathbb{N}$, since the function ψ_k is bounded and positive in \mathbb{R} , there is then $\xi_k \in \mathbb{R}$ such that

$$(4.1) \quad \psi_k(\xi_k) \geq \left(1 - \frac{1}{k+1}\right) \|\psi_k\|_{L^\infty(\mathbb{R})}.$$

In particular, $\psi_k(\xi_k) \rightarrow +\infty$ as $k \rightarrow +\infty$. Furthermore, one has

$$\psi'_k(\xi) \geq \frac{d}{c^*} \psi_k(\xi+1) - \frac{2d + \mu + \gamma}{c^*} \psi_k(\xi) \quad \text{in } \mathbb{R}$$

for all $k \in \mathbb{N}$. Since each ψ_k is positive, it follows from Lemma 3.2 that the functions $\xi \mapsto \psi_k(\xi \pm 1)/\psi_k(\xi)$ are globally bounded in \mathbb{R} independently of $k \in \mathbb{N}$, and so are the functions $\xi \mapsto \psi'_k(\xi)/\psi_k(\xi)$, from the equation (1.3) satisfied with speed $c_k \in (c^*, c^* + 1]$ (remember also that $0 < \phi_k < 1$ in \mathbb{R}). As a consequence,

$$\psi_k(\xi + \xi_k) \xrightarrow[k \rightarrow +\infty]{} +\infty \quad \text{locally uniformly in } \xi \in \mathbb{R}.$$

Lemma 3.3 then implies that

$$\Phi_k(\xi) := \phi_k(\xi + \xi_k) \rightarrow 0$$

as $k \rightarrow +\infty$ locally uniformly in $\xi \in \mathbb{R}$.

From the boundedness of the sequence $\{\psi'_k/\psi_k\}$ in $L^\infty(\mathbb{R})$, one also infers that the functions

$$\xi \mapsto \Psi_k(\xi) = \frac{\psi_k(\xi + \xi_k)}{\psi_k(\xi_k)}$$

are locally bounded independently of k (in the sense that $\sup_{k \in \mathbb{N}} \|\Psi_k\|_{L^\infty(K)} < +\infty$ for any compact set $K \subset \mathbb{R}$). Each function Ψ_k obeys

$$c_k \Psi'_k = d D[\Psi_k] - (\mu + \gamma) \Psi_k + \beta \Phi_k \Psi_k \quad \text{in } \mathbb{R},$$

whence the functions Ψ'_k are locally bounded too. From Arzela-Ascoli theorem, the positive functions Ψ_k converge locally uniformly in \mathbb{R} , up to extraction of a subsequence, to a continuous nonnegative function Ψ_∞ . Furthermore, from the above equation and the fact that $\Phi_k \rightarrow 0$ as $k \rightarrow +\infty$ locally uniformly in \mathbb{R} (together with $c_k \rightarrow c^* > 0$), the functions Ψ'_k converge locally uniformly in \mathbb{R} too. Therefore, the functions Ψ_k converge in $C^1_{loc}(\mathbb{R})$ to Ψ_∞ and the function Ψ_∞ satisfies

$$(4.2) \quad c^* \Psi'_\infty = d D[\Psi_\infty] - (\mu + \gamma) \Psi_\infty \quad \text{in } \mathbb{R}.$$

Notice that this function Ψ_∞ is thus automatically of class $C^\infty(\mathbb{R})$. Furthermore, Ψ_∞ is nonnegative and $\Psi_\infty(0) = \lim_{k \rightarrow +\infty} \Psi_k(0) = 1$. As in the proof of Lemma 3.8 for the solution of (3.12), one then infers that Ψ_∞ is positive in \mathbb{R} .

