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THE HEAT FLOW WITH A CRITICAL EXPONENTIAL NONLINEARITY

TOBIAS LAMM, FRÉDÉRIC ROBERT, AND MICHAEL STRUWE

ABSTRACT. We analyze the possible concentration behavior of heat flows related to the Moser-Trudinger energy and derive quantization results completely analogous to the quantization results for solutions of the corresponding elliptic equation. As an application of our results we obtain the existence of critical points of the Moser-Trudinger energy in a supercritical regime.

1. INTRODUCTION

On any bounded domain $\Omega \subset \mathbb{R}^2$ the Moser-Trudinger energy functional

$$E(u) = \frac{1}{2} \int_{\Omega} (e^{u^2} - 1) dx$$

for any $\alpha \leq 4\pi$ admits a maximizer in the space

$$(1) \quad M_{\alpha} = \{u \in H_0^1(\Omega); u \geq 0, \|\nabla u\|_{L^2}^2 = \alpha\},$$

corresponding to a solution $0 < u \in H_0^1(\Omega)$ of the equation

$$(2) \quad -\Delta u = \lambda u e^{u^2} \text{ in } \Omega$$

for some $\lambda > 0$; see [7]. Moreover, when Ω is a ball numerical evidence [9] shows that for small $\alpha > 4\pi$ there exists a pair of critical points of E in M_{α} , corresponding to a relative maximizer and a saddle point of E , respectively. However, standard variational techniques fail in this “supercritical” energy range and ad hoc methods devised to remedy the situation so far have only been partially successful in producing the expected existence results; compare [11], [12]. As in various other geometric variational problems a flow method might turn out to be more useful in this regard.

Given a smooth function $0 \leq u_0 \in H_0^1(\Omega)$, we consider smooth solutions $u = u(t, x)$ to the equation

$$(3) \quad u_t e^{u^2} = \Delta u + \lambda u e^{u^2} \text{ in } [0, \infty[\times \Omega$$

with initial and boundary data

$$(4) \quad u(0) = u_0, \quad u = 0 \text{ on } [0, \infty[\times \partial\Omega.$$

The function $\lambda = \lambda(t)$ may be determined so that the Dirichlet integral of u is preserved along the flow. As we shall see, also the case where the volume of the evolving metric $g = e^{u^2} g_{\mathbb{R}^2}$ is fixed gives rise to interesting applications, and both constraints can easily be analyzed in parallel.

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1.1. **Fixed volume.** Fixing the volume is equivalent to the constraint

$$(5) \quad E(u(t)) = E(u_0) =: c_0 \text{ for all } t,$$

which can be achieved by imposing the condition

$$(6) \quad \frac{d}{dt}E(u(t)) = \int_{\Omega} uu_t e^{u^2} dx = \lambda \int_{\Omega} u^2 e^{u^2} dx - \int_{\Omega} |\nabla u|^2 dx = 0.$$

Clearly, we may assume that u_0 does not vanish identically and that $c_0 > 0$; otherwise $u \equiv 0$ is the unique smooth solution to (3) - (5) for any choice of $\lambda(t)$.

Note that when we multiply (3) with u_t and use (6), upon integrating by parts we obtain the relation

$$(7) \quad \int_{\Omega} u_t^2 e^{u^2} dx + \frac{1}{2} \frac{d}{dt} \left(\int_{\Omega} |\nabla u|^2 dx \right) = \lambda \frac{d}{dt} E(u(t)) = 0;$$

that is, the flow (3) - (5) may be regarded as the gradient flow (with respect to the metric g) for the Dirichlet energy with the critical exponential constraint (5).

Equation (6) and the energy inequality (7) imply the uniform bound

$$(8) \quad \lambda \int_{\Omega} u^2 e^{u^2} dx = \int_{\Omega} |\nabla u|^2 dx \leq \int_{\Omega} |\nabla u_0|^2 dx =: \Lambda_0.$$

Since we can easily estimate $e^a \leq 1 + 4a$ for $0 \leq a \leq 1/4$, we have

$$(9) \quad \begin{aligned} \int_{\Omega} u^2 e^{u^2} dx &= \int_{\Omega} u^2 (e^{u^2} - 1) dx + \int_{\Omega} u^2 dx \\ &\geq \frac{1}{4} \int_{\Omega} (e^{u^2} - 1) dx - \frac{1}{4} \int_{\{x \in \Omega; u \leq 1/2\}} (e^{u^2} - 1) dx + \int_{\Omega} u^2 dx \geq \frac{E(u)}{2} \geq \frac{c_0}{2}, \end{aligned}$$

for all t . Therefore, recalling that $c_0 > 0$, from (8) we deduce that with the constant $\lambda_0 = 2\Lambda_0/c_0 > 0$ there holds

$$(10) \quad 0 < \lambda(t) \leq \lambda_0 \text{ for all } t \geq 0.$$

Finally, the maximum principle yields that $u \geq 0$.

1.2. **Constant Dirichlet integral.** If, on the other hand, we choose λ so that

$$(11) \quad \frac{1}{2} \frac{d}{dt} \left(\int_{\Omega} |\nabla u|^2 dx \right) = - \int_{\Omega} u_t \Delta u dx = \lambda \int_{\Omega} |\nabla u|^2 dx - \int_{\Omega} |\Delta u|^2 e^{-u^2} dx = 0,$$

for a solution of (3), (4) satisfying (11) the Dirichlet integral is preserved; that is,

$$(12) \quad \int_{\Omega} |\nabla u|^2 dx = \int_{\Omega} |\nabla u_0|^2 dx = \Lambda_0,$$

In this case, from (7) we find the equation

$$(13) \quad \int_{\Omega} u_t^2 e^{u^2} dx = \lambda \frac{d}{dt} E(u(t)),$$

and (3), (4) with the constraint (12) turns into the (positive) gradient flow for the Moser-Trudinger energy with prescribed Dirichlet integral. Again clearly we may assume that $\Lambda_0 > 0$.

Recalling the identity

$$\frac{d}{dt} E(u(t)) = \lambda \int_{\Omega} u^2 e^{u^2} dx - \int_{\Omega} |\nabla u|^2 dx$$

(13) and (9), for any t we have

$$(14) \quad \frac{c_0}{2} \int_0^t \lambda dt \leq \Lambda_0 t + E(u(t)) - E(u_0),$$

where $c_0 = E(u_0) \leq E(u(t))$ for all $t \geq 0$. Similarly, from (13) we obtain

$$(15) \quad \int_0^t (\lambda^{-1} \int_{\Omega} u_t^2 e^{u^2} dx) dt = E(u(t)) - E(u_0).$$

Hence we can hope to obtain bounds for solutions of (3), (4), (12) whenever the Moser-Trudinger energy is bounded along the flow.

1.3. Results. Building on previous results from [2], [6], and [14], in this paper we establish the following result for the flow (3), (4) with either the constraint (5) or the constraint (12).

Theorem 1.1. *For any $c_0 > 0$ and any smooth initial data $0 \leq u_0 \in H_0^1(\Omega)$ satisfying (5) the evolution problem (3) - (5) admits a unique smooth solution $u \geq 0$ for all $t > 0$. Likewise, for any smooth $0 \leq u_0 \in H_0^1(\Omega)$ satisfying (12) for a given $\Lambda_0 > 0$ the evolution problem (3), (4), (12) admits a unique smooth solution $u \geq 0$ for small $t > 0$ which can be continued smoothly for all $t > 0$, provided that $E(u(t))$ remains bounded. In both cases, for a suitable sequence $t_k \rightarrow \infty$ the functions $u(t_k) \rightarrow u_{\infty}$ weakly in $H_0^1(\Omega)$, where $u_{\infty} \in H_0^1(\Omega)$ is a solution to the problem (2) for some constant $\lambda_{\infty} \geq 0$. Moreover, either $u(t_k) \rightarrow u_{\infty}$ strongly in $H_0^1(\Omega)$, $\lambda_{\infty} > 0$, and $0 < u_{\infty} \in H_0^1(\Omega)$ satisfies, respectively, (5) or (12), or there exist $i_* \in \mathbb{N}$ and points $x^{(i)} \in \bar{\Omega}$, $l_i \in \mathbb{N}$, $1 \leq i \leq i_*$, such that as $k \rightarrow \infty$ we have*

$$|\nabla u(t_k)|^2 dx \xrightarrow{w^*} |\nabla u_{\infty}|^2 dx + \sum_{i=1}^{i_*} 4\pi l_i \delta_{x^{(i)}}$$

weakly in the sense of measures. By (8) or (12) then necessarily $4\pi \sum_{i=1}^{i_} l_i \leq \Lambda_0$.*

The quantization result in the case of divergence of the flow relies on the precise microscopic description of blow-up given in Sections 4 and 5; see in particular Theorems 4.2 and 5.1. Their derivation will take up the major part of this paper. These results are in complete analogy with the results of Adimurthi-Struwe [2] and Druet [6] for solutions of the corresponding elliptic equation (2). In the final Section 6 we apply Theorem 1.1 to obtain saddle-point solutions for (2) in the supercritical regime of large energy.

2. GLOBAL EXISTENCE

Let $u(t)$ be a solution of (3), (4) with either the constraint (5) or (12). In the latter case we also assume that $E(u(t))$ remains bounded. For any $t \geq 0$ let $m(t) = \|u(t)\|_{L^{\infty}}$. Writing equation (3) in the form

$$u_t - e^{-u^2} \Delta u = \lambda u \text{ in } [0, \infty[\times \Omega$$

and observing that $\Delta u \leq 0$ at any point where $u(t)$ achieves its maximum, we conclude that the supremum of the function $\tilde{u}(t) = e^{-\int_0^t \lambda(s) ds} u(t)$ is non-increasing in time. That is, for any $0 \leq t_0 \leq t < \infty$ we have

$$(16) \quad m(t) \leq e^{\int_{t_0}^t \lambda(s) ds} m(t_0).$$

Together with (10), (14) this immediately gives the following result.

Lemma 2.1. *Suppose that $E(u(t))$ is uniformly bounded. Then there exist constants $\lambda_1 > 0$, C_1 depending on u_0 such that for any $t \geq 0$ we have*

$$\|u(t)\|_{L^\infty} \leq e^{\int_0^t \lambda(s) ds} \|u_0\|_{L^\infty} \leq C_1 e^{\lambda_1 t} \|u_0\|_{L^\infty}.$$

Existence of a unique smooth solution on any finite time interval now follows from standard results on uniformly parabolic equations.

3. ASYMPTOTIC BEHAVIOR

3.1. Weak subconvergence. First consider the constraint (5). Integrating in time, from (7) we then obtain

$$(17) \quad \int_0^\infty \int_\Omega u_t^2 e^{u^2} dx dt \leq \frac{1}{2} \int_\Omega |\nabla u_0|^2 dx.$$

Hence we can find a sequence $t_k \rightarrow \infty$ such that

$$(18) \quad \int_{\{t_k\} \times \Omega} u_t^2 e^{u^2} dx \rightarrow 0 \text{ as } k \rightarrow \infty.$$

In view of (10) and (8) from any such sequence (t_k) we may extract a subsequence such that $\lambda_\infty = \lim_{k \rightarrow \infty} \lambda(t_k)$ exists and such that, in addition, $u_k = u(t_k) \rightharpoonup u_\infty$ weakly in $H_0^1(\Omega)$ and pointwise almost everywhere as $k \rightarrow \infty$. From (8) by means of the Vitali convergence theorem we then deduce that for a further subsequence the terms $\lambda u e^{u^2}$, evaluated at $t = t_k$, converge to $\lambda_\infty u_\infty e^{u_\infty^2}$ in $L^1(\Omega)$. Thus, upon passing to the limit $k \rightarrow \infty$ in (3) we see that u_∞ is a (weak) solution to equation (2). But since $u_\infty \in H_0^1(\Omega)$, from the Moser-Trudinger inequality it follows that $u_\infty e^{u_\infty^2} \in L^p(\Omega)$ for any $p < \infty$, and u_∞ is, in fact, smooth.

Similarly, in the case of the the constraint (12), assuming that $E(u(t))$ is uniformly bounded from above along the flow (3), (4), from (15) we obtain the bound

$$(19) \quad \int_0^\infty \left(\lambda^{-1} \int_\Omega u_t^2 e^{u^2} dx \right) dt \leq \lim_{t \rightarrow \infty} E(u(t)) - E(u_0) < \infty,$$

and we can find a sequence $t_k \rightarrow \infty$ such that

$$(20) \quad \lambda(t_k)^{-1} \int_{\{t_k\} \times \Omega} u_t^2 e^{u^2} dx \rightarrow 0 \text{ as } k \rightarrow \infty.$$

Necessarily the sequence $(\lambda(t_k))$ is bounded. Indeed, upon multiplying (3) by u we infer that at time t_k with error $o(1) \rightarrow 0$ we have

$$\lambda \int_\Omega u^2 e^{u^2} dx = \int_\Omega |\nabla u|^2 dx + \int_\Omega uu_t e^{u^2} dx$$

But by (20) and Hölder's inequality, at time $t = t_k$ with error $o(1) \rightarrow 0$ as $k \rightarrow \infty$ we can estimate

$$(21) \quad \left| \int_\Omega uu_t e^{u^2} dx \right|^2 \leq \lambda \int_\Omega u^2 e^{u^2} dx \cdot \lambda^{-1} \int_\Omega u_t^2 e^{u^2} dx = o(1) \lambda \int_\Omega u^2 e^{u^2} dx$$

and we have

$$(22) \quad (1 + o(1)) \lambda \int_\Omega u^2 e^{u^2} dx = \int_\Omega |\nabla u|^2 dx = \Lambda_0.$$

Our claim now follows from (9). Note that, in particular, the approximate identity (8) thus also holds in the case of the constraint (12).

3.2. The case when u is bounded. If in addition we assume that the function u is uniformly bounded we find that any sequence (u_k) as above is bounded in $H^2(\Omega)$ and hence possesses a subsequence such that $u_k \rightarrow u_\infty$ strongly in $H_0^1(\Omega)$ as $k \rightarrow \infty$. Hence $u_\infty \in H_0^1(\Omega)$ satisfies, respectively, (5) or (12), and $u_\infty > 0$ by the maximum principle.

In the case of the constraint (5), and provided that u is bounded, we can even show relative compactness of the sequence $u_k = u(t_k)$ for *any* sequence $t_k \rightarrow \infty$.

Proposition 3.1. *Let u solve (3) - (5). Suppose that there exists a uniform constant $M > 0$ such that $u(t, x) \leq M$ for all $x \in \Omega$ and all $t \geq 0$. Then any sequence $u_k = u(t_k)$ with $t_k \rightarrow \infty$ has a strongly convergent subsequence.*

Proof. It suffices to show that under the assumptions of the Proposition the convergence in (18) can be improved to be uniform in time. To show this we use (3) to calculate

$$\begin{aligned} u_{tt} &= \lambda_t u + \lambda u_t - 2uu_t e^{-u^2} \Delta u + e^{-u^2} \Delta u_t \\ &= \lambda_t u + \lambda u_t + e^{-u^2} \Delta u_t - 2u u_t^2 + 2\lambda u^2 u_t. \end{aligned}$$

Thus we obtain

$$\begin{aligned} \frac{1}{2} \frac{d}{dt} \left(\int_{\Omega} u_t^2 e^{u^2} dx \right) &= \int_{\Omega} u_t u_{tt} e^{u^2} dx + \int_{\Omega} u_t^3 u e^{u^2} dx \\ &= \lambda_t \int_{\Omega} u u_t e^{u^2} dx + \lambda \int_{\Omega} u_t^2 e^{u^2} dx + \int_{\Omega} u_t \Delta u_t dx \\ &\quad - 2 \int_{\Omega} u u_t^3 e^{u^2} dx + 2\lambda \int_{\Omega} u^2 u_t^2 e^{u^2} dx \end{aligned}$$

By (6) the first term on the right vanishes. Moreover, we may use the fact $u_t = 0$ on $\partial\Omega$ to integrate by parts in the third term. Also using Hölders inequality and Sobolev's embedding $W^{1,2} \hookrightarrow L^4$ then with constants $C = C(M)$ we find

$$\begin{aligned} &\int_{\Omega} |\nabla u_t|^2 dx + \frac{1}{2} \frac{d}{dt} \left(\int_{\Omega} u_t^2 e^{u^2} dx \right) \\ (23) \quad &\leq C \int_{\Omega} u_t^2 e^{u^2} dx + C \left(\int_{\Omega} u_t^2 e^{u^2} dx \right)^{\frac{1}{2}} \left(\int_{\Omega} u_t^4 dx \right)^{\frac{1}{2}} \\ &\leq C \int_{\Omega} u_t^2 e^{u^2} dx + C_1 \left(\int_{\Omega} u_t^2 e^{u^2} dx \right)^{\frac{1}{2}} \int_{\Omega} (|\nabla u_t|^2 + u_t^2 e^{u^2}) dx. \end{aligned}$$

To proceed, we use an argument similar to [13], p. 271. Given any number $\varepsilon_0 > 0$, by (17) there exist arbitrary large times t_0 such that

$$(24) \quad \int_{\{t_0\} \times \Omega} u_t^2 e^{u^2} dx < \varepsilon_0.$$

For any such t_0 we may choose a maximal $t_0 \leq t_1 \leq \infty$ such that

$$(25) \quad \sup_{t_0 \leq t \leq t_1} \int_{\{t\} \times \Omega} u_t^2 e^{u^2} dx \leq 2\varepsilon_0.$$

If we now fix $\varepsilon_0 = \frac{1}{16C_1^2}$, from (23) at any time $t \in [t_0, t_1]$ we obtain

$$(26) \quad \frac{1}{2} \frac{d}{dt} \left(\int_{\Omega} u_t^2 e^{u^2} dx \right) \leq C \int_{\Omega} u_t^2 e^{u^2} dx.$$

Integrating from t_0 to t and using (17), for any $t \in [t_0, t_1]$ we get

$$(27) \quad \int_{\{t\} \times \Omega} u_t^2 e^{u^2} dx \leq \int_{\{t_0\} \times \Omega} u_t^2 e^{u^2} dx + C \int_{t_0}^{\infty} \int_{\Omega} u_t^2 e^{u^2} dx < 2\varepsilon_0,$$

if t_0 is large enough. For such t_0 then $t_1 = \infty$, and we conclude

$$(28) \quad \begin{aligned} & \limsup_{t \rightarrow \infty} \int_{\{t\} \times \Omega} u_t^2 e^{u^2} dx \\ & \leq \liminf_{t_0 \rightarrow \infty} \left(\int_{\{t_0\} \times \Omega} u_t^2 e^{u^2} dx + C \int_{t_0}^{\infty} \int_{\Omega} u_t^2 e^{u^2} dx \right) = 0. \end{aligned}$$

Using again the assumption that u is uniformly bounded this directly implies that

$$(29) \quad \limsup_{t \rightarrow \infty} \|u(t)\|_{H^2} < \infty$$

and hence the claim. \square

4. BLOW-UP ANALYSIS

It remains to analyze the blow-up behavior of a solution u to (3), (4) satisfying either (5) or (12) in the case when u is unbounded. As we shall see, this can be done in complete analogy with the corresponding time-independent problem. The key is the following lemma, which refines our above choice of (t_k) .

