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## **Non-intrusive Diagnosis of a PEMFC**

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*Abstract***— Many reasons may cause malfunction of a proton exchange membrane fuel cell stack. Moreover, a single defect cell can affect the whole stack functioning, therefore the localization of this cell is crucial as well as identifying the root causes for such malfunctioning. The diagnosis tool has to be non intrusive and to be easy to replicate to a new fuel cell system. This paper explores paths to achieve such monitoring based on the use of the magnetic field induced by the fuel cell stack internal currents. Designing and using a three–dimensional simulation tool allows to establish a cause and effect relationship between fuel cell defects and specific magnetic signatures. This preliminary work enables to determine the required magnetic sensor accuracy, the sensor orientation and location as well as the experimental precautions for relevant measurements.** 

*Index Terms***— Proton exchange membrane fuel cell; Maxwell-Stefan diffusion; magnetic field; 3D PDE modeling, diagnosis.** 

#### I. AIM

Using polymer electrolyte membrane (PEM) fuel cells (FCs) is an attractive way to produce energy without greenhouse gas. A defect in the FC stack, such as drying, flooding, starvation, active layer deterioration, induces major negative effect on the operation of the overall energy conversion system (FC and its auxiliaries): performance decreases and may cause the FC system to shut down, with possible damages to the PEMFC or to one of the auxiliaries. Identification of the defect and understanding the physical processes involved are essential to diagnose accurately the reduced performance causes. The aim of this work is to compute the magnetic field created by a PEMFC in order to identify a specific signature in case of malfunctioning cell.

### II. PEMFC MODELING

To predict the FC performance, the unknowns,  $W_{H2}$ ,  $W_{H2O,a}$ , wH2O,c, wO2, wN2, pa, pc, φm, φ, and **v** have to be computed by solving the equations of mathematical physics that describe the overall cell; computational domain is given in Fig. 1.

It is assumed that:

- Three species oxygen, nitrogen and water are present on the cathode side, while hydrogen and water are considered in the anode side.
- The PEMFC is operating in steady state.
- Flow is considered laminar.
- There is no condensation in the cell. The single phase model described here is sufficient for modelling the mass transport in the porous anode and cathode up to moderately high current densities [1] [2].
- The FC temperature remains constant and homogeneous all over the cell.
- The membrane is gas-tight.
- The channels are approximated by straight channels.



#### A. *Mass and Momentum balances (coupled)*

In anode and cathode sides, the transport phenomena are similar, only the species differ. In channels, the velocity **v** and the pressure of the gas mixtures are described by Navier-Stokes equations, i.e. the conservation of momentum (1) and the conservation of mass known as mass continuity

TA BLEI NOMENCLATURE

ID 169 2

 equation (2). As GDL and CL are porous media, the Navier-Stokes equations have to be modified that leads to the socalled Brinkman equations [3] (3).



Fig. 1. Cell detail: zoom on the GDLs and the MEA.

In the channel, the following system has to be solved

$$
\begin{cases}\n\nabla \cdot (\rho_{a,c} \mathbf{v}_{a,c} \otimes \mathbf{v}_{a,c} - \mu_{a,c} \nabla \mathbf{v}_{a,c}) = \nabla \cdot [\mu_{a,c} (\nabla \mathbf{v}_{a,c})^T] \\
-\nabla \cdot (\rho_{a,c} \mathbf{v}_{a,c}) = 0\n\end{cases}
$$
\n(1)\n
$$
\nabla \cdot (\rho_{a,c} \mathbf{v}_{a,c}) = 0
$$
\n(2)

while in the GDLs and in the CLs the below-stated system of equations needs to be solved:

$$
\begin{cases}\n\nabla \left( \frac{\rho_{a,c}}{\varepsilon_{\text{GDL,CL}}} v_{a,c} \otimes v_{a,c} - \frac{\mu_{a,c}}{\varepsilon_{\text{GDL,CL}}} \nabla v_{a,c} \right) = \\
\nabla \left[ \frac{\mu_{a,c}}{\varepsilon_{\text{GDL,CL}}} (\nabla v_{a,c})^{\text{T}} \right] - \nabla \left( p_{a,c} + \frac{2}{3} \frac{\mu_{a,c}}{\varepsilon_{\text{GDL,CL}}} \nabla v_{a,c} \right) - \frac{\mu_{a,c}}{\kappa_{\text{GDL,Cl}}} = 0 \\
\nabla \left( \rho_{a,c} \mathbf{v}_{a,c} \right) = 0\n\end{cases} \tag{3}
$$

