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EXISTENCE AND SYMMETRY OF LEAST ENERGY NODAL SOLUTIONS FOR HAMILTONIAN ELLIPTIC SYSTEMS

DENIS BONHEURE, EDERSON MOREIRA DOS SANTOS, MIGUEL RAMOS,
AND HUGO TAVARES

ABSTRACT. In this paper we prove existence of least energy nodal solutions for the Hamiltonian elliptic system with Hénon-type weights

$$-\Delta u = |x|^\beta |v|^{q-1} v, \quad -\Delta v = |x|^\alpha |u|^{p-1} u \quad \text{in } \Omega, \quad u = v = 0 \text{ on } \partial\Omega,$$

where Ω is a bounded smooth domain in \mathbb{R}^N , $N \geq 1$, $\alpha, \beta \geq 0$ and the nonlinearities are superlinear and subcritical, namely

$$1 > \frac{1}{p+1} + \frac{1}{q+1} > \frac{N-2}{N}.$$

When Ω is either a ball or an annulus centered at the origin and $N \geq 2$, we show that these solutions display the so-called foliated Schwarz symmetry. It is natural to conjecture that these solutions are not radially symmetric. We provide such a symmetry breaking in a range of parameters where the solutions of the system behave like the solutions of a single equation. Our results on the above system are new even in the case of the Lane-Emden system (i.e. without weights). As far as we know, this is the first paper that contains results about least energy nodal solutions for strongly coupled elliptic systems and their symmetry properties.

1. INTRODUCTION

We consider the Hamiltonian elliptic system with Hénon-type weights

$$-\Delta u = |x|^\beta |v|^{q-1} v, \quad -\Delta v = |x|^\alpha |u|^{p-1} u \quad \text{in } \Omega, \quad u = v = 0 \text{ on } \partial\Omega, \quad (1.1)$$

where Ω is a bounded domain in \mathbb{R}^N , $N \geq 1$, and $\alpha, \beta \geq 0$. We consider superlinear and subcritical nonlinearities, namely

$$1 > \frac{1}{p+1} + \frac{1}{q+1} > \frac{N-2}{N}. \quad (\text{H})$$

Observe that the first condition is also equivalent to $pq > 1$.

The system (1.1) is strongly coupled in the sense that $u \equiv 0$ if and only if $v \equiv 0$. Moreover, u changes sign if and only if v changes sign.

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We recall that a strong solution to this problem corresponds to a pair (u, v) with

$$u \in W^{2,(q+1)/q}(\Omega) \cap W_0^{1,(q+1)/q}(\Omega), \quad v \in W^{2,(p+1)/p}(\Omega) \cap W_0^{1,(p+1)/p}(\Omega)$$

satisfying the system in (1.1) for a.e. $x \in \Omega$. By using a bootstrap method (see [37, Theorem 1(a)]), it can be shown that strong solutions are actually classical solutions.

Consider the energy functional

$$E(u, v) = \int_{\Omega} \nabla u \cdot \nabla v \, dx - \frac{1}{p+1} \int_{\Omega} |x|^{\alpha} |u|^{p+1} \, dx - \frac{1}{q+1} \int_{\Omega} |x|^{\beta} |v|^{q+1} \, dx, \quad (1.2)$$

which is well defined for strong solutions thanks to assumption (H).

One can use various variational settings to deal with the system (1.1), see for instance the surveys [10, 21, 34]. Once the existence of at least one critical point is proved, a natural question is that of the existence of a least energy one, by which we mean a critical point at the level

$$c = \inf\{E(u, v) : (u, v) \text{ is a nonzero strong solution of (1.1)}\}.$$

The solutions at this energy are usually referred to as ground state solutions and in many problems, they are of special interest. In our setting the existence of such solutions is clear and rely on a simple compactness argument. On the other hand, it is useful to get a variational characterisation of these solutions to derive qualitative properties, see for example [8, 9, 10]. In our setting, any solution at level c is positive (or negative). This can be established using a Nehari type characterization of the level c . We emphasize that the adequate associated Nehari manifold is then of infinite codimension. We refer to [10] for more details.

Existence of sign-changing solutions has been obtained in [32, Theorem 4], under the extra assumption $p > 1$ and $q > 1$, where it is proved that there exists an unbounded sequence of solutions (u_k, v_k) such that both $(u_k + v_k)^+ \neq 0$ and $(u_k + v_k)^- \neq 0$ for every k . In fact, for a pair of solutions (u, v) , $u + v$ changes sign if and only if u and v change sign. Our results therefore present some improvement of the result given in [32, Theorem 4] since, by imposing the mere super linearity condition $pq > 1$, we are able to prove the existence of a (least energy) nodal solution to (1.1).

Define then the *least energy nodal level* as

$$c_{\text{nod}} = \inf\{E(u, v) : (u, v) \text{ is a nonzero strong solution of (1.1) and } u^{\pm}, v^{\pm} \neq 0\}.$$

It is not obvious that this level is achieved since this no more follows from a simple minimisation argument. Indeed, even if we have enough compactness to extract a converging subsequence, the limit could be a critical point (u, v) such that both u and v are positive (or negative). The existence of a least energy nodal solution for the scalar Lane-Emden equation [14, 3, 4] follows from the minimisation of the functional over a nodal Nehari set. It is not clear at all how such a nodal Nehari

set associated to the energy functional E could be defined. Anyhow, our first main result shows c_{nod} is achieved.

Theorem 1.1. *Let $N \geq 1$, $\alpha \geq 0$, $\beta \geq 0$ and suppose that (H) is satisfied. Then the level c_{nod} is achieved, that is, there exists a strong solution (u, v) of (1.1) such that $u^\pm, v^\pm \neq 0$ and $E(u, v) = c_{\text{nod}}$.*

Our proof relies on a dual method as in Clément and van der Vorst [15] or Alves and Soares [2]¹ who deal with the singularly perturbed system

$$-\varepsilon^2 \Delta u + V(x)u = |v|^{q-1}v, \quad -\varepsilon^2 \Delta v + V(x)v = |u|^{p-1}u \quad \text{in } \mathbb{R}^N$$

assuming the extra assumptions $p > 1$ and $q > 1$. With respect to [2], the hypothesis (H) includes more general powers, namely $pq > 1$ is enough. This means in particular that we cover the biharmonic operator with Hénon weight, that is $q = 1$ and $\beta = 0$, with Navier boundary conditions. In this context, the problem (1.1) reads as

$$\Delta^2 u = |x|^\alpha |u|^{p-1}u \quad \text{in } \Omega, \quad u = \Delta u = 0 \quad \text{on } \partial\Omega, \quad (1.3)$$

with $\alpha \geq 0$, $\frac{1}{2} > \frac{1}{p+1} > \frac{N-4}{2N}$ and Theorem 1.1 applies.

Next we investigate the symmetry of these solutions in case the domain is radial. Let Ω be either a ball or an annulus centred at the origin. Recall that a function $u : \Omega \rightarrow \mathbb{R}$ is called foliated Schwarz symmetric with respect to some unit vector $p \in \mathbb{R}^N$ if, for a.e. $r > 0$ such that $\partial B_r(0) \subset \Omega$ and for every $c \in \mathbb{R}$, the restricted superlevel set $\{x \in \partial B_r(0) : u(x) \geq c\}$ is either equal to $\partial B_r(0)$ or to a geodesic ball in $\partial B_r(0)$ centred at rp . In other words, u is foliated Schwarz symmetric if u is axially symmetric with respect to the axis $\mathbb{R}p$ and non increasing in the polar angle $\theta = \arccos(x \cdot p) \in [0, \pi]$.

In the past thirty years the study on the symmetry properties of positive or least energy solutions of strongly coupled elliptic systems has been an active research field, see for instance [42, 35, 20, 22, 45, 13, 31, 8, 9, 10, 17, 18, 19]. The basic tools used to derive the symmetry of the solutions are the method of moving planes [25], symmetrization or polarization and Morse index arguments. All these techniques were originally developed for second order elliptic equations and their use in the context of strongly coupled elliptic system requires more involved arguments.

Our second main result is the following.

Theorem 1.2. *Let $\Omega \subset \mathbb{R}^N$, $N \geq 2$, be either a ball or an annulus centred at the origin. Let (u, v) be a least energy nodal solution of (1.1). Then there exists $p \in \partial B_1(0)$ such that both u and v are foliated Schwarz symmetric with respect to p .*

As mentioned above, our results cover the biharmonic operator complementing therefore some of the results in [43].

¹In order to apply the dual variational method, the two potentials on the left hand sides of [2, eq. (1.3)] must be equal as follows from [2, line 5 p. 114].

Corollary 1.3. *Let $\Omega \subset \mathbb{R}^N$, $N \geq 1$ and assume that $\frac{1}{2} > \frac{1}{p+1} > \frac{N-4}{2N}$. Then the fourth order problem (1.3) admits a least energy nodal solution. Moreover, if Ω is either a ball or an annulus centred at the origin, $N \geq 2$, then any least energy nodal solution of (1.3) is such that u and $-\Delta u$ are foliated Schwarz symmetric with respect to the the same unit vector $p \in \mathbb{R}^N$.*

For the scalar Lane-Emden equation (i.e. without weights), it is known that any least energy nodal solution has Morse index 2. Combined with the analysis of the Morse index of the sign changing radial solutions when Ω is either a ball or an annulus, this leads to the conclusion that whereas least energy solutions are radially symmetric, least energy nodal solutions are not. The foliated Schwarz symmetry is thus somehow optimal.

For the Hénon-Lane-Emden system (1.1), it is not clear how to compute (or even define) the Morse index of the solutions. Although we conjecture that for any p, q satisfying (H) and $\alpha, \beta \geq 0$ every least energy nodal solutions of (1.1) are non radial, we are not able to prove it. Symmetry breaking occurs at least for $p \sim q$, $\alpha \sim 0$ and $\beta \sim 0$.

Theorem 1.4. *Assume $N \geq 2$ and $\Omega \subset \mathbb{R}^N$ is either a ball or an annulus centred at the origin. Let q_0 satisfy*

$$q_0 > 1 \quad \text{and} \quad q_0 + 1 < 2N/(N - 2) \quad \text{if } N \geq 3.$$

Then there exists $\delta_0 > 0$ such that, if $p, q \in [q_0 - \delta_0, q_0 + \delta_0]$, $\alpha, \beta \in [0, \delta_0]$, then any least energy nodal solution (u, v) of (1.1) is such that both u and v are non radially symmetric.

At this point, we emphasize that when $\Omega \subset \mathbb{R}^N$ is either a ball or an annulus, we can work in a functional framework of radially symmetric functions yielding at least one radial sign-changing solution having least energy among all radial nodal solutions. The previous theorem gives therefore a range of coexistence of both radially symmetric and non radially symmetric sign-changing solutions. When $p > 1$ and $q > 1$, the existence of infinitely many radial sign-changing solutions follows also from applying the method of [32, Theorem 4] in a functional framework of radially symmetric functions. Although these solutions can a priori coincide with the solutions obtained in [32, Theorem 4], we clearly do not expect that to happen. A difficult question seems to be that of providing a precise information on the number of nodes of radial nodal solutions. Both the gluing approach [6] and an optimal partition method [16], reminiscent of the original approach of Nehari, seem out of reach for the Lane-Emden system (without weight). Therefore, the existence of radially symmetric solutions with a prescribed number of nodes is a challenging open question.

Finally, we show that all of our results apply to the Hénon equation

$$-\Delta u = |x|^\alpha |u|^{p-1} u \quad \text{in } \Omega, \quad u = 0 \quad \text{on } \partial\Omega. \quad (1.4)$$

For this purpose we prove a kind of symmetry theorem for the components, which guarantees that when $p = q$, $\alpha = \beta$, then any solution (u, v) of (1.1) is such that

$u = v$ and whence u solves (1.4). For related classifications results we mention [31, Theorem 1.2] and [24].

Theorem 1.5. *Assume that $u, v \in H_0^1(\Omega)$ solves the system (1.1)*

$$-\Delta u = |x|^\beta |v|^{q-1} v, \quad -\Delta v = |x|^\beta |u|^{q-1} u \quad \text{in } \Omega, \quad u = v = 0 \quad \text{on } \partial\Omega,$$

where Ω is a bounded domain in \mathbb{R}^N , $N \geq 1$, $q > 1$ and $q + 1 \leq 2N/(N - 2)$ if $N \geq 3$. Then $u = v$.

Combining Theorems 1.1, 1.2, 1.4 and 1.5, we get the following results about least energy nodal solutions of the Hénon equation (1.4).

Corollary 1.6. *Let $N \geq 1$, $\alpha \geq 0$, $p > 1$ and $p + 1 < 2N/(N - 2)$ if $N \geq 3$. Then:*

- i) There exist least energy nodal solutions of (1.4).*
- ii) Let $\Omega \subset \mathbb{R}^N$, $N \geq 2$, be either a ball or an annulus centred at the origin. Let u be a least energy nodal solution of (1.4). Then there exists $p \in \partial B_1(0)$ such that u is foliated Schwarz symmetric with respect to p .*
- iii) Assume that $\Omega \subset \mathbb{R}^N$, $N \geq 2$, is either a ball or an annulus centred at the origin. Then there exists $\delta_0 > 0$ such that, if $\alpha \in [0, \delta_0]$, then any least energy nodal solution u of (1.4) is such that u is non radially symmetric.*

Items i) and ii) are known cf. [5, 14, 30]. To our knowledge, Corollary 1.6 iii) is new, though expected from a perturbation analysis for small α . However, we stress that the approach of [1] to symmetric breaking of any least energy nodal solution of the autonomous equation

$$-\Delta u = f(u) \quad \text{in } \Omega, \quad u = 0 \quad \text{on } \partial\Omega,$$

where $\Omega \subset \mathbb{R}^N$, $N \geq 2$, that stands either for a ball or an annulus centred at the origin, cannot be extended to the non autonomous case. In particular it cannot be extended to the Hénon equation (1.4). We expect however that least energy nodal solutions of (1.4) are non radial. We provide a proof for α small whereas this should follow from an asymptotic analysis as in [39] for α large. The general case seems more delicate.

The structure of the paper is the following. In Section 2, we introduce the variational setting corresponding to the dual method and prove Theorem 1.1, showing the existence of a least energy nodal solution, providing as well alternative characterizations of the level c_{nod} . In Section 3 we prove the Schwarz foliated symmetry of these solutions when Ω is a radial bounded domain, namely Theorem 1.2. Finally in Section 4 we provide some examples of symmetry breaking, proving Theorems 1.4 and 1.5.

