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Fe_xGa_{1-x}BO₃ single crystals: synthesis and characterization

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Abstract A series of Fe_xGa_{1-x}BO₃ single crystals in the concentration range $0 \leq x \leq 1$ has been synthesized by solution in the melt method. In order to determine optimal crystallization regimes, two settings have been worked out and applied in the present study: the first one, for precise differential thermal analysis and the second one, for the probe method. X-ray fluorescence and X-ray diffraction analysis have allowed accurate determination of iron contents and lattice parameters for synthesized crystals with different x . Computer-assisted EPR studies of Fe³⁺ have revealed a high perfection of the crystals: low degree of disorder and the absence of twinning.

1 Introduction

Iron borate single-crystal FeBO₃, as a model compound, has been used in numerous studies, in particular, in magnetism. This is explained by an extraordinary set of its physical characteristics: magnetic, resonance, magneto-acoustic, optical and magneto-optical, etc., e.g., see [1–15]. A simultaneous occurrence of some of these characteristics is quite outstanding; e.g., FeBO₃ is transparent for visible light and, at the same time, magnetically ordered. This crystal has calcite structure, rhombohedral, of D_{3d}^6 space group. In the hexagonal system, unit cell parameters are $a = 4.626$ and $c = 14.493$ Å, and the unit cell volume is

$V = 268.57$ Å³ [16]. From the viewpoint of magnetic structure, it is a two sublattice easy-plane antiferromagnet with a weak ferromagnetism, the Néel temperature being $T_N = 348$ K [16]. In spite of a great number of studies, the exact nature of certain physical characteristics of this crystal still remains unclear. For instance, magnetic anisotropy (volume and surface [6, 7]) in FeBO₃ is known to include two contributions: single-ion anisotropy and magnetic dipole interactions. However, their respective contributions have not yet been reliably evaluated.

Isomorphous substitution of a part of *paramagnetic* iron in FeBO₃ by *diamagnetic* gallium is a very powerful technique allowing a comprehensive description of various intracrystalline interactions. It makes use of the fact that different mechanisms contributing in the physical characteristics have different concentration and temperature dependences. The series of Fe_xGa_{1-x}BO₃ single crystals, synthesized in a wide concentration range, allows studying in detail the transformation of various physical characteristics under the transition from magnetically ordered to paramagnetic state.

Iron borate crystals can be synthesized by two routes: (1) from the gas phase and (2) from the solution in the melt. Technology of synthesis of pure iron borate, FeBO₃, single crystals by both routes has been developed by our group [17, 18]. Using the gas-phase technique, we have obtained *bulk* single crystals of iron borate with large *non-basal* faces of optical quality. The existence of natural non-basal faces has allowed finding out and describing surface magnetism caused by specific surface anisotropy [6, 7]. Another effect observed only in bulk crystals is magnetic birefringence of sound [4]. The solution in the melt technique allows obtaining *thin*, 0.05–0.1 mm along the trigonal C₃ axis [16] and up to 10 mm in the basal plane (0001) [16], crystals. These crystals are of high structural

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perfection and much less costly in manufacturing. For the purposes of the present work, this technique has proved to be most appropriate. Thus the solution in the melt technique has been used for isomorphous substitution of Fe^{3+} by Ga^{3+} in $\text{Fe}_x\text{Ga}_{1-x}\text{BO}_3$.

Using the solution-melt technique simplifies the crystal synthesis, allows obtaining materials with complex compositions and reduces the crystallization temperature of refractory compounds. Usually, crystallization regimes suitable for obtaining desired crystal phases are determined empirically, requiring a significant time and financial resources. Hence, express methods of determining the ranges of stability of crystalline phases and temperature borders of the phase formation are desirable. In the present work for this purpose, we have used differential thermal analysis (DTA) and probe method.

After successful $\text{Fe}_x\text{Ga}_{1-x}\text{BO}_3$ synthesis, a detailed characterization is required for further studies and applications of the samples. The exact crystal composition and parameters of the crystal lattice have been determined by X-ray fluorescence analysis (XRF) and X-ray diffraction (XRD), respectively. Besides, electron paramagnetic resonance (EPR) has been used to estimate local disorder in the $\text{Fe}_x\text{Ga}_{1-x}\text{BO}_3$ crystals at low x values.