Lemma 4.1. *Suppose that $\limsup_{t \rightarrow \infty} \|u(t)\|_{L^\infty} = \infty$ and that $E(u(t)) \leq E_\infty$ for some constant $E_\infty < \infty$. Then there is a sequence $t_k \rightarrow \infty$ with associated numbers $\lambda_k = \lambda(t_k) \rightarrow \lambda_\infty \geq 0$ such that $u(t_k) \rightharpoonup u_\infty$ weakly in $H_0^1(\Omega)$ as $k \rightarrow \infty$ and*

$$\|u(t_k)\|_{L^\infty} \rightarrow \infty, \quad \lambda_k^{-1} \int_{\{t_k\} \times \Omega} |u_t|^2 e^{u^2} dx dt \rightarrow 0.$$

Proof. Suppose by contradiction that there exist $t_0 \geq 0$ and a constant $C_0 > 0$ such that for all $t \geq t_0$ either there holds

$$m(t) = \|u(t)\|_{L^\infty} \leq C_0,$$

or

$$(30) \quad \lambda(t) \leq C_0 \int_{\{t\} \times \Omega} |u_t|^2 e^{u^2} dx.$$

Consider first the constraint (5). If $m(t) > C_0$ for all $t \geq t_0$, then (30) holds for all such t and upon integrating in time from (7) for any $t \geq t_0$ we obtain

$$(31) \quad \int_{t_0}^t \lambda(s) ds \leq C_0 \int_0^\infty \int_{\Omega} |u_t|^2 e^{u^2} dx dt \leq \frac{C_0 \Lambda_0}{2} =: C_1 < \infty.$$

Applying (16) to the shifted flow $u(t - t_0)$ we find $\sup_{t \geq t_0} m(t) \leq m(t_0) e^{C_1} < \infty$, contrary to assumption.

If for some $t_0 \leq t_1 < t_2 \leq \infty$ and all $t_1 < t < t_2$ we have $m(t_1) = C_0 < m(t)$, then (30) holds for all such t and we obtain (31) with t_1 replacing t_0 for all $t \in [t_1, t_2]$. Applying (16) to the shifted flow $u(t - t_1)$, for any such $t_0 \leq t_1 < t_2 \leq \infty$ we obtain the bound $\sup_{t_1 < t \leq t_2} m(t) \leq C_0 e^{C_1} < \infty$, again contradicting our hypotheses.

In case of the constraint (12), whenever for some $t_0 \leq t_1 < t_2 \leq \infty$ and all $t_1 < t < t_2$ there holds $m(t) > C_0$ from (30) and (15) we obtain

$$(32) \quad t_2 - t_1 \leq C_0 \int_0^\infty \left(\lambda(t)^{-1} \int_\Omega |u_t|^2 e^{u^2} dx \right) dt \leq C_0 E_\infty =: T_0 < \infty.$$

By (32) the length of any interval $I =]t_1, t_2[$ with $m(t) > C_0$ for $t \in I$ is uniformly bounded. Since $\limsup_{t \rightarrow \infty} m(t) = \infty$, we may then assume that $m(t_1) = C_0$. Applying (16) to the shifted flow $u(t - t_1)$, by (14) for any such interval we find $\sup_{t_1 < t \leq t_2} m(t) \leq C_0 e^{C_2}$, where $C_2 = 2c_0^{-1}(\Lambda_0 T_0 + E_\infty) < \infty$. Thus we also have $\limsup_{t \rightarrow \infty} m(t) \leq C_0 e^{C_2}$, contrary to hypothesis. \square

For a sequence (t_k) as determined in Lemma 4.1 above we let $u_k = u(t_k)$, $k \in \mathbb{N}$ and set $\dot{u}_k = u_t(t_k)$. The symbols t , t_k then no longer explicitly appear and we may use these letters for other purposes. Also let $\eta = \log \left(\frac{2}{1+|x|^2} \right)$ be the standard solution of Liouville's equation

$$(33) \quad -\Delta \eta = e^{2\eta} \text{ on } \mathbb{R}^2$$

induced by stereographic projection from S^2 , with

$$(34) \quad \int_{\mathbb{R}^2} e^{2\eta} dx = 4\pi =: \Lambda_1.$$

Similar to [2], [6] the following result now holds.

Theorem 4.2. *There exist a number $i_* \in \mathbb{N}$ and points $x^{(i)} \in \bar{\Omega}$, $1 \leq i \leq i_*$, such that as $k \rightarrow \infty$ suitably for each i with suitable points $x_k = x_k^{(i)} \rightarrow x^{(i)}$ and scale factors $0 < r_k = r_k^{(i)} \rightarrow 0$ satisfying*

$$(35) \quad \lambda_k r_k^2 u_k^2(x_k) e^{u_k^2(x_k)} = 4$$

we have

$$(36) \quad \eta_k(x) = \eta_k^{(i)}(x) := u_k(x_k)(u_k(x_k + r_k x) - u_k(x_k)) \rightarrow \eta_0 = \log \left(\frac{1}{1+|x|^2} \right)$$

locally uniformly on \mathbb{R}^2 , where $\eta_0 = \eta - \log 2$ satisfies

$$(37) \quad -\Delta \eta_0 = 4e^{2\eta_0} \text{ on } \mathbb{R}^2,$$

and there holds

$$(38) \quad \lim_{L \rightarrow \infty} \lim_{k \rightarrow \infty} \lambda_k \int_{B_{Lr_k}(x_k)} u_k^2 e^{u_k^2} dx = 4 \int_{\mathbb{R}^2} e^{2\eta_0} dx = \Lambda_1.$$

Equality $x^{(i)} = x^{(j)}$ may occur, but we have

$$(39) \quad \frac{\text{dist}(x_k^{(i)}, \partial\Omega)}{r_k^{(i)}}, \frac{|x_k^{(i)} - x_k^{(j)}|}{r_k^{(i)}} \rightarrow \infty \text{ for all } 1 \leq i \neq j \leq i_*,$$

and there holds the uniform pointwise estimate

$$(40) \quad \lambda_k \inf_i |x - x_k^{(i)}|^2 u_k^2(x) e^{u_k^2(x)} \leq C,$$

for all $x \in \Omega$ and all $k \in \mathbb{N}$.

Finally, $u_k \rightarrow u_\infty$ in $H_{loc}^2(\Omega \setminus \{x_1, \dots, x_{i_*}\})$ as $k \rightarrow \infty$.

Proof. Choose $x_k = x_k^{(1)} \in \Omega$ such that $u_k(x_k) = \sup_{x \in \Omega} u_k$ and let $r_k = r_k^{(1)}$ be given by (35). We claim that $r_k \rightarrow 0$ as $k \rightarrow \infty$. Otherwise, (35) gives $\lambda_k u_k^2(x_k) e^{u_k^2(x_k)} \leq C < \infty$, and with the help of Lemma 4.1 we can estimate

$$\int_{\Omega} |u_k(x_k) \dot{u}_k e^{u_k^2}|^2 dx \leq \lambda_k u_k^2(x_k) e^{u_k^2(x_k)} \left(\lambda_k^{-1} \int_{\Omega} \dot{u}_k^2 e^{u_k^2} dx \right) \rightarrow 0$$

as $k \rightarrow \infty$. By (3) then the sequence $(u_k(x_k) \Delta u_k)$ is bounded in L^2 and it follows that $u_k \rightarrow 0$ uniformly as $k \rightarrow \infty$ contradicting our assumption that $u_k(x_k) \rightarrow \infty$. Therefore $r_k \rightarrow 0$ as $k \rightarrow \infty$.

Suppose that we already have determined points $x_k^{(1)}, \dots, x_k^{(i-1)}$ such that (36) and (39) hold and let $x_k = x_k^{(i)} \in \Omega$ be such that

$$(41) \quad \lambda_k \inf_{j < i} |x_k - x_k^{(j)}|^2 u_k^2(x_k) e^{u_k^2(x_k)} = \sup_{x \in \Omega} \left(\lambda_k \inf_{j < i} |x - x_k^{(j)}|^2 u_k^2(x) e^{u_k^2(x)} \right) \rightarrow \infty$$

as $k \rightarrow \infty$. If no such $x_k = x_k^{(i)}$ exists the induction terminates, establishing (40).

Choose $r_k = r_k^{(i)} \rightarrow 0$ satisfying (35). In view of (41) we have $|x_k - x_k^{(j)}|/r_k \rightarrow \infty$ for all $j < i$; that is, half of (39). Moreover, denoting as $v_k(x) = u_k(x_k + r_k x)$ the scaled function u_k on the domain

$$\Omega_k = \{x; x_k + r_k x \in \Omega\},$$

with error $o(1) \rightarrow 0$ as $k \rightarrow \infty$ for any $L > 0$ we can estimate

$$(42) \quad \sup_{x \in \Omega_k, |x| \leq L} v_k^2(x) e^{v_k^2(x)} \leq (1 + o(1)) v_k^2(0) e^{v_k^2(0)} = (1 + o(1)) u_k^2(x_k) e^{u_k^2(x_k)}.$$

Let $\eta_k(x) = \eta_k^{(i)}(x)$ be defined as in (36). Also denoting as $\dot{v}_k(x) = \dot{u}_k(x_k + r_k x)$ the scaled function $\dot{u}_k = u_t(t_k)$, then we have

$$-\Delta \eta_k = \lambda_k r_k^2 v_k(0) v_k e^{v_k^2} - r_k^2 \dot{v}_k v_k(0) e^{v_k^2} =: I_k + II_k \text{ on } \Omega_k.$$

Observe that for any $L > 0$ the bound (42) implies the uniform estimate

$$(43) \quad \begin{aligned} 0 < I_k &= \lambda_k r_k^2 v_k(0) v_k e^{v_k^2} \leq \lambda_k r_k^2 \sup\{v_k^2(0) e^{v_k^2(0)}, v_k^2 e^{v_k^2}\} \\ &\leq (1 + o(1)) \lambda_k r_k^2 v_k^2(0) e^{v_k^2(0)} = (4 + o(1)) \text{ on } B_L(0); \end{aligned}$$

moreover, with (35) and Lemma 4.1 for the second term we have

$$(44) \quad \begin{aligned} \int_{\Omega_k \cap B_L(0)} |II_k|^2 dx &\leq (1 + o(1)) \lambda_k r_k^2 v_k^2(0) e^{v_k^2(0)} \left(\lambda_k^{-1} \int_{\Omega_k \cap B_L(0)} r_k^2 \dot{v}_k^2 e^{v_k^2} dx \right) \\ &= (4 + o(1)) \lambda_k^{-1} \int_{\Omega \cap B_{Lr_k}(x_k)} |u_t(t_k)|^2 e^{u_k^2} dx \rightarrow 0 \end{aligned}$$

with error $o(1) \rightarrow 0$ as $k \rightarrow \infty$.

Note that (41) forces $v_k(0) \rightarrow \infty$. Since (42) also implies the bound

$$(45) \quad 2\eta_k = v_k^2 - v_k^2(0) - (v_k - v_k(0))^2 \leq o(1) \text{ on } \Omega_k \cap B_L(0),$$

it follows that

$$\text{dist}(0, \partial\Omega_k) = \frac{\text{dist}(x_k, \partial\Omega)}{r_k} \rightarrow \infty.$$

Otherwise, by (43) - (45), the mean value property of harmonic functions and the fact that $\eta_k \rightarrow -\infty$ on $\partial\Omega_k$ as $k \rightarrow \infty$ we have locally uniform convergence $\eta_k \rightarrow -\infty$ in Ω_k , which contradicts the fact that $\eta_k(0) = 0$. By the same reasoning we also may assume that as $k \rightarrow \infty$ a subsequence $\eta_k \rightarrow \eta_\infty$ in H_{loc}^2 and locally uniformly. Recalling that $v_k(0) \rightarrow \infty$, then we also have

$$(46) \quad (v_k - v_k(0)) \rightarrow 0, \quad \rho_k := \frac{v_k}{v_k(0)} \rightarrow 1, \quad a_k := 1 + \frac{\eta_k}{2v_k^2(0)} \rightarrow 1$$

locally uniformly. Observing that $e^{v_k^2 - v_k^2(0)} = e^{2a_k \eta_k}$ and using (35), we conclude

$$I_k = \lambda_k r_k^2 v_k(0) v_k e^{v_k^2} = 4\rho_k e^{2a_k \eta_k} \rightarrow 4e^{2\eta_\infty}$$

locally uniformly. Thus, η_∞ solves (37); moreover, for any $L > 1$ by (8) or (22) we have

$$4 \int_{B_L(0)} e^{2\eta_\infty} dx = \lim_{k \rightarrow \infty} \int_{B_L(0)} 4\rho_k^2 e^{2a_k \eta_k} dx = \lim_{k \rightarrow \infty} \int_{B_{Lr_k}(x_k)} \lambda_k u_k^2 e^{u_k^2} dx \leq \Lambda_0.$$

By Fatou's lemma, upon letting $L \rightarrow \infty$ we find $\int_{\mathbb{R}^2} e^{2\eta_\infty} dx < \infty$. In view of the equation $\eta(0) = \lim_{k \rightarrow \infty} \eta_k(0) = 0$ together with (45), the classification of Chen-Li [4] then yields that $\eta_\infty = \eta - \log 2 = \eta_0$, as claimed, which completes the induction step. In view of (38) the induction must terminate when $i > \Lambda_0/\Lambda_1$.

Finally, to see the asserted local H^2 -convergence away from x_i , $1 \leq i \leq i_*$, observe that by (40) and estimates similar to (43), (44) for any x_0 with

$$\inf_{1 \leq i \leq i_*} |x_0 - x_k^{(i)}| \geq 3R_0 > 0$$

the sequence (Δu_k) is bounded in $L^2(B_{2R_0}(x_0))$. Boundedness of (u_k) on $B_{R_0}(x_0)$ and convergence $u_k \rightarrow u_\infty$ in $H^2(B_{R_0}(x_0))$ then follow from boundedness of $(E(u_k))$ and elliptic regularity. \square

5. QUANTIZATION

Throughout this section we continue to assume that $\limsup_{t \rightarrow \infty} \|u(t)\|_{L^\infty} = \infty$ and for a sequence (t_k) as determined in Lemma 4.1 we let $u_k = u(t_k) \rightharpoonup u_\infty$ weakly in $H_0^1(\Omega)$ as $k \rightarrow \infty$, and $\dot{u}_k = u_t(t_k)$ as above. By (8) or (22), respectively, with error $o(1) \rightarrow 0$ there holds

$$(47) \quad \int_{\Omega} |\nabla u_k|^2 dx = (1 + o(1)) \lambda_k \int_{\Omega} u_k^2 e^{2u_k^2} dx \rightarrow \Lambda$$

for some $\Lambda < \infty$. By Theorem 4.2, moreover, we may assume that

$$|\nabla u(t_k)|^2 dx \xrightarrow{w^*} |\nabla u_\infty|^2 dx + \sum_{i=1}^{i_*} \Lambda^{(i)} \delta_{x^{(i)}}$$

weakly in the sense of measures, where $\Lambda^{(i)} \geq \Lambda_1 = 4\pi$ on account of (38). Similar to [6] and [14] we now obtain the following quantization result for the ‘‘defect’’ $\Lambda^{(i)}$ at each $x^{(i)}$.

Theorem 5.1. *We have $\Lambda^{(i)} = 4\pi l_i = l_i \Lambda_1$ for some $l_i \in \mathbb{N}$, $1 \leq i \leq i_*$.*

For the proof we argue as in [14]. We first consider the radial case.

5.1. The radial case. Let $\Omega = B_R(0) =: B_R$ and assume that $u(t, x) = u(t, |x|)$. In this case by Theorem 4.2 for any $i \leq i_*$ we have $r_k^{-1}x_k \rightarrow 0$ as $k \rightarrow \infty$, where $x_k = x_k^{(i)}$ and $r_k = r_k^{(i)}$ is given by (35); otherwise, the blow-up limit $\eta_0 = \lim_{k \rightarrow \infty} \eta_k^{(i)}$ could not be radially symmetric. In particular, from (39) it follows that $i_* = 1$; moreover, by (36) we have $u_k^2(x_k) = \sup_{\Omega} u_k^2 = u_k^2(0) + o(1)$. Thus, up to an error $o(1) \rightarrow 0$ locally uniformly as $k \rightarrow \infty$ we may replace the original function $\eta_k = \eta_k^{(1)}$ defined in (36) by the function

$$\eta_k(x) = u_k(0)(u_k(r_k x) - u_k(0)).$$

Observe that by radial symmetry or Theorem 4.2 we also have convergence $u_k \rightarrow u_{\infty}$ locally uniformly away from $x = 0$ as $k \rightarrow \infty$.