2. EXISTENCE OF A LEAST ENERGY NODAL LEVEL

Let us now introduce in a precise way the variational setting corresponding to the dual method. Given $r \geq 1$ and $\gamma \geq 0$, we denote

$$L^r(\Omega, |x|^{-\gamma}) := \{u : \Omega \rightarrow \mathbb{R} \text{ measurable} : \int_{\Omega} |u|^r |x|^{-\gamma} dx < \infty\},$$

which is a Banach space equipped with the norm

$$\|u\|_{r,\gamma} := \left(\int_{\Omega} |u|^r |x|^{-\gamma} dx \right)^{1/r}.$$

Observe that, since Ω is bounded and $\gamma \geq 0$, we have the inclusions $L^r(\Omega, |x|^{-\gamma}) \subset L^r(\Omega)$, where the last is the usual $L^r(\Omega)$ -space. In fact, it is easy to check that there exists a constant $C(\Omega)$ such that

$$\|u\|_r \leq C(\Omega)^{\frac{\gamma}{r}} \|u\|_{r,\gamma} \quad \forall u \in L^r(\Omega, |x|^{-\gamma}), \quad r > 1, \quad \gamma \geq 0. \quad (2.1)$$

In an informal basis, the method consists in taking the inverse of the Laplace operator, rewriting the system as

$$(-\Delta)^{-1}(|x|^{\beta}|v|^{q-1}v) = u, \quad (-\Delta)^{-1}(|x|^{\alpha}|u|^{p-1}u) = v.$$

and defining $w_1 = |x|^{\alpha}|u|^{p-1}u$, $w_2 = |x|^{\beta}|v|^{q-1}v$, which leads to

$$(-\Delta)^{-1}w_2 = |x|^{-\frac{\alpha}{p}}|w_1|^{\frac{1}{p}-1}w_1, \quad (-\Delta)^{-1}w_1 = |x|^{-\frac{\beta}{q}}|\omega_2|^{\frac{1}{q}-1}\omega_2. \quad (2.2)$$

We will work in the product space

$$X := L^{\frac{p+1}{p}}(\Omega, |x|^{-\frac{\alpha}{p}}) \times L^{\frac{q+1}{q}}(\Omega, |x|^{-\frac{\beta}{q}}),$$

$$\|(w_1, w_2)\| := \|w_1\|_{\frac{p+1}{p}, \frac{\alpha}{p}} + \|w_2\|_{\frac{q+1}{q}, \frac{\beta}{q}} \quad \forall w = (w_1, w_2) \in X$$

and will use the map $T : X \rightarrow L^1(\Omega)$ given by

$$Tw = w_1Kw_2 + w_2Kw_1 \quad w = (w_1, w_2) \in X$$

where, with some abuse of notations, K denotes the inverse of the minus Laplace operator with zero Dirichlet boundary condition. We observe that we use the same letter K to denote both the operators $K_{(p+1)/p} : L^{\frac{p+1}{p}}(\Omega) \rightarrow W^{2, \frac{p+1}{p}}(\Omega) \cap W_0^{1, \frac{p+1}{p}}(\Omega)$ and $K_{(q+1)/q} : L^{\frac{q+1}{q}}(\Omega) \rightarrow W^{2, \frac{q+1}{q}}(\Omega) \cap W_0^{1, \frac{q+1}{q}}(\Omega)$. Later on, we will use the fact that

$$u \in L^t(\Omega) \cap L^s(\Omega) \Rightarrow K_t u = K_s u \quad (t, s > 1),$$

which is a consequence of the uniqueness of the Dirichlet problem. Thanks to the subcriticality, namely the second inequality at (H), we have compact embeddings

$$W^{2, \frac{q+1}{q}}(\Omega) \subset L^{p+1}(\Omega), \quad W^{2, \frac{p+1}{p}}(\Omega) \subset L^{q+1}(\Omega).$$

This, together with standard elliptic estimates, yields

$$\begin{aligned} \int_{\Omega} |w_1Kw_2| dx &\leq \|w_1\|_{\frac{p+1}{p}} \|Kw_2\|_{p+1} \leq C_1 \|w_1\|_{\frac{p+1}{p}} \|Kw_2\|_{W^{2, \frac{q+1}{q}}} \\ &\leq C_2 \|w_1\|_{\frac{p+1}{p}} \|w_2\|_{\frac{q+1}{q}} \leq C_3 \|w_1\|_{\frac{p+1}{p}, \frac{\alpha}{p}} \|w_2\|_{\frac{q+1}{q}, \frac{\beta}{q}}, \end{aligned}$$

and an analogous estimate holds for w_2Kw_1 . Thus

$$\int_{\Omega} |Tw| dx \leq c \|w_1\|_{\frac{p+1}{p}, \frac{\alpha}{p}} \|w_2\|_{\frac{q+1}{q}, \frac{\beta}{q}} \quad \forall w = (w_1, w_2) \in X. \quad (2.3)$$

Also, using integration by parts,

$$\int_{\Omega} w_1 K w_2 dx = \int_{\Omega} w_2 K w_1 dx \quad \forall w = (w_1, w_2) \in X. \quad (2.4)$$

Now, let $I : X \rightarrow \mathbb{R}$ be the functional defined by

$$I(w_1, w_2) = \frac{p}{p+1} \int_{\Omega} |w_1|^{\frac{p+1}{p}} |x|^{-\frac{\alpha}{p}} dx + \frac{q}{q+1} \int_{\Omega} |w_2|^{\frac{q+1}{q}} |x|^{-\frac{\beta}{q}} dx - \frac{1}{2} \int_{\Omega} T w dx. \quad (2.5)$$

It is easy to see that I is a C^1 functional and, thanks to (2.4), that its derivative is given by

$$I'(w)(\varphi, \psi) = \int_{\Omega} |w_1|^{\frac{1}{p}-1} w_1 \varphi |x|^{-\frac{\alpha}{p}} dx + \int_{\Omega} |w_2|^{\frac{1}{q}-1} w_2 \psi |x|^{-\frac{\beta}{q}} dx - \int_{\Omega} (\varphi K w_2 + \psi K w_1) dx,$$

for every $(w_1, w_2), (\varphi, \psi) \in X$. In particular, (w_1, w_2) is a critical point of I if and only if (2.2) holds, and so $(u, v) := (|x|^{-\frac{\alpha}{p}} |w_1|^{\frac{1}{p}-1} w_1, |x|^{-\frac{\beta}{q}} |w_2|^{\frac{1}{q}-1} w_2)$ is a strong solution of the original system (1.1). In this case we have that

$$I(w_1, w_2) = E(u_1, u_2) = \frac{pq-1}{(p+1)(q+1)} \int_{\Omega} |u|^{p+1} |x|^{\alpha} dx > 0.$$

Working in this framework, we can rewrite the least energy nodal level as

$$c_{\text{nod}} = \inf\{I(w_1, w_2) : w_1^{\pm}, w_2^{\pm} \neq 0, I'(w_1, w_2) = 0\}.$$

In the following, we will adapt some of the ideas of [2, 15] into our situation. A first novelty in our arguments consists in introducing the following constants λ and μ (in view of evening the different powers of w_1 and w_2 in the functional I), as well as the introduction of the map θ , in view of obtaining Proposition 2.4 ahead. Let

$$\lambda := \frac{2p(q+1)}{p+q+2pq}, \quad \mu = \frac{2q(p+1)}{p+q+2pq},$$

so that

$$\gamma := \lambda \frac{p+1}{p} = \mu \frac{q+1}{q} = \frac{2(p+1)(q+1)}{p+q+2pq} \in (1, 2) \quad \text{and} \quad \lambda + \mu = 2.$$

We start by introducing the Nehari type set:

$$\mathcal{N}_{\text{nod}} = \{(w_1, w_2) \in X : w_1^{\pm} \neq 0, w_2^{\pm} \neq 0 \text{ and } I'(w)(\lambda w_1^+, \mu w_2^+) = I'(w)(\lambda w_1^-, \mu w_2^-) = 0\}$$

and the level

$$\tilde{c}_{\text{nod}} = \inf_{w \in \mathcal{N}_{\text{nod}}} I(w),$$

which, we will check later on that, coincides with c_{nod} and that I is positive on \mathcal{N}_{nod} ; cf. Theorem 2.6 and (2.18) respectively. The study of this problem will be done by means of a fiber-type map: given $(w_1, w_2) \in X$ such that $w_1^\pm \neq 0$ and $w_2^\pm \neq 0$, define $\theta = \theta_w : [0, \infty) \times [0, \infty) \rightarrow \mathbb{R}$ by

$$\theta(t, s) = I(t^\lambda w_1^+ - s^\lambda w_1^-, t^\mu w_2^+ - s^\mu w_2^-).$$

Observe that if $t, s > 0$, then

$$\nabla \theta(t, s) = (0, 0) \quad \text{if and only if} \quad (t^\lambda w_1^+ - s^\lambda w_1^-, t^\mu w_2^+ - s^\mu w_2^-) \in \mathcal{N}_{\text{nod}}. \quad (2.6)$$

In what follows it will be important to prove that θ_w admits a critical point. However, this seems not to hold for every $w \in X$. Indeed, for some w it turns out that θ_w has no critical points and its supremum is plus infinity. For that reason, we introduce the following auxiliary set

$$\begin{aligned} \mathcal{N}_0 &= \left\{ w \in X : \begin{array}{l} \lambda \int_{\Omega} w_1^+ K w_2 dx + \mu \int_{\Omega} w_1 K w_2^+ dx > 0 \\ \lambda \int_{\Omega} w_1^- K w_2 dx + \mu \int_{\Omega} w_1 K w_2^- dx < 0 \end{array} \right\} \\ &= \{ w \in X : \lambda C_1 + \mu C_2 < 2B^+, \mu C_1 + \lambda C_2 < 2B^- \}, \end{aligned}$$

where we have used the fact that $\lambda + \mu = 2$, and the notations

$$B^+ = \int_{\Omega} w_1^+ K w_2^+ dx > 0, \quad B^- = \int_{\Omega} w_1^- K w_2^- dx > 0, \quad (2.7)$$

and

$$C_1 = \int_{\Omega} w_1^+ K w_2^- dx > 0, \quad C_2 = \int_{\Omega} w_1^- K w_2^+ dx > 0.$$

We point out that, by the strong maximum principle, if $w \in \mathcal{N}_0$ then $w_1^\pm \neq 0$, $w_2^\pm \neq 0$.

Lemma 2.1. *The set \mathcal{N}_0 is nonempty and $\mathcal{N}_{\text{nod}} \subseteq \mathcal{N}_0$.*

Proof. *Claim 1* – $\mathcal{N}_0 \neq \emptyset$. Let φ_2 be the second eigenfunction of $-\Delta$ with zero Dirichlet boundary condition, and denote its eigenvalue by λ_2 . Then $(\varphi_2, \varphi_2) \in \mathcal{N}_0$, as

$$\lambda \int_{\Omega} \varphi_2^+ K \varphi_2 dx + \mu \int_{\Omega} \varphi_2 K \varphi_2^+ dx = \frac{2}{\lambda_2} \int_{\Omega} (\varphi_2^+)^2 > 0$$

and

$$\lambda \int_{\Omega} \varphi_2^- K \varphi_2 dx + \mu \int_{\Omega} \varphi_2 K \varphi_2^- dx = -\frac{2}{\lambda_2} \int_{\Omega} (\varphi_2^-)^2 < 0.$$

Claim 2 – $\mathcal{N}_{\text{nod}} \subseteq \mathcal{N}_0$. This is an immediate consequence of the equalities that define \mathcal{N}_{nod} . Indeed, since $w \neq 0$,

$$\begin{aligned} \lambda \int_{\Omega} w_1^+ K w_2 dx + \mu \int_{\Omega} w_1 K w_2^+ dx \\ = \lambda \int_{\Omega} |w_1^+|^{\frac{p+1}{p}} |x|^{-\frac{\alpha}{p}} dx + \mu \int_{\Omega} |w_2^+|^{\frac{q+1}{q}} |x|^{-\frac{\beta}{q}} dx > 0 \end{aligned}$$

and

$$\begin{aligned} \lambda \int_{\Omega} w_1^- K w_2 dx + \mu \int_{\Omega} w_1 K w_2^- dx \\ = -\lambda \int_{\Omega} |w_1^-|^{\frac{p+1}{p}} |x|^{-\frac{\alpha}{p}} dx + \mu \int_{\Omega} |w_2^-|^{\frac{q+1}{q}} |x|^{-\frac{\beta}{q}} dx < 0. \quad \square \end{aligned}$$

Let us now study in detail the map θ , following a standard procedure. We recall that $\theta = \theta_w$ for $w = (w_1, w_2) \in X$ such that $w_1^{\pm} \neq 0$ and $w_2^{\pm} \neq 0$. After a few computations, we can rewrite it as

$$\theta(t, s) = A^+ t^\gamma + A^- s^\gamma - B^+ t^2 - B^- s^2 + C_1 t^\lambda s^\mu + C_2 t^\mu s^\lambda,$$

with

$$A^\pm = \frac{p}{p+1} \int_{\Omega} |w_1^\pm|^{\frac{p+1}{p}} |x|^{-\frac{\alpha}{p}} dx + \frac{q}{q+1} \int_{\Omega} |w_2^\pm|^{\frac{q+1}{q}} |x|^{-\frac{\beta}{q}} dx > 0. \quad (2.8)$$

Our goal is to prove that θ_w , for $w \in \mathcal{N}_0$, admits a unique critical point, which is a global maximum, attained at a pair with positive components. We divide the proof of this fact in several lemmas.

Lemma 2.2. *Let $w \in \mathcal{N}_0$ and take $\theta = \theta_w$. Then θ has a global maximum at some (t_0, s_0) with $t_0, s_0 > 0$. Moreover, every local maximum must have positive components.*

Proof. Young's inequality yields

$$t^\lambda s^\mu \leq \frac{\lambda t^2}{2} + \frac{\mu s^2}{2} \quad \text{and} \quad t^\mu s^\lambda \leq \frac{\mu t^2}{2} + \frac{\lambda s^2}{2} \quad \forall t, s \geq 0,$$

and thus

$$\theta(t, s) \leq A^+ t^\gamma + A^- s^\gamma + \left(\frac{1}{2}(\lambda C_1 + \mu C_2) - B^+ \right) t^2 + \left(\frac{1}{2}(\mu C_1 + \lambda C_2) - B^- \right) s^2.$$

As $w \in \mathcal{N}_0$, the coefficients of the quadratic terms are negative, hence $\theta(t, s) \rightarrow -\infty$ as $|s| + |t| \rightarrow +\infty$, and θ admits a global maximum (t_0, s_0) with *nonnegative* components.

To conclude, let us prove that it cannot happen that either $t_0 = 0$ or $s_0 = 0$. This is an immediate consequence of the fact that

$$\theta(t, s) = A^+ t^\gamma - B^+ t^2 + s^\gamma (A^- - B^- s^{2-\gamma}) + C_1 t^\lambda s^\mu + C_2 t^\mu s^\lambda > \theta(t, 0)$$

for $s > 0$ sufficiently small. Analogously, $\theta(t, s) > \theta(0, s)$ for $t > 0$ sufficiently small. \square

Lemma 2.3. *If (t, s) is a critical point of θ with $t, s > 0$, then (t, s) is a non degenerate local maximum.*

Proof. If (t, s) is a critical point of θ , then

$$2B^+ = \gamma A^+ t^{\gamma-2} + \lambda C_1 t^{\lambda-2} s^\mu + \mu C_2 t^{\mu-2} s^\lambda,$$

and

$$2B^- = \gamma A^- s^{\gamma-2} + \mu C_1 t^\lambda s^{\mu-2} + \lambda C_2 t^\mu s^{\lambda-2}.$$

Thus

$$\begin{aligned} \theta_{tt}(t, s) &= \gamma(\gamma-1)A^+ t^{\gamma-2} - 2B^+ + \lambda(\lambda-1)C_1 t^{\lambda-2} s^\mu + \mu(\mu-1)C_2 t^\mu s^{\lambda-2} \\ &= \gamma(\gamma-2)A^+ t^{\gamma-2} + \lambda(\lambda-2)C_1 t^{\lambda-2} s^\mu + \mu(\mu-2)C_2 t^\mu s^{\lambda-2} < 0, \end{aligned}$$

$$\begin{aligned} \theta_{ss}(t, s) &= \gamma(\gamma-1)A^- s^{\gamma-2} - 2B^- + \mu(\mu-1)C_1 t^\lambda s^{\mu-2} + \lambda(\lambda-1)C_2 t^\mu s^{\lambda-2} \\ &= \gamma(\gamma-2)A^- s^{\gamma-2} + \mu(\mu-2)C_1 t^\lambda s^{\mu-2} + \lambda(\lambda-2)C_2 t^\mu s^{\lambda-2} < 0, \end{aligned}$$

and

$$\theta_{ts}(t, s) = \lambda\mu C_1 t^{\lambda-1} s^{\mu-1} + \lambda\mu C_2 t^{\mu-1} s^{\lambda-1}.$$

The proof is complete as soon as we prove that $\theta_{ts}^2 < \theta_{tt}\theta_{ss}$, which is equivalent to

$$\begin{aligned} &\lambda^2 \mu^2 C_1^2 t^{2(\lambda-1)} s^{2(\mu-1)} + \lambda^2 \mu^2 C_2^2 t^{2(\mu-1)} s^{2(\lambda-1)} + 2\lambda^2 \mu^2 C_1 C_2 \\ &< \lambda(\lambda-2)\mu(\mu-2)C_1^2 t^{2(\lambda-1)} s^{2(\mu-1)} + \lambda(\lambda-2)\mu(\mu-2)C_2^2 t^{2(\mu-1)} s^{2(\lambda-1)} \\ &\quad + [\lambda^2(\lambda-2)^2 + \mu^2(\mu-2)^2]C_1 C_2 + \gamma^2(\gamma-2)^2 A^+ A^- t^{\gamma-2} s^{\gamma-2} \\ &\quad + \gamma(\gamma-2)\mu(\mu-2)A^+ C_1 t^{\lambda+\gamma-2} s^{\mu-2} + \gamma(\gamma-2)\lambda(\lambda-2)A^- C_1 t^{\lambda-2} s^{\mu+\gamma-2} \\ &\quad + \gamma(\gamma-2)\lambda(\lambda-2)A^+ C_2 t^{\mu+\gamma-2} s^{\lambda-2} + \mu(\mu-2)\gamma(\gamma-2)C_2 A^- s^{\lambda+\gamma-2} t^\mu \end{aligned} \tag{2.9}$$

Now $\lambda + \mu = 2$ is equivalent to $\lambda^2 \mu^2 = \lambda(\lambda-2)\mu(\mu-2)$, which in turn implies that

$$2\lambda^2 \mu^2 \leq \lambda^2(\lambda-2)^2 + \mu^2(\mu-2)^2.$$

Since the last five terms in (2.9) are positive, combining all of these we prove that the desired inequality holds true. \square

Proposition 2.4. *Let $w \in \mathcal{N}_0$. The map θ_w admits a unique critical point (t_0, s_0) with $t_0, s_0 > 0$, which corresponds to its unique global maximum. Moreover, the pair (t_0, s_0) can be characterized as being the (unique) solution to the system:*

$$\begin{cases} 2B^+ = \gamma A^+ t^{\gamma-2} + \lambda C_1 (s/t)^\mu + \mu C_2 (s/t)^\lambda, \\ 2B^- = \gamma A^- s^{\gamma-2} + \mu C_1 (t/s)^\lambda + \lambda C_2 (t/s)^\mu. \end{cases} \tag{2.10}$$

Proof. The only think left to prove is the uniqueness statement. One could argue as in [40, Proposition 3.2], but here instead we present a shorter argument, which makes use of the Poincaré-Hopf Theorem [28]. Recall that this result states that given M a smooth manifold with boundary, and X a vector field having only isolated zeros x_i ($i \in I$) and such that it points outward on ∂M , then

$$\chi(M) = \sum_{i \in I} \text{index}(X, x_i),$$

where $\chi(\cdot)$ is the Euler characteristic. Let (t_0, s_0) be a global maximum as in Lemma 2.2. Take M to be a bounded regular set containing (t_0, s_0) , which coincides with the square $[\varepsilon, L] \times [\varepsilon, L]$ except at the corners, where it is smooth. Then

$\chi(M) = 1$, and $X = -\nabla\theta$ points outward on ∂M for sufficiently small ε and L large enough. Lemma 2.3 on the other hand implies that $\text{index}(-\nabla\theta, (s, t)) = 1$ at each critical point (s, t) , and thus we prove that (t_0, s_0) is indeed the unique critical point of θ_w . \square

Lemma 2.5. *Let $w \in \mathcal{N}_{\text{nod}}$ be such that $I(w) = \tilde{c}_{\text{nod}}$. Then $I'(w) = 0$.*

Proof. We will argue as in [23]; see also [5, 16].

