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# Block-Based Adaptive Mesh Refinement scheme using numerical density of entropy production for three-dimensional two-fluid flows

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## Abstract

In this work, we present a fast and parallel finite volume scheme on unstructured meshes applied to complex fluid flow. The mathematical model is based on a three-dimensional compressible low Mach two-phase flows model, combined with a linearized "artificial pressure" law. This hyperbolic system of conservation laws allows an explicit scheme, improved by a block-based adaptive mesh refinement scheme. Following a previous one-dimensional work, the useful numerical density of entropy production is used as mesh refinement criterion. Moreover, the computational time is preserved using a local time stepping method. Finally, we show through several test cases the efficiency of the present scheme on two and three-dimensional dam-break problems over an obstacle.

**Keywords:** low Mach model, dam-break problem, two-phase flow, artificial compressibility, finite volume, block-based mesh refinement, entropy production, local time stepping.

## 1 Introduction

The understanding of wave hydrodynamics is of primary interest for countless ocean and naval engineering applications: dynamics of ships and floating structures, stability of offshore structures, mass, momentum and energy fluxes between ocean and atmosphere, coastal erosion and submersion processes, etc. Despite significant progress in numerical, experimental and theoretical works during the last decades, the wave dynamics remains a fairly open research field. Many difficulties arise when trying to describe the violent transformations observed during wave breaking and/or impact on rigid structures. The involved physical processes, such as splash-ups or gas pockets entrapment, are quite complex and can hardly be characterized by field or laboratory experiments or analytical approaches. Boosted by the continuous improvement of computer technology, the Computational Fluid Dynamics (CFD) simulations are expected to provide a unique insight, in terms of spatio-temporal resolution and controlled conditions, into such wave dynamics. Therefore, the numerical simulation of breaking and impacting waves is both an attractive research topic and a challenging task for coastal and environmental engineering.

The most comprehensive numerical approach, in terms of physical relevancy, would be the solving of the full set of Navier-Stokes (NS) equations for both air and water phases down to the Kolmogorov dissipation scale, together with non-linear dynamic and kinematic free surface boundary conditions and other boundary conditions to represent solid surfaces or open boundaries. Such an ambition is still out of reach for now in terms of computing power. Consequently, most of the studies based on the NS equations are generally performed in the two dimensional case (see *e.g.* [19, 18]). Even if impressive results have been recently obtained using finite volume scheme [24], Lattice-Boltzmann model [6], finite element model [8, 15] or smoothed particle hydrodynamics model [5], three-dimensional simulations still demand a significant effort

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in software development and mesh refinement technique with the use of powerful computers (see *e.g.* [20, 11, 1, 30]). Turbulence models are often included in NS wave models whereas, despite of its importance in wave breaking and impact problems, the behaviour of the air-water micro-bubbles mixture is hardly taken into account [26].

Due to their much smaller requirements in cpu-time, simplified models based on potential theory or shallow-water equations are also widely used in research and engineering applications (see *e.g.* [14, 9, 28]). Unfortunately, these models are not able to accurately describe the powerful rotational two-phases processes occurring during wave breaking or impact, such as splash-up, overturning and free surface reconnection.

The overall objective of our research work is to propose an "intermediate" three-dimensional model aiming, on one hand, to be physically relevant in the context of highly dynamical and aerated flows associated with breaking and impacting waves and, on the other hand, to be less cpu-time consuming than the full three-dimensional Navier-Stokes equations. The presented model has been designed and continuously improved to this end [17, 13, 27], keeping in mind that an accurate capture of interface dynamics is a key issue for simulating wave transformation near breaking and after impacting. The assumption is made that, at least for the early stages, the physics of impacting flows is dominated by pressure forces and large scale overturning motions rather than small-scale friction and dissipation processes. Viscous effects are thus neglected and the two-phase flows model can be based on the compressible Euler equations. An artificial linearized pressure law is used to compute low Mach flows (see for instance [?]) with a preconditioned physical pressure coefficient avoiding too constraining CFL conditions. The obtained system is hyperbolic and, under some assumptions, the Riemann problem is well-posed and can be solved using explicit parallel Godunov finite volume solver (as described in [13]). The physical relevancy of this compressible two phase flow model has been successively tested on experimental and numerical test cases [17], improved by the use of an isothermal model [13] and recently validated on breaking wave problem over a non flat bottom with and without macro-roughness [27].

The purpose of this paper is to present the implementation and to test the efficiency of a new block-based adaptive mesh refinement (BB-AMR) combined with an efficient parallel algorithm. Moreover, based on recent developments in one-dimensional case [10], the use of the numerical density of entropy production as mesh refinement criterion is extended to our three-dimensional model. Following [10], the time step constrain imposed by the mesh refinement is outweighed by the implementation a local time stepping method together with first and second order Adams-Bashforth time integration schemes.

The first section of the paper is dedicated to the presentation of the model governing equations (section 2.1) and the finite volume approximation (section 2.2). A particular attention is paid on the new BB-AMR scheme and the related adaptation of the space grid using the numerical density of entropy production. Combining both adaptive mesh refinement and local time-stepping, we construct a three-dimensional parallel solver on unstructured meshes. The model confrontation with experiments in both 2D and 3D dam-break problems is presented in section 3. Finally, concluding remarks and prospects are given in section 4.

## 2 The mathematical and the numerical multi-fluid model

The hydrodynamics of breaking and impacting waves is extremely complex. As exposed in the introductory section, the most relevant approach would be to solve the full set of Navier-Stokes equations. Applied to three-dimensional cases and accounting for the time and space scales required by turbulence description, such computation is virtually inaccessible. The definition of a simplified mathematical and numerical model able to properly describe the physics of breaking and impacting waves at a decent cpu-time cost is a challenging problem which is the core of our research work. The purpose of this section is to briefly present the mathematical model with governing equations and finite volume approximation already described and validated in [17], [13] and [27], and to detail the implementation of both the time integration and the three-dimensional extension of the adaptive multi-scale scheme introduced in [10] on unstructured meshed using a new BB-AMR.

## 2.1 Governing equation

We consider a compressible two-fluids flows problem in three space dimension. Viscosity, surface tension and heat conduction are neglected. The incompressibility condition is relaxed using a low Mach approach in order to lead to an hyperbolic system of conservation laws. Thus, based on [13], the following three-dimensional isothermal hyperbolic and compressible Euler equations system is applied to a mixture fluid of air and water:

$$\begin{cases} \frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{u}) & = 0 \\ \frac{\partial (\rho \mathbf{u})}{\partial t} + \nabla \cdot (\rho \mathbf{u} \otimes \mathbf{u}) + \nabla p & = \rho \mathbf{g} \end{cases} \quad (1)$$

where the unknowns depend on spatial coordinates  $(x, y, z)$  and time  $t$ . The unknowns are the density  $\rho$ , the three components of the velocity  $\mathbf{u} = (u, v, w)$ , the pressure  $p$ . Here,  $\mathbf{g}$  stands for gravitational acceleration.

Air and water fractions within the mixture are defined by the volume fraction function  $\varphi \in [0, 1]$  ( $\varphi = 0$  in the water, and  $\varphi = 1$  in the air). With this definition of  $\varphi$ , the pressure of the two-phase flow problem is a function of the density  $\rho$  and the volume fraction  $\varphi$ , where  $\varphi$  solves the following non-conservative transport equation:

$$\frac{\partial \varphi}{\partial t} + \mathbf{u} \cdot \nabla \varphi = 0. \quad (2)$$

It is usually admitted that a flow is incompressible if the Mach number  $M = \|\mathbf{u}\|/c$  is lower than 0.1 ( $c$  is the sound speed), keeping in mind that the real (physical) Mach number is generally much smaller (of the order of  $1/400 \sim 1/1600$ ). In particular, this is constraining for explicit finite volume solver in which the time step  $\Delta t$  needs to satisfy a CFL (Courant-Friedrichs-Lewy) condition. Note also that the numerical scheme efficiency is expected to decrease with the Mach number. Therefore, an artificial pressure law (isothermal equation of state) is used to close the system :

$$p = c_0^2(\rho - (\varphi \rho_A + (1 - \varphi)\rho_W)) + p_0. \quad (3)$$

In this expression,  $\rho_A$  and  $\rho_W$  stand for air and water densities, respectively,  $c_0$  is the artificial sound speed (defined below) and  $p_0$  a reference pressure. For further details about the EOS choice, the reader is referred to [13]. The value of  $c_0$  is chosen as a compromise between the limits of compressible effects, the rate of numerical diffusion and a reasonable CFL constraint. The optimal choice for the artificial sound speed is still an open problem and would require a large amount of comparative test cases in a wide range of hydrodynamical conditions. In the present context, i.e. for flow velocity of the order of 1 m/s, an optimized value  $c_0 = 20\text{m/s}$  is used. Further improvements such as variable artificial sound speed can be considered. Finally, it is emphasized that in the boundary mixture region  $0 < \varphi < 1$  related to numerical diffusion processes, the proposed pressure law has no physical meanings.

