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Vincent Monchiet, Djimedo Kondo. Combined voids size and shape effects on the macroscopic criterion of ductile nanoporous materials. *International Journal of Plasticity*, 2013, 43, pp.20-41. <10.1016/j.ijplas.2012.10.007>. <hal-01165815>

HAL Id: hal-01165815

<https://hal.science/hal-01165815v1>

Submitted on 18 Mar 2021

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Combined voids size and shape effects on the macroscopic criterion of ductile nanoporous materials

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Abstract

In this paper, we investigate the interfacial stress effects on the macroscopic yield function of ductile porous media containing nanosized spheroidal cavities. The solid matrix is assumed rigid-ideal plastic and of von Mises type with associated flow rule. We then perform limit analysis of a spheroidal unit cell containing a confocal spheroidal (prolate or oblate) cavity and subjected to arbitrary mechanical loadings. Voids size effects are captured by considering at the interface between the matrix and the cavity a surface stress model which relates the jump of the traction vector to the interfacial residual stress and interfacial plastic strain. This accounts for a thin shell of material in which occurs a strong plastic strain accumulation. We then provide a closed-form two-field based estimate of the overall dissipation which contains additional terms related to the interfacial plasticity. By taking advantage of this result, we derive parametric equations of the macroscopic yield surface of the nanoporous plastic material. The obtained estimates are assessed through comparisons with numerical data. Finally, it is shown that the resulting macroscopic criterion of the nanoporous material exhibits specific features such as (i) a dependence of the yield stress on the size of the spheroidal nanovoids, (ii) asymmetry between the yield stress in uniaxial tension and compression, (iii) more pronounced size effects for oblate voids than for prolate ones.

Key words: Ductile Materials, Yield Criterion, Voids Size Effects, Prolate or Oblate Cavities, Interfacial Stress.

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1 Introduction

The elastic properties of solids are significantly affected by the presence of surfaces and interfaces (see [68,13,10,41] for experimental evidence and [48,14,11,59,66] for numerical results). Surface effects are attributed to the presence of few layers of atoms which experience a different local environment than atoms in the bulk and have a different equilibrium positions and energy. The classical three-dimensional linear elasticity theory generally neglects these effects since the considered objects and structural elements are, at less, micro-sized but never of the size of one or few nanometers. For such nanosized objects, the interfacial effects become predominant since the area of surface per unit of volume is very high. Surface effects in standard elasticity theories are modeled by means of the Gurtin and Murdoch stress interface model [36,37] which assumes a jump of the traction vector while the displacement field is considered continuous across the surface (see also Povstenko [58], Steigman and Ogden [62] for various extensions). The jump condition may comply with a generalized Young-Laplace equation which extends to solid-solid interface the well known equation that describes capillarity effects in fluid mechanics. Thus, the discontinuity of the traction vector consists of two parts, the first one is attributed to the presence of interfacial residual stresses, independent of the deformation, and the second one being related to elastic deformation of the interface. Note that various studies, based on *ab initio* calculations and molecular dynamics, have been recently proposed for the identification of interfacial elastic coefficients and the residual stress that occur at the free surface of a monocrystalline material [48,61,38,71]. Mention must be also done of studies evaluating the effects of surface free energy on the elastic behavior of nano-sized particles (see for instance [17]).

Gurtin type stress surface model has been used to investigate the inclusions size dependency of elastic properties of nanocomposites. For instance, [19,21,22] have proposed to generalize the fundamental framework of micromechanical procedure to take into account the surface/interface stress effect at the nanoscale. Sharma and Ganti [60] generalize the Eshelby inclusion problem by incorporating the interface stress. Other works have proposed to derive bounds for the size dependent effective properties of nanocomposites [42,6,7]. However, it can be noted that all these studies generally concern spherical or cylindrical nano-inhomogeneities since the solution of the corresponding Eshelby type problem is only known for this kind of inclusion shape. Mention has to be made of the study by Ou et al. [54,55] who, by using displacement harmonic potential functions methods, attempted to derive the elastic fields around a nanosized spheroidal inclusion or cavity with surface effect. However, their studies do not deliver closed form or semi analytic expressions. In fact, a general micromechanics theory of elastic materials containing nanoinclusions, based on the solution of the Eshelby ellipsoidal inhomogeneity problem with

interfacial stress, is still lacking and is not the subject of the present study.

Surface effects on the yield strength and plastic behavior of nano objects and nanostructured materials have also been observed (see for instance [46,69] for experimental evidence and [15,26,56,57,16,70,34] for numerical simulations). Few recent studies consider stress interfaces effects in the context of non-linear composites or nanoporous ductile materials [72–74,18,34,50]. In particular, by performing a limit analysis of a hollow sphere, Dormieux and Kondo [18] generalize the Gurson model [35] in order to predict void size effects. To this end, they make use of a plastic version of the Gurtin stress interface model (see for instance [50]) which relates the interfacial stress to the plastic deformation at the cavity surface. The resulting model shows a void size dependency of the macroscopic yield strength of nanoporous media: for nanosized spherical cavities, the strength domain appears to be significantly larger than that predicted by the Gurson model.

Alternatively, for porous media, non-local plasticity theories have been used in [64,45,67,40,43,44] for the plastic behavior of the solid matrix instead of the von Mises model. In [45,67], the authors extend the Rice-Tracey or the Gurson model by considering, for the solid matrix, the Taylor dislocation based theory of plasticity introduced by [28]. In [40,43,44], the size effect is captured by the Fleck & Hutchinson phenomenological strain gradient plasticity theory [24] (see also [25] for a study of voids size effect on the plastic behavior of porous ductile media). It is interesting to note that the result obtained by Dormieux et al. [18] are qualitatively comparable to that earlier works mentioned above, which also predict an increase of the elastic domain by decreasing the size of the cavity. Still in the context of non-local plasticity theories, several studies are devoted to the study of combined voids shape and voids size effects (see for instance [39,40,43,44]). In all these studies, the void size effect is interpreted as the dependence of the solid matrix plasticity with an internal length scale which is physically attributed to the generation and the storage of geometrically necessary dislocations (Nye [53], Cottrell [12], Ashby [2], Fleck et al. [23], Gao et al. [27]). A priori, there is no direct link between the "interfacial stress" and the "non local approaches". However, and along the lines of Fleck and Hutchinson [24], microstructural effects of plastic media occur when the strain gradient are sufficiently large compared to dislocation spacing. Still, according to these authors, *"the expanding void is surrounded by a shell of hardened material due to the presence of both strain and strain gradient"*. Note that, based on [50], the use of the plastic version of the stress interface model can be physically interpreted as the existence of a thin region surrounding the interface in which intense gradient of plastic deformation will occur. Interface stresses model can be seen as an alternative approach to account for plastic strain accumulation at the cavity surface.

Such approach has been developed by [18] who restrict their analysis to the

case of spherical cavities. In the present study, we provide an extension of their approach to the case of spheroidal (prolate or oblate) cavities. By doing so, we also extend the classical work of Gologanu et al. [31,32] by simultaneously accounting for combined voids shape and void size effects. The paper is organized as follows. In section 2, we present the limit analysis of the hollow spheroid in the presence of stress interface. In section 3, we recall the trial velocity field already considered in [31,32] from which is computed a compact expression of the interfacial plastic strain. A two-field estimate of the macroscopic dissipation, including new terms related to the stress interface, is derived in section 4 from which we determine the parametric equations of the macroscopic yield surface of the nanoporous medium (see section 5). Finally, in section 6, we provide various illustrations of the combined voids shape and size dependency of the macroscopic criterion. In particular, the results of our two-field estimate are compared to data obtained by numerical computation of the macroscopic dissipation.

2 The limit analysis accounting for interfacial stress

2.1 Description of the unit cell

Let us consider a spheroidal cavity of semi-axes a_1 and b_1 embedded in a confocal spheroid of semi-axes a_2 and b_2 (see figure 1). The axis of the spheroids are aligned with $0x_3$, where $(0, x_1, x_2, x_3)$ is a cartesian coordinate system of orthonormal basis $(\underline{e}_1, \underline{e}_2, \underline{e}_3)$.

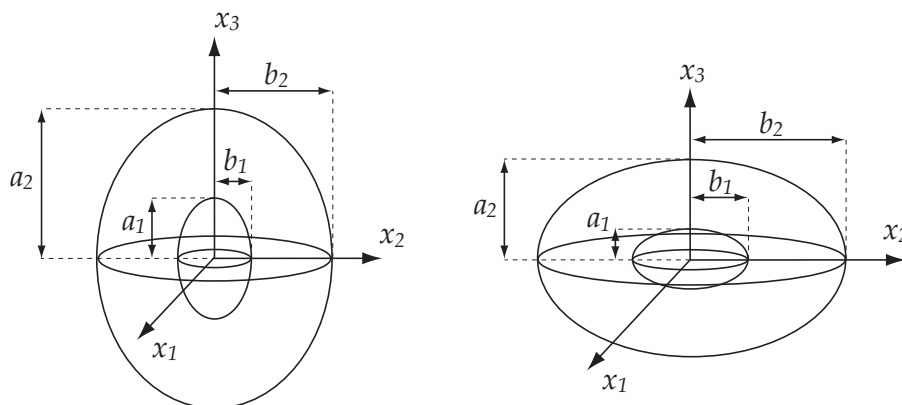


Fig. 1. The spheroidal unit cell with a confocal spheroidal void. Case of prolate cavity (at the left) and prolate cavity (at the right).

The assumption of confocality has been made in order to simplify analytical computation of the macroscopic criterion in the framework of limit analysis

methods. Note that it is the unique one that has been considered in literature dealing with Gurson type approach. Despite its simplicity, the confocal cell allows to represent microstructures of the type "hollow spheroid assemblages" (see for instance Benveniste and Milton [3]).

The volume of the cavity is $V_1 = 4\pi a_1 b_1^2/3$ while the total volume of the unit-cell is $V_2 = 4\pi a_2 b_2^2/3$. The shape of the cavity is described by the aspect ratio a_1/b_1 , with $a_1 > b_1$ corresponding to a prolate cavity while $b_1 > a_1$ to an oblate one. Let us denote by c the focal distance and by e_1 the void eccentricity, defined by:

$$\begin{aligned} c = \sqrt{a_1^2 - b_1^2} = \sqrt{a_2^2 - b_2^2} \quad e_1 = \frac{c}{a_1} \quad e_2 = \frac{c}{a_2} \quad (\text{prolate}) \\ c = \sqrt{b_1^2 - a_1^2} = \sqrt{b_2^2 - a_2^2} \quad e_1 = \frac{c}{b_1} \quad e_2 = \frac{c}{b_2} \quad (\text{oblate}) \end{aligned} \quad (1)$$

We will use both cylindrical coordinates ρ, θ, z and $(\underline{e}_\rho, \underline{e}_\theta, \underline{e}_z)$ the associated orthonormal basis and the classical spheroidal coordinates λ, φ, θ (associated orthogonal basis $(\underline{e}_\lambda, \underline{e}_\varphi, \underline{e}_\theta)$ defined by:

$$\begin{cases} x_1 = b \sin(\varphi) \cos(\theta) \\ x_2 = b \sin(\varphi) \sin(\theta) \\ x_3 = a \cos(\varphi) \end{cases} \quad \begin{cases} \rho = b \sin(\varphi) \\ \theta = \theta \\ z = x_3 = a \cos(\varphi) \end{cases} \quad (2)$$

and

$$\begin{cases} \underline{e}_\lambda = \frac{1}{L_\lambda} \{a \sin(\varphi) \underline{e}_\rho + b \cos(\varphi) \underline{e}_3\} \\ \underline{e}_\varphi = \frac{1}{L_\lambda} \{b \cos(\varphi) \underline{e}_\rho - a \sin(\varphi) \underline{e}_3\} \\ \underline{e}_\theta = \underline{e}_\theta \end{cases} \quad (3)$$

with $L_\lambda = \sqrt{a^2 \sin^2(\varphi) + b^2 \cos^2(\varphi)}$, $\lambda \in [0, +\infty[$, $\varphi \in [0, \pi]$ and $\underline{e}_\rho = \cos(\theta)\underline{e}_1 + \sin(\theta)\underline{e}_2$, $\theta \in [0, 2\pi]$. In the above equations: $a = c \cosh(\lambda)$ and $b = c \sinh(\lambda)$ for a prolate void, while for the case of an oblate void $a = c \sinh(\lambda)$ and $b = c \cosh(\lambda)$. The iso- λ surfaces define confocal spheroids with foci $c = \sqrt{|a^2 - b^2|}$ and eccentricity $e = c/a$, for a prolate void while $e = c/b$ for an oblate one. The porosity f is defined by: $f = (a_1 b_1^2)/(a_2 b_2^2)$. The matrix of the spheroidal unit cell is made of a rigid-plastic material obeying to the von Mises yield criterion, $\sigma_{eq} \leq \sigma_0$, (σ_0 being the yield stress in tension and σ_{eq} the von Mises

equivalent stress $\sigma_{eq} = \sqrt{\frac{3}{2}\bar{\boldsymbol{\sigma}} : \bar{\boldsymbol{\sigma}}}$ and the associated flow rule:

$$\bar{\boldsymbol{\sigma}} = \frac{2\sigma_0}{3} \frac{\mathbf{d}}{d_{eq}}, \quad \text{with} \quad d_{eq} = \sqrt{\frac{2}{3}\mathbf{d} : \mathbf{d}}, \quad \mathbf{d} = \frac{1}{2}(\nabla \otimes \underline{\mathbf{v}} + \underline{\mathbf{v}} \otimes \nabla) \quad (4)$$

where $\underline{\mathbf{v}}$ is the velocity field, \mathbf{d} the strain rate tensor and d_{eq} the von Mises equivalent strain rate. $\bar{\boldsymbol{\sigma}}$ represents the deviatoric part of the local stress tensor.