Step 1 – Assume by contradiction that $I'(w) \neq 0$. Then there exists $v \in X$ such that $I'(w)v = -2$. By continuity, there exists a small $\varepsilon > 0$ such that

$$I'((t^\lambda w_1^+ - s^\lambda w_1^-, t^\mu w_2^+ - s^\mu w_2^-) + rv)v < -1 \quad \forall 0 \leq r \leq \varepsilon, |t-1|, |s-1| \leq \varepsilon.$$

Set $D := [1-\varepsilon, 1+\varepsilon] \times [1-\varepsilon, 1+\varepsilon]$. We fix a smooth function $\eta : D \rightarrow [0, \varepsilon]$ such that $\eta(1, 1) = \varepsilon$ and $\eta = 0$ on ∂D , and denote

$$h(t, s) = (h_1(t, s), h_2(t, s)) := (t^\lambda w_1^+ - s^\lambda w_1^-, t^\mu w_2^+ - s^\mu w_2^-) + \eta(t, s)v,$$

$$H(t, s) = \left(I'(h(t, s))(\lambda h_1^+(t, s), \mu h_2^+(t, s)), I'(h(t, s))(\lambda h_1^-(t, s), \mu h_2^-(t, s)) \right).$$

By possibly taking a smaller ε , we can insure by continuity that

$$h(t, s) \in \mathcal{N}_0 \quad \forall (t, s) \in D.$$

Step 2 – We claim that there exists $(t_0, s_0) \in D$ such that $H(t_0, s_0) = (0, 0)$. To prove this, we use the classical Miranda's Theorem [29]. We will need to compute H on ∂D , where as $\eta = 0$,

$$H(t, s) = \begin{pmatrix} I'(t^\lambda w_1^+ - s^\lambda w_1^-, t^\mu w_2^+ - s^\mu w_2^-)(\lambda t^\lambda w_1^+, \mu t^\mu w_2^+) \\ I'(t^\lambda w_1^+ - s^\lambda w_1^-, t^\mu w_2^+ - s^\mu w_2^-)(\lambda s^\lambda w_1^-, \mu s^\mu w_2^-) \end{pmatrix} = \begin{pmatrix} t \theta_t(t, s) \\ -s \theta_s(t, s) \end{pmatrix}.$$

We have $\nabla\theta(1, 1) = (0, 0)$, which tells us that

$$\gamma A^+ - 2B^+ + \lambda C_1 + \mu C_2 = 0, \quad \gamma A^- - 2B^- + \mu C_1 + \lambda C_2 = 0.$$

For $s \in [1-\varepsilon, 1+\varepsilon]$ we have that, if $t = 1 + \varepsilon$, then

$$\begin{aligned} \theta_t &= \gamma A^+(1+\varepsilon)^{\gamma-1} - 2B^+(1+\varepsilon) + \lambda C_1(1+\varepsilon)^{\lambda-1} s^\mu + \mu C_2(1+\varepsilon)^{\mu-1} s^\lambda \\ &\leq \gamma A^+(1+\varepsilon)[(1+\varepsilon)^{\gamma-2} - 1] < 0; \end{aligned}$$

while if $t = 1 - \varepsilon$

$$\theta_t \geq \gamma A^+(1-\varepsilon)[(1-\varepsilon)^{\gamma-2} - 1] > 0.$$

Analogously, for $t \in [1-\varepsilon, 1+\varepsilon]$, we have

$$\theta_s < 0 \quad \text{for } s = 1 + \varepsilon, \quad \theta_s > 0 \quad \text{for } s = 1 - \varepsilon$$

and the claim follows.

Conclusion – By the previous point, which shows that $h(t_0, s_0) \in \mathcal{N}_{\text{nod}}$, it follows that

$$\begin{aligned} \tilde{c}_{\text{nod}} &\leq I(h(t_0, s_0)) \\ &= I(t_0^\lambda w_1^+ - s_0^\lambda w_1^-, t_0^\mu w_2^+ - s_0^\mu w_2^-) \\ &\quad + \int_0^{\eta(t_0, s_0)} I'((t_0^\lambda w_1^+ - s_0^\lambda w_1^-, t_0^\mu w_2^+ - s_0^\mu w_2^-) + rv) v dr \\ &\leq \theta_w(t_0, s_0) - \eta(t_0, s_0) \leq \theta_w(1, 1) - \eta(t_0, s_0) = \tilde{c}_{\text{nod}} - \eta(t_0, s_0), \end{aligned}$$

and so $\eta(t_0, s_0) = 0$ and, in particular, $\theta_w(t_0, s_0) = \theta_w(1, 1)$. By the uniqueness of maximum provided by Proposition 2.4 we must have $(t_0, s_0) = (1, 1)$ while, by construction, $\eta(1, 1) = \varepsilon > 0$, a contradiction. \square

Our purpose in the remainder of this section is to prove the following result.

Theorem 2.6. *The number \tilde{c}_{nod} is attained by a function $w \in \mathcal{N}_{\text{nod}}$. Moreover,*

$$c_{\text{nod}} = \tilde{c}_{\text{nod}} = \inf_{w \in \mathcal{N}_0} \sup_{t, s > 0} I(t^\lambda w_1^+ - s^\lambda w_1^-, t^\mu w_2^+ - s^\mu w_2^-) > 0. \quad (2.11)$$

Observe that this implies our first main result, namely Theorem 1.1. From (2.6), Lemma 2.1 and Proposition 2.4, we have

$$\tilde{c}_{\text{nod}} = \inf_{w \in \mathcal{N}_0} \sup_{t, s > 0} I(t^\lambda w_1^+ - s^\lambda w_1^-, t^\mu w_2^+ - s^\mu w_2^-). \quad (2.12)$$

So the only thing left to prove is that \tilde{c}_{nod} is achieved. To this aim, we rely on an indirect argument. The difficulty of a direct approach is that the weak limit of a minimizing sequence of \tilde{c}_{nod} does not belong necessarily to \mathcal{N}_0 , and we cannot project it in \mathcal{N}_{nod} . Indeed, even if we can bound from below the norms of the positive and negative parts of a minimizing sequence, the weak convergence in X does not imply the weak convergence of the positive and negative part of the sequence. Observe also that the lack of regularity of the nodal Nehari set makes rather tricky the use of Ekeland's principle to build a Palais-Smale sequence from a minimizing sequence.

To overpass this difficulty, we regularize the problem by introducing the auxiliary functional $I_\varepsilon : \tilde{X} \rightarrow \mathbb{R}$ defined by

$$\begin{aligned} I_\varepsilon(w_1, w_2) &= \frac{p}{p+1} \left(\varepsilon \int_\Omega |\nabla w_1|^{\frac{p+1}{p}} dx + \int_\Omega |w_1|^{\frac{p+1}{p}} |x|^{-\frac{\alpha}{p}} dx \right) \\ &\quad + \frac{q}{q+1} \left(\varepsilon \int_\Omega |\nabla w_2|^{\frac{q+1}{q}} dx + \int_\Omega |w_2|^{\frac{q+1}{q}} |x|^{-\frac{\beta}{q}} dx \right) - \frac{1}{2} \int_\Omega Tw dx, \end{aligned}$$

where $\varepsilon > 0$ and $\tilde{X} := (W^{1, \frac{p+1}{p}}(\Omega) \times W^{1, \frac{q+1}{q}}(\Omega)) \cap X$.

Such a regularization is a standard approach to regularize non uniformly elliptic operators, such as the curvature operator, see for instance [41, 27, 7]. It is surprisingly useful in our context to regularize a zero order term. Actually, the main utility of this approach is that it somehow provides a regularized minimizing

sequence which solves an approximating system. The key point is then that the regularization does not affect the geometry of the original functional, while the presence of the gradient terms give rise to Euler-Lagrange equations in which we can pass to the limit. For that, we will exploit the fact that \tilde{X} is dense in X .

Lemma 2.7. *Let $\tilde{X} = (W^{1, \frac{p+1}{p}}(\Omega) \times W^{1, \frac{q+1}{q}}(\Omega)) \cap X$ be endowed with the norm*

$$\|(w_1, w_2)\| = \|\nabla w_1\|_{\frac{p+1}{p}} + \|\nabla w_2\|_{\frac{q+1}{q}} + \|w_1\|_{\frac{p+1}{p}, \frac{\alpha}{p}} + \|w_2\|_{\frac{q+1}{q}, \frac{\beta}{q}}.$$

Then \tilde{X} is a reflexive Banach space which is continuously embedded in $W^{1, \frac{p+1}{p}}(\Omega) \times W^{1, \frac{q+1}{q}}(\Omega)$. Moreover \tilde{X} is a dense subspace of X .

Proof. The first two statements are obvious. To prove the density of \tilde{X} in X , first observe that

$$T : L^{\frac{p+1}{p}}(\Omega) \times L^{\frac{q+1}{q}}(\Omega) \longrightarrow L^{\frac{p+1}{p}}(\Omega, |x|^{-\frac{\alpha}{p}}) \times L^{\frac{q+1}{q}}(\Omega, |x|^{-\frac{\beta}{q}}),$$

defined by $T(f, g) = (f|x|^{\frac{\alpha}{p+1}}, g|x|^{\frac{\beta}{q+1}})$, is an isometric isomorphism. For each $\delta > 0$, fix $\varphi_\delta \in C_c^\infty(\mathbb{R}^N)$ such that

$$0 \leq \varphi_\delta \leq 1, \quad \varphi_\delta(x) = 1 \text{ if } |x| \geq 2\delta \text{ and } \varphi_\delta(x) = 0 \text{ if } |x| \leq \delta.$$

Then observe that $A = \{(f\varphi_\delta, g\varphi_\delta); f, g \in C_c^\infty(\Omega), \delta > 0\}$ is dense in $L^{\frac{p+1}{p}}(\Omega) \times L^{\frac{q+1}{q}}(\Omega)$ and that $T(A) \subset \tilde{X} \subset X$. \square

One can check easily that the functional I_ε belongs to $C^1(\tilde{X})$, and

$$\begin{aligned} I'_\varepsilon(w)(\varphi, \psi) &= \varepsilon \int_\Omega |\nabla w_1|^{\frac{1}{p}-1} \nabla w_1 \cdot \nabla \varphi \, dx + \int_\Omega |w_1|^{\frac{1}{p}-1} w_1 \varphi |x|^{-\frac{\alpha}{p}} \, dx \\ &+ \varepsilon \int_\Omega |\nabla w_2|^{\frac{1}{q}-1} \nabla w_2 \cdot \nabla \psi \, dx + \int_\Omega |w_2|^{\frac{1}{q}-1} w_2 \psi |x|^{-\frac{\beta}{q}} \, dx - \int_\Omega (\varphi K w_2 + \psi K w_1) \, dx, \end{aligned}$$

for every $(w_1, w_2), (\varphi, \psi) \in \tilde{X}$. In particular, (w_1, w_2) is a critical point of I_ε if and only if

$$\begin{cases} -\varepsilon \operatorname{div}(|\nabla w_1|^{\frac{1}{p}-1} \nabla w_1) + |w_1|^{\frac{1}{p}-1} w_1 |x|^{-\frac{\alpha}{p}} = K w_2 \\ -\varepsilon \operatorname{div}(|\nabla w_2|^{\frac{1}{q}-1} \nabla w_2) + |w_2|^{\frac{1}{q}-1} w_2 |x|^{-\frac{\beta}{q}} = K w_1 \end{cases} \quad (2.13)$$

in \tilde{X}^* . Define now

$$\tilde{\mathcal{N}}_0 = \left\{ w \in \tilde{X} : \begin{aligned} &\lambda \int_\Omega w_1^+ K w_2 \, dx + \mu \int_\Omega w_1 K w_2^+ \, dx > 0 \\ &\lambda \int_\Omega w_1^- K w_2 \, dx + \mu \int_\Omega w_1 K w_2^- \, dx < 0 \end{aligned} \right\}$$

and, for each $\varepsilon > 0$,

$$\begin{aligned} \mathcal{N}_{\text{nod}}^\varepsilon &= \{(w_1, w_2) \in \tilde{X} : w_1^\pm \neq 0, w_2^\pm \neq 0 \text{ and} \\ &\quad I'_\varepsilon(w)(\lambda w_1^+, \mu w_2^+) = I'_\varepsilon(w)(\lambda w_1^-, \mu w_2^-) = 0\}. \end{aligned}$$

Observe that

$$\mathcal{N}_{\text{nod}}^\varepsilon \subset \tilde{\mathcal{N}}_0 \subset \mathcal{N}_0.$$

In the statement of the next lemma and its proof, we keep the definitions of A^\pm and B^\pm as given in (2.8) and (2.7).

Lemma 2.8. *Let $w \in \tilde{\mathcal{N}}_0$. Then the map $\mathbb{R}^+ \times \mathbb{R}^+ \rightarrow \mathbb{R}$ defined by*

$$(t, s) \mapsto I_\varepsilon(t^\lambda w_1^+ - s^\lambda w_1^-, t^\mu w_2^+ - s^\mu w_2^-)$$

admits a unique critical point (t_0, s_0) , which is a global maximum. Moreover, the pair (t_0, s_0) can be characterized univocally by the condition

$$(t^\lambda w_1^+ - s^\lambda w_1^-, t^\mu w_2^+ - s^\mu w_2^-) \in \mathcal{N}_{\text{nod}}^\varepsilon$$

or equivalently through the system

$$\begin{cases} 2B^+ = \gamma \left(A^+ + \varepsilon \frac{p}{p+1} \int_{\Omega} |\nabla w_1^+|^{\frac{p+1}{p}} + \varepsilon \frac{q}{q+1} \int_{\Omega} |\nabla w_2^+|^{\frac{q+1}{q}} \right) t^{\gamma-2} \\ \quad + \lambda C_1 (s/t)^\mu + \mu C_2 (s/t)^\lambda, \\ 2B^- = \gamma \left(A^- + \varepsilon \frac{p}{p+1} \int_{\Omega} |\nabla w_1^-|^{\frac{p+1}{p}} + \varepsilon \frac{q}{q+1} \int_{\Omega} |\nabla w_2^-|^{\frac{q+1}{q}} \right) s^{\gamma-2} \\ \quad + \mu C_1 (t/s)^\lambda + \lambda C_2 (t/s)^\mu. \end{cases} \quad (2.14)$$

Proof. Since the functional I_ε has exactly the same shape and geometry of I , it is enough to repeat the proofs of Lemmas 2.1-2.3 and of Proposition 2.4, replacing only \mathcal{N}_0 and \mathcal{N}_{nod} by $\tilde{\mathcal{N}}_0$ and $\mathcal{N}_{\text{nod}}^\varepsilon$ respectively, and A^\pm by

$$A^\pm + \varepsilon \frac{p}{p+1} \int_{\Omega} |\nabla w_1^\pm|^{\frac{p+1}{p}} + \varepsilon \frac{q}{q+1} \int_{\Omega} |\nabla w_2^\pm|^{\frac{q+1}{q}}. \quad \square$$

Define the levels

$$c_{\text{nod}}^\varepsilon = \inf \{ I_\varepsilon(w) : w \in \tilde{X}, w_1^\pm, w_2^\pm \neq 0, I'_\varepsilon(w) = 0 \}$$

and

$$\tilde{c}_{\text{nod}}^\varepsilon = \inf_{\mathcal{N}_{\text{nod}}^\varepsilon} I_\varepsilon.$$

Lemma 2.9. *Given $\varepsilon > 0$, let $w \in \mathcal{N}_{\text{nod}}^\varepsilon$ be such that $I_\varepsilon(w) = \tilde{c}_{\text{nod}}^\varepsilon$. Then $I'_\varepsilon(w) = 0$.*

Proof. The proof follows the lines of that of Lemma 2.5 with obvious changes as in the proof of Lemma 2.8. \square

Proposition 2.10. *Given $\varepsilon > 0$, the number $c_{\text{nod}}^\varepsilon$ is attained by a function $w_\varepsilon \in \mathcal{N}_{\text{nod}}^\varepsilon$. Moreover, we have*

$$c_{\text{nod}}^\varepsilon = \tilde{c}_{\text{nod}}^\varepsilon = \inf_{w \in \tilde{\mathcal{N}}_0} \sup_{t, s > 0} I(t^\lambda w_1^+ - s^\lambda w_1^-, t^\mu w_2^+ - s^\mu w_2^-) > 0. \quad (2.15)$$

Proof. Step 1 – $\mathcal{N}_{\text{nod}}^\varepsilon$ is not empty for every $\varepsilon > 0$. This clearly follows from Lemma 2.8.

