



HAL
open science

THE FOKKER-PLANCK EQUATION WITH SUBCRITICAL CONFINEMENT FORCE

Otared Kavian, Stéphane Mischler, Mamadou Ndao

► **To cite this version:**

Otared Kavian, Stéphane Mischler, Mamadou Ndao. THE FOKKER-PLANCK EQUATION WITH SUBCRITICAL CONFINEMENT FORCE. *Journal de Mathématiques Pures et Appliquées*, 2021, 151, pp.171-211. 10.1016/j.matpur.2021.04.007. hal-01241680v3

HAL Id: hal-01241680

<https://hal.science/hal-01241680v3>

Submitted on 8 Oct 2020

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

THE FOKKER-PLANCK EQUATION WITH SUBCRITICAL CONFINEMENT FORCE

O. KAVIAN, S. MISCHLER, M. NDAO

ABSTRACT. We consider the Fokker-Planck equation with subcritical confinement force field which may not derive from a potential function. We prove the existence of an equilibrium (in the case of a general force) and we establish some subgeometric rate of convergence to the equilibrium (depending on the space to which belongs the initial datum) in many spaces. Our results generalize similar results introduced by Toscani, Villani [33] and Röckner, Wang [31] for some forces associated to a potential and extended by Douc, Fort, Guillin [12] and Bakry, Cattiaux, Guillin [4] for some general forces: the spaces are more general, the rates are sharper.

Version of September 16, 2019

AMS Subject Classification (2000): 47D06 One-parameter semigroups and linear evolution equations [See also 34G10, 34K30], 35P15 Estimation of eigenvalues, upper and lower bounds, 47H20 Semigroups of nonlinear operators [See also 37L05, 47J35, 54H15, 58D07], 35Q84 Fokker-Planck equations.

Keywords: Fokker-Planck equation; semigroup; weak Poincaré inequality; weak dissipativity; Krein-Rutman theorem; spectral mapping theorem.

CONTENTS

1.	Introduction	1
2.	The splitting $\mathcal{L} = \mathcal{A} + \mathcal{B}$ and growth estimates on $S_{\mathcal{B}}$	9
3.	Boundedness of $S_{\mathcal{L}}$	19
4.	The Case 1 when a Poincaré inequality holds	20
5.	The general case - the stationary problem	23
6.	The general case - decay estimate	28
	References	37

1. INTRODUCTION

In the present work, we consider the Fokker-Planck equation

$$(1.1) \quad \partial_t f = \mathcal{L}f = \Delta f + \operatorname{div}(f \mathbf{F})$$

on the density function $f = f(t, x)$, $t > 0$, $x \in \mathbb{R}^d$, $d \geq 1$, in the case of a subcritical confinement force \mathbf{F} , and which is complemented with an initial condition

$$(1.2) \quad f(0, x) = f_0(x), \quad \forall x \in \mathbb{R}^d.$$

More precisely, we will always assume that the force field $\mathbf{F} \in C^1(\mathbb{R}^d, \mathbb{R}^d)$ satisfies

$$(1.3) \quad x \cdot \mathbf{F}(x) \geq |x| \langle x \rangle^{\gamma-1}, \quad \operatorname{div}(\mathbf{F}(x)) \leq C_F |x|^{\gamma-2}, \quad \forall x \in B_{R_0}^c,$$

as well as

$$(1.4) \quad |D\mathbf{F}(x)| \leq C'_F \langle x \rangle^{\gamma-2}, \quad \forall x \in \mathbb{R}^d,$$

for some constants $C_F \geq d$, $R_0 > 0$, $C'_F > 0$ and an exponent

$$(1.5) \quad \gamma \in (0, 1).$$

Here and below, we denote $\langle x \rangle := (1 + |x|^2)^{1/2}$ for any $x \in \mathbb{R}^d$.

It is worth mentioning that we have made two normalization hypotheses by taking a diffusion coefficient equal to 1 in (1.1) as well as a lower bound constant equal to 1 in the first condition in (1.3). Of course, these two normalization hypotheses can be removed by standard scaling arguments (in the time and position variables) and thus do not restrict the generality of our analysis, but on the other hand noticeably simplify the presentation.

A typical example of a force field is the one associated to a confinement potential

$$(1.6) \quad \mathbf{F}(x) := \nabla V(x), \quad V(x) := \frac{\langle x \rangle^\gamma}{\gamma} + V_0, \quad V_0 \in \mathbb{R}.$$

In this case, we may observe that

$$(1.7) \quad G(x) := e^{-V(x)} \in L^1(\mathbb{R}^d) \cap C^2(\mathbb{R}^d),$$

is a stationary solution of (1.1), and even an equilibrium state. We may assume that G is a probability measure, by choosing the constant V_0 adequately. We recall that when \mathbf{F} is given by (1.6) with $\gamma \geq 1$, the following Poincaré inequality

$$\exists c > 0, \quad \int_{\mathbb{R}^d} |f(x)|^2 \exp(-V(x)) dx \leq c \int_{\mathbb{R}^d} |\nabla f(x)|^2 \exp(-V(x)),$$

holds for any f such that $\int_{\mathbb{R}^d} f(x) \exp(-V(x)) dx = 0$. Such a Poincaré inequality does not hold when $\gamma \in (0, 1)$, which is the case studied in this paper, but only a weak version of this inequality remains true (see [31], and below (1.14) and section 4.1). In particular, there is no spectral gap for the associated operator \mathcal{L} , nor is there an exponential trend to the equilibrium for the associated semigroup. Similarly, the classical *logarithmic Sobolev inequality* does not hold but only a *modified* version of it, see the discussion in [33, Section 2].

In the general case of a force field which is not the gradient of a potential, one may see easily that the above Fokker-Planck equation preserves positivity, that is

$$f(t, \cdot) \geq 0, \quad \forall t \geq 0, \quad \text{if } f_0 \geq 0,$$

and that it conserves mass, that is

$$(1.8) \quad \mathbb{M}(f(t, \cdot)) = \mathbb{M}(f_0), \quad \forall t \geq 0, \quad \text{with } \mathbb{M}(g) := \int_{\mathbb{R}^d} g(x) dx.$$

Moreover, the Fokker-Planck operator \mathcal{L} generates a (Markov) semigroup in many Lebesgue spaces which has a unique positive normalized steady state (or invariant probability measure).

Before stating our existence result, let us introduce some notation. For any exponent $p \in [1, \infty]$, we define the polynomial and exponential weight functions $m : \mathbb{R}^d \rightarrow \mathbb{R}_+$, by

$$(1.9) \quad m(x) := \langle x \rangle^k, \quad \text{for some } k > k^*, \text{ with } k^* := \max(d, C_F)/p',$$

where $p' := p/(p-1)$ is the conjugated exponent associated to p ,

$$(1.10) \quad m(x) := \exp(\kappa \langle x \rangle^s), \quad \text{for some } 0 < s < \gamma \text{ and } \kappa > 0,$$

or

$$(1.11) \quad m(x) := \exp(\kappa \langle x \rangle^\gamma), \quad \text{for some } \kappa \in (0, 1/\gamma),$$

as well as the associated Lebesgue spaces

$$L^p(m) = \{f \in L^1_{\text{loc}}(\mathbb{R}^d); \|f\|_{L^p(m)} := \|fm\|_{L^p} < \infty\}.$$

We also use the shorthands $L^p_k = L^p(m)$ when $m(x) = \langle x \rangle^k$. It is noteworthy that, for such a choice of weights m , we have $L^p(m) \subset L^1(\mathbb{R}^d)$.

As a first step, we have the following existence and uniqueness result.

Theorem 1.1. *For any exponent $p \in [1, \infty]$, any weight function m satisfying either of definitions (1.9), (1.10) or (1.11), and any initial datum $f_0 \in L^p(m)$, there exists a unique global solution f to the Fokker-Planck equation (1.1)–(1.2), such that for any $T > 0$,*

$$f \in C([0, T]; L^1(\mathbb{R}^d)) \cap L^\infty(0, T; L^p(m)).$$

Moreover the associated flow preserves positivity and conserves mass. Also, the operator \mathcal{L} generates a strongly continuous semigroup $S_{\mathcal{L}}(t)$ in $L^p(m)$ when $p \in [1, \infty)$.

On the other hand, there exists a unique positive, unit mass, stationary solution G such that

$$G \in L^1_1, \quad \mathbb{M}(G) = 1 \quad \text{and} \quad \Delta G + \text{div}(G\mathbf{F}) = 0.$$

More precisely, there exists $\kappa_1^* \geq \kappa_0^* := 1/\gamma$ such that for any $\kappa_0 \in (0, \kappa_0^*)$, $\kappa_1 \in (\kappa_1^*, \infty)$, there holds

$$(1.12) \quad C_1 e^{-\kappa_1 \langle x \rangle^\gamma} \leq G \leq C_0 e^{-\kappa_0 \langle x \rangle^\gamma} \quad \text{on } \mathbb{R}^d,$$

for some constructive constants $C_0, C_1 > 0$.

The well-posedness for the evolution equation is not surprising and the proof follows classical dissipativity arguments. Due to the lack of compactness of the associated semigroup, the standard Krein-Rutman theory does not apply directly in the case $\gamma \in (0, 1)$, and the existence of a stationary solution is not straightforward. We follow here a similar fixed point theorem strategy as in the recent work [30] (see also [16, 15] and the references therein) where the Fokker-Planck equation with general force field (1.1)–(1.5) in the case $\gamma \geq 1$ is considered. Our approach provides an alternative deterministic proof of the existence of a stationary solution which may also be established using the probabilistic approach developed by Douc, Fort and Guillin in [12]. Another deterministic approach is also presented in [9] (which applies to much more general Markov semigroups).

Once the existence of a stationary solution is known, we are interested in the long time behaviour of the solution $f(t, \cdot)$ to the Fokker-Planck equation (1.1). For the sake of clarity, we consider separately the following two cases.

Case 1. Following [31, Example 1.4(c)], we consider the case when furthermore the above steady state G fulfills a *weak weighted Poincaré-Wirtinger inequality*. More precisely, in this case we assume that there exist some constants $R_0, c_1, c_2 > 0$ such that the function $V := -\log G \in C^1(\mathbb{R}^d)$ satisfies

$$(1.13) \quad \forall x \in B_{R_0}^c, \quad c_1 |x|^\gamma \leq V(x) \leq c_2 |x|^\gamma,$$

and also that there exists $\mu > 0$ such that for any $f \in \mathcal{D}(\mathbb{R}^d)$ with $\mathbb{M}(f) = 0$, the following *weak weighted Poincaré-Wirtinger inequality* holds

$$(1.14) \quad \int |\nabla(f/G)|^2 G dx \geq \mu \int f^2 \langle x \rangle^{2\gamma-2} G^{-1} dx.$$

Inequality (1.14) holds when $c_1 = c_2$ in (1.13), see Lemma 4.1. Although this inequality probably belongs to folklore, we were not able to find a precise reference. We however refer to [31, Theorem 3.3] and [11, Theorem 2.18] where related inequalities are established.

The weak Poincaré inequality (1.14) is a consequence of a “*local Poincaré inequality*” (or “*Poincaré-Wirtinger inequality*”) together with the fact that the following Lyapunov condition (see for instance [4, 3])

$$(1.15) \quad \Delta w - \nabla V \cdot \nabla w \leq (-\zeta(x) + M\chi_R)w, \quad \forall x \in \mathbb{R}^d,$$

holds for some well chosen function $w : \mathbb{R}^d \rightarrow [1, \infty)$. Here it is assumed that M and R are two positive constants, $\chi_R(x) := \chi(x/R)$ is a truncation function defined through a certain $\chi \in \mathcal{D}(\mathbb{R}^d)$ such that $\mathbf{1}_{\{|x| \leq 1\}} \leq \chi \leq \mathbf{1}_{\{|x| \leq 2\}}$ and $\zeta(x) := \zeta_0 \langle x \rangle^{2(1-\gamma)}$ for a constant $\zeta_0 > 0$. It is worth emphasizing that, in this case, the force field \mathbf{F} can be written as

$$(1.16) \quad \mathbf{F} = \nabla V + \mathbf{F}_0, \quad \operatorname{div}(e^{-V} \mathbf{F}_0) = 0,$$

with no other specific condition on \mathbf{F}_0 except that \mathbf{F} still satisfies conditions (1.3)–(1.5). Under these circumstances, we can give a simpler proof than in the general case.

Case 2. This corresponds to the general case when \mathbf{F} satisfies only conditions (1.3)–(1.5), without any further assumption on the stationary state G , which in general cannot be determined explicitly. Using the above notations for M, R, χ_R and ζ introduced in the inequality (1.15), the assumptions (1.3)–(1.5) made on \mathbf{F} imply in particular the following inequality

$$(1.17) \quad \mathcal{L}^* m^p := \Delta m^p - \mathbf{F} \cdot \nabla m^p \leq (-\zeta(x) + M\chi_R)m^p, \quad \forall x \in \mathbb{R}^d,$$

which is another version of the Lyapunov condition (1.15).

The main and fundamental difference between these two cases is that the first one involves assumptions on the equilibrium state $G = e^{-V}$ while the second one only involves an assumption on the force field \mathbf{F} .

In the sequel, when $a(t) \geq 0$ and $b(t) \geq 0$ are two functions of time $t > 0$, we write $a(t) \lesssim b(t)$ to mean that there exists a positive constant c_0 independent of t such that one has $a(t) \leq c_0 b(t)$ for all $t > 0$.

Our main result is as follows:

Theorem 1.2. *Let \mathbf{F} satisfy (1.3)–(1.5), and let $p \in [1, \infty]$. For a weight function m satisfying either of definitions (1.9), (1.10) or (1.11), we define the subgeometric*

rate function Θ_m as follows: when $m(x) = \langle x \rangle^k$, we take $\beta \in (0, (k - k^*)/(2 - \gamma))$ arbitrary and we set

$$(1.18) \quad \Theta_m(t) := (1 + t)^{-\beta}.$$

When $m(x) = \exp(\kappa \langle x \rangle^s)$, we take $\sigma := s/(2 - \gamma)$ and we set for a $\lambda > 0$

$$(1.19) \quad \Theta_m(t) := \exp(-\lambda t^\sigma).$$

Then for any initial datum $f_0 \in L^p(m)$, the associated solution $f = f(t, x)$ to the Fokker-Planck equation (1.1) satisfies

$$(1.20) \quad \|f(t, \cdot) - \mathbb{M}(f_0)G\|_{L^p} \lesssim \Theta_m(t) \|f_0 - \mathbb{M}(f_0)G\|_{L^p(m)},$$

where we recall that $\mathbb{M}(f_0)$ denotes the mass of f_0 defined in (1.8).

Remark 1.3. We believe that Theorem 1.2 is new in the sense that the decay estimate (1.18)-(1.20) is optimal and constructive and that the Lebesgue spaces involved appear with the same exponent $p \in [1, \infty]$. This last fact is of main importance when one is interested in the stability of nonlinear evolution PDEs, see for instance [10].

We discuss below in details some of the existent literature about the same kind of results, some possible extensions and the general strategy of our approach.

Remark 1.4. When the force field \mathbf{F} and the equilibrium G satisfy similar assumptions as described in **Case 1**, several previous results are known. A less accurate rate of decay than the one given by (1.18)-(1.20), actually a decay rate of order $\mathcal{O}(t^{-(k-2)/(2(2-\gamma))})$ in L^1 -norm, has been proved by G. Toscani, C. Villani in [33, Theorem 3] under the additional assumptions that the initial datum f_0 is nonnegative, has finite energy and finite Boltzmann entropy. The proof of this first result relies on a modified and weak version of the log-Sobolev inequality. (It is known that the standard (and stronger) version of the log-Sobolev inequality only holds when $\gamma \geq 2$). In [31, Example 1.4 (c) and Section 3 for the proof], M. Röckner, F.Y. Wang have established the exponential rate

$$(1.21) \quad \|f(t, \cdot) - \mathbb{M}(f_0)G\|_{L^2} \lesssim e^{-\lambda t^{\sigma^\sharp}} \|f_0 - \mathbb{M}(f_0)G\|_{L^\infty(G^{-1})},$$

with $\sigma^\sharp := \gamma/(4 - 3\gamma)$, some $\lambda > 0$ and any $f_0 \in L^\infty(G^{-1})$. The proof of (1.21) is based on a weak Poincaré inequality and the existence of a family of entropy functionals.

Estimate (1.18)-(1.20) is more accurate than (1.21) since, when the weight function $m(x) := \exp(\kappa \langle x \rangle^\gamma)$ we have $\sigma = \gamma/(2 - \gamma)$, and the function $\Theta_m(t) := \exp(-\lambda t^{\gamma/(2-\gamma)})$ decays faster than $\exp(-\lambda t^{\sigma^\sharp})$. Also, more importantly, the same Lebesgue exponent p on both sides of the inequality is involved in (1.20) while it is not the case in (1.21). Also, as we shall see below, from (1.18)-(1.20) one easily deduces by interpolation that for any $\theta \in (0, 1)$ one has

$$\|f(t, \cdot) - \mathbb{M}(f_0)G\|_{L^p(m^\theta)} \lesssim \Theta_m(t) \|f_0 - \mathbb{M}(f_0)G\|_{L^p(m)}.$$

Remark 1.5. D. Bakry, P. Cattiaux & A. Guillin in [4] have extended the approach of M. Röckner, F.Y. Wang to the case of a general force field. More precisely, they establish an inequality, which they call *weak Lyapunov-Poincaré inequality*, in [4, Theorem 3.10 and Section 4.2.1] from which they deduce the same rate of convergence (1.21). That work is based on the Lyapunov condition method, which we have already discussed in the presentation of **Case 1** and which is closely related

to the method we use here. The Lyapunov condition method has then been extensively studied during the last decade, and we refer for instance to [4, 3, 11] and the references therein for more details.

Remark 1.6. A completely different approach has been developed by R. Douc, G. Fort & A. Guillin in [12] generalizing the classical Harris-Meyn-Tweedie theory to the weak confinement setting and adapted to both the potential case and the general force case. We also refer to the notes on the convergence of Markov processes by M. Hairer, and more precisely to [19, Theorem 4.1], for a more explicit formulation of the subgeometric convergence result and a simplified proof (starting from a slightly different set of hypotheses). It provides estimate (1.20) in the case $p := 1$ and $m := G^{-1}$ (with same rate), see [19, (7.3)]. The drawback is that the approach is definitively probabilistic and the constants are maybe not so explicit. In a forthcoming paper [9], J.A. Cañizo and S. Mischler provide a deterministic and elementary proof of the above mentioned theorem. We also refer to A. Eberle, A. Guillin & R. Zimmze [14] for recent related results.

Remark 1.7. The same kind of decay estimates remains true when, on the left hand side of (1.20), the Lebesgue norm L^p is replaced by a weighted Lebesgue norm $L^p(m^\theta)$ with $0 \leq \theta < 1$. More precisely, when $m(x) = \langle x \rangle^k$ and θ is such that $k^*/k < \theta < 1$, we choose $\beta := k(1 - \theta)/(2 - \gamma)$, and if $0 \leq \theta \leq k^*/k$, we choose $\beta \in (0, (k - k^*)/(2 - \gamma))$ arbitrary. In both cases, we define Θ_m through (1.18). When $m(x) = \exp(\kappa \langle x \rangle^s)$ the definition of the decay rate Θ_m is unchanged.

Remark 1.8. When an exponential weight function $m(x) := \exp(\kappa \langle x \rangle^\gamma)$ is considered, one could have a field force \mathbf{F} satisfying the first condition of (1.3) and $\operatorname{div}(\mathbf{F}) \leq C'_F \langle x \rangle^{\gamma' - 2}$, with $\gamma' < 2\gamma$, or $\gamma' = 2\gamma$ but with C'_F small enough, and obtain similar results. However we do not push our investigations in that direction, since the general ideas of the proof are essentially the same.