The surface between the void and the solid matrix, denoted Γ , is assumed to be described by a stress-type interface model, which was introduced by [36] in the context of elasticity and extended by [18,50] to plasticity. Stress interfaces ensure the continuity of the velocity field while the traction vector, $\underline{\mathbf{t}} = \boldsymbol{\sigma} \cdot \underline{\mathbf{n}}$, undergoes a jump which is governed by the Young-Laplace equation:

$$[\underline{\mathbf{t}}]_{\Gamma} = -\text{div}_s(\boldsymbol{\tau}) \quad (5)$$

In the above equation, $\text{div}_s(\boldsymbol{\tau})$ denotes the surface divergence of tensor $\boldsymbol{\tau}$ such that vector $\text{div}_s(\boldsymbol{\tau}) = (\boldsymbol{\tau} \otimes \underline{\mathbf{v}}) : \mathbf{P}(\underline{\mathbf{n}})$, with $\mathbf{P}(\underline{\mathbf{n}}) = \mathbf{I} - \underline{\mathbf{n}} \otimes \underline{\mathbf{n}}$, \mathbf{I} being the second order identity tensor and $\underline{\mathbf{n}}$ the normal unit vector taken on the interface Γ and oriented from the void to the solid matrix. In (5), the notation $[\underline{\mathbf{t}}]_{\Gamma}$ represents the jump of $\underline{\mathbf{t}}$ across Γ defined as follows: $[\underline{\mathbf{t}}]_{\Gamma} = \underline{\mathbf{t}}^+ - \underline{\mathbf{t}}^-$, where $\underline{\mathbf{t}}^+$ and $\underline{\mathbf{t}}^-$ are the values of the traction vector calculated on both sides of Γ . When the interface Γ is the boundary between a void and a solid, $\underline{\mathbf{t}}^- = 0$ and relation (5) reduces to $\underline{\mathbf{t}}^+ = -\text{div}_s(\boldsymbol{\tau})$.

In the Gurtin model [36], the interfacial stress $\boldsymbol{\tau}$ has two contributions, the first one is attributed to the presence of interfacial residual stresses, independent of the deformation, while the second one is related to elastic deformation of the interface. This model has been later extended in the context of plasticity by Monchiet and Bonnet [51] along a method formalized by Benveniste [4] in the context of linear elasticity. This method is based on the concept of equivalent elastic interphase which accounts for the presence of a few layers of atoms having a different local environment than atoms in the bulk. By considering the appropriate values for the elasticity coefficients of the interphase, it can be replaced by an idealized imperfect interface involving the jump of the traction vector which is described by the Gurtin model. Later, [51] have extended the work of [4] to the case of a plastic interphase. The interfacial model obtained in this work induces a jump of the traction vector which is still given by equation (5). However, in (5), the interfacial stress $\boldsymbol{\tau}$ is related to the plastic deformation of the interface. Its expression is:

$$\boldsymbol{\tau} = \tau_r \mathbf{P}(\underline{\mathbf{n}}) + \frac{2\tau_0}{3d_{eq}^s} (\mathbf{d}_s + \text{tr}(\mathbf{d}_s) \mathbf{P}(\underline{\mathbf{n}})) \quad (6)$$

where τ_r is the interfacial residual stress, τ_0 is a material parameter. Although limited, the particular form $\tau_r \mathbf{P}(\underline{n})$ (isotropic in the plane of the interface), considered for the interfacial residual stress is able to represent some cases of practical interest such as capillary forces. Note that recent ab initio based calculations by [71] indicate a slight anisotropy for these residual stresses. In (6), \mathbf{d}_s is the interfacial plastic strain rate which is defined as the restriction of the total strain rate to the tangent plane of normal unit vector \underline{n} . In (6), d_{eq}^s is the surface equivalent strain rate. The expressions of \mathbf{d}_s and d_{eq}^s are:

$$\begin{aligned} \mathbf{d}_s &= \mathbf{P}(\underline{n}) \cdot \mathbf{d} \cdot \mathbf{P}(\underline{n}) \\ d_{eq}^s &= \left[\frac{2}{3} \left(\mathbf{d}_s : \mathbf{d}_s + (\text{tr } \mathbf{d}_s)^2 \right) \right]^{1/2} \end{aligned} \quad (7)$$

Note that for a spheroidal surface $\underline{n} = \underline{e}_\lambda$, $\mathbf{P}(\underline{n}) = \underline{e}_\theta \otimes \underline{e}_\theta + \underline{e}_\varphi \otimes \underline{e}_\varphi$ and the quantities \mathbf{d}_s and d_{eq}^s read:

$$\begin{aligned} \mathbf{d}_s &= d_{\theta\theta} \underline{e}_\theta \otimes \underline{e}_\theta + d_{\varphi\varphi} \underline{e}_\varphi \otimes \underline{e}_\varphi + d_{\theta\varphi} (\underline{e}_\theta \otimes \underline{e}_\varphi + \underline{e}_\varphi \otimes \underline{e}_\theta) \\ d_{eq}^s &= \left[\frac{4}{3} (d_{\theta\theta}^2 + d_{\varphi\varphi}^2 + d_{\theta\varphi}^2 + d_{\theta\theta} d_{\varphi\varphi}) \right]^{1/2} \end{aligned} \quad (8)$$

2.2 Limit analysis taking into consideration interfacial stresses

As already mentioned, Hill-Mandel kinematic homogenization approach will be used to derive the overall plastic potential of the porous solid. Thus, uniform strain rate boundary conditions are considered on external surface of the hollow spheroid:

$$\underline{v}(\lambda = \lambda_2) = \mathbf{D} \cdot \underline{x} \quad (9)$$

where \mathbf{D} is the macroscopic strain rate tensor.

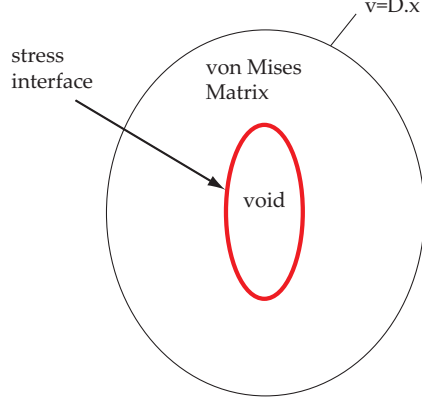


Fig. 2. Schematic showing the hollow spheroid with a stress interface.

The limit stress at the macroscopic scale, Σ , is:

$$\Sigma = \frac{\partial \Pi}{\partial \mathbf{D}}(\mathbf{D}) \quad (10)$$

where $\Pi(\mathbf{D})$, is the macroscopic dissipation given by:

$$\Pi(\mathbf{D}) = \inf_{\underline{v} \in \mathcal{K}} \frac{1}{V_2} \left[\int_{\Omega - \omega} \sigma_0 d_{eq} dV + \int_{\Gamma} [\underline{t}]_{\Gamma} \cdot \underline{v} dS \right] \quad (11)$$

where the infimum is taken on \mathcal{K} which is the space of admissible velocity fields, i.e. of continuous and differentiable velocity fields that comply with condition (9). The first integral in (11) is computed over the solid matrix, defined by $\Omega - \omega$, where Ω is the domain corresponding to the unit cell, while ω is the domain occupied by the void. The second integral in (11) is performed over the surface Γ of the void. It is recalled that $V_2 = 4\pi a_2 b_2^2/3$ is the volume of Ω . Using the generalized Young-Laplace equation (5), it is readily seen that the second integral is related to the interfacial residual stress and interfacial plastic strain. It then follows that:

$$\int_{\Gamma} [\underline{t}]_{\Gamma} \cdot \underline{v} dS = - \int_{\Gamma} \text{div}_s(\boldsymbol{\tau}) \cdot \underline{v} dS \quad (12)$$

Moreover, for any continuously differentiable two order tensor $\boldsymbol{\tau}$ and vector \underline{v} ,

$$\text{div}_s(\boldsymbol{\tau} \cdot \underline{v}) = \text{div}_s(\boldsymbol{\tau}) \cdot \underline{v} + \boldsymbol{\tau} : \mathbf{d}_s \quad (13)$$

The integral over Γ of the quantity $\text{div}_s(\boldsymbol{\tau} \cdot \underline{v})$ being null over any closed surface, it follows that:

$$\int_{\Gamma} \operatorname{div}_s(\boldsymbol{\tau}) \cdot \underline{v} dS = - \int_{\Gamma} \boldsymbol{\tau} : \mathbf{d}_s dS \quad (14)$$

Finally, replacing $\boldsymbol{\tau}$ by its expression (6), one obtains:

$$\Pi(\mathbf{D}) = \inf \left[\frac{1}{V_2} \int_{\Omega-\omega} \sigma_0 d_{eq} dV + \frac{1}{V_2} \int_{\Gamma} \tau_0 d_{eq}^s dS + \frac{1}{V_2} \int_{\Gamma} \tau_r \operatorname{tr}(\mathbf{d}^s) dS \right] \quad (15)$$

Note that this expression of $\Pi(\mathbf{D})$ contains three terms: the first is classic in the context of limit analysis of ductile porous media; the last two terms are associated to interfacial plastic dissipations.

Note that, in the spheroidal coordinate system, the surface integral is defined by:

$$\int_{\Gamma} \bullet dS = \int_{\varphi=0}^{\varphi=\pi} \int_{\theta=0}^{\theta=2\pi} \bullet b_1 L_{\lambda_1} \sin(\varphi) d\varphi d\theta \quad (16)$$

where L_{λ_1} is the value of L_{λ} for $\lambda = \lambda_1$, i.e. $L_{\lambda_1} = \sqrt{a_1^2 \sin^2(\varphi) + b_1^2 \cos^2(\varphi)}$; for the definition of the volume integral in the expression of $\Pi(\mathbf{D})$, the reader is referred to [31,32].

3 The choice of the trial velocity fields

A crucial step in the derivation of an approximate closed form expression of the macroscopic potential of the ductile porous material lies in the choice of the trial velocity field. Generally, this field is composed in two parts: one involving a constant traceless tensor \mathbf{A} , one, denoted \underline{v}^E , which is heterogeneous, i.e.

$$\underline{v} = \mathbf{A} \cdot \underline{x} + B \underline{v}^E \quad (17)$$

We propose here to adopt, for \underline{v}^E , the velocity field considered by Gologanu et al. [31,32] and Monchiet et al. [49]. This field has the property to comply with uniform strain rate conditions on any iso- $\lambda = cst$ spheroid. Gologanu et al. [33], Monchiet et al. [51] considered more sophisticated trial velocity fields in order to derive an improved estimate of the macroscopic yield locus of ductile porous media. In these approaches, the trial velocity field introduce a larger number of parameters which cannot be entirely determined by the boundary conditions. Due to additional difficulties related to the consideration of interfacial stress effects, we propose to restrict our analysis with the trial velocity fields first considered in [31,32]. Its expression is:

$$\begin{cases} v_\lambda^E = \frac{c^3}{bL_\lambda} \left[1 + \frac{1}{2}(1 - 3\alpha)(1 - 3\cos^2(\varphi)) \right] \\ v_\varphi^E = \frac{3c^3}{4ab^2L_\lambda} [(1 - \alpha)b^2 - 2\alpha a^2] \sin(2\varphi) \\ v_\theta^E = 0 \end{cases} \quad (18)$$

In which α depends on λ or equivalently e and is given by:

$$\alpha = \begin{cases} \frac{ab^2}{c^3} \operatorname{arctanh} \left\{ \frac{c}{a} \right\} - \frac{b^2}{c^2} = \frac{1 - e^2}{e^3} (\operatorname{arctanh}(e) - e) & \text{(prolate void)} \\ -\frac{ab^2}{c^3} \operatorname{arctan} \left\{ \frac{c}{a} \right\} + \frac{b^2}{c^2} = \frac{e - \arcsin(e)\sqrt{1 - e^2}}{e^3} & \text{(oblate void)} \end{cases} \quad (19)$$

For the explicit dependence of a and b on the coordinates λ and e , the reader is referred to section 2.1. The verification of uniform strain rate boundary

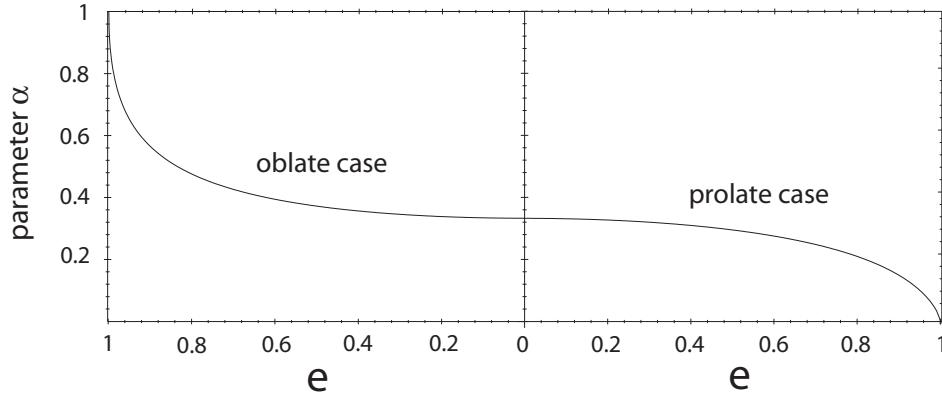


Fig. 3. α as function of e for a prolate and an oblate void.