For $|x| = r$ let $u_k(r) = u_k(x)$ and set

$$\lambda_k u_k^2 e^{u_k^2} =: e_k \text{ in } \Omega.$$

We also denote as

$$w_k(x) = u_k(0)(u_k(x) - u_k(0))$$

the unscaled function η_k , satisfying the equation

$$-\Delta w_k = \lambda_k u_k(0) u_k e^{u_k^2} - d_k,$$

where the term $d_k = u_k(0) \dot{u}_k e^{u_k^2}$ for any $L > 0$ can be estimated

$$(48) \quad \int_{B_{Lr_k}} |d_k| dx \leq \sup_{B_{Lr_k}} \left(\frac{u_k(0)}{u_k} \right) \left(\lambda_k \int_{B_{Lr_k}} u_k^2 e^{u_k^2} dx \cdot \lambda_k^{-1} \int_{B_{Lr_k}} \dot{u}_k^2 e^{u_k^2} dx \right)^{1/2}.$$

Hence by Theorem 4.2, Lemma 4.1, and (47) we conclude that $d_k \rightarrow 0$ in $L^1(B_{Lr_k})$ for any $L > 0$ as $k \rightarrow \infty$. Finally, we set

$$\lambda_k u_k(0) u_k e^{u_k^2} =: f_k \text{ in } \Omega = B_R$$

and for $0 < r < R$ let

$$\Lambda_k(r) = \int_{B_r} e_k dx, \quad \sigma_k(r) = \int_{B_r} f_k dx,$$

Observe that with error $o(1) \rightarrow 0$ as $k \rightarrow \infty$ we have $e_k \leq (1 + o(1))f_k$, $\Lambda_k(r) \leq \sigma_k(r) + o(1)$; moreover, Theorem 4.2 implies

$$(49) \quad \lim_{L \rightarrow \infty} \lim_{k \rightarrow \infty} \Lambda_k(Lr_k) = \lim_{L \rightarrow \infty} \lim_{k \rightarrow \infty} \sigma_k(Lr_k) = \lim_{L \rightarrow \infty} 4 \int_{B_L} e^{2\eta_0} dx = \Lambda_1.$$

We can now show our first decay estimate. Let $u'_k = \frac{\partial u_k}{\partial r}$, and so on.

Lemma 5.2. *For any $0 < \varepsilon < 1$, letting $T_k > 0$ be minimal such that $u_k(T_k) = \varepsilon u_k(0)$, for any constant $b < 2$ and sufficiently large k there holds*

$$w_k(r) \leq b \log \left(\frac{r_k}{r} \right) \text{ on } B_{T_k}$$

and we have

$$\lim_{k \rightarrow \infty} \Lambda_k(T_k) = \lim_{k \rightarrow \infty} \sigma_k(T_k) = \Lambda_1 = 4\pi.$$

Proof. Note that $T_k \rightarrow 0$ as $k \rightarrow \infty$ in view of the locally uniform convergence $u_k \rightarrow u_\infty$ away from 0.

Since $u_k(t) \geq \varepsilon u_k(0)$ for $Lr_k \leq t \leq T_k$, from (49) and an estimate similar to (48) for all such $t = t_k$ we obtain

$$(50) \quad \begin{aligned} 2\pi t w'_k(t) &= \int_{\partial B_t} \partial_\nu w_k \, do = \int_{B_t} \Delta w_k \, dx \\ &= -\sigma_k(t) + o(1) \leq -\Lambda_1 + o(1) \end{aligned}$$

with error $o(1) \rightarrow 0$ uniformly in t , if first $k \rightarrow \infty$ and then $L \rightarrow \infty$. For any $b < 2$ and sufficiently large $L \geq L(b)$, for $k \geq k_0(L)$ we thus obtain that

$$w'_k(t) \leq -\frac{b}{t}$$

for all $Lr_k \leq t \leq T_k$. Since $\eta_0(L) < -b \log L$ for all $L > 0$, in view of Theorem 4.2 clearly we may choose $k_0(L)$ such that $\eta_k(L) < -b \log L$ for all $k \geq k_0(L)$. For any such k and any $r \in [Lr_k, T_k]$, upon integrating from Lr_k to r then we find

$$(51) \quad \begin{aligned} w_k(r) &\leq w_k(Lr_k) - b \log \left(\frac{r}{Lr_k} \right) \\ &= \eta_k(L) + b \log L + b \log \left(\frac{r_k}{r} \right) \leq b \log \left(\frac{r_k}{r} \right), \end{aligned}$$

as claimed. For $r \leq Lr_k$ the asserted bound already follows from Theorem 4.2.

Inserting (51) in the definition of f_k and recalling (35), for $Lr_k \leq r \leq T_k$ with sufficiently large $L > 0$ and $k \geq k_0(L)$ then we obtain

$$\begin{aligned} f_k &= \lambda_k (u_k^2(0) + w_k) e^{u_k^2(0)} e^{2(1 + \frac{w_k}{2u_k^2(0)})w_k} \\ &\leq \lambda_k r_k^2 u_k^2(0) e^{u_k^2(0)} r_k^{-2} e^{(1+\varepsilon)w_k} \leq 4r_k^{-2} \left(\frac{r_k}{r} \right)^{(1+\varepsilon)b}. \end{aligned}$$

Choosing $b < 2$ such that $(1 + \varepsilon)b = 2 + \varepsilon$, upon integrating over B_{T_k} we obtain

$$\begin{aligned} \sigma_k(T_k) &= \int_{B_{T_k}} f_k \, dx \leq \Lambda_1 + \int_{B_{T_k} \setminus B_{Lr_k}} f_k \, dx \\ &\leq \Lambda_1 + C r_k^{-2} \int_{B_{T_k} \setminus B_{Lr_k}} \left(\frac{r_k}{r} \right)^{2+\varepsilon} \, dx \leq \Lambda_1 + C \varepsilon^{-1} \left(\frac{r_k}{Lr_k} \right)^\varepsilon \leq \Lambda_1 + \varepsilon, \end{aligned}$$

if first $L > L_0(\varepsilon)$ and then $k \geq k_0(L)$ is chosen sufficiently large. Since $\varepsilon > 0$ is arbitrary, the proof is complete. \square

If we now choose $\varepsilon_k \downarrow 0$ such that with $s_k = T_k(\varepsilon_k)$ we have $u_k(s_k) \rightarrow \infty$, by Theorem 4.2 we also have $r_k/s_k \rightarrow 0$, $s_k \rightarrow 0$ as $k \rightarrow \infty$. That is, we can achieve that

$$(52) \quad \lim_{k \rightarrow \infty} \Lambda_k(s_k) = \Lambda_1, \quad \lim_{k \rightarrow \infty} \frac{u_k(s_k)}{u_k(r_k)} = \lim_{k \rightarrow \infty} \frac{r_k}{s_k} = \lim_{k \rightarrow \infty} s_k = 0.$$

In addition, from (49) we obtain that

$$(53) \quad \lim_{L \rightarrow \infty} \lim_{k \rightarrow \infty} (\Lambda_k(s_k) - \Lambda_k(Lr_k)) = 0.$$

Let $r_k = r_k^{(1)}$, $s_k = s_k^{(1)}$. We now proceed by iteration. Suppose that for some integer $l \geq 1$ we already have determined numbers $r_k^{(1)} < s_k^{(1)} < \dots < r_k^{(l)} < s_k^{(l)}$ such that

$$(54) \quad \lim_{k \rightarrow \infty} \Lambda_k(s_k^{(l)}) = l\Lambda_1$$

and

$$(55) \quad \lim_{L \rightarrow \infty} \lim_{k \rightarrow \infty} (\Lambda_k(s_k^{(l)}) - \Lambda_k(Lr_k^{(l)})) = \lim_{k \rightarrow \infty} \frac{u_k(s_k^{(l)})}{u_k(r_k^{(l)})} = \lim_{k \rightarrow \infty} \frac{r_k^{(l)}}{s_k^{(l)}} = \lim_{k \rightarrow \infty} s_k^{(l)} = 0.$$

For $0 < s < t < R$ let

$$N_k(s, t) = \int_{B_t \setminus B_s} e_k \, dx = \int_{B_t \setminus B_s} \lambda_k u_k^2 e^{u_k^2} \, dx = 2\pi \int_s^t \lambda_k r u_k^2 e^{u_k^2} \, dr$$

and define

$$P_k(t) = t \frac{\partial}{\partial t} N_k(s, t) = t \int_{\partial B_t} e_k \, d\sigma = 2\pi \lambda_k t^2 u_k^2(t) e^{u_k^2(t)}.$$

Note that (40) implies the uniform bound $P_k \leq C$; moreover, with a uniform constant C_0 for any t we have

$$(56) \quad \inf_{t/2 \leq t' \leq t} P_k(t') \leq C_0 N_k(t/2, t).$$

A preliminary quantization now can be achieved, as follows.

Lemma 5.3. *i) Suppose that for some $t_k > s_k^{(l)}$ there holds*

$$\sup_{s_k^{(l)} < t < t_k} P_k(t) \rightarrow 0 \text{ as } k \rightarrow \infty.$$

Then we have

$$\lim_{k \rightarrow \infty} N_k(s_k^{(l)}, t_k) = 0.$$

ii) Conversely, if for some $t_k > s_k^{(l)}$ and a subsequence (u_k) there holds

$$\lim_{k \rightarrow \infty} N_k(s_k^{(l)}, t_k) = \nu_0 > 0, \quad \lim_{k \rightarrow \infty} t_k = 0,$$

then either $\nu_0 \geq \pi$, or we have

$$\liminf_{k \rightarrow \infty} P_k(t_k) \geq \nu_0$$

and

$$\lim_{L \rightarrow \infty} \liminf_{k \rightarrow \infty} N_k(s_k^{(l)}, Lt_k) \geq \pi, \quad \lim_{L \rightarrow \infty} \limsup_{k \rightarrow \infty} N_k(s_k^{(l)}, t_k/L) = 0.$$

Proof. *i) For $s = s_k^{(l)} < t$ we integrate by parts to obtain*

$$(57) \quad \begin{aligned} 2N_k(s, t) &= \int_{B_t \setminus B_s} e_k \operatorname{div} x \, dx = P_k(t) - P_k(s) - \int_{B_t \setminus B_s} x \cdot \nabla e_k \, dx \\ &\leq P_k(t) - 4\pi \int_s^t \lambda_k r^2 u_k'(1 + u_k^2) u_k e^{u_k^2} \, dr. \end{aligned}$$

In order to further estimate the right hand side we observe that (3) for any $t < R$ yields the identity

$$(58) \quad -2\pi t u_k(t) u'_k(t) = \int_{B_t} \lambda_k u_k(t) u_k e^{u_k^2} dx - \int_{B_t} u_k(t) \dot{u}_k e^{u_k^2} dx.$$

Estimating $u_k^2(t) e^{u_k^2} \leq \max\{u_k^2(t) e^{u_k^2(t)}, u_k^2 e^{u_k^2}\}$, by Lemma 4.1, (40), and (47), we can easily bound the contribution from the second integral

$$(59) \quad \begin{aligned} \left(\int_{B_t} u_k(t) |\dot{u}_k| e^{u_k^2} dx \right)^2 &\leq \lambda_k \int_{B_t} u_k^2(t) e^{u_k^2} dx \cdot \lambda_k^{-1} \int_{B_t} \dot{u}_k^2 e^{u_k^2} dx \\ &\leq o(1) (\pi \lambda_k t^2 u_k^2(t) e^{u_k^2(t)} + \lambda_k \int_{B_t} u_k^2 e^{u_k^2} dx) = o(1), \end{aligned}$$

where $o(1) \rightarrow 0$ as $k \rightarrow \infty$. From (58) we then obtain that at any sequence of points $t = t_k$ where $u'_k(t) \geq 0$ there holds

$$(60) \quad \int_{B_t} \lambda_k u_k(t) u_k e^{u_k^2} dx = o(1).$$

On the other hand, if for $t_{k0} = t_0 \leq r \leq t = t_k$ there holds $u'_k(r) \leq 0 = u'_k(t_0)$, by (60) we can estimate

$$(61) \quad \begin{aligned} \int_{B_t} \lambda_k u_k(t) u_k e^{u_k^2} dx &\leq \int_{B_t \setminus B_{t_0}} \lambda_k u_k^2 e^{u_k^2} dx + \int_{B_{t_0}} \lambda_k u_k(t_0) u_k e^{u_k^2} dx \\ &= N_k(t_0, t) + o(1). \end{aligned}$$

In view of (59)-(61) and (55), for $s = s_k^{(l)} \leq r \leq t = t_k$ and with $r_k = r_k^{(l)}$ we then can estimate

$$(62) \quad \begin{aligned} -2\pi r u_k(r) u'_k(r) &= \int_{B_r} \lambda_k u_k(r) u_k e^{u_k^2} dx + o(1) \\ &\leq N_k(s, r) + \int_{B_s} \lambda_k u_k(s) u_k e^{u_k^2} dx + o(1) \\ &\leq N_k(s, r) + N_k(Lr_k, s) + \frac{u_k(s)}{u_k(Lr_k)} \Lambda_k(Lr_k) + o(1) \\ &= N_k(s, r) + o(1), \end{aligned}$$

where $o(1) \rightarrow 0$ when first $k \rightarrow \infty$ and then $L \rightarrow \infty$. Indeed, the first inequality is clear when $u'_k \leq 0$ in $[s, r]$, and otherwise follows from (60), (61). The second inequality may be seen in a similar way. Recalling (57) we thus arrive at the estimate

$$(63) \quad \begin{aligned} 2N_k(s, t) &\leq P_k(t) + 2 \int_s^t \lambda_k r (1 + u_k^2) e^{u_k^2} N_k(s, r) dr + o(1) \\ &\leq P_k(t) + \pi^{-1} N_k(s, t)^2 + o(1). \end{aligned}$$

If we now assume that

$$\sup_{s < t < t_k} P_k(t) \rightarrow 0 \text{ as } k \rightarrow \infty,$$

upon letting t increase from $t = s = s_k^{(l)}$ to t_k we find

$$\lim_{k \rightarrow \infty} N_k(s_k^{(l)}, t_k) = 0,$$

as claimed.

ii) On the other hand, if we suppose that for some $t_k > s_k^{(l)}$ we have

$$(64) \quad 0 < \lim_{k \rightarrow \infty} N_k(s_k^{(l)}, t_k) = \nu_0 < \pi,$$

from (63) with error $o(1) \rightarrow 0$ as $k \rightarrow \infty$ we conclude that

$$(65) \quad \nu_0 + o(1) \leq (2 - \nu_0/\pi)N_k(s_k^{(l)}, t_k) \leq P_k(t_k) + o(1).$$

It then also follows that

$$\lim_{L \rightarrow \infty} \liminf_{k \rightarrow \infty} N_k(s_k^{(l)}, Lt_k) \geq \pi.$$

Otherwise, (56) and (65) for a subsequence (u_k) yield the uniform bound

$$C_0 \liminf_{k \rightarrow \infty} N_k(Lt_k/2, Lt_k) \geq \liminf_{k \rightarrow \infty} \inf_{Lt_k/2 \leq t \leq Lt_k} P_k(t) \geq \nu_0$$

for all $L \geq 2$. Choosing $L = 2^m$, where $m \in \mathbb{N}$, and summing over $1 \leq m \leq M$, we obtain

$$C_0 \liminf_{k \rightarrow \infty} \Lambda_k(2^M t_k) \geq C_0 \liminf_{k \rightarrow \infty} N_k(t_k, 2^M t_k) \geq \nu_0 M \rightarrow \infty \text{ as } M \rightarrow \infty,$$

contrary to assumption (47). Upon replacing t_k by t_k/L in the previous argument and recalling our assumption (64), by the same reasoning we also arrive at the estimate

$$\lim_{L \rightarrow \infty} \liminf_{k \rightarrow \infty} N_k(s_k^{(l)}, t_k/L) = 0.$$

This completes the proof. \square

Suppose that for some $t_k > s_k^{(l)}$ with $t_k \rightarrow 0$ as $k \rightarrow \infty$ there holds

$$\liminf_{k \rightarrow \infty} N_k(s_k^{(l)}, t_k) > 0.$$

Then we can find a subsequence (u_k) and numbers $r_k^{(l+1)} \in]s_k^{(l)}, t_k[$ such that

$$(66) \quad \lim_{k \rightarrow \infty} N_k(s_k^{(l)}, r_k^{(l+1)}) = \nu_0 > 0.$$

Replacing our original choice of $r_k^{(l+1)}$ by a smaller number, if necessary, we may assume that $\nu_0 < \pi$. Lemma 5.3 then implies that

$$(67) \quad \lim_{L \rightarrow \infty} \liminf_{k \rightarrow \infty} N_k(s_k^{(l)}, Lr_k^{(l+1)}) \geq \pi, \quad \lim_{L \rightarrow \infty} \limsup_{k \rightarrow \infty} N_k(s_k^{(l)}, r_k^{(l+1)}/L) = 0,$$

and that

$$(68) \quad \liminf_{k \rightarrow \infty} P_k(r_k^{(l+1)}) > 0.$$

In particular, since $r_k^{(l+1)} \leq t_k \rightarrow 0$ we then conclude that $u_k(r_k^{(l+1)}) \rightarrow \infty$.

The desired precise quantization result at the scale $r_k^{(l+1)}$ is a consequence of the following Proposition.

Proposition 5.4. *There exist a subsequence (u_k) such that*

$$\eta_k^{(l+1)}(x) := u_k(r_k^{(l+1)})(u_k(r_k^{(l+1)}(x) - u_k(r_k^{(l+1)})) \rightarrow \eta(x)$$

locally uniformly on $\mathbb{R}^2 \setminus \{0\}$ as $k \rightarrow \infty$, where $\eta(x) = \log(\frac{2}{1+|x|^2})$.

Postponing the details of the proof of Proposition 5.4 to the next section, we now complete the proof of Theorem 5.1.

