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1 Tomography of core-mantle boundary and lowermost mantle
2 coupled by geodynamics: joint models of shear and
3 compressional velocity

4 Gaia Soldati¹, Lapo Boschi^{2,3}, Steve Della Mora^{4,5}, and Alessandro M. Forte⁶

5 ¹Istituto Nazionale di Geofisica e Vulcanologia, Roma, Italy.

6 ²Sorbonne Universités, UPMC Univ Paris 06, UMR 7193, Institut des Sciences de la
7 Terre Paris (iSTeP), F-75005 Paris, France.

8 ³CNRS, UMR 7193, Institut des Sciences de la Terre Paris (iSTeP), F-75005 Paris,
9 France.

10 ⁴Institute of Geophysics, Department of Earth Sciences, ETH Zürich, Switzerland.

11 ⁵Zürich Insurance, Whiteley, U.K.

12 ⁶GEOTOP, Université du Québec à Montréal, Canada

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14 Abstract

15 We conduct joint tomographic inversions of P and S travel time observations to obtain
16 models of δv_P and δv_S in the entire mantle. We adopt a recently published method which
17 takes into account the geodynamic coupling between mantle heterogeneity and core-mantle
18 boundary (CMB) topography by viscous flow, where sensitivity of the seismic travel times to
19 the CMB is accounted for implicitly in the inversion. The approach involved scaling P -wave
20 velocity (more sensitive to the CMB) to density anomalies, in the assumption that mantle
21 heterogeneity has a purely thermal origin, so that velocity and density heterogeneity are
22 proportional to one another. On the other hand, it has sometimes been suggested that S -
23 wave velocity might be more directly sensitive to temperature, while P heterogeneity is more
24 strongly influenced by chemical composition. In the present study, we use only S -, and not
25 P -velocity, to estimate density heterogeneity through linear scaling, and hence the sensitivity
26 of core-reflected P phases to mantle structure. The seismic maps of the Earth's mantle and
27 CMB topography that we derive can explain the inverted seismic data while being physically
28 consistent with each other. Their similarity to the ones obtained by scaling P -velocity to
29 density suggests that compositional anomaly has a limited impact on viscous flow in the deep
30 mantle.

31 1 Introduction

32 Our understanding of the dynamics of Earth's mantle is largely based on a precise imaging
33 of its velocity structure. Combining tomographic inversion with geodynamic modelling, and
34 using seismically-inferred density variations as a proxy for CMB topography, *Soldati et al.*
35 (2012) obtained mantle v_P models which are physically sound (their geodynamic coupling to
36 CMB topography is accounted for) and which fit the seismic data (ISC data sets of P -wave
37 arrivals) at least as well as models obtained from seismic data alone. Their method requires
38 that a tomography model of mantle seismic velocity be interpreted in terms of equivalent
39 density anomalies via a constant or radially varying scaling factor. However, density and
40 velocity heterogeneity are proportional to one another only if mantle heterogeneity is of
41 purely thermal origin, and no compositional heterogeneity is present. This is certainly not
42 strictly true, and it is still questioned to what extent it is a valid approximation of the real
43 Earth (e.g. *Karato, 2003; Deschamps and Trampert, 2003; Trampert et al., 2004; Della Mora*
44 *et al., 2011*), at least for the uppermost and lowermost regions of the mantle.

45 Shear and compressional velocity are in principle sensitive to both composition and tem-
 46 perature, but we do not know a priori the relative importance of the two effects. Since
 47 *Soldati et al.* (2012) assume a linear relationship between temperature T and v_P anomalies,
 48 we investigate here the other end-member, scaling v_S to T . This is motivated by the stronger
 49 sensitivity of S -wave anomalies to temperature (and density) variations (*Goes et al.*, 2000).