Remark 1.9. When a polynomial weight function $m(x) = \langle x \rangle^k$ is considered, the decay rate in (1.20) is given by (1.18), which is better than the decay rate one might obtain by a mere interpolation argument between $L^2(\exp(\kappa \langle x \rangle^\gamma))$ and $L^1(\mathbb{R}^d)$. More precisely, assume that when the weight function $m(x) = \exp(\kappa \langle x \rangle^\gamma)$, the function Θ_m being given by (1.19) with $s := \gamma/(2 - \gamma)$, we have (1.20), as well as the estimate $\|f(t)\|_{L^1} \lesssim \|f_0\|_{L^1}$ for any $f_0 \in L^1(\exp(\kappa \langle x \rangle^\gamma))$ such that $\mathbb{M}(f_0) = 0$. Then for any $R > 0$ we have $\mathbb{M}(f_0 \mathbf{1}_{B_R}) = -\mathbb{M}(f_0 \mathbf{1}_{B_R^c})$, and thus we may write $f_0 = f_{01} + f_{02} + f_{03}$ where

$$f_{01} := (f_0 - \mathbb{M}(f_0 \mathbf{1}_{B_R})) \mathbf{1}_{B_R}, \quad f_{02} := f_0 \mathbf{1}_{B_R^c}, \quad f_{03} := \mathbb{M}(f_0 \mathbf{1}_{B_R^c}) \mathbf{1}_{B_R}.$$

Therefore for $t > 0$, denoting by $f_j(t)$ the solution of (1.1) with initial datum f_{0j} , one has

$$\begin{aligned} \|f(t)\|_{L^1} &\leq \|f_1(t)\|_{L^1} + \|f_2(t)\|_{L^1} + \|f_3(t)\|_{L^1} \\ &\lesssim \exp(-\lambda t^{\gamma/(2-\gamma)}) \|f_{01}\|_{L^1(\exp(\kappa \langle x \rangle^\gamma))} + \|f_{02}\|_{L^1} + \|f_{03}\|_{L^1} \\ &\lesssim \exp(-\lambda t^{\gamma/(2-\gamma)}) \|(f_0 - \mathbb{M}(f_0 \mathbf{1}_{B_R})) \mathbf{1}_{B_R}\|_{L^1(\exp(\kappa \langle x \rangle^\gamma))} \\ &\quad + \|\mathbb{M}(f_0 \mathbf{1}_{B_R^c}) \mathbf{1}_{B_R}\|_{L^1} + \|f_0 \mathbf{1}_{B_R^c}\|_{L^1} \\ &\lesssim \exp(-\lambda t^{\gamma/(2-\gamma)}) e^{\kappa \langle R \rangle^\gamma} \|f_0\|_{L^1} + (R^{d-k} + R^{-k}) \|f\|_{L_k^1} \\ &\lesssim \left(\exp(-\lambda t^{\gamma/(2-\gamma)} + \kappa \langle R \rangle^\gamma) + R^{d-k} \right) \|f\|_{L_k^1}. \end{aligned}$$

Assuming that $t > (2\kappa/\lambda)^{(2-\gamma)/\gamma}$, we may choose R so that $\kappa\langle R \rangle^\gamma = \lambda t^{\gamma/(2-\gamma)}/2$, we find that when $k > d$, for any $t > (2\kappa/\lambda)^{(2-\gamma)/\gamma}$ we have

$$\|f(t)\|_{L^1} \lesssim t^{-\frac{(k-d)}{2-\gamma}} \|f_0\|_{L^1_k}.$$

This decay estimate is not as sharp as the one given by Theorem 1.2 when the weight function is $m := \langle x \rangle^k$, since according to the definition (1.18) in this case we have actually a decay rate of $\mathcal{O}(t^{-K})$ for any $K \in (0, k/(2-\gamma))$. \square

Remark 1.10. In a few previous papers, due in particular to R.E. Caflisch [7, 8], T.M. Liggett [23], G. Toscani & C. Villani [33], Y. Guo [18] or K. Aoki & F. Golse [1], a certain number of models, arising from statistical physics, has been considered for which only subgeometric (but not geometric) rate of decay to the equilibrium can be established. As it is the present case, one can associate to each of these models a linear operator which does not enjoy any spectral gap in its spectrum set and that is the reason why exponential rate of convergence fails.

An abstract theory for non-uniformly exponentially stable semigroups (with non exponential decay rate) has also been recently developed and we refer the interested reader to [5, 6] and the references therein. We finally refer to K. Carrapatoso & S. Mischler [10] where similar semigroup analysis as here is developed and applied in order to establish the well-posedness of the Landau equation in large spaces.

Let us briefly explain the main ideas behind our method of proof. In **Case 1**, and as a first step, we may use the argument introduced in M. Röckner & F.Y. Wang [31] (see also O. Kaviani [20, Lemma 1.3]) which we briefly recall now. We consider three Banach spaces E_2 , E_1 and E_0 , such that $E_2 \subset E_1 \subset E_0 \subset L^1$, and more precisely E_1 is an interpolation space of order $1 - 1/\alpha$ between E_0 and E_2 for some $\alpha \in (1, \infty)$, that is

$$(1.22) \quad \|f\|_{E_1} \leq C_\alpha \|f\|_{E_0}^{1/\alpha} \|f\|_{E_2}^{1-1/\alpha}, \quad \forall f \in E_2,$$

and such that the semigroup $S_{\mathcal{L}}(t)$ associated to the Fokker-Planck equation can be solved in each of these spaces. Moreover, assume that for any $f_0 \in E_2$, the solution $S_{\mathcal{L}}(t)f_0 = f(t)$ to the Fokker-Planck equation (1.1)–(1.2) satisfies the following two differential inequalities

$$(1.23) \quad \frac{d}{dt} \|f(t)\|_{E_1} \leq -\lambda \|f(t)\|_{E_0}, \quad \frac{d}{dt} \|f(t)\|_{E_2} \leq 0,$$

for some constant $\lambda > 0$. Using the fact that $\|f(t)\|_{E_2} \leq \|f_0\|_{E_2}$, as a consequence of the above second differential inequality, together with (1.22), we obtain the closed differential inequality

$$\frac{d}{dt} \|f(t)\|_{E_1} \leq -\lambda C_\alpha^{-\alpha} \|f_0\|_{E_2}^{-(\alpha-1)} \|f(t)\|_{E_1}^\alpha.$$

We may readily integrate this inequality and we obtain the estimate

$$(1.24) \quad \|f(t)\|_{E_1} \lesssim t^{-1/(\alpha-1)} \|f_0\|_{E_2}.$$

Now, choosing $E_1 = L^2(G^{-1/2})$, $E_0 := L^2(G^{-1/2}\langle x \rangle^{\gamma-1})$ and $E_2 = L^\infty(G^{-1})$, one may see that the first differential inequality in (1.23) is an immediate consequence of the weak Poincaré inequality (1.14). The second differential inequality is a kind of generalized relative entropy principle (see [31, 24]). The above estimate (1.24) is a somewhat rough variant of estimate (1.19). It is noteworthy that for $\alpha \in (1, 2)$,

we get that the associated semigroup $S_{\mathcal{L}}$ defined by $S_{\mathcal{L}}(t)f_0 = f(t)$ satisfies in particular $\|S_{\mathcal{L}}\|_{E_2 \rightarrow E_1} \in L^1(0, \infty)$.

We then generalize the decay estimate to a wider class of Banach spaces by adapting the extension theory introduced and developed in C. Mouho [29], M.P. Gualdani, S. Mischler & C. Mouhot [17] and S. Mischler & C. Mouhot [26], and used in M. Ndao [30] for the case $1 \leq \gamma \leq 2$. In order to do so, we consider the two Banach spaces $\mathcal{E}_2 := L^p(m)$ and $\mathcal{E}_1 := L^p$, and we introduce a splitting $\mathcal{L} = \mathcal{A} + \mathcal{B}$, where \mathcal{A} is an appropriately defined bounded operator so that \mathcal{B} becomes a dissipative operator. Then, we show that for $i = 1$ and $i = 2$

$$\|S_{\mathcal{B}\mathcal{A}}\|_{\mathcal{E}_i \rightarrow \mathcal{E}_i} \in L^1(\mathbb{R}_+), \quad \|S_{\mathcal{B}}\|_{\mathcal{E}_2 \rightarrow \mathcal{E}_1} \in L^1(\mathbb{R}_+), \quad \|\mathcal{A}S_{\mathcal{B}}\|_{\mathcal{E}_1 \rightarrow \mathcal{E}_1} \in L^1(\mathbb{R}_+).$$

If \mathcal{T}_i , with $i = 1, 2$, are two given operator valued, measurable functions, defined on $(0, \infty)$, we denote by

$$(\mathcal{T}_1 * \mathcal{T}_2)(t) := \int_0^t \mathcal{T}_1(\tau)\mathcal{T}_2(t - \tau) d\tau$$

their convolution on \mathbb{R}_+ . We then set $\mathcal{T}^{(*0)} := I$, $\mathcal{T}^{(*1)} := \mathcal{T}$ and, for any $k \geq 2$, $\mathcal{T}^{(*k)} := \mathcal{T}^{*(k-1)} * \mathcal{T}$. We may show that for $n \in \mathbb{N}$ sufficiently large (actually $n \geq 1 + (d/2)$ is enough), we have

$$(1.25) \quad \|(S_{\mathcal{B}\mathcal{A}})^{(*n)}\|_{E_1 \rightarrow E_1} \in L^1(\mathbb{R}_+), \quad \|(\mathcal{A}S_{\mathcal{B}})^{(*n)}\|_{E_1 \rightarrow E_1} \in L^1(\mathbb{R}_+).$$

Then, from the usual Duhamel formula, the solution of (1.1) can be written as

$$f(t) = S_{\mathcal{B}}(t)f_0 + \int_0^t S_{\mathcal{B}}(t - \tau)\mathcal{A}S_{\mathcal{L}}(\tau)f_0 d\tau.$$

Thus, using the above notations for the convolution of operator valued functions, we have $S_{\mathcal{L}} = S_{\mathcal{B}} + S_{\mathcal{B}} * (\mathcal{A}S_{\mathcal{L}})$, and interchanging the role played by \mathcal{L} and \mathcal{B} in this expression, we get the following operator versions of Duhamel formulas

$$(1.26) \quad S_{\mathcal{L}} = S_{\mathcal{B}} + S_{\mathcal{B}} * (\mathcal{A}S_{\mathcal{L}}) = S_{\mathcal{B}} + (S_{\mathcal{B}\mathcal{A}}) * S_{\mathcal{L}}$$

$$(1.27) \quad = S_{\mathcal{B}} + S_{\mathcal{L}} * (\mathcal{A}S_{\mathcal{B}}) = S_{\mathcal{B}} + (S_{\mathcal{L}\mathcal{A}}) * S_{\mathcal{B}}.$$

Upon replacing recursively $S_{\mathcal{L}}$ in either of the expressions on the right hand side by either of the Duhamel's formula, we get, for instance:

$$\begin{aligned} S_{\mathcal{L}} &= S_{\mathcal{B}} + S_{\mathcal{B}} * \mathcal{A}\{S_{\mathcal{B}} + (S_{\mathcal{B}\mathcal{A}}) * S_{\mathcal{L}}\} \\ &= S_{\mathcal{B}} + (S_{\mathcal{B}\mathcal{A}}) * S_{\mathcal{B}} + (S_{\mathcal{B}\mathcal{A}})^{(*2)} * S_{\mathcal{L}}. \end{aligned}$$

By induction on the integers $n_1 \geq 0$, $n_2 \geq 0$ and $n_1 + n_2 \geq 1$, we thus obtain

$$(1.28) \quad S_{\mathcal{L}} = \sum_{k=0}^{n_1+n_2-1} S_{\mathcal{B}} * (\mathcal{A}S_{\mathcal{B}})^{(*k)} + (S_{\mathcal{B}\mathcal{A}})^{(*n_1)} * S_{\mathcal{L}} * (\mathcal{A}S_{\mathcal{B}})^{(*n_2)}.$$

Using the above formulas (1.28) and estimates (1.25), as well as the decay estimate (1.24) for initial data in the space E_2 , we conclude that $\|S_{\mathcal{L}}\|_{E_2 \rightarrow E_1} \in L^1(\mathbb{R}_+)$ which is nothing but a rough version of the estimates presented in Theorem 1.2. While the method leading to (1.24) in E_i can be performed only in very specific (Hilbert) spaces, the above extension method is very general and can be used in a large class of Banach spaces \mathcal{E}_i (once we already know the decay in one pair of spaces (E_1, E_2)).

On the other hand, in **Case 2**, we start proving an equivalent to estimate (1.24) in one appropriate pair of (small) spaces. We then argue similarly as in the previous case. We present three different strategies in order to establish an equivalent to estimate (1.24). On the one hand, we may adapt the strategy of M. Röckner and F.Y. Wang introduced to deal with **Case 1** by using a generalization of the weak Poincaré inequality (1.14) which also holds in the case of a general force field and which has been established by D. Bakry, P. Cattiaux & A. Guillin in [4].

A second approach consists in using the generalization of the well-known Doebelin-Harris-Meyn-Tweedie theory by R. Douc, G. Fort & A. Guillin in [12] to the so-called *subgeometric* framework, namely to the case when exponential trend to the equilibrium is not available.

A third and more original way consists in adapting the Krein-Rutman theory to the present context. On the one hand, it is a simple version of the Krein-Rutman theory because the equation is mass conserving, a property which implies that the largest eigenvalue of \mathcal{L} is $\lambda_1 = 0$. On the other hand, it is not a classical version because the operator \mathcal{L} does not have a compact resolvent (however it has power-compact resolvent in the sense of J. Voigt [34]) and, more importantly, 0 is not necessarily an isolated point in the spectrum. First adapting (from [15, 28] for instance) some more or less standard arguments, we prove that there exists G , a unique stationary solution of (1.1) which is positive, has unit mass and is such that $G \in L^\infty(\exp(\kappa\langle x \rangle^\gamma))$, for all $\kappa \in (0, 1/\gamma)$. Next, we prove an estimate similar to (1.24) by establishing a set of accurate estimates on the resolvent operators $\mathcal{R}_\mathcal{B}(z)$, $R_\mathcal{L}(z)$ and by using the iterated Duhamel formula

$$S_\mathcal{L} = \sum_{k=0}^5 S_\mathcal{B} * (\mathcal{A}S_\mathcal{B})^{(*k)} + S_\mathcal{L} * (\mathcal{A}S_\mathcal{B})^{(*6)},$$

together with the inverse Laplace formula

$$S_\mathcal{L} * (\mathcal{A}S_\mathcal{B})^{(*6)}(t) = \frac{i}{2\pi} \frac{1}{t^n} \int_{-i\infty}^{+i\infty} e^{zt} \frac{d^n}{dz^n} [R_\mathcal{L}(z)(\mathcal{A}R_\mathcal{B}(z))^6] dz,$$

which holds true for any time $t > 0$ and any integer n .

To finish this introduction, let us describe the plan of the paper. In Section 2, we introduce an appropriate splitting $\mathcal{L} = \mathcal{A} + \mathcal{B}$ and present the main estimates on the semigroup $S_\mathcal{B}$. In Section 3, we deduce that the semigroup $S_\mathcal{L}$ is bounded in the spaces $L^p(m)$. In Section 4, the proof of Theorem 1.2 is carried out in the case when a weak Poincaré inequality (1.14) is satisfied (**Case 1**). In Section 5, the second part of the proof of Theorem 1.1 on the stationary problem is presented in the general case (**Case 2**). Finally, Section 6 is devoted to the proof of Theorem 1.2 in the general case.

Acknowledgements. The second author's work is supported by the french "ANR blanche" project Stab: ANR-12-BS01-0019. We thank M. Hairer and A. Guillin for enlightening discussions on several results in relation with the present work.

2. THE SPLITTING $\mathcal{L} = \mathcal{A} + \mathcal{B}$ AND GROWTH ESTIMATES ON $S_\mathcal{B}$

We introduce the splitting of the operator \mathcal{L} defined by

$$(2.1) \quad \mathcal{A}f := M\chi_R f, \quad \mathcal{B}f := \mathcal{L}f - M\chi_R f$$

where M is positive constant, and for a fixed truncation function $\chi \in \mathcal{D}(\mathbb{R}^d)$ such that $1_{B(0,1)} \leq \chi \leq 1_{B(0,2)}$, and for $R > 1$ which will be chosen appropriately as well as M , we set $\chi_R(x) := \chi(x/R)$. We establish several growth and regularity estimates on the semigroup $S_{\mathcal{B}}$ and the family of operators $\mathcal{A}S_{\mathcal{B}}$ which will be of fundamental importance in the sequel.

2.1. Basic growth estimates.

Lemma 2.1. *For any exponent $p \in [1, \infty]$ and any polynomial or exponential weight function m given by (1.9), (1.10) or (1.11), we can choose R, M large enough in the definition (2.1) of \mathcal{B} such that the operator \mathcal{B} is dissipative in $L^p(m)$, namely*

$$(2.2) \quad \|S_{\mathcal{B}}(t)\|_{L^p(m) \rightarrow L^p(m)} \leq 1, \quad \forall t \geq 0.$$

Moreover, if $m(x) = \langle x \rangle^k$, set $\beta := k(1 - \theta)/(2 - \gamma)$ for $k^*/k < \theta < 1$, and $\beta \in (0, (k - k^*)/(2 - \gamma))$ arbitrary when $\theta \leq k^*/k$. Then the function Θ_m being defined by (1.18), we have

$$(2.3) \quad \|S_{\mathcal{B}}(t)\|_{L^p(m) \rightarrow L^p(m^\theta)} \lesssim \Theta_m(t).$$

If $m(x) = \exp(\kappa \langle x \rangle^s)$ satisfies (1.10) or (1.11), the above inequality holds, provided the function Θ_m is defined by

$$\Theta_m(t) := \exp(-\lambda t^{s/(2-\gamma)}),$$

where $\lambda > 0$ can be chosen arbitrarily when $s < \gamma$, and $\lambda < \lambda_*$, with

$$\lambda_* := (\kappa(1 - \theta))^{(2-2\gamma)/(2-\gamma)} (\kappa\gamma(1 - \kappa\gamma))^{\gamma/(2-\gamma)},$$

when $s = \gamma$.

Proof of Lemma 2.1. The proof is similar to the proof of [17, Lemma 3.8] and [26, Lemma 3.8].

Step 1. We first fix $p \in [1, \infty)$, assuming m is as in the statement of the Lemma. We start recalling an identity satisfied by the operator \mathcal{B} (see the proof of [26, Lemma 3.8]). For any smooth, rapidly decaying and positive function f , we have

$$\begin{aligned} \int_{\mathbb{R}^d} (\mathcal{B}f) f^{p-1} m^p dx &= \\ &= -(p-1) \int_{\mathbb{R}^d} |\nabla(mf)|^2 (mf)^{p-2} dx + \int_{\mathbb{R}^d} f^p m^p \psi_{m,p}^0(x) dx, \end{aligned}$$

with

$$(2.4) \quad \psi_{m,p}^0(x) := \frac{(2-p)}{p} \frac{\Delta m}{m} + \frac{2}{p'} \frac{|\nabla m|^2}{m^2} + \frac{1}{p'} \operatorname{div}(\mathbf{F}) - \mathbf{F} \cdot \frac{\nabla m}{m} - M \chi_R.$$

Observe that

$$\begin{aligned} \frac{\nabla m}{m} &= k\kappa x \langle x \rangle^{s-2} \\ \frac{\Delta m}{m} &= k\kappa d \langle x \rangle^{s-2} + s(s-2)\kappa |x|^2 \langle x \rangle^{s-4} + \nu |x|^2 \langle x \rangle^{2s-4}, \end{aligned}$$

where we have set

$$\begin{aligned} s &:= 0, & \kappa &:= 1, & \nu &:= k(k-2), & \text{when } m(x) &= \langle x \rangle^k, \\ k &:= s, & \nu &:= (s\kappa)^2, & & & \text{when } m(x) &= \exp(\kappa \langle x \rangle^s). \end{aligned}$$

In this latter case, for $s \in (0, \gamma]$, the third term in the definition of $\psi_{m,p}^0$ is negligible with respect to the first and second terms, and thus

$$\begin{aligned} \psi_{m,p}^0(x) |x|^{2-\gamma-s} &\xrightarrow{|x| \rightarrow \infty} -a^* := (\kappa\gamma)^2 - \kappa\gamma < 0 \quad \text{if } s = \gamma, \\ \psi_{m,p}^0(x) |x|^{2-\gamma-s} &\xrightarrow{|x| \rightarrow \infty} -a^* := -\infty \quad \text{if } 0 < s < \gamma. \end{aligned}$$

When $m = \langle x \rangle^k$, and $k > k^*(p) := C_F/p'$, the first and second terms are negligible with respect to the third term, and then

$$\limsup_{|x| \rightarrow \infty} \psi_{m,p}^0(x) |x|^{2-\gamma} \leq -a^* := \left(1 - \frac{1}{p}\right) C_F - k < 0.$$

We deduce that for any $a \in (0, a^*(m, p))$, we can choose $R > 1$ and M large enough in such a way that $\psi_{m,p}^0(x) \leq -a \langle x \rangle^{\gamma+s-2}$ for all $x \in \mathbb{R}^d$, and then

$$(2.5) \quad \int (\mathcal{B}f) f^{p-1} m^p \leq -a \int |f|^p m^p \langle x \rangle^{\gamma+s-2} - (p-1) \int |\nabla(fm)|^2 (fm)^{p-1}.$$

In particular, using only the fact that the RHS term is negative, we conclude that the operator \mathcal{B} is dissipative and we classically deduce that the semigroup $S_{\mathcal{B}}$ is well-defined on $L^p(m)$ for $p \in [1, \infty)$ and that it is a strongly continuous contraction semigroup, in other words, (2.2) holds for any $p \in [1, \infty)$. Since we may choose R, M such that the above inequality holds true for any $p \in [1, \infty)$ when $a \in (0, a^*(m, \infty))$, we may pass to the limit as $p \rightarrow \infty$ in (2.2) and we conclude that $S_{\mathcal{B}}$ is a contraction semigroup in $L^p(m)$, for any $p \in [1, \infty]$.

