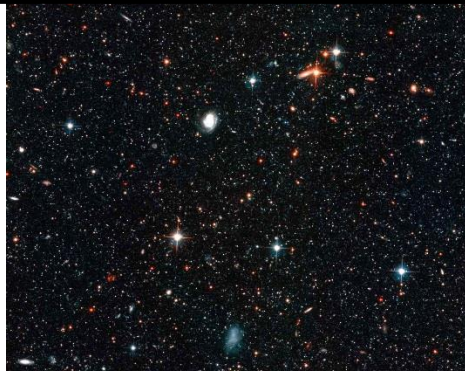


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An energy and determinist approach of quantum mechanics



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2015

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First part

Classical approach of Dirac equation and its solutions

I

Introduction

This manuscript describes a set of reflections on the formalism describing the quantum mechanics in its status at the beginning of 21st century. The first chapters are devoted to the resumption of the physical concepts that accompany the theory through the prism of the DIRAC equation. Abstraction is made of some strong mathematical concepts useful in the synthesis of the theory, but not necessary for the physical understanding of phenomena. The following chapters are entirely devoted to the implementation of extensive work on solutions in the form of stationary modes of the DIRAC equation.

A century of research, both theoretical and experimental helped to significantly refine the knowledge of particle physics. If the experimental tools regularly lead to significant advances in the observation and measurement of phenomena, the theoretical framework seems to be frozen since tens of years, and without a compelling perspective on a next evolution.

This theoretical framework is facing a problem that appears insurmountable: the particles behave both in the manner of a wave, and in the manner of a body of material. Unable to account for this phenomenon, the theory is reduced to treat experimental observations on a statistical and probabilistic point of view. It succeeds in a remarkable way, but at the price of a mathematical complexity that is made necessary to overcome the fact that the physics underlying the observed phenomena is not known to us with sufficient accuracy to lighten the mathematical formalism.

To advance in a reflection that allows physically account for the wave-particle duality, we must develop elements which, while being in perfect coherence with all of the existing formalisms, are breaking with the probabilistic vision of this part of physics. This can be achieved only by a deterministic approach to clarify how the material moves between its wavelike and its corpuscular aspects.

The work presented in this document is based on a set of exact, but new solutions to this day, of the DIRAC equation. This approach ensures therefore, intrinsically, consistency with all the theoretical properties built around this equation. It led, ultimately, to a description deterministic and no more probabilistic, of the wave-particle duality.

I - An energy approach

In the classical wave equations, the dimension of the quantity that spreads under the form of a wave is usually set by the second member. One can illustrate this remark by one example chosen in electromagnetism, concerning the vector and scalar potentials:

$$\vec{\nabla}^2 \cdot \vec{A} - \epsilon_0 \mu_0 \cdot \frac{\partial^2 \vec{A}}{\partial t^2} = -\mu_0 \cdot \vec{J} \quad (\text{I-1})$$

$$\vec{\nabla}^2 \cdot \varphi - \epsilon_0 \mu_0 \cdot \frac{\partial^2 \varphi}{\partial t^2} = -\frac{\rho}{\epsilon_0} \quad (\text{I-2})$$

In the wave equation of Schrödinger or Klein-Gordon, the solution ψ function is a quantity without information on its dimensions, because it is present in the second part of the equation:

$$j\hbar \frac{\partial}{\partial t} (\psi) = -\frac{\hbar^2}{2m} \nabla^2 (\psi) + E_p \psi \quad (\text{I-3})$$

$$\nabla^2 (\psi) - \frac{1}{c^2} \frac{\partial^2}{\partial t^2} (\psi) = \frac{m^2 c^2}{\hbar^2} \psi \quad (\text{I-4})$$

This means that the ψ solution reflects the function of propagation of the wave. The spreading quantity can therefore be chosen on the basis of physical considerations.

The fact that the function ψ is complex, coupled with a probabilistic view of quantum mechanics, have led to give to the quantity $\psi\psi^*$ the meaning of a volumetric density of probability, involving a normalization on any volume where the particle is located with certainty:

$$\iiint_V \psi\psi^* dv = 1 \quad (\text{I-5})$$

If we adopt a deterministic vision of the phenomena, there is no more probability density. The physical quantity that spreads can be considered to be an energy, and consistency with the probabilistic vision suggests to give to the quantity $\psi\psi^*$ the meaning of a volumetric energy density. The condition of standardization is obtained by expressing that integration on a volume V where the particle extends must give the total energy E of the particle:

$$\iiint_V \psi\psi^* dv = E \quad (\text{I-6})$$

Wave functions issued from the probabilistic theory and from the deterministic theory are therefore proportional in a ratio of square root of E .

The wave function ψ becomes homogeneous with the square root of energy. And since there are negative energies, the complex number $j = \sqrt{-1}$ appears in a natural way and with a clear physical meaning in the solutions of the Dirac equation describing both positive energy particles and particles of negative energies.

I - 1 The different kinds of energy

In special relativity, four-vectors are identified as invariant physical quantities by change of frame. It follows that their pseudo-norm is constant and does not depends on the frame in which it is evaluated.

For a mass m moving at speed v , four-vector momentum-energy is expressed as follows:

$$\tilde{\mathbf{P}} = \begin{pmatrix} m \frac{dx}{dt} \\ m \frac{dy}{dt} \\ m \frac{dz}{dt} \\ \frac{E}{c} \end{pmatrix} = \begin{pmatrix} mv_x \\ mv_y \\ mv_z \\ \frac{E}{c} \end{pmatrix} = \begin{pmatrix} p_x \\ p_y \\ p_z \\ \frac{E}{c} \end{pmatrix} \quad (\text{I-7})$$

In this expression, E means the whole energy of the particle of mass m , moving with a speed v , and therefore linked to rest mass by the relationship:

$$m = \frac{m_0}{\sqrt{1 - \frac{v^2}{c^2}}} \quad (\text{I-8})$$

The pseudo-norm raised to the square of the four-vector pulse energy is written:

$$\left(m \frac{dx}{dt}\right)^2 + \left(m \frac{dy}{dt}\right)^2 + \left(m \frac{dz}{dt}\right)^2 - \left(\frac{E}{c}\right)^2 = cte \quad (\text{I-9})$$

or again:

$$E^2 - p^2c^2 = cte \quad (\text{I-10})$$

where p is the pulse module:

$$p = \sqrt{\left(m \frac{dx}{dt}\right)^2 + \left(m \frac{dy}{dt}\right)^2 + \left(m \frac{dz}{dt}\right)^2} = m \sqrt{\left(\frac{dx}{dt}\right)^2 + \left(\frac{dy}{dt}\right)^2 + \left(\frac{dz}{dt}\right)^2} = mv \quad (\text{I-11})$$

One determines the constant by writing that, according to the relation which links energy and mass, the total energy of a particle of mass m is also equal to:

$$E = mc^2 = \frac{m_0c^2}{\sqrt{1 - \frac{v^2}{c^2}}} \quad (\text{I-12})$$

where m_0 is the rest mass.

We can deduce the value of the constant:

$$E^2 - p^2 c^2 = \left(\frac{m_0 c^2}{\sqrt{1 - \frac{v^2}{c^2}}} \right)^2 - \left(\frac{m_0 v}{\sqrt{1 - \frac{v^2}{c^2}}} \right)^2 c^2 = (m_0 c^2)^2 \quad (\text{I-13})$$

and finally the expression of the conservation of energy in special relativity:

$$E^2 = (pc)^2 + (m_0 c^2)^2 \quad (\text{I-14})$$

We accept as a postulate from the expression above, that within a particle of mass m and moving at speed v , there may be only three specific forms of energy:

- A form of energy representative of the whole energy E , which we will subsequently justify the designation of wave energy
- A form of energy representing pulse energy: pc
- A form of energy representing the mass energy at rest: $m_0 c^2$

Let us also assume that, in the world of the infinitely small, these three forms of energy are distinct, in the same way as electrical and magnetic energy in electromagnetism, or potential and kinetic energy in mechanics.

I - 2 Stationary energy

Consider a particle at rest, and therefore which has no pulse energy. One is led to the assumption that its mass energy is located in a finite spatial extent. This energy is constant in time, and we therefore call it stationary energy. Since there is no pulse energy, the total energy in this particular case, that is to say the wave energy, is equal to the mass energy.

We then ask about what we know on the stationarity of the energy in the physics of waves in general. This stationarity is a property of systems that retain the trapped wave in a limited portion of the space. If we assume the system without loss, and if we introduce the wave energy in this system, it is preserved in time.

The fact of importance is that this energy never stays at rest in the portion of the space where it is confined, but settles in stationary modes that include at least two types of energies that are exchanged in general in time quadrature.

In mechanics, the waves can propagate in a medium with certain elasticity, and exchange occurs between the elastic potential energy and kinetic energy of the moving elements.

In electromagnetism, the exchange is between the electric and magnetic energy.

If we now return to the situation of a particle at rest that contains both of the wave energy and the mass energy, we postulate that these two forms of energy exchange in the portion of the space where they are confined. This assumption stems directly from the particle behavior under the form of a wave or a particle.

If the particle is moving, it appears a third form of energy which is the impulse energy. We postulate that this energy will participate in exchanges with the two previous ones in a form that remains to be determined, but which verifies the equation of conservation of energy (I-14).

Later in this document, we will not wonder on the way in which the different types of energy are confined in a region of space: the answer to this question is not known to us. But we will show that if we make the assumption that the different kinds of energy exchange in stationary modes, this led to a quantum physics deterministic and consistent with the actual probabilistic theory.

II

DIRAC equation

The DIRAC equation is part of some fundamental equations of contemporary physics. It was obtained in the wake of two other very important equations of quantum mechanics: the SCHROEDINGER equation and the KLEIN-GORDON equation. It has supplanted the latter two because it describes with more precision the reality of observed phenomena.

The objective of this short chapter of presentation is not to discuss the contribution of this equation on quantum physics that the reader will find in books and websites related to this subject and to which we will return later. It's just, in a first time, to retrace the approach leading to the DIRAC equation, and underline some difficulties on the physical interpretation of its solution.

I - The SCHROEDINGER equation

Any linear physical phenomenon $\psi(x, y, z, t)$ which propagates at speed v in a three dimensions space as time flows can be described by a wave equation:

$$\nabla^2 \psi(x, y, z, t) - \frac{1}{v^2} \frac{\partial^2 \psi(x, y, z, t)}{\partial t^2} = 0 \quad (\text{II-1})$$

When dealing with a signal with a sinusoidal temporal variation, we can put:

$$\psi(x, y, z, t) = \varphi(x, y, z) \exp(\pm j\omega t) \quad (\text{II-2})$$

and the wave equation for the quantity φ independent of time becomes:

$$\nabla^2 \varphi(x, y, z) + \frac{\omega^2}{v^2} \varphi(x, y, z) = 0 \quad (\text{II-3})$$

Following Louis De BROGLIE assuming that one can associate to any particle of mass m moving at a velocity v a wave of wavelength $\lambda = h/mv$, where h is the PLANCK constant, the time-independent wave equation attached to the particle becomes:

$$\nabla^2 \varphi(x, y, z) + \frac{\omega^2}{v^2} \varphi(x, y, z) = \nabla^2 \varphi(x, y, z) + \left(\frac{2\pi}{\lambda}\right)^2 \varphi(x, y, z) = 0 \quad (\text{II-4})$$

On the other hand, in the context of classical mechanics, energy total E of this particle is the sum of its kinetic energy E_c and its potential energy E_p , which allows to write:

$$E_c = \frac{1}{2}mv^2 = E - E_p \quad (\text{II-5})$$

from which is deducting successively:

$$mv = \sqrt{2m(E - E_p)} \quad (\text{II-6})$$

$$\lambda = \frac{h}{mv} = \frac{h}{\sqrt{2m(E - E_p)}} \quad (\text{II-7})$$

This last relationship put in the wave equation (I-4) provides the time-independent SCHROEDINGER equation:

$$\nabla^2 \varphi(x, y, z) + \frac{2m}{\hbar^2} (E - E_p) \varphi(x, y, z) = 0 \quad \text{with } \hbar = \frac{h}{2\pi} \quad (\text{II-8})$$

The heuristic approach that has been proposed for this relationship can be completed to bring up the time dependence.

For a locally flat, monochromatic wave going away towards infinity, we can write the wave function in a general manner under the form:

$$\psi(\vec{r}, t) = \psi_0 \exp[j(\vec{k}\vec{r} - \omega t)] \quad \text{with } j^2 = -1 \quad (\text{II-9})$$

In the quantum world where $\lambda = h/mv$ and where the whole energy E is related to the frequency of the wave by the relation:

$$E = hv = \frac{h}{2\pi} \omega = \hbar\omega \quad (\text{II-10})$$

we obtain by substitution:

$$\psi(\vec{r}, t) = \psi_0 \exp\left[\frac{j}{\hbar}(\vec{p}\vec{r} - Et)\right] \quad \text{with } \vec{p} = m\vec{v} \quad (\text{II-11})$$

In deriving this expression over time, it follows:

$$\frac{\partial}{\partial t}(\psi(\vec{r}, t)) = -\frac{j}{\hbar} E \psi_0 \exp\left[\frac{j}{\hbar}(\vec{p}\vec{r} - Et)\right] = -\frac{j}{\hbar} E \psi(\vec{r}, t) \quad (\text{II-12})$$

and taking the gradient of this same expression (II-11):

$$\vec{\nabla}(\psi(\vec{r}, t)) = \frac{j}{\hbar} \vec{p} \psi(\vec{r}, t) \quad (\text{II-13})$$

From the relations (II-12) and (II-13), are deduced the quantization rules of classical or relativistic mechanics equations which allow to obtain similar equations in the quantum field:

$$j\hbar \frac{\partial}{\partial t}(\psi) = E\psi \quad (\text{II-14})$$

$$-j\hbar \vec{\nabla}(\psi) = \vec{p}\psi \quad (\text{II-15})$$

For a particle whose total energy is given in classical mechanics by the sum of kinetic energy and potential energy:

$$E = E_c + E_p = \frac{1}{2}mv^2 + E_p = \frac{p^2}{2m} + E_p \quad (\text{II-16})$$

It is deduced after multiplication by the wave function ψ :

$$E\psi = \frac{p^2}{2m}\psi + E_p\psi \quad (\text{II-17})$$

and finally by using the rules of quantification (II-14) (II-15):

$$j\hbar \frac{\partial}{\partial t}(\psi) = \frac{(-j\hbar \vec{\nabla})^2}{2m}(\psi) + E_p\psi = -\frac{\hbar^2}{2m}\nabla^2(\psi) + E_p\psi \quad (\text{II-18})$$

which is the time dependent SCHRÖDINGER equation.

This equation allows to find the main series describing the emission lines of the hydrogen atom as well as other parameters like the diameters of the BOHR orbit, but it is in default when there is interest in more subtle phenomena such as the fine structure levels for a hydrogen atom or taking into account of the spin of the electron.

II - The KLEIN-GORDON equation

The inadequacies of the SCHROEDINGER equation are attributed to the fact that the quantization rules have been applied to an energy balance using classical mechanics.

Relativity introduces a relationship whose scope is much broader because it is invariant under the LORENTZ transformation:

$$E^2 = p^2c^2 + m_0^2c^4 \quad (\text{II-19})$$

and we obtain immediately by applying the rules of quantification (II-14) (II-15):

$$-\hbar^2 \frac{\partial^2}{\partial t^2}(\psi) = -\hbar^2c^2\nabla^2(\psi) + m_0^2c^4\psi \quad (\text{II-20})$$

which is the KLEIN-GORDON equation. We can present it in a form that make the wave equation appearing with a second member:

$$\nabla^2(\psi) - \frac{1}{c^2} \frac{\partial^2}{\partial t^2}(\psi) = \frac{m_0^2c^2}{\hbar^2}\psi \quad (\text{II-21})$$

A detailed review of its solutions allows to show some inadequacies, particularly in his inability to describe the spin of the electron which is of order 1/2 as well as some fine levels of energy of the hydrogen atom.

III - The DIRAC equation

Beginning with the relativistic energy conservation equation:

$$E^2 = p^2 c^2 + m_0^2 c^4 \quad (\text{II-22})$$

and applying the rules of quantification of quantum mechanics:

$$E \rightarrow j\hbar \frac{\partial}{\partial t} \quad p_x \rightarrow -j\hbar \frac{\partial}{\partial x} \quad p_y \rightarrow -j\hbar \frac{\partial}{\partial y} \quad p_z \rightarrow -j\hbar \frac{\partial}{\partial z} \quad (\text{II-23})$$

It transforms the KLEIN-GORDON equation in the formula:

$$\left(j\hbar \frac{\partial}{\partial t} \right)^2 (\psi) = \left(-j\hbar c \frac{\partial}{\partial x} \right)^2 (\psi) + \left(-j\hbar c \frac{\partial}{\partial y} \right)^2 (\psi) + \left(-j\hbar c \frac{\partial}{\partial z} \right)^2 (\psi) + m_0^2 c^4 \quad (\text{II-24})$$

According to the formalism proposed by DIRAC, we have to find coefficients α_i such that the above equation is verified in the form of a first order partial differential equation raised to square:

$$\left(j\hbar \frac{\partial}{\partial t} \right)^2 (\psi) = \left[\alpha_1 \left(-j\hbar c \frac{\partial}{\partial x} \right) + \alpha_2 \left(-j\hbar c \frac{\partial}{\partial y} \right) + \alpha_3 \left(-j\hbar c \frac{\partial}{\partial z} \right) + \alpha_0 (m_0 c^2) \right]^2 (\psi) \quad (\text{II-25})$$

If the relationship above is checked through the presence of coefficients α_i , the solution can be obtained by solving the equation obtained by removing the squares:

$$\left(j\hbar \frac{\partial}{\partial t} \right) (\psi) = \left[\alpha_1 \left(-j\hbar c \frac{\partial}{\partial x} \right) + \alpha_2 \left(-j\hbar c \frac{\partial}{\partial y} \right) + \alpha_3 \left(-j\hbar c \frac{\partial}{\partial z} \right) + \alpha_0 (m_0 c^2) \right] (\psi) \quad (\text{II-26})$$

It is impossible to find a α_i real or complex in response to conditions (II-25) and (II-26), but you can find matrices α_i who meet the following, necessary and sufficient conditions to transform equation (II-24) into equation (II-25):

$$\alpha_i^2 = 1 \quad (\text{II-27})$$

$$\alpha_i \alpha_j + \alpha_j \alpha_i = 0 \quad (\text{II-28})$$

These matrices are not unique, but a simple choice is that proposed by Dirac:

$$\alpha_0 = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix} \quad (\text{II-29})$$

The following three are defined from the called PAULI matrices σ_i :

$$\alpha_i = \begin{pmatrix} 0 & \sigma_i \\ \sigma_i & 0 \end{pmatrix} \quad \sigma_1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \quad \sigma_2 = \begin{pmatrix} 0 & -j \\ j & 0 \end{pmatrix} \quad \sigma_3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \quad (\text{II-30})$$

or still explicitly:

$$\alpha_1 = \begin{pmatrix} 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \\ 0 & 1 & 0 & 0 \\ 1 & 0 & 0 & 0 \end{pmatrix} \quad \alpha_2 = \begin{pmatrix} 0 & 0 & 0 & -j \\ 0 & 0 & j & 0 \\ 0 & -j & 0 & 0 \\ j & 0 & 0 & 0 \end{pmatrix} \quad \alpha_3 = \begin{pmatrix} 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & -1 \\ 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \end{pmatrix} \quad (\text{II-31})$$

The structure of the matrixes α_i requires the search for the solution ψ in the form of a column vector:

$$\psi = \begin{pmatrix} \Psi_0 \\ \Psi_1 \\ \Psi_2 \\ \Psi_3 \end{pmatrix} \quad (\text{II-32})$$

In order to put the system of differential equations in a presentation such as the mass energy appears in factor with the identity matrix, we multiply the equation (II-26) by the matrix α_0 . It is known that its square is equal to the identity matrix:

$$\alpha_0 \left(j\hbar \frac{\partial}{\partial t} \right) (\psi) = \left[\alpha_0 \alpha_1 \left(-j\hbar c \frac{\partial}{\partial x} \right) + \alpha_0 \alpha_2 \left(-j\hbar c \frac{\partial}{\partial y} \right) + \alpha_0 \alpha_3 \left(-j\hbar c \frac{\partial}{\partial z} \right) + (m_0 c^2) \right] (\psi) \quad (\text{II-33})$$

We put then:

$$\gamma_0 = \alpha_0 \quad \gamma_1 = \alpha_0 \alpha_1 \quad \gamma_2 = \alpha_0 \alpha_2 \quad \gamma_3 = \alpha_0 \alpha_3 \quad (\text{II-34})$$

and we gather all terms containing partial derivatives:

$$\left[\gamma_0 \left(\frac{\partial}{\partial t} \right) + \gamma_1 \left(c \frac{\partial}{\partial x} \right) + \gamma_2 \left(c \frac{\partial}{\partial y} \right) + \gamma_3 \left(c \frac{\partial}{\partial z} \right) \right] (\psi) = \frac{m_0 c^2}{j\hbar} (\psi) \quad (\text{II-35})$$

The matrix identity, implicit in the right term will be omitted in the rest of the document.

The matrices γ_i are obtained explicitly as follows:

$$\gamma_0 = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix} \quad (\text{II-36})$$

$$\gamma_1 = \alpha_0 \alpha_1 = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix} \begin{pmatrix} 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \\ 0 & 1 & 0 & 0 \\ 1 & 0 & 0 & 0 \end{pmatrix} = \begin{pmatrix} 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \\ 0 & -1 & 0 & 0 \\ -1 & 0 & 0 & 0 \end{pmatrix} \quad (\text{II-37})$$

$$\gamma_2 = \alpha_0 \alpha_2 = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix} \begin{pmatrix} 0 & 0 & 0 & -j \\ 0 & 0 & j & 0 \\ 0 & -j & 0 & 0 \\ j & 0 & 0 & 0 \end{pmatrix} = \begin{pmatrix} 0 & 0 & 0 & -j \\ 0 & 0 & j & 0 \\ 0 & j & 0 & 0 \\ -j & 0 & 0 & 0 \end{pmatrix} \quad (\text{II-38})$$

$$\gamma_3 = \alpha_0 \alpha_3 = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix} \begin{pmatrix} 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & -1 \\ 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \end{pmatrix} = \begin{pmatrix} 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & -1 \\ -1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \end{pmatrix} \quad (\text{II-39})$$

The arrangement of these matrices can be synthesized on the basis of the PAULI matrices:

$$\gamma_i = \begin{pmatrix} 0 & \sigma_i \\ -\sigma_i & 0 \end{pmatrix} \quad (\text{II-40})$$

and we can check the following property:

$$(\gamma_0)^2 = 1, (\gamma_1)^2 = (\gamma_2)^2 = (\gamma_3)^2 = -1 \quad (\text{II-41})$$

Considering now that the quantity ψ of the initial wave equation arises in the form of a vector with 4 column, complete writing of the matrix system is given by:

$$\begin{aligned}
 \frac{mc}{\hbar} \begin{pmatrix} \Psi_0 \\ \Psi_1 \\ \Psi_2 \\ \Psi_3 \end{pmatrix} &= \frac{j}{c} \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix} \begin{pmatrix} \frac{\partial \Psi_0}{\partial t} \\ \frac{\partial \Psi_1}{\partial t} \\ \frac{\partial \Psi_2}{\partial t} \\ \frac{\partial \Psi_3}{\partial t} \end{pmatrix} + j \begin{pmatrix} 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \\ 0 & -1 & 0 & 0 \\ -1 & 0 & 0 & 0 \end{pmatrix} \begin{pmatrix} \frac{\partial \Psi_0}{\partial x} \\ \frac{\partial \Psi_1}{\partial x} \\ \frac{\partial \Psi_2}{\partial x} \\ \frac{\partial \Psi_3}{\partial x} \end{pmatrix} \\
 &+ j \begin{pmatrix} 0 & 0 & 0 & -j \\ 0 & 0 & j & 0 \\ 0 & j & 0 & 0 \\ -j & 0 & 0 & 0 \end{pmatrix} \begin{pmatrix} \frac{\partial \Psi_0}{\partial y} \\ \frac{\partial \Psi_1}{\partial y} \\ \frac{\partial \Psi_2}{\partial y} \\ \frac{\partial \Psi_3}{\partial y} \end{pmatrix} + j \begin{pmatrix} 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & -1 \\ -1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \end{pmatrix} \begin{pmatrix} \frac{\partial \Psi_0}{\partial z} \\ \frac{\partial \Psi_1}{\partial z} \\ \frac{\partial \Psi_2}{\partial z} \\ \frac{\partial \Psi_3}{\partial z} \end{pmatrix}
 \end{aligned} \tag{II-42}$$

or again, in a more condensed manner, under the form of a partial derivatives equations system of 4 equations, after putting $\eta = \frac{mc}{\hbar}$:

$$\begin{aligned}
 \eta \Psi_0 &= j \frac{\partial \Psi_0}{\partial(ct)} + j \frac{\partial \Psi_3}{\partial x} + \frac{\partial \Psi_3}{\partial y} + j \frac{\partial \Psi_2}{\partial z} \\
 \eta \Psi_1 &= j \frac{\partial \Psi_1}{\partial(ct)} + j \frac{\partial \Psi_2}{\partial x} - \frac{\partial \Psi_2}{\partial y} - j \frac{\partial \Psi_3}{\partial z} \\
 \eta \Psi_2 &= -j \frac{\partial \Psi_2}{\partial(ct)} - j \frac{\partial \Psi_1}{\partial x} - \frac{\partial \Psi_1}{\partial y} - j \frac{\partial \Psi_0}{\partial z} \\
 \eta \Psi_3 &= -j \frac{\partial \Psi_3}{\partial(ct)} - j \frac{\partial \Psi_0}{\partial x} + \frac{\partial \Psi_0}{\partial y} + j \frac{\partial \Psi_1}{\partial z}
 \end{aligned} \tag{II-43}$$

These equations show that the 4 quantity associated with the wave of material ψ interact, without it is possible to specify in a clear and detailed manner the physical nature of their interaction. The representation in terms of spinors provides a rigorous framework on which are based of multiple developments, but the notion of spinor remains abstract, despite the important efforts of the scientific community to give concrete illustrations.

III

DIRAC bi-spinors

In the previous chapter, we concluded that solutions of the Dirac equation are presented in the form of a quantity with 4 components:

$$\psi = \begin{pmatrix} \psi_0 \\ \psi_1 \\ \psi_2 \\ \psi_3 \end{pmatrix} \quad (\text{III-1})$$

A careful review of the system of Dirac:

$$\begin{aligned} \eta\psi_0 &= j \frac{\partial\psi_0}{\partial(ct)} + j \frac{\partial\psi_3}{\partial x} + \frac{\partial\psi_3}{\partial y} + j \frac{\partial\psi_2}{\partial z} \\ \eta\psi_1 &= j \frac{\partial\psi_1}{\partial(ct)} + j \frac{\partial\psi_2}{\partial x} - \frac{\partial\psi_2}{\partial y} - j \frac{\partial\psi_3}{\partial z} \\ \eta\psi_2 &= -j \frac{\partial\psi_2}{\partial(ct)} - j \frac{\partial\psi_1}{\partial x} - \frac{\partial\psi_1}{\partial y} - j \frac{\partial\psi_0}{\partial z} \\ \eta\psi_3 &= -j \frac{\partial\psi_3}{\partial(ct)} - j \frac{\partial\psi_0}{\partial x} + \frac{\partial\psi_0}{\partial y} + j \frac{\partial\psi_1}{\partial z} \end{aligned} \quad (\text{III-2})$$

shows that these solutions can be grouped by 2 in a behavior with analogies. To display these analogies, it is opportune to introduce the system in a slightly different arrangement:

$$\begin{aligned} j \frac{\partial\psi_0}{\partial(ct)} &= \eta\psi_0 - j \frac{\partial\psi_3}{\partial x} - \frac{\partial\psi_3}{\partial y} - j \frac{\partial\psi_2}{\partial z} \\ j \frac{\partial\psi_1}{\partial(ct)} &= \eta\psi_1 - j \frac{\partial\psi_2}{\partial x} + \frac{\partial\psi_2}{\partial y} + j \frac{\partial\psi_3}{\partial z} \\ j \frac{\partial\psi_2}{\partial(ct)} &= -\eta\psi_2 - j \frac{\partial\psi_1}{\partial x} - \frac{\partial\psi_1}{\partial y} - j \frac{\partial\psi_0}{\partial z} \\ j \frac{\partial\psi_3}{\partial(ct)} &= -\eta\psi_3 - j \frac{\partial\psi_0}{\partial x} + \frac{\partial\psi_0}{\partial y} + j \frac{\partial\psi_1}{\partial z} \end{aligned} \quad (\text{III-3})$$

Therefore, if we put:

$$\Psi = \begin{pmatrix} \Psi_0 \\ \Psi_1 \\ \Psi_2 \\ \Psi_3 \end{pmatrix} = \begin{pmatrix} \varphi \\ \chi \end{pmatrix} \quad \text{with } \varphi = \begin{pmatrix} \Psi_0 \\ \Psi_1 \end{pmatrix} \quad \text{and} \quad \chi = \begin{pmatrix} \Psi_2 \\ \Psi_3 \end{pmatrix} \quad (\text{III-4})$$

Dirac system breaks down into two coupled systems:

$$j \frac{\partial}{\partial(ct)} \begin{pmatrix} \Psi_0 \\ \Psi_1 \end{pmatrix} = \eta \begin{pmatrix} \Psi_0 \\ \Psi_1 \end{pmatrix} - j \begin{pmatrix} \frac{\partial}{\partial z} & \frac{\partial}{\partial x} - j \frac{\partial}{\partial y} \\ \frac{\partial}{\partial x} + j \frac{\partial}{\partial y} & -\frac{\partial}{\partial z} \end{pmatrix} \begin{pmatrix} \Psi_2 \\ \Psi_3 \end{pmatrix} \quad (\text{III-5})$$

$$j \frac{\partial}{\partial(ct)} \begin{pmatrix} \Psi_2 \\ \Psi_3 \end{pmatrix} = -\eta \begin{pmatrix} \Psi_2 \\ \Psi_3 \end{pmatrix} - j \begin{pmatrix} \frac{\partial}{\partial z} & \frac{\partial}{\partial x} - j \frac{\partial}{\partial y} \\ \frac{\partial}{\partial x} + j \frac{\partial}{\partial y} & -\frac{\partial}{\partial z} \end{pmatrix} \begin{pmatrix} \Psi_0 \\ \Psi_1 \end{pmatrix} \quad (\text{III-6})$$

in which it can be shown that the quantities to both φ and χ components behave as mathematical objects known as spinors.

A relief of writing is proposed by calling M the matrix:

$$M = -j \begin{pmatrix} \frac{\partial}{\partial z} & \frac{\partial}{\partial x} - j \frac{\partial}{\partial y} \\ \frac{\partial}{\partial x} + j \frac{\partial}{\partial y} & -\frac{\partial}{\partial z} \end{pmatrix} \quad (\text{III-7})$$

The coupled system (III-5, 6) in terms of spinors then takes the form:

$$\begin{aligned} j \frac{\partial}{\partial(ct)} (\varphi) &= \eta(\varphi) + M(\chi) \\ j \frac{\partial}{\partial(ct)} (\chi) &= -\eta(\chi) + M(\varphi) \end{aligned} \quad (\text{III-8})$$

or still:

$$j \frac{\partial}{\partial(ct)} \begin{pmatrix} \varphi \\ \chi \end{pmatrix} = \begin{pmatrix} \eta & M \\ M & -\eta \end{pmatrix} \begin{pmatrix} \varphi \\ \chi \end{pmatrix} \quad (\text{III-9})$$

In the absence of additional assumptions, it is impossible to go forward in the writing of the cross-relationships between two spinors.

On the other hand, we can seek special solutions with a temporal dependence in $\exp(-j\omega t)$ with a total energy $E = \hbar\omega$, hence a time in dependence $\exp\left(-j \frac{E}{\hbar} t\right)$.

It is important to note that the j which appears in the exponential has nothing to do with the complex representation of a physical wave in $\cos(\omega t)$. Its physical meaning already discussed in chapter I, will be detailed in the following chapters.

After taking account of the time derivation in the relationship (III-9), the coupled spinors system has the form:

$$\frac{E}{\hbar c} \begin{pmatrix} \varphi \\ \chi \end{pmatrix} = \begin{pmatrix} \eta & M \\ M & -\eta \end{pmatrix} \begin{pmatrix} \varphi \\ \chi \end{pmatrix} \quad (\text{III-10})$$

or again taking account of $\eta = \frac{mc}{\hbar}$:

$$E \begin{pmatrix} \varphi \\ \chi \end{pmatrix} = \begin{pmatrix} mc^2 & \hbar c M \\ \hbar c M & -mc^2 \end{pmatrix} \begin{pmatrix} \varphi \\ \chi \end{pmatrix} \quad (\text{III-11})$$

A rigorous writing should reveal the 2 dimension identity matrix multiplied by η for the relationship (III-10) terms and the terms in mc^2 for relationship (III-11). The use wants when this matrix is implied, as necessary for the coherence of the dimensions, it is not necessarily indicated for relief of writing.

I - Spinors for a not moving particle

It is possible to show solutions with positive energy that describe the particles of mass m , and negative energy solutions that describe the same anti-particles mass. To show this property, we treat the case of immobile particles: if x , y , z are fixed and constant, then the derivatives with respect to x , y , and z are zero, and the matrix M is zero. We derive from (III-11):

$$E \begin{pmatrix} \varphi \\ \chi \end{pmatrix} = \begin{pmatrix} mc^2 & 0 \\ 0 & -mc^2 \end{pmatrix} \begin{pmatrix} \varphi \\ \chi \end{pmatrix} \quad (\text{III-12})$$

or still:

$$\begin{aligned} E(\varphi) &= mc^2(\varphi) \\ E(\chi) &= -mc^2(\chi) \end{aligned} \quad (\text{III-13})$$

We identify that the energy associated with the spinor φ is equal to its mass energy, while the one associated with the spinor χ is equal to its opposite, and thus represents the energy associated with the anti-particle. We can deduce that the spinor φ is associated with the particle, while the spinor χ is associated with the antiparticle. The solution to the Dirac equation, which includes two spinors simultaneously, allows to describe both the behavior of the particle and its antiparticle.

In a very general manner, any spinor φ can be decomposed on the canonical basis in the following manner:

$$\phi = \begin{pmatrix} \psi_0 \\ \psi_1 \end{pmatrix} = \psi_0 \begin{pmatrix} 1 \\ 0 \end{pmatrix} + \psi_1 \begin{pmatrix} 0 \\ 1 \end{pmatrix} \quad (\text{III-14})$$

where the two independent spinors of the canonical basis are related to two possible spin for an electron States. The practice is that the spinor $\begin{pmatrix} 1 \\ 0 \end{pmatrix}$ is associated with the state of spin said «up» (\uparrow), and the spinor $\begin{pmatrix} 0 \\ 1 \end{pmatrix}$ is associated with the state of spin said «down» (\downarrow).

As for the particle with negative energy, we adopt the following decomposition:

$$\chi = \begin{pmatrix} \psi_1 \\ \psi_2 \end{pmatrix} = \psi_2 \begin{pmatrix} 1 \\ 0 \end{pmatrix} + \psi_3 \begin{pmatrix} 0 \\ 1 \end{pmatrix} \quad (\text{III-15})$$

and identify the two spin states described above.

Gathering these partial results, we can represent four distinct states of the DIRAC bi-spinors in the frame where the particle is at rest:

$$\psi = \begin{pmatrix} \psi_0 \\ \psi_1 \\ \psi_2 \\ \psi_3 \end{pmatrix} = \begin{pmatrix} 1 \\ 0 \\ 0 \\ 0 \end{pmatrix} \text{ Particle at positive energy in spin "up".}$$

$$\psi = \begin{pmatrix} \psi_0 \\ \psi_1 \\ \psi_2 \\ \psi_3 \end{pmatrix} = \begin{pmatrix} 0 \\ 1 \\ 0 \\ 0 \end{pmatrix} \text{ "Particle at positive energy in spin "down".}$$

$$\psi = \begin{pmatrix} \psi_0 \\ \psi_1 \\ \psi_2 \\ \psi_3 \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \\ 1 \\ 0 \end{pmatrix} \text{ Particle with negative energy in spin "up".}$$

$$\psi = \begin{pmatrix} \psi_0 \\ \psi_1 \\ \psi_2 \\ \psi_3 \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \\ 0 \\ 1 \end{pmatrix} \text{ Particle with negative energy in spin "down".}$$

II - Spinors for a particle in motion

When the particle is moving, the matrix M is not null, and we use the relationship (III-11) called for memory:

$$E \begin{pmatrix} \varphi \\ \chi \end{pmatrix} = \begin{pmatrix} mc^2 & \hbar c M \\ \hbar c M & -mc^2 \end{pmatrix} \begin{pmatrix} \varphi \\ \chi \end{pmatrix} \quad (\text{III-16})$$

This relationship shows that two spinors interact and that they are not independent of one another. As a necessary consequence, we can think that particle and antiparticle form a whole, and that one cannot move without the other is associated with this movement.

It comes, developing the matrix relationship (III-16):

$$\begin{aligned} E(\varphi) &= mc^2(\varphi) + \hbar c M(\chi) \\ E(\chi) &= \hbar c M(\varphi) - mc^2(\chi) \end{aligned} \quad (\text{III-17})$$

or still

$$\begin{aligned} (E - mc^2)(\varphi) &= \hbar c M(\chi) \\ (E + mc^2)(\chi) &= \hbar c M(\varphi) \end{aligned} \quad (\text{III-18})$$

or still:

$$\begin{aligned} (\varphi) &= \frac{\hbar c}{E - mc^2} M(\chi) \\ (\chi) &= \frac{\hbar c}{E + mc^2} M(\varphi) \end{aligned} \quad (\text{III-19})$$

These relationships indicate that in the special case where we fit, that is to say in the case where the 4 elements of the Dirac bispinors have a time in dependence $\exp\left(-j \frac{E}{\hbar} t\right)$, from the knowledge of one of the spinors, we can deduce the other.

It is possible to develop a more advanced formalism of relations (III-19) above. We must take again the matrix M and describe it on the basis of the PAULI matrices called for memory:

$$\sigma_1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \quad \sigma_2 = \begin{pmatrix} 0 & -j \\ j & 0 \end{pmatrix} \quad \sigma_3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \quad (\text{III-20})$$

It gives:

$$M = -j \begin{pmatrix} \frac{\partial}{\partial z} & \frac{\partial}{\partial x} - j \frac{\partial}{\partial y} \\ \frac{\partial}{\partial x} + j \frac{\partial}{\partial y} & -\frac{\partial}{\partial z} \end{pmatrix} = -j \left(\frac{\partial}{\partial x} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} + \frac{\partial}{\partial y} \begin{pmatrix} 0 & -j \\ j & 0 \end{pmatrix} + \frac{\partial}{\partial z} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \right) \quad (\text{III-21})$$

We introduce by substituting the partial derivatives (III-21), the operator pulse of quantum mechanics, whose three components are explicitly:

$$\hat{p}_x = -j\hbar \frac{\partial}{\partial x} \quad \hat{p}_y = -j\hbar \frac{\partial}{\partial y} \quad \hat{p}_z = -j\hbar \frac{\partial}{\partial z} \quad (\text{III-22})$$

This allows to present the M matrix in the form:

$$\mathbf{M} = \frac{1}{\hbar} \left\{ \hat{p}_x \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} + \hat{p}_y \begin{pmatrix} 0 & -j \\ j & 0 \end{pmatrix} + \hat{p}_z \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \right\} \quad (\text{III-23})$$

It follows an abuse of writing that would give cold sweats to more than a teacher of mathematics:

We call:

$$\vec{\sigma} = \begin{pmatrix} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \\ \begin{pmatrix} 0 & -j \\ j & 0 \end{pmatrix} \\ \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \end{pmatrix} \quad (\text{III-24})$$

a "vector" containing the three PAULI matrices, and:

$$\vec{p} = \begin{pmatrix} -j\hbar \frac{\partial}{\partial x} \\ -j\hbar \frac{\partial}{\partial y} \\ -j\hbar \frac{\partial}{\partial z} \end{pmatrix} \quad (\text{III-25})$$

a "vector" which represents the pulse operator of the quantum mechanics.

Therefore the relation expressing the matrix M (III-23) may come under the following condensed form, using the usual formalism of the scalar product between two vectors:

$$\mathbf{M} = \frac{\vec{\sigma} \cdot \vec{p}}{\hbar} \quad (\text{III-26})$$

We can deduce the condensed writing which expresses the relationship between the two spinors which constitute the solution of DIRAC, in the event of a time in dependence

$$\exp\left(-j \frac{E}{\hbar} t\right) :$$

$$(\varphi) = \frac{\vec{\sigma} \cdot \vec{p} c}{E - mc^2} (\chi) \quad (\text{III-27})$$

$$(\chi) = \frac{\vec{\sigma} \cdot \vec{p} c}{E + mc^2} (\varphi)$$

By substitution of one relationship in the other, for example the second in the first, we obtain:

$$(\varphi) = \frac{\vec{\sigma} \cdot \vec{p}c}{E - mc^2} \frac{\vec{\sigma} \cdot \vec{p}c}{E + mc^2} (\varphi) = \frac{(\vec{\sigma} \cdot \vec{p}c)^2}{E^2 - (mc^2)^2} (\varphi) \quad (\text{III-28})$$

Using the property of the PAULI matrices whose square gives the identity matrix, yields finally:

$$(\varphi) = \frac{(\vec{p}c)^2}{E^2 - (mc^2)^2} (\varphi) \quad (\text{III-29})$$

This relationship will be checked only if:

$$(\vec{p}c)^2 = E^2 - (mc^2)^2 \quad (\text{III-30})$$

or again:

$$(\vec{p}c)^2 + (mc^2)^2 = E^2 \quad (\text{III-31})$$

This relationship confirms the consistency of solutions expressed with relativistic energy conservation, but also with the possibility of solutions with negative energies of the form:

$$E = \pm \sqrt{(\vec{p}c)^2 + (mc^2)^2} \quad (\text{III-32})$$

IV

Spin 1/2 of the electron

The solution in terms of spinors mentioned in the previous chapter suggests that there is within the electron something that spins, without knowing precisely what. We will therefore focus in this chapter on the kinematics of rotation, and try to put in relation with the solution to the system of DIRAC.

I - The concept of angular momentum

For a material point of mass m , located in M , and moving around an origin O , the kinetic momentum that will be designated by L is expressed as the pulse momentum:

$$\vec{L} = \overrightarrow{OM} \wedge \vec{p} = \overrightarrow{OM} \wedge m\vec{v} \quad (\text{IV-1})$$

It is a vector quantity which is brought by the axis of rotation.

If one denotes by x, y, z , the components of the position vector, and p_x, p_y, p_z , the components of the momentum vector:

$$\overrightarrow{OM} = \begin{pmatrix} x \\ y \\ z \end{pmatrix} \quad \vec{p} = \begin{pmatrix} p_x \\ p_y \\ p_z \end{pmatrix} \quad (\text{IV-2})$$

we get the components of the angular momentum in the form:

$$\begin{aligned} L_x &= yp_z - zp_y \\ L_y &= zp_x - xp_z \\ L_z &= xp_y - yp_x \end{aligned} \quad (\text{IV-3})$$

The transition to quantum mechanics imposes one overrides the "position" quantities and "quantity of motion" with operators "position" and "quantity of motion", which allows to define the components of the angular momentum operator in the form:

$$\begin{aligned} \hat{L}_x &= \hat{y}\hat{p}_z - \hat{z}\hat{p}_y \\ \hat{L}_y &= \hat{z}\hat{p}_x - \hat{x}\hat{p}_z \\ \hat{L}_z &= \hat{x}\hat{p}_y - \hat{y}\hat{p}_x \end{aligned} \quad (\text{IV-4})$$

where the definition of the position and pulse momentum, applied to a wave function ψ operators is recalled below:

$$\begin{aligned}\hat{p}_x(\psi) &= -j\hbar \frac{\partial\psi}{\partial x} & \hat{p}_y(\psi) &= -j\hbar \frac{\partial\psi}{\partial y} & \hat{p}_z(\psi) &= -j\hbar \frac{\partial\psi}{\partial z} \\ \hat{x}(\psi) &= x\psi & \hat{y}(\psi) &= y\psi & \hat{z}(\psi) &= z\psi\end{aligned}\quad (IV-5)$$

The description of the above angular momentum operator allows to treat without problem the orbital angular momentum of a particle of mass m which revolves around an origin, but she is distressed to describe the angular momentum of spin. This is an internal movement to the particle which is poorly known and we don't know how to express it in terms of the position and momentum operators.

To overcome this difficulty, it was necessary to look for a property of the kinetic moment which do not involve the position and momentum operators, and that contains sufficient information to describe the rotation.

This property is constructed from a function called switch whose we can give the following definition in quantum mechanics:

If A and B are two operators, the switch of A and B is defined by:

$$[A, B] = AB - BA \quad (IV-6)$$

For example, we can show that the operators position and speed applied to a wave function do not switch, which is to say that these two operators switch is non-zero:

$$\left[x, \frac{\partial}{\partial x} \right] (\psi) = x \left(\frac{\partial\psi}{\partial x} \right) - \frac{\partial}{\partial x} (x\psi) = x \left(\frac{\partial\psi}{\partial x} \right) - \psi - x \left(\frac{\partial\psi}{\partial x} \right) = -\psi \quad (IV-7)$$

If we built switches of angular momentum operators (IV-4), they have the following property:

$$\begin{aligned}\left[\hat{L}_x, \hat{L}_y \right] (\psi) &= (\hat{L}_x \hat{L}_y - \hat{L}_y \hat{L}_x) (\psi) = j\hbar \hat{L}_z (\psi) \\ \left[\hat{L}_z, \hat{L}_x \right] (\psi) &= (\hat{L}_z \hat{L}_x - \hat{L}_x \hat{L}_z) (\psi) = j\hbar \hat{L}_y (\psi) \\ \left[\hat{L}_y, \hat{L}_z \right] (\psi) &= (\hat{L}_y \hat{L}_z - \hat{L}_z \hat{L}_y) (\psi) = j\hbar \hat{L}_x (\psi)\end{aligned}\quad (IV-8)$$

The demonstration is obtained directly from the definition of angular momentum operators (IV-4). On the example of the first line of (IV - 8) yields successively:

$$\hat{L}_x (\psi) = -j\hbar \left(y \frac{\partial\psi}{\partial z} - z \frac{\partial\psi}{\partial y} \right) \quad \hat{L}_y (\psi) = -j\hbar \left(z \frac{\partial\psi}{\partial x} - x \frac{\partial\psi}{\partial z} \right) \quad (IV-9)$$

$$\hat{L}_y (\hat{L}_x (\psi)) = -\hbar^2 \left\{ z \frac{\partial}{\partial x} \left(y \frac{\partial\psi}{\partial z} - z \frac{\partial\psi}{\partial y} \right) - x \frac{\partial}{\partial z} \left(y \frac{\partial\psi}{\partial z} - z \frac{\partial\psi}{\partial y} \right) \right\} \quad (IV-10)$$

$$\hat{L}_y(\hat{L}_x(\psi)) = -\hbar^2 \left\{ zy \frac{\partial^2 \psi}{\partial x \partial z} - z^2 \frac{\partial^2 \psi}{\partial x \partial y} - xy \frac{\partial^2 \psi}{\partial z^2} + x \frac{\partial \psi}{\partial y} + xz \frac{\partial^2 \psi}{\partial z \partial y} \right\} \quad (\text{IV-11})$$

and by reversing the order of operators:

$$\hat{L}_x(\hat{L}_y(\psi)) = -\hbar^2 \left\{ y \frac{\partial}{\partial z} \left(z \frac{\partial \psi}{\partial x} - x \frac{\partial \psi}{\partial z} \right) - z \frac{\partial}{\partial y} \left(z \frac{\partial \psi}{\partial x} - x \frac{\partial \psi}{\partial z} \right) \right\} \quad (\text{IV-12})$$

$$\hat{L}_x(\hat{L}_y(\psi)) = -\hbar^2 \left\{ y \frac{\partial \psi}{\partial x} + yz \frac{\partial^2 \psi}{\partial x \partial z} - yx \frac{\partial^2 \psi}{\partial z^2} - z^2 \frac{\partial^2 \psi}{\partial y \partial x} + xz \frac{\partial^2 \psi}{\partial z \partial y} \right\} \quad (\text{IV-13})$$

By subtraction of (IV-13) from (IV-11), yields the result presented in (IV-8):

$$\hat{L}_x(\hat{L}_y(\psi)) - \hat{L}_y(\hat{L}_x(\psi)) = -\hbar^2 \left\{ y \frac{\partial \psi}{\partial x} - x \frac{\partial \psi}{\partial y} \right\} = j\hbar \hat{L}_z \quad (\text{IV-14})$$

The property (IV-8) no longer depends on the position and momentum operators. Even if the image should be taken with care, one has somehow built a system of three equations and three unknowns which presents for these unknown similar properties to those that are generated by relations (IV-4) that define the angular momentum in quantum mechanics.