Denote as $v_k^{(l+1)}(x) = u_k(r_k^{(l+1)}x)$, $\dot{v}_k^{(l+1)}(x) = \dot{u}_k(r_k^{(l+1)}x)$ the scaled functions u_k and \dot{u}_k , respectively. Omitting the superscript $(l+1)$ for brevity, similar to the proof of Theorem 4.2 for $\eta_k := \eta_k^{(l+1)}$ we have

$$-\Delta\eta_k = \lambda_k r_k^2 v_k(1) v_k e^{v_k^2} - r_k^2 \dot{v}_k v_k(1) e^{v_k^2} =: I_k + II_k,$$

where $II_k \rightarrow 0$ in $L_{loc}^2(\mathbb{R}^2 \setminus \{0\})$ as $k \rightarrow \infty$. Moreover, letting $\rho_k = \rho_k^{(l+1)} := \frac{v_k}{v_k(1)}$, $a_k = a_k^{(l+1)} = 1 + \frac{\eta_k}{2v_k^2(1)}$, by Proposition 5.4 we have $a_k \rightarrow 1$, $\rho_k \rightarrow 1$ as $k \rightarrow \infty$ locally uniformly away from $x = 0$, and

$$I_k = \lambda_k r_k^2 v_k(1) v_k e^{v_k^2} = \lambda_k r_k^2 v_k^2(1) e^{v_k^2(1)} \rho_k e^{v_k^2 - v_k^2(1)} = (2\pi)^{-1} P_k(r_k) \rho_k e^{2a_k \eta_k}.$$

Now observe that η solves equation (33) on \mathbb{R}^2 with

$$\int_{\mathbb{R}^2} e^{2\eta} dx = 4\pi = \Lambda_1.$$

We therefore conclude that $P_k(r_k) \rightarrow 2\pi$ and

(69)

$$\begin{aligned} \lim_{L \rightarrow \infty} \lim_{k \rightarrow \infty} N_k(r_k^{(l+1)}/L, Lr_k^{(l+1)}) &= \lim_{L \rightarrow \infty} \lim_{k \rightarrow \infty} \int_{B_L \setminus B_{1/L}} \lambda_k r_k^2 v_k^2 e^{v_k^2} dx \\ &= \lim_{L \rightarrow \infty} \lim_{k \rightarrow \infty} \int_{B_L \setminus B_{1/L}} (2\pi)^{-1} P_k(r_k) \rho_k^2 e^{2a_k \eta_k} dx = \lim_{L \rightarrow \infty} \int_{B_L \setminus B_{1/L}} e^{2\eta} dx = \Lambda_1. \end{aligned}$$

From (67) then we obtain that

$$\begin{aligned} \lim_{L \rightarrow \infty} \lim_{k \rightarrow \infty} N_k(s_k^{(l)}, Lr_k^{(l+1)}) \\ = \lim_{L \rightarrow \infty} \lim_{k \rightarrow \infty} (N_k(s_k^{(l)}, r_k^{(l+1)}/L) + N_k(r_k^{(l+1)}/L, Lr_k^{(l+1)})) = \Lambda_1, \end{aligned}$$

and our induction hypothesis (54) yields

$$(70) \quad \lim_{L \rightarrow \infty} \lim_{k \rightarrow \infty} \Lambda_k(Lr_k^{(l+1)}) = \lim_{L \rightarrow \infty} \lim_{k \rightarrow \infty} (\Lambda_k(s_k^{(l)}) + N_k(s_k^{(l)}, Lr_k^{(l+1)})) = (l+1)\Lambda_1.$$

Moreover, $r_k^{(l+1)}/s_k^{(l)} \rightarrow \infty$ as $k \rightarrow \infty$. Indeed, if we assume that $r_k^{(l+1)} \leq Ls_k^{(l)}$ for some L by Proposition 5.4 we have $N_k(s_k^{(l)}/2, s_k^{(l)}) \geq \nu_0$ for some constant $\nu_0 = \nu_0(L) > 0$, contradicting (55).

In order to obtain decay analogous to Lemma 5.2 and then also the analogue of (55) at the scale $r_k^{(l+1)}$, denote as

$$w_k^{(l+1)}(x) = u_k(r_k^{(l+1)})(u_k(x) - u_k(r_k^{(l+1)}))$$

the unscaled function $\eta_k^{(l+1)}$, satisfying the equation

$$(71) \quad -\Delta w_k^{(l+1)} = \lambda_k u_k(r_k^{(l+1)}) u_k e^{u_k^2} - u_k(r_k^{(l+1)}) \dot{u}_k e^{u_k^2} =: f_k^{(l+1)} - d_k^{(l+1)}$$

in $\Omega = B_R$. We then have the analogue of Lemma 5.2, which may be proved in the same fashion.

Lemma 5.5. *For any $\varepsilon > 0$, letting $T_k = T_k^{(l+1)} > r_k^{(l+1)}$ be minimal such that $u_k(T_k) = \varepsilon u_k(r_k^{(l+1)})$, for any constant $b < 2$ and sufficiently large k and L there holds*

$$w_k^{(l+1)}(r) \leq b \log \left(\frac{T_k^{(l+1)}}{r} \right) \text{ on } B_{T_k} \setminus B_{Lr_k^{(l+1)}}$$

and we have

$$\lim_{k \rightarrow \infty} N_k(s_k^{(l)}, T_k) = \Lambda_1.$$

Proof. Denote $w_k^{(l+1)} = w_k$, $r_k^{(l+1)} = r_k$, $d_k^{(l+1)} = d_k$ for simplicity. Coupled with the uniform bound $u_k(t) \geq \varepsilon u_k(r_k)$ for $r_k \leq t \leq T_k$, the estimate (59) yields decay of $\int_{B_{T_k}} |d_k| dx$. Thus, for $Lr_k \leq t = t_k \leq T_k$ from (69) and Proposition 5.4 we have

$$(72) \quad \begin{aligned} 2\pi t w_k'(t) &= \int_{\partial B_t} \partial_\nu w_k \, d\sigma = \int_{B_t} \Delta w_k \, dx \leq - \int_{B_{Lr_k}} \frac{u_k(r_k)}{u_k} e_k \, dx + o(1) \\ &\leq -N_k(r_k/L, Lr_k) + o(1) \leq -\Lambda_1 + o(1), \end{aligned}$$

with error $o(1) \rightarrow 0$ uniformly in t , if first $k \rightarrow \infty$ and then $L \rightarrow \infty$. For any $b < 2$ and sufficiently large $L \geq L(b)$ for $k \geq k_0(L)$, we thus obtain that

$$w_k'(t) \leq -\frac{b}{t}$$

for all $Lr_k \leq t \leq T_k$. For such t it then follows that

$$\begin{aligned} e_k &\leq \lambda_k u_k^2(r_k) e^{u_k^2(r_k)} e^{2(1 + \frac{w_k}{2u_k^2(r_k)})w_k} \\ &\leq (2\pi)^{-1} P(r_k) r_k^{-2} e^{(1+\varepsilon)w_k} \leq C r_k^{-2} \left(\frac{r_k}{r} \right)^{(1+\varepsilon)b}, \end{aligned}$$

and the proof may be completed as in Lemma 5.2. \square

For suitable numbers $s_k^{(l+1)} = T_k^{(l+1)}(\varepsilon_k)$, where $\varepsilon_k \downarrow 0$ is chosen such that $u_k(s_k^{(l+1)}) = \varepsilon_k u_k(r_k^{(l+1)}) \rightarrow \infty$ as $k \rightarrow \infty$, then we have

$$(73) \quad \lim_{k \rightarrow \infty} \Lambda_k(s_k^{(l+1)}) = (l+1)\Lambda_1$$

and

$$(74) \quad \begin{aligned} &\lim_{L \rightarrow \infty} \lim_{k \rightarrow \infty} (\Lambda_k(s_k^{(l+1)}) - \Lambda_k(Lr_k^{(l+1)})) \\ &= \lim_{k \rightarrow \infty} \frac{r_k^{(l+1)}}{s_k^{(l+1)}} = \lim_{k \rightarrow \infty} \frac{u_k(s_k^{(l+1)})}{u_k(r_k^{(l+1)})} = \lim_{k \rightarrow \infty} s_k^{(l+1)} = 0, \end{aligned}$$

completing the induction step. In view of (47) and Lemma 5.3 the iteration must terminate after finitely many steps $1 \leq l \leq l_*$, after which

$$N_k(s_k^{(l_*)}, t_k) \rightarrow 0 \text{ as } k \rightarrow \infty$$

for any sequence $t_k \rightarrow 0$ as $k \rightarrow \infty$. This concludes the proof of Theorem 5.1 in the radial case.

5.2. Proof of Proposition 5.4. Throughout this section we let $r_k = r_k^{(l+1)}$, etc., and we set $r_k^- = r_k^{(l)}$, $s_k^- = s_k^{(l)}$. Again denote as $v_k(x) = u_k(r_k x)$, $\dot{v}_k(x) = \dot{u}_k(r_k x)$ the scaled functions u_k , \dot{u}_k , respectively. As usual we write $v_k(x) = v_k(r)$ for $r = |x|$. Recall that (68) implies that $v_k(1) = u_k(r_k) \rightarrow \infty$.

Lemma 5.6. *As $k \rightarrow \infty$ we have $v_k(x) - v_k(1) \rightarrow 0$ locally uniformly on $\mathbb{R}^2 \setminus \{0\}$.*

Proof. The function $\tilde{v}_k(x) = v_k(x) - v_k(1)$ satisfies the equation

$$-\Delta \tilde{v}_k = g_k - l_k,$$

where $g_k = \lambda_k r_k^2 v_k e^{v_k^2}$ and with $l_k = r_k^2 \dot{v}_k e^{v_k^2}$.

We claim that $g_k \rightarrow 0$ locally uniformly away from 0. Indeed, since $r_k \rightarrow 0$, for any x where $g_k(x) \geq r_k$ we have $v_k(x) = u_k(r_k x) \geq \gamma_k$ with constants $\gamma_k \rightarrow \infty$ independent of x . Hence for any $L > 0$ and any $1/L \leq |x| \leq L$ we either can bound $g_k(x) \leq r_k \rightarrow 0$, or

$$\begin{aligned} g_k(x) &= \lambda_k r_k^2 v_k(x) e^{v_k^2(x)} = \lambda_k r_k^2 u_k(r_k x) e^{u_k^2(r_k x)} \\ &= (2\pi)^{-1} |x|^{-2} P_k(r_k |x|) / u_k(r_k x) \leq CL^2 \gamma_k^{-1} \rightarrow 0 \end{aligned}$$

as $k \rightarrow \infty$. Moreover, (40) and Lemma 4.1 imply

$$(75) \quad \begin{aligned} \int_{B_L \setminus B_{1/L}(0)} |l_k|^2 dx &\leq \lambda_k r_k^2 \sup_{1/L \leq |x| \leq L} e^{v_k^2(x)} \left(\lambda_k^{-1} \int_{B_{Lr_k}(x_k)} \dot{u}_k^2 e^{u_k^2} dx \right) \\ &\leq (2\pi)^{-1} L^2 \sup_{1/L \leq |x| \leq L} \frac{P_k(r_k |x|)}{u_k^2(r_k |x|)} \left(\lambda_k^{-1} \int_{\Omega} \dot{u}_k^2 e^{u_k^2} dx \right) \rightarrow 0 \end{aligned}$$

for any fixed $L > 1$ as $k \rightarrow \infty$.

Since from (7) or (12), respectively, we also have the uniform L^2 -bound

$$\|\nabla \tilde{v}_k\|_{L^2} = \|\nabla u_k\|_{L^2} \leq C,$$

we may extract a subsequence (u_k) such that $\tilde{v}_k \rightarrow \tilde{v}$ weakly in $H_{loc}^1(\mathbb{R}^2)$, where \tilde{v} is harmonic away from the origin. In addition, $\nabla \tilde{v} \in L^2(\mathbb{R}^2)$; since the point $x = 0$ has vanishing H^1 -capacity, we then have $\Delta \tilde{v} = 0$ in the distribution sense on all of \mathbb{R}^2 and \tilde{v} is a smooth, everywhere harmonic function. Again invoking the fact that $\nabla \tilde{v} \in L^2(\mathbb{R}^2)$, and recalling that $\tilde{v}(1) = \tilde{v}_k(1) = 0$, then we see that \tilde{v} vanishes identically; that is, $\tilde{v}_k \rightarrow 0$ weakly in $H_{loc}^1(\mathbb{R}^2)$.

Recalling that for radially symmetric functions weak H^1 -convergence implies locally uniform convergence away from the origin, we obtain the claim. \square

Now $\eta_k(x) = v_k(1)(v_k(x) - v_k(1))$ satisfies the equation

$$(76) \quad -\Delta \eta_k = \lambda_k r_k^2 v_k(1) v_k e^{v_k^2} - r_k^2 v_k(1) \dot{v}_k e^{v_k^2} =: I_k + II_k.$$

By Lemma 5.6 for any $L > 1$ we can bound $\sup_{B_L \setminus B_{1/L}} v_k(1)/v_k \leq 2$ for sufficiently large k . Lemma 4.1, (47), and (59) then yield

$$(77) \quad \begin{aligned} \int_{B_L} |II_k| dx &\leq \int_{B_1} |II_k| dx + \int_{B_L \setminus B_1} |II_k| dx \\ &\leq o(1) + 2 \left(\lambda_k \int_{B_{Lr_k}} u_k^2 e^{u_k^2} dx \cdot \lambda_k^{-1} \int_{B_{Lr_k}} \dot{u}_k^2 e^{u_k^2} dx \right)^{1/2} \rightarrow 0, \end{aligned}$$

with error $o(1) \rightarrow 0$ as $k \rightarrow \infty$ for any fixed $L > 1$. Upon estimating $v_k(1)v_k e^{v_k^2} \leq \max\{v_k^2(1)e^{v_k^2(1)}, v_k^2 e^{v_k^2}\}$, for $1/L \leq |x| \leq L$ by (40) we can bound the remaining term

$$(78) \quad I_k(x) \leq (2\pi)^{-1} \max\{P_k(r_k), |x|^{-2}P_k(r_k|x|)\} \leq C(1 + L^2)$$

Moreover, letting $\hat{v}_k = v_k/v_k(1) \rightarrow 1$ in $B_L \setminus B_{1/L}$, we have

$$(79) \quad I_k = \lambda_k r_k^2 v_k^2(1) e^{v_k^2(1)} \hat{v}_k e^{v_k^2 - v_k^2(1)} = p_k \hat{v}_k e^{\eta_k(1 + \hat{v}_k)},$$

where $p_k = (2\pi)^{-1}P_k(r_k) \geq p_0 > 0$ by (68).

Finally, similar to (62) and in view of (55) we find

$$(80) \quad \begin{aligned} \int_{B_{1/L}(0)} I_k dx &= \int_{B_{r_k/L}(0)} \lambda_k u_k(r_k) u_k e^{u_k^2} dx \\ &\leq N_k(Lr_k^-, r_k/L) + C\Lambda \frac{u_k(s_k^-)}{u_k(Lr_k^-)} \rightarrow 0, \end{aligned}$$

if we first let $k \rightarrow \infty$ and then pass to the limit $L \rightarrow \infty$.

Lemma 5.7. *There exist a subsequence (u_k) such that $\eta_k \rightarrow \eta_\infty$ locally uniformly on $\mathbb{R}^2 \setminus \{0\}$ as $k \rightarrow \infty$.*

Proof. For any $L > 1$ decompose $\eta_k = h_k + n_k$ on $B_L \setminus B_{1/L}(0)$, where $\Delta h_k = 0$ in $B_L \setminus B_{1/L}(0)$, and where $n_k = 0$ on $\partial(B_L \setminus B_{1/L}(0))$. In view of (77), (78), and passing to a subsequence, if necessary, we may assume that $n_k \rightarrow n$ as $k \rightarrow \infty$ in $W^{1,q}$ on $B_L \setminus B_{1/L}(0)$ for any $q < 2$ and therefore also uniformly by radial symmetry.

On the other hand, letting $h_k^+ = \max\{0, h_k\}$, from (78) - (79) for sufficiently large k we obtain the estimate

$$\begin{aligned} \int_{B_L \setminus B_{1/L}(0)} h_k^+ dx &\leq \int_{B_L \setminus B_{1/L}(0)} (\eta_k^+ + |n_k|) dx \\ &\leq \int_{B_L \setminus B_{1/L}(0)} e^{(1 + \hat{v}_k)\eta_k} dx + C(L) \leq C(L) < \infty. \end{aligned}$$

From the mean value property of harmonic functions and Harnack's inequality we conclude that either $h_k \rightarrow h$ locally uniformly on $B_L \setminus B_{1/L}(0)$, or $h_k \rightarrow -\infty$ and hence $\eta_k \rightarrow -\infty$ locally uniformly on $B_L \setminus B_{1/L}(0)$ as $k \rightarrow \infty$. But the identity $\eta_k(1) = 0$ excludes the latter case, and the assertion follows. \square

Now we can complete the proof of Proposition 5.4. Since $\Delta\eta_k$ by (76) - (80) is uniformly bounded in $L^1(B_L(0))$, the sequence (η_k) is bounded in $W^{1,q}(B_L(0))$ for any $q < 2$ and any $L > 1$ and we may assume that $\eta_k \rightarrow \eta_0$ also weakly locally in $W^{1,q}$ on \mathbb{R}^2 as $k \rightarrow \infty$.

By Lemmas 5.6 and 5.7 we may then pass to the limit $k \rightarrow \infty$ in equation (76) to see that η_∞ solves the equation

$$(81) \quad -\Delta\eta_\infty = p_\infty e^{2\eta_\infty} \text{ on } \mathbb{R}^2 \setminus \{0\},$$

for some constant $p_\infty = \lim_{k \rightarrow \infty} p_k > 0$. Moreover, by Lemma 5.7, and (78) we have

$$p_\infty e^{2\eta_\infty} = \lim_{k \rightarrow \infty} p_k \hat{v}_k^2 e^{\eta_k(\hat{v}_k(x)+1)} = \lim_{k \rightarrow \infty} \hat{v}_k I_k = \lim_{k \rightarrow \infty} r_k^2 e_k(r_k \cdot)$$

locally uniformly on $\mathbb{R}^2 \setminus \{0\}$. Thus, with a uniform constant C for any $L > 1$ we have

$$p_\infty \int_{B_L \setminus B_{1/L}(0)} e^{2\eta_\infty} dx \leq \liminf_{k \rightarrow \infty} \int_{B_{Lr_k} \setminus B_{r_k/L}(0)} e_k dx \leq C\Lambda.$$

Passing to the limit $L \rightarrow \infty$, we see that $e^{2\eta_\infty} \in L^1(\mathbb{R}^2)$. By (77) and (80) we also have

$$\limsup_{k \rightarrow \infty} \int_{B_{1/L}(0)} |\Delta \eta_k| dx \rightarrow 0$$

as $L \rightarrow \infty$. Hence η_∞ extends to a distribution solution of (81) on all of \mathbb{R}^2 . Our claim then follows from the Chen-Li [4] classification of all solutions η_∞ to equation (81) on \mathbb{R}^2 with $e^{2\eta_\infty} \in L^1(\mathbb{R}^2)$ in view of radial symmetry of η_∞ together with the fact that $\eta_\infty(1) = \eta_k(1) = 0$.