#### B. *Mass and specifies balance*

The mass and molar fractions obey the Maxwell-Stefan equations which describe the multi-species diffusion and convection in channels (5-6) and electrodes. Pair of species (i,j) interactions are characterized by the binary diffusion coefficient  $D_{ij}$  [4]: to account for the electrodes' porosity, the effective diffusion coefficient  $D_{ij}^{eff}$  (9) substitute the  $D_{ij}$  coefficient in relation (8).

$$
\nabla \left[ -\rho \cdot w_i \sum_{j=1}^{e} D_{ij} \left( \nabla x_j + (x_j - w_j) \frac{\nabla p_{a,c}}{p_{a,c}} \right) + w_i \rho v_{a,c} \right] = R_i
$$
\n
$$
\mathbf{N}_i = \left[ -\rho \cdot w_i \sum_{j=1}^{e} D_{ij} \left( \nabla x_j + (x_j - w_j) \frac{\nabla p_{a,c}}{p_{a,c}} \right) + \rho v_{a,c} w_i \right]
$$
\n
$$
(6)
$$

Where  $\rho$  is the mixture density and  $D_{ij}$  is defined by:

$$
\rho = \left(\sum_{i=1}^{e} x_i M_i\right) \frac{p_{a,c}}{RT}
$$
\n(7)

$$
D_{ij} = D_{ij}^0 (P^0, T^0) \frac{P^0}{T^0} \left(\frac{T}{T^0}\right)^{\frac{3}{2}}
$$
 (8)

In the GDLs and CLs into account, the effective diffusion coefficients are computed by the Bruggeman relation:

$$
D_{ij}^{\text{eff}} = D_{ij} \varepsilon_{\text{GDL,CL}}^{3/2}
$$
 (9)

Using the Onsager reciprocal relation [5], [6] also leads to the following relation:

$$
D_{ij}^{\text{eff}} = D_{ji}^{\text{eff}} \tag{10}
$$

In the continuity equation (5), species production and consumption are represented by a source term (water production) and a sink term (reactant consumption) respectively. The corresponding Ri terms is expressed in  $kg/(m^3.s)$ . Obviously  $R_i=0$  in flow channels.

$$
R_{O_2} = \frac{j_c}{4F} M_{O_2}
$$
 (11)

 $(1)$ 

$$
R_{H_2} = -\frac{j_a}{2F} M_{H_2}
$$
 (12)

$$
R_{H_2O} = -\frac{j_c}{2F} M_{H_2O}
$$
 (13)

Where  $j_a$  and  $j_c$  are the transfer current densities at CLs.

(2) TABLE II GEOMETRY PARAMETERS

.	
<b>Parameters</b>	Value
Active area	$100 \text{ cm}^2$
Anode GDL thickness	$380 \mu m$
Cathode GDL thickness	$380 \mu m$
Membrane thickness	$100 \mu m$
Anode catalyst layer thickness	$50 \mu m$
Cathode catalyst layer thickness	$50 \mu m$
<b>GDL</b> porosity	0.4
Catalyst layer porosity	0.3
Anode and cathode channel width	1 <sub>mm</sub>
Anode and cathode channel depth	1 mm
Cog height of bipolar plates	$0.9 \text{ mm}$
<b>GDL</b> conductivity	$220$ S/m
Membrane conductivity	$10 \text{ S/m}$

#### C. *Redox reaction and produced current*

Indeed, hydrogen oxidation leads to charges' displacements: a proton one occurs through the membrane inducing a ionic current  $j_{H+}$  while electrons travel are conducted by the electrodes resulting in an electronic current, je. Current conservation is expressed as follows:

$$
\nabla \cdot \mathbf{j}_{cell} = \nabla \cdot \mathbf{j}_{e} + \nabla \cdot \mathbf{j}_{H+} = 0 \tag{14}
$$

Ohm's law links current j<sub>i</sub> densities to potentials  $\varphi_i$ :

$$
\nabla \left( -\sigma_e^{\text{eff}} \nabla \varphi_e \right) = -j_{a,c} \tag{15}
$$

$$
\nabla \left( -\sigma_{\rm H+}^{\rm eff} \nabla \varphi_{\rm H+} \right) = \mathbf{j}_{\rm a,c} \tag{16}
$$