II - The spin angular momentum operator

The reasoning of the previous paragraph has enabled us to clarify the relationships that define an operator of angular momentum for a wave function ψ . This operator is defined from the relationships of commutation recalled to memory:

$$\begin{aligned} [\hat{L}_x, \hat{L}_y](\psi) &= (\hat{L}_x \hat{L}_y - \hat{L}_y \hat{L}_x)(\psi) = j\hbar \hat{L}_z(\psi) \\ [\hat{L}_z, \hat{L}_x](\psi) &= (\hat{L}_z \hat{L}_x - \hat{L}_x \hat{L}_z)(\psi) = j\hbar \hat{L}_y(\psi) \\ [\hat{L}_y, \hat{L}_z](\psi) &= (\hat{L}_y \hat{L}_z - \hat{L}_z \hat{L}_y)(\psi) = j\hbar \hat{L}_x(\psi) \end{aligned} \quad (\text{IV-15})$$

Solutions issue from the formulation of DIRAC appear in the form of spinors, that is to say in the form of a mathematical being containing two wave functions. In the previous chapter, we showed for example that the solution to positive energy was represented by the spinor:

$$\varphi = \begin{pmatrix} \psi_0 \\ \psi_1 \end{pmatrix} \quad (\text{IV-16})$$

The question that arises is: how can we generalize the (IV-15) relationships that define the operator of angular momentum for a wave function, in order to define an operator of angular momentum for a spinor that contains not one but two wave functions?

In other words, if S is one such operator, we want to be able to write:

$$\begin{aligned}
 [\hat{S}_x, \hat{S}_y](\varphi) &= (\hat{S}_x \hat{S}_y - \hat{S}_y \hat{S}_x)(\varphi) = j\hbar \hat{S}_z(\varphi) \\
 [\hat{S}_z, \hat{S}_x](\varphi) &= (\hat{S}_z \hat{S}_x - \hat{S}_x \hat{S}_z)(\varphi) = j\hbar \hat{S}_y(\varphi) \\
 [\hat{S}_y, \hat{S}_z](\varphi) &= (\hat{S}_y \hat{S}_z - \hat{S}_z \hat{S}_y)(\varphi) = j\hbar \hat{S}_x(\varphi)
 \end{aligned} \tag{IV-17}$$

or still:

$$\begin{aligned}
 [\hat{S}_x, \hat{S}_y] \begin{pmatrix} \Psi_0 \\ \Psi_1 \end{pmatrix} &= (\hat{S}_x \hat{S}_y - \hat{S}_y \hat{S}_x) \begin{pmatrix} \Psi_0 \\ \Psi_1 \end{pmatrix} = j\hbar \hat{S}_z \begin{pmatrix} \Psi_0 \\ \Psi_1 \end{pmatrix} \\
 [\hat{S}_z, \hat{S}_x] \begin{pmatrix} \Psi_0 \\ \Psi_1 \end{pmatrix} &= (\hat{S}_z \hat{S}_x - \hat{S}_x \hat{S}_z) \begin{pmatrix} \Psi_0 \\ \Psi_1 \end{pmatrix} = j\hbar \hat{S}_y \begin{pmatrix} \Psi_0 \\ \Psi_1 \end{pmatrix} \\
 [\hat{S}_y, \hat{S}_z] \begin{pmatrix} \Psi_0 \\ \Psi_1 \end{pmatrix} &= (\hat{S}_y \hat{S}_z - \hat{S}_z \hat{S}_y) \begin{pmatrix} \Psi_0 \\ \Psi_1 \end{pmatrix} = j\hbar \hat{S}_x \begin{pmatrix} \Psi_0 \\ \Psi_1 \end{pmatrix}
 \end{aligned} \tag{IV-18}$$

It appears that operators $\hat{S}_x, \hat{S}_y, \hat{S}_z$ can no longer be defined as simple operators used for a single wave function: they must be made by a matrix 2 X 2 of operators and these matrices must check the characteristic relations of angular momentum operators:

$$\begin{aligned}
 [\hat{S}_x, \hat{S}_y] &= (\hat{S}_x \hat{S}_y - \hat{S}_y \hat{S}_x) = j\hbar \hat{S}_z \\
 [\hat{S}_z, \hat{S}_x] &= (\hat{S}_z \hat{S}_x - \hat{S}_x \hat{S}_z) = j\hbar \hat{S}_y \\
 [\hat{S}_y, \hat{S}_z] &= (\hat{S}_y \hat{S}_z - \hat{S}_z \hat{S}_y) = j\hbar \hat{S}_x
 \end{aligned} \tag{IV-19}$$

The PAULI matrices, mentioned below for memory, are good candidates for this role:

$$\sigma_1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \quad \sigma_2 = \begin{pmatrix} 0 & -j \\ j & 0 \end{pmatrix} \quad \sigma_3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \tag{IV-20}$$

By multiplying by the quantity \hbar , the first switch (IV-19) is written:

$$\begin{aligned}
 [\hbar\sigma_1, \hbar\sigma_2] &= \hbar \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} j\hbar \begin{pmatrix} 0 & -j \\ j & 0 \end{pmatrix} - \hbar \begin{pmatrix} 0 & -j \\ j & 0 \end{pmatrix} j\hbar \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \\
 [\hbar\sigma_1, \hbar\sigma_2] &= \hbar^2 \left\{ \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 0 & -j \\ j & 0 \end{pmatrix} - \begin{pmatrix} 0 & -j \\ j & 0 \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \right\} \\
 [\hbar\sigma_1, \hbar\sigma_2] &= \hbar^2 \left\{ \begin{pmatrix} j & 0 \\ 0 & -j \end{pmatrix} - \begin{pmatrix} -j & 0 \\ 0 & j \end{pmatrix} \right\} = 2j\hbar^2 \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}
 \end{aligned} \tag{IV-21}$$

or still:

$$[\hbar\sigma_1, \hbar\sigma_2] = 2j\hbar(\hbar\sigma_3) \tag{IV-22}$$

We don't find exactly the relationship of switching expressed in (IV - 19), since there is a factor of 2 which takes place. If we want to find exactly the relationship that defines the components of an angular momentum, we must necessarily introduce a factor $\frac{1}{2}$ in the PAULI matrices. It is this factor that will induce the spin $\frac{1}{2}$ of the electron.

In summary, the components $\hat{S}_x, \hat{S}_y, \hat{S}_z$ sought, and which therefore verifies (IV-19), are as follows:

$$\begin{aligned}\hat{S}_x &= \frac{\hbar}{2} \sigma_1 = \frac{\hbar}{2} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \\ \hat{S}_y &= \frac{\hbar}{2} \sigma_2 = \frac{\hbar}{2} \begin{pmatrix} 0 & -j \\ j & 0 \end{pmatrix} \\ \hat{S}_z &= \frac{\hbar}{2} \sigma_3 = \frac{\hbar}{2} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}\end{aligned}\tag{IV-23}$$

If we have well advanced in the formalism for describing the spin angular momentum, we are away considerably from its physical meaning. The passage, which seems obvious, from relations (IV-3) to (IV-4), is a passage where relations (IV-3) have a real physical meaning, while relations (IV-4) have already no more because they are supported by operators. Their generalization to spinors even increases the level of abstraction, and distance with the physical representation of the angular momentum that is used in classical mechanics.

III - The interpretation of the angular momentum in terms of rotation operator

The description of rotation takes meaning only if it allows to specify the angle of rotation. We will therefore focus in this part to show how one can express this angle on the basis of the elements that were used to characterize this rotation in the preceding paragraphs, and in particular the components of the angular momentum.

The problem is not simple, because it is necessary to characterize not a classic rotation in a three-dimensional space, but a rotation operator that acts on a wave function ψ in a first time, and on a spinor φ with two components in a second time.

We must first establish the matrix of a rotation of angle θ in an Euclidean space, whose axis is chosen arbitrarily toward Oz. This rotation belong to the xOy plane, and we represent below the rotation between two points M and M' with coordinates M(x,y) and M'(x',y').

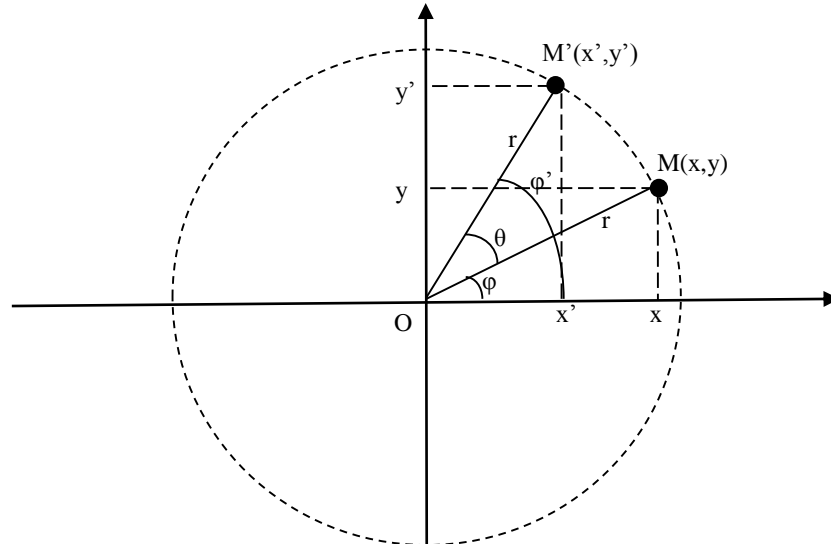


Figure (IV-1): Rotation in the xOy plane

The relationship between the coordinates (x, y) and (x', y') may be established geometrically as follows:

$$\begin{aligned} x &= r \cdot \cos(\varphi) \\ y &= r \cdot \sin(\varphi) \end{aligned} \tag{IV - 24}$$

$$\begin{aligned} x' &= r \cdot \cos(\varphi') = r \cdot \cos(\theta + \varphi) = r (\cos(\theta) \cos(\varphi) - \sin(\theta) \sin(\varphi)) = x \cdot \cos(\theta) - y \cdot \sin(\theta) \\ y' &= r \cdot \sin(\varphi') = r \cdot \sin(\theta + \varphi) = r (\sin(\theta) \cos(\varphi) + \cos(\theta) \sin(\varphi)) = x \cdot \sin(\theta) + y \cdot \cos(\theta) \end{aligned} \tag{IV-25}$$

or still with a matrix written in three dimensions:

$$\begin{pmatrix} x' \\ y' \\ z' \end{pmatrix} = \begin{pmatrix} \cos \theta & -\sin \theta & 0 \\ \sin \theta & \cos \theta & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} x \\ y \\ z \end{pmatrix} \tag{IV-26}$$

For an opposite angle rotation $-\theta$, the matrix is obtained by changing the sign of sinus:

$$\begin{pmatrix} x' \\ y' \\ z' \end{pmatrix} = \begin{pmatrix} \cos \theta & \sin \theta & 0 \\ -\sin \theta & \cos \theta & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} x \\ y \\ z \end{pmatrix} \tag{IV-27}$$

In the description of the rotation operator applied to a given function, if we assimilate the function to an object to make a concrete representation, a rotation of the axes to the right can be considered a rotation equivalent of the object to the left.

For an operator of angle rotation θ , everything happens as if the coordinates used by the function were under a rotation of angle $-\theta$.

Taking into account these elements, we can specify what is the rotation operator applied to a function ψ , and we will adopt the following definition for an Oz axis rotation of angle θ and applied to a wave function ψ :

$$\hat{R}_{z,\theta}\psi(x, y, z) = \psi(x', y', z') \quad (\text{IV-28})$$

In this relationship, in accordance with the previous comments, we use the transformation of coordinates (IV-21).

The next step is to establish a link between this rotation operator, and the angular momentum operator $\hat{L} = (\hat{L}_x, \hat{L}_y, \hat{L}_z)$ which was used to characterize the movement of rotation in (IV-4), inspired from relations (IV-3) of classical mechanics. This link will be formalized in a first time, for an infinitesimal rotation.

Starting from the definition of the operator rotation (IV-28) given above:

$$\hat{R}_{z,\theta}\psi(x, y, z) = \psi(x \cdot \cos \theta + y \cdot \sin \theta, -x \cdot \sin \theta + y \cdot \cos \theta, z) \quad (\text{IV-29})$$

the introduction of the developments limited to first order of $\sin(\theta)$ and $\cos(\theta)$ leads to the infinitesimal rotation operator of angle $\delta\theta$:

$$\hat{R}_{z,\delta\theta}\psi(x, y, z) = \psi(x + y \cdot \delta\theta, -x \cdot \delta\theta + y, z) \quad (\text{IV-30})$$

Making use of the development limited to first order of a function of two variables x and y which vary from an infinitesimal amount respectively equal to ε and η :

$$f(x + \varepsilon, y + \eta) = f(x, y) + \varepsilon \frac{\partial f(x, y)}{\partial x} + \eta \frac{\partial f(x, y)}{\partial y} \quad (\text{IV-31})$$

the infinitesimal rotation operator (IV-30) goes in the form:

$$\hat{R}_{z,\delta\theta}\psi(x, y, z) = \psi(x, y, z) + y(\delta\theta) \frac{\partial \psi}{\partial x} - x(\delta\theta) \frac{\partial \psi}{\partial y} \quad (\text{IV-32})$$

It is interpreted using the operators "position" and "pulse" recalled to memory below:

$$\begin{aligned} \hat{p}_x(\psi) &= -j\hbar \frac{\partial \psi}{\partial x} & \hat{p}_y(\psi) &= -j\hbar \frac{\partial \psi}{\partial y} & \hat{p}_z(\psi) &= -j\hbar \frac{\partial \psi}{\partial z} \\ \hat{x}(\psi) &= x\psi & \hat{y}(\psi) &= y\psi & \hat{z}(\psi) &= z\psi \end{aligned} \quad (\text{IV-33})$$

to give the following expression:

$$\hat{R}_{z,\delta\theta}\psi(x, y, z) = \psi(x, y, z) - j \frac{\delta\theta}{\hbar} (\hat{x}\hat{p}_y - \hat{y}\hat{p}_x) \psi(x, y, z) \quad (\text{IV-34})$$

We recognize the operator of angular momentum (IV-4):

$$\hat{L}_z = \hat{x}\hat{p}_y - \hat{y}\hat{p}_x \quad (\text{IV-35})$$

We deduce the expression of the infinitesimal rotation operator around Oz axis and angle θ , applied to a wave function ψ , depending on the angular momentum operator:

$$\hat{R}_{z,\delta\theta}\psi(x,y,z) = \psi(x,y,z) - j\frac{\delta\theta}{\hbar}\hat{L}_z\psi(x,y,z) = \left\{ \hat{I} - j\frac{\delta\theta}{\hbar}\hat{L}_z \right\} \psi(x,y,z) \quad (IV-36)$$

where \hat{I} denotes the identity operator.

We can now deduce the intrinsic expression of the infinitesimal rotation operator:

$$\hat{R}_{z,\delta\theta} = \hat{I} - j\frac{\delta\theta}{\hbar}\hat{L}_z \quad (IV-37)$$

This relationship can be generalized to a rotation of angle θ , in any of several ways. One of the simplest is to divide this angle by an integer N , which tends towards infinity: we can assimilate the infinitely small angle $\delta\theta$ in the relationship (IV-37) with the angle θ/N .

$$\hat{R}_{z,\frac{\theta}{N}} = \hat{I} - \frac{1}{N} \left(j\frac{\theta}{\hbar}\hat{L}_z \right) \quad (IV-38)$$

One then writes that for a rotation of angle θ , we must applied N times the infinitesimal rotation of angle θ/N :

$$\hat{R}_{z,\theta} = \left(\hat{R}_{z,\frac{\theta}{N}} \right)^N = \left(\hat{I} - \frac{1}{N} \left(j\frac{\theta}{\hbar}\hat{L}_z \right) \right)^N \quad (IV-39)$$

And we can now write a transition to the limit:

$$\hat{R}_{z,\theta} = \lim_{N \rightarrow \infty} \left(\hat{R}_{z,\frac{\theta}{N}} \right)^N = \lim_{N \rightarrow \infty} \left(\hat{I} - \frac{1}{N} \left(j\frac{\theta}{\hbar}\hat{L}_z \right) \right)^N \quad (IV-40)$$

Using the known result:

$$\lim_{N \rightarrow \infty} \left(1 + \frac{x}{N} \right)^N = \exp(x) \quad (IV-41)$$

we finally get the expression of the rotation of axis Oz and angle θ operator, according to the angular momentum operator \hat{L}_z :

$$\hat{R}_{z,\theta} = \exp \left(j\frac{\theta}{\hbar}\hat{L}_z \right) \quad (IV-42)$$

It remains to conclude this chapter, to generalize this operator to rotation of spinors which are functions of waves in two dimensions.

It has been shown above that the operator $\hat{L} = (\hat{L}_x, \hat{L}_y, \hat{L}_z)$ should be replaced by a matrix operator $\hat{S} = (\hat{S}_x, \hat{S}_y, \hat{S}_z)$ able to act on mathematical two-dimensional beings, and that this operator is inferred from the PAULI matrices following relationship (IV-23) called for memory:

$$\begin{aligned}\hat{S}_x &= \frac{\hbar}{2} \sigma_1 = \frac{\hbar}{2} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \\ \hat{S}_y &= \frac{\hbar}{2} \sigma_2 = \frac{\hbar}{2} \begin{pmatrix} 0 & -j \\ j & 0 \end{pmatrix} \\ \hat{S}_z &= \frac{\hbar}{2} \sigma_3 = \frac{\hbar}{2} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}\end{aligned}\tag{IV-43}$$

The operator of rotation of axis Oz and angle θ , able to act on a spinor, therefore has the following form:

$$\hat{R}_{z,\theta} = \exp\left(j \frac{\theta}{\hbar} \hat{S}_z\right)\tag{IV-44}$$

or still:

$$\hat{R}_{z,\theta} = \exp\left(j \frac{\theta}{2} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}\right)\tag{IV-45}$$

From a practical point of view, we can rarely put any exponential matrix in the form of a 2 X 2 matrix whose each term is exactly known. We may nevertheless obtained an approximate solution using the development in series of the exponential function. For any matrix M, this gives:

$$\exp(M) = \sum_{k=0}^{\infty} \frac{M^k}{k!} = 1 + M + \frac{M^2}{2!} + \frac{M^3}{3!} + \dots\tag{IV-46}$$

However, it is possible to obtain exactly the operator matrix of the rotation around the main axis Ox, Oy, and Oz.

Operator of rotation around the Oz axis:

The operator matrix is diagonal, which allows to give the explicit form of the exponential matrix:

$$\hat{R}_{z,\theta} = \exp\left(j \frac{\theta}{2} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}\right) = \begin{pmatrix} \exp\left(j \frac{\theta}{2}\right) & 0 \\ 0 & \exp\left(-j \frac{\theta}{2}\right) \end{pmatrix} = \begin{pmatrix} \cos\left(\frac{\theta}{2}\right) + j \sin\left(\frac{\theta}{2}\right) & 0 \\ 0 & \cos\left(\frac{\theta}{2}\right) - j \sin\left(\frac{\theta}{2}\right) \end{pmatrix}\tag{IV-47}$$

Operator of rotation around the Oy axis:

$$\hat{R}_{y,\theta} = \exp\left(j\frac{\theta}{\hbar}\hat{S}_y\right) = \exp\left(j\frac{\theta}{2}\begin{pmatrix} 0 & -j \\ j & 0 \end{pmatrix}\right) = \exp\left(\frac{\theta}{2}\begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}\right) \quad (\text{IV-48})$$

The particular shape of the 2 X 2 matrix allows to show that series of the exponential expansion gives the development in series of cos and sin functions in terms of the Matrix result following the relationship:

$$A = \begin{pmatrix} 0 & x \\ -x & 0 \end{pmatrix} \quad \exp(A) = \begin{pmatrix} \cos(x) & \sin(x) \\ -\sin(x) & \cos(x) \end{pmatrix} \quad (\text{IV-49})$$

Yields the expression of the operator of rotation about the Oy axis:

$$\hat{R}_{y,\theta} = \begin{pmatrix} \cos\left(\frac{\theta}{2}\right) & \sin\left(\frac{\theta}{2}\right) \\ -\sin\left(\frac{\theta}{2}\right) & \cos\left(\frac{\theta}{2}\right) \end{pmatrix} \quad (\text{IV-50})$$

Operator of rotation around Ox axis:

$$\hat{R}_{x,\theta} = \exp\left(j\frac{\theta}{\hbar}\hat{S}_x\right) = \exp\left(j\frac{\theta}{2}\begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}\right) \quad (\text{IV-51})$$

As in the previous example, the particular shape of the 2 X 2 matrix allows to show that the series of the exponential expansion give the hyperbolic functions cosh and sinh series expansion in the terms of the Matrix result following the relationship:

$$A = \begin{pmatrix} 0 & x \\ x & 0 \end{pmatrix} \quad \exp(A) = \begin{pmatrix} \cosh(x) & \sinh(x) \\ \sinh(x) & \cosh(x) \end{pmatrix} \quad (\text{IV-52})$$

Yields the expression of the operator of rotation around Ox axis:

$$\hat{R}_{x,\theta} = \begin{pmatrix} \text{ch}\left(j\frac{\theta}{2}\right) & \text{sh}\left(j\frac{\theta}{2}\right) \\ \text{sh}\left(j\frac{\theta}{2}\right) & \text{ch}\left(j\frac{\theta}{2}\right) \end{pmatrix} = \begin{pmatrix} \cos\left(\frac{\theta}{2}\right) & j\sin\left(\frac{\theta}{2}\right) \\ j\sin\left(\frac{\theta}{2}\right) & \cos\left(\frac{\theta}{2}\right) \end{pmatrix} \quad (\text{IV-53})$$

V

Covariance of the DIRAC equation

The DIRAC equation on which we relied to establish the system of differential equations has the form:

$$j \left[\gamma_0 \left(\frac{\partial}{\partial(ct)} \right) + \gamma_1 \left(\frac{\partial}{\partial x} \right) + \gamma_2 \left(\frac{\partial}{\partial y} \right) + \gamma_3 \left(\frac{\partial}{\partial z} \right) \right] (\psi) = \frac{m_0 c}{\hbar} (\psi) \quad (V-1)$$

By adopting the notation:

$$x^0 = ct, x^1 = x, x^2 = y, x^3 = z \quad (V-2)$$

we can condense the writing of (V-1) using the summation over indices rule:

$$\left(j \gamma^\mu \frac{\partial}{\partial x^\mu} - \frac{m_0 c}{\hbar} \right) \psi(x^\mu) = 0 \quad \mu = 0, 1, 2, 3 \quad (V-3)$$

It turns out that this equation is covariant, it means that she keeps the same shape when changing frame as it is given by special relativity. The demonstration, a bit laborious, may be found in references.

The invariance of the laws of physics by change of frame is probably the most important criterion on the general validity of a physical law. It can be summarized as follows:

Let us consider a frame in which the DIRAC equation is written in the form (V-3) above. If we change the frame, it will induce a change of spatial coordinates and time x^μ in each of the spacetime points that will become x'^μ ; and this will induce a change in the wave function ψ which will become ψ' .

The invariance of the laws of physics by frame change requires that the DIRAC equation is written in the frame (R'), under the form:

$$\left(j \gamma^\mu \frac{\partial}{\partial x'^\mu} - \frac{m_0 c}{\hbar} \right) \psi'(x'^\mu) = 0 \quad (V-4)$$

Another very general law of physics is the law of local conservation. It expresses the fact that when a physical quantity evolves in time and space, the conservation of this quantity is expressed using a four divergence equal to zero. One of the best-known examples is the conservation of the electric load that is obtained from the four-vector current density ($\rho c, j_x, j_y, j_z$) under the form:

$$\frac{\partial(\rho c)}{\partial(ct)} + \frac{\partial j_x}{\partial x} + \frac{\partial j_y}{\partial y} + \frac{\partial j_z}{\partial z} = 0 \quad (\text{V-5})$$

It is possible to show that there is a quantity kept by the solutions of the DIRAC equation, and this quantity is called current of DIRAC in reference to the relationship (V-5) above.

If we call J this quantity, the condensed writing of a four-divergence equal to zero is as follows:

$$\frac{\partial}{\partial x^\mu} J^\mu = \partial_\mu J^\mu = 0 \quad (\text{V-6})$$

A few non-trivial manipulations detailed in many courses allow to switch from the DIRAC equation into the following relationship:

$$\frac{\partial}{\partial x^\mu} (\bar{\psi} \gamma^\mu \psi) = 0 \quad (\text{V-7})$$

in which:

$$\bar{\psi} = (\psi^*)^T \gamma^0 = (\psi_0^*, \psi_1^*, \psi_2^*, \psi_3^*) \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix} = (\psi_0^*, \psi_1^*, -\psi_2^*, -\psi_3^*) \quad (\text{V-8})$$

From (V-6) and (V-7), we can deduce the four-vector current of DIRAC:

$$J^\mu = \bar{\psi} \gamma^\mu \psi \quad (\text{V-9})$$

where γ^μ represent matrices of DIRAC introduced in chapter II.

We can give the explicit expressions of each of the components of the four-vector J^μ representing s the DIRAC's currents:

index component 0:

$$J^0 = \bar{\psi} \gamma^0 \psi = (\psi_0^*, \psi_1^*, -\psi_2^*, -\psi_3^*) \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix} \begin{pmatrix} \psi_0 \\ \psi_1 \\ \psi_2 \\ \psi_3 \end{pmatrix} = (\psi_0^*, \psi_1^*, \psi_2^*, \psi_3^*) \begin{pmatrix} \psi_0 \\ \psi_1 \\ \psi_2 \\ \psi_3 \end{pmatrix} \quad (\text{V-10})$$

$$J^0 = \psi_0^* \psi_0 + \psi_1^* \psi_1 + \psi_2^* \psi_2 + \psi_3^* \psi_3 \quad (\text{V-11})$$

index component 1:

$$J^1 = \bar{\psi}\gamma^1\psi = (\psi_0^*, \psi_1^*, -\psi_2^*, -\psi_3^*) \begin{pmatrix} 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \\ 0 & -1 & 0 & 0 \\ -1 & 0 & 0 & 0 \end{pmatrix} \begin{pmatrix} \psi_0 \\ \psi_1 \\ \psi_2 \\ \psi_3 \end{pmatrix} = (\psi_3^*, \psi_2^*, \psi_1^*, \psi_0^*) \begin{pmatrix} \psi_0 \\ \psi_1 \\ \psi_2 \\ \psi_3 \end{pmatrix} \quad (V-12)$$

$$J^1 = \psi_3^*\psi_0 + \psi_2^*\psi_1 + \psi_1^*\psi_2 + \psi_0^*\psi_3 \quad (V-13)$$

index component 2:

$$J^2 = \bar{\psi}\gamma^2\psi = (\psi_0^*, \psi_1^*, -\psi_2^*, -\psi_3^*) \begin{pmatrix} 0 & 0 & 0 & -j \\ 0 & 0 & j & 0 \\ 0 & j & 0 & 0 \\ -j & 0 & 0 & 0 \end{pmatrix} \begin{pmatrix} \psi_0 \\ \psi_1 \\ \psi_2 \\ \psi_3 \end{pmatrix} = (j\psi_3^*, -j\psi_2^*, j\psi_1^*, -j\psi_0^*) \begin{pmatrix} \psi_0 \\ \psi_1 \\ \psi_2 \\ \psi_3 \end{pmatrix} \quad (V-14)$$

$$J^2 = j\psi_3^*\psi_0 - j\psi_2^*\psi_1 + j\psi_1^*\psi_2 - j\psi_0^*\psi_3 \quad (V-15)$$

index component 3:

$$J^3 = \bar{\psi}\gamma^3\psi = (\psi_0^*, \psi_1^*, -\psi_2^*, -\psi_3^*) \begin{pmatrix} 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & -1 \\ -1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \end{pmatrix} \begin{pmatrix} \psi_0 \\ \psi_1 \\ \psi_2 \\ \psi_3 \end{pmatrix} = (\psi_2^*, -\psi_3^*, \psi_0^*, -\psi_1^*) \begin{pmatrix} \psi_0 \\ \psi_1 \\ \psi_2 \\ \psi_3 \end{pmatrix} \quad (V-16)$$

$$J^3 = \psi_2^*\psi_0 - \psi_3^*\psi_1 + \psi_0^*\psi_2 - \psi_1^*\psi_3 \quad (V-17)$$

Hence in summary:

$$\begin{aligned} J^0 &= \psi_0^*\psi_0 + \psi_1^*\psi_1 + \psi_2^*\psi_2 + \psi_3^*\psi_3 \\ J^1 &= \psi_3^*\psi_0 + \psi_2^*\psi_1 + \psi_1^*\psi_2 + \psi_0^*\psi_3 \\ J^2 &= j\psi_3^*\psi_0 - j\psi_2^*\psi_1 + j\psi_1^*\psi_2 - j\psi_0^*\psi_3 \\ J^3 &= \psi_2^*\psi_0 - \psi_3^*\psi_1 + \psi_0^*\psi_2 - \psi_1^*\psi_3 \end{aligned} \quad (V-18)$$

The J^μ represent the density currents of probability of presence of the particle.

Second part

Energy approach of Dirac equation and its exact solutions in Cartesian coordinates

VI

Stationary modes

In a general manner, one presents the possible solutions to the DIRAC equations under the form of plane waves of the type:

$$\psi(x, y, z, t) = A \exp \left[j(\omega t - \vec{k} \vec{r}) \right] \quad (\text{VI-1})$$

This relationship is characteristic of a wave which propagates to the angular frequency ω and along the wave vector \vec{k} . It does not describe exchanges of energy such as those that can arise in an electromagnetic cavity for example. It is also subject to question about the number $j = \sqrt{-1}$ which is present and has nothing to do with the complex formalism usual for the description of waves.

We may think to put into evidence stationary modes by summing two or more travelling solutions, but solutions (VI-1) to DIRAC system are never fully explained, which makes this possibility inoperative.

We have therefore to formalize a method of searching for stationary solutions to the DIRAC system recalled to memory:

$$\begin{aligned} \eta \psi_0 &= j \frac{\partial \psi_0}{\partial (ct)} + j \frac{\partial \psi_3}{\partial x} + \frac{\partial \psi_3}{\partial y} + j \frac{\partial \psi_2}{\partial z} \\ \eta \psi_1 &= j \frac{\partial \psi_1}{\partial (ct)} + j \frac{\partial \psi_2}{\partial x} - \frac{\partial \psi_2}{\partial y} - j \frac{\partial \psi_3}{\partial z} \\ \eta \psi_2 &= -j \frac{\partial \psi_2}{\partial (ct)} - j \frac{\partial \psi_1}{\partial x} - \frac{\partial \psi_1}{\partial y} - j \frac{\partial \psi_0}{\partial z} \\ \eta \psi_3 &= -j \frac{\partial \psi_3}{\partial (ct)} - j \frac{\partial \psi_0}{\partial x} + \frac{\partial \psi_0}{\partial y} + j \frac{\partial \psi_1}{\partial z} \end{aligned} \quad (\text{VI-2})$$

and we hypothesize that any wave function can be written as a linear combination of stationary modes which may be present in a three-dimensional cavity:

$$\begin{aligned}
 \psi_0 &= \{a_0 S_x S_y S_z + b_0 S_x C_y S_z + c_0 C_x S_y S_z + d_0 C_x C_y S_z + e_0 S_x S_y C_z + f_0 S_x C_y C_z + g_0 C_x S_y C_z + h_0 C_x C_y C_z\} C_t \\
 &+ \{i_0 S_x S_y S_z + j_0 S_x C_y S_z + k_0 C_x S_y S_z + l_0 C_x C_y S_z + m_0 S_x S_y C_z + n_0 S_x C_y C_z + o_0 C_x S_y C_z + p_0 C_x C_y C_z\} S_t \\
 \psi_1 &= \{a_1 S_x S_y S_z + b_1 S_x C_y S_z + c_1 C_x S_y S_z + d_1 C_x C_y S_z + e_1 S_x S_y C_z + f_1 S_x C_y C_z + g_1 C_x S_y C_z + h_1 C_x C_y C_z\} C_t \\
 &+ \{i_1 S_x S_y S_z + j_1 S_x C_y S_z + k_1 C_x S_y S_z + l_1 C_x C_y S_z + m_1 S_x S_y C_z + n_1 S_x C_y C_z + o_1 C_x S_y C_z + p_1 C_x C_y C_z\} S_t \\
 \psi_2 &= \{a_2 S_x S_y S_z + b_2 S_x C_y S_z + c_2 C_x S_y S_z + d_2 C_x C_y S_z + e_2 S_x S_y C_z + f_2 S_x C_y C_z + g_2 C_x S_y C_z + h_2 C_x C_y C_z\} C_t \\
 &+ \{i_2 S_x S_y S_z + j_2 S_x C_y S_z + k_2 C_x S_y S_z + l_2 C_x C_y S_z + m_2 S_x S_y C_z + n_2 S_x C_y C_z + o_2 C_x S_y C_z + p_2 C_x C_y C_z\} S_t \\
 \psi_3 &= \{a_3 S_x S_y S_z + b_3 S_x C_y S_z + c_3 C_x S_y S_z + d_3 C_x C_y S_z + e_3 S_x S_y C_z + f_3 S_x C_y C_z + g_3 C_x S_y C_z + h_3 C_x C_y C_z\} C_t \\
 &+ \{i_3 S_x S_y S_z + j_3 S_x C_y S_z + k_3 C_x S_y S_z + l_3 C_x C_y S_z + m_3 S_x S_y C_z + n_3 S_x C_y C_z + o_3 C_x S_y C_z + p_3 C_x C_y C_z\} S_t
 \end{aligned} \tag{VI-3}$$

In this definition, the following abbreviated notation has been used:

$$S_x = \sin(k_x x) \quad S_y = \sin(k_y y) \quad S_z = \sin(k_z z) \quad S_t = \sin(k_t ct) \tag{VI-4}$$

$$C_x = \cos(k_x x) \quad C_y = \cos(k_y y) \quad C_z = \cos(k_z z) \quad C_t = \cos(k_t ct) \tag{VI-5}$$

The wave vector is represented by its k_x, k_y, k_z components, while for a homogeneous notation and consistent with relativity, the product ωt has been replaced by the expression $k_t ct$, which allows to highlight the two four-vectors:

$$\begin{array}{l}
 \text{4-vector position:} \\
 \left(\begin{array}{c} x \\ y \\ z \\ ct \end{array} \right)
 \end{array}
 \quad
 \begin{array}{l}
 \text{4-vector wave:} \\
 \left(\begin{array}{c} k_x \\ k_y \\ k_z \\ k_t = \frac{\omega}{c} \end{array} \right)
 \end{array}
 \tag{VI-6}$$

Coefficients $a_i, b_i, c_i, d_i, e_i, f_i, g_i, h_i, i_i, j_i, k_i, l_i, m_i, n_i, o_i, p_i$, for $i = 0, 1, 2, 3$, are real constants that weigh each of the modes and will serve as an unknown in the search for the wave functions $\psi_0, \psi_1, \psi_2, \psi_3$ solutions of the DIRAC system.

This leads, for each equation of the DIRAC system, to express the partial derivatives of the wave functions $\psi_0, \psi_1, \psi_2, \psi_3$ and to formulate a homogeneous system of 16 equations for the coefficients $a_i, b_i, c_i, d_i, e_i, f_i, g_i, h_i, i_i, j_i, k_i, l_i, m_i, n_i, o_i, p_i$. The obtained global system will therefore be a homogeneous system of 64 equations with 64 unknowns.

Calculations are a bit laborious but without difficulties. They are presented in their entirety in the following lines where we remember having put $k_t = \omega/c$.

First equation of the DIRAC system: $j \frac{\partial \psi_0}{\partial(ct)} + j \frac{\partial \psi_3}{\partial x} + \frac{\partial \psi_3}{\partial y} + j \frac{\partial \psi_2}{\partial z} - \eta \psi_0 = 0$

$$\begin{aligned}
 j \frac{\partial \psi_0}{\partial(ct)} &= -jk_t \{a_0 S_x S_y S_z + b_0 S_x C_y S_z + c_0 C_x S_y S_z + d_0 C_x C_y S_z + e_0 S_x S_y C_z + f_0 S_x C_y C_z + g_0 C_x S_y C_z + h_0 C_x C_y C_z\} S_t \\
 &+ jk_t \{i_0 S_x S_y S_z + j_0 S_x C_y S_z + k_0 C_x S_y S_z + l_0 C_x C_y S_z + m_0 S_x S_y C_z + n_0 S_x C_y C_z + o_0 C_x S_y C_z + p_0 C_x C_y C_z\} C_t \\
 j \frac{\partial \psi_3}{\partial x} &= jk_x \{a_3 C_x S_y S_z + b_3 C_x C_y S_z - c_3 S_x S_y S_z - d_3 S_x C_y S_z + e_3 C_x S_y C_z + f_3 C_x C_y C_z - g_3 S_x S_y C_z - h_3 S_x C_y C_z\} C_t \\
 &+ jk_x \{i_3 C_x S_y S_z + j_3 C_x C_y S_z - k_3 S_x S_y S_z - l_3 S_x C_y S_z + m_3 C_x S_y C_z + n_3 C_x C_y C_z - o_3 S_x S_y C_z - p_3 S_x C_y C_z\} S_t \\
 \frac{\partial \psi_3}{\partial y} &= k_y \{a_3 S_x C_y S_z - b_3 S_x S_y S_z + c_3 C_x C_y S_z - d_3 C_x S_y S_z + e_3 S_x C_y C_z - f_3 S_x S_y C_z + g_3 C_x C_y C_z - h_3 C_x S_y C_z\} C_t \\
 &k_y \{i_3 S_x C_y S_z - j_3 S_x S_y S_z + k_3 C_x C_y S_z - l_3 C_x S_y S_z + m_3 S_x C_y C_z - n_3 S_x S_y C_z + o_3 C_x C_y C_z - p_3 C_x S_y C_z\} S_t \\
 j \frac{\partial \psi_2}{\partial z} &= jk_z \{a_2 S_x S_y C_z + b_2 S_x C_y C_z + c_2 C_x S_y C_z + d_2 C_x C_y C_z - e_2 S_x S_y S_z - f_2 S_x C_y S_z - g_2 C_x S_y S_z - h_2 C_x C_y S_z\} C_t \\
 &+ jk_z \{i_2 S_x S_y C_z + j_2 S_x C_y C_z + k_2 C_x S_y C_z + l_2 C_x C_y C_z - m_2 S_x S_y S_z - n_2 S_x C_y S_z - o_2 C_x S_y S_z - p_2 C_x C_y S_z\} S_t \\
 \eta \psi_0 &= \eta \{a_0 S_x S_y S_z + b_0 S_x C_y S_z + c_0 C_x S_y S_z + d_0 C_x C_y S_z + e_0 S_x S_y C_z + f_0 S_x C_y C_z + g_0 C_x S_y C_z + h_0 C_x C_y C_z\} C_t \\
 &+ \eta \{i_0 S_x S_y S_z + j_0 S_x C_y S_z + k_0 C_x S_y S_z + l_0 C_x C_y S_z + m_0 S_x S_y C_z + n_0 S_x C_y C_z + o_0 C_x S_y C_z + p_0 C_x C_y C_z\} S_t
 \end{aligned} \tag{VI-7}$$

We can deduce the homogeneous system associated with the first equation of DIRAC system:

$$\begin{aligned}
 -jk_t a_0 - jk_x k_3 - k_y j_3 - jk_z m_2 - \eta i_0 &= 0 \\
 -jk_t b_0 - jk_x l_3 + k_y i_3 - jk_z n_2 - \eta j_0 &= 0 \\
 -jk_t c_0 + jk_x i_3 - k_y l_3 - jk_z o_2 - \eta k_0 &= 0 \\
 -jk_t d_0 + jk_x j_3 + k_y k_3 - jk_z p_2 - \eta l_0 &= 0 \\
 -jk_t e_0 - jk_x o_3 - k_y n_3 + jk_z i_2 - \eta m_0 &= 0 \\
 -jk_t f_0 - jk_x p_3 + k_y m_3 + jk_z j_2 - \eta n_0 &= 0 \\
 -jk_t g_0 + jk_x m_3 - k_y p_3 + jk_z k_2 - \eta o_0 &= 0 \\
 -jk_t h_0 + jk_x n_3 + k_y o_3 + jk_z l_2 - \eta p_0 &= 0 \\
 jk_t i_0 - jk_x c_3 - k_y b_3 - jk_z e_2 - \eta a_0 &= 0 \\
 jk_t j_0 - jk_x d_3 + k_y a_3 - jk_z f_2 - \eta b_0 &= 0 \\
 jk_t k_0 + jk_x a_3 - k_y d_3 - jk_z g_2 - \eta c_0 &= 0 \\
 jk_t l_0 + jk_x b_3 + k_y c_3 - jk_z h_2 - \eta d_0 &= 0 \\
 jk_t m_0 - jk_x g_3 - k_y f_3 + jk_z a_2 - \eta e_0 &= 0 \\
 jk_t n_0 - jk_x h_3 + k_y e_3 + jk_z b_2 - \eta f_0 &= 0 \\
 jk_t o_0 + jk_x e_3 - k_y h_3 + jk_z c_2 - \eta g_0 &= 0 \\
 jk_t p_0 + jk_x f_3 + k_y g_3 + jk_z d_2 - \eta h_0 &= 0
 \end{aligned} \tag{VI-8}$$

Second equation of DIRAC system: $j \frac{\partial \psi_1}{\partial(ct)} + j \frac{\partial \psi_2}{\partial x} - \frac{\partial \psi_2}{\partial y} - j \frac{\partial \psi_3}{\partial z} - \eta \psi_1 = 0$

$$\begin{aligned}
 j \frac{\partial \psi_1}{\partial(ct)} &= -jk_t \{a_1 S_x S_y S_z + b_1 S_x C_y S_z + c_1 C_x S_y S_z + d_1 C_x C_y S_z + e_1 S_x S_y C_z + f_1 S_x C_y C_z + g_1 C_x S_y C_z + h_1 C_x C_y C_z\} S_t \\
 &+ jk_t \{i_1 S_x S_y S_z + j_1 S_x C_y S_z + k_1 C_x S_y S_z + l_1 C_x C_y S_z + m_1 S_x S_y C_z + n_1 S_x C_y C_z + o_1 C_x S_y C_z + p_1 C_x C_y C_z\} C_t \\
 j \frac{\partial \psi_2}{\partial x} &= jk_x \{a_2 C_x S_y S_z + b_2 C_x C_y S_z - c_2 S_x S_y S_z - d_2 S_x C_y S_z + e_2 C_x S_y C_z + f_2 C_x C_y C_z - g_2 S_x S_y C_z - h_2 S_x C_y C_z\} C_t \\
 &+ jk_x \{i_2 C_x S_y S_z + j_2 C_x C_y S_z - k_2 S_x S_y S_z - l_2 S_x C_y S_z + m_2 C_x S_y C_z + n_2 C_x C_y C_z - o_2 S_x S_y C_z - p_2 S_x C_y C_z\} S_t \\
 - \frac{\psi_2}{\partial y} &= -k_y \{a_2 S_x C_y S_z - b_2 S_x S_y S_z + c_2 C_x C_y S_z - d_2 C_x S_y S_z + e_2 S_x C_y C_z - f_2 S_x S_y C_z + g_2 C_x C_y C_z - h_2 C_x S_y C_z\} C_t \\
 &- k_y \{i_2 S_x C_y S_z - j_2 S_x S_y S_z + k_2 C_x C_y S_z - l_2 C_x S_y S_z + m_2 S_x C_y C_z - n_2 S_x S_y C_z + o_2 C_x C_y C_z - p_2 C_x S_y C_z\} S_t \\
 - j \frac{\partial \psi_3}{\partial z} &= -jk_z \{a_3 S_x S_y C_z + b_3 S_x C_y C_z + c_3 C_x S_y C_z + d_3 C_x C_y C_z - e_3 S_x S_y S_z - f_3 S_x C_y S_z - g_3 C_x S_y S_z - h_3 C_x C_y S_z\} C_t \\
 &- jk_z \{i_3 S_x S_y C_z + j_3 S_x C_y C_z + k_3 C_x S_y C_z + l_3 C_x C_y C_z - m_3 S_x S_y S_z - n_3 S_x C_y S_z - o_3 C_x S_y S_z - p_3 C_x C_y S_z\} S_t \\
 \eta \psi_1 &= \eta \{a_1 S_x S_y S_z + b_1 S_x C_y S_z + c_1 C_x S_y S_z + d_1 C_x C_y S_z + e_1 S_x S_y C_z + f_1 S_x C_y C_z + g_1 C_x S_y C_z + h_1 C_x C_y C_z\} C_t \\
 &+ \eta \{i_1 S_x S_y S_z + j_1 S_x C_y S_z + k_1 C_x S_y S_z + l_1 C_x C_y S_z + m_1 S_x S_y C_z + n_1 S_x C_y C_z + o_1 C_x S_y C_z + p_1 C_x C_y C_z\} S_t
 \end{aligned}
 \tag{VI-9}$$

We can deduce the homogeneous system associated with the second equation of DIRAC system:

$$\begin{aligned}
 -jk_t a_1 - jk_x k_2 + k_y j_2 + jk_z m_3 - \eta i_1 &= 0 \\
 -jk_t b_1 - jk_x l_2 - k_y i_2 + jk_z n_3 - \eta j_1 &= 0 \\
 -jk_t c_1 + jk_x i_2 + k_y l_2 + jk_z o_3 - \eta k_1 &= 0 \\
 -jk_t d_1 + jk_x j_2 - k_y k_2 + jk_z p_3 - \eta l_1 &= 0 \\
 -jk_t e_1 - jk_x o_2 + k_y n_2 - jk_z i_3 - \eta m_1 &= 0 \\
 -jk_t f_1 - jk_x p_2 - k_y m_2 - jk_z j_3 - \eta n_1 &= 0 \\
 -jk_t g_1 + jk_x m_2 + k_y p_2 - jk_z k_3 - \eta o_1 &= 0 \\
 -jk_t h_1 + jk_x n_2 - k_y o_2 - jk_z l_3 - \eta p_1 &= 0 \\
 jk_t i_1 - jk_x c_2 + k_y b_2 + jk_z e_3 - \eta a_1 &= 0 \\
 jk_t j_1 - jk_x d_2 - k_y a_2 + jk_z f_3 - \eta b_1 &= 0 \\
 jk_t k_1 + jk_x a_2 + k_y d_2 + jk_z g_3 - \eta c_1 &= 0 \\
 jk_t l_1 + jk_x b_2 - k_y c_2 + jk_z h_3 - \eta d_1 &= 0 \\
 jk_t m_1 - jk_x g_2 + k_y f_2 - jk_z a_3 - \eta e_1 &= 0 \\
 jk_t n_1 - jk_x h_2 - k_y e_2 - jk_z b_3 - \eta f_1 &= 0 \\
 jk_t o_1 + jk_x e_2 + k_y h_2 - jk_z c_3 - \eta g_1 &= 0 \\
 jk_t p_1 + jk_x f_2 - k_y g_2 - jk_z d_3 - \eta h_1 &= 0
 \end{aligned}
 \tag{VI-10}$$

Third equation of DIRAC system: $-j \frac{\partial \psi_2}{\partial(ct)} - j \frac{\partial \psi_1}{\partial x} - \frac{\partial \psi_1}{\partial y} - j \frac{\partial \psi_0}{\partial z} - \eta \psi_2 = 0$

$$\begin{aligned}
 & -j \frac{\partial \psi_2}{\partial(ct)} = jk_t \{a_2 S_x S_y S_z + b_2 S_x C_y S_z + c_2 C_x S_y S_z + d_2 C_x C_y S_z + e_2 S_x S_y C_z + f_2 S_x C_y C_z + g_2 C_x S_y C_z + h_2 C_x C_y C_z\} S_t \\
 & -jk_t \{a_2 S_x S_y S_z + b_2 S_x C_y S_z + c_2 C_x S_y S_z + d_2 C_x C_y S_z + e_2 S_x S_y C_z + f_2 S_x C_y C_z + g_2 C_x S_y C_z + h_2 C_x C_y C_z\} C_t \\
 & -j \frac{\partial \psi_1}{\partial x} = -jk_x \{a_1 C_x S_y S_z + b_1 C_x C_y S_z - c_1 S_x S_y S_z - d_1 S_x C_y S_z + e_1 C_x S_y C_z + f_1 C_x C_y C_z - g_1 S_x S_y C_z - h_1 S_x C_y C_z\} C_t \\
 & -jk_x \{i_1 C_x S_y S_z + j_1 C_x C_y S_z - k_1 S_x S_y S_z - l_1 S_x C_y S_z + m_1 C_x S_y C_z + n_1 C_x C_y C_z - o_1 S_x S_y C_z - p_1 S_x C_y C_z\} S_t \\
 & -\frac{\partial \psi_1}{\partial y} = -k_y \{a_1 S_x C_y S_z - b_1 S_x S_y S_z + c_1 C_x C_y S_z - d_1 C_x S_y S_z + e_1 S_x C_y C_z - f_1 S_x S_y C_z + g_1 C_x C_y C_z - h_1 C_x S_y C_z\} C_t \\
 & -k_y \{i_1 S_x C_y S_z - j_1 S_x S_y S_z + k_1 C_x C_y S_z - l_1 C_x S_y S_z + m_1 S_x C_y C_z - n_1 S_x S_y C_z + o_1 C_x C_y C_z - p_1 C_x S_y C_z\} S_t \\
 & -j \frac{\partial \psi_0}{\partial z} = -jk_z \{a_0 S_x S_y C_z + b_0 S_x C_y C_z + c_0 C_x S_y C_z + d_0 C_x C_y C_z - e_0 S_x S_y S_z - f_0 S_x C_y S_z - g_0 C_x S_y S_z - h_0 C_x C_y S_z\} C_t \\
 & -jk_z \{i_0 S_x S_y C_z + j_0 S_x C_y C_z + k_0 C_x S_y C_z + l_0 C_x C_y C_z - m_0 S_x S_y S_z - n_0 S_x C_y S_z - o_0 C_x S_y S_z - p_0 C_x C_y S_z\} S_t \\
 & \eta \psi_2 = \eta \{a_2 S_x S_y S_z + b_2 S_x C_y S_z + c_2 C_x S_y S_z + d_2 C_x C_y S_z + e_2 S_x S_y C_z + f_2 S_x C_y C_z + g_2 C_x S_y C_z + h_2 C_x C_y C_z\} C_t \\
 & + \eta \{i_2 S_x S_y S_z + j_2 S_x C_y S_z + k_2 C_x S_y S_z + l_2 C_x C_y S_z + m_2 S_x S_y C_z + n_2 S_x C_y C_z + o_2 C_x S_y C_z + p_2 C_x C_y C_z\} S_t
 \end{aligned} \tag{VI-11}$$