For the ease of reference, equations (1) are written into a more compact form :

$$\frac{\partial \mathbf{w}(t)}{\partial t} + \nabla \cdot \mathbf{f}(t, \mathbf{w}) = \mathbf{G} \quad (4)$$

where  $\mathbf{w}$ ,  $\mathbf{f}$ ,  $\mathbf{G}$  stands respectively for conservative variables, flux and source:

$$\mathbf{w} = \begin{pmatrix} \rho \\ \rho \mathbf{u} \end{pmatrix}, \quad \mathbf{f} = \begin{pmatrix} \rho \mathbf{u} \\ \rho \mathbf{u} \otimes \mathbf{u} + pI \end{pmatrix} \quad \mathbf{G} = \begin{pmatrix} 0 \\ \rho \mathbf{g} \end{pmatrix}.$$

Here  $I$  stands for the  $3 \times 3$  identity matrix.

## 2.2 Finite volume approximation

In this section, we recall the well-known semi-discrete finite volume approximation of eqs. (4). The system is solved in  $\mathbb{R}^3$  and the source term is omitted for the sake of simplicity. The computational domain  $\Omega \subset \mathbb{R}^3$

is split into a set of control volumes, also referred as cells,  $\Omega = \cup_k C_k$ . We start by integrating Equations (4), by means of Green formula, on a cell  $C_k$ :

$$\int_{C_k} \frac{\partial \mathbf{w}(t)}{\partial t} + \sum_a \int_{\partial C_{k/a}} \mathbf{f}(t, \mathbf{w}) \cdot \mathbf{n}_{k/a} = 0 \quad (5)$$

where  $\mathbf{n}_{k/a}$  denotes the unit normal vector on the boundary  $\partial C_{k/a}$ . Then, setting

$$\mathbf{w}_k(t) \approx \frac{1}{|C_k|} \int_{C_k} \mathbf{w}(\mathbf{x}, t) d\Omega$$

where  $|C_k|$  stands for the volume of the cell  $C_k$ , we approximate (5) by

$$\frac{\partial \mathbf{w}_k(t)}{\partial t} + \sum_a \mathbf{F}(\mathbf{w}_k(t), \mathbf{w}_a(t); \mathbf{n}_{k/a}) = 0 \quad (6)$$

where the numerical flux

$$\mathbf{F}(\mathbf{w}_k(t), \mathbf{w}_a(t); \mathbf{n}_{k/a}) \approx \frac{1}{|C_{k/a}|} \int_{\partial C_{k/a}} \mathbf{f}(t, \mathbf{w}) \cdot \mathbf{n}_{k/a} ds$$

is defined by the Godunov solver. More precisely, the numerical flux  $\mathbf{F}(\mathbf{w}_k(t), \mathbf{w}_a(t); \mathbf{n}_{k/a}) := \mathbf{f}(w_{k/a}^*(t))$ , noted for the sake of simplicity  $\mathbf{F}_{k/a}(t)$ , is computed with the exact solution of the Riemann problem at the interface  $k/a$ :  $w^* = \mathfrak{R}(0, \mathbf{w}_k, \mathbf{w}_a)$  (for further details, see for instance, [29] or [13]).

Pressure oscillations observed in the multi-fluid case are avoided using the Abgrall's method [2] leading to a non-conservative discretisation of Equation (2). It is computed through the contact discontinuity velocity of the Riemann problems  $u^*$  solved at the interface  $k/a$ . As a consequence, the numerical flux for Equation (2) simply reads  $\min(u^*, 0)(\phi_k - \phi_a)$ .

Alternatively, a full conservative scheme with another pressure law can be used [3] but the less time consuming approach of Abgrall is retained here. The accurate interface capture is ensured by the use of a second order approximation (MUSCL) with Barth limiter. Further details can be found in [13].

### 2.3 Time integration and local time stepping method

This section is dedicated to the presentation of the principles of the Adams-Bashforth time integration and its implementation in our model. This approach has been shown to significantly decrease the computational cost compared to classical Runge-Kutta time integration which needs intermediate computations. The Runge-Kutta approach is deeply related to the size of the smallest control volume and therefore the time step is severely constrained when the mesh is refined. By contrast, even if the Adams-Bashforth is known to be less stable and less accurate, it can be easily handled in the framework of local time stepping (as described below) to save computational time (see for instance, [4] or [10]).

Equations (6) are first integrated over the time interval  $[0, T]$  subdivided such that  $t_{n+1} = t_n + \delta t_n$  with  $\delta t_n$  the time step.

The Adams-Bashforth method of order  $m$  consists in replacing the numerical flux of eq. (6) by a Lagrange polynomial interpolation of the same order [16]. Note that this explicit approximation is built with the fluxes previously computed and stored and then the integration of eq. (4), in the case of interest i.e. the second order Adams-Bashforth, leads to:

$$\mathbf{w}_k(t_{n+1}) = \mathbf{w}_k(t_n) - \sum_a \left( \frac{\delta t_n}{h_k} \mathbf{F}_k(t_n) - \frac{\delta t_n^2}{2\delta t_{n-1} h_k} (\mathbf{F}_k(t_n) - \mathbf{F}_k(t_{n-1})) \right) \quad (7)$$

where  $h_k = \frac{|C_k|}{\sum_a |\partial C_{k/a}|}$ . The main advantage of the Adams-Bashforth method is to avoid the computation of intermediate fluxes and thus facilitate a local time stepping approach (see [10] for details).

## 2.4 Mesh refinement criterion and BB-AMR scheme

### 2.4.1 Entropy production as a mesh refinement parameter

As pointed out in [10], many works are based on *a posteriori* error estimates, which are constructed from mathematical arguments for mesh refinement [34]. But paradoxically, to our knowledge, very few works use a refinement criterion based on physical consideration, such as entropy [12]. Based on this latter as a mesh refinement parameter, [10] proposed an efficient adaptive numerical scheme. The aim of the present work is to detail the extension of this scheme to three-dimensional cases and to demonstrate its efficiency.

It is well-known that solving equations system (4) with high accuracy is a challenging problem since solutions can and will breakdown at a finite time, even if the initial data are smooth, and develop complex structure (shock wave interactions). In such a situation, the uniqueness of the (weak) solution is lost and is recovered by completing the system (4) with an entropy inequality of the form:

$$\frac{\partial s(\mathbf{w})}{\partial t} + \nabla \cdot \psi(s(\mathbf{w})) \leq 0, \quad (8)$$

where  $(s, \psi)$  stands for a convex entropy-entropy flux pair satisfying the relation

$$(\nabla_w \psi(s(\mathbf{w})))^t = (\nabla_w s(\mathbf{w}))^t D_w f(\mathbf{w}).$$

This inequality in eq. (8) is used to select the physical relevant solution. Moreover, the entropy satisfies a conservation equation only in regions where the solution is smooth and an inequality when the solution develops shocks. Therefore, the entropy production can be seen as a "smoothness indicator" at the numerical level.