50 We present here a new method to conduct joint inversions of P - and S -waves, which are no
 51 longer coupled via any assumption on their scaling; they are now indirectly coupled through
 52 the CMB, which is computed by integration of the δv_S heterogeneity structure weighted by
 53 the geodynamic sensitivity kernels (*Forte and Peltier*, 1991). This is different from previous
 54 joint P - S tomographic models (e.g., *Su and Dziewonski*, 1997; *Vasco and Johnson*, 1998;
 55 *Kennett et al.*, 1998; *Saltzer et al.*, 2001; *Kennett and Gorbato*, 2004; *Houser et al.*, 2008), in
 56 that the inversion is now also constrained by the expected physical coupling between mantle
 57 and CMB, and no a-priori value for δv_S -to- δv_P scaling is prescribed. Furthermore, only v_S ,
 58 and not v_P , is attached to density structure, while v_P heterogeneities are completely free
 59 parameters. The assumption that δv_S (rather than δv_P) be proportional to density anomaly
 60 is motivated by the observation of, e.g., *Boschi et al.* (2007, 2008) that the distribution of
 61 deep plume roots correlates much better with lowermost-mantle δv_S than δv_P . We take this
 62 as an indication that δv_S is more sensitive than δv_P to thermal variation, while δv_P is more
 63 strongly affected by chemical heterogeneity.

64 The mantle and CMB models we find following this new hybrid approach after scaling
 65 v_S velocity to density are almost coincident with the ones obtained by *Soldati et al.* (2012)
 66 scaling v_P to density; this is a strong indication that the potential presence of compositional
 67 heterogeneity in the lowermost mantle, while it may have important local effects, does not
 68 heavily affect the viscous convective flow (*Simmons et al.*, 2009).

69 2 Method

70 We assume a linear relationship between travel time data and seismic velocities (*Boschi and*
 71 *Dziewonski*, 2000), and use the LSQR method (*Paige and Saunders*, 1982) to iteratively
 72 invert for δv_P and δv_S the following system of equations

$$\delta t_S = - \int_{\text{path}} \frac{\delta v_S(r(s), \theta(s), \phi(s))}{v_S^2(r)} ds, \quad (1)$$

$$\delta t_P = - \int_{\text{path}} \frac{\delta v_P(r(s), \theta(s), \phi(s))}{v_P^2(r)} ds, \quad (2)$$

$$\delta t_{PcP} = - \int_{\text{path}} \frac{\delta v_P(r(s), \theta(s), \phi(s))}{v_P^2(r)} ds + K_{PcP} \frac{\delta c(\theta_b, \phi_b)}{c}, \quad (3)$$

where $r = r(s), \theta = \theta(s), \phi = \phi(s)$ is the ray-path equation, (θ_b, ϕ_b) are the coordinates at which the PcP raypath is reflected off the CMB, and K_{PcP} the sensitivity of δt to CMB undulations (defined e.g. by *Dziewonski and Gilbert (1976)*).

The CMB topography δc is the other unknown function to be determined; *Forte and Peltier (1991)* show that its spherical harmonic coefficients δc_{lm} coincide with the harmonic coefficients $\delta \rho_{lm}$ of density perturbation modulated by the viscosity-dependent CMB sensitivity kernels $B_l(r)$

$$\delta c_{lm} = \frac{1}{\Delta \rho_{cmb}} \int_c^a B_l(r) \delta \rho_{lm}(r) dr, \quad (4)$$

with c and a denoting the reference, mean radii of the CMB and Earth's surface, respectively, and $\Delta \rho_{cmb}$ the density jump across the CMB according to PREM (*Dziewonski and Anderson, 1981*). The sensitivity kernels B_l are computed adopting the radial viscosity profile selected by *Mitrovica and Forte (1997)* on the basis of the fit to geoid and post-glacial rebound data, neglecting the effect of lateral viscosity variations (*Moucha et al., 2007*).

Replacing δc in equation 3 with its harmonic expansion 4, and expressing $\delta v_P/v_P$ and $\delta v_S/v_S$ as a linear combination of voxels (*Soldati et al., 2012*), the system of equations above may be summarized in the compact formula

$$\begin{pmatrix} \delta t_S \\ \delta t_P \\ \delta t_{PcP} \end{pmatrix} = \begin{pmatrix} 0 & A_{mantle}^S \\ A_{mantle}^P & 0 \\ A_{mantle}^{PcP} & A_{CMB}^{PcP} K_{PcP} \end{pmatrix} \cdot \begin{pmatrix} \delta v_P \\ \delta v_S \end{pmatrix}, \quad (5)$$

where the submatrices A_{mantle}^p and A_{CMB}^p represent the sensitivity of the seismic phase p (S, P, PcP) to mantle and CMB structure, respectively. Taking into account the mechanical relationship between deep mantle heterogeneity and CMB deflections allows the equations for δv_P and δv_S to be coupled via the sensitivity kernels K_{PcP} , and the resulting velocity models to be physically consistent with each other and with the CMB topography, obtained from equation 4 after scaling δv_S to $\delta \rho$. Relative density anomalies are indeed assumed to be proportional to shear-velocity ones through the constant factor $\delta \ln \rho / \delta \ln v_S = 0.27$ (average

of the profile proposed by Karato (1993) on the basis of mineralogical experiments). This simplistic assumption has been proven to be a good approximation of the mineral properties of the mantle through a series of experiments using the depth-dependent scaling factors employed by *Simmons et al.* (2009) and shown in their Figure 3.