We can deduce the homogeneous system associated with the third equation of DIRAC system:

$$\begin{aligned}
 & jk_t a_2 + jk_x k_1 + k_y j_1 + jk_z m_0 - \eta i_2 = 0 \\
 & jk_t b_2 + jk_x l_1 - k_y i_1 + jk_z n_0 - \eta j_2 = 0 \\
 & jk_t c_2 - jk_x i_1 + k_y l_1 + jk_z o_0 - \eta k_2 = 0 \\
 & jk_t d_2 - jk_x j_1 - k_y k_1 + jk_z p_0 - \eta l_2 = 0 \\
 & jk_t e_2 + jk_x o_1 + k_y n_1 - jk_z i_0 - \eta m_2 = 0 \\
 & jk_t f_2 + jk_x p_1 - k_y m_1 - jk_z j_0 - \eta n_2 = 0 \\
 & jk_t g_2 - jk_x m_1 + k_y p_1 - jk_z k_0 - \eta o_2 = 0 \\
 & jk_t h_2 - jk_x n_1 - k_y o_1 - jk_z l_0 - \eta p_2 = 0 \\
 & -jk_t i_2 + jk_x c_1 + k_y b_1 + jk_z e_0 - \eta a_2 = 0 \\
 & -jk_t j_2 + jk_x d_1 - k_y a_1 + jk_z f_0 - \eta b_2 = 0 \\
 & -jk_t k_2 - jk_x a_1 + k_y d_1 + jk_z g_0 - \eta c_2 = 0 \\
 & -jk_t l_2 - jk_x b_1 - k_y c_1 + jk_z h_0 - \eta d_2 = 0 \\
 & -jk_t m_2 + jk_x g_1 + k_y f_1 - jk_z a_0 - \eta e_2 = 0 \\
 & -jk_t n_2 + jk_x h_1 - k_y e_1 - jk_z b_0 - \eta f_2 = 0 \\
 & -jk_t o_2 - jk_x e_1 + k_y h_1 - jk_z c_0 - \eta g_2 = 0 \\
 & -jk_t p_2 - jk_x f_1 - k_y g_1 - jk_z d_0 - \eta h_2 = 0
 \end{aligned} \tag{VI-12}$$

Fourth equation of DIRAC system: $-j \frac{\partial \Psi_3}{\partial(ct)} - j \frac{\partial \Psi_0}{\partial x} + \frac{\partial \Psi_0}{\partial y} + j \frac{\partial \Psi_1}{\partial z} - \eta \Psi_3 = 0$

$$\begin{aligned}
 & -j \frac{\partial \Psi_3}{\partial(ct)} = jk_t \{a_3 S_x S_y S_z + b_3 S_x C_y S_z + c_3 C_x S_y S_z + d_3 C_x C_y S_z + e_3 S_x S_y C_z + f_3 S_x C_y C_z + g_3 C_x S_y C_z + h_3 C_x C_y C_z\} S_t \\
 & - jk_t \{i_3 S_x S_y S_z + j_3 S_x C_y S_z + k_3 C_x S_y S_z + l_3 C_x C_y S_z + m_3 S_x S_y C_z + n_3 S_x C_y C_z + o_3 C_x S_y C_z + p_3 C_x C_y C_z\} C_t \\
 & - j \frac{\partial \Psi_0}{\partial x} = -jk_x \{a_0 C_x S_y S_z + b_0 C_x C_y S_z - c_0 S_x S_y S_z - d_0 S_x C_y S_z + e_0 C_x S_y C_z + f_0 C_x C_y C_z - g_0 S_x S_y C_z - h_0 S_x C_y C_z\} C_t \\
 & - jk_x \{i_0 C_x S_y S_z + j_0 C_x C_y S_z - k_0 S_x S_y S_z - l_0 S_x C_y S_z + m_0 C_x S_y C_z + n_0 C_x C_y C_z - o_0 S_x S_y C_z - p_0 S_x C_y C_z\} S_t \\
 & \frac{\partial \Psi_0}{\partial y} = k_y \{a_0 S_x C_y S_z - b_0 S_x S_y S_z + c_0 C_x C_y S_z - d_0 C_x S_y S_z + e_0 S_x C_y C_z - f_0 S_x S_y C_z + g_0 C_x C_y C_z - h_0 C_x S_y C_z\} C_t \\
 & + k_y \{i_0 S_x C_y S_z - j_0 S_x S_y S_z + k_0 C_x C_y S_z - l_0 C_x S_y S_z + m_0 S_x C_y C_z - n_0 S_x S_y C_z + o_0 C_x C_y C_z - p_0 C_x S_y C_z\} S_t \\
 & j \frac{\partial \Psi_1}{\partial z} = jk_z \{a_1 S_x S_y C_z + b_1 S_x C_y C_z + c_1 C_x S_y C_z + d_1 C_x C_y C_z - e_1 S_x S_y S_z - f_1 S_x C_y S_z - g_1 C_x S_y S_z - h_1 C_x C_y S_z\} C_t \\
 & + jk_z \{i_1 S_x S_y C_z + j_1 S_x C_y C_z + k_1 C_x S_y C_z + l_1 C_x C_y C_z - m_1 S_x S_y S_z - n_1 S_x C_y S_z - o_1 C_x S_y S_z - p_1 C_x C_y S_z\} S_t \\
 & \eta \Psi_3 = \eta \{a_3 S_x S_y S_z + b_3 S_x C_y S_z + c_3 C_x S_y S_z + d_3 C_x C_y S_z + e_3 S_x S_y C_z + f_3 S_x C_y C_z + g_3 C_x S_y C_z + h_3 C_x C_y C_z\} C_t \\
 & + \eta \{i_3 S_x S_y S_z + j_3 S_x C_y S_z + k_3 C_x S_y S_z + l_3 C_x C_y S_z + m_3 S_x S_y C_z + n_3 S_x C_y C_z + o_3 C_x S_y C_z + p_3 C_x C_y C_z\} S_t
 \end{aligned} \tag{VI-13}$$

We can deduce the homogeneous system associated with the fourth equation of DIRAC system:

$$\begin{aligned}
 & jk_t a_3 + jk_x k_0 - k_y j_0 - jk_z m_1 - \eta i_3 = 0 \\
 & jk_t b_3 + jk_x l_0 + k_y i_0 - jk_z n_1 - \eta j_3 = 0 \\
 & jk_t c_3 - jk_x i_0 - k_y l_0 - jk_z o_1 - \eta k_3 = 0 \\
 & jk_t d_3 - jk_x j_0 + k_y k_0 - jk_z p_1 - \eta l_3 = 0 \\
 & jk_t e_3 + jk_x o_0 - k_y n_0 + jk_z i_1 - \eta m_3 = 0 \\
 & jk_t f_3 + jk_x p_0 + k_y m_0 + jk_z j_1 - \eta n_3 = 0 \\
 & jk_t g_3 - jk_x m_0 - k_y p_0 + jk_z k_1 - \eta o_3 = 0 \\
 & jk_t h_3 - jk_x n_0 + k_y o_0 + jk_z l_1 - \eta p_3 = 0 \\
 & - jk_t i_3 + jk_x c_0 - k_y b_0 - jk_z e_1 - \eta a_3 = 0 \\
 & - jk_t j_3 + jk_x d_0 + k_y a_0 - jk_z f_1 - \eta b_3 = 0 \\
 & - jk_t k_3 - jk_x a_0 - k_y d_0 - jk_z g_1 - \eta c_3 = 0 \\
 & - jk_t l_3 - jk_x b_0 + k_y c_0 - jk_z h_1 - \eta d_3 = 0 \\
 & - jk_t m_3 + jk_x g_0 - k_y f_0 + jk_z a_1 - \eta e_3 = 0 \\
 & - jk_t n_3 + jk_x h_0 + k_y e_0 + jk_z b_1 - \eta f_3 = 0 \\
 & - jk_t o_3 - jk_x e_0 - k_y h_0 + jk_z c_1 - \eta g_3 = 0 \\
 & - jk_t p_3 - jk_x f_0 + k_y g_0 + jk_z d_1 - \eta h_3 = 0
 \end{aligned} \tag{VI-14}$$

The complete system of 64 equations relating to the coefficients $a_i, b_i, c_i, d_i, e_i, f_i, g_i, h_i, i_i, j_i, k_i, l_i, m_i, n_i, o_i, p_i$ for $i = 0, 1, 2, 3$ can now be summarized on 2 columns:

$$\begin{aligned}
 & -jk_t a_0 - jk_x k_3 - k_y j_3 - jk_z m_2 - \eta i_0 = 0 & -jk_t a_1 - jk_x k_2 + k_y j_2 + jk_z m_3 - \eta i_1 = 0 \\
 & -jk_t b_0 - jk_x l_3 + k_y i_3 - jk_z n_2 - \eta j_0 = 0 & -jk_t b_1 - jk_x l_2 - k_y i_2 + jk_z n_3 - \eta j_1 = 0 \\
 & -jk_t c_0 + jk_x i_3 - k_y l_3 - jk_z o_2 - \eta k_0 = 0 & -jk_t c_1 + jk_x i_2 + k_y l_2 + jk_z o_3 - \eta k_1 = 0 \\
 & -jk_t d_0 + jk_x j_3 + k_y k_3 - jk_z p_2 - \eta l_0 = 0 & -jk_t d_1 + jk_x j_2 - k_y k_2 + jk_z p_3 - \eta l_1 = 0 \\
 & -jk_t e_0 - jk_x o_3 - k_y n_3 + jk_z i_2 - \eta m_0 = 0 & -jk_t e_1 - jk_x o_2 + k_y n_2 - jk_z i_3 - \eta m_1 = 0 \\
 & -jk_t f_0 - jk_x p_3 + k_y m_3 + jk_z j_2 - \eta n_0 = 0 & -jk_t f_1 - jk_x p_2 - k_y m_2 - jk_z j_3 - \eta n_1 = 0 \\
 & -jk_t g_0 + jk_x m_3 - k_y p_3 + jk_z k_2 - \eta o_0 = 0 & -jk_t g_1 + jk_x m_2 + k_y p_2 - jk_z k_3 - \eta o_1 = 0 \\
 & -jk_t h_0 + jk_x n_3 + k_y o_3 + jk_z l_2 - \eta p_0 = 0 & -jk_t h_1 + jk_x n_2 - k_y o_2 - jk_z l_3 - \eta p_1 = 0 \\
 & jk_t i_0 - jk_x c_3 - k_y b_3 - jk_z e_2 - \eta a_0 = 0 & jk_t i_1 - jk_x c_2 + k_y b_2 + jk_z e_3 - \eta a_1 = 0 \\
 & jk_t j_0 - jk_x d_3 + k_y a_3 - jk_z f_2 - \eta b_0 = 0 & jk_t j_1 - jk_x d_2 - k_y a_2 + jk_z f_3 - \eta b_1 = 0 \\
 & jk_t k_0 + jk_x a_3 - k_y d_3 - jk_z g_2 - \eta c_0 = 0 & jk_t k_1 + jk_x a_2 + k_y d_2 + jk_z g_3 - \eta c_1 = 0 \\
 & jk_t l_0 + jk_x b_3 + k_y c_3 - jk_z h_2 - \eta d_0 = 0 & jk_t l_1 + jk_x b_2 - k_y c_2 + jk_z h_3 - \eta d_1 = 0 \\
 & jk_t m_0 - jk_x g_3 - k_y f_3 + jk_z a_2 - \eta e_0 = 0 & jk_t m_1 - jk_x g_2 + k_y f_2 - jk_z a_3 - \eta e_1 = 0 \\
 & jk_t n_0 - jk_x h_3 + k_y e_3 + jk_z b_2 - \eta f_0 = 0 & jk_t n_1 - jk_x h_2 - k_y e_2 - jk_z b_3 - \eta f_1 = 0 \\
 & jk_t o_0 + jk_x e_3 - k_y h_3 + jk_z c_2 - \eta g_0 = 0 & jk_t o_1 + jk_x e_2 + k_y h_2 - jk_z c_3 - \eta g_1 = 0 \\
 & jk_t p_0 + jk_x f_3 + k_y g_3 + jk_z d_2 - \eta h_0 = 0 & jk_t p_1 + jk_x f_2 - k_y g_2 - jk_z d_3 - \eta h_1 = 0 \\
 \\
 & jk_t a_2 + jk_x k_1 + k_y j_1 + jk_z m_0 - \eta i_2 = 0 & jk_t a_3 + jk_x k_0 - k_y j_0 - jk_z m_1 - \eta i_3 = 0 \\
 & jk_t b_2 + jk_x l_1 - k_y i_1 + jk_z n_0 - \eta j_2 = 0 & jk_t b_3 + jk_x l_0 + k_y i_0 - jk_z n_1 - \eta j_3 = 0 \\
 & jk_t c_2 - jk_x i_1 + k_y l_1 + jk_z o_0 - \eta k_2 = 0 & jk_t c_3 - jk_x i_0 - k_y l_0 - jk_z o_1 - \eta k_3 = 0 \\
 & jk_t d_2 - jk_x j_1 - k_y k_1 + jk_z p_0 - \eta l_2 = 0 & jk_t d_3 - jk_x j_0 + k_y k_0 - jk_z p_1 - \eta l_3 = 0 \\
 & jk_t e_2 + jk_x o_1 + k_y n_1 - jk_z i_0 - \eta m_2 = 0 & jk_t e_3 + jk_x o_0 - k_y n_0 + jk_z i_1 - \eta m_3 = 0 \\
 & jk_t f_2 + jk_x p_1 - k_y m_1 - jk_z j_0 - \eta n_2 = 0 & jk_t f_3 + jk_x p_0 + k_y m_0 + jk_z j_1 - \eta n_3 = 0 \\
 & jk_t g_2 - jk_x m_1 + k_y p_1 - jk_z k_0 - \eta o_2 = 0 & jk_t g_3 - jk_x m_0 - k_y p_0 + jk_z k_1 - \eta o_3 = 0 \\
 & jk_t h_2 - jk_x n_1 - k_y o_1 - jk_z l_0 - \eta p_2 = 0 & jk_t h_3 - jk_x n_0 + k_y o_0 + jk_z l_1 - \eta p_3 = 0 \\
 & -jk_t i_2 + jk_x c_1 + k_y b_1 + jk_z e_0 - \eta a_2 = 0 & -jk_t i_3 + jk_x c_0 - k_y b_0 - jk_z e_1 - \eta a_3 = 0 \\
 & -jk_t j_2 + jk_x d_1 - k_y a_1 + jk_z f_0 - \eta b_2 = 0 & -jk_t j_3 + jk_x d_0 + k_y a_0 - jk_z f_1 - \eta b_3 = 0 \\
 & -jk_t k_2 - jk_x a_1 + k_y d_1 + jk_z g_0 - \eta c_2 = 0 & -jk_t k_3 - jk_x a_0 - k_y d_0 - jk_z g_1 - \eta c_3 = 0 \\
 & -jk_t l_2 - jk_x b_1 - k_y c_1 + jk_z h_0 - \eta d_2 = 0 & -jk_t l_3 - jk_x b_0 + k_y c_0 - jk_z h_1 - \eta d_3 = 0 \\
 & -jk_t m_2 + jk_x g_1 + k_y f_1 - jk_z a_0 - \eta e_2 = 0 & -jk_t m_3 + jk_x g_0 - k_y f_0 + jk_z a_1 - \eta e_3 = 0 \\
 & -jk_t n_2 + jk_x h_1 - k_y e_1 - jk_z b_0 - \eta f_2 = 0 & -jk_t n_3 + jk_x h_0 + k_y e_0 + jk_z b_1 - \eta f_3 = 0 \\
 & -jk_t o_2 - jk_x e_1 + k_y h_1 - jk_z c_0 - \eta g_2 = 0 & -jk_t o_3 - jk_x e_0 - k_y h_0 + jk_z c_1 - \eta g_3 = 0 \\
 & -jk_t p_2 - jk_x f_1 - k_y g_1 - jk_z d_0 - \eta h_2 = 0 & -jk_t p_3 - jk_x f_0 + k_y g_0 + jk_z d_1 - \eta h_3 = 0
 \end{aligned}$$

(VI-15)

It is a homogeneous system that allows non-zero solution only if its determinant is zero. But the literal expression of the determinant of a system of 64 equations with 64 unknowns is not trivial to obtain.

However, one can try to identify it by some physical considerations. If there is a relationship between k_x , k_y , k_z , k_t and η which allows to obtain solutions to this system, this relationship must express the conservation of energy.

It was recalled in the introduction that the pseudo-norm of the pulse energy four-vector:

$$\tilde{\mathbf{P}} = \begin{pmatrix} p_x \\ p_y \\ p_z \\ \frac{E}{c} \end{pmatrix} \quad (\text{VI-16})$$

does not depend on the frame in which it is expressed and it has been shown that its consistency is used to establish the relationship of conservation of energy:

$$E^2 = (pc)^2 + (m_0c^2)^2 = (p_xc)^2 + (p_yc)^2 + (p_zc)^2 + (m_0c^2)^2 \quad (\text{VI-17})$$

The counterpart in (VI-16) four-vector in quantum mechanics is obtained by multiplying the four-vector of wave by the barred PLANCK's constant:

$$\tilde{\mathbf{P}}_Q = \hbar \begin{pmatrix} k_x \\ k_y \\ k_z \\ k_t = \frac{\omega}{c} \end{pmatrix} \quad (\text{VI-18})$$

For the same reasons as before, the pseudo-norm of this four-vector is constant and this constant is necessarily the rest mass energy divided by c^2 . We can deduce:

$$\hbar^2(k_t^2 - k_x^2 - k_y^2 - k_z^2) = (m_0c)^2 \quad (\text{VI-19})$$

or again, by making use of the notation used in the expression of the DIRAC system recalled in (V-2):

$$\eta = \frac{m_0c}{\hbar} \quad (\text{VI-20})$$

$$k_t^2 = k_x^2 + k_y^2 + k_z^2 + \eta^2 \quad (\text{VI-21})$$

In summary, if DIRAC system has solutions, these must necessarily be consistent with the equation of conservation of energy (VI-21).

It is now possible to show that a solution in the form of a linear combination of stationary modes is solution of the KLEIN-GORDON equation recalled below:

$$\left[\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} - \frac{1}{c^2} \frac{\partial^2}{\partial t^2} \right] (\psi) = \frac{m^2 c^2}{\hbar^2} (\psi) \quad (\text{VI-22})$$

Consider a wave function ψ representing one any cavity three-dimensional patterns expressed in (VI-3), for example to fix ideas:

$$\psi(x, y, z, t) = A \sin(k_x x) \sin(k_y y) \cos(k_z z) \cos(k_t ct) \quad (\text{VI-23})$$

By substituting the function ψ from (VI-23) in (VI-22), we get:

$$\left(-k_x^2 - k_y^2 - k_z^2 + k_t^2 \right) (\psi) = \eta^2 (\psi) \quad (\text{VI-24})$$

which suggests that DIRAC system admits solutions in the form of stationary modes provided that the equation of conservation of energy (VI-21) is satisfied.

Based on these assumptions, it can be shown that the determinant of the entire system of 64 equations in 64 unknowns has the determinant:

$$\left(-k_t^2 + k_x^2 + k_y^2 + k_z^2 + \eta^2 \right)^{32} \quad (\text{VI-35})$$

It is concluded definitively that when this determinant is zero, that is, when the energy conservation equation is verified, there are solutions to the DIRAC system in the form of standing waves.

VII

Exact solutions to the DIRAC system

The previous chapter has allowed to show that there were stationary solutions to this system. To be convincing, we must be able to clarify them.

A detailed analysis of the system shows that when the determinant is zero, that is, when the following condition occurs:

$$k_t^2 = k_x^2 + k_y^2 + k_z^2 + \eta^2 \quad (\text{VII-1})$$

the choice of one of spinors is arbitrary, and the other follows.

Taking into account this observation, two tables of solutions have been built for spinors defined in previous chapters:

$$\varphi = \begin{pmatrix} \psi_0 \\ \psi_1 \end{pmatrix} \quad \text{and} \quad \chi = \begin{pmatrix} \psi_2 \\ \psi_3 \end{pmatrix} \quad (\text{VII-2})$$

I - Solutions with a single mode excited on one of the components of spinors

The first table (table VII-1) was developed by choosing the spinor φ and calculating the corresponding spinor χ to be solution for the DIRAC system. In order to sweep the set of solutions, the choice of the spinor φ has been made by setting the component ψ_0 successively with all the possible modes while maintaining $\psi_1 = 0$, then by setting the component ψ_1 successively all possible modes while keeping now $\psi_0 = 0$.

	a	b	c	d	e	f	g	h	i	j	k	l	m	n	o	p
	SxSySzCt	SxCySzCt	CxSySzCt	CxCySzCt	SxSyCzCt	SxCyCzCt	CxSyCzCt	CxCyCzCt	SxSySzSt	SxCySzSt	CxSySzSt	CxCySzSt	SxSyCzSt	SxCyCzSt	CxSyCzSt	CxCyCzSt
ψ_0	a0	b0	c0	d0	e0	f0	g0	h0	i0	j0	k0	l0	m0	n0	o0	p0
ψ_1	a1	b1	c1	d1	e1	f1	g1	h1	i1	j1	k1	l1	m1	n1	o1	p1
$\psi_2(\eta^2-k^2)$	j.e0.η.Kz m0.Kt.Kz b1.η.Ky j.c1.η.Kx -j.j1.Kt.Ky k1.Kt.Kx	j.f0.η.Kz n0.Kt.Kz -a1.η.Ky j.d1.η.Kx j.i1.Kt.Ky l1.Kt.Kx	j.g0.η.Kz o0.Kt.Kz -j.a1.η.Kx d1.η.Ky -i1.Kt.Kx -j.l1.Kt.Ky	j.h0.η.Kz p0.Kt.Kz -j.b1.η.Kx -c1.η.Ky -j1.Kt.Kx j.k1.Kt.Ky	-j.a0.η.Kz -i0.Kt.Kz f1.η.Ky j.g1.η.Kx -j.n1.Kt.Ky o1.Kt.Kx	-j.b0.η.Kz -j0.Kt.Kz -e1.η.Ky j.h1.η.Kx j.m1.Kt.Ky p1.Kt.Kx	-j.c0.η.Kz -k0.Kt.Kz -j.e1.η.Kx h1.η.Ky -m1.Kt.Kx -j.p1.Kt.Ky	-j.d0.η.Kz -l0.Kt.Kz -j.f1.η.Kx -g1.η.Ky -n1.Kt.Kx j.o1.Kt.Ky	-e0.Kt.Kz j.m0.η.Kz j.b1.Kt.Ky -c1.Kt.Kx j1.η.Ky j.k1.η.Kx	-f0.Kt.Kz j.n0.η.Kz -j.a1.Kt.Ky -d1.Kt.Kx -i1.η.Ky j.l1.η.Kx	-g0.Kt.Kz j.o0.η.Kz a1.Kt.Kx j.d1.Kt.Ky -j.i1.η.Kx l1.η.Ky	-h0.Kt.Kz j.p0.η.Kz b1.Kt.Kx -j.c1.Kt.Ky -j.j1.η.Kx -k1.η.Ky	a0.Kt.Kz -j.i0.η.Kz j.f1.Kt.Ky -g1.Kt.Kx n1.η.Ky j.o1.η.Kx	b0.Kt.Kz -j.j0.η.Kz -j.e1.Kt.Ky -h1.Kt.Kx -m1.η.Ky j.p1.η.Kx	c0.Kt.Kz -j.k0.η.Kz e1.Kt.Kx j.h1.Kt.Ky -j.m1.η.Kx p1.η.Ky	d0.Kt.Kz -j.l0.η.Kz f1.Kt.Kx -j.g1.Kt.Ky -j.n1.η.Kx -o1.η.Ky
$\psi_3(\eta^2-k^2)$	-b0.η.Ky j.c0.η.Kx j.j0.Kt.Ky k0.Kt.Kx -j.e1.η.Kz -m1.Kt.Kz	a0.η.Ky j.d0.η.Kx -j.i0.Kt.Ky l0.Kt.Kx -j.f1.η.Kz -n1.Kt.Kz	-j.a0.η.Kx -d0.η.Ky -i0.Kt.Kx j.l0.Kt.Ky -j.g1.η.Kz -o1.Kt.Kz	-j.b0.η.Kx c0.η.Ky -j0.Kt.Kx -j.k0.Kt.Ky -j.h1.η.Kz -p1.Kt.Kz	-f0.η.Ky j.g0.η.Kx j.n0.Kt.Ky o0.Kt.Kx j.a1.η.Kz i1.Kt.Kz	e0.η.Ky j.h0.η.Kx -j.m0.Kt.Ky p0.Kt.Kx j.b1.η.Kz j1.Kt.Kz	-j.e0.η.Kx -h0.η.Ky -m0.Kt.Kx j.p0.Kt.Ky j.c1.η.Kz k1.Kt.Kz	-j.f0.η.Kx g0.η.Ky -n0.Kt.Kx -j.o0.Kt.Ky j.d1.η.Kz l1.Kt.Kz	-j.b0.Kt.Ky -c0.Kt.Kx -j0.η.Ky j.k0.η.Kx e1.Kt.Kz -j.m1.η.Kz	j.a0.Kt.Ky -d0.Kt.Kx i0.η.Ky j.l0.η.Kx f1.Kt.Kz -j.n1.η.Kz	a0.Kt.Kx -j.d0.Kt.Ky -j.i0.η.Kx -l0.η.Ky g1.Kt.Kz -j.o1.η.Kz	b0.Kt.Kx j.c0.Kt.Ky -j.j0.η.Kx k0.η.Ky h1.Kt.Kz -j.p1.η.Kz	j.f0.Kt.Ky -g0.Kt.Kx -n0.η.Ky j.o0.η.Kx -a1.Kt.Kz j.i1.η.Kz	j.e0.Kt.Ky -h0.Kt.Kx m0.η.Ky j.p0.η.Kx -b1.Kt.Kz j.j1.η.Kz	e0.Kt.Kx -j.h0.Kt.Ky -j.m0.η.Kx -p0.η.Ky -c1.Kt.Kz j.k1.η.Kz	f0.Kt.Kx j.g0.Kt.Ky -j.n0.η.Kx o0.η.Ky -d1.Kt.Kz j.l1.η.Kz

Table VII-1 : Stationary solutions of DIRAC equation for ψ_2 and ψ_3 versus the excited modes ψ_0 and ψ_1 .

The reading of this table is made in the following way: ψ_0 and ψ_1 being chosen as a amplitude for a stationary mode x_0 or x_1 (x represents any letter included between a and p), ψ_2 and ψ_3 are determined by identifying in each columns the modes related to amplitudes x_0 or x_1 .

The wave functions expressed in these solutions are dimensionless. From a purely mathematical point of view, these wave functions may have a multiplicative constant, which allows to express them in different units. We will use this property later.

Examples are proposed in order to familiarize themselves with the reading of the table (VII-1). We put for homogeneity of notation, $x_t = ct$ which allows to write the term of temporal phase under the form $\omega t = k_t x_t$.

Example 1:

$$\psi_0 = a_0 \sin(k_x x) \sin(k_y y) \sin(k_z z) \cos(k_t x_t) \quad (\text{VII-3})$$

$$\psi_1 = 0$$

$$\psi_2 = a_0 \frac{k_t k_z}{\eta^2 - k_t^2} \sin(k_x x) \sin(k_y y) \cos(k_z z) \sin(k_t x_t) - ja_0 \frac{\eta k_z}{\eta^2 - k_t^2} \sin(k_x x) \sin(k_y y) \cos(k_z z) \cos(k_t x_t)$$

$$\psi_3 = a_0 \frac{\eta k_y}{\eta^2 - k_t^2} \sin(k_x x) \cos(k_y y) \sin(k_z z) \cos(k_t x_t) + ja_0 \frac{k_t k_y}{\eta^2 - k_t^2} \sin(k_x x) \cos(k_y y) \sin(k_z z) \sin(k_t x_t)$$

$$+ a_0 \frac{k_t k_x}{\eta^2 - k_t^2} \cos(k_x x) \sin(k_y y) \sin(k_z z) \sin(k_t x_t) - ja_0 \frac{\eta k_x}{\eta^2 - k_t^2} \cos(k_x x) \sin(k_y y) \sin(k_z z) \cos(k_t x_t)$$

Example 2:

$$\psi_0 = 0 \quad (\text{VII-4})$$

$$\psi_1 = a_1 \sin(k_x x) \sin(k_y y) \sin(k_z z) \cos(k_t x_t)$$

$$\psi_2 = -a_1 \frac{\eta k_y}{\eta^2 - k_t^2} \sin(k_x x) \cos(k_y y) \sin(k_z z) \cos(k_t x_t) - ja_1 \frac{k_t k_y}{\eta^2 - k_t^2} \sin(k_x x) \cos(k_y y) \sin(k_z z) \sin(k_t x_t)$$

$$+ a_1 \frac{k_t k_x}{\eta^2 - k_t^2} \cos(k_x x) \sin(k_y y) \sin(k_z z) \sin(k_t x_t) - ja_1 \frac{\eta k_x}{\eta^2 - k_t^2} \cos(k_x x) \sin(k_y y) \sin(k_z z) \cos(k_t x_t)$$

$$\psi_3 = -a_1 \frac{k_t k_z}{\eta^2 - k_t^2} \sin(k_x x) \sin(k_y y) \cos(k_z z) \sin(k_t x_t) + ja_1 \frac{\eta k_z}{\eta^2 - k_t^2} \sin(k_x x) \sin(k_y y) \cos(k_z z) \cos(k_t x_t)$$

Example 3:

$$\psi_0 = b_0 \sin(k_x x) \cos(k_y y) \sin(k_z z) \cos(k_t x_t) \quad (\text{VII-5})$$

$$\psi_1 = 0$$

$$\psi_2 = b_0 \frac{k_t k_z}{\eta^2 - k_t^2} \sin(k_x x) \cos(k_y y) \cos(k_z z) \sin(k_t x_t) - jb_0 \frac{\eta k_z}{\eta^2 - k_t^2} \sin(k_x x) \cos(k_y y) \cos(k_z z) \cos(k_t x_t)$$

$$\psi_3 = -b_0 \frac{\eta k_y}{\eta^2 - k_t^2} \sin(k_x x) \sin(k_y y) \sin(k_z z) \cos(k_t x_t) - jb_0 \frac{k_t k_y}{\eta^2 - k_t^2} \sin(k_x x) \sin(k_y y) \sin(k_z z) \sin(k_t x_t)$$

$$+ b_0 \frac{k_t k_x}{\eta^2 - k_t^2} \cos(k_x x) \cos(k_y y) \sin(k_z z) \sin(k_t x_t) - jb_0 \frac{\eta k_x}{\eta^2 - k_t^2} \cos(k_x x) \cos(k_y y) \sin(k_z z) \cos(k_t x_t)$$

Example 4:

$$\psi_0 = 0 \quad (VII-6)$$
$$\psi_1 = b_1 \sin(k_x x) \cos(k_y y) \sin(k_z z) \cos(k_t x_t)$$

$$\psi_2 = b_1 \frac{\eta k_y}{\eta^2 - k_t^2} \sin(k_x x) \sin(k_y y) \sin(k_z z) \cos(k_t x_t) + j b_1 \frac{k_t k_y}{\eta^2 - k_t^2} \sin(k_x x) \sin(k_y y) \sin(k_z z) \sin(k_t x_t)$$
$$+ b_1 \frac{k_t k_x}{\eta^2 - k_t^2} \cos(k_x x) \cos(k_y y) \sin(k_z z) \sin(k_t x_t) - j b_1 \frac{\eta k_x}{\eta^2 - k_t^2} \cos(k_x x) \cos(k_y y) \sin(k_z z) \cos(k_t x_t)$$

$$\psi_3 = -b_1 \frac{k_t k_z}{\eta^2 - k_t^2} \sin(k_x x) \cos(k_y y) \cos(k_z z) \sin(k_t x_t) + j b_1 \frac{\eta k_z}{\eta^2 - k_t^2} \sin(k_x x) \cos(k_y y) \cos(k_z z) \cos(k_t x_t)$$

The second table (table VII-2) has been developed by choosing the spinor χ and calculating the spinor ϕ corresponding to be solution of the DIRAC system. As in the previous table, the choice of the spinor χ was made by setting the component ψ_2 successively with all the possible modes while maintaining $\psi_3 = 0$, then by setting the component ψ_3 successively in all possible modes while keeping now $\psi_2 = 0$.

	a	b	c	d	e	f	g	h	i	j	k	l	m	n	o	p
	SxSySzSt	SxSySzSt	CxSySzSt	CxSySzSt	SxSyCzCt	SxSyCzCt	CxSyCzCt	CxSyCzCt	SxSySzSt	SxSySzSt	CxSySzSt	CxSySzSt	SxSyCzSt	SxSyCzSt	CxSyCzSt	CxSyCzSt
	-j.e2.η.Kz m2.Kt.Kz -b3.η.Ky -j.c3.η.Kx -j.j3.Kt.Ky k3.Kt.Kx	-j.f2.η.Kz n2.Kt.Kz a3.η.Ky -j.d3.η.Kx j.i3.Kt.Ky L3.Kt.Kx	-j.g2.η.Kz o2.Kt.Kz j.a3.η.Kx -d3.η.Ky -i3.Kt.Kx -j.L3.Kt.Ky	-j.h2.η.Kz p2.Kt.Kz j.b3.η.Kx c3.η.Ky -j3.Kt.Kx j.k3.Kt.Ky	j.a2.η.Kz -i2.Kt.Kz -f3.η.Ky -j.g3.η.Kx -j.n3.Kt.Ky o3.Kt.Kx	j.b2.η.Kz -j2.Kt.Kz e3.η.Ky -j.h3.η.Kx j.m3.Kt.Ky p3.Kt.Kx	j.c2.η.Kz -k2.Kt.Kz j.e3.η.Kx -h3.η.Ky -m3.Kt.Kx -j.p3.Kt.Ky	j.d2.η.Kz -L2.Kt.Kz j.f3.η.Kx g3.η.Ky -n3.Kt.Kx j.o3.Kt.Ky	-e2.Kt.Kz -j.m2.η.Kz j.b3.Kt.Ky -c3.Kt.Kx -j3.η.Ky -j.k3.η.Kx	-f2.Kt.Kz -j.n2.η.Kz -j.a3.Kt.Ky -d3.Kt.Kx i3.η.Ky -j.L3.η.Kx	-g2.Kt.Kz -j.o2.η.Kz a3.Kt.Kx j.d3.Kt.Ky j.i3.η.Kx -L3.η.Ky	-h2.Kt.Kz -j.p2.η.Kz b3.Kt.Kx -j.c3.Kt.Ky j.j3.η.Kx k3.η.Ky	a2.Kt.Kz j.i2.η.Kz j.f3.Kt.Ky -g3.Kt.Kx -n3.η.Ky -j.o3.η.Kx	b2.Kt.Kz j.j2.η.Kz -j.e3.Kt.Ky m3.η.Ky -j.p3.η.Kx	c2.Kt.Kz j.k2.η.Kz e3.Kt.Kx j.h3.Kt.Ky j.m3.η.Kx -p3.η.Ky	d2.Kt.Kz j.l2.η.Kz f3.Kt.Kx -j.g3.Kt.Ky j.n3.η.Kx o3.η.Ky
ψ_0																
	b2.η.Ky -j.c2.η.Kx j.j2.Kt.Ky k2.Kt.Kx j.e3.η.Kz -m3.Kt.Kz	-a2.η.Ky -j.d2.η.Kx -j.i2.Kt.Ky L2.Kt.Kx j.f3.η.Kz -n3.Kt.Kz	j.a2.η.Kx d2.η.Ky -i2.Kt.Kx j.L2.Kt.Ky j.g3.η.Kz -o3.Kt.Kz	j.b2.η.Kx -c2.η.Ky -j2.Kt.Kx -j.k2.Kt.Ky j.h3.η.Kz -p3.Kt.Kz	f2.η.Ky -j.g2.η.Kx j.n2.Kt.Ky o2.Kt.Kx -j.a3.η.Kz i3.Kt.Kz	-e2.η.Ky -j.h2.η.Kx -j.m2.Kt.Ky p2.Kt.Kx -j.b3.η.Kz j3.Kt.Kz	h2.η.Ky -m2.Kt.Kx j.p2.Kt.Ky -j.c3.η.Kz k3.Kt.Kz	j.f2.η.Kx -g2.η.Ky -n2.Kt.Kx -j.o2.Kt.Ky -j.d3.η.Kz L3.Kt.Kz	-j.b2.Kt.Ky -c2.Kt.Kx j2.η.Ky -j.k2.η.Kx e3.Kt.Kz j.m3.η.Kz	j.a2.Kt.Ky -d2.Kt.Kx -i2.η.Ky -j.L2.η.Kx f3.Kt.Kz j.n3.η.Kz	a2.Kt.Kx -j.d2.Kt.Ky j.i2.η.Kx L2.η.Ky g3.Kt.Kz j.o3.η.Kz	b2.Kt.Kx j.c2.Kt.Ky j.j2.η.Kx -k2.η.Ky h3.Kt.Kz j.p3.η.Kz	-j.f2.Kt.Ky -g2.Kt.Kx n2.η.Ky -j.o2.η.Kx -a3.Kt.Kz -j.i3.η.Kz	j.e2.Kt.Ky -h2.Kt.Kx -m2.η.Ky -j.p2.η.Kx -b3.Kt.Kz -j.j3.η.Kz	e2.Kt.Kx -j.h2.Kt.Ky j.m2.η.Kx p2.η.Ky -c3.Kt.Kz -j.k3.η.Kz	f2.Kt.Kx j.g2.Kt.Ky j.n2.η.Kx -o2.η.Ky -d3.Kt.Kz -j.l3.η.Kz
$\psi_2(\eta^2-k^2)$	a2	b2	c2	d2	e2	f2	g2	h2	i2	j2	k2	L2	m2	n2	o2	p2
$\psi_3(\eta^2-k^2)$	a3	b3	c3	d3	e3	f3	g3	h3	i3	j3	k3	L3	m3	n3	o3	p3

Table VII-2 : Stationary solutions of DIRAC equation for ψ_0 and ψ_1 versus the excited modes ψ_2 and ψ_3 .

As previously, some examples are proposed in order to familiarize themselves with the reading of the table (VII-2). It is recalled that we adopted the notation $x_t = ct$ which allows to write the term of temporal phase $\omega t = k_t x_t$.

Example 5:

$$\begin{aligned}\psi_0 &= a_2 \frac{k_t k_z}{\eta^2 - k_t^2} \sin(k_x x) \sin(k_y y) \cos(k_z z) \sin(k_t x_t) + ja_2 \frac{\eta k_z}{\eta^2 - k_t^2} \sin(k_x x) \sin(k_y y) \cos(k_z z) \cos(k_t x_t) \\ \psi_1 &= -a_2 \frac{\eta k_y}{\eta^2 - k_t^2} \sin(k_x x) \cos(k_y y) \sin(k_z z) \cos(k_t x_t) + ja_2 \frac{k_t k_y}{\eta^2 - k_t^2} \sin(k_x x) \cos(k_y y) \sin(k_z z) \sin(k_t x_t) \\ &+ a_2 \frac{k_t k_x}{\eta^2 - k_t^2} \cos(k_x x) \sin(k_y y) \sin(k_z z) \sin(k_t x_t) + ja_2 \frac{\eta k_x}{\eta^2 - k_t^2} \cos(k_x x) \sin(k_y y) \sin(k_z z) \cos(k_t x_t) \\ \psi_2 &= a_2 \sin(k_x x) \sin(k_y y) \sin(k_z z) \cos(k_t x_t) \\ \psi_3 &= 0\end{aligned}\tag{VII-7}$$

Example 6:

$$\begin{aligned}\psi_0 &= a_3 \frac{\eta k_y}{\eta^2 - k_t^2} \sin(k_x x) \cos(k_y y) \sin(k_z z) \cos(k_t x_t) - ja_3 \frac{k_t k_y}{\eta^2 - k_t^2} \sin(k_x x) \cos(k_y y) \sin(k_z z) \sin(k_t x_t) \\ &+ a_3 \frac{k_t k_x}{\eta^2 - k_t^2} \cos(k_x x) \sin(k_y y) \sin(k_z z) \sin(k_t x_t) + ja_3 \frac{\eta k_x}{\eta^2 - k_t^2} \cos(k_x x) \sin(k_y y) \sin(k_z z) \cos(k_t x_t) \\ \psi_1 &= -a_3 \frac{k_t k_z}{\eta^2 - k_t^2} \sin(k_x x) \sin(k_y y) \cos(k_z z) \sin(k_t x_t) - ja_3 \frac{\eta k_z}{\eta^2 - k_t^2} \sin(k_x x) \sin(k_y y) \cos(k_z z) \cos(k_t x_t) \\ \psi_2 &= 0 \\ \psi_3 &= a_3 \sin(k_x x) \sin(k_y y) \sin(k_z z) \cos(k_t x_t)\end{aligned}\tag{VII-8}$$

Example 7:

$$\begin{aligned}\psi_0 &= -h_2 \frac{k_t k_z}{\eta^2 - k_t^2} \cos(k_x x) \cos(k_y y) \sin(k_z z) \sin(k_t x_t) - jh_2 \frac{\eta k_z}{\eta^2 - k_t^2} \cos(k_x x) \cos(k_y y) \sin(k_z z) \cos(k_t x_t) \\ \psi_1 &= h_2 \frac{\eta k_y}{\eta^2 - k_t^2} \cos(k_x x) \sin(k_y y) \cos(k_z z) \cos(k_t x_t) - jh_2 \frac{k_t k_y}{\eta^2 - k_t^2} \cos(k_x x) \sin(k_y y) \cos(k_z z) \sin(k_t x_t) \\ &- h_2 \frac{k_t k_x}{\eta^2 - k_t^2} \sin(k_x x) \cos(k_y y) \cos(k_z z) \sin(k_t x_t) - jh_2 \frac{\eta k_x}{\eta^2 - k_t^2} \sin(k_x x) \cos(k_y y) \cos(k_z z) \cos(k_t x_t) \\ \psi_2 &= h_2 \cos(k_x x) \cos(k_y y) \cos(k_z z) \cos(k_t x_t) \\ \psi_3 &= 0\end{aligned}\tag{VII-9}$$

Example 8:

$$\begin{aligned}
 \psi_0 &= -h_3 \frac{\eta k_y}{\eta^2 - k_t^2} \cos(k_x x) \sin(k_y y) \cos(k_z z) \cos(k_t x_t) + j h_3 \frac{k_t k_y}{\eta^2 - k_t^2} \cos(k_x x) \sin(k_y y) \cos(k_z z) \sin(k_t x_t) \\
 &- h_3 \frac{k_t k_x}{\eta^2 - k_t^2} \sin(k_x x) \cos(k_y y) \cos(k_z z) \sin(k_t x_t) - j h_3 \frac{\eta k_x}{\eta^2 - k_t^2} \sin(k_x x) \cos(k_y y) \cos(k_z z) \cos(k_t x_t) \\
 \psi_1 &= h_3 \frac{k_t k_z}{\eta^2 - k_t^2} \cos(k_x x) \cos(k_y y) \sin(k_z z) \sin(k_t x_t) + j h_3 \frac{\eta k_z}{\eta^2 - k_t^2} \cos(k_x x) \cos(k_y y) \sin(k_z z) \cos(k_t x_t) \\
 \psi_2 &= 0 \\
 \psi_3 &= h_3 \cos(k_x x) \cos(k_y y) \cos(k_z z) \cos(k_t x_t)
 \end{aligned} \tag{VII-10}$$

Any linear combination of solutions to the system of DIRAC is still a solution of the system. One can thus construct alternative solutions, of which some examples are given below.

II - Travelling Solutions

A wave that spreads can be seen as the sum of two standing waves, allowing travelling solutions from previous tables (table VI-1 and 2).

One can for example choose the following modes, in which we reminded the notation of the phase time $\omega t = k_t x_t$:

$$\cos(k_x x) \cos(k_y y) \cos(k_z z) \cos(k_t x_t) + \cos(k_x x) \cos(k_y y) \sin(k_z z) \sin(k_t x_t) \tag{VII-11}$$

what gives after factorisation:

$$\cos(k_x x) \cos(k_y y) \{ \cos(k_z z) \cos(k_t x_t) + \sin(k_z z) \sin(k_t x_t) \} \tag{VII-12}$$

a stationary wave in x, y, and propagative wave along the z axis:

$$\cos(k_x x) \cos(k_y y) \cos(k_t x_t - k_z z) \tag{VII-13}$$

The approach is as follows: we write the solution corresponding to each of the modes we want to add in (VII-11), which gives for the first mode:

$$\psi_0 = h_0 \cos(k_x x) \cos(k_y y) \cos(k_z z) \cos(k_t x_t) \tag{VII-14}$$

$$\psi_1 = 0$$

$$\psi_2 = j h_0 \frac{\eta k_z}{\eta^2 - k_t^2} \cos(k_x x) \cos(k_y y) \sin(k_z z) \cos(k_t x_t) - h_0 \frac{k_t k_z}{\eta^2 - k_t^2} \cos(k_x x) \cos(k_y y) \sin(k_z z) \sin(k_t x_t)$$

$$\psi_3 = -h_0 \frac{\eta k_y}{\eta^2 - k_t^2} \cos(k_x x) \sin(k_y y) \cos(k_z z) \cos(k_t x_t) - j h_0 \frac{k_t k_y}{\eta^2 - k_t^2} \cos(k_x x) \sin(k_y y) \cos(k_z z) \sin(k_t x_t)$$

$$-h_0 \frac{k_t k_x}{\eta^2 - k_t^2} \sin(k_x x) \cos(k_y y) \cos(k_z z) \sin(k_t x_t) + j h_0 \frac{\eta k_x}{\eta^2 - k_t^2} \sin(k_x x) \cos(k_y y) \cos(k_z z) \cos(k_t x_t)$$

and for the second:

$$\psi_0 = l_0 \cos(k_x x) \cos(k_y y) \sin(k_z z) \sin(k_t x_t) \tag{VII-15}$$

$$\psi_1 = 0$$

$$\begin{aligned}\Psi_2 &= -jI_0 \frac{\eta k_z}{\eta^2 - k_t^2} \cos(k_x x) \cos(k_y y) \cos(k_z z) \sin(k_t x_t) - I_0 \frac{k_t k_z}{\eta^2 - k_t^2} \cos(k_x x) \cos(k_y y) \cos(k_z z) \cos(k_t x_t) \\ \Psi_3 &= -I_0 \frac{\eta k_y}{\eta^2 - k_t^2} \cos(k_x x) \sin(k_y y) \sin(k_z z) \sin(k_t x_t) + jI_0 \frac{k_t k_y}{\eta^2 - k_t^2} \cos(k_x x) \sin(k_y y) \sin(k_z z) \cos(k_t x_t) \\ &+ I_0 \frac{k_t k_x}{\eta^2 - k_t^2} \sin(k_x x) \cos(k_y y) \sin(k_z z) \cos(k_t x_t) + jI_0 \frac{\eta k_x}{\eta^2 - k_t^2} \sin(k_x x) \cos(k_y y) \sin(k_z z) \sin(k_t x_t)\end{aligned}$$

We put then: $h_0 = l_0 = A$, we sum term-to-term sum the solutions (VII-14) and (VII-15), in order to obtain the travelling solution along z after reduction:

$$\begin{aligned}\Psi_0 &= A \cos(k_x x) \cos(k_y y) \cos(k_t x_t - k_z z) \\ \Psi_1 &= 0 \\ \Psi_2 &= jA \frac{\eta k_z}{\eta^2 - k_t^2} \cos(k_x x) \cos(k_y y) \sin(k_t x_t - k_z z) - A \frac{k_t k_z}{\eta^2 - k_t^2} \cos(k_x x) \cos(k_y y) \cos(k_t x_t - k_z z) \\ \Psi_3 &= -A \frac{\eta k_y}{\eta^2 - k_t^2} \cos(k_x x) \sin(k_y y) \cos(k_t x_t - k_z z) + jA \frac{k_t k_y}{\eta^2 - k_t^2} \cos(k_x x) \sin(k_y y) \sin(k_t x_t - k_z z) \\ &+ A \frac{k_t k_x}{\eta^2 - k_t^2} \sin(k_x x) \cos(k_y y) \sin(k_t x_t - k_z z) + jA \frac{\eta k_x}{\eta^2 - k_t^2} \sin(k_x x) \cos(k_y y) \cos(k_t x_t - k_z z)\end{aligned}\tag{VII-16}$$

In reiterating this process with the other two directions, we can develop travelling solutions in x, y, z whose example is given below:

$$\begin{aligned}\Psi_0 &= \eta \cos(k_t x_t - k_x x - k_y y - k_z z) - jk_t \sin(k_t x_t - k_x x - k_y y - k_z z) \\ \Psi_1 &= 0 \\ \Psi_2 &= -jk_z \sin(k_t x_t - k_x x - k_y y - k_z z) \\ \Psi_3 &= k_y \sin(k_t x_t - k_x x - k_y y - k_z z) - jk_x \sin(k_t x_t - k_x x - k_y y - k_z z)\end{aligned}\tag{VII-17}$$

III - Other solutions

Basic solutions expressed in tables 1 and 2 above may also be combined to get solutions whose shape is a little different. Two examples are proposed.