conditions at the outer spheroid $\lambda = \lambda_2$ leads to:

$$\mathbf{A} = \mathbf{D} - D_m \mathbf{T}_2; \quad B = \frac{a_2 b_2^2}{c^3} D_m \quad (20)$$

where $\mathbf{T}_2 = \mathbf{T}(\lambda_2)$, with

$$\mathbf{T}(\lambda) = 3 \frac{a_2 b_2^2}{ab^2} \left[\frac{1}{2}(1 - \alpha) \mathbf{I}_t + \alpha \mathbf{I}_n \right] \quad (21)$$

and

$$\mathbf{I}_t = \mathbf{e}_1 \otimes \mathbf{e}_1 + \mathbf{e}_2 \otimes \mathbf{e}_2, \quad \mathbf{I}_n = \mathbf{e}_3 \otimes \mathbf{e}_3 \quad (22)$$

These tensors are the two projectors associated to the transverse isotropy related to the Ox_3 axis of the spheroidal cavity. They comply to: $\mathbf{I}_n \cdot \mathbf{I}_n = \mathbf{I}_n$, $\mathbf{I}_t \cdot \mathbf{I}_t = \mathbf{I}_t$, $\mathbf{I}_n \cdot \mathbf{I}_t = 0$ and $\mathbf{I} = \mathbf{I}_n + \mathbf{I}_t$ where is recalled that \mathbf{I} is the second order identity tensor. From equation (18), it is readily seen that the strain rate field in the matrix is the sum of a homogenous deviatoric field \mathbf{A} and a non-homogeneous field \mathbf{d}^E :

$$\mathbf{d} = \mathbf{A} + B\mathbf{d}^E \quad (23)$$

with \mathbf{d}^E given by:

$$\begin{cases} d_{\lambda\lambda}^E = -\frac{3c^3}{ab^2}(1-\alpha) + \frac{3c^3a}{2b^2L_\lambda^2}(1-3\alpha)\sin^2(\varphi) \\ d_{\varphi\varphi}^E = \frac{3c^3}{2ab^2}(1-\alpha) - \frac{3c^3a}{2b^2L_\lambda^2}(1-3\alpha)\sin^2(\varphi) \\ d_{\theta\theta}^E = \frac{3c^3}{2ab^2}(1-\alpha) \\ d_{\lambda\varphi}^E = \frac{3c^3}{4bL_\lambda^2}(1-3\alpha)\sin(2\varphi) \end{cases} \quad (24)$$

Note that the strain field \mathbf{d}^E can be put into the form the compact form:

$$\mathbf{d}^E = \frac{c^3}{a_2b_2^2} \left[\mathbf{T}(\lambda) - 3\frac{a_2b_2^2}{ab^2} \underline{e}_\lambda \otimes \underline{e}_\lambda \right] \quad (25)$$

In relation (23), \mathbf{A} and B are replaced by expressions (20) while relation (25) is used for \mathbf{d}^E . The total strain rate, \mathbf{d} , reads then:

$$\mathbf{d} = \mathbf{D} + D_m \left[\mathbf{T}(\lambda) - \mathbf{T}_2 - 3\frac{a_2b_2^2}{ab^2} \underline{e}_\lambda \otimes \underline{e}_\lambda \right] \quad (26)$$

where it is recalled that $\mathbf{T}_2 = \mathbf{T}(\lambda_2)$ and $\mathbf{T}(\lambda)$ is given by (21). Taking into account this expression, the surface strain rate tensor, \mathbf{d}_s , defined by the first relation in (7) reads:

$$\mathbf{d}_s = \mathbf{P}(\theta, \varphi) \cdot \left[\mathbf{D} + D_m(\mathbf{T}_1 - \mathbf{T}_2) \right] \cdot \mathbf{P}(\theta, \varphi) \quad (27)$$

for which account has been taken of the identity $\mathbf{P}(\varphi, \theta) \cdot \underline{e}_\lambda = 0$ and $\mathbf{T}_1 = \mathbf{T}(\lambda_1)$. As it will be shown in the next section, the representation of the interfacial plastic strain rate tensor, in the form given by (27), is required for the

computation of the integrals which enter into the definition of $\Pi(\mathbf{D})$ given by (11).

4 Closed-form expression of the macroscopic dissipation

The trial velocity field being entirely determined by the boundary conditions, no minimization procedure is required for $\Pi(\mathbf{D})$ which, according to (15), can be put into the form:

$$\Pi(\mathbf{D}) = \Pi_1(\mathbf{D}) + \Pi_2(\mathbf{D}) + \Pi_3(\mathbf{D}) \quad (28)$$

where the quantities $\Pi_1(\mathbf{D})$, $\Pi_2(\mathbf{D})$ and $\Pi_3(\mathbf{D})$ read:

$$\begin{aligned} \Pi_1(\mathbf{D}) &= \frac{1}{V_2} \int_{\Omega-\omega} \sigma_0 d_{eq} dV \\ \Pi_2(\mathbf{D}) &= \frac{1}{V_2} \int_{\Gamma} \tau_r \operatorname{tr}(\mathbf{d}_s) dS \\ \Pi_3(\mathbf{D}) &= \frac{1}{V_2} \int_{\Gamma} \tau_0 d_{eq}^s dS \end{aligned} \quad (29)$$

Expressions of these quantities will be provided in the following sections, 4.1, 4.2 and 4.3, respectively.

4.1 Approximate expression of $\Pi_1(\mathbf{D})$

Based on the trial velocity field (17) (together with (18)), Gologanu et al. [31,32] derived closed-form expressions of $\Pi_1(\mathbf{D})$ and then a two-field estimate of the macroscopic criterion in the cases of a prolate and an oblate cavity. Although the two approximate expressions obtained by these authors prove to be accurate, comparatively to the exact two-field solution (evaluated numerically without any approximations), we propose to use the unified expression established in [49], valid for both cases of a prolate and an oblate cavity. The latter, derived in the more general case of an anisotropic Hill plastic porous medium with spheroidal cavities, reduced in the limit case of the von Mises matrix to:

$$\Pi_1(\mathbf{D}) = \sigma_0 f \left[Y \operatorname{arcsinh} \left\{ \frac{uY}{X} \right\} - \frac{\sqrt{X^2 + u^2 Y^2}}{u} \right]_{u=u_1}^{u=u_2} \quad (30)$$

in which the quantities X and Y depend on the macroscopic strain rate tensor as:

$$\begin{aligned} X &= \left[(1 - \zeta)(D_q + (1 - 3\alpha_2)D_m)^2 + D_s^2 + D_t^2 \right]^{1/2}, \\ Y &= \frac{1}{f\kappa} [3 + \eta(1 - 3\alpha_2)] D_m + \frac{\eta}{f\kappa} D_q \end{aligned} \quad (31)$$

Expressions of coefficients η , κ and ζ as well as u_1 and u_2 are reported in appendix A. The invariant quantities, D_q , D_s and D_t , in agreement with the void-induced transverse isotropy (with a symmetry axis Ox_3), are defined as:

$$\begin{cases} D_q = \frac{2}{3}D_{33} - \frac{1}{3}(D_{11} + D_{22}) \\ D_s = \frac{1}{\sqrt{3}}\sqrt{(D_{22} - D_{11})^2 + 4D_{12}^2} \\ D_t = \frac{2}{\sqrt{3}}\sqrt{D_{13}^2 + D_{23}^2} \end{cases} \quad (32)$$

The macroscopic equivalent strain rate can then be decomposed as $D_{eq}^2 = D_q^2 + D_s^2 + D_t^2$.

4.2 Exact expression for $\Pi_2(\mathbf{D})$

We now propose to compute the exact expression of the dissipation $\Pi_2(\mathbf{D})$ associated to the interfacial residual stresses. Introducing expression (27) in the second relation in (29) leads to:

$$\Pi_2(\mathbf{D}) = \frac{\tau_r S_1}{V_2} [\mathbf{D} + D_m(\mathbf{T}_1 - \mathbf{T}_2)] : \overline{\mathbf{P}} \quad (33)$$

where the second order tensor $\overline{\mathbf{P}}$ is defined by:

$$\overline{\mathbf{P}} = \frac{1}{S_1} \int_{\Gamma} \mathbf{P}(\theta, \varphi) dS = \frac{1}{2}(1 + \gamma)\mathbf{I}_t + (1 - \gamma)\mathbf{I}_n \quad (34)$$

and S_1 is the area of the spheroid of coordinate $\lambda = \lambda_1$:

$$S_1 = \begin{cases} \frac{2\pi a_1^2 b_1}{c} \arctan \left\{ \frac{c}{b_1} \right\} + 2\pi b_1^2 & \text{(prolate void)} \\ \frac{2\pi a_1^2 b_1}{c} \operatorname{arctanh} \left\{ \frac{c}{b_1} \right\} + 2\pi b_1^2 & \text{(oblate void)} \end{cases} \quad (35)$$

The parameter γ , which depends on the cavity shape, reads:

$$\gamma = \frac{b_1^2}{a_1^2 - b_1^2} \left(1 - \frac{4b_1^2\pi}{S_1} \right) \quad (36)$$

For more details concerning the computation of the integral in (34), one may refer to appendix B.

The value of γ for the particular cases of a spherical, cylindrical and penny-shaped crack are provided in table 1. The variation of γ with the void eccentricity e_1 are given on figure 4. It can be noted that, for any value of eccentricity, γ lies in the range $[0, 1]$.

	sphere	cylinder	crack
γ	1/3	0	1

Table 1

Values of γ for some particular cases of a spheroidal cavity.

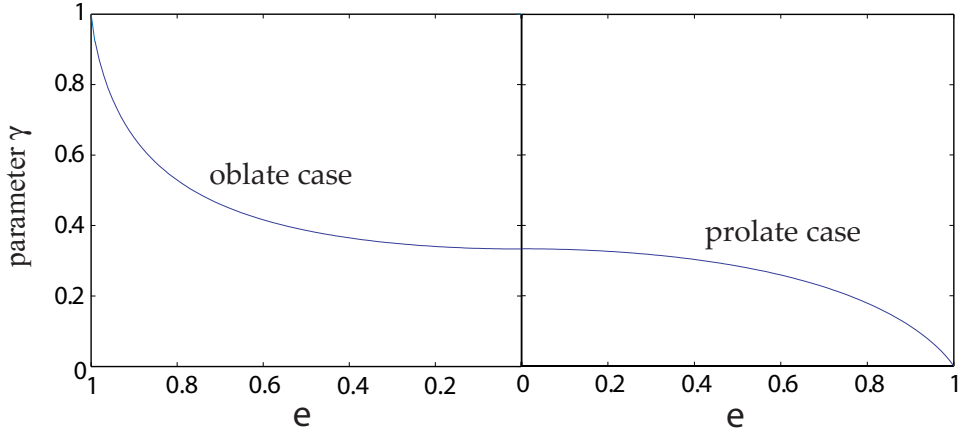


Fig. 4. γ as function of e_1 for a prolate and an oblate void.

Interestingly, $\Pi_2(\mathbf{D})$, given by (33), can be put into the form:

$$\Pi_2(\mathbf{D}) = \frac{\tau_r S_1}{V_1} \mathbf{D} : \mathbf{N} \quad (37)$$

where $V_1 = 4\pi a_1 b_1^2/3$ is the volume of the spheroidal cavity while the second order tensor \mathbf{N} is given by:

$$\mathbf{N} = \left[\mathbb{I} + \frac{1}{3} \mathbf{I} \otimes (\mathbf{T}_1 - \mathbf{T}_2) \right] : \bar{\mathbf{P}} \quad (38)$$

\mathbb{I} , whose components are $I_{ijkl} = (\delta_{ik}\delta_{jl} + \delta_{il}\delta_{jk})/2$, is the fourth order identity tensor, δ_{ij} being the Kronecker symbol. From the definition (21) of $\mathbf{T}(\lambda)$ and the expression (34) of $\overline{\mathbf{P}}$, one obtains:

$$\begin{aligned} \mathbf{N} = & \frac{1}{2} \left[1 + \gamma + (\alpha_1 - f\alpha_2)(1 - 3\gamma) \right] \mathbf{I}_t \\ & + \left[1 - \gamma - \frac{1}{2}(1 - \alpha_1 - f(1 - \alpha_2)(1 - 3\gamma)) \right] \mathbf{I}_n \end{aligned} \quad (39)$$

This tensor depends on the eccentricity e_1, e_2 (through α_1, α_2 and γ in (19) and (36)) and of the porosity f . In the particular case of spherical void, \mathbf{N} reduces to $\frac{2}{3}\mathbf{I}$. In the general case, it is convenient to express $\Pi_2(\mathbf{D})$ as function of D_m, D_q, D_s and D_t , given in (32):

$$\Pi_2(\mathbf{D}) = \frac{\tau_r S_1}{V_1} (g_1 D_m + g_2 D_q) \quad (40)$$

with:

$$g_1 = 2 - \frac{1}{2}(1 - 3\gamma)(1 - 3\alpha_1 - f(1 - 3\alpha_2)), \quad g_2 = \frac{f}{2}(1 - 3\gamma) \quad (41)$$

4.3 An upper bound for $\Pi_3(\mathbf{D})$

We are now interested by the derivation of the dissipation $\Pi_3(\mathbf{D})$ associated to the plastic deformation of the surface Γ . It is defined by the last relation in (29), in which the interfacial equivalent strain rate, d_{eq}^s , is given by (7). The required expressions, (7) and (27), for the interfacial plastic strain rate leads to:

$$d_{eq}^s = \left[\frac{2}{3} (\mathbf{D} + D_m(\mathbf{T}_1 - \mathbf{T}_2)) : \mathbb{P}(\theta, \varphi) : (\mathbf{D} + D_m(\mathbf{T}_1 - \mathbf{T}_2)) \right]^{1/2} \quad (42)$$

where the fourth order tensor $\mathbb{P}(\theta, \varphi)$ is given by:

$$\mathbb{P}(\theta, \varphi) = \mathbf{P}(\theta, \varphi) \overline{\otimes} \mathbf{P}(\theta, \varphi) + \mathbf{P}(\theta, \varphi) \otimes \mathbf{P}(\theta, \varphi) \quad (43)$$

for which the following notation has been used (for any couple of second order tensors \mathbf{a} and \mathbf{b}):

$$(\mathbf{a} \overline{\otimes} \mathbf{b})_{ijkl} = \frac{1}{2} (a_{ik} b_{jl} + a_{il} b_{jk}) \quad (44)$$

At this stage, it is not possible to derive a closed-form expression of $\Pi_3(\mathbf{D})$, even in the simplest case of a spherical void (see [18]). However, based on the Cauchy-Schwartz inequality, one can establish an upper bound for $\Pi_3(\mathbf{D})$. Indeed, for the positive and square integrable function d_{eq}^s , one has:

$$\frac{1}{S_1} \int_{\Gamma} d_{eq}^s dS \leq \sqrt{\frac{1}{S_1} \int_{\Gamma} (d_{eq}^s)^2 dS} \quad (45)$$

The following approximation is then used:

Approximation A1 : In the expression of $\Pi_3(\mathbf{D})$ the integral of the interfacial equivalent strain rate over the surface of the spheroidal cavity, i.e. $\int_{\Gamma} d_{eq}^s dS$, is replaced by $S_1 \left[\frac{1}{S_1} \int_{\Gamma} (d_{eq}^s)^2 dS \right]^{1/2}$.