To this end and to be consistent with the discretisation of eq. (7), a discrete version of the entropy production, called *numerical density of entropy production*, is defined as follows:

$$S_k^n = \frac{s_k^{n+1} - s_k^n}{\Delta t_n} + \sum_a \left( \frac{\delta t_n}{h_k} \psi_{k/a}(t_n) + \frac{\delta t_n^2}{2\delta t_{n-1} h_k} (\psi_{k/a}(t_n) - \psi_{k/a}(t_{n-1})) \right) \cdot \mathbf{n}_{k/a}. \quad (9)$$

In this expression,  $\psi_{k/a}$  represents the entropy flux calculated from the resolution of the Riemann problem at the interface of cells  $k$  and  $a$ .

From a theoretical viewpoint, except during a shock, the numerical density of entropy production given by (9) is zero. But, at the numerical level, one can observe that it does not vanish and hence can be used as mesh refinement criterion. More precisely, [10] have shown that, in the case of one-dimensional gas dynamics equations for ideal gas, the support of relative error coincides with the support of numerical density of entropy production. As a consequence, the mesh is automatically refined only in regions where errors are identified.

The local entropy production must be compared to the total entropy  $\bar{S} = \frac{1}{|\Omega|} \sum_k S_k^n$ . Two coefficients  $0 \leq \alpha_{\min} \leq \alpha_{\max} \leq 1$  are thus defined to determine the ratio of numerical production of entropy leading to mesh refinement or mesh coarsening.

For each cell  $C_k$ :

- if  $S_k^n > \bar{S} \alpha_{\max}$ , the mesh is refined and split and,
- if  $S_k^n < \bar{S} \alpha_{\min}$  the mesh is coarsened whenever it is possible following the rule defined hereinafter.

**Remark 2.1** *The threshold parameters  $\alpha_{\max}$  and  $\alpha_{\min}$  are determined empirically for each problem to get a good compromise between computational cost and accuracy as done previously in the 1D framework [10]. More precisely,  $\alpha_{\min}$  and  $\alpha_{\max}$  allow to set a percentage of mesh refinement and mesh coarsening with respect to the quantity  $\bar{S}$ . It is not surprising that these settings will deteriorate or improve the accuracy of the numerical solution. For instance, the more  $\alpha_{\min}$  and  $\alpha_{\max}$  are small, the more accurate are the results to the expense of the computational time.*

Finally, in the present case, for the two-fluid model, the expression of entropy and entropy flux are

$$s = \frac{1}{2}\rho\mathbf{u}^2 + c_0^2\rho\ln\rho - c_0^2(\rho_W - \rho_A)\varphi, \quad \psi = \left(\frac{1}{2}\rho\mathbf{u}^2 + c_0^2\rho(\ln\rho + 1)\right)\mathbf{u}.$$

#### 2.4.2 BB-AMR scheme

Keeping in mind the following requirements: parallel treatment, fast computation and hierarchical grid (mesh level cannot exceed 2 for stability reasons as pointed out in [10]), the main and challenging difficulty is to develop a suitable mesh refinement tool which is related on the design of software data structures and numerical scheme.

Moreover, as pointed out by several authors, the extension toward three-dimensional case is not an easy task. Interesting works have been presented for 2D Cartesian grid or quad-tree [7, 32, 25, 33], octree for 3D simulations [23, 11], and anisotropic AMR [8, 15].

The present study aims to generalize the scheme proposed by [10] with a finite volume solver described in the previous section. This naturally leads to quad-tree meshing in two dimensional problem and octree meshing in three dimensional problem.

The presence of a complex moving air-water interface implies to re-mesh at each time step, which is clearly time consuming. Keeping in mind our willingness to find a relevant compromise between the contradictory aims of physical relevancy and numerical performance, we introduce a Cartesian block-based mesh approach (somehow like in [31, 35]). The computational domain is split in several blocks. In practice, each block corresponds to the initial unstructured mesh composed of hexahedral cells. According to the mesh refinement procedure (with respect to  $S_k^n$ ,  $\bar{S}$ , the mesh refinement level  $N$ ), each block is sub-divided in a Cartesian way ( $2^{N-1}n_x, 2^{N-1}n_y, 2^{N-1}n_z$ ) where  $(n_x, n_y, n_z)$  stands for the initial block discretisation.

As done in [10], for each refined cells (or blocks), averaged values at time  $t^n$  are projected on each sub-cell and fluxes are computed as simply as possible to avoid heavy computation. Consequently, the level of two adjacent blocks cannot exceeds 2 to avoid oscillations (see [10] for further details). Thus, the interface defining two blocks is most of the time is a non conforming one (as displayed in figure 2(b)).

Next, in order to spare proportionally the balanced distribution of cpu load, the cells of each block are re-distributed in a fixed number of domains according to the Cuthill-McKee numbering. The number of domain being fixed, each domain are loaded in a given MPI process. The procedure is certainly not optimal but the synchronization time between each domain are almost the same, hence overall more efficient. Finally, the re-numbering and re-meshing being expensive, the mesh is kept constant on a time interval, called AMR time-step, given by *the smallest block* (and not by the smallest cell) and the maximum velocity. This procedure provides a good compromise between computational time and accuracy.

Finally, let us illustrate the algorithm to define unstructured meshes using the BB-AMR scheme in the two-dimensional case (*i.e.*  $n_z = 1$  always). During the refinement process, the number of blocks and domains are kept always constant. As displayed in figure 1(a), we consider an initial mesh composed of 9 blocks, assumed to be of level 1 (the coarsest level with  $n_x = n_y = 1$ ), and decomposed on 3 domains according to a Cuthill-McKee numbering of the blocks (3 + 3 + 3 internal cells). For example the first domain (blue) is composed of three internal cells and three adjacent "ghost" cells in order to share to send/receive information between domains. After some computations, at the AMR-time  $T_1$ , let us suppose that blocks 1 and 4 have to be refined. As displayed on figure 1(b), the blocks distribution into the domains has been modified (5 + 5 + 5 internal cells). One can observed that the new mesh is a non-conforming mesh. For example, the cell belonging to the block 2 is bordered by 7 faces. If, at the AMR-time  $T_2$ , only the block 1 exceeds the mesh refinement criterion, it has to be refined. But, as the level of two adjacent blocks cannot exceeds 2 in order to avoid oscillations, the block 2 must also be refined. As displayed in figure 1(c), the blocks distribution has been modified (16 + 7 + 7 internal cells). The distribution is not well balanced because we distribute the blocks and not the cells for sake of simplicity. This drawback is reduced as the number of blocks is much bigger than the number of domains.

The three dimensional case is dealt in the same way. For illustration purpose, figure 2 shows an example of 3D block-based mesh decomposition (64 + 39 + 39 internal cells) with 3 domains and 27 blocks, where the

first block in the corner is of mesh refinement level 3.

### 3 Numerical results

The model is here tested with classical two and three dimensional test cases on dam-break problems over an obstacle. The 2D case is used to validate the mesh decomposition strategy while the 3D case quantitatively tests the overall performance of the presented model with and without BB-AMR scheme confronted to experimental and state-of-the-art numerical results.

#### 3.1 A 2D dam-break problem

This test case was initially introduced to study a particle method for simulating incompressible viscous flow with a particular focus on water dispersion and droplets after wave impact over an obstacle [21]. The experimental data obtained by [21] is here used to test the ability of our adaptive strategy to represent the complex structure of the air-water interface.

The test case configuration is shown in Fig. 3. The domain size is  $4L \times 2L$  with  $L = 14.6$  cm. A water column ( $2L$  high and  $L$  wide) and a rigid obstacle ( $2h$  high and  $h$  wide, with  $h=2.4$ cm) are initially located at the left and center of the domain, respectively. Air and water are at rest at the start of the experiment. The selected numerical parameters are  $CFL = 0.8$ , simulation time = 1.5 (s), 321 blocks, 120 cores, 321 domains or MPI process, 5 as the maximum level of mesh refinement,  $\alpha_{\max}$  (mesh refinement parameter) = 0.2,  $\alpha_{\min}$  (mesh coarsening parameter) = 0.02. Slip boundary conditions are imposed with no penetration. The simulation time is divided in 300 AMR time intervals.