5.3. The general case. For the proof of Theorem 5.1 in the general case fix an index $1 \leq i \leq i_*$ and let $x_k = x_k^{(i)} \rightarrow x^{(i)}$, $0 < r_k = r_k^{(i)} \rightarrow 0$ be determined as in Theorem 4.2 so that $u_k(x_k) = \max_{|x-x_k| \leq Lr_k} u_k(x)$ for any $L > 0$ and sufficiently large k and such that

$$(82) \quad \eta_k(x) = \eta_k^{(i)}(x) := u_k(x_k)(u_k(x_k + r_k x) - u_k(x_k)) \rightarrow \log \left(\frac{1}{1 + |x|^2} \right)$$

as $k \rightarrow \infty$. For each k we may shift the origin so that henceforth we may assume that $x_k = 0$ for all k . Denote as $\Omega_k = \Omega_k^{(i)}$ the shifted domain Ω . We also extend u_k by 0 outside Ω_k to obtain $u_k \in H^1(\mathbb{R}^2)$, still satisfying (47).

Again we let $e_k = \lambda_k u_k^2 e^{u_k^2}$, $f_k = \lambda_k u_k(0) u_k e^{u_k^2}$, and for $0 < r < R$ we set

$$\Lambda_k(r) = \int_{B_r} e_k dx, \quad \sigma_k(r) = \int_{B_r} f_k dx,$$

satisfying (49).

Also introduce the spherical mean $\bar{u}_k(r) = \int_{\partial B_r} u_k do$ of u_k on ∂B_r , and so on, and set $\tilde{e}_k = \lambda_k \bar{u}_k^2 e^{\bar{u}_k^2}$.

The spherical mean \bar{w}_k of the function

$$w_k(x) = u_k(0)(u_k(x) - u_k(0)),$$

satisfies the equation

$$(83) \quad -\Delta \bar{w}_k = \bar{f}_k - \bar{d}_k,$$

where $\bar{f}_k = \lambda_k u_k(0) \overline{u_k e^{u_k^2}}$ and where

$$\bar{d}_k = u_k(0) \overline{u_k e^{u_k^2}} \rightarrow 0 \text{ in } L^1(B_{Lr_k})$$

for any $L > 0$ as $k \rightarrow \infty$ similar to (59).

Note that by Jensen's inequality we have

$$(84) \quad \tilde{e}_k \leq \bar{e}_k;$$

hence

$$\tilde{\Lambda}_k(r) := \int_{B_r} \tilde{e}_k dx \leq \Lambda_k(r), \quad \int_{B_r} \bar{f}_k dx = \sigma_k(r).$$

Observe that in analogy with (49) Theorem 4.2 implies

$$(85) \quad \lim_{L \rightarrow \infty} \lim_{k \rightarrow \infty} \tilde{\Lambda}_k(Lr_k) = \lim_{L \rightarrow \infty} \lim_{k \rightarrow \infty} \Lambda_k(Lr_k) = \lim_{L \rightarrow \infty} \lim_{k \rightarrow \infty} \sigma_k(Lr_k) = \Lambda_1.$$

To proceed, we need the following estimate similar to the gradient estimate of Druet [6], Proposition 2. For any $k \in \mathbb{N}$, $x \in \Omega$ we let

$$R_k(x) = \inf_{1 \leq j \leq i_*} |x - x_k^{(j)}|.$$

Proposition 5.8. *There exists a uniform constant C such that for all $y \in \Omega$ there holds*

$$\sup_{z \in B_{R_k(y)/2}(y)} |u_k(y) - u_k(z)| u_k(y) \leq C,$$

uniformly in $k \in \mathbb{N}$.

The proof of Proposition 5.8 is given in the next section.

Recalling that $x_k^{(i)} = 0$, we let

$$\rho_k = \rho_k^{(i)} = \frac{1}{2} \inf_{j \neq i} |x_k^{(j)}|,$$

and we set $\rho_k = \text{diam}(\Omega)$ if $\{j; j \neq i\} = \emptyset$, that is, if there is no other concentration point but $x_k^{(i)}$. We now use Proposition 5.8 to deal with concentrations around the point $x_k^{(i)}$ at scales which are small with respect to ρ_k .

Indeed, for $|x| \leq \rho_k$ we have $|x| = R_k(x)$; therefore, by Proposition 5.8 and Lemma 5.14 below for any $0 < r \leq \rho_k$ with a uniform constant C there holds

$$(86) \quad \sup_{r/2 \leq |x| \leq r} u_k^2(x) - \inf_{r/2 \leq |x| \leq r} u_k^2(x) \leq C.$$

Hence, in particular, there holds

$$(87) \quad \sup_{r/2 \leq |x| \leq r} e^{u_k^2(x)} \leq C e^{\bar{u}_k^2(r)},$$

and we conclude the estimate

$$(88) \quad \frac{1}{C_3} \sup_{r/2 \leq |x| \leq r} u_k^2(x) e^{u_k^2(x)} \leq (1 + \bar{u}_k^2(r)) e^{\bar{u}_k^2(r)} \leq C_3 \inf_{r/2 \leq |x| \leq r} (1 + u_k^2(x)) e^{u_k^2(x)}$$

with a uniform constant C_3 . In the following we proceed as in [14]; therefore we only sketch the necessary changes we have to perform in the present case.

Because of our choice of origin $x_k^{(i)} = 0$ there holds $u_k(x) \leq u_k(0)$ for all $|x| \leq Lr_k$, $k \geq k_0(L)$; hence at this scale there also holds the inequality $e_k \leq f_k$.

Similar to Lemma 5.2 with the help of (88) we obtain

Lemma 5.9. *For any $\varepsilon > 0$, if there is a minimal number $0 < T_k \leq \rho_k$ such that $\bar{u}_k(T_k) = \varepsilon u_k(0)$, then for any constant $b < 2$ and sufficiently large k there holds*

$$\bar{w}_k(r) \leq b \log \left(\frac{r_k}{r} \right) \quad \text{on } B_{T_k}$$

and we have

$$\lim_{k \rightarrow \infty} \tilde{\Lambda}_k(T_k) = \lim_{k \rightarrow \infty} \Lambda_k(T_k) = \lim_{k \rightarrow \infty} \sigma_k(T_k) = 4\pi.$$

Next we define for $0 \leq s < t \leq \rho_k$

$$N_k(s, t) = \int_{B_t \setminus B_s} e_k dx = \lambda_k \int_{B_t \setminus B_s} u_k^2 e^{u_k^2} dx,$$

and

$$\tilde{N}_k(s, t) = \int_{B_t \setminus B_s} \tilde{e}_k dx = 2\pi\lambda_k \int_s^t r \bar{u}_k^2 e^{\bar{u}_k^2} dr \leq N_k(s, t),$$

where we used Jensen's inequality for the last estimate. Moreover we let

$$P_k(t) = t \frac{\partial}{\partial t} N_k(s, t) = t \int_{\partial B_t} e_k do$$

and

$$\tilde{P}_k(t) = t \frac{\partial}{\partial t} \tilde{N}_k(s, t) = t \int_{\partial B_t} \tilde{e}_k do = 2\pi t^2 \lambda_k \bar{u}_k^2 e^{\bar{u}_k^2} \leq P_k(t).$$

The estimate (88) implies

$$(89) \quad N_k(s, t) \leq C_3 \tilde{N}_k(s, t) + o(1) \quad \text{and} \quad P_k(t) \leq C_3 \tilde{P}_k(t) + o(1),$$

with error $o(1) \rightarrow 0$ as $k \rightarrow \infty$, uniformly in $s \leq t \leq \rho_k$. Moreover, similar to [14], estimate (26), by (88) with uniform constants C_4, C_5 we have

$$(90) \quad P_k(t) \leq C_4 N_k(t/2, t) + o(1) \leq C_5 P_k(t/2) + o(1).$$

If for some $\varepsilon > 0$ there is no $T_k = T_k(\varepsilon) \leq \rho_k$ as in Lemma 5.9 we continue our argument as described in Case 1 after Proposition 5.11. Otherwise, we proceed by iteration as in the radially symmetric case. Choose a sequence $\varepsilon_k \downarrow 0$ such that with corresponding numbers $s_k = T_k(\varepsilon_k) \leq \rho_k$ we have $\bar{u}_k(s_k) \rightarrow \infty$ as $k \rightarrow \infty$. Then there holds

$$\lim_{k \rightarrow \infty} \Lambda_k(s_k) = \Lambda_1 = 4\pi$$

and

$$\lim_{L \rightarrow \infty} \lim_{k \rightarrow \infty} (\Lambda_k(s_k) - \Lambda_k(Lr_k)) = \lim_{k \rightarrow \infty} \frac{\bar{u}_k(s_k)}{\bar{u}_k(r_k)} = \lim_{k \rightarrow \infty} \frac{r_k}{s_k} = \lim_{k \rightarrow \infty} s_k = 0.$$

By a slight abuse of notation we let $r_k = r_k^{(1)}, s_k = s_k^{(1)}$. Suppose that for some $l \geq 0$ we already have determined numbers $r_k^{(1)} < s_k^{(1)} < \dots < s_k^{(l)} \leq \rho_k$ such that

$$(91) \quad \lim_{L \rightarrow \infty} \lim_{k \rightarrow \infty} \Lambda_k(s_k^{(l)}) = \Lambda_1 l = 4\pi l$$

and

$$(92) \quad \lim_{L \rightarrow \infty} \lim_{k \rightarrow \infty} (\Lambda_k(s_k^{(l)}) - \Lambda_k(Lr_k^{(l)})) = \lim_{k \rightarrow \infty} \frac{\bar{u}_k(s_k^{(l)})}{\bar{u}_k(r_k^{(l)})} = \lim_{k \rightarrow \infty} \frac{r_k^{(l)}}{s_k^{(l)}} = \lim_{k \rightarrow \infty} s_k^{(l)} = 0.$$

Similar to Lemma 5.3 we now have the following result.

Lemma 5.10. *i) Suppose that for some $s_k^{(l)} < t_k \leq \rho_k$ there holds*

$$\sup_{s_k^{(l)} < t < t_k} P_k(t) \rightarrow 0 \quad \text{as } k \rightarrow \infty.$$

Then we have

$$\lim_{k \rightarrow \infty} N_k(s_k^{(l)}, t_k) = 0.$$

ii) Conversely, if for some $s_k^{(l)} < t_k$ and a subsequence (u_k) there holds

$$\lim_{k \rightarrow \infty} N_k(s_k^{(l)}, t_k) = \nu_0 > 0, \quad \lim_{k \rightarrow \infty} \frac{t_k}{\rho_k} = 0,$$

then either $\nu_0 \geq \pi$, or we have

$$\liminf_{k \rightarrow \infty} P_k(t_k) \geq \nu_0$$

and

$$\lim_{L \rightarrow \infty} \liminf_{k \rightarrow \infty} N_k(s_k^{(l)}, Lt_k) \geq \pi, \quad \lim_{L \rightarrow \infty} \limsup_{k \rightarrow \infty} N_k(s_k^{(l)}, t_k/L) = 0.$$

Proof. i) Because of the estimate (89) it is enough to prove the Lemma with $N_k(s, t)$ and $P_k(t)$ replaced by $\tilde{N}_k(s, t)$ and $\tilde{P}_k(t)$. For $s = s_k^{(l)} < t$ we integrate by parts as before to obtain

$$(93) \quad 2\tilde{N}_k(s, t) \leq \tilde{P}_k(t) - 4\pi \int_s^t \lambda_k r^2 \bar{u}'_k(1 + \bar{u}_k^2) \bar{u}_k e^{\bar{u}_k^2} dr.$$

As in the proof of Lemma 5.3 equation (3) yields the identity

$$(94) \quad -2\pi t \bar{u}_k(t) \bar{u}'_k(t) = \int_{B_t} \lambda_k \bar{u}_k(t) \overline{u_k e^{u_k^2}} dx - \int_{B_t} \bar{u}_k(t) \overline{\dot{u}_k e^{u_k^2}} dx.$$

for any $t \leq \rho_k$. Arguing as in (59) we get that

$$\int_{B_t} \bar{u}_k(t) \overline{\dot{u}_k e^{u_k^2}} dx \rightarrow 0$$

as $k \rightarrow \infty$. In view of (94) and Jensen's inequality at any sequence of points $t = t_k$ where $\bar{u}'_k(t) \geq 0$ then there holds

$$(95) \quad 0 \leq \int_{B_t} \lambda_k \bar{u}_k(t) \bar{u}_k e^{\bar{u}_k^2} dx \leq \int_{B_t} \lambda_k \bar{u}_k(t) \overline{u_k e^{u_k^2}} dx = o(1).$$

Conversely, if $\bar{u}'_k(r) \leq 0 = \bar{u}'_k(t_0)$ for $t_{k0} = t_0 \leq r \leq t = t_k$, by (95) we can estimate

$$(96) \quad \begin{aligned} \int_{B_t} \lambda_k \bar{u}_k(t) \bar{u}_k e^{\bar{u}_k^2} dx &\leq \int_{B_t \setminus B_{t_0}} \lambda_k \bar{u}_k^2 e^{\bar{u}_k^2} dx + \int_{B_{t_0}} \lambda_k \bar{u}_k(t_0) \bar{u}_k e^{\bar{u}_k^2} dx \\ &= \tilde{N}_k(t_0, t) + o(1). \end{aligned}$$

Combining the above estimates, similar to (62) for $s = s_k^{(l)} \leq r \leq t = t_k$ we get

$$(97) \quad \begin{aligned} -2\pi r \bar{u}_k(r) \bar{u}'_k(r) &= \int_{B_r} \lambda_k \bar{u}_k(r) \bar{u}_k e^{\bar{u}_k^2} dx + o(1) \\ &\leq \tilde{N}_k(s, r) + \int_{B_s} \lambda_k \bar{u}_k(s) \bar{u}_k e^{\bar{u}_k^2} dx + o(1) \\ &\leq \tilde{N}_k(s, r) + \tilde{N}_k(Lr_k, s) + \frac{\bar{u}_k(s)}{\bar{u}_k(Lr_k)} \Lambda_k(Lr_k) + o(1) \\ &= \tilde{N}_k(s, r) + o(1), \end{aligned}$$

where $o(1) \rightarrow 0$ when first $k \rightarrow \infty$ and then $L \rightarrow \infty$. As in (62) the first inequality is clear when $\bar{u}'_k \leq 0$ in $[s, r]$, and otherwise follows from (95), (96). The second

inequality is proved similarly. Thus we conclude the estimate

$$(98) \quad \begin{aligned} 2\tilde{N}_k(s, t) &\leq \tilde{P}_k(t) + 2 \int_s^t \lambda_k r (1 + \bar{u}_k^2) e^{\bar{u}_k^2} \tilde{N}_k(s, r) dr + o(1) \\ &\leq \tilde{P}_k(t) + \pi^{-1} \tilde{N}_k(s, t)^2 + o(1). \end{aligned}$$

If we now assume that

$$\sup_{s < t < t_k} \tilde{P}_k(t) \leq C_3 \sup_{s < t < t_k} P_k(t) + o(1) \rightarrow 0 \text{ as } k \rightarrow \infty,$$

as in Lemma 5.3 we find the desired decay

$$\lim_{k \rightarrow \infty} \tilde{N}_k(s_k^{(l)}, t_k) = 0$$

when we let t increase from $t = s = s_k^{(l)}$ to t_k .

ii) In view of (98) the second assertion can be proved as in Lemma 5.3. \square

By the preceding result it now suffices to consider the following two cases. In **Case A** for any sequence $t_k = o(\rho_k)$ we have

$$\sup_{s_k^{(l)} < t < t_k} P_k(t) \rightarrow 0 \text{ as } k \rightarrow \infty,$$

and then in view of Lemma 5.10 also

$$(99) \quad \lim_{L \rightarrow \infty} \lim_{k \rightarrow \infty} N_k(s_k^{(l)}, \rho_k/L) = 0,$$

thus completing the concentration analysis at scales up to $o(\rho_k)$.

In **Case B** for some $s_k^{(l)} < t_k \leq \rho_k$ there holds

$$\limsup_{k \rightarrow \infty} N_k(s_k^{(l)}, t_k) > 0, \quad \lim_{k \rightarrow \infty} \frac{t_k}{\rho_k} = 0.$$

Then, as in the radial case, from Lemma 5.10 we infer that for a subsequence (u_k) and suitable numbers $r_k^{(l+1)} \in]s_k^{(l)}, t_k[$ we have

$$(100) \quad \lim_{L \rightarrow \infty} \lim_{k \rightarrow \infty} N_k(s_k^{(l)}, Lr_k^{(l+1)}) \geq \pi, \quad \liminf_{k \rightarrow \infty} P_k(r_k^{(l+1)}) > 0;$$

in particular, $\bar{u}_k(r_k^{(l+1)}) \rightarrow \infty$ as $k \rightarrow \infty$. Moreover, as in Lemma 5.10 the bound (100) implies that $r_k^{(l+1)}/s_k^{(l)} \rightarrow \infty$ as $k \rightarrow \infty$. Indeed, assume by contradiction that $r_k^{(l+1)} \leq Ls_k^{(l)}$ for some $L > 0$. Then from (88), (90), and recalling that $N_k(s_k^{(l)}/2, s_k^{(l)}) \rightarrow 0$ as $k \rightarrow \infty$ we obtain that $P_k(r_k^{(l+1)}) \rightarrow 0$ contrary to (100). Also note that

$$(101) \quad \lim_{L \rightarrow \infty} \limsup_{k \rightarrow \infty} N_k(s_k^{(l)}, r_k^{(l+1)}/L) = \lim_{k \rightarrow \infty} \frac{r_k^{(l+1)}}{\rho_k} = \lim_{k \rightarrow \infty} \frac{t_k}{\rho_k} = 0.$$

Moreover, we have the following analogue of Proposition 5.4.