With this approximation, an upper bound of the plastic dissipation $\Pi_3(\mathbf{D})$ is:

$$\begin{aligned} \Pi_3(\mathbf{d}) &= \frac{S_1}{V_2} \left[\frac{1}{S_1} \int_{\Gamma} (d_{eq}^s)^2 dS \right]^{1/2} \\ &= \frac{S_1}{V_2} \left[\frac{2}{3} (\mathbf{D} + D_m(\mathbf{T}_1 - \mathbf{T}_2)) : \bar{\mathbb{P}} : (\mathbf{D} + D_m(\mathbf{T}_1 - \mathbf{T}_2)) \right]^{1/2} \end{aligned} \quad (46)$$

where the fourth order tensor $\bar{\mathbb{P}}$ is defined by:

$$\begin{aligned} \bar{\mathbb{P}} &= \frac{1}{S_1} \int_{\Gamma} \mathbb{P}(\theta, \varphi) dS = (3\gamma + \mu)\mathbb{E}_1 + 2\mu\mathbb{E}_2 + \frac{2\gamma + \mu}{2}\mathbb{E}_3 \\ &+ \left[\frac{5}{2}(1 - \gamma) - 2\mu \right] \mathbb{E}_4 + \left[\frac{3}{2}(1 - \gamma) - \mu \right] (\mathbb{E}_5 + \mathbb{E}_6) \end{aligned} \quad (47)$$

\mathbb{E}_n for $n = 1..6$ are the fourth order tensors of the Walpole basis [65] (the elements of this basis are recalled in appendix B). Parameter μ is given by:

$$\mu = \frac{a^2}{a^2 - b^2} (1 - 3\gamma) \quad (48)$$

The details concerning the computation of the integral in (47) can be found in appendix B.

The values of μ for a spherical or a cylindrical void, as well as for a penny-shaped crack are provided in table 2. For any value of the eccentricity, coefficient μ lies in the range $[0, 1]$. Its variation with respect to e_1 is shown on figure 5.

Table 2

Values of μ for particular cases of a spheroidal cavity.

	sphere	cylinder	crack
μ	8/15	1	0

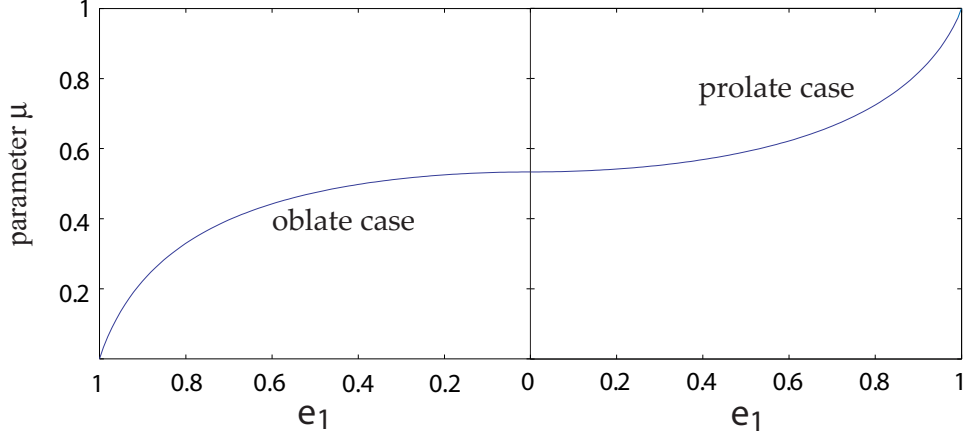


Fig. 5. Variations of μ with the eccentricity e_1 in the case of prolate and an oblate void.

Finally, $\Pi_3(\mathbf{D})$ can be put into the form:

$$\Pi_3(\mathbf{D}) = \frac{\tau_0 S_1}{V_1} \sqrt{\frac{2}{3} \mathbf{D} : \mathbb{S} : \mathbf{D}} \quad (49)$$

where the fourth order tensor \mathbb{S} is defined by:

$$\mathbb{S} = f^2 \left[\mathbb{I} + \mathbf{I} \otimes (\mathbf{T}_1 - \mathbf{T}_2) \right] : \bar{\mathbb{P}} : \left[\mathbb{I} + (\mathbf{T}_1 - \mathbf{T}_2) \otimes \mathbf{I} \right] \quad (50)$$

and the components of \mathbb{S} are provided in appendix B. As already done for $\Pi_1(\mathbf{D})$ and $\Pi_2(\mathbf{D})$, we now express $\Pi_3(\mathbf{D})$ as function of the strain rate invariants introduced in (32):

$$\Pi_3(\mathbf{D}) = \frac{\tau_0 S_1}{V_1} \sqrt{h_1 D_m^2 + h_2 D_q^2 + 2h_3 D_m D_q + h_4 D_s^2 + h_5 D_t^2} \quad (51)$$

where the coefficients h_n are given by:

$$\left\{ \begin{array}{l} h_1 = (3\gamma + 3\mu - 2)U^2 + 2(3\gamma - 1)U + 4 \\ h_2 = f^2(3\gamma + 3\mu - 2) \\ h_3 = f(3\gamma + 3\mu - 2)U + f(3\gamma - 1) \\ h_4 = \frac{f^2}{2}(2\gamma + \mu) \\ h_5 = \frac{f^2}{2}(5 - 4\mu - 5\gamma) \end{array} \right. \quad (52)$$

with $U = 1 - 3\alpha_1 - f(1 - 3\alpha_2)$, γ being given by (36) and μ by (48).

5 An estimate of the macroscopic yield surface

5.1 The parametric equations

The macroscopic yield surface of the nanoporous medium is obtained from equation (10) in which the macroscopic dissipation $\Pi(\mathbf{D})$ is obtained as the sum of $\Pi_1(\mathbf{D})$, $\Pi_2(\mathbf{D})$ and $\Pi_3(\mathbf{D})$ whose expressions are given by (30), (40) and (51) respectively.

Equation (10) defines a hyper-surface of dimension 5 in the space of macroscopic stress. Note that $\Pi_1(\mathbf{D})$, $\Pi_2(\mathbf{D})$ and $\Pi_3(\mathbf{D})$ and then $\Pi(\mathbf{D})$ depends on the macroscopic strain rate tensor, \mathbf{D} , only through the invariants D_m , D_q , D_s and D_t . The yield surface can then be deduced from the following parametric equations:

$$\left\{ \begin{array}{l} \Sigma_m = \frac{1}{3} \frac{\partial \Pi}{\partial D_m} \\ \Sigma_q = \frac{\partial \Pi}{\partial D_q} \\ \Sigma_s = \frac{\partial \Pi}{\partial D_s} \\ \Sigma_t = \frac{\partial \Pi}{\partial D_t} \end{array} \right. \quad (53)$$

where Σ_m , Σ_q , Σ_s and Σ_t are the invariants of the macroscopic stress defined by:

$$\left\{ \begin{array}{l} \Sigma_m = \frac{1}{3}(\Sigma_{11} + \Sigma_{22} + \Sigma_{33}) \\ \Sigma_q = \Sigma_{33} - \frac{1}{2}(\Sigma_{11} + \Sigma_{22}) \\ \Sigma_s = \frac{\sqrt{3}}{2}\sqrt{(\Sigma_{22} - \Sigma_{11})^2 + 4\Sigma_{12}^2} \\ \Sigma_t = \sqrt{3}\sqrt{\Sigma_{13}^2 + \Sigma_{23}^2} \end{array} \right. \quad (54)$$

Relations (53) corresponds to a surface of dimension 3 in the space of dimension 4 defined by $(\Sigma_m, \Sigma_q, \Sigma_s, \Sigma_t)$. Unfortunately, for general loading case, it is not possible to derive a closed-form solution of the system given by (53), even for cylindrical or spherical voids. However, we succeed to identify closed-form expressions in some particular cases of macroscopic stress loading or values of the material parameters σ_0 , τ_0 and τ_r . These solutions are presented in the next section.

The impossibility to obtain an explicit expression of the macroscopic yield surface does not matter by itself since relations (53) provide a set of parametric equations which are sufficient for the representation of the yield surface. Introducing the following three non dimensional variables $\xi_1 = D_q/D_m$, $\xi_2 = D_s/D_m$ and $\xi_3 = D_t/D_m$, relations (53) can be rewritten as:

$$\left\{ \begin{array}{l} \Sigma_m = \mathcal{G}_1(\xi_1, \xi_2, \xi_3) \\ \Sigma_q = \mathcal{G}_2(\xi_1, \xi_2, \xi_3) \\ \Sigma_s = \mathcal{G}_3(\xi_1, \xi_2, \xi_3) \\ \Sigma_t = \mathcal{G}_4(\xi_1, \xi_2, \xi_3) \end{array} \right. \quad (55)$$

Functions $\mathcal{G}_n(\xi_1, \xi_2, \xi_3)$ for $n = 1, 2, 3, 4$ are explicitly given in appendix C. The macroscopic yield surface can then be obtained by varying ξ_1 , ξ_2 and ξ_3 from $-\infty$ to $+\infty$.

Remark 1: The term corresponding to $\Pi_2(\mathbf{D})$ is linear with respect to the macroscopic strain rate tensor, \mathbf{D} ; this introduces the constant tensor $\frac{\tau_r S_1}{V_1} \mathbf{N}$ in equation (10). The derivatives of $\Pi_1(\mathbf{D})$ and $\Pi_3(\mathbf{D})$ according to \mathbf{D} are null for $\mathbf{D} = 0$. Consequently, the interfacial residual stress, through the term $\frac{\tau_r S_1}{V_1} \mathbf{N}$, induces a change in position of the center of the macroscopic yield surface. Since tensor \mathbf{N} is axisymmetric (it has only two components along \mathbf{I}_t and \mathbf{I}_n), the residual stress translates the macroscopic yield locus along the direction related to the macroscopic stress defined by Σ_m and Σ_q . The position of the center in the space $(\Sigma_m, \Sigma_q, \Sigma_s, \Sigma_t)$ is denoted $(C_m, C_q, 0, 0)$ where C_m

and C_q are given by:

$$C_m = \frac{\tau_r S_1}{3V_1} g_1, \quad C_q = \frac{\tau_r S_1}{V_1} g_2 \quad (56)$$

Coefficients g_1 and g_2 are given by relations (41). Note also that C_q is null in the case of a spherical void (corresponding to $\gamma = 1/3$) which suggests that the interfacial residual stress τ_r only induces a translation of the macroscopic yield locus along the axis related to Σ_m .

Remark 2: By replacing ξ_1 by $-\xi_1$ in the parametric equations (55), Σ_m and Σ_q are changed in $-\Sigma_m$ and $-\Sigma_q$ when $\tau_r = 0$. This suggests that the macroscopic yield surface, in the plane (Σ_m, Σ_q) , is symmetric with respect to its center (C_m, C_q) . The representation of this yield surface can then be restricted to the half plane $(\Sigma_m > 0, \Sigma_q)$ in the absence of residual stresses.

Remark 3: By replacing ξ_2 by $-\xi_2$ in the parametric equations (55), Σ_m is changed in $-\Sigma_m$ while Σ_s remains unchanged. It follows that the yield locus is symmetric with respect to axis $\Sigma_s = 0$ in the plane (Σ_m, Σ_s) . Noting also that Σ_s is positive by definition, the yield locus is represented in the quart plane $(\Sigma_m > 0, \Sigma_s > 0)$. This remark also holds for the component Σ_t at the place of Σ_s .

Remark 4: In the particular case of spherical voids, expressions of $\mathcal{G}_1, \mathcal{G}_2, \mathcal{G}_3, \mathcal{G}_4$ are given in appendix C. This leads to a parametric equation for the mean stress Σ_m and equivalent stress Σ_{eq} :

$$\left\{ \begin{array}{l} \Sigma_m = \frac{2\sigma_0}{3} \left[\operatorname{arcsinh} \left(\frac{2u}{f\xi} \right) \right]_f^1 + \frac{2\tau_r}{a_1} + \frac{\tau_0}{a_1} \frac{4}{\sqrt{4 + \frac{3f^2}{5}\xi^2}} \\ \Sigma_{eq} = -\sigma_0 \left[\frac{\sqrt{4u^2 + f^2\xi^2}}{u\xi} \right]_f^1 + \frac{\tau_0}{a_1} \frac{9\xi f^2}{5\sqrt{4 + \frac{3f^2}{5}\xi^2}} \end{array} \right. \quad (57)$$

where $\xi = D_{eq}/D_m$ and a_1 is the radius of the nanovoid.