Step 2. Take $p \in [1, \infty)$ and $k > k^*(p) = C_F/p'$, and finally, assuming first that $\theta > k^*/k$, set $\ell := \theta k \in (k^*, k)$. If $f_0 \in L^p(m)$ with $m := \langle x \rangle^k$, denote $f(t) := S_{\mathcal{B}}(t)f_0$. Dropping the last term in (2.5), we have for $a \in (0, a^*(m, p))$

$$\frac{d}{dt} \int |f|^p \langle x \rangle^{p\ell} \leq -ap \int |f|^p \langle x \rangle^{p\ell+\gamma-2}.$$

Using Hölder's inequality

$$\int f^p \langle x \rangle^{p\ell} \leq \left(\int f^p \langle x \rangle^{p\ell+\gamma-2} \right)^\eta \left(\int f^p \langle x \rangle^{pk} \right)^{1-\eta}$$

with $\eta := (k - \ell)/[k - \ell + (2 - \gamma)/p] \in (0, 1)$, and the fact that the semigroup $S_{\mathcal{B}}$ is a contraction semigroup in L_k^p by (2.2), upon denoting $\alpha := \eta/(1 - \eta) = p(k - \ell)/(2 - \gamma)$, we get

$$\frac{d}{dt} Y_\theta(t) \leq -ap Y_\theta(t)^{(\alpha+1)/\alpha} Y_1(0)^{-1/\alpha}, \quad \text{where } Y_\tau(t) := \int f^p \langle x \rangle^{p\tau}.$$

Integrating the above differential inequality yields

$$Y_\theta(t) \leq \left(\frac{\alpha}{apt} \right)^\alpha Y_1(0),$$

which in turn implies (2.3) with $\Theta_m(t)$ replaced with $\left(\frac{(k-\ell)/(2-\gamma)}{at} \right)^{\frac{k-\ell}{2-\gamma}}$. Since for $0 \leq t \leq 1$ we have clearly $Y_\theta(t) \lesssim Y_1(0)$ the proof of (2.3) is complete when $p < \infty$ and $m(x) = \langle x \rangle^k$ and $\ell := k\theta > k^*$.

In the case where $\ell = k\theta \leq k^*$, it is enough to pick $\theta_0 > \theta$ so that $k\theta_0 > k^*$ and observe that we have $Y_\theta \leq Y_{\theta_0}$: in this way one is convinced that (2.3) holds for all $p < \infty$ and $0 \leq \theta < 1$.

We deduce the same estimate for $p = \infty$ by letting $p \rightarrow \infty$ in (2.3).

Step 3. Similarly, when the weight function m is an exponential as defined in (1.10) or (1.11), take $p \in [1, \infty)$. Given an initial datum $f_0 \in L^p(m)$, denote $f(t) := S_{\mathcal{B}}(t)f_0$, and set $Y_\theta(t) := \|f(t)\|_{L^p(m^\theta)}^p$. Thanks to the above Step 1 we have for all $t \geq 0$ and $0 < \theta \leq 1$

$$Y_\theta(t) \leq Y_\theta(0).$$

For $\rho > 0$ denote by B_ρ the ball of \mathbb{R}^d centered at the origin with radius ρ . Using the estimate (2.5) with the weight function m^θ , neglecting the last term of that inequality, we have successively

$$\begin{aligned} \frac{d}{dt} Y_\theta(t) &= p \int (\mathcal{B}f) f^{p-1} m^{p\theta} \\ &\leq -ap \int_{B_\rho} |f|^p m^{p\theta} \langle x \rangle^{\gamma+s-2} \\ &\leq -ap \langle \rho \rangle^{\gamma+s-2} \int_{B_\rho} |f|^p m^{p\theta} \\ &\leq -ap \langle \rho \rangle^{\gamma+s-2} Y_\theta + ap \langle \rho \rangle^{\gamma+s-2} \int_{B_\rho^c} |f|^p m^{p\theta} \\ &\leq -ap \langle \rho \rangle^{\gamma+s-2} Y_\theta + ap \langle \rho \rangle^{\gamma+s-2} m(\rho)^{-p(1-\theta)} \int_{B_\rho^c} |f_0|^p m^p \\ &\leq -ap \langle \rho \rangle^{\gamma+s-2} Y_\theta + ap \langle \rho \rangle^{\gamma+s-2} m(\rho)^{-p(1-\theta)} \int |f_0|^p m^p. \end{aligned}$$

Integrating this differential inequality we deduce

$$\begin{aligned} Y_\theta(t) &\leq \exp(-apt \langle \rho \rangle^{\gamma+s-2}) Y_\theta(0) + m(\rho)^{-p(1-\theta)} Y_\theta(0). \\ &\leq (\exp(-apt \langle \rho \rangle^{\gamma+s-2}) + \exp(-p(1-\theta)\rho^s)) Y_\theta(0). \end{aligned}$$

We may choose ρ such that $a \langle \rho \rangle^{\gamma+s-2} t = (1-\theta)\rho^s$, that is we may take ρ of order $t^{1/(2-\gamma)}$, which allows us to conclude that (2.3) also holds in the exponential case. As indicated above, the estimate (2.3) for $p = \infty$ is obtained by letting $p \rightarrow \infty$. \square

The following two lemmas state that when the weight function is exponential, that is $m(x) := \exp(\kappa \langle x \rangle^\gamma)$, the semigroup $S_{\mathcal{B}}$ is ultracontractive in the spaces $L^p(m)$, that is it maps $L^1(m)$ into $L^\infty(m)$ for $t > 0$. As it is pointed out in [21] (see Remark 2.2 of this reference for a proof based on Probability arguments, and Remark 5.2 for a simple proof based on comparison theorems for parabolic equations), when one considers an operator of the type $Lf := \Delta f + \nabla V \cdot \nabla f$ with V satisfying, for some constants $R > 0$ and $c_0 \geq 0$,

$$\forall x \in B_R^c, \quad \frac{\Delta V^{1/2}}{V^{1/2}} + c_0 \geq 0,$$

and if there exists a positive constant $c_1 > 0$ such that $V(x) \geq c_1$ for all $x \in B_R^c$, then the semigroups $S_L(t)$ and $S_{L^*}(t)$ are ultracontractive in the spaces $L^p(\exp(V))$ (the above condition on $\Delta V^{1/2}/V^{1/2}$ is a sort of convexity condition at infinity). Here the operator \mathcal{B} is not exactly of the same type as L , but nevertheless the ultracontractivity of the semigroup $S_{\mathcal{B}}(t)$, as well as that of $(S_{\mathcal{B}})^*$, holds. (Recall also that the ultracontractivity of the semigroup $S_{\mathcal{B}}$ is equivalent to an appropriate form of Nash inequality for the operator \mathcal{B}).

Lemma 2.2. *Consider the weight function $m_0 := \exp(\kappa\langle x \rangle^\gamma)$, for $0 < \kappa\gamma < 1$. Then there exists $R_0, M_0 > 0$ such that for $M \geq M_0$ and $R \geq R_0$ we have*

$$(2.6) \quad \forall t > 0, \quad \|S_{\mathcal{B}}(t)\|_{L^1(m_0) \rightarrow L^2(m_0)} \lesssim t^{-d/4}.$$

Proof of Lemma 2.2. The proof is similar to the proof of [17, Lemma 3.9], see also [26, Section 3], [27, Section 2] and [20]. For the sake of completeness we sketch it below.

Consider $f_0 \in L^2(m_0)$ and denote $f(t) = S_{\mathcal{B}}(t)f_0$. From (2.4) with $p = 2$ and throwing out the last term in that inequality, we find

$$\frac{d}{dt} \frac{1}{2} \int_{\mathbb{R}^d} f(t)^2 m_0^2 dx \leq - \int_{\mathbb{R}^d} |\nabla(f(t)m_0)|^2 dx.$$

Using Nash's inequality for $g := f(t)m_0$ ([22, Chapter 8]) stating that for some constant $c > 0$

$$(2.7) \quad \int_{\mathbb{R}^d} g^2 dx \leq c \left(\int_{\mathbb{R}^d} |\nabla g|^2 dx \right)^{\frac{d}{d+2}} \left(\int_{\mathbb{R}^d} |g| dx \right)^{\frac{4}{d+2}}$$

we get (for another constant $c > 0$)

$$(2.8) \quad X'(t) \leq -2c Y(t)^{-4/d} X(t)^{1+\frac{2}{d}},$$

where for brevity of notations we have set

$$X(t) := \|f(t)\|_{L^2(m_0)}^2, \quad Y(t) := \|f(t)\|_{L^1(m_0)}.$$

Since according to (2.2) we have $Y(t) \leq Y_0$ for $t > 0$, we may integrate the differential inequality (2.8) and obtain (2.6). \square

The next result states that the adjoint of \mathcal{B} generates also an ultracontractive semigroup in the spaces $L^p(m_0)$.

Lemma 2.3. *Consider the weight function $m_0 := \exp(\kappa\langle x \rangle^\gamma)$, for $0 < \kappa\gamma < 1$. Then there exists $R_1 \geq R_0$ and $M_1 \geq M_0$ (where M_0 and R_0 are defined in the previous lemma) such that for $M \geq M_1$ and $R \geq R_1$, the semigroup generated by \mathcal{B}_* , the formal adjoint of \mathcal{B} , satisfies*

$$(2.9) \quad \forall t > 0, \quad \|S_{\mathcal{B}_*}(t)\|_{L^1(m_0) \rightarrow L^2(m_0)} \lesssim t^{-d/4}.$$

Consequently, for $M \geq M_1$ and $R \geq R_1$, we have

$$(2.10) \quad \forall t > 0, \quad \|S_{\mathcal{B}}(t)\|_{L^2(m_0) \rightarrow L^\infty(m_0)} \lesssim t^{-d/4}.$$

Proof of Lemma 2.3. We first observe that if the operator B is of the form

$$Bf = \Delta f + \mathbf{b}(x) \cdot \nabla f + a(x)f,$$

and we make the transform $h := fm$, then the corresponding operator $B_m h := mB(m^{-1}h)$ is of the same type and is given by

$$B_m h = \Delta h + \left[\mathbf{b}(x) - 2 \frac{\nabla m}{m} \right] \cdot \nabla h + \left[-\frac{\Delta m}{m} + 2 \frac{|\nabla m|^2}{m^2} + a - \mathbf{b}(x) \cdot \frac{\nabla m}{m} \right] h.$$

Observe also that the formal adjoint of B , denoted by B_* to avoid any misunderstanding, is given by

$$B_* g = \Delta g - \mathbf{b}(x) \cdot \nabla g + (a(x) - \operatorname{div}(\mathbf{b}(x)))g.$$

Applying these observations to $\mathcal{B}f = \Delta f + \mathbf{F} \cdot \nabla f + (\operatorname{div}(\mathbf{F}) - M\chi_R)f$, we get that for $h := g m_0$ the operator $\mathcal{B}_{*,m}$, associated to the formal adjoint \mathcal{B}_* , is given by

$$(2.11) \quad \mathcal{B}_{*,m}h = \Delta h - \left[\mathbf{F} - 2 \frac{\nabla m}{m} \right] \cdot \nabla h + \left[\frac{\Delta m}{m} - M\chi_R - \mathbf{F} \cdot \frac{\nabla m}{m} \right] h.$$

Thus, if $g_0 \geq 0$ is a smooth initial datum, then the solution g of

$$\partial_t g = \mathcal{B}_*g, \quad g(0, x) = g_0(x),$$

yields the function $h := g m_0$ which satisfies the evolution equation

$$(2.12) \quad \partial_t h = \mathcal{B}_{*,m_0}h, \quad h(0, x) = h_0(x).$$

Now, one can verify easily that for $h_0 \in C_c^\infty(\mathbb{R}^d)$ and $h_0 \geq 0$, the solution h to the equation

$$\partial_t h = \Delta h + \mathbf{b}(x) \cdot \nabla h + a(x)h, \quad h(0, x) = h_0(x)$$

satisfies $h(t, x) \geq 0$ and for $1 \leq p < \infty$ we have the identity

$$\frac{d}{dt} \frac{1}{p} \int h(t, x)^p dx = -(p-1) \int |\nabla h|^2 h^{p-2} dx + \frac{1}{p} \int (p a(x) - \operatorname{div}(\mathbf{b}(x))) h^p dx.$$

As a consequence, applying this to the operator \mathcal{B}_{*,m_0} , we have that the solution h of equation (2.12) verifies

$$\frac{d}{dt} \frac{1}{p} \int h(t, x)^p dx \leq -(p-1) \int |\nabla h|^2 h^{p-2} dx + \int h^p \psi_{*,m_0,p} dx,$$

where, for convenience, we have set

$$(2.13) \quad \psi_{*,p,m_0} := \frac{(p-2)}{p} \frac{\Delta m_0}{m_0} + \frac{2}{p} \frac{|\nabla m_0|^2}{m_0^2} + \frac{1}{p} \operatorname{div}(\mathbf{F}) - \mathbf{F} \cdot \frac{\nabla m_0}{m_0} - M\chi_R.$$

Proceeding as we did above in the study of the function defined in (2.4), we may choose, if necessary, M and R large enough (in particular larger than M_0, R_0 given by Lemma 2.2), so that for all $x \in \mathbb{R}^d$ we have $\psi_{*,m_0} \leq 0$. Therefore we conclude that

$$(2.14) \quad \frac{d}{dt} \frac{1}{p} \|h(t)\|_{L^p}^p \leq -(p-1) \int |\nabla h|^2 h^{p-2} dx.$$

On the one hand, taking $p := 1$, we deduce that the semigroup generated by \mathcal{B}_{*,m_0} is a contraction semigroup in $L^1(\mathbb{R}^d)$, that is $\|h(t)\|_{L^1} \leq \|h_0\|_{L^1}$ for all $t \geq 0$.

On the other hand, taking $p := 2$ and using Nash's inequality (2.7), together with the fact that $\|h(t)\|_{L^1}$ is non increasing, we deduce that if we set $X(t) := \|h(t)\|_{L^2}^2$, then for some constant $c > 0$, the function $X(t)$ satisfies the differential inequality

$$\frac{d}{dt} X(t) \leq -c \|h_0\|_{L^1}^{-4/d} X(t)^{(2+d)/d}.$$

Integrating this, we get that for all $t > 0$

$$\|h(t)\|_{L^2} \lesssim t^{-d/4} \|h_0\|_{L^1}.$$

From this, by a density argument and the splitting of any initial datum as the difference of two nonnegative functions, we conclude that for any $g_0 \in L^1(m_0)$ the associated solution to $\partial_t g = \mathcal{B}_*g$ satisfies

$$(2.15) \quad \|S_{\mathcal{B}_*}(t)g_0\|_{L^2(m_0)} = \|g(t)\|_{L^2(m_0)} \lesssim t^{-d/4} \|g_0\|_{L^1(m_0)}, \quad \forall t > 0,$$

which is precisely (2.9). To conclude the proof of the Lemma, observe that for $f, g \in C_c^\infty(\mathbb{R}^d)$ we have

$$\begin{aligned} (\mathcal{B}f|g)_{L^2(m_0)} &= \int \mathcal{B}f g m_0^2 dx = \int \mathcal{B}_{m_0}(f m_0)(g m_0) dx \\ &= \int (f m_0) \mathcal{B}_{*,m_0}(g m_0) dx = \int (f m_0) m_0^{-1} \mathcal{B}_{*,m_0}(g m_0) m_0(x)^2 dx \\ &= (f|\mathcal{B}_*g)_{L^2(m_0)}. \end{aligned}$$

This allows one to verify that $(S_{\mathcal{B}}(t))^* = S_{\mathcal{B}_*}(t)$, the adjoint being taken in the sense of the Hilbert space $L^2(m_0)$, where we assume that this space is identified with its dual. Therefore, since with these conventions we have $(L^1(m_0))' = L^\infty(m_0)$, thanks to (2.15), we conclude (2.10). \square

Putting together the previous estimates, we get the following ultracontractivity result on the semigroup $S_{\mathcal{B}}$ and on the iterated convolution family of operators $(\mathcal{A}S_{\mathcal{B}})^{(*n)}$.

Lemma 2.4. *Consider the weight function $m_0 := \exp(\kappa\langle x \rangle^\gamma)$, for $0 < \kappa\gamma < 1$. Then, M, R being large enough as in Lemma 2.3, there exists $\lambda_* \in (0, \infty)$ such that for any $p, q \in [1, \infty]$, $p \leq q$, and for any $0 \leq \theta_2 < \theta_1 \leq 1$, the semigroup $S_{\mathcal{B}}$ satisfies*

$$(2.16) \quad \|S_{\mathcal{B}}(t)\|_{L^p(m_0^{\theta_2}) \rightarrow L^q(m_0^{\theta_1})} \lesssim t^{-(d/2)(1/p-1/q)} e^{-\lambda_* t^{\gamma/(2-\gamma)}}, \quad \forall t > 0.$$

Moreover, if $n \geq d/2$ is an integer, for all $\lambda < \lambda_*$ and all $t > 0$, we have

$$(2.17) \quad \|(\mathcal{A}S_{\mathcal{B}})^{(*n)}(t)\|_{L^1(m_0) \rightarrow L^\infty(m_0^2)} \lesssim e^{-\lambda t^{\gamma/(2-\gamma)}}.$$

Proof of Lemma 2.4. Step 1. Writing $S_{\mathcal{B}}(t) = S_{\mathcal{B}}(t/2)S_{\mathcal{B}}(t/2)$, and using (2.6) together with (2.10), we deduce that for any $t > 0$

$$(2.18) \quad \begin{aligned} \|S_{\mathcal{B}}(t)\|_{L^1(m_0) \rightarrow L^\infty(m_0)} &\leq \|S_{\mathcal{B}}(t/2)\|_{L^2(m_0) \rightarrow L^\infty(m_0)} \|S_{\mathcal{B}}(t/2)\|_{L^1(m_0) \rightarrow L^2(m_0)} \\ &\lesssim t^{-d/2}. \end{aligned}$$

Since on the other hand we have also $\|S_{\mathcal{B}}(t)f_0\|_{L^p(m_0)} \leq \|f_0\|_{L^p(m_0)}$, a classical interpolation argument yields that for $1 \leq p \leq q \leq \infty$, we have

$$(2.19) \quad \|S_{\mathcal{B}}(t)\|_{L^p(m_0) \rightarrow L^q(m_0)} \lesssim t^{-(d/2)(1/p-1/q)}.$$

Using Lemma 2.1, since

$$\|S_{\mathcal{B}}(t)\|_{L^q(m_0) \rightarrow L^q(m_0^2)} \lesssim \exp(-\lambda t^{\gamma/(2-\gamma)}),$$

one sees that the proof of (2.16) with $\theta_1 = 1$ is complete. To see that (2.16) holds with $\theta_1 < 1$, it is enough to observe that $m_0^{\theta_1}(x)$ is of the same type as $m_0(x) = \exp(\kappa\langle x \rangle^\gamma)$ provided κ is replaced with $\theta_1\kappa$.