Step 2 – boundedness and convergence of minimizing sequences. Let $(w_n)_n \subset \mathcal{N}_{\text{nod}}^\varepsilon$ be a minimizing sequence for $\tilde{c}_{\text{nod}}^\varepsilon$. Denote, by simplicity,

$$a_n = \int_{\Omega} |w_{1,n}|^{\frac{p+1}{p}} |x|^{-\frac{\alpha}{p}} dx, \quad a_n^\pm = \int_{\Omega} |w_{1,n}^\pm|^{\frac{p+1}{p}} |x|^{-\frac{\alpha}{p}} dx$$

and

$$b_n = \int_{\Omega} |w_{2,n}|^{\frac{q+1}{q}} |x|^{-\frac{\beta}{q}} dx, \quad b_n^\pm = \int_{\Omega} |w_{2,n}^\pm|^{\frac{q+1}{q}} |x|^{-\frac{\beta}{q}} dx.$$

One has

$$\begin{aligned} \lambda \left(a_n^+ + \varepsilon \int_{\Omega} |\nabla w_{1,n}^+|^{\frac{p+1}{p}} dx \right) + \mu \left(b_n^+ + \varepsilon \int_{\Omega} |\nabla w_{2,n}^+|^{\frac{q+1}{q}} dx \right) \\ = \lambda \int_{\Omega} w_{1,n}^+ K w_{2,n} dx + \mu \int_{\Omega} w_{2,n}^+ K w_{1,n} dx, \end{aligned} \quad (2.16)$$

$$\begin{aligned} \lambda \left(a_n^- + \varepsilon \int_{\Omega} |\nabla w_{1,n}^-|^{\frac{p+1}{p}} dx \right) + \mu \left(b_n^- + \varepsilon \int_{\Omega} |\nabla w_{2,n}^-|^{\frac{q+1}{q}} dx \right) \\ = -\lambda \int_{\Omega} w_{1,n}^- K w_{2,n} dx - \mu \int_{\Omega} w_{2,n}^- K w_{1,n} dx. \end{aligned} \quad (2.17)$$

By adding (2.16) and (2.17), we obtain

$$\lambda \left(a_n + \varepsilon \int_{\Omega} |\nabla w_{1,n}|^{\frac{p+1}{p}} dx \right) + \mu \left(b_n + \varepsilon \int_{\Omega} |\nabla w_{2,n}|^{\frac{q+1}{q}} dx \right) = 2 \int_{\Omega} w_{1,n} K w_{2,n} dx$$

and we deduce that

$$\begin{aligned} I_\varepsilon(w_n) &= \frac{(pq-1)p}{(p+1)(2pq+p+q)} \left(a_n + \varepsilon \int_{\Omega} |\nabla w_{1,n}|^{\frac{p+1}{p}} dx \right) \\ &\quad + \frac{(pq-1)q}{(q+1)(2pq+p+q)} \left(b_n + \varepsilon \int_{\Omega} |\nabla w_{2,n}|^{\frac{q+1}{q}} dx \right) > 0. \end{aligned} \quad (2.18)$$

Observe that this shows that I_ε is positive on $\mathcal{N}_{\text{nod}}^\varepsilon$. Therefore $(w_n)_n$ is bounded in \tilde{X} , and up to a subsequence, we have that $w_n \rightharpoonup w$ weakly in \tilde{X} , strongly in $L^{\frac{p+1}{p}}(\Omega) \times L^{\frac{q+1}{q}}(\Omega)$. In particular, $(w_{1,n}^\pm, w_{2,n}^\pm) \rightarrow (w_1^\pm, w_2^\pm)$ in $L^{\frac{p+1}{p}}(\Omega) \times L^{\frac{q+1}{q}}(\Omega)$.

Step 3 – $w \in \tilde{\mathcal{N}}_0$. We need to show that

$$\lambda \int_{\Omega} w_1^+ K w_2 dx + \mu \int_{\Omega} w_1 K w_2^+ dx > 0$$

and

$$\lambda \int_{\Omega} w_1^- K w_2 dx + \mu \int_{\Omega} w_1 K w_2^- dx < 0.$$

From Step 2 and the continuity of K , we infer that the right-hand side in (2.16) and (2.17) do converge. We now show that the left-hand side in (2.16) and (2.17) are bounded away from zero. Starting from (2.16), we get

$$\lambda a_n^+ + \mu b_n^+ \leq \lambda \int_{\Omega} w_{1,n}^+ K w_{2,n}^+ dx + \mu \int_{\Omega} w_{2,n}^+ K w_{1,n}^+ dx = 2 \int_{\Omega} w_{1,n}^+ K w_{2,n}^+,$$

which by (2.3) yields, for any $\delta > 0$,

$$\lambda a_n^+ + \mu b_n^+ \leq c \|w_{1,n}\|_{\frac{p+1}{p}, \frac{p}{p}} \|w_{2,n}\|_{\frac{q+1}{q}, \frac{p}{q}} \leq \delta a_n^+ + \frac{C}{\delta} (b_n^+)^{\frac{q(p+1)}{q+1}}.$$

Since $q(p+1)/(q+1) > 1$, we deduce that $b_n^+ \geq \bar{\delta} > 0$ for some $\bar{\delta} > 0$. The inequalities

$$b_n^- \geq \bar{\delta} > 0, \quad a_n^\pm \geq \bar{\delta} > 0$$

follow by arguing in a similar way.

Conclusion – By Lemma 2.8, we can take (t_0, s_0) such that

$$(t_0^\lambda w_1^+ - s_0^\lambda w_1^-, t_0^\mu w_2^+ - s_0^\mu w_2^-) \in \mathcal{N}_{\text{nod}}^\varepsilon.$$

By the uniqueness assertion in the same lemma and the weak lower semicontinuity of the norm, we infer that

$$\begin{aligned} \tilde{c}_{\text{nod}}^\varepsilon &\leq I_\varepsilon(t_0^\lambda w_1^+ - s_0^\lambda w_1^-, t_0^\mu w_2^+ - s_0^\mu w_2^-) \\ &\leq \liminf I_\varepsilon(t_0^\lambda w_{1,n}^+ - s_0^\lambda w_{1,n}^-, t_0^\mu w_{2,n}^+ - s_0^\mu w_{2,n}^-) \leq \liminf I_\varepsilon(w_n) = \tilde{c}_{\text{nod}}^\varepsilon. \end{aligned}$$

Hence $(t_0^\lambda w_1^+ - s_0^\lambda w_1^-, t_0^\mu w_2^+ - s_0^\mu w_2^-) \in \mathcal{N}_{\text{nod}}^\varepsilon$ achieves $\tilde{c}_{\text{nod}}^\varepsilon$.

At last, the characterization (2.15) of the critical level can be proved in a straightforward way. \square

Our strategy to prove Theorem 2.6 now essentially consists in passing to the limit in (2.13) when $\varepsilon \rightarrow 0$. As a first step, we prove the convergence of the critical level, namely $\tilde{c}_{\text{nod}}^\varepsilon \rightarrow \tilde{c}_{\text{nod}}$ as $\varepsilon \rightarrow 0$. We start with two preliminary lemmas.

Lemma 2.11. *Take $(w_1, w_2) \in \tilde{\mathcal{N}}_0 \subset \mathcal{N}_0$ and let (t_0, s_0) be the unique pair such that*

$$t_0, s_0 > 0, \quad (t_0^\lambda w_1^+ - s_0^\lambda w_1^-, t_0^\mu w_2^+ - s_0^\mu w_2^-) \in \mathcal{N}_{\text{nod}}$$

while, for each $\varepsilon > 0$, let $(t_\varepsilon, s_\varepsilon)$ be the unique pair such that

$$t_\varepsilon, s_\varepsilon > 0, \quad (t_\varepsilon^\lambda w_1^+ - s_\varepsilon^\lambda w_1^-, t_\varepsilon^\mu w_2^+ - s_\varepsilon^\mu w_2^-) \in \mathcal{N}_{\text{nod}}^\varepsilon.$$

Then

$$(t_\varepsilon, s_\varepsilon) \rightarrow (t_0, s_0) \quad \text{as } \varepsilon \rightarrow 0.$$

Proof. The pair $(t_\varepsilon, s_\varepsilon)$ solves (2.14), so that

$$2B^+ t_\varepsilon^{2-\gamma} \geq \gamma A^+ \quad \text{and} \quad 2B^- s_\varepsilon^{\gamma-2} \geq \gamma A^-.$$

Since $2-\gamma > 0$, we infer that $t_\varepsilon, s_\varepsilon \geq a > 0$ for some constant a independent of ε .

For the sake of contradiction, assume that $\{(t_\varepsilon, s_\varepsilon)\}$ is unbounded as $\varepsilon \rightarrow 0$.

Case 1 – there exists $b > 0$ such that $a \leq s_\varepsilon \leq b$ or $a \leq t_\varepsilon \leq b$ for $\varepsilon \in (0, 1]$. In the first alternative, taking the limit in the first equation of (2.14), we obtain $2B^+ = 0$ which is a contradiction. In the second alternative, taking the limit in the second equation of (2.14) leads to $2B^- = 0$ which is still a contradiction.

Case 2 – both $t_\varepsilon, s_\varepsilon \rightarrow +\infty$ as $\varepsilon \rightarrow 0$. We divide this case in two subcases.

Case 2.1 – $t_\varepsilon/s_\varepsilon \rightarrow +\infty$ or $t_\varepsilon/s_\varepsilon \rightarrow 0$. This case leads again to either $2B^+ = 0$ or $2B^- = 0$.

Case 2.2 – $t_\varepsilon/s_\varepsilon \rightarrow l \in \mathbb{R}^+$. Taking the limit in (2.14) gives

$$\begin{cases} 2B^+ = \lambda C_1(1/l)^\mu + \mu C_2(1/l)^\lambda \\ 2B^- = \mu C_1 l^\lambda + \lambda C_2 l^\mu. \end{cases}$$

If $l \leq 1$, then $2B^- \leq \mu C_1 + \lambda C_2$ whereas $2B^+ < \lambda C_1 + \mu C_2$ if $l > 1$. In both cases, we obtain an inequality which contradicts the fact that $(w_1, w_2) \in \tilde{\mathcal{N}}_0$.

We now conclude that, up to a subsequence, $t_\varepsilon \rightarrow \bar{t} > 0$, $s_\varepsilon \rightarrow \bar{s} > 0$, which satisfy (2.10). Hence, the uniqueness assertion in Proposition 2.4 implies $(\bar{t}, \bar{s}) = (t_0, s_0)$. \square

We just proved the continuity of the projection on $\mathcal{N}_{\text{nod}}^\varepsilon$ when $\varepsilon \rightarrow 0$. We will need also the continuity of the projection on \mathcal{N}_{nod} with respect to strong convergence in X .

Lemma 2.12. *Take $(w_1, w_2) \in \mathcal{N}_0$ and $(w_{1,n}, w_{2,n}) \in \tilde{\mathcal{N}}_0$ such that*

$$(w_{1,n}, w_{2,n}) \rightarrow (w_1, w_2) \quad \text{in } X, \text{ as } n \rightarrow \infty.$$

Let (t_n, s_n) and (t_0, s_0) be the unique pairs of positive components such that

$$(t_n^\lambda w_{1,n}^+ - s_n^\lambda w_{1,n}^-, t_n^\mu w_{2,n}^+ - s_n^\mu w_{2,n}^-), (t_0^\lambda w_1^+ - s_0^\lambda w_1^-, t_0^\mu w_2^+ - s_0^\mu w_2^-) \in \mathcal{N}_{\text{nod}}$$

Then

$$(t_n, s_n) \rightarrow (t_0, s_0) \quad \text{as } n \rightarrow \infty.$$

Proof. We have

$$\begin{cases} 2B_n^+ = \gamma A_n^+ t_n^{\gamma-2} + \lambda C_{1,n} (s_n/t_n)^\mu + \mu C_{2,n} (s_n/t_n)^\lambda, \\ 2B_n^- = \gamma A_n^- s_n^{\gamma-2} + \mu C_{1,n} (t_n/s_n)^\lambda + \lambda C_{2,n} (t_n/s_n)^\mu. \end{cases}$$

From the strong convergence in X and the continuity of K , we deduce that

$$\begin{aligned} A_n^\pm &:= \frac{p}{p+1} \int_\Omega |w_{1,n}^\pm|^{\frac{p+1}{p}} |x|^{-\frac{\alpha}{p}} dx + \frac{q}{q+1} \int_\Omega |w_{2,n}^\pm|^{\frac{q+1}{q}} |x|^{-\frac{\beta}{q}} dx \\ &\rightarrow \frac{p}{p+1} \int_\Omega |w_1^\pm|^{\frac{p+1}{p}} |x|^{-\frac{\alpha}{p}} dx + \frac{q}{q+1} \int_\Omega |w_2^\pm|^{\frac{q+1}{q}} |x|^{-\frac{\beta}{q}} dx =: A^\pm > 0, \end{aligned}$$

(recall that $w_i^\pm \not\equiv 0$, $i = 1, 2$ whenever $w \in \mathcal{N}_0$),

$$B_n^\pm := \int_\Omega w_{1,n}^\pm K w_{2,n}^\pm dx \rightarrow \int_\Omega w_{1,n}^\pm K w_{2,n}^\pm dx =: B^\pm > 0$$

and

$$\begin{aligned} C_{1,n} &:= \int_\Omega w_{1,n}^+ K w_{2,n}^- dx \rightarrow \int_\Omega w_1^+ K w_2^- dx =: C_1 > 0, \\ C_{2,n} &:= \int_\Omega w_{2,n}^+ K w_{1,n}^- dx \rightarrow \int_\Omega w_2^+ K w_1^- dx =: C_2 > 0. \end{aligned}$$

Using Proposition 2.4 and arguing exactly as in the proof of Lemma 2.11, we can infer that $(t_n, s_n)_n$ is a bounded sequence which actually converges to (t_0, s_0) . \square

We can now turn to the convergence of the critical level which implies that the extension of the map $\mathbb{R}^+ \ni \varepsilon \mapsto \tilde{c}_{\text{nod}}^\varepsilon$ by $\tilde{c}_{\text{nod}}^0 = \tilde{c}_{\text{nod}}$ is right-continuous at zero.

Proposition 2.13. *We have $\tilde{c}_{\text{nod}}^\varepsilon \rightarrow \tilde{c}_{\text{nod}}$, as $\varepsilon \rightarrow 0$.*

Proof. We deal successively with the upper and lower-semicontinuity.

Step 1 – Upper semi-continuity.