The aim of the present paper is to synthesize series of $\text{Fe}_x\text{Ga}_{1-x}\text{BO}_3$ crystals in a wide range of substitutions and to obtain their principal characteristics, such as exact crystal composition, parameters of the crystal lattice and quality of the crystals. The synthesized crystals provide the basis for future fundamental magnetic, magneto-optical and resonance studies of the influence of diamagnetic dilution on the properties of the magnetically concentrated iron borate. Besides, from the practical viewpoint, $\text{Fe}_x\text{Ga}_{1-x}\text{BO}_3$ crystals are promising for applications as memory elements, magneto-optical and magneto-acoustic transducers, etc., and our developed crystal synthesis technique is well adapted to fine-tuning the characteristics of such materials by varying the concentration of magnetic ions.

2 Experimental

$\text{Fe}_x\text{Ga}_{1-x}\text{BO}_3$ single crystals have been synthesized by solution in the melt technique. Crystallizations were carried out in the Ga_2O_3 Fe_2O_3 B_2O_3 PbO PbF_2 system. Some examples of charge compositions used in the synthesis are shown in Table 1.

In order to determine exact contents of iron in the synthesized crystals, x_{crystal} , we have carried out XRF analysis.

The parameters of crystal lattice in the series of $\text{Fe}_x\text{Ga}_{1-x}\text{BO}_3$ crystals have been determined by XRD with

Table 1 Charge compositions used for synthesizing $\text{Fe}_x\text{Ga}_{1-x}\text{BO}_3$ crystals for some x values

x	Ga_2O_3 (Mass %)	Fe_2O_3 (Mass %)	B_2O_3 (Mass %)	PbO (Mass %)	PbF_2 (Mass %)
0.02	18.52	0.08			
0.2	16.99	1.61	42.4	27.3	11.7
0.3	14.49	4.11			

a Rigaku SmartLab diffractometer using copper radiation $\text{CuK}\alpha$ in the angle range of $5^\circ < 2\theta < 100^\circ$ with the step of 0.02° .

For estimating the quality of the synthesized crystals, we have chosen EPR characterization technique, particularly sensitive to imperfections of local structure. At low x values, only the EPR of diluted Fe^{3+} in GaBO_3 lattice is observed; therefore, this ion can serve as a high-sensitive probe for evaluating the degree of structural disorder in crystals. A single crystal with $x = 0.003$ has been studied by EPR with an X-band Bruker spectrometer at 4 K and magnetizing fields up to 1 T.