In  $(15)$  and  $(16)$ , the source terms j<sub>i</sub> derive from the electrochemical reactions described by Butler-Volmer equation.

$$
j_a = a_v j_o^{ref} \left(\frac{C_{H_2}}{C_{H_2}^{ref}}\right)^{1/2} \left[\exp\left(\frac{\alpha_a F}{RT}\right) \eta_{act} - \exp\left(-\frac{\alpha_c F}{RT}\right) \eta_{act}\right] \quad (17)
$$

$$
j_c = a_v j_0^{ref} \left( \frac{C_{O_2}}{C_{O_2}^{ref}} \right) \left[ exp \left( \frac{\alpha_a F}{RT} \right) \eta_{act} - exp \left( -\frac{\alpha_c F}{RT} \right) \eta_{act} \right]
$$
(18)

The reference exchange current  $j_o^{ref}$  depends on several parameters such as operating temperature and requires a large number of experiments [7-8] to be properly identified.  $\eta_{\text{act}}$  is the voltage drop between electrodes and electrolyte:

At the anode side 
$$
\eta_{\text{act}} = \varphi_e - \varphi_{H+}
$$
 (19)

At the cathode side  $\eta_{\text{act}} = \varphi_e - \varphi_{H+} - V_{oc}$  (20)

#### D. *Boundary conditions*

The boundary conditions have to be specified for all variables in each domain. In anode and cathode channels, inflows are given by their velocity and their inlet mixtures:

$$
v_{a,c} = v_{a,c,in} \tag{21}
$$

$$
w_a = w_{H_{2,in}} \tag{22}
$$

$$
\mathbf{w}_{\rm c} = \begin{cases} \mathbf{w}_{\mathbf{O}_{2,\rm in}} \\ \mathbf{w}_{\rm H_2O_{in}} \end{cases}
$$
 (23)

The outlet boundaries are given by:

$$
p_{a,c} = p_{out} \tag{24}
$$

and

$$
-\mathbf{n}\left[\rho.w_i\sum_{j=1}^{e}D_{ij}\left(\nabla x_j + (x_j - w_j)\frac{\nabla p_{a,c}}{p_{a,c}}\right)\right] = 0
$$
\n(25)

Elsewhere no slip conditions apply and no flux go out. Thus:

$$
-\mathbf{n}.\mathbf{N}_{i} = 0 \tag{26}
$$

The potential is fixed at gas channel-gas diffusion layer interfaces:

At the cathode side: 
$$
\varphi_e = V_{cell}
$$
 (27)

At the anode side: 
$$
\varphi_e = 0
$$
 (28)

At the membrane: 
$$
\varphi_{H+} = 0
$$
 (29)

At the top and at the bottom and on the back and on the front of the FC unique cell, no current goes out. Hence:

$$
-\mathbf{n}.\mathbf{j}_e = 0\tag{30}
$$

$$
-\mathbf{n}.\mathbf{j}_{H+} = 0 \tag{31}
$$

#### E. *Magnetic formulation*

As discussed earlier, the main aim of this study is to propose a set of guiding principles that will support the implementation of non intrusive fuel cell malfunctioning detection. A default in the heart of the fuel cell might modify the current density repartition and consequently the resulting magnetic field. After computing the fuel cell electrical operating point, the magnetic field is calculated by means of the magnetic vector potential (32).

In this part, the studied domain is the single fuel cell  $\Omega_{FC}$ and the air around it  $\Omega$ <sub>a</sub>. The air domain dimensions are chosen much larger than the cell ones.

$$
\Delta A - \mu j = 0 \quad \text{in } \Omega_a \cup \Omega_{FC} \tag{32}
$$

To ensure the uniqueness of the solution, a Coulomb gauge (33) has to be coupled to Poisson equation (32).

$$
\nabla \mathbf{A} = 0 \tag{33}
$$

#### F. *Boundary conditions*

The boundary surface is denoted by  $\Gamma_{air}$ . Zero Dirichlet boundary conditions can be applied.

$$
\mathbf{A} = \mathbf{0} \text{ on } \Gamma_{\text{air}} \tag{34}
$$

To obtain the PEMFC operating point (V, I), 3D coupling PDE, describing the mixture gas diffusion and the redox reactions, have to be solved under a given operating conditions. Subsequently, the resulting magnetic field is computed by means of MVP. All the equations have been implemented in FEM software.