Figure 4 represents the comparison between experimental and numerical snapshots during the dam-break process on the left and right columns, respectively. The selected times are  $t = 0.2, 0.3, 0.4, 0.5$  and 1s. At  $t = 0.2$ s, the wave has already impacted the obstacle and a thick water jet is ejected. At  $t = 0.3$ s and 0.4s, the jet extends and hits the right wall, inducing entrapment of a large air pocket which remains significant during the whole backwash and retreat process observed at  $t = 0.5$  and 1s.

The macroscopic features of such complex hydrodynamics, such as impact times against the obstacle and wall or jet shape and thickness, are rather well reproduced by the model. Moreover, even though our model was not designed to capture small scale interface splitting and reconnection processes, some of droplets or packets of droplets are captured as displayed on the right column of figure 4 at  $t = 0.4, 0.5$ s and 1s. Increasing the level of refinement, leading to more cells, will improve the accuracy of the numerical description of such processes, but this is not the aim of the present work.

The initial configuration for the numerical test case is shown in Fig. 5. Five levels of mesh refinement are used (see Fig. 5, right part). The mesh is initially denser in the areas of interest, i.e. the water phase and around the interface and then progressively coarsens in the air phase when moving away from the water column. Figures 6 and 7 represent the behavior of the dynamic mesh driven by the numerical density of entropy production at time  $t = 0.2$ s and 0.4s. The mesh is observed to remain and/or become the finest in the most active regions of the flow, i.e. interfacial, impact or sheared areas. This overall compared evolution of the collapsing water column with experiments demonstrates the relevancy of our refinement criterion, in particular when tracking the air-water interface. Note that the number of cells varies from 70 000 and 100 000 during all the simulation for an elapsed computing time about 5 hours. The present case shows the performance of the automatic mesh refinement strategy. Here, the numerical scheme is parallel and we have used more MPI process than cores based on a simple meshing strategy, one block is put in one domain. Such a simple procedure is of course rather time consuming as the load of the processors is not well balanced because of the required synchronization at each time step. As explained previously and illustrated in the following three-dimensional cases, this synchronization issue can be addressed by the use of BB-AMR scheme.



## 3.2 A 3D dam-break problem without AMR block-based mesh decomposition

In order to test the efficiency of the model, we focus on the well-known 3D dam-break experiments carried out by [20]. This test case being both highly dynamic and strongly three-dimensional as well as well-documented by pressure sensors and wave gauges, it has been widely used to evaluate the performance of several numerical models such as VOF method [20], SPH method [22] or Eulerian-Lagrangian method [30].

The initial configuration of the problem as well as sensors location are shown in Figure 8. The setup consists in a hexahedral block of fluid collapsing in a three-dimensional pool and impacting a rectangular rigid obstacle on a flat bottom (see [20] for detailed description). The computational domain is 3.22m long, 1m high and 1m width. The obstacle is located from  $x = 0.6635\text{m}$  of the tank. The base fixed mesh is splitted into 24 domains refined around the obstacle. Each domain is devoted to one process on the cluster. The numerical results are first qualitatively compared to Kleefsman’s video recording in Fig. 9 and quantitatively compared in Fig. 10 to the experimental data and published numerical results mentioned hereinbefore.

Figure 9 shows a detailed qualitative comparison between Kleefsman’s numerical results (left column), experimental video recording (central column) and our numerical results (right column) at times  $t = 0.4, 0.6, 1, 1.8, 2$  and  $4.8\text{s}$ . The identification of the exact position of the experimental free surface is limited by the flow aeration observed in video recordings. One notes however the very satisfying overall agreement between both experimental and numerical Kleefsman’s data and our results. The successive steps of the experiment are qualitatively well described by our numerical model, including the dambreak before the impact at  $t = 0.4\text{s}$ , the first impact on the obstacle  $t = 0.6\text{s}$ , the reflection on the wall  $t = 1\text{s}$ , the formation of the first reflected wave  $t = 1\text{s}$  and the flow return after the second reflection on the right boundary  $t = 4.8\text{s}$ . The comparison between numerical results reveals the fundamental difference between models. The Kleefsman’s model [20] produces a variety of small scale structures, such as centimetric secondary surface waves or deformations and water fragmentation in drops, while our model shows a much regular aspect of the free surface. At this stage of the qualitative analysis, and recalling that none of the models properly resolves the air-water mixture at the capillary scale, it appears difficult to identify which one of the numerical methods produces the better approximation of the real physical phenomenon.

The qualitative analysis showed that our model nicely compares with existing experimental and numerical results for a fully 3D wave impact case. Our numerical results are now quantitatively compared to experimental data of Kleefsman [20] in terms of pressure (sensor  $P1$ ) and free surface height (gauge  $H2$ ). Similar comparisons with experiments performed on gauge  $H4$  and sensor  $P7$  data (not shown here) showed the same level of agreement. Figure 10 shows the temporal evolution of pressure and free surface elevation for a 6s run comparing experimental data and numerical predictions provided by our model and a set of other published numerical results. Let us first focus on the comparison between the experiments (blue dotted line) and our model (red solid lines). At  $t = 0\text{s}$ , the dam breaks but, as the  $P1$  location is not reached by water, both pressures correspond to atmospheric pressure (Fig. 10, left plot). Around  $t = 0.45\text{s}$ , the water impacts the obstacle. One notes the simultaneous violent increase of pressure for both numerical and experimental data. The numerical pressure peak is less intense than the experimental one. As the water flows around and above the obstacle and then reflects on the left wall, the pressure at  $P1$  continuously decreases. After  $t = 1.8\text{s}$ , the obstacle is submerged by the reflecting wave which induces a slight increase of pressure observed on both experimental and numerical results. One notes here the appearance of a small delay between experiments and model. The pressure then still decreases until the arrival of the second reflected wave. The wave amplitude is rather well reproduced by the model but it comes with a delay around  $0.3\text{s}$  and a smoother shape. Further comparisons are performed with experimental wave height for gauge  $H2$  (Fig. 10, left plot). The agreement between numerical and experimental wave heights is quite satisfactory. Both incoming and reflected waves are nicely reproduced by the model. Some differences are observed, in particular for the return wave which arrives with a slight delay around  $0.3\text{s}$  and a loss of amplitude.

Despite some discrepancies between model and experimental results, figure 10 demonstrates the ability of the present numerical in describing the whole dambreak process. Kleefsman’s experiments have also recently been used as a benchmark test for several other numerical models. A comparison is here proposed between the presented model, the truly incompressible Smoothed Particle Hydrodynamics of [22], the NS-VOFSM model with hybrid Eulerian-Lagrangian resolution of the transport equations by [30] and the Kleefsman’s NS-

VOF model [20]. Figure 10 shows a quantitative comparison between the four models and the experimental data for water height at H2 gauge (left graph) and pressure at P1 sensor (right graph). One notes first that both water level and pressure evolution are macroscopically well described by all numerical models, excepted the pressure peak during the first impact which is poorly described by the SPH method. Some differences are observed between models but their detailed description is not the aim of the present work. The main conclusion of the proposed comparison is to highlight that our mathematical and numerical model nicely compares with state-of-the-art heavier numerical methods in a quantitative point of view.

### 3.3 A 3D dam-break problem with AMR block-based mesh decomposition

The previous section demonstrated that our simple two-fluid model is, at the least, competitive with heavier numerical models. The Block Based AMR method described in this paper is now introduced and tested against the same Kleefsman’s experiment [20]. Numerical parameters are  $CFL = 0.8$ , 3628 blocks, 48 cores, 48 domains or Mpi process, 4 levels of mesh refinement,  $\alpha_{max} = 0.2$  and  $\alpha_{min} = 0.02$  as mesh refinement and coarsening parameters, respectively. Slip boundary conditions are imposed with no penetration. The simulation time (4.8s) is divided in 240 AMR time intervals

The model accuracy in describing flow dynamics being unchanged, we restrict our comments to the numerical aspects of the computation. Figure 11 shows the evolution of the total number of cells. It is observed to evolve from 800 000 cells up to about 1 500 000 cells. By contrast with [8, 15], the number of cells is not prescribed, but prescribing the maximum of mesh refinement level and the parameters  $\alpha_{min}$  and  $\alpha_{max}$  is sufficient to prevent a number of cells out of control. As seen in figure 12 (left), the distribution of cells on the 48 domains follows the flow evolution (Fig. 12, right plot) as explained in section 2.4.2 and maintains a well-balanced processors load. Finally, the present simulation, performed with the standard Runge Kutta time integration on 48 Intel X5675 cores, is achieved in 10 hours. The same test case simulated without BB-AMR but with a faster Adams-Bashforth time integration (3 levels) on 12 Itanium II cores is nearly five time longer.