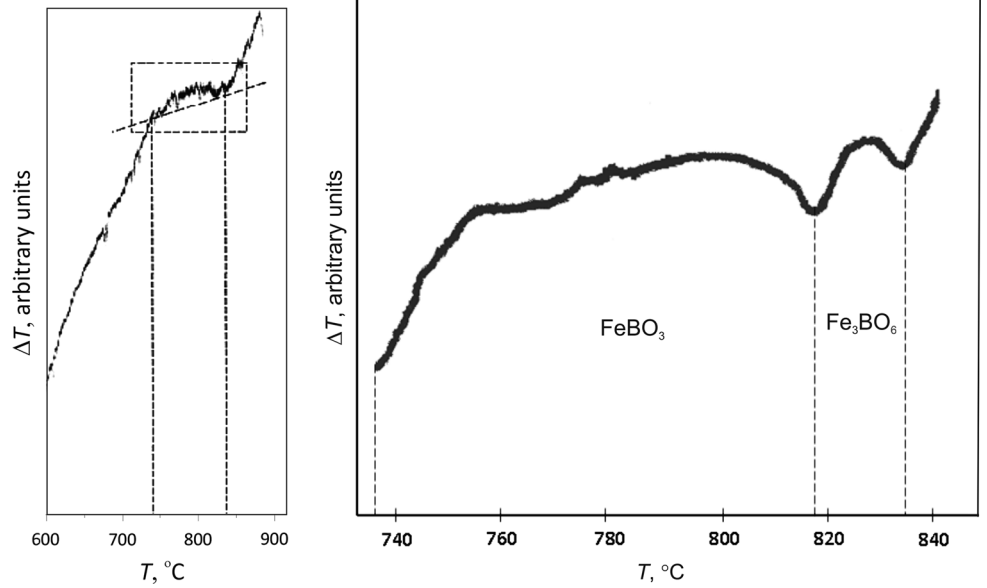
3 Results and discussion

DTA can be applied to a wide range of substances; meanwhile, there have been only few DTA studies of the solution-melt systems. This can be explained by the fact that for diluted solutions, sensitivity of this technique is greatly reduced because of the smallness of the exothermic peak on the thermograms [19]. Therefore, we have developed a highly sensitive setting for rapid analysis of small amounts of the solution melts in the temperature range from 300 to 1150 °C. This setting allows detecting weak thermal effects caused by the formation of crystalline phases and determining temperature borders of crystallization with sufficient accuracy [20].

As a differential thermocouple, we have used two platinum crucibles having a volume of 1.5 cm^3 . This allows recording on the thermograms the temperature difference ΔT between the crucibles of the order of 0.01°C .

We have carried out DTA studies of a number of compositions for the crystal synthesis with different concentrations of the components. As an example, the DTA cooling curve for the composition used to synthesize FeBO_3 is shown in Fig. 1 (left). As one can see, in the range of 835–737 °C, a flat exothermic peak is observed. A detailed analysis of this temperature range has revealed the existence of two exothermic effects, above and below ca. 817 °C, see Fig. 1 (right). The first, at 835–817 °C, and the second, at 817–737 °C, exothermic effects correspond to the emergence of Fe_3BO_6 and FeBO_3 phases, respectively.

Fig. 1 DTA cooling curve for the composition used for the synthesis of FeBO_3 in wide temperature range (*left*) and in a more restricted range of emergence of FeBO_3 and Fe_3BO_6 phases (*right*)



The obtained DTA thermograms have proved highly efficient for studying phase transitions in the solution-melt systems. In particular, they allowed determining temperature borders of crystallization for $\text{Fe}_x\text{Ga}_{1-x}\text{BO}_3$ crystals in the whole concentration range, $0 \leq x \leq 1$.

Certain shortcomings of the DTA technique, namely high crystallization rate and impossibility of in situ control of the emerging phases, require a complementary analysis method. With this aim, we have used the probe method; it was applied to studying the same compositions as in the case of the DTA. As a probe, a platinum wire of 0.5 mm in diameter, fixed on the lifting device, has been used. The probe was immersed into the melt, then the temperature was modified following a distinct regime in accordance with a predetermined program, and finally, the probe was removed. The microcrystals formed on the probe have been studied by optical and scanning electron microscopies, see Fig. 2. In this way, emergence of the same crystalline phases in the same temperature ranges as with the DTA technique has been confirmed.

Figure 3 shows a typical temperature regime used in synthesizing crystals of a definite composition by solution in the melt technique. It includes the following steps: (1) heating of the furnace, (2) homogenization of the melt, (3) sharp temperature dropping, (4) nucleation and crystal growth, (5) and (6) cooling the furnace.

The synthesized $\text{Fe}_x\text{Ga}_{1-x}\text{BO}_3$ crystals are shown in Fig. 4. All crystals have the shape of hexagonal plates with the dimensions of a few millimeters in the basal plane and ca. 0.1 mm in the perpendicular direction. Their aspect is gradually modified with the iron contents, from almost transparent to green and dark green.

Figure 5 shows the results of XRF analysis; as one can see, x_{crystal} substantially differs from those in the charge x_{charge} ; besides, a considerable difference in iron contents occurs for different crystals extracted from the same crucible. For instance, for $x_{\text{charge}} = 0.01$, x_{crystal} is in the range from 0.002 to 0.04.