We employ a database of ~ 630000 summary rays travel times of P waves and ~ 63000 of PcP waves extracted from the International Seismological Centre (ISC) bulletin, as corrected by *Antolik et al.* (2001), plus the ~ 170000 S waves arrivals computed by Houser et al. (2008) via a cross correlation technique. With respect to our previous study (*Soldati et al.*, 2012), we neglect here the data set of PKP travel times. The combination of the P - and S -wave data sets provides a better coverage throughout most of the mantle and allows us to get a significant increase in information. The velocity models are parametrized in terms of 15 layers (200 km-thick) of 1656 equal area voxels each (plus 1656 pixels to describe the CMB topography), measuring $5^\circ \times 5^\circ$ at the equator.

Multi-parameter inversions may be strongly influenced by factors like the weighting associated with different data or the misfit functions. Following *Soldati et al.* (2012), we assign to each travel time a weight (exponential function) depending on its deviation from PREM predictions. We select different values for the cutoff of our subsets of data, depending on their standard deviation, while we assume the same relative weight for each data set, despite the difference in the number of phase arrivals. The LSQR linearized inversion is regularized via radial and lateral roughness damping only (e.g. *Boschi and Dziewonski*, 1999).

3 Results

3.1 Mantle velocity and CMB topography models

We first compute δv_P and δv_S models of the Earth’s mantle and the corresponding CMB topography map via classic tomographic (T) inversions of the entire data set. In this case, the inverse problem 5 is totally decoupled (the geodynamic part of the matrix being null) and therefore equivalent to separate inversions of the S , P , PcP data sets. The solution model δv_P , shown in Figure 1 (column 1), is consistent with previously published ones (e.g. *Boschi and Dziewonski*, 1999, 2000; *Soldati et al.*, 2003). The same holds true for the model of mantle δv_S anomalies (column 3), in close agreement with the results found by *Kennett et al.* (1998); *Ritsema et al.* (1999); *Houser et al.* (2008); *Soldati et al.* (2012); *Auer et al.* (2014). Compared to the δv_P solutions, δv_S anomalies are considerably larger and the slow structures

126 beneath Southern Africa and Pacific Ocean emerge much more clearly; this suggests a stronger
 127 sensitivity of v_S to thermal anomalies.

128 We also show in Figure 1 the results of inverting the entire data set with the tomographic-
 129 geodynamic (TG) approach described in Section 2, assuming the radial viscosity profile by
 130 *Mitrovica and Forte* (1997). The δv_P and δv_S anomaly maps so obtained (column 2, 4, respec-
 131 tively) are almost coincident with the corresponding T models, derived without considering
 132 the mantle-CMB coupling.

133 As shown in Table 1, the TG model of δv_P achieves similar or higher variance reduction
 134 to the different subsets of data than that associated with the purely tomographic T one, and
 135 similarly the TG model of δv_S fits the S data better than the corresponding T model. This
 136 is a nontrivial result, since the number of free parameters is reduced in the TG inversion (e.g.
 137 *Soldati et al.*, 2012). That a model achieves a higher variance reduction with a lower number
 138 of free parameters is an indication that the added regularization provided by geodynamic
 139 constraints helps the solution to converge to a better model. The δv_P model of Figure 1 is in
 140 agreement with those found by *Soldati et al.* (2012) using the same approach to invert solely
 141 the ISC P -wave data set, and scaling δv_P (instead of δv_S) to $\delta \rho$ heterogeneity.

142 We show in Figure 2a our model of CMB topography, obtained directly from the T
 143 inversion, and in 2b the model we obtained stepwise from the TG inversion and integration
 144 of δv_S as described by *Soldati et al.* (2012). Both models are dominated by harmonic degree 2,
 145 corresponding to systematically negative CMB topography under the circumpacific ring. This
 146 is generally consistent with the observations of, e.g., *Morelli and Dziewonski* (1987); *Boschi*
 147 *and Dziewonski* (1999); *Rodgers and Wahr* (1993); *Forte et al.* (1995); *Soldati et al.* (2003).
 148 With respect to the T model, the TG one is characterized by stronger CMB depression at very
 149 high and low latitudes, and an overall smaller amplitude (with a depth-to-valley amplitude
 150 of 13.3 km, vs. a value of 14.6 km for the T model).