Finally, for every $\xi \in \mathbb{R}$, there holds $\psi_k(\xi + \xi_k) \leq \|\psi_k\|_{L^\infty(\mathbb{R})} \leq (1 + 1/k) \psi_k(\xi_k)$ from (4.1). In other words, $\Psi_k(\xi) \leq 1 + 1/k$ for every $\xi \in \mathbb{R}$ and $k \in \mathbb{N}$ with $k \geq 1$, whence $\Psi_\infty(\xi) \leq 1$ for every $\xi \in \mathbb{R}$. Therefore, since $\Psi_\infty(0) = 1$, 0 is a global maximum of the function Ψ_∞ , and $\Psi'_\infty(0) = 0$, $D[\Psi_\infty](0) \leq 0$. The equation (4.2) evaluated at 0 leads to a contradiction, since μ and γ are positive. The proof of Lemma 4.2 is thereby complete. \square

End of the proof of Theorem 1.1 in case $c = c^$.* First of all, Lemma 3.8 applied with $\underline{c} = c^* > 0$ and $\bar{c} = c^* + 1$ yields the existence of $\varepsilon > 0$ such that $\Psi'(\xi) > 0$ for every $\xi \in \mathbb{R}$ with $\Psi(\xi) \leq \varepsilon$, and for every solution (Φ, Ψ) of (1.3) and (1.5) with speed $c \in [c^*, c^* + 1]$. Without loss of generality, one can assume that

$$(4.3) \quad 0 < \varepsilon \leq e^* = \frac{\mu}{\beta} (\sigma - 1).$$

Coming back to our solutions (ϕ_k, ψ_k) of (1.3) (with speed c_k) satisfying (1.5), (1.6) and (1.7), it follows from Lemma 4.1 and the positivity of each ψ_k that one can also assume without loss of generality that

$$0 < \varepsilon < \inf_{k \in \mathbb{N}} \|\psi_k\|_{L^\infty(\mathbb{R})}.$$

Therefore, for each $k \in \mathbb{N}$, since $\psi_k(-\infty) = 0$ and $\psi_k > 0$, there is $\xi_k \in \mathbb{R}$ such that

$$\psi_k(\xi_k) = \varepsilon.$$

Shift the origin at ξ_k and denote

$$\tilde{\phi}_k(\xi) = \phi_k(\xi + \xi_k) \quad \text{and} \quad \tilde{\psi}_k(\xi) = \psi_k(\xi + \xi_k).$$

From Lemma 4.2, the sequence $\{\tilde{\psi}_k\}$ is bounded in $L^\infty(\mathbb{R})$. Remember also that $0 < \tilde{\phi}_k < 1$ in \mathbb{R} and $c_k \rightarrow c^* > 0$ as $k \rightarrow +\infty$. Therefore, up to extraction of a subsequence, the functions $\tilde{\phi}_k$ and $\tilde{\psi}_k$ converge in $C_{loc}^\infty(\mathbb{R})$ to some bounded $C^\infty(\mathbb{R})$ functions ϕ and ψ solving (1.3) with speed c^* . Furthermore, $0 \leq \phi \leq 1$ and $\psi \geq 0$ in \mathbb{R} , while

$$\psi(0) = \varepsilon > 0.$$

In order to complete the proof of Theorem 1.1 in case $c = c^*$, one shall show that the pair (ϕ, ψ) is non-trivial and satisfies the desired limiting conditions at $\pm\infty$, that is, the conditions (1.5), (1.6) and (1.7) hold.

Let us first show that

$$\psi > 0 \quad \text{in } \mathbb{R}.$$

Indeed, if there is $\xi^* \in \mathbb{R}$ such that $\psi(\xi^*) = 0$, then $\psi'(\xi^*) = 0$ and equation (1.3) at ξ^* yields $\psi(\xi^* \pm 1) = 0$, whence $\psi(\xi^* + m) = 0$ for all $m \in \mathbb{Z}$ by immediate induction. But $c^*\psi' \geq -(2d + \mu + \gamma)\psi$ in \mathbb{R} , whence the function $\xi \mapsto \psi(\xi) e^{(2d + \mu + \gamma)\xi/c^*}$ is nondecreasing. Since $\psi \geq 0$ in \mathbb{R} and $\psi(\xi^* + m) = 0$ for all $m \in \mathbb{Z}$, one infers that $\psi = 0$ in \mathbb{R} , which is impossible since $\psi(0) = \varepsilon > 0$. Thus, $\psi > 0$ in \mathbb{R} . Once the positivity of ψ is known, it follows as in the proof of Lemma 3.1 that

$$0 < \phi < 1 \quad \text{in } \mathbb{R}.$$

In other words, the pair (ϕ, ψ) fulfills (1.5).