Step 2. In order to show (2.17), first note that the operator \mathcal{A} consisting simply in a multiplication by a smooth compactly supported function, thanks to the above lemmas, we clearly have, for all $t > 0$,

$$(2.20) \quad \|\mathcal{A}S_{\mathcal{B}}(t)\|_{L^{p_1}(m_0) \rightarrow L^{p_2}(m_0^2)} \lesssim t^{-\alpha} e^{-\lambda t^{\sigma^*}},$$

with $\sigma^* := \gamma/(2-\gamma) \leq 1$ and where it is understood that $\alpha := d/2$ if $(p_1, p_2) := (1, \infty)$, and $\alpha := 0$ when $p_1 = p_2$. We claim that the three estimates for three

choices $(p_1, p_2) = (1, \infty)$, and $p_1 = p_2 = 1$, as well as $p_1 = p_2 = \infty$, imply that for all integers $n \geq 1$ we have

$$(2.21) \quad \|(\mathcal{A}S_{\mathcal{B}})^{*(n)}(t)\|_{L^{p_1}(m_0) \rightarrow L^{p_2}(m_0^2)} \leq C_n t^{n-1-\alpha} e^{-\lambda t^{\sigma^*}},$$

from which one readily deduces (2.17). We prove (2.21) by induction. Estimates (2.21) are clearly true for $n = 1$. Let us assume that $p_1 = 1$ and $p_2 = \infty$, for which (2.21) is true for a certain $n \geq 1$. Introducing the shorthand notation $u := \mathcal{A}S_{\mathcal{B}}$ and $\|\cdot\|_{p_1 \rightarrow p_2} = \|\cdot\|_{L^{p_1}(m_0) \rightarrow L^{p_2}(m_0^2)}$, we have

$$\begin{aligned} \|u^{*(n+1)}(t)\|_{1 \rightarrow \infty} &\leq \int_0^{t/2} \|u^{(n)}(t-s)\|_{1 \rightarrow \infty} \|u(s)\|_{1 \rightarrow 1} ds \\ &\quad + \int_{t/2}^t \|u^{(n)}(t-s)\|_{\infty \rightarrow \infty} \|u(s)\|_{1 \rightarrow \infty} ds \\ &\leq C_n C_1 e^{-\lambda t^{\sigma^*}} \int_0^{t/2} (t-s)^{-\alpha+n-1} ds \\ &\quad + C_n C_1 e^{-\lambda t^{\sigma^*}} \int_{t/2}^t (t-s)^{n-1} s^{-\alpha} ds \\ &\leq C_n C_1 e^{-\lambda t^{\sigma^*}} t^{-\alpha+n} \left\{ \int_0^{1/2} (1-\tau)^{-\alpha+n-1} d\tau \right. \\ &\quad \left. + \int_{1/2}^1 (1-\tau)^{n-1} \tau^{-\alpha} d\tau \right\}, \end{aligned}$$

where we have used the fact that $t^{\sigma^*} \leq (t-s)^{\sigma^*} + s^{\sigma^*}$ for any $0 \leq s \leq t$, since $0 < \sigma^* \leq 1$. This proves estimate (2.21) at rank $n+1$ and $(p_1, p_2) = (1, \infty)$.

The proof of the other cases $(p_1, p_2) = (1, 1)$ and $(p_1, p_2) = (\infty, \infty)$ is similar, if not much simpler, and can be left to the reader. \square

2.2. Additional growth estimates. In order to deal with the general case in Section 6, we will need a more accurate version of the previous estimates.

Lemma 2.5. *Consider $m_0 := e^{2\kappa\langle x \rangle^\gamma}$ with $\kappa \in (0, 1/(4\gamma))$ and define the sequence of spaces*

$$(2.22) \quad X_k := L^2(m_k), \quad m_k := \frac{m_0}{\nu_k}, \quad \nu_k(x) := \sum_{\ell=0}^k \frac{(\kappa \langle x \rangle^\gamma)^\ell}{\ell!},$$

for any $k \in \mathbb{N}$. There exist some constants R and M in the definition of \mathcal{B} and some constant $\beta > 0$ such that for any $k, j \in \mathbb{N}$, $k \geq j$ and any $\alpha \in (0, \alpha^*)$, $\alpha^* := 1/2(1-\gamma)$, the semigroup $S_{\mathcal{B}}$ satisfies the growth estimate

$$(2.23) \quad \|S_{\mathcal{B}}(t)\|_{X_{k-j} \rightarrow X_k} \lesssim e^{-\lambda \langle t^\alpha \rangle^{2(\gamma-1)} t} + \left(1 \wedge \frac{k^j}{\kappa^j \langle t^\alpha \rangle^{\gamma j}}\right) \quad \forall t > 0.$$

Proof of Lemma 2.5. We easily compute

$$\frac{\nabla m_k}{m_k} = \gamma \kappa x \langle x \rangle^{\gamma-2} \left[2 - \frac{1 + \dots + (\kappa \langle x \rangle^\gamma)^{k-1} / (k-1)!}{1 + \dots + (\kappa \langle x \rangle^\gamma)^k / k!} \right],$$

from which we deduce

$$\left| \frac{\nabla m_k}{m_k} \right| \leq 2\gamma \kappa |x| \langle x \rangle^{\gamma-2},$$

and

$$-\mathbf{F} \cdot \frac{\nabla m_k}{m_k} \leq -\gamma \kappa |x|^\gamma \langle x \rangle^{\gamma-2} \quad \forall x \in B_{R_0}^c.$$

As a consequence, for $f(t) := S_{\mathcal{B}}(t)f_0$, we have

$$\frac{d}{dt} \int f^2 m_k^2 \leq - \int |\nabla f|^2 m_k^2 + \int f^2 m_k^2 \psi_k,$$

with

$$\begin{aligned} \psi_k &= \frac{|\nabla m_k|^2}{m_k^2} + \frac{1}{2} \operatorname{div}(\mathbf{F}) - \mathbf{F} \cdot \frac{\nabla m_k}{m_k} - M \chi_R \\ &\leq 4\gamma^2 \kappa^2 |x|^2 \langle x \rangle^{2\gamma-4} + \frac{1}{2} C'_F \langle x \rangle^{\gamma-2} \\ &\quad - \gamma \kappa |x|^\gamma \langle x \rangle^{\gamma-2} \mathbf{1}_{B_{R_0}^c} - M \chi_R \\ &\leq -2\lambda \langle x \rangle^{2(\gamma-1)} \end{aligned}$$

for any $x \in \mathbb{R}^d$ and $k \in \mathbb{N}$, by fixing $\kappa > 0$ small enough (as we did) and then R and M large enough. We deduce

$$(2.24) \quad \frac{d}{dt} \int f^2 m_k^2 \leq - \int |\nabla f|^2 m_k^2 - 2\lambda \int f^2 m_k^2 \langle x \rangle^{2(\gamma-1)},$$

and in particular

$$Y_k(t) := \int f^2 m_k^2 \leq Y_k(0) \quad \text{for any } k \geq 0.$$

We now observe that for any $j \in \mathbb{N}$, $0 \leq j \leq k$, there hold $m_k \leq m_{k-j}$ as well as

$$m_k(x) \leq \frac{m_0(x)}{\nu_{k-j}(x) (\kappa \langle x \rangle^\gamma / k)^j} = \frac{k^j}{\kappa^j \langle x \rangle^{\gamma j}} m_{k-j}(x) \quad \forall x \in \mathbb{R}^d.$$

The two inequalities together, we have proved

$$(2.25) \quad \forall j \leq k, \forall x \in \mathbb{R}^d \quad m_k(x) \leq \left(1 \wedge \frac{k^j}{\kappa^j \langle x \rangle^{\gamma j}}\right) m_{k-j}(x).$$

As a consequence, for any $\rho > 0$, we have

$$\begin{aligned} Y_k &= \int_{B_\rho} f^2 m_k^2 + \int_{B_\rho^c} f^2 m_k^2 \\ &\leq \langle \rho \rangle^{2(1-\gamma)} \int f^2 m_k^2 \langle x \rangle^{2(\gamma-1)} + \left(1 \wedge \frac{k^j}{\kappa^j \langle \rho \rangle^{\gamma j}}\right)^2 Y_{k-j}. \end{aligned}$$

Coming back to (2.24), we deduce

$$\frac{d}{dt} Y_k \leq -2\lambda \langle \rho \rangle^{2(\gamma-1)} Y_k + 2\lambda \langle \rho \rangle^{2(\gamma-1)} \left(1 \wedge \frac{k^j}{\kappa^j \langle \rho \rangle^{\gamma j}}\right)^2 Y_{k-j}(0),$$

which in turn implies

$$Y_k(t) \leq \left\{ e^{-2\lambda \langle \rho \rangle^{2(\gamma-1)} t} + \left(1 \wedge \frac{k^j}{\kappa^j \langle \rho \rangle^{\gamma j}}\right)^2 \right\} Y_{k-j}(0) \quad \forall t \geq 0, \rho > 0.$$

We conclude by making the choice $\rho = t^\alpha$ for $\alpha \in (0, \alpha^*)$. \square

Lemma 2.6. *Consider $m_0 := e^{\kappa \langle x \rangle^\gamma}$ with $\kappa \in (0, 1/(2\gamma))$. There exist two constants $C, \lambda > 0$ such that $S_{\mathcal{B}}$ satisfies (recall that $\sigma^* = \gamma/(2 - \gamma)$)*

$$(2.26) \quad \|S_{\mathcal{B}}(t)\|_{L^2(m_0) \rightarrow H^1} \leq \frac{C}{t^{1/2}} e^{-\lambda t^{\sigma^*}} \quad \forall t > 0.$$

Proof of Lemma 2.6. The function $f = S_{\mathcal{B}}(t)f_0$ satisfies

$$\partial_i \partial_i f = \Delta \partial_i f + (\partial_i \operatorname{div}(\mathbf{F}))f + \operatorname{div}(\mathbf{F}) \partial_i f + \partial_i \mathbf{F}_j \partial_j f + \mathbf{F} \cdot \nabla \partial_i f - M \partial_i \chi_R f - M \chi_R \partial_i f,$$

so that

$$\begin{aligned} \frac{1}{2} \frac{d}{dt} \int |\nabla f|^2 m_0^2 &= - \int |D^2 f|^2 m_0^2 + \frac{1}{2} \int |\nabla f|^2 \Delta m_0^2 - \int (\operatorname{div} \mathbf{F}) f (\Delta f) m_0^2 \\ &\quad - \int (\operatorname{div}(\mathbf{F})) f \nabla f \cdot \nabla m_0^2 - \int \partial_i \mathbf{F}_j \partial_i f \partial_j f m_0^2 - \frac{1}{2} \int (\operatorname{div}(\mathbf{F})) |\nabla f|^2 m_0^2 \\ &\quad - \frac{1}{2} \int |\nabla f|^2 \mathbf{F} \cdot \nabla m_0^2 + M \int \operatorname{div}(\nabla \chi_R m_0^2) f^2 - \int M \chi_R |\nabla f|^2 \\ &\leq -\frac{1}{2} \int |D^2 f|^2 m_0^2 dx + \int |\nabla f|^2 m_0^2 \psi_1 dx + \int f^2 m_0^2 \psi_2 dx, \end{aligned}$$

with

$$\psi_1(x) := \frac{1}{2} \frac{\Delta m_0^2}{m_0^2} + \frac{1}{2} |\operatorname{div}(\mathbf{F})| \frac{|\nabla m_0^2|}{m_0^2} + \frac{3}{2} |D\mathbf{F}| - \mathbf{F} \cdot \frac{\nabla m_0}{m_0} - M \chi_R$$

and

$$\psi_2(x) := (\operatorname{div}(\mathbf{F}))^2 + \frac{1}{2} \frac{\Delta m_0^2}{m_0^2} + \frac{1}{2} |\operatorname{div} \mathbf{F}| \frac{|\nabla m_0^2|}{m_0^2} + M |\Delta \chi_R| + M |\nabla \chi_R| \frac{|\nabla m_0^2|}{m_0^2}.$$

Choosing M and R large enough, we have for some constants $a > 0$ and $C \in \mathbb{R}$

$$\psi_1(x) \leq -a \langle x \rangle^{2(\gamma-1)}, \quad \psi_2(x) \leq C \langle x \rangle^{2(\gamma-1)},$$

and, choosing then $\eta > 0$ small enough, we get

$$\begin{aligned} \frac{d}{dt} \int (f^2 + \eta |\nabla f|^2) m_0^2 &\leq -\frac{1}{2} \int (|\nabla f|^2 + \eta |D^2 f|^2) m_0^2 \\ &\quad - a \int (f^2 + \eta |\nabla f|^2) m_0^2 \langle x \rangle^{2(\gamma-1)}. \end{aligned}$$

On the one hand, keeping only the second term and arguing as in the proof of Lemma 2.1, we get

$$(2.27) \quad \|S_{\mathcal{B}}(t)\|_{\mathcal{B}(H^1(m_0), H^1)} \leq \Theta_m(t), \quad \forall t \geq 0.$$

On the other hand, using the elementary inequality

$$\int |\nabla f|^2 m_0^2 = - \int f D^2 f m_0^2 + \frac{1}{2} \int f^2 \Delta m_0^2 \leq \|f\|_{L^2(m_0)} \|f\|_{H^2(m_0)} + C \|f\|_{L^2(m_0)}^2,$$

the differential inequality

$$\frac{d}{dt} \|\nabla f\|_{L^2(m_0)}^2 \leq -\|D^2 f\|_{L^2(m_0)}^2 + \|\psi^2\|_{L^\infty} \|f\|_{L^2(m_0)}^2$$

and the contraction in $L^2(m_0)$, we then obtain

$$\frac{d}{dt} \|\nabla f\|_{L^2(m_0)}^2 \leq -\|f_0\|_{L^2(m_0)}^{-2} \|\nabla f\|_{L^2(m_0)}^4 + \|\psi^2\|_{L^\infty} \|f_0\|_{L^2(m_0)}^2.$$

As in the proof of Lemma 2.1, we deduce

$$(2.28) \quad \|\nabla f\|_{L^2(m_0)}^2 \leq \frac{C}{t} \|f_0\|_{L^2(m_0)}^2 \quad \forall t \in (0, 1).$$

We easily deduce (2.26) by gathering (2.27) and (2.28). \square

3. BOUNDEDNESS OF $S_{\mathcal{L}}$

In this section, we establish some estimates on $S_{\mathcal{L}}$ which are simple results yielded by the iterated Duhamel formula (1.28) and the previous estimates on $S_{\mathcal{B}}$.

Lemma 3.1. *For any exponent $p \in [1, \infty]$ and any weight function m given by (1.9), (1.10) or (1.11), there exists $C(m, p)$ such that*

$$\|S_{\mathcal{L}}(t)\|_{\mathcal{B}(L^p(m))} \leq C(m, p), \quad \forall t \geq 0.$$

Proof of Lemma 3.1. The estimate

$$(3.1) \quad \|S_{\mathcal{L}}(t)\|_{\mathcal{B}(L^1)} \leq 1, \quad \forall t \geq 0,$$

is clear and is a consequence of the fact that $S_{\mathcal{L}}$ is a positive and mass conserving semigroup.

Step 1. We consider first the case $p = 1$ and we write the Duhamel formula

$$S_{\mathcal{L}} = S_{\mathcal{B}} + S_{\mathcal{B}} * \mathcal{A}S_{\mathcal{L}},$$

which allows one to deduce that

$$\|S_{\mathcal{L}}\|_{L_k^1 \rightarrow L_k^1} \leq \|S_{\mathcal{B}}\|_{L_k^1 \rightarrow L_k^1} + \|S_{\mathcal{B}}\|_{L_{k+\ell}^1 \rightarrow L_k^1} * \|\mathcal{A}S_{\mathcal{L}}\|_{L^1 \rightarrow L_{k+\ell}^1} \leq C,$$

because $t \mapsto \|\mathcal{A}S_{\mathcal{L}}(t)\|_{L^1 \rightarrow L_{k+\ell}^1}$ is a bounded function thanks to (3.1), and, upon choosing $\ell := 2(2 - \gamma)$ and applying Lemma 2.1, we have $\|S_{\mathcal{B}}\|_{L_{k+\ell}^1 \rightarrow L_k^1} \lesssim \langle t \rangle^{-2}$, showing that $\|S_{\mathcal{B}}(t)\|_{L_{k+\ell}^1 \rightarrow L_k^1}$ belongs to $L^1(0, \infty)$.

Step 2. In order to study the case $1 < p \leq \infty$, we write the iterated Duhamel formula (1.28) with $n_1 := \ell + 1$ and $n_2 := 0$ and we get

$$S_{\mathcal{L}} = S_{\mathcal{B}} + \dots + S_{\mathcal{B}} * (\mathcal{A}S_{\mathcal{B}})^{*(\ell-1)} + S_{\mathcal{B}} * (\mathcal{A}S_{\mathcal{B}})^{*(\ell)} * (\mathcal{A}S_{\mathcal{L}})$$

with $\ell := d/2 + 1$ and then we infer that

$$\begin{aligned} \|S_{\mathcal{L}}\|_{\mathcal{B}(L^p(m))} &\leq \sum_{j=0}^{\ell-1} \|S_{\mathcal{B}}\|_{\mathcal{B}(L^p(m))} * \|\mathcal{A}S_{\mathcal{B}}\|_{\mathcal{B}(L^p(m))}^{*(j)} \\ &\quad + \|S_{\mathcal{B}}\|_{L^p(\varpi) \rightarrow L^p(m)} * \|(\mathcal{A}S_{\mathcal{B}})^{*(\ell)}\|_{L^1(\varpi) \rightarrow L^\infty(\varpi^2)} * \|\mathcal{A}S_{\mathcal{L}}\|_{L^p(m) \rightarrow L^1(\varpi)}, \end{aligned}$$

where $\varpi := e^{\nu \langle x \rangle^\gamma}$ with $\nu \in (0, 1/\gamma)$ large enough. Using (2.20) and (2.17), we see that the RHS term is uniformly bounded as the sum of $\ell + 1$ functions, each one being the convolution of one L^∞ function with integrable functions. \square

Remark 3.2. When $p = 1$, $m = \langle x \rangle^k$ and $f(t) \geq 0$, the above estimate means exactly that for $k > 2 - \gamma > 1$, there exists a constant C_k such that

$$M_k(S_{\mathcal{L}}(t)f_0) \leq C_k M_k(f_0), \quad M_k(f) := \int_{\mathbb{R}^d} f \langle x \rangle^k dx.$$

We then recover (with a simpler proof) and generalize (to a wider class of confinement force fields) a result obtained by Toscani and Villani in [33, section 2].

4. THE **Case 1** WHEN A POINCARÉ INEQUALITY HOLDS

In this section we restrict our analysis to **Case 1**, that is when furthermore the steady state G to the Fokker-Planck equation (1.1) is such that the *weak weighted Poincaré-Wirtinger inequality* (1.14) holds true and the function $V := -\log G$ is of order $\langle x \rangle^\gamma$, that is it satisfies (1.13). We start by showing the probably classical fact that (1.14) holds in the case when $c_1 = c_2$ in (1.13). The proof uses some of the arguments exposed in [3, 11].

Lemma 4.1. *Consider a potential function V such that (1.13) holds with $c_1 = c_2$ and denote $G := e^{-V}$ (that we assume to have mass 1). There exists $\mu \in (0, \infty)$ such that the following weak weighted Poincaré inequality*

$$\int |\nabla h|^2 G dx \geq \mu \int h^2 \langle x \rangle^{2(\gamma-1)} G dx$$

holds for any $h \in C_b^1(\mathbb{R}^d)$ such that $\mathbb{M}(hG) = 0$.