151 3.2 Comparison of v_P and v_S models

152 We show in Figure 3a the correlation between shear and compressional velocity anomalies
 153 and that between shear and bulk sound velocity (v_ϕ) anomalies for T and TG mantle mod-
 154 els. We observe a positive correlation between δv_P and δv_S throughout the mantle, and
 155 a corresponding decorrelation between δv_ϕ and δv_S , with little or no difference between T
 156 and TG inversions. All the correlation curves tend to decrease in the midmantle and at D''
 157 depth, as also found by *Della Mora et al.* (2011) inverting direct P and S waves with a clas-

158 sic tomographic approach and with finer radial parameterization. The simultaneous drop of
 159 correlation δv_P - δv_S and anticorrelation δv_ϕ - δv_S at the base of the mantle has been already
 160 observed by several studies (e.g. *Su and Dziewonski*, 1997; *Kennett et al.*, 1998; *Becker and*
 161 *Boschi*, 2002; *Antolik et al.*, 2003; *Simmons et al.*, 2010), and may indicate a non-thermal
 162 origin for the velocity heterogeneities (*Karato and Karki*, 2001; *Saltzer et al.*, 2001; *Hirose*,
 163 2006).

164 Another indicator of potential compositional heterogeneity is the ratio $R_{S,P}$ between shear
 165 and compressional velocity anomalies, which may be computed following various approaches.
 166 Figure 3b shows in red the depth dependence of $R_{S,P}$ computed from Equation 1 of *Della*
 167 *Mora et al.* (2011), while blue and green curves refer to the ratio of the RMS of models δv_S and
 168 δv_P , multiplied/not multiplied by the Pearson's correlation coefficient. This coefficient acts as
 169 a weight: wherever the correlation between δv_P and δv_S is poor, it is meaningless to assume
 170 that they are proportional, and to compute their ratio. Independently from the formula
 171 applied, and from the approach used to invert the data (T/TG), $R_{S,P}$ is positive throughout
 172 the mantle and increases (of different amounts depending on the formula used to compute
 173 it) below 2100 km depth, with a relative maximum around 2600 km depth. This trend is
 174 consistent with previous studies (*Saltzer et al.*, 2001; *Della Mora et al.*, 2011), and suggests
 175 at least a partial influence of chemical composition on the seismic velocity heterogeneities
 176 (*Karato*, 2003). Again, the importance of compositional effects does not appear to be altered
 177 by the inclusion of geodynamic constraints (TG) in the inverse problem.

178 4 Discussion

179 Using a large data set of P , PcP and S travel time data and the fact that CMB deflections
 180 should be gravitationally related to velocity structure of the deep mantle, we apply the
 181 method by *Soldati et al.* (2012) to the case of joint inversion for both compressional and
 182 shear velocity anomalies. In this approach, CMB topography is not explicitly inverted for,
 183 but required to be coupled to mapped mantle heterogeneities according to the viscous flow
 184 theory of *Forte and Peltier* (1991). This is a entirely new approach to joint inversion, in that
 185 it does not require any assumption on the scaling between v_P and v_S anomalies, which are
 186 instead only coupled by the CMB. This study represents a significant extension to that of
 187 *Soldati et al.* (2012), in that no scaling was assumed between v_P and density anomalies, but
 188 rather between v_S and density anomalies. Interestingly, despite this important difference in

189 approach, the results of *Soldati et al.* (2012) are confirmed: neither the CMB topography, nor
 190 the pattern of mapped low/high velocity are strongly perturbed. The robustness of this result
 191 suggests that, while considerable compositional heterogeneity exists in the lowermost mantle,
 192 a flow model controlled by temperature/density heterogeneity alone is consistent with the
 193 mapped CMB topography and pattern of low/high seismic velocity in the mantle.

194 Despite the fundamental differences between the results presented here and those of *Sol-*
 195 *dati et al.* (2012), this is an indication that the global pattern of convective currents in the
 196 lower mantle is governed by thermal, rather than thermochemical variations.