Let us then show that the pair (ϕ, ψ) satisfies the limiting conditions (1.6) at $-\infty$. Since the pair (ϕ, ψ) solves (1.3) and (1.5) with speed c^* , the choice of $\varepsilon > 0$ above and the property $\psi(0) = \varepsilon$ imply that $\psi' > 0$ in $(-\infty, 0]$. In particular, the limit $L = \lim_{\xi \rightarrow -\infty} \psi(\xi)$ exists, and $L \in [0, \varepsilon)$. If $L > 0$, then the same arguments as in the proof of Lemma 3.9 imply that $\phi(-\infty)$ exists and $\phi(-\infty) = (\mu + \gamma)/\beta = 1/\sigma \in (0, 1)$. The same arguments also

yield $L = \psi(-\infty) = \mu(1 - \phi(-\infty))/(\beta\phi(-\infty)) = (\mu/\beta)(\sigma - 1) = e^*$. Hence, $e^* = L < \varepsilon$, contradicting (4.3). Therefore,

$$L = \psi(-\infty) = 0.$$

Furthermore, for any sequence $\{\tilde{\xi}_k\}$ converging to $-\infty$, the functions $\xi \mapsto \phi(\xi + \tilde{\xi}_k)$ and $\xi \mapsto \psi(\xi + \tilde{\xi}_k)$ converge in $C_{loc}^\infty(\mathbb{R})$, up to extraction of a subsequence, to a pair $(\phi_\infty, 0)$, for some function $0 \leq \phi_\infty \leq 1$ solving (3.11) with speed $c_\infty = c^*$. It follows as in the proof of Lemma 3.8 that $\phi_\infty = 1$ in \mathbb{R} . Since the limit does not depend on the choice the sequence $\{\tilde{\xi}_k\}$, one gets that the limit $\lim_{\xi \rightarrow -\infty} \phi(\xi)$ exists, and

$$\phi(-\infty) = 1.$$

In other words, the pair (ϕ, ψ) satisfies (1.6).

Let us finally show that the non-triviality conditions (1.7) hold at $+\infty$. Firstly, as in the proof of Lemma 3.7, there holds $\inf_{\mathbb{R}} \phi > 0$. Secondly, Lemma 3.8 and (1.5) imply at once that $\liminf_{\xi \rightarrow +\infty} \psi(\xi) > 0$. Thirdly, one concludes that $\limsup_{\xi \rightarrow +\infty} \phi(\xi) < 1$ as in the proof of (3.6) and that (3.7) holds as in the case $c > c^*$. The solution (ϕ, ψ) thus fulfills all desired properties and the proof of Theorem 1.1 in case $c = c^*$ is thereby complete. \square

5. NON-EXISTENCE OF TRAVELING WAVES FOR $c < c^*$

In this section, (ϕ, ψ) denotes a classical solution of (1.3) satisfying (1.5) and (1.6), with a speed $c \in \mathbb{R}$. By classical, we mean that ϕ and ψ are of class $C^1(\mathbb{R})$ (and then of class $C^\infty(\mathbb{R})$) if $c \neq 0$, and that ϕ and ψ are continuous if $c = 0$. We shall prove that, necessarily, $c \geq c^*$. To do so, we consider separately the cases $c > 0$, $c < 0$ and $c = 0$.

First case: $c > 0$. Since the positive function ψ satisfies

$$\psi'(\xi) \geq \frac{d}{c} \psi(\xi + 1) - \frac{2d + \mu + \gamma}{c} \psi(\xi)$$

for all $\xi \in \mathbb{R}$, Lemma 3.2 implies that the functions $\xi \mapsto \psi(\xi \pm 1)/\psi(\xi)$ are bounded, and then so is the function $\xi \mapsto \psi'(\xi)/\psi(\xi)$. Consider now any sequence $\{\xi_k\}$ converging to $-\infty$. The positive functions

$$\xi \mapsto \psi_k(\xi) := \frac{\psi(\xi + \xi_k)}{\psi(\xi_k)}$$

are locally bounded and they satisfy

$$c \psi_k' = d D[\psi_k] - (\mu + \gamma) \psi_k + \beta \phi(\cdot + \xi_k) \psi_k \quad \text{in } \mathbb{R}.$$