In the first example, two modes of the wave function Ψ_1 are excited, one weighted by (-1) , and the other weighted by the term $(j\eta/k_t)$:

$$\begin{aligned}\Psi_0 &= 0 \\ \Psi_1 &= -\sin(k_x x) \cos(k_y y) \cos(k_z z) \cos(k_t x_t) + j \frac{\eta}{k_t} \sin(k_x x) \cos(k_y y) \cos(k_z z) \sin(k_t x_t) \\ \Psi_2 &= j \frac{k_y}{k_t} \sin(k_x x) \sin(k_y y) \cos(k_z z) \sin(k_t x_t) + \frac{k_x}{k_t} \cos(k_x x) \cos(k_y y) \cos(k_z z) \sin(k_t x_t) \\ \Psi_3 &= \frac{k_z}{k_t} \sin(k_x x) \cos(k_y y) \sin(k_z z) \sin(k_t x_t)\end{aligned}\tag{VII-18}$$

In the second example, these are two modes of the wave function ψ_2 who are excited, one weighted by (-1), and the other-weighted term ($j\eta/k_t$):

$$\psi_0 = -\frac{k_z}{k_t} \cos(k_x x) \cos(k_y y) \cos(k_z z) \cos(k_t x_t)$$

$$\psi_1 = \frac{k_x}{k_t} \sin(k_x x) \cos(k_y y) \sin(k_z z) \cos(k_t x_t) + j \frac{k_y}{k_t} \cos(k_x x) \sin(k_y y) \sin(k_z z) \cos(k_t x_t)$$

$$\psi_2 = j \frac{\eta}{k_t} \cos(k_x x) \cos(k_y y) \sin(k_z z) \cos(k_t x_t) - \cos(k_x x) \cos(k_y y) \sin(k_z z) \sin(k_t x_t)$$

$$\psi_3 = 0$$

(VII-19)

VIII

The wave-particle duality

When the particle comes in the form of material, it obeys the conservation equation of energy from special relativity:

$$E^2 = p^2 c^2 + m_0^2 c^4 \quad (\text{VIII-1})$$

When the particle is in wave form, its total energy and its wave vector are such as:

$$E = \hbar\omega \quad p_x = \hbar k_x \quad p_y = \hbar k_y \quad p_z = \hbar k_z \quad (\text{VIII-2})$$

By introducing these relations in (VIII-1), we can deduce that wave quantities must respect the following relationship for compatibility with energy conservation imposed by relativity:

$$\hbar^2 \omega^2 = \hbar^2 (k_x^2 + k_y^2 + k_z^2) c^2 + m_0^2 c^4 \quad (\text{VIII-3})$$

or again:

$$\frac{\omega^2}{c^2} = (k_x^2 + k_y^2 + k_z^2) + \frac{m_0^2 c^2}{\hbar^2} \quad (\text{VIII-4})$$

One found exactly the relationship required to get solutions to the system of DIRAC.

We can deduce that this relationship, associated with relationships (VIII-2) expresses the conservation of energy, both if the particle presents itself in the form of material or in the wave form.

The quantum relationship of conservation of energy:

$$\hbar^2 \omega^2 = \hbar^2 c^2 (k_x^2 + k_y^2 + k_z^2) + (m_0 c^2)^2 \quad (\text{VIII-5})$$

will play a fundamental role in the analysis of stationary solutions which will be proposed.

As for the relationship (VIII-1) issue of relativity, we distinguish three types of energy:

$\hbar\omega$: Wave energy in reference to the pulse ω that appears in this expression.

$\hbar k_x, \hbar k_y, \hbar k_z$: Impulse energy following the directions x, y, z.

m_0c^2 : Mass energy

Among the different forms of solution, we choose one that allows a direct interpretation of the role of these energies in the DIRAC bispinor wave functions.

We take as starting point a solution expressed in the previous chapter:

$$\begin{aligned}
 \psi_0 &= 0 \\
 \psi_1 &= -\sin(k_x x) \cos(k_y y) \cos(k_z z) \cos(k_t x_t) + j \frac{\eta}{k_t} \sin(k_x x) \cos(k_y y) \cos(k_z z) \sin(k_t x_t) \\
 \psi_2 &= j \frac{k_y}{k_t} \sin(k_x x) \sin(k_y y) \cos(k_z z) \sin(k_t x_t) + \frac{k_x}{k_t} \cos(k_x x) \cos(k_y y) \cos(k_z z) \sin(k_t x_t) \\
 \psi_3 &= \frac{k_z}{k_t} \sin(k_x x) \cos(k_y y) \sin(k_z z) \sin(k_t x_t)
 \end{aligned}
 \tag{VIII-6}$$

In this solution, the wave functions are unitless. But we have, from a mathematical perspective, any freedom to multiply all of these wave functions by a constant, and one that seems indicated in this case is equal to k_t to get:

$$\begin{aligned}
 \psi_0 &= 0 \\
 \psi_1 &= -k_t \sin(k_x x) \cos(k_y y) \cos(k_z z) \cos(k_t x_t) + j\eta \sin(k_x x) \cos(k_y y) \cos(k_z z) \sin(k_t x_t) \\
 \psi_2 &= jk_y \sin(k_x x) \sin(k_y y) \cos(k_z z) \sin(k_t x_t) + k_x \cos(k_x x) \cos(k_y y) \cos(k_z z) \sin(k_t x_t) \\
 \psi_3 &= k_z \sin(k_x x) \cos(k_y y) \sin(k_z z) \sin(k_t x_t)
 \end{aligned}
 \tag{VIII-7}$$

Substituting the expression $\eta = \frac{m_0c}{\hbar}$, and multiplying again all wave functions by the constant quantity $\hbar c$, we get:

$$\begin{aligned}
 \psi_0 &= 0 \\
 \psi_1 &= -\hbar\omega \sin(k_x x) \cos(k_y y) \cos(k_z z) \cos(k_t x_t) + jm_0c^2 \sin(k_x x) \cos(k_y y) \cos(k_z z) \sin(k_t x_t) \\
 \psi_2 &= j\hbar ck_y \sin(k_x x) \sin(k_y y) \cos(k_z z) \sin(k_t x_t) + \hbar ck_x \cos(k_x x) \cos(k_y y) \cos(k_z z) \sin(k_t x_t) \\
 \psi_3 &= \hbar ck_z \sin(k_x x) \cos(k_y y) \sin(k_z z) \sin(k_t x_t)
 \end{aligned}
 \tag{VIII-8}$$

Each wave function now has an energy dimension, and one make the observation that in this solution, each term contains an energy of different nature, considering that two pulse energies in orthogonal directions are necessarily differentiated.

If we look at a point where the mass energy is maximum, we must have at this point $|\sin(k_x x)| = |\cos(k_y y)| = |\cos(k_z z)| = 1$ which returns to put $\cos(k_x x) = \sin(k_y y) = \sin(k_z z) = 0$. The solution then takes the form:

$$\begin{aligned}
 \psi_0 &= 0 \\
 \psi_1 &= -\hbar\omega \sin(k_x x) \cos(k_y y) \cos(k_z z) \cos(k_t x_t) + jm_0 c^2 \sin(k_x x) \cos(k_y y) \cos(k_z z) \sin(k_t x_t) \\
 \psi_2 &= 0 \\
 \psi_3 &= 0
 \end{aligned}
 \tag{VIII-9}$$

The wave function ψ_1 must retain special attention: when the mass energy is maximum, it takes the form:

$$\psi_1 = \left\{ -\hbar\omega \cos(k_t x_t) + jm_0 c^2 \sin(k_t x_t) \right\}
 \tag{VIII-10}$$

where it is recognized the wave energy $\hbar\omega$ and the mass energy $m_0 c^2$. But the remarkable result that teaches us the relationship (VIII-9) is that these energies are evolving in time quadrature, and that when one is maximum, the other is minimum.

In other words, when the particle is in its total mass form, it has no wave energy, and when it occurs in its total wave form, it presents no mass energy. Energy present in the particle so alternates between mass and wave forms to the pulse ω defined by the equation of conservation of quantum energy, that for energy pulse equal to zero is simply written:

$$\hbar^2 \omega^2 = (m_0 c^2)^2
 \tag{VIII-11}$$

It can be assumed that it is in this ongoing exchange of energy that lies the mystery of the wave-particle duality which appears, in the light of the relationship (VIII-10), sometimes in the form of mass, sometimes in wave form.

In the general case, it is still the wave function ψ_1 which brings these energy exchanges. The terms that carry the impulse energy are those of the second spinor:

$$\begin{aligned}
 \psi_2 &= j\hbar c k_y \sin(k_x x) \sin(k_y y) \cos(k_z z) \sin(k_t x_t) + \hbar c k_x \cos(k_x x) \cos(k_y y) \cos(k_z z) \sin(k_t x_t) \\
 \psi_3 &= \hbar c k_z \sin(k_x x) \cos(k_y y) \sin(k_z z) \sin(k_t x_t)
 \end{aligned}
 \tag{VIII-12}$$

It is recalled that all these energy exchanges must agree with the law of energy conservation:

$$\hbar^2 \omega^2 = \hbar^2 c^2 (k_x^2 + k_y^2 + k_z^2) + (m_0 c^2)^2
 \tag{VIII-13}$$

The fact that the pulse energy is carried by the second spinor seem to be understood by imagining that the antiparticle corresponds to the situation in which the second spinor deals with exchanges between mass energy and wave energy. On the basis of this hypothesis, one which is detected in an experiment is that whose spinor contains the energy of mass in the solution of DIRAC.

The minimum pulse ω_0 at which this exchange of energy is performed, is given by the relationship:

$$\omega_0 = \frac{m_0 c^2}{\hbar} \quad (\text{VIII-14})$$

The numerical application for an electron gives:

$$\omega_0 = \frac{m_0 c^2}{\hbar} = \frac{(9,11 \cdot 10^{-31}) (3 \cdot 10^8)^2}{1,05 \cdot 10^{-34}} \approx 7,8 \cdot 10^{20} \text{ rd/s} \quad (\text{VIII-15})$$

The great value of this pulse could explain the great difficulty to see these energy exchanges from an experimental point of view.

IX

The currents of DIRAC (1)

The expression of these currents has been given in one of the previous chapters. It is recalled for memory:

$$\begin{aligned} J^0 &= \psi_0^* \psi_0 + \psi_1^* \psi_1 + \psi_2^* \psi_2 + \psi_3^* \psi_3 \\ J^1 &= \psi_3^* \psi_0 + \psi_2^* \psi_1 + \psi_1^* \psi_2 + \psi_0^* \psi_3 \\ J^2 &= j\psi_3^* \psi_0 - j\psi_2^* \psi_1 + j\psi_1^* \psi_2 - j\psi_0^* \psi_3 \\ J^3 &= \psi_2^* \psi_0 - \psi_3^* \psi_1 + \psi_0^* \psi_2 - \psi_1^* \psi_3 \end{aligned} \quad (\text{IX-1})$$

These currents check the local conservation of energy equation:

$$\frac{\partial J^0}{\partial(ct)} + \frac{\partial J^1}{\partial x} + \frac{\partial J^2}{\partial y} + \frac{\partial J^3}{\partial z} = 0 \quad (\text{IX-2})$$

I - DIRAC currents for a stationary solution

A practical calculation of these currents will be performed for a stationary solution. We choose for this the solution discussed in the previous chapter:

$$\begin{aligned} \psi_0 &= 0 \\ \psi_1 &= -k_t \sin(k_x x) \cos(k_y y) \cos(k_z z) \cos(k_t x_t) + j\eta \sin(k_x x) \cos(k_y y) \cos(k_z z) \sin(k_t x_t) \\ \psi_2 &= jk_y \sin(k_x x) \sin(k_y y) \cos(k_z z) \sin(k_t x_t) + k_x \cos(k_x x) \cos(k_y y) \cos(k_z z) \sin(k_t x_t) \\ \psi_3 &= k_z \sin(k_x x) \cos(k_y y) \sin(k_z z) \sin(k_t x_t) \end{aligned} \quad (\text{IX-3})$$

Calculations are a little long, but without difficulties. Finally, we get the following expressions:

$$\begin{aligned} J^0 &= k_t^2 \sin^2(k_x x) \cos^2(k_y y) \cos^2(k_z z) \cos^2(k_t x_t) + \eta^2 \sin^2(k_x x) \cos^2(k_y y) \cos^2(k_z z) \sin^2(k_t x_t) \\ &+ k_y^2 \sin^2(k_x x) \sin^2(k_y y) \cos^2(k_z z) \sin^2(k_t x_t) + k_x^2 \cos^2(k_x x) \cos^2(k_y y) \cos^2(k_z z) \sin^2(k_t x_t) \\ &+ k_z^2 \sin^2(k_x x) \cos^2(k_y y) \sin^2(k_z z) \sin^2(k_t x_t) \end{aligned} \quad (\text{IX-4})$$

$$J^1 = \begin{pmatrix} \eta k_y \sin^2(k_x x) \sin(2k_y y) \cos^2(k_z z) \sin^2(k_t x_t) \\ -k_x k_t \sin(2k_x x) \cos^2(k_y y) \cos^2(k_z z) \sin(k_t x_t) \cos(k_t x_t) \end{pmatrix} \quad (IX-5)$$

$$J^2 = \begin{pmatrix} k_y k_t \sin^2(k_x x) \sin(2k_y y) \cos^2(k_z z) \sin(k_t x_t) \cos(k_t x_t) \\ k_x \eta \sin(2k_x x) \cos^2(k_y y) \cos^2(k_z z) \sin^2(k_t x_t) \end{pmatrix} \quad (IX-6)$$

$$J^3 = k_t k_z \sin^2(k_x x) \cos^2(k_y y) \sin(2k_z z) \sin(k_t x_t) \cos(k_t x_t) \quad (IX-7)$$

We can then verify the conservation equation (IX-2). The details of the calculations is not given since it is not hard. Yields:

$$\frac{\partial J^0}{\partial x_t} + \frac{\partial J^1}{\partial x} + \frac{\partial J^2}{\partial y} + \frac{\partial J^3}{\partial z} = k_t \sin(2k_t x_t) \sin^2(k_x x) \cos^2(k_y y) \cos^2(k_z z) (-k_t^2 + \eta^2 + k_x^2 + k_y^2 + k_z^2) \quad (IX-8)$$

which, taking into account the relationship of conservation of energy:

$$k_t^2 = k_x^2 + k_y^2 + k_z^2 + \eta^2 \quad (IX-9)$$

leads to the expected result:

$$\frac{\partial J^0}{\partial x_t} + \frac{\partial J^1}{\partial x} + \frac{\partial J^2}{\partial y} + \frac{\partial J^3}{\partial z} = 0 \quad (IX-10)$$

We propose to reconsider the above results in a more physical approach. This leads to multiply the wave functions by the amount $\hbar c$ and DIRAC current by the amount $(\hbar c)^2$. In this description, the wave function has the form:

$$\begin{aligned} \psi_0 &= 0 \\ \psi_1 &= -\hbar \omega \sin(k_x x) \cos(k_y y) \cos(k_z z) \cos(k_t x_t) + j m_0 c^2 \sin(k_x x) \cos(k_y y) \cos(k_z z) \sin(k_t x_t) \\ \psi_2 &= j \hbar c k_y \sin(k_x x) \sin(k_y y) \cos(k_z z) \sin(k_t x_t) + \hbar c k_x \cos(k_x x) \cos(k_y y) \cos(k_z z) \sin(k_t x_t) \\ \psi_3 &= \hbar c k_z \sin(k_x x) \cos(k_y y) \sin(k_z z) \sin(k_t x_t) \end{aligned} \quad (IX-11)$$

From a mathematical point of view, we know that this wave function is defined to a multiplicative constant close, which we will call C, so that it can be put in a more general form:

$$\begin{aligned} \psi_0 &= 0 \\ \psi_1 &= C \left\{ -\hbar \omega \sin(k_x x) \cos(k_y y) \cos(k_z z) \cos(k_t x_t) + j m_0 c^2 \sin(k_x x) \cos(k_y y) \cos(k_z z) \sin(k_t x_t) \right\} \\ \psi_2 &= C \left\{ j \hbar c k_y \sin(k_x x) \sin(k_y y) \cos(k_z z) \sin(k_t x_t) + \hbar c k_x \cos(k_x x) \cos(k_y y) \cos(k_z z) \sin(k_t x_t) \right\} \\ \psi_3 &= C \left\{ \hbar c k_z \sin(k_x x) \cos(k_y y) \sin(k_z z) \sin(k_t x_t) \right\} \end{aligned} \quad (IX-12)$$

The constant C must allow:

- 1 – to give to the quantity $\psi\psi^*$ the dimension of a volumetric energy density.
- 2 – to ensure that the totality of the particle energy is confined in a parallelepiped with the dimension along x is between X_1 and X_2 , along y between Y_1 and Y_2 , and along z between Z_1 , Z_2 , so its volume V is equal to:

$$V = (X_2 - X_1) (Y_2 - Y_1) (Z_2 - Z_1) \quad (\text{IX-13})$$

The term J^0 of the currents of DIRAC represents the total energy volume density included in the box. It has for expression:

$$\begin{aligned} J^0 = & C^2 (\hbar\omega)^2 \sin^2(k_x x) \cos^2(k_y y) \cos^2(k_z z) \cos^2(k_t x_t) \\ & + C^2 (m_0 c^2)^2 \sin^2(k_x x) \cos^2(k_y y) \cos^2(k_z z) \sin^2(k_t x_t) \\ & + C^2 (\hbar c k_y)^2 \sin^2(k_x x) \sin^2(k_y y) \cos^2(k_z z) \sin^2(k_t x_t) \\ & + C^2 (\hbar c k_x)^2 \cos^2(k_x x) \cos^2(k_y y) \cos^2(k_z z) \sin^2(k_t x_t) \\ & + C^2 (\hbar c k_z)^2 \sin^2(k_x x) \cos^2(k_y y) \sin^2(k_z z) \sin^2(k_t x_t) \end{aligned} \quad (\text{IX-14})$$

To get the total energy E in the parallelepiped, we need to integrate on the volume of this latter, hence:

$$E = \iiint_V \left\{ \begin{aligned} & C^2 (\hbar\omega)^2 \sin^2(k_x x) \cos^2(k_y y) \cos^2(k_z z) \cos^2(k_t x_t) \\ & + C^2 (m_0 c^2)^2 \sin^2(k_x x) \cos^2(k_y y) \cos^2(k_z z) \sin^2(k_t x_t) \\ & + C^2 (\hbar c k_y)^2 \sin^2(k_x x) \sin^2(k_y y) \cos^2(k_z z) \sin^2(k_t x_t) \\ & + C^2 (\hbar c k_x)^2 \cos^2(k_x x) \cos^2(k_y y) \cos^2(k_z z) \sin^2(k_t x_t) \\ & + C^2 (\hbar c k_z)^2 \sin^2(k_x x) \cos^2(k_y y) \sin^2(k_z z) \sin^2(k_t x_t) \end{aligned} \right\} dx dy dz \quad (\text{IX-15})$$

We will do the classical hypothesis that stationary modes settle in conditions such as the dimensions of the box are multiples of the half-wavelength along each of the x, y, z direction.

$$\begin{aligned} (X_2 - X_1) &= n \frac{\lambda_x}{2} = n \frac{\pi}{k_x} \\ (Y_2 - Y_1) &= m \frac{\lambda_y}{2} = m \frac{\pi}{k_y} \\ (Z_2 - Z_1) &= p \frac{\lambda_z}{2} = p \frac{\pi}{k_z} \end{aligned} \quad (\text{IX-16})$$

where m, n, p are positive or negative integers. Under these conditions, one has the following property:

$$\begin{aligned}
 \int_{X_1}^{X_2} \sin^2(k_x x) dx &= \int_{X_1}^{X_2} \cos^2(k_x x) dx = \frac{X_2 - X_1}{2} \\
 \int_{Y_1}^{Y_2} \sin^2(k_y y) dy &= \int_{Y_1}^{Y_2} \cos^2(k_y y) dy = \frac{Y_2 - Y_1}{2} \\
 \int_{Z_1}^{Z_2} \sin^2(k_z z) dz &= \int_{Z_1}^{Z_2} \cos^2(k_z z) dz = \frac{Z_2 - Z_1}{2}
 \end{aligned} \tag{IX-17}$$

and the total energy calculated according to the relationship (IX-15) takes the form:

$$E = C^2 \frac{(X_2 - X_1)(Y_2 - Y_1)(Z_2 - Z_1)}{8} \left\{ \begin{aligned} &(\hbar\omega)^2 \cos^2(k_t x_t) \\ &+ (m_0 c^2)^2 \sin^2(k_t x_t) \\ &+ (\hbar c k_y)^2 \sin^2(k_t x_t) \\ &+ (\hbar c k_x)^2 \sin^2(k_t x_t) \\ &+ (\hbar c k_z)^2 \sin^2(k_t x_t) \end{aligned} \right\} \tag{IX-18}$$

or again:

$$E = C^2 \frac{V}{8} \left\{ (\hbar\omega)^2 \cos^2(k_t x_t) + \left[(m_0 c^2)^2 + (\hbar c k_y)^2 + (\hbar c k_x)^2 + (\hbar c k_z)^2 \right] \sin^2(k_t x_t) \right\} \tag{IX-19}$$

From relationship:

$$(\hbar\omega)^2 = (m_0 c^2)^2 + (\hbar c k_y)^2 + (\hbar c k_x)^2 + (\hbar c k_z)^2 \tag{IX-20}$$

It is deduced that the total energy in the cavity is either under the form of a wave energy, either in the form of a combination of mass and impulse energy. These energy exchanges are in time quadrature, as shown in the (IX-19) relationship: when one is maximum, the other is null and vice versa.

From relations (IX-19) and (IX-20) are deduced:

$$E = C^2 \frac{V}{8} (\hbar\omega)^2 \tag{IX-21}$$

and since the total energy E is equal to $\hbar\omega$, this imposes to the C constant the following value:

$$C = \sqrt{\frac{8}{(\hbar\omega)V}} \tag{IX-22}$$

After taking into account of this multiplicative constant, DIRAC currents become homogeneous to a volumetric energy density, and the total energy is normalized to the energy of the particle.

The four-divergence:

$$\frac{\partial J^0}{\partial x_t} + \frac{\partial J^1}{\partial x} + \frac{\partial J^2}{\partial y} + \frac{\partial J^3}{\partial z} = 0 \quad (\text{IX-23})$$

is interpreted in the same manner as in electromagnetism (POYNTING theorem) or in general relativity. It expresses the fact that if there is a change in energy in a volume $dV = dx dy dz$ during a time element dt , it's because this variation has crossed the border defined by the closed surface bounding the volume element.

II - DIRAC currents for a propagative solution

As an example, we choose a propagative solution in x, y, z :

$$\psi_0 = \eta \cos(k_t x_t - k_x x - k_y y - k_z z) - j k_t \sin(k_t x_t - k_x x - k_y y - k_z z) \quad (\text{IX-24})$$

$$\psi_1 = 0$$

$$\psi_2 = -j k_z \sin(k_t x_t - k_x x - k_y y - k_z z)$$

$$\psi_3 = k_y \sin(k_t x_t - k_x x - k_y y - k_z z) - j k_x \sin(k_t x_t - k_x x - k_y y - k_z z)$$

Calculations give the following DIRAC currents:

$$\begin{aligned} J^0 &= \eta^2 \cos^2(k_t x_t - k_x x - k_y y - k_z z) + k_t^2 \sin^2(k_t x_t - k_x x - k_y y - k_z z) \\ &+ k_x^2 \sin^2(k_t x_t - k_x x - k_y y - k_z z) + k_y^2 \sin^2(k_t x_t - k_x x - k_y y - k_z z) \\ &+ k_z^2 \sin^2(k_t x_t - k_x x - k_y y - k_z z) \end{aligned} \quad (\text{IX-25})$$

$$\begin{aligned} J^1 &= 2k_y \eta \sin(k_t x_t - k_x x - k_y y - k_z z) \cos(k_t x_t - k_x x - k_y y - k_z z) \\ &+ 2k_x k_t \sin^2(k_t x_t - k_x x - k_y y - k_z z) \end{aligned} \quad (\text{IX-26})$$

$$\begin{aligned} J^2 &= -2k_x \sin(k_t x_t - k_x x - k_y y - k_z z) \cos(k_t x_t - k_x x - k_y y - k_z z) + \\ &(2k_y k_t \sin^2(k_t x_t - k_x x - k_y y - k_z z)) \end{aligned} \quad (\text{IX-27})$$

$$J^3 = (2k_t k_z \sin^2(k_t x_t - k_x x - k_y y - k_z z)) \quad (\text{IX-28})$$

The current J^0 is of particular interest because it contains, to a multiplicative constant close, the total energy of the particle. It may come in the form:

$$J^0 = \eta^2 \cos^2(k_t x_t - k_x x - k_y y - k_z z) + (k_t^2 + k_x^2 + k_y^2 + k_z^2) \sin^2(k_t x_t - k_x x - k_y y - k_z z) \quad (\text{IX-29})$$

After multiplication by the constant $(\hbar c)^2$ and standardization by the C^2 constant defined in the previous paragraph, it represents the total volume density attached to the particle.

Using the relation of energy conservation:

$$k_t^2 = k_x^2 + k_y^2 + k_z^2 + \eta^2 \quad (\text{IX-30})$$

it becomes:

$$J^0 = k_t^2 - (k_x^2 + k_y^2 + k_z^2) \left\{ \cos^2(k_t x_t - k_x x - k_y y - k_z z) - \sin^2(k_t x_t - k_x x - k_y y - k_z z) \right\} \quad (\text{IX-31})$$

or again:

$$J^0 = k_t^2 - (k_x^2 + k_y^2 + k_z^2) \cos[2(k_t x_t - k_x x - k_y y - k_z z)] \quad (\text{IX-32})$$

Total energy volume density is given by the term $C^2(\hbar c k_t)^2$. It fluctuates around this value with a spatial and temporal variation in average which is zero.

The local conservation of energy equation:

$$\frac{\partial J^0}{\partial x_t} + \frac{\partial J^1}{\partial x} + \frac{\partial J^2}{\partial y} + \frac{\partial J^3}{\partial z} = 0 \quad (\text{IX-33})$$

gives the following result:

$$\frac{\partial J^0}{\partial x_t} + \frac{\partial J^1}{\partial x} + \frac{\partial J^2}{\partial y} + \frac{\partial J^3}{\partial z} = (k_t^2 - k_x^2 - k_y^2 - k_z^2 - \eta^2) \sin[2(k_t x_t - k_x x - k_y y - k_z z)] = 0 \quad (\text{IX-34})$$

III - Uniform DIRAC currents

We will name uniform currents, currents that are not dependent neither on time nor on space. There are exact solutions to the DIRAC equation which have this property, of which an example is provided below:

$$\begin{aligned} \psi_0 &= k_z \exp \{j(-k_t x_t + k_x x + k_y y + k_z z)\} \\ \psi_1 &= k_x \exp \{j(-k_t x_t + k_x x + k_y y + k_z z)\} + j k_y \exp \{j(-k_t x_t + k_x x + k_y y + k_z z)\} \\ \psi_2 &= (k_t - \eta) \exp \{j(-k_t x_t + k_x x + k_y y + k_z z)\} \\ \psi_3 &= 0 \end{aligned} \quad (\text{IX-35})$$

Because products of wave functions by conjugated wave functions eliminate the exponential, we find:

$$\begin{aligned} J^0 &= k_x^2 + k_y^2 + k_z^2 + (k_t - \eta)^2 \\ J^1 &= 2(k_t - \eta)k_x \\ J^2 &= 2(k_t - \eta)k_y \\ J^3 &= 2(k_t - \eta)k_z \end{aligned} \quad (\text{IX-36})$$

One observe that the volume density of total energy $C^2(\hbar c)^2 J^0$ presents itself in the form of a combination of all energies working into the cavity, and it is uniform, that is to say independent of space and time.

This excitement of some modes making it uniform energy within the particle density can occur only under specific conditions. Indeed, we know that the total energy within the particle density is equal to $C^2(\hbar c k_t)^2$, hence the relationship:

$$J^0 = k_x^2 + k_y^2 + k_z^2 + (\eta - k_t)^2 = k_x^2 + k_y^2 + k_z^2 + \eta^2 + k_t^2 - 2\eta k_t = k_t^2 \quad (\text{IX-37})$$

By introducing the relationship of energy conservation, we can deduce:

$$2k_t^2 - 2\eta k_t = k_t^2 \quad (\text{IX-38})$$

still, by substituting expressions of k_t and η :

$$\eta = \frac{m_0 c}{\hbar} = \frac{\omega}{2c} \quad (\text{IX-39})$$

The condition for obtaining a uniform density is therefore given by the relationship:

$$\hbar c k_t = 2m_0 c^2 \quad (\text{IX-40})$$

This relationship expresses the fact it needs a total internal energy equal to twice the mass energy particle to allow the installation of such modes.

We can connect this observation to the fact that all modes expressed in the solution (IX-35) exchange all kinds of energy in time and space quadrature between positive and negative energies identified by factor $j = \text{square root}(-1)$. The antiparticle can appear only if the total energy is at least twice the mass energy of the particle.

X

Principle of indeterminacy

This principle, enunciated by HEISENBERG, during the early days of quantum mechanics, was popularized in the expression: "it is impossible to know both the position and momentum of a particle". From a physical point of view, it is whole contained in a relationship that connects the uncertainty on the Δx position and the uncertainty on momentum Δp_x of a particle in the quantum world:

$$\Delta x \cdot \Delta p_x > \frac{\hbar}{2} \quad (\text{X-1})$$

We can deduce an alternative formulation by noting that the fundamental principle of the dynamic allows to write that the variation of the amount of movement Δp_x is done through an outdoor action F_x called force, while a duration of Δt :

$$\Delta p_x = F_x \cdot \Delta t \quad (\text{X-2})$$

Yields, noting that energy can be seen as the product of a force by displacement:

$$\Delta x \cdot \Delta p_x = \Delta x \cdot F_x \cdot \Delta t = \Delta E \cdot \Delta t \geq \frac{\hbar}{2} \quad (\text{X-3})$$

This principle has solid theoretical foundations, based on the fact that the position and the momentum of quantum operators do not commute. Since the result of the measurement of position and pulse, made at the same place and at the same time, depends on the order in which it performs this measure, this indicates that there is necessarily an uncertainty on the result of these measures.

If we now consider an exact stationary solution of the DIRAC equation such as that which has been chosen as an example in the previous chapters:

$$\begin{aligned} \psi_0 &= 0 \\ \psi_1 &= C \left\{ -\hbar \omega \sin(k_x x) \cos(k_y y) \cos(k_z z) \cos(k_t x_t) + j m_0 c^2 \sin(k_x x) \cos(k_y y) \cos(k_z z) \sin(k_t x_t) \right\} \\ \psi_2 &= C \left\{ j \hbar c k_y \sin(k_x x) \sin(k_y y) \cos(k_z z) \sin(k_t x_t) + \hbar c k_x \cos(k_x x) \cos(k_y y) \cos(k_z z) \sin(k_t x_t) \right\} \\ \psi_3 &= C \left\{ \hbar c k_z \sin(k_x x) \cos(k_y y) \sin(k_z z) \sin(k_t x_t) \right\} \end{aligned} \quad (\text{X-4})$$

there is a question that naturally comes to mind. The solution (X-4) above is perfectly deterministic: each type of energy is known, in theory, with infinite precision for a position (x,

y, z) and an instant (t) given. This state indeed seems in contradiction with the HEISENBERG uncertainty principle.

To remove this contradiction, we must first admit in the form of postulate the following conclusion: an observer can obtain information from a physical system only if it exchanges energy with this system. A corollary of this assumption is that two systems that do not exchange energy ignore each other and do not interact: they can work simultaneously at the same time and in the same place.

On the basis of this assumption, we examine, for the above solution (X-4), the volumetric energy density present in the particle which has been calculated in the previous chapter:

$$\begin{aligned}
 J^0 &= C^2(\hbar\omega)^2 \sin^2(k_x x) \cos^2(k_y y) \cos^2(k_z z) \cos^2(k_t x_t) \\
 &+ C^2(m_0 c^2)^2 \sin^2(k_x x) \cos^2(k_y y) \cos^2(k_z z) \sin^2(k_t x_t) \\
 &+ C^2(\hbar c k_y)^2 \sin^2(k_x x) \sin^2(k_y y) \cos^2(k_z z) \sin^2(k_t x_t) \\
 &+ C^2(\hbar c k_x)^2 \cos^2(k_x x) \cos^2(k_y y) \cos^2(k_z z) \sin^2(k_t x_t) \\
 &+ C^2(\hbar c k_z)^2 \sin^2(k_x x) \cos^2(k_y y) \sin^2(k_z z) \sin^2(k_t x_t)
 \end{aligned} \tag{X-5}$$

There are different types of energy outlined in the previous chapters:

- the total energy or wave energy: $\hbar\omega$
- the energy of mass: $m_0 c^2$
- pulse energy: $\hbar c k_x, \hbar c k_y, \hbar c k_z$

Let us place on a point in space (x, y, z) where the volume density of mass energy of the particle is maximum. Let us assume that the position where this mass energy is maximum can be determined with any precision desired. To ensure that this condition is achieved, the coordinates x, y, z must check:

$$\sin^2(k_x x) = \cos^2(k_y y) = \cos^2(k_z z) = 1 \tag{X-6}$$

What requires:

$$\cos^2(k_x x) = \sin^2(k_y y) = \sin^2(k_z z) = 0 \tag{X-7}$$

It appears the following remarkable result: all impulse energy densities present in the particle are zero at this location.

In other words, if we move to a point where we can, through an exchange of energy with the energy of mass, know with precision the position of the particle, we cannot get any information on its momentum at this point because its impulse energy is zero at this place.

The reciprocal is expressed in the following way: if one moves to a place where the impulse energy following x is maximum, then mass energy and impulse energy along y and z are zero. A similar property is checked by permutation on the variables x, y, z.

These observations allow to understand how a completely deterministic theory built on exact stationary solutions to the DIRAC equation remains compatible with the HEISENBERG uncertainty principle. This principle is based on the hypothesis that measurements of position and speed are made pointwise in the same place, while the energy approach shows that the energies corresponding to these two quantities are shifted in the space. If this approach proves to be correct, it can be concluded that it is possible to know the position and velocity of a particle with arbitrary precision, provided you locate in the place where these characteristics are present in the particle

The second relation of indeterminacy (X-3) which deals with the energy and time:

$$\Delta E \cdot \Delta t \geq \frac{\hbar}{2} \quad (\text{X-8})$$

gives rise to a somewhat different interpretation. It concerns total energy or wave energy, whose volume density is given by (X-5):

$$C^2(\hbar\omega)^2 \sin^2(k_x x) \cos^2(k_y y) \cos^2(k_z z) \cos^2(k_t x_t) \quad (\text{X-9})$$

The points of the coordinate (x, y, z) in space where this energy is maximum are the same as those where the mass energy is at a maximum, they obey therefore relations (X-6), and the volume density of the wave energy is written in these points:

$$C^2(\hbar\omega)^2 \cos^2(k_t x_t) \quad (\text{X-10})$$

It appears that the measure of this energy depends on the moment in which it is carried out, in the same way as previously impulse energy or mass energy depended on the place where they were measured, and then we have similar uncertainty relations.

XI

DE BROGLIE wavelength

This chapter aims to show that the stationary solutions of the DIRAC equation are fully compatible with the conclusions of Louis DE BROGLIE on the wavelength associated with the motion of each particle. It allows to make the link between the wave description of quantum mechanics and Relativistic description of the motion of a particle.

The particle is assumed to have a straight trajectory along the Oz axis with constant velocity v . However, it is likely to have stationary modes according to Ox and Oy directions. We fall in a similar situation well known in electromagnetism, which is that of a wave in a perfectly conducting rectangular waveguide. The elements presented in this chapter have a greater analogy to those involving guided propagation.

The relationship of conservation of energy requires, to a multiplicative constant close:

$$k_t^2 = k_x^2 + k_y^2 + k_z^2 + \eta^2 = \frac{\omega^2}{c^2} \quad (\text{XI-1})$$

Assuming that the particle moves along Oz, it is natural to consider the space pulse k_z deduced from (XI-1):

$$k_z^2 = \frac{\omega^2}{c^2} - (k_x^2 + k_y^2 + \eta^2) \quad (\text{XI-2})$$

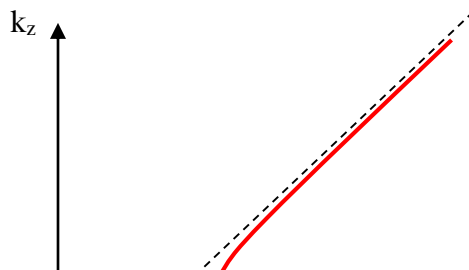
We name temporal cut pulse, the pulse ω_c such as:

$$\omega_c^2 = (k_x^2 + k_y^2 + \eta^2)c^2 \quad (\text{XI-3})$$

Which allows to express the space pulsation k_z in the form:

$$k_z^2 = \frac{\omega^2 - \omega_c^2}{c^2} \quad (\text{XI-4})$$

The curve representative $k_z = f(\omega)$ has the following look:



1/c

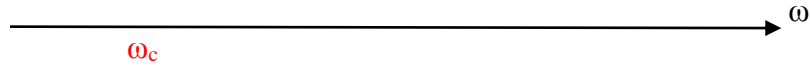


Figure (XI-1): *representation of the space pulsation versus temporal pulsation*

For a wave that will be named phase wave, and which propagates with a phase:

$$\varphi = \omega t - k_z z \quad (\text{XI-5})$$

we define the phase velocity as the velocity of the sliding of the phase wave:

$$v_\varphi = \frac{\omega}{k_z} \quad (\text{XI-6})$$

We can express this phase velocity depending on the cut pulse ω_c defined above, using the relationship (XI-4):

$$v_\varphi = \frac{c}{\sqrt{1 - \left(\frac{\omega_c}{\omega}\right)^2}} \quad (\text{XI-7})$$

Since $\omega > \omega_c$, this speed is greater than the speed of light. It has a physical reality as it represents the sliding of the phase velocity, but it cannot represent the speed of energy which must remain below c .

We define the speed of propagation of energy, or group velocity v_g by the derivative of the temporal pulse over the space pulsation:

$$v_g = \frac{d\omega}{dk_z} \quad (\text{XI-8})$$

We can notice on the figure (XI-1) that the speed of propagation of energy is zero for $\omega = \omega_c$, and that it tends to c for $\omega \gg \omega_c$. Energy cannot spread for $\omega < \omega_c$.

Group velocity can be expressed using the cut pulse, as has been done for the phase velocity. As a first step, one differentiates the relationship of energy conservation (XI-1) to obtain:

$$k_z \cdot dk_z = \frac{\omega \cdot d\omega}{c^2} \quad (\text{XI-9})$$

From which it takes an immediate relationship between phase velocity and group Velocity:

$$\frac{\omega}{k_z} \frac{d\omega}{dk_z} = c^2 \quad \text{and hence: } v_\phi v_g = c^2 \quad (\text{XI-10})$$

We can deduce:

$$v_g = \frac{d\omega}{dk_z} = \frac{k_z}{\omega} c^2 = \frac{c^2}{v_\phi} = c \sqrt{1 - \left(\frac{\omega_c}{\omega}\right)^2} \quad (\text{XI-11})$$

To make the link with the mechanical relativistic displacement of a particle, and following Louis DE BROGLIE, we admit that a particle of mass m_0 at rest has an energy at wave pulse ω_c since there is no displacement of energy:

$$\hbar\omega_c = m_0 c^2 \quad (\text{XI-12})$$

When this particle is moving at speed v , it presents an increase in its total energy given by the theory of relativity, which wave representation is going to the angular frequency ω :

$$\hbar\omega = \frac{m_0 c^2}{\sqrt{1 - \left(\frac{v}{c}\right)^2}} \quad (\text{XI-13})$$

On the basis of the assumptions (XI-12) and (XI-13), we can deduce that ω and ω_c pulse should check between them relations:

$$\left(\frac{\omega_c}{\omega}\right)^2 = 1 - \left(\frac{v}{c}\right)^2 \quad (\text{XI-14})$$

From the definition of group velocity (XI-11), we derived:

$$\left(\frac{\omega_c}{\omega}\right)^2 = 1 - \left(\frac{v_g}{c}\right)^2 \quad (\text{XI-15})$$

It follows from (XI-14) and (XI-15) that one can identify the speed of mass v to group velocity v_g of the wave's phase attached to the moveable mass. This identification ensures a representation of relativistic mass energy and quantum wave energy that is fully compatible.

The last step is to express the wavelength λ associated to phase wave of frequency ν and pulse ω that moves at the speed v_ϕ . From the previous paragraph, we deduced that it is associated, through its group velocity, to a mass m moving at speed v . By making use of the relationship $v_\phi v_g = c^2$ that establishes a relationship between phase velocity and group velocity or speed of the particle, one obtains:

$$\lambda = \frac{v_\phi}{v} = v_\phi \frac{2\pi}{\omega} = \frac{c^2}{v} \frac{2\pi}{\omega} = \frac{c^2}{v} \frac{2\pi}{\frac{m_0 c^2}{\hbar \sqrt{1-\beta^2}}} = \frac{h}{mv} \quad (\text{XI-16})$$

XII

Generalized DIRAC equation

This part deals with the DIRAC equation for a charged particle in an electromagnetic field characterized by a scalar potential ϕ and a vector potential (A_x, A_y, A_z). These potentials are considered constant and uniform, i.e. independent of x, y, z and t . We are looking for, as previously, a solution in the form of linear combinations of stationary modes which would settle into a rectangular cavity.

Following an approach similar to chapter II, the formalism of DIRAC leads to find the solutions of the new equation in which the electromagnetic potential four-vector is introduced:

$$\left\{ \gamma_0 \left(j \frac{\partial}{\partial x_t} - \frac{q\phi}{c\hbar} \right) + \gamma_1 \left(j \frac{\partial}{\partial x} - \frac{qA_x}{\hbar} \right) + \gamma_2 \left(j \frac{\partial}{\partial y} - \frac{qA_y}{\hbar} \right) + \gamma_3 \left(j \frac{\partial}{\partial z} - \frac{qA_z}{\hbar} \right) \right\} (\psi) = \frac{m_0 c}{\hbar} (\psi) \quad (\text{XII-1})$$

The wave function ψ is a bi-spinor with four components:

$$\psi = \begin{pmatrix} \psi_0 \\ \psi_1 \\ \psi_2 \\ \psi_3 \end{pmatrix} \quad (\text{XII-2})$$

The matrices $\gamma_0, \gamma_1, \gamma_2, \gamma_3$, are given in chapter II, and in order to simplify expressions, we put:

$$\eta = \frac{m_0 c}{\hbar} \quad (\text{XII-3})$$

Injecting matrices γ_i in the above equation (XII-3), we obtain the system of differential equations:

$$\begin{aligned}
 \eta\psi_0 &= j\left(\frac{\partial\psi_0}{\partial x_t} + j\frac{q\phi}{\hbar}\psi_0\right) + j\left(\frac{\partial\psi_3}{\partial x} + j\frac{qA_x}{\hbar}\psi_3\right) + \left(\frac{\partial\psi_3}{\partial y} + j\frac{qA_y}{\hbar}\psi_3\right) + j\left(\frac{\partial\psi_2}{\partial z} + j\frac{qA_z}{\hbar}\psi_2\right) \\
 \eta\psi_1 &= j\left(\frac{\partial\psi_1}{\partial x_t} + j\frac{q\phi}{\hbar}\psi_1\right) + j\left(\frac{\partial\psi_2}{\partial x} + j\frac{qA_x}{\hbar}\psi_2\right) - \left(\frac{\partial\psi_2}{\partial y} + j\frac{qA_y}{\hbar}\psi_2\right) - j\left(\frac{\partial\psi_3}{\partial z} + j\frac{qA_z}{\hbar}\psi_3\right) \\
 \eta\psi_2 &= -j\left(\frac{\partial\psi_2}{\partial x_t} + j\frac{q\phi}{\hbar}\psi_2\right) - j\left(\frac{\partial\psi_1}{\partial x} + j\frac{qA_x}{\hbar}\psi_1\right) - \left(\frac{\partial\psi_1}{\partial y} + j\frac{qA_y}{\hbar}\psi_1\right) - j\left(\frac{\partial\psi_0}{\partial z} + j\frac{qA_z}{\hbar}\psi_0\right) \\
 \eta\psi_3 &= -j\left(\frac{\partial\psi_3}{\partial x_t} + j\frac{q\phi}{\hbar}\psi_3\right) - j\left(\frac{\partial\psi_0}{\partial x} + j\frac{qA_x}{\hbar}\psi_0\right) + \left(\frac{\partial\psi_0}{\partial y} + j\frac{qA_y}{\hbar}\psi_0\right) + j\left(\frac{\partial\psi_1}{\partial z} + j\frac{qA_z}{\hbar}\psi_1\right)
 \end{aligned} \tag{XII-4}$$

or again, by isolating the differential system of the free particle:

$$\begin{aligned}
 \left(\eta + \frac{q\phi}{\hbar c}\right)\psi_0 + \frac{qA_x}{\hbar}\psi_3 - j\frac{qA_y}{\hbar}\psi_3 + \frac{qA_z}{\hbar}\psi_2 &= j\frac{\partial\psi_0}{\partial x_t} + j\frac{\partial\psi_3}{\partial x} + \frac{\partial\psi_3}{\partial y} + j\frac{\partial\psi_2}{\partial z} \\
 \left(\eta + \frac{q\phi}{\hbar c}\right)\psi_1 + \frac{qA_x}{\hbar}\psi_2 + j\frac{qA_y}{\hbar}\psi_2 - \frac{qA_z}{\hbar}\psi_3 &= j\frac{\partial\psi_1}{\partial x_t} + j\frac{\partial\psi_2}{\partial x} - \frac{\partial\psi_2}{\partial y} - j\frac{\partial\psi_3}{\partial z} \\
 \left(\eta - \frac{q\phi}{\hbar c}\right)\psi_2 - \frac{qA_x}{\hbar}\psi_1 + j\frac{qA_y}{\hbar}\psi_1 - \frac{qA_z}{\hbar}\psi_0 &= -j\frac{\partial\psi_2}{\partial x_t} - j\frac{\partial\psi_1}{\partial x} - \frac{\partial\psi_1}{\partial y} - j\frac{\partial\psi_0}{\partial z} \\
 \left(\eta - \frac{q\phi}{\hbar c}\right)\psi_3 - \frac{qA_x}{\hbar}\psi_0 - j\frac{qA_y}{\hbar}\psi_0 + \frac{qA_z}{\hbar}\psi_1 &= -j\frac{\partial\psi_3}{\partial x_t} - j\frac{\partial\psi_0}{\partial x} + \frac{\partial\psi_0}{\partial y} + j\frac{\partial\psi_1}{\partial z}
 \end{aligned} \tag{XII-5}$$

Again for a relief of notation, we put:

$$\eta_t = \frac{q\phi}{\hbar c} \quad \eta_x = \frac{qA_x}{\hbar} \quad \eta_y = \frac{qA_y}{\hbar} \quad \eta_z = \frac{qA_z}{\hbar} \tag{XII-6}$$

which leads to the system representing the DIRAC equation:

$$\begin{aligned}
 (\eta + \eta_t)\psi_0 + \eta_x\psi_3 - j\eta_y\psi_3 + \eta_z\psi_2 &= j\frac{\partial\psi_0}{\partial x_t} + j\frac{\partial\psi_3}{\partial x} + \frac{\partial\psi_3}{\partial y} + j\frac{\partial\psi_2}{\partial z} \\
 (\eta + \eta_t)\psi_1 + \eta_x\psi_2 + j\eta_y\psi_2 - \eta_z\psi_3 &= j\frac{\partial\psi_1}{\partial x_t} + j\frac{\partial\psi_2}{\partial x} - \frac{\partial\psi_2}{\partial y} - j\frac{\partial\psi_3}{\partial z} \\
 (\eta - \eta_t)\psi_2 - \eta_x\psi_1 + j\eta_y\psi_1 - \eta_z\psi_0 &= -j\frac{\partial\psi_2}{\partial x_t} - j\frac{\partial\psi_1}{\partial x} - \frac{\partial\psi_1}{\partial y} - j\frac{\partial\psi_0}{\partial z} \\
 (\eta - \eta_t)\psi_3 - \eta_x\psi_0 - j\eta_y\psi_0 + \eta_z\psi_1 &= -j\frac{\partial\psi_3}{\partial x_t} - j\frac{\partial\psi_0}{\partial x} + \frac{\partial\psi_0}{\partial y} + j\frac{\partial\psi_1}{\partial z}
 \end{aligned} \tag{XII-7}$$

It appears as a very complicated differential system.