197 References

- 198 Antolik, M., G. Ekström, and A. M. Dziewonski (2001), Global event location with full and
 199 sparse data sets using three-dimensional models of mantle P-wave velocity, *Pure Appl.*
 200 *Geophysics*, *158*(1-2), 291–317.
- 201 Antolik, M., Y. J. Gu, G. Ekström, and A. M. Dziewonski (2003), J362d28: a new joint model
 202 of compressional and shear velocity in the Earth’s mantle, *Geophysical J. Int.*, *153*(2), 443–
 203 466.
- 204 Auer, L., L. Boschi, T. W. Becker, T. Nissen-Meyer, and D. Giardini (2014), Savani: A
 205 variable-resolution whole-mantle model of anisotropic shear-velocity variations based on
 206 multiple datasets, *J. Geophys. Res.*, doi:10.1002/2013JB010773.
- 207 Becker, T. W., and L. Boschi (2002), A comparison of tomographic and geodynamic mantle
 208 models, *Geochemistry Geophysics Geosystems*, *3*, 10.129/2001GC000168.
- 209 Boschi, L., and A. M. Dziewonski (1999), High- and low-resolution images of the Earth’s
 210 mantle: Implications of different approaches to tomographic modeling, *J. Geophys. Res.*,
 211 *104*(B11), 25,567–25,594.
- 212 Boschi, L., and A. M. Dziewonski (2000), Whole Earth tomography from delay times of P,
 213 PcP, and PKP phases: Lateral heterogeneities in the outer core or radial anisotropy in the
 214 mantle?, *J. Geophys. Res.*, *105*(B6), 13,675–13,696.
- 215 Boschi, L., T. W. Becker, and B. Steinberger (2007), Mantle plumes: Dynamic mod-
 216 els and seismic images, *Geochemistry Geophysics Geosystems*, *8*, Q10,006, doi:10.1029/
 217 2007GC001733.

218 Boschi, L., T. W. Becker, and B. Steinberger (2008), On the statistical significance of cor-
219 relations between synthetic mantle plumes and tomographic models, *Phy. Earth. Planet.*
220 *Int.*, *167*(3-4), 230–238.

221 Della Mora, S., L. Boschi, P. J. Tackley, T. Nakagawa, and D. Giardini (2011), Low seismic
222 resolution cannot explain S/P velocity decorrelation in the lower mantle, *Geophys. Res.*
223 *Lett.*, doi:10.1029/2011GL047559.

224 Deschamps, F., and J. Trampert (2003), Mantle tomography and its relation to temperature
225 and composition, *Phys. Earth Planet. Inter.*, *140*(4), 277–291.

226 Dziewonski, A. M., and D. L. Anderson (1981), Preliminary reference earth model, *Phy.*
227 *Earth. Planet. Int.*, *25*(4), 297–356.

228 Dziewonski, A. M., and F. Gilbert (1976), Effect of small, aspherical perturbations on travel
229 times and a re-examination of corrections for ellipticity, *Geophys. J. Roy. Astron. Soc.*,
230 *44*(1), 7–17.

231 Forte, A. M., and W. L. Peltier (1991), Viscous-flow models of global geophysical observables
232 1: Forward problems, *J. Geophys. Res.*, *96*, 20,131–20,159.

233 Forte, A. M., J. X. Mitrovica, and R. L. Woodward (1995), Seismic-geodynamic determination
234 of the origin of excess ellipticity of the core-mantle boundary, *Geophys. Res. Lett.*, *22*(9),
235 1013–1016.

236 Goes, S., R. Govers, and P. Vacher (2000), Shallow mantle temperatures under Europe from
237 P and S wave tomography, *J. Geophys. Res.*, *105*(B5), 11,153–11,169.

238 Hirose, K. (2006), Postperovskite phase transition and its geophysical implications, *Rev.*
239 *Geoph.*, *44*.

240 Houser, C., G. Masters, P. Shearer, and G. Laske (2008), Shear and compressional velocity
241 models of the mantle from cluster analysis of long-period waveforms, *Geophysical J. Int.*,
242 *174*(1), 195–212.

243 Karato, S. I. (2003), *The dynamic structure of the deep Earth*, Princeton University Press.

244 Karato, S. I., and B. B. Karki (2001), Origin of lateral variation of seismic wave velocities
245 and density in the deep mantle, *J. Geophys. Res.*, *106*(B10), 21,771–21,783.