## 4 Concluding remarks and perspectives

This paper reports on numerical simulation of two-phases dam-break flows over an obstacle. The presented mathematical model is based on previous works of [17, 13, 27]. The model has been specifically developed to simulate violent three-dimensional air and water flows, such as breaking or impacting waves. A set of assumptions is made to establish the best compromise between computing costs and physical relevancy in the considered physical framework. As such, this can be considered as an alternative “intermediate“ approach between full Navier Stokes solvers and simplified Shallow Water models. It is expected that such approach, which provides access to fully three-dimensional simulations at a rather reasonable computing cost, can bring a great insight into a variety of coastal and marine engineering problems.

The model is based on an artificial linearized pressure law for low Mach regime. This artificial pressure law induces several interesting effects from mathematical and numerical viewpoints described above and in previous related papers. The novelty of the present work is to extend the multi-scale adaptive scheme proposed by [10] to our three-dimensional hyperbolic equations. The automatic mesh refinement process is performed thanks to an original criterion: the numerical density of entropy production. The increase of computational costs caused by the re-meshing procedure is outweighed by the introduction of an efficient Block-Based Adaptive Mesh Refinement method (BB-AMR) allowing an easy and fast parallel implementation and computation. Through several 2D and 3D wave impact test cases, the overall numerical approach is validated. In particular, the entropy production appears to be a relevant criterion for automatic mesh refinement process, which leads to better describe the regions of interest (interfaces and impact zones) and to drastically reduce computational time.

On-going developments of the model focus on the optimization of the artificial sound celerity, with the prospect of introducing a spatially and temporally variable value. As this approach can virtually be applied

to any hyperbolic system, longer term extensions of the method will concern the application to shallow water models and fluid-structure interaction issues.

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## References

- [1] S. Abadie, D. Morichon, S. Grilli, and S. Glockner. Numerical simulation of waves generated by landslides using a multiple-fluid navierstokes model. *Coastal Engineering*, 57(9):779 – 794, 2010.
- [2] R. Abgrall. How to prevent pressure oscillations in multicomponent flow calculations: a quasi-conservative approach. *J. Comp Phys.*, 125:150–160, 1996.
- [3] G. Allaire, S. Clerc, and S. Kokh. A five-equation model for the simulation of interfaces between compressible fluids. *J. Comp Phys.*, 181:577–616, 2002.
- [4] C. Altmann, T. Belat, M. Gutnic, P. Helluy, H. Mathis, E. Sonnendrücker, W. Angulo, and J.M. Hérard. A local time-stepping discontinuous galerkin algorithm for the mhd system. *ESAIM*, 28:33–54, 2009.
- [5] C. Altomare, A.J.C. Crespo, B.D. Rogers, J.M. Dominguez, X. Gironella, and M. Gmez-Gesteira. Numerical modelling of armour block sea breakwater with smoothed particle hydrodynamics. *Computers and Structures*, 130(0):34 – 45, 2014.
- [6] A. Banari, C. Janen, S. T. Grilli, and M. Krafczyk. Efficient gpgpu implementation of a lattice boltzmann model for multiphase flows with high density ratios. *Computers & Fluids*, 93(0):1 – 17, 2014.
- [7] M.J. Berger and P. Colella. Local adaptive mesh refinement for shock hydrodynamics. *Journal of Computational Physics*, 82(1):64 – 84, 1989.
- [8] T. Coupez and E. Hachem. Solution of high-reynolds incompressible flow with stabilized finite element and adaptive anisotropic meshing. *Computer Methods in Applied Mechanics and Engineering*, 267(0):65 – 85, 2013.
- [9] D. Dutykh, R. Poncet, and F. Dias. The volna code for the numerical modeling of tsunami waves: Generation, propagation and inundation. *European Journal of Mechanics - B/Fluids*, 30(6):598 – 615, 2011.
- [10] M. Ersoy, F. Golay, and L. Yushchenko. Adaptive multiscale scheme based on numerical density of entropy production for conservation laws. *Central European Journal of Mathematics*, 11(8):1392–1415, 2013.
- [11] D. Fuster, G. Agbaglah, C. Josserand, S. Popinet, and S. Zaleski. Numerical simulation of droplets, bubbles and waves: state of the art. *Fluid Dynamics Research*, 41(6):065001, 2009.
- [12] F. Golay. Numerical entropy production and error indicator for compressible flows. *Comptes Rendus Mecanique*, 337(4):233 – 237, 2009.
- [13] F. Golay and P. Helluy. Numerical schemes for low mach wave breaking. *International Journal of Computational Fluid Dynamics*, 21(2):69–86, 2007.

- [14] S. Grilli, P. Guyenne, and F. Dias. A fully non-linear model for three-dimensionnal overturning waves over an arbitrary bottom. *Int. J. Num. Meth. in Fluids*, 35:829–867, 2001.
- [15] E. Hachem, S. Feghali, R. Codina, and T. Coupez. Immersed stress method for fluid structure interaction using anisotropic mesh adaptation. *International Journal for Numerical Methods in Engineering*, 94(9):805–825, 2013.
- [16] E. Hairer, S. P. Nørsett, and G. Wanner. *Solving ordinary differential equations. I*, volume 8 of *Springer Series in Computational Mathematics*. Springer-Verlag, Berlin, second edition, 1993. Nonstiff problems.
- [17] P. Helluy, F. Golay, J.P. Caltagirone, P. Lubin, S. Vincent, D. Drevrard, R. Marcer, P. Fraunié, N. Seguin, S. Grilli, A.N. Lesage, A. Dervieux, and O. Allain. Numerical simulation of wave breaking. *M2AN*, 39(3):591–607, 2005.
- [18] S.C. Hsiao and T.C. Lin. Tsunami-like solitary waves impinging and overtopping an impermeable seawall: Experiment and rans modeling. *Coastal Engineering*, 57(1):1 – 18, 2010.
- [19] T.W. Hsu, C.M. Hsieh, and R. R. Hwang. Using rans to simulate vortex generation and dissipation around impermeable submerged double breakwaters. *Coastal Engineering*, 51(7):557 – 579, 2004.
- [20] K.M.T. Kleefsman, G. Fekken, A.E.P. Veldman, B. Iwanowski, and B. Buchner. A volume-of-fluid based simulation method for wave impact problems. *Journal of Computational Physics*, 206(1):363 – 393, 2005.
- [21] S. Koshizuka, H. Tamako, and Y. Oka. A particle method for incompressible viscous flow with fluid fragmentations. *Computational Fluid Dynamics Journal*, 4(1):29–46, 1995.
- [22] E.S. Lee, D. Violeau, R. Issa, and S. Ploix. Application of weakly compressible and truly incompressible sph to 3-d water collapse in waterworks. *Journal of Hydraulic Research*, 48(sup1):50–60, 2010.
- [23] F. Losasso, F. Gibou, and R. Fedkiw. Simulating water and smoke with an octree data structure. *ACM Trans. Graph.*, 23(3):457–462, August 2004.
- [24] P. Lubin and S. Glockner. detailed numerical investigation of the three-dimensional flow structures under breaking waves. In *7th international conference on Coastal Dynamics, Arcachon*, 2013.
- [25] C. Min and F. Gibou. A second order accurate level set method on non-graded adaptive cartesian grids. *Journal of Computational Physics*, 225(1):300 – 321, 2007.
- [26] L.-R. Plumerault, D. Astruc, P. Villedieu, and P. Maron. A numerical model for aerated-water wave breaking. *International Journal for Numerical Methods in Fluids*, pages n/a–n/a, 2011.
- [27] A. N. Sambe, D. Sous, F. Golay, P. Fraunié, and R. Marcer. Numerical wave breaking with macro-roughness. *European Journal of Mechanics - B/Fluids*, 30(6):577 – 588, 2011.
- [28] F. Shi, J.-T. Kirby, J.-H. Harris, J.-D. Geiman, and S.-T. Grilli. A high-order adaptive time-stepping tvd solver for boussinesq modeling of breaking waves and coastal inundation. *Ocean Modelling*, 4344(0):36 – 51, 2012.
- [29] E.F. Toro. *Riemann solvers and numerical methods for fluid dynamics*. 2nd ed., Springer: New York, 1999.
- [30] S. Vincent, G. Balmigère, J.-P. Caltagirone, and E. Meillot. Eulerian-lagrangian multiscale methods for solving scalar equations - application to incompressible two-phase flows. *Journal of Computational Physics*, 229(1):73 – 106, 2010.