Proposition 5.11. *There exist a subsequence (u_k) such that*

$$\eta_k^{(l+1)}(x) := \bar{u}_k(r_k^{(l+1)})(u_k(r_k^{(l+1)}x) - \bar{u}_k(r_k^{(l+1)})) \rightarrow \eta(x)$$

locally uniformly on $\mathbb{R}^2 \setminus \{0\}$ as $k \rightarrow \infty$, where η solves (33), (34).

Proposition 5.11 is a special case of Proposition 5.12 below, whose proof will be presented in Section 5.5.

From Proposition 5.11 the desired energy quantization result at the scale $r_k^{(l+1)}$ follows as in the radial case. If $\rho_k \geq \rho_0 > 0$ we can argue as in [14], p. 416, to obtain numbers $s_k^{(l+1)}$ satisfying (91), (92) for $l+1$ and such that $\bar{u}_k(s_k^{(l+1)}) \rightarrow \infty$ as $k \rightarrow \infty$. By iteration we then establish (91), (92) up to $l = l_0$ for some maximal index $l_0 \geq 0$ and thus complete the concentration analysis near the point $x^{(i)}$.

If $\rho_k \rightarrow 0$ as $k \rightarrow \infty$, we distinguish the following two cases. In **Case 1** for some $\varepsilon_0 > 0$ and all $t \in [r_k^{(l+1)}, \rho_k]$ there holds $\bar{u}_k(t) \geq \varepsilon_0 \bar{u}_k(r_k^{(l+1)})$. The decay estimate that we established in Lemma 5.9 then remains valid throughout this range and (91) holds true for any choice $s_k^{(l+1)} = o(\rho_k)$ for $l = l+1$. Again the concentration analysis at scales up to $o(\rho_k)$ is complete. In **Case 2**, for any $\varepsilon > 0$ there is a minimal $T_k = T_k(\varepsilon) \in [r_k^{(l+1)}, \rho_k]$ as in Lemma 5.9 such that $\bar{u}_k(T_k) = \varepsilon \bar{u}_k(r_k^{(l+1)})$. Then as before we can define numbers $s_k^{(l+1)} < \rho_k$ with $\bar{u}_k(s_k^{(l+1)}) \rightarrow \infty$ as $k \rightarrow \infty$ so that (91), (92) also hold true for $l+1$, and we proceed by iteration up to some maximal index $l_0 \geq 0$ where either Case 1 or Case A holds with final radius $r^{(l_0)}$.

For the concentration analysis at the scale ρ_k first assume that for some number $L \geq 1$ there is a sequence (x_k) such that $\rho_k/L \leq R_k(x_k) \leq |x_k| \leq L\rho_k$ and

$$(102) \quad \lambda_k |x_k|^2 u_k^2(x_k) e^{u_k^2(x_k)} \geq \nu_0 > 0.$$

By Proposition 5.8 we may assume that $|x_k| = \rho_k$. As in [14], Lemma 4.6, we then have $\bar{u}_k(\rho_k)/\bar{u}_k(r_k^{(l_0)}) \rightarrow 0$ as $k \rightarrow \infty$, ruling out Case 1; that is, at scales up to $o(\rho_k)$ we end with Case A. The desired quantization result at the scale ρ_k then is a consequence of the following result that we demonstrate in Section 5.5 below.

Proposition 5.12. *Assuming (102), there exists a finite set $S_0 \subset \mathbb{R}^2$ and a subsequence (u_k) such that*

$$\eta_k(x) := u_k(x_k)(u_k(\rho_k x) - u_k(x_k)) \rightarrow \eta(x)$$

locally uniformly on $\mathbb{R}^2 \setminus S_0$ as $k \rightarrow \infty$, where η solves (33), (34).

By Proposition 5.12 in case of (102) there holds

$$\lim_{L \rightarrow \infty} \limsup_{k \rightarrow \infty} \int_{\{x \in \Omega; \frac{\rho_k}{L} \leq R_k(x) \leq |x| \leq L\rho_k\}} e_k dx = \Lambda_1 = 4\pi.$$

Letting

$$X_{k,1} = X_{k,1}^{(i)} = \{x_k^{(j)}; \exists C > 0 : |x_k^{(j)}| \leq C\rho_k \text{ for all } k\}$$

and carrying out the above blow-up analysis up to scales of order $o(\rho_k)$ also on all balls of center $x_k^{(j)} \in X_{k,1}$, then from (92) we have

$$\lim_{L \rightarrow \infty} \lim_{k \rightarrow \infty} \Lambda_k(L\rho_k) = \Lambda_1(1 + I_1) = 4\pi(1 + I_1),$$

where I_1 is the total number of bubbles concentrating at the points $x_k^{(j)} \in X_{k,1}^{(i)}$ at scales $o(\rho_k)$.

On the other hand, if (102) fails to hold clearly we have

$$(103) \quad \lim_{L \rightarrow \infty} \limsup_{k \rightarrow \infty} \int_{\{x \in \Omega; \frac{\rho_k}{L} \leq R_k(x) \leq |x| \leq L\rho_k\}} e_k dx = 0,$$

and the energy estimate at the scale ρ_k again is complete.

In order to deal with secondary concentrations around $x_k^{(i)} = 0$ at scales exceeding ρ_k , with $X_{k,1}$ defined as above we let

$$\rho_{k,1} = \rho_{k,1}^{(i)} = \frac{1}{2} \inf_{\{j; x_k^{(j)} \notin X_{k,1}\}} |x_k^{(j)}|;$$

again we set $\rho_{k,1} = \text{diam}(\Omega)$, if $\{j; x_k^{(j)} \notin X_{k,1}\} = \emptyset$. From this definition it follows that $\rho_{k,1}/\rho_k \rightarrow \infty$ as $k \rightarrow \infty$. Then either we have

$$\lim_{L \rightarrow \infty} \limsup_{k \rightarrow \infty} N_k(L\rho_k, \frac{\rho_{k,1}}{L}) = 0,$$

and we iterate to the next scale; or there exist radii $t_k \leq \rho_{k,1}$ such that $t_k/\rho_k \rightarrow \infty$, $t_k/\rho_{k,1} \rightarrow 0$ as $k \rightarrow \infty$ and a subsequence (u_k) such that

$$(104) \quad P_k(t_k) \geq \nu_0 > 0 \quad \text{for all } k.$$

The argument then depends on whether (102) or (103) holds. In case of (102), as in [14], Lemma 4.6, the bound (104) and Proposition 5.12 imply that $\bar{u}_k(t_k)/\bar{u}_k(\rho_k) \rightarrow 0$ as $k \rightarrow \infty$. Then all the previous results remain true for $r \in [L\rho_k, \rho_{k,1}]$ for sufficiently large L , and we can continue as before to resolve concentrations in this range of scales.

In case of (103) we further need to distinguish whether Case A or Case 1 holds at the final stage of our analysis at scales $o(\rho_k)$. In fact, for the following estimates we also consider all points $x_k^{(j)} \in X_{k,1}^{(i)}$ in place of $x_k^{(i)}$. Recalling that in Case A we have (92) and (99), and arguing as above in Case 1, on account of (103) for a suitable sequence of numbers $s_{k,1}^{(0)}$ such that $s_{k,1}^{(0)}/\rho_k \rightarrow \infty$, $t_k/s_{k,1}^{(0)} \rightarrow \infty$ as $k \rightarrow \infty$ we find

$$\lim_{L \rightarrow \infty} \lim_{k \rightarrow \infty} \left(\Lambda(s_{k,1}^{(0)}) - \sum_{x_k^{(j)} \in X_{k,1}^{(i)}} \Lambda_k^{(j)}(Lr_k^{(l_0^{(j)})}) \right) = 0,$$

where $\Lambda_k^{(j)}(r)$ and $r_k^{(l_0^{(j)})}$ are computed as above with respect to the concentration point $x_k^{(j)}$. In particular, with such a choice of $s_{k,1}^{(0)}$ we find the intermediate quantization result

$$\lim_{k \rightarrow \infty} \Lambda_k(s_{k,1}^{(0)}) = \Lambda_1 I_1 = 4\pi I_1$$

analogous to (91), where I_1 is defined as above. Moreover, in Case 1 we can argue as in [14], Lemma 4.8, to conclude that $\bar{u}_k(t_k)/\bar{u}_k(r_k^{(l_0^{(j)})}) \rightarrow 0$ as $k \rightarrow \infty$; therefore, similar to (92) in Case A, we can achieve that

$$\lim_{k \rightarrow \infty} \frac{\bar{u}_k(s_{k,1}^{(0)})}{\bar{u}_k(r_k^{(l_0^{(j)})})} = \lim_{k \rightarrow \infty} \frac{r_k^{(l_0^{(j)})}}{s_{k,1}^{(0)}} = 0$$

for all $x_k^{(j)} \in X_{k,1}^{(i)}$ where Case 1 holds.

We then finish the argument by iteration. For $l \geq 2$ we inductively define the sets

$$X_{k,l} = X_{k,l}^{(i)} = \{x_k^{(j)}; \exists C > 0 : |x_k^{(j)}| \leq C\rho_{k,l-1} \text{ for all } k\}$$

and we let

$$\rho_{k,l} = \rho_{k,l}^{(i)} = \frac{1}{2} \inf_{\{j; x_k^{(j)} \notin X_{k,l}^{(i)}\}} |x_k^{(j)}|;$$

as before, we set $\rho_{k,l} = \text{diam}(\Omega)$, if $\{j; x_k^{(j)} \notin X_{k,l}^{(i)}\} = \emptyset$. Iteratively performing the above analysis at all scales $\rho_{k,l}$, thereby exhausting all concentration points $x_k^{(j)}$, upon passing to further subsequences, we finish the proof of Theorem 1.1.

5.4. Proof of Proposition 5.8. We argue as in [14], thereby closely following the proof of Druet [6], Proposition 2. Suppose by contradiction that

$$(105) \quad L_k := \sup_{y \in \Omega} \left(\sup_{z \in B_{R_k(y)/2}(y)} |u_k(y) - u_k(z)| u_k(y) \right) \rightarrow \infty \text{ as } k \rightarrow \infty.$$

Let $y_k \in \Omega$, $z_k \in B_{R_k(y_k)/2}(y_k)$ satisfy

$$(106) \quad |u_k(y_k) - u_k(z_k)| u_k(y_k) \geq L_k/2.$$

Lemma 5.13. *We have $u_k(y_k) \rightarrow \infty$ as $k \rightarrow \infty$.*

Proof. Suppose by contradiction that $u_k(y_k) \leq C < \infty$. From (106) we then find that $u_k(z_k) \rightarrow \infty$ as $k \rightarrow \infty$. Also letting $\hat{z}_k = (y_k + z_k)/2$, we now observe that

$$R_k(z_k), R_k(\hat{z}_k) \geq R_k(y_k)/2 > |y_k - z_k| = 2|y_k - \hat{z}_k| = 2|\hat{z}_k - z_k|;$$

hence

$$y_k \in B_{R_k(\hat{z}_k)/2}(\hat{z}_k), \hat{z}_k \in B_{R_k(z_k)/2}(z_k).$$

But then the estimate

$$\frac{L_k}{2u_k(y_k)} \leq |u_k(z_k) - u_k(y_k)| \leq |u_k(z_k) - u_k(\hat{z}_k)| + |u_k(\hat{z}_k) - u_k(y_k)|,$$

our assumption that $u_k(y_k) \leq C$, and our choice of y_k, z_k imply

$$\frac{1}{L_k} (|u_k(\hat{z}_k) - u_k(y_k)| u_k(\hat{z}_k) + |u_k(\hat{z}_k) - u_k(z_k)| u_k(z_k)) \rightarrow \infty$$

as $k \rightarrow \infty$, and a contradiction to (105) results. \square

A similar reasoning also yields the following result.

Lemma 5.14. *There exists an absolute constant C such that*

$$\sup_{z \in B_{R_k(y)/2}(y)} |u_k^2(y) - u_k^2(z)| \leq CL_k,$$

uniformly in $y \in \Omega$. In fact, we may take $C = 6$.

Proof. From the identity

$$\begin{aligned} u_k^2(y) - u_k^2(z) &= (u_k(y) - u_k(z))(u_k(y) + u_k(z)) \\ &= 2(u_k(y) - u_k(z))u_k(y) - (u_k(y) - u_k(z))^2 \end{aligned}$$

we conclude the bound

$$|u_k^2(y) - u_k^2(z)| \leq 2L_k + (u_k(y) - u_k(z))^2$$

for all $y \in \Omega$, $z \in B_{R_k(y)/2}(y)$, and we are done unless for some such points y and z there holds $(u_k(y) - u_k(z))^2 \geq 4L_k$. Suppose we are in this case. From (105) we then obtain the estimate $u_k(y) \leq \sqrt{L_k}/2$ and hence $u_k(z) \geq 2\sqrt{L_k}$.

Letting $\hat{z} = (y + z)/2$, as in the proof of Lemma 5.13 above we observe that $R_k(z), R_k(\hat{z}) \geq R_k(y)/2 \geq |y - z|$ and

$$y \in B_{R_k(\hat{z})/2}(\hat{z}), \hat{z} \in B_{R_k(z)/2}(z).$$

Since $u_k(z) \geq 2\sqrt{L_k}$, the bound (105) implies that $u_k(\hat{z}) \geq 3\sqrt{L_k}/2$. But then, upon estimating

$$\begin{aligned} 2L_k &\geq |u_k(y) - u_k(\hat{z})|u_k(\hat{z}) + |u_k(\hat{z}) - u_k(z)|u_k(z) \\ &\geq 3\sqrt{L_k}|u_k(y) - u_k(z)|/2 \geq 3L_k \end{aligned}$$

we arrive at the desired contradiction. \square

From Theorem 4.2 and Lemma 5.13 it follows that $s_k := R_k(y_k) \rightarrow 0$ as $k \rightarrow \infty$. Set

$$\Omega_k = \{y; y_k + s_k y \in \Omega\}$$

and scale

$$v_k(y) = u_k(y_k + s_k y), \dot{v}_k(y) = \dot{u}_k(y_k + s_k y), y \in \Omega_k.$$

Letting $x_k^{(i)}$ be as in the statement of Theorem 4.2, we set

$$y_k^{(i)} = \frac{x_k^{(i)} - y_k}{s_k}, 1 \leq i \leq i_*,$$

and let

$$S_k = \{y_k^{(i)}; 1 \leq i \leq i_*\}.$$

Note that in the scaled coordinates we have

$$\text{dist}(0, S_k) = \inf\{|y_k^{(i)}|; 1 \leq i \leq i_*\} = 1.$$

Also let

$$p_k = \frac{z_k - y_k}{s_k} \in B_{1/2}(0).$$

Then there holds

$$\begin{aligned} (107) \quad L_k/2 &\leq |v_k(p_k) - v_k(0)|v_k(0) \\ &\leq \sup_{y \in \Omega_k} \left(\sup_{z \in B_{\text{dist}(y, S_k)/2}(y)} |v_k(y) - v_k(z)|v_k(y) \right) = L_k; \end{aligned}$$

moreover, from Lemma 5.14 we have

$$(108) \quad \sup_{y \in \Omega_k} \left(\sup_{z \in B_{\text{dist}(y, S_k)/2}(y)} |v_k^2(y) - v_k^2(z)| \right) \leq CL_k.$$

Since $s_k = R_k(y_k) \rightarrow 0$ we may assume that as $k \rightarrow \infty$ the domains Ω_k exhaust the domain

$$\Omega_0 = \mathbb{R} \times] - \infty, R_0[,$$

where $0 < R_0 \leq \infty$. We also may assume that as $k \rightarrow \infty$ either $|y_k^{(i)}| \rightarrow \infty$ or $y_k^{(i)} \rightarrow y^{(i)}$, $1 \leq i \leq i_*$, and we let S_0 be the set of accumulation points of S_k , satisfying $\text{dist}(0, S_0) = 1$. For $R > 0$ denote as

$$K_R = K_{k,R} = \Omega_k \cap B_R(0) \setminus \bigcup_{y \in S_k} B_{1/R}(y).$$

Note that we have

$$R_k(y_k + s_k y) = s_k \operatorname{dist}(y, S_k) \geq s_k/R \quad \text{for all } y \in K_R.$$

Thus (40) in Theorem 4.2 implies the bound

$$(109) \quad \lambda_k s_k^2 v_k^2(y) e^{v_k^2(y)} \leq C = C(R) \quad \text{for all } y \in K_R.$$

Finally, letting

$$(110) \quad -v_k \Delta v_k = \lambda_k s_k^2 v_k^2 e^{v_k^2} - s_k^2 v_k v_k e^{v_k^2} =: I_k + II_k,$$

by (47) we can estimate

$$(111) \quad \|I_k\|_{L^1(\Omega_k)} = \lambda_k \int_{\Omega_k} s_k^2 v_k^2 e^{v_k^2} dy = \lambda_k \int_{\Omega} u_k^2 e^{u_k^2} dx \leq C;$$

moreover, by Hölder's inequality and Lemma 4.1 we have

$$(112) \quad \|II_k\|_{L^1(\Omega_k)}^2 \leq \left(\lambda_k \int_{\Omega} u_k^2 e^{u_k^2} dx \right) \cdot \left(\lambda_k^{-1} \int_{\Omega} u_k^2 e^{u_k^2} dx \right) \rightarrow 0$$

as $k \rightarrow \infty$. In view of (109) we also have the local L^2 -bounds

$$(113) \quad \begin{aligned} \|I_k\|_{L^2(K_R)}^2 &\leq C \sup_{K_R} (\lambda_k s_k^2 v_k^2 e^{v_k^2}) \cdot \left(\lambda_k \int_{\Omega_k} s_k^2 v_k^2 e^{v_k^2} dy \right) \\ &\leq C(R) \lambda_k \int_{\Omega} u_k^2 e^{u_k^2} dx \leq C(R), \end{aligned}$$

while Lemma 4.1 implies

$$(114) \quad \|II_k\|_{L^2(K_R)}^2 \leq C \sup_{K_R} (\lambda_k s_k^2 v_k^2 e^{v_k^2}) \cdot \left(\lambda_k^{-1} \int_{\Omega} u_k^2 e^{u_k^2} dx \right) \rightarrow 0$$

as $k \rightarrow \infty$, for any $R > 0$. Similarly, for any $R > 0$ we find

$$(115) \quad \|\Delta v_k\|_{L^2(K_R)} \rightarrow 0 \quad (k \rightarrow \infty).$$