Figure 6 shows powder XRD patterns for FeBO_3 and $\text{Fe}_{0.05}\text{Ga}_{0.95}\text{BO}_3$ crystals. In the XRD pattern of pure iron borate FeBO_3 , we have observed lines at $2\theta = 25.43^\circ, 33.38^\circ, 37.23^\circ, 38.95^\circ, 43.39^\circ, 47.11^\circ, 52.21^\circ, 55.09^\circ, 55.78^\circ, 61.58^\circ, 62.73^\circ, 67.04^\circ, 70.12^\circ, 70.54^\circ, 70.72^\circ, 79.31^\circ, 82.51^\circ, 83.12^\circ, 83.65^\circ$ and 89.19° , identifiable as corresponding to the planes (012), (104), (006), (110), (113), (202), (024), (116), (018), (211), (122), (214), (208), (300), (119), (0 0 12), (306), (128), (223) and (312) in the hexagonal system [16].

Isomorphous substitution of iron with gallium produces a change of the lattice parameters manifesting itself in a shift of the XRD peaks from their positions in pure iron borate, see Fig. 6. The hexagonal lattice parameters a and c can be calculated using the following expression [21]:

$$\frac{1}{d^2} = \frac{4h^2 + hk + k^2}{3a^2} + \frac{l^2}{c^2} \quad (1)$$

where d is the interplanar spacing calculated using the Bragg's formula $d = \lambda/2\sin\theta$, and h , k and l represent the lattice planes.

The lattice parameters in $\text{Fe}_x\text{Ga}_{1-x}\text{BO}_3$ crystals have been calculated using the positions of the most "convenient" peaks, e.g., (300) for a and (006) and (0 0 12) for c . Next, these parameters have been refined by graphical extrapolation for all diffraction peaks in the angle range of

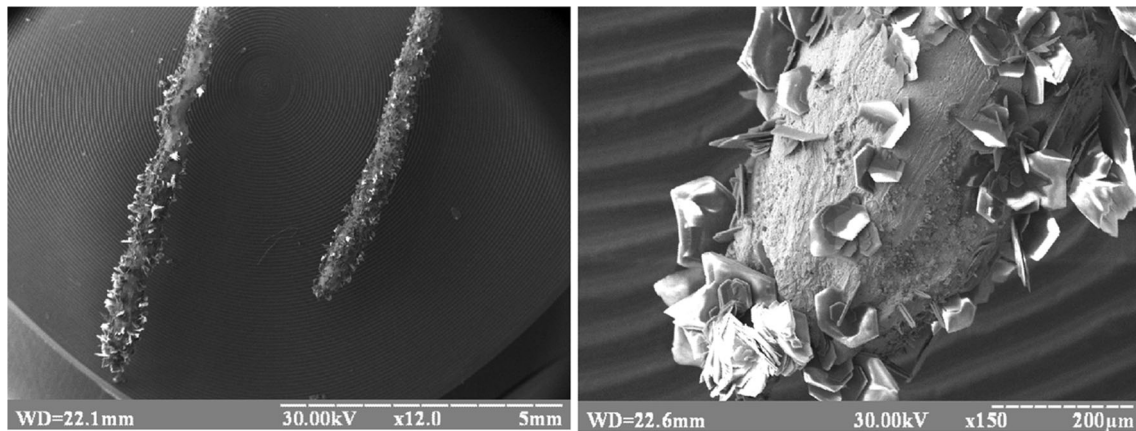


Fig. 2 Electron microscope pictures showing FeBO_3 single crystals on the probes for two different magnifications

Fig. 3 Temperature mode of crystallization used in synthesizing $\text{Fe}_x\text{Ga}_{1-x}\text{BO}_3$ crystals of a definite composition