#### III. SOME RESULTS

### A. *Defect cathode channel*

The considered fuel cell features a 100 cm<sup>2</sup> active area. As one of its cathode channels is blocked, O<sub>2</sub> does no more diffuse correctly (Fig. 2) and consequently the electrochemical reaction as well as the current distribution are locally disturbed; finally the resulting magnetic field is modified (Fig. 3 and 4).

#### B. *Defect anode channel*

The same default is simulated at the anode side; a stopper at anode channel is located at the top or the middle of the single fuel cell as shown in Fig 5.

Indeed, due to the stopper, hydrogen diffuses hardly in the default channel. However for a specified current of 27A, as shown in Fig. 5, the same default at anode side does not alter significantly the magnetic field distribution compared to cathode side.

TABLE III OPERATING CONDITIONS

<b>Parameters</b>	Value
Fuel cell temperature	60 °C
Hydrogen flow rate	$624 \text{ mL/min}$
Air flow rate	$1.56$ L/min
Inlet $H_2$ mass fraction (anode)	0.743
Inlet $O_2$ mass fraction (cathode)	0.228
Inlet $H_2O$ mass fraction (cathode)	0.023
Anode inlet flow velocity	$0.2 \text{ m/s}$
Cathode inlet flow velocity	$0.5 \text{ m/s}$
Anode viscosity	$1.19\ 10^{-5}$ Pa/s
Cathode viscosity	2.46 $10^{-5}$ Pa/s
Porous electrode permeability	$1.18 \times 10^{-11}$ m <sup>2</sup>
GDL permeability	$2.36 \times 10^{-12}$ m <sup>2</sup>
$H_2-H_2O$ binary $D_{ii}$	$1.055$ 10 <sup>-4</sup> m <sup>2</sup> /s
$O_2$ -N <sub>2</sub> binary $D_{ii}$	$2.751\ 10^{-5}$ m <sup>2</sup> /s
$N_2-H_2O$ Binary $D_{ii}$	$2.951\,10^{-5}$ m/s
$O_2$ -H <sub>2</sub> O binary $D_{ii}$	$3.233\ 10^{-5}$ m/s
Fuel cell voltage	From 0.9 V to $0.6$ V



Fig. 2. O2 diffusion: (a) healthy cell and (b) defective cathode channel

 $(a)$  $(b)$ 

Fig. 3. Magnetic induction tangential component  $B<sub>x</sub>$  at 0.6 V-31A.: xz view of (a) a healthy cell and (b) a defective cathode channel cell



Fig. 4. Comparison between the tangential component of the magnetic induction for a healthy and defective cathode channel



Fig. 5. Comparison between the tangential component of the magnetic induction for a healthy and defective anode channel

#### IV. CONCLUSIONS

This study has allowed to build a 3-dimensional computer model of a fuel cell whose originality is to be able to compute the induced magnetic field inside and around the fuel cell. This model has been used in different steady state scenarios.

Among other things, it permits to compare the magnetic field maps of a good operating cell and a defect one. The final purpose is obviously to determine which defect can be easily detected by magnetic measurement means and understand the precautions to be taken to achieve relevant measurements and diagnosis.

At a few tens of percent of rated current (a few amps to a few tens of amps for the studied cell), a defect at the cathode can be detected by measuring the tangential component of B with a sensor located closed to the cell. By contrast, a defect placed at the anode side produces a weak effect on the magnetic field B; it would be difficult to achieve a relevant diagnosis of such situation.

Experimental facility has to be led to demonstrate the scientific and technical feasibility of detecting a cell default among a fuel cell stack. Specifically, the present study shows that, regarding the rated cell current density (approximately 0.5 to 0.7 A.cm-2), the resulted magnetic field does not exceed the  $10^{-6}$  T range which can be challenging to measure especially considering that FC system auxiliaries may distort the measuring results. It is therefore essential to choose a Hall probe featuring a high sensitivity (in the order of mG) with relatively small active area due to the thickness of a single fuel cell.

According to all these considerations, the experimental setup has to be carried out so that the two current collectors do not disturb the measurement of the magnetic field induced by the fuel cell. One way to do so is to shield magnetically the fuel cell. Another promising way consists in correcting the measurements using the computation of auxiliaries' noises.

In sum, the simulation results allow to check the feasibility of diagnosis by means of magnetic field measurement and size the test bench and the associated sensors. This experimental stage is a key step and is the next phase of this work.

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