- [31] M. J. Williamschen and C.P.T. Groth. Parallel anisotropic block-based adaptive mesh refinement algorithm for three-dimensional flows. In *21st AIAA Computational Fluid Dynamics Conference, June 24-27, San Diego, CA*, pages 1–22, 2013.
- [32] K.F.C. Yiu, D.M. Greaves, S. Cruz, A. Saalehi, and A.G.L. Borthwick. Quadtree grid generation: Information handling, boundary fitting and cfd applications. *Computers & Fluids*, 25(8):759 – 769, 1996.
- [33] M. Zhang and W.M. Wu. A two dimensional hydrodynamic and sediment transport model for dam break based on finite volume method with quadtree grid. *Applied Ocean Research*, 33(4):297 – 308, 2011.
- [34] X.D. Zhang, J.Y. Trépanier, and R. Camarero. A posteriori error estimation for finite-volume solutions of hyperbolic conservation laws. *Comp. Meth. App. Eng.*, 185:1–19, 2000.
- [35] Z.X.J. Zheng and C.P.T. Groth. Block-based adaptive mesh refinement finite-volume scheme for hybrid multi-block meshes. In *7st conference on Computational Fluid Dynamics (ICCFD7), July 9-13, Hawaii*, pages 1–19, 2012.

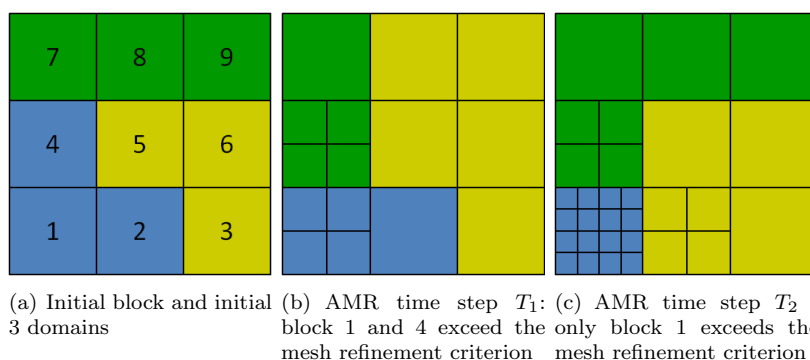


Figure 1: Example of two dimensional BB-AMR with 3 domains (blue-yellow-green) and 9 blocks

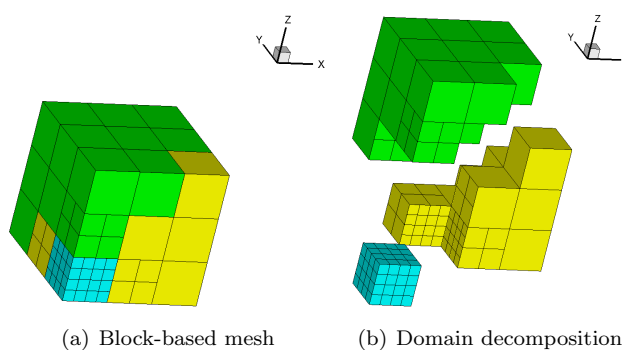


Figure 2: Example of three dimensional block-based mesh with 3 domains and 27 blocks

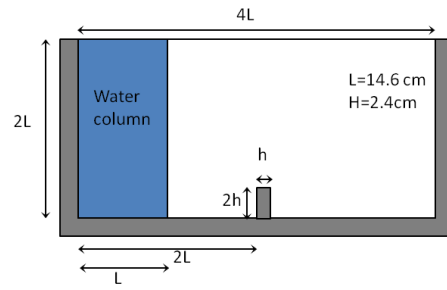


Figure 3: Collapse of water column [21]

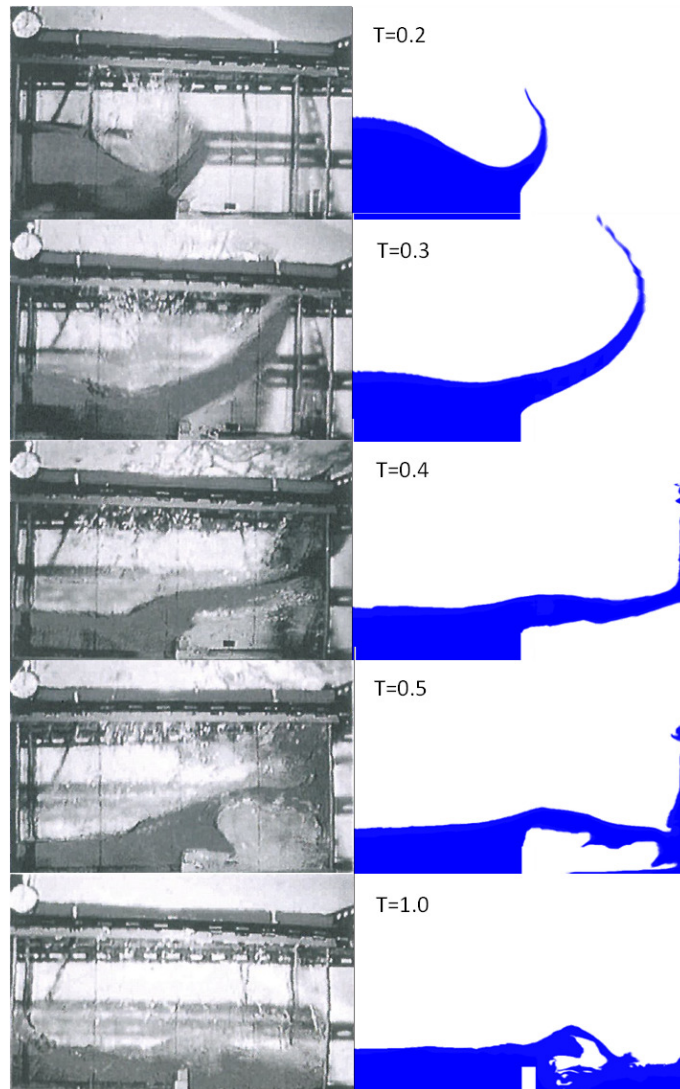


Figure 4: 2D dam-break problem, confrontation with experiments: experimental simulation [21] (left) – numerical simulation (right).