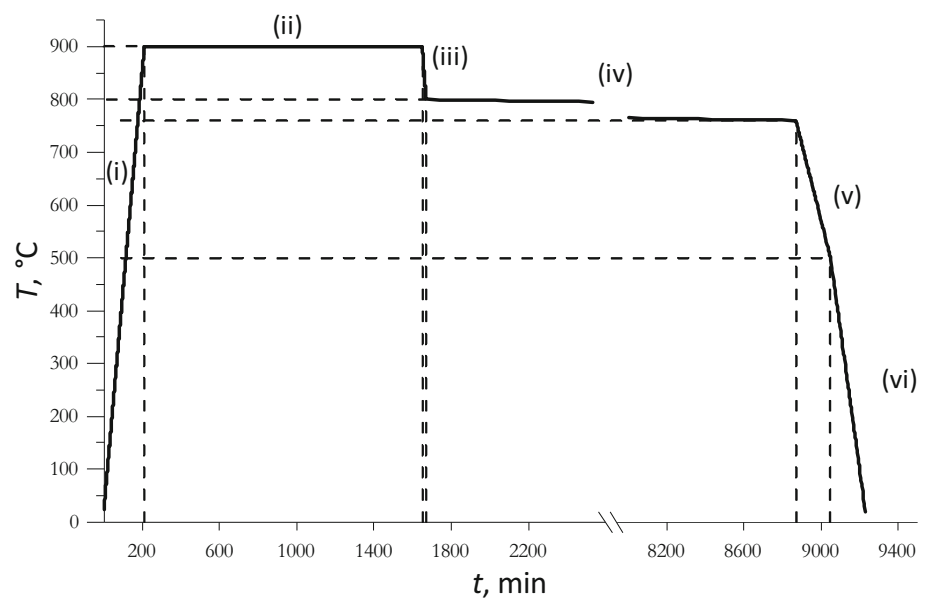


Fig. 4 $\text{Fe}_x\text{Ga}_{1-x}\text{BO}_3$ single crystals with different x (left) and GaBO_3 crystals on the seed holder (right)



$40^\circ < 2\theta < 100^\circ$. The dependences of a and c on the crystal composition are shown in Fig. 7.

According to the Vegard's law [22], suggesting a linear relationship between the lattice parameters of a solid solution and concentrations of its constituent components,

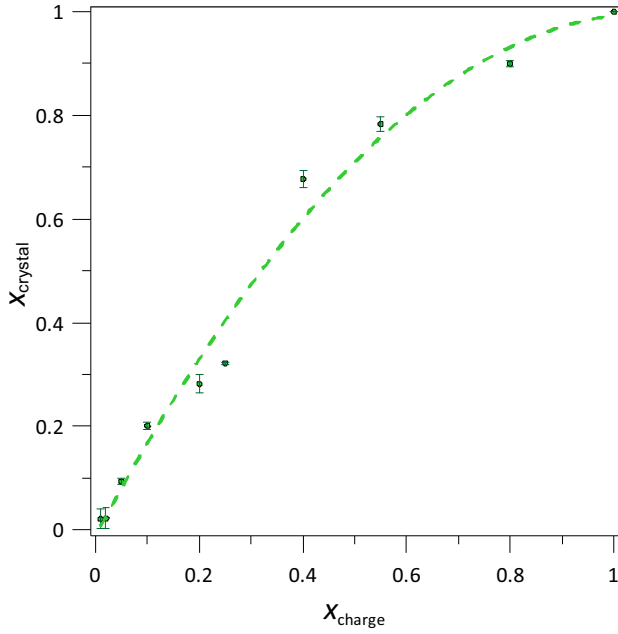
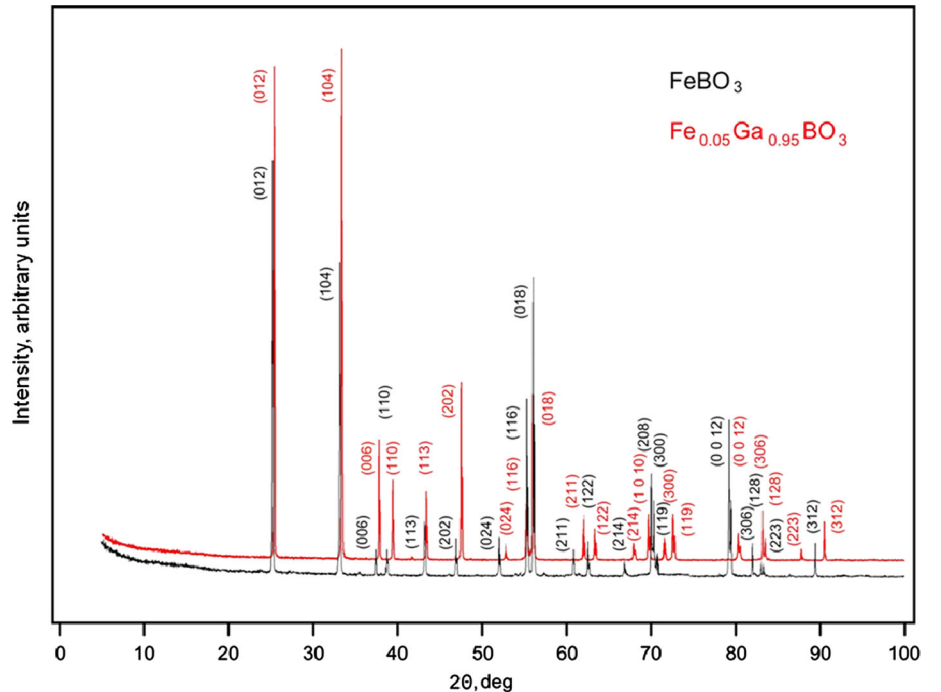