Proof of Lemma 4.1. We observe that $|\nabla V(x)|^2 \sim c_1^2 \gamma^2 |x|^{2\gamma-2}$ and $|\Delta V(x)| \sim c_1 \gamma |\gamma + d - 2| |x|^{\gamma-2}$, so that

$$(4.1) \quad \forall \theta \in (0, 1), \exists R_1 > 0, \quad \theta |\nabla V|^2 \geq \Delta V \text{ on } B_{R_1}^c.$$

We fix $h \in \mathcal{D}(\mathbb{R}^d)$ such that $\mathbb{M}(hG) = 0$. On the one hand, defining $g = h e^{-aV}$ and using the identity

$$\int |\nabla h|^2 e^{-V} = \int |\nabla g|^2 e^{(2a-1)V} + a \int h^2 [(1-a)|\nabla V|^2 - \Delta V] e^{-V}$$

with $a := (1 - \theta)/2$ together with (4.1), we obtain a first estimate

$$(4.2) \quad \frac{(1-\theta)^2}{4} \int h^2 |\nabla V|^2 e^{-V} \leq \int |\nabla h|^2 e^{-V} + b \int_{B_{R_0}} h^2 e^{-V}.$$

On the other hand, using the Poincaré-Wirtinger inequality in $H^1(B_R, Gdx)$, we get

$$(4.3) \quad \int_{B_R} h^2 \frac{G}{\mathbb{M}(G\mathbf{1}_{B_R})} \leq C_R \int_{B_R} |\nabla h|^2 \frac{G}{\mathbb{M}(G\mathbf{1}_{B_R})} + \left(\int_{B_R} h \frac{G}{\mathbb{M}(G\mathbf{1}_{B_R})} \right)^2,$$

with possibly $C_R \rightarrow \infty$ when $R \rightarrow \infty$. Using the vanishing moment condition $\mathbb{M}(hG) = 0$ and the Cauchy-Schwarz inequality, we also have

$$(4.4) \quad \left(\int_{B_R} h G \right)^2 = \left(\int_{B_R^c} h G \right)^2 \leq \varepsilon_R \int_{B_R^c} h^2 |\nabla V|^2 G,$$

with

$$\varepsilon_R := \int_{B_R^c} |\nabla V|^{-2} e^{-V} \rightarrow 0 \quad \text{as } R \rightarrow \infty.$$

Gathering (4.3) and (4.4) for $R \geq R_2$, we obtain the second estimate

$$(4.5) \quad \begin{aligned} \int_{B_R} h^2 e^{-V} &\leq C_R \int_{B_R} |\nabla h|^2 e^{-V} dx + 2 \left(\int_{B_R} h e^{-V} dx \right)^2 \\ &\leq C_R \int_{B_R} |\nabla h|^2 e^{-V} dx + 2\varepsilon_R \int_{B_R^c} h^2 |\nabla V|^2 e^{-V}. \end{aligned}$$

Inserting (4.5) into (4.2) with $R \geq \max(R_1, R_2)$, we obtain

$$\frac{(1-\theta)^2}{4} \int h^2 |\nabla V|^2 e^{-V} \leq (1+C_R) \int |\nabla h|^2 e^{-V} + 2\varepsilon_R \int_{B_R^c} h^2 |\nabla V|^2 e^{-V},$$

and we conclude by taking R large enough in such a way that $2\varepsilon_R < (1-\theta)^2/4$. \square

4.1. Decay estimate in a small space.

We first recall the following decay estimate in a small space. We define

$$E_1 := L^2(G^{-1/2}), \quad E_2 := L^\infty(G^{-1}).$$

Theorem 4.2 (Röckner-Wang [31]). *Under assumptions (1.14)–(1.13) on the steady state G , set*

$$(4.6) \quad \Theta_{G^{-1}}(t) := \exp(-\lambda t^{\gamma/(2-\gamma)}).$$

Then there exists λ^ such that for $\lambda < \lambda_*$ the semigroup $S_{\mathcal{L}}$ satisfies for all $t > 0$*

$$(4.7) \quad \|S_{\mathcal{L}}(t)f_0 - \mathbb{M}(f_0)G\|_{E_1} \lesssim \Theta_{G^{-1}}(t) \|f_0 - \mathbb{M}(f_0)G\|_{E_2}.$$

We briefly sketch a proof of Theorem 4.2 which uses as a cornerstone the weak Poincaré inequality (1.14). We refer to [31] for a more detailed proof which uses closely related (but different) arguments and more precisely which are based on the so-called *weak Poincaré inequality* (see [31, 11, 4]).

Proof of Theorem 4.2. We define $V := -\log G$, we set $\mathbf{F}_0 := \mathbf{F} - \nabla V$ and we observe that $\operatorname{div}(G\mathbf{F}_0) = 0$. The Fokker-Planck equation (1.1) can be rewritten as

$$(4.8) \quad \partial_t f = \operatorname{div}[G\nabla(fG^{-1}) + f\mathbf{F}_0].$$

Now, since $f\mathbf{F}_0 = fG^{-1} \cdot G\mathbf{F}_0$ and $\operatorname{div}(G\mathbf{F}_0) = 0$, we have

$$\operatorname{div}(f\mathbf{F}_0) = \nabla(fG^{-1}) \cdot G\mathbf{F}_0.$$

Thus if $j : \mathbb{R} \rightarrow \mathbb{R}$ is a locally Lipschitz function such that $j(fG^{-1}) \in L^1$, we see that $j'(fG^{-1})\operatorname{div}(f\mathbf{F}_0) = \nabla j(fG^{-1}) \cdot G\mathbf{F}_0$ and therefore

$$(4.9) \quad \int j'(fG^{-1})\operatorname{div}(f\mathbf{F}_0) = \int \nabla j(fG^{-1}) \cdot G\mathbf{F}_0 = - \int j(fG^{-1})\operatorname{div}(G\mathbf{F}_0) = 0.$$

In particular, taking $j(s) := s^2/2$ we see that

$$\int fG^{-1}\operatorname{div}(f\mathbf{F}_0) = 0,$$

and as a consequence, multiplying (4.8) by fG^{-1} we obtain

$$(4.10) \quad \begin{aligned} \frac{1}{2} \frac{d}{dt} \int f^2 G^{-1} &= - \int G |\nabla(fG^{-1})|^2 + \int \operatorname{div}[f\mathbf{F}_0] fG^{-1} \\ &\leq -\mu \int f^2 \langle x \rangle^{2\gamma-2} G^{-1}, \end{aligned}$$

thanks to the weak Poincaré inequality (1.14).

More generally, for any convex function $j : \mathbb{R} \rightarrow \mathbb{R}$, locally in $W^{2,\infty}$, using (4.9) one sees easily that the following generalized relative entropy inequality holds (see [24], and also [31] and the references therein)

$$\begin{aligned} \frac{d}{dt} \int j(fG^{-1})G &= \int j'(fG^{-1})\operatorname{div}[G\nabla(fG^{-1})] + \int j'(fG^{-1})\operatorname{div}[f\mathbf{F}_0] \\ &= - \int j''(fG^{-1})|\nabla(fG^{-1})|^2 \leq 0. \end{aligned}$$

In particular, taking $j(s) = |s|^p$, we obtain

$$\|f(t)G^{-1}\|_{L^p} \leq \|f_0G^{-1}\|_{L^p},$$

and passing to the limit $p \rightarrow \infty$, we get

$$(4.11) \quad \|f(t)G^{-1}\|_{L^\infty} \leq \|f_0G^{-1}\|_{L^\infty}, \quad \forall t \geq 0.$$

We obtain (4.7) by gathering the two estimates (4.10) and (4.11). More precisely, we write, with $k := 2 - 2\gamma$,

$$\begin{aligned} \int f^2 G^{-1} &\leq \int_{B_R} f^2 G^{-1} + \int_{B_R^c} f^2 G^{-1} \\ &\leq R^k \int_{B_R} f^2 G^{-1} \langle x \rangle^{-k} + \|fG^{-1}\|_{L^\infty}^2 \int_{B_R^c} G. \end{aligned}$$

Together with (4.10) and (4.11), we get for $R \geq R_0$

$$\frac{d}{dt} \int f^2 G^{-1} \leq -R^{-k} \int f^2 G^{-1} + \|f_0G^{-1}\|_{L^\infty}^2 \int_{B_R^c} e^{-c_1|x|^\gamma}.$$

Integrating in time, for any $\theta \in (0, 1)$, we find a constant $C_\theta > 0$ such that

$$\begin{aligned} \int f^2 G^{-1} &\leq e^{-tR^{-k}} \int f_0^2 G^{-1} + C_\theta e^{-\theta c_1 R^\gamma} \|f_0G^{-1}\|_{L^\infty}^2 \\ &\leq e^{-\lambda t^{\gamma/(2-\gamma)}} \|f_0G^{-1}\|_{L^\infty}^2, \end{aligned}$$

by choosing $R := (t/(\theta c_1))^{1/(2-\gamma)}$ for t large enough and defining $\lambda := (\theta c_1)^{-1/(2-\gamma)}$, which is nothing but (4.7). \square

4.2. Rate of decay in a large space. We now present the proof of our main Theorem 1.2 in the **Case 1** when furthermore G satisfies the weak Poincaré inequality (1.14) and the asymptotic estimates (1.13). We recall that the projection operator Π is defined by

$$\Pi f := \mathbb{M}(f) G.$$

Step 1. Here we consider the case $p \in [1, 2]$. We begin by fixing m a polynomial or exponential weight function, and we introduce a stronger confinement exponential weight $m_0(x) := \exp(\kappa_0 \langle x \rangle^\gamma)$, with $\kappa_0 \in (0, 1/\gamma)$. We denote by $\Theta_m(t)$, $\Theta_{m_0}(t)$ and $\Theta_{G^{-1}}(t)$ the associated rate of decay defined in (1.18), (1.19) and in the statement of Theorem 4.2, and we may assume that $\Theta_{m_0}/\Theta_m, \Theta_{G^{-1}}/\Theta_m \in L^1(0, \infty)$. We split the semigroup on invariant spaces

$$S_{\mathcal{L}} = \Pi S_{\mathcal{L}} + (I - \Pi) S_{\mathcal{L}},$$

and together with the iterated Duhamel formula (1.28) with $n_1 = 0$ and $n_2 := n \geq d/2 + 1$, we can write

$$S_{\mathcal{L}} - \Pi = S_{\mathcal{L}}(I - \Pi) = (I - \Pi)S_{\mathcal{L}} = T_1 + T_2,$$

where

$$(4.12) \quad \begin{aligned} T_1 &:= (I - \Pi) \left(S_{\mathcal{B}} + \sum_{\ell=1}^{n-1} (S_{\mathcal{B}} \mathcal{A})^{(*\ell)} * S_{\mathcal{B}} \right) \\ T_2 &:= ((I - \Pi)S_{\mathcal{L}}) * (\mathcal{A} S_{\mathcal{B}})^{(*n)}. \end{aligned}$$

In order to estimate the first term T_1 , we recall that

$$\|S_{\mathcal{B}}(t)\|_{L^p(m) \rightarrow L^p} \leq \Theta_m(t), \quad \|S_{\mathcal{B}}(t)\mathcal{A}\|_{L^p \rightarrow L^p} \leq \Theta_{m_0}(t),$$

thanks to Lemma 2.1. We write

$$\|T_1\|_{\mathcal{B}(L^p(m), L^p)} \leq u_0 + \dots + u_{n-1}$$

where, for $\ell \in \{1, \dots, n-1\}$, we set

$$u_\ell := C \|S_{\mathcal{B}}\mathcal{A}\|_{\mathcal{B}(L^p)}^{(*\ell)} * \|S_{\mathcal{B}}\|_{\mathcal{B}(L^p(m), L^p)}.$$

Since clearly $\Theta_m^{-1}(t) \leq \Theta_m^{-1}(s) \Theta_m^{-1}(t-s)$ for all $0 \leq s \leq t$, we deduce

$$\|\Theta_m^{-1}u_\ell\|_{L_t^\infty} \leq \|\Theta_m^{-1}S_{\mathcal{B}}\mathcal{A}\|_{L_t^1(\mathcal{B}(L^p))}^\ell \|\Theta_m^{-1}S_{\mathcal{B}}\|_{L_t^\infty(\mathcal{B}(L^p(m), L^p))},$$

and then

$$\forall t \geq 0, \quad \|T_1(t)\|_{\mathcal{B}(L^p(m), L^p)} \leq C \Theta_m(t).$$

In order to estimate the second term T_2 , we recall that from Lemma 2.1, Lemma 2.4 and Theorem 4.2, we have

$$\begin{aligned} \|\mathcal{A}S_{\mathcal{B}}(t)\|_{\mathcal{B}(L^p(m), L^1(m_0))} &\leq \Theta_m(t), \\ \|(\mathcal{A}S_{\mathcal{B}})^{*(n-1)}(t)\|_{\mathcal{B}(L^1(m_0), L^\infty(m_1))} &\leq \Theta_{m_0}(t), \\ \|S_{\mathcal{L}}(t)(I - \Pi)\|_{\mathcal{B}(L^\infty(G^{-1}), L^2(G^{-1/2}))} &\leq \Theta_{G^{-1}}(t), \end{aligned}$$

where $m_1 := m_0 + G^{-1}$. We thus obtain that T_2 satisfies the same estimate as T_1 . We deduce (1.20) by gathering the two estimates obtained on T_1 and T_2 .

Step 2. We consider next the case $p \in [2, \infty]$. Thanks to the iterated Duhamel formula (1.28) with $n_1 := n \geq d/2 + 1$ and $n_2 = 1$, and using the fact that

$$(S_{\mathcal{B}}\mathcal{A})^{(*n)} = S_{\mathcal{B}} * (S_{\mathcal{B}}\mathcal{A})^{*(n-1)} * \mathcal{A},$$

we can write

$$S_{\mathcal{L}} - \Pi = T_1 + T_3$$

where T_1 is defined again by (4.12) and

$$T_3 := (I - \Pi)S_{\mathcal{B}} * (\mathcal{A}S_{\mathcal{B}})^{*(n-1)} * \mathcal{A}S_{\mathcal{L}} * (\mathcal{A}S_{\mathcal{B}})$$

and we conclude in a similar way as in Step 1. \square

5. THE GENERAL CASE - THE STATIONARY PROBLEM

In this section, we present the proof of the existence and uniqueness of the solution to the stationary problem as claimed in Theorem 1.1 in the general case, that is assuming that the force field satisfies conditions (1.3)–(1.5), without any further assumption on the existence and particular properties of a positive stationary state.

In order to do so, we define the “small” Hilbert space

$$X := L^2(m_0), \quad m_0 := \exp(\kappa \langle x \rangle^\gamma), \quad \kappa \in (0, 1/\gamma),$$

and we study some of the spectral properties of the operator \mathcal{L} acting in X . We consider below the eigenvalue problem on the imaginary axis. We start recalling some standard notions and notations. For an operator Λ acting in a Banach space X , we denote $N(\Lambda) := \Lambda^{-1}(\{0\})$ its null space, $\Sigma(\Lambda)$ its spectrum, $\Sigma_P(\Lambda)$ its point spectrum set, that is the set of the eigenvalues of Λ , while $\rho(\Lambda) := \mathbb{C} \setminus \Sigma(\Lambda)$ will denote its resolvent set and $R_\Lambda(z)$, for $z \in \rho(\Lambda)$, the resolvent operator defined by

$$R_\Lambda(z) := (\Lambda - zI)^{-1}.$$

We denote X_+ the positive cone of X . In this section we shall prove

Theorem 5.1. *There exists a unique steady state $G \in X$, positive, normalized with total mass $\mathbb{M}(G) = 1$, such that*

$$(5.1) \quad N(\mathcal{L}^n) = \text{span}(G), \quad \forall n \geq 1.$$

Moreover, there holds

$$(5.2) \quad \Sigma_P(\mathcal{L}) \cap \{z \in \mathbb{C}; \Re z \geq 0\} = \{0\}.$$

Finally, estimate (1.12) holds true.

We start by recalling some elementary properties regarding the positivity of the semigroup $S_{\mathcal{L}}(t)$.

Proposition 5.2. *The operator \mathcal{L} and its semigroup $S_{\mathcal{L}}$ satisfy the following properties:*

- (a) *The semigroup $S_{\mathcal{L}}$ is positive, namely $S_{\mathcal{L}}(t)f \geq 0$ for any $f \geq 0$ and $t \geq 0$.*
- (b) *The operators \mathcal{L} and \mathcal{L}^* satisfy Kato's inequality, that is if $f \in D(\mathcal{L})$ is complex valued, then we have*

$$(5.3) \quad \mathcal{L}|f| \geq \frac{1}{|f|} \Re(\bar{f} \mathcal{L}f), \quad \mathcal{L}^*|g| \geq \frac{1}{|g|} \Re(\bar{g} \mathcal{L}^*g) \quad \text{in } \mathcal{D}'(\mathbb{R}^d).$$

The above inequality means in particular that for all $\psi \in D(\mathcal{L}^*) \cap X_+$ and all $\varphi \in D(\mathcal{L}) \cap X_+$, we have

$$\langle |f|, \mathcal{L}^* \psi \rangle \geq \Re \langle |f|^{-1} \bar{f} \mathcal{L}f, \psi \rangle, \quad \langle \mathcal{L}\varphi, |g| \rangle \geq \Re \langle \varphi, |g|^{-1} \bar{g} \mathcal{L}g \rangle.$$

- (c) *The operator $-\mathcal{L}$ satisfies a “weak maximum principle”, namely for any $a > 0$ and $g \in X_+$, there holds*

$$(5.4) \quad f \in D(\mathcal{L}) \text{ and } (-\mathcal{L} + a)f = g \text{ imply } f \geq 0.$$

- (d) *The opposite of the resolvent operator is a positive operator, namely for any $a > 0$ and $g \in X_+$, there holds $-R_{\mathcal{L}}(a)g \in X_+$.*

Proof of Proposition 5.2. The argument to establish (5.3) is a classical one, but we outline it briefly. For a smooth complex valued function f if we set $f_\varepsilon := (\varepsilon^2 + |f|^2)^{1/2} - \varepsilon$, then for $1 \leq j \leq d$ we have

$$(5.5) \quad \partial_j f_\varepsilon = \frac{\Re(\bar{f} \partial_j f)}{(\varepsilon^2 + |f|^2)^{1/2}},$$

$$(5.6) \quad \partial_{jj} f_\varepsilon = \frac{\Re(\bar{f} \partial_{jj} f)}{(\varepsilon^2 + |f|^2)^{1/2}} + \frac{|\partial_j f|^2}{(\varepsilon^2 + |f|^2)^{1/2}} - \frac{(\Re(\bar{f} \partial_j f))^2}{(\varepsilon^2 + |f|^2)^{3/2}}.$$

Observe that $f_\varepsilon \rightarrow |f|$ in $H_{\text{loc}}^1(\mathbb{R}^d)$ as $\varepsilon \rightarrow 0$, we have $\partial_j |f| = |f|^{-1} \Re(\bar{f} \partial_j f)$ in $H_{\text{loc}}^1(\mathbb{R}^d)$. Also since

$$\frac{|\partial_j f|^2}{(\varepsilon^2 + |f|^2)^{1/2}} - \frac{(\Re(\bar{f} \partial_j f))^2}{(\varepsilon^2 + |f|^2)^{3/2}} \geq 0,$$

we conclude that $\partial_{jj} f_\varepsilon \geq f_\varepsilon^{-1} \Re(\bar{f} \partial_{jj} f)$. Passing to the limit in the sense of the distributions we obtain $\partial_{jj} |f| \geq |f|^{-1} \Re(\bar{f} \partial_{jj} f)$, and using the expressions of the operators \mathcal{L} and \mathcal{L}^* we obtain (5.3).

The properties (a), (c) and (d) are classical consequences of (5.3), applied to real valued functions, see for instance [2] or [25]. \square

Remark 5.3. For later use, we point out that if $f \in H_{\text{loc}}^2(\mathbb{R}^d)$ is complex valued and moreover is such that $|f| > 0$ on \mathbb{R}^d , then as $\varepsilon \rightarrow 0$ one may pass to the limit in (5.5) and (5.6) and obtain the following equalities

$$(5.7) \quad \partial_j |f| = \frac{\Re(\bar{f} \partial_j f)}{|f|},$$

$$(5.8) \quad \partial_{jj} |f| = \frac{\Re(\bar{f} \partial_{jj} f)}{|f|} + \frac{|\partial_j f|^2}{|f|} - \frac{(\Re(\bar{f} \partial_j f))^2}{|f|^3}.$$

Proposition 5.4. *The operator $-\mathcal{L}$ satisfies a “strong maximum principle”, namely for any given real valued function $f \in X \setminus \{0\}$, there holds*

$$f, |f| \in D(\mathcal{L}), \quad \mathcal{L}f = 0 \text{ and } \mathcal{L}|f| = 0 \text{ imply } f > 0 \text{ or } f < 0.$$

Proof of Proposition 5.4. Consider $f \in X \setminus \{0\}$ such that $f \in D(\mathcal{L})$ and $\mathcal{L}f = 0$. By a bootstrap regularization argument, we classically have $f \in C(\mathbb{R}^d)$. By assumption, there exist then $x_0 \in \mathbb{R}^d$, and two constants $c, r > 0$ such that $|f| \geq c$ on $B(x_0, r)$. From Lemma 2.1, we also have that the operator $\mathcal{L} - aI$ is dissipative for a large enough, in the sense that

$$(5.9) \quad \forall f \in D(\mathcal{L}) \quad ((\mathcal{L} - a)f|f|)_X \leq -\|f\|_X^2.$$

For instance one can take $a := M + 1$, where M is the constant entering in the definition of \mathcal{B} , using the fact that \mathcal{B} is then dissipative and that the same holds for $\mathcal{L} - MI$ because $0 \leq \mathcal{A} \leq M$.