246 Kennett, B., and A. Gorbatov (2004), Seismic heterogeneity in the mantle - strong shear
 247 wave signature of slabs from joint tomography, *Phys. Earth Planet. Inter.*, *146*(1-2), 87–
 248 100, doi:10.1016/j.pepi.2003.07.033.

249 Kennett, B. L. N., S. Widiyantoro, and R. D. van der Hilst (1998), Joint seismic tomography
 250 for bulk sound and shear wave speed in the Earth’s mantle, *J. Geophys. Res.*, *103*(B6),
 251 12,469–12,493.

252 Mitrovica, J., and A. M. Forte (1997), The radial profile of mantle viscosity: results from the
 253 joint inversion of convection and post-glacial rebound observables, *J. Geophys. Res.*, *102*,
 254 2751–2769.

255 Morelli, A., and A. M. Dziewonski (1987), Topography of the core-mantle boundary and
 256 lateral homogeneity of the liquid core, *Nature*, *325*, 678–683.

257 Moucha, R., A. M. Forte, J. X. Mitrovica, and A. Daradich (2007), Lateral variations in man-
 258 tle rheology: implications for convection related surface observables and inferred viscosity
 259 models, *Geophys. J. Int.*, *169*(1), 113–135.

260 Paige, C., and M. Saunders (1982), LSQR - an algorithm for sparse linear-equations and
 261 sparse least-squares, *ACM Transactions on Mathematical Software*, *8*.

262 Ritsema, J., H. J. V. Heijst, and J. H. Woodhouse (1999), Complex shear wave velocity
 263 structure imaged beneath Africa and iceland, *Science*, *286*(5446), 1925–1928.

264 Rodgers, A., and J. Wahr (1993), Inference of core-mantle boundary topography from ISC
 265 PcP and PKP traveltimes, *Geophys. J. Int.*, *115*, 991–1011.

266 Saltzer, R. L., R. D. van der Hilst, and H. Karason (2001), Comparing P and S wave hetero-
 267 geneity in the mantle, *Geophys. Res. Lett.*, *28*(7), 1335–1338.

268 Simmons, N. A., A. M. Forte, and S. P. Grand (2009), Joint seismic, geodynamic and mineral
 269 physical constraints on three-dimensional mantle heterogeneity: Implications for the rela-
 270 tive importance of thermal versus compositional heterogeneity, *Geophysical J. Int.*, *177*(3),
 271 1284–1304.

272 Simmons, N. A., A. M. Forte, L. Boschi, and S. P. Grand (2010), GyPSuM: A joint tomo-
 273 graphic model of mantle density and seismic wave speeds., *J. Geophys. Res.*, *115*, B12,310,
 274 doi:10.1029/2010JB007,631.

- 275 Soldati, G., L. Boschi, and A. Piersanti (2003), Outer core density heterogeneity and the
276 discrepancy between PKP and PcP travel time observations, *Geophys. Res. Lett.*, *30*(4),
277 1190.
- 278 Soldati, G., L. Boschi, and A. Forte (2012), Tomography of core-mantle boundary and low-
279 ermost mantle coupled by geodynamics, *Geophys. J. Int.*, doi:10.1111/j.1365-246X.2012.
280 05413.x.
- 281 Su, W. J., and A. M. Dziewonski (1997), Simultaneous inversion for 3-D variations in shear
282 and bulk velocity in the mantle, *Phy. Earth. Planet. Int.*, *100*(1-4), 135–156.
- 283 Trampert, J., F. Deschamps, J. Resovsky, and D. Yuen (2004), Probabilistic tomography
284 maps chemical heterogeneities throughout the lower mantle, *Science*, *306*(5697), 853–856.
- 285 Vasco, D. W., and L. Johnson (1998), Whole earth structure estimated from seismic arrival
286 times, *J. Geophys. Res.*, *103*, 2633–2671.

Mantle/CMB model	S (%)	P (%)	PcP (%)
δv_P (T)	/	25.4	10.2
δv_S (T)	60.3	/	/
$\delta v_P, \delta v_S$ (TG)	61.3	25.3	10.5

Table 1: Variance reductions of different databases (columns) achieved by models T and TG , as indicated. Note that model T consists of the results of two entirely decoupled inversions (first and second row), as in *Della Mora et al.* (2011), while model TG (third row) is an individual, joint model including both v_S and v_P anomalies.