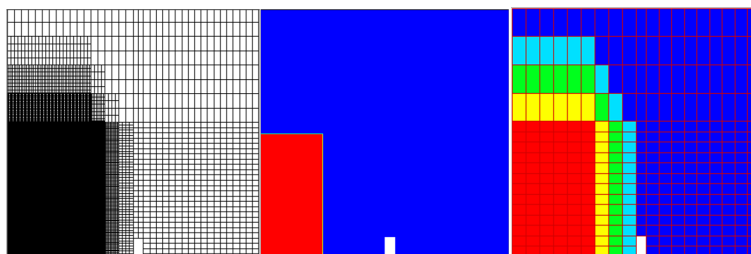


Figure 5: Initial configuration for 2D dam-break problem: mesh (left), density with blue and red corresponding to air and water, respectively (center), mesh refinement level (1 to 5) per block (right)

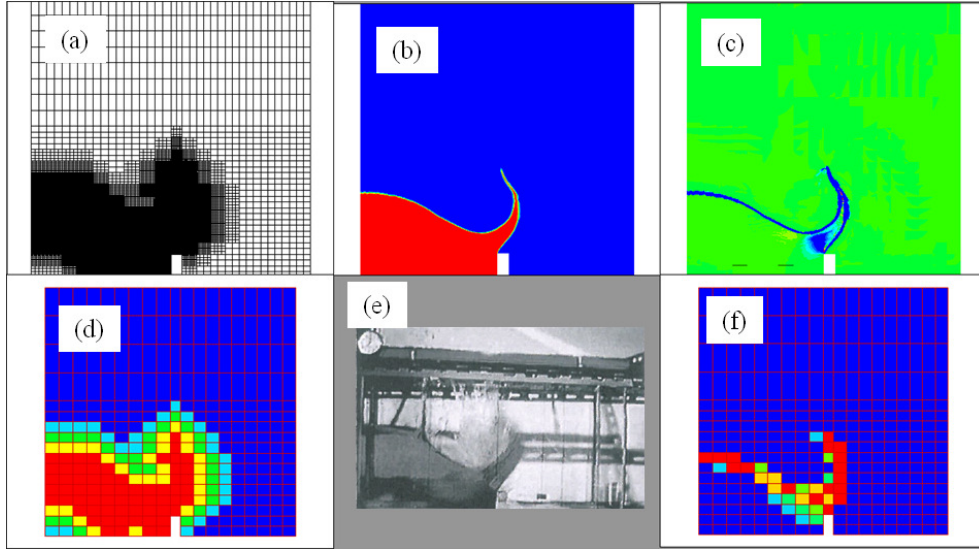


Figure 6: 2D dam-break problem at  $t=0.2s$ .

(a) Mesh; (b) Density (air-blue, water-red); (c) Density of numerical entropy production (equation (9): green-zero, blue-negative values); (d) Mesh refinement level per block (1 to 5); (e) Experiment [21]; (f) Mesh refinement criterion per block.

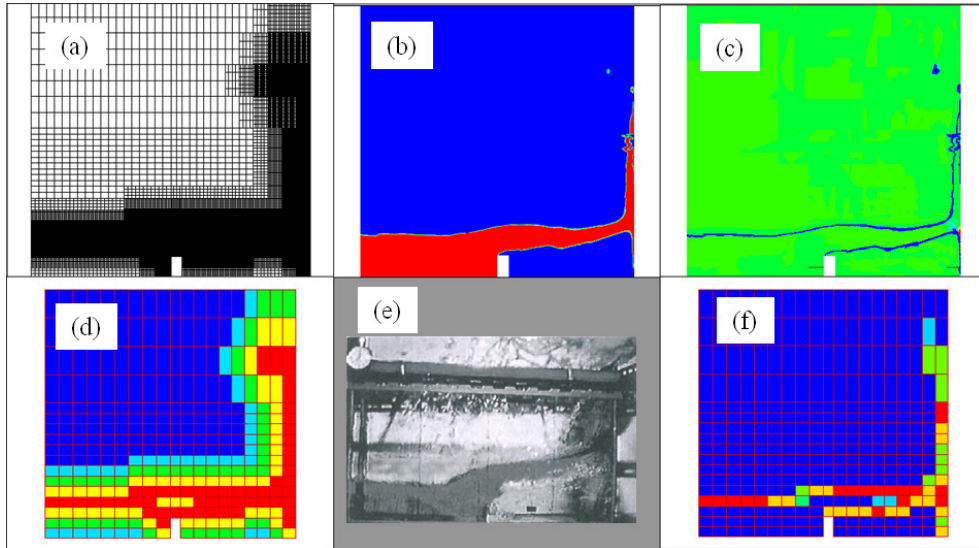


Figure 7: 2D dam-break problem at  $t=0.4s$ .

(a) Mesh; (b) Density (air-blue, water-red); (c) Density of numerical entropy production (equation (9): green-zero, blue-negative values); (d) Mesh refinement level per block (1 to 5); (e) Experiment [21]; (f) Mesh refinement criterion per block



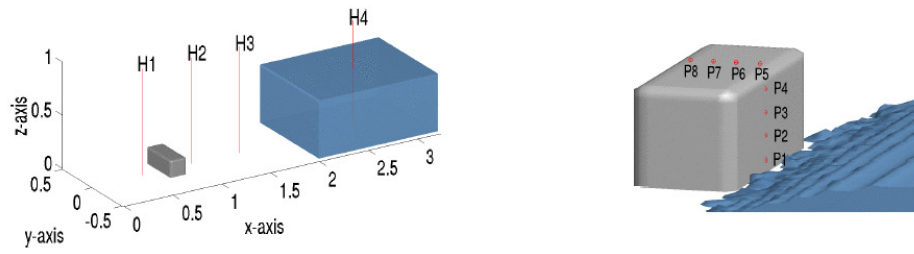
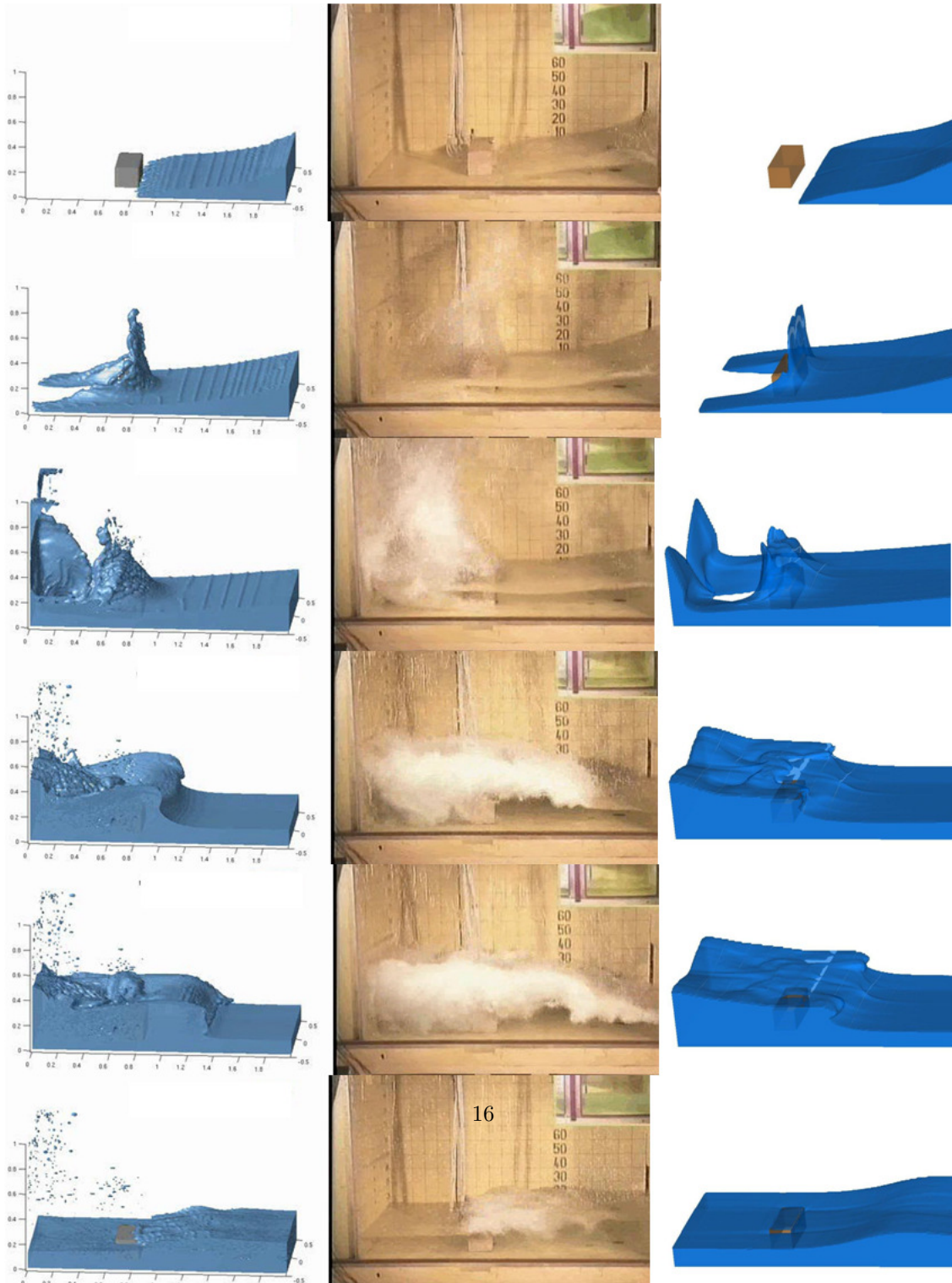


Figure 8: Dam break problem: domain geometry and sensors points.