We next observe that for $\sigma > 0$ large enough, the function

$$g(x) := c \exp(\sigma r^2/2 - \sigma|x - x_0|^2/2)$$

satisfies $g = c$ on $\partial B(x_0, r)$ and

$$(-\mathcal{L} + a)g = (a + d\sigma + \sigma \mathbf{F} \cdot (x - x_0) - \text{div} \mathbf{F} - \sigma^2|x - x_0|^2)g \leq 0 \quad \text{on } B(x_0, r)^c.$$

We define $h := (g - |f|)^+$ and $\Omega := \mathbb{R}^d \setminus \bar{B}(x_0, r)$. We have $h \in H_0^1(\Omega, m dx)$ and

$$\begin{aligned} (\mathcal{L} - a)h &\geq \theta'(g - |f|) \mathcal{L}(g - |f|) - ah \\ &= \theta'(g - |f|) [(\mathcal{L} - a)g + a|f|] \geq 0, \end{aligned}$$

where we have used the notation $\theta(s) = s^+$. Thanks to a straightforward generalization of (5.9) to $H_0^1(\Omega, m)$, we deduce

$$0 \leq ((\mathcal{L} - a)h|h)_{L^2(\Omega, m)} \leq -\|h\|_{L^2(\Omega, m)}^2,$$

and then $h = 0$. This implies that we have $|f| \geq g$ on $\Omega = \bar{B}(x_0, r)^c$, and thus $|f| > 0$ on \mathbb{R}^d . Since $f \in C(\mathbb{R}^d)$, we must have either $f \geq g > 0$ or $f \leq -g < 0$. \square

Proof of Theorem 5.1. We split the proof into six steps.

Step 1. We prove that there exists $G \in N(\mathcal{L})$ such that $G > 0$, which yields in particular that $0 \in \Sigma_P(\mathcal{L})$. In other words, we prove that there exists a positive and normalized (with mass 1) steady state $G \in X$. For the equivalent norm $\|\cdot\|$ defined on X by

$$\|f\| := \sup_{t>0} \|S_{\mathcal{L}}(t)f\|_X,$$

we have $\|S_{\mathcal{L}}(t)f\| \leq \|f\|$ for all $t \geq 0$, that is the semigroup $S_{\mathcal{L}}$ is a contraction semigroup on $(X, \|\cdot\|)$. There exists $R > 0$ large enough such that the intersection of the closed hyperplane $H := \{f; \mathbb{M}(f) = 1\}$ and the closed ball of radius R in

$(X, \|\cdot\|)$ is a convex, non empty subset. Then consider the closed, weakly compact convex set

$$\mathbb{K} := \{f \in X_+ ; \|f\| \leq R, \quad \mathbb{M}(f) = 1\}.$$

Since $S_{\mathcal{L}}(t)$ is a linear, weakly continuous, contraction in $(X, \|\cdot\|)$ and $\mathbb{M}(S_{\mathcal{L}}(t)f) = \mathbb{M}(f)$ for all $t \geq 0$, we see that \mathbb{K} is stable under the action of the semigroup. Therefore we apply the Markov–Kakutani fixed point theorem (see for instance [13, Theorem 6, chapter V, § 10.5, page 456] or its (possibly nonlinear) variant [15, Theorem 1.2]) and we conclude that there exists $G \in \mathbb{K}$ such that $S_{\mathcal{L}}(t)G = G$. Therefore we have in particular $G \in D(\mathcal{L})$ and $\mathcal{L}G = 0$. Moreover since $G \geq 0$, $G \neq 0$ and $G \in C(\mathbb{R}^d)$, we may conclude that $G > 0$ on \mathbb{R}^d by the strong maximum principle.

Step 2. In this step we prove that $N(\mathcal{L}^2) = N(\mathcal{L})$, which implies that $N(\mathcal{L}^n) = N(\mathcal{L})$ for all $n \geq 1$. Otherwise, there would exist $g_1 \in D(\mathcal{L})$ with $\|g_1\| = 1$ and $g_2 \in D(\mathcal{L})$ such that $\mathcal{L}g_1 = 0$ and $\mathcal{L}g_2 = g_1$. Since $S_{\mathcal{L}}(t)g_1 = g_1$ for all $t \geq 0$, one sees easily that one must have $S_{\mathcal{L}}(t)g_2 = g_2 + tg_1$. However, since $g_1 \neq 0$, this is in contradiction with the fact that the semigroup $S_{\mathcal{L}}$ is bounded on X as stated in Lemma 3.1.

Step 3. We prove here that $N(\mathcal{L}) = \text{span}(G)$. Since \mathcal{L} is a real operator, we may restrict ourselves to real valued eigenfunctions. Consider a real valued eigenfunction $f \in N(\mathcal{L})$ with $f \neq 0$. First we observe that thanks to Kato’s inequality we have

$$0 = (\mathcal{L}f) \text{sgn}(f) \leq \mathcal{L}|f|.$$

Actually this inequality must be an equality, since otherwise we would have

$$0 \neq \langle \mathcal{L}|f|, 1 \rangle = \langle |f|, \mathcal{L}^*1 \rangle = 0,$$

which is a contradiction. As a consequence, we have also $|f| \in N(\mathcal{L})$, so that the strong maximum principle, Proposition 5.4, implies that either $f > 0$ or $f < 0$, and without loss of generality we may assume that $f > 0$, and up to a multiplication by a normalization factor, we may also assume that $\mathbb{M}(f) = 1$. Now, using once more Kato’s inequality we have

$$0 = \mathbf{1}_{[f-G>0]} \mathcal{L}(f-G) \leq \mathcal{L}(f-G)^+,$$

and due to the same reasons as above, we may conclude that this last inequality is in fact an equality, that is $(f-G)^+ \in N(\mathcal{L})$. The strong maximum principle implies that either $(f-G)^+ \equiv 0$ or $(f-G)^+ > 0$ on \mathbb{R}^d . This means that either we have $f \leq G$ or $f > G$ on \mathbb{R}^d . Thanks to the normalization hypothesis $\mathbb{M}(f) = \mathbb{M}(G) = 1$, the second possibility must be excluded and thus we have $f \leq G$ on \mathbb{R}^d . Repeating the same argument with $(G-f)^+$ we deduce that $G \leq f$ and we conclude that $f = G$.

Step 4. We prove here that $i\mathbb{R} \cap \Sigma_P(\mathcal{L}) = \{0\}$: that is that the only eigenvalue with vanishing real part is zero, or in other words, (5.2) holds. We consider a couple (f, μ) of eigenfunction and eigenvalue, with $\mu := i\omega \in i\mathbb{R}$, and normalized so that $\mathbb{M}(|f|) = 1$. Using the complex version of Kato’s inequality (5.3), we have

$$(5.10) \quad \mathcal{L}|f| \geq \frac{1}{|f|} \Re(\bar{f} \mathcal{L}f) = 0.$$

Computing $\langle \mathcal{L}|f|, 1 \rangle$, thanks to the above inequality, we get

$$0 \leq \langle \mathcal{L}|f|, 1 \rangle = \langle |f|, \mathcal{L}^*1 \rangle = 0,$$

which implies that the inequality is in fact an equality, that is $\mathcal{L}|f| = 0$, and since $\mathbb{M}(G) = \mathbb{M}(|f|) = 1$, we conclude that $|f| = G > 0$.

Next, using Remark 5.3 and (5.7)–(5.8), we have

$$\begin{aligned} |f| \mathcal{L}|f| &= \Re e(\bar{f} \Delta f) + \Re e(\bar{f} \mathbf{F} \cdot \nabla f) + \operatorname{div}(\mathbf{F})|f|^2 + |\nabla f|^2 - \frac{1}{|f|^2} |\Re e(\bar{f} \nabla f)|^2 \\ &= \Re e(\bar{f} \mathcal{L}f) + |\nabla f|^2 - \frac{1}{|f|^2} |\Re e(\bar{f} \nabla f)|^2. \end{aligned}$$

Together with $\mathcal{L}|f| = 0$ and $\Re e(\bar{f} \mathcal{L}f) = 0$ (since $\mathcal{L}f = i\omega f$), the above identity implies

$$(5.11) \quad |\nabla f|^2 - \frac{1}{|f|^2} |\Re e(\bar{f} \nabla f)|^2 = 0.$$

From this, as explained below, we infer that $f = \exp(i\theta)G$ for some constant $\theta \in [0, 2\pi]$, and thus $\mathcal{L}f = \exp(i\theta)\mathcal{L}G = 0$ and $\omega = 0$.

Indeed, in order to see that $f = \exp(i\theta)G$, for some $\theta \in [0, 2\pi]$, let us write

$$f = u + iv$$

for two real valued functions u and v . Then, since $\Re e(\bar{f} \nabla f) = u\nabla u + v\nabla v$, relation (5.11) means that

$$(u^2 + v^2) (|\nabla u|^2 + |\nabla v|^2) = |u\nabla u + v\nabla v|^2,$$

which yields $v\nabla u - u\nabla v = 0$. Since u and v are both continuous functions and are not both identically equal to zero, we may assume for instance that there exists $x_0 \in \mathbb{R}^d$ such that $u(x_0) > 0$. Thus, if we denote by Ω the connected component of $\{x \in \mathbb{R}^d; u(x) > 0\}$ containing x_0 , we have $\nabla(v/u) = 0$ on Ω . Hence $v = \alpha u$ on Ω for some $\alpha \in \mathbb{R}$. However we must have $\Omega = \mathbb{R}^d$, since otherwise we would have $|f| = 0$ on $\partial\Omega$, which would be a contradiction with the fact that $G = |f| > 0$ in \mathbb{R}^d . Thus, setting $a := 1 + i\alpha$ we get $f = au$ and thus $u = |u| = |f|/|a|$ is a positive steady state, so that $|u| = G/|a| > 0$ from Step 3. We conclude that $f = aG/|a| = \exp(i\theta)G$ for some $\theta \in [0, 2\pi]$.

Step 5. Denoting $m := e^{\kappa\langle x \rangle^\gamma}$, with $\kappa \in (0, 1/\gamma)$, and $\alpha := \kappa\gamma(1 - \kappa\gamma)/2 > 0$, the function $\psi_{m,1}^0$ being defined by (2.4) for $p = 1$, $M, R > 0$ large enough so that

$$m\psi_{m,1}^0 = \mathcal{L}^*m - M\chi_{Rm} = \Delta m - \mathbf{F} \cdot \nabla m - M\chi_{Rm} \leq -\alpha m,$$

with $\alpha > 0$, we have

$$0 = \int (\mathcal{L}G)m = \int G(\mathcal{L}^*m) = \int G(m\psi_{m,1}^0 + M\chi_{Rm}),$$

and therefore

$$\int Gm \leq \frac{1}{\alpha} \int GM\chi_{Rm} \leq \frac{M}{\alpha} \mathbb{M}(G).$$

We have proved $G \in L^1(m)$. Using the regularization property of $S_{\mathcal{B}}$ established in Lemma 2.3 and Lemma 2.4 as well as the same Duhamel formula as in Step 1 in the proof of Lemma 3.1, we deduce $G \in L^\infty(m)$.

Step 6. We define $\kappa_1^* := \gamma^{-1} \limsup_{|x| \rightarrow \infty} x \cdot \mathbf{F}(x)/|x|^\gamma$ and we know that $\kappa^* \in [1/\gamma, \infty)$ from assumptions (1.3) and (1.4). We then denote $g := ce^{\kappa|x|^\gamma}$, with

$\kappa > \kappa_1^*$, and we compute

$$\begin{aligned} g^{-1}[(-\mathcal{L} + a)g] &= -\kappa^2 \gamma^2 |x|^{2(\gamma-1)} + \kappa \gamma \mathbf{F} \cdot x |x|^{\gamma-2} \\ &\quad + a - \operatorname{div} \mathbf{F} + \kappa \gamma (d + \gamma - 2) |x|^{\gamma-2} \leq 0 \quad \text{on } B(0, R)^c, \end{aligned}$$

for some $R > 0$ large enough. Arguing as in the proof of Proposition 5.4, we deduce that $G \geq g$ on $B(0, R)^c$, and this yields the claimed lower bound (see also the proof of Lemma 6.2 below for similar arguments). \square

6. THE GENERAL CASE - DECAY ESTIMATE

In this section we present the proof of our main Theorem 1.2 in the general case, that is assuming that the force field satisfies conditions (1.3)–(1.5). We denote by X_0 the “small” Banach space

$$X_0 := L^p(m_0), \quad m_0 := \exp(\kappa \langle x \rangle^\gamma), \quad \kappa \in (0, \kappa_0^*), \quad p = 1 \text{ or } 2,$$

and we will establish a semigroup decay in that space using three different strategies. The proof of the decay of the semigroup in a general Banach space $L^p(m)$, with a weight function m which is a polynomial or an exponential function, then follows the same arguments as the ones which have been developed in section 4.2 and may be skipped here.

6.1. Modified Poincaré inequality approach. In this section we fix the weight function $m = m_0$ as defined above. We recall the so-called *Lyapunov condition*

$$(6.1) \quad \mathcal{L}^* m_0 \leq -\xi_0 + M \mathbf{1}_{B_{R_0}}, \quad \xi_0(x) := \zeta_0 m_0(x) \langle x \rangle^{2(\gamma-1)},$$

which holds for some appropriate constants $\zeta_0, M, R_0 \in (0, \infty)$ and which has been established in Step 1 in the proof of Lemma 2.1. We also recall that the stationary state G given by Theorem (5.1) satisfies estimate (1.12), and thus a local Poincaré (or Poincaré-Wirtinger) inequality on any ball with constants which may be estimated explicitly. Under both above assumptions, we may adapt the proof of [4, Theorem 3.6] by Bakry, Cattiaux and Guillin in order to get the following “*weak Lyapunov-Poincaré inequality*” which is a generalisation of the *weak weighted Poincaré-Wirtinger inequality* (1.14) established in Lemma 4.1.

Lemma 6.1. *There exists a constant $\beta \in (0, \infty)$ such that denoting $w_0 := 1 + \beta m_0$ and recalling that the function ξ_0 has been defined by (6.1), the following inequality*

$$(6.2) \quad \frac{1}{4} \int h^2 \xi_0 G \leq \int \left(w_0 |\nabla h|^2 - \frac{1}{2} h^2 \mathcal{L}^* w_0 \right) G,$$

holds for all functions $h \in C_b^1(\mathbb{R}^d)$ such that $\mathbb{M}(hG) = 0$.

Proof of Lemma 6.1. We take $h \in C_b^1(\mathbb{R}^d)$ such that $\mathbb{M}(hG) = 0$. From (6.1), we have

$$\int h^2 \xi_0 G \leq \int h^2 (M \mathbf{1}_{B_{R_0}} - \mathcal{L}^* m_0) G.$$

For any $R \geq R_0$, the local Poincaré inequality writes

$$\begin{aligned} \int h^2 \mathbf{1}_{B_R} G &\leq \kappa_R \int_{B_R} |\nabla h|^2 G + \frac{1}{\mathbb{M}(G \mathbf{1}_{B_R})} \left(\int_{B_R^c} h G \right)^2 \\ &\leq \frac{\kappa_R}{1 + \beta} \int |\nabla h|^2 w_0 G + \frac{\mathbb{M}(\xi_0^{-1} G \mathbf{1}_{B_R^c})}{\mathbb{M}(G \mathbf{1}_{B_R})} \int h^2 \xi_0 G, \end{aligned}$$

where we have used the the zero mass condition $\mathbb{M}(hG) = 0$ in the first line and the Cauchy-Schwarz inequality in the second line. We then deduce

$$\left(1 - M \frac{\mathbb{M}(\xi_0^{-1} G \mathbf{1}_{B_R^c})}{\mathbb{M}(G \mathbf{1}_{B_R})}\right) \int h^2 \xi_0 G \leq \frac{M \kappa_R}{1 + \beta} \int |\nabla h|^2 w_0 G + \int h^2 (-\mathcal{L}^* w_0) G.$$

We may then first fix $R > R_0$ large enough and next $\beta > 0$ large enough in such a way that (6.2) holds. \square

We introduce the weight function $m := w_0^{1/2} G^{-1/2}$ and the associated Hilbert space $L^2(m)$. Despite the fact that in general the operator \mathcal{L} is not symmetric in $L^2(m)$, the associated Dirichlet form has a nice positivity property. More precisely, in order to make the discussion simpler, we shall consider the conjugate operator $h \mapsto Lh := G^{-1} \mathcal{L}(Gh)$, which amounts to make the change of function $h := G^{-1} f$, and thus we define

$$Lh := G^{-1} \mathcal{L}(Gh) = \Delta h + \left(2 \frac{\nabla G}{G} + \mathbf{F}\right) \cdot \nabla h.$$

For any smooth function f , defining the quadratic form $\mathcal{E}(f) := (-\mathcal{L}f|f)_{L^2(m)}$, we compute, upon integrating by parts

$$\begin{aligned} \mathcal{E}(f) &:= (-\mathcal{L}f|f)_{L^2(m)} = (-Lh|h)_{L^2(w_0^{1/2} G^{1/2})} \\ &= - \int \left[\Delta h + \left(2 \frac{\nabla G}{G} + \mathbf{F}\right) \cdot \nabla h \right] h w_0 G \\ &= \int |\nabla h|^2 G w_0 - \frac{1}{2} \int \nabla(h^2) [-G \nabla w_0 + w_0 \nabla G + \mathbf{F} w_0 G] \\ &= \int |\nabla h|^2 G w_0 + \frac{1}{2} \int h^2 [-G \Delta w_0 + w_0 \Delta G + \operatorname{div}(w_0 G \mathbf{F})] \\ &= \int |\nabla h|^2 G w_0 - \frac{1}{2} \int h^2 G \mathcal{L}^* w_0, \end{aligned}$$

where we have used the fact that G is a steady state in the last line. For any smooth function f such that $\mathbb{M}(f) = 0$, we deduce from the weak Lyapunov-Poincaré inequality (6.2) that

$$\mathcal{E}(f) \geq \frac{1}{4} \int f^2 \xi_0 G^{-1}.$$

As a consequence, denoting $f(t) := S_{\mathcal{L}}(t) f_0$ for $f_0 \in X_0$ with $\mathbb{M}(f_0) = 0$, we compute

$$\frac{d}{dt} \int f^2 w_0 G^{-1} = -\mathcal{E}(f) \leq -\frac{1}{4} \frac{\zeta_0}{1 + \beta} \int f^2 \langle x \rangle^{2(\gamma-1)} w_0 G^{-1}.$$

This differential inequality is similar to (4.10) and we thus may conclude as in Section 4 that

$$\|f_t\|_{L^2} \lesssim \Theta_{G^{-1}}(t) \|f_0\|_{L^2(m)}, \quad \forall t \geq 0.$$

6.2. Harris-Meyn-Tweedie approach. We present now a second way to get a decay estimate in one small space. We start establishing a strict positivity estimate which will be the main technical estimate in our analysis.

Lemma 6.2. *For any $R > 0$, there exist $T > 0$ and $\psi \in L^\infty(\mathbb{R}^d)$, $\psi \geq 0$, $\psi \not\equiv 0$, such that*

$$S_{\mathcal{L}}(T)g_0 \geq \psi \int_{B_R} g_0 dx, \quad \forall g_0 \geq 0.$$

Proof of Lemma 6.2. We split the proof into three parts.