Also observe that (47) yields the uniform bound

$$(116) \quad \|\nabla v_k\|_{L^2(\Omega_k)} \leq C.$$

Lemma 5.15. *We have $R_0 = \infty$; that is, $\Omega_0 = \mathbb{R}^2$.*

Proof. Suppose by contradiction that $R_0 < \infty$. Choosing $R = 2R_0$, from (4) and (108) we conclude the uniform bound

$$\sup_{y \in K_R} v_k^2(y) \leq CL_k$$

with $C = C(R)$. Letting $w_k = \frac{v_k}{\sqrt{L_k}}$, we then have $0 \leq w_k \leq C$, while (115) and (116) give

$$\|\nabla w_k\|_{L^2(\Omega_k)} + \|\Delta w_k\|_{L^2(K_R)} \rightarrow 0 \quad \text{as } k \rightarrow \infty.$$

Since $w_k = 0$ on $\partial\Omega_k \cap K_R$, it follows that $w_k \rightarrow 0$ locally uniformly, contradicting the fact that $|w_k(p_k) - w_k(0)|w_k(0) \geq 1/2$. \square

Lemma 5.16. *As $k \rightarrow \infty$ we have*

$$\frac{v_k}{v_k(0)} \rightarrow 1 \quad \text{locally uniformly in } \mathbb{R}^2 \setminus S_0.$$

Proof. Recall from Lemma 5.13 that

$$c_k := u_k(y_k) = v_k(0) \rightarrow \infty \text{ as } k \rightarrow \infty.$$

Letting $w_k = c_k^{-1}v_k$, from (115) and (116) for any $R > 0$ then we have

$$\|\nabla w_k\|_{L^2(\Omega_k)} + \|\Delta w_k\|_{L^2(K_R)} \rightarrow 0 \text{ as } k \rightarrow \infty,$$

and we conclude that w_k converges locally uniformly on $\mathbb{R}^2 \setminus S_0$ to a constant limit function w . Recalling that $\text{dist}(0, S_0) = 1$, we obtain that $w \equiv w(0) = 1$, as claimed. \square

Define

$$\tilde{v}_k(y) = \frac{1}{L_k}(v_k(y) - v_k(0))v_k(0).$$

We claim that \tilde{v}_k grows at most logarithmically. To see this, let $s_0 \geq 2 \sup_i |y^{(i)}|$ and fix $q = 3/2$. For any fixed $R > 0$, any $y \in K_R$ with $|y| \geq q^L s_0$ let $y_l = q^{l-L}y$, $0 \leq l \leq L$, so that $y_{l-1} \in B_{\text{dist}(y_l, S_k)/2}(y_l)$ for all $l \geq 1$ and sufficiently large k . Note that we have $|v_k(y_0) - v_k(0)|v_k(0) \leq CL_k$. By Lemma 5.16 with error $o(1) \rightarrow 0$ as $k \rightarrow \infty$ then we can estimate

$$\begin{aligned} |\tilde{v}_k(y)| &\leq \frac{1}{L_k} \sum_{l=1}^L |v_k(y_l) - v_k(y_{l-1})|v_k(0) + C \\ (117) \quad &\leq \frac{1+o(1)}{L_k} \sum_{l=1}^L |v_k(y_l) - v_k(y_{l-1})|v_k(y_l) + C \\ &\leq C + (1+o(1))L \leq C + (C+o(1))\log|y|. \end{aligned}$$

Moreover, from (113), (114) and Lemma 5.16 for any $R > 0$ with a constant $C = C(R)$ we obtain

$$(118) \quad \|\Delta \tilde{v}_k\|_{L^2(K_R)} \leq \sup_{K_R} \left(\frac{v_k(0)}{L_k v_k} \right) \|v_k \Delta v_k\|_{L^2(K_R)} \leq C \sup_{K_R} \left(\frac{v_k(0)}{L_k v_k} \right) \rightarrow 0$$

as $k \rightarrow \infty$. Thus we may assume that $\tilde{v}_k \rightarrow \tilde{v}$ locally uniformly away from S_0 , where \tilde{v} satisfies

$$(119) \quad \Delta \tilde{v} = 0, \tilde{v}(0) = 0, \sup_{B_{1/2}(0)} \tilde{v} \geq 1/2, |\tilde{v}(y)| \leq C + C \log(1 + |y|).$$

Fix any point $x_0 \in S_0$. For any $r > 0$ upon estimating $v_k(0)v_k e^{v_k^2} \leq \max\{v_k^2(0)e^{v_k^2(0)}, v_k^2 e^{v_k^2}\}$ we have

$$L_k \int_{B_r(x_0)} |\Delta \tilde{v}_k| dx = \int_{B_r(x_0)} v_k(0) |\Delta v_k| dx = I_k + II_k,$$

where

$$\begin{aligned} I_k &= \int_{B_r(x_0)} \lambda_k s_k^2 v_k(0) v_k e^{v_k^2} dx \leq C \lambda_k s_k^2 v_k^2(0) e^{v_k^2(0)} + \lambda_k \int_{B_r(x_0)} s_k^2 v_k^2 e^{v_k^2} dx \\ &\leq C \lambda_k R_k^2(y_k) u_k^2(y_k) e^{u_k^2(y_k)} + \lambda_k \int_{\Omega} u_k^2 e^{u_k^2} dx \leq C \end{aligned}$$

by Theorem 4.2 and (88). Similarly, by Hölder's inequality

$$|II_k|^2 = \left| \int_{B_r(x_0)} s_k^2 v_k(0) |\dot{v}_k| e^{v_k^2} dx \right|^2 \leq C \lambda_k^{-1} \int_{B_r(x_0)} s_k^2 \dot{v}_k^2 e^{v_k^2} dx \rightarrow 0.$$

as $k \rightarrow \infty$. It follows that $\Delta \tilde{v}_k \rightarrow 0$ in $L^1_{loc}(\mathbb{R}^2)$ as $k \rightarrow \infty$. The sequence (\tilde{v}_k) therefore is uniformly locally bounded in $W^{1,q}$ for any $q < 2$ and the limit $\tilde{v} \in W^1_{loc}(\mathbb{R}^2)$ extends as a weakly harmonic function to all of \mathbb{R}^2 . The mean value property together with the logarithmic growth condition (119) then implies that \tilde{v} is a constant; see for instance [3], Theorem 2.4. That is, $\tilde{v} \equiv \tilde{v}(0) = 0$. But by (119) we have $\sup_{B_{1/2}(0)} |\tilde{v}| \geq 1/2$, which is the desired contradiction and completes the proof of Proposition 5.8.

5.5. Proof of Proposition 5.12. We follow closely the proof of Proposition 4.7 in [14]. Fix an index $i \in \{1, \dots, i_\star\}$ and write $r_k = \rho_k$. Define

$$v_k(y) = u_k(x_k^{(i)} + r_k y)$$

where $y \in \Omega_k = \Omega_k^{(i)} = \{y; x_k^{(i)} + r_k y \in \Omega\}$. Also let

$$y_k^{(j)} = \frac{x_k^{(j)} - x_k^{(i)}}{r_k}$$

and

$$S_k = S_k^{(i)} = \{y_k^{(j)}; 1 \leq j \leq i_\star\}.$$

By choosing a subsequence we may assume that as $k \rightarrow \infty$ either $|y_k^{(j)}| \rightarrow \infty$ or $y_k^{(j)} \rightarrow y^{(j)}$, $1 \leq j \leq i_\star$, and we let $S_0 = S_0^{(i)}$ be the set of accumulation points of S_k . Note that $0 \in S_0$. Finally we let

$$y_k^{(0)} = \frac{x_k - x_k^{(i)}}{r_k}$$

be the scaled points x_k for which (102) holds and which satisfy $|y_k^{(0)}| = 1$. Choosing another subsequence we may assume that $y_k^{(0)} \rightarrow y^{(0)}$ as $k \rightarrow \infty$.

Recalling that $v_k(y_k^{(0)}) \rightarrow \infty$ by (102) and observing that $\mathbb{R}^2 \setminus S_0$ is connected, from Proposition 5.8 and a standard covering argument we obtain that

$$(120) \quad v_k - v_k(y_k^{(0)}) \rightarrow 0 \quad \text{as } k \rightarrow \infty$$

locally uniformly on $\mathbb{R}^2 \setminus S_0$. Moreover, as $k \rightarrow \infty$, the sets Ω_k exhaust all of \mathbb{R}^2 .

Next we note that η_k satisfies the equation

$$(121) \quad -\Delta \eta_k = \lambda_k r_k^2 v_k(y_k^{(0)}) v_k e^{v_k^2} - r_k^2 v_k(y_k^{(0)}) \dot{v}_k e^{v_k^2} = I_k + II_k$$

on Ω_k . For $L > 1$ set $K_L = B_L(0) \setminus (\cup_{y_0 \in S_0} B_{1/L}(y_0))$. Another covering argument together with (88) allows to bound $e^{v_k^2} \leq C e^{v_k^2(y_k^{(0)})} = C e^{u_k^2(x_k)}$ on K_L , where $C = C(L)$. By (40) and Lemma 4.1 for any $L > 0$ we then obtain

$$\begin{aligned} \int_{K_L} |II_k|^2 dx &\leq C \lambda_k r_k^2 v_k^2(y_k^{(0)}) e^{v_k^2(y_k^{(0)})} \cdot (\lambda_k^{-1} \int_{B_L(0)} r_k^2 \dot{v}_k^2 e^{v_k^2} dx) \\ &= C \lambda_k r_k^2 u_k^2(x_k) e^{u_k^2(x_k)} \cdot (\lambda_k^{-1} \int_{B_{Lr_k}(x_k^{(i)})} \dot{u}_k^2 e^{u_k^2} dx) \rightarrow 0 \end{aligned}$$

as $k \rightarrow \infty$. Next rewrite I_k as

$$I_k = \lambda_k r_k^2 v_k^2(y_k^{(0)}) e^{v_k^2(y_k^{(0)})} \dot{v}_k e^{\eta_k(\dot{v}_k+1)},$$

where $\hat{v}_k = \frac{v_k}{v_k(y_k^{(0)})}$. From (120) we get that $\hat{v}_k \rightarrow 1$ locally uniformly on $\mathbb{R}^2 \setminus S_0$ while from (102) we conclude that

$$\lambda_k r_k^2 v_k^2(y_k^{(0)}) e^{v_k^2(y_k^{(0)})} = \lambda_k r_k^2 u_k^2(x_k) e^{u_k^2(x_k)} \rightarrow \mu_0$$

for some $\mu_0 > 0$ as $k \rightarrow \infty$. Since by Proposition 5.8 η_k is locally uniformly bounded, from (121) and the above considerations via standard L^2 -theory we obtain that η_k is uniformly locally bounded in H^2 away from S_0 . Hence we conclude that η_k converges locally uniformly away from S_0 and weakly locally in H^2 to some limit $\eta_0 \in H_{loc}^2(\mathbb{R}^2 \setminus S_0)$ which is smooth away from S_0 and which satisfies the equation

$$(122) \quad -\Delta \eta_0 = \mu_0 e^{2\eta_0}$$

on $\mathbb{R}^2 \setminus S_0$. Recalling that $\hat{v}_k I_k = \lambda_k r_k^2 v_k^2 e^{v_k^2}$, from (47) we can estimate

$$\begin{aligned} \int_{\mathbb{R}^2} e^{2\eta_0} dx &\leq \lim_{L \rightarrow \infty} \liminf_{k \rightarrow \infty} \int_{K_L} \hat{v}_k^2 e^{\eta_k(\hat{v}_k+1)} dx = \lim_{L \rightarrow \infty} \liminf_{k \rightarrow \infty} \int_{K_L} \mu_0^{-1} \hat{v}_k I_k dx \\ &\leq \mu_0^{-1} \limsup_{k \rightarrow \infty} \int_{\Omega} \lambda_k u_k^2 e^{u_k^2} dx \leq C\Lambda \end{aligned}$$

as before, and $e^{2\eta_0} \in L^1(\mathbb{R}^2)$.

Similar to (77) we can moreover estimate for every $L \geq 1$

$$\int_{B_L(y_0)} |II_k| dx \rightarrow 0 \quad \text{as } k \rightarrow \infty,$$

and analogous to (80) we have

$$\int_{B_{1/L}(y_0)} I_k dx \rightarrow 0$$

for any $y_0 \in S_0$ if we let first $k \rightarrow \infty$ and then $L \rightarrow \infty$. Hence for such y_0 we conclude that

$$\limsup_{k \rightarrow \infty} \int_{B_{1/L}(0)} |\Delta \eta_k| dx \rightarrow 0 \quad \text{as } L \rightarrow \infty.$$

This shows that η_0 extends as a distribution solution of (122) on all of \mathbb{R}^2 . The claim then follows from the classification result of Chen-Li [4].

In the case of Proposition 5.11 we argue similarly by scaling with $r_k = r_k^{(l+1)}$. Note that in this case $S_0 = \{0\}$.

6. APPLICATIONS

In this final section we will use Theorem 1.1 to obtain solutions to (2) in the supercritical high energy regime.

Let Ω be a bounded domain in \mathbb{R}^2 . Recall the Moser-Trudinger inequality

$$(123) \quad \sup_{u \in H_0^1(\Omega); \|\nabla u\|_{L^2(\Omega)}^2 \leq 1} \int_{\Omega} e^{4\pi u^2} dx < \infty;$$

see [10], [15]. The exponent $\alpha = 4\pi$ is critical for this Orlicz space embedding in the sense that for any $\alpha > 4\pi$ there holds

$$(124) \quad \sup_{u \in H_0^1(\Omega); \|\nabla u\|_{L^2(\Omega)}^2 \leq 1} \int_{\Omega} e^{\alpha u^2} dx = \infty.$$

Indeed, suppose that $B_R(0) \subset \Omega$. Following Moser [10], for $0 < \rho < R$ consider the functions

$$m_{\rho,R}(x) = \frac{1}{\sqrt{2\pi}} \begin{cases} \sqrt{\log\left(\frac{R}{\rho}\right)}, & 0 \leq |x| \leq \rho, \\ \log\left(\frac{R}{r}\right) / \sqrt{\log\left(\frac{R}{\rho}\right)}, & \rho \leq |x| = r < R, \\ 0, & R \leq |x|. \end{cases}$$

Note that $\|\nabla m_{\rho,R}\|_{L^2(\Omega)}^2 = 1$, and for any $\alpha > 4\pi$ we have

$$(125) \quad \int_{\Omega} e^{\alpha m_{\rho,R}^2} dx \rightarrow \infty \text{ as } \rho \rightarrow 0.$$

After scaling, (123) gives

$$(126) \quad c_{\alpha} = c_{\alpha}(\Omega) := \sup_{u \in H_0^1(\Omega); \|\nabla u\|_{L^2(\Omega)}^2 \leq \alpha} E(u) < \infty$$

for any $\alpha \leq 4\pi$, while for any $\alpha > 4\pi$ from (124) we have

$$(127) \quad \sup_{u \in H_0^1(\Omega); \|\nabla u\|_{L^2(\Omega)}^2 \leq \alpha} E(u) = \infty.$$

If we normalize $\text{vol}(\Omega) = \pi$, the constant $c_{4\pi}(\Omega)$ is maximal when $\Omega = B_1(0) =: B$, as can be seen by symmetrization. Let $c_* = c_{4\pi}(B)$.

6.1. Solutions with “large” Moser-Trudinger energy on non-contractible domains. We obtain the following analogue of a result by Coron [5]; our result also is related to Theorem 1.1 in [12].

Theorem 6.1. *For any $c^* > c_*$ there are numbers $R_1 > R_2 > 0$ with the following property. Given any domain $\Omega \subset \mathbb{R}^2$ with $\text{vol}(\Omega) = \pi$ containing the annulus $B_{R_1} \setminus B_{R_2}(0)$ and such that $0 \notin \overline{\Omega}$, for any constant c_0 with $c_{4\pi}(\Omega) < c_0 < c^*$ problem (2) admits a positive solution u with $E(u) = c_0$.*

The proof of Theorem 6.1 relies on the following observation.

Lemma 6.2. *Let (u_k) be a sequence in $H_0^1(\Omega)$ such that*

$$E(u_k) \geq c > c_{4\pi}(\Omega), \quad \int_{\Omega} |\nabla u_k|^2 dx \rightarrow 4\pi \text{ as } k \rightarrow \infty.$$

Then there exists a point $x_0 \in \overline{\Omega}$ such that $|\nabla u_k|^2 dx \xrightarrow{w^} 4\pi\delta_{x_0}$ weakly in the sense of measures as $k \rightarrow \infty$ suitably.*

Proof. We may assume that $u_k \xrightarrow{w} u$ weakly in $H_0^1(\Omega)$ and pointwise almost everywhere as $k \rightarrow \infty$. Negating our claim, there exist $\alpha_1, r_1 > 0$ with $\alpha_1 < 4\pi$ such that

$$\sup_{k \in \mathbb{N}, x_1 \in \Omega} \int_{B_{r_1}(x_1) \cap \Omega} |\nabla u_k|^2 dx \leq \alpha_1.$$

But then by a reasoning as in the proof of Lemma 3.3 in [2] we conclude that the functions $e^{u_k^2}$ are uniformly bounded in L^q for some $q > 1$, and by Vitali's convergence theorem we have

$$E(u) = \lim_{k \rightarrow \infty} E(u_k) \geq c > c_{4\pi}(\Omega).$$

Since $\int_{\Omega} |\nabla u|^2 dx \leq 4\pi$, the latter contradicts (126), which proves our claim. \square

The proof of Theorem 6.1 now is achieved via a saddle-point construction similar to Section 3.4 in [12]. We may assume that $0 < R_1 < 1/2$. Given such R_1 , fix $R = R_1/4$. For each $R_2 < R_1/8 = R/2$, moreover, we let $\tau = \tau_{R_2} \in C_0^\infty(B_R(0))$ be a cut-off function $0 \leq \tau \leq 1$ satisfying $\tau \equiv 1$ on $B_{R_2}(0)$ and such that $\tau \rightarrow 0$ in $H^1(\mathbb{R}^2)$ as $R_2 \rightarrow 0$.