By incorporating the effects of interfacial residual stress, (57) extends the results established by Dormieux and Kondo [18] who investigated the case of spherical voids.

5.2 Closed form expression of the yield function in some particular cases

- *Case 1:* $\tau_0 = 0$

In this case, the term related to the interfacial plastic strain dissipation, $\Pi_3(\mathbf{D})$, vanishes in the expression of $\Pi(\mathbf{D})$. The macroscopic criterion takes then the form (see appendix D for more details):

$$\begin{aligned} \frac{(\Sigma_q - C_q - \eta T)^2}{(1 - \zeta)\sigma_0^2} + \frac{\Sigma_s^2}{\sigma_0^2} + \frac{\Sigma_t^2}{\sigma_0^2} \\ + 2(1 + g)(f + g) \cosh\left(\frac{\kappa T}{\sigma_0}\right) - (1 + g)^2 - (f + g)^2 = 0 \end{aligned} \quad (58)$$

with $T = \Sigma_m - C_m - \frac{1}{3}(1 - 3\alpha_2)(\Sigma_q - C_q)$. It is recalled that coefficients η , ζ , κ and g are given in appendix A.

For the particular case of spherical voids (for which $a_1 = b_1$, $\zeta = 0$, $\eta = 0$, $\kappa = 3/2$ and $\alpha_2 = 1/3$), note that the above expression coincides with that given by [18] when the pore space is filled by a fluid and capillary effects develop in the solid/fluid interface (here $\tau_r = \gamma^{sf}$, the quantity γ^{sf} denoting the surface tension):

$$\frac{\Sigma_{eq}^2}{\sigma_0^2} + 2f \cosh\left(3 \frac{\Sigma_m - 2\tau_r/a_1}{2\sigma_0}\right) - 1 - f^2 = 0 \quad (59)$$

Note that (59) reduces to the well-known result of Gurson [35] when $\tau_r = 0$ or for large values of cavity radius a_1 . In the same vein, when $\tau_r = 0$ and cavities are spheroidal, (58) reduces to the one established in [49] for a von Mises Matrix (obtained as a particular case of a plastically anisotropic matrix)¹.

- *Case 2:* $\frac{\partial \Pi_1(\mathbf{D})}{\partial(\mathbf{D})} \ll \frac{\partial \Pi_2(\mathbf{D})}{\partial(\mathbf{D})}$ and $\frac{\partial \Pi_1(\mathbf{D})}{\partial(\mathbf{D})} \ll \frac{\partial \Pi_3(\mathbf{D})}{\partial(\mathbf{D})}$

Introducing (37) and (49) into relation (10) yields:

$$\boldsymbol{\Sigma} = \frac{S_1}{V_1} \left[\tau_r \mathbf{N} + \frac{2\tau_0}{3} \frac{\mathbb{S} : \mathbf{D}}{\sqrt{\frac{2}{3}} \mathbf{D} : \mathbb{S} : \mathbf{D}} \right] \quad (60)$$

with \mathbf{N} given by (39) and \mathbb{S} is provided in appendix B. (60) leads to the following macroscopic elliptic criterion:

$$\frac{3}{2} \left[\boldsymbol{\Sigma} - \frac{\tau_r S_1}{V_1} \mathbf{N} \right] : \mathbb{S}^{-1} : \left[\boldsymbol{\Sigma} - \frac{\tau_r S_1}{V_1} \mathbf{N} \right] = \left(\frac{\tau_0 S_1}{V_1} \right)^2 \quad (61)$$

which is anisotropic due to the presence of directional tensors \mathbf{N} and \mathbb{S} .

In the case of a hollow sphere, one has $S_1 = 4\pi a_1^2$, $V_1 = 4\pi a_1^3/3$, $\mathbf{N} = 2/3 \mathbf{I}$

¹ Note that the latter has the same form as the GLD-model [31,32] but with other definitions of the coefficients η , ζ , κ , g .

and $\mathbb{S} = 2\mathbb{J} + \frac{3f^2}{5}\mathbb{K}$ in which $\mathbb{J} = \frac{1}{3}\mathbf{I} \otimes \mathbf{I}$, $\mathbb{K} = \mathbb{I} - \mathbb{J}$, \mathbb{I} being the fourth order symmetric identity. The macroscopic criterion then reduces to:

$$\frac{5}{27f^2}\Sigma_{eq}^2 + \frac{1}{4}\left(\Sigma_m - \frac{2\tau_r}{a_1}\right)^2 = \left(\frac{\tau_0}{a_1}\right)^2 \quad (62)$$

- *Case 3:* $D_m = D_q = D_t = 0$

The hollow spheroid is subjected to the following macroscopic shear strain rate loading: $\mathbf{D} = D_{12}\underline{e}_1 \overset{s}{\otimes} \underline{e}_2$ or $\mathbf{D} = D_{11}(\underline{e}_1 \otimes \underline{e}_1 - \underline{e}_2 \otimes \underline{e}_2)$. In that cases, $D_s = 2|D_{12}|/\sqrt{3}$ or $D_s = 2|D_{11}|/\sqrt{3}$, respectively. The macroscopic criterion corresponding to this shear loading reads:

$$\Sigma_s = (1-f)\sigma_0 + \frac{\tau_0 S_1}{V_1} f \sqrt{\gamma + \frac{\mu}{2}} \quad (63)$$

It is noted that this result provides an additional interfacial term, $\frac{\tau_0 S_1}{V_1} f \sqrt{\gamma + \frac{\mu}{2}}$, comparatively to the GLD-model [31,32] or result by [49]. Note also that the interfacial residual stress τ_r does not affect the above shear strength Σ_s .

- *Case 4:* $D_m = D_q = D_s = 0$

This is similar to the above one and corresponds to the following macroscopic shear strain rate loading, $\mathbf{D} = D_{13}\underline{e}_1 \overset{s}{\otimes} \underline{e}_3$ or $\mathbf{D} = D_{23}\underline{e}_2 \overset{s}{\otimes} \underline{e}_3$, for which $D_t = 2|D_{13}|/\sqrt{3}$ and $D_t = 2|D_{23}|/\sqrt{3}$ respectively. The macroscopic shear strength corresponding to this loading is then given by:

$$\Sigma_t = (1-f)\sigma_0 + \frac{\tau_0 S_1}{V_1} f \sqrt{\frac{5}{2}(1-\gamma) - 2\mu} \quad (64)$$

6 Illustrations and validations with numerical calculations

In this section, we propose to illustrate the salient features of the new macroscopic criterion of a ductile material containing spheroidal nanocavities. In particular, the combined effects of voids size and voids shape will be analyzed. These illustrations are considered for an axisymmetric loading represented by the values of the stress components $\Sigma_q = \Sigma_{33} - \Sigma_{11}$ and $\Sigma_m = (2\Sigma_{11} + \Sigma_{33})/3$. We assume that $\Sigma_s = \Sigma_t = 0$, which consists to impose $D_s = D_t = 0$ in the expressions of the macroscopic dissipation and then to put $\xi_2 = \xi_3 = 0$ in the parametric equations (55). The macroscopic criterion is then described by two parametric equations evolving the macroscopic stress invariants Σ_m and

Σ_t and the parameter ξ_1 .

We will also check the accuracy of the approximate criterion by comparing its predictions to numerical solutions. Indeed, equations (55) have been derived from (53) in which a number of approximations has been used in order to obtain a closed-form expression of the macroscopic dissipation $\Pi(\mathbf{D})$. The accuracy of the approximate criterion needs then to be assessed through comparisons with the exact two-field criterion computed numerically. We will first deal with a spherical and then a cylindrical void; the cases of an arbitrary prolate and oblate cavity is thereafter investigated. Finally, we will focus on the case of a penny-shaped crack for which particular attention is needed when computing the limit of $\Pi_2(\mathbf{D})$ and $\Pi_3(\mathbf{D})$ when the aspect ratio tends to zero.

6.1 Case of spherical nano-voids

6.1.1 Assessment from atomistic simulations

In order to investigate the role of the interfacial stress on the macroscopic yielding of plastic materials containing nanovoids we propose a qualitative evaluation of the the interfacial parameters τ_0 and τ_r . To this end, we refer to the recent study of Mi et al. [47] who quantify the effect of void radius on the yield strength of a porous cubic single crystal aluminium by making use of atomistic simulations (see also Traiviratana et al. [63] who provide similar results). In this work, they consider a unit cell containing a spherical cavity submitted to a uniaxial tension E_{33} , other components of the macroscopic strain remaining equal to zero. Four size for the cavity radius has been considered while the volume fraction of cavity is kept constant and equal to $f = 0.0042$. The numerical data in [47] have been reported on figure 6 (the circles) and show the values of the yield strength on the macroscopic yield stress for Σ_{11} as function of the radius of the cavity. It must be noted that these results show a strong effect of the cavity radius on the macroscopic yield strength.

We now propose to compare the results obtained by atomistic simulations with the macroscopic model defined by the parametric equations (57). Under uniaxial tension, the strain rate tensor and the stress tensor are on the form:

$$\mathbf{D} = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & D_{33} \end{pmatrix}, \quad \boldsymbol{\Sigma} = \begin{pmatrix} \Sigma_{11} & 0 & 0 \\ 0 & \Sigma_{11} & 0 \\ 0 & 0 & \Sigma_{33} \end{pmatrix} \quad (65)$$

For such loading case, $D_m = D_{33}/3$ while $D_{eq} = 2D_{33}/3$. Parameter ξ , which appears in equations (57), takes then the value $\xi = 2$. In other hand, the yield stress Σ_{33} is given by $\Sigma_{33} = \Sigma_m + 2\Sigma_{eq}/3$ where Σ_m and Σ_{eq} are given by (57). It must be noted that the dependence of Σ_{33} with the void radius is on the form $\Sigma_{33} = P + Q/a_1$ for which the two constants P and Q are given by:

$$\begin{cases} P = \frac{2\sigma_0}{3} \left[\operatorname{arcsinh} \left(\frac{u}{f} \right) - \frac{\sqrt{u^2 + f^2}}{u} \right]_f^1 \\ Q = \frac{2\tau_r}{a_1} + \frac{2\tau_0}{a_1} \sqrt{1 + \frac{3f^2}{5}} \end{cases} \quad (66)$$

and which are evaluated with the numerical data provided in the paper of Mi et al. [47]. This computation leads to: $P = 2.266 \text{ GPa}$ and $Q = 1.089 \text{ GPa}$. The volume fraction is known and equal to $f = 0.0042$, the value of σ_0 is deduced from the first relation in (66), it gives $\sigma_0 \simeq 0.59 \text{ GPa}$. In (66), the value of Q depends on two unknown material parameters, τ_0 and τ_r . In order to obtain a qualitative evaluation of τ_0 and τ_r , we first consider the particular case $\tau_r = 0$ and then $\tau_0 = 0$. The values of the porosity being very small, it is possible, for this qualitative evaluation of the surface parameters, to neglect the term $3f^2/4$ under the square root in the expression of Q . This coefficient is then approximated by:

$$Q = \frac{2\tau_r}{a_1} + \frac{2\tau_0}{a_1} \quad (67)$$

Note that in the above expression the effect of τ_0 and τ_r on coefficient Q is the same. If $\tau_r = 0$, one has $\tau_0 = 0.545 \text{ N/m}$. The second case, $\tau_0 = 0$, leads to the same value for τ_r . Note that the computation of both τ_0 and τ_r could be done by performing the atomistic simulation in the case of uniaxial compression, indeed, it will be recalled that the term containing τ_r introduces an asymmetry between tension and compression.

In figure 6, the numerical data obtained from atomistic simulations (the circles) are compared to the parametric criterion (57) with the calibrated material parameters τ_0 and τ_r (the full line). Just as information, we also represent the prediction of the Gurson model, which do not account for the size effect, and which corresponds to $\Sigma_{33} = P$. It must be noted that the model can reproduce

the dependence of the yield strength with the cavity radius in the particular case of spherical shape.

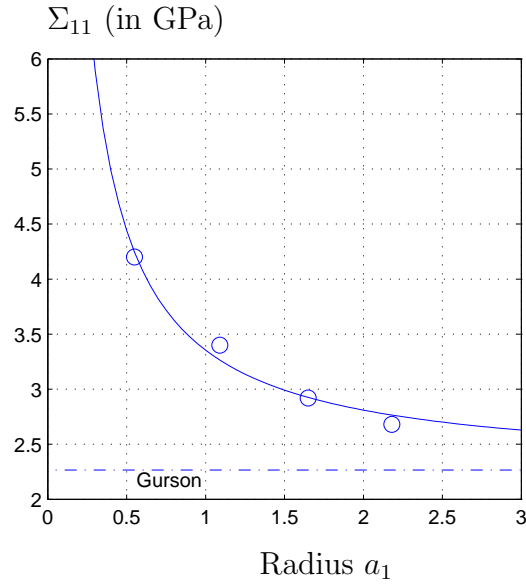


Fig. 6. Yield strength for Σ_{11} as function of the radius of the cavity. Comparison between the numerical data obtained from atomistic simulation by Mi et al. [47] (circles), the macroscopic model $\Sigma_{11} = P + Q/a_1$ with optimized values for P and Q (the full line) and the Gurson prediction (dashed line).