Step 1. We prove that for any $p \in [1, \infty)$ and $m := \langle x \rangle^k$, $k > 0$ large enough, there exist $C, \alpha > 1$, such that for any $g_0 \in L^1(m)$, the function $g(t) := S_{\mathcal{L}}(t)g_0$ satisfies

$$(6.3) \quad \|\nabla g(t)\|_{L^p} \leq \frac{C}{t^\alpha} \|g_0\|_{L^1(m)}, \quad \forall t > 0.$$

First, as a consequence of Lemma 2.4 and the iterated Duhamel formula

$$S_{\mathcal{L}} = S_{\mathcal{B}} + \cdots + (S_{\mathcal{B}}\mathcal{A})^{*(n-1)} * S_{\mathcal{B}} + (S_{\mathcal{B}}\mathcal{A})^{*n} * S_{\mathcal{L}},$$

we have that $S_{\mathcal{L}}$ is ultra-contractive, that is

$$\|g(t)\|_{L^\infty} \leq \frac{C}{t^{d/2}} \|g_0\|_{L^1(m)}, \quad \forall t > 0.$$

Next, we differentiate the Fokker-Planck equation and proceeding in a similar way as in the proof of Lemma 2.6, we get

$$\|g(t)\|_{H^1} \leq \frac{C}{t^{1/2}} \|g_0\|_{L^2(m)}, \quad \forall t > 0.$$

Repeating the proof of Lemma 2.4 for the derivative of $g(t)$, we may conclude in a quite standard way to (6.3).

Step 2. We prove that for any $t > 0, R > 0$, there exist $r, \lambda > 0$ and $x_0 \in B_R$ such that

$$S_{\mathcal{L}}(t)g_0 \geq \lambda \mathbf{1}_{B(x_0, r)} \int_{B_R} g_0 dx.$$

We fix $R > 0$ and $g_0 \in L^1(\mathbb{R}^d)$ with $\text{supp } g \subset B_R$. With the same notations as in the previous step, we have

$$\int g(t, x) dx = \int g_0(x) dx = \int_{B_R} g_0 dx.$$

Moreover, thanks to Lemma 3.1, there exists $A > 0$, such that

$$\int g(t, x) |x| dx \leq A \int g_0 |x| dx \leq AR \int_{B_R} g_0 dx.$$

For any $\rho > 0$, , we write

$$\begin{aligned} \int_{B_\rho} g(t, x) dx &= \int g(t, x) dx - \int_{B_\rho^c} g(t, x) dx \\ &\geq \int g_0 dx - \frac{1}{\rho} \int g(t, x) |x| dx \\ &\geq \left(1 - \frac{AR}{\rho}\right) \int_{B_R} g_0(x) dx \geq \frac{1}{2} \int_{B_R} g_0(x) dx, \end{aligned}$$

by choosing $\rho > 0$ large enough. As a consequence we infer that there exists $x_0(t) \in B_\rho$ such that

$$g(t, x_0(t)) \geq 2\lambda := \frac{1}{2|B_\rho|} \int_{B_\rho} g(t, x) dx \geq \frac{1}{4|B_\rho|} \int_{B_R} g_0(x) dx.$$

Thanks to (6.3) and Morrey embedding $W^{1,p}(\mathbb{R}^d) \subset C^{0,1/2}(\mathbb{R}^d)$, which holds true for $p \in (1, \infty)$ large enough, we know that there exists a constant $c_0 > 0$ such that for all $t > 0$ and $x, y \in \mathbb{R}^d$

$$|g(t, x) - g(t, y)| \leq c_0 \|\nabla g(t, \cdot)\|_{L^p} \cdot |x - y|^{1/2} \leq c_0 C \|g_0\|_{L^1} t^{-\alpha} |x - y|^{1/2}.$$

Therefore, we have

$$g(t, x) \geq g(t, x_0(t)) - c_0 C \|g_0\|_{L^1} t^{-\alpha} |x - x_0(t)|^{1/2} \geq \lambda$$

for $x \in B(x_0(t), r)$, where r is chosen by setting $c_0 C \|g_0\|_{L^1} t^{-\alpha} r^{1/2} = \lambda$.

Step 3. (Spreading of the positivity.) We prove that for any $r_0, r_1 > 0$ and $x_0 \in \mathbb{R}^d$, there exist $t_1, \kappa_1 > 0$ (computable) such that

$$g_0 \geq \mathbf{1}_{B(x_0, r_0)} \quad \Rightarrow \quad g(t_1, \cdot) \geq \kappa_1 \mathbf{1}_{B(x_0, r_1)}.$$

We first observe that $h(t, x) := g(t, x) e^{t\|\operatorname{div}\mathbf{F}\|_\infty}$ is a super solution of the Fokker-Planck equation with initial data g_0 , that is $\partial_t h - \mathcal{L}h \geq 0$ and

$$\partial_t h \geq \mathcal{L}^\sharp h := \Delta h + \mathbf{F} \cdot \nabla h.$$

We introduce the function

$$\varphi(t, x) := \varphi_0 - \varepsilon, \quad \varphi_0 := (t + t_0)^{-\alpha} e^{-\mu|x-x_0|^2/(t+t_0)},$$

for some parameters $\alpha, \mu, \varepsilon, t_0 > 0$ to be specified latter. We compute

$$\begin{aligned} \partial_t \varphi &= \left(\mu \frac{|x-x_0|^2}{(t+t_0)^2} - \frac{\alpha}{t+t_0} \right) \varphi_0, \quad \nabla_x \varphi = -2\mu \frac{x-x_0}{t+t_0} \varphi_0, \\ \Delta_x \varphi &= \left(4\mu^2 \frac{|x-x_0|^2}{(t+t_0)^2} - 2 \frac{\mu d}{t+t_0} \right) \varphi_0. \end{aligned}$$

To simplify notations, we only consider the case $x_0 = 0$ (the general case can be dealt exactly in the same way) and we denote $r = |x|$ and $\tau = t + t_0$, and we consider φ as function of τ and r . For $\mu \geq 1$, $T > 0$ and $t \in (0, T)$, we deduce

$$\begin{aligned} \frac{\partial_t \varphi - \mathcal{L}^\sharp \varphi}{\varphi_0} &= \mu \frac{r^2}{\tau^2} - \frac{\alpha}{\tau} + 2d \frac{\mu}{\tau} - 4\mu^2 \frac{r^2}{\tau^2} + 2\mu \frac{\mathbf{F} \cdot x}{\tau} \\ &\leq -\frac{\alpha}{\tau} - 2\mu^2 \frac{r^2}{\tau^2} + 2d \frac{\mu}{\tau} + 4\|\mathbf{F}\|_{L^\infty}, \end{aligned}$$

and then

$$(6.4) \quad \partial_t \varphi \leq \mathcal{L}^\sharp \varphi \quad \text{on} \quad (0, T) \times \mathbb{R}^d$$

for any $\alpha \geq 2d\mu + 4T\|\mathbf{F}\|_{L^\infty}$. We fix ε such that

$$\varphi(0, r_0) = t_0^{-\alpha} e^{-\mu r_0^2/t_0} - \varepsilon = 0.$$

We have built φ such that

$$\partial_t (h - \varphi) \geq \mathcal{L}^\sharp (h - \varphi) \quad \text{on} \quad (0, T) \times \mathbb{R}^d, \quad (h - \varphi)(0) \geq 0 \quad \text{on} \quad \mathbb{R}^d.$$

Because $(h - \varphi)_- \in L^1((0, T) \times \mathbb{R}^d)$, by the weak maximum principle we deduce that $h \geq \varphi$ on $(0, T) \times \mathbb{R}^d$, so that

$$g(t, x) \geq e^{-t\|\operatorname{div}\mathbf{F}\|_{L^\infty}} \varphi(t, x) \quad \text{on} \quad (0, T) \times \mathbb{R}^d.$$

We now define $r_1 := r_0 t_0^{-\beta} \tau^\beta$ as well as

$$\psi(t) := \varphi(t, r_1) = \tau^{-\alpha} e^{-\mu r_0^2 t_0^{-2\beta} \tau^{2\beta-1}} - t_0^{-\alpha} e^{-\mu r_0^2/t_0}.$$

We have (recall that $\tau = t + t_0$)

$$(6.5) \quad \psi(0) = 0, \quad \psi'(t) = \left(\frac{\mu r_0^2 t_0^{-2\beta} (1-2\beta)}{\tau^{1-2\beta}} - \alpha \right) \tau^{-\alpha-1} e^{-\mu r_0^2 \tau^{2\beta-1}} \geq 0$$

whenever $\alpha \leq \frac{\mu r_0^2 t_0^{-2\beta} (1-2\beta)}{\tau^{1-2\beta}}$. We choose $\mu = 1 + r_0^{-2}$, $T = 2t_0 \leq 1$ to be fixed below, $\beta = 1/4$ and $\alpha := 2d + 4\|\mathbf{F}\|_{L^\infty} \geq 2$. We observe that (6.4) is true on the time interval $(0, T)$ and (6.5) is true on $(0, 2t_0)$ by choosing $t_0 := 1/(4\alpha) \leq 1/2$. In particular, $\psi(2t_0) \geq 0$. As a consequence,

$$\varphi(2t_0, x) \geq \bar{\varphi}_1 := \varphi(2t_0, r_{t_0}) = (3t_0)^{-\alpha} e^{-\mu r_0^2 2^{-1/2}/t_0} - t_0^{-\alpha} e^{-\mu r_0^2/t_0} > 0,$$

for any $|x| < r_{t_0} := r_0 2^{1/4}$ and then

$$g(2t_0, x) \geq \bar{\varphi}_1 e^{-2t_0 \|\operatorname{div} \mathbf{F}\|_{L^\infty}} \mathbf{1}_{B(x_0, r_0 2^{1/4})},$$

by the above weak maximum principle. Iterating this argument, we get the announced estimate with

$$t_1 := 2nt_0, \quad n := \left(1 + \left(4 \frac{\log r_1/r_0}{\log 2} \right) \right), \quad \kappa_1 := e^{-2t_1 \|\operatorname{div} \mathbf{F}\|_{L^\infty}} \prod_{k=1}^n \bar{\varphi}_k$$

and

$$\bar{\varphi}_k := (3t_0)^{-\alpha} e^{-\mu \rho_k^2 2^{-1/2}/t_0} - t_0^{-\alpha} e^{-\mu \rho_k^2/t_0}, \quad \rho_k := r_0 2^{k/4}.$$

Step 4 and conclusion. From Step 2, we know that we may find $\tau_0, r_0, R, \lambda > 0$ and $x_0 \in B(0, R)$ such that

$$g(\tau_0, \cdot) \geq c_0 \mathbf{1}_{B(x_0, r_0)}, \quad c_0 := \lambda \int_{B_R} g_0 dx.$$

Using Step 3 with $r_1 := 2R$, we obtain

$$g(\tau_0 + t_1, \cdot) \geq \lambda \kappa_1 \mathbf{1}_{B(0, R)} \int_{B_R} g_0 dx,$$

which is nothing but the announced Harris condition. \square

We recall the following so called “*abstract subgeometric Harris-Meyn-Tweedie theorem*”.

Theorem 6.3 (subgeometric Harris-Meyn-Tweedie). *Consider a semigroup S in $L^1(m_0)$ which is positive and mass conservative. Assume furthermore that the following two additional conditions hold:*

(1) *the Lyapunov condition*

$$\mathcal{L}^* m \leq -\varphi(m) + b,$$

for some weight function $m : \mathbb{R}^d \rightarrow [1, \infty)$ converging to infinity and some positive, concave and converging to infinity function $\varphi : \mathbb{R}_+ \rightarrow \mathbb{R}_+$ and some constant $b > 0$;

(2) *the strict positivity (or irreducibility) condition*

$$S_T f \geq \nu \int_{B_R} f, \quad \forall f \geq 0,$$

for any $R > 0$ and some $T > 0$ and $\nu \geq 0$, $\nu \neq 0$. Then, there exists some constant $C > 0$ such that

$$(6.6) \quad \|S_t f_0\|_{L^1} \leq \frac{C}{(\varphi \circ H_\varphi^{-1})(t)} \|f_0\|_{L^1(m)},$$

for any $f_0 \in L^1(m)$, $\mathbb{M}(f_0) = 0$, where H_φ is the function defined by

$$H_\varphi(u) := \int_1^u \frac{ds}{\varphi(s)}.$$

Theorem 6.3 is stated in [4, Theorem 1.2] and it is proved in [12, (3.5) in Theorem 3.2] (see also [12, Section 5.1]) by the mean of probabilistic tools. A simple semigroup proof is presented in [9]. A variant of [4, Theorem 1.2] is also presented in [19, Section 4]. Gathering the Lyapunov condition (6.1), the strict positivity property established in Lemma 6.2 and the subgeometric Harris-Meyn-Tweedie Theorem 6.3, we deduce the following decay estimate.

Corollary 6.4. *For any $f_0 \in L^1(m_0)$ such that $\mathbb{M}(f_0) = 0$, the associated solution $f(t, \cdot) = S_{\mathcal{L}}(t)f_0$ to the Fokker-Planck equation (1.1) satisfies*

$$(6.7) \quad \|f(t, \cdot)\|_{L^1} \lesssim \Theta_{m_0}(t) \|f_0\|_{L^1(m_0)},$$

which is nothing but the conclusion of Theorem 1.2 in $L^1(m_0)$.

Proof of Corollary 6.4. We write (1.17) (or (6.1)) as

$$(6.8) \quad \mathcal{L}^* m_0 \leq -\varphi(m_0) + M,$$

with

$$\varphi(y) := C \frac{y}{(\log y)^\alpha}, \quad \alpha := 2 \frac{1-\gamma}{\gamma}.$$

We easily compute

$$H_\varphi(u) = C \int_1^u (\log y)^\alpha \frac{dy}{y} = C \int_1^{\log u} z^\alpha dz = C (\log u)^{\alpha+1},$$

so that $H_\varphi^{-1}(v) = \exp(cv^{\frac{1}{\alpha+1}})$. From (6.6), we conclude (6.7) with $\Theta_{m_0}(t) := K \exp(-\lambda t^{\frac{\gamma}{2-\gamma}})$, $K \geq 1$, $\lambda > 0$. \square

6.3. Krein-Rutman type approach. We finally present a third approach which is mainly an adaptation of an argument used in the proof of [28, Theorem 2.1]. The key argument is an accurate estimate on the confinement process of the semigroup and it does not deeply use the positivity estimates. The drawback comes from the fact that the rate of decay in the small space X_0 is lower than in the small space $L^2(G^{-1/2})$ (see section 4.2 for the notations).

We consider the sequence of spaces $(X_k)_{k \in \mathbb{N}}$, as defined in (2.22), and $X_\infty := L^2(m_0^{1/2})$. We observe that $X_k \subset X_{k+1} \subset X_\infty$ for any $k \in \mathbb{N}$. For $0 \leq \eta \leq 1$ we also denote $X_{k,\eta}$ the space defined by Hilbertian interpolation between $X_{k,0} = X_k$ and $X_{k,1} := \{f \in X_k; \mathcal{L}f \in X_k\}$, that is, with the notations of L. Tartar [32, Chapter 22, page 109],

$$X_{k,\eta} := (X_k, X_{k,1})_{\eta,2}.$$

Lemma 6.5. *Let us fix an integer $j > 2(1-\gamma)/\gamma > 0$. There exists a constant C such that for any $\ell_1, \ell_2, k \in \mathbb{N}$, $k \geq j$, $\ell_i \geq 1$, we have for all $z \in \mathbb{C}$ with $\Re z > 0$*

$$(6.9) \quad \|R_{\mathcal{B}}(z)\|_{X_{k-j} \rightarrow X_k} \leq \mathcal{C}_k := C k^j,$$

$$(6.10) \quad \|\mathcal{A}R_{\mathcal{B}}(z)^{\ell_1} \mathcal{A}R_{\mathcal{B}}(z)^{\ell_2}\|_{X_0 \rightarrow X_0} \leq (\ell_1! \ell_2!)^j C^{\ell_1 + \ell_2} / \langle y \rangle^{1/2},$$

where $z = x + iy$, $x, y \in \mathbb{R}$.

Proof of Lemma 6.5. We use the representation formula

$$R_{\mathcal{B}}(z) = - \int_0^\infty e^{-zt} S_{\mathcal{B}}(t) dt$$

together with the estimates established in Lemmas 2.1, 2.4 and 2.5 in the following way. In order to simplify the presentation, we only consider in the sequel the boundary case $\Re z = x = 0$.

On the one hand, we define, as in Lemma 2.5, $\alpha^* := 1/2(1-\gamma)$, so that $\alpha^* \gamma j > 1$, and we observe that the LHS term of (2.23) belongs to $L^1(\mathbb{R}_+)$. Therefore, with the notations of Lemma 2.5, we have for any $y \in \mathbb{R}$

$$\|R_{\mathcal{B}}(iy)\|_{X_{k-j} \rightarrow X_k} \leq C_1 + \left(\frac{k}{\kappa}\right)^j \frac{1}{\alpha \gamma j - 1} \leq C k^j.$$

On the other hand, from (2.26), we have

$$(6.11) \quad \sup_{y \in \mathbb{R}} \|R_{\mathcal{B}}(iy)\|_{X_\infty \rightarrow H^1} \leq \int_0^\infty \|S_{\mathcal{B}}(t)\|_{X_\infty \rightarrow H^1} dt \leq C.$$

The latter estimate together with (6.9) yield that, for any $y \in \mathbb{R}$ and $\ell_2 \in \mathbb{N}^*$, we have

$$(6.12) \quad \begin{aligned} \|\mathcal{A}R_{\mathcal{B}}(iy)^{\ell_2}\|_{X_0 \rightarrow X_{0,1/2}} &\leq \\ &\leq \|\mathcal{A}R_{\mathcal{B}}(iy)\|_{X_\infty \rightarrow X_{0,1/2}} \|R_{\mathcal{B}}(iy)\|_{X_{(\ell_2-2)j} \rightarrow X_{(\ell_2-1)j}} \cdots \|R_{\mathcal{B}}(iy)\|_{X_0 \rightarrow X_j} \\ &\leq (\ell_2!)^j C^{\ell_2}. \end{aligned}$$

On the other hand, from the identity

$$R_{\mathcal{B}}(z) = z^{-1}(\mathcal{R}_{\mathcal{B}}(z)\mathcal{B} - I)$$

and an interpolation argument, we deduce that

$$\|R_{\mathcal{B}}(iy)\|_{X_{0,1/2} \rightarrow X_j} \leq C/\langle y \rangle^{1/2},$$

and therefore for any $\ell_1 \in \mathbb{N}$

$$\|\mathcal{A}R_{\mathcal{B}}(iy)^{\ell_1}\|_{X_{0,1/2} \rightarrow X_0} \leq (\ell_1!)^j C^{\ell_1} / \langle y \rangle^{1/2}.$$

It is now clear that the above estimate together with (6.12) completes the proof of estimate (6.10). \square

Lemma 6.6. *Let us fix again an integer $j > 2(1-\gamma)/\gamma > 0$. There exists a constant C such that for any $\ell \in \mathbb{N}^*$, denoting by Π the projection on G , that is $\Pi(f) := \mathbb{M}(f)G$, we have*

$$(6.13) \quad \sup_{\Re z > 0} \|(I - \Pi)R_{\mathcal{L}}(z)^\ell\|_{X_0 \rightarrow X_\infty} \leq C^\ell (\ell!)^j.$$

Proof of Lemma 6.6. Since the operator $\mathcal{L} - aI$ is dissipative for any $a > 0$, we clearly have

$$C_{\mathcal{L},a} := \sup_{\Re z \geq a} \|(I - \Pi)R_{\mathcal{L}}(z)\|_{X_0 \rightarrow X_\infty} \lesssim \sup_{\Re z \geq a} \|R_{\mathcal{L}}(z)\|_{X_0 \rightarrow X_\infty} < \infty,$$

and thus we have only to prove that the constant $C_{\mathcal{L},a}$ does not blow up when $a \rightarrow 0^+$.

