



HAL
open science

The amplitude equation for weakly nonlinear reversible phase boundaries

Sylvie Benzoni-Gavage, Jean-François Coulombel

► **To cite this version:**

Sylvie Benzoni-Gavage, Jean-François Coulombel. The amplitude equation for weakly nonlinear reversible phase boundaries. [Research Report] CNRS. 2015. hal-01208192

HAL Id: hal-01208192

<https://hal.science/hal-01208192v1>

Submitted on 2 Oct 2015

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

The amplitude equation for weakly nonlinear reversible phase boundaries

Sylvie BENZONI-GAVAGE[†] & Jean-François COULOMBEL[‡]

[†] Université de Lyon, Université Claude Bernard Lyon 1,
CNRS, UMR5208, Institut Camille Jordan, 43 boulevard du 11 novembre 1918
F-69622 Villeurbanne-Cedex, France

[‡] CNRS, Université de Nantes, Laboratoire de Mathématiques Jean Leray (CNRS UMR6629)
2 rue de la Houssinière, BP 92208, 44322 Nantes Cedex 3, France

Emails: benzoni@math.univ-lyon1.fr, jean-francois.coulombel@univ-nantes.fr

October 2, 2015

Abstract

This technical note is a complement to an earlier paper [Benzoni-Gavage & Rosini, *Comput. Math. Appl.* 2009], which aims at a deeper understanding of a basic model for propagating phase boundaries that was proved to admit surface waves [Benzoni-Gavage, *Nonlinear Anal.* 1998]. The amplitude equation governing the evolution of weakly nonlinear surface waves for that model is computed explicitly, and is eventually found to have enough symmetry properties for the associated Cauchy problem to be locally well-posed.

1 Introduction

The reader is assumed to be familiar with the weakly nonlinear theory developed in [3], which is very much inspired from the seminal work [4]. In particular, we keep the same notation as in [3, Section 2], and consider the liquid-vapor phase transition model described in [3, Paragraphs 3.1, 3.2]. Our goal is to make explicit the amplitude equation, which corresponds to [3, Equation (2.20)] in an abstract framework. We adopt slightly different conventions compared with [3, Paragraph 3.3]. Namely, the eigenmodes with nonzero real part for the state ahead of the phase transition are¹:

$$\beta_1^- := \frac{a_\ell - i u_\ell \eta_0}{c_\ell^2 - u_\ell^2}, \quad \beta_1^+ := \frac{-a_\ell - i u_\ell \eta_0}{c_\ell^2 - u_\ell^2} = -\overline{\beta_1^-}, \quad a_\ell := -c_\ell \sqrt{(c_\ell^2 - u_\ell^2) |\check{\eta}|^2 - \eta_0^2}. \quad (1)$$

The eigenmodes with nonzero real part for the state behind the phase transition are:

$$\beta_2^- := \frac{-a_r + i u_r \eta_0}{c_r^2 - u_r^2}, \quad \beta_2^+ := \frac{a_r + i u_r \eta_0}{c_r^2 - u_r^2} = -\overline{\beta_2^-}, \quad a_r := c_r \sqrt{(c_r^2 - u_r^2) |\check{\eta}|^2 - \eta_0^2}. \quad (2)$$

¹The main difference here with [3, Paragraph 3.3] is our sign convention for a_ℓ , the latter being denoted by α_ℓ in [3, Paragraph 3.3].

The remaining, purely imaginary, eigenmodes are:

$$\beta_3^+ = \dots = \beta_{d+1}^+ := \frac{i \eta_0}{u_\ell}, \quad \beta_3^- = \dots = \beta_{d+1}^- := -\frac{i \eta_0}{u_r}.$$

The corresponding right eigenvectors are:

$$R_1^\pm := \begin{pmatrix} r_1^\pm \\ 0 \end{pmatrix}, \quad r_1^- := \begin{pmatrix} -i \eta_0 + u_\ell \beta_1^- \\ i c_\ell^2 \check{\eta} \\ -a_\ell \end{pmatrix}, \quad r_1^+ := \begin{pmatrix} i \eta_0 - u_\ell \beta_1^+ \\ -i c_\ell^2 \check{\eta} \\ -a_\ell \end{pmatrix} = \overline{r_1^-}, \quad (3)$$

$$R_2^\pm := \begin{pmatrix} 0 \\ r_2^\pm \end{pmatrix}, \quad r_2^- := \begin{pmatrix} -i \eta_0 - u_r \beta_2^- \\ i c_r^2 \check{\eta} \\ -a_r \end{pmatrix}, \quad r_2^+ := \begin{pmatrix} i \eta_0 + u_r \beta_2^+ \\ -i c_r^2 \check{\eta} \\ -a_r \end{pmatrix} = \overline{r_2^-}, \quad (4)$$

$$R_j^\pm := \begin{pmatrix} r_j^\pm \\ 0 \end{pmatrix}, \quad R_j^- := \begin{pmatrix} 0 \\ r_j^- \end{pmatrix}, \quad j = 3, \dots, d+1, \\ r_j^+ := \begin{pmatrix} 0 \\ \eta_0 \check{e}_{j-2} \\ u_\ell \check{\eta} \cdot \check{e}_{j-2} \end{pmatrix}, \quad r_j^- := \begin{pmatrix} 0 \\ \eta_0 \check{e}_{j-2} \\ u_r \check{\eta} \cdot \check{e}_{j-2} \end{pmatrix}, \quad (5)$$

with $\check{e}_1 := \check{\eta}$ and the $d-2$ vectors $\check{e}_2, \dots, \check{e}_{d-1} \in \mathbb{R}^{d-1}$ span $\check{\eta}^\perp$.

The left eigenvectors are:

$$L_1^\pm := \begin{pmatrix} \ell_1^\pm \\ 0 \end{pmatrix}, \quad \ell_1^- := \frac{c_\ell^2 - u_\ell^2}{2 a_\ell (u_\ell a_\ell + i c_\ell^2 \eta_0)} \begin{pmatrix} i \eta_0 - 2 u_\ell \beta_1^+ \\ -i \check{\eta} \\ \beta_1^+ \end{pmatrix}, \\ \ell_1^+ := \frac{c_\ell^2 - u_\ell^2}{2 a_\ell (u_\ell a_\ell - i c_\ell^2 \eta_0)} \begin{pmatrix} -i \eta_0 + 2 u_\ell \beta_1^- \\ i \check{\eta} \\ -\beta_1^- \end{pmatrix} = \overline{\ell_1^-},$$

$$L_2^\pm := \begin{pmatrix} 0 \\ \ell_2^\pm \end{pmatrix}, \quad \ell_2^- := \frac{c_r^2 - u_r^2}{2 a_r (u_r a_r + i c_r^2 \eta_0)} \begin{pmatrix} -i \eta_0 - 2 u_r \beta_2^+ \\ i \check{\eta} \\ \beta_2^+ \end{pmatrix}, \\ \ell_2^+ := \frac{c_r^2 - u_r^2}{2 a_r (u_r a_r - i c_r^2 \eta_0)} \begin{pmatrix} i \eta_0 + 2 u_r \beta_2^- \\ -i \check{\eta} \\ -\beta_2^- \end{pmatrix} = \overline{\ell_2^-},$$

$$L_j^\pm := \begin{pmatrix} \ell_j^\pm \\ 0 \end{pmatrix}, \quad L_j^- := \begin{pmatrix} 0 \\ \ell_j^- \end{pmatrix}, \quad j = 3, \dots, d+1, \\ \ell_3^+ := \frac{1}{\eta_0^2 + u_\ell^2 |\check{\eta}|^2} \begin{pmatrix} u_\ell \\ -\eta_0 / (u_\ell |\check{\eta}|^2) \check{\eta} \\ -1 \end{pmatrix}, \quad \ell_3^- := \frac{1}{\eta_0^2 + u_r^2 |\check{\eta}|^2} \begin{pmatrix} -u_r \\ \eta_0 / (u_r |\check{\eta}|^2) \check{\eta} \\ 1 \end{pmatrix}, \\ \ell_j^+ := -\frac{1}{u_\ell \eta_0} \begin{pmatrix} 0 \\ \check{e}'_{j-2} \\ 0 \end{pmatrix}, \quad \ell_j^- := \frac{1}{u_r \eta_0} \begin{pmatrix} 0 \\ \check{e}'_{j-2} \\ 0 \end{pmatrix}, \quad j = 4, \dots, d+1,$$

where the $d - 2$ vectors $\check{e}'_2, \dots, \check{e}'_{d-1}$ belong to $\check{\eta}^\perp$ and form the dual basis of $\check{e}_2, \dots, \check{e}_{d-1}$. Unlike the choice in [3, page 1476], the left eigenvectors here satisfy the normalization property:

$$(L_i^\pm)^* \check{\mathbb{A}}^d(\underline{v}) R_j^\pm = \delta_{i,j}, \quad (L_i^\pm)^* \check{\mathbb{A}}^d(\underline{v}) R_j^\mp = 0.$$

After linearizing the jump conditions we are left with the matrices

$$H(\underline{v}) := \begin{pmatrix} 0 & 0 & 1 & 0 & 0 & -1 \\ 0 & u_\ell I_{d-1} & 0 & 0 & -u_r I_{d-1} & 0 \\ c_\ell^2 - u_\ell^2 & 0 & 2u_\ell & -(c_r^2 - u_r^2) & 0 & -2u_r \\ (c_\ell^2 - u_\ell^2)u_\ell & 0 & u_\ell^2 + \mu & -(c_r^2 - u_r^2)u_r & 0 & -u_r^2 - \mu \end{pmatrix}, \quad (6)$$

where

$$\mu := \frac{1}{2} u_\ell^2 + g(\rho_\ell) = \frac{1}{2} u_r^2 + g(\rho_r), \quad (7)$$

and

$$J(\underline{v}) \eta := \begin{pmatrix} [\rho] \eta_0 \\ [p] \check{\eta} \\ 0 \\ (\mu [\rho] - [p]) \eta_0 \end{pmatrix}. \quad (8)$$

Here we have used the relation

$$\frac{1}{2} j [u] + [\rho f] = \mu [\rho] - [p],$$

where μ is defined in (7), in order to simplify the last entry of $J(\underline{v}) \eta$.

The Lopatinskii determinant is defined by

$$\Delta(\eta) := \det (J(\underline{v}) \eta \quad H(\underline{v}) R_1^- \quad \dots \quad H(\underline{v}) R_{d+1}^-). \quad (9)$$

Each column vector in the above determinant is computed by using (8), (6) and (3), (4), (5) for the definition of the right eigenvectors. Then some simple manipulations on the rows and columns of the above determinant yield (see [1] for similar computations):

$$\Delta(\eta) = (-u_r \eta_0)^{d-2} u_r \det (\check{e}_1 \quad \dots \quad \check{e}_{d-1}) \det \begin{pmatrix} [\rho] \eta_0 & a_\ell & a_r & |\check{\eta}|^2 \\ 0 & u_\ell a_\ell + i \eta_0 c_\ell^2 & u_r a_r + i \eta_0 c_r^2 & 2 u_r |\check{\eta}|^2 \\ -[p] \eta_0 & i \eta_0 u_\ell c_\ell^2 & i \eta_0 u_r c_r^2 & u_r^2 |\check{\eta}|^2 \\ [p] & -i u_\ell c_\ell^2 & -i u_r c_r^2 & \eta_0 \end{pmatrix}.$$

For future use, we introduce the quantity

$$\Upsilon := (-u_r \eta_0)^{d-2} u_r \det (\check{e}_1 \quad \dots \quad \check{e}_{d-1}) \in \mathbb{R} \setminus \{0\}. \quad (10)$$

The expression of the Lopatinskii determinant then reduces to

$$\Delta(\eta) = -[\rho] [u] \Upsilon (\eta_0^2 + u_r^2 |\check{\eta}|^2) (u_\ell u_r a_\ell a_r + c_\ell^2 c_r^2 \eta_0^2). \quad (11)$$

We fix a root $\underline{\eta}$ of the Lopatinskii determinant (see [1] for the properties of such roots, in particular the location of $|\underline{\eta}_0|/|\underline{\check{\eta}}|$ with respect to various velocities associated with the phase transition). From now on, an underline refers to evaluation at the root $\underline{\eta}$ of the Lopatinskii determinant.

We now define a vector $\sigma \in \mathbb{C}^{d+2}$ by computing some of the minors of $\Delta(\eta)$. More precisely, the vectors $H(\underline{v}) \underline{R}_1^-, \dots, H(\underline{v}) \underline{R}_{d+1}^-$ are linearly independent and we can thus define a vector $\sigma \in \mathbb{C}^{d+2} \setminus \{0\}$ satisfying

$$\forall X \in \mathbb{C}^{d+2}, \quad \det \begin{pmatrix} X & H(\underline{v}) \underline{R}_1^- & \dots & H(\underline{v}) \underline{R}_{d+1}^- \end{pmatrix} = \sigma^* X.$$

This vector σ can be computed explicitly by performing some elementary manipulations on each minor of (9). We do not give the detailed calculations but rather give the expression of σ . We find

$$\sigma^* = \underline{\Upsilon} \begin{pmatrix} D_1 & \check{D} \check{\eta}^T & D_{d+1} & D_{d+2} \end{pmatrix}, \quad (12)$$

where $\underline{\Upsilon}$ denotes the quantity Υ in (10) evaluated at the frequency η , and²

$$\begin{aligned} D_1 + \mu D_{d+2} &:= -(\eta_0^2 + u_r^2 |\check{\eta}|^2) ([u] c_\ell^2 c_r^2 \eta_0 - i u_\ell u_r (c_r^2 \underline{a}_\ell - c_\ell^2 \underline{a}_r)), \\ \check{D} &:= -[u] u_r (\underline{a}_\ell (u_r \underline{a}_r - i c_r^2 \eta_0) + \underline{a}_r (u_\ell \underline{a}_\ell - i c_\ell^2 \eta_0)), \\ D_{d+1} &:= -i (\eta_0^2 + u_r^2 |\check{\eta}|^2) (u_r c_r^2 \underline{a}_\ell - u_\ell c_\ell^2 \underline{a}_r), \\ D_{d+2} &:= [u] \eta_0 (\underline{a}_\ell \underline{a}_r + c_\ell^2 c_r^2 |\check{\eta}|^2) \\ &\quad + i \eta_0^2 (c_r^2 \underline{a}_\ell - c_\ell^2 \underline{a}_r) - i |\check{\eta}|^2 (u_r (u_\ell - 2 u_r) c_r^2 \underline{a}_\ell + u_\ell u_r c_\ell^2 \underline{a}_r). \end{aligned} \quad (13)$$

It remains to compute the coefficients γ_1, γ_2 satisfying:

$$J(\underline{v}) \eta + \gamma_1 H(\underline{v}) \underline{R}_1^- + \gamma_2 H(\underline{v}) \underline{R}_2^- = 0.$$

Observe that our convention differs from that in [3, page 1477]. We get:

$$\gamma_1 := \frac{[\rho] u_r \eta_0}{u_r \underline{a}_\ell - i c_\ell^2 \eta_0} = \frac{-i [\rho] c_r^2 \eta_0^2}{\underline{a}_\ell (u_\ell \underline{a}_r - i c_r^2 \eta_0)}, \quad \gamma_2 := \frac{-[\rho] u_\ell \eta_0}{u_\ell \underline{a}_r - i c_r^2 \eta_0} = \frac{i [\rho] c_\ell^2 \eta_0^2}{\underline{a}_r (u_r \underline{a}_\ell - i c_\ell^2 \eta_0)}, \quad (14)$$

where the equalities follow from the relation $u_\ell u_r \underline{a}_\ell \underline{a}_r + c_\ell^2 c_r^2 \eta_0^2 = 0$ that is satisfied by the root η of the Lopatinskii determinant. For notational convenience, we also set

$$\gamma_3 = \dots = \gamma_{d+1} := 0.$$

Following [3, Proposition 2.2], the evolution of a weakly nonlinear phase transition is governed by a scalar amplitude w obeying a nonlocal Burgers equation:

$$a_0(k) \partial_\tau \widehat{w}(\tau, k) + \int_{\mathbb{R}} a_1(k - k', k') \widehat{w}(\tau, k - k') \widehat{w}(\tau, k') dk' = 0, \quad (15)$$

where a_0 and a_1 are given by Equations (2.24) and (2.25) in [3, page 1471]. With the present notation, this yields

$$a_0(k) = \begin{cases} \alpha_0 / (i k) & \text{if } k > 0, \\ \overline{\alpha_0} / (i k) & \text{if } k < 0, \end{cases}$$

and α_0 is a complex number whose definition is recalled in Equation (17) below. The expression of the kernel a_1 is recalled and made explicit in Section 3 below.

²Since \underline{a}_ℓ is negative and \underline{a}_r is positive, D_{d+1} is nonzero and we can thus check that σ is a nonzero vector.

2 Computation of the coefficient α_0

Proposition 1. *The coefficient α_0 in the expression of a_0 is given by*

$$\alpha_0 = -\frac{[\rho] [u] \Upsilon}{\eta_0} (\eta_0^2 + u_r^2 |\check{\eta}|^2) \left\{ u_\ell^2 u_r^2 \left(\frac{a_\ell^2}{c_\ell^2} + \frac{a_r^2}{c_r^2} \right) + 2 c_\ell^2 c_r^2 \eta_0^2 \right\}, \quad (16)$$

and it coincides with the derivative of the Lopatinskii determinant Δ with respect to η_0 at its root $\underline{\eta}$. In particular, α_0 is a nonzero real number.

Proof. We recall that the expression of α_0 is

$$\alpha_0 = \sigma^* [\tilde{f}_0(\underline{v})] + i \sigma^* H(\underline{v}) \underline{R}_p^+ \frac{(\underline{L}_p^+)^* \gamma_q \underline{R}_q^-}{\underline{\beta}_p^+ - \underline{\beta}_q^-}, \quad (17)$$

where we use Einstein's summation convention over repeated indices. We first observe that the Hermitian product $(\underline{L}_p^+)^* \underline{R}_q^-$, $q = 1, 2$, vanishes as soon as p is larger than 4. In the same way, the products $(\underline{L}_1^+)^* \underline{R}_2^-$, $(\underline{L}_3^+)^* \underline{R}_2^-$ and $(\underline{L}_2^+)^* \underline{R}_1^-$ vanish so the expression of α_0 reduces to the sum of four terms:

$$\alpha_0 = \sigma^* [\tilde{f}_0(\underline{v})] + i \sigma^* H(\underline{v}) \underline{R}_2^+ \frac{(\underline{L}_2^+)^* \gamma_2 \underline{R}_2^-}{\underline{\beta}_2^+ - \underline{\beta}_2^-} + i \sigma^* H(\underline{v}) \underline{R}_1^+ \frac{(\underline{L}_1^+)^* \gamma_1 \underline{R}_1^-}{\underline{\beta}_1^+ - \underline{\beta}_1^-} + i \sigma^* H(\underline{v}) \underline{R}_3^+ \frac{(\underline{L}_3^+)^* \gamma_1 \underline{R}_1^-}{\underline{\beta}_3^+ - \underline{\beta}_1^-}. \quad (18)$$

We now compute each of these four quantities separately. Using

$$[\tilde{f}_0(\underline{v})] = \frac{1}{\eta_0} J(\underline{v}) \underline{\eta} - \frac{1}{\eta_0} \begin{pmatrix} 0 \\ [p] \check{\eta} \\ 0 \\ 0 \end{pmatrix},$$

and the orthogonality relation $\sigma^* J(\underline{v}) \underline{\eta} = 0$, we get

$$\begin{aligned} \sigma^* [\tilde{f}_0(\underline{v})] &= -[\rho] \Upsilon \frac{u_\ell u_r |\check{\eta}|^2}{\eta_0} \check{D} \\ &= [\rho] [u] \Upsilon \left\{ \frac{u_\ell u_r^2 |\check{\eta}|^2}{\eta_0} (\underline{a}_\ell (u_r \underline{a}_r - i c_r^2 \eta_0) + \underline{a}_r (u_\ell \underline{a}_\ell - i c_\ell^2 \eta_0)) \right\}. \end{aligned} \quad (19)$$

We now turn to the second term on the right in (18). There holds

$$\frac{(\underline{L}_2^+)^* \gamma_2 \underline{R}_2^-}{\underline{\beta}_2^+ - \underline{\beta}_2^-} = -\frac{c_r^2 |\check{\eta}|^2 (u_r \underline{a}_r - i c_r^2 \eta_0)}{2 \underline{a}_r^2 (\eta_0^2 + u_r^2 |\check{\eta}|^2)} \gamma_2,$$

and we also compute

$$\begin{aligned} i \sigma^* H(\underline{v}) \underline{R}_2^+ &= i \sigma^* \overline{H(\underline{v}) \underline{R}_2^-} = 2 \Upsilon c_r^2 \left\{ \eta_0 (D_{d+1} + u_r D_{d+2}) - u_r |\check{\eta}|^2 \check{D} \right\} \\ &= 2 [u] \Upsilon c_r^2 \underline{a}_r (\eta_0^2 + u_r^2 |\check{\eta}|^2) (u_r \underline{a}_\ell - i c_\ell^2 \eta_0), \end{aligned} \quad (20)$$

where the latter relation is obtained by using the expressions (13). Recalling the expression (14) of γ_2 , we obtain

$$i \sigma^* H(\underline{v}) \underline{R}_2^+ \frac{(\underline{L}_2^+)^* \gamma_2 \underline{R}_2^-}{\underline{\beta}_2^+ - \underline{\beta}_2^-} = [\rho] [u] \Upsilon \left\{ i \frac{u_\ell u_r \underline{a}_\ell}{\underline{a}_r} c_r^2 |\check{\eta}|^2 (u_r \underline{a}_r - i c_r^2 \eta_0) \right\}. \quad (21)$$

We now examine the third term on the right in (18). There holds:

$$\frac{(\underline{L}_1^+)^* \gamma_1 \underline{R}_1^-}{\underline{\beta}_1^+ - \underline{\beta}_1^-} = -\frac{c_\ell^2 |\check{\eta}|^2 (u_\ell \underline{a}_\ell - i c_\ell^2 \eta_0)}{2 \underline{a}_\ell^2 (\eta_0^2 + u_\ell^2 |\check{\eta}|^2)} \gamma_1,$$

and we also compute

$$\begin{aligned} i \sigma^* H(\underline{v}) \underline{R}_1^+ &= i \sigma^* \overline{H(\underline{v}) \underline{R}_1^-} = 2 \Upsilon c_\ell^2 \left\{ u_\ell |\check{\eta}|^2 \check{D} - \eta_0 (D_{d+1} + u_\ell D_{d+2}) \right\} \\ &= -2 [u] \Upsilon c_\ell^2 \underline{a}_\ell (\eta_0^2 + u_r^2 |\check{\eta}|^2) (u_\ell \underline{a}_r - i c_r^2 \eta_0). \end{aligned} \quad (22)$$

Using the expression (14) of γ_1 , we obtain

$$i \sigma^* H(\underline{v}) \underline{R}_1^+ \frac{(\underline{L}_1^+)^* \gamma_1 \underline{R}_1^-}{\underline{\beta}_1^+ - \underline{\beta}_1^-} = [\rho] [u] \Upsilon \left\{ i \frac{u_\ell u_r \underline{a}_r}{\underline{a}_\ell} c_\ell^2 |\check{\eta}|^2 \frac{\eta_0^2 + u_r^2 |\check{\eta}|^2}{\eta_0^2 + u_\ell^2 |\check{\eta}|^2} (u_\ell \underline{a}_\ell - i c_\ell^2 \eta_0) \right\}. \quad (23)$$

It remains to compute the last term on the right in (18) and to add the four expressions. First we compute

$$\frac{(\underline{L}_3^+)^* \gamma_1 \underline{R}_1^-}{\underline{\beta}_3^+ - \underline{\beta}_1^-} = -\frac{c_\ell^2}{\eta_0^2 + u_\ell^2 |\check{\eta}|^2} \gamma_1,$$

and we also compute

$$\sigma^* H(\underline{v}) \underline{R}_3^+ = -[u] \Upsilon u_\ell |\check{\eta}|^2 (2 D_{d+1} + (u_\ell + u_r) D_{d+2}), \quad (24)$$

where we have used the relation (which amounts to $\sigma^* H(\underline{v}) \underline{R}_3^- = 0$):

$$D_1 + \eta_0 \check{D} + 2 u_r D_{d+1} + (\mu + u_r^2) D_{d+2} = 0.$$

The expression (24) can be factorized by using the definitions (13) of D_{d+1} , D_{d+2} , and we obtain

$$\sigma^* H(\underline{v}) \underline{R}_3^+ = -\frac{[u]^2 \Upsilon u_\ell |\check{\eta}|^2}{\eta_0} (\eta_0^2 - u_\ell u_r |\check{\eta}|^2) (\underline{a}_r (u_r \underline{a}_\ell - i c_\ell^2 \eta_0) + \underline{a}_\ell (u_\ell \underline{a}_r - i c_r^2 \eta_0)). \quad (25)$$

Using (14), we derive the expression

$$i \sigma^* H(\underline{v}) \underline{R}_3^+ \frac{(\underline{L}_3^+)^* \gamma_1 \underline{R}_1^-}{\underline{\beta}_3^+ - \underline{\beta}_1^-} = [\rho] [u] \Upsilon \left\{ i [u] u_\ell c_\ell^2 |\check{\eta}|^2 \frac{\eta_0^2 - u_\ell u_r |\check{\eta}|^2}{\eta_0^2 + u_\ell^2 |\check{\eta}|^2} (u_r \underline{a}_r - i c_r^2 \eta_0) \right\}. \quad (26)$$

According to the decomposition (18), the coefficient α_0 is the sum of the four quantities in (19), (21), (23) and (26). Factorizing $[\rho] [u] \Upsilon$ in each term, we first observe that the imaginary part of the sum equals zero. We can thus simplify α_0 by retaining only the real part of each term. This leads to the expression

$$\begin{aligned} \frac{\eta_0 \alpha_0}{[\rho] [u] \Upsilon} &= -c_\ell^2 c_r^2 \eta_0^2 (u_r^2 + u_\ell u_r) |\check{\eta}|^2 + u_\ell u_r \underline{a}_\ell \frac{c_r^2 \eta_0^2}{\underline{a}_r} c_r^2 |\check{\eta}|^2 \\ &\quad + u_\ell u_r \underline{a}_r \frac{c_\ell^2 \eta_0^2}{\underline{a}_\ell} c_\ell^2 |\check{\eta}|^2 \frac{\eta_0^2 + u_r^2 |\check{\eta}|^2}{\eta_0^2 + u_\ell^2 |\check{\eta}|^2} + (u_\ell u_r - u_\ell^2) |\check{\eta}|^2 c_\ell^2 c_r^2 \eta_0^2 \frac{\eta_0^2 - u_\ell u_r |\check{\eta}|^2}{\eta_0^2 + u_\ell^2 |\check{\eta}|^2}. \end{aligned}$$

At this stage, some elementary manipulations lead to the expression (16) of α_0 .

The link between α_0 and the partial derivative $\partial_{\eta_0} \Delta(\underline{\eta})$ comes from the relation

$$\eta_0 \frac{\partial}{\partial \eta_0} (u_\ell u_r a_\ell a_r + c_\ell^2 c_r^2 \eta_0^2) \Big|_{\underline{\eta}} = u_\ell^2 u_r^2 \left(\frac{a_\ell^2}{c_\ell^2} + \frac{a_r^2}{c_r^2} \right) + 2 c_\ell^2 c_r^2 \eta_0^2,$$

which is obtained by differentiating (11) and the expressions (1), (2) with respect to η_0 and then evaluating at $\underline{\eta}$. \square

Since the coefficient α_0 is real, we obtain $a_0(k) = \alpha_0/(ik)$ for all $k \neq 0$. In particular, the amplitude equation (15) reduces to

$$\partial_\tau \widehat{w}(\tau, k) + \frac{ik}{\alpha_0} \int_{\mathbb{R}} a_1(k - k', k') \widehat{w}(\tau, k - k') \widehat{w}(\tau, k') dk' = 0.$$

As far as smooth solutions are concerned, the Cauchy problem associated with this kind of nonlocal Burgers equation is known to be locally well-posed under rather simple algebraic conditions (see [2]). These conditions are invariant under multiplication by a nonzero real constant, so that it is sufficient to investigate whether they are satisfied by the slightly simpler kernel $4\pi a_1$. This is the purpose of the next section.

3 Computation of the quadratic kernel

We define the kernel $q(k, k') := 4\pi a_1(k, k')$. Following [3], we can decompose q as follows:

$$q(k, k') = \sum_{j=1}^5 q_j(k, k'),$$

where the kernels q_1, \dots, q_5 are given by³

$$q_1(k, k') := \sigma(k + k')^* \sum_{j=0}^{d-1} \underline{\eta}_j \left\{ d\tilde{f}^j(v_r) \cdot (\widehat{r}_+(k, 0) + \widehat{r}_+(k', 0)) - d\tilde{f}^j(v_\ell) \cdot (\widehat{r}_-(k, 0) + \widehat{r}_-(k', 0)) \right\}, \quad (27)$$

$$q_2(k, k') := -\sigma(k + k')^* \left\{ d^2 \tilde{f}^d(v_r) \cdot (\widehat{r}_+(k, 0), \widehat{r}_+(k', 0)) - d^2 \tilde{f}^d(v_\ell) \cdot (\widehat{r}_-(k, 0), \widehat{r}_-(k', 0)) \right\}, \quad (28)$$

$$q_3(k, k') := i(k + k') \int_0^{+\infty} L(k + k', z) d\mathbb{A}(v, \underline{\eta}) \cdot \widehat{r}(k, z) \cdot \widehat{r}(k', z) dz, \quad (29)$$

$$q_4(k, k') := \int_0^{+\infty} L(k + k', z) \frac{\partial}{\partial z} \left(d\check{\mathbb{A}}^d(v) \cdot \widehat{r}(k, z) \cdot \widehat{r}(k', z) \right) dz, \quad (30)$$

$$q_5(k, k') := - \int_0^{+\infty} L(k + k', z) \check{\mathbb{A}}(v, \underline{\eta}) \left(\frac{\partial \widehat{r}}{\partial z}(k, z) + \frac{\partial \widehat{r}}{\partial z}(k', z) \right) dz. \quad (31)$$

We first examine the kernel q_1 and derive its expression for all values of (k, k') .

Lemma 1. *Let us define the quantity*

$$Q := 2[\rho][u] \underline{\Upsilon} (\eta_0^2 + u_r^2 |\underline{\eta}|^2) (\underline{\beta}_1^- + \underline{\beta}_2^-) i u_\ell u_r \underline{a}_\ell \underline{a}_r \frac{u_\ell \underline{a}_r + i c_r^2 \underline{\eta}_0}{u_\ell \underline{a}_r - i c_r^2 \underline{\eta}_0}. \quad (32)$$

³We keep the notation of [3] for the functions $\widehat{r}_\pm, \widehat{r}$ and so on.

Then the kernel q_1 in (27) satisfies

$$q_1(k, k') = \begin{cases} 0 & \text{if } k > 0 \text{ and } k' > 0, \\ \overline{Q} & \text{if } k > 0, k' < 0 \text{ and } k + k' > 0. \end{cases}$$

Proof. The function g^0 is known to be an entropy for the isothermal Euler equations with corresponding flux (g^1, \dots, g^d) . We thus have the relations

$$\forall j = 0, \dots, d, \quad dg^0(v) A^j(v) = dg^j(v),$$

where we use the convention $A^0(v) = I$ for all v . Using this relation between the Jacobian matrices, we get

$$\sum_{j=0}^{d-1} \eta_j d\tilde{f}^j(v_r) \cdot \underline{r}_2^- = \left(\begin{array}{c} \sum_{j=0}^{d-1} \eta_j A^j(v_r) \underline{r}_2^- \\ dg^0(v_r) \sum_{j=0}^{d-1} \eta_j A^j(v_r) \underline{r}_2^- \end{array} \right) = i \beta_2^- \left(\begin{array}{c} A^d(v_r) \underline{r}_2^- \\ dg^d(v_r) \cdot \underline{r}_2^- \end{array} \right) = -i \beta_2^- H(\underline{v}) \underline{R}_2^-.$$

Similarly we have

$$\sum_{j=0}^{d-1} \eta_j d\tilde{f}^j(v_\ell) \cdot \underline{r}_1^- = -i \beta_1^- H(\underline{v}) \underline{R}_1^-.$$

For $k > 0$ and $k' > 0$, we thus get

$$\begin{aligned} q_1(k, k') &= 2\sigma^* \sum_{j=0}^{d-1} \gamma_2 \eta_j d\tilde{f}^j(v_r) \cdot \underline{r}_2^- - 2\sigma^* \sum_{j=0}^{d-1} \gamma_1 \eta_j d\tilde{f}^j(v_\ell) \cdot \underline{r}_1^- \\ &= -2i \gamma_2 \beta_2^- \sigma^* H(\underline{v}) \underline{R}_2^- + 2i \gamma_1 \beta_1^- \sigma^* H(\underline{v}) \underline{R}_1^- = 0, \end{aligned}$$

because σ is orthogonal to both $H(\underline{v}) \underline{R}_1^-$ and $H(\underline{v}) \underline{R}_2^-$.

Let us now consider the case $k > 0, k' < 0$ and $k + k' > 0$. Using the same relations as above for the differentials dg^j , we obtain

$$\begin{aligned} q_1(k, k') &= \sigma^* \sum_{j=0}^{d-1} \eta_j d\tilde{f}^j(v_r) \cdot (\gamma_2 \underline{r}_2^- + \overline{\gamma}_2 \underline{r}_2^+) - \sigma^* \sum_{j=0}^{d-1} \eta_j d\tilde{f}^j(v_\ell) \cdot (\gamma_1 \underline{r}_1^- + \overline{\gamma}_1 \underline{r}_1^+) \\ &= -i \overline{\gamma}_2 \beta_2^+ \sigma^* H(\underline{v}) \underline{R}_2^+ + i \overline{\gamma}_1 \beta_1^+ \sigma^* H(\underline{v}) \underline{R}_1^+. \end{aligned}$$

The Hermitian products $\sigma^* H(\underline{v}) \underline{R}_2^+$ and $\sigma^* H(\underline{v}) \underline{R}_1^+$ have already been computed in the proof of Proposition 1, see (20) and (22). We then obtain

$$\begin{aligned} q_1(k, k') &= -2[u] \Upsilon (\eta_0^2 + u_r^2 |\check{\eta}|^2) c_r^2 \underline{a}_r (u_r \underline{a}_\ell - i c_\ell^2 \eta_0) \overline{\gamma}_2 \beta_2^+ \\ &\quad - 2[u] \Upsilon (\eta_0^2 + u_r^2 |\check{\eta}|^2) c_\ell^2 \underline{a}_\ell (u_\ell \underline{a}_r - i c_r^2 \eta_0) \overline{\gamma}_1 \beta_1^+. \end{aligned}$$

We use the definition (14) to obtain

$$c_r^2 \underline{a}_r (u_r \underline{a}_\ell - i c_\ell^2 \eta_0) \overline{\gamma}_2 = c_\ell^2 \underline{a}_\ell (u_\ell \underline{a}_r - i c_r^2 \eta_0) \overline{\gamma}_1 = i [\rho] c_\ell^2 c_r^2 \eta_0^2 \frac{u_\ell \underline{a}_r - i c_r^2 \eta_0}{u_\ell \underline{a}_r + i c_r^2 \eta_0},$$

and the claim follows using the relation $c_\ell^2 c_r^2 \eta_0^2 = -u_\ell u_r \underline{a}_\ell \underline{a}_r$. \square

Deriving the expression of the kernels q_3, q_4, q_5 requires the expression of the row vector $L(k + k', z)$, which we derive right now.

Lemma 2. *For $k > 0$, there holds*

$$L(k, z) = \left(\frac{\omega_1}{\gamma_1} \exp(-k \underline{\beta}_1^+ z) \tilde{\ell}_1 + \frac{\omega_3}{\gamma_1} \exp(-k \underline{\beta}_3^+ z) \tilde{\ell}_3 - \frac{\omega_2}{\gamma_2} \exp(-k \underline{\beta}_2^+ z) \tilde{\ell}_2 \right),$$

where we have set

$$\begin{aligned} \tilde{\ell}_1 &:= (i \eta_0 - 2 u_\ell \underline{\beta}_1^+ & -i \check{\eta}^T & \underline{\beta}_1^+), & \tilde{\ell}_3 &:= (-u_\ell^2 |\check{\eta}|^2 & \eta_0 \check{\eta}^T & u_\ell |\check{\eta}|^2), \\ \tilde{\ell}_2 &:= (-i \eta_0 - 2 u_r \underline{\beta}_2^+ & i \check{\eta}^T & \underline{\beta}_2^+), \end{aligned}$$

and

$$\begin{aligned} \omega_1 &:= [\rho] [u] \underline{\Upsilon} \frac{\eta_0^2 + u_r^2 |\check{\eta}|^2}{\eta_0^2 + u_\ell^2 |\check{\eta}|^2} i u_\ell \eta_0 (u_r \underline{a}_r - i c_r^2 \eta_0), \\ \omega_3 &:= [\rho] [u]^2 \underline{\Upsilon} \frac{\eta_0^2 - u_\ell u_r |\check{\eta}|^2}{\eta_0^2 + u_\ell^2 |\check{\eta}|^2} (u_r \underline{a}_r - i c_r^2 \eta_0), & \omega_2 &:= [\rho] [u] \underline{\Upsilon} i u_r \eta_0 (u_\ell \underline{a}_\ell - i c_\ell^2 \eta_0). \end{aligned}$$

Proof. We first observe that σ is orthogonal to the vectors $H(\underline{v}) \underline{R}_p^+$ for $p \geq 4$, so for $k > 0$ the expression of $L(k, z)$ reduces to

$$L(k, z) = \sum_{p=1}^3 \sigma^* H(\underline{v}) \underline{R}_p^+ \exp(-k \underline{\beta}_p^+ z) (\underline{L}_p^+)^*.$$

The expression of the products $\sigma^* H(\underline{v}) \underline{R}_p^+$, $p = 1, 2, 3$ can be found in (20), (22), (25), and we use the definitions of the left eigenvectors \underline{L}_p^+ to derive the expressions given in Lemma 2. \square

We now examine the kernel q_5 , which, as q_1 but unlike q_2, q_3, q_4 , does not contain any term in $p''(\rho_{\ell, r})$.

Lemma 3. *With Q defined in (32), the kernel q_5 in (31) satisfies*

$$q_5(k, k') = \begin{cases} 0 & \text{if } k > 0 \text{ and } k' > 0, \\ \frac{Q}{k} k' & \text{if } k > 0, k' < 0 \text{ and } k + k' > 0. \end{cases}$$

Proof. For $k > 0$ and $k' > 0$, we compute

$$\check{\mathbb{A}}(v, \underline{\eta}) \left(\frac{\partial \hat{r}}{\partial z}(k, z) + \frac{\partial \hat{r}}{\partial z}(k', z) \right) = \begin{pmatrix} i \gamma_1 (\underline{\beta}_1^-)^2 (k \exp(k \underline{\beta}_1^- z) + k' \exp(k' \underline{\beta}_1^- z)) A^d(v_\ell) \underline{r}_1^- \\ i \gamma_2 (\underline{\beta}_2^-)^2 (k \exp(k \underline{\beta}_2^- z) + k' \exp(k' \underline{\beta}_2^- z)) A^d(v_r) \underline{r}_2^- \end{pmatrix}.$$

Using the orthogonality properties $\tilde{\ell}_1 A^d(v_\ell) \underline{r}_1^- = \tilde{\ell}_3 A^d(v_\ell) \underline{r}_1^- = \tilde{\ell}_3 A^d(v_r) \underline{r}_2^- = 0$, we get $q_5(k, k') = 0$ if $k > 0$ and $k' > 0$ because the integrand in (31) vanishes.

Let us now consider the case $k > 0$, $k' < 0$ and $k + k' > 0$. From the previous argument, we still find that the term $L(k + k', z) \check{\mathbb{A}}(v, \underline{\eta}) \partial_z \hat{r}(k, z)$ vanishes. We thus get

$$\begin{aligned} q_5(k, k') &= - \int_0^{+\infty} L(k + k', z) \check{\mathbb{A}}(v, \underline{\eta}) \frac{\partial \hat{r}}{\partial z}(k', z) dz \\ &= -i k' \int_0^{+\infty} L(k + k', z) \left(\frac{\gamma_1 (\underline{\beta}_1^+)^2}{\gamma_2 (\underline{\beta}_2^+)^2} \exp(k' \underline{\beta}_1^+ z) A^d(v_\ell) \underline{r}_1^+ \right) dz. \end{aligned}$$

We now use the expression of $L(k+k', z)$ in Lemma 2. The expression of $q_5(k, k')$ is simplified by recalling the orthogonality property $\tilde{\ell}_3 A^d(v_\ell) \underline{r}_1^+ = 0$ and we get

$$\begin{aligned} q_5(k, k') &= -i k' \int_0^{+\infty} \frac{\overline{\gamma}_1}{\gamma_1} \omega_1 \tilde{\ell}_1 A^d(v_\ell) \underline{r}_1^+ (\underline{\beta}_1^+)^2 \exp(-k \underline{\beta}_1^+ z) dz \\ &\quad - i k' \int_0^{+\infty} \frac{\overline{\gamma}_2}{\gamma_2} \omega_2 \tilde{\ell}_2 A^d(v_r) \underline{r}_2^+ (\underline{\beta}_2^+)^2 \exp(-k \underline{\beta}_2^+ z) dz, \\ &= -i \frac{k'}{k} \left\{ \frac{\overline{\gamma}_1}{\gamma_1} \omega_1 \tilde{\ell}_1 A^d(v_\ell) \underline{r}_1^+ \underline{\beta}_1^+ + \frac{\overline{\gamma}_2}{\gamma_2} \omega_2 \tilde{\ell}_2 A^d(v_r) \underline{r}_2^+ \underline{\beta}_2^+ \right\}. \end{aligned}$$

The conclusion of Lemma 3 then follows from the relations

$$\begin{aligned} \frac{\overline{\gamma}_1}{\gamma_1} &= -\frac{\overline{\gamma}_2}{\gamma_2} = -\frac{u_\ell \underline{a}_r - i c_r^2 \eta_0}{u_\ell \underline{a}_r + i c_r^2 \eta_0}, \\ \omega_1 \tilde{\ell}_1 A^d(v_\ell) \underline{r}_1^+ &= -\omega_2 \tilde{\ell}_2 A^d(v_r) \underline{r}_2^+ = 2[\rho][u] \Upsilon (\eta_0^2 + u_r^2 |\check{\eta}|^2) u_\ell u_r \underline{a}_\ell \underline{a}_r. \end{aligned}$$

□

We immediately get

Corollary 1. *With Q defined in (32), the kernels q_1, q_5 in (27), (31) satisfy*

$$(q_1 + q_5)(k, k') = \begin{cases} 0 & \text{if } k > 0 \text{ and } k' > 0, \\ \overline{Q} \left(1 + \frac{k'}{k} \right) & \text{if } k > 0, k' < 0 \text{ and } k + k' > 0. \end{cases}$$

Our goal now is to derive an explicit expression for the kernels q_2, q_3, q_4 in (28), (29), (30). This is more intricate because these kernels are quadratic with respect to the vector $\hat{r}(k, z)$ and there is more algebra involved to obtain a factorized expression in each region of the (k, k') -plane. We begin with some preliminary computations that will be useful in Propositions 2 and 3 below.

Lemma 4. *The coordinates (13) of the vector σ satisfy*

$$\begin{aligned} \gamma_1 \check{D} &= -[\rho][u] u_r \eta_0 (u_r \underline{a}_r - i c_r^2 \eta_0), \\ \gamma_2 \check{D} &= [\rho][u] u_r \eta_0 (u_\ell \underline{a}_\ell - i c_\ell^2 \eta_0), \\ \gamma_1 D_{d+1} &= -[\rho] u_r (\eta_0^2 + u_r^2 |\check{\eta}|^2) (u_\ell \underline{a}_r + i c_r^2 \eta_0), \\ \gamma_2 D_{d+1} &= -[\rho] u_\ell (\eta_0^2 + u_r^2 |\check{\eta}|^2) (u_r \underline{a}_\ell + i c_\ell^2 \eta_0), \\ \gamma_1 D_{d+2} &= [\rho] (\eta_0^2 + u_r^2 |\check{\eta}|^2) (u_r \underline{a}_r + i c_r^2 \eta_0) - [\rho][u] u_r |\check{\eta}|^2 (u_r \underline{a}_r - i c_r^2 \eta_0), \\ \gamma_2 D_{d+2} &= [\rho] (\eta_0^2 + u_r^2 |\check{\eta}|^2) (u_\ell \underline{a}_\ell + i c_\ell^2 \eta_0) + [\rho][u] u_r |\check{\eta}|^2 (u_\ell \underline{a}_\ell - i c_\ell^2 \eta_0). \end{aligned}$$

Proof. From the definition (13), there holds

$$\check{D} = [u] u_r (\underline{a}_r (u_r \underline{a}_\ell - i c_\ell^2 \eta_0) + \underline{a}_\ell (u_\ell \underline{a}_r - i c_r^2 \eta_0)),$$

and we then use the expression (14) of γ_1, γ_2 to compute $\gamma_1 \check{D}$ and $\gamma_2 \check{D}$.

The expressions of $\gamma_{1,2} D_{d+1}$ follow from the observation that D_{d+1} satisfies

$$\begin{aligned} \eta_0 D_{d+1} &= -(\eta_0^2 + u_r^2 |\check{\eta}|^2) (u_r \underline{a}_\ell - i c_\ell^2 \eta_0) (u_\ell \underline{a}_r + i c_r^2 \eta_0) \\ &= (\eta_0^2 + u_r^2 |\check{\eta}|^2) (u_r \underline{a}_\ell + i c_\ell^2 \eta_0) (u_\ell \underline{a}_r - i c_r^2 \eta_0). \end{aligned}$$

We then use again (14).

To compute the product $\gamma_1 D_{d+2}$, we recall the relation (22) which we found in the proof of Proposition 1. It reads

$$\mathfrak{Y}_0 (D_{d+1} + u_\ell D_{d+2}) - u_\ell |\check{\mathfrak{Y}}|^2 \check{D} = [u] \underline{a}_\ell (\mathfrak{Y}_0^2 + u_r^2 |\check{\mathfrak{Y}}|^2) (u_\ell \underline{a}_r - i c_r^2 \mathfrak{Y}_0).$$

We multiply the latter relation by γ_1 and use the previous expressions of $\gamma_1 \check{D}$ and $\gamma_1 D_{d+1}$ to obtain that of $\gamma_1 D_{d+2}$. The expression of $\gamma_2 D_{d+2}$ is then easily deduced by using $\gamma_2/\gamma_1 = i c_\ell^2 \mathfrak{Y}_0^2/(u_r \underline{a}_r)$. \square

We now compute the kernels q_2, q_3, q_4 in the region $\{k > 0, k' > 0\}$.

Proposition 2. *Let us define the quantities*

$$\begin{aligned} Q_\ell &:= [\rho] [u] \underline{\Upsilon} u_\ell u_r \frac{\underline{a}_r}{\underline{a}_\ell} (\mathfrak{Y}_0^2 + u_r^2 |\check{\mathfrak{Y}}|^2) (\mathfrak{Y}_0^2 + u_\ell^2 |\check{\mathfrak{Y}}|^2) \gamma_1 (i \mathfrak{Y}_0 - u_\ell \underline{\beta}_1^-), \\ Q_r &:= [\rho] [u] \underline{\Upsilon} u_\ell u_r \frac{\underline{a}_\ell}{\underline{a}_r} (\mathfrak{Y}_0^2 + u_r^2 |\check{\mathfrak{Y}}|^2)^2 \gamma_2 (i \mathfrak{Y}_0 + u_r \underline{\beta}_2^-), \\ Q_\# &:= 2 [\rho] \underline{\Upsilon} (\mathfrak{Y}_0^2 + u_r^2 |\check{\mathfrak{Y}}|^2) (\mathfrak{Y}_0^2 + u_\ell u_r |\check{\mathfrak{Y}}|^2) i c_\ell^2 c_r^2 \mathfrak{Y}_0 \left(\frac{c_r^2 \gamma_2}{\rho_r u_r} - \frac{c_\ell^2 \gamma_1}{\rho_\ell u_\ell} \right). \end{aligned} \quad (33)$$

Then the kernels q_2, q_3, q_4 defined in (28), (29) and (30) satisfy

$$(q_2 + q_3 + q_4)(k, k') = \left(\frac{p''(\rho_\ell)}{2} + \frac{c_\ell^2}{\rho_\ell} \right) Q_\ell + \left(\frac{p''(\rho_r)}{2} + \frac{c_r^2}{\rho_r} \right) Q_r + Q_\#,$$

for all $k > 0$ and $k' > 0$.

Proof. We first recall the expressions of the second differentials that are involved in the kernels q_2, q_3, q_4 , see [3, page 1480]:

$$\sum_{k=1}^{d-1} \check{\mathfrak{Y}}_k d^2 f^k(v_{\ell,r}) \cdot (v, v) = p''(\rho_{\ell,r}) \begin{pmatrix} 0 \\ \rho^2 \check{\mathfrak{Y}} \\ 0 \end{pmatrix} + \frac{2}{\rho_{\ell,r}} \begin{pmatrix} 0 \\ \check{\mathfrak{Y}} \cdot \check{\mathfrak{J}} \check{\mathfrak{J}} \\ \check{\mathfrak{Y}} \cdot \check{\mathfrak{J}} (j_d - u_{\ell,r} \rho) \end{pmatrix}, \quad (34)$$

$$\begin{aligned} d^2 \tilde{f}^d(v_{\ell,r}) \cdot (v, v) &= \left(\frac{d^2 f^d(v_{\ell,r}) \cdot (v, v)}{d^2 g^d(v_{\ell,r}) \cdot (v, v)} \right) = p''(\rho_{\ell,r}) \begin{pmatrix} 0 \\ 0 \\ \rho^2 \\ u_{\ell,r} \rho^2 \end{pmatrix} + \frac{2}{\rho_{\ell,r}} \begin{pmatrix} 0 \\ (j_d - u_{\ell,r} \rho) \check{\mathfrak{J}} \\ (j_d - u_{\ell,r} \rho)^2 \\ 0 \end{pmatrix} \\ &+ \frac{1}{\rho_{\ell,r}} \begin{pmatrix} 0 \\ 0 \\ 0 \\ 3 u_{\ell,r} (j_d - u_{\ell,r} \rho)^2 - u_{\ell,r} c_{\ell,r}^2 \rho^2 + 2 c_{\ell,r}^2 \rho j_d + u_{\ell,r} \check{\mathfrak{J}} \cdot \check{\mathfrak{J}} \end{pmatrix}. \end{aligned} \quad (35)$$

The proof of Proposition 2 splits in several steps. We assume from now on that k and k' are both positive and we wish to compute the expression of the kernel $q_2 + q_3 + q_4$.

• Step 1: computation of the $p''(\rho_\ell)$ factor. In this first step, we collect all the terms that involve $p''(\rho_\ell)$ in $q_2 + q_3 + q_4$. The contribution of the kernel q_2 equals

$$\begin{aligned} &\underline{\Upsilon} (D_{d+1} + u_\ell D_{d+2}) \gamma_1^2 (i \mathfrak{Y}_0 - u_\ell \underline{\beta}_1^-)^2 \\ &= -[\rho] [u] \underline{\Upsilon} \gamma_1 (i \mathfrak{Y}_0 - u_\ell \underline{\beta}_1^-)^2 \left\{ (\mathfrak{Y}_0^2 + u_r^2 |\check{\mathfrak{Y}}|^2) i c_r^2 \mathfrak{Y}_0 + (u_r \underline{a}_r - i c_r^2 \mathfrak{Y}_0) u_\ell u_r |\check{\mathfrak{Y}}|^2 \right\}, \end{aligned} \quad (36)$$

where we have used Lemma 4 to compute the product $(D_{d+1} + u_\ell D_{d+2}) \gamma_1$.

The contribution of the kernel q_3 equals

$$\begin{aligned} & i(k+k') \int_0^{+\infty} \frac{\omega_1}{\gamma_1} (-i \check{\eta}^T) e^{-(k+k') \underline{\beta}_1^+ z} \gamma_1^2 (i \eta_0 - u_\ell \underline{\beta}_1^-)^2 \check{\eta} e^{(k+k') \underline{\beta}_1^- z} dz \\ & + i(k+k') \int_0^{+\infty} \frac{\omega_3}{\gamma_1} (\eta_0 \check{\eta}^T) e^{-(k+k') \underline{\beta}_3^+ z} \gamma_1^2 (i \eta_0 - u_\ell \underline{\beta}_1^-)^2 \check{\eta} e^{(k+k') \underline{\beta}_1^- z} dz \\ & = \gamma_1 (i \eta_0 - u_\ell \underline{\beta}_1^-)^2 \left\{ \omega_1 \frac{|\check{\eta}|^2}{\underline{\beta}_1^+ - \underline{\beta}_1^-} + \omega_3 |\check{\eta}|^2 \frac{i \eta_0}{\underline{\beta}_3^+ - \underline{\beta}_1^-} \right\}. \end{aligned}$$

Similarly, the contribution of the kernel q_4 reads

$$\begin{aligned} & - \int_0^{+\infty} \frac{\omega_1}{\gamma_1} \underline{\beta}_1^+ e^{-(k+k') \underline{\beta}_1^+ z} \gamma_1^2 (i \eta_0 - u_\ell \underline{\beta}_1^-)^2 (k+k') \underline{\beta}_1^- e^{(k+k') \underline{\beta}_1^- z} dz \\ & - \int_0^{+\infty} \frac{\omega_3}{\gamma_1} u_\ell |\check{\eta}|^2 e^{-(k+k') \underline{\beta}_3^+ z} \gamma_1^2 (i \eta_0 - u_\ell \underline{\beta}_1^-)^2 (k+k') \underline{\beta}_1^- e^{(k+k') \underline{\beta}_1^- z} dz \\ & = - \gamma_1 (i \eta_0 - u_\ell \underline{\beta}_1^-)^2 \left\{ \omega_1 \frac{\underline{\beta}_1^+ \underline{\beta}_1^-}{\underline{\beta}_1^+ - \underline{\beta}_1^-} + \omega_3 |\check{\eta}|^2 \frac{u_\ell \underline{\beta}_1^-}{\underline{\beta}_3^+ - \underline{\beta}_1^-} \right\}. \end{aligned}$$

Adding the contributions of q_3 and q_4 gives the term

$$\gamma_1 (i \eta_0 - u_\ell \underline{\beta}_1^-)^2 \left\{ -\frac{1}{2 \underline{a}_\ell} \left(\frac{\underline{a}_\ell^2}{c_\ell^2} + c_\ell^2 |\check{\eta}|^2 \right) \omega_1 + u_\ell |\check{\eta}|^2 \omega_3 \right\}. \quad (37)$$

We now use the definitions of ω_1 and ω_3 , see Lemma 2, and add the contributions in (36) and (37) in order to obtain the $p''(\rho_\ell)$ term in $q_2 + q_3 + q_4$. We first obtain that the sum of the right hand side of (36) and the expression in (37) equals

$$-[\rho][u] \Upsilon \gamma_1 (i \eta_0 - u_\ell \underline{\beta}_1^-)^2 \frac{\eta_0^2 + u_r^2 |\check{\eta}|^2}{2 c_\ell^2 (\eta_0^2 + u_\ell^2 |\check{\eta}|^2)} (u_\ell \underline{a}_\ell + i c_\ell^2 \eta_0) \left(c_r^2 \eta_0^2 + \frac{u_r \underline{a}_r}{u_\ell \underline{a}_\ell} u_\ell^2 c_\ell^2 |\check{\eta}|^2 \right).$$

This last expression is simplified a little further by observing that we have

$$(i \eta_0 - u_\ell \underline{\beta}_1^-) (u_\ell \underline{a}_\ell + i c_\ell^2 \eta_0) = \frac{-1}{c_\ell^2 - u_\ell^2} (u_\ell^2 \underline{a}_\ell^2 + c_\ell^4 \eta_0^2) = -c_\ell^2 (\eta_0^2 + u_\ell^2 |\check{\eta}|^2),$$

and

$$c_r^2 \eta_0^2 + \frac{u_r \underline{a}_r}{u_\ell \underline{a}_\ell} u_\ell^2 c_\ell^2 |\check{\eta}|^2 = u_\ell u_r \frac{\underline{a}_r}{\underline{a}_\ell} \left(c_\ell^2 |\check{\eta}|^2 - \frac{\underline{a}_\ell^2}{c_\ell^2} \right) = u_\ell u_r \frac{\underline{a}_r}{\underline{a}_\ell} (\eta_0^2 + u_\ell^2 |\check{\eta}|^2).$$

Eventually, we find that the sum of all the terms that involve $p''(\rho_\ell)$ factorizes as $Q_\ell/2$ where Q_ℓ is defined in (33).

• Step 2: computation of the $p''(\rho_r)$ factor. We follow the same strategy as in the first step and compute the contribution that involves $p''(\rho_r)$ in each kernel. The contribution of the kernel q_2 equals

$$- \Upsilon (D_{d+1} + u_r D_{d+2}) \gamma_2^2 (i \eta_0 + u_r \underline{\beta}_2^-)^2 = [\rho][u] \Upsilon \gamma_2 (i \eta_0 + u_r \underline{\beta}_2^-)^2 \frac{u_\ell \underline{a}_\ell}{c_r^2} \left\{ i \eta_0 u_r \underline{a}_r - u_r^2 c_r^2 |\check{\eta}|^2 \right\}, \quad (38)$$

where we have used Lemma 4 to simplify $(D_{d+1} + u_r D_{d+2}) \gamma_2$. The contribution of the kernel q_3 equals

$$\begin{aligned} i(k+k') \int_0^{+\infty} \frac{\omega_2}{\gamma_2} (i \check{\eta}^T) e^{-(k+k') \underline{\beta}_2^+ z} \gamma_2^2 (i \eta_0 + u_r \underline{\beta}_2^-)^2 \check{\eta} e^{(k+k') \underline{\beta}_2^- z} dz \\ = -\gamma_2 (i \eta_0 + u_r \underline{\beta}_2^-)^2 \omega_2 \frac{|\check{\eta}|^2}{\underline{\beta}_2^+ - \underline{\beta}_2^-}, \end{aligned}$$

and the contribution of the kernel q_4 reads

$$\begin{aligned} \int_0^{+\infty} \frac{\omega_2}{\gamma_2} \underline{\beta}_2^+ e^{-(k+k') \underline{\beta}_2^+ z} \gamma_2^2 (i \eta_0 + u_r \underline{\beta}_2^-)^2 (k+k') \underline{\beta}_2^- e^{(k+k') \underline{\beta}_2^- z} dz \\ = \gamma_2 (i \eta_0 + u_r \underline{\beta}_2^-)^2 \omega_2 \frac{\underline{\beta}_2^+ \underline{\beta}_2^-}{\underline{\beta}_2^+ - \underline{\beta}_2^-}. \end{aligned}$$

Adding the contributions of q_3 and q_4 gives the term

$$-\gamma_2 (i \eta_0 + u_r \underline{\beta}_2^-)^2 \frac{\omega_2}{2 \underline{a}_r} \left(\frac{\underline{a}_r^2}{c_r^2} + c_r^2 |\check{\eta}|^2 \right).$$

When we add the latter term with the expression in (38), we obtain

$$-[\rho] [u] \Upsilon \gamma_2 (i \eta_0 + u_r \underline{\beta}_2^-)^2 (\eta_0^2 + u_r^2 |\check{\eta}|^2) \frac{i u_r \eta_0}{2 \underline{a}_r} (u_\ell \underline{a}_\ell + i c_\ell^2 \eta_0),$$

and this quantity is further simplified by using the relation

$$(i \eta_0 + u_r \underline{\beta}_2^-) i \eta_0 (u_\ell \underline{a}_\ell + i c_\ell^2 \eta_0) = -u_\ell \underline{a}_\ell \frac{u_r^2 \underline{a}_r^2 + c_r^4 \eta_0^2}{c_r^2 (c_r^2 - u_r^2)} = -u_\ell \underline{a}_\ell (\eta_0^2 + u_r^2 |\check{\eta}|^2).$$

Eventually, we find that the sum of all the terms that involve $p''(\rho_r)$ factorizes as $Q_r/2$.

• Step 3: computation of the remaining terms. In order to prove Proposition 2, we can assume from now on, and without loss of generality that $p''(\rho_\ell) = p''(\rho_r) = 0$ in (34) and (35). With this simplification, we compute

$$\sum_{k=1}^{d-1} \check{\eta}_k d^2 f^k(v_\ell) \cdot (x_1^-, x_1^-) = -\frac{2 i c_\ell^4 |\check{\eta}|^2}{\rho_\ell} \begin{pmatrix} 0 \\ -i \check{\eta} \\ \underline{\beta}_1^- \end{pmatrix}, \quad (39)$$

$$d^2 \tilde{f}^d(v_\ell) \cdot (x_1^-, x_1^-) = \frac{2 c_\ell^2}{\rho_\ell} \begin{pmatrix} 0 \\ -i c_\ell^2 \underline{\beta}_1^- \check{\eta} \\ c_\ell^2 |\check{\eta}|^2 + (i \eta_0 - u_\ell \underline{\beta}_1^-)^2 \\ i c_\ell^2 \eta_0 \underline{\beta}_1^- + u_\ell (i \eta_0 - u_\ell \underline{\beta}_1^-)^2 \end{pmatrix}, \quad (40)$$

$$\sum_{k=1}^{d-1} \check{\eta}_k d^2 f^k(v_r) \cdot (x_2^-, x_2^-) = \frac{2 i c_r^4 |\check{\eta}|^2}{\rho_r} \begin{pmatrix} 0 \\ i \check{\eta} \\ \underline{\beta}_2^- \end{pmatrix}, \quad (41)$$

$$d^2 \tilde{f}^d(v_r) \cdot (x_2^-, x_2^-) = \frac{2 c_r^2}{\rho_r} \begin{pmatrix} 0 \\ i c_r^2 \underline{\beta}_2^- \check{\eta} \\ c_r^2 |\check{\eta}|^2 + (i \eta_0 + u_r \underline{\beta}_2^-)^2 \\ -i c_r^2 \eta_0 \underline{\beta}_2^- + u_r (i \eta_0 + u_r \underline{\beta}_2^-)^2 \end{pmatrix}, \quad (42)$$

where in (40) and (42), we have used the relations

$$c_\ell^2 (\underline{\beta}_1^-)^2 = c_\ell^2 |\check{y}|^2 + (i \eta_0 - u_\ell \underline{\beta}_1^-)^2, \quad c_r^2 (\underline{\beta}_2^-)^2 = c_r^2 |\check{y}|^2 + (i \eta_0 + u_r \underline{\beta}_2^-)^2.$$

With these expressions, let us look first at the kernel q_2 in (28). Using (40), we compute

$$\begin{aligned} \sigma^* d^2 \tilde{f}^d(v_\ell) \cdot (x_1^-, x_1^-) \gamma_1^2 &= \frac{2c_\ell^2}{\rho_\ell} \underline{\Upsilon} \gamma_1 \left\{ i c_\ell^2 \underline{\beta}_1^- (\eta_0 \gamma_1 D_{d+2} - |\check{y}|^2 \gamma_1 \check{D}) + c_\ell^2 |\check{y}|^2 \gamma_1 D_{d+1} \right. \\ &\quad \left. + (i \eta_0 - u_\ell \underline{\beta}_1^-)^2 \gamma_1 (D_{d+1} + u_\ell D_{d+2}) \right\}, \end{aligned}$$

and Lemma 4 turns this expression into

$$\begin{aligned} \sigma^* d^2 \tilde{f}^d(v_\ell) \cdot (x_1^-, x_1^-) \gamma_1^2 &= \frac{2c_\ell^2}{\rho_\ell} [\rho] \underline{\Upsilon} (\eta_0^2 + u_r^2 |\check{y}|^2) \gamma_1 \left\{ u_r \underline{a}_r (u_\ell \underline{a}_\ell + i c_\ell^2 \eta_0) \underline{\beta}_1^- - u_r c_\ell^2 |\check{y}|^2 (u_\ell \underline{a}_r + i c_r^2 \eta_0) \right\} \\ &\quad - \frac{2c_\ell^2}{\rho_\ell} [\rho] [u] \underline{\Upsilon} \gamma_1 (i \eta_0 - u_\ell \underline{\beta}_1^-)^2 \left\{ (\eta_0^2 + u_r^2 |\check{y}|^2) i c_r^2 \eta_0 + u_\ell u_r |\check{y}|^2 (u_r \underline{a}_r - i c_r^2 \eta_0) \right\}. \quad (43) \end{aligned}$$

Similarly, we derive the relation

$$\begin{aligned} \sigma^* d^2 \tilde{f}^d(v_r) \cdot (x_2^-, x_2^-) \gamma_2^2 &= -\frac{2c_r^2}{\rho_r} [\rho] \underline{\Upsilon} (\eta_0^2 + u_r^2 |\check{y}|^2) \gamma_2 \left\{ u_\ell \underline{a}_\ell (u_r \underline{a}_r + i c_r^2 \eta_0) \underline{\beta}_2^- + u_\ell c_r^2 |\check{y}|^2 (u_r \underline{a}_\ell + i c_\ell^2 \eta_0) \right\} \\ &\quad + \frac{2c_r^2}{\rho_r} [\rho] [u] \underline{\Upsilon} \gamma_2 (i \eta_0 + u_r \underline{\beta}_2^-)^2 \left\{ (\eta_0^2 + u_r^2 |\check{y}|^2) i c_\ell^2 \eta_0 + u_r^2 |\check{y}|^2 (u_\ell \underline{a}_\ell - i c_\ell^2 \eta_0) \right\}. \quad (44) \end{aligned}$$

The kernels q_3, q_4 both read as a sum of two contributions, one from the system ahead of the phase boundary, and one from the system behind the phase boundary. We compute each of these contributions separately in order to combine them with either (43) or (44). Using (39), the ‘left’ contribution of q_3 equals

$$\begin{aligned} &i(k+k') \int_0^{+\infty} \frac{\omega_1}{\gamma_1} \tilde{\ell}_1 e^{-(k+k') \underline{\beta}_1^+ z} \gamma_1^2 \frac{-2i c_\ell^4 |\check{y}|^2}{\rho_\ell} \begin{pmatrix} 0 \\ -i \check{y} \\ \underline{\beta}_1^- \end{pmatrix} e^{(k+k') \underline{\beta}_1^- z} dz \\ &+ i(k+k') \int_0^{+\infty} \frac{\omega_3}{\gamma_1} \tilde{\ell}_3 e^{-(k+k') \underline{\beta}_3^+ z} \gamma_1^2 \frac{-2i c_\ell^4 |\check{y}|^2}{\rho_\ell} \begin{pmatrix} 0 \\ -i \check{y} \\ \underline{\beta}_1^- \end{pmatrix} e^{(k+k') \underline{\beta}_1^- z} dz \\ &= \frac{2c_\ell^4}{\rho_\ell} \gamma_1 \omega_1 \frac{|\check{y}|^2}{\underline{\beta}_1^+ - \underline{\beta}_1^-} \tilde{\ell}_1 \begin{pmatrix} 0 \\ -i \check{y} \\ \underline{\beta}_1^- \end{pmatrix} + \frac{2c_\ell^4}{\rho_\ell} \gamma_1 \omega_3 \frac{|\check{y}|^2}{\underline{\beta}_3^+ - \underline{\beta}_1^-} \tilde{\ell}_3 \begin{pmatrix} 0 \\ -i \check{y} \\ \underline{\beta}_1^- \end{pmatrix}, \quad (45) \end{aligned}$$

and, similarly (using now (41) rather than (39)), the ‘right’ contribution of q_3 equals

$$i(k+k') \int_0^{+\infty} \frac{\omega_2}{\gamma_2} \tilde{\ell}_2 e^{-(k+k') \underline{\beta}_2^+ z} \gamma_2^2 \frac{2i c_r^4 |\check{\eta}|^2}{\rho_r} \begin{pmatrix} 0 \\ i \check{\eta} \\ \underline{\beta}_2^- \end{pmatrix} e^{(k+k') \underline{\beta}_2^- z} dz$$

$$= -\frac{2c_r^4}{\rho_r} \gamma_2 \omega_2 \frac{|\check{\eta}|^2}{\underline{\beta}_2^+ - \underline{\beta}_2^-} \tilde{\ell}_2 \begin{pmatrix} 0 \\ i \check{\eta} \\ \underline{\beta}_2^- \end{pmatrix}. \quad (46)$$

The kernel q_4 is computed by first observing that the vectors $d^2 f^d(v_\ell) \cdot (\underline{r}_1^-, \underline{r}_1^-)$, and $d^2 f^d(v_r) \cdot (\underline{r}_2^-, \underline{r}_2^-)$ are obtained by retaining only the three first coordinates in (40) and (42):

$$d^2 f^d(v_\ell) \cdot (\underline{r}_1^-, \underline{r}_1^-) = \frac{2c_\ell^4 \underline{\beta}_1^-}{\rho_\ell} \begin{pmatrix} 0 \\ -i \check{\eta} \\ \underline{\beta}_1^- \end{pmatrix}, \quad d^2 f^d(v_r) \cdot (\underline{r}_2^-, \underline{r}_2^-) = \frac{2c_r^4 \underline{\beta}_2^-}{\rho_r} \begin{pmatrix} 0 \\ i \check{\eta} \\ \underline{\beta}_2^- \end{pmatrix}.$$

Consequently, the ‘left’ contribution of q_4 equals

$$-\int_0^{+\infty} \frac{\omega_1}{\gamma_1} \tilde{\ell}_1 e^{-(k+k') \underline{\beta}_1^+ z} \gamma_1^2 \frac{2c_\ell^4 \underline{\beta}_1^-}{\rho_\ell} \begin{pmatrix} 0 \\ -i \check{\eta} \\ \underline{\beta}_1^- \end{pmatrix} (k+k') \underline{\beta}_1^- e^{(k+k') \underline{\beta}_1^- z} dz$$

$$-\int_0^{+\infty} \frac{\omega_3}{\gamma_1} \tilde{\ell}_1 e^{-(k+k') \underline{\beta}_3^+ z} \gamma_1^2 \frac{2c_\ell^4 \underline{\beta}_1^-}{\rho_\ell} \begin{pmatrix} 0 \\ -i \check{\eta} \\ \underline{\beta}_1^- \end{pmatrix} (k+k') \underline{\beta}_1^- e^{(k+k') \underline{\beta}_1^- z} dz$$

$$= -\frac{2c_\ell^4}{\rho_\ell} \gamma_1 \omega_1 \frac{(\underline{\beta}_1^-)^2}{\underline{\beta}_1^+ - \underline{\beta}_1^-} \tilde{\ell}_1 \begin{pmatrix} 0 \\ -i \check{\eta} \\ \underline{\beta}_1^- \end{pmatrix} - \frac{2c_\ell^4}{\rho_\ell} \gamma_1 \omega_3 \frac{(\underline{\beta}_1^-)^2}{\underline{\beta}_3^+ - \underline{\beta}_1^-} \tilde{\ell}_3 \begin{pmatrix} 0 \\ -i \check{\eta} \\ \underline{\beta}_1^- \end{pmatrix}, \quad (47)$$

and the ‘right’ contribution of q_4 equals

$$\frac{2c_r^4}{\rho_r} \gamma_2 \omega_2 \frac{(\underline{\beta}_2^-)^2}{\underline{\beta}_2^+ - \underline{\beta}_2^-} \tilde{\ell}_2 \begin{pmatrix} 0 \\ i \check{\eta} \\ \underline{\beta}_2^- \end{pmatrix}. \quad (48)$$

The ‘left’ contribution of $q_3 + q_4$ is obtained by adding the expressions in (45) and (47), which gives

$$-\frac{2c_\ell^4}{\rho_\ell} \gamma_1 \omega_1 \frac{(\underline{\beta}_1^-)^2 - |\check{\eta}|^2}{\underline{\beta}_1^+ - \underline{\beta}_1^-} \tilde{\ell}_1 \begin{pmatrix} 0 \\ -i \check{\eta} \\ \underline{\beta}_1^- \end{pmatrix} - \frac{2c_\ell^4}{\rho_\ell} \gamma_1 \omega_3 \frac{(\underline{\beta}_1^-)^2 - |\check{\eta}|^2}{\underline{\beta}_3^+ - \underline{\beta}_1^-} \tilde{\ell}_3 \begin{pmatrix} 0 \\ -i \check{\eta} \\ \underline{\beta}_1^- \end{pmatrix}$$

$$= -\frac{2c_\ell^2}{\rho_\ell} \gamma_1 (i \eta_0 - u_\ell \underline{\beta}_1^-)^2 \left\{ \omega_1 \tilde{\ell}_1 \begin{pmatrix} 0 \\ -i \check{\eta} \\ \underline{\beta}_1^- \end{pmatrix} + \omega_3 \tilde{\ell}_3 \begin{pmatrix} 0 \\ -i \check{\eta} \\ \underline{\beta}_1^- \end{pmatrix} \right\}$$

$$= \frac{c_\ell^2}{\rho_\ell \underline{a}_\ell} \gamma_1 (i \eta_0 - u_\ell \underline{\beta}_1^-)^2 \left\{ (\eta_0^2 + u_\ell^2 |\check{\eta}|^2 - 2c_\ell^2 |\check{\eta}|^2) \omega_1 + 2u_\ell \underline{a}_\ell |\check{\eta}|^2 \omega_3 \right\}. \quad (49)$$

The ‘right’ contribution of $q_3 + q_4$ is obtained by adding the expressions in (46) and (48), which gives

$$\begin{aligned} \frac{2c_r^4}{\rho_r} \gamma_2 \omega_2 \frac{(\underline{\beta}_2^-)^2 - |\check{\eta}|^2}{\underline{\beta}_2^+ - \underline{\beta}_2^-} \tilde{\ell}_2 \begin{pmatrix} 0 \\ i \check{\eta} \\ \underline{\beta}_2^- \end{pmatrix} &= \frac{2c_r^2}{\rho_r} \gamma_2 \omega_2 (i \eta_0 + u_r \underline{\beta}_2^-)^2 \frac{\underline{\beta}_2^+ \underline{\beta}_2^- - |\check{\eta}|^2}{\underline{\beta}_2^+ - \underline{\beta}_2^-} \\ &= \frac{c_r^2}{\rho_r \underline{a}_r} \gamma_2 (i \eta_0 + u_r \underline{\beta}_2^-)^2 (\eta_0^2 + u_r^2 |\check{\eta}|^2 - 2c_r^2 |\check{\eta}|^2) \omega_2. \end{aligned} \quad (50)$$

We can now compute the ‘left’ contribution of the full kernel $q_2 + q_3 + q_4$ by combining the expression in (43) with the one in (49). We use here the expressions of ω_1 and ω_3 given in Lemma 2. The sum of (43) and (49) reads

$$\begin{aligned} A_\ell := & \frac{2c_\ell^2}{\rho_\ell} [\rho] \underline{\Upsilon} (\eta_0^2 + u_r^2 |\check{\eta}|^2) \gamma_1 \left\{ u_r \underline{a}_r (u_\ell \underline{a}_\ell + i c_\ell^2 \eta_0) \underline{\beta}_1^- - u_r c_\ell^2 |\check{\eta}|^2 (u_\ell \underline{a}_r + i c_r^2 \eta_0) \right\} \\ & - \frac{2c_\ell^2}{\rho_\ell} [\rho] [u] \underline{\Upsilon} \gamma_1 (i \eta_0 - u_\ell \underline{\beta}_1^-)^2 \left\{ (\eta_0^2 + u_r^2 |\check{\eta}|^2) i c_r^2 \eta_0 + u_\ell u_r |\check{\eta}|^2 (u_r \underline{a}_r - i c_r^2 \eta_0) \right\} \\ & + \frac{c_\ell^2}{\rho_\ell \underline{a}_\ell} [\rho] [u] \underline{\Upsilon} \gamma_1 (i \eta_0 - u_\ell \underline{\beta}_1^-)^2 \left\{ 2 u_\ell \underline{a}_\ell |\check{\eta}|^2 [u] \frac{\eta_0^2 - u_\ell u_r |\check{\eta}|^2}{\eta_0^2 + u_\ell^2 |\check{\eta}|^2} (u_r \underline{a}_r - i c_r^2 \eta_0) \right. \\ & \left. + (\eta_0^2 + u_\ell^2 |\check{\eta}|^2 - 2c_\ell^2 |\check{\eta}|^2) i u_\ell \eta_0 (u_r \underline{a}_r - i c_r^2 \eta_0) \frac{\eta_0^2 + u_r^2 |\check{\eta}|^2}{\eta_0^2 + u_\ell^2 |\check{\eta}|^2} \right\}. \end{aligned}$$

The three last rows in the definition of A_ℓ are factorized after a little bit of algebra, and we get

$$\begin{aligned} A_\ell = & \frac{2c_\ell^2}{\rho_\ell} [\rho] \underline{\Upsilon} (\eta_0^2 + u_r^2 |\check{\eta}|^2) \gamma_1 \left\{ u_r \underline{a}_r (u_\ell \underline{a}_\ell + i c_\ell^2 \eta_0) \underline{\beta}_1^- - u_r c_\ell^2 |\check{\eta}|^2 (u_\ell \underline{a}_r + i c_r^2 \eta_0) \right\} \\ & - \frac{1}{\rho_\ell} [\rho] [u] \underline{\Upsilon} \gamma_1 (i \eta_0 - u_\ell \underline{\beta}_1^-)^2 u_\ell u_r \frac{\underline{a}_r}{\underline{a}_\ell} (u_\ell \underline{a}_\ell + i c_\ell^2 \eta_0) (\eta_0^2 + u_r^2 |\check{\eta}|^2). \end{aligned}$$

We have thus shown that the ‘left’ contribution A_ℓ of $q_2 + q_3 + q_4$ reads

$$A_\ell = \frac{2c_\ell^2}{\rho_\ell} [\rho] \underline{\Upsilon} (\eta_0^2 + u_r^2 |\check{\eta}|^2) \gamma_1 \left\{ u_r \underline{a}_r (u_\ell \underline{a}_\ell + i c_\ell^2 \eta_0) \underline{\beta}_1^- - u_r c_\ell^2 |\check{\eta}|^2 (u_\ell \underline{a}_r + i c_r^2 \eta_0) \right\} + \frac{c_\ell^2}{\rho_\ell} Q_\ell. \quad (51)$$

The ‘right’ contribution of $q_2 + q_3 + q_4$ by combining the expression in (44) (with a minus sign, recall the definition (28) of the kernel q_2) with the one in (50):

$$\begin{aligned} A_r := & \frac{2c_r^2}{\rho_r} [\rho] \underline{\Upsilon} (\eta_0^2 + u_r^2 |\check{\eta}|^2) \gamma_2 \left\{ u_\ell \underline{a}_\ell (u_r \underline{a}_r + i c_r^2 \eta_0) \underline{\beta}_2^- + u_\ell c_r^2 |\check{\eta}|^2 (u_r \underline{a}_\ell + i c_\ell^2 \eta_0) \right\} \\ & - \frac{2c_r^2}{\rho_r} [\rho] [u] \underline{\Upsilon} \gamma_2 (i \eta_0 + u_r \underline{\beta}_2^-)^2 \left\{ (\eta_0^2 + u_r^2 |\check{\eta}|^2) i c_\ell^2 \eta_0 + u_r^2 |\check{\eta}|^2 (u_\ell \underline{a}_\ell - i c_\ell^2 \eta_0) \right\} \\ & + \frac{c_r^2}{\rho_r \underline{a}_r} \gamma_2 (i \eta_0 + u_r \underline{\beta}_2^-)^2 (\eta_0^2 + u_r^2 |\check{\eta}|^2 - 2c_r^2 |\check{\eta}|^2) \omega_2. \end{aligned}$$

Once again, the last two rows are factorized after a few calculations that we skip, and we get

$$\begin{aligned} A_r = & \frac{2c_r^2}{\rho_r} [\rho] \underline{\Upsilon} (\eta_0^2 + u_r^2 |\check{\eta}|^2) \gamma_2 \left\{ u_\ell \underline{a}_\ell (u_r \underline{a}_r + i c_r^2 \eta_0) \underline{\beta}_2^- + u_\ell c_r^2 |\check{\eta}|^2 (u_r \underline{a}_\ell + i c_\ell^2 \eta_0) \right\} \\ & - \frac{1}{\rho_r} [\rho] [u] \underline{\Upsilon} \gamma_2 (i \eta_0 + u_r \underline{\beta}_2^-)^2 u_\ell u_r \frac{\underline{a}_\ell}{\underline{a}_r} (u_r \underline{a}_r + i c_r^2 \eta_0) (\eta_0^2 + u_r^2 |\check{\eta}|^2). \end{aligned}$$

We have thus shown that the ‘left’ contribution A_ℓ of $q_2 + q_3 + q_4$ reads

$$A_r = \frac{2c_r^2}{\rho_r} [\rho] \underline{\Upsilon} (\eta_0^2 + u_r^2 |\check{\eta}|^2) \gamma_2 \left\{ u_\ell \underline{a}_\ell (u_r \underline{a}_r + i c_r^2 \eta_0) \underline{\beta}_2^- + u_\ell c_r^2 |\check{\eta}|^2 (u_r \underline{a}_\ell + i c_\ell^2 \eta_0) \right\} + \frac{c_r^2}{\rho_r} Q_r. \quad (52)$$

The final step of the proof consists in simplifying the remaining terms in A_ℓ and A_r . More specifically, the first factor in the expression (51) of A_ℓ can be simplified as follows:

$$\begin{aligned} & u_\ell \left\{ u_r \underline{a}_r (u_\ell \underline{a}_\ell + i c_\ell^2 \eta_0) \underline{\beta}_1^- - u_r c_\ell^2 |\check{\eta}|^2 (u_\ell \underline{a}_r + i c_r^2 \eta_0) \right\} \\ &= u_r \underline{a}_r (u_\ell \underline{a}_\ell + i c_\ell^2 \eta_0) (u_\ell \underline{\beta}_1^- - i \eta_0) - c_\ell^2 \eta_0^2 (u_r \underline{a}_r + i c_r^2 \eta_0) - u_\ell u_r c_\ell^2 |\check{\eta}|^2 (u_\ell \underline{a}_r + i c_r^2 \eta_0) \\ &= u_r \underline{a}_r c_\ell^2 (\eta_0^2 + u_\ell^2 |\check{\eta}|^2) - c_\ell^2 \eta_0^2 (u_r \underline{a}_r + i c_r^2 \eta_0) - u_\ell u_r c_\ell^2 |\check{\eta}|^2 (u_\ell \underline{a}_r + i c_r^2 \eta_0) \\ &= -i c_\ell^2 c_r^2 \eta_0 (\eta_0^2 + u_\ell u_r |\check{\eta}|^2). \end{aligned}$$

Similarly, the first factor in the expression (52) of A_r can be simplified by using

$$u_r \left\{ u_\ell \underline{a}_\ell (u_r \underline{a}_r + i c_r^2 \eta_0) \underline{\beta}_2^- + u_\ell c_r^2 |\check{\eta}|^2 (u_r \underline{a}_\ell + i c_\ell^2 \eta_0) \right\} = i c_\ell^2 c_r^2 \eta_0 (\eta_0^2 + u_\ell u_r |\check{\eta}|^2).$$

Using these two last simplifications, we can add (51) and (52) and obtain

$$(q_2 + q_3 + q_4)(k, k') = \frac{c_\ell^2}{\rho_\ell} Q_\ell + \frac{c_r^2}{\rho_r} Q_r + Q_\sharp,$$

with Q_\sharp defined in (33). This completes the proof of Proposition 2. \square

We now compute the kernel $q_2 + q_3 + q_4$ in the domain $\{k > 0, k' > 0, k + k' > 0\}$.

Proposition 3. *Let Q_ℓ and Q_r be defined in (33) and let us define*

$$Q_b := -2 [\rho] [u] \underline{\Upsilon} (\eta_0^2 + u_r^2 |\check{\eta}|^2) u_\ell u_r |\check{\eta}|^2 \left\{ \frac{c_\ell^4 \underline{a}_r}{\rho_\ell \underline{a}_\ell} \overline{\gamma}_1 (i \eta_0 - u_\ell \underline{\beta}_1^+) + \frac{c_r^4 \underline{a}_\ell}{\rho_r \underline{a}_r} \overline{\gamma}_2 (i \eta_0 + u_r \underline{\beta}_2^+) \right\}.$$

Then the kernels q_2, q_3, q_4 defined in (28), (29) and (30) satisfy

$$(q_2 + q_3 + q_4)(k, k') = \left\{ \left(\frac{p''(\rho_\ell)}{2} - \frac{c_\ell^2}{\rho_\ell} \right) \overline{Q}_\ell + \left(\frac{p''(\rho_r)}{2} - \frac{c_r^2}{\rho_r} \right) \overline{Q}_r + Q_b \right\} \left(1 + \frac{k'}{k} \right),$$

for all (k, k') such that $k > 0, k' < 0$ and $k + k' > 0$.

Proof. We split again the proof in several steps, as was done in the proof of Proposition 2.

• Step 1: computation of the $p''(\rho_\ell)$ factor. We collect again the contributions that involve $p''(\rho_\ell)$. The contribution of the kernel q_2 equals

$$\begin{aligned} & \underline{\Upsilon} (D_{d+1} + u_\ell D_{d+2}) |\gamma_1|^2 (-i \eta_0 + u_\ell \underline{\beta}_1^-) (i \eta_0 - u_\ell \underline{\beta}_1^+) \\ &= -[\rho] [u] \underline{\Upsilon} \overline{\gamma}_1 |i \eta_0 - u_\ell \underline{\beta}_1^-|^2 \left\{ (\eta_0^2 + u_r^2 |\check{\eta}|^2) i c_r^2 \eta_0 + (u_r \underline{a}_r - i c_r^2 \eta_0) u_\ell u_r |\check{\eta}|^2 \right\}, \quad (53) \end{aligned}$$

where the product $(D_{d+1} + u_\ell D_{d+2}) \gamma_1$ has already been computed when deriving (36).

The contribution of the kernel q_3 equals

$$\begin{aligned} & i(k+k') \int_0^{+\infty} \frac{\omega_1}{\gamma_1} (-i \check{\eta}^T) e^{-(k+k') \underline{\beta}_1^+ z} |\gamma_1|^2 |i \eta_0 - u_\ell \underline{\beta}_1^-|^2 \check{\eta} e^{k \underline{\beta}_1^- z} e^{k' \underline{\beta}_1^+ z} dz \\ & + i(k+k') \int_0^{+\infty} \frac{\omega_3}{\gamma_1} (\eta_0 \check{\eta}^T) e^{-(k+k') \underline{\beta}_3^+ z} |\gamma_1|^2 |i \eta_0 - u_\ell \underline{\beta}_1^-|^2 \check{\eta} e^{k \underline{\beta}_1^- z} e^{k' \underline{\beta}_1^+ z} dz \\ & = \overline{\gamma_1} |i \eta_0 - u_\ell \underline{\beta}_1^-|^2 \left\{ \left(1 + \frac{k'}{k}\right) \frac{\omega_1 |\check{\eta}|^2}{\underline{\beta}_1^+ - \underline{\beta}_1^-} + \frac{i \eta_0 (k+k') u_\ell |\check{\eta}|^2 \omega_3}{k(i \eta_0 - u_\ell \underline{\beta}_1^-) + k'(i \eta_0 - u_\ell \underline{\beta}_1^+)} \right\}. \end{aligned}$$

Similarly, the contribution of the kernel q_4 reads

$$\begin{aligned} & - \int_0^{+\infty} \frac{\omega_1}{\gamma_1} \underline{\beta}_1^+ e^{-(k+k') \underline{\beta}_1^+ z} |\gamma_1|^2 |i \eta_0 - u_\ell \underline{\beta}_1^-|^2 (k \underline{\beta}_1^- + k' \underline{\beta}_1^+) e^{k \underline{\beta}_1^- z} e^{k' \underline{\beta}_1^+ z} dz \\ & - \int_0^{+\infty} \frac{\omega_3}{\gamma_1} u_\ell |\check{\eta}|^2 e^{-(k+k') \underline{\beta}_3^+ z} |\gamma_1|^2 |i \eta_0 - u_\ell \underline{\beta}_1^-|^2 (k \underline{\beta}_1^- + k' \underline{\beta}_1^+) e^{k \underline{\beta}_1^- z} e^{k' \underline{\beta}_1^+ z} dz \\ & = -\overline{\gamma_1} |i \eta_0 - u_\ell \underline{\beta}_1^-|^2 \left\{ \frac{\omega_1}{\underline{\beta}_1^+ - \underline{\beta}_1^-} \left(\underline{\beta}_1^+ \underline{\beta}_1^- + \frac{k'}{k} (\underline{\beta}_1^+)^2 \right) + \frac{(k u_\ell \underline{\beta}_1^- + k' u_\ell \underline{\beta}_1^+) u_\ell |\check{\eta}|^2 \omega_3}{k(i \eta_0 - u_\ell \underline{\beta}_1^-) + k'(i \eta_0 - u_\ell \underline{\beta}_1^+)} \right\}. \end{aligned}$$

Adding the contributions of q_3 and q_4 gives the term

$$\overline{\gamma_1} |i \eta_0 - u_\ell \underline{\beta}_1^-|^2 \left\{ -\frac{1}{2 \underline{a}_\ell} \left(\frac{\underline{a}_\ell^2}{c_\ell^2} + c_\ell^2 |\check{\eta}|^2 \right) \omega_1 + u_\ell |\check{\eta}|^2 \omega_3 + \frac{k'}{k} \omega_1 \frac{|\check{\eta}|^2 - (\underline{\beta}_1^+)^2}{\underline{\beta}_1^+ - \underline{\beta}_1^-} \right\}. \quad (54)$$

We now add the contributions in (53) and (54) in order to obtain the $p''(\rho_\ell)$ term in $q_2 + q_3 + q_4$. There is first a constant term that is independent of (k, k') , and this term is entirely similar to the one derived when adding (36) and (37) (see Step 1 in the proof of Proposition 2). Namely, the constant term in the sum of (53) and (54) equals

$$\begin{aligned} & - [\rho] [u] \underline{\Upsilon} \overline{\gamma_1} |i \eta_0 - u_\ell \underline{\beta}_1^-|^2 \frac{\eta_0^2 + u_r^2 |\check{\eta}|^2}{2 c_\ell^2 (\eta_0^2 + u_\ell^2 |\check{\eta}|^2)} (u_\ell \underline{a}_\ell + i c_\ell^2 \eta_0) \left(c_r^2 \eta_0^2 + \frac{u_r \underline{a}_r}{u_\ell \underline{a}_\ell} u_\ell^2 c_\ell^2 |\check{\eta}|^2 \right) \\ & = \frac{1}{2} [\rho] [u] \underline{\Upsilon} \overline{\gamma_1} (-i \eta_0 + u_\ell \underline{\beta}_1^+) (\eta_0^2 + u_r^2 |\check{\eta}|^2) \left(c_r^2 \eta_0^2 + \frac{u_r \underline{a}_r}{u_\ell \underline{a}_\ell} u_\ell^2 c_\ell^2 |\check{\eta}|^2 \right) = \frac{1}{2} \overline{Q}_\ell. \end{aligned}$$

The last contribution in the $p''(\rho_\ell)$ term is the one that depends on (k, k') in (54), that is

$$\begin{aligned} & \frac{k'}{k} \overline{\gamma_1} |i \eta_0 - u_\ell \underline{\beta}_1^-|^2 \omega_1 \frac{|\check{\eta}|^2 - (\underline{\beta}_1^+)^2}{\underline{\beta}_1^+ - \underline{\beta}_1^-} \\ & = \frac{k'}{k} [\rho] [u] \underline{\Upsilon} \overline{\gamma_1} |i \eta_0 - u_\ell \underline{\beta}_1^-|^2 \frac{\eta_0^2 + u_r^2 |\check{\eta}|^2}{\eta_0^2 + u_\ell^2 |\check{\eta}|^2} \frac{(-u_\ell u_r \underline{a}_r)}{c_\ell^2} (u_\ell \underline{a}_\ell - i c_\ell^2 \eta_0) \frac{(i \eta_0 - u_\ell \underline{\beta}_1^+)^2}{2 c_\ell^2 \underline{a}_\ell} (c_\ell^2 - u_\ell^2) \\ & = \frac{k'}{2k} [\rho] [u] \underline{\Upsilon} \overline{\gamma_1} (-i \eta_0 + u_\ell \underline{\beta}_1^+) (\eta_0^2 + u_r^2 |\check{\eta}|^2) u_\ell u_r \frac{\underline{a}_r}{\underline{a}_\ell} \frac{|i \eta_0 - u_\ell \underline{\beta}_1^-|^2}{c_\ell^2} (c_\ell^2 - u_\ell^2) = \frac{k'}{2k} \overline{Q}_\ell. \end{aligned}$$

We have thus shown that the $p''(\rho_\ell)$ term in $(q_2 + q_3 + q_4)(k, k')$ is as claimed in Proposition 3.

• Step 2: computation of the $p''(\rho_r)$ factor. The contribution of the kernel q_2 equals

$$-\underline{\Upsilon} (D_{d+1} + u_r D_{d+2}) |\gamma_2|^2 |i \eta_0 + u_r \underline{\beta}_2^-|^2 = [\rho] [u] \underline{\Upsilon} \overline{\gamma_2} |i \eta_0 + u_r \underline{\beta}_2^-|^2 \frac{u_\ell \underline{a}_\ell}{c_r^2} \left\{ i \eta_0 u_r \underline{a}_r - u_r^2 c_r^2 |\check{\eta}|^2 \right\}. \quad (55)$$

The contribution of the kernel q_3 equals

$$\begin{aligned} i(k+k') \int_0^{+\infty} \frac{\omega_2}{\gamma_2} (i\check{\eta}^T) e^{-(k+k')\beta_2^+ z} |\gamma_2|^2 |i\eta_0 + u_r \beta_2^-|^2 \check{\eta} e^{k\beta_2^- z} e^{k'\beta_2^+ z} dz \\ = -\overline{\gamma_2} |i\eta_0 + u_r \beta_2^-|^2 \omega_2 \frac{|\check{\eta}|^2}{\beta_2^+ - \beta_2^-} \left(1 + \frac{k'}{k}\right), \end{aligned}$$

and the contribution of the kernel q_4 reads

$$\begin{aligned} \int_0^{+\infty} \frac{\omega_2}{\gamma_2} \beta_2^+ e^{-(k+k')\beta_2^+ z} |\gamma_2|^2 |i\eta_0 + u_r \beta_2^-|^2 (k\beta_2^- + k'\beta_2^+) e^{k\beta_2^- z} e^{k'\beta_2^+ z} dz \\ = \overline{\gamma_2} |i\eta_0 + u_r \beta_2^-|^2 \frac{\omega_2}{\beta_2^+ - \beta_2^-} \left(\beta_2^+ \beta_2^- + \frac{k'}{k} (\beta_2^+)^2 \right). \end{aligned}$$

Adding the contributions of q_3 and q_4 gives the term

$$\overline{\gamma_2} |i\eta_0 + u_r \beta_2^-|^2 \left\{ -\frac{\omega_2}{2\underline{a}_r} \left(\frac{\underline{a}_r^2}{c_r^2} + c_r^2 |\check{\eta}|^2 \right) + \frac{k'}{k} \omega_2 \frac{(\beta_2^+)^2 - |\check{\eta}|^2}{\beta_2^+ - \beta_2^-} \right\}. \quad (56)$$

When we add the latter term with the expression in (55), we obtain the constant term (independent of (k, k')):

$$-[\rho][u] \underline{\Upsilon} \overline{\gamma_2} |i\eta_0 + u_r \beta_2^-|^2 (\eta_0^2 + u_r^2 |\check{\eta}|^2) \frac{i u_r \eta_0}{2\underline{a}_r} (u_\ell \underline{a}_\ell + i c_\ell^2 \eta_0) = \frac{1}{2} \overline{Q}_r.$$

The only term that depends on (k, k') arises in (56) and equals

$$\frac{k'}{k} \overline{\gamma_2} |i\eta_0 + u_r \beta_2^-|^2 \omega_2 \frac{(i\eta_0 + u_r \beta_2^+)^2}{2c_r^2 \underline{a}_r} (c_r^2 - u_r^2) = \frac{k'}{2k} \overline{\gamma_2} (i\eta_0 + u_r \beta_2^+)^2 \frac{\eta_0^2 + u_r^2 |\check{\eta}|^2}{\underline{a}_r} \omega_2 = \frac{k'}{2k} \overline{Q}_r.$$

The sum of all the terms that involve $p''(\rho_r)$ factorizes as claimed in Proposition 3.

• Step 3: computation of the remaining terms. In order to prove Proposition 3, we can assume from now on, and without loss of generality that $p''(\rho_\ell) = p''(\rho_r) = 0$ in (34) and (35). With this simplification, we compute

$$\sum_{k=1}^{d-1} \check{\eta}_k d^2 f^k(v_\ell) \cdot (r_1^-, r_1^+) = -i \frac{c_\ell^4 |\check{\eta}|^2}{\rho_\ell} \begin{pmatrix} 0 \\ 2i\check{\eta} \\ -(\beta_1^+ + \beta_1^-) \end{pmatrix}, \quad (57)$$

$$d^2 \tilde{f}^d(v_\ell) \cdot (r_1^-, r_1^+) = \frac{c_\ell^4}{\rho_\ell} \begin{pmatrix} 0 \\ i(\beta_1^+ + \beta_1^-) \check{\eta} \\ -2\beta_1^+ \beta_1^- \\ -2u_\ell \beta_1^+ \beta_1^- \end{pmatrix}, \quad (58)$$

$$\sum_{k=1}^{d-1} \check{\eta}_k d^2 f^k(v_r) \cdot (r_2^-, r_2^+) = -i \frac{c_r^4 |\check{\eta}|^2}{\rho_r} \begin{pmatrix} 0 \\ 2i\check{\eta} \\ \beta_2^+ + \beta_2^- \end{pmatrix}, \quad (59)$$

$$d^2 \tilde{f}^d(v_r) \cdot (r_2^-, r_2^+) = -\frac{c_r^4}{\rho_r} \begin{pmatrix} 0 \\ i(\beta_2^+ + \beta_2^-) \check{\eta} \\ 2\beta_2^+ \beta_2^- \\ 2u_r \beta_2^+ \beta_2^- \end{pmatrix}, \quad (60)$$

As was done in Step 3 of the proof of Proposition 2, we are going to split the computations between the ‘left’ and ‘right’ contributions. Let us first concentrate on all ‘left’ terms. The ‘left’ contribution of q_2 is computed from (58) and equals

$$\begin{aligned} \sigma^* d^2 \tilde{f}^d(v_\ell) \cdot (x_1^-, x_1^+) |\gamma_1|^2 &= \sigma^* \frac{c_\ell^4}{\rho_\ell} \begin{pmatrix} 0 \\ i(\beta_1^+ + \beta_1^-) \tilde{y} \\ -2\beta_1^+ \beta_1^- \\ -2u_\ell \beta_1^+ \beta_1^- \end{pmatrix} |\gamma_1|^2 \\ &= \frac{2c_\ell^4}{\rho_\ell(c_\ell^2 - u_\ell^2)} \Upsilon \overline{\gamma_1} \left\{ u_\ell \eta_0 |\tilde{y}|^2 \gamma_1 \check{D} - (\eta_0^2 - c_\ell^2 |\tilde{y}|^2) \gamma_1 (D_{d+1} + u_\ell D_{d+2}) \right\}. \end{aligned}$$

Using Lemma 4, we find that the ‘left’ contribution of q_2 is given by

$$\frac{2c_\ell^4}{\rho_\ell(c_\ell^2 - u_\ell^2)} [\rho] [u] \Upsilon \overline{\gamma_1} \left\{ i c_r^2 \eta_0 (\eta_0^2 - c_\ell^2 |\tilde{y}|^2) (\eta_0^2 + u_r^2 |\tilde{y}|^2) - u_\ell u_r c_\ell^2 |\tilde{y}|^4 (u_r \underline{a}_r - i c_r^2 \eta_0) \right\}. \quad (61)$$

Similarly, we use (60) to compute the ‘right’ contribution of q_2 , which yields

$$\begin{aligned} -\sigma^* d^2 \tilde{f}^d(v_\ell) \cdot (x_2^-, x_2^+) |\gamma_2|^2 &= \frac{2c_r^4}{\rho_r(c_r^2 - u_r^2)} [\rho] [u] \Upsilon \overline{\gamma_2} \\ &\quad \times \left\{ (i c_\ell^2 \eta_0^3 + u_\ell \underline{a}_\ell u_r^2 |\tilde{y}|^2) (\eta_0^2 - c_r^2 |\tilde{y}|^2) - \eta_0^2 u_r^2 |\tilde{y}|^2 (u_\ell \underline{a}_\ell - i c_\ell^2 \eta_0) \right\}. \end{aligned} \quad (62)$$

Using (57), the ‘left’ contribution of q_3 equals

$$\begin{aligned} &i(k+k') \int_0^{+\infty} \frac{\omega_1}{\gamma_1} \tilde{\ell}_1 e^{-(k+k')\beta_1^+ z} \frac{-i c_\ell^4 |\tilde{y}|^2}{\rho_\ell} \begin{pmatrix} 0 \\ 2i \tilde{y} \\ -(\beta_1^+ + \beta_1^-) \end{pmatrix} |\gamma_1|^2 e^{k\beta_1^- z} e^{k'\beta_1^+ z} dz \\ &+ i(k+k') \int_0^{+\infty} \frac{\omega_3}{\gamma_1} \tilde{\ell}_3 e^{-(k+k')\beta_3^+ z} \frac{-i c_\ell^4 |\tilde{y}|^2}{\rho_\ell} \begin{pmatrix} 0 \\ 2i \tilde{y} \\ -(\beta_1^+ + \beta_1^-) \end{pmatrix} |\gamma_1|^2 e^{k\beta_1^- z} e^{k'\beta_1^+ z} dz \\ &= \frac{c_\ell^4}{\rho_\ell} \overline{\gamma_1} \omega_1 \tilde{\ell}_1 \begin{pmatrix} 0 \\ 2i \tilde{y} \\ -(\beta_1^+ + \beta_1^-) \end{pmatrix} \frac{|\tilde{y}|^2}{\beta_1^+ - \beta_1^-} \left(1 + \frac{k'}{k}\right) \\ &\quad + \underbrace{\frac{c_\ell^4}{\rho_\ell} \overline{\gamma_1} \omega_3 \tilde{\ell}_1 \begin{pmatrix} 0 \\ 2i \tilde{y} \\ -(\beta_1^+ + \beta_1^-) \end{pmatrix} \frac{(k+k') |\tilde{y}|^2}{k(\beta_3^+ - \beta_1^-) + k'(\beta_3^+ - \beta_1^+)}}_{\diamond}, \end{aligned} \quad (63)$$

and, similarly (using now (59) rather than (57)), the ‘right’ contribution of q_3 equals

$$\begin{aligned} &i(k+k') \int_0^{+\infty} \frac{\omega_2}{\gamma_2} \tilde{\ell}_2 e^{-(k+k')\beta_2^+ z} \frac{-i c_r^4 |\tilde{y}|^2}{\rho_r} \begin{pmatrix} 0 \\ 2i \tilde{y} \\ \beta_2^+ + \beta_2^- \end{pmatrix} |\gamma_2|^2 e^{k\beta_2^- z} e^{k'\beta_2^+ z} dz \\ &= \frac{c_r^4}{\rho_r} \overline{\gamma_2} \omega_2 \tilde{\ell}_2 \begin{pmatrix} 0 \\ 2i \tilde{y} \\ \beta_2^+ + \beta_2^- \end{pmatrix} \frac{|\tilde{y}|^2}{\beta_2^+ - \beta_2^-} \left(1 + \frac{k'}{k}\right). \end{aligned} \quad (64)$$

The ‘left’ contribution of q_4 is computed by retaining the three first rows in (58):

$$\begin{aligned}
& - \int_0^{+\infty} \frac{\omega_1}{\gamma_1} \tilde{\ell}_1 e^{-(k+k') \underline{\beta}_1^+ z} \frac{c_\ell^A}{\rho_\ell} \begin{pmatrix} 0 \\ i(\underline{\beta}_1^+ + \underline{\beta}_1^-) \check{y} \\ -2 \underline{\beta}_1^+ \underline{\beta}_1^- \end{pmatrix} |\gamma_1|^2 (k \underline{\beta}_1^- + k' \underline{\beta}_1^+) e^{k \underline{\beta}_1^- z} e^{k' \underline{\beta}_1^+ z} dz \\
& - \int_0^{+\infty} \frac{\omega_3}{\gamma_1} \tilde{\ell}_1 e^{-(k+k') \underline{\beta}_3^+ z} \frac{c_\ell^A}{\rho_\ell} \begin{pmatrix} 0 \\ i(\underline{\beta}_1^+ + \underline{\beta}_1^-) \check{y} \\ -2 \underline{\beta}_1^+ \underline{\beta}_1^- \end{pmatrix} |\gamma_1|^2 (k \underline{\beta}_1^- + k' \underline{\beta}_1^+) e^{k \underline{\beta}_1^- z} e^{k' \underline{\beta}_1^+ z} dz \\
& = - \frac{c_\ell^A}{\rho_\ell} \frac{\omega_1}{\gamma_1} \tilde{\ell}_1 \begin{pmatrix} 0 \\ i(\underline{\beta}_1^+ + \underline{\beta}_1^-) \check{y} \\ -2 \underline{\beta}_1^+ \underline{\beta}_1^- \end{pmatrix} \frac{1}{\underline{\beta}_1^+ - \underline{\beta}_1^-} \left(\underline{\beta}_1^- + \frac{k'}{k} \underline{\beta}_1^+ \right) \\
& \quad - \underbrace{\frac{c_\ell^A}{\rho_\ell} \frac{\omega_3}{\gamma_1} \tilde{\ell}_1 \begin{pmatrix} 0 \\ i(\underline{\beta}_1^+ + \underline{\beta}_1^-) \check{y} \\ -2 \underline{\beta}_1^+ \underline{\beta}_1^- \end{pmatrix} \frac{k \underline{\beta}_1^- + k' \underline{\beta}_1^+}{k(\underline{\beta}_3^+ - \underline{\beta}_1^-) + k'(\underline{\beta}_3^+ - \underline{\beta}_1^+)}}_{\diamond}, \tag{65}
\end{aligned}$$

and the ‘right’ contribution of q_4 equals

$$\begin{aligned}
& - \frac{c_r^A}{\rho_r} \frac{\omega_2}{\gamma_2} \tilde{\ell}_2 \begin{pmatrix} 0 \\ i(\underline{\beta}_2^+ + \underline{\beta}_2^-) \check{y} \\ 2 \underline{\beta}_2^+ \underline{\beta}_2^- \end{pmatrix} \frac{1}{\underline{\beta}_2^+ - \underline{\beta}_2^-} \left(\underline{\beta}_2^- + \frac{k'}{k} \underline{\beta}_2^+ \right) \\
& = \frac{c_r^A}{\rho_r} \frac{\omega_2}{\gamma_2} \tilde{\ell}_2 \begin{pmatrix} 0 \\ i(\underline{\beta}_2^+ + \underline{\beta}_2^-) \check{y} \\ 2 \underline{\beta}_2^+ \underline{\beta}_2^- \end{pmatrix} - \frac{c_r^A}{\rho_r} \frac{\omega_2}{\gamma_2} \tilde{\ell}_2 \begin{pmatrix} 0 \\ i(\underline{\beta}_2^+ + \underline{\beta}_2^-) \check{y} \\ 2 \underline{\beta}_2^+ \underline{\beta}_2^- \end{pmatrix} \frac{\underline{\beta}_2^+}{\underline{\beta}_2^+ - \underline{\beta}_2^-} \left(1 + \frac{k'}{k} \right). \tag{66}
\end{aligned}$$

The ‘left’ contribution of $q_3 + q_4$ is computed as follows. We first observe that the difference between the diamond term in (63) and the diamond term in (65) reads

$$\frac{c_\ell^2}{\rho_\ell} \frac{\omega_3}{\gamma_1} u_\ell |\check{y}|^2 \left\{ c_\ell^2 |\check{y}|^2 - c_\ell^2 \underline{\beta}_1^+ \underline{\beta}_1^- - (i \eta_0 - u_\ell \underline{\beta}_1^+) (i \eta_0 - u_\ell \underline{\beta}_1^-) \right\} = \frac{2 c_\ell^6}{\rho_\ell (c_\ell^2 - u_\ell^2)} u_\ell |\check{y}|^4 \overline{\gamma_1} \omega_3.$$

Consequently the ‘left’ contribution of $q_3 + q_4$, which corresponds to the sum of (63) and (65), equals

$$\begin{aligned}
& \frac{c_\ell^A \overline{\gamma_1} \omega_1}{\rho_\ell (\underline{\beta}_1^+ - \underline{\beta}_1^-)} \left\{ |\check{y}|^2 \tilde{\ell}_1 \begin{pmatrix} 0 \\ 2i \check{y} \\ -(\underline{\beta}_1^+ + \underline{\beta}_1^-) \end{pmatrix} - \underline{\beta}_1^+ \tilde{\ell}_1 \begin{pmatrix} 0 \\ i(\underline{\beta}_1^+ + \underline{\beta}_1^-) \check{y} \\ -2 \underline{\beta}_1^+ \underline{\beta}_1^- \end{pmatrix} \right\} \left(1 + \frac{k'}{k} \right) \\
& + \frac{c_\ell^A}{\rho_\ell} \overline{\gamma_1} \omega_1 \tilde{\ell}_1 \begin{pmatrix} 0 \\ i(\underline{\beta}_1^+ + \underline{\beta}_1^-) \check{y} \\ -2 \underline{\beta}_1^+ \underline{\beta}_1^- \end{pmatrix} + \frac{2 c_\ell^6}{\rho_\ell (c_\ell^2 - u_\ell^2)} u_\ell |\check{y}|^4 \overline{\gamma_1} \omega_3 \\
& = \frac{2 c_\ell^A}{\rho_\ell (\underline{\beta}_1^+ - \underline{\beta}_1^-)} \overline{\gamma_1} \omega_1 (\underline{\beta}_1^+ \underline{\beta}_1^- - |\check{y}|^2) ((\underline{\beta}_1^+)^2 - |\check{y}|^2) \left(1 + \frac{k'}{k} \right) \\
& \quad - \frac{2 c_\ell^A}{\rho_\ell (c_\ell^2 - u_\ell^2)} \overline{\gamma_1} \omega_1 \left\{ (\eta_0^2 - c_\ell^2 |\check{y}|^2) \underline{\beta}_1^+ + i u_\ell \eta_0 |\check{y}|^2 \right\} + \frac{2 c_\ell^6}{\rho_\ell (c_\ell^2 - u_\ell^2)} u_\ell |\check{y}|^4 \overline{\gamma_1} \omega_3.
\end{aligned}$$

After a little bit of algebra, the ‘left’ contribution of $q_3 + q_4$ is found to be equal to

$$-\frac{c_\ell^2}{\rho_\ell} [\rho] [u] \underline{\Upsilon} u_\ell u_r \frac{a_r}{a_\ell} (\eta_0^2 + u_r^2 |\check{\eta}|^2) \left(\frac{a_\ell^2}{c_\ell^2} + c_\ell^2 |\check{\eta}|^2 \right) \overline{\gamma}_1 (i \eta_0 - u_\ell \beta_1^+) \left(1 + \frac{k'}{k} \right) \\ + \frac{2c_\ell^4}{\rho_\ell (c_\ell^2 - u_\ell^2)} [\rho] [u] \underline{\Upsilon} \overline{\gamma}_1 \left\{ -i c_r^2 \eta_0 (\eta_0^2 - c_\ell^2 |\check{\eta}|^2) (\eta_0^2 + u_r^2 |\check{\eta}|^2) + u_\ell u_r c_\ell^2 |\check{\eta}|^4 (u_r a_r - i c_r^2 \eta_0) \right\}.$$

When combined with (61), we have thus shown that the ‘left’ contribution of the kernel $q_2 + q_3 + q_4$ reads

$$-\frac{c_\ell^2}{\rho_\ell} [\rho] [u] \underline{\Upsilon} u_\ell u_r \frac{a_r}{a_\ell} (\eta_0^2 + u_r^2 |\check{\eta}|^2) \left(\frac{a_\ell^2}{c_\ell^2} + c_\ell^2 |\check{\eta}|^2 \right) \overline{\gamma}_1 (i \eta_0 - u_\ell \beta_1^+) \left(1 + \frac{k'}{k} \right) \\ = - \left(\frac{c_\ell^2}{\rho_\ell} \overline{Q}_\ell + 2 [\rho] [u] \underline{\Upsilon} u_\ell u_r |\check{\eta}|^2 (\eta_0^2 + u_r^2 |\check{\eta}|^2) \frac{c_\ell^4 a_r}{\rho_\ell a_\ell} \overline{\gamma}_1 (i \eta_0 - u_\ell \beta_1^+) \right) \left(1 + \frac{k'}{k} \right). \quad (67)$$

Let us now compute the ‘right’ contribution of the kernel $q_3 + q_4$. We add the expressions in (64) and (66), which gives

$$\frac{2c_r^4}{\rho_r} \overline{\gamma}_2 \omega_2 \left\{ (\beta_2^+ \beta_2^-) \beta_2^+ - \frac{\beta_2^+ + \beta_2^-}{2} |\check{\eta}|^2 \right\} - \frac{2c_r^4}{\rho_r} \overline{\gamma}_2 \omega_2 ((\beta_2^+)^2 - |\check{\eta}|^2) \frac{\beta_2^+ \beta_2^- - |\check{\eta}|^2}{\beta_2^+ - \beta_2^-} \left(1 + \frac{k'}{k} \right),$$

or equivalently

$$\frac{2c_r^4}{\rho_r (c_r^2 - u_r^2)} [\rho] [u] \underline{\Upsilon} \overline{\gamma}_2 i u_r \eta_0 (u_\ell a_\ell - i c_\ell^2 \eta_0) \left\{ (\eta_0^2 - c_r^2 |\check{\eta}|^2) \beta_2^+ - i u_r \eta_0 |\check{\eta}|^2 \right\} \\ - \frac{c_r^4}{\rho_r} [\rho] [u] \underline{\Upsilon} u_\ell u_r \frac{a_\ell}{a_r} (\eta_0^2 + u_r^2 |\check{\eta}|^2) \left(\frac{a_r^2}{c_r^2} + c_r^2 |\check{\eta}|^2 \right) \overline{\gamma}_2 (i \eta_0 + u_r \beta_2^+) \left(1 + \frac{k'}{k} \right). \quad (68)$$

The first row in (68) is exactly the opposite of (62), that is of the ‘right’ contribution of q_2 . In other words, we have found that the ‘right’ contribution of $q_2 + q_3 + q_4$ is given by the second row in (68), which reads

$$- \left(\frac{c_r^2}{\rho_r} \overline{Q}_r + 2 [\rho] [u] \underline{\Upsilon} u_\ell u_r |\check{\eta}|^2 (\eta_0^2 + u_r^2 |\check{\eta}|^2) \frac{c_r^4 a_\ell}{\rho_r a_r} \overline{\gamma}_2 (i \eta_0 + u_r \beta_2^+) \right) \left(1 + \frac{k'}{k} \right).$$

The kernel $q_2 + q_3 + q_4$ is the sum of the latter quantity and the right hand side of (67). Collecting the terms, this completes the proof of Proposition 3. \square

Corollary 2. *With Q defined in (32), $Q_\ell, Q_r, Q_\sharp, Q_b$ defined in Propositions 2 and 3, the kernel q satisfies*

$$q(k, k') = \begin{cases} \left(\frac{p''(\rho_\ell)}{2} + \frac{c_\ell^2}{\rho_\ell} \right) Q_\ell + \left(\frac{p''(\rho_r)}{2} + \frac{c_r^2}{\rho_r} \right) Q_r + Q_\sharp & \text{if } k > 0 \text{ and } k' > 0, \\ \left\{ \left(\frac{p''(\rho_\ell)}{2} - \frac{c_\ell^2}{\rho_\ell} \right) \overline{Q}_\ell + \left(\frac{p''(\rho_r)}{2} - \frac{c_r^2}{\rho_r} \right) \overline{Q}_r + Q_b + \overline{Q} \right\} \left(1 + \frac{k'}{k} \right) & \text{if } k > 0, k' < 0, k + k' > 0. \end{cases}$$

4 Conclusion

There remains a final simplification in order to achieve the final form of the kernel q . Our main result reads as follows.

Proposition 4. *With Q_ℓ, Q_r, Q_\sharp defined in Propositions 2, let us define*

$$Q_\natural := \left(\frac{p''(\rho_\ell)}{2} + \frac{c_\ell^2}{\rho_\ell} \right) Q_\ell + \left(\frac{p''(\rho_r)}{2} + \frac{c_r^2}{\rho_r} \right) Q_r + Q_\sharp.$$

Then the kernel $q = 4\pi a_1$ satisfies

$$q(k, k') = \begin{cases} Q_\natural & \text{if } k > 0 \text{ and } k' > 0, \\ \frac{Q_\natural}{1 + \frac{k'}{k}} & \text{if } k > 0, k' < 0 \text{ and } k + k' > 0. \end{cases}$$

In particular, q satisfies Hunter's stability condition $q(1, 0^+) = \overline{q(1, 0^-)}$, and (15) is locally well-posed in $H^2(\mathbb{R})$.

Proof. In view of the expression of q given in Corollary 2, we only need to show the relation

$$\frac{c_\ell^2}{\rho_\ell} Q_\ell + \frac{c_r^2}{\rho_r} Q_r + \frac{1}{2} Q_\sharp = \frac{1}{2} (Q + \overline{Q_b}),$$

with Q_ℓ, Q_r, Q_\sharp given in Proposition 2, Q_b in Proposition 3, and Q in Lemma 1. Simplifying by the common factor $[\rho][u]\Upsilon(\varpi_0^2 + u_r^2 |\check{\varpi}|^2) u_\ell u_r$, we are led to showing that the following relation holds:

$$\begin{aligned} & \frac{\underline{a}_r}{\rho_\ell \underline{a}_\ell} c_\ell^2 (\varpi_0^2 + u_\ell^2 |\check{\varpi}|^2) \gamma_1 (i \varpi_0 - u_\ell \underline{\beta}_1^-) + \frac{\underline{a}_\ell}{\rho_r \underline{a}_r} c_r^2 (\varpi_0^2 + u_r^2 |\check{\varpi}|^2) \gamma_2 (i \varpi_0 + u_r \underline{\beta}_2^-) \\ & - \frac{1}{\mathbf{j}[u] \varpi_0} (\varpi_0^2 + u_\ell u_r |\check{\varpi}|^2) i \underline{a}_\ell \underline{a}_r (c_r^2 \gamma_2 - c_\ell^2 \gamma_1) = (\underline{\beta}_1^- + \underline{\beta}_2^-) i \underline{a}_\ell \underline{a}_r \frac{u_\ell \underline{a}_r + i c_r^2 \varpi_0}{u_\ell \underline{a}_r - i c_r^2 \varpi_0} \\ & + \frac{\underline{a}_r}{\rho_\ell \underline{a}_\ell} c_\ell^4 |\check{\varpi}|^2 \gamma_1 (i \varpi_0 - u_\ell \underline{\beta}_1^-) + \frac{\underline{a}_\ell}{\rho_r \underline{a}_r} c_r^4 |\check{\varpi}|^2 \gamma_2 (i \varpi_0 + u_r \underline{\beta}_2^-), \quad (69) \end{aligned}$$

where we use the notation $\mathbf{j} := \rho_\ell u_\ell = \rho_r u_r$ to denote the mass flux across the phase boundary.

The verification of (69) proceeds as follows. We first combine the first and third row in (69) by recalling the definitions of $\underline{a}_\ell, \underline{a}_r$, see (1) and (2). Thus verifying (69) amounts to showing

$$\begin{aligned} & (\underline{\beta}_1^- + \underline{\beta}_2^-) \frac{u_\ell \underline{a}_r + i c_r^2 \varpi_0}{u_\ell \underline{a}_r - i c_r^2 \varpi_0} - i \frac{\gamma_1}{\rho_\ell} (i \varpi_0 - u_\ell \underline{\beta}_1^-) - i \frac{\gamma_2}{\rho_r} (i \varpi_0 + u_\ell \underline{\beta}_2^-) \\ & + \frac{1}{\mathbf{j}[u] \varpi_0} (\varpi_0^2 + u_\ell u_r |\check{\varpi}|^2) (c_r^2 \gamma_2 - c_\ell^2 \gamma_1) = 0. \quad (70) \end{aligned}$$

Let us now define

$$\begin{aligned} B_\ell & := \underline{\beta}_1^- \frac{u_\ell \underline{a}_r + i c_r^2 \varpi_0}{u_\ell \underline{a}_r - i c_r^2 \varpi_0} - i \frac{\gamma_1}{\rho_\ell} (i \varpi_0 - u_\ell \underline{\beta}_1^-) - \frac{1}{\mathbf{j}[u] \varpi_0} (\varpi_0^2 + u_\ell u_r |\check{\varpi}|^2) c_\ell^2 \gamma_1, \\ B_r & := \underline{\beta}_2^- \frac{u_\ell \underline{a}_r + i c_r^2 \varpi_0}{u_\ell \underline{a}_r - i c_r^2 \varpi_0} - i \frac{\gamma_2}{\rho_r} (i \varpi_0 + u_\ell \underline{\beta}_2^-) + \frac{1}{\mathbf{j}[u] \varpi_0} (\varpi_0^2 + u_\ell u_r |\check{\varpi}|^2) c_r^2 \gamma_2, \end{aligned}$$

so that (70), which is the relation we wish to prove, reads $B_\ell + B_r = 0$.

Using the relations

$$\frac{u_\ell \underline{a}_r + i c_r^2 \eta_0}{u_\ell \underline{a}_r - i c_r^2 \eta_0} = -\frac{u_r \underline{a}_\ell + i c_\ell^2 \eta_0}{u_r \underline{a}_\ell - i c_\ell^2 \eta_0}, \quad (u_r \underline{a}_\ell - i c_\ell^2 \eta_0) \gamma_1 = -\rho_\ell [u] \eta_0,$$

we get

$$u_\ell (u_r \underline{a}_\ell - i c_\ell^2 \eta_0) B_\ell = -i \eta_0 (u_\ell \underline{a}_\ell + i c_\ell^2 \eta_0). \quad (71)$$

Similarly, we compute

$$u_r (u_\ell \underline{a}_r - i c_r^2 \eta_0) B_r = -i \eta_0 (u_r \underline{a}_r + i c_r^2 \eta_0). \quad (72)$$

Combining the relations (71) and (72), we get

$$\begin{aligned} & u_\ell u_r (u_r \underline{a}_\ell - i c_\ell^2 \eta_0) (u_\ell \underline{a}_r - i c_r^2 \eta_0) (B_\ell + B_r) \\ &= -i \eta_0 u_r (u_\ell \underline{a}_r - i c_r^2 \eta_0) (u_\ell \underline{a}_\ell + i c_\ell^2 \eta_0) - i \eta_0 u_\ell (u_r \underline{a}_\ell - i c_\ell^2 \eta_0) (u_r \underline{a}_r + i c_r^2 \eta_0) \\ &= -i \eta_0 (u_\ell + u_r) (u_\ell u_r \underline{a}_\ell \underline{a}_r + c_\ell^2 c_r^2 \eta_0^2) = 0. \end{aligned}$$

This means that (70) holds, and consequently the expression of the kernel q is as claimed in Proposition 4. The verification of Hunter's stability condition $q(1, 0^+) = \overline{q(1, 0^-)}$ is then straightforward, and local well-posedness in $H^2(\mathbb{R})$ follows from the main result in [2]. \square

References

- [1] S. Benzoni-Gavage. Stability of multi-dimensional phase transitions in a van der Waals fluid. *Nonlinear Anal.*, 31(1-2):243–263, 1998.
- [2] S. Benzoni-Gavage. Local well-posedness of nonlocal Burgers equations. *Differential Integral Equations*, 22(3-4):303–320, 2009.
- [3] S. Benzoni-Gavage and M. Rosini. Weakly nonlinear surface waves and subsonic phase boundaries. *Comput. Math. Appl.*, 57(3-4):1463–1484, 2009.
- [4] J. K. Hunter. Nonlinear surface waves. In *Current progress in hyperbolic systems: Riemann problems and computations (Brunswick, ME, 1988)*, volume 100 of *Contemp. Math.*, pages 185–202. Amer. Math. Soc., 1989.