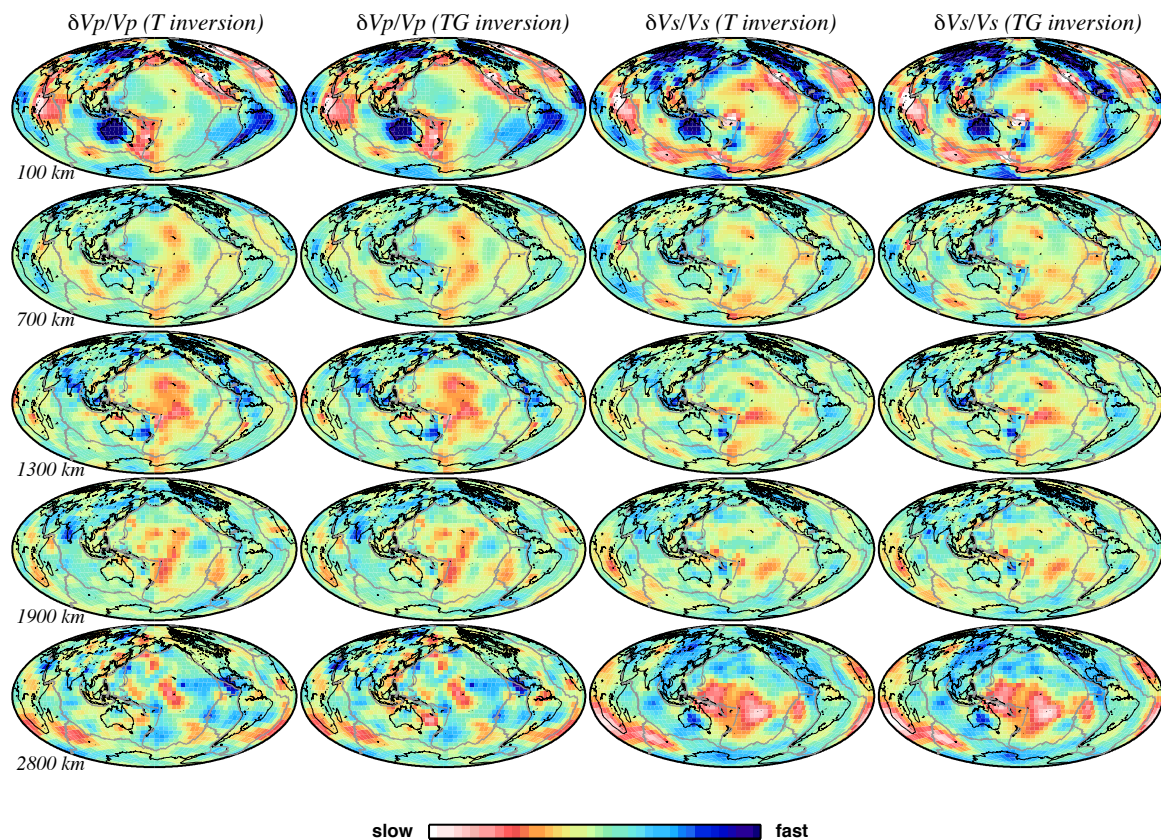


Figure 1: Maps of relative v_P and v_S velocity variations (%) at (top to bottom) five different depths in the mantle, obtained from the entire P , PcP , S data set. The maps are derived by jointly inverting ISC and *Houser et al.* (2008) data with a purely tomographic approach T (columns 1, 3), and with a tomographic-geodynamic approach TG (columns 2, 4). The scale for the δv_P maps is $\pm 2\%$ at 100 km depth and $\pm 1\%$ elsewhere, that for the δv_S maps is $\pm 4\%$ at 100 km depth and $\pm 3\%$ elsewhere; blue regions denote higher than average velocity, and red regions denote lower than average velocity.