Therefore, the functions ψ_k' are locally bounded too (remember that $\phi(-\infty) = 1$). From Arzela-Ascoli theorem, up to extraction of a subsequence, the functions ψ_k converge locally uniformly (and then in $C_{loc}^1(\mathbb{R})$ from the above equation) to a function ψ_∞ solving

$$(5.1) \quad c \psi_\infty' = d D[\psi_\infty] + (\beta - \mu - \gamma) \psi_\infty \quad \text{in } \mathbb{R}.$$

Furthermore, $\psi_\infty \geq 0$ in \mathbb{R} and $\psi_\infty(0) = 1$. As in the proof of Lemma 3.8 for the function Ψ_∞ solving (3.12), it follows that $\psi_\infty > 0$ in \mathbb{R} . Now, the function $z = \psi_\infty'/\psi_\infty$ solves

$$(5.2) \quad \frac{c}{d} z(x) = e^{\int_x^{x+1} z(s) ds} + e^{\int_x^{x-1} z(s) ds} - 2 + \frac{\beta - \mu - \gamma}{d} \quad \text{in } \mathbb{R}.$$

Proposition 3.5 implies that the limits $z(\pm\infty)$ exist in \mathbb{R} and are roots ω of the equation

$$c\omega = d(e^\omega + e^{-\omega} - 2) + \beta - \mu - \gamma.$$

Since $c > 0$ and $\beta > \mu + \gamma$, the roots must be positive and $c \geq c^*$ by definition of c^* in (1.4). The proof of the necessity condition is thereby complete in the case $c > 0$.

Second case: $c < 0$. Denote $\Phi(\xi) = \phi(-\xi)$ and $\Psi(\xi) = \psi(-\xi)$. The functions Φ and Ψ satisfy (1.3) and (1.5) with speed $|c| > 0$, together with the limiting conditions $(\Phi(+\infty), \Psi(+\infty)) = (1, 0)$. Furthermore, since the positive function Ψ satisfies

$$\Psi'(\xi) \geq \frac{d}{|c|} \Psi(\xi + 1) - \frac{2d + \mu + \gamma}{|c|} \Psi(\xi)$$

and (1.3) with speed $|c|$, it follows as in the above case $c > 0$ that the function Ψ'/Ψ is bounded. Since $\Psi > 0$ in \mathbb{R} and $\Psi(+\infty) = 0$, one can consider a sequence $\{\xi_k\}$ converging to $+\infty$ such that

$$\Psi'(\xi_k) \leq 0 \quad \text{for all } k \in \mathbb{N}.$$

As above, up to extraction of a subsequence, the functions $\xi \mapsto \Psi_k(\xi) := \Psi(\xi + \xi_k)/\Psi(\xi_k)$ converge in $C_{loc}^1(\mathbb{R})$ to a positive solution Ψ_∞ of (5.1) with $|c|$ instead of c , and such that $\Psi_\infty(0) = 1$. Furthermore, here, $\Psi'_\infty(0) \leq 0$. The function $Z := \Psi'_\infty/\Psi_\infty$ satisfies (5.2) with $|c|$ instead of c and it follows from Proposition 3.5 that the limits $Z(\pm\infty)$ exist in \mathbb{R} and are roots Ω of the equation

$$|c|\Omega = d(e^\Omega + e^{-\Omega} - 2) + \beta - \mu - \gamma.$$

Since $\beta > \mu + \gamma$, the roots are positive (and $|c| \geq c^*$). In particular, Z is positive at $\pm\infty$. But $Z(0) = \Psi'_\infty(0)/\Psi_\infty(0) = \Psi'_\infty(0) \leq 0$. Hence, the continuous function Z has a minimum Ξ in \mathbb{R} , that is $Z(\Xi) \leq Z(\xi)$ for all $\xi \in \mathbb{R}$. By differentiating the equation satisfied by Z , one gets as in (3.14) that

$$|c|Z'(\xi) = d(Z(\xi + 1) - Z(\xi)) \frac{\Psi_\infty(\xi + 1)}{\Psi_\infty(\xi)} + d(Z(\xi - 1) - Z(\xi)) \frac{\Psi_\infty(\xi - 1)}{\Psi_\infty(\xi)} \quad \text{in } \mathbb{R}.$$