Fix $w \in \mathcal{N}_0$. Since \tilde{X} is dense in X and \mathcal{N}_0 is open, there exists $(w_n)_n \subset \tilde{\mathcal{N}}_0$ such that $w_n \rightarrow w$ strongly in X . Given $\varepsilon > 0$ and $n \in \mathbb{N}$, according to Lemma 2.8, there exist unique $t_{n,\varepsilon}, s_{n,\varepsilon} > 0$ such that

$$(t_{n,\varepsilon}^\lambda w_{1,n}^+ - s_{n,\varepsilon}^\lambda w_{1,n}^-, t_{n,\varepsilon}^\mu w_{2,n}^+ - s_{n,\varepsilon}^\mu w_{2,n}^-) \in \mathcal{N}_{\text{nod}}^\varepsilon.$$

Therefore, we have

$$\begin{aligned} \tilde{c}_{\text{nod}}^\varepsilon &\leq I_\varepsilon(t_{n,\varepsilon}^\lambda w_{1,n}^+ - s_{n,\varepsilon}^\lambda w_{1,n}^-, t_{n,\varepsilon}^\mu w_{2,n}^+ - s_{n,\varepsilon}^\mu w_{2,n}^-) \\ &= I(t_{n,\varepsilon}^\lambda w_{1,n}^+ - s_{n,\varepsilon}^\lambda w_{1,n}^-, t_{n,\varepsilon}^\mu w_{2,n}^+ - s_{n,\varepsilon}^\mu w_{2,n}^-) \\ &\quad + \varepsilon t_{n,\varepsilon}^\gamma \left(\frac{p}{p+1} \int_\Omega |\nabla w_{1,n}^+|^{\frac{p+1}{p}} + \frac{q}{q+1} \int_\Omega |\nabla w_{2,n}^+|^{\frac{q+1}{q}} \right) \\ &\quad + \varepsilon s_{n,\varepsilon}^\gamma \left(\frac{p}{p+1} \int_\Omega |\nabla w_{1,n}^-|^{\frac{p+1}{p}} + \frac{q}{q+1} \int_\Omega |\nabla w_{2,n}^-|^{\frac{q+1}{q}} \right). \end{aligned} \tag{2.20}$$

Now observe that by Lemma 2.11, for each fixed $n \in \mathbb{N}$, we have

$$(t_{n,\varepsilon}, s_{n,\varepsilon}) \rightarrow (t_{n,0}, s_{n,0}) \quad \text{as } \varepsilon \rightarrow 0,$$

where $t_{n,0}, s_{n,0} > 0$ is the unique pair of positive components such that

$$(t_{n,0}^\lambda w_{1,n}^+ - s_{n,0}^\lambda w_{1,n}^-, t_{n,0}^\mu w_{2,n}^+ - s_{n,0}^\mu w_{2,n}^-) \in \mathcal{N}_{\text{nod}}.$$

Taking the limit in (2.20) as $\varepsilon \rightarrow 0$, we obtain

$$\limsup_{\varepsilon \rightarrow 0} \tilde{c}_{\text{nod}}^\varepsilon \leq I(t_{n,0}^\lambda w_{1,n}^+ - s_{n,0}^\lambda w_{1,n}^-, t_{n,0}^\mu w_{2,n}^+ - s_{n,0}^\mu w_{2,n}^-).$$

On the other hand, by Lemma 2.12, we have

$$(t_{n,0}, s_{n,0}) \rightarrow (t_0, s_0) \quad \text{as } n \rightarrow \infty,$$

where $t_0, s_0 > 0$ is the unique pair of positive components such that

$$(t_0^\lambda w_1^+ - s_0^\lambda w_1^-, t_0^\mu w_2^+ - s_0^\mu w_2^-) \in \mathcal{N}_{\text{nod}}.$$

Hence we deduce that

$$\begin{aligned} \limsup_{\varepsilon \rightarrow 0} \tilde{c}_{\text{nod}}^\varepsilon &\leq I(t_0^\lambda w_1^+ - s_0^\lambda w_1^-, t_0^\mu w_2^+ - s_0^\mu w_2^-) \\ &= \sup_{t,s>0} I(t^\lambda w_1^+ - s^\lambda w_1^-, t^\mu w_2^+ - s^\mu w_2^-). \end{aligned}$$

Since this holds for every $w \in \mathcal{N}_0$, (2.12) implies

$$\limsup_{\varepsilon \rightarrow 0} \tilde{c}_{\text{nod}}^\varepsilon \leq \inf_{w \in \mathcal{N}_0} \sup_{t,s>0} I(t^\lambda w_1^+ - s^\lambda w_1^-, t^\mu w_2^+ - s^\mu w_2^-) = \tilde{c}_{\text{nod}}.$$

Step 2 – Lower semi-continuity.

Take $w_\varepsilon \in \mathcal{N}_{\text{nod}}^\varepsilon$ such that $I_\varepsilon(w_\varepsilon) = \tilde{c}_{\text{nod}}^\varepsilon$ and $I'_\varepsilon(w_\varepsilon) = 0$. Since $\mathcal{N}_{\text{nod}}^\varepsilon \subset \tilde{\mathcal{N}}_0 \subset \mathcal{N}_0$, there exists unique $(t_\varepsilon, s_\varepsilon) \in \mathbb{R}^+ \times \mathbb{R}^+$ such that

$$(t_\varepsilon^\lambda w_{1,\varepsilon}^+ - s_\varepsilon^\lambda w_{1,\varepsilon}^-, t_\varepsilon^\mu w_{2,\varepsilon}^+ - s_\varepsilon^\mu w_{2,\varepsilon}^-) \in \mathcal{N}_{\text{nod}}$$

Therefore, we have

$$\begin{aligned} \tilde{c}_{\text{nod}} &\leq I(t_\varepsilon^\lambda w_{1,\varepsilon}^+ - s_\varepsilon^\lambda w_{1,\varepsilon}^-, t_\varepsilon^\mu w_{2,\varepsilon}^+ - s_\varepsilon^\mu w_{2,\varepsilon}^-) \\ &\leq I(t_\varepsilon^\lambda w_{1,\varepsilon}^+ - s_\varepsilon^\lambda w_{2,\varepsilon}^-, t_\varepsilon^\mu w_{2,\varepsilon}^+ - s_\varepsilon^\mu w_{2,\varepsilon}^-) \\ &\quad + \varepsilon t_\varepsilon^\gamma \left(\frac{p}{p+1} \int_\Omega |\nabla w_{1,\varepsilon}^+|^{\frac{p+1}{p}} + \frac{q}{q+1} \int_\Omega |\nabla w_{2,\varepsilon}^+|^{\frac{q+1}{q}} \right) \\ &\quad + \varepsilon s_\varepsilon^\gamma \left(\frac{p}{p+1} \int_\Omega |\nabla w_{1,\varepsilon}^-|^{\frac{p+1}{p}} + \frac{q}{q+1} \int_\Omega |\nabla w_{2,\varepsilon}^-|^{\frac{q+1}{q}} \right) \\ &= I_\varepsilon(t_\varepsilon^\lambda w_{1,\varepsilon}^+ - s_\varepsilon^\lambda w_{2,\varepsilon}^-, t_\varepsilon^\mu w_{2,\varepsilon}^+ - s_\varepsilon^\mu w_{2,\varepsilon}^-) \\ &\leq \sup_{t,s>0} I_\varepsilon(t^\lambda w_{1,\varepsilon}^+ - s^\lambda w_{2,\varepsilon}^-, t^\mu w_{2,\varepsilon}^+ - s^\mu w_{2,\varepsilon}^-) = I_\varepsilon(w_\varepsilon) = \tilde{c}_{\text{nod}}^\varepsilon. \quad \square \end{aligned}$$

Consider now the family of approximating minimizers

$$W_\varepsilon := \{(w_\varepsilon, \varepsilon) \in \mathcal{N}_{\text{nod}}^\varepsilon \times \mathbb{R}^+ \mid I_\varepsilon(w_{1,\varepsilon}, w_{2,\varepsilon}) = \tilde{c}_{\text{nod}}^\varepsilon\}.$$

Our subsequent step is to prove that, given a sequence $(\varepsilon_n)_n$ converging to zero and a sequence $(w_{\varepsilon_n})_n$ such that $(w_{\varepsilon_n}, \varepsilon_n) \in W_{\varepsilon_n}$, $(w_{\varepsilon_n})_n$ converges strongly in X and achieves \tilde{c}_{nod} .

Arguing exactly as in the proof of Proposition 2.10, we infer that

$$\begin{aligned} I_\varepsilon(w_\varepsilon) &= \frac{(pq-1)p}{(p+1)(2pq+p+q)} \left(\int_\Omega |w_{1,\varepsilon}|^{\frac{p+1}{p}} |x|^{-\frac{\alpha}{p}} dx + \varepsilon \int_\Omega |\nabla w_{1,\varepsilon}|^{\frac{p+1}{p}} \right) \\ &\quad + \frac{(pq-1)q}{(q+1)(2pq+p+q)} \left(\int_\Omega |w_{2,\varepsilon}|^{\frac{q+1}{q}} |x|^{-\frac{\beta}{q}} dx + \varepsilon \int_\Omega |\nabla w_{2,\varepsilon}|^{\frac{q+1}{q}} \right). \end{aligned}$$

Combining this identity with Proposition 2.13, we deduce the existence of $C > 0$ such that

$$\sup_{\varepsilon \in (0,1]} \left(\int_\Omega |w_{1,\varepsilon}|^{\frac{p+1}{p}} |x|^{-\frac{\alpha}{p}} dx, \int_\Omega |w_{2,\varepsilon}|^{\frac{q+1}{q}} |x|^{-\frac{\beta}{q}} dx \right) \leq C. \quad (2.22)$$

Moreover, arguing as in Step 2 of Proposition 2.10, we deduce that

$$\inf_{\varepsilon \in (0,1]} \left(\int_\Omega |w_{1,\varepsilon}^\pm|^{\frac{p+1}{p}} |x|^{-\frac{\alpha}{p}} dx, \int_\Omega |w_{2,\varepsilon}^\pm|^{\frac{q+1}{q}} |x|^{-\frac{\beta}{q}} dx \right) \geq \bar{\delta} > 0, \quad (2.23)$$

together with the lower estimates

$$\lambda \int_\Omega w_{1,\varepsilon}^+ K w_{2,\varepsilon} dx + \mu \int_\Omega w_{2,\varepsilon}^+ K w_{1,\varepsilon} dx \geq 2\bar{\delta}, \quad (2.24)$$

$$-\lambda \int_\Omega w_{1,\varepsilon}^- K w_{2,\varepsilon} dx - \mu \int_\Omega w_{2,\varepsilon}^- K w_{1,\varepsilon} dx \geq 2\bar{\delta}, \quad (2.25)$$

which hold for every $\varepsilon \in (0, 1]$.

Next we prove that the gradient terms disappear when taking the limit in $I'_\varepsilon(w_\varepsilon)$ as $\varepsilon \rightarrow 0$.

Proposition 2.14. *Let $(w_\varepsilon, \varepsilon) \in W_\varepsilon$. We have*

$$\max \left(\varepsilon \int_{\Omega} |\nabla w_{1,\varepsilon}|^{\frac{p+1}{p}}, \quad \varepsilon \int_{\Omega} |\nabla w_{2,\varepsilon}|^{\frac{q+1}{q}} \right) \rightarrow 0, \quad \text{as } \varepsilon \rightarrow 0.$$

Proof. From the inequalities in (2.21), we actually deduce that

$$\lim_{\varepsilon \rightarrow 0} I(t_\varepsilon^\lambda w_{1,\varepsilon}^+ - s_\varepsilon^\lambda w_{2,\varepsilon}^-, t_\varepsilon^\mu w_{2,\varepsilon}^+ - s_\varepsilon^\mu w_{2,\varepsilon}^-) = \tilde{c}_{\text{nod}},$$

$$\lim_{\varepsilon \rightarrow 0} \varepsilon t_\varepsilon^\gamma \left(\frac{p}{p+1} \int_{\Omega} |\nabla w_{1,\varepsilon}^+|^{\frac{p+1}{p}} + \frac{q}{q+1} \int_{\Omega} |\nabla w_{2,\varepsilon}^+|^{\frac{q+1}{q}} \right) = 0,$$

and

$$\lim_{\varepsilon \rightarrow 0} \varepsilon s_\varepsilon^\gamma \left(\frac{p}{p+1} \int_{\Omega} |\nabla w_{1,\varepsilon}^-|^{\frac{p+1}{p}} + \frac{q}{q+1} \int_{\Omega} |\nabla w_{2,\varepsilon}^-|^{\frac{q+1}{q}} \right) = 0.$$

The conclusion is an obvious consequence of the following claim.

Claim – $t_\varepsilon, s_\varepsilon \not\rightarrow 0$. We argue as in the proof of Lemmas 2.11 and 2.12, though we proceed with extra care since we have no information yet about the convergence of (w_ε) in X . The pair $(t_\varepsilon, s_\varepsilon)$ satisfies

$$2t_\varepsilon^{2-\gamma} B_\varepsilon^+ \geq \gamma A_\varepsilon^+ \quad 2s_\varepsilon^{2-\gamma} B_\varepsilon^- \geq \gamma A_\varepsilon^- \quad (2.26)$$

with

$$A_\varepsilon^\pm := \frac{p}{p+1} \int_{\Omega} |w_{1,\varepsilon}^\pm|^{\frac{p+1}{p}} |x|^{-\frac{\alpha}{p}} dx + \frac{q}{q+1} \int_{\Omega} |w_{2,\varepsilon}^\pm|^{\frac{q+1}{q}} |x|^{-\frac{\beta}{q}} dx$$

and

$$B_\varepsilon^\pm := \int_{\Omega} w_{1,\varepsilon}^\pm K w_{2,\varepsilon}^\pm dx.$$

From (2.22), we deduce the existence of $w, (f_+, g_+), (f_-, g_-) \in X$ such that, up to a subsequence, $w_\varepsilon \xrightarrow{X} w$,

$$w_{1,\varepsilon}^\pm \rightharpoonup f_\pm \geq 0, \quad \text{weakly in } L^{\frac{p+1}{p}}(\Omega, |x|^{-\frac{\alpha}{p}})$$

and

$$w_{2,\varepsilon}^\pm \rightharpoonup g_\pm \geq 0, \quad \text{weakly in } L^{\frac{q+1}{q}}(\Omega, |x|^{-\frac{\beta}{q}}).$$

Taking the limit as $\varepsilon \rightarrow 0$ in (2.24)–(2.25), we obtain

$$\begin{aligned} \lambda \int_{\Omega} f_+ K(g_+ - g_-) dx + \mu \int_{\Omega} g_+ K(f_+ - f_-) dx &\geq 2\bar{\delta}, \\ -\lambda \int_{\Omega} f_- K(g_+ - g_-) dx - \mu \int_{\Omega} g_- K(f_+ - f_-) dx &\geq 2\bar{\delta}, \end{aligned}$$

and it is clear that $f_\pm, g_\pm \not\equiv 0$. Going back to (2.26), we see that the lower estimate (2.23) yields $A_\varepsilon^\pm \geq \tilde{\delta} > 0$, where $\tilde{\delta}$ is independent of ε . Since moreover

$$B_\varepsilon^\pm \rightarrow \int_{\Omega} f_\pm K g_\pm dx > 0,$$

we have that $t_\varepsilon, s_\varepsilon \not\rightarrow 0$, as claimed. \square

We are now ready to conclude this section by proving its main result.

Proof of Theorem 2.6. The key point is the next claim.

Claim – up to a subsequence, $w_\varepsilon \rightarrow w$ strongly in X for some $w \in X$.