One may attempt to find the wave functions ψ_i ($i = 0, 1, 2, 3$) solutions of this system as it was made in chapter VI, in the form:

$$\begin{aligned} \Psi_i = & \{a_i S_x S_y S_z + b_i S_x C_y S_z + c_i C_x S_y S_z + d_i C_x C_y S_z + e_i S_x S_y C_z + f_i S_x C_y C_z + g_i C_x S_y C_z + h_i C_x C_y C_z\} C_t \\ & + \{i_i S_x S_y S_z + j_i S_x C_y S_z + k_i C_x S_y S_z + l_i C_x C_y S_z + m_i S_x S_y C_z + n_i S_x C_y C_z + o_i C_x S_y C_z + p_i C_x C_y C_z\} S_t \end{aligned} \quad (\text{XII-8})$$

with the usual notation:

$$S_x = \sin(k_x x) \quad S_y = \sin(k_y y) \quad S_z = \sin(k_z z) \quad S_t = \sin(k_t ct) \quad (\text{XII-9})$$

$$C_x = \cos(k_x x) \quad C_y = \cos(k_y y) \quad C_z = \cos(k_z z) \quad C_t = \cos(k_t ct) \quad (\text{XII-10})$$

Using the methodology set out in chapter VI, we are led to build a homogeneous system of 64 equations with 64 unknowns which is described below:

$$\begin{aligned} -jka_0 - jk_x k_3 - k_y j_3 - jk_z m_2 - (\eta + \eta_t) i_0 - (\eta_x - j\eta_y) i_3 - \eta_z i_2 &= 0 \\ -jkb_0 - jk_x l_3 + k_y i_3 - jk_z n_2 - (\eta + \eta_t) j_0 - (\eta_x - j\eta_y) j_3 - \eta_z j_2 &= 0 \\ -jkc_0 + jk_x i_3 - k_y l_3 - jk_z o_2 - (\eta + \eta_t) k_0 - (\eta_x - j\eta_y) k_3 - \eta_z k_2 &= 0 \\ -jkd_0 + jk_x j_3 + k_y k_3 - jk_z p_2 - (\eta + \eta_t) l_0 - (\eta_x - j\eta_y) l_3 - \eta_z l_2 &= 0 \\ -jke_0 - jk_x o_3 - k_y n_3 + jk_z i_2 - (\eta + \eta_t) m_0 - (\eta_x - j\eta_y) m_3 - \eta_z m_2 &= 0 \\ -jkf_0 - jk_x p_3 + k_y m_3 + jk_z j_2 - (\eta + \eta_t) n_0 - (\eta_x - j\eta_y) n_3 - \eta_z n_2 &= 0 \\ -jkg_0 + jk_x m_3 - k_y p_3 + jk_z k_2 - (\eta + \eta_t) o_0 - (\eta_x - j\eta_y) o_3 - \eta_z o_2 &= 0 \\ -jkh_0 + jk_x n_3 + k_y o_3 + jk_z l_2 - (\eta + \eta_t) p_0 - (\eta_x - j\eta_y) p_3 - \eta_z p_2 &= 0 \\ jki_0 - jk_x c_3 - k_y b_3 - jk_z e_2 - (\eta + \eta_t) a_0 - (\eta_x - j\eta_y) a_3 - \eta_z a_2 &= 0 \\ jkj_0 - jk_x d_3 + k_y a_3 - jk_z f_2 - (\eta + \eta_t) b_0 - (\eta_x - j\eta_y) b_3 - \eta_z b_2 &= 0 \\ jkk_0 + jk_x a_3 - k_y d_3 - jk_z g_2 - (\eta + \eta_t) c_0 - (\eta_x - j\eta_y) c_3 - \eta_z c_2 &= 0 \\ jkl_0 + jk_x b_3 + k_y c_3 - jk_z h_2 - (\eta + \eta_t) d_0 - (\eta_x - j\eta_y) d_3 - \eta_z d_2 &= 0 \\ jkm_0 - jk_x g_3 - k_y f_3 + jk_z a_2 - (\eta + \eta_t) e_0 - (\eta_x - j\eta_y) e_3 - \eta_z e_2 &= 0 \\ jkn_0 - jk_x h_3 + k_y e_3 + jk_z b_2 - (\eta + \eta_t) f_0 - (\eta_x - j\eta_y) f_3 - \eta_z f_2 &= 0 \\ jko_0 + jk_x e_3 - k_y h_3 + jk_z c_2 - (\eta + \eta_t) g_0 - (\eta_x - j\eta_y) g_3 - \eta_z g_2 &= 0 \\ jkp_0 + jk_x f_3 + k_y g_3 + jk_z d_2 - (\eta + \eta_t) h_0 - (\eta_x - j\eta_y) h_3 - \eta_z h_2 &= 0 \end{aligned} \quad (\text{XII-11})$$

$$\begin{aligned}
& -jka_1 - jk_x k_2 + k_y j_2 + jk_z m_3 - (\eta + \eta_t) i_1 - (\eta_x + j\eta_y) i_2 + \eta_z i_3 = 0 \\
& -jkb_1 - jk_x l_2 - k_y i_2 + jk_z n_3 - (\eta + \eta_t) j_1 - (\eta_x + j\eta_y) j_2 + \eta_z j_3 = 0 \\
& -jkc_1 + jk_x i_2 + k_y l_2 + jk_z o_3 - (\eta + \eta_t) k_1 - (\eta_x + j\eta_y) k_2 + \eta_z k_3 = 0 \\
& -jkd_1 + jk_x j_2 - k_y k_2 + jk_z p_3 - (\eta + \eta_t) l_1 - (\eta_x + j\eta_y) l_2 + \eta_z l_3 = 0 \\
& -jke_1 - jk_x o_2 + k_y n_2 - jk_z i_3 - (\eta + \eta_t) m_1 - (\eta_x + j\eta_y) m_2 + \eta_z m_3 = 0 \\
& -jkf_1 - jk_x p_2 - k_y m_2 - jk_z j_3 - (\eta + \eta_t) n_1 - (\eta_x + j\eta_y) n_2 + \eta_z n_3 = 0 \\
& -jkg_1 + jk_x m_2 + k_y p_2 - jk_z k_3 - (\eta + \eta_t) o_1 - (\eta_x + j\eta_y) o_2 + \eta_z o_3 = 0 \\
& -jkh_1 + jk_x n_2 - k_y o_2 - jk_z l_3 - (\eta + \eta_t) p_1 - (\eta_x + j\eta_y) p_2 + \eta_z p_3 = 0 \\
& jki_1 - jk_x c_2 + k_y b_2 + jk_z e_3 - (\eta + \eta_t) a_1 - (\eta_x + j\eta_y) a_2 + \eta_z a_3 = 0 \\
& jkj_1 - jk_x d_2 - k_y a_2 + jk_z f_3 - (\eta + \eta_t) b_1 - (\eta_x + j\eta_y) b_2 + \eta_z b_3 = 0 \\
& jkk_1 + jk_x a_2 + k_y d_2 + jk_z g_3 - (\eta + \eta_t) c_1 - (\eta_x + j\eta_y) c_2 + \eta_z c_3 = 0 \\
& jkl_1 + jk_x b_2 - k_y c_2 + jk_z h_3 - (\eta + \eta_t) d_1 - (\eta_x + j\eta_y) d_2 + \eta_z d_3 = 0 \\
& jkm_1 - jk_x g_2 + k_y f_2 - jk_z a_3 - (\eta + \eta_t) e_1 - (\eta_x + j\eta_y) e_2 + \eta_z e_3 = 0 \\
& jkn_1 - jk_x h_2 - k_y e_2 - jk_z b_3 - (\eta + \eta_t) f_1 - (\eta_x + j\eta_y) f_2 + \eta_z f_3 = 0 \\
& jko_1 + jk_x e_2 + k_y h_2 - jk_z c_3 - (\eta + \eta_t) g_1 - (\eta_x + j\eta_y) g_2 + \eta_z g_3 = 0 \\
& jkp_1 + jk_x f_2 - k_y g_2 - jk_z d_3 - (\eta + \eta_t) h_1 - (\eta_x + j\eta_y) h_2 + \eta_z h_3 = 0
\end{aligned} \tag{XII-12}$$

$$\begin{aligned}
& jka_2 + jk_x k_1 + k_y j_1 + jk_z m_0 - (\eta - \eta_t) i_2 + (\eta_x - j\eta_y) i_1 + \eta_z i_0 = 0 \\
& jkb_2 + jk_x l_1 - k_y i_1 + jk_z n_0 - (\eta - \eta_t) j_2 + (\eta_x - j\eta_y) j_1 + \eta_z j_0 = 0 \\
& jkc_2 - jk_x i_1 + k_y l_1 + jk_z o_0 - (\eta - \eta_t) k_2 + (\eta_x - j\eta_y) k_1 + \eta_z k_0 = 0 \\
& jkd_2 - jk_x j_1 - k_y k_1 + jk_z p_0 - (\eta - \eta_t) l_2 + (\eta_x - j\eta_y) l_1 + \eta_z l_0 = 0 \\
& jke_2 + jk_x o_1 + k_y n_1 - jk_z i_0 - (\eta - \eta_t) m_2 + (\eta_x - j\eta_y) m_1 + \eta_z m_0 = 0 \\
& jkf_2 + jk_x p_1 - k_y m_1 - jk_z j_0 - (\eta - \eta_t) n_2 + (\eta_x - j\eta_y) n_1 + \eta_z n_0 = 0 \\
& jkg_2 - jk_x m_1 + k_y p_1 - jk_z k_0 - (\eta - \eta_t) o_2 + (\eta_x - j\eta_y) o_1 + \eta_z o_0 = 0 \\
& jkh_2 - jk_x n_1 - k_y o_1 - jk_z l_0 - (\eta - \eta_t) p_2 + (\eta_x - j\eta_y) p_1 + \eta_z p_0 = 0 \\
& -jki_2 + jk_x c_1 + k_y b_1 + jk_z e_0 - (\eta - \eta_t) a_2 + (\eta_x - j\eta_y) a_1 + \eta_z a_0 = 0 \\
& -jkj_2 + jk_x d_1 - k_y a_1 + jk_z f_0 - (\eta - \eta_t) b_2 + (\eta_x - j\eta_y) b_1 + \eta_z b_0 = 0 \\
& -jkk_2 - jk_x a_1 + k_y d_1 + jk_z g_0 - (\eta - \eta_t) c_2 + (\eta_x - j\eta_y) c_1 + \eta_z c_0 = 0 \\
& -jkl_2 - jk_x b_1 - k_y c_1 + jk_z h_0 - (\eta - \eta_t) d_2 + (\eta_x - j\eta_y) d_1 + \eta_z d_0 = 0 \\
& -jkm_2 + jk_x g_1 + k_y f_1 - jk_z a_0 - (\eta - \eta_t) e_2 + (\eta_x - j\eta_y) e_1 + \eta_z e_0 = 0 \\
& -jkn_2 + jk_x h_1 - k_y e_1 - jk_z b_0 - (\eta - \eta_t) f_2 + (\eta_x - j\eta_y) f_1 + \eta_z f_0 = 0 \\
& -jko_2 - jk_x e_1 + k_y h_1 - jk_z c_0 - (\eta - \eta_t) g_2 + (\eta_x - j\eta_y) g_1 + \eta_z g_0 = 0 \\
& -jkp_2 - jk_x f_1 - k_y g_1 - jk_z d_0 - (\eta - \eta_t) h_2 + (\eta_x - j\eta_y) h_1 + \eta_z h_0 = 0
\end{aligned} \tag{XII-13}$$

$$\begin{aligned}
& jka_3 + jk_x k_0 - k_y j_0 - jk_z m_1 - (\eta - \eta_t) i_3 + (\eta_x + j\eta_y) i_0 - \eta_z i_1 = 0 \\
& jkb_3 + jk_x l_0 + k_y i_0 - jk_z n_1 - (\eta - \eta_t) j_3 + (\eta_x + j\eta_y) j_0 - \eta_z j_1 = 0 \\
& jkc_3 - jk_x i_0 - k_y l_0 - jk_z o_1 - (\eta - \eta_t) k_3 + (\eta_x + j\eta_y) k_0 - \eta_z k_1 = 0 \\
& jkd_3 - jk_x j_0 + k_y k_0 - jk_z p_1 - (\eta - \eta_t) l_3 + (\eta_x + j\eta_y) l_0 - \eta_z l_1 = 0 \\
& jke_3 + jk_x o_0 - k_y n_0 + jk_z i_1 - (\eta - \eta_t) m_3 + (\eta_x + j\eta_y) m_0 - \eta_z m_1 = 0 \\
& jkf_3 + jk_x p_0 + k_y m_0 + jk_z j_1 - (\eta - \eta_t) n_3 + (\eta_x + j\eta_y) n_0 - \eta_z n_1 = 0 \\
& jkg_3 - jk_x m_0 - k_y p_0 + jk_z k_1 - (\eta - \eta_t) o_3 + (\eta_x + j\eta_y) o_0 - \eta_z o_1 = 0 \\
& jkh_3 - jk_x n_0 + k_y o_0 + jk_z l_1 - (\eta - \eta_t) p_3 + (\eta_x + j\eta_y) p_0 - \eta_z p_1 = 0 \\
& - jki_3 + jk_x c_0 - k_y b_0 - jk_z e_1 - (\eta - \eta_t) a_3 + (\eta_x + j\eta_y) a_0 - \eta_z a_1 = 0 \\
& - jkj_3 + jk_x d_0 + k_y a_0 - jk_z f_1 - (\eta - \eta_t) b_3 + (\eta_x + j\eta_y) b_0 - \eta_z b_1 = 0 \\
& - jkk_3 - jk_x a_0 - k_y d_0 - jk_z g_1 - (\eta - \eta_t) c_3 + (\eta_x + j\eta_y) c_0 - \eta_z c_1 = 0 \\
& - jkl_3 - jk_x b_0 + k_y c_0 - jk_z h_1 - (\eta - \eta_t) d_3 + (\eta_x + j\eta_y) d_0 - \eta_z d_1 = 0 \\
& - jkm_3 + jk_x g_0 - k_y f_0 + jk_z a_1 - (\eta - \eta_t) e_3 + (\eta_x + j\eta_y) e_0 - \eta_z e_1 = 0 \\
& - jkn_3 + jk_x h_0 + k_y e_0 + jk_z b_1 - (\eta - \eta_t) f_3 + (\eta_x + j\eta_y) f_0 - \eta_z f_1 = 0 \\
& - jko_3 - jk_x e_0 - k_y h_0 + jk_z c_1 - (\eta - \eta_t) g_3 + (\eta_x + j\eta_y) g_0 - \eta_z g_1 = 0 \\
& - jkp_3 - jk_x f_0 + k_y g_0 + jk_z d_1 - (\eta - \eta_t) h_3 + (\eta_x + j\eta_y) h_0 - \eta_z h_1 = 0
\end{aligned}$$

(XII-14)

The first idea is to build on the results of Chapter VI, it means to assign a mode to ψ_0 , and search patterns that are solutions of the system for wave functions ψ_1, ψ_2, ψ_3 . A long and tedious mathematical work has not allowed to express an exact solution for ψ_1, ψ_2, ψ_3 . This work has not led to show that such a solution does not exist, and so the problem remains open. Complementary indications will be given in chapter XXII.

Progress towards a possible solution will come from a physical analysis of this system.

In first place, the determinant of this system should be null. We hypothesize that the condition of nullity is provided by the equation of energy conservation.

We must therefore establish this new equation of conservation on the basis of the energy provided to the charged particle by the presence of the electromagnetic potential. This energy is of two kinds:

- the energy provided by the scalar potential ϕ allows to increase the kinetic energy of the charged particle. Relativity suggests that this increase in energy is transformed in mass energy.
- the energy provided by the vector potential (A_x, A_y, A_z) allows only to change the direction of the trajectory of the particle, without kinetic energy supply: this is therefore a purely impulse energy.

Outside the presence of the electromagnetic field, the equation of conservation of energy is written:

$$-k_t^2 + k_x^2 + k_y^2 + k_z^2 + \eta^2 = 0 \quad (\text{XII-15})$$

Inputs of energy of the electromagnetic field in this relationship can be introduced in many ways without being trivial to make a priori choice among all possible formulations.

Reflection led to conclude that the correct form is as follows:

$$-(k_t + \eta_t)^2 + (k_x + \eta_x)^2 + (k_y + \eta_y)^2 + (k_z + \eta_z)^2 + \eta^2 = 0 \quad (\text{XII-16})$$

It should be noted in particular that the kinetic energy η_t provided by the scalar potential ϕ is not associated with the mass energy η as it could think intuitively.

The rigorous justification is that the sum of the wave four-vector and the electromagnetic potential four-vector gives a four-vector too:

$$\begin{pmatrix} k_x \\ k_y \\ k_z \\ k_t \end{pmatrix} + \begin{pmatrix} \eta_x \\ \eta_y \\ \eta_z \\ \eta_t \end{pmatrix} = \begin{pmatrix} k_x + \eta_x \\ k_y + \eta_y \\ k_z + \eta_z \\ k_t + \eta_t \end{pmatrix} \quad (\text{XII-17})$$

And since the pseudo-norm of the four-vector sum is constant, this leads directly to the equation of conservation of energy (XII-16).

Therefore, we hypothesize that the determinant of the overall system of 64 equations with 64 unknowns is zero when the equation of conservation of energy (XII-16) is checked.

But, it does not provide track to move towards a solution of the generalized electromagnetic interacting DIRAC equation system.

A detailed analysis of the system of DIRAC allows however to see if one excite a wave function with a mode in time quadrature for the scalar potential and and in space quadrature for the vector potential, it manages to get some solutions relative to the stationary modes (XII-11,12,13,14). Quadrature modes must also express exchanges of energy between positive and negative energy which introduced so the quantity $j = \text{sqrt}(-1)$ between expressions of these modes.

In summary, the presence of the electromagnetic field don't excite wave functions in the form of independent stationary modes, but in the form of combinations of modes related in time and space quadrature, and they reflect exchanges of energy between positive energy and negative energy.

These modes will therefore present themselves in the form of combinations of functions of x, y, z, t with amplitude A , which have the following expression:

$$A \exp(\pm j k_x x) \exp(\pm j k_y y) \exp(\pm k_z z) \exp(\pm j k_t x_t) = A \exp \left\{ j (\pm k_t x_t \pm k_x x \pm k_y y \pm k_z z) \right\} \quad (\text{XII-18})$$

The signs + and - present in this expression will affect the relationship of conservation of energy which should take them into account in the form:

$$-(\eta_t \pm k_t)^2 + (\eta_x \pm k_x)^2 + (\eta_y \pm k_y)^2 + (\eta_z \pm k_z)^2 + \eta^2 = 0 \quad (\text{XII-19})$$

These considerations are illustrated on this particular example, which corresponds to an exact solution of the DIRAC equation generalized to electromagnetic interacting.

$$\begin{aligned} \psi_0 &= -(k_z + \eta_z) \exp \{j(-k_t x_t + k_x x + k_y y + k_z z)\} \\ \psi_1 &= -(k_x + \eta_x) \exp \{j(-k_t x_t + k_x x + k_y y + k_z z)\} - j(k_y + \eta_y) \exp \{j(-k_t x_t + k_x x + k_y y + k_z z)\} \\ \psi_2 &= (\eta - k_t + \eta_t) \exp \{j(-k_t x_t + k_x x + k_y y + k_z z)\} \\ \psi_3 &= 0 \end{aligned} \quad (\text{XII-20})$$

This solution is associated with the conservation of energy equation:

$$-(\eta_t - k_t)^2 + (\eta_x + k_x)^2 + (\eta_y + k_y)^2 + (\eta_z + k_z)^2 + \eta^2 = 0 \quad (\text{XII-21})$$

It is suitable to be convincing, to detail checks of these properties. We skip in the calculations below exponential coming in factor with all terms.

First generalized DIRAC equation:

$$(\eta + \eta_t) \psi_0 + \eta_x \psi_3 - j \eta_y \psi_3 + \eta_z \psi_2 = j \frac{\partial \psi_0}{\partial x_t} + j \frac{\partial \psi_3}{\partial x} + \frac{\partial \psi_3}{\partial y} + j \frac{\partial \psi_2}{\partial z} \quad (\text{XII-22})$$

$$-(\eta + \eta_t)(k_z + \eta_z) + 0 + 0 + \eta_z(\eta - k_t + \eta_t) = -(k_z + \eta_z)k_t + 0 + 0 - (\eta - k_t + \eta_t)k_z \quad (\text{XII-23})$$

Second generalized DIRAC equation:

$$(\eta + \eta_t) \psi_1 + \eta_x \psi_2 + j \eta_y \psi_2 - \eta_z \psi_3 = j \frac{\partial \psi_1}{\partial x_t} + j \frac{\partial \psi_2}{\partial x} - \frac{\partial \psi_2}{\partial y} - j \frac{\partial \psi_3}{\partial z} \quad (\text{XII-24})$$

$$\begin{aligned} (\eta + \eta_t) \{-(k_x + \eta_x) - j(k_y + \eta_y)\} + \eta_x(\eta - k_t + \eta_t) + j \eta_y(\eta - k_t + \eta_t) - 0 = \\ k_t \{-(k_x + \eta_x) - j(k_y + \eta_y)\} - k_x(\eta - k_t + \eta_t) - j k_y(\eta - k_t + \eta_t) - 0 \end{aligned} \quad (\text{XII-25})$$

Third generalized DIRAC equation:

$$(\eta - \eta_t) \psi_2 - \eta_x \psi_1 + j \eta_y \psi_1 - \eta_z \psi_0 = -j \frac{\partial \psi_2}{\partial x_t} - j \frac{\partial \psi_1}{\partial x} - \frac{\partial \psi_1}{\partial y} - j \frac{\partial \psi_0}{\partial z} \quad (\text{XII-26})$$

$$\begin{aligned} (\eta - \eta_t)(\eta - k_t + \eta_t) + \eta_x \{(k_x + \eta_x) + j(k_y + \eta_y)\} - j \eta_y \{(k_x + \eta_x) + j(k_y + \eta_y)\} + \eta_z(k_z + \eta_z) = \\ -(\eta - k_t + \eta_t)k_t - k_x \{(k_x + \eta_x) + j(k_y + \eta_y)\} + j k_y \{(k_x + \eta_x) + j(k_y + \eta_y)\} - k_z(k_z + \eta_z) \end{aligned} \quad (\text{XII-27})$$

The imaginary terms vanish, and it remains:

$$-(\eta_t - k_t)^2 + (\eta_x + k_x)^2 + (\eta_y + k_y)^2 + (\eta_z + k_z)^2 + \eta^2 = 0 \quad (\text{XII-28})$$

or so the equation of conservation of energy.

Fourth generalized DIRAC equation:

$$(\eta - \eta_t)\psi_3 - \eta_x\psi_0 - j\eta_y\psi_0 + \eta_z\psi_1 = -j\frac{\partial\psi_3}{\partial x_t} - j\frac{\partial\psi_0}{\partial x} + \frac{\partial\psi_0}{\partial y} + j\frac{\partial\psi_1}{\partial z} \quad (\text{XII-29})$$

$$\begin{aligned} 0 + \eta_x(k_z + \eta_z) + j\eta_y(k_z + \eta_z) - \eta_z\{(k_x + \eta_x) + j(k_y + \eta_y)\} = \\ 0 - k_x(k_z + \eta_z) - jk_y(k_z + \eta_z) + k_z\{(k_x + \eta_x) + j(k_y + \eta_y)\} \end{aligned} \quad (\text{XII-30})$$

Verification that the (XII-20) solution is an exact solution to the generalized DIRAC system is completed.

We would think that on the basis of the exact solution (XII-20), it is possible to construct purely real stationary solutions for example by summing two solutions in $\exp(j\omega t)$ and $\exp(-j\omega t)$, which would be in contradiction with the previous statement on the impossibility to obtain such solutions in the generalized DIRAC system.

In fact, it is impossible to sum these solutions, because they are relative each to a different energy conservation equation, and thus to a different condition of nullity of the determinant of the system.

One can however mix real stationary modes and complex stationary modes, as shown in the exact solution below, in which the presence of the electromagnetic potential is reduced to the scalar potential:

$$\begin{aligned} \psi_0 &= 0 \\ \psi_1 &= -(\eta - k_t - \eta_t)\sin(k_x x)\cos(k_y y)\cos(k_z z)\exp(jk_t x_t) \\ \psi_2 &= -k_y \sin(k_x x)\sin(k_y y)\cos(k_z z)\exp(jk_t x_t) + jk_x \cos(k_x x)\cos(k_y y)\cos(k_z z)\exp(jk_t x_t) \\ \psi_3 &= jk_z \sin(k_x x)\cos(k_y y)\sin(k_z z)\exp(jk_t x_t) \end{aligned} \quad (\text{XII-31})$$

It is associated with the conservation of energy equation:

$$(k_t + \eta_t)^2 - k_x^2 - k_y^2 - k_z^2 - \eta^2 = 0 \quad (\text{XII-32})$$

XIII

The currents of DIRAC (2)

DIRAC currents related to the solutions of the generalized equation to the presence of an electromagnetic field have a remarkable property which will be illustrated on the example of solution proposed earlier and recalled to memory:

$$\begin{aligned}\psi_0 &= -(k_z + \eta_z) \exp \{j(-k_t x_t + k_x x + k_y y + k_z z)\} \\ \psi_1 &= -(k_x + \eta_x) \exp \{j(-k_t x_t + k_x x + k_y y + k_z z)\} - j(k_y + \eta_y) \exp \{j(-k_t x_t + k_x x + k_y y + k_z z)\} \\ \psi_2 &= (\eta - k_t + \eta_t) \exp \{j(-k_t x_t + k_x x + k_y y + k_z z)\} \\ \psi_3 &= 0\end{aligned}\tag{XIII-1}$$

The expression of these currents is recalled below. Multiplied by a constant adequate, they become homogeneous to a volumetric energy density, and the term J^0 represents the total energy volume density.

$$\begin{aligned}J^0 &= \psi_0^* \psi_0 + \psi_1^* \psi_1 + \psi_2^* \psi_2 + \psi_3^* \psi_3 \\ J^1 &= \psi_3^* \psi_0 + \psi_2^* \psi_1 + \psi_1^* \psi_2 + \psi_0^* \psi_3 \\ J^2 &= j\psi_3^* \psi_0 - j\psi_2^* \psi_1 + j\psi_1^* \psi_2 - j\psi_0^* \psi_3 \\ J^3 &= \psi_2^* \psi_0 - \psi_3^* \psi_1 + \psi_0^* \psi_2 - \psi_1^* \psi_3\end{aligned}\tag{XIII-2}$$

In addition, these currents must check the local conservation of energy equation:

$$\frac{\partial J^0}{\partial x_t} + \frac{\partial J^1}{\partial x} + \frac{\partial J^2}{\partial y} + \frac{\partial J^3}{\partial z} = 0\tag{XIII-3}$$

One obtains, in a straightforward way, from wave functions (XIII-1) and currents (XIII-2), the following expressions:

$$\begin{aligned}J^0 &= (k_x + \eta_x)^2 + (k_z + \eta_z)^2 + (k_z + \eta_z)^2 + (\eta - k_t + \eta_t)^2 \\ J^1 &= -2(\eta - k_t + \eta_t)(k_x + \eta_x) \\ J^2 &= -2(\eta - k_t + \eta_t)(k_y + \eta_y) \\ J^3 &= -2(\eta - k_t + \eta_t)(k_z + \eta_z)\end{aligned}\tag{XIII-4}$$

The surprising result that appears in the expression of these currents is that they depend on neither time nor space. In other words, the volume density of total energy represented by J^0 is uniform inside the block that contains the energy of the particle.

This result presents a greater analogy to the example of a particle which has twice its mass energy which is presented in chapter VI, with however a difference subject to the assumptions made in the preparation of the solutions in the presence of electromagnetic field.

If we accept that there is no purely real solutions for the wave functions obtained in the presence of an electromagnetic field, this particular scheme of uniform energy within the particle density is imposed by the presence of the electromagnetic field, while nothing requires it for the free particle.

The normalization constant C is obtained by writing that J^0 represents the volume density of total energy of the particle placed in an electromagnetic field, and then the total energy contained in volume V that delimits the particle:

$$E = (\hbar^2 c^2) (\mathbf{k}_t + \boldsymbol{\eta}_t)^2 V \quad (\text{XIII-5})$$

In equating this relationship with (XIII-4) multiplied by the normalizing constant C^2 , one obtains:

$$E = C^2 \left\{ (\mathbf{k}_x + \boldsymbol{\eta}_x)^2 + (\mathbf{k}_y + \boldsymbol{\eta}_y)^2 + (\mathbf{k}_z + \boldsymbol{\eta}_z)^2 + (\eta - k_t + \eta_t)^2 \right\} = (\hbar^2 c^2) (\mathbf{k}_t + \boldsymbol{\eta}_t)^2 V \quad (\text{XIII-6})$$

or still, by introducing the equation of conservation of energy called for memory:

$$(\mathbf{k}_t + \boldsymbol{\eta}_t)^2 = (\mathbf{k}_x + \boldsymbol{\eta}_x)^2 + (\mathbf{k}_y + \boldsymbol{\eta}_y)^2 + (\mathbf{k}_z + \boldsymbol{\eta}_z)^2 + \eta^2 \quad (\text{XIII-7})$$

$$C^2 = \frac{(\hbar^2 c^2) (\mathbf{k}_t + \boldsymbol{\eta}_t)^2 V}{(\mathbf{k}_t + \boldsymbol{\eta}_t)^2 - \eta^2 + (\eta - k_t + \eta_t)^2} \quad (\text{XIII-8})$$

XIV

Conclusion of the second part

We know that solutions to the DIRAC equation which correspond to the reality of the observations are solutions developed using spherical modes and in this context, one may wonder what is the interest to work on solutions obtained in Cartesian coordinates.

It appears that Cartesian solutions are more quickly attainable, and that they can learn to us valuable information about the behavior of energy which the particle is constituted.

On the basis of an equation of conservation of energy which is a fundamental physical reasoning and on the assumption of an internal evolution of energy based on exchanges between stationary modes, the exact solutions to the DIRAC equation deliver new elements likely to describe the physics of the infinitely small.

There is no assumption on the spatial extent of the modes that are supposed to describe the behavior of the particle, but it is legitimate to think that this scope exceeds the size given in classical physics if you want for example to be able to explain the phenomena of interference. This interpretation was already present in the thought of Louis DE BROGLIE during his thesis: "Do we assume the localized periodic phenomenon *inside* the piece of energy? This is not necessary and will result in paragraph (III) it is probably spread in a large region of space."

In opposition to the Copenhagen school, the energy interpretation of stationary solutions is perfectly deterministic, but it does not contradict the experiences of the probabilistic vision of quantum mechanics.

It justifies the wave particle duality in indicating in what manner the internal energy to the particle alternately goes in the form of mass energy and wave energy.

It shows how the HEISENBERG uncertainty principle is interpreted by indicating how the mass energy and impulse energy are not simultaneously present in the same place.

Finally, it is fully compatible with the interpretation of a wave phenomenon associated with the particle following the DE BROGLIE wavelength.

Third part

Energy approach of Dirac equation and its exact solutions in spherical coordinates

XV

DIRAC equation in spherical coordinates

The analysis of stationary solutions of the DIRAC equation in Cartesian coordinates allowed to highlight the properties which, while being backed by a perfectly deterministic theory, are in agreement with all the results obtained in the statistical interpretation of the Copenhagen school.

If stationary modes are able to represent the exchange of energy within particles, there are little chance that it is in the shape of a parallelepiped. Everything indicates, in particular solutions of the SCHRÖDINGER equation, that the coordinate system the most suitable, one that provides solutions in agreement with experimental observations, is the system of spherical coordinates (Figure XV-1)

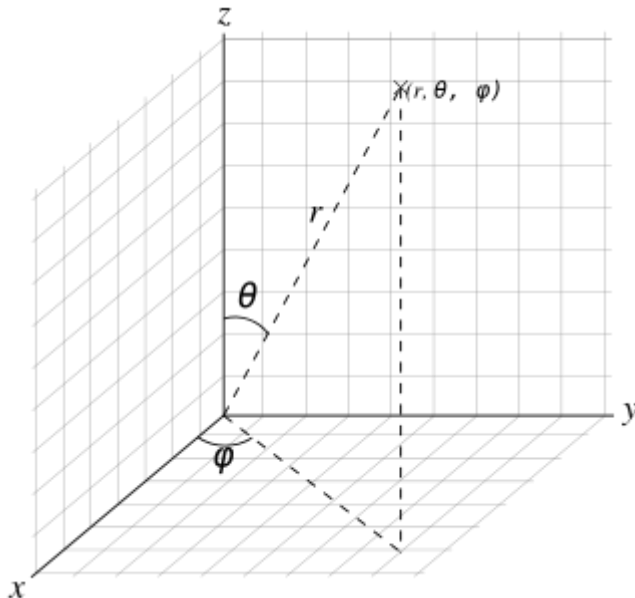


Figure XV-1: *representation of the spherical coordinate system*

It is expected to appear in exact solutions to the DIRAC equation in spherical coordinates, informations that allow to better understand how the spin of the electron is related to internal rotation of energy.

Even before discussing the search for solutions, we must transform the DIRAC equation in spherical coordinates.

The starting point is given by the link relations between the cartesian and spherical coordinates:

$$\begin{aligned} x &= r \sin \theta \cos \varphi \\ y &= r \sin \theta \sin \varphi \\ z &= r \cos \theta \end{aligned} \tag{XV-1}$$

We deduce the differential relations:

$$\begin{aligned} dx &= dr \sin \theta \cos \varphi + r \cos \theta \cos \varphi d\theta - r \sin \theta \sin \varphi d\varphi \\ dy &= dr \sin \theta \sin \varphi + r \cos \theta \sin \varphi d\theta + r \sin \theta \cos \varphi d\varphi \\ dz &= dr \cos \theta - r \sin \theta d\theta \end{aligned} \tag{XV-2}$$

Or again using matrix writing:

$$\begin{pmatrix} dx \\ dy \\ dz \end{pmatrix} = \begin{pmatrix} \sin \theta \cos \varphi & r \cos \theta \cos \varphi & -r \sin \theta \sin \varphi \\ \sin \theta \sin \varphi & r \cos \theta \sin \varphi & r \sin \theta \cos \varphi \\ \cos \theta & -r \sin \theta & 0 \end{pmatrix} \begin{pmatrix} dr \\ d\theta \\ d\varphi \end{pmatrix} \tag{XV-3}$$

And then, by matrix inversion:

$$\begin{pmatrix} dr \\ d\theta \\ d\varphi \end{pmatrix} = \begin{pmatrix} \frac{\sin \theta \cos \varphi}{\cos \theta \cos \varphi} & \frac{\sin \theta \sin \varphi}{\cos \theta \sin \varphi} & \frac{\cos \theta}{-\sin \theta} \\ \frac{r}{-r \sin \varphi} & \frac{r}{\cos \varphi} & \frac{r}{0} \\ \frac{r}{r \sin \theta} & \frac{r}{r \sin \theta} & 0 \end{pmatrix} \begin{pmatrix} dx \\ dy \\ dz \end{pmatrix} \tag{XV-4}$$

The DIRAC equation in Cartesian coordinates is recalled below:

$$j \left[\gamma_0 \left(\frac{\partial}{\partial x_t} \right) + \gamma_1 \left(\frac{\partial}{\partial x} \right) + \gamma_2 \left(\frac{\partial}{\partial y} \right) + \gamma_3 \left(\frac{\partial}{\partial z} \right) \right] (\psi) = \eta (\psi) \tag{XV-5}$$

where matrices γ_i are of the form:

$$\gamma_0 = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix} \gamma_1 = \begin{pmatrix} 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \\ 0 & -1 & 0 & 0 \\ -1 & 0 & 0 & 0 \end{pmatrix} \gamma_2 = \begin{pmatrix} 0 & 0 & 0 & -j \\ 0 & 0 & j & 0 \\ 0 & j & 0 & 0 \\ -j & 0 & 0 & 0 \end{pmatrix} \gamma_3 = \begin{pmatrix} 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & -1 \\ -1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \end{pmatrix} \tag{XV-6}$$

The objective is to transform this equation to write it with the partial derivatives with respect to spherical variables r, θ, φ :

$$j \left[\gamma_0 \left(\frac{\partial}{\partial x_i} \right) + \gamma_r \left(\frac{\partial}{\partial r} \right) + \gamma_\theta \frac{1}{r} \left(\frac{\partial}{\partial \theta} \right) + \gamma_\varphi \frac{1}{r \sin \theta} \left(\frac{\partial}{\partial \varphi} \right) \right] (\psi) = \eta(\psi) \quad (\text{XV-7})$$

The issue is the determination of the new matrices $\gamma_r, \gamma_\theta, \gamma_\varphi$. This requires, as a first step, to formalize the passage of the partial derivatives with respect to x, y, z into partial derivatives with respect to r, θ, φ .

To establish these relationships, we can use the total differential, which is a constant of the transformation:

$$d\psi = \frac{\partial \psi}{\partial r} dr + \frac{\partial \psi}{\partial \theta} d\theta + \frac{\partial \psi}{\partial \varphi} d\varphi \quad (\text{XV-8})$$

By introducing in this relation the differentials $dr, d\theta, d\varphi$ given in (XV-4), we obtain:

$$\begin{aligned} d\psi &= \frac{\partial \psi}{\partial r} (\sin \theta \cos \varphi dx + \sin \theta \sin \varphi dy + \cos \theta dz) + \\ &\frac{\partial \psi}{\partial \theta} \frac{1}{r} (\cos \theta \cos \varphi dx + \cos \theta \sin \varphi dy - \sin \theta dz) \\ &+ \frac{\partial \psi}{\partial \varphi} \frac{1}{r \sin \theta} (-\sin \varphi dx + \cos \varphi dy) \end{aligned} \quad (\text{XV-9})$$

We gather terms which are linked to Cartesian differentials dx, dy, dz :

$$\begin{aligned} d\psi &= dx \left\{ \frac{\partial \psi}{\partial r} \sin \theta \cos \varphi + \frac{\partial \psi}{\partial \theta} \frac{1}{r} \cos \theta \cos \varphi - \frac{\partial \psi}{\partial \varphi} \frac{1}{r \sin \theta} \sin \varphi \right\} \\ &+ dy \left\{ \frac{\partial \psi}{\partial r} \sin \theta \sin \varphi + \frac{\partial \psi}{\partial \theta} \frac{1}{r} \cos \theta \sin \varphi + \frac{\partial \psi}{\partial \varphi} \frac{1}{r \sin \theta} \cos \varphi \right\} \\ &+ dz \left\{ \frac{\partial \psi}{\partial r} \cos \theta - \frac{\partial \psi}{\partial \theta} \frac{1}{r} \sin \theta \right\} \end{aligned} \quad (\text{XV-10})$$

We then identify this expression with the total differential in Cartesian coordinates, which is a constant of the transformation:

$$d\psi = \frac{\partial \psi}{\partial x} dx + \frac{\partial \psi}{\partial y} dy + \frac{\partial \psi}{\partial z} dz \quad (\text{XV-11})$$

This identification provides the searched transformation between the partial derivatives in spherical coordinates and the partial derivatives in Cartesian coordinates:

$$\begin{aligned}
 \frac{\partial \psi}{\partial x} &= \frac{\partial \psi}{\partial r} \sin \theta \cos \varphi + \frac{\partial \psi}{\partial \theta} \frac{1}{r} \cos \theta \cos \varphi - \frac{\partial \psi}{\partial \varphi} \frac{1}{r \sin \theta} \sin \varphi \\
 \frac{\partial \psi}{\partial y} &= \frac{\partial \psi}{\partial r} \sin \theta \sin \varphi + \frac{\partial \psi}{\partial \theta} \frac{1}{r} \cos \theta \sin \varphi + \frac{\partial \psi}{\partial \varphi} \frac{1}{r \sin \theta} \cos \varphi \\
 \frac{\partial \psi}{\partial z} &= \frac{\partial \psi}{\partial r} \cos \theta - \frac{\partial \psi}{\partial \theta} \frac{1}{r} \sin \theta
 \end{aligned}
 \tag{XV-12}$$

It is then possible to change the coordinate system by reporting these equalities in DIRAC equation in Cartesian coordinates:

$$\mathbf{j} \left[\gamma_0 \left(\frac{\partial}{\partial x_t} \right) + \gamma_1 \left(\frac{\partial}{\partial x} \right) + \gamma_2 \left(\frac{\partial}{\partial y} \right) + \gamma_3 \left(\frac{\partial}{\partial z} \right) \right] (\psi) = \frac{m_0 c}{\hbar} (\psi)
 \tag{XV-13}$$

The explicit formulation is fully developed below:

$$\begin{aligned}
 \eta \begin{pmatrix} \psi_0 \\ \psi_1 \\ \psi_2 \\ \psi_3 \end{pmatrix} &= \mathbf{j} \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix} \begin{pmatrix} \frac{\partial \psi_0}{\partial x_t} \\ \frac{\partial \psi_1}{\partial x_t} \\ \frac{\partial \psi_2}{\partial x_t} \\ \frac{\partial \psi_3}{\partial x_t} \end{pmatrix} \\
 + \mathbf{j} \begin{pmatrix} 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \\ 0 & -1 & 0 & 0 \\ -1 & 0 & 0 & 0 \end{pmatrix} &\begin{pmatrix} \frac{\partial \psi_0}{\partial r} \sin \theta \cos \varphi + \frac{\partial \psi_0}{\partial \theta} \frac{1}{r} \cos \theta \cos \varphi - \frac{\partial \psi_0}{\partial \varphi} \frac{1}{r \sin \theta} \sin \varphi \\ \frac{\partial \psi_1}{\partial r} \sin \theta \cos \varphi + \frac{\partial \psi_1}{\partial \theta} \frac{1}{r} \cos \theta \cos \varphi - \frac{\partial \psi_1}{\partial \varphi} \frac{1}{r \sin \theta} \sin \varphi \\ \frac{\partial \psi_2}{\partial r} \sin \theta \cos \varphi + \frac{\partial \psi_2}{\partial \theta} \frac{1}{r} \cos \theta \cos \varphi - \frac{\partial \psi_2}{\partial \varphi} \frac{1}{r \sin \theta} \sin \varphi \\ \frac{\partial \psi_3}{\partial r} \sin \theta \cos \varphi + \frac{\partial \psi_3}{\partial \theta} \frac{1}{r} \cos \theta \cos \varphi - \frac{\partial \psi_3}{\partial \varphi} \frac{1}{r \sin \theta} \sin \varphi \end{pmatrix} \\
 + \mathbf{j} \begin{pmatrix} 0 & 0 & 0 & -\mathbf{j} \\ 0 & 0 & \mathbf{j} & 0 \\ 0 & \mathbf{j} & 0 & 0 \\ -\mathbf{j} & 0 & 0 & 0 \end{pmatrix} &\begin{pmatrix} \frac{\partial \psi_0}{\partial r} \sin \theta \sin \varphi + \frac{\partial \psi_0}{\partial \theta} \frac{1}{r} \cos \theta \sin \varphi + \frac{\partial \psi_0}{\partial \varphi} \frac{1}{r \sin \theta} \cos \varphi \\ \frac{\partial \psi_1}{\partial r} \sin \theta \sin \varphi + \frac{\partial \psi_1}{\partial \theta} \frac{1}{r} \cos \theta \sin \varphi + \frac{\partial \psi_1}{\partial \varphi} \frac{1}{r \sin \theta} \cos \varphi \\ \frac{\partial \psi_2}{\partial r} \sin \theta \sin \varphi + \frac{\partial \psi_2}{\partial \theta} \frac{1}{r} \cos \theta \sin \varphi + \frac{\partial \psi_2}{\partial \varphi} \frac{1}{r \sin \theta} \cos \varphi \\ \frac{\partial \psi_3}{\partial r} \sin \theta \sin \varphi + \frac{\partial \psi_3}{\partial \theta} \frac{1}{r} \cos \theta \sin \varphi + \frac{\partial \psi_3}{\partial \varphi} \frac{1}{r \sin \theta} \cos \varphi \end{pmatrix} \\
 + \mathbf{j} \begin{pmatrix} 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & -1 \\ -1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \end{pmatrix} &\begin{pmatrix} \frac{\partial \psi_0}{\partial r} \cos \theta - \frac{\partial \psi_0}{\partial \theta} \frac{1}{r} \sin \theta \\ \frac{\partial \psi_1}{\partial r} \cos \theta - \frac{\partial \psi_1}{\partial \theta} \frac{1}{r} \sin \theta \\ \frac{\partial \psi_2}{\partial r} \cos \theta - \frac{\partial \psi_2}{\partial \theta} \frac{1}{r} \sin \theta \\ \frac{\partial \psi_3}{\partial r} \cos \theta - \frac{\partial \psi_3}{\partial \theta} \frac{1}{r} \sin \theta \end{pmatrix}
 \end{aligned} \tag{XV-14}$$

We organize terms around the partial derivative in r , θ , φ , in order to obtain a matrix relation:

$$\mathbf{j} \left[\gamma_0 \left(\frac{\partial}{\partial x_t} \right) + \gamma_r \left(\frac{\partial}{\partial r} \right) + \gamma_\theta \frac{1}{r} \left(\frac{\partial}{\partial \theta} \right) + \gamma_\varphi \frac{1}{r \sin \theta} \left(\frac{\partial}{\partial \varphi} \right) \right] (\psi) = \eta(\psi) \tag{XV-15}$$

This grouping led by identification to the searched matrices γ_r , γ_θ , γ_φ :

$$\begin{aligned}
 \gamma_0 &= \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix} & \gamma_r &= \begin{pmatrix} 0 & 0 & \cos \theta & \sin \theta e^{-j\varphi} \\ 0 & 0 & \sin \theta e^{j\varphi} & -\cos \theta \\ -\cos \theta & -\sin \theta e^{-j\varphi} & 0 & 0 \\ -\sin \theta e^{j\varphi} & \cos \theta & 0 & 0 \end{pmatrix} \\
 \gamma_\theta &= \begin{pmatrix} 0 & 0 & -\sin \theta & \cos \theta e^{-j\varphi} \\ 0 & 0 & \cos \theta e^{j\varphi} & \sin \theta \\ \sin \theta & -\cos \theta e^{-j\varphi} & 0 & 0 \\ -\cos \theta e^{j\varphi} & -\sin \theta & 0 & 0 \end{pmatrix} & \gamma_\varphi &= \begin{pmatrix} 0 & 0 & 0 & -j e^{-j\varphi} \\ 0 & 0 & j e^{j\varphi} & 0 \\ 0 & j e^{-j\varphi} & 0 & 0 \\ -j e^{j\varphi} & 0 & 0 & 0 \end{pmatrix} \\
 & & & (XV-16)
 \end{aligned}$$

It should be noted in particular that the obtained matrices check the general properties of DIRAC matrices:

$$(\gamma_0)^2 = 1, \quad (\gamma_r)^2 = (\gamma_\theta)^2 = (\gamma_\varphi)^2 = -1 \quad (XV-17)$$

Developed system of DIRAC matrix writing takes the form:

$$\begin{aligned}
 \eta \begin{pmatrix} \psi_0 \\ \psi_1 \\ \psi_2 \\ \psi_3 \end{pmatrix} &= j \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix} \begin{pmatrix} \frac{\partial \psi_0}{\partial x_t} \\ \frac{\partial \psi_1}{\partial x_t} \\ \frac{\partial \psi_2}{\partial x_t} \\ \frac{\partial \psi_3}{\partial x_t} \end{pmatrix} \\
 + j \begin{pmatrix} 0 & 0 & \cos \theta & \sin \theta e^{-j\varphi} \\ 0 & 0 & \sin \theta e^{j\varphi} & -\cos \theta \\ -\cos \theta & -\sin \theta e^{-j\varphi} & 0 & 0 \\ -\sin \theta e^{j\varphi} & \cos \theta & 0 & 0 \end{pmatrix} \begin{pmatrix} \frac{\partial \psi_0}{\partial r} \\ \frac{\partial \psi_1}{\partial r} \\ \frac{\partial \psi_2}{\partial r} \\ \frac{\partial \psi_3}{\partial r} \end{pmatrix} \\
 + j \frac{1}{r} \begin{pmatrix} 0 & 0 & -\sin \theta & \cos \theta e^{-j\varphi} \\ 0 & 0 & \cos \theta e^{j\varphi} & \sin \theta \\ \sin \theta & -\cos \theta e^{-j\varphi} & 0 & 0 \\ -\cos \theta e^{j\varphi} & -\sin \theta & 0 & 0 \end{pmatrix} \begin{pmatrix} \frac{\partial \psi_0}{\partial \theta} \\ \frac{\partial \psi_1}{\partial \theta} \\ \frac{\partial \psi_2}{\partial \theta} \\ \frac{\partial \psi_3}{\partial \theta} \end{pmatrix} \\
 + j \frac{1}{r \sin \theta} \begin{pmatrix} 0 & 0 & 0 & -j e^{-j\varphi} \\ 0 & 0 & j e^{j\varphi} & 0 \\ 0 & j e^{-j\varphi} & 0 & 0 \\ -j e^{j\varphi} & 0 & 0 & 0 \end{pmatrix} \begin{pmatrix} \frac{\partial \psi_0}{\partial \varphi} \\ \frac{\partial \psi_1}{\partial \varphi} \\ \frac{\partial \psi_2}{\partial \varphi} \\ \frac{\partial \psi_3}{\partial \varphi} \end{pmatrix} +
 \end{aligned}
 \tag{XV-18}$$

We can deduce the 4 equations with partial derivatives expressing the DIRAC system in spherical coordinates:

$$\begin{aligned}
 \eta_{\psi_0} &= j \left\{ \frac{\partial \psi_0}{\partial x_t} + \cos \theta \frac{\partial \psi_2}{\partial r} + \sin \theta e^{-j\varphi} \frac{\partial \psi_3}{\partial r} + \frac{1}{r} \left(-\sin \theta \frac{\partial \psi_2}{\partial \theta} + \cos \theta e^{-j\varphi} \frac{\partial \psi_3}{\partial \theta} \right) - \frac{j e^{-j\varphi}}{r \sin \theta} \frac{\partial \psi_3}{\partial \varphi} \right\} \\
 \eta_{\psi_1} &= j \left\{ \frac{\partial \psi_1}{\partial x_t} + \sin \theta e^{j\varphi} \frac{\partial \psi_2}{\partial r} - \cos \theta \frac{\partial \psi_3}{\partial r} + \frac{1}{r} \left(\cos \theta e^{j\varphi} \frac{\partial \psi_2}{\partial \theta} + \sin \theta \frac{\partial \psi_3}{\partial \theta} \right) + \frac{j e^{j\varphi}}{r \sin \theta} \frac{\partial \psi_2}{\partial \varphi} \right\} \\
 \eta_{\psi_2} &= j \left\{ -\frac{\partial \psi_2}{\partial x_t} - \cos \theta \frac{\partial \psi_0}{\partial r} - \sin \theta e^{-j\varphi} \frac{\partial \psi_1}{\partial r} + \frac{1}{r} \left(\sin \theta \frac{\partial \psi_0}{\partial \theta} - \cos \theta e^{-j\varphi} \frac{\partial \psi_1}{\partial \theta} \right) + \frac{j e^{-j\varphi}}{r \sin \theta} \frac{\partial \psi_1}{\partial \varphi} \right\} \\
 \eta_{\psi_3} &= j \left\{ -\frac{\partial \psi_3}{\partial x_t} - \sin \theta e^{j\varphi} \frac{\partial \psi_0}{\partial r} + \cos \theta \frac{\partial \psi_1}{\partial r} + \frac{1}{r} \left(-\cos \theta e^{j\varphi} \frac{\partial \psi_0}{\partial \theta} - \sin \theta \frac{\partial \psi_1}{\partial \theta} \right) - \frac{j e^{j\varphi}}{r \sin \theta} \frac{\partial \psi_0}{\partial \varphi} \right\}
 \end{aligned}
 \tag{XV-19}$$

It is a non-linear system, and the method used with the Cartesian coordinates is no longer applicable.