Hence, $Z(\Xi \pm 1) = Z(\Xi)$, and $Z(\Xi + m) = Z(\Xi) = \min_{\mathbb{R}} Z$ for all $m \in \mathbb{Z}$ by immediate induction. Therefore, $Z(\pm\infty) = \min_{\mathbb{R}} Z \leq Z(0) \leq 0$, a contradiction with the positivity of $Z(\pm\infty)$. As a consequence, the case $c < 0$ is ruled out.

Third case: $c = 0$. Here, the function ψ satisfies $dD[\psi] + (\beta\phi - \mu - \gamma)\psi = 0$ in \mathbb{R} . Since $d > 0$, $\beta > \mu + \gamma$, $\phi(-\infty) = 1$ and $\psi > 0$ in \mathbb{R} , it follows that there exists $\xi_0 \in \mathbb{R}$ such that $D[\psi](\xi) < 0$ for all $\xi \leq \xi_0$. Denote

$$\theta(\xi) = \psi(\xi) - \psi(\xi + 1).$$

The condition $D[\psi] < 0$ in $(-\infty, \xi_0]$ means that $\theta(\xi - 1) < \theta(\xi)$ for all $\xi \leq \xi_0$. Furthermore, since $\psi > 0$ in \mathbb{R} and $\psi(-\infty) = 0$, there is $\xi_1 \leq \xi_0$ such that $\theta(\xi_1) < 0$. Since $\theta(\xi_1 - m) < \theta(\xi_1)$ for all $m \in \mathbb{N}$ with $m \geq 1$, one infers that

$$\psi(\xi_1 - m) - \psi(\xi_1) = \sum_{j=1}^m \theta(\xi_1 - j) < m\theta(\xi_1)$$

for all $m \in \mathbb{N}$ with $m \geq 1$. Thus, $\psi(\xi_1 - m) < \psi(\xi_1) + m\theta(\xi_1) \rightarrow -\infty$ as $m \rightarrow +\infty$ since $\theta(\xi_1) < 0$. This contradicts the positivity of ψ . As a consequence, the case $c = 0$ is ruled out too and the proof of Theorem 1.1 is thereby complete. \square

Remark 5.1. We give here another proof of the positivity of c when ψ is bounded (cf. [25]). Since $\phi(-\infty) = 1$ and $\beta > \mu + \gamma$, there is a sufficiently large K such that

$$\beta \phi(\xi) - \mu - \gamma > \frac{\beta - \mu - \gamma}{2} > 0 \quad \text{for } \xi \in (-\infty, -K).$$

Integrating the second equation of (1.3) from $-\infty$ to $\xi < -K$, using $\psi(-\infty) = 0$ and the positivity and boundedness of ψ , we obtain

$$\begin{aligned} c\psi(\xi) &= d \left\{ \int_{\xi}^{\xi+1} \psi(s) ds - \int_{\xi-1}^{\xi} \psi(s) ds \right\} + \int_{-\infty}^{\xi} [\beta \phi(s) - \mu - \gamma] \psi(s) ds \\ &\geq -d \left\{ \sup_{s \in \mathbb{R}} \psi(s) \right\} + \frac{\beta - \mu - \gamma}{2} \int_{-\infty}^{\xi} \psi(s) ds \end{aligned}$$

for all $\xi < -K$. It follows that the integral

$$R(\xi) := \int_{-\infty}^{\xi} \psi(s) ds$$

is well-defined for all $\xi < -K$ (and then for all $\xi \in \mathbb{R}$ by continuity of ψ). Integrating the second equation of (1.3) twice, we obtain

$$cR(x) = d \left\{ \int_x^{x+1} R(\xi) d\xi - \int_{x-1}^x R(\xi) d\xi \right\} + \int_{-\infty}^x \int_{-\infty}^{\xi} [\beta \phi(s) - \mu - \gamma] \psi(s) ds d\xi$$

for all $x \in \mathbb{R}$. Since $R(\xi)$ is strictly increasing, we conclude that $c > 0$.