Fig. 5 Iron contents in $\text{Fe}_x\text{Ga}_{1-x}\text{BO}_3$ crystals extracted from one and the same crucible versus initial iron contents in the charge. The dashed line is a guide for the eyes

Fig. 6 XRD patterns of FeBO_3 (bottom curve) and $\text{Fe}_{0.05}\text{Ga}_{0.95}\text{O}_3$ (top curve)



the following relation should hold for $\text{Fe}_x\text{Ga}_{1-x}\text{BO}_3$ crystals:

$$\begin{aligned} a_{\text{Fe}_x\text{Ga}_{1-x}\text{BO}_3} &= x_{\text{crystal}}a_{\text{FeBO}_3} + (1 - x_{\text{crystal}})a_{\text{GaBO}_3}, \\ c_{\text{Fe}_x\text{Ga}_{1-x}\text{BO}_3} &= x_{\text{crystal}}c_{\text{FeBO}_3} + (1 - x_{\text{crystal}})c_{\text{GaBO}_3}. \end{aligned} \quad (2)$$

The dashed lines in Fig. 7 are obtained by least square fitting of the experimental data to this relation. One can see that in the limits of errors, this relation quite adequately describes the experimental tendencies.

The experimental EPR spectrum, curve a in Fig. 8, is typical of isolated Fe^{3+} in oxygen environment in the case where the Zeeman interaction is comparable with the quadrupole fine structure. The spin Hamiltonian parameters have been determined by computer simulating the EPR spectra with laboratory-made code based on a generalized spin Hamiltonian [23]. Besides, no structurally non-equivalent iron sites have been detected, proving the absence of twinning in the crystals.

In the presence of certain local disorder in the crystals, all coordinates of the atoms surrounding the paramagnetic ions undergo random site-to-site distributions. Consequently, the spin Hamiltonian parameters also are distributed, resulting in broadening of certain resonance lines and concomitant decrease in their amplitudes. To estimate the degree of local disorder, we have put forward a simulation code on the basis of the general spin Hamiltonian expressed by means of tesseral spherical tensor operators $\mathcal{T}_{lm}^{ls}(\mathbf{n}, \mathbf{S})$ [23–26]:

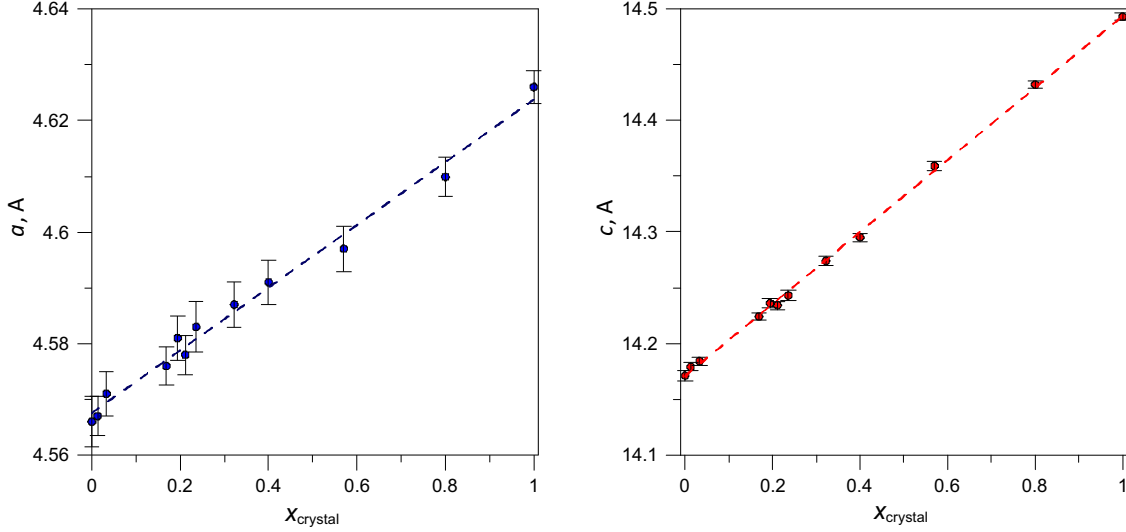


Fig. 7 Concentration dependences of lattice parameters a (left) and c (right). Dashed lines indicate least square fittings

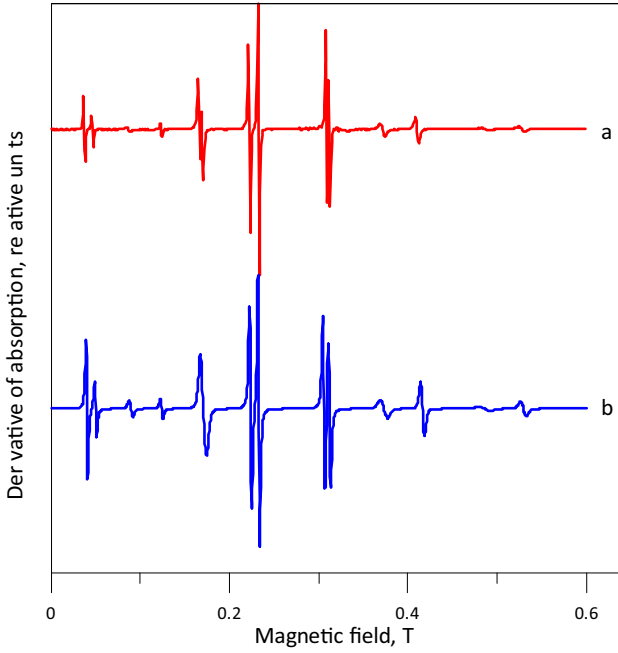


Fig. 8 Experimental (a) and computer generated (b) EPR spectra. The simulation has been made with $\sigma = 0.0003 \text{ \AA}$

$$\mathcal{H} = \sum_{l=2,4} \sum_{m=l}^l \mathcal{B}_{lm}^{0l} \mathcal{T}_{lm}(\mathbf{S}) + g_e \beta \mathbf{B} \sum_{l_s=1,3,5} \sum_{l=\begin{cases} l_s & 1 \\ l_s+1 \end{cases}}^l \sum_{m=l}^l \mathcal{B}_{lm}^{1l_s} \mathcal{T}_{lm}^{1l_s}(\mathbf{n}, \mathbf{S}) \quad (3)$$

where the first and second terms on the right-hand side are, respectively, zero-field and linear Zeeman spin Hamiltonians; the \mathcal{B}_{lm}^{0l} parameters are proportional to the

corresponding Stevens parameters B_l^n [24], β is the Bohr magneton, \mathbf{n} is the unit vector of the magnetizing field \mathbf{B} , \mathbf{S} is the electron spin, and g_e is the free electron g -factor.