The complexity of this system is such that it is difficult (impossible?) to consider a purely mathematical method allowing to lead to an exact solution. We will show in the following chapters that a physical approach based on exchanges of energy between spherical modes allows to progress towards such solutions.

XVI

KLEIN-GORDON equation in spherical coordinates

In the process of finding exact solutions to the DIRAC equation in spherical coordinates, we hypothesize that the conservation equation of energy between the mass energy, wave energy, and impulse energy will be given to us by the KLEIN-GORDON equation in spherical coordinates:

$$\left[\frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2}{\partial \varphi^2} - \frac{\partial^2}{\partial x_t^2} \right] (\psi) = \eta^2 (\psi) \quad (\text{XVI-1})$$

The modal solutions of this equation are presented in the form of a product of three separable functions in (r) , (x_t) , and (θ, φ) .

I - Separable solution (θ, φ)

This solution is based on spherical harmonics $Y_{\ell m}$ that are functions of two parameters:

- ℓ is called the harmonic degree of $Y_{\ell m}$, and it is a natural number.
- m is known as the spherical harmonic order: it is an integer such that $|m| \leq \ell$ and therefore: $m = -\ell, -\ell+1, \dots, \ell$

$$Y_{\ell m}(\theta, \varphi) = (-1)^{\frac{m+|m|}{2}} \sqrt{\frac{2\ell+1}{4\pi}} \sqrt{\frac{(\ell-|m|)!}{(\ell+|m|)!}} P_{\ell}^{|m|}(\cos \theta) \exp(jm\varphi) \quad (\text{XVI-2})$$

$$0 \leq \theta \leq \pi \quad \text{et} \quad 0 \leq \varphi \leq 2\pi$$

P_{ℓ}^m is a LEGENDRE polynomial raised to power m .

Spherical harmonics are solutions of the eigenvalue equation:

$$\left[\frac{1}{\sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial}{\partial \theta} \right) + \frac{1}{\sin^2 \theta} \frac{\partial^2}{\partial \varphi^2} \right] Y_{\ell m}(\theta, \varphi) = -\ell(\ell+1) Y_{\ell m}(\theta, \varphi) \quad (\text{XVI-3})$$

The first standard spherical harmonics are the following:

$$\begin{aligned}
 Y_{00} &= \frac{1}{\sqrt{4\pi}} \\
 Y_{10} &= \sqrt{\frac{3}{4\pi}} \cos \theta & Y_{1-1} &= \sqrt{\frac{3}{8\pi}} \sin \theta e^{-j\varphi} & Y_{11} &= -\sqrt{\frac{3}{8\pi}} \sin \theta e^{j\varphi} \\
 Y_{20} &= \sqrt{\frac{5}{16\pi}} (3\cos^2 \theta - 1) & Y_{2-1} &= \sqrt{\frac{15}{8\pi}} \sin \theta \cos \theta e^{-j\varphi} & Y_{21} &= -\sqrt{\frac{15}{8\pi}} \sin \theta \cos \theta e^{j\varphi} \\
 & & Y_{2-2} &= \sqrt{\frac{15}{32\pi}} \sin^2 \theta e^{-j2\varphi} & Y_{22} &= \sqrt{\frac{15}{32\pi}} \sin^2 \theta e^{j2\varphi} \\
 Y_{30} &= \sqrt{\frac{7}{16\pi}} (5\cos^2 \theta - 1) \cos \theta & Y_{3\mp 1} &= \pm \sqrt{\frac{35}{64\pi}} (5\cos^2 \theta - 1) \sin \theta e^{\mp j\varphi} \\
 Y_{3\mp 2} &= \sqrt{\frac{105}{32\pi}} \sin^2 \theta \cos \theta e^{\mp j2\varphi} & Y_{3\mp 3} &= \pm \sqrt{\frac{35}{64\pi}} \sin^3 \theta e^{\mp j3\varphi}
 \end{aligned} \tag{XVI-4}$$

II – The separable solution in r

This solution is given by any linear combination of spherical BESSEL functions $j_n(r)$ and $y_n(r)$ where the index n is a natural number. These functions can be defined from the BESSEL functions:

$$\begin{aligned}
 j_n(r) &= \sqrt{\frac{\pi}{2r}} J_{n+\frac{1}{2}}(r) \\
 y_n(r) &= \sqrt{\frac{\pi}{2r}} N_{n+\frac{1}{2}}(r)
 \end{aligned} \tag{XVI-5}$$

or from generators, more convenient to determine the explicit formulations:

$$\begin{aligned}
 j_n(r) &= +(-1)^n r^n \left(\frac{1}{r} \frac{d}{dr} \right)^n \left(\frac{\sin r}{r} \right) \\
 y_n(r) &= -(-1)^n r^n \left(\frac{1}{r} \frac{d}{dr} \right)^n \left(\frac{\cos r}{r} \right)
 \end{aligned} \tag{XVI-6}$$

The first spherical BESSEL functions are as follows:

$$\begin{aligned}
 j_0(r) &= \frac{\sin r}{r} & y_0(r) &= -\frac{\cos r}{r} \\
 j_1(r) &= \frac{\sin r}{r^2} - \frac{\cos r}{r} & y_1(r) &= -\frac{\cos r}{r^2} - \frac{\sin r}{r} \\
 j_2(r) &= \left(\frac{3}{r^3} - \frac{1}{r} \right) \sin r - 3 \frac{\cos r}{r^2} & y_2(r) &= -\left(\frac{3}{r^3} - \frac{1}{r} \right) \cos r - 3 \frac{\sin r}{r^2} \\
 j_3(r) &= \left(\frac{15}{r^4} - \frac{6}{r^2} \right) \sin r - \left(\frac{15}{r^3} - \frac{1}{r} \right) \cos r & y_3(r) &= -\left(\frac{15}{r^4} - \frac{6}{r^2} \right) \cos r - \left(\frac{15}{r^3} - \frac{1}{r} \right) \sin r
 \end{aligned} \tag{XVI-7}$$

If we refer to $f_n(r)$ the general function representative of $j_n(r)$ or $y_n(r)$, $f_n(r)$ is solution of the spherical BESSEL differential equation:

$$\frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial}{\partial r} \right) f_n(r) + \left[1 - \frac{n(n+1)}{r^2} \right] f_n(r) = 0 \quad (\text{XVI-8})$$

Or in an equivalent way:

$$\frac{\partial^2 f_n(r)}{\partial r^2} + \frac{2}{r} \frac{\partial f_n(r)}{\partial r} + \left[1 - \frac{n(n+1)}{r^2} \right] f_n(r) = 0 \quad (\text{XVI-9})$$

Spherical cavity modes dependent only from r are called the pulsed modes. They are functions of a radial propagation constant or radial space pulsation that we refer to by k_r and who plays a role analogous to propagation constants k_x , k_y , or k_z for rectangular cavities.

By introducing the variable $R = k_r r$ in spherical BESSEL equation (XVI-9), we get successively:

$$\begin{aligned} \frac{\partial^2 f_n(R)}{\partial R^2} + \frac{2}{R} \frac{\partial f_n(R)}{\partial R} + \left[1 - \frac{n(n+1)}{R^2} \right] f_n(R) &= 0 \\ \frac{1}{k_r^2} \frac{\partial^2 f_n(k_r r)}{\partial r^2} + \frac{1}{k_r^2} \frac{2}{r} \frac{\partial f_n(k_r r)}{\partial r} + \left[1 - \frac{n(n+1)}{(k_r r)^2} \right] f_n(k_r r) &= 0 \\ \frac{\partial^2 f_n(k_r r)}{\partial r^2} + \frac{2}{r} \frac{\partial f_n(k_r r)}{\partial r} + \left[k_r^2 - \frac{n(n+1)}{r^2} \right] f_n(k_r r) &= 0 \\ \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial}{\partial r} \right) f_n(k_r r) + \left[k_r^2 - \frac{n(n+1)}{r^2} \right] f_n(k_r r) &= 0 \end{aligned} \quad (\text{XVI-10})$$

III – The separable solution in x_t

If we refer to this solution by u , it is solution of the eigenvalue equation:

$$\frac{\partial^2}{\partial x_t^2} u(k_t x_t) + k_t^2 u(k_t x_t) = 0 \quad (\text{XVI-11})$$

It is thus constituted by any linear combination of the trigonometric functions $\cos(k_t x_t)$ and $\sin(k_t x_t)$.

IV - The complete solution in stationary modes separated in $(r, \theta, \varphi$ and x_t)

The KLEIN-GORDON equation in spherical coordinates is recalled for memory:

$$\left[\frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2}{\partial \varphi^2} - \frac{\partial^2}{\partial x_t^2} \right] (\psi) = \eta^2 (\psi) \quad (\text{XVI-12})$$

Solutions in stationary modes appear under the form of a product of independent functions in r , (θ, φ) , and x_t :

$$\psi(r, \theta, \varphi, t) = f_r(k_r r) Y_{\ell m}(\theta, \varphi) u(k_t x_t) \quad (\text{XVI-13})$$

The introduction of this form of solution in the KLEIN-GORDON equation in spherical coordinates leads to the following substitutions:

$$\begin{aligned} \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial}{\partial r} \right) \psi(r, \theta, \varphi, t) &= - \left[k_r^2 - \frac{\ell(\ell+1)}{r^2} \right] f_r(k_r r) Y_{\ell m}(\theta, \varphi) u(k_t x_t) \\ \left[\frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2}{\partial \varphi^2} \right] \psi(r, \theta, \varphi, t) &= - \frac{\ell(\ell+1)}{r^2} f_r(k_r r) Y_{\ell m}(\theta, \varphi) u(k_t x_t) \\ - \frac{\partial^2}{\partial x_t^2} \psi(r, \theta, \varphi, t) &= k_t^2 f_r(k_r r) Y_{\ell m}(\theta, \varphi) u(k_t x_t) \end{aligned} \quad (\text{XVI-14})$$

By substitution of these 3 relationships in (XVI-12), yields the relationship of energy:

$$k_t^2 = \eta^2 + k_r^2 \quad (\text{XVI-15})$$

After multiplication by the constant $(\hbar c)^2$, this relation becomes:

$$(\hbar c k_t)^2 = (\hbar c \eta)^2 + (\hbar c k_r)^2 \quad (\text{XVI-16})$$

or again :

$$(\hbar \omega)^2 = (m_0 c^2)^2 + (\hbar c k_r)^2 \quad (\text{XVI-17})$$

As in Cartesian coordinates, it identifies in this relationship three kinds of energy: wave energy $(\hbar \omega)$, masse energy $(m_0 c^2)$ and impulse energy $(\hbar c k_r)$.

It should be noted that this relationship of conservation is independent of the excited modes depending on θ and φ . This means that the distribution between mass energy, wave energy, and impulse energy is not depending on modes excited in rotation following θ and φ .

As in the case of Cartesian coordinates, this relationship of energy conservation will play a fundamental role in obtaining stationary solutions to the DIRAC equation in spherical coordinates.

XVII

Exact solutions of the DIRAC equation in spherical coordinates

In previous chapters, we have shown that the exact solutions of the Dirac equation in spherical coordinates must check the system:

$$\begin{aligned}\eta\psi_0 &= j \left\{ \frac{\partial\psi_0}{\partial x_t} + \cos\theta \frac{\partial\psi_2}{\partial r} + \sin\theta e^{-j\varphi} \frac{\partial\psi_3}{\partial r} + \frac{1}{r} \left(-\sin\theta \frac{\partial\psi_2}{\partial\theta} + \cos\theta e^{-j\varphi} \frac{\partial\psi_3}{\partial\theta} \right) - \frac{j e^{-j\varphi}}{r \sin\theta} \frac{\partial\psi_3}{\partial\varphi} \right\} \\ \eta\psi_1 &= j \left\{ \frac{\partial\psi_1}{\partial x_t} + \sin\theta e^{j\varphi} \frac{\partial\psi_2}{\partial r} - \cos\theta \frac{\partial\psi_3}{\partial r} + \frac{1}{r} \left(\cos\theta e^{j\varphi} \frac{\partial\psi_2}{\partial\theta} + \sin\theta \frac{\partial\psi_3}{\partial\theta} \right) + \frac{j e^{j\varphi}}{r \sin\theta} \frac{\partial\psi_2}{\partial\varphi} \right\} \\ \eta\psi_2 &= j \left\{ -\frac{\partial\psi_2}{\partial x_t} - \cos\theta \frac{\partial\psi_0}{\partial r} - \sin\theta e^{-j\varphi} \frac{\partial\psi_1}{\partial r} + \frac{1}{r} \left(\sin\theta \frac{\partial\psi_0}{\partial\theta} - \cos\theta e^{-j\varphi} \frac{\partial\psi_1}{\partial\theta} \right) + \frac{j e^{-j\varphi}}{r \sin\theta} \frac{\partial\psi_1}{\partial\varphi} \right\} \\ \eta\psi_3 &= j \left\{ -\frac{\partial\psi_3}{\partial x_t} - \sin\theta e^{j\varphi} \frac{\partial\psi_0}{\partial r} + \cos\theta \frac{\partial\psi_1}{\partial r} + \frac{1}{r} \left(-\cos\theta e^{j\varphi} \frac{\partial\psi_0}{\partial\theta} - \sin\theta \frac{\partial\psi_1}{\partial\theta} \right) - \frac{j e^{j\varphi}}{r \sin\theta} \frac{\partial\psi_0}{\partial\varphi} \right\}\end{aligned}\tag{XVII-1}$$

We have also formulated the hypothesis that if exact solutions exist in the form of stationary modes describing energy exchange within the particle, these solutions must be compatible with the equation of conservation of energy established using the KLEIN-GORDON equation:

$$k_t^2 = \eta^2 + k_r^2\tag{XVII-2}$$

However, these elements are insufficient to advance in the search for solutions to the above system (XVII-1).

We must therefore find new features able, in a heuristic approach, to restrict the field of possible solutions.

We are going to do this using two observations of the exact solutions obtained in Cartesian coordinates. Consider for example the following solution:

$$\begin{aligned}\psi_0 &= 0 \\ \psi_1 &= -(\hbar\omega)\sin(k_x x)\cos(k_y y)\cos(k_z z)\cos(k_t x_t) + j(m_0 c^2)\sin(k_x x)\cos(k_y y)\cos(k_z z)\sin(k_t x_t) \\ \psi_2 &= j\hbar c k_y \sin(k_x x)\sin(k_y y)\cos(k_z z)\sin(k_t x_t) + \hbar c k_x \cos(k_x x)\cos(k_y y)\cos(k_z z)\sin(k_t x_t) \\ \psi_3 &= \hbar c k_z \sin(k_x x)\cos(k_y y)\sin(k_z z)\sin(k_t x_t)\end{aligned}\tag{XVII-3}$$

The first observation to be noted is that the bi-spinor representing some of these solutions can be broken down as follows:

- the first spinor expresses on one of its components exchange between mass energy and wave energy.
- the second spinor expresses on its two components impulse energy exchange.

The second observation is linked to the wave-particle duality. It appears in exact solutions in the form of a same spatial mode, one with energy-positive, the other with negative energy (because the presence of $j = \sqrt{-1}$), excited in time quadrature, which indicates that when the mass energy is maximum, the wave energy is zero, and vice versa. Modes that are present in this exchange of energy must be solution of the KLEIN-GORDON equation that appears systematically in a system of DIRAC equations.

Based on these two observations, empirical tests have shown that there is indeed exact spherical solutions to DIRAC system (XVII-1) in the form of stationary modes, with respect to the equation of conservation of energy (XVII-2)

Some spherical modes do not allow to obtain exact solutions to the DIRAC equation. The empirical methodology proposed here does not say whether there is a mathematical absence of these solutions, or if these modes-related solutions are achievable by other methods

The first spherical mode $Y_{00}(\theta, \varphi)$ solution of the KLEIN-GORDON equation is equal to a constant, so it is independent of θ , and φ . It doesn't lead to an exact solution of the Dirac equation.

We will develop in detail in the following chapters, exact solutions for some modes of the spherical solutions of the KLEIN-GORDON equation.

XVIII

Exact solutions of DIRAC equation on the modes Y_{n-n} et Y_{nn}

We are interested in this chapter to modes solutions of the KLEIN-GORDON equation for which the parameters l and $|m|$ are equal, and we shall put $l = |m| = n \neq 0$.

The angular description of these modes is given by the following spherical harmonics, in which the normalization constant, which plays no role to establish the validity of the solutions has been omitted:

$$Y_{n-n} = \sin^n \theta e^{-jn\varphi} \quad Y_{nn} = \sin^n \theta e^{jn\varphi} \quad (\text{XVIII-1})$$

We will work on the mode Y_{n-n} , before deriving, by simple considerations, the expression of the Y_{nn} mode solutions.

The spherical harmonic Y_{n-n} is associated with the radial function given by two spherical BESSEL functions of order n :

$$f_n(k_r r) = j_n(k_r r) = \sqrt{\frac{\pi}{2r}} J_{n+\frac{1}{2}}(k_r r) \quad (\text{XVIII-2})$$

or

$$f_n(k_r r) = y_n(k_r r) = \sqrt{\frac{\pi}{2r}} N_{n+\frac{1}{2}}(k_r r) \quad (\text{XVIII-3})$$

Modal solutions of the KLEIN-GORDON equation expressed under the form:

$$\psi(r, \theta, \varphi, t) = f_n(k_r r) Y_{n-n}(\theta, \varphi) u(k_t x_t) \quad (\text{XVIII-4})$$

where the time dependence is given by any linear combination of functions: $u(k_t x_t) = \cos(k_t x_t)$ ou $u(k_t x_t) = \sin(k_t x_t)$

The approach discussed in the previous chapter incentive to propose the following solution which will prove to be an exact solution:

$$\begin{aligned}
 \psi_0 &= \eta f_n(k_r r) \sin^n \theta e^{-nj\varphi} \sin(k_t x_t) + j k_t f_n(k_r r) \sin^n \theta e^{-nj\varphi} \cos(k_t x_t) \\
 \psi_1 &= 0 \\
 \psi_2 &= j \cos \theta \sin^n \theta e^{-nj\varphi} \sin(k_t x_t) \left\{ -k_r f_n'(k_r r) + n \frac{f_n(k_r r)}{r} \right\} \\
 \psi_3 &= j \sin(k_t x_t) \sin^{n-1} \theta e^{-(n-1)j\varphi} \left\{ -\sin^2 \theta k_r f_n'(k_r r) - n(\cos^2 \theta + 1) \frac{f_n(k_r r)}{r} \right\}
 \end{aligned} \tag{XVIII-5}$$

in which we have adopted the light notation: $f_n'(k_r r) = df_n(k_r r)/d(k_r r)$.

The first spinor describes the duality wave-corpucle through the wave function ψ_0 . Mass energy is exchanged with the wave energy on the spatial mode defined by the spherical harmonic Y_{n-n} . It responds to the approach discussed in the previous chapter.

The second spinor exchanges impulse energy to which we will return later. It is inferred from the spherical DIRAC system after substitution of the first spinor.

The equation of conservation of energy associated with this solution is recalled for memory:

$$k_t^2 = \eta^2 + k_r^2 \tag{XVIII-6}$$

This relationship is independent of the nature of the excited spherical modes.

In previous chapters, we have shown that the exact solutions of the DIRAC equation in spherical coordinates must check the system:

$$\begin{aligned}
 \eta \psi_0 &= j \left\{ \frac{\partial \psi_0}{\partial x_t} + \cos \theta \frac{\partial \psi_2}{\partial r} + \sin \theta e^{-j\varphi} \frac{\partial \psi_3}{\partial r} + \frac{1}{r} \left(-\sin \theta \frac{\partial \psi_2}{\partial \theta} + \cos \theta e^{-j\varphi} \frac{\partial \psi_3}{\partial \theta} \right) - \frac{j e^{-j\varphi}}{r \sin \theta} \frac{\partial \psi_3}{\partial \varphi} \right\} \\
 \eta \psi_1 &= j \left\{ \frac{\partial \psi_1}{\partial x_t} + \sin \theta e^{j\varphi} \frac{\partial \psi_2}{\partial r} - \cos \theta \frac{\partial \psi_3}{\partial r} + \frac{1}{r} \left(\cos \theta e^{j\varphi} \frac{\partial \psi_2}{\partial \theta} + \sin \theta \frac{\partial \psi_3}{\partial \theta} \right) + \frac{j e^{j\varphi}}{r \sin \theta} \frac{\partial \psi_2}{\partial \varphi} \right\} \\
 \eta \psi_2 &= j \left\{ -\frac{\partial \psi_2}{\partial x_t} - \cos \theta \frac{\partial \psi_0}{\partial r} - \sin \theta e^{-j\varphi} \frac{\partial \psi_1}{\partial r} + \frac{1}{r} \left(\sin \theta \frac{\partial \psi_0}{\partial \theta} - \cos \theta e^{-j\varphi} \frac{\partial \psi_1}{\partial \theta} \right) + \frac{j e^{-j\varphi}}{r \sin \theta} \frac{\partial \psi_1}{\partial \varphi} \right\} \\
 \eta \psi_3 &= j \left\{ -\frac{\partial \psi_3}{\partial x_t} - \sin \theta e^{j\varphi} \frac{\partial \psi_0}{\partial r} + \cos \theta \frac{\partial \psi_1}{\partial r} + \frac{1}{r} \left(-\cos \theta e^{j\varphi} \frac{\partial \psi_0}{\partial \theta} - \sin \theta \frac{\partial \psi_1}{\partial \theta} \right) - \frac{j e^{j\varphi}}{r \sin \theta} \frac{\partial \psi_0}{\partial \varphi} \right\}
 \end{aligned} \tag{XVIII-7}$$

This system verification requires only basic calculations, but it is laborious. It is however a decisive argument to convince of the validity of the proposed solution, and aid is proposed below giving the explicit formulation of each of the terms of these equations.

The first equation, more complex because it contains the KLEIN-GORDON equation, will be treated as a last, and we propose to give first the terms relating to equations 2, 3, and 4. Verification of the sums is left to the reader care.

Verification elements of equation 2 of the spherical DIRAC system:

$$\eta\psi_1 = j \left\{ \frac{\partial\psi_1}{\partial x_t} + \sin\theta e^{j\varphi} \frac{\partial\psi_2}{\partial r} - \cos\theta \frac{\partial\psi_3}{\partial r} + \frac{1}{r} \left(\cos\theta e^{j\varphi} \frac{\partial\psi_2}{\partial\theta} + \sin\theta \frac{\partial\psi_3}{\partial\theta} \right) + \frac{j e^{j\varphi}}{r \sin\theta} \frac{\partial\psi_2}{\partial\varphi} \right\} \quad (\text{XVIII-8})$$

$$\eta\psi_1 = 0$$

$$j \frac{\partial\psi_1}{\partial x_t} = 0$$

$$\begin{aligned} j \sin\theta e^{j\varphi} \frac{\partial\psi_2}{\partial r} &= -\cos\theta \sin^{n+1}\theta \sin(k_t x_t) e^{-(n-1)j\varphi} \left\{ -k_r^2 f_n''(k_r r) + n \left(\frac{k_r f_n'(k_r r)}{r} - \frac{f_n(k_r r)}{r^2} \right) \right\} \\ -j \cos\theta \frac{\partial\psi_3}{\partial r} &= \cos\theta \sin^{n-1}\theta \sin(k_t x_t) e^{-(n-1)j\varphi} \left\{ -\sin^2\theta k_r^2 f_n''(k_r r) - n \left(\cos^2\theta + 1 \right) \left(\frac{k_r f_n'(k_r r)}{r} - \frac{f_n(k_r r)}{r^2} \right) \right\} \\ j \frac{\cos\theta e^{j\varphi}}{r} \frac{\partial\psi_2}{\partial\theta} &= -\cos\theta \sin^{n-1}\theta \left(n \cos^2\theta - \sin^2\theta \right) \sin(k_t x_t) e^{-(n-1)j\varphi} \left\{ -\frac{k_r f_n'(k_r r)}{r} + n \frac{f_n(k_r r)}{r^2} \right\} \\ j \frac{\sin\theta}{r} \frac{\partial\psi_3}{\partial\theta} &= \sin(k_t x_t) \sin^{n-1}\theta e^{-(n-1)j\varphi} \left\{ \begin{aligned} &(n+1) \sin^2\theta \cos\theta \frac{k_r f_n'(k_r r)}{r} \\ &+ n \left((n-1) \cos^3\theta - 2 \sin^2\theta \cos\theta + (n-1) \cos\theta \right) \frac{f_n(k_r r)}{r^2} \end{aligned} \right\} \\ -\frac{e^{j\varphi}}{r \sin\theta} \frac{\partial\psi_2}{\partial\varphi} &= -n \cos\theta \sin^{n-1}\theta \sin(k_t x_t) e^{-(n-1)j\varphi} \left\{ -\frac{k_r f_n'(k_r r)}{r} + n \frac{f_n(k_r r)}{r^2} \right\} \end{aligned}$$

(XVIII-9)

Verification elements of equation 3 of the spherical DIRAC system:

$$\eta\psi_2 = j \left\{ -\frac{\partial\psi_2}{\partial x_t} - \cos\theta \frac{\partial\psi_0}{\partial r} - \sin\theta e^{-j\varphi} \frac{\partial\psi_1}{\partial r} + \frac{1}{r} \left(\sin\theta \frac{\partial\psi_0}{\partial\theta} - \cos\theta e^{-j\varphi} \frac{\partial\psi_1}{\partial\theta} \right) + \frac{j e^{-j\varphi}}{r \sin\theta} \frac{\partial\psi_1}{\partial\varphi} \right\} \quad (\text{XVIII-10})$$

$$\begin{aligned} \eta\psi_2 &= j \eta \cos\theta \sin^n\theta e^{-nj\varphi} \sin(k_t x_t) \left\{ -k_r f_n'(k_r r) + n \frac{f_n(k_r r)}{r} \right\} \\ -j \frac{\partial\psi_2}{\partial x_t} &= k_t \cos\theta \sin^n\theta e^{-nj\varphi} \cos(k_t x_t) \left\{ -k_r f_n'(k_r r) + n \frac{f_n(k_r r)}{r} \right\} \\ -j \cos\theta \frac{\partial\psi_0}{\partial r} &= -j \eta k_r f_n'(k_r r) \cos\theta \sin^n\theta e^{-nj\varphi} \sin(k_t x_t) + k_t k_r f_n'(k_r r) \cos\theta \sin^n\theta e^{-nj\varphi} \cos(k_t x_t) \\ -j \sin\theta e^{-j\varphi} \frac{\partial\psi_1}{\partial r} &= 0 \\ j \frac{\sin\theta}{r} \frac{\partial\psi_0}{\partial\theta} &= j \eta \frac{f_n(k_r r)}{r} n \sin^n\theta \cos\theta e^{-nj\varphi} \sin(k_t x_t) - k_t \frac{f_n(k_r r)}{r} n \sin^n\theta \cos\theta e^{-nj\varphi} \cos(k_t x_t) \\ -j \frac{\cos\theta}{r} e^{-j\varphi} \frac{\partial\psi_1}{\partial\theta} &= 0 \\ -\frac{e^{-j\varphi}}{r \sin\theta} \frac{\partial\psi_1}{\partial\varphi} &= 0 \end{aligned}$$

(XVIII-11)

Verification elements of equation 4 of the spherical DIRAC system:

$$\eta\psi_3 = j \left\{ -\frac{\partial\psi_3}{\partial x_t} - \sin\theta e^{j\varphi} \frac{\partial\psi_0}{\partial r} + \cos\theta \frac{\partial\psi_1}{\partial r} + \frac{1}{r} \left(-\cos\theta e^{j\varphi} \frac{\partial\psi_0}{\partial\theta} - \sin\theta \frac{\partial\psi_1}{\partial\theta} \right) - \frac{j e^{j\varphi}}{r \sin\theta} \frac{\partial\psi_0}{\partial\varphi} \right\} \quad (\text{XVIII-12})$$

$$\begin{aligned} \eta\psi_3 &= j\eta \sin(k_t x_t) \sin^{n-1} \theta e^{-(n-1)j\varphi} \left\{ -\sin^2 \theta k_r f_n'(k_r r) - n(\cos^2 \theta + 1) \frac{f_n(k_r r)}{r} \right\} \\ -j \frac{\partial\psi_3}{\partial x_t} &= k_t \cos(k_t x_t) \sin^{n-1} \theta e^{-(n-1)j\varphi} \left\{ -\sin^2 \theta k_r f_n'(k_r r) - n(\cos^2 \theta + 1) \frac{f_n(k_r r)}{r} \right\} \\ -j \sin\theta e^{j\varphi} \frac{\partial\psi_0}{\partial r} &= -j\eta k_r f_n'(k_r r) \sin^{n+1} \theta e^{-(n-1)j\varphi} \sin(k_t x_t) + k_r k_t f_n'(k_r r) \sin^{n+1} \theta e^{-(n-1)j\varphi} \cos(k_t x_t) \\ j \cos\theta \frac{\partial\psi_1}{\partial r} &= 0 \\ -j \frac{\cos\theta e^{j\varphi}}{r} \frac{\partial\psi_0}{\partial\theta} &= \left\{ \begin{array}{l} -j\eta \frac{f_n(k_r r)}{r} n \sin^{n-1} \theta \cos^2 \theta e^{-(n-1)j\varphi} \sin(k_t x_t) \\ + k_t \frac{f_n(k_r r)}{r} n \sin^{n-1} \theta \cos^2 \theta e^{-(n-1)j\varphi} \cos(k_t x_t) \end{array} \right\} \\ -j \frac{\sin\theta}{r} \frac{\partial\psi_1}{\partial\theta} &= 0 \\ \frac{e^{j\varphi}}{r \sin\theta} \frac{\partial\psi_0}{\partial\varphi} &= -jn\eta \frac{f_n(k_r r)}{r} \sin^{n-1} \theta e^{-(n-1)j\varphi} \sin(k_t x_t) + nk_t \frac{f_n(k_r r)}{r} \sin^{n-1} \theta e^{-(n-1)j\varphi} \cos(k_t x_t) \end{aligned} \quad (\text{XVIII-13})$$

Verification elements of equation 1 of the spherical DIRAC system:

As stated previously, this verification justifies special attention because it contains the equation of conservation of energy derived from the KLEIN-GORDON equation.

$$\eta\psi_0 = j \left\{ \frac{\partial\psi_0}{\partial x_t} + \cos\theta \frac{\partial\psi_2}{\partial r} + \sin\theta e^{-j\varphi} \frac{\partial\psi_3}{\partial r} + \frac{1}{r} \left(-\sin\theta \frac{\partial\psi_2}{\partial\theta} + \cos\theta e^{-j\varphi} \frac{\partial\psi_3}{\partial\theta} \right) - \frac{j e^{-j\varphi}}{r \sin\theta} \frac{\partial\psi_3}{\partial\varphi} \right\} \quad (\text{XVIII-14})$$

The explicit calculation of each of the terms is given below:

$$\begin{aligned}
\eta \Psi_0 &= \eta^2 f_n(k_r, r) \sin^n \theta e^{-nj\varphi} \sin(k_t x_t) + j \eta k_t f_n(k_r, r) \sin^n \theta e^{-nj\varphi} \cos(k_t x_t) \\
j \frac{\partial \Psi_0}{\partial x_t} &= j \eta k_t f_n(k_r, r) \sin^n \theta e^{-nj\varphi} \cos(k_t x_t) + k_t^2 f_n(k_r, r) \sin^n \theta e^{-nj\varphi} \sin(k_t x_t) \\
j \cos \theta \frac{\partial \Psi_2}{\partial r} &= -\cos^2 \theta \sin^n \theta \sin(k_t x_t) e^{-nj\varphi} \left\{ -k_r^2 f_n''(k_r, r) + n \left(\frac{k_r f_n'(k_r, r)}{r} - \frac{f_n(k_r, r)}{r^2} \right) \right\} \\
j \sin \theta e^{-j\varphi} \frac{\partial \Psi_3}{\partial r} &= -\sin^n \theta \sin(k_t x_t) e^{-nj\varphi} \left\{ -\sin^2 \theta k_r^2 f_n''(k_r, r) - n(\cos^2 \theta + 1) \left(\frac{k_r f_n'(k_r, r)}{r} - \frac{f_n(k_r, r)}{r^2} \right) \right\} \\
-j \frac{\sin \theta}{r} \frac{\partial \Psi_2}{\partial \theta} &= \sin^n (n \cos^2 \theta - \sin^2 \theta) \sin(k_t x_t) e^{-nj\varphi} \left\{ -\frac{k_r f_n'(k_r, r)}{r} + n \frac{f_n(k_r, r)}{r^2} \right\} \\
j \frac{\cos \theta}{r} e^{-j\varphi} \frac{\partial \Psi_3}{\partial \theta} &= \sin(k_t x_t) \cos \theta e^{-nj\varphi} \left\{ (n+1) \sin^n \theta \cos \theta \frac{k_r f_n'(k_r, r)}{r} \right. \\
&\quad \left. + n((n-1) \sin^{n-2} \theta \cos^3 \theta - 2 \sin^n \theta \cos \theta + (n-1) \sin^{n-2} \theta \cos \theta) \frac{f_n(k_r, r)}{r} \right\} \\
\frac{e^{-j\varphi}}{r \sin \theta} \frac{\partial \Psi_3}{\partial \varphi} &= (n-1) \sin(k_t x_t) \sin^{n-2} \theta e^{-nj\varphi} \left\{ -\sin^2 \theta k_r \frac{f_n'(k_r, r)}{r} - n \frac{\cos^2 \theta + 1}{r^2} f_n(k_r, r) \right\}
\end{aligned} \tag{XVIII-15}$$

After simplification by $\sin(k_t x_t) \cdot e^{-j\varphi}$, the first equation of the DIRAC system is written:

$$\begin{aligned}
\eta^2 f_n(k_r, r) \sin^n \theta &= \sin^n \theta k_t^2 f_n(k_r, r) + \sin^n \theta k_r^2 f_n''(k_r, r) \\
+ \frac{k_r f_n'(k_r, r)}{r} &\left\{ -n \cos^2 \theta \sin^n \theta + n \sin^n \theta (\cos^2 \theta + 1) - \sin^n (n \cos^2 \theta - \sin^2 \theta) \right\} \\
+ \frac{f_n(k_r, r)}{r^2} &\left\{ + (n+1) \sin^n \theta \cos^2 \theta - (n-1) \sin^n \theta \right. \\
&\left. \left\{ n \cos^2 \theta \sin^n \theta - n \sin^n \theta (\cos^2 \theta + 1) + n \sin^n (n \cos^2 \theta - \sin^2 \theta) \right\} \right. \\
&\left. \left\{ + n((n-1) \sin^{n-2} \theta \cos^4 \theta - 2 \sin^n \theta \cos^2 \theta + (n-1) \sin^{n-2} \theta \cos^2 \theta) \right\} \right\} \\
&\left\{ -n(n-1) \sin^{n-2} \theta (\cos^2 \theta + 1) \right\}
\end{aligned} \tag{XVIII-16}$$

Terms within brackets reduce in a remarkable manner to obtain:

$$\eta^2 f_n(k_r, r) = k_t^2 f_n(k_r, r) + k_r^2 f_n''(k_r, r) + 2 \frac{k_r f_n'(k_r, r)}{r} - n(n+1) \frac{f_n(k_r, r)}{r^2} \tag{XVIII-17}$$

We can rewrite this relationship by inserting the partial derivatives with respect to r :

$$\eta^2 f_n(k_r, r) = k_t^2 f_n(k_r, r) + \frac{\partial^2}{\partial r^2} f_n(k_r, r) + \frac{2}{r} \frac{\partial}{\partial r} f_n(k_r, r) - n(n+1) \frac{f_n(k_r, r)}{r^2} \tag{XVIII-18}$$

Or again:

$$\eta^2 f_n(k_r, r) = k_t^2 f_n(k_r, r) + \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial}{\partial r} \right) f_n(k_r, r) - n(n+1) \frac{f_n(k_r, r)}{r^2} \tag{XVIII-19}$$

Using the property of the spherical BESSEL function established in (XVI-10) and recalled for memory:

$$\frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial}{\partial r} \right) f_\ell(k_r r) + \left[k_r^2 - \frac{\ell(\ell+1)}{r^2} \right] f_\ell(k_r r) = 0 \quad (\text{XVIII-20})$$

we obtain for $\ell = n$:

$$\begin{aligned} \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial}{\partial r} \right) f_n(k_r r) + \left[k_r^2 - \frac{n(n+1)}{r^2} \right] f_n(k_r r) &= 0 \\ \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial}{\partial r} \right) f_n(k_r r) - \frac{n(n+1)}{r^2} f_n(k_r r) &= -k_r^2 f_n(k_r r) \end{aligned} \quad (\text{XVIII-21})$$

By postponing this last result in the below mentioned (XVIII-19) relationship:

$$\eta^2 f_n(k_r r) = k_t^2 f_n(k_r r) + \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial}{\partial r} \right) f_n(k_r r) - n(n+1) \frac{f_n(k_r r)}{r^2} \quad (\text{XVIII-22})$$

We obtain the relationship of conservation of energy:

$$\begin{aligned} \eta^2 f_n(k_r r) &= k_t^2 f_n(k_r r) - k_r^2 f_n(k_r r) \\ k_t^2 &= \eta^2 + k_r^2 \end{aligned} \quad (\text{XVIII-23})$$

Verification of the exact solution to the DIRAC equation in spherical coordinates is completed. As in Cartesian coordinates, to be valid, this solution must be associated with the equation of conservation of energy in spherical coordinates.

At the conclusion of this chapter, we look at the similar mode in Y_{nn} :

$$Y_{nn} = \sin^n \theta e^{jn\varphi} \quad (\text{XVIII-24})$$

If it ignores the normalization constant that is not involved in the calculations, the only difference from the spherical harmonic Y_{n-n} of the preceding paragraph is in the sign of φ .

This induces an immediate consequence: we can no longer find exact solution by shifting the duality wave-corpuscule by the wave function ψ_0 . The exact solution to the spherical DIRAC system can only be constructed by carrying this duality by the wave function ψ_1 . We then get a similar solution to the previous one, in which the wave functions of the spinor that carries the pulse energy is exchanged:

$$\begin{aligned} \psi_0 &= 0 \\ \psi_1 &= \eta f_n(k_r r) \sin^n \theta e^{nj\varphi} \sin(k_t x_t) + j k_t f_n(k_r r) \sin^n \theta e^{nj\varphi} \cos(k_t x_t) \\ \psi_2 &= j \sin(k_t x_t) \sin^{n-1} \theta e^{(n-1)j\varphi} \left\{ -\sin^2 \theta k_r f_n'(k_r r) - n(\cos^2 \theta + 1) \frac{f_n(k_r r)}{r} \right\} \\ \psi_3 &= j \cos \theta \sin^n \theta e^{nj\varphi} \sin(k_t x_t) \left\{ k_r f_n'(k_r r) - n \frac{f_n(k_r r)}{r} \right\} \end{aligned} \quad (\text{XVIII-25})$$

We will see in a next chapter that the sign of φ determines the direction of rotation of the energy. Using the conventional rules of orientation of the spin (rule of the corkscrew for example), we can associate a spin direction to each of the proposed solutions.

The spinor:

$$\begin{aligned}\psi_0 &= \eta f_n(k, r) \sin^n \theta e^{-nj\varphi} \sin(k_t x_t) + j k_t f_n(k, r) \sin^n \theta e^{-nj\varphi} \cos(k_t x_t) \\ \psi_1 &= 0\end{aligned}\tag{XVIII-26}$$

can be associated with the positive z-oriented spin and it is usually called spin "up".

The spinor:

$$\begin{aligned}\psi_0 &= 0 \\ \psi_1 &= \eta f_n(k, r) \sin^n \theta e^{nj\varphi} \sin(k_t x_t) + j k_t f_n(k, r) \sin^n \theta e^{nj\varphi} \cos(k_t x_t)\end{aligned}\tag{XVIII-27}$$

can be associated with the negative z-oriented spin and it is usually called spin "down".

XIX

Other exact solutions

I - The rotation modes

Basic exact solutions presented in the previous chapter do not highlight the rotation of power.

In doing similarly looking for propagative solution during the study in Cartesian coordinates, it is possible to find exact solutions that express this rotation.

Spherical DIRAC system is linear with respect to the sine and cosine functions of x_t , then we can build new exact solutions by summing solutions whose modes variations cover x_t .

We choose to work in a way representing a 'down' spin and we call (solution 1) the solution obtained in (XVIII-25):

$$\begin{aligned}\psi_0 &= 0 \\ \psi_1 &= \eta f_n(k_r, r) \sin^n \theta e^{nj\phi} \sin(k_t x_t) + j k_t f_n(k_r, r) \sin^n \theta e^{nj\phi} \cos(k_t x_t) \\ \psi_2 &= j \sin(k_t x_t) \sin^{n-1} \theta e^{(n-1)j\phi} \left\{ -\sin^2 \theta k_r f_n'(k_r, r) - n(\cos^2 \theta + 1) \frac{f_n(k_r, r)}{r} \right\} \\ \psi_3 &= j \cos \theta \sin^n \theta e^{nj\phi} \sin(k_t x_t) \left\{ k_r f_n'(k_r, r) - n \frac{f_n(k_r, r)}{r} \right\}\end{aligned} \quad (\text{XIX-1})$$

We call (solution 2) the exact solution obtained by exchanging the sine and cosine in the wave function ψ_1 :

$$\begin{aligned}\psi_0 &= 0 \\ \psi_1 &= \eta f_n(k_r, r) \sin^n \theta e^{nj\phi} \cos(k_t x_t) - j k_t f_n(k_r, r) \sin^n \theta e^{nj\phi} \sin(k_t x_t) \\ \psi_2 &= j \cos(k_t x_t) \sin^{n-1} \theta e^{(n-1)j\phi} \left\{ -\sin^2 \theta k_r f_n'(k_r, r) - n(\cos^2 \theta + 1) \frac{f_n(k_r, r)}{r} \right\} \\ \psi_3 &= j \cos \theta \sin^n \theta e^{nj\phi} \cos(k_t x_t) \left\{ k_r f_n'(k_r, r) - n \frac{f_n(k_r, r)}{r} \right\}\end{aligned} \quad (\text{XIX-2})$$

We build the exact solution obtained by the linear combination: (solution 2) + j (solution 1):

$$\psi_0 = 0$$

$$\psi_1 = \eta f_n(k_r r) \sin^n \theta \exp j(k_t x_t + n\varphi) - k_t f_n(k_r r) \sin^n \theta \exp j(k_t x_t + n\varphi)$$

$$\psi_2 = j \exp j(k_t x_t + (n-1)\varphi) \sin^{n-1} \theta \left\{ -\sin^2 \theta k_r f_n'(k_r r) - n(\cos^2 \theta + 1) \frac{f_n(k_r r)}{r} \right\} \quad (\text{XIX-3})$$

$$\psi_3 = j \cos \theta \sin^n \theta \exp j(k_t x_t + n\varphi) \left\{ k_r f_n'(k_r r) - n \frac{f_n(k_r r)}{r} \right\}$$

Terms with $\exp j(k_t x_t + n\varphi)$ and $\exp j(k_t x_t + (n-1)\varphi)$ represent both temporal rotation energy depending on the angle φ and exchanges between positive and negative energy shown by the presence of the imaginary term j .

The direction of rotation can be evaluated in the same way that the meaning of a progressive wave propagation.

A wave as $(\omega t - kx)$ is moving towards the positive x : a wave as $(k_t x_t - n\varphi)$ is moving towards positive φ .

A wave as $(\omega t + kx)$ is progressing towards the negative x : a wave as $(k_t x_t + n\varphi)$ moves towards the negative φ , which justifies the name of spin 'down' for the solution (XIX-3).

It appears impossible to construct exact solutions in rotation as $(k_t x_t + \theta)$ because there are no similar solutions in $\sin\theta$ and $\cos\theta$ which are necessary to obtain by combination with temporal $\sin(k_t x_t)$ and $\cos(k_t x_t)$ functions of rotation by θ .

II – Other solutions

Among the modal solutions of the KLEIN-GORDON equation, for an order l given, $|m|$ can take all values between 0 and l . The previous chapters study show that the modes such as $l = |m| = n \neq 0$ lead to exact solutions of the DIRAC equation.

It turns out that all modes do not result in exact solutions.

A specific test on the Y_{2-1} and Y_{21} modes shows that these modes allow exact solutions which are reported below.

II-1 - mode Y_{2-1}

It is associated with the spherical harmonic Y_{2-1} :

$$Y_{2-1} = \sqrt{\frac{15}{8\pi}} \sin \theta \cos \theta e^{-j\varphi} \quad (\text{XIX-4})$$

The radial function is given by two spherical BESSEL functions of order 2:

$$f_2(k_r r) = j_2(k_r r) = \left(\frac{3}{(k_r r)^3} - \frac{1}{(k_r r)} \right) \sin(k_r r) - 3 \frac{\cos(k_r r)}{(k_r r)^2} \quad (\text{XIX-5})$$

ou

$$f_2(k_r r) = y_2(k_r r) = - \left(\frac{3}{(k_r r)^3} - \frac{1}{(k_r r)} \right) \cos(k_r r) - 3 \frac{\sin(k_r r)}{(k_r r)^2} \quad (\text{XIX-6})$$

Modal solutions of the KLEIN-GORDON equation expressed in the form:

$$\psi(r, \theta, \varphi, t) = f_2(k_r r) Y_{2-1}(\theta, \varphi) u(k_t x_t) \quad (\text{XIX-7})$$

where the time dependence is given by any linear combination of functions: $u(k_t x_t) = \cos(k_t x_t)$ or $u(k_t x_t) = \sin(k_t x_t)$

The way proposed in the previous chapters allows to get the following exact solution:

$$\begin{aligned} \psi_0 &= \eta f_2(k_r r) \sin \theta \cos \theta e^{-j\varphi} \sin(k_t x_t) + j k_t f_2(k_r r) \sin \theta \cos \theta e^{-j\varphi} \cos(k_t x_t) \\ \psi_1 &= 0 \\ \psi_2 &= j \sin \theta e^{-j\varphi} \sin(k_t x_t) \left\{ -\cos^2 \theta k_r f_2'(k_r r) + (1 - 2 \sin^2 \theta) \frac{f_2(k_r r)}{r} \right\} \\ \psi_3 &= j \sin(k_t x_t) \cos \theta \left\{ -\sin^2 \theta k_r f_2'(k_r r) - 2 \cos^2 \theta \frac{f_2(k_r r)}{r} \right\} \end{aligned} \quad (\text{XIX-8})$$

in which we adopted the notation: $f_2'(k_r r) = df_2(k_r r)/d(k_r r)$.

II-2 - mode Y₂₁

It is associated with the spherical harmonic Y_{21} recalled below:

$$Y_{21} = -\sqrt{\frac{15}{8\pi}} \sin \theta \cos \theta e^{j\varphi} \quad (\text{XIX-9})$$

The normalizing constant being not involved in calculations, the only difference from the spherical harmonic Y_{2-1} of the preceding paragraph is in the sign of φ . This indicates that the direction of rotation of the energy is reversed from the Y_{2-1} mode.