6.1.2 Representation of the macroscopic criterion

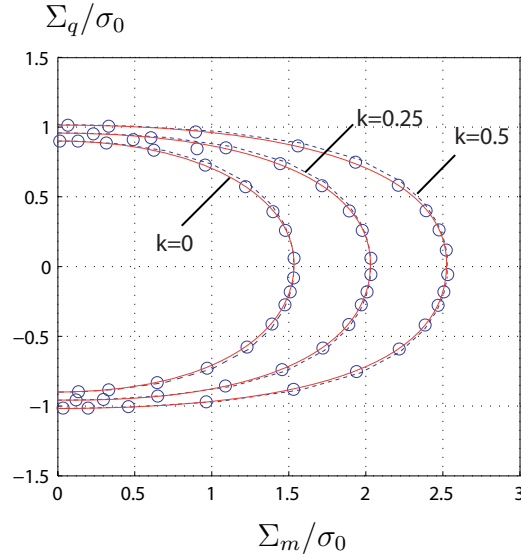


Fig. 7. Yield loci for spherical nanovoids. Comparison between the approximate criterion (dashed line) and the exact, numerical two-field criteria (full line with circles) for a porosity $f = 0.1$ and for various values of $k = \tau_0/(a_1\sigma_0)$.

Figure 7 illustrates the particular case of spherical nano-cavities for which the approximate yield surface is given by the parametric equation (55). Two values of the non dimensional parameter $k = \tau_0/(a_1\sigma_0)$, namely $k = 0.25$ and $k = 0.5$, as well as a fixed porosity $f = 0.1$ and $\tau_r = 0$ are considered. By using the values $\tau_0 = 0.54N/m$ and $\sigma_0 = 0.59GPa$, the cases $k = 0.25$ and $k = 0.5$ correspond to a cavity of radius $1.83nm$ and $3.66nm$ respectively. The results show an important effect of the parameter k on the macroscopic yield locus: the yield strength domain increases when the voids size decreases. In this figure, a very good agreement is also observed between the approximate criterion and the results obtained by numerically evaluating the macroscopic dissipation (15).

6.2 Case of cylindrical nano-cavities

The case of a hollow cylinder corresponds to the limits $a_1 \rightarrow \infty$ and $a_2 \rightarrow \infty$, that is $e_1 \rightarrow 1$ and $e_2 \rightarrow 1$. On figure 8 is represented the macroscopic yield function in the plane (Σ_m, Σ_q) for a porosity $f = 0.1$ and for various values of the non dimensional parameter $k = \tau_0/(b_1\sigma_0)$; the residual stress $\tau_r = 0$ for this illustration. These results show an important effect of the cavity size on the macroscopic yield locus, similar to that observed for spherical voids.

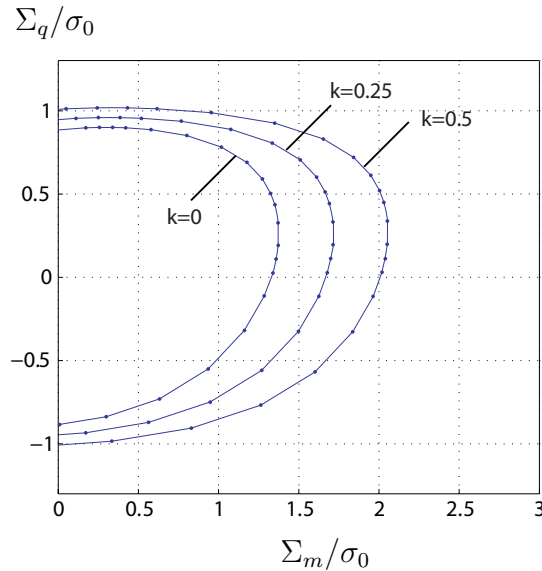


Fig. 8. Yield loci for the ductile material with a cylindrical nanovoid for a porosity $f = 0.1$ and for various values of $k = \tau_0/(b_1\sigma_0)$.

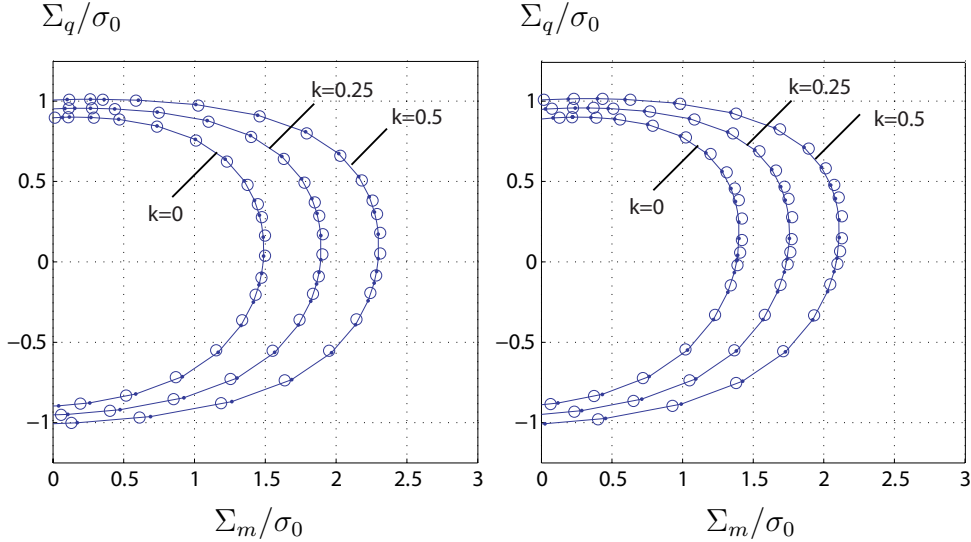


Fig. 9. Yield loci for a prolate cavity with an aspect ratio $a_1/b_1 = 2$ (at the left) $a_1/b_1 = 5$ (at the right) and a porosity $f = 0.1$. Comparison between the approximate (full line with points) and the exact, numerical, two-field criteria (circles) for various values of $k = \tau_0/(b_1\sigma_0)$.

6.3 Ductile materials with spheroidal nanovoids

On Figure 9, is represented the macroscopic yield surface for the case of a prolate void. On the left, this yield surface is plotted for an aspect ratio $a_1/b_1 = 2$ and a porosity $f = 0.1$, while, at the right, the figure corresponds to $a_1/b_1 = 5$ and the same porosity. k is the non dimensional parameter defined as $k = \tau_0/(b_1\sigma_0)$. To check the accuracy of the approximate criterion (full line with the points), we compare the results with numerical data of the exact two-field criterion (circles): a very good agreement is observed. Figure 10 displays similar results for the case of an oblate cavity. At the right, the aspect ratio is $a_1/b_1 = 1/2$ while at the left, an aspect ratio of $a_1/b_1 = 1/5$ is considered. A very good agreement between the "approximate" and the "exact" two-field criterion is observed for low values of the parameter $k = \tau_0/(b_1\sigma_0)$. When increasing k , some slight differences are noted in the region of positive values of Σ_q and high values of the mean stress Σ_m . Beyond that differences, it can be also noted that the approximate criterion seems to preserve the upper bound character of the limit analysis-based approach, since the predictions of the approximate two-field criterion are larger than the numerical ones.

We now represent on figure 11 the yield strength of the nanoporous medium as function of the parameter $k = \tau_0/(b_1\sigma_0)$. At the right, we represent the deviatoric yield strength, Σ_q (other components of the macroscopic stress are null), obtained by putting $\xi_2 = \xi_3 = 0$ and $\xi_1 = +\infty$ in the parametric equations (55). A porosity $f = 0.1$ is considered and we compare the predictions

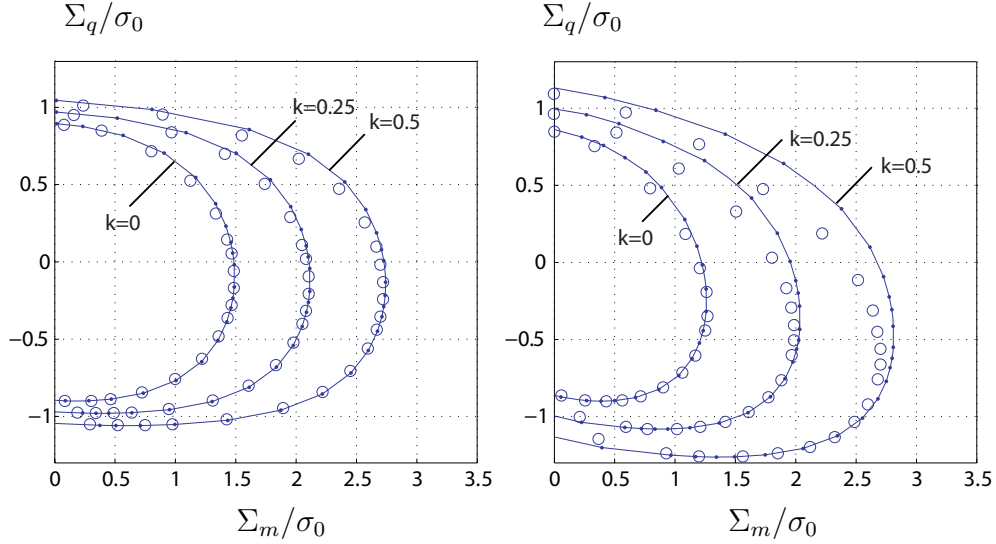


Fig. 10. Yield loci for a oblate cavity with an aspect ratio $a_1/b_1 = 1/2$ (at the left) $a_1/b_1 = 1/5$ (at the right) and a porosity $f = 0.1$. Comparison between the approximate (full line with points) and the exact, numerical, two-field criteria (circles), for various values of $k = \tau_0/(b_1\sigma_0)$.

for a prolate cavity (with the aspect ratio $a_1/b_1 = 5$) with those obtained for an oblate one (with the aspect ratio $a_1/b_1 = 1/5$). Also, in each case, the approximate value for the plastic limit is compared to the numerical exact one. The symbols used on figure 11 are the same as that of figures 9 and 10: the full line with the points corresponds to the approximate solution, whereas the circles are associated to the numerical two-field estimate. On the left, we present similar results for Σ_m (other components of the macroscopic stress are null), obtained by putting $\xi_1 = \xi_2 = \xi_3 = 0$ in the parametric equations (55). As expected, the yield strength increase when decreasing the voids size. It is worth noticing that the voids size effects seem to be more important in the case of an oblate cavity than for a prolate one. The reason is that for the same void porosity, the area of the void surface is greater for an oblate cavity than for a prolate one.

For completeness, we now propose to evaluate the effect of the interfacial residual stresses on the macroscopic yield surface of the nanoporous material. To this end, we come back to the case $\tau_0 = 0$ which leads to the closed-form expression (58). On figure 12, on the left, are represented the yield surfaces for a prolate cavity with an aspect ratio $a_1/b_1 = 1/2$, a porosity $f = 0.1$ and two values for the non dimensional parameter $k_r = \tau_r/(b_1\sigma_0)$. On the right, are represented similar results for the case of an oblate cavity having an aspect ratio $a_1/b_1 = 1/5$. A significant influence of the residual stress on the yield strength is observed. More precisely, the residual stresses induce a translation of the center of the macroscopic yield surface along both axis related to the macroscopic stress invariants Σ_m and Σ_q . This imply an asymmetry of the macroscopic yield strength between tension and compression. Again, it is

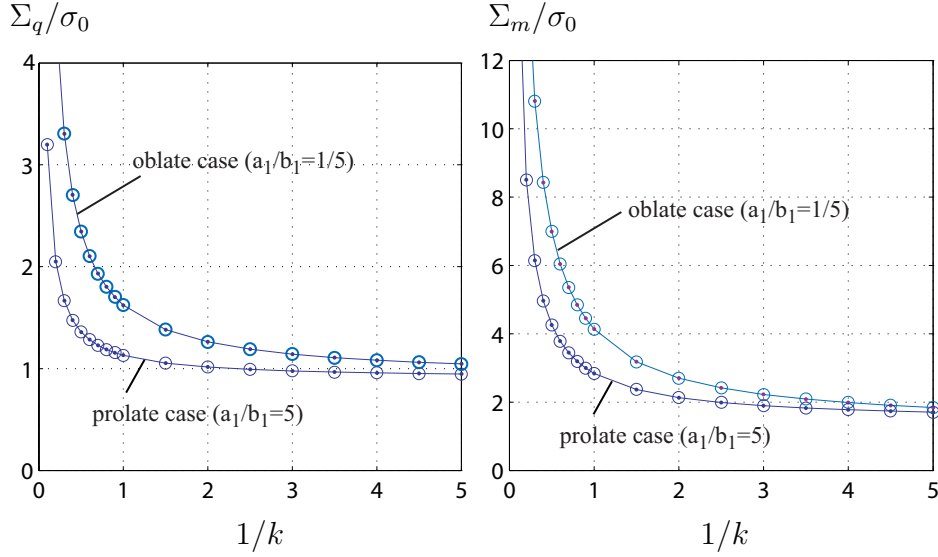


Fig. 11. On the left: variations of the shear yield strength for Σ_q as function of $1/k = b_1\sigma_0/\tau_0$. At the right: variations of the yield strength for Σ_m as function of $1/k = b_1\sigma_0/\tau_0$. The considered porosity is $f = 0.1$ and the aspect ratios are $a_1/b_1 = 5$ (prolate case) and $b_1/a_1 = 5$ (oblate case). The full line represents the approximate solution, the discrete points (circles) the numerical exact two-field solution.