REFERENCES

- [1] M. Alfaro, J. Coville, *Rapid traveling waves in the nonlocal Fisher equation connect two unstable states*, Appl. Math. Letters 25 (2012), 2095–2099.
- [2] N. Apreutesei, N. Bessonov, V. Volpert, V. Vougalter, *Spatial structures and generalized travelling waves for an integro-differential equation*, Disc. Cont. Dyn. Syst. B 13 (2010), 537–557.
- [3] H. Berestycki, F. Hamel, A. Kiselev, L. Ryzhik, *Quenching and propagation in KPP reaction-diffusion equations with a heat loss*, Arch. Ration. Mech. Anal. 178 (2005), 57–80.
- [4] H. Berestycki, G. Nadin, B. Perthame, L. Ryzhik, *The non-local Fisher-KPP equation: traveling waves and steady states*, Nonlinearity 22 (2009), 2813–2844.
- [5] J.W. Cahn, J. Mallet-Paret, E.S. van Vleck, *Travelling wave solutions for systems of ODE's on a two-dimensional spatial lattice*, SIAM J. Appl. Math. 59 (1999), 455–493.
- [6] X. Chen, S.-C. Fu, J.-S. Guo, *Uniqueness and asymptotics of traveling waves of monostable dynamics on lattices*, SIAM J. Math. Anal. 38 (2006), 233–258.
- [7] X. Chen, J.-S. Guo, *Existence and asymptotic stability of traveling waves of discrete quasilinear monostable equations*, J. Diff. Equations 184 (2002), 549–569.
- [8] X. Chen, J.-S. Guo, *Uniqueness and existence of traveling waves for discrete quasilinear monostable dynamics*, Math. Ann. 326 (2003), 123–146.
- [9] X. Chen, J.-S. Guo, C.-C. Wu, *Traveling waves in discrete periodic media for bistable dynamics*, Arch. Ration. Mech. Anal. 189 (2008), 189–236.