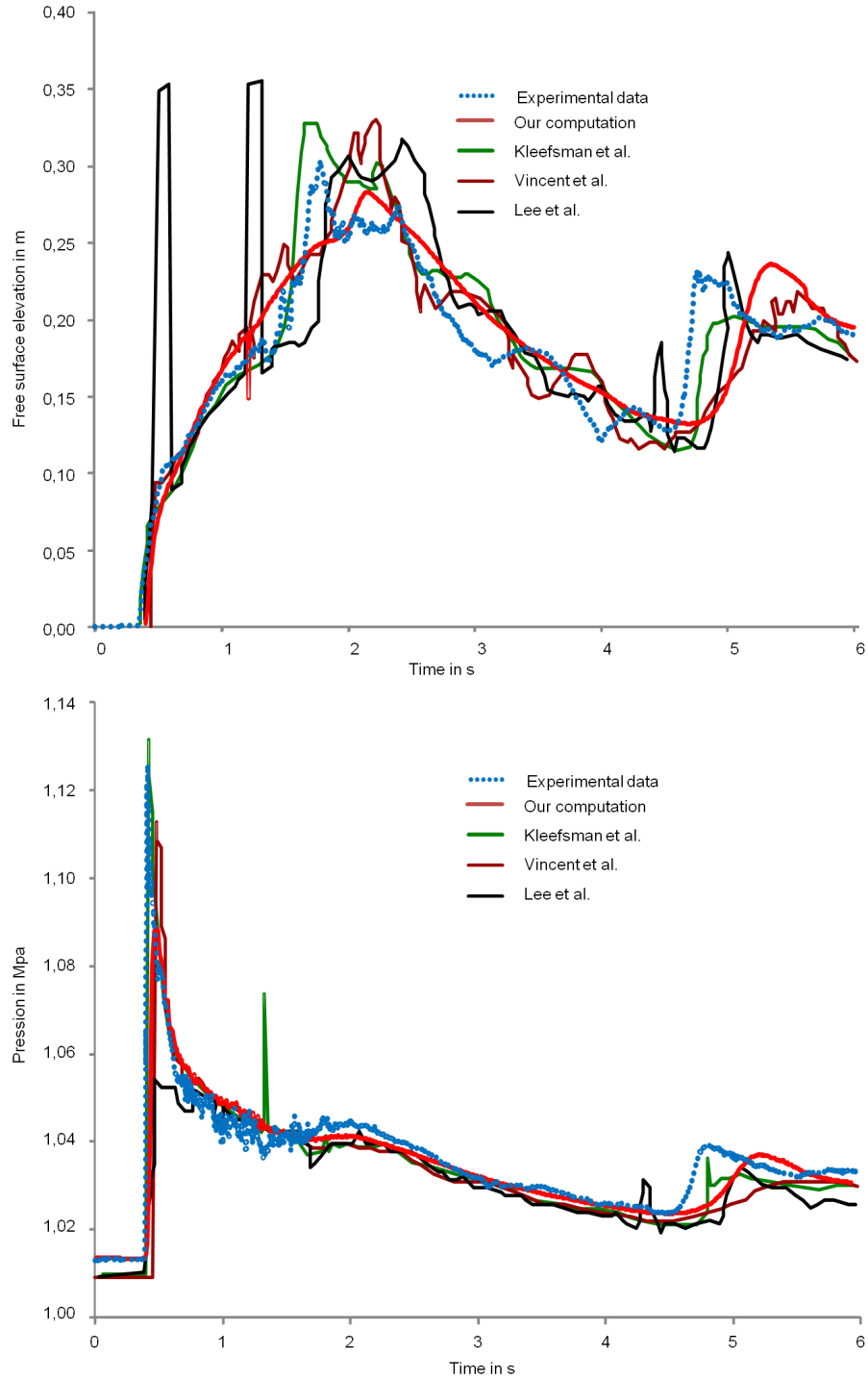


Figure 10: Comparison between three numerical models [20, 30, 22] and the present one. Top: free surface elevation (H2 gauge). Bottom: pressure at P1 sensor.

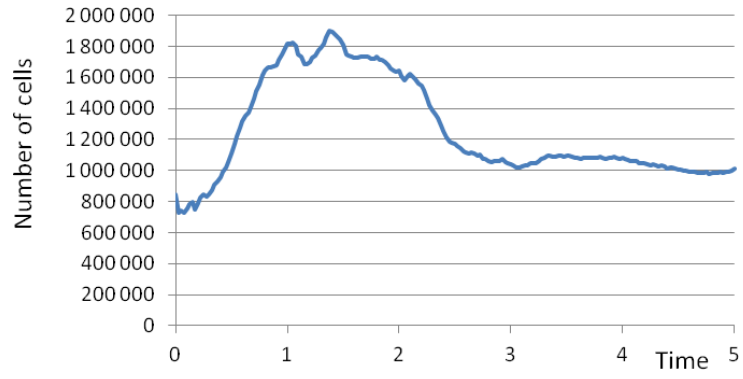


Figure 11: Number of cells during the computation.

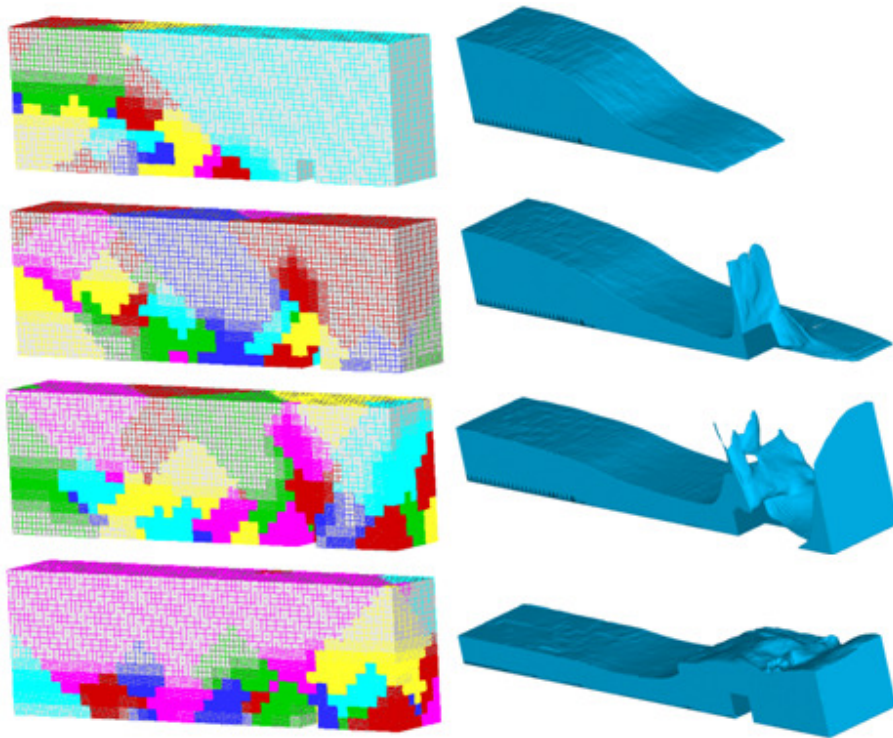


Figure 12: Domains due to the BB-AMR scheme (left) and air-water interface (right) at time  $0.4s, 0.6s, 1.0s, 2s$ .