The following exact solution is obtained:

$$\begin{aligned} \psi_0 &= 0 \\ \psi_1 &= \eta f_2(k_r r) \sin \theta \cos \theta e^{j\varphi} \sin(k_t x_t) + j k_t f_2(k_r r) \sin \theta \cos \theta e^{j\varphi} \cos(k_t x_t) \\ \psi_2 &= -j \sin(k_t x_t) \cos \theta \left\{ \sin^2 \theta k_r f_2'(k_r r) + 2 \cos^2 \theta \frac{f_2(k_r r)}{r} \right\} \\ \psi_3 &= j \sin(k_t x_t) \sin \theta e^{j\varphi} \left\{ \cos^2 \theta k_r f_2'(k_r r) - (1 - 2 \sin^2 \theta) \frac{f_2(k_r r)}{r} \right\} \end{aligned} \quad (\text{XIX-10})$$

II-3 – Other modes

It has not been discovered simple rule for determining if a mode is an exact solution or not when $l > |m|$, forcing a search to specific checks which becomes particularly laborious when

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the order of the modes increases. It is an open problem: If some modes do not lead to exact solutions of the DIRAC equation, there is necessarily a physical reason that prevents the installation of these modes and which should be elucidated.

XX

Some considerations on exact solutions in spherical coordinates

The passage of the Cartesian coordinates to spherical coordinates allows to retrieve properties that have been highlighted in the analysis of exact solutions in Cartesian coordinates, in a completely deterministic approach.

The major advantage of a formulation of exact solutions to the DIRAC equation in spherical coordinates lies in the fact that we'll be able express analytically the rotation of power.

Some first elements of reflection are proposed below. These elements seem to confirm that the developed solutions are in agreement with the main observed physical phenomena. However we will have to move further in this reflection before concluding or not to the final validity of the proposed solutions.

I – Particle of spin 1

Let us consider an exact solution obtained on mode Y_{1-1} :

$$\begin{aligned}\psi_0 &= \eta f_1(k_r r) \sin \theta e^{-j\varphi} \sin(k_t x_t) + j k_t f_1(k_r r) \sin \theta e^{-j\varphi} \cos(k_t x_t) \\ \psi_1 &= 0 \\ \psi_2 &= j \cos \theta \sin \theta e^{-j\varphi} \sin(k_t x_t) \left\{ -k_r f_1'(k_r r) + \frac{f_1(k_r r)}{r} \right\} \\ \psi_3 &= j \sin(k_t x_t) \left\{ -\sin^2 \theta k_r f_1'(k_r r) - (\cos^2 \theta + 1) \frac{f_1(k_r r)}{r} \right\}\end{aligned}\tag{XX-1}$$

where f_1 is a spherical BESSEL function of order 1:

$$j_1(k_r r) = \frac{\sin(k_r r)}{(k_r r)^2} - \frac{\cos(k_r r)}{(k_r r)} \quad y_1(k_r r) = -\frac{\cos(k_r r)}{(k_r r)^2} - \frac{\sin(k_r r)}{(k_r r)}\tag{XX-2}$$

and where we adopted the notation $f_1'(k_r r) = df_1(k_r r)/d(k_r r)$.

After multiplication by the constant $\hbar c$, each of the wave functions has the dimension of energy:

$$\begin{aligned}
 \psi_0 &= (m_0 c^2) f_1(k_r, r) \sin \theta e^{-j\varphi} \sin(k_t x_t) + j(\hbar \omega) f_1(k_r, r) \sin \theta e^{-j\varphi} \cos(k_t x_t) \\
 \psi_1 &= 0 \\
 \psi_2 &= j(\hbar c k_r) \cos \theta \sin \theta e^{-j\varphi} \sin(k_t x_t) \left\{ -f_1'(k_r, r) + \frac{f_1(k_r, r)}{k_r r} \right\} \\
 \psi_3 &= j(\hbar c k_r) \sin(k_t x_t) \left\{ -\sin^2 \theta f_1'(k_r, r) - (\cos^2 \theta + 1) \frac{f_1(k_r, r)}{k_r r} \right\}
 \end{aligned} \tag{XX-3}$$

The functions above have a finite value in the vicinity of $r = 0$ only when f_1 is a spherical Bessel function of the first kind J_1 . We will maintain despite all subsequently the name f_1 in order to discuss solutions in their greater generality.

We ignore in the discussion that follows the multiplicative constant for standardization. It is recalled that this constant has a dual role: make so that each wave function has the dimension of a square root of volume energy density, and that the integral of the density on the volume containing the energy gives the total energy contained in the particle.

The first spinor formed by ψ_0 and ψ_1 carries the exchange of energy between mass energy and wave energy.

The second spinor formed by ψ_2 and ψ_3 carries spatial and temporal pulse energy exchanges, which is confirmed by the presence of the expression $(\hbar c k_r)$. This impulse energy can be firstly assigned to each of the variables of space r , θ , φ .

The movement of rotation which is highlighted in exact solutions takes place around the Oz axis: it induces an impulse energy along φ .

The variable space r plays a role analogous to variables space x , y and z in Cartesian coordinates: it induces an impulse energy following r .

The θ variable appears to play no role in impulse terms, and we will make the assumption that the impulse energy associated with this variable is zero.

The wave function ψ_3 cannot be associated with a rotating motion along φ : let us assume that it is relative to r impulse energy. We will therefore associate the ψ_2 wave function to impulse energy of rotation. This hypothesis is supported by the fact that the ψ_2 wave function changes sign when passing from a spin "up" for a spin "down". It is also confirmed by the spatial distribution of energy depending on θ , as we shall see later.

When performing a rotation by an angle $\varphi = 2\pi$, the bispinor (XX-1) found its initial position, which suggests that we can associate this mode with a particle of spin 1. We know that in this case the kinetic angular momentum has the maximum chance of making an angle $\theta = 45^\circ$ and $\theta = 135^\circ$ with the Oz axis.

We can link this property to the pulse rotational energy carried by the ψ_2 wave function:

$$\psi_2 = j(\hbar c k_r) \cos \theta \sin \theta e^{-j\varphi} \sin(k_t x_t) \left\{ -f_1'(k_r, r) + \frac{f_1(k_r, r)}{k_r r} \right\} \tag{XX-4}$$

From the point of view of the volumetric energy density, the θ dependence is given by the function:

$$A(\theta) = (\cos \theta \sin \theta)^2 \quad (\text{XX-5})$$

This function is shown in figure (XX-1) below:

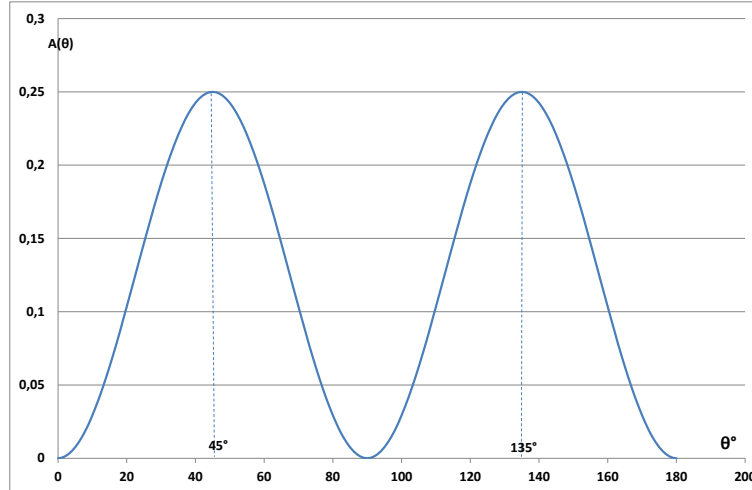


Figure (XX-1): Representation of the distribution range of volumetric density of impulse energy versus θ for the Y_{11} or Y_{1-1} mode likely to represent a particle of spin 1

The rotational impulse energy density shows a maximum in the $\theta = 45^\circ$ and $\theta = 135^\circ$ directions. It is recalled that in an energy and deterministic approach, distribution around these values no longer represents a probability density, but an energy volume density.

Impulse energy along r is carried by the wave ψ_3 function:

$$\psi_3 = j(\hbar c k_r) \sin(k_t x_t) \left\{ -\sin^2 \theta f_1'(k_r r) - (\cos^2 \theta + 1) \frac{f_1(k_r r)}{k_r r} \right\} \quad (\text{XX-6})$$

None of the terms separately seems to be subject to a simple physical interpretation.

The photon is an important special case of particle of spin 1. It's a particle whose rest mass is null and in this energy approach, wave functions which describe this particle obey the following relationships to a multiplicative constant close:

$$\begin{aligned} \psi_0 &= (\hbar \omega) f_1(k_r r) \sin \theta e^{-j\phi} \cos(k_t x_t) \\ \psi_1 &= 0 \\ \psi_2 &= (\hbar c k_r) \cos \theta \sin \theta e^{-j\phi} \sin(k_t x_t) \left\{ -f_1'(k_r r) + \frac{f_1(k_r r)}{k_r r} \right\} \\ \psi_3 &= (\hbar c k_r) \sin(k_t x_t) \left\{ -\sin^2 \theta f_1'(k_r r) - (\cos^2 \theta + 1) \frac{f_1(k_r r)}{k_r r} \right\} \end{aligned} \quad (\text{XX-7})$$

The energy exchanges occur between the wave energy brought by the first spinor, and impulse energy carried by the second spinor: when the first is maximum ($|\cos(k_t x_t)| = 1$), the second is null ($\sin(k_t x_t) = 0$) and vice versa.

In the energy approach, the mass energy of the photon being zero, it is impossible to know its position. During its meeting with other particles, it can only share pulse energy.

II – Particle of spin 1/2

We know that the DIRAC equation is the equation whose solutions are the closest to the physical behavior of the electron. We must therefore find in exact solutions to this equation in spherical coordinates, some key properties highlighting both theoretical and experimental behavior of this particle.

The solution which appears to have the greatest analogy is one based on mode Y_{22} or $Y_{2,-2}$:

$$\begin{aligned}\psi_0 &= \eta f_2(k_r r) \sin^2 \theta e^{-2j\varphi} \sin(k_t x_t) + j k_t f_2(k_r r) \sin^2 \theta e^{-2j\varphi} \cos(k_t x_t) \\ \psi_1 &= 0 \\ \psi_2 &= j k_r \cos \theta \sin^2 \theta e^{-2j\varphi} \sin(k_t x_t) \left\{ -f_2'(k_r r) + 2 \frac{f_2(k_r r)}{k_r r} \right\} \\ \psi_3 &= j k_r \sin(k_t x_t) \sin \theta e^{-j\varphi} \left\{ -\sin^2 \theta f_2'(k_r r) - 2(\cos^2 \theta + 1) \frac{f_2(k_r r)}{k_r r} \right\}\end{aligned}\tag{XX-8}$$

where f_2 is a spherical BESSEL function of order 2:

$$j_2(k_r r) = \left(\frac{3}{(k_r r)^3} - \frac{1}{(k_r r)} \right) \sin(k_r r) - 3 \frac{\cos(k_r r)}{(k_r r)^2}\tag{XX-9}$$

$$y_2(k_r r) = - \left(\frac{3}{(k_r r)^3} - \frac{1}{(k_r r)} \right) \cos(k_r r) - 3 \frac{\sin(k_r r)}{(k_r r)^2}\tag{XX-10}$$

and where we adopted the notation $f_2'(k_r r) = df_2(k_r r)/d(k_r r)$.

After multiplication by the constant $\hbar c$, each of the wave functions has the dimension of energy, and we will do as previously, abstraction of the constant of standardization in the analysis of this solution:

$$\begin{aligned}
 \psi_0 &= (m_0 c^2) f_2(k_r r) \sin^2 \theta e^{-2j\varphi} \sin(k_t x_t) + j(\hbar\omega) f_2(k_r r) \sin^2 \theta e^{-2j\varphi} \cos(k_t x_t) \\
 \psi_1 &= 0 \\
 \psi_2 &= j(\hbar c k_r) \cos \theta \sin^2 \theta e^{-2j\varphi} \sin(k_t x_t) \left\{ -f_2'(k_r r) + 2 \frac{f_2(k_r r)}{k_r r} \right\} \\
 \psi_3 &= j(\hbar c k_r) \sin(k_t x_t) \sin \theta e^{-j\varphi} \left\{ -\sin^2 \theta f_2'(k_r r) - 2(\cos^2 \theta + 1) \frac{f_2(k_r r)}{k_r r} \right\}
 \end{aligned} \tag{XX-11}$$

The first important element concerns the rotation angle φ : if the wave function ψ_3 rotates of 2π , then ψ_0 and ψ_2 wave functions are rotated 4π .

The spin $1/2$ of the electron is often presented by explaining that the bi-spinor must rotate 4π before returning to its original position. It appears in the solution (XX-8) new elements which indicate that all components of a spinor rotation do not necessarily vary with the same angular range.

Seen under this aspect, it is necessary that ψ_0 and ψ_2 components turn of 4π so that the component ψ_3 returned to its initial state.

The second element which suggests that the (XX-8) solution may characterize the inner workings of the electron is contained within the ψ_2 wave function:

$$\psi_2 = j(\hbar c k_r) \cos \theta \sin^2 \theta e^{-2j\varphi} \sin(k_t x_t) \left\{ -f_2'(k_r r) + 2 \frac{f_2(k_r r)}{k_r r} \right\} \tag{XX-12}$$

This wave function carries impulse energy along φ . Since the θ direction is separable, we can find out in which direction this energy is maximum along φ , hence we have to solve:

$$\frac{d}{d\theta} (\cos \theta \sin^2 \theta) = -\sin^3 \theta + 2 \sin \theta \cos^2 \theta = 0 \tag{XX-13}$$

what gives an angle θ_{\max} such as:

$$\text{tg}(\theta_{\max}) = \sqrt{2} \tag{XX-14}$$

And then:

$$\theta_{\max} = 54,73^\circ \tag{XX-15}$$

From the point of view of the volumetric energy density, the θ dependence is given by the function:

$$A(\theta) = (\cos \theta \sin^2 \theta)^2 \tag{XX-16}$$

This function is shown in figure (XX-2) below:

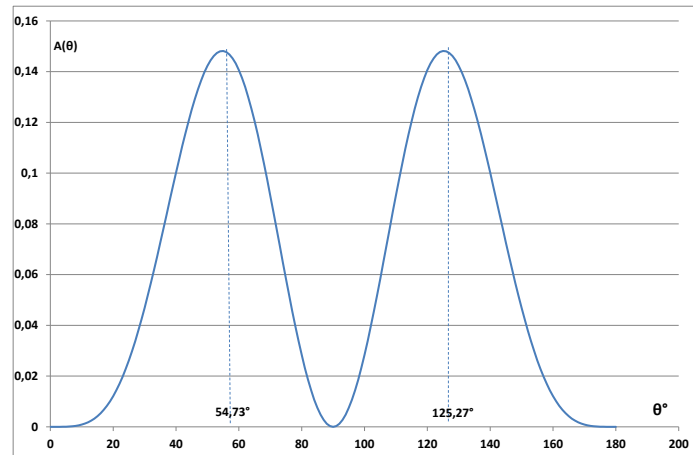


Figure (XX-2): Representation of the distribution of volumetric density of impulse energy versus θ for Y_{22} or Y_{2-2} modes likely to represent a particle of spin $1/2$

It appears that these directions correspond to the directions of the angular momentum of the electron.

III – Spin particle « $1/n$ »

One can search in the exact solutions to the DIRAC equation on modes Y_{nn} or Y_{n-n} , the general expression of the direction in which the angular momentum of spin is maximum.

Impulse energy of spin is carried by the wave function ψ_2 who is remembered for memory:

$$\psi_2 = j \cos \theta \sin^n \theta e^{-nj\phi} \sin(k_t x_t) \left\{ -k_r f_n'(k_r r) + n \frac{f_n(k_r r)}{r} \right\} \quad (\text{XX-17})$$

Since this expression is separable in θ , one can determine the θ_{\max} value for which this impulse energy is maximum. This is equivalent to solving the equation:

$$\frac{d}{d\theta} (\cos \theta \sin^n \theta) = -\sin^{n+1} \theta + n \sin^{n-1} \theta \cos^2 \theta = 0 \quad (\text{XX-18})$$

What gives an angle θ_{\max} such:

$$\text{tg}^2(\theta_{\max}) = n \quad (\text{XX-19})$$

It turns out that this direction is identical to that provided by the quantization of angular momentum of spin in classical mechanics for a spin in “ $1/n$ ”. The quotation marks mean that we enter here in a field for which there is, to the knowledge of the author, no particle known to date with this property for n different from 1 and 2.

To show it, we adopt a quantization of angular momentum of spin along the Oz axis in the form:

$$S_z = \frac{1}{n} \hbar \quad (\text{XX-20})$$

Standard angular momentum S in quantum mechanics is written:

$$S = \sqrt{\frac{1}{n} \left(\frac{1}{n} + 1 \right)} \hbar \quad (\text{XX-21})$$

And the angle θ_{\max} is given by the relationship:

$$\cos(\theta_{\max}) = \frac{S_z}{S} = \frac{\frac{1}{n}}{\sqrt{\frac{1}{n} \left(\frac{1}{n} + 1 \right)}} = \frac{1}{\sqrt{n+1}} \quad (\text{XX-22})$$

Equality with the angle (XX-19) provided by the energy approach and the angle (XX-22) provided by classical quantization of the spin angular momentum is given by the trigonometric relationship:

$$\frac{1}{\cos^2 \theta} = \text{tg}^2 \theta + 1 \quad (\text{XX-23})$$

XXI

Spherical DIRAC currents

Exact solutions to the DIRAC equation can only represent physical solutions if they comply with the local conservation of energy, which the mathematical translation is carried out by expressing that the four-divergence of volume energy density must be null.

$$\frac{\partial}{\partial x^\mu} J^\mu = \partial_\mu J^\mu = 0 \quad (\text{XXI-1})$$

We have seen in Cartesian coordinates that one could extract from DIRAC equations the following relationship:

$$\frac{\partial}{\partial x^\mu} (\bar{\psi} \gamma^\mu \psi) = 0 \quad (\text{XXI-2})$$

In which:

$$\bar{\psi} = (\psi^*)^T \gamma^0 = (\psi_0^*, \psi_1^*, \psi_2^*, \psi_3^*) \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix} = (\psi_0^*, \psi_1^*, -\psi_2^*, -\psi_3^*) \quad (\text{XXI-3})$$

By identifying (XXI-2) with (XXI-1), we infer the four-vector current of DIRAC:

$$J^\mu = \bar{\psi} \gamma^\mu \psi \quad (\text{XXI-4})$$

where γ^μ are the DIRAC matrices.

To establish expressions of these currents in spherical coordinates, we must use the DIRAC matrices obtained in this coordinate system. These matrices are index by the letters t, r, θ, ϕ which represent differential to which they apply.

$$\gamma_t = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix} \quad (\text{XXI-5})$$

$$\gamma_r = \begin{pmatrix} 0 & 0 & \cos \theta & \sin \theta e^{-j\varphi} \\ 0 & 0 & \sin \theta e^{j\varphi} & -\cos \theta \\ -\cos \theta & -\sin \theta e^{-j\varphi} & 0 & 0 \\ -\sin \theta e^{j\varphi} & \cos \theta & 0 & 0 \end{pmatrix} \quad (\text{XXI-6})$$

$$\gamma_\theta = \begin{pmatrix} 0 & 0 & -\sin \theta & \cos \theta e^{-j\varphi} \\ 0 & 0 & \cos \theta e^{j\varphi} & \sin \theta \\ \sin \theta & -\cos \theta e^{-j\varphi} & 0 & 0 \\ -\cos \theta e^{j\varphi} & -\sin \theta & 0 & 0 \end{pmatrix} \quad (\text{XXI-7})$$

$$\gamma_\varphi = \begin{pmatrix} 0 & 0 & 0 & -je^{-j\varphi} \\ 0 & 0 & je^{j\varphi} & 0 \\ 0 & je^{-j\varphi} & 0 & 0 \\ -je^{j\varphi} & 0 & 0 & 0 \end{pmatrix} \quad (\text{XXI-8})$$

Details of the analytical expressions of these currents are shown below:

index component t:

$$J^t = \bar{\Psi} \gamma_0 \Psi = (\Psi_0^*, \Psi_1^*, -\Psi_2^*, -\Psi_3^*) \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix} \begin{pmatrix} \Psi_0 \\ \Psi_1 \\ \Psi_2 \\ \Psi_3 \end{pmatrix} = (\Psi_0^*, \Psi_1^*, \Psi_2^*, \Psi_3^*) \begin{pmatrix} \Psi_0 \\ \Psi_1 \\ \Psi_2 \\ \Psi_3 \end{pmatrix}$$

$$J^t = \Psi_0^* \Psi_0 + \Psi_1^* \Psi_1 + \Psi_2^* \Psi_2 + \Psi_3^* \Psi_3 \quad (\text{XXI-9})$$

index component r:

$$\mathbf{J}^r = \bar{\Psi} \gamma_r \Psi = (\psi_0^*, \psi_1^*, -\psi_2^*, -\psi_3^*) \begin{pmatrix} 0 & 0 & \cos \theta & \sin \theta e^{-j\varphi} \\ 0 & 0 & \sin \theta e^{j\varphi} & -\cos \theta \\ -\cos \theta & -\sin \theta e^{-j\varphi} & 0 & 0 \\ -\sin \theta e^{j\varphi} & \cos \theta & 0 & 0 \end{pmatrix} \begin{pmatrix} \Psi_0 \\ \Psi_1 \\ \Psi_2 \\ \Psi_3 \end{pmatrix}$$

$$= (\psi_2^* \cos \theta + \psi_3^* \sin \theta e^{j\varphi}, \psi_2^* \sin \theta e^{-j\varphi} - \psi_3^* \cos \theta, \psi_0^* \cos \theta + \psi_1^* \sin \theta e^{j\varphi}, \psi_0^* \sin \theta e^{-j\varphi} - \psi_1^* \cos \theta) \begin{pmatrix} \Psi_0 \\ \Psi_1 \\ \Psi_2 \\ \Psi_3 \end{pmatrix}$$

$$\mathbf{J}^r = \psi_0 (\psi_2^* \cos \theta + \psi_3^* \sin \theta e^{j\varphi}) + \psi_1 (\psi_2^* \sin \theta e^{-j\varphi} - \psi_3^* \cos \theta) + \psi_2 (\psi_0^* \cos \theta + \psi_1^* \sin \theta e^{j\varphi}) + \psi_3 (\psi_0^* \sin \theta e^{-j\varphi} - \psi_1^* \cos \theta)$$

(XXI-10)

index component θ :

$$\mathbf{J}^\theta = \bar{\Psi} \frac{\gamma_\theta}{r} \Psi = (\psi_0^*, \psi_1^*, -\psi_2^*, -\psi_3^*) \frac{1}{r} \begin{pmatrix} 0 & 0 & -\sin \theta & \cos \theta e^{-j\varphi} \\ 0 & 0 & \cos \theta e^{j\varphi} & \sin \theta \\ \sin \theta & -\cos \theta e^{-j\varphi} & 0 & 0 \\ -\cos \theta e^{j\varphi} & -\sin \theta & 0 & 0 \end{pmatrix} \begin{pmatrix} \Psi_0 \\ \Psi_1 \\ \Psi_2 \\ \Psi_3 \end{pmatrix}$$

$$= \frac{1}{r} (-\psi_2^* \sin \theta + \psi_3^* \cos \theta e^{j\varphi}, \psi_2^* \cos \theta e^{-j\varphi} + \psi_3^* \sin \theta, -\psi_0^* \sin \theta + \psi_1^* \cos \theta e^{j\varphi}, \psi_0^* \cos \theta e^{-j\varphi} + \psi_1^* \sin \theta) \begin{pmatrix} \Psi_0 \\ \Psi_1 \\ \Psi_2 \\ \Psi_3 \end{pmatrix}$$

$$r\mathbf{J}^\theta = \psi_0 (-\psi_2^* \sin \theta + \psi_3^* \cos \theta e^{j\varphi}) + \psi_1 (\psi_2^* \cos \theta e^{-j\varphi} + \psi_3^* \sin \theta) + \psi_2 (-\psi_0^* \sin \theta + \psi_1^* \cos \theta e^{j\varphi}) + \psi_3 (\psi_0^* \cos \theta e^{-j\varphi} + \psi_1^* \sin \theta)$$

(XXI-11)

index component φ :

$$\mathbf{J}^\varphi = \bar{\Psi} \frac{\gamma_\varphi}{r \sin \theta} \Psi = (\psi_0^*, \psi_1^*, -\psi_2^*, -\psi_3^*) \frac{1}{r \sin \theta} \begin{pmatrix} 0 & 0 & 0 & -j e^{-j\varphi} \\ 0 & 0 & j e^{j\varphi} & 0 \\ 0 & j e^{-j\varphi} & 0 & 0 \\ -j e^{j\varphi} & 0 & 0 & 0 \end{pmatrix} \begin{pmatrix} \Psi_0 \\ \Psi_1 \\ \Psi_2 \\ \Psi_3 \end{pmatrix}$$

$$\mathbf{J}^\varphi = \frac{j}{r \sin \theta} (\psi_3^* e^{j\varphi}, -\psi_2^* e^{-j\varphi}, \psi_1^* e^{j\varphi}, -\psi_0^* e^{-j\varphi}) \begin{pmatrix} \Psi_0 \\ \Psi_1 \\ \Psi_2 \\ \Psi_3 \end{pmatrix}$$

(XXI-12)

$$r \sin \theta \mathbf{J}^\varphi = j (\psi_0 \psi_3^* e^{j\varphi} - \psi_1 \psi_2^* e^{-j\varphi} + \psi_2 \psi_1^* e^{j\varphi} - \psi_3 \psi_0^* e^{-j\varphi})$$

These analytical expressions allow to check how exact solutions obtained comply with the local conservation of energy. The calculations may be particularly laborious. Some elements of calculations are presented below, on the example of the rotating solution obtained in a previous chapter:

$$\begin{aligned}
 \psi_0 &= 0 \\
 \psi_1 &= (\eta - k_t) f_n(k_r r) \sin^n \theta \exp j(k_t x_t + n\varphi) \\
 \psi_2 &= j \exp j(k_t x_t + (n-1)\varphi) \sin^{n-1} \theta \left\{ -\sin^2 \theta k_r f_n'(k_r r) - n(\cos^2 \theta + 1) \frac{f_n(k_r r)}{r} \right\} \\
 \psi_3 &= j \cos \theta \sin^n \theta \exp j(k_t x_t + n\varphi) \left\{ k_r f_n'(k_r r) - n \frac{f_n(k_r r)}{r} \right\}
 \end{aligned} \tag{XXI-13}$$

Calculation of J^t :

$$\begin{aligned}
 J^t &= \psi_0^* \psi_0 + \psi_1^* \psi_1 + \psi_2^* \psi_2 + \psi_3^* \psi_3 \\
 J^t &= (\eta - k_t)^2 \left\{ f_n(k_r r) \sin^n \theta \right\}^2 + k_r^2 \left\{ -\sin^{n+1} \theta f_n'(k_r r) - n \sin^{n-1} \theta (\cos^2 \theta + 1) \frac{f_n(k_r r)}{k_r r} \right\}^2 \\
 &+ (\cos \theta \sin^n \theta)^2 k_r^2 \left\{ f_n'(k_r r) - n \frac{f_n(k_r r)}{k_r r} \right\}^2
 \end{aligned} \tag{XXI-14}$$

On this particular mode, the volume density of total energy depends neither time nor of the space variable φ

Calculation of J^r :

$$\begin{aligned}
 J^r &= \psi_0^* (\psi_2^* \cos \theta + \psi_3^* \sin \theta e^{j\varphi}) + \psi_1^* (\psi_2^* \sin \theta e^{-j\varphi} - \psi_3^* \cos \theta) \\
 &+ \psi_2^* (\psi_0^* \cos \theta + \psi_1^* \sin \theta e^{j\varphi}) + \psi_3^* (\psi_0^* \sin \theta e^{-j\varphi} - \psi_1^* \cos \theta)
 \end{aligned} \tag{XXI-15}$$

For the reporting solution, we have $\psi_0 = 0$:

$$J^r = \psi_1^* (\psi_2^* \sin \theta e^{-j\varphi} - \psi_3^* \cos \theta) + \psi_2^* (\psi_1^* \sin \theta e^{j\varphi}) + \psi_3^* (-\psi_1^* \cos \theta) \tag{XXI-16}$$

The exponential terms vanish by conjugation in the products of the wave functions. After this cancellation, it remains in ψ_2 and ψ_3 purely imaginary terms. We can deduce:

$$J^r = 0 \tag{XXI-17}$$

Calculation of J^θ :

$$\begin{aligned}
 rJ^\theta &= \psi_0^* (-\psi_2^* \sin \theta + \psi_3^* \cos \theta e^{j\varphi}) + \psi_1^* (\psi_2^* \cos \theta e^{-j\varphi} + \psi_3^* \sin \theta) \\
 &+ \psi_2^* (-\psi_0^* \sin \theta + \psi_1^* \cos \theta e^{j\varphi}) + \psi_3^* (\psi_0^* \cos \theta e^{-j\varphi} + \psi_1^* \sin \theta)
 \end{aligned} \tag{XXI-18}$$

For the reporting solution, we have $\psi_0 = 0$:

$$rJ^0 = \psi_1(\psi_2^* \cos \theta e^{-j\varphi} + \psi_3^* \sin \theta) + \psi_2(\psi_1^* \cos \theta e^{j\varphi}) + \psi_3(\psi_1^* \sin \theta) \quad (\text{XXI-19})$$

The exponential terms vanish by conjugation in the products of the wave functions. After this cancellation, it remains in ψ_2 and ψ_3 purely imaginary terms. We can deduce:

$$J^0 = 0 \quad (\text{XXI-20})$$

Calculation of J^φ :

$$r \sin \theta J^\varphi = j(\psi_0 \psi_3^* e^{j\varphi} - \psi_1 \psi_2^* e^{-j\varphi} + \psi_2 \psi_1^* e^{j\varphi} - \psi_3 \psi_0^* e^{-j\varphi}) \quad (\text{XXI-21})$$

For the reporting solution, we have $\psi_0 = 0$:

$$r \sin \theta J^\varphi = j(-\psi_1 \psi_2^* e^{-j\varphi} + \psi_2 \psi_1^* e^{j\varphi}) \quad (\text{XXI-22})$$

We obtain:

$$J^\varphi = 2(k_t - \eta) \frac{f_n(k_r r)}{r} \sin^{2n-2} \theta \left\{ \sin^2 \theta k_r f_n'(k_r r) + n(\cos^2 \theta + 1) \frac{f_n(k_r r)}{r} \right\} \quad (\text{XXI-23})$$

Since we are in possession of the currents of DIRAC, we can now check if these currents that can be associated to the evolution of the density of energy in time and space, well check the relationship of local conservation of energy.

This relationship is written in spherical coordinates:

$$\frac{\partial J^t}{\partial x_t} + \frac{1}{r^2} \frac{\partial(r^2 J^r)}{\partial r} + \frac{1}{r \sin \theta} \frac{\partial(\sin \theta J^\theta)}{\partial \theta} + \frac{1}{r \sin \theta} \frac{\partial J^\varphi}{\partial \varphi} = 0 \quad (\text{XXI-24})$$

It is verified, because each of the terms of this divergence, taken separately, is zero. A somewhat more detailed analysis shows that about these modes in rotation:

- Total volume energy density J^t depends neither on time, nor on variable spatial φ . On the other hand, it varies following r and θ , and is depending on the modes that are excited.
- Currents of DIRAC following r and θ being zero, there is no flow of energy following these directions.
- The Dirac current along φ is non-zero, reflecting the existence of a flow of energy linked with the angular rotation of energy in that angular direction. This flow of energy is uniform because it does not depend on φ . It can be attached to the rotation, with a constant speed, of energy following this direction.

XXII

Generalized DIRAC equation in spherical coordinates

Obtaining solutions to the Dirac equation in the form of spherical modes leads naturally to wonder about the existence of such solutions when the particle is immersed in a constant and uniform four-potential.

This problem has already been processed in Cartesian coordinates in chapter XII. We propose to resume it in spherical coordinates in a somewhat different approach.

We must first determine the equation of conservation of energy between two frames linked by relativity.

We adopt as a starting point a Cartesian coordinate system and the form of the KLEIN-GORDON equation invariant under the LORENTZ transformation:

$$\left\{ \left(j \frac{\partial}{\partial x^\mu} - \frac{q}{\hbar} A_\mu \right) \left(j \frac{\partial}{\partial x^\mu} - \frac{q}{\hbar} A_\mu \right) \right\} (\psi) = \left(\frac{m_0 c}{\hbar} \right)^2 (\psi) \quad (\text{XXII-1})$$

In this relationship, the A_μ represent the components of the four-potential (ϕ/c , A_x , A_y , A_z). In order to lighten writing, we put, as in chapter XII:

$$\eta_t = \frac{q\phi}{\hbar c} \quad \eta_x = \frac{qA_x}{\hbar} \quad \eta_y = \frac{qA_y}{\hbar} \quad \eta_z = \frac{qA_z}{\hbar} \quad \eta = \frac{m_0 c}{\hbar} \quad (\text{XXII-2})$$

The development of the equation (XXII-1) in a metric (+,-,-,-) gives us:

$$\left\{ \left(-\frac{\partial^2}{\partial x_t^2} - 2j\eta_t \frac{\partial}{\partial x_t} + \eta_t^2 \right) - \left(-\frac{\partial^2}{\partial x^2} - 2j\eta_x \frac{\partial}{\partial x} + \eta_x^2 \right) \right\} \left\{ -\left(-\frac{\partial^2}{\partial y^2} - 2j\eta_y \frac{\partial}{\partial y} + \eta_y^2 \right) - \left(-\frac{\partial^2}{\partial z^2} - 2j\eta_z \frac{\partial}{\partial z} + \eta_z^2 \right) \right\} (\psi) = \eta^2 (\psi) \quad (\text{XXII-3})$$

Or still :

$$\left(-\frac{\partial^2}{\partial x_t^2} - 2j\eta_t \frac{\partial}{\partial x_t} + \eta_t^2\right)(\psi) = \left\{ \begin{array}{l} \left(-\frac{\partial^2}{\partial x^2} - 2j\eta_x \frac{\partial}{\partial x} + \eta_x^2\right) + \\ \left(-\frac{\partial^2}{\partial y^2} - 2j\eta_y \frac{\partial}{\partial y} + \eta_y^2\right) + \\ \left(-\frac{\partial^2}{\partial z^2} - 2j\eta_z \frac{\partial}{\partial z} + \eta_z^2\right) \end{array} \right\} (\psi) + \eta^2(\psi) \quad (\text{XXII-4})$$

The scalar potential affects the term on the left of the equal sign, whereas the vector potential affects the right term. We will consider these two cases separately.

I – The scalar potential

We are interested in this part to the effect of the scalar potential, and we are working with a null vector potential, which leads to put $\eta_x = \eta_y = \eta_z = 0$ in the equation (XXII-4):

$$\left(-\frac{\partial^2}{\partial x_t^2} - 2j\eta_t \frac{\partial}{\partial x_t} + \eta_t^2\right)(\psi) = -\left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2}\right)(\psi) + \eta^2(\psi) \quad (\text{XXII-5})$$

The introduction of spherical coordinates provides the Laplacian in the form:

$$\left(-\frac{\partial^2}{\partial x_t^2} - 2j\eta_t \frac{\partial}{\partial x_t} + \eta_t^2\right)(\psi) = -\left\{ \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2}{\partial \varphi^2} \right\} (\psi) + \eta^2(\psi) \quad (\text{XXII-6})$$

It appears from the analysis of spherical modes of Chapter XVI that stationary modes should check the relationship:

$$\left(-\frac{\partial^2}{\partial x_t^2} - 2j\eta_t \frac{\partial}{\partial x_t} + \eta_t^2\right)(\psi) = (k_r^2 + \eta^2)(\psi) \quad (\text{XXII-7})$$

It appears a note already formulated in another way in chapter XII on the basis of observations relating to the linear system XII-11, 12, 13, 14: it is impossible to find a stationary solution in $\sin(k_t x_t)$ or $\cos(k_t x_t)$ that allows to obtain an equation of conservation of energy derived from (XXII-7).

Possible solutions appear in the form:

$$\cos(k_t x_t) \pm j \sin(k_t x_t) = \exp(\pm j k_t x_t) \quad (\text{XXII-8})$$

They reflect the excitement of modes in quadrature which exchange positive and negative energy.

By postponing this relationship (XXII-8) in (XXII-7), the relationship of conservation of energy which must be verified by the system of Dirac in spherical coordinates in a scalar potential is obtained:

$$(k_t^2 \pm 2k_t \eta_t + \eta_t^2) = (k_r^2 + \eta^2) \quad (\text{XXII-9})$$

Or still:

$$(k_t \pm \eta_t)^2 = (k_r^2 + \eta^2) \quad (\text{XXII-10})$$

It's the relation obtained in chapter XII, transposed in spherical coordinates.

The system of Dirac generalized to a scalar potential in spherical coordinates is written:

$$\begin{aligned} (\eta + \eta_t) \psi_0 &= j \left\{ \frac{\partial \psi_0}{\partial x_t} + \cos \theta \frac{\partial \psi_2}{\partial r} + \sin \theta e^{-j\varphi} \frac{\partial \psi_3}{\partial r} + \frac{1}{r} \left(-\sin \theta \frac{\partial \psi_2}{\partial \theta} + \cos \theta e^{-j\varphi} \frac{\partial \psi_3}{\partial \theta} \right) - \frac{j e^{-j\varphi}}{r \sin \theta} \frac{\partial \psi_3}{\partial \varphi} \right\} \\ (\eta + \eta_t) \psi_1 &= j \left\{ \frac{\partial \psi_1}{\partial x_t} + \sin \theta e^{j\varphi} \frac{\partial \psi_2}{\partial r} - \cos \theta \frac{\partial \psi_3}{\partial r} + \frac{1}{r} \left(\cos \theta e^{j\varphi} \frac{\partial \psi_2}{\partial \theta} + \sin \theta \frac{\partial \psi_3}{\partial \theta} \right) + \frac{j e^{j\varphi}}{r \sin \theta} \frac{\partial \psi_2}{\partial \varphi} \right\} \\ (\eta - \eta_t) \psi_2 &= j \left\{ -\frac{\partial \psi_2}{\partial x_t} - \cos \theta \frac{\partial \psi_0}{\partial r} - \sin \theta e^{-j\varphi} \frac{\partial \psi_1}{\partial r} + \frac{1}{r} \left(\sin \theta \frac{\partial \psi_0}{\partial \theta} - \cos \theta e^{-j\varphi} \frac{\partial \psi_1}{\partial \theta} \right) + \frac{j e^{-j\varphi}}{r \sin \theta} \frac{\partial \psi_1}{\partial \varphi} \right\} \\ (\eta - \eta_t) \psi_3 &= j \left\{ -\frac{\partial \psi_3}{\partial x_t} - \sin \theta e^{j\varphi} \frac{\partial \psi_0}{\partial r} + \cos \theta \frac{\partial \psi_1}{\partial r} + \frac{1}{r} \left(-\cos \theta e^{j\varphi} \frac{\partial \psi_0}{\partial \theta} - \sin \theta \frac{\partial \psi_1}{\partial \theta} \right) - \frac{j e^{j\varphi}}{r \sin \theta} \frac{\partial \psi_0}{\partial \varphi} \right\} \end{aligned} \quad (\text{XXII-11})$$

Using the exact solution (XXI-13) for example, the exact solution in a scalar potential is obtained in the form

$$\begin{aligned} \psi_0 &= 0 \\ \psi_1 &= (\eta - k_t - \eta_t) f_n(k_r r) \sin^n \theta \exp j(k_t x_t + n\varphi) \\ \psi_2 &= -j \exp j(k_t x_t + (n-1)\varphi) \sin^{n-1} \theta \left\{ \sin^2 \theta k_r f_n'(k_r r) + n(\cos^2 \theta + 1) \frac{f_n(k_r r)}{r} \right\} \\ \psi_3 &= j \cos \theta \sin^n \theta \exp j(k_t x_t + n\varphi) \left\{ k_r f_n'(k_r r) - n \frac{f_n(k_r r)}{r} \right\} \end{aligned} \quad (\text{XXII-12})$$

It is associated with the relation of energy conservation:

$$(k_t + \eta_t)^2 = (k_r^2 + \eta^2) \quad (\text{XXII-13})$$

II – The vector potential

From (XXII-4), we deduce the KLEIN-GORDON equation written in Cartesian coordinates for a scalar potential equal to zero and a constant vector potential:

$$-\frac{\partial^2}{\partial x_i^2}(\psi) = \left\{ \left(-\frac{\partial^2}{\partial x^2} - 2j\eta_x \frac{\partial}{\partial x} + \eta_x^2 \right) + \left(-\frac{\partial^2}{\partial y^2} - 2j\eta_y \frac{\partial}{\partial y} + \eta_y^2 \right) + \left(-\frac{\partial^2}{\partial z^2} - 2j\eta_z \frac{\partial}{\partial z} + \eta_z^2 \right) + \eta^2 \right\} (\psi) \quad (\text{XXII-14})$$

Or still, by rearranging the terms:

$$-\frac{\partial^2}{\partial x_i^2}(\psi) = \left(-\frac{\partial^2}{\partial x^2} - \frac{\partial^2}{\partial y^2} - \frac{\partial^2}{\partial z^2} - 2j \frac{\partial}{\partial x} \eta_x - 2j \frac{\partial}{\partial y} \eta_y - 2j \frac{\partial}{\partial z} \eta_z + \eta_x^2 + \eta_y^2 + \eta_z^2 + \eta^2 \right) (\psi) \quad (\text{XXII-15})$$

The passage in spherical coordinates is made by substituting the Laplacian and the partial derivatives already encountered in Chapter XV

$$\begin{aligned} \frac{\partial \psi}{\partial x} &= \frac{\partial \psi}{\partial r} \sin \theta \cos \varphi + \frac{\partial \psi}{\partial \theta} \frac{1}{r} \cos \theta \cos \varphi - \frac{\partial \psi}{\partial \varphi} \frac{1}{r \sin \theta} \sin \varphi \\ \frac{\partial \psi}{\partial y} &= \frac{\partial \psi}{\partial r} \sin \theta \sin \varphi + \frac{\partial \psi}{\partial \theta} \frac{1}{r} \cos \theta \sin \varphi + \frac{\partial \psi}{\partial \varphi} \frac{1}{r \sin \theta} \cos \varphi \\ \frac{\partial \psi}{\partial z} &= \frac{\partial \psi}{\partial r} \cos \theta - \frac{\partial \psi}{\partial \theta} \frac{1}{r} \sin \theta \end{aligned} \quad (\text{XXII-16})$$

The vector potential components are transformed following relations:

$$\begin{aligned} \eta_x &= \frac{qA_x}{\hbar} = \frac{qA_r}{\hbar} \sin \theta \cos \varphi = \eta_r \sin \theta \cos \varphi \\ \eta_y &= \frac{qA_y}{\hbar} = \frac{qA_r}{\hbar} \sin \theta \sin \varphi = \eta_r \sin \theta \sin \varphi \\ \eta_z &= \frac{qA_z}{\hbar} = \frac{qA_r}{\hbar} \cos \theta = \eta_r \cos \theta \end{aligned} \quad (\text{XXII-17})$$

We can deduce the KLEIN-GORDON equation in spherical coordinates in a vector potential

$$-\frac{\partial^2}{\partial x_i^2}(\psi) = \left(\begin{aligned} &-\frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial}{\partial r} \right) - \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial}{\partial \theta} \right) - \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2}{\partial \varphi^2} \\ &- 2j \left(\frac{\partial}{\partial r} \sin \theta \cos \varphi + \frac{\partial}{\partial \theta} \frac{1}{r} \cos \theta \cos \varphi - \frac{\partial}{\partial \varphi} \frac{1}{r \sin \theta} \sin \varphi \right) \eta_r \sin \theta \cos \varphi \\ &- 2j \left(\frac{\partial}{\partial r} \sin \theta \sin \varphi + \frac{\partial}{\partial \theta} \frac{1}{r} \cos \theta \sin \varphi + \frac{\partial}{\partial \varphi} \frac{1}{r \sin \theta} \cos \varphi \right) \eta_r \sin \theta \sin \varphi \\ &- 2j \left(\frac{\partial}{\partial r} \cos \theta - \frac{\partial}{\partial \theta} \frac{1}{r} \sin \theta \right) \eta_r \cos \theta \\ &+ (\eta_r \sin \theta \cos \varphi)^2 + (\eta_r \sin \theta \sin \varphi)^2 + (\eta_r \cos \theta)^2 + \eta^2 \end{aligned} \right) (\psi) \quad (\text{XXII-18})$$

After reduction of the terms which vanish, we get:

$$-\frac{\partial^2}{\partial x_t^2}(\psi) = \left(\begin{array}{l} -\frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial}{\partial r} \right) - \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial}{\partial \theta} \right) - \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2}{\partial \varphi^2} \\ -2j \frac{\partial}{\partial r} \eta_r + \eta_r^2 + \eta^2 \end{array} \right) (\psi) \quad (\text{XXII-19})$$

If we require that the wave function ψ comes in the form of spherical stationary modes, this leads to the following expression of conservation of energy:

$$k_t^2(\psi) = \left(k_r^2 - 2j \frac{\partial}{\partial r} \eta_r + \eta_r^2 + \eta^2 \right) (\psi) \quad (\text{XXII-20})$$

This relationship is completely analogous to the relationship (XXII-7) obtained for the scalar potential.

It becomes a problem that could not be overcome by the author: dependence in r of the spherical BESSEL functions in the stationary solution r does not make the relationship of conservation of energy (XXII-20) independent of r . It follows that no formulation of an exact solution to the DIRAC equation in spherical coordinates and under a constant vector potential could be formulated.

XXIII

Conclusion of the third part

I – On the conservation of energy

The importance of this conservation was repeatedly recalled in this manuscript. It is of two kinds: one is conservation of energy during one change of inertial frame, the other on local conservation of energy. Both contribute to the coherence of the presented solutions.

I-1 Conservation of energy for change of frame

It is linked to the fact that the pseudo-norm of the pulse energy four vector:

$$\tilde{\mathbf{P}} = \begin{pmatrix} m \frac{dx}{dt} \\ m \frac{dy}{dt} \\ m \frac{dz}{dt} \\ \frac{E}{c} \end{pmatrix} = \begin{pmatrix} mv_x \\ mv_y \\ mv_z \\ \frac{E}{c} \end{pmatrix} = \begin{pmatrix} p_x \\ p_y \\ p_z \\ \frac{E}{c} \end{pmatrix} \quad (\text{XXIII-1})$$

expressed explicitly under the form :

$$E^2 - p^2c^2 = \text{cte} \quad (\text{XXIII-2})$$

is a constant that does not depend on the inertial frame in which it is assessed. This constant is equal to the energy of the rest mass squared, it is inferred that in any inertial frame in which the particle is likely to evolve, the following relationship holds:

$$E^2 = (pc)^2 + (m_0c^2)^2 \quad (\text{XXIII-3})$$

In the frame in which it is at rest, its impulse energy is zero, and its total energy is equal to its mass energy. In a frame where it is no more at rest, its impulse energy takes a finite value that increases its total energy in the report given by the respect for the relationship (XXIII-2) or energy (XXIII-3) above.

When we transpose these concepts to quantum mechanics, one obtains the relationship that translates the same phenomena of energy conservation by changing frame:

$$\hbar^2 c^2 k_t^2 = \hbar^2 c^2 (k_x^2 + k_y^2 + k_z^2) + (m_0 c^2)^2 \quad (\text{XXIII-4})$$

or still after division of the two members of equality by $\hbar^2 c^2$;

$$k_t^2 = (k_x^2 + k_y^2 + k_z^2) + \eta^2 \quad (\text{XXIII-5})$$

When switching from one inertial frame to another Galilean frame, this relationship expresses the fact that during changes in 3 kinds of different energies, the mass energy remains constant, while the impulse energy and wave energy vary in the proportions given by the relationship (XXIII-4)

Because the relativistic DIRAC equation:

$$\left(j\gamma^\mu \frac{\partial}{\partial x^\mu} - \frac{m_0 c}{\hbar} \right) \psi(x^\mu) = 0 \quad \mu = 0, 1, 2, 3 \quad (\text{XXIII-6})$$

is invariant under change of inertial frame, this imposes that solutions that are expressed on the basis of k_t, k_x, k_y, k_z, m_0 , check the equation of conservation of energy (XXIII-4).

II – On the imaginary form of the DIRAC equation solutions

The solutions of the DIRAC equation are represented by a bispinor. It is recognized that each of these spinors represents for one, a particle, for the other, its antiparticle.

Once the particle is moving, it appears that the two spinors are no more independent, but closely correlated in a relationship whose physical meaning escapes to classical quantum mechanics. In particular, the presence of imaginary terms found no satisfactory explanation, if it is accepted that the squared wave functions represent probability densities.

In an energy approach, the squared wave functions are homogeneous to a volumetric energy density. All physics is built around a signed energy representation, which can be positive or negative. For example, the total energy of a satellite in orbit around the Earth is negative, and the negative sign is justified there because he cannot escape the gravity. The analysis presented in the first part also shows that two spinors which are solutions of the DIRAC equation have energy of mass opposite when the particle is at rest. Therefore, if we consider a quantity that represents the square root of this energy, or more exactly the square root of the energy volume density, there is no problem to see the complex number $j = \sqrt{-1}$, which simply indicates that it is a quantity linked with a negative energy. Thus, in the energy exchanges that reflect the solutions in the form of stationary modes, the presence of the complex number $j = \sqrt{-1}$ simply refers to an energy that must be counted negatively.