- [10] S.N. Chow, J. Mallet-Paret, W. Shen, *Travelling waves in lattice dynamical systems*, J. Diff. Equations 149 (1998), 249–291.
- [11] R. Coutinho, B. Fernandez, *Extended symbolic dynamics in bistable CML: existence and stability of fronts*, Physica D 108 (1997), 60–80.
- [12] R. Coutinho, B. Fernandez, *Fronts and interfaces in bistable extended mappings*, Nonlinearity 11 (1998), 1407–1433.
- [13] R. Coutinho, B. Fernandez, *Fronts in extended systems of bistable maps coupled via convolutions*, Nonlinearity 17 (2004), 23–47.
- [14] R. Coutinho, B. Fernandez, *Fronts between periodic patterns for bistable recursions on lattices*, J. Dyn. Diff. Equations 25 (2013), 17–31.
- [15] J. Fang, J. Wei, X.-Q. Zhao, *Spreading speeds and travelling waves for non-monotone time-delayed lattice equations*, Proc. Royal Soc. A 466 (2010), 1919–1934.
- [16] J. Fang, J. Wei, X.-Q. Zhao, *Uniqueness of traveling waves for nonlocal lattice equations*, Proc. Amer. Math. Soc. 139 (2011), 1361–1373.
- [17] J. Fang, X.-Q. Zhao, *Monotone wavefronts of the nonlocal Fisher-KPP equation*, Nonlinearity 24 (2011), 3043–3054.
- [18] G. Faye, M. Holzer, *Modulated traveling fronts for a nonlocal Fisher-KPP equation: a dynamical systems approach*, J. Diff. Equations 258 (2015), 2257–2289.
- [19] S.-C. Fu, J.-S. Guo, C.-C. Wu, *Traveling wave solutions for a discrete diffusive epidemic model*, J. Nonlinear Convex Anal. (to appear).
- [20] S. Genieys, V. Volpert, P. Auger, *Pattern and waves for a model in population dynamics with nonlocal consumption of resources*, Math. Model. Nat. Phenom. 1 (2006), 63–80.
- [21] J.-S. Guo, F. Hamel, *Front propagation for discrete periodic monostable equations*, Math. Ann. 335 (2006), 489–525.
- [22] J.-S. Guo, Y.-C. Lin, *Traveling wave solution for a lattice dynamical system with convolution type nonlinearity*, Disc. Cont. Dyn. Syst. 32 (2012), 101–124.
- [23] J.-S. Guo, C.-H. Wu, *Existence and uniqueness of traveling waves for a monostable 2-D lattice dynamical system*, Osaka J. Math. 5 (2008), 327–346.
- [24] J.-S. Guo, C.-H. Wu, *Wave propagation for a two-component lattice dynamical system arising in strong competition models*, J. Differential Equations 250 (2011), 3504–3533.
- [25] J.-S. Guo, C.-H. Wu, *Traveling wave front for a two-component lattice dynamical system arising in competition models*, J. Differential Equations 252 (2012), 4357–4391.
- [26] D. Hankerson, B. Zinner, *Wavefronts for a cooperative tridiagonal system of differential equations*, J. Dyn. Diff. Equations 5 (1993), 359–373.
- [27] H.W. Hethcote, *The mathematics of infectious diseases*, SIAM Review 42 (2000), 599–653.
- [28] A. Hoffman, H.J. Hupkes, E.S. Van Vleck, *Multi-dimensional stability of waves travelling through rectangular lattices in rational directions*, Trans. Amer. Math. Soc. 367 (2015), 8757–8808.
- [29] A. Hoffman, B. Kennedy, *Existence and uniqueness of traveling waves in a class of unidirectional differential equations*, Disc. Cont. Dyn. Syst. 30 (2011), 137–167.
- [30] W. Hudson, B. Zinner, *Existence of traveling waves for a generalized discrete Fisher’s equation*, Comm. Appl. Nonlinear Anal. 1 (1994), 23–46.
- [31] J.P. Keener, *Propagation and its failure in coupled systems of discrete excitable cells*, SIAM J. Appl. Math. 47 (1987), 556–572.
- [32] J.P. Keener, *Homogenization and propagation in the bistable equation*, Physica D 136 (2000), 1–17.
- [33] S.-G. Liao, C.-C. Wu, *Propagation failure in discrete periodic media*, J. Diff. Equ. Appl. 19 (2013), 1268–1275.
- [34] S. Ma, P. Weng, X. Zou, *Asymptotic speed of propagation and traveling wavefronts in a non-local delayed lattice differential equation*, Nonlinear Anal. TMA 65 (2006), 1858–1890.

- [35] S. Ma, X. Zou, *Existence, uniqueness and stability of travelling waves in a discrete reaction-diffusion monostable equation with delay*, J. Diff. Equations 217 (2005), 54–87.
- [36] G. Nadin, B. Perthame, M. Tang, *Can a traveling wave connect two unstable states? The case of the nonlocal Fisher equation*, C. R. Acad. Sci. Paris I 349 (2010), 553–557.
- [37] G. Nadin, L. Rossi, L. Ryzhik, B. Perthame, *Wave-like solutions for nonlocal reaction-diffusion equations: a toy model*, Math. Model. Nat. Phenom. 8 (2013), 33–41.
- [38] Z. Xu, J. Ma, *Monotonicity, asymptotics and uniqueness of travelling wave solution of a non-local delayed lattice dynamical system*, Disc. Cont. Dyn. Syst. 35 (2015), 5107–5131.
- [39] B. Zinner, *Stability of traveling wavefronts for the discrete Nagumo equation*, SIAM J. Math. Anal. 22 (1991), 1016–1020.
- [40] B. Zinner, *Existence of traveling wavefront solutions for the discrete Nagumo equation*, J. Diff. Equations 96 (1992), 1–27.
- [41] B. Zinner, G. Harris, W. Hudson, *Traveling wavefronts for the discrete Fisher’s equation*, J. Diff. Equations 105 (1993), 46–62.

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