Step 1. We claim that for any fixed M , there holds

$$(6.14) \quad \sup_{z \in \Omega_M} \|(I - \Pi)R_{\mathcal{L}}(z)\|_{X_{k-j} \rightarrow X_k} \leq C_M \mathcal{C}_k,$$

where we define $\Omega_M := \{z = x + iy \in \mathbb{C}, 0 < x \leq 1, |y| \leq M\}$ and we recall that \mathcal{C}_k is defined in (6.9). We argue by contradiction, assuming that there exists $y \in [-M, M]$ and a sequence (z_n) in Ω_M such that

$$z_n \rightarrow z := iy \quad \text{and} \quad \mathcal{C}_n^{-1} \|R_{\mathcal{L}_1}(z_n)\|_{\mathcal{B}(X_{n-j}, X_n)} \rightarrow \infty,$$

with $\mathcal{L}_1 := \Pi^\perp \mathcal{L}$, where for brevity we note $\Pi^\perp := I - \Pi$, despite the fact that Π is not an orthogonal projection. The last family of blowing up estimates means that there exist sequences $(\tilde{f}_n)_n$ and $(\tilde{g}_n)_n$ such that

$$\mathcal{C}_n^{-1} \|\tilde{f}_n\|_{X_n} \rightarrow \infty, \quad \|\tilde{g}_n\|_{X_{n-j}} = 1, \quad \tilde{f}_n = R_{\mathcal{L}_1}(z_n) \tilde{g}_n,$$

or, equivalently, that there exist $(f_n)_n$ in X_n and $(g_n)_n$ in X_{n-j} satisfying

$$\|f_n\|_{X_n} = 1, \quad \mathcal{C}_n \|g_n\|_{X_{n-j}} \rightarrow 0, \quad g_n = (\mathcal{L}_1 - z_n) f_n.$$

This in turn would imply that

$$(6.15) \quad R_{\mathcal{B}}(z_n) \mathcal{A} \Pi^\perp f_n + \Pi^\perp f_n - z_n R_{\mathcal{B}}(z_n) \Pi f_n = R_{\mathcal{B}}(z_n) g_n,$$

with

$$\|R_{\mathcal{B}}(z_n) g_n\|_{X_n} \leq \mathcal{C}_n \|g_n\|_{X_{n-j}} \rightarrow 0,$$

by using (6.9). Since $(f_n)_n$ is bounded in $X_n \subset X_\infty = L^2(e^{\kappa \langle x \rangle^\gamma})$, by weak compactness of this sequence in X_∞ , we find $f \in X_\infty$ and a subsequence denoted again by $(f_n)_n$ such that $f_n \rightharpoonup f$ weakly in X_∞ , and then $\mathcal{A} \Pi^\perp f_n \rightharpoonup \mathcal{A} \Pi^\perp f$ weakly in X_0 . Together with (6.11), we deduce that

$$(6.16) \quad R_{\mathcal{B}}(z_n) \mathcal{A} \Pi^\perp f_n \rightarrow R_{\mathcal{B}}(z) \mathcal{A} \Pi^\perp f \quad \text{strongly in } X_0.$$

Now, passing (weakly) to the limit in (6.15), we have

$$R_{\mathcal{B}}(z) \mathcal{A} \Pi^\perp f + \Pi^\perp f - z R_{\mathcal{B}}(z) \Pi f = 0.$$

We claim that $f \neq 0$. If not, we get from (6.16) that

$$R_{\mathcal{B}}(z_n) \mathcal{A} \Pi^\perp f_n \rightarrow 0 \quad \text{and} \quad \Pi f_n \rightarrow 0 \quad \text{strongly in } X_0$$

and then together with (6.15) that $\|\Pi^\perp f_n\|_{X_n} \rightarrow 0$. Thus we would have

$$1 = \|f_n\|_{X_n} \leq \|\Pi f_n\|_{X_0} + \|\Pi^\perp f_n\|_{X_n} \rightarrow 0,$$

which is a contradiction. As a consequence, we have exhibited an $f \in X_\infty \setminus \{0\}$ such that $(\mathcal{L}_1 - zI)f = 0$. This means that f is an eigenvector for $\mathcal{L}_1 = \Pi^\perp \mathcal{L}$ associated to an eigenvalue $z \in i\mathbb{R}$; however this is in contradiction with the fact that $\Sigma_P(\Pi^\perp \mathcal{L}) \cap i\mathbb{R} = \emptyset$. Thus the proof of (6.14) is complete.

Step 2. In this step we complete the proof of the Lemma. We begin by recalling that $\mathcal{L} = \mathcal{A} + \mathcal{B}$ and then we write

$$R_{\mathcal{L}}(z) = R_{\mathcal{B}}(z) - R_{\mathcal{L}}(z) \mathcal{A} R_{\mathcal{B}}(z)$$

and

$$R_{\mathcal{L}}(z)(1 - \mathcal{V}(z)) = R_{\mathcal{B}}(z) - R_{\mathcal{B}}(z) \mathcal{A} R_{\mathcal{B}}(z), \quad \text{where } \mathcal{V}(z) := (\mathcal{A} R_{\mathcal{B}}(z))^2.$$

Let \tilde{X}_k be defined as X_k but with a coefficient $\tilde{\kappa} > \kappa$ so that $\tilde{X}_k \subset \tilde{X}_\infty \subset X_0 \subset X_k$ with embedding constants uniformly bounded with respect to k . First we may fix M large enough such that $\|\mathcal{V}(z)\|_{\mathcal{B}(X_0)} \leq 1/2$ and $\|\mathcal{V}(z)\|_{\mathcal{B}(\tilde{X}_0)} \leq 1/2$, for any

$z = x + iy$, with $|y| \geq M$: this is indeed possible thanks to (6.10) by choosing $\ell_1 = \ell_2 = 1$. Next, we write the expansion

$$R_{\mathcal{L}}(z) = R_{\mathcal{B}}(z) - (R_{\mathcal{B}}(z) - R_{\mathcal{B}}(z) \mathcal{A} R_{\mathcal{B}}(z)) \left(\sum_{j=0}^{\infty} \mathcal{V}(z)^j \right) \mathcal{A} R_{\mathcal{B}}(z),$$

where the series converges normally in $\mathcal{B}(X_0)$ and in $\mathcal{B}(\tilde{X}_0)$. More precisely, for $z = x + iy$ and $|y| \geq M$ we have

$$\begin{aligned} \|R_{\mathcal{L}}(z)\|_{X_{k-j} \rightarrow X_k} &\leq \|R_{\mathcal{B}}(z)\|_{X_{k-j} \rightarrow X_k} \\ &\quad + \|(R_{\mathcal{B}} - R_{\mathcal{B}} \mathcal{A} R_{\mathcal{B}})(z)\|_{\tilde{X}_0 \rightarrow X_0} \left(\sum_{j=0}^{\infty} \|\mathcal{V}(z)\|_{\mathcal{B}(\tilde{X}_0)}^j \right) \|\mathcal{A} R_{\mathcal{B}}\|_{X_{k-j} \rightarrow \tilde{X}_0}, \\ (6.17) \quad &\leq \|R_{\mathcal{B}}(z)\|_{X_{k-j} \rightarrow X_k} + 2 \|(R_{\mathcal{B}} - R_{\mathcal{B}} \mathcal{A} R_{\mathcal{B}})(z)\|_{\tilde{X}_0 \rightarrow X_0} \|\mathcal{A} R_{\mathcal{B}}\|_{X_{k-j} \rightarrow \tilde{X}_0}. \end{aligned}$$

The right hand side of the above inequality being bounded by a constant $C C_k$, we conclude that

$$\sup_{|\Re z| \geq M} \|(I - \Pi) R_{\mathcal{L}}(z)\|_{X_{k-j} \rightarrow X_k} \leq C C_k.$$

Together with (6.14), we conclude the proof of (6.13) in the case when $\ell = 1$. For the general case $\ell \geq 1$, we argue similarly as we did in the proof of (6.10). \square

Theorem 6.7. *Let $\sigma_{\mathcal{L}}^* := 1/\lfloor 2/\gamma \rfloor$, where $\lfloor s \rfloor$ stands for the integer part of the real number s , and let $m_0(x) := \exp(\kappa \langle x \rangle^\gamma)$ and $0 < \kappa \gamma < 1/8$. Then for any $\sigma \in (0, \sigma_{\mathcal{L}}^*]$ and $\theta < 1$ there exist $\lambda > 0$ such that for all $t > 0$*

$$\|f(t) - \mathbb{M}(f_0) G\|_{L^2(m_0^\theta)} \lesssim \exp(-\lambda t^\sigma) \|f_0 - \mathbb{M}(f_0) G\|_{L^2(m_0)}.$$

Proof of Theorem 6.7. We write the representation formulas (taken from [28])

$$(6.18) \quad S_{\mathcal{L}}(t)f = \Pi f + \sum_{\ell=0}^5 (I - \Pi) S_{\mathcal{B}} * (\mathcal{A} S_{\mathcal{B}})^{(*\ell)}(t) f + \mathcal{T}(t) f,$$

where $\mathcal{T}(t) := \lim_{M \rightarrow \infty} \frac{i}{2\pi} \int_{a-iM}^{a+iM} e^{zt} (I - \Pi) \mathcal{R}_{\mathcal{L}}(z) (\mathcal{A} \mathcal{R}_{\mathcal{B}}(z))^6 dz,$

for any $f \in D(\mathcal{L})$, $t \geq 0$ and $a > 0$.

Thanks to Lemma 2.4, we know that each term $\|(I - \Pi) S_{\mathcal{B}} * (\mathcal{A} S_{\mathcal{B}})^{(*\ell)}(t) f\|$ in the above expression of $S_{\mathcal{L}}(t)$ is bounded by $\mathcal{O}(\Theta_{m_0}(t))$. In order to conclude, we have to estimate the last term.

We introduce the shorthands $\Phi_1 := R_{\mathcal{L}}(I - \Pi)$, $\Phi_\ell = \mathcal{A} R_{\mathcal{B}}$, for $2 \leq \ell \leq 7$, and we perform n integration by part in the formula giving $\mathcal{T}(t)$ to get

$$(6.19) \quad \mathcal{T}(t) = \frac{i}{2\pi} \frac{1}{t^n} \int_{a-i\infty}^{a+i\infty} e^{zt} \frac{d^n}{dz^n} \left(\prod_{i=1}^7 \Phi_i(z) \right) dz.$$

Using the fact that all the functions appearing in the integral are bounded on the imaginary axis, together with the resolvent identity

$$R_{\Lambda}^n(z) := \frac{d^n}{dz^n} R_{\Lambda}(z) = n! R_{\Lambda}(z)^n,$$

we find in $\mathcal{B}(X_0)$, thanks to Leibniz formula and for any $z = x + iy \in \mathbb{C}$ with $0 \leq y \leq 1$,

$$\begin{aligned} \left\| \frac{d^n}{dz^n} \left(\prod_{\ell=1}^7 \Phi_\ell(z) \right) \right\| &\leq 7^n \sup_{\alpha \in \mathbb{N}^n, |\alpha|=n} \left\| \prod_{\ell=1}^7 \frac{d^{\alpha_\ell}}{dz^{\alpha_\ell}} \Phi_\ell(z) \right\| \\ &\leq 7^n \sup_{\alpha \in \mathbb{N}^7, |\alpha|=n} \prod_{i=1}^7 \|\alpha_i! (I - \Pi) R_{\mathcal{L}}^{1+\alpha_i} \alpha_2! \mathcal{A} R_{\mathcal{B}}^{1+\alpha_2} \dots \alpha_7! \mathcal{A} R_{\mathcal{B}}^{1+\alpha_7}(z)\| \\ &\leq C^n (n!)^{1+j} \langle y \rangle^{-3/2}, \end{aligned}$$

where in the last step we have used Lemma 6.6 for some integer j which will be fixed later. Next, using the bound $n! \leq (Cn)^n$, we get

$$\|\mathcal{T}(t)\| \leq C^n n^{(1+j)n} t^{-n} \quad \forall t > 0, \forall k \geq 0.$$

For any $t \geq t^*$, where t^* is large enough and depends on j , we choose $n \in \mathbb{N}$ such that

$$t \geq 2Cn^{1+j} \geq t/2,$$

and we obtain

$$\|\mathcal{T}(t)\| \leq (Cn^{1+j}t^{-1})^n \leq (1/2)^n \leq (1/2)^{(t/4C)^{\frac{1}{j+1}}} \quad \forall t > 0.$$

As a consequence, with the choice $j := \lfloor [2(1-\gamma)/\gamma] + 1 \rfloor$, we have proved that for all $t \geq 0$ we have

$$\|\mathcal{T}(t)\| \leq \exp(-\lambda t^{\frac{1}{1+j}}),$$

which clearly ends the proof. \square

REFERENCES

- [1] AOKI, K., AND GOLSE, F. On the speed of approach to equilibrium for a collisionless gas. *Kinet. Relat. Models* 4, 1 (2011), 87–107.
- [2] ARENDT, W., GRABOSCH, A., GREINER, G., GROH, U., LOTZ, H. P., MOUSTAKAS, U., NAGEL, R., NEUBRANDER, F., AND SCHLOTTERBECK, U. *One-parameter semigroups of positive operators*, vol. 1184 of *Lecture Notes in Mathematics*. Springer-Verlag, Berlin, 1986.
- [3] BAKRY, D., BARTHE, F., CATTIAUX, P., AND GUILLIN, A. A simple proof of the Poincaré inequality for a large class of probability measures including the log-concave case. *Electron. Commun. Probab.* 13 (2008), 60–66.
- [4] BAKRY, D., CATTIAUX, P., AND GUILLIN, A. Rate of convergence for ergodic continuous Markov processes: Lyapunov versus Poincaré. *J. Funct. Anal.* 254, 3 (2008), 727–759.
- [5] BÁTKAI, A., ENGEL, K.-J., PRÜSS, J., AND SCHNAUBELT, R. Polynomial stability of operator semigroups. *Math. Nachr.* 279, 13-14 (2006), 1425–1440.
- [6] BATTY, C. J. K., AND DUYCKAERTS, T. Non-uniform stability for bounded semi-groups on Banach spaces. *J. Evol. Equ.* 8, 4 (2008), 765–780.
- [7] CAFLISCH, R. E. The Boltzmann equation with a soft potential. I. Linear, spatially-homogeneous. *Comm. Math. Phys.* 74, 1 (1980), 71–95.
- [8] CAFLISCH, R. E. The Boltzmann equation with a soft potential. II. Nonlinear, spatially-periodic. *Comm. Math. Phys.* 74, 2 (1980), 97–109.
- [9] CAÑIZO, J. A., AND MISCHLER, S. A semigroup proof of Harris' theorem. In progress.
- [10] CARRAPATOSO, K., AND MISCHLER, S. Landau equation for very soft and Coulomb potentials near Maxwellian. *Ann. PDE* 3, 1 (2017), 1–65.
- [11] CATTIAUX, P., GOZLAN, N., GUILLIN, A., AND ROBERTO, C. Functional inequalities for heavy tailed distributions and application to isoperimetry. *Electron. J. Probab.* 15 (2010), no. 13, 346–385.
- [12] DOUC, R., FORT, G., AND GUILLIN, A. Subgeometric rates of convergence of f -ergodic strong Markov processes. *Stochastic Process. Appl.* 119, 3 (2009), 897–923.

- [13] DUNFORD, N., AND SCHWARTZ, J. T. *Linear Operators. I. General Theory*. With the assistance of W. G. Bade and R. G. Bartle. Pure and Applied Mathematics, Vol. 7. Interscience Publishers, Inc., New York; Interscience Publishers, Ltd., London, 1958.
- [14] EBERLE, A., GUILLIN, A., AND ZIMMZE, R. Quantitative Harris type theorems for diffusions and McKean-Vlasov processes. *To appear in Transactions of the American Mathematical Society*.
- [15] ESCOBEDO, M., MISCHLER, S., AND RODRIGUEZ RICARD, M. On self-similarity and stationary problem for fragmentation and coagulation models. *Ann. Inst. H. Poincaré Anal. Non Linéaire* 22, 1 (2005), 99–125.
- [16] GAMBA, I. M., PANFEROV, V. AND VILLANI, C. On the Boltzmann equation for diffusively excited granular media. *Comm. Math. Phys.* 246, (2004), 503–541.
- [17] GUALDANI, M. P., MISCHLER, S., AND MOUHOT, C. Factorization of non-symmetric operators and exponential H -Theorem. *Mém. Soc. Math. Fr.* 153, (2017), 1-137.
- [18] GUO, Y. The Landau equation in a periodic box. *Comm. Math. Phys.* 231 (2002), 391–434.
- [19] HAIRER, M. Convergence of Markov Processes, (2016), available on <http://www.hairer.org/notes/Convergence.pdf>
- [20] KAVIAN, O. Remarks on the Kompaneets equation, a simplified model of the Fokker-Planck equation. In *Nonlinear partial differential equations and their applications. Collège de France Seminar, Vol. XIV (Paris, 1997/1998)*, vol. 31 of *Stud. Math. Appl.* North-Holland, Amsterdam, 2002, pp. 469–487.
- [21] KAVIAN, O., KERKYACHARIAN, G., AND ROYNETTE, B. Quelques remarques sur l’ultracontractivité. *J. Funct. Anal.* 111, 1 (1993), 155–196.
- [22] LIEB, E. H., AND LOSS, M. *Analysis*, second ed., vol. 14 of *Graduate Studies in Mathematics*. American Mathematical Society, Providence, RI, 2001.
- [23] LIGGETT, T. M. L_2 rates of convergence for attractive reversible nearest particle systems: the critical case. *Ann. Probab.* 19, 3 (1991), 935–959.
- [24] MICHEL, P., MISCHLER, S., AND PERTHAME, B. General relative entropy inequality: an illustration on growth models. *J. Math. Pures Appl. (9)* 84, 9 (2005), 1235–1260.
- [25] MISCHLER, S. Semigroups in Banach spaces - factorization approach for spectral analysis and asymptotic estimates. In preparation.
- [26] MISCHLER, S., AND MOUHOT, C. Exponential stability of slowly decaying solutions to the Kinetic-Fokker-Planck equation. *Arch. Ration. Mech. Anal.* 221, 2 (2016), 677–723.
- [27] MISCHLER, S., QUIÑINAO, C., AND TOUBOUL, J. On a kinetic FitzHugh-Nagumo model of neuronal network. *Comm. Math. Phys.* 342, 3 (2016), 1001–1042.
- [28] MISCHLER, S., AND SCHER, J. Spectral analysis of semigroups and growth-fragmentation equations. (*Ann. Inst. H. Poincaré Anal. Non Linéaire* 33, 3 (2016), 849–898.
- [29] MOUHOT, C. Rate of convergence to equilibrium for the spatially homogeneous Boltzmann equation with hard potentials. *Comm. Math. Phys.* 261, 3 (2006), 629–672.
- [30] NDAO, M. Convergence to equilibrium for the Fokker-Planck equation with a general force field. PhD thesis, Université de Versailles Saint-Quentin (2018).
- [31] RÖCKNER, M., AND WANG, F.-Y. Weak Poincaré inequalities and L^2 -convergence rates of Markov semigroups. *J. Funct. Anal.* 185, 2 (2001), 564–603.
- [32] TARTAR, L. *An introduction to Sobolev spaces and interpolation spaces*, vol. 3 of *Lecture Notes of the Unione Matematica Italiana*. Springer, Berlin; UMI, Bologna, 2007.
- [33] TOSCANI, G., AND VILLANI, C. On the trend to equilibrium for some dissipative systems with slowly increasing a priori bounds. *J. Statist. Phys.* 98, 5-6 (2000), 1279–1309.
- [34] VOIGT, J. A perturbation theorem for the essential spectral radius of strongly continuous semigroups. *Monatsh. Math.* 90, 2 (1980), 153–161.

(Otarred Kavian) UNIVERSITÉ PARIS-SACLAY; UVSQ & CNRS, UMR 8100; LABORATOIRE DE MATHÉMATIQUES DE VERSAILLES; 45 AVENUE DES ETATS UNIS; 78035 VERSAILLES CEDEX, FRANCE.
Email address: kavian@math.uvsq.fr

(Stéphane Mischler) UNIVERSITÉ PARIS-DAUPHINE, INSTITUT UNIVERSITAIRE DE FRANCE (IUF), PSL RESEARCH UNIVERSITY, CNRS, UMR [7534], CEREMADE, PLACE DU MARÉCHAL DE LATRE DE TASSIGNY 75775 PARIS CEDEX 16, FRANCE.
Email address: mischler@ceremade.dauphine.fr

(Mamadou Ndao) UNIVERSITÉ PARIS-SACLAY; UVSQ & CNRS, UMR 8100; LABORATOIRE DE MATHÉMATIQUES DE VERSAILLES; 45 AVENUE DES ETATS UNIS; 78035 VERSAILLES CEDEX, FRANCE.
Email address: `nda075019@yahoo.fr`