In an energy approach, DIRAC currents take a clear physical meaning: they are homogeneous to a volumetric energy density, and thus naturally check the local conservation of energy equation.

III – On the notion of negative energy

If we can be satisfied by the previous paragraph that it is possible to use a representation of signed energy, it appears that the concept of negative energy that is used in this energy approach must be explained more because its physical meaning is not clear.

In this document, we found it for the first time in the classic formalism with relations (III-12) and (III-13) that connect two spinors when the particle is at rest. These two relationships are obtained in the case where the time dependence of spinors is in $\exp(-j\omega t)$. They are recalled below:

$$E \begin{pmatrix} \varphi \\ \chi \end{pmatrix} = \begin{pmatrix} mc^2 & 0 \\ 0 & -mc^2 \end{pmatrix} \begin{pmatrix} \varphi \\ \chi \end{pmatrix} \quad (\text{XXIII-7})$$

or still :

$$\begin{aligned} E(\varphi) &= mc^2(\varphi) \\ E(\chi) &= -mc^2(\chi) \end{aligned} \quad (\text{XXIII-8})$$

The minus sign which appears with the energy associated with the second spinor leads to think that this spinor describes the behavior of the antiparticle. However, the experimental behavior of the antiparticle of the electron shows that its mass energy is identical to that of the electron: its load alone has changed sign.

Then, the sign - which is present before the mass energy seems not to be able to be associated with the description of the mass energy of a particle of antimatter.

Consequently, arises clearly, the physical meaning to be given to negative energy in this energy approach.

We propose to develop an interpretation based on a simple, concrete example of classical physics. This interpretation should be considered with caution as long as it has not received additional evidence in his favour on the part of the scientific community.

The example that will be developed concerns electrical energy.

Let us consider a classical RLC circuit. If we load the capacity before the closure of this circuit, the energy stored by the capacity performs round-trips between capacity and self. At each round trip, a fraction of the energy is dissipated in the resistance R until the total disappearance of energy and thus oscillations.

One ways for the continuation of the oscillations is to include in this circuit a negative resistance (-R). This resistance can be constructed using an operational amplifier that delivers a voltage proportional to the intensity in a fraction (-R): it is therefore a very practical device.

The power P_R dissipated in the resistance R at each time t is simply written:

$$P_R(t) = R I^2(t) \quad (\text{XXIII-9})$$

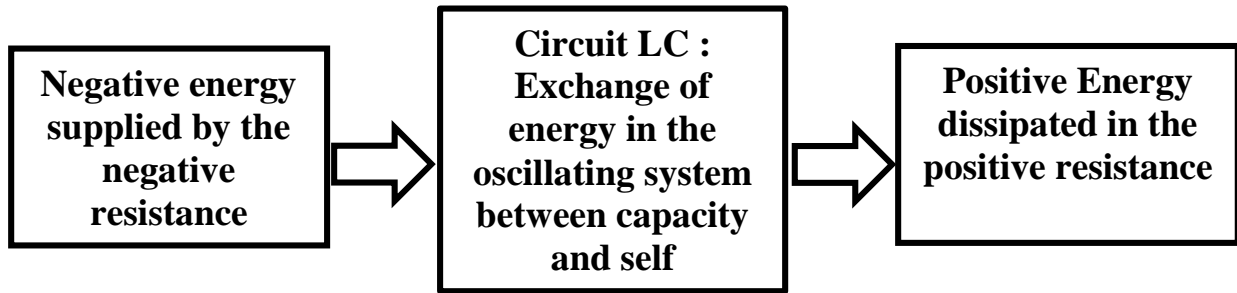
and it is apparent that this power is counted positively.

Power $P_{(-R)}$ provided by the negative resistance at each time t is written thus:

$$P_{(-R)}(t) = -R I^2(t) \quad (\text{XXIII-10})$$

and it is apparent that this power is counted negatively.

Such a device can be summarised as follows:



The physical interpretation that can be made is as follows: a positive power expresses a power that disappears from the system that exchanges of energy, while a negative power is a power that is fed into this system. In this simple case, the conservation of energy introduced in the system at any moment implies:

$$P_R(t) + P_{(-R)}(t) = R I^2(t) - R I^2(t) = 0 \quad (\text{XXIII-11})$$

If now, by analogy with the wave functions, we are interested in quantities ψ_+ and ψ_- that represent the square root of power, these quantities are defined by the relationship:

$$\begin{aligned} \psi_+ &= \sqrt{P_R(t)} = \sqrt{R I^2(t)} \\ \psi_- &= \sqrt{P_{(-R)}(t)} = \sqrt{-R I^2(t)} = j\sqrt{R I^2(t)} \end{aligned} \quad (\text{XXIII-12})$$

One can no longer simply express the relationship of conservation of energy above using ψ_+ and ψ_- quantities. If one wants to nevertheless express the sum or the difference between these quantities that are exchanged in the system, one is led to consider expressions of the type:

$$\begin{aligned} \psi &= \psi_+ + \psi_- = \sqrt{P_R(t)} + \sqrt{P_{(-R)}(t)} = \sqrt{R I^2(t)} + j\sqrt{R I^2(t)} \\ \psi^* &= \psi_+ - \psi_- = \sqrt{P_R(t)} - \sqrt{P_{(-R)}(t)} = \sqrt{R I^2(t)} - j\sqrt{R I^2(t)} \end{aligned} \quad (\text{XXIII-13})$$

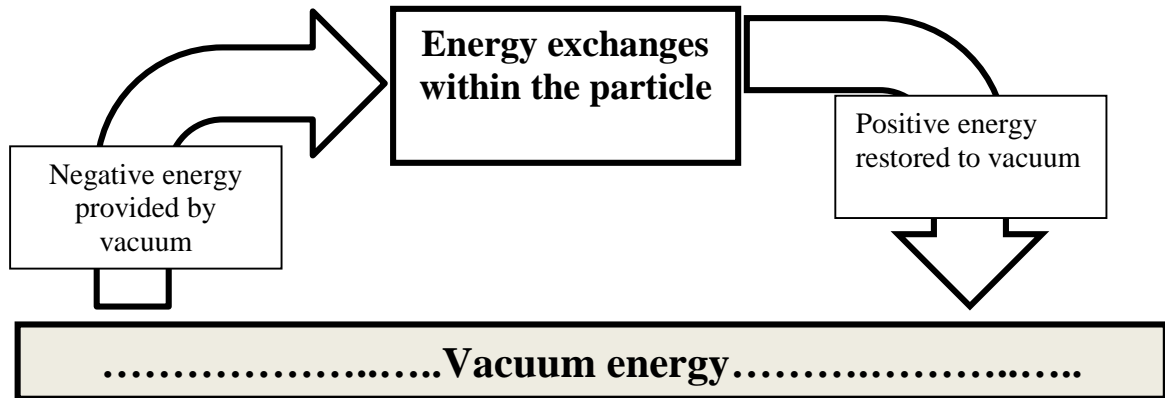
The exchanged power $P(t)$ may be expressed by multiplying the previous imaginary quantity by its conjugate:

$$P(t) = \frac{1}{2} \psi \psi^* = R I^2(t) \quad (\text{XXIII-14})$$

By looking at this expression, we observe that the calculation of this exchanged power is always positive: we get a quantity whose behavior is analogous to Dirac currents.

When one attempt to transpose the energy interpretation above towards quantum energy exchanges, it immediately comes an inevitable question: where is the origin of the energy introduced in the particle (negative energy), and where is going the energy that escapes (positive energy)?

It seems that the only possible response is an exchanged with the energy of the vacuum, whose existence seems to be confirmed. Under this hypothesis, the exchange of energy could be represented schematically as follows:



A scheme of this nature seems consistent with the current knowledge of the quantum fluctuations of the vacuum, and gives a physical meaning clear and unambiguous to the notion of positive energy and negative energy. It expresses the fact that in the quantum world, particles exist and propagate by means of permanent exchange of energy with the energy of the vacuum. We could see in these exchanges support spread allowing photons to propagate in vacuum over distances of several light-years.

IV – general conclusion

Quantum mechanics cannot be satisfied a very long time yet the impasse in which it is maintained by the ignorance of the underlying physical phenomena to the wave-particle duality.

Because the DIRAC equation is the equation that describes the best to date, the behavior of the particles which compose the infinitely small world, we can think learn significant informations from it, if one is able to extract exact solutions.

The general treatment of this equation as it is adopted today, and as it is developed in the first part, does not give all usable informations from this equation.

Some solutions are achievable only at the cost of additional conditions relating to the conservation of energy. Combining these conditions to the hypothesis of solutions in the form of stationary modes, it is possible to show that one can construct a deterministic physics and energy vision of the infinitely small physics.

The analysis of the solutions in Cartesian coordinates has allowed to check the complete consistency of these solutions with the vision of the Copenhagen school.

Solutions in spherical coordinates confirms this consistency with various theoretical and experimental elements.

The wave-particle duality is confirmed as an exchange of energy between mass energy and wave energy. The presence of imaginary terms in these exchanges shows that there is continuously transfer of energy between positive energy and negative energy. It is this exchange, taken as a plausible hypothesis, which allowed to access by a heuristic reasoning to exact solutions in spherical coordinates.

The concept of spin $\frac{1}{2}$, and its strange rotation of 4π to recover the initial state has an extremely simple physical explanation in one of exact solutions: when the component that carries the impulse energy following r rotates 2π to return to its starting point, the component that carries the impulse energy depending on φ and the component carrying the wave-particle duality turn 4π .

The uncertainty principle is not as exclusive as in Cartesian coordinates, but it may be noted that at any given time, mass energy and impulse energy are never maximum at the same place, which is sufficient to show that in a point in space we cannot have all of the information concerning these two kinds of energy. Since we can only have a measure of the momentum or position by an exchange of energy with these two quantities, these measures are necessarily tainted uncertainty when they occur in the same point of space.

Fourth part

Exact stationary Solutions to DIRAC equation in a variable potential

XXIV

The DIRAC equation in a variable potential: classical approach

Modal solutions to the generalized DIRAC equation proposed in previous chapters have been developed in a very simplistic framework that is a uniform and constant potential, i.e. independent of space and time. We know that in these circumstances the associated electromagnetic field is null. So far, the presence of the potential causes changes in the particle energy that are highlighted in the exact solutions that have been developed.

Now, we want to move towards exact solutions when the particle is immersed in a variable potential in space and in time. Such solutions should enable us to apprehend the modal changes that occur in the particle in the presence of an electric and magnetic field, uniform or variable in space and time.

To move towards such solutions, it is necessary to lean on the existing classical formalism, in which the constraints imposed by the modal solutions will be introduced. This formalism is most often presented in a very condensed form which enables an overall vision of the phenomena, but often mask elements of great complexity. These elements are listed in detail below, in order to introduce changes induced by the stationary solutions in the following chapters.

We take as starting point the system of DIRAC for a particle in a variable potential, which means that each component of the potential may depend on the variables of space and time:

$$\begin{aligned}\eta\psi_0 &= j\left(\frac{\partial\psi_0}{\partial x_t} + j\frac{q\phi}{\hbar}\psi_0\right) + j\left(\frac{\partial\psi_3}{\partial x} + j\frac{qA_x}{\hbar}\psi_3\right) + \left(\frac{\partial\psi_3}{\partial y} + j\frac{qA_y}{\hbar}\psi_3\right) + j\left(\frac{\partial\psi_2}{\partial z} + j\frac{qA_z}{\hbar}\psi_2\right) \\ \eta\psi_1 &= j\left(\frac{\partial\psi_1}{\partial x_t} + j\frac{q\phi}{\hbar}\psi_1\right) + j\left(\frac{\partial\psi_2}{\partial x} + j\frac{qA_x}{\hbar}\psi_2\right) - \left(\frac{\partial\psi_2}{\partial y} + j\frac{qA_y}{\hbar}\psi_2\right) - j\left(\frac{\partial\psi_3}{\partial z} + j\frac{qA_z}{\hbar}\psi_3\right) \\ \eta\psi_2 &= -j\left(\frac{\partial\psi_2}{\partial x_t} + j\frac{q\phi}{\hbar}\psi_2\right) - j\left(\frac{\partial\psi_1}{\partial x} + j\frac{qA_x}{\hbar}\psi_1\right) - \left(\frac{\partial\psi_1}{\partial y} + j\frac{qA_y}{\hbar}\psi_1\right) - j\left(\frac{\partial\psi_0}{\partial z} + j\frac{qA_z}{\hbar}\psi_0\right) \\ \eta\psi_3 &= -j\left(\frac{\partial\psi_3}{\partial x_t} + j\frac{q\phi}{\hbar}\psi_3\right) - j\left(\frac{\partial\psi_0}{\partial x} + j\frac{qA_x}{\hbar}\psi_0\right) + \left(\frac{\partial\psi_0}{\partial y} + j\frac{qA_y}{\hbar}\psi_0\right) + j\left(\frac{\partial\psi_1}{\partial z} + j\frac{qA_z}{\hbar}\psi_1\right)\end{aligned}$$

(XXIV-1)

As in the rest of the document, it retains the abridged notation:

$$\eta_t = \frac{q\phi}{\hbar c} \quad \eta_x = \frac{qA_x}{\hbar} \quad \eta_y = \frac{qA_y}{\hbar} \quad \eta_z = \frac{qA_z}{\hbar} \quad (XXIV-2)$$

In this notation, the components of the potential are related to the electromagnetic field by the following relationships:

- For the magnetic field:

$$\vec{B} = \begin{pmatrix} B_x \\ B_y \\ B_z \end{pmatrix} = \overrightarrow{\text{Rot}}(\vec{A}) = \begin{pmatrix} \frac{\partial A_z}{\partial y} - \frac{\partial A_y}{\partial z} \\ \frac{\partial A_x}{\partial z} - \frac{\partial A_z}{\partial x} \\ \frac{\partial A_y}{\partial x} - \frac{\partial A_x}{\partial y} \end{pmatrix} \quad (XXIV-3)$$

For the electric field:

$$\vec{E} = \begin{pmatrix} E_x \\ E_y \\ E_z \end{pmatrix} = -\overrightarrow{\text{Grad}}(\phi) - \frac{\partial \vec{A}}{\partial t} = \begin{pmatrix} -\frac{\partial \phi}{\partial x} - \frac{\partial A_x}{\partial t} \\ -\frac{\partial \phi}{\partial y} - \frac{\partial A_y}{\partial t} \\ -\frac{\partial \phi}{\partial z} - \frac{\partial A_z}{\partial t} \end{pmatrix} \quad (XXIV-4)$$

In order to highlight the behaviour of the bi-spinor, we rearrange the system (XXIV-1) above in the form:

$$\begin{aligned} j \frac{\partial \psi_0}{\partial x_t} - \eta_t \psi_0 &= \eta \psi_0 - j \frac{\partial \psi_3}{\partial x} + \eta_x \psi_3 - \frac{\partial \psi_3}{\partial y} - j \eta_y \psi_3 - j \frac{\partial \psi_2}{\partial z} + \eta_z \psi_2 \\ j \frac{\partial \psi_1}{\partial x_t} - \eta_t \psi_1 &= \eta \psi_1 - j \frac{\partial \psi_2}{\partial x} + \eta_x \psi_2 + \frac{\partial \psi_2}{\partial y} + j \eta_y \psi_2 + j \frac{\partial \psi_3}{\partial z} - \eta_z \psi_3 \\ j \frac{\partial \psi_2}{\partial x_t} - \eta_t \psi_2 &= -\eta \psi_2 - j \frac{\partial \psi_1}{\partial x} + \eta_x \psi_1 - \frac{\partial \psi_1}{\partial y} - j \eta_y \psi_1 - j \frac{\partial \psi_0}{\partial z} + \eta_z \psi_0 \\ j \frac{\partial \psi_3}{\partial x_t} - \eta_t \psi_3 &= -\eta \psi_3 - j \frac{\partial \psi_0}{\partial x} + \eta_x \psi_0 + \frac{\partial \psi_0}{\partial y} + j \eta_y \psi_0 + j \frac{\partial \psi_1}{\partial z} - \eta_z \psi_1 \end{aligned} \quad (XXIV-5)$$

Therefore, if we put:

$$\Psi = \begin{pmatrix} \Psi_0 \\ \Psi_1 \\ \Psi_2 \\ \Psi_3 \end{pmatrix} = \begin{pmatrix} \varphi \\ \chi \end{pmatrix} \quad \text{with } \varphi = \begin{pmatrix} \Psi_0 \\ \Psi_1 \end{pmatrix} \quad \text{et} \quad \chi = \begin{pmatrix} \Psi_2 \\ \Psi_3 \end{pmatrix} \quad (\text{XXIV-6})$$

Dirac system breaks down into two coupled systems:

$$\begin{aligned} j \left(\frac{\partial}{\partial x_t} + j\eta_t \right) \begin{pmatrix} \Psi_0 \\ \Psi_1 \end{pmatrix} &= \eta \begin{pmatrix} \Psi_0 \\ \Psi_1 \end{pmatrix} - j \begin{pmatrix} \left(\frac{\partial}{\partial z} + j\eta_z \right) & \left(\frac{\partial}{\partial x} + j\eta_x \right) - j \left(\frac{\partial}{\partial y} + j\eta_y \right) \\ \left(\frac{\partial}{\partial x} + j\eta_x \right) + j \left(\frac{\partial}{\partial y} + j\eta_y \right) & - \left(\frac{\partial}{\partial z} + j\eta_z \right) \end{pmatrix} \begin{pmatrix} \Psi_2 \\ \Psi_3 \end{pmatrix} \\ j \left(\frac{\partial}{\partial x_t} + j\eta_t \right) \begin{pmatrix} \Psi_2 \\ \Psi_3 \end{pmatrix} &= -\eta \begin{pmatrix} \Psi_2 \\ \Psi_3 \end{pmatrix} - j \begin{pmatrix} \left(\frac{\partial}{\partial z} + j\eta_z \right) & \left(\frac{\partial}{\partial x} + j\eta_x \right) - j \left(\frac{\partial}{\partial y} + j\eta_y \right) \\ \left(\frac{\partial}{\partial x} + j\eta_x \right) + j \left(\frac{\partial}{\partial y} + j\eta_y \right) & - \left(\frac{\partial}{\partial z} + j\eta_z \right) \end{pmatrix} \begin{pmatrix} \Psi_0 \\ \Psi_1 \end{pmatrix} \end{aligned} \quad (\text{XXIV-7})$$

We put for easy writing:

$$\mathbf{M} = -j \begin{pmatrix} \left(\frac{\partial}{\partial z} + j\eta_z \right) & \left(\frac{\partial}{\partial x} + j\eta_x \right) - j \left(\frac{\partial}{\partial y} + j\eta_y \right) \\ \left(\frac{\partial}{\partial x} + j\eta_x \right) + j \left(\frac{\partial}{\partial y} + j\eta_y \right) & - \left(\frac{\partial}{\partial z} + j\eta_z \right) \end{pmatrix} \quad (\text{XXIV-8})$$

The coupled system takes a simplified form:

$$\begin{aligned} j \left(\frac{\partial}{\partial x_t} + j\eta_t \right) (\varphi) &= \eta(\varphi) + \mathbf{M}(\chi) \\ j \left(\frac{\partial}{\partial x_t} + j\eta_t \right) (\chi) &= -\eta(\chi) + \mathbf{M}(\varphi) \end{aligned} \quad (\text{XXIV-9})$$

Or still:

$$\begin{aligned} \left(j \frac{\partial}{\partial x_t} - (\eta_t + \eta) \right) (\varphi) &= \mathbf{M}(\chi) \\ \left(j \frac{\partial}{\partial x_t} - (\eta_t - \eta) \right) (\chi) &= \mathbf{M}(\varphi) \end{aligned} \quad (\text{XXIV-10})$$

From which we deduce formally:

$$\begin{aligned} \left(j \frac{\partial}{\partial x_t} - (\eta_t + \eta) \right) \left(j \frac{\partial \varphi}{\partial x_t} - (\eta_t - \eta) \varphi \right) &= M^2 \varphi \\ \left(j \frac{\partial}{\partial x_t} - (\eta_t - \eta) \right) \left(j \frac{\partial \chi}{\partial x_t} - (\eta_t + \eta) \chi \right) &= M^2 \chi \end{aligned} \quad (\text{XXIV-11})$$

Yields by developing the left member of equality:

$$\begin{aligned} -\frac{\partial^2 \varphi}{\partial x_t^2} - 2j \frac{\partial \varphi}{\partial x_t} \eta_t + \eta_t^2 \varphi - j \frac{\partial \eta_t}{\partial x_t} \varphi - \eta^2 \varphi &= M^2 \varphi \\ -\frac{\partial^2 \chi}{\partial x_t^2} - 2j \frac{\partial \chi}{\partial x_t} \eta_t + \eta_t^2 \chi - j \frac{\partial \eta_t}{\partial x_t} \chi - \eta^2 \chi &= M^2 \chi \end{aligned} \quad (\text{XXIV-12})$$

To make progress towards a possible solution, we must now establish the expression of matrix M^2 .

If we adopt the following notation for PAULI matrixes and the operators involved in matrix computations:

$$\bar{\sigma} = \begin{pmatrix} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \\ \begin{pmatrix} 0 & -j \\ j & 0 \end{pmatrix} \\ \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \end{pmatrix} \quad \bar{p} = \begin{pmatrix} -j\hbar \frac{\partial}{\partial x} \\ -j\hbar \frac{\partial}{\partial y} \\ -j\hbar \frac{\partial}{\partial z} \end{pmatrix} \quad \bar{\pi} = (\bar{p} + q\bar{A}) = \begin{pmatrix} -j\hbar \frac{\partial}{\partial x} + qA_x \\ -j\hbar \frac{\partial}{\partial y} + qA_y \\ -j\hbar \frac{\partial}{\partial z} + qA_z \end{pmatrix} = \begin{pmatrix} -j\hbar \left(\frac{\partial}{\partial x} + j\eta_x \right) \\ -j\hbar \left(\frac{\partial}{\partial y} + j\eta_y \right) \\ -j\hbar \left(\frac{\partial}{\partial z} + j\eta_z \right) \end{pmatrix} \quad (\text{XXIV-13})$$

the matrix M established in (XXIV-8) is expressed as follows:

$$\begin{aligned} \hbar M(\varphi) &= [\bar{\sigma} \cdot (\bar{p} + q\bar{A})](\varphi) = -j\hbar \left[\begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \left(\frac{\partial}{\partial x} + j\eta_x \right) + \begin{pmatrix} 0 & -j \\ j & 0 \end{pmatrix} \left(\frac{\partial}{\partial y} + j\eta_y \right) + \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \left(\frac{\partial}{\partial z} + j\eta_z \right) \right] (\varphi) \\ M &= -j \begin{pmatrix} \left(\frac{\partial}{\partial z} + j\eta_z \right) & \left(\frac{\partial}{\partial x} + j\eta_x \right) - j \left(\frac{\partial}{\partial y} + j\eta_y \right) \\ \left(\frac{\partial}{\partial x} + j\eta_x \right) + j \left(\frac{\partial}{\partial y} + j\eta_y \right) & - \left(\frac{\partial}{\partial z} + j\eta_z \right) \end{pmatrix} \end{aligned} \quad (\text{XXIV-14})$$

The development of the matrix M^2 is a bit laborious. The reader will take care that each of the terms of the matrix represents an operator and must therefore be treated as such in the operations of derivations.

If we point out each of the terms of this matrix by:

$$M^2 = \begin{pmatrix} M_{11} & M_{12} \\ M_{21} & M_{22} \end{pmatrix} \quad (\text{XXIV-15})$$

$$M^2 = \begin{pmatrix} -\frac{\partial^2 \varphi}{\partial x_t^2} - 2j \frac{\partial \varphi}{\partial x_t} \eta_t + \eta_t^2 \varphi - j \frac{\partial \eta_t}{\partial x_t} \varphi - \eta_t^2 \varphi & 0 \\ 0 & -\frac{\partial^2 \varphi}{\partial x_t^2} - 2j \frac{\partial \varphi}{\partial x_t} \eta_t + \eta_t^2 \varphi - j \frac{\partial \eta_t}{\partial x_t} \varphi - \eta_t^2 \varphi \end{pmatrix} \quad (\text{XXIV-19})$$

It appears that equality above can be checked only if:

$$\left(\frac{\partial \eta_z}{\partial y} - \frac{\partial \eta_y}{\partial z} \right) \varphi + j \left(\frac{\partial \eta_x}{\partial z} - \frac{\partial \eta_z}{\partial x} \right) \varphi = 0 \quad (\text{XXIV-20})$$

For this relationship to be true regardless of any spinor φ , this requires:

$$\begin{aligned} \left(\frac{\partial \eta_z}{\partial y} - \frac{\partial \eta_y}{\partial z} \right) &= 0 \\ \left(\frac{\partial \eta_x}{\partial z} - \frac{\partial \eta_z}{\partial x} \right) &= 0 \end{aligned} \quad (\text{XXIV-21})$$

When the particle is immersed in a magnetic field, it is deduced from (XXIV-21) and (XXIV-3) that the system of DIRAC have only solutions when the magnetic field is directed along the Oz axis.

Since angular momentum is also oriented along the axis Oz, this leads to a quantification of the spin when the particle is immersed in a magnetic field, with only two possible States.

Specifically, when a free particle is introduced in a STERN and GERLACH device for example, this necessarily induced a spatial evolution of the internal modes to the particle. This spatial evolution will put in alignment the magnetic moment of the electron with the magnetic field generated by the device.

It can be assumed that this alignment will be done on an extremely low duration. If it is estimated empirically on a time τ equal to a hundred period, it gets an order of magnitude of time τ as followed:

$$\tau = 100 \cdot \frac{2\pi}{\omega_0} = 200\pi \frac{\hbar}{m_0 c^2} = 200\pi \frac{1,05 \cdot 10^{-34}}{(9,11 \cdot 10^{-31})(3 \cdot 10^8)^2} \approx 8,09 \cdot 10^{-19} \text{s} \quad (\text{XXIV-22})$$

By comparing this value to the duration t of the path of the electron on a distance of 1 meter, the speed limit for light:

$$t = \frac{1}{3 \cdot 10^8} = 0,33 \cdot 10^{-8} \text{s} \quad (\text{XXIV-23})$$

It can be concluded that the alignment of the magnetic moment of the electron with the magnetic field generated by the device of STERN and GERLACH is almost instantly to its

entry into the device. It thus gets a description of the concept of spatial quantization, a concept which is extremely difficult to have a physical representation when addressing quantum mechanics.

We propose to finish this chapter to remember the relationship established by DIRAC, which highlights the role of the magnetic field on spinors:

$$\hbar^2 M^2(\varphi) = [\vec{\sigma} \cdot (\vec{p} + q\vec{A})]^2(\varphi) = (\vec{p} + q\vec{A})^2(\varphi) + \vec{\sigma} \cdot \overrightarrow{\text{Rot}}(q\vec{A})(\varphi) = (\vec{p} + q\vec{A})^2(\varphi) + \vec{\sigma} \cdot q\vec{B}(\varphi) \quad (\text{XXIV-24})$$

Matrix M^2 left of equality has already been expressed in (XXIV-18). The calculation of the terms located to the right of equality is detailed below:

$$\begin{aligned} (\vec{p} + q\vec{A})^2(\varphi) &= \begin{pmatrix} -j\hbar \frac{\partial}{\partial x} + qA_x \\ -j\hbar \frac{\partial}{\partial y} + qA_y \\ -j\hbar \frac{\partial}{\partial z} + qA_z \end{pmatrix} \begin{pmatrix} -j\hbar \frac{\partial \varphi}{\partial x} + qA_x \varphi \\ -j\hbar \frac{\partial \varphi}{\partial y} + qA_y \varphi \\ -j\hbar \frac{\partial \varphi}{\partial z} + qA_z \varphi \end{pmatrix} = -\hbar^2 \begin{pmatrix} \frac{\partial}{\partial x} - j\eta_x \\ \frac{\partial}{\partial y} - j\eta_y \\ \frac{\partial}{\partial z} - j\eta_z \end{pmatrix} \begin{pmatrix} \frac{\partial \varphi}{\partial x} - j\eta_x \varphi \\ \frac{\partial \varphi}{\partial y} - j\eta_y \varphi \\ \frac{\partial \varphi}{\partial z} - j\eta_z \varphi \end{pmatrix} \\ (\vec{p} + q\vec{A})^2(\varphi) &= -\hbar^2 \left\{ \begin{aligned} &\left(\frac{\partial^2 \varphi}{\partial x^2} + 2j\eta_x \frac{\partial \varphi}{\partial x} - \eta_x^2 \varphi + j\varphi \frac{\partial \eta_x}{\partial x} \right) + \left(\frac{\partial^2 \varphi}{\partial y^2} + 2j\eta_y \frac{\partial \varphi}{\partial y} - \eta_y^2 \varphi + j\varphi \frac{\partial \eta_y}{\partial y} \right) \\ &+ \left(\frac{\partial^2 \varphi}{\partial z^2} + 2j\eta_z \frac{\partial \varphi}{\partial z} - \eta_z^2 \varphi + j\varphi \frac{\partial \eta_z}{\partial z} \right) \end{aligned} \right\} \quad (\text{XXIV-25}) \end{aligned}$$

$$\begin{aligned} \vec{\sigma} \cdot \overrightarrow{\text{Rot}}(q\vec{A}) &= \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} \frac{\partial qA_z}{\partial y} - \frac{\partial qA_y}{\partial z} \\ \frac{\partial qA_x}{\partial z} - \frac{\partial qA_z}{\partial x} \end{pmatrix} + \begin{pmatrix} 0 & -j \\ j & 0 \end{pmatrix} \begin{pmatrix} \frac{\partial qA_x}{\partial z} - \frac{\partial qA_z}{\partial x} \\ \frac{\partial qA_y}{\partial x} - \frac{\partial qA_x}{\partial y} \end{pmatrix} + \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \begin{pmatrix} \frac{\partial qA_y}{\partial x} - \frac{\partial qA_x}{\partial y} \\ \frac{\partial qA_z}{\partial y} - \frac{\partial qA_y}{\partial z} \end{pmatrix} \\ \vec{\sigma} \cdot \overrightarrow{\text{Rot}}(q\vec{A}) &= \begin{pmatrix} \frac{\partial qA_y}{\partial x} - \frac{\partial qA_x}{\partial y} & \left(\frac{\partial qA_z}{\partial y} - \frac{\partial qA_y}{\partial z} \right) - j \left(\frac{\partial qA_x}{\partial z} - \frac{\partial qA_z}{\partial x} \right) \\ \left(\frac{\partial qA_z}{\partial y} - \frac{\partial qA_y}{\partial z} \right) + j \left(\frac{\partial qA_x}{\partial z} - \frac{\partial qA_z}{\partial x} \right) & - \left(\frac{\partial qA_y}{\partial x} - \frac{\partial qA_x}{\partial y} \right) \end{pmatrix} \\ \vec{\sigma} \cdot \overrightarrow{\text{Rot}}(q\vec{A}) &= \hbar^2 \begin{pmatrix} \frac{\partial \eta_y}{\partial x} - \frac{\partial \eta_x}{\partial y} & \left(\frac{\partial \eta_z}{\partial y} - \frac{\partial \eta_y}{\partial z} \right) - j \left(\frac{\partial \eta_x}{\partial z} - \frac{\partial \eta_z}{\partial x} \right) \\ \left(\frac{\partial \eta_z}{\partial y} - \frac{\partial \eta_y}{\partial z} \right) + j \left(\frac{\partial \eta_x}{\partial z} - \frac{\partial \eta_z}{\partial x} \right) & - \left(\frac{\partial \eta_y}{\partial x} - \frac{\partial \eta_x}{\partial y} \right) \end{pmatrix} \quad (\text{XXIV-26}) \end{aligned}$$

The relationship:

$$\hbar^2 M^2(\varphi) = [\vec{\sigma} \cdot (\vec{p} + q\vec{A})]^2(\varphi) = (\vec{p} + q\vec{A})^2(\varphi) + \vec{\sigma} \cdot \overrightarrow{\text{Rot}}(q\vec{A})(\varphi) \quad (\text{XXIV-27})$$

is checked by introducing the identity matrix in factor with the term $(\vec{p} + q\vec{A})^2$.

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XXV

Exact solutions of the Dirac equation in a variable potential in Cartesian coordinates

We discuss in this chapter looking for exact solutions to the DIRAC equation in a variable potential. We know from the previous chapter that this potential can have a magnetic field, but only directed along the Oz axis.

The presence of the electromagnetic field will induce changes on stationary modes, and so on the components of the wave vector k_t , k_x , k_y , k_z , which will depend on the components of the potential, and there are several ways to take into account these changes.

We adopt the following search strategy: on a point of space and time, the potential will depend on the spatial and temporal variables x , y , z , t . We seek modal solutions in the form of imaginary exponentials $\exp[j(k_t x_t + k_x x + k_y y + k_z z)]$.

The changes with respect to the modes of free space (without potential) will carry on modifications on the amplitude of each wave functions, and on the components of the wave vector k_t , k_x , k_y , k_z , through the conservation of energy equation.

In the previous chapter, we have established the following equality which must be valid for any spinor:

$$M^2 = \begin{pmatrix} -\frac{\partial^2 \varphi}{\partial x_t^2} - 2j \frac{\partial \varphi}{\partial x_t} \eta_t + \eta_t^2 \varphi - j \frac{\partial \eta_t}{\partial x_t} \varphi - \eta_t^2 \varphi & 0 \\ 0 & -\frac{\partial^2 \varphi}{\partial x_t^2} - 2j \frac{\partial \varphi}{\partial x_t} \eta_t + \eta_t^2 \varphi - j \frac{\partial \eta_t}{\partial x_t} \varphi - \eta_t^2 \varphi \end{pmatrix} \quad (\text{XXV-1})$$

The particular shape of relation (XXV-1) shows that the two spinors are decoupled. Developing writing, we obtain the equation of conservation of energy in its most general form by gathering the influence of the vector potential under the two sign + and -:

$$\begin{aligned} \frac{\partial^2 \varphi}{\partial x_t^2} + 2j \frac{\partial \varphi}{\partial x_t} \eta_t - \eta_t^2 \varphi + j \frac{\partial \eta_t}{\partial x_t} \varphi + \eta^2 \varphi = \\ \left\{ \left(\frac{\partial^2 \varphi}{\partial x^2} + 2j \eta_x \frac{\partial \varphi}{\partial x} - \eta_x^2 \varphi + j \varphi \frac{\partial \eta_x}{\partial x} \right) + \left(\frac{\partial^2 \varphi}{\partial y^2} + 2j \eta_y \frac{\partial \varphi}{\partial y} - \eta_y^2 \varphi + j \varphi \frac{\partial \eta_y}{\partial y} \right) \right\} \\ \left\{ + \left(\frac{\partial^2 \varphi}{\partial z^2} + 2j \eta_z \frac{\partial \varphi}{\partial z} - \eta_z^2 \varphi + j \varphi \frac{\partial \eta_z}{\partial z} \right) \pm \left\{ \frac{\partial \eta_y}{\partial x} \varphi - \frac{\partial \eta_x}{\partial y} \varphi \right\} \right\} \end{aligned} \quad (\text{XXV-2})$$

There are solutions only if we are in a magnetic field oriented along Oz, which imposes:

$$\begin{aligned} \left(\frac{\partial \eta_z}{\partial y} - \frac{\partial \eta_y}{\partial z} \right) = 0 \\ \left(\frac{\partial \eta_x}{\partial z} - \frac{\partial \eta_z}{\partial x} \right) = 0 \end{aligned} \quad (\text{XXV-3})$$

We will restrict the search of solutions field placing us in the following simplifying assumptions:

$$\frac{\partial \eta_t}{\partial x_t} = \frac{\partial \eta_x}{\partial x} = \frac{\partial \eta_y}{\partial y} = \eta_z = 0 \quad (\text{XXV-4})$$

We obtain after simplification:

$$\frac{\partial^2 \varphi}{\partial x_t^2} + 2j \frac{\partial \varphi}{\partial x_t} \eta_t - \eta_t^2 \varphi + \eta^2 \varphi = \left\{ \left(\frac{\partial^2 \varphi}{\partial x^2} + 2j \eta_x \frac{\partial \varphi}{\partial x} - \eta_x^2 \varphi \right) + \left(\frac{\partial^2 \varphi}{\partial y^2} + 2j \eta_y \frac{\partial \varphi}{\partial y} - \eta_y^2 \varphi \right) \right\} \\ \left\{ + \frac{\partial^2 \varphi}{\partial z^2} \pm \left\{ \frac{\partial \eta_y}{\partial x} \varphi - \frac{\partial \eta_x}{\partial y} \varphi \right\} \right\} \quad (\text{XXV-5})$$

We introduce the particular form of the solution:

$$\varphi = \exp j(k_t x_t + k_x x + k_y y + k_z z) \quad (\text{XXV-6})$$

That gives the energy conservation equation relative to (XXV-6) in a variable potential:

$$(k_t + \eta_t)^2 = (k_x + \eta_x)^2 + (k_y + \eta_y)^2 + k_z^2 + \eta^2 \pm \left(\frac{\partial \eta_y}{\partial x} - \frac{\partial \eta_x}{\partial y} \right) \quad (\text{XXV-7})$$

It is now possible to begin the search for solutions itself. For a null potential, DIRAC system is written:

$$\begin{aligned}
 \eta\psi_0 &= j\frac{\partial\psi_0}{\partial x_t} + j\frac{\partial\psi_3}{\partial x} + \frac{\partial\psi_3}{\partial y} + j\frac{\partial\psi_2}{\partial z} \\
 \eta\psi_1 &= j\frac{\partial\psi_1}{\partial x_t} + j\frac{\partial\psi_2}{\partial x} - \frac{\partial\psi_2}{\partial y} - j\frac{\partial\psi_3}{\partial z} \\
 \eta\psi_2 &= -j\frac{\partial\psi_2}{\partial x_t} - j\frac{\partial\psi_1}{\partial x} - \frac{\partial\psi_1}{\partial y} - j\frac{\partial\psi_0}{\partial z} \\
 \eta\psi_3 &= -j\frac{\partial\psi_3}{\partial x_t} - j\frac{\partial\psi_0}{\partial x} + \frac{\partial\psi_0}{\partial y} + j\frac{\partial\psi_1}{\partial z}
 \end{aligned}
 \tag{XXV-8}$$

A possible solution to this system in an exponential form has already been developed in the previous chapters:

$$\begin{aligned}
 \psi_0 &= (\eta - k_t)\exp j(k_t x_t + k_x x + k_y y + k_z z) \\
 \psi_1 &= 0 \\
 \psi_2 &= k_z \exp j(k_t x_t + k_x x + k_y y + k_z z) \\
 \psi_3 &= (k_x + jk_y)\exp j(k_t x_t + k_x x + k_y y + k_z z)
 \end{aligned}
 \tag{XXV-9}$$

In a variable potential, DIRAC system is amended as follows:

$$\begin{aligned}
 (\eta + \eta_t)\psi_0 + \eta_x\psi_3 - j\eta_y\psi_3 + \eta_z\psi_2 &= j\frac{\partial\psi_0}{\partial x_t} + j\frac{\partial\psi_3}{\partial x} + \frac{\partial\psi_3}{\partial y} + j\frac{\partial\psi_2}{\partial z} \\
 (\eta + \eta_t)\psi_1 + \eta_x\psi_2 + j\eta_y\psi_2 - \eta_z\psi_3 &= j\frac{\partial\psi_1}{\partial x_t} + j\frac{\partial\psi_2}{\partial x} - \frac{\partial\psi_2}{\partial y} - j\frac{\partial\psi_3}{\partial z} \\
 (\eta - \eta_t)\psi_2 - \eta_x\psi_1 + j\eta_y\psi_1 - \eta_z\psi_0 &= -j\frac{\partial\psi_2}{\partial x_t} - j\frac{\partial\psi_1}{\partial x} - \frac{\partial\psi_1}{\partial y} - j\frac{\partial\psi_0}{\partial z} \\
 (\eta - \eta_t)\psi_3 - \eta_x\psi_0 - j\eta_y\psi_0 + \eta_z\psi_1 &= -j\frac{\partial\psi_3}{\partial x_t} - j\frac{\partial\psi_0}{\partial x} + \frac{\partial\psi_0}{\partial y} + j\frac{\partial\psi_1}{\partial z}
 \end{aligned}
 \tag{XXV-10}$$

A possible solution to the system of DIRAC in a variable potential appears in the following form:

$$\begin{aligned}
 \psi_0 &= (\eta - k_t - \eta_t)\exp j(k_t x_t + k_x x + k_y y + k_z z) \\
 \psi_1 &= 0 \\
 \psi_2 &= k_z \exp j(k_t x_t + k_x x + k_y y + k_z z) \\
 \psi_3 &= \{k_x + \eta_x + j(k_y + \eta_y)\}\exp j(k_t x_t + k_x x + k_y y + k_z z)
 \end{aligned}
 \tag{XXV-11}$$

It made the observation that the form of the solution remains identical to that which was obtained for a uniform potential. The introduction of a variable potential comes only to change the equation of energy conservation.

We give below a few elements of check. For this solution, the first equation of DIRAC system provides energy conservation equation.

First equation of the DIRAC system:

$$\begin{aligned}
 (\eta + \eta_t)\psi_0 + \eta_x\psi_3 - j\eta_y\psi_3 + \eta_z\psi_2 &= j\frac{\partial\psi_0}{\partial x_t} + j\frac{\partial\psi_3}{\partial x} + \frac{\partial\psi_3}{\partial y} + j\frac{\partial\psi_2}{\partial z} \\
 (\eta + \eta_t)\psi_0 + \eta_x\psi_3 - j\eta_y\psi_3 &= (\eta + \eta_t)(\eta - k_t - \eta_t) + \eta_x\{k_x + \eta_x + j(k_y + \eta_y)\} - j\eta_y\{k_x + \eta_x + j(k_y + \eta_y)\} \\
 j\frac{\partial\psi_0}{\partial x_t} &= -k_t(\eta - k_t - \eta_t) \\
 j\frac{\partial\psi_3}{\partial x} &= -k_x\{k_x + \eta_x + j(k_y + \eta_y)\} - \frac{\partial\eta_y}{\partial x} \\
 \frac{\partial\psi_3}{\partial y} &= jk_y\{k_x + \eta_x + j(k_y + \eta_y)\} + \frac{\partial\eta_x}{\partial y} \\
 j\frac{\partial\psi_2}{\partial z} &= -k_z^2
 \end{aligned}
 \tag{XXV-12}$$

The sums are left to the care of the reader. We get to the first equation:

$$(k_t + \eta_t)^2 = (k_x + \eta_x)^2 + (k_y + \eta_y)^2 + k_z^2 + \eta^2 + \left(\frac{\partial\eta_y}{\partial x} - \frac{\partial\eta_x}{\partial y} \right)
 \tag{XXV-13}$$

Second equation of the DIRAC system:

$$\begin{aligned}
 (\eta + \eta_t)\psi_1 + \eta_x\psi_2 + j\eta_y\psi_2 - \eta_z\psi_3 &= j\frac{\partial\psi_1}{\partial x_t} + j\frac{\partial\psi_2}{\partial x} - \frac{\partial\psi_2}{\partial y} - j\frac{\partial\psi_3}{\partial z} \\
 \eta_x\psi_2 + j\eta_y\psi_2 &= \eta_x k_z + j\eta_y k_z \\
 j\frac{\partial\psi_2}{\partial x} &= -k_x k_z \\
 -\frac{\partial\psi_2}{\partial y} &= -jk_y k_z \\
 -j\frac{\partial\psi_3}{\partial z} &= k_z\{k_x + \eta_x + j(k_y + \eta_y)\}
 \end{aligned}
 \tag{XXV-14}$$

Third equation of the DIRAC system:

$$\begin{aligned}
 (\eta - \eta_t)\psi_2 - \eta_x\psi_1 + j\eta_y\psi_1 - \eta_z\psi_0 &= -j\frac{\partial\psi_2}{\partial x_t} - j\frac{\partial\psi_1}{\partial x} - \frac{\partial\psi_1}{\partial y} - j\frac{\partial\psi_0}{\partial z} \\
 (\eta - \eta_t)\psi_2 &= (\eta - \eta_t)k_z \\
 -j\frac{\partial\psi_2}{\partial x_t} &= k_t k_z \\
 -j\frac{\partial\psi_0}{\partial z} &= k_z(\eta - k_t - \eta_t)
 \end{aligned}
 \tag{XXV-15}$$

Fourth equation of the DIRAC system:

$$\begin{aligned}(\eta - \eta_t)\psi_3 - \eta_x\psi_0 - j\eta_y\psi_0 + \eta_z\psi_1 &= -j\frac{\partial\psi_3}{\partial x_t} - j\frac{\partial\psi_0}{\partial x} + \frac{\partial\psi_0}{\partial y} + j\frac{\partial\psi_1}{\partial z} \\(\eta - \eta_t)\psi_3 - \eta_x\psi_0 - j\eta_y\psi_0 &= (\eta - \eta_t)\{k_x + \eta_x + j(k_y + \eta_y)\} - \eta_x(\eta - k_t - \eta_t) - j\eta_y(\eta - k_t - \eta_t) \\-j\frac{\partial\psi_3}{\partial x_t} &= k_t\{k_x + \eta_x + j(k_y + \eta_y)\} \\-j\frac{\partial\psi_0}{\partial x} &= k_x(\eta - k_t - \eta_t) \\\frac{\partial\psi_0}{\partial y} &= jk_y(\eta - k_t - \eta_t)\end{aligned}$$

(XXV-16)

XXVI

General conclusion

The energy approach to quantum mechanics proposed in this document is in opposition with the probabilistic approach of the Copenhagen school that dominates today, practically in the form of a single thought, the vision of quantum physics.

This probabilistic vision is installed during the discovery of particular physical properties of the infinitely small world. It is imposed under the constraint of the HEISENBERG uncertainty principle, supported by the equivalence of the SCHRÖDINGER equation with this principle. This vision was the only one that allows to account for all experimental results in consistency with quantum uncertainty principle of indeterminacy. It became inescapable as soon as it remained in agreement with experimental results more accurate and refined. Gradually, the physicists accepted it as a safe theory able to predict all of the properties of the quantum world.

When we adopt a critical point of view, the main problem that arises in the evolution of this theory is that it is based on a postulate given by the indeterminacy theorem. Therefore, everything that is developed using this theory is seen through the filter of this indeterminacy and can only be developed in a probabilistic approach. It becomes impossible to get out of this framework to explain specific phenomena of quantum mechanics.

One can only notice the powerlessness of the Copenhagen school to provide an explanation for certain phenomena like the wave-particle duality. After a century of research, and an investment of best physicists in understanding this phenomenon, it can be estimated without too much risk that there is an intrinsic impossibility to explain this duality by using a probabilistic theory.

Other complex phenomena such as entanglement have probabilistic explanation only by using contortions that come out of the classical laws of physics, and so raises some questions.

The theorem of uncertainty itself is accepted as a property of the quantum world, without any attempt of explanation of the physical nature of its interpretation.

Finally, the imaginary nature of the equation of SCHRÖDINGER and its solutions remains Enigma full, and can find no satisfactory explanation in this formalism.

However, to offer an alternative or an evolution of the Copenhagen school is a delicate challenge because any development must remain consistent with all of the achievements of this school, which is an unmovable base of quantum mechanics.

The energy approach presented in this document meets this challenge.

If we consider the DIRAC equation as a refinement of SCHRÖDINGER and KLEIN-GORDON equations, and if we remember the equivalence between the mechanics of the SCHRÖDINGER equation and HEISENBERG matrices, then, exact solutions to the DIRAC equation must be consistent with the overall results built on the principle of HEISENBERG and SCHRÖDINGER equation.

An analysis of exact solutions in the form of standing waves leads to reconsider the physical meaning of the wave function. The Copenhagen school gives their conjugated product the meaning of a volume density of probability of presence of the particle, appellation which hides the ignorance of the underlying physical phenomena to the concept of presence of a particle and which prevents any progress towards a more subtle understanding of the functioning of the particle.

By adopting the meaning of a volumetric energy density, we offer a way in which, without no other hypothesis that energy conservation and the evolution of energy in the form of standing waves, all issues not resolved by the Copenhagen school find a simple explanation.

The internal energy of the particle is broken down into three energies of different nature: energy of mass, impulse, and wave energy. These energies are exchanged between them and in an ultimate way, with the energy of the vacuum. This exchange of energy is signed and takes an opposite sign depending on whether the particle receives or renders the vacuum energy. DIRAC currents gather volume densities of all these energies according to each direction of space, allowing to apply the local conservation of energy principle in calculating their divergence. The imaginary nature of the wave function has no more any problem of interpretation.

There is no longer need to make ad hoc assumptions, or construct convoluted theories to explain the duality wave-matter, since it appears naturally in the stationary solutions as a very fast exchange between mass energy and wave energy.

The mystery of the principle of indeterminacy disappears. The energy exchanges within the particle show that it is impossible to have simultaneously and to the same place, all of the information on the impulse energy, and mass energy that is relative to the position operator.

It remains only a completely deterministic approach in agreement with all achievements of the Copenhagen school.

This research work is proposed to the scientific community interested in this topic to be shared, criticized, amended and/or may be strengthened.

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Second part

The second part deals with searching of exact solutions to the DIRAC equation in the form of stationary modes. It was not found any work having a direct link with the objective of this document. Consulted publications relate to the search for exact solutions to the DIRAC equation associated with diverse potentials or associated with gravitational problems

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Third part

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