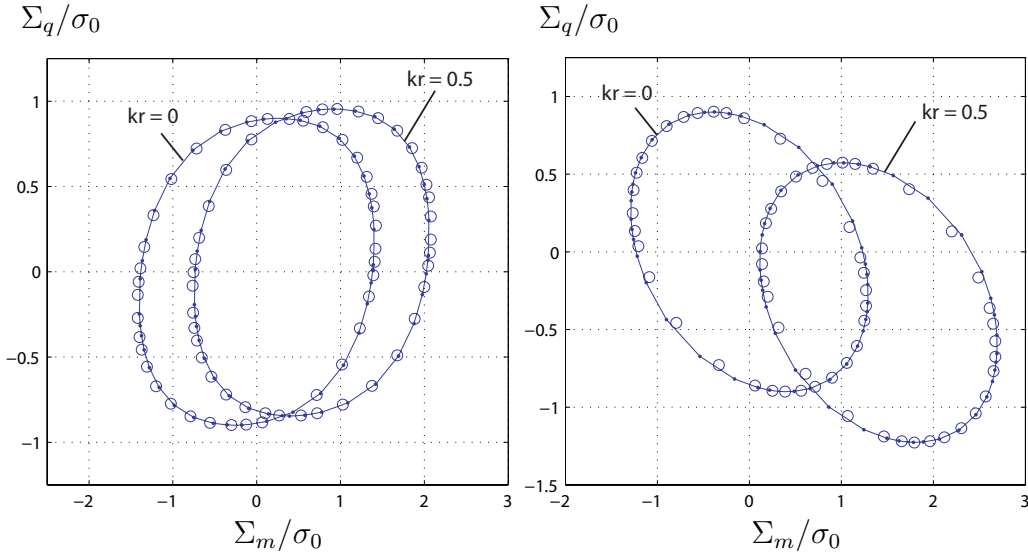


Fig. 12. Yield loci for a prolate cavity with an aspect ratio $a_1/b_1 = 5$ (on the left) for an oblate cavity with $a_1/b_1 = 1/5$ (on the right), a porosity $f = 0.1$ and for two values of the parameter $k_r = \tau_r/(b_1\sigma_0)$.

observed that the interfacial stress effects are more pronounced for an oblate cavity than for a prolate one (for the same reasons mentioned above).

6.4 The particular case of nano penny-shaped cracks

We consider now the case of a nanosized penny-shaped crack, corresponding to an oblate spheroidal nanocavity with an aspect ratio $\epsilon = a_1/b_1 \rightarrow 0$ or equivalently with an eccentricity $e_1 \rightarrow 1$. Since the volume V_1 of the cavity tends to zero, the ratio $\frac{\tau_0 S_1}{V_1}$ which appears in expression of $\Pi_2(\mathbf{D})$ and $\Pi_3(\mathbf{D})$, respectively given by (40) and (51), diverge. In the other hand, coefficients h_n for $n = 1..5$ which enters under the square root in (51) (and given by (52)) as well as the coefficients g_n which appears in (40) (and given by (41)) tend to zero when the limit $e_1 \rightarrow 1$ is considered. However, it can be shown that $\Pi_2(\mathbf{D})$ and $\Pi_3(\mathbf{D})$ have a finite limit when $e_1 \rightarrow 1$. Indeed, in the case of a penny-shaped crack, coefficients h_n are proportional to ϵ^2 and g_n are proportional to ϵ while the ratio $\frac{\tau_0 S_1}{V_1}$ is proportional to $\frac{1}{\epsilon}$. By computing carefully the expressions of the dissipation $\Pi_2(\mathbf{D})$ and $\Pi_3(\mathbf{D})$, we get:

$$\begin{aligned}\Pi_2(\mathbf{D}) &= \frac{3\tau_r}{2b_1}(\tilde{g}_1 D_m + \tilde{g}_2 D_q) \\ \Pi_3(\mathbf{D}) &= \frac{3\tau_0}{2b_1} \sqrt{\tilde{h}_1 D_m^2 + \tilde{h}_2 D_q^2 + 2\tilde{h}_3 D_m D_q + \tilde{h}_4 D_s^2}\end{aligned}\tag{68}$$

where coefficients \tilde{g}_n for $n = 1, 2$ are obtained by taking the limit of g_n/ϵ for $\epsilon \rightarrow 0$, while \tilde{h}_n for $n = 1, 2, 3, 4$ are obtained as the limit of h_n/ϵ^2 for $\epsilon \rightarrow 0$. Note that term involving D_t does not appear in $\Pi_3(\mathbf{D})$ (68); this is due to the fact that the limit h_5/ϵ^2 for $\epsilon \rightarrow 0$ is zero. The final expressions for \tilde{g}_n and \tilde{h}_n are:

$$\left\{\begin{array}{l} \tilde{g}_1 = \frac{3\pi}{2} - (1 - 3\alpha_2)d \\ \tilde{g}_2 = d \\ \tilde{h}_1 = 24 + \frac{9\pi^2}{4} - 3\pi(1 - 3\alpha_2)d + (1 - 3\alpha_2)^2 d^2 \\ \tilde{h}_2 = d^2 \\ \tilde{h}_3 = \frac{1}{2}(3\pi - 2d(1 - 3\alpha_2))d \\ \tilde{h}_4 = d^2 \end{array}\right.\tag{69}$$

in which d is the crack density parameter, $d = b_1^3/(a_2 b_2^2)$, introduced first by Bristow [8] and later considered by Budiansky and O'Connell [9] when studying elastic properties of cracked media. In figure 13, we represent the macroscopic criterion for a penny-shaped crack having a density parameter $d = 0.2$; various values of the non dimensional parameter $k = \tau_0/(b_1 \sigma_0)$ and

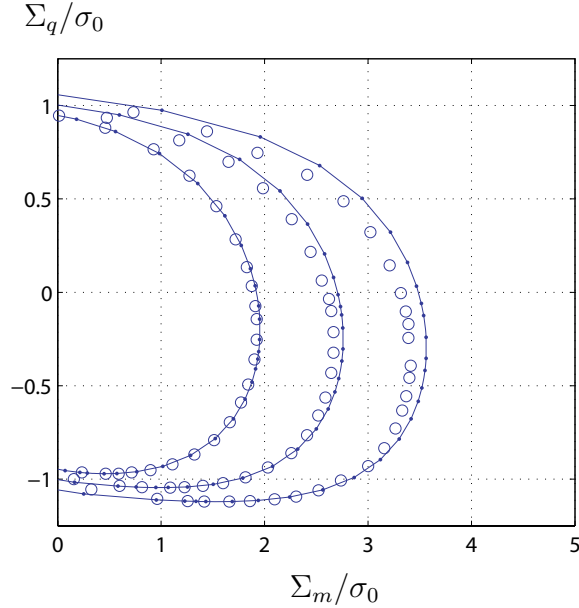


Fig. 13. Yield loci for a nano penny-shaped crack. Comparison between the approximate (full line with points) and the exact numerical two-field criteria (circles) for a crack density parameter $d = 0.2$ and for various values of $k = \tau_0/(b_1\sigma_0)$.

$\tau_r = 0$ are considered. Again, it can be observed that the approximate criterion (full line with points) compare well with the exact two-field criterion (circles).

7 Conclusion

The present paper has been devoted to the study of anisotropic ductile nanoporous materials. Particular emphasis has been put on interfacial stress and its influence on the overall yield criterion. The combined effects of void shape and size are investigated by considering prolate and oblate cavities with surface stresses. The theoretical development has been carried out by making use of limit analysis of a spheroidal unit cell made up of a von Mises solid matrix containing a confocal spheroidal cavity. To this end, the trial velocity field, introduced by Gologanu et al. [31,32], has been considered in order to derive the macroscopic yield criterion.

An original aspect of the analysis lies in the incorporation of a stress interface (between void and matrix) which accounts for a jump of the traction vector across the surface of the void. To this end, we took advantage of the plastic version of the Gurtin model, recently proposed by Monchiet and Bonnet [50], in which the jump of the traction vector at the void-matrix interface results from two contributions: surface residual stresses and plastic deformation at the void surface. In addition to the classical volumetric dissipation in the matrix, two different contributions to the macroscopic dissipation have been

identified; they correspond to interfacial residual stresses and plastic strain rates, respectively. Closed-form expressions of these interfacial contributions have been derived. This allowed to determine approximate parametric equations of the macroscopic yield function which, interestingly, points out the crucial role of the ratio between the void surface and the cell volume. These results extend those recently established by [18] which deal with ductile materials containing spherical nanovoids and in which voids size effects have been also demonstrated. In addition to the above results, numerical calculations, based on the considered two-field trial velocity have been performed in order to derive an upper bound for the macroscopic yield criterion. Illustrations has been provided for various shape of the cavity and various values of the material parameters which enter into the stress interface model. It has been shown that the stress interface has a significant effect on the macroscopic yield surface of the ductile material containing spheroidal nanocavities. In particular, the surface plastic deformation induces, at the macroscopic level, an increase of the yield strength. Concerning the interfacial residual stress, it leads to a change in position of the macroscopic yield surface, resulting in an asymmetry between tension and compression. Note that a good agreement has been noted between the predictions of the approximate parametric equations and that of the numerical upper bound. Finally, it has been shown that the surface effects on the macroscopic yield strength are more pronounced for an oblate cavity than for a prolate one.

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A Coefficients in $\Pi_1(\mathbf{D})$

In this appendix are provided the coefficients which enter in the closed-form expression of the plastic dissipation $\Pi_1(\mathbf{D})$ given by (30). Expressions of coefficients η and ζ are

$$\eta = \frac{\kappa^2(1+g)(f+g)(\alpha_2 - \alpha_1)}{(1-f)}; \quad \zeta = \frac{\kappa^2(1+g)(f+g)(\alpha_2 - \alpha_1)^2}{(1-f)^2} \quad (\text{A.1})$$

where g is defined by:

$$g = \frac{c^3}{a_2 b_2^2} = \frac{e_2^3}{\sqrt{1-e_2^2}} \quad (\text{A.2})$$

and coefficient κ is defined by:

$$\begin{aligned} \kappa &= \frac{1}{\ln(f)} \int_{e_1}^{e_2} \left[3(1+3\alpha^2) + \frac{6(1-3\alpha)}{3-2e^2} \right] \frac{3-2e^2}{e(1-e^2)} de \quad (\text{prolate}) \\ \kappa &= \frac{1}{\ln\left(\frac{g+f}{1+g}\right)} \int_{e_1}^{e_2} \left[3(1+3\alpha^2) + \frac{6(1-3\alpha)(1-e^2)}{3-e^2} \right] \frac{3-e^2}{e(1-e^2)} de \quad (\text{oblate}) \end{aligned} \quad (\text{A.3})$$

u_1 and u_2 are given by:

$$u_1 = \frac{f}{1+g}, \quad u_2 = \frac{f}{f+g} \quad (\text{A.4})$$

where $g = 0$ for the case of a prolate cavity and g is given by (A.2) for the case of an oblate cavity.

B Computation of \bar{P} , $\bar{\mathbb{P}}$ and \mathbb{S}

Performing the integral for θ over $[0, 2\pi]$ in (34), it comes:

$$\int_{\theta=0}^{\theta=2\pi} \mathbf{P}(\theta, \varphi) = \pi(1+w^2)\mathbf{I}_t + \pi(1-w^2)\mathbf{I}_n \quad (\text{B.1})$$

where w is defined by:

$$w = \frac{b \cos(\varphi)}{L_\lambda} \quad (\text{B.2})$$

For computing the second integral in (34), according to φ over the interval $[0, \pi]$, one can use the following result:

$$\frac{2\pi}{S_1} \int_{\varphi=0}^{\varphi=\pi} w^2 b_1 L_{\lambda_1} \sin(\varphi) d\varphi = \gamma \quad (\text{B.3})$$

Tensor $\bar{\mathbb{P}}$ is a transversely isotropic (with respect to the Ox_3 axis) and symmetric fourth order tensor. This tensor is then a priori defined by 5 independent coefficients. Moreover, by considering the following equalities:

$$\begin{aligned}
\bar{P}_{1111} + \bar{P}_{1122} + \bar{P}_{1133} &= \frac{1}{3}\bar{P}_{11} = \frac{1-\gamma}{6} \\
2\bar{P}_{1133} + \bar{P}_{3333} &= \frac{1}{3}\bar{P}_{33} = \frac{\gamma}{3} \\
\bar{P}_{1111} + \bar{P}_{1212} + \bar{P}_{1313} &= \frac{2}{5}\bar{P}_{11} = \frac{1-\gamma}{5} \\
2\bar{P}_{1313} + \bar{P}_{3333} &= \frac{2}{5}\bar{P}_{33} = \frac{2\gamma}{5}
\end{aligned} \tag{B.4}$$

One has only to compute one component of $\bar{\mathbb{P}}$ for having all other ones. For instance we compute \bar{P}_{3333} and we choose to define it as 2μ :

$$\bar{P}_{3333} = 2\mu = \frac{1}{S_1} \int_{\varphi=0}^{\varphi=\pi} \int_{\theta=0}^{\theta=2\pi} 2(1-w^2)^2 b_1 L_{\lambda_1} \sin(\varphi) d\varphi \tag{B.5}$$

Expression of μ is given by (48). After some algebraic calculations, it is easy to compute all the components of tensor $\bar{\mathbb{P}}$ from (B.4) and to put the results in the form given by (47) where the tensor of the Walpole basis are recall below:

$$\begin{cases}
\mathbb{E}_1 = \frac{1}{2}\mathbf{I}_t \otimes \mathbf{I}_t, & \mathbb{E}_2 = \mathbf{I}_n \otimes \mathbf{I}_n \\
\mathbb{E}_3 = \mathbf{I}_t \underline{\otimes} \mathbf{I}_t - \mathbb{E}_1, & \mathbb{E}_4 = \mathbf{I}_t \underline{\otimes} \mathbf{I}_n + \mathbf{I}_n \underline{\otimes} \mathbf{I}_t \\
\mathbb{E}_5 = \mathbf{I}_n \otimes \mathbf{I}_t, & \mathbb{E}_6(\underline{\xi}) = \mathbf{I}_t \otimes \mathbf{I}_n
\end{cases} \tag{B.6}$$