We first deduce from (2.22) the existence of $w \in X$, $g_1 \in L^{p+1}(\Omega, |x|^{-\frac{\alpha}{p}})$, $g_2 \in L^{q+1}(\Omega, |x|^{-\frac{\beta}{q}})$ such that $w_\varepsilon \rightharpoonup w$ weakly in X ,

$$|w_{1,\varepsilon}|^{\frac{1}{p}-1} w_{1,\varepsilon} \rightharpoonup g_1 \text{ weakly in } L^{p+1}(\Omega, |x|^{-\frac{\alpha}{p}})$$

and

$$|w_{2,\varepsilon}|^{\frac{1}{q}-1} w_{2,\varepsilon} \rightharpoonup g_2 \text{ weakly in } L^{q+1}(\Omega, |x|^{-\frac{\beta}{q}}).$$

In particular, taking $\varepsilon \rightarrow 0$ in the approximating system (2.13) and using Proposition 2.14, we conclude that

$$g_1 |x|^{-\frac{\alpha}{p}} = K w_2, \quad g_2 |x|^{-\frac{\beta}{q}} = K w_1. \quad (2.27)$$

This implies in particular that

$$g_1 |x|^{-\frac{\alpha}{p}} \in W^{2, \frac{q+1}{q}}(\Omega) \cap W_0^{1, \frac{q+1}{q}}(\Omega) \subset L^{p+1}(\Omega)$$

and

$$g_2 |x|^{-\frac{\beta}{q}} \in W^{2, \frac{p+1}{p}}(\Omega) \cap W_0^{1, \frac{p+1}{p}}(\Omega) \subset L^{q+1}(\Omega).$$

Writing $g_1 = |f_1|^{\frac{1}{p}-1} f_1$, with $f_1 = |g_1|^{p-1} g_1$, we observe that $f_1 \in L^{\frac{p+1}{p}}(\Omega, |x|^{-\frac{\alpha}{q}})$ because

$$\int_{\Omega} |f_1|^{\frac{p+1}{p}} |x|^{-\frac{\alpha}{p}} dx = \int_{\Omega} |g_1|^{p+1} |x|^{-\frac{\alpha}{p}} dx < \infty.$$

Now take $\delta > 0$ and fix $h_\delta \in W^{1, \frac{p+1}{p}}(\Omega) \cap L^{\frac{p+1}{p}}(\Omega, |x|^{-\frac{\alpha}{p}})$ such that

$$\|f_1 - h_\delta\|_{\frac{p+1}{p}, \frac{\alpha}{p}} < \delta. \quad (2.28)$$

Using the test function $w_{1,\varepsilon} - f_1$ in the equation

$$-\varepsilon \operatorname{div}(|\nabla w_{1,\varepsilon}|^{\frac{1}{p}-1} \nabla w_{1,\varepsilon}) + (|w_{1,\varepsilon}|^{\frac{1}{p}-1} w_{1,\varepsilon} - |f_1|^{\frac{1}{p}-1} f_1) |x|^{-\frac{\alpha}{p}} = K w_{2,\varepsilon} - K w_2,$$

we deduce that

$$\begin{aligned} & \int_{\Omega} (|w_{1,\varepsilon}|^{\frac{1}{p}-1} w_{1,\varepsilon} - |f_1|^{\frac{1}{p}-1} f_1) (w_{1,\varepsilon} - h_\delta) |x|^{-\frac{\alpha}{p}} dx \\ &= -\varepsilon \int_{\Omega} |\nabla w_{1,\varepsilon}|^{\frac{1}{p}-1} \nabla w_{1,\varepsilon} \cdot \nabla (w_{1,\varepsilon} - h_\delta) dx + \int_{\Omega} (K w_{2,\varepsilon} - K w_2) (w_{1,\varepsilon} - h_\delta) dx. \end{aligned}$$

Proposition 2.14, the upper bounds (2.22) and the compactness of the operator K now imply that the right hand side of the last equation converges to 0 as $\varepsilon \rightarrow 0$. Consequently, we can take $\bar{\varepsilon} > 0$ such that

$$\left| \int_{\Omega} (|w_{1,\varepsilon}|^{\frac{1}{p}-1} w_{1,\varepsilon} - |f_1|^{\frac{1}{p}-1} f_1) (w_{1,\varepsilon} - h_\delta) |x|^{-\frac{\alpha}{p}} dx \right| < \delta, \text{ for every } \varepsilon < \bar{\varepsilon}. \quad (2.29)$$

Combining (2.22), (2.28) and (2.29), we infer that

$$\begin{aligned}
& \left| \int_{\Omega} (|w_{1,\varepsilon}|^{\frac{1}{p}-1} w_{1,\varepsilon} - |f_1|^{\frac{1}{p}-1} f_1)(w_{1,\varepsilon} - f_1) |x|^{-\frac{\alpha}{p}} dx \right| \\
& \leq \left| \int_{\Omega} (|w_{1,\varepsilon}|^{\frac{1}{p}-1} w_{1,\varepsilon} - |f_1|^{\frac{1}{p}-1} f_1)(w_{1,\varepsilon} - h_{\delta}) |x|^{-\frac{\alpha}{p}} dx \right| \\
& \quad + \left| \int_{\Omega} (|w_{1,\varepsilon}|^{\frac{1}{p}-1} w_{1,\varepsilon} - |f_1|^{\frac{1}{p}-1} f_1)(h_{\delta} - f_1) |x|^{-\frac{\alpha}{p}} dx \right| \\
& \leq \delta + \|h_{\delta} - f_1\|_{\frac{p+1}{p}, \frac{\alpha}{p}} \| |w_{1,\varepsilon}|^{\frac{1}{p}-1} w_{1,\varepsilon} - |f_1|^{\frac{1}{p}-1} f_1 \|_{p+1, \frac{\alpha}{p}} \leq \delta + C\delta,
\end{aligned}$$

for some $C > 0$ (independent of ε and δ). We therefore conclude that

$$\lim_{\varepsilon \rightarrow 0} \int_{\Omega} (|w_{1,\varepsilon}|^{\frac{1}{p}-1} w_{1,\varepsilon} - |f_1|^{\frac{1}{p}-1} f_1)(w_{1,\varepsilon} - f_1) |x|^{-\frac{\alpha}{p}} dx = 0. \quad (2.30)$$

Using the classical pointwise estimate

$$(|\xi|^{\frac{1}{p}-1} \xi - |\eta|^{\frac{1}{p}-1} \eta) \cdot (\xi - \eta) \geq 2^{\frac{p-1}{p}} |\xi - \eta|^{\frac{p+1}{p}} \quad \text{if } 0 < p < 1,$$

see for instance [36], we easily deduce from (2.30) that $w_{1,\varepsilon} \rightarrow f_1 = w_1$ strongly in $L^{\frac{p+1}{p}}(\Omega, |x|^{-\frac{\alpha}{p}})$ when $p \leq 1$. In the complementary case $p > 1$, using the pointwise estimate

$$(|\xi|^{\frac{1}{p}-1} \xi - |\eta|^{\frac{1}{p}-1} \eta) \cdot (\xi - \eta) \geq \frac{1}{p} |\xi - \eta|^2 (|\xi| + |\eta|)^{\frac{1}{p}-1} \quad \text{if } p \geq 1,$$

see again [36], we observe that

$$\begin{aligned}
\int_{\Omega} |w_{1,\varepsilon} - f_1|^{\frac{p+1}{p}} |x|^{-\frac{\alpha}{p}} dx &= \int_{\Omega} \frac{|w_{1,\varepsilon} - f_1|^{\frac{p+1}{p}}}{(|w_{1,\varepsilon}| + |f_1|)^{\frac{p^2-1}{2p^2}}} (|w_{1,\varepsilon}| + |f_1|)^{\frac{p^2-1}{2p^2}} |x|^{-\frac{\alpha}{p}} dx \\
&\leq \left(\int_{\Omega} \frac{|w_{1,\varepsilon} - f_1|^2}{(|w_{1,\varepsilon}| + |f_1|)^{1-\frac{1}{p}}} |x|^{-\frac{\alpha}{p}} dx \right)^{\frac{p+1}{2p}} \left(\int_{\Omega} (|w_{1,\varepsilon}| + |f_1|)^{\frac{p+1}{p}} |x|^{-\frac{\alpha}{p}} dx \right)^{\frac{p-1}{2p}} \\
&\leq C \left(\int_{\Omega} (|w_{1,\varepsilon}|^{\frac{1}{p}-1} w_{1,\varepsilon} - |f_1|^{\frac{1}{p}-1} f_1)(w_{1,\varepsilon} - f_1) |x|^{-\frac{\alpha}{p}} dx \right)^{\frac{p+1}{2p}}
\end{aligned}$$

and we reach the same conclusion as for $p \leq 1$. Obviously, the convergence of the component $w_{2,\varepsilon}$ follows in an analogous way.

The previous claim guarantees that $w_{\varepsilon} \rightarrow w$ with $w_1^{\pm}, w_2^{\pm} \neq 0$ since (2.23) now implies

$$\min \left(\int_{\Omega} |w_1^{\pm}|^{\frac{p+1}{p}} |x|^{-\frac{\alpha}{p}} dx, \int_{\Omega} |w_2^{\pm}|^{\frac{q+1}{q}} |x|^{-\frac{\beta}{q}} dx \right) \geq \bar{\delta} > 0.$$

From Propositions 2.13 and 2.14 together with the strong convergence in X , we conclude that $I(w) = \tilde{c}_{\text{nod}}$. Moreover, the equations in (2.27) tells that w is actually a critical point of I so that we have indeed proved that \tilde{c}_{nod} is achieved by a critical point of the functional I .

At last, the characterization (2.11) of the critical level follows in a straightforward way as previously mentioned. \square

3. LEAST ENERGY NODAL SOLUTIONS ARE FOLIATED SCHWARZ SYMMETRIC

Let Ω be a bounded radial domain centred at the origin, namely a ball or an annulus. The purpose of this section is to prove Theorem 1.2 via polarization methods, in the spirit of [5, 40]. First, we introduce some definitions and recall some known results. Define the sets

$$\mathcal{H}_0 = \{H \subset \mathbb{R}^N : H \text{ is a closed half-space in } \mathbb{R}^N \text{ with } 0 \in \partial H\}$$

and, for $p \neq 0$,

$$\mathcal{H}_0(p) = \{H \in \mathcal{H}_0 : p \in \text{int}(H)\}.$$

For each $H \in \mathcal{H}_0$ we denote the reflection in \mathbb{R}^N with respect to the hyperplane ∂H by $\sigma_H : \mathbb{R}^N \rightarrow \mathbb{R}^N$, and define the polarization of a function $u : \Omega \rightarrow \mathbb{R}$ with respect to $H \in \mathcal{H}_0$ by

$$u_H(x) = \begin{cases} \max\{u(x), u(\sigma_H(x))\} & x \in H \cap \Omega, \\ \min\{u(x), u(\sigma_H(x))\} & x \in \Omega \setminus H. \end{cases}$$

As far as we know the link between polarization and foliated Schwarz symmetry appeared first in [38]; cf. [5, Theorem 2.6] for further results about the foliated Schwarz symmetry of least energy solutions of some second order elliptic equations with radial data. We recall from [12, Lemma 4.2], see also [44, Proposition 2.7], the following equivalent characterization of foliated Schwarz symmetry which involves polarization.

Proposition 3.1. *Let $u : \Omega \rightarrow \mathbb{R}$ be a continuous function and take $p \in \partial B_1(0)$. Then the following statements are equivalent:*

- i) u is foliated Schwarz symmetric with respect to p ;
- ii) $u_H(x) = u(x) \forall x \in \Omega \cap H$, whenever $H \in \mathcal{H}_0(p)$.

Moreover, the next lemma collects some known properties about polarization; cf. [44, Lemma 3.1] for the first property and [5, Lemma 2.1] for the second.

Lemma 3.2. *Let $u : \Omega \rightarrow \mathbb{R}$ be a measurable function and $H \in \mathcal{H}_0$.*

- i) *If $F : \Omega \times \mathbb{R} \rightarrow \mathbb{R}$ is a continuous function such that $F(x, t) = F(y, t)$ for every $x, y \in \Omega$ such that $|x| = |y|$ and $t \in \mathbb{R}$ and $\int_{\Omega} |F(x, u(x))| dx < +\infty$, then $\int_{\Omega} F(x, u_H) dx = \int_{\Omega} F(x, u) dx$.*
- ii) $(u_H)^+ = (u^+)_H$, $(u_H)^- = -(-u^-)_H$.

Observe that the second statement of the previous result implies that:

$$(au^+ - bu^-)_H = a(u_H)^+ - b(u_H)^-, \quad \forall a, b > 0. \quad (3.1)$$

Finally, before we head to the proof of Theorem 1.2, we recall the following key estimate from [9, Lemma 3.7].

Lemma 3.3. *Given $u \in L^{\frac{p+1}{p}}(\Omega)$, $v \in L^{\frac{q+1}{q}}(\Omega)$ and any $H \in \mathcal{H}_0$, we have that*

$$\int_{\Omega} uKv \leq \int_{\Omega} u_H K(v_H).$$

We are now ready to prove our second main result.

Proof of Theorem 1.2. Let (u, v) be a least energy nodal solution of (1.1) and take the corresponding pair $(w_1, w_2) \in X$. Fix any $r > 0$ such that $\partial B_r(0) \subset \Omega$ and take $p \in \partial B_1(0)$ such that $w_1(rp) = \max_{\partial B_r(0)} w_1$. Given $H \in \mathcal{H}_0(p)$, we aim at proving that $(w_1)_H(x) = w_1(x)$ and $(w_2)_H(x) = w_2(x)$ for $x \in \Omega \cap H$. As

$$\begin{aligned} 0 &< \lambda \int_{\Omega} w_1^+ K w_2 \, dx + \mu \int_{\Omega} w_1 K w_2^+ \, dx \\ &\leq \lambda \int_{\Omega} (w_1^+)_H K (w_2)_H \, dx + \mu \int_{\Omega} (w_1)_H K (w_2^+)_H \, dx \end{aligned}$$

and

$$\begin{aligned} 0 &< \lambda \int_{\Omega} (-w_1^-) K w_2 \, dx + \mu \int_{\Omega} w_1 K (-w_2^-) \, dx \\ &\leq \lambda \int_{\Omega} (-w_1^-)_H K (w_2)_H \, dx + \mu \int_{\Omega} (w_1)_H K (-w_2^-)_H \, dx \\ &= -\lambda \int_{\Omega} ((w_1)_H)^- K (w_2)_H \, dx - \mu \int_{\Omega} (w_1)_H K ((w_2)_H)^- \, dx \end{aligned}$$

then $((w_1)_H, (w_2)_H) \in \mathcal{N}_0$, and from Proposition 2.4 we know there exist $t_0, s_0 > 0$ such that $(t_0^\lambda (w_1)_H^+ - s_0^\lambda ((w_1)_H)^-, t_0^\mu (w_2)_H^+ - s_0^\mu ((w_2)_H)^-) \in \mathcal{N}_{\text{nod}}$. Thus, by putting together Lemma 3.3 with (3.1) and with the uniqueness of global maximum,

$$\begin{aligned} c_{\text{nod}} &\leq I(t_0^\lambda (w_1)_H^+ - s_0^\lambda ((w_1)_H)^-, t_0^\mu (w_2)_H^+ - s_0^\mu ((w_2)_H)^-) \\ &= I((t_0^\lambda w_1^+ - s_0^\lambda w_1^-)_H, (t_0^\mu w_2^+ - s_0^\mu w_2^-)_H) \\ &\leq I(t_0^\lambda w_1^+ - s_0^\lambda w_1^-, t_0^\mu w_2^+ - s_0^\mu w_2^-) = \theta_w(t_0, s_0) \\ &\leq \sup_{t, s > 0} \theta_w(t, s) = \theta_w(1, 1) = I(w) = c_{\text{nod}}. \end{aligned}$$

Thus $(t_0, s_0) = (1, 1)$, $((w_1)_H, (w_2)_H) \in \mathcal{N}_{\text{nod}}$ and $I((w_1)_H, (w_2)_H) = c_{\text{nod}}$. By Lemma 2.5, $I'((w_1)_H, (w_2)_H) = 0$. Going bak to (u, v) , we have that both this pair as well as (u_H, v_H) solve (1.1). Thus

$$-\Delta(u_H - u) = |x|^\beta (|v_H|^{q-1} v_H - |v|^{q-1} v), \quad -\Delta(v_H - v) = |x|^\alpha (|u_H|^{p-1} u_H - |u|^{p-1} u), \quad (3.2)$$

in $\Omega \cap H$, and $u_H - u = v_H - v = 0$ on $\partial\Omega \cup (\Omega \cap \partial H)$. As $v_H \geq v$ in $\Omega \cap H$, then $-\Delta(u_H - u) \geq 0$ and by the maximum principle we have that either $u_H \equiv u$ or $u_H > u$. Since $u_H(rp) = u(rp)$, then $u_H \equiv u$ in $\Omega \cap H$. Going back to (3.2), we have $-\Delta(v_H - v) = 0$, and thus also $v_H \equiv v$. \square

4. SYMMETRY BREAKING

We start by proving Theorem 1.5.

Proof of Theorem 1.5. Multiplying the first equation of (1.1) by u , the second of (1.1) by v and integrating both gives

$$\int_{\Omega} |\nabla u|^2 dx = \int_{\Omega} |v|^{q-1} v u |x|^{\beta} dx \leq \left(\int_{\Omega} |v|^{q+1} |x|^{\beta} dx \right)^{\frac{q}{q+1}} \left(\int_{\Omega} |u|^{q+1} |x|^{\beta} dx \right)^{\frac{1}{q+1}}$$

and

$$\int_{\Omega} |\nabla v|^2 dx = \int_{\Omega} |u|^{q-1} u v |x|^{\beta} dx \leq \left(\int_{\Omega} |u|^{q+1} |x|^{\beta} dx \right)^{\frac{q}{q+1}} \left(\int_{\Omega} |v|^{q+1} |x|^{\beta} dx \right)^{\frac{1}{q+1}}.$$

Multiplying the first equation of (1.1) by v , the second of (1.1) by u and integrating both gives

$$\int_{\Omega} \nabla u \cdot \nabla v dx = \int_{\Omega} |v|^{q+1} |x|^{\beta} dx = \int_{\Omega} |u|^{q+1} |x|^{\beta} dx.$$

Putting these estimates together, we infer that

$$\int_{\Omega} |\nabla u|^2 dx + \int_{\Omega} |\nabla v|^2 dx \leq 2 \int_{\Omega} \nabla u \cdot \nabla v dx,$$

which obviously implies $u = v$. \square

Remember that for the single equation

$$-\Delta u = |u|^{q-1} u, \quad \text{in } \Omega, \quad u = 0 \text{ on } \partial\Omega,$$

it is known, cf. [1, Theorem 1.3], that any least energy nodal solution is non radial when $\Omega \subset \mathbb{R}^N$, $N \geq 2$, is either a ball or an annulus centred at the origin. We will show that when (p, q) is close to some couple (q_0, q_0) , and (α, β) is close to $(0, 0)$, this property is also true.