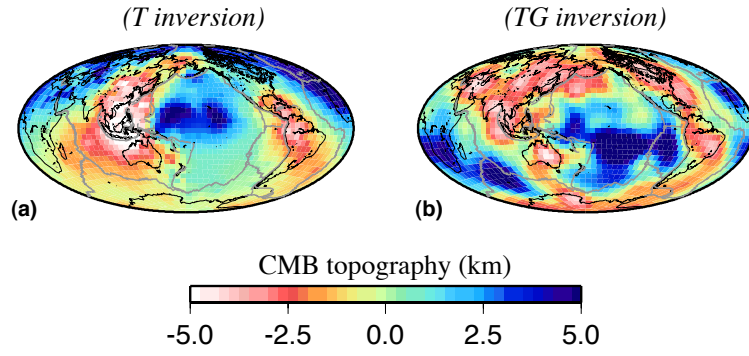


Figure 2: Maps of CMB topography (km) obtained from the entire P , PcP , S data set. The maps are obtained by jointly inverting ISC and *Houser et al.* (2008) data with a purely tomographic approach T (a), and with a tomographic-geodynamic approach TG (b), integrating the mantle shear velocity anomaly modulated by the sensitivity kernels as in equation 4.

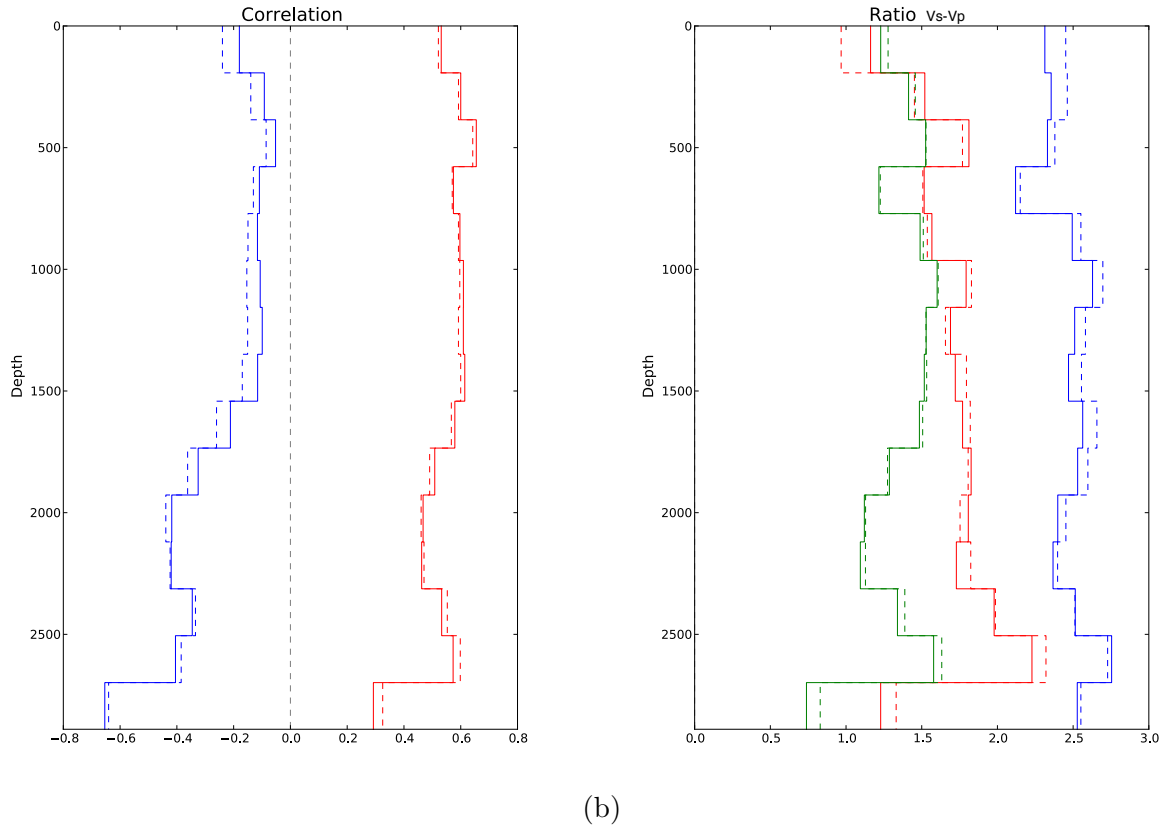


Figure 3: (a) Radial correlation between v_S and v_P heterogeneities (red) and between bulk sound velocity and v_P heterogeneities (blue), for T (solid) and TG (dashed) mantle models. (b) Ratio between v_S and v_P heterogeneities as a function of depth (red); same ratio computed via the RMS of both models as a function of depth (blue); same as the latter, but multiplied by the v_S - v_P correlation coefficients (panel a, red lines) (green). Solid and dashed lines as in panel a.