Finally, the components of tensor \mathbb{S} , given by (50), in the Walpole basis, are:

$$\mathbb{S} = \sum_n s_n \mathbb{S}_n \tag{B.7}$$

with:

$$\left\{ \begin{array}{l}
s_1 = 3\gamma[1 - 3\alpha_1 + 3f\alpha_2][1 - \alpha_1 + f\alpha_2] \\
\quad + \mu(1 - 3\alpha_1 + 3f\alpha_2)^2 \\
\quad + 6[1 - \alpha_1 + f\alpha_2][\alpha_1 - f\alpha_2] \\
s_2 = \frac{3\gamma}{2}[1 - 3\alpha_1 - 3f(1 - \alpha_2)][1 - \alpha_1 - f(1 - \alpha_2)] \\
\quad + \frac{\mu}{2}(1 - 3\alpha_1 - 3f(1 - \alpha_2))^2 \\
\quad + 3[\alpha_1 + (1 - \alpha_2)f][1 - \alpha_1 - f(1 - \alpha_2)] \\
s_3 = f^2 \left(\gamma + \frac{\mu}{2} \right) \\
s_4 = f^2 \left(\frac{5}{2}(1 - \gamma) - 2\mu \right) \\
s_5 = \frac{3\gamma}{2}[1 - 3\alpha_1 - 3f(1 - \alpha_2)][1 - \alpha_1 + f\alpha_2] + \frac{3f\gamma}{2} \\
\quad + \frac{\mu}{2}[1 - 3\alpha_1 + 3f\alpha_2][1 - 3\alpha_1 - 3f(1 - \alpha_2)] \\
\quad + 3[\alpha_1 + f(1 - \alpha_2)][1 - \alpha_1 + f\alpha_2] - \frac{3f}{2} \\
s_6 = s_5
\end{array} \right. \quad (\text{B.8})$$

C Expression of functions $\mathcal{G}_n(\xi_1, \xi_2, \xi_3)$

Starting from:

$$\left\{ \begin{array}{l}
\frac{\partial \Pi}{\partial D_n}(\mathbf{D}) = \sigma_0 f \left[\operatorname{arcsinh} \left(\frac{uY}{X} \right) \right]_{u=u_1}^{u=u_2} \frac{\partial Y}{\partial D_n} \\
\quad - \sigma_0 f \left[\frac{\sqrt{X^2 + u^2 Y^2}}{X} \right]_{u=u_1}^{u=u_2} \frac{\partial X}{\partial D_n}, \\
\quad + \frac{\tau_r S_1}{V_1} \frac{\partial W}{\partial D_n} + \frac{\tau_0 S_1}{V_1} \frac{\partial Z}{\partial D_n}
\end{array} \right. \quad (\text{C.1})$$

with $n = m, q, s, t$ and:

$$W = g_1 D_m + g_2 D_q \quad (C.2)$$

$$Z = \sqrt{h_1 D_m^2 + h_2 D_q^2 + 2h_3 D_m D_q + h_4 D_s^2 + h_5 D_t^2}$$

and X, Y are given by (31). By denoting ξ_n the triplet (ξ_1, ξ_2, ξ_3) , it follows that:

$$\left\{ \begin{array}{l} \mathcal{G}_1(\xi_n) = \frac{\sigma_0 \mathcal{U}(\xi_n)}{f\kappa} \left[1 + \frac{\eta}{3}(1 - 3\alpha_2) \right] \\ \quad - \frac{1}{3} \sigma_0 (1 - \zeta) \mathcal{V}(\xi_n) \left[\xi_1 + (1 - 3\alpha_2) \right] (1 - 3\alpha_2) \\ \quad + \frac{\tau_r S_1}{3V_1} g_1 + \frac{\tau_0 S_1}{3V_1} \frac{h_1 + h_3 \xi_1}{\mathcal{Z}(\xi_n)} \\ \mathcal{G}_2(\xi_n) = \frac{\sigma_0 \eta \mathcal{U}(\xi_n)}{f\kappa} - \sigma_0 (1 - \zeta) \mathcal{V}(\xi_n) \left[\xi_1 + (1 - 3\alpha_2) \right] \\ \quad + \frac{\tau_r S_1}{V_1} g_2 + \frac{\tau_0 S_1}{V_1} \frac{h_2 \xi_1 + h_3}{\mathcal{Z}(\xi_n)} \\ \mathcal{G}_3(\xi_n) = -\sigma_0 \mathcal{V}(\xi_n) \xi_2 + \frac{\tau_0 S_1}{V_1} \frac{h_4}{\mathcal{Z}(\xi_n)} \xi_2 \\ \mathcal{G}_4(\xi_n) = -\sigma_0 \mathcal{V}(\xi_n) \xi_3 + \frac{\tau_0 S_1}{V_1} \frac{h_5}{\mathcal{Z}(\xi_n)} \xi_3 \end{array} \right. \quad (C.3)$$

with:

$$\left\{ \begin{array}{l} \mathcal{U}(\xi_n) = f \left[\operatorname{arcsinh} \left(\frac{u \mathcal{Y}(\xi_n)}{\mathcal{X}(\xi_n)} \right) \right]_{u_1}^{u_2} \\ \mathcal{V}(\xi_n) = f \left[\frac{\sqrt{\mathcal{X}^2(\xi_n) + u^2 \mathcal{Y}^2(\xi_n)}}{u \mathcal{X}^2(\xi_n)} \right]_{u_1}^{u_2} \end{array} \right. \quad (C.4)$$

and:

$$\left\{ \begin{array}{l} \mathcal{X}(\xi_n) = \left[(1 - \zeta) \left(\xi_1 + (1 - 3\alpha_2) \right)^2 + \xi_2^2 + \xi_3^2 \right]^{1/2} \\ \mathcal{Y}(\xi_n) = \frac{1}{f\kappa} [3 + \eta(1 - 3\alpha_2)] + \frac{\eta}{f\kappa} \xi_1 \\ \mathcal{Z}(\xi_n) = \sqrt{h_1 + h_2 \xi_1^2 + 2h_3 \xi_1 + h_4 \xi_2^2 + h_5 \xi_3^2} \end{array} \right. \quad (C.5)$$

Consider the particular case of a spherical nanovoids for which the different coefficients, which enters into the definition of functions $\mathcal{G}_1(\xi_n), \mathcal{G}_2(\xi_n), \mathcal{G}_3(\xi_n), \mathcal{G}_4(\xi_n)$,

read:

$$\begin{aligned} \eta = \zeta = 0, \quad \kappa = \frac{3}{2}, \quad \alpha_2 = \frac{1}{3}, \quad g_1 = 2, \quad g_2 = 0, \\ h_1 = 4, \quad h_2 = h_3 = h_4 = \frac{3f^2}{5}, \quad S_1 = 4\pi a_1^2, \quad V_1 = \frac{4}{3}\pi a_1^3 \end{aligned} \quad (\text{C.6})$$

Consequently:

$$\left\{ \begin{aligned} \mathcal{G}_1(\xi_n) &= \frac{2\sigma_0 \mathcal{U}(\xi)}{3f} + \frac{2\tau_r S_1}{3V_1} + \frac{\tau_0}{a_1} \frac{4}{\sqrt{4 + \frac{3f^2}{5}\xi^2}} \\ \mathcal{G}_2(\xi_n) &= -\sigma_0 \mathcal{V}(\xi) \xi_1 + \frac{\tau_0}{a_1} \frac{9f^2}{5\sqrt{4 + \frac{3f^2}{5}\xi^2}} \xi_1 \\ \mathcal{G}_3(\xi_n) &= -\sigma_0 \mathcal{V}(\xi) \xi_2 + \frac{\tau_0}{a_1} \frac{9f^2}{5\sqrt{4 + \frac{3f^2}{5}\xi^2}} \xi_2 \\ \mathcal{G}_4(\xi_n) &= -\sigma_0 \mathcal{V}(\xi) \xi_3 + \frac{\tau_0}{a_1} \frac{9f^2}{5\sqrt{4 + \frac{3f^2}{5}\xi^2}} \xi_3 \end{aligned} \right. \quad (\text{C.7})$$

in which:

$$\left\{ \begin{aligned} \mathcal{U}(\xi) &= f \left[\operatorname{arcsinh} \left(\frac{2u}{f\xi} \right) \right]_f \\ \mathcal{V}(\xi) &= f \left[\frac{\sqrt{\xi^2 + \frac{4u^2}{f^2}}}{u\xi^2} \right]_f \end{aligned} \right. \quad (\text{C.8})$$

with $\xi = \sqrt{\xi_1^2 + \xi_2^2 + \xi_3^2}$. The mean stress Σ_m is given $\Sigma_m = \mathcal{G}_1(\xi_n)$ ($\mathcal{G}_1(\xi_n)$ being given by the first relation in (C.7)). The equivalent stress Σ_{eq} is obtained as a combination of the last three equations in (C.7):

$$\begin{aligned} \Sigma_{eq} &= \sqrt{[\mathcal{G}_2(\xi_n)]^2 + [\mathcal{G}_3(\xi_n)]^2 + [\mathcal{G}_4(\xi_n)]^2} \\ &= -\sigma_0 \mathcal{V}(\xi) \xi + \frac{\tau_0}{a_1} \frac{9f^2}{5\sqrt{4 + \frac{3f^2}{5}\xi^2}} \xi \end{aligned} \quad (\text{C.9})$$

D Derivation of the macroscopic criterion for $\tau_0 = 0$

We put $\tau_0 = 0$ in the expressions of functions $\mathcal{G}_1(\xi_n), \mathcal{G}_2(\xi_n), \mathcal{G}_3(\xi_n), \mathcal{G}_4(\xi_n)$ given by (C.3). It leads to:

$$\left\{ \begin{array}{l} \mathcal{G}_1(\xi_n) = \frac{\sigma_0 \mathcal{U}(\xi_n)}{f\kappa} \left[1 + \frac{\eta}{3}(1 - 3\alpha_2) \right] \\ \quad - \frac{1}{3}\sigma_0(1 - \zeta)\mathcal{V}(\xi_n) \left[\xi_1 + (1 - 3\alpha_2) \right] (1 - 3\alpha_2) + \frac{\tau_r S_1}{3V_1} g_1 \\ \mathcal{G}_2(\xi_n) = \frac{\sigma_0 \eta \mathcal{U}(\xi_n)}{f\kappa} - \sigma_0(1 - \zeta)\mathcal{V}(\xi_n) \left[\xi_1 + (1 - 3\alpha_2) \right] + \frac{\tau_r S_1}{V_1} g_2 \\ \mathcal{G}_3(\xi_n) = -\sigma_0 \mathcal{V}(\xi_n) \xi_2 \\ \mathcal{G}_4(\xi_n) = -\sigma_0 \mathcal{V}(\xi_n) \xi_3 \end{array} \right. \quad (\text{D.1})$$

Let us introduce:

$$\begin{aligned} \mathcal{G}_1^*(\xi_n) &= \mathcal{G}_1(\xi_n) - \frac{\tau_r S_1}{3V_1} g_1 \\ \mathcal{G}_2^*(\xi_n) &= \mathcal{G}_2(\xi_n) - \frac{\tau_r S_1}{V_1} g_2 \end{aligned} \quad (\text{D.2})$$

From the two first relations in (D.1), it can be easily shown that:

$$\mathcal{U}(\xi_n) = \frac{f\kappa}{\sigma_0} \left[\mathcal{G}_1^*(\xi_n) - \frac{1}{3}(1 - 3\alpha_2)\mathcal{G}_2^*(\xi_n) \right] \quad (\text{D.3})$$

Note also that, from equations (55) and (56), one has:

$$\mathcal{G}_1^*(\xi_n) = \Sigma_m - C_m, \quad \mathcal{G}_2^*(\xi_n) = \Sigma_q - C_q \quad (\text{D.4})$$

It follows that (D.3) can be rewritten as:

$$\mathcal{U}(\xi_n) = \frac{f\kappa T}{\sigma_0} \quad (\text{D.5})$$

with $T = \Sigma_m - C_m - \frac{1}{3}(1 - 3\alpha_2)(\Sigma_q - C_q)$. Using again relations (55) and (56), the last three equations in (D.1) can be read:

$$\begin{cases} \Sigma_q - C_q - \eta T = -\sigma_0(1 - \zeta)\mathcal{V}(\xi_n)\left[\xi_1 + (1 - 3\alpha_2)\right] \\ \Sigma_s = -\sigma_0\mathcal{V}(\xi_n)\xi_2 \\ \Sigma_t = -\sigma_0\mathcal{V}(\xi_n)\xi_3 \end{cases} \quad (\text{D.6})$$

which, by combination, leads to (see equations (C.4) and (C.5) for the definitions of $\mathcal{V}(\xi_n)$ and $\mathcal{X}(\xi_n)$):

$$\mathcal{V}(\xi_n)\mathcal{X}(\xi_n) = \frac{1}{\sigma_0} \left[\frac{(\Sigma_q - C_q - \eta T)^2}{1 - \zeta} + \Sigma_s^2 + \Sigma_t^2 \right]^{1/2} \quad (\text{D.7})$$

The last step of the demonstration consist to use the following identity:

$$\mathcal{V}^2(\xi_n)\mathcal{X}^2(\xi_n) + \frac{2f^2}{u_1 u_2} \cosh\left(\frac{\mathcal{U}(\xi_n)}{f}\right) - \frac{f^2}{u_1^2} - \frac{f^2}{u_2^2} = 0 \quad (\text{D.8})$$

which is obtained by eliminating $\mathcal{V}(\xi_n)$ in (C.4). Replacing u_1 and u_2 by their expressions given by (A.4) and considering relations (D.5) and (D.7), it easy to recover (58).