For $x_0 \in \mathbb{R}^2$ let $m_{\rho,R,x_0}(x) = m_{\rho,R}(x - x_0)$. With a suitable number $0 < \rho < R$ to be determined, for any x_0 with $|x_0| = 3R$, any $0 \leq s < 1$ then we define

$$v_{s,x_0}(x) = m_{s\rho,R,(1-s)x_0}(x)(1 - \tau(x)) \in H_0^1(B_{R_1} \setminus B_{R_2}(0)).$$

Provided that Ω contains the annulus $B_{R_1} \setminus B_{R_2}(0)$, these functions then also belong to $H_0^1(\Omega)$.

Given $c^* > c_*$, we fix the numbers $0 < \rho < R$, $0 < R_2 < R/2$ so that

$$(128) \quad \inf_{0 < s \leq 1, |x_0|=3R} \left(\frac{1}{2} \int_{\Omega} (e^{8\pi v_{s,x_0}^2} - 1) dx \right) > c^*$$

for all such domains Ω . This is possible by (125). Fixing such a domain Ω , finally, for any given $c_{4\pi}(\Omega) < c_0 < c^*$ we let

$$w_{s,x_0} = \sqrt{\alpha_{s,x_0}} v_{s,x_0},$$

where for each s, x_0 the number α_{s,x_0} is uniquely determined such that

$$E(w_{s,x_0}) = \frac{1}{2} \int_{\Omega} (e^{\alpha_{s,x_0} v_{s,x_0}^2} - 1) dx = c_0.$$

Observe that (126) and (128) imply the bounds $4\pi < \alpha_{s,x_0} < 8\pi$ for each s, x_0 , and

$$(129) \quad \alpha_{s,x_0} \rightarrow 4\pi \text{ as } s \rightarrow 0$$

uniformly in $|x_0| = 3R$ by (125).

Let $u_{s,x_0}(t)$ be the solution to the initial value problem (3) - (5) with initial data $u_{s,x_0}(0) = w_{s,x_0} \geq 0$.

Lemma 6.3. *With a uniform constant $\alpha_0 > 4\pi$ there holds*

$$(130) \quad \sup_{0 < s \leq 1, |x_0|=3R} \int_{\Omega} |\nabla u_{s,x_0}(t)|^2 dx \geq \alpha_0$$

for all $0 \leq t < \infty$.

Proof. Otherwise by (7) we have $\|\nabla u_{s,x_0}(t)\|_{L^2}^2 \rightarrow 4\pi$ as $t \rightarrow \infty$, uniformly in s and x_0 , and from Lemma 6.2 we conclude that

$$\sup_{0 < s \leq 1, |x_0|=3R} \text{dist}(m(u_{s,x_0}(t)), \Omega) \rightarrow 0$$

as $t \rightarrow \infty$, where

$$m(u) = \frac{\int_{\Omega} x |\nabla u|^2 dx}{\int_{\Omega} |\nabla u|^2 dx}$$

is the center of mass. Moreover, by (129), (8), and Lemma 6.2 we have

$$\sup_{0 < s \leq s_0, |x_0|=3R} \text{dist}(m(u_{s,x_0}(t)), \Omega) \rightarrow 0$$

as $s_0 \rightarrow 0$, uniformly in $t \geq 0$. Recall that $0 \notin \overline{\Omega}$. Thus, for some sufficiently small number $0 < s_0 < 1$ and sufficiently large $T > 0$ with a uniform constant $\delta > 0$ we have

$$\inf_{|x_0|=3R} |m(u_{s,x_0}(t))| \geq \delta > 0,$$

provided that either $0 < s \leq s_0$ or $t \geq T$. Identifying $\partial B_{3R}(0)$ with S^1 and letting $\pi_{S^1}(p) = p/|p|$ for $p \in \mathbb{R}^2 \setminus \{0\}$, then for sufficiently small $0 < s_0 < 1$ and sufficiently large $T > 0$ we can define a homotopy $H = H(\cdot, r): S^1 \times]0, T+1] \rightarrow S^1$ by letting

$$H(x_0, r) = \begin{cases} \pi_{S^1}(m(u_{r,x_0}(0))), & 0 < r \leq s_0, \\ \pi_{S^1}(m(u_{s_0,x_0}(r-s_0))), & s_0 \leq r \leq T+s_0, \\ \pi_{S^1}(m(u_{r-T,x_0}(T))), & T+s_0 \leq r \leq T+1. \end{cases}$$

Then clearly $H(\cdot, T+1) \equiv \text{const}$, whereas $H(x_0, r) \rightarrow x_0/|x_0|$ as $r \rightarrow 0$, uniformly in x_0 , which is impossible. The contradiction proves the claim. \square

Proof of Theorem 6.1 For any $t > 0$ by Lemma 6.3 there are $0 < s(t) \leq 1$, $x_0(t)$ with $|x_0(t)| = 3R$ such that

$$(131) \quad \int_{\Omega} |\nabla u_{s(t),x_0(t)}(t)|^2 dx \geq \alpha_0 > 4\pi.$$

Let (s_1, x_0) be a point of accumulation of $(s(t), x_0(t))$ as $t \rightarrow \infty$. Note that by (7) for any fixed time t_0 we have

$$(132) \quad \begin{aligned} 8\pi > \alpha_{s_1, x_0} &= \int_{\Omega} |\nabla u_{s_1, x_0}(0)|^2 dx \geq \int_{\Omega} |\nabla u_{s_1, x_0}(t_0)|^2 dx \\ &\geq \liminf_{t \rightarrow \infty} \int_{\Omega} |\nabla u_{s(t), x_0(t)}(t_0)|^2 dx \geq \liminf_{t \rightarrow \infty} \int_{\Omega} |\nabla u_{s(t), x_0(t)}(t)|^2 dx \geq \alpha_0 > 4\pi. \end{aligned}$$

Fix $u_0 = u_{s_1, x_0}(0) \geq 0$ and let $u(t)$ be the solution to the initial value problem (3) - (5) with initial data $u(0) = u_0$ with associated parameter $\lambda(t)$. We claim that $u(t)$ is uniformly bounded and hence converges to a solution $u_{\infty} > 0$ of (2) with

$$\int_{\Omega} |\nabla u_{\infty}|^2 dx > 4\pi \quad \text{and} \quad E(u_{\infty}) = c_0.$$

This will finish the proof of the Theorem.

Indeed, suppose by contradiction that $u(t)$ blows up as $t \rightarrow \infty$. For a sequence of numbers $t_k \rightarrow \infty$ as constructed in Lemma 4.1 then as $k \rightarrow \infty$ we have $\lambda_k := \lambda(t_k) \rightarrow \lambda_{\infty} \geq 0$; moreover, we may assume that $u_k := u(t_k) \xrightarrow{w} u_{\infty}$ in $H_0^1(\Omega)$ and pointwise almost everywhere, where u_{∞} solves (2). Finally, Theorem 1.1 and (132) also give the bound

$$(133) \quad \int_{\Omega} |\nabla u_{\infty}|^2 dx < 4\pi.$$

It then follows that $\lambda_\infty = 0$. Indeed, if we assume $\lambda_\infty > 0$, from (8) and the dominated convergence theorem we infer

$$E(u_\infty) = \lim_{k \rightarrow \infty} E(u_k) = c_0 > c_{4\pi}(\Omega),$$

which is impossible in view of (133) and (126). But with $\lambda_\infty = 0$ in view of (2) also u_∞ must vanish identically, and from Theorem 1.1 it follows that

$$(134) \quad \lim_{k \rightarrow \infty} \int_{\Omega} |\nabla u_k|^2 dx = 4\pi l,$$

for some $l \in \mathbb{N}$, contradicting (132). The proof is complete. \square

6.2. Saddle points of the Moser-Trudinger energy. By [7], Corollary 7, on any bounded domain $\Omega \subset \mathbb{R}^2$ the Moser-Trudinger energy E attains its maximum $\beta_{4\pi}^* := c_{4\pi}(\Omega)$ in the set $M_{4\pi}$ defined in (1). Moreover, we have

Lemma 6.4. *The set $K_{4\pi}$ of maximizers of E in $M_{4\pi}$ is compact.*

Proof. Any $u \in K_{4\pi}$ solves (2). Given a sequence $(u_k) \subset K_{4\pi}$, we may assume that $u_k \rightharpoonup u_\infty$ weakly in $H_0^1(\Omega)$ as $k \rightarrow \infty$ while by (9) the associated numbers $\lambda_k \rightarrow \lambda_\infty \geq 0$. If $\lambda_\infty > 0$, from (8) and the dominated convergence theorem as above we conclude that $E(u_k) \rightarrow E(u_\infty)$, so that $E(u_\infty) = \beta_{4\pi}^*$ and $u_\infty \neq 0$. But by a result of P.-L. Lions [8], Theorem I.6, this implies that the functions $e^{u_k^2}$ are uniformly bounded in L^q for some $q > 1$, and $u_k \rightarrow u_\infty$ strongly in $H_0^1(\Omega)$, as claimed. On the other hand, if $\lambda_\infty = 0$, from (2) we conclude that also u_∞ must vanish and $E(u_\infty) = 0$. Theorem I.6 in [8] then implies weak convergence

$$(135) \quad |\nabla u_k|^2 dx \xrightarrow{w^*} 4\pi \delta_{x_0}$$

for some $x_0 \in \bar{\Omega}$ in the sense of measures, and by Flucher [7], Lemma 4 and Theorem 5, we have $E(u_k) < \beta_{4\pi}^*$ for large k , contradicting our choice of (u_k) . \square

In view of Lemma 6.4 now Lemma 5.3 from [11] remains valid for a general domain and there exist numbers $\alpha^* > 4\pi$, $\varepsilon > 0$ such that for any $4\pi < \alpha < \alpha^*$ there holds

$$\beta_\alpha^* := \sup_{N_{\alpha,\varepsilon}} E > \sup_{N_{\alpha,2\varepsilon} \setminus N_{\alpha,\varepsilon}} E$$

where

$$N_{\alpha,\varepsilon} = \{u \in M_\alpha; \exists v \in K_{4\pi} : \|\nabla(u-v)\|_{L^2} < \varepsilon\}.$$

Moreover, for any such α there exists $\bar{u} \in N_{\alpha,\varepsilon}$ where $\beta_\alpha^* = E(\bar{u})$ is attained, and \bar{u} solves (2) for some $\bar{\lambda} \geq 0$. By (127) the set

$$\Gamma_\alpha = \{\gamma \in C^0([0, 1]; M_\alpha); \gamma(0) = \bar{u}, E(\gamma(1)) > \beta_\alpha^*\},$$

then is non-void for any $4\pi < \alpha < \alpha^*$. Since any $\gamma \in \Gamma_\alpha$ necessarily passes through the set $N_{\alpha,2\varepsilon} \setminus N_{\alpha,\varepsilon}$ we have

$$(136) \quad \beta_\alpha := \sup_{\gamma \in \Gamma_\alpha} \inf_{0 < s < 1} E(\gamma(s)) < \beta_\alpha^*.$$

Finally, observing that

$$c_\alpha = \sup_{u \in M_\alpha} E(u) \rightarrow 0 \text{ as } \alpha \rightarrow 0,$$

we can choose $4\pi < \alpha_1 < \alpha^*$ such that

$$(137) \quad c_{\alpha-4\pi} < \beta_\alpha \text{ for all } \alpha \in]4\pi, \alpha_1[.$$

Clearly, we may assume that $\alpha_1 \leq 8\pi$.

Theorem 6.5. *For any $4\pi < \alpha < \alpha_1$ there exists a pair of solutions $\underline{u}, \bar{u} \in M_\alpha$ of (2) with $0 < E(\underline{u}) < E(\bar{u}) = \beta_\alpha^*$.*

This result completes Theorem 1.8 from [11] where the existence of a pair of solutions of (2) only was shown for almost every $4\pi < \alpha < \alpha_1$.

Proof of Theorem 6.5 Let $4\pi < \alpha < \alpha_1$. It remains to find \underline{u} . Fix some $\gamma \in \Gamma_\alpha$ with

$$\inf_{0 < s < 1} E(\gamma(s)) > c_{\alpha-4\pi}.$$

Fix a number β with $\beta_\alpha < \beta < \beta_\alpha^*$. As long as $E(u(s, t)) \leq \beta$ let $u(s, t) \geq 0$ be the solution to the initial value problem (3), (4), (12) with initial data $u(s, 0) = \gamma(s) \geq 0$, and let $u(s, t) = u(s, t(s))$ for all $t \geq t(s)$ if there is some first $t(s) \geq 0$ where $E(u(s, t(s))) = \beta$. Note that by the implicit function theorem the family $u(s, t)$ thus defined depends continuously both on s and t unless $u_t(s, t(s)) = 0$ for some s with $E(u(s, t(s))) = \beta$, that is, unless there is a solution $0 < u \in M_\alpha$ of (2) with $E(u) = \beta$, in which case the proof is complete.

For $t > 0$ let $0 \leq s(t) < 1$ be such that

$$E(u(s(t), t)) = \inf_{0 < s < 1} E(u(s, t)) \leq \beta_\alpha$$

and let s_1 be a point of accumulation of $(s(t))_{t>0}$ as $t \rightarrow \infty$. Note that similar to (132) by (13) for any fixed time t_0 we have

$$E(u(s_1, t_0)) \leq \limsup_{t \rightarrow \infty} E(u(s(t), t_0)) \leq \limsup_{t \rightarrow \infty} E(u(s(t), t)) \leq \beta_\alpha.$$

Fix $u_0 = \gamma(s_1) \geq 0$ and let $u(t)$ with associated parameter $\lambda(t)$ be the solution to the initial value problem (3), (4), (12) with initial data $u(0) = u_0$, satisfying

$$(138) \quad c_{\alpha-4\pi} < E(\gamma(s_1)) = E(u(0)) \leq E(u(t)) \leq \beta_\alpha < \beta < \beta_\alpha^* \text{ for all } t.$$

We claim that $u(t)$ is uniformly bounded and thus converges to a solution $0 < u_\infty \in M_\alpha$ of (2) with $0 < E(u_\infty) < \beta_\alpha^*$. For this we argue as in the proof of Theorem 6.1.

Indeed, suppose by contradiction that $u(t)$ blows up as $t \rightarrow \infty$. For a sequence of numbers $t_k \rightarrow \infty$ as constructed in Lemma 4.1 then as $k \rightarrow \infty$ we have $\lambda_k := \lambda(t_k) \rightarrow \lambda_\infty \geq 0$; moreover, we may assume that $u_k := u(t_k) \xrightarrow{w} u_\infty$ in $H_0^1(\Omega)$ and pointwise almost everywhere, where $u_\infty \geq 0$ solves (2) with $\|\nabla u_\infty\|_{L^2}^2 \leq \alpha - 4\pi$ in view of Theorem 1.1. But then $\lambda_\infty = 0$. Indeed, if $\lambda_\infty > 0$, from (22) and the dominated convergence theorem we infer $E(u_k) \rightarrow E(u_\infty) \leq c_{\alpha-4\pi}$, contradicting (138). But with $\lambda_\infty = 0$ in view of (2) also u_∞ must vanish identically, and Theorem 1.1 yields the contradiction $\alpha = 4\pi$. The proof is complete. \square

REFERENCES

- [1] Adimurthi, O. Druet: *Blow-up analysis in dimension 2 and a sharp form of Trudinger-Moser inequality*, Comm. Partial Differential Equations 29 (2004), no. 1-2, 295–322.
- [2] Adimurthi, M. Struwe: *Global compactness properties of semilinear elliptic equations with critical exponential growth*, J. Funct. Analysis 175 (2000), 125–167.

- [3] Adimurthi, F. Robert, M. Struwe: *Concentration phenomena for Liouville's equation in dimension four*, J. Eur. Math. Soc. (JEMS) 8 (2006), no. 2, 171–180.
- [4] Wen Xiong Chen, Congming Li: *Classification of solutions of some nonlinear elliptic equations*, Duke Math. J. 63 (1991), no. 3, 615–622.
- [5] J.-M. Coron: *Topologie et cas limite des injections de Sobolev*, C.R. Acad. Sc. Paris **299**, Ser. I (1984), 209–212.
- [6] O. Druet: *Multibumps analysis in dimension 2: quantification of blow-up levels*, Duke Math. J. 132 (2006), no. 2, 217–269.
- [7] M. Flucher: *Extremal functions for the Trudinger-Moser inequality in 2 dimensions*. Comment. Math. Helv. 67 (1992), no. 3, 471–497.
- [8] P.-L. Lions: *The concentration-compactness principle in the calculus of variations. The limit case. I*. Rev. Mat. Iberoamericana 1 (1985), no. 1, 145–201.
- [9] J. P. Monahan: *Numerical solution of a non-linear boundary value problem*, Thesis, Princeton Univ., 1971.
- [10] J. Moser: *A sharp form of an inequality by N. Trudinger*, Indiana Univ. Math. J. **20** (1971), 1077–1092.
- [11] M. Struwe: *Critical points of embeddings of $H_0^{1,n}$ into Orlicz spaces*, Ann. Inst. H. Poincaré, Anal. Non Linéaire, 5 (1984), 425–464.
- [12] M. Struwe: *Positive solutions of critical semilinear elliptic equations on non-contractible planar domains*. J. Eur. Math. Soc. (JEMS) 2 (2000), no. 4, 329–388.
- [13] M. Struwe: *Curvature flows on surfaces*, Annali Sc. Norm. Sup. Pisa, Ser. V, 1 (2002), 247–274.
- [14] M. Struwe: *Quantization for a fourth order equation with critical exponential growth*, Math. Z. 256 (2007), 397–424.
- [15] N.S. Trudinger: *On embeddings into Orlicz spaces and some applications*, J. Math. Mech. **17** (1967), 473–484.

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