The advantage of using the spin Hamiltonian (3) is that it allows explicitly relating the disorder-caused distributions of atomic positions with those of the spin Hamiltonian parameters. As far as $\mathcal{B}_{lm}^{l_B l_S}$ are components of irreducible tensors of rank l [24], they can be consistently expressed through the superposition model [27], as follows, for details, see [25, 26]:

$$\mathcal{B}_{lm}^{l_B l_S} = \sum_{j=1}^n \bar{b}_{lm}^{l_B l_S}(r_j) C_l^m(\vartheta_j, \varphi_j) \quad (4)$$

where j enumerates the nearest neighbors of the paramagnetic ion (ligands) with spherical coordinates r_j , ϑ_j and φ_j and $C_l^m(\vartheta_j, \varphi_j)$ are tesseral coordination factors. For simplicity, we assume that the radial functions $\bar{b}_{lm}^{l_B l_S}(r_j)$ depend only on l_B and l_S and have the radial dependence given by the relation $\bar{b}_{lm}^{l_B l_S}(r_j) = \bar{b}_{l_B l_S} \times (r_0/r_j)^{t_{l_B l_S}}$, with the intrinsic parameters $\bar{b}_{l_B l_S}$ and power exponents $t_{l_B l_S}$ considered as fitting parameters.

In order to numerically estimate the amount of local disorder, we have attributed different normal random site-to-site distributions to the Cartesian coordinates of Fe^{3+} ligands, and we have calculated the parameter distributions of the spin Hamiltonian (3) by means of the superposition model. Figure 9 shows computer simulations of a chosen EPR line, for different amounts of local disorder estimated from the width, σ of the distributions of the ligand coordinates. As one can see, with increasing the disorder, the width of this line drastically increases, and its amplitude decreases, so that the integrated line intensity remains

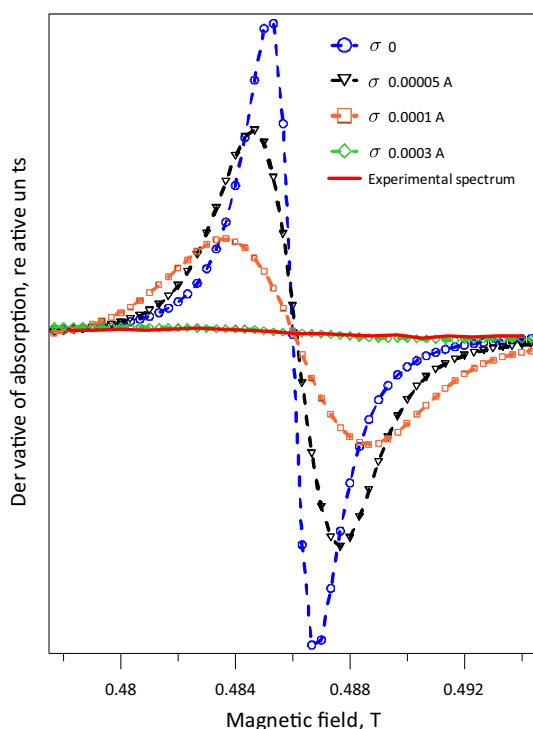


Fig. 9 Experimental EPR line at ca. 0.486 T and the corresponding part of computer generated spectrum with different distribution widths of the ligand coordinates

roughly the same. The best agreement between the experimental and computer-generated line profile is attained at $\sigma \approx 0.0003 \text{ \AA}$. Besides, Fig. 8 shows that with such parameter distributions, a good fitting, curve b, is achieved for the whole experimental spectrum.

4 Conclusions

We have succeeded in obtaining high-quality $\text{Fe}_x\text{Ga}_{1-x}\text{BO}_3$ single crystals in the whole range of concentrations, using the solution in the melt technique for isomorphous substitution in iron-gallium borate system. By means of differential thermal analysis and the probe method, we have determined the component ratios in the charge and the corresponding temperature modes optimal for the crystal synthesis for different x .

The XRF analysis has revealed a certain dispersion of the contents of Fe^{3+} and Ga^{3+} ions in the crystals extracted from the same crucible. The XRD analysis has shown that diamagnetic dilution modifies the parameters of crystal lattice of gallium borate according to the Vegard's law.

The EPR studies have allowed evaluating the degree of local disorder, and its value (the distribution width of atom coordinates estimated as $\sigma \approx 0.0003 \text{ \AA}$) confirms a high quality of the crystals.

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