Take q_0 satisfying

$$q_0 > 1 \quad \text{such that} \quad q_0 + 1 < 2N/(N-2) \text{ if } N \geq 3, \quad (4.1)$$

and δ_0 such that

$$q_0 - \delta_0 > 1 \quad \text{and} \quad q_0 + 1 + \delta_0 < 2N/(N-2) \text{ if } N \geq 3, \quad (4.2)$$

that is, such that the square $[q_0 - \delta, q_0 + \delta_0]^2$ is contained in the region of the points (p, q) such that (H) holds.

The proof of Theorem 1.4 consists in doing some asymptotic estimates of the least energy nodal solutions and levels as $p, q \rightarrow q_0$ and $\alpha, \beta \rightarrow 0$, combined with the known fact that, at the diagonal point (q_0, q_0) and $\alpha = \beta = 0$, least energy nodal solutions are non radial. Having this in mind, let us introduce some notations. Given (p, q) satisfying (H), $\alpha, \beta \geq 0$, we denote by $c_{\text{nod}}^{p,q,\alpha,\beta}$ the least energy nodal level of (1.1), and by $E_{p,q,\alpha,\beta}$ its associated energy (1.2). We will also use the variational framework introduced in Section 2, denoting by $I_{p,q,\alpha,\beta}$

the energy functional (2.5). Recall that $E_{p,q,\alpha,\beta}(u, v) = I_{p,q,\alpha,\beta}(w_1, w_2)$ at critical points, under the relation

$$(u, v) := (|x|^{-\frac{\alpha}{p}} |w_1|^{\frac{1}{p}-1} w_1, |x|^{-\frac{\beta}{q}} |w_2|^{\frac{1}{q}-1} w_2).$$

Finally, recall the characterizations (cf. Theorem 2.6):

$$c_{\text{nod}}^{p,q,\alpha,\beta} = \inf_{\mathcal{N}_{\text{nod}}^{p,q,\alpha,\beta}} I_{p,q,\alpha,\beta} = \inf_{w \in \mathcal{N}_0^{p,q,\alpha,\beta}} \sup_{t,s>0} I_{p,q,\alpha,\beta}(t^\lambda w_1^+ - s^\lambda w_1^-, t^\mu w_2^+ - s^\mu w_2^-),$$

where

$$\lambda = \lambda(p, q) := \frac{2p(q+1)}{p+q+2pq}, \quad \mu = \mu(p, q) := \frac{2q(p+1)}{p+q+2pq},$$

and

$$\mathcal{N}_{\text{nod}}^{p,q,\alpha,\beta} = \left\{ (w_1, w_2) \in X_{p,q,\alpha,\beta} : w_1^\pm \neq 0, w_2^\pm \neq 0 \text{ and } I'_{p,q,\alpha,\beta}(w)(\lambda w_1^+, \mu w_2^+) = I'_{p,q,\alpha,\beta}(w)(\lambda w_1^-, \mu w_2^-) = 0 \right\},$$

$$\mathcal{N}_0^{p,q,\alpha,\beta} := \left\{ w \in X_{p,q,\alpha,\beta} : \begin{array}{l} \lambda \int_\Omega w_1^+ K w_2 dx + \mu \int_\Omega w_1 K w_2^+ dx > 0 \\ \lambda \int_\Omega w_1^- K w_2 dx + \mu \int_\Omega w_1 K w_2^- dx < 0 \end{array} \right\},$$

with $X_{p,q,\alpha,\beta} = L^{\frac{p+1}{p}}(\Omega, |x|^{-\frac{\alpha}{p}}) \times L^{\frac{q+1}{q}}(\Omega, |x|^{-\frac{\beta}{q}})$.

For simplicity, when $p = q$ and $\alpha = \beta = 0$ we will use the notation c_{nod}^p for $c_{\text{nod}}^{p,p,0,0}$. First we prove a uniform lower bound for the positive and negative parts of elements of $\mathcal{N}_{\text{nod}}^{p,q,\alpha,\beta}$.

Lemma 4.1. *Given q_0 satisfying (4.1) there exists $\delta_0 > 0$ and $\varepsilon > 0$ such that*

$$\int_\Omega |w_1^\pm|^{\frac{p+1}{p}} |x|^{-\frac{\alpha}{p}} dx \geq \varepsilon \quad \text{and} \quad \int_\Omega |w_2^\pm|^{\frac{q+1}{q}} |x|^{-\frac{\beta}{q}} dx \geq \varepsilon$$

for every $(w_1, w_2) \in \mathcal{N}_{\text{nod}}^{p,q,\alpha,\beta}$ with $p, q \in [q_0 - \delta_0, q_0 + \delta_0]$ and $\alpha, \beta \in [0, \delta_0]$.

Proof. We use the estimates in the proof of Theorem 2.6 - step 3, this time keeping a better track of the constants. We split the proof in several steps.

1) There exists C_1 (independent of p and q) such that

$$\|u\|_{p+1} \leq C_1 \|u\|_{W^{2, \frac{q+1}{q}}} \quad \forall u \in W^{2, \frac{q+1}{q}}(\Omega), \quad p, q \in [q_0 - \delta_0, q_0 + \delta_0].$$

Since Ω has finite measure and $p \leq q_0 + \delta_0$, from Hölder's estimates we deduce that

$$\|u\|_{p+1} \leq |\Omega|^{\frac{q_0 + \delta_0 - p}{(p+1)(q_0 + \delta_0 + 1)}} \|u\|_{q_0 + \delta_0 + 1} \leq \kappa_1 |\Omega|^{\frac{q_0 + \delta_0 - p}{(p+1)(q_0 + \delta_0 + 1)}} \|u\|_{W^{2, \frac{q_0 + \delta_0 + 1}{q_0 + \delta_0}}},$$

where κ_1 is a constant associated to the embedding $W^{2, \frac{q_0 + \delta_0 + 1}{q_0 + \delta_0}} \hookrightarrow L^{q_0 + \delta_0 + 1}$; recall that δ_0 is such that (4.2) holds. Moreover, again by using Hölder estimates and

also that $(q_0 + \delta_0 + 1)/(q_0 + \delta_0) \leq (q + 1)/q$,

$$\begin{aligned} \|u\|_{W^{2, \frac{q_0 + \delta_0 + 1}{q_0 + \delta_0}}} &= \left(\sum_{|\alpha| \leq 2} \int_{\Omega} |D^\alpha u|^{\frac{q_0 + \delta_0 + 1}{q_0 + \delta_0}} dx \right)^{\frac{q_0 + \delta_0}{q_0 + \delta_0 + 1}} \\ &\leq \left(|\Omega|^{\frac{q_0 + \delta_0 - q}{(q+1)(q_0 + \delta_0)}} \sum_{|\alpha| \leq 2} \|D^\alpha u\|_{\frac{q+1}{q}}^{\frac{q_0 + \delta_0 + 1}{q_0 + \delta_0}} \right)^{\frac{q_0 + \delta_0}{q_0 + \delta_0 + 1}} \\ &\leq \left(\frac{N(N-1)}{2} + N + 1 \right)^{\frac{q_0 + \delta_0}{q_0 + \delta_0 + 1}} |\Omega|^{\frac{q_0 + \delta_0 - q}{(q+1)(q_0 + \delta_0 + 1)}} \|u\|_{W^{2, \frac{q+1}{q}}}, \end{aligned}$$

and thus $\|u\|_{p+1} \leq \kappa(p, q) \|u\|_{W^{2, \frac{q+1}{q}}}$, with

$$\kappa(p, q) = \kappa_1 \left(\frac{N(N-1)}{2} + N + 1 \right)^{\frac{q_0 + \delta_0}{q_0 + \delta_0 + 1}} |\Omega|^{\frac{q_0 + \delta_0 - q}{(q+1)(q_0 + \delta_0 + 1)}} |\Omega|^{\frac{q_0 + \delta_0 - p}{(p+1)(q_0 + \delta_0 + 1)}},$$

which is bounded from above by some C_1 , for every $p, q \in [q_0 - \delta_0, q_0 + \delta_0]$.

2) There exists C_2 such that, for all $u \in W^{2, \frac{q_0 + \delta_0 + 1}{q_0 + \delta_0}}(\Omega) \cap W_0^{1, \frac{q_0 + \delta_0 + 1}{q_0 + \delta_0}}(\Omega)$

$$\|Ku\|_{W^{2, \frac{q_0 + \delta_0 + 1}{q_0 + \delta_0}}} \leq C_2 \|u\|_{\frac{q_0 + \delta_0 + 1}{q_0 + \delta_0}},$$

cf. [26, Lemma 9.17].

3) As $(w_1, w_2) \in \mathcal{N}_{\text{nod}}^{p, q, \alpha, \beta}$, from steps 1) and 2) above there exists $C > 0$ independent of p, q, α, β such that

$$\begin{aligned} \lambda \int_{\Omega} |w_1^+|^{\frac{p+1}{p}} |x|^{-\frac{\alpha}{p}} dx + \mu \int_{\Omega} |w_2^+|^{\frac{q+1}{q}} |x|^{-\frac{\beta}{q}} dx &\leq 2 \int_{\Omega} w_1^+ K w_2^+ dx \\ &\leq 2 \|w_1^+\|_{\frac{p+1}{p}} \|K w_2^+\|_{p+1} \leq 2C_1 \|w_1^+\|_{\frac{p+1}{p}} \|K w_2^+\|_{W^{2, \frac{q_0 + \delta_0 + 1}{q_0 + \delta_0}}} \\ &\leq 2C_1 C_2 \|w_1^+\|_{\frac{p+1}{p}} \|w_2^+\|_{\frac{q_0 + \delta_0 + 1}{q_0 + \delta_0}} \leq \tilde{C} \|w_1^+\|_{\frac{p+1}{p}} \|w_2^+\|_{\frac{q+1}{q}} \\ &\leq C \|w_1^+\|_{\frac{p+1}{p}, \frac{\alpha}{p}} \|w_2^+\|_{\frac{q+1}{q}, \frac{\beta}{q}}, \end{aligned}$$

where we have used estimate (2.1) and the fact that $q \leq q_0 + \delta_0$. By using the Young's inequality

$$Cab \leq \frac{\lambda}{2} a^{\frac{p+1}{p}} + \frac{(2p)^p C^{p+1}}{\lambda^p (p+1)^{p+1}} b^{p+1} \quad \forall a, b \geq 0,$$

we have

$$\mu \int_{\Omega} |w_2^+|^{\frac{q+1}{q}} |x|^{-\frac{\beta}{q}} dx \leq \frac{(2p)^p C^{p+1}}{\lambda^p (p+1)^{p+1}} \left(\int_{\Omega} |w_2^+|^{\frac{q+1}{q}} |x|^{-\frac{\beta}{q}} dx \right)^{\frac{q(p+1)}{q+1}},$$

and thus

$$\int_{\Omega} |w_2^+|^{\frac{q+1}{q}} |x|^{-\frac{\beta}{q}} dx \geq \left(\frac{\mu \lambda^p (p+1)^{p+1}}{(2p)^p C^{p+1}} \right)^{\frac{q+1}{pq-1}} =: K(p, q).$$

As $K(q_0, q_0) > 0$, then from sufficiently small δ_0 we have $K(p, q) \geq \varepsilon > 0$ for every $p, q \in [q_0 - \delta_0, q_0 + \delta_0]$.

The lower bounds for the remaining integrals follow in an analogous way. \square

Lemma 4.2. *We have*

$$\limsup c_{\text{nod}}^{p,q,\alpha,\beta} \leq c_{\text{nod}}^{q_0} \quad \text{as } p, q \rightarrow q_0, \alpha, \beta \rightarrow 0.$$

In particular, there exists δ_0 and $\kappa > 0$ such that

$$0 < c_{\text{nod}}^{p,q,\alpha,\beta} \leq \kappa, \quad \forall p, q \in [q_0 - \delta_0, q_0 + \delta_0], \alpha, \beta \in [0, \delta_0].$$

Proof. Take $p_n, q_n \rightarrow q_0, \alpha_n, \beta_n \rightarrow 0$.

1) We adapt some ideas from [11, Lemma 3], where a different problem is considered. Let (w_1, w_2) be such that $w_1^\pm \neq 0, w_2^\pm \neq 0$,

$$I_{q_0}(w_1, w_2) = c_{\text{nod}}^{q_0}, \quad I'_{q_0}(w_1, w_2) = 0,$$

where $I_{q_0} = I_{q_0, q_0, 0, 0}$. Denote $\lambda_n := \lambda(p_n, q_n)$ and $\mu_n := \mu(p_n, q_n)$. Since $\lambda_n, \mu_n \rightarrow 1$ as $n \rightarrow \infty$ and

$$\int_{\Omega} w_1^+ K w_2 dx + \int_{\Omega} w_1 K w_2^+ dx > 0, \quad \int_{\Omega} w_1^- K w_2 dx + \int_{\Omega} w_1 K w_2^- dx < 0,$$

then $(w_1, w_2) \in \mathcal{N}_0^{p_n, q_n, \alpha_n, \beta_n}$ for large n , and

$$c_{\text{nod}}^{p_n, q_n, \alpha_n, \beta_n} \leq \sup_{t, s > 0} I_{p_n, q_n, \alpha_n, \beta_n}(t^{\lambda_n} w_1^+ - s^{\lambda_n} w_1^-, t^{\mu_n} w_2^+ - s^{\mu_n} w_2^-).$$

Assume that the supremum at the right hand side is achieved at $(t, s) = (t_n, s_n)$.

2) We claim that $t_n, s_n \rightarrow 1$.

2a) First observe that t_n, s_n are bounded. In fact, repeating the computations of Lemma 2.2, we have this time that

$$\begin{aligned} I_{p_n, q_n, \alpha_n, \beta_n}(t_n^{\lambda_n} w_1^+ - s_n^{\lambda_n} w_1^-, t_n^{\mu_n} w_2^+ - s_n^{\mu_n} w_2^-) &\leq A_n^+ t_n^{\gamma_n} + A_n^- s_n^{\gamma_n} + \\ &+ \left(\frac{1}{2}(\lambda_n C_1 + \mu_n C_2) - B^+ \right) t_n^2 + \left(\frac{1}{2}(\mu_n C_1 + \lambda_n C_2) - B^- \right) s_n^2, \end{aligned}$$

with

$$A_n^\pm = \frac{p_n}{p_n + 1} \int_{\Omega} |w_1^\pm|^{\frac{p_n+1}{p_n}} |x|^{-\frac{\alpha_n}{p_n}} dx + \frac{q_n}{q_n + 1} \int_{\Omega} |w_2^\pm|^{\frac{q_n+1}{q_n}} |x|^{-\frac{\beta_n}{q_n}} dx,$$

(positive and bounded in n), B^\pm and C_i are as in Lemma 2.2, and

$$\gamma_n = \lambda_n \frac{p_n + 1}{p_n} = \mu_n \frac{q_n + 1}{q_n} \rightarrow \frac{q_0 + 1}{q_0} \in (1, 2).$$

Since moreover

$$\begin{aligned}\frac{1}{2}(\lambda_n C_1 + \mu_n C_2) - B^+ &\rightarrow \frac{1}{2}(C_1 + C_2) - B^+ < 0, \\ \frac{1}{2}(\mu_n C_1 + \lambda_n C_2) - B^- &\rightarrow \frac{1}{2}(C_1 + C_2) - B^- < 0,\end{aligned}$$

then if $|s_n| + |t_n| \rightarrow \infty$ we would have

$$0 < c_{\text{nod}}^{p_n, q_n, \alpha_n, \beta_n} \leq I_{p_n, q_n, \alpha_n, \beta_n}(t_n^{\lambda_n} w_1^+ - s_n^{\lambda_n} w_1^-, t_n^{\mu_n} w_2^+ - s_n^{\mu_n} w_2^-) \rightarrow -\infty,$$

a contradiction.

2b) We have $t_n, s_n \not\rightarrow 0$. In fact,

$$(t_n^{\lambda_n} w_1^+ - s_n^{\lambda_n} w_1^-, t_n^{\mu_n} w_2^+ - s_n^{\mu_n} w_2^-) \in \mathcal{N}_{\text{nod}}^{p_n, q_n, \alpha_n, \beta_n},$$

hence by Lemma 4.1

$$\int_{\Omega} |t_n^{\lambda_n} w_1^+|^{\frac{p_n+1}{p_n}} |x|^{-\frac{\alpha_n}{p_n}} dx, \int_{\Omega} |s_n^{\lambda_n} w_1^-|^{\frac{p_n+1}{p_n}} |x|^{-\frac{\alpha_n}{p_n}} dx \geq \varepsilon > 0,$$

which proves the statement.

2c) The claim of 2) now follows. We have $t_n \rightarrow \bar{t} \neq 0$, $s_n \rightarrow \bar{s} \neq 0$. Since

$$\begin{aligned}I_{p_n, q_n, \alpha_n, \beta_n}(w_1, w_2) &\leq \sup_{t, s > 0} I_{p_n, q_n, \alpha_n, \beta_n}(t^{\lambda_n} w_1^+ - s^{\lambda_n} w_1^-, t^{\mu_n} w_2^+ - s^{\mu_n} w_2^-) \\ &= I_{p_n, q_n, \alpha_n, \beta_n}(t_n^{\lambda_n} w_1^+ - s_n^{\lambda_n} w_1^-, t_n^{\mu_n} w_2^+ - s_n^{\mu_n} w_2^-).\end{aligned}$$

by passing to the limit,

$$\sup_{t, s > 0} I_{q_0}(t w_1^+ - s w_1^-, t w_2^+ - s w_2^-) = I_{q_0}(w_1, w_2) \leq I_{q_0}(\bar{t} w_1^+ - \bar{s} w_1^-, \bar{t} w_2^+ - \bar{s} w_2^-).$$

By the uniqueness provided by Proposition 2.4, we have $\bar{t} = \bar{s} = 1$.

3) Finally, by making $n \rightarrow \infty$ in the inequality

$$c_{\text{nod}}^{p_n, q_n, \alpha_n, \beta_n} \leq I_{p_n, q_n, \alpha_n, \beta_n}(t_n^{\lambda_n} w_1^+ - s_n^{\lambda_n} w_1^-, t_n^{\mu_n} w_2^+ - s_n^{\mu_n} w_2^-),$$

we obtain

$$\limsup c_{\text{nod}}^{p_n, q_n, \alpha_n, \beta_n} \leq I_{q_0}(w_1, w_2) = c_{\text{nod}}^{q_0}.$$

□

As a consequence, we have the following a priori bound.

Lemma 4.3. *Given q_0 satisfying (4.1) there exists $\delta_0 > 0$ and $\kappa > 0$ such that*

$$\|(u, v)\|_{\infty} \leq \kappa$$

for every (u, v) least energy nodal solution of (1.1) with $p, q \in [q_0 - \delta_0, q_0 + \delta_0]$, $\alpha, \beta \in [0, \delta_0]$.

Proof. Having the uniform upper bound of the energy levels coming from the previous lemma, and since the nonlinearities in (1.1) satisfy

$$||x|^\alpha |s|^{p_n} s|, ||x|^\beta |s|^{q_n} s| \leq C(1 + |s|^{q_0+1+\delta}) \quad \forall s \in \mathbb{R},$$

for $p, q \in [q_0 - \delta_0, q_0 + \delta_0]$, $\delta \in [0, \delta_0]$, C independent of p, q, α, β , then one can reason exactly as in the proof of [33, Lemma 5.4] (see also [10, Theorem 5.18]) to obtain uniform L^∞ bounds. \square

Lemma 4.4. *Take q_0 satisfying (4.1). Then*

$$c_{\text{nod}}^{p,q,\alpha,\beta} \rightarrow c_{\text{nod}}^{q_0} \quad \text{as } p, q \rightarrow q_0, \alpha, \beta \rightarrow 0.$$

Moreover, the corresponding least energy nodal solutions converge: if $(u_{p,q,\alpha,\beta}, v_{p,q,\alpha,\beta})$ is a sign changing solution of (1.1) with

$$E_{p,q,\alpha,\beta}(u_{p,q,\alpha,\beta}, v_{p,q,\alpha,\beta}) = c_{\text{nod}}^{p,q,\alpha,\beta},$$

then

$$u_{p,q,\alpha,\beta} \rightarrow u, \quad v_{p,q,\alpha,\beta} \rightarrow v \quad \text{in } C^{1,\gamma}(\overline{\Omega}) \text{ for every } 0 < \gamma < 1,$$

where (u, v) solves (1.1) for $p = q = q_0$, $\alpha = \beta = 0$, and $E_{q_0}(u, v) = c_{\text{nod}}^{q_0}$.

Proof. Take $p_n, q_n \rightarrow q_0$, $\alpha_n, \beta_n \rightarrow 0$, and let (u_n, v_n) be the corresponding least energy nodal solution of (1.1) with $(p, q, \alpha, \beta) = (p_n, q_n, \alpha_n, \beta_n)$. Then $\|(u_n, v_n)\|_\infty \leq \kappa$ and, by elliptic estimates, the sequence (u_n, v_n) is uniformly bounded in $W^{2,s} \times W^{2,t}$ for every $s, t > 1$. Thus there exists u, v such that $u_n \rightarrow u, v_n \rightarrow v$ in $C^{1,\gamma}(\overline{\Omega})$, and (u, v) solves

$$-\Delta u = |v|^{q_0-1} v, \quad -\Delta v = |u|^{q_0-1} u \quad \text{in } \Omega, \quad u = v = 0 \text{ on } \partial\Omega. \quad (4.3)$$

Defining

$$w_{1n} = |x|^{\alpha_n} |u_n|^{p_n-1} u_n, \quad w_{2n} = |x|^{\beta_n} |v_n|^{q_n-1} v_n,$$

we deduce from Lemma 4.1 that

$$\begin{aligned} \int_{\Omega} |u_n^\pm|^{p_n+1} |x|^{\alpha_n} dx &= \int_{\Omega} |w_{1n}^\pm|^{\frac{p_n+1}{p_n}} |x|^{-\frac{\alpha_n}{p_n}} dx \geq \varepsilon, \\ \int_{\Omega} |v_n^\pm|^{q_n+1} |x|^{\beta_n} dx &= \int_{\Omega} |w_{2n}^\pm|^{\frac{q_n+1}{q_n}} |x|^{-\frac{\beta_n}{q_n}} dx \geq \varepsilon, \end{aligned}$$

for some $\varepsilon > 0$ independent of n . Thus (u, v) is a sign changing solution of (4.3), and

$$c_{\text{nod}}^{q_0} \leq E_{q_0}(u, v) = \lim_n E_{p_n, q_n, \alpha_n, \beta_n}(u_n, v_n) = \lim_n c_{\text{nod}}^{p_n, q_n, \alpha_n, \beta_n}.$$

Combining this information with Lemma 4.2 yields the desired result. \square

Proof of Theorem 1.4. This is now an easy consequence of Lemma 1.5 and Lemma 4.4. Arguing by contradiction, we would get a *radial* least energy nodal solution of

$$-\Delta u = |u|^{q-1} u, \quad \text{in } \Omega, \quad u = 0 \text{ on } \partial\Omega,$$

contradicting [1, Theorem 1.3]. \square

Remark 4.5. Reasoning as in this section, we can prove that the map $(p, q, \alpha, \beta) \mapsto c_{\text{nod}}^{p,q,\alpha,\beta}$ is continuous for (p, q) satisfying (H), $\alpha, \beta \geq 0$, and that the corresponding least energy nodal solutions converge.

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REFERENCES

- [1] Amandine Aftalion and Filomena Pacella. Qualitative properties of nodal solutions of semilinear elliptic equations in radially symmetric domains. *C. R. Math. Acad. Sci. Paris*, 339(5):339–344, 2004.
- [2] Claudianor O. Alves and Sérgio H. M. Soares. Singularly perturbed elliptic systems. *Nonlinear Anal.*, 64(1):109–129, 2006.
- [3] Thomas Bartsch and Tobias Weth. A note on additional properties of sign changing solutions to superlinear elliptic equations. *Topol. Methods Nonlinear Anal.*, 22(1), 1–14, 2003.
- [4] Thomas Bartsch and Tobias Weth. Three nodal solutions of singularly perturbed elliptic equations on domains without topology. *Ann. Inst. H. Poincaré Anal. Non Linéaire*, 22(3), 259–281, 2005
- [5] Thomas Bartsch, Tobias Weth, and Michel Willem. Partial symmetry of least energy nodal solutions to some variational problems. *J. Anal. Math.*, 96:1–18, 2005.
- [6] Thomas Bartsch and Michel Willem. Infinitely many radial solutions of a semilinear elliptic problem on \mathbf{R}^N . *Arch. Rational Mech. Anal.*, 124(3):261–276, 1993.
- [7] Denis Bonheure, Patrick Habets, Franco Obersnel, and Pierpaolo Omari. Classical and non-classical solutions of a prescribed curvature equation. *J. Differential Equations*, 243(2) :208–237, 2007.
- [8] Denis Bonheure, Ederson Moreira dos Santos, and Miguel Ramos. Ground state and non-ground state solutions of some strongly coupled elliptic systems. *Trans. Amer. Math. Soc.*, 364(1):447–491, 2012.
- [9] Denis Bonheure, Ederson Moreira dos Santos, and Miguel Ramos. Symmetry and symmetry breaking for ground state solutions of some strongly coupled elliptic systems. *J. Funct. Anal.*, 264(1):62–96, 2013.
- [10] Denis Bonheure, Ederson Moreira dos Santos, and Hugo Tavares. Hamiltonian elliptic systems: a guide to variational methods. *Port. Math.* 71(3-4):301–395, 2014.
- [11] Denis Bonheure and Jean Van Schaftingen. Bound state solutions for a class of nonlinear Schrödinger equations. *Rev. Mat. Iberoam.*, 24(1):297–351, 2008.

- [12] Friedemann Brock. Symmetry and monotonicity of solutions to some variational problems in cylinders and annuli. *Electron. J. Differential Equations*, pages No. 108, 20 pp. (electronic), 2003.
- [13] Jérôme Busca and Boyan Sirakov. Symmetry results for semilinear elliptic systems in the whole space. *J. Differential Equations*, 163(1):41–56, 2000.
- [14] Alfonso Castro, Jorge Cossio, and John M. Neuberger. A sign-changing solution for a superlinear Dirichlet problem. *Rocky Mountain J. Math.*, 27(4):1041–1053, 1997.
- [15] Ph. Clément and R. C. A. M. Van der Vorst. On a semilinear elliptic system. *Differential Integral Equations*, 8(6):1317–1329, 1995.
- [16] M. Conti, S. Terracini, and G. Verzini. Nehari’s problem and competing species systems. *Ann. Inst. H. Poincaré Anal. Non Linéaire*, 19(6):871–888, 2002.
- [17] Lucio Damascelli and Filomena Pacella. Symmetry results for cooperative elliptic systems via linearization. *SIAM J. Math. Anal.*, 45(3):1003–1026, 2013.
- [18] Lucio Damascelli, Francesca Gladiali, and Filomena Pacella. A symmetry result for semilinear cooperative elliptic systems. In *Recent trends in nonlinear partial differential equations. II. Stationary problems*, volume 595 of *Contemp. Math.*, pages 187–204. Amer. Math. Soc., Providence, RI, 2013.
- [19] Lucio Damascelli, Filomena Pacella, and Francesca Gladiali. Symmetry results for cooperative elliptic systems in unbounded domains. *Indiana Univ. Math. J.*, 63(3):615–649, 2014.
- [20] Djairo G. de Figueiredo. Monotonicity and symmetry of solutions of elliptic systems in general domains. *NoDEA Nonlinear Differential Equations Appl.*, 1(2):119–123, 1994.
- [21] Djairo G. de Figueiredo. Semilinear elliptic systems: existence, multiplicity, symmetry of solutions. In *Handbook of differential equations: stationary partial differential equations. Vol. V*, Handb. Differ. Equ., pages 1–48. Elsevier/North-Holland, Amsterdam, 2008.
- [22] Djairo G. De Figueiredo and Jianfu Yang. Decay, symmetry and existence of solutions of semilinear elliptic systems. *Nonlinear Anal.*, 33(3):211–234, 1998.
- [23] Sébastien de Valeriola and Michel Willem. Existence of nodal solutions for some nonlinear elliptic problems. In *Nonlinear elliptic partial differential equations*, volume 540 of *Contemp. Math.*, pages 231–240. Amer. Math. Soc., Providence, RI, 2011.
- [24] Alberto Farina. Symmetry of components, liouville-type theorems and classification results for some nonlinear elliptic systems. *arXiv:1307.6949*, 2013.
- [25] B. Gidas, Wei Ming Ni, and L. Nirenberg. Symmetry and related properties via the maximum principle. *Comm. Math. Phys.*, 68(3):209–243, 1979.
- [26] David Gilbarg and Neil S. Trudinger. *Elliptic partial differential equations of second order*. Classics in Mathematics. Springer-Verlag, Berlin, 2001. Reprint of the 1998 edition.
- [27] O. Ladyzhenskaya and N. Ural’tseva. Local estimates for the gradients of solutions to the simplest regularization of a class of nonuniformly elliptic equations. (Russian) *Zap. Nauchn. Sem. S.-Peterburg. Otdel. Mat. Inst. Steklov. (POMI)* 213 (1994), Kraev. Zadachi Mat. Fiz. Smezh. Voprosy Teor. Funktsii. 25, 75–92, 225; translation in *J. Math. Sci.* 84(1), 862–872, 1997.
- [28] John W. Milnor. *Topology from the differentiable viewpoint*. Princeton Landmarks in Mathematics. Princeton University Press, Princeton, NJ, 1997. Based on notes by David W. Weaver, Revised reprint of the 1965 original.
- [29] Carlo Miranda. Un’osservazione su un teorema di Brouwer. *Boll. Un. Mat. Ital. (2)*, 3:5–7, 1940.
- [30] Filomena Pacella and Tobias Weth. Symmetry of solutions to semilinear elliptic equations via Morse index. *Proc. Amer. Math. Soc.*, 135(6):1753–1762 (electronic), 2007.
- [31] Pavol Quittner and Philippe Souplet. Symmetry of components for semilinear elliptic systems. *SIAM J. Math. Anal.*, 44(4):2545–2559, 2012.
- [32] M. Ramos, H. Tavares, and W. Zou. A Bahri-Lions theorem revisited. *Adv. Math.*, 222(6):2173–2195, 2009.

- [33] Miguel Ramos and Hugo Tavares. Solutions with multiple spike patterns for an elliptic system. *Calc. Var. Partial Differential Equations*, 31(1):1–25, 2008.
- [34] Bernhard Ruf. Superlinear elliptic equations and systems. In *Handbook of differential equations: stationary partial differential equations. Vol. V*, Handb. Differ. Equ., pages 211–276. Elsevier/North-Holland, Amsterdam, 2008.
- [35] A. W. Shaker. On symmetry in elliptic systems. *Appl. Anal.*, 41(1-4):1–9, 1991.
- [36] Jacques Simon. Régularité de la solution d’une équation non linéaire dans R^N . In *Journées d’Analyse Non Linéaire (Proc. Conf., Besançon, 1977)*, volume 665 of *Lecture Notes in Math.*, pages 205–227. Springer, Berlin, 1978.
- [37] Boyan Sirakov. On the existence of solutions of Hamiltonian elliptic systems in \mathbf{R}^N . *Adv. Differential Equations*, 5(10-12):1445–1464, 2000.
- [38] Didier Smets and Michel Willem. Partial symmetry and asymptotic behavior for some elliptic variational problems. *Calc. Var. Partial Differential Equations*, 18(1):57–75, 2003.
- [39] Didier Smets, Michel Willem, and Jiabao Su. Non-radial ground states for the Hénon equation. *Commun. Contemp. Math.*, 4(3):467–480, 2002.
- [40] H. Tavares and T. Weth. Existence and symmetry results for competing variational systems. *NoDEA Nonlinear Differential Equations Appl.*, 20(3):715–740, 2013.
- [41] R. Temam. Solutions généralisées de certaines équations du type hypersurfaces minima. *Arch. Rational Mech. Anal.*, 44: 121–156, 1971.
- [42] William C. Troy. Symmetry properties in systems of semilinear elliptic equations. *J. Differential Equations*, 42(3):400–413, 1981.
- [43] Tobias Weth. Nodal solutions to superlinear biharmonic equations via decomposition in dual cones. *Topol. Methods Nonlinear Anal.*, 28(1):33–52, 2006.
- [44] Tobias Weth. Symmetry of solutions to variational problems for nonlinear elliptic equations via reflection methods. *Jahresber. Dtsch. Math.-Ver.*, 112(3):119–158, 2010.
- [45] Henghui Zou. Symmetry of ground states for a semilinear elliptic system. *Trans. Amer. Math. Soc.*, 352(3):1217